



华南师范大学
SOUTH CHINA NORMAL UNIVERSITY



LPC

Lattice Parton
Collaboration

TMD Wave Functions of Pion from Lattice QCD

arXiv:2302.09961; JHEP.08.172(2023)

华俊 华南师范大学

2023.10.08

第三届中国格点量子色动力学研讨会

Out Line



01 Motivation & Recent Progress

02 Soft Function

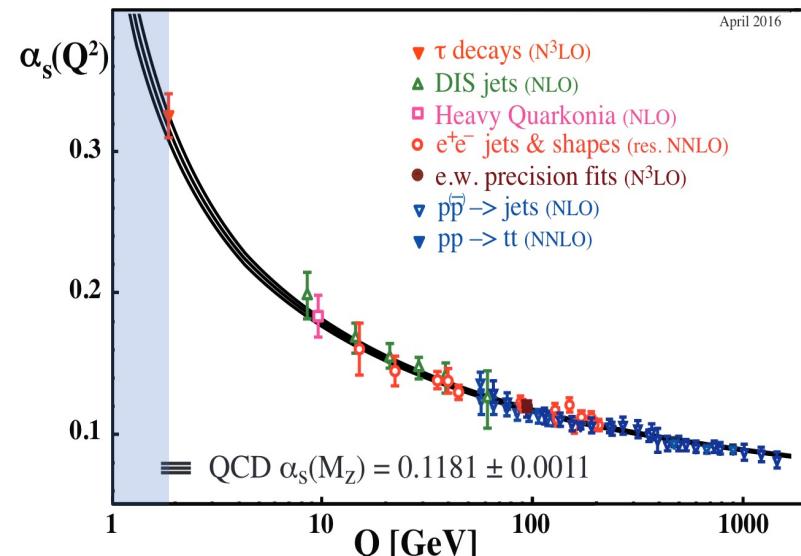
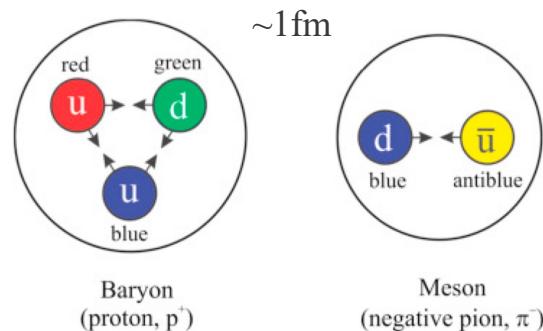
03 TMDWFs by LaMET

04 Outlook and Summary





Limited by QCD confinement



➤ QCD factorization (1982)

$$\left| e(l') \cdot e(l) \cdot h(p) \right|^2 \approx \left| e(l') \cdot e(l) \cdot xp, k_T \right|^2 \otimes \left| xp, k_T \cdot h(p) \right|^2$$

Measured cross section

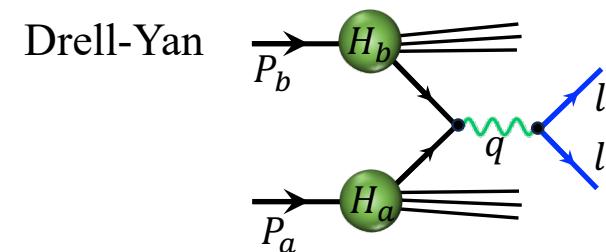
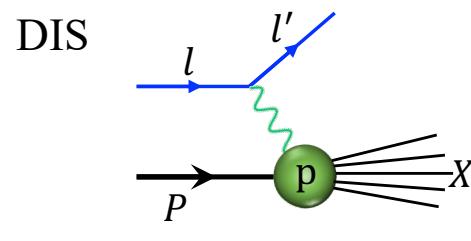
Asymptotic freedom
(calculated hard kernel)

Color confinement
(Parton in hadron)

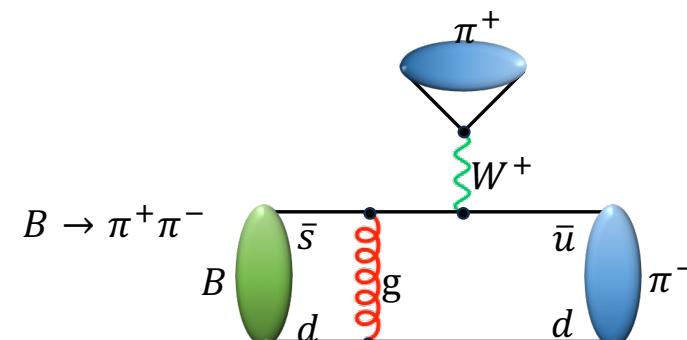
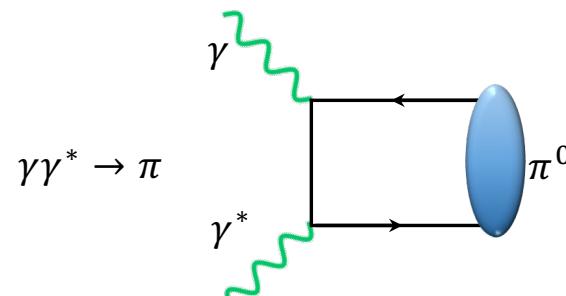
parton density
amplitude
(PDF, LCDA)



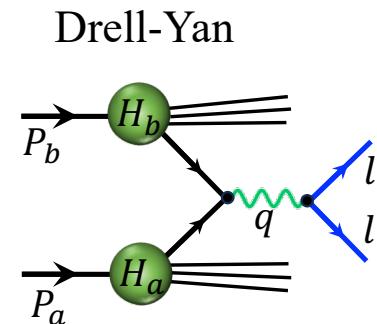
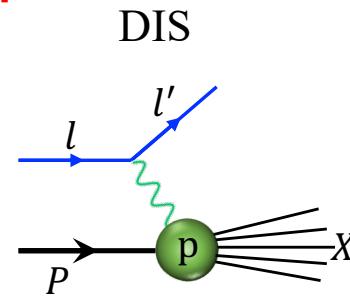
- ◆ PDFs: the probability distribution of partons (quarks and gluons) within a hadron — Inclusive process



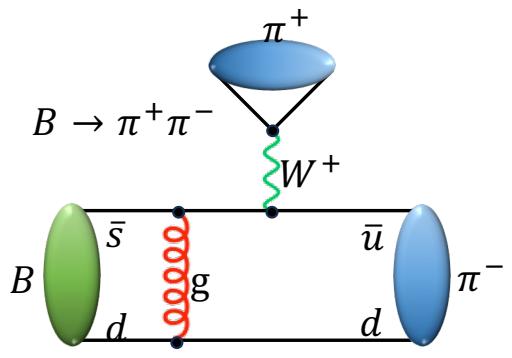
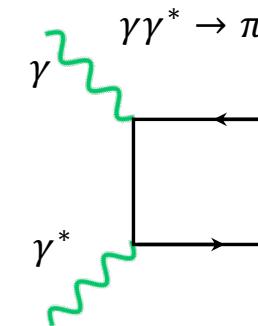
- ◆ LCDAs: the probability amplitude for partons within a hadron — Exclusive process



