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arXiv: 2302.10965, 2211.02028  
in collaboration with  
Vedran Brdar, Andr e de Gouv ea  
and Pedro A. N. Machado

# Neutrino magnetic moment portal and supernova: new constraints and multi-messenger opportunities

July 7, 2023 WIN2023

# Heavy Neutral Lepton - Sterile Neutrino

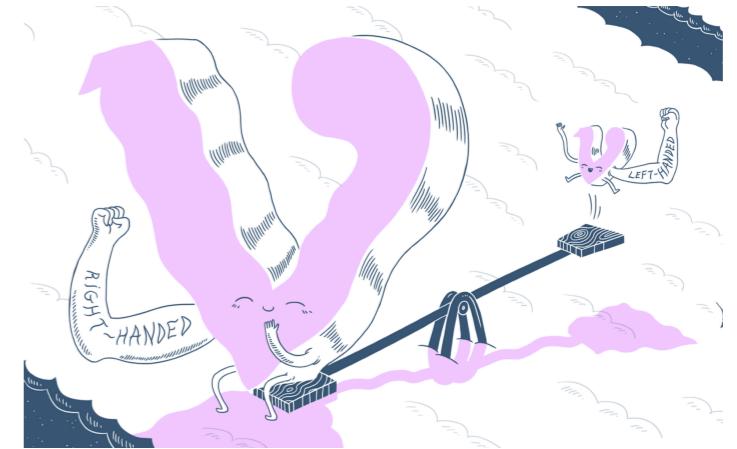
[Minkowski, Mohapatra, Senjanović, Gell-Mann, Ramond, Slansky, Yanagida]

$$\mathcal{L} \supset \frac{1}{2} \overline{N^c} M_R N + \bar{L} Y_\nu \tilde{H} N + \text{h.c.}$$

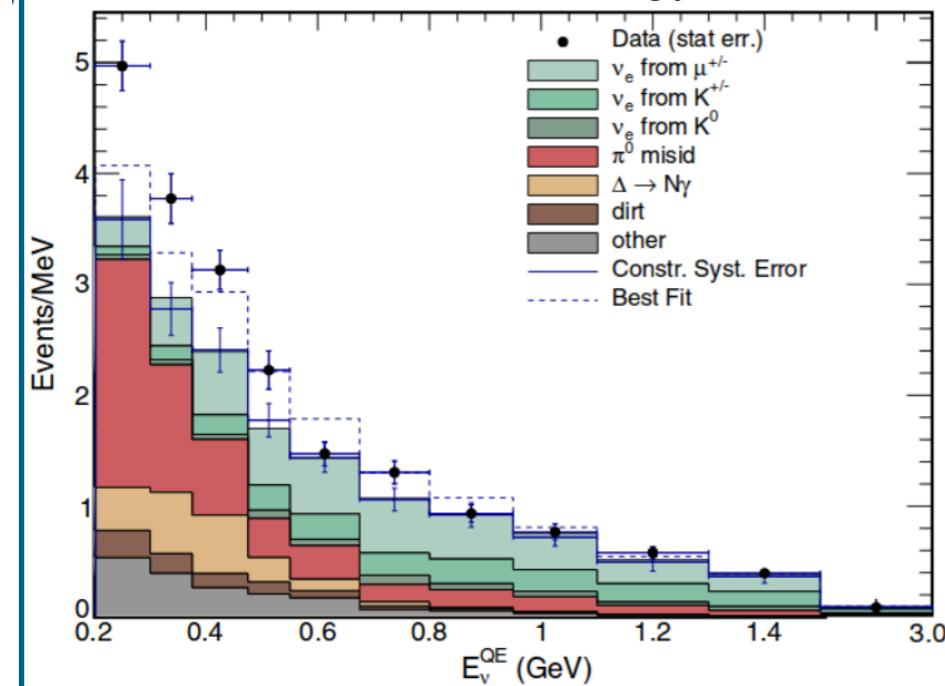
Neutrino mass:  
type-I seesaw  
mechanism

$$\mathcal{M} = \begin{pmatrix} 0 & M_D \\ M_D & M_R \end{pmatrix}$$

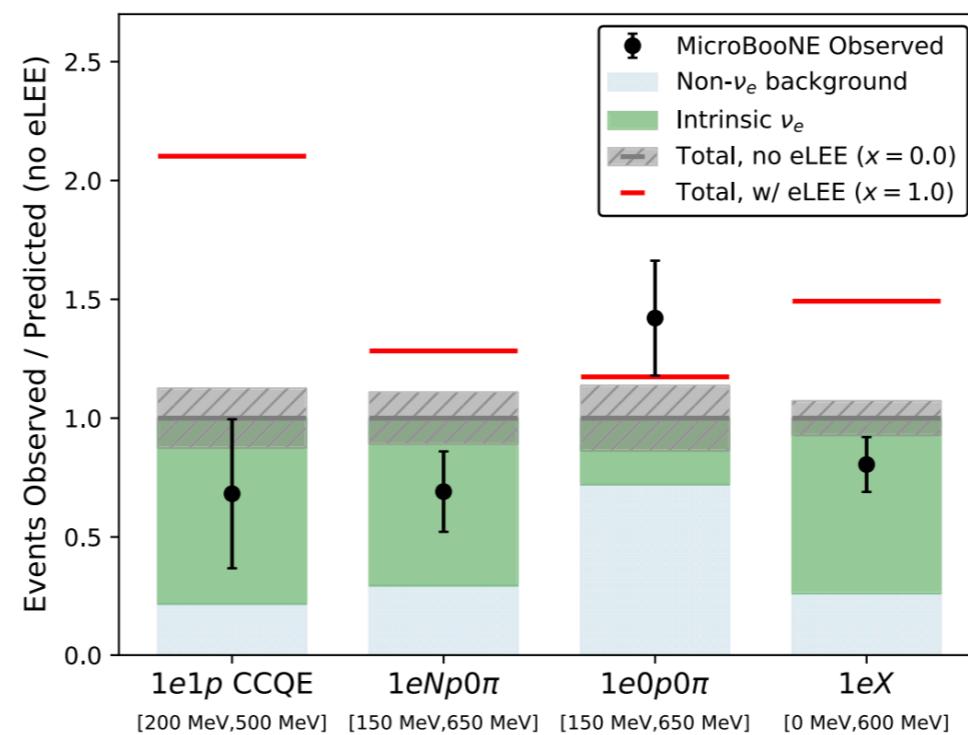
$$m_\nu = -M_D M_R^{-1} M_D^T$$



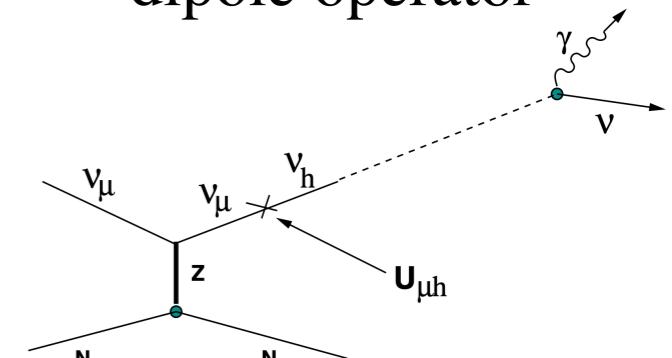
MinibooNE low energy excess



[PRL 128, 241801 (2022)]



Via mixing and  
dipole operator



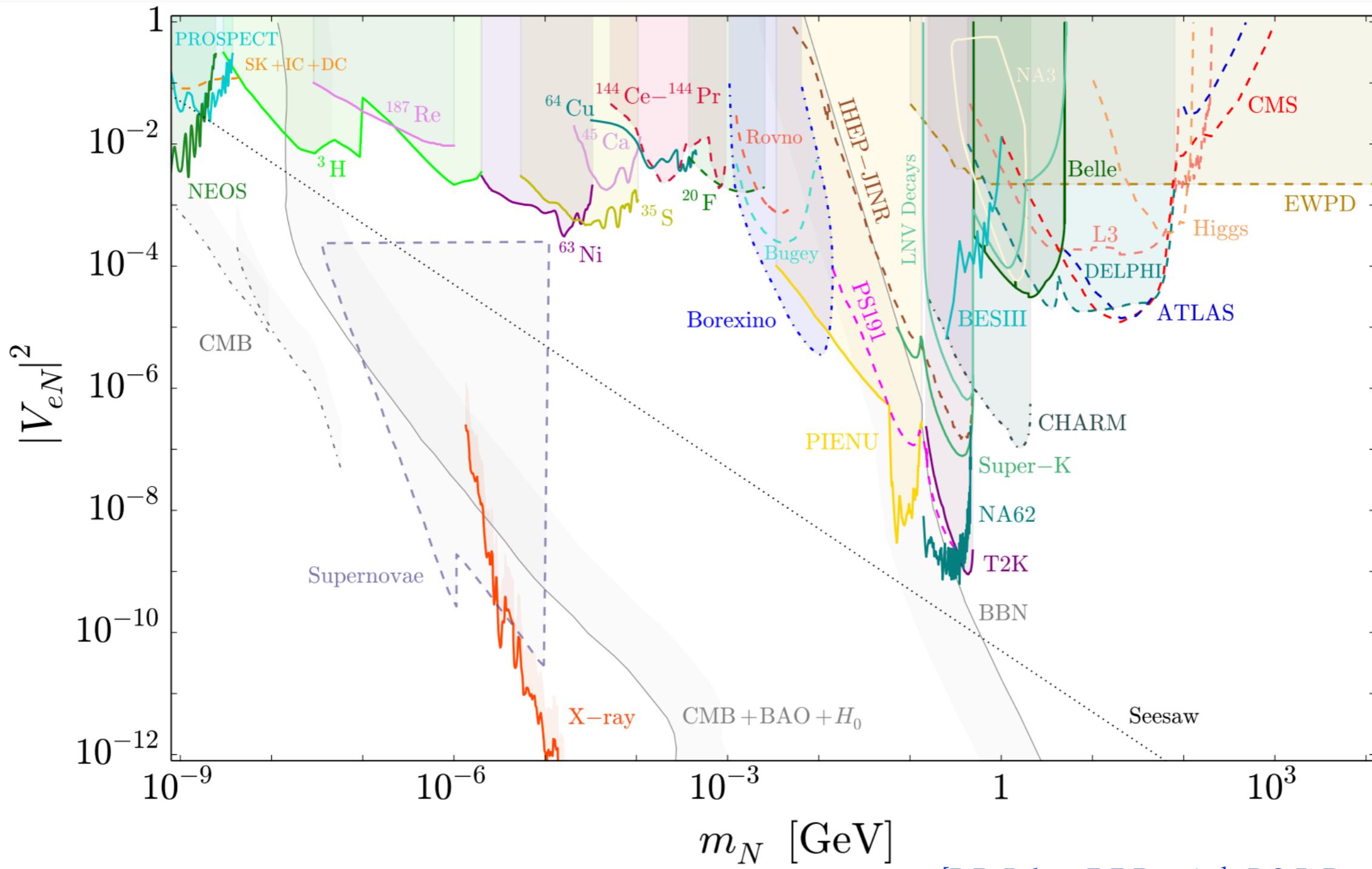
$$\theta^2 \sim 10^{-3}$$

$$\mu \sim 10^{-9} \mu_B$$

[S. N. Gninenko, PRL 103, 241802 (2009)]

# Heavy Neutral Lepton - Mixing Portal

$$\Gamma_{\nu_s \rightarrow \nu_a \gamma} = 1.38 \cdot 10^{-29} \text{ sec}^{-1} \left( \frac{\sin^2 2\theta}{10^{-7}} \right) \left( \frac{m_s}{1 \text{ keV}} \right)^5.$$



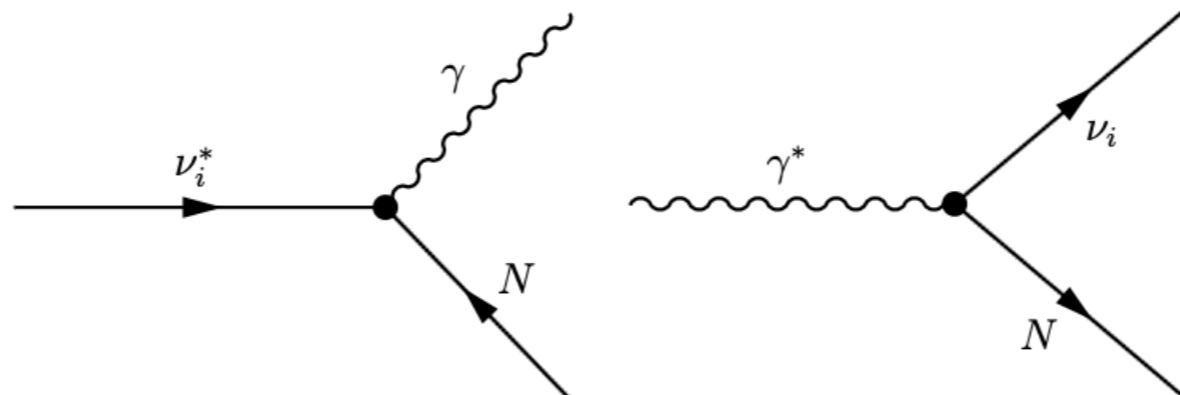
[P. D. Bolton, F. F. Deppisch, P. S. B. Dev, arXiv: 1912.03058]

# Heavy Neutral Lepton - Dipole Portal

When the mixing effect is subdominant, ...

$$\mathcal{L} \supset \frac{1}{2} \mu_\nu \bar{\nu}_L^\alpha \sigma^{\mu\nu} N F_{\mu\nu}$$

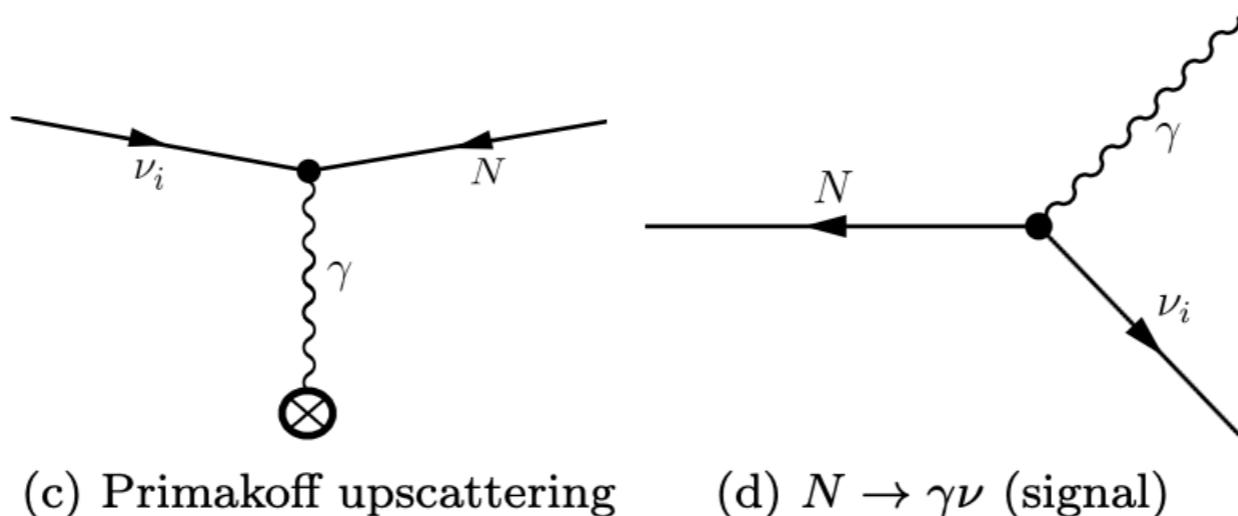
$$\pi^\pm, K^\pm \rightarrow \mu^\pm \left( \overset{(-)}{\nu_\mu}{}^* \rightarrow \gamma \overset{(-)}{N} \right) \quad \pi^0, \eta \rightarrow \gamma (\gamma^* \rightarrow \nu_a N)$$



