The symbology of Feynman integrals from twistor geometries I



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8.22

2203.16112, 2207.13482, 22XX.XXXXX

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- Symbology for Amplitudes & Integrals
 - For multi-polylogarithmic (MPL) functions, symbol letters are most basic building blocks
 - Canonical DE method for general integrals (very efficient; mathematical structure?)
 - Possible mathematical structures of symbol letters help us make prediction for the alphabet \rightarrow bootstrap amplitudes/individual integrals
 - In this talk (1st part: N=4 SYM, 2nd part: general cases), we introduce a geometric way, which recovers the symbol letters from twistor geometries

One-loop integrals and Schubert problems [N. Arkani-Hamed, 21']

- For dual conformal invariant integrals, symbol letters are DCI invariants → geometrical invariants in momentum twistor space
- E.g.1:one-loop four-mass box (weight-2 MPL)

$$i - 1 \quad i \quad j - 1$$

$$l \quad \downarrow \quad j \quad j - 1$$

$$l \quad \downarrow \quad \downarrow \quad j \quad j - 1$$

$$l \quad \downarrow \quad \downarrow \quad j \quad j = \int_{AB} \frac{\langle i - 1ik - 1k \rangle \langle j - 1jl - 1l \rangle}{\langle ABi - 1i \rangle \langle ABk - 1k \rangle \langle ABl - 1l \rangle}$$

$$\frac{1}{2\Delta_{i,j,k,l}} \left(v \otimes \frac{z_{i,j,k,l}}{\bar{z}_{i,j,k,l}} + u \otimes \frac{1 - \bar{z}_{i,j,k,l}}{1 - z_{i,j,k,l}} \right)$$

$$u = \frac{\langle i - 1ij - 1j \rangle \langle k - 1kl - 1l \rangle}{\langle i - 1ik - 1k \rangle \langle j - 1jl - 1l \rangle}, \quad v = \frac{\langle i - 1il - 1l \rangle \langle j - 1jk - 1k \rangle}{\langle i - 1ik - 1k \rangle \langle j - 1jl - 1l \rangle}$$

$$z_{i,j,k,l} \bar{z}_{i,j,k,l} = u, \quad (1 - z_{i,j,k,l})(1 - \bar{z}_{i,j,k,l}) = v.$$

$$\Delta_{i,j,k,l} = \sqrt{(1 - u - v)^2 - 4uv}$$

One-loop integrals and Schubert problems [N. Arkani-Hamed, 21']

 For dual conformal invariant integrals, symbol letters are DCI invariants → geometrical invariants in momentum twistor space

$$i-1 \quad i \\ l \longrightarrow j \\ k \quad k-1 \qquad = \int_{AB} \frac{\langle i-1ik-1k \rangle \langle j-1jl-1l \rangle}{\langle ABj-1j \rangle \langle ABk-1k \rangle \langle ABl-1l \rangle} \\ \frac{1}{2\Delta_{i,j,k,l}} \left(v \otimes \frac{z_{i,j,k,l}}{\bar{z}_{i,j,k,l}} + u \otimes \frac{1-\bar{z}_{i,j,k,l}}{1-z_{i,j,k,l}} \right)$$

• LS of box: AB intersects with 4 external lines simultaneously [Schubert, 19th century; N. Arkani-Hamed et al, 10']

$$i-1 \qquad j-1 \qquad k-1 \qquad l-1$$

$$\alpha_{1} \qquad \beta_{1} \qquad \gamma_{1} \qquad \delta_{1} \qquad (AB)_{1} \qquad \frac{(\alpha_{1},\beta_{1})(\gamma_{1},\delta_{1})}{(\alpha_{1},\gamma_{1})(\beta_{1},\delta_{1})} = z_{i,j,k,l}, \\ \frac{(\alpha_{1},\beta_{1})(\gamma_{1},\beta_{1})}{(\alpha_{1},\gamma_{1})(\beta_{1},\delta_{1})} = 1-z_{i,j,k,l} \qquad \frac{(X_{1},X_{3})(X_{2},X_{4})}{(X_{1},X_{4})(X_{2},X_{3})} := \frac{\langle X_{1}X_{3}I \rangle \langle X_{2}X_{4}I \rangle}{\langle X_{1}X_{4}I \rangle \langle X_{2}X_{3}I \rangle}$$

$$\alpha_{2} \qquad \beta_{2} \qquad \gamma_{2} \qquad \delta_{2} \qquad (AB)_{2} \qquad \frac{(\alpha_{2},\beta_{2})(\gamma_{2},\delta_{2})}{(\alpha_{2},\gamma_{2})(\beta_{2},\delta_{2})} = \bar{z}_{i,j,k,l}, \\ \frac{(\alpha_{2},\delta_{2})(\gamma_{2},\beta_{2})}{(\alpha_{2},\gamma_{2})(\beta_{2},\delta_{2})} = 1-\bar{z}_{i,j,k,l} \qquad \text{I: arbitrary reference line}$$

