

# CP-symmetry studies of baryon-antibaryon pairs with the BESIII experiment

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EDP Seminar, June 9, 2022

with Karin Schönning and Andrzej Kupsc

Introduction

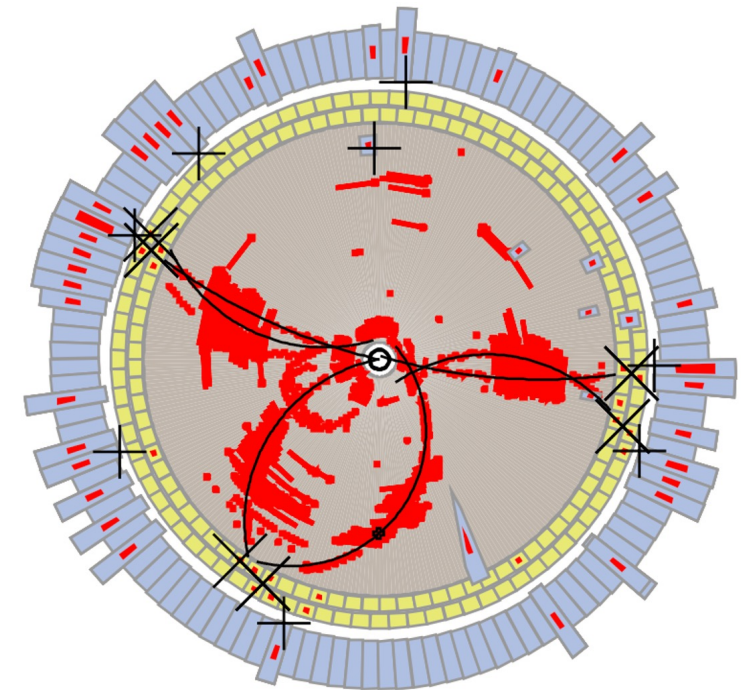
BESIII Experiment

Results Single Weak Decays

**New Result: *First Weak Phase Measurement in Baryon Decays***

Future prospects

Summary and Outlook



Display of simulated  
 $e^-e^+ \rightarrow \Lambda\pi^-\bar{\Lambda}\pi^+ \rightarrow p\pi^-\pi^-\bar{p}\pi^+\pi^+$



Universe likely began with equal abundance of matter and antimatter

Matter and antimatter are created and annihilate pair wise

Yet almost no existing antimatter in Universe

Strong observational evidence against symmetric Universe

Initial condition? Not likely. Standard Model washes out asymmetry

Local fraction of anti-stars over normal stars  
in galactic disk to be  $< 2.5 \times 10^{-6}$  at 95% CL

\*A. D. Sakharov, J. Exp. Theor. Phys. Lett. 5, 24

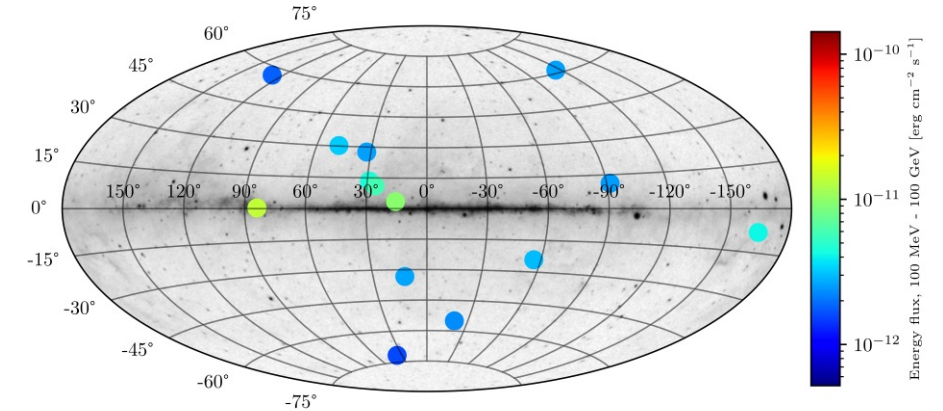


FIG. 1. Positions and energy flux in the 100 MeV–100 GeV range of antistar candidates selected in 4FGL-DR2. Galactic coordinates. The background image shows the *Fermi* 5-year all-sky photon counts above 1 GeV (image credit: NASA/DOE/*Fermi* LAT Collaboration).

*Fourteen anti-star candidates in Fermi 10 year data  
PRD 103, 083016(2021)*



The dynamical mechanism, *baryogenesis*, not understood

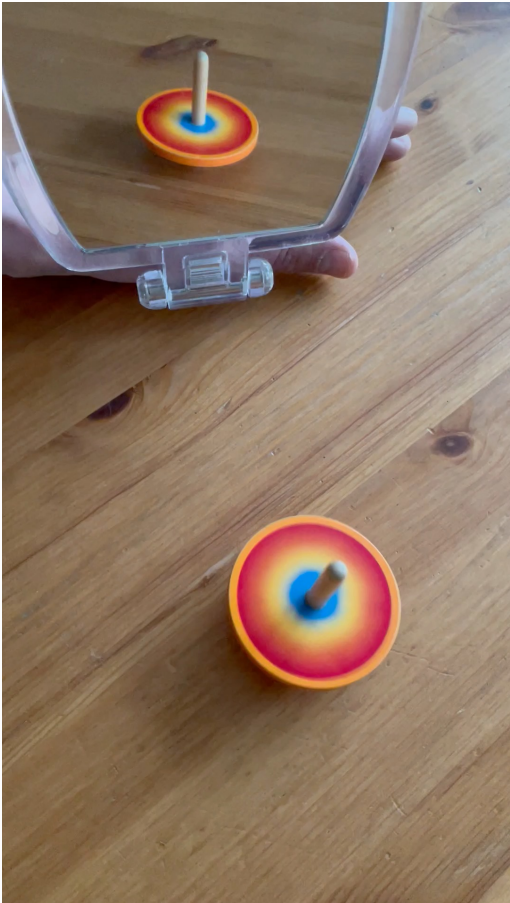
Sakharov's three criteria \*

- 1) Baryon number violation
- 2) Charge,  $C$ , and Charge conjugation Parity,  $P$ , violating processes
- 3) Departure from thermal equilibrium



CP violation is subtle effect requiring precision studies of many particle physics processes using complementary methods

\*A. D. Sakharov, *J. Exp. Theor. Phys. Lett.* 5, 24



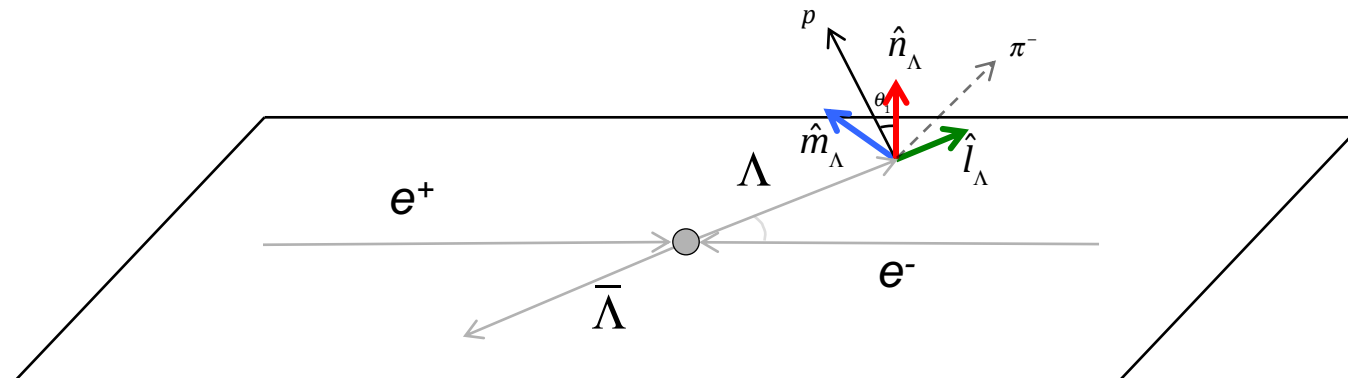
No CP violation detected for *baryons*

Additional degree of freedom for baryons compared to mesons : spin

Spin behaves differently compared to momentum when inverting spatial coordinates

Studying baryons provides complementary path in understanding CP symmetry

Focus on *hyperons*, strange quark systems and heavier cousins of proton and neutron



# Strangeness $\Delta S = 1$ mesons

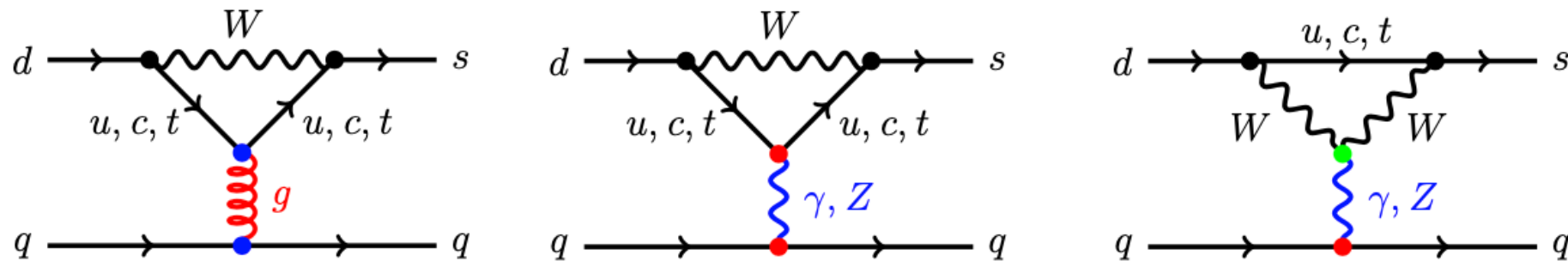
In strange sector most precise probe is  $\Delta S = 1$  direct CPV ( $\varepsilon'$ ) relative to indirect CPV ( $\varepsilon$ ) in  $K_{S,L} \rightarrow \pi\pi$  decays

CPV mechanism in SM requires penguin diagrams involving all three quark families

$$(\varepsilon'/\varepsilon)_{EXP} = (16.6 \pm 2.3) \times 10^{-4} *$$

$$(\varepsilon'/\varepsilon)_{SM} = (17.4 \pm 6.1) \times 10^{-4} + (\varepsilon'/\varepsilon)_{BSM} = (-4 - +10) \times 10^{-4} **$$

SM calculation involves partial cancellation of QCD and EW penguins which posed challenge until recently



QCD (left) and EW penguin diagrams (middle, right)\*\*\*

\* Phys. Lett. B544 (2002) 97–112; 0909.2555 [hep-ex]

\*\* Eur. Phys. J. C 80 (2020) 8, 705

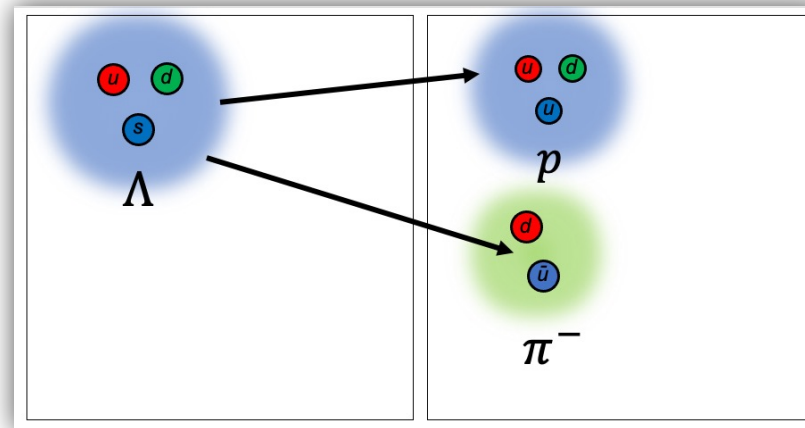
\*\*\* arXiv: 2203.03035

# Strangeness $\Delta S = 1$ baryons

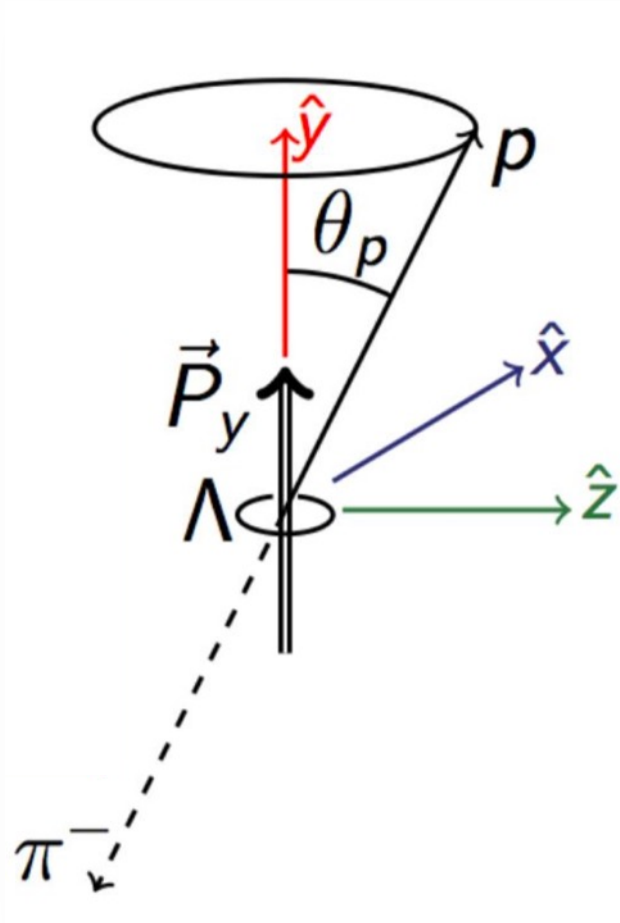
For  $K_{S,L} \rightarrow \pi\pi$  the direct CPV  $\varepsilon'$  comes from  $|\Delta I| = 1/2$  and  $|\Delta I| = 3/2$

Hyperon non-leptonic two-body weak decays tests  $\Delta S = 1$  CP

Hyperon weak decays offer complementary approach to meson decays as mainly  $|\Delta I| = 1/2$  transition occur



# Asymmetry parameters and Polarization



Polarization of hyperons experimentally accessible in weak parity violating decays

They are *self analysing*: daughter particles are emitted according to polarization of mother hyperon

Example: Angular distribution of  $\Lambda \rightarrow p\pi^-$

$$I(\cos \theta_p) \propto 1 + \alpha P_\Lambda \cos \theta_p$$

Asymmetry parameter  
CP-observable

Polarization





# Asymmetry parameters and Polarization

Asymmetry parameters describe relationship between the amplitudes  $S$  (parity violating) and  $P$  (parity conserving)

$$\alpha = 2\text{Re}(A_{S^*} \cdot A_P), \quad \beta = 2\text{Im}(A_{S^*} \cdot A_P) \quad \gamma = |A_S|^2 - |A_P|^2$$

$$\beta = \sqrt{(1 - \alpha^2)} \sin(\phi_Y), \quad \gamma = \sqrt{(1 - \alpha^2)} \cos(\phi_Y)$$

