

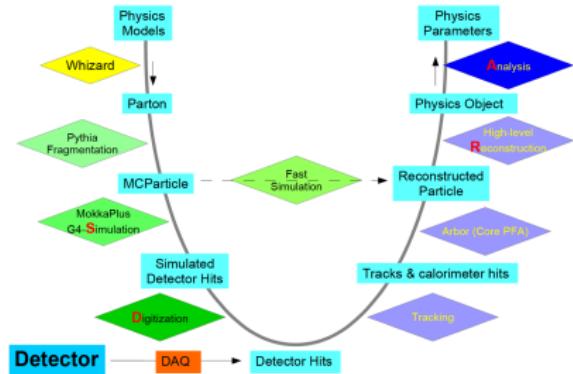
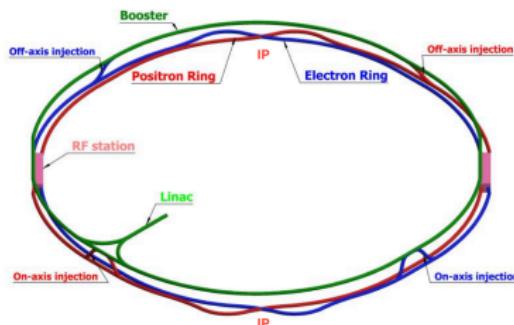
The measurement of Higgs decay (bb/cc/gg, rare, and FCNC) at the CEPC

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Higgs 2023

12.01.2023

CEPC



| Operation mode | \sqrt{s} GeV | Running time year | $\int L$ ab^{-1} | Event yields |
|----------------|-------------------|----------------------|-----------------------|---------------------------------------|
| H | 240 | 7/10 | 5.6/20 | $1 \times 10^6 / 4 \times 10^6$ |
| Z | 91.2 | 2 | 16/100 | $7 \times 10^{11} / 3 \times 10^{12}$ |
| W^-W^+ | 160 | 1 | 2.6/6 | $2 \times 10^7 / 1 \times 10^8$ |
| $t\bar{t}$ | 360 | 5 | 1 | 5×10^5 |

Motivation

- Measuring the branching fractions of the $H \rightarrow b\bar{b}/c\bar{c}/gg$ decays is one of the core CEPC physics objectives.
- The rare and FCNC hadronic decays of the Higgs boson are of great interest to NP.

Contents

- The relative accuracy of signal strength measurement of $\nu\nu H(H \rightarrow b\bar{b}, c\bar{c}, gg)$ and $qqH(H \rightarrow b\bar{b}, c\bar{c}, gg)$. ($\sqrt{s} = 240$ GeV, $\int L = 5.6 ab^{-1}$)
- key performance: flavor tagging performance and color singlet identification
- The upper limits on branching ratio of Higgs rare and FCNC hadronic decays at the CEPC. ($\sqrt{s} = 240$ GeV, $\int L = 20 ab^{-1}$)
- summary

The analysis of $H \rightarrow b\bar{b}/c\bar{c}/gg$ in $\nu\nu H$ channel.