Two-loop Symbology and twistor geometries 9-point double-box integral (weight-4 MPL)

$$S(I_9) = \sum A_1 \otimes A_2 \otimes B \otimes C$$



 $A_1 \otimes A_2$ (first-two-entries) : letters from one-loop intersections on internal lines

B (third entries): ? ? (rational letters & algebraic letters containing four-mass square roots)

C (last entries): ? ? (4 letters containing two-loop square roots)

• <u>B: cross-ratios from external lines</u>

After two one-loop Schubert problems are solved, each shared external lines have four intersections as well



degeneration: algebraic letters & rational letters ALL 3rd entry of 9-point double-box integral can be generated by this way (n<9 letters of MHV/NMHV amplitudes in N=4 SYM 6pt: 9 rational; 7pt: 42 rational, 8pt: 272 rational+18 algebraic)



<u>C: At 2 loops, new topologies lead to new intersections and correspondingly new letters</u>

$$E = \alpha_1 Z_5 + \beta_1 Z_6 + Z_4, \ F = \alpha_2 Z_5 + \beta_2 Z_6 + Z_7$$

$$G = \gamma_1 Z_7 + \delta_1 Z_8 + Z_9, \ H = \gamma_2 Z_7 + \delta_2 Z_8 + Z_6$$

$$\alpha_2 = \delta_2 = 0, \beta_2 = \frac{1}{\gamma_2}$$

• Solving 7 on-shell conditions

Jacobian

factor $\langle 7 \ (89) \cap (612) \ (645) \cap (623) \rangle - \langle 6 \ (89) \cap (712) \ (745) \cap (723) \rangle \gamma_2^2 + (\langle 6 \ (45) \cap (236) \ (789) \cap (712) \rangle - \langle 7 \ (45) \cap (237) \ (689) \cap (612) \rangle) \gamma_2$



 γ_{12}

• (composite) LS: jacobian=0, two solutions for γ_2 , two intersections ϵ_1 and ϵ_2 a new square root Δ_9 $(CD)_3 (EF)_1 (GH)_1 (CD)_1 (CD)_4 (EF)_2 (GH)_2 (CD)_2$

 γ_{21}

 γ_{11}

 γ_{22}

<u>C: At 2 loops, new topologies lead to new intersections and correspondingly new letters</u>

Using two intersections, we can construct:

 $\frac{(\epsilon_1,\gamma_{11})(\epsilon_2,\gamma_{12})}{(\epsilon_2,\gamma_{11})(\epsilon_1,\gamma_{12})} = \frac{(1+az_{2,5,7,9})(1+b\bar{z}_{2,5,7,9})}{(1+a\bar{z}_{2,5,7,9})(1+bz_{2,5,7,9})}, \quad \frac{(\epsilon_1,\gamma_{21})(\epsilon_2,\gamma_{22})}{(\epsilon_2,\gamma_{21})(\epsilon_1,\gamma_{22})} = \frac{(\epsilon_1,\gamma_{11})(\epsilon_2,\gamma_{12})}{(\epsilon_2,\gamma_{11})(\epsilon_1,\gamma_{12})}\Big|_{reflection}$

a & b: combinations involving Δ_9 last entries of the db integral!





9-point double-box integral (weight-4 MPL)

$$S(I_9) = \sum A_1 \otimes A_2 \otimes B \otimes C$$



 $A_1 \otimes A_2$ (first-two-entries) : letters from one-loop intersections on internal lines

B (third entries): letters from one-loop intersections on external lines

C (last entries): letters from two-loop intersections on external lines

More general integrals?...

Elliptic Schubert problems and letters [in progress w. M.Wilhelm et al]

Cannot take residues in all 4L variables→ rigidity
 Rigidity=1: elliptic curve, Rigidity=2, K3 surface …
 e.g. 10-pt double-box integral

$$E = \alpha_1 Z_6 + \beta_1 Z_7 + Z_5, \ F = \gamma_1 Z_6 + \delta_1 Z_7 + Z_8$$
$$G = \alpha_2 Z_{10} + \beta_2 Z_1 + Z_2, \ H = \gamma_2 Z_{10} + \delta_2 Z_1 + Z_3$$







Elliptic Schubert problems and letters [in progress w. M.Wilhelm et al]

e.g. 10-pt double-box integral[2106.14902]

Elliptic "cross-ratios:"

 $\frac{1}{\omega_1} \int_a^\omega \frac{\mathrm{d}x}{\sqrt{P(x)}}$

a & b: two intersections from another Schubert problem ω_1 : elliptic period





Successfully reproduce the last entries of 10-point db integrals and give a predictions of 12-point db integrals \rightarrow bootstrapping

More general cases: K3 last entries?...