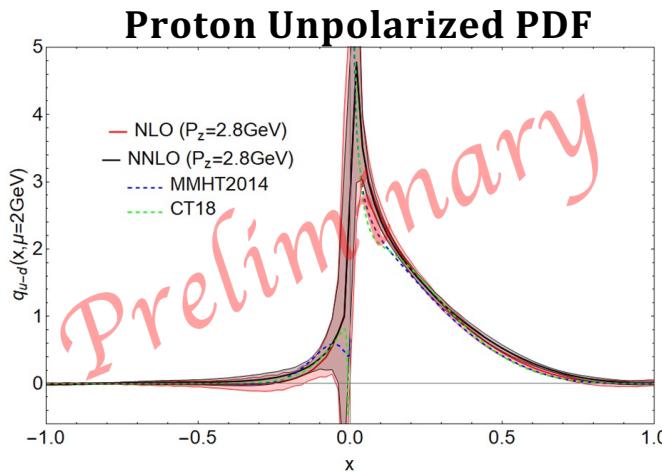
➤ PDF



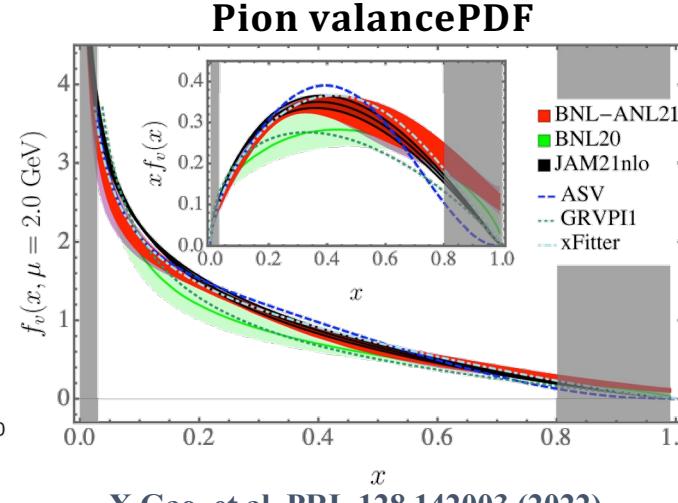
➤ LCDA



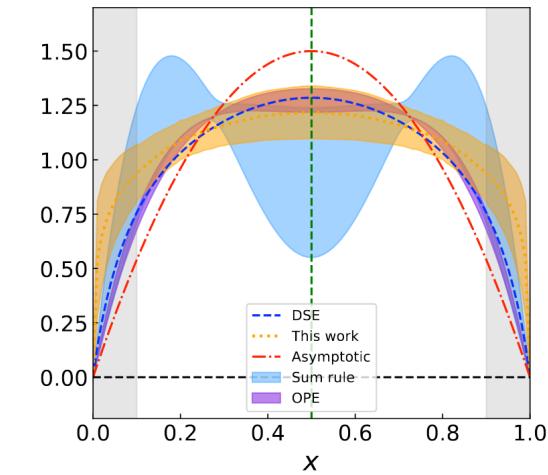
Pion LCDA



Preliminary result from LPC



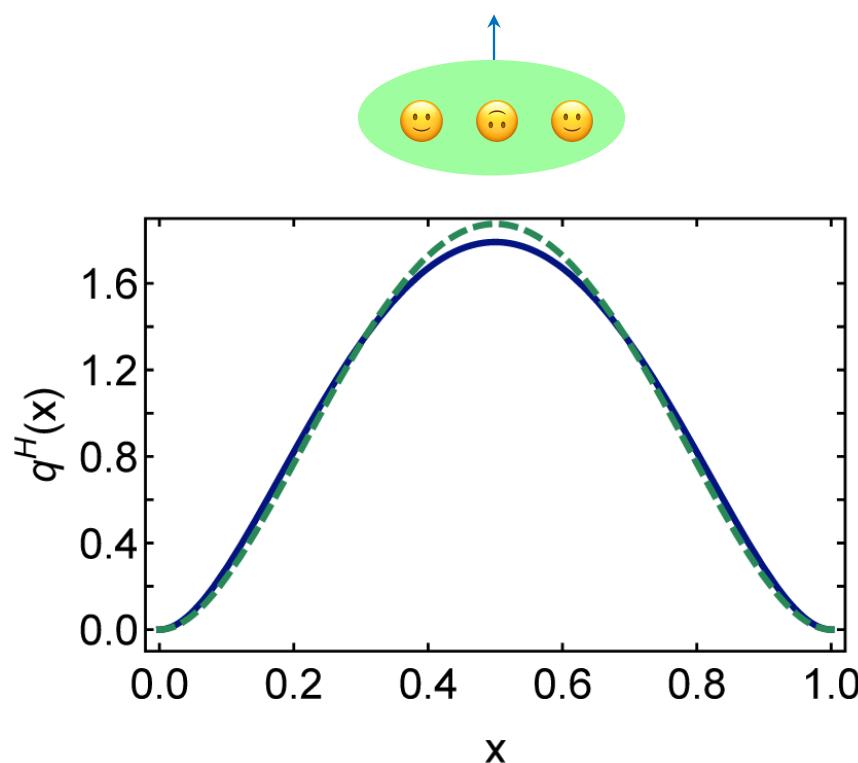
X.Gao, et.al. PRL.128.142003 (2022)



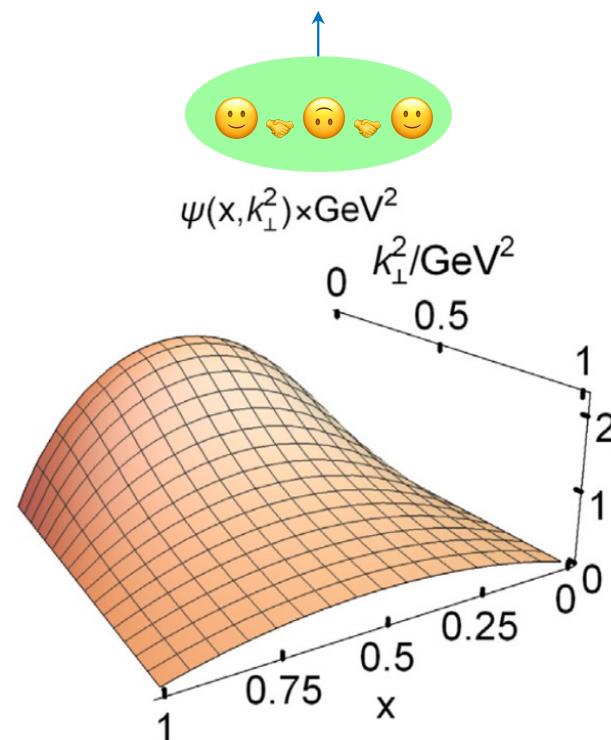
(LPC) J.Hua et.al. PRL.129,132001 (2022)



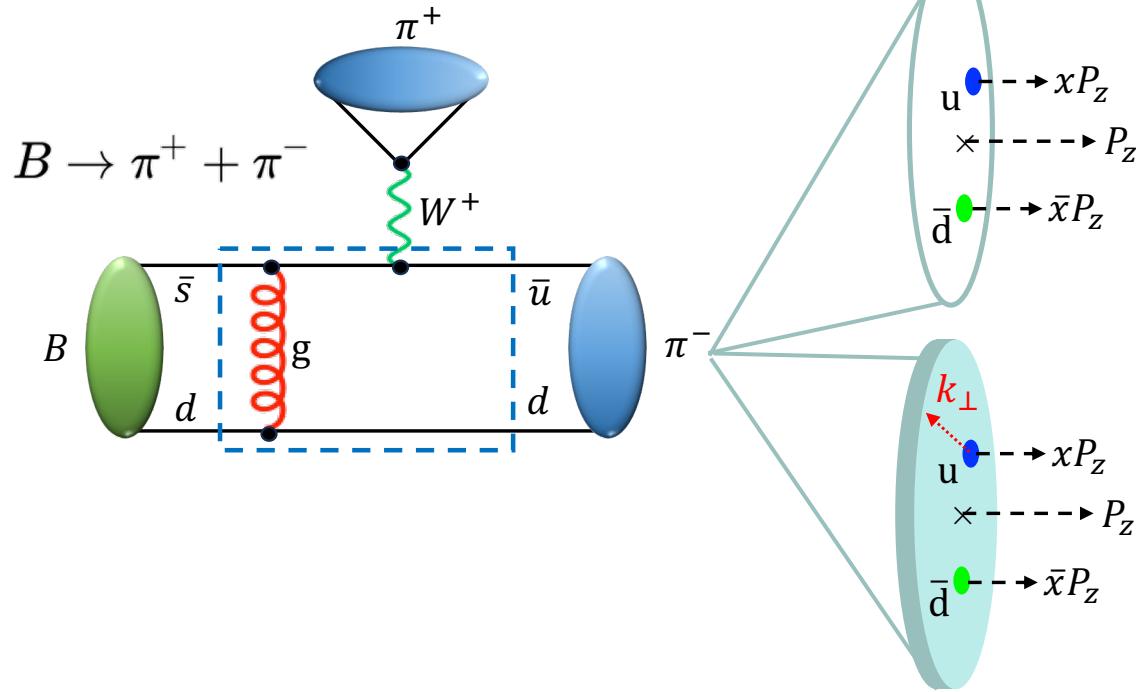
➤ One dimensional LCDA



➤ Three dimensional TMDWF

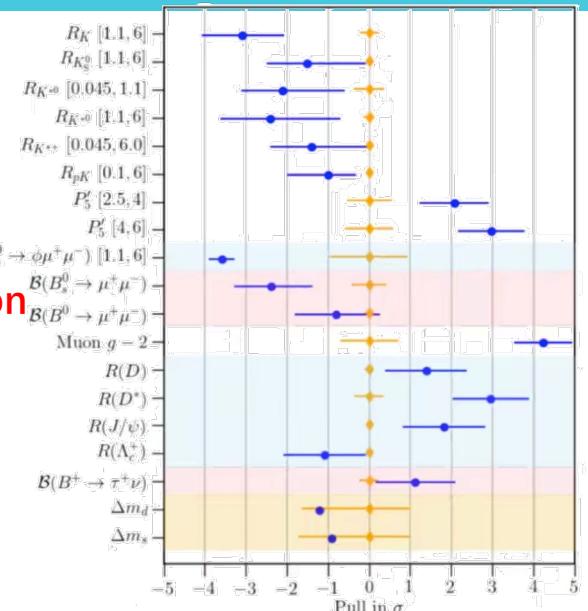


C.D.Roberts et.al. PPNP.120, 138883 (2021)



$$\mathcal{A} = \langle \pi^+ \pi^- | H | B \rangle \sim \int d^4 k_1 d^4 k_2 d^4 k_3 \text{Tr}[C(t) \psi_B(p_1) \psi_{\pi^+}(p_2) \psi_{\pi^-}(p_3) H(k_1, k_2, k_3, t)]$$

