UPC Physics Workshop 2023

Nuclear excitation by electron capture in electron-ion collisions



南开大学物理科学学院

2023年05月27日

NEEC

NEEC: Nuclear Excitation by Electron Capture



- 1976年在理论上提出, 2018年首个实验迹象被报道
- 核结构的研究、核天体物理的研究
- 通过调控电子与原子来调控原子核
- 核能触发 (Isomer triggering): 新型核能
- 核钟 (Nuclear clock) 相关研究

Why NEEC?

 NEEC有望用于触发同核异能素 能量释放:新型核能





Isomer triggering

Isomer: long-lived excited state of nuclei

Gunst, Litvinov, Keitel, Pálffy, Phys. Rev. Lett. 112, 082501(2014) Gunst, Wu, Kumar, Keitel, Pálffy, Phys. Plasmas 22, 112706 (2015) Wu, Gunst, Keitel, Pálffy, Phys. Rev. Lett. 120, 052504 (2018) Gunst, Wu, Keitel, Pálffy, Phys. Rev. E 97, 063205 (2018) Wu, Keitel, Pálffy, Phys. Rev. A 100, 063420 (2019)

Why NEEC?

首个NEEC实验迹象在2018年被报道

◎离子束-固体靶相互作用

●^{93m}Mo同核异能素衰变

C. J. Chiara et al., Nature 554, 216 (2018)

◇ 首个NEEC实验迹象的理论分析?◇ 通过电子波函数的改变来影响NEEC?一 涡旋电子束



☑研究背景

I首个NEEC实验迹象的讨论

凶涡旋电子束







I首个NEEC实验迹象的讨论

区涡旋电子束





Wu, Keitel, Pálffy, Phys. Rev. Lett. 122, 212501 (2019)

首次实验观测NEEC

- ^{93m}Mo 衰变
- $P_{\rm exc} = 0.01$
- 没有相关的NEEC理论计算

C. J. Chiara et al., Nature 554, 216 (2018)



NEEC 截面

$$\sigma_q^{\alpha}(E) = S_q^{\alpha} \frac{\Gamma_{q,\alpha}/(2\pi)}{(E - E_{q,\alpha})^2 + \frac{1}{4}\Gamma_{q,\alpha}^2}$$

 S_q^{α} : resonance strength

$$S_q^{\alpha} \propto \left| \left\langle \Psi_g^N \right| \left\langle \Psi_g^e \right| H_N | \Psi_i^e, \psi_s \rangle \left| \Psi_i^N \right\rangle \right|^2$$

• Coulomb interaction (E transitions)

$$H_{en} = \int d^3 r_n \frac{\rho_n(\vec{r}_n)}{\left| \vec{r}_e - \vec{r}_n \right|}$$

• Magnetic Hamiltonian (M transitions)

$$H_{magn} = -\frac{1}{c} \overrightarrow{\alpha} \int d^3 r_n \frac{\overrightarrow{j}_n(\overrightarrow{r}_n)}{\left| \overrightarrow{r} - \overrightarrow{r}_n \right|}$$

Matrix elements

- Nuclear via reduced transition probability *B*(*E*/*ML*)
- Electronic wavefunctions Bound electrons: GRASP92

NEEC几率

$$P = \sum_{q,\alpha} \int f_q \phi \sigma_q^{\alpha} dt$$

NEEC截面

$$\sigma_q^{\alpha}(E) = S_q^{\alpha}(E) \frac{\Gamma_{q,\alpha}/(2\pi)}{(E - E_{q,\alpha})^2 + \frac{1}{4}\Gamma_{q,\alpha}^2}$$

Γ很小

$$P = \sum_{q,\alpha} f_q(E_{q,\alpha}^{\text{ion}}) P_q^{\alpha}(E_{q,\alpha}^{\text{ion}})$$
$$P_q^{\alpha} = n_e S_q^{\alpha}(E_{q,\alpha}) \frac{m_i}{m_e} \frac{1}{-\left(\frac{E^{\text{ion}}}{dx}\right)\Big|_{E_{q,\alpha}^{\text{ion}}}}$$

Wu, Keitel, Pálffy, Phys. Rev. Lett. 122, 212501 (2019)



Wu, Keitel, Pálffy, Phys. Rev. Lett. 122, 212501 (2019)



Wu, Keitel, Pálffy, Phys. Rev. Lett. 122, 212501 (2019)



^{93m}Mo产生和衰变观测区域没有有效分离从而可能引起污染?

S. Guo *et al.*, Nature 594, E1 (2021)

C. J. Chiara et al., Nature 594, E3 (2021)

⊘使用Isomer Beam重新观测NEEC



- **O** ¹²C(⁸⁶Kr, 5*n*)^{93m}Mo
- **O** 93mMo离子能量: 460 MeV
- 分离^{93m}Mo产生和衰变观测区域
- o 低伽马本底
- 激发几率上限: $P_{\text{exc}} < 2 \times 10^{-5}$

O NEEC几率理论比值: *P*(460 MeV)/*P*(840 MeV)~8%

S. Guo, B. Ding, X. H. Zhou, Y. B. Wu,, Y. H. Zhang, Phys. Rev. Lett. 128, 242502 (2022)



☑研究背景

☑首个NEEC实验迹象的讨论

☑涡旋电子束



NEEC-涡旋电子束

通过调节电子波函数来操纵原子核?



