



CP violation in Higgs sector

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Why CPV interesting?

- Baryon genesis: why more matter than anti-matter?
- Sakharov's conditions
 Standard Model
 Baryon number violation
 C and CP violations
 Departure from the equilibrium

In SM

- Sphaleron process can lead B violation. But,
- CKM, only CPV source in SM, is not enough
- Higgs mechanism in SM is too smooth

Higgs boson open the door

- The discovery of Higgs boson open a new sector
 - Many BSM, e.g. 2HDM, provide new CP violation sources and strong phase transition at Higgs sector



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Higgs boson open the door

- The discovery of Higgs boson open a new sector
 - Many BSM, e.g. 2HDM, provide new CP violation sources and strong phase transition at Higgs sector

- LHC Higgs physics program
 - Is there CP violation in Higgs coupling
 - EW symmetry breaking the 1st or 2nd order phase transition

$$\begin{aligned} \mathcal{L} &= -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ &+ i \overline{\psi} \overline{\psi} \psi + h.c. \\ &+ \overline{\psi} \overline{\psi} \overline{\psi} + h.c. \\ &+ \overline{\psi} \overline{\psi} \overline{\psi} \overline{\psi} - V(\phi) \end{aligned}$$

Higgs boson open the door

- Higgs boson is found a CP-even scalar, as SM prediction
 - J^P=0⁺ compared with alternative spin-model



But, mixing between CP-even and CP-odd, which could lead to CP violation, is still allowed

Indirect constraint

 Low energy experiments, e.g. electron EDM, can constrain the Higgs CP indirectly



$$\mathcal{L} \supset -\frac{y_f}{\sqrt{2}} \left(\kappa_f \, \bar{f} f + i \kappa_f \, \bar{f} \gamma_5 f \right) h$$

ACME collaboration: **eEDM<1.1×10⁻²⁹ e·cm** $|\tilde{\kappa}_e| \lesssim 1.7 \times 10^{-2}$ $|\tilde{\kappa}_t| \lesssim 1.0 \times 10^{-2}$

- But, very model dependent
 - Gauge-dependent contributions, UV-divergent diagrams, etc.

CP violation in Higgs sector

- General methodology:
 - using event topology to build some CP sensitive angle



- Indirect measurements not discussed here
 - e.g. ggf loop, cross section, etc
- ATLAS public results: https://twiki.cern.ch/twiki/bin/view/AtlasPublic

New results from ATLAS

- CP structure in Higgs-Vector boson coupling
 - Vector Boson Fusion (VBF) with $H \rightarrow \gamma \gamma$
 - arXiv:2208.02338, accepted by PRL
 - Main collaborators: Y. Fang, B. Liu, F. Guo, H. Chen, A. Tarek, etc.
- CP structure in Higgs-fermion Yukawa coupling
 - Higgs decay to lepton: $H \rightarrow \tau^+ \tau^-$
 - arXiv:2208.02338, accepted by EPJC
 - Main collaborators: H. Li, A. De Maria, K. Lie, T. Cheng, etc.

Study HVV CP in VBF production

• HVV vertex in VBF Higgs production



- Independent from Higgs decay: here, use $H \rightarrow \gamma \gamma$ (SM Br)
- EFT interpretation: two bases used

 $\begin{aligned} & \text{Warsaw: } \mathcal{L}_{\text{SMEFT}}^{\text{CP-odd}} \supset \frac{c_{H\tilde{W}}}{\Lambda^2} H^{\dagger} H \tilde{W}_{\mu\nu}^I W^{\mu\nu I} + \frac{c_{H\tilde{B}}}{\Lambda^2} H^{\dagger} H \tilde{B}_{\mu\nu}^A B^{\mu\nu} + \frac{c_{H\tilde{W}B}}{\Lambda^2} H^{\dagger} \sigma^I H \tilde{W}_{\mu\nu}^I B^{\mu\nu} \\ \\ & \text{HISZ: } \mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \tilde{g}_{H\gamma\gamma} H \tilde{A}_{\mu\nu} A^{\mu\nu} + \tilde{g}_{H\gammaZ} H \tilde{A}_{\mu\nu} Z^{\mu\nu} \\ & + \tilde{g}_{HZZ} H \tilde{Z}_{\mu\nu} Z^{\mu\nu} + \tilde{g}_{HWW} H \tilde{W}_{\mu\nu}^+ W^{-\mu\nu} \end{aligned} \qquad \tilde{d} = -\frac{m_W^2}{\Lambda^2} c_{\tilde{W}W} \end{aligned}$

