

Synchrotron radiation background

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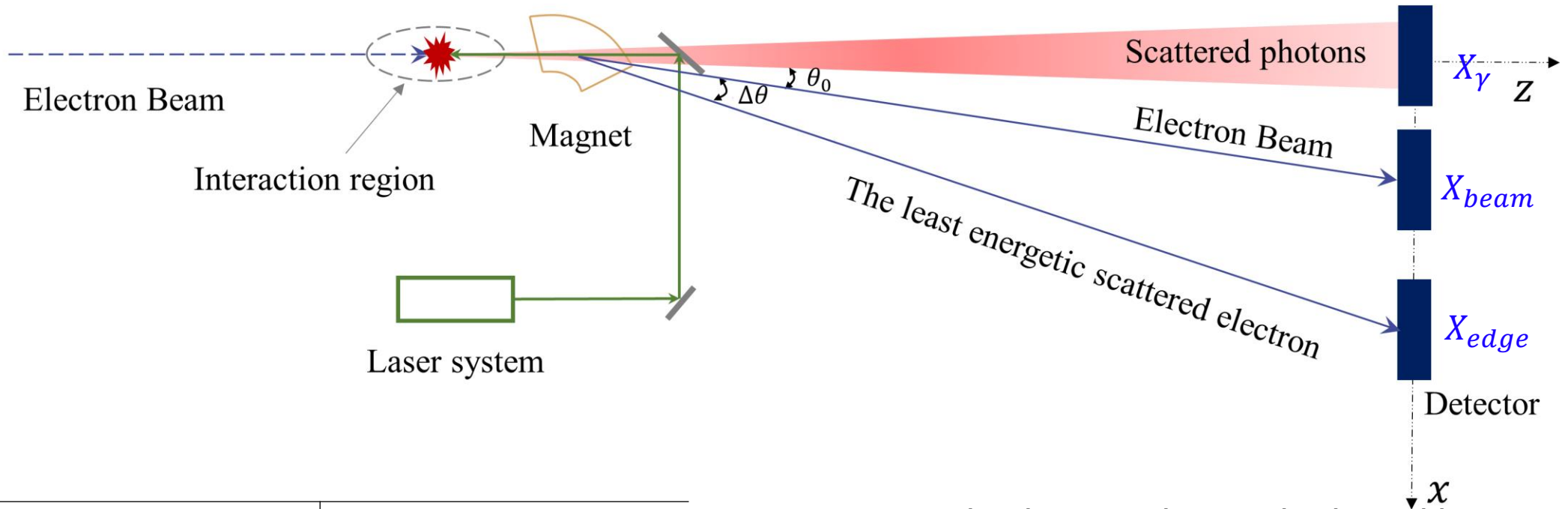
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Outline

- Calibration of beam energy
- Properties of scattered photons
- Synchrotron radiation

Method of energy calibration

- **Method:** Compton back-scattering combining a bending magnet



Electron beam		Nd:YAG Laser system	
Energy (GeV)	120	$\lambda(\text{nm})$	532
N_e	15×10^{10}	Energy(J)	0.1
Collision angle α	$\sim 2.35 \text{ mrad}$		
Compton scattering cross section	202 mb		

- Compton back-scattering method used in BEPC by measuring the energy of scattered photons with accuracy is 2×10^5 .

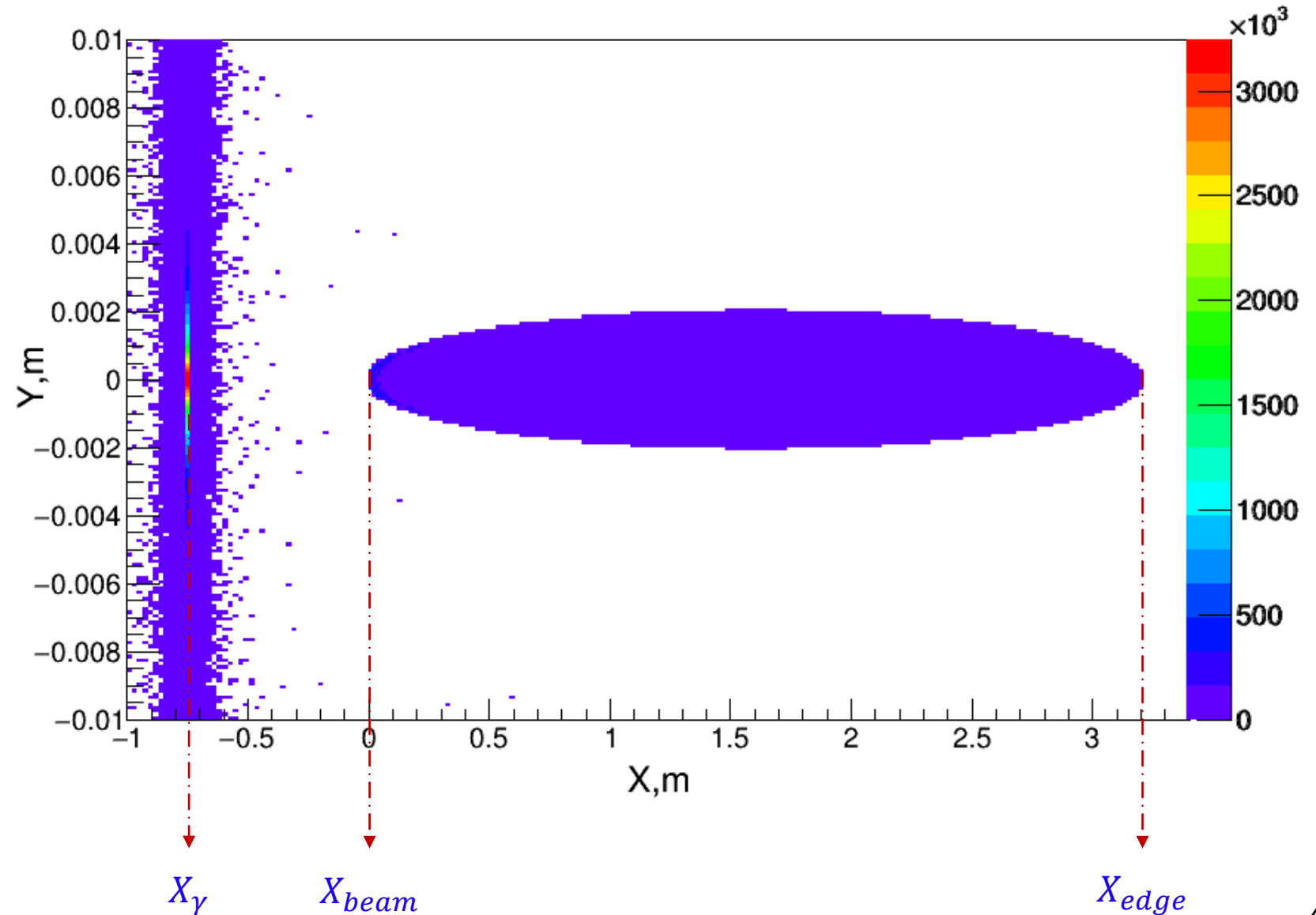
<https://doi.org/10.1016/j.nima.2011.08.050>

