Collaborative research center CRC 110

"Symmetries and the emergence of structure in QCD"



Hi-Lites from the German side

- Third funding period -

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Deutsche Forschungsgemeinschaft



CRC 110 - 3rd FP - Rizhao meeting – July 21, 2023 \cdot O < \wedge ∇ > \triangleright (

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A.1: Analysis of rescattering effects in 3π final states

Stamen, Isken, Kubis, Mikhasenko, Niehus, Eur. Phys. J. C 83 (2023) 510

• nontrivial rescattering effects in 3π decays dominated by $\pi\rho$:

 $J^{PC} = 0^{--}$ [exotic] 1^{--} [ω, ϕ] 1^{-+} [π_1] 2^{++} [a_2]

 \rightarrow 1⁻⁺ and 2⁺⁺ relevant for COMPASS

tool: Khuri–Treiman equations

necess. events in Dalitz plot N

to reject naïve isobar model

(\doteq undistorted ρ lineshape)

at 5σ significance?



related experimental analysis: in progress

- first observation of avoided level crossing in 3-particle resonance on the lattice
- model system: complex φ^4 theory

$$\begin{split} \mathcal{L} &= \sum_{i=0,1} \left[\frac{1}{2} \partial^{\mu} \varphi_{i}^{\dagger} \partial_{\mu} \varphi_{i} + \frac{1}{2} m_{i}^{2} \varphi_{i}^{\dagger} \varphi_{i} + \lambda_{i} (\varphi_{i}^{\dagger} \varphi_{i})^{2} \right] \\ &+ \frac{g}{2} \varphi_{1}^{\dagger} \varphi_{0}^{3} + \text{h.c.} \,. \end{split}$$

- comparison: resonance (red) versus no resonance (gray) scenario at g = 8.87
- excellent agreement between different finite volume formalism



[[]M. Garofalo et al., JHEP 02 (2023) 252, arXiv:2211.05605]

A3: BORN OPPENHEIMER EFT (BOEFT) CALCULATION of HYBRIDS DECAY



BOEFT prediction for hybrids multiplets (boxes) in comparison to neutral exotic charmonium states (crosses)

Results are obtained also for bottomonia exotics

This opens the way to a treatment of ALL XYZ exotics in the BOEFT framework: we are currently addressing **Tetraquarks**



BOEFT predictions for the hybrids to quarkonium decays in comparison to the total decay widths of the neutral exotic charmonium states

N. Brambilla, W.K. Lai, A. Mohapatra, A. Vairo *Phys.Rev.D* 107 (2023) 5, 054034





A5: θ -dependence of nuclei & Big Bang nucleosynthesis

Lee, UGM, Olive, Shifman, Vonk, Phys. Rev. Res. 2 (2020) 033392

- Does BBN give bounds on the elusive θ -parameter of QCD?
- Use chiral EFT and meson-exchange to θ -dependent study nuclear properties



 \hookrightarrow di-proton and di-neutron get bound for heta=0.7, 0.2

 \hookrightarrow triple-lpha reaction affected, lack of ¹⁶O for $heta \gtrsim 0.1$

Constraints on θ , but no anthropics: Universe looks similar as long as $\theta \leq 0.1$

Project A7: H. Dreiner & C.D. Lü

Detector



- Proton decays to Kaon + light long-lived neutralino via UDD
- Neutralino flys and decays in detector via LQD
- Signature: Kaon + Neutralino decay

Light long-lived neutralinos at colliders:

(1) *Phys.Rev.D* 103 (2021) 7,075013; (2) *JHEP* 03 (2021) 148; (3) *JHEP* 04 (2022) 057;

(4) SciPost Phys. 14 (2023) 134; (5) JHEP 02 (2023) 120; (6) e-Print: 2306.11803

• Supersymmetric light long-lived neutralino: novel proton decay mode



(to appear shortly)





A.8 Charmless Exclusive B Decays — Highlights M. Beneke (TUM)

