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Improved radiative corrections for $\tau \rightarrow \pi\ (K)\ \nu_\tau\ [\gamma]$ and reliable new physics tests

See next talk by Alberto Lusiani

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3) $\tau \to P\, \nu_\tau\, [\gamma]$ ($P=\pi,K$)

4) Calculation of $R_{\tau/P} \equiv \dfrac{\Gamma(\tau \to P\nu_\tau[\gamma])}{\Gamma(P \to \mu\nu_\mu[\gamma])}$

5) Results

6) Applications

7) Conclusions

1. Motivation

- ✓ Lepton Universality (LU) as a basic tenet of the Standard Model (SM).
  - ✓ A few anomalies observed in semileptonic B meson decays*.
  - ✓ Lower energy observables currently provide the most precise test of LU**.
- ✓ We aim to test muon-tau lepton universality through the ratio (P = π, K)***:

$$R_{\tau/P} \equiv \frac{\Gamma(\tau \to P\nu_\tau[\gamma])}{\Gamma(P \to \mu\nu_\mu[\gamma])} = \left|\frac{g_\tau}{g_\mu}\right|_P^2 R^{(0)}_{\tau/P}\left(1+\delta R_{\tau/P}\right)$$

  - ✓ $g_\tau = g_\mu$ according to LU.
  - ✓ $R_{\tau/P}^{(0)}$ is the LO result $R^{(0)}_{\tau/P} = \frac{1}{2}\frac{M_\tau^3}{m_\mu^2 m_P}\frac{(1-m_P^2/M_\tau^2)^2}{(1-m_\mu^2/m_P^2)^2}$.
  - ✓ $\delta R_{\tau/P}$ encodes the radiative corrections.
- ✓ $\delta R_{\tau/P}$ was calculated by Decker & Finkemeier (DF'95) ^ :
  - ✓ $\delta R_{\tau/\pi}$ = (0.16 ± 0.14)% and $\delta R_{\tau/K}$ = (0.90 ± 0.22)%.
- ✓ Important phenomenological and theoretical reasons to address the analysis again.

* Albrecht et al.'21
** Bryman et al.'21

*** Marciano & Sirlin'93
^ Decker & Finkemeier'95

1. Motivation

- ✓ Phenomenological disagreement in LU tests:

See next talk by Alberto Lusiani, HFLAV22

  - ✓ Using $\frac{\Gamma(\tau \to P\nu_\tau[\gamma])}{\Gamma(P \to \mu\nu_\mu[\gamma])}$ and DF'95*, HFLAV** reported:
    - ✓ $|g_\tau/g_\mu|_\pi = 0.9958 \pm 0.0026$ (at 1.6σ of LU)
    - ✓ $|g_\tau/g_\mu|_K = 0.9879 \pm 0.0063$ (at 1.9σ of LU)
  - ✓ Using $\frac{\Gamma(\tau \to e\bar{\nu}_e\nu_\tau[\gamma])}{\Gamma(\mu \to e\bar{\nu}_e\nu_\mu[\gamma])}$, HFLAV** reported:
    - ✓ $|g_\tau/g_\mu| = 1.0010 \pm 0.0014$ (at 0.7σ of LU)
  - ✓ Using $\frac{\Gamma(W \to \tau\nu_\tau)}{\Gamma(W \to \mu\nu_\mu)}$, CMS and ATLAS*** and reported:
    - ✓ $|g_\tau/g_\mu| = 0.995 \pm 0.006$ (at 0.8σ of LU)

* Decker & Finkemeier'95
** HFLAV'21
*** CMS'21, ATLAS'21

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✓ Theoretical issues within DF'95*:

- ✓ Hadronic form factors are different for real- and virtual-photon corrections, do not satisfy the correct QCD short-distance behavior, violate unitarity, analicity and the chiral limit at leading non-trivial orders.
- ✓ A cutoff to regulate the loop integrals (separating long- and short-distance corrections)
- ✓ Unrealistic uncertainties (purely $O(e^2p^2)$ ChPT size).

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✓ Unrealistic uncertainties (purely $O(e^2p^2)$ ChPT size).

✓ By-products of the project:

✓ Radiative corrections in $\Gamma(\tau \to P\nu_\tau[\gamma])$.