$\sqrt{s} = 240$ GeV, $\int L = 5.6 ab^{-1}$

| | $v\nu H q\bar{q}/gg$ | 2f | SW | SZ | WW | ZZ | Mixed | ZH | $\frac{\sqrt{S+B}}{S} (%)$ |
|---------------------------------------|----------------------|--------|--------|--------|--------|--------|--------|--------|----------------------------|
| total | 178890 | 8.01E8 | 1.95E7 | 9.07E6 | 5.08E7 | 6.39E6 | 2.18E7 | 961606 | 16.86 |
| recoilMass (GeV) $\in (74, 131)$ | 157822 | 5.11E7 | 2.17E6 | 1.38E6 | 4.78E6 | 1.30E6 | 1.08E6 | 74991 | 4.99 |
| visEn (GeV) $\in (109, 143)$ | 142918 | 2.37E7 | 1.35E6 | 8.81E5 | 3.60E6 | 1.03E6 | 6.29E5 | 50989 | 3.92 |
| leadLepEn (GeV) $\in (0, 42)$ | 141926 | 2.08E7 | 3.65E5 | 7.24E5 | 2.81E6 | 9.72E5 | 1.34E5 | 46963 | 3.59 |
| multiplicity $\in (40, 130)$ | 139545 | 1.66E7 | 2.36E5 | 5.24E5 | 2.62E6 | 9.07E5 | 4977 | 42751 | 3.29 |
| leadNeuEn (GeV) $\in (0, 41)$ | 138653 | 1.46E7 | 2.24E5 | 4.72E5 | 2.49E6 | 8.69E5 | 4552 | 42303 | 3.12 |
| Pt (GeV) $\in (20, 60)$ | 121212 | 248715 | 1.56E5 | 2.48E5 | 1.51E6 | 4.31E5 | 999 | 35453 | 1.37 |
| Pl (GeV) $\in (0, 50)$ | 118109 | 52784 | 1.05E5 | 74936 | 7.30E5 | 1.13E5 | 847 | 34279 | 0.94 |
| -log10(Y23) $\in (3.375, +\infty)$ | 96156 | 40861 | 26088 | 60349 | 2.25E5 | 82560 | 640 | 10691 | 0.76 |
| InvMass (GeV) $\in (116, 134)$ | 71758 | 22200 | 11059 | 6308 | 77912 | 13680 | 248 | 6915 | 0.64 |
| BDT $\in (-0.02, 1)$ | 60887 | 9140 | 266 | 2521 | 3761 | 3916 | 58 | 1897 | 0.47 |

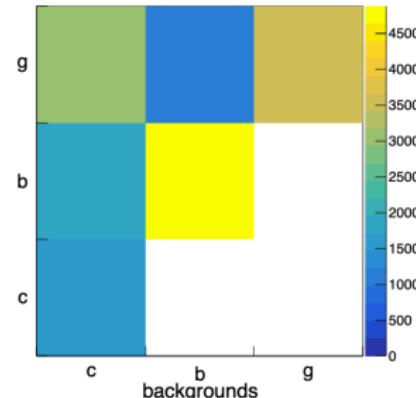
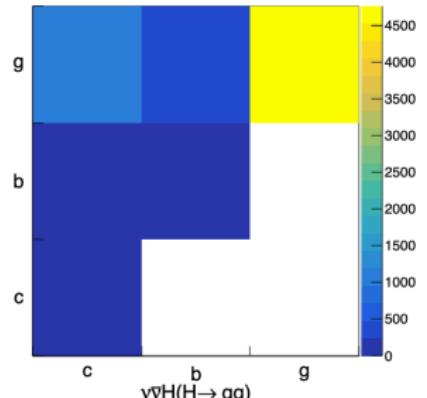
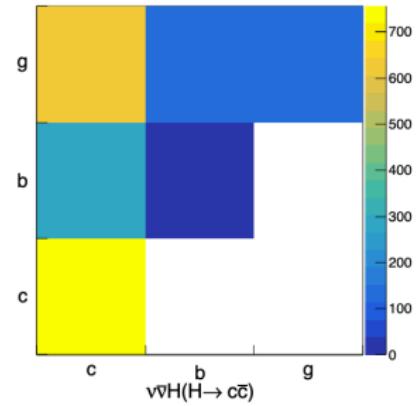
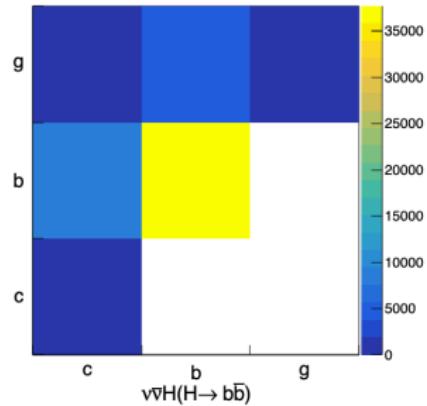
- ① The final state particles of remaining events are clustered into two jets.
- ② The b-likeness and c-likeness of two jets were calculated with the LCFIPlus package.
- ③ The cut on b-likeness and c-likeness can be found to maximize the value of $eff(b \rightarrow b) + eff(c \rightarrow c) + eff(udsg \rightarrow udsg)$.

migration matrix: the elements represent the identification efficiency and misidentification rate of each quark species.