Optimal Observable

- Single CP sensitive observable
 - combine kinematics information of VBF jets and Higgs, to constrain Wilsonian coefficient, e.g. \tilde{d} and $c_{h\widetilde{W}}$

$$|\mathcal{M}|^{2} = |\mathcal{M}_{SM}|^{2} + \left(\tilde{d} \cdot 2\operatorname{Re}(\mathcal{M}_{SM}^{*}\mathcal{M}_{CP-odd})\right) + \tilde{d}^{2} \cdot |\mathcal{M}_{CP-odd}|^{2}$$

$$\mathcal{O}_{opt} = \frac{2\operatorname{Re}(\mathcal{M}_{SM}^{*}\mathcal{M}_{CP-odd})}{|\mathcal{M}_{SM}|^{2}}$$

- Evaluated with
 - momentum fraction of initial-state parton x1(x2), four-momenta of Higgs boson and two VBF jets

• At rec. level,
$$x_{1,2}^{\text{reco}} = \frac{m_{Hjj}}{\sqrt{s}} e^{\pm y_{Hjj}}$$

 $m_{Hjj} (y_{Hjj})$: invariant mass (rapidity) of Higgs and VBF jet system

Optimal Observable

- Single CP sensitive observable
 - combine kinematics information of VBF jets and Higgs, to constrain Wilsonian coefficient, e.g. \tilde{d} and $c_{h\widetilde{W}}$

$$|\mathcal{M}|^{2} = |\mathcal{M}_{\rm SM}|^{2} + \left[\tilde{d} \cdot 2\operatorname{Re}(\mathcal{M}_{\rm SM}^{*}\mathcal{M}_{\rm CP-odd})\right] + \tilde{d}^{2} \cdot |\mathcal{M}_{\rm CP-odd}|^{2}$$

$$\mathcal{O}_{\text{opt}} = \frac{2 \operatorname{Re}(\mathcal{M}_{\text{SM}}^* \mathcal{M}_{\text{CP-odd}})}{|\mathcal{M}_{\text{SM}}|^2}$$

- Phenomenon
 - Inter. : asymmetry, direct CPV
 - Quad. : over cross section



VBF signature at detector



Illustrative figure

VBF: H→γγ

- Events selection
 - $N_{jets} \ge 2$ and $H \rightarrow \gamma \gamma$ final states
 - Two BDTs: VBF v.s. ggF and VBF $H \rightarrow \gamma \gamma$ v.s. continuum



Optimal Observable

- Events further categorized into six OO bins
 - VBF Higgs signal extracted by a simultaneous fit on diphoton mass



$$|\mathcal{M}|^{2} \neq |\mathcal{M}_{\rm SM}|^{2} + 2 \cdot c_{i} \cdot \operatorname{Re}(\mathcal{M}_{\rm SM}^{*}\mathcal{M}_{\rm CP-odd})|^{2}$$
$$+ c_{i}^{2} \cdot |\mathcal{M}_{\rm CP-odd}|^{2}.$$

Interpretation

1. Interference only:

factor out the impact on overall cross section and only sensitive to the CP violation.

Optimal Observable

- Events further categorized into 6 OO bins
 - VBF Higgs signal extracted by a simultaneous fit on diphoton mass



$$|\mathcal{M}|^{2} = |\mathcal{M}_{SM}|^{2} + 2 \cdot c_{i} \cdot \operatorname{Re}(\mathcal{M}_{SM}^{*}\mathcal{M}_{CP\text{-odd}}) + c_{i}^{2} \cdot |\mathcal{M}_{CP\text{-odd}}|^{2}.$$

