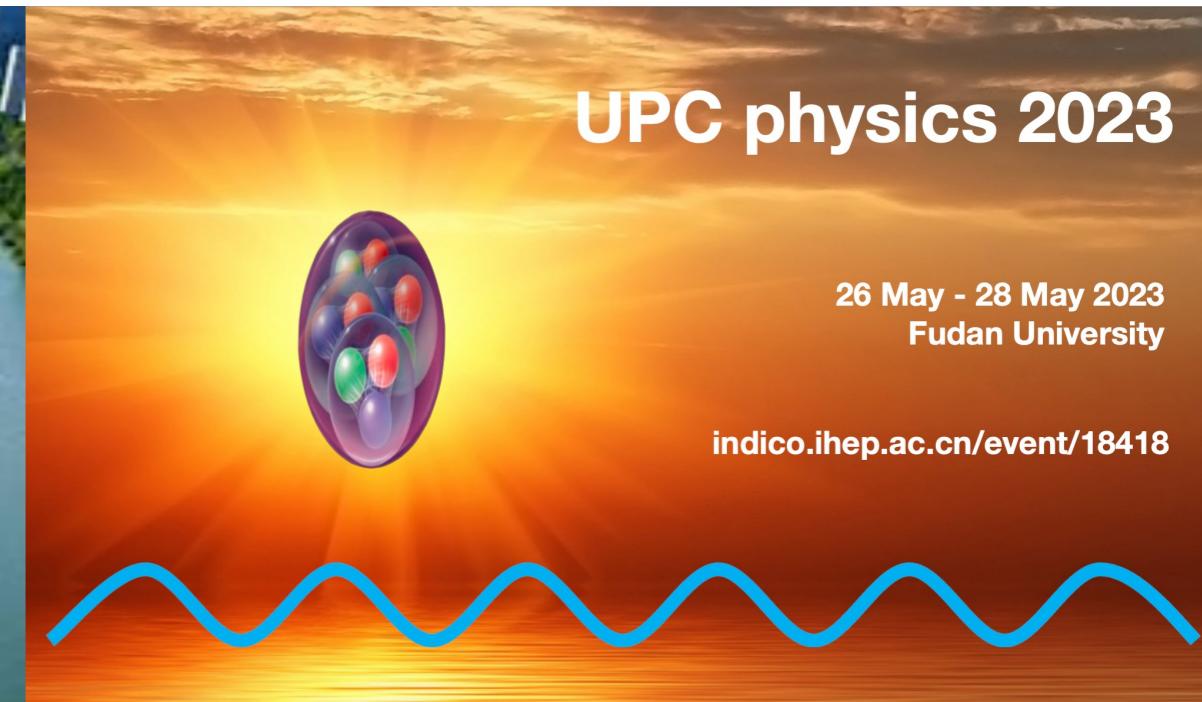




Impact parameter dependence of photon-induced processes in heavy-ion collisions

杨 帅

华南师范大学



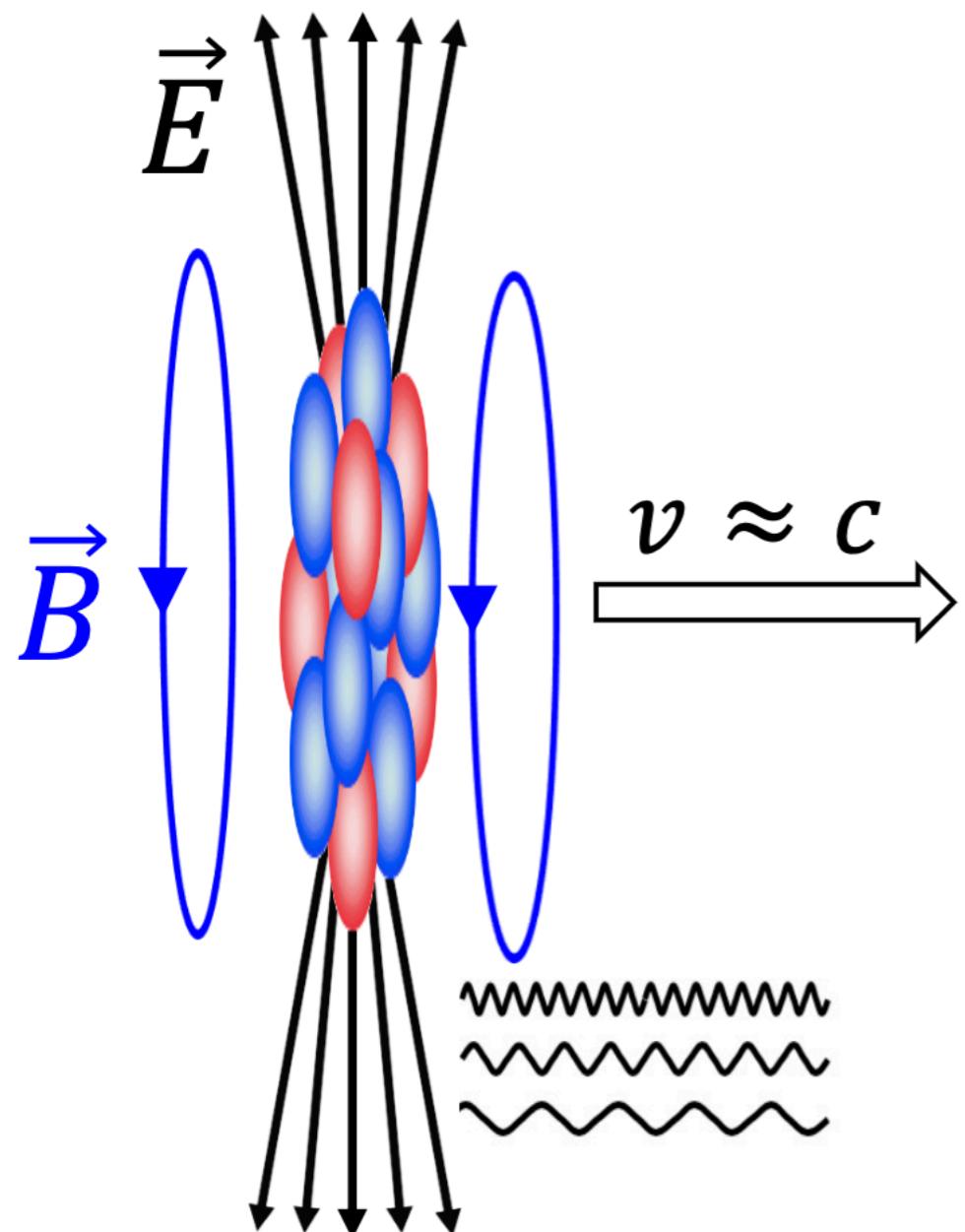
Local organizers:
Xu-Guang Huang
Guo-Liang Ma
Ding-Yu Shao
Jie Zhao



Equivalent photon

- Equivalent Photon Approximation

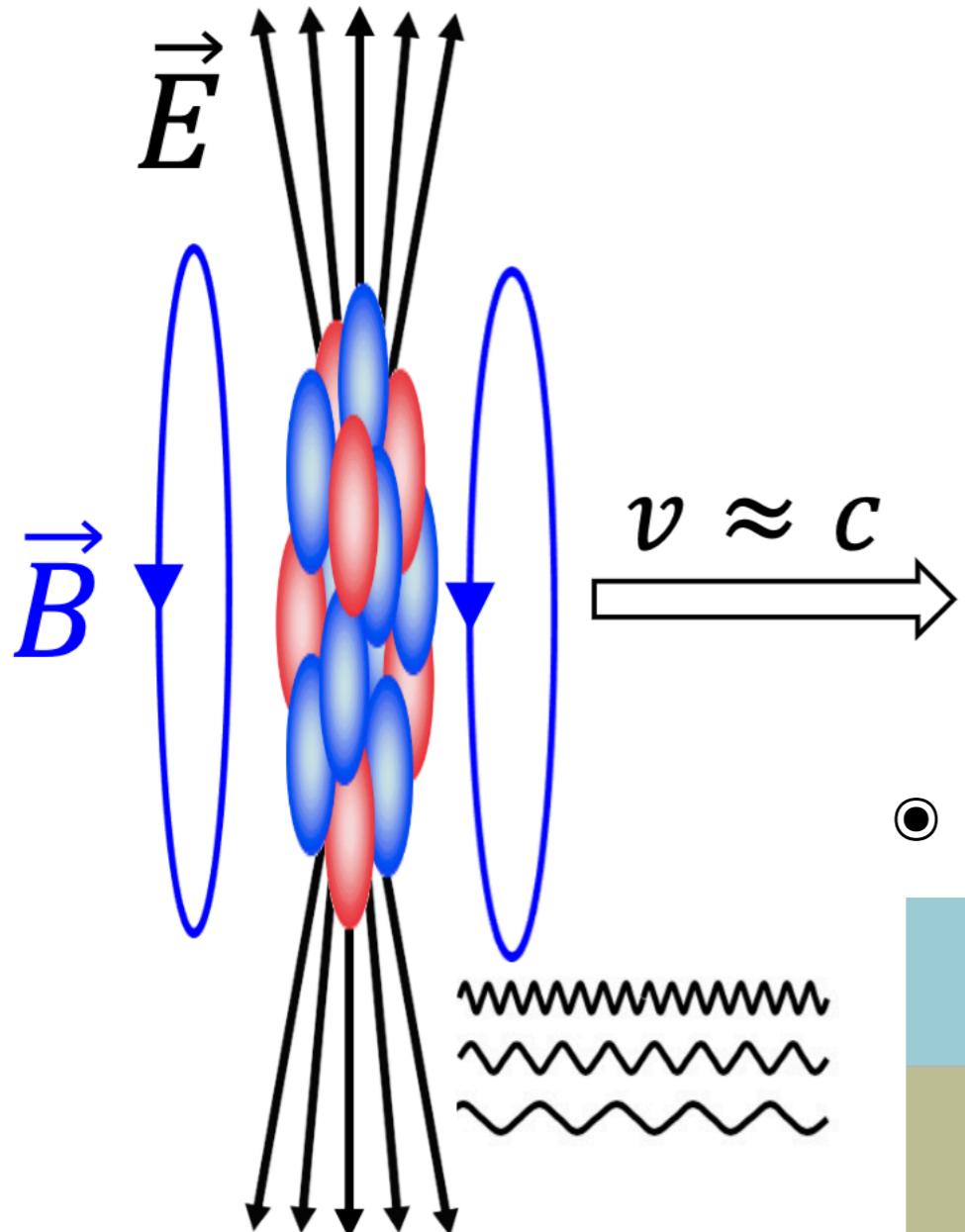
- Photon Flux $\propto Z^2$
- $p_T < \sim 30 \text{ MeV}$ ($Q^2 \sim 0$)



Equivalent photon

- Equivalent Photon Approximation

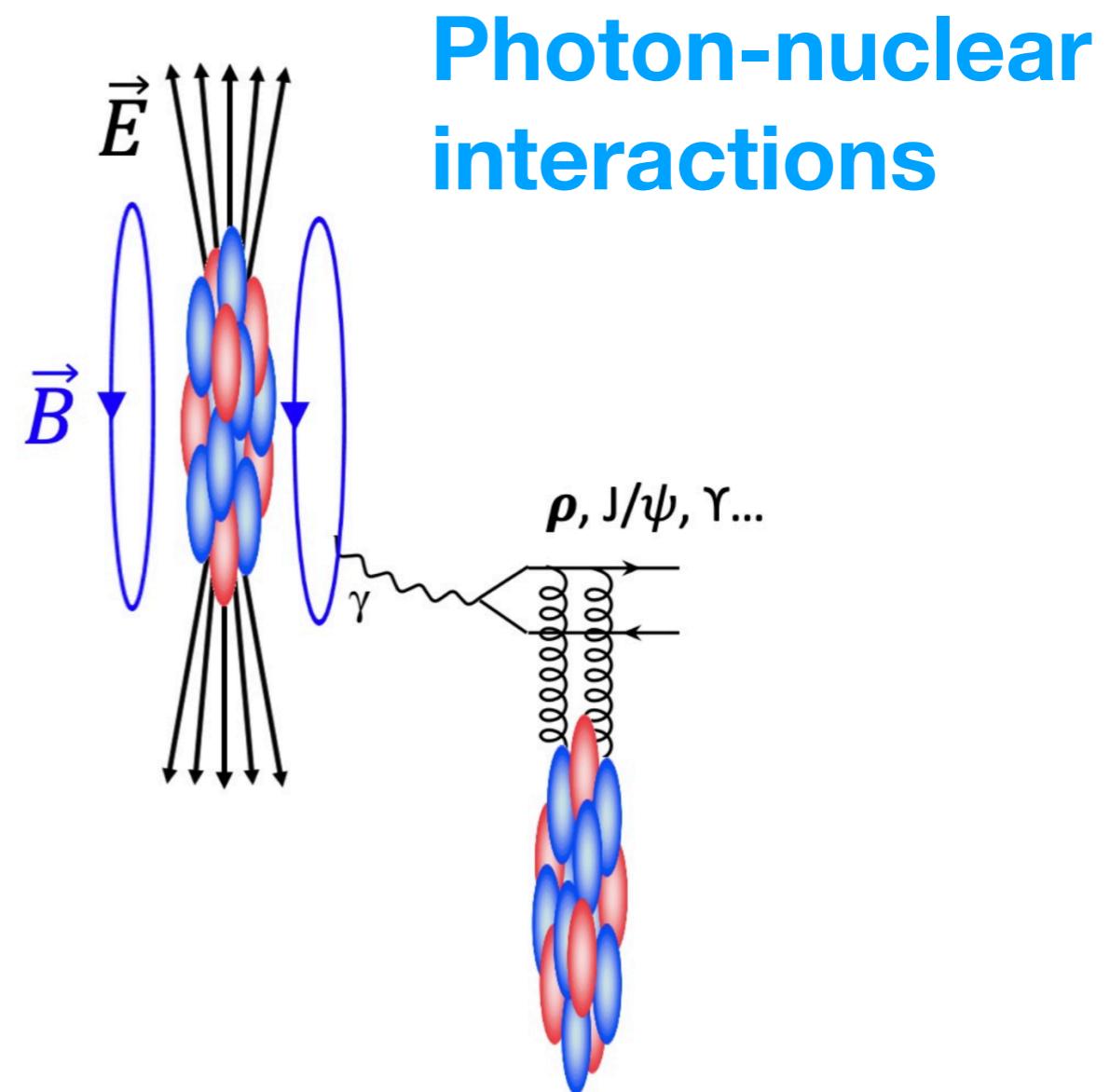
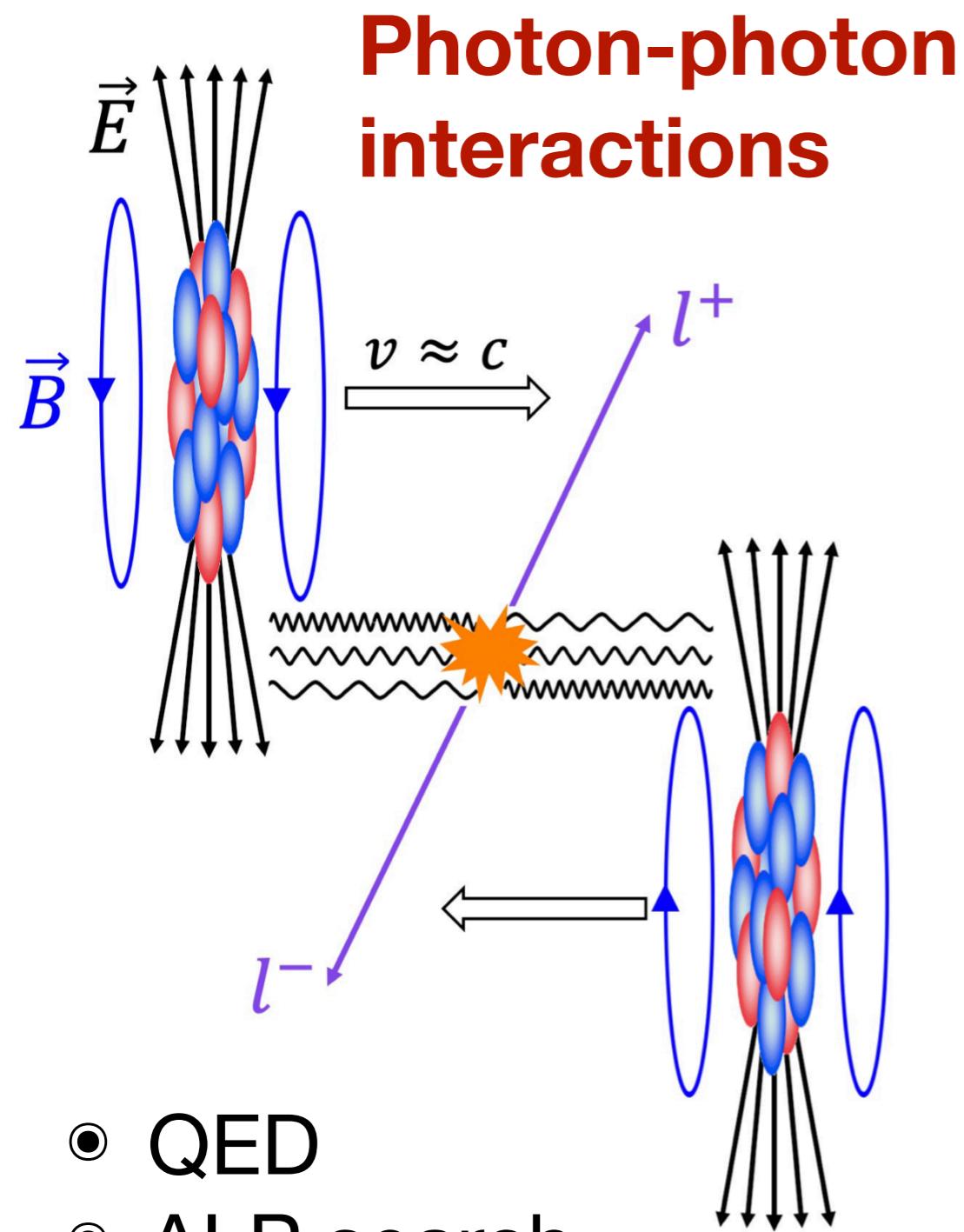
- Photon Flux $\propto Z^2$
- $p_T < \sim 30 \text{ MeV}$ ($Q^2 \sim 0$)



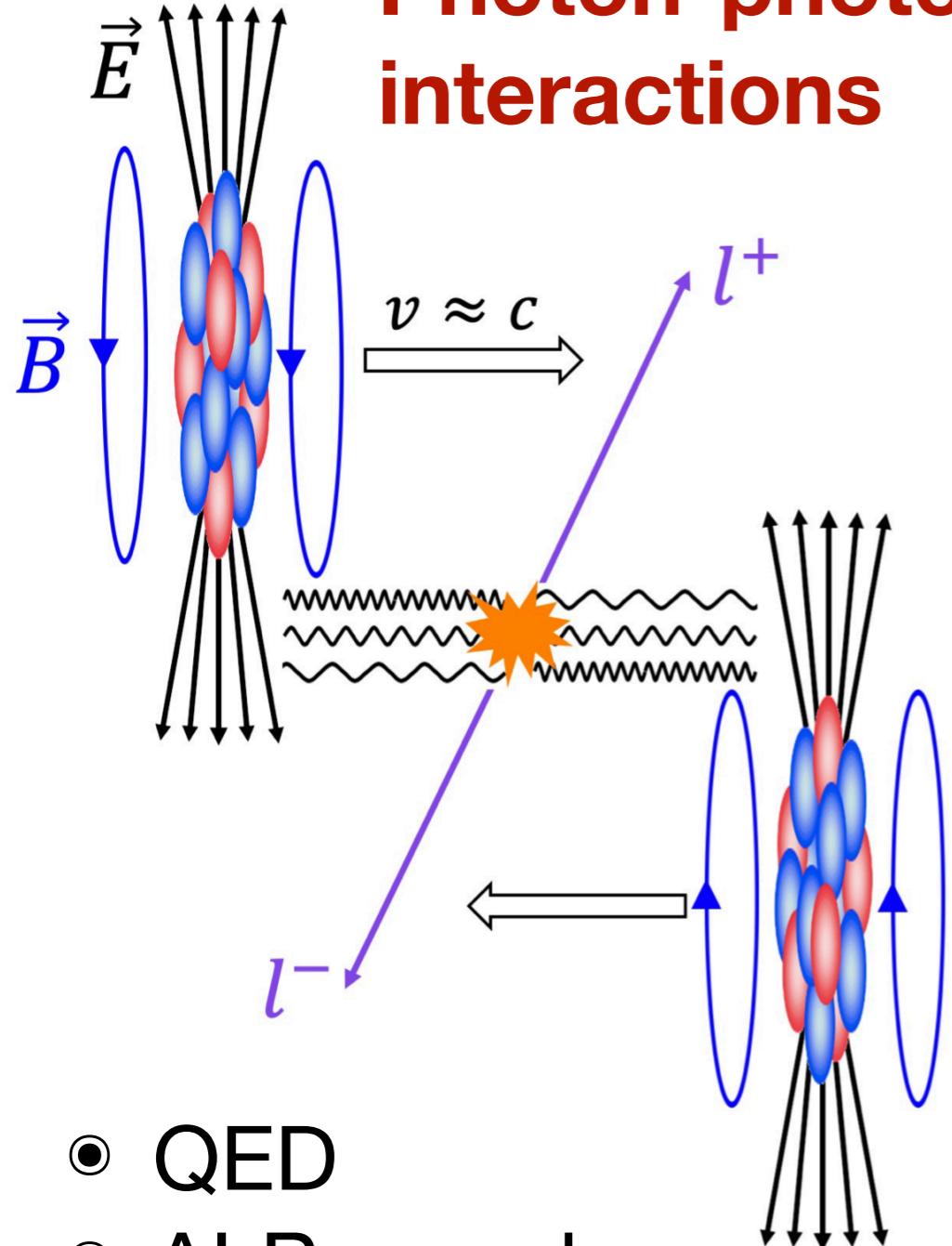
- Photon kinematics

maximum energy $E_{\gamma,\max} \sim \gamma(\hbar c/R)$	80 GeV in Pb+Pb@LHC 3 GeV in Au+Au@RHIC
typical p_T (& virtuality) $p_{T\max} \sim \hbar c/R$	O(30) MeV @ RHIC & LHC
Coherent strengths (rates) scale as Z^2 : nuclei >> protons	Flux of photons on other nucleus $\sim Z^2$, flux of photons on photons $\sim Z^4$ (45M!)

