Searching for heavy neutral leptons using  $\tau$  decays at BABAR

# Claudia Patrignani representing the BABAR Collaboration

Claudia.Patrignani3@unibo.it

The 13th International Workshop on  $e^+e^-$  collisions from Phi to Psi Fudan University, Shanghai, China. August 15 - 19, 2022



#### Standard Model

Standard Model extremely successful

very large number of high precision tests – almost all passed with flying colours yet some intriguing discrepancies:

#### SM anomalies

#### Astrophysical evidence for Dark matter



#### anomalous muon magnetic moment



W mass



Gravitational lensing



Galactic rotation



Bullet Cluster





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### Physics beyond Standard Model: neutrino portal

Dark matter firmly established – "SM anomalies" tantalizing search for new physics incorporating new terms in Lagrangian



This talk focused on the "Neutrino portal"

 $\it BABAR$  is exploring also LFV, dark photon and axion portals (see talks by N. Tasneem , B. Echenard, and S. Ngan Nguyen)



# Neutrino Oscillations

Neutrino oscillation firmly established:

neutrinos have mass

PMNS matrix parametrization:

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{\rm CP}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{\rm CP}} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} P$$

 $\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$ 



but a number of questions unanswered

- Neutrino mass is small: what are the origins of this mass?
- Why is neutrino mixing so different from quark mixing?
- CP violation?
- Nature of neutrinos?

# Dirac /Majorana neutrinos?



coupling to  $\bar{\nu}s \Longrightarrow$  Majorana mass

in "standard" Standard model massless neutrinos

 $\implies$  no mixing

masses and mass hierarchy much different from charged leptons



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### Heavy Neutral Leptons?

Heavy Neutral leptons are additional neutrino states which have (some) mass but no weak hyper-charge, electric charge, weak isospin and color charge.

HNLs proposed by many BSM models to help explaining

Neutrino Oscillations, Baryon Asymmetry

Dark Matter

If neutrinos get their mass from Higgs, Yukawa couplings must be exceedingly small

Not so in see saw models with 5-dim operator and additional Majorana neutrinos

PRD 23, 165 (1981) https://doi.org/10.3389/fphy.2018.00040

Sterile neutrinos could also help explaining			
Reactor Anti-neutrino anomaly PRD 83, 073006 (2011),			
Gallium anomaly PRC 80, 015807 (2009)			
Accelerator anomalies: LSND PRD 64, 112007 (2001)			
MiniBooNE PRL 110, 161801 (2013)			
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# Predicted mass ranges for HNL

Different models may introduce HNL in a broad mass range:

- $m_4 pprox \mathcal{O}(\,\mathrm{eV}/\,c^2)$  address the so-called "oscillation anomalies"
- $m_4 pprox \mathcal{O}(\,\mathrm{keV}\!/c^2)$  provide warm dark matter candidate.
- $m_4 \approx \mathcal{O}(\text{ GeV}/c^2)$  explain deviations in SM decays.
- m<sub>4</sub> ≈ O(TeV/c<sup>2</sup>) explain Baryonic Asymmetry via low-scale scenarios of leptogenesis without conflict with other cosmological observations.

e.g.:  $\nu$ MSM model PLB 620, 17 (2005) introduces three right-handed singlet HNLs:

- Two  $\,{\rm GeV}/c^2$ -scale particles to solve origin and smallness of SM neutrino mass with see saw mechanism
- a third HNL  $keV/c^2$ -scale dark matter candidate, that also provides leptogenesis (Majorana mass term)

different methods/techniques needed to test such a variety of models

HNL in the  $\,{\rm MeV}/c^2{\rm -}\,{\rm GeV}/c^2$  scale can be searched for at existing accelerator experiments.



