

Charmed baryons results at Belle

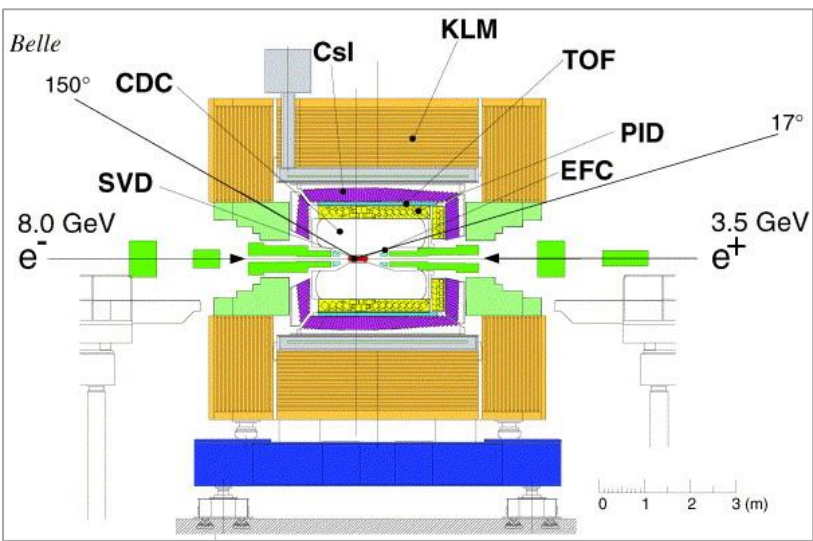
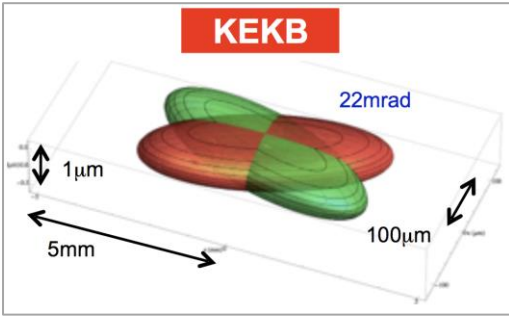
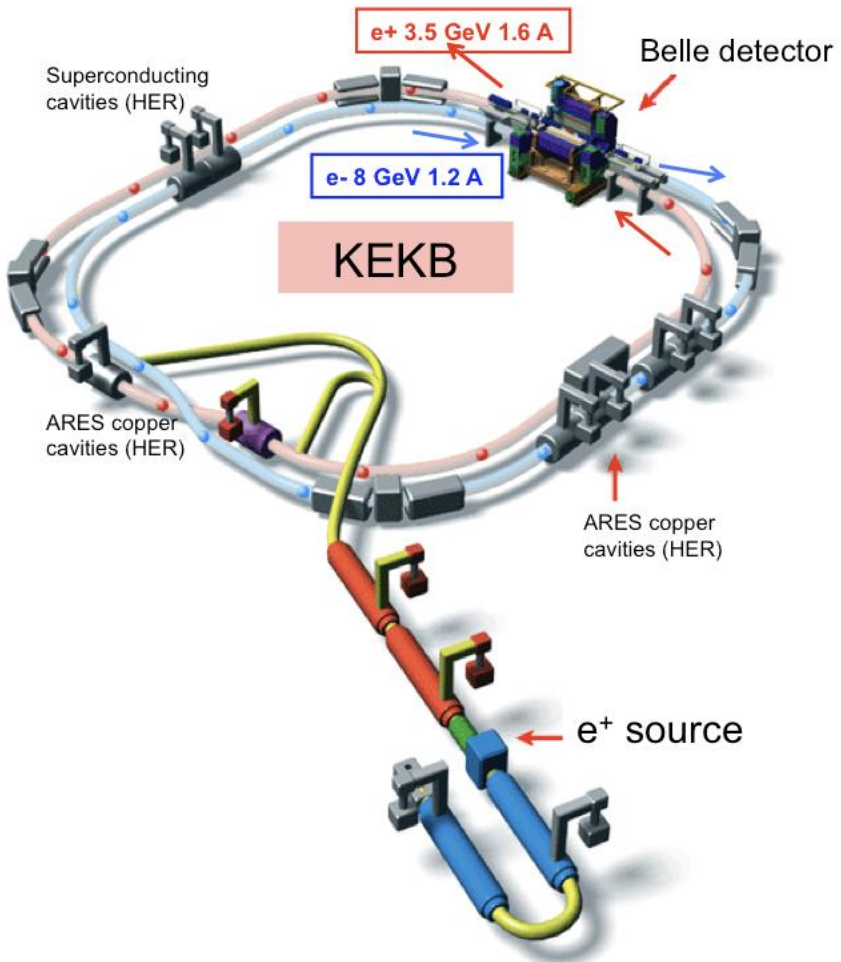
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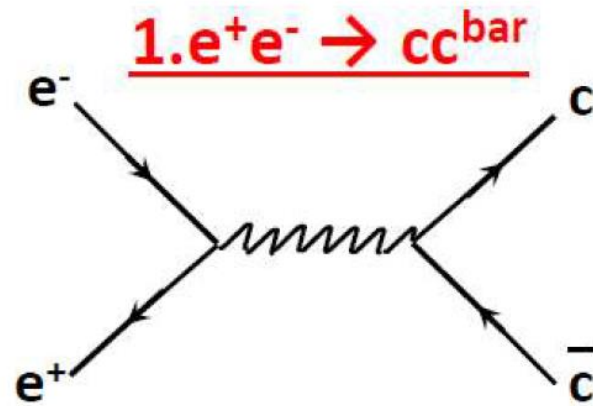
The 13th International Workshop on e^+e^- collisions from Phi to Psi

2022-08-15

Belle experiment

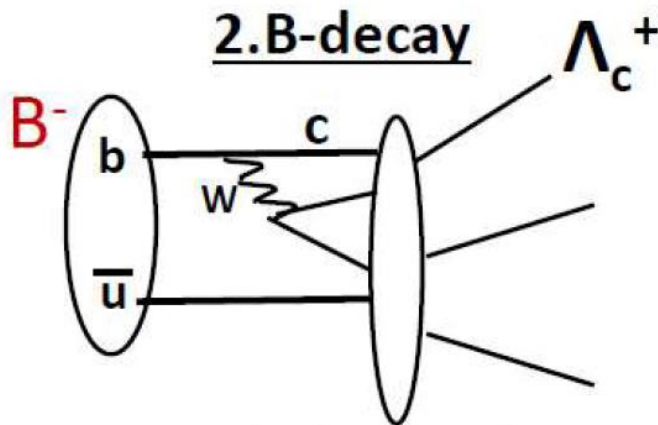


Charmed baryon production at B-factory



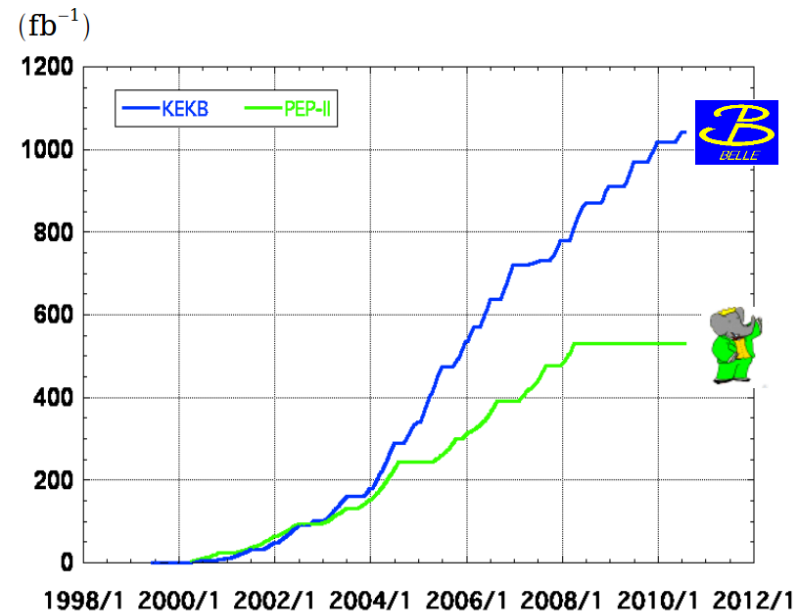
Baryons produced via fragmentation

- Charmed baryons – rather direct
- Hyperons – later stage of fragmentation



- B meson is efficiently produced via $\Upsilon(4S)$
- Once bottom is produced, it favorably decays into charm

Integrated luminosity of B factories



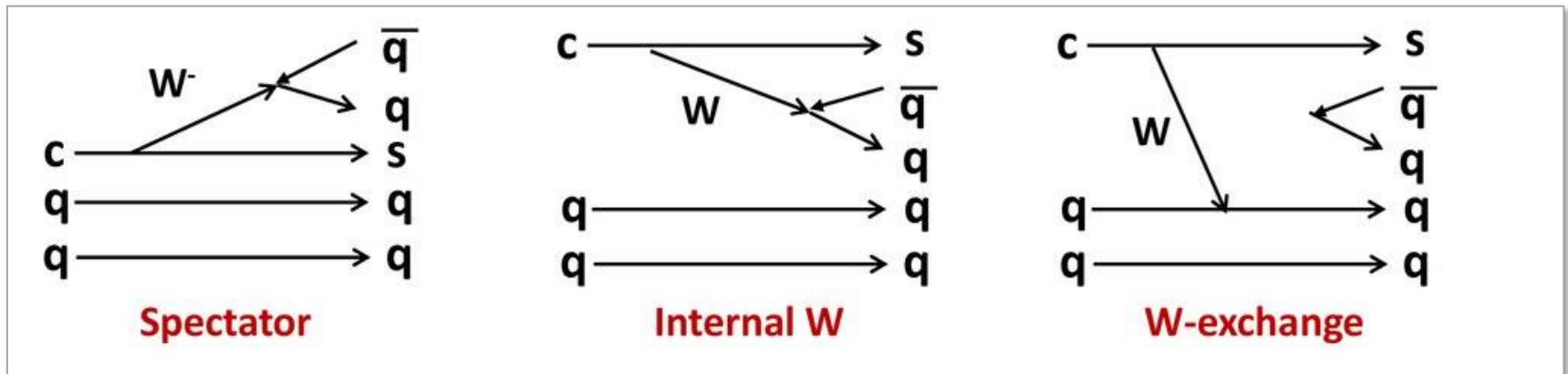
> 1 ab⁻¹
On resonance:
 $\Upsilon(5S)$: 121 fb⁻¹
 $\Upsilon(4S)$: 711 fb⁻¹
 $\Upsilon(3S)$: 3 fb⁻¹
 $\Upsilon(2S)$: 25 fb⁻¹
 $\Upsilon(1S)$: 6 fb⁻¹
Off reson./scan:
 ~ 100 fb⁻¹

~ 550 fb⁻¹
On resonance:
 $\Upsilon(4S)$: 433 fb⁻¹
 $\Upsilon(3S)$: 30 fb⁻¹
 $\Upsilon(2S)$: 14 fb⁻¹
Off resonance:
 ~ 54 fb⁻¹

Huge statistics

The charmed baryon physics

- The weak decay of charmed baryon is not understood well.
 - Three diagrams contribute to in the tree level, strength of each is not known
- Ground state charmed baryon is a good laboratory to study strange baryons as decay proceed via $c \rightarrow s$ transition.



- Belle has collected $\sim 1 \text{ ab}^{-1} e^+e^-$ data samples (mainly at $\Upsilon(4S)$).
 - $10^9 e^+e^- \rightarrow c\bar{c}$ samples
 - $7.7 \times 10^8 B\bar{B}$ samples
- Huge data sample enable to study various charmed baryons.