✓ CKM unitarity test via $\Gamma(\tau \to K\nu_\tau[\gamma])$ or via the ratio $\Gamma(\tau \to K\nu_\tau[\gamma]) / \Gamma(\tau \to \pi\nu_\tau[\gamma])$.

✓ Constraints on possible non-standard interactions in $\Gamma(\tau \to P\nu_\tau[\gamma])$^.

* Decker & Finkemeier'95
** HFLAV'21
*** CMS'21, ATLAS'21

^ Cirigliano et al.'10 '19
^ González-Alonso & Martín-Camalich '16
^ Gonzàlez-Solís et al. '20

2. $P \to \mu\, \nu_\mu\, [\gamma]$ ($P=\pi,K$)

✓ Calculated unambigously within the Standard Model (Chiral Perturbation Theory, ChPT*).

✓ Notation by Marciano & Sirlin** and numbers by Cirigliano & Rosell*** (D=d,s for π,K and $F_\pi \approx 92.2$ MeV):

LO result

short-distance EW correction ≈ 1.0232**

structure independent (SI) contributions (point-like approximation)^

$$\Gamma(P\to\mu\nu_\mu[\gamma]) = \frac{G_F^2|V_{uD}|^2F_P^2}{4\pi}\, m_P\, m_\mu^2 \left(1-\frac{m_\mu^2}{m_P^2}\right)^2 S_{\rm EW}\left\{1+\frac{\alpha}{\pi}F(m_\mu^2/m_P^2)\right\}\times$$

$$\left\{1-\frac{\alpha}{\pi}\left[\frac{3}{2}\log\frac{m_\rho}{m_P}+c_1^{(P)}+\frac{m_\mu^2}{m_\rho^2}\left(c_2^{(P)}\log\frac{m_\rho^2}{m_\mu^2}+c_3^{(P)}+c_4^{(P)}(m_\mu/m_P)\right)-\frac{m_P^2}{m_\rho^2}\tilde{c}_2^{(P)}\log\frac{m_\rho^2}{m_\mu^2}\right]\right\}$$

structure-dependent (SD) contributions
[coefficents reported in Cirigliano & IR'07]

✓ The only model-dependence is the determination of the counterterms in $c_1^{(P)}$ and $c_3^{(P)}$:

✓ Large-N_C expansion of QCD: ChPT is enlarged by including the lightest multiplets of spin-one resonances such that the relevant Green functions are well-behaved at high energies†.

* Weinberg'79
* Gasser & Leutwyler'84 '85
** Marciano & Sirlin'93

*** Cirigliano & IR'07
^ Kinoshita'59

† Ecker et al.'89
† Cirigliano et al.'06

3. $\tau \to P\, \nu_\tau\, [\gamma]$ ($P=\pi,K$)

- ✓ Calculated within an effective approach encoding the hadronization:
  - ✓ Large-N_C expansion of QCD: ChPT is enlarged by including the lightest multiplets of spin-one resonances such that the relevant Green functions are well-behaved at high energies*.
- ✓ We follow a similar notation to $P\to\mu\nu_\mu[\gamma]$ ($D=d,s$ for π,K and $F_\pi \approx 92.2$ MeV):

LO result

short-distance EW correction ≈ 1.0232**

structure independent (SI) contributions (point-like approximation)***

$$\Gamma(\tau \to P\nu_\tau[\gamma]) = \frac{G_F^2 |V_{uD}|^2 F_P^2}{8\pi} M_\tau^3 \left(1-\frac{m_P^2}{M_\tau^2}\right)^2 S_{\mathrm{EW}} \left\{1+\frac{\alpha}{\pi}G(m_P^2/M_\tau^2)\right\}\times$$
$$\left\{1-\frac{3\alpha}{2\pi}\log\frac{m_\rho}{M_\tau}+\delta_{\tau P}\big|_{\mathrm{rSD}}+\delta_{\tau P}\big|_{\mathrm{vSD}}\right\}$$

real-photon structure-dependent (rSD) contributions

virtual-photon structure-dependent (vSD) contributions

  - ✓ Real-photon structure-dependent (rSD) contributions from Guo & Roig'10^.
  - ✓ Virtual-photon structure-dependent (vSD) contributions not calculated in the literature.