I) Factorization theorem for charmless and *B* to charm hadronic two-body decays $B \rightarrow M_1 M_2(\gamma)$ including QED effects + analysis of branching fractions and CP asymmetry sum rules in πK final states

[2008.10615 (charmless) and 2107.03819 (heavy-light + semi-leptonic)]

II) Definition, renormalization and evolution of light-cone distribution amplitudes in QCD+QED for light mesons (π , . . .) and heavy mesons (*B*) [2108.05589 (light mesons) + 2204.09091 (heavy mesons)]

$$\Gamma(u, v; \mu) = -\frac{\alpha_{\rm em}Q_M}{\pi} \,\delta(u - v) \left(Q_M \left(\ln \frac{\mu}{2E} + \frac{3}{4} \right) - Q_d \ln u + Q_u \ln \bar{u} \right) \\ - \left(\frac{\alpha_s C_F}{\pi} + \frac{\alpha_{\rm em}}{\pi} Q_u Q_d \right) \left[\text{ ERBL } \right]_+ \qquad (Evolution \, kernel)$$

(III)

Mass-dependent LCDA of a heavy meson from the universal leading-twist HQET LCDA

[2305.06401]



A9: EFT for Nuclear Electroweak Currents

N³LO EM current operator: isoscalar part

Application: ⁴He charge radius: effective field theory and experiment

$$r_{C}(^{4}\text{He}) = r_{str}^{2}(^{4}\text{He}) + \left(r_{p}^{2} + \frac{3}{4m_{p}^{2}}\right) + r_{m}^{2}$$

Our prediction for ⁴He **charge** radius

 $r_C(^4$ **He**) = (1.6798 ± 0.0035) fm



Preliminary, using CODATA 2018 r_{p} and own determination of $r_{n}\ PDG22$



Our prediction for ⁴He charge radius is fully consistent with the muonic-atom spectroscopy

A10: Parity-violating Pion-Nucleon coupling h_{π}^1 from Lattice QCD



Exploratory Lattice QCD computation towards Parity-Violating NN π coupling h_{π}^1

• first-time calculation of all diagrams, including quark-loops "B" and "D", from all light and strange flavored 4-quark operators



• bare coupling from connected "W" diagram at $M_\pi \approx 260 \,\mathrm{MeV}$, lattice spacing $a \approx 0.9 \,\mathrm{fm}$ and box size $L \approx 3 \,\mathrm{fm}$

$$h_{\pi}^{1}(W, \text{bare}) = 8.08(98) \cdot 10^{-7}$$
 arXiv: 2306.03211 [hep - latt]

A.11: Dispersion relations for $B^- o \ell^- ar{ u}_\ell \ell' ar{\ell'}$

Kürten, Zanke, Kubis, van Dyk, Phys. Rev. D 107 (2023) 053006

- $B^- \rightarrow \ell^- \bar{\nu}_{\ell} \gamma$: access *B*-meson LCDA e.g. Beneke, Rohrwild 2011
- $B^- \to \ell^- \bar{\nu}_\ell \ell' \bar{\ell'}$: off-shell dependence of $B \to \gamma^*$ form factors LHCb
- goal: relate $B \to \gamma^*$ to $B \to \rho, \omega$ via dispersion relations
 - \rightarrow need form factor basis free of kinematic singularities / zeros

Bardeen, Tung 1968; Tarrach 1975

 \longrightarrow BTT generalised from electromagnetic to weak form factors

- *z*-expansion of $B \rightarrow V$ form factors Bharucha, Straub, Zwicky 2016
- branching ratios, forward–backward asymmetry, form factors:

Process	$\mathcal{B} imes 10^8$	A_{FB}	3	
$B^- \to e^- \bar{\nu}_e \mu^- \mu^+$	$3.19(43)_N(25)_{V_{ub}}$	$-0.358(31)_N$	$[d_{1}^{2}, q^{2})]$	
$B^- \to \mu^- \bar{\nu}_\mu e^- e^+$	$3.78(47)_N(30)_{V_{ub}}$	$-0.398(38)_N$	$\mathcal{F}_1(k_2)$	10
$B^- \to \tau^- \bar{\nu}_\tau e^- e^+$	$2.75(27)_N(22)_{V_{ub}}$	$-0.500(18)_N$		
$B^- \to \tau^- \bar{\nu}_\tau \mu^- \mu^+$	$1.77(23)_N(14)_{V_{ub}}$	$-0.458(15)_N$	$\frac{1}{g^2 [C_0 V^2]}$	

A.11: Dispersion relations for $B^- \rightarrow \ell^- \bar{\nu}_\ell \ell' \bar{\ell'}$ – p. 1

B1. New insights into the interpretation of hadronic form factors

 $\rho(r) = \int d^3q F(-\mathbf{q}^2) e^{-i\mathbf{q}\cdot\mathbf{r}} \quad \longleftarrow \text{ The standard "Breit-frame" density can$ **not**be interpreted as a charge density for light hadrons Burkardt, Miller, Jaffe,

\Rightarrow no consistent definition of local spatial densities possible?

Define the charge density of a spin-0 system in terms of a generic wave packet state:

$$\rho_{\phi}(\mathbf{r}) = \langle \phi, \mathbf{0} | \hat{\rho}(\mathbf{r}, 0) | \phi, \mathbf{0} \rangle = \int \frac{d^3 p \, d^3 p'}{2(2\pi)^3 \sqrt{EE'}} \, \phi^{\star}(\mathbf{p}') \underbrace{\langle p' | \hat{\rho}(\mathbf{r}, 0) | p \rangle}_{e^{-i(\mathbf{p}' - \mathbf{p}) \cdot \mathbf{r}} (E + E') F(q^2)} \phi(\mathbf{p})$$

For localized packets, $\rho_{\phi}(\mathbf{r})$ becomes independent of the wave packet state ϕ :

$$\rho(r) = \int \frac{d^3 q}{(2\pi)^3} e^{-i\mathbf{q}\cdot\mathbf{r}} \int_{-1}^{+1} d\alpha \, \frac{1}{2} F\left[(\alpha^2 - 1)\,\mathbf{q}^2\right] \quad \longleftarrow \text{ valid in the frame with } \langle \mathbf{p} \rangle = 0$$

Epelbaum, Gegelia, Lange, Meißner, Polyakov, PRL129 (2022)

The new definition is valid for any particle's mass, is consistent with special relativity and coincides with the known 2d-densities in the infinite-momentum frame.

Higher-spin systems and gravitational FFs: Panteleeva, Epelbaum, Gegelia, Meißner, PRD 106 (2022) 056019; e-Print: 2211.09596 (to appear in EPJC); e-Print: 2305.01491; Alharazin, Sun, Epelbaum, Gegelia, Meißner, JHEP 02 (2023) 163.





 \implies Three-body dynamics and isospin violation must matter

We solve
$$T = V + VGT$$
 with $V_{LO} = {}^{3}S_{1}$ $I_{3D_{1}}$ I_{3

@ LO: One free parameter (for each cut off) \rightarrow Width fixed by mass

• Excellent description of the data

B.3: $T_{cc}(3875)$

- Precision needs 3 body dynamics (problem: experimental resolution)
- $r = -2.4 \pm 0.9$ fm, but $r_0 = +1.4 \pm 0.9$ fm (corrected for isospin viol.)

 \implies T_{cc}^+ qualifies as isoscalar DD^* molecule

• OPE gives prominent left-hand cut $@M_{\pi} = 280 \text{ MeV}$

Effect should be visible in recent lattice data







B4: Three-particle analog of the Lellouch-Lüscher formula

F. Müller and A. Rusetsky, JHEP 03 (2021) 152, F. Müller, J.-Y. Pang, A. Rusetsky and J.-J. Wu, JHEP 02 (2023) 214.

• A three-particle analog of the Lüscher-Lellouch formula has beed derived, relating the finite- and infinite-volume three-particle decay matrix elements.