Photon interactions

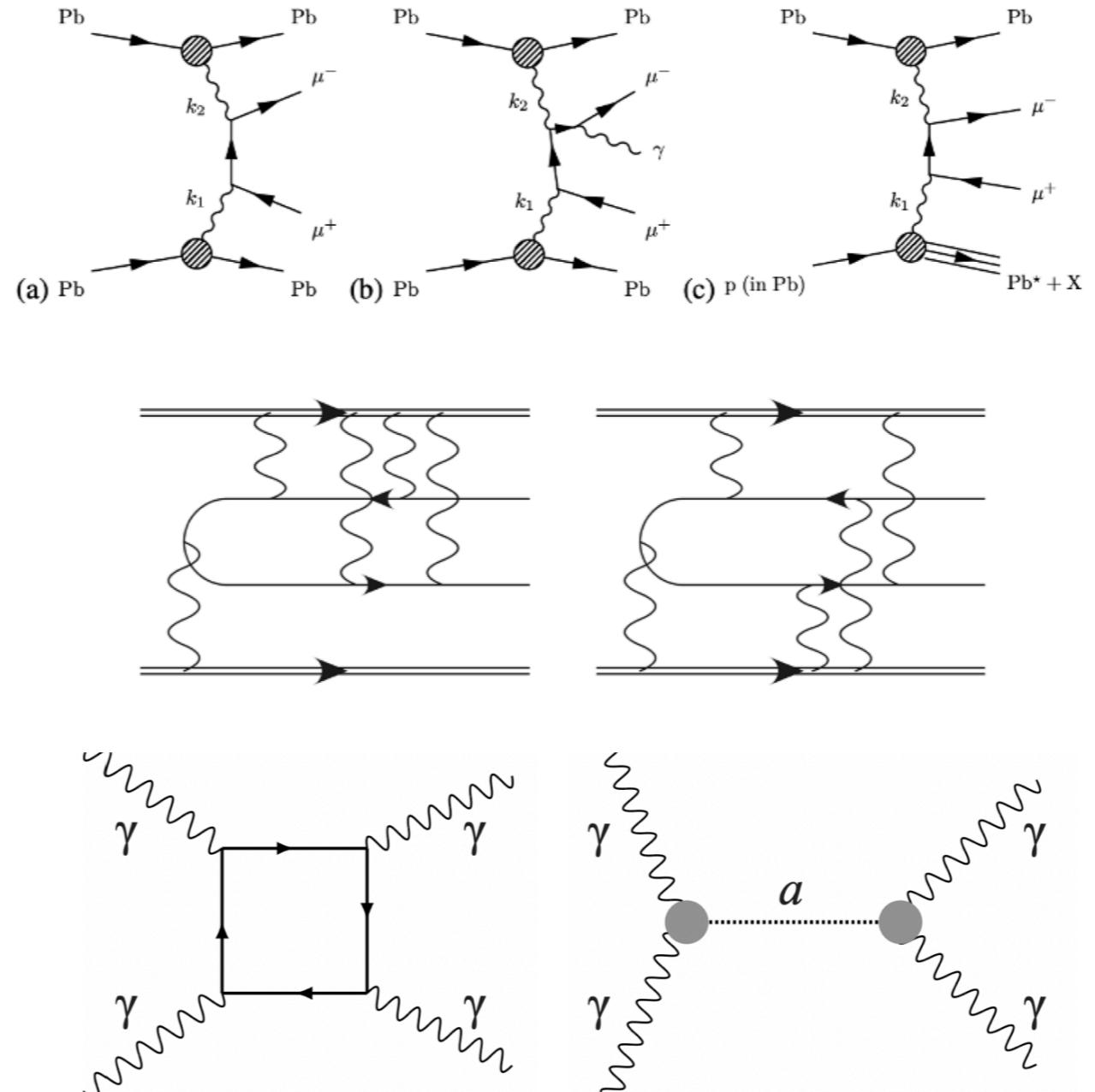


Photon interactions



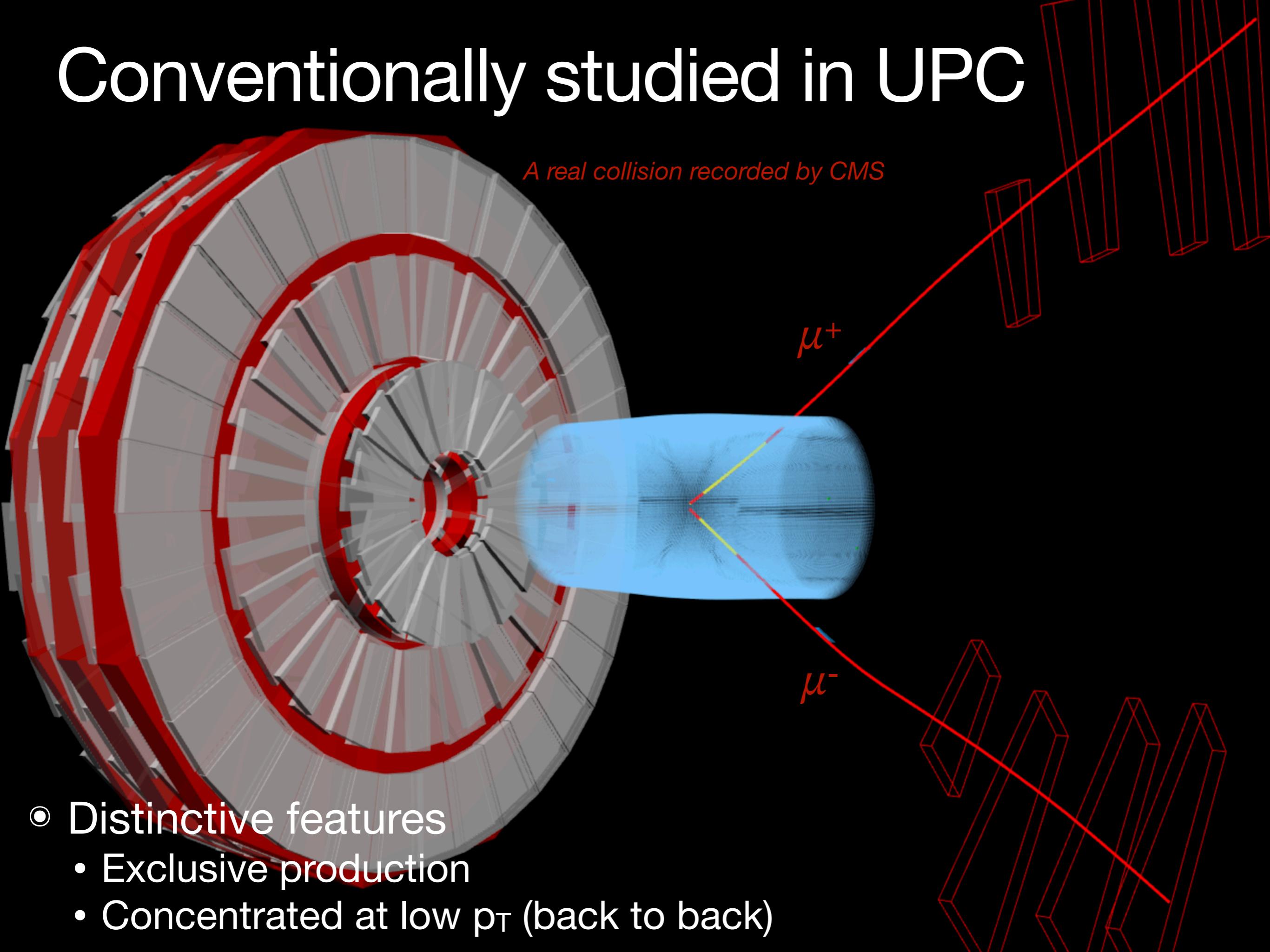
Shuai Yang

- QED
- ALP search



STAR, *PRC* 70 (2004) 031902; *PRL* 121 (2018) 132301;
PRL 127 (2021) 052302
ATLAS, *Nat. Phys.* 13 (2017) 852; *PRL* 121 (2018) 212301;
PRL 123 (2019) 052001; *PRC* 104 (2021) 024906
CMS, *PRL* 127 (2021) 122001

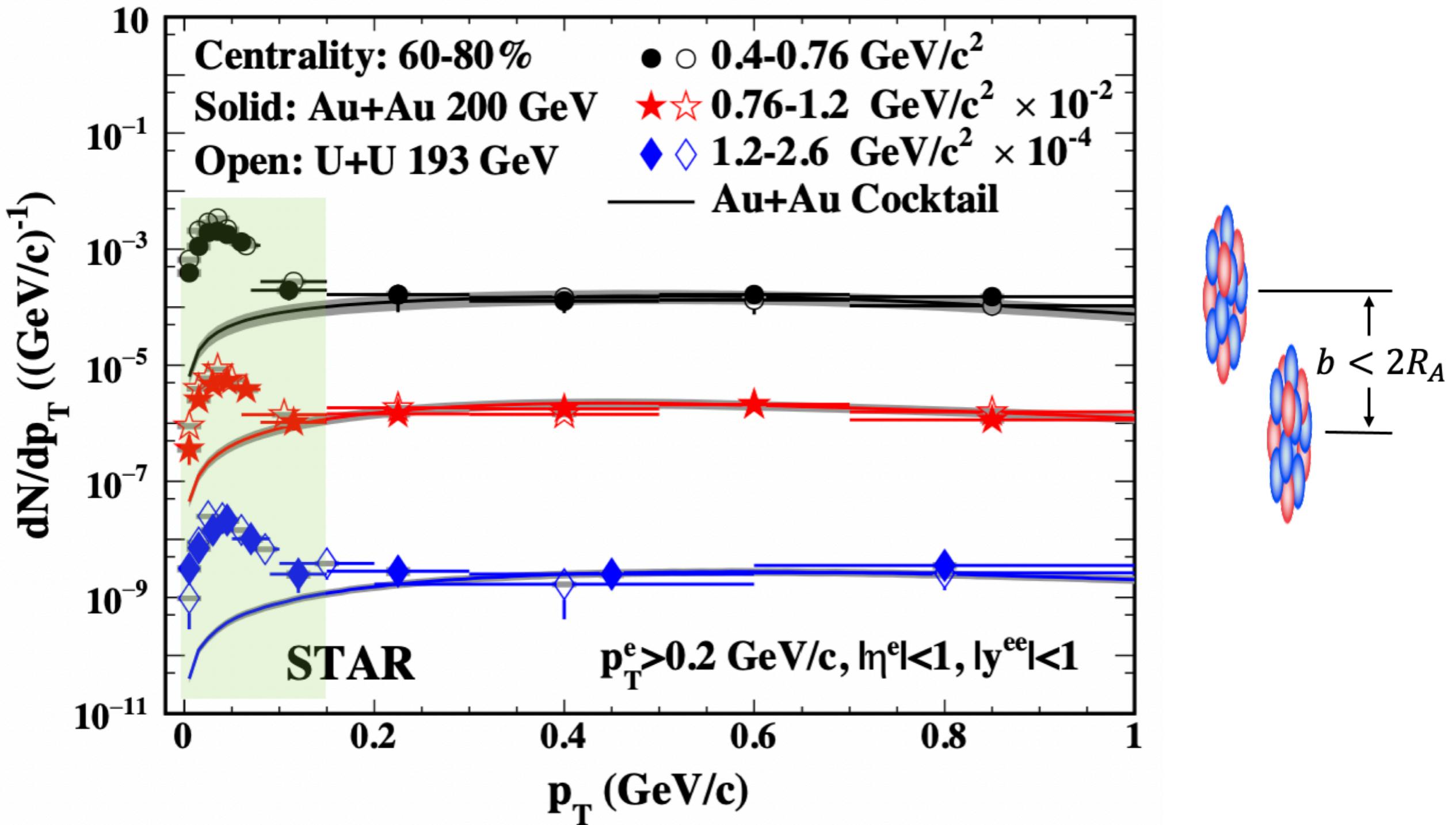
Conventionally studied in UPC



- Distinctive features
 - Exclusive production
 - Concentrated at low p_T (back to back)

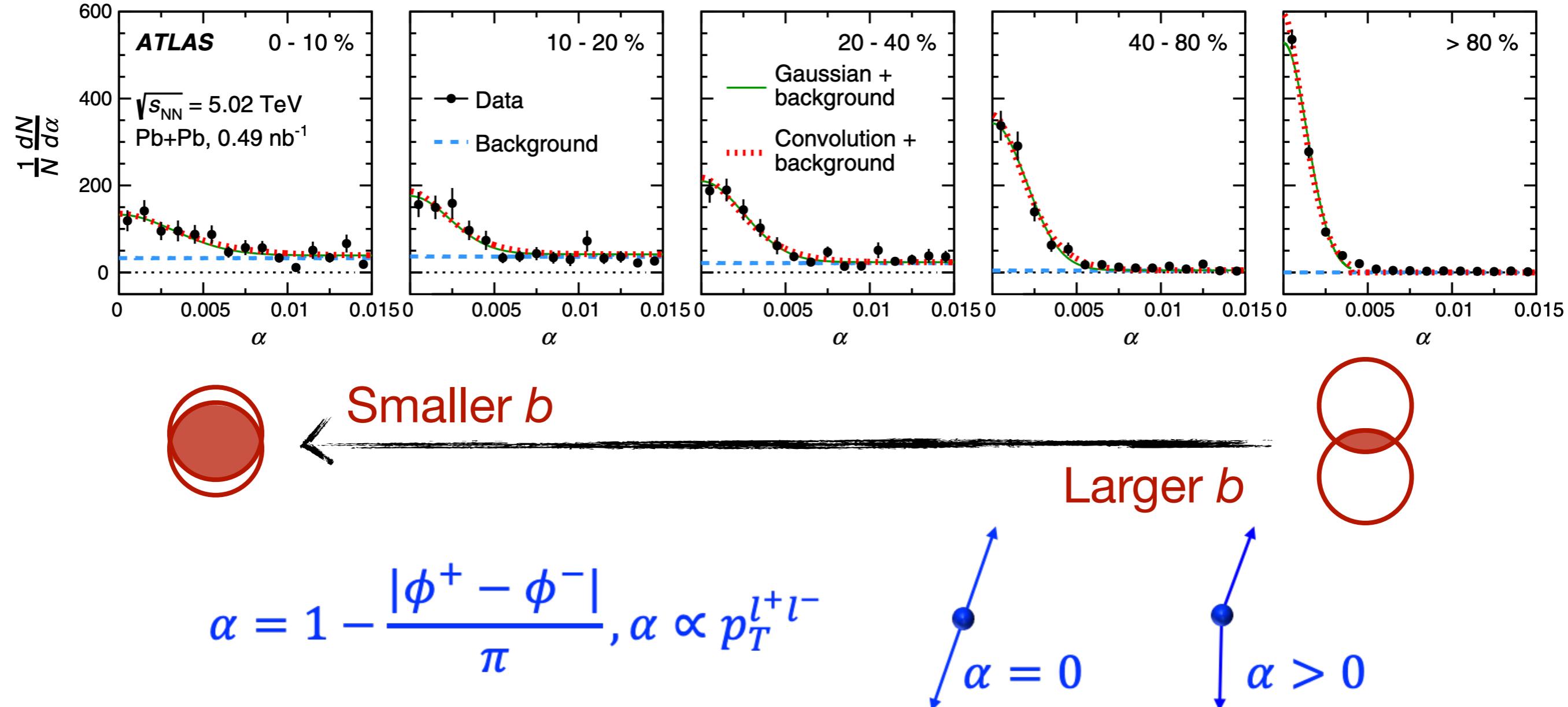
Observed in hadronic collisions

STAR, PRL 121 (2018) 132301



Modification of lepton pairs

ATLAS, PRL 121 (2018) 212301

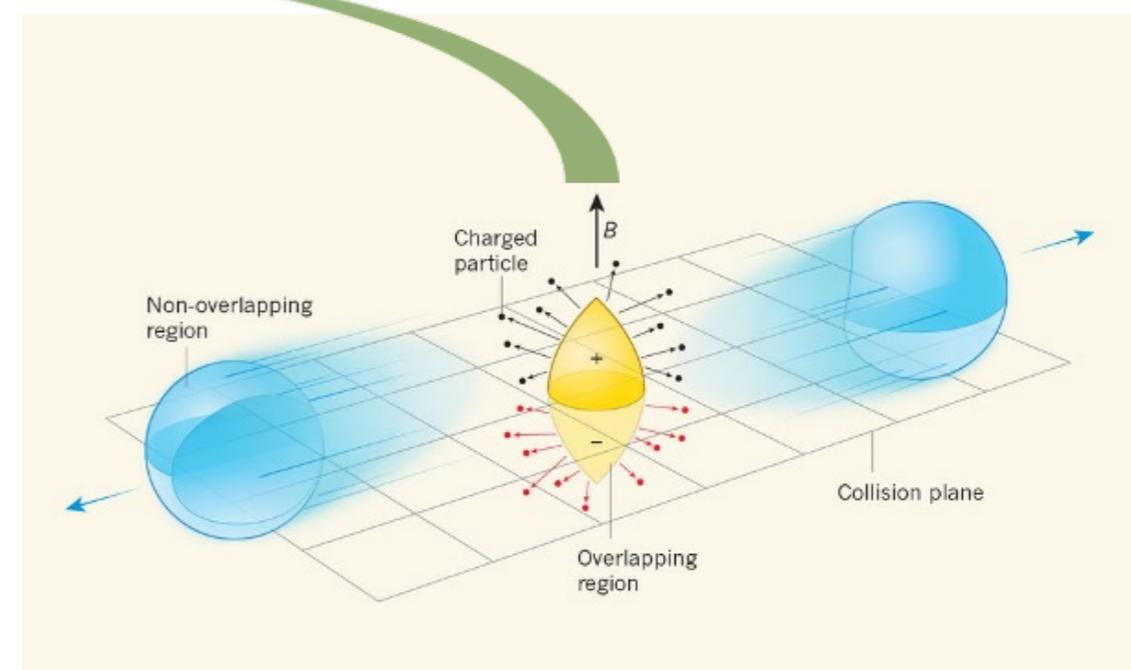
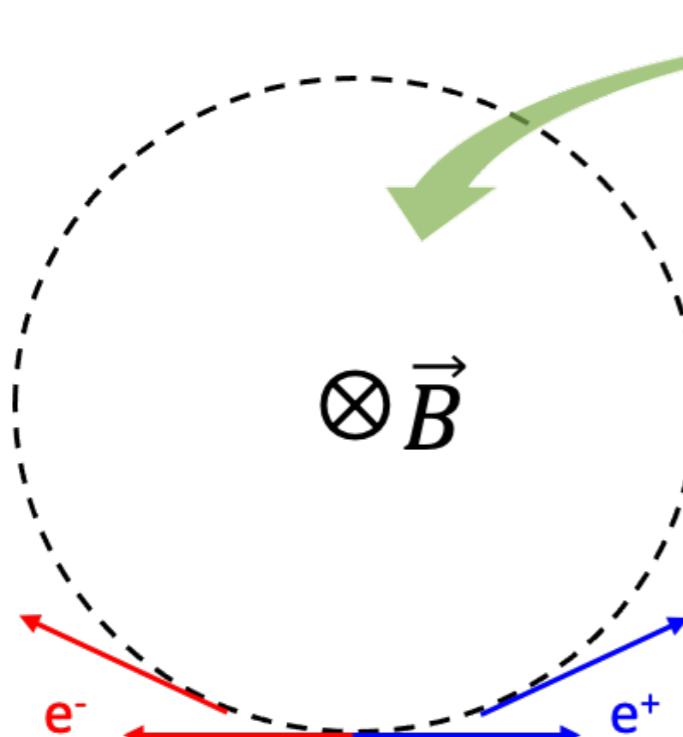
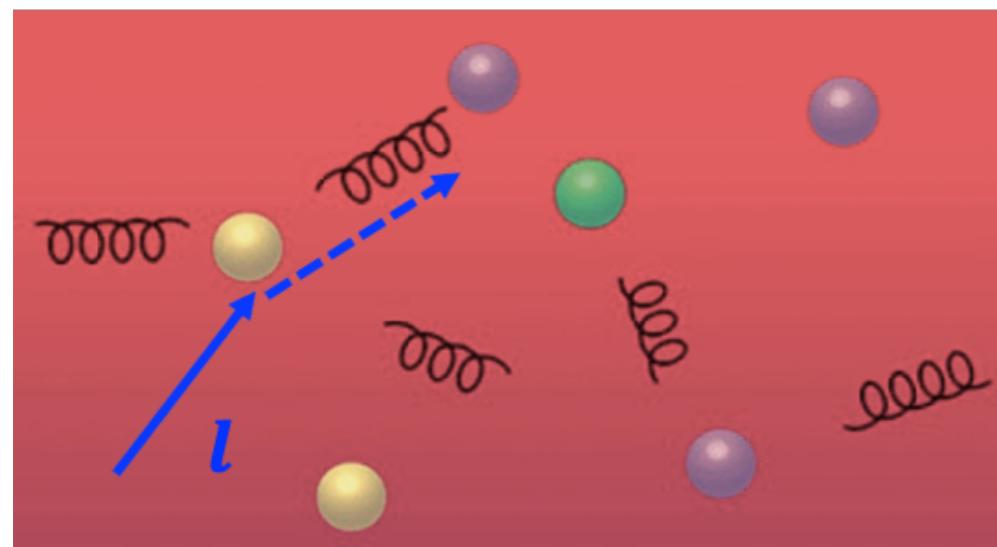


- Back-to-back correlation becomes weaker towards central collisions

Puzzle of the physics origin

STAR, PRL 121 (2018) 132301
ATLAS, PRL 121 (2018) 212301
Klein et al., PRL 122 (2019) 132301

Final-state effect?

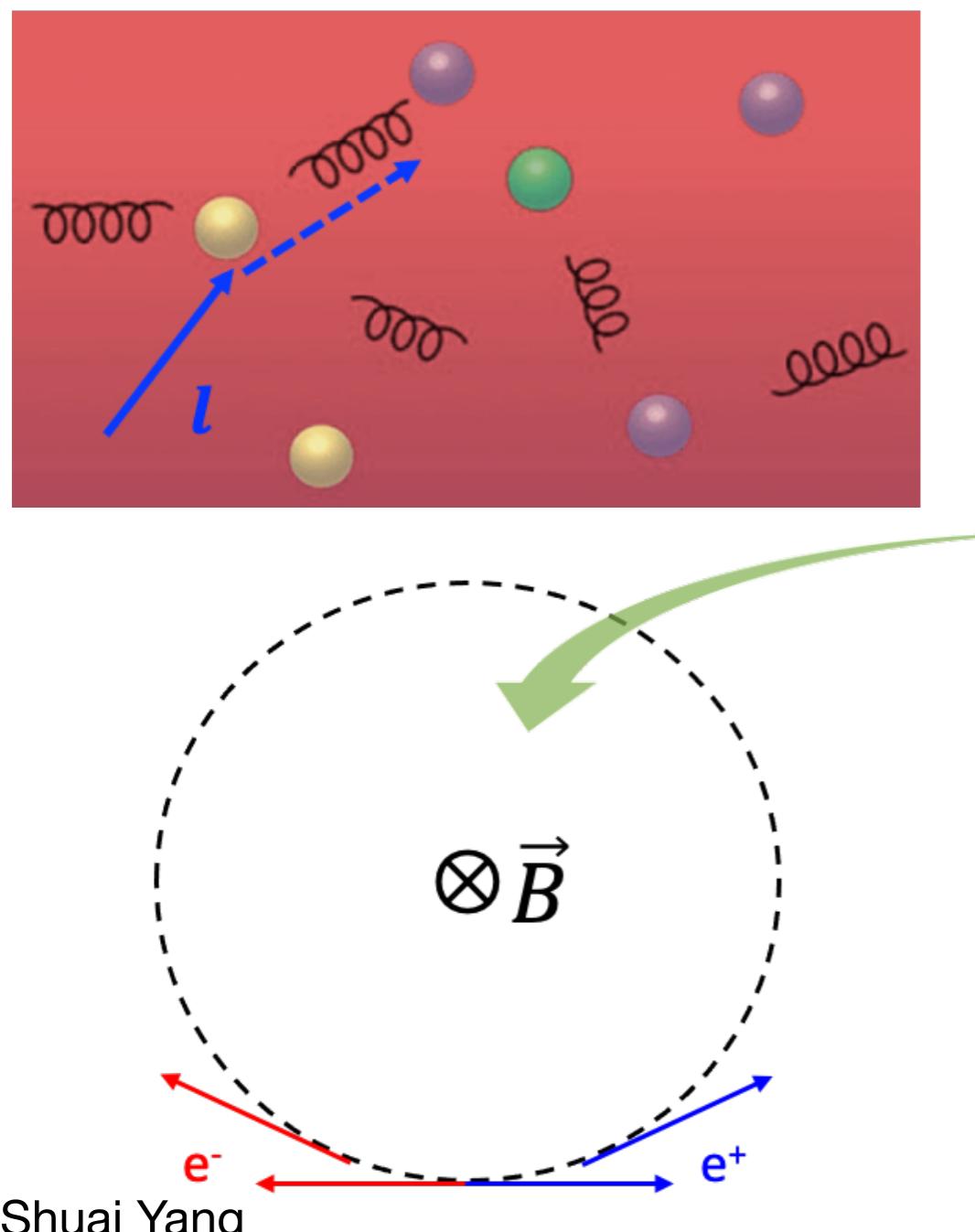


Puzzle of the physics origin

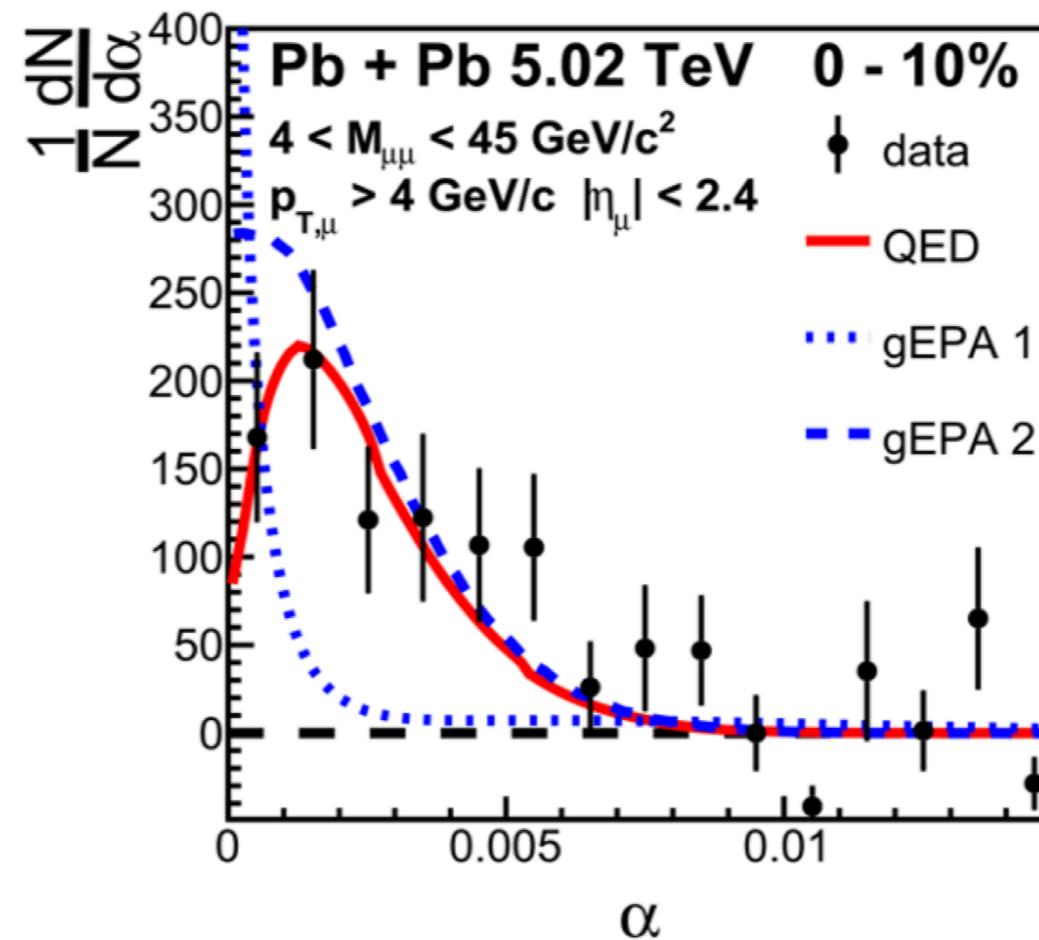
STAR, PRL 121 (2018) 132301
ATLAS, PRL 121 (2018) 212301
Klein et al., PRL 122 (2019) 132301

Zha et al., PLB 800 (2020) 135089
Li et al., PRD 101 (2020) 034015
Wang et al., PRD 104 (2021) 056011

Final-state effect?

