

RHIC-BES Seminar
April 25, 2023

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Rice University

Outline

- QCD Phase Diagram
- What role can dileptons play?
- Some theoretical considerations
- Some experimental considerations
- The Dilepton Multitool
 - spectrometer, chronometer, thermometer, barometer, polarimeter, multimeter
- Future prospects



QCD phase diagram

Experimentally, one can access different regions of phase diagram by varying centre-of-mass energy

- experimental data over 3-4 orders of magnitude in $\sqrt{s_{NN}}$

LHC, RHIC, and FAIR provide access to low μ_B region

- cross-over region

Several experiments/facilities give access to μ_B regions that both cover cross-over, possible 1st order PT and a conjectured CP

- AGS, SPS (NA61/SHINE)
- SIS18/FAIR (Hades)
- RHIC beam energy scan (BES)

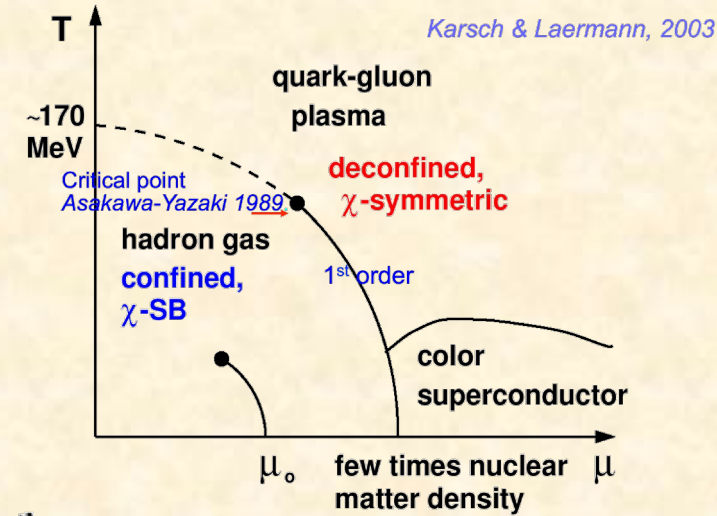
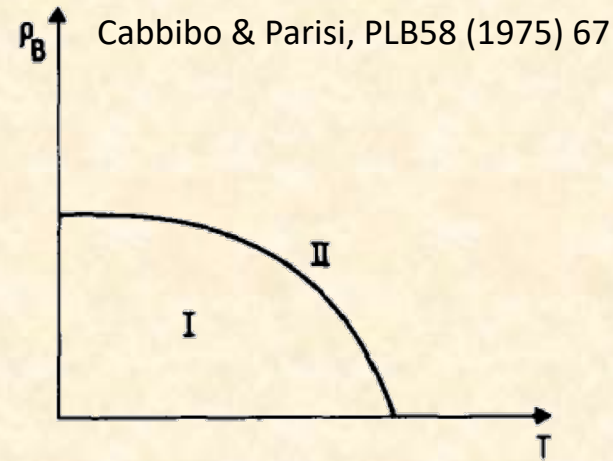
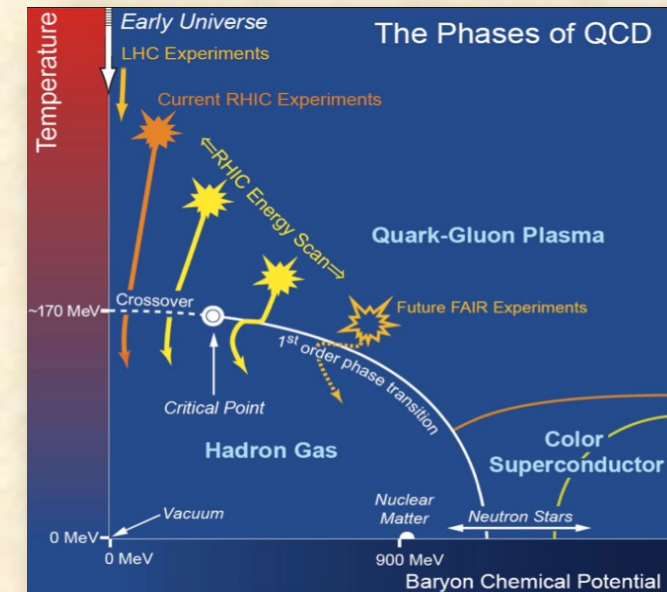
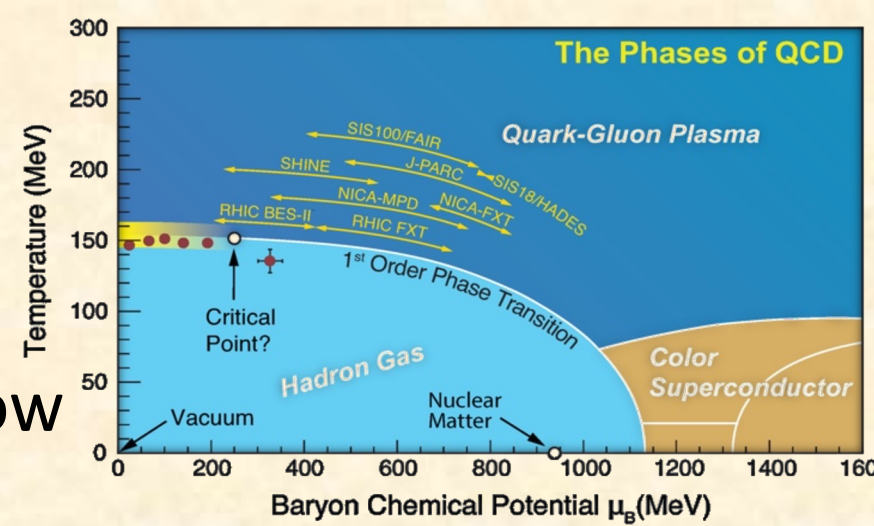


Fig. 1. Schematic phase diagram of hadronic matter. ρ_B is the density of baryonic number. Quarks are confined in phase I and unconfined in phase II.



BES-2 White Paper – STAR Note 598

Charting the QCD phase diagram



- Turn-off of QGP signatures - suppression, elliptic flow
- First-order phase transition - changes in EoS due to attractive force (softest point)
 - “step” in mean transverse mass of identified particles
 - non-monotonic behavior of directed flow slope at mid-rapidity ($dv_1/dy|_{y=0}$)
- Critical point - divergence of the correlation length \Rightarrow non-monotonic behavior of higher moments of conserved quantities
 - experimentally, skewness S , and kurtosis κ of event-by-event net-particle distributions

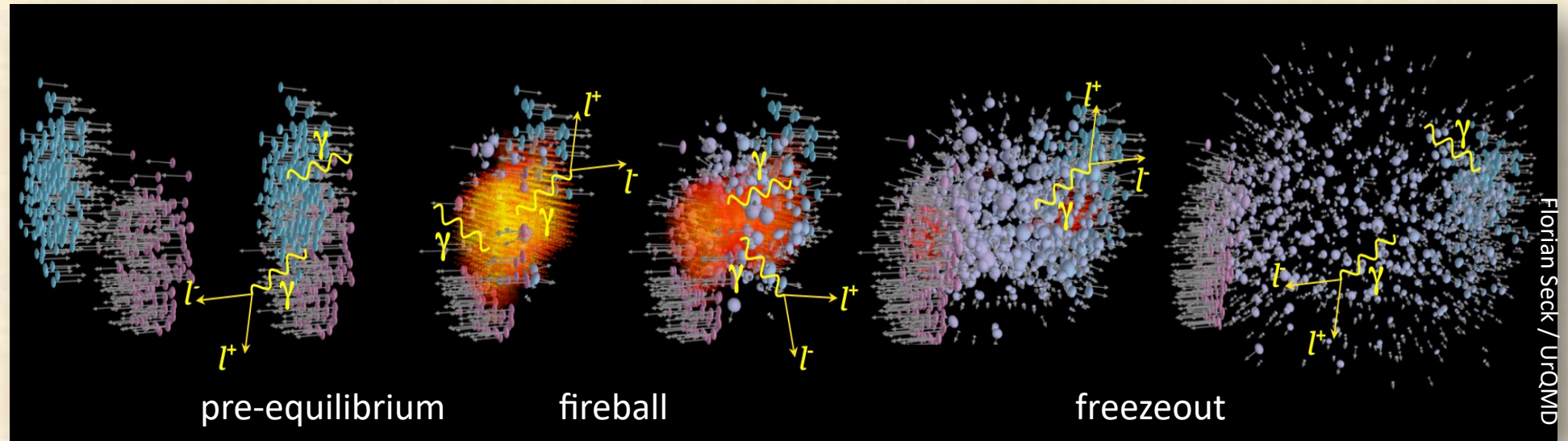
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What can dileptons do?

© ParticleZoo.net

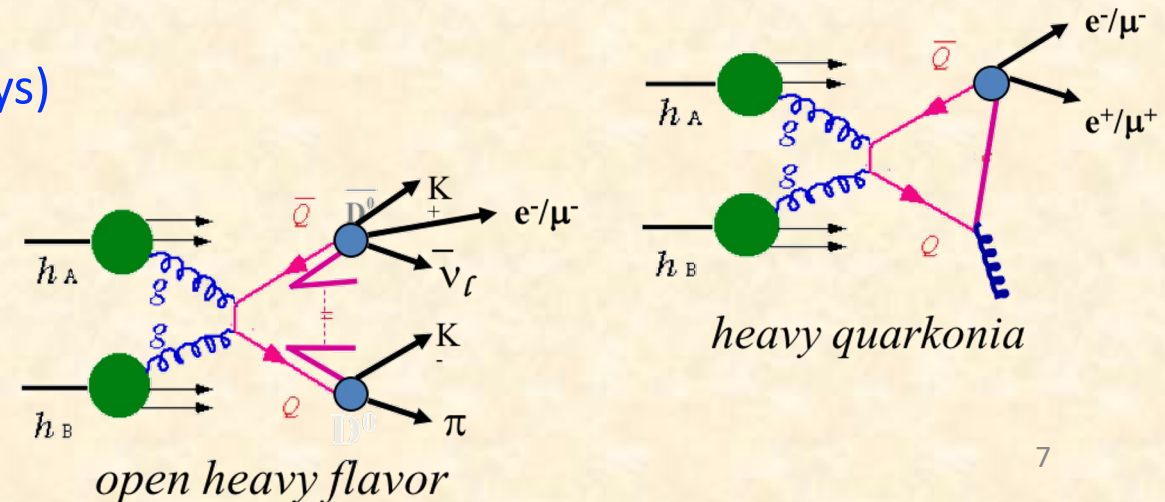
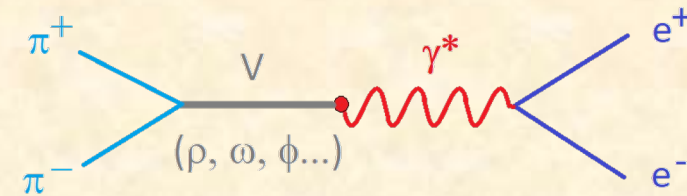
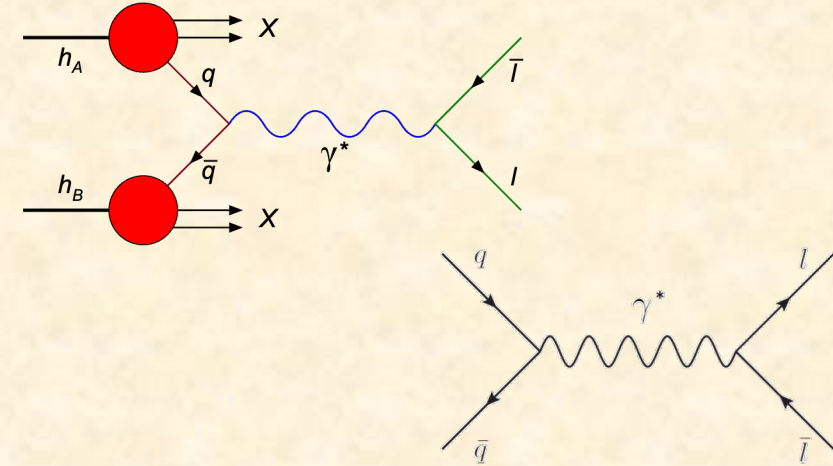


Dileptons are excellent penetrating probes

- colorless objects \therefore no coupling to strongly interacting matter
- produced in various ways throughout the system's evolution
- long mean free paths
- experimentally provide an additional “knob”: invariant mass

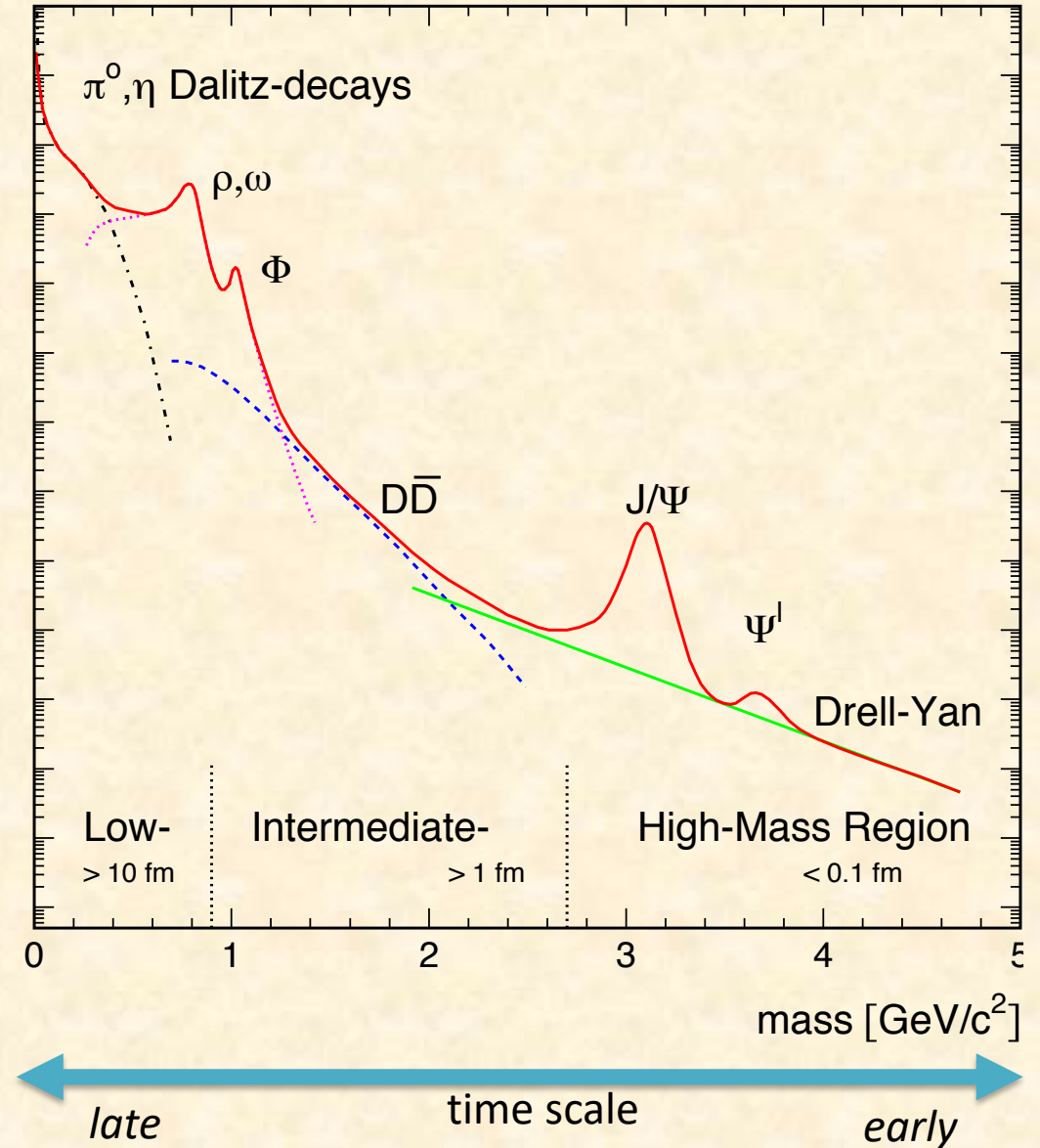
Dilepton invariant mass spectrum

- Primordial emissions, pre-equilibrium
 - Drell-Yan, $NN \rightarrow e^+e^- X$
 - heavy flavor production ($c\bar{c}, b\bar{b}$), quarkonia & open charm
- Thermal radiation from QGP/hadronic matter
 - QGP thermal radiation $q\bar{q} \rightarrow e^+e^-$
 - HG thermal radiation $\pi^+\pi^- \rightarrow e^+e^-$
 - in-medium ρ
 - other 4π , multi-meson interactions, incl. $\rho - a_1$ mixing
- Long-lived hadron and resonance decays
 - decays of light mesons of $\pi^0, \eta, \omega, \phi$ (incl. Dalitz decays)
 - in-medium modification of vector mesons
 - decays of quarkonia $J/\Psi, \Psi'$ and correlated $D\bar{D}$ pairs
 - nuclear modification effects