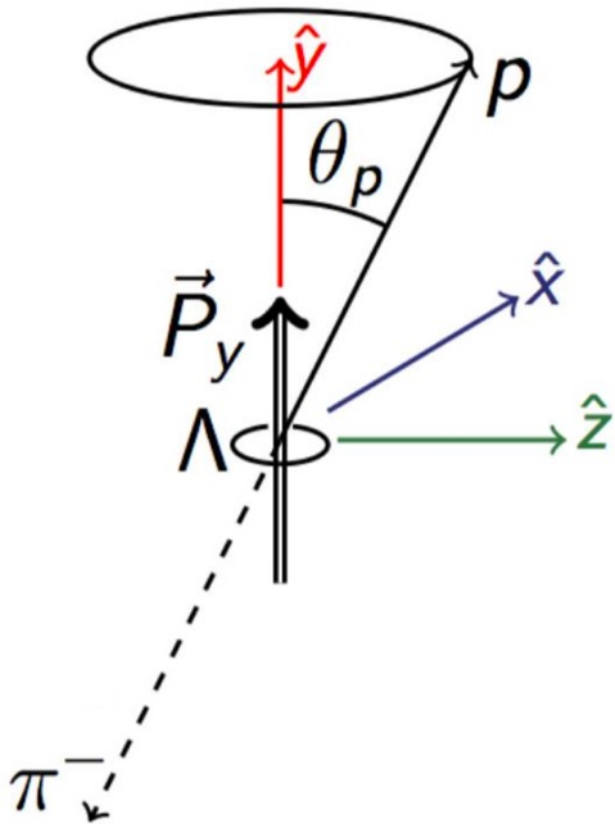
$\alpha$  and  $\phi_Y$  are determined experimentally

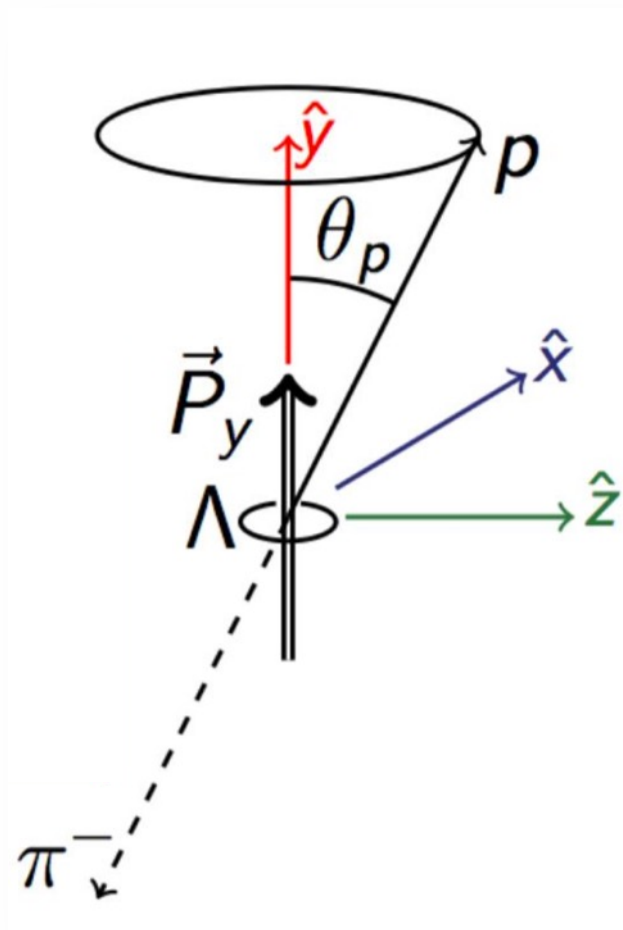
$\phi_Y$  more challenging: polarimeter or via sequential weak decays, e.g.  $\Xi$

$\alpha$  and  $\phi_Y$  are CP-odd observables

$$-1 \leq \alpha \leq 1$$

$$-\pi \leq \phi \leq \pi$$





We determine  $\alpha$  from  $\Lambda \rightarrow p\pi^-$  and  $\bar{\alpha}$  from  $\bar{\Lambda} \rightarrow \bar{p}\pi^+$

If CP symmetry holds then  $\alpha = -\bar{\alpha}$

$$A_{CP} = \frac{\alpha + \bar{\alpha}}{\alpha - \bar{\alpha}}$$

This test not limited only to  $\Lambda \rightarrow p\pi^-$  but all non-leptonic two-body weak decays

$$\begin{aligned}
 S &= |S| e^{i\delta_S} e^{i\xi_S} & P &= |P| e^{i\delta_P} e^{i\xi_P} \\
 \bar{S} &= -|\bar{S}| e^{i\delta_S} e^{-i\xi_S} & \bar{P} &= |\bar{P}| e^{i\delta_P} e^{-i\xi_P}
 \end{aligned}$$

$\delta$  strong baryon pion phase shift at cm energy of hyperon mass

$\xi$  weak CP-odd phase for  $\Delta I = 1/2$

Under assumption that isospin  $1/2$  transitions dominate

$$A_{CP} = \frac{\alpha + \bar{\alpha}}{\alpha - \bar{\alpha}} \approx -\tan(\delta_P - \delta_S) \tan(\xi_P - \xi_S)^*$$

strong phase diff      weak phase diff  
non-zero iff FSI

Corrections from isospin  $3/2$  are only a few percent, hence sufficiently good approximation

\* Phys. Rev Lett 55 162 (1985)

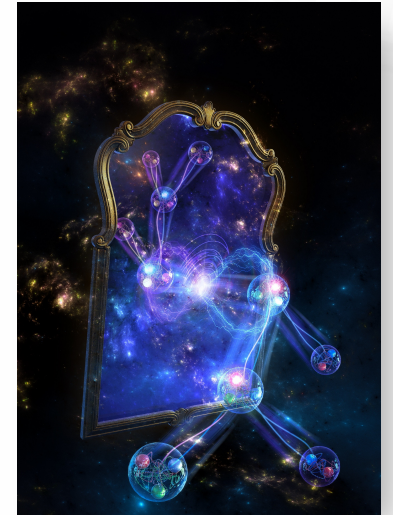
$$\Xi^- \rightarrow \Lambda \pi^-, \Lambda \rightarrow p \pi^-$$

$$A_{CP}^{\Xi} = \frac{\alpha_{\Xi} + \bar{\alpha}_{\Xi}}{\alpha_{\Xi} - \bar{\alpha}_{\Xi}} \approx -\sin\langle\phi_{\Xi}\rangle \frac{\sqrt{1-\alpha_{\Xi}^2}}{\alpha_{\Xi}} \tan(\zeta_P - \zeta_S)_{\Xi} *$$

$$\Delta\phi_{CP} = \frac{\phi_{\Xi} + \bar{\phi}_{\Xi}}{2} \approx \cos\langle\phi_{\Xi}\rangle \frac{\alpha_{\Xi}}{\sqrt{1-\alpha_{\Xi}^2}} \tan(\zeta_P - \zeta_S)_{\Xi} *$$

strong contribution  $\phi_{\Xi} \approx 0$  weak phase diff - potentially CPV

$\Delta\phi_{CP}$  more sensitive to CP-violating effects of  $A_{CP}^{\Xi} *$



\* Phys. Rev Lett 55 162 (1985)

# Strangeness $\Delta S = 1$ BSM

$$K_{S,L} \rightarrow \pi\pi$$

$$(\epsilon'/\epsilon)_{EXP} = (16.6 \pm 2.3) \times 10^{-4} *$$

$$(\epsilon'/\epsilon)_{SM} = (17.4 \pm 6.1) \times 10^{-4} + (\epsilon'/\epsilon)_{BSM} = (-4 - +10) \times 10^{-4} **$$

$$Y \rightarrow B\pi$$

$$(\xi_P - \xi_S)_{BSM} = \frac{C'_B}{B_G} \left( \frac{\epsilon'}{\epsilon} \right)_{BSM} + \frac{C_B}{\kappa} \epsilon_{BSM}$$

Wolfenstein parameters  $\eta \lambda^5 A^2$   
 $1.36(7) \times 10^{-4}$

***	$\xi_P - \xi_S$		$C_B$	$C'_B$
	$(\eta \lambda^5 A^2)$	$[10^{-4} \text{ rad}]$		
	SM		BSM	
$\Lambda \rightarrow p\pi^-$	$-0.1 \pm 1.5$	$-0.2 \pm 2.2$	$0.9 \pm 1.8$	$0.4 \pm 0.9$
$\Xi^- \rightarrow \Lambda\pi^-$	$-1.5 \pm 1.2$	$-2.1 \pm 1.7$	$-0.5 \pm 1.0$	$0.4 \pm 0.7$

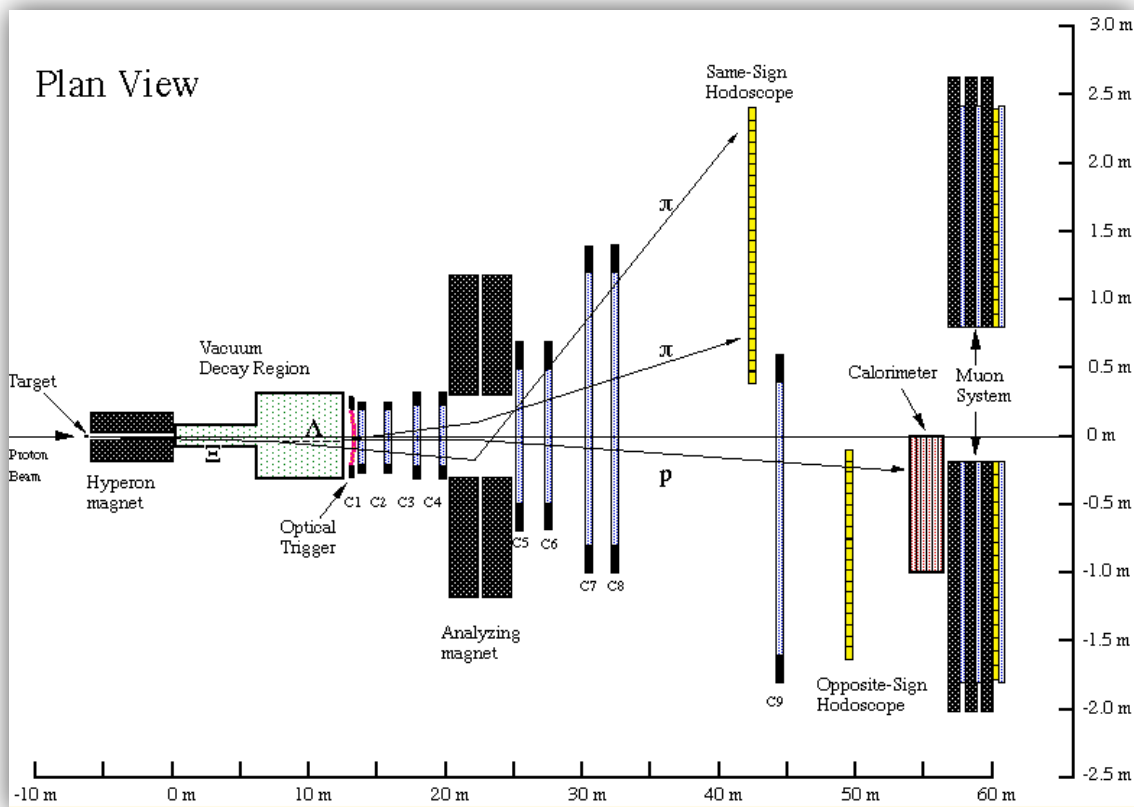
Phys. Rev. D 67, 056001 (2003)  
 Phys. Rev. D 69, 076008 (2004)



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\* Phys. Lett. B544 (2002) 97–112; 0909.2555 [hep-ex]  
 \*\* Eur. Phys. J. C 80 (2020) 8, 705  
 \*\*\* arXiv: 2203.03035

Chromomagnetic BSM penguin operators



HyperCP (E871) Fermilab

800 GeV/c proton on fixed target Cu

Best  $A_{CP}$  limits obtained so far 117M  $\Xi$  and 41M  $\bar{\Xi}$

$$A_{CP \Xi \Lambda} = \frac{\alpha_{\Xi} \alpha_{\Lambda} + \bar{\alpha}_{\Xi} \bar{\alpha}_{\Lambda}}{\alpha_{\Xi} \alpha_{\Lambda} - \bar{\alpha}_{\Xi} \bar{\alpha}_{\Lambda}} = (0.0 \pm 5.1 \pm 4.7) \times 10^{-4} *$$

$$|A_{SM \Xi \Lambda}| \leq 5 \times 10^{-5} **$$

144M polarized  $\Xi$  (~5%)  $\phi_{\Xi, \text{HyperCP}} = -0.042 \pm 0.011 \pm 0.011$

\*PRL 93, 262001 (2004)

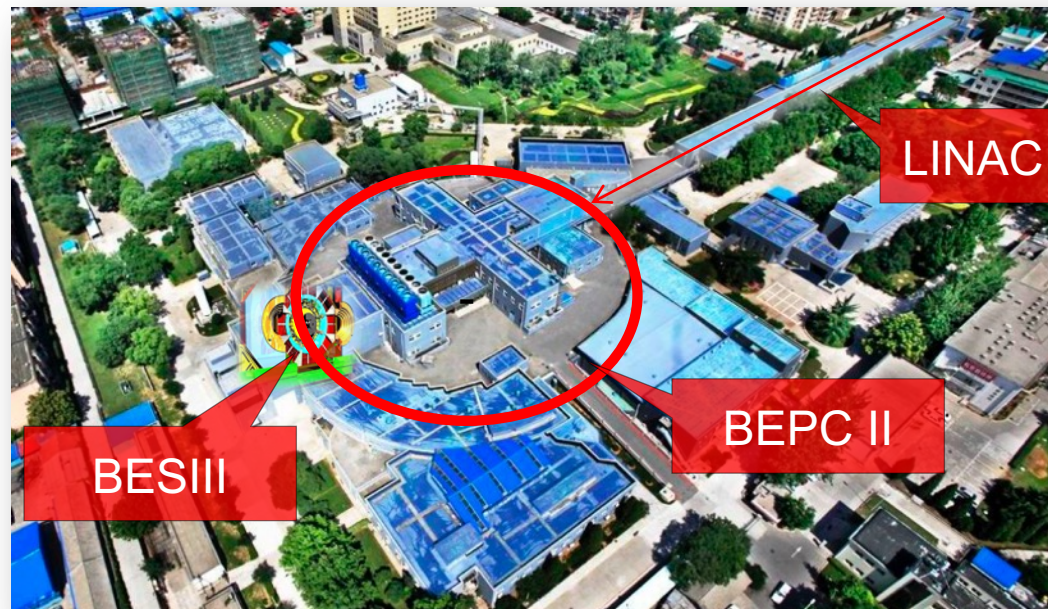
\*\* PRD 67, 056001 (2003)

\*\*\* NPB, Proc Suppl 187, 208 (2009)

$$862M \Xi \text{ \& } 230M \bar{\Xi} \quad A_{CP \Xi \Lambda} = \frac{\alpha_{\Xi} \alpha_{\Lambda} + \bar{\alpha}_{\Xi} \bar{\alpha}_{\Lambda}}{\alpha_{\Xi} \alpha_{\Lambda} - \bar{\alpha}_{\Xi} \bar{\alpha}_{\Lambda}} = (-6.0 \pm 2.1 \pm 2.0) \times 10^{-4} ***$$



