



电子核子碰撞实验

第二节:实验方法详解

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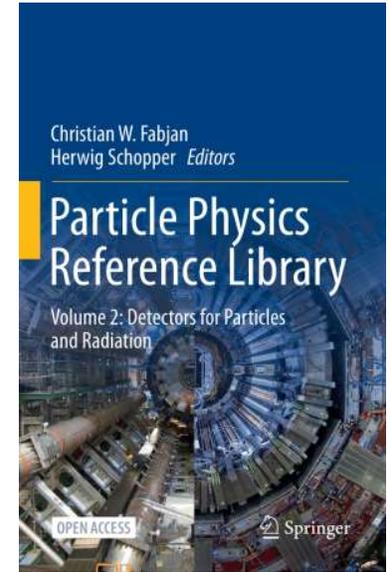
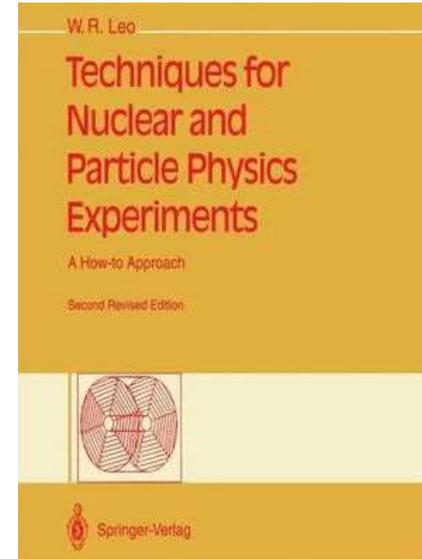
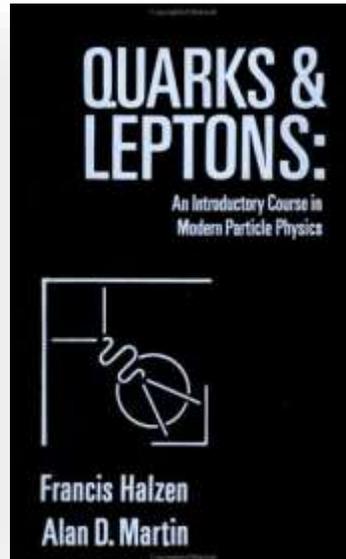
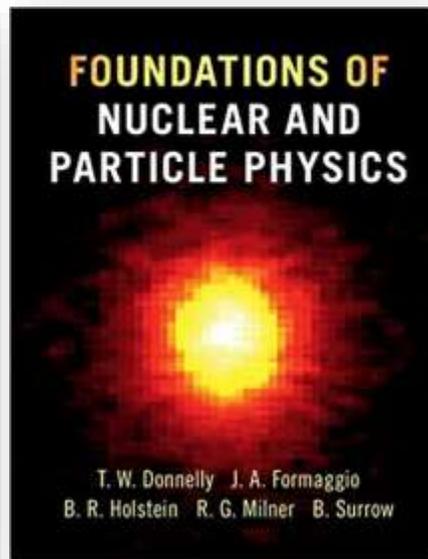
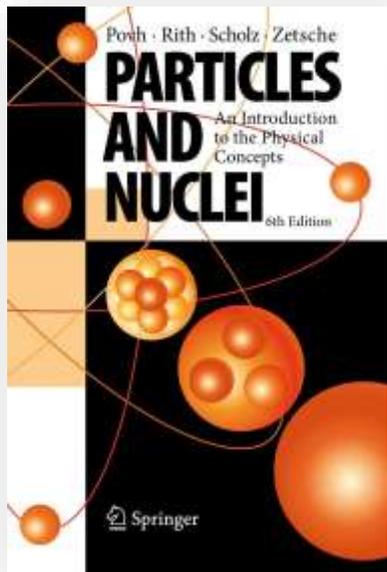
复旦大学2022年“优秀学生培养计划”粒子物理与核物理暑期学校
08/20/2022



课堂内容

- 第一节：物质结构概要
- 第二节：实验方法详解和实验实例(Hall-A Tritium Experiments)
- 第三节：未来展望

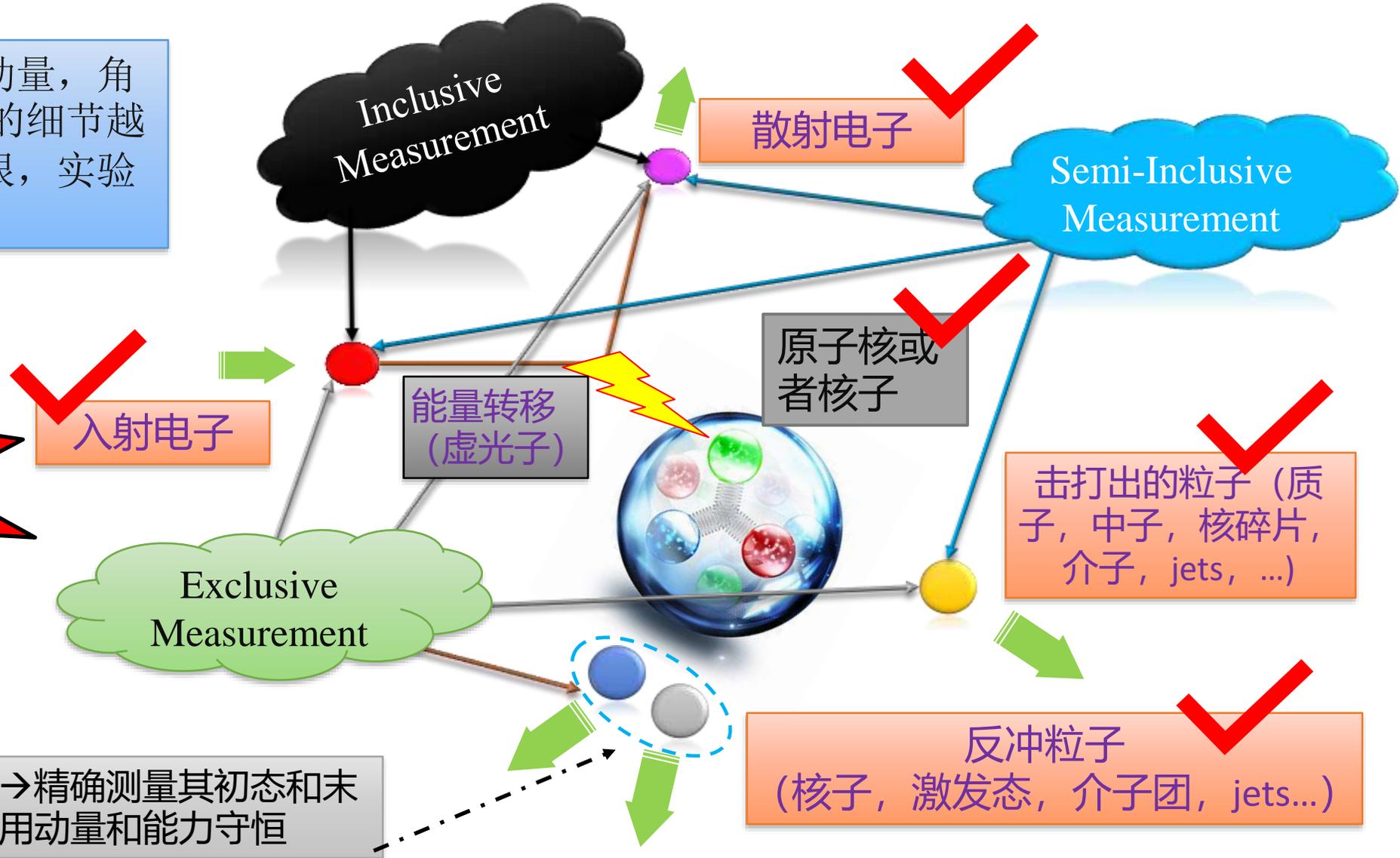
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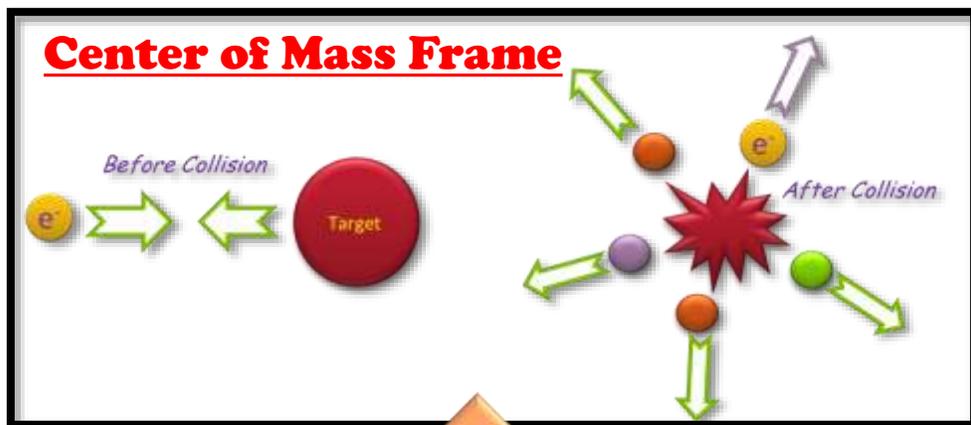
► 高能电子核子碰撞过程:

探测越多的终态粒子（动量，角度，粒子类型），了解的细节越多，但获取的信息越有限，实验也越有挑战！

如何观测？

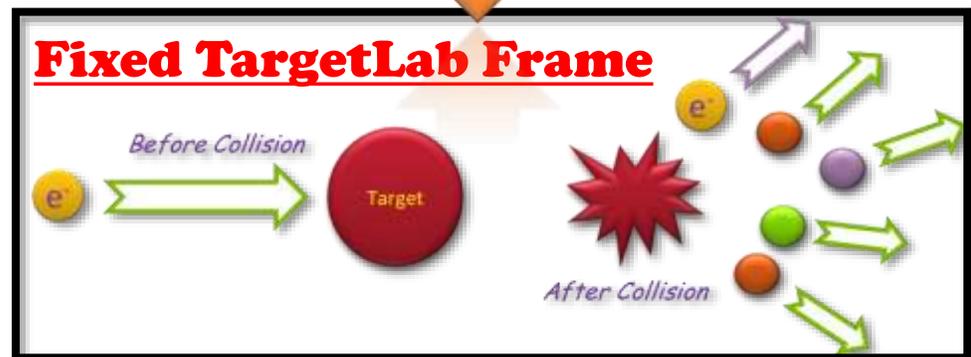


实验环境:

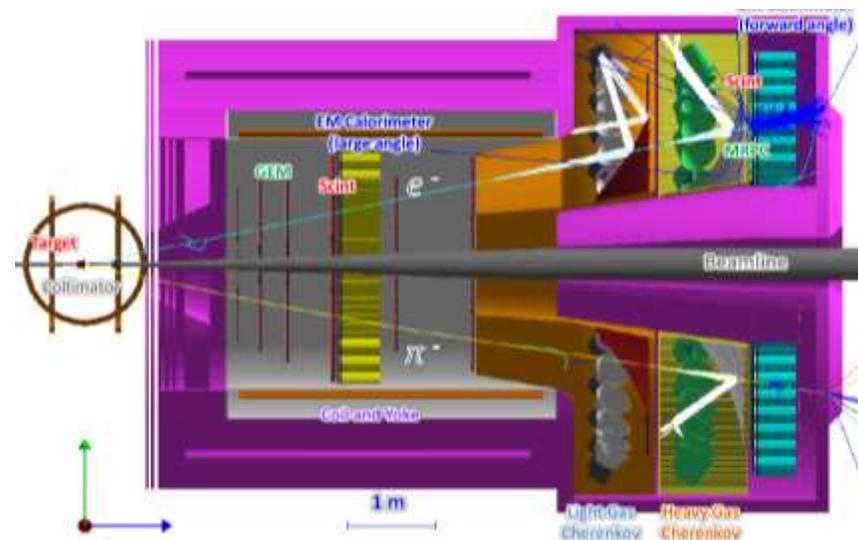
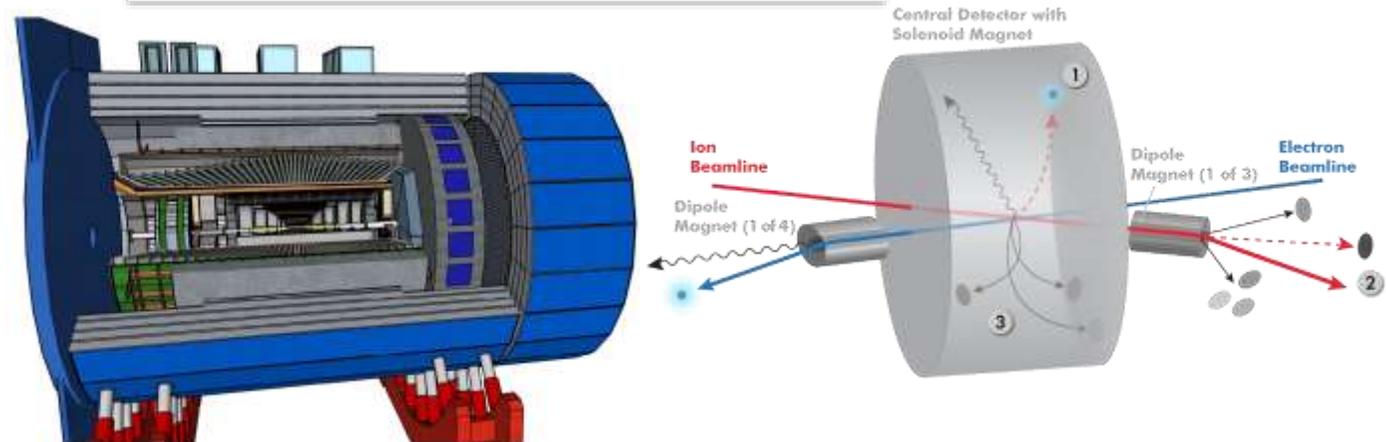


Lorentz Boost

Collider Frame



Electron-Ion Collider (EIC)



SoLID

➤ 从物理量到实验观测量: 如SIDIS

□ 入射和出射电子的动量参量 (Kinematic Variables)

- ✓ 入射和出射动量 ($P_{e0}, \theta_{e0}, \phi_{e0}, P_{e1}, \theta_{e1}, \phi_{e1}$)
- ✓ 高能电子加速器上: ($E_0, 0, 0, E', \theta_{e1}, \phi_{e1}$)
- ✓ 入射电子的极化率和极化方向

$$Q^2 = 4E_0 E' \sin^2(\theta/2) \quad \nu = E_0 - E'$$

$$y = \frac{E_l - E_l'}{E_l} = \frac{q \cdot P}{l \cdot P} \quad x_{bj} = \frac{Q^2}{2m_p \nu}$$

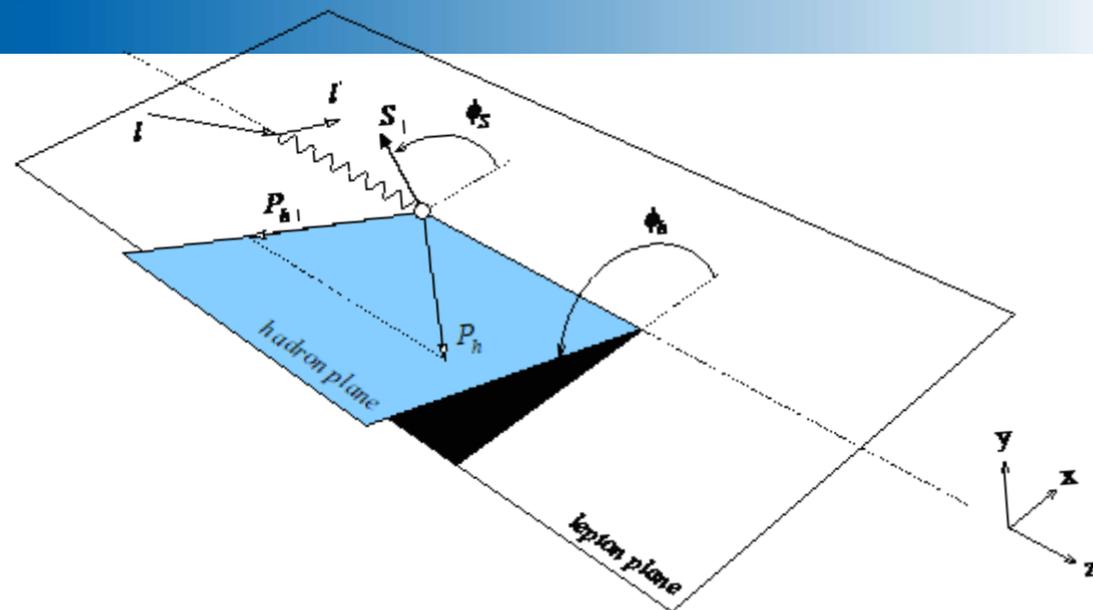
□ 核子(“靶”)的动量参量 (Kinematic Variables)

- ✓ 质量和动量 ($M_N, P_N, \theta_N, \phi_N$)
- ✓ 固定靶实验 ($M_N, 0, 0, 0$)
- ✓ 极化率和极化方向 (ϕ_T)

□ 末态强子 (介子, 核子) 的动量参量

- ✓ 质量和动量 ($M_h, P_h, \theta_h, \phi_h$)
- ✓ 强子极化方向

$$P_T = \frac{p \cdot P_h}{|q^2|} \quad z = \frac{P \cdot P_h}{P \cdot q}$$



□ 物理观测量: 如单举微分反应截面

探测到的散射电子数

Pi介子掺杂修正

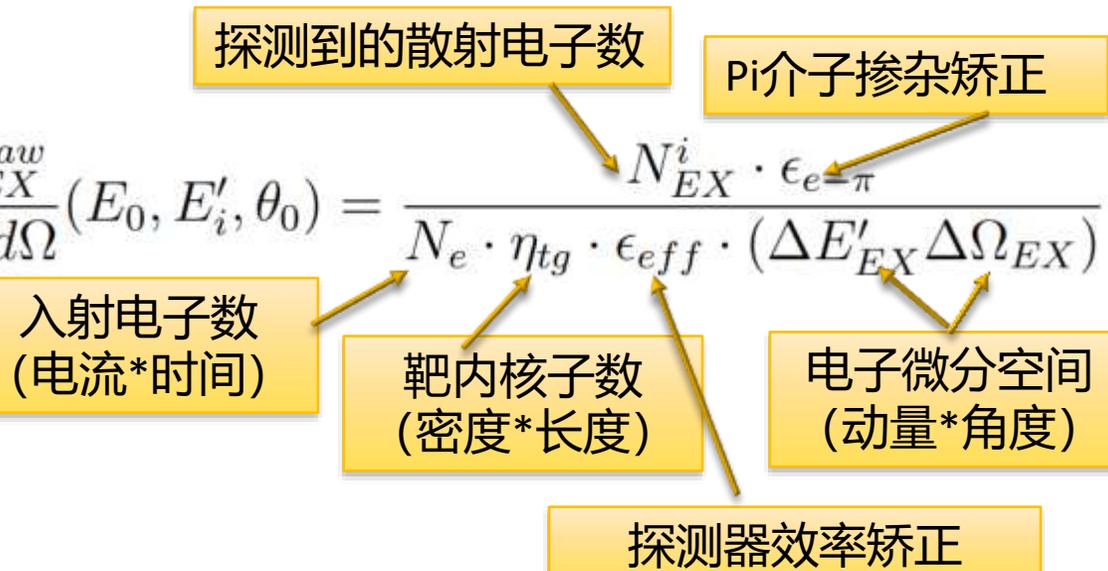
$$\frac{d\sigma_{EX}^{raw}}{dE' d\Omega}(E_0, E'_i, \theta_0) = \frac{N_{EX}^i \cdot \epsilon_{e=\pi}}{N_e \cdot \eta_{tg} \cdot \epsilon_{eff} \cdot (\Delta E'_{EX} \Delta \Omega_{EX})}$$

入射电子数
(电流*时间)

靶内核子数
(密度*长度)

电子微分空间
(动量*角度)

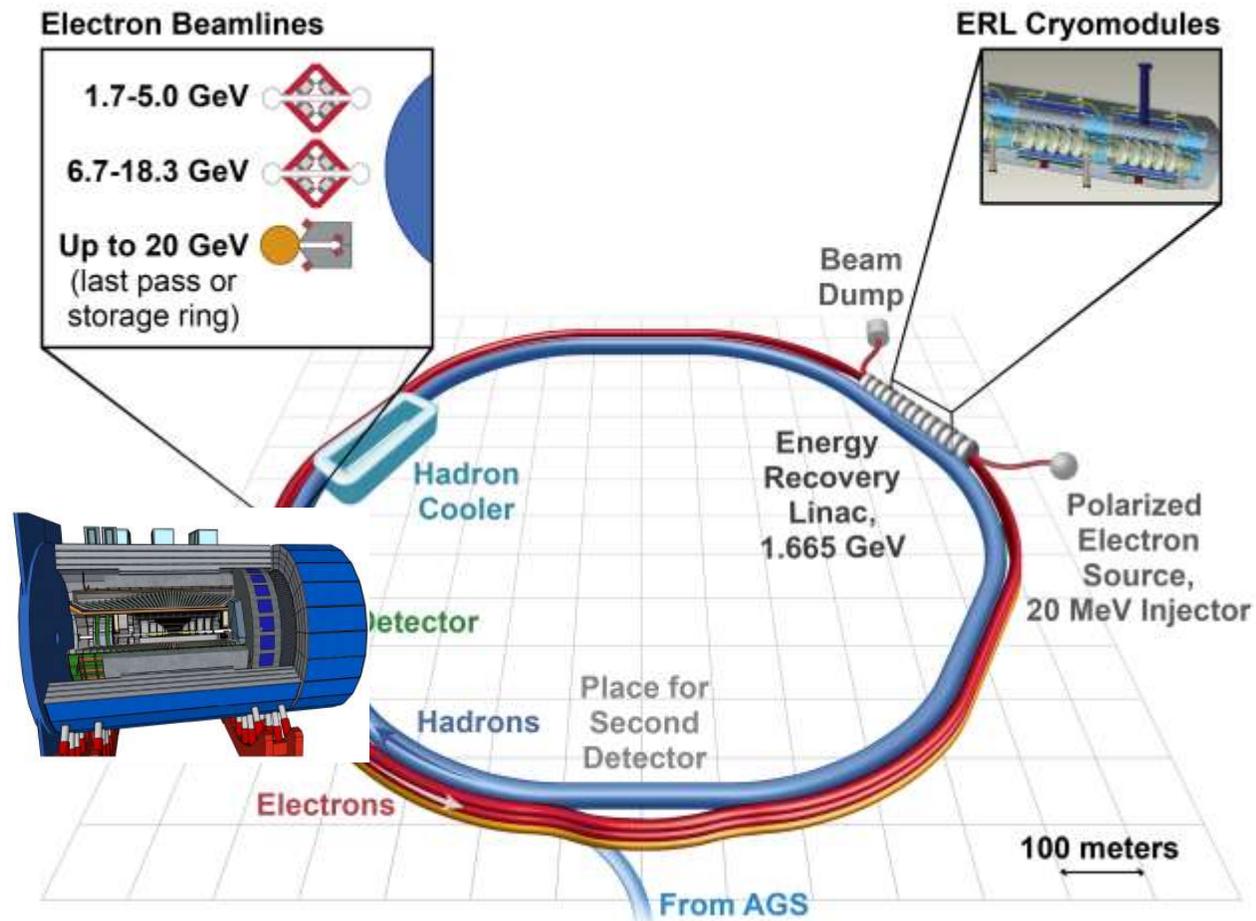
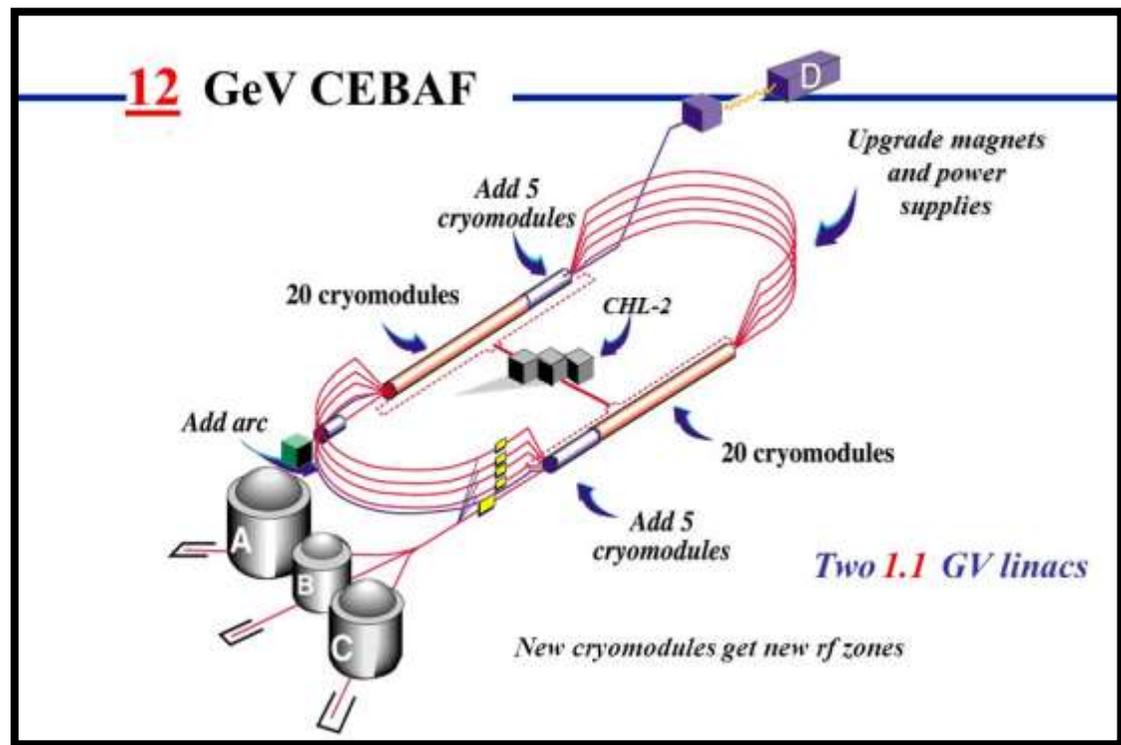
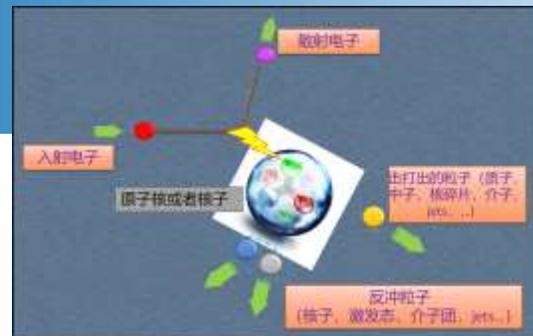
探测器效率修正



➤ 入射电子:

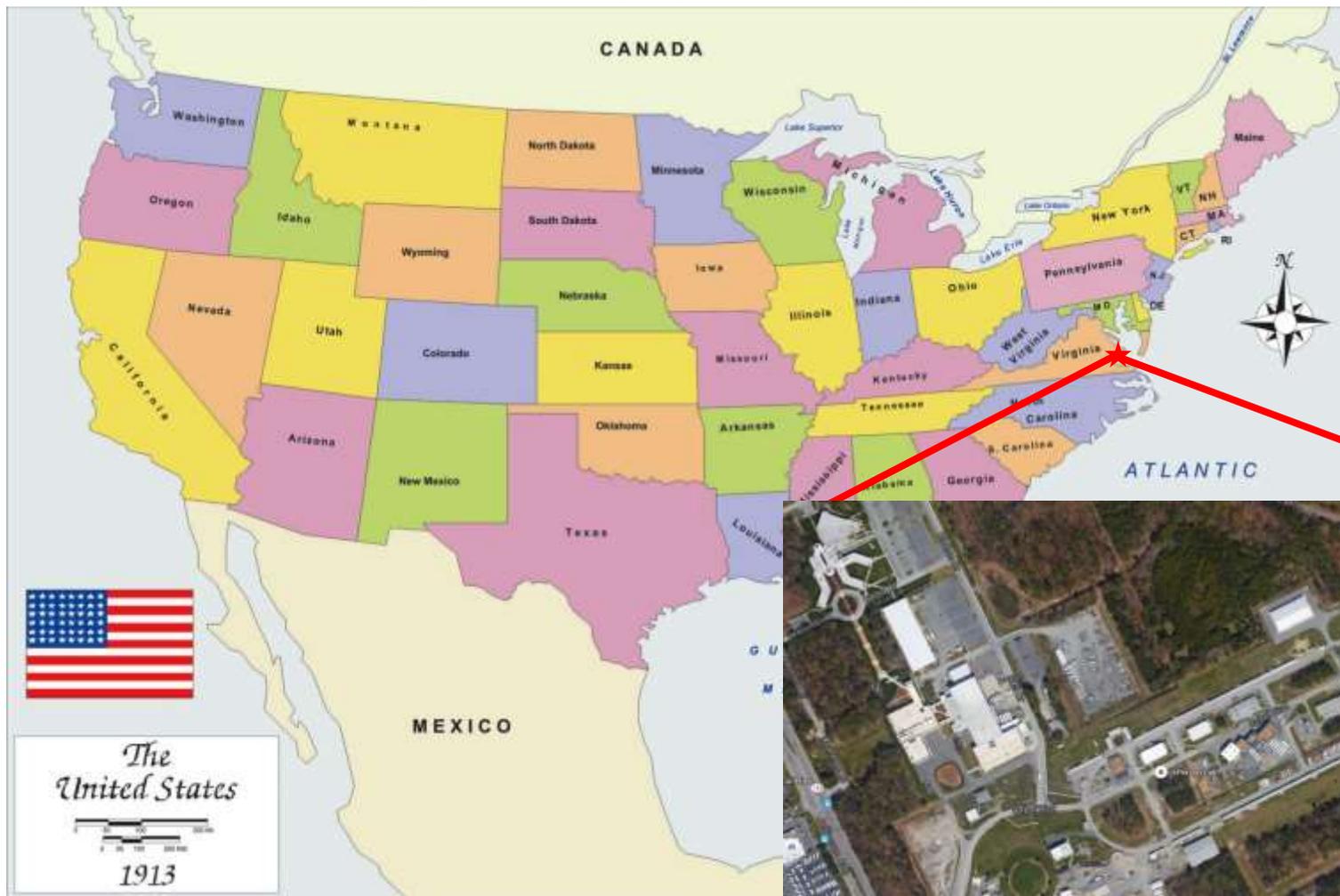
□ 入射电子的能量, 极化率和流强都通过调试加速器参数而固定

□ EIC中, 质子或者离子也由加速器提供

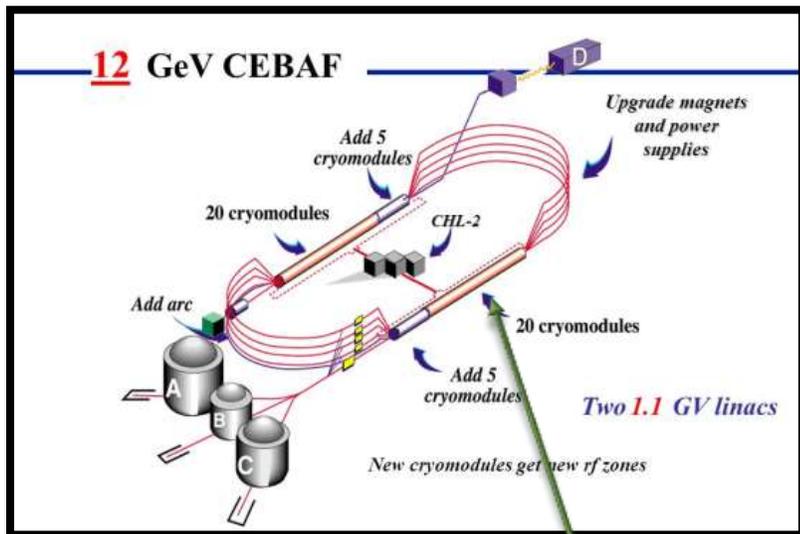


➤ Jefferson Lab坐落于美国弗吉尼亚州Newport News：

Located at Newport News, Virginia; Funded by Department of Energy; First operation in 1990s



➤ 电子加速器技术(Jefferson Lab):



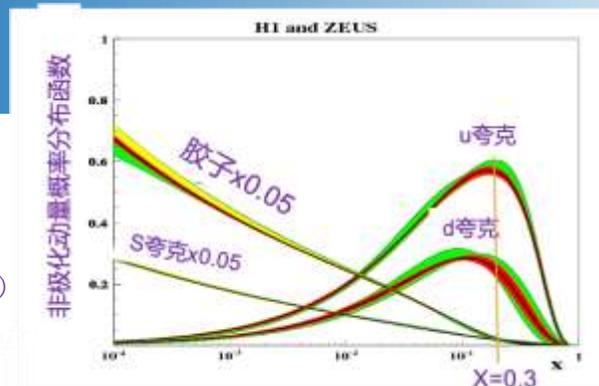
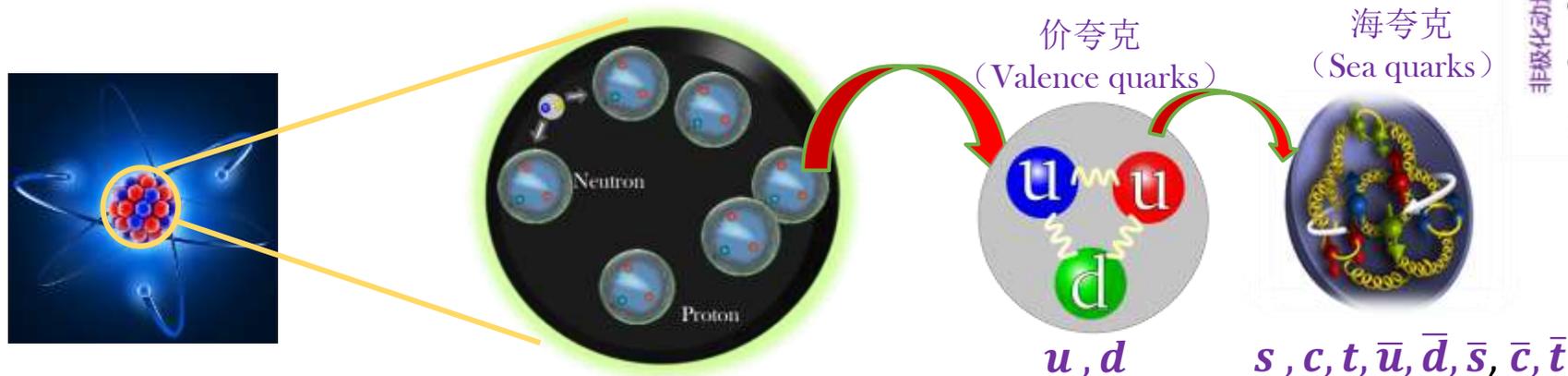
ARC



Superconducting Radio-Frequency Cavity

- ❖ 高流量电子加速器
- ❖ 低温超导加速腔
- ❖ 纵向和横向极化可控
- ❖ 同时输送连续200uA电子
- ❖ 无线电快频率分离技术可将不同能量，不同电流的电子输送到4个实验大厅同时做实验

➤如何研究夸克？：



□ 现实: 无法直接观测到自由存在的夸克!

- 通过测量 质子 (uud) 和中子 (udd) 来测量 u 和 d 夸克的性质
价夸克: 在原子中, “u” 和 “d” 夸克被间接观测到的几率最大 (携带高动能的比例最大)

□ 然而: 实验上没有自由中子可以使用 (中子的衰变周期是15分钟)

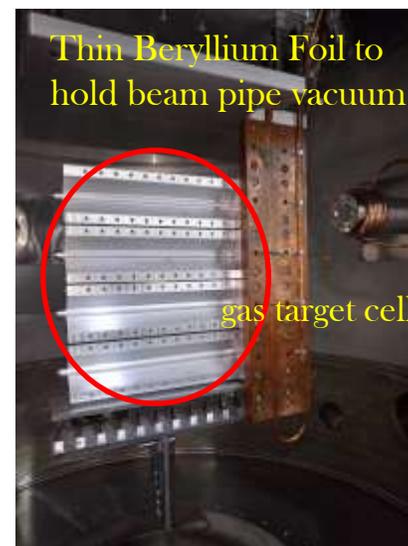
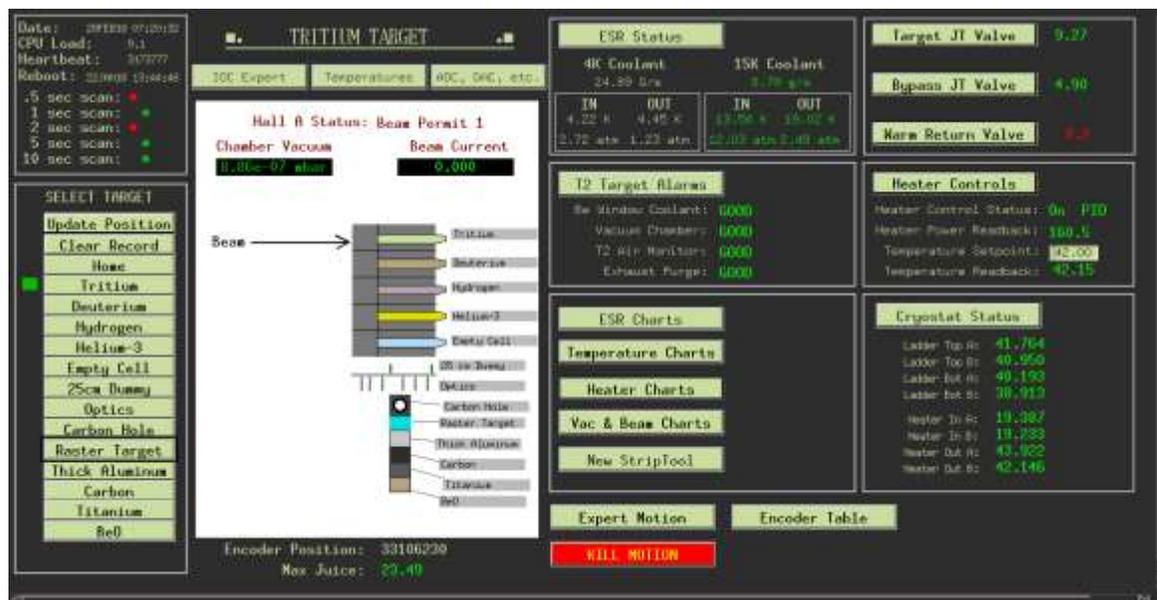
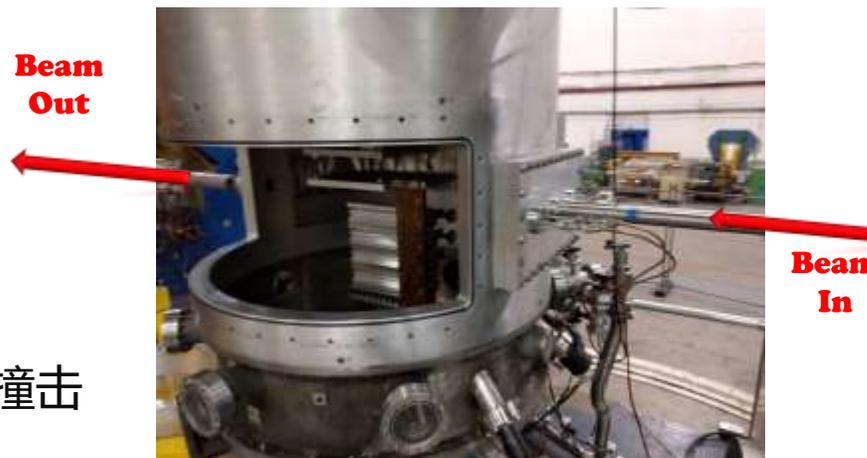
- 一般用Deuterium (p+n) 或者 Helium-3 (2p+n) 作为有效中子
- 然而: 核子里面的中子是时刻运动的 (Fermi-Motion); 可能和自由中子不一样 (EMC effect).

□ 海夸克: 比u和d重的夸克, 观测几率小 (携带高动量的比例小), 通过真空能和胶子衰变涨落

- 可以通过观测重介子 (e.g., $K^+(u-sbar)$, $K^-(ubar-s)$), 或者亚稳定的重强子衰变

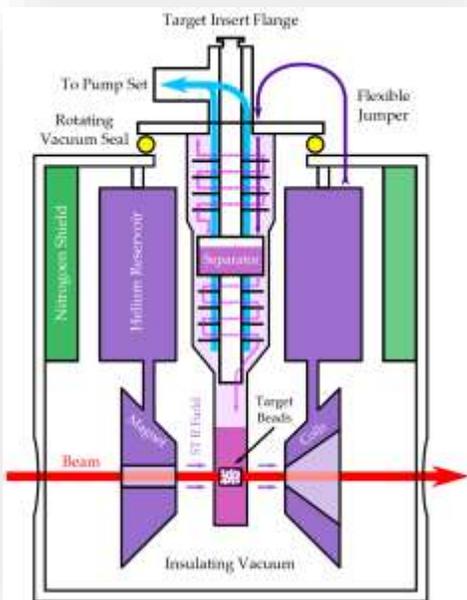
➤ 非极化固定靶

- ❖ 超低温真空腔存储多种气体，液体和固体靶
- ❖ 电子输送管道直接穿过真空腔
- ❖ 避免靶的氧化和高温气化
- ❖ 远程控制不同靶（或者不同原子核）到电子束流线上被撞击



➤极化固定靶:

Polarized NH₃ Targets



“proton”

重要参数

- ✓ 系统半径: 1.2m(W) x 2m(H)
- ✓ 磁场强度: 5 Tesla
- ✓ 靶长: ~3 cm
- ✓ 亮度: $10^{35} \text{ cm}^{-2} \text{ s}^{-1}$
- ✓ 最高极化率: >90%
- ✓ 平均极化率: 70%
- ✓ 稀释效应: ~0.13
- ✓ 运行温度: 1K

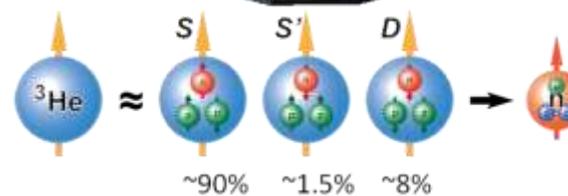
- 经过多年的实验，非常成熟的技术
- 达到很高极化率，但是实验环境要求也很高
- JLab的主要极化靶
- UVA的极化靶组是最顶尖的

<https://userweb.jlab.org/%7Eckeith/Frozen/Frozen.html>

<http://twist.phys.virginia.edu/>

Polarized He₃ Targets

“neutron”



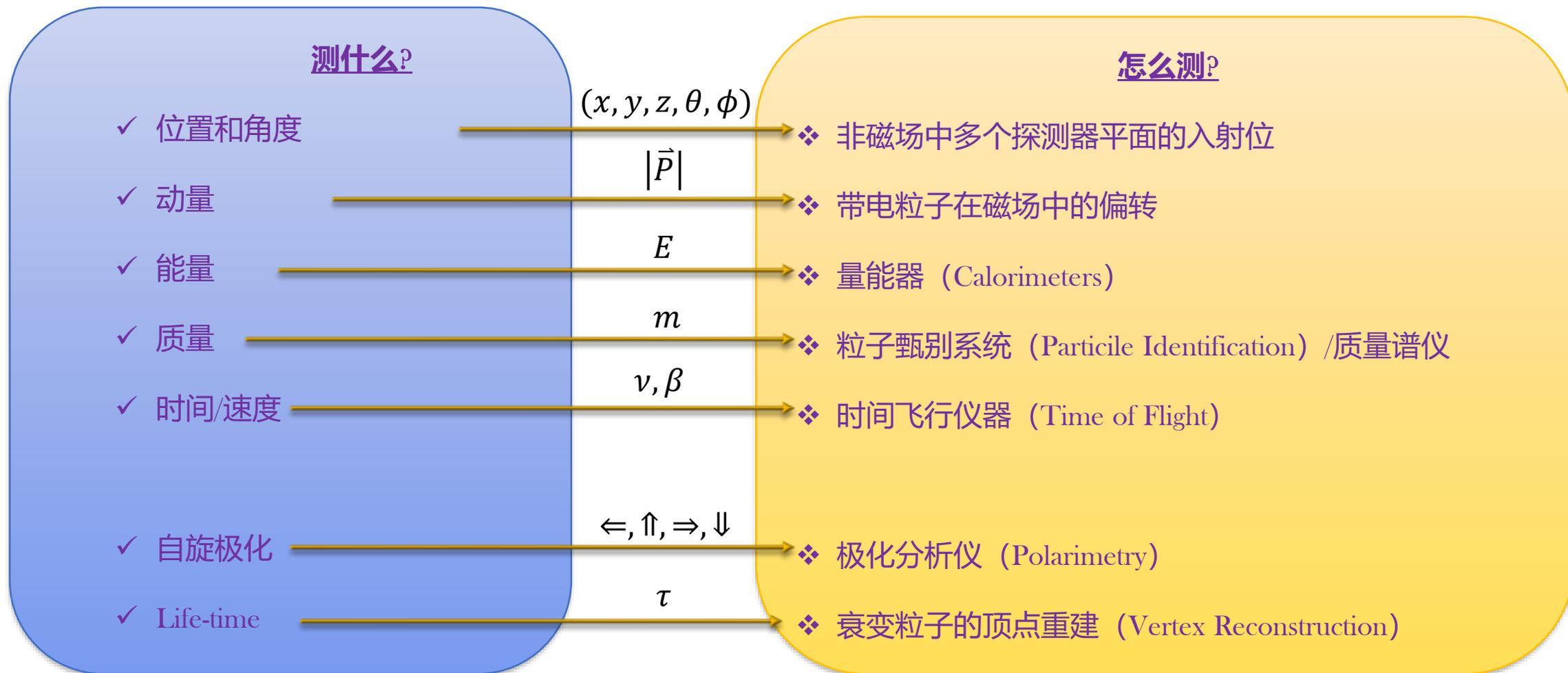
重要参数

- ✓ 系统半径: R=1.83m(outer)
- ✓ 磁场强度: 25~30 Gauss
- ✓ 靶长: 40 cm
- ✓ 亮度: $10^{36} \text{ cm}^{-2} \text{ s}^{-1}$
- ✓ 最高极化率: >70%
- ✓ 平均极化率: ~60% @ 15uA
- ✓ 稀释效应: 0.3
- ✓ 运行温度: ~20C

- 多个极化方向 (longitudinally and transversely polarized)
- 1998年 开始在Jlab使用 (Jian-Ping Chen)
- 已经在13个实验使用过
- 将会在7个以上已经批准的12GeV实验上面使用
- 还有不少改进空间
- 有很宽的医学物理和其他方面的应用

http://hallaweb.jlab.org/equipment/targets/polhe3/polhe3_tgt.html

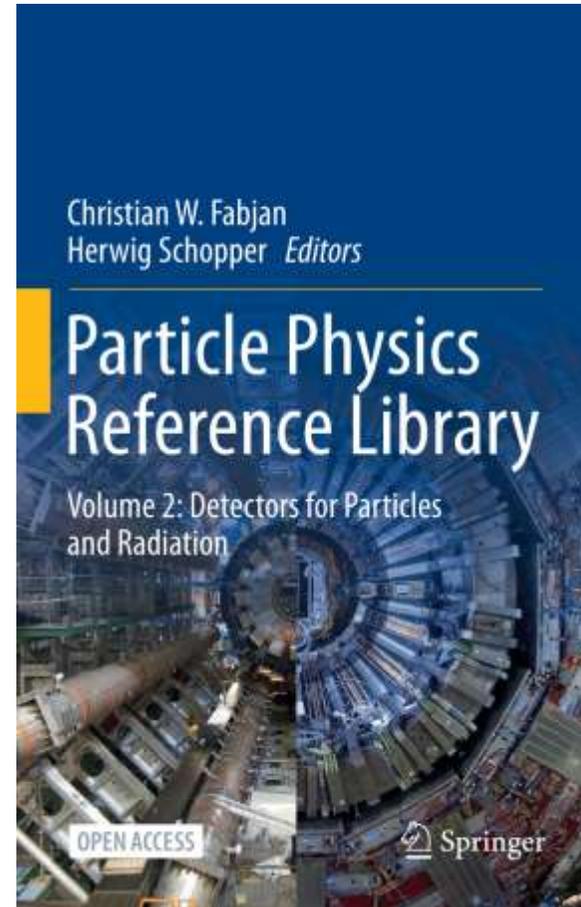
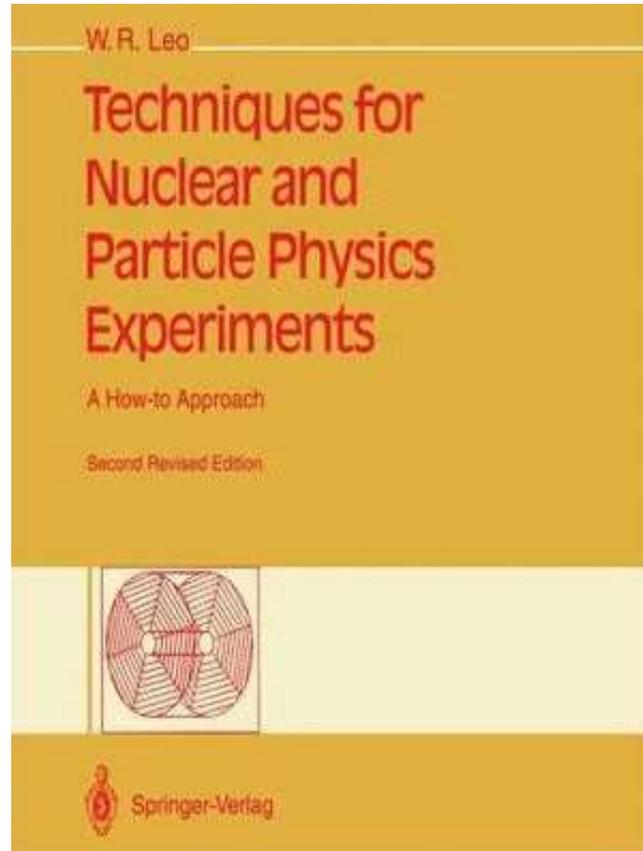
➤ 末态粒子信息



❑ 需要不同高能粒子探测器的相互配合

➤ 部分探测器举例 → **Backup#42~#57**

□ 关于探测器, 模拟, 数据分析的推荐书籍

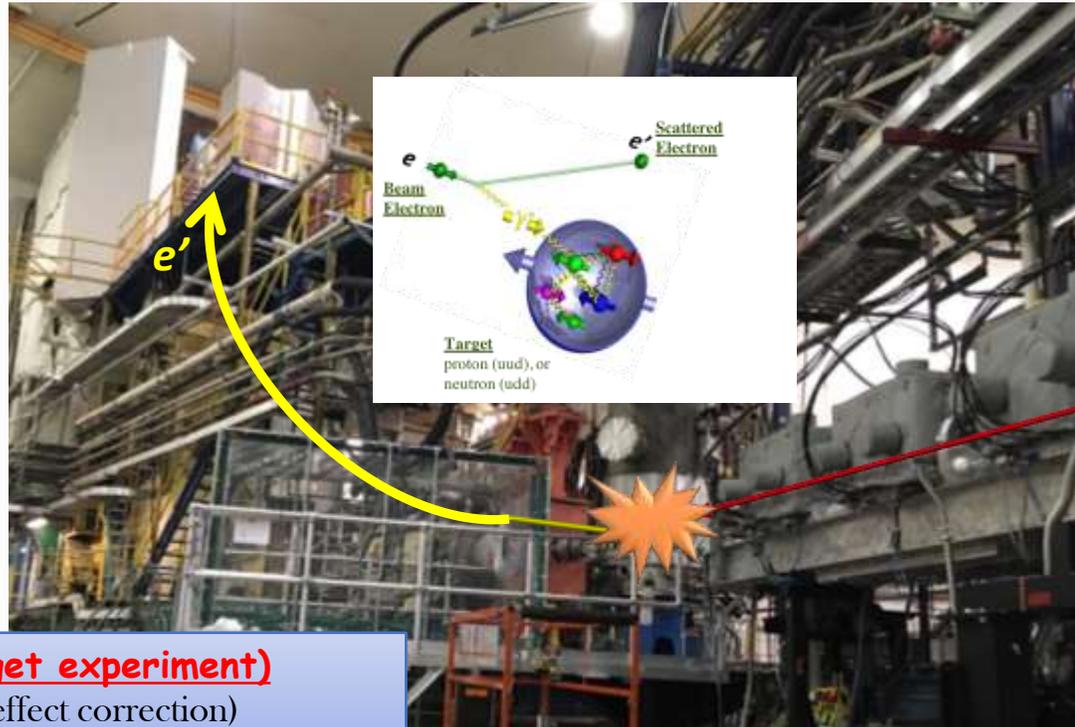


➤ 单粒子探测

- ❖ 参与碰撞反应的粒子一般要离开碰撞点飞行一段时间/距离再由探测器测量
- ❖ 探测器信号的产生时间和粒子到达时间可能有不同时（探测器尺寸，反应时间，信号传输线路等）
- ❖ 需要触发系统判断碰撞发生时间，并通知数据采集系统收集并同步所有探测器信号

Detectors

- ✓ Position
- ✓ Angle
- ✓ Time
- ✓ Speed
- ✓ Momentum
- ✓ Energy
- ✓ Particle Type (Mass, Charge)



e Beam

- ✓ Current
- ✓ Energy
- ✓ Position
- ✓ Polarization
- ✓ Arrival Time

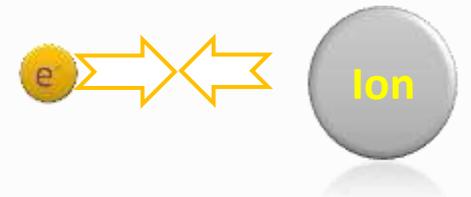
Target (fixed target experiment)

- ✓ In-Beam Density (boiling effect correction)
- ✓ Polarization (for a polarized target only)
- ✓ Target Atoms (if mixed materials)
- ✓ Reaction Point (for a long target)

Ion Beam (EIC)

- ✓ Current
- ✓ Energy
- ✓ Position
- ✓ Angle
- ✓ Polarization
- ✓ Arrival Time

Collider Mode



➤ 单粒子探测

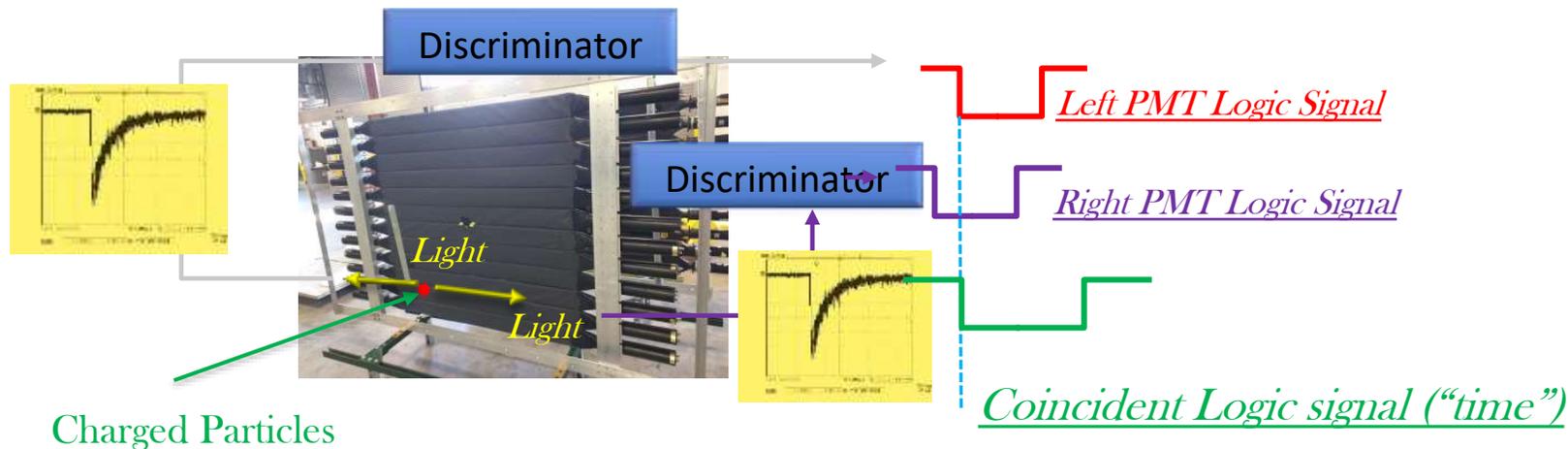
❖ 数据采集系统包括：

- ✓ 触发系统 → 判断碰撞事件发生
- ✓ 前端电子学 → 收集探测器信号（如PMT的脉冲信号）
- ✓ 电脑系统 → 收集所有系统信号（束流，靶信息等），判断触发信号的有效性，存储前端电子学输送的探测器信号

