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# **Running quark masses at future** $e^+e^-$ colliders

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#### The running of the Standard Model parameters

• Standard Model parameters are subject to renormalization, so they depend on the renormalization scale  $\mu_R$  according to the Renormalization Group Equations:

$$\beta(\alpha) = \mu_R^2 \frac{\partial \alpha}{\partial \mu_R^2} = -\sum_{i=0}^{\infty} \left(\beta_i \alpha^{i+2}\right)$$
$$\gamma(\alpha) = -\mu_R \frac{1}{m} \frac{\partial m}{\partial \mu_R^2} = \sum_{k=0}^{\infty} \left(\gamma_k \alpha^{k+1}\right)$$

- They must be determined experimentally, and the observed value is tied to the characteristic energy of the process it is involved with
- Calculations at 5-loop precision enable a precise determination of the evolution of the strong coupling



#### Quark masses running: the top quark

- First experimental investigation of the running of the top  $\overline{MS}$  mass with pp collision data at  $\sqrt{s} = 13$  TeV, collected by CMS experiment at the LHC
- Running extracted from measurements of the differential  $\sigma_{t\bar{t}}$ lacksquareas a function of  $m_{t\bar{t}}$
- NLO QCD calculations to simultaneously extract  $m_t(\mu)$  at each bin at parton level
- The running hypothesis is tested **considering**  $\mu = m_{t\bar{t}}$

 $\rightarrow$  Top mass running in agreement with SM RGE within  $1.1\sigma$ 

 $\rightarrow$  No-running scenario ruled out at above %95 CL



### Top quark mass running in $e^+e^-$ colliders

- Radiative  $e^+e^- \rightarrow t\bar{t} + X + \gamma$  events above the  $t\bar{t}$  production threshold  $> 2m_t$  allow:
  - Precise determination of  $m_t$  in the MSR (R) mass scheme (with  $R < m_t$ ) with  $N^3LO \ QCD + N^2LL$  prediction
  - Measurements of  $m_t$  at different renormalization scales
- Study performed in two operating and realistic scenarios: ILC and CLIC

cms energy	CLIC, $\sqrt{s}$	= 380  GeV	ILC, $\sqrt{s}$ :	= 500  GeV
luminosity $[fb^{-1}]$	500	1000	500	4000
statistical	$140{ m MeV}$	$90\mathrm{MeV}$	$350\mathrm{MeV}$	$110\mathrm{MeV}$
theory	$46\mathrm{MeV}$		$55{ m MeV}$	
lum. spectrum	$20{ m MeV}$		$20{ m MeV}$	
photon response	$16{ m MeV}$		$85{ m MeV}$	
total	$150\mathrm{MeV}$	$110\mathrm{MeV}$	$360\mathrm{MeV}$	$150\mathrm{MeV}$

arXiv:1912.01275



R-evolution of top quark mass demonstrated for  $R < m_t$  in the 500 GeV ILC program



#### Bottom quark mass measurements

Running of the bottom quark mass also studied from low-energy measurements ullet



- Bottom quark mass production threshold
- Most precise determinations from mass of bottomonium bound states in  $e^+e^- \rightarrow$  hadrons
- World combination from PDG in  $\overline{MS}$  scheme:

$$m_b(m_b) = 4.18^{+0.03}_{-0.02}$$



- Hadronic decays of Z bosons
- Event-shapes and jet-rates sensitive to subleading mass effects
- A combination of the most precise measurements in *MS* scheme yields:

$$m_b(m_Z) = 2.82 \pm 0.28 \ GeV$$

#### $m_h$ from Higgs boson decay rates (I)

Higgs discovery and observation of its decay into bottom quarks enabled a new measurement at the Large Hadron Collider

We profit from the ratio of branching ratios:



\* Dependence of  $\Gamma(H \to b\bar{b})$  with  $m_b$  obtained assuming SM Yukawa couplings and  $m_b \ll m_H$ :

$$\Gamma(H \to b\bar{b}) = \frac{3G_F m_H}{4\sqrt{2}\pi} m_b(\mu)^2 \left(1 + \delta_{ew}\right) \times \left(1 + \delta_{QCD} + \delta_t + \delta_{mix}\right)$$

**Excellent convergence** of the perturbative calculation at  $\mu_R = m_H$ , so small higher-order corrections expected

$$T(H \to b\bar{b})$$
  
 $T(H \to ZZ)$ 



#### $m_h$ from Higgs boson decay rates (II)

several values of  $m_b(m_H)$ :

$$\frac{\Gamma(H \to b\bar{b})}{\Gamma(H \to ZZ)} = 2.82 \frac{m_b^2}{GeV^2} - 0.0014 \frac{m_b^4}{GeV^4} + \mathcal{O}(m_b^6)$$

