

Introduction to Heavy Ion Collisions – Experimental

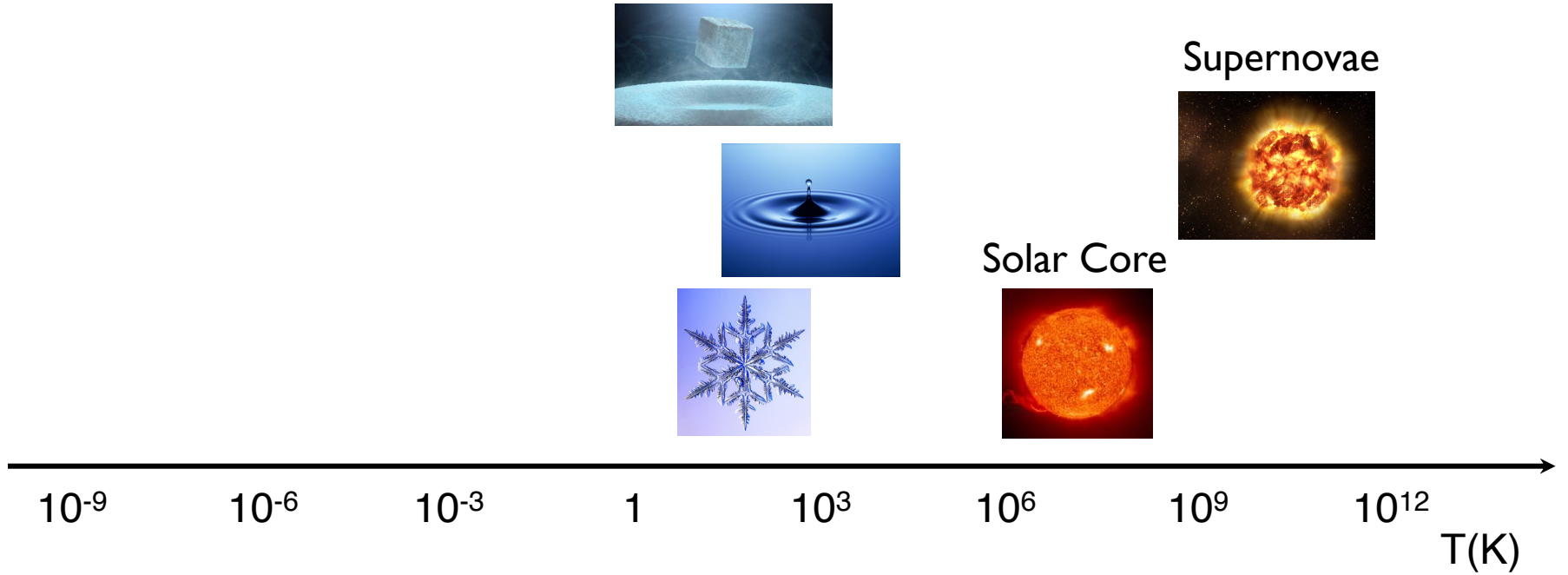
Part1 Collective Flow

陈震宇
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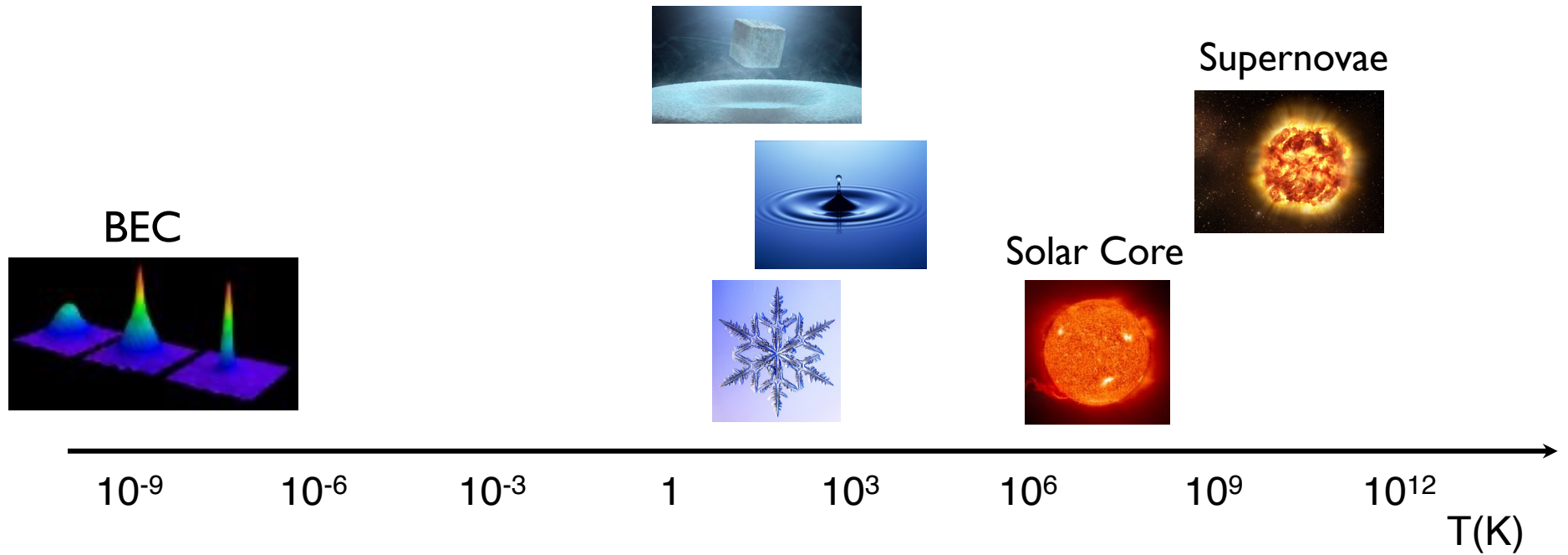


山东大学
SHANDONG UNIVERSITY

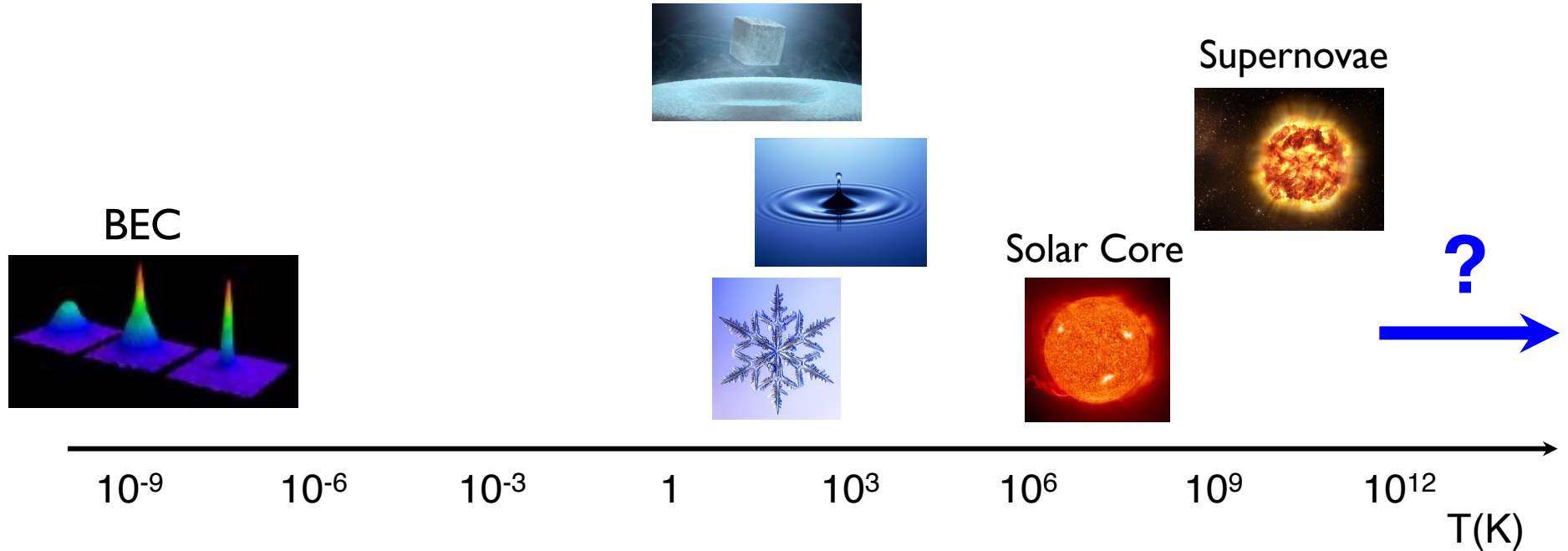
Phases of matter



Phases of matter

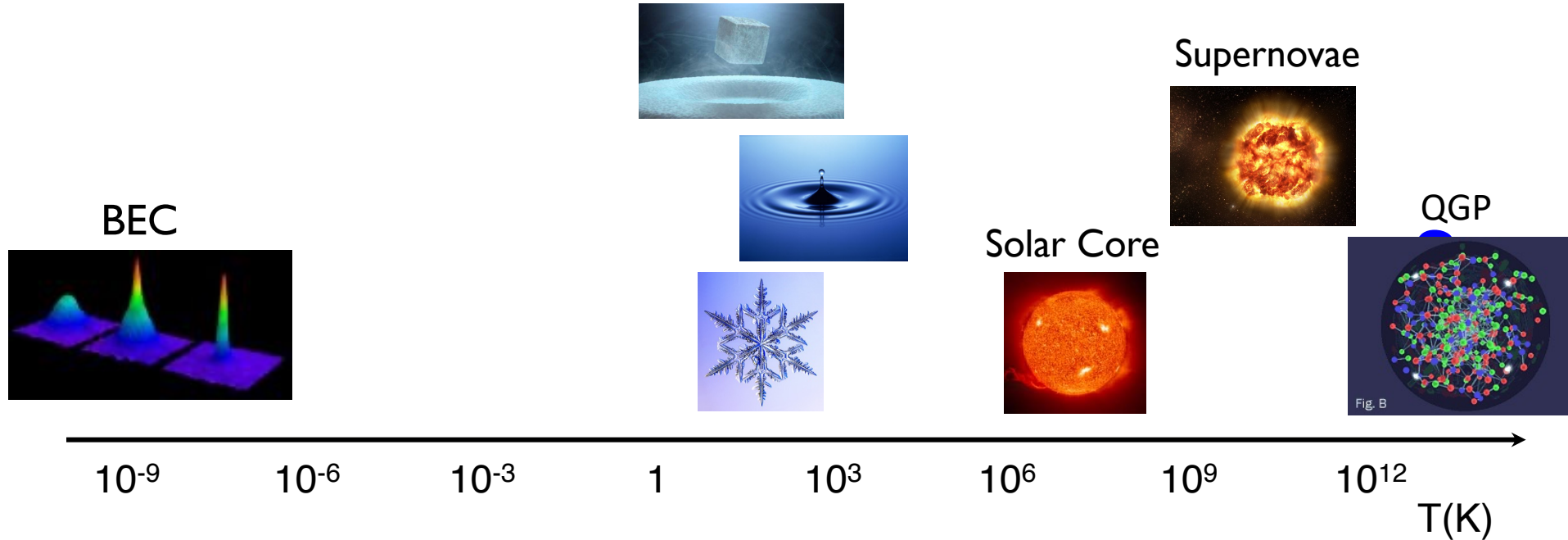


Phases of matter



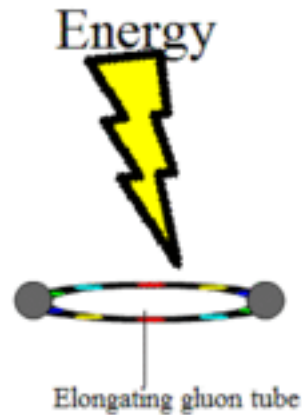
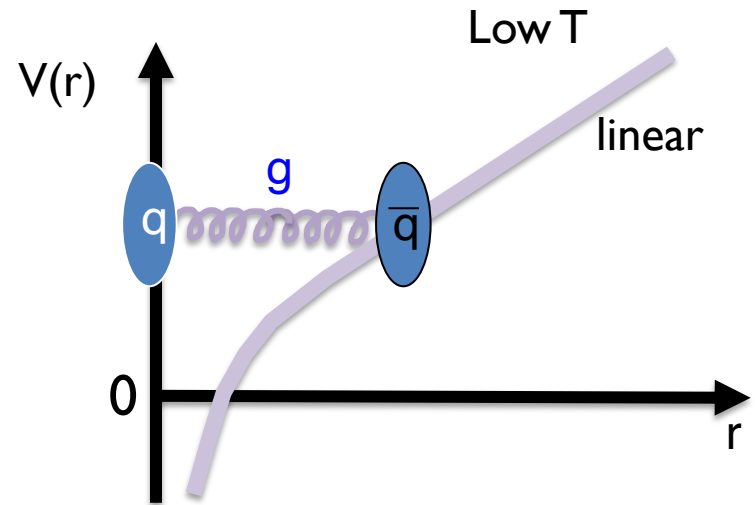
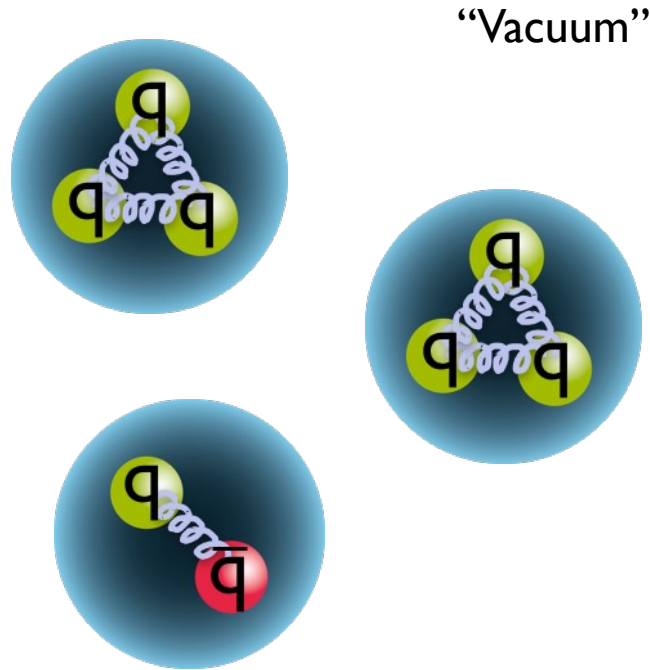
What are the properties of matter at extremely high temperatures ($> 10^{12}$ K)?

Phases of matter



What are the properties of matter at extremely high temperatures ($> 10^{12}$ K)?

Confinement \rightarrow De-confinement

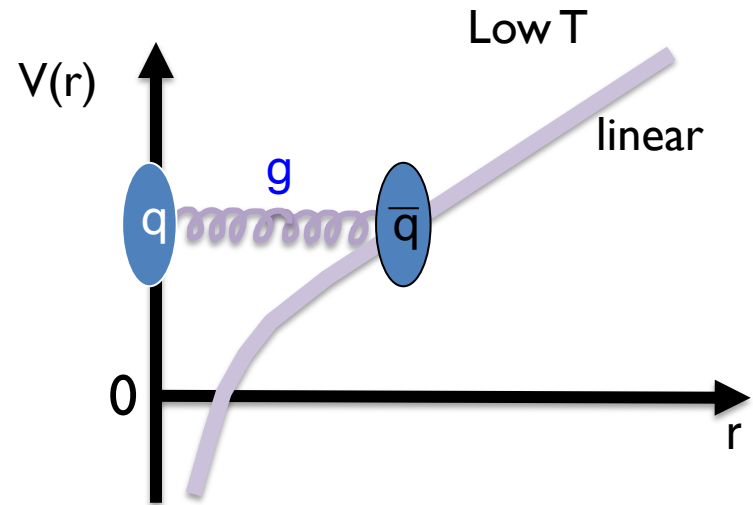


Confinement \rightarrow De-confinement

“Vacuum”



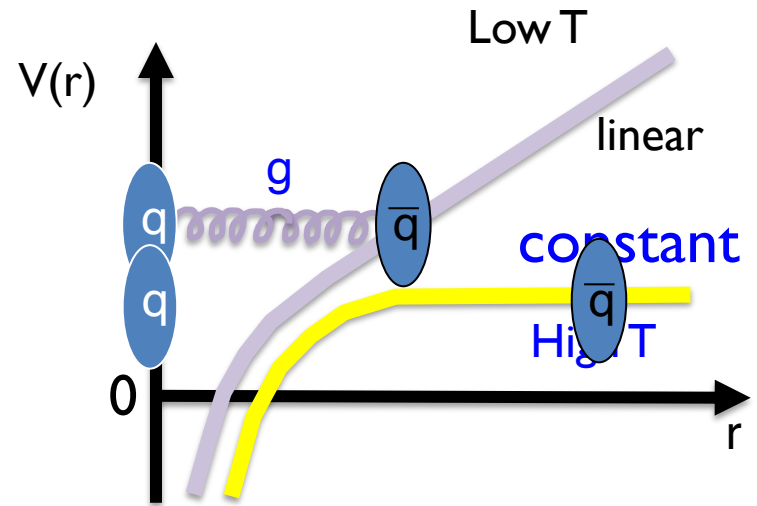
Taken from “Cooking Master Boy”



Confinement \rightarrow De-confinement



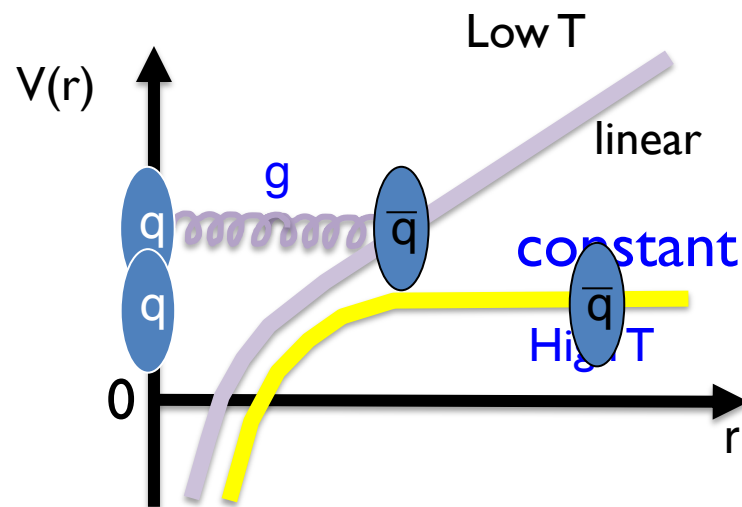
Taken from “Cooking Master Boy”



Confinement → De-confinement



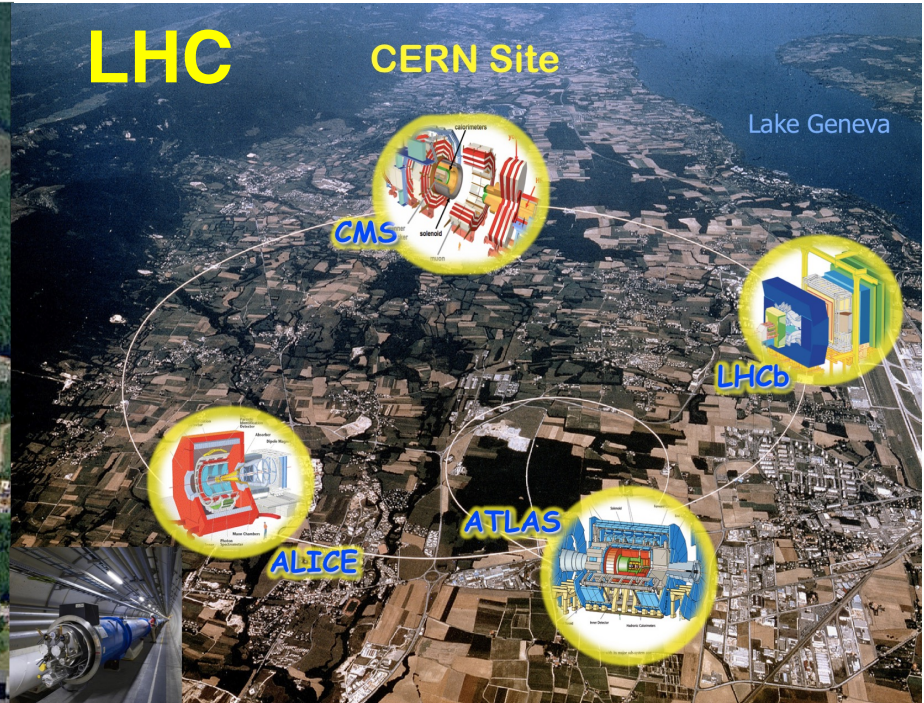
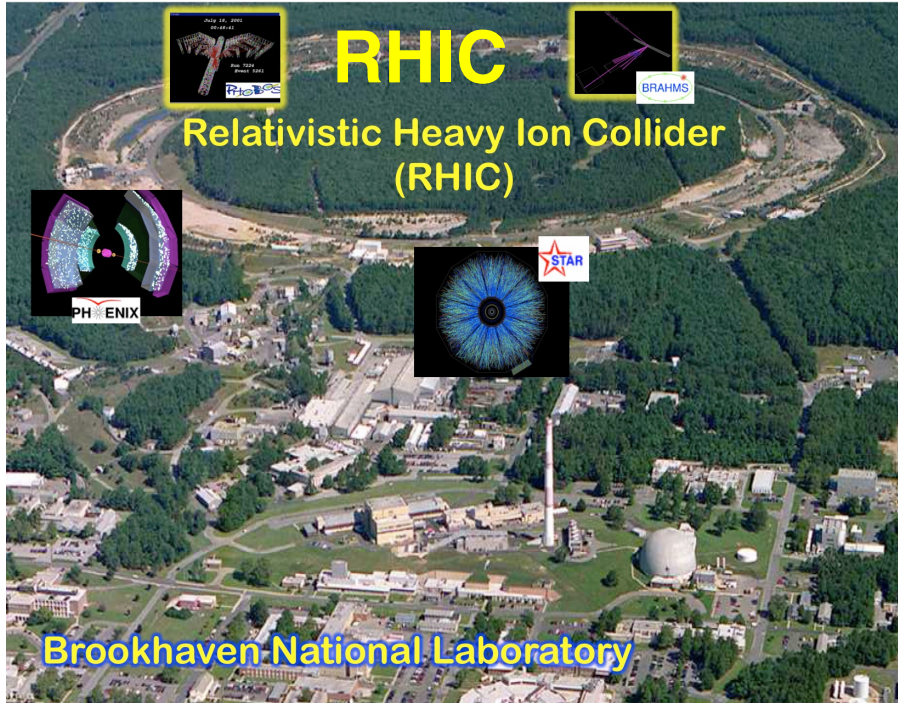
Taken from “Cooking Master Boy”



“Quark-Gluon Plasma”
(QGP)

$$T_c \sim 154 \text{ MeV} \sim 2 \times 10^{12} \text{ K}$$

Creating QGP in lab



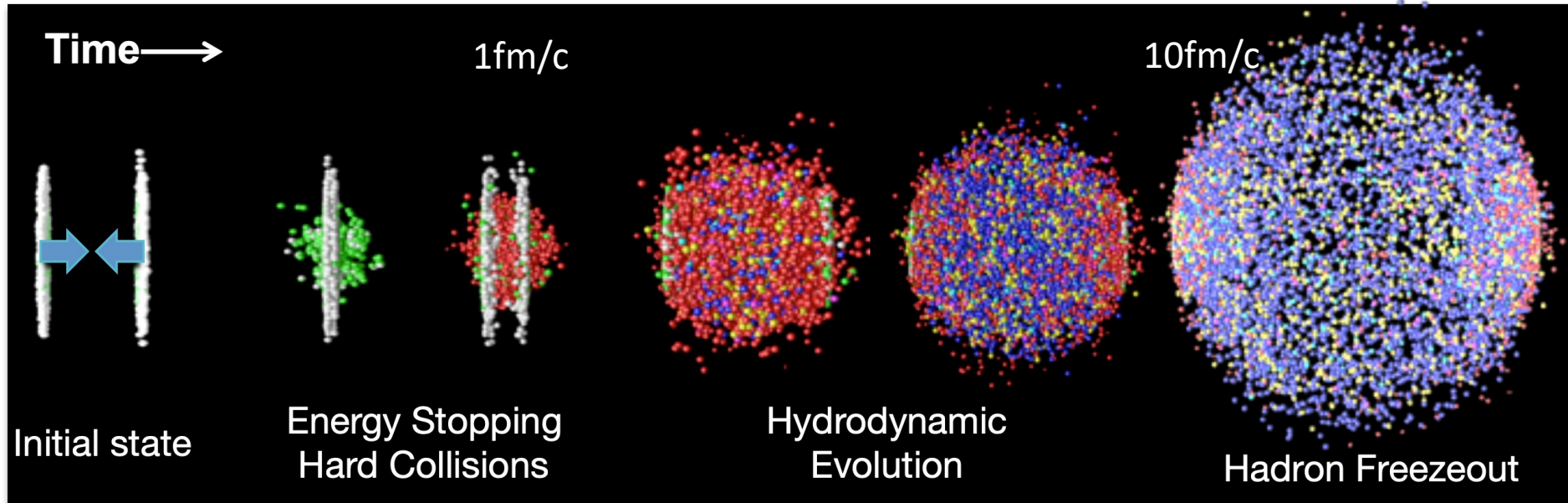
First collisions in 2000

- Diameter 1.2 km
- pp, dAu, CuCu, AuAu, UU, OO, ...
- $\sqrt{s_{NN}} \sim \mathbf{0.007 - 0.2 \text{ TeV}}$
- 99.995% speed of light

First collisions in 2010

- Diameter 8.6 km
- pp, PbPb, pPb, XeXe
- $\sqrt{s_{NN}} \sim \mathbf{5 - 8 \text{ TeV}}$
- 99.9999991% speed of light

QGP evolution

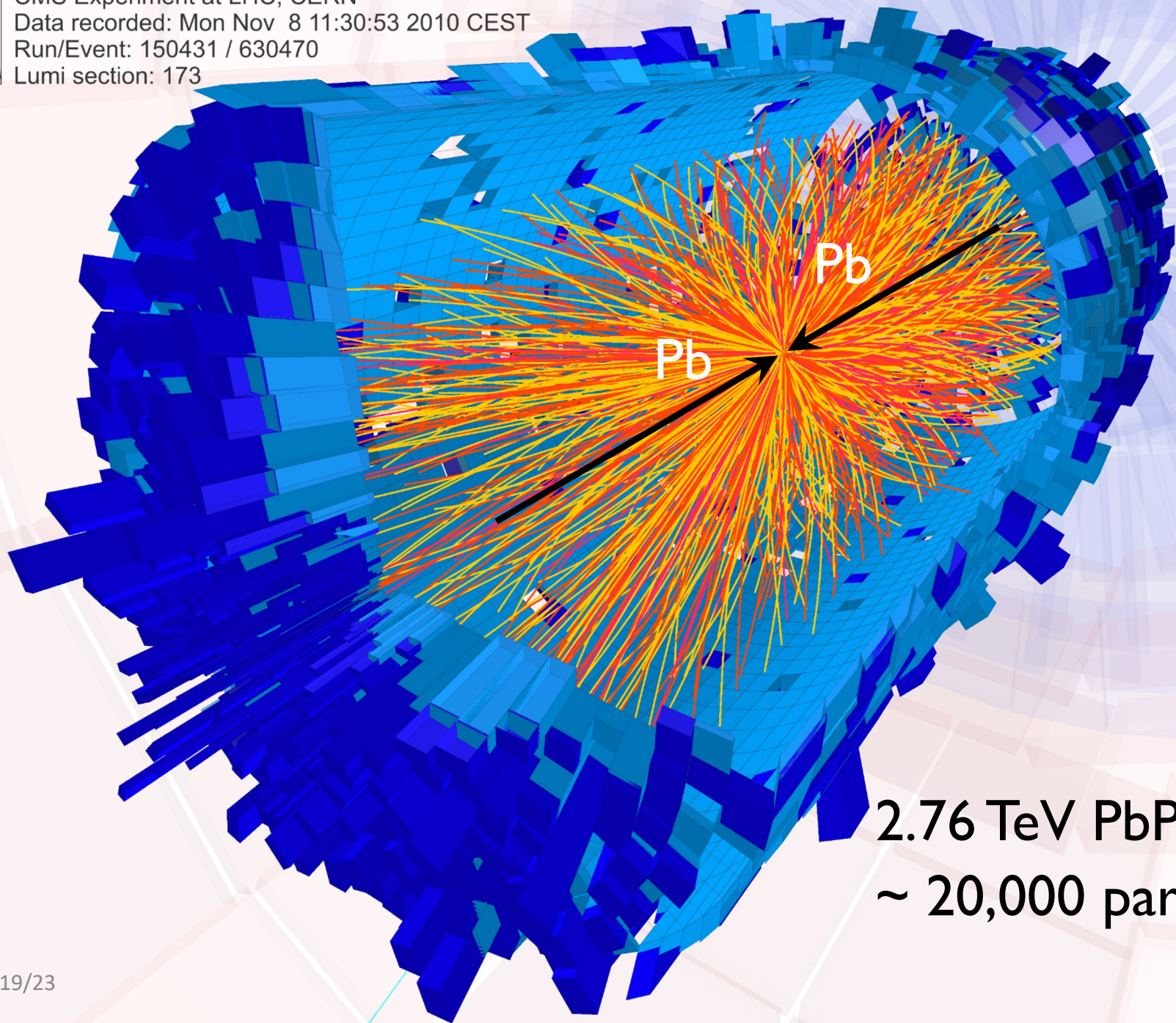


$1\text{ fm}/c \approx 3 \times 10^{-24}\text{ s}$

Extremely short-lived matter



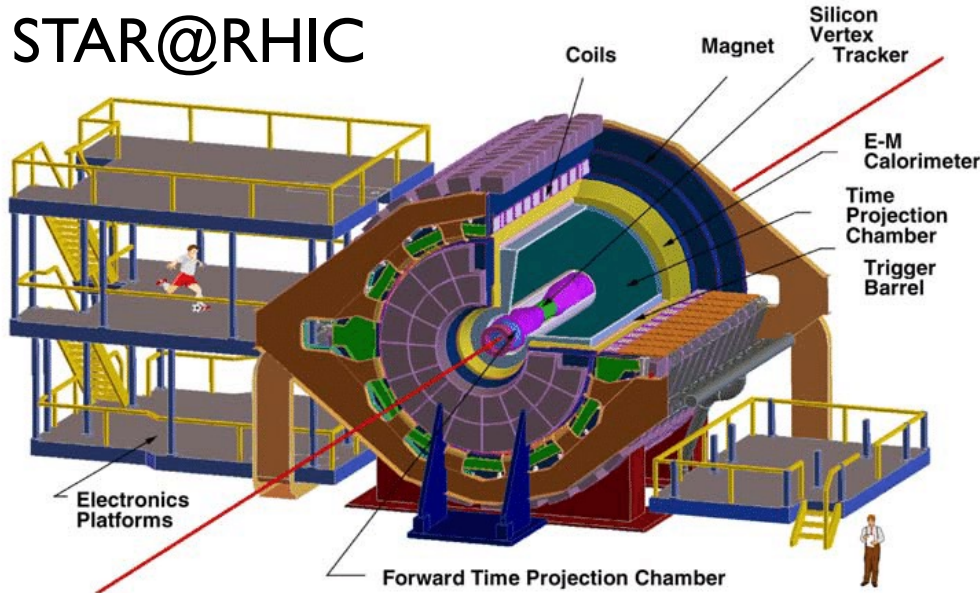
CMS Experiment at LHC, CERN
Data recorded: Mon Nov 8 11:30:53 2010 CEST
Run/Event: 150431 / 630470
Lumi section: 173



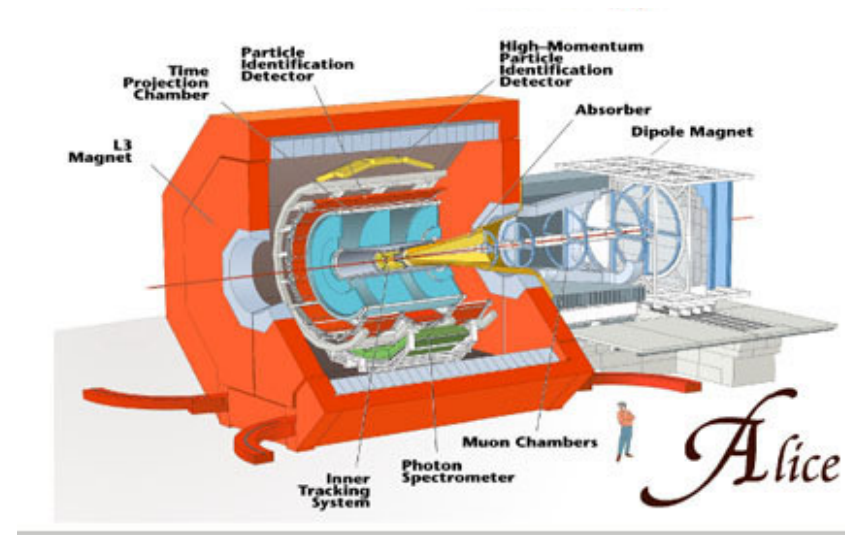
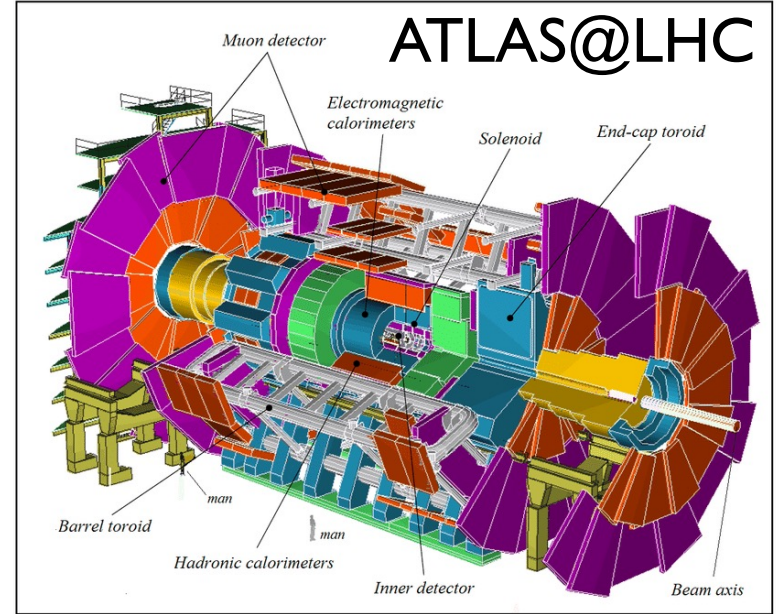
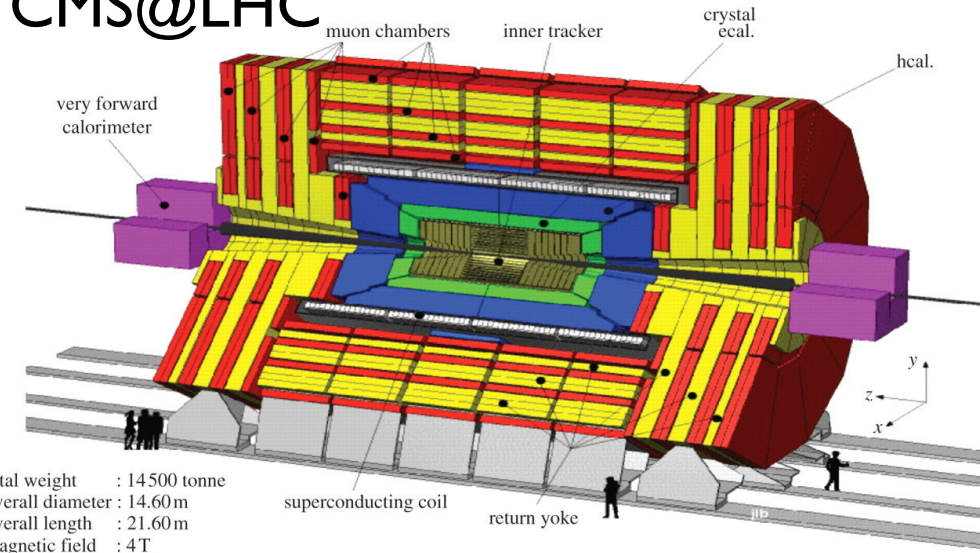
2.76 TeV PbPb
~ 20,000 particles