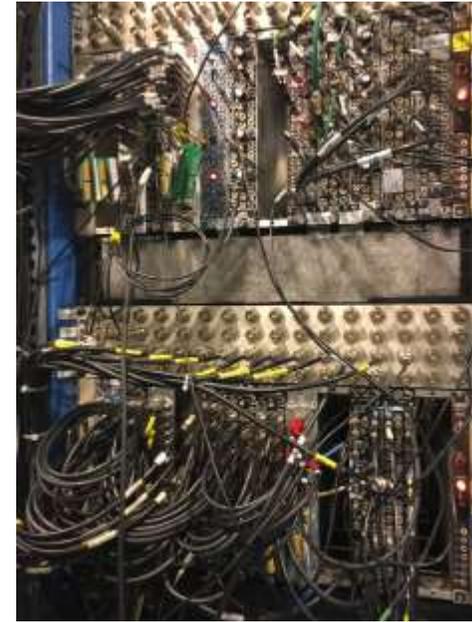
❖ 不同的探测器系统需要不同前端电子学

❖ 不同实验，不同粒子探测，需要设计不同触发系统

例子：如何设计触发信号，判断带电粒子穿过塑料闪烁体探测器 时间？

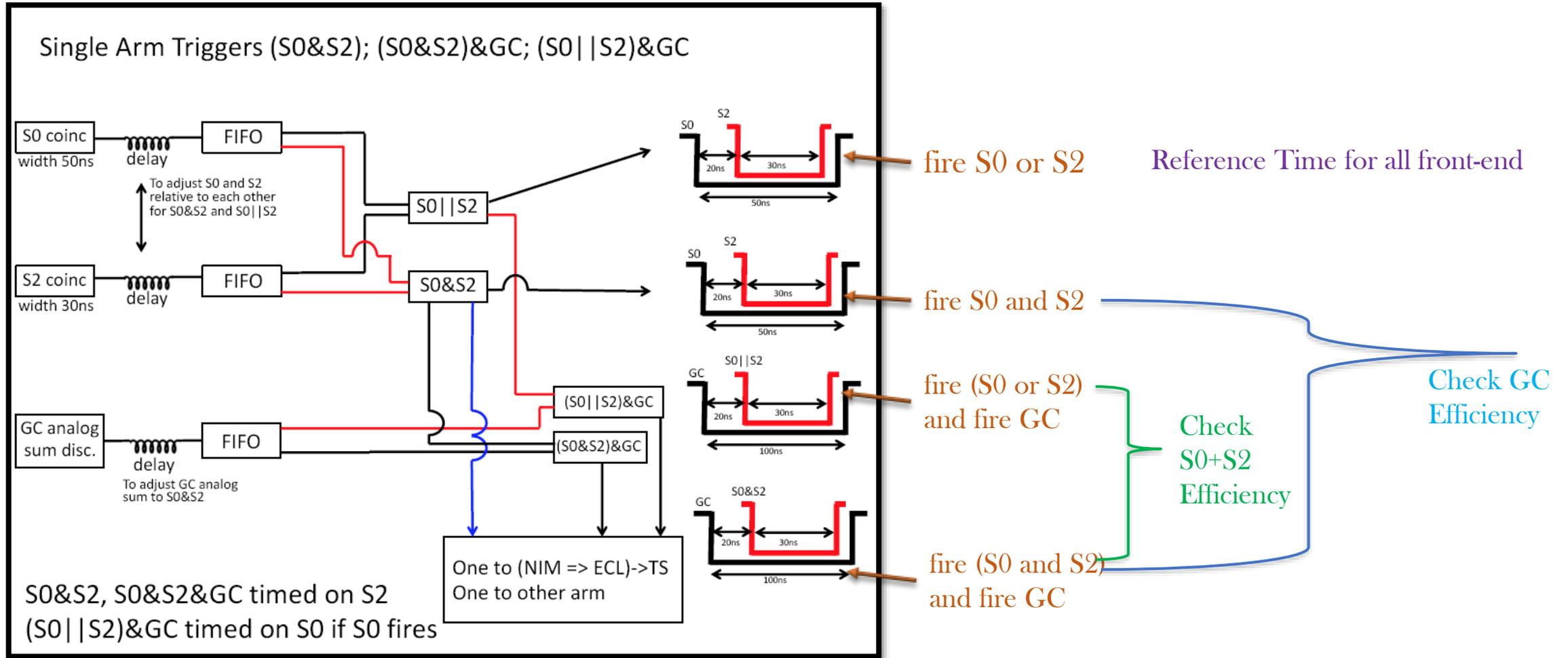


Note: delay time due to different cable lengths should be considered



➤ 触发系统设计:

- ❖ 排除背景信号和偶然触发信号, 需要两组探测器“同时”看到同一信号
- ❖ 排除其他粒子信号, 需要加入粒子甄别探测器的触发信号 (如CER可以保留电子, 去除Pi介子)
- ❖ 根据实验需求, 设计各种复杂的触发系统

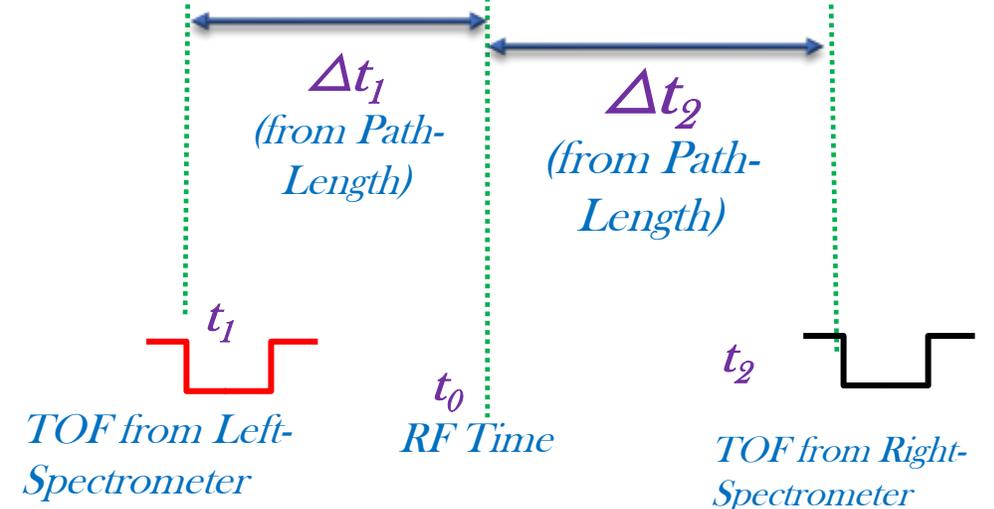
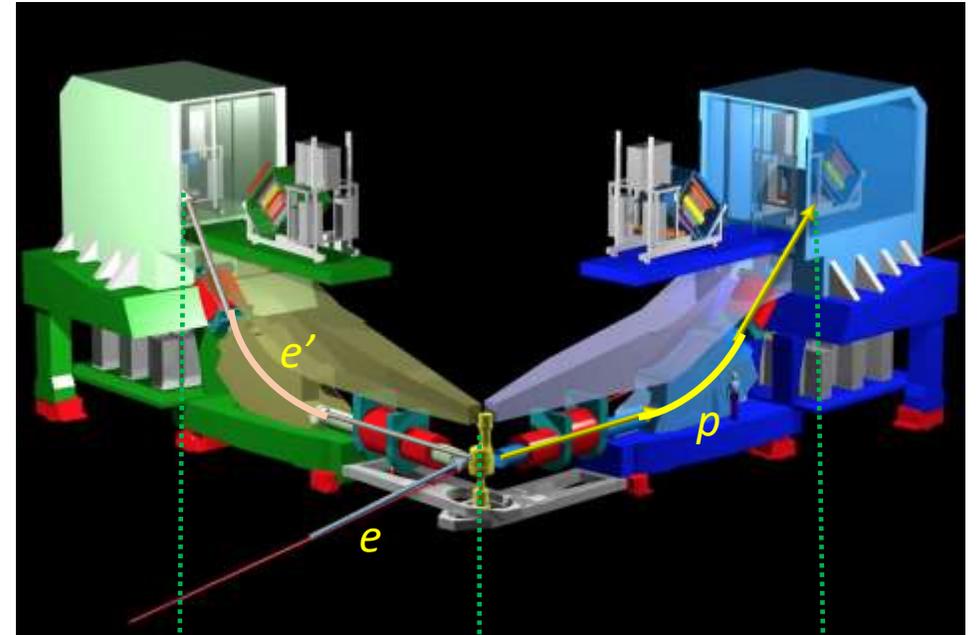
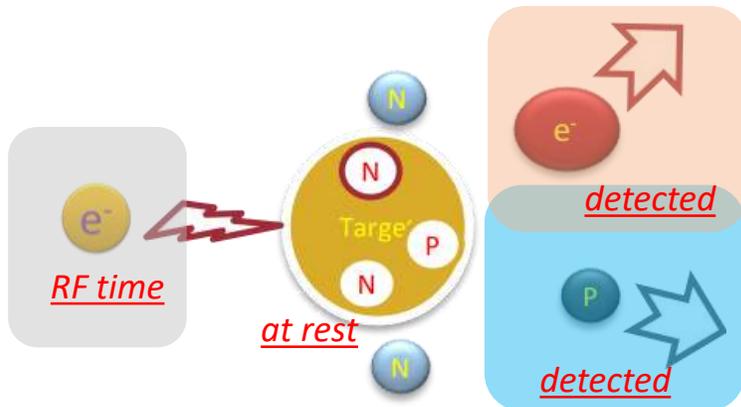


➤ 双（多）粒子探测

- ❑ SIDIS & Exclusive Measurements → 需要判断测量的末态粒子来自同一个碰撞反应
- ❑ 入射电子束流非绝对连续, 如Jlab, 每2ns来一束电子团簇 (beam bunch)
- ❑ 一个团簇一般之多一个电子会与原子核碰撞 → 团簇到达时间 (RF Time) == 碰撞发生时间



- ❑ 末态粒子由探测器测量, 时间根据飞行距离反推到反应点



➤ 实验控制室：

□ 实验运行是24/7，持续几个月到几年→需要合作者在控制室轮流值班



□ 值班人员(Jlab为8 hrs一班)：

- 采集数据；
- 监控探测状态
- 初步数据分析
- 修改实验设置
- 记录各种信息
- 与多部门之间沟通



□ 专家（主要为博士生和博士后）

- 为值班人员提供支持
- 深入数据分析
- 维修探测器
- 解决其他问题

❖ 其他人员：

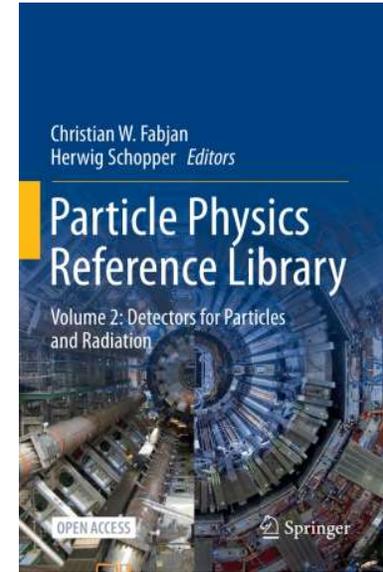
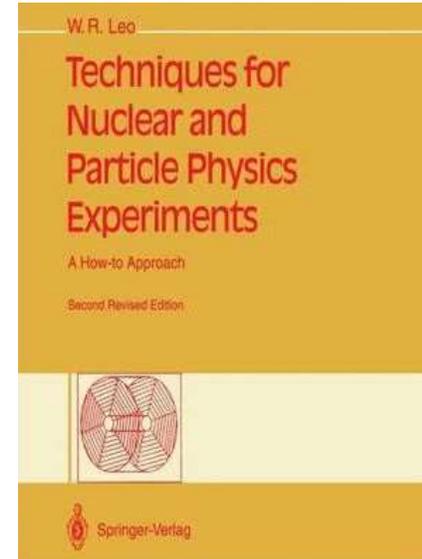
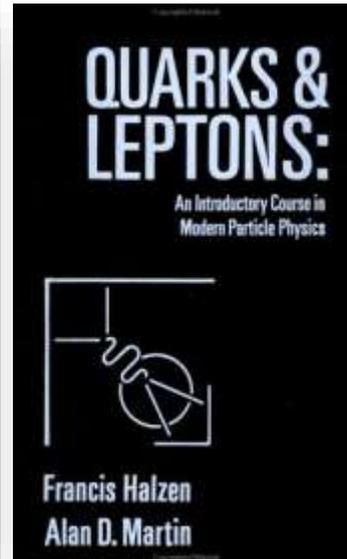
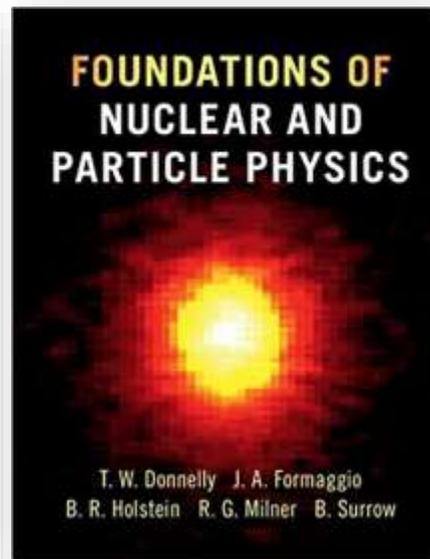
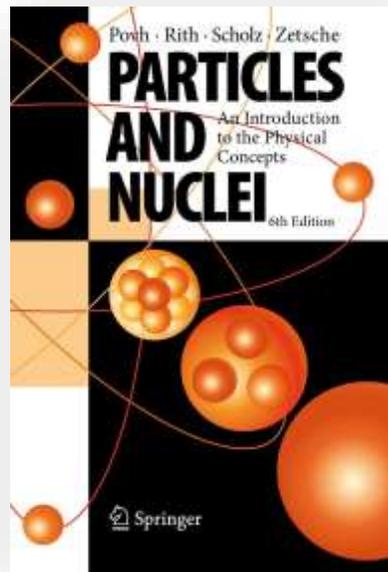
- 实验发言人，实验室管理层→定制实验计划，多部门沟通协调
- 实验协调员：所有信息沟通的节点
- 支持团队（加速器，靶，电子学，探测器，低温制冷，辐射防护等）



课堂内容

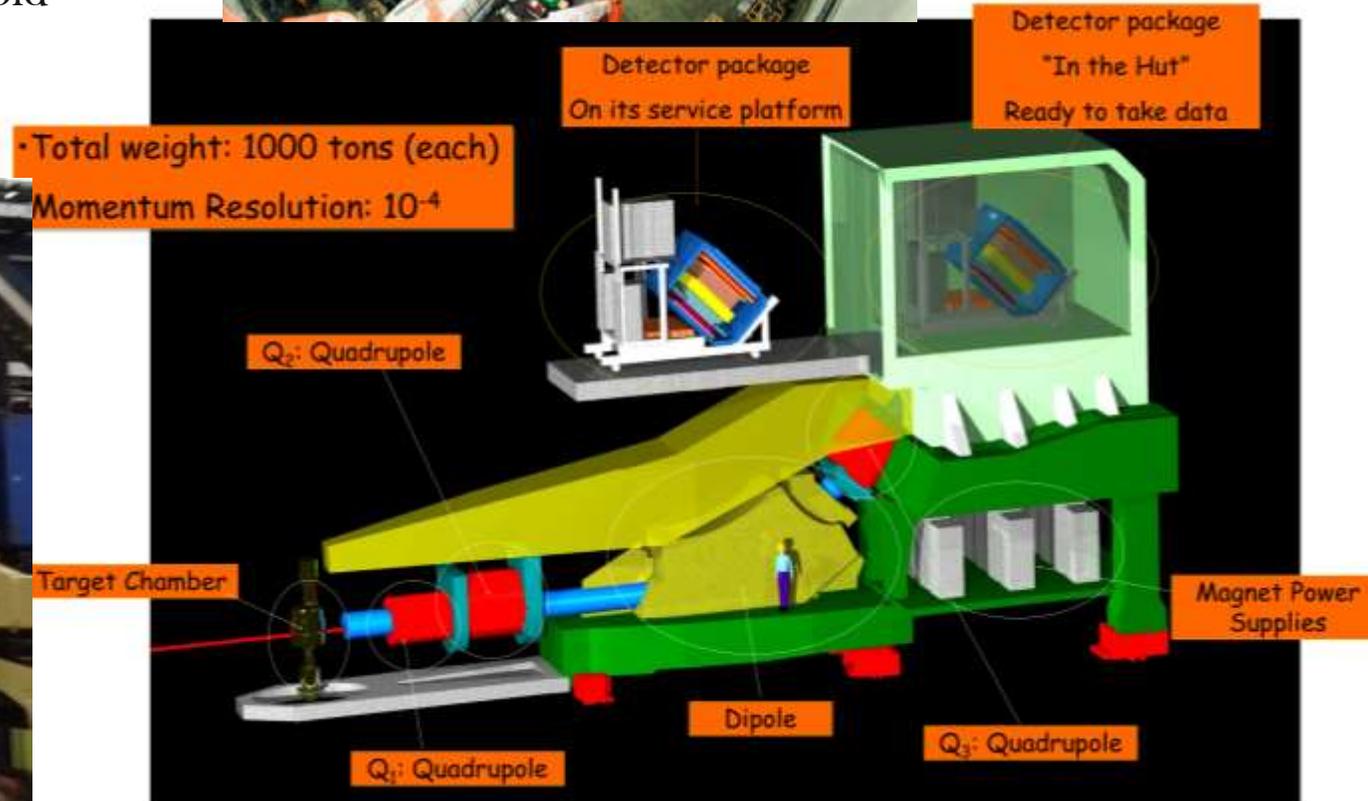
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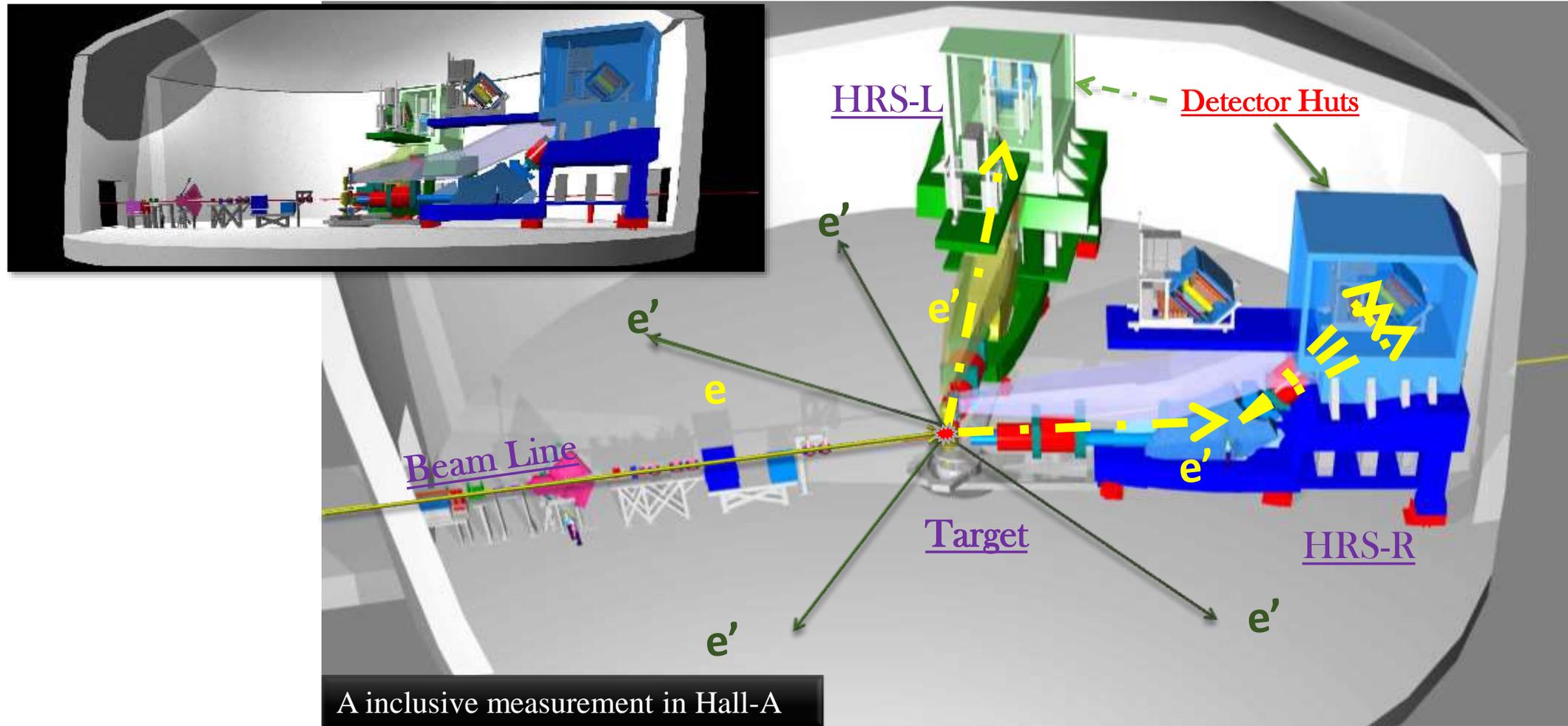
➤ Hall-A实验大厅

- ❑ Two identical high resolution spectrometers (HRS)
- ❑ Each HRS includes 3 Quadrupoles (Q) + 1 Dipole
QQDQ setup to obtain high momentum resolution (10^{-4})
- ❑ A ~ 20 meter flight path allowing to precise measure the angles and positions
- ❑ The detector-hut is ~ 10 meter above the floor to avoid radiation damage and background.



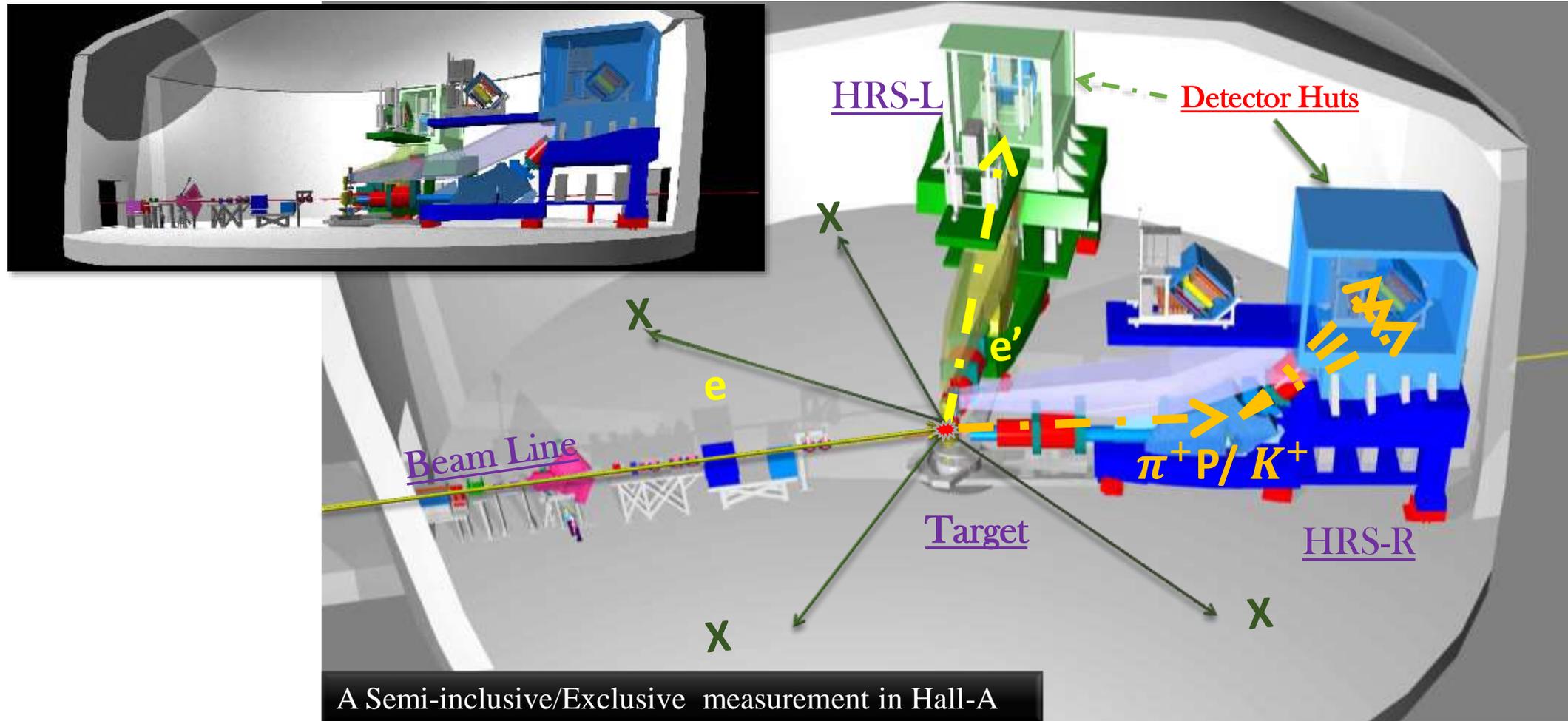
➤ High Resolution Magnetic Spectrometer (HRS)

- ❑ 极小接收角度 ($\pm 60\text{mrad}$)和极小接收动量 ($\pm 5\%$) 但极高分辨率
- ❑ 可以接收极高流量的反应 (高入射电流, 高靶厚度, 极高粒子计数率)
- ❑ 一般一组实验设置只能测量单个物理反应



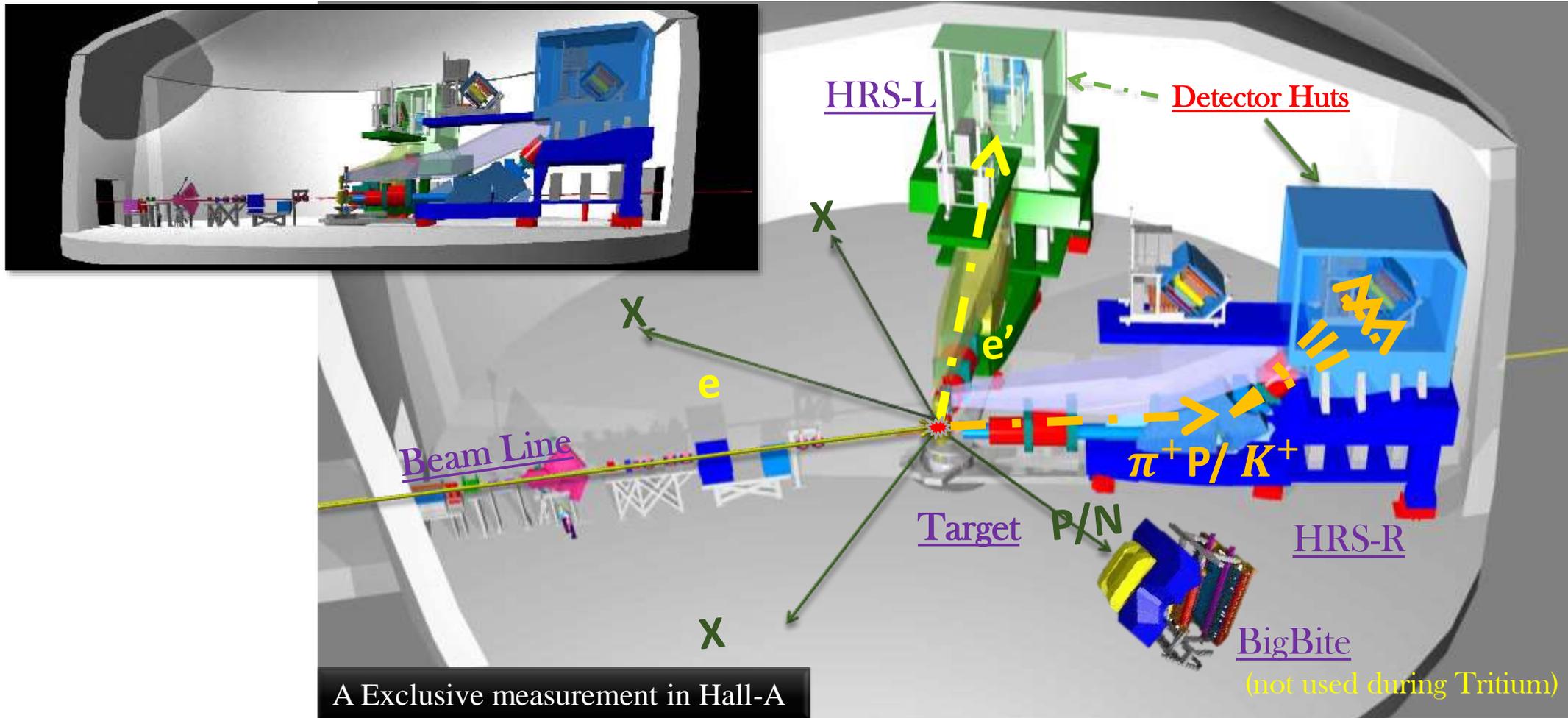
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- ❑ 一般一组实验设置只能测量单个物理反应



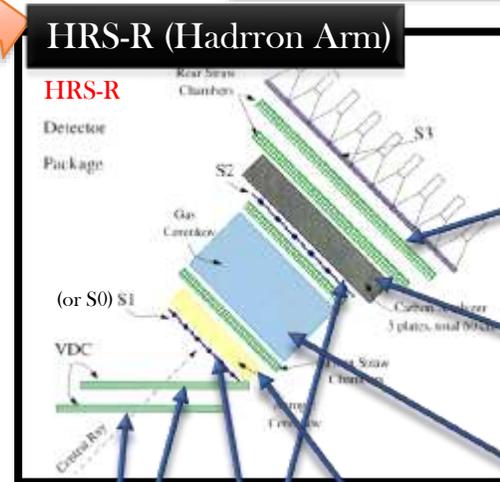
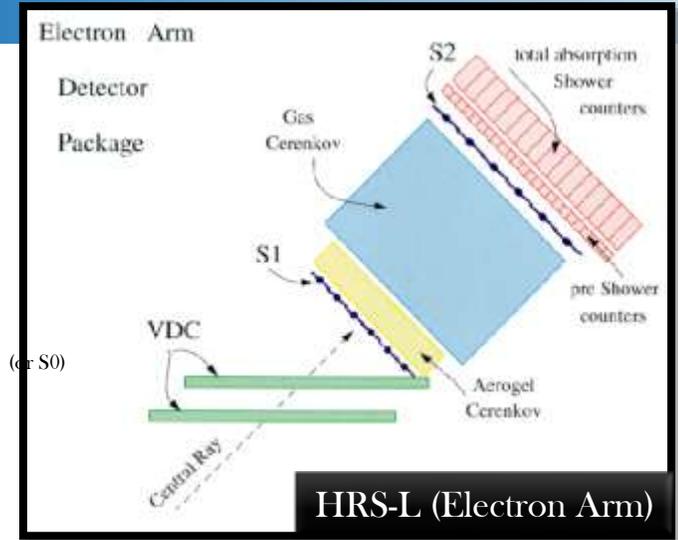
➤ High Resolution Magnetic Spectrometer (HRS)

- ❑ 极小接收角度 ($\pm 60\text{mrad}$)和极小接收动量 ($\pm 5\%$) 但极高分辨率
- ❑ 可以接收极高流量的反应 (高入射电流, 高靶厚度, 极高粒子计数率)
- ❑ 一般一组实验设置只能测量单个物理反应



➤ Hall-A实验大厅

The detector hut contains a set of detectors hold in a retractable rack, protected by a thick concrete door from radiations (e.g. neutrons).



