

Electron scattering from nucleon and nuclei

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Lecture 1 : electron scattering in general

What is electron scattering?

what can we learn by electron scattering?

what have been revealed so far for stable nuclei?

Lecture 2 : selected topics related to NUSYS

electron scattering for

proton charge radius

exotic nuclei

neutron distribution in nuclei

A series of lectures on electron scattering at Beihang Univ., Aug. 3, 4

Electron scattering from nucleon and nuclei

Introduction to electron scattering

what is electron scattering?

what can we learn from electron scattering?

Selected topics of electron scattering for nucleon and nuclei

exotic nuclei

proton charge radius (+ neutron charge radius)

neutron distribution in nuclei

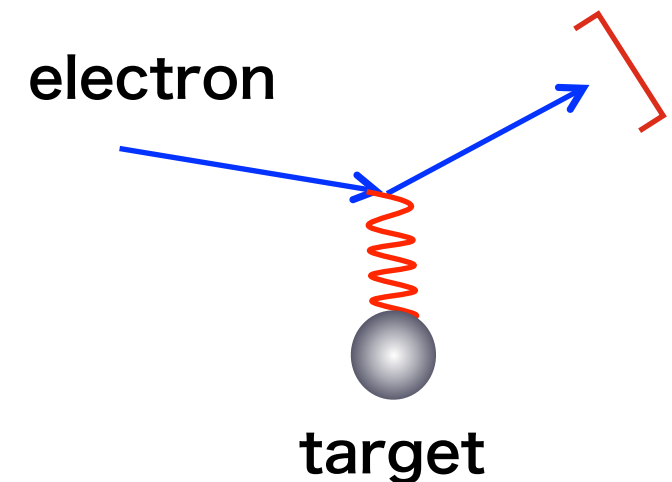
email me if you would like to have. (suda@lns.tohoku.ac.jp)

<https://www.dropbox.com/scl/fo/irxbkzkd5pffbkrh5s234/h?rlkey=nf4xb3prcrg2pch65ck8ry1aj&dl=0>

Electron scattering has consistently played an essential role to reveal detailed structures of nucleon and nuclei

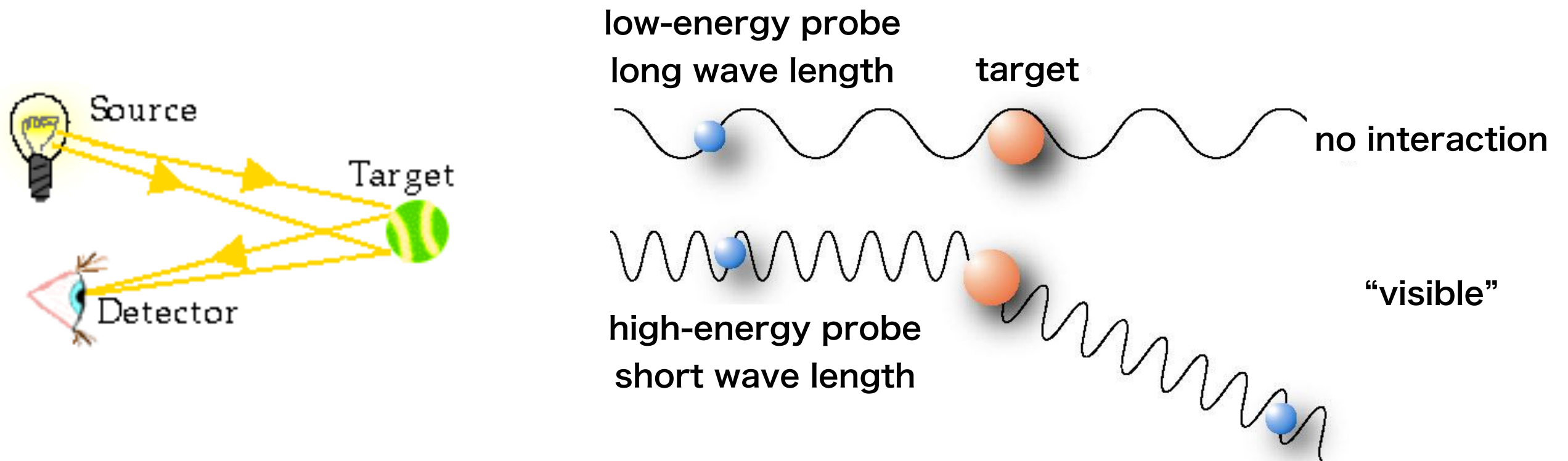
one detects only scattered electrons

➔ very “simple” measurements



1. elementary particle - structure-less -
2. electro-weak interaction - best understood -
3. “relatively” weak - probing deep inside of the target -

Electrons do not experience the nuclear strong force.

