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- Primordial black holes (PBHs)
 - 1 PBHs as Dark Matter
 - 2 Hawking Radiation
 - **3** Constraints on PBHs by $\bar{\nu}_e$ Fluxes Upper limits
 - **4** The Sensitivity Floor for Primordial Black Holes Neutrino

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Evidence for Dark Matter

There is compelling astrophysical and cosmological evidence supporting the existence of dark matter (DM).

- Rotation Curve
- Cosmic Microwave Background
- Collisions of Galaxy Clusters
- Large Structure Formation



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Experimental Motivation

Experimental Anomalies

• The discrepancies between observation and theoretical prediction could be the hint for new physics

Null Observations

• Due to the increasing sensitivity, we can improve the measurement limits and exclude a significant region of the model parameter space

Theoretical Motivation

The SM cannot be the complete theory of the fundamental constituents of the universe.

- The SM does not contain gravity
- The neutrinos have mass
- The dark matter and dark energy
- Matter-antimatter asymmetry
- Strong CP problem



Primordial Black Hole Dark Matter

Dark matter is non-baryonic, stable, cold, and weakly interacting.

- PBHs form due to the gravitational collapse of overdensities
- Lifetime of PBHs can be longer than the age of the Universe
- PBHs can be non-relativistic
- PBHs are expected to only interact gravitationally

Primordial Black Hole Dark Matter

- PBHs could have formed with a broad range of masses in the early universe at different time after the big bang.
- The lifetime of a black hole is approximately given by,

$$t_{\rm evap} \ = t_{\rm Universe} \ \left(\frac{M_{\rm PBH}}{5 \times 10^{14} g} \right)^3, \label{eq:tevap}$$

where $M_{\rm PBH}$ is the PBH initial mass, and $t_{\rm Universe} \approx 13.8$ Gyr is the age of the Universe. We can infer that PBHs weighing less than 5×10^{14} g would have evaporated by now.

Primordial Black Hole Dark Matter

If the PBHs exist today, it would be a part of the DM, and we represent the PBHs abundance, a fraction of the DM denstity, as $f_{\rm PBH} \equiv \frac{\Omega_{\rm PBH}}{\Omega_{\rm DM}}$.



[https://github.com/bradkav/PBHbounds.]

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Hawking Radiation

• The curvature of spacetime near the event horizon of a black hole causes the vacuum state of the quantum fields to become excited, leading to the emission of particles.

$$\left. \frac{d^2 N_i(E,t)}{dt dE} \right|_{\text{primary}} = \frac{g_i}{2\pi} \frac{\Gamma_i(E, M_{\text{PBH}})}{e^{E/T_{\text{BH}}} \pm 1},$$

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• The differential neutrino flux can be separated into the galactic and extragalactic contribution

$$\frac{d\Phi_{\rm MW}}{dE_{\nu}} = \frac{d^2 N_{\nu}}{dE_{\nu} dt} \frac{f_{\rm PBH}}{M_{\rm PBH}} \int \frac{d\Omega}{4\pi} \int_0^{\ell_{\rm max}} d\ell \rho_{\rm NFW}(r(\ell,\phi)),$$
$$\frac{d\Phi_{\rm EG}\left(E_o\right)}{dE_o} = \frac{cf_{\rm PBH}\rho_{\rm DM}}{M_{\rm PBH}} \int_{z_{\rm min}}^{z_{\rm max}} \frac{dz}{H} \frac{d^2 N_{\nu}((1+z)E_o,t)}{dE_{\nu} dt}.$$

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Hawking Radiation



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Figure: Three dominant antineutrino backgrounds, and galactic (dotted), extragalactic (dashed) and total (solid) $\bar{\nu}_e$ fluxes from PBHs evaporation as a function of energy with $f_{\rm PBH} = 1$, assuming a monochromatic mass distribution.

$\bar{\nu}_e$ Fluxes Upper Limits

The primary electron antineutrino detection channel in most experiments is inverse-beta decay, $\bar{\nu}_e + p \rightarrow n + e^+$.



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Upper Limit on Dark Matter Fraction of PBHs $f_{\rm PBH}$



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Since these various experiments are independent, all the $\bar{\nu}_e$ fluxes upper limits can be taken into account to obtain a global constraint: $\exists \bar{\nu} \equiv \neg \land \land$ Qishan Liu WIN 2023 July 5, 2023 12/18



arXiv:1307.5458

• In the DM direct detection experiment, the recoil signal would be limited by coherent elastic neutrino-nucleus scattering

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arXiv:1307.5458

- In the DM direct detection experiment, the recoil signal would be limited by coherent elastic neutrino-nucleus scattering
- The DSNB serves as an irreducible background that forms a sensitivity floor in PBHs parameter space

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• As neutrino experiments become more sensitive, DSNB would become a background for PBHs neutrino searches.

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- Additionally, their energy distributions are different, but the flux of DSNB is uncertain.

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- The DSNB has a different angular distribution from the signal produced by PBHs but the detector cannot distinguish between them.
- Additionally, their energy distributions are different, but the flux of DSNB is uncertain.
- Both of reactor neutrinos and atmospheric neutrinos can be measured accurately, so we do not take them into account by assuming they can be eliminated by other methods.

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• To obtain the sensitivity floor, we estimate the number of expected events for PBHs and DSNB in the Hyper-Kamiokande,

$$N_i^{\text{PBH(DSNB)}} = \varepsilon N_t t \int F_i^{\text{PBH(DSNB)}} \sigma_i(E) dE,$$

where $\varepsilon = 67\%$ is the detector efficiency, $N_t = 2.5 \times 10^{34}$ is the number of targets, $\sigma_i(E)$ is the inverse beta decay cross-section.

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• We use the log-likelihood ratio,

$$\chi^2 = -2\ln\frac{L(f_{\rm PBH})}{L(f_{\rm PBH}=0)},$$

where L are given by Poisson likelihood.



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Figure: Current and prospective upper bounds on PBHs, and the sensitivity floor for PBHs neutrino.

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- The current data excludes PBHs as the sole component of dark matter up to masses of $6.4 \times 10^{15} \, {\rm g}$. This represents a significant improvement of approximately 20% in comparison to the previous upper limits obtained from SK's data, which was around $5.2 \times 10^{15} \, {\rm g}$.
- In addition, since the DSNB is an unavoidable isotropic background, we thus estimate the sensitivity floor and show that it is difficult for neutrino detectors to detect the PBH DM above 9×10^{15} g.

Thank you for your attention!

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