Calorimeters

Polarimetry
(Proton
Polarization,
not in used)

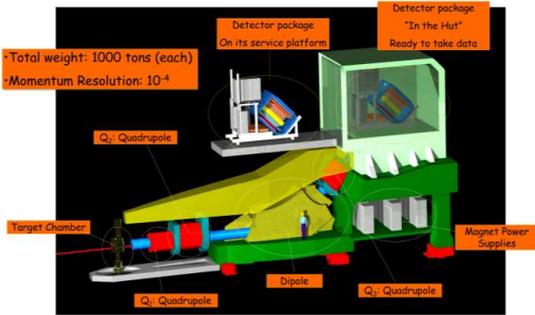
$e^{+/-}$ & $\pi^{+/-}$
Separation

Timing

$\pi^{+/-}$ and $K^{+/-}$ Separation

Front-End Electrons + Trigger System

Positions and Angles Tracking



➤ 同一探测器, 四组不同物理测量(2017~2018)

E12-10-103 (MARATHON)

- ❖ d/u Ratio at $x \rightarrow 1$
- ❖ Isospin-dependent EMC Effect



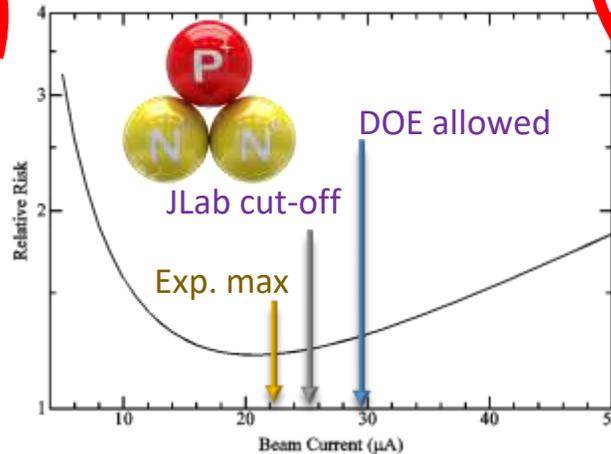
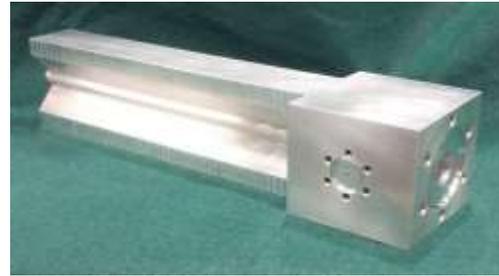
Spokespeople: G (Makis) Petratos, J. Gomez, A. Katramatou, R. Holt (J. Arrington), D. Meekins, R. Ransome,
 PhD Students: T. Hague, M. Mycz, T. Su (Kent), J. Bane (Tennessee), Tyler Kutz (Stony Brooks), H. Liu (Columbia)

E12-17-003 (Hypernucleus)

- ❖ Search for Lambda-N-N Hypernuclear



Spokespeople: F. Garibaldi, P. Markowitz, S. Nakamura, J. Reinhold, L. Tang, G. Urciuoli
 PhD Student: Bishnu Pendey (Hampton)



Tritium Specs:

- 1099 Ci
- 200 psia
- Cool with 40K
- Beam 22.5mA (~15W heat)

E12-11-112 (Inclusive SRC)

- ❖ Study Isospin Effect in 2N-SRC and 3N-SRC
- ❖ Measure GMn at small Q^2



Spokespeople: J. Arrington, D. Day, D. Higinbotham, P. Solgignon*, Z. Ye
 PhD Students: Shujie Li, Nathaly Santiesteban (UNH)

E12-14-009 (Exclusive SRC)

- ❖ Measure proton mom. dis. in ^3H & ^3He
- ❖ Verify in neutron-rich nuclei: $n_p(k) > n_n(k)$



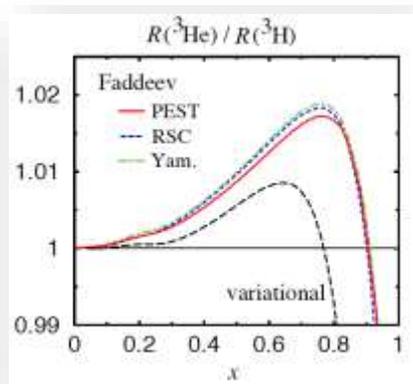
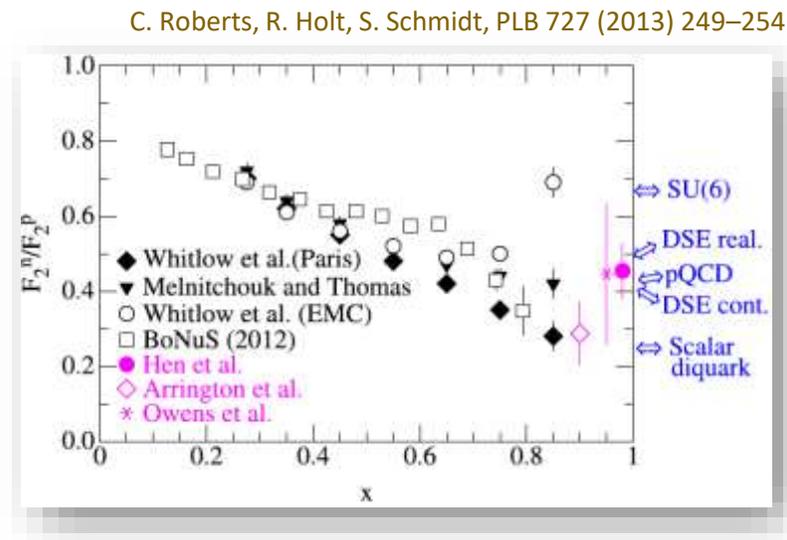
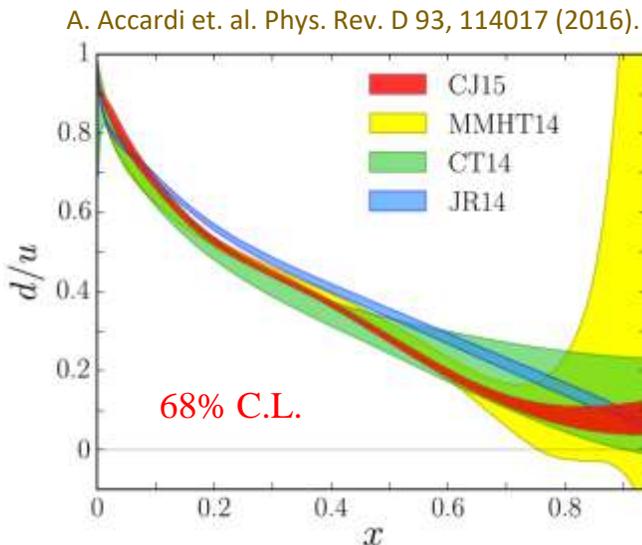
Spokespeople: L. Weinstein, W. Boeglin, F. Hauenstein, O. Hen, S. Gilad
 PhD Students: Reynier Cruz Torres (MIT)

➤ MARATHON Experiment

❖ 精确测量 $x \rightarrow 1$ 的 F_2^n/F_2^p 可以判断不同QCD模型对 d/u 的预测

$$F_2 = \sum_q e_q^2 f_1^q(x)$$

❖ 传统实验用 H_2 测 F_2^p , 用 D_2 作为有效中子测 $F_2^n \rightarrow$ 具有超大的核效应修正



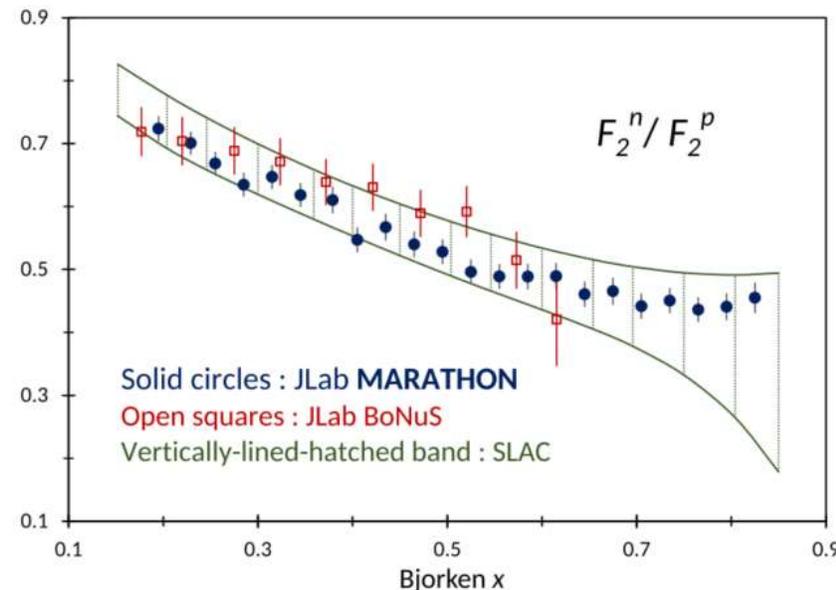
❖ 全新方法: 利用镜像原子核 H_3 和 He_3 (核效应几乎一样), 提取 F_2^n/F_2^p

$$\frac{F_2^n}{F_2^p} = \frac{2\mathcal{R} - F_2^{3He}/F_2^{3H}}{2F_2^{3He}/F_2^{3H} - \mathcal{R}}$$

Super-Ratio in EMC (DIS)

$$\mathcal{R} = \frac{R(3He)}{R(3H)}$$

$$R(3He) = \frac{F_{3He}}{2F_p + F_n}, \quad R(3H) = \frac{F_{3H}}{F_p + 2F_n}$$

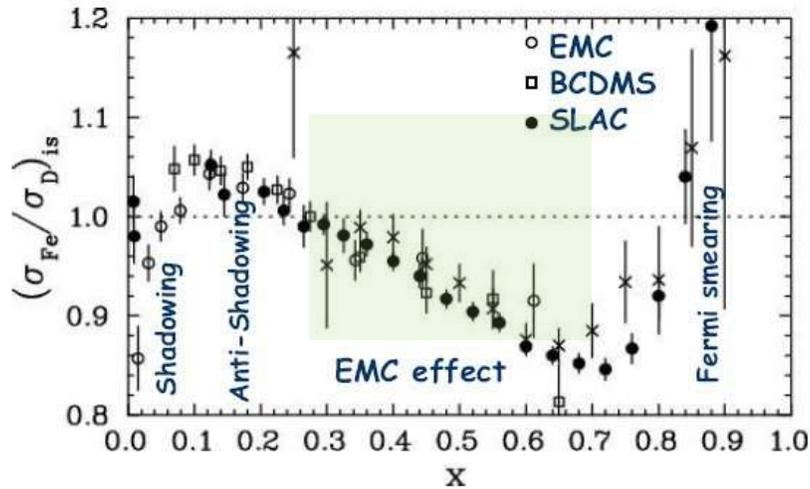


❖ 结果已发表

PRL128, 132003 (2022)

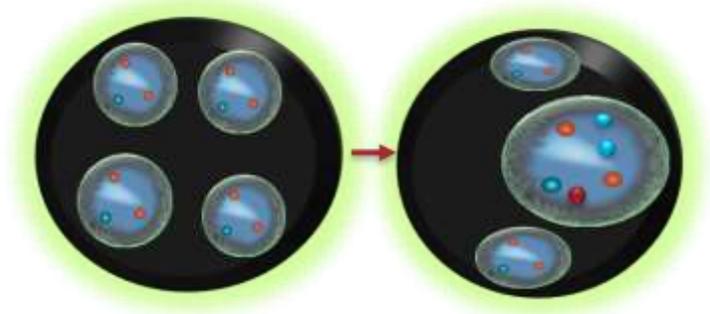
➤ MARATHON Experiment

❖ EMC Effect: Per-nucleon DIS cross-section ratio between a nucleus-A to the deuteron decreases linearly in $0.3 < x < 0.7$

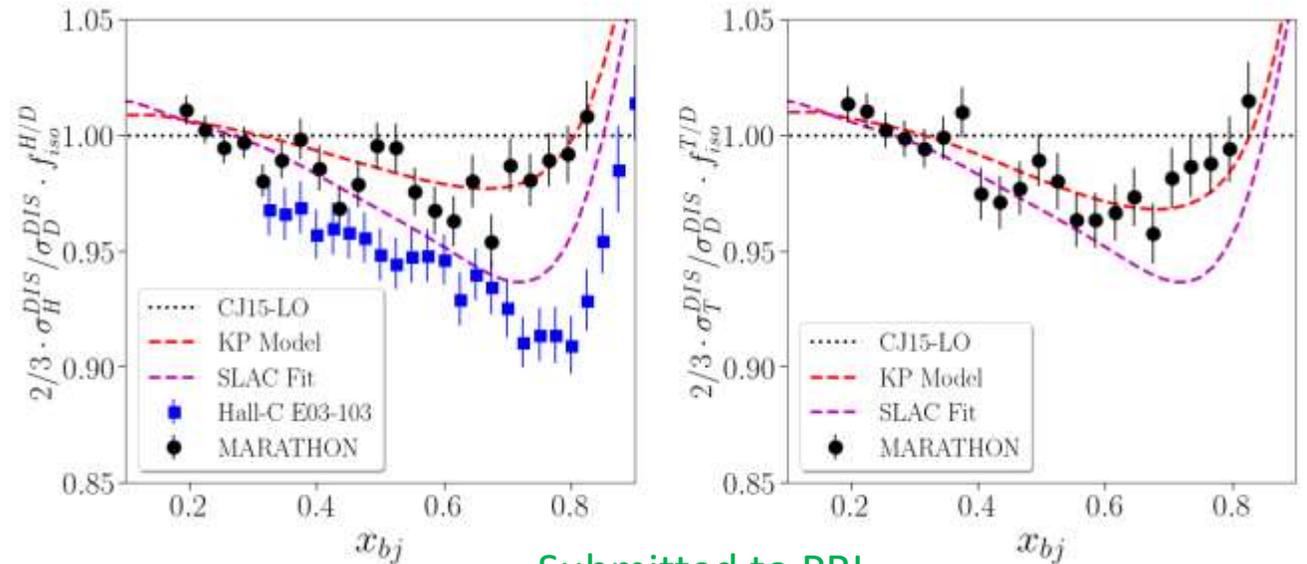


- No accepted explanation
- Nucleon must be modified
 - Which nucleons are modified?
 - Isospin Dependent?
 - Flavor dependent?

○ 自由的核子等于束缚的核子?



- Tritium/D2 and He3/D2 gave maximum isospin asymmetry in DIS

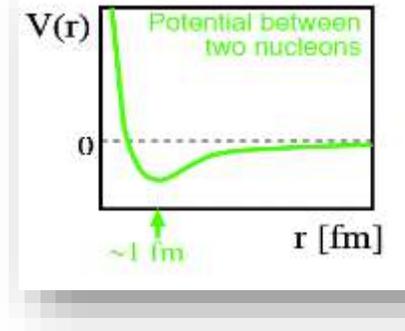


Submitted to PRL

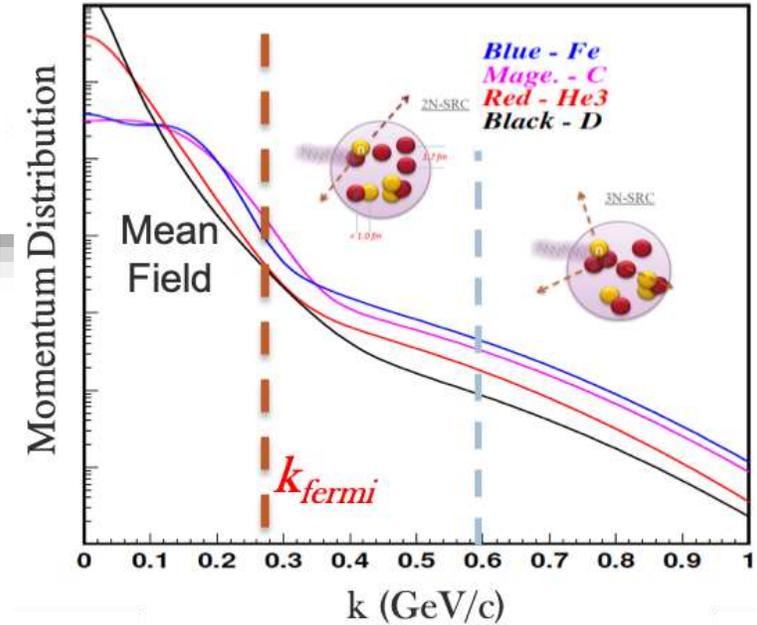
- But no flavor sensitive → Need new experiment !

➤ E12-11-112 (inclusive SRC):

- ❖ 核子之间80%的概率是长距“弱”作用 (平均场模型)
- ❖ 短程关联 (Short-range correlations, SRC):
 - ✓ 重核内2N-SRC (3N-SRC) 类似于²D (³He&³H) 原子核
 - ✓ 不同原子核内的核子动量分布具有相似尾巴
 - ✓ 高密度, 高动量 (研究中子星内部, 原子核的夸克结构等)

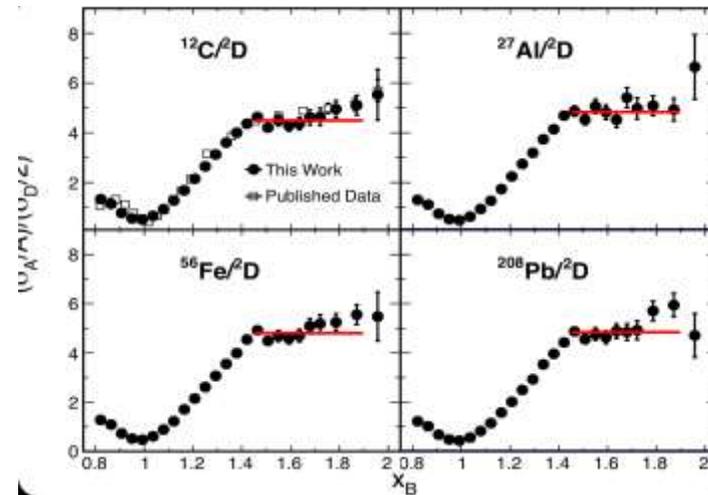
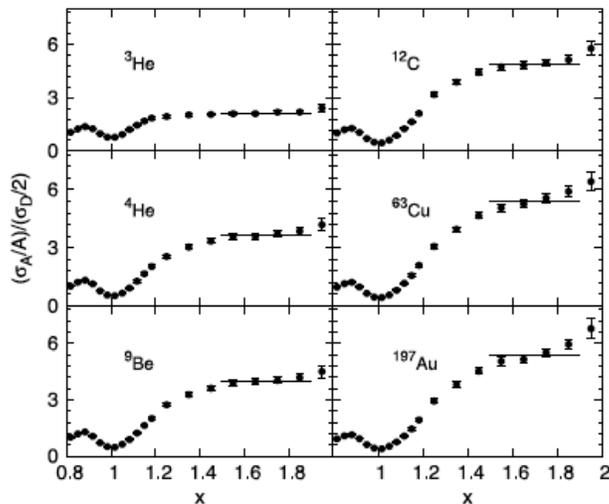


C. Ciofi degli Atti and S. Simula, Phys. Rev. C 53 (1996).



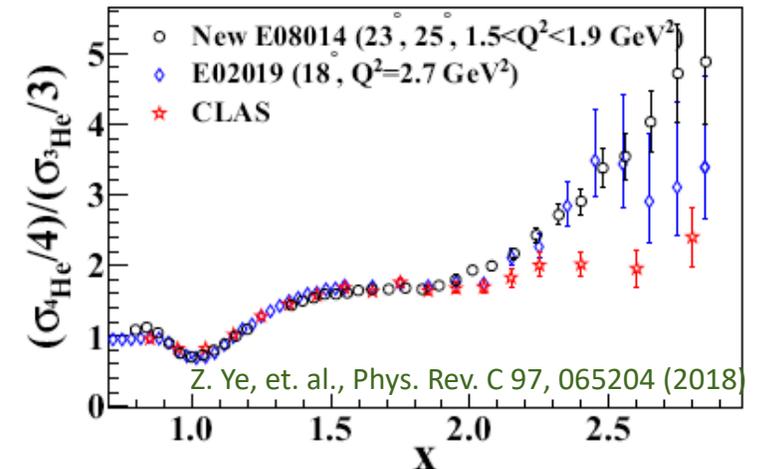
❖ 2N-SRC has been observed by inclusive and exclusive QE scattering

N. Fomin, E02-019 results, arXiv:0812:2144



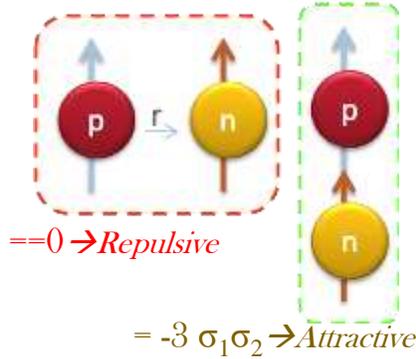
Schmookler et al., Nature (2019)

■ No signal about 3N-SRC!



➤ E12-11-112 (inclusive SRC):

❖ 2N-SRC is strongly Isospin-Dependent

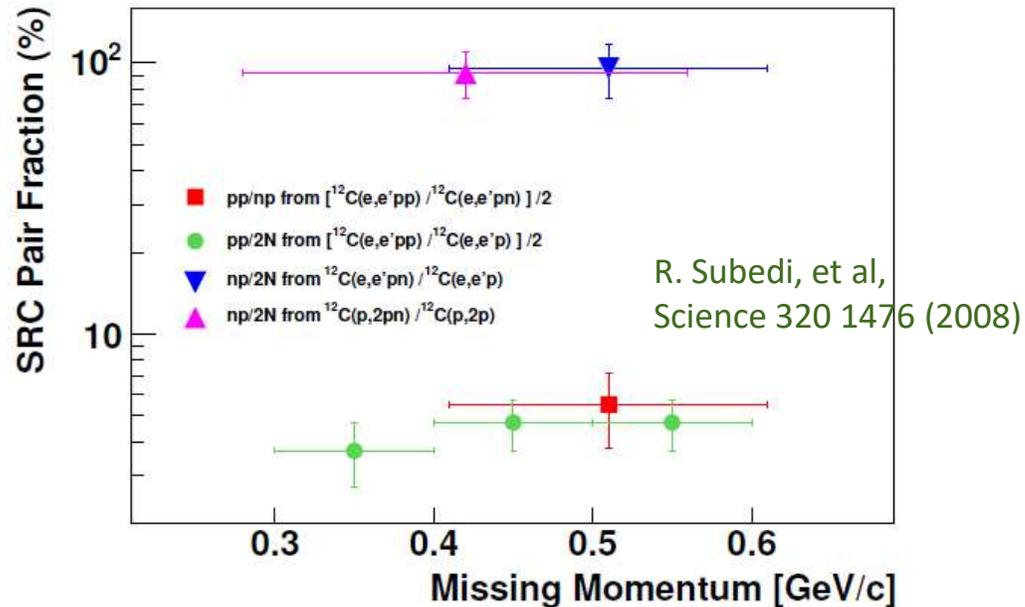


Spin component in V(r)

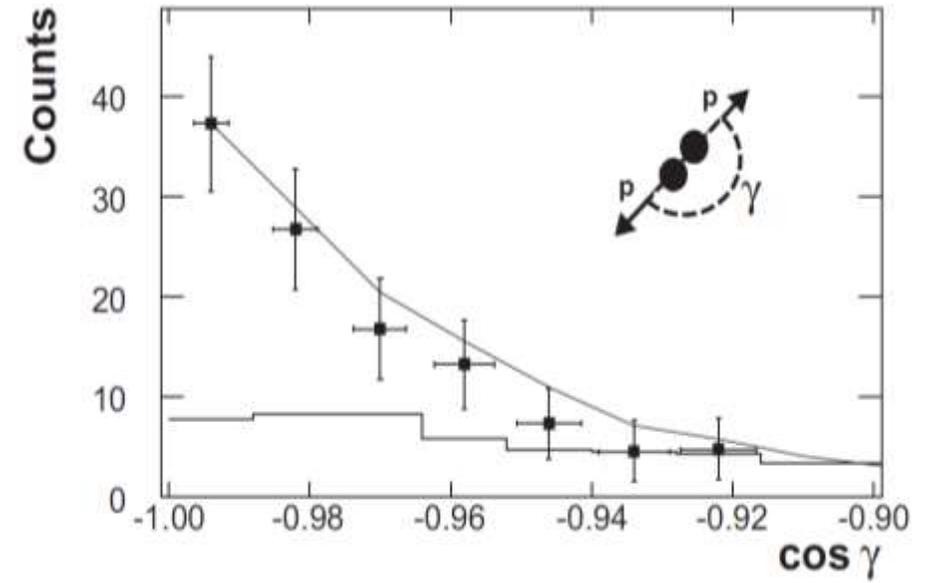
$$-S_{12} = -3(\vec{\sigma}_1 \cdot \hat{r})(\vec{\sigma}_2 \cdot \hat{r}) + (\vec{\sigma}_1 \cdot \vec{\sigma}_2)$$

At small r, tensor-force dominates:

■ Back-to-back in SRC w/ A(e,e'pN)A-2 reaction in Hall-A, JLab

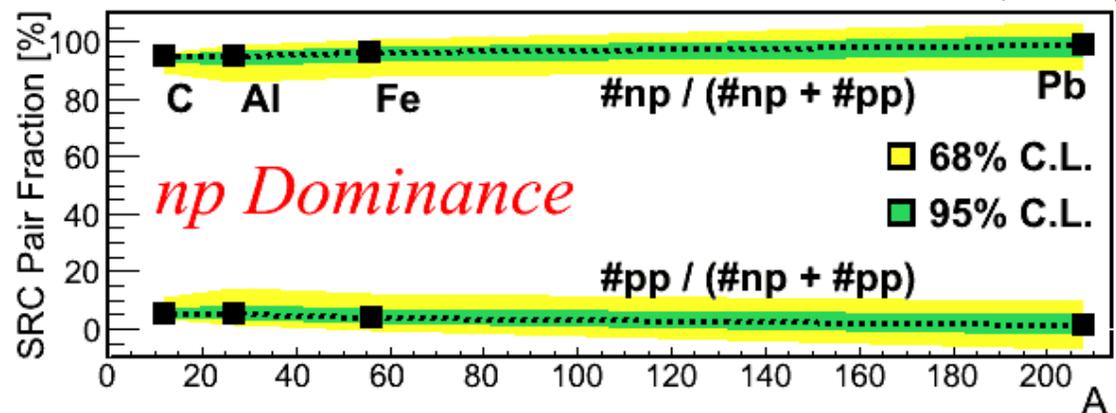


R. Shneor et al, PRL 99 072501 (2007)



■ np pairs are 90% in 2N-SRC!