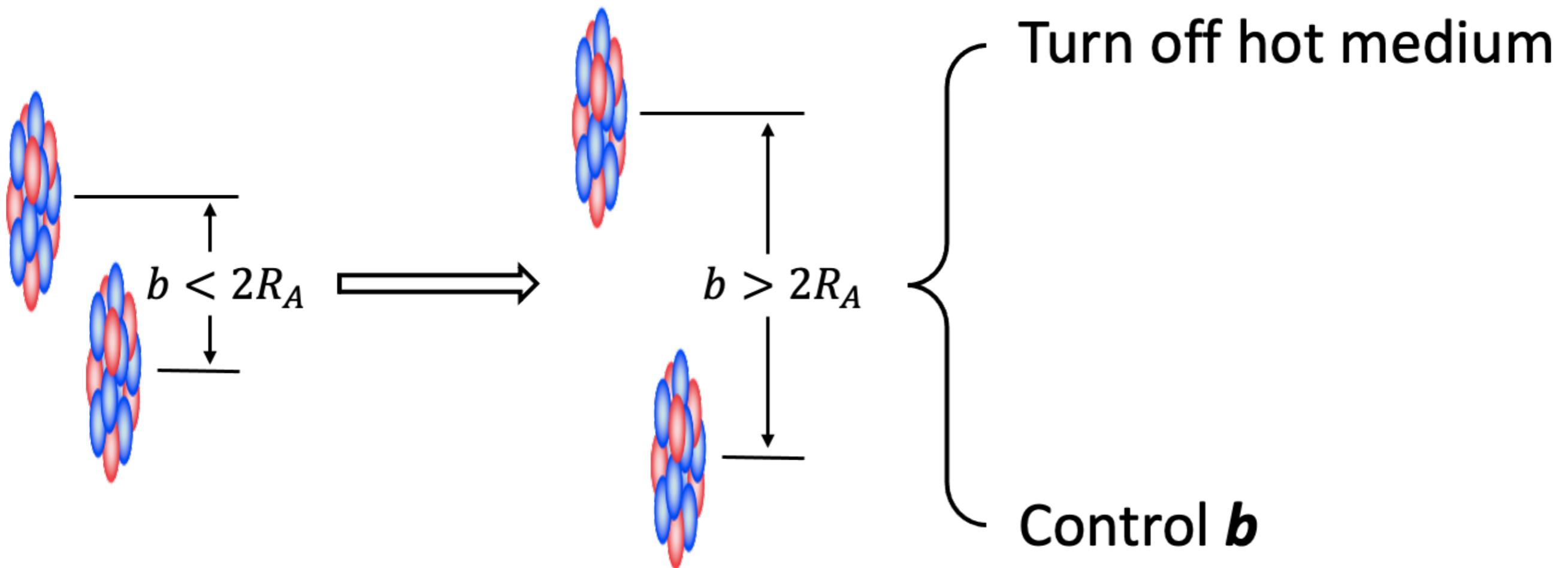


Initial-state effect?



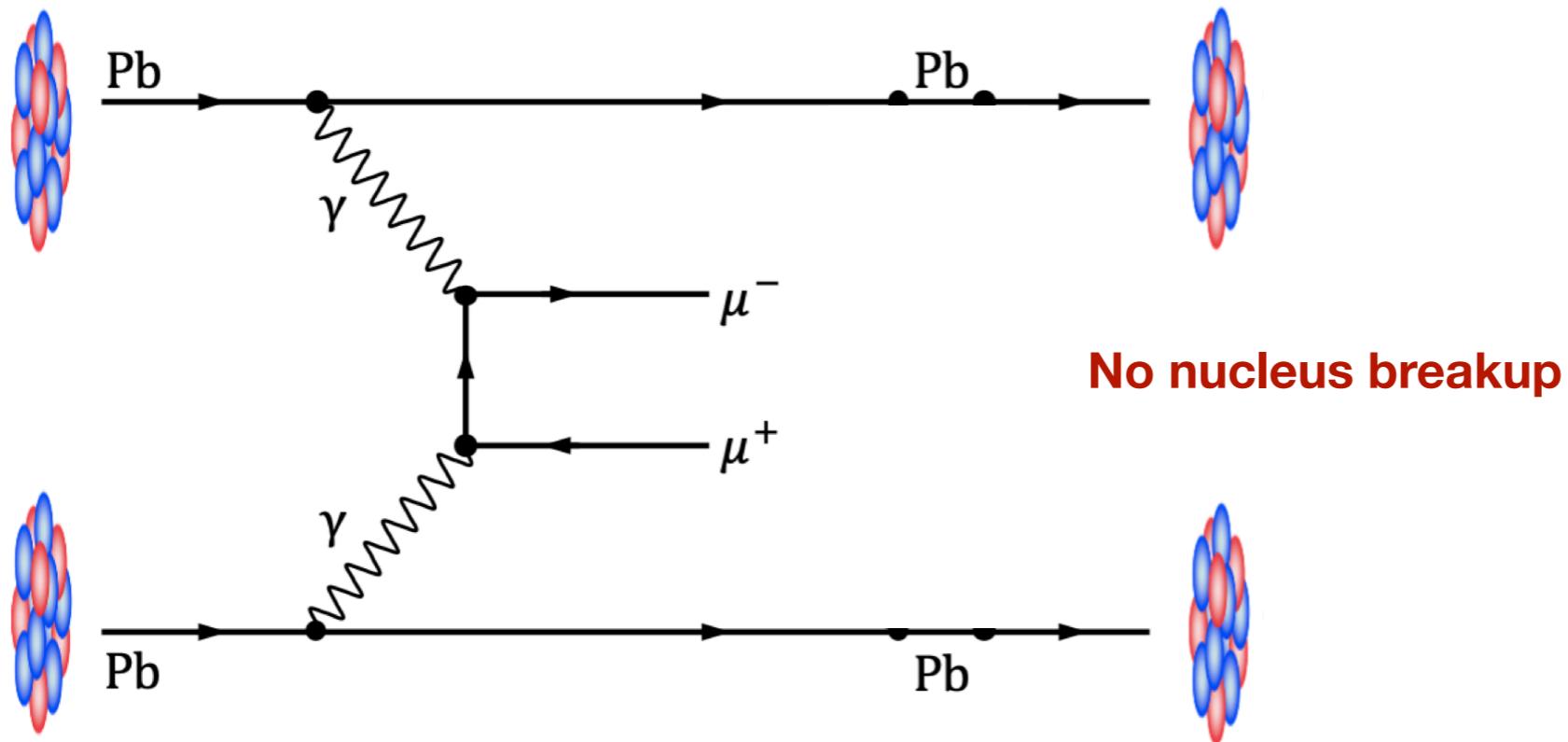
- Described by lowest-order QED without medium effect
 - **b** dependence of initial photon p_T

Experimentally explore the puzzle



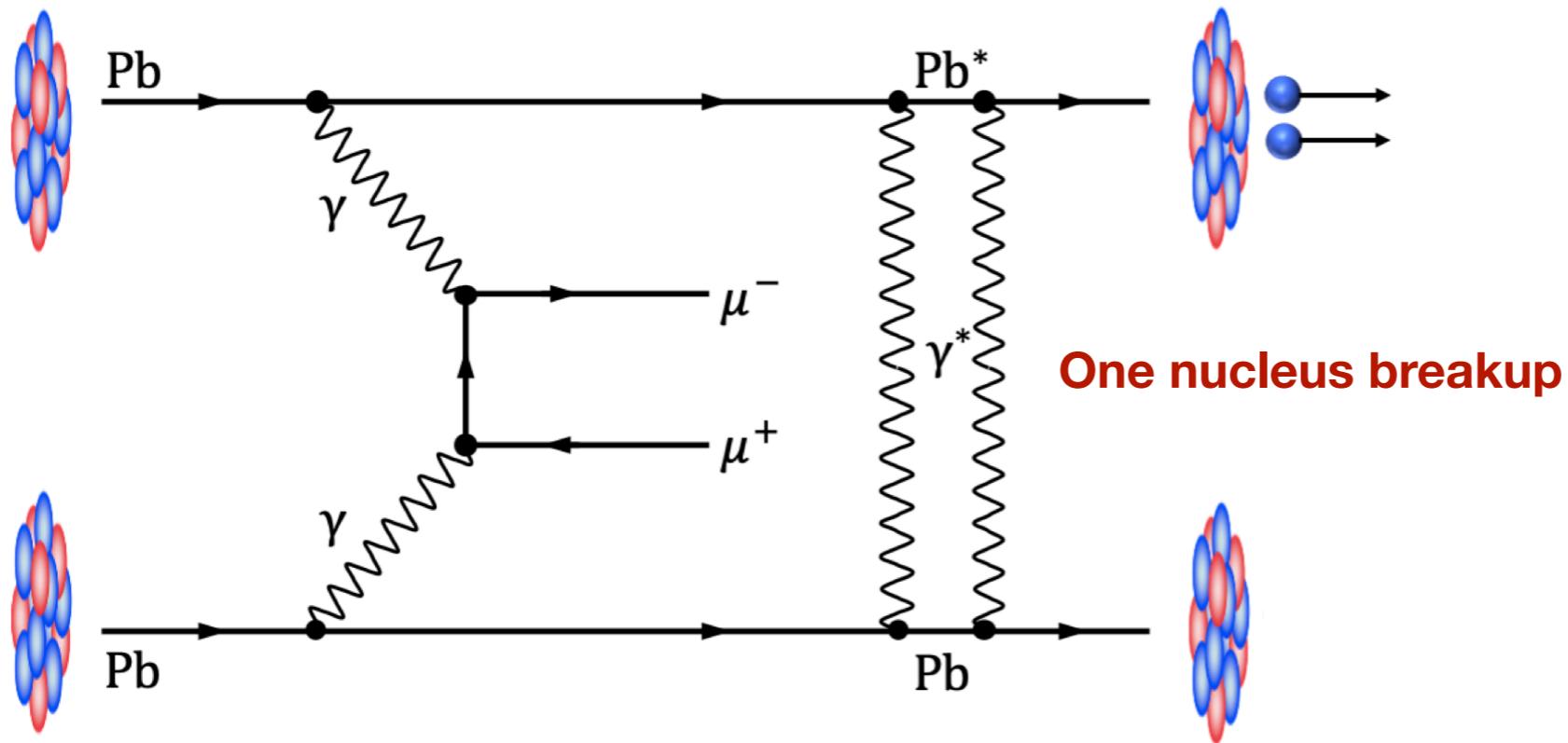
Control b in UPC

Nuclei **may** exchange soft photon(s) \Rightarrow nuclear dissociation



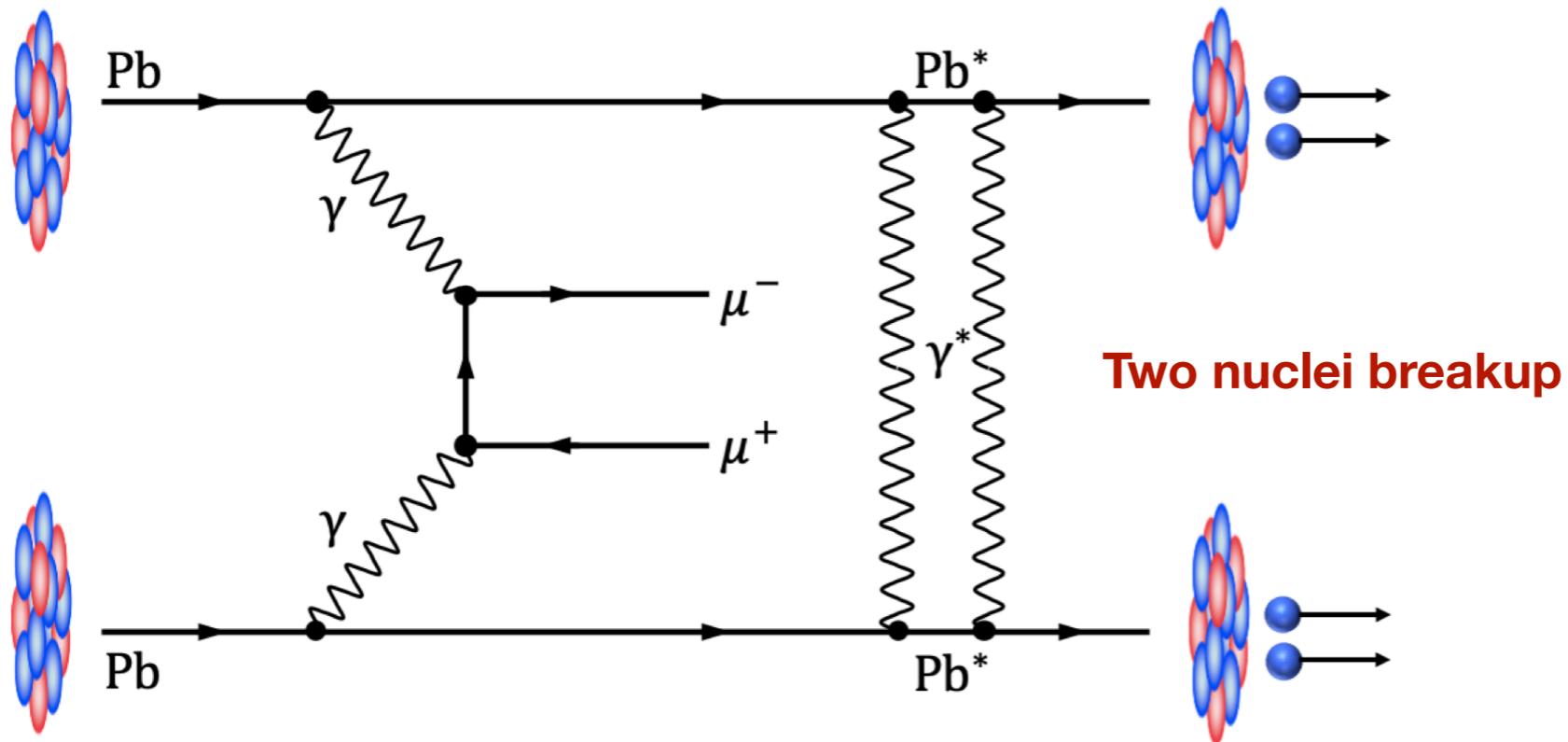
Control b in UPC

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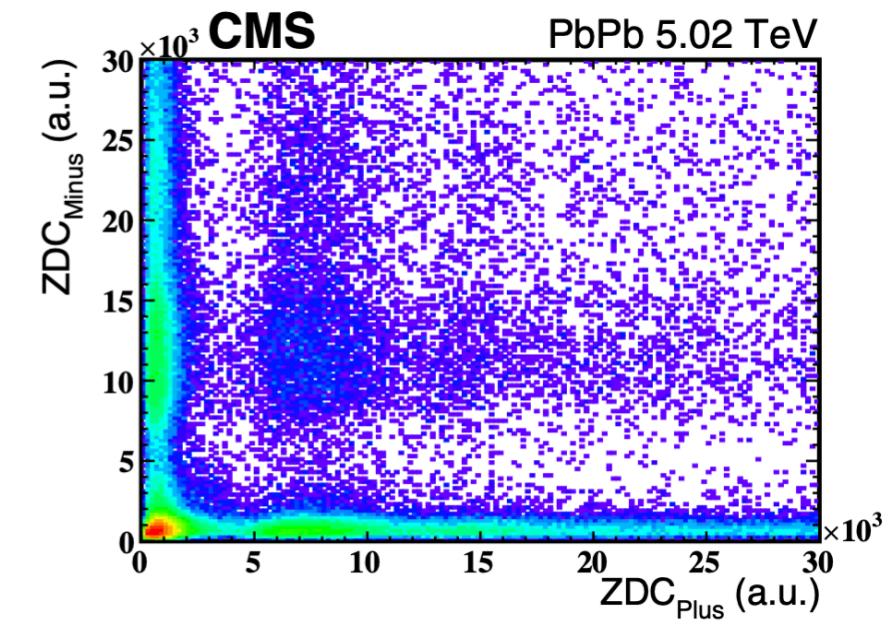
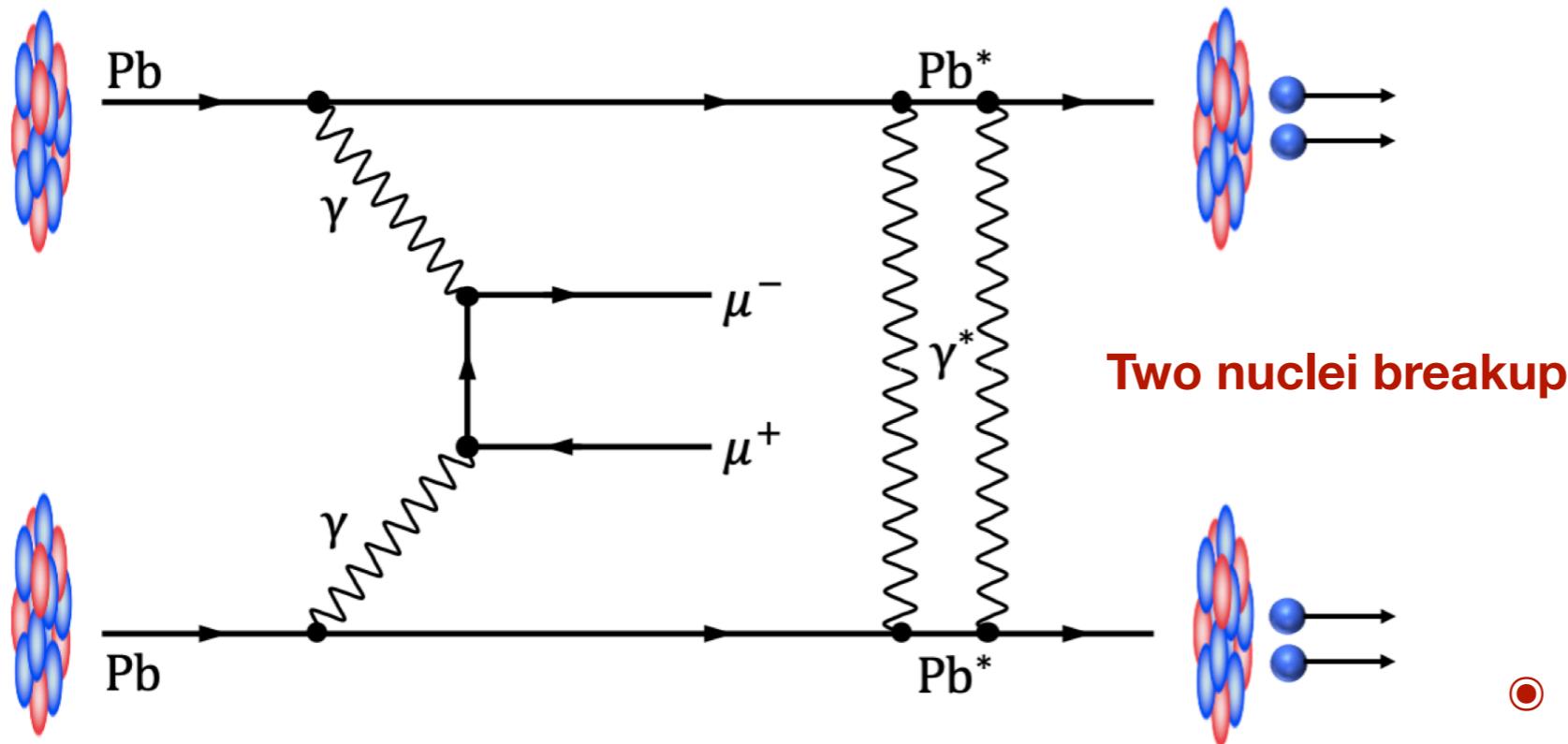
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Control b in UPC

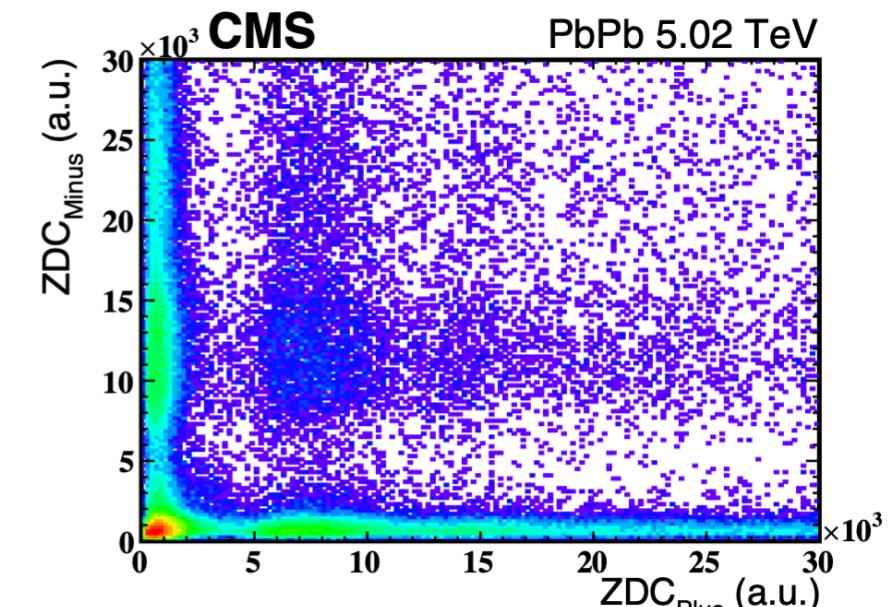
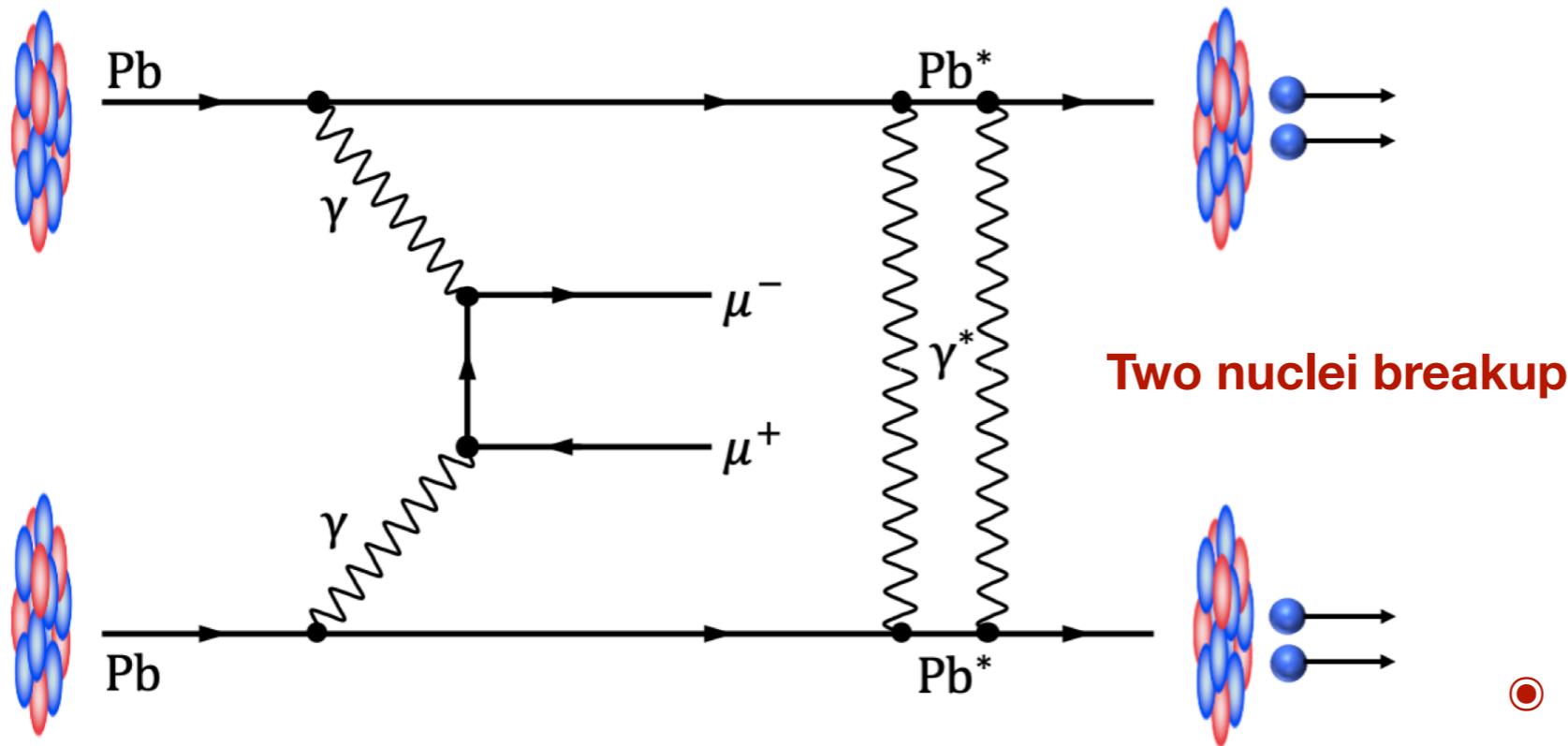
Nuclei **may** exchange soft photon(s) \Rightarrow nuclear dissociation