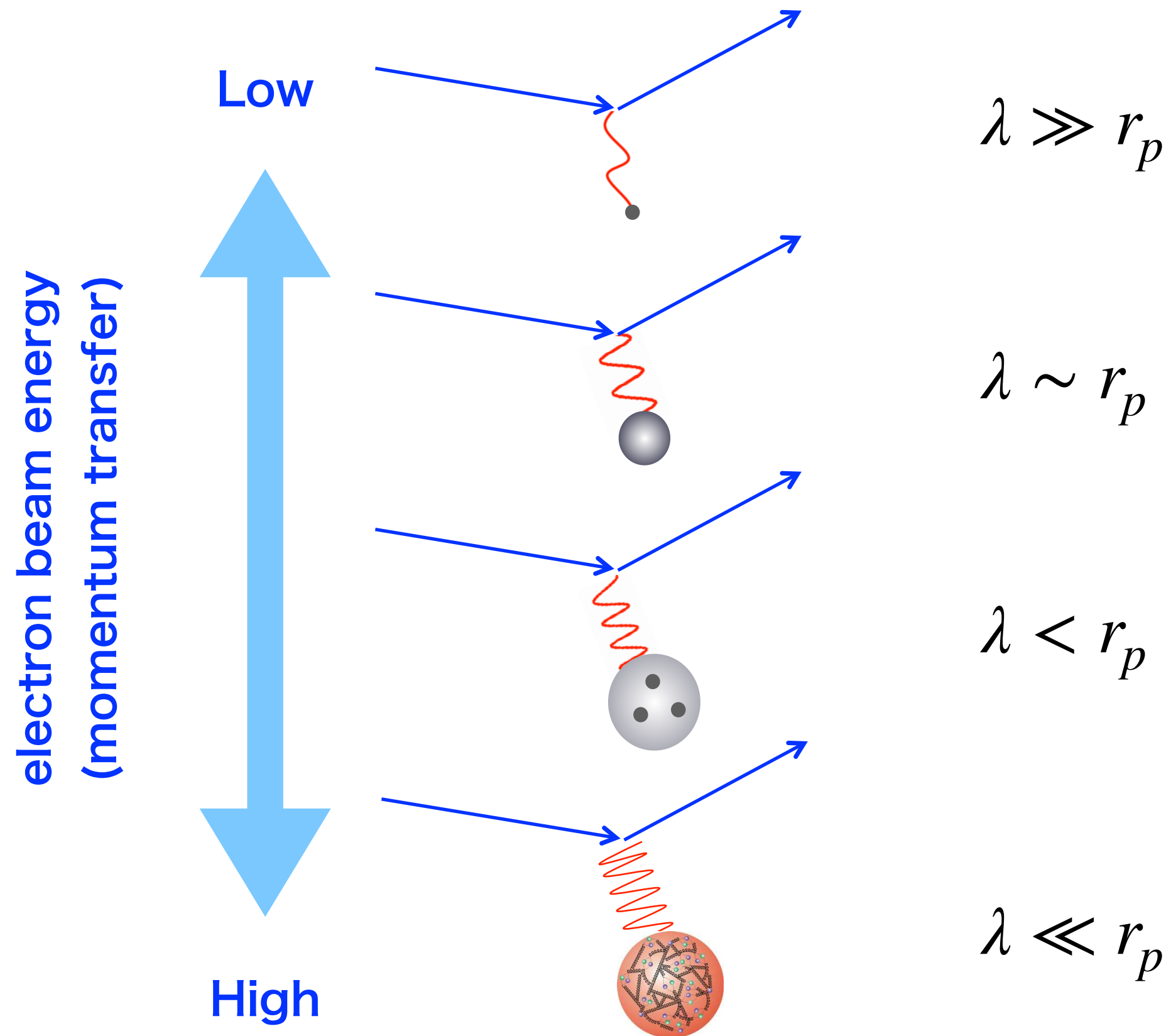


de Broglie wave length of electron \sim target size (\sim fm)