* Ecker et al.'89
* Cirigliano et al.'06
** Erler'02

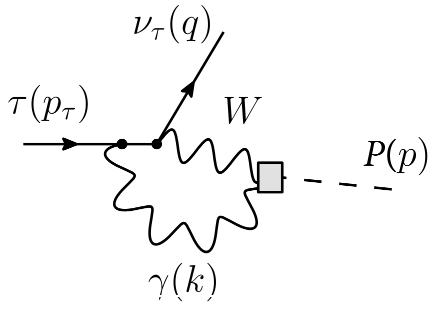
*** Kinoshita'59
^ Guo & Roig'10

3. τ → P ν_τ [γ]   (P=π,K)

✓ Virtual-photon structure-dependent contribution (vSD):

$$i\mathcal{M}[\tau\to P\nu_\tau]|_{\rm SD} = G_F V_{uD} e^2 \int \frac{{\rm d}^d k}{(2\pi)^d} \frac{\ell^{\mu\nu}}{k^2[(p_\tau+k)^2-M_\tau^2]} \left[i\epsilon_{\mu\nu\lambda\rho}k^\lambda p^\rho F_V^P(W^2,k^2) + F_A^P(W^2,k^2)\lambda_{1\mu\nu} + 2B(k^2)\lambda_{2\mu\nu}\right]$$

$$\begin{aligned} \ell^{\mu\nu} &= \bar{u}(q)\gamma^\mu(1-\gamma_5)[(\not p_\tau + \not k) + M_\tau]\gamma^\nu u(p_\tau) \\ \lambda_{1\mu\nu} &= \left[(p+k)^2 + k^2 - m_P^2\right] g_{\mu\nu} - 2k_\mu p_\nu \\ \lambda_{2\mu\nu} &= k^2 g_{\mu\nu} - \frac{k^2(p+k)_\mu p_\nu}{(p+k)^2 - m_P^2} \end{aligned}$$



✓ Form factors from Guo & Roig'10 and Guevara et al.'13,'21*:

$$\begin{aligned} F_V^P(W^2,k^2) &= \frac{-N_C M_V^4}{24\pi^2 F_P (k^2-M_V^2)(W^2-M_V^2)} \\ F_A^P(W^2,k^2) &= \frac{F_P}{2}\frac{M_A^2 - 2M_V^2 - k^2}{(M_V^2-k^2)(M_A^2-W^2)} \\ B(k^2) &= \frac{F_P}{M_V^2-k^2} \end{aligned}$$

✓ Well-behaved two- and three-point Green functions.

✓ Chiral and U(3) limits.

✓ M_V and M_A vector- and axial-vector resonance mass: $M_V=M_\rho$ and $M_A=M_{a1}$ (π case); $M_V=M_{K^*}$ and $M_A \approx M_{f1}$ (K case).

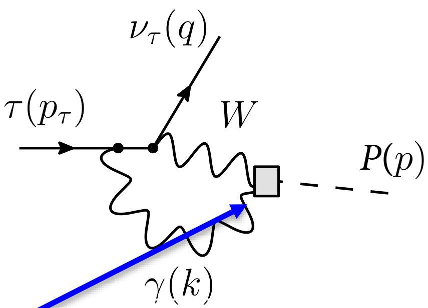
* Guo & Roig'10
* Guevara et al.'13,'21

3. $\tau \to P\,\nu_\tau\,[\gamma]$ $(P=\pi,K)$

✓ Virtual-photon structure-dependent contribution (vSD):

$$i\mathcal{M}[\tau\to P\nu_\tau]|_{\rm SD}=G_F V_{uD} e^2\int\frac{{\rm d}^d k}{(2\pi)^d}\frac{\ell^{\mu\nu}}{k^2[(p_\tau+k)^2-M_\tau^2]}\left[i\epsilon_{\mu\nu\lambda\rho}k^\lambda p^\rho F_V^P(W^2,k^2)+F_A^P(W^2,k^2)\lambda_{1\mu\nu}+2B(k^2)\lambda_{2\mu\nu}\right]$$

$$\begin{aligned}\ell^{\mu\nu} &= \bar{u}(q)\gamma^\mu(1-\gamma_5)[(\not{p}_\tau+\not{k})+M_\tau]\gamma^\nu u(p_\tau)\\ \lambda_{1\mu\nu} &= \left[(p+k)^2+k^2-m_P^2\right]g_{\mu\nu}-2k_\mu p_\nu\\ \lambda_{2\mu\nu} &= k^2 g_{\mu\nu}-\frac{k^2(p+k)_\mu p_\nu}{(p+k)^2-m_P^2}\end{aligned}$$