O. Hen et al., Science 346, 614 (2014)



➤ E12-11-112 (inclusive SRC):

- ❖ Isospin-Dependence was seen in (e,e'p) with big errors
→ but not fully conformed by (e,e')
- ❖ Use Inclusive QE scattering on He3/H3 to study isospin dependence
✓ Much clearer than Ca48/Ca40 ratio

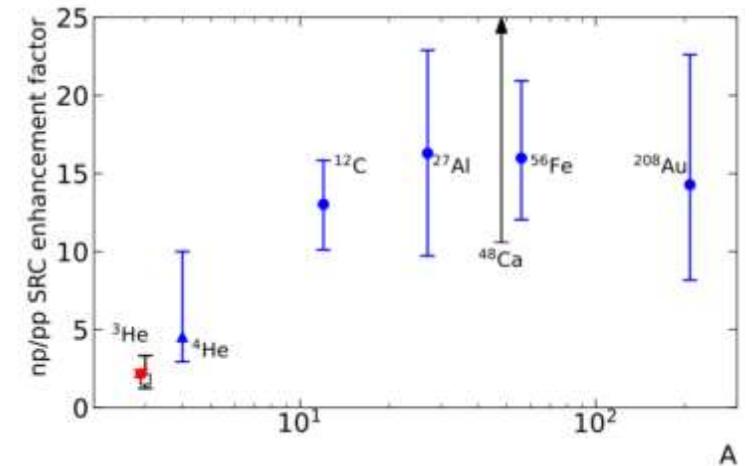
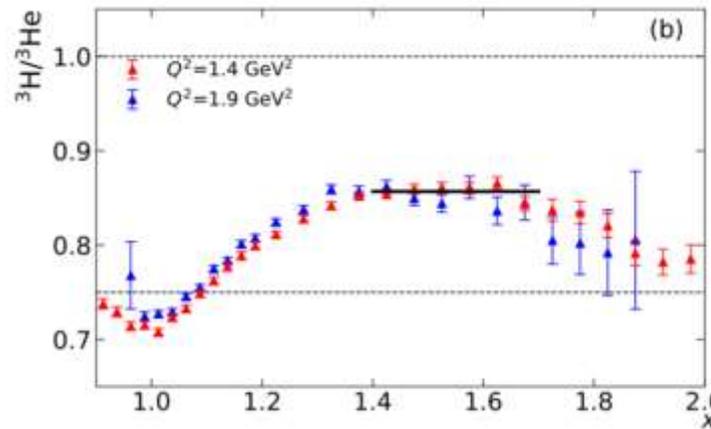
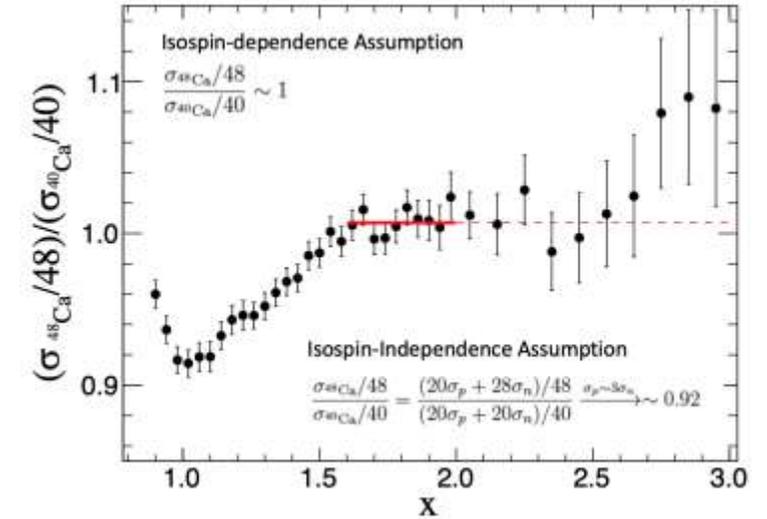
$$R = \frac{\sigma_{H3/3}}{\sigma_{He3/3}} = \frac{(\sigma_p + 1\sigma_n)/3}{(2\sigma_p + \sigma_n)/3} \xrightarrow{\sigma_p \approx 3\sigma_n} 0.71$$

$$R = \frac{\sigma_{H3/3}}{\sigma_{He3/3}} = \frac{(2\sigma_{np} + \sigma_{nn})/3}{(2\sigma_{np} + \sigma_{pp})/3} \xrightarrow{T=0, np-domin} 1.0$$

- ◆ Very high precision in (e,e') with H3&He3 to study isospin-SRC
- ◆ Less np dominated in light nuclei vs a more universal feature in heavy nuclei

S. Li, et. al, accepted by Nature (to be published on Sep. 1st)

D. Nguyen, Z. Ye, et al, RPC **102**, 064004 (2020)

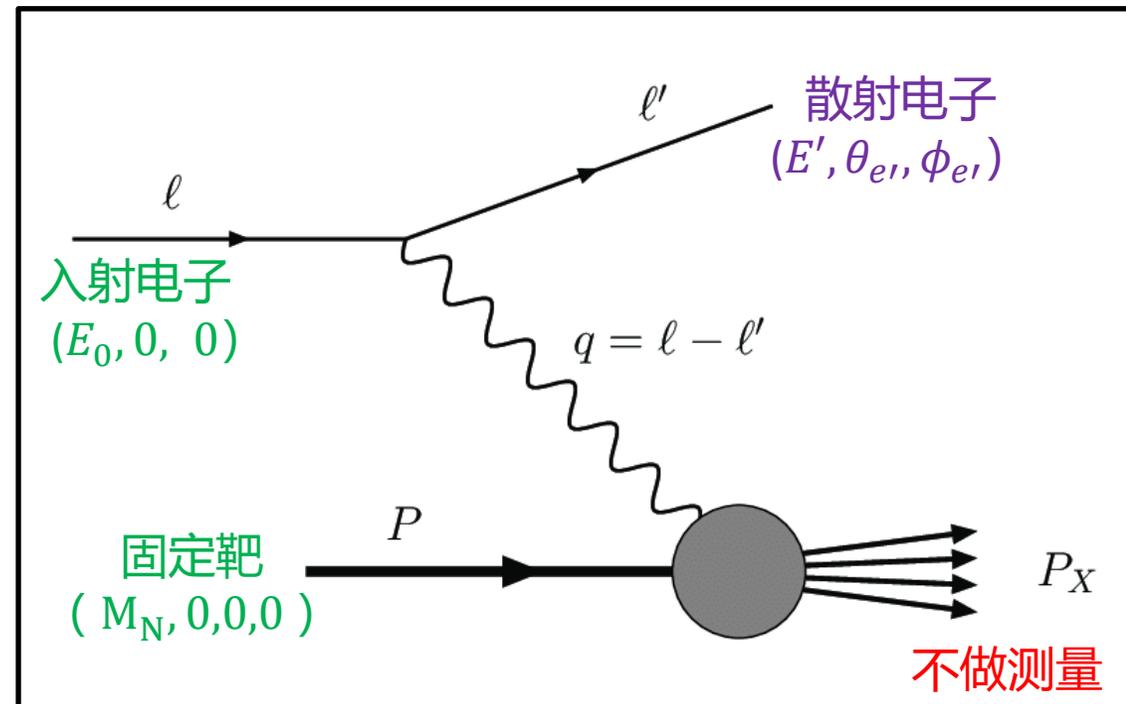


➤ Unpolarized Inclusive Cross-Section Measurements:

❑ 实验测量单举微分反应截面

$$\frac{d\sigma_{EX}^{raw}}{dE' d\Omega}(E_0, E'_i, \theta_0) = \frac{N_{EX}^i \cdot \epsilon_{e=\pi}}{N_e \cdot \eta_{tg} \cdot \epsilon_{eff} \cdot (\Delta E'_{EX} \Delta \Omega_{EX})}$$

探测到的散射电子数 π介子掺杂修正
入射电子数 (电流*时间) 靶内核子数 (密度*长度) 探测器效率修正 电子微分空间 (动量*角度)



❑ 动量参数

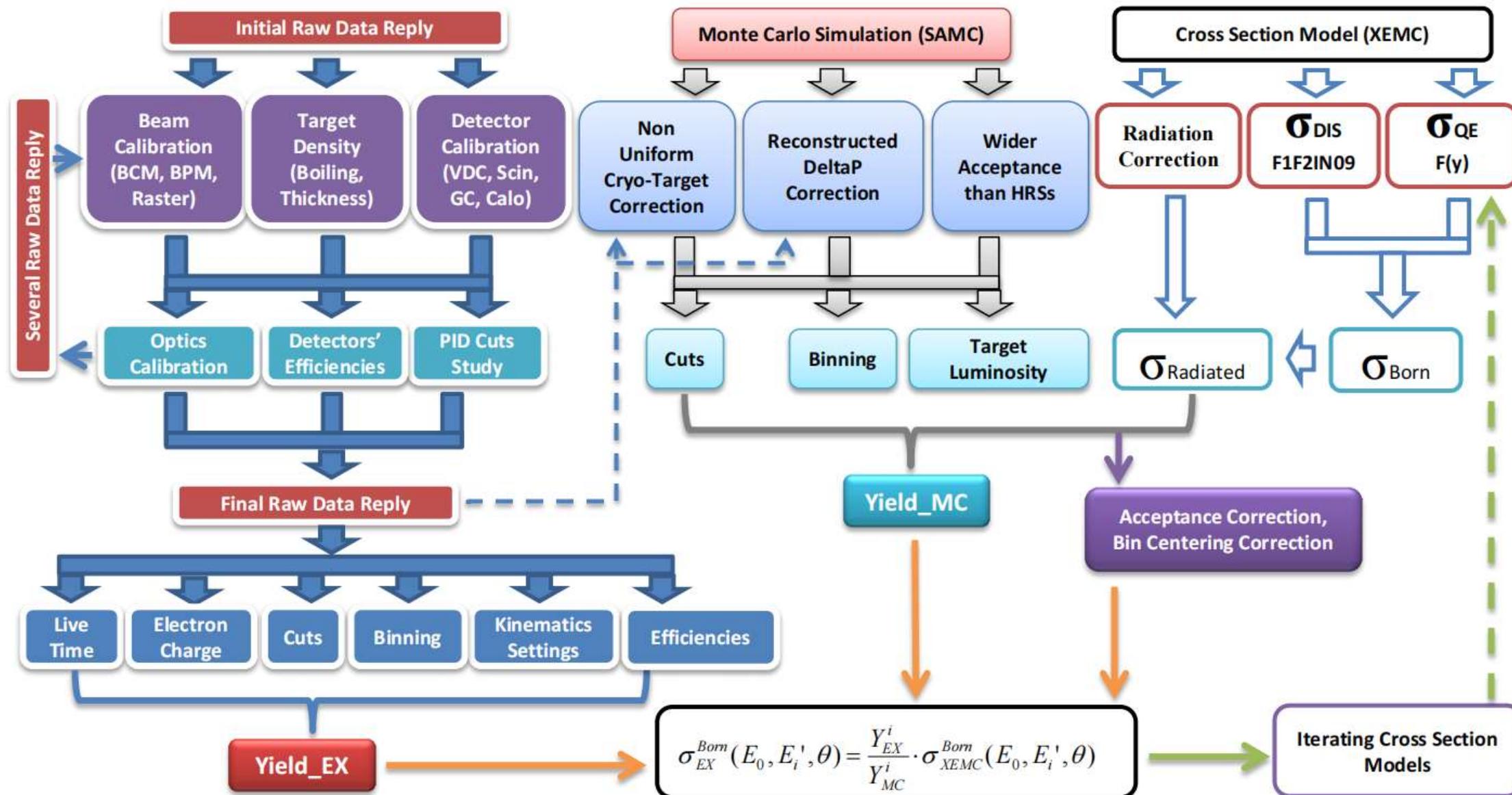
$$Q^2 = 4E_0 E' \sin^2(\theta/2) \quad \nu = E_0 - E' \quad x_{bj} = \frac{Q^2}{2m_p \nu}$$

❑ 从观测量到物理量:

$$d\sigma_{EX}^{raw}(E_0, E'_{EX}, \theta_{EX}) \xrightarrow{\text{theory corrections}} d\sigma_{EX}^{born}(E_0, E'_{Real}, \theta_{real}) \xrightarrow{\text{Jacobian Trans.}} d\sigma_{EX}^{born}(Q^2, x)$$

$$\xrightarrow{\text{Leading-Order}} \frac{2\pi\alpha^2}{(xS^2)} \frac{1-y + \frac{1}{2}y^2}{y^2} \sum_q e_q^2 f_1^q(x)$$

实验分析流程:

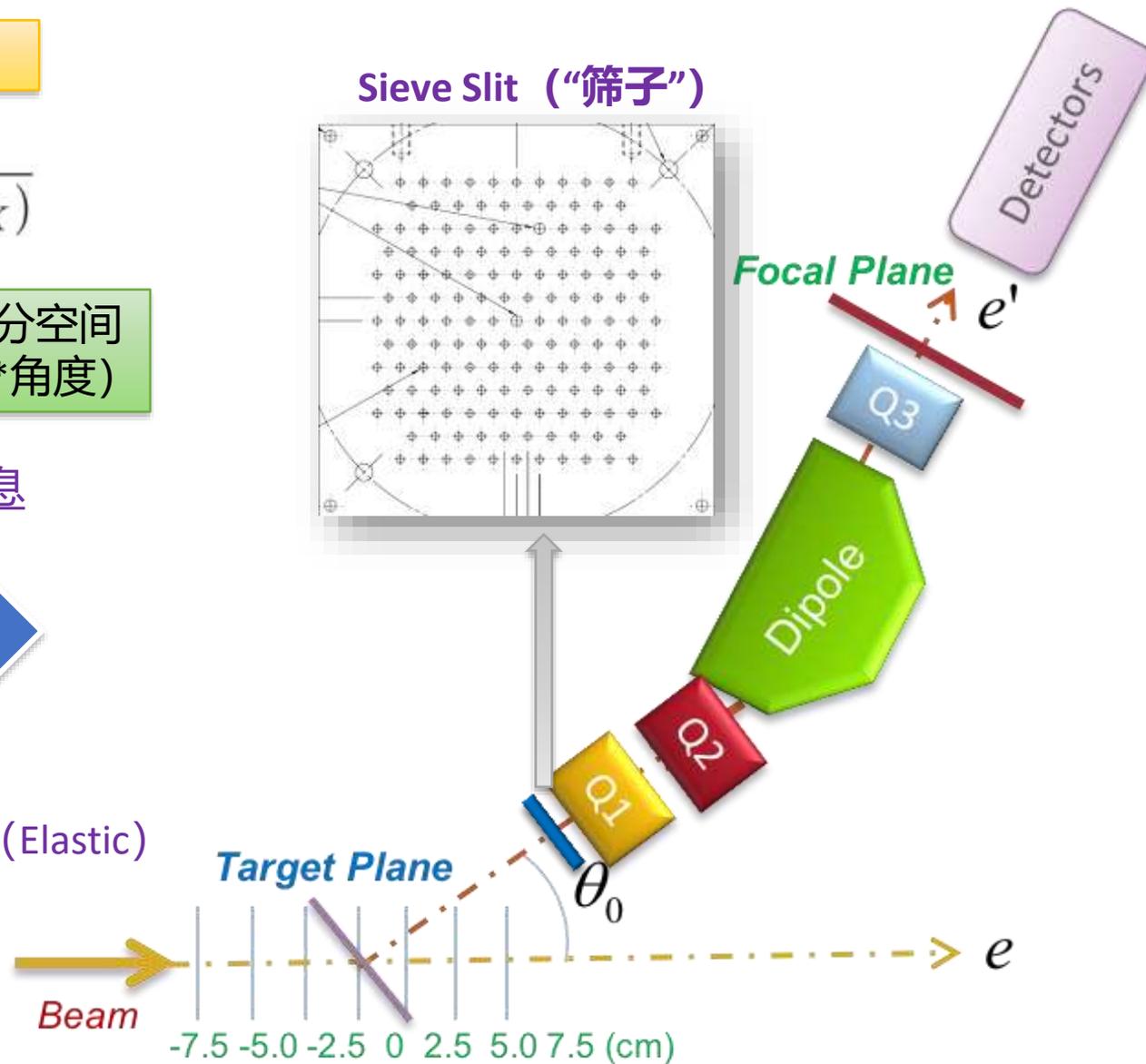
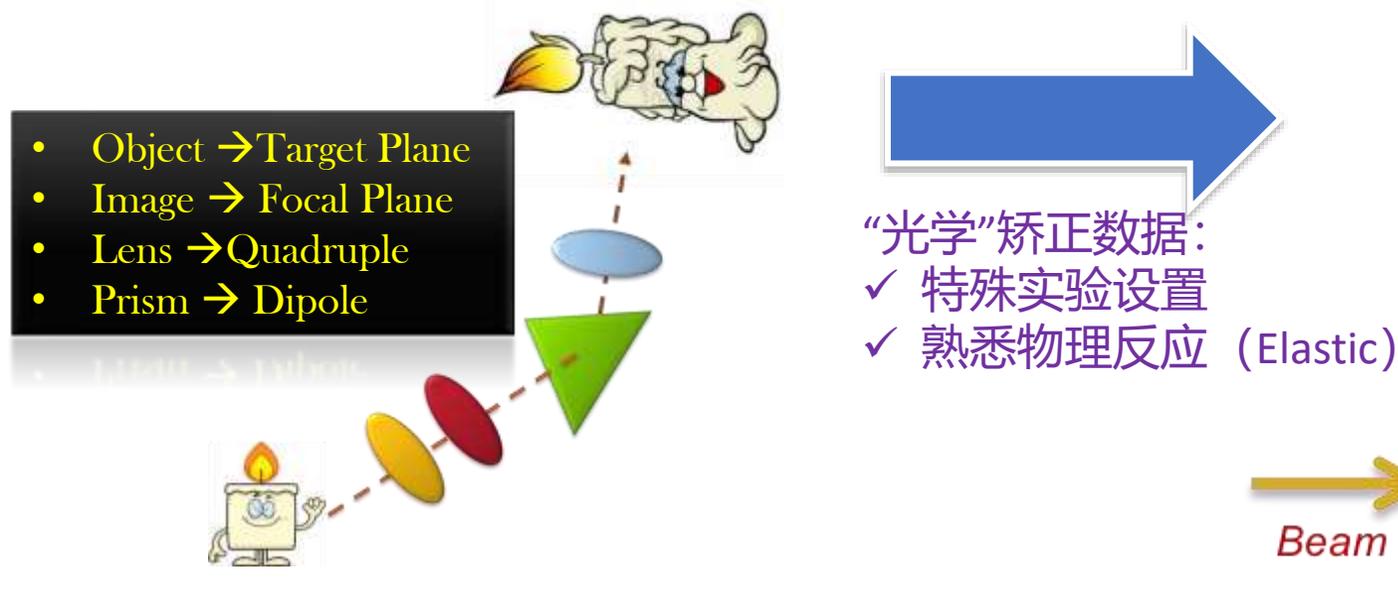


➤ Unpolarized Inclusive Cross-Section Measurements:

$$\frac{d\sigma_{EX}^{raw}}{dE'd\Omega}(E_0, E'_i, \theta_0) = \frac{N_{EX}^i \cdot \epsilon_{e=\pi}}{N_e \cdot \eta_{tg} \cdot \epsilon_{eff} \cdot (\Delta E'_{EX} \Delta \Omega_{EX})}$$

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“Optics” Reconstruction → 重建反应平面的动量信息



➤ Unpolarized Inclusive Cross-Section Measurements:

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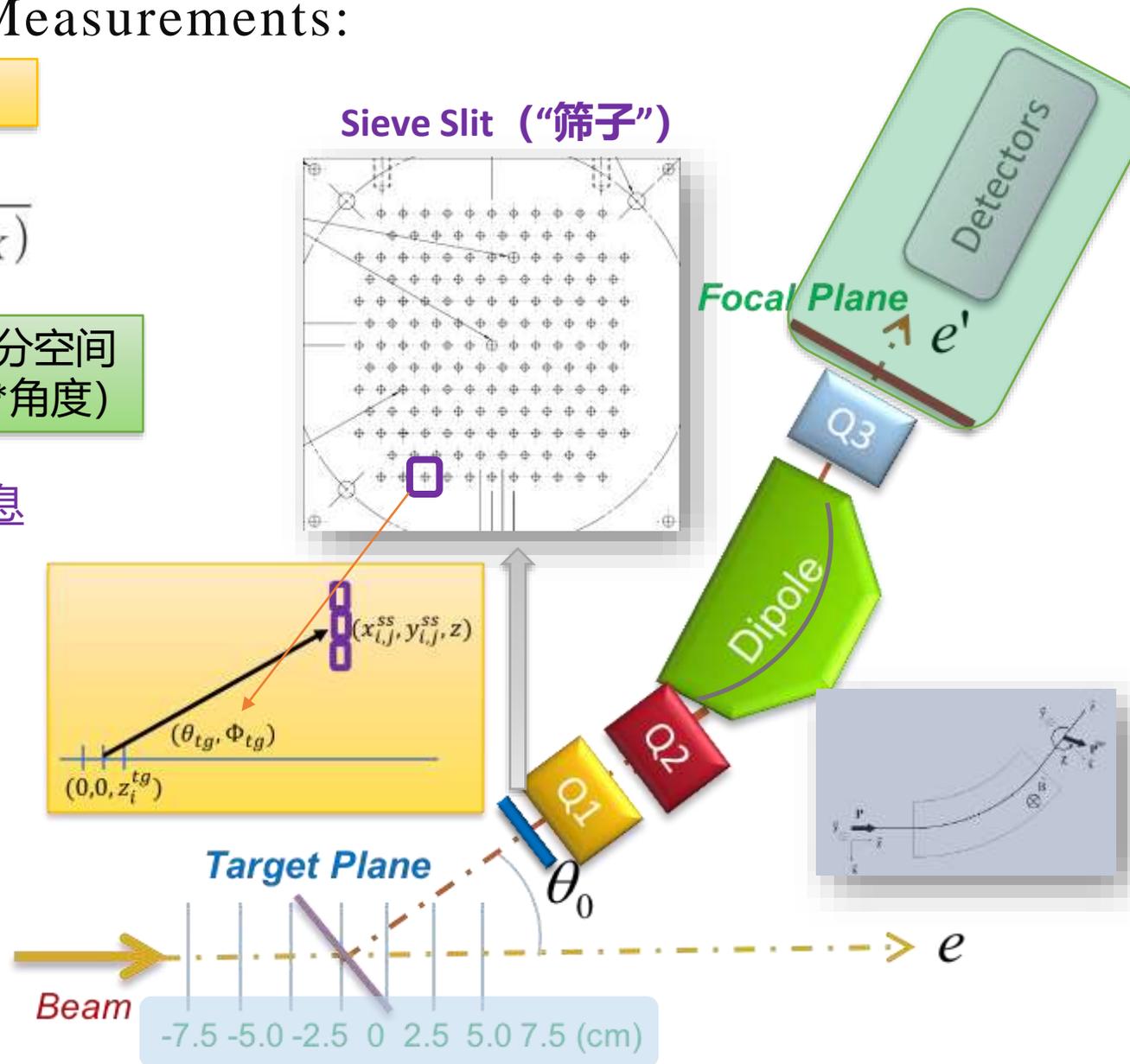
“Optics” Reconstruction → 重建反应平面的动量信息

$$\begin{pmatrix} \delta p \\ y_{tg} \\ \theta_{tg} \\ \phi_{tg} \end{pmatrix} = \begin{bmatrix} Optics \\ Matrix \end{bmatrix} \begin{pmatrix} x_{fp} \\ y_{fp} \\ \theta_{fp} \\ \phi_{fp} \end{pmatrix}$$

$$\begin{cases} \delta p = \sum_{i,j,k,l} C_{ijk}^D x_{fp}^i y_{fp}^j \theta_{fp}^k \phi_{fp}^l \\ y_{tg} = \sum_{i,j,k,l} C_{ijk}^Y x_{fp}^i y_{fp}^j \theta_{fp}^k \phi_{fp}^l \\ \theta_{tg} = \sum_{i,j,k,l} C_{ijk}^T x_{fp}^i y_{fp}^j \theta_{fp}^k \phi_{fp}^l \\ \phi_{tg} = \sum_{i,j,k,l} C_{ijk}^P x_{fp}^i y_{fp}^j \theta_{fp}^k \phi_{fp}^l \end{cases}$$

□ 光学系数通过已知实验数据拟合

Measured in Focal Plane



➤ Unpolarized Inclusive Cross-Section Measurements:

$$\frac{d\sigma_{EX}^{raw}}{dE'd\Omega}(E_0, E'_i, \theta_0) = \frac{N_{EX}^i \cdot \epsilon_{e=\pi}}{N_e \cdot \eta_{tg} \cdot \epsilon_{eff} \cdot (\Delta E'_{EX} \Delta \Omega_{EX})}$$

探测到的散射电子数 → N_{EX}^i
π介子掺杂修正 → $\epsilon_{e=\pi}$

入射电子数 (电流*时间) → N_e
靶内核子数 (密度*长度) → η_{tg}
探测器效率修正 → ϵ_{eff}
电子微分空间 (动量*角度) → $(\Delta E'_{EX} \Delta \Omega_{EX})$

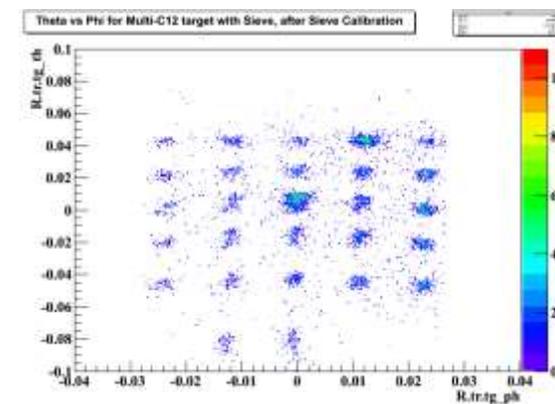
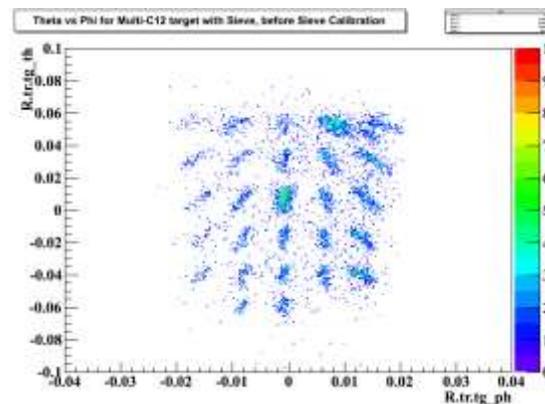
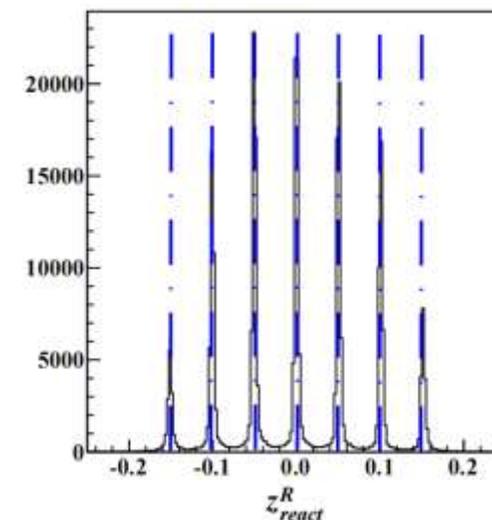
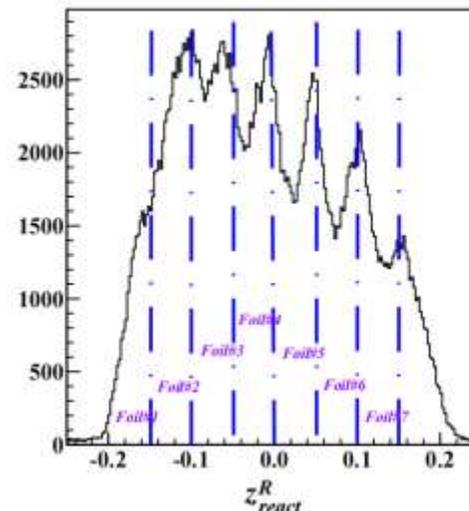
“Optics” Reconstruction → 重建反应平面的动量信息

$$\begin{pmatrix} \delta p \\ y_{tg} \\ \theta_{tg} \\ \phi_{tg} \end{pmatrix} = \begin{bmatrix} \text{Optics} \\ \text{Matrix} \end{bmatrix} \begin{pmatrix} x_{fp} \\ y_{fp} \\ \theta_{fp} \\ \phi_{fp} \end{pmatrix}$$