Charmed baryon resent results at Belle

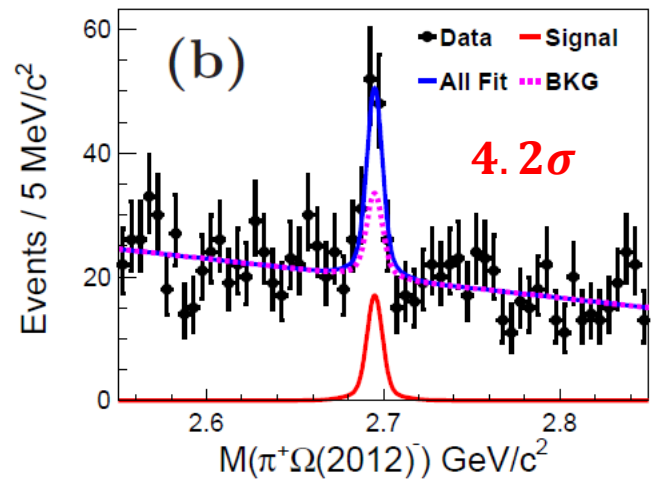
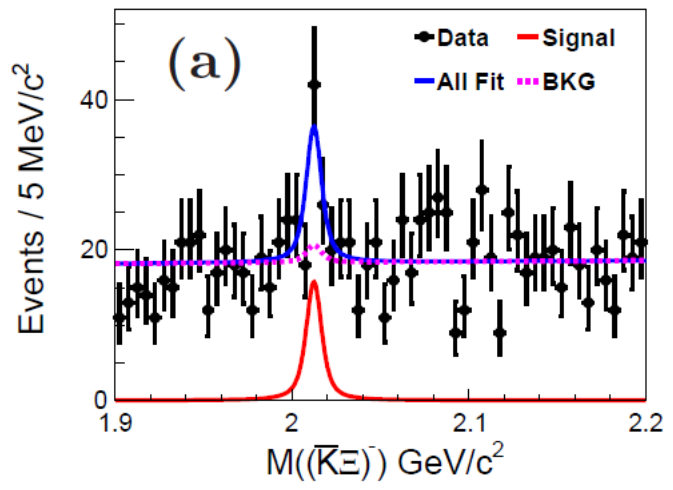
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Process	Reference	Physics Covered
$\Omega_c^0 \rightarrow \pi^+ [\Omega(2012)^- \rightarrow \bar{K}\Xi]$	PRD 104, 052005 (2021)	Evidence: 4.2σ
$\Lambda_c^+ \rightarrow p\pi^0/p\eta$	PRD 103, 072004 (2021)	SCS decays
$\Lambda_c^+ \rightarrow p\omega$	PRD 104, 072008 (2021)	SCS decay, 3pi mode
$\Lambda_c^+ \rightarrow p\eta'$	JHEP 03, 090 (2022)	SCS decay, 5.4σ
$\Lambda_c^+ \rightarrow \eta\Lambda\pi^+$	PRD 103, 052005 (2021)	$\Lambda(1670), \Sigma(1385)^+$
$\Lambda_c^+ \rightarrow \Lambda h^+/\Sigma^0 h^+$	Belle note #625	Direct CPV search, Br, Acp
$\Lambda_c^+ \rightarrow \Sigma^+\pi^0/\Sigma^+\eta/\Lambda_c^+ \rightarrow \Sigma^+\eta'$	Belle note #626	Br, Acp
$\Lambda_c^+ \rightarrow \Sigma^+\gamma/\Xi_c^0 \rightarrow \Xi^0\gamma$	arXiv: 2206.12517	Weak radiative decays, UL
$\Xi_c^0 \rightarrow \Lambda K_S^0/\Sigma^0 K_S^0/\Sigma^+ K^-$	PRD 105, L011102 (2022)	Br, CF decays
$\Xi_c^0 \rightarrow \Xi^- \ell^+ \nu_\ell$	PRL 127, 121803 (2021)	Br, Acp, LFU
$\Xi_c^0 \rightarrow \Lambda_c^+ \pi^-$	arXiv: 2206.08527	Heavy-flavor-conserving decay
$\Omega_c^0 \rightarrow \Omega^- \ell^+ \nu_\ell$	PRD 105, L091101 (2022)	Br, LFU
$\Sigma_c(2455)^+, \Sigma(2520)^+$	PRD 104, 052003 (2021)	mass and widths
$\Sigma_c(2455)^{0,++}\pi^\pm$	arXiv: 2206.08822	New excited charmed baryon, 4.2σ

Evidence for $\Omega_c^0 \rightarrow \pi^+ [\Omega(2012)^- \rightarrow \bar{K}\Xi^-]$

- Searching for new production model is very important to understand the nature of $\Omega(2012)^-$
- A theoretical study of the $\Omega(2012)^-$ in the nonleptonic weak decays of $\Omega_c^0 \rightarrow \pi^+ (\bar{K}\Xi)^-$ was reported [1]
- We do the search in both $K_S^0 \Xi^-$ and $K^- \Xi^0$ final states

PRD 104, 052005 (2021)



Combined

$$\begin{aligned}
 &= \frac{\mathcal{B}(\Omega_c^0 \rightarrow \pi^+ \Omega(2012)^-) \times \mathcal{B}(\Omega(2012)^- \rightarrow (\bar{K}\Xi)^-)}{\mathcal{B}(\Omega_c^0 \rightarrow \pi^+ \Omega^-)} \\
 &= \frac{N_{\text{sig.}}^{\text{obs}} \times \epsilon_{\pi^+ \Omega^-}}{N_{\pi^+ \Omega^-}^{\text{obs}} \times (f_1 \times \epsilon_1 \times \mathcal{B}_1 + f_2 \times \epsilon_2 \times \mathcal{B}_2)} \\
 &= 0.220 \pm 0.059(\text{stat.}) \pm 0.035(\text{syst.}),
 \end{aligned}$$

[1] PRD 102, 076009

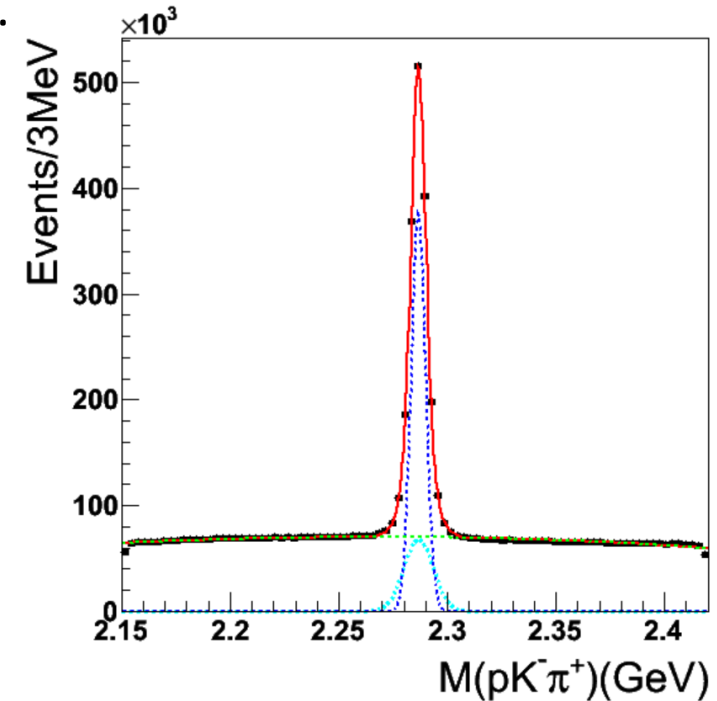
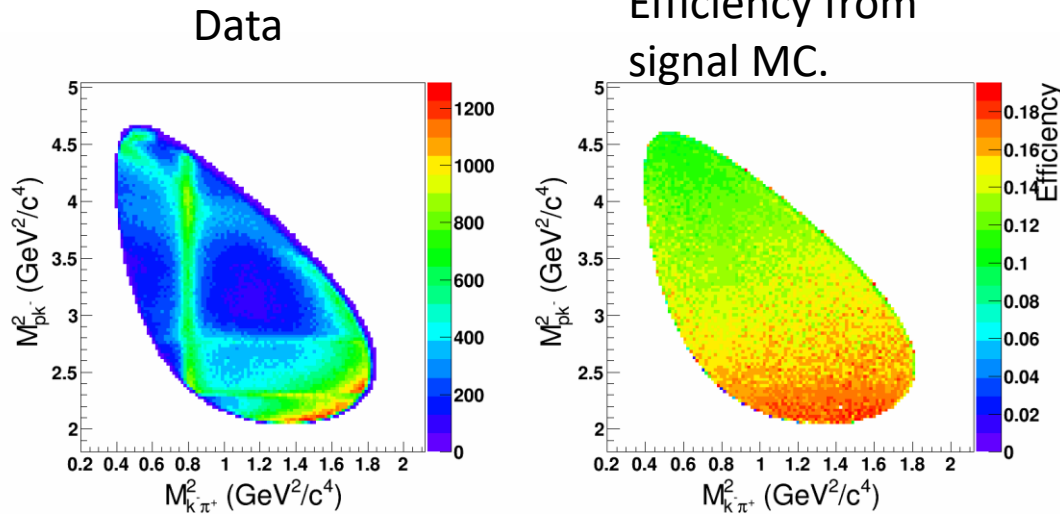
Measurement of $\Lambda_c^+ \rightarrow p\pi^0/p\eta$

A method of branching ratio with respect to Cabibbo-Favored decay $\Lambda_c^+ \rightarrow pK^-\pi^+$ (reference mode) is applied to measure the branching fractions of signal decay.

PRD 103, 052005 (2021)

$$\frac{B(\text{Signal})}{B(\text{CF})} = \frac{N^{\text{obs}}(\text{Signal})}{\epsilon^{\text{MC}}(\text{Signal})} \times \frac{\epsilon^{\text{MC}}(\text{CF})}{N^{\text{obs}}(\text{CF})}$$

$\Lambda_c^+ \rightarrow pK^-\pi^+$ efficiency estimation: Dalitz method.



