

Discussion of EPOL

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极化讨论会议

Outline

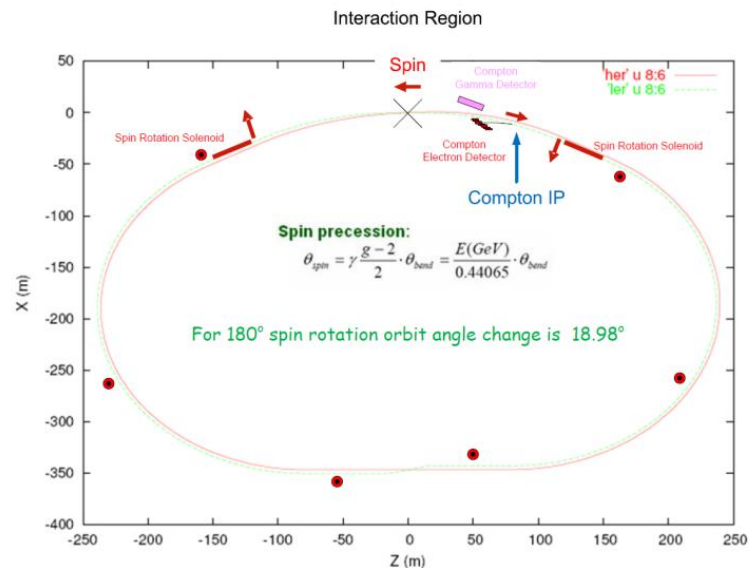
□ 1. Super-B

<https://cds.cern.ch/record/1296190/files/arXiv:1009.6178.pdf>

□ 2. FCC

<https://arxiv.org/pdf/2208.00585.pdf>

Super-B (Longitudinal polarization)		
Physical motivation ($\Delta P/P$)	<1%	
Method	Compton polarimeter	
Location	Upstream of the IR	
Physical IP/Spin projection in Compton IP	1/0.968	
Systematical uncertainties	1 mrad in the orbit	0.25%
	Beam energy uncertainty 20MeV at the 4.18 GeV	0.2%



- 不放在下游的原因：
 - the space at this location is minimal to locate the Compton IP
 - severe backgrounds from the $e^+ e^-$ collisions would be intolerable in the Compton gamma and electron detectors.
- The [orbit angle change](#) for π spin rotation is -0.3312 radians at 4.18 GeV. The selected location of the Compton IP in a magnetic field free region has [an orbit angle change](#) of -0.3580 radians between the Compton IP and the Interaction Region resulting in the spin direction ~ 14 degrees from longitudinal. At the Compton IP the longitudinal spin projection is 0.968. The longitudinal polarization at the IR will be larger by 1/0.968 than that measured in the Compton polarimeter

