

# Modern Methods for Multi-Scale Scattering Amplitudes

Ben Page

CERN, Theoretical Physics Department

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# Precise Perturbative Predictions

- **Perturbation theory** is our major tool for making predictions.

$$\sigma \sim \sigma_{\text{LO}} + \alpha \delta\sigma_{\text{NLO}} + \alpha^2 \delta\sigma_{\text{NNLO}} + \mathcal{O}(\alpha^3).$$

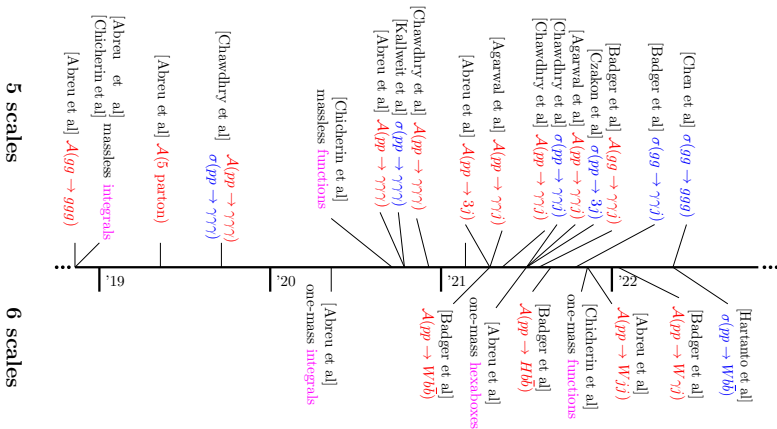
- At fixed order  $\sigma$  is built from phase space integrals over amplitudes:

$$\sigma \sim \int d\Phi |\mathcal{A}|^2.$$

For-high precision, we need **multi-loop** scattering amplitudes.

# Large Progress in Multi-Scale NNLO Calculations

In hadron-collider context, we have witnessed large advances:



Can we also apply these advances to lepton colliders?

# Calculating Multi-Scale Loop Amplitudes

# The Modern Perspective and Challenges

**Aim:** efficient, stable evaluation of amplitude over phase space.

(e.g. at worst 4 digits precision).

- Too many integrals  $\Rightarrow$  write  $A^{(2)}$  in terms of **master integrals**:

$$A^{(2)}(p_1, \dots, p_5) = \sum_k \underbrace{\mathcal{C}_k(p_1, \dots, p_5)}_{\text{rational functions}} \underbrace{\mathcal{I}_k(s_{12}, s_{23}, s_{34}, s_{45}, s_{51})}_{\text{master integrals}}.$$

- Coefficients  $\mathcal{C}_k$  are process specific, integrals  $\mathcal{I}_k$  are universal.

## Challenges

- $\mathcal{I}_k$ : analytic complexity  $\Rightarrow$  **efficient, stable** evaluation difficult.
- $\mathcal{C}_k$ : intermediate expression swell, **poorly understood structure**.

# Master Integrals

# Master Integrals – the Problems

**Aim:** As  $\mathcal{I}_k$  must be combined with  $\mathcal{C}_k$ , need strong control of **precision**.

[Monte Carlo hard as low control over precision.]

## Roadblocks to a “good representation”:

- Integrals are special functions with complicated **branch cuts**.  
⇒ Analytic continuation to physical region is complex.
- With many scales, one finds a myriad of **square root** branch cuts.  
⇒ Limited use of known special functions.
- With massive fields, **special functions** are poorly understood.

# The Workhorse – Differential Equations

## Framework:

- Feynman integrals can be computed from their differential equations

$$d\mathcal{I}_k = \mathbf{M}_{kl}(\epsilon, \vec{s})\mathcal{I}_l.$$

[Gehrmann, Remiddi '01]

- For (pure) polylogarithmic integrals,  $\mathbf{M}$  in “canonical form” [Henn '13].

$$\mathbf{M}_{kl}(\epsilon, \vec{s}) = \underbrace{\epsilon}_{\text{regulator}} \sum_{\alpha, l} \underbrace{M_{kl}^{\alpha}}_{\text{rational numbers}} d \log \left( \underbrace{W_{\alpha}[\vec{s}]}_{\text{algebraic functions}} \right).$$

## The steps:

- 1 Find pure  $\mathcal{I}_k$ . (Many approaches, though unsolved problem).

- 2 Construct alphabet from cut differential equations.

[Abreu, Ita, Moriello, BP, Tschernow, Zeng '20]

- 3 Compute  $M_{kl}^{\alpha}$  via Ansatz, boundary condition from physics.