Interpretation

2. Interference + quadrature:

For completeness, to compare with other analyses

Constrain CP-odd in HVV coupling

- CP-odd operators in two EFT bases constrained
 - HISZ basis: \tilde{d} ; Warsaw basis: $c_{h\tilde{W}}$
 - \tilde{d} further constrained by combining with H $\rightarrow \tau \tau$ previous results



- Results compatible with SM and no CP violation observed
- The most stringent constraints on CP-violating effects in the HVV coupling so far

Constrain CP-odd in HVV coupling

	68% (exp.)	95% (exp.)	68% (obs.)	95% (obs.)
\tilde{d} (inter. only)	[-0.027, 0.027]	[-0.055, 0.055]	[-0.011, 0.036]	[-0.032, 0.059]
\tilde{d} (inter.+quad.)	[-0.028, 0.028]	[-0.061, 0.060]	[-0.010, 0.040]	[-0.034, 0.071]
\tilde{d} from $H \to \tau \tau$	[-0.038, 0.036]	_	[-0.090, 0.035]	-
Combined \tilde{d}	[-0.022, 0.021]	[-0.046, 0.045]	[-0.012, 0.030]	[-0.034, 0.057]
$c_{H\tilde{W}}$ (inter. only)	[-0.48, 0.48]	[-0.94, 0.94]	[-0.16, 0.64]	[-0.53, 1.02]
$c_{H\tilde{W}}$ (inter.+quad.)	[-0.48, 0.48]	[-0.95, 0.95]	[-0.15, 0.67]	[-0.55, 1.07]

- CP-odd operators in two EFT bases constrained
 - \tilde{d} in HISZ basis: 95% confidence interval(CI), achieved for the first time
 - $c_{h\widetilde{W}}$ in Warsaw basis: 95% CI, five times better than $H \rightarrow \gamma \gamma$ differential measurement (arXiv:2202.00487)
- Combination with previous $H \rightarrow \tau \tau$ results on \tilde{d}

CPV in Yukawa coupling

• CP-odd HVV coupling is high dimension operators, usually arisen at one loop, e.g. 2HDM prediction:

$$\left[\mathcal{L}_{\text{eff}} \supset -\frac{\tilde{g}_{hZZ}}{2} hZ_{\mu\nu}\tilde{Z}^{\mu\nu} - \tilde{g}_{hWW} hW^+_{\mu\nu}\tilde{W}^{-\mu\nu} \right] \quad \begin{bmatrix} \tilde{g}_{hZZ} \simeq -\frac{\sin\alpha_2}{\tan\beta} \frac{1}{6 \times 10^5 \text{ GeV}} \\ \tilde{g}_{\mu\nu\nu} \simeq -\frac{\sin\alpha_2}{10} \frac{1}{6 \times 10^5 \text{ GeV}} \end{bmatrix}$$

- LHC constraint
 - (137 fb⁻¹, CMS PAS HIG-19-009) $ilde{g}_{hZZ} \lesssim rac{1}{3 imes 10^3 {
 m GeV}}$
- CP-odd in Higgs Yukawa coupling to fermion Hff coupling
 - Can be arisen at tree level in, e.g. 2HDM

$${\cal L}_{
m Yuk} \supset -rac{m_f}{v} \left(\kappa_f ar{f}f + i ilde{\kappa}_f ar{f} \gamma_5 f
ight) h$$

 $g_{nWW} = \tan \beta 5 \times 10^5 \text{ GeV}$

CP structure of H→ττ coupling

- CP-mixing of Higgs to τ lepton yukawa coupling
 - ϕ^*_{CP} : angle between **two t decay planes** encodes spin correlation
 - Propagated to the tau decay products

$$\mathcal{L}_{Y} = -\frac{m_{\tau}}{V} \kappa_{\tau} (\cos \phi_{\tau} \bar{\tau} \tau + \sin \phi_{\tau} \bar{\tau} i \gamma_{5} \tau) h \qquad \text{1. CP mixing parametrized by } \varphi_{\tau}$$

$$\frac{d\Gamma}{(H \to \tau^{+} \tau^{-}) \sim 1 - s_{\tau} \bar{s}_{\tau} + R(2\alpha^{H\tau\tau}) \cdot s_{+} \bar{s}_{+} \sim 1 - \frac{\pi^{2}}{2} b(E^{+}) b(E^{-}) \cdot \cos(\phi_{CP} - 2\alpha^{H\tau\tau})}$$

16



 $d\phi_{CP}$

2. ϕ_{τ} measured as modulation, or phase shift, in ϕ^*_{CP} distribution.