Detectors

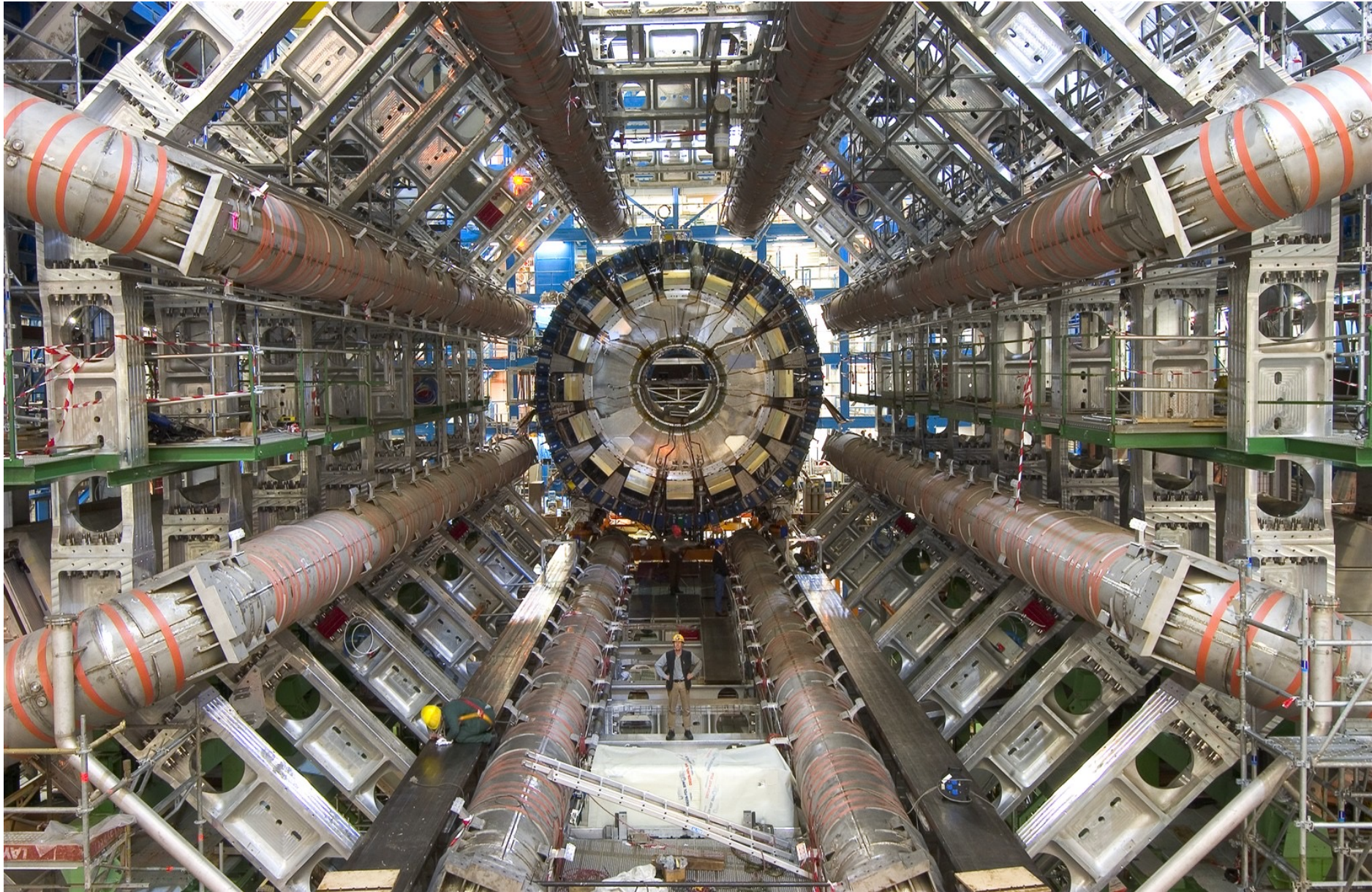
STAR@RHIC



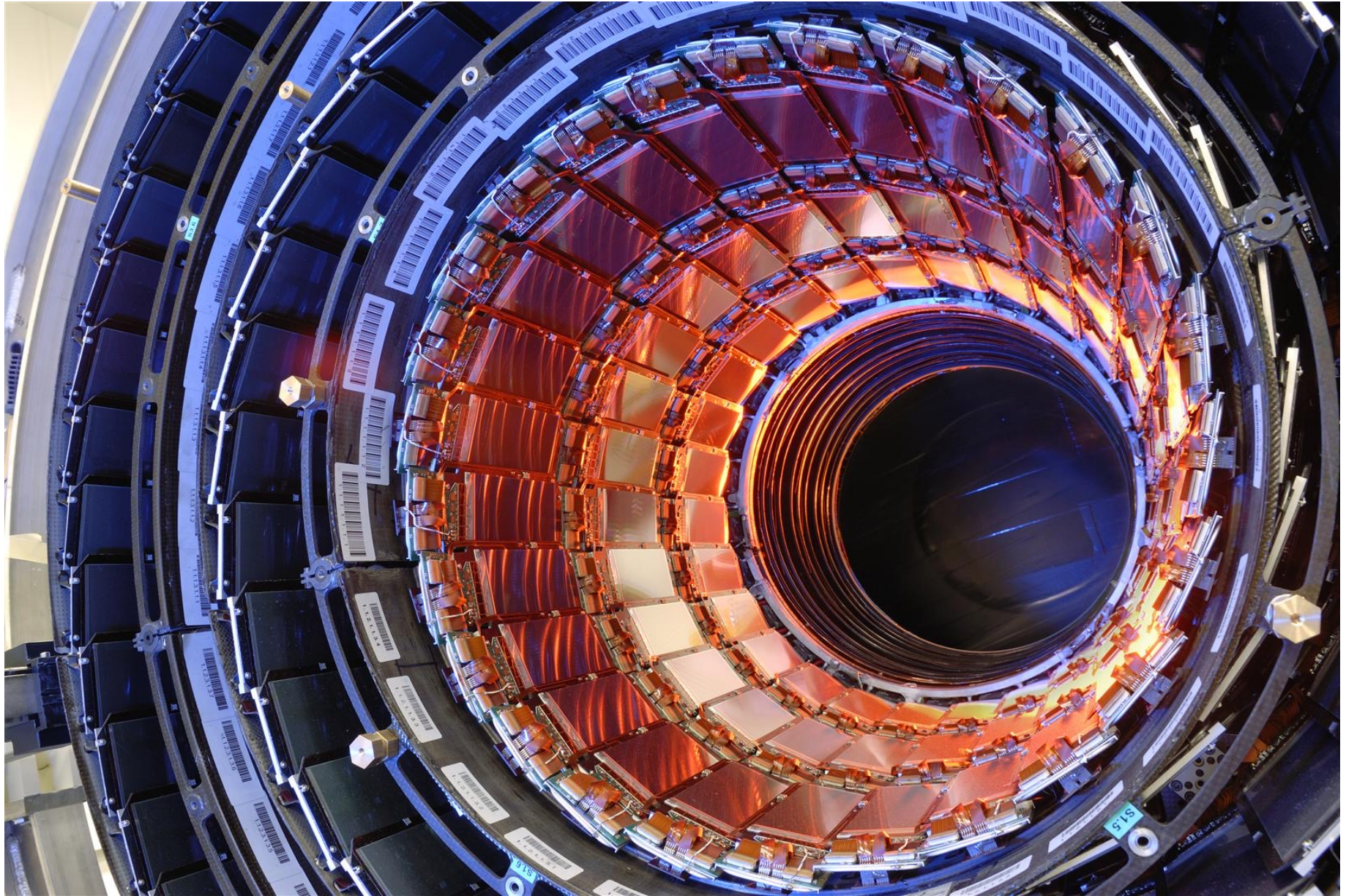
CMS@LHC



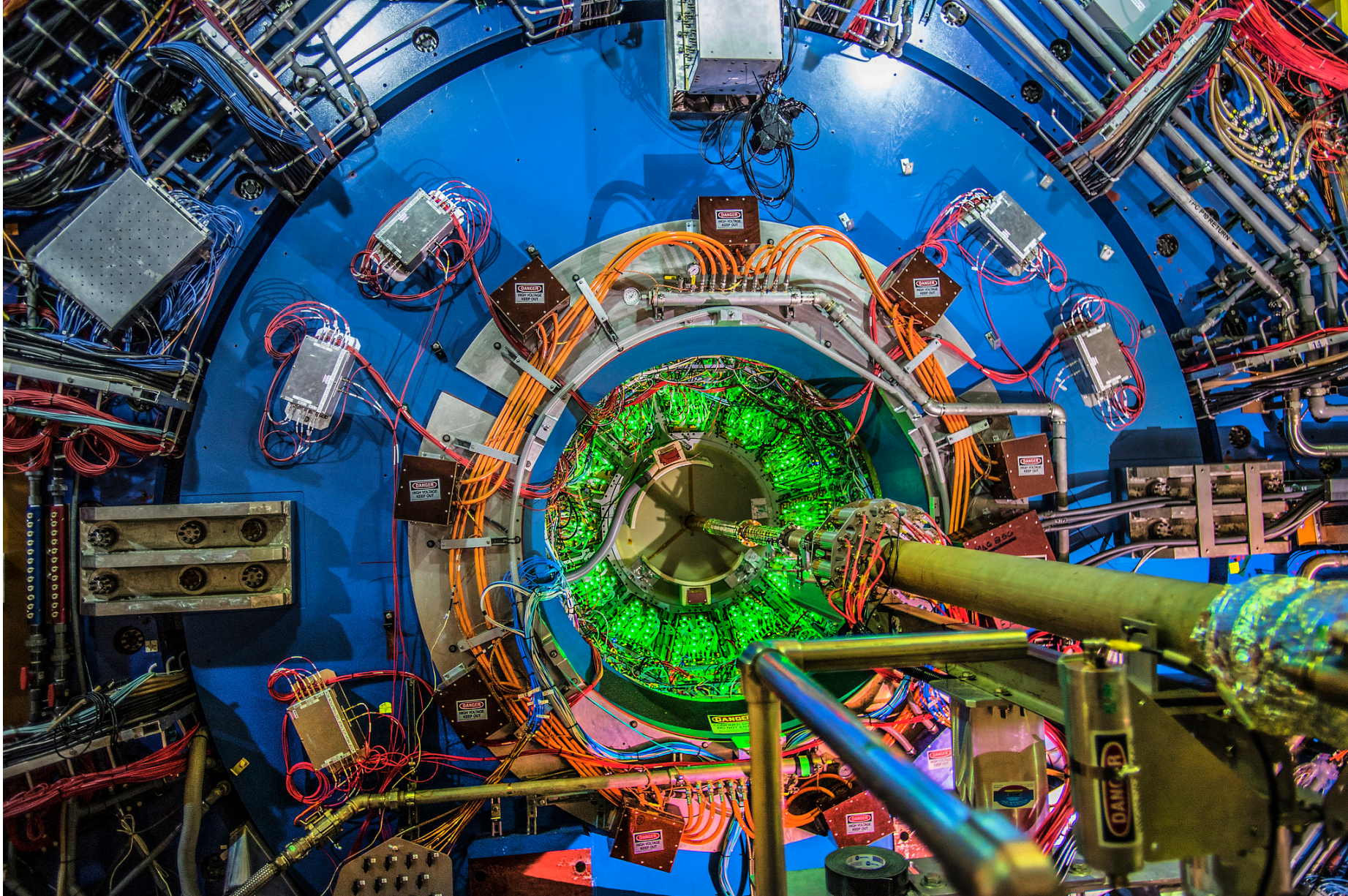
Detectors



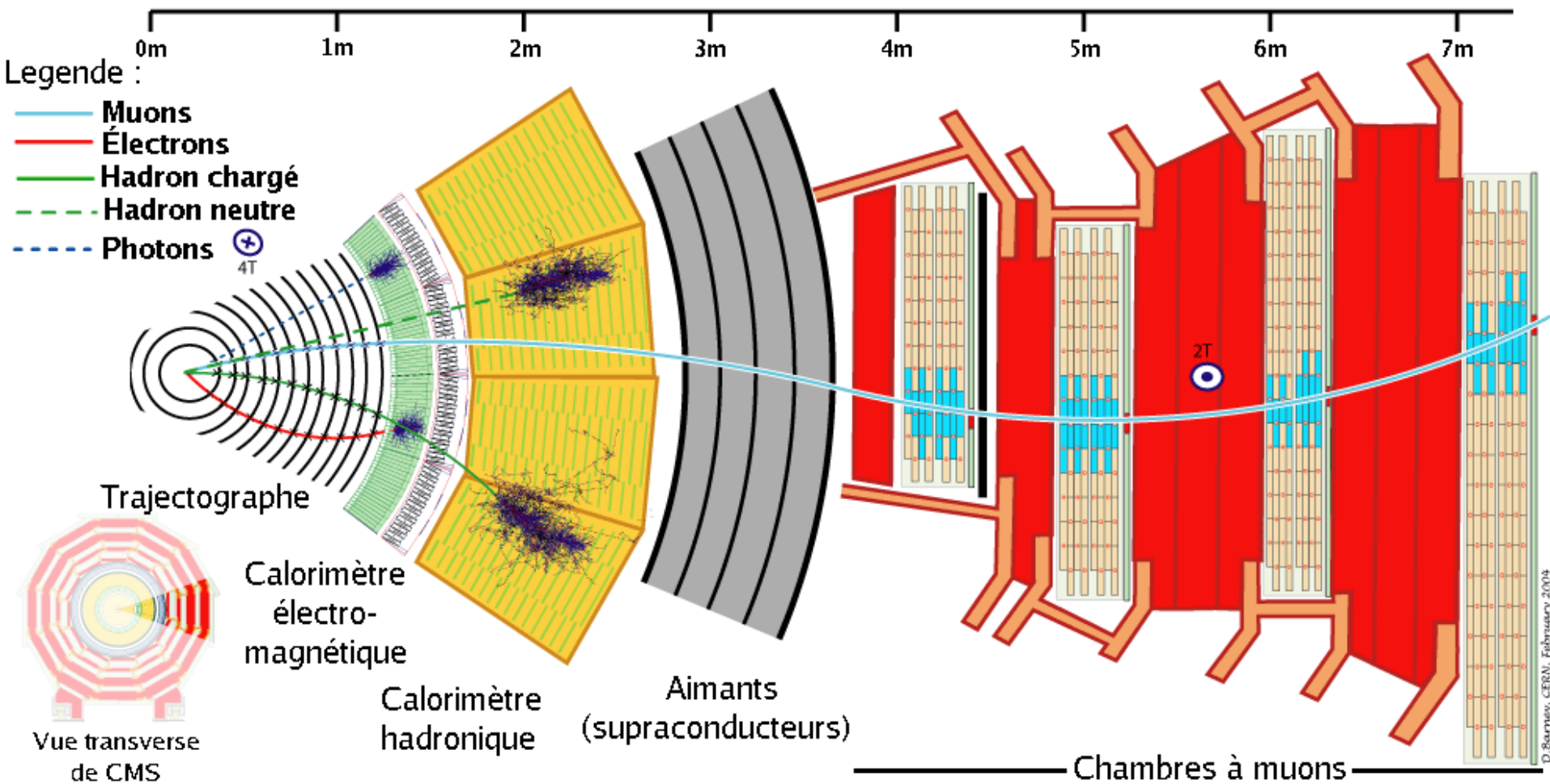
Detectors



Detectors



Particle journey



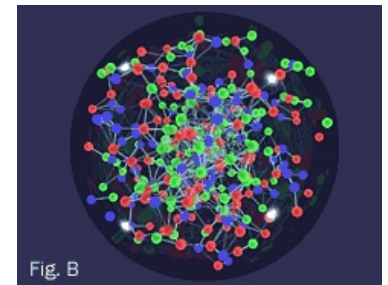
Discovery of a strongly coupled QGP

RHIC Scientists Serve Up "Perfect" Liquid (2005)

New state of matter more remarkable than predicted -- raising many new questions

Monday, April 18, 2005

TAMPA, FL -- The four detector groups conducting research at the [Relativistic Heavy Ion Collider](#) (RHIC) -- a giant atom "smasher" located at the U.S. Department of Energy's Brookhaven National Laboratory -- say they've created a new state of hot, dense matter out of the quarks and gluons that are the basic particles of atomic nuclei, but it is a state quite different and even more remarkable than had been predicted. In [peer-reviewed papers](#) summarizing the first three years of RHIC findings, the scientists say that instead of behaving like a gas of free quarks and gluons, as was expected, the matter created in RHIC's heavy ion collisions appears to be more like a *liquid*.



Discovery of a strongly coupled QGP

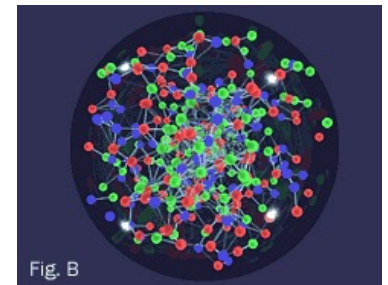
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Strongly coupled like a *liquid*, instead of gas



Discovery of a strongly coupled QGP

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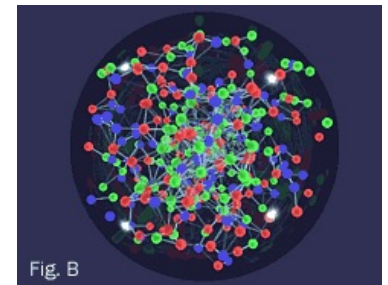
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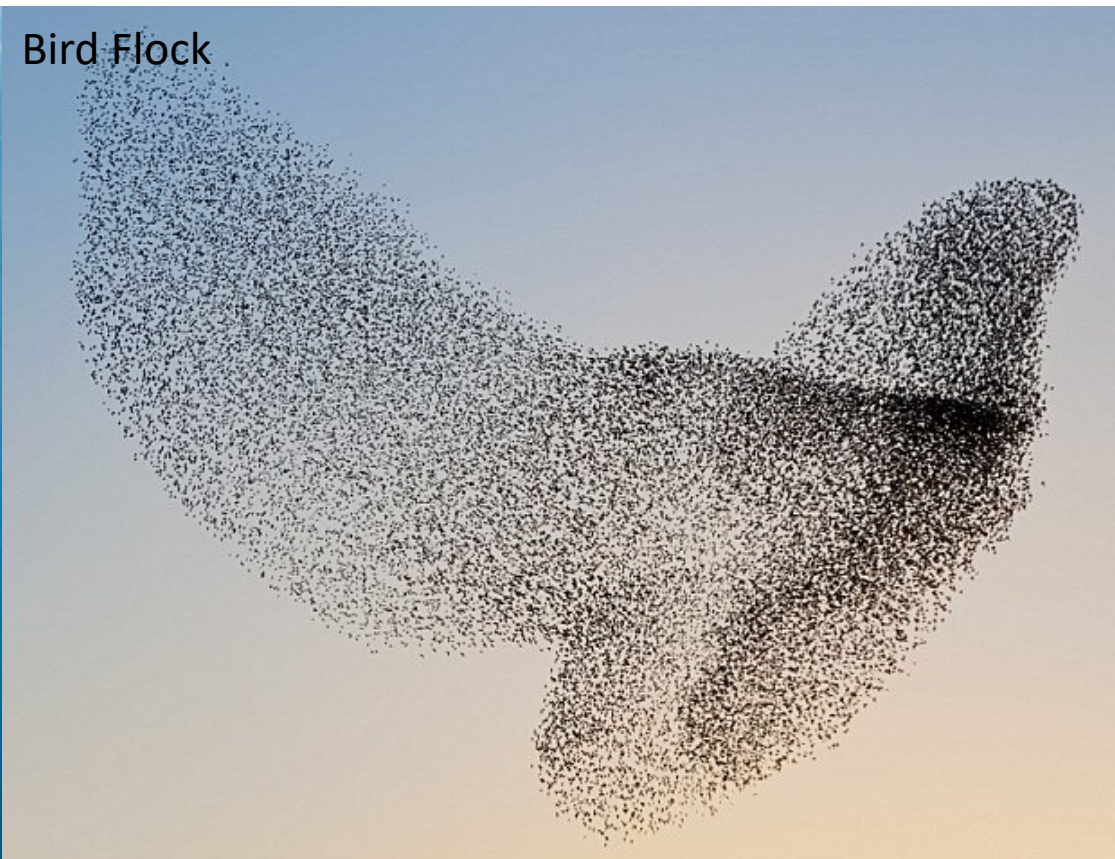
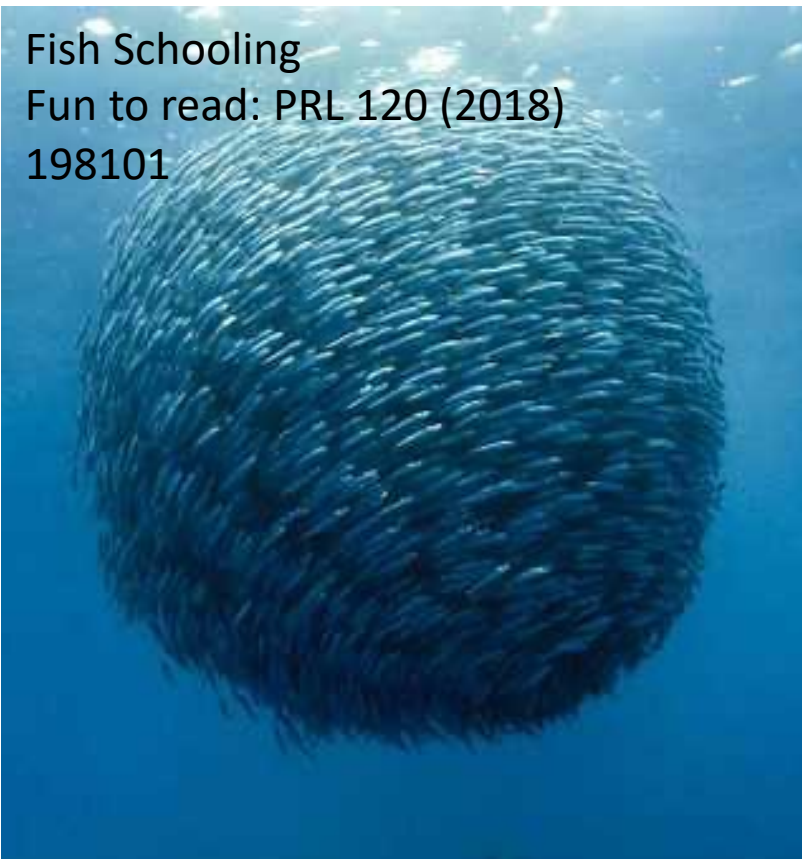
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Strongly coupled like a *liquid*, instead of gas

Important evidence from collectivity



Collectivity



Collectivity

Fish Schooling

Fun to read: PRL 120 (2018)

198101



Bird Flock



Tourists on the Great Wall during national holiday



Collectivity

Fish Schooling

Fun to read: PRL 120 (2018)

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Bird Flock



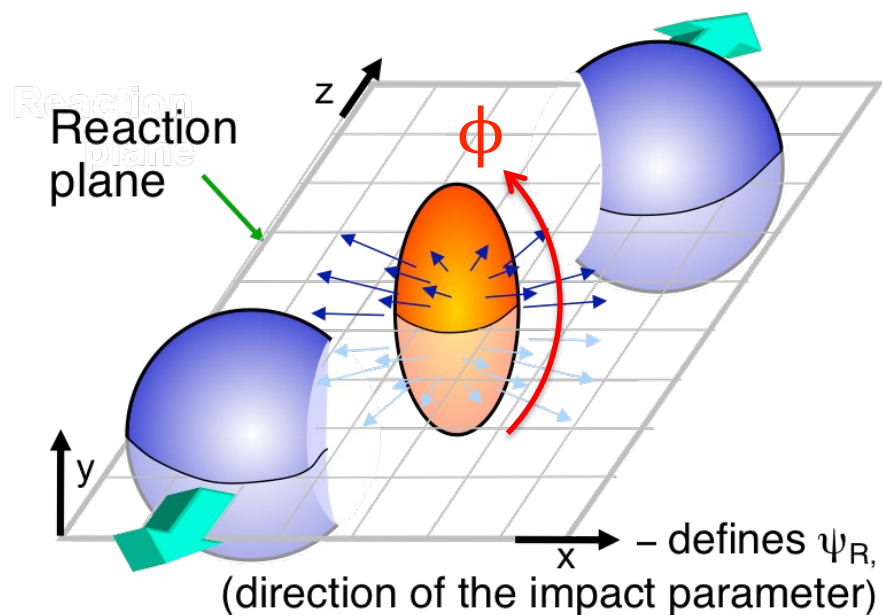
Tourists on the Great Wall during national holiday



Group of objects interact with each other frequently and move as a whole

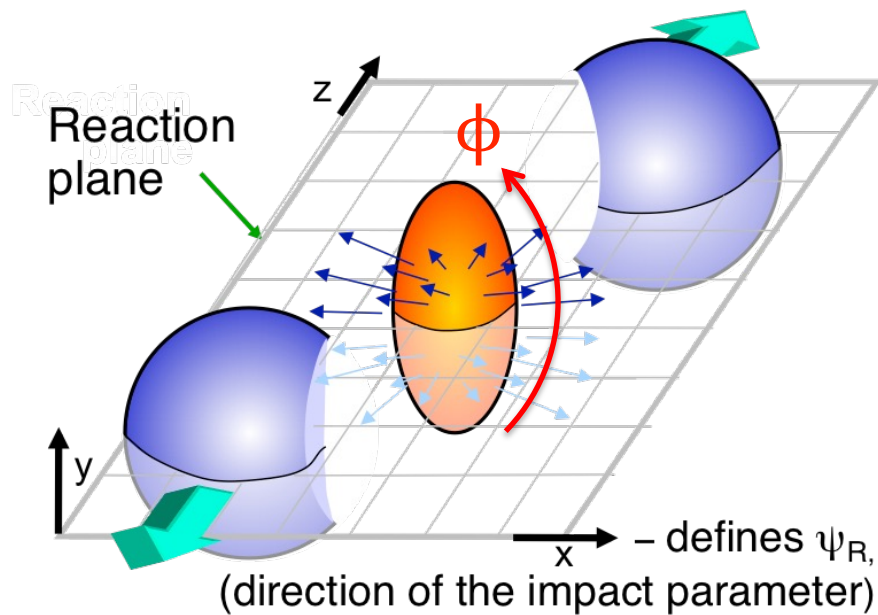
Collectivity – Evidence of fluidity

Off-center AA collision



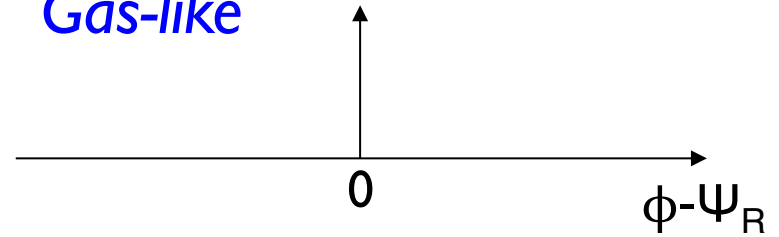
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Off-center AA collision



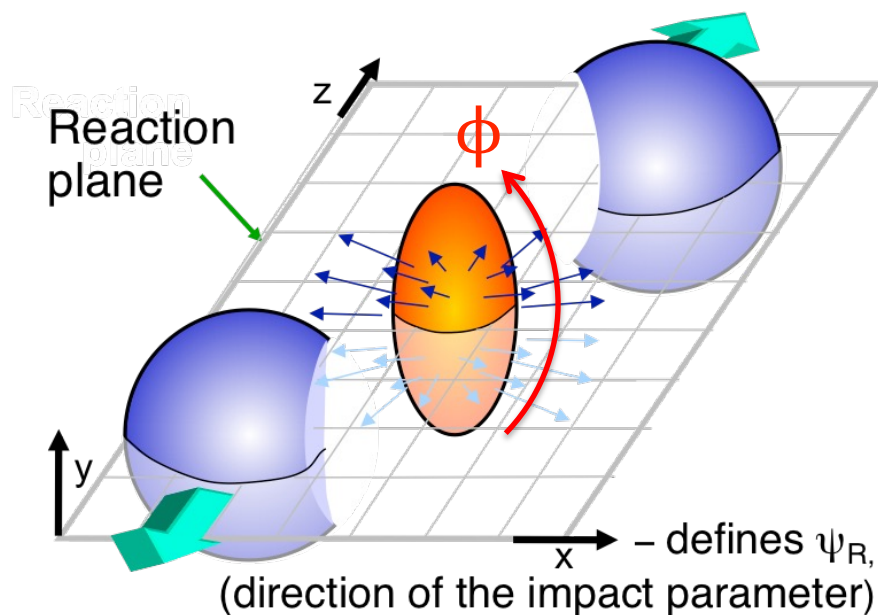
Momentum anisotropy

Gas-like

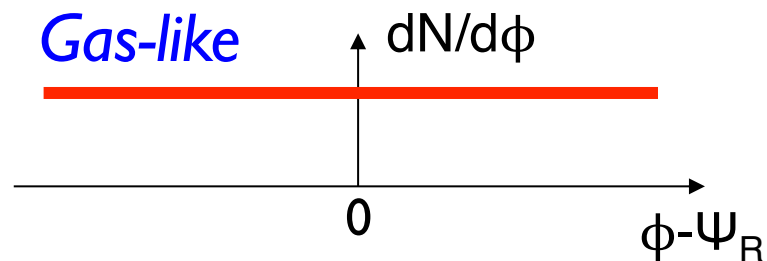


Collectivity – Evidence of fluidity

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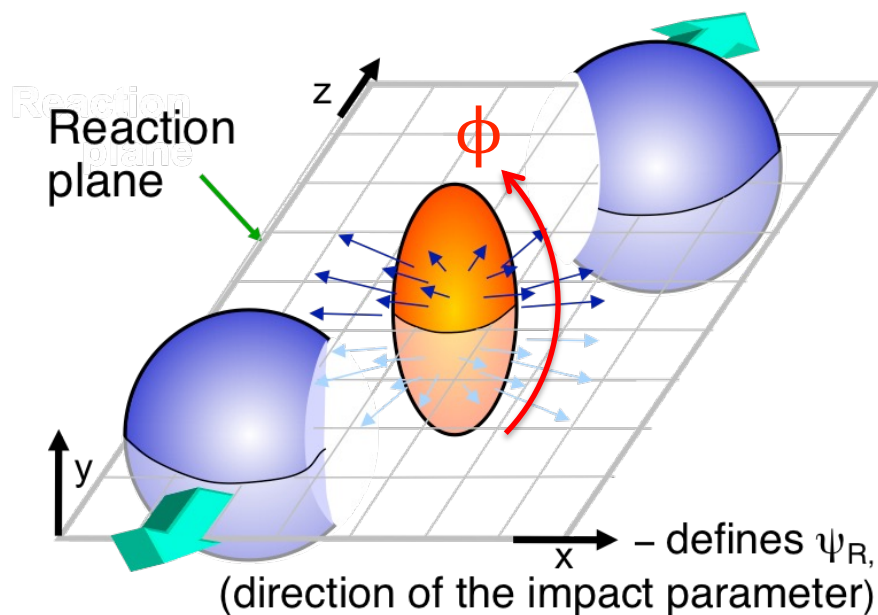


Momentum anisotropy



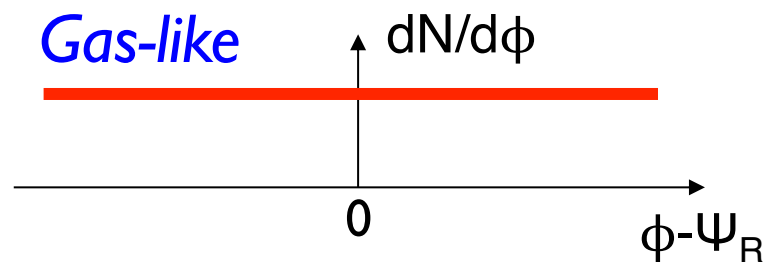
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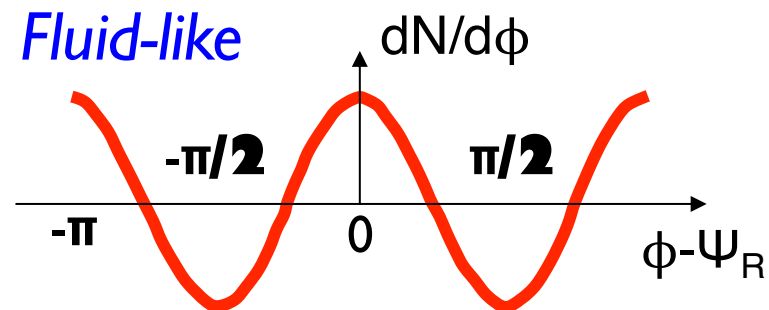


Momentum anisotropy

Gas-like



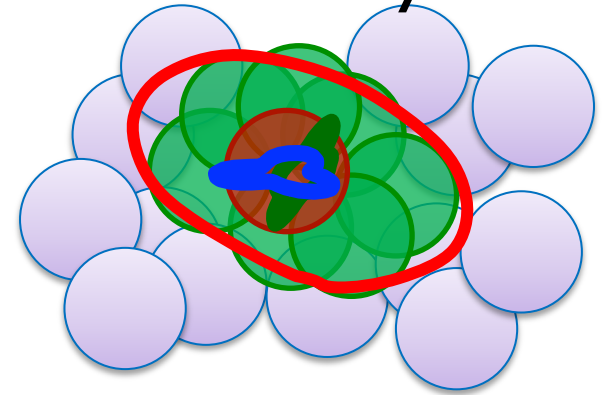
Fluid-like



Collectivity – Evidence of fluidity

$$\begin{array}{c} \text{Orange Oval} \\ \varepsilon_2 \\ \cos 2\Delta\phi \end{array} + \begin{array}{c} \text{Orange Triangle} \\ \varepsilon_3 \\ \cos 3\Delta\phi \end{array} + \begin{array}{c} \text{Orange Square} \\ \varepsilon_4 \\ \cos 4\Delta\phi \end{array} + \dots =$$

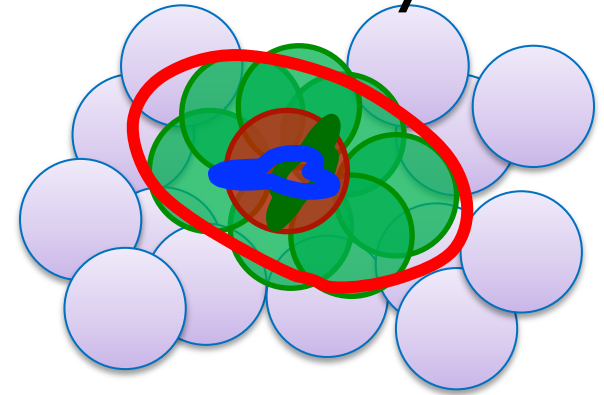
Initial Geometry



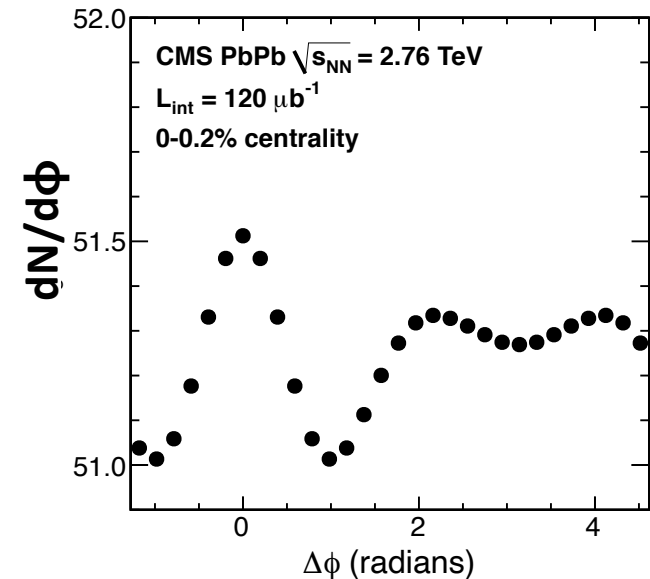
Collectivity – Evidence of fluidity

$$\begin{array}{c} \text{orange oval} \\ \varepsilon_2 \\ \cos 2\Delta\phi \end{array} + \begin{array}{c} \text{orange triangle} \\ \varepsilon_3 \\ \cos 3\Delta\phi \end{array} + \begin{array}{c} \text{orange rounded square} \\ \varepsilon_4 \\ \cos 4\Delta\phi \end{array} + \dots =$$

Initial Geometry



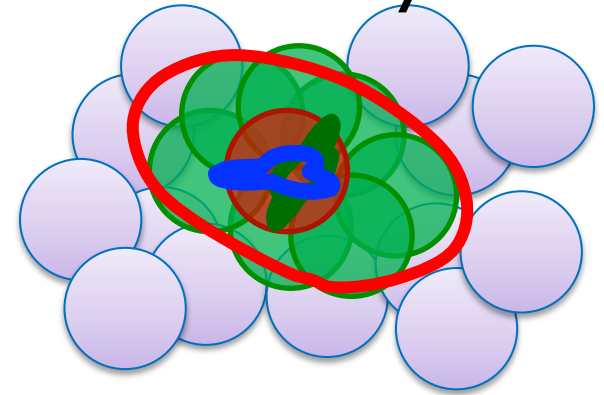
Particle distribution



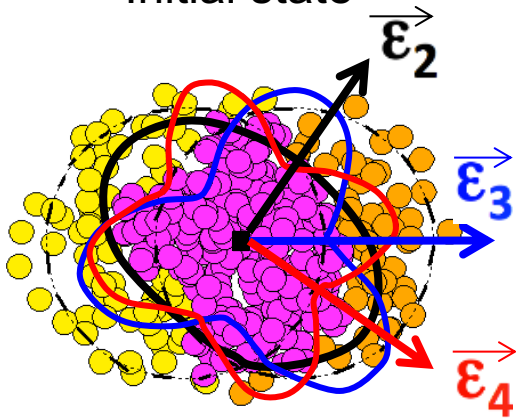
Collectivity – Initial to final

$$\begin{array}{c} \text{Oval} \\ \epsilon_2 \\ \cos 2\Delta\phi \end{array} + \begin{array}{c} \text{Triangle} \\ \epsilon_3 \\ \cos 3\Delta\phi \end{array} + \begin{array}{c} \text{Square} \\ \epsilon_4 \\ \cos 4\Delta\phi \end{array} + \dots =$$

Initial Geometry



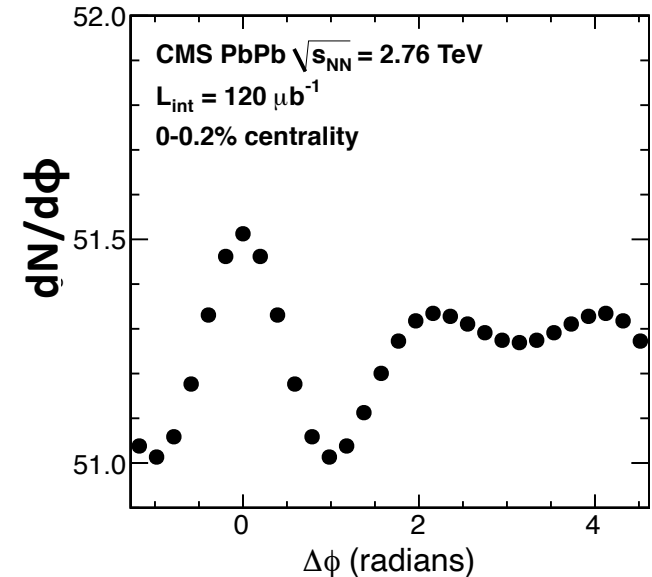
Initial state



Hydro-response

Space-time dynamics

Particle distribution



$$\vec{\epsilon}_n \equiv \epsilon_n e^{in\Phi_n^*} \equiv -\frac{\langle r^n e^{in\phi} \rangle}{\langle r^n \rangle}$$

Energy-momentum conservation

$$\partial_{\mu} T^{\mu\nu} = 0$$

Energy-momentum conservation

$$\partial_{\mu} T^{\mu\nu} = 0$$

Ideal hydro: System always in local equilibrium

$$T^{\mu\nu} = (\epsilon + P)u^{\mu}u^{\nu} - P g^{\mu\nu}$$

Hydro response controlled by QCD Equation of State $P(\epsilon)$

Energy-momentum conservation

$$\partial_{\mu} T^{\mu\nu} = 0$$

Viscous hydro: Including near-equilibrium corrections

$$T^{\mu\nu} = \epsilon u^{\mu} u^{\nu} - (P + \Pi) \Delta^{\mu\nu} + \pi^{\mu\nu}$$

← Bulk Pressure

← Shear Tensor

Energy-momentum conservation

$$\partial_\mu T^{\mu\nu} = 0$$

Viscous hydro: Including near-equilibrium corrections

$$T^{\mu\nu} = \epsilon u^\mu u^\nu - (P + \Pi)\Delta^{\mu\nu} + \pi^{\mu\nu}$$

Bulk Pressure Shear Tensor

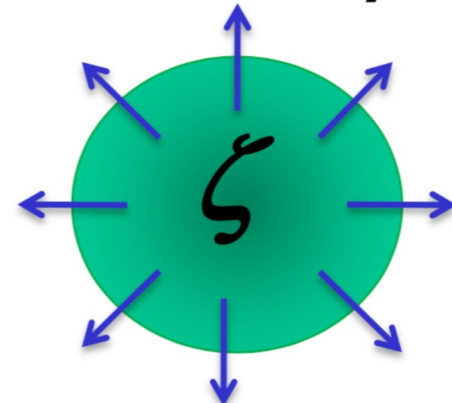
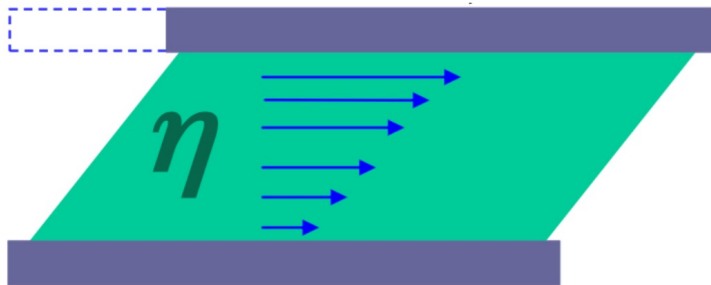
With 1st order gradient expansion

$$\pi^{\mu\nu} = -\eta\sigma^{\mu\nu}$$

η : Shear viscosity

$$\Pi = -\zeta\nabla\frac{1}{\lambda}u^\lambda$$

ζ : Bulk viscosity



Energy-momentum conservation

$$\partial_\mu T^{\mu\nu} = 0$$

Viscous hydro: Including near-equilibrium corrections

$$T^{\mu\nu} = \epsilon u^\mu u^\nu - (P + \Pi) \Delta^{\mu\nu} + \pi^{\mu\nu}$$

Bulk Pressure Shear Tensor

With 2nd order gradient expansion

$$\begin{aligned} \pi^{\mu\nu} = & -\eta \sigma^{\mu\nu} + \eta \tau_\pi \left[\langle D \sigma^{\mu\nu} \rangle + \frac{\nabla_\lambda^\perp u^\lambda}{3} \sigma^{\mu\nu} \right] + \kappa \left[R^{\langle \mu\nu \rangle} - 2u_\lambda u_\rho R^{\lambda \langle \mu\nu \rangle \rho} \right] + \lambda_1 \sigma^{\langle \mu}_\lambda \sigma^{\nu \rangle \lambda} \\ & + \lambda_2 \sigma^{\langle \mu}_\lambda \Omega^{\nu \rangle \lambda} + \lambda_3 \Omega^{\langle \mu}_\lambda \Omega^{\nu \rangle \lambda} + \kappa^* 2u_\lambda u_\rho R^{\lambda \langle \mu\nu \rangle \rho} + \eta \tau_\pi^* \frac{\nabla_\lambda^\perp u^\lambda}{3} \sigma^{\mu\nu} + \bar{\lambda}_4 \nabla_\perp^{\langle \mu} \ln \epsilon \nabla_\perp^{\nu \rangle} \ln \epsilon \\ \Pi = & -\zeta (\nabla_\lambda^\perp u^\lambda) + \zeta \tau_\Pi D (\nabla_\lambda^\perp u^\lambda) + \xi_1 \sigma^{\mu\nu} \sigma_{\mu\nu} + \xi_2 (\nabla_\lambda^\perp u^\lambda)^2 \\ & + \xi_3 \Omega^{\mu\nu} \Omega_{\mu\nu} + \bar{\xi}_4 \nabla_\mu^\perp \ln \epsilon \nabla_\perp^\mu \ln \epsilon + \xi_5 R + \xi_6 u^\lambda u^\rho R_{\lambda\rho}. \end{aligned}$$

Energy-momentum conservation

$$\partial_\mu T^{\mu\nu} = 0$$

Viscous hydro: Including near-equilibrium corrections

$$T^{\mu\nu} = \epsilon u^\mu u^\nu - (P + \Pi) \Delta^{\mu\nu} + \pi^{\mu\nu}$$

← Bulk Pressure ← Shear Tensor

With 2nd order gradient expansion

$$\begin{aligned} \pi^{\mu\nu} = & -\eta \sigma^{\mu\nu} + \eta \tau_\pi \left[\langle D \sigma^{\mu\nu} \rangle + \frac{\nabla_\lambda^\perp u^\lambda}{3} \sigma^{\mu\nu} \right] + \kappa \left[R^{\langle \mu\nu \rangle} - 2u_\lambda u_\rho R^{\lambda \langle \mu\nu \rangle \rho} \right] + \lambda_1 \sigma^{\langle \mu}_\lambda \sigma^{\nu \rangle \lambda} \\ & + \lambda_2 \sigma^{\langle \mu}_\lambda \Omega^{\nu \rangle \lambda} + \lambda_3 \Omega^{\langle \mu}_\lambda \Omega^{\nu \rangle \lambda} + \kappa^* 2u_\lambda u_\rho R^{\lambda \langle \mu\nu \rangle \rho} + \eta \tau_\pi^* \frac{\nabla_\lambda^\perp u^\lambda}{3} \sigma^{\mu\nu} + \bar{\lambda}_4 \nabla_\perp^{\langle \mu} \ln \epsilon \nabla_\perp^{\nu \rangle} \ln \epsilon \\ \Pi = & -\zeta (\nabla_\lambda^\perp u^\lambda) + \zeta \tau_\Pi D (\nabla_\lambda^\perp u^\lambda) + \xi_1 \sigma^{\mu\nu} \sigma_{\mu\nu} + \xi_2 (\nabla_\lambda^\perp u^\lambda)^2 \\ & + \xi_3 \Omega^{\mu\nu} \Omega_{\mu\nu} + \bar{\xi}_4 \nabla_\mu^\perp \ln \epsilon \nabla_\perp^\mu \ln \epsilon + \xi_5 R + \xi_6 u^\lambda u^\rho R_{\lambda\rho}. \end{aligned}$$

ab.initio calc. for QGP not easy, relies on **model/data** comparison

Collectivity – how to measure

- Single particle distribution

$$\frac{dN}{d\phi} = \frac{N}{2\pi} \left[1 + 2 \sum_n v_n \cos n(\phi - \Phi_n) \right]$$

Event Plane

Anisotropic Flow

Collectivity – how to measure

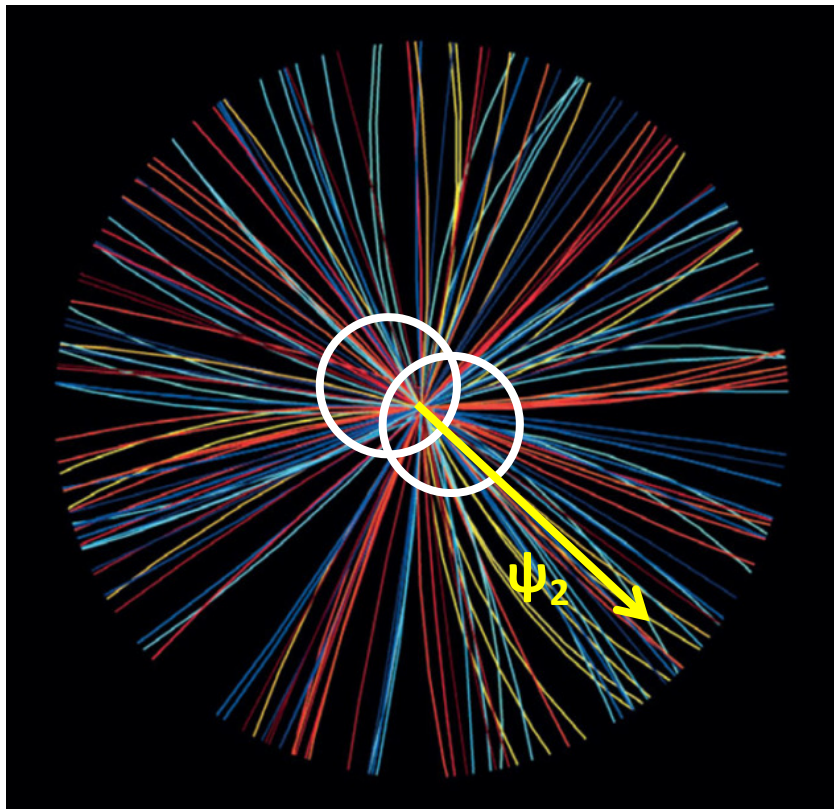
- Single particle distribution

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Event Plane

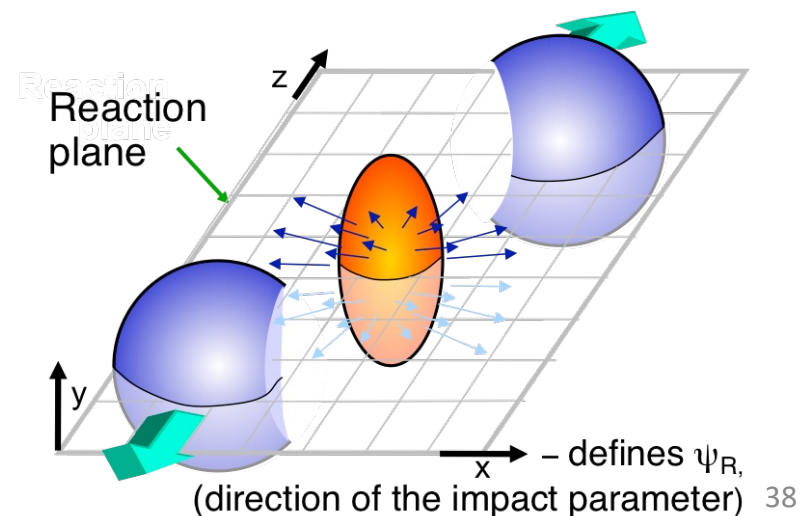
Anisotropic Flow

- Event Plane method



Event Plane from emitted particles

$$\psi_n = \frac{1}{n} \tan^{-1} \frac{\sum_i \sin(n\phi_i)}{\sum_i \cos(n\phi_i)}$$



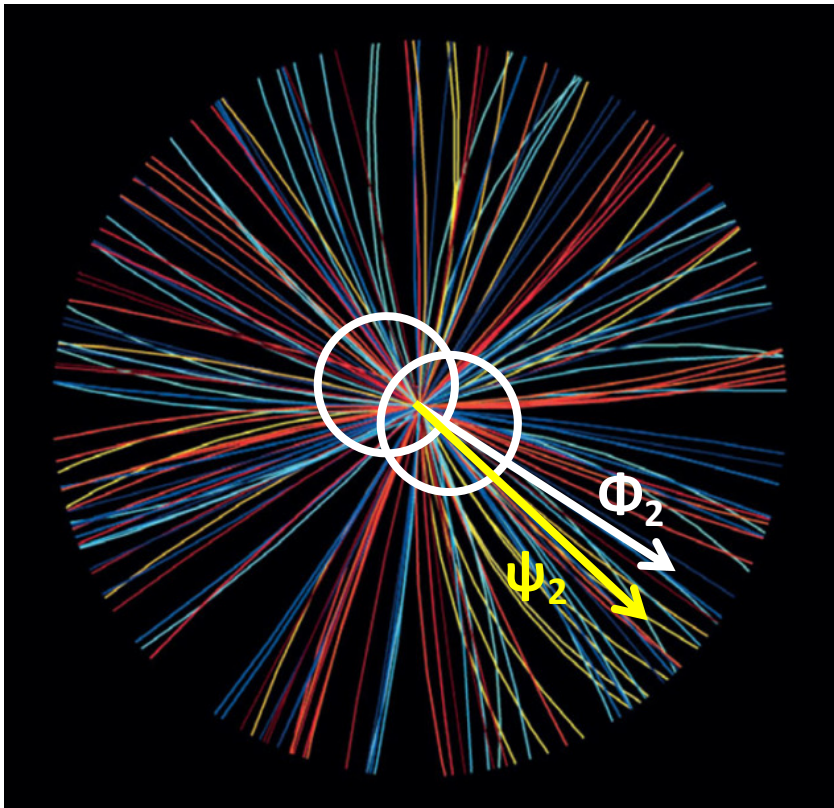
Collectivity – how to measure

- Single particle distribution

$$\frac{dN}{d\phi} = \frac{N}{2\pi} \left[1 + 2 \sum_n v_n \cos n(\phi - \Phi_n) \right]$$

Event Plane
↙
↘
Anisotropic Flow

- Event Plane method



Event Plane from emitted particles

$$\psi_n = \frac{1}{n} \tan^{-1} \frac{\sum_i \sin(n\phi_i)}{\sum_i \cos(n\phi_i)}$$

Event Plane Resolution

$$\mathcal{R}_n \propto \cos n(\Phi_n - \psi_n)$$

$$v_n = v_n^{obs} / \mathcal{R}_n$$

\mathcal{R}_n cannot be perfectly determined!

