

Heavy-quark-pair production at lepton colliders at NNNLO in QCD

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In collaboration with Xin Guan, Chuan-Qi He, Xiao Liu, Yan-Qing Ma
Based on: [arXiv: 2209.14259]



1. Introduction: Motivation and Background

2. Computational Techniques

3. Numerical Results

4. Summary and Outlook

Motivation and Background

- Top quark plays a crucial role
 - Electroweak(EW) physics
 - physics beyond the Standard Model
- Unprecedented experiment precision
 - e^+e^- colliders offer a clean environment
 - Precision measurement: top-quark mass, cross section, A_{FB}^t
 - A few ‰ to % precision at the on-going colliders are possible.

The International Linear Collider. (2013)
The CEPC Study Group. (2018)
FCC-ee, A. Abada *et al.* (2019)
ILD Concept Group. (2020)

Theoretical efforts in literature

➤ NNLO

- Near-threshold expansion and high energy expansion of cross section are known for more than two decades

B. H. Smith, M. B. Voloshin. *Phys. Lett. B*(1994)

K. G. Chetyrkin, A. Kwiatkowski. *Z. Phys. C*(1994)

K. G. Chetyrkin, J. H. Kuhn, M. Steinhauser. *Nucl. Phys. B*(1996)

- Cross section and fully differential distributions in the continuum

Jun Gao, Hua Xing Zhu. *Phys. Rev. D*(2014)

Jun Gao, Hua Xing Zhu. *Phys. Rev. Lett.* (2014)

Long Chen, Oliver Dekkers, et al. *JHEP*(2016)

\sqrt{s} [GeV]	360	381.3	400	500
Δ_1	0.627	0.352	0.266	0.127
Δ_2	0.281	0.110	0.070	0.020

$$\sigma_{NNLO} = \sigma_{LO}(1 + \Delta_1 + \Delta_2)$$

From Long Chen, Oliver Dekkers, et al. *JHEP*(2016)

- The theoretical uncertainty is about **1%** From Jun Gao, Hua Xing Zhu. *Phys. Rev. D*(2014)
- Full computation of NNNLO QCD correction is **necessary!**

➤ NNNLO(partial results)

- Near-threshold expansion M. Beneke, et al. *Phys. Rev. Lett.* (2015)
- Massive form factors M. Fael, F. Lange, K. Schonwald, M. Steinhauser. *Phys. Rev. Lett.* (2022)
M. Fael, F. Lange, K. Schonwald, M. Steinhauser. *Phys. Rev. D* (2022)

Outline

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The computational work-flow

- Generating Feynman diagrams ($\mathcal{O}(\text{minutes})$)
- Dealing with Lorentz and color algebras ($\mathcal{O}(\text{few minutes to hours})$)
- IBP(integration-by-parts) reduction of loop integrals ($\mathcal{O}(\text{hours to days})$)
- Evaluating master integrals ($\mathcal{O}(\text{days})$ typically in frontier)

Dealing with Feynman Amplitudes

➤ Generating Feynman diagrams

- Tools: **QGRAF** [P. Nogueira, J. Comput. Phys. \(1993\)](#)
- Sum all virtual and real correction at NNNLO
- Focus: **vector contribution**, i.e. $e^+e^- \rightarrow \gamma^* \rightarrow t\bar{t}$