| | b | c | g |
|---|--------|--------|--------|
| b | 0.8675 | 0.0887 | 0.0437 |
| c | 0.1136 | 0.6263 | 0.2601 |
| g | 0.0411 | 0.1007 | 0.8582 |

identified as

events distribution based on migration matrix :



get the signal strength accuracy by the Log-likelihood function:

$$-2 \cdot \log(\ell) = \sum_{i=1}^{i=6} \frac{[S_b \cdot N_{b,i} + S_c \cdot N_{c,i} + S_{light} \cdot N_{light,i} + N_{bkg,i} - N_i]^2}{N_i}$$

- S_b : the signal strength of $\nu\nu H b\bar{b}$
- $N_{b,i}$: the event number of $\nu\nu H b\bar{b}$ in i th bin
- N_i : the total event number in i 'th bin of $\nu\nu H b\bar{b}$, $\nu\nu H/c\bar{c}$, $\nu\nu H gg$ and backgrounds
- $N_{bkg,i}$ is the expected event number in i th bin of backgrounds,
- similar for S_c , S_{light} , $N_{c,i}$, and $N_{light,i}$

hessian matrix =
$$\begin{bmatrix} \frac{\partial^2 \log(\ell)}{\partial S_g \partial S_c} & \frac{\partial^2 \log(\ell)}{\partial S_g \partial S_b} & \frac{\partial^2 \log(\ell)}{\partial S_g \partial S_g} \\ \frac{\partial^2 \log(\ell)}{\partial S_b \partial S_c} & \frac{\partial^2 \log(\ell)}{\partial S_b \partial S_b} & \frac{\partial^2 \log(\ell)}{\partial S_b \partial S_g} \\ \frac{\partial^2 \log(\ell)}{\partial S_c \partial S_c} & \frac{\partial^2 \log(\ell)}{\partial S_c \partial S_b} & \frac{\partial^2 \log(\ell)}{\partial S_c \partial S_g} \end{bmatrix}$$

- The error covariance is obtained from the Hessian matrix.
- The relative accuracy of signal strength is the square roots of the diagonal elements of the covariance matrix, it is **0.49%/5.75%/1.82%** for $\nu\nu H b\bar{b}/c\bar{c}/gg$.

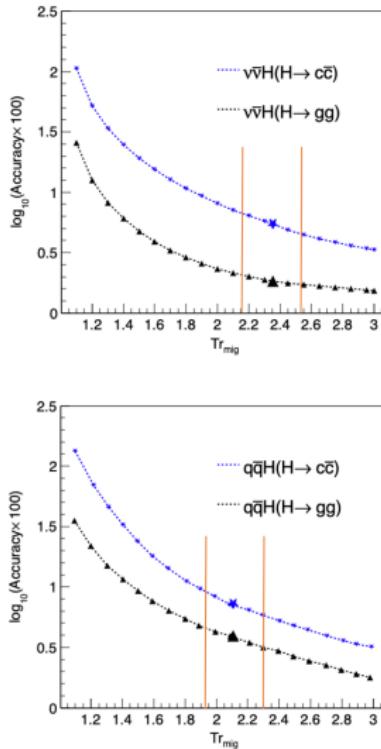
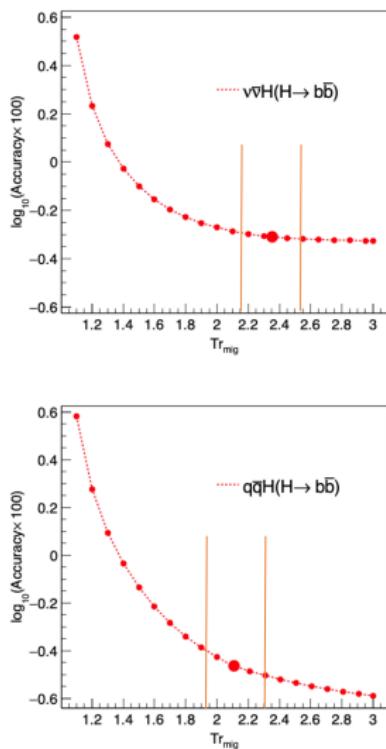
The similar method can be applied to $q\bar{q}Hb\bar{b}/c\bar{c}/gg$.