Fit to $M(pK^-\pi^+)$ from data using double Gaussian + second-order polynomial

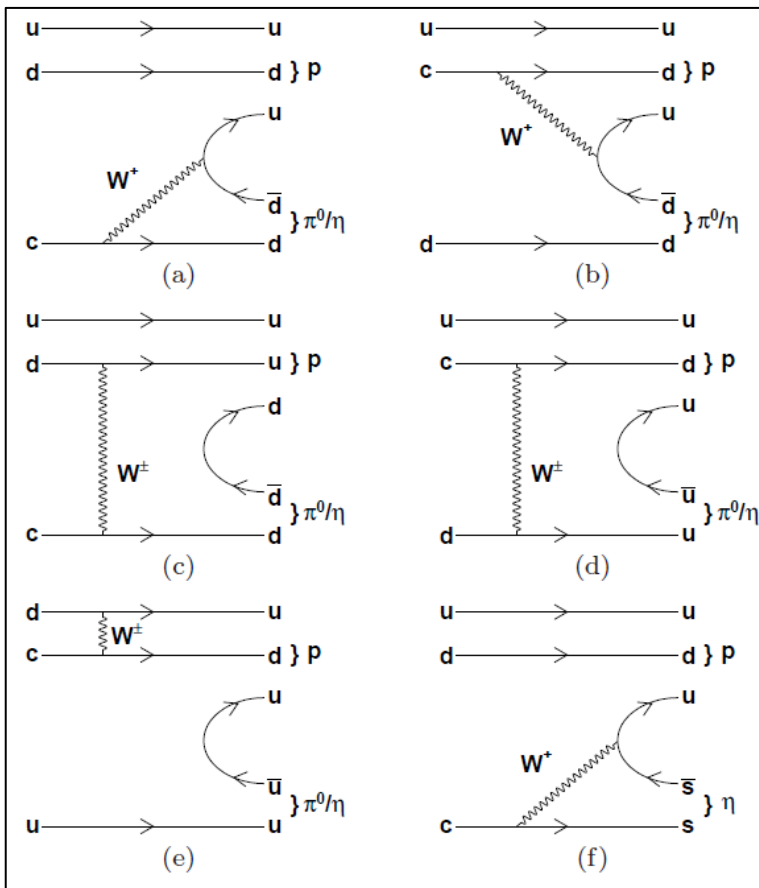
$$\epsilon = \sum s_i / \sum_j (s_j / \epsilon_j) = (14.06 \pm 0.01)\%$$

Yield: 1476200 ± 1560
 $\chi^2/ndf=1.06$

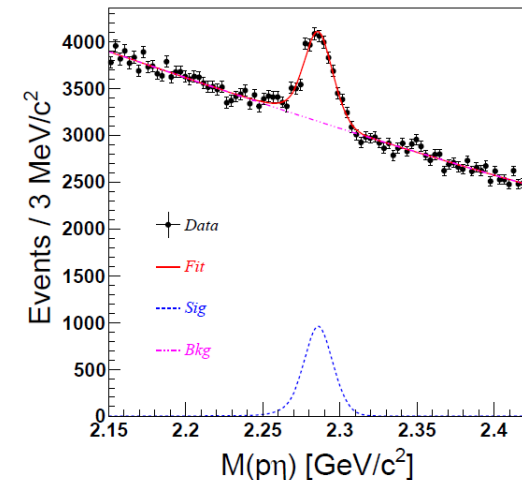
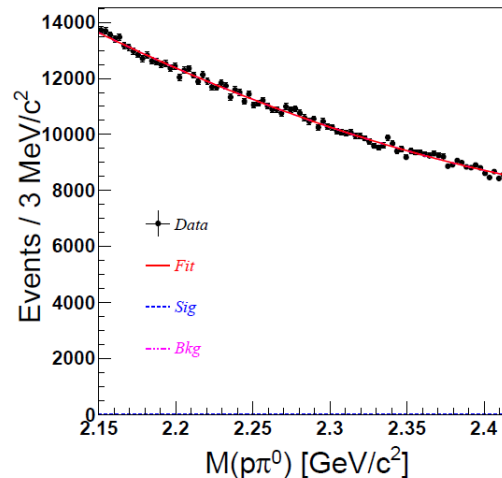
Measurement of $\Lambda_c^+ \rightarrow p\pi^0 / p\eta$

PRD 103, 072004 (2021)

- The W-boson exchange mechanism plays an important role in the decay of charmed baryons.
- Thus, measuring the branching fractions of these two Singly Cabibbo-Suppressed decays will help elucidate the decay mechanism of charmed baryons.



(a, b) Internal W emission
(c, d, e) W exchange
(f) Internal W emission



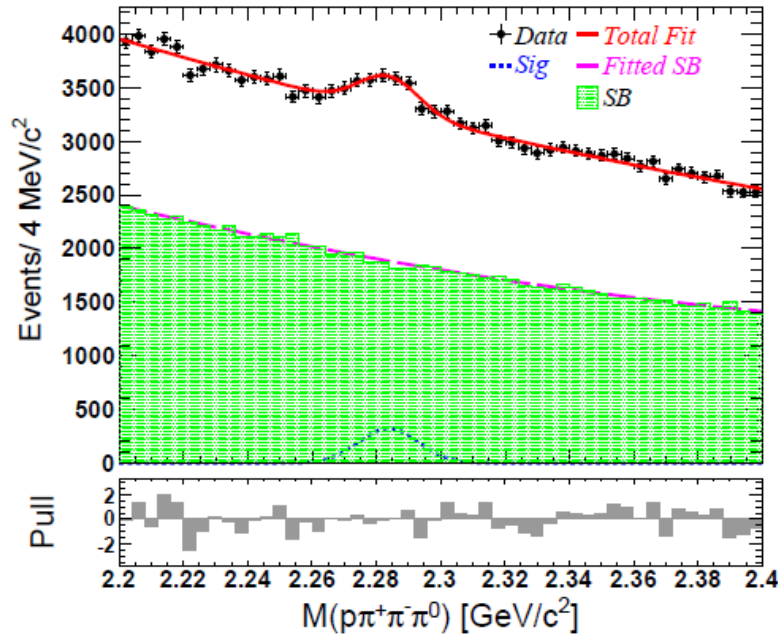
$$Br(\Lambda_c^+ \rightarrow p\pi^0) < 8 \times 10^{-5} \text{ @ 90\% C.L.}$$

$$Br(\Lambda_c^+ \rightarrow p\eta) = (1.42 \pm 0.05 \pm 0.11) \times 10^{-3}$$

Observation of $\Lambda_c^+ \rightarrow p\omega$, $\Lambda_c^+ \rightarrow p\eta'$

- LHCb reported the first observation of a SCS decay $\Lambda_c^+ \rightarrow p\omega[\rightarrow \mu^+\mu^-]$, $\text{Br}(\Lambda_c^+ \rightarrow p\omega) = (9.4 \pm 3.9) \times 10^{-4}$
- We perform same measurement by reconstructing $\omega \rightarrow \pi^+\pi^-\pi^0$ decay

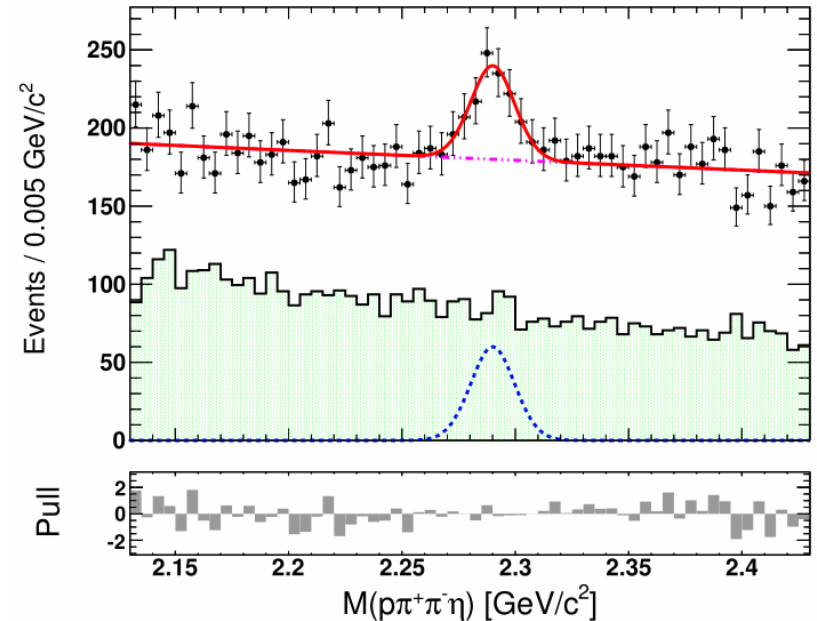
PRD 104, 072008 (2021)



$$\text{Br}(\Lambda_c^+ \rightarrow p\omega) = (8.27 \pm 0.75 \pm 0.62 \pm 0.42) \times 10^{-4}$$

- Most precise measurement
- Consistent with LCHb result

JHEP 03, 090 (2022)



$$\text{Br}(\Lambda_c^+ \rightarrow p\eta') = (4.73 \pm 0.82 \pm 0.47 \pm 0.24) \times 10^{-4}$$

- First observation of Λ_c^+ in $\Lambda_c^+ \rightarrow p\eta'$
- Consistent with the $SU(3)_F$ calculation

Measurements of $\Lambda_c^+ \rightarrow \Lambda h^+, \Sigma^0 h^+ (h^+ = K, \pi)$

- SCS decays of charmed hadrons provide an ideal laboratory for studying CPV [1,2]
- The only observation of CPV in charm sector was made by LHCb $D^0 \rightarrow h^+ h^-$ [3]
- Direct CPV searches in two-body SCS:
 - No good theoretical predication, no experimental measurements

$$A_{\text{raw}}^{\text{corr}}(\Lambda_c^+ \rightarrow \Lambda K^+) - A_{\text{raw}}^{\text{corr}}(\Lambda_c^+ \rightarrow \Lambda \pi^+) = A_{CP}^{\text{dir}}(\Lambda_c^+ \rightarrow \Lambda K^+) - \cancel{A_{CP}^{\text{dir}}(\Lambda_c^+ \rightarrow \Lambda \pi^+)}. \quad A_{\text{raw}} = \frac{N(\Lambda_c^+ \rightarrow f) - N(\bar{\Lambda}_c^- \rightarrow \bar{f})}{N(\Lambda_c^+ \rightarrow f) + N(\bar{\Lambda}_c^- \rightarrow \bar{f})}.$$

- We can also measure the decay asymmetry parameters

$$A_{CP}^{\alpha}(\text{total}) = \frac{\alpha_{\Lambda_c^+} \alpha_- - \alpha_{\bar{\Lambda}_c^-} \alpha_+}{\alpha_{\Lambda_c^+} \alpha_- + \alpha_{\bar{\Lambda}_c^-} \alpha_+}, \quad \frac{dN}{d \cos \theta_{\Lambda}} \propto 1 + \alpha_{\Lambda_c^+} \alpha_- \cos \theta_{\Lambda}$$

