

## **A.1 Flavor symmetries and final-state interactions in hadronic decays**

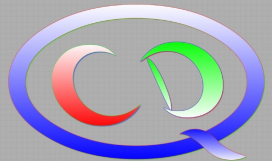
**PLs: Feng-Kun Guo, Bastian Kubis, Stephan Paul**

## **A.5 Quark mass dependence of hadronic observables**

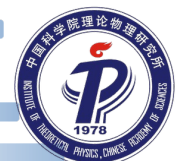
**PLs: Feng-Kun Guo, Ulf-G. Meißner**

**The 7th Symposium on Symmetries and the Emergence of Structure in QCD**

19-22 July, 2023, Rizhao



# Accomplishments of the projects



Chinese side + collaborating with German nodes

## ● A.1

### □ Three-body decays

- $D\pi$  FSI from  $B \rightarrow D\pi\pi$
- Probing structure of  $D_{s1}(2460)$  using  $D_{s1} \rightarrow D_s\pi^+\pi^-$

### □ Triangle singularities

- estimate of the  $X(3872)\gamma$  production rates via the TS mechanism
- analysis of updated BESIII data for the reactions involving  $Z_c(3900)$  with TS contributions

### □ Threshold cusps

- extraction of S-wave  $ND$  scattering lengths from  $\Lambda_b \rightarrow \pi^- p D^0$
- mechanism of  $J/\psi p$  photoproduction

□ ...

## ● A.5

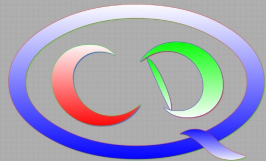
### □ Quark mass dependence

- Analyzing  $D\pi-D\eta-D_s\bar{K}$  finite volume energy levels from HadSpec at unphysical quark mass

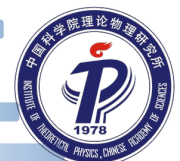
### □ $\theta$ -dependence $\Rightarrow$ QCD axion properties

- $\Delta$  contribution to the axioproduction of pion  $aN \rightarrow \pi N$

□ ...



# $D\pi$ FSI in $B \rightarrow D\pi\pi$



M.-L. Du, FKG, C. Hanhart, B. Kubis, U.-G. Meißner, PRL 126 (2021) 192001

- Lightest  $D_0^*$ :  $D_0^*(2300)$  in PDG since 2004

- Resonance parameters from fitting to  $D\pi$  spectrum using BW

PDG2023

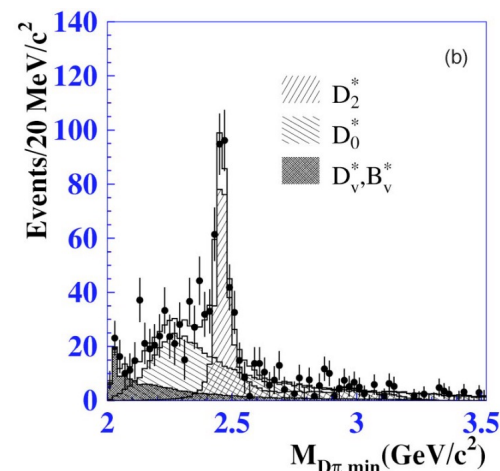
$D_0^*(2300)$  MASS  $2343 \pm 10$  MeV

$D_0^*(2300)$  WIDTH  $229 \pm 16$  MeV

- Two poles in unitarized CHPT

$(M, \Gamma/2)$	Lower (MeV)	Higher (MeV)
$D_0^*$	$(2105_{-8}^{+6}, 102_{-11}^{+10})$	$(2451_{-26}^{+36}, 134_{-8}^{+7})$
$D_1$	$(2247_{-6}^{+5}, 107_{-10}^{+11})$	$(2555_{-30}^{+47}, 203_{-9}^{+8})$

Albaladejo et al., PLB 767 (2017) 465;  
M.-L. Du et al., PRD 98 (2018) 094018



Belle, PRD 69 (2004) 112002

SU(3) irreps:  $\mathbf{3} \otimes \mathbf{8} = \mathbf{15} \oplus \mathbf{6} \oplus \mathbf{3}$

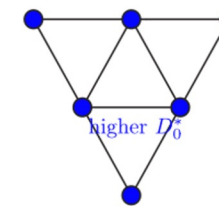
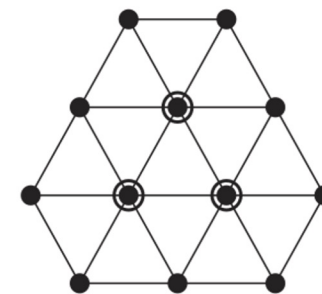
Albaladejo et al., PLB767(2017)465

$S = 2$

$S = 1$

$S = 0$

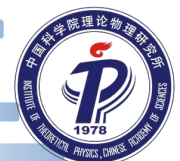
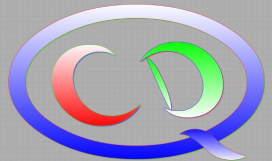
$S = -1$



WT term:  $\overline{\mathbf{15}}$ : repulsive;  $\mathbf{6}$ : attractive;  $\overline{\mathbf{3}}$ : most attractive

- Question

➤ What can be learned from precise LHCb data of  $B^- \rightarrow D^+ \pi^- \pi^-$ ?