| | qqHqq | 2f | SW | SZ | WW | ZZ | Mixed | ZH | $\frac{\sqrt{s} + B}{\sqrt{s}} (\%)$ |
|---|--------|--------|--------|--------|--------|--------|--------|--------|--------------------------------------|
| total | 527488 | 8.01E8 | 1.95E7 | 9.07E6 | 5.08E7 | 6.39E6 | 2.18E7 | 613008 | 5.71 |
| multiplicity $\in (27, +\infty)$ | 527488 | 3.04E8 | 1.46E7 | 3.37E6 | 4.85E7 | 6.00E6 | 1.81E7 | 577930 | 3.77 |
| <i>leadLepEn</i> $\in (0, 59)$ | 527036 | 2.98E8 | 6.76E6 | 2.44E6 | 3.93E7 | 5.40E6 | 1.79E7 | 531411 | 3.65 |
| <i>visEn</i> $\in (199, 278)$ | 510731 | 1.21E8 | 1.29E6 | 551105 | 2.14E7 | 3.06E6 | 1.71E7 | 180571 | 2.52 |
| <i>leadNeuEn</i> $\in (0, 57)$ | 509623 | 5.68E7 | 716161 | 168030 | 2.04E7 | 2.93E6 | 1.65E7 | 176387 | 1.94 |
| <i>thrust</i> $\in (0, 0.86)$ | 460535 | 7.81E6 | 473732 | 132126 | 1.88E7 | 2.60E6 | 1.54E7 | 167863 | 1.47 |
| $-\log(Y_{34})$ $\in (0, 5.8875)$ | 451468 | 4.90E6 | 181432 | 119836 | 1.74E7 | 2.40E6 | 1.45E7 | 165961 | 1.40 |
| <i>HiggsJetsA</i> $\in (2.18, 2\pi)$ | 326207 | 2.83E6 | 110156 | 58613 | 4.54E6 | 870276 | 3.74E6 | 96560 | 1.08 |
| <i>ZJetsA</i> $\in (1.97, 2\pi)$ | 279030 | 1.37E6 | 33491 | 37101 | 2.39E6 | 496611 | 2.00E6 | 74005 | 0.93 |
| <i>ZHiggsA</i> $\in (2.32, 2\pi)$ | 274530 | 1.32E6 | 17026 | 33847 | 2.28E6 | 468340 | 1.91E6 | 69620 | 0.92 |
| <i>circle</i> BDT | 268271 | 1.20E6 | 10193 | 31567 | 2.13E6 | 424514 | 1.79E6 | 65434 | 0.90 |
| $\epsilon(0.02, 1)$ | 192278 | 378300 | 40 | 307 | 271436 | 141446 | 244126 | 30022 | 0.57 |

With log-likelihood function,

$-2 \cdot \log(\ell) = \sum_{i=1}^{i=36} \frac{[S_b \cdot N_{b,i} + S_c \cdot N_{c,i} + S_{light} \cdot N_{light,i} + N_{bkg,i} - N_i]^2}{N_i}$, the signal strength accuracy is **0.35%/7.74%/3.96%** for $q\bar{q}Hb\bar{b}/c\bar{c}/gg$.

key performance: flavor tagging



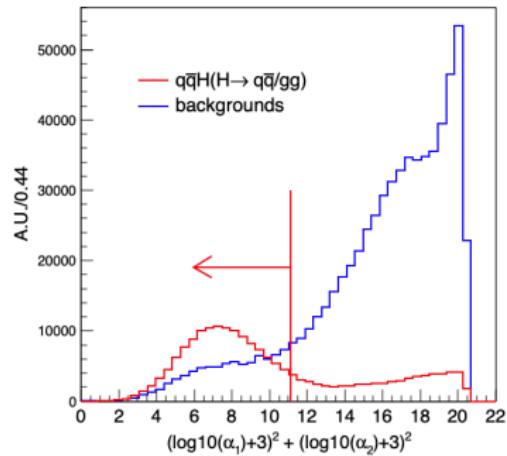
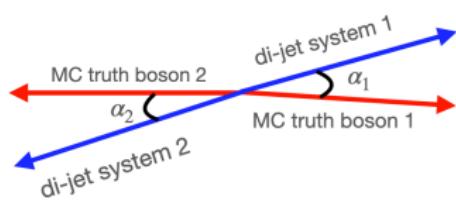
Tr_{mig} : the trace of migration matrix