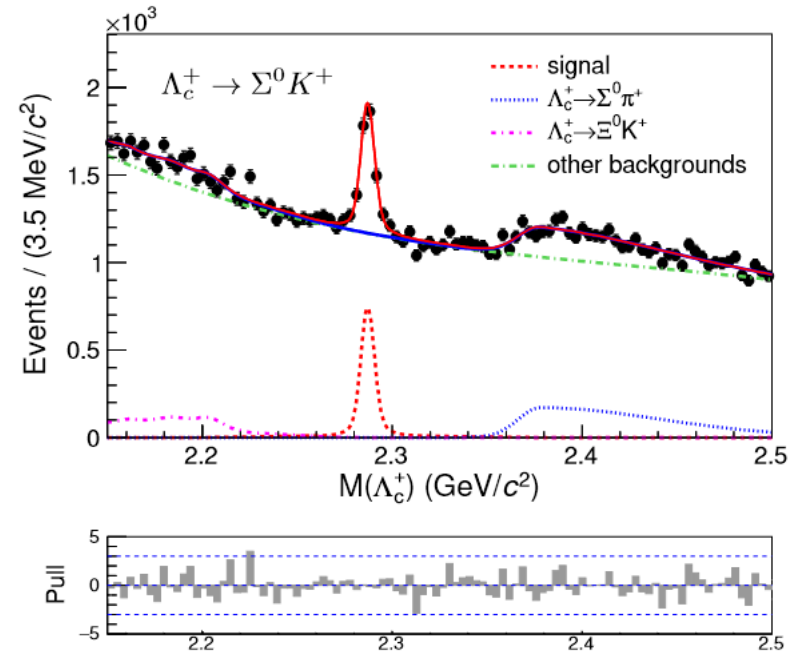
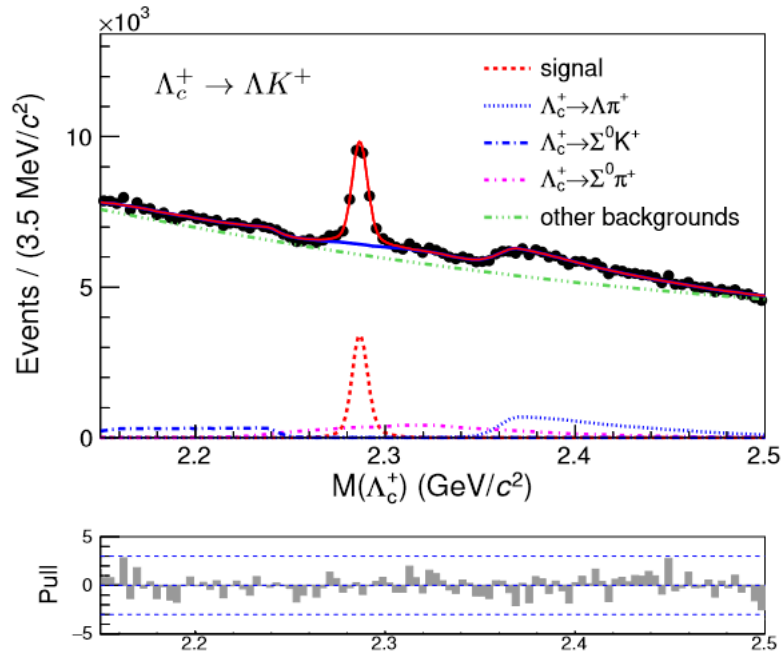
$$A_{CP}^{\tilde{\alpha}}(\text{total}) = A_{CP}^{\alpha}(\Lambda \rightarrow p \pi^-) \quad \frac{dN}{d \cos \theta_{\Sigma^0} d \cos \theta_{\Lambda}} \propto 1 - \alpha_{\Lambda_c^+} \alpha_- \cos \theta_{\Sigma^0} \cos \theta_{\Lambda}$$

θ_{Λ} (θ_{Σ^0}) is the angle between the proton (Λ) momentum vector and the opposite of the Σ^0 (Λ_c^+) momentum vector in the Λ (Σ^0) rest frame; $\alpha_- = 0.7542 \pm 0.0022$ from world average value measured by BESIII.

- And branching fractions $\frac{\mathcal{B}_{\text{sig}}}{\mathcal{B}_{\text{ref}}} = \frac{N_{\text{sig}}/\epsilon_{\text{sig}}}{N_{\text{ref}}/\epsilon_{\text{ref}}}$,

[1] PRD 86, 616 014023 (2012)
 [2] PRD 85, 034036618 (2012)
 [3] PRL 627 122, 211803 (2019)

Measurements of $\Lambda_c^+ \rightarrow \Lambda h^+, \Sigma^0 h^+ (h^+ = K, \pi)$



Direct CPV measurements:

$$A_{CP}^{dir}(\Lambda_c^+ \rightarrow \Lambda K^+) = (+2.1 \pm 2.6 \pm 0.1) \times 10^{-2}$$

$$A_{CP}^{dir}(\Lambda_c^+ \rightarrow \Sigma^0 K^+) = (+2.5 \pm 5.4 \pm 0.4) \times 10^{-2}$$

The first direct CP asymmetry measurement for SCS two-body decays of charmed baryons.

Branching fraction measurements:

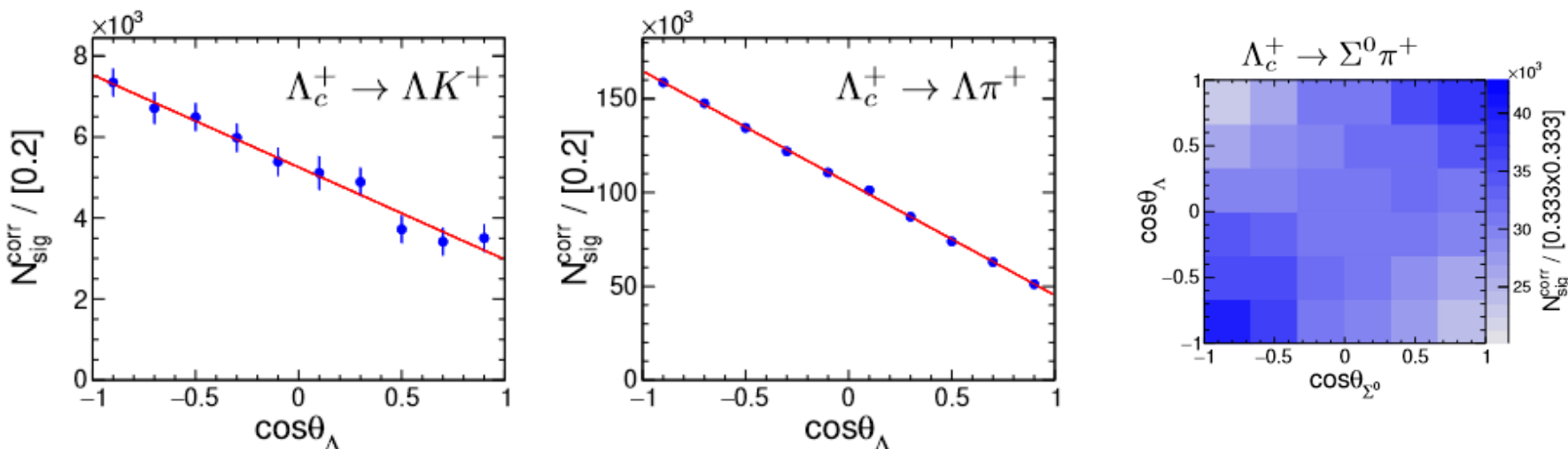
$$Br(\Lambda_c^+ \rightarrow \Lambda K^+) = (6.57 \pm 0.17 \pm 0.11 \pm 0.35) \times 10^{-4}$$

$$Br(\Lambda_c^+ \rightarrow \Sigma^0 K^+) = (3.58 \pm 0.19 \pm 0.06 \pm 0.19) \times 10^{-4}$$

Significantly improved precision

Preliminary result

Measurements of $\Lambda_c^+ \rightarrow \Lambda h^+, \Sigma^0 h^+ (h^+ = K, \pi)$



Averaged decay asymmetry parameters:

- **First measurements** of K^+ modes
- Improved precision of π^+ modes

$$\begin{aligned} \alpha_{\text{avg}}(\Lambda_c^+ \rightarrow \Lambda K^+) &= -0.585 \pm 0.049 \pm 0.018, \\ \alpha_{\text{avg}}(\Lambda_c^+ \rightarrow \Lambda \pi^+) &= -0.755 \pm 0.005 \pm 0.003, \\ \alpha_{\text{avg}}(\Lambda_c^+ \rightarrow \Sigma^0 K^+) &= -0.55 \pm 0.18 \pm 0.09, \\ \alpha_{\text{avg}}(\Lambda_c^+ \rightarrow \Sigma^0 \pi^+) &= -0.463 \pm 0.016 \pm 0.008, \end{aligned}$$

α -induced CP asymmetry:

- **First measurements** of A_{CP}^α for SCS decays of charmed baryons

Channel	A_{CP}^α
$\Lambda_c^+ \rightarrow \Lambda K^+$	$-0.023 \pm 0.086 \pm 0.071$
$\Lambda_c^+ \rightarrow \Lambda \pi^+$	$+0.020 \pm 0.007 \pm 0.013$
$\Lambda_c^+ \rightarrow \Sigma^0 K^+$	$+0.08 \pm 0.35 \pm 0.14$
$\Lambda_c^+ \rightarrow \Sigma^0 \pi^+$	$-0.023 \pm 0.034 \pm 0.030$

Λ -hyperon CP violation:

- **First measurement** of hyperon CPV searches in CF charm decays.

$$A_{CP}^\alpha = +0.013 \pm 0.007 \pm 0.011$$

No evidence of baryon CPV is found

Preliminary result

Measurements of $\Lambda_c^+ \rightarrow \Sigma^+ \pi^0$, $\Lambda_c^+ \rightarrow \Sigma^+ \eta$, $\Lambda_c^+ \rightarrow \Sigma^+ \eta'$

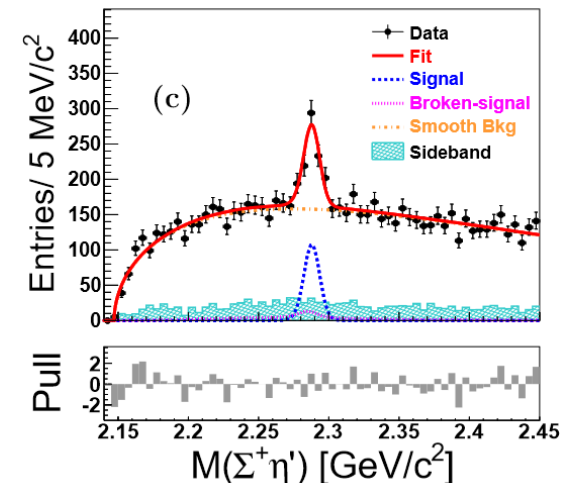
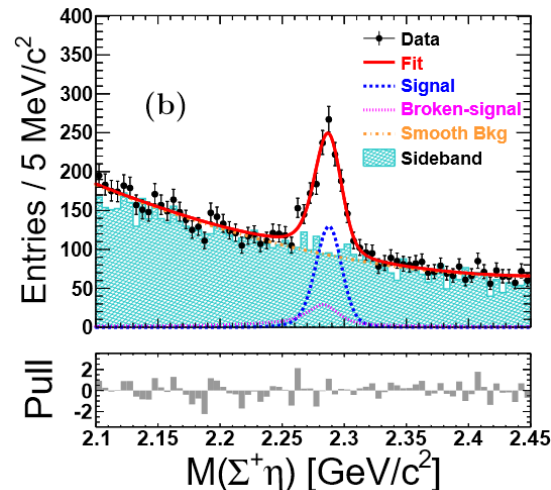
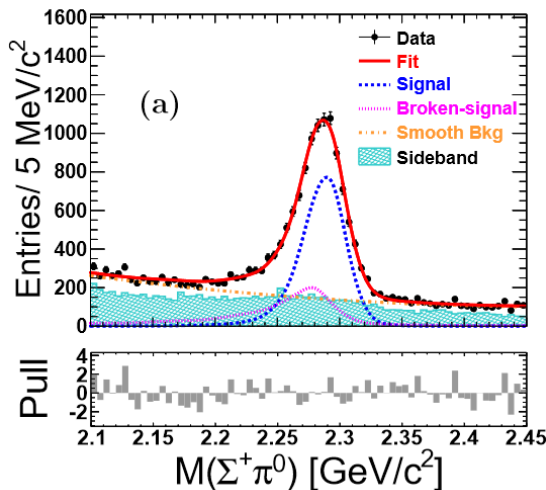
- Measure branching fractions and asymmetry parameters of charmed baryon weak decays: $B_c \rightarrow B + M$ (M is pseudoscalar or vector meson)

Decay	Körner [1]	Ivanov [2]	Żenczykowski [7]	Sharma [8]	Zou [10]	Geng [11]	Experiment [12]
$\Lambda_c^+ \rightarrow \Sigma^+ \eta$	0.16	0.11	0.90	0.57	0.74	0.32 ± 0.13	0.44 ± 0.20
$\Lambda_c^+ \rightarrow \Sigma^+ \eta'$	1.28	0.12	0.11	0.10	–	1.44 ± 0.56	1.5 ± 0.6
$\Lambda_c^+ \rightarrow \Sigma^+ \pi^0$	0.70	0.43	0.39	–0.31	–0.76	-0.35 ± 0.27	-0.55 ± 0.11
$\Lambda_c^+ \rightarrow \Sigma^+ \eta$	0.33	0.55	0.00	–0.91	–0.95	-0.40 ± 0.47	–
$\Lambda_c^+ \rightarrow \Sigma^+ \eta'$	–0.45	–0.05	–0.91	0.78	–	$1.00^{+0.00}_{-0.17}$	–

Branching fraction

Asymmetry parameter

- The rates of $\Lambda_c^+ \rightarrow \Sigma^+ \eta$ and $\Lambda_c^+ \rightarrow \Sigma^+ \eta'$ are naively comparable or the former is larger than the latter. However, the branching fraction of $\Lambda_c^+ \rightarrow \Sigma^+ \eta'$ measured by BESIII was found to be larger than the $\Lambda_c^+ \rightarrow \Sigma^+ \eta$ mode.
$$\frac{\mathcal{B}(\Lambda_c^+ \rightarrow \Sigma^+ \eta')}{\mathcal{B}(\Lambda_c^+ \rightarrow \Sigma^+ \eta)} = 3.5 \pm 2.1 \pm 0.4.$$
- Compare the asymmetry parameters of $\Lambda_c^+ \rightarrow \Sigma^+ \pi^0$ and $\Lambda_c^+ \rightarrow \Sigma^0 \pi^+$ could test the isospin symmetry.