(a) Weak meson decays

(b) Dalitz-like decay

$$\text{Br}(M \rightarrow N) \propto \mu_\nu^2 m_M^2$$



(c) Primakoff upscattering

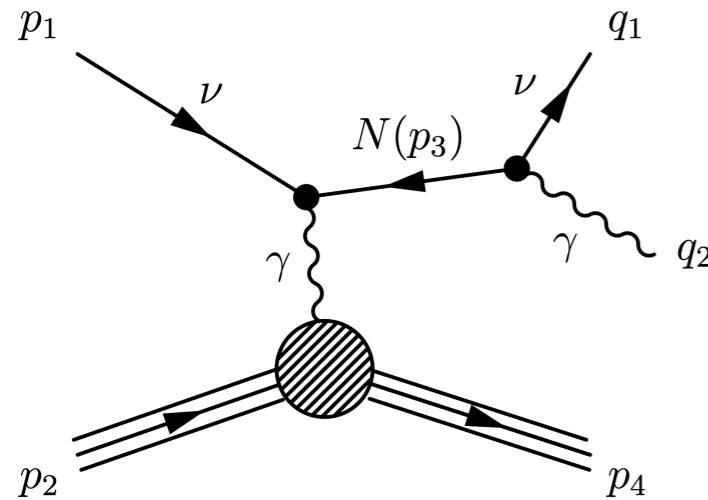
(d)  $N \rightarrow \gamma\nu$  (signal)

$$\Gamma_N = \frac{6}{4\pi} \mu_\nu^2 M_N^3$$

(flavor universal)

# Current Probes: terrestrial experiments

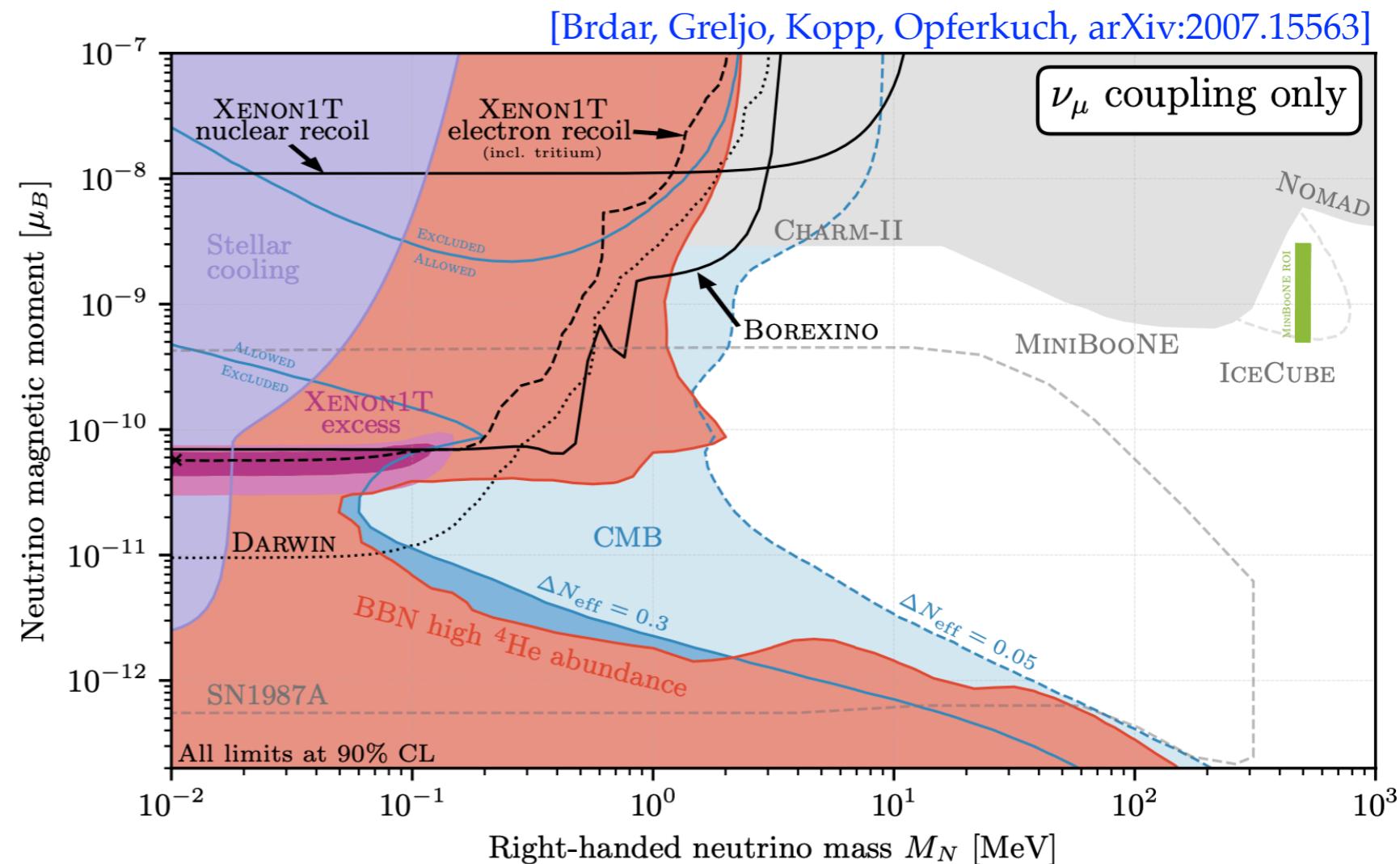
Beam dump experiments:  
MiniBoone, NONAD



$$\pi^0, \eta \rightarrow \gamma (\gamma^* \rightarrow \nu_a N)$$

$$\pi^\pm, K^\pm \rightarrow \mu^\pm \left( \bar{\nu}_\mu^* \rightarrow \gamma N \right)$$

relevant for transition magnetic  
moments between  $\nu_\mu$  and  $N$

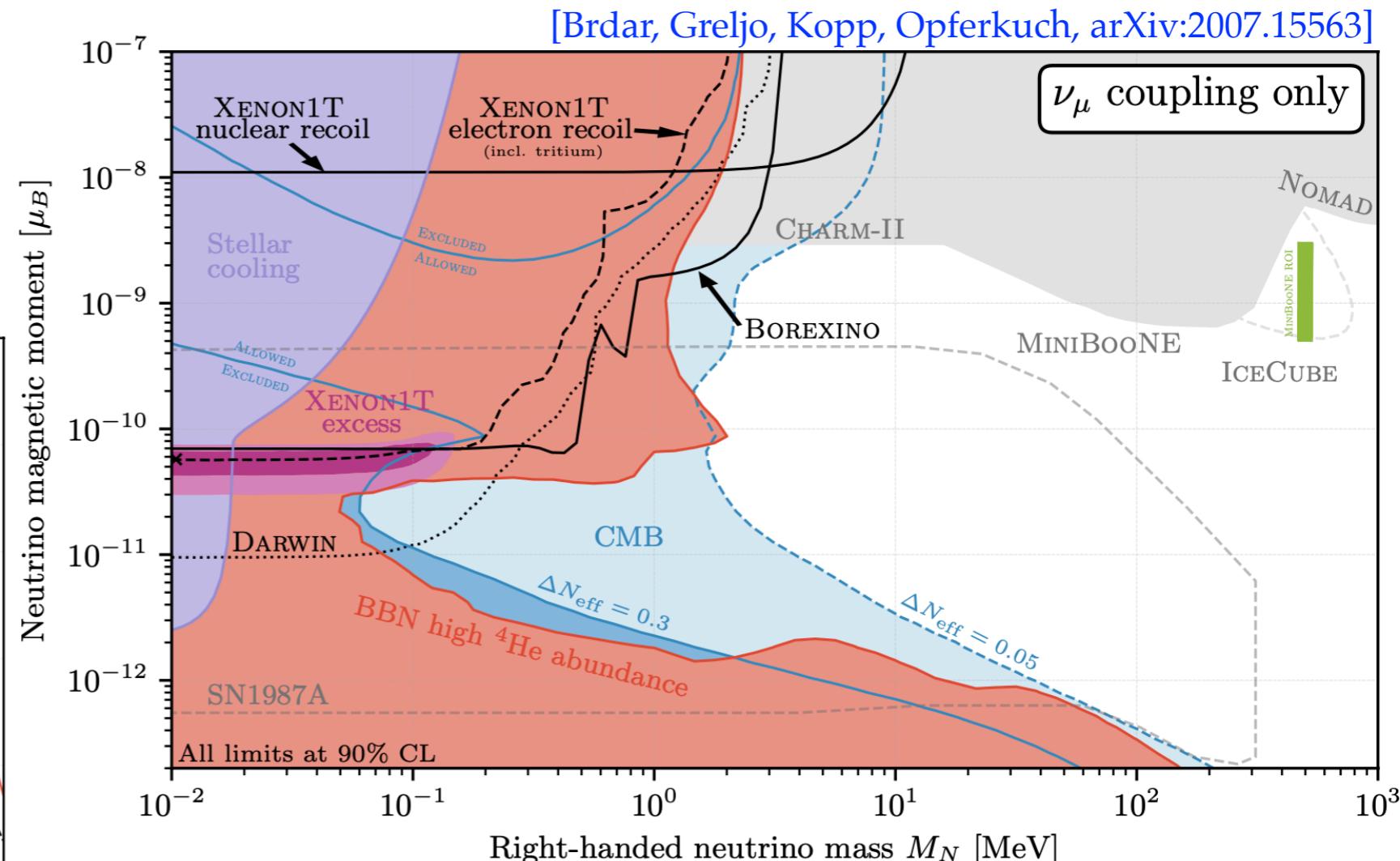
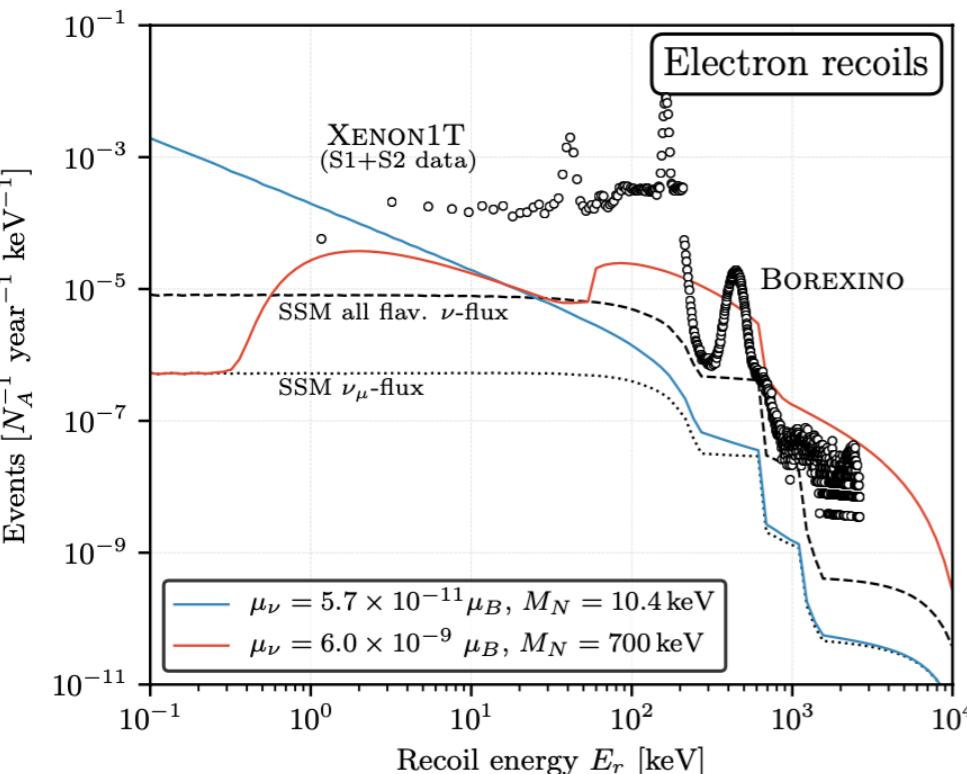


# Current Probes: terrestrial experiments

Solar neutrino spectrum:  
Xenon1T, Borexino

$$\nu_L + e^- \rightarrow N + e^-$$

$$\nu_L + X_Z^A \rightarrow N + X_Z^A$$



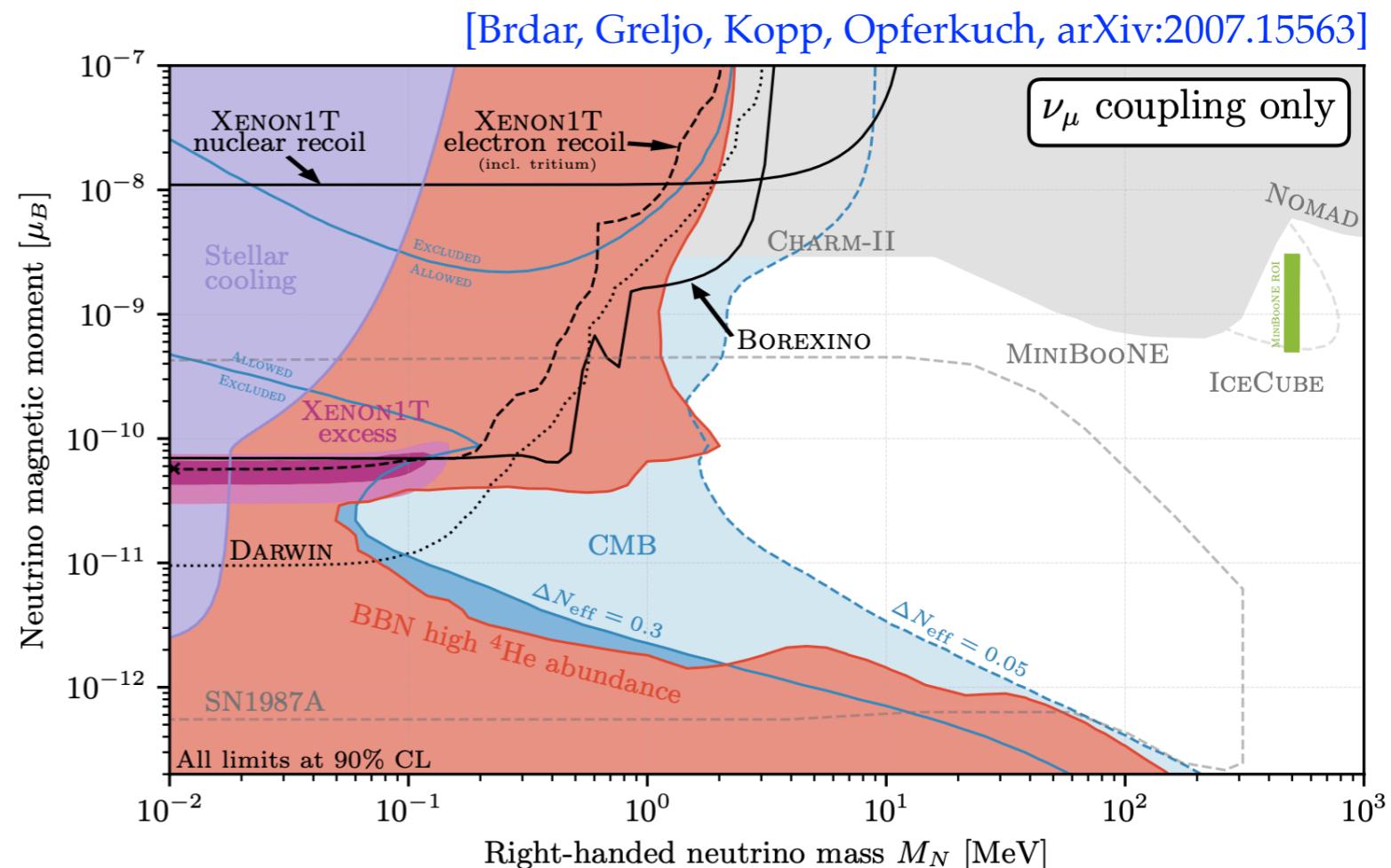
$M_N \gtrsim 6 \text{ GeV}$  (LHC, LEP):  $e^+ e^- (q\bar{q}) \rightarrow (N \rightarrow \gamma\nu)\bar{\nu} + h.c.$