The perfect flavor tagging performance can improve the $v\bar{v}H(H \rightarrow b\bar{b}/c\bar{c}/gg)$ signal strength accuracy by **2%/63%/13%**

and

$q\bar{q}H(H \rightarrow b\bar{b}/c\bar{c}/gg)$ signal strength accuracy by **35%/122%/181%**.

key performance: Color Singlet Identification (CSI)



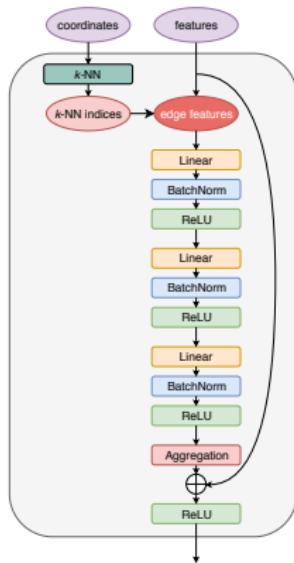
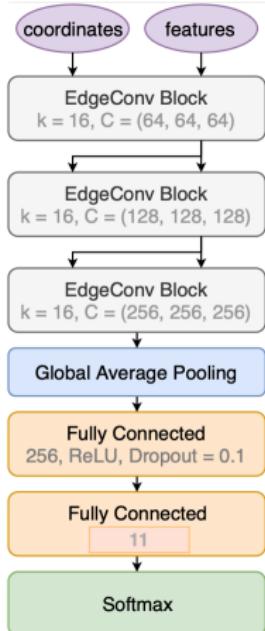
| | $qqHb\bar{b}$ | $qqHc\bar{c}$ | $qqHg\bar{g}$ |
|----------------|---------------|---------------|---------------|
| w.o. alpha cut | 0.35% | 7.74% | 3.96% |
| w.i. alpha cut | 0.33% | 4.37% | 2.08% |
| improved by | 6% | 77% | 90% |

note: The CSI in this report is just a demonstrator to illustrate the importance of CSI, we need to construct a CSI evaluator at the reconstruction level in the future.

The rare and FCNC hadronic decays of the Higgs.

ParticleNet (arXiv:1902.08570): Graph Neural Network based flavor tagging algorithm

Jet origin identification:



| | | b | \bar{b} | c | \bar{c} | s | \bar{s} | u | \bar{u} | d | \bar{d} | G |
|-----------|-----------|-------|-----------|-------|-----------|-------|-----------|-------|-----------|-------|-----------|-------|
| True | b | 0.745 | 0.163 | 0.033 | 0.025 | 0.004 | 0.003 | 0.002 | 0.003 | 0.002 | 0.002 | 0.017 |
| | \bar{b} | 0.170 | 0.737 | 0.026 | 0.033 | 0.003 | 0.004 | 0.003 | 0.002 | 0.002 | 0.003 | 0.018 |
| c | 0.015 | 0.014 | 0.743 | 0.055 | 0.036 | 0.031 | 0.025 | 0.009 | 0.009 | 0.009 | 0.018 | 0.043 |
| \bar{c} | 0.016 | 0.015 | 0.056 | 0.739 | 0.032 | 0.037 | 0.009 | 0.026 | 0.017 | 0.010 | 0.010 | 0.043 |
| s | 0.003 | 0.002 | 0.020 | 0.018 | 0.543 | 0.102 | 0.030 | 0.080 | 0.063 | 0.045 | 0.092 | |
| \bar{s} | 0.003 | 0.003 | 0.018 | 0.020 | 0.102 | 0.542 | 0.084 | 0.028 | 0.045 | 0.062 | 0.094 | |
| u | 0.002 | 0.003 | 0.020 | 0.011 | 0.044 | 0.131 | 0.367 | 0.055 | 0.080 | 0.174 | 0.111 | |
| \bar{u} | 0.003 | 0.003 | 0.011 | 0.019 | 0.132 | 0.043 | 0.062 | 0.356 | 0.178 | 0.081 | 0.111 | |
| d | 0.003 | 0.003 | 0.012 | 0.019 | 0.112 | 0.092 | 0.082 | 0.207 | 0.277 | 0.079 | 0.112 | |
| \bar{d} | 0.003 | 0.003 | 0.020 | 0.012 | 0.092 | 0.112 | 0.219 | 0.076 | 0.079 | 0.272 | 0.113 | |
| G | 0.015 | 0.014 | 0.024 | 0.024 | 0.052 | 0.052 | 0.043 | 0.041 | 0.034 | 0.034 | 0.667 | |