- The technique is “non-destructive”: $\sim 10^6$ Compton scattered particles in one collision.

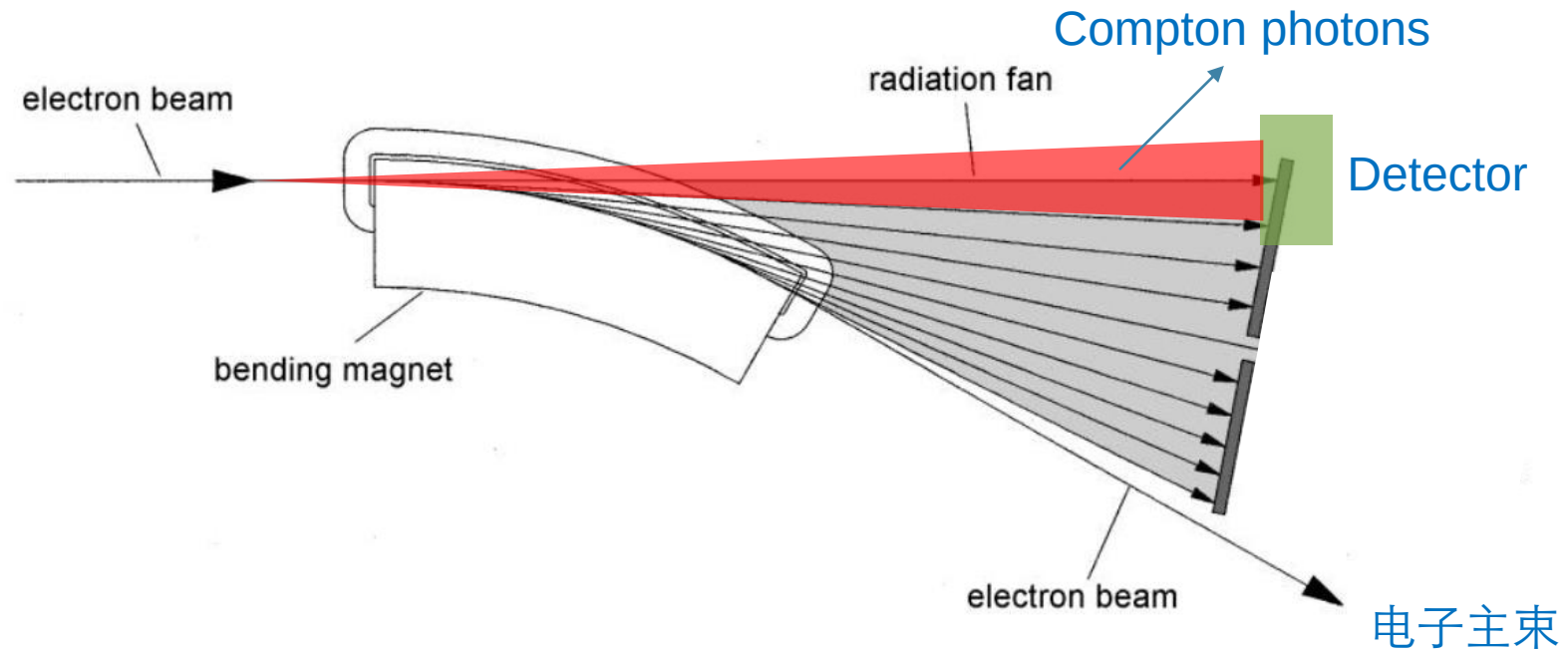
Spatial distribution of scattered particles

- Beam energy can be calibrated by:
 - Position of the main electron beam particles(X_{beam}).
 - Position of scattered photons(X_γ).
 - Position of the scattered electrons with the least energy(X_{edge}).

$$E_{beam} = \frac{(m_e c^2)^2}{4W_0} \frac{X_{edge} - X_{beam}}{X_{beam} - X_\gamma}$$