Measurements of $\Lambda_c^+ \rightarrow \Sigma^+ \pi^0$, $\Lambda_c^+ \rightarrow \Sigma^+ \eta$, $\Lambda_c^+ \rightarrow \Sigma^+ \eta'$

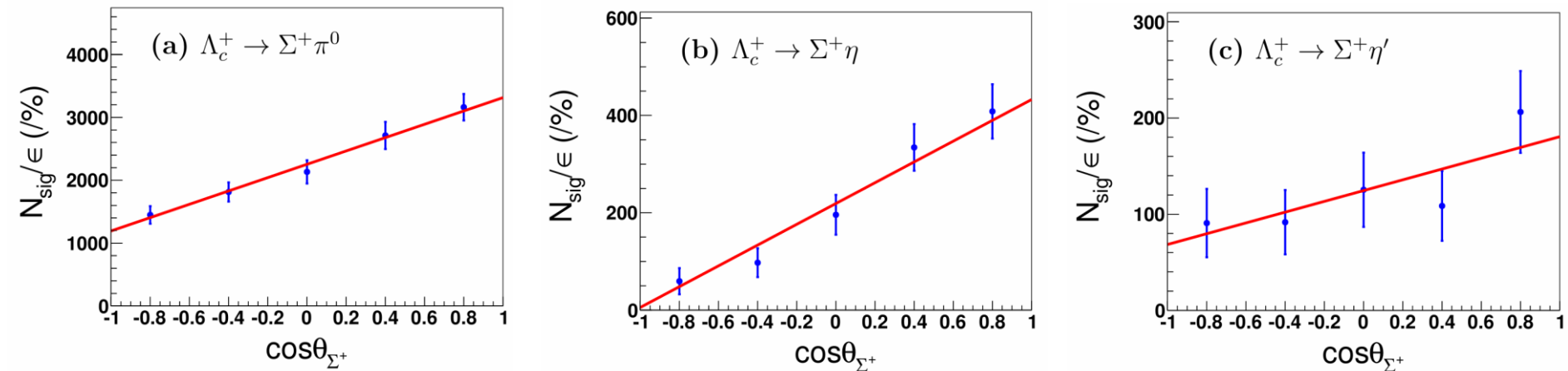
$$Br(\Lambda_c^+ \rightarrow \Sigma^+ \eta) = (3.14 \pm 0.35 \pm 0.11 \pm 0.25) \times 10^{-3}$$

$$Br(\Lambda_c^+ \rightarrow \Sigma^+ \eta') = (4.16 \pm 0.75 \pm 0.21 \pm 0.33) \times 10^{-3}$$

Most precise
results to date

$$\frac{\mathcal{B}_{\Lambda_c^+ \rightarrow \Sigma^+ \eta'}}{\mathcal{B}_{\Lambda_c^+ \rightarrow \Sigma^+ \eta}} = 1.34 \pm 0.28 \pm 0.06,$$

BESIII: $\frac{\mathcal{B}(\Lambda_c^+ \rightarrow \Sigma^+ \eta')}{\mathcal{B}(\Lambda_c^+ \rightarrow \Sigma^+ \eta)} = 3.5 \pm 2.1 \pm 0.4.$



$$\alpha_{\Sigma^+ \pi^0} = -0.48 \pm 0.02 \pm 0.02$$

Agrees with the world average value: -0.55 ± 0.11 in much improved precision.

Consistent with $\alpha_{\Sigma^0 \pi^+} = -0.463 \pm 0.016 \pm 0.008$, no isospin symmetry broken.

$$\alpha_{\Sigma^+ \eta} = -0.99 \pm 0.03 \pm 0.05 \text{ and } \alpha_{\Sigma^+ \eta'} = -0.46 \pm 0.06 \pm 0.03$$

Measured for the first time.

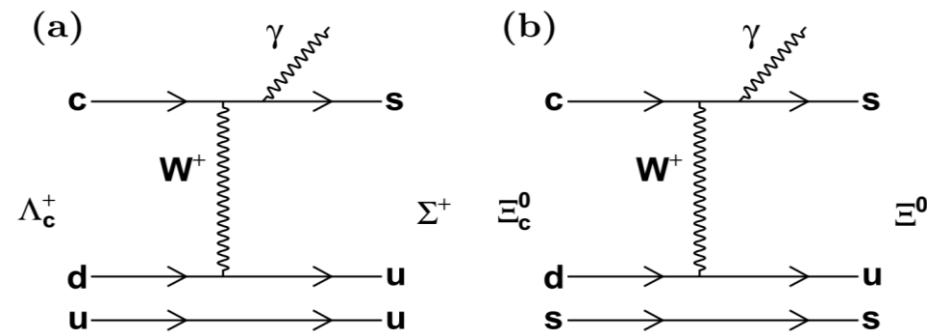
Preliminary result

First search for the weak radiative decays

$$\Lambda_c^+ \rightarrow \Sigma^+ \gamma \text{ and } \Xi_c^0 \rightarrow \Xi^0 \gamma$$

arXiv:2206.12517

- Weak radiative decays of charmed hadrons are dominated by the **long-range nonperturbative processes** that can enhance the branching fractions up to 10^{-4} , whereas short-range interactions are predicted to yield rates at the level of 10^{-8} [1,2]
- At the **CF level**, there are two decay modes for the weak radiative decays of anti-triplet charmed baryons induced from $cd \rightarrow us + \gamma$, i.e., $\Lambda_c^+ \rightarrow \Sigma^+ \gamma$ and $\Xi_c^0 \rightarrow \Xi^0 \gamma$ decays.
- The theoretical estimates of branching fractions cover ranges of $(4.5 - 29.1) \times 10^{-5}$ and $(3.0 - 19.5) \times 10^{-5}$ for $\Lambda_c^+ \rightarrow \Sigma^+ \gamma$ and $\Xi_c^0 \rightarrow \Xi^0 \gamma$ decays, respectively [3-6].



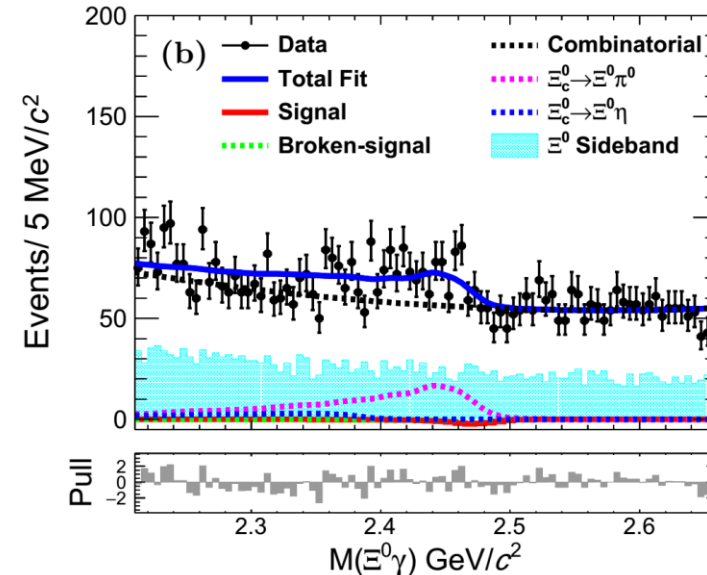
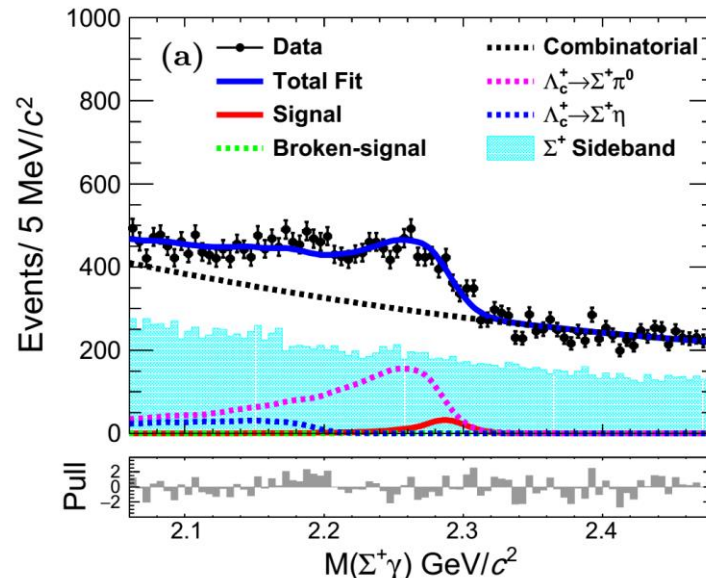
- [1] PRD 52, 6383 (1995).
- [2] PLB 382, 415 (1996).
- [3] PRD 28, 2176 (1983).
- [4] PRD 47, 2858 (1993).
- [5] PRD 51, 1199 (1995).
- [6] arXiv:2109.01216.

First search for the weak radiative decays

$$\Lambda_c^+ \rightarrow \Sigma^+ \gamma \text{ and } \Xi_c^0 \rightarrow \Xi^0 \gamma$$

arXiv:2206.12517

- There are no evident $\Lambda_c^+ \rightarrow \Sigma^+ \gamma$ or $\Xi_c^0 \rightarrow \Xi^0 \gamma$ signals. The signal significance of $\Lambda_c^+ \rightarrow \Sigma^+ \gamma$ decay is only 2.2σ , after considering the systematic uncertainty.