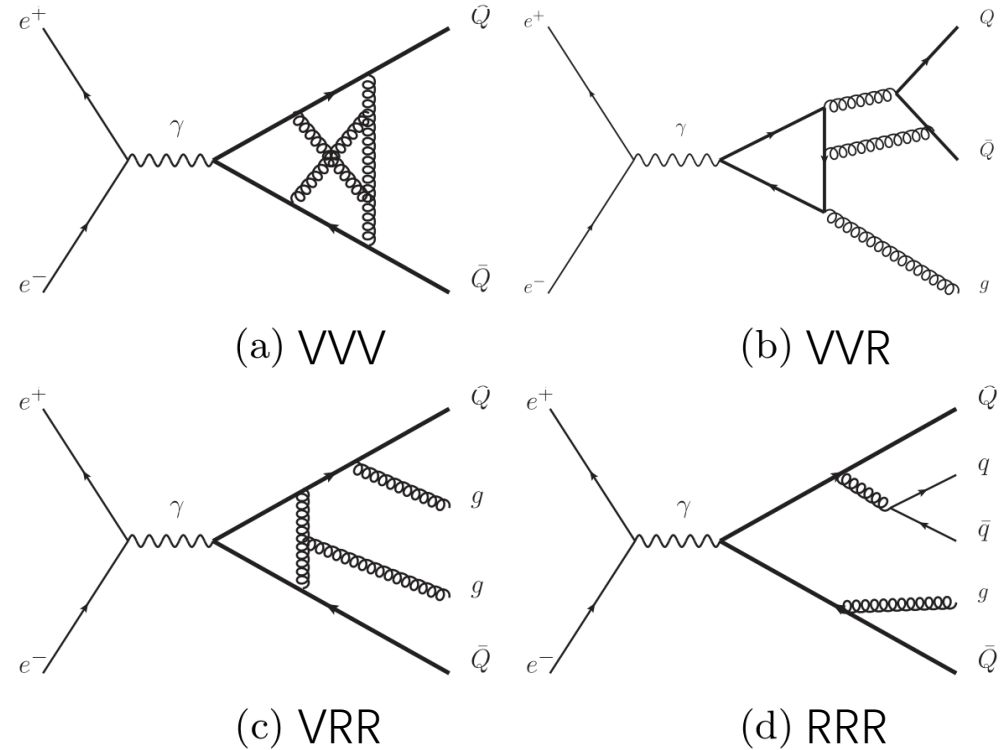
➤ Dealing with Feynman amplitudes

- Lorentz and color algebra
- Phase-space integrals \Rightarrow reverse unitarity

$$\delta(x) \rightarrow \frac{1}{2\pi i} \left(\frac{1}{x - i0} - \frac{1}{x + i0} \right)$$

- Express squared amplitudes as linear combinations of scalar integrals
- Total **475** families (VVV, VVR, VRR, RRR)

$$|M|^2 = \sum_{i=1}^{O(10^4)} a_i \times I_i$$



“V” means “Virtual”, “R” means “Real”

Integrals reduction

➤ Strategy:

- Laporta's algorithm [S. Laporta. *Int. J. Mod. Phys. A* \(2000\)](#)
- Tools: **LiteRed** and **FiniteFlow** \Rightarrow IBP system (over finite field) [R. N. Lee. *J. Phys. Conf. Ser.* \(2014\)](#)
[T. Peraro. *JHEP* \(2019\)](#)
- Numerical samplings \Rightarrow reconstruct the reduction coefficients



Block-triangular relations [X. Guan, X. Liu, Y.-Q. Ma. *Chin. Phys. C* \(2020\)](#)

- Block-triangular form is an improved linear system
 - ✓ Contains as many equations as target integrals
 - ✓ **Several orders of magnitude less** than plain IBP system
- Applied to this work and our recent another work [X. Chen, X. Guan, C.-Q. He, Z. Li, X. Liu, Y.-Q. Ma. *arXiv:2209.14953*](#)
[\rightarrow Talk by Xin Guan]

Integrals reduction computing resources

➤ Computing resources:

- Total 475 families

Contribution	Number of families	Computing resources (the most expensive family)
VVV	62	10 hours on 16 cores
VVR	147	5 hours on 16 cores
VRR	167	4 hours on 16 cores
RRR	99	2 hours on 16 cores

- The computational time of the most expensive family of VVV contribution is about **one order of magnitude shorter** than that reported in literature

[M. Fael, F. Lange, K. Schonwald, M. Steinhauser. *Phys. Rev. D* \(2022\)](#)

- **Block-triangular form reduction is much more efficient!**

Evaluating master integrals

R. N. Lee, V. A. Smirnov. *JHEP*(2018)

R. Bonciani, G. Degrossi, P. P. Giardino, R. Grober. *Comput. Phys. Commun* (2019)

H. Frellesvig, M. Hidding, L. Maestri, F. Moriello, G. Salvatori. *JHEP*(2020)

➤ Method: using **numerical differential equations** based on power series expansion

Long Chen, M. Czakon, M. Niggetiedt. *JHEP*(2021)

Martijin Hidding. *Comput. Phys. Commun.* (2021)