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* Guo & Roig'10
* Guevara et al.'13,'21

4. Calculation of $R_{\tau/P} = R_{\tau/P}^{(0)}(1 + \delta R_{\tau/P}) = R_{\tau/P}^{(0)}(1 + \delta_{\tau P} - \delta_{P\mu})$

1. Structure-independent contribution (point-like approximation): SI.

✓ We confirm the results by DF'95*.

$$\delta R_{\tau/P}\big|_{\rm SI} = \frac{\alpha}{2\pi}\left\{\frac{3}{2}\log\frac{M_\tau^2 m_P^2}{m_\mu^4} + \frac{3}{2} + g\left(\frac{m_P^2}{M_\tau^2}\right) - f\left(\frac{m_\mu^2}{m_P^2}\right)\right\}$$

$$f(x) = 2\left(\frac{1+x}{1-x}\log x - 2\right)\log(1-x) - \frac{x(8-5x)}{2(1-x)^2}\log x + 4\frac{1+x}{1-x}{\rm Li}_2(x) - \frac{x}{1-x}\left(\frac{3}{2} + \frac{4}{3}\pi^2\right)$$

$$g(x) = 2\left(\frac{1+x}{1-x}\log x - 2\right)\log(1-x) - \frac{x(2-5x)}{2(1-x)^2}\log x + 4\frac{1+x}{1-x}{\rm Li}_2(x) + \frac{x}{1-x}\left(\frac{3}{2} - \frac{4}{3}\pi^2\right)$$

$\delta R_{\tau/\pi}|_{SI} = 1.05\%$ and $\delta R_{\tau/K}|_{SI} = 1.67\%$

2. Real-photon structure-dependent contribution: rSD.

✓ $\delta_{P\mu}|_{rSD}$ from Cirigliano & IR'07**: $\delta_{\pi\mu}|_{rSD} = -1.3\cdot 10^{-8}$ and $\delta_{K\mu}|_{rSD} = -1.7\cdot 10^{-5}$.

✓ $\delta_{\tau P}|_{rSD}$ from Guo & Roig'10***: $\delta_{\tau\pi}|_{rSD} = 0.15\%$ and $\delta_{\tau K}|_{rSD} = (0.18 \pm 0.05)\%$.

$\delta R_{\tau/\pi}|_{rSD} = 0.15\%$ and $\delta R_{\tau/K}|_{rSD} = (0.18 \pm 0.15)\%$

* Decker & Finkemeier'95
** Cirigliano & Rosell'07
*** Guo & Roig'10

4. Calculation of $R_{\tau/P} = R_{\tau/P}^{(0)} (1 + \delta R_{\tau/P}) = R_{\tau/P}^{(0)} (1 + \delta_{\tau P} - \delta_{P\mu})$

3. Virtual-photon structure-dependent contribution: vSD.

- ✓ $\delta_{P\mu}|_{vSD}$ from Cirigliano & IR'07*: $\delta_{\pi\mu}|_{vSD} = (0.54 \pm 0.12)\%$ and $\delta_{K\mu}|_{vSD} = (0.43 \pm 0.12)\%$.
- ✓ $\delta_{\tau P}|_{vSD}$, **new calculation**: $\delta_{\tau\pi}|_{vSD} = (-0.48 \pm 0.56)\%$ and $\delta_{\tau K}|_{vSD} = (-0.45 \pm 0.57)\%$.

$$\delta R_{\tau/\pi}|_{vSD} = (-1.02 \pm 0.57)\% \text{ and } \delta R_{\tau/K}|_{vSD} = (-0.88 \pm 0.58)\%$$

* Cirigliano & IR'07

4. Calculation of $R_{\tau/P} = R_{\tau/P}^{(0)}(1 + \delta R_{\tau/P}) = R_{\tau/P}^{(0)}(1 + \delta_{\tau P} - \delta_{P\mu})$

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$$\delta R_{\tau/\pi}|_{vSD} = (-1.02 \pm 0.57)\% \text{ and } \delta R_{\tau/K}|_{vSD} = (-0.88 \pm 0.58)\%$$