# Current Probes: cosmology

CMB, BBN:  $N_{\text{eff}}$

- Relativistic:  $N_{\text{eff}}$
- Inject extra photons  $N \rightarrow \nu\gamma$

$$\begin{aligned}\tau_N &= \frac{16\pi}{\mu_\nu^2 M_N^3} \\ &= 3760 \text{ sec} \times \left( \frac{1 \times 10^{-11} \mu_B}{\mu_\nu} \right)^2 \left( \frac{\text{MeV}}{M_N} \right)^3\end{aligned}$$

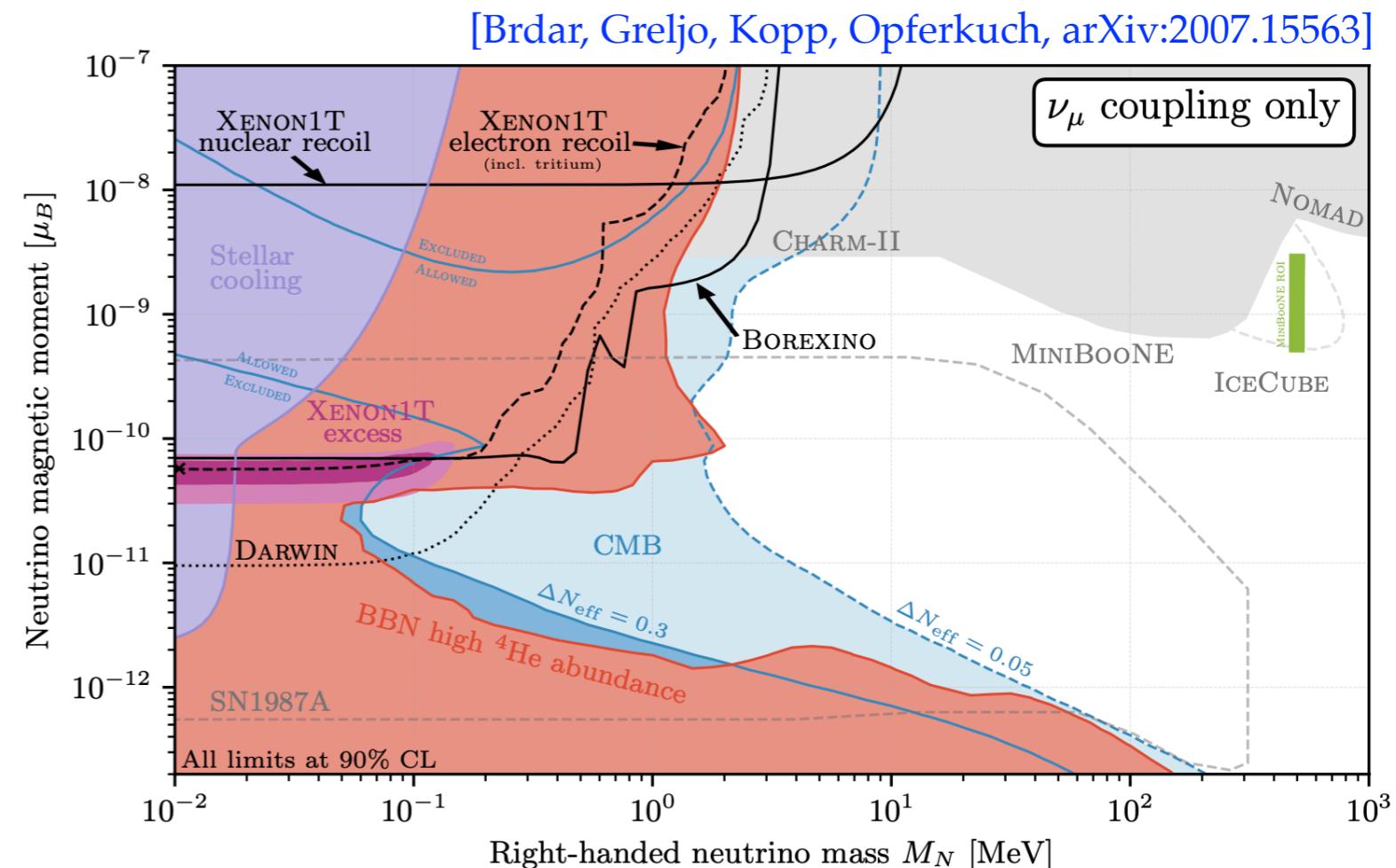


$$T_{\text{dec}} \sim 1.28 \text{ GeV} \left( \frac{10^{-11} \mu_B}{\mu_\nu} \right)^2$$

# Current Probes: Supernova

10% of energy loss  
to sterile neutrino

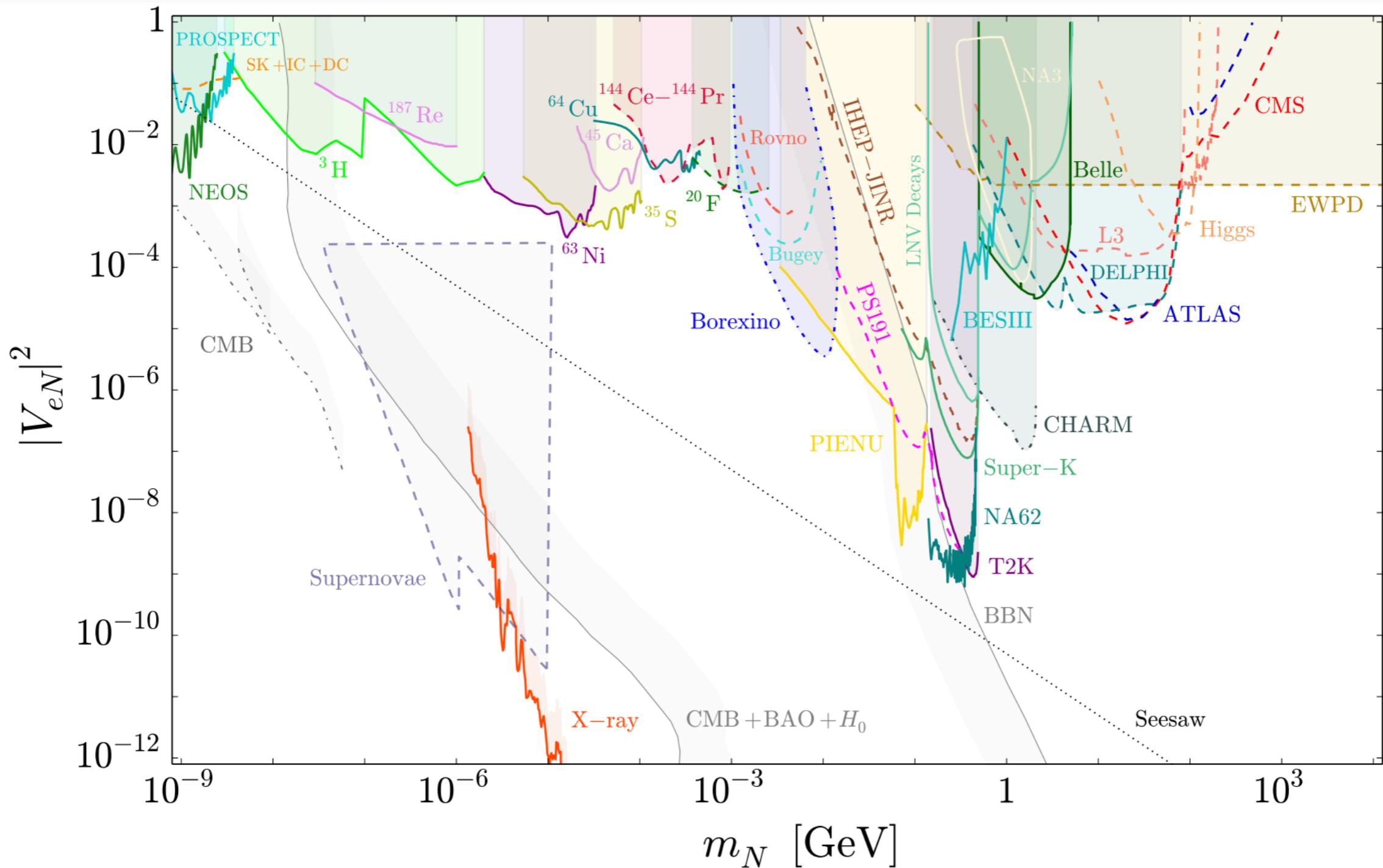
$$\begin{aligned}\nu + e^\pm &\rightarrow N + e^\pm \\ \nu + p &\rightarrow N + p \\ e^+ + e^- &\rightarrow \bar{\nu} + N \\ \gamma + \nu &\rightarrow N\end{aligned}$$



# UV completion

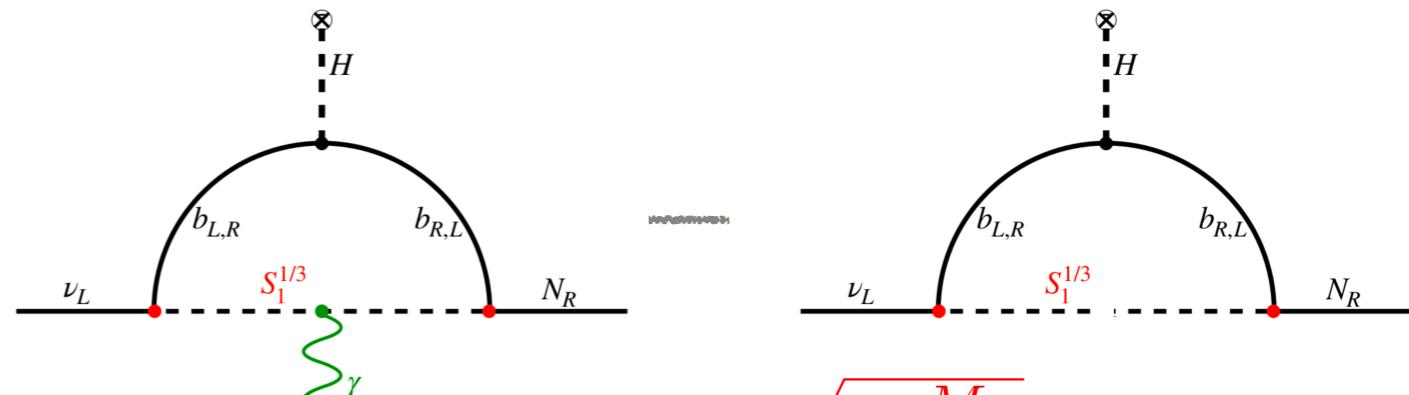
In what theory the effect of mixing is subdominant...

$$\theta^2 \sim \frac{m_\nu}{M_N}, \text{ keV} < M_N < 100 \text{ MeV}$$



# UV completion

- Voloshin-type symmetry  $SU(2)_\nu$



$$\mu_\nu \sim 10^{-12} \mu_B \frac{\sqrt{m_\nu M_N}}{\text{MeV}}$$

$$(N_R^C, \nu_L)^T \in \mathbf{2}$$

$$SU(2)_L \rightarrow SU(3)_L$$

$$\bar{\nu}_L N_R \rightarrow -\bar{\nu}_L \nu_R$$

$$\bar{\nu}_L \sigma_{\mu\nu} N_R F^{\mu\nu} \rightarrow \bar{\nu}_L \sigma_{\mu\nu} N_R F^{\mu\nu}$$

$$m_{\nu N} \sim \frac{\mu_\nu}{\mu_B} \frac{\alpha}{4\pi} \frac{m_V^2}{2m_e}$$

$$\mu_\nu \sim 10^{-8} \mu_B \frac{\sqrt{m_\nu M_N}}{\text{MeV}}$$

$SU(2)_\nu$  Symmetry-breaking scale  $m_V$  at the TeV scale

$$m_\nu \sim 0.1 \text{eV}, \mu_\nu \sim 10^{-11} \mu_B$$

# Multimessenger Signals

$$\nu + e^\pm \rightarrow N + e^\pm$$

$$\nu + p \rightarrow N + p$$



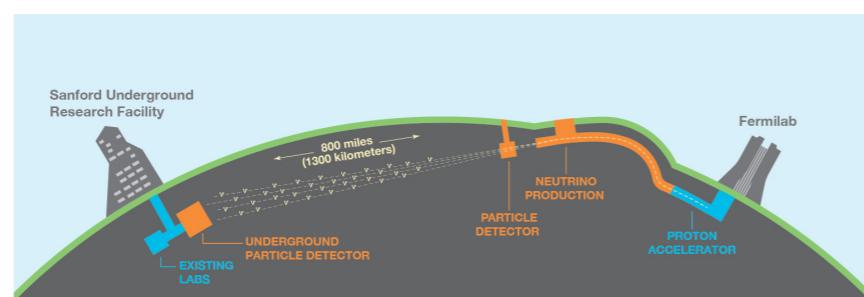
$$\mathcal{L} \supset \frac{1}{2} \mu_\nu \bar{\nu}_L^\alpha \sigma^{\mu\nu} N F_{\mu\nu}$$

$$e^+ + e^- \rightarrow \bar{\nu} + N$$

$$\gamma + \nu \rightarrow N$$



$$N \rightarrow \nu + \gamma$$



[V. Brdar, A. D. Gouv^ea, YYL, P. A. N. Machado, arXiv:2302.10965]