Collectivity – how to measure

- Particle pair distribution

$$\frac{dN^{pair}}{d\Delta\phi} = \frac{N^{pair}}{2\pi} \left[1 + 2 \sum_n V_{n\Delta,ab} \cos n(\phi_a - \phi_b) \right]$$

$V_{n\Delta,ab} = v_{n,a}v_{n,b}$ when $\Phi_n^a = \Phi_n^b$

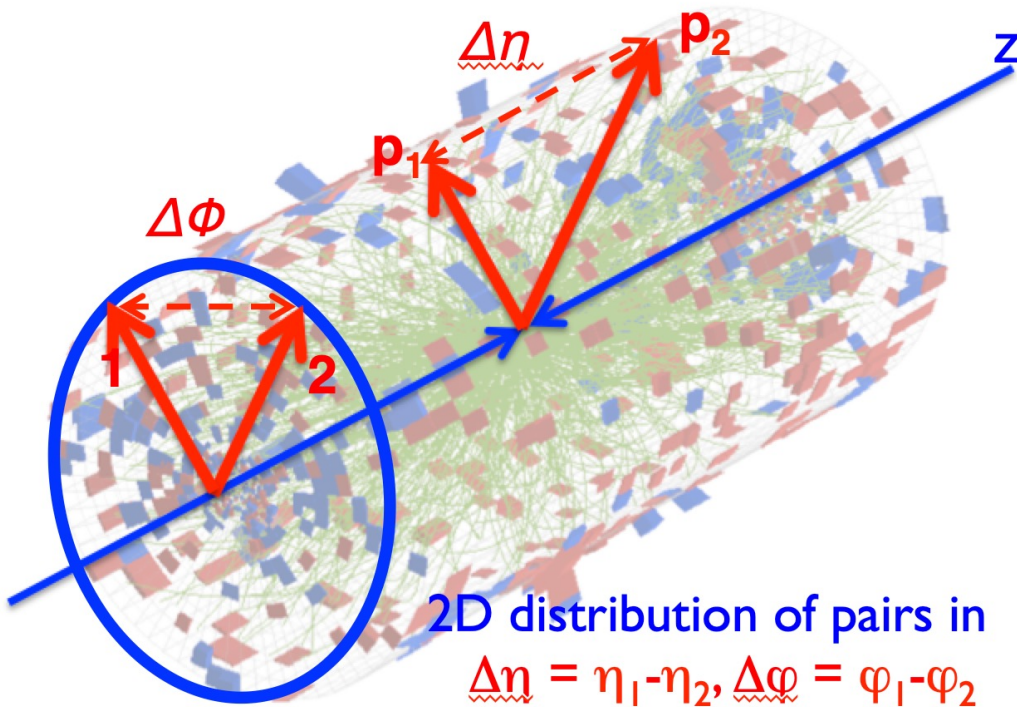
Collectivity – how to measure

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$V_{n\Delta,ab} = v_{n,a}v_{n,b}$ when $\Phi_n^a = \Phi_n^b$

- Two particle correlation method



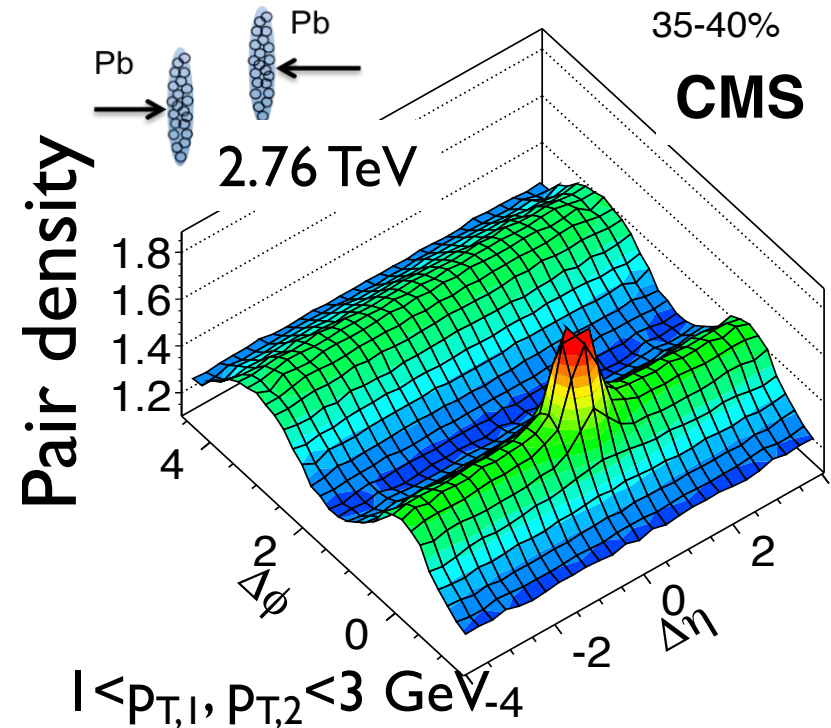
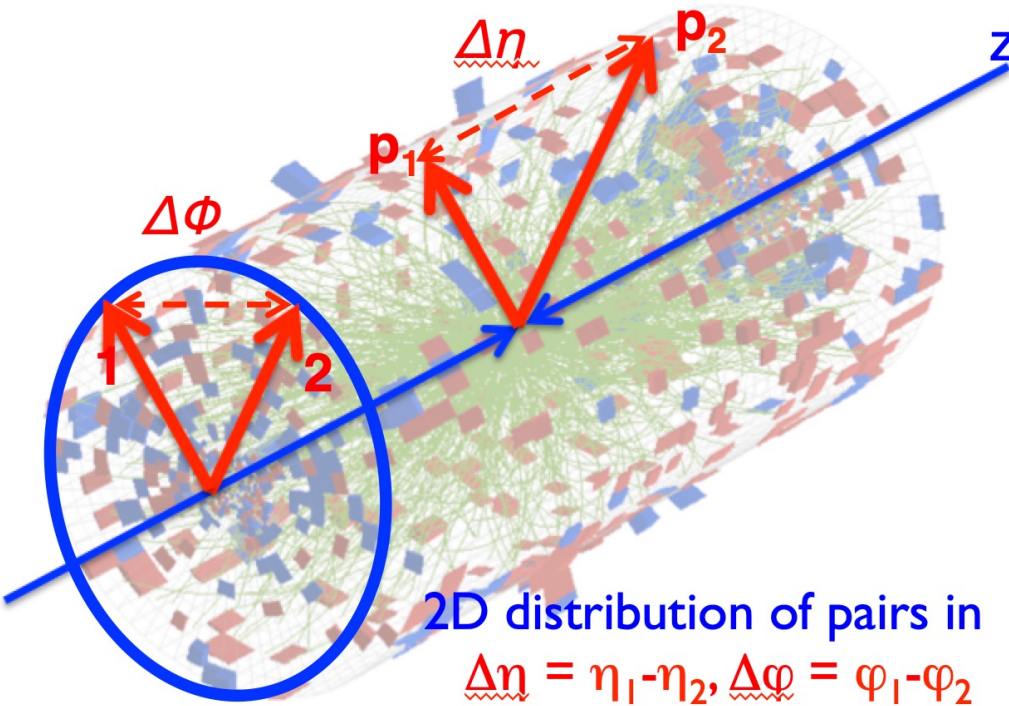
Collectivity – how to measure

- Particle pair distribution

$$\frac{dN^{pair}}{d\Delta\phi} = \frac{N^{pair}}{2\pi} \left[1 + 2 \sum_n V_{n\Delta,ab} \cos n(\phi_a - \phi_b) \right]$$

$$V_{n\Delta,ab} = v_{n,a} v_{n,b} \quad \text{when } \Phi_n^a = \Phi_n^b$$

- Two particle correlation method



Collectivity – how to measure

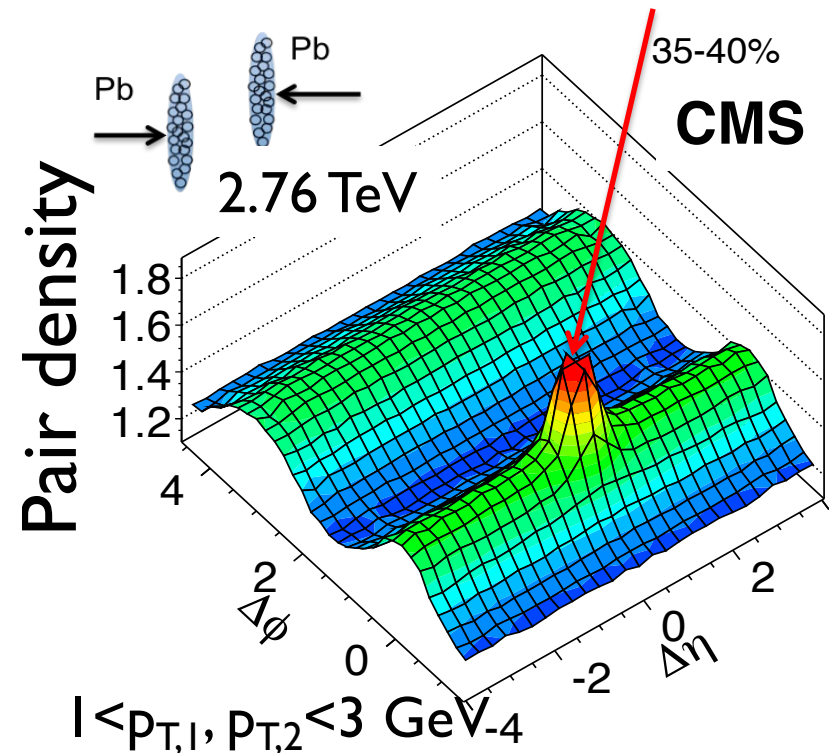
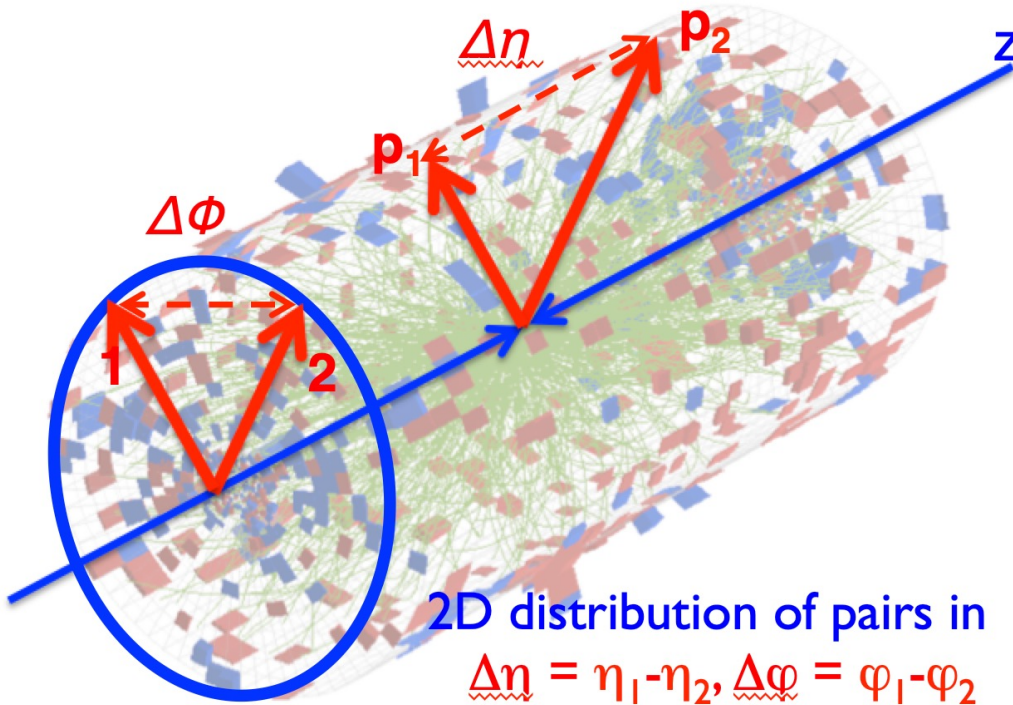
- Particle pair distribution

$$\frac{dN^{pair}}{d\Delta\phi} = \frac{N^{pair}}{2\pi} \left[1 + 2 \sum_n V_{n\Delta,ab} \cos n(\phi_a - \phi_b) \right]$$

$$V_{n\Delta,ab} = v_{n,a} v_{n,b} \text{ when } \Phi_n^a = \Phi_n^b$$

- Two particle correlation method

Short-range correlations



Collectivity – how to measure

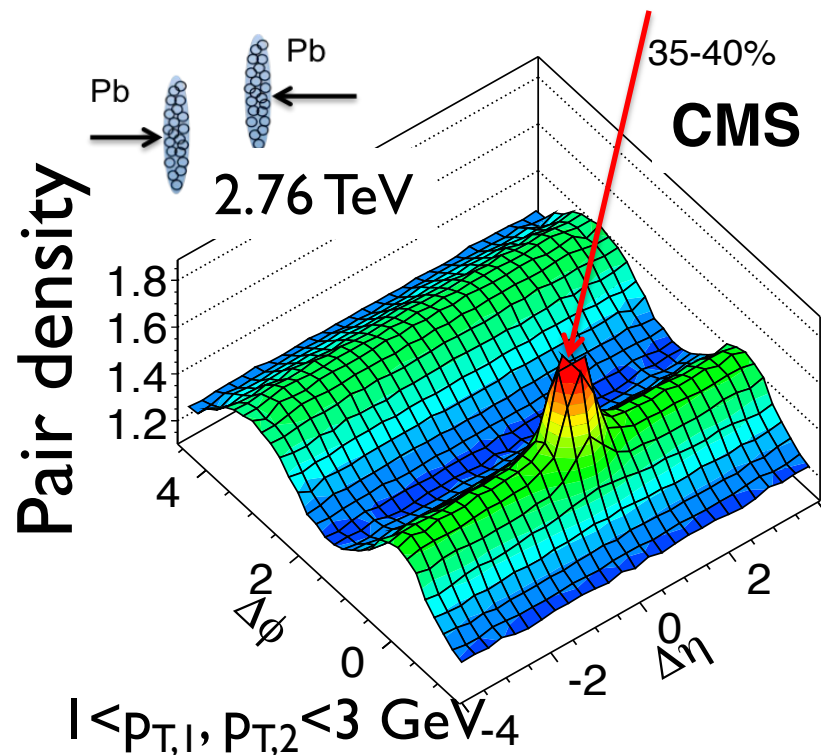
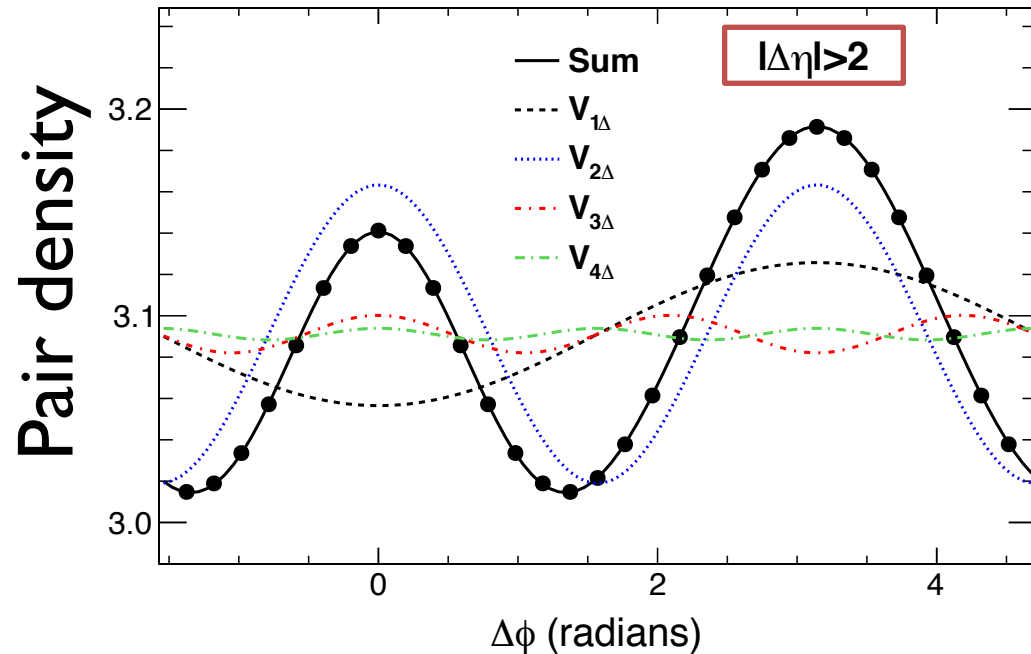
- Particle pair distribution

$$\frac{dN^{pair}}{d\Delta\phi} = \frac{N^{pair}}{2\pi} \left[1 + 2 \sum_n V_{n\Delta,ab} \cos n(\phi_a - \phi_b) \right]$$

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- Two particle correlation method

Short-range correlations



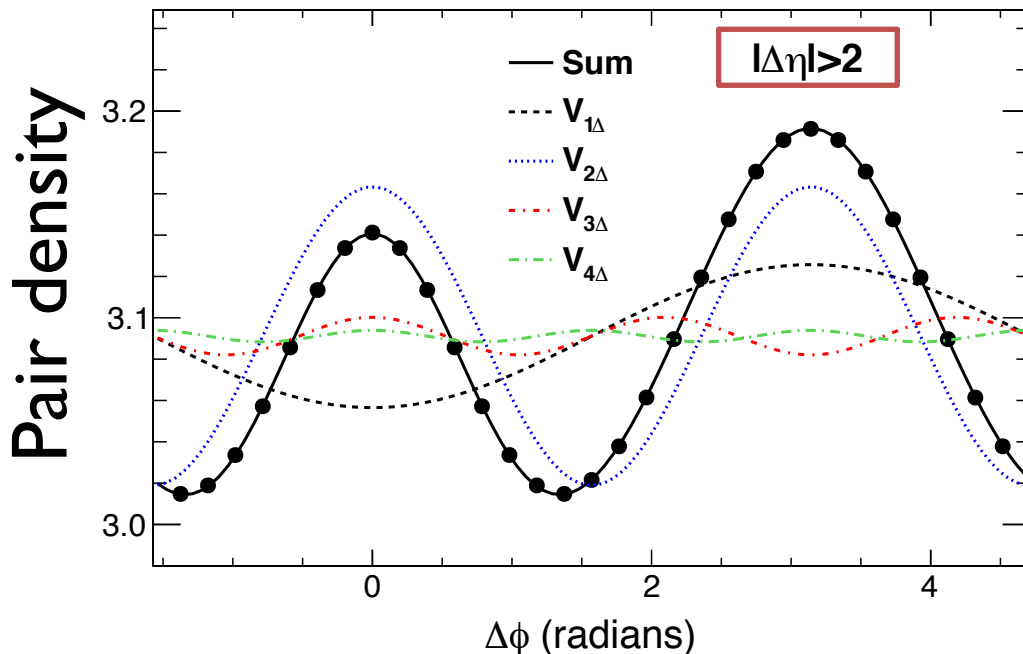
Collectivity – how to measure

- Particle pair distribution

$$V_{n\Delta,ab} = v_{n,a}v_{n,b} \quad \text{when } \Phi_n^a = \Phi_n^b$$

$$\frac{dN^{pair}}{d\Delta\phi} = \frac{N^{pair}}{2\pi} \left[1 + 2 \sum_n V_{n\Delta,ab} \cos n(\phi_a - \phi_b) \right]$$

- Two particle correlation method



Fourier analysis:

$$\frac{dN^{pair}}{d\Delta\phi} \sim 1 + 2V_{1\Delta} \cos(\Delta\phi) + 2V_{2\Delta} \cos(2\Delta\phi) + 2V_{3\Delta} \cos(3\Delta\phi) + 2V_{4\Delta} \cos(4\Delta\phi) + \dots$$

$$V_{n\Delta} = v_{n,a}v_{n,b}$$

Collectivity – how to measure

- Multi-particle distribution

$$\frac{dN_1}{d\phi_1} \cdots \frac{dN_m}{d\phi_m} \Rightarrow v_{n_1} v_{n_2} \cdots v_{n_m} \cos(n_1\phi_1 + n_2\phi_2 + \cdots + n_m\phi_m)$$

Collectivity – how to measure

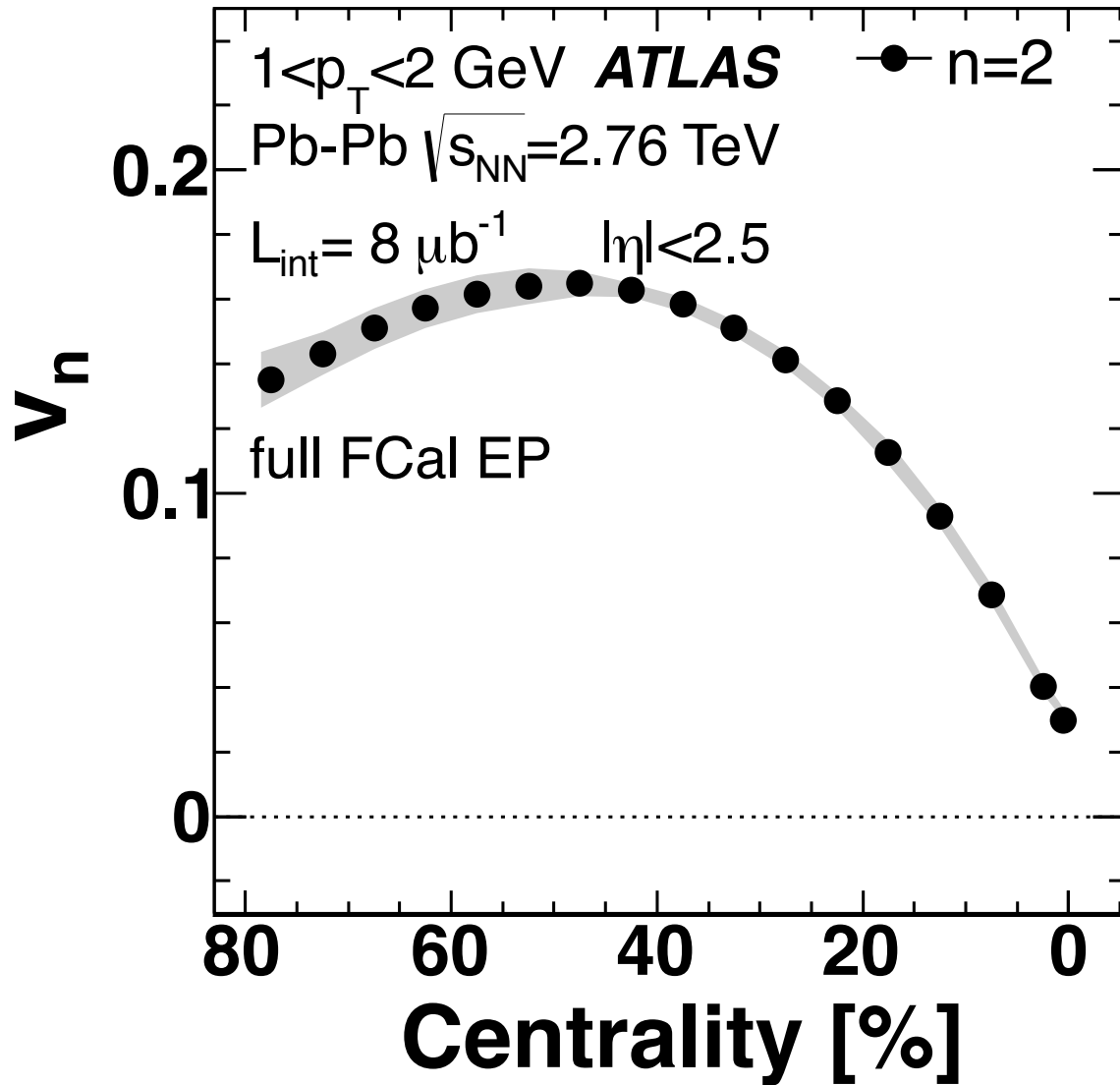
- Multi-particle distribution

$$\frac{dN_1}{d\phi_1} \cdots \frac{dN_m}{d\phi_m} \Rightarrow v_{n_1} v_{n_2} \cdots v_{n_m} \cos(n_1\phi_1 + n_2\phi_2 + \cdots + n_m\phi_m)$$

- Cumulant method
 - PRC 83 (2011) 044913, $v_n\{m\}$ with multi-particle correlation
 - PRC 89 (2014) 064904, Symmetric Cumulant (*Standard Candles*)
 - PRC 90 (2014) 024905, Asymmetric Cumulant
- $v_n\{m\}$ – Less sensitive to few particle correlation
- SC – Event-by-event correlation between v_n and v_m
- AC – Initial geometry fluctuation & non-linear mixing between v_n, v_m

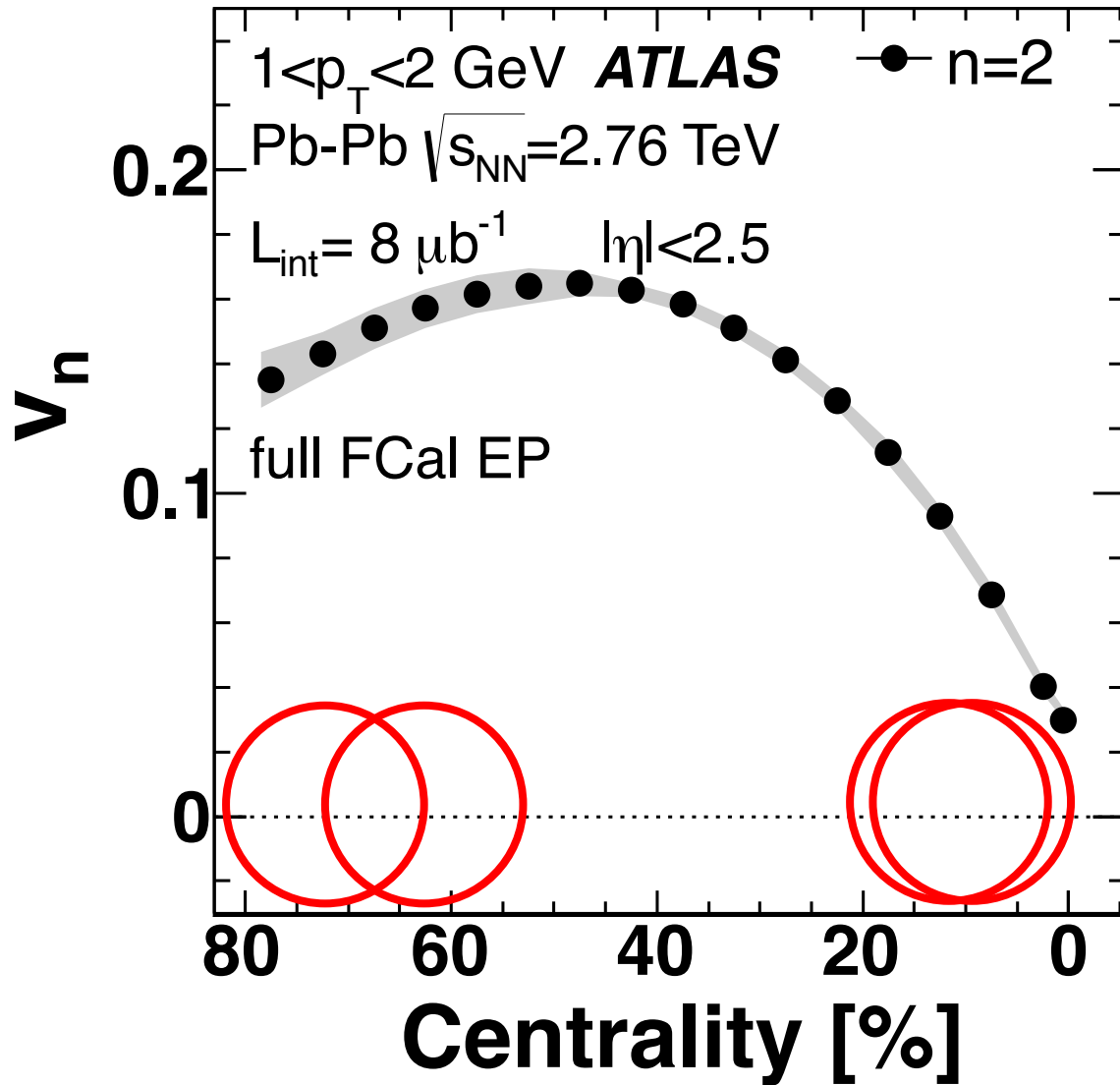
v_2 vs centrality

PRC 86 (2012) 014907



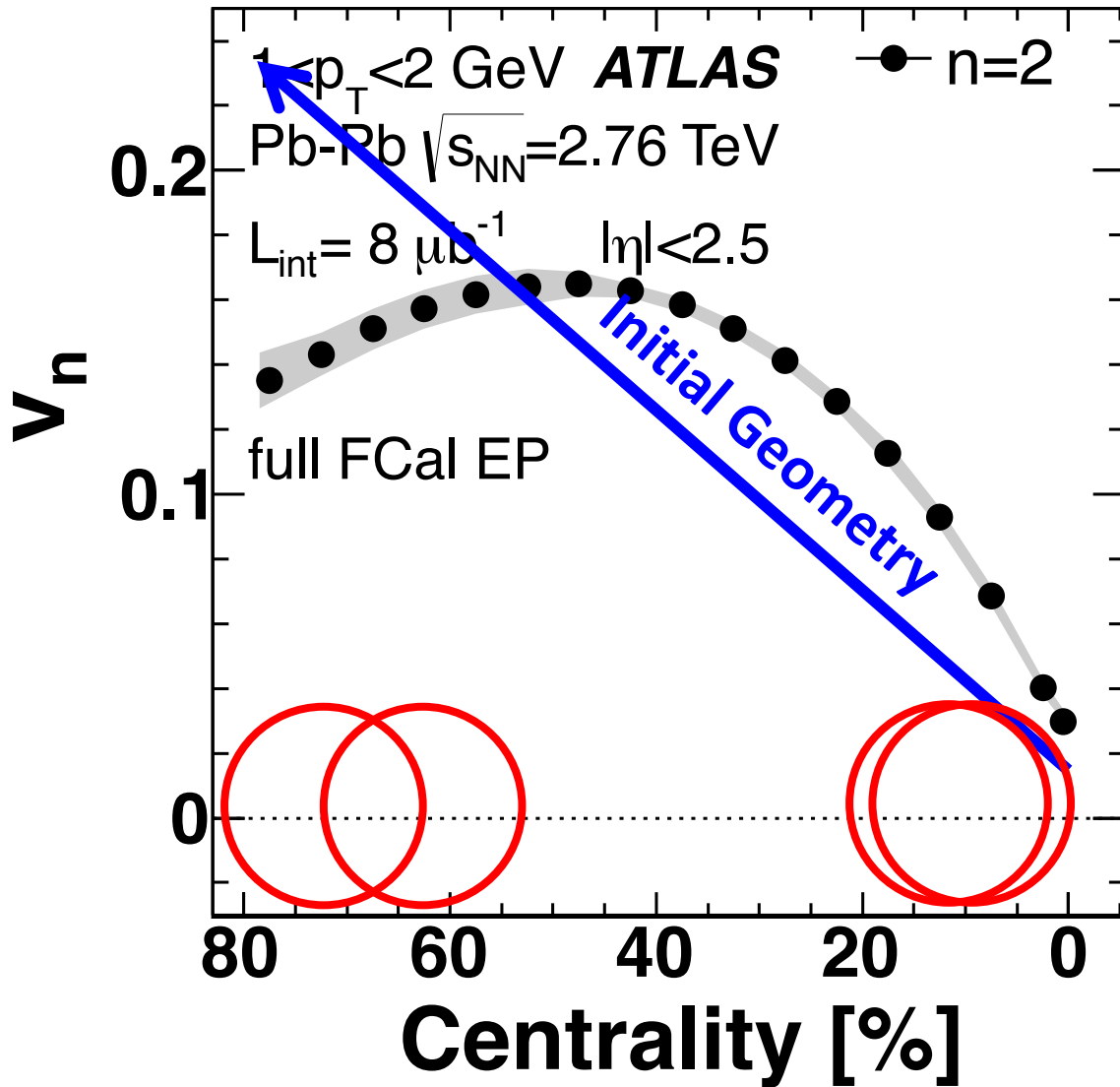
v_2 vs centrality

PRC 86 (2012) 014907



v_2 vs centrality

PRC 86 (2012) 014907

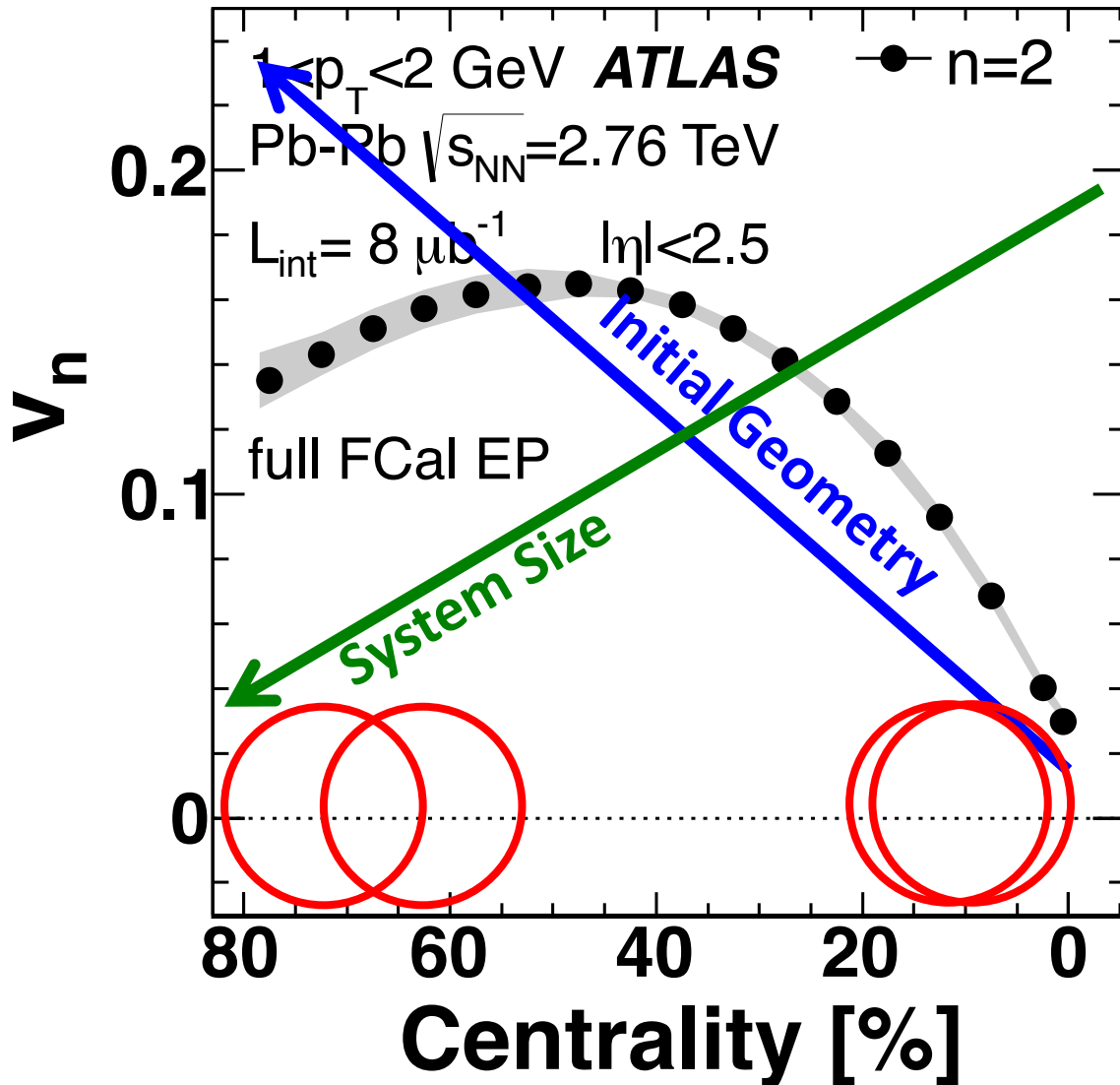


Geometry response:

$$v_n \propto A \varepsilon_n$$

v_2 vs centrality

PRC 86 (2012) 014907



Geometry response:

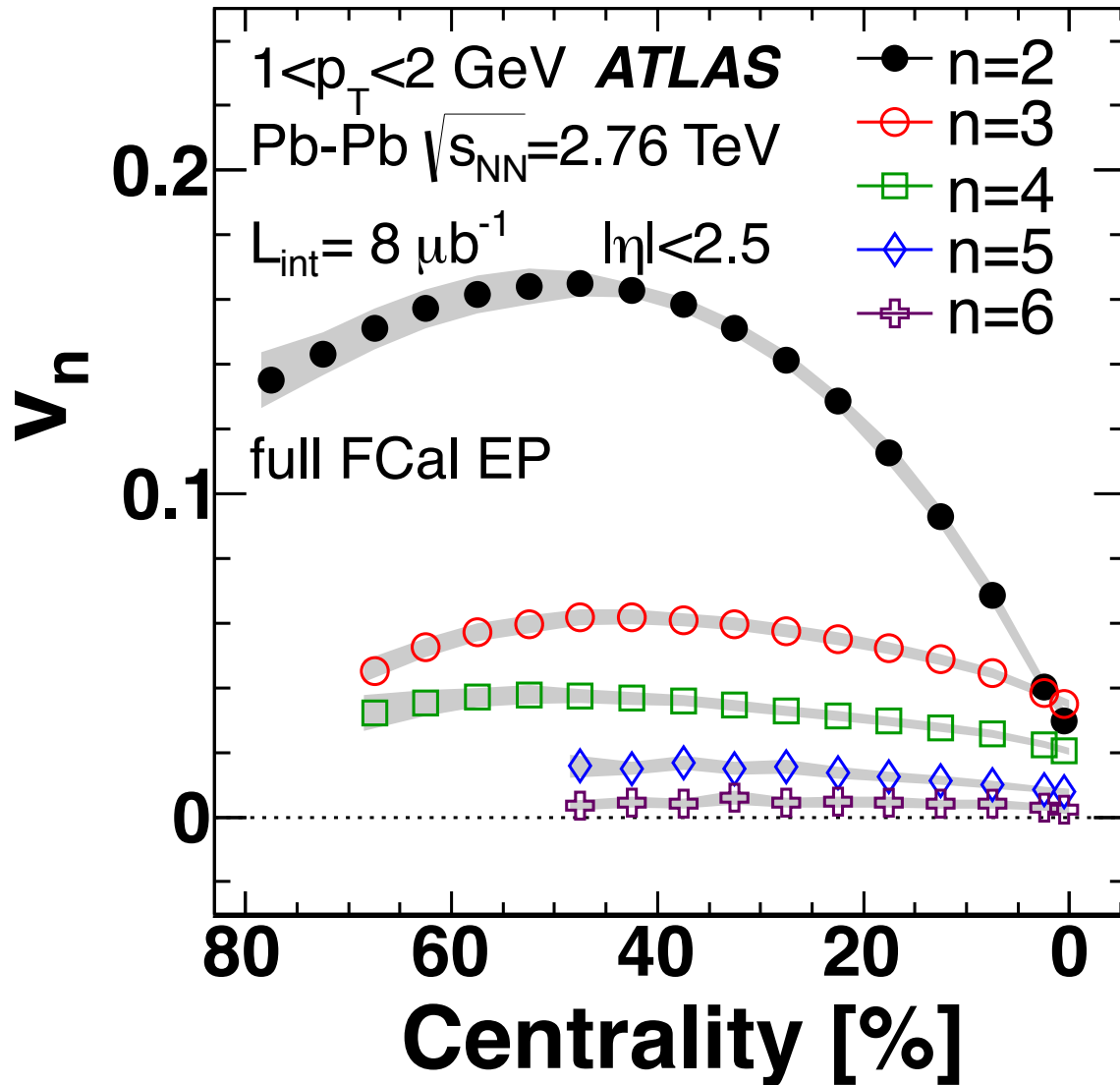
$$v_n \propto A \varepsilon_n$$

System Size:

Less interactions
develop less flow

v_n vs centrality

PRC 86 (2012) 014907



Geometry response:

$$v_n \propto A \varepsilon_n$$

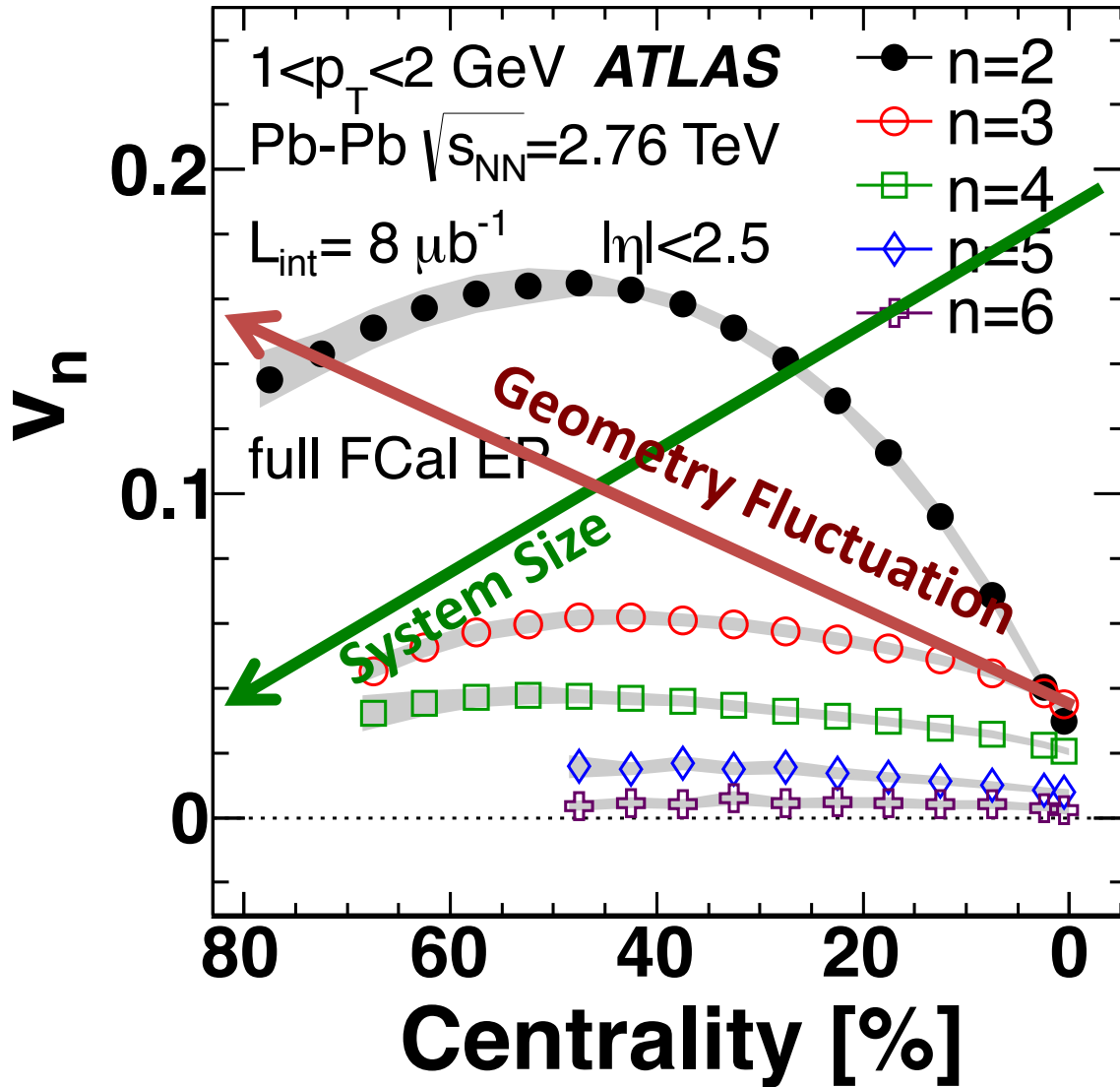
System Size:

Less interactions
develop less flow

Eccentricity from
Geometry Fluctuation

v_n vs centrality

PRC 86 (2012) 014907



Geometry response:

$$v_n \propto A \varepsilon_n$$

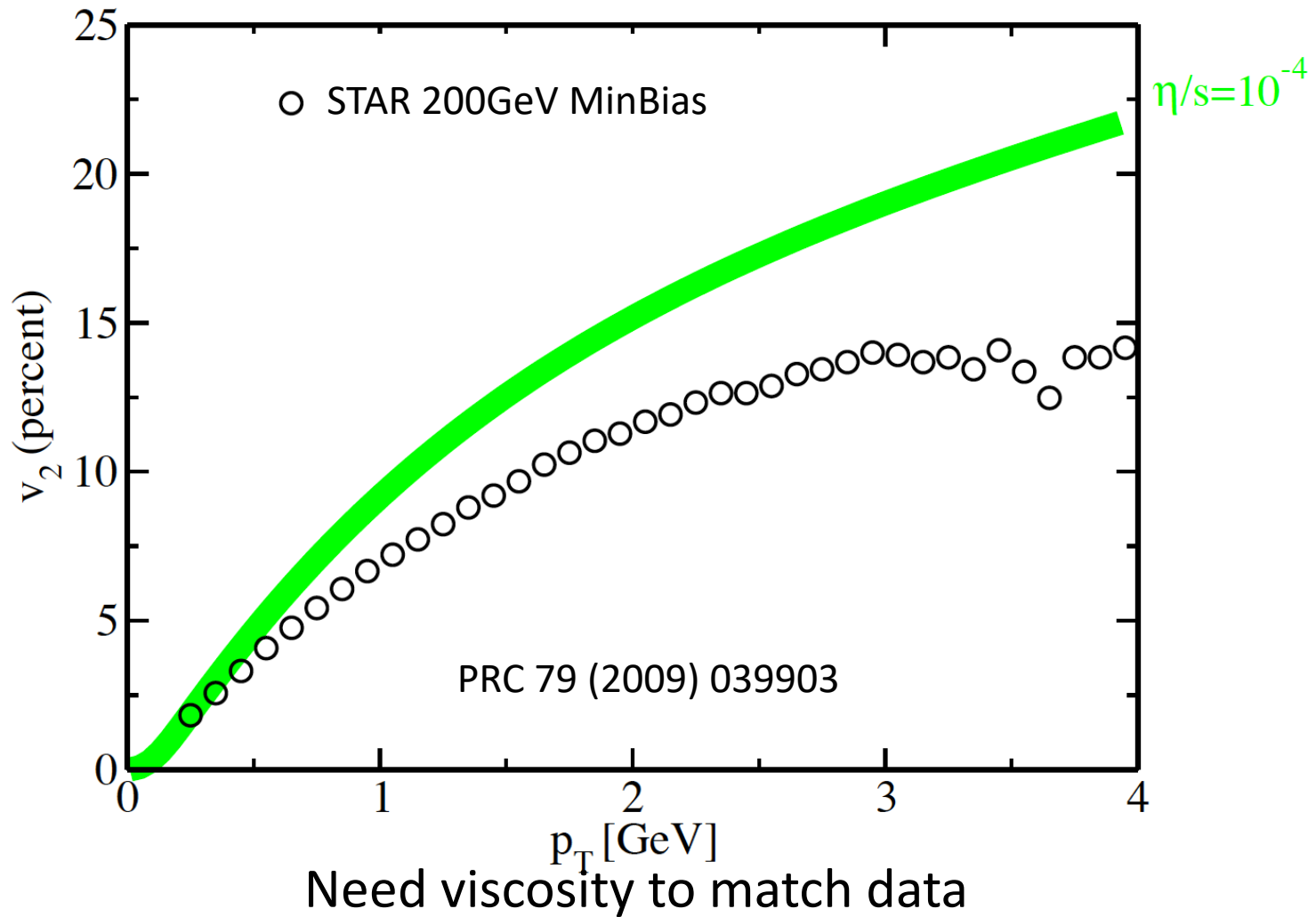
System Size:

Less interactions
 develop less flow

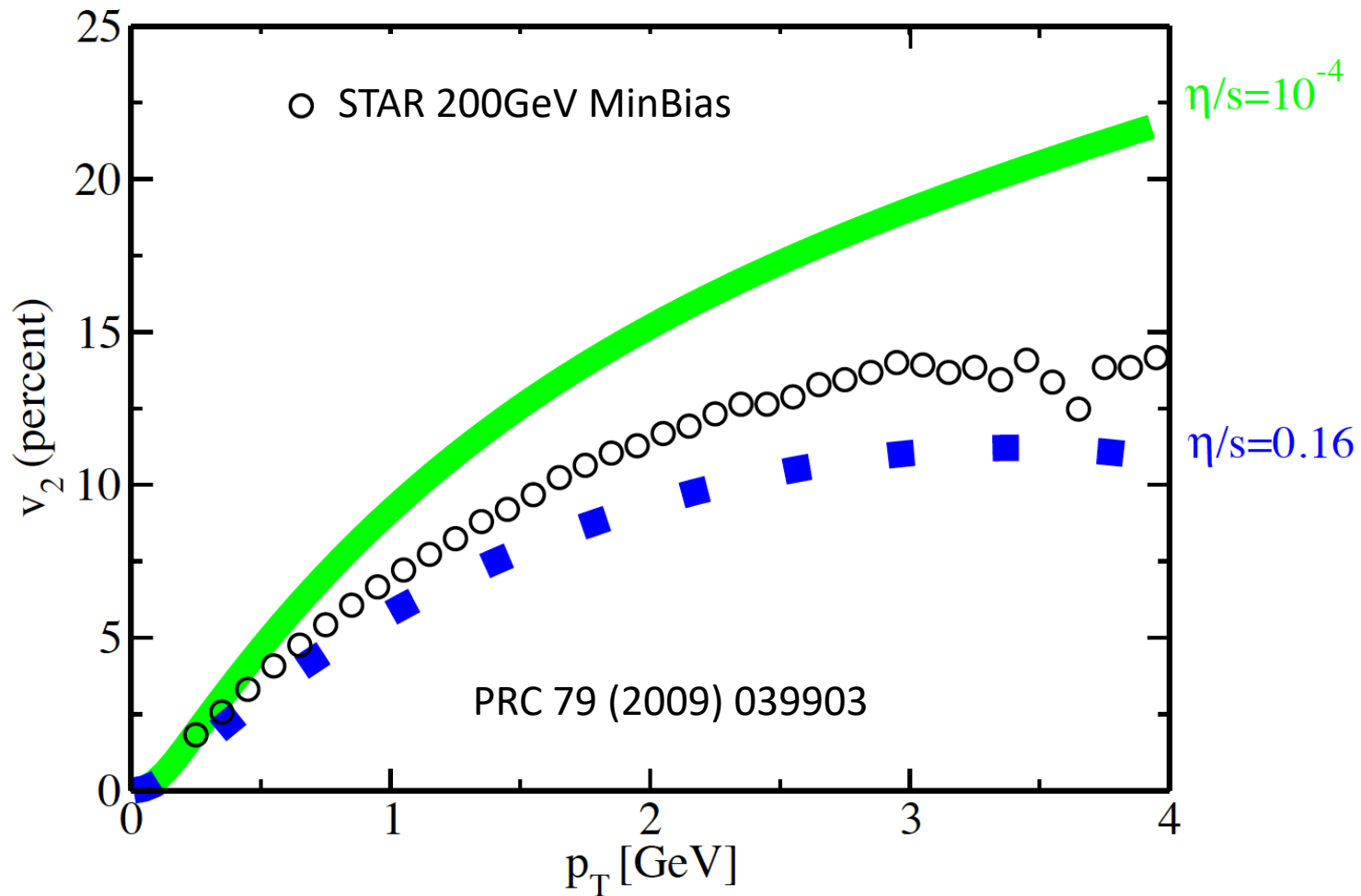
Eccentricity from
 Geometry Fluctuation:

$$GF \propto 1/N_{par}$$

v_2 vs p_T – constrain viscosity

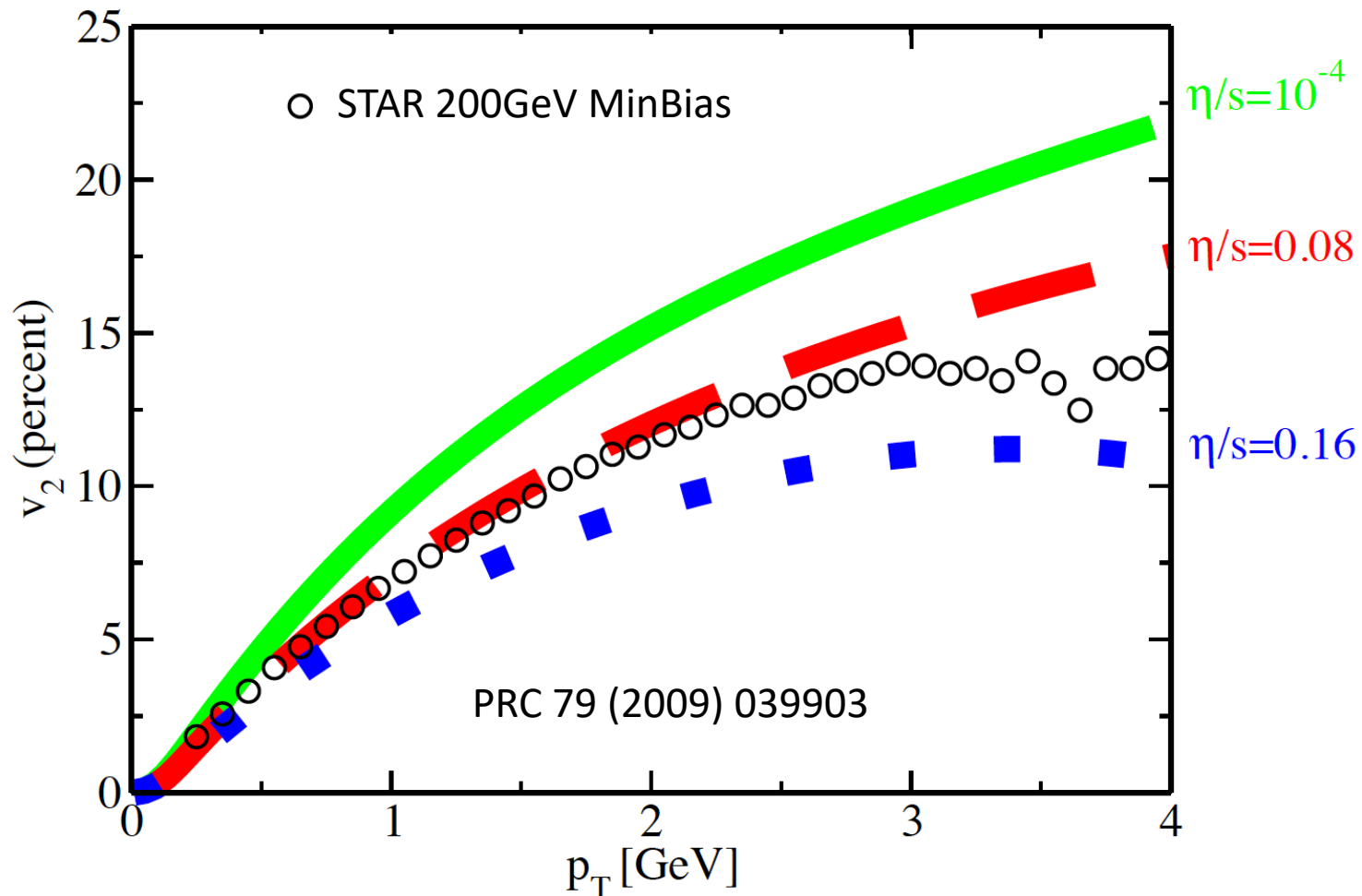


v_2 vs p_T – constrain viscosity



Strong constraint on viscosity

v_2 vs p_T – constrain viscosity

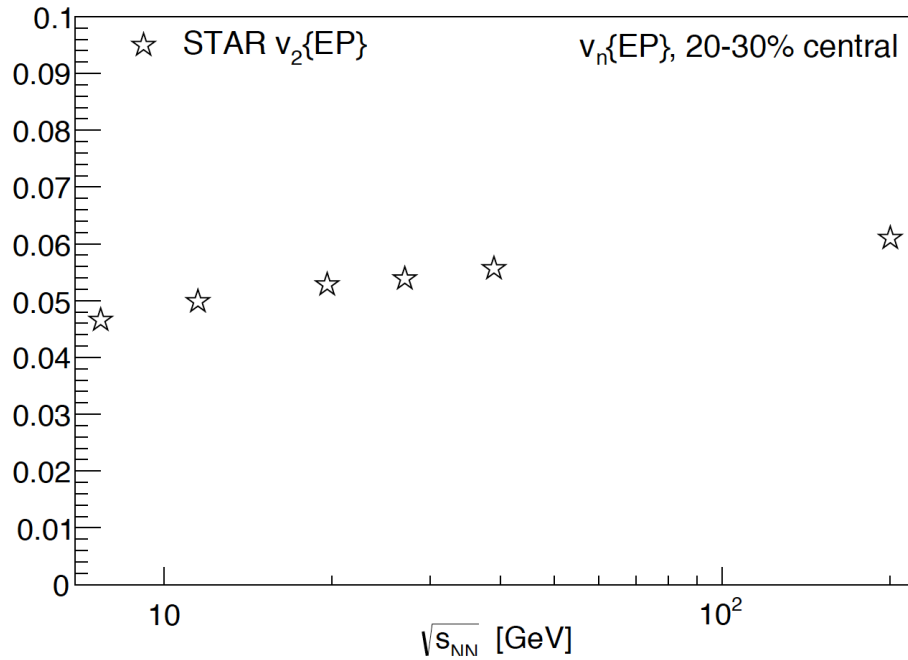


Strong constraint on viscosity

Reveal the nature of “perfect” fluid

Ads/CFT lower limit $\eta/s = 1/4\pi \approx 0.078$, PRL 94 (2005) 11601

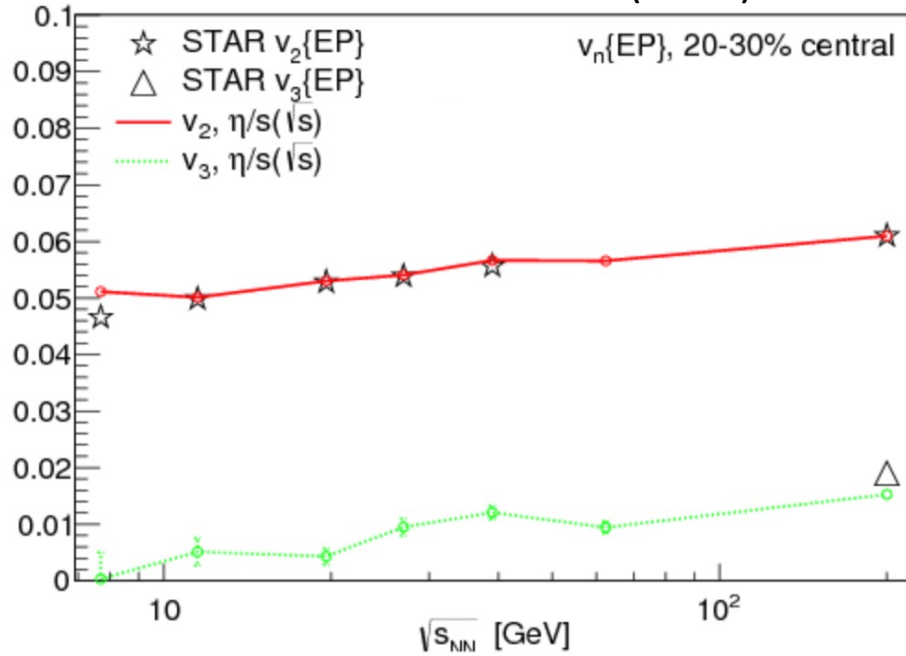
v_n vs. energy – viscosity vs. energy



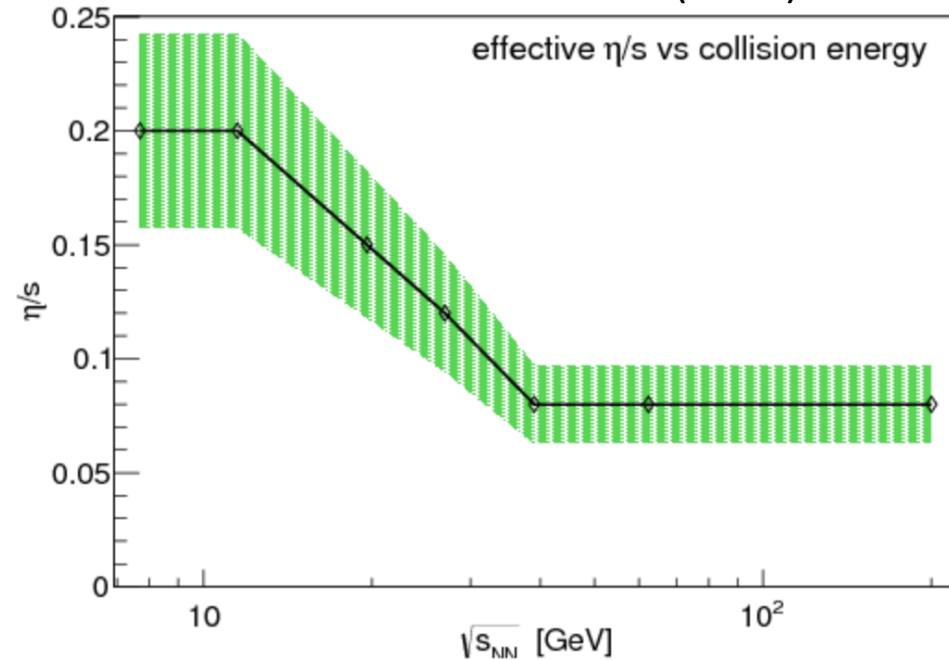
v_2 increases with collision energy

v_n vs. energy – viscosity vs. energy

PRC 91 (2015) 064901

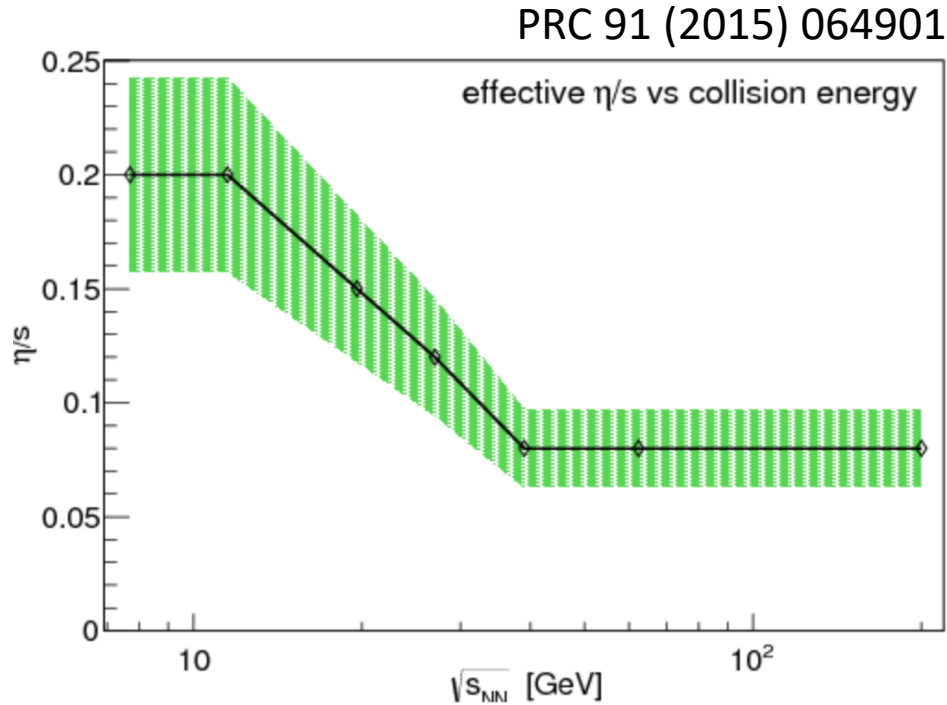
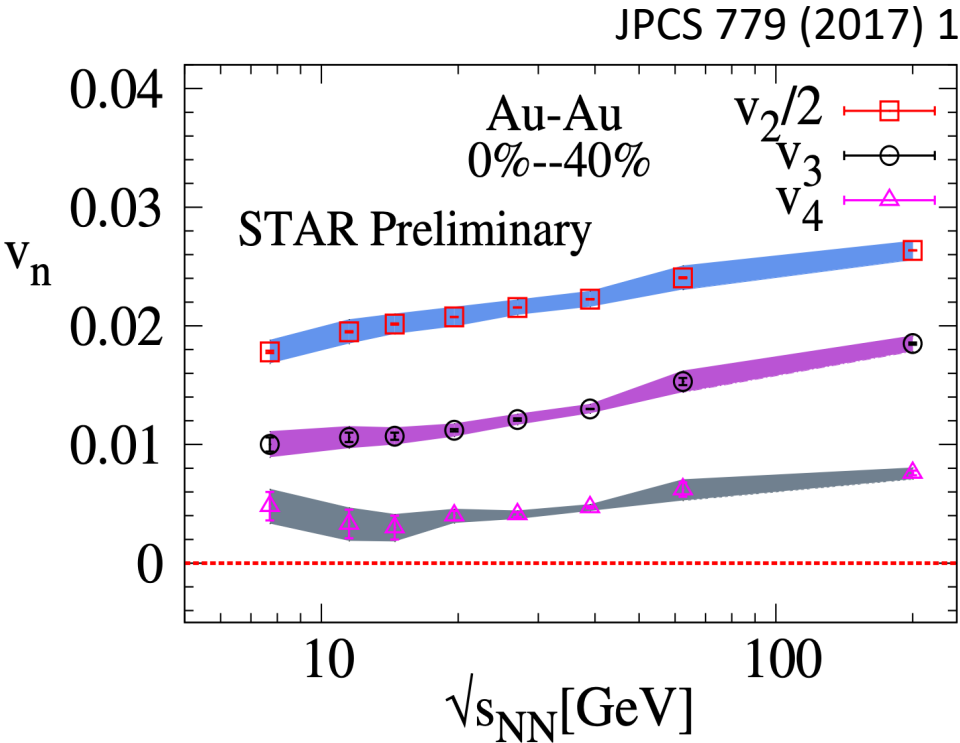


PRC 91 (2015) 064901



v_2 increases with collision energy
Indication of collision energy dependent viscosity

v_n vs. energy – viscosity vs. energy

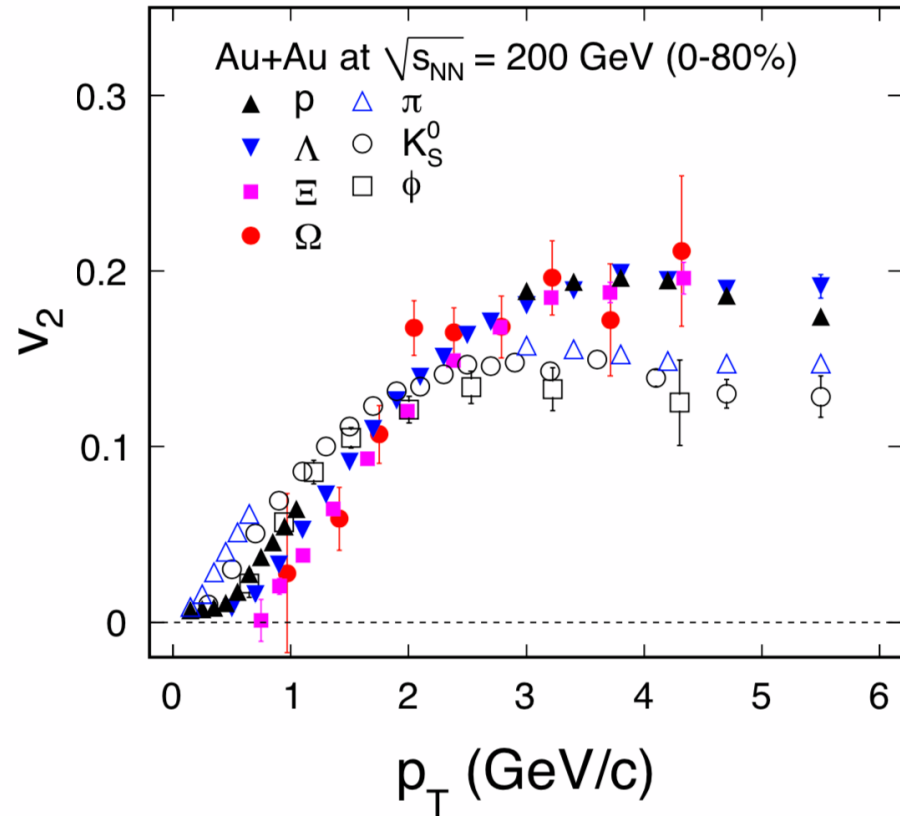
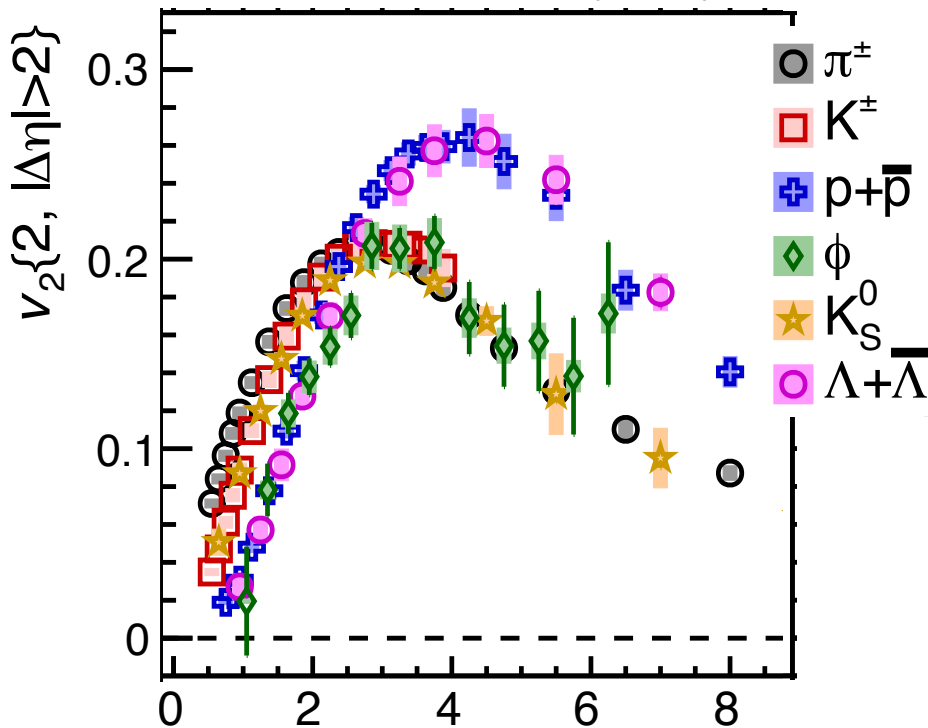


v_2 increases with collision energy
Indication of collision energy dependent viscosity
New results available for further investigation

v_n vs PID

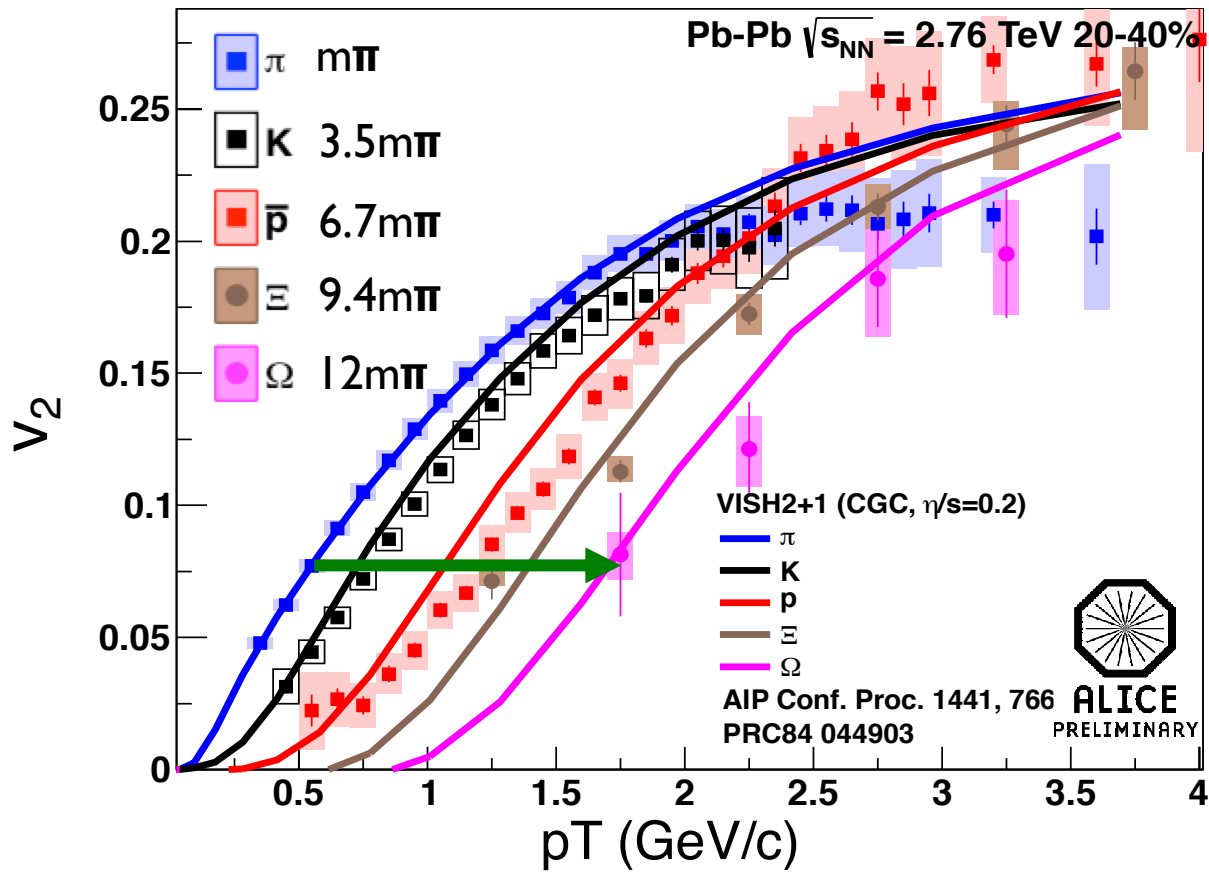
v_n vs PID – mass ordering

JHEP 09 (2018) 006



Mass ordering at low $p_T < \approx 3\text{GeV}$
 Heavier particle has smaller v_n at same p_T

v_n vs PID – mass ordering

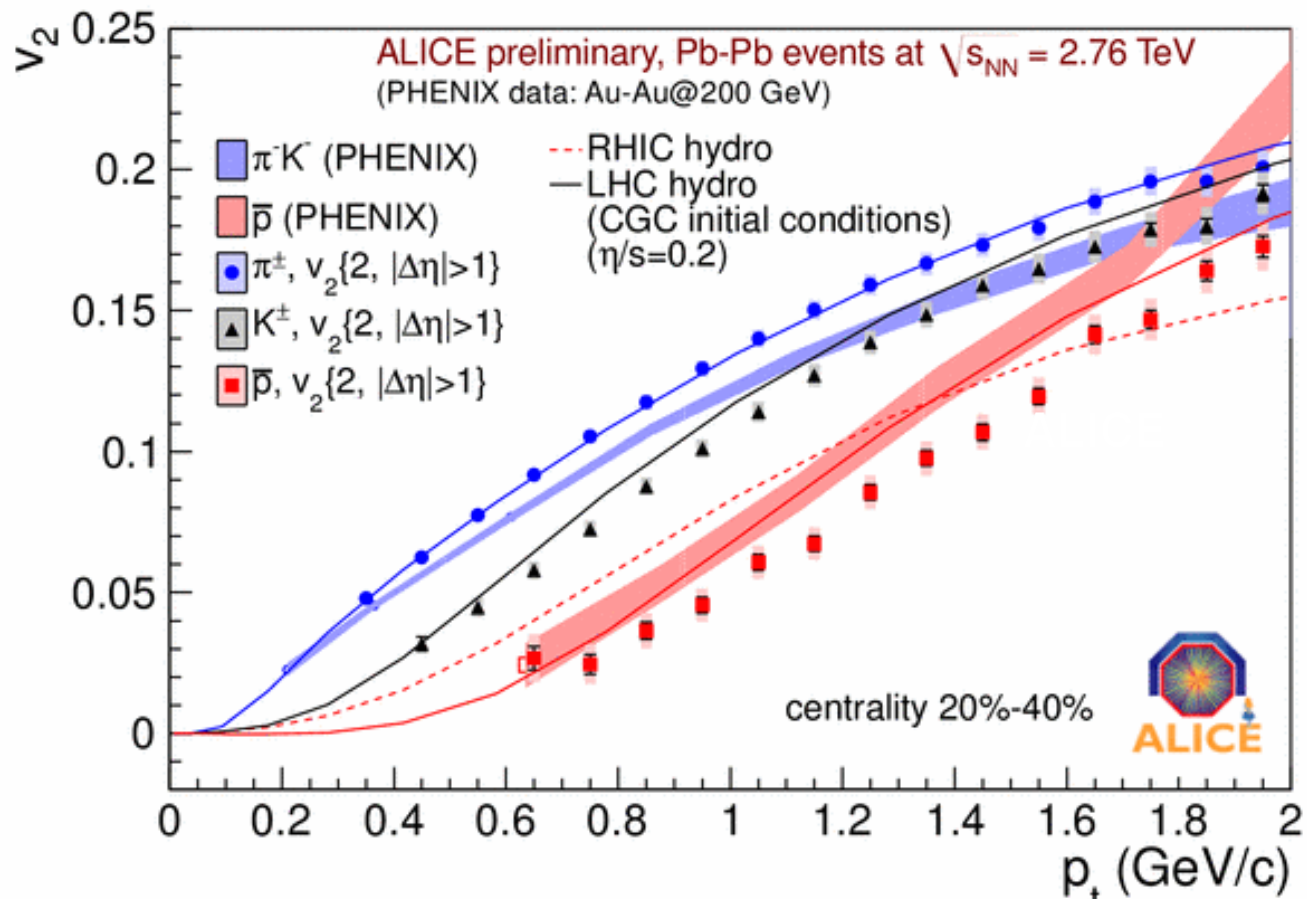


Mass ordering at low $p_T < \approx 3\text{GeV}$

Heavier particle has smaller v_n at same p_T

Predicted by hydro: common boost from pressure gradient

v_n vs PID – mass ordering



Mass ordering at low $p_T < \approx 3\text{GeV}$

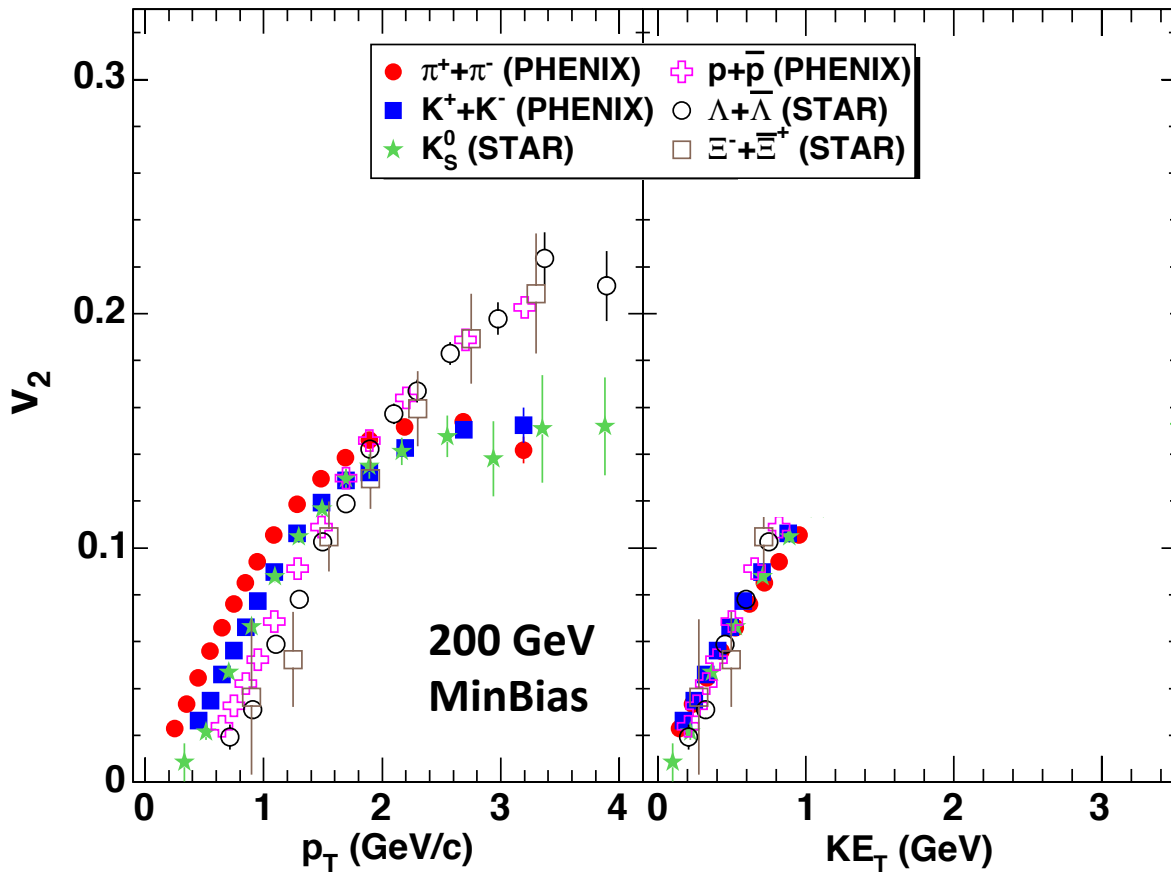
Heavier particle has smaller v_n at same p_T