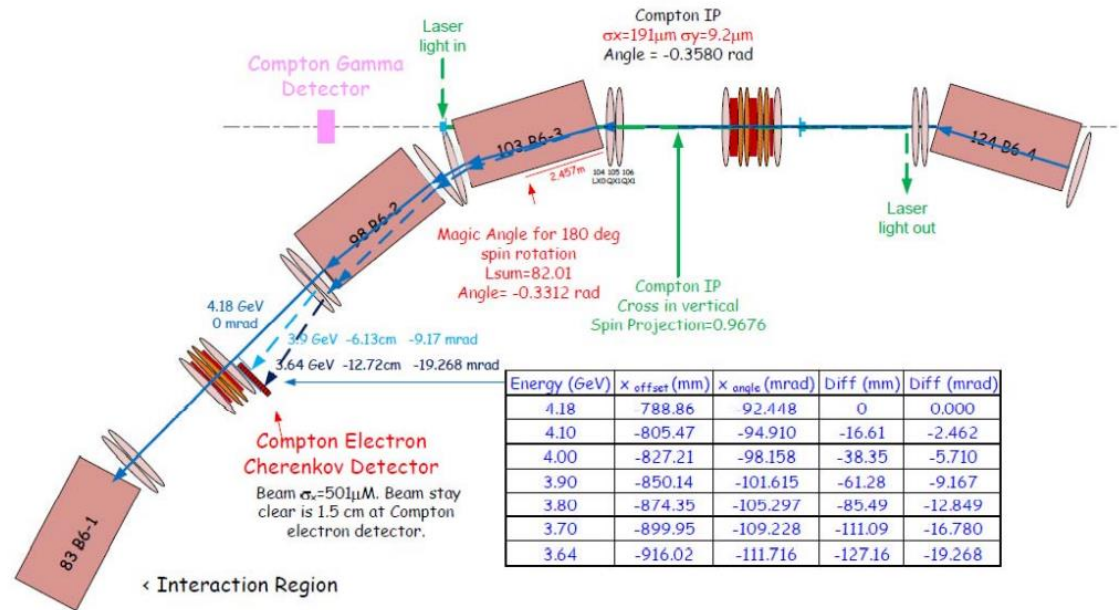


Figure 16.9: Layout of the Compton polarimeter.

Compton polarimeter

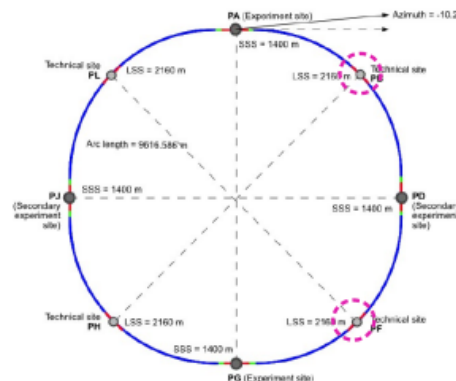
Laser	Nd:YLF circularly polarized laser
Laser length	10 ps
Laser photons	2.3×10^{10}
Electron beam	4.06 GeV
Compton edge electrons	3.64 GeV
Compton rates	$\sim 196980 / \text{sec}$

Table 16.4: Systematic errors expected for the polarization measurement.

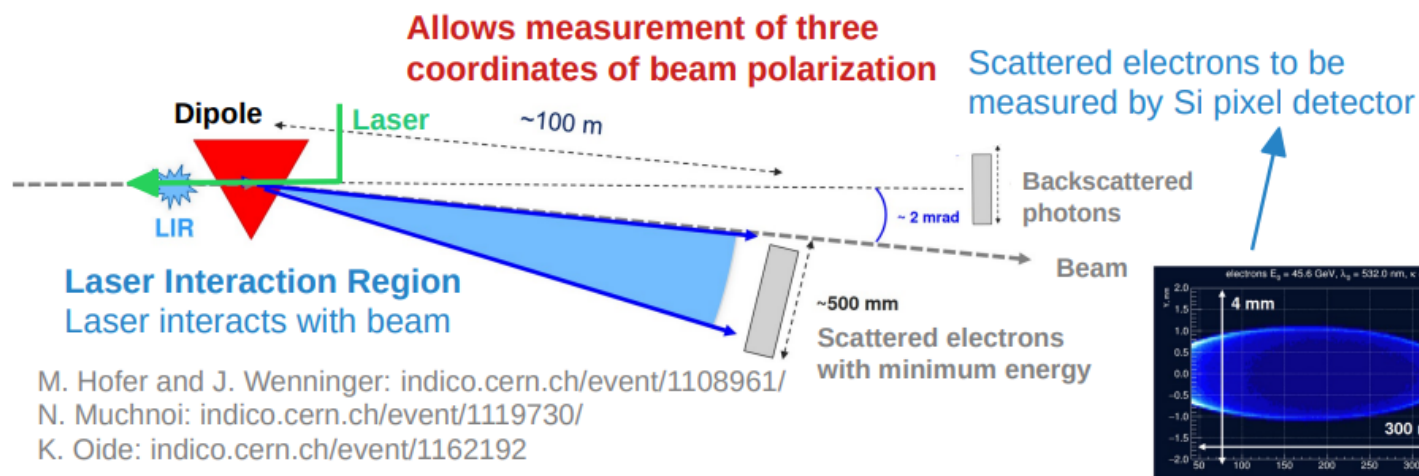
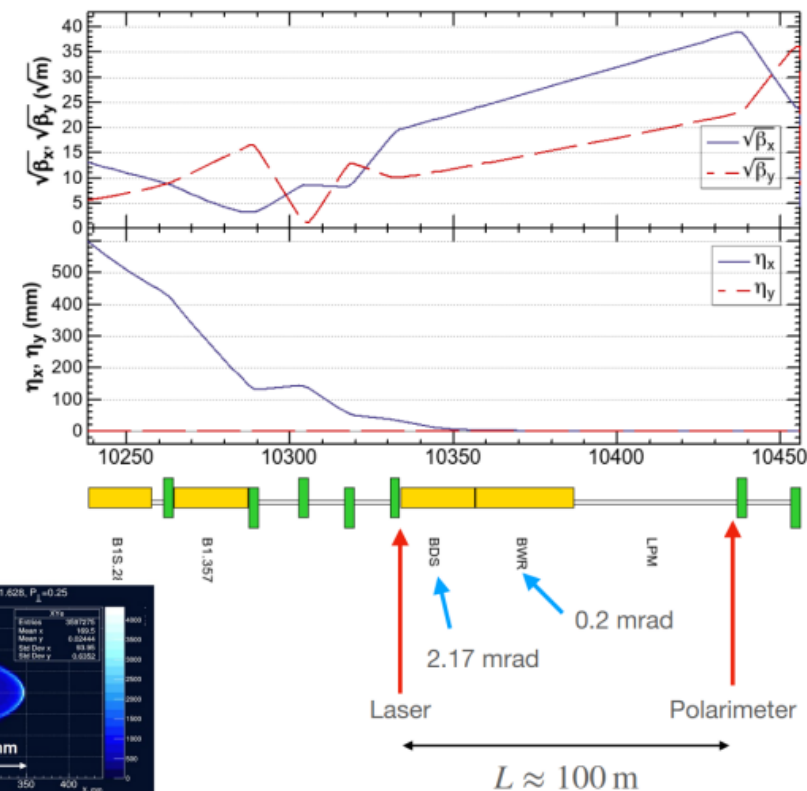
Item	$\delta P/P$
Laser Polarization	$< 0.1\%$
Background uncertainty	$< 0.25\%$
Linearity of phototube response	$< 0.25\%$
Uncertainty in dP (Difference between the luminosity weighted polarization and the Compton IP polarization. Includes uncertainties due to beam energy and direction uncertainties.)	$< 0.4\%$
Uncertainty in asymmetry analyzing power	$\sim 0.5\%$
Total Systematic Error	$< 1.0\%$