Dilepton invariant mass spectrum

- High Mass Range (HMR)
 - $M_{ee} > 3 \text{ GeV}/c^2$
 - primordial emission, Drell-Yan
 - Heavy quarkonia: J/ψ and Υ suppression
- Intermediate Mass Range (IMR)
 - $1.1 < M_{ee} < 3 \text{ GeV}/c^2$
 - QGP thermal radiation
 - Semi-leptonic decay of correlated charm heavy-flavor modification
- Low Mass Range (LMR)
 - $M_{ee} < 1.1 \text{ GeV}/c^2$
 - in-medium modification of vector mesons
 - fireball lifetime measurement
 - transport coefficients (electrical conductivity)



EM production rates

From thermal field theory†, using EM current-current correlation function:

$$\Pi_{em}^{\mu\nu}(q_0, q) = -i \int d^4x e^{iqx} \Theta(x^0) \langle\langle [j^\mu(x), j^\nu(0)] \rangle\rangle$$

$$j^\mu = \sum_q e_q \bar{q} \gamma^\mu q = \frac{2}{3} \bar{u} \gamma^\mu u - \frac{1}{3} \bar{d} \gamma^\mu d - \frac{1}{3} \bar{s} \gamma^\mu s$$

with the thermal emission rates

- photons:

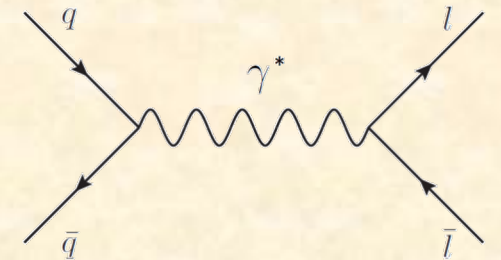
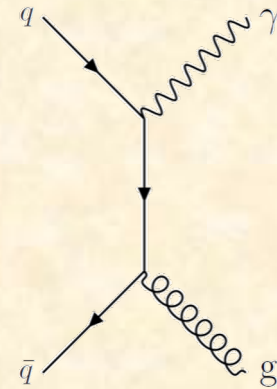
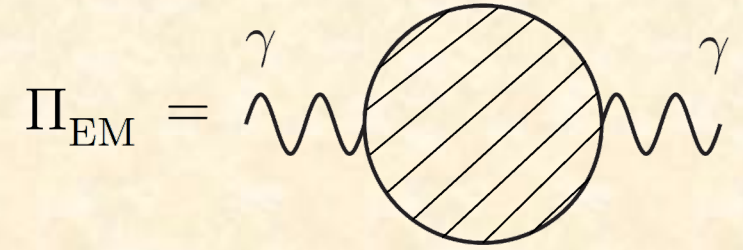
$$p_0 \frac{dN_\gamma}{d^4x d^3p} = -\frac{\alpha_{em}}{\pi^2} f^B(p_0; T) \frac{1}{2} g_{\mu\nu} \text{Im} \Pi_{em}^{\mu\nu}(M=0, p; \mu_B, T)$$

- dileptons:

$$\frac{dN_{ll}}{d^4x d^4p} = -\frac{\alpha_{em}^2}{\pi^3 M^2} L(M) f^B(p_0; T) \frac{1}{3} g_{\mu\nu} \text{Im} \Pi_{em}^{\mu\nu}(M, p; \mu_B, T)$$

$L(M)$ is lepton space factor and $f^B(p; T)$ is the thermal Bose distribution

- both governed by same underlying spectral functions
 - but different kinematic regimes (lightlike and timelike)



† see e.g. Friman, et al., Lecture Notes in Phys. 814 (2011) 1

Connection with vector mesons

For lightest quarks

$$j^\mu = \sum_q e_q \bar{q} \gamma^\mu q = \frac{2}{3} \bar{u} \gamma^\mu u - \frac{1}{3} \bar{d} \gamma^\mu d - \frac{1}{3} \bar{s} \gamma^\mu s$$

or grouping into isospin states $I = 1(\rho), 0(\omega, \phi)$:

$$\begin{aligned} j^\mu &= \frac{1}{2} (\bar{u} \gamma^\mu u - \bar{d} \gamma^\mu d) + \frac{1}{6} \bar{d} \gamma^\mu d (\bar{u} \gamma^\mu u + \bar{d} \gamma^\mu d) - \frac{1}{3} \bar{s} \gamma^\mu s \\ &= \frac{1}{\sqrt{2}} j_\rho^\mu + \frac{1}{3\sqrt{2}} j_\omega^\mu + \frac{1}{3} j_\phi^\mu \end{aligned}$$

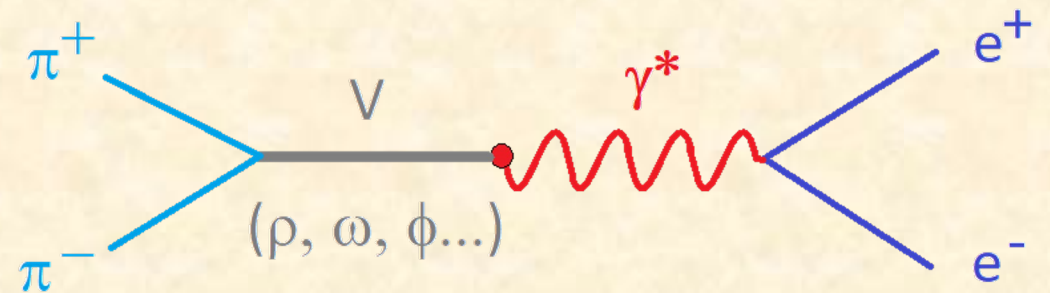
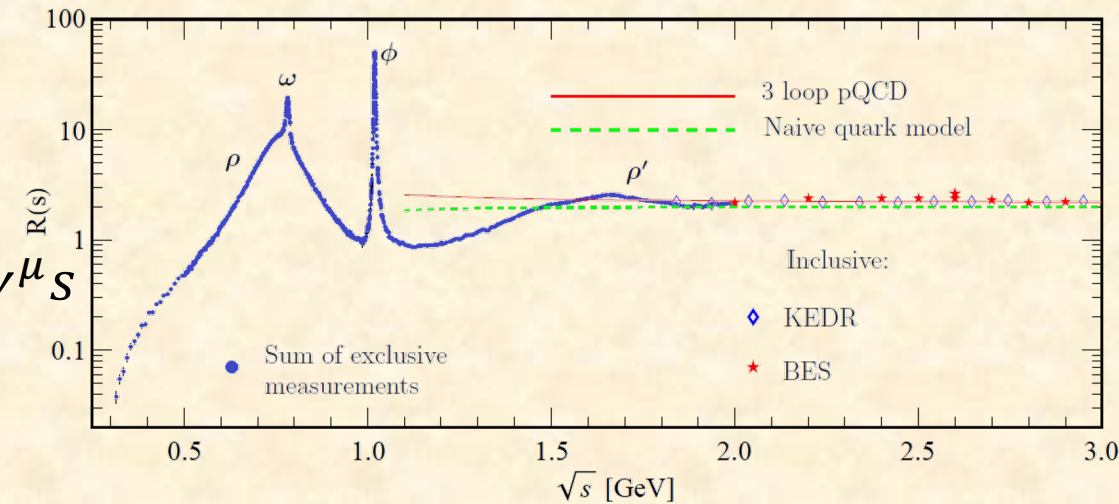
which leads at low M :

$$\text{Im } \Pi_{em} \sim D_\rho + \frac{1}{9} D_\omega + \frac{2}{9} D_\phi$$

vector meson dominance

- carry same quantum numbers as photons
- can directly decay into dileptons
- $\rho(770)$ dominant source

$\text{Im } \Pi_{em}$ is well understood in vacuum:



In-medium vector mesons (1)

ρ meson will interact with hadrons in the medium

propagator will have various contributions to the self-energy

$$D_\rho(M, q; T, \mu_B) = \frac{1}{(M^2 - m_\rho^2 - \Sigma_{\rho\pi\pi} - \Sigma_{\rho M} - \Sigma_{\rho B})}$$

$$\Sigma_{\rho\pi\pi} = \text{diagram showing } \rho \text{ meson interacting with a pion cloud via } \pi \text{ mesons}$$

in-medium pion cloud

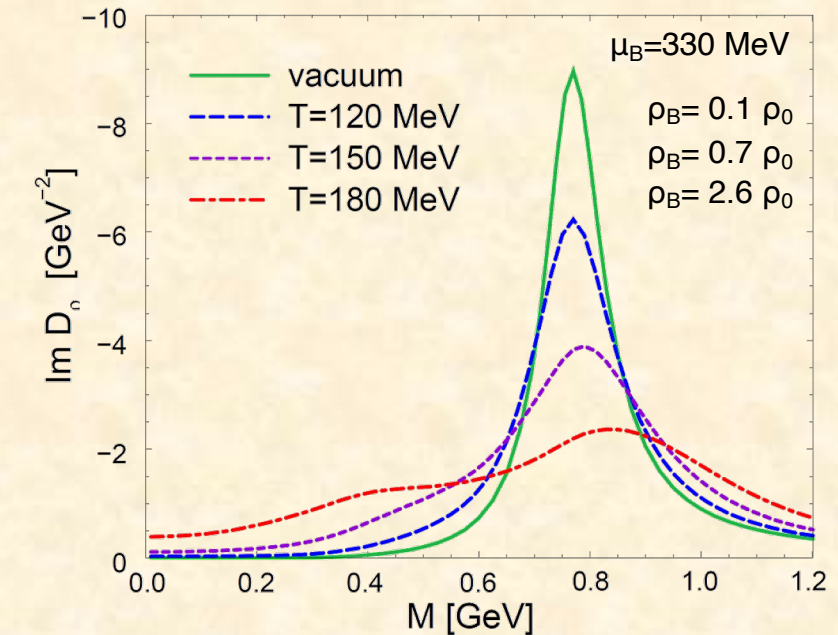
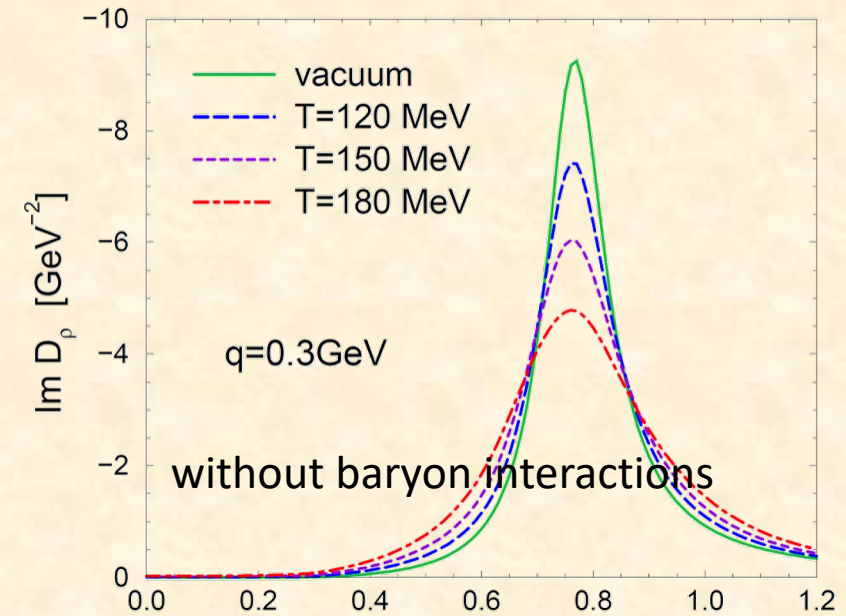
$$\Sigma_{\rho B, M} = \text{diagram showing } \rho \text{ meson interacting with a hadron } h \text{ via } N^*(1520) \text{ resonance}$$

direct ρ -hadron scattering

$$\text{diagram showing } \rho \text{ meson interacting with a hadron } h \text{ via } a_1(1260) \text{ resonance}$$

strong broadening of ρ spectral function

→ baryons are important



In-medium vector mesons (2)

QCD langrangian contains subgroup $SU_L(n_f) \times SU_R(n_f)$

- chiral symmetric in limit of vanishing quark masses
 - lattice QCD: dynamical formation of $\langle q\bar{q} \rangle \sim \Delta_{l,s}$
breaks chiral symmetry
 - profound effect on chiral partners

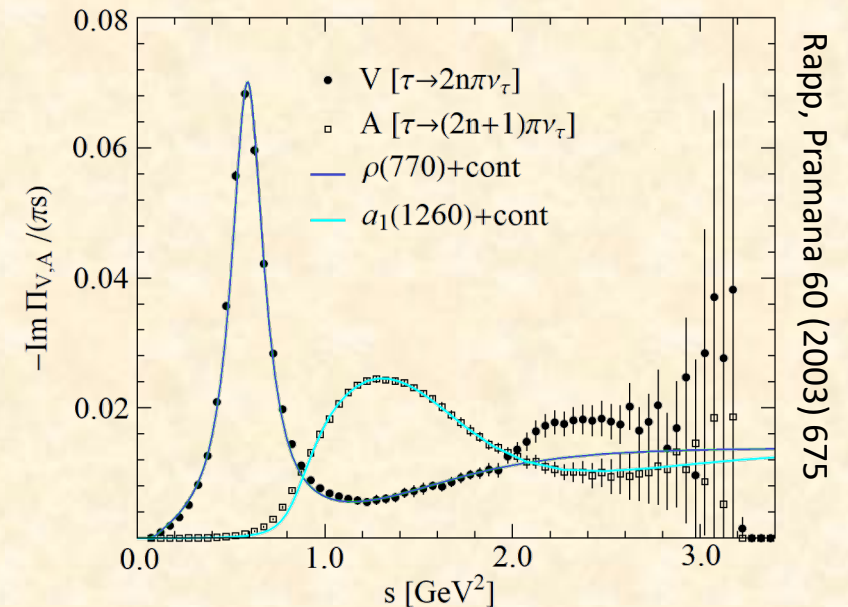
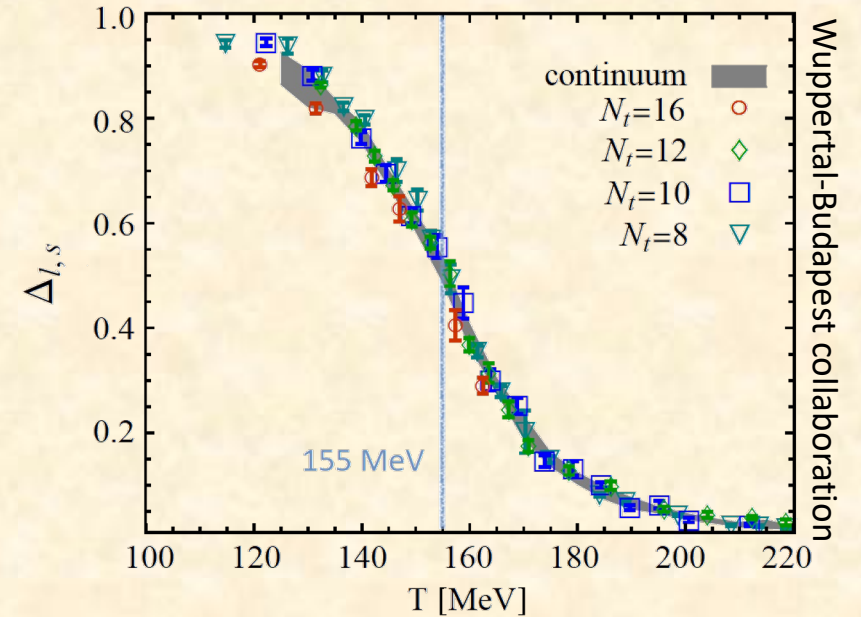
$$\langle q\bar{q} \rangle = \langle q_L \bar{q}_R + q_R \bar{q}_L \rangle$$

significant mass splitting between chiral partners

$\rho(770)$ - $a_1(1260)$, nucleon(940) - N(1535), σ - π

- Weinberg (chiral) sum rules connect SFs to condensates:

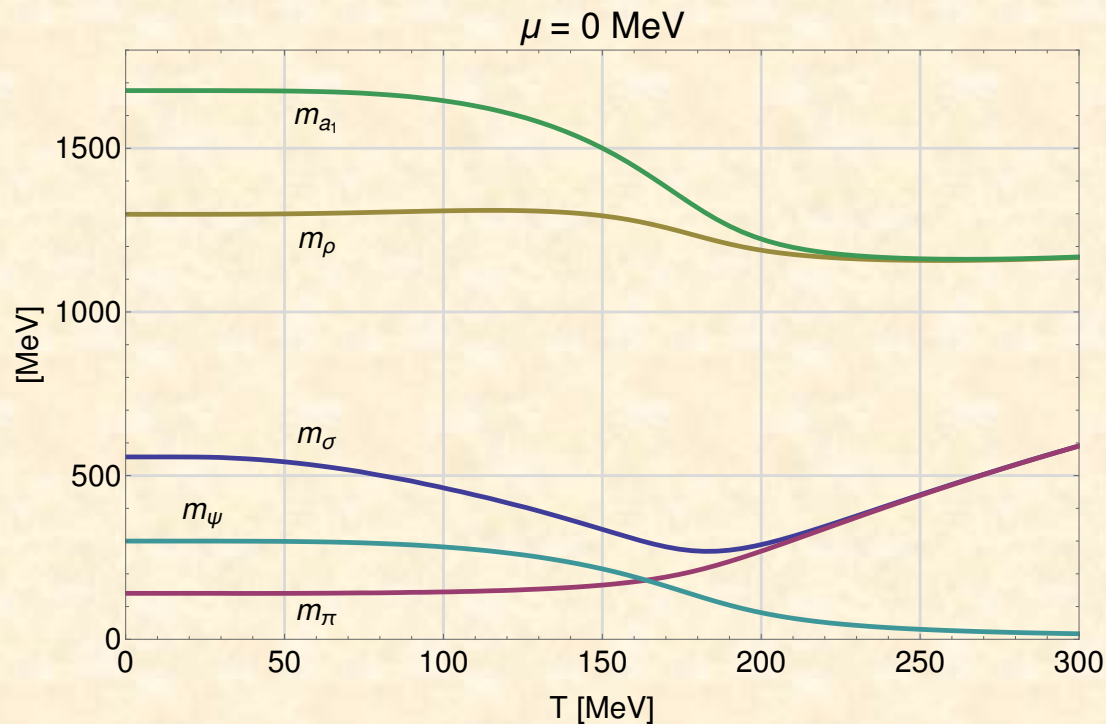
$$\int_0^\infty \frac{ds}{\pi} (\Pi_V(s) - \Pi_A(s)) = m_\pi^2 f_\pi^2 = -2 m_q \langle q\bar{q} \rangle$$



Chiral symmetry restoration

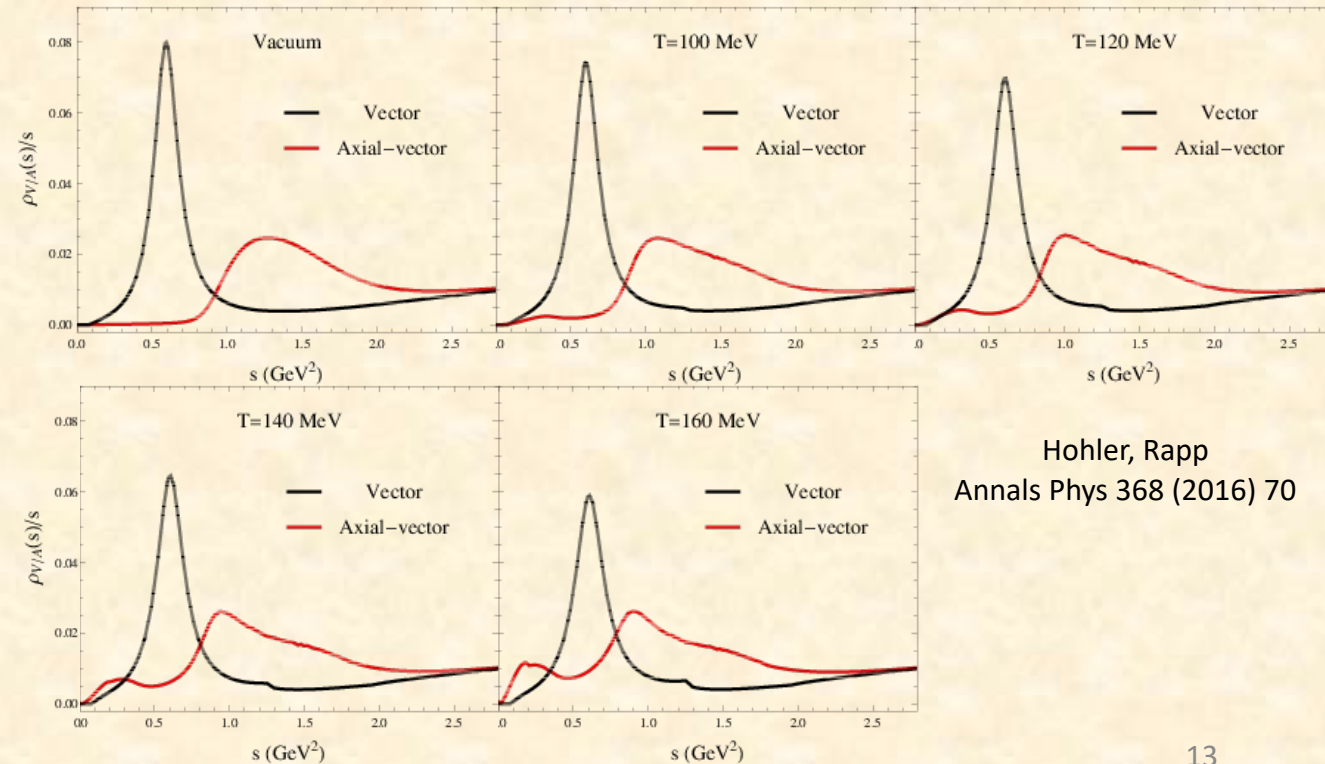
- restoration of chiral symmetry manifests itself in mixing of V and A correlators
- ρ mesons melts in hot matter while a_1 decreases and degenerates

➤ chiral mass splitting "burns off"



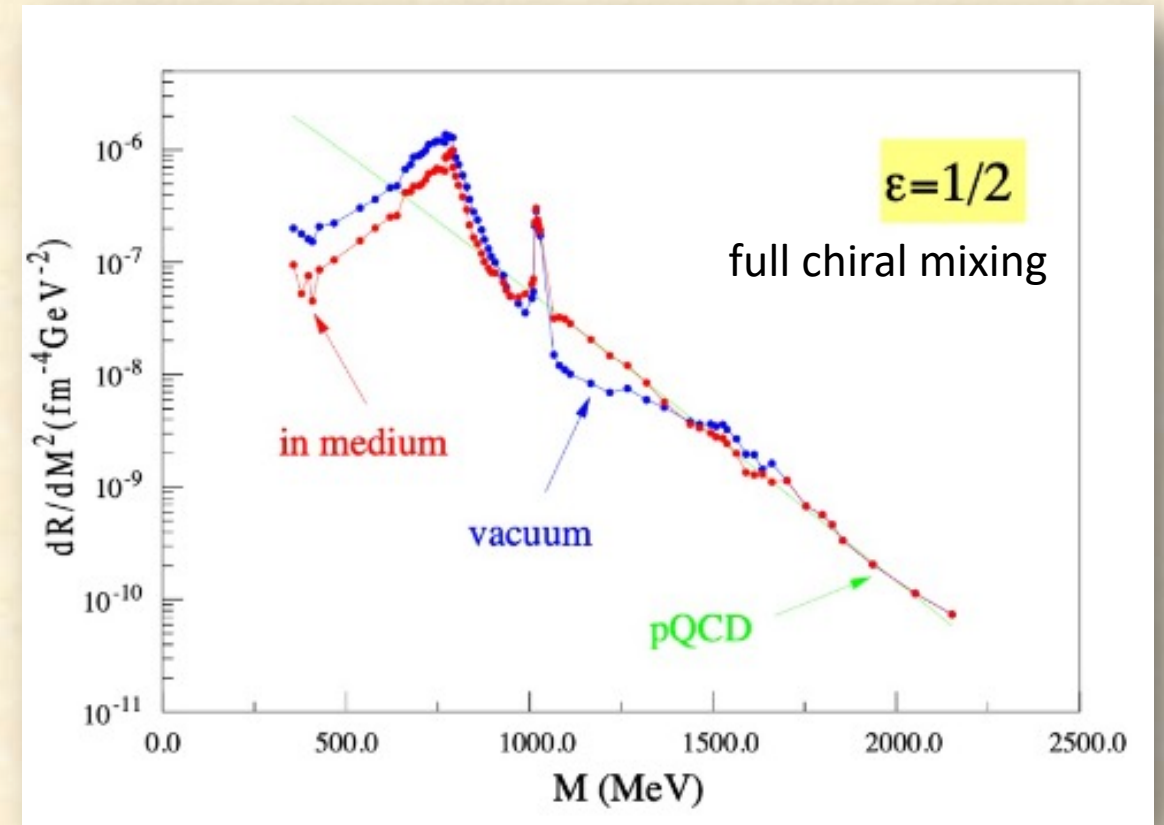
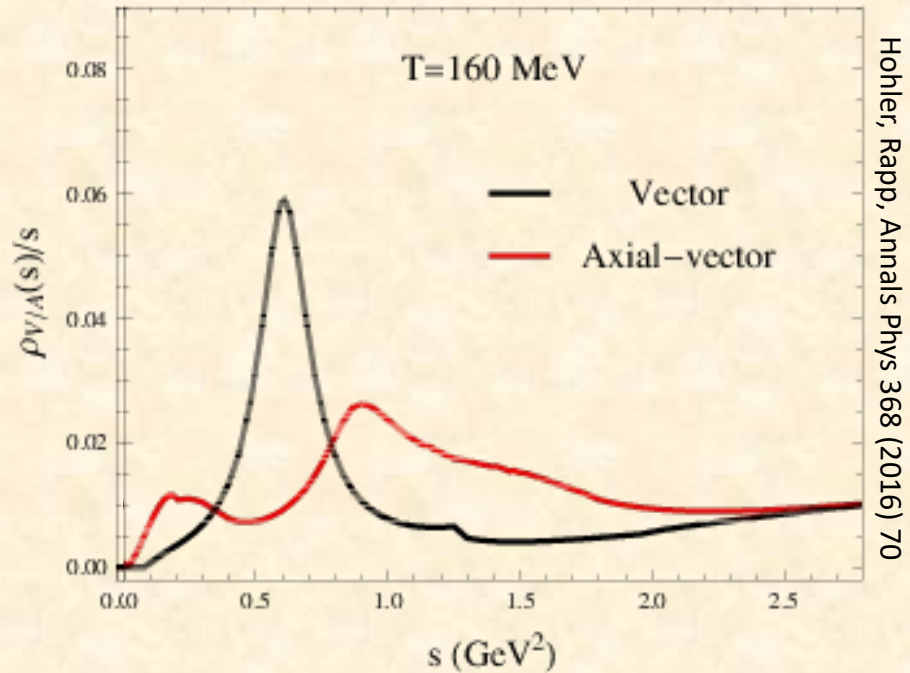
Jung, et al. PRD 95 (2017) 036020

Massive Yang-Mills in hot pion gas



Hohler, Rapp
Annals Phys 368 (2016) 70

Chiral symmetry restoration: ρ - a_1 mixing



$$\pi a_1 \rightarrow \rho' \rightarrow l^+ l^-$$

- mixing “moves strength from the axialvector to the vector channel” Rapp, Wambach, Adv. Nucl.Phys. 25 (2000) 1

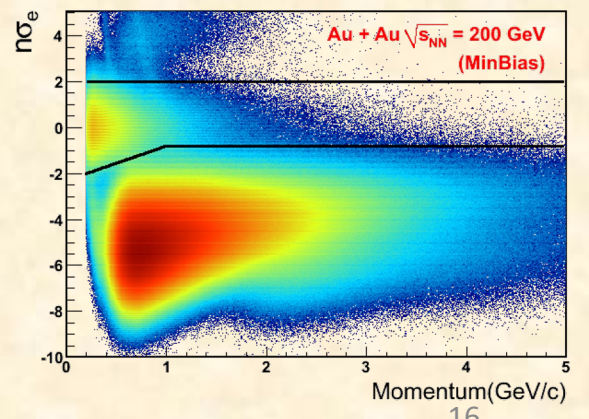
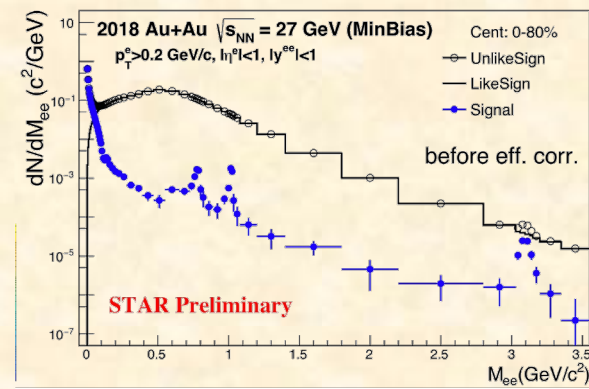
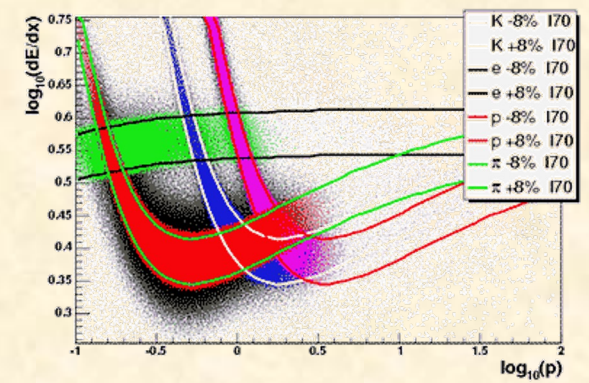
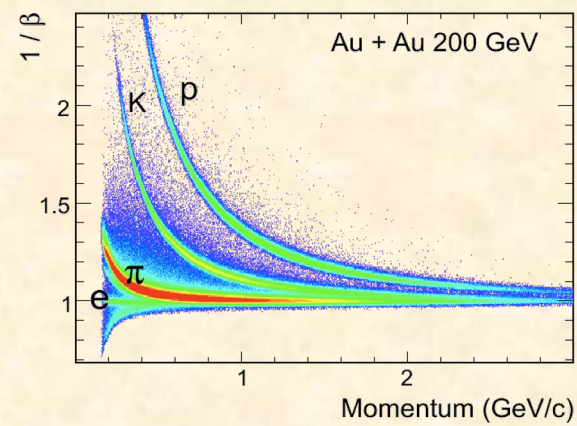
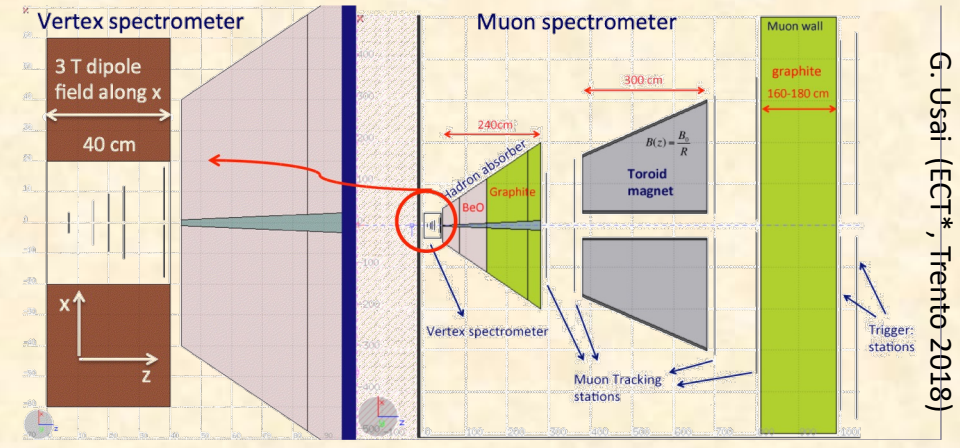
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The experimental challenge