Collinear Factorization



TMD Factorization

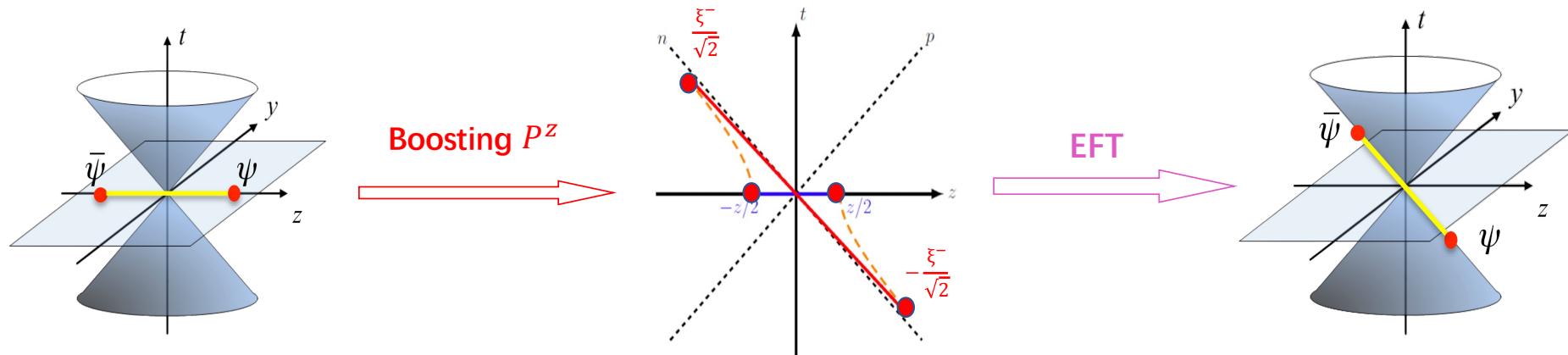
Y.Y.Keum et.al. PLB. 504, 6(2001)

C.D.Lu et.al. PRD.63,074009(2001)

A.Ali et.al. PRL.76, 074018(2007)



➤ Large Momentum Effective Theory:



LaMET is capable for Entire x dependence distributions



Recent Progress on LCDA

R.Zhang et.al. PRD. 125, 094519(2020)

Pion LCDA with 3 lattice spacings

X.Ji et.al. NPB.964,115311(2021)

A hybrid renormalization scheme

(LPC) J.Hua et.al. PRL.127, 062002(2021)

K^*, ϕ LCDA at physical with hybrid

(LPC) Y.K.Huo et.al. NPB. 969, 115443 (2021)

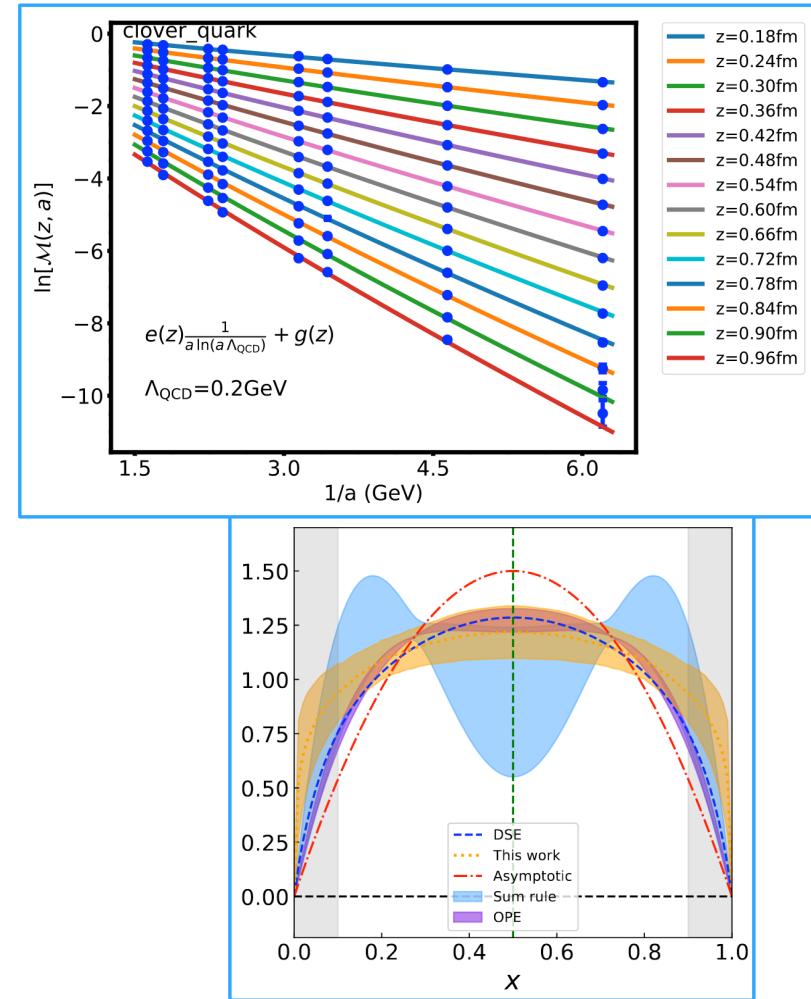
Solve linear divergence by self renormalization

(LPC) J.Hua et.al. PRL.129,132001 (2022)

π, K LCDA with self renormalization

R.Zhang et.al. RPB.844,138081(2023)

Resummation to improve endpoint region





Recent Progress on TMDWF

(LPC) Q.A. Zhang et.al. PRL. 125, 192001 (2020)
Soft function and CS kernel (First)

M. Schlemmer et.al. JHEP.08,004(2021)
CS kernel by different TMDs

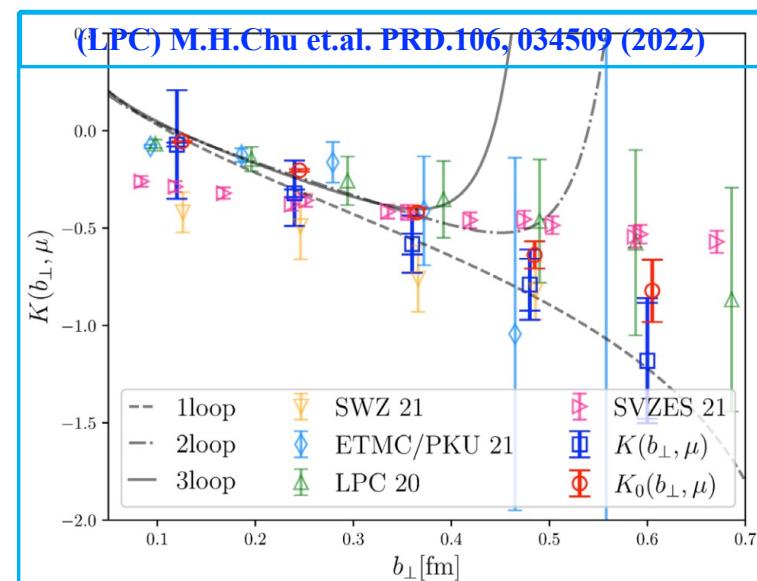
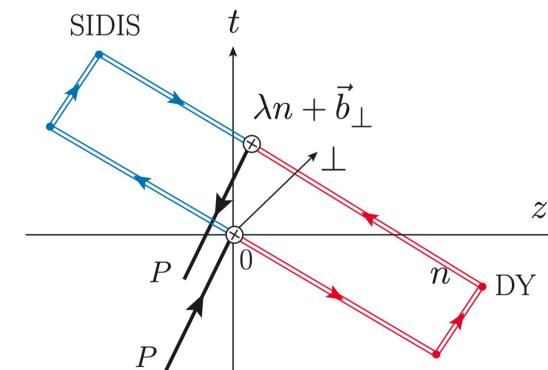
P. Shanahan et.al. PRD.104, 114502(2021)
CS kernel from quasi-TMDPDFs (1-loop)

L.Yuan, X.Feng et.al. PRL. 128, 062002 (2022)
Twists' effects on soft function

(LPC) K.Zhang PRL.129,082002 (2022)
Renormalization of TMDs on lattice

(LPC) M.H.Chu et.al. PRD.106, 034509 (2022)
CS kernel from quasi-TMDWFs (1-loop)

A. Avkhadieva et.al. arXiv2307.12359 (2023)
CS kernel at physical pion mass





☰ TMD Factorization in LaMET

➤ Multiplicative factorization of quasi-TMDWF in LaMET

$$\frac{\tilde{\Psi}^\pm(x, b_\perp, \mu, \zeta^z) S_I^{\frac{1}{2}}(b_\perp, \mu)}{H^\pm(x, \zeta^z, \mu) \exp\left[\frac{1}{2} K(b_\perp, \mu) \ln \frac{\pm \zeta^z + i\epsilon}{\zeta}\right]} = \Psi^\pm(x, b_\perp, \mu, \zeta) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{x \zeta}, \frac{M^2}{(P^z)^2}, \frac{1}{b_\perp^2 \zeta_z}\right)$$

X.D.Ji et.al. Rev.Mod.Phys. 93, 035005 (2021)

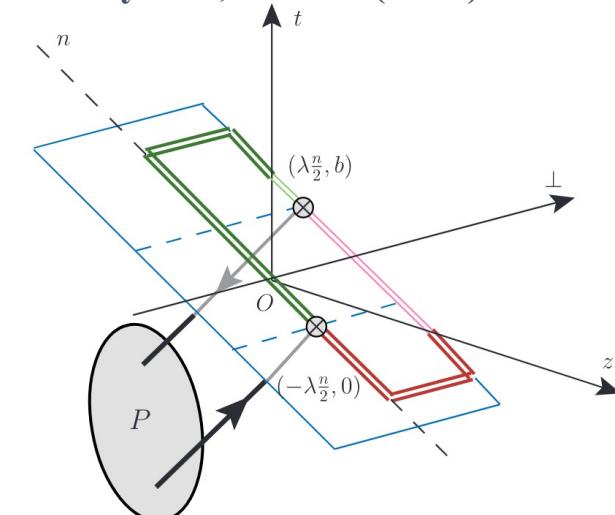
$\tilde{\Psi}^\pm(x, b_\perp, \mu, \zeta_z)$: Quasi-TMDWF,

$S_r(b_\perp, \mu)$: Intrinsic soft function,

$H^\pm(\zeta_z, \bar{\zeta}_z, \mu^2)$: Matching coefficient,

$K(b_\perp, \mu)$: Collins-Soper kernel,

$\Psi^\pm(x, b_\perp, \mu, \zeta)$: TMDWF.