Polarimeter

- One polarimeter per beam
- First definition of specifications
 - 2 mrad angle
 - 100 m drift space
 - 2 m space for LIR (monitoring of location to be designed)



Polarimeter implemented in straight section without IP or RF



M. Hofer and J. Wenninger: indico.cern.ch/event/1108961/
 N. Muchnoi: indico.cern.ch/event/1119730/
 K. Oide: indico.cern.ch/event/1162192/

A Stokes parameters

The Stokes parameters describe the polarization state of electromagnetic radiation. Their definition is slightly different in different sources, so below are the definition that we use here.

- $\xi_0 = E_x^2 + E_y^2$ is the intensity of light.

With normalization $E_x^2 + E_y^2 = 1$ for 100% polarized laser radiation $\xi_0 = \sqrt{\xi_1^2 + \xi_2^2 + \xi_3^2} = 1$.

- $\xi_1 = E_x^2 - E_y^2$.

$$E_x = 1, E_y = 0, \xi_1 = +1:$$

100% linear polarization along x-axis.

$$E_x = 0, E_y = 1, \xi_1 = -1:$$

100% linear polarization along y-axis.

- $\xi_2 = 2E_x E_y \cos(\delta)$.

$$E_x = E_y, \delta = 0, \xi_2 = +1:$$

100% linear polarization along $\varphi = +\pi/4$.

$$E_x = E_y, \delta = \pi, \xi_2 = -1:$$

100% linear polarization along $\varphi = -\pi/4$.

- $\xi_3 = 2E_x E_y \sin(\delta)$.

$$E_x = E_y, \delta = +\pi/2, \xi_3 = +1:$$

100% right circular polarization.

$$E_x = E_y, \delta = -\pi/2, \xi_3 = -1:$$

100% left circular polarization.