Electron vortex beams carry orbital angular momentum (OAM)

Wu, Gargiulo, Carbone, Keitel, Pálffy, Phys. Rev. Lett. 128, 162501 (2022)

NEEC-涡旋电子束

电子波函数

$$\psi(\mathbf{r}) = \int \frac{d^2 \mathbf{p}_{\perp}}{(2\pi)^2} a_{\zeta m}(\mathbf{p}_{\perp}) u_{\mathbf{p}} e^{i\mathbf{p}\cdot\mathbf{r}}$$

$$a_{\zeta m}(\mathbf{p}_{\perp}) = (-i)^m e^{im\alpha_p} \sqrt{2\pi/\zeta} \delta(|\mathbf{p}_{\perp}| - \zeta)$$

m: vortex quantum number related to OAM $\mathbf{p} = (\mathbf{p}_{\perp}, p_z) = (\zeta \cos \alpha_p, \zeta \sin \alpha_p, p_z)$



$$\begin{split} \psi_{\pm\frac{1}{2}}(\mathbf{r}) \propto e^{ip_{z}z} \begin{bmatrix} \begin{pmatrix} \sqrt{E_{\pm}} \\ 0 \\ \cos\theta_{0}\sqrt{E_{-}} \\ 0 \end{bmatrix} e^{im\varphi}J_{m}(\zeta r_{\perp}) + i \begin{pmatrix} 0 \\ 0 \\ 0 \\ \sin\theta_{0}\sqrt{E_{-}} \end{pmatrix} e^{i(m+1)\varphi}J_{m+1}(\zeta r_{\perp}) \end{bmatrix} \\ \psi_{\pm\frac{1}{2}}(\mathbf{r}) \propto e^{ip_{z}z} \begin{bmatrix} \begin{pmatrix} 0 \\ \sqrt{E_{\pm}} \\ 0 \\ -\cos\theta_{0}\sqrt{E_{-}} \end{pmatrix} e^{im\varphi}J_{m}(\zeta r_{\perp}) - i \begin{pmatrix} 0 \\ 0 \\ \sin\theta_{0}\sqrt{E_{-}} \\ 0 \end{bmatrix} e^{i(m-1)\varphi}J_{m-1}(\zeta r_{\perp}) \end{bmatrix} \end{split}$$

NEEC-涡旋电子束



NEEC - 涡旋电子

$$\sigma_{neec} = \frac{2\pi^2}{p^2} \frac{2p}{J_z} Y_{neec} L_d(E - E_d)$$

$$Y_{neec} = \frac{\zeta b^2}{2} \int_0^{2\pi} \int_0^{2\pi} \frac{d\alpha_p}{2\pi} \frac{d\alpha_k}{2\pi} e^{im(\alpha_p - \alpha_k)} Y_{neec}^{p,k} {}_0F_1(2;u)$$

$$Y_{neec}^{p,k} = Y_a \sum_{\kappa,m_l} \frac{Y_b}{2l+1} Y_{lm_l}^*(\theta_k, \varphi_k) Y_{lm_l}(\theta_p, \varphi_p)$$

$$u = -b^2 \zeta^2 [1 - \cos(\alpha_k - \alpha_p)]/2$$

Transitions of electric multipolarity L

$$Y_a = \frac{4\pi^2 (2J_d + 1)}{(2J_i + 1)(2L + 1)^2} \frac{B\rho_i}{R_0^{2(L+2)}}; \quad Y_b = \left[C(j_d \ L \ j; \frac{1}{2} \ 0 \ \frac{1}{2})\right]^2 |R_{L,\kappa_d,\kappa}|^2$$

Transitions of magnetic multipolarity L

$$Y_{a} = \frac{4\pi^{2}(2J_{d}+1)}{(2J_{i}+1)L^{2}(2L+1)^{2}}B\rho_{i}; \quad Y_{b} = (2j+1)(\kappa_{d}+\kappa)^{2} \binom{j_{d}}{\frac{1}{2}} \frac{j}{-\frac{1}{2}} \frac{L}{0}^{2} \left|R_{L,\kappa_{d},\kappa}\right|^{2}$$

Wu, Gargiulo, Carbone, Keitel, Pálffy, Phys. Rev. Lett. 128, 162501 (2022)

NEEC-涡旋电子束



Wu, Gargiulo, Carbone, Keitel, Pálffy, Phys. Rev. Lett. 128, 162501 (2022)



☑研究背景

☑首个NEEC实验迹象的讨论

☑涡旋电子束



总结

首个NEEC实验迹象 — 离子束-固体靶

- NEEC几率 $P \sim 10^{-11} \ll P_{exc} = 0.01$ (首个NEEC实验迹象)
- 离子能量偏低、低伽马本底条件下重新测量: $P_{exc} < 2 \times 10^{-5}$

NEEC — 涡旋电子束-离子束

- •相比于平面波,涡旋电子束可以使NEEC截面提高2-6个数量级
- 通过电子波函数的改变影响原子核能级的跃迁

谢谢大家!