# Beijing Electron Positron Collider BEPC II



*Aerial view of BEPC II and BESIII*

$e^+e^-$  collider in CMS range 2.0 – 4.95 GeV

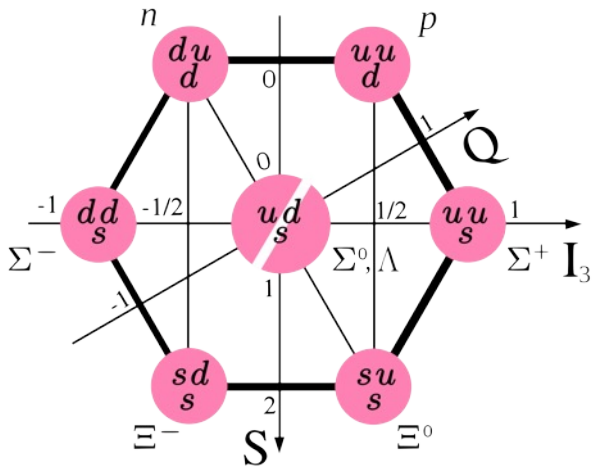
Optimized in  $\tau$  - charm region

Data taking since 2009, peak luminosity  $10^{33} \text{ cm}^{-2}\text{s}^{-1}$

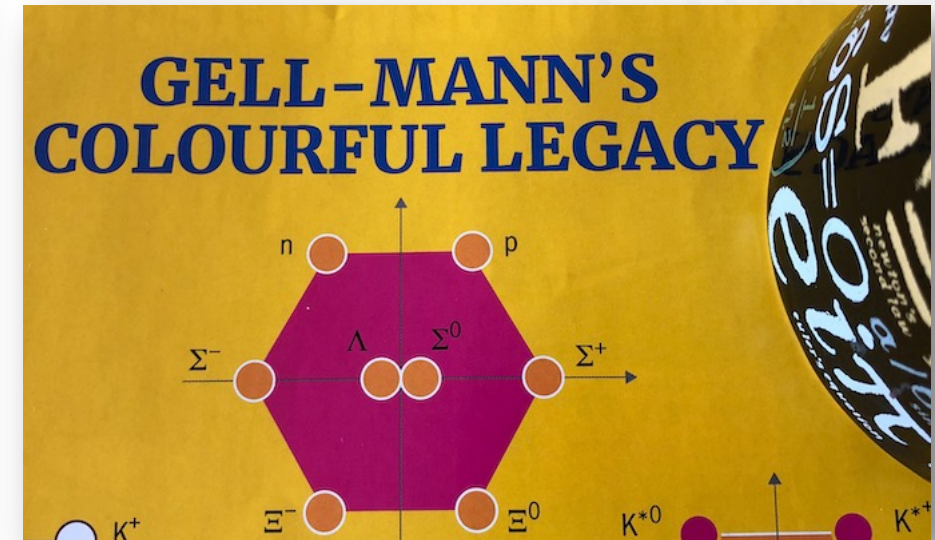


# Baryon Octet at BESIII

From CERN Courier cover July-August 2019



hyperon	Mass [GeV/c <sup>2</sup> ]	$c\tau$ [cm]	decay (BF)
$\Lambda(uds)$	1.116	7.9	$p\pi^-$ (63.9%) $n\pi^0$ (35.8%)
$\Sigma^-(dds)$	1.197	4.4	$n\pi^-$ (99.8%)
$\Sigma^+(uus)$	1.189	2.4	$p\pi^0$ (51.6%) $n\pi^+$ (48.3%)
$\Xi^0(uss)$	1.315	8.7	$\Lambda\pi^0$ (99.5%)
$\Xi^-(dss)$	1.321	5.1	$\Lambda\pi^-$ (99.8%)



**Thresholds:**

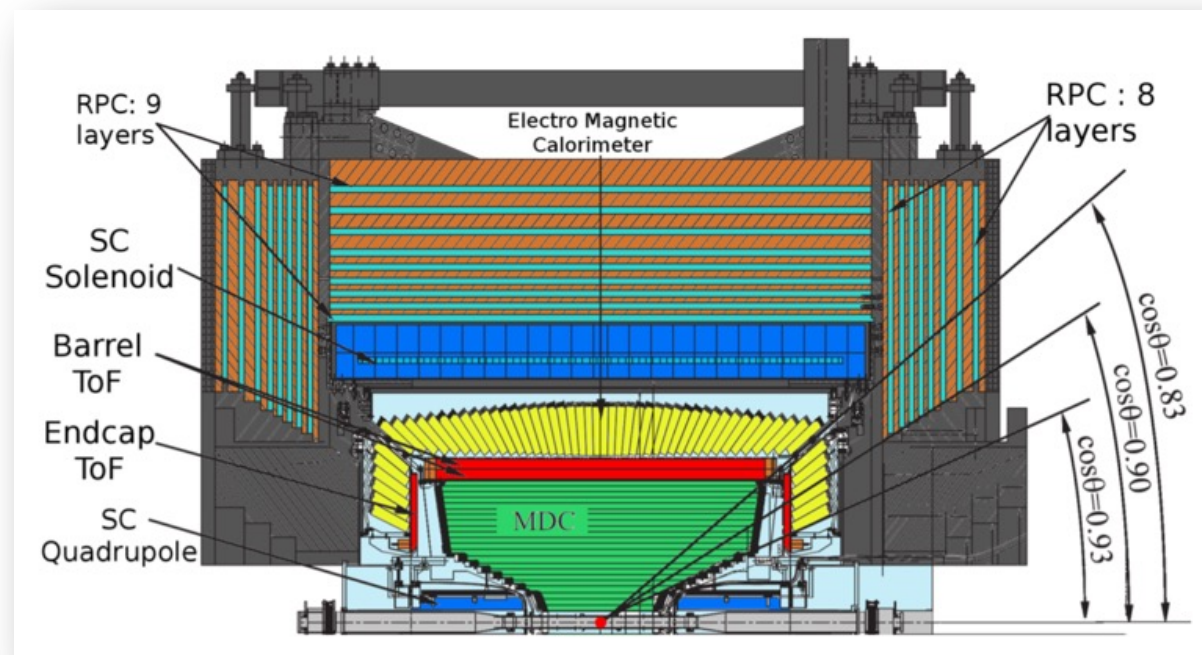
$\Sigma^+\bar{\Sigma}^-$ 2.379 GeV	$\Lambda\bar{\Lambda}$ 2.231 GeV	$\Sigma^-\bar{\Sigma}^+$ 2.395 GeV
$\Xi^0\bar{\Xi}^0$ 2.630 GeV	$\Sigma^0\bar{\Sigma}^0$ 2.385 GeV	$\Xi^-\bar{\Xi}^+$ 2.643 GeV

Full baryon octet kinematically accessible at  $J/\psi$  resonance





# BESII Hyperons



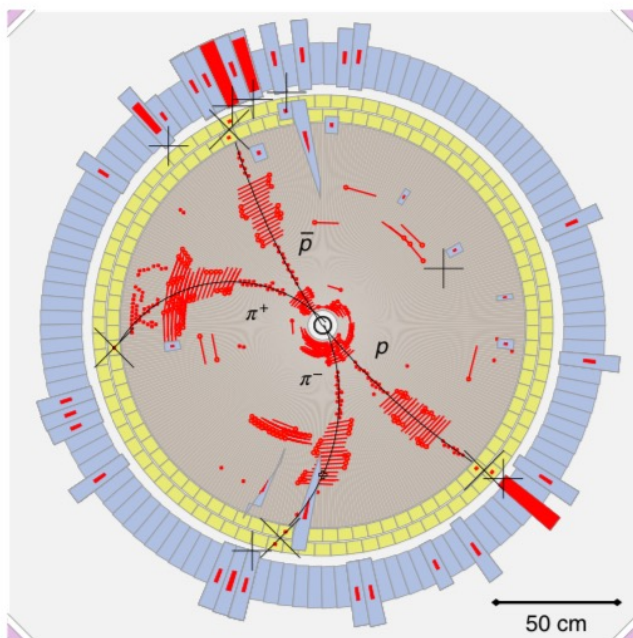
Multipurpose detector with very good resolution, near  $4\pi$  angular coverage

Symmetric particle – anti-particle conditions and produced in a quantum entangled state

$e^+e^-$  experiment low hadronic background and controlled systematic uncertainties

World's largest charmonia data sample collected  $N_{J/\psi} = (1.0087 \pm 0.0044) \times 10^{10}$





**Fig. 2 | An example  $J/\psi \rightarrow (\Lambda \rightarrow p\pi^-)(\bar{\Lambda} \rightarrow \bar{p}\pi^+)$  event in the BESIII detector.** Cross-section of the detector in the plane perpendicular to the colliding electron-positron beams and a schematic representation of the information collected for the event. The mean decay length of the neutral  $\Lambda(\bar{\Lambda})$  is 5 cm. The curved tracks of the charged particles from the subsequent  $\Lambda(\bar{\Lambda})$  decays are registered in the drift chamber, indicated by the brown region of the display. The momenta of (anti-)baryons are greater than  $750 \text{ MeV } c^{-1}$  and pions are less than  $300 \text{ MeV } c^{-1}$ .

***BESIII, Nature Physics 15 (2019) 631***

Charged track coverage  $|\cos\theta| < 0.93$

Mom. res of charged tracks 0.5% at 1 GeV/c

Neutrals  $|\cos\theta| < 0.8$  and  $0.86 < |\cos\theta| < 0.92$

Energy resolution 2.5% (5%) at 1 GeV for  
barrell (end cap)

ToF can be used together with  $dE/dx$  MDC for PID

But for fully charged modes e.g.  $\Lambda$  and  $\Xi$  momentum requirements enough to separate protons from pions



Production parameters of spin 1/2 baryons at cbar : angular distribution parameter  $\alpha_\psi$  and relative phase  $\Delta\Phi$

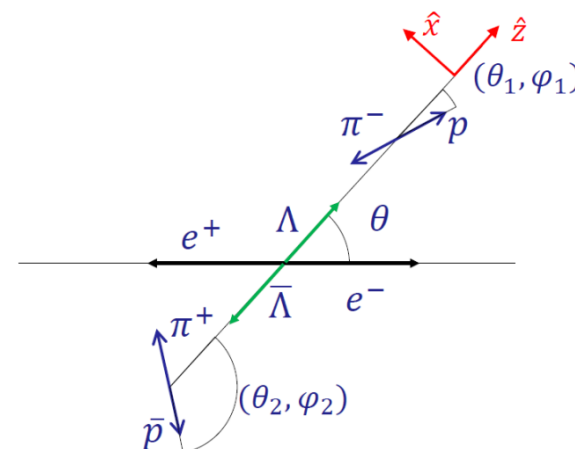
Decay parameters for 2-body decays:  $\alpha$  and  $\bar{\alpha}$

$\mathcal{T}_0 - \mathcal{T}_6$  are functions with experimentally measured observables

$$W(\xi) = \underbrace{\mathcal{T}_0(\xi) + \alpha_\psi \mathcal{T}_5(\xi)}_{\text{Unpolarised part}} - \underbrace{\alpha \bar{\alpha} [\mathcal{T}_1(\xi) + \sqrt{1 - \alpha_\psi^2} \cos(\Delta\Phi) \mathcal{T}_2(\xi) + \alpha_\psi \mathcal{T}_6(\xi)]}_{\text{Polarised part}} + \underbrace{\sqrt{1 - \alpha_\psi^2} \sin(\Delta\Phi) [\alpha \mathcal{T}_3(\xi) - \bar{\alpha} \mathcal{T}_4(\xi)]}_{\text{Spin correlated part}}$$

Polarization necessary to "disentangle"  $\alpha$  from  $\bar{\alpha}$

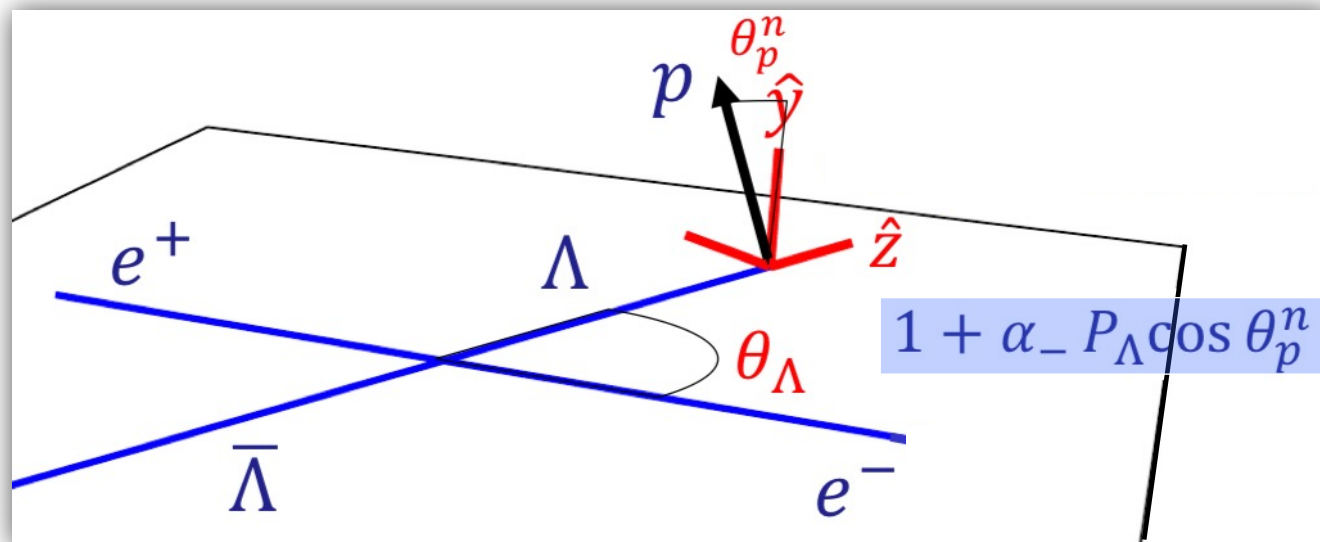
$$\begin{aligned} \mathcal{T}_0(\xi) &= 1 \\ \mathcal{T}_1(\xi) &= \sin^2 \theta \sin \theta_1 \sin \theta_2 \cos \phi_1 \cos \phi_2 + \cos^2 \theta \cos \theta_1 \cos \theta_2 \\ \mathcal{T}_2(\xi) &= \sin \theta \cos \theta (\sin \theta_1 \cos \theta_2 \cos \phi_1 + \cos \theta_1 \sin \theta_2 \cos \phi_2) \\ \mathcal{T}_3(\xi) &= \sin \theta \cos \theta \sin \theta_1 \sin \phi_1 \\ \mathcal{T}_4(\xi) &= \sin \theta \cos \theta \sin \theta_2 \sin \phi_2 \\ \mathcal{T}_5(\xi) &= \cos^2 \theta \\ \mathcal{T}_6(\xi) &= \cos \theta_1 \cos \theta_2 - \sin^2 \theta \sin \theta_1 \sin \theta_2 \sin \phi_1 \sin \phi_2 \end{aligned}$$



When initial state is unpolarized and process is parity conserving, hyperons polarized perpendicular to production plane

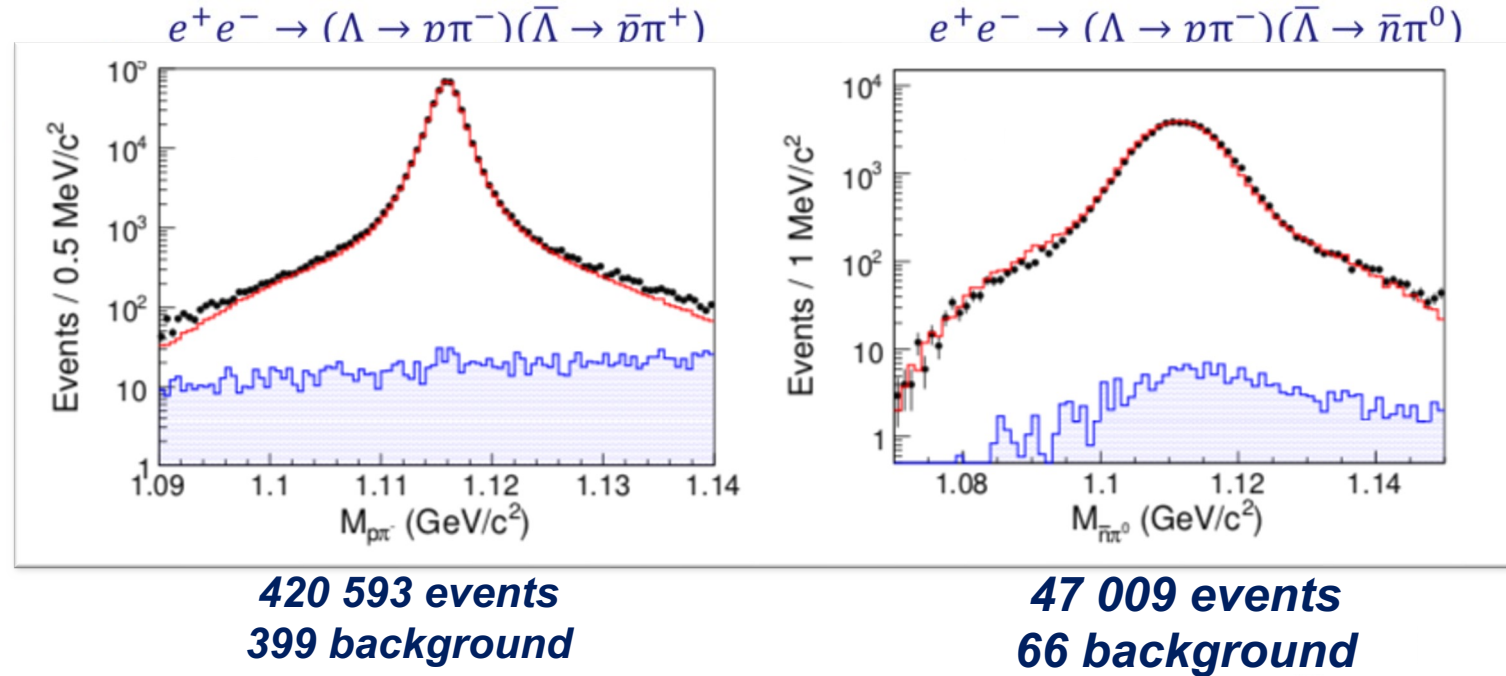
Phase is production related, depending on CMS energy and scattering angle  $\Delta\Phi \neq 0$  from interfering amplitudes (e.g. s- and d- waves)  $\Delta\Phi = 0$  threshold