FCC-Compton polarimeter

1.1 Cross section

ICS cross section depends on polarization states of all initial and final particles. In case of averaging over the polarizations of the final states **the cross section depends solely from the initial photon and electron polarizations**. Stokes parameters ξ_1, ξ_2, ξ_3 describe the polarization of laser light as it is explained in Appendix A. The electron beam polarization has three components $\zeta_x, \zeta_y, \zeta_z$, the total degree of polarization $\sqrt{\zeta_x^2 + \zeta_y^2 + \zeta_z^2} \in [0 : 1]$. We take differential cross section from ref. [34] and after Lorentz transformations it is represented in u and φ variables by the sum of the six terms:

$$\begin{aligned}
 &\text{Unpolarized polarized} \longrightarrow \frac{1}{r_e^2} \frac{d\sigma_0}{du d\varphi} = \frac{1}{\kappa(1+u)^3} \left[1 + (1+u)^2 - 4\frac{u}{\kappa} \left(1 - \frac{u}{\kappa} \right) (1+u) \right], \\
 &\text{Laser linearly polarized} \left\{ \begin{aligned} \frac{1}{r_e^2} \frac{d\sigma_{\xi_1}}{du d\varphi} &= \frac{4\xi_1}{\kappa(1+u)^2} \frac{u}{\kappa} \left(1 - \frac{u}{\kappa} \right) \cos(2\varphi), \\ \frac{1}{r_e^2} \frac{d\sigma_{\xi_2}}{du d\varphi} &= \frac{4\xi_2}{\kappa(1+u)^2} \frac{u}{\kappa} \left(1 - \frac{u}{\kappa} \right) \sin(2\varphi), \end{aligned} \right. \\
 &\text{Laser circularly polarized} \left\{ \begin{aligned} \frac{1}{r_e^2} \frac{d\sigma_x}{du d\varphi} &= \frac{-2\xi_3\zeta_x}{(1+u)^3} \frac{u}{\kappa} \sqrt{\frac{u}{\kappa} \left(1 - \frac{u}{\kappa} \right)} \cos(\varphi), \\ \frac{1}{r_e^2} \frac{d\sigma_y}{du d\varphi} &= \frac{-2\xi_3\zeta_y}{(1+u)^3} \frac{u}{\kappa} \sqrt{\frac{u}{\kappa} \left(1 - \frac{u}{\kappa} \right)} \sin(\varphi), \\ \frac{1}{r_e^2} \frac{d\sigma_z}{du d\varphi} &= \frac{\xi_3\zeta_z}{(1+u)^3} \frac{u}{\kappa} (u+2) \left(1 - 2\frac{u}{\kappa} \right). \end{aligned} \right. \tag{1.7} \\
 &\text{electron three polarization}
 \end{aligned}$$

FCC-Compton polarimeter

check the fitting procedure. The subject of this study is the trivial case when the Stokes vector of laser polarization is $[\xi_1, \xi_2, \xi_3] = [0, 0, 1]$ and the electron beam is unpolarized $\zeta_x, \zeta_y, \zeta_z = [0, 0, 0]$.

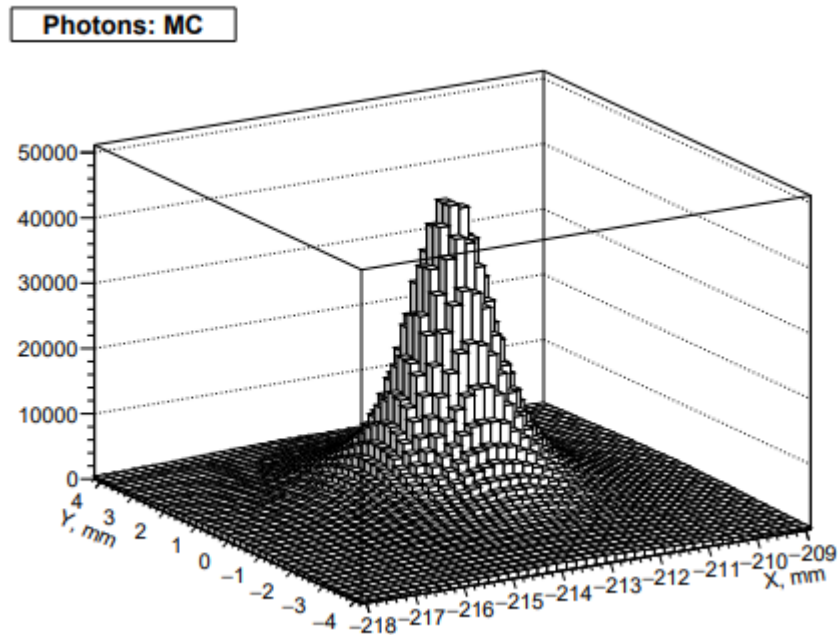


Table 4. Photon spot fit results.