# Multimessenger Signals

$$\nu + e^\pm \rightarrow N + e^\pm$$

$$\nu + p \rightarrow N + p$$

$$e^+ + e^- \rightarrow \bar{\nu} + N$$

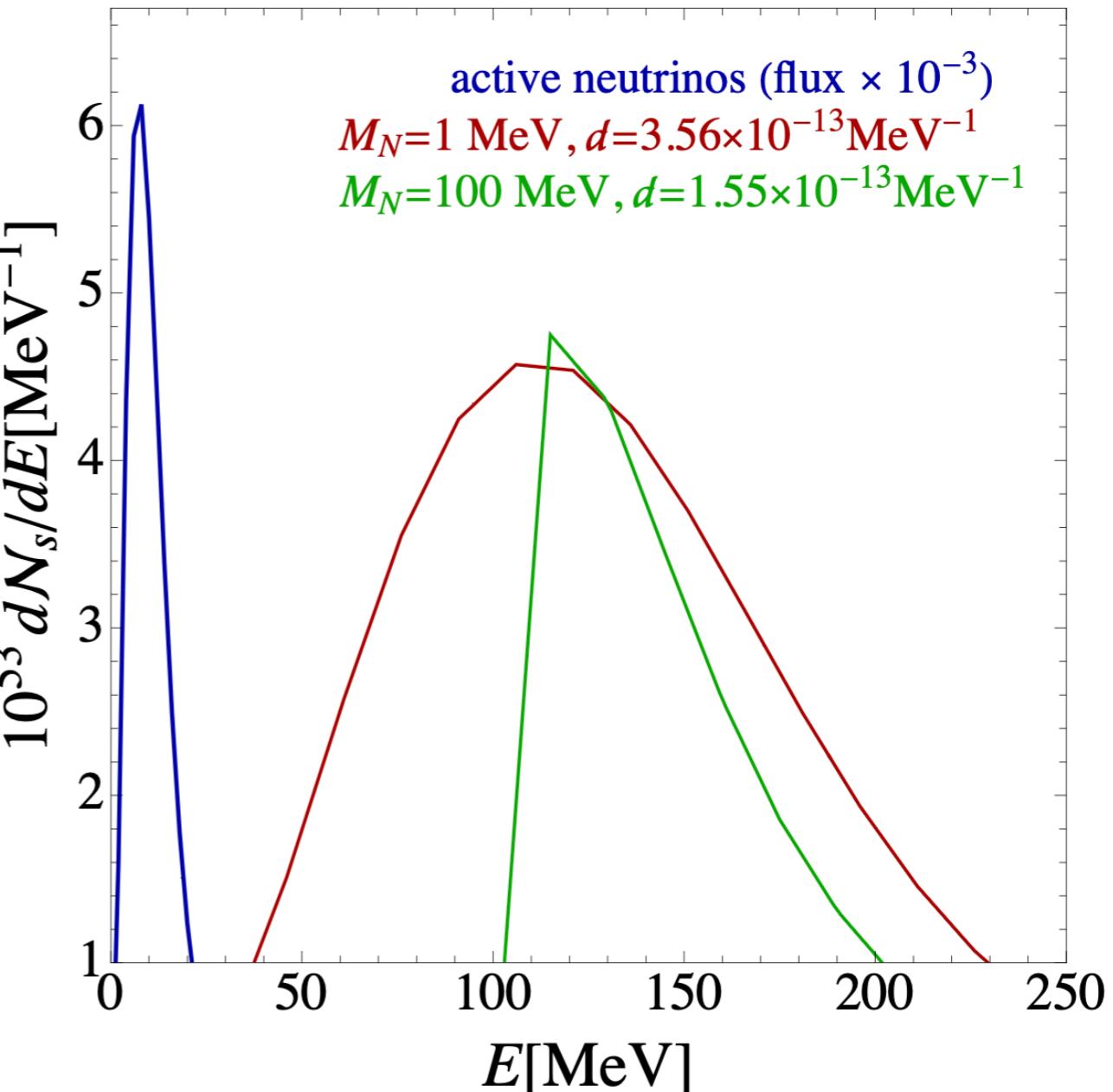
$$\gamma + \nu \rightarrow N$$

$$\frac{1}{4\pi r^2} \frac{\partial^2}{\partial r \partial t} \left( \frac{d\mathcal{N}_s}{dE_N} \right) = \sigma n_e \frac{dn_\nu}{dE}$$

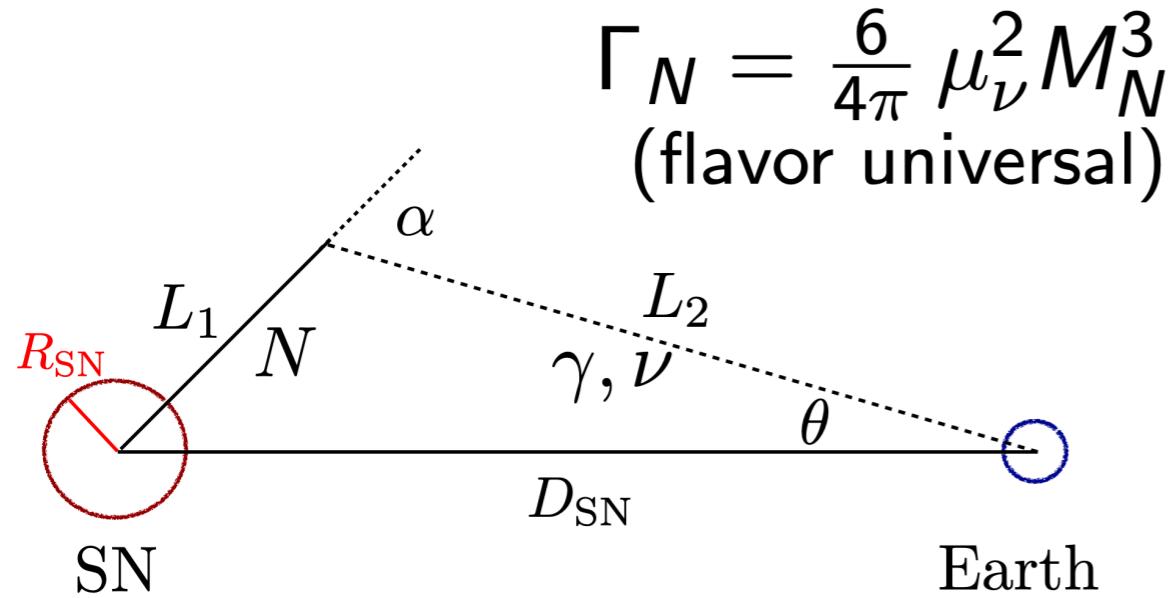
using data associated to  
the simulation performed  
by the Garching group of  
an  $8.8M_\odot$  progenitor star

sterile neutrinos are produced in a dense  $T \sim 100$  MeV core and leave  
subsequently the star without further interactions

[V. Brdar, A. D. Gouv<sup>ea</sup>, YYL, P. A. N. Machado, arXiv:2302.10965]



# Multimessenger Signals



$$\Delta t = L_1/\beta + L_2 - D_{\text{SN}}$$

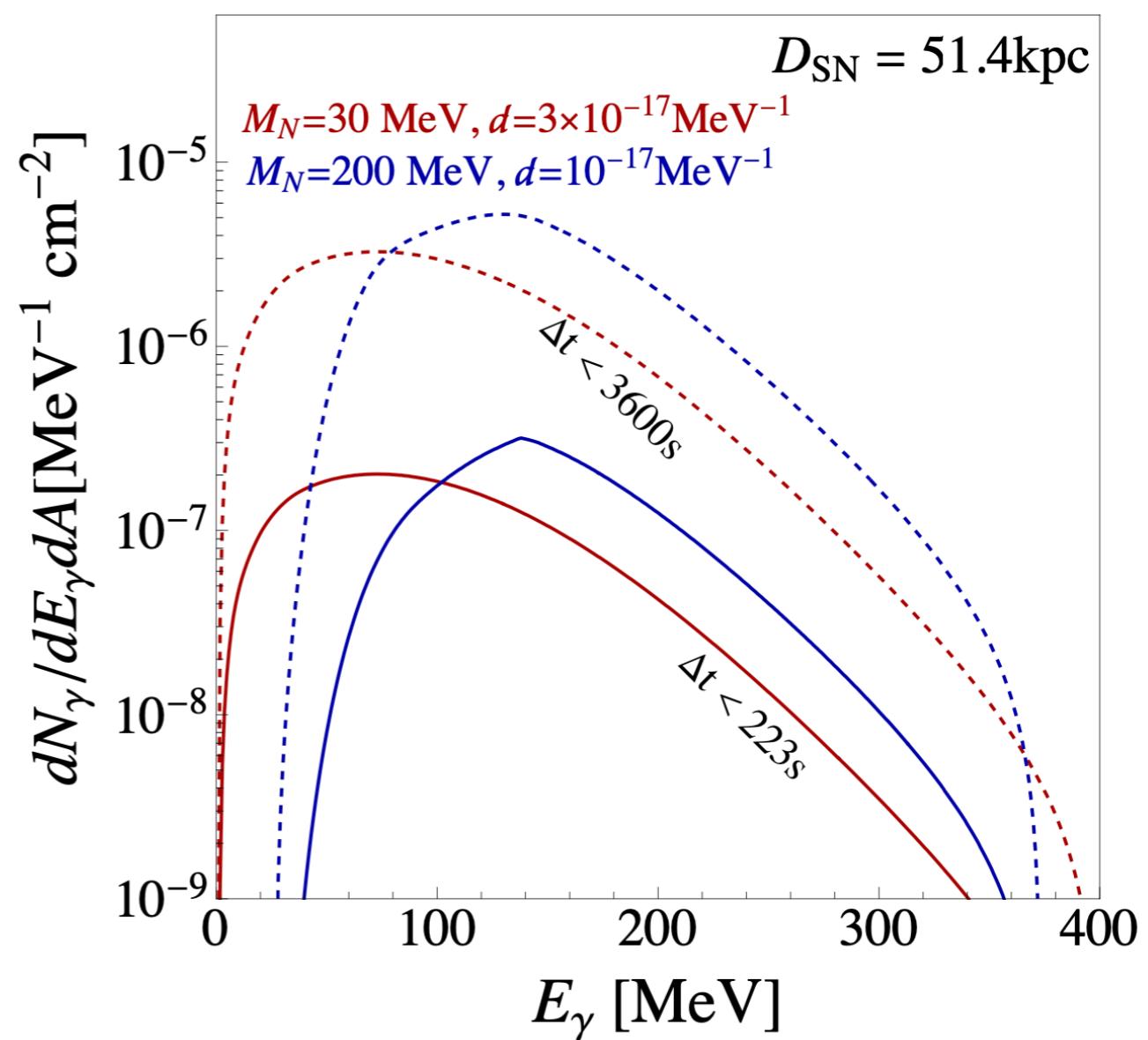
$$\beta = \sqrt{E_N^2 - M_N^2}/E_N$$

$$R_{\text{SN}}^{\gamma/\nu} \leq L_1 \leq L_1^{\max}$$

$$R_{\text{SN}}^\gamma = 3 \times 10^{10} \text{ m}$$

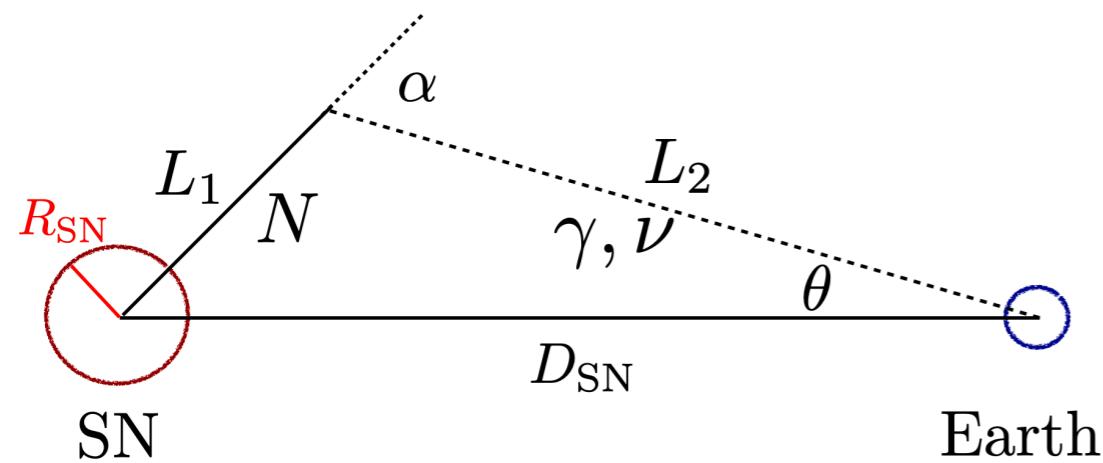
$$R_{\text{SN}}^\nu = 30 \text{ km}$$

$$\cos \alpha = \frac{2E_N E_{\gamma/\nu} - M_N^2}{2E_{\gamma/\nu} \sqrt{E_N^2 - M_N^2}}$$

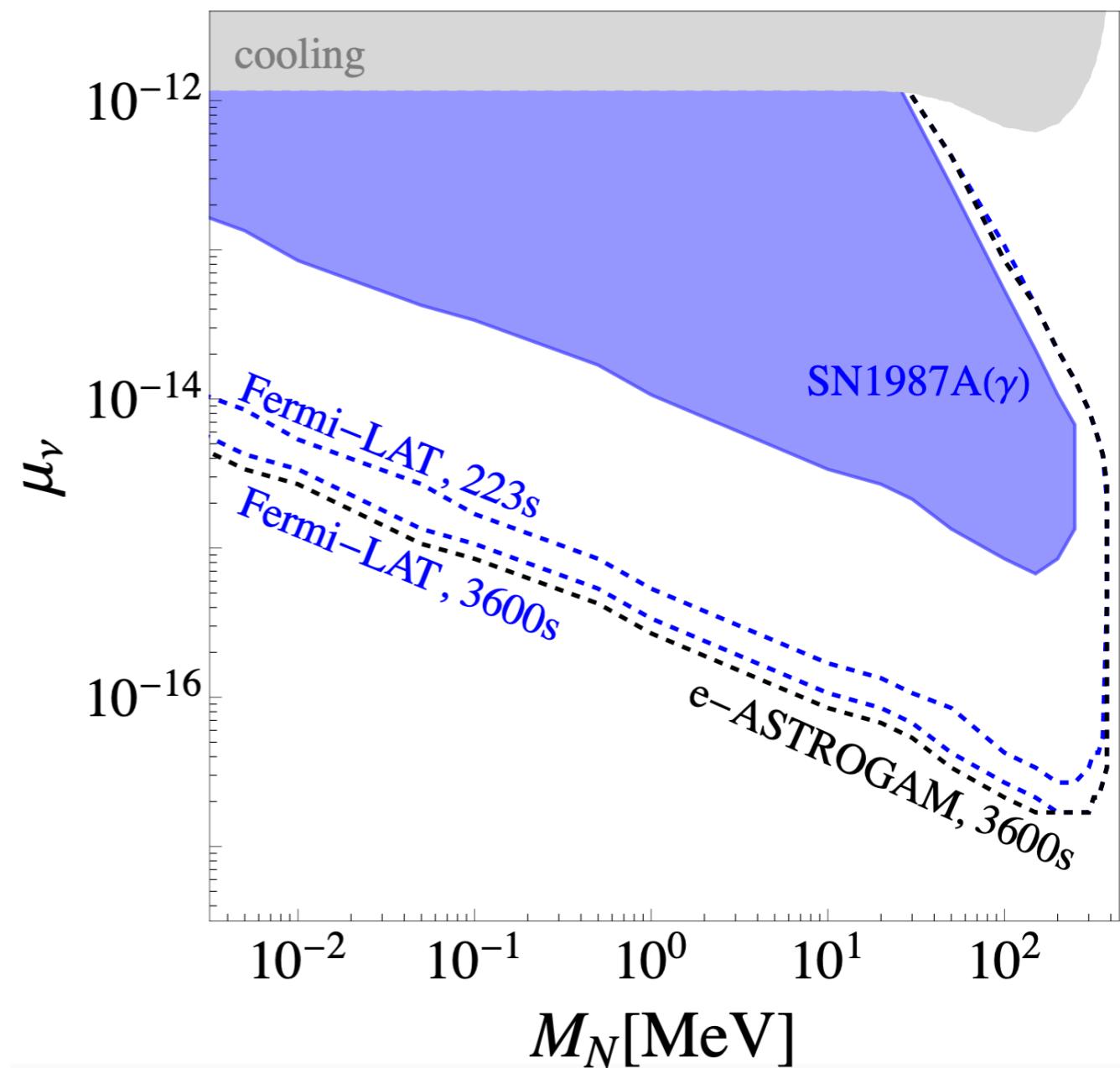


[V. Brdar, A. D. Gouvea, YYL, P. A. N. Machado, arXiv:2302.10965]