- di-leptons need large acceptance + high purity PID
 - good momentum resolution
- electron PID from a combination of a tracker with time-of-flight (velocity), energy loss measurements (dE/dx), or RICH (very high γ thresholds) information.
- muon PID from employing hadron absorber with tracking before and after the absorber
- dileptons are rare probes: production rate is very low
 - for example: $\frac{\rho \rightarrow e^+e^-}{\rho \rightarrow \pi^+\pi^-} \sim 5 \times 10^{-5}$
- large combinatorial background
 - photon conversions from detector materials
 - Dalitz decays from light mesons
 - purity of muons, “fake” muons from weak decays
 - signal/background can be as low as 1%



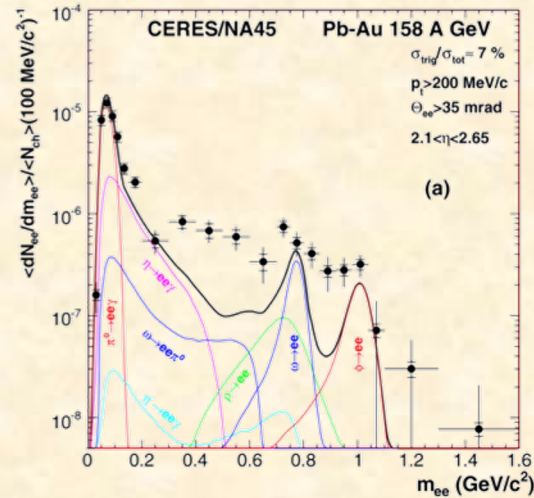
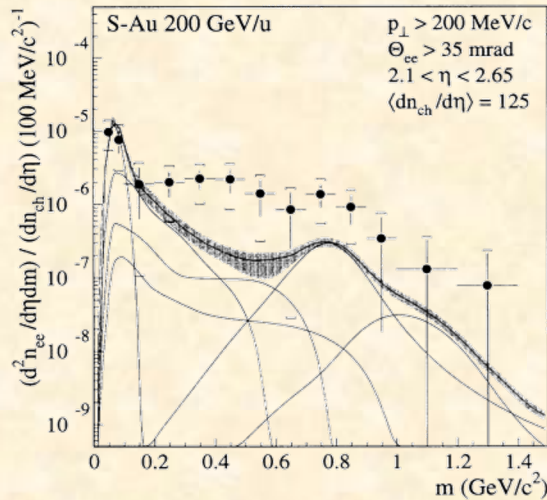
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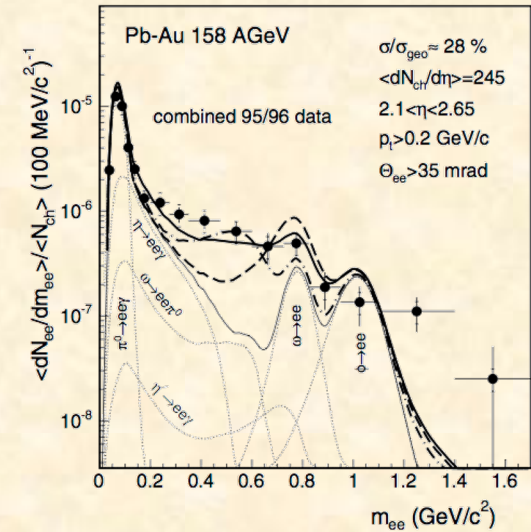


SPS dielectrons spectra (CERES)

First observation of a significant LMR enhancement – PRL 75 (1995) 1272



PLB 666 (2008) 425



EPL C41 (2005) 475

Vacuum ρ unable to describe this data

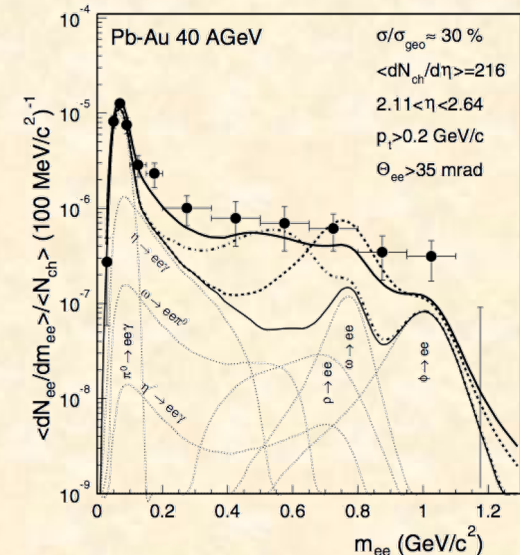
➤ Introduce in-medium modifications

- decrease of ρ mass (Brown-Rho)
 - mass expected to scale with q-qbar condensate
- broadening of ρ spectral function (Rapp-Wambach)
 - hadronic (baryons) scattering

Both rely on high baryon densities

Both showed good agreement with 158 and 40 AGeV

dashed = vacuum ρ ; dash-dotted = DM; solid = RB

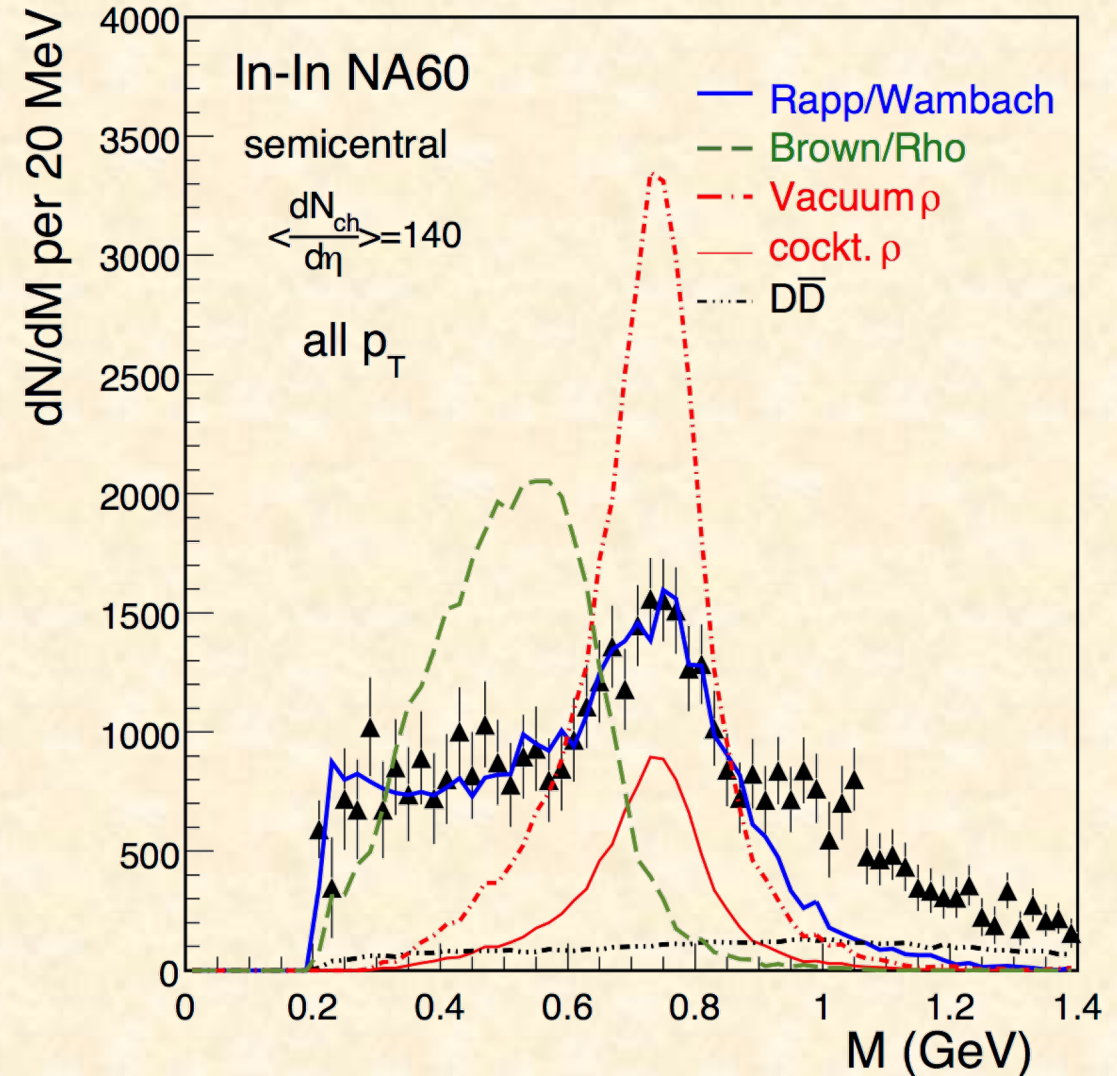


PRL 91 (2003) 042301

SPS dimuons spectra (NA60)

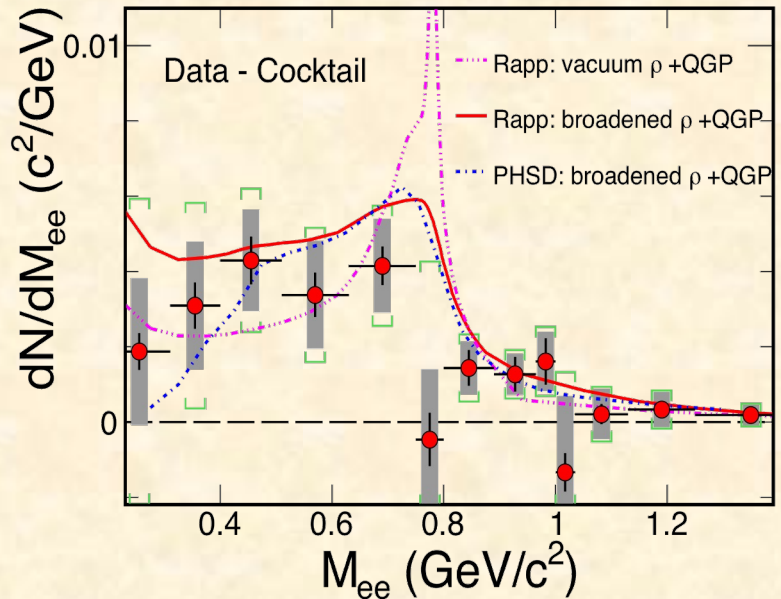
Excess in LMR $\mu^+\mu^-$ – EPJ C61 (2009) 711

- rules out: dropping-mass scenario
- very good agreement with Resonance Width Broadening for $M_{\mu\mu} < 0.9 \text{ GeV}/c^2$



RHIC dielectron spectra at 200GeV

✓ STAR



PRC 92 (2015) 024912

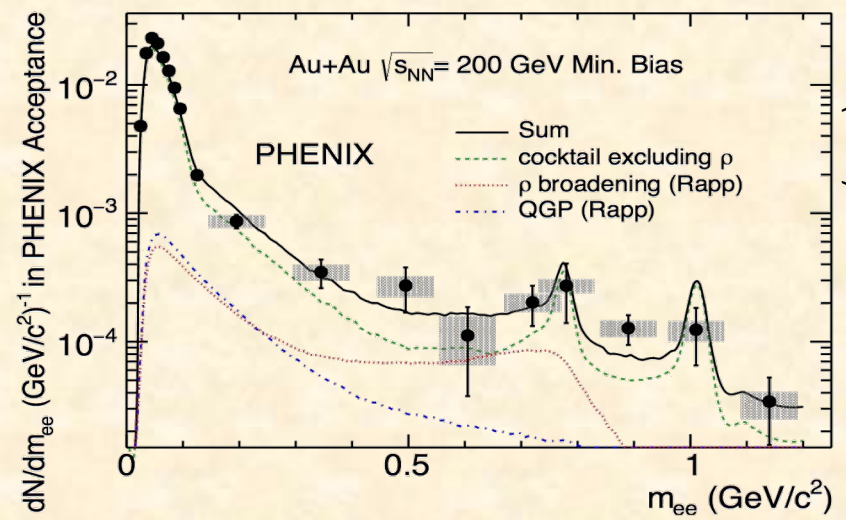
R. Rapp, PRC 63 (2001) 054907
 O. Linnyk et al., PRC 85 024910 (2012)

- Data does not support vacuum ρ
- Within uncertainties agreement between experiment and theory

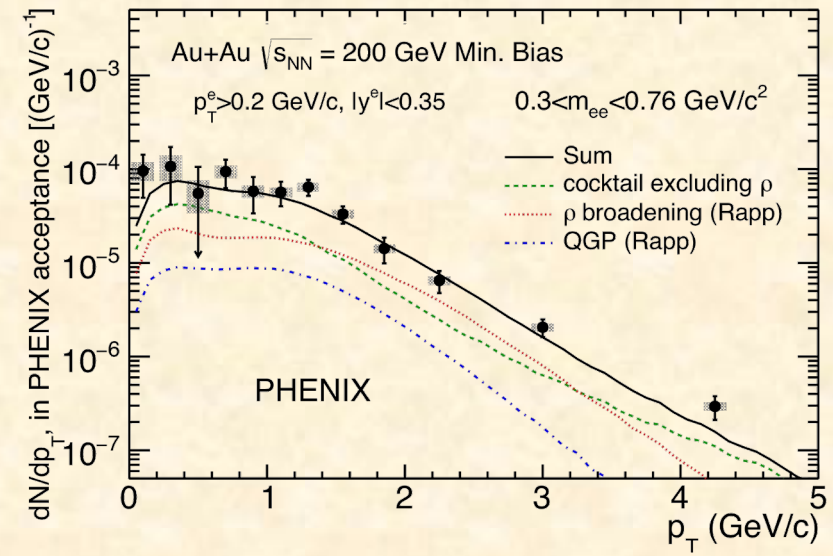
TABLE VIII. Reduced χ^2 for model calculations compared to the excess data in the invariant-mass region of 0.3–1.0 GeV/c².