- The irregular volume-dependence in the matrix element is studied within the non-relativistic EFT. Short-distance effects are encoded in the effective couplings that feature only exponentially suppressed finite-volume corrections. The lattice calulations enable one to extract these couplings from the fit to the data.
- The framework can be formulated in a manifestly Lorentz-invariant form by choosing the quantization axis along the total four-momentum of the three-particle system: important for moving frames.

B.6: Charge-Symmetry Breaking of A=7 and 8 hypernuclei

Large CSB contributions due to long-ranged pion exchange

- use ${}^4_{\Lambda}He {}^4_{\Lambda}H$ to determine LECs
- first prediction of Λn scattering lengths Haidenbauer et al., FBS62, 105 (2021)
- significant uncertainty in input data

Application to A = 7 and 8

- benchmark confirms reliable predictions for SRG-evolved interactions up to 3-baryon level
- spin-dependence of CSB differs for $A = 7.8^{-1}$
- comparison of two scenarios CSB/CSB*
- consistent description of A = 4,7,8
- hypernuclei provide constraints on YN



Hoai Le et al., PRC107, 024002 (2023)





Evidence against a first-order phase transition in neutron star cores

CRC110, Project B7

Len Brandes, Wolfram Weise and Norbert Kaiser [arXiv:2306.06218]

► Bayesian inference of sound speed $c_s^2 = \partial P / \partial \varepsilon$ inside neutron stars based on available data:

[Brandes, Weise and Kaiser, PRD 107 (2023)]

extreme evidence

- Shapiro time-delays
- NICER X-ray measurements
- ► Gravitational waves (GW) of binary neutron star mergers

- ChEFT results at small densities
- Perturbative QCD calculations at asymptotically high densities
- ▶ New heavy-mass measurement ($M = 2.35 \pm 0.17 M_{\odot}$) of black-widow PSR J0952-0607
- ▶ Posterior credible bands for sound speed (*left*) and mass-radius relation (*center*) & Bayes factor against small sound speeds ($c_s^2 \le 0.1$) in neutron star cores (*right*):

 10^{2} verv strong Bayes factor 95% 2.51.00 median strong 0.75 $[{}^{\odot} 2.0]{W}_{M}$ moderate ~<u>~</u> 0.50anecdotal PSR J0740+6620 0.25 10° 1.0 PSR J0030+0451 no evidence 4U 1702-429 0.00^{L}_{0} 250 750 5001000 12 14 10 2.02.1 2.2 2.3 1.9 $\varepsilon \, [{\rm MeV \, fm^{-3}}]$ $R \, [\mathrm{km}]$ M/M_{\odot}

- ▶ With black-widow pulsar: strong evidence for $c_s^2 > 0.1$ in core of neutron stars with masses $M \le 2.1 M_{\odot}$
 - \rightarrow First-order phase transition unlikely

B8: Quarkonium production cross sections @ LHC, $\sqrt{s} = 7$ TeV

Using pNRQCD factorization, the inclusive J/ψ , $\psi(2S)$, $\Upsilon(2S)$, $\Upsilon(3S)$ production cross sections can be fitted with only 3 nonperturbative unknown instead of the 12 in NRQCD. The quality of the fit is very good. Also computed: photo-, EW ass. production, polarization.