#### Extended PMNS Matrix and current limits on $|U_{\ell 4}|^2$

1 ... >



$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \\ \nu_L \end{pmatrix} = \begin{pmatrix} U_{e1} \\ U_{\mu1} \\ U_{\tau1} \\ U_{L1} \end{pmatrix}$$

$$\begin{array}{ccccc} U_{e2} & U_{e3} & U_{e4} \\ U_{\mu2} & U_{\mu3} & U_{\mu4} \\ U_{\tau2} & U_{\tau3} & U_{\tau4} \\ U_{L2} & U_{L3} & U_{L4} \end{array} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \\ \nu_4 \end{pmatrix}$$



limits on  $|U_{\ell n}|^2$ : current (filled); projected (solid)

less stringent limits on  $|U_{\tau 4}|^2$ can be improved by studying  $\tau$  decays



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Exploit the peculiar topology provided by the boost in  $E_{CM}$  to select  $\tau$  pairs charged daughters from each  $\tau$  in one hemisphere



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### Outline of the method

 $\tau^- o \pi^- \pi^- \pi^+ \nu_x$ determines the 2D distribution of Mass of the missing neutrino  $(m_4)$ hadronic energy  $E_h$  vs invariant mass  $m_h$ 1.0 In the center of mass:  $E_{\tau} = \frac{1}{2} E_{cm}$ Energy Fraction 0.8  $3m_{\pi^{\pm}} < m_h < m_{\pi} - m_A$ 0.6  $E_{\tau} - \sqrt{m_4^2 + q_+^2} < \frac{E_h}{E_h} < E_{\tau} - \sqrt{m_4^2 + q_-^2}$ 0.4 $m_4 = 500 \, MeV/c^2$  $m_4 = 0 MeV/c^2$ 0.2 $q_{\pm} = \frac{m_{\tau}}{2} \left( \frac{m_{h}^{2} - m_{\tau}^{2} - m_{4}^{2}}{m^{2}} \right) \sqrt{\frac{E_{\tau}^{2}}{m^{2}}} - 1 \pm \frac{E_{\tau}}{2} \sqrt{\left( 1 - \frac{(m_{h} + m_{4})^{2}}{m^{2}} \right) \left( 1 - \frac{(m_{h} - m_{4})^{2}}{m^{2}} \right)}$  $\tau^+ \rightarrow 3\pi^{\pm} + \nu_4$ 0.20.4 0.0 0.6 0.8 1.0 2D distribution as sum of  $\nu_{\tau}$  (SM,m=0) Mass Fraction component and a HNL component  $\frac{\mathrm{d}\Gamma(\tau^- \to \nu h^-)}{\mathrm{d}m_h \mathrm{d}E_h} \bigg|_{\mathrm{Total}} = \left| U_{\tau 4} \right|^2 \frac{\mathrm{d}\Gamma(\tau^- \to \nu h^-)}{\mathrm{d}m_h \mathrm{d}E_h} \bigg|_{\mathrm{HNL}} + (1 - |U_{\tau 4}|^2) \frac{\mathrm{d}\Gamma(\tau^- \to \nu h^-)}{\mathrm{d}m_h \mathrm{d}E_h} \bigg|_{\mathrm{SM}}$ 

NEW

# Background and signal simulation

Event sample has contribution from a number of sources



NEW

 $\implies$  MC simulations





# Signal MC samples





# Fit model

arXiv:2207.09575 NEW

Minimize likelihood for the observed # of events in each bin to extract 
$$|U_{\tau 4}|^2$$
  