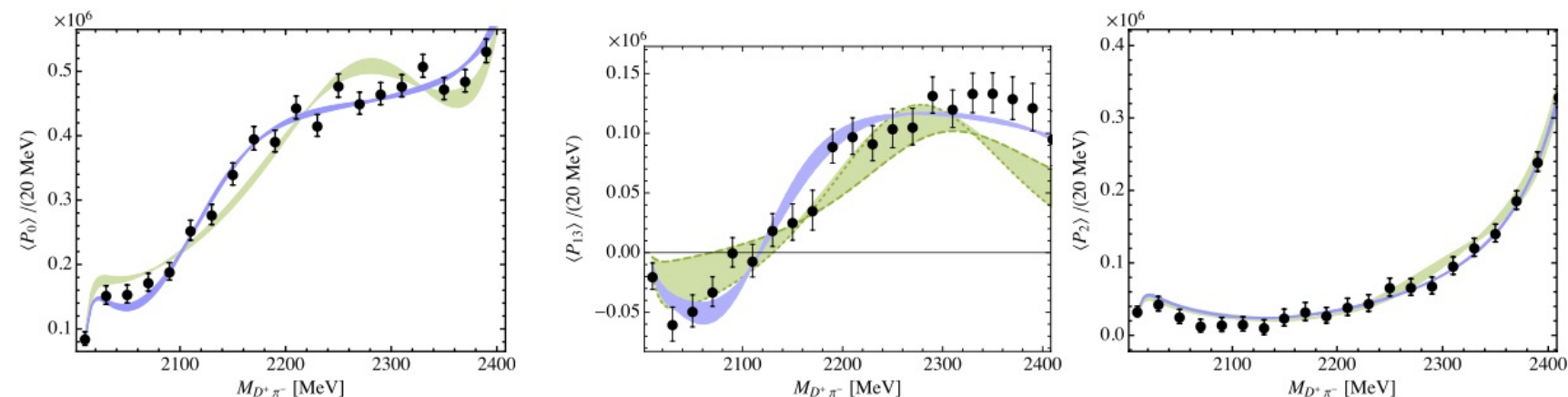


# $D\pi$ FSI in $B \rightarrow D\pi\pi$

M.-L. Du, FKG, C. Hanhart, B. Kubis, U.-G. Meißner, PRL 126 (2021) 192001

- Fits of the  $B^- \rightarrow D^+ \pi^- \pi^-$  data with KT equation considering 3-body unitarity: using  $S$ -wave  $D\pi$  scattering phase from UCHPT ( $\chi^2/\text{d.o.f.} = 1.2$ ) and from BW ( $\chi^2/\text{d.o.f.} = 2.0$ )

Data: LHCb, PRD94(2016)072001



$D_0^*(2300)$

$I(J^P) = 1/2(0^+)$

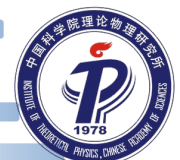
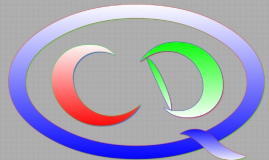


was  $D_0^*(2400)$

There is a strong evidence that recent data on  $B \rightarrow D\pi\pi$  (AAJ 2015Y, AAJ 2016AH) and  $B \rightarrow D\pi K$  (AAJ 2014BH, AAJ 2015V, AAJ 2015X) call for two poles in the scalar  $I = 1/2 \pi D$  amplitude in this mass range. The data are consistent with a lower pole at  $(2105_{-8}^{+6}) - i(102_{-11}^{+10})$  MeV and a higher pole at  $(2451_{-26}^{+35}) - i(134_{-9}^{+7})$  MeV (DU 2018A, DU 2019, DU 2021). For details see review on "Heavy Non- $\bar{q}q$  Mesons."

- The LHCb data are well described with UCHPT amplitude with two  $D_0^*$  states; the lowest has a mass about 2.1 GeV
- Support from recent lattice results: L. Gayer et al. [HadSpec], JHEP07(2021)123  
 $D_0^*$  with  $M \approx 2.2$  GeV and  $\Gamma \approx 0.4$  GeV obtained using  $M_\pi \approx 239$  MeV

Independent analyses of data from other groups are called for to change the PDG entry!!



# $D_{s1} \rightarrow D_s \pi^+ \pi^-$

M.-N. Tang, Y.-H. Lin, FKG, C. Hanhart, U.-G. Meißner, CTP 75 (2023) 055203

- Experimental measurement:

$$\frac{\Gamma(D_{s1}(2460)^+ \rightarrow D_s^+ \pi^+ \pi^-)}{\Gamma(D_{s1}(2460)^+ \rightarrow D_s^{*+} \pi^0)} = \begin{cases} 0.14 \pm 0.04 \pm 0.02 & \text{Belle, PRL 92(2004)012002} \\ 0.09 \pm 0.02 & \text{PDG fit} \end{cases}$$

- Isospin breaking  $D_{s1} \rightarrow D_s^* \pi^0$ :  $(111 \pm 15)$  keV H.-L. Fu et al., EPJA58(2022)70

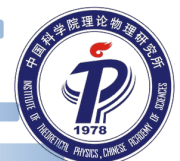
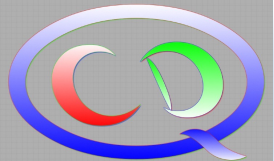
- How about  $D_{s1}(2460)^+ \rightarrow D_s^+ \pi^+ \pi^-$ ?

- CHPT result from S. Fajfer, A. Prapotnik Brdnik, PRD92(2015)074047

$$\Gamma(D_{s1}(2460)^+ \rightarrow D_s^+ \pi^+ \pi^-) = 0.25(4)(7) \begin{pmatrix} +2 \\ -4 \end{pmatrix} \text{ keV}$$

Both  $D_s$  and  $D_{s1}$  were treated statically,  $P$ -wave happens between  $\pi^+$  and  $\pi^- \Rightarrow I(\pi^+ \pi^-) = 1$ , isospin breaking

- But **isospin is conserved (!!)** for  $P$ -wave between  $D_s$  and isoscalar  $\pi^+ \pi^-$