• Brambilla Chung Vairo Wang PRD 105 (2022) 11, JHEP 03 (2023) 242

Investigation of the Lightest Hybrid Meson Candidate π_1

Kopf et al, Eur. Phys J. C(2021) 81, 1056



Table 1 Obtained masses, total widths and ratios of partial widths for the pole of the spin-exotic π_1 -wave and for the two poles in the a_2 -wave, the $a_2(1320)$ and the $a_2(1700)$. The first uncertainty is the statistical and the second the systematic one

Name	Pole mass (MeV/ c^2)	Pole width (MeV)	$\Gamma_{\pi\eta'}/\Gamma_{\pi\eta}$ (%)	$\Gamma_{KK}/\Gamma_{\pi\eta}$ (%)
<i>a</i> ₂ (1320)	$1318.7 \pm 1.9 ^{+1.3}_{-1.3}$	$107.5 \pm 4.6 {+3.3 \atop -1.8}$	$4.6 \pm 1.5 {}^{+7.0}_{-0.6}$	$31 \pm 22 {+9 \atop -11}$
$a_2(1700)$	$1686\pm22^{+19}_{-7}$	$412\pm75{}^{+64}_{-57}$	$3.5 \pm 4.4 {}^{+6.9}_{-1.2}$	$2.9 \pm 4.0 {}^{+1.1}_{-1.2}$
π_1	$1623 \pm 47 {+24 \atop -75}$	$455\pm88{}^{+144}_{-175}$	$554 \pm 110 {}^{+180}_{-27}$	_

Highlights B.11 "Coupled-channel dynamics":

PLs: D. Rönchen and B.-S. Zou

Extension of JüBo dynamical coupled-channel model to $K\Sigma$ photoproduction on the proton: EPI A 58, 229 (2022)

- Simultaneous analysis of $\pi N \rightarrow \pi N$, ηN , $K\Lambda$, $K\Sigma$ and $\gamma p \rightarrow \pi N$, ηN , $K\Lambda$, $K^+\Sigma^0 \& K^0\Sigma^+$
- almost 72,000 data points in total, $W_{max} = 2.4 \text{ GeV}$

Resonance analysis:

- all 4-star N and Δ states up to J = 9/2 are seen + some states rated less than 4 stars (exception: $N(1895)1/2^{-}$)
- no additional s-channel diagrams had to be included, but indications for new dyn. gen. poles
- \Rightarrow resonance parameters to be included in PDG averages

Further highlights:

- Prediction of dynamically generated states in coupled $\bar{D}^{(*)}\Lambda_c \bar{D}^{(*)}\Sigma_c^{(*)}$ system Zheng-Li Wang *et al.* EPJ C 82 (2022) Observation of several bound states, some close to LHCb pentaquarks
- Inclusion of the $\pi N \rightarrow \omega N$ channel in JüBo DCC model Yu-Fei Wang et al. PRD 106, 094031 (2022)
- In progress: JüBo DCC analysis of $\bar{K}N$ system (PhD S. Rawat): S = -1 hyperon resonance spectrum



Computing transverse momentum dependent PDFs

$$f^{TMD}(x, b_{\perp}, \mu, \zeta) = H(\zeta_{z}, \mu)e^{-\ln(\frac{\zeta_{z}}{\zeta})K(b_{\perp},\mu)}S_{r}^{\frac{1}{2}}(b_{\perp}, \mu)f^{qTMD}(x, b_{\perp}, \mu, \zeta_{z}), \qquad \zeta : \text{Collins-Soper scale}$$

$$K(b_{\perp}, \mu): \text{Collins-Soper kernel}$$

$$K(b_{\perp}, \mu): \text{Collins-Soper kernel}$$
Computed in lattice QCD, with staple-shaped link
Reduced Soft function,
computed in LQCD
$$\int_{1.6}^{1.6} \int_{1.4}^{1.2} \int_{1.6}^{1.6} \int_{1.6}^{1.6} \int_{1.4}^{1.2} \int_{1.6}^{1.6} \int_{1.6}^{1.6} \int_{1.4}^{1.2} \int_{1.6}^{1.6} \int_{$$

Missing ingredient: the quasi-TMD PDF, $f^{qTMD}(x, b_{\perp}, \mu, \zeta_z)$ Its computation is underway, arXiv: 2305.11824