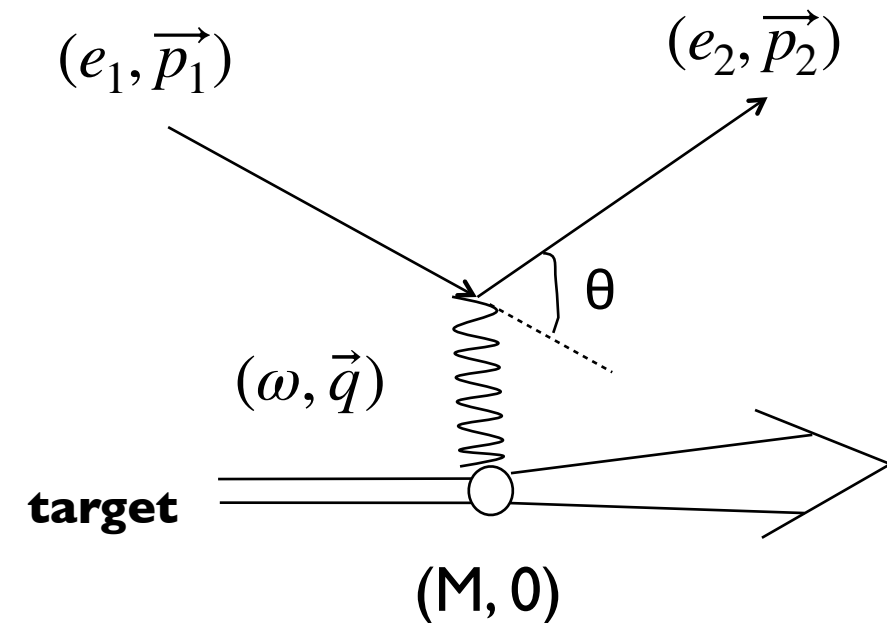
$$\lambda = \frac{\lambda}{2\pi} = \frac{\hbar}{p} = \frac{\hbar c}{pc}$$

$$= \frac{197[\text{MeV} \cdot \text{fm}]}{pc[\text{MeV}/c \cdot c]} \sim 1\text{fm}$$

$$p \sim 200 \text{ MeV}/c$$



Lab. frame (target at rest)



kinematical valuables under URL

$$\omega = e_1 - e_2$$

energy transfer

$$\vec{q} = \vec{e}_1 - \vec{e}_2$$

momentum transfer

$$Q^2 = \omega^2 - \vec{q}^2 = 4e_1e_2 \sin^2 \frac{\theta}{2}$$

4-momentum transfer

$$e^2 = \vec{p}_e^2 + m_e^2$$

URL (Ultra-relativistic limit) : $e \gg m_e \Rightarrow |\vec{p}_{1,2}| = e_{1,2}$

Mott cross section

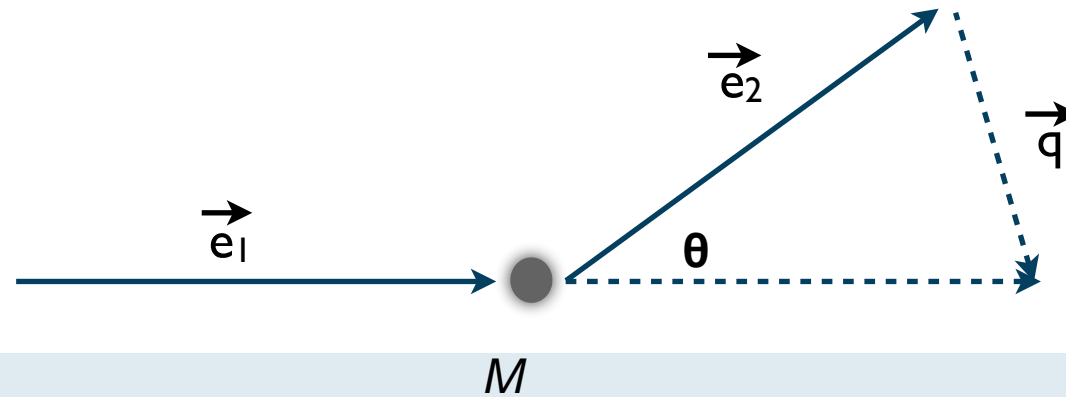
$$\beta \sim 1 \quad e \sim p$$

$$\frac{d\sigma_{Mott}}{d\Omega} = \frac{Z^2 \alpha^2}{4e^2} \frac{\cos^2 \frac{\theta}{2}}{\sin^4 \frac{\theta}{2}}$$

Rutherford cross section

$$\beta \ll 1 \quad e = \frac{p^2}{2m}$$

$$\frac{d\sigma_{Rutherford}}{d\Omega} = \frac{Z^2 \alpha^2}{16e^2} \frac{1}{\sin^4 \frac{\theta}{2}}$$



kinematics

$$e_2 = \frac{e_1}{1 + \frac{2e_1}{M} \sin^2 \frac{\theta}{2}}$$

$$M \gg e_1 \Rightarrow e_2 \sim e_1$$

$$|e_1| = |e_2| (= e)$$

$$\omega = e_1 - e_2 = 0$$

$$|q| = 2e \sin(\theta/2)$$

$$q = e_1 (= e_2) \quad (\theta = 60^\circ)$$

cross section (under one-photon-exchange approximation)