# $D_{s1} \rightarrow D_s \pi^+ \pi^-$

M.-N. Tang, Y.-H. Lin, FKG, C. Hanhart, U.-G. Meißner, CTP 75 (2023) 055203

- Hadronic molecule v.s. other (compact) components

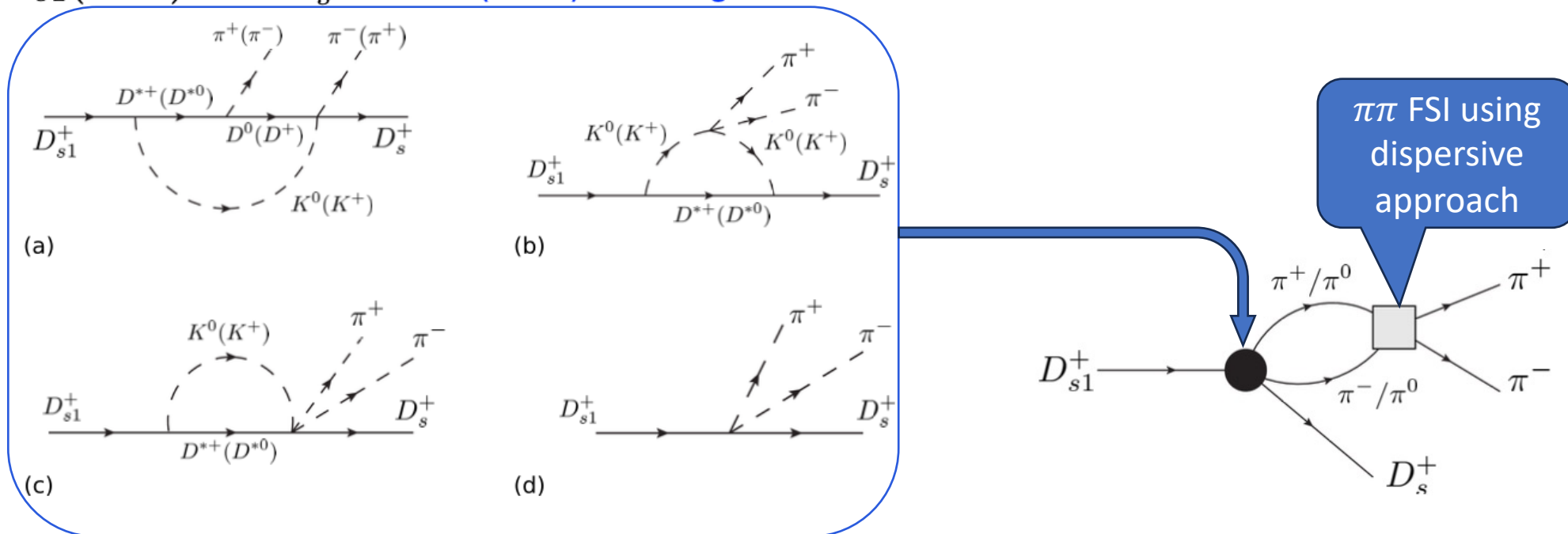
Effective coupling contains crucial information

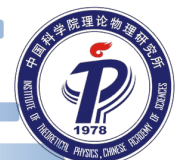
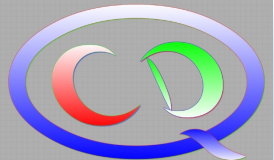
S. Weinberg, PR137(1965)B672

$$g^2 \propto (1 - Z) \sqrt{2\mu E_B}$$

is maximized for a pure molecular state

- Diagrams for  $D_{s1}(2460)^+ \rightarrow D_s^+ \pi^+ \pi^-$ : (a,b,c): leading for molecular state





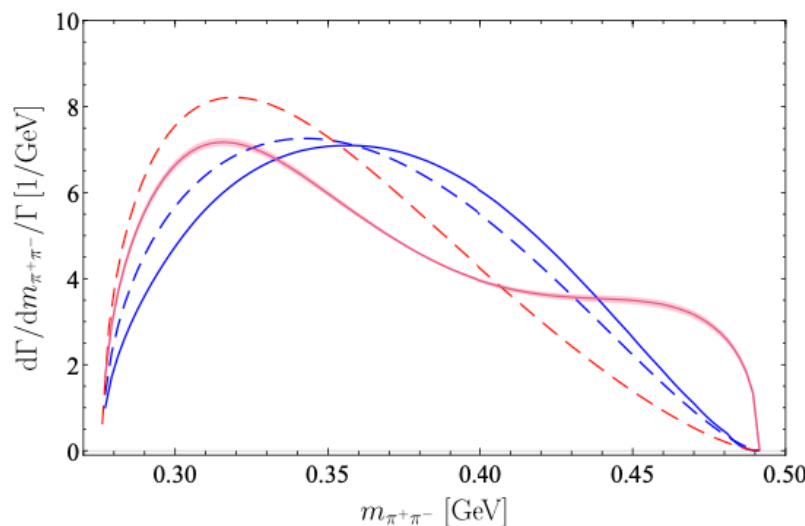
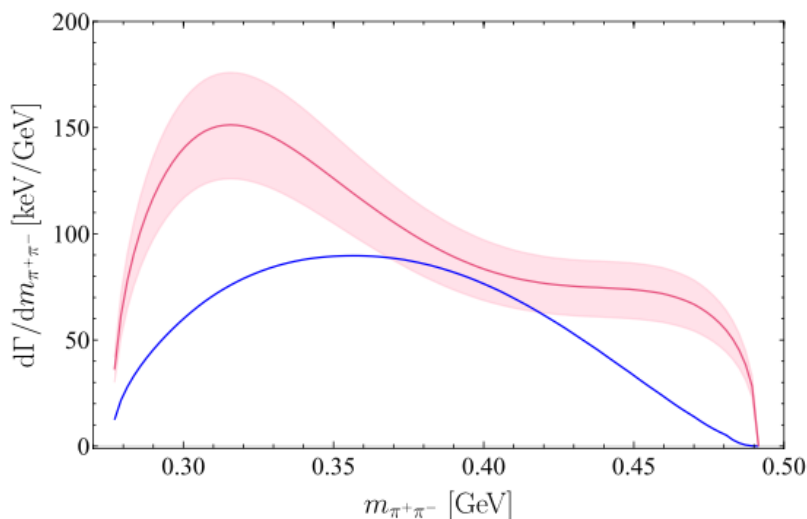
# $D_{s1} \rightarrow D_s \pi^+ \pi^-$