Predicted by hydro: common boost from pressure gradient

Indicate larger boost @ LHC energies

v_n vs PID – KE_T scaling

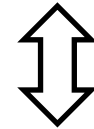
PRL 98 (2007) 162301



KE_T scaling @ $KE_T \lesssim 1\text{GeV}$: Hydro pressure gradient

$$KE_T = \sqrt{m^2 + p_T^2} - m$$

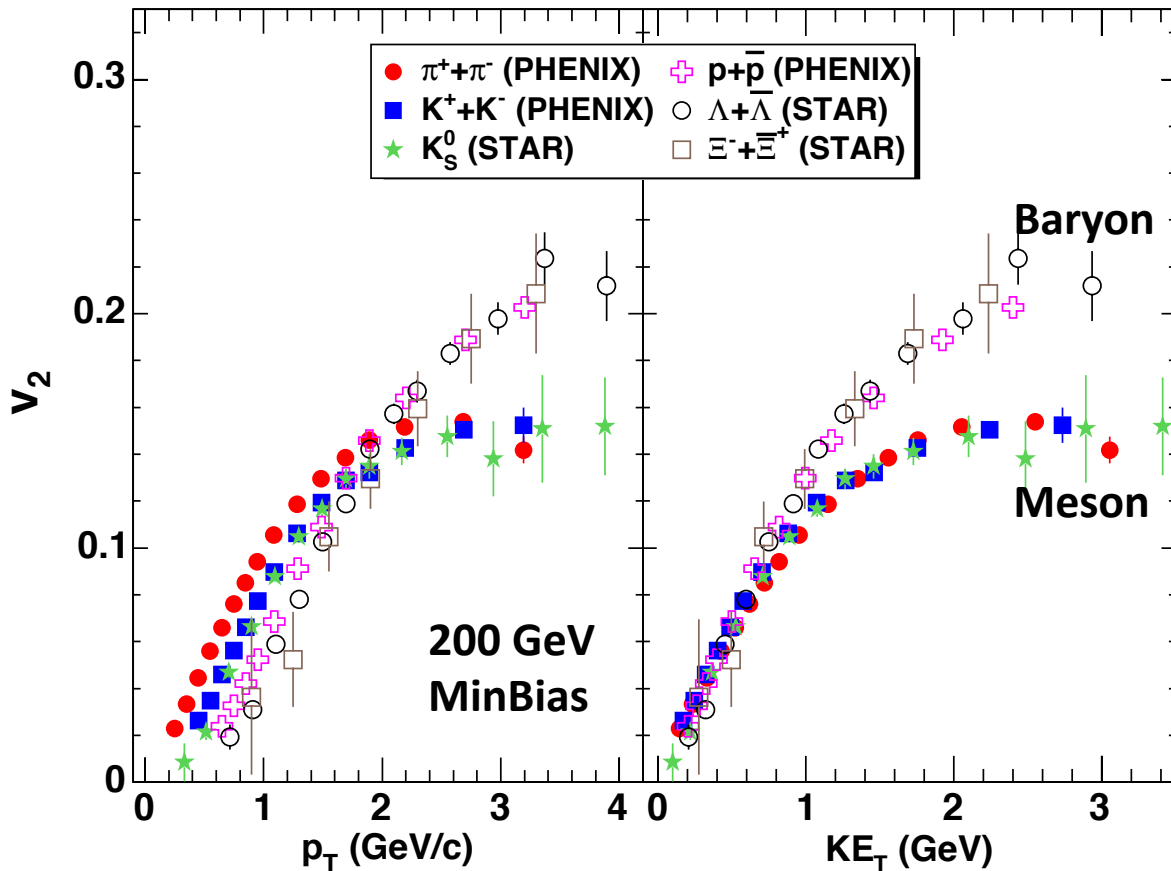
Pressure gradient



collective kinetic energy

v_n vs PID – Quark Coalescence

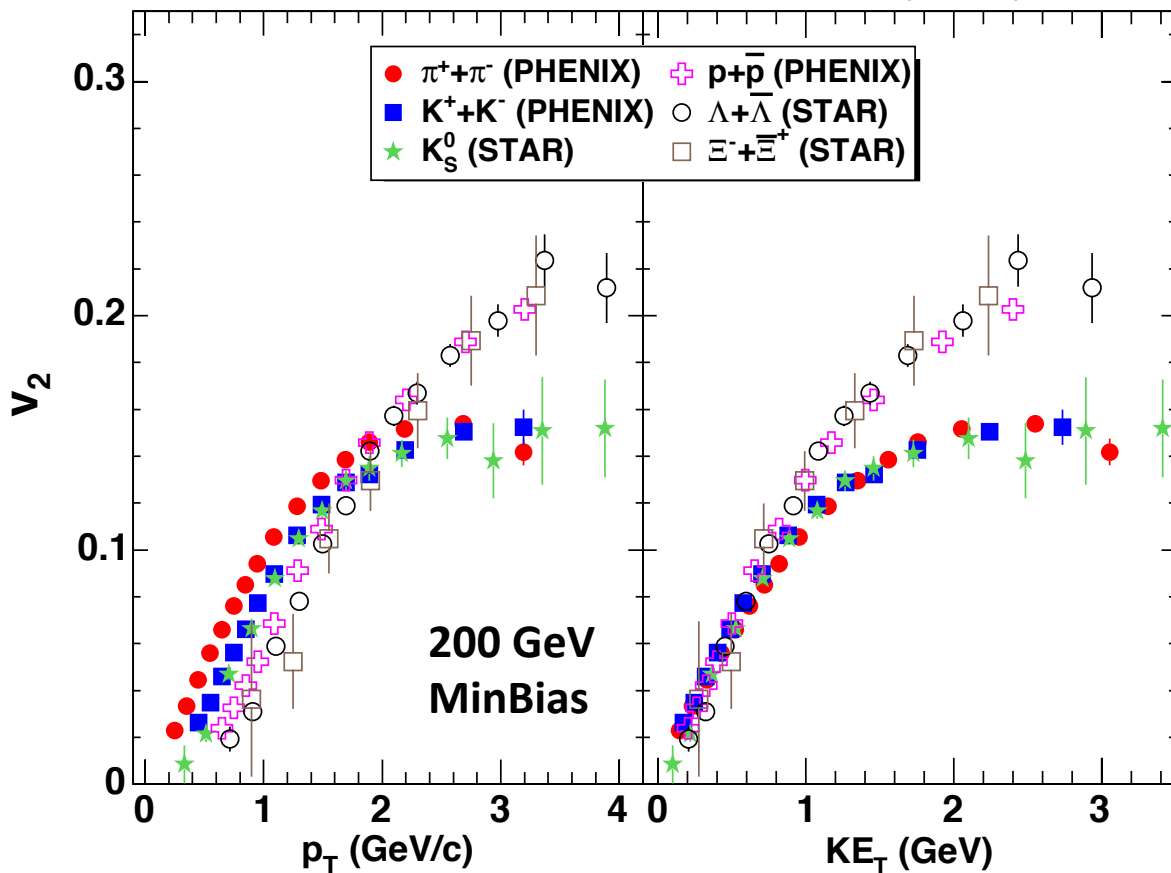
PRL 98 (2007) 162301



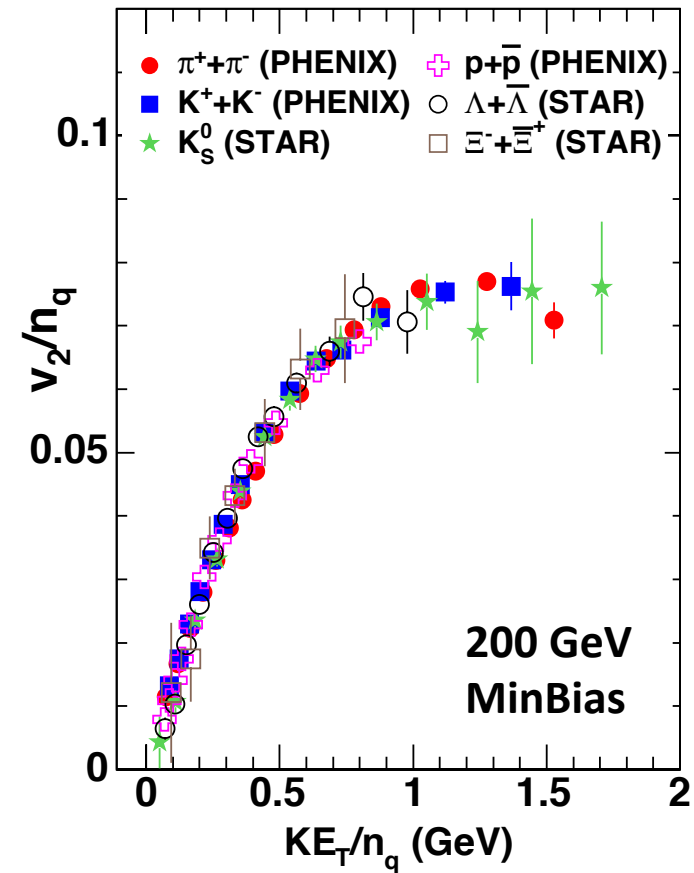
KE_T scaling @ $KE_T < \approx 1\text{GeV}$: Hydro pressure gradient
Meson/Baryon splitting: Quark coalescence

v_n vs PID – NCQ scaling

PRL 98 (2007) 162301



PRL 98 (2007) 162301



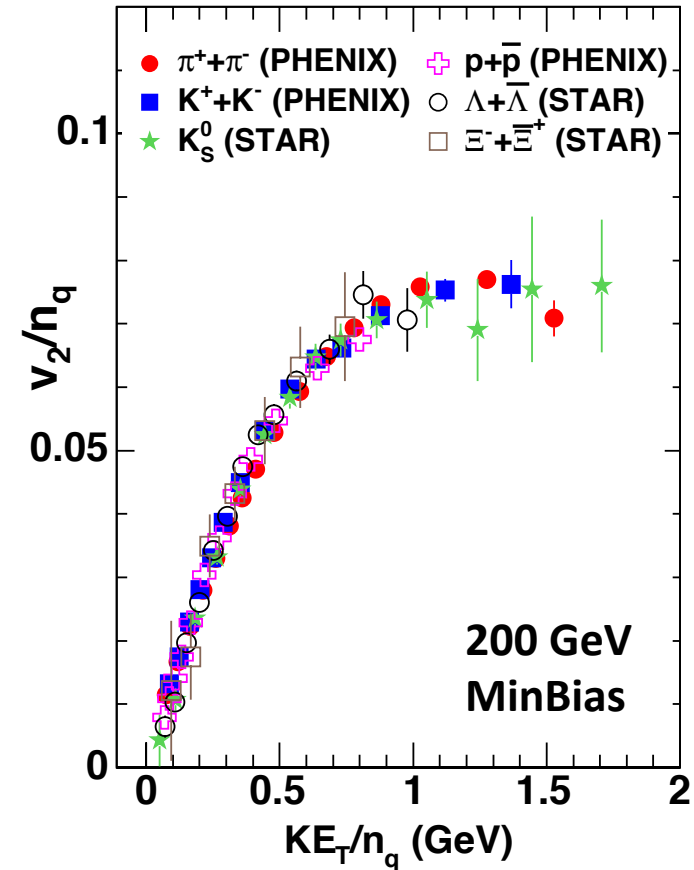
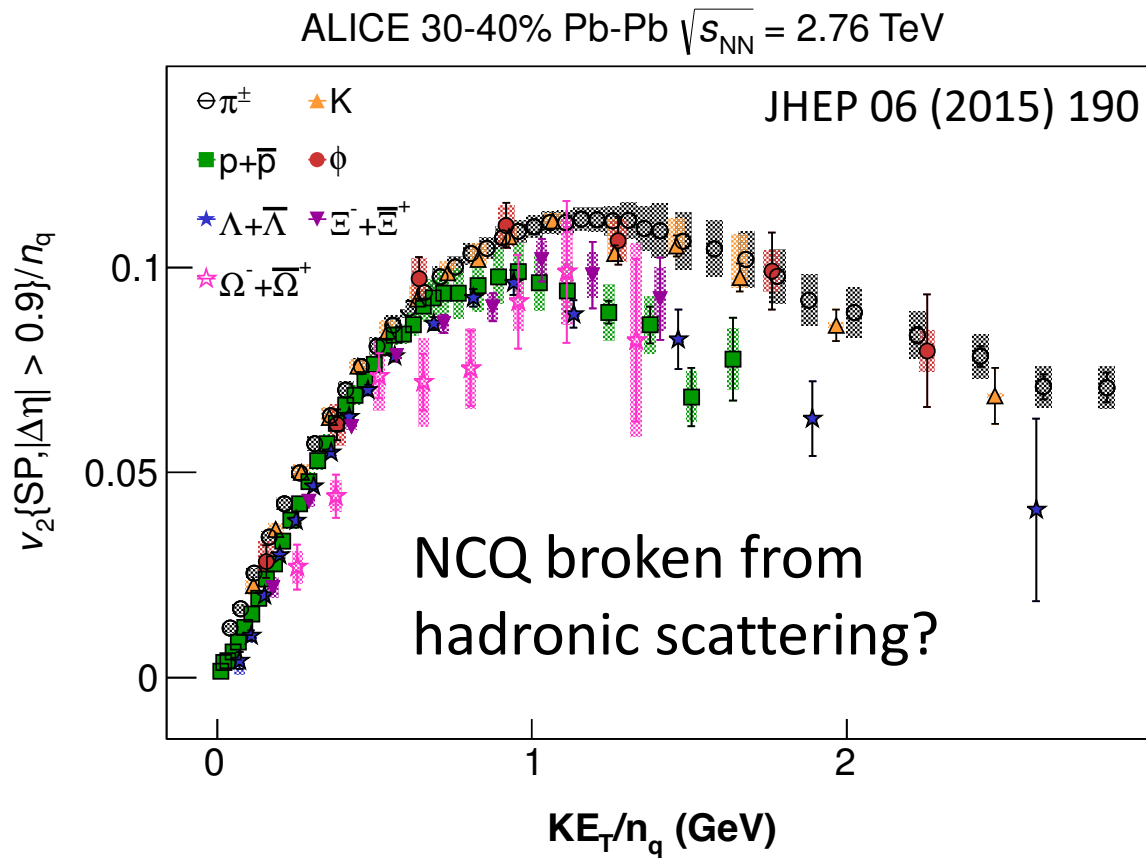
KE_T scaling @ $KE_T < \approx 1\text{GeV}$: Hydro pressure gradient

Meson/Baryon splitting: Quark coalescence

Number of Constituent Quark scaling: Parton degree of freedom

v_n vs PID – NCQ scaling

PRL 98 (2007) 162301

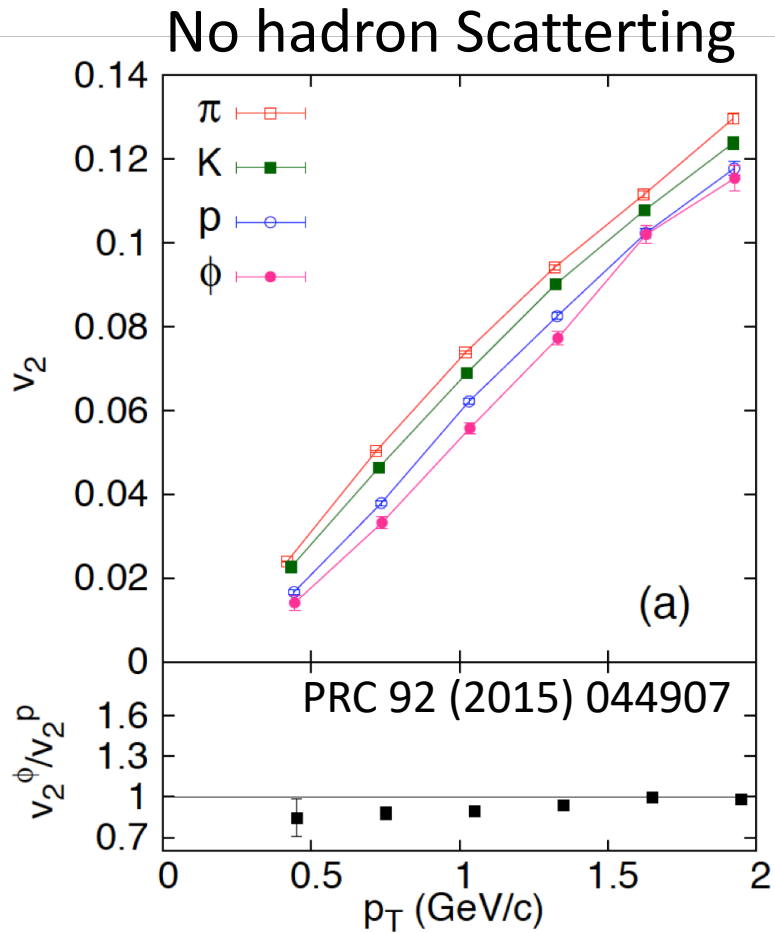


KE_T scaling @ $KE_T < \approx 1$ GeV: Hydro pressure gradient

Meson/Baryon splitting: Quark coalescence

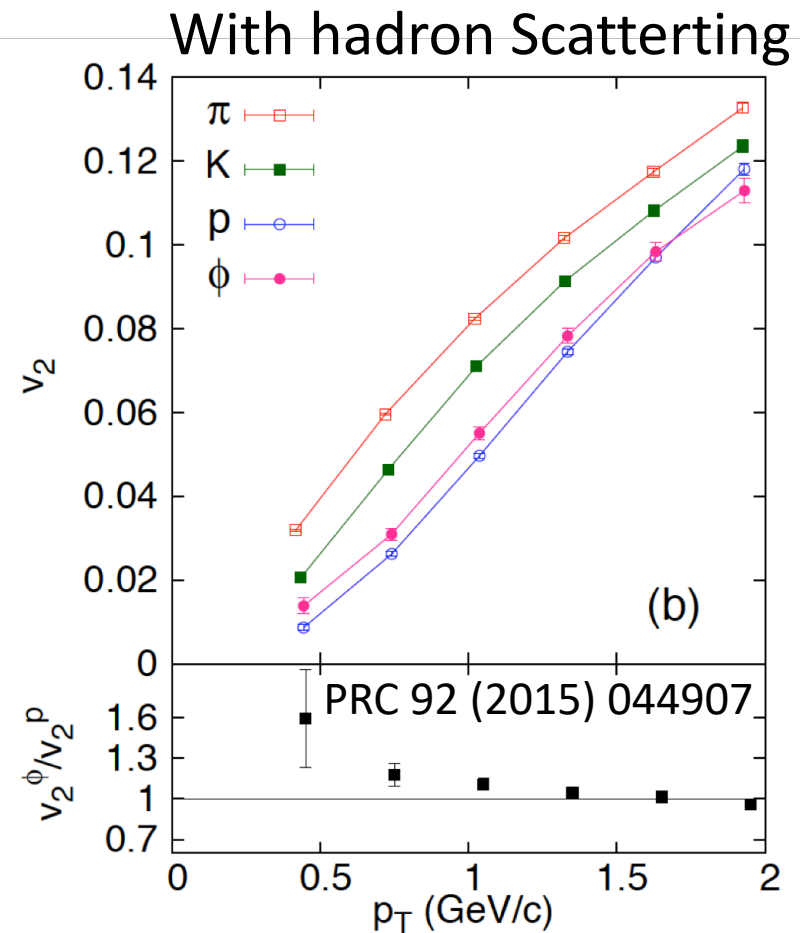
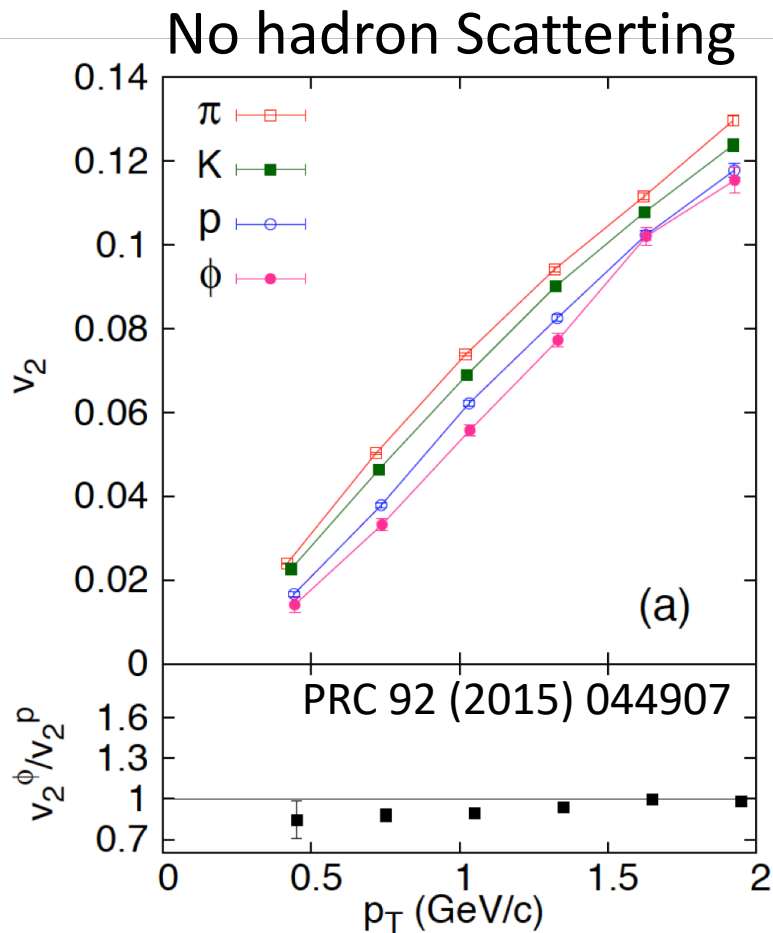
Number of Constituent Quark scaling: Parton degree of freedom

v_n vs PID – impact of hadronic phase



Ideal hydro + hadron cascade (JAM)

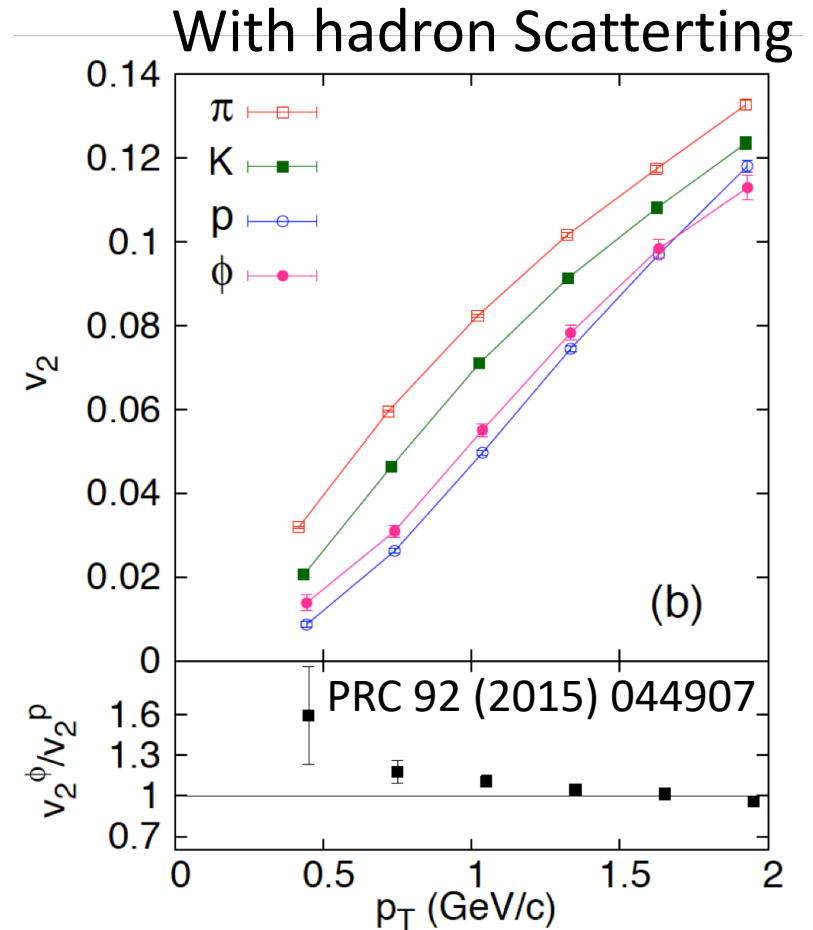
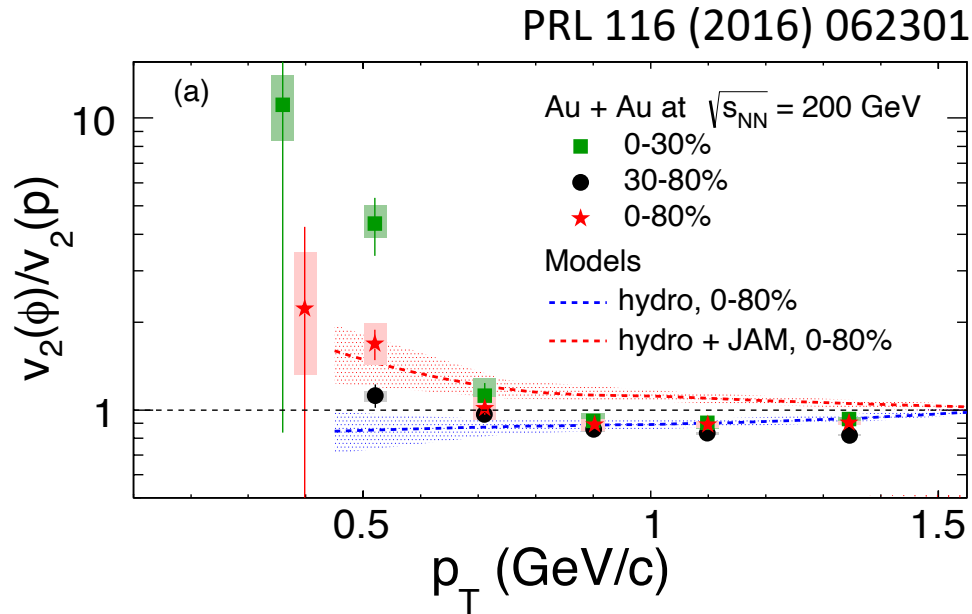
v_n vs PID – impact of hadronic phase



Ideal hydro + hadron cascade (JAM)

Small hadron cross section breaks mass ordering for ϕ

v_n vs PID – impact of hadronic phase

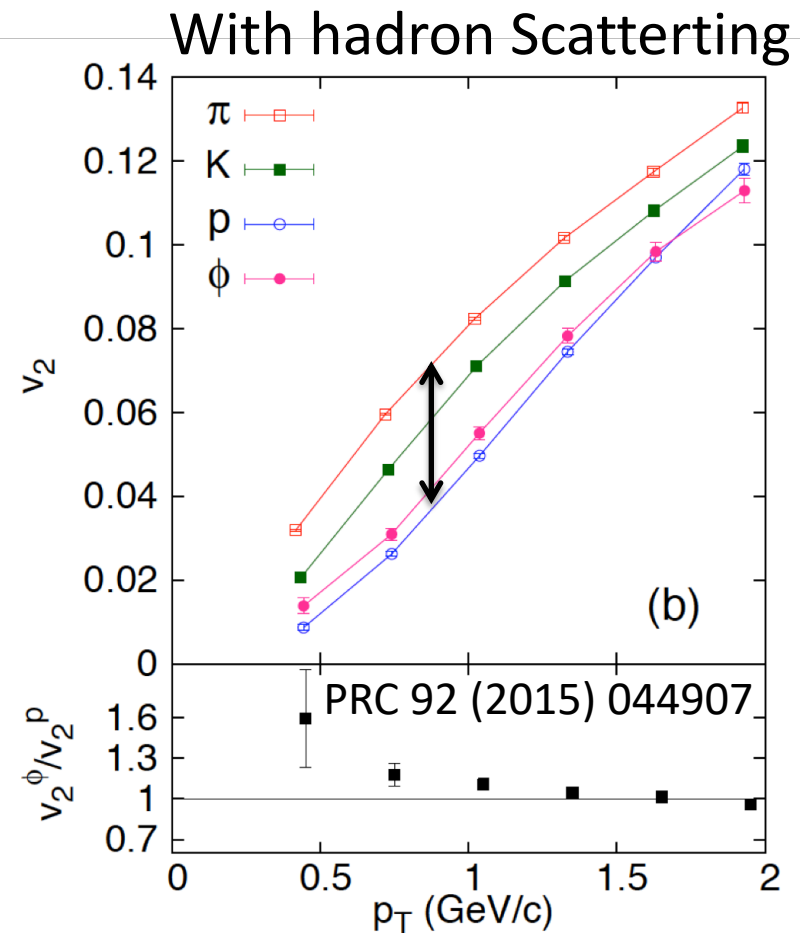
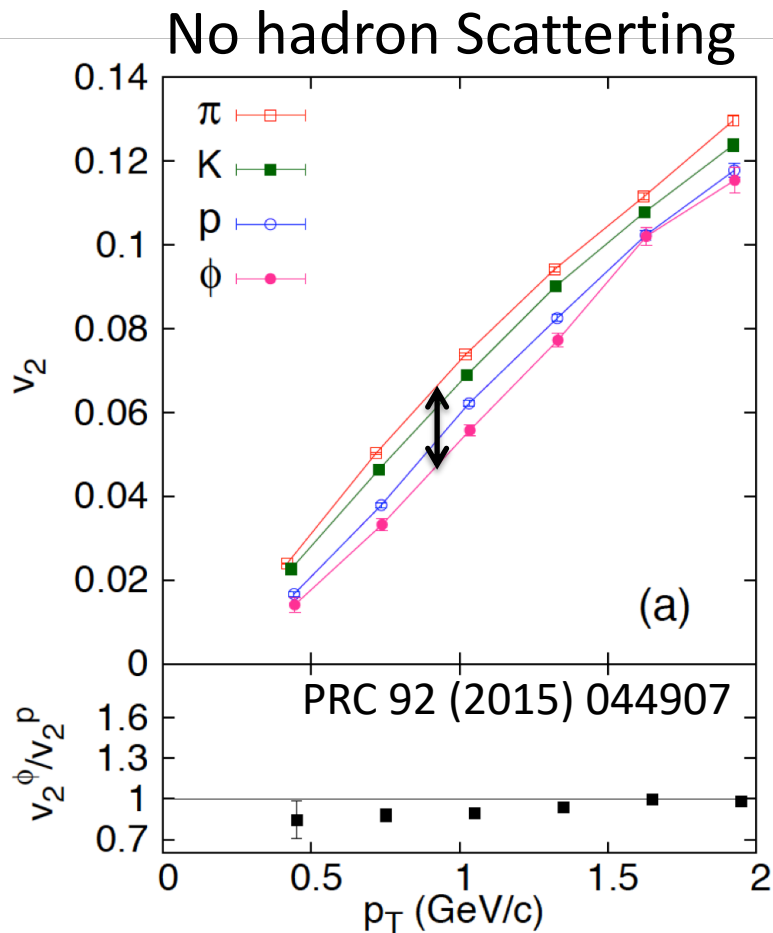


Ideal hydro + hadron cascade (JAM)

Small hadron cross section breaks mass ordering for ϕ

Confirmed by data

v_n vs PID – impact of hadronic phase



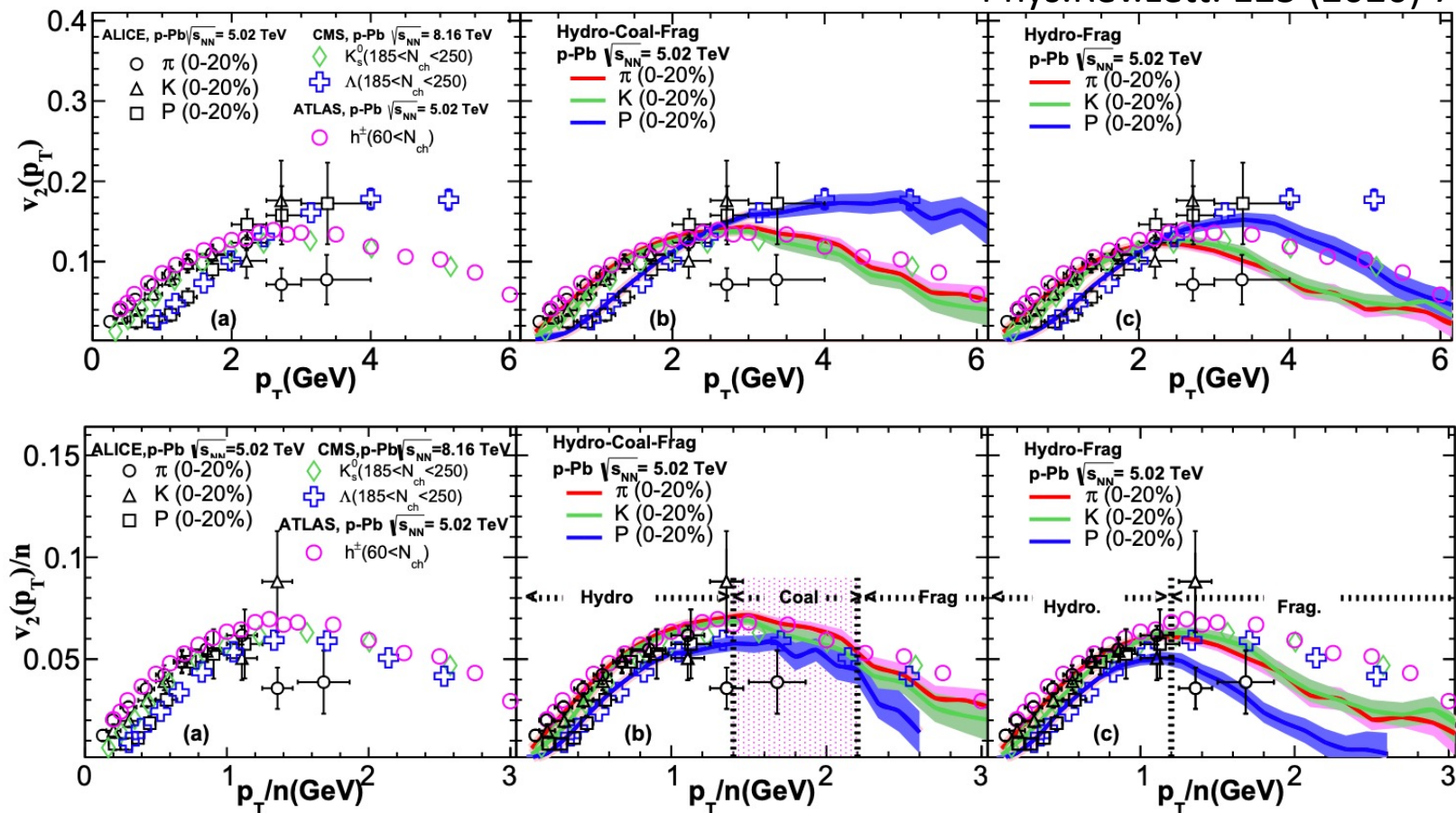
Ideal hydro + hadron cascade (JAM)

Small hadron cross section breaks mass ordering for phi

Confirmed by data; **implication on NCQ scaling?**

v_n vs PID – impact from coalescence?