Lattice ensembles

$L^3 \times T$	a (fm)	m_π^{sea} (MeV)	m_π^v (MeV)
$24^3 \times 64$	0.12	310	670
measurement			
1053×4			

- 2+1+1 flavors of HISQ action (MILC)
- Momenta: 1.72, 2.15, 2.58, 3.01GeV
- Coulomb gauge fixed wall source

$L^3 \times T$	a (fm)	m_π^{sea} (MeV)	m_π^v (MeV)
$48^3 \times 48$	0.098	333	662
measurement			
952×4			

- 2+1 flavors of Symanzik gauge action (CLS)
- Momenta: 1.58, 2.11, 2.64, 3.16GeV
- Coulomb gauge fixed wall source





≡ Soft Function by LaMET

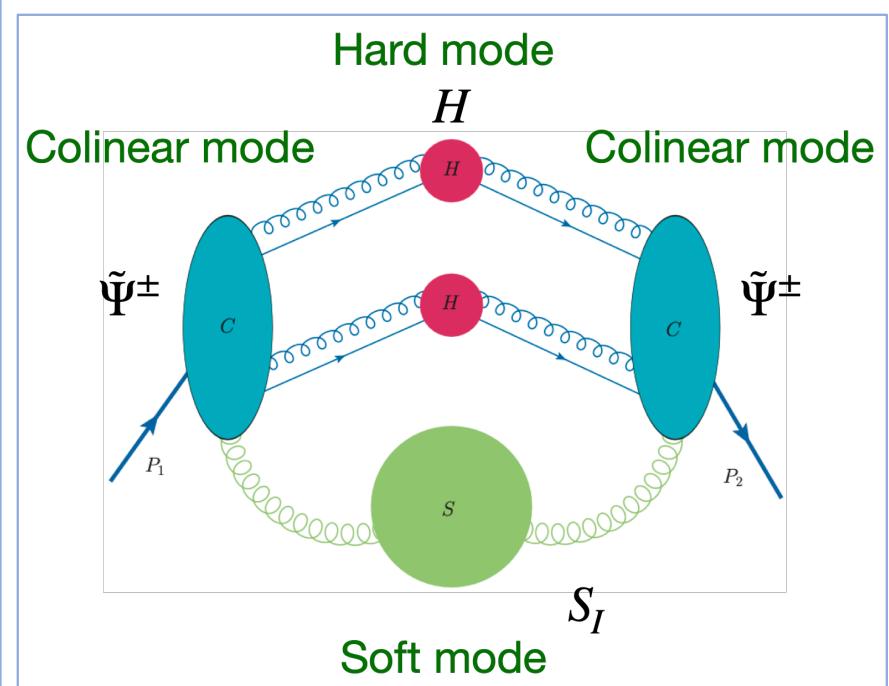
➤ Four quark form factor:

$$F(b_\perp, P_1, P_2, \Gamma, \mu) = \frac{\langle P_2 | \bar{q}(b_\perp) \Gamma q(b_\perp) \bar{q}(0) \Gamma' q(0) | P_1 \rangle}{\langle 0 | \bar{q}(0) \gamma^\mu \gamma^5 q(0) | P_1 \rangle \langle P_2 | \bar{q}(0) \gamma_\mu \gamma^5 q(0) | 0 \rangle}$$

Normalization factor: $f_\pi^2 P_1 P_2$

➤ Factorization of form factor:

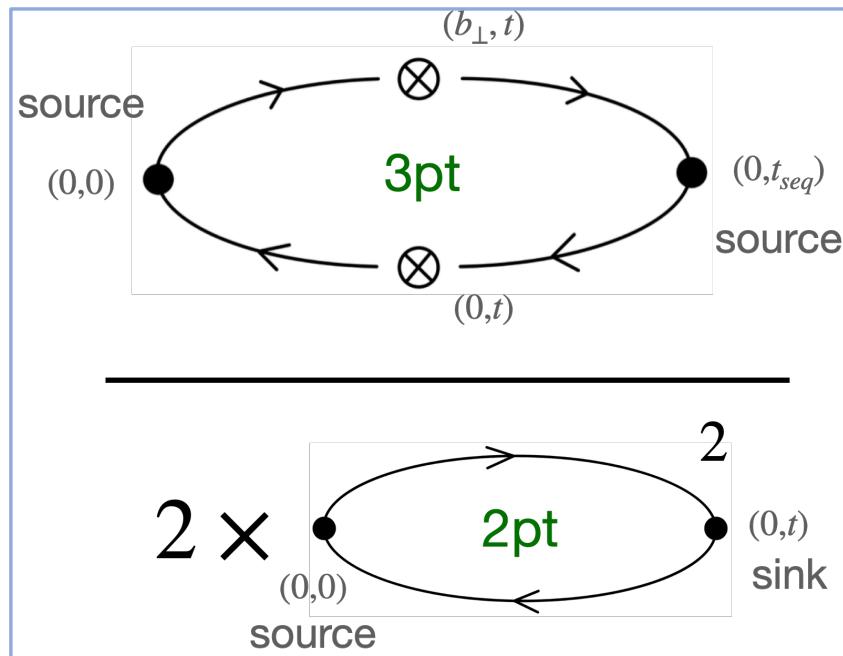
$$S_I(b_\perp, \mu) = \frac{F(b_\perp, P_1, P_2, \Gamma, \mu)}{\int dx_1 dx_2 H(x_1, x_2, \Gamma) \tilde{\Psi}^{\pm*}(x_2, b_\perp, \zeta^z) \tilde{\Psi}^\pm(x_1, b_\perp, \zeta^z)}$$



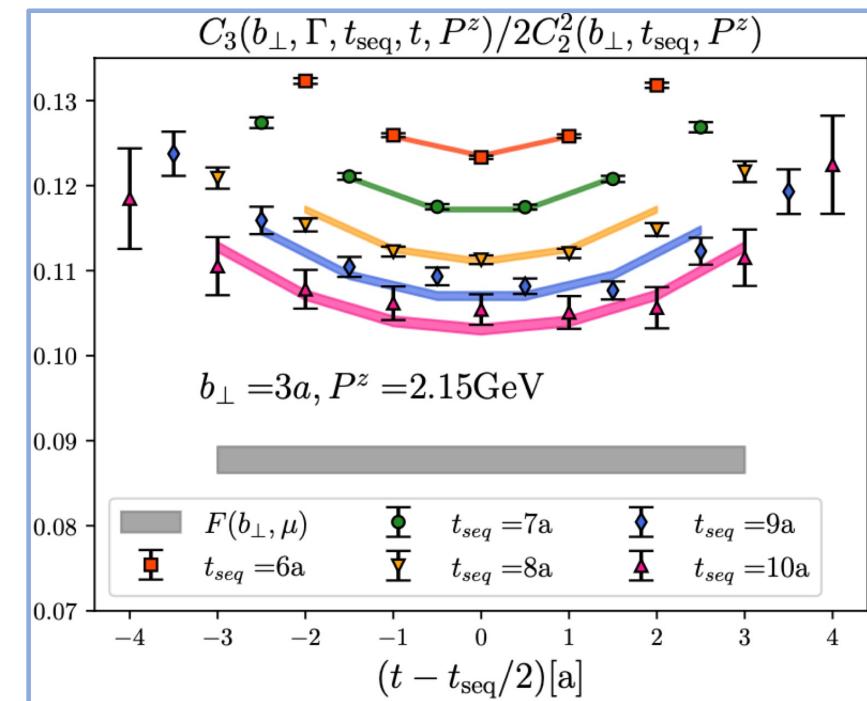


Soft Function by LaMET

$$\frac{C_3(b_\perp, \Gamma, t_{\text{sep}}, t, P^z)}{2C_2^2(t_{\text{sep}}/2, P^z)} = F(b_\perp, \Gamma, P^z) \frac{1 + c_1(e^{-\Delta Et} + e^{-\Delta E(t_{\text{sep}}-t)})}{1 + c_2 e^{-\Delta Et_{\text{sep}}/2}}$$



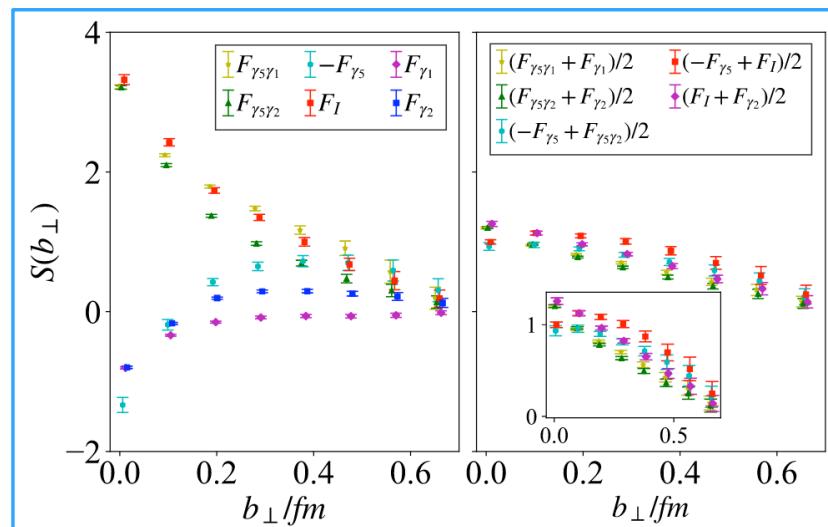
Joint fit for $t_{\text{seq}} = (7,8,9,10)$