Synchrotron radiation background



Bending magnet

B	0.5 T
length	3 m

Theory

- For a horizontally accelerated electron, per solid angle

$$d\Omega = \sin\vartheta d\vartheta d\varphi$$

ϑ is polar angle;
 φ is azimuth angle

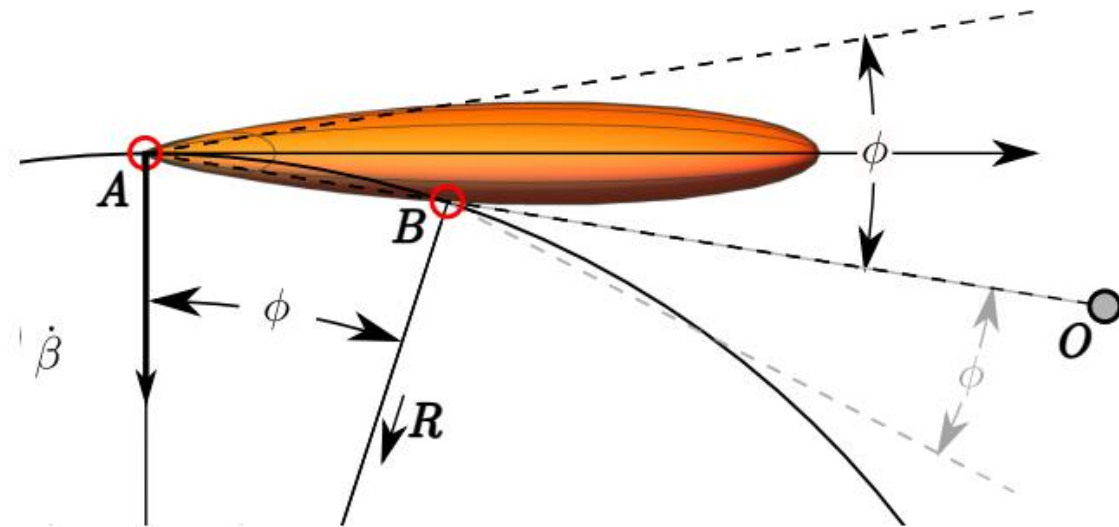
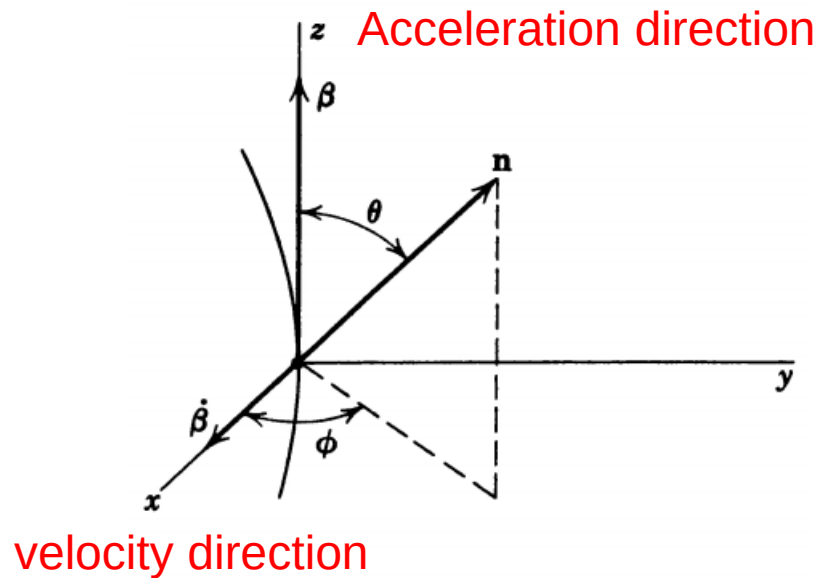


Figure 2.4: Synchrotron radiation cone of width $\phi = 1/\gamma$ sweeping across an observation point O while the photon source is moving from points A to B on the circular trajectory.

SR properties

- Bending radius

Bending radius of charged particle path in homogeneous field B at energy E is

$$\frac{1}{\rho}[\text{m}] = \frac{eB}{\beta E} = 0.2998 \frac{B[\text{T}]}{\beta E[\text{GeV}]} \quad (1)$$

Magnet parameter: $B = 0.5 \text{ T}$, length = 3 m;

CEPC Higgs mode: $E_{beam} = 120 \text{ GeV}$

$$\frac{1}{\rho}[\text{m}] = 0.2998 \frac{0.5}{120} = 0.0012$$

$$\rho = 800.5337 \text{ m}$$

1. Critical energy

发射的光子的特征能量 E_c 与特征波长 λ_c 是:

$$E_c = \hbar \frac{3}{2} \frac{c\gamma^3}{R} = 4.7948 \text{ MeV}$$

连续光谱的特征波长是:

$$\lambda_c = \frac{4\pi}{3} \rho \gamma^{-3} = 2.5876 \text{ e-13 m}$$

Theory

当一个带电粒子在一个弯转磁铁中作圆周运动时，若在圆形轨道外一个固定点上观察，可以接受到电子轨道上的一小段圆弧上来的同步辐射，这一小段辐射脉冲可以表示同步辐射的频谱。

同步辐射的完全谱分布：

$$\frac{d^2 P}{d\Omega d\omega} = \frac{P_\gamma \gamma}{\omega_c} F(\omega, \psi), \quad (5.11)$$

其中 Ω 是 solid angle, ω_c 为临界频率, $F(\omega, \psi)$ 为谱的角分布函数, 用 $F_\sigma(\omega, \psi)$ 和 $F_\pi(\omega, \psi)$ 代表水平和垂直偏振分量:

$$P_\gamma = \frac{2}{3} \frac{e^2 c}{4\pi\epsilon_0} \frac{\beta^4 \gamma^4}{\rho^2} \quad \omega_c = \frac{3c\gamma^3}{2\rho}, \quad (5.12)$$

$$F(\omega, \psi) = F_\sigma(\omega, \psi) + F_\pi(\omega, \psi), \quad (5.13)$$

$$F(\omega, \psi) = \left(\frac{3}{2\pi}\right)^3 \left(\frac{\omega}{2\omega_c}\right)^2 (1 + \gamma^2 \psi^2)^2 \left[K_{2/3}^2(\xi) + \frac{\gamma^2 \psi^2}{1 + \gamma^2 \psi^2} K_{1/3}^2(\xi) \right], \quad (5.14)$$

$$F_\sigma(\omega, \psi) = \left(\frac{3}{2\pi}\right)^3 \left(\frac{\omega}{2\omega_c}\right)^2 (1 + \gamma^2 \psi^2)^2 K_{2/3}^2(\xi), \quad (5.15)$$

$$F_\pi(\omega, \psi) = \left(\frac{3}{2\pi}\right)^3 \left(\frac{\omega}{2\omega_c}\right)^2 (1 + \gamma^2 \psi^2) \gamma^2 \psi^2 K_{1/3}^2(\xi)$$