X. Liu, Y.-Q. Ma. *arXiv:2201.11669*(2021)

M. Fael, F. Lange, K. Schonwald, M. Steinhauser. *Phys. Rev. Lett.* (2022)

➤ General procedure

- Set up the differential equations with respect to $x = \frac{4m^2}{s}$ by IBP reducing derivatives

$$\frac{dI_i(\epsilon, x)}{dx} = \sum_j A_{ij} I_j(\epsilon, x)$$

- Obtain boundary condition (**non-trivial**)

We use **AMFlow** to calculate master integrals at $x = \frac{4}{23}$ (arbitrary regular point) [X. Liu, Y.-Q. Ma. arXiv:2201.11669\(2021\)](#) [→ Talk by Xiao Liu]

- Generalized series expansions around singularities

$$I(\epsilon, x) = \sum_{\mu, k, n} c_{\mu, k, n}(\epsilon) x^{\mu(\epsilon)} \log^k x x^n$$

} ⇒ Cover physical region

- Taylor expansions around regular points
- Match between two neighboring expansions

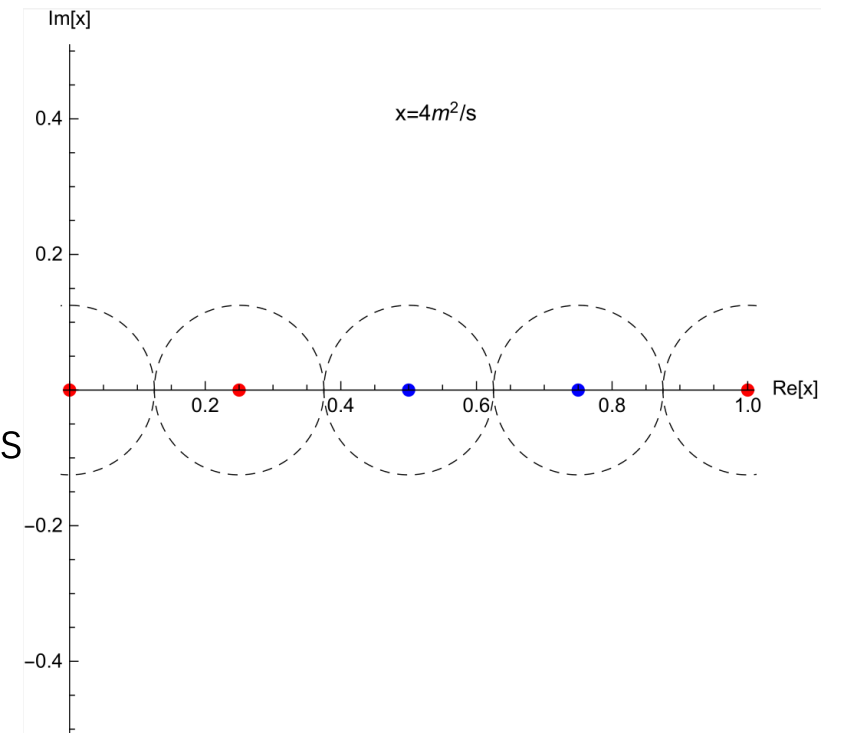
Results for master integrals

➤ Constructing a piecewise function of $x(= \frac{4m^2}{s})$ for each master integral

- 5 deeply expanded power series at the following points

$$x \rightarrow \left\{ 0, \frac{1}{4}, \frac{1}{2}, \frac{3}{4}, 1 \right\}$$

- $0, \frac{1}{4}, 1$ are physical singularities, $\frac{1}{2}$ and $\frac{3}{4}$ are regular points
- The expansion at each x_0 is **valid** in $x \in \left(x_0 - \frac{1}{8}, x_0 + \frac{1}{8}\right)$
- The whole physical region $x \in (0,1)$ is covered



Renormalization

- α_s in the \overline{MS} scheme: multiplicative
- Wave-function renormalization in the **on-shell** scheme: multiplicative
- Quark-mass renormalization in the **on-shell** scheme:
 - Replace bare mass by renormalized mass via $m^b = m^{OS} + \delta m$
 - Keep on-shell condition of external top quark momentum $p^2 = (m^{OS})^2$ untouched
 - In spirit of perturbation theory, δm is a small quantity and can be expanded
$$\sigma(m^b) = \sigma(m^{OS}) + \left. \frac{\partial \sigma(m)}{\partial m} \right|_{m=m^{OS}} \delta m + \dots, \quad \delta m = (Z_m^{OS} - 1)m^{OS}$$
 - Substitute renormalization constant