[Abreu, BP, Zeng '18] [Abreu, Ita, Moriello, BP, Tschernow, Zeng '20] [Liu, Ma '21]

Open question: How to canonicalize beyond polylogarithms?



# Numerically Solving the Differential Equation

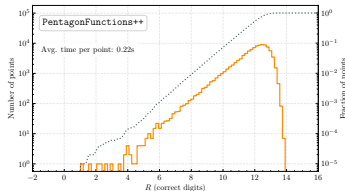
**Aim:** numerically evaluate  $\mathcal{I}_j$  in  $\epsilon$  expansion,  $\mathcal{I}_j(\vec{p}, \epsilon) = \sum_{k=-4}^0 d_{j,k} \epsilon^k h_k + \mathcal{O}(\epsilon)$ .

**Pentagon functions:** [Gehrmann, Henn, Ito Presti '18]

- **Dedicated** iterated integral code:

$$h_k \sim \int_0^1 d \log(W_n[t_n]) \cdots \int_0^{t_2} d \log(W_1[t_1]).$$

- Relations between  $h_k$  known.
- Very efficient, only dynamical scales.



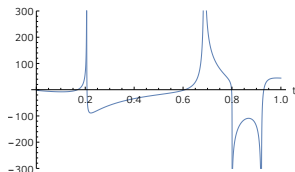
[Chicherin, Sotnikov '20]

**Series expansions:**

- Patch together  $h_k$  from series solutions:

$$h_k \sim \sum_{j_1, j_2} (t - t_0)^{j_1/2} \log(t - t_0)^{j_2}.$$

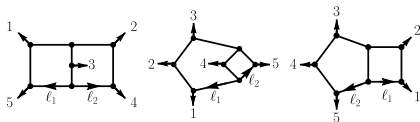
- Implementations: DiffExp/AMFlow.  
[Hidding '20] [Liu, Ma '22]
- Proof-of-concept, promising for masses.



[Abreu, Ita, Moriello, BP, Tschernow, Zeng '20]

# Master Integrals – State of the Art (Two-Loop Five-Point)

- Five-point Massless



[Papadopoulos, Tommasini, Wever '15]

[Gehrmann, Henn, lo Presti '18]

[Abreu, Page, Zeng '18]

[Chicherin, Gehrmann, Henn, lo Presti, Mitev, Wasser '18]

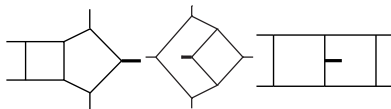
[Abreu, Dixon, Herrmann, Page, Zeng '18]

[Chicherin, Gehrmann, Henn, Wasser, Zhang, Zoia '18]

[Gehrmann, Henn, lo Presti '18]

[Chicherin, Sotnikov '20]

- Five-point One-Mass



[Abreu, Ita, Moriello, Page Tschernow, Zeng '20]

[Abreu, Ita, Page, Tschernow '21]

[Papadopoulos, Tommasini, Wever 15]

[Canko Papadopoulos, Syrrakos 20]

[Kardos, Papadopoulos, Smirnov, Syrrakos, Wever 19]

[Chicherin, Sotnikov, Zoia '21]

[Abreu, Chicherin, Ita, Page, Sotnikov, Tschernow, Zoia. to appear]

Great progress for many dynamical scales, no fixed scales.

# Integral Coefficients

# The Analytic Reconstruction Approach

- Given known  $h_k$ , must now determine  $C_{j,k}$ .

$$\mathcal{A} = \sum_{j=-4}^0 e^j \sum_{k \in B} C_{j,k} h_k + O(\epsilon).$$

- Strategy: reconstruct analytic  $C_{j,k}$  from numerical samples (mod  $p$ ).

[Schabinger, von Manteuffel '14; Peraro '16]

$$C_{j,k}(p_1, \dots, p_n) = \sum_{l=1}^N c_{j,k,l} a_{j,k,l}(p_1, \dots, p_n), \quad c_{j,k,l} \in \mathbb{Q}.$$

Two steps:

- 1 Construct linear Ansatz.
- 2 Gather numerical data and fix unknowns from numerical data.

Runtime cost  $\sim N \times T_{\text{eval}}$ .

# Ansätze for Rational Functions

- Coefficients  $C_{j,k}(p_1, \dots, p_n)$  are rational functions of kinematics.

$$C_{j,k}(p_1, \dots, p_n) = \frac{\mathcal{N}_{j,k}(p_1, \dots, p_n)}{D_{j,k}(p_1, \dots, p_n)}$$

- Denominators simple when we know the integrals: product of letters.

$$D_{j,k} = \prod_l W_l^{q_{jkl}}$$

[Abreu, Dormans, Ita, Febres Cordero, Page '18]

- Simpler problem – reconstruct polynomial  $\mathcal{N}_{j,k}$ . Multiple techniques:

- Newton reconstruction.