b/c/s/g tagging efficiency can be higher than 91%/80%/64%/67%

$$\sqrt{s} = 240 \text{ GeV}, \int L = 20ab^{-1}$$

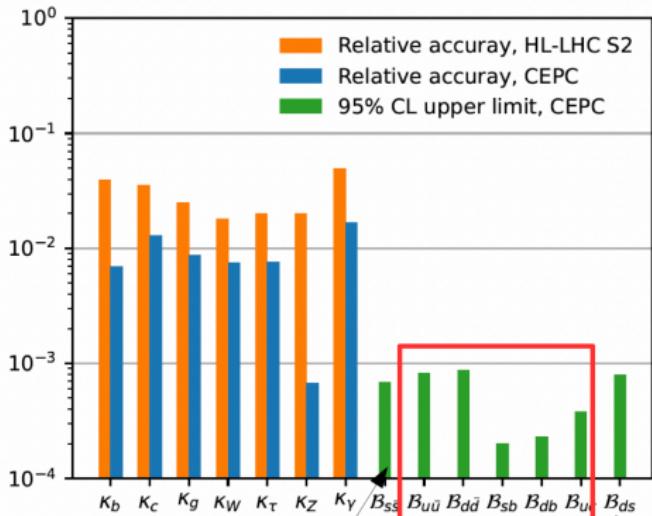


TABLE I: Summary of background events of $H \rightarrow b\bar{b}/c\bar{c}/gg$, Z , and W prior to flavor-based event selection, along with the expected upper limits on Higgs decay branching ratios at 95% CL. Expectations are derived based on the background-only hypothesis.

| | Bkg. (10^3) | | | | Upper limit (10^{-3}) | | | | | | |
|-----------------|-----------------|-----|-----|------|---------------------------|------------|------------|------|------|------|------|
| | | H | Z | W | $s\bar{s}$ | $u\bar{u}$ | $d\bar{d}$ | sb | db | uc | ds |
| $\nu\bar{\nu}H$ | 151 | 20 | 2.1 | 0.81 | 0.95 | 0.99 | 0.26 | 0.27 | 0.46 | 0.93 | |
| $\mu^+\mu^-H$ | 50 | 25 | 0 | 2.6 | 3.0 | 3.2 | 0.5 | 0.6 | 1.0 | 3.0 | |
| e^+e^-H | 26 | 16 | 0 | 4.1 | 4.6 | 4.8 | 0.7 | 0.9 | 1.6 | 4.3 | |
| Comb. | - | - | - | 0.75 | 0.91 | 0.95 | 0.22 | 0.23 | 0.39 | 0.86 | |

- [28] J. Duarte-Campos, G. Perez, M. Schlaffer, and A. Soffer. Probing the Higgs-strange-quark coupling at e^+e^- colliders using light-jet flavor tagging. *Phys. Rev. D*, 101(11):115005, 2020.
- [50] Alexander Albert et al. Strange quark as a probe for new physics in the Higgs sector. In *Snowmass 2021*, 3 2022.
- [59] J. de Blas et al. Higgs Boson Studies at Future Particle Colliders. *JHEP*, 01:139, 2020.
- [60] Jorge De Blas, Gauthier Durieux, Christophe Grojean, Jiayin Gu, and Ayan Paul. On the future of Higgs, electroweak and diboson measurements at lepton colliders. *JHEP*, 12:117, 2019.

Summary :