Verification

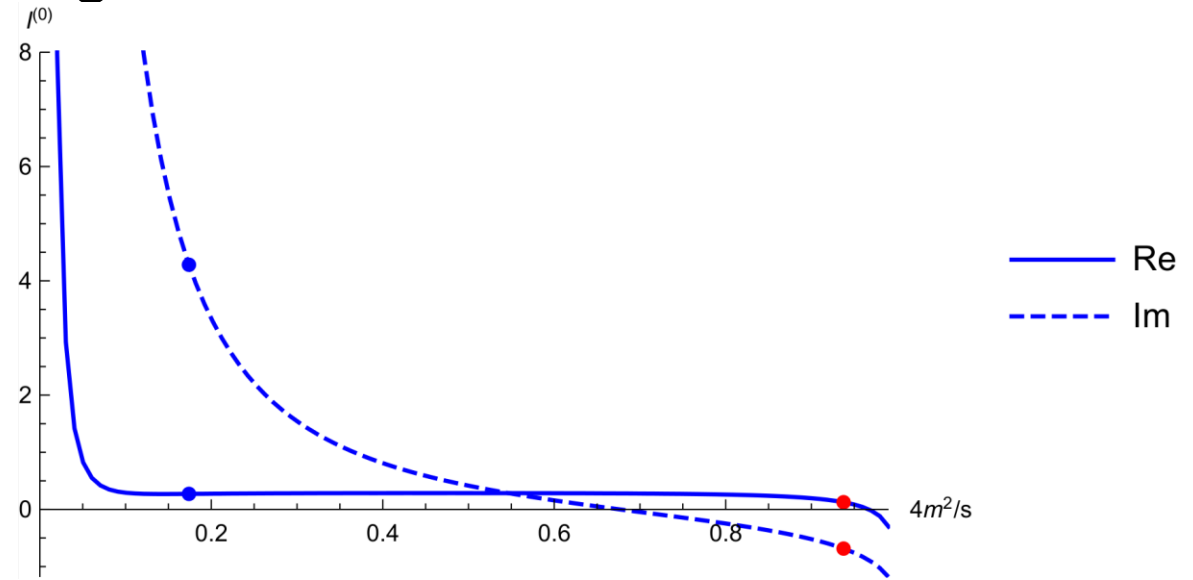
➤ Free of ultraviolet and infrared divergences

➤ Check master integrals

All MIs obtained by solving DEs with boundary conditions $I\left(x = \frac{4}{23}, \epsilon\right)$

|| checked

Evaluating master integrals by **AMFlow** at $x = \frac{15}{16}$
(another arbitrary point)



The blue lines correspond to the results of differential equations, the blue points correspond to boundary condition, the red points correspond to check point.

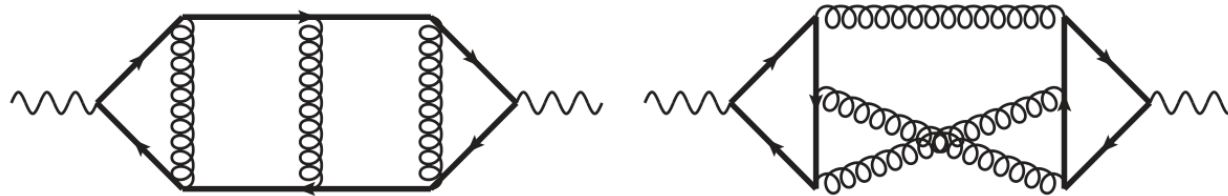
➤ Three-loop contributions agree with massive form factors

[M. Fael, F. Lange, K. Schonwald, M. Steinhauser. *Phys. Rev. Lett.* \(2022\)](#)

[M. Fael, F. Lange, K. Schonwald, M. Steinhauser. *Phys. Rev. D* \(2022\)](#)

Verification

- We calculate the cross section in another totally different way
 - **Optical theorem**: imaginary part of forward scattering amplitude \Rightarrow cross section
 - Same computational strategy
 - Sample Feynman diagrams:



- Using **block-triangular** relations is expected to be **2 orders of magnitude faster** than plain IBP system
- Contributions from **massless final states** without any top quark must be subtracted
- Contributions from **four-heavy-quark production** should be subtracted when we consider the region above its threshold, i.e. $x < 1/4$