● Zero Degree Calorimeter

Control b in UPC

Nuclei **may** exchange soft photon(s) \Rightarrow nuclear dissociation

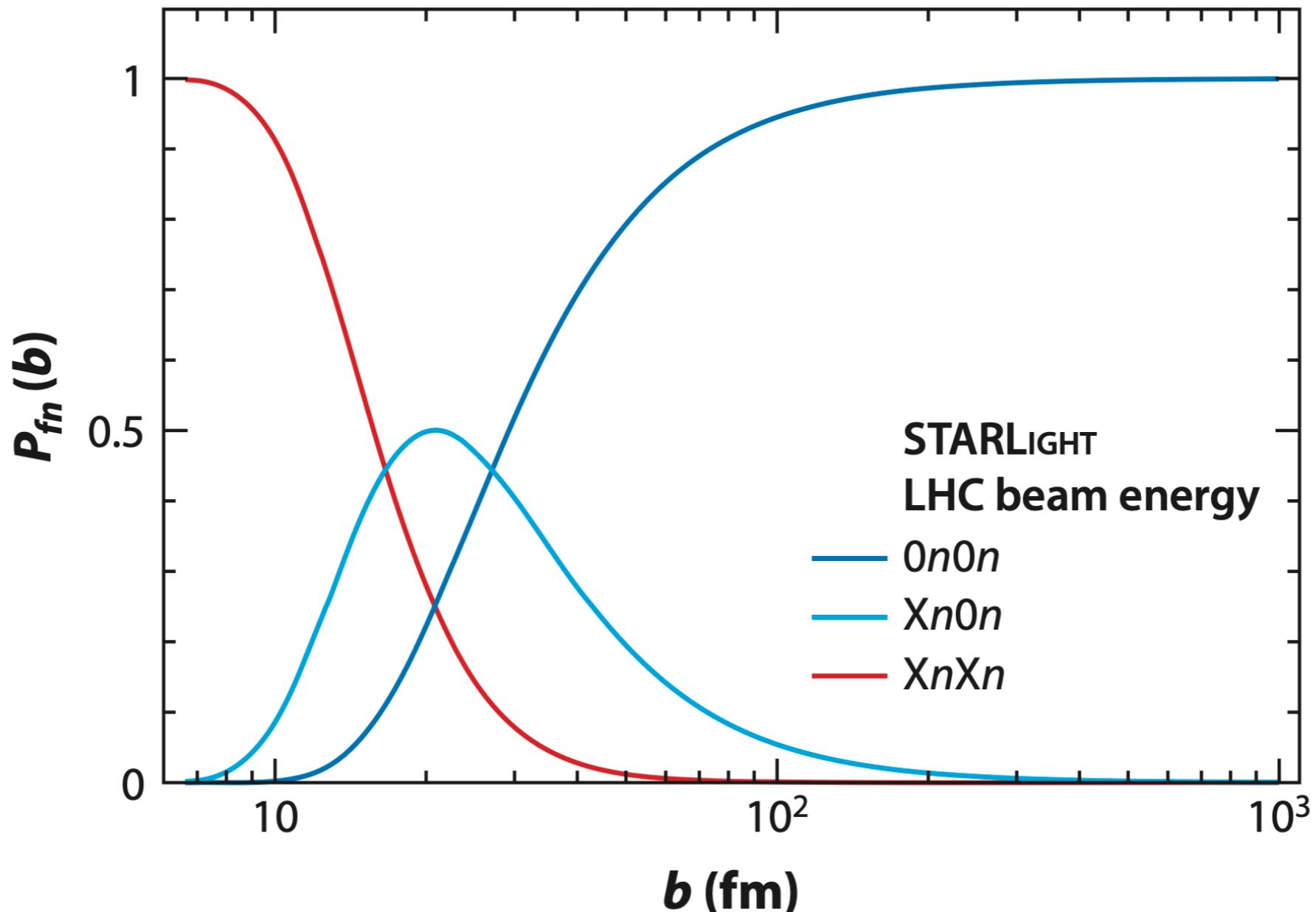


○ Zero Degree Calorimeter

$$N(k) = \int d^2 b N(k, b) P_{0\text{had}}(b) P_1(b) P_2(b)$$

, where $P_i(b) \propto 1/b^2$

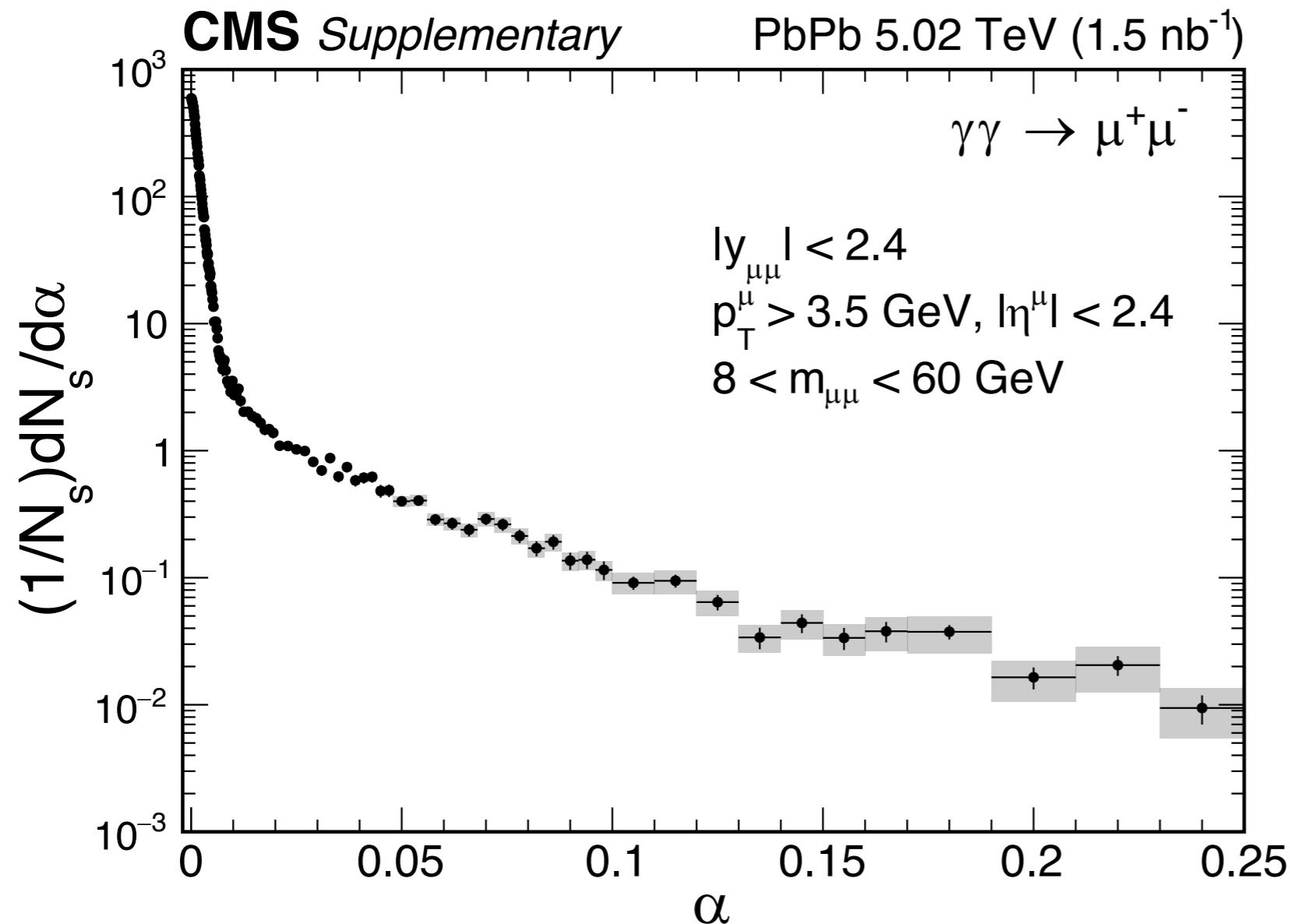
Control b in UPC



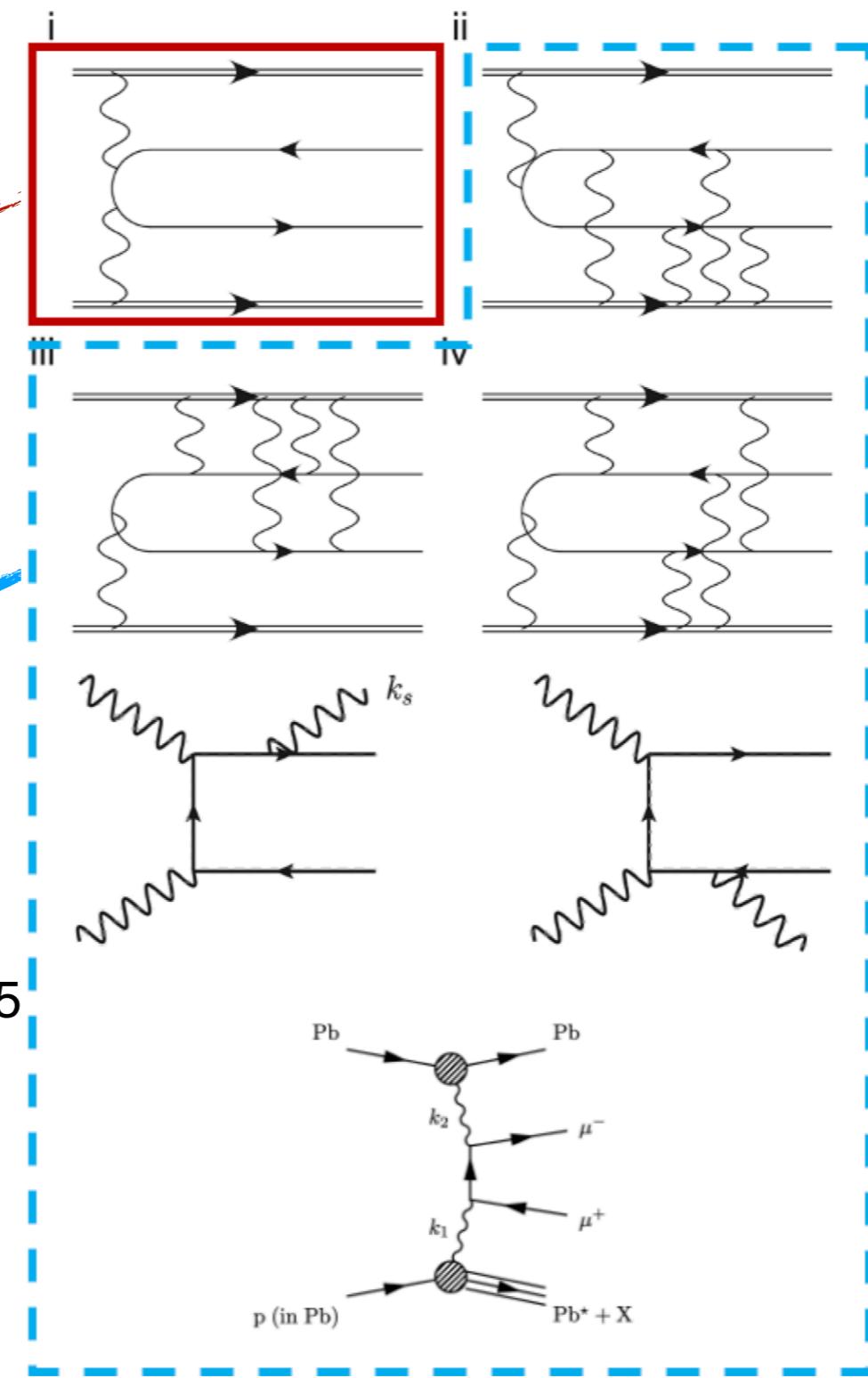
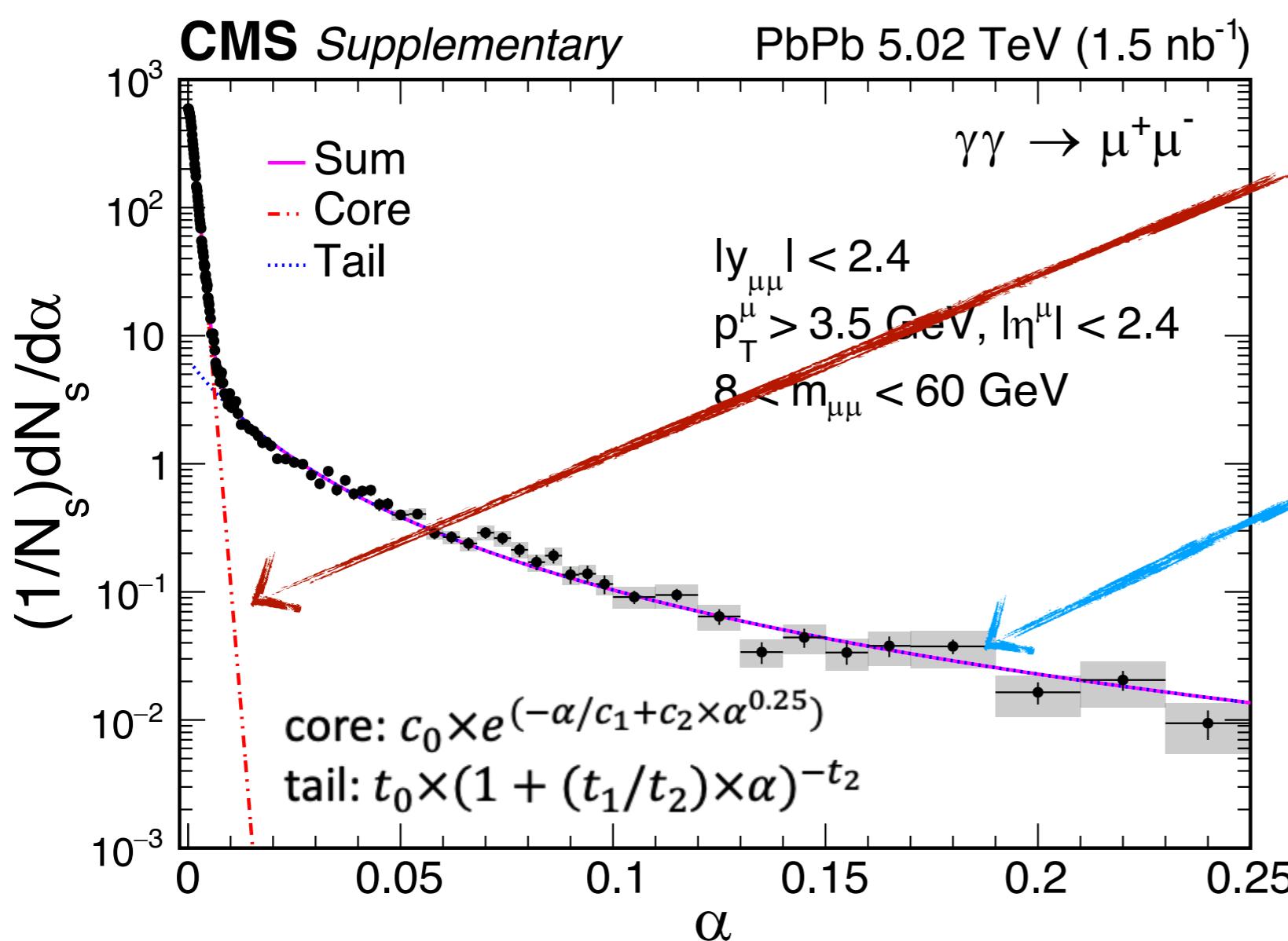
Klein and Steinberg, Ann. Rev. Nucl. Part. Sci. 70 (2020) 323

$$\mathbf{b}_{XnXn} < \mathbf{b}_{0nXn} < \mathbf{b}_{0n0n}$$

α spectrum in UPC

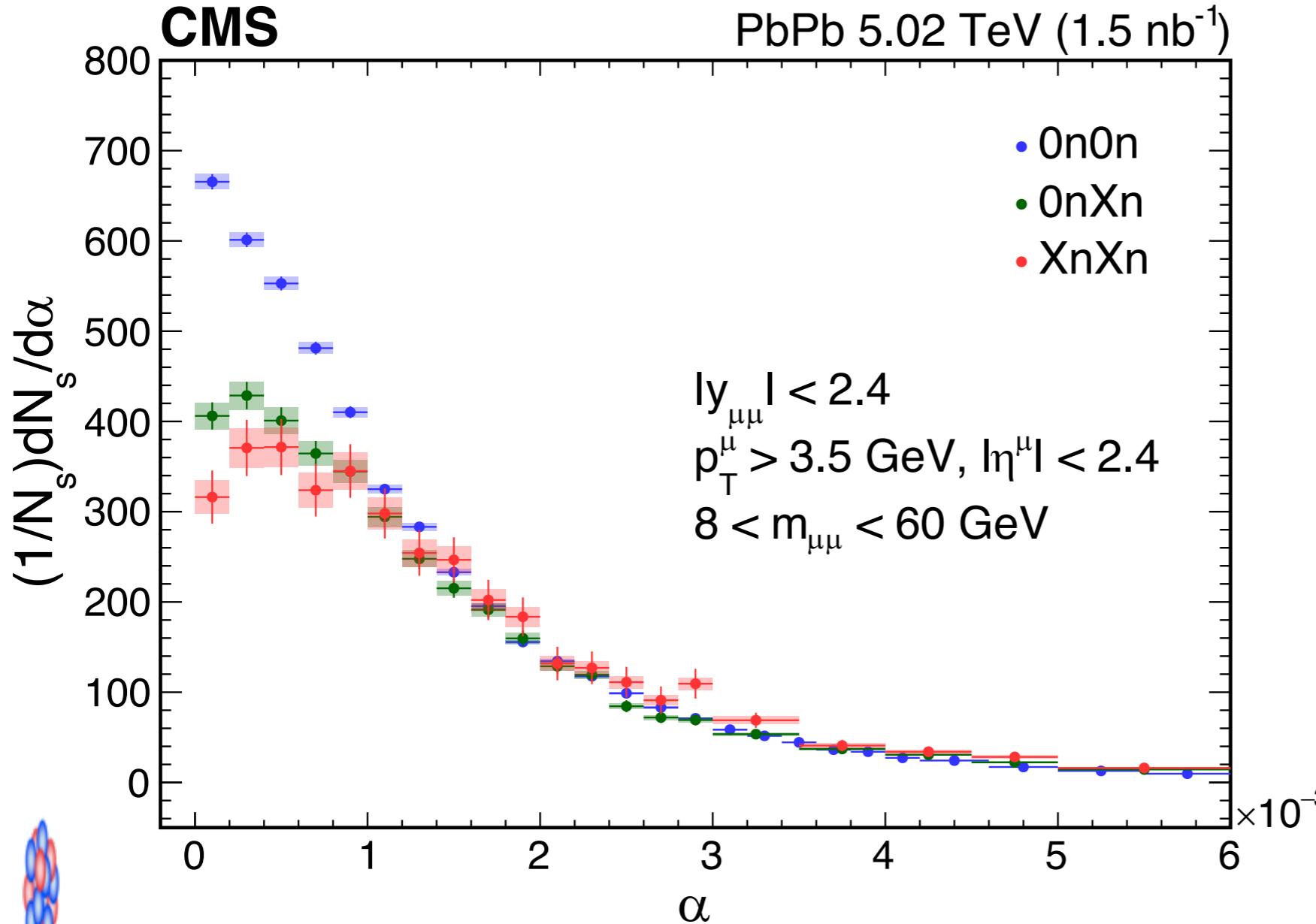


α spectrum in UPC



Decouple α spectrum with empirical function

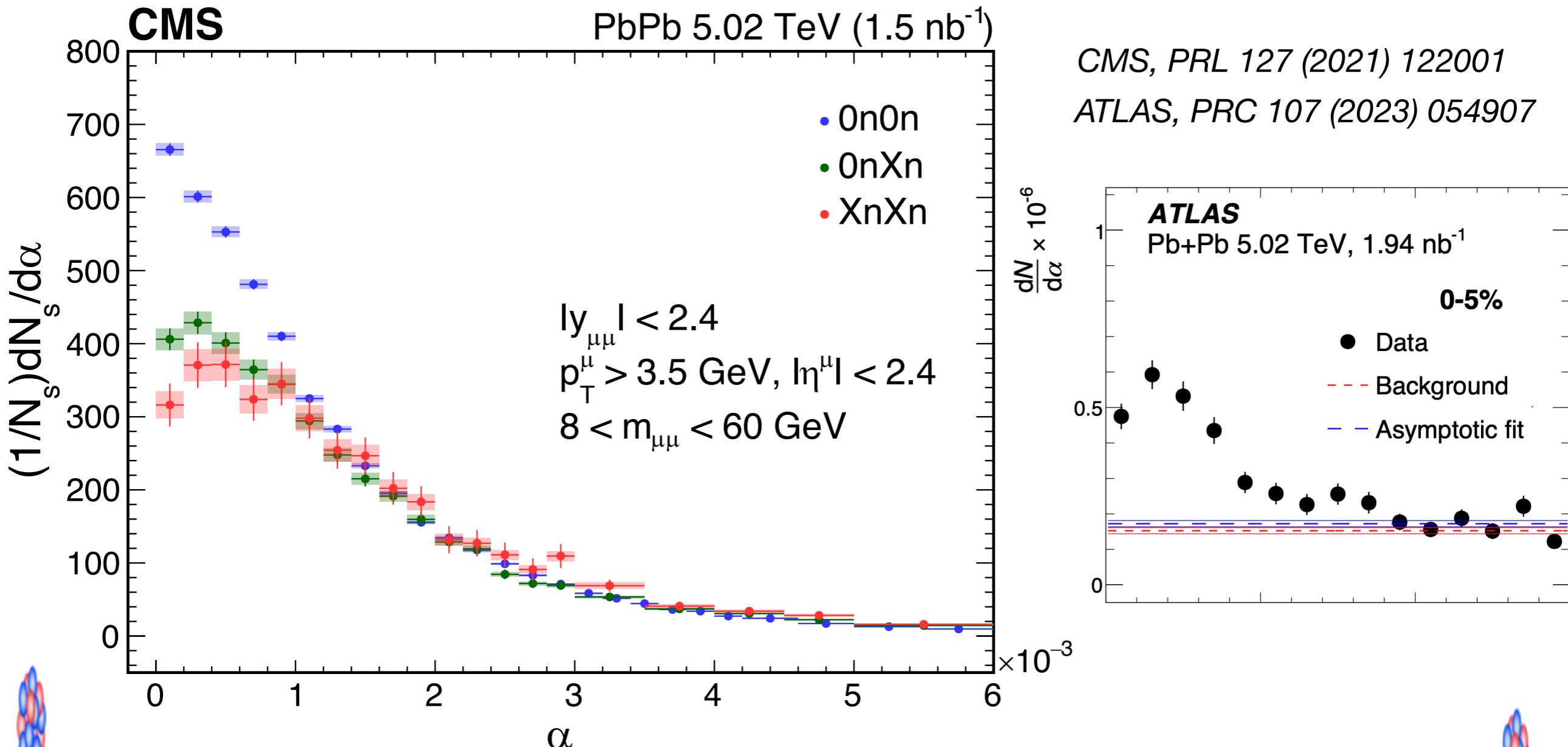
α spectrum vs. neutron multiplicity



CMS, PRL 127 (2021) 122001

- 0n0n (fewer neutrons) \Rightarrow XnXn (more neutrons)
 - α spectrum becomes broader