Model	χ^2 /ndf	p value
Rapp: vacuum ρ + QGP	41.3/8	2.4×10^{-7}
Rapp: broadened ρ + QGP	8.0/8	0.32
PHSD: broadened ρ + QGP	16.5/8	0.040

✓ PHENIX

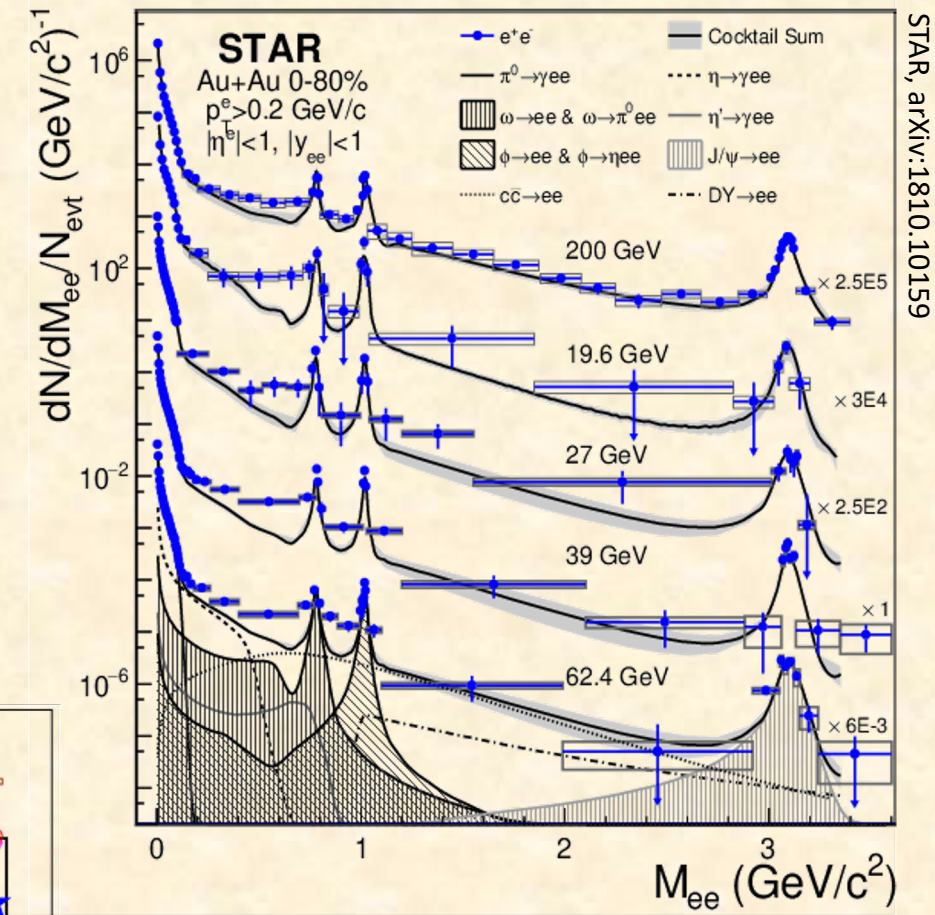
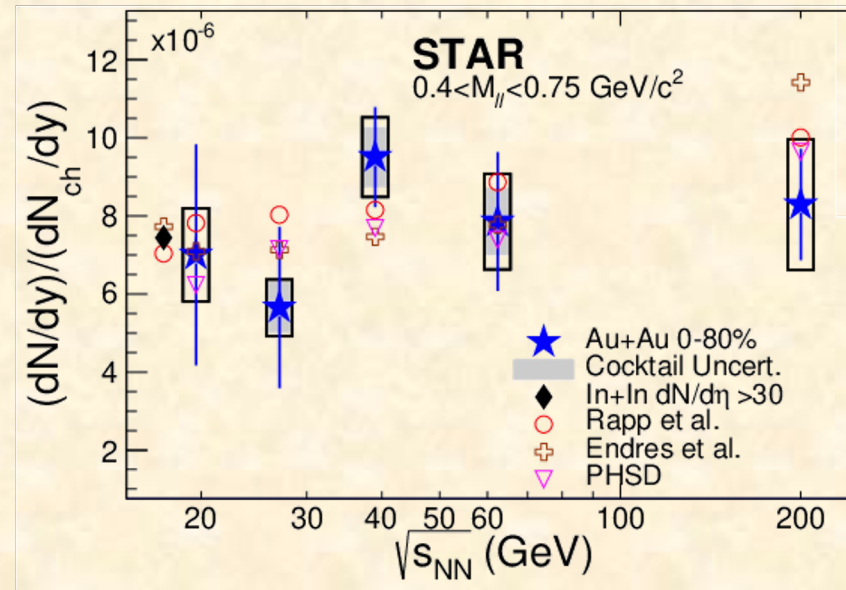
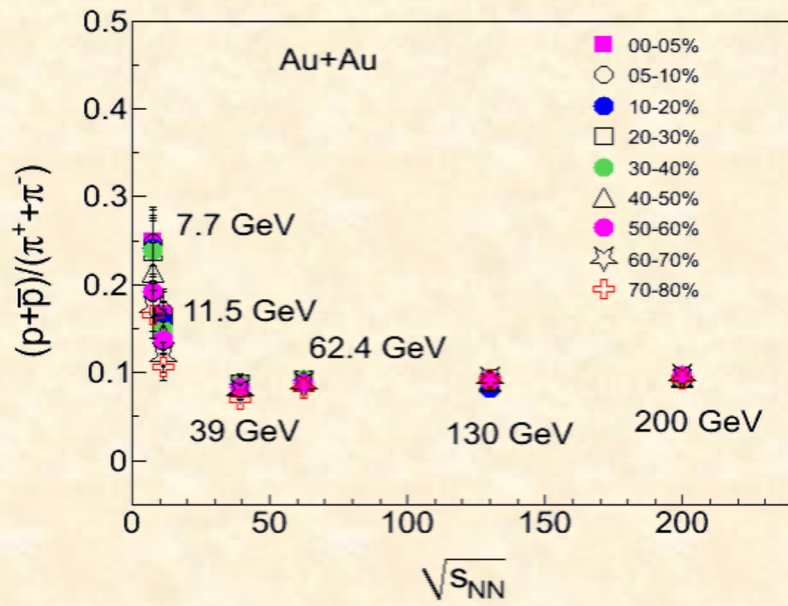


PRC 93 (2016) 014904



RHIC dielectron production from BES

- Excess established at RHIC by PHENIX & STAR
 - in-medium modification?
 - indications of chiral symmetry restoration?
- RHIC Beam Energy Scan
 - explore low-mass range down to SPS energies
 - opportunity to determine excitation function
 - dependence on temperature, total baryon density, and medium lifetime
 - normalized excess yield shows no significant $v_{s_{NN}}$ dependence
 - nor does the total baryon density
- BES Phase 1: limited precision to constrain model assumptions
 - especially for $v_{s_{NN}} < 19$ GeV



STAR, arXiv:1810.10159

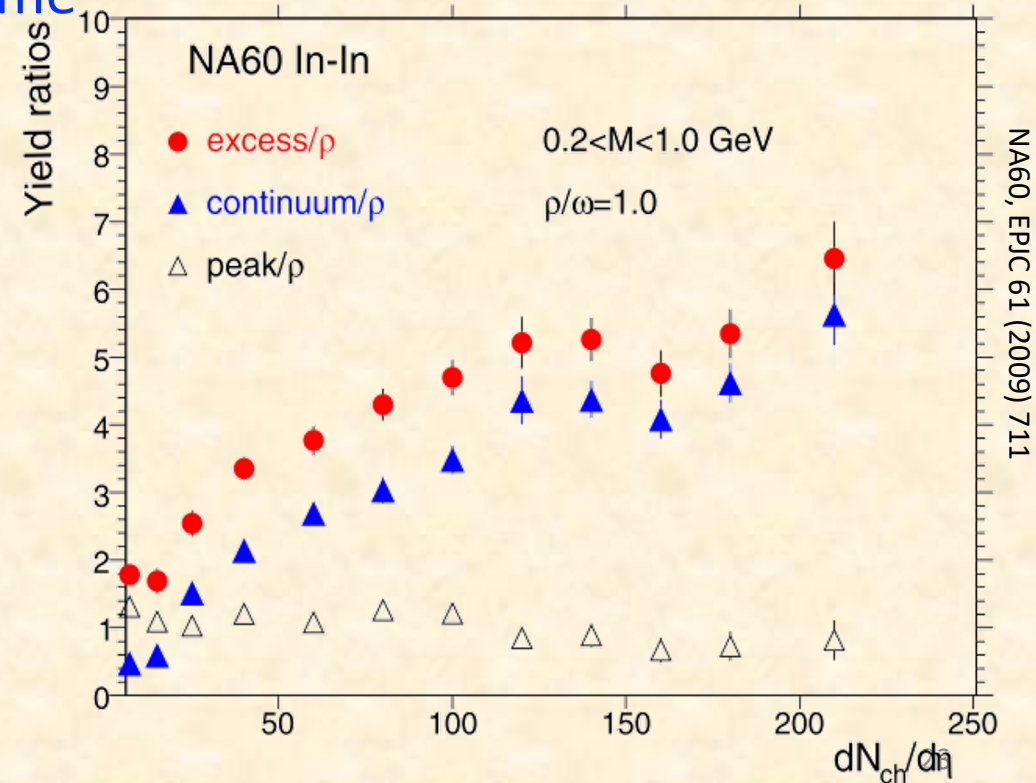
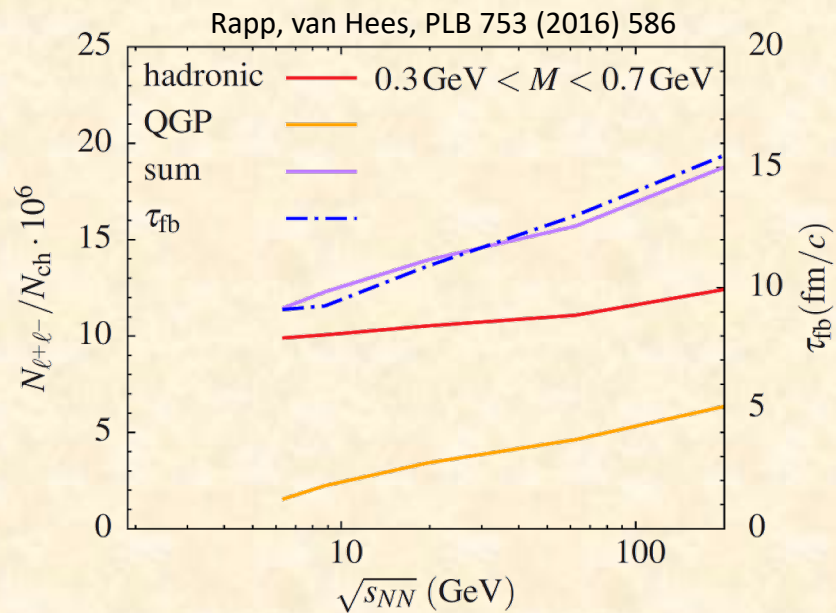
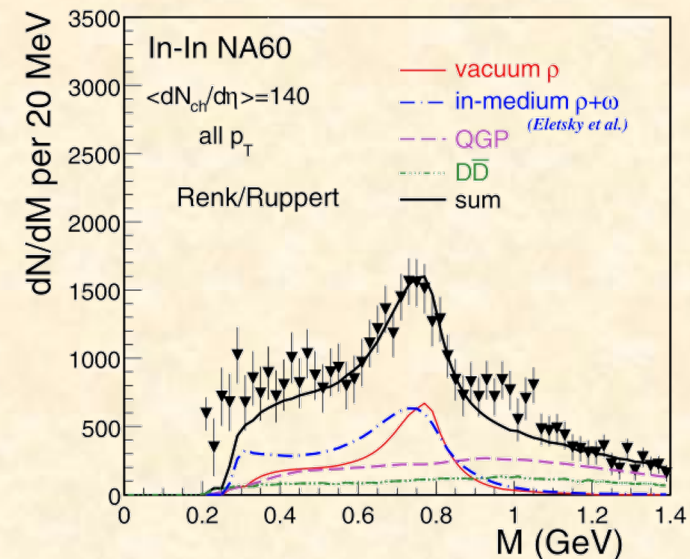
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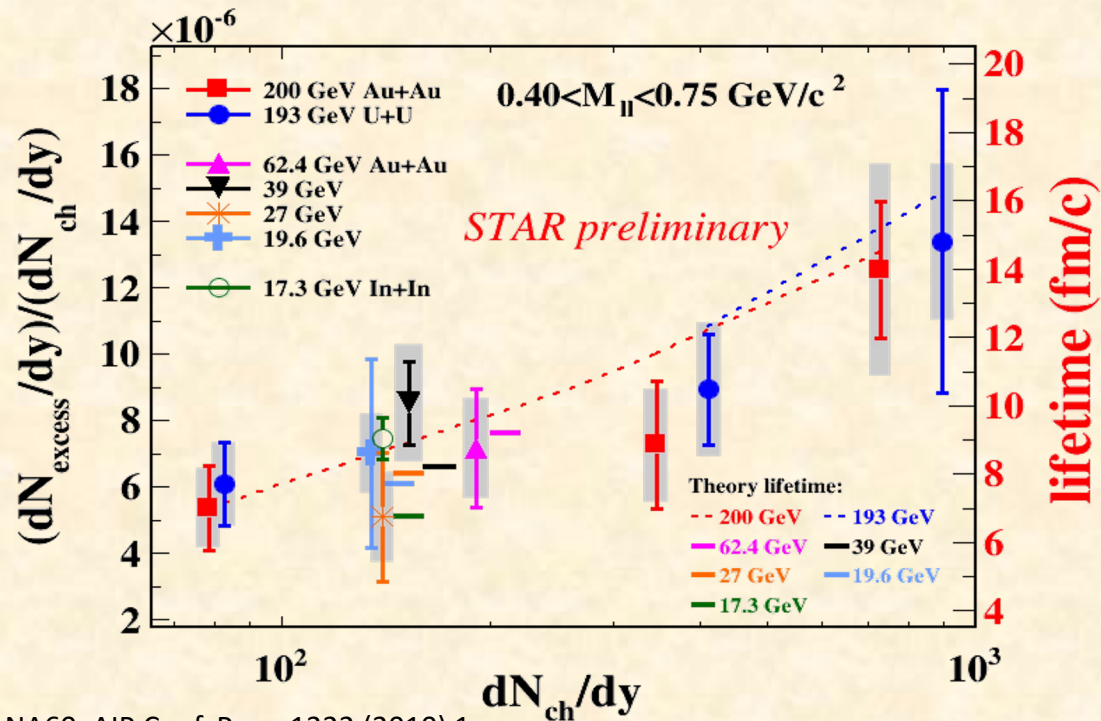
Dileptons as chronometer

- ρ peak as a clock for fireball lifetimes
 - see e.g., U.Heinz, KS.Lee, PLB 259 (1991) 162
- NA60: “ ρ clock”
 - centrality dependence of excess yield
 - reaches up to 6 generations
- Normalized excess yields in LMR track medium lifetime
 - sensitive to onset of 1st order phase transition?
 - sensitive to anomalous variations in lifetime in vicinity of CP?



Dileptons as chronometer

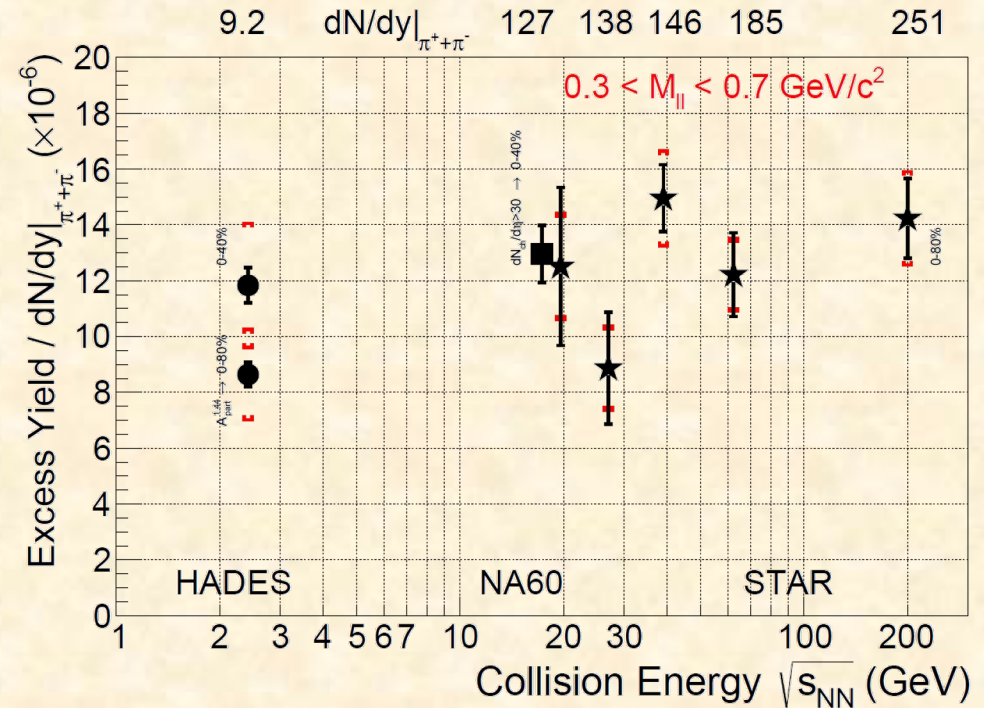
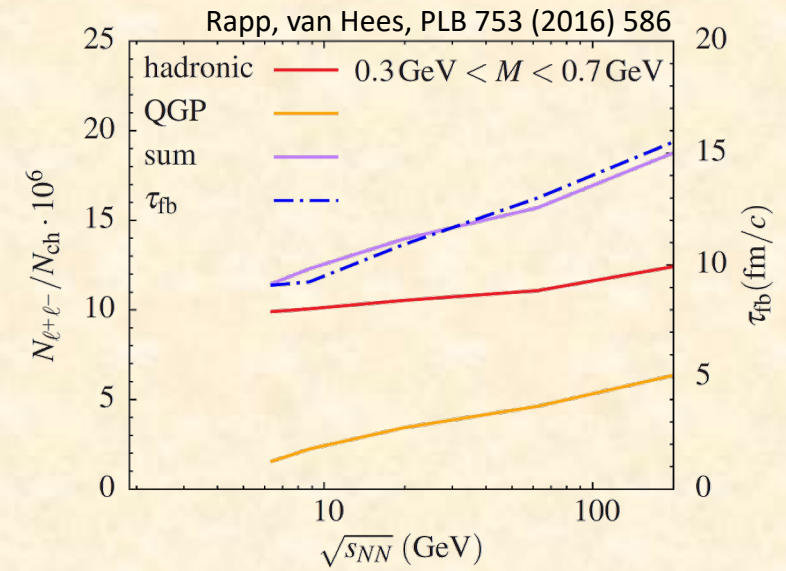
- Integrated excess radiation
 - measured below free ρ/ω mass
 - results from HADES, NA60, STAR look promising
- Experimental uncertainties are large
 - high statistics measurements needed