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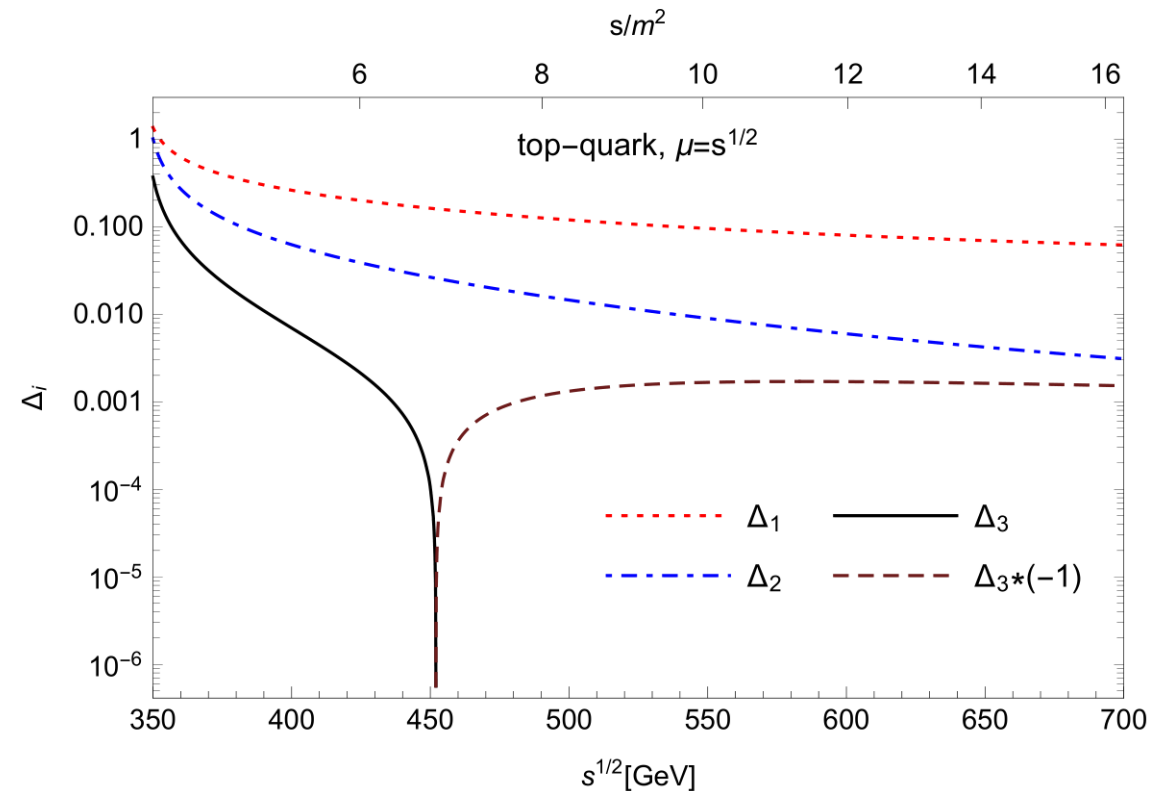
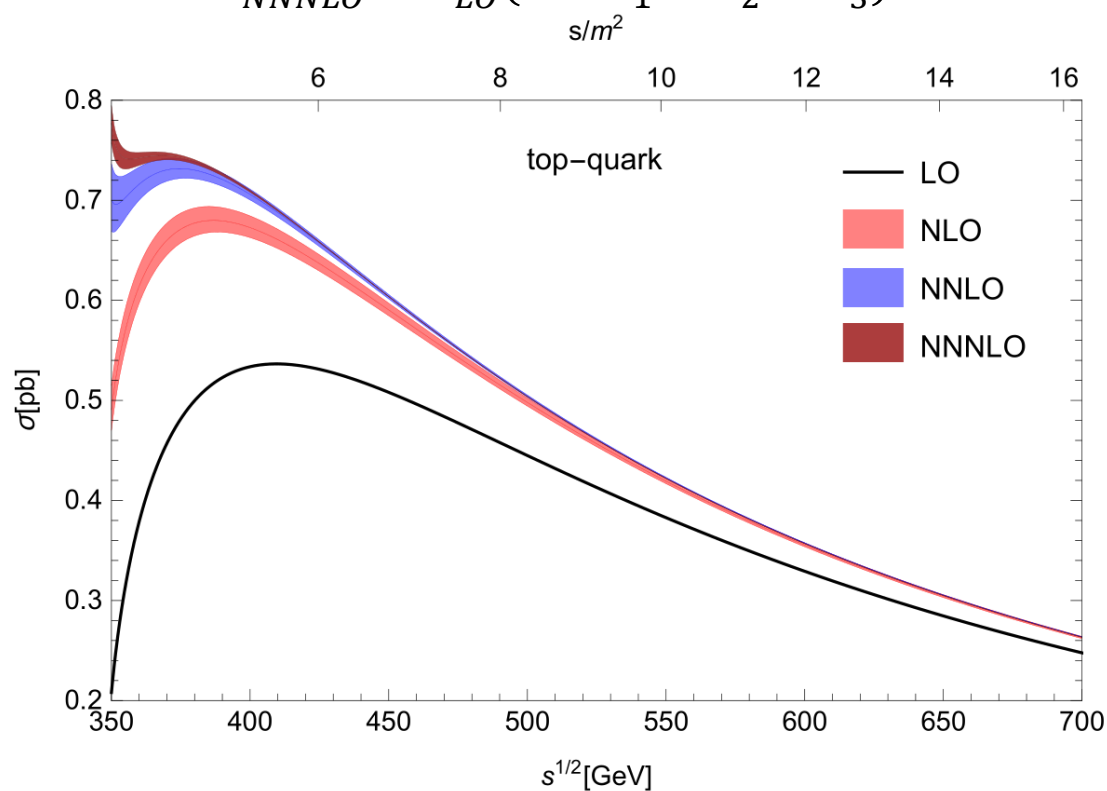
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Numerical Results for NNNLO cross section