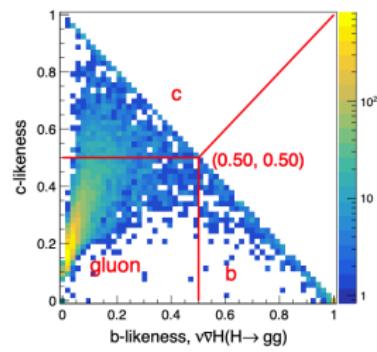
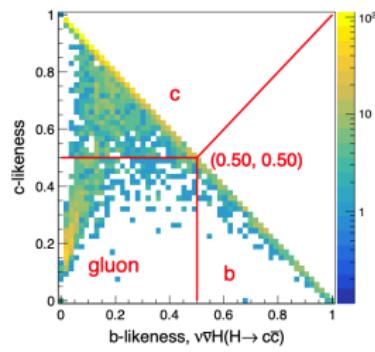
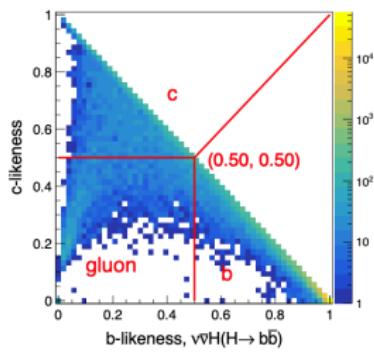
- Combine all four different channels of $\mu\mu H$, eeH (from Chin. Phys. C, 44(1):013001, 2020), $\nu\nu H$ and qqH , the total signal strength of $H \rightarrow b\bar{b}, c\bar{c}, gg$ can be measured to a relative accuracy of 0.27%/4.03%/1.56%.
JHEP, 11:100, 2022

| Z decay mode | $H \rightarrow b\bar{b}$ | $H \rightarrow c\bar{c}$ | $H \rightarrow gg$ |
|------------------------------|--------------------------|--------------------------|--------------------|
| $Z \rightarrow e^+e^-$ | 1.57% | 14.43% | 10.31% |
| $Z \rightarrow \mu^+\mu^-$ | 1.06% | 10.16% | 5.23% |
| $Z \rightarrow q\bar{q}$ | 0.35% | 7.74% | 3.96% |
| $Z \rightarrow \nu\bar{\nu}$ | 0.49% | 5.35% | 1.77% |
| combination | 0.27% | 4.03% | 1.56% |

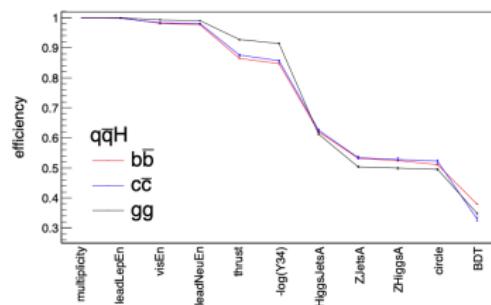
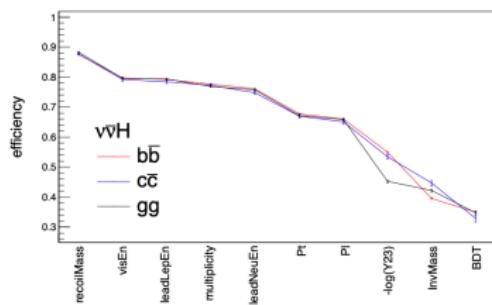
- The flavor tagging and color singlet identification (CSI) are the critical performances for these benchmarks.
- With jet origin identification, the branching ratio upper limit for $H \rightarrow s\bar{s}$ ($H \rightarrow u\bar{u}/d\bar{d}$) improved by a factor of 3 (~ 10) compared to previous studies.

Many thanks !

Backup



systematic uncertainty



systematic uncertainties are categorized into three groups:

- The first group are those that are significantly smaller than the statistical uncertainties, including the reconstructed energy/momentum scale of the physics objects.
- The second group are those comparable to the statistical uncertainty, especially the integrated luminosity.
- The third group are those that can be significantly larger than the statistical uncertainty, including CSI and the jet configuration.

$$M_{mig} = \frac{Tr_{mig} - Tr_{opt}}{Tr_I - Tr_{opt}} \cdot (M_I - M_{opt}) + M_{opt}$$

$$M_{mig} = \frac{Tr_{mig} - Tr_{opt}}{Tr_{1/3} - Tr_{opt}} \cdot (M_{1/3} - M_{opt}) + M_{opt}$$

| | b | c | g |
|---|---|---|---|
| b | 1 | 0 | 0 |
| c | 0 | 1 | 0 |
| g | 0 | 0 | 1 |

identified as
perfect

| | b | c | g |
|------|-----|-----|-----|
| true | 1/3 | 1/3 | 1/3 |
| c | 1/3 | 1/3 | 1/3 |
| g | 1/3 | 1/3 | 1/3 |

identified as
none