这里, K 为分数阶修正 Bessel 函数, ξ 则为

$$\xi = \frac{\omega}{2\omega_c} (1 + \gamma^2 \psi^2)^{3/2}. \quad (5.16)$$

Theory

$$\frac{d^2 P}{d\Omega d\omega} = \frac{P_\gamma \gamma}{\omega_c} F(\omega, \psi),$$

将谱的角功率分布对 Ω 进行积分，得到同步辐射的谱分布：

$$\frac{dP}{d\omega} = \frac{P_\gamma}{\omega_c} S\left(\frac{\omega}{\omega_c}\right), \quad (5.17)$$

其中

$$\begin{aligned} S\left(\frac{\omega}{\omega_c}\right) &= S_\sigma\left(\frac{\omega}{\omega_c}\right) + S_\pi\left(\frac{\omega}{\omega_c}\right) \\ &= \frac{9\sqrt{3}\omega}{8\pi\omega_c} \int_{\omega/\omega_c}^{\infty} K_{5/3}(\zeta) d\zeta, \end{aligned} \quad (5.18)$$

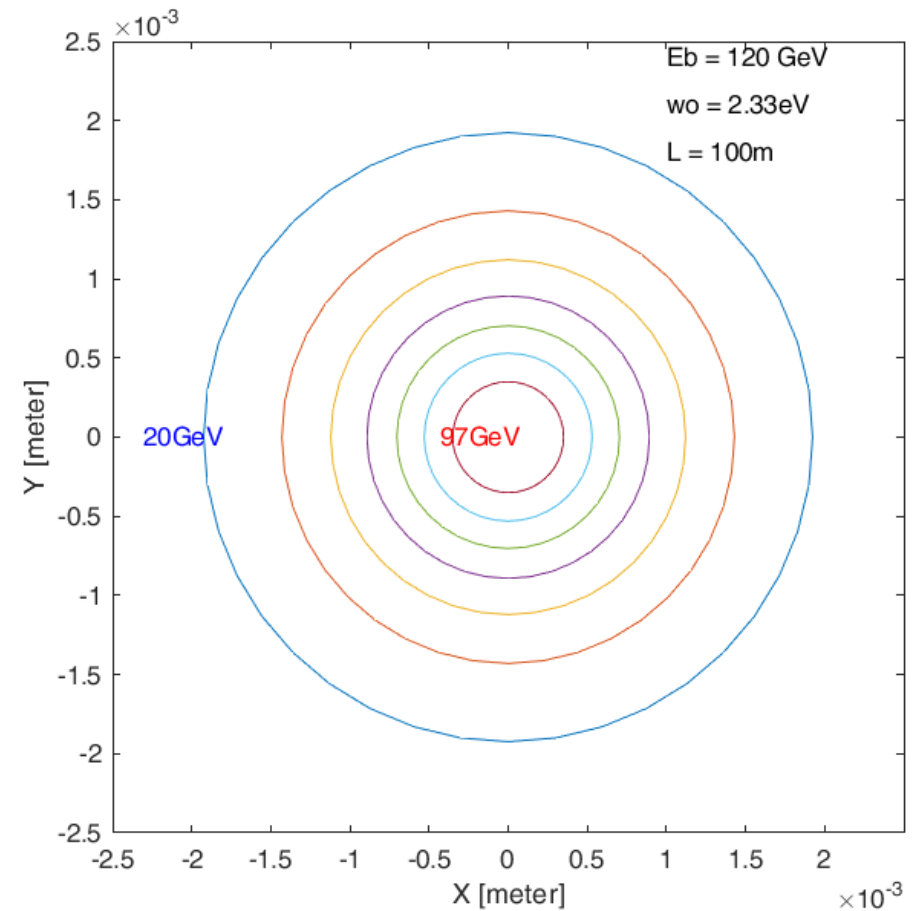
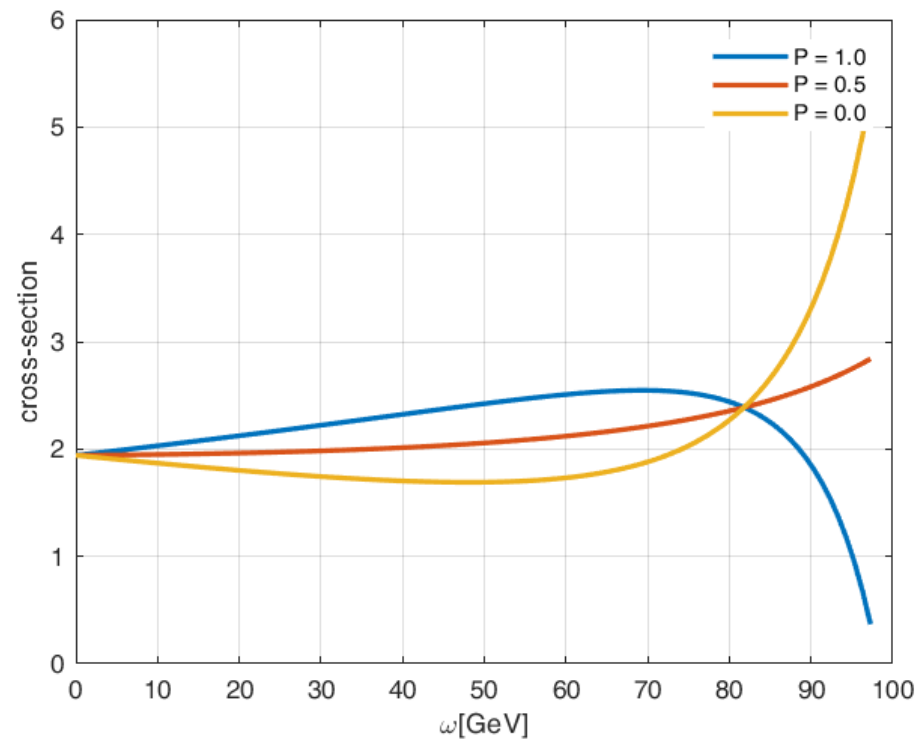
$$\begin{aligned} S_\sigma\left(\frac{\omega}{\omega_c}\right) &= \frac{9\sqrt{3}\omega}{16\pi\omega_c} \left[\int_{\omega/\omega_c}^{\infty} K_{5/3}(\zeta) d\zeta + K_{2/3}\left(\frac{\omega}{\omega_c}\right) \right] \\ S_\pi\left(\frac{\omega}{\omega_c}\right) &= \frac{9\sqrt{3}\omega}{16\pi\omega_c} \left[\int_{\omega/\omega_c}^{\infty} K_{5/3}(\zeta) d\zeta - K_{2/3}\left(\frac{\omega}{\omega_c}\right) \right]. \end{aligned} \quad (5.19)$$