# Multimessenger Signals : $\gamma$ – ray detection



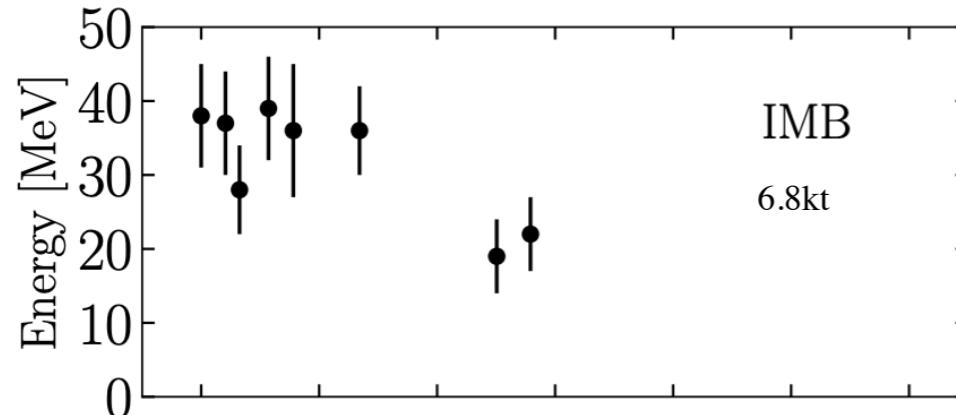
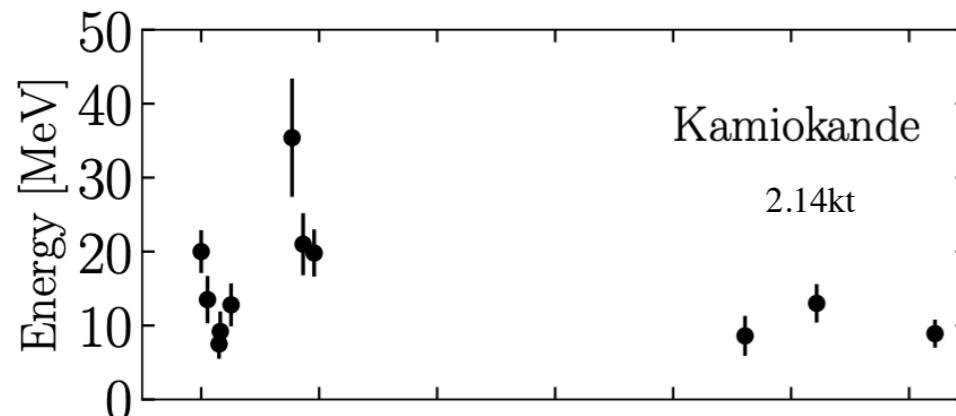
- At the time of SN1987A, the **Gamma-Ray Spectrometer (GRS)** observed  $N_{\text{obs}} = 1393$  photons with energy  $25\text{-}100 \text{ MeV}$  at  $\Delta t < 223s$
- Assuming a SN event happens in the galaxy at a distance of  $D_{\text{SN}} = 10\text{kpc}$ ,  
**Fermi-LAT**:  $E_\gamma > 100\text{MeV}, \theta < 5^\circ$   
**e-ASTROGAM**:  $E_\gamma > 1\text{MeV}, \theta < 1.25^\circ$



[V. Brdar, A. D. Gouveia, YYL, P. A. N. Machado, arXiv:2302.10965]

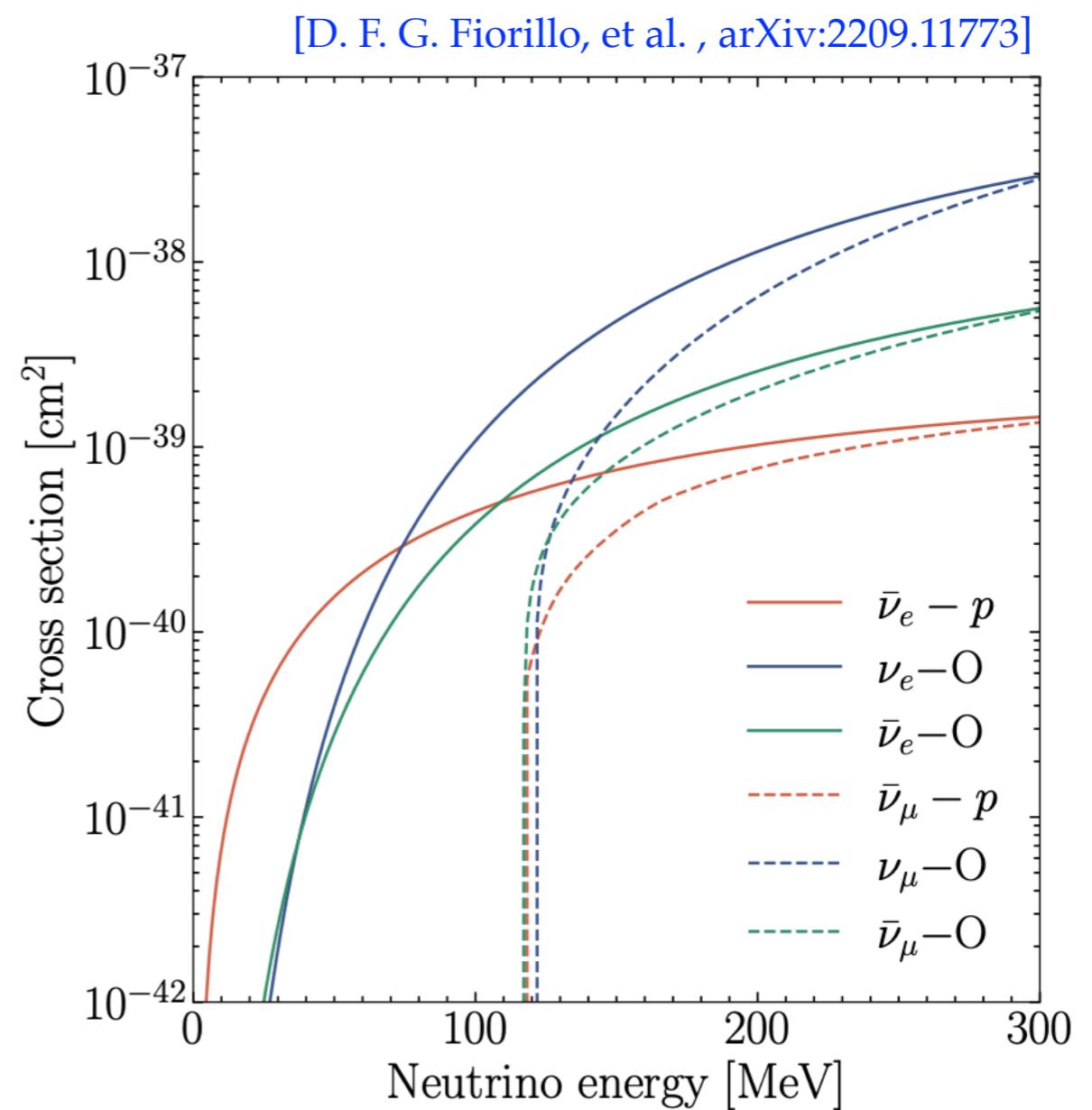
# Multimessenger Signals : neutrino detection

## SN1987A, neutrino events water-Cherenkov detectors



- No significant excess was observed by Kamiokande-II and IMB for  $E_\nu > 50\text{MeV}, \Delta t < 2 \text{ days}$

$$N_\nu^{\text{BSM}} = N_{\text{tgt}} \int dE_\nu \frac{dN_\nu}{dE_\nu dA}(E_\nu) \sigma_{\text{IBD}}(E_\nu) \varepsilon(E_\nu)$$



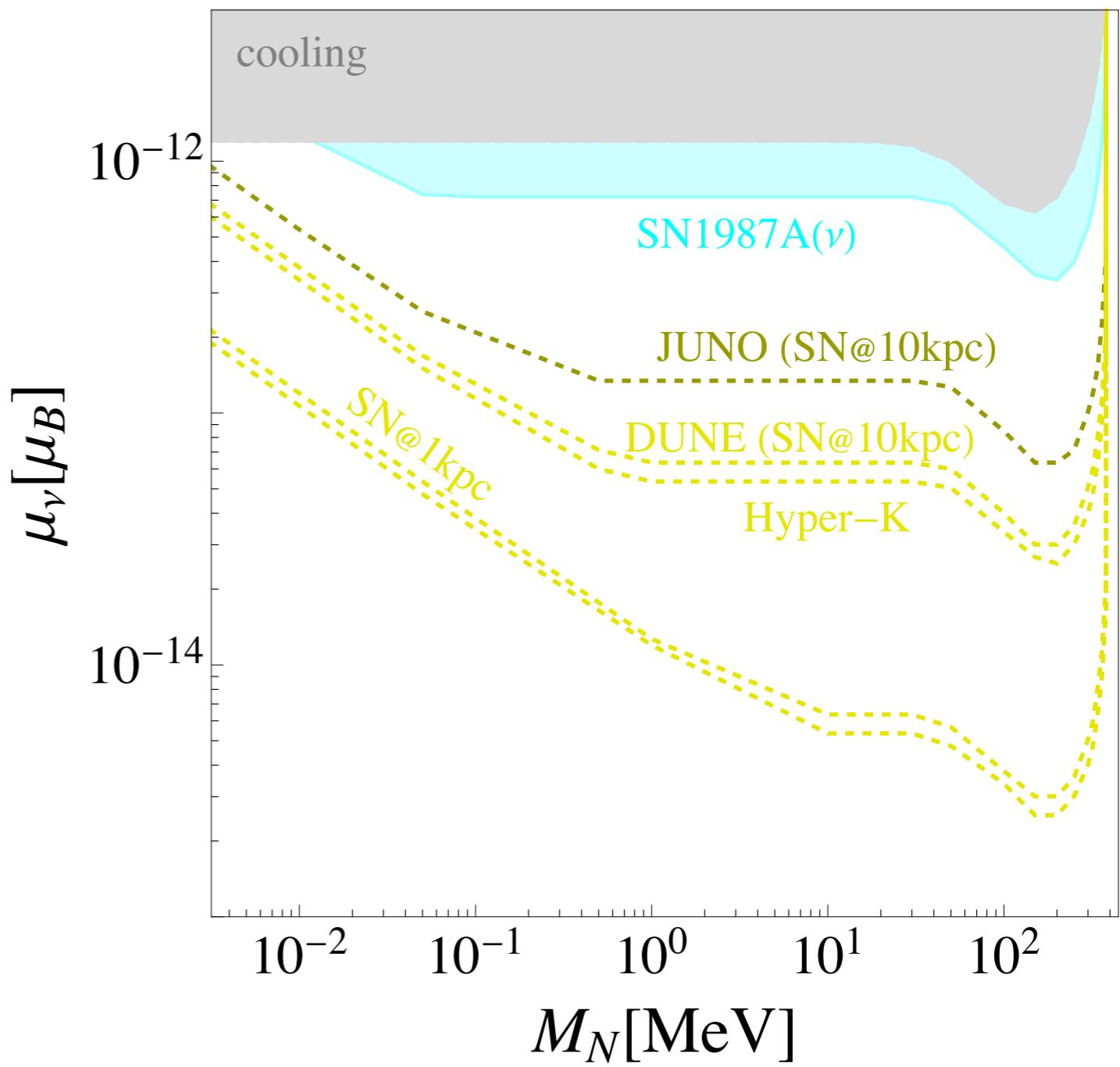
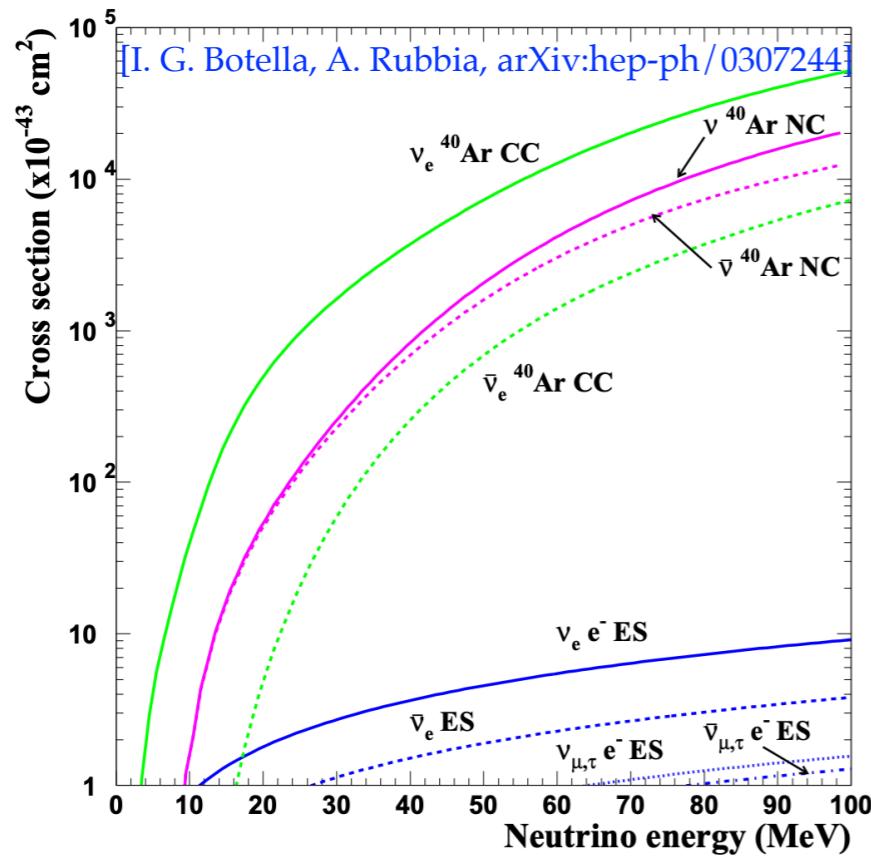
# Multimessenger Signals : neutrino detection