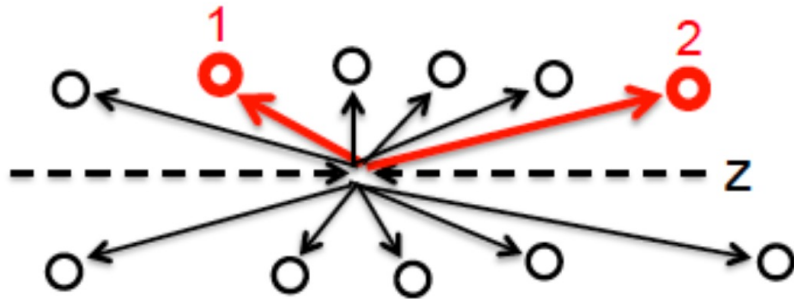
Phys.Rev.Lett. 125 (2020) 7



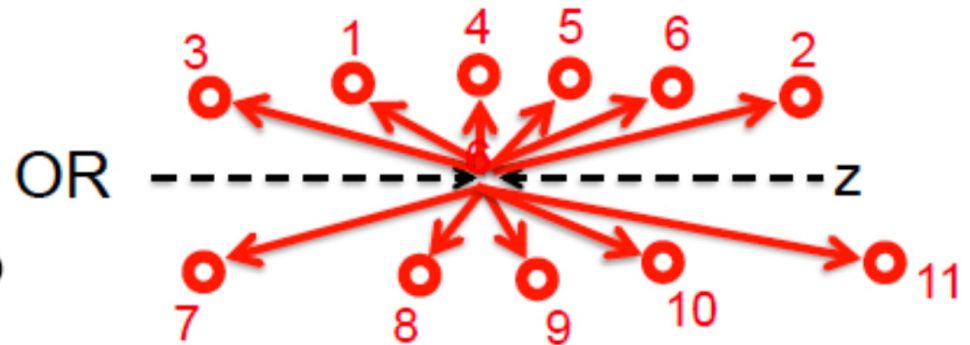
Coalescence of
quenched jet particle and bulk particle
can break the scaling

v_n from multi-particle correlation

Two-particle correlation



Multi-particle correlation



Multi-particle (>2) cumulants:

$$\langle\langle 6 \rangle\rangle = \langle\langle e^{in(\phi_1 + \phi_2 + \phi_3 - \phi_4 - \phi_5 - \phi_6)} \rangle\rangle$$

$$c_n\{6\} = \langle\langle 6 \rangle\rangle - 9 \cdot \langle\langle 4 \rangle\rangle \langle\langle 2 \rangle\rangle + 12 \cdot \langle\langle 2 \rangle\rangle^3$$

$$v_n\{4\} = \sqrt[4]{-c_n\{4\}}$$

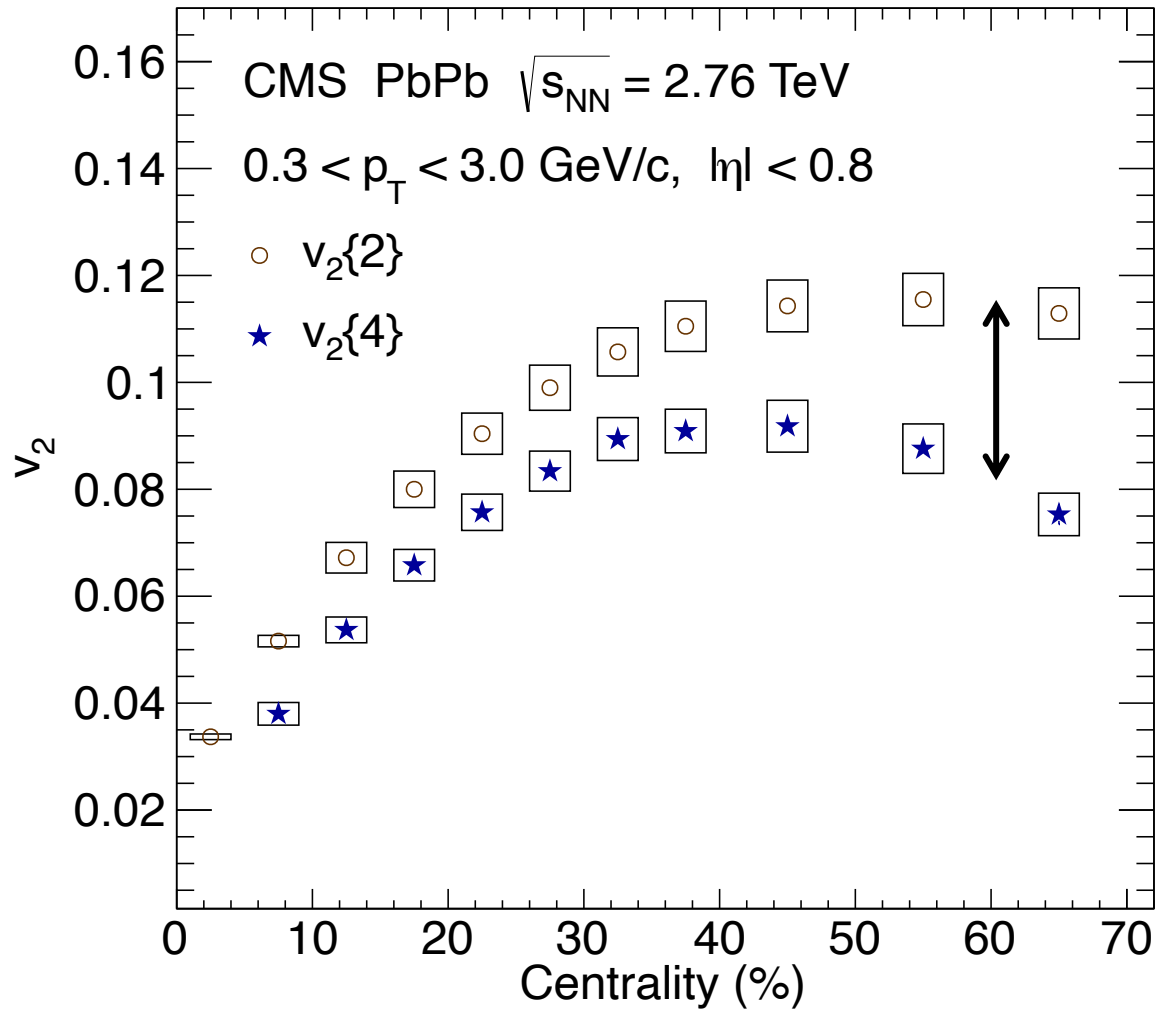
$$v_n\{6\} = \sqrt[4]{\frac{1}{4}c_n\{6\}}$$

Q-cumulant, PRC 83 (2011) 044913

$v_n\{4\} \approx v_n\{6\} \approx \dots \rightarrow$ system is collective

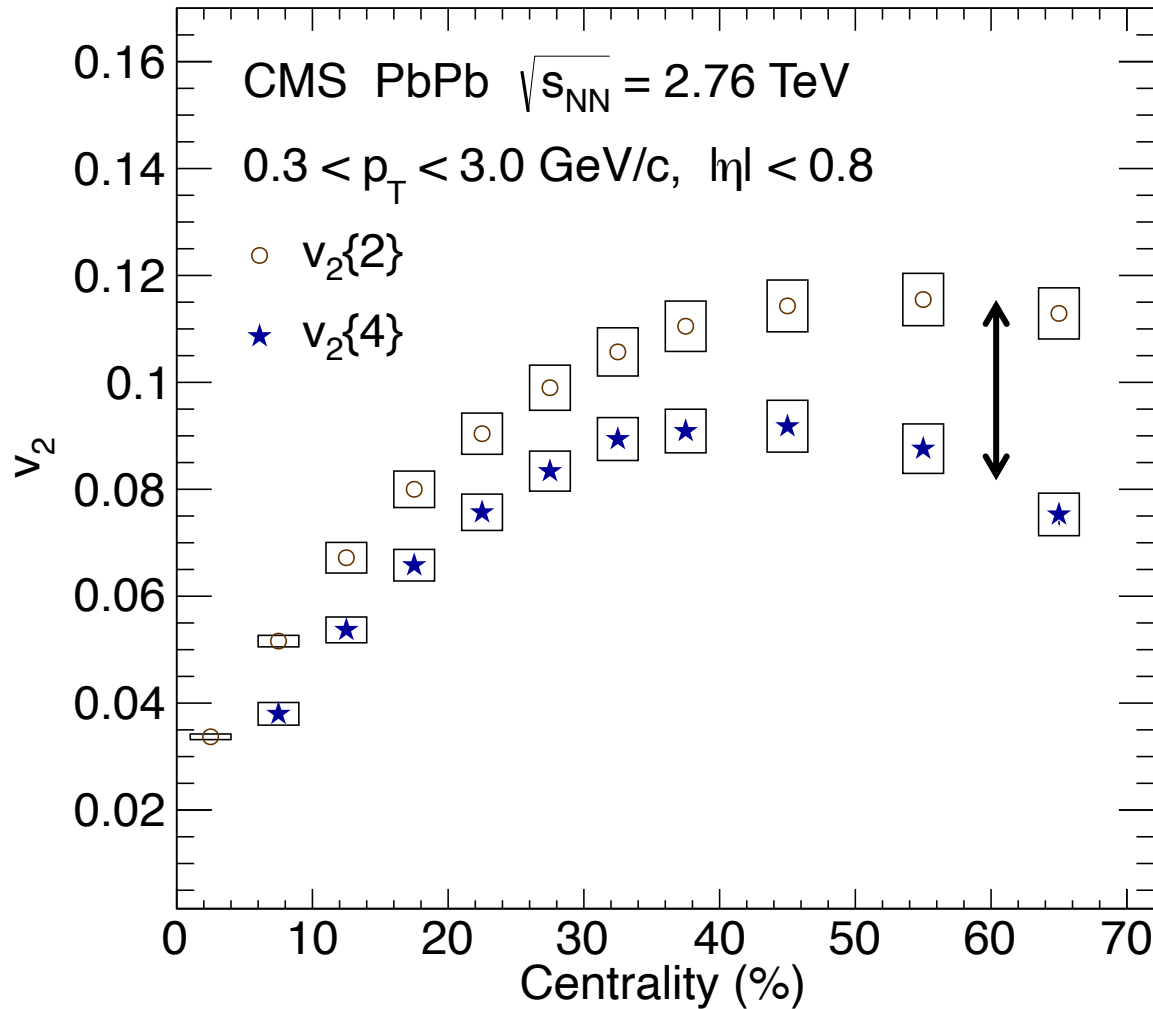
v_n from multi-particle correlation

PRC 87 (2013) 014902

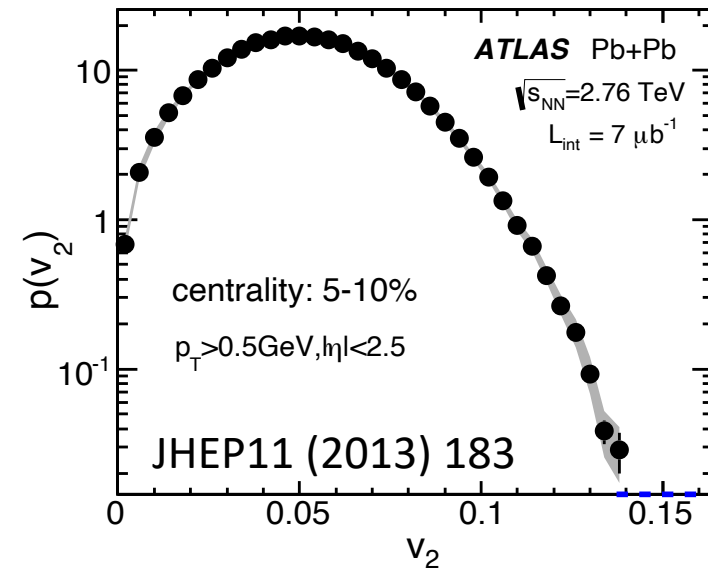


v_n from multi-particle – flow fluctuation

PRC 87 (2013) 014902

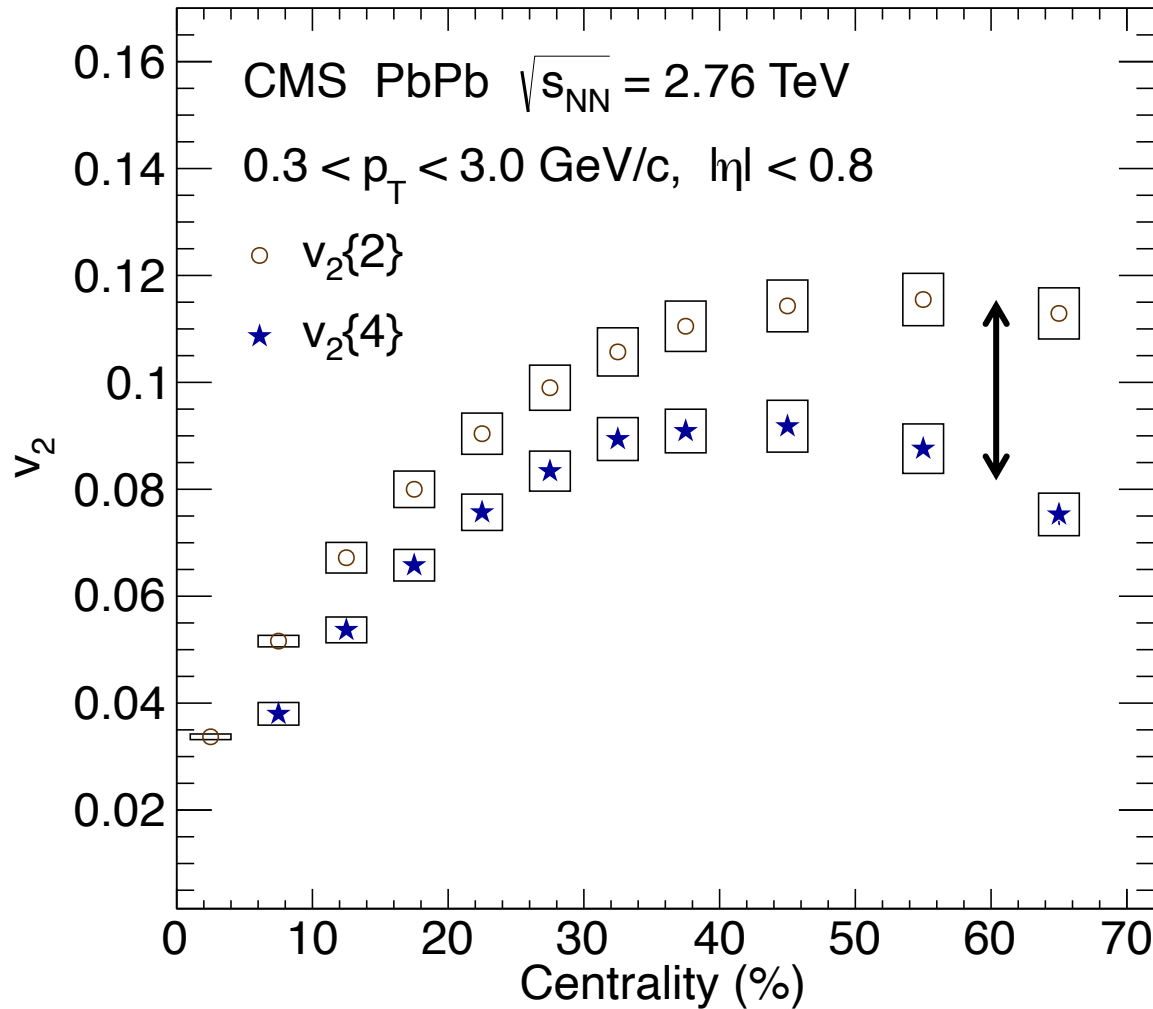


Event-by-Event Flow Fluctuation



v_n from multi-particle – flow fluctuation

PRC 87 (2013) 014902



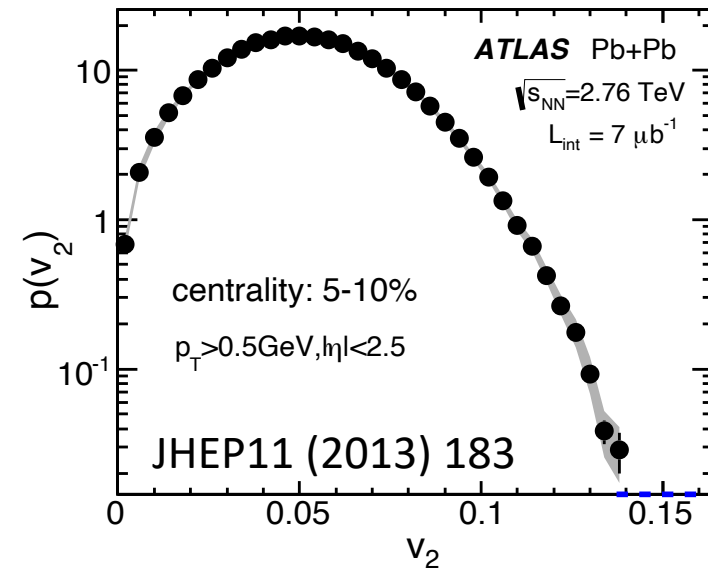
Event-by-Event Flow Fluctuation

$$v_n\{2\}^2 = \langle v_n^2 \rangle = \langle v_n \rangle^2 + \sigma_{v_n}^2$$

$$v_n\{4\}^2 = (2\langle v_n^2 \rangle^2 - \langle v_n^4 \rangle)^{\frac{1}{2}}$$

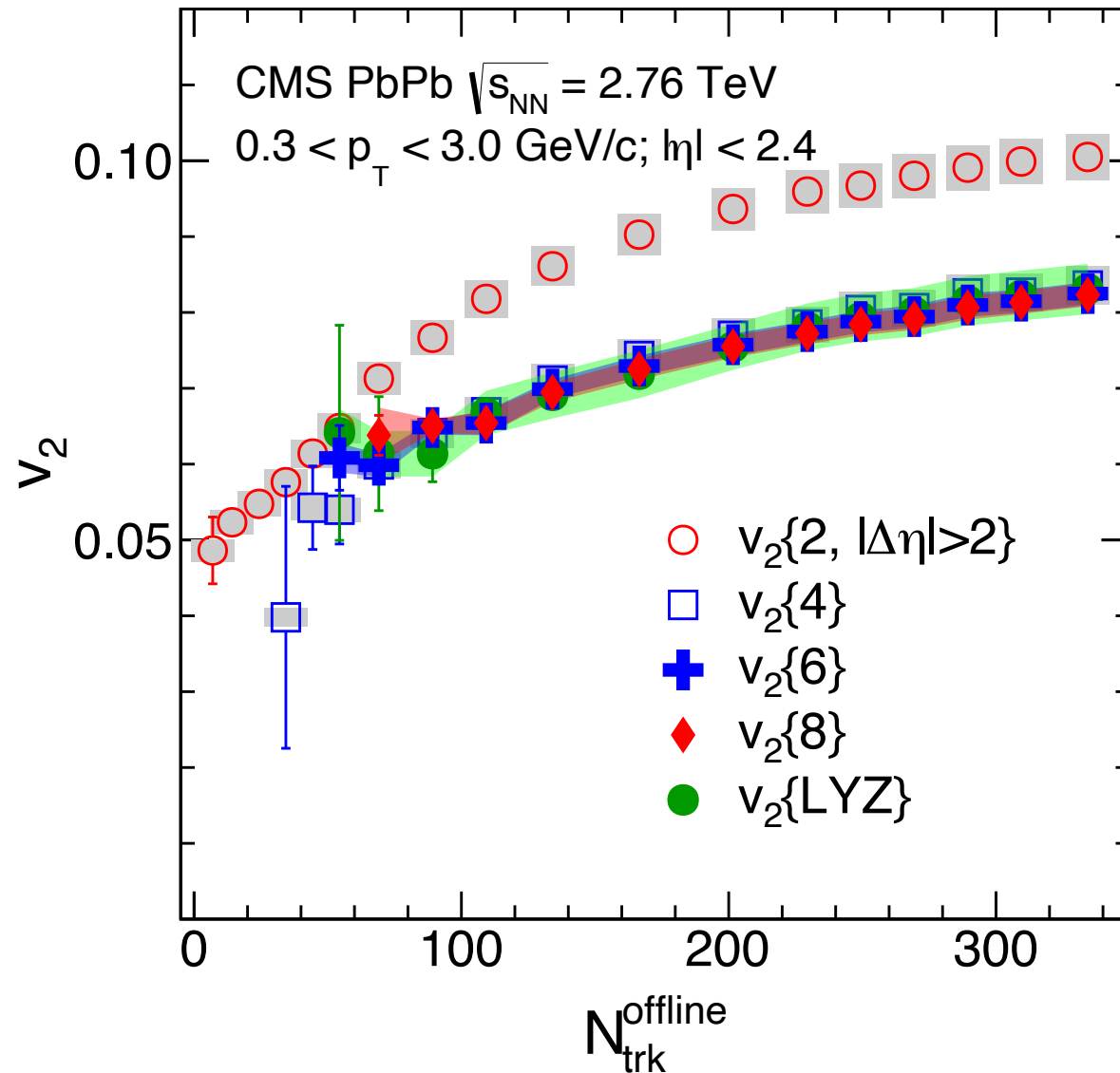
$$\approx \langle v_n \rangle^2 - \sigma_{v_n}^2$$

PRC 80 (2009) 014904



v_n from multi-particle – flow fluctuation

PRL 115 (2015) 012301



Event-by-Event
Flow Fluctuation

$$v_n\{2\}^2 = \langle v_n^2 \rangle = \langle v_n \rangle^2 + \sigma_{v_n}^2$$

$$v_n\{4\}^2 = (2\langle v_n^2 \rangle^2 - \langle v_n^4 \rangle)^{\frac{1}{2}}$$

$$\approx \langle v_n \rangle^2 - \sigma_{v_n}^2$$

PRC 80 (2009) 014904

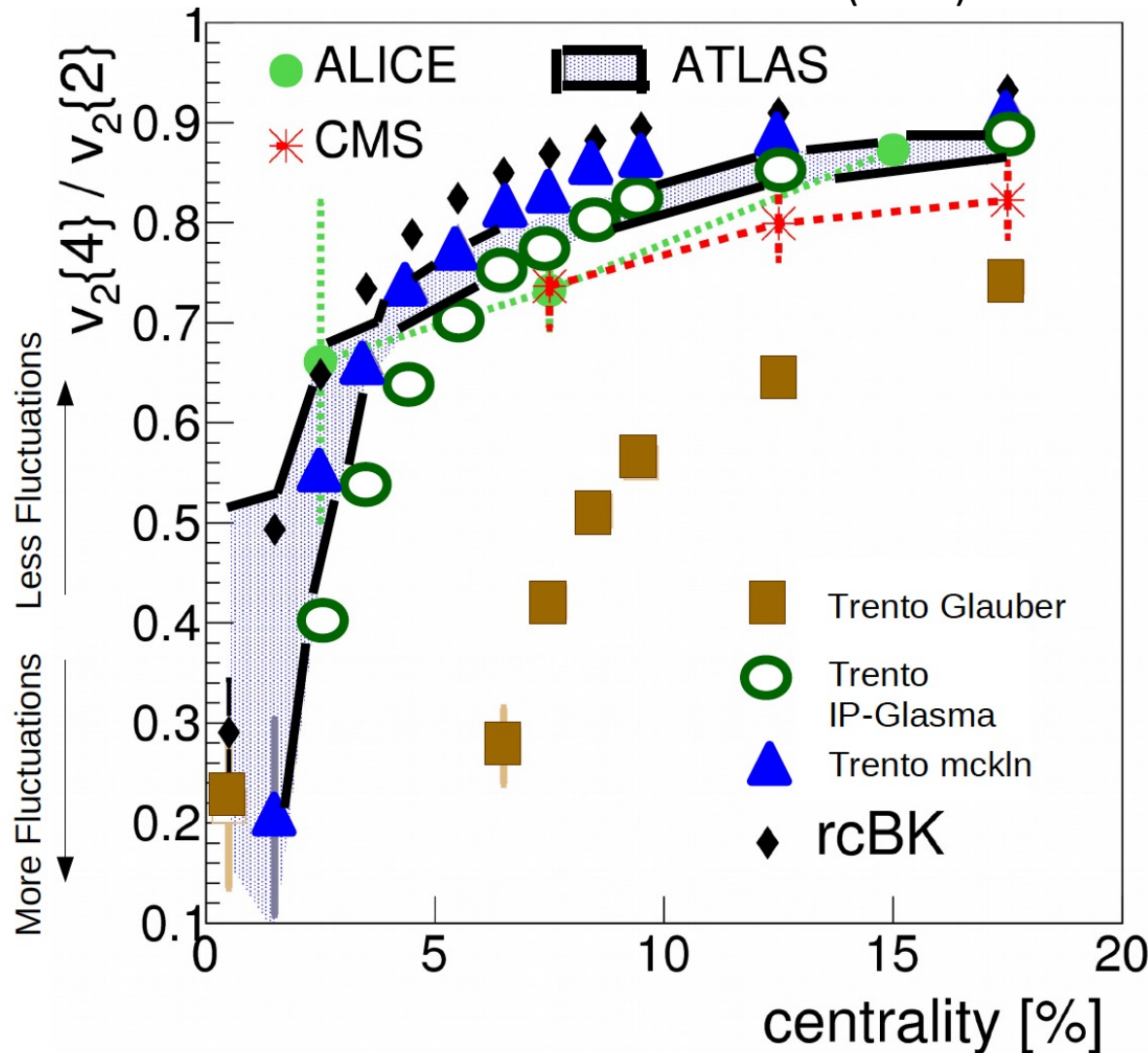
$v_2\{4,6,8,\text{LYZ}\}$

indicate multi-particle
nature of correlation

More implications
on flow fluctuation
see EPJC 74 (2014) 3157

v_n from multi-particle – constrain initial-state

PRC 95 (2017) 054910



Event-by-Event
Flow Fluctuation

$$v_n\{2\}^2 = \langle v_n^2 \rangle = \langle v_n \rangle^2 + \sigma_{v_n}^2$$

$$v_n\{4\}^2 = (2\langle v_n^2 \rangle^2 - \langle v_n^4 \rangle)^{\frac{1}{2}}$$

$$\approx \langle v_n \rangle^2 - \sigma_{v_n}^2$$

PRC 80 (2009) 014904

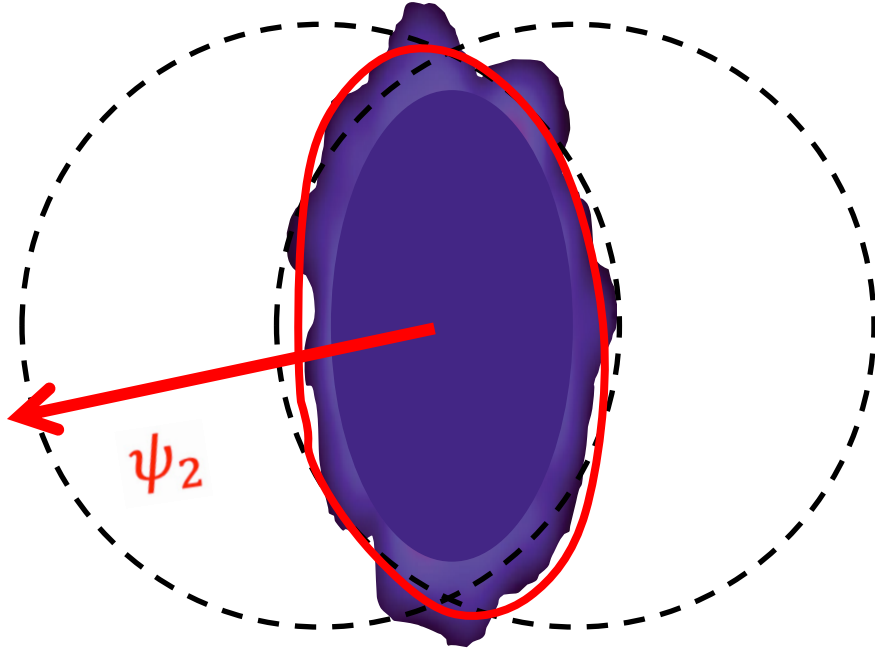
Strong constraints on
initial state models

MC Glauber fails!

TRENTO: PRC 92(2015) 011901

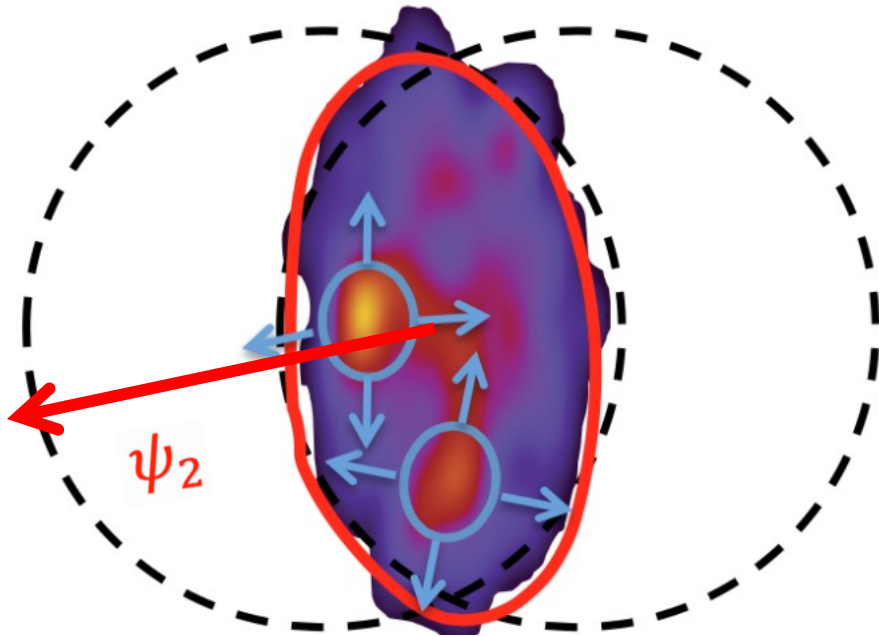
rcBK: arXiv:1011.5161

Decorrelation – Transverse



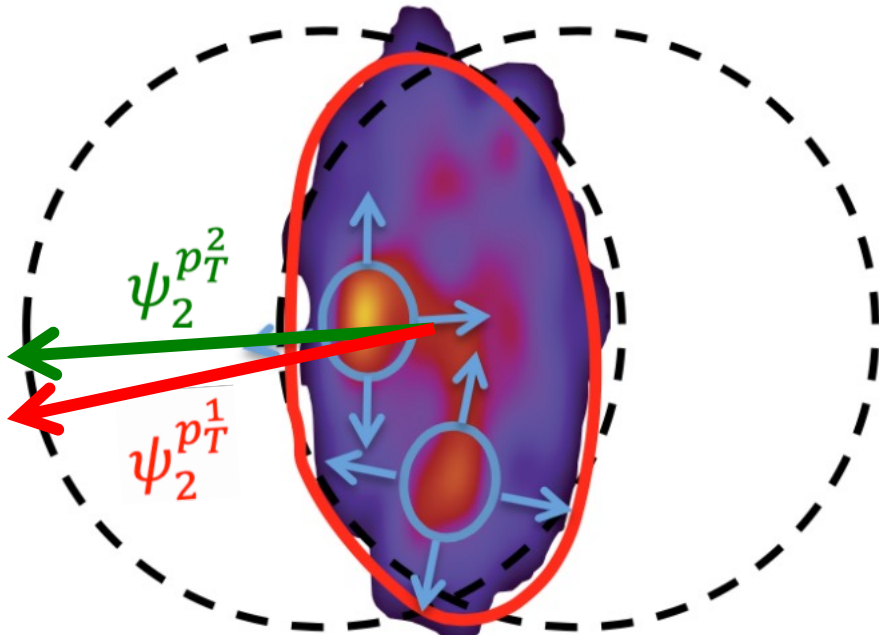
All particles share the same event plane

Decorrelation – Transverse



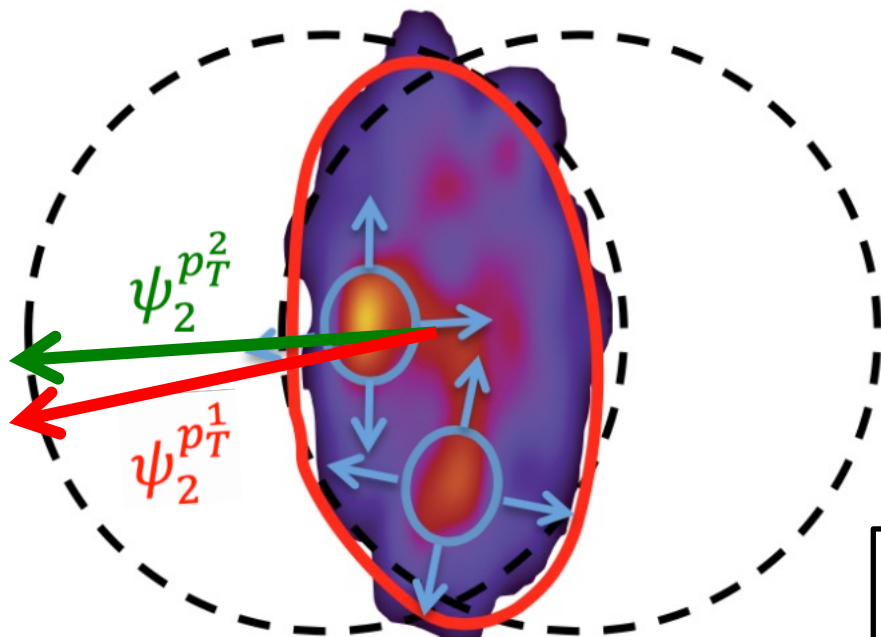
Local hot spots push particles
to higher p_T

Decorrelation – Transverse



Local hot spots push particles
to higher p_T
AND
result in different event planes
for different p_T range

Decorrelation – Transverse



Local hot spots push particles
to higher p_T
AND
result in different event planes
for different p_T range

Factorization breaking

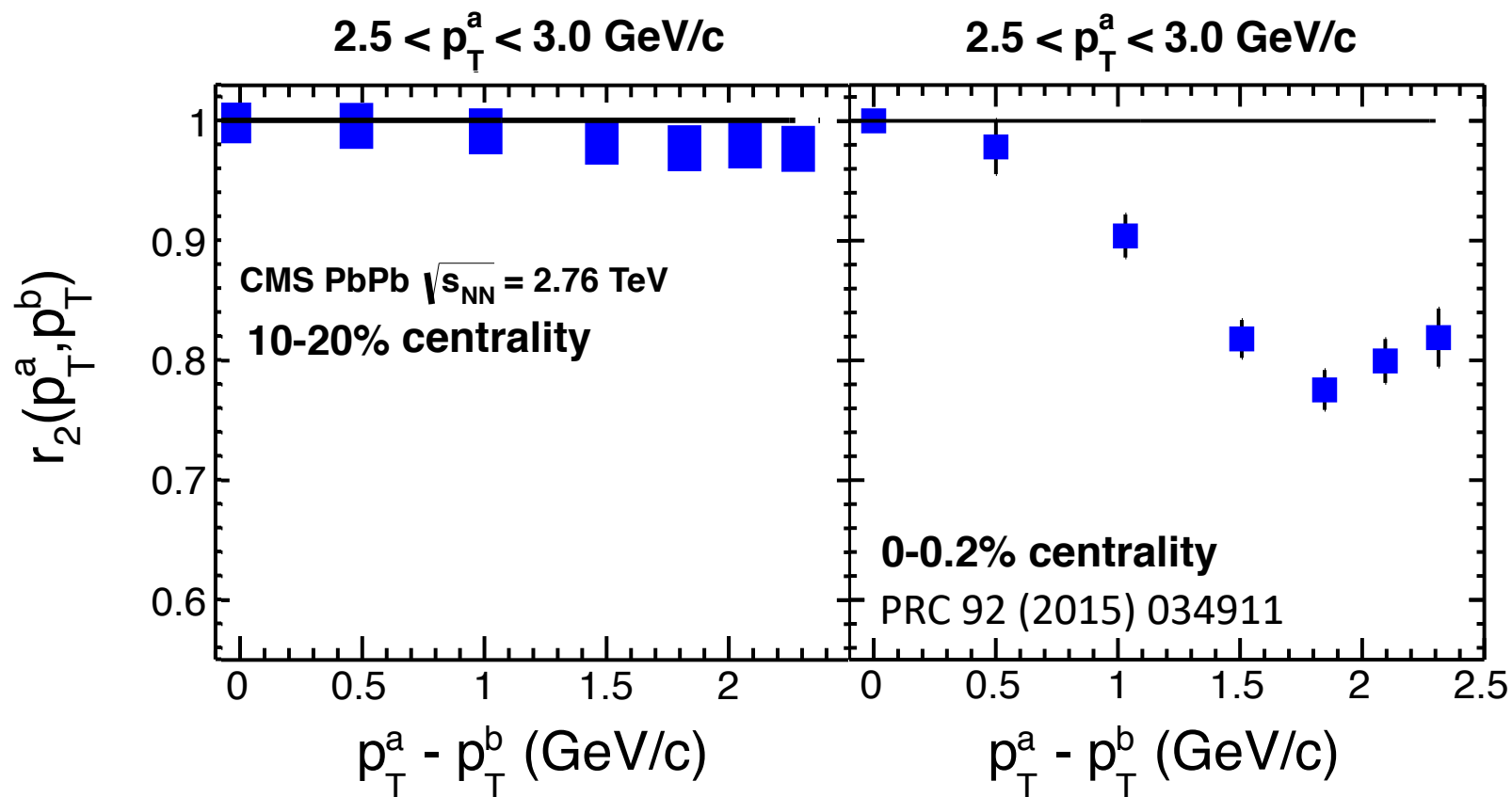
$$V_{n\Delta}(p_T^a, p_T^b) \neq v_n(p_T^a) \times v_n(p_T^b)$$

Effects measurable by

$$r_n(p_T^a, p_T^b) \equiv \frac{V_{n\Delta}(p_T^a, p_T^b)}{\sqrt{V_{n\Delta}(p_T^a, p_T^a)} \sqrt{V_{n\Delta}(p_T^b, p_T^b)}} \sim \langle \cos [n(\Psi_n(p_T^a) - \Psi_n(p_T^b))] \rangle$$

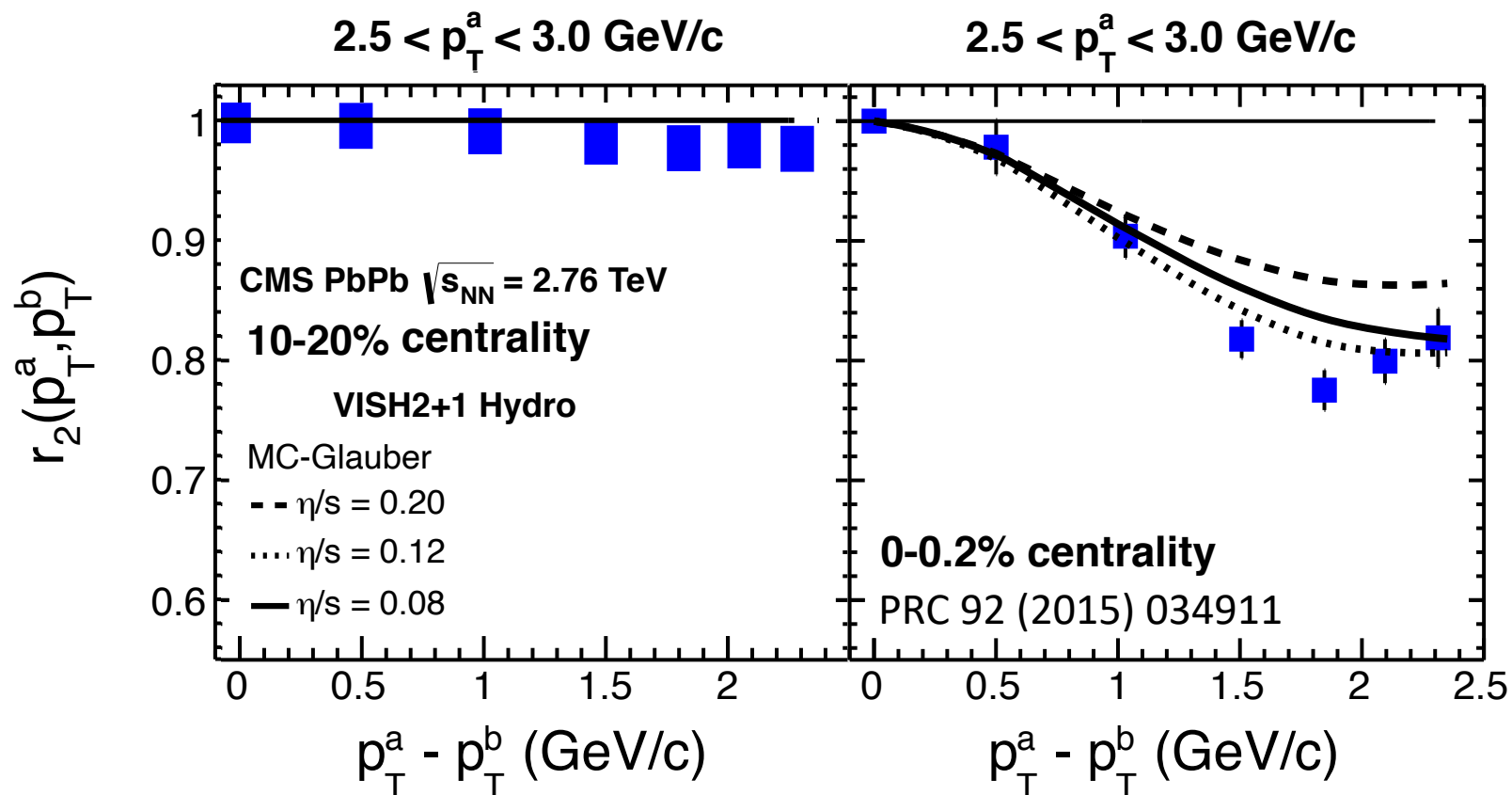
PRC 92 (2015) 034911

Decorrelation – p_T



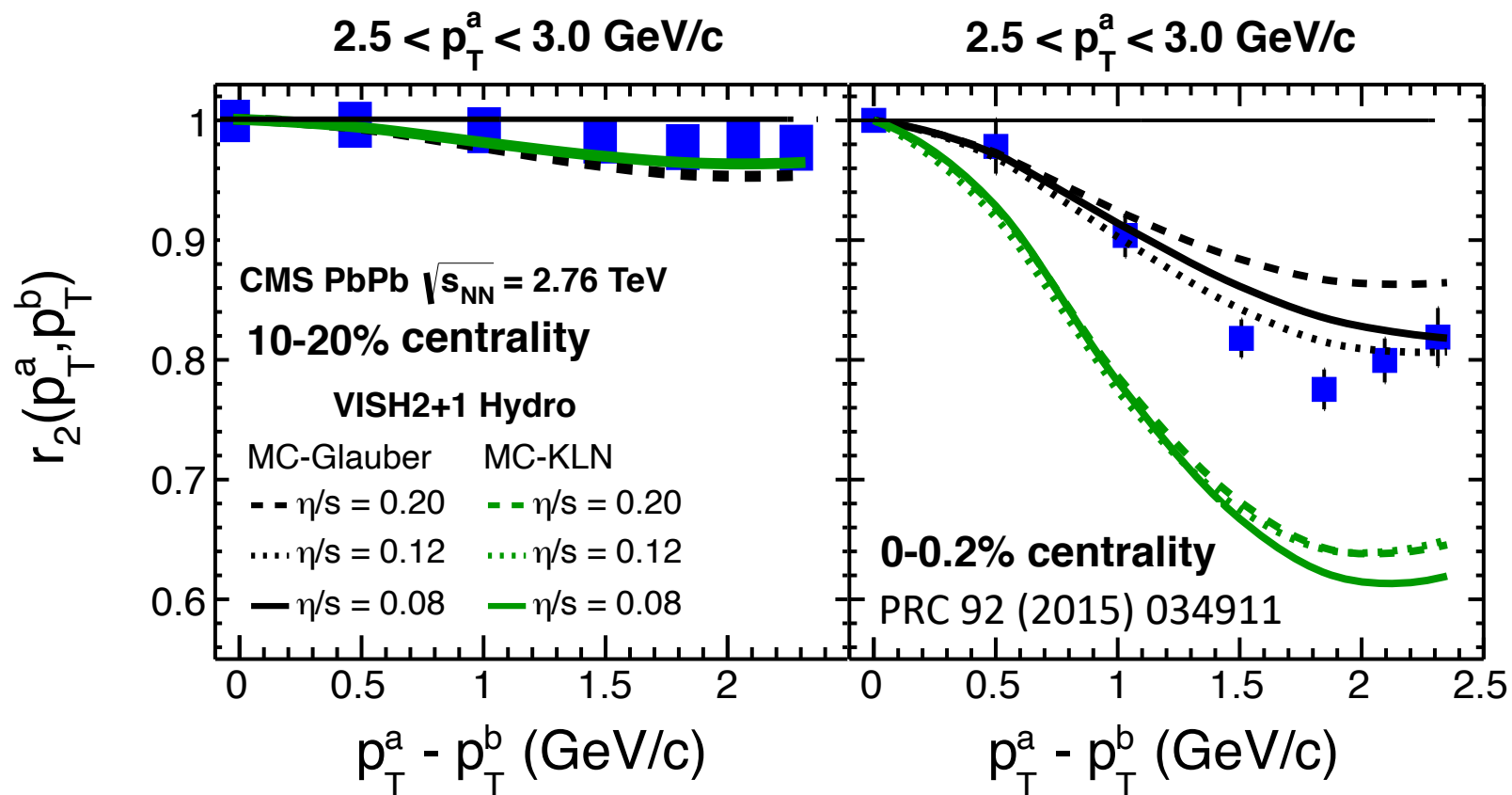
Most prominent at central collisions where fluctuation dominates

Decorrelation – p_T



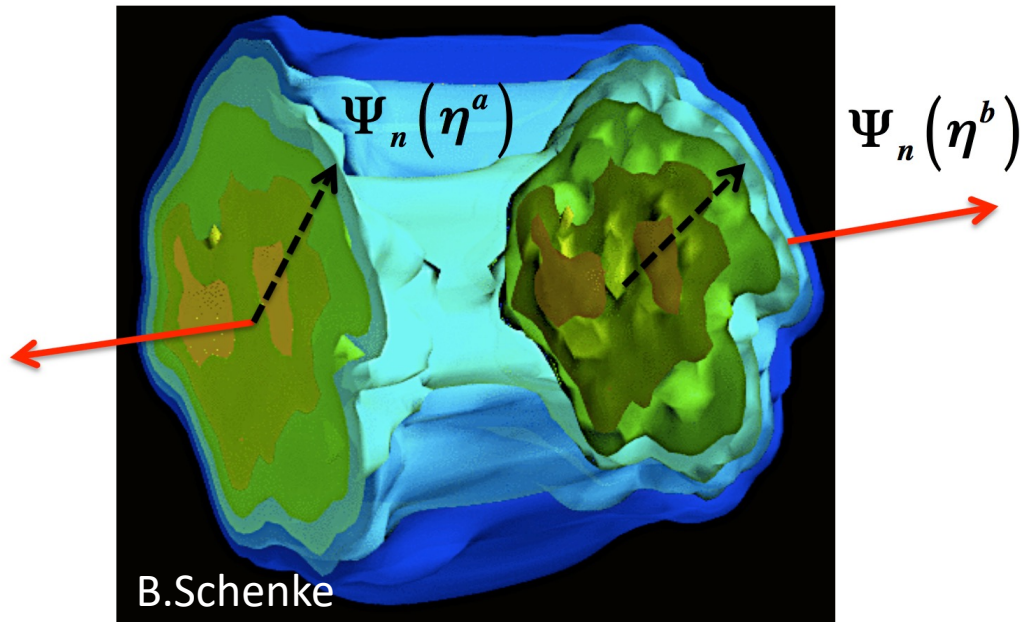
Most prominent at central collisions where fluctuation dominates
Some dependence on η/s

Decorrelation – p_T

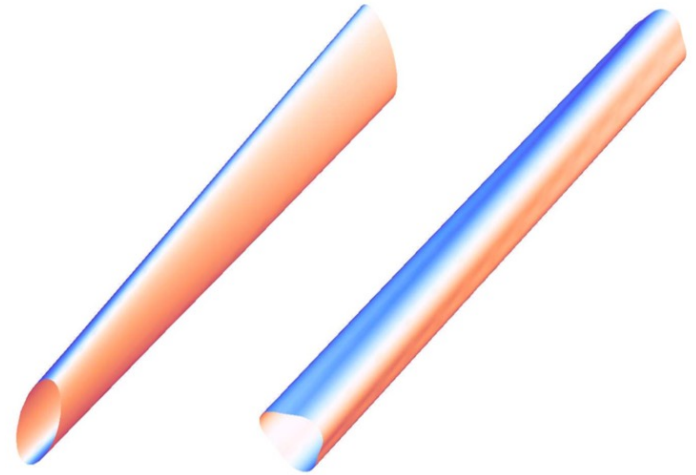


Most prominent at central collisions where fluctuation dominates
 Some dependence on η/s
 Very sensitive to initial state geometry