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma_{Mott}}{d\Omega} |F_c(q)|^2$$

$$F_c(\vec{q}) = \int \rho(\vec{r}) e^{-i\vec{q} \cdot \vec{r}} d\vec{r}$$

form factor

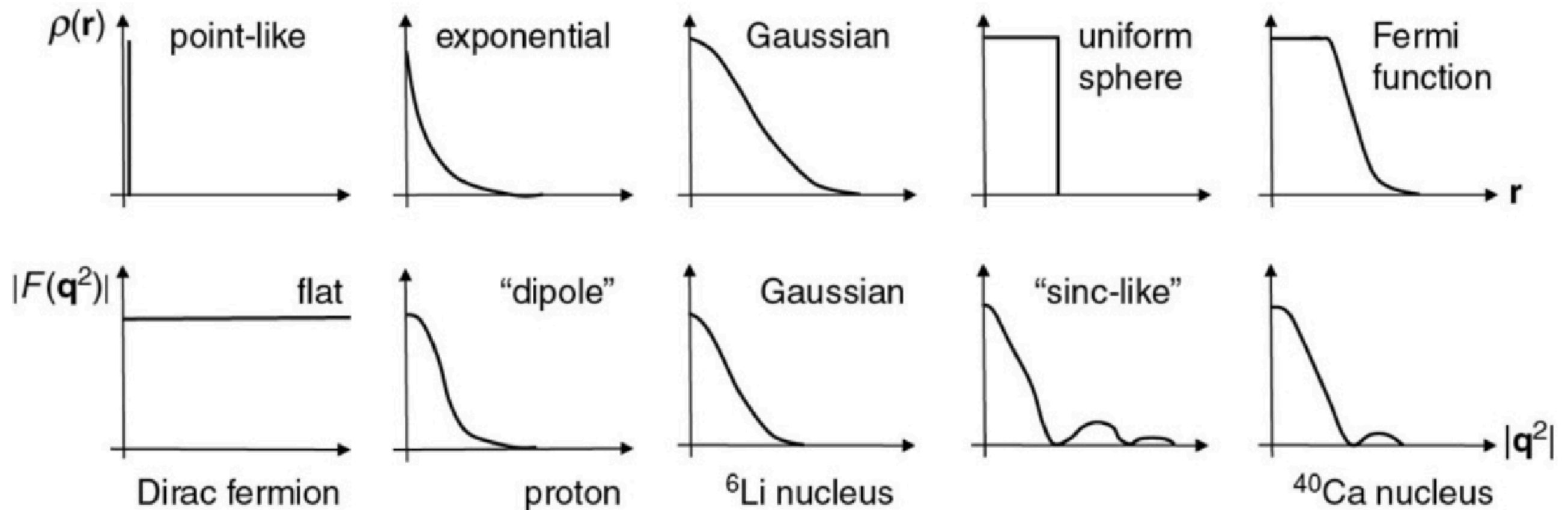
including target structure information