Soft Function by LaMET

Operator mixing in Soft function:



Y.Li et.al. PRL.128, 062002 (2022)

By Fierz rearrangement analysis, these combination can suppress high twists contribution:

$$\begin{aligned}
 & \sum F(\Gamma = \gamma^\mu) + F(\Gamma = \gamma^\mu \gamma_5) \\
 &= (\bar{\psi}_a \gamma^{x,y} \psi_b)(\bar{\psi}_c \gamma_{x,y} \psi_d) + (\bar{\psi}_a \gamma^{x,y} \gamma_5 \psi_b)(\bar{\psi}_c \gamma_{x,y} \gamma_5 \psi_d) \\
 &= \bar{\psi}_c \gamma^\mu \gamma_5 \psi_b \bar{\psi}_a \gamma_\mu \gamma_5 \psi_d + \bar{\psi}_c \gamma^\mu \psi_b \bar{\psi}_a \gamma_\mu \psi_d \\
 & F(\Gamma = I) - F(\Gamma = \gamma_5) \\
 &= (\bar{\psi}_a \psi_b)(\bar{\psi}_c \psi_d) - (\bar{\psi}_a \gamma_5 \psi_b)(\bar{\psi}_c \gamma_5 \psi_d) \\
 &= \frac{1}{2} \bar{\psi}_c \gamma^\mu \gamma_5 \psi_b \bar{\psi}_a \gamma_\mu \gamma_5 \psi_d - \frac{1}{2} \bar{\psi}_c \gamma^\mu \psi_b \bar{\psi}_a \gamma_\mu \psi_d
 \end{aligned}$$

- The UV divergence in the I and γ_5 form factor. can be removed by the renormalization constant of scalar density operator Z.F.Deng et.al. JHEP.09, 046 (2022)

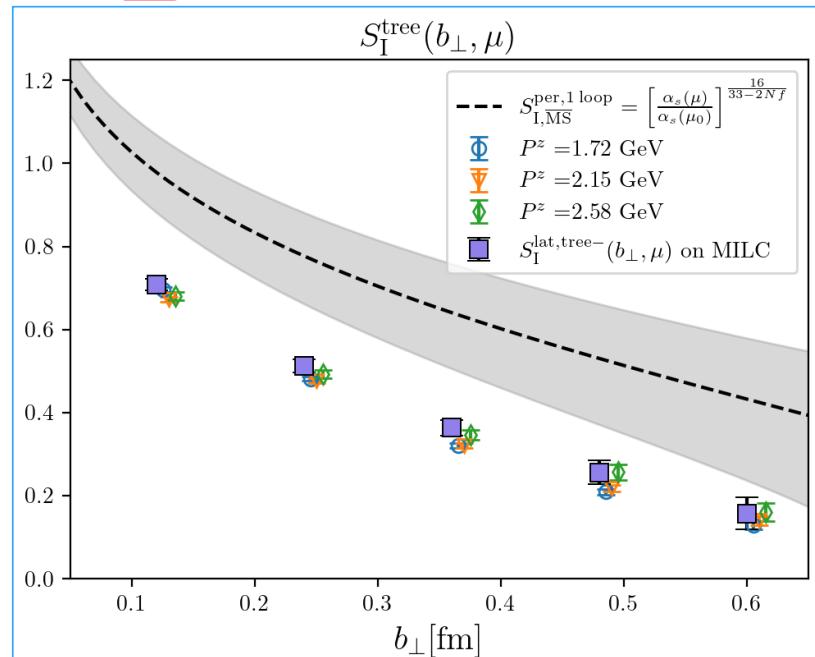
$$Z_S = 1 + \frac{\alpha_s C_F}{4\pi} \frac{3}{\epsilon_{UV}}. \quad (59)$$



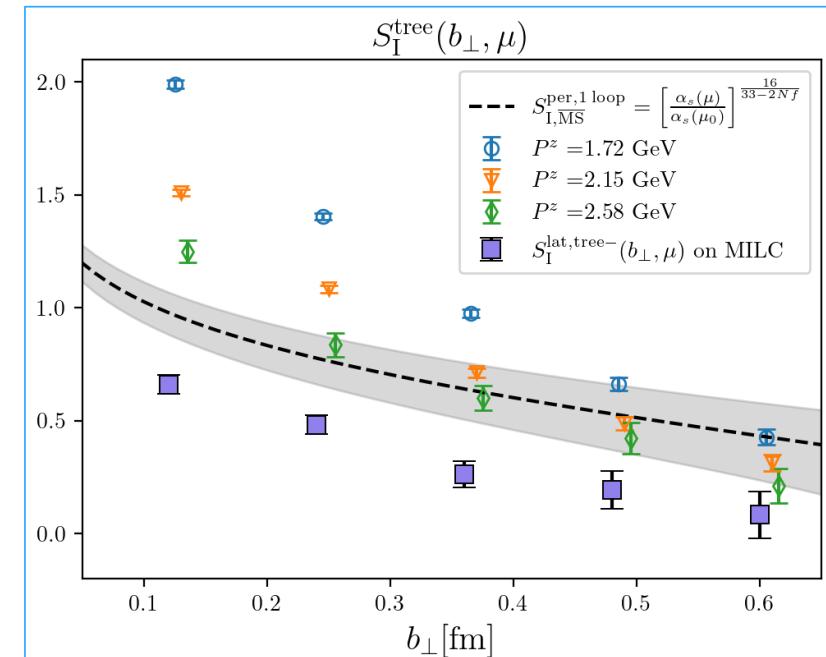
☰ Soft Function by LaMET

➤ Pz dependence of soft function for 2 combination:

$$\sum F(\Gamma = \gamma^\mu) + F(\Gamma = \gamma^\mu \gamma_5)$$



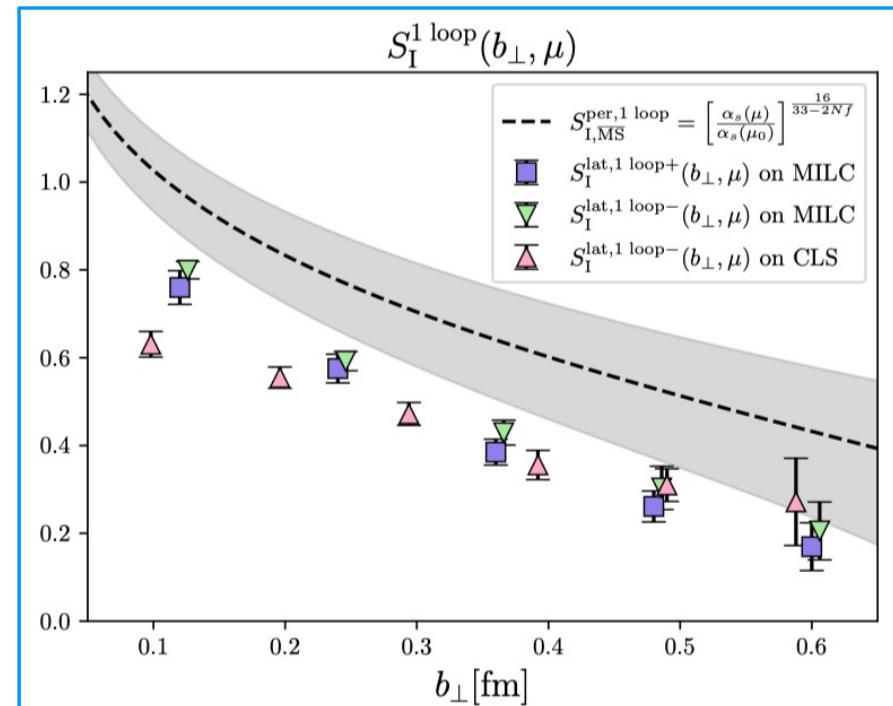
$$F(\Gamma = I) - F(\Gamma = \gamma_5)$$





☰ Soft Function by LaMET

- 1-loop matching soft function extracted by MILC and CLS ensemble
- Consistent for ‘+/-’ cases: soft function is universal
- Discrete effects are significant