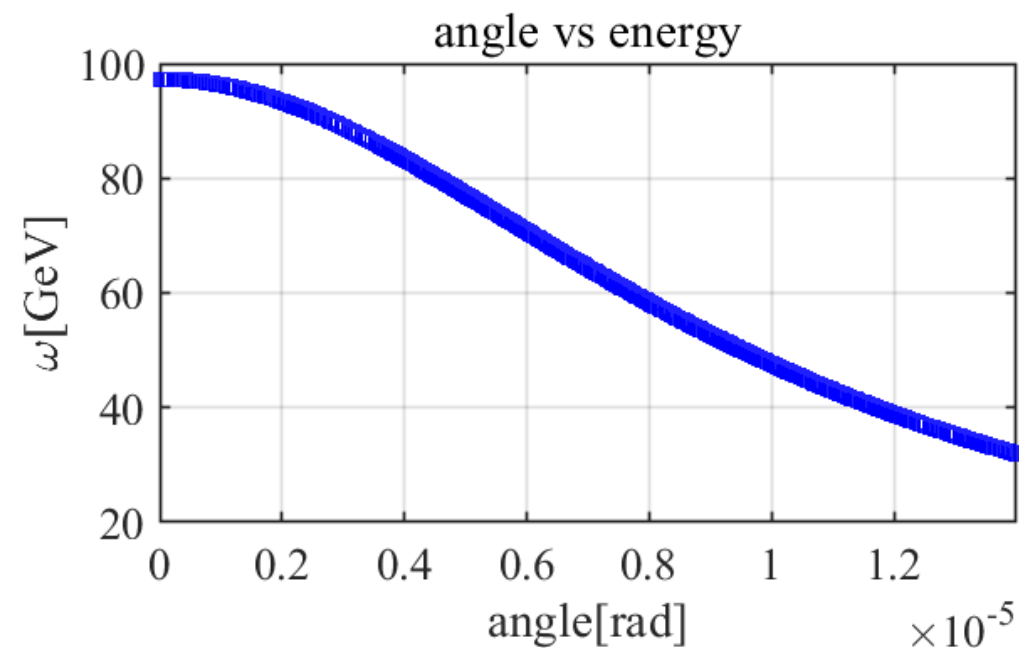
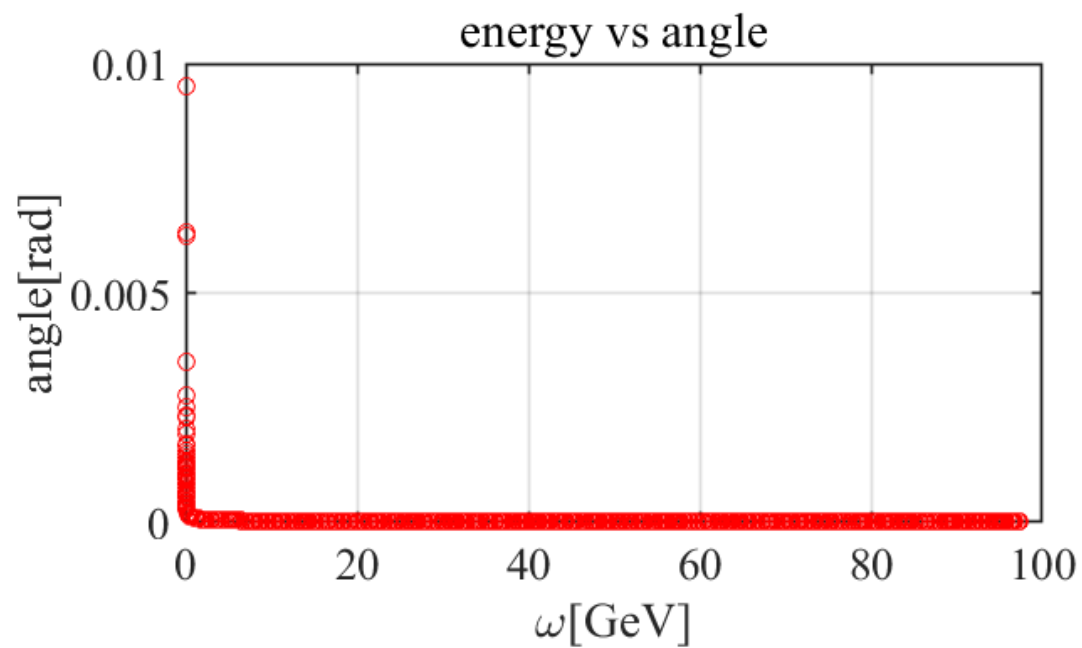