- Assuming a SN event happens in the galaxy at a distance of  $D_{SN} = 10\text{kpc}$ ,

**JUNO:** 20kt fiducial volume, liquid scintillator detector

**DUNE:** 40kt, liquid argon

**Hyper-K:** 188kt fiducial volume  
water Cherenkov

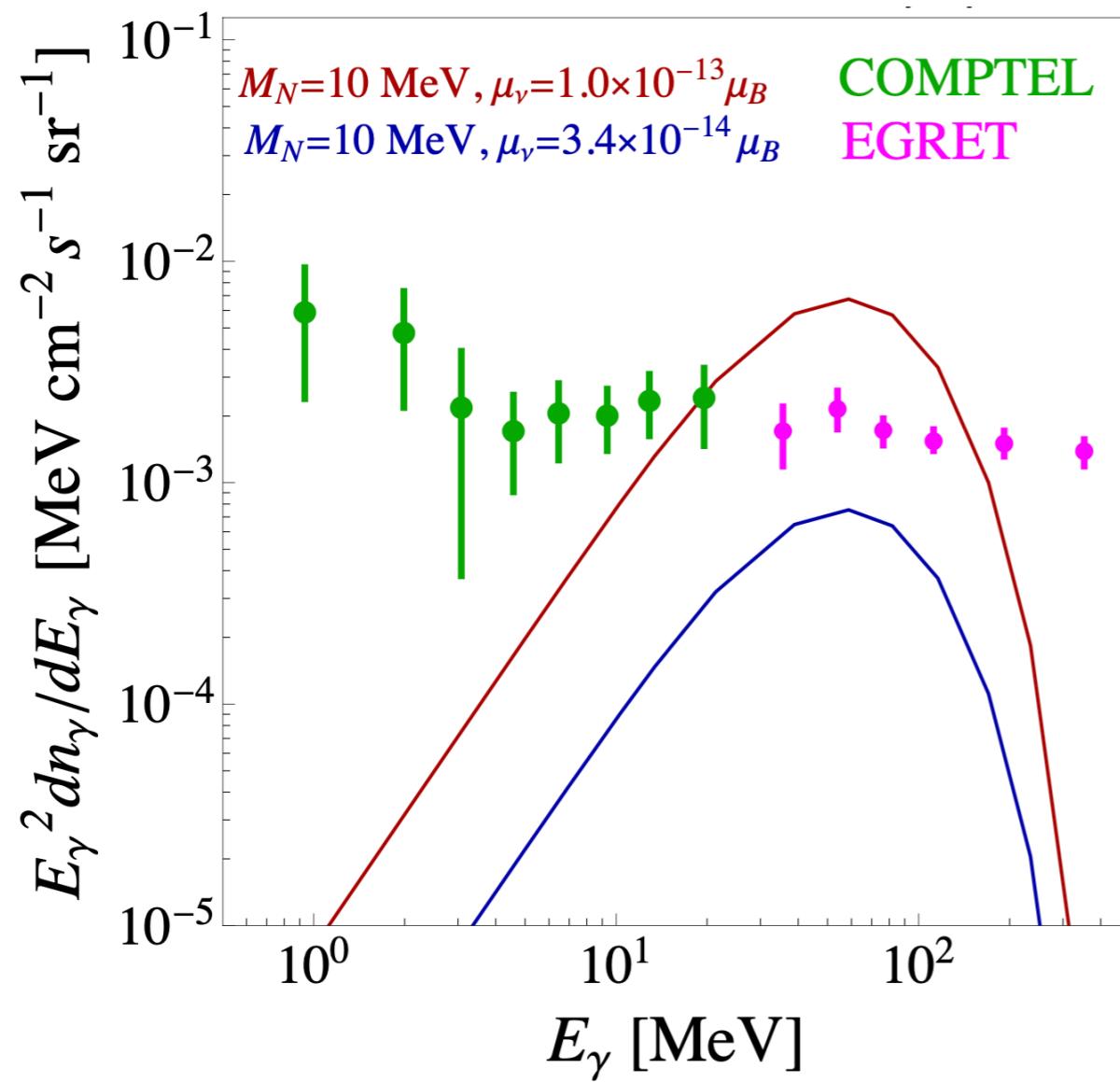


[V. Brdar, A. D. Gouveia, YYL, P. A. N. Machado, arXiv:2302.10965]

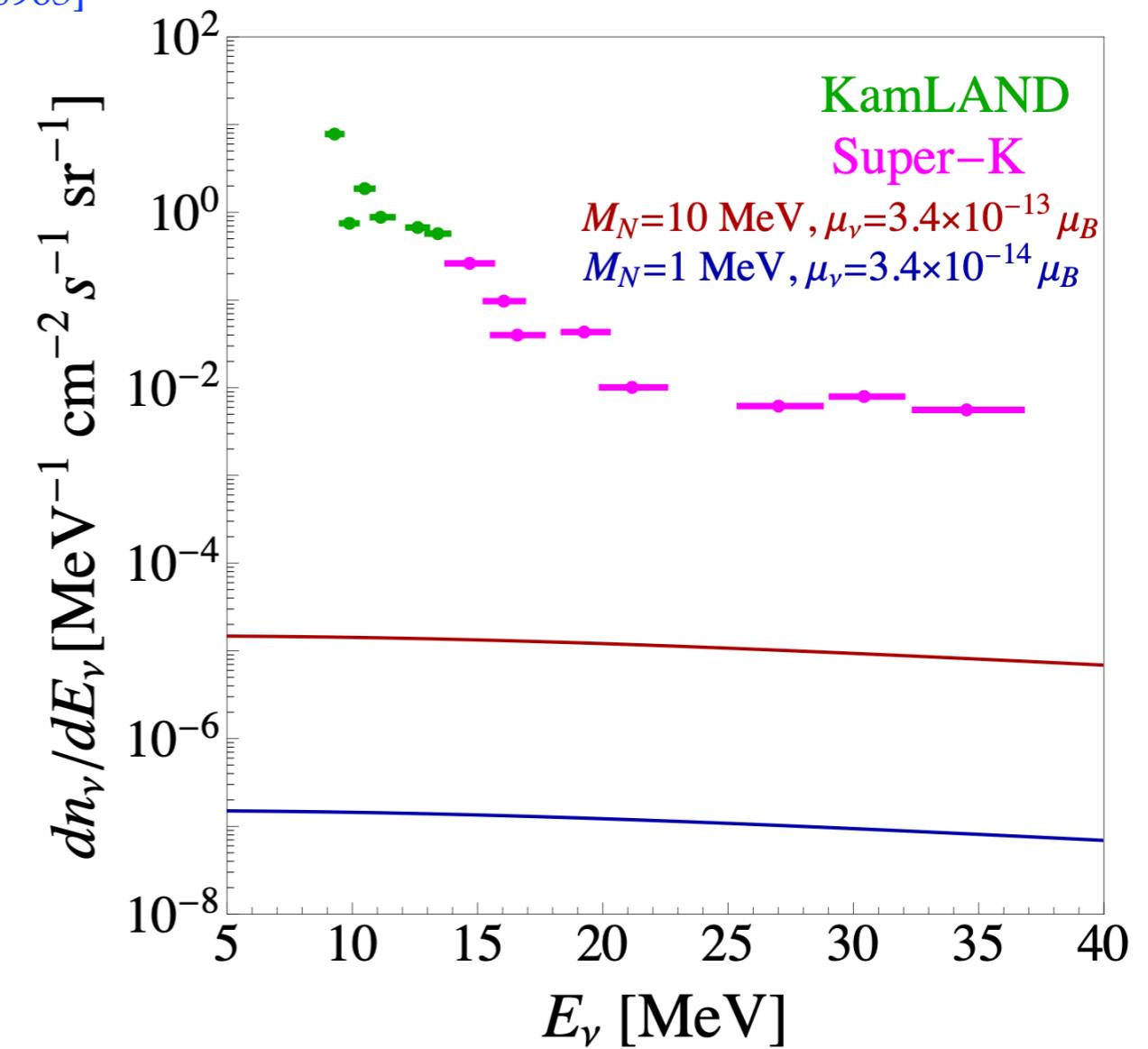
# Multimessenger Signals : diffused BSM Photon and neutrino background

$$\frac{dn_N}{dE} = \frac{c}{4\pi} \int_0^\infty dz(1+z) n'_{cc}(z) \frac{d\mathcal{N}_s}{dE}(E_z)$$

[V. Brdar, A. D. Gouv^ea, YYL, P. A. N. Machado, arXiv:2302.10965]



extragalactic background light

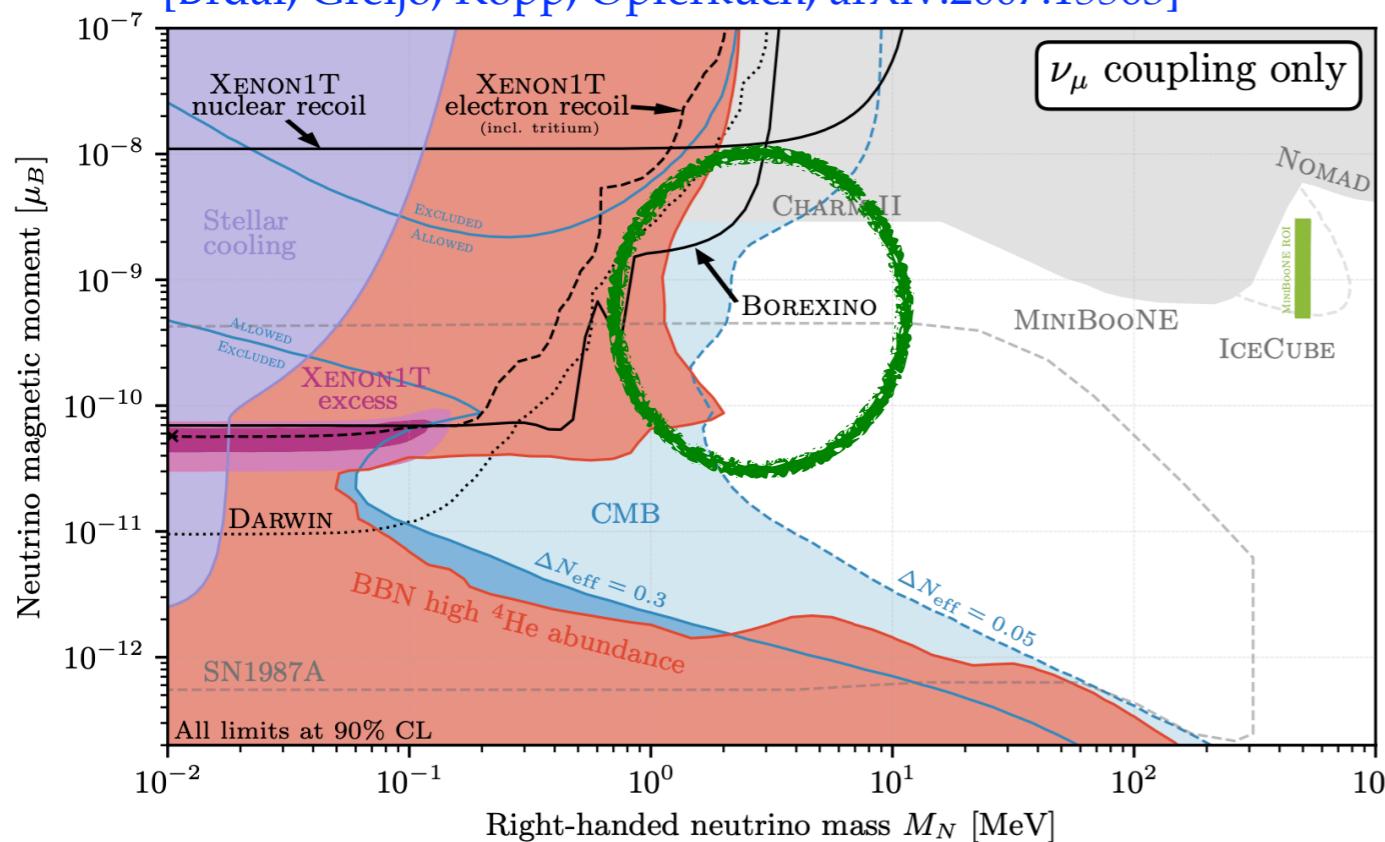


# Conclusion? Outlooks

## Solar neutrino —

Can JUNO cover the green circle region, given that its bigger than Borexino, also probes electron recoils below 100keV.