Analyticity requires that SL FF  $\sim$  TL FF as  $|q^2|$  approaches  $\infty$   $\Delta\Phi = 0$



$$e^+e^- \rightarrow J/\psi \rightarrow \Lambda\bar{\Lambda} \rightarrow p\pi^-\bar{p}\pi^+(n\pi^0)$$

BESIII, Nature Physics 15 (2019) 631



Results based on  $1.3 \times 10^9$  J/ψ events

Excellent agreement between data and simulation

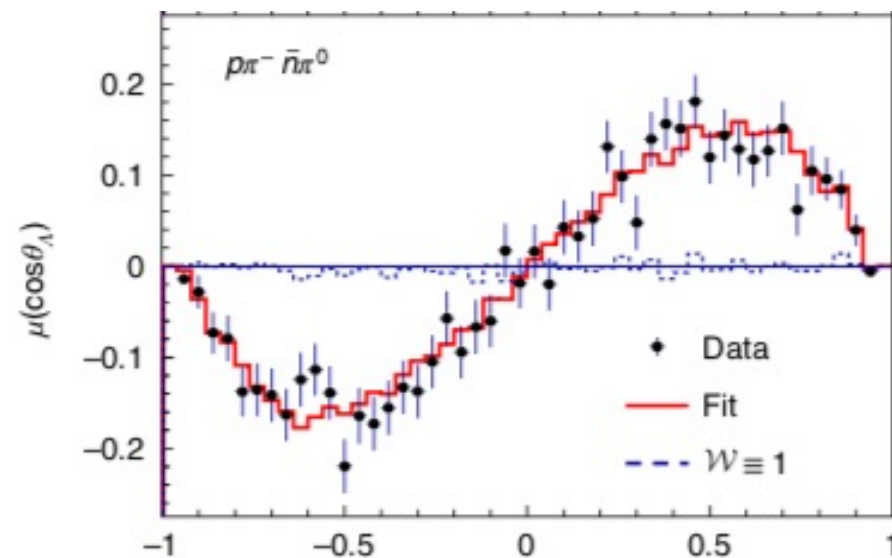
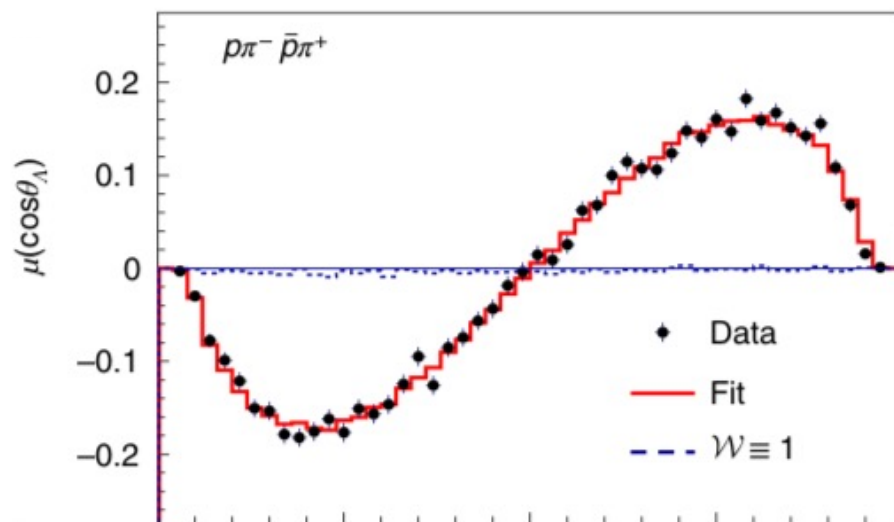
Background levels at fraction of percent

Unbinned maximum log-likelihood approach used for estimation of parameters



$$e^+e^- \rightarrow J/\psi \rightarrow \Lambda\bar{\Lambda} \rightarrow \rho\pi^-\bar{\rho}\pi^+ (n\pi^0)$$

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$$\bar{P}_Y(\cos\theta_\Lambda) = \frac{\sqrt{1 - \alpha_\psi^2} \cos\theta_\Lambda \sin\theta_\Lambda}{1 + \alpha_\psi \cos^2\theta_\Lambda} \sin(\Delta\Phi)$$

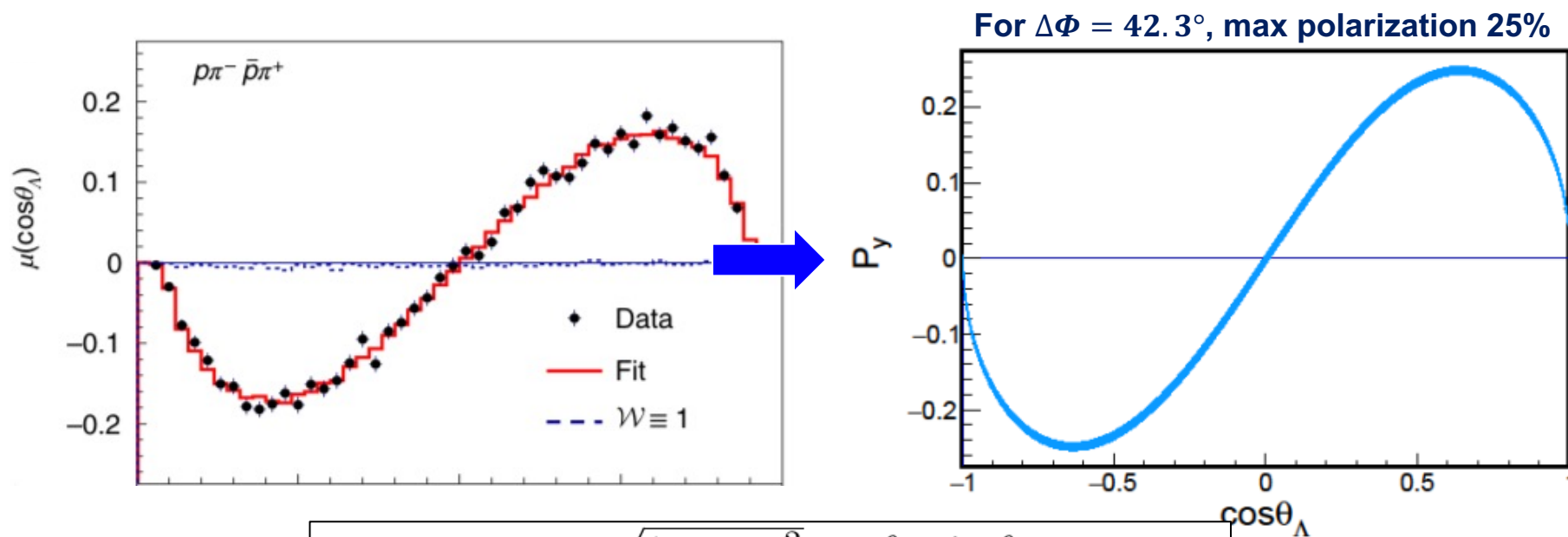
First measurement of hyperon polarization at  $J/\psi$  resonance

Non-zero  $\Delta\Phi$  allows for direct and precise measurements of asymmetry parameters





BESIII, Nature Physics 15 (2019) 631



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## BESIII, Nature Physics 15 (2019) 631

## 2019 Review of Particle Physics.

M. Tanabashi *et al.* (Particle Data Group), Phys. Rev. D **98**, 030001 (2018) and 2019 update. $\Lambda$  DECAY PARAMETERS

See the "Note on Baryon Decay Parameters" in the neutron Listings. Some early results have been omitted.

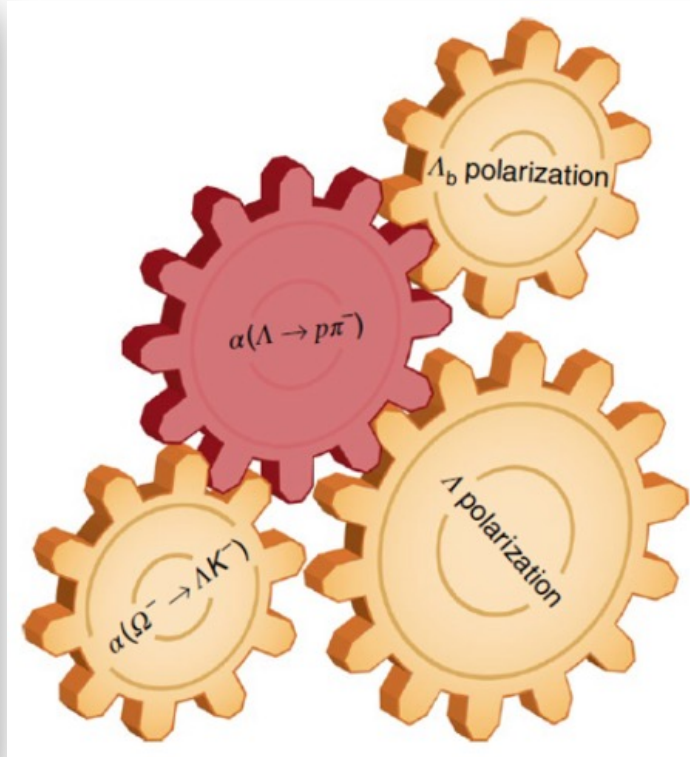
 $\alpha_{\Lambda \rightarrow p\pi^-}$ 

INSPIRE search

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
$0.750 \pm 0.009 \pm 0.004$	420k	ABLIKIM 2018AG	BES3	$J/\psi$ to $\Lambda\bar{\Lambda}$
... We do not use the following data for averages, fits, limits, etc. ...				
$0.584 \pm 0.046$	8500	ASTBURY 1975	SPEC	
$0.649 \pm 0.023$	10325	CLELAND 1972	OSPK	
$0.67 \pm 0.06$	3520	DAUBER 1969	HBC	From $\Xi$ decay
$0.645 \pm 0.017$	10130	OVERSETH 1967	OSPK	$\Lambda$ from $\pi^-p$
$0.62 \pm 0.07$	1156	CRONIN 1963	CNTR	$\Lambda$ from $\pi^-p$

## References:

ABLIKIM 2018AG	arXiv:1808.08917	
ASTBURY 1975	NP B99 30	Measurement of the Differential Cross Section and the Spin Correlation Parameters $P$ , $A$ , and $R$ in the Backward Peak of $\pi^- p \rightarrow K^0 \Lambda$ at 5 GeV/c
CLELAND 1972	NP B40 221	A Measurement of the $\beta$ -Parameter in the Charged Nonleptonic Decay of the $\Lambda^0$ Hyperon
DAUBER 1969	PR 179 1262	Production and Decay of Cascade Hyperons
OVERSETH 1967	PRL 19 391	Time Reversal Invariance in $\Lambda$ Decay
CRONIN 1963	PR 129 1795	Measurement of the Decay Parameters of the $\Lambda$ Particle



## news &amp; views

## PARTICLE PHYSICS

## Anomalous asymmetry

A measurement based on quantum entanglement of the parameter describing the asymmetry of the  $\Lambda$  hyperon decay is inconsistent with the current world average. This shows that relying on previous measurements can be hazardous.

Ulrik Egede

*"Relying on previous measurements can be hazardous"*

$$\alpha(\Lambda \rightarrow p\pi^-) = 0.750(9)_{\text{stat}}(4)_{\text{syst}}$$

$17 \pm 3$  % off from old  $\alpha_{\text{PDG}}(\Lambda \rightarrow p\pi^-) = 0.642(13)$  established in 1978. Sets the new PDG standard

Consequences for all results relying on 2018 PDG  $\alpha_{\text{PDG}}$ . Measurements re-scaled or re-measured





2019 Review of Particle Physics.  
M. Tanabashi *et al.* (Particle Data Group), Phys. Rev. D **98**, 030001 (2018) and 2019 update.

**$\Lambda$  DECAY PARAMETERS**

See the "Note on Baryon Decay Parameters" in the neutron Listings. Some early results have been omitted.

$\alpha_-$  FOR  $\Lambda \rightarrow p\pi^-$  INSPIRE search

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ASTBURY 1975 NP B99 30 Measurement of the Differential Cross Section and the Spin Correlation Parameters  $P$ ,  $A$ , and  $R$  in the Backward Peak of  $\pi^- p \rightarrow K^0 \Lambda$  at 5 GeV/c

CLELAND 1972 NP B40 221 A Measurement of the  $\beta$ -Parameter in the Charged Nonleptonic Decay of the  $\Lambda^0$  Hyperon

DAUBER 1969 PR 179 1262 Production and Decay of Cascade Hyperons

OVERSETH 1967 PRL 19 391 Time Reversal Invariance in  $\Lambda$  Decay

CRONIN 1963 PR 129 1795 Measurement of the Decay Parameters of the  $\Lambda$  Particle

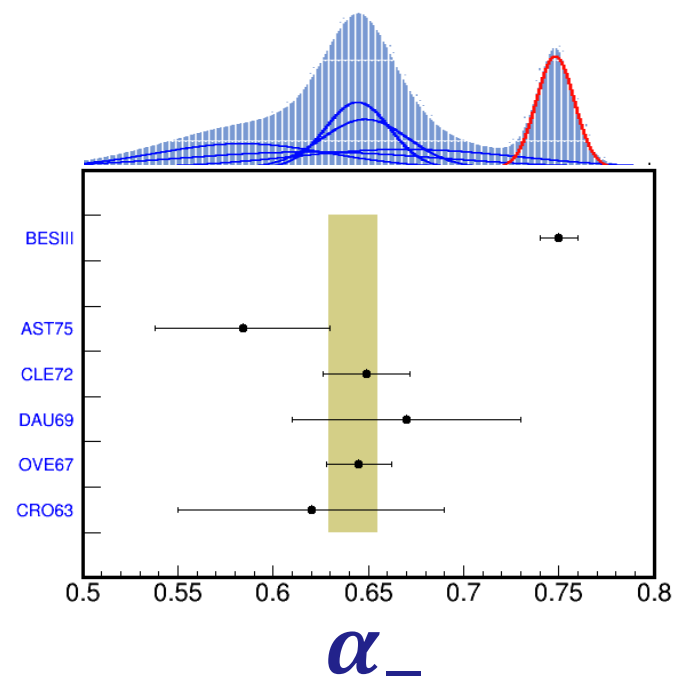


Fig. A. Kupsc

Re-measurement  $\alpha(\Lambda \rightarrow p\pi^-) = 0.721(6)_{\text{stat}}(5)_{\text{syst}}^{**}$

Even this result not compatible;  $\langle \alpha(\Lambda \rightarrow p\pi^-) \rangle_{\text{BESIII}} = 0.754(3)(2)$ , assuming  $\alpha_- = -\alpha_+$ ,  $\rho = 0.82$

$\alpha_{\Lambda, \text{PDG}} = 0.732 \pm 0.014_{\text{tot}}$  based on the two mutually incompatible values

More input was needed!