Summary

- Twistor geometries provide a geometric explanations for the symbol letters, both for L-loop amplitudes and integrals in N=4 SYM
- Through this way, we can predict symbol letters for individual integrals and finally bootstrap them instead of direct computations
- This method can be generalized to elliptic cases, and proves to be useful for 10-/12-pt double-box integral
- Relations to cut integrals/diagrammatic coactions...
- K3 integrals?...

Thanks!



Symbology of Feynman Integrals from Twistor Geometries (II)

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August 24, 2022

Based on [2207.13482], also see [2203.16112] [2206.04609]

Momentum Twistors and D_{ual}**C**_{onformal}**I**_{nvariance} SO(6) = SL(4)



Expressions independent of I_{∞} are DCl, e.g., cross-ratios. $\frac{(x_a - x_b)^2 (x_c - x_d)^2}{(x_a - x_d)^2 (x_b - x_c)^2} = \frac{\langle a - 1ab - 1b \rangle \langle c - 1cd - 1d \rangle}{\langle a - 1ad - 1d \rangle \langle b - 1bc - 1c \rangle}$

Momentum Twistors and D_{ual}**C**_{onformal}**I**_{nvariance} SO(6) = SL(4)



Expressions independent of I_{∞} are DCl, e.g., cross-ratios. $\frac{(x_a - x_b)^2 (x_c - x_d)^2}{(x_a - x_d)^2 (x_b - x_c)^2} = \frac{\langle a - 1ab - 1b \rangle \langle c - 1cd - 1d \rangle}{\langle a - 1ad - 1d \rangle \langle b - 1bc - 1c \rangle}$

Momentum Twistors and Dual**C**onformal Invariance SO(6) = SL(4)



$$\frac{(x_a - x_b)^2 (x_c - x_d)^2}{(x_a - x_d)^2 (x_b - x_c)^2} = \frac{\langle a - 1ab - 1b \rangle \langle c - 1cd - 1d \rangle}{\langle a - 1ad - 1d \rangle \langle b - 1bc - 1c \rangle}$$

Planar Integrals (DCI Kinematics)



$$\frac{dx_{0}}{(x_{0} - x_{2})^{2}(x_{0} - x_{4})^{2}(x_{0} - x_{6})^{2}(x_{0} - x_{8})^{2}} = \int \frac{d\mu_{AB} \langle 121_{\otimes} / \langle 0 | 1_{\otimes} / \langle 0 | 1_{\otimes}$$

2

Planar Integrals (Generic Kinematics)



Generically, an integral may depend on $x_{\infty} \sim I_{\infty}$.

$$\frac{1}{(x_0 - x_a)^2} = \frac{\langle ABI_{\infty} \rangle \langle a - 1aI_{\infty} \rangle}{\langle ABa - 1a \rangle}$$
$$d^4 x_0 = \frac{d^4 A d^4 B}{GL(2)} \frac{1}{\langle ABI_{\infty} \rangle^4} =: \frac{d\mu_{AB}}{\langle ABI_{\infty} \rangle^4}$$

$$\int \frac{\mathrm{d}^4 x_0}{(x_0 - x_2)^2 (x_0 - x_4)^2 (x_0 - x_6)^2} = \int \frac{\mathrm{d}\mu_{AB} \langle 12I_{\infty} \rangle \langle 34I_{\infty} \rangle \langle 56I_{\infty} \rangle}{\langle AB12 \rangle \langle AB34 \rangle \langle AB56 \rangle \langle ABI_{\infty} \rangle}$$

Non-DCI Alphabets from DCI



DCI alphabet from kinematic cluster algebras, missing only a few letters. [Henn et al., 2012.12285]



DCI alphabet from DCI Schubert [see YQL's talk], covering the entire alphabet. [He et al., 2206.04609]

Non-DCI Alphabets from DCI



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Massive non-DCI equivalent to DCI

Study non-DCI integrals on their own?

One-Loop Building Blocks: Triangles



One-Loop Building Blocks: Triangles



One-Loop Building Blocks: Triangles







One-Loop: 0m Pentagon



One-Loop: 0m Pentagon



Two-Loop: 2me Slashed-Box

















Surprising! "Clever" routing only?