[Brdar, Greljo, Kopp, Opferkuch, arXiv:2007.15563]

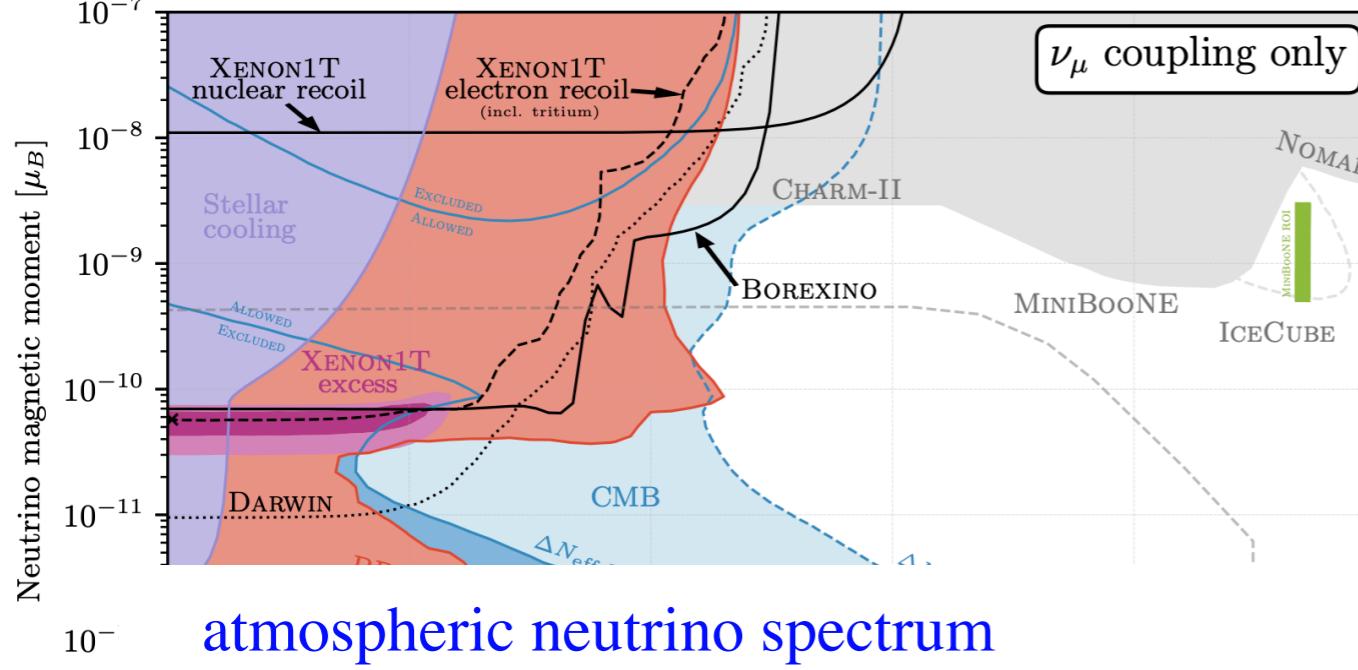


[Z. Ye, F. Zhang, D. Xu, J. Liu, arXiv:2103.11771]

$\mu\nu < 10^{-11} \mu\text{B}$  can be reached at an energy threshold greater than 40keV (SM case)

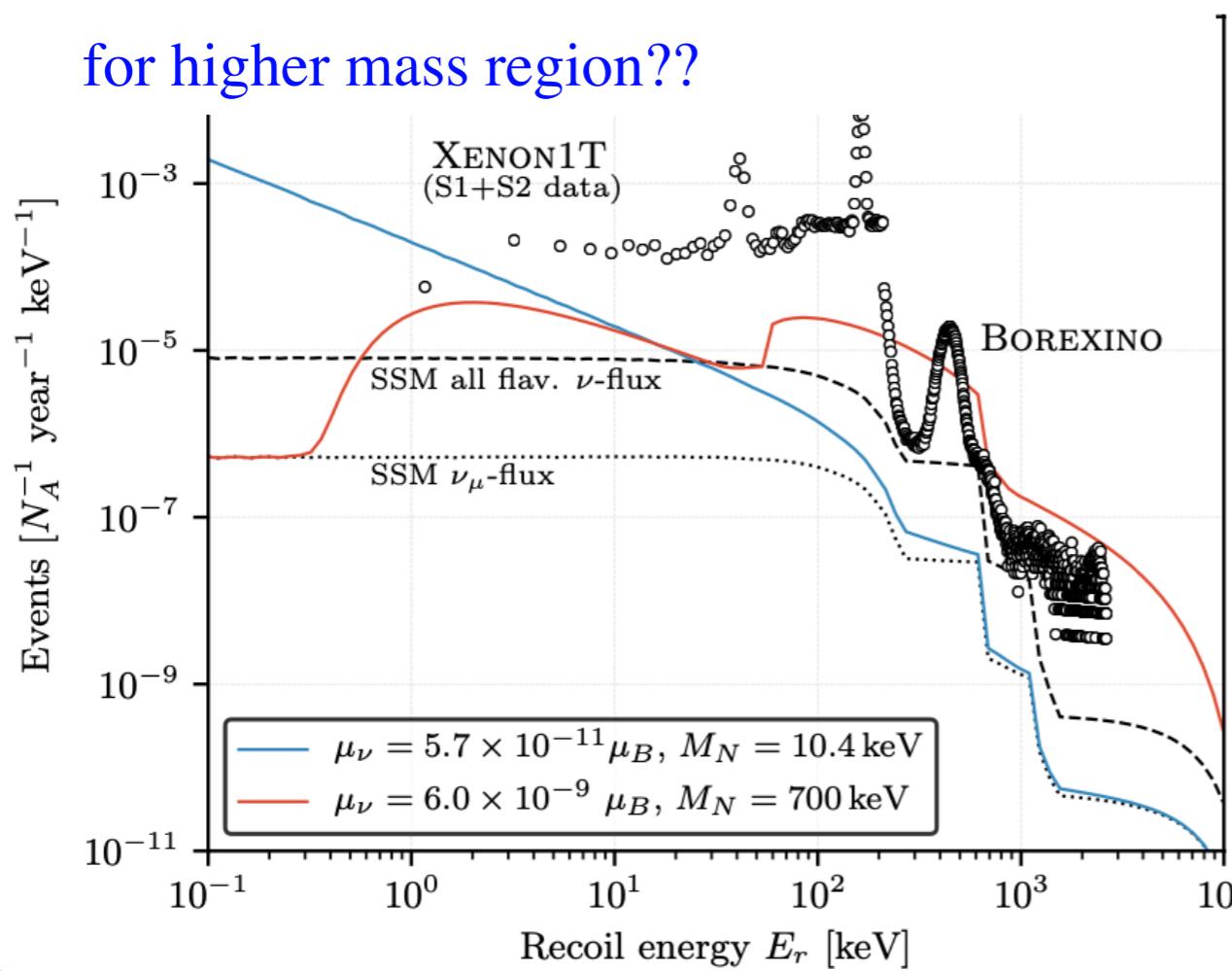
# Conclusion? Outlooks

[Brdar, Greljo, Kopp, Opferkuch, arXiv:2007.15563]

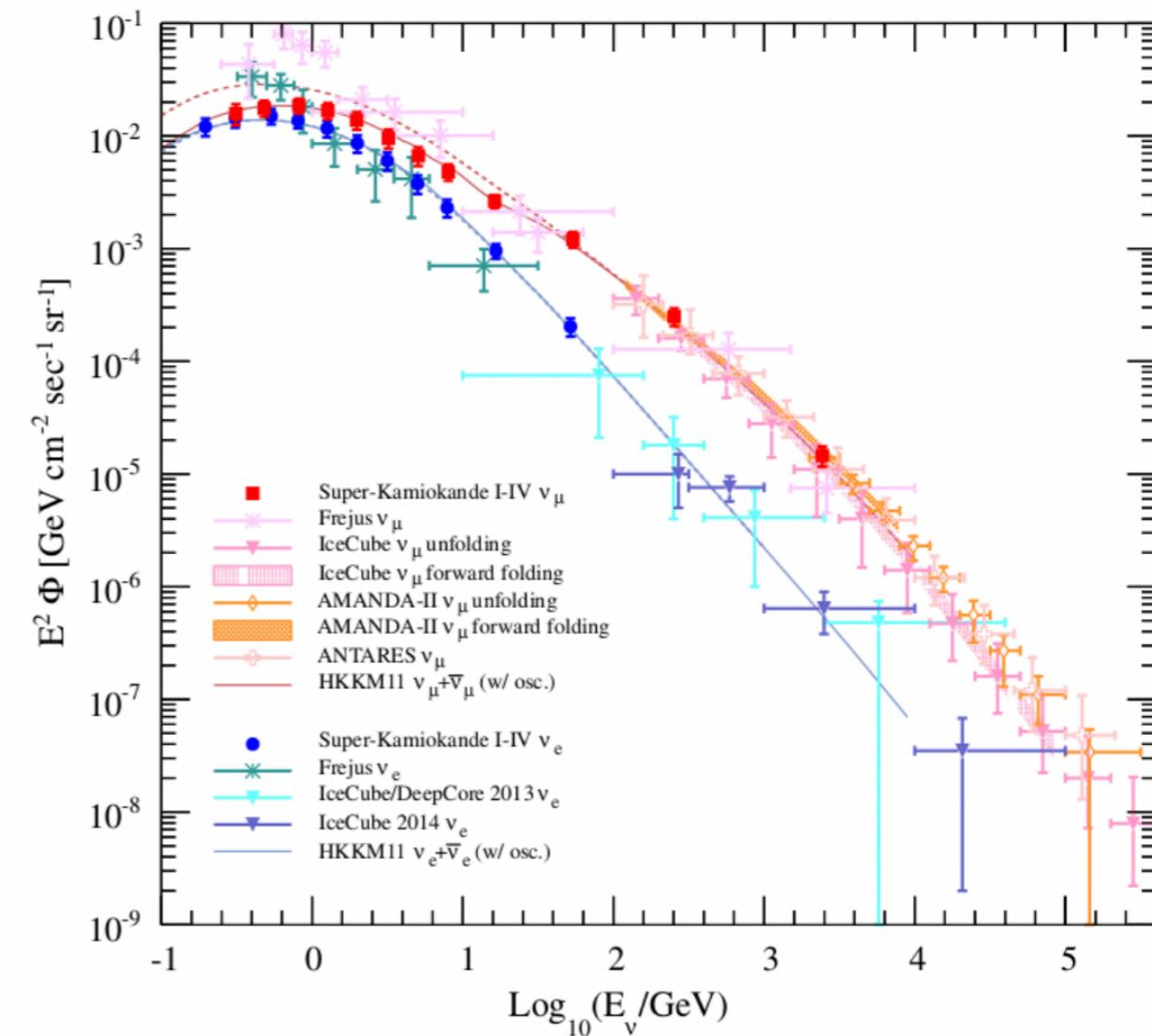


atmospheric neutrino spectrum

for higher mass region??



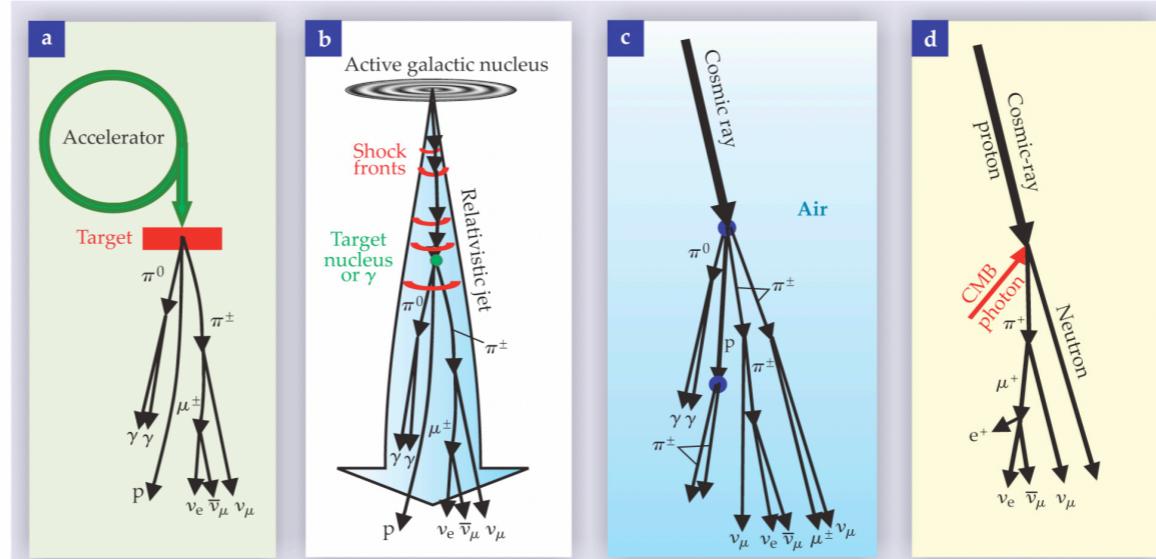
[JUNO, arXiv:2103.09908]



# Sterile Neutrino: Multi-messenger signals for Photon detection and Neutrino experiments

Halzen Klein 2008

- ◆ Astrophysical neutrinos

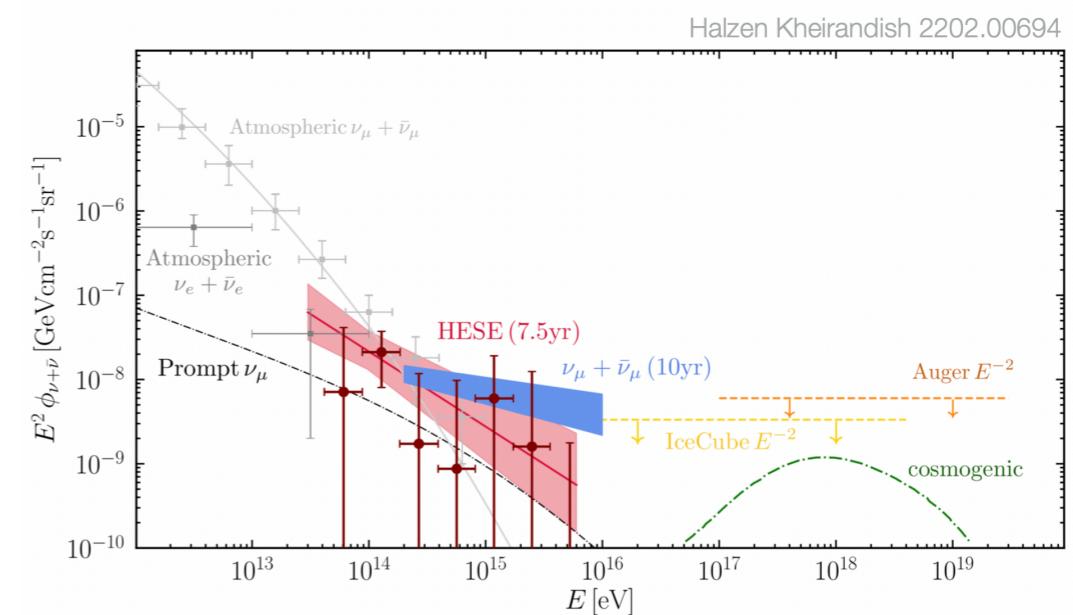


Sterile neutrino  $N$  production via mixing:  $\nu \rightarrow N$

$N$  decaying to photon that can be observed by photon detectors:  $N \rightarrow \nu + \gamma$

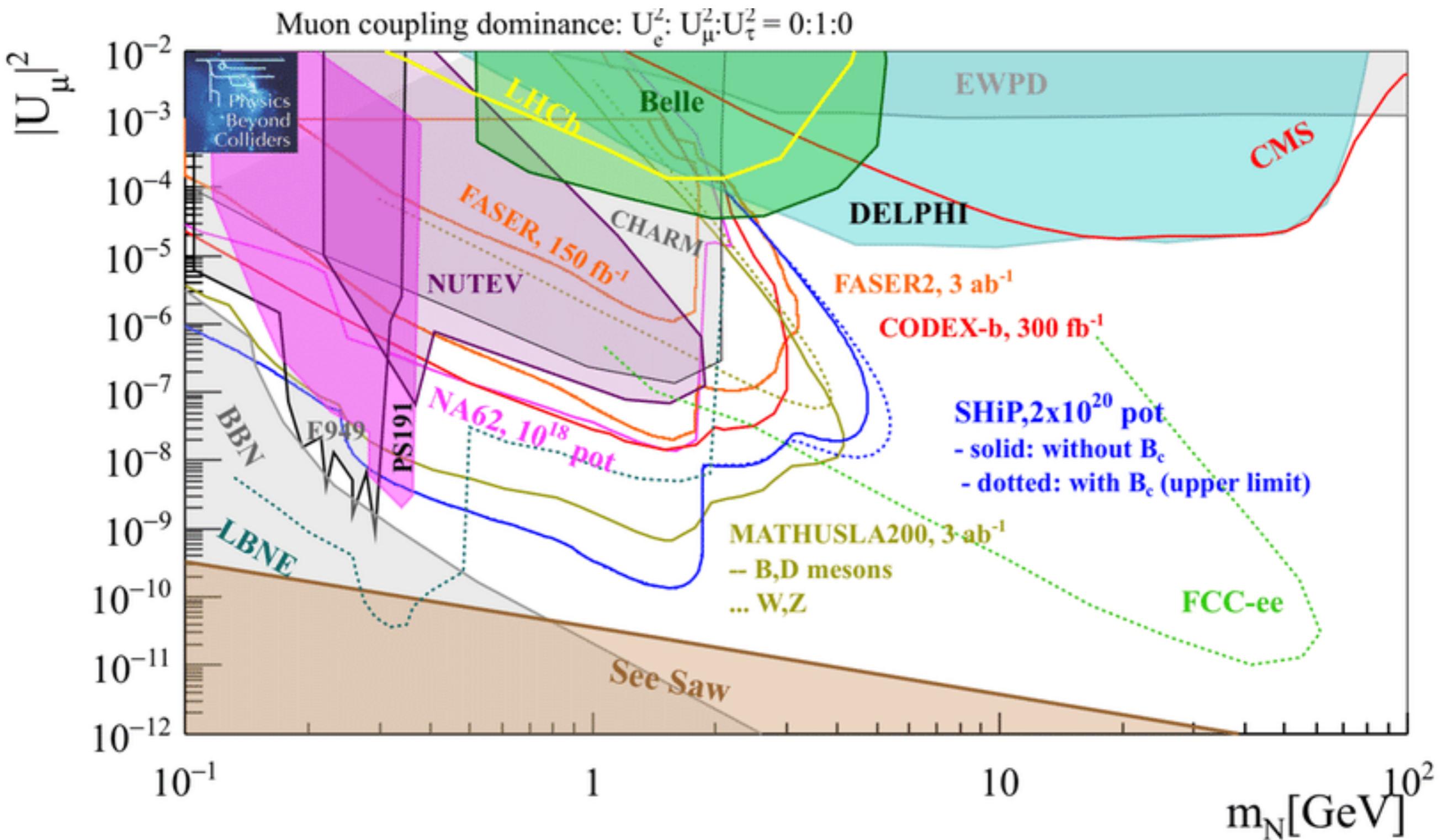
- ◆ Diffused neutrinos fluxes being observed or bounded

LHAASO's measurements of diffused TeV photon flux might already constrains parameter space of sterile neutrino?