\* BESIII, Nature Physics 15 (2019) 631

\*\* Phys. Rev. Lett. 123 (2019) 18, 182301

BESIII, Nature Physics 15 (2019) 631

$$A_{CP, \Lambda} = \frac{\alpha_{\Lambda^+} + \bar{\alpha}_{\Lambda^-}}{\alpha_{\Lambda^-} - \bar{\alpha}_{\Lambda^+}} = -0.006 \pm 0.012_{stat} \pm 0.007_{syst}$$

$$-3 \times 10^{-5} \leq A_{\Lambda SM} \leq 4 \times 10^{-5}^*$$

$$A_{CP, \Lambda prev} = 0.013 \pm 0.021_{tot}^{**}$$

Most precise test of CP for  $\Lambda$  and compatible with SM expectations

Table 1 | Summary of the results

Parameters	This work	Previous results
$\alpha_{\psi}$	$0.461 \pm 0.006 \pm 0.007$	$0.469 \pm 0.027$ (ref. <sup>14</sup> )
$\Delta\Phi$	$42.4 \pm 0.6 \pm 0.5^\circ$	-
$\alpha_-$	$0.750 \pm 0.009 \pm 0.004$	$0.642 \pm 0.013$ (ref. <sup>6</sup> )
$\alpha_+$	$-0.758 \pm 0.010 \pm 0.007$	$-0.71 \pm 0.08$ (ref. <sup>6</sup> )
$\bar{\alpha}_0$	$-0.692 \pm 0.016 \pm 0.006$	-
$A_{CP}$	$-0.006 \pm 0.012 \pm 0.007$	$0.006 \pm 0.021$ (ref. <sup>6</sup> )
$\bar{\alpha}_0/\alpha_+$	$0.913 \pm 0.028 \pm 0.012$	-

Parameters:  $J/\psi \rightarrow \Lambda \bar{\Lambda}$  angular distribution parameter  $\alpha_{\psi}$ , helicity phase  $\Delta\Phi$ , asymmetry parameters for the  $\Lambda \rightarrow p\pi^-$  ( $\alpha_-$ ),  $\bar{\Lambda} \rightarrow \bar{p}\pi^+$  ( $\alpha_+$ ) and  $\bar{\Lambda} \rightarrow \bar{n}\pi^0$  ( $\bar{\alpha}_0$ ) decays, CP asymmetry  $A_{CP}$  and ratio  $\bar{\alpha}_0/\alpha_+$ . The first uncertainty is 1 s.d. statistical, and the second is systematic, calculated as described in the Methods.

\*. Phys. Rev. D67, 056001 (2003)

\*\* Phys Rev C54, 1877 (1996)



Results based on 83k events

Parameter	Measured value
$\alpha_{J/\psi}$	$-0.508 \pm 0.006 \pm 0.004$
$\Delta\Phi_{J/\psi}$	$-0.270 \pm 0.012 \pm 0.009$
$\alpha_{\psi'}$	$0.682 \pm 0.03 \pm 0.011$
$\Delta\Phi_{\psi'}$	$0.379 \pm 0.07 \pm 0.014$
$\alpha_0$	$-0.998 \pm 0.037 \pm 0.009$
$\bar{\alpha}_0$	$0.990 \pm 0.037 \pm 0.011$

$$\langle \alpha_\Sigma \rangle = -0.994 \pm 0.004 \pm 0.002$$

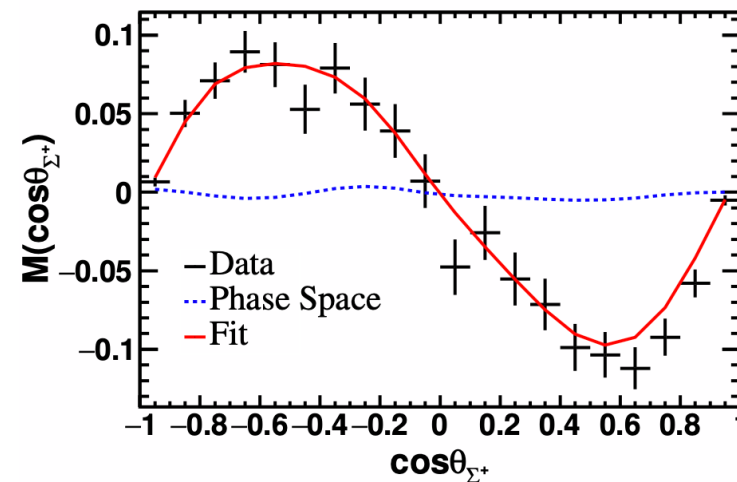
First CP measurement for any  $\Sigma$  decay

$$A_{CP\Sigma} = \frac{\alpha_\Sigma + \alpha_{\bar{\Sigma}}}{\alpha_\Sigma - \alpha_{\bar{\Sigma}}} = -0.004 \pm 0.037_{stat} \pm 0.010_{syst} *$$

$$A_{CP\Sigma SM} 3.6 \times 10^{-6} **$$

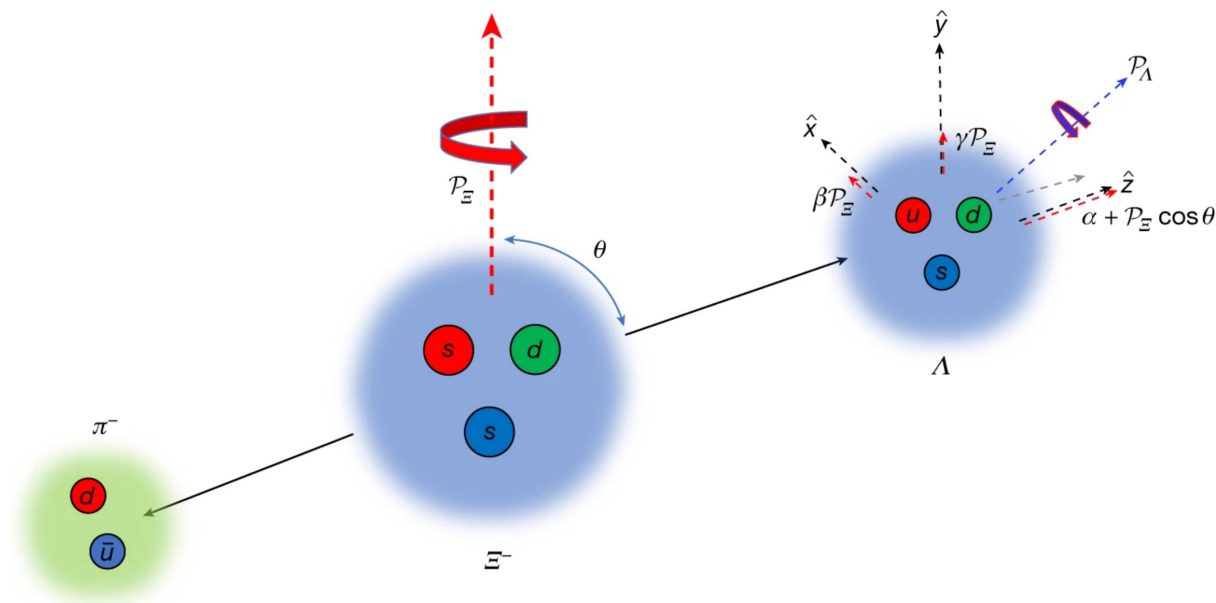
\* Phys.Rev.Lett. 125 (2020) 5, 052004

\*\* Phys. Rev. D67, 056001 (2003)



$$J/\psi \rightarrow \Xi^- \bar{\Xi}^+ \rightarrow \Lambda \pi^- \bar{\Lambda} \pi^+$$

Nature 606 64-69 (2022)



*Longitudinal polarization  $\alpha$  and  $P$  ; transversal polarization  $\phi$*

**Probing CP symmetry and weak phases with entangled double-strange baryons**



$$A_{CP}^E = \frac{\alpha_E + \bar{\alpha}_E}{\alpha_E - \bar{\alpha}_E} \approx -\sin\langle\phi_E\rangle \frac{\sqrt{1-\alpha_E^2}}{\alpha_E} \tan(\xi_P - \xi_S)_E *$$

$$\Delta\phi_{CP} = \frac{\phi_E + \bar{\phi}_E}{2} \approx \cos\langle\phi_E\rangle \frac{\alpha_E}{\sqrt{1-\alpha_E^2}} \tan(\xi_P - \xi_S)_E *$$

strong contribution  $\phi_E \approx 0$       weak phase diff - potentially CPV

$\Delta\phi_{CP}$  more sensitive to CP-violating effects of  $A_{CP}^E$ .  
Proposed more 35 years ago but not measured until now!

\* Phys. Rev Lett 55 162 (1985)

- The formalism exploits polarisation, entanglement and sequential decays \* \*\*



$$\mathcal{W}(\xi; \omega) = \sum_{\mu, \nu=0}^3 \textcircled{C_{\mu\nu}} \sum_{\mu', \nu'=0}^3 \textcircled{a_{\mu\mu'}^{\bar{E}} a_{\nu\nu'}^{\bar{E}} a_{\mu'0}^{\Lambda} a_{\nu'0}^{\bar{\Lambda}}}$$

- Nine-dimensional phase space given by nine helicity angles
- Eight free parameters determined by maximum log likelihood method:

$$\alpha_{\psi}, \Delta\Phi, \alpha_E, \bar{\alpha}_E, \phi_E, \bar{\phi}_E, \alpha_{\Lambda}, \bar{\alpha}_{\Lambda}$$

↑ ↑ ↑ ↑  
not measured before

Formalism developed in Uppsala

\* Phys. Rev. D 99, 056008 (2019)  
\*\* Phys. Rev. D 100, 114005 (2019)

- The formalism exploits polarisation, entanglement and sequential decays \* \*\*

$$\mathcal{W}(\xi; \omega) = \sum_{\mu, \nu=0}^3 \mathbb{C}_{\mu\nu} \sum_{\mu', \nu'=0}^3 a_{\mu\mu'}^{\Xi} a_{\nu\nu'}^{\Xi} a_{\mu'0}^{\Lambda} a_{\nu'0}^{\bar{\Lambda}}$$

$C_{\mu\nu}$  4x4 real matrix representing polarizations and spin correlations

Spin density matrices  $a_{\mu\nu}$

$$C_{\mu\nu} = (1 + \alpha_{\psi} \cos^2 \theta) \begin{pmatrix} 1 & 0 & P_y & 0 \\ 0 & C_{xx} & 0 & C_{xz} \\ -P_y & 0 & C_{yy} & 0 \\ 0 & -C_{xz} & 0 & C_{zz} \end{pmatrix}$$

$$\begin{pmatrix} 1 & 0 & 0 & \alpha_D \\ \alpha_D \sin \theta \cos \varphi & \gamma_D \cos \theta \cos \varphi - \beta_D \sin \varphi & -\beta_D \cos \theta \cos \varphi - \gamma_D \sin \varphi & \sin \theta \cos \varphi \\ \alpha_D \sin \theta \sin \varphi & \beta_D \cos \varphi + \gamma_D \cos \theta \sin \varphi & \gamma_D \cos \varphi - \beta_D \cos \theta \sin \varphi & \sin \theta \sin \varphi \\ \alpha_D \cos \theta & -\gamma_D \sin \theta & \beta_D \sin \theta & \cos \theta \end{pmatrix}$$

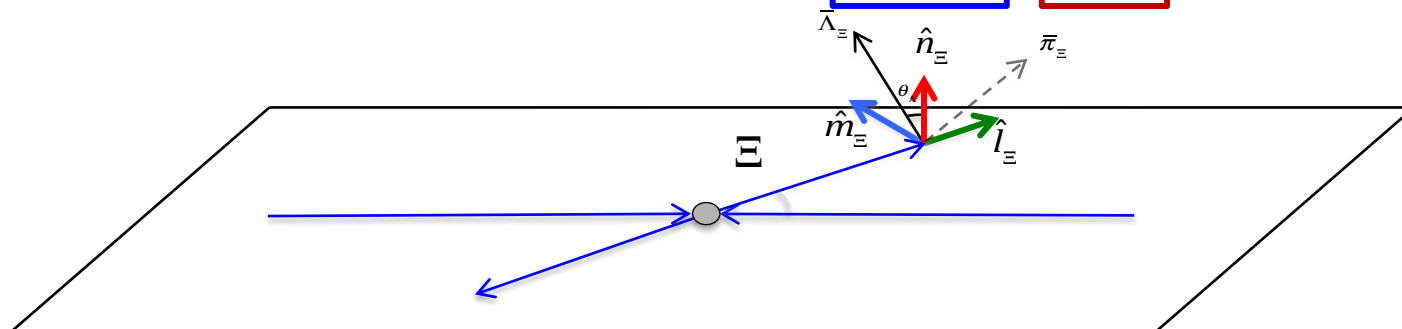
\* Phys. Rev. D 99, 056008 (2019)  
\*\* Phys. Rev. D 100, 114005 (2019)

# Formalism $J/\psi \rightarrow \Xi\bar{\Xi} \rightarrow \Lambda(\rightarrow p\pi)\bar{\Lambda}\pi(\rightarrow \bar{p}\pi^+)$

Here  $\Delta\Phi \neq 0$  is not needed to measure decay parameters! \*, \*\*

$$\Delta\Phi \neq 0 : \quad M = \begin{array}{|c|} \hline \Xi^- \bar{\Xi}^+ \\ \hline \end{array} \quad \begin{array}{|c|} \hline \Lambda\bar{\Lambda} \\ \hline \end{array} \quad (7)$$

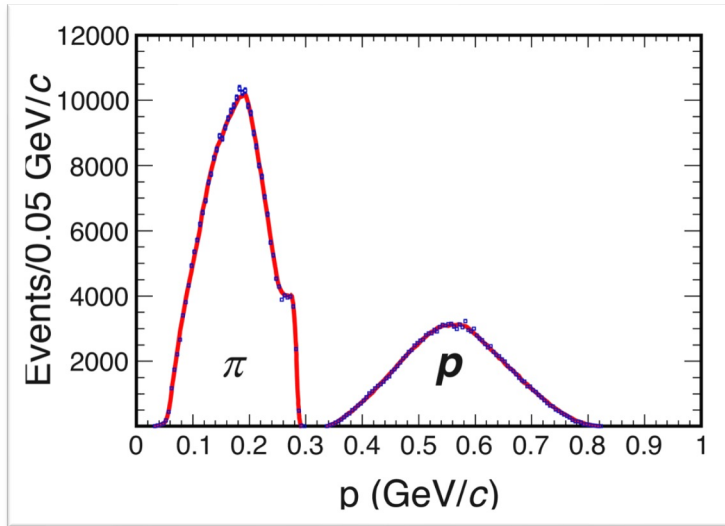
$$\Delta\Phi = 0 : \quad M = \begin{array}{|c|} \hline \Xi^- \bar{\Xi}^+ \\ \hline \end{array} \quad \begin{array}{|c|} \hline \Lambda\bar{\Lambda} \\ \hline \end{array} \quad (5)$$