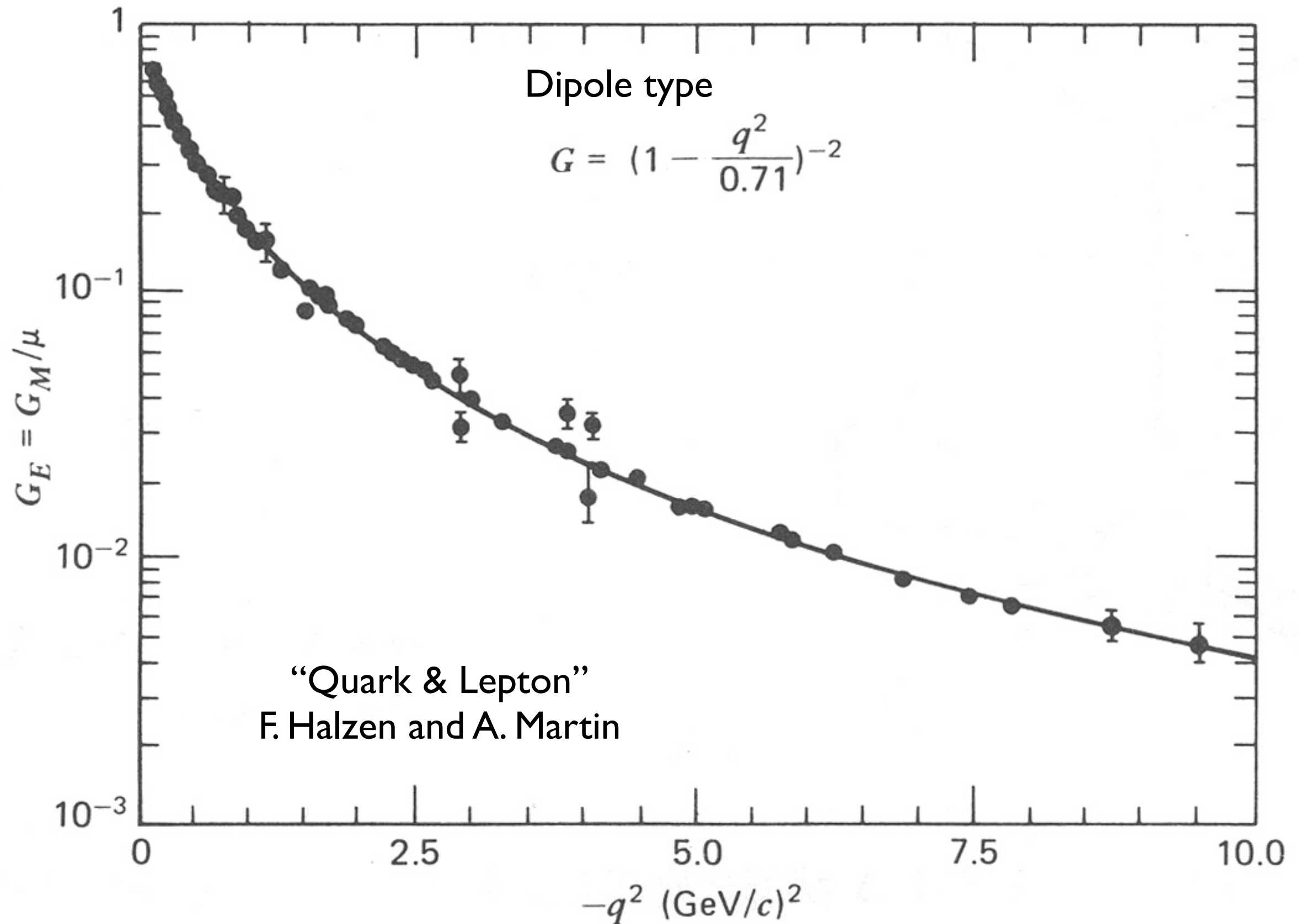
(=1 if the target is a point particle having the charge Z)

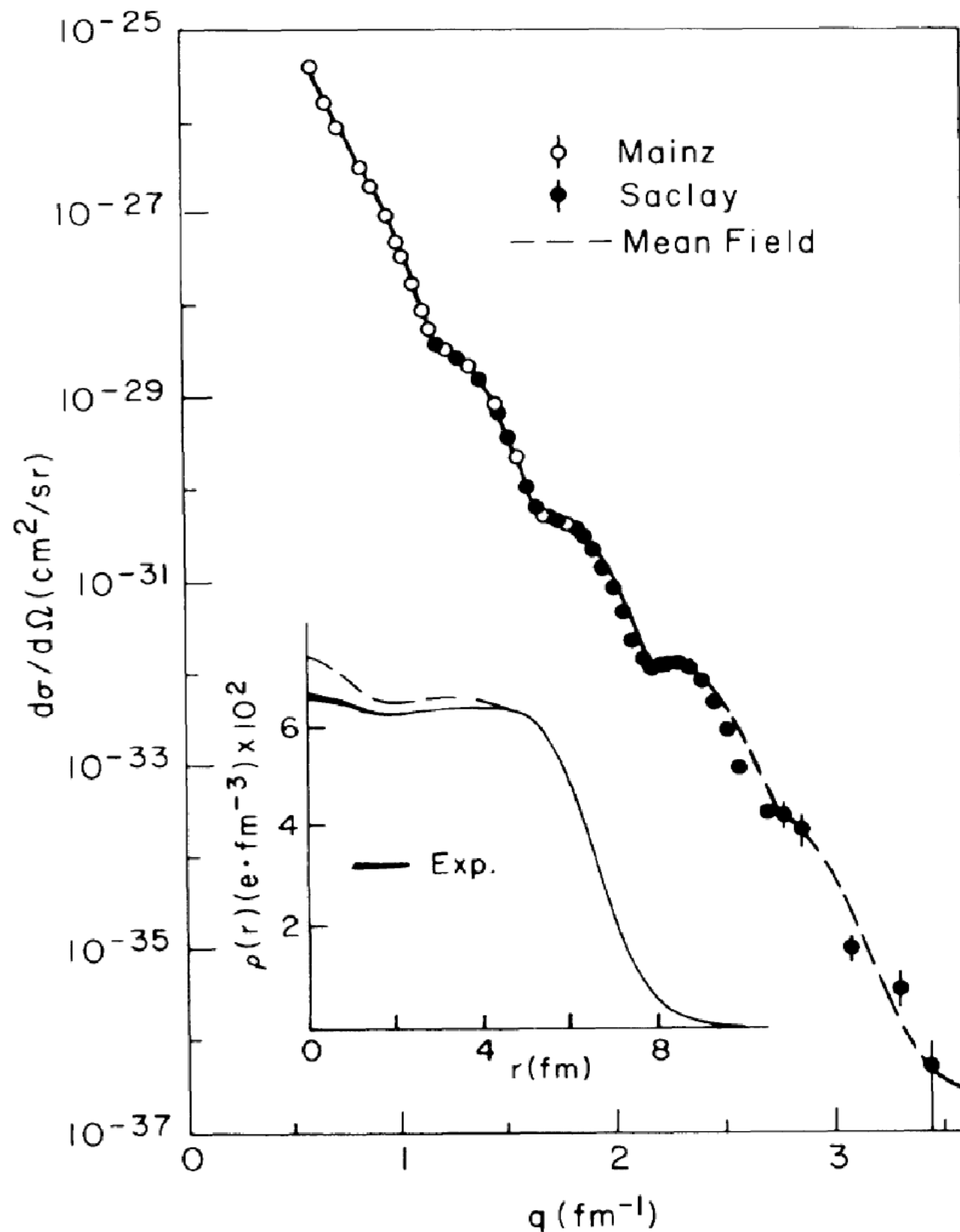
$$F_c(\vec{q}) = \int \rho(\vec{r}) e^{-i\vec{q} \cdot \vec{r}} d\vec{r}$$