- ✓ Uncertainties dominated by $\delta_{\tau P}|_{vSD}$:
  - ✓ P decays within ChPT [counterterms can be determined by matching ChPT with the resonance effective approach at higher energies], whereas τ decays within resonance effective approach [no matching to determine the counterterms].
  - ✓ Estimation of the model-dependence by comparing our results with a less general scenario where only well-behaved two-point Green functions and a reduced resonance Lagrangian is used: ±0.22% and ±0.24% for the pion and the kaon case.
  - ✓ Estimation of the counterterms by considering the running between 0.5 and 1.0 GeV: ±0.52% (similar procedure in Marciano & Sirlin'93). **Conservative estimate**, since vSD counterterms affecting in P decays imply similar corrections to our estimation of the vSD counterterms in τ decays.

* Cirigliano & Rosell'07

5. Results

Contribution	$\delta R_{\tau/\pi}$	$\delta R_{\tau/K}$	Ref.
SI	$+1.05\%$	$+1.67\%$	*
rSD	$+0.15\%$	$+(0.18 \pm 0.05)\%$	**
vSD	$-(1.02 \pm 0.57)\%$	$-(0.88 \pm 0.58)\%$	new
Total	$+(0.18 \pm 0.57)\%$	$+(0.97 \pm 0.58)\%$	new

Errors are not reported if they are lower than 0.01%.

- ✓ Central values agree remarkably with DF'95, merely a coincidence: $\delta R_{\tau/\pi} = (0.16 \pm 0.14)\%$ and $\delta R_{\tau/K} = (0.90 \pm 0.22)\%$, **but** in that work:
  - ✓ problematic hadronization: form factors are different for real- and virtual-photon corrections, do not satisfy the correct QCD short-distance behavior, violate unitarity, analicity and the chiral limit at leading non-trivial orders.
  - ✓ a cutoff to regulate the loop integrals, splitting unphysically long- and short-distance regimes.
  - ✓ unrealistic uncertainties (purely $O(e^2p^2)$ ChPT size).

* Decker & Finkemeier'95
** Cirigliano & Rosell'07
** Guo & Roig'10

6. Application I: Radiative corrections in $\Gamma(\tau \to P\nu_\tau[\gamma])$

short-distance
EW correction
≈ 1.0201*

$$\Gamma(\tau \to P\nu_\tau[\gamma]) = \frac{G_F^2 |V_{uD}|^2 F_P^2}{8\pi} M_\tau^3 \left(1 - \frac{m_P^2}{M_\tau^2}\right)^2 S_{\rm EW} (1 + \delta_{\tau P})$$

✓ $\delta_{\tau P}$ includes SI and SD radiative corrections.

$$\delta_{\tau P} = \frac{\alpha}{2\pi}\left(g\left(\frac{m_P^2}{M_\tau^2}\right) + \frac{19}{4} - \frac{2\pi^2}{3} - 3\log\frac{m_\rho}{M_\tau}\right) + \delta_{\tau P}\big|_{\rm rSD} + \delta_{\tau P}\big|_{\rm vSD} = \begin{cases} \delta_{\tau\pi} = (-0.24 \pm 0.56)\% \\ \delta_{\tau K} = (-0.15 \pm 0.57)\% \end{cases}$$

* Erler'02

6. Application II: lepton universality test

$$R_{\tau/P} \equiv \frac{\Gamma(\tau \to P\nu_\tau[\gamma])}{\Gamma(P \to \mu\nu_\mu[\gamma])} = \left|\frac{g_\tau}{g_\mu}\right|_P^2 \frac{1}{2}\frac{M_\tau^3}{m_\mu^2 m_P}\frac{(1-m_P^2/M_\tau^2)^2}{(1-m_\mu^2/m_P^2)^2}\left(1+\delta R_{\tau/P}\right)$$

$$\left|\frac{g_\tau}{g_\mu}\right|_\pi = 0.9964 \pm 0.0028_{\rm th} \pm 0.0025_{\rm exp} = 0.9964 \pm 0.0038$$

$$\left|\frac{g_\tau}{g_\mu}\right|_K = 0.9857 \pm 0.0028_{\rm th} \pm 0.0072_{\rm exp} = 0.9857 \pm 0.0078$$