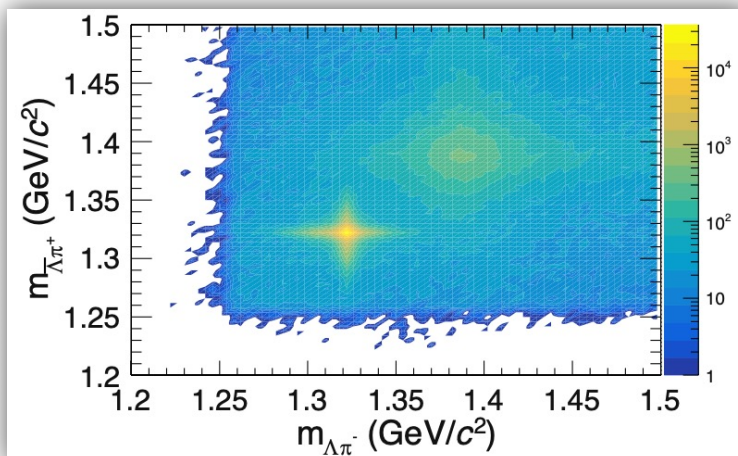
\* Phys. Rev. D 99, 056008 (2019)  
 \*\* Phys. Rev. D 100, 114005 (2019)



Nature 606 64-69 (2022)



at least one proton, one anti-proton, two positively and two negatively charged pion candidates  
 momentum criteria used to select proton ( $p > 0.32 \text{ GeV}/c$ ) and pion ( $p < 0.30 \text{ GeV}/c$ ) candidates  
 $\Lambda$  and  $\Xi$  candidates formed with succesful vertex fits



Mass windows  $|m(p\pi) - m_{\Lambda}| < 11.5 \text{ MeV}/c^2$  and  $|m(\Lambda\pi) - m_{\Xi}| < 12.0 \text{ MeV}/c^2$

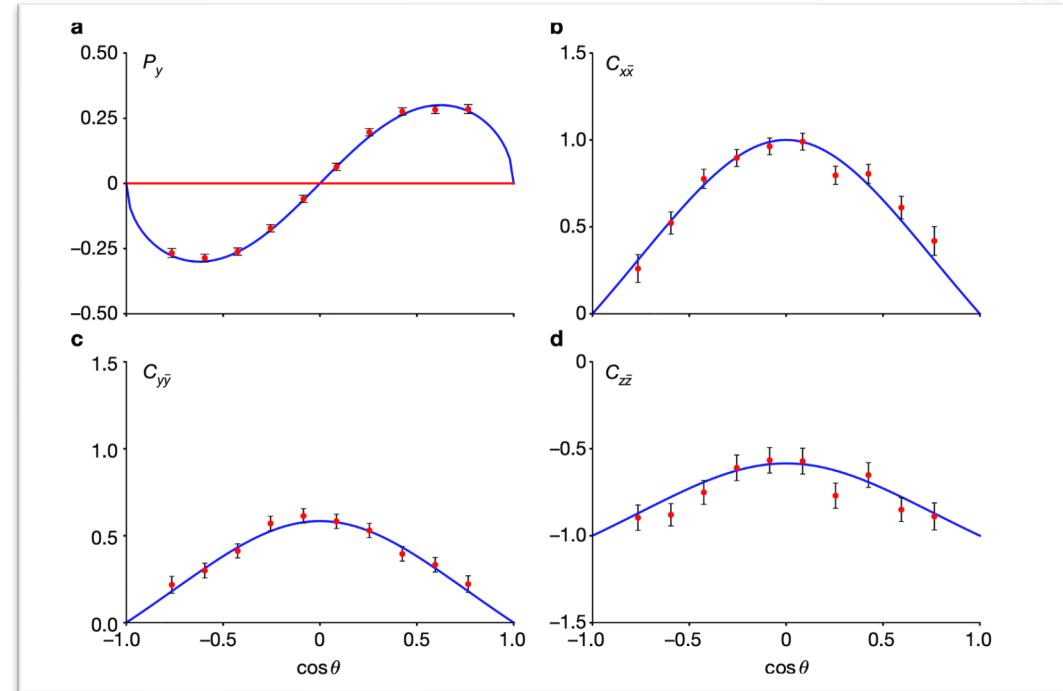
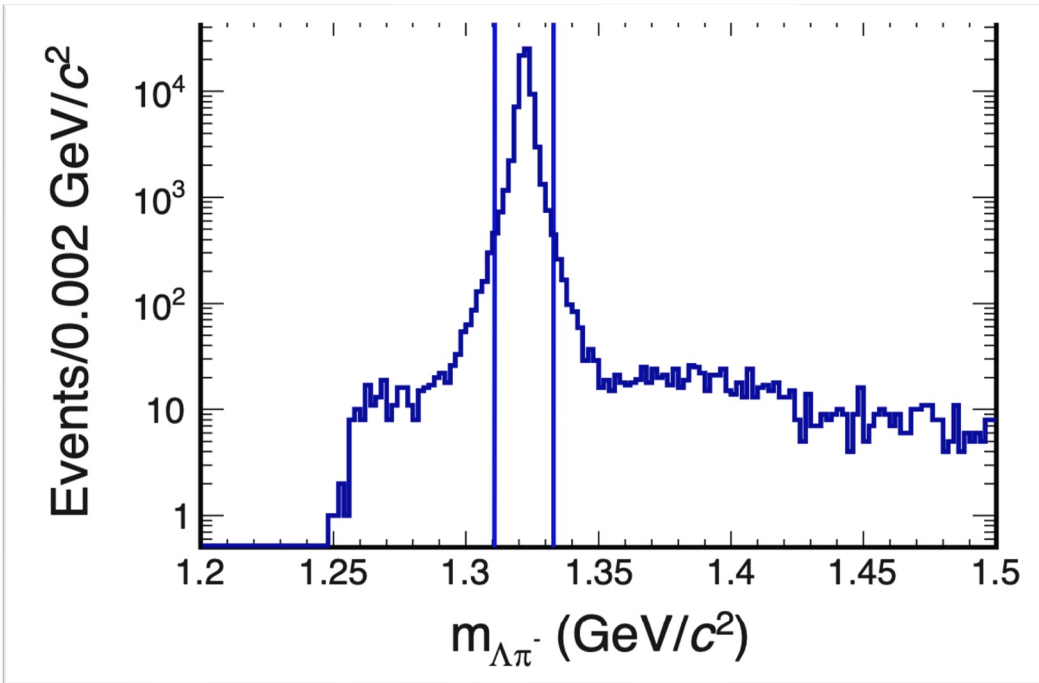
4C-kinematic fit on the hypothesis  $e^+e^- \rightarrow J/\Xi \rightarrow \Xi^-\bar{\Xi}^+$  is used as veto

The decay lengths of  $\Lambda$  and  $\Xi$  candidates greater than 0.

For improved data-MC consistency only events with  $|\cos\theta| < 0.84$



Nature 606 64-69 (2022)



73 200 exclusively measured  $\Xi^- \bar{\Xi}^+ \rightarrow \Lambda\pi^- \bar{\Lambda}\pi^+$  events

Very low level of background,  $199 \pm 17$  events

Here *entanglement* from spin correlations allows us to “disentangle” the weak and strong contributions



Table 1 | Summary of results

Parameter	This work	Previous result	Reference	
$\alpha_{\psi}$	$0.586 \pm 0.012 \pm 0.010$	$0.58 \pm 0.04 \pm 0.08$	Ref. <sup>49</sup>	*
$\Delta\Phi$	$1.213 \pm 0.046 \pm 0.016 \text{ rad}$	–		
$\alpha_{\Xi}$	$-0.376 \pm 0.007 \pm 0.003$	$-0.401 \pm 0.010$	Ref. <sup>26</sup>	**
$\phi_{\Xi}$	$0.011 \pm 0.019 \pm 0.009 \text{ rad}$	$-0.037 \pm 0.014 \text{ rad}$	Ref. <sup>26</sup>	**
$\bar{\alpha}_{\Xi}$	$0.371 \pm 0.007 \pm 0.002$	–		
$\bar{\phi}_{\Xi}$	$-0.021 \pm 0.019 \pm 0.007 \text{ rad}$	–		
$\alpha_{\Lambda}$	$0.757 \pm 0.011 \pm 0.008$	$0.750 \pm 0.009 \pm 0.004$	Ref. <sup>4</sup>	***
$\bar{\alpha}_{\Lambda}$	$-0.763 \pm 0.011 \pm 0.007$	$-0.758 \pm 0.010 \pm 0.007$	Ref. <sup>4</sup>	***
$\xi_p - \xi_s$	$(1.2 \pm 3.4 \pm 0.8) \times 10^{-2} \text{ rad}$	–		
$\delta_p - \delta_s$	$(-4.0 \pm 3.3 \pm 1.7) \times 10^{-2} \text{ rad}$	$(10.2 \pm 3.9) \times 10^{-2} \text{ rad}$	Ref. <sup>3</sup>	****
$A_{\text{CP}}^{\Xi}$	$(6 \pm 13 \pm 6) \times 10^{-3}$	–		
$\Delta\phi_{\text{CP}}^{\Xi}$	$(-5 \pm 14 \pm 3) \times 10^{-3} \text{ rad}$	–		
$A_{\text{CP}}^{\Lambda}$	$(-4 \pm 12 \pm 9) \times 10^{-3}$	$(-6 \pm 12 \pm 7) \times 10^{-3}$	Ref. <sup>4</sup>	***
$\langle\phi_{\Xi}\rangle$	$0.016 \pm 0.014 \pm 0.007 \text{ rad}$			

The  $J/\psi \rightarrow \Xi^- \bar{\Xi}^+$  angular distribution parameter  $\alpha_{\psi}$ , the hadronic form factor phase  $\Delta\Phi$ , the decay parameters for  $\Xi^- \rightarrow \Lambda \pi^-$  ( $\alpha_{\Xi}, \phi_{\Xi}$ ),  $\bar{\Xi}^+ \rightarrow \bar{\Lambda} \pi^+$  ( $\bar{\alpha}_{\Xi}, \bar{\phi}_{\Xi}$ ),  $\Lambda \rightarrow p \pi^-$  ( $\alpha_{\Lambda}$ ) and  $\bar{\Lambda} \rightarrow \bar{p} \pi^+$  ( $\bar{\alpha}_{\Lambda}$ ); the CP asymmetries  $A_{\text{CP}}^{\Xi}$ ,  $\Delta\phi_{\text{CP}}^{\Xi}$  and  $A_{\text{CP}}^{\Lambda}$ , and the average  $\langle\phi_{\Xi}\rangle$ . The first and second uncertainties are statistical and systematic, respectively.

First measurement of polarisation

First direct determination of all  $\Xi^- \bar{\Xi}^+$  decay parameters

Previous experiments determined product  $\alpha_{\Xi} \alpha_{\Lambda}$

Independent measurement of  $\Lambda$  decay parameters. Excellent agreement with previous BESIII results. Similar precision despite 6x smaller data sample

\* PRD 93, 072003 (2018)  
 \*\* PDG 2020  
 \*\*\* Nat. Ph. 15, 631 (2019)  
 \*\*\*\* PRL 93, 011802 (2004)



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$A_{\text{CP}}^{\Xi^-}$	$(6 \pm 13 \pm 6) \times 10^{-3}$	-		
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First extraction of weak phase diff  
for any weakly decaying baryon

$$(\xi_p - \xi_s) = (1.2 \pm 3.4 \pm 0.8) \times 10^{-2} \text{ rad}$$

Consistent with SM expectation

$$(\xi_p - \xi_s)_{\text{SM}} = (1.8 \pm 1.5) \times 10^{-4} \text{ rad}$$

New method for direct weak phase extraction!

Three CP-tests, (two independent)  
in single measurement

- \* PRD 93, 072003 (2018)
- \*\* PDG 2020
- \*\*\* Nat. Ph. 15, 631 (2019)
- \*\*\*\* PRL 93, 011802 (2004)



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$\delta_P - \delta_S$	$(-4.0 \pm 3.3 \pm 1.7) \times 10^{-2} \text{ rad}$	$(10.2 \pm 3.9) \times 10^{-2} \text{ rad}$	Ref. <sup>3</sup>	****
$A_{CP}^{\Xi^-}$	$(6 \pm 13 \pm 6) \times 10^{-3}$	-		
$\Delta\phi_{CP}^{\Xi^-}$	$(-5 \pm 14 \pm 3) \times 10^{-3} \text{ rad}$	-		
$A_{CP}^\Lambda$	$(-4 \pm 12 \pm 9) \times 10^{-3}$	$(-6 \pm 12 \pm 7) \times 10^{-3}$	Ref. <sup>4</sup>	***
$\langle\phi_{\Xi^-}\rangle$	$0.016 \pm 0.014 \pm 0.007 \text{ rad}$			

The  $J/\psi \rightarrow \Xi^- \bar{\Xi}^+$  angular distribution parameter  $\alpha_\psi$ , the hadronic form factor phase  $\Delta\Phi$ , the decay parameters for  $\Xi^- \rightarrow \Lambda \pi^-$  ( $\alpha_{\Xi^-}, \phi_{\Xi^-}$ ),  $\Xi^- \rightarrow \bar{\Lambda} \pi^+$  ( $\bar{\alpha}_{\Xi^-}, \bar{\phi}_{\Xi^-}$ ),  $\Lambda \rightarrow p \pi^-$  ( $\alpha_\Lambda$ ) and  $\bar{\Lambda} \rightarrow \bar{p} \pi^+$  ( $\bar{\alpha}_\Lambda$ ); the CP asymmetries  $A_{CP}^{\Xi^-}$ ,  $\Delta\phi_{CP}^{\Xi^-}$  and  $A_{CP}^\Lambda$ , and the average  $\langle\phi_{\Xi^-}\rangle$ . The first and second uncertainties are statistical and systematic, respectively.

We obtain the same precision for  $\phi$  as HyperCP with **three orders of magnitude** smaller data sample!

$$\phi_{\Xi^-, \text{HyperCP}} = -0.042 \pm 0.011 \pm 0.011$$

$$\langle\phi_{\Xi^-}\rangle = 0.016 \pm 0.014 \pm 0.007$$

Strong phase measurement compatible with SM  $(1.9 \pm 4.9) \times 10^{-2}$  but in tension with HyperCP  $2.6\sigma$

PRD 67 056001 (2004)

- \* PRD 93, 072003 (2018)
- \*\* PDG 2020
- \*\*\* Nat. Ph. 15, 631 (2019)
- \*\*\*\* PRL 93, 011802 (2004)



## Article

Extended Data Table 1 | Correlation coefficients for the production and asymmetry decay parameters

	$\alpha_\Psi$	$\Delta\Phi$	$\alpha_\Xi$	$\phi$	$\alpha_\Lambda$	$\bar{\alpha}_\Xi$	$\bar{\alpha}_\Lambda$	$\bar{\phi}_\Xi$
$\alpha_\Psi$	1.0	0.414	-0.008	-0.006	0.107	0.014	0.120	0.003
$\Delta\Phi$		1.0	-0.016	0.016	0.133	0.008	0.138	-0.029
$\alpha_\Xi$			1.0	-0.000	0.280	0.024	0.071	0.010
$\phi_\Xi$				1.0	-0.002	-0.010	-0.010	0.013
$\alpha_\Lambda$					1.0	0.070	0.401	0.014
$\bar{\alpha}_\Xi$						1.0	0.269	0.001
$\bar{\alpha}_\Lambda$							1.0	0.006
$\bar{\phi}_\Xi$								1.0

Correlation between  $\alpha_\Lambda$  and  $\bar{\alpha}_\Lambda$   
lower compared to  $\Lambda\bar{\Lambda}$

Low correlation between  $\phi_\Xi$   
and all other parameters



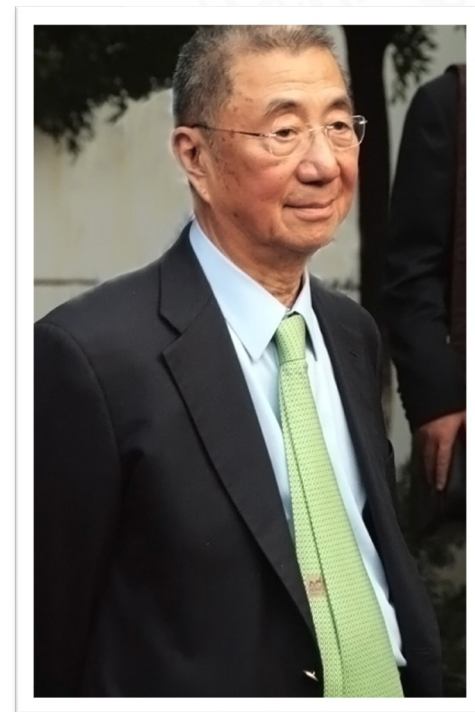
We have presented a novel *model-independent method that uses spin entanglement in sequential weak decay chain  $E^- \rightarrow \Lambda\pi^-, \Lambda \rightarrow p\pi^-$*

First measurement of weak phase difference for any baryon decay published yesterday. First Nature publication of BESIII, but why does it matter?