$X_0 = -213.538 \pm 0.001 \text{ mm}$	$Y_0 = -0.002 \pm 0.001 \text{ mm}$
$\sigma_x = 255 \pm 3 \text{ } \mu\text{m}$	$\sigma_y = 30 \pm 18 \text{ } \mu\text{m}$
$\xi_1 = 0.000 \pm 0.002$	$\xi_2 = -0.001 \pm 0.001$
$\xi_3 \zeta_x = 0.004 \pm 0.006$	$\xi_3 \zeta_y = -0.008 \pm 0.006$
$\xi_3 \zeta_z = 0.000 \pm 0.002$	$\chi^2/\text{NDF} = 9796.9/9990$

FCC-Compton polarimeter

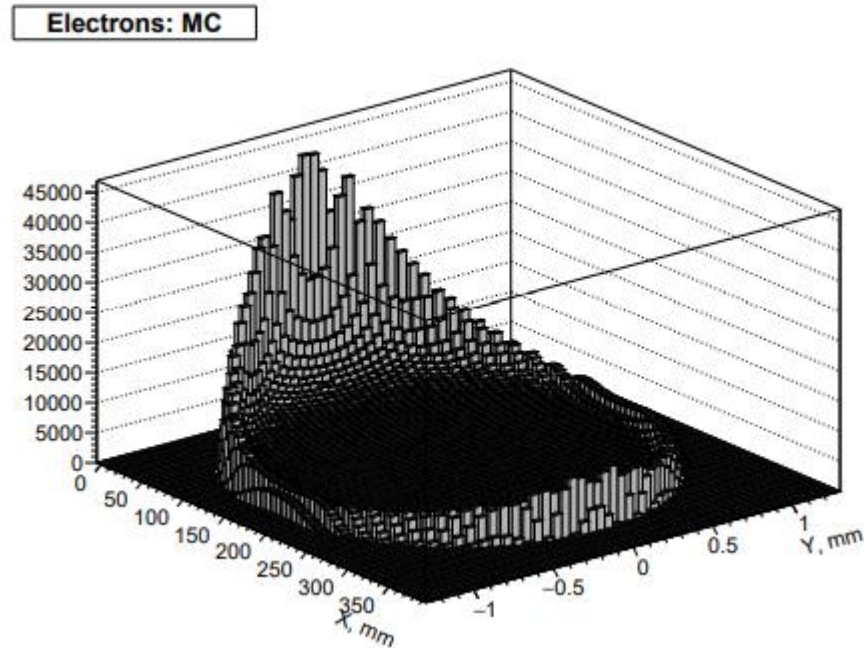


Table 5. Scattered electrons ellipse fit results.

$X_1 = 0.0035 \pm 0.0016$ mm	$X_2 = 347.635 \pm 0.003$ mm
$Y_1 = -1.0682 \pm 0.0001$ mm	$Y_2 = 1.0682 \pm 0.0001$ mm
$\sigma_x = 320.7 \pm 1.5$ μ m	$\sigma_y = 27.06 \pm 0.03$ μ m
$\xi_1 = 0.001 \pm 0.001$	$\xi_2 = 0.432 \pm 0.198$
$\xi_3 \zeta_x = 1.000 \pm 0.195$	$\xi_3 \zeta_y = -0.001 \pm 0.002$
$\xi_3 \zeta_z = 0.000 \pm 0.001$	$\chi^2/\text{NDF} = 51568.9/52270$
$\varepsilon_0 = 45.5997 \pm 0.0007$ GeV	