Thank you

# BACK UP



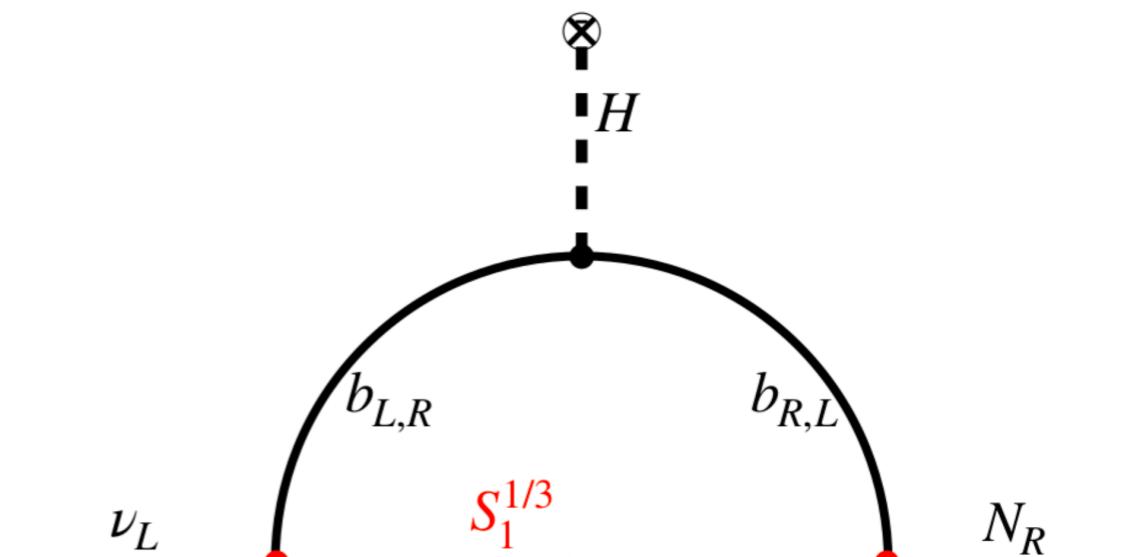
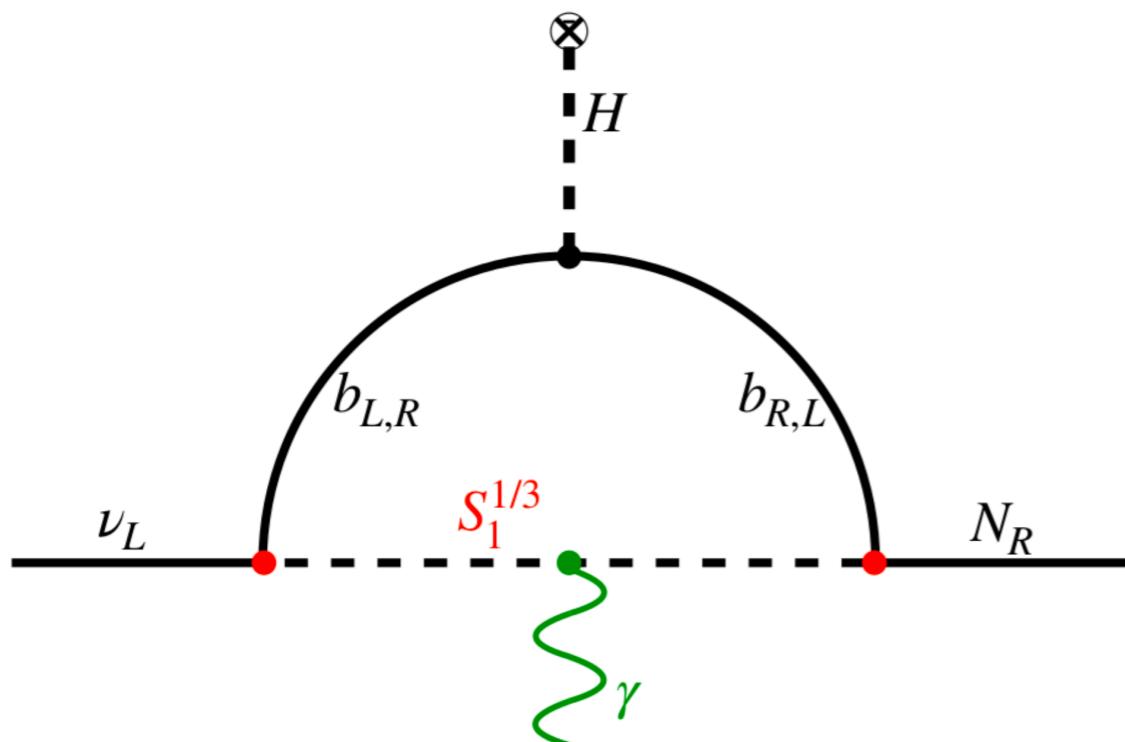
# UV completion

In what theory the effect of mixing is subdominant...

$$\theta^2 \sim \frac{m_\nu}{M_N}, \text{keV} < M_N < 100\text{MeV}$$

- Consider heavy scalar lepto-quark  $S_1 \sim (\bar{3}, 1, 1/3)$

$$\mathcal{L} \supset y_1 \bar{b}_R^c N S_1 + y_2 \bar{Q}_L^3 L_L^i {}^c S_1^\dagger + h.c.$$



$$\mu_\nu \approx \frac{e y_1 y_2}{8\pi^2 m_{LQ}^2} m_b \log \frac{m_b^2}{m_{LQ}^2}$$

$$m_{\nu N} \sim \frac{\mu_\nu}{\mu_B} \frac{m_{LQ}^2}{2m_e}$$

# UV completion

In what theory the effect of mixing is subdominant...

$$\theta^2 \sim \frac{m_\nu}{M_N}, \text{keV} < M_N < 100\text{MeV}$$

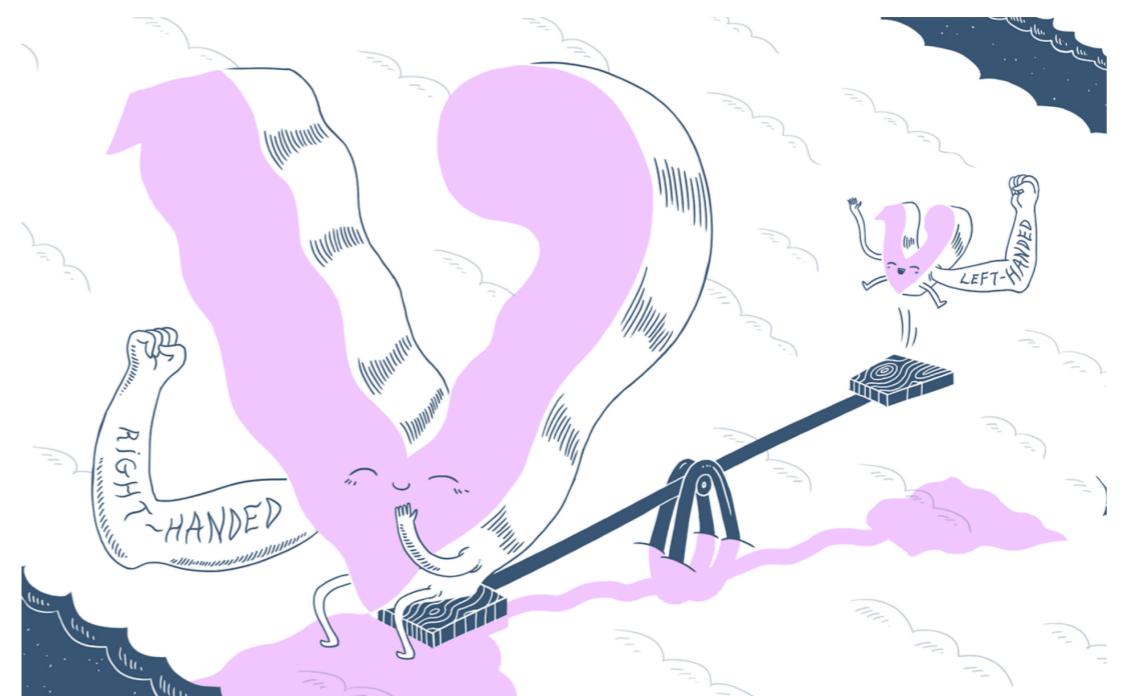
- Consider heavy scalar lepto-quark  $S_1 \sim (\bar{3}, 1, 1/3)$

$$\mathcal{L} \supset y_1 \bar{b}_R^c N S_1 + y_2 \bar{Q}_L^3 L_L^i {}^c S_1^\dagger + h.c.$$

$$m_{\nu N} \sim \frac{\mu_\nu}{\mu_B} \frac{m_{LQ}^2}{2m_e}$$

$$m_\nu \sim \frac{m_{\nu N}^2}{M_N}$$

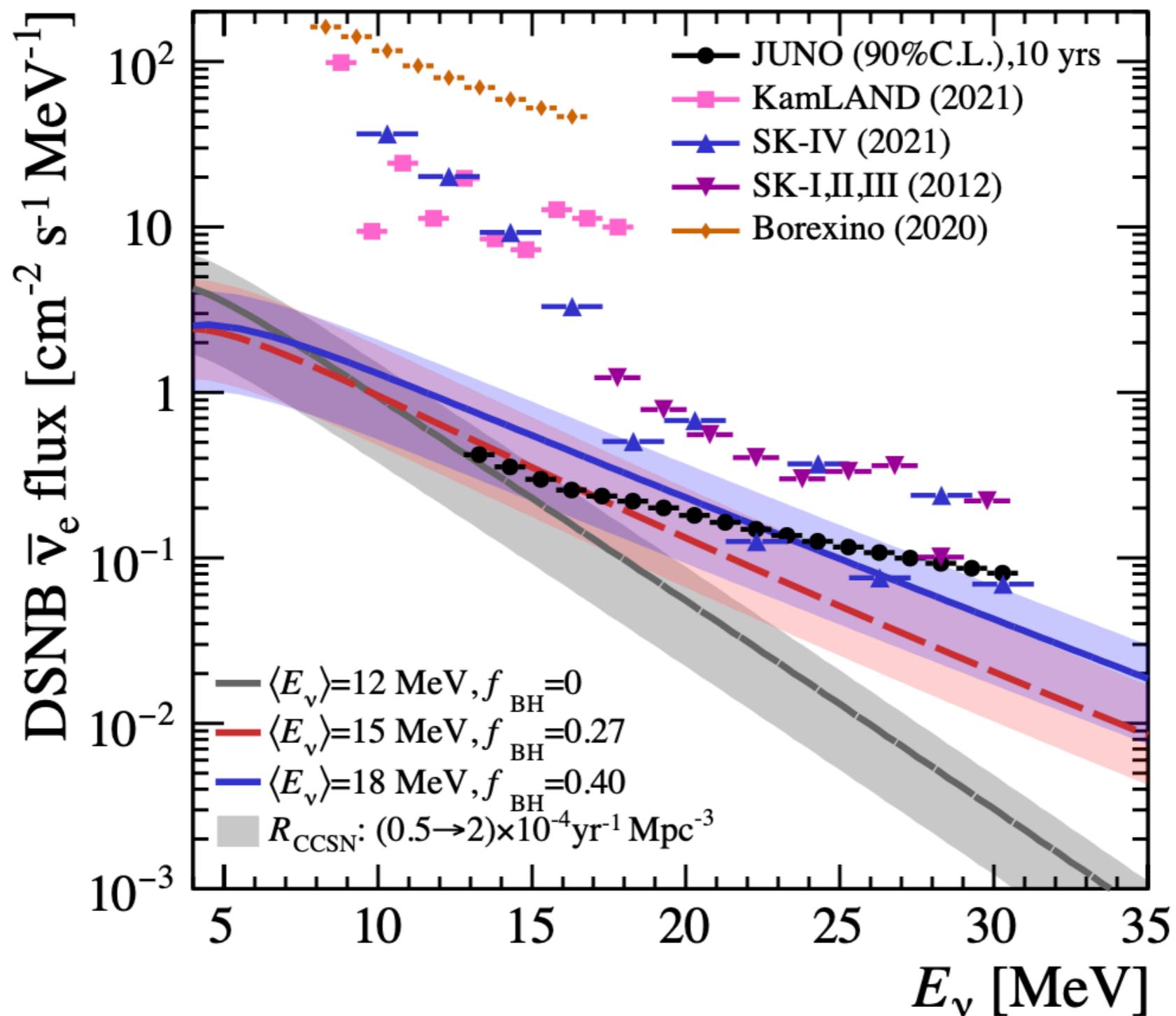
$$\mu_\nu \sim 10^{-12} \mu_B \frac{\sqrt{m_\nu M_N}}{\text{MeV}}$$



for lepto-quark masses at the TeV scale

$$m_\nu \sim 0.1\text{eV}, \mu_\nu \sim 10^{-15} \mu_B$$

# Multimessenger Signals : diffused BSM Photon and neutrino background



[A. Abusleme, arXiv:2205.08830]