6. Application II: lepton universality test

$$R_{\tau/P} \equiv \frac{\Gamma(\tau \to P\nu_\tau[\gamma])}{\Gamma(P \to \mu\nu_\mu[\gamma])} = \left|\frac{g_\tau}{g_\mu}\right|_P^2 \frac{1}{2}\frac{M_\tau^3}{m_\mu^2 m_P}\frac{(1-m_P^2/M_\tau^2)^2}{(1-m_\mu^2/m_P^2)^2}\left(1+\delta R_{\tau/P}\right)$$

PDG

$\delta R_{\tau/\pi} = (0.18 \pm 0.57)\%$
$\delta R_{\tau/K} = (0.97 \pm 0.58)\%$

$$\left|\frac{g_\tau}{g_\mu}\right|_\pi = 0.9964 \pm 0.0028_{\mathrm{th}} \pm 0.0025_{\mathrm{exp}} = 0.9964 \pm 0.0038$$

$$\left|\frac{g_\tau}{g_\mu}\right|_K = 0.9857 \pm 0.0028_{\mathrm{th}} \pm 0.0072_{\mathrm{exp}} = 0.9857 \pm 0.0078$$

- ✓ π case: at 0.9σ of LU vs. 1.6σ of LU in HFLAV'21* using DF'95**
- ✓ K case: at 1.8σ of LU vs. 1.9σ of LU in HFLAV'21* using DF'95**

* HFLAV'21
** Decker & Finkemeier'95

6. Application III: CKM unitarity test in the ratio $\Gamma(\tau \to K\nu_\tau[\gamma]) / \Gamma(\tau \to \pi\nu_\tau[\gamma])$

$$\frac{\Gamma(\tau \to K\nu_\tau[\gamma])}{\Gamma(\tau \to \pi\nu_\tau[\gamma])} = \frac{|V_{us}|^2}{|V_{ud}|^2} \frac{F_K^2}{F_\pi^2} \frac{(1-m_K^2/M_\tau^2)^2}{(1-m_\pi^2/M_\tau^2)^2} (1+\delta)$$

$$\left|\frac{V_{us}}{V_{ud}}\right| = 0.2288 \pm 0.0010_{\rm th} \pm 0.0017_{\rm exp} = 0.2288 \pm 0.0020$$

6. Application III: CKM unitarity test in the ratio $\Gamma(\tau \to K\nu_\tau[\gamma])$ / $\Gamma(\tau \to \pi\nu_\tau[\gamma])$

FLAG'20*:
$F_K/F_\pi = 1.1932 \pm 0.0019$

$$\delta = \delta_{\tau K} - \delta_{\tau\pi} = +(0.10 \pm 0.80)\%$$

PDG

$$\frac{\Gamma(\tau \to K\nu_\tau[\gamma])}{\Gamma(\tau \to \pi\nu_\tau[\gamma])} = \frac{|V_{us}|^2}{|V_{ud}|^2}\frac{F_K^2}{F_\pi^2}\frac{(1-m_K^2/M_\tau^2)^2}{(1-m_\pi^2/M_\tau^2)^2}(1+\delta)$$

$$\left|\frac{V_{us}}{V_{ud}}\right| = 0.2288 \pm 0.0010_{\mathrm{th}} \pm 0.0017_{\mathrm{exp}} = 0.2288 \pm 0.0020$$

- ✓ 2.1σ away from CKM unitarity, considering $|V_{ud}|$=0.97373±0.00031**.
- ✓ To be compared with $|V_{us}/V_{ud}|$=0.2291±0.0009***, obtained with kaon semileptonic decays. Our error does not reach this level due to lack of statistics in τ decays.

* FLAG'20
** Hardy & Towner'20
*** Seng et al.'21

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PDG

$$\frac{\Gamma(\tau \to K\nu_\tau[\gamma])}{\Gamma(\tau \to \pi\nu_\tau[\gamma])} = \frac{|V_{us}|^2}{|V_{ud}|^2} \frac{F_K^2}{F_\pi^2} \frac{(1-m_K^2/M_\tau^2)^2}{(1-m_\pi^2/M_\tau^2)^2} (1+\delta)$$

Conservative estimation of the errors in δ, since we have directly propagated the uncertainties of $\delta_{\tau K}$ and $\delta_{\tau\pi}$.

$$\left|\frac{V_{us}}{V_{ud}}\right| = 0.2288 \pm 0.0010_{\mathrm{th}} \pm 0.0017_{\mathrm{exp}} = 0.2288 \pm 0.0020$$

- ✓ 2.1σ away from CKM unitarity, considering $|V_{ud}|=0.97373\pm0.00031$**.
- ✓ To be compared with $|V_{us}/V_{ud}|=0.2291\pm0.0009$***, obtained with kaon semileptonic decays. Our error does not reach this level due to lack of statistics in τ decays.