NA60, AIP Conf. Proc. 1322 (2010) 1

STAR, PLB 750 (2015) 64

STAR, arXiv 1810.10159



Galatyuk, JPC Conf:32 (2020) 010079

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Dileptons as thermometer

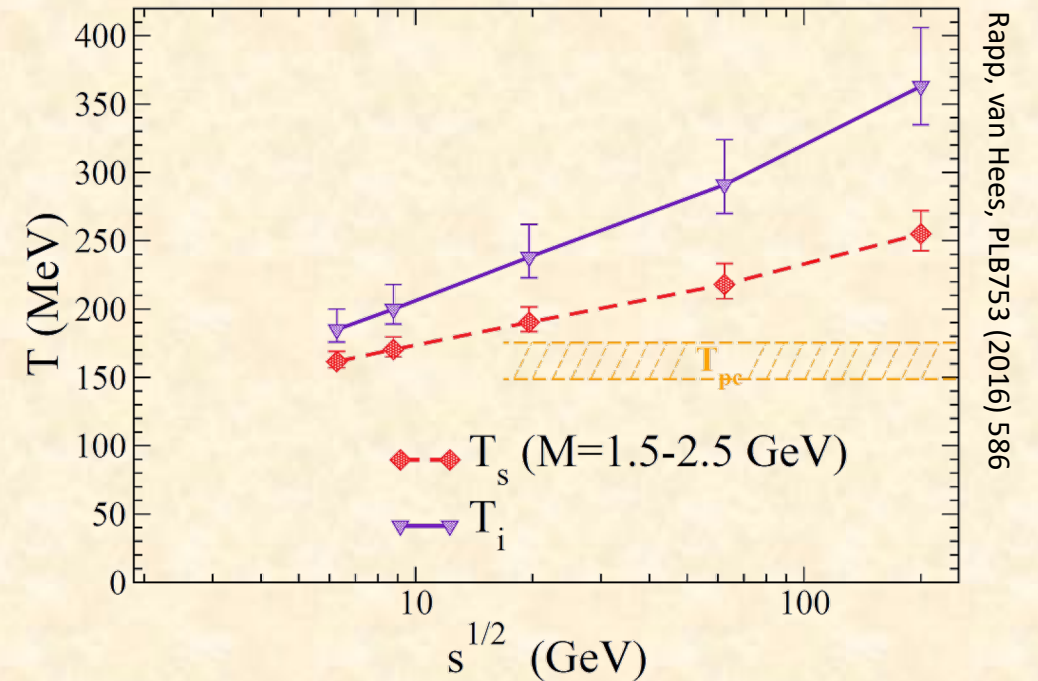
Real thermal dilepton radiation:

- LMR - dilepton spectra saturated by light vector mesons
- IMR - quark-antiquark continuum

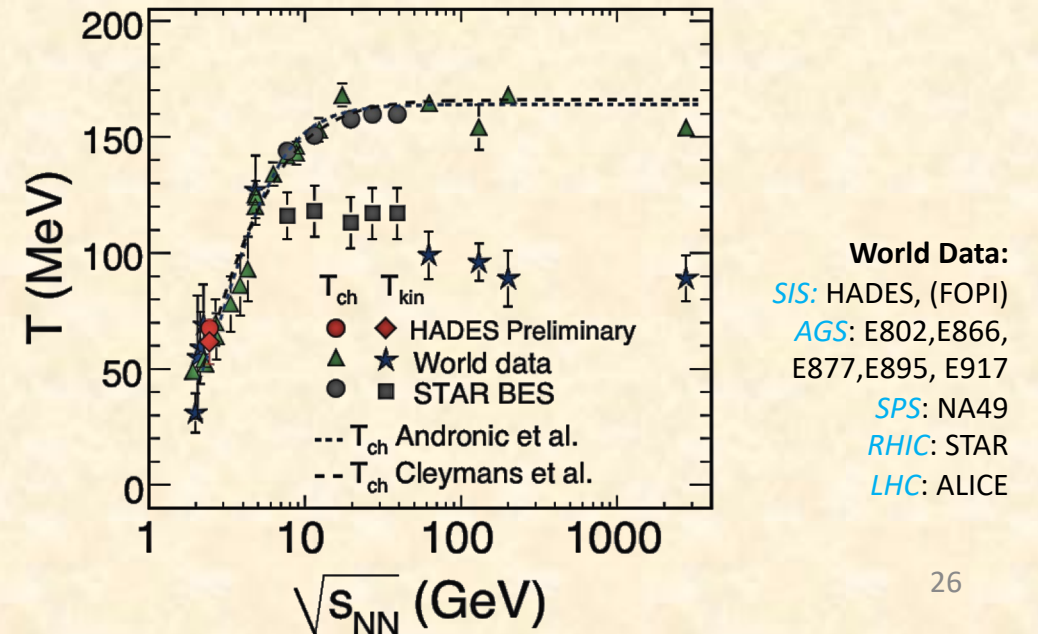
IMR dilepton rate

$$\frac{dR_{ll}}{dM} \propto \left(\frac{M}{T}\right)^{\frac{3}{2}} \exp\left(-\frac{M}{T}\right)$$

- M by construction Lorentz-invariant
- independent of flow \rightarrow no blue-shift effects
- average over the system evolution
- Other bulk temperature measurements rely on hadron yields and spectra
 - chemical and kinetic freezeout
 - separation between T_{chem} and T_{kin} grows with $\sqrt{s_{\text{NN}}}$



Rapp, van Hees, PLB753 (2016) 586

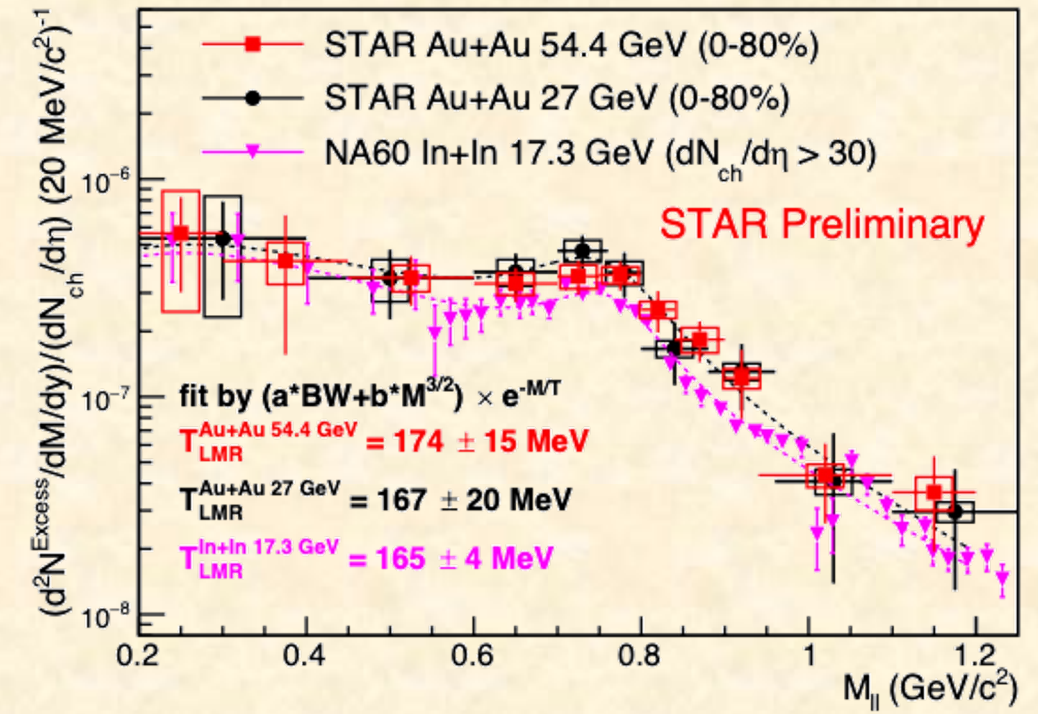
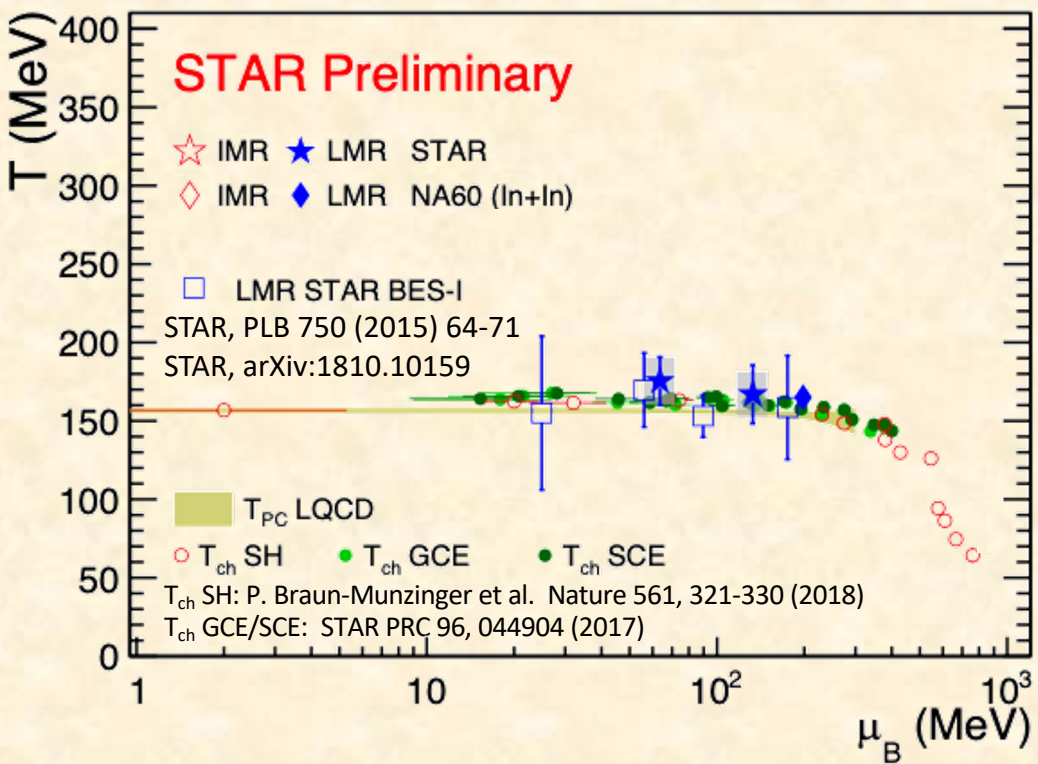


World Data:
 SIS: HADES, (FOPI)
 AGS: E802, E866, E877, E895, E917
 SPS: NA49
 RHIC: STAR
 LHC: ALICE

LMR temperature measurements

At SPS/RHIC energies

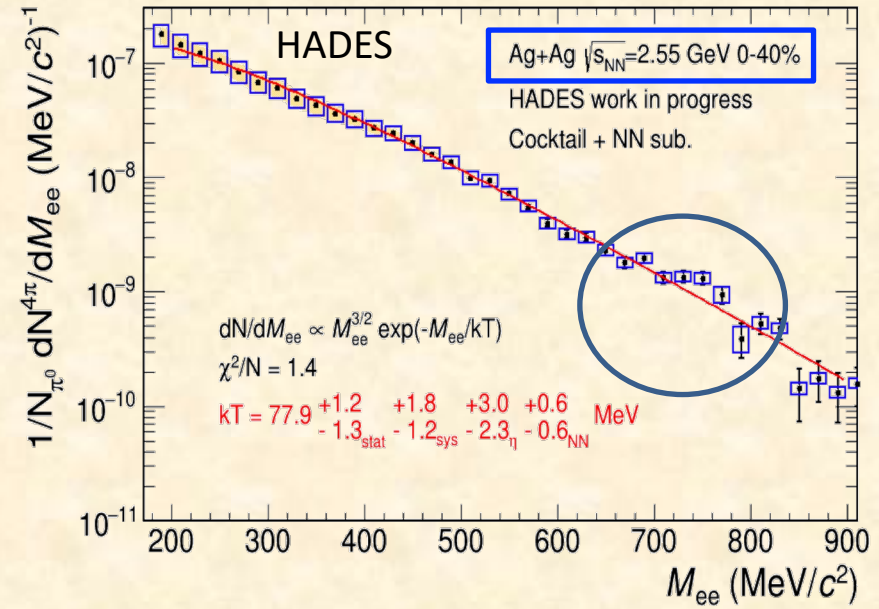
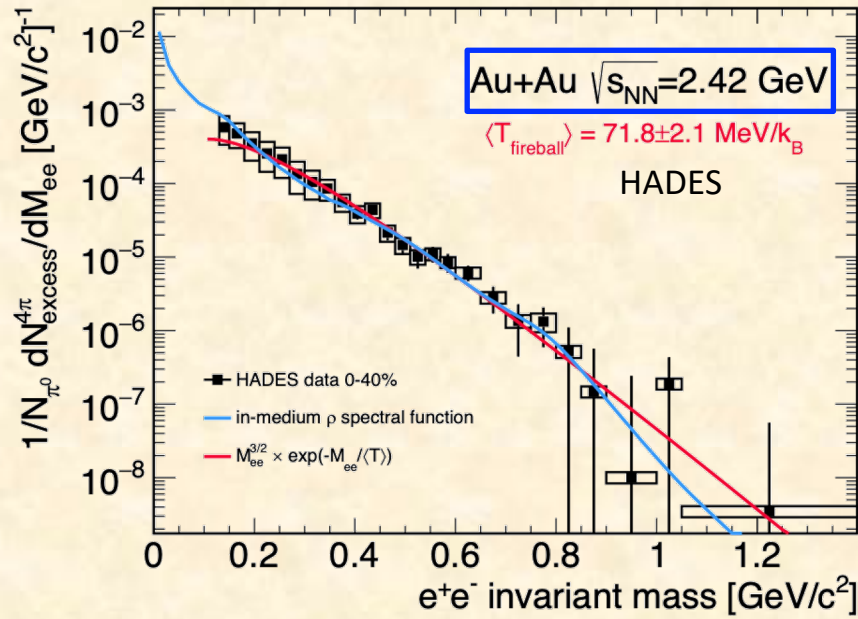
- predominantly thermal dileptons from in-medium ρ
- include Breit-Wigner in T_{LMR} fit
- recent STAR results at $\sqrt{s_{\text{NN}}} = 27$ and 54 GeV show similar mass spectra and extracted T_{LMR}
 - compared with NA60 at 17.3 GeV



- temperatures close to T_{ch} and T_{pc}
 - emitted from hadronic phase
 - predominantly around phase transition

LMR temperature measurements at higher μ_B

Nat.Phys.15 (2019) 1040

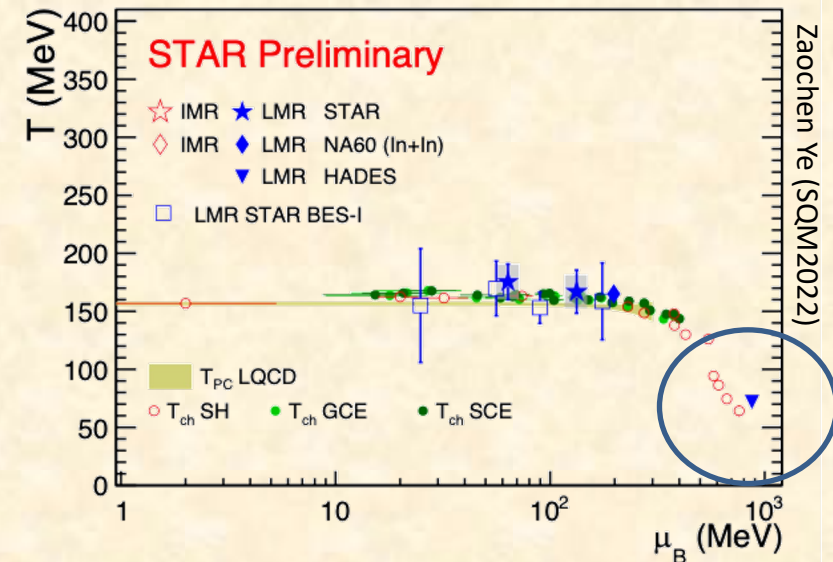


At lower energies

- higher baryon density ($\mu_B \sim 700 - 900$ MeV)
- in-medium ρ substantially modified through frequent scattering with baryons
 - almost structureless exponential distribution