➤ Top quark

- Input values: $m_t = 173.34\text{GeV}$, $\alpha = 1/132.2$, $\alpha_s^{n_f=5}(m_Z = 91.1876\text{GeV}) = 0.1181$
- $\sigma_{NNNLO} = \sigma_{LO}(1 + \Delta_1 + \Delta_2 + \Delta_3)$

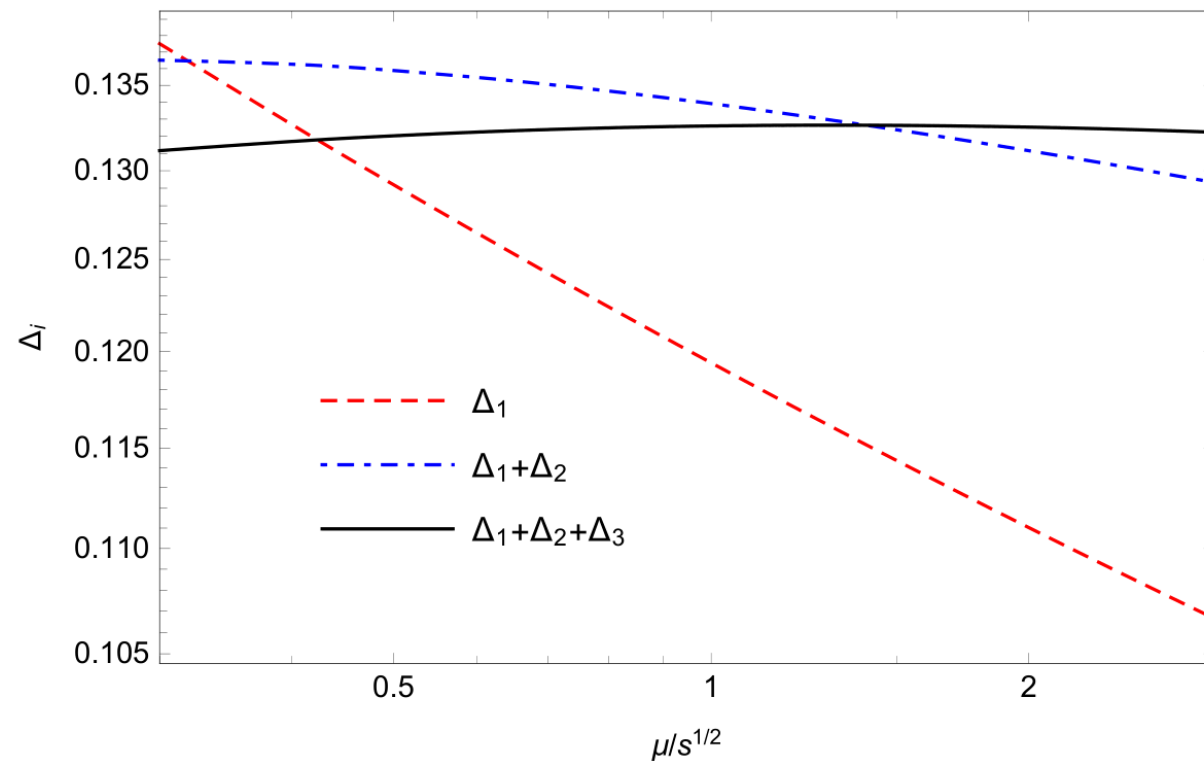


The middle lines correspond to $\mu = \sqrt{s}$, the upper and lower lines correspond to $\mu = 2\sqrt{s}$ and $\mu = \sqrt{s}/2$

Numerical Results for NNNLO cross section

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top-quark, $s^{1/2}=500\text{GeV}$



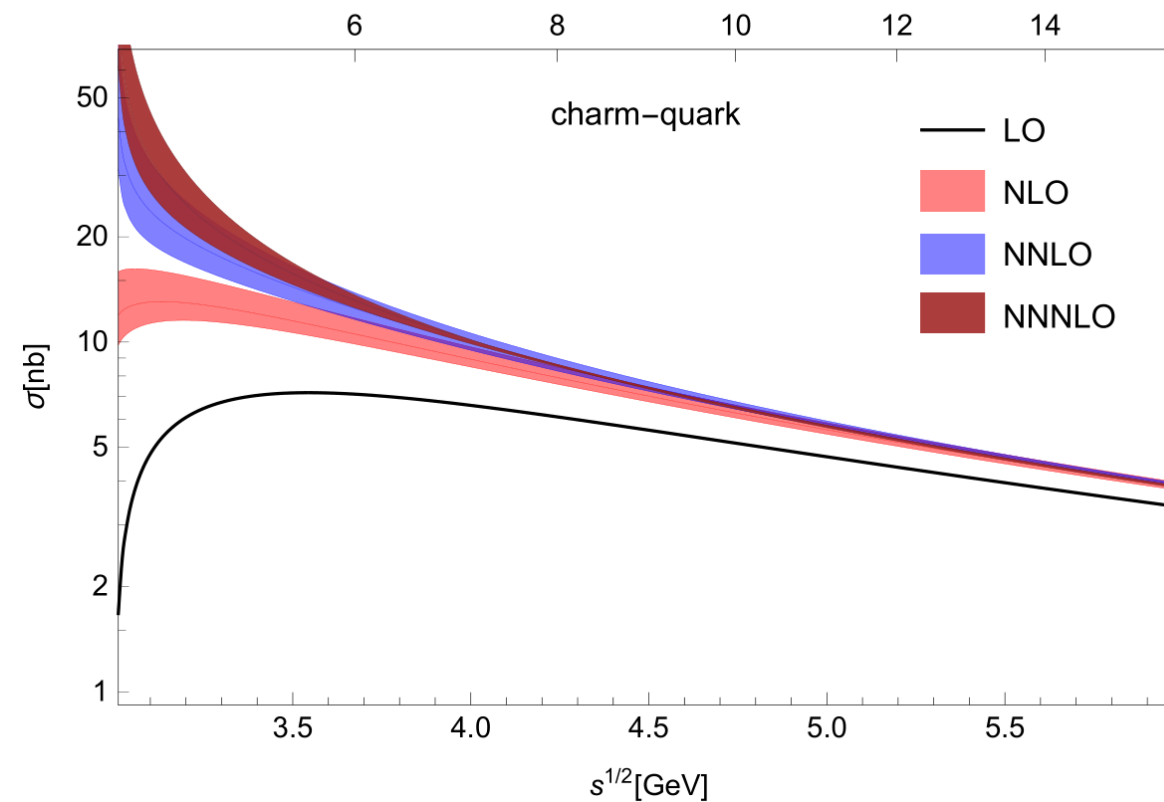
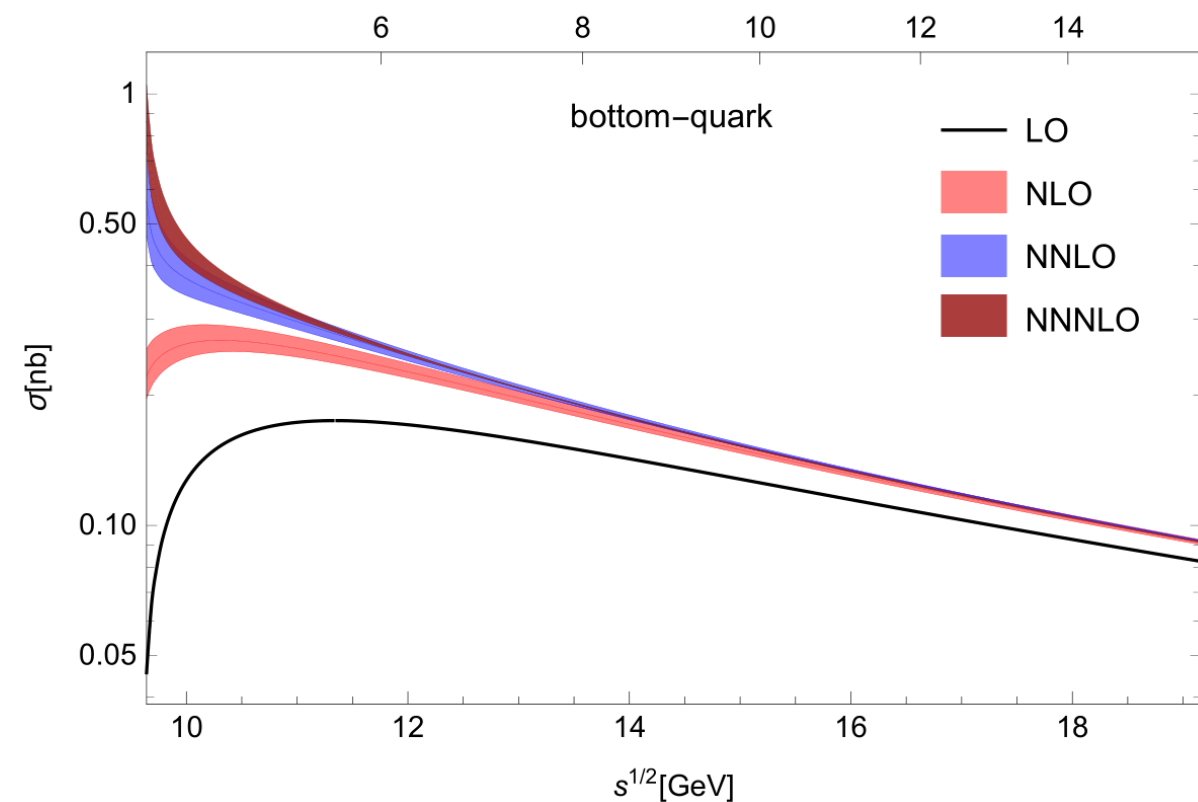