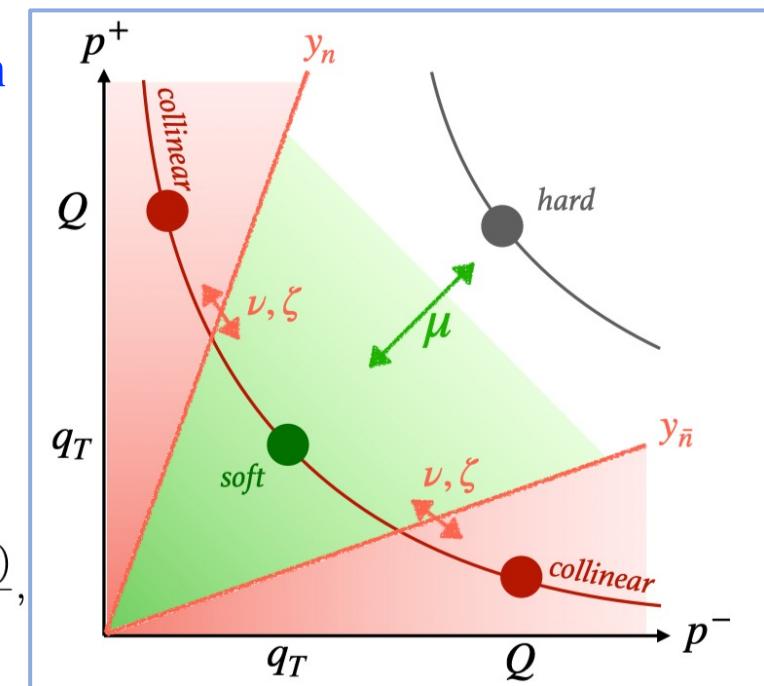
≡ CS-Kernel by LaMET

- Collins-Soper kernel: describe the evolution for rapidity scale:

$$2\zeta \frac{d}{d\zeta} \ln \Psi(x, b_\perp, \mu, \zeta) = K(b_\perp, \mu),$$

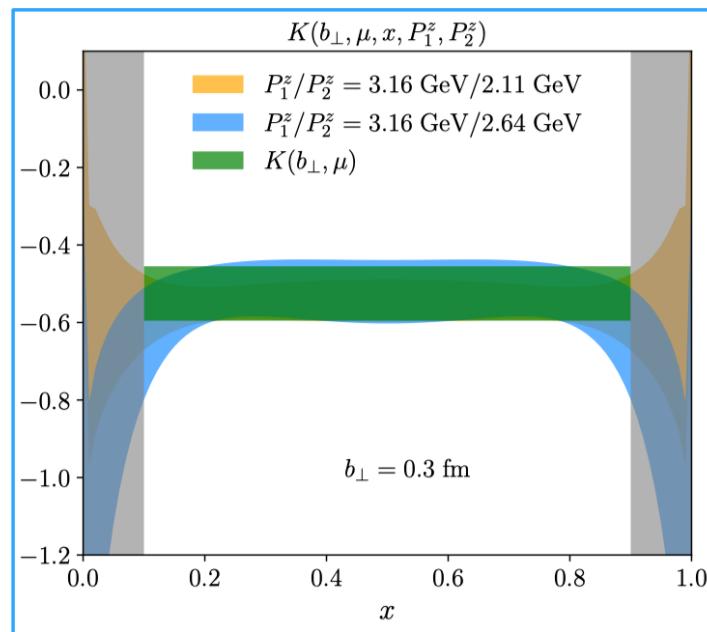
- In LaMET factorization, CS-kernel can be extracted by ratio:

$$K(b_\perp, \mu, x, P_1^z, P_2^z) = \frac{1}{\ln(P_1^z/P_2^z)} \ln \frac{H^\pm(xP_2^z, \mu)\tilde{\Psi}^\pm(x, b_\perp, \mu, P_1^z)}{H^\pm(xP_1^z, \mu)\tilde{\Psi}^\pm(x, b_\perp, \mu, P_2^z)},$$

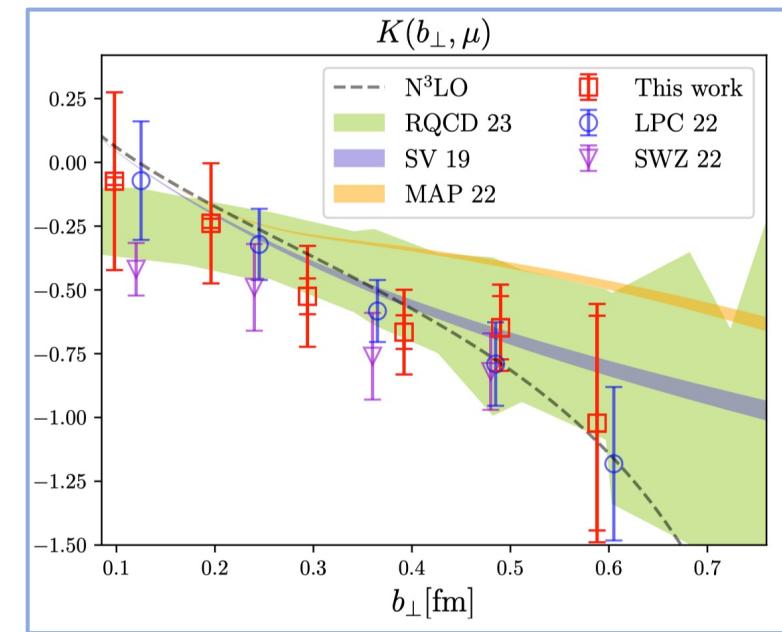


≡ CS-Kernel by LaMET

$$K(b_\perp, \mu, x, P_1^z, P_2^z) = \frac{1}{\ln(P_1^z/P_2^z)} \ln \frac{H^\pm(xP_2^z, \mu)\tilde{\Psi}^\pm(x, b_\perp, \mu, P_1^z)}{H^\pm(xP_1^z, \mu)\tilde{\Psi}^\pm(x, b_\perp, \mu, P_2^z)},$$



1-loop CS-kernel on CLS ensemble

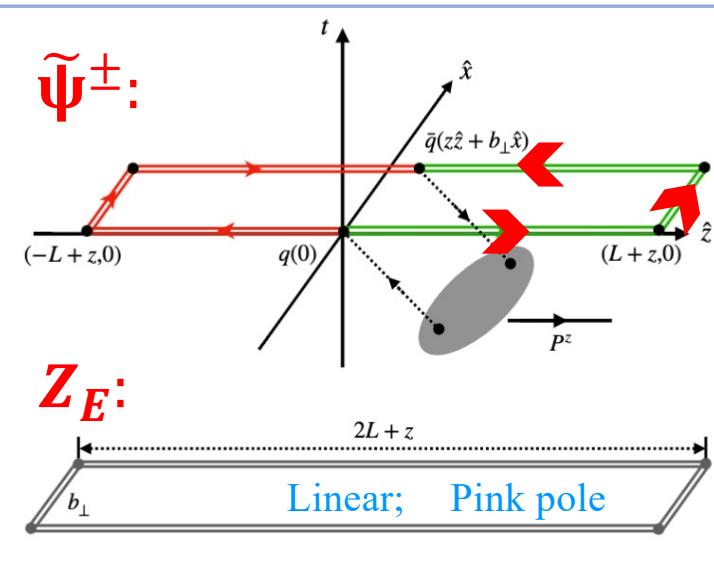




☰ Quasi TMDWFs

➤ Quasi TMDWF in Euclidean lattice:

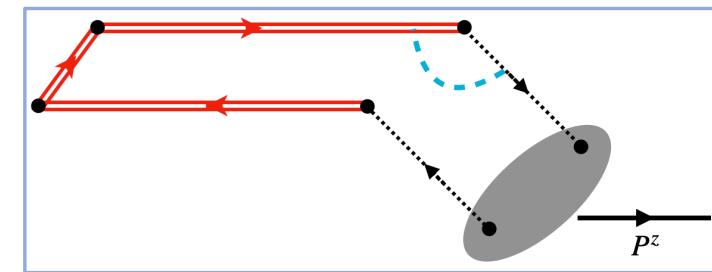
$$\tilde{\Psi}^{\pm}(x, b_{\perp}, \mu, \zeta^z) = \lim_{L \rightarrow \infty} \int \frac{P^z dz}{2\pi} e^{ixzP^z} \times \frac{\langle 0 | \bar{q}(z\hat{n}_z + b_{\perp}\hat{n}_{\perp}) \gamma^t \gamma_5 U_{c\pm} q(0) | \pi(P^z) \rangle}{\sqrt{Z_E(2L \pm z, b_{\perp}, \mu)} Z_O(1/a, \mu, \Gamma)}.$$



➤ Staple-shaped gauge-link:

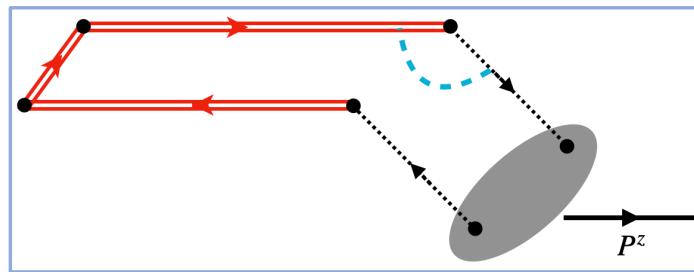
$$U_{c\pm} = U_z^\dagger(z\hat{n}_z + b_{\perp}\hat{n}_{\perp}; L) U_{\perp}(\pm L\hat{n}_z + z\hat{n}_z; b_{\perp}) \times U_z(0\hat{n}_z; \pm L + z).$$