How to detect tau lepton

- Tau lepton
 - Heaviest lepton: 1.777 GeV, Decay length: ct \approx 87 μ m
 - ~ 35% to e or μ, ~ 65% to hadrons
- Experimental signature
 - Tau decay to e or μ looks like prompt e or μ
 - τ_{had} has a narrower energy deposit at calorimeter and finite number of tracks



τ_{had} identification

• Multi-Variate Analysis to combine all detector information



ATL-PHYS-PUB-2019-033

Tau decay mode classification

• Two methods to build τ decay plane



Tau decay mode classification

• Identify π^0 in hadron collider



Deep NN: Deep Set





Efficiency migration matrix



True Tau Decay Mode 23

Overview of $H \rightarrow \tau \tau$ analysis

- $Br(H \rightarrow \tau \tau) \approx 6\%$
- 2 final states: $\tau_{lep}\tau_{had}$, $\tau_{had}\tau_{had}$



23%, $\tau_h \tau_\mu$

q

Production modes: ggF and VBF



$H \rightarrow \tau \tau$: VBF

- VBF: tag two forward high p_T jets
 - High di-jet invariant mass: m_{ii}
 - Large separation: $\Delta \eta_{ii}$





Illustrative Figure

$H \rightarrow \tau \tau$: ggH

- VBF: tag two forward high p_T jets
- ggH: larger color factor than quark initiated $Z \rightarrow \tau \tau$
 - more high p_T gluon radiation



Color factor:
$$C_g = 3 \text{ v.s. } C_q = 4/3$$

Event categorization

• Improve $H \rightarrow \tau \tau$ significance

VBF		Boost			
$\begin{array}{l} p_{\mathrm{T}}^{j_2} > 30 \ \mathrm{GeV} \\ m_{jj} > 400 \ \mathrm{GeV} \\ \Delta \eta_{jj} > 3.0 \\ \eta_{j_1} \cdot \eta_{j_2} < 0 \\ \mathrm{Central} \ \tau\text{-leptons} \end{array}$		Not VBF $p_{\rm T}^{\tau\tau} > 100 {\rm ~GeV}$			
Signal region $(110 < m_{\tau\tau}^{\rm MMC} < 150 {\rm GeV})$					
VBF_1	VBF_1 VBF_0		Boost_0		
BDT(VBF) > 0 BDT(VBF) < 0		$ \begin{vmatrix} \Delta R_{\tau\tau} < 1.5 \text{ and} \\ p_{\rm T}^{\tau\tau} > 140 \text{ GeV} \end{vmatrix} $	$\begin{vmatrix} \Delta R_{\tau\tau} > 1.5 \text{ or} \\ p_{\rm T}^{\tau\tau} < 140 \text{ GeV} \end{vmatrix}$		
$Z \to \tau \tau$ control regions (60 < $m_{\tau \tau}^{\rm MMC}$ < 110 GeV)					
VBF_1 Z CR	VBF_0 Z CR	Boost_1 Z CR	Boost_0 Z CR		

CP sensitive observable

• ϕ^*_{CP} : angle between two tau decay planes, sensitive to ϕ_{τ}



Event categorization

- Improve significance on ϕ_{τ}
 - Impact Parameter significance: $d_0^{sig} = \frac{d_0}{\sigma}$
 - τ spin-analyzing functions: y_{+}^{ρ/a_1} from ρ or a1 meson decays