M.-N. Tang, Y.-H. Lin, FKG, C. Hanhart, U.-G. Meißner, CTP 75 (2023) 055203

- **Double bump structure in the  $\pi\pi$  invariant mass distribution** as a feature of the hadronic molecular picture

★ **Red: molecular**, assuming the  $D_{s1} D_s \pi\pi$  contact term to vanish

★ **Blue: compact**, without  $D^* K$  loops

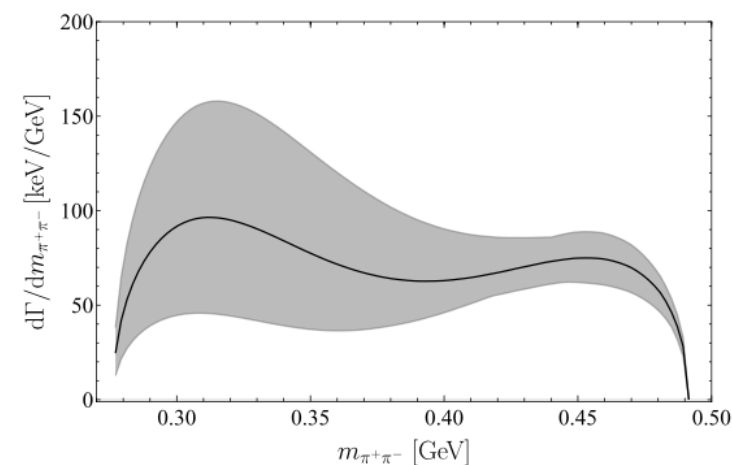


Dashed: without  $\pi\pi$  FSI

- agree with Belle measurement  $0.14 \pm 0.04 \pm 0.02$ :

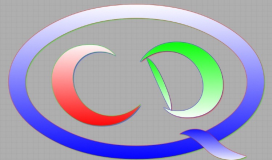
$$\left. \frac{\Gamma(D_{s1}(2460)^+ \rightarrow D_s^+ \pi^+ \pi^-)}{\Gamma(D_{s1}(2460)^+ \rightarrow D_s^{*+} \pi^0)} \right|_{\text{mol.}} = 0.19^{+0.07}_{-0.05}$$

- With contact term (consistent with zero) fixed from the Belle measurement

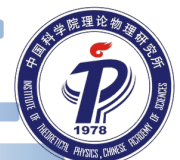


- Prediction for the bottom analogue:

$$\Gamma(B_{s1}^0 \rightarrow B_s^0 \pi^+ \pi^-) = (3 \pm 1) \text{ keV}$$



# Threshold cusps: $ND$ scattering length



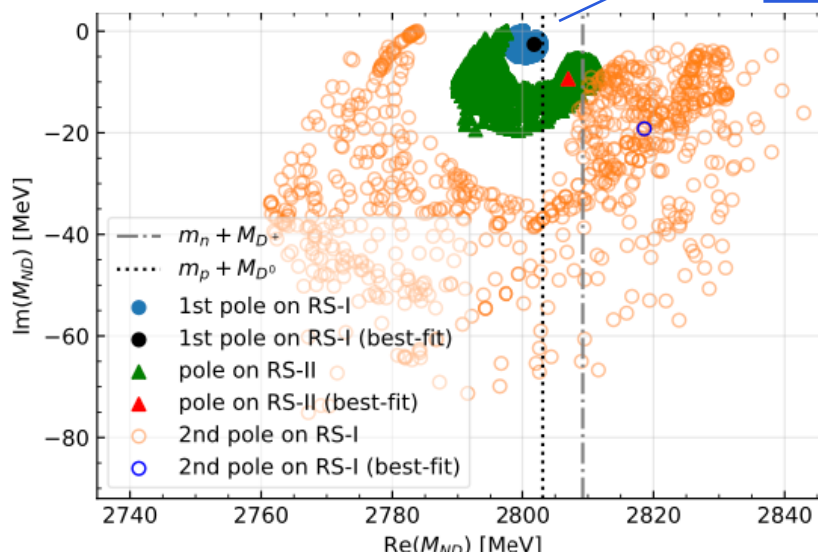
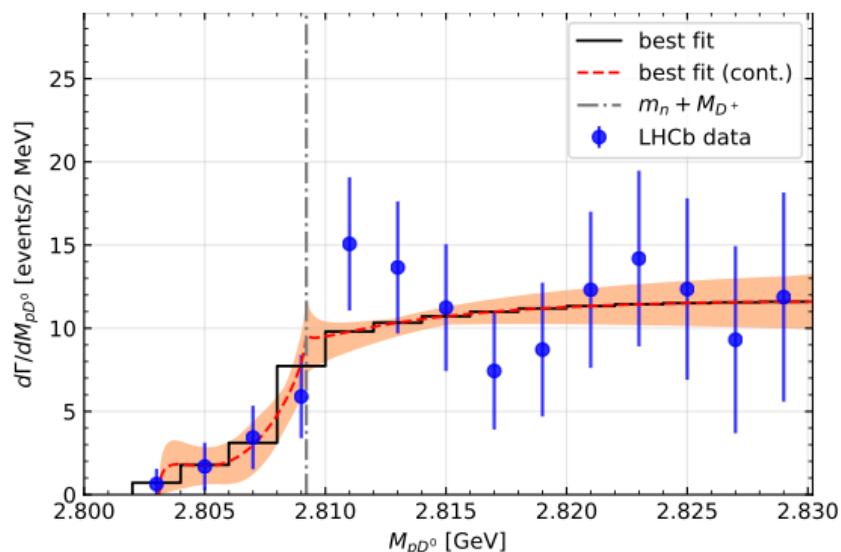
S. Sakai, FKG, B. Kubis, PLB 808 (2020) 135623

- $ND$  near-threshold interaction is relevant for understanding charmed baryons around 2.8 GeV, e.g.  $\Sigma_c(2800)$  ( $J^P = ??$ )
- The LHCb data on  $\Lambda_b \rightarrow \pi^- p D^0$  show peculiar behavior around  $nD^+$  threshold

LHCb, JHEP05(2017)030

- Analysis with nonrelativistic EFT with two channels:  $pD^0$  and  $nD^+$ .