• ATLAS and CMS measurements at 139  $fb^{-1}$  and 35  $fb^{-1}$ :

$$\Gamma^{b\bar{b}}/\Gamma^{ZZ}$$

 $0.87^{+0.22}_{-0.17} (stat)^{+0.18}_{-0.12} (syst) \qquad 2.61^{+0.32}_{-0.27} (stat)^{+0.26}_{-0.19} (syst) GeV$ 

 $0.84^{+0.27}_{-0.21} (stat)^{+0.26}_{-0.17} (syst) \qquad 2.57^{+0.39}_{-0.35} (stat)^{+0.37}_{-0.28} (syst) GeV$ 

\* Individual partial widths  $\Gamma(H \rightarrow b\bar{b}, ZZ)$  estimated separately. The ratio is built and parametrized for

 $m_b(m_H)$ 

$$m_b(m_H) = 2.60^{+0.36}_{-0.31} GeV$$







✤ A modified RGE SM evolution through x parameter:

$$m(\mu; x, m_b(m_b)) = x \left[ m_b^{RGE}(\mu, m_b(m_b)) - m_b(m_b) \right] + x \left[ m_b^{RGE}(\mu, m_b) \right] + x \left[ m_b^{RGE}(\mu$$

 $x = 0 \rightarrow$  No-running scenario

 $x = 1 \rightarrow SM$  scenario

 $m_b(m_b)$ 



✤ A modified RGE SM evolution through x parameter:

$$m(\mu; x, m_b(m_b)) = x \left[ m_b^{RGE}(\mu, m_b(m_b)) - m_b(m_b) \right] + 1$$

\* Taking  $\mu_0 = m_b$ , calculate  $m_b(m_Z)$  and  $m_b(m_H)$  from different starting values of  $m_b(m_b)$  and x at 5-loop precision

$$\chi 2(m_b(m_b), x) = \frac{\sum_{\mu_i} (m_b^{exp}(\mu_i) - m(\mu_i; x, m_b(m_b)))^2}{\sigma_i^2}$$





$$m(\mu; x, m_b(m_b)) = x \left[ m_b^{RGE}(\mu, m_b(m_b)) - m_b(m_b) \right] + 1$$

precision





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\* Taking  $\mu_0 = m_b$ , calculate  $m_b(m_Z)$  and  $m_b(m_H)$  from different starting values of  $m_b(m_b)$  and x at 5-loop precision

$$\begin{split} m_b(m_b) &= 4.18^{+0.03}_{-0.02} \; GeV \; , \\ x &= 1.08 \pm 0.15 \; (exp) \; \pm 0.05 \; (\alpha_S) \end{split}$$

- No-running scenario ruled out with  $> 5\sigma$
- SM running confirmed within  $1\sigma$



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#### Prospects in future colliders

Collider	Channel	Expected experimental unc. on channel meas.	Expected experimental unc. on $m_b(m_H)$
HL-LHC	$BR(H \rightarrow b\bar{b})$	4 %	±63 MeV
ILC:250	$\frac{BR(H \to b\bar{b})}{BR(H \to WW)}$	0.86 %	$\pm 12 MeV$
ILC:250+500		0.47 %	$\pm 6 MeV$

 $\rightarrow$  Very competitive measurements are possible with this method. The prospects for theory uncertainties need to be carefully assessed.



#### Summary

Quark masses evolution with the energy scale subject of study at hadron and  $e^+e^-$  colliders

- Running of the top quark mass inspected at the LHC with NLO QCD precision
- Study from radiative events at  $e^+e^-$  colliders shows a better precision in top quark mass measurement and provides observation of its running
- magnitud in the ILC program

First determination of the bottom quark mass at the Higgs mass scale at the LHC presented. The running is observed with  $> 5\sigma$  including this measurement to the existing low-energy ones

• Excellent prospects for this method at future  $e^+e^-$  colliders: uncertainty reduced by one order of

### **Bonus slides**

## $m_h(m_7)$ combination from LEP and SLC

experiment

 $m_b(m_Z)$  |GeV| ALEPH[14]  $3.27 \pm 0.22$  (stat.)  $\pm 0.44$  (syst.)  $\pm 0.16$  (theo.) DELPHI[16]  $2.85 \pm 0.18$  (stat.)  $\pm 0.23$  (syst.)  $\pm 0.12$  (theo.) OPAL[15]  $2.67 \pm 0.03$  (stat.)  $^{+0.29}_{-0.37}$  (syst.)  $\pm 0.19$  (theo.) SLD[12, 13]  $2.56 \pm 0.27$  (stat.)  $^{+0.28}_{-0.38}$  (syst.)  $^{+0.49}_{-1.48}$  (theo.)