$$\mathcal{L} = \prod_{ij} f(n_{ij}; n_{obs}, \vec{\theta}) = \prod_{ij} \frac{|\mathcal{V}_{HNL} + \nu_{\tau-SM} + \nu_{BKG}\rangle_{ij}^{(n_{obs})_{ij}} e^{-(\mathcal{V}_{HNL} + \nu_{BKG} + \nu_{\tau-SM})_{ij}}{(n_{obs})_{ij}!} \times \prod_{k} f(\theta_{k}, \tilde{\theta}_{k})$$
as a sum of the expected number of events in each bin for  
• HNL 3-prong decays  $\tau^{-} \rightarrow \pi^{+}\pi^{-}\pi^{-}\nu_{X}$ :  
 $\hat{\mathcal{V}}_{HNL,ij} = n_{HNL,ij}^{reco} = N_{\tau,gen} \cdot (|U_{\tau 4}|^2) \cdot p_{HNL,ij}$   
• SM 3-prong decays  $\tau^{-} \rightarrow \pi^{+}\pi^{-}\pi^{-}\nu_{\tau}$ :  
 $\hat{\mathcal{V}}_{\tau-SM,ij} = n_{\tau-SM,ij}^{reco} = N_{\tau,gen} \cdot (1 - |U_{\tau 4}|^2) \cdot p_{\tau-SM,ij}$   
• other sources  
 $\hat{\mathcal{V}}_{BKG,ij} = n_{BKG,ij}^{reco} = n_{\tau-other,ij}^{reco} + n_{non-\tau,ij}^{reco}$   
 $p_{i,j}$  accounts for acceptance × efficiency × Ph.Sp.  
Upper limits (Wick's theorem):  $q = -2\ln(\frac{\mathcal{L}_{H_0}(|U_{\tau 4}|_{i}^2; \hat{\theta}_{0}, data)}{\mathcal{L}_{H_1}(|\hat{U}_{\tau 4}|_{i}^2; \hat{\theta}, data)}) = -2\ln(\Delta \mathcal{L})$ 

#### Event selection

Selection optimized for



		•
Cut		Purpose
	Number of tracks	Ensure 1+3 prong topology
	Total charge on all 4 charged tracks is 0	Charge conservation
	$p_{CM}^{miss} > 0.9\% \sqrt{s}$	Suppresses non-tau backgrounds
	All tracks: $p_{trans} > 250 {\rm MeV/c}$	To reach DIRC <sup>1</sup>
	All tracks: $-0.76 < \cos(\theta) < 0.9$	Acceptance of DIRC <sup>1</sup>
	1 prong: $\frac{2p}{E} < 0.9\%$	Consistent with tau decay
	PID Requirements	Uses Electron and Muon ID algorithms
	C. Patrignani – Claudia.Patrignani3@unibo.it	PHIPSI 2022 August 15 - 19, 2022, Shanghai 14

### 2D plots: examples (data and MC)

electron-tag  $\tau^+$  and  $\tau^-$  samples studied separately for consistency

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(similarly for muon-tag)
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C. Patrignani - Claudia.Patrignani3@unibo.it

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Electron Tag

NFW/

Mass Fraction

BABAR

### Systematic uncertainties: normalization



Uncertainty	Yield Change $(\pm)$
Luminosity	0.44%
$\sigma(ee \rightarrow \tau \tau)$	0.31%
Branching Fractions (1 prong)	e: 0.22%
	$\mu: 0.22\%$
Branching Fractions (3 prong)	$3\pi: 0.57\%$
PID Efficiency	e: 2%
	$\mu: 1\%$
	$\pi: 3\%$
Bhabha Contamination	0.2%
$q\bar{q}$ Contamination (data)	0.1%
Tracking Efficiency	negligible
Detector Modeling	negligible
Beam Energy	negligible
Tau Mass	negligible



#### Systematic uncertainties: 2D shape templates

The  $\tau \to 3\pi \nu$  decay is dominated by  $\tau \to a_1 \nu$ The shape of the 3h template depends on the  $a_1$  mass and width parameters  $\implies$  systematic based on uncertainties from PDG



# Limits on $|U_{\tau 4}|^2$ vs HNL mass





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#### Summary and outlook

HNLs offer ways of explaining several observational phenomena.

The possible masses of the HNLs is model dependent and can range from  $eV/c^2$  up to very heavy masses.

In the last few years several new results have been published including results from collider-based experiments and neutrino experiments.

A new technique has been adopted by *BABAR* to search for HNLs obtaining new preliminary upper limits on  $|U_{\tau 4}|^2$  between 100 and 1300 MeV/ $c^2$ 

currently the most stringent between 300 and 1300  $\mathrm{MeV}/c^2$ 

submitted to PRD

 $\ensuremath{\textit{BABAR}}$  result is competitive with expected limits from planned experiments in the next decade

Improvements could be obtained by the same technique by Belle-II



# BACKUP



#### The BABAR experiment



#### Data samples

As of 2008/04/11 00:00