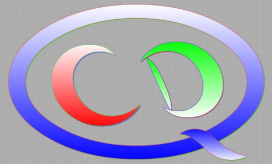
Could be  $\Sigma_c(2800)$   
 $\Rightarrow J^P = 1/2^-$



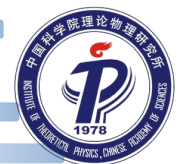
- Scattering lengths:

$a_{ND}$ [fm]	Our result	SU(4) [10]	SU(4) [11]	SU(8) [12]	Meson-exchange model [15]
$I = 0$	$-0.79^{+0.66}_{-0.61}$	-0.43	$-0.57 + i 0.001$	$0.004 + i 0.002$	$-0.41 + i 0.04$
$I = 1$	$-3.8^{+1.4}_{-2.0} + i 2.7^{+1.6}_{-2.7}$	-0.41	$-1.47 + i 0.65$	$0.33 + i 0.05$	$-2.07 + i 0.57$






# Mechanism of $J/\psi p$ photoproduction



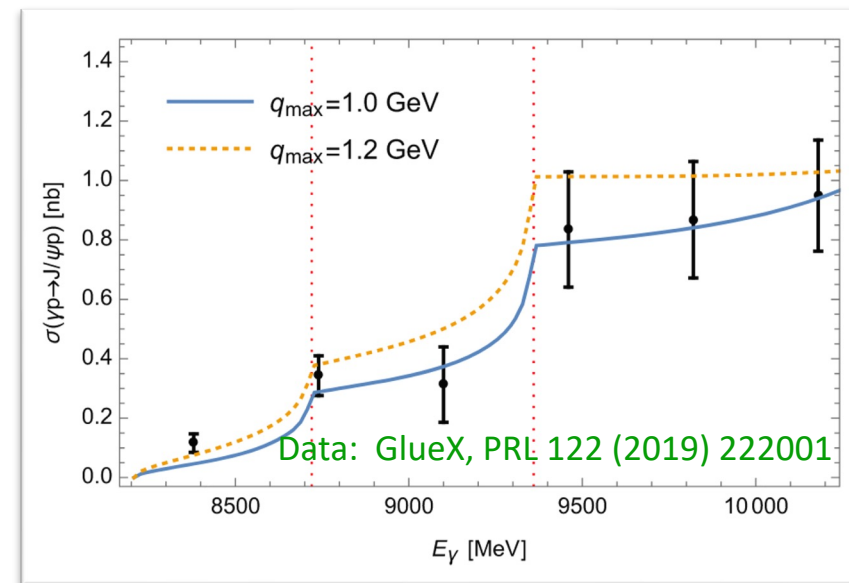
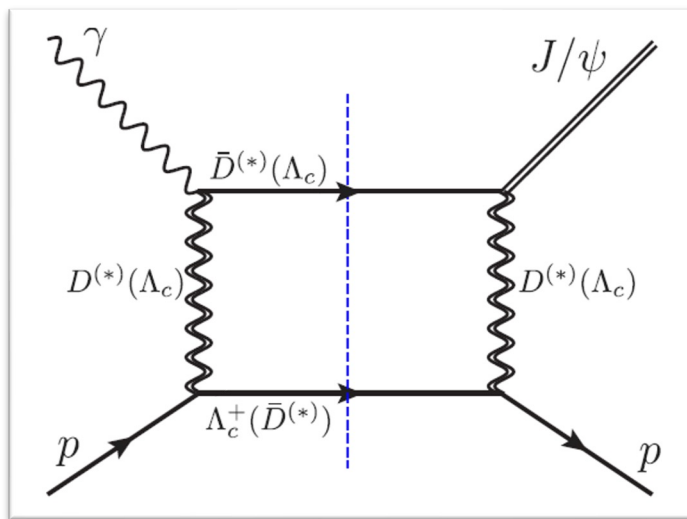
M.-L. Du, V. Baru, FKG, C. Hanhart, U.-G. Meißner, A. Nefediev, I. Strakovsky, EPJC 80 (2020) 1053

- Open-charm channels easier to be produced than  $J/\psi p$ ; thresholds nearby



Unitarity:  $J/\psi p \rightarrow J/\psi p$  enters w/o VMD, but cannot be singled out

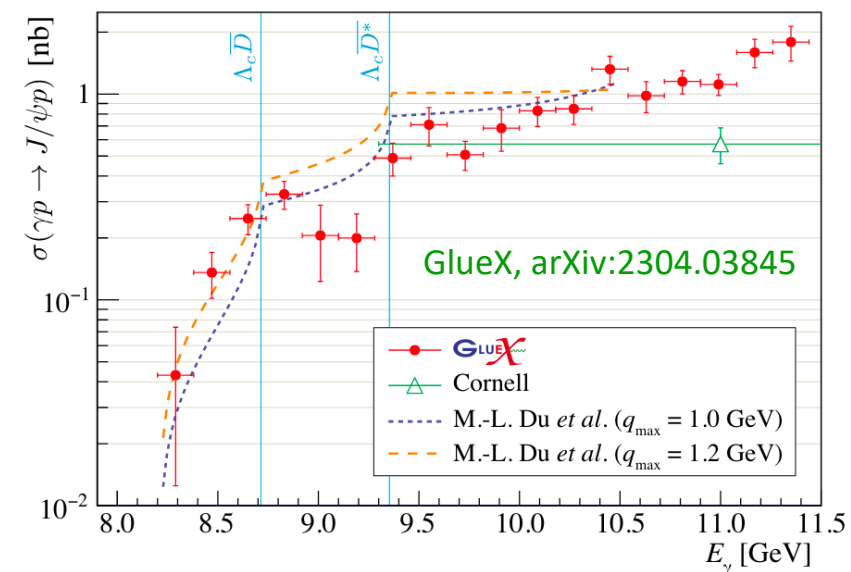
$\Lambda_c^+ + \bar{D}^-$	$2286 + 1865 = 4151 \text{ MeV}$
$J/\psi + p$	$3097 + 938 = 4035 \text{ MeV}$