for spherically symmetric distributions

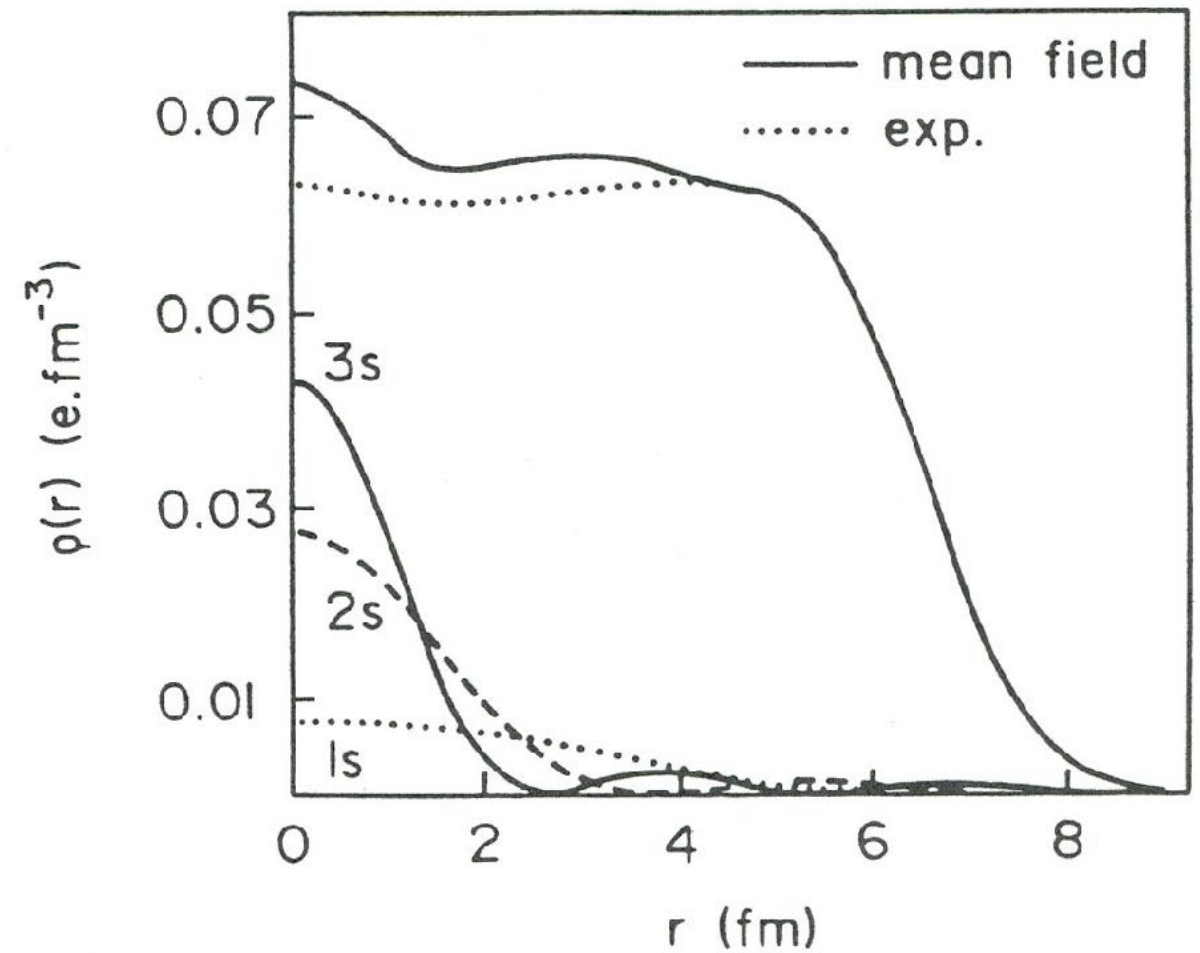
$$F_c(q) = 4\pi \int_0^\infty \rho(r) \frac{\sin(qr)}{qr} r^2 dr$$







J. M. Cavedon et al. PRL 58 (1987) 195 -198.



Phys. Rev. Lett. 38(1977)153.