$$\varepsilon_0 = \frac{(mc^2)^2}{4\omega_0} \frac{X_2 - X_1}{X_1 - X_0}.$$

FCC-Compton polarimeter

the electron beam. The Stokes vector of laser polarization is chosen as $[\xi_1, \xi_2, \xi_3] = [0.1, 0.1, 0.99]$ with small amount of vestigial linear polarization which may exist due to imperfect polarization control. The set of the electron beam polarization parameters is $[\zeta_x, \zeta_y, \zeta_z] = [0.1, 0.25, 0.1]$. The expected polarization of the pilot electron bunches at FCC-ee (averaged over thousands beam revolutions) has ζ_y component only. The ζ_x and ζ_z components are added to the simulations in order to investigate the possibility of measuring the electron beam polarization in general case.

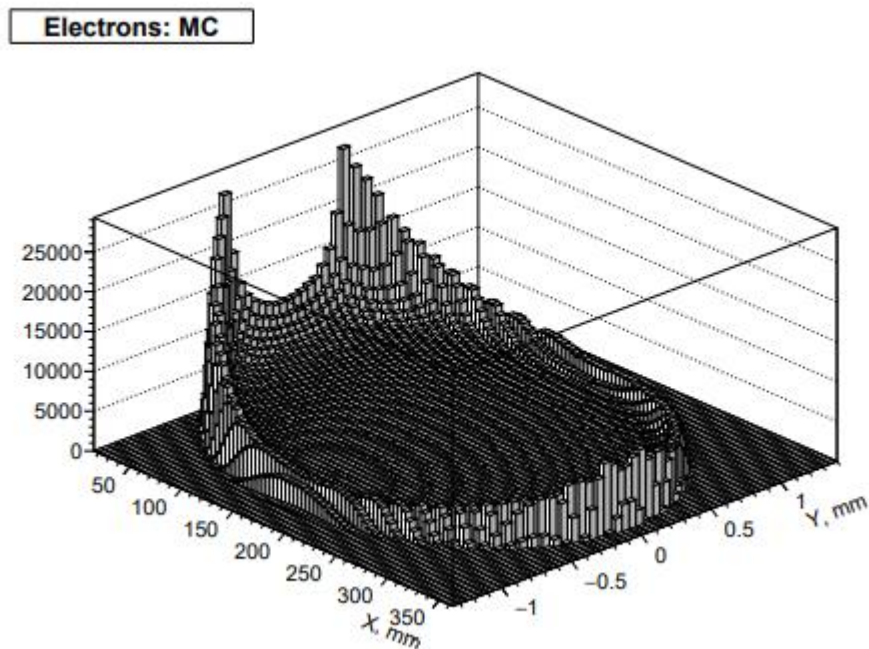


Table 6. Photon spot fit results.

$X_0 = -213.539 \pm 0.002$ mm	$Y_0 = 0.000 \pm 0.001$ mm
$\sigma_x = 246 \pm 4$ μ m	$\sigma_y = 13 \pm 70$ μ m
$\xi_1 = 0.102 \pm 0.002$	$\xi_2 = 0.100 \pm 0.001$
$\xi_3 \zeta_x = 0.095 \pm 0.007$	$\xi_3 \zeta_y = 0.247 \pm 0.006$
$\xi_3 \zeta_z = 0.105 \pm 0.002$	$\chi^2/\text{NDF} = 9935.8/9990$

Table 7. Scattered electrons ellipse fit results.

$X_1 = 0.013 \pm 0.007$ mm	$X_2 = 347.632 \pm 0.004$ mm
$Y_1 = -1.0682 \pm 0.0001$ mm	$Y_2 = 1.0684 \pm 0.0001$ mm
$\sigma_x = 319.6 \pm 4.3$ μ m	$\sigma_y = 27.15 \pm 0.03$ μ m
$\xi_1 = 0.100 \pm 0.001$	$\xi_2 = 0.100$
$\xi_3 \zeta_x = 0.099$	$\xi_3 \zeta_y = 0.246 \pm 0.002$
$\xi_3 \zeta_z = 0.099 \pm 0.001$	$\chi^2/\text{NDF} = 50152.7/51245$
$\varepsilon_0 = 45.5959 \pm 0.0025$ GeV	

Discussion and Summary

1. 从物理角度推导截面公式
2. 探测器角度摆放