- Taking $\Lambda_c^+ \rightarrow pK^- \pi^+$ and $\Xi_c^0 \rightarrow \Xi^- \pi^+$ as normalization channels, the upper limits at 90% confidence level (C.L.) on the ratios of branching fractions are

$$\frac{\text{Br}(\Lambda_c^+ \rightarrow \Sigma^+ \gamma)}{\text{Br}(\Lambda_c^+ \rightarrow pK^- \pi^+)} < 3.99 \times 10^{-3}$$

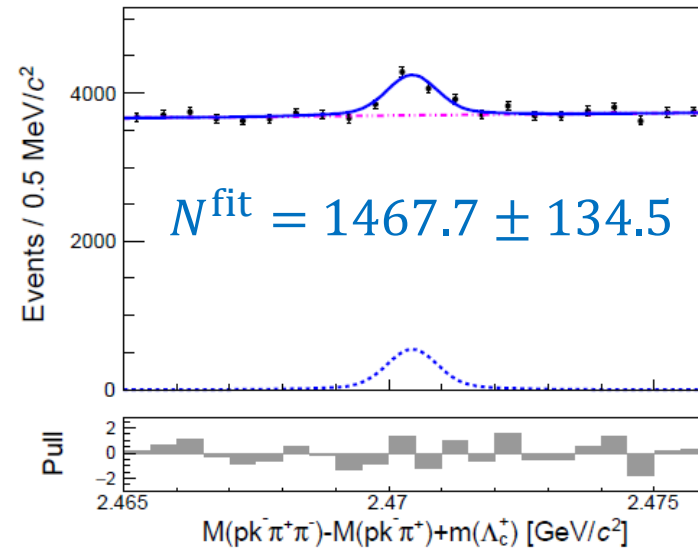
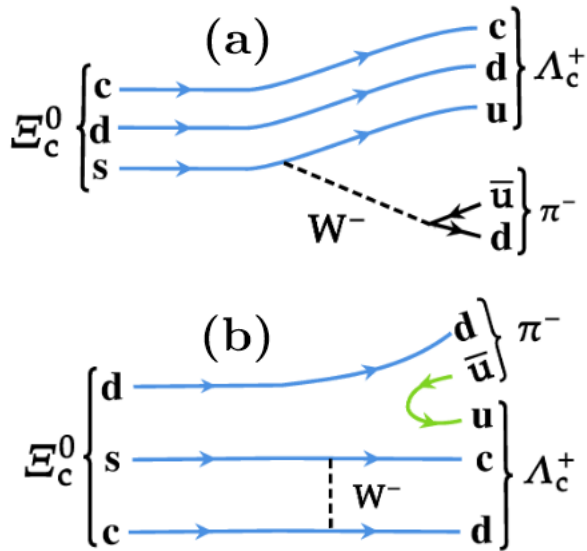
$$\frac{\text{Br}(\Xi_c^0 \rightarrow \Xi^0 \gamma)}{\text{Br}(\Xi_c^0 \rightarrow \Xi^- \pi^+)} < 1.15 \times 10^{-2}$$

- The upper limits at 90% C.L. on the absolute branching fractions are determine:

$$\text{Br}(\Lambda_c^+ \rightarrow \Sigma^+ \gamma) < 2.55 \times 10^{-4}$$

$$\text{Br}(\Xi_c^0 \rightarrow \Xi^0 \gamma) < 1.73 \times 10^{-4}$$

- The heavy-flavor-conserving nonleptonic decays in charmed baryons containing both an s and a c quark, that proceeds via the decay of the s quark, i.e. $\Xi_c^0 \rightarrow \Lambda_c^+ \pi^-$.
- The decay width of $\Xi_c^0 \rightarrow \Lambda_c^+ \pi^-$ is based on the size of the s quark decay amplitude of $s \rightarrow u(\bar{u}d)$ and the weak scattering amplitude $cs \rightarrow dc$.



$$Br(\Xi_c^0 \rightarrow \Lambda_c^+ \pi^-) = (6.8 \pm 0.6 \pm 0.6 \pm 2.0) \times 10^{-3}$$

- The result is consistent with previous measurement by LHCb experiment and recent theoretical calculation, and larger than the previous theoretical predictions.

Measurements of $\Omega_c^0 \rightarrow \Omega^- l^+ \nu$ and $\Xi_c^0 \rightarrow \Xi^- l^+ \nu$

Semi-leptonic decays of charmed baryons:

- Ideal test of QCD in transition region of (non-)perturbative.
- The cleanest processes among charm decays
- Test lepton flavor universality (LFU).

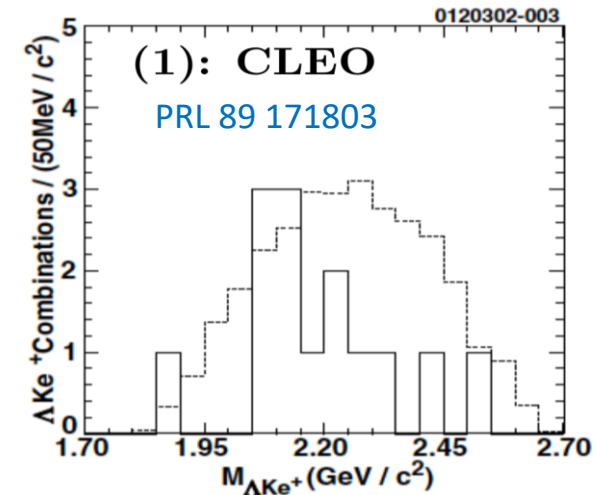
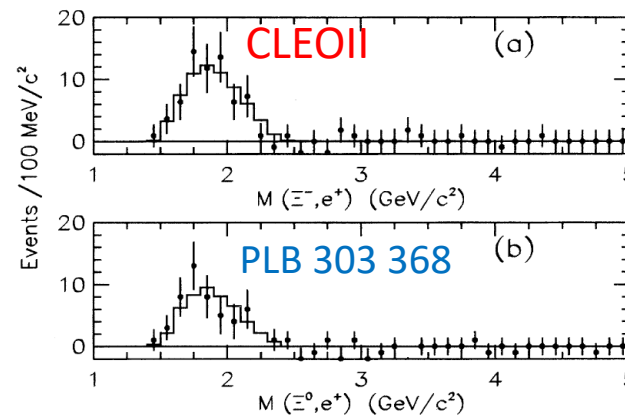
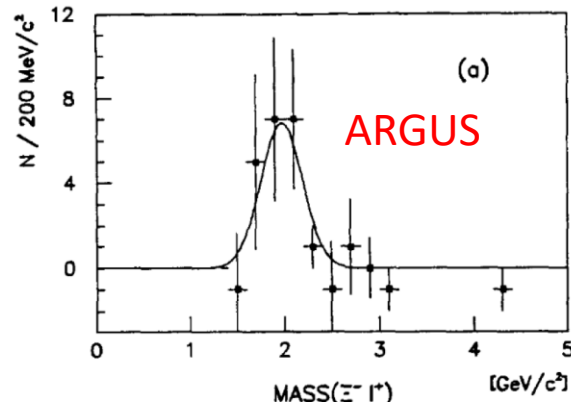
Experimentally:

- BESIII measured the $\mathcal{B}(\Lambda_c^+ \rightarrow \Lambda l^+ \nu)$
- ARGUS and CLEOII measured $\mathcal{B}(\Xi_c \rightarrow \Xi l^+ \nu)$
- CLEO measured $\mathcal{B}(\Omega_c^0 \rightarrow \Omega^- e^+ \nu)$

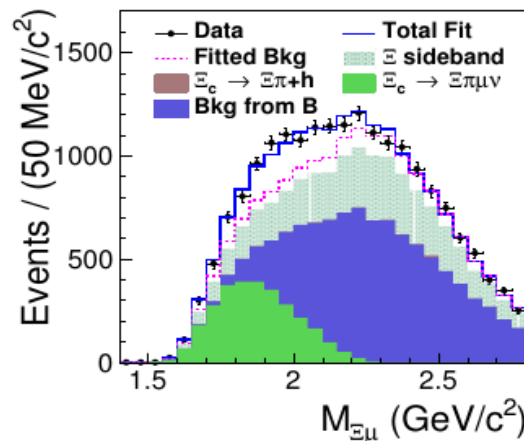
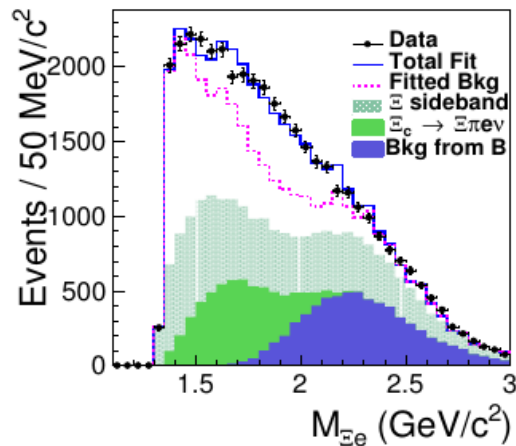
$$\mathcal{B}(\Lambda_c^+ \rightarrow \Lambda e^+ \nu_e) = (3.6 \pm 0.4)\% \text{ PRL 115, 221805(2015)}$$

$$\mathcal{B}(\Lambda_c^+ \rightarrow \Lambda \mu^+ \nu_e) = (3.5 \pm 0.4)\% \text{ PLB 767, 42 (2017)}$$

} large uncertainty



Measurements of $\Omega_c^0 \rightarrow \Omega^- l^+ \nu$ and $\Xi_c^0 \rightarrow \Xi^- l^+ \nu$



Data-driven method for bkg

- Mis-selected l^+
- Wrongly constructed Ξ
- $\Xi_c \rightarrow \Xi \pi l^+ \nu_\ell$
- $\Xi_c \rightarrow \Xi \pi + h$
- B decay

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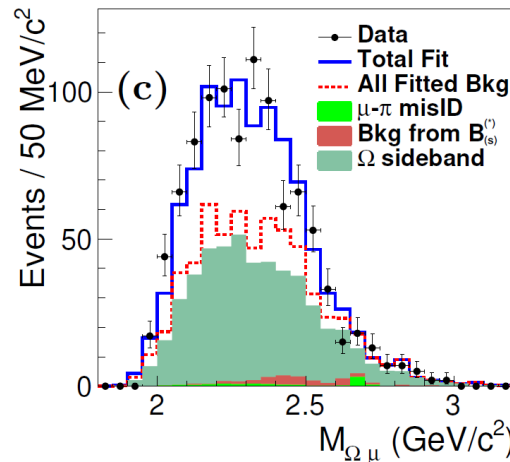
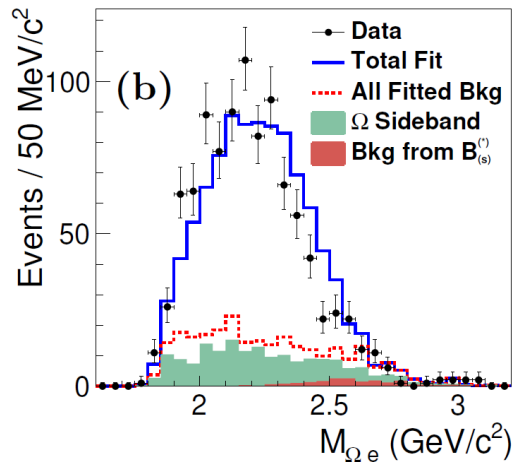
$$Br(\Xi_c^0 \rightarrow \Xi^- e^+ \nu_e) = (1.31 \pm 0.39)\%$$

Previous: $(2.34 \pm 1.59)\%$

$$Br(\Xi_c^0 \rightarrow \Xi^- \mu^+ \nu_\mu) = (1.27 \pm 0.39)\%$$

Consistent with LFU

PRD 105, L091101 (2022)



$$\frac{\mathcal{B}(\Omega_c^0 \rightarrow \Omega^- e^+ \nu)}{\mathcal{B}(\Omega_c^0 \rightarrow \Omega^- \pi^+)} = 1.98 \pm 0.15$$

$$\frac{\mathcal{B}(\Omega_c^0 \rightarrow \Omega^- \mu^+ \nu)}{\mathcal{B}(\Omega_c^0 \rightarrow \Omega^- \pi^+)} = 1.94 \pm 0.21$$