* FLAG'20
** Hardy & Towner'20
*** Seng et al.'21

6. Application IV: CKM unitarity test in $\Gamma(\tau \to K\nu_\tau[\gamma])$

$$\Gamma(\tau \to K\nu_\tau[\gamma]) = \frac{G_F^2 F_K^2}{8\pi} |V_{us}|^2 M_\tau^3 \left(1 - \frac{m_K^2}{M_\tau^2}\right)^2 S_{\rm EW} \, (1 + \delta_{\tau K})$$

$$|V_{us}| = 0.2220 \pm 0.0008_{\rm th} \pm 0.0016_{\rm exp} = 0.2220 \pm 0.0018$$

6. Application IV: CKM unitarity test in $\Gamma(\tau \to K\nu_\tau[\gamma])$

FLAG'20*: $\sqrt{2}F_K = (155.7 \pm 0.3)$ MeV

short-distance EW correction ≈ 1.0201**

$\delta_{\tau K} = (-0.15 \pm 0.57)\%$

PDG

$$\Gamma(\tau \to K\nu_\tau[\gamma]) = \frac{G_F^2 F_K^2}{8\pi} |V_{us}|^2 M_\tau^3 \left(1 - \frac{m_K^2}{M_\tau^2}\right)^2 S_{\mathrm{EW}} (1 + \delta_{\tau K})$$

$$|V_{us}| = 0.2220 \pm 0.0008_{\mathrm{th}} \pm 0.0016_{\mathrm{exp}} = 0.2220 \pm 0.0018$$

- ✓ 2.6σ away from CKM unitarity, considering $|V_{ud}|$=0.97373±0.00031***.
- ✓ To be compared with $|V_{us}|$=0.2234±0.0015^ or $|V_{us}|$=0.2231±0.0006†, obtained this last one with kaon semileptonic decays. Our error does not reach this level due to lack of statistics in τ decays.

* FLAG'20
** Erler'02
*** Hardy & Towner'20
^ HFLAV'21
† Seng et al.'21

6. Application V: constraining non-standard interactions in $\Gamma(\tau \to P\nu_\tau[\gamma])$

$$\Gamma(\tau \to P\nu_\tau[\gamma]) = \frac{G_F^2 |\widetilde{V}_{uD}|^2 F_P^2}{8\pi} M_\tau^3 \left(1 - \frac{m_P^2}{M_\tau^2}\right)^2 S_{\rm EW} \left(1 + \delta_{\tau P} + 2\Delta^{\tau P}\right)$$

Values of $\Delta^{\tau P}$ reported in the $\overline{\rm MS}$-scheme and at a scale of μ=2 GeV.

$$\Delta^{\tau P} = \epsilon_L^\tau - \epsilon_L^e - \epsilon_R^\tau - \epsilon_R^e - \frac{m_P^2}{M_\tau(m_u + m_D)}\epsilon_P^\tau = \begin{cases} \Delta^{\tau\pi} = -(0.15 \pm 0.72) \cdot 10^{-2} \\ \Delta^{\tau K} = -(0.36 \pm 1.18) \cdot 10^{-2} \end{cases}$$

6. Application V: constraining non-standard interactions in $\Gamma(\tau \to P\nu_\tau[\gamma])$

$|V_{ud}|$ = 0.97373 ± 0.00031*
$|V_{us}/V_{ud}|$ = 0.2288 ± 0.0020

FLAG'20*:
$\sqrt{2}F_\pi$ = (130.2 ± 0.8) MeV
$\sqrt{2}F_K$ = (155.7 ± 0.3) MeV

short-distance EW correction ≈ 1.0201**

$$\delta_{\tau\pi} = (-0.24 \pm 0.56)\%$$
$$\delta_{\tau K} = (-0.15 \pm 0.57)\%$$