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Measured in Focal Plane



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$$\frac{d\sigma_{EX}^{raw}}{dE'd\Omega}(E_0, E'_i, \theta_0) = \frac{N_{EX}^i \cdot \epsilon_{e=\pi}}{N_e \cdot \eta_{tg} \cdot \epsilon_{eff} \cdot (\Delta E'_{EX} \Delta \Omega_{EX})}$$

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□ 入射电子数:

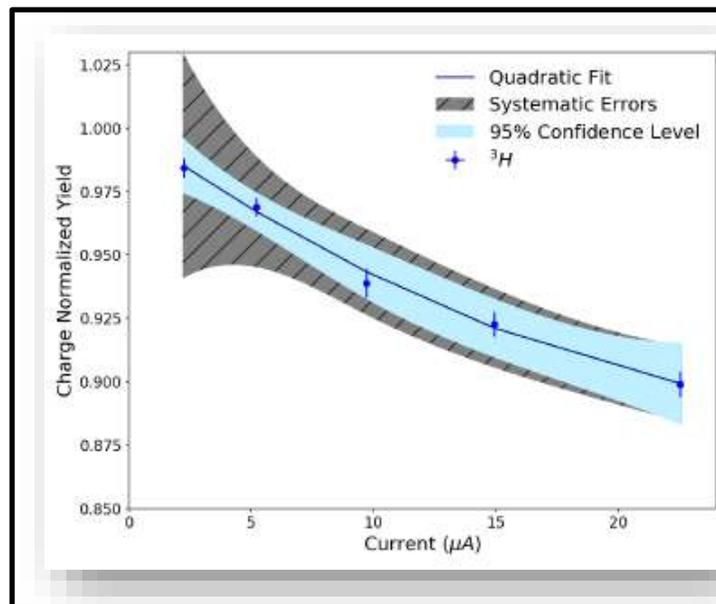
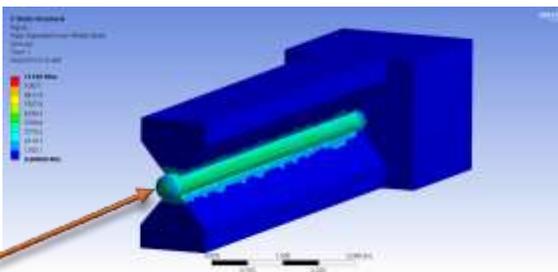
- ✓ 实验向加速器请求特定电流
- ✓ 实验大厅实时测量真实电流涨落
- ✓ 记录电子束流打靶的时间

$$N_e = \sum_i I_i \cdot \Delta T_i / e$$

□ 靶核子数: $\eta_{tg} = \frac{\rho \cdot l \cdot N_a}{A}$,

- ✓ 靶的长度为实验已知量
- ✓ 固体靶的密度为已知量
- ✓ 液体和气体的密度随温度涨落 (电子束流加热作用)

Electrons



Target Boiling Correction

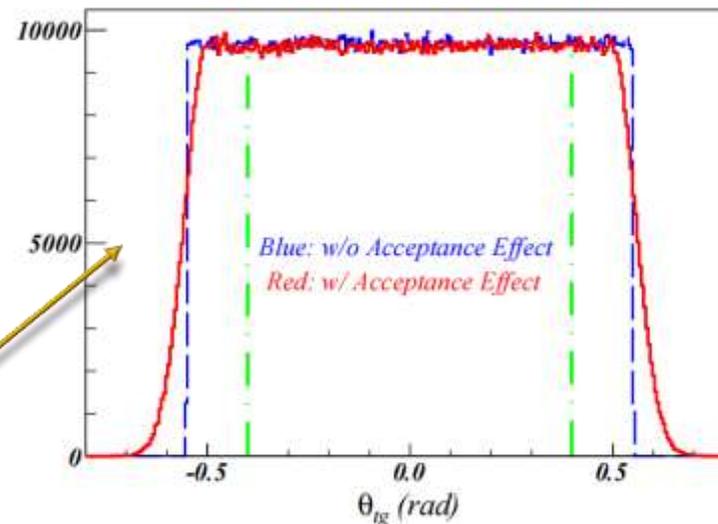
- ✓ 向加速器要求多组电流 (2.5uA, 5uA, 10uA, 15uA, 22.5uA)
- ✓ 测量电子轰击液体/气体靶的散射电子数
- ✓ 同时测量电子轰击固体靶作为对照组
- ✓ 分析散射电子数的变化 (修正电流后)
- ✓ 拟合修正函数

$$\rho = \rho_0 \cdot (1.0 - B \cdot I/100)$$

➤ Unpolarized Inclusive Cross-Section Measurements:

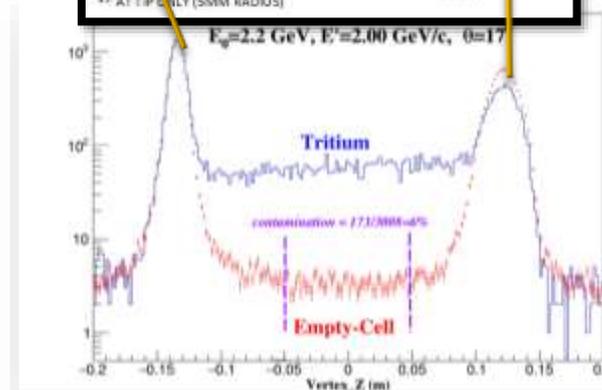
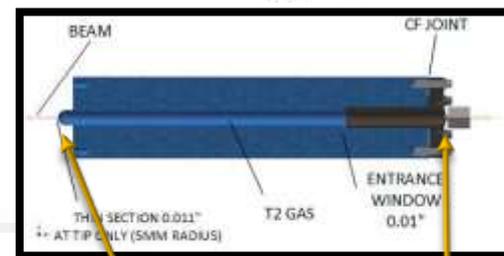
$$\frac{d\sigma_{EX}^{raw}}{dE'd\Omega}(E_0, E'_i, \theta_0) = \frac{N_{EX}^i \cdot \epsilon_{e=\pi}}{N_e \cdot \eta_{tg} \cdot \epsilon_{eff} \cdot (\Delta E'_{EX} \Delta \Omega_{EX})}$$

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❑ 粒子甄别 (Particle Identification) :

- ✓ 来自碰撞点的散射电子
- ✗ 物理信息被改变的散射电子 (如能损, 次级碰撞等)
- ✗ 电子束流被靶外其他材料散射 (气体容器, 空气, 加速器管道材料等) → 取中间一半的气体

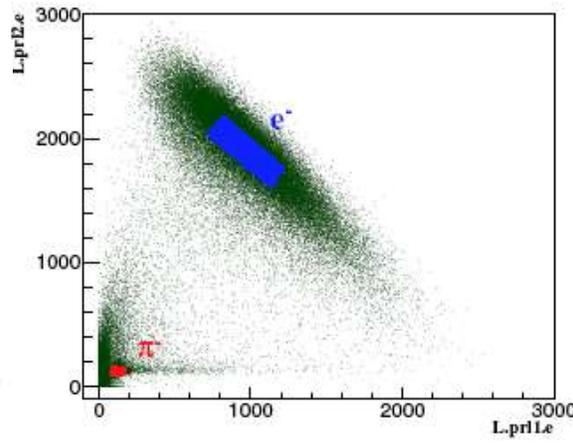
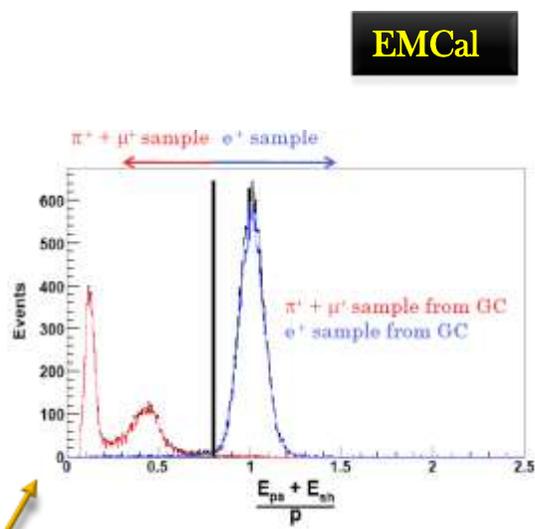


➤ Unpolarized Inclusive Cross-Section Measurements:

$$\frac{d\sigma_{EX}^{raw}}{dE'd\Omega}(E_0, E'_i, \theta_0) = \frac{N_{EX}^i \cdot \epsilon_{e=\pi}}{N_e \cdot \eta_{tg} \cdot \epsilon_{eff} \cdot (\Delta E'_{EX} \Delta \Omega_{EX})}$$

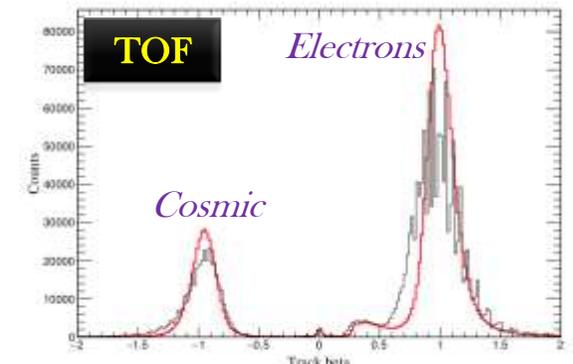
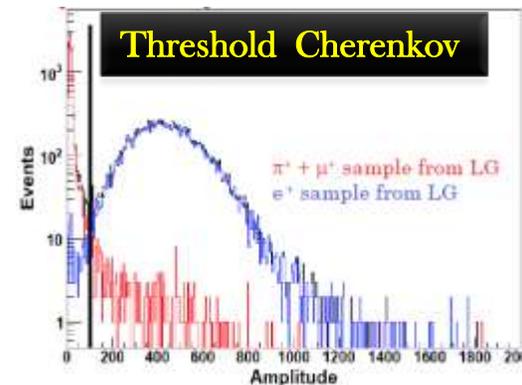
探测到的散射电子数 (绿色框) π 介子掺杂修正 (黄色框)

入射电子数 (电流*时间) (黄色框) 靶内核子数 (密度*长度) (黄色框) 探测器效率修正 (黄色框) 电子微分空间 (动量*角度) (黄色框)



❑ 粒子甄别 (Particle Identification) :

- ✓ 来自碰撞点的散射电子
- ✗ 物理信息被改变的散射电子 (如能损, 次级碰撞等)
- ✗ 电子束流被靶外其他材料散射 (气体容器, 空气, 加速器管道材料等) → 取中间一半的气体
- ✗ 来自高能光子衰败的正负电子对 (极少, 实验验证)
- ✗ 误判为散射电子的其他粒子, π/K 介子, 宇宙射线等

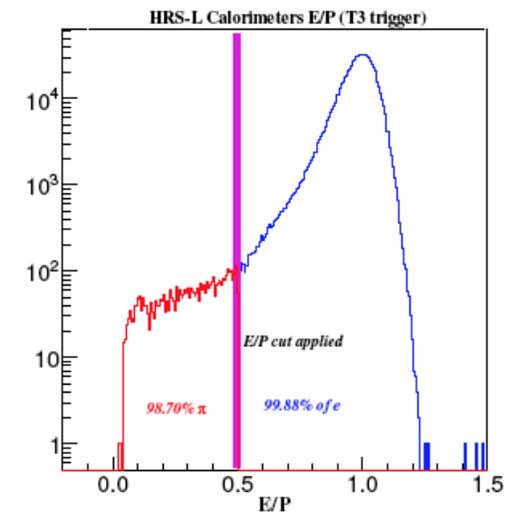
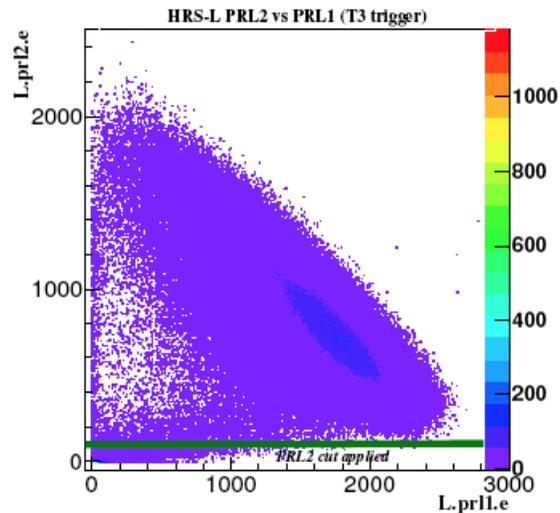


➤ Unpolarized Inclusive Cross-Section Measurements:

$$\frac{d\sigma_{EX}^{raw}}{dE'd\Omega}(E_0, E'_i, \theta_0) = \frac{N_{EX}^i \cdot \epsilon_{e=\pi}}{N_e \cdot \eta_{tg} \cdot \epsilon_{eff} \cdot (\Delta E'_{EX} \Delta \Omega_{EX})}$$

探测到的散射电子数 → N_{EX}^i
 π 介子掺杂修正 → $\epsilon_{e=\pi}$

入射电子数 (电流*时间) → N_e
靶内核子数 (密度*长度) → η_{tg}
探测器效率修正 → ϵ_{eff}
电子微分空间 (动量*角度) → $(\Delta E'_{EX} \Delta \Omega_{EX})$



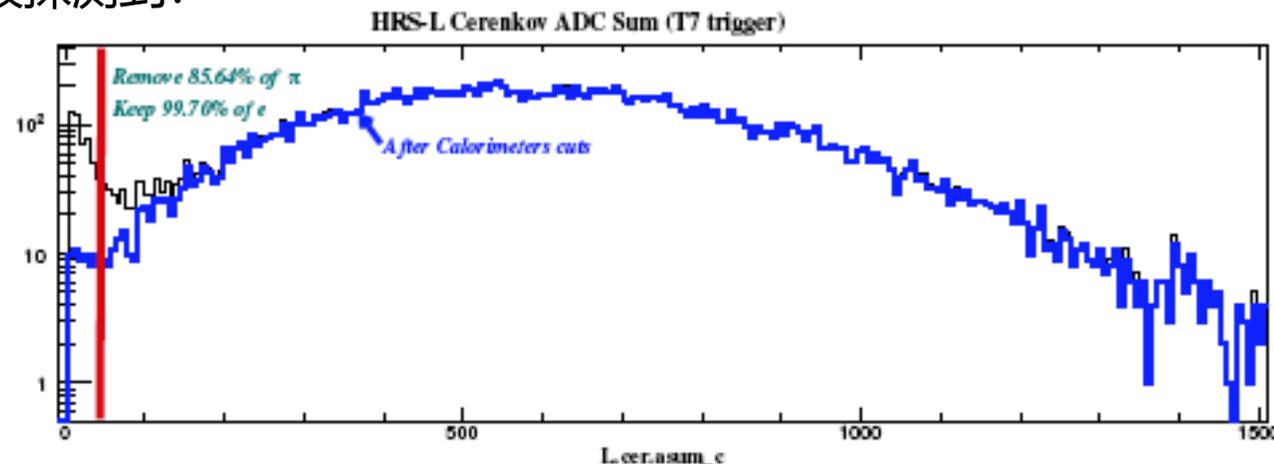
❑ 探测器效率 → 穿过探测器的粒子 (电子) 多少比例被探测到?

受以下硬件问题影响

- ✓ 粒子达到探测器的死区
- ✓ 没留下足够大信号
- ✓ 无法进行有效轨迹重建
- ✓ 无法产生有效触发信号
- ✓ 前端电子学继续数据过慢
- ✓ 电脑存储数据过慢

❑ PID效率 (如介子掺杂修正)

→ 数据分析过程中人为判断的准确度



❑ 通过两组探测器探测同一种粒子, 相互判断

(如: N个被探测器#1测到的电子, 多少被探测器#2同时看到?)

➤ Unpolarized Inclusive Cross-Section Measurements:

$$\frac{d\sigma_{EX}^{raw}}{dE'd\Omega}(E_0, E'_i, \theta_0) = \frac{N_{EX}^i \cdot \epsilon_{e=\pi}}{N_e \cdot \eta_{tg} \cdot \epsilon_{eff} \cdot (\Delta E'_{EX} \Delta \Omega_{EX})}$$

探测到的散射电子数

π 介子掺杂修正

入射电子数 (电流*时间)

靶内核子数 (密度*长度)

探测器效率修正

电子微分空间 (动量*角度)

□ 其他修正

- ✓ Energy Lose
- ✓ Radiative Correction
- ✓ Bing-Centering Correction
- ✓ Acceptance Correction

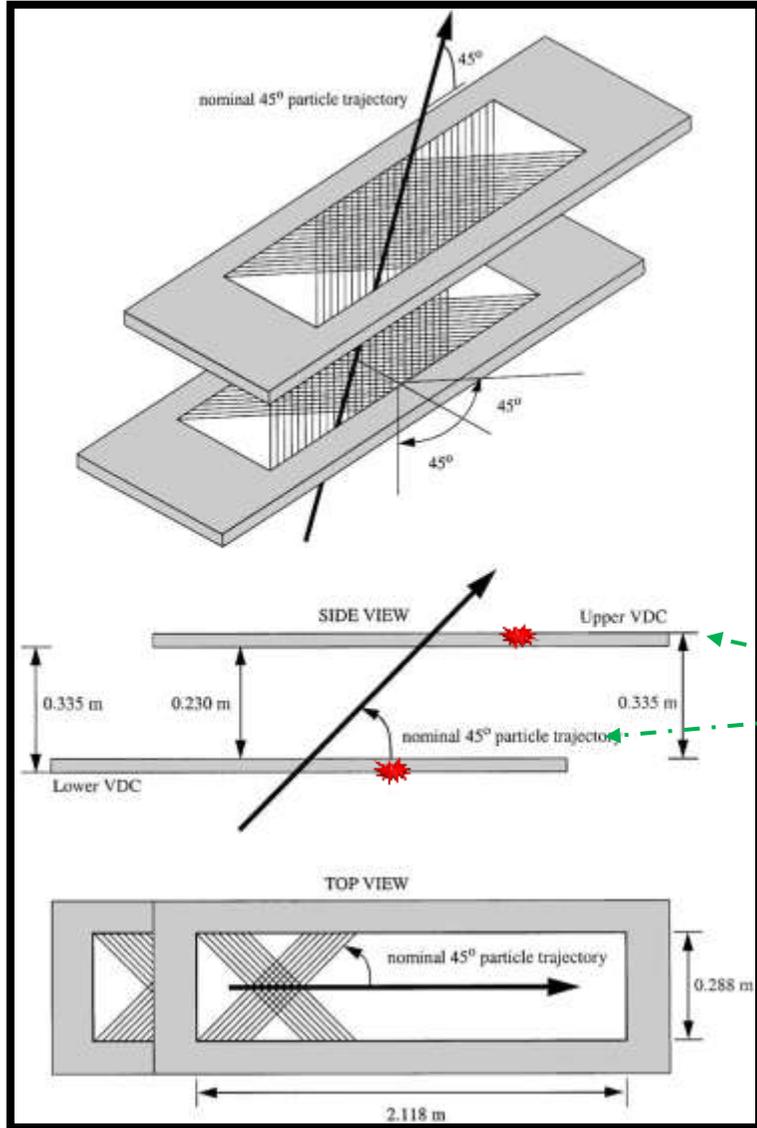
□ 误差分析

- 电子散射实验综述
- 各种粒子探测器
- 数据采集系统量
- 实验分析实例

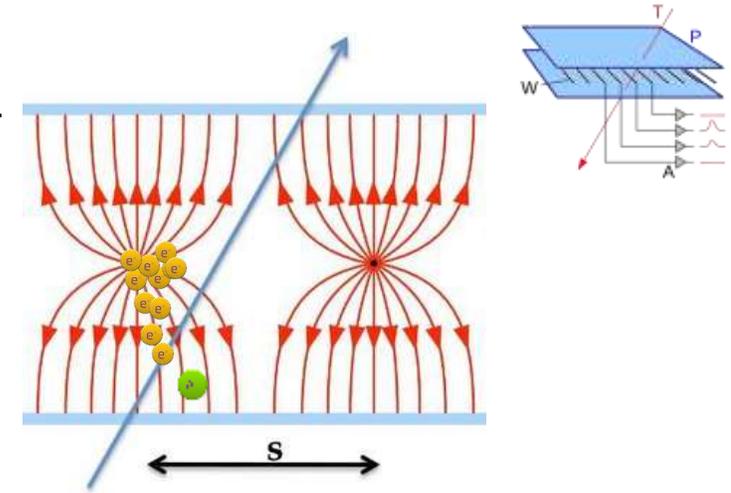
下节预告：未来展望

Backup

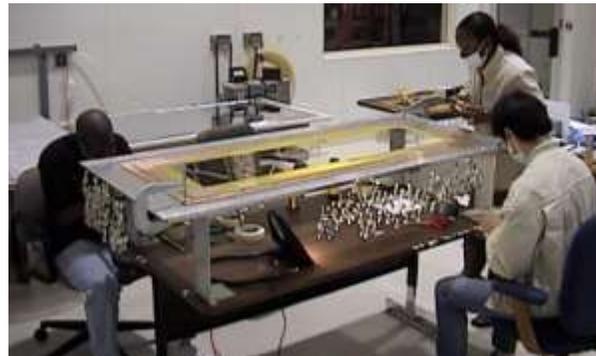
➤ Drift Chambers (DC)



- ❑ 高能带电粒子经过气体电离出外围电子
- ❑ 强电场中，电离电子产生雪崩效应产生大量次级电子；
- ❑ 电子漂移到阳极丝以脉冲信号被前端电子学记录

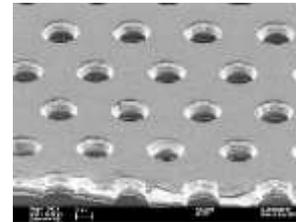
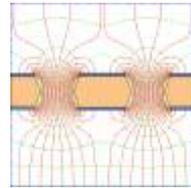
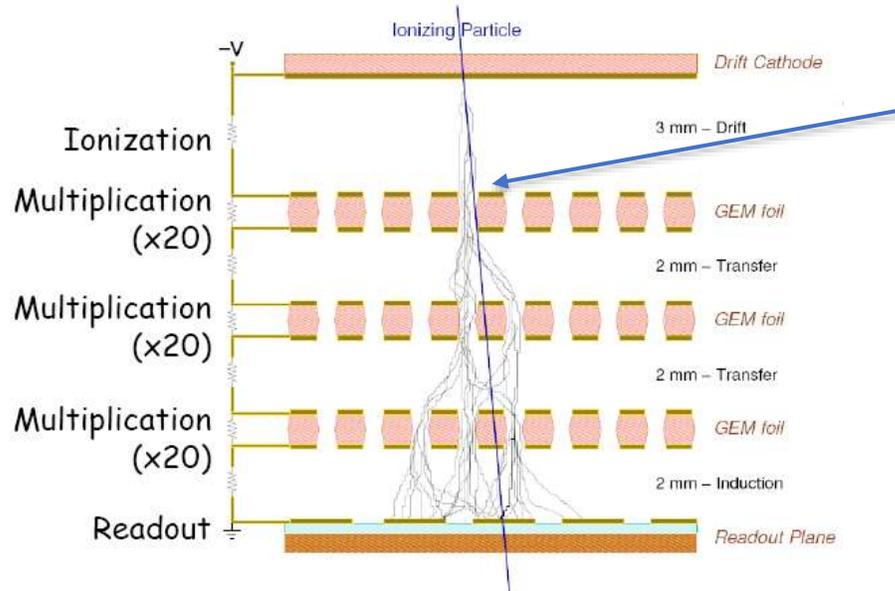
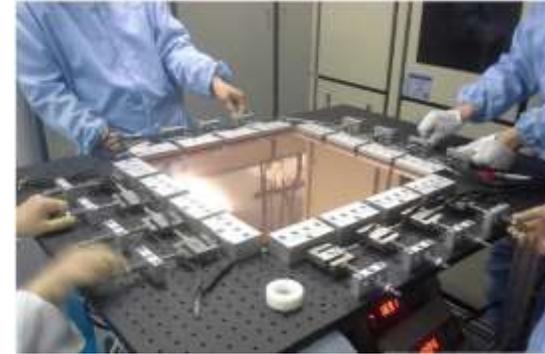


- ❖ 每个DC至少两个不同方向的阳极丝平面（45度夹角）；通过判断每个平面的哪根阳极丝有脉冲信号，即可判断入射位置
- ❖ 非电磁场环境，两组DC提供两组 (x, y) 信息，即可判断粒子穿过的位置 (x, y) and 角度 (θ, ϕ)

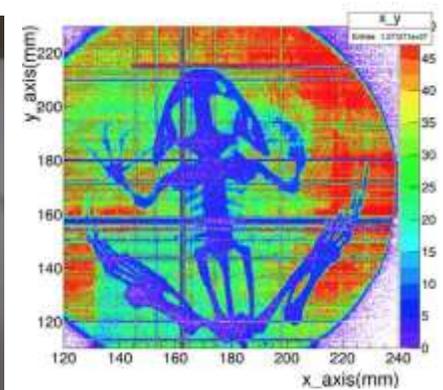


➤ Gas Electric Multiplier (GEM)

- ❑ 在多层薄膜上精细打洞
- ❑ 洞内部产生强电场
- ❑ 电离电子经过洞产生雪崩放大，末端记录大信号
- ❑ 可在强磁场中工作
- ❑ 高计数率 ($50\text{MHz}/\text{mm}^2$)



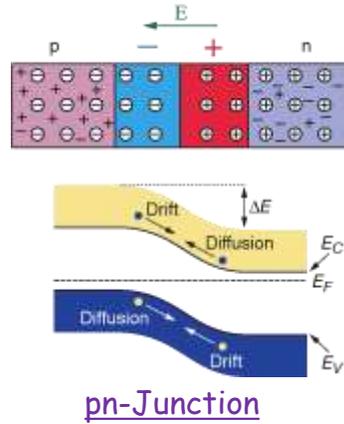
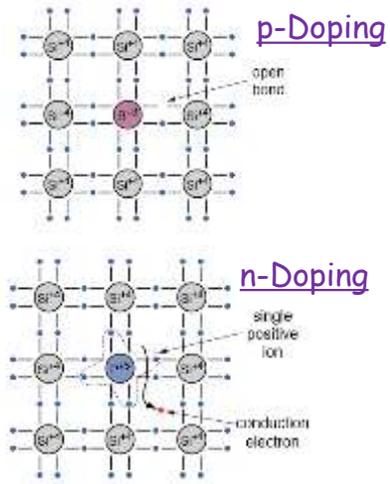
❖ 极高的位置分辨率



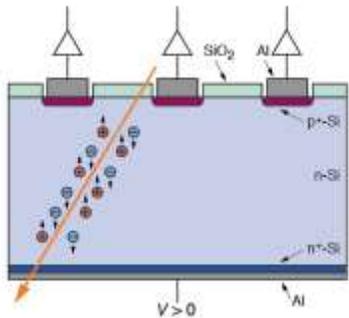
- ❑ 电子脉冲信号被底部读出条读出
- ❑ 其他类似MPGD技术 (muRELL)