[Peraro '16]

- Univariate partial fractions.

[Badger, Hartanto, Zoia '21]

- Vandermonde reconstruction.

[Abreu et al '21]

- $p$ -adic numbers.

[de Laurentis, BP '22]

Major challenge: Constructing **compact** Ansätze.

# Gathering Numerical Data

Two major paths:

## “The Feynman way”

- Feynman diagrams via QGRAF.  
Integrands via FORM.
- Projectors + form factors.  
[Chen '19; Peraro, Tancredi '19 '20]
- Finite field IBP reduction.  
FiniteFlow [Peraro '19]

## “Numerical Unitarity”

- Parameterize integrand with  
master/surface integrands.  
[Ita '15]
- Generate integrand from  
unitarity cuts.

Use to compute coefficients numerically on a phase-space point  $\vec{p}_0$ :

$$\mathcal{A}(\vec{p}_0) = \sum_i C_j(\vec{p}_0, \epsilon) \mathcal{I}_j(\vec{p}, \epsilon) = \sum_{j=-4}^0 \epsilon^j \sum_{k \in B} \underbrace{C_{j,k}(\vec{p}_0)}_{\text{numerical data}} h_k(\vec{p}) + \mathcal{O}(\epsilon).$$

Principal bottleneck: IBP reduction.

# $pp \rightarrow Vjj$ Case Study

$$\mathcal{A}_g = \text{Diagram 1} + \dots, \quad \mathcal{A}_Q = \text{Diagram 2} + \dots$$

- Leading-colour helicity-amplitude finite remainders.
- Planar pentagon functions from [Chicherin, Sotnikov, Zoia '21].
- Ansatz: Univariate partial fractions, Vandermonde sampling.
- $\mathcal{O}(10^6)$  samples for most complex  $\mathcal{A}_n$ . Sample time:  $\mathcal{O}(3)$  minutes.

Analytic results at <http://www.hep.fsu.edu/~ffebres/W4partons>



# Towards $e^+e^-$ Colliders

# From $pp \rightarrow Vjj$ to $e^+e^- \rightarrow 4j$ (QCD corrections)

$pp \rightarrow Vjj$  amplitudes computed in terms of QCD current  $\mathcal{A}^\mu$ :

$$\mathcal{A} = \mathcal{A}^\mu J_\mu, \quad J_\mu \sim \frac{\bar{u}(p_6)\gamma_\mu v(p_5)}{s_{56} - M_V^2 + i\Gamma_V M_V}.$$

## Steps to achieve crossing:

- Attach appropriate leptonic current.
- **Analytically continue** from  $2 \rightarrow 3$  to  $1 \rightarrow 4$  kinematics.

## Nonplanar contributions:

- Terms subleading in  $N_c$  not included, likely suppressed.
- Contribution proportional to  $\sum_f Q_f$  potentially required.  
(Non-planar integrals: [Abreu, Chicherin, Ita, Tschernow, BP, Sotnikov, Zoia. to appear])

## No progress required:

- $e^+e^- \rightarrow 4j$  at  $\mathcal{O}(\alpha_s^2)$ . [just crossing!]

## Incremental progress required:

- $e^+e^- \rightarrow 3j$  at  $\mathcal{O}(\alpha_s^3)$ . [+ 1 dynamical scales]
- $e^+e^- \rightarrow QQj$  at  $\mathcal{O}(\alpha_s^2)$  [elliptic, low # dynamical scales].
- $e^+e^- \rightarrow 5j$  at  $\mathcal{O}(\alpha_s^2)$  [+ 1 leg, + 3 dynamical scales].

## Breakthrough required? [Elliptic integrals, many dynamical scales.]

- $e^+e^- \rightarrow \nu\nu H$  at  $\mathcal{O}(\alpha^2)$ .
- $e^+e^- \rightarrow l^+l^- H$  at  $\mathcal{O}(\alpha^2)$ .

## Status

- Loop amplitudes remain a major bottleneck to precise predictions.
- Large progress in multi-scale amplitudes will benefit  $e^+e^-$  colliders.

## Lessons learned

- Large final-state phase space  $\Rightarrow$  need efficient/stable amplitudes.
- Calculations based on structural understanding of Feynman integrals.

## The Next Frontier?

- Large improvements desired for cases with massive particles (elliptics).
- Time to go fully numerical?