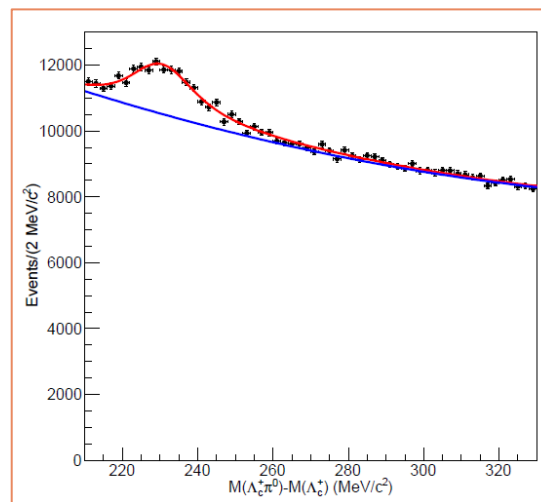
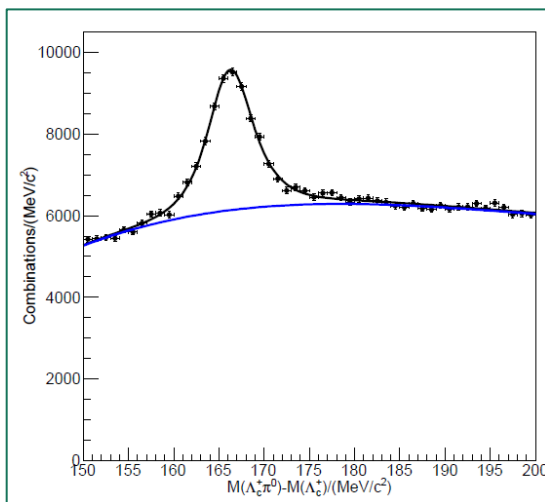
Previous: 2.4 ± 1.2

Consistent with LFU

Measurement of the masses and widths of the $\Sigma_c(2455)^+$ and $\Sigma_c(2520)^+$ baryons

PRD 104, 052003 (2021)

- The doubly-charged and neutral Σ_c states are well studied.
- However, there is comparatively little experimental information on the singly-charged Σ_c^+ states due to π^0 transition.
- Good tests of models of isospin mass splittings, most mass models predict Σ_c^+ states have masses a little lower than doubly-charged and neutral analogs.
- Reconstruct Σ_c^+ baryons from $\Lambda_c^+\pi^0$ decays. **Fits performed to limited-range subsets**



$\Sigma_c(2455)^+$:

- $\Delta(M) = 166.17 \pm 0.05^{+0.16}_{-0.07}$ MeV
- $\Gamma = 2.3 \pm 0.3 \pm 0.3$ MeV

$\Sigma_c(2520)^+$:

- $\Delta(M) = 230.9 \pm 0.5^{+0.5}_{-0.1}$ MeV
- $\Gamma = 17.2^{+2.3+3.1}_{-2.1-0.7}$ MeV

First non-zero measurements of the intrinsic widths.

No deviation from the expectations of their isospin partners

- Charmed baryons are good platform to study related theories such as HQET.
- High mass excited states are still not well understood.
- Most of Λ_c^+ excited states decay (via Σ_c) to $\Lambda_c^+ \pi \pi$ final states

Experimentally

$B^0 \rightarrow \Lambda_c \pi \pi \bar{p}$ is a promising channel to study

- Low background contributions
- High branching fraction $\sim 10^{-4}$
- Large $B\bar{B}$ sample: 772×10^6

B^0 Decay Modes

Γ_{500}	$\bar{\Lambda}_c^- p \pi^+ \pi^-$ (nonresonant)	$(5.5 \pm 1.0) \times 10^{-4}$
Γ_{501}	$\bar{\Sigma}_c(2520)^{--} p \pi^+$	$(1.02 \pm 0.18) \times 10^{-4}$
Γ_{502}	$\bar{\Sigma}_c(2520)^0 p \pi^-$	$< 3.1 \times 10^{-5}$
Γ_{503}	$\bar{\Sigma}_c(2455)^0 p \pi^-$	$(1.08 \pm 0.16) \times 10^{-4}$

$\Lambda_c(2880)^+$ Decay Modes

$$\Gamma_1 \quad \Lambda_c^+ \pi^+ \pi^-$$

$\Lambda_c(2595)^+$ Decay Modes

$$\Gamma_1 \quad \Lambda_c^+ \pi^+ \pi^-$$

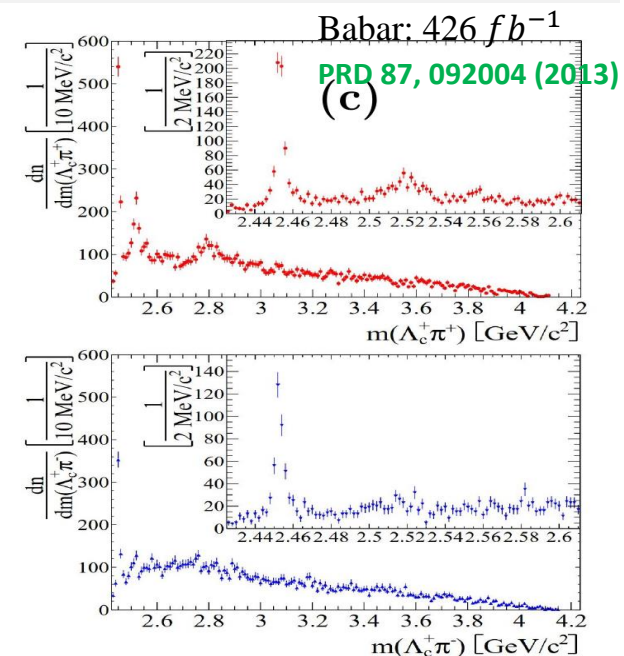
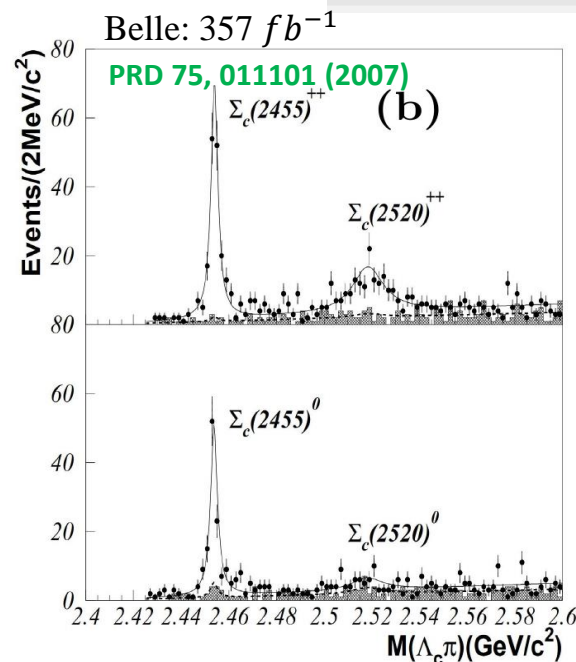
$\Lambda_c(2625)^+$ Decay Modes

$$\Gamma_1 \quad \Lambda_c^+ \pi^+ \pi^-$$

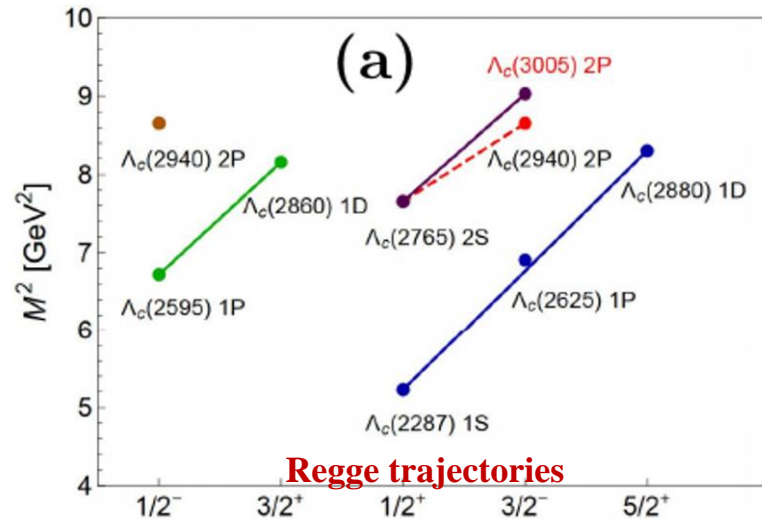
$\Lambda_c(2940)^+$ Decay Modes

$$\Gamma_1 \quad p D^0$$

$$\Gamma_2 \quad \Sigma_c(2455)^{0,++} \pi^\pm$$



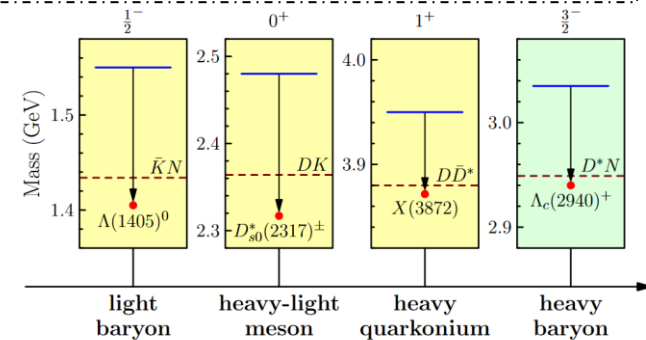
- $\Lambda_c(2940) \rightarrow \Sigma_c(2455)\pi$ is the current highest mass Λ_c excited state
- LHCb result prefers $J^P = \frac{3}{2}^-$, but doesn't exclude other values [JHEP 05, 030 (2017)].
- The mass of $\Lambda_c(2940)$ is lower than predication of contradictory quark model.