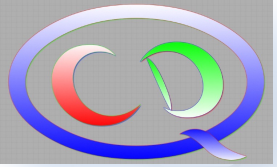


- Unique feature: cusps at  $\Lambda_c \bar{D}^{(*)}$  thresholds

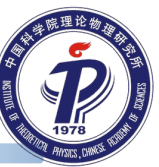
- The same mechanism for  $J/\psi p \rightarrow J/\psi p$  leads to small scattering length; need to compare with the scattering length from gluon exchanges (ongoing):

$$|a^{J=1/2}| = 0.2 \dots 3.1 \text{ mfm}, \quad |a^{J=3/2}| = 0.2 \dots 3.0 \text{ mfm},$$





# $D\pi$ energy levels in finite volume



A. Asokan, M.-N. Tang, FKG, C. Hanhart, Y. Kamiya, U.-G. Meißner, arXiv:2212.07856

- $D_0^*$  from lattice

- one bound state pole just below  $D\pi$  threshold with  $M_\pi = 390$  MeV HadSpec, JHEP 10 (2016) 011

- evolving to a resonance with  $M \approx 2.2$  GeV,  $\Gamma \approx 0.4$  GeV when  $M_\pi = 239$  MeV HadSpec, JHEP 08 (2021) 123

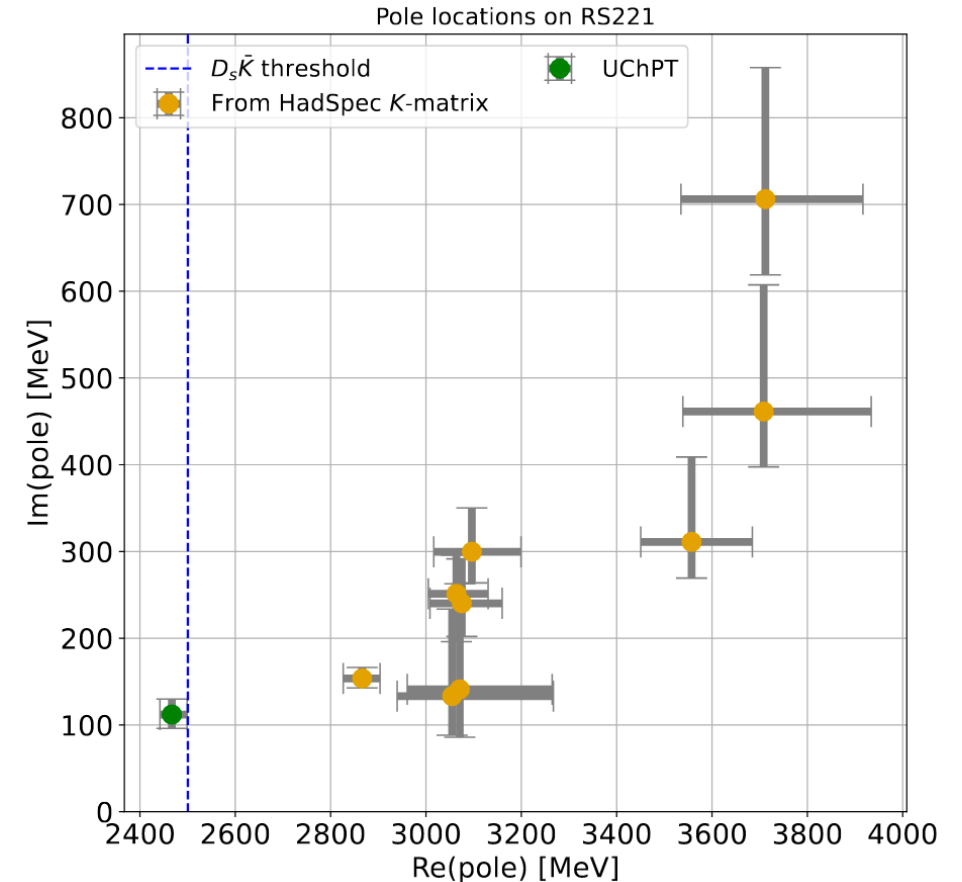
- K-matrix parametrization of the T-matrix

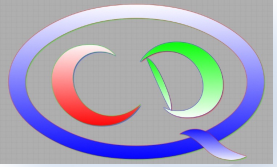
- the one used by HadSpec

$$T_{ij}^{-1}(s) = K_{ij}^{-1}(s) + I_{ij}(s)$$

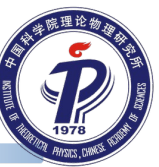
$$K_{ij}(s) = \left( g_i^{(0)} + g_i^{(1)} s \right) \left( g_j^{(0)} + g_j^{(1)} s \right) \frac{1}{m^2 - s} + \gamma_{ij}^{(0)} + \gamma_{ij}^{(1)} s$$