$y_{\pm}^{\rho} = (E_{\pi^{\pm}} - E_{\pi^{0}}) / (E_{\pi^{\pm}} + E_{\pi^{0}})$		$y_{\pm}^{a_1} = \frac{E_{2\pi} - E_{\pi_1^{\pm}}}{E_{2\pi} + E_{\pi_1^{\pm}}} - \frac{E_{2\pi} - E_{\pi_1^{\pm}}}{E_{2\pi} - E_{\pi_1^{\pm}}}} - \frac{E_{2\pi} - E_{\pi_1^{\pm}}}{E_{\pi_1^{\pm}}} - \frac{E_{\pi_1^{\pm}}}{E_{\pi_1^{\pm}}} - \frac{E_{\pi_1^{\pm}}}}{E_{\pi_1^{\pm}}} - \frac{E_{\pi_1^{\pm}}}{E_{\pi_1^{\pm}}} - \frac{E_{\pi_1^{\pm}}}{E_{\pi_1^{\pm}}}} - \frac{E_{\pi_1^{\pm}}}{E_{\pi_1^{\pm}}} - \frac{E_{\pi_1^{\pm}}}}{E_{\pi_1^{\pm}}} - \frac{E_{\pi_1^{\pm}}}{E_{\pi_1^{\pm}}} - \frac{E_{\pi_1^{\pm}}}}{E_{\pi_1^{\pm}}}} - \frac{E_{\pi_1^{\pm}}}{E_{\pi_1^{\pm}}} - \frac{E_{\pi_1^{\pm}}}}{E_{\pi_1^{\pm}}} - \frac{E_{\pi_1^{\pm}}}{E_{\pi_1^{\pm}}}} - \frac{E_{\pi_1^{\pm}}}{E_{\pi_1^{\pm}}}} - \frac{E_{\pi_1^{\pm}}}}{E_{\pi_1^{\pm}}}} - \frac{E_{\pi_1^{\pm}}}{E_{\pi_1^{\pm}}}} - \frac{E_{\pi_1^{\pm}}}}{E_{\pi_1^{\pm}}}} - E$	$-\frac{m_{3\pi}^2 - m_{\pi_1^\pm} + m_{2\pi}^2}{2m_{3\pi}^2}$	
Channel	Signal region	Decay mode combination	Selection criteria	
$\rm High \\ \tau_{had} \tau_{had} \\ Medium$		1 p0n - 1 p0n	$\begin{aligned} d_0^{\text{sig}}(\tau_1) &> 1.5\\ d_0^{\text{sig}}(\tau_2) &> 1.5 \end{aligned}$	
	High	1 p0n - 1 p1n	$\begin{split} d_0^{\rm sig}(\tau_{\rm 1p0n}) &> 1.5 \\ y^{\rho}(\tau_{\rm 1p1n}) &> 0.1 \end{split}$	
		1p1n–1p1n	$ y^{\rho}(\tau_1)y^{\rho}(\tau_2) > 0.2$	
		1p0n–1pXn	$\begin{aligned} d_0^{\rm sig}(\tau_{\rm 1p0n}) &> 1.5\\ y^{\rho}(\tau_{\rm 1pXn}) &> 0.1 \end{aligned}$	
	Medium	1p1n–1pXn	$ y^{\rho}(\tau_{1\text{pln}})y^{\rho}(\tau_{1\text{pXn}}) > 0.2$	
		1p1n-3p0n	$ert y^ ho(au_{1 ext{pln}}) ert > 0.1 \ ert y^{a_1}(au_{3 ext{p0n}}) ert > 0.6$	
	Low	All above	Not satisfying selection criteria	20

100² 100² 1 100²

CP structure of $H \rightarrow \tau \tau$ coupling

- Analysis strategy:
 - Simultaneous fit on $\phi^*_{\ CP}$ to extract the CP mixing angle



CP structure of $H \rightarrow \tau \tau$ coupling



- Observed (expected) $\phi_{\tau} \approx 9^{\circ} \pm 16^{\circ}$ ($0^{\circ} \pm 28^{\circ}$) at the 68% CL
 - Results compatible with SM expectation within uncertainties
 - Excluded pure CP-odd state at 3.4σ significance

Summary

- CP violation is one of the most interesting and profound topics in particle physics, where Higgs sector is a new field to look into
- Two new results from ATLAS experiment presented
 - for CP in the Higgs coupling to vector bosons and fermions.
 - New techniques: decay plane building, Optimal Observable (Matrix Element method), MVA method widely used in CPV
- Diverse theoretical frameworks: EFT based, Amplitude, etc.
- Both analyses still statistically limited. More data in Run 3 will definitely help