Non-planar root Σ_5 and 5 associated algebraic letters





Non-planar root Σ_5 and 5 associated algebraic letters













Summary & Outlook

- Full coverage of alphabets Geometric handle on symbology
- Exact matches Predicting letters and bootstrapping integrals

- Systematic construction (no more & no less)
- Elliptic, K3, non-planar, non-planar elliptic, ...
- Detailed structure of symbols / alphabets



Back-up Slides



For reviews, see [1012.6032] [1204.3894]

Parametrization of Kinematics



3-dim Projective Geometry

No notion of "parallel"

- Lines and planes always intersect $(ab) \cap (cde) = Z_a \langle bcde \rangle Z_b \langle acc \rangle$
- Planes and planes always intersect

$$\begin{aligned} (ab) \cap (cde) &= Z_a \langle bcde \rangle - Z_b \langle acde \rangle \\ (abc) \cap (def) &= (ab) \langle cdef \rangle + (bc) \langle adef \rangle \\ &+ (ca) \langle bdef \rangle \end{aligned}$$

$\frac{m(m-3)}{2}$ cross-ratios for m points on a line

- Write {Z₁ ~ Z_m} as linear combinations of any two points {P, Q} on the line.
- Compute cross-ratios using minors of C, or using any reference line I_{ref}.

$$(Z_1 \quad \cdots \quad Z_m)_{4 \times m} = (P \quad Q)_{4 \times 2} \mathbf{C}_{2 \times m}$$

$$u_{ij} = \frac{(i,j-1)(i-1,j)}{(i,j)(i-1,j-1)} = \frac{\langle ij-1I_{\text{ref}} \rangle \langle i-1jI_{\text{ref}} \rangle}{\langle ijI_{\text{ref}} \rangle \langle i-1j-1I_{\text{ref}} \rangle}$$

2-loop Schubert

Maximal cut (with Jacobian) (7+1) equations, 8 unknowns

$$\begin{pmatrix} A \\ B \end{pmatrix} = \begin{pmatrix} 1 & 0 & \alpha_1 & \beta_1 \\ 0 & 1 & \gamma_1 & \delta_1 \end{pmatrix} \begin{pmatrix} Z_1 \\ P \\ Z_2 \\ Q \end{pmatrix}, \qquad \begin{pmatrix} C \\ D \end{pmatrix} = \begin{pmatrix} 1 & 0 & \alpha_2 & \beta_2 \\ 0 & 1 & \gamma_2 & \delta_2 \end{pmatrix} \begin{pmatrix} Z_2 \\ P \\ Z_3 \\ Q \end{pmatrix}$$

1. Solve cutting conditions for any 7 unknowns $\langle AB12 \rangle = \langle AB56 \rangle = \langle ABPQ \rangle = \langle ABCD \rangle = \langle CD23 \rangle = \langle CD45 \rangle = \langle CDPQ \rangle = 0$ $\implies \begin{cases} \beta_1 = \gamma_1 = \beta_2 = \gamma_2 = 0, \ \alpha_1 = \alpha_1^{(1)}(\delta_2), \ \delta_1 = \delta_1^{(1)}(\delta_2), \ \alpha_2 = \alpha_2^{(1)}(\delta_2) \\ \beta_1 = \gamma_1 = \beta_2 = \gamma_2 = 0, \ \alpha_1 = \alpha_1^{(2)}(\delta_2), \ \delta_1 = \delta_1^{(2)}(\delta_2), \ \alpha_2 = \alpha_2^{(2)}(\delta_2) \end{cases}$

2. Solve the Jacobian for the remaining unknown

 $\frac{\partial(\operatorname{cutting eqns.})}{\partial(\alpha_1, \beta_1, \gamma_1, \delta_1, \alpha_2, \beta_2, \gamma_2)}\Big|_{\operatorname{cutting sol.}^{(1)}} = 0 \implies \delta_2^{(1,1)} = \cdots \text{ or } \delta_2^{(1,2)} = \cdots$ $\frac{\partial(\operatorname{cutting eqns.})}{\partial(\alpha_1, \beta_1, \gamma_1, \delta_1, \alpha_2, \beta_2, \gamma_2)}\Big|_{\operatorname{cutting sol.}^{(2)}} = 0 \implies \delta_2^{(2,1)} = \cdots \text{ or } \delta_2^{(2,2)} = \cdots$



Generically, 4 Schubert solutions.

In this example, it happens that $\delta_1 = \delta_2$. 2 intersections on I_{∞} , corresponding to $\delta_{1,2} = \frac{\cdots \pm \Delta_{nc}}{2}$.