➤ Silicon Tracker

□ 硅半导体pn-Junction



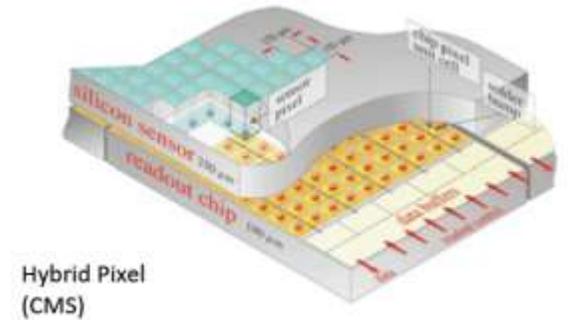
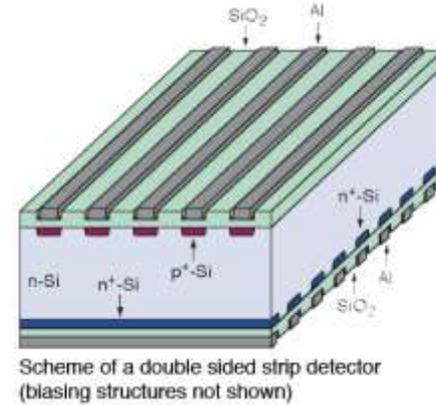
□ 高能带电粒子硅半导体的电子并产生雪崩信号



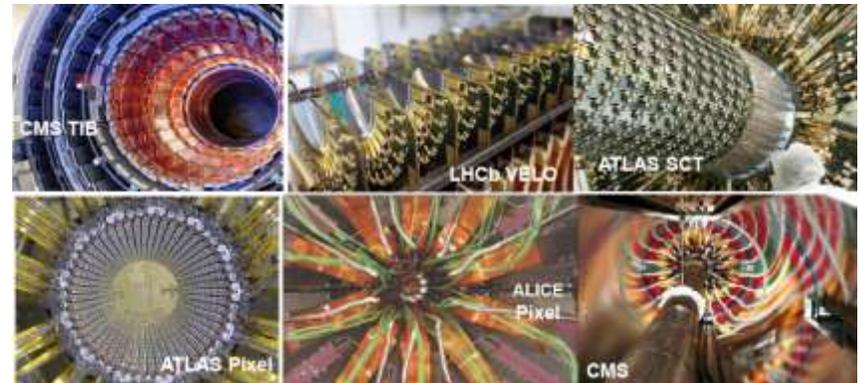
- ✓ Voltage < 200 V
- ✓ Amplification ~10⁷

□ 电子脉冲信号被底部读出条读出

❖ 长条读出 (时间探测) 和像素读出 (时间和位置探测)



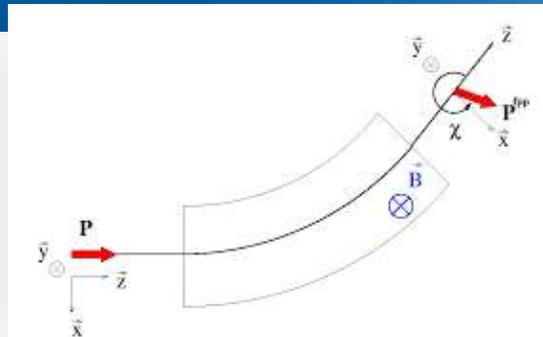
- ❖ 可在强磁场工作 (顶点测量) ; 可塑性高
- ❖ 怕温度变化; 怕中子辐射



轨迹探测器

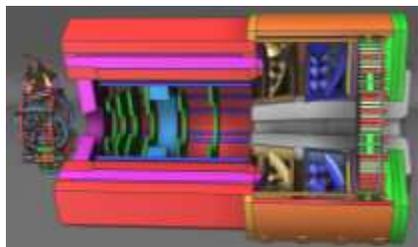
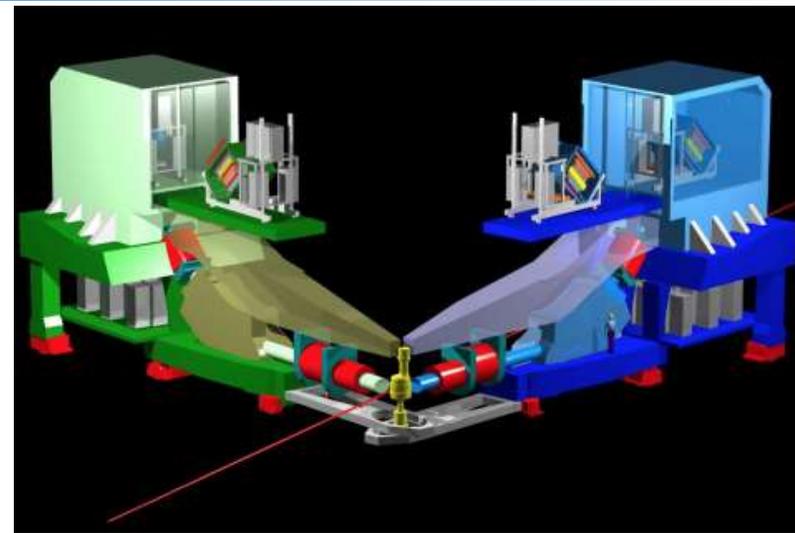
➤ 带电粒子动量

❑ 在均匀强磁场中，带电粒子的偏转半径由其动量决定

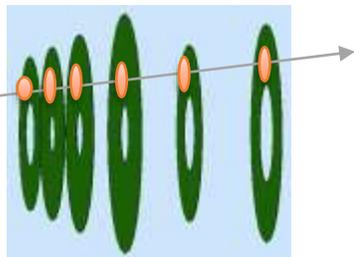


❑ 高精度磁谱仪（双极子+四极子磁铁组合）→可测量0.1%精度的带电粒子动量

❑ In a Solenoid magnet (e.g., SoLID, EIC, or other collider detectors), 多组tracking探测器可决定粒子的偏转轨迹，从而获得去动量

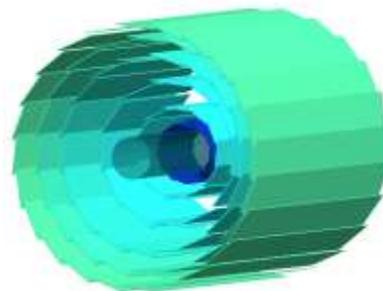
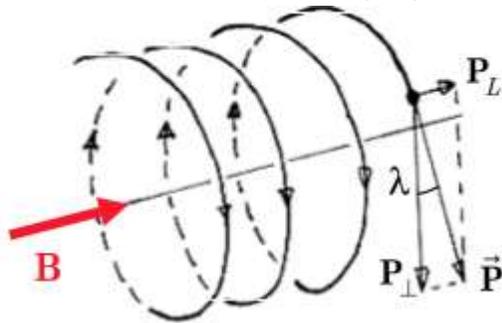


SoLID

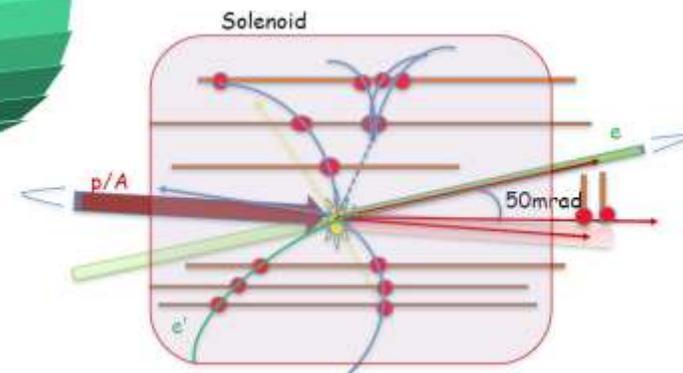


固定靶实验高能粒子趋
近于朝前飞行

$$R(m) = \frac{P_T(\text{GeV})}{0.3 \cdot B(T)}$$



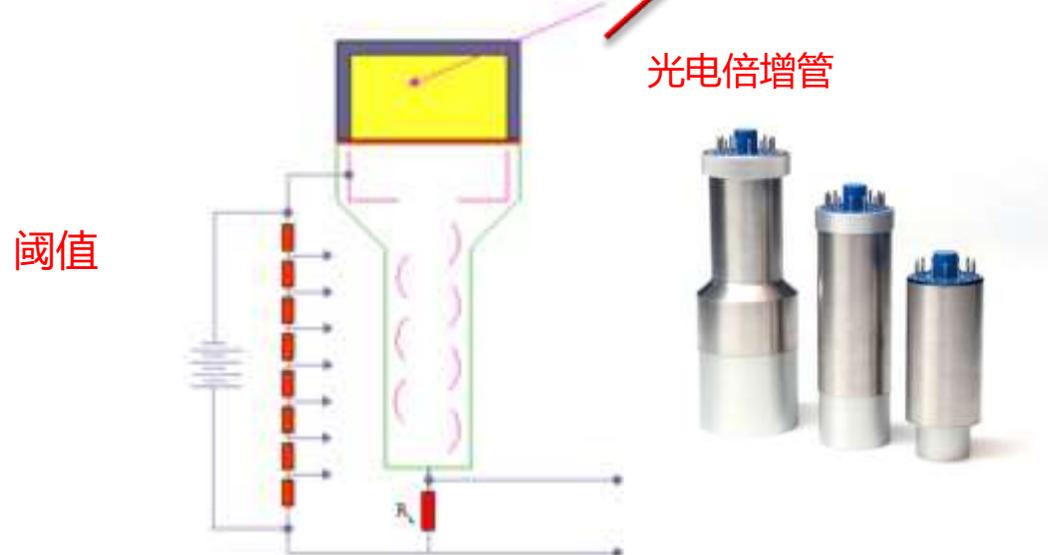
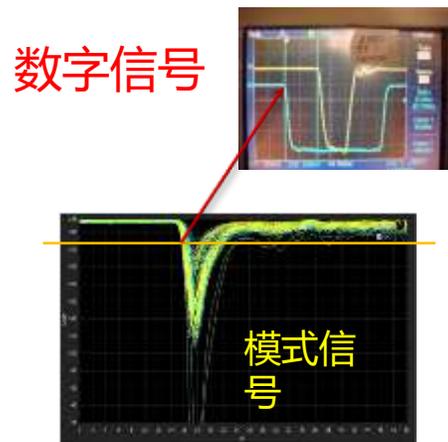
对撞机探测器



对撞实验中粒子“均匀”飞向各个方向

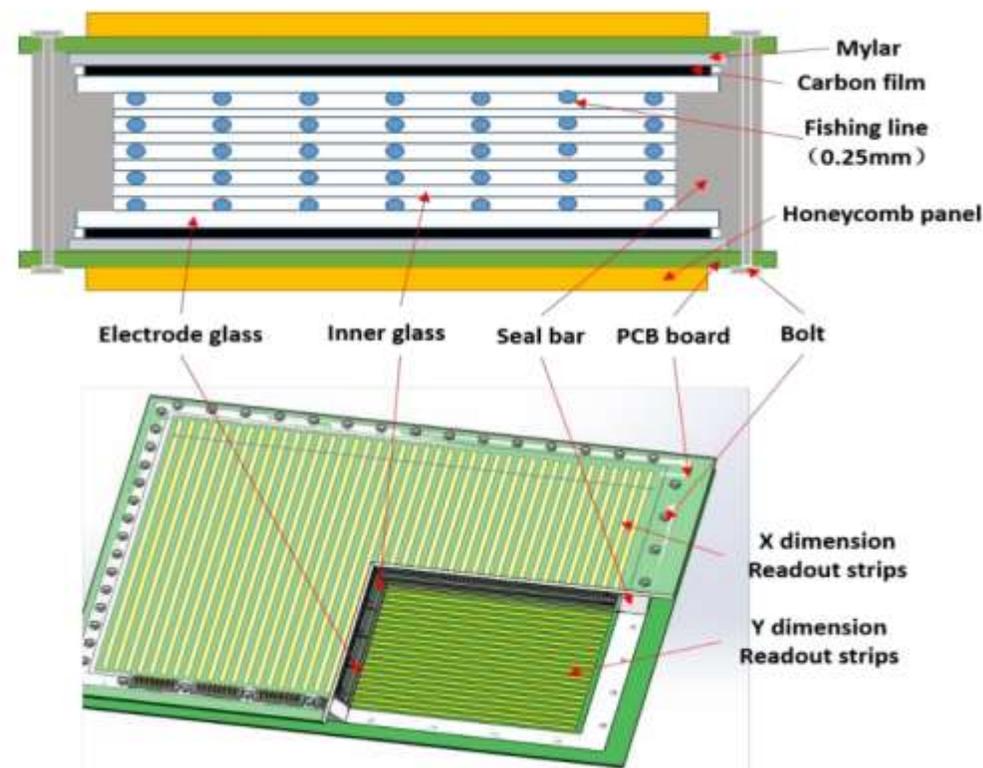
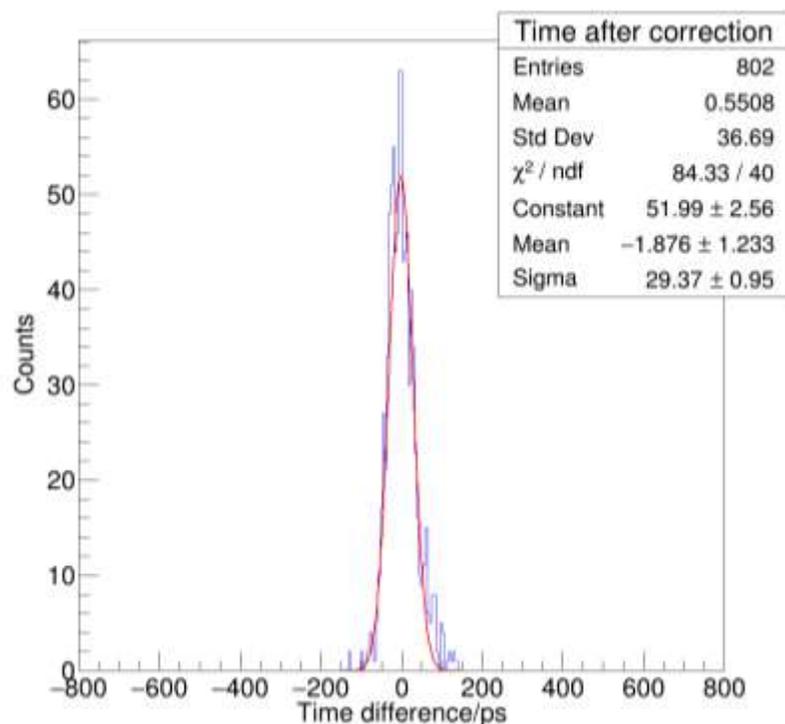
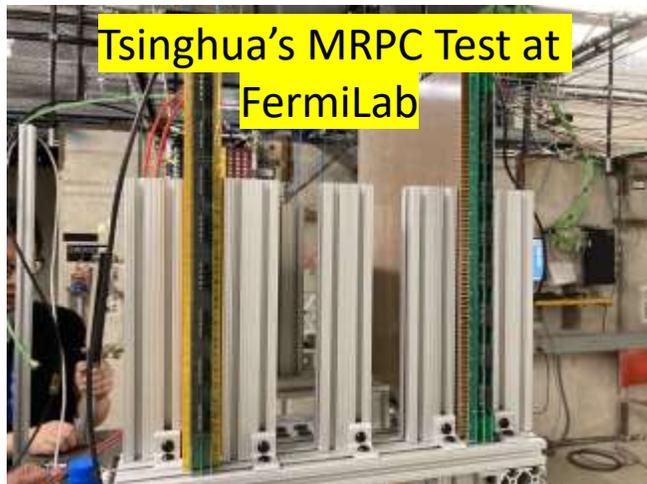
➤ 塑料闪烁体

- ❖ 塑料掺杂荧光材料
- ❖ 带电粒子穿过塑料并电离产生光信号
- ❖ 光电倍增管 (Photon-Multiplier Tube, PMT) 将光信号转变为“模式”电信号
- ❖ 以超过某一阈值的模式信号为时间起点 (粒子穿过时间)
- ❖ 时间信息以数字信号方式传递到数据采集系统



➤ Multi-gap Resistive Plate Chamber (MRPC)

- ❖ 气体由多层低电阻玻璃分割成极薄气隙
- ❖ 5000~15000V高压
- ❖ 电离电子在强电场的气隙中雪崩放大并被迅速收集脉冲信号
- ❖ 极高时间分辨率 (<20ps)
- ❖ 极高计数率- 50 KHz/cm²
- ❖ 耐辐射; 耐强磁场



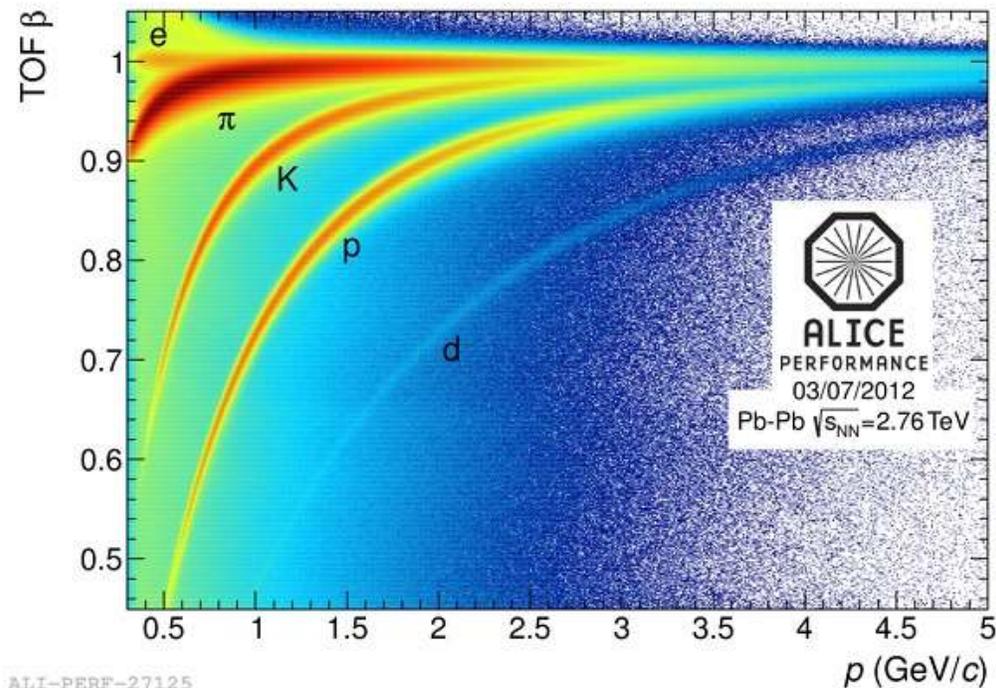
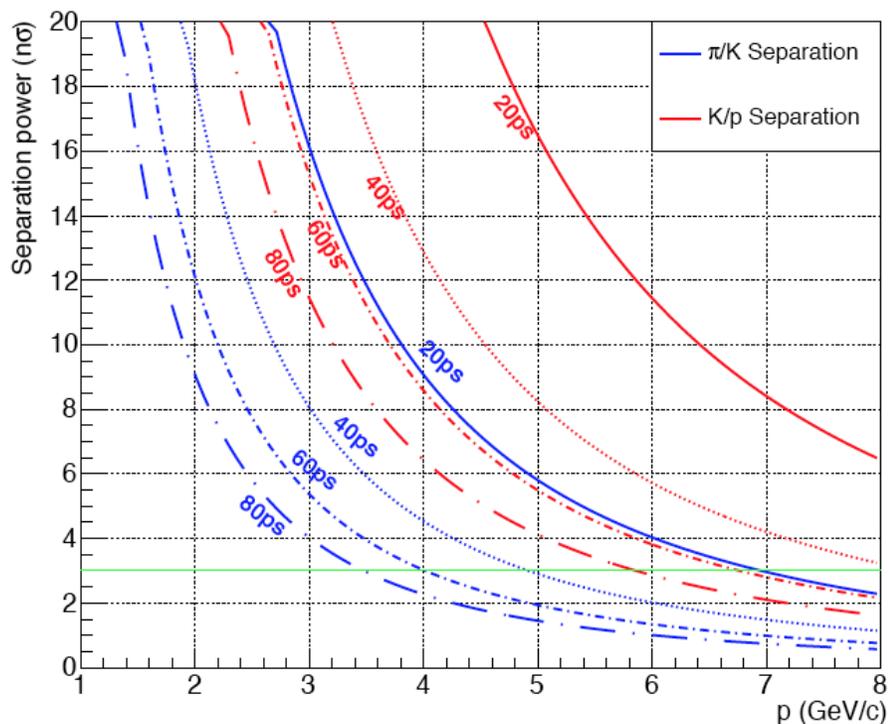
Time-of-Flight (TOF):

带电粒子的速度可由两个固定距离时间飞行仪探测器测量:

$$\beta = \frac{v}{c} = \frac{L}{(t_1 - t_2) \cdot c}$$

TOF-Beta 关联于粒子动量 (Tracking 获得) 和粒子种类 (质量):

$$\beta = \frac{p}{\sqrt{p^2 + m^2 c^2}}$$

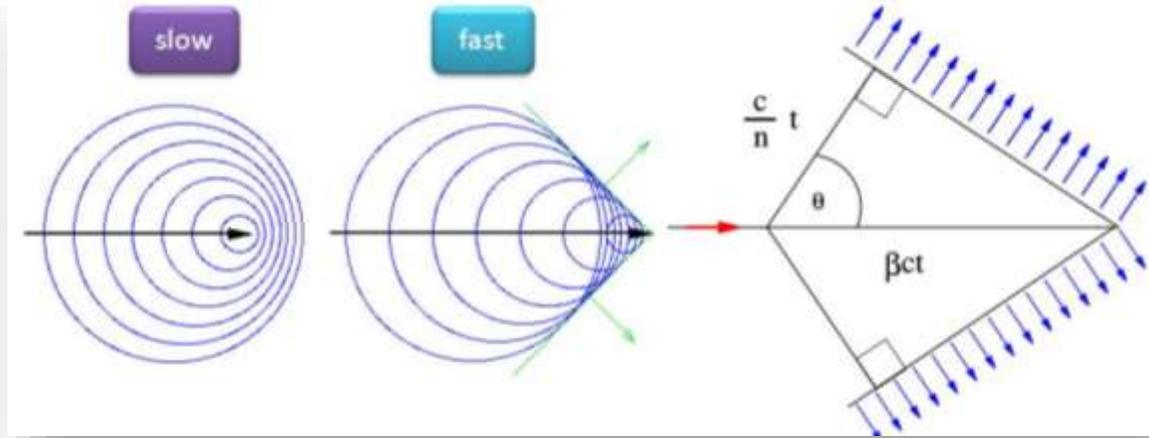


穿过同一组时间飞行仪的带电粒子, 如动量相同, 其时间差反应其质量之差:

$$\Delta t = t_1 - t_2 \simeq \frac{Lc}{2p^2} (m_1^2 - m_2^2)$$

➤ 奇伦科夫辐射原理

- 当带电粒子穿过某介质，且其速度高于介质内光速，产生奇伦科夫辐射
 - 介质内分子被带电粒子极化变为双极子 (dipoles)
 - 双极子震荡产生奇伦科夫幅射光



- 奇伦科夫光的辐射角取决于带电粒子速度和介质的折射率：

$$\cos\theta = \frac{1}{\beta n}$$

- 介质的折射率决定了特定粒子发生奇伦科夫光辐射的动量阈值：

$$P_{threshold} = \frac{mc}{\sqrt{n^2 - 1}}$$

奇伦科夫光阈值探测器

- 已知介质，已知动量，通过测量角度可获取粒子质量（种类）：

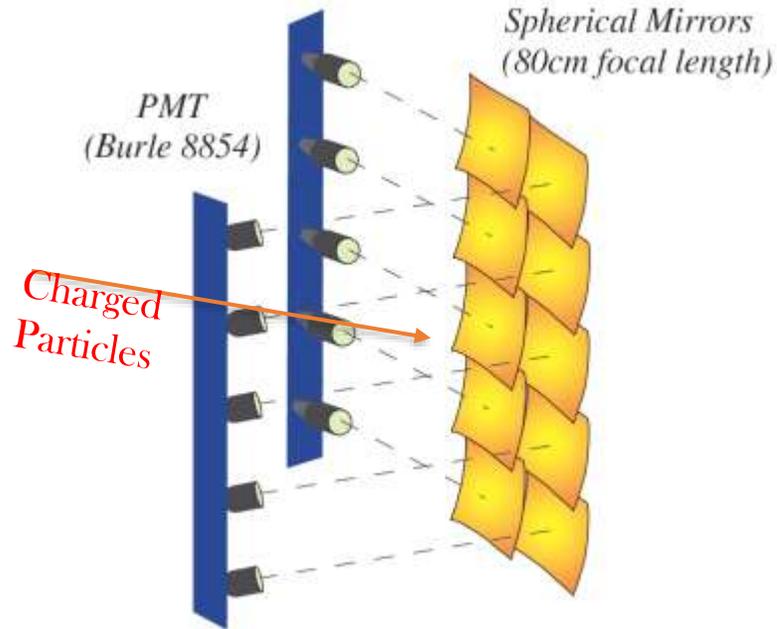
$$\beta = \frac{p}{\sqrt{p^2 + m^2 c^2}}$$

奇伦科夫光环成像探测器 (RICH)

➤ 奇伦科夫辐射原理

❑ Hall-A Gas Cherenkov Detector was tuned to allow electrons with 18 MeV/c to emit light, while pions requires 4.3 GeV/c

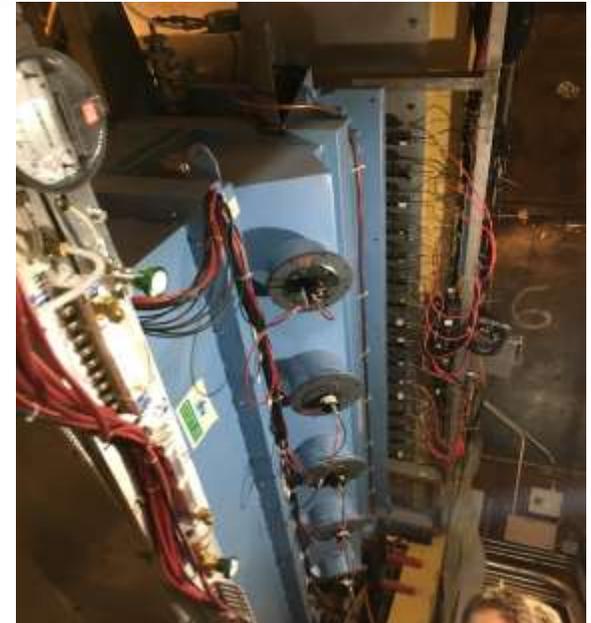
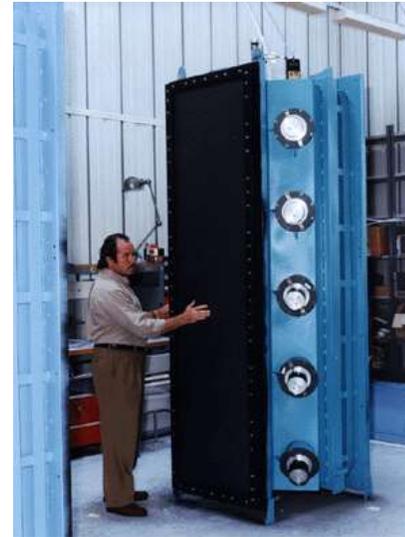
$$P_{threshold} = \frac{mc}{\sqrt{n^2 - 1}}$$



Cherenkov lights are focused by mirrors and reflected to the PMT

radiator	index	Threshold (GeV/c)			
		e	π	K	p
quartz (DIRC)	1.473	0.00048	0.13	0.47	0.88
aerogel (mRICH)	1.03	0.00207	0.57	2.00	3.80
aerogel (dRICH)	1.02	0.00245	0.69	2.46	4.67
C ₂ F ₆ (dRICH)	1.0008	0.01277	3.49	12.34	23.45
CF ₄ (gRICH)	1.00056	0.01527	4.17	14.75	28.03

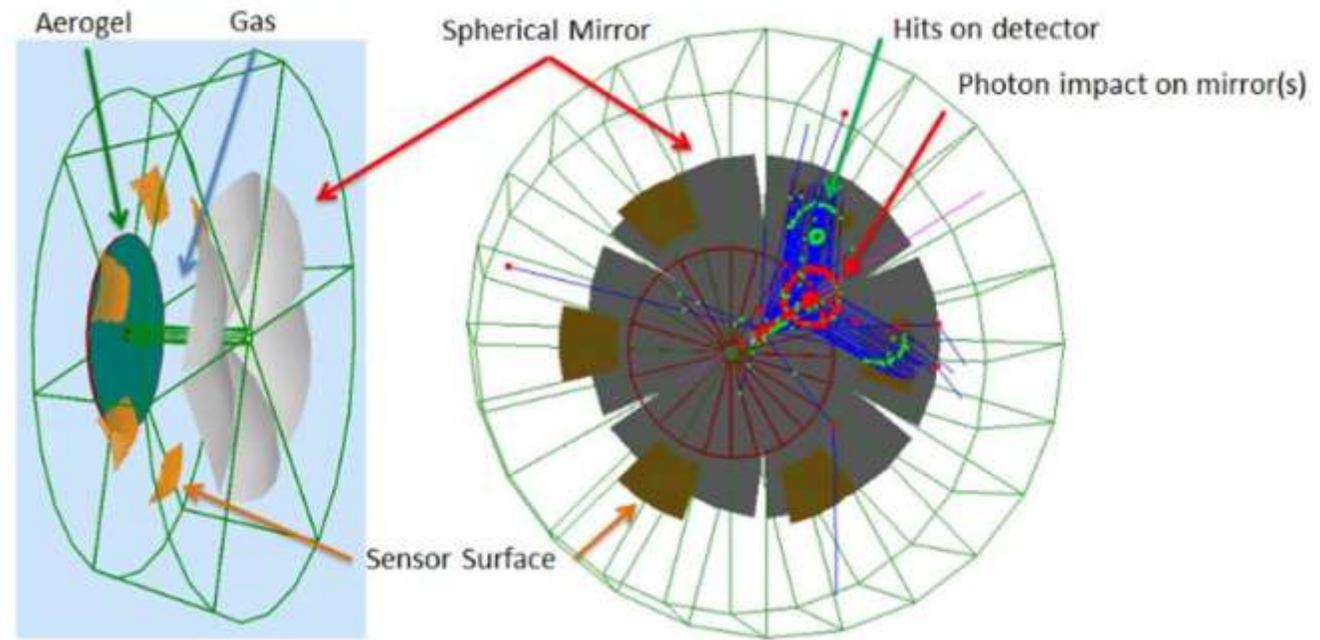
Table 11.40: Table of Cherenkov thresholds for various media.