Backup



Study HVV CP in $H \rightarrow ZZ$ decay

- Parameterize in terms of cross section fractions, f_{ai}
- For the V=W,Z,

$$f_{ai} = \frac{|a_i|^2 \sigma_i}{|a_1|^2 \sigma_1 + |a_2|^2 \sigma_2 + |a_3|^2 \sigma_3 + |\kappa_1|^2 \sigma_{\Lambda 1} + |\kappa_1^{Z\gamma}|^2 \sigma_{\Lambda 1}^{Z\gamma}} \operatorname{sgn}\left(\frac{a_i}{a_1}\right)$$

• Four fractions: f_{a2} , f_{a3} , $f_{\Lambda 1}$, and $f_{\Lambda 1}^{Z\gamma}$



Studies of ggH vertex in $H \rightarrow WW$

CP properties of the Higgs boson in the effective Higgs – gluon coupling

- Production mode: ggH
 - In the large top quark mass limit: the CP structure of the topquark Yukawa coupling is inherited by the effective Higgs – gluon interaction
 - Assumption: SM-like HVV coupling
- Theoretical framework: EFT with Higgs Characterization

$$\mathcal{L}_{0}^{\text{loop}} = -\frac{g_{Hgg}}{4} \left(\kappa_{gg} \cos(\alpha) G^{a}_{\mu\nu} G^{a,\mu\nu} + \kappa_{gg} \sin(\alpha) G^{a}_{\mu\nu} \tilde{G}^{a,\mu\nu} \right) H$$

Effective coupling CP-mixing angle

- Target: constrain the CP-mixing angle *α*
- Sensitive observable: signed $-\Delta \phi_{jj}$ Signed difference in ϕ of the leading and subleading jets for events with at least two jets
 - Shape sensitive to CP effects



Chiara Arcangeletti

10

H

a

gb

Studies of ggH vertex in $H \rightarrow WW$

- Select ggF +2 jets events
- BDT to discriminate $H \rightarrow WW$ signal from the main backgrounds
- Build 12 categories in the 2D space BDT vs $\Delta \eta_{ii}$
 - BDT score split: maximize the signal/bkg ratio
 - $\Delta \eta_{ii}$ split: separation between different CP hypotheses for $\Delta \phi_{ii}$ increase at high $\Delta \eta_{ii}$
 - Perform a fit on signed- $\Delta \phi_{ii}$



arXiv:2109.13808v1





Reconstruct di-tau mass

units

Arbitrary

- Likelihood based method to recover neutrinos
 - Simulate kinematics: $Prob(\Delta R, p_{\tau})$
 - Include measured information

$$\mathcal{P}(\not\!\!\!E_{\Gamma_{x,y}}) = \exp\left(-\frac{(\Delta \not\!\!\!E_{\Gamma_{x,y}})^2}{2\sigma^2}\right)$$

Final LLH



CP structure of $H \rightarrow \tau \tau$ coupling

ATLAS Run: 350144 Event: 1545345207 2018-05-13 02:47:13 CEST

HVV CP in $H \rightarrow ZZ$



As already mentioned above, we analyze all major 1- and 3-prong tau decay modes:

$$\tau \to l + \mathbf{v}_l + \mathbf{v}_\tau \,, \tag{3}$$

$$\tau \to a_1 + \nu_\tau \to \pi + 2\pi^0 + \nu_\tau \,, \tag{4}$$

$$\tau \to a_1^{L,T} + \nu_\tau \to 2\pi^\pm + \pi^\mp + \nu_\tau \,, \tag{5}$$

$$\tau \to \rho + \nu_{\tau} \to \pi + \pi^0 + \nu_{\tau} \,, \tag{6}$$

$$\tau \to \pi + \nu_{\tau} \,. \tag{7}$$

We call the decay mode $\tau \rightarrow a_1^{L,T} + v_{\tau}$ in (5) also '1-prong', because the 4-momentum of a_1^{\pm} can be obtained from the measured 4-momenta of the 3 charged pions. The longitudinal (*L*) and transverse (*T*) helicity states of the a_1 resonance can be separated by using known kinematic distributions [22–25].