Decorrelation – Longitudinal



Torqued fireball



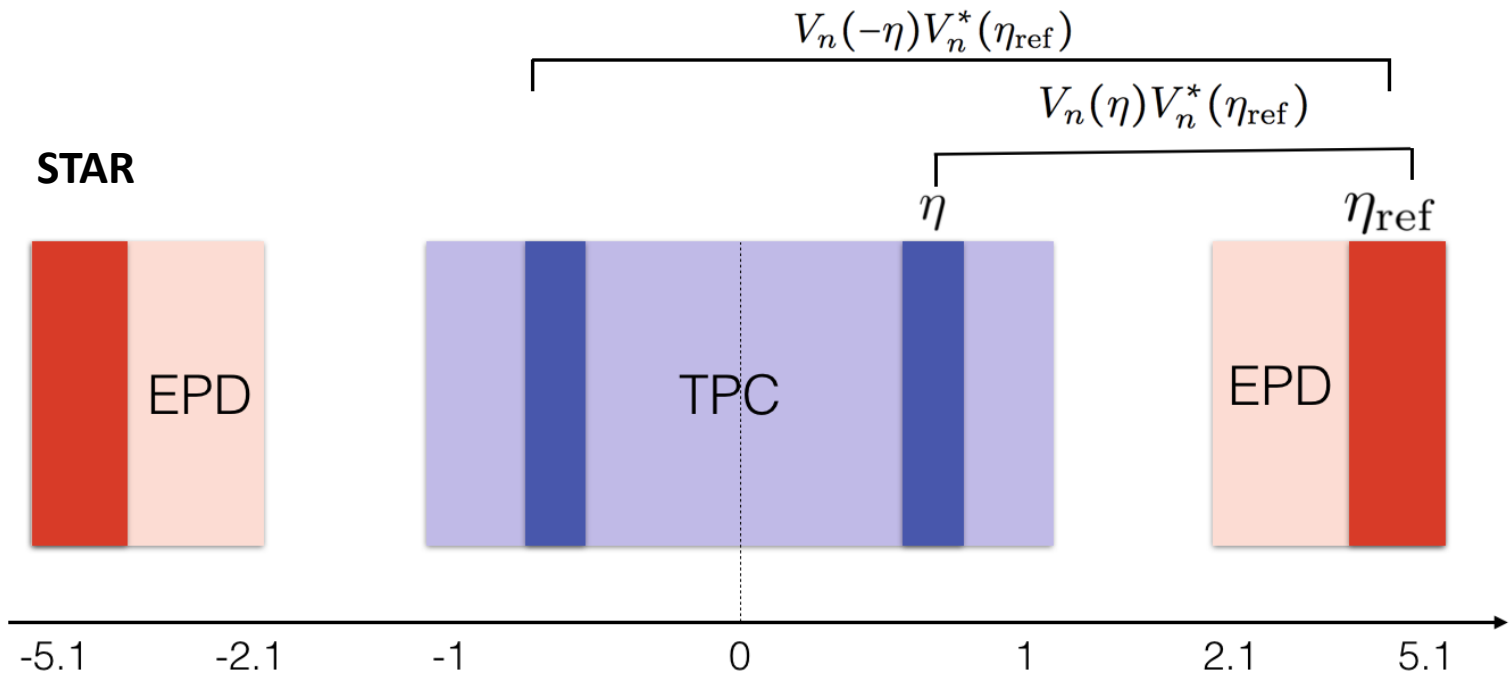
PRC 83 (2011) 034911

Event Plane decorrelation along longitudinal direction
Provide constraints on how system evolves in 3D

Decorrelation – Longitudinal

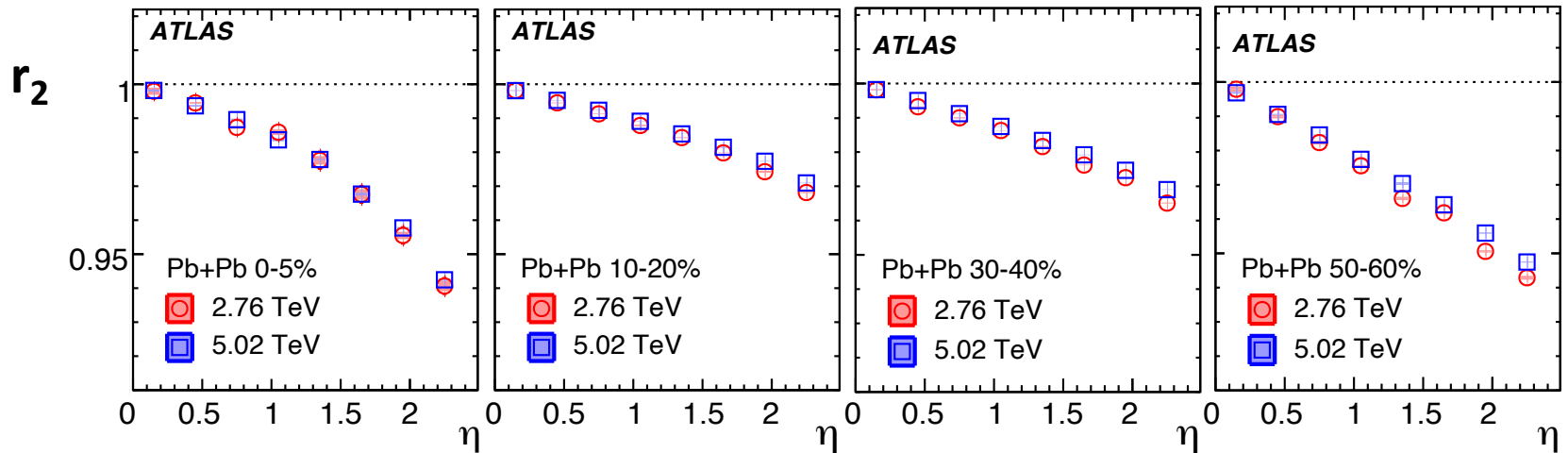
Effects measurable by

$$\begin{aligned}
 r_n(\eta) &= \frac{\langle V_n(-\eta)V_n^*(\eta_{\text{ref}}) \rangle}{\langle V_n(\eta)V_n^*(\eta_{\text{ref}}) \rangle} \\
 &= \frac{\langle v_n(-\eta)v_n(\eta_{\text{ref}}) \cos n(\Psi_n(-\eta) - \Psi_n(\eta_{\text{ref}})) \rangle}{\langle v_n(\eta)v_n(\eta_{\text{ref}}) \cos n(\Psi_n(\eta) - \Psi_n(\eta_{\text{ref}})) \rangle}
 \end{aligned}$$



Decorrelation – Longitudinal

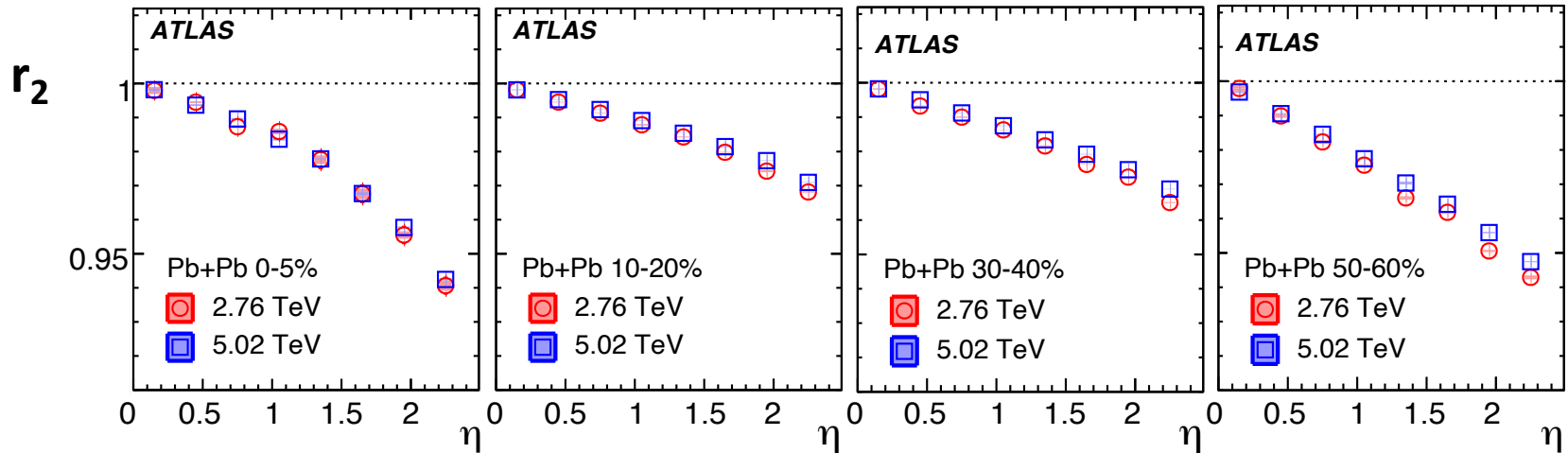
EPJC 76 (2018) 142



r_2 : Centrality dependence; Initial geometry + fluctuation

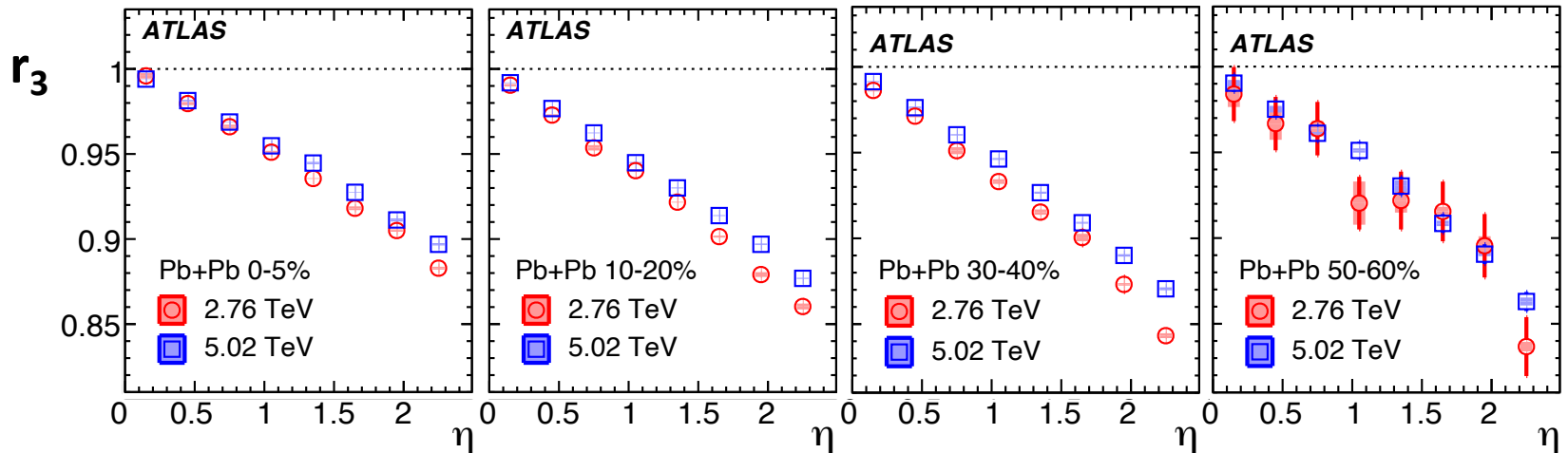
Decorrelation – Longitudinal

EPJC 76 (2018) 142

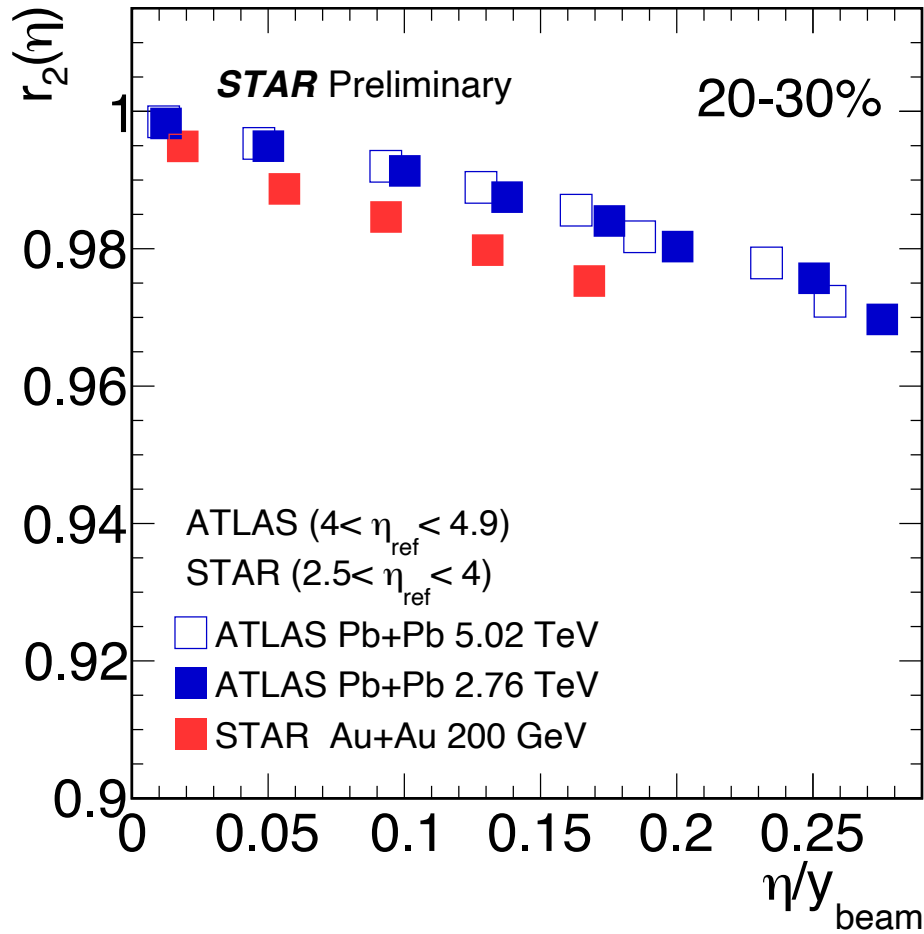


r_2 : Centrality dependence; Initial geometry + fluctuation

r_3 : No centrality dependence; Pure fluctuation

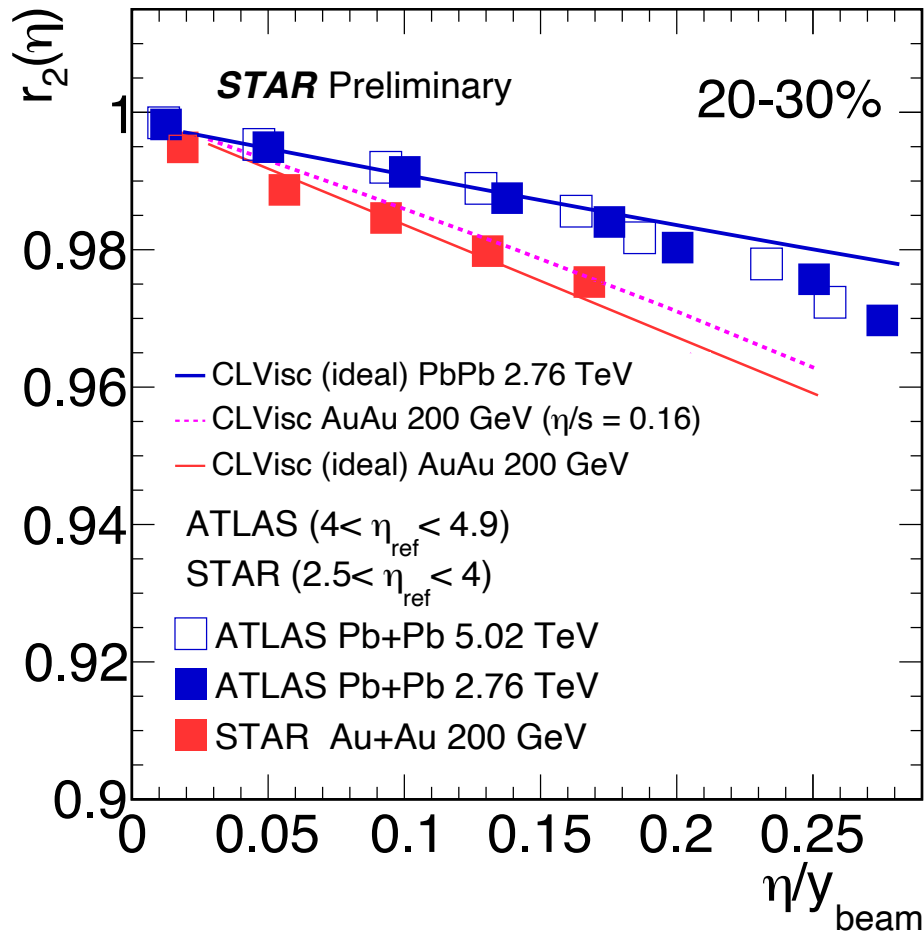


Decorrelation – Longitudinal



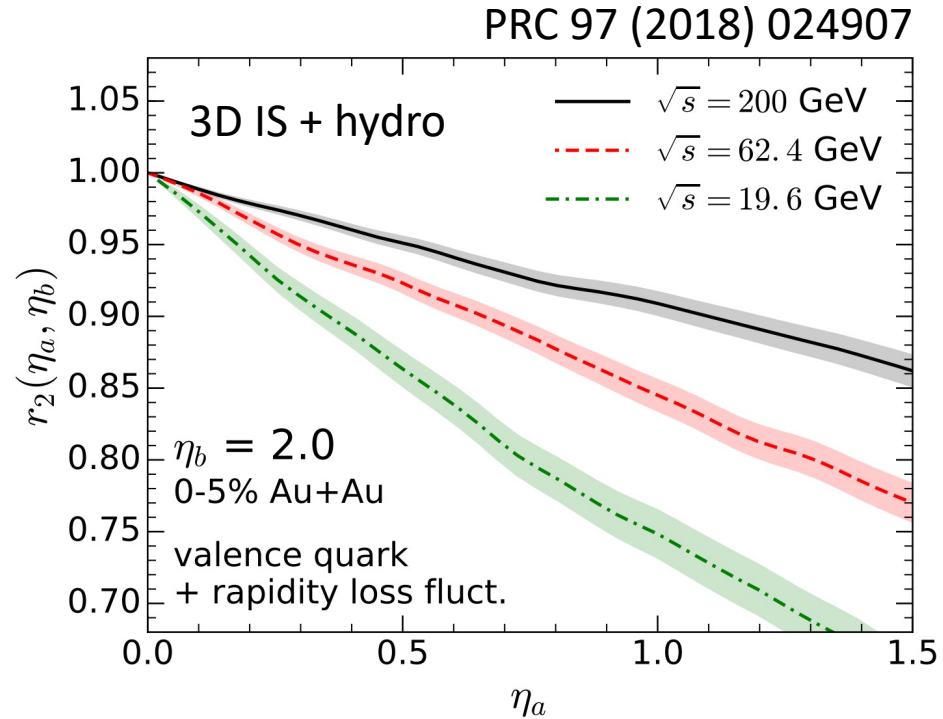
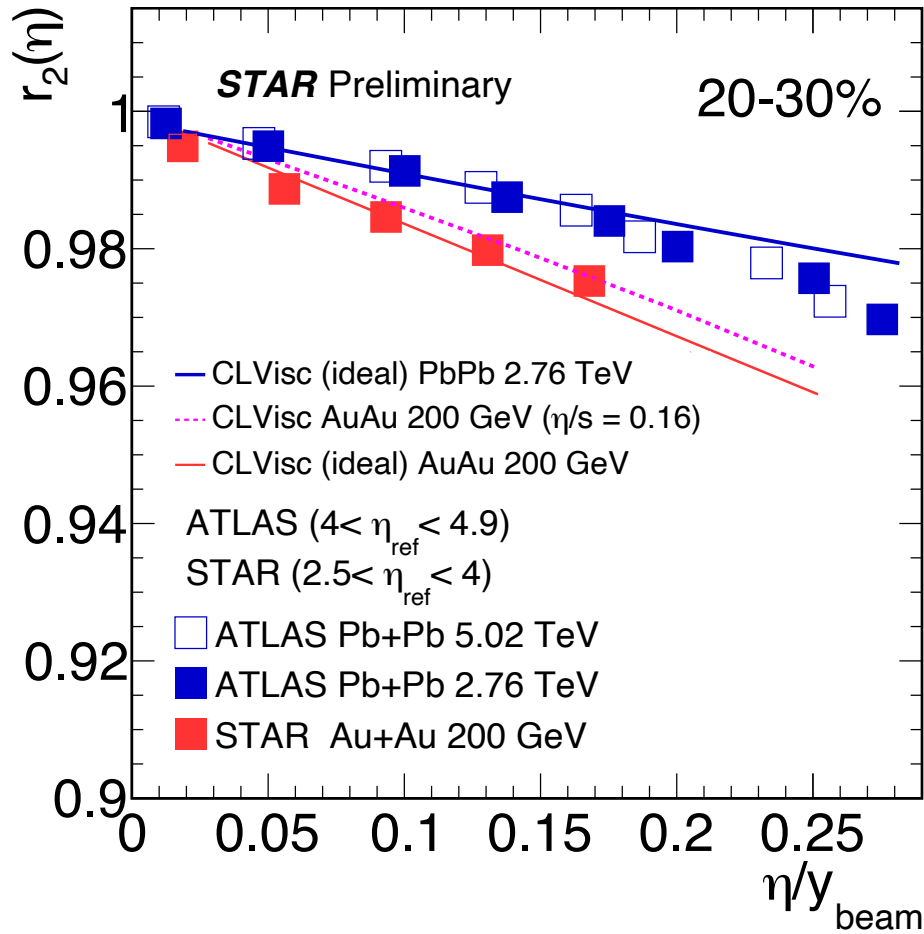
Larger effect seen at RHIC than LHC

Decorrelation – Longitudinal



Larger effect seen at RHIC than LHC

Decorrelation – Longitudinal



Larger effect seen at RHIC than LHC

3D IS + Hydro predicts larger decorrelation at lower energies

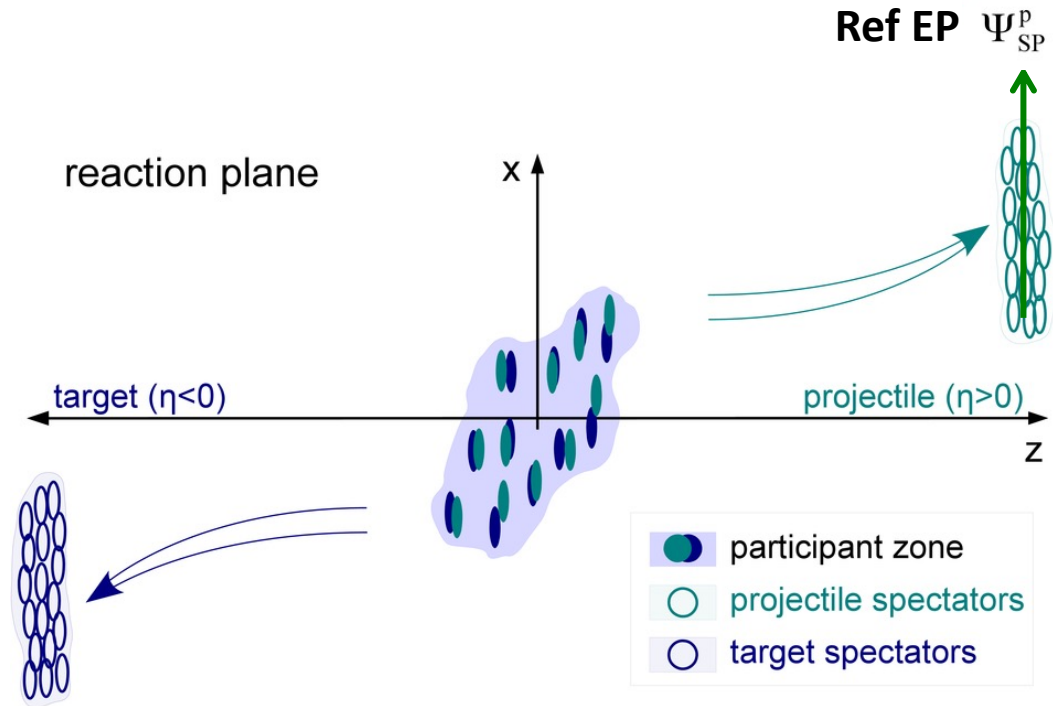
Directed flow v_1

$$v_1 = \langle \cos(\phi - \psi_{sp}) \rangle$$

Preferred moving direction
at different rapidity

Collision symmetry

$$v_1(\eta) = -v_1(-\eta)$$

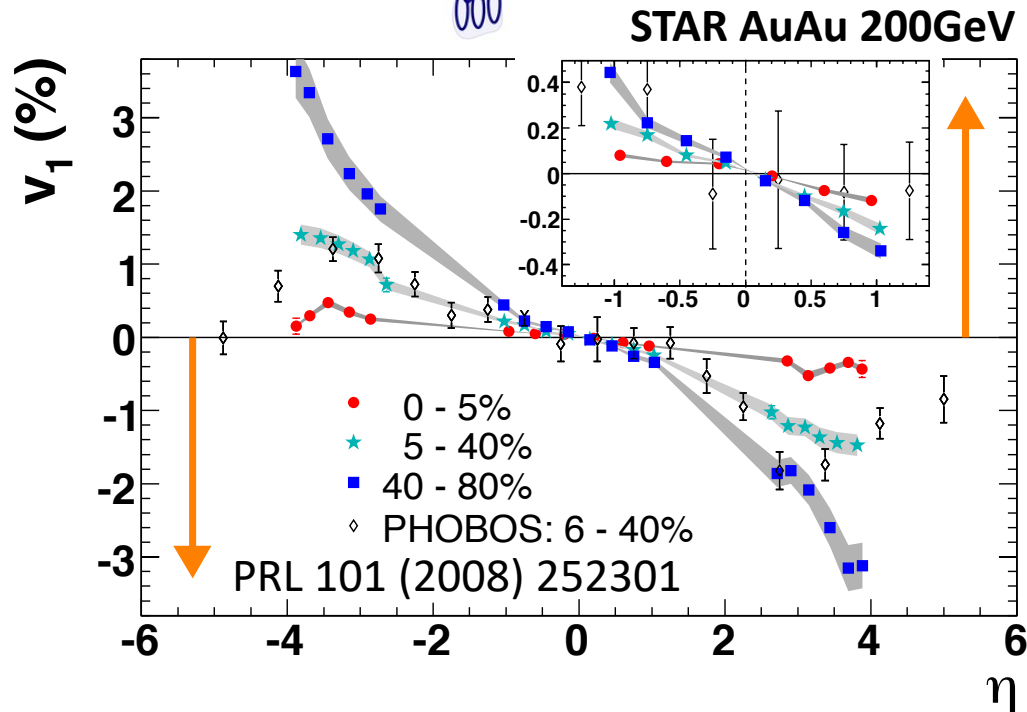
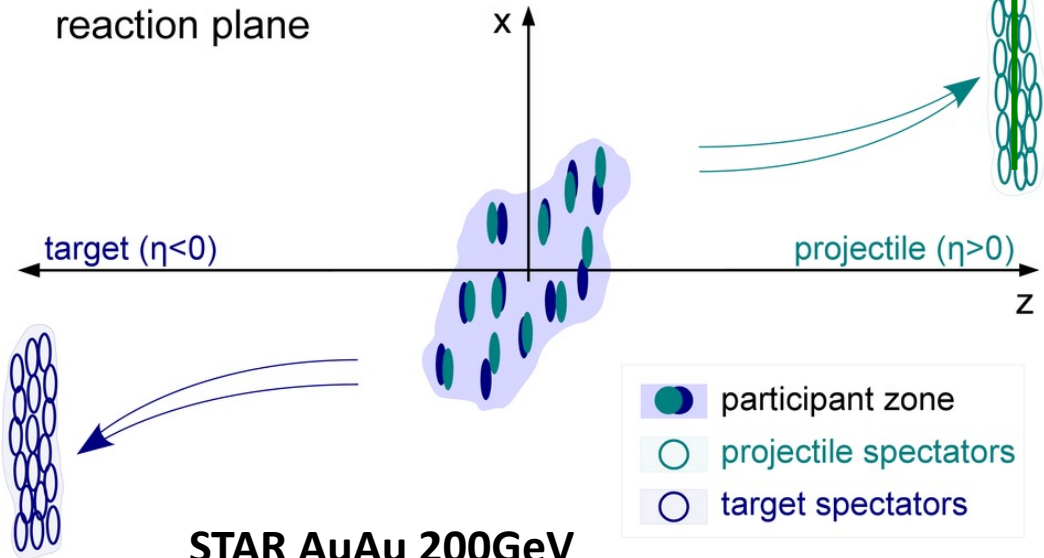


Directed flow v_1

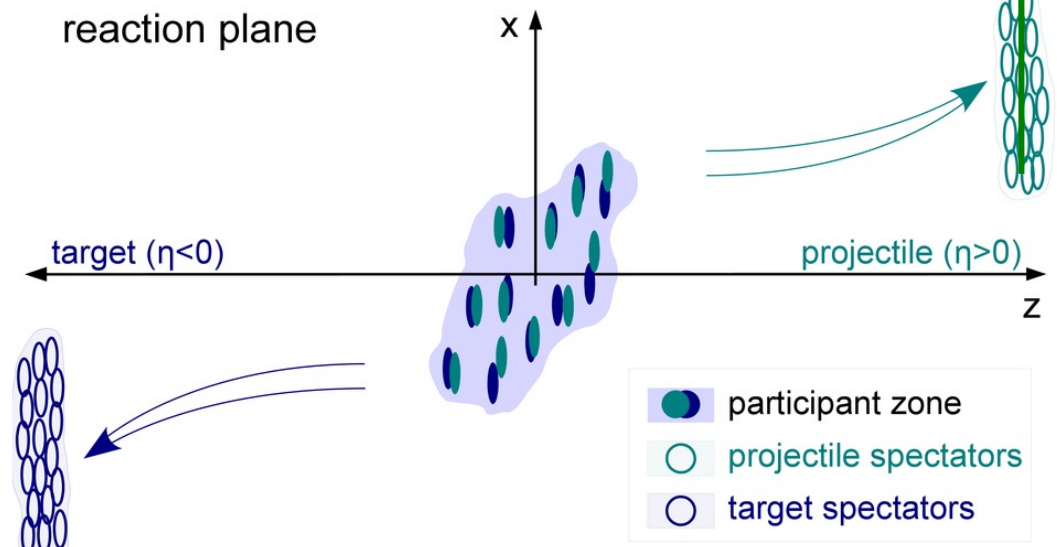
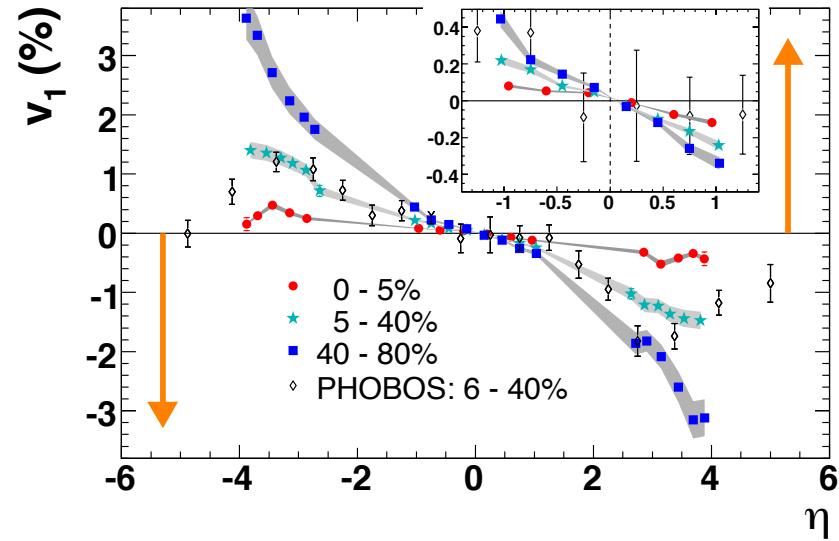
$$v_1 = \langle \cos(\phi - \psi_{sp}) \rangle$$

Preferred moving direction
at different rapidity

Collision symmetry
 $v_1(\eta) = -v_1(-\eta)$



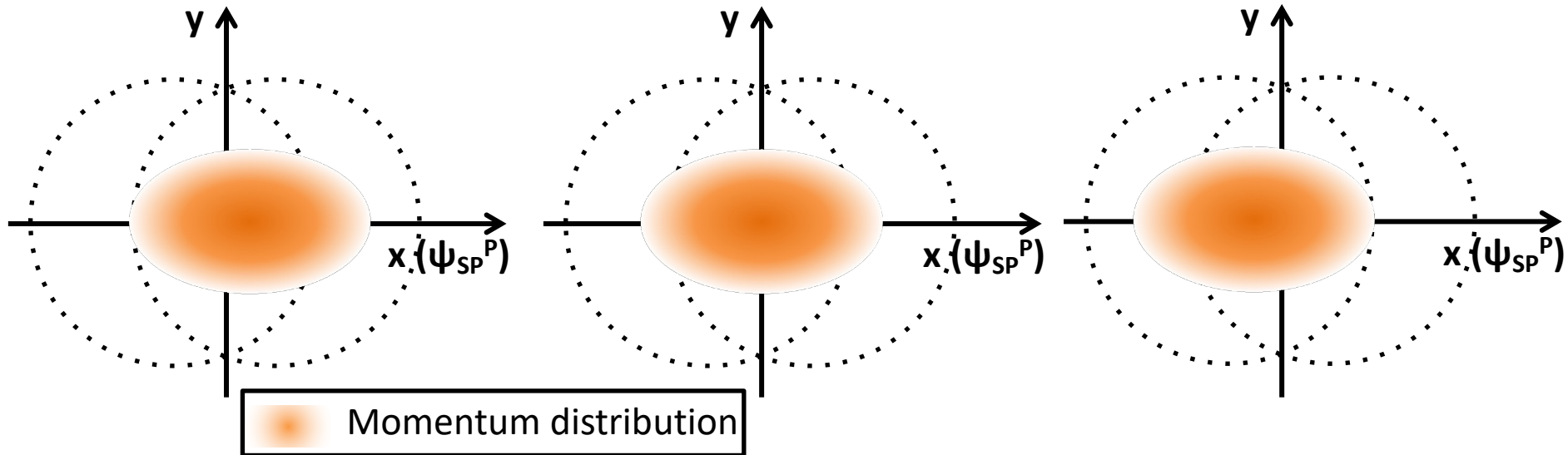
Directed flow v_1



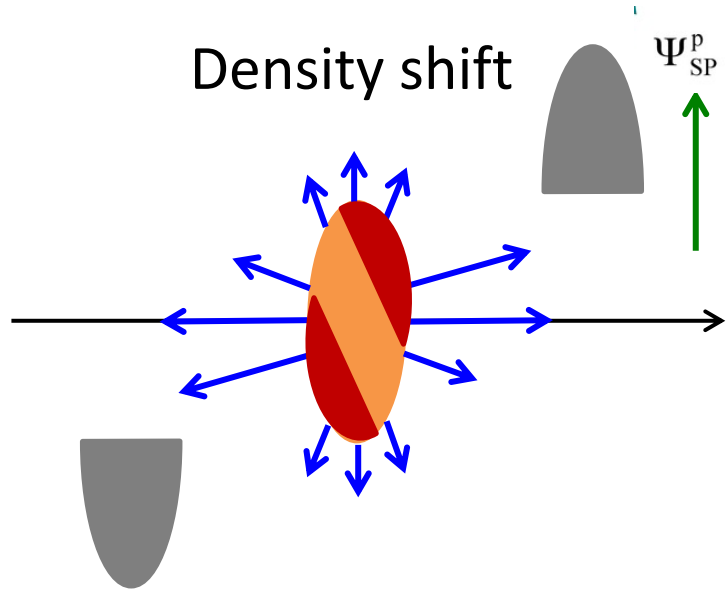
$\eta < 0, v_1 > 0$

$\eta = 0, v_1 = 0$

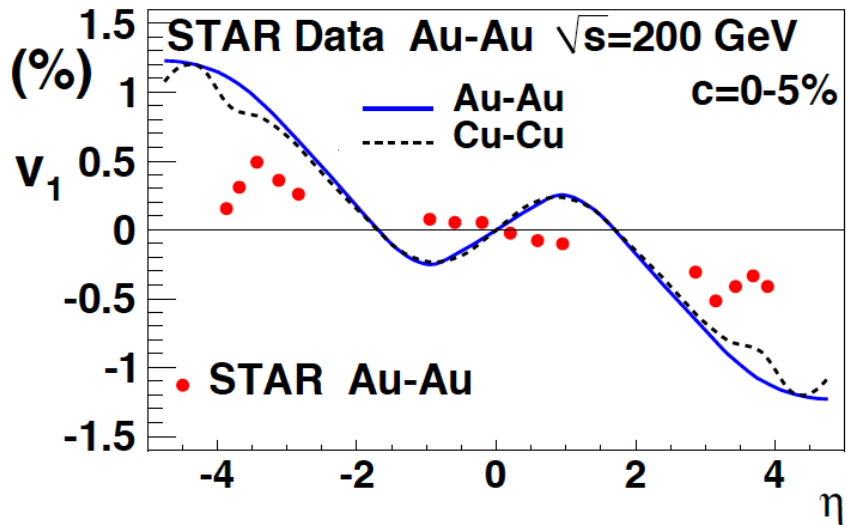
$\eta > 0, v_1 < 0$



Directed flow v_1

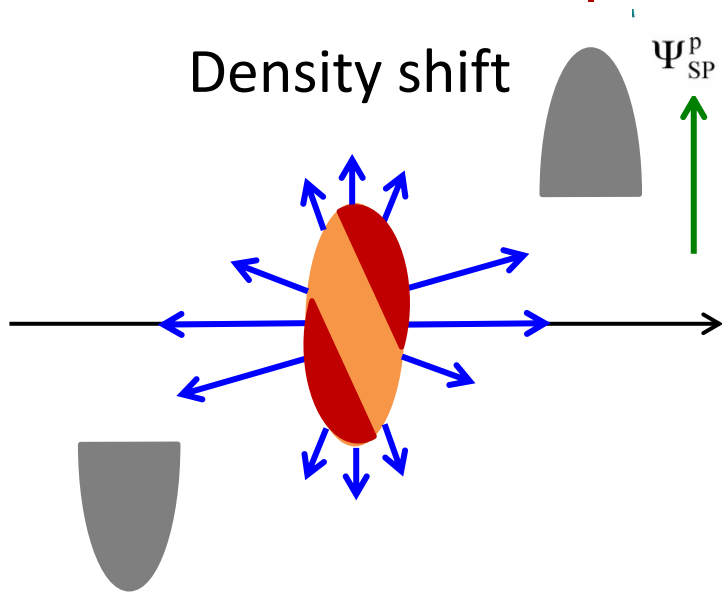


PRC 81 (2010) 054902

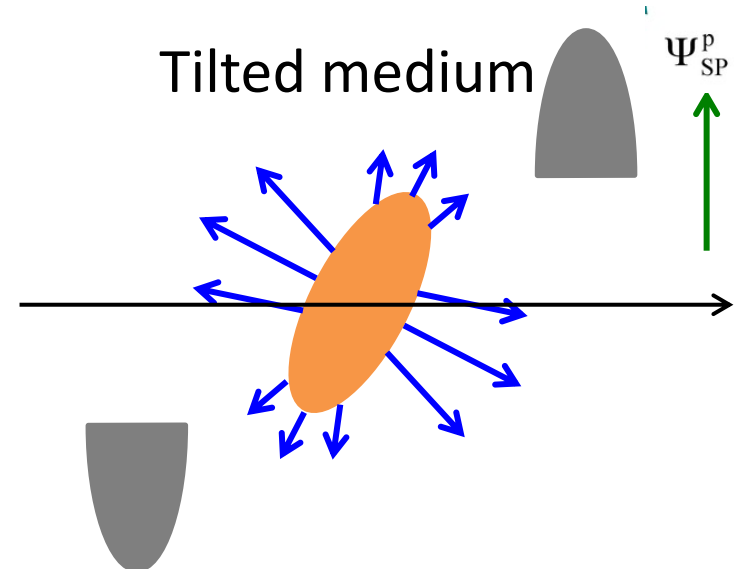


Directed flow v_1

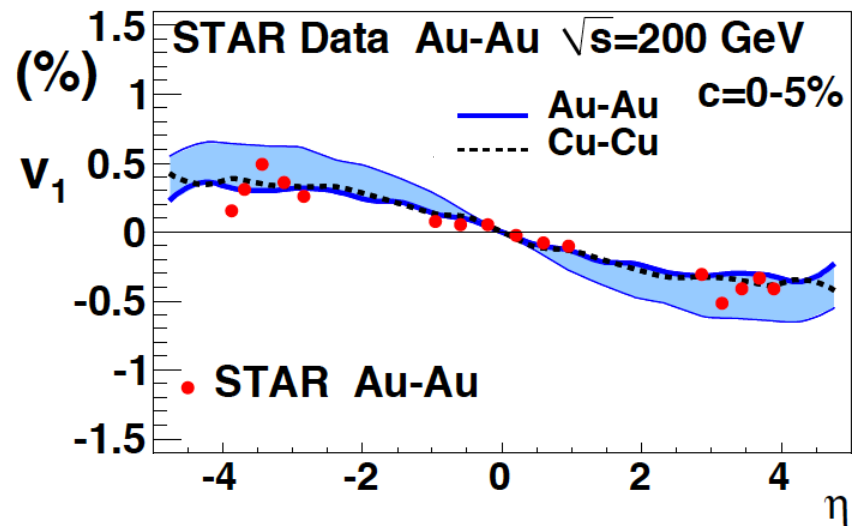
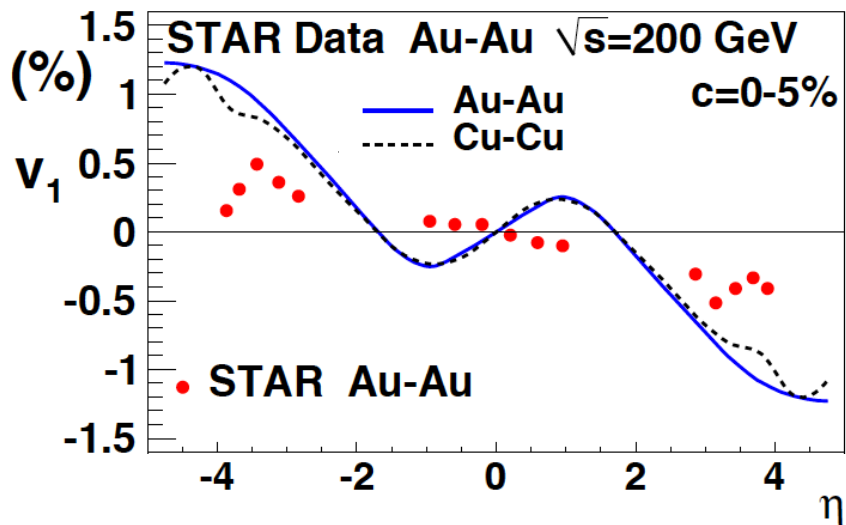
Density shift