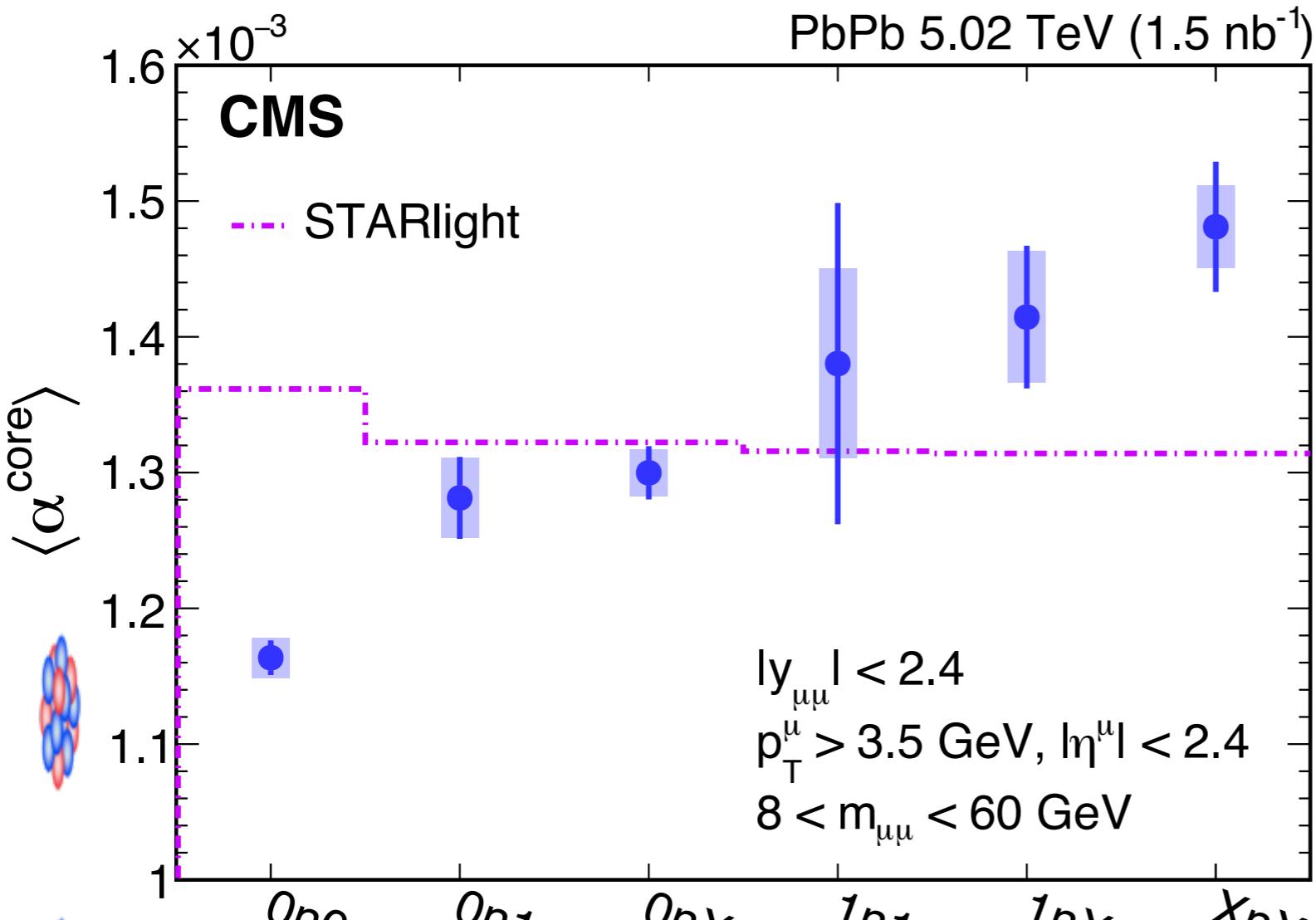
α spectrum vs. neutron multiplicity



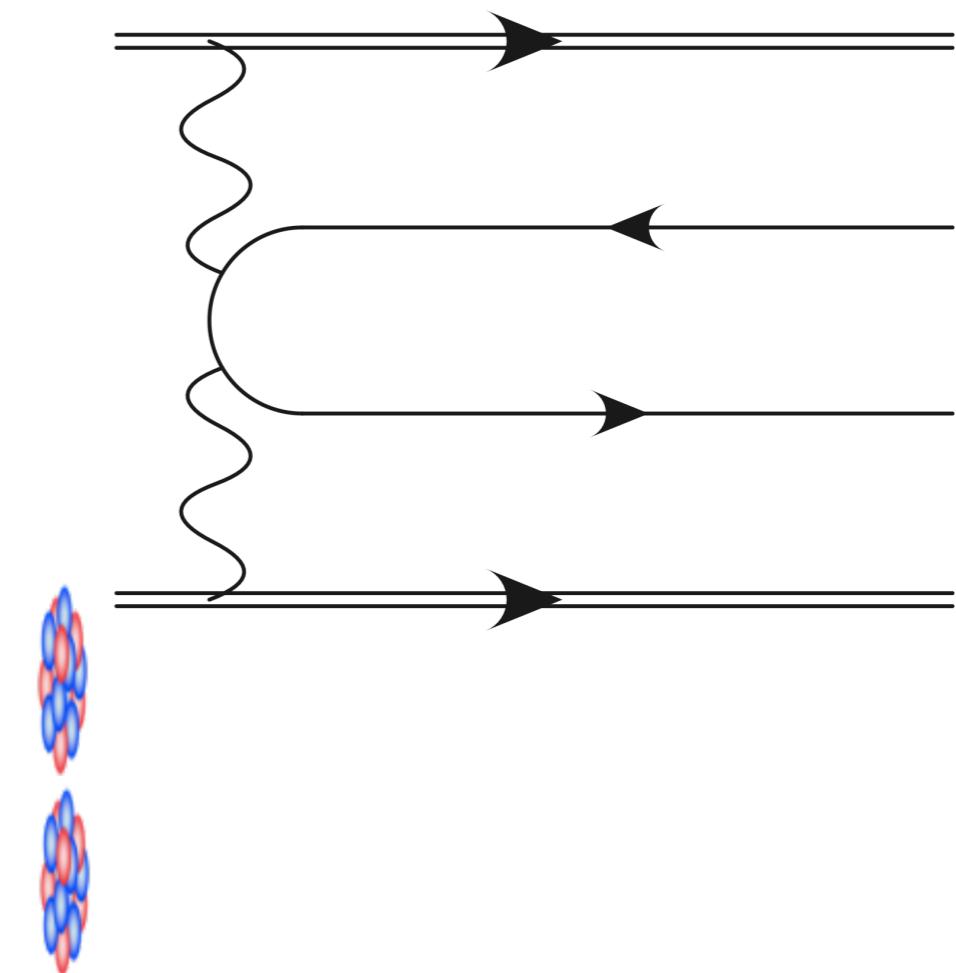
● 0n0n (fewer neutrons) \Rightarrow XnXn (more neutrons)

- α spectrum becomes broader
- Seems has depletion in the very small α

$\langle \alpha^{\text{core}} \rangle$ vs. neutron multiplicity

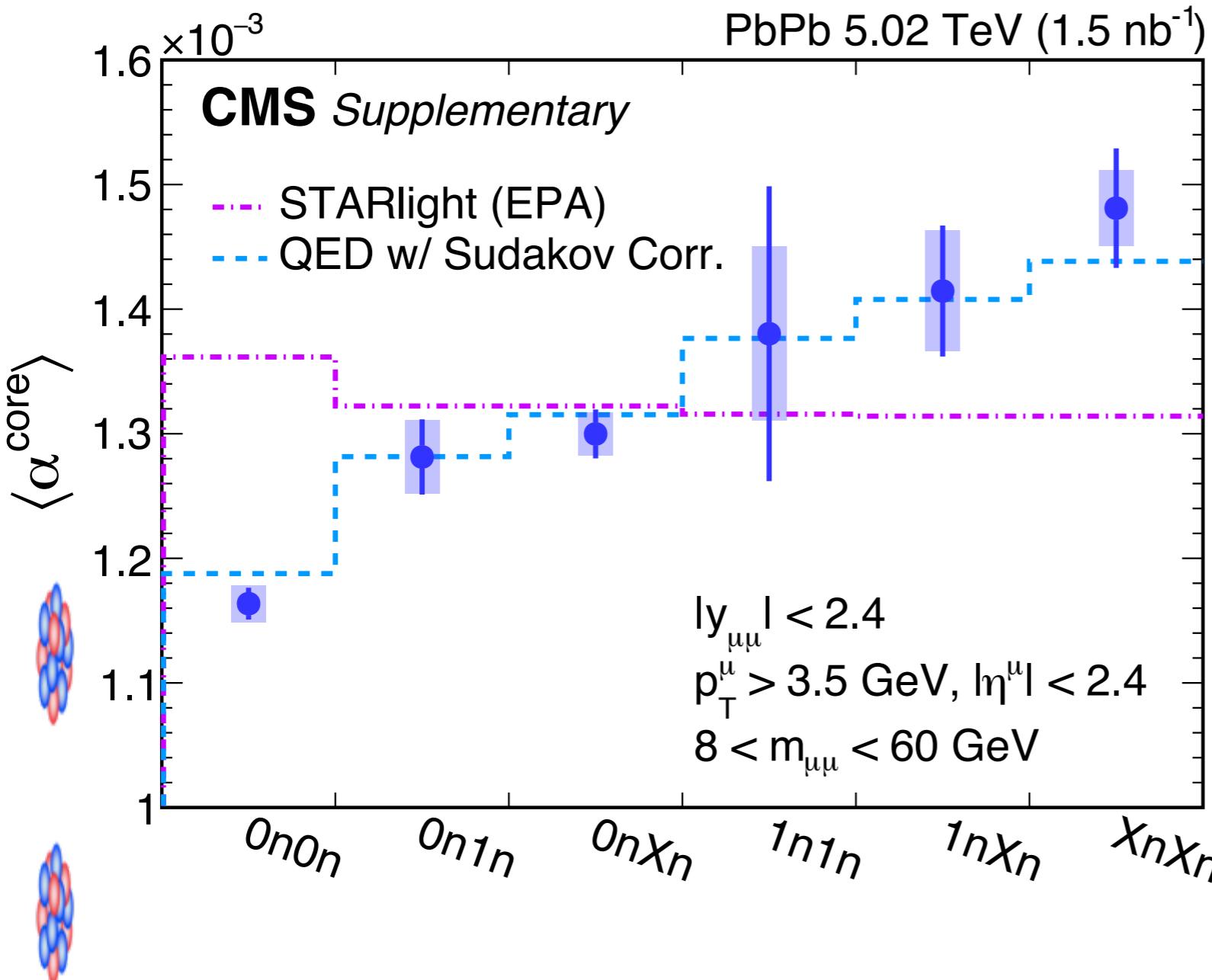


Klein et al., Comput. Phys. Commun. 212 (2017) 258
Klein et al., PRL 122 (2019) 132301
Brandenburg et al., arXiv:2006.07365

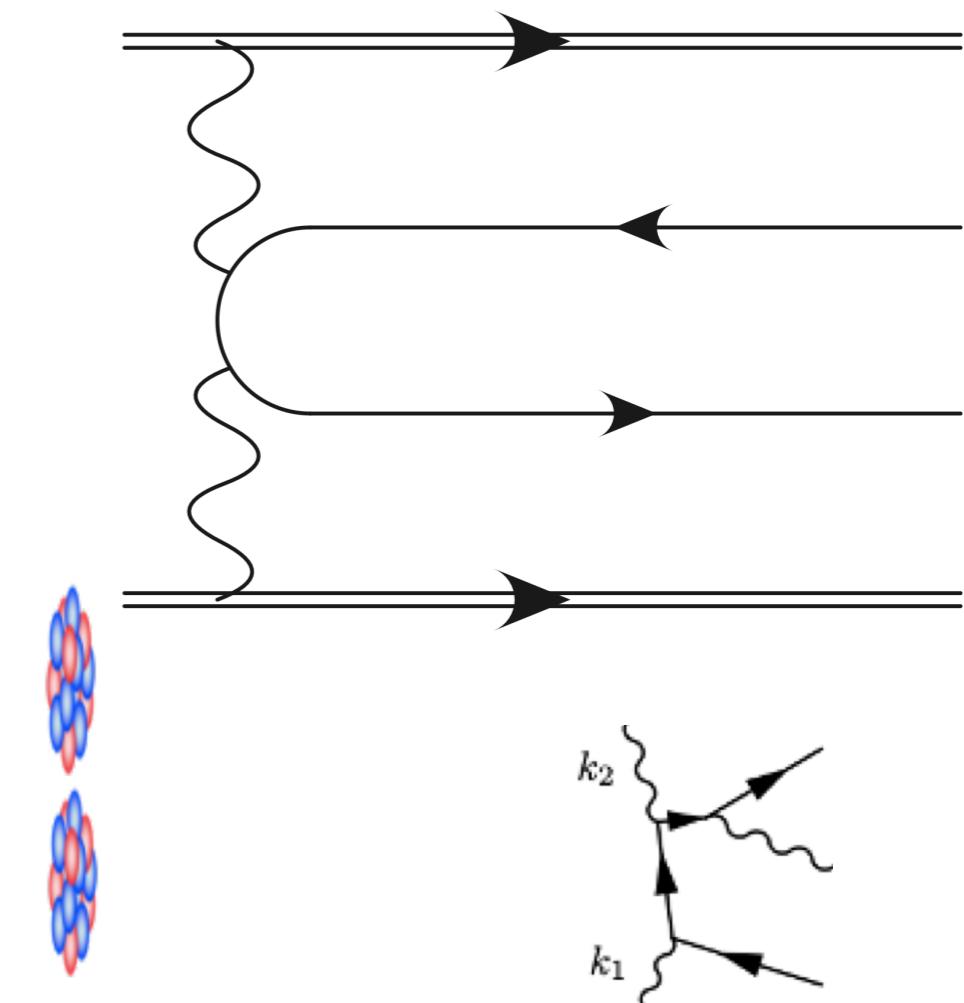


- Strong neutron multiplicity dependence of $\langle \alpha^{\text{core}} \rangle$
 - b** dependence of initial photon p_T

$\langle \alpha^{\text{core}} \rangle$ vs. neutron multiplicity

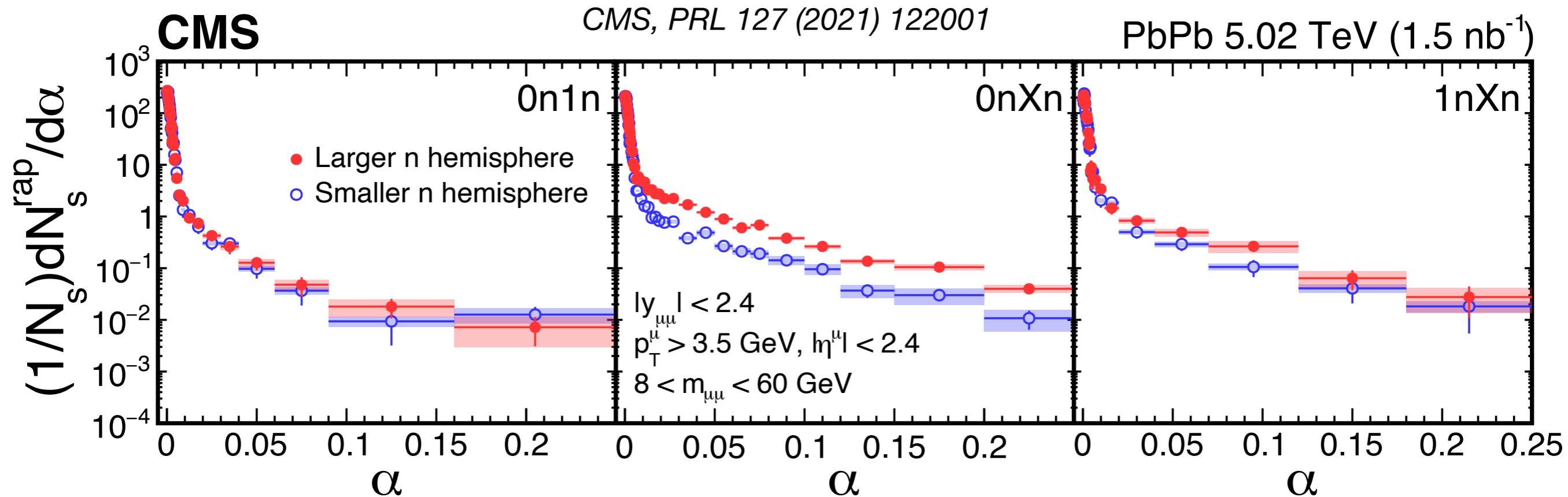


Klein et al., Comput. Phys. Commun. 212 (2017) 258
Klein et al., PRL 122 (2019) 132301
Brandenburg et al., arXiv:2006.07365



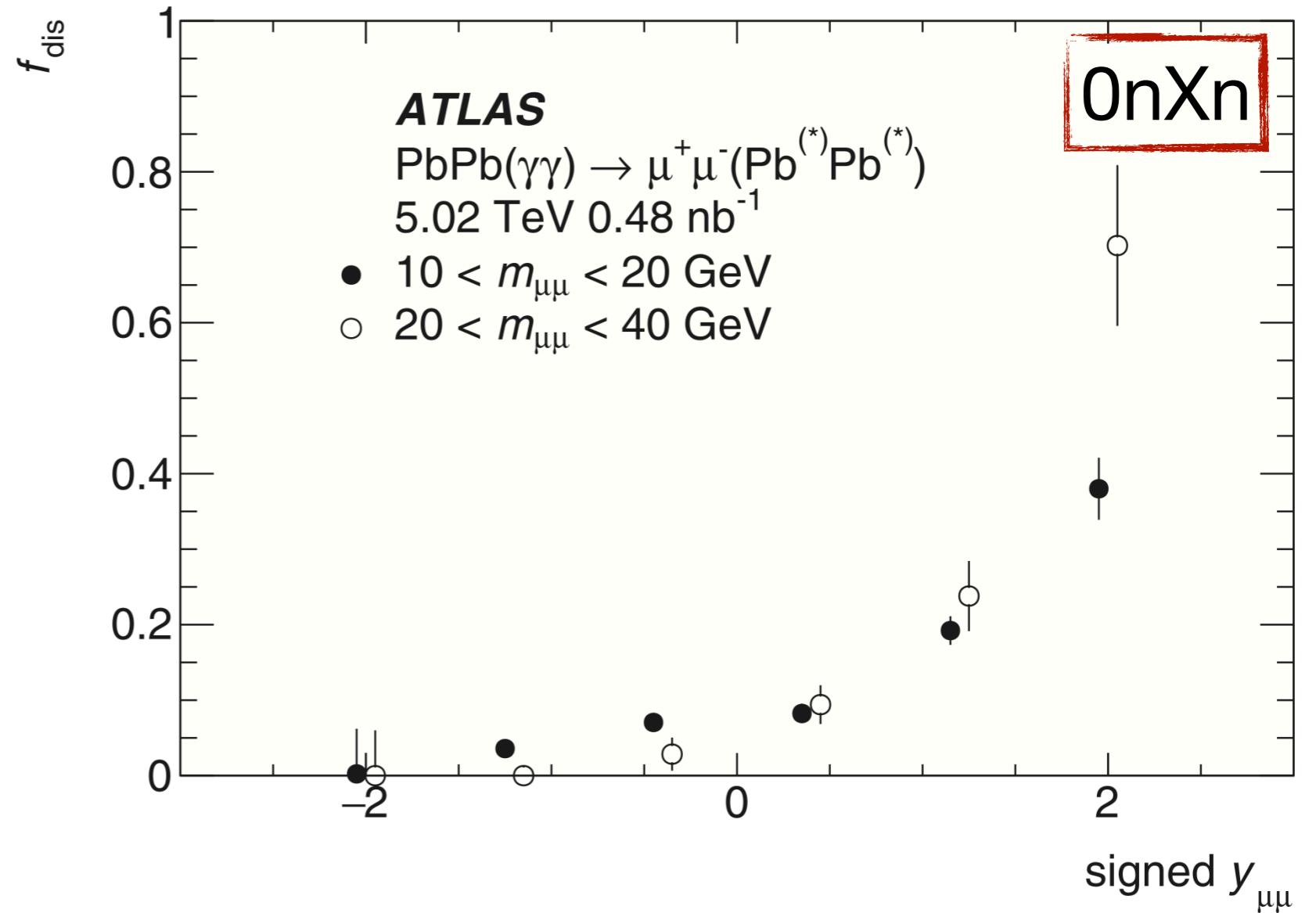
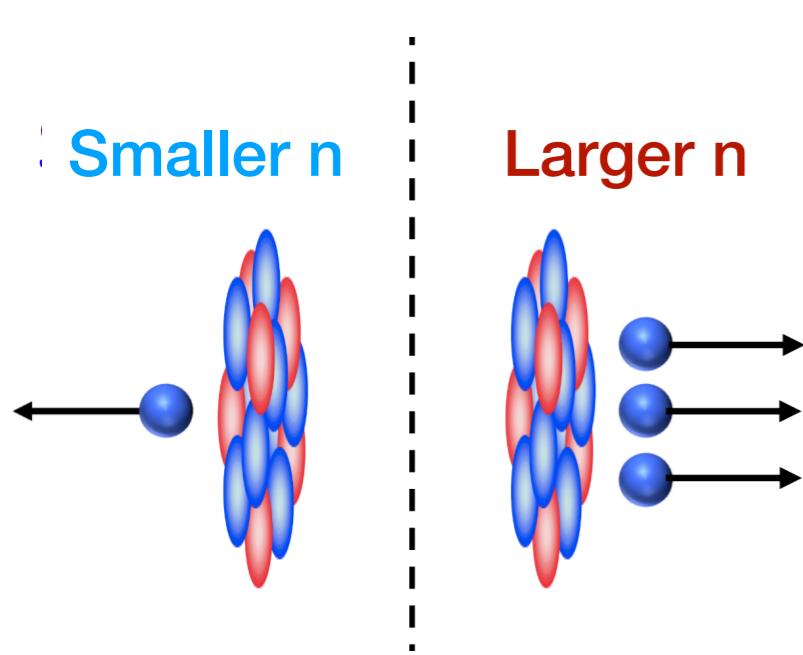
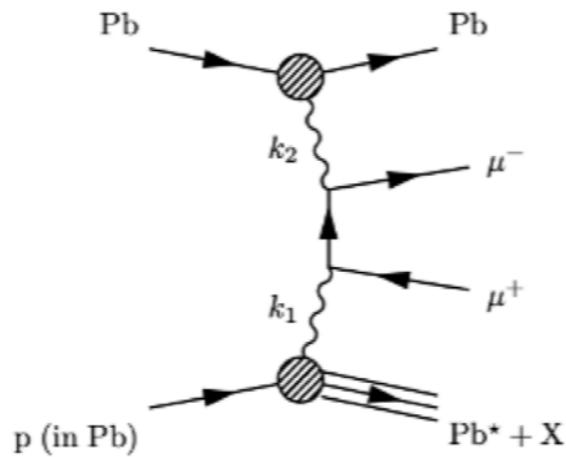
- Qualitatively described by a leading order QED model

Rapidity dependence of α spectra



Rapidity dependence of α spectra

ATLAS, PRC 104 (2021) 024906

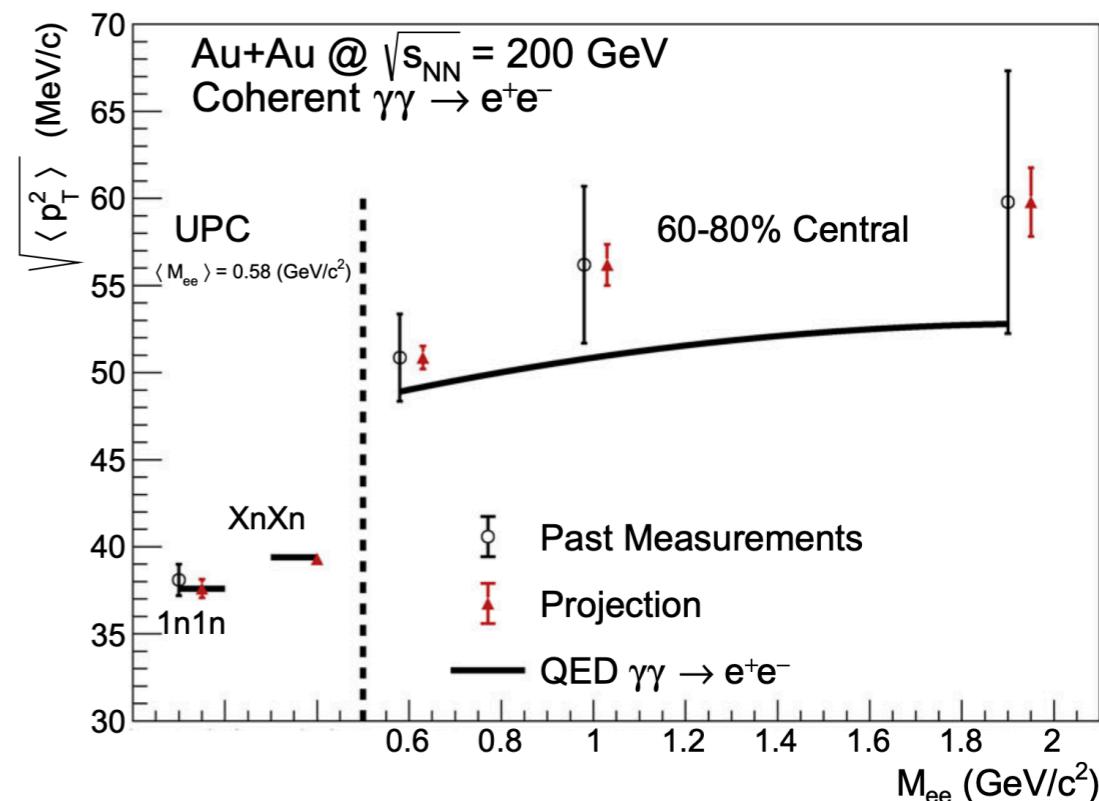


- Similar observation by ATLAS

Impact to explore QGP EM properties

- The b dependence of photon p_T should be considered to explore QGP EM properties
 - RHIC run 2023-2025
 - LHC run3 & 4