Magic numbers and stability

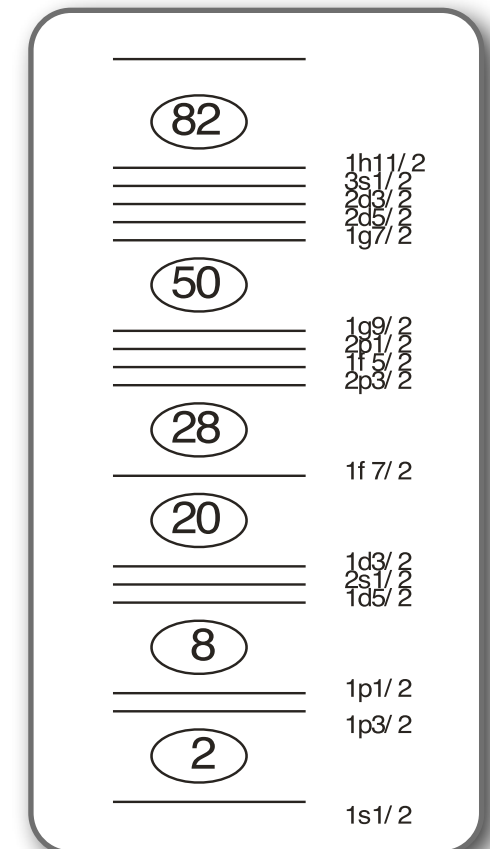
2, 8, 20, 28, 50, 82, 126

doubly magic nuclei

^4He , ^{16}O , $^{40,48}\text{Ca}$, ^{208}Pb



M.G.Mayer and J.Jensen



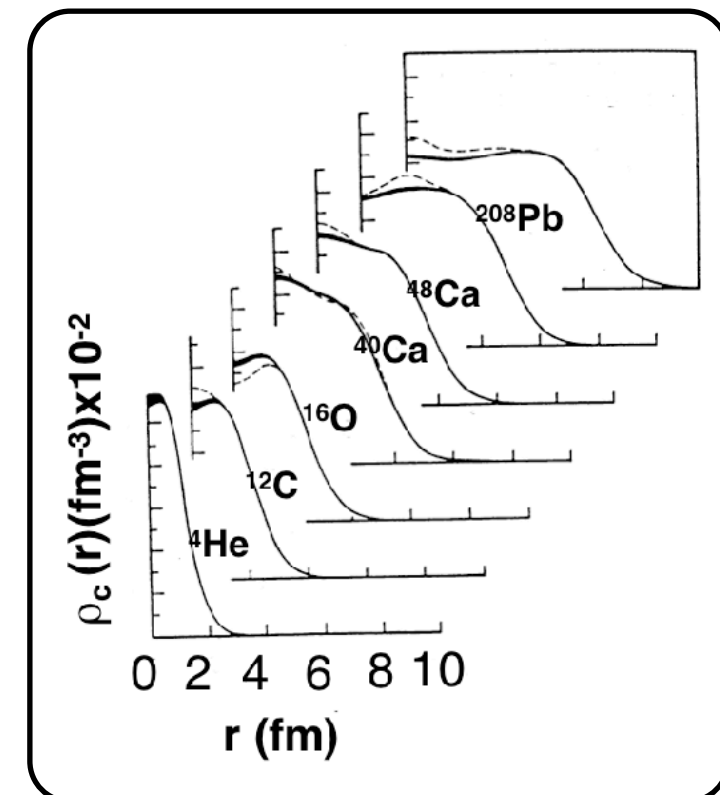
Density distributions

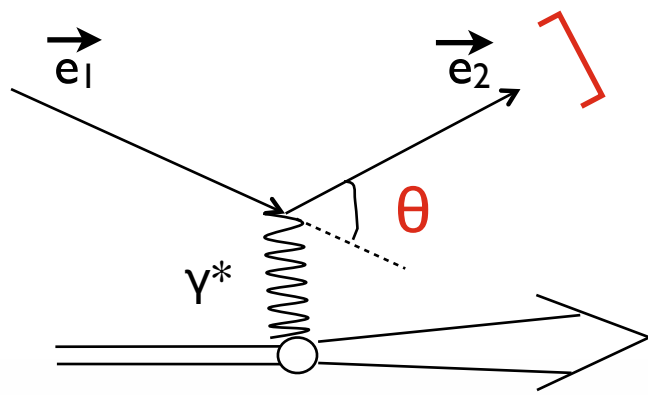
electron scattering

- 1) radius and mass number $r = r_0 A^{1/3}$
- 2) density saturation
- 3) surface diffuseness $\sim \text{const.}$
- 4) $\rho_p(r) \sim \rho_n(r)$