the second pole present, but badly constrained





# $D\pi$ energy levels in finite volume



A. Asokan, M.-N. Tang, FKG, C. Hanhart, Y. Kamiya, U.-G. Meißner, arXiv:2212.07856

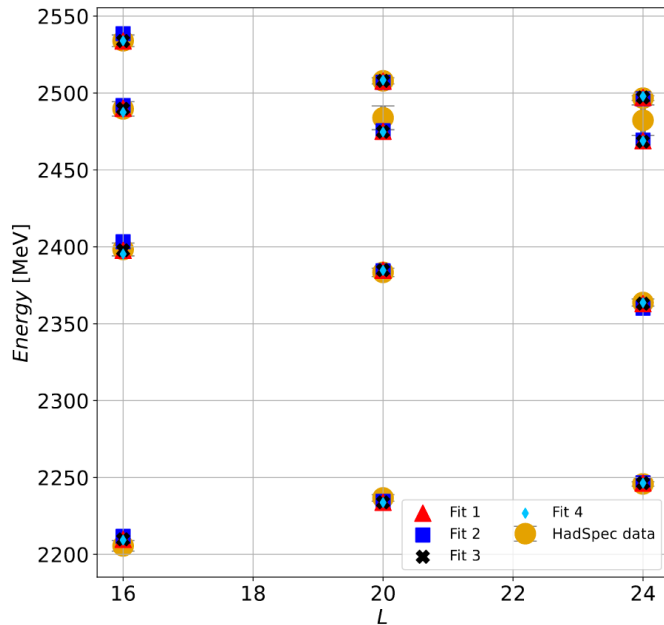
- SU(3) constrained K-matrix

- contact terms in all three irreps

- bare poles in  $\bar{3}$  and  $6$

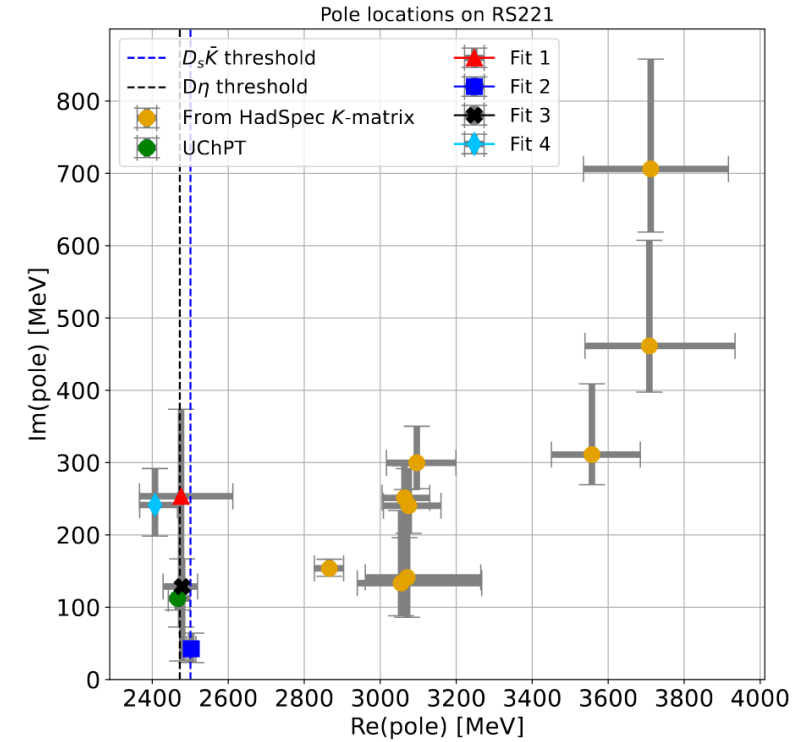
$$K = \left( \frac{g_3^2}{m_3^2 - s} + c_3 \right) C_{\bar{3}} + \left( \frac{g_6^2}{m_6^2 - s} + c_6 \right) C_6 + c_{15} C_{15}.$$

- Fit to energy levels by HadSpec

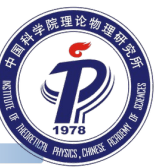
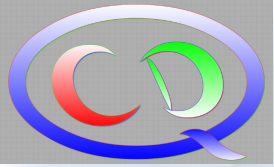


- Existence of the second  $D_0^*$  pole

- mass consistent with the one in UCHPT



	$g_3$ [GeV]	$m_3$ [MeV]	$g_6$ [GeV]	$m_6$ [MeV]	$c_3$	$c_6$	$c_{15}$	$\chi^2$	$\chi^2/\text{dof}$
Fit 1	$2.92 \pm 0.39$	$2275.1 \pm 0.6$	$0.32 \pm 0.32$	$2542 \pm 50$	$4 \pm 3$	$0.7 \pm 0.4$	$-0.6 \pm 0.2$	7.1	1.4
Fit 2	$2.31 \pm 0.14$	$2274.5 \pm 0.8$	$0.66 \pm 0.17$	$2560 \pm 37$	-	-	$-0.6 \pm 0.2$	14	1.9
Fit 3	$2.91 \pm 0.39$	$2275.1 \pm 0.6$	$1.20 \pm 0.62$	$2735 \pm 266$	$4 \pm 2$	-	$-0.6 \pm 0.2$	7.3	1.2
Fit 4	$3.16 \pm 0.38$	$2275.3 \pm 0.6$	-	-	$5 \pm 2$	$1.0 \pm 0.2$	$-0.4 \pm 0.2$	8.2	1.2



# Axioproduction of pion

T. Vonk, F.-K. Guo and U.-G. Meißner, PRD 105 (2022) 054029

- The pion axioproduction  $aN \rightarrow \pi N$  process was proposed to be employed to detect axions