Brandenburg et al., EPJA 57 (2021) 299



Impact to explore QGP EM properties

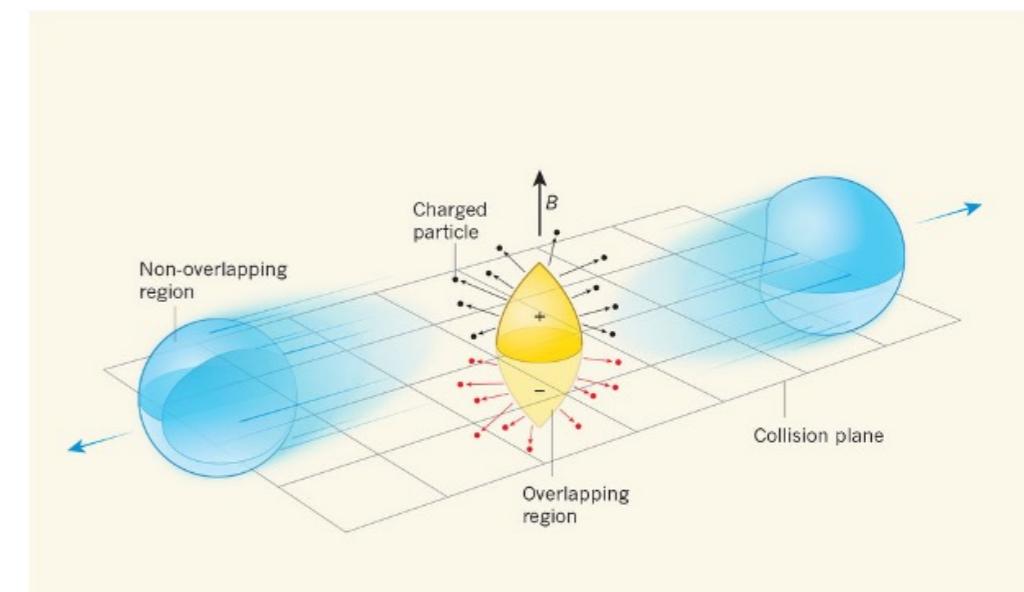
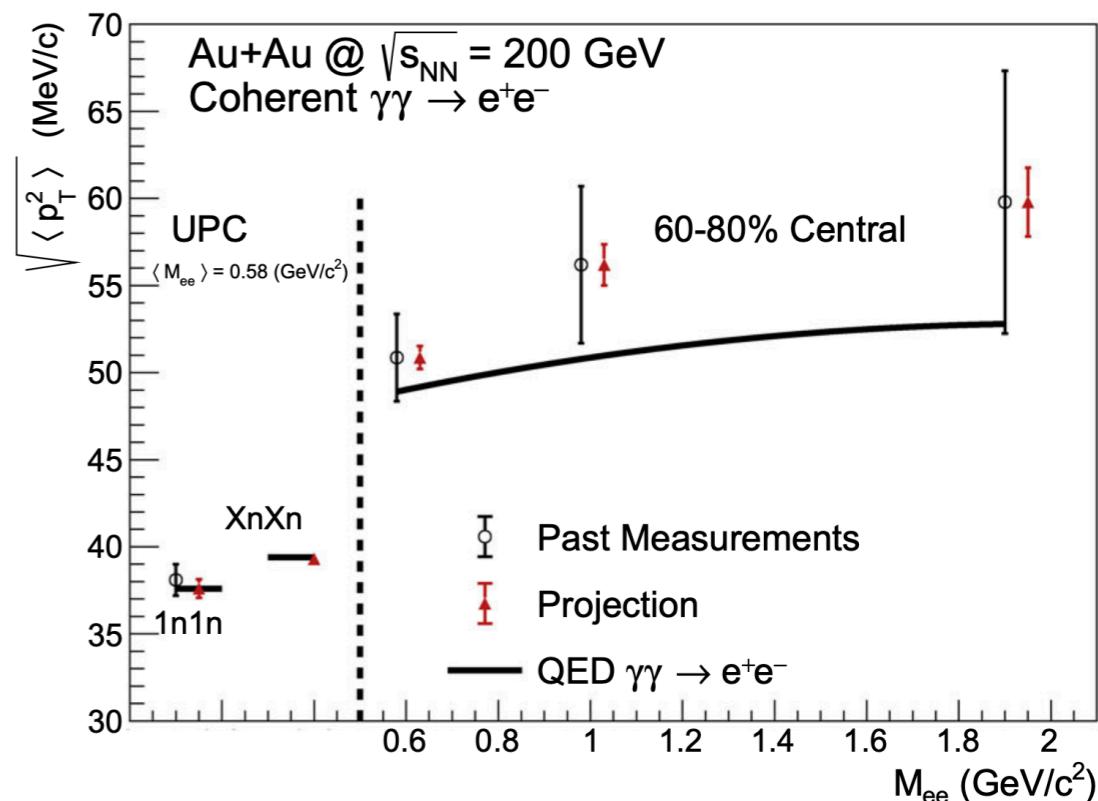
- The b dependence of photon p_T should be considered to explore QGP EM properties

- RHIC run 2023-2025
- LHC run3 & 4

- $\langle p_T \rangle$ or $\langle \alpha \rangle$ w.r.t. event plane

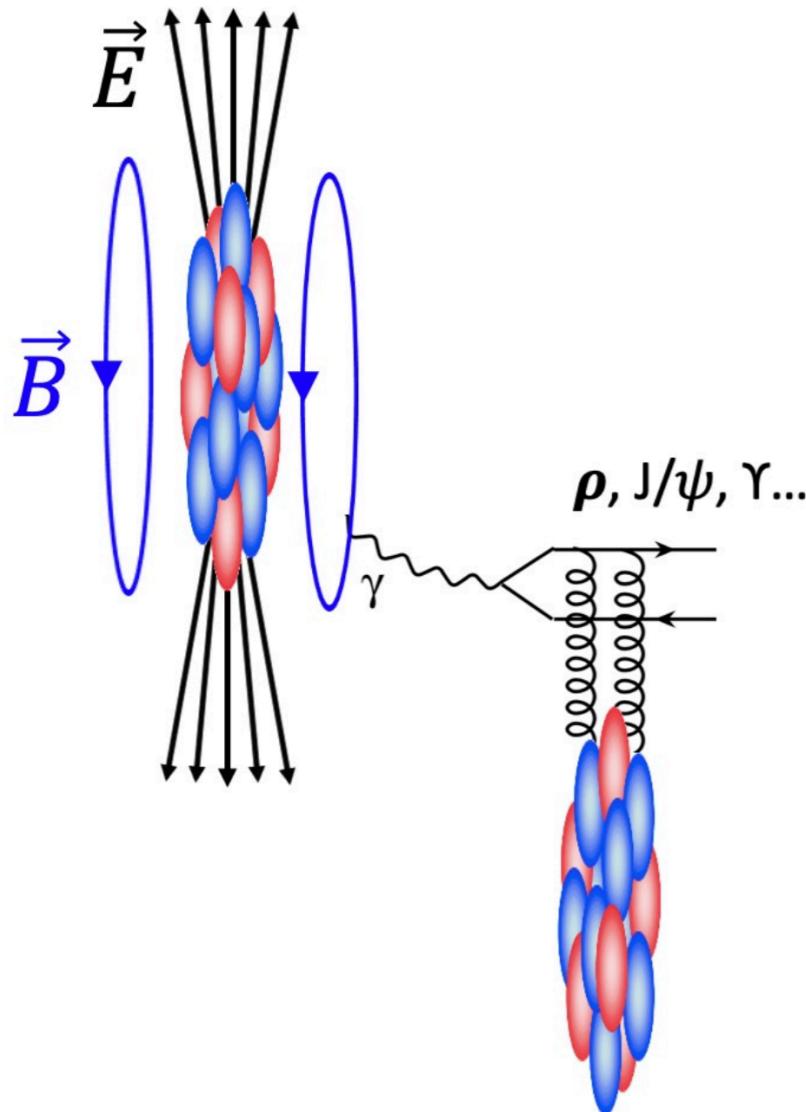
- In plane $>$ out of plane \Rightarrow Magnetic field
- In plane $<$ out of plane \Rightarrow Multiple scattering

Brandenburg et al., EPJA 57 (2021) 299

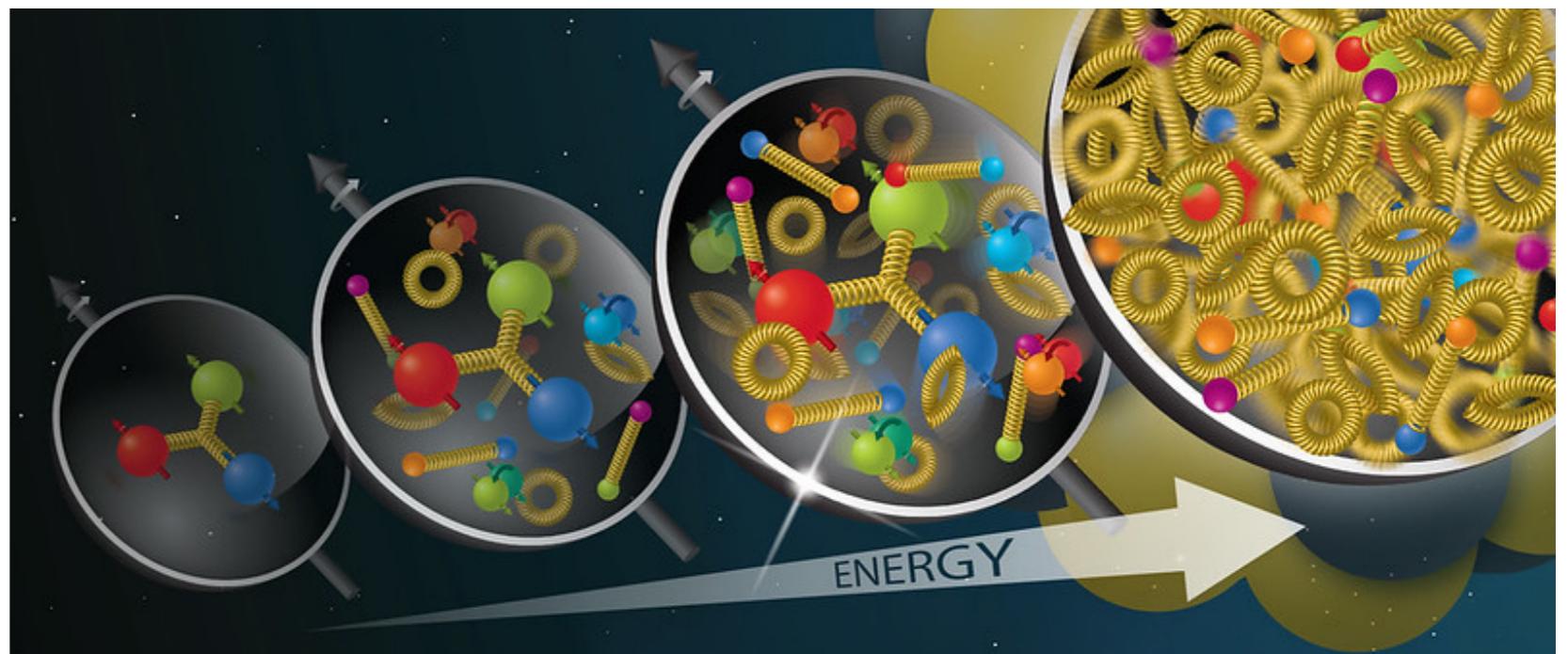


ATLAS, PRC 107 (2023) 054907

Impact to image the nuclear



$$\frac{d\sigma(\gamma A \rightarrow VA)}{dt} |_{t=0} \propto [xG_A(x, Q^2)]^2$$



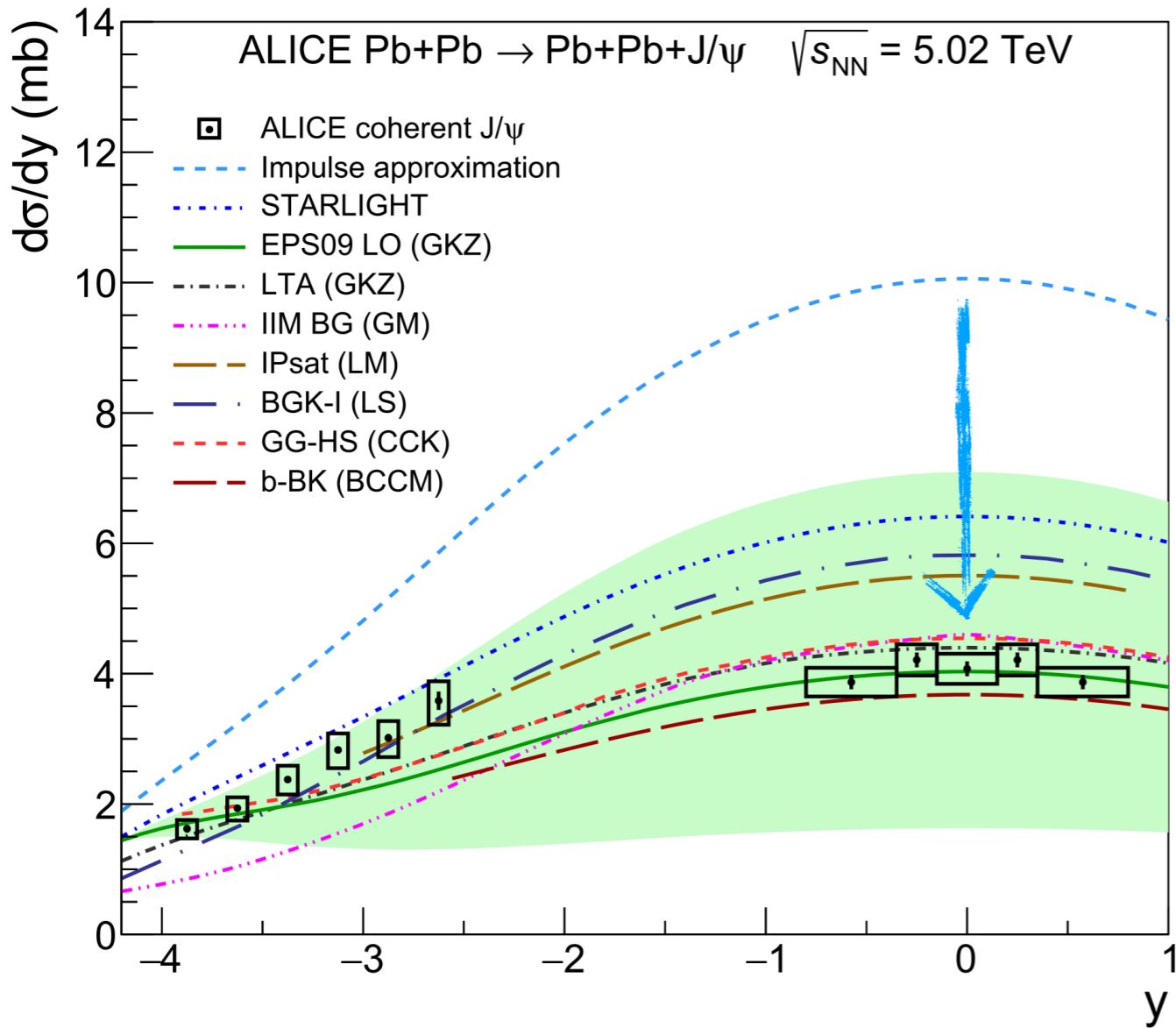
Middle Rapidity

Forward Rapidity

Impact to image the nuclear

ALICE, PLB 798 (2019) 134926

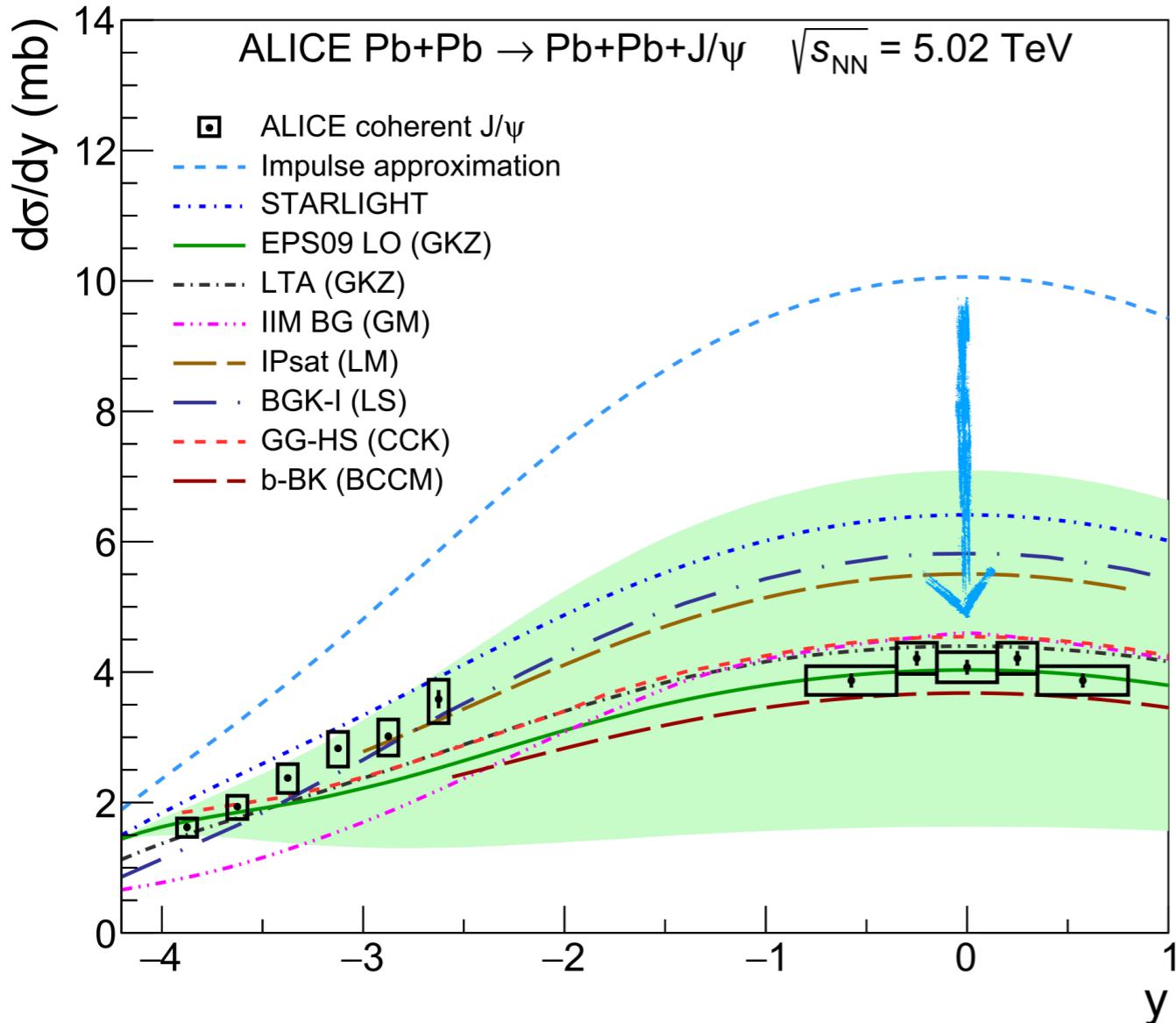
ALICE, EPJC 81 (2021) 712



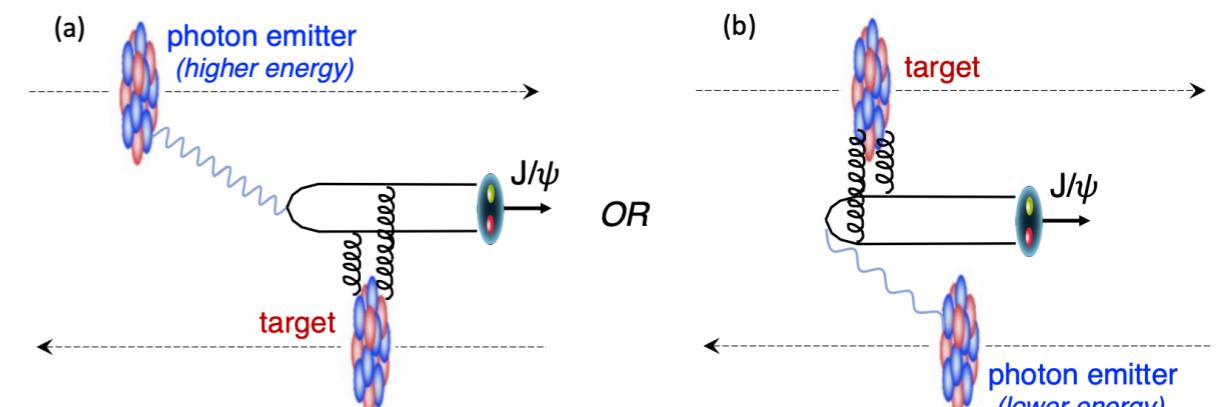
- Nuclear gluon suppression factor $R_g^{Pb} = 0.64 \pm 0.04$ at $x \sim 10^{-3}$

Impact to image the nuclear

ALICE, PLB 798 (2019) 134926
 ALICE, EPJC 81 (2021) 712



- Nuclear gluon suppression factor $R_g^{Pb} = 0.64 \pm 0.04$ at $x \sim 10^{-3}$
- Two-way ambiguity in A+A UPC