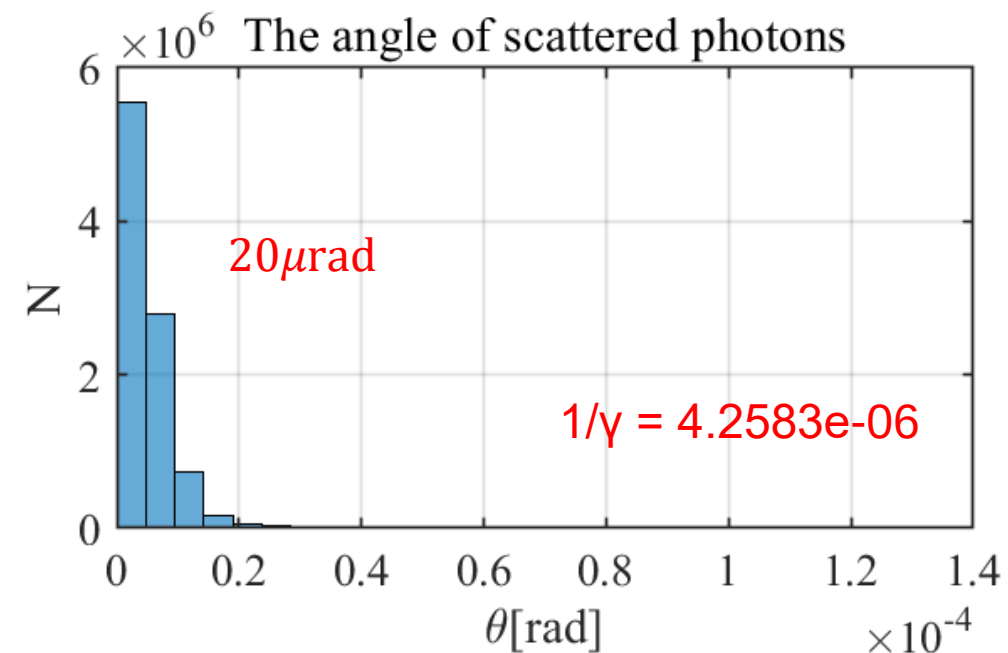
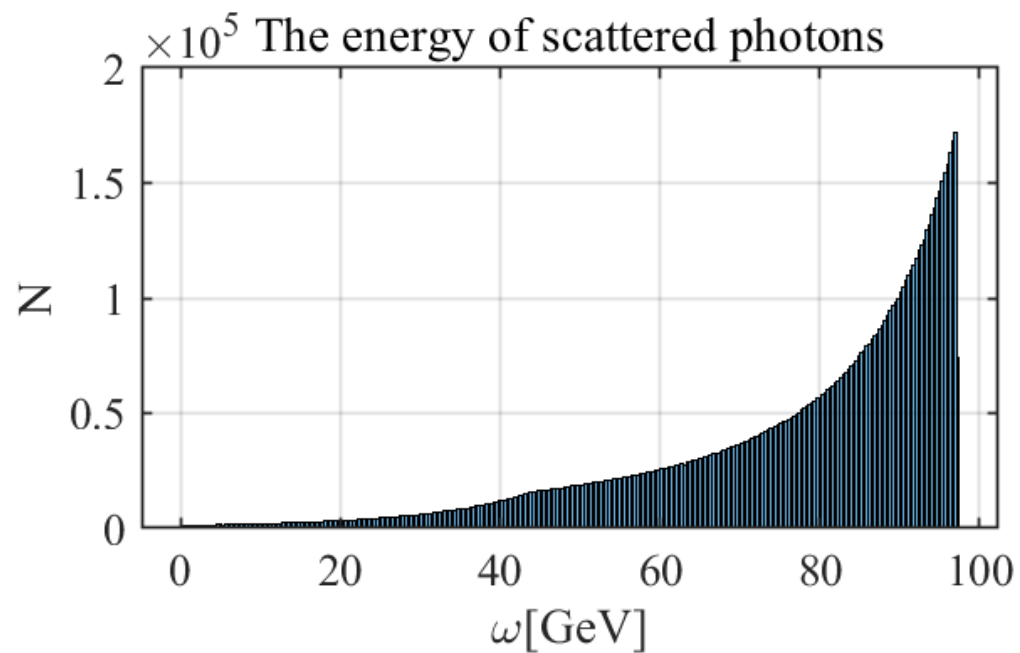
将谱的角功率分布对频率积分,得到:

$$\frac{dP}{d\Omega} = \frac{21}{32} \frac{P_\gamma}{2\pi} \frac{\gamma}{(1 + \gamma^2 \psi^2)^{5/2}} \left[1 + \frac{5}{7} \frac{\gamma^2 \psi^2}{1 + \gamma^2 \psi^2} \right].$$

散射光子分布计算

Properties of scattered photons





Properties of scattered photons

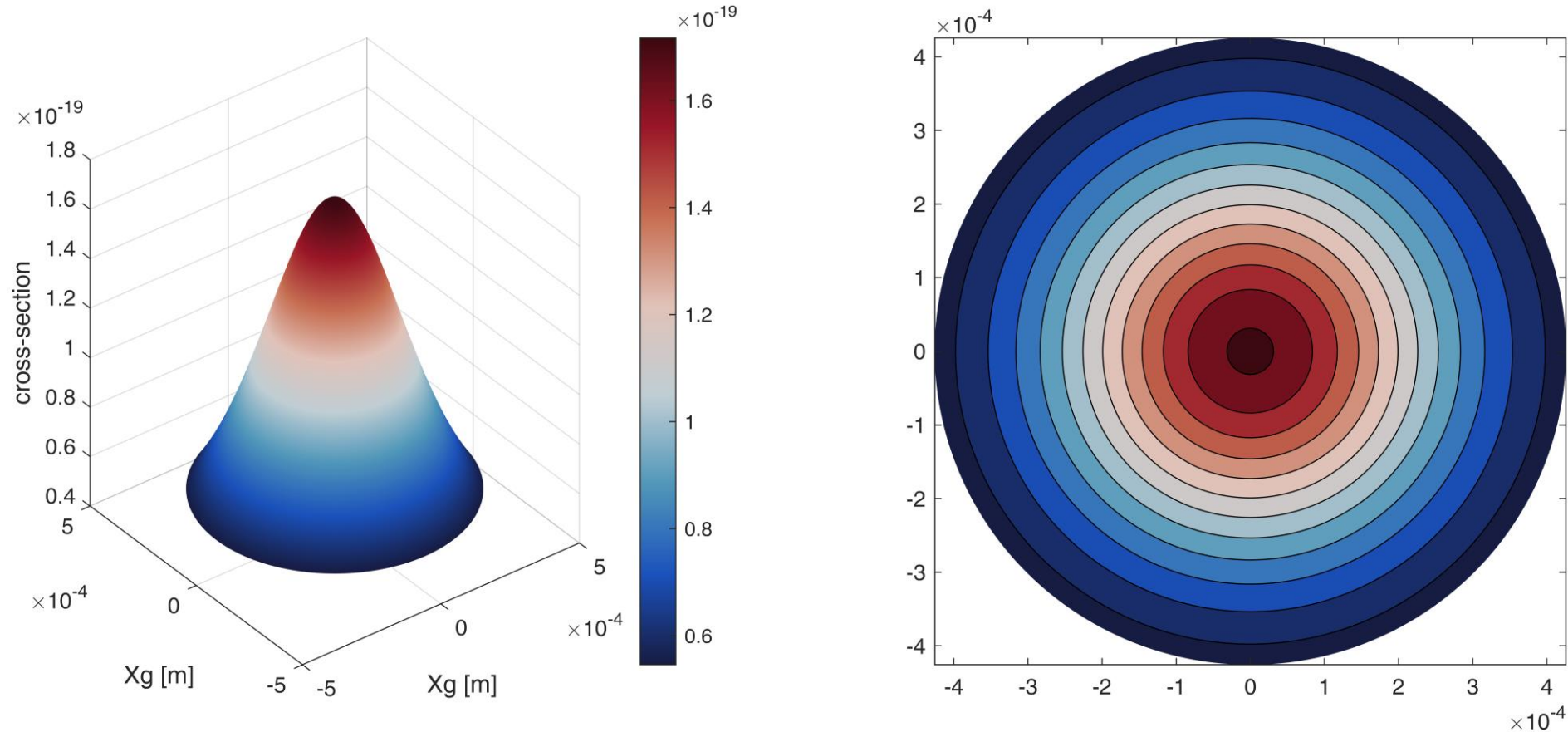


Figure. The computed spatial photons of Compton gamma-ray photons by a head-on a circularly polarized 532 nm laser with an 120 GeV electron beam.

Compton photons

➤ Samples

$$X_\gamma = L_1 \theta_\gamma \cos \varphi$$

$$Y_\gamma = L_1 \theta_\gamma \sin \varphi$$

➤ Fit function

$$d\sigma_0 = \frac{r_e^2}{\kappa^2(1+u)^3} \left(\kappa(1+(1+u)^2) - 4\frac{u}{\kappa}(1+u)(\kappa-u) \left[1 - \xi_\perp \cos(2(\varphi - \varphi_\perp)) \right] \right) du d\varphi,$$

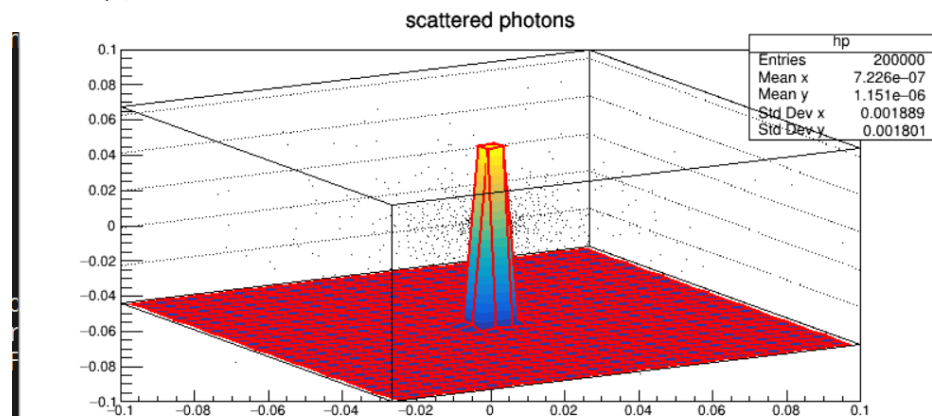
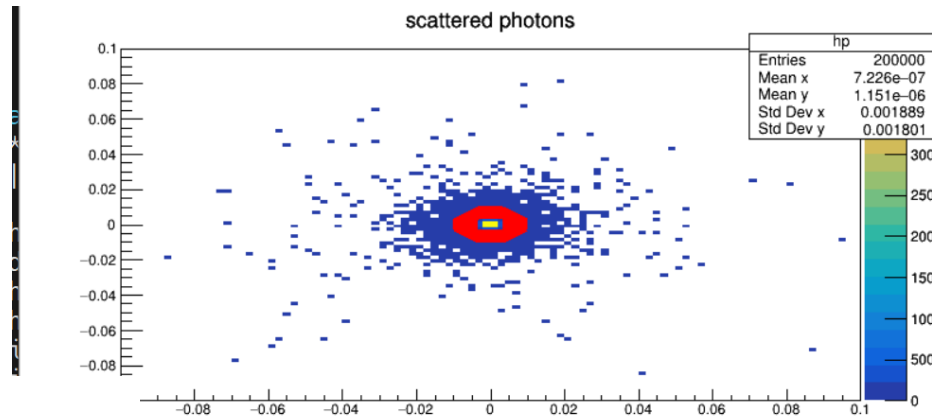
$$\frac{d\sigma_0}{dX_\gamma dY_\gamma} = \text{Jacobian} \cdot \frac{d\sigma_0}{du d\varphi}$$

$$X_\gamma = X - X_0$$

$$Y_\gamma = Y - Y_0$$

拟合参数

➤ Fit results



1	#normalization	1.49214e+26	3.41415e+23	7.83505e+19	7.38839e-29
2	position x	-3.90393e-07	1.77336e-06	-8.16127e-09	-3.78509e-02
3	position y	-2.44924e-07	1.01461e-06	-4.20812e-08	-1.97616e-01