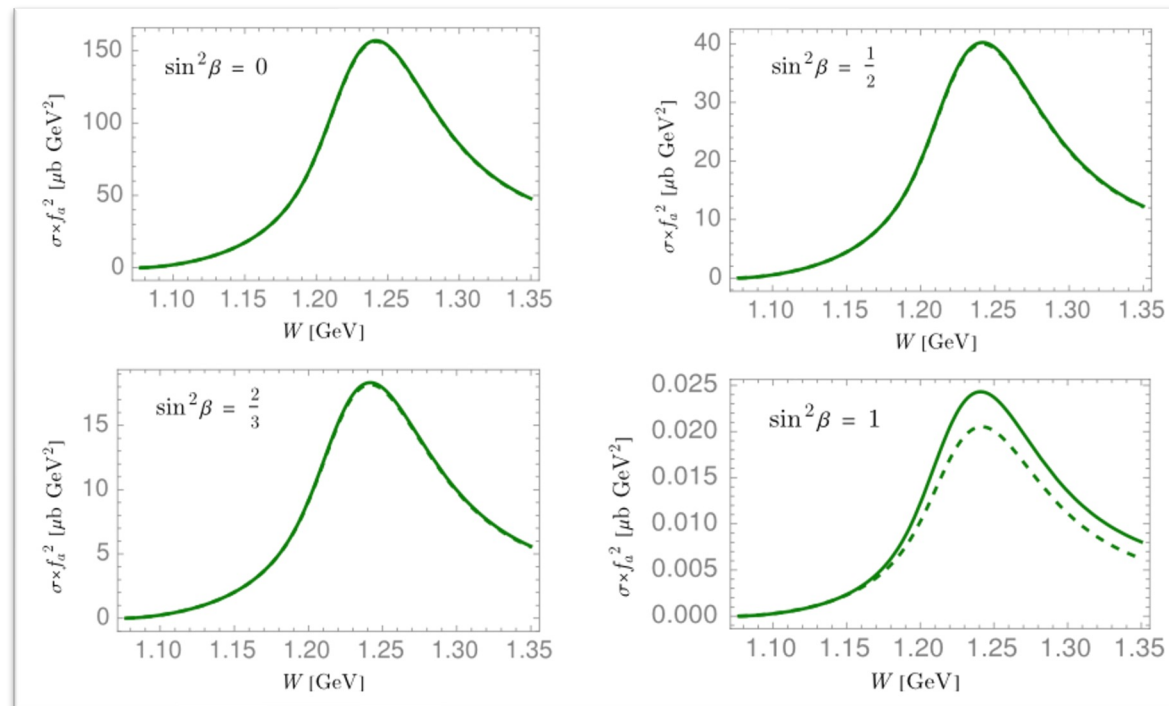
□ with naive estimate around the  $\Delta$  resonance

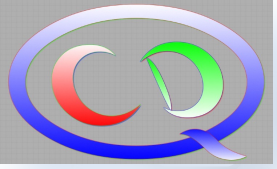
P. Carena, B. Fore, M. Giannotti, A. Mirizzi, S. Reddy, PRL 126 (2021) 071102

$$\sigma_{aN \rightarrow \pi N} \approx \frac{F^2}{f_a^2} \sigma_{\pi N \rightarrow \pi N} = \mathcal{O}\left(\frac{1 \text{ mb GeV}^2}{f_a^2}\right)$$

- However, **isospin breaking!** The  $\Delta$  contribution to the  $aN \rightarrow \pi N$  using UCHPT

□ feature: isospin breaking in amplitude  $\propto \frac{m_d - m_u}{m_d + m_u}$  v.s. normal isospin breaking  $\propto \frac{m_d - m_u}{\Lambda_{\text{QCD}}}$  or  $\frac{m_d - m_u}{m_s}$

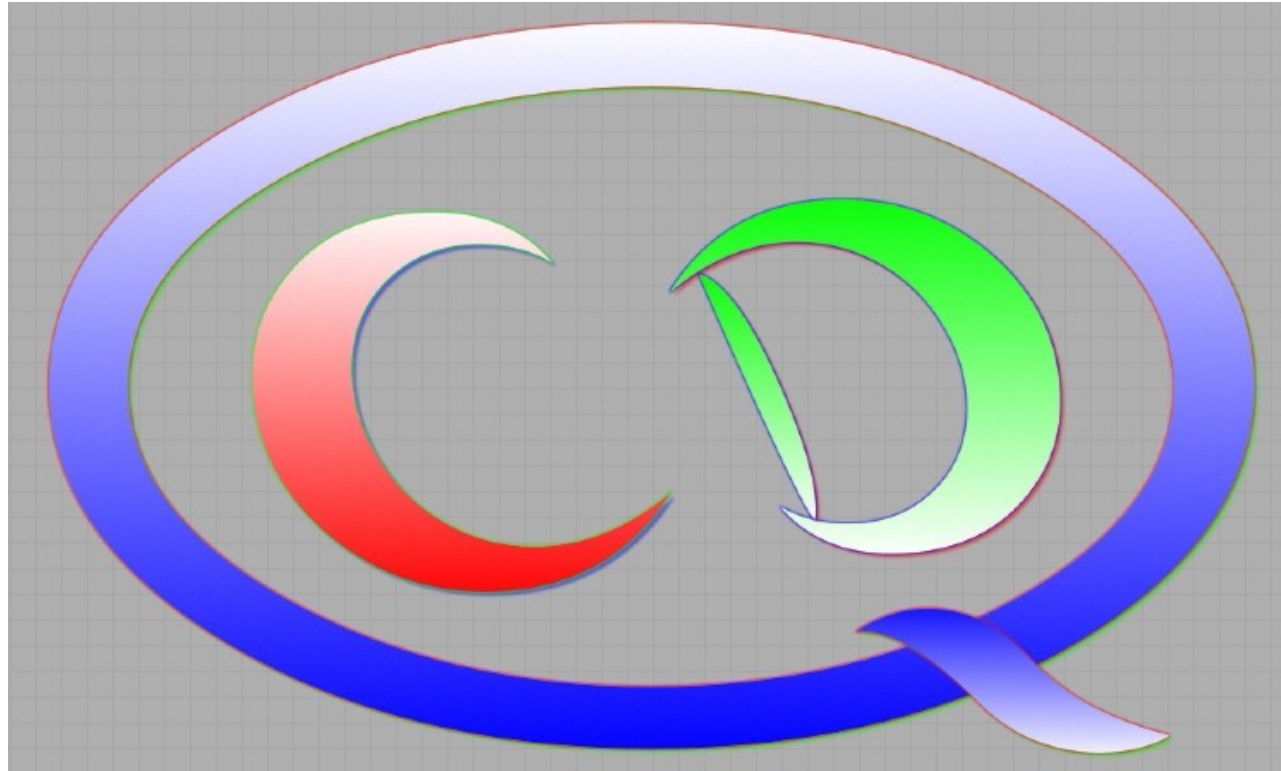




# Summary



- Good progress has been made
- Intensive collaborations between the Chinese and German nodes



**Thank you for your attention !**