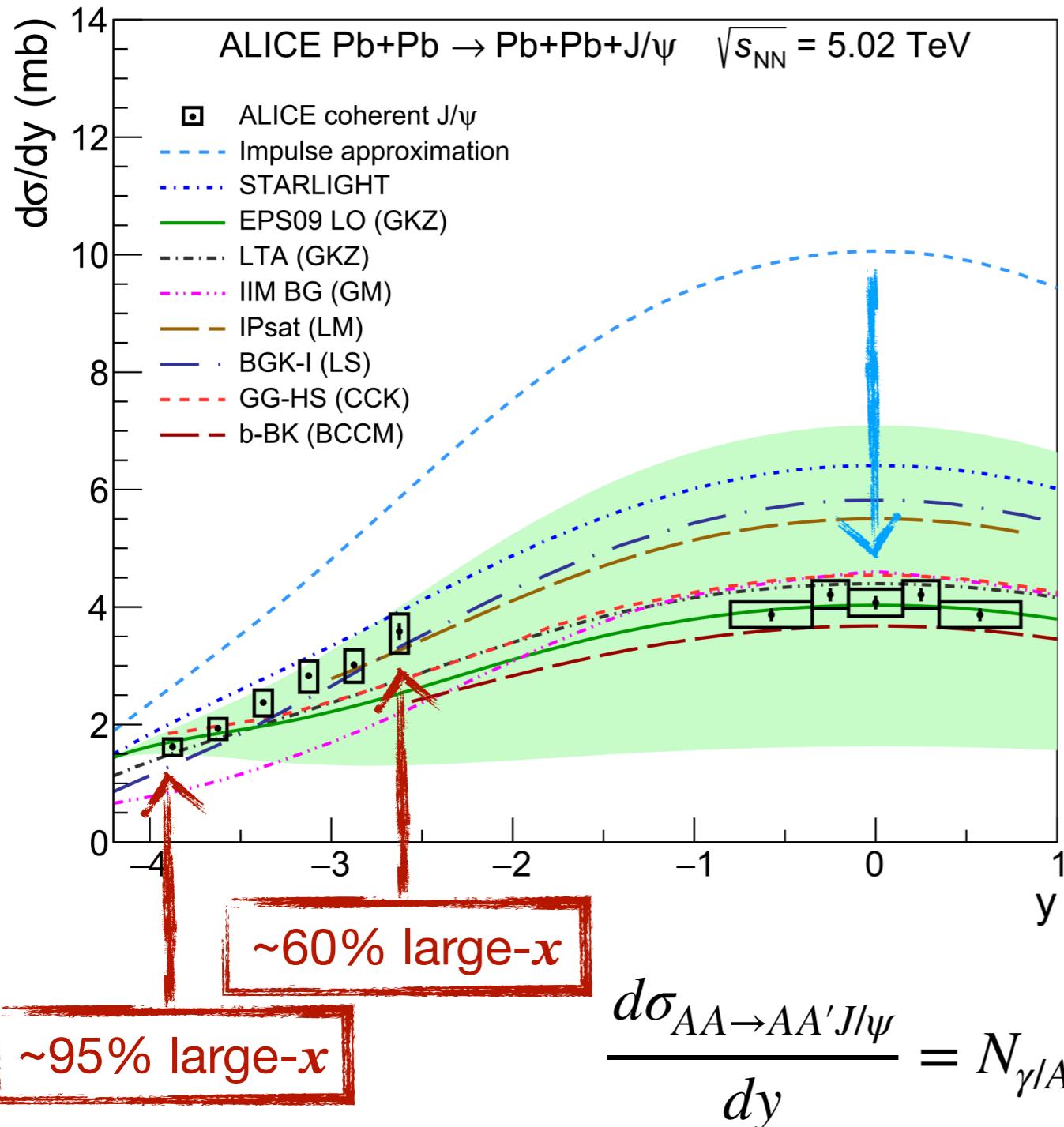
$$x = \frac{M_{J/\psi}}{2E_{beam}} e^{-y}$$

$$x = \frac{M_{J/\psi}}{2E_{beam}} e^{+y}$$

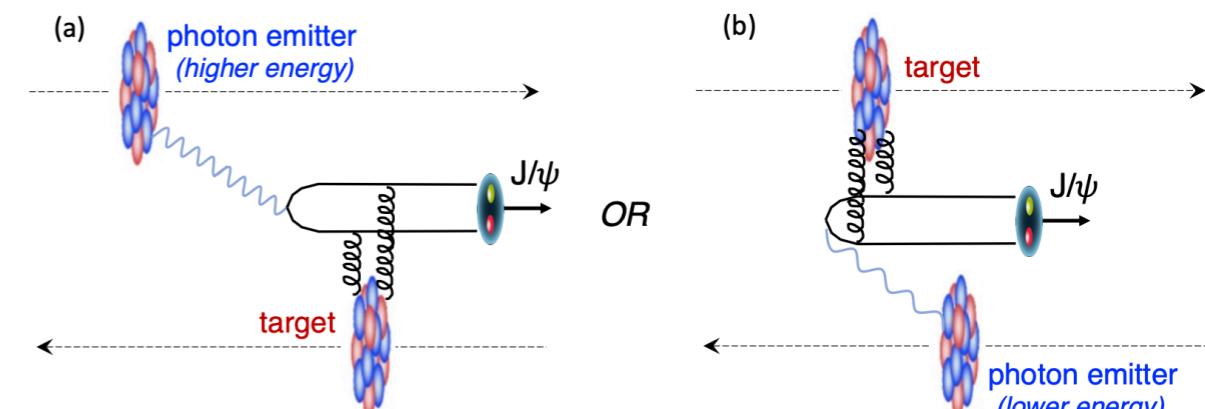
$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}}{dy} = N_{\gamma/A}(\omega_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_1)} + N_{\gamma/A}(\omega_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_2)}$$

Impact to image the nuclear

ALICE, PLB 798 (2019) 134926
 ALICE, EPJC 81 (2021) 712



- Nuclear gluon suppression factor $R_g^{Pb} = 0.64 \pm 0.04$ at $x \sim 10^{-3}$
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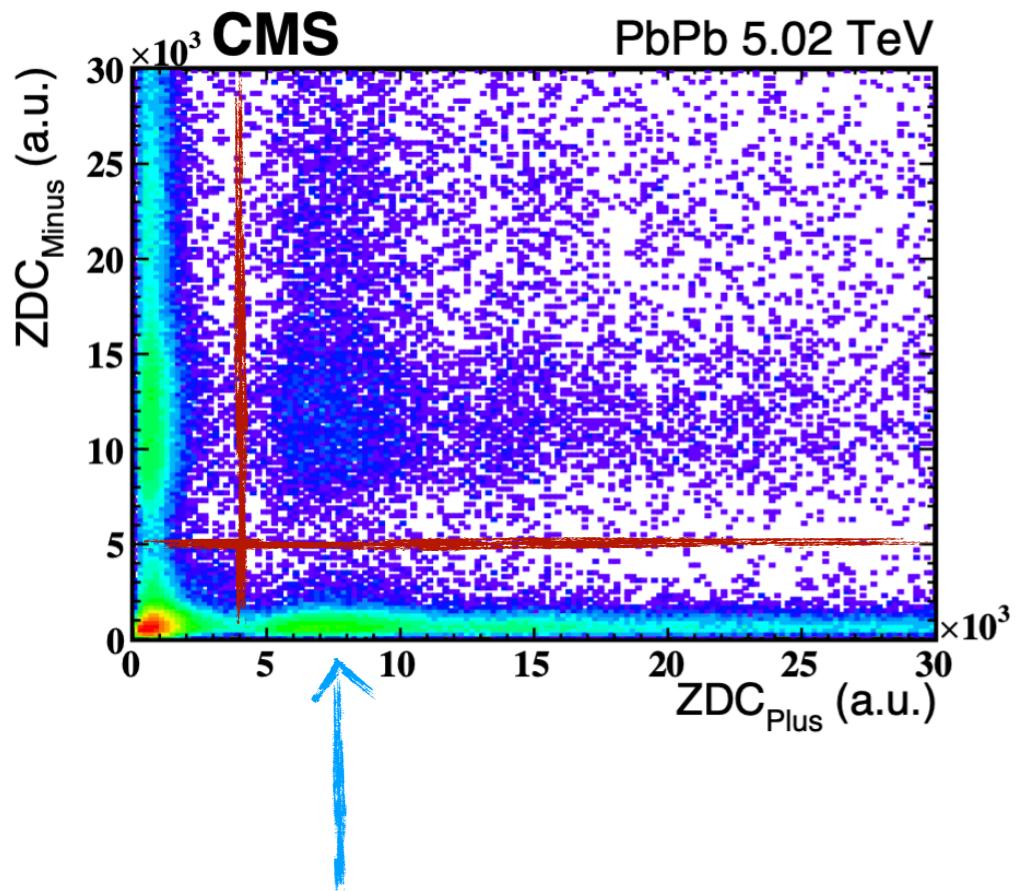
$$x = \frac{M_{J/\psi}}{2E_{beam}} e^{-y}$$

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Impact to image the nuclear

Guzey *et al.*, EPJC 74 (2014) 2942

CMS, PRL 127 (2021) 122001



Single neutron peak

$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}^{0n0n}}{dy} = N_{\gamma/A}^{0n0n}(\omega_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_1)} + N_{\gamma/A}^{0n0n}(\omega_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_2)}$$

$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}^{0nXn}}{dy} = N_{\gamma/A}^{0nXn}(\omega_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_1)} + N_{\gamma/A}^{0nXn}(\omega_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_2)}$$

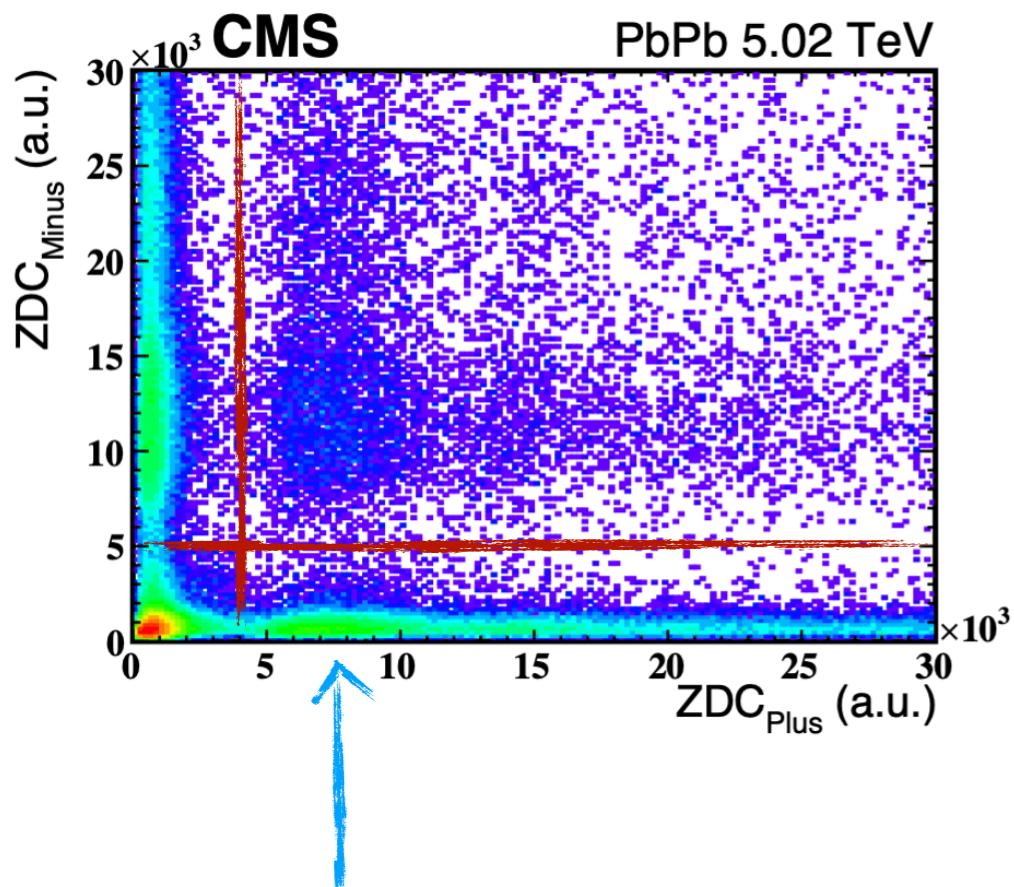
$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}^{XnXn}}{dy} = N_{\gamma/A}^{XnXn}(\omega_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_1)} + N_{\gamma/A}^{XnXn}(\omega_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_2)}$$

Impact to image the nuclear

Guzey et al., EPJC 74 (2014) 2942

CMS, PRL 127 (2021) 122001

Experimental measurements



$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}^{0n0n}}{dy} = N_{\gamma/A}^{0n0n}(\omega_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_1)} + N_{\gamma/A}^{0n0n}(\omega_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_2)}$$

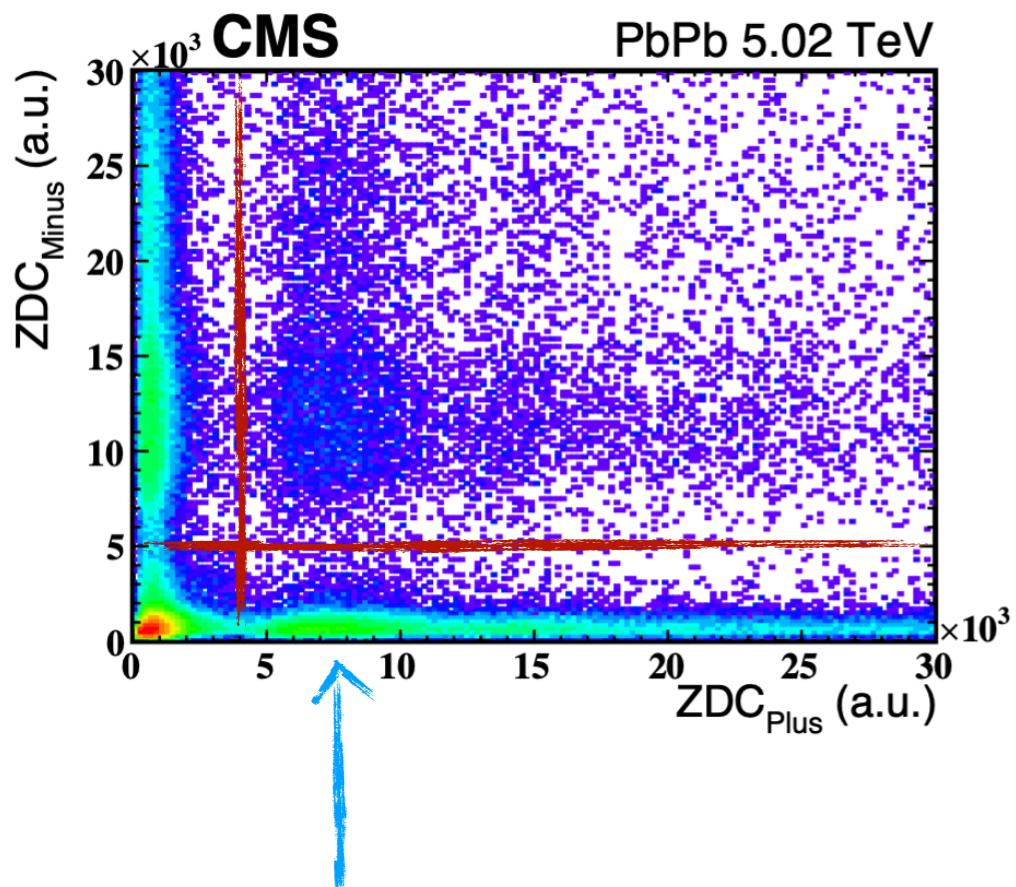
$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}^{0nXn}}{dy} = N_{\gamma/A}^{0nXn}(\omega_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_1)} + N_{\gamma/A}^{0nXn}(\omega_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_2)}$$

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Impact to image the nuclear

Guzey et al., EPJC 74 (2014) 2942

CMS, PRL 127 (2021) 122001



Experimental measurements

Photon flux from theory

$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}^{0n0n}}{dy} = N_{\gamma/A}^{0n0n}(\omega_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_1)} + N_{\gamma/A}^{0n0n}(\omega_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_2)}$$

$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}^{0nXn}}{dy} = N_{\gamma/A}^{0nXn}(\omega_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_1)} + N_{\gamma/A}^{0nXn}(\omega_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_2)}$$

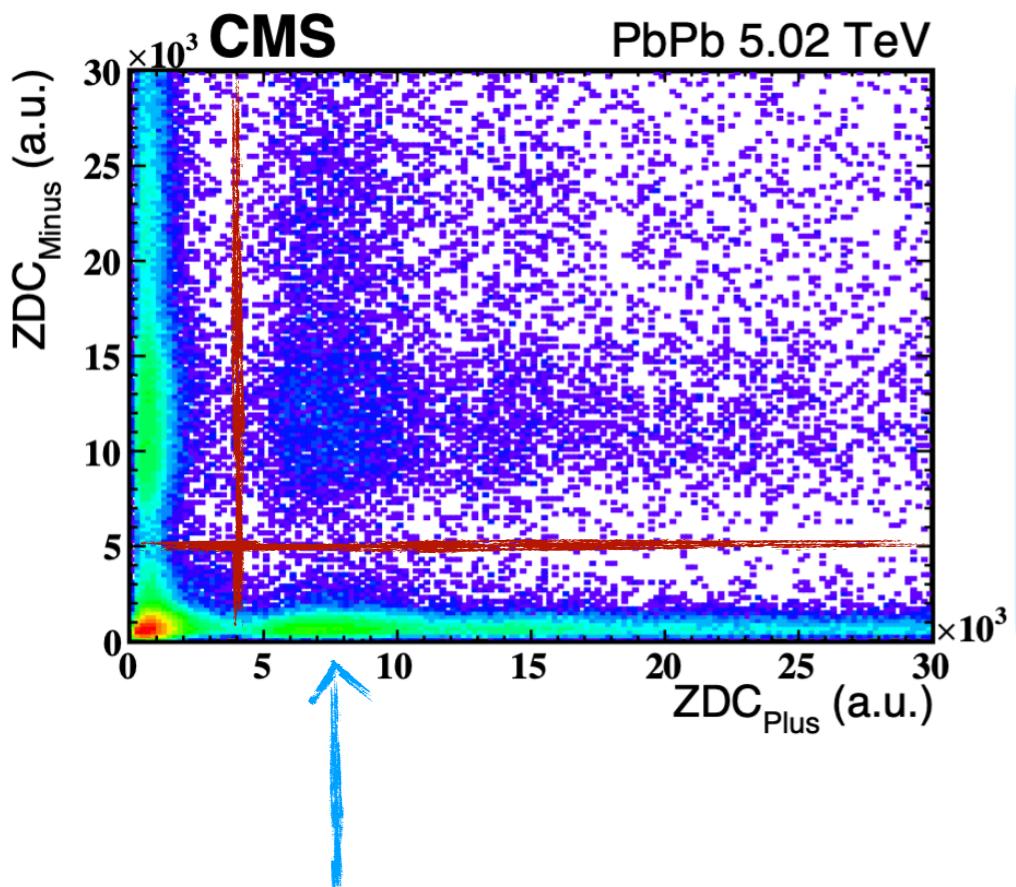
$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}^{XnXn}}{dy} = N_{\gamma/A}^{XnXn}(\omega_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_1)} + N_{\gamma/A}^{XnXn}(\omega_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_2)}$$

Single neutron peak

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Guzey et al., EPJC 74 (2014) 2942

CMS, PRL 127 (2021) 122001



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Photon flux from theory

$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}^{0n0n}}{dy} = N_{\gamma/A}^{0n0n}(\omega_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_1)} + N_{\gamma/A}^{0n0n}(\omega_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_2)}$$

$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}^{0nXn}}{dy} = N_{\gamma/A}^{0nXn}(\omega_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_1)} + N_{\gamma/A}^{0nXn}(\omega_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_2)}$$

$$\frac{d\sigma_{AA \rightarrow AA' J/\psi}^{XnXn}}{dy} = N_{\gamma/A}^{XnXn}(\omega_1) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_1)} + N_{\gamma/A}^{XnXn}(\omega_2) \cdot \sigma_{\gamma A \rightarrow J/\psi A'(\omega_2)}$$

What we need!

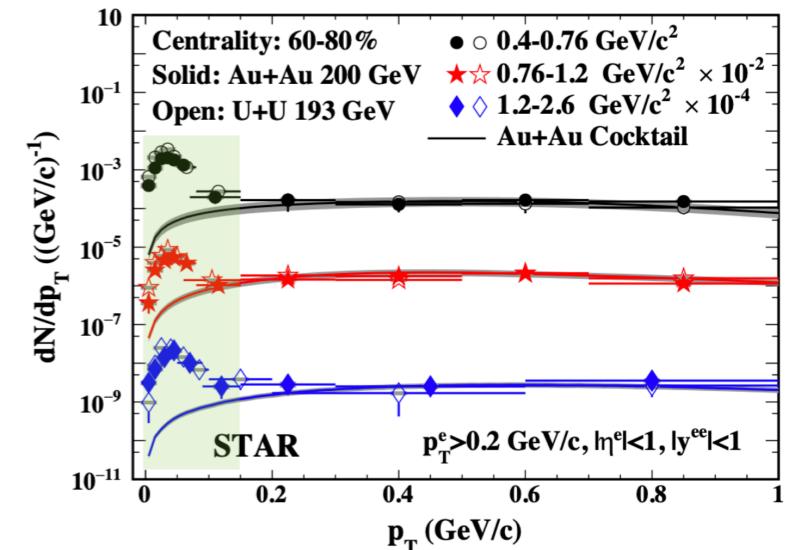
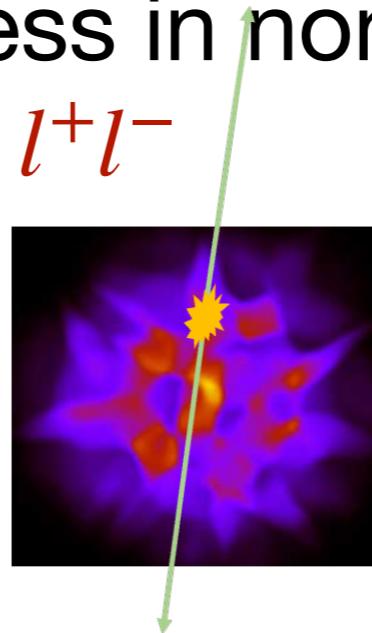
Solve the “two-way ambiguity”

Probe gluons at *small x* in heavy nucleus!

Summary

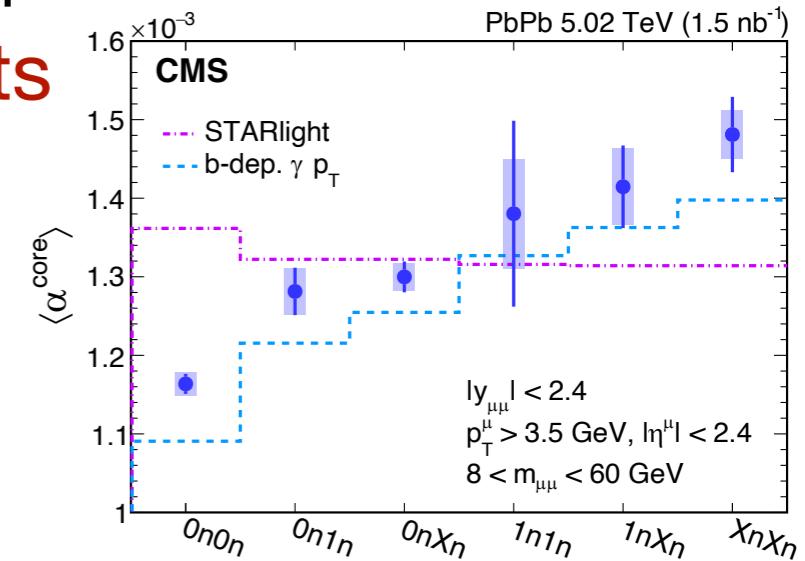
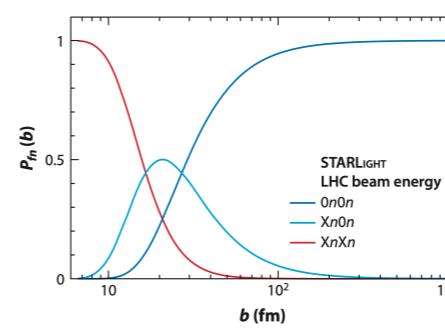
- Observed Breit-Wheeler process in non-UPC

- Probe QGP medium using $\gamma\gamma \rightarrow l^+l^-$



- Observed of b dependence of photon p_T

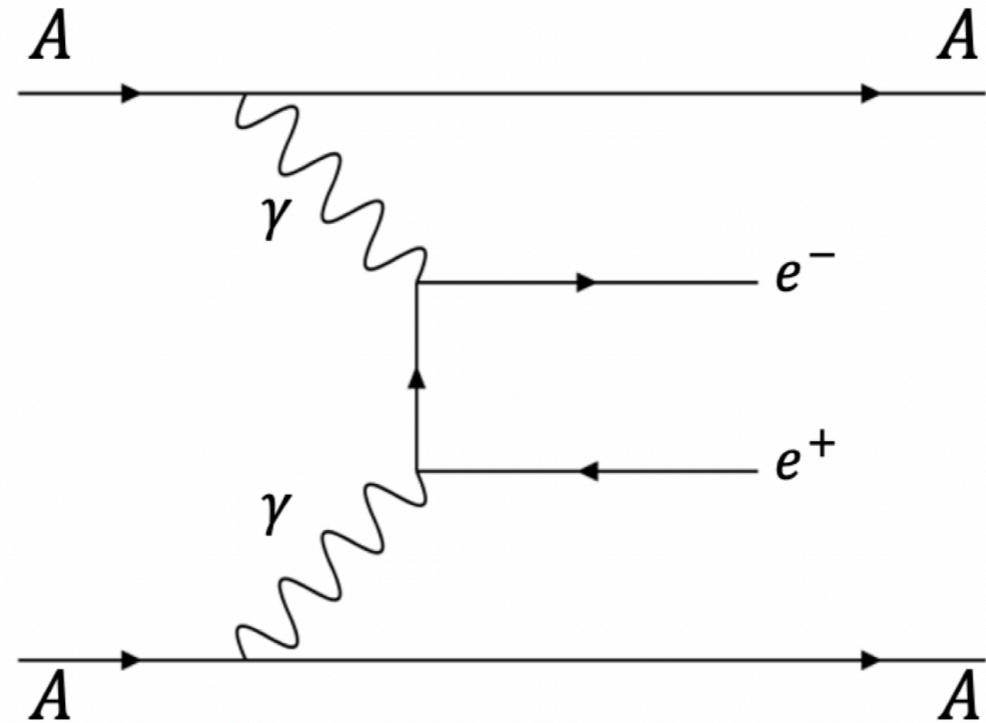
- Precise reference for probing QGP EM effects



- New tool of exploring small x physics from b dependence of photoproduced vector mesons