Logarithm divergence





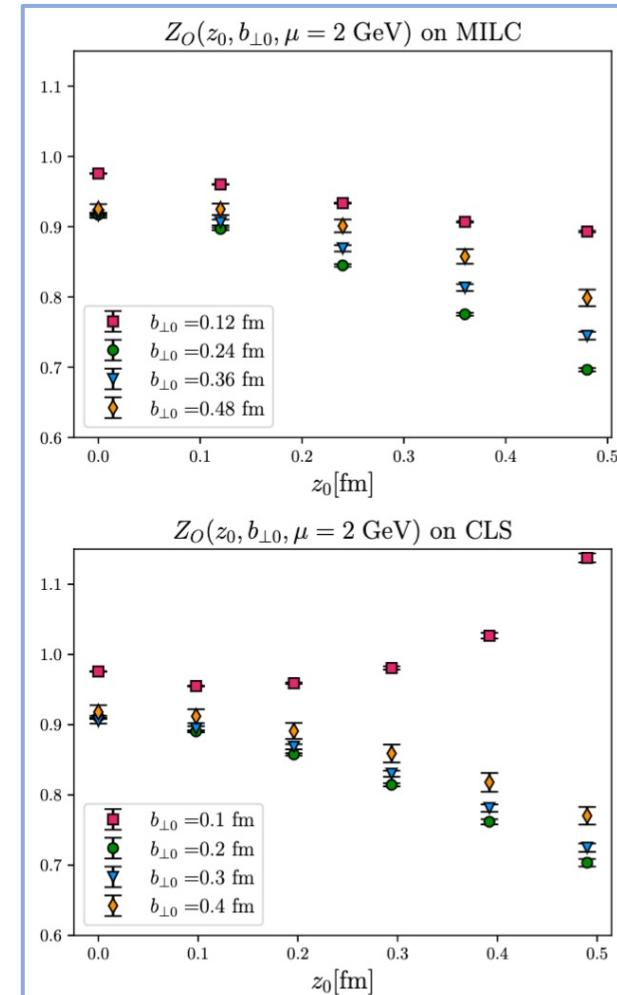
☰ Quasi TMDWFs



➤ Quark Wilson line vertex renormalization:

$$Z_O(1/a, \mu) = \frac{\tilde{\Psi}^{\pm,0}(z_0, b_{\perp 0}, \zeta^z = 0, L)}{\sqrt{Z_E(2L + |z_0|, b_{\perp 0}, \mu)} \tilde{\psi}^{\overline{\text{MS}}}(z_0, b_{\perp 0}, \mu)}.$$

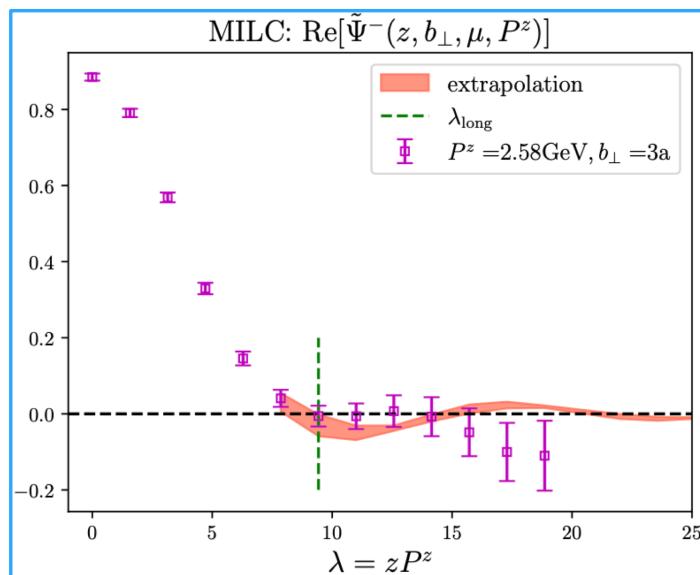
(LPC) K.Zhang PRL.129,082002 (2022)





☰ Quasi TMDWFs

$$\tilde{\Psi}^\pm(x, b_\perp, \mu, \zeta^z) = \lim_{L \rightarrow \infty} \int \frac{P^z dz}{2\pi} e^{ixzP^z} \frac{\langle 0 | \bar{q}(z\hat{n}_z + b_\perp \hat{n}_\perp) \gamma^t \gamma_5 U_{c\pm} q(0) | \pi(P^z) \rangle}{\sqrt{Z_E(2L \pm z, b_\perp, \mu)} Z_O(1/a, \mu, \Gamma)}$$



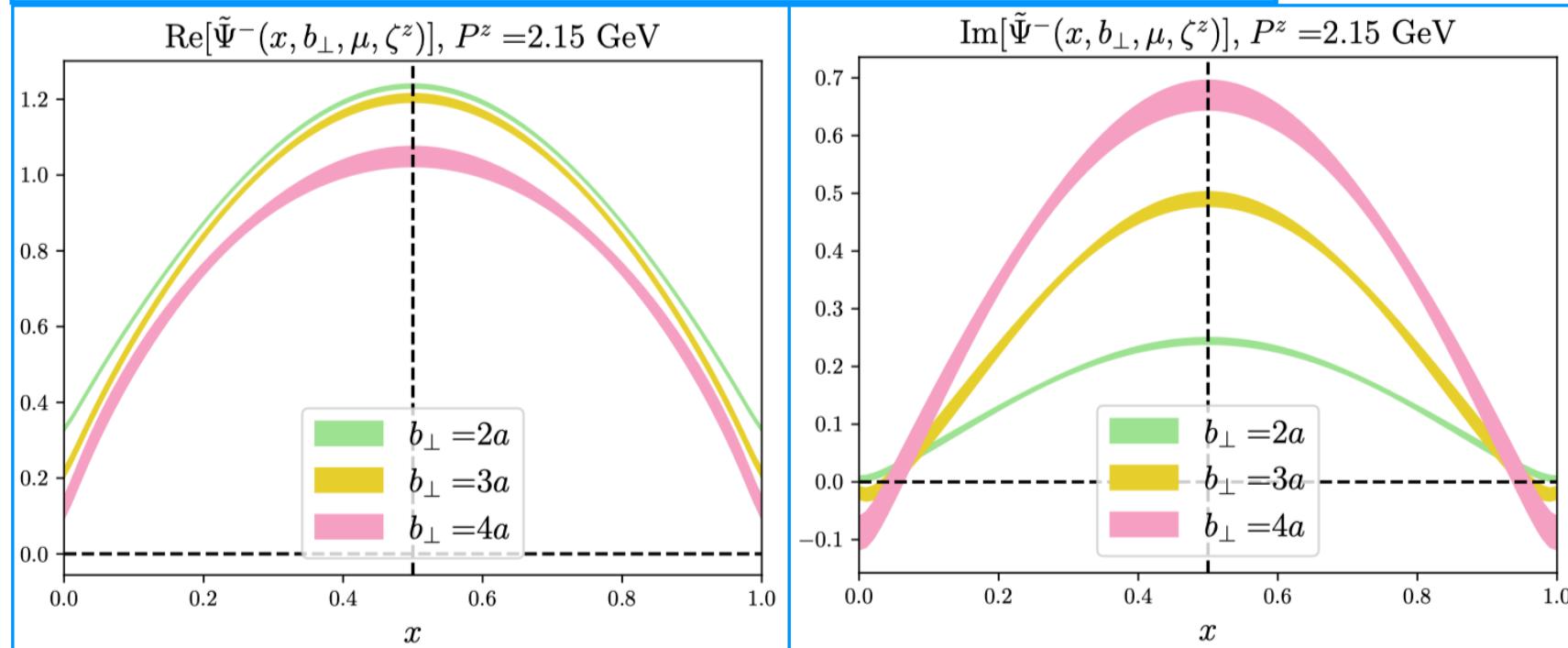
- Quasi TMDWF in coordinate space and extrapolation in large λ
- Physical based parameterization :

$$\tilde{\Psi}(z, b_\perp, \mu, P^z) = f(b_\perp) \left[\frac{c_1}{(-i\lambda)^d} + e^{i\lambda} \frac{c_2}{(i\lambda)^d} \right] e^{-\frac{\lambda}{\lambda_0}}$$



☰ Quasi TMDWFs

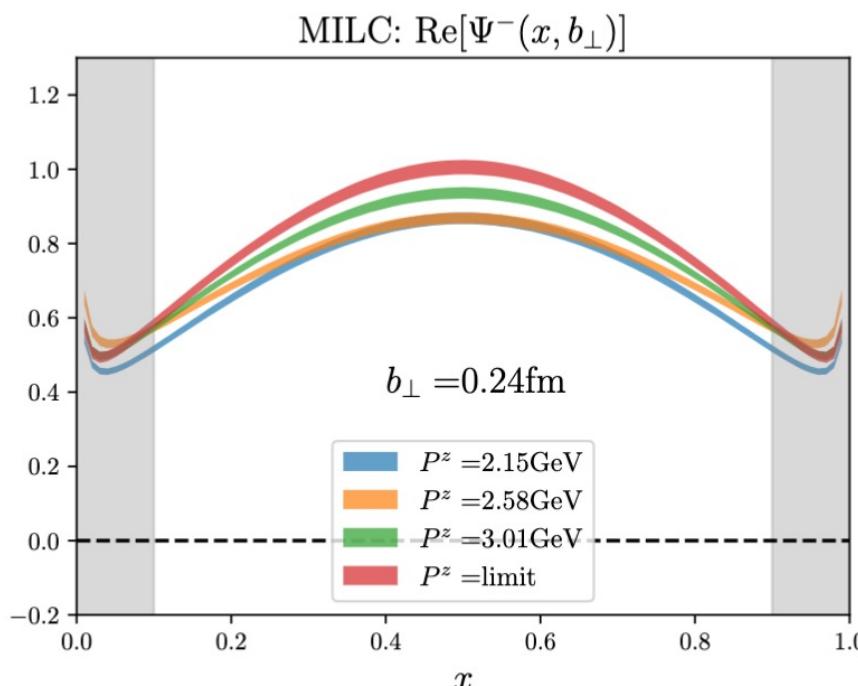
$$\tilde{\Psi}^\pm(x, b_\perp, \mu, \zeta^z) = \lim_{L \rightarrow \infty} \int \frac{P^z dz}{2\pi} e^{ixzP^z} \frac{\langle 0 | \bar{q}(z \hat{n}_z + b_\perp \hat{n}_\perp) \gamma^t \gamma_5 U_{c\pm} q(0) | \pi(P^z) \rangle}{\sqrt{Z_E(2L \pm z, b_\perp, \mu)} Z_O(1/a, \mu, \Gamma)}$$