➤ HADES: $T_{\text{LMR}} = 70-80$ MeV

- Au+Au @ 2.42 GeV and Ag+Ag @ 2.55 GeV



IMR temperature measurements

➤ Access to $q\bar{q}$ radiation

- NA60[†] first $\mu^+\mu^-$ measurement: $T_{IMR} = 205 \pm 12 \text{ MeV}$

- range $1.2 < M < 2.0 \text{ GeV}/c^2$

- Recent STAR IMR e^+e^- results: $T_{IMR} \sim 320 \text{ MeV}$

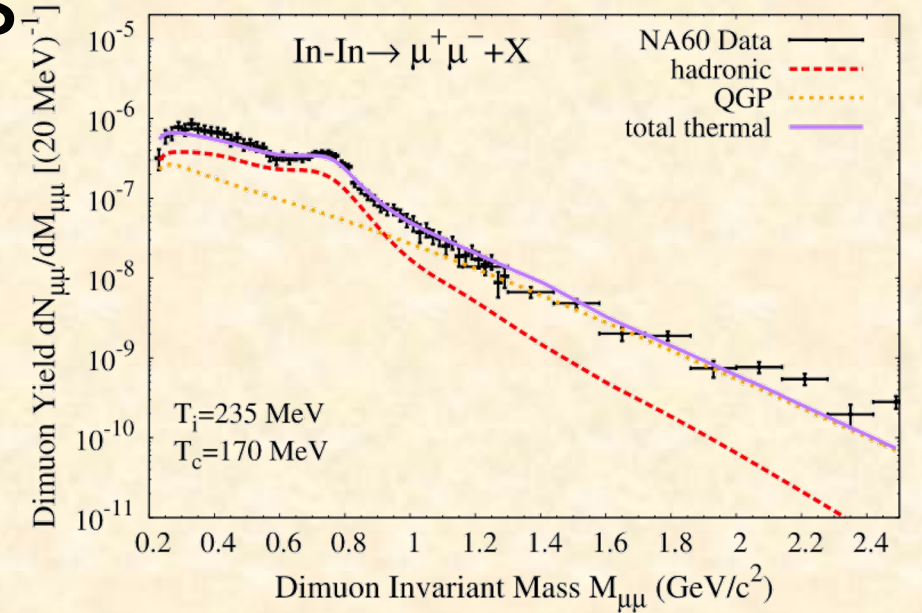
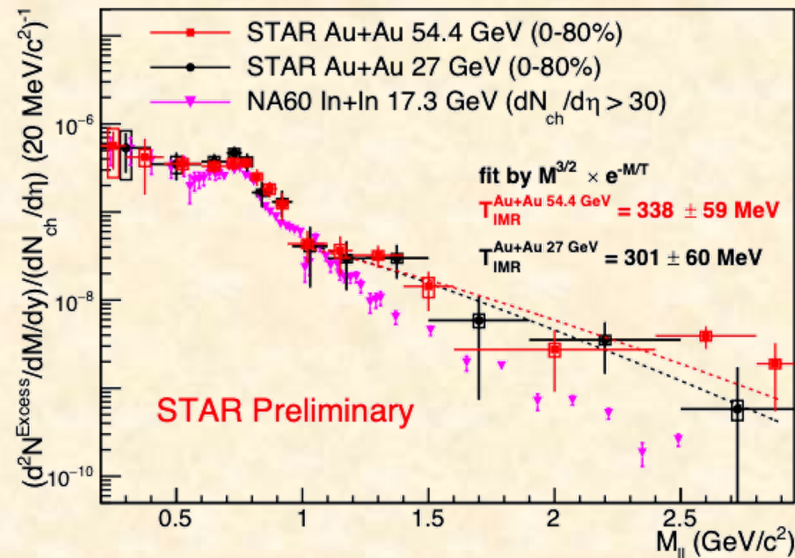
- compare with $T_{IMR}(NA60)^{++} = 246 \pm 15 \text{ MeV}$

- range $1.2 < M < 2.5 \text{ GeV}/c^2$

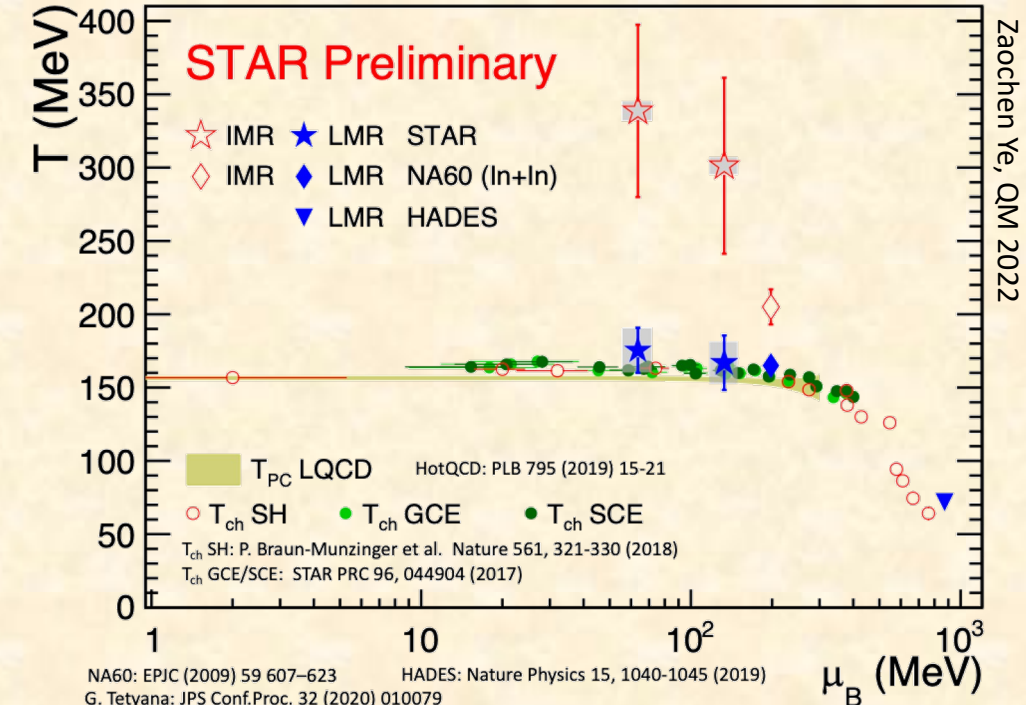
– average T_{IMR} higher due to longer lifetime?

– supported by generally higher yields

✓ Average T_{IMR} well above T_{pc}



Rapp, van Hees, PLB 753 (2016) 586



† NA60, AIP Conf. Proc. 1322 (2010) 1
 †† B. Mueller/Z. Ye, priv. comm

NA60: EPJC (2009) 59 607-623
 HADES: Nature Physics 15, 1040-1045 (2019)
 G. Tetyana: JPS Conf. Proc. 32 (2020) 010079

Outline

- QCD Phase Diagram
- What role can dileptons play?
- Some theoretical considerations
- Some experimental considerations
- **The Dilepton Multitool**
 - spectrometer, chronometer, thermometer, barometer, polarimeter, multimeter
- Future prospects



Dileptons as barometer

m_T distributions:

take medium flow into account

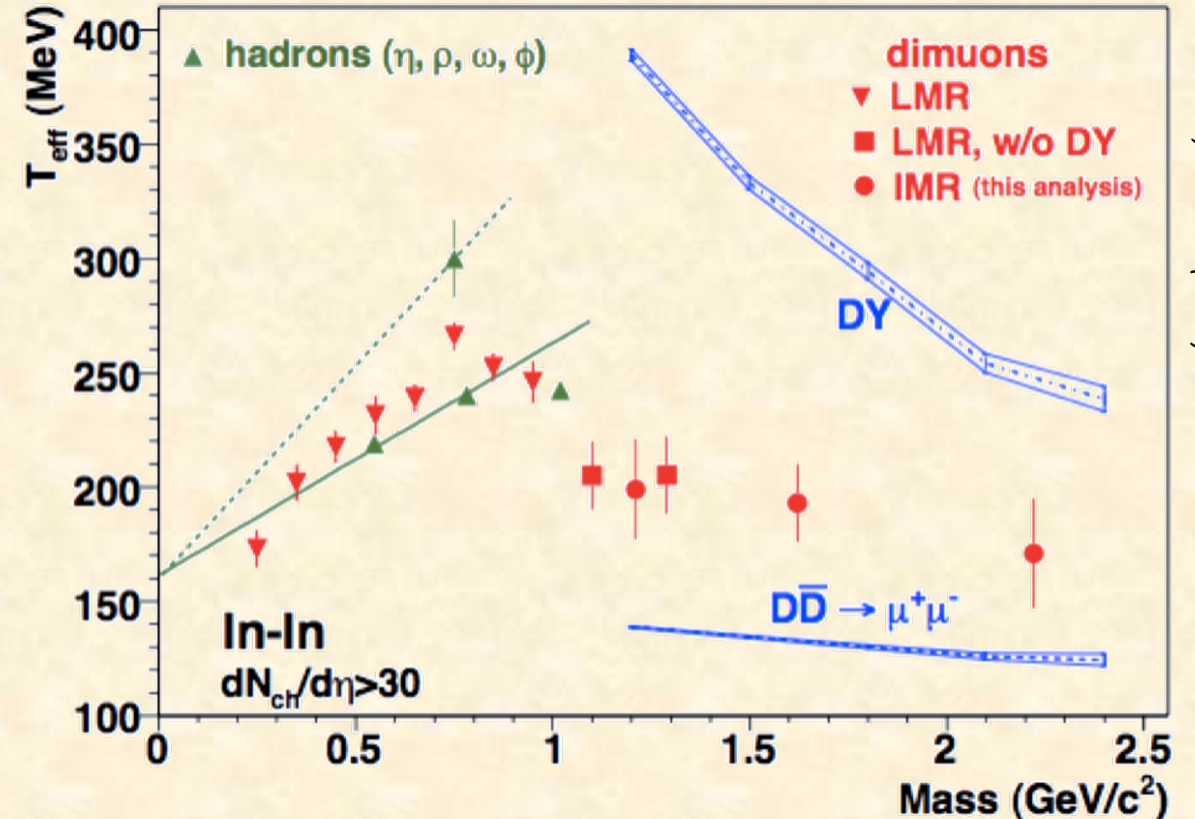
$$\frac{1}{m_T} \frac{dN}{dm_T} \propto \exp\left(-\frac{m_T}{T_{eff}}\right)$$

- LMR: inverse slopes show mass dependence

- T_{eff} linearly rises, and peaks at m_ρ
- radiation pushed by radial flow

- IMR: no indication of mass dependence

- sudden drop of T_{eff} by ~ 50 MeV
- dominant source from hadronic to partonic matter
- $T_{eff} \sim 200$ MeV

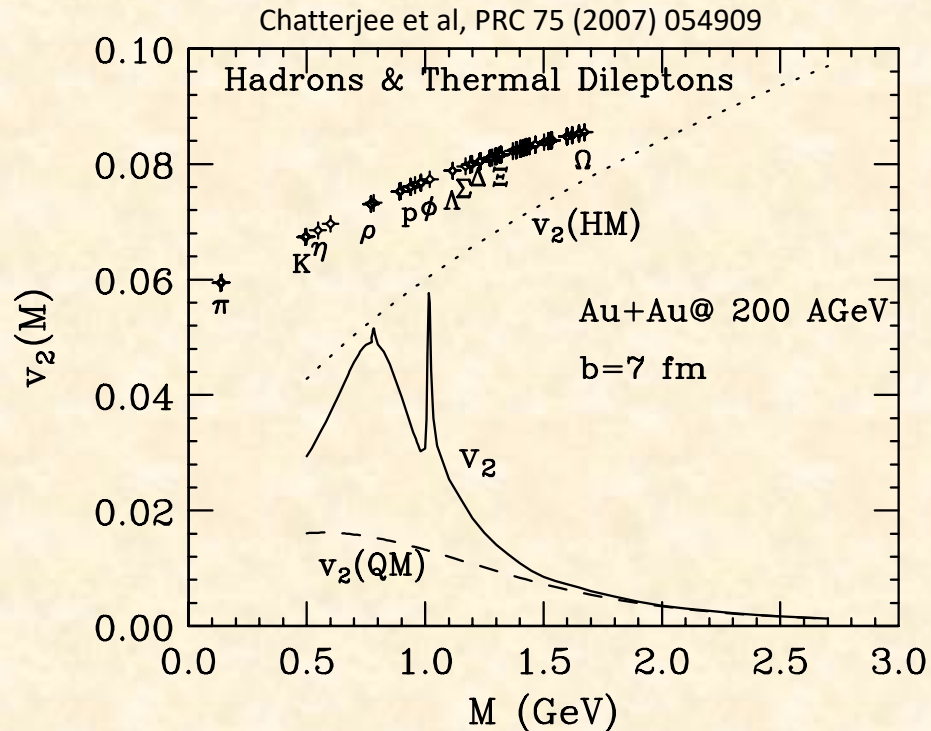


Dileptons as barometer

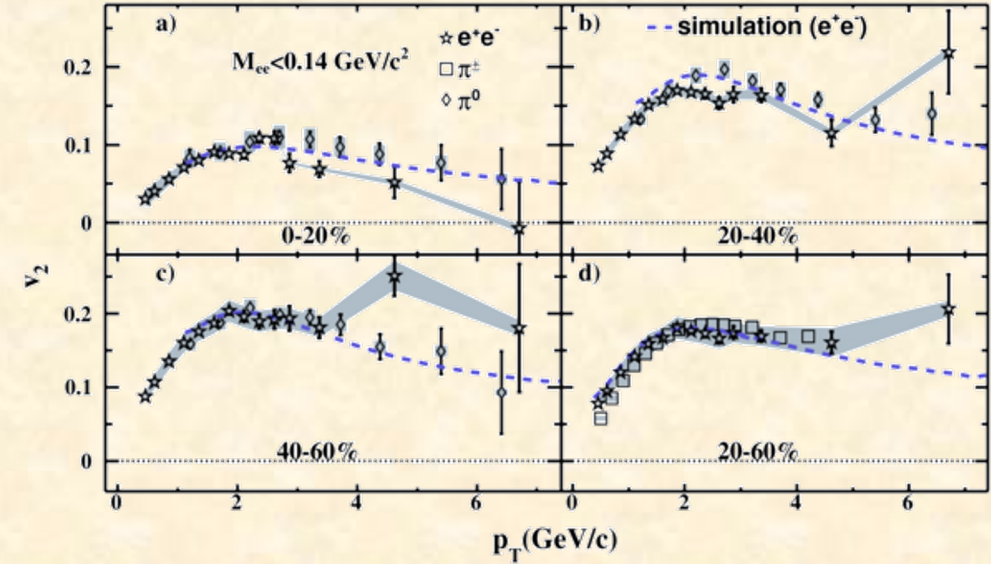
Azimuthal anisotropy

challenge: isolate v_2 of excess dielectrons

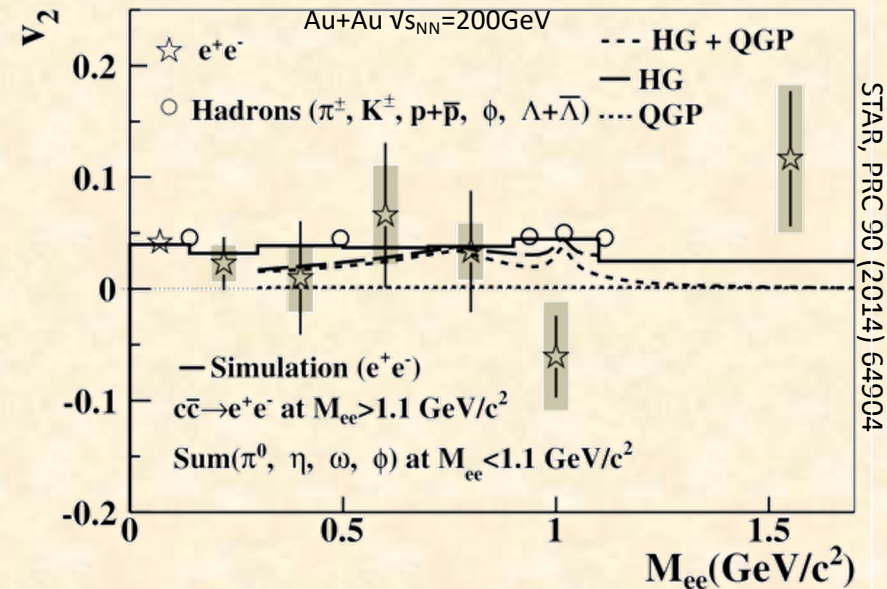
➤ to distinguish between HG and QGP need uncertainties $<4\%$...