*“...the most striking finding in the current paper is the demonstrated improvement that the use of entanglement, the imprimatur of quantum mechanics, gives in leveraging the sensitivity of tests of CP violation” - Anonymous Nature referee 1*

*“In reality, a theory in natural science cannot be without experimental foundations; physics, in particular, comes from experimental work.” - 丁肇中*

Consider how far we have come from the original discovery of J/ψ!

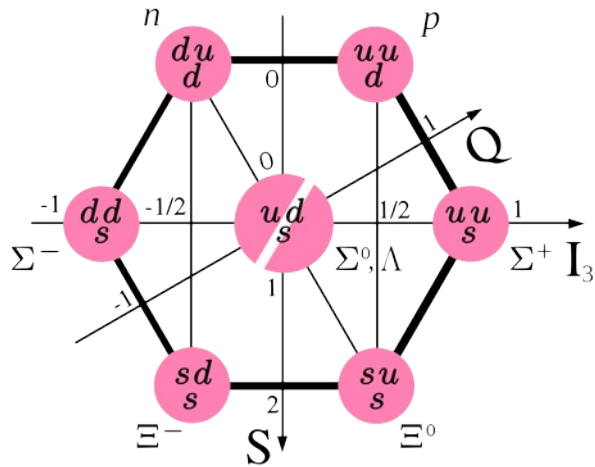


**Samuel Chao Chung Ting 丁肇中**  
Discovered J/ψ Nobel prize 1976



Based on these discoveries and  $10^{10}$   $J/\psi$  data set there are many more modes which can be analyzed by BESIII

In particular reactions with neutral final state particles  $n, \gamma$



	$D$	$B$	$\langle \alpha_D \rangle$	$\langle \phi_D \rangle$ [rad]	$A_{CP}$	Comment
$\Lambda \rightarrow p\pi^-$ [ $\Lambda p$ ]	64%	<b>0.755(03)<sup>a</sup></b>	<b>-0.113(61)<sup>b</sup></b>	<b>-0.005(10)<sup>a</sup></b>		
		0.754(3)(2)	-	-0.006(12)(7)		BESIII [31]
		0.721(6)(5) <sup>*</sup>	-	-		CLAS [47]
		0.760(6)(3)	-	-0.004(12)(9)		BESIII [32]
$\Lambda \rightarrow n\pi^0$ [ $\Lambda n$ ]	36%	<b>0.692(17)<sup>c</sup></b>	-	-		BESIII [31]
$\Sigma^+ \rightarrow p\pi^0$ [ $\Sigma p$ ]	52%	<b>-0.994(04)<sup>d</sup></b>	<b>0.63(59)<sup>g</sup></b>	<b>-0.004(37)<sup>d</sup></b>		
$\Sigma^+ \rightarrow n\pi^+$ [ $\Sigma n$ ]	48%	<b>0.068(13)<sup>*</sup></b>	<b>2.91(35)<sup>*</sup></b>	-		PDG [28]
$\Sigma^- \rightarrow n\pi^-$ [ $\Sigma^-$ ]	100%	<b>-0.068(08)<sup>*</sup></b>	<b>0.17(26)<sup>*</sup></b>	-		PDG [28]
$\Xi^0 \rightarrow \Lambda\pi^0$ [ $\Xi 0$ ]	100%	<b>-0.345(08)<sup>e</sup></b>	<b>0.36(21)<sup>*</sup></b>	-		AVG [48, 49]
$\Xi^- \rightarrow \Lambda\pi^-$ [ $\Xi^-$ ]	100%	<b>-0.379(04)<sup>f</sup></b>	<b>-0.042(16)<sup>*</sup></b>	-		AVG [28, 50]
		-0.373(5)(2)	0.016(14)(7)	<b>0.006(13)(6)</b>		BESIII [32]

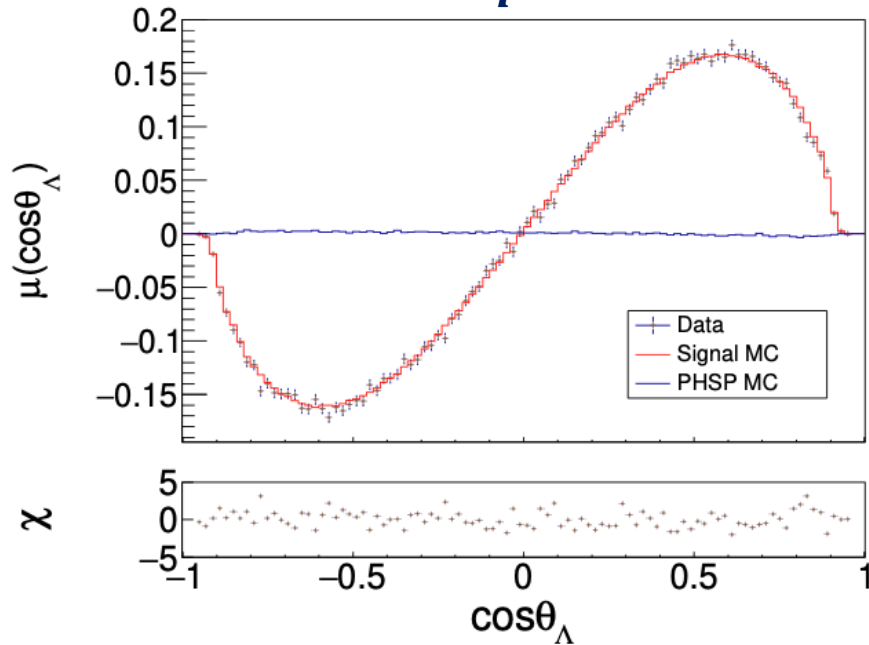
$\Xi^0 \rightarrow \Lambda\pi^0, \Lambda \rightarrow p\pi^-$  only few hundred events required to match the current PDG precision of  $\phi_{\Xi^0}$

Only indirect determination of  $\alpha_{\Xi}$  exists

$$\phi_{\Xi^0 \rightarrow \Lambda\pi^0} = 21(12)_{tot}^\circ$$





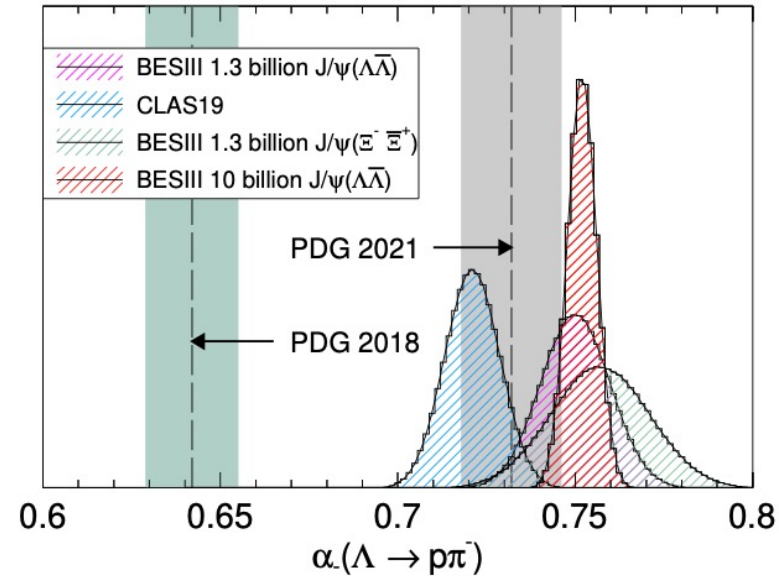
$\Lambda \rightarrow p\pi^-$ 


Based on 3.23M  $\bar{\Lambda}\Lambda$

Par.	This work	Previous results [8]
$\alpha_{J/\psi}$	$0.4748 \pm 0.0022 \pm 0.0024$	$0.461 \pm 0.006 \pm 0.007$
$\Delta\Phi$	$0.7521 \pm 0.0042 \pm 0.0080$	$0.740 \pm 0.010 \pm 0.009$
$\alpha_-$	$0.7519 \pm 0.0036 \pm 0.0019$	$0.750 \pm 0.009 \pm 0.004$
$\alpha_+$	$-0.7559 \pm 0.0036 \pm 0.0029$	$-0.758 \pm 0.010 \pm 0.007$
$A_{CP}$	$-0.0025 \pm 0.0046 \pm 0.0011$	$0.006 \pm 0.012 \pm 0.007$
$\alpha_{\text{avg}}$	$0.7542 \pm 0.0010 \pm 0.0020$	-

arXiv: 2204.11058

arXiv: 2204.11058



$$\langle \alpha(\Lambda \rightarrow p\pi^-) \rangle_{\Lambda} = 0.754(1)(2)$$

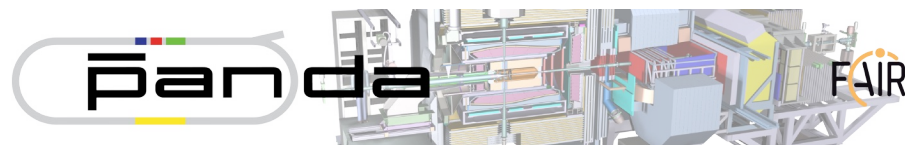
$$\langle \alpha(\Lambda \rightarrow p\pi^-) \rangle_{\Xi} = 0.760(6)(3)$$

Precision determined by this new result  
Most precise  $A_{CP, \Lambda}$ !

Indication of what precision to expect

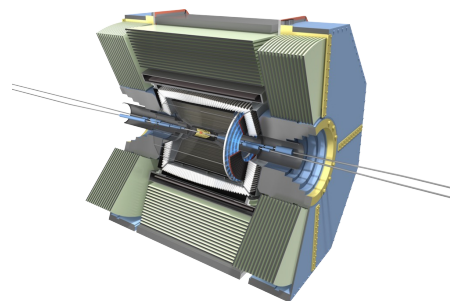


PANDA@FAIR



The potential of  $\Lambda$  and  $\Xi-\Xi-$  studies with PANDA at FAIR  
 Eur. Phys. J. A 57 No. 154 (2021), [arXiv: 2009.11582](https://arxiv.org/abs/2009.11582),

BELLE-II

Super  $\tau$  charm factories

# Case study Super $\tau$ charm factories

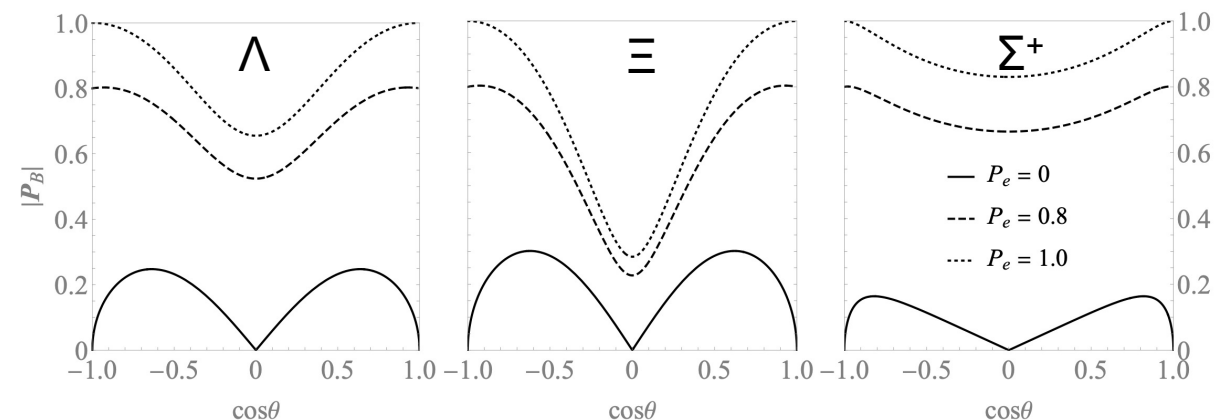
Using polarized electron beam can greatly enhance sensitivity!

Non-polarized beam

$$C_{\mu\nu} = (1 + \alpha_\psi \cos^2\theta) \begin{pmatrix} 1 & 0 & P_y & 0 \\ 0 & C_{xx} & 0 & C_{xz} \\ -P_y & 0 & C_{yy} & 0 \\ 0 & -C_{xz} & 0 & C_{zz} \end{pmatrix}$$

polarized beam

$$\frac{3}{3 + \alpha_\psi} \cdot \begin{pmatrix} 1 + \alpha_\psi \cos^2\theta & \gamma_\psi P_e \sin\theta & \beta_\psi \sin\theta \cos\theta & (1 + \alpha_\psi) P_e \cos\theta \\ \gamma_\psi P_e \sin\theta & \sin^2\theta & 0 & \gamma_\psi \sin\theta \cos\theta \\ -\beta_\psi \sin\theta \cos\theta & 0 & \alpha_\psi \sin^2\theta & -\beta_\psi P_e \sin\theta \\ -(1 + \alpha_\psi) P_e \cos\theta & -\gamma_\psi \sin\theta \cos\theta & -\beta_\psi P_e \sin\theta & -\alpha_\psi - \cos^2\theta \end{pmatrix}$$