PDG

$$\Gamma(\tau \to P\nu_\tau[\gamma]) = \frac{G_F^2 |\widetilde{V}_{uD}|^2 F_P^2}{8\pi} M_\tau^3 \left(1 - \frac{m_P^2}{M_\tau^2}\right)^2 S_{\rm EW} \left(1 + \delta_{\tau P} + 2\Delta^{\tau P}\right)$$

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- ✓ To be compared with $\Delta^{\tau\pi} = -(0.15 \pm 0.67)\cdot 10^{-2}$ of Cirigliano et al.'19^.
- ✓ To be compared with $\Delta^{\tau\pi} = -(0.12 \pm 0.68)\cdot 10^{-2}$ and $\Delta^{\tau K} = (-0.41 \pm 0.93)\cdot 10^{-2}$ of González-Solís et al.'20†.

* Hardy & Towner'20
** FLAG'20
*** Erler'02

^ Cirigliano et al.'19
† Gonzàlez-Solís et al. '20

7. Conclusions

✓ The observable and **our result**:

$$R_{\tau/P} \equiv \frac{\Gamma(\tau \to P\nu_\tau[\gamma])}{\Gamma(P \to \mu\nu_\mu[\gamma])} = \left|\frac{g_\tau}{g_\mu}\right|_P^2 R^{(0)}_{\tau/P}\left(1+\delta R_{\tau/P}\right) \quad \longrightarrow \quad \begin{cases} \delta R_{\tau/\pi} = (0.18 \pm 0.57)\% \\ \delta R_{\tau/K} = (0.97 \pm 0.58)\% \end{cases}$$

✓ Framework: ChPT for π decays and a resonance extension of ChPT for τ decays.

✓ Consistent with DF'95*, but with more robust assumptions and yielding a reliable uncertainty.

✓ Applications:

- ✓ Theoretical determination of radiative corrections in $\Gamma(\tau \to P\nu_\tau[\gamma])$.
- ✓ $|g_\tau/g_\mu|_P$ at 0.9σ (π) and 1.8σ (K) of LU, reducing HFLAV'21** disagreement with LU.
- ✓ CKM unitarity in $\Gamma(\tau\to K\nu_\tau[\gamma])/\Gamma(\tau\to\pi\nu_\tau[\gamma])$: $|V_{us}/V_{ud}| = 0.2288 \pm 0.0020$, at 2.1σ from unitarity.
- ✓ CKM unitarity in $\Gamma(\tau\to K\nu_\tau[\gamma])$: $|V_{us}| = 0.2220 \pm 0.0018$, at 2.6σ from unitarity.
- ✓ Constraining non-standard interactions in $\Gamma(\tau \to P\nu_\tau[\gamma])$: update of $\Delta^{\tau P}$.

✓ Our results have been incorporated in the very recent **HFLAV'22**.

See next talk by Alberto Lusiani

* Decker & Finkemeier'95
** HFLAV'21

7. Conclusions

Reliable NP tests for present & future exps.

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* Decker & Finkemeier'95

** HFLAV'21

Comparison with Decker & Finkemeier'95 (DF'95) in the π case

Contribution	$\delta R_{\tau\pi}$ by DF'95 [$\mu_{cut} = 1.5$ GeV]	our $\delta R_{\tau\pi}$
SI	$+0.84\%^*$	$+1.05\%$
rSD	$+0.05\%$	$+0.15\%$
vSD	$-0.49\%^*$	$-(1.02 \pm 0.57)\%$
short-distance	$-0.25\%^*$	0
Total	$+(0.16 \pm 0.14)\%^*$	$+(0.18 \pm 0.57)\%$

- ✓ Virtual corrections by DF'95 are μ_{cut}-dependent, since long- and short-distance photonic contributions were separated unphysically: figures with an asterisk are cutoff-dependent.
- ✓ The quoted error in the radiative correction of DF'95 arises from uncertainties in hadronic parameters of SD contributions and from variations in the cutoff parameter, μ_{cut}.
- ✓ For the SI contribution in DF'95 we have added to the result obtained in the point-like approximation (1.05%) the term coming from cutting off the loops at μ_{cut} (−0.21%).
- ✓ Different contributions of $\delta R_{\tau/K}$ are not provided in DF'95, which prevents a comparison.
- ✓ Although central values for the sum of all the corrections agree remarkably, this is a coincidence, since central values for the SD corrections are largely different within both approaches.