同步辐射计算

Synchrotron radiation background

更实用的是弯铁中每单位转角中的光子数,即光通量为

$$\begin{aligned}\frac{dN_{ph}}{d\psi} &= \frac{1}{\hbar\omega} \frac{d^2P}{d\omega d\psi} \\ &= \frac{1}{\hbar\omega} \frac{1}{2\pi} \frac{dP}{d\omega} \Delta\omega \\ &= \frac{P_r}{2\pi\hbar\omega_c} \cdot \frac{\Delta\omega}{\omega} \cdot S(\omega/\omega_c) \quad (\text{II.73})\end{aligned}$$

上式中的 ψ 为水平偏转角,绕储存环一圈为 2π , $\hbar = \frac{h}{2\pi}$ 为简约化的普朗克常数。把 $P_r = \frac{2}{3} r_e m_0 c^2 \frac{c\gamma^4}{\rho^2}$ 和 $\hbar\omega_c = \frac{3}{2} \hbar c \frac{\gamma^3}{\rho}$ 代入式(II.73), 则

$$\frac{dN_{ph}}{d\psi} = \frac{4\alpha}{9} \gamma f_{rev} \frac{\Delta\omega}{\omega} S(\omega/\omega_c) \quad (\text{II.74})$$

式(II.74)中的 $\alpha = \frac{e^2}{\hbar c}$ 为精细结构常数,在实用单位制中 $\alpha = 7.29735 \times 10^{-3}$, $f_{rev} = \frac{c}{2\pi\rho}$ 是电子的回旋频率。式(II.74)表示一个电子在 $d\psi$ 角度内产生的光子数,如果有 N_e 个电子在储存环中回旋,它产生的光子通量为式(II.74)的 N_e 倍,那末一个电流为 I ($I = f_{rev} e N_e$) 的循环束流,产生的光子数则为^[48]

$$\begin{aligned}\frac{dN_{ph}}{d\psi} &= \frac{4\alpha}{9} \frac{E}{m_0 c^2} \frac{I}{e} \frac{\Delta\omega}{\omega} S(\omega/\omega_c) \\ &= C_\psi I E \frac{\Delta\omega}{\omega} S(\omega/\omega_c) \quad (\text{I.75})\end{aligned}$$

这就是光通量公式(7.13),它是单位时间、单位水平偏转角内辐射的光子数。此处 I 为储存环中的循环束流,单位为安培; E 为电子的能量,单位是 GeV; $\frac{\Delta\omega}{\omega}$ 是带宽,通常取 0.1%; $S(\omega/\omega_c)$ 是光谱函数,数值见表 7.1; C_ψ 是一个常数

$$C_\psi = \frac{4\alpha}{9m_0 c^2 e} = 3.96139 \times 10^{16} \frac{\text{光子数}}{\text{秒} \cdot \text{毫弧度} \cdot \text{安培} \cdot \text{GeV}}$$

$$N_{ph} = 3.3416e+13 [\text{光子数}/(\text{秒} \cdot \text{毫弧度} \cdot \text{安培} \cdot \text{GeV})]$$

Synchrotron radiation background

回旋频率： $f_{eq} = 3000$

束团数目： $N = 242$

单发束团一次的SR flux: $N_{ph} = 4.6028e+07$ [光子数/(秒·毫弧度·安培·GeV)]

Detector acceptance angle = $2/\gamma = 0.0085$ mrad



$N_{ph} = 391238$

Synchrotron radiation flux

$$\begin{aligned}\frac{dN_{\text{ph}}}{d\psi} &= \frac{4\alpha}{9} \frac{E}{m_0 c^2} \frac{I}{e} \frac{\Delta\omega}{\omega} S(\omega/\omega_c) \\ &= C_\psi I E \frac{\Delta\omega}{\omega} S(\omega/\omega_c)\end{aligned}$$

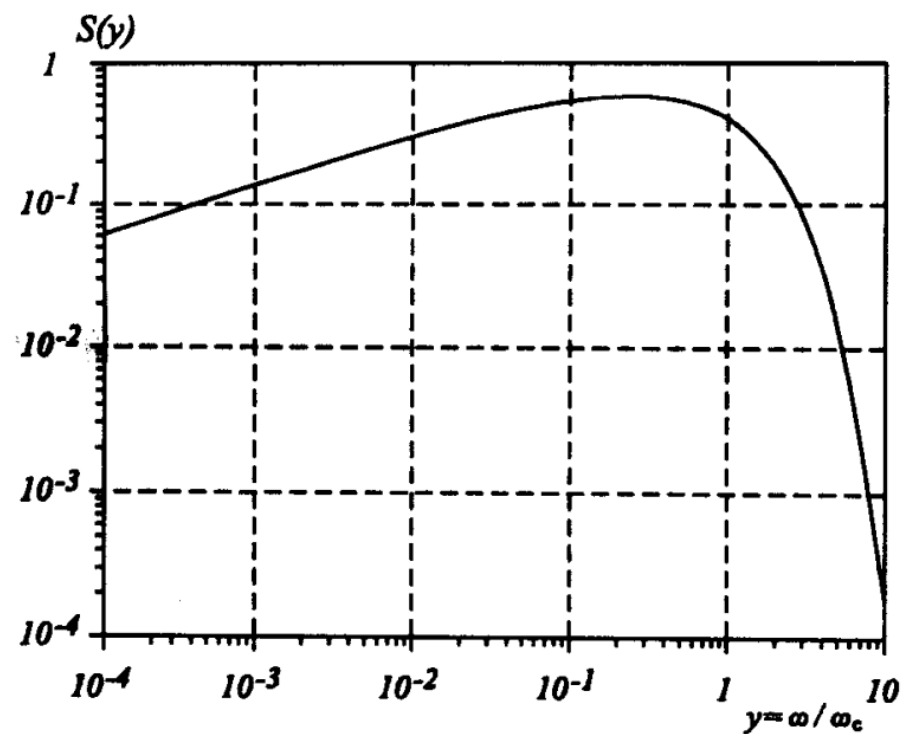


图 7.1 同步辐射光谱的谱适函数 $S(y)$ 曲线

Back-up