v_2 from π^0 Dalitz decay consistent with simulations based on published $\pi^0 v_2$



STAR inclusive e^+e^-

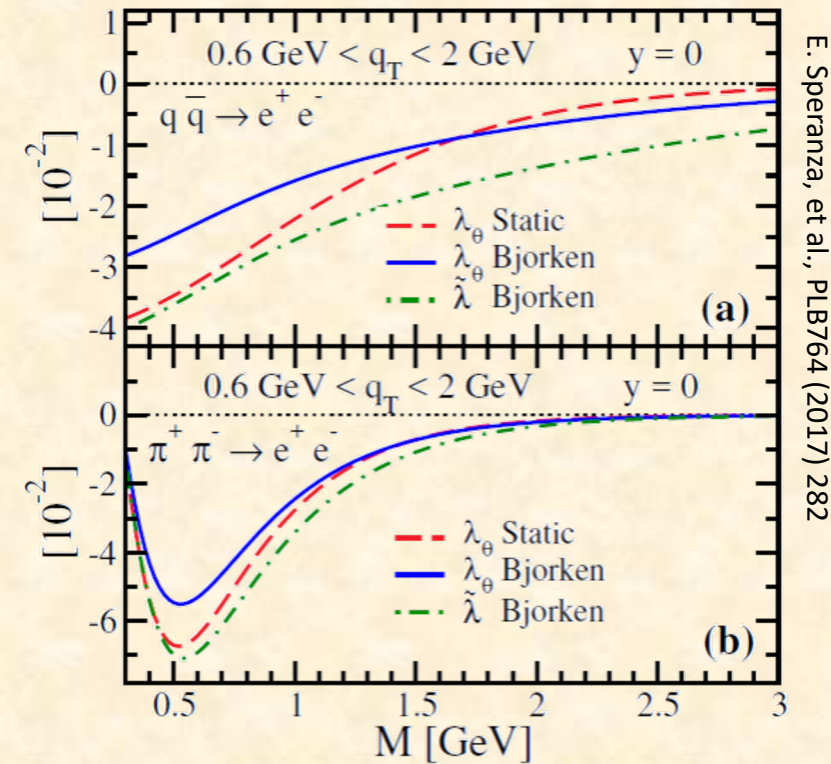


Dileptons as a polarimeter

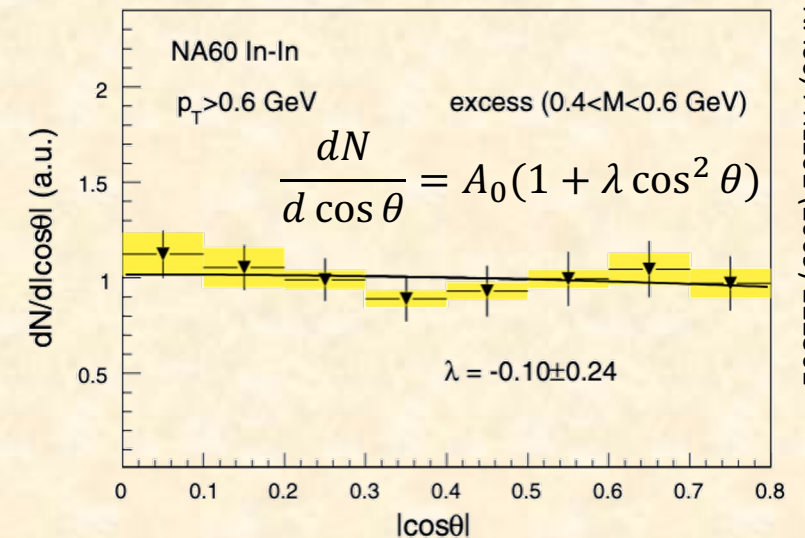
Use the angular distribution of dilepton rates

$$\frac{dR}{d^4q d\Omega_l} = N(1 + \lambda_\theta \cos^2 \theta_l + \lambda_\phi \sin^2 \theta_l \cos 2\varphi_l + \dots)$$

- anisotropy coefficients λ :
 - give info on γ^* polarization
 - relate to production mechanisms
 - e.g., $\lambda_\theta = +1$ (T) for DY, and -1 (L) in $\pi\pi$ annihilation
- integrated over M, q_T, y coefficients $\lambda_\theta \lesssim 1\%$
 - expect small, but finite polarization in a thermal system
 - consistent NA60's null finding within uncertainties
- need high-statistics future experiments
 - systematic study of all relevant process required



E. Speranza, et al., PLB764 (2017) 282



NA60, PRL102 (2009) 222301

Dileptons as multimeter: electrical conductivity

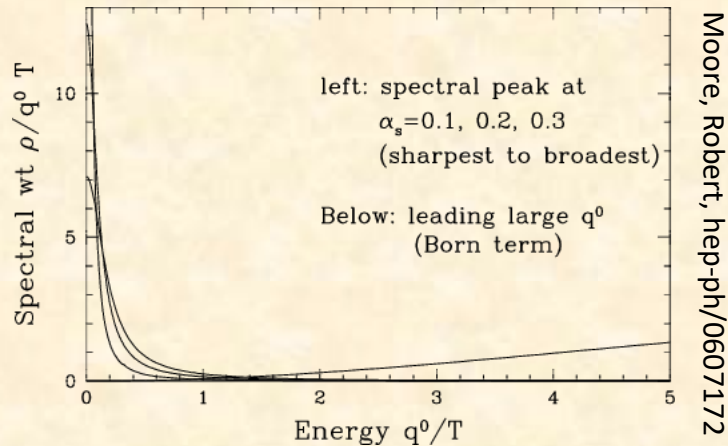
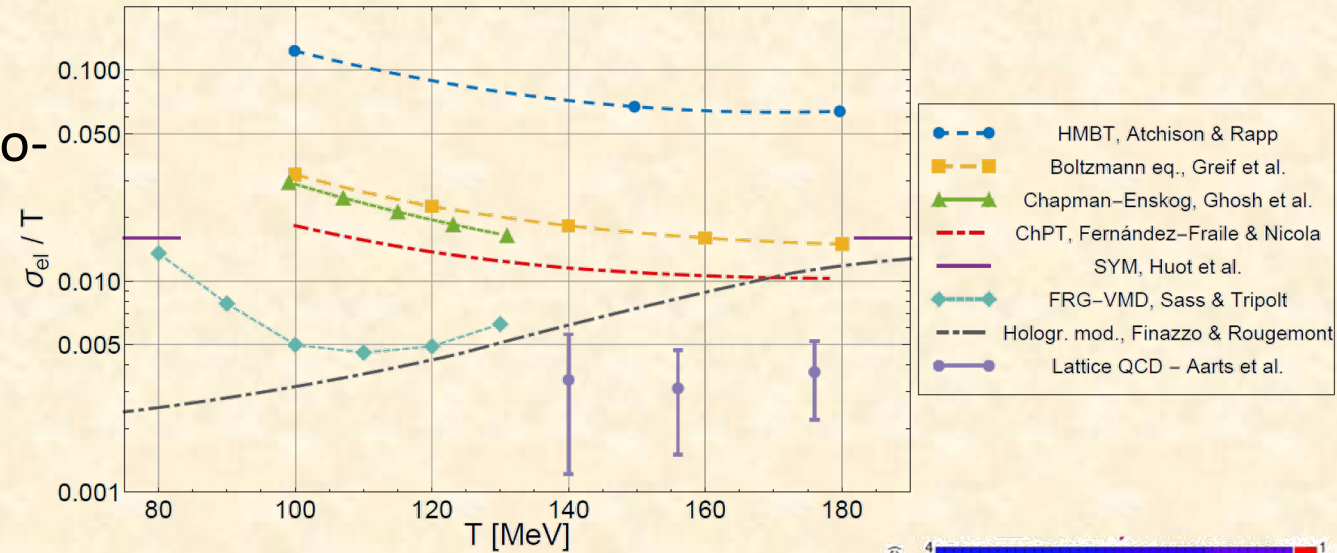
- Electrical conductivity defined as

$$J^\mu = \sigma_{el} E^\mu$$

- Can be extracted from EM correlator in the zero-momentum, low-energy limit:

$$\sigma_{el}(T) = -e^2 \lim_{q_0 \rightarrow 0} \frac{\text{Im}\Pi_{em}(q_0, \vec{q} = 0, T)}{q_0}$$

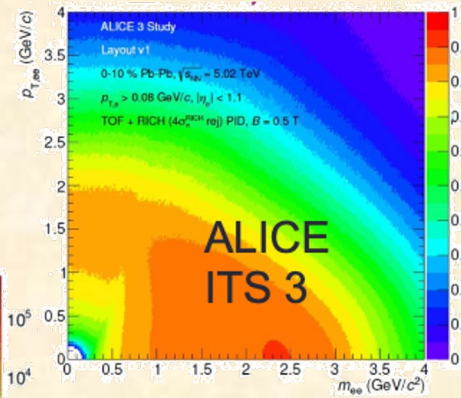
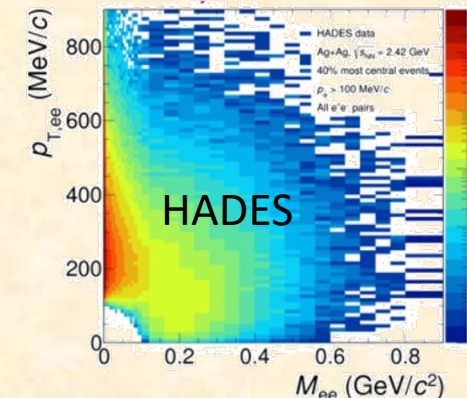
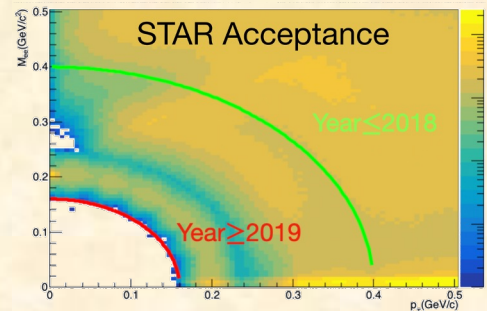
- wide range of theory predictions



Moore, Robert, hep-ph/0607172

Experimental challenge:

- low invariant mass and low p_T
- precise knowledge of (elastic) cross sections among hadrons

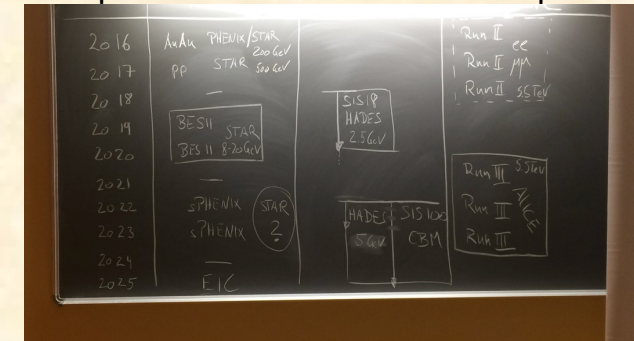


e.g. Greif, Greiner, Denicol, PRD93 (2016) 096012
Atchinson, Rapp, J.Phys.Conf.Ser.832 (2017) 012057

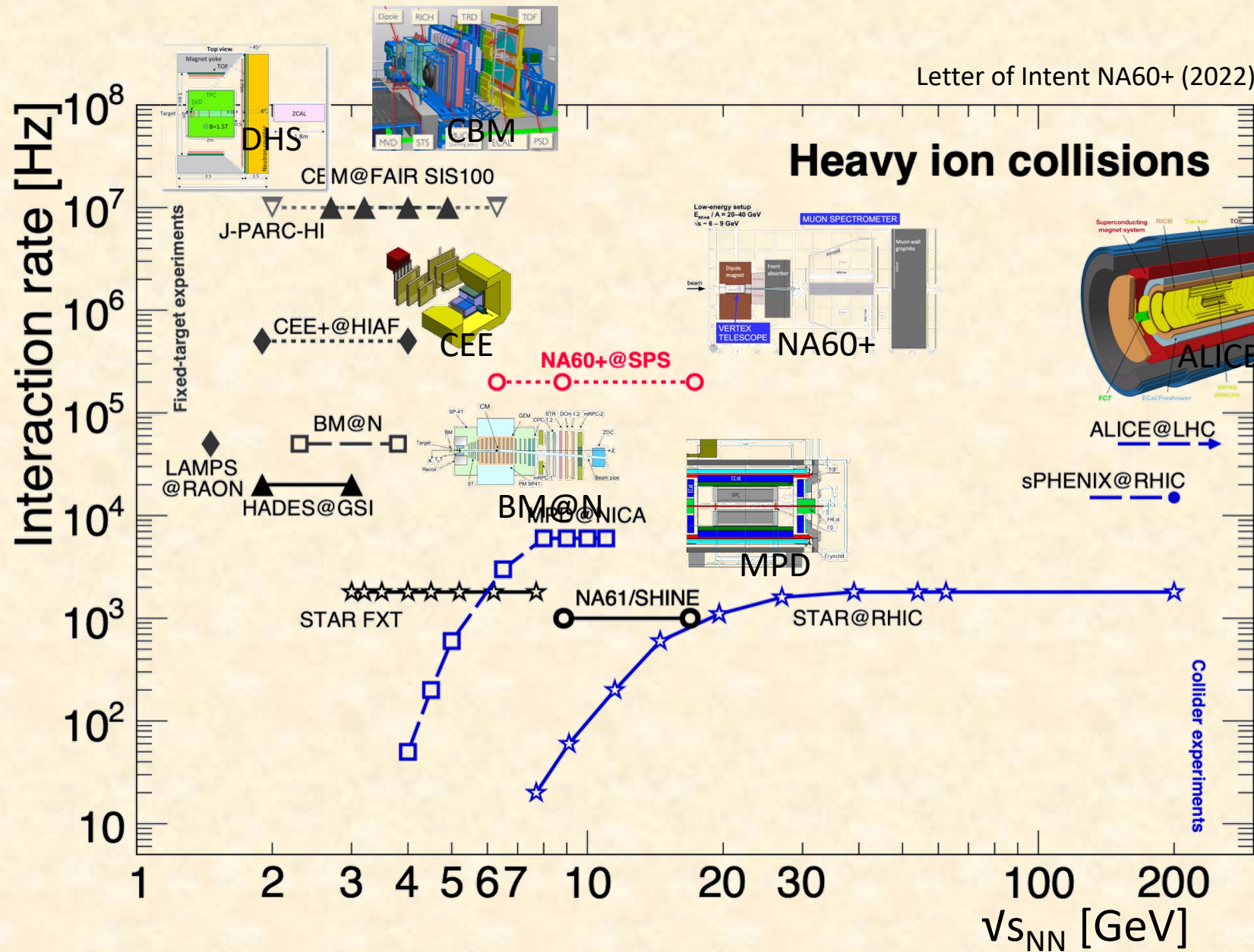
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Future Prospects



Future Measurements

- high statistics
 - high interaction rates
 - large acceptance
- precise references
 - cocktail (mesons, HF, DY)
 - multipurpose detectors
- good control on backgrounds
 - materials: conversion rejection
 - $e^+/- \mu^+/-$ purity

Summary



- Dilepton measurements provide access to wide range of unique physical observables
 - lifetime, temperature, transport properties, chiral symmetry restoration, ...
- Potential of accurate dilepton measurements is well demonstrated at SPS, SIS18, RHIC, and LHC energies
 - combined with new theoretical developments and insights
- For future experimental progress high-statistics data is key
 - an increasing world-wide effort to map out the QCD phase map and deliver its landmarks