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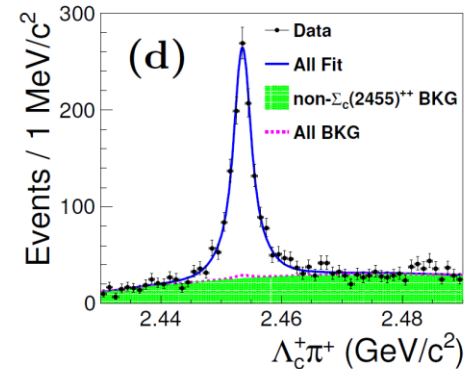
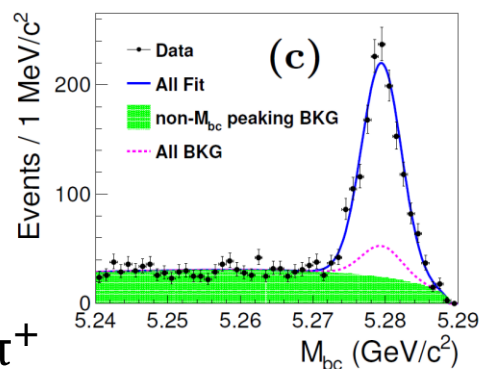
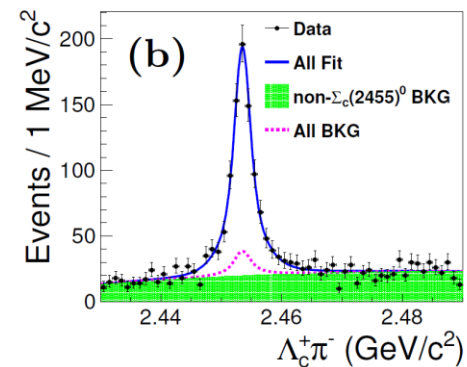
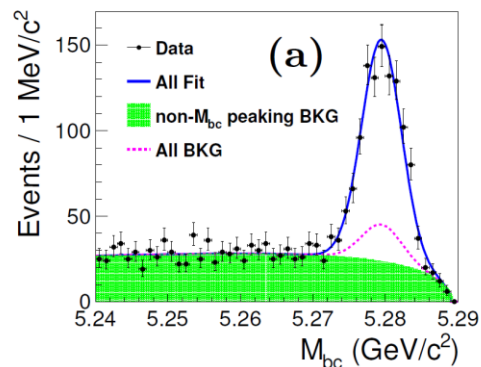
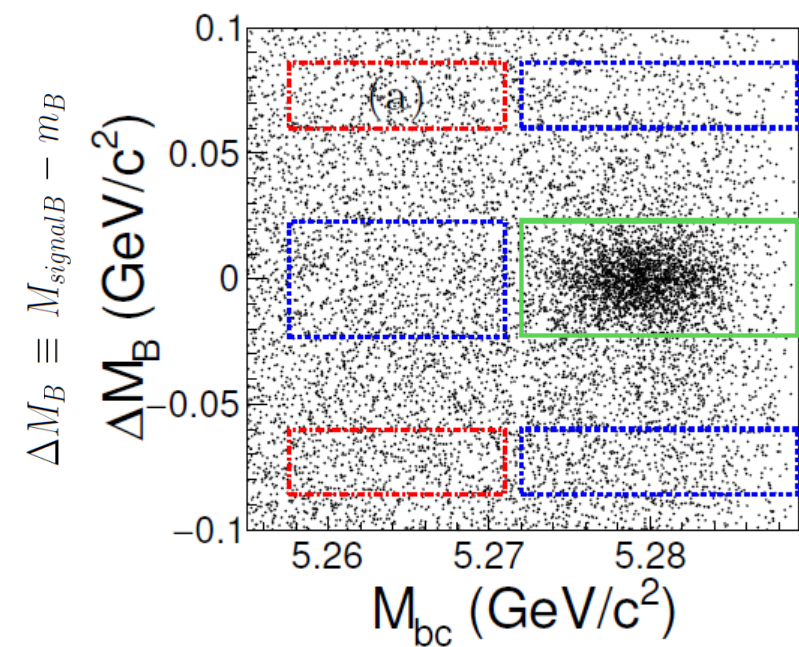
State	J^P	Present	[9]	[10]	[11]	[16]	[20]	[21]	PDG [1]
$2P$	$\frac{1}{2}^-$	2.978	3.017	2.983	2.989	2.890		2.980	
$2P$	$\frac{3}{2}^-$	2.970	3.034	3.005	3.000	2.917		3.004	$2.9396^{+0.0014}_{-0.0015}$

[9]. Phys. Lett. B 659, 612 (2008). [16]. Phys. Rev. D 92, 114029 (2015).
 [10]. Phys. Rev. D 84, 014025 (2011). [21]. Eur. Phys. J. C 77, 154 (2017).
 [11]. Phys. Rev. D 91, 054034 (2015).



Theoretically:

- $\Lambda_c(2940)$ has $J^P = \frac{1}{2}^-$, belongs to another Regge trajectories
- Search for $\Lambda_c(3005)$
- $\Lambda_c(2940)$ has D^*N contribution is similar to (1405) , $D_s(2317)$, $X(3872)$,
- $\Lambda_c\left(\frac{1}{2}^-, 2P\right)$ has reversed mass which is larger than $\Lambda_c\left(\frac{3}{2}^-, 2P\right)$



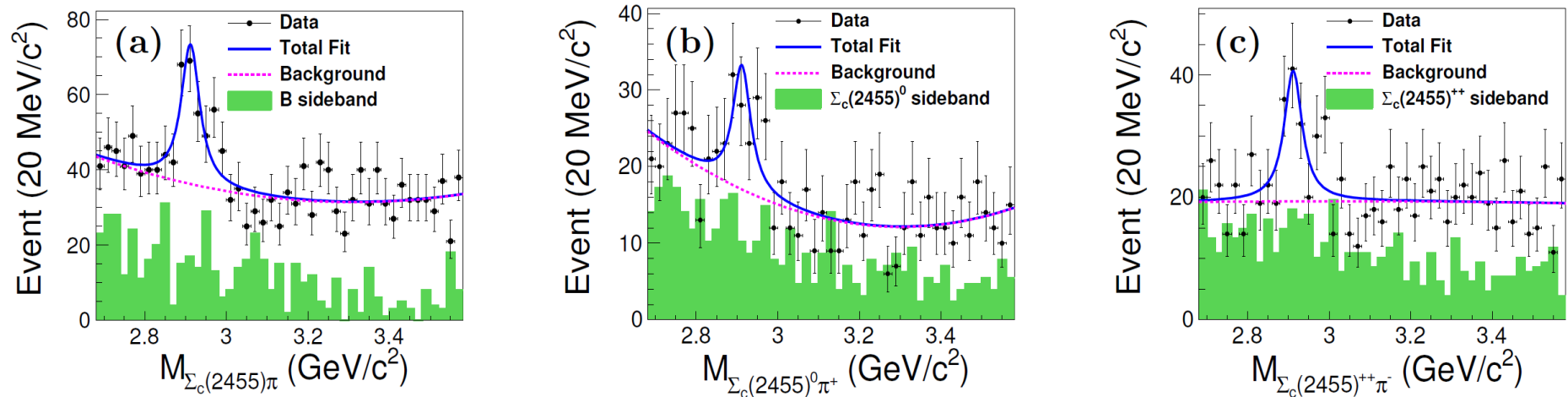
Reconstruct Λ_c^+ via $pK^-\pi^+$, pK_S , and $\Lambda\pi^+$

Significant signal of $\bar{B}^0 \rightarrow \Sigma_c(2455)\pi\bar{p}$

Measured branching fraction consistent with previous study at improved precision

B decay mode	yield	Br(10^{-4})	PDG (10^{-4})
$\Sigma_c(2455)^0 \pi \bar{p}$	767 ± 44	$1.09 \pm 0.06 \pm 0.07$	1.08 ± 0.16
$\Sigma_c(2455)^{++} \pi \bar{p}$	1213 ± 73	$1.84 \pm 0.11 \pm 0.12$	1.88 ± 0.24

Search for excited Λ_c states



- A resonant structure in $\Sigma_c(2455)\pi$ invariant mass spectra, the mass and width are measured to be:

$$m = (2913.8 \pm 5.6 + 3.7) \text{ MeV}/c^2, \Gamma = (51.8 \pm 20.0 \pm 18.8) \text{ MeV}$$

- The significance in the most conservative case is 4.2σ
- The J^p prefers to be $\frac{1}{2}^-$, $\Lambda_c\left(\frac{1}{2}^-, 2P\right)$, thus name it $\Lambda_c(2910)$
- Need further efforts from both experimental and theoretical sides to understand $\Lambda_c(2910)$ and $\Lambda_c(2940)$

- Although Belle has stopped data taking for >10 years ago, we are still producing exciting results.
- The Belle II experiment at SuperKEKB aims to find New Physics beyond the SM with ultimate precision measurement (a few %, typically) of heavy flavor decays.
- Belle II started data taking on 25 March 2019 with its full detector and is performing as expected and obtained early physics results.
- Belle II can achieve lots of important results on baryons and charmed baryons near future.

Thanks!

Measurements of $\Omega_c^0 \rightarrow \Omega^- l^+ \nu$ and $\Xi_c^0 \rightarrow \Xi^- l^+ \nu$

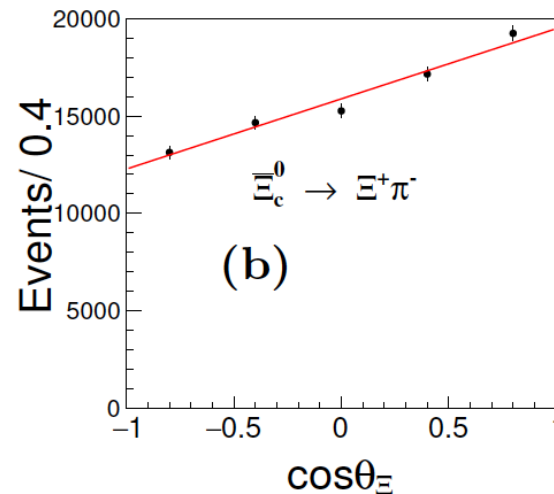
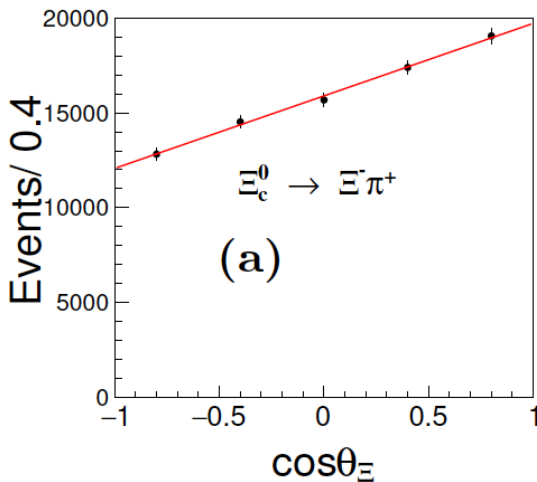
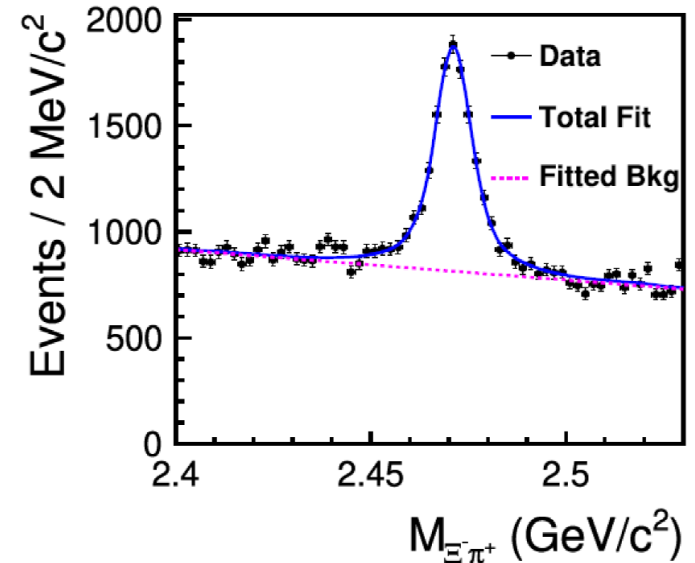
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\mathcal{A}_{CP} of $\Xi_c^0 \rightarrow \Xi^- \pi^+$: search for CP violation in charmed baryons section

$$\frac{dN}{d \cos \theta_{\Xi}} \propto 1 + \alpha_{\Xi^- \pi^+} + \alpha_{\Xi^-} \cos \theta_{\Xi}$$

θ_{Ξ} : angle between the \vec{p}_{Λ} and $-\vec{p}_{\Xi_c^0}$
in the Ξ^- rest frame

$$\mathcal{A}_{CP} = (\alpha_{\Xi^- \pi^+} + \alpha_{\Xi^+ \pi^-}) / (\alpha_{\Xi^- \pi^+} - \alpha_{\Xi^+ \pi^-})$$



$$\alpha_{\Xi^- \pi^+} = -0.60 \pm 0.045$$

$$\alpha_{\Xi^+ \pi^-} = 0.58 \pm 0.045$$

$$\mathcal{A}_{CP} = 0.015 \pm 0.056$$

Previous: -0.60 ± 0.4