☰ TMDWFs by LaMET

$$\tilde{\Psi}^\pm(x, b_\perp, \mu, \zeta^z) S_I^{\frac{1}{2}}(b_\perp, \mu) = H^\pm(x, \zeta^z, \mu) \exp\left[\frac{1}{2}K(b_\perp, \mu) \ln \frac{\pm \zeta^z + i\epsilon}{\zeta}\right] \Psi^\pm(x, b_\perp, \mu, \zeta)$$



$$\begin{aligned} H^\pm(x, \zeta^z, \mu) \\ = 1 + \frac{\alpha_s C_F}{4\pi} \left(-\frac{5\pi^2}{6} - 4 + l_\pm + \bar{l}_\pm - \frac{1}{2} (l_\pm^2 + \bar{l}_\pm^2) \right) \\ l_\pm = \ln[(-x\zeta^z \pm i\epsilon)/\mu^2] \end{aligned}$$

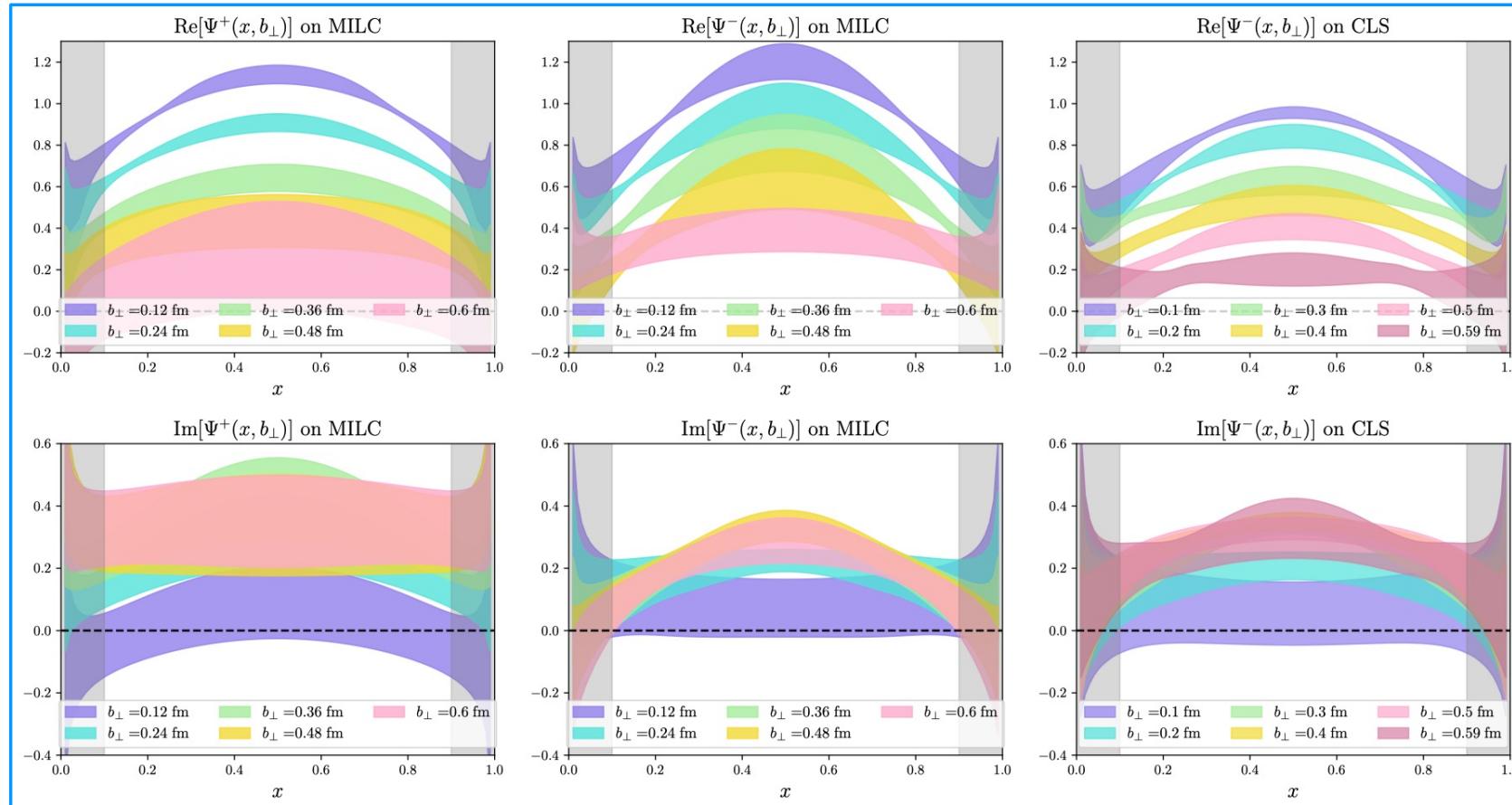
➤ Pz dependence of TMDWF after matting

- **Pz extrapolation:**

$$\Psi^\pm(x, P_z) = \Psi^\pm(x, P_z \rightarrow \infty) + \frac{c_2(x)}{P_z^2} + \mathcal{O}\left(\frac{1}{P_z^4}\right)$$

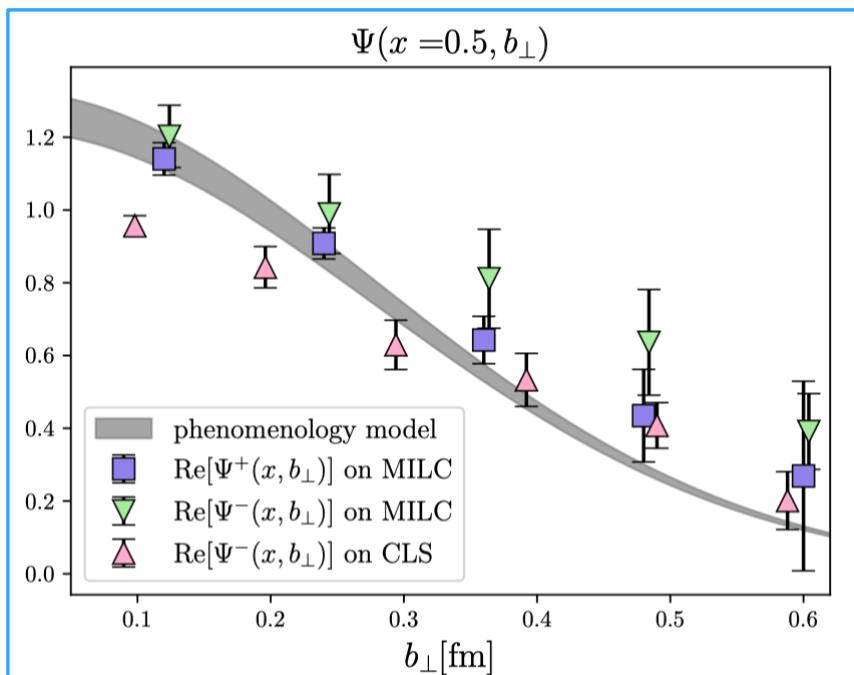


TMDWFs by LaMET





☰ TMDWFs by LaMET



- Decay behavior of b_\perp at $x = 0.5$
- A comparison with a phenomenological model at $x = 0.5$

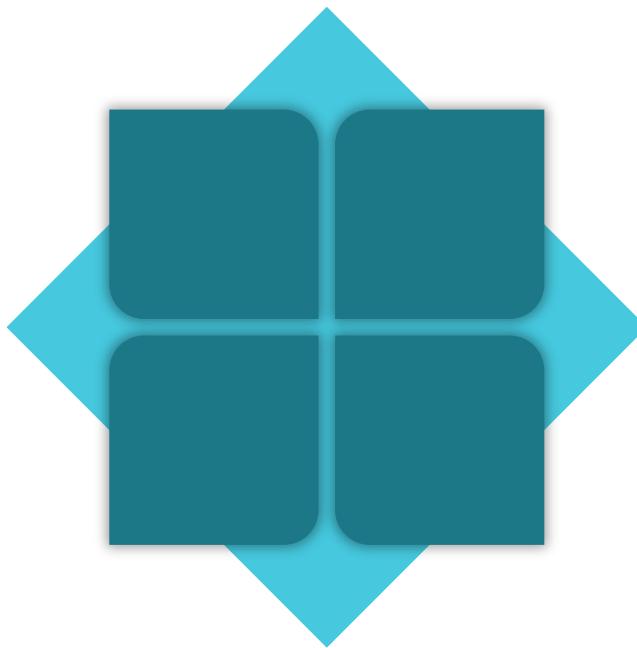
$$\Psi(x, b_\perp) = 6x(1-x) \left[1 + \frac{3}{2} a_2^\pi (5(2x-1)^2 - 1) \right] \exp \left[-\frac{x(1-x)b_\perp^2}{4a^2} \right]$$

C.D.Lv et.al. PRD75,094020 (2007)

Summary

*. We calculate the one-loop intrinsic soft function and TMDWF with LaMET on MILC and CLS ensembles.

*. The MILC and CLS results show good agreement, but discrete errors are still relatively significant in current results.



*. Fierz rearrangement analysis can be adopted to suppress high twist's effect in soft function.

*. Future calculations with more b_\perp on smaller lattice spacings are necessary to get more complete TMDWF results.

Thanks for your attentions!