Backups

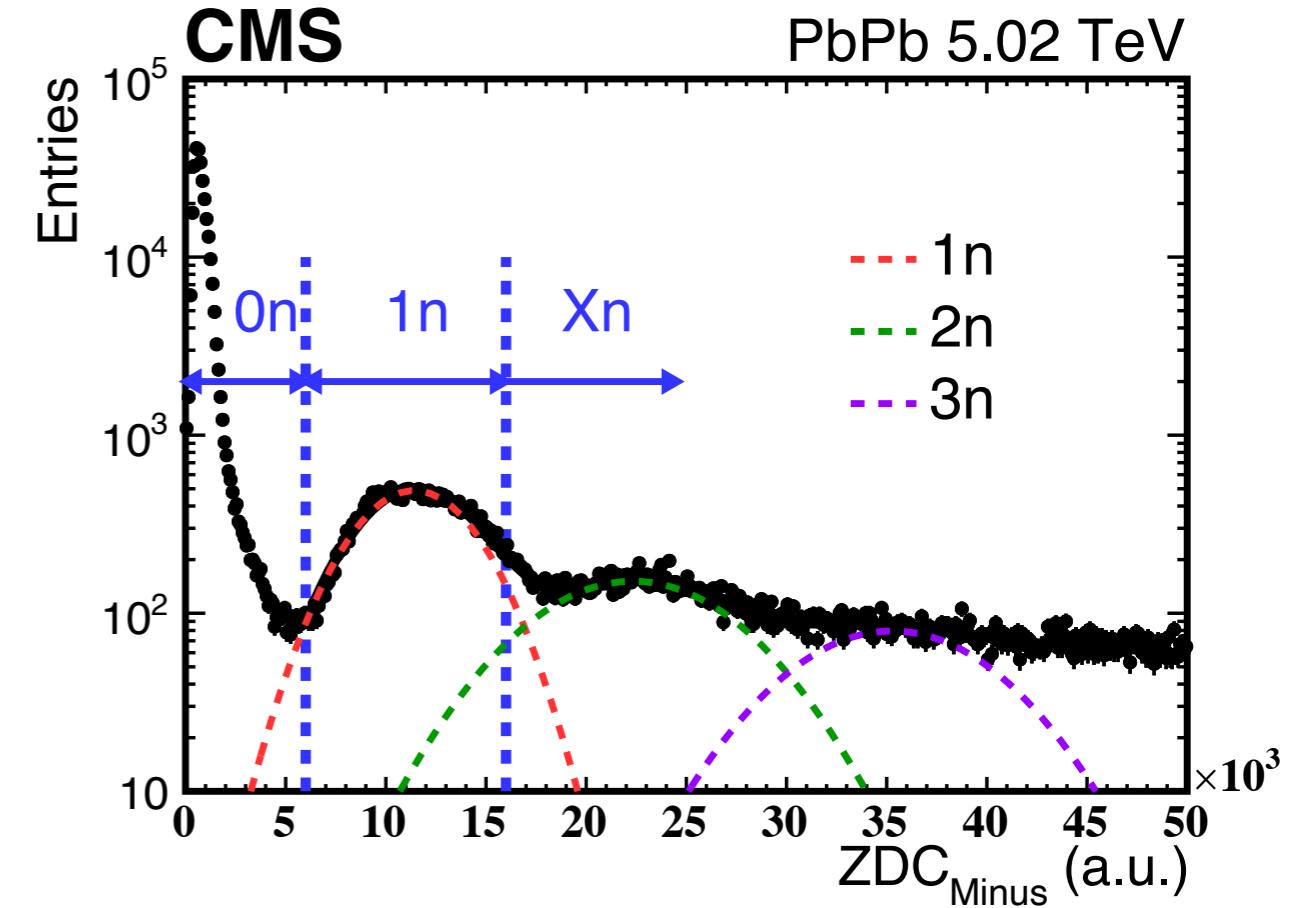
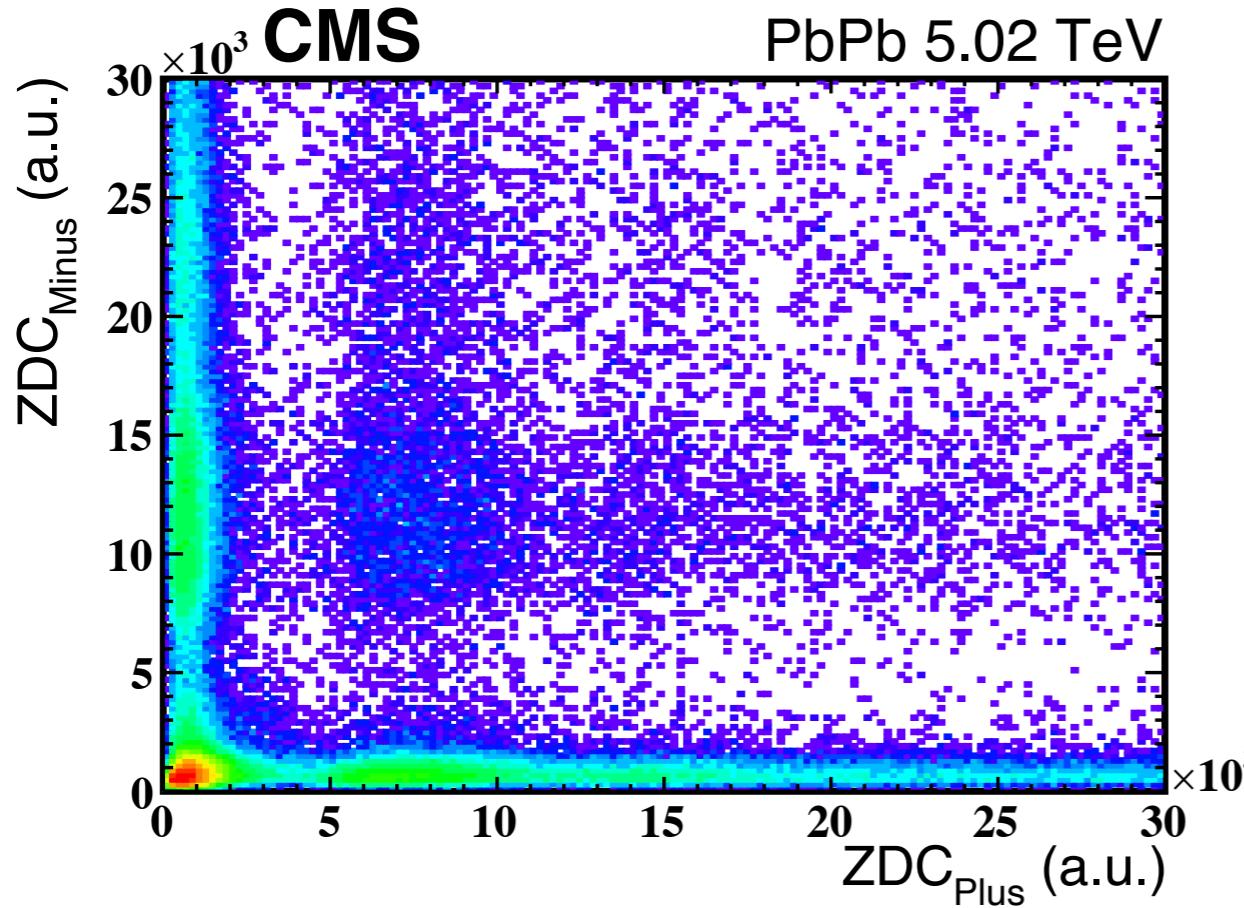
Breit-Wheeler process



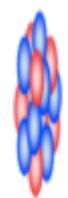
- **Breit-Wheeler process:** converting **real** photon into e^+e^-
 - Proposed in 1934 [Breit & Wheeler, Phys. Rev. 46 (1934) 1087]

$Q^2 < (\hbar/R_A)^2$ in UPC \Rightarrow almost real

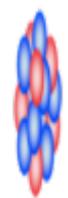
Determine neutron multiplicity



- Straight cuts to disentangle neutrons
 - 0n0n, 0n1n, 0nXn, 1n1n, 1nXn, XnXn ($X \geq 2$)

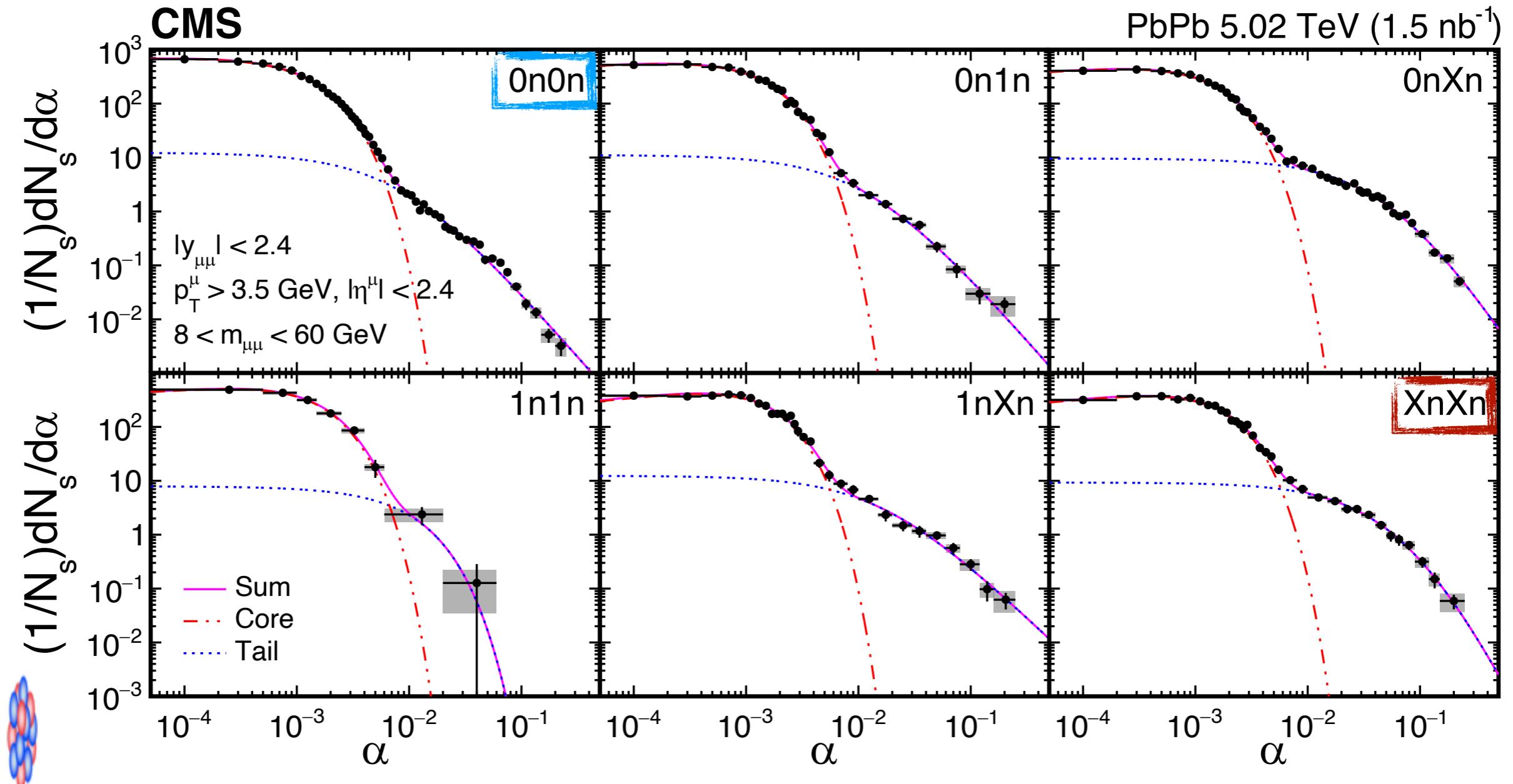


Fewer neutrons



More neutrons

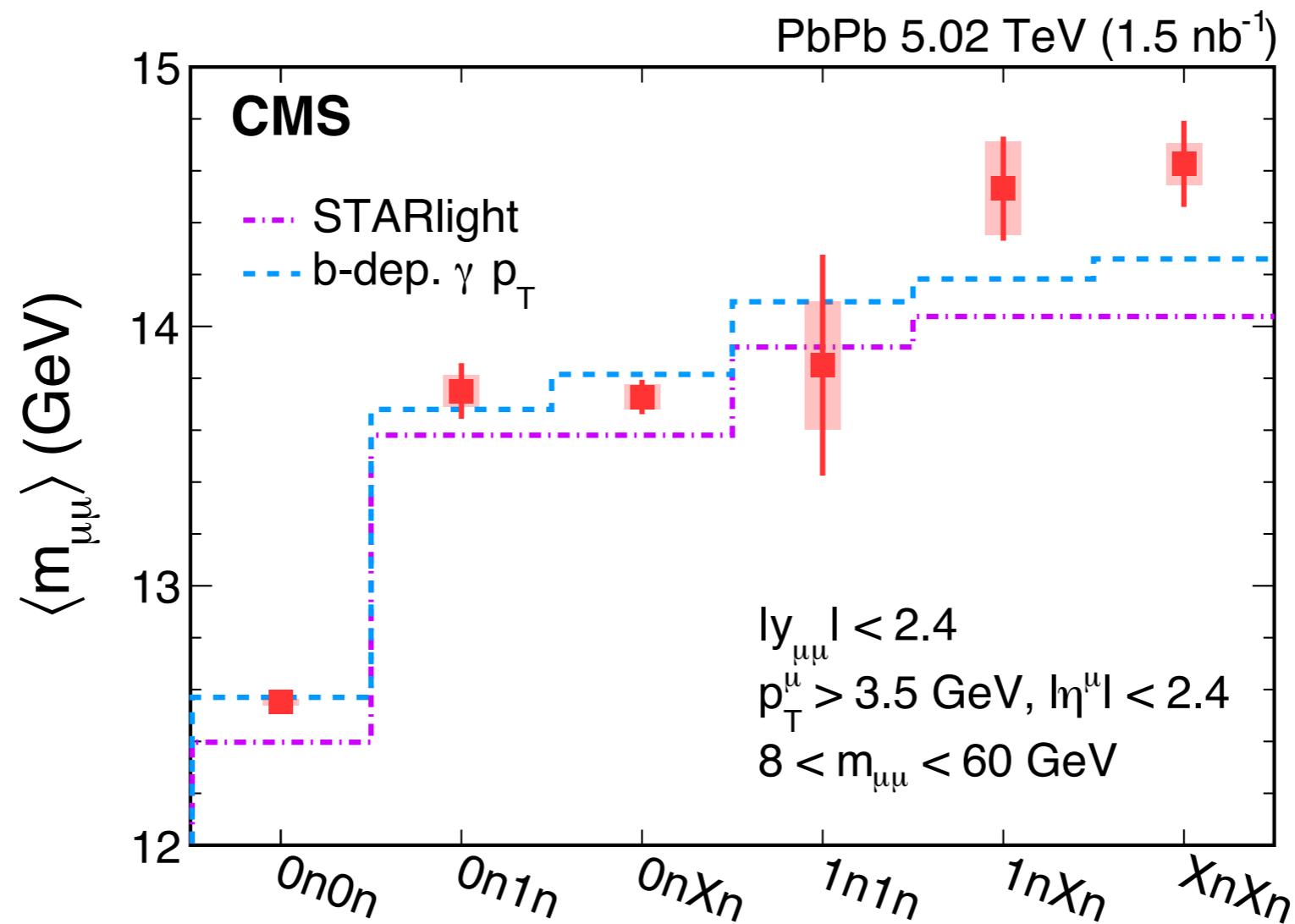
α spectrum vs. neutron multiplicity



- 0n0n (fewer neutrons) \Rightarrow XnXn (more neutrons)
 - Tail contribution becomes larger

$\langle M_{\mu\mu} \rangle$ vs. neutron multiplicity

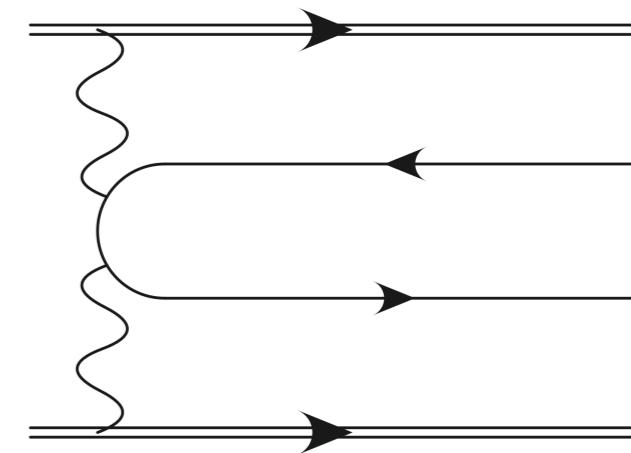
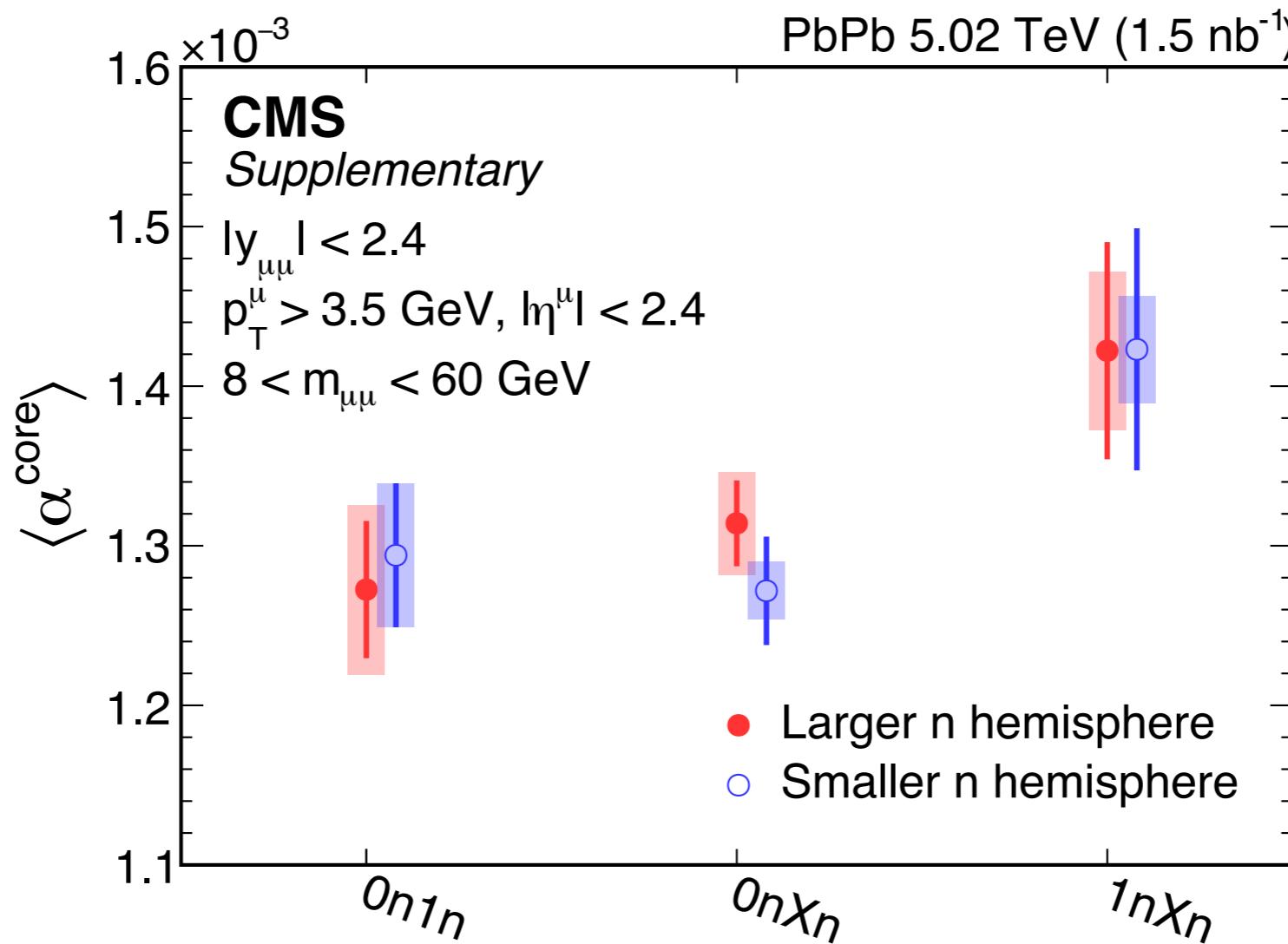
CMS, PRL 127 (2021) 122001



- Strong neutron multiplicity dependence of $\langle M_{\mu\mu} \rangle$
 - Deviation from constant: $\gg 5\sigma$
 - **b** dependence of initial photon energy

Rapidity dependence of $\langle \alpha^{\text{core}} \rangle$

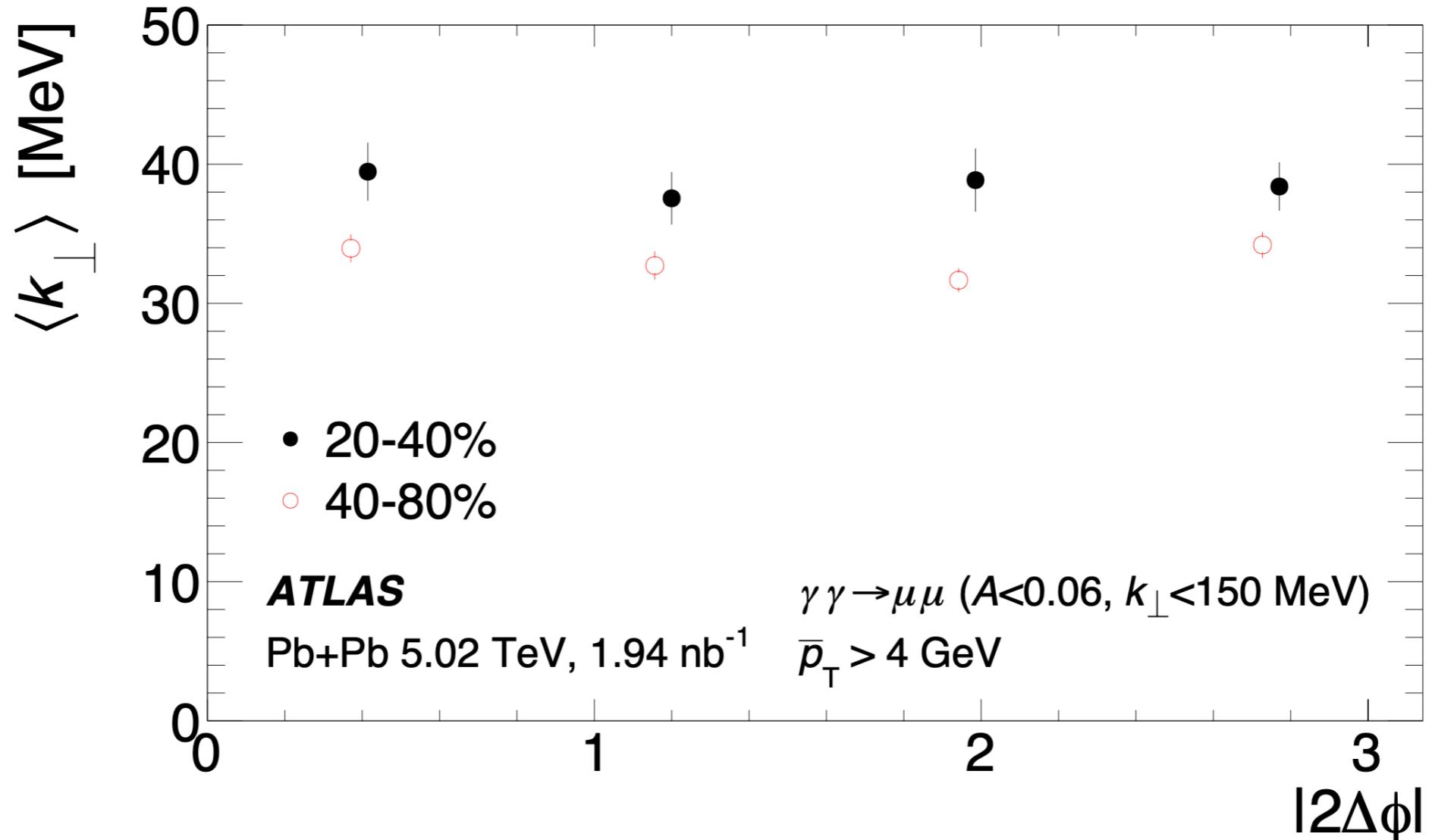
CMS, PRL 127 (2021) 122001



- $\langle \alpha^{\text{core}} \rangle$ has no rapidity dependence
 - Core dominantly comes from LO $\gamma\gamma$ scatterings

Event plane dependence

ATLAS, arXiv:2206.12594



$$|2\Delta\phi| \equiv |2(\phi_{\mu\mu} - \Psi_2)|$$

$$\phi_{\mu\mu} = \frac{1}{2} (\phi_1 + \pi + \phi_2)$$

$$k_{\perp} \equiv \frac{1}{2} (p_{T1} + p_{T2}) (\pi - |\phi_1 - \phi_2|) = \pi \alpha \bar{p}_T$$