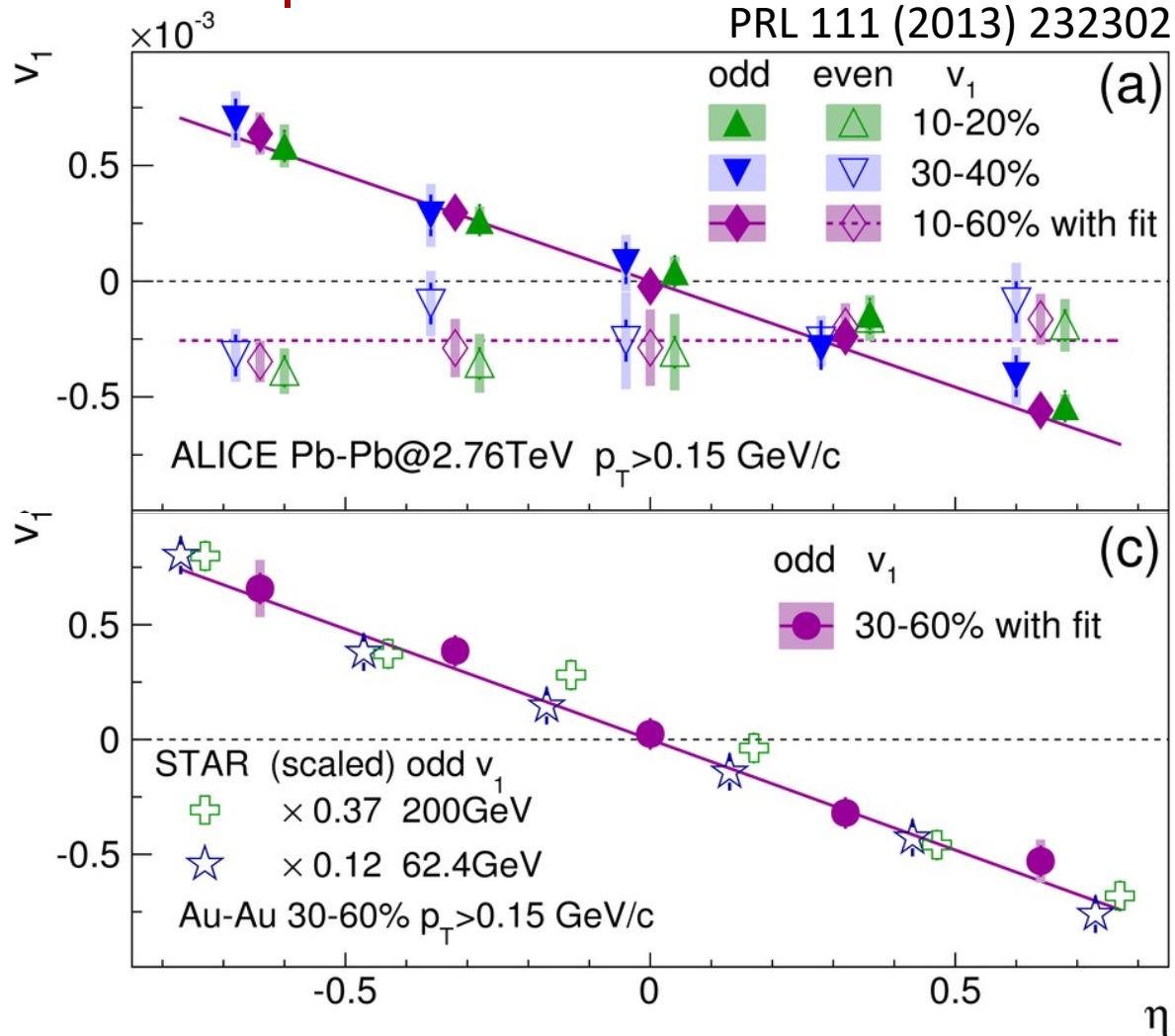
Tilted medium



PRC 81 (2010) 054902



Directed flow v_1



Larger slope of v_1^{odd} at RHIC energies: smaller tilt at LHC

Directed flow v_1

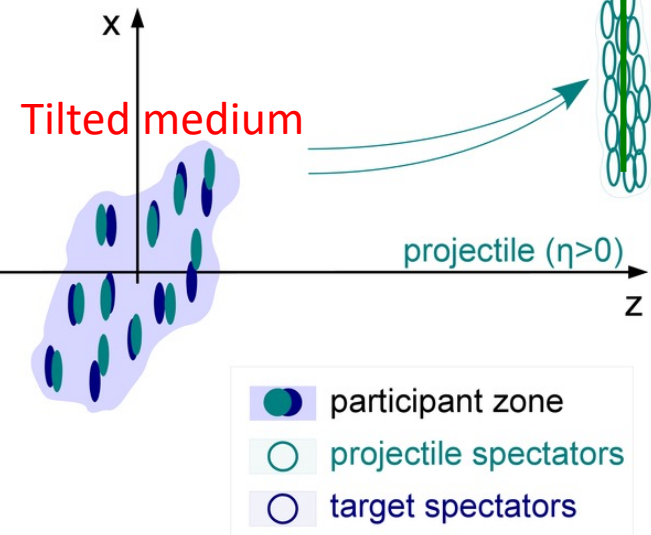
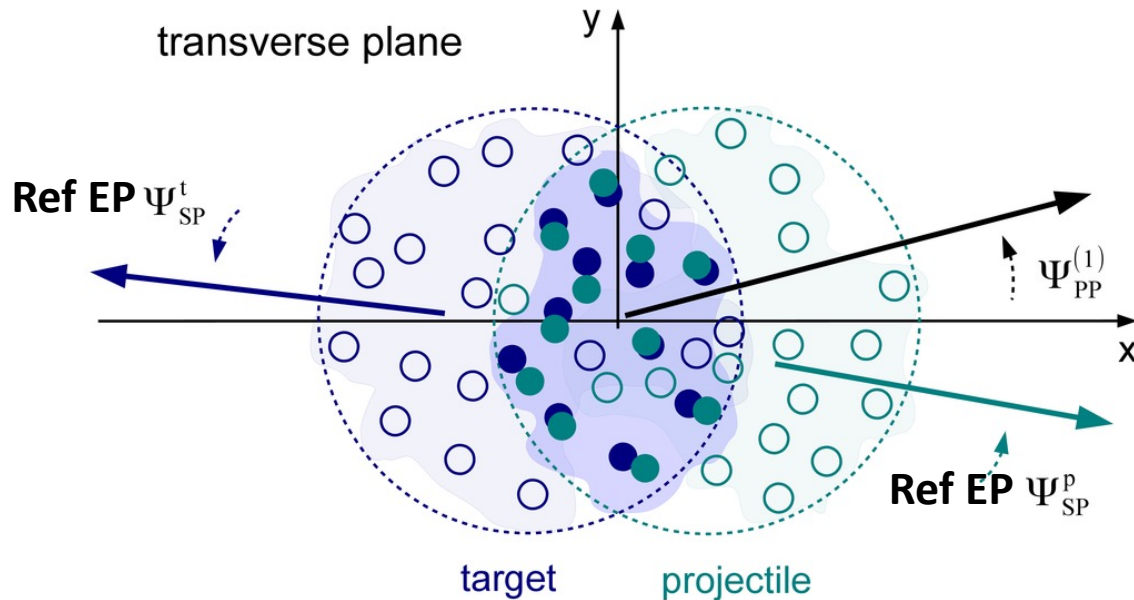
$$v_1 = \langle \cos(\phi - \psi_{sp}) \rangle$$

Preferred moving direction
at different rapidity

Collision symmetry
 $v_1(\eta) = -v_1(-\eta)$

v_1^{odd}

transverse plane

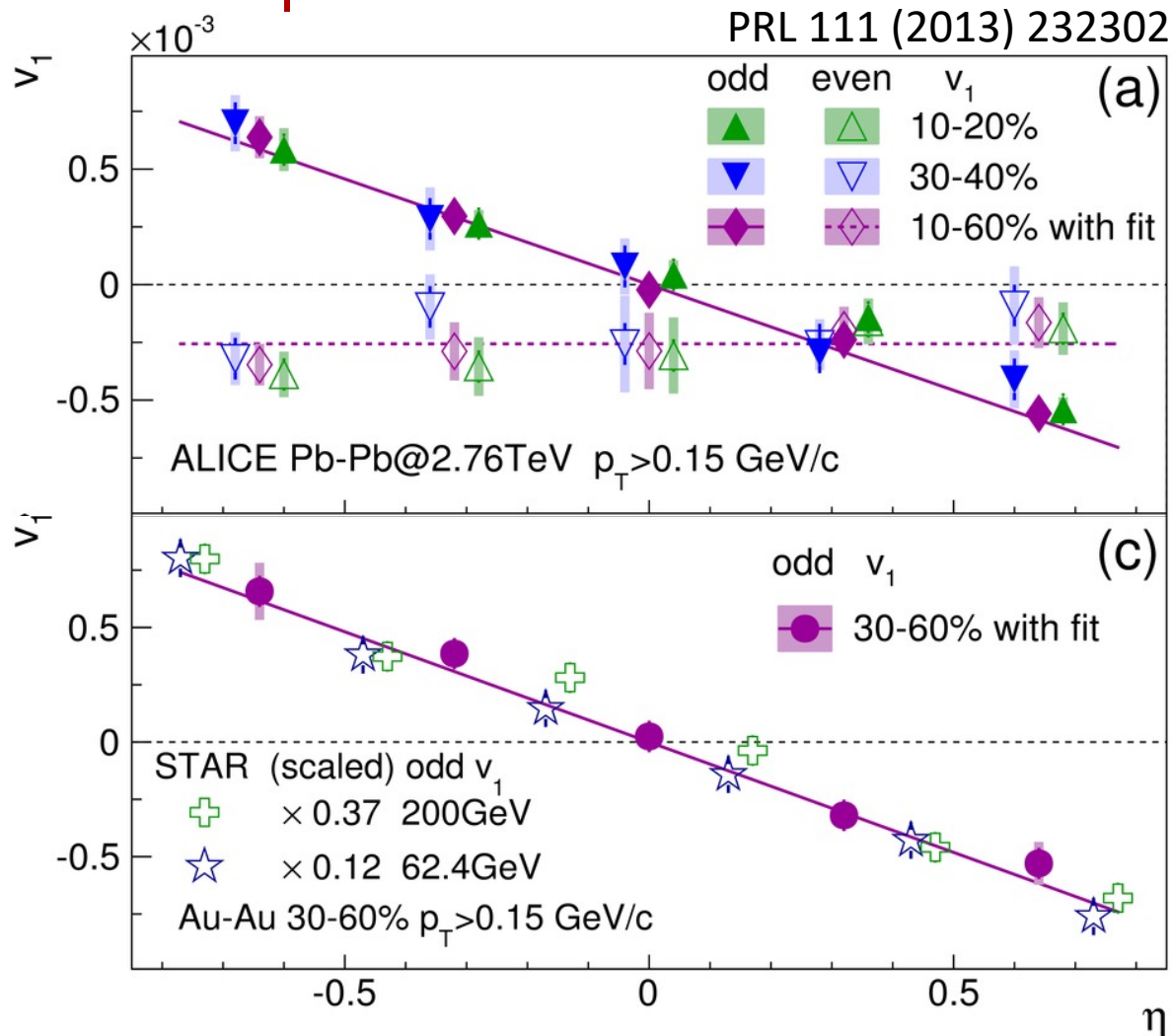


Transverse plane
geometry fluctuation

$$v_1(\eta) = v_1(-\eta)$$

v_1^{even}

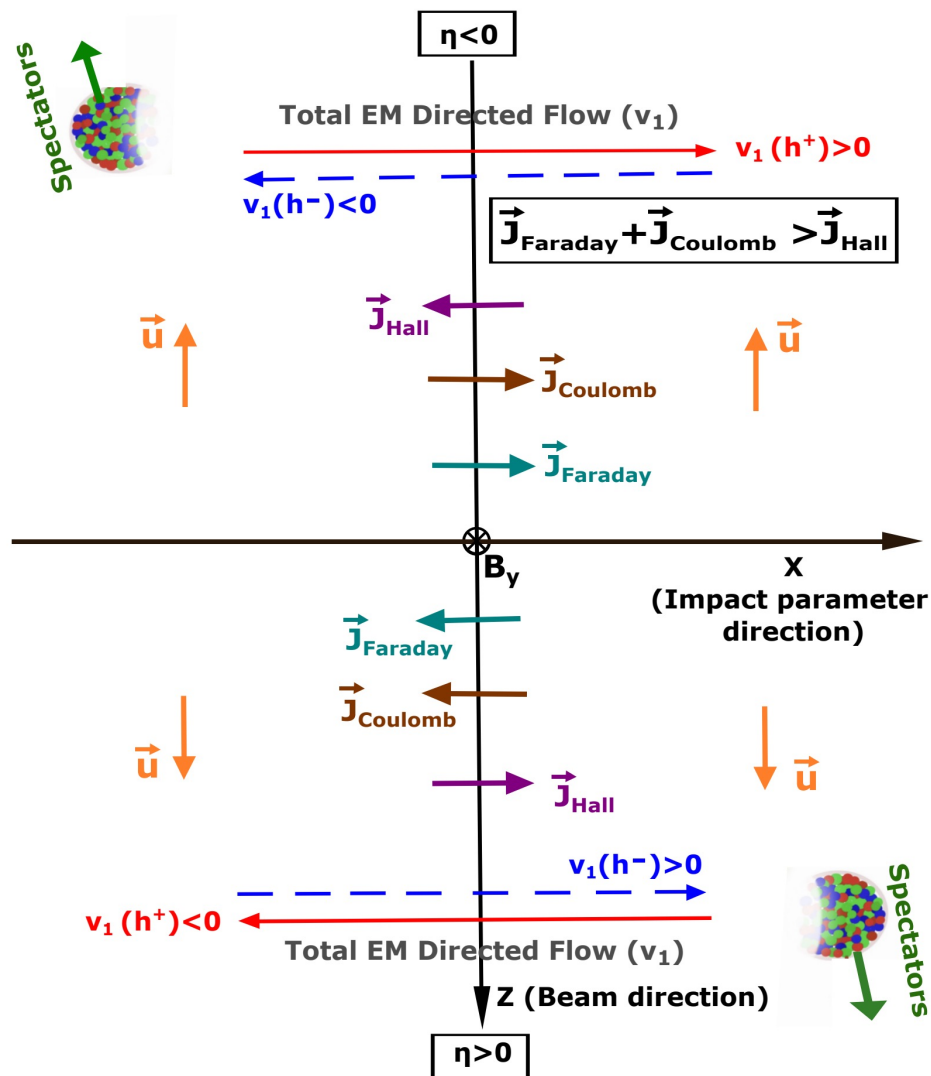
Directed flow v_1



Larger slope of v_1^{odd} at RHIC energies: smaller tilt at LHC

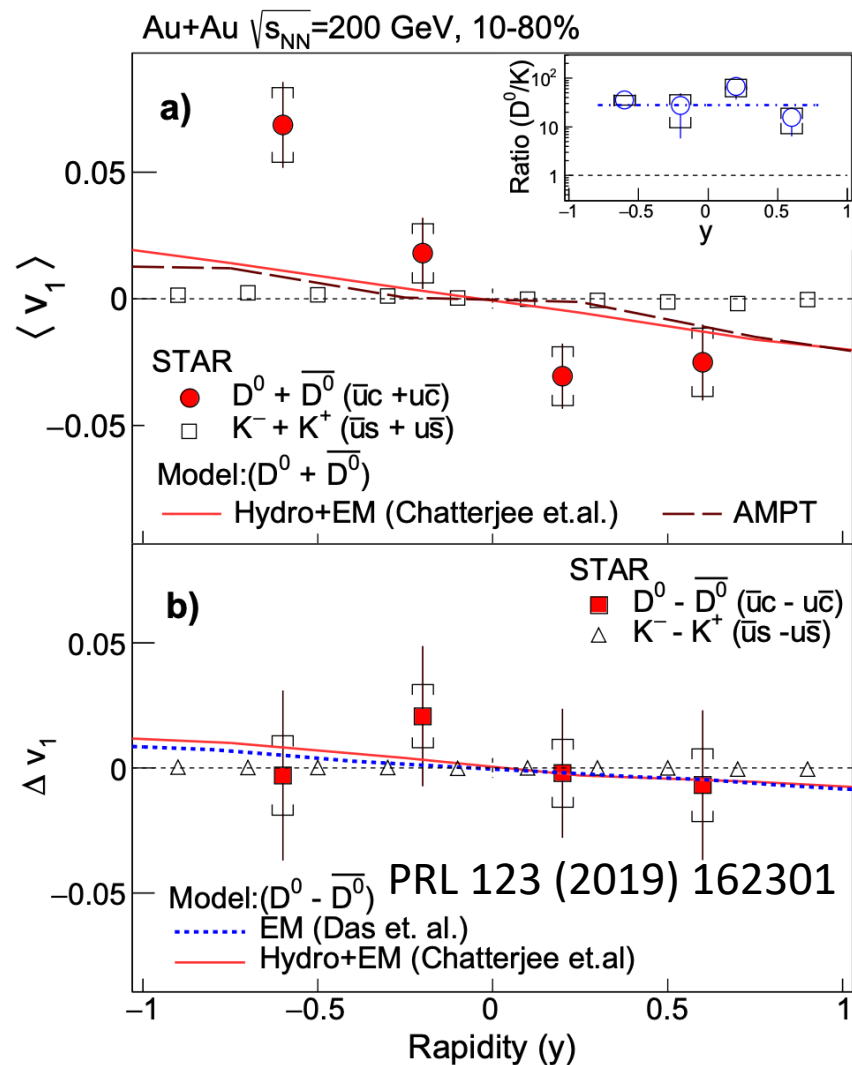
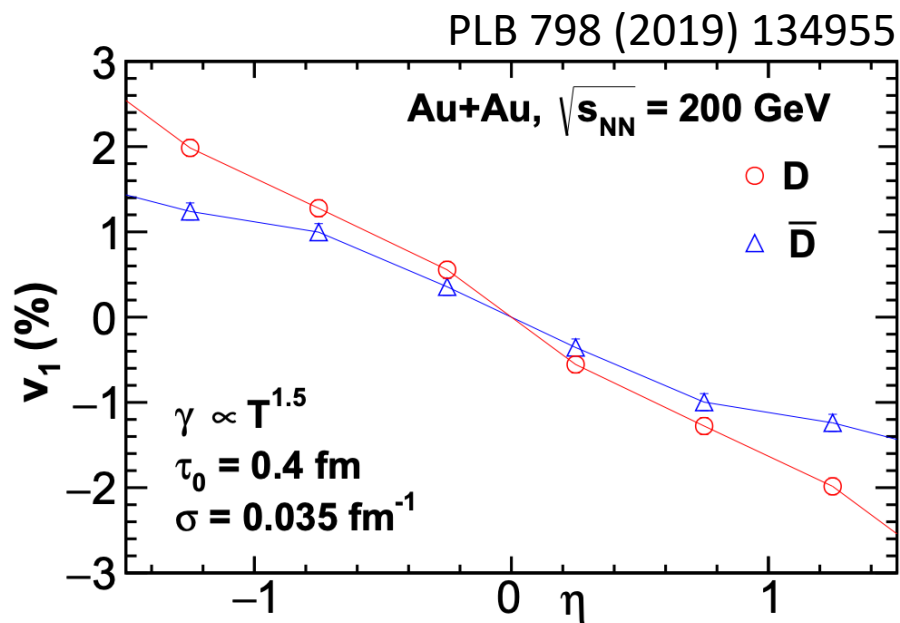
Non-zero v_1^{even} : sizable geometry fluctuation

Directed flow v_1 – B & E field effects



Main electro-magnetic effects on v_1

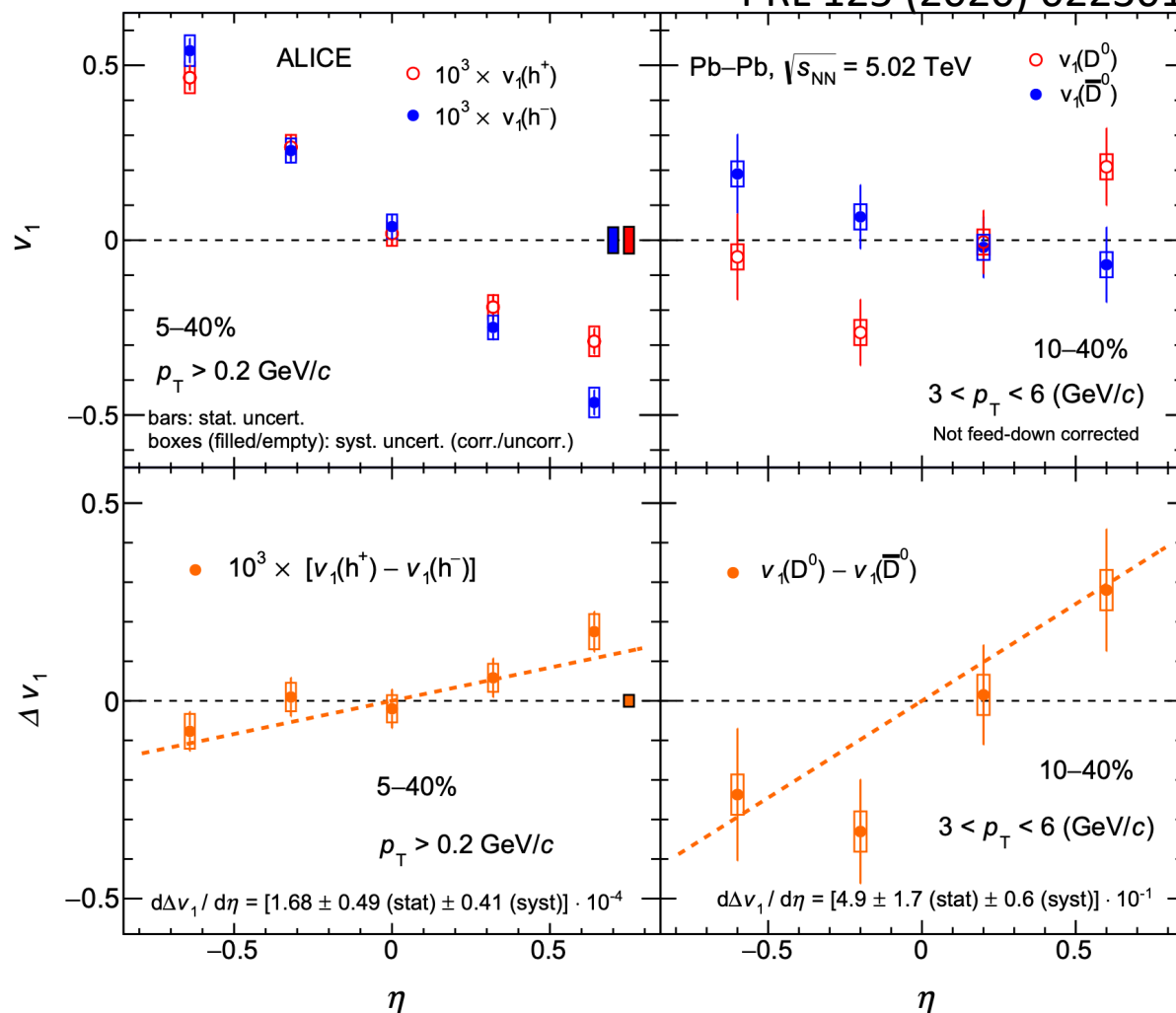
Directed flow v_1 – B & E field effects



Models predict negative slope of Δv_1 @ RHIC
Faraday + Coulomb > Hall

Directed flow v_1 – B & E field effects

PRL 125 (2020) 022301

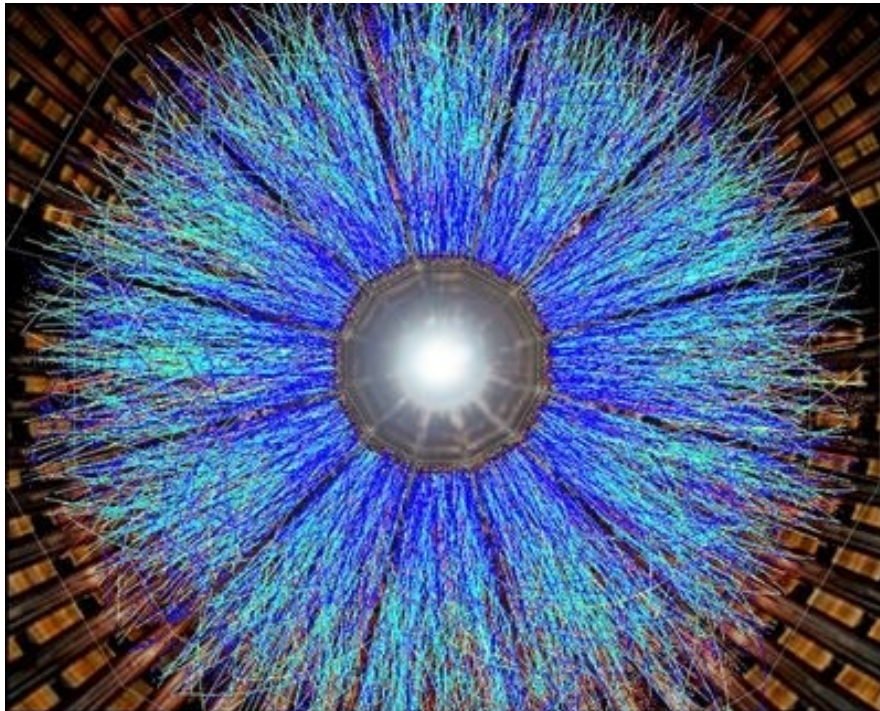


Hint of positive slope of Δv_1 @ LHC
Faraday + Coulomb < Hall

Summary

Collective flow measurements are versatile tools to study all stages of heavy ion collision

- QGP viscosity $\Leftrightarrow v_n$ vs centrality; v_n vs energy; v_n vs p_T ; ...
- Initial state \Leftrightarrow multi-particle v_n ; decorrelation; v_1 ; ...
- Hadronic phase \Leftrightarrow PID v_n ; ...
- ...

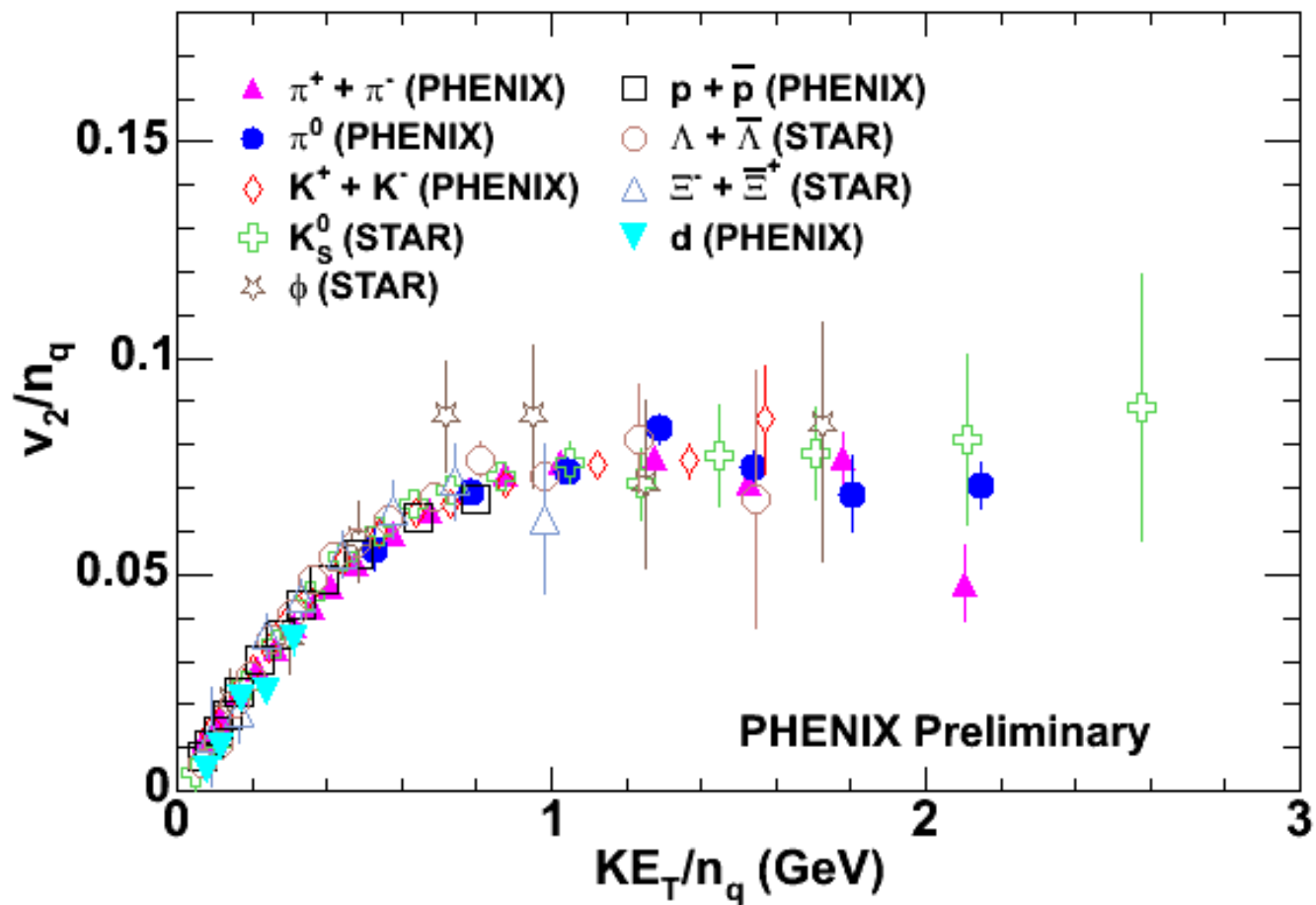


Questions to
zhenyuchen@sdu.edu.cn

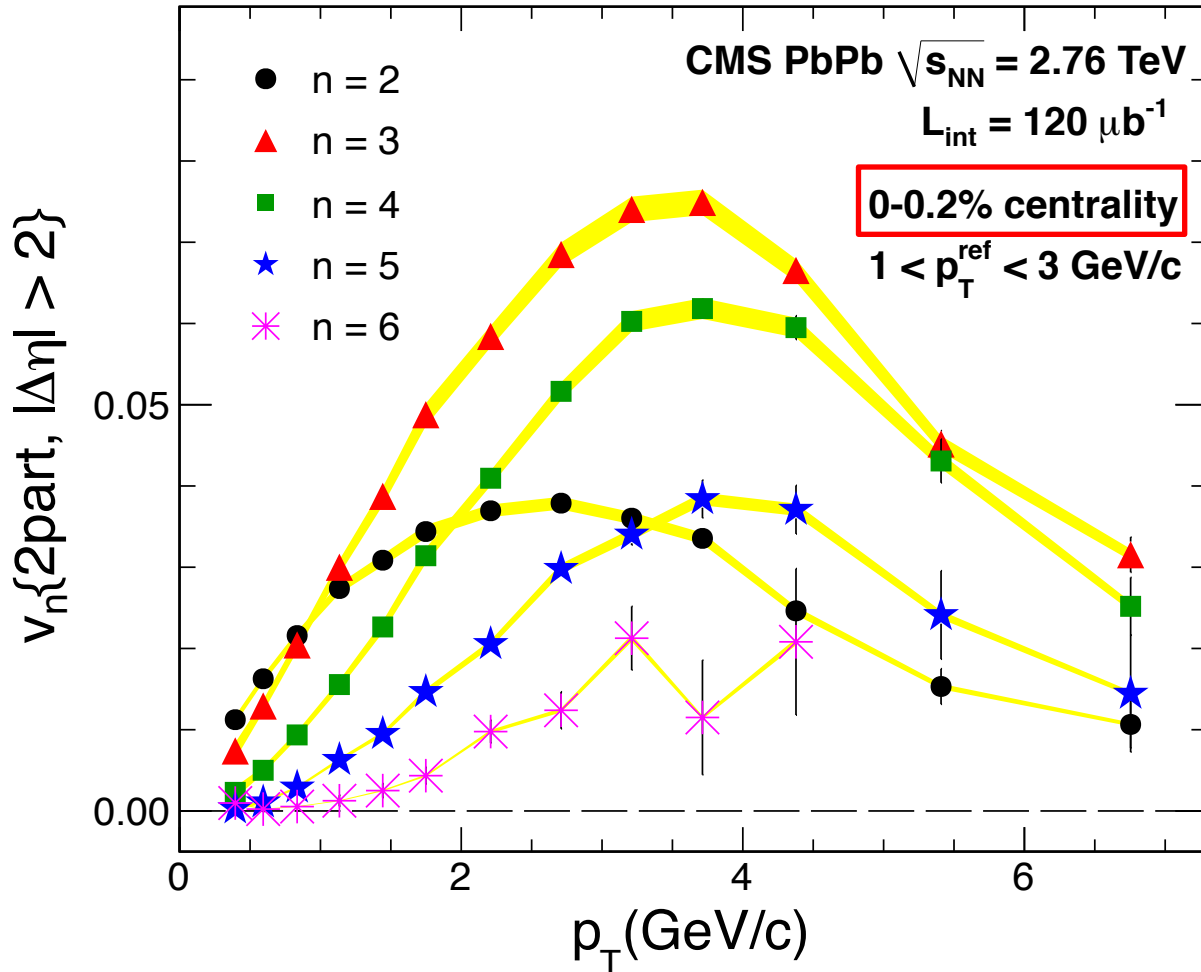


Back up

NCQ



v_n vs centrality



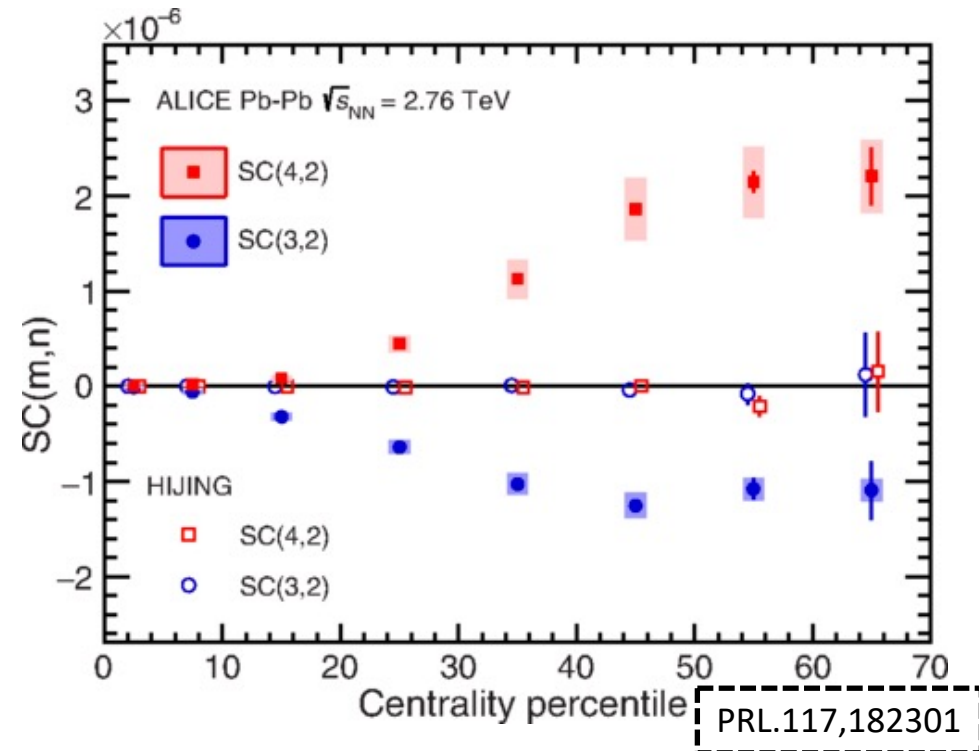
Probe geometry fluctuation with Ultra-Central-Collision

Symmetric cumulants

- Correlation between harmonics:

$$SC(n,m) = \langle v_n^2 v_m^2 \rangle - \langle v_n^2 \rangle \langle v_m^2 \rangle$$

- Symmetric Cumulant (SC) developed by ALICE
 - Based on 4-particle cumulant technique
 - Non-flow highly suppressed at high multiplicity



$SC(2,3) < 0 \rightarrow v_2$ and v_3 are **anti-correlated**

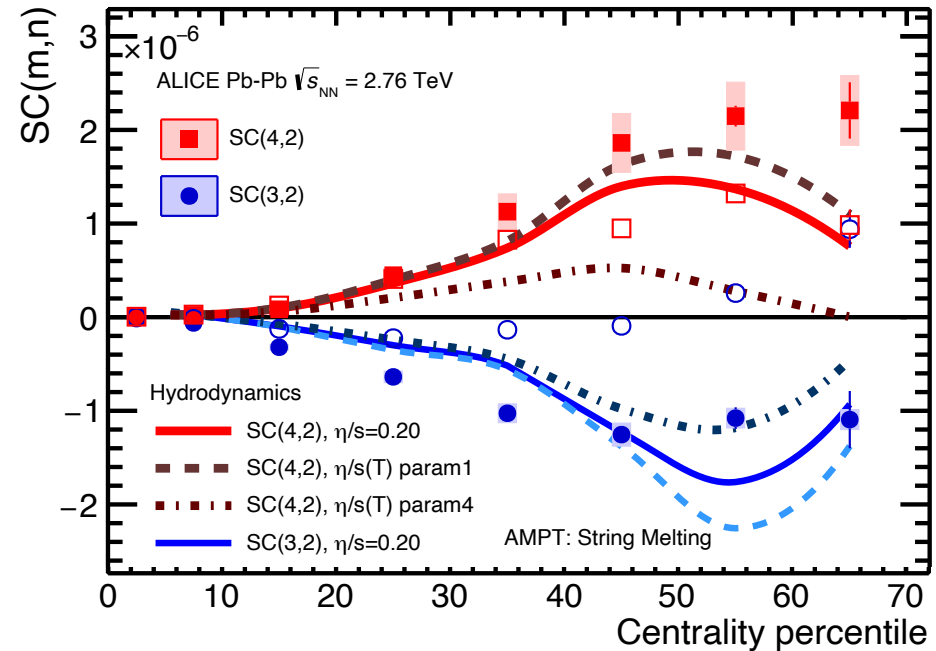
$SC(2,4) > 0 \rightarrow v_2$ and v_4 are **correlated**

Symmetric cumulants

- Correlation between harmonics:

$$SC(n,m) = \langle v_n^2 v_m^2 \rangle - \langle v_n^2 \rangle \langle v_m^2 \rangle$$

- Symmetric Cumulant (SC) developed by ALICE
 - Based on 4-particle cumulant technique
 - Non-flow highly suppressed at high multiplicity

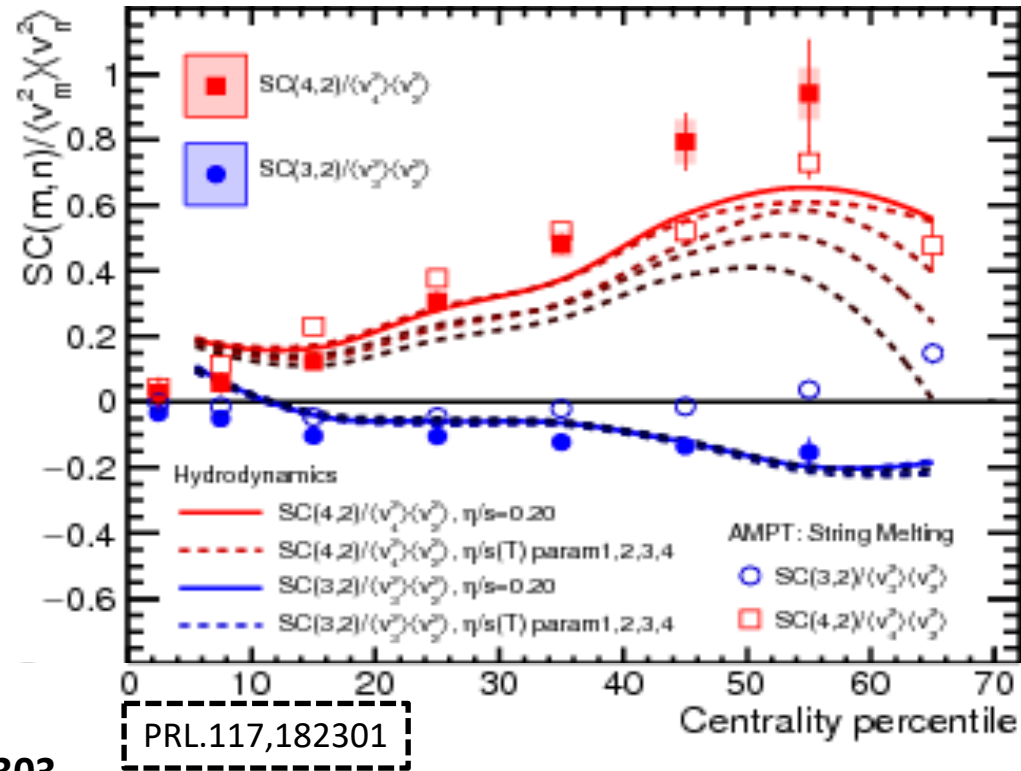


PRL.117,182301

$SC(2,3) < 0 \rightarrow v_2$ and v_3 are **anti-correlated**
 $SC(2,4) > 0 \rightarrow v_2$ and v_4 are **correlated**

Symmetric cumulants

- SC normalized by $\langle v_n^2 \rangle \langle v_m^2 \rangle$
- Normalized SC(2,4)
 - Probe Initial State fluctuation
 - Sensitive to medium response
- Normalized SC(2,3)
 - Probe Initial State fluctuation
 - **Insensitive** to medium response



- Constraints on models
 - [Giacalone et al.](#) **arxiv 1605.08303**
 - [Gardim et al.](#) **arxiv 1608.02982**
 - [Norhona-Holster et al.](#) **arxiv 1609.05171**
 - [Welsh et al.](#) **arxiv 1609.09418**
 - [Zhu et al.](#) **arxiv 1608.05305**

Mass ordering broken for phi