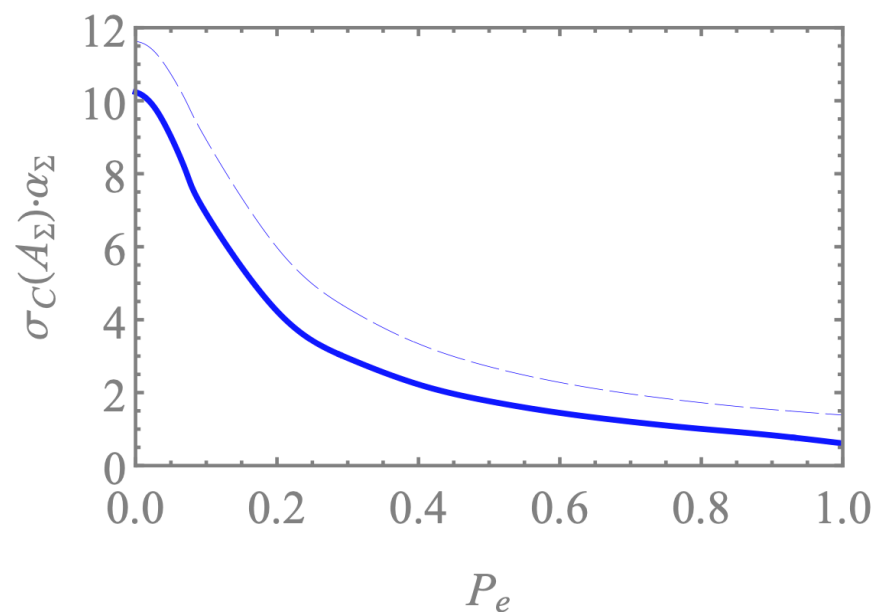
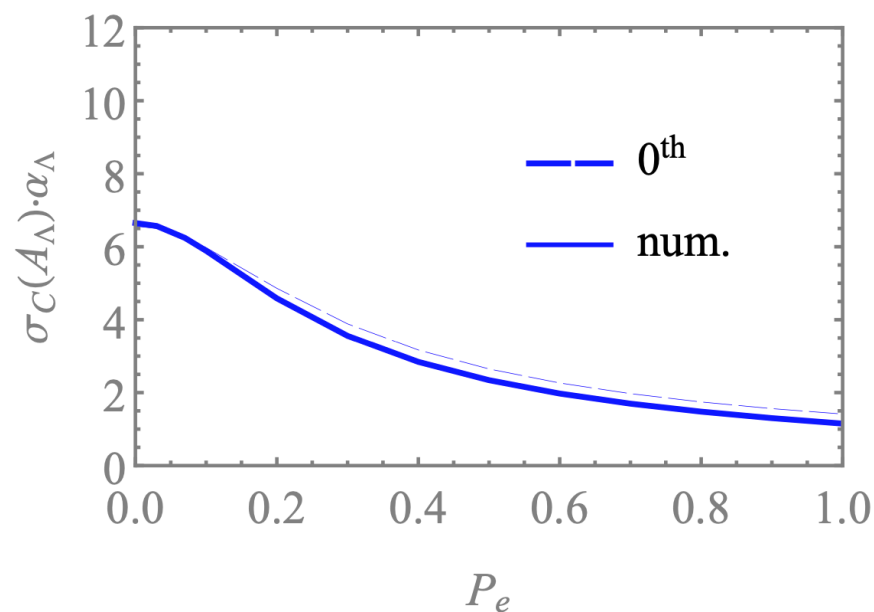


arXiv 2203:03035



S $\tau$ CF: single weak decay

$$\sigma_C(\omega_k) := \sigma(\omega_k) \times \sqrt{N}$$



Double tag

0<sup>th</sup> order approximation  $\sigma(A_{CP})\sqrt{N} = \sigma_C(A_{CP}) \approx \sqrt{\frac{3}{2}} \frac{1}{\alpha_D \sqrt{\langle \mathbf{P}_B^2 \rangle}}$

arXiv 2203:03035

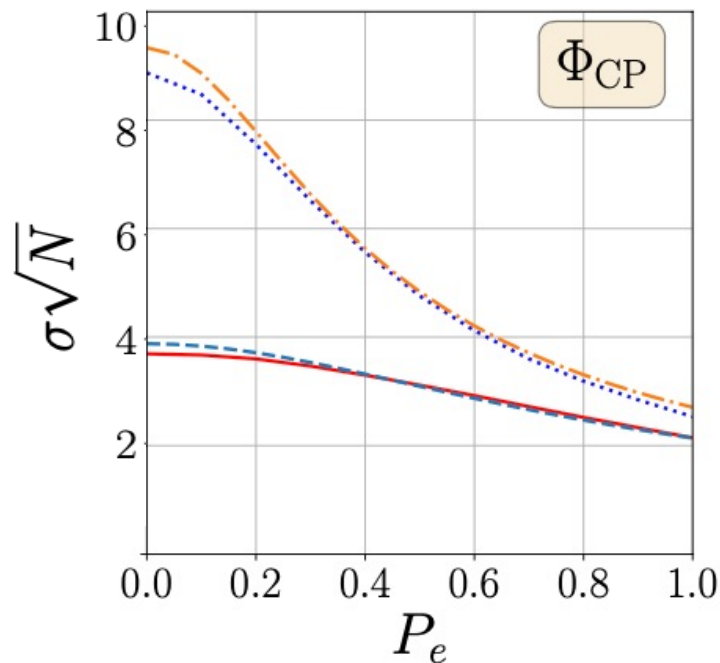


# S $\tau$ CF: sequential weak decay

Using polarized electron beam can greatly enhance sensitivity!

$$\mathcal{I}_0(\Phi_{\text{CP}}) = \frac{2N}{27} (1 - \alpha_{\Xi}^2) \alpha_{\Lambda}^2 \left[ (3 + \alpha_{\Xi}^2 \alpha_{\Lambda}^2) \langle \mathbb{P}_{\Xi}^2 \rangle + \frac{2}{3} (\alpha_{\Xi}^2 (3 - 2\alpha_{\Lambda}^2) + 3\alpha_{\Lambda}^2) \langle \mathbb{S}_{\Xi\Xi}^2 \rangle \right]$$

$$\sigma(\Phi_{\text{CP}}) = 1/\sqrt{\mathcal{I}(\Phi_{\text{CP}})}$$



↑  
Spin correlation contributions

Dotted: ST

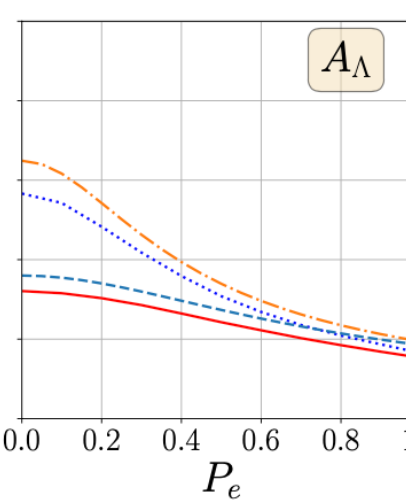
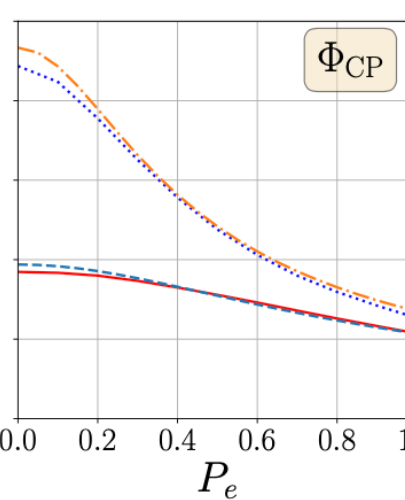
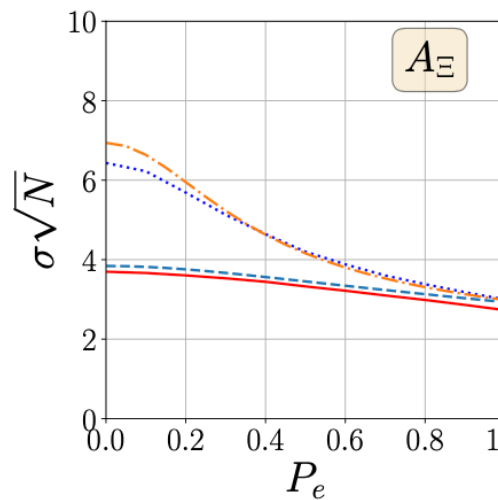
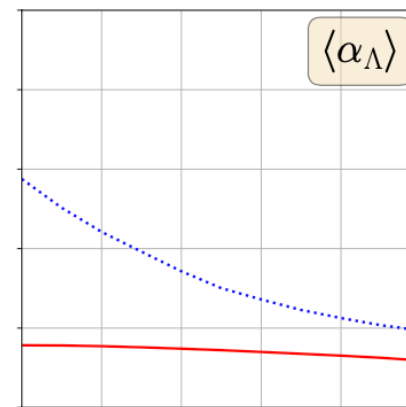
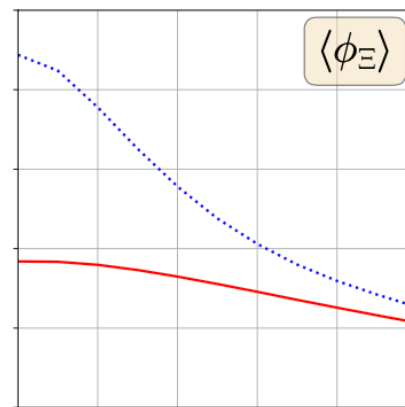
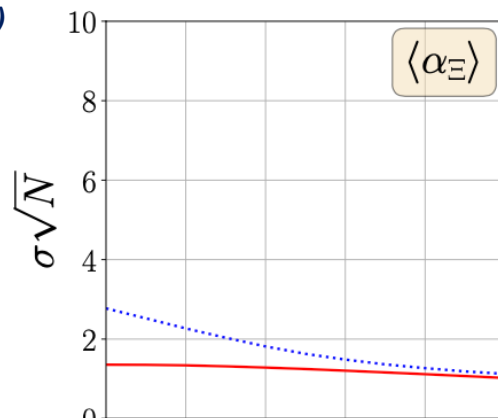
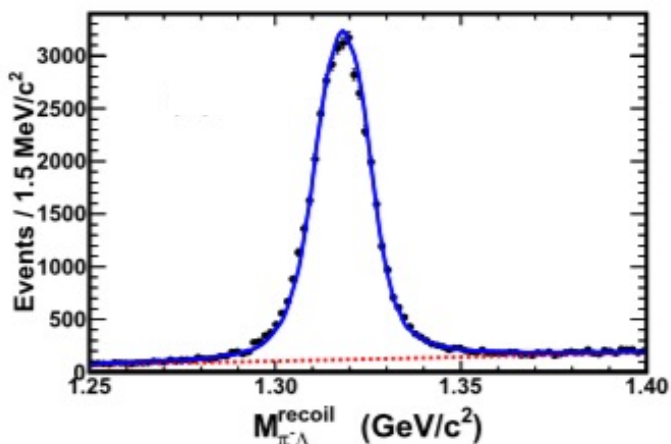
Solid: DT

Analytic approximation: dashed dotted  
dashed

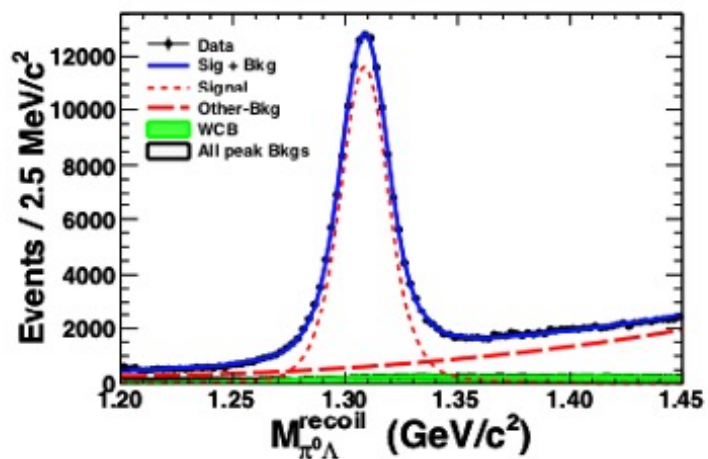


It becomes beneficial to also include single-tag events

(BESIII) PHYSICAL REVIEW D 93, 072003 (2016)



(BESIII) Phys Lett B 770, 217 (2017)

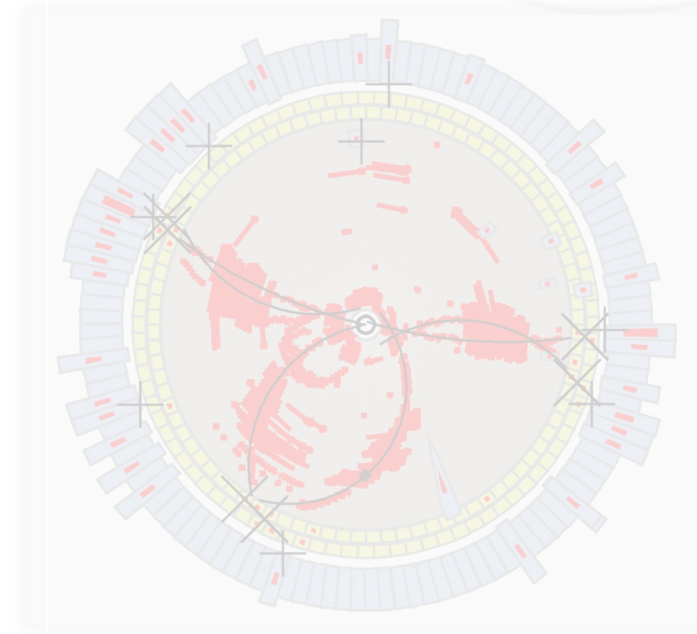


We have presented a novel *model-independent method that uses spin entanglement in sequential weak decay chain*  $E^- \rightarrow \Lambda\pi^-, \Lambda \rightarrow p\pi^-$

First measurement of weak phase difference for any baryon decay published yesterday. First Nature publication of BESIII

The benefits of using entangled pairs can be adopted by other experiments e.g. PANDA, BELLE-II and Super-tau Charm factories. Polarization of 0.8 possible?

BESIII recently collected  $1.0 \times 10^{10} J/\psi$  events. More results to be expected in future!



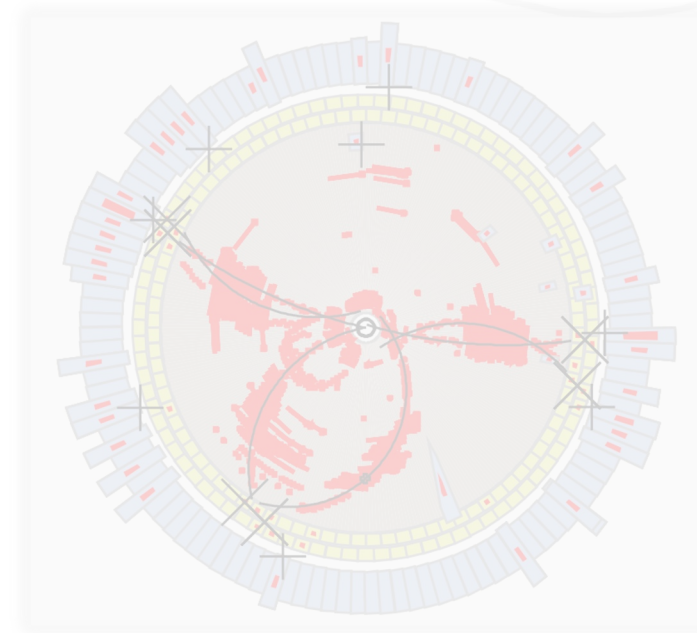
**Thank you for your attention!**



Extended Data Table 2 | Contributing systematic uncertainties, and the sum in quadrature

$\times 10^2$	$\alpha_\psi$	$\Delta\Phi$	$\alpha_\Xi$	$\bar{\alpha}_\Xi$	$\alpha_\Lambda$	$\bar{\alpha}_\Lambda$	$\phi_\Xi$	$\bar{\phi}_\Xi$
Statistical	1.2	4.6	0.70	0.70	1.05	1.06	1.91	1.93
Kin. fit	0.36	1.5	0.18	0.17	0.21	0.43	0.77	0.44
mass win $\Lambda$	0.44	0.44	0.07	0.02	0.56	0.33	0.17	0.46
mass win $\Xi$	0.25	-	-	-	0.36	-	0.46	-
dec. length $\Lambda$	-	-	-	-	0.30	0.40	-	-
Track. eff.	0.80	0.41	0.27	0.05	0.21	0.14	0.16	0.16
Sum syst.	1.0	1.6	0.33	0.18	0.79	0.69	0.93	0.66

First row: statistical uncertainty as reference. The uncertainties of  $\Delta\Phi$  and  $\phi$  are given in radians. All values multiplied by a factor  $10^2$ , as indicated at top left.





$\times 10^2$	$A_{\Lambda,CP}$	$A_{\Xi,CP}$	$\Delta\phi_{CP}$ (rad)	$\delta_P - \delta_S$ (rad)	$\zeta_P - \zeta_S$ (rad)
Statistical	1.17	1.34	1.37	3.3	3.4
Kin. fit	0.32	0.47	0.16	1.3	0.4
mass win. $\Lambda$	0.59	0.07	0.14	0.8	0.4
mass win. $\Xi$	0.38	-	0.20	0.7	0.5
dec. length $\Lambda$	0.46	-	-	-	-
Track. eff.	0.05	0.29	0.003	0.4	$2 \cdot 10^{-3}$
Sum syst.	0.90	0.56	0.29	1.7	0.75

First row: statistical uncertainty as reference. All values multiplied by a factor  $10^2$ , as indicated at top left.