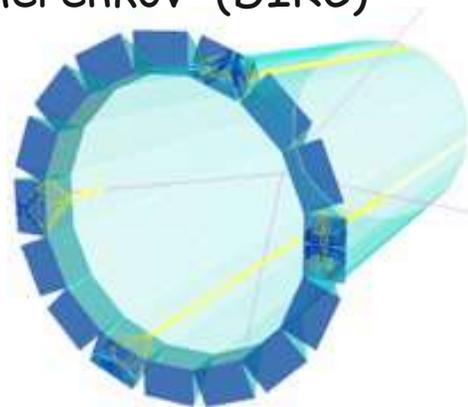
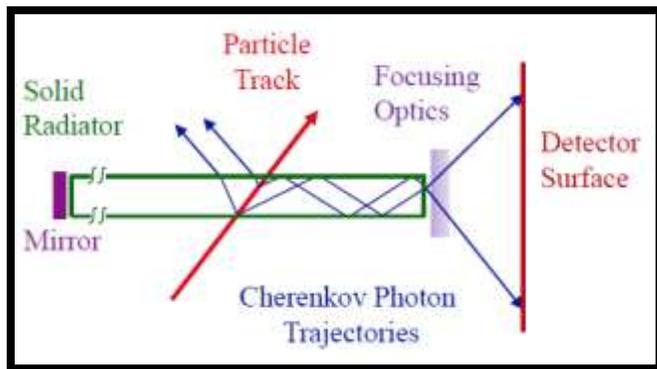
➤ 奇伦科夫辐射原理

□ Ring Image Cherenkov Detector

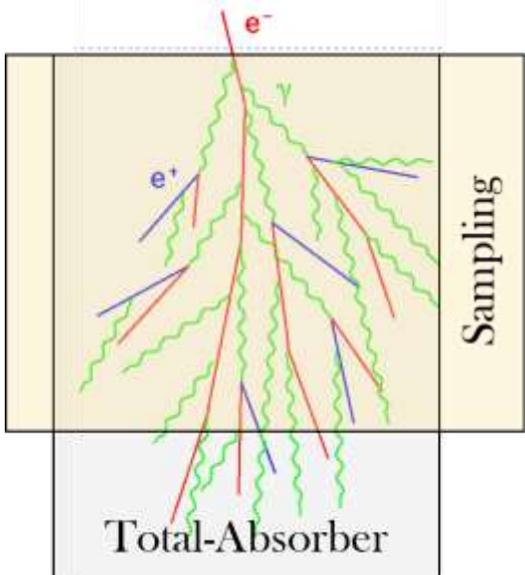
$$\cos\theta = \frac{1}{\beta n} \quad \beta = \frac{p}{\sqrt{p^2 + m^2 c^2}}$$



□ Detector of Internal Reflective Cherenkov (DIRC)

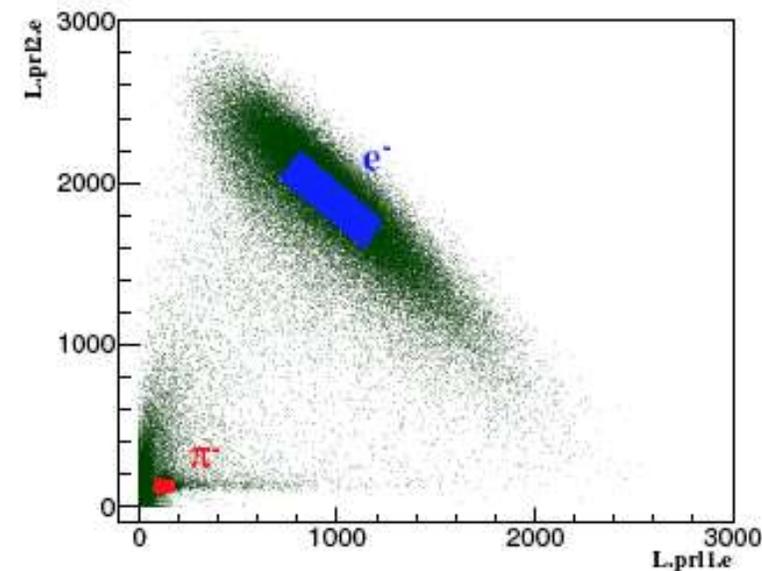
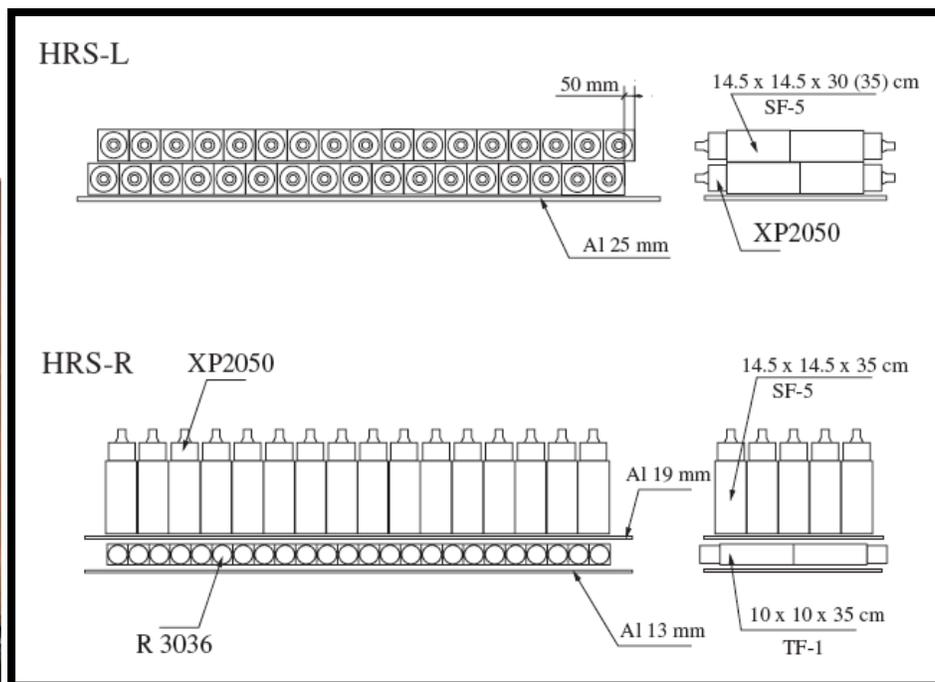


➤电磁量能器 (EM Calorimeter)



- ❑ 光子, 正负电子穿过致密材料会释放大部分能量 (电磁雪崩过程, 包括电离, 韧致辐射, 电子对产生等)
- ❑ 强子主要在EMCal里留下电离能; 需要更厚的致密材料才发生雪崩过程.
- ❑ 要求: 致密+透明, 如PbWO₄ Crystal or Pb-Glass ; 光由底部光电倍增管读出.
- ❑ 一般使用双层结构 (PreShower + Shower)

Electrons are more easily develop EM cascade even in the first layer



Pions most likely develop their energy only via ionization in the first layer;

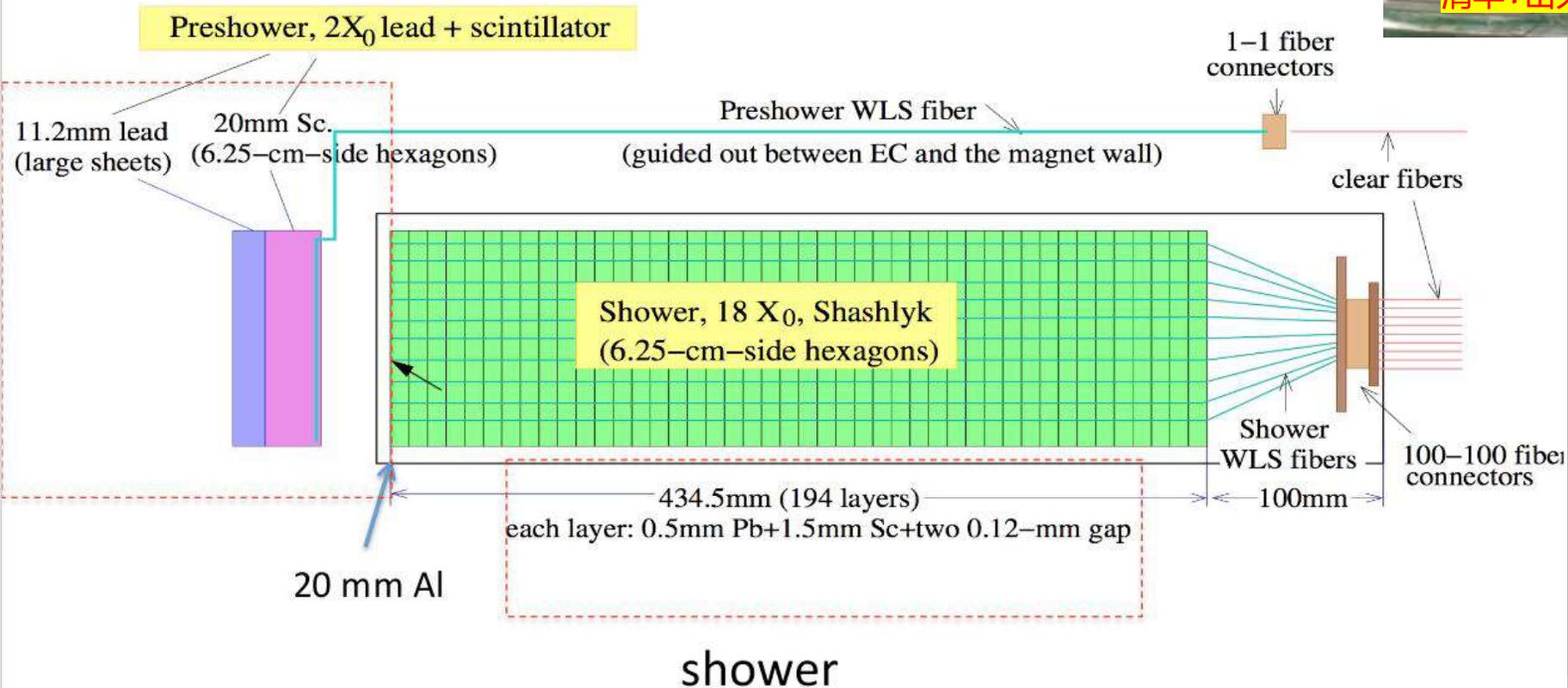
➤ Shashlyk 采样行电磁量能器 (Shashlyk EMCal)

- ❑ Shashlyk Type: (铅板+塑闪...) 叠加; 透明光纤引出光+光电倍增管读出信号
- ❑ 能量分辨率低
- ❑ 便宜; 维护容易



清华+山大研发的**Shshlyk EMCal**在**Jlab**测试

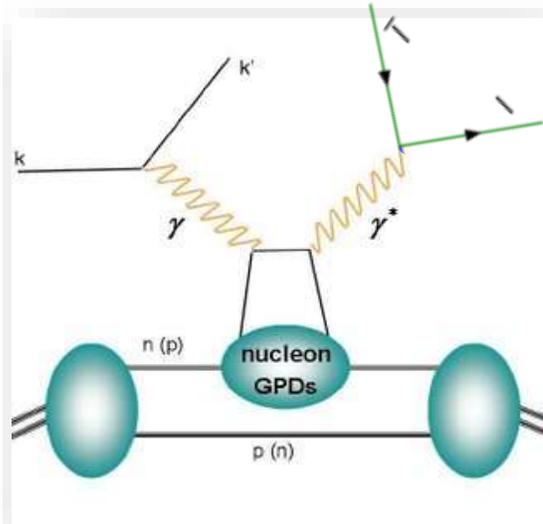
Preshower+ Al



➤ 缪子探测器系统

❑ 缪子 (Muon) 是“重电子”，物理性质几乎一样，只是更重 (105.66 MeV) 和不稳定 (2.2 us)

❑ 一些物理测量里面，测量正负缪子比测量正负电子更加干净



✓如Double-DVCS对GPD的测量

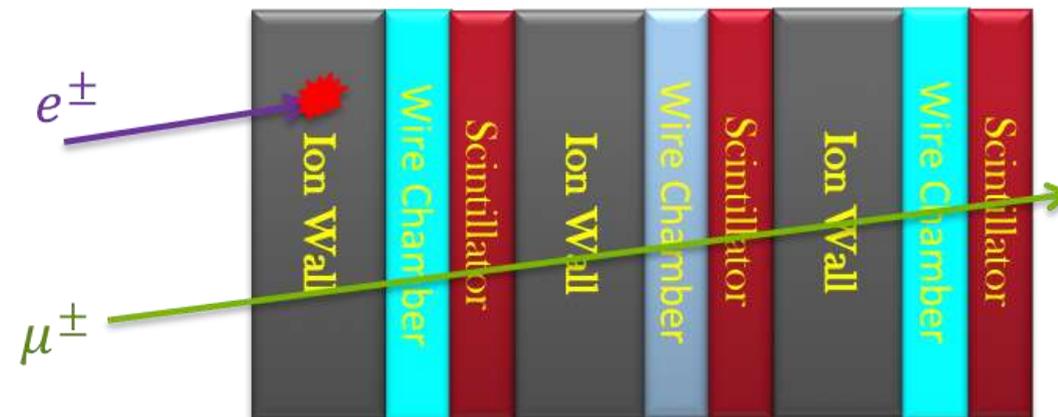
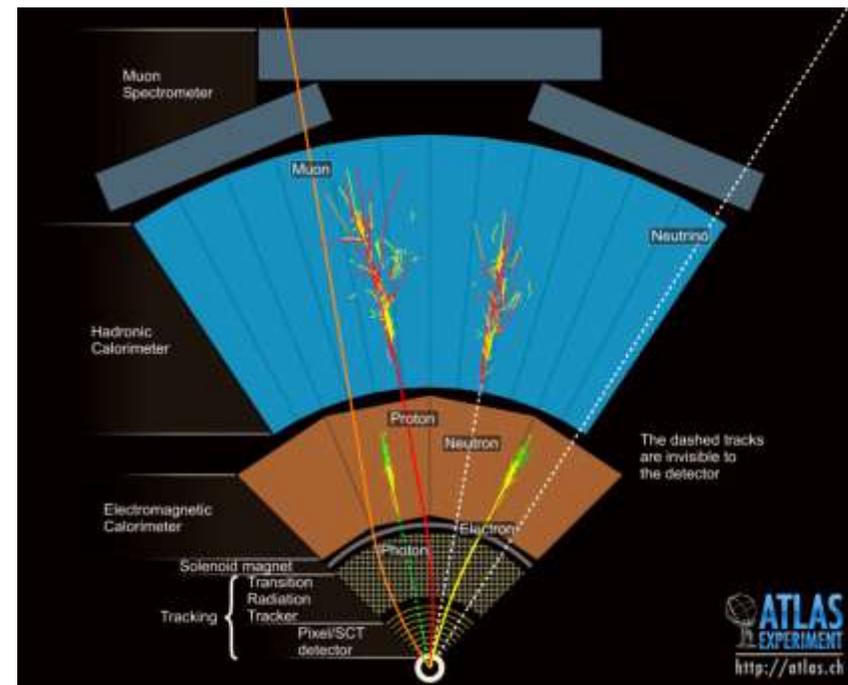
✓电子容易与其他背景混淆：

✓光子及其容易产生正负电子对

✓各种电离过程容易产生电子

❑ 缪子可以产生电离信号，但是很难产生韧致辐射 (电磁雪崩及其困难)

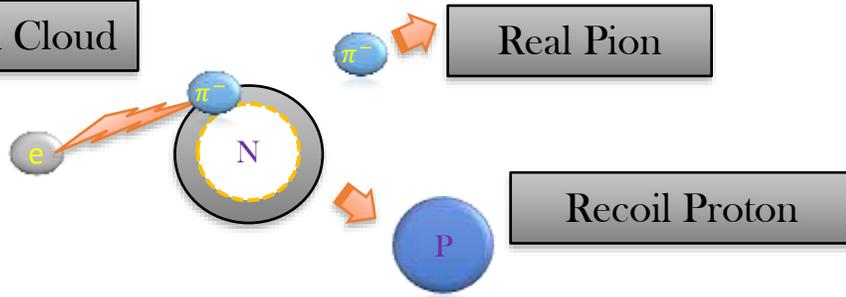
❑ 利用其极强的穿透能力，设计特殊的缪子探测器系统进行精确测量



➤反冲核子探测器系统（固定靶）

- ❑ 利用D2或He3作为“有效中子靶”，有时需确定电子击中的是中子并测量其动量
- ❑ 测量Pi介子和K介子的结构，电子击中的核子外部的Pi和K，留下低能反冲核子

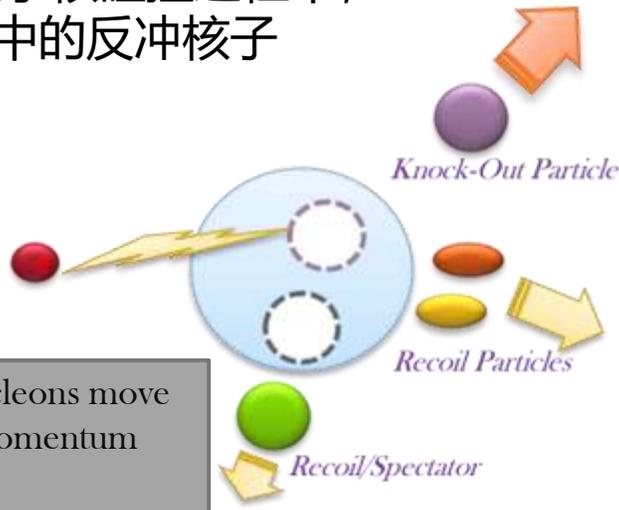
Virtual Pion Cloud



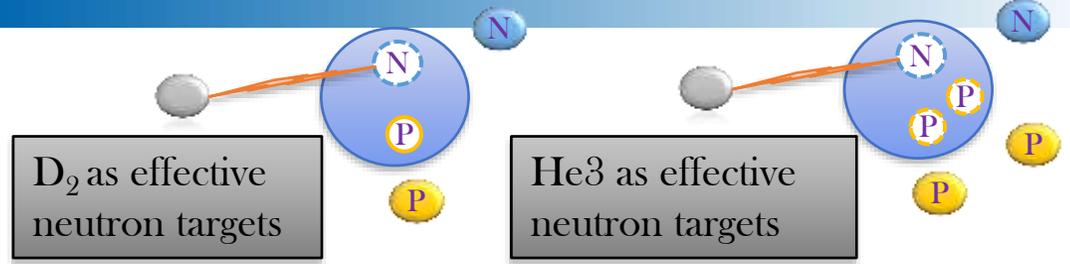
Real Pion

Recoil Proton

- ❑ 电子-原子核碰撞过程中，未被击中的反冲核子



Spectator Nucleons move with Fermi-Momentum (<math><200 \text{ MeV}/c</math>)



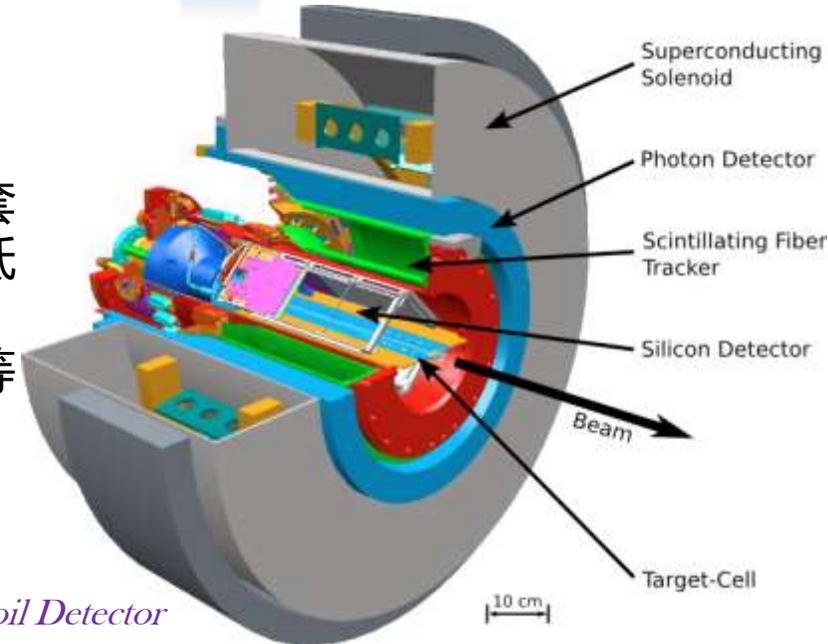
D₂ as effective neutron targets

He₃ as effective neutron targets

- ❑ 反冲核子一般能量极低，难以穿透靶区



- ❑ 针对气体靶，设计一套微型探测器系统，对低能核子进行轨迹重建，时间测量，粒子甄别等



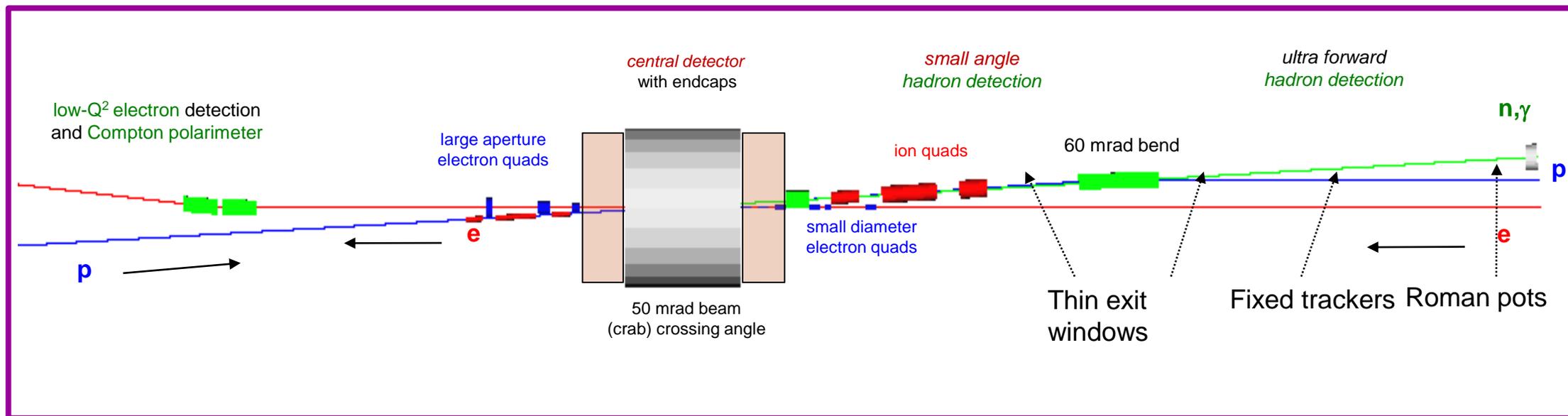
HERMES Recoil Detector

➤反冲核子探测器系统（对撞机）

❖ 在对撞机系统里，反冲核子拥有高动量（来自加速器）→更容易探测

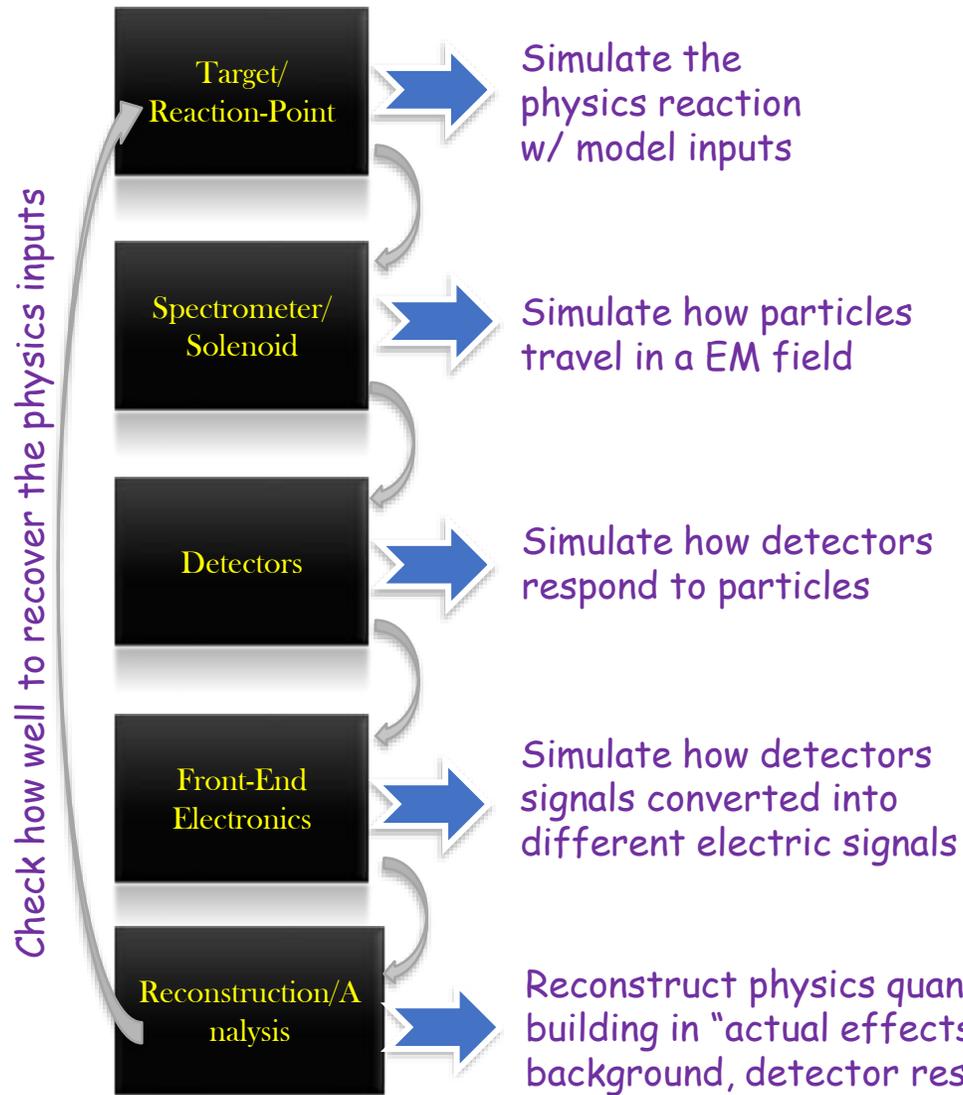


❖ 一般在对撞点远端（100m），探测器系统结合加速器系统，进行反冲核子的测量



➤ General Idea:

- ❖ A Monte-Carlo Simulation is basically like running an experiment in your computer



- ❖ Many application
 - ✓ Design experiments
 - ✓ Design detectors
 - ✓ Evaluate detector/electronics performance
 - ✓ Help pre-processing of the online experimental data
 - ✓ Apply correction and evaluation uncertainties to real data
- ❖ Software and Tools
 - ✓ Event Generators for dedicated physics processes (w/ theory models)
 - ✓ Fast simulation to study particles propagating through EM fields and materials (DYI codes)
 - ✓ Sopesticated MC simulation, using Geant4