The scale dependence has been reduced from **0.72%** at **NNLO** to **0.15%** at **NNNLO** for a collider energy of **500GeV**, which **meets the precision (1‰ to 1%)** requested by future lepton colliders!

ILD Concept Group. (2020)

Numerical Results for NNNLO cross section

➤ Bottom quark and Charm quark

- Input values: $m_b = 4.8\text{GeV}$, $m_c = 1.5\text{GeV}$, $\alpha = 1/132.2$, $\alpha_s^{n_f=5}(m_Z = 91.1876\text{GeV}) = 0.1181$



The middle lines correspond to $\mu = \sqrt{s}$, the upper and lower lines correspond to $\mu = 2\sqrt{s}$ and $\mu = \sqrt{s}/2$

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Summary and Outlook

➤ Summary

- We provide the **first NNNLO perturbative correction** to $t\bar{t}$ production cross section at lepton colliders mediated by a **virtual photon**.
- We combine several powerful techniques:
 - ✓ **Block-triangular relations**
 - ✓ **Auxiliary mass flow**
 - ✓ Numerical sampling of ϵ
 - ✓ A novel way to perform renormalization
- The calculation are done **in two totally different way** for cross check.
- Our result significantly **reduce theoretical uncertainty**, which **meets the precision** requested by future lepton colliders.

➤ Outlook

- Contribution of top-quark-pair production mediated by a virtual **Z boson** will be calculated in the near future.
- Our strategy is **applicable** for many other process.