R. Hofstadter





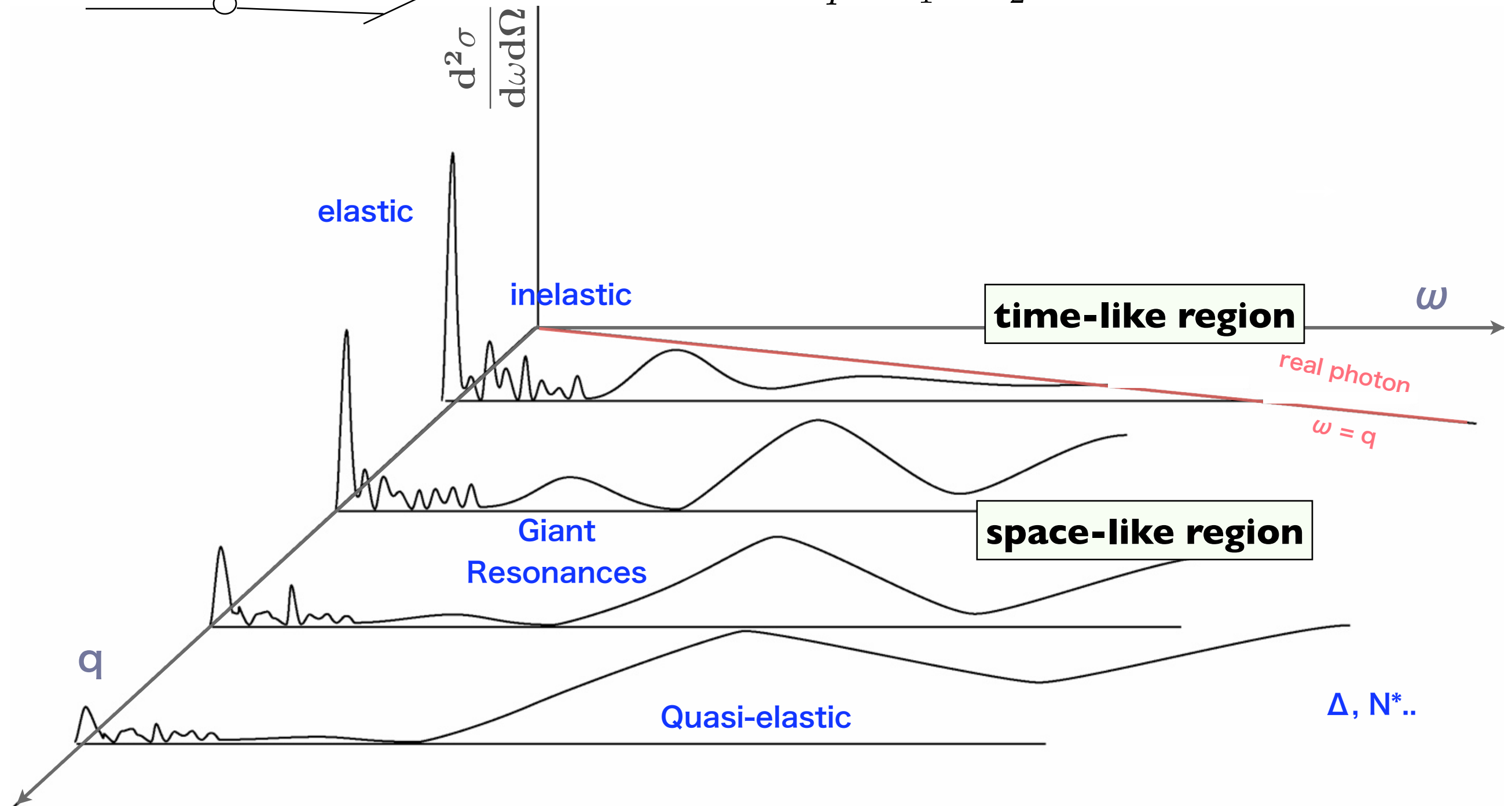
kinematical variables

$$\omega = e_1 - e_2$$

energy transfer

$$\vec{q} = \vec{e}_1 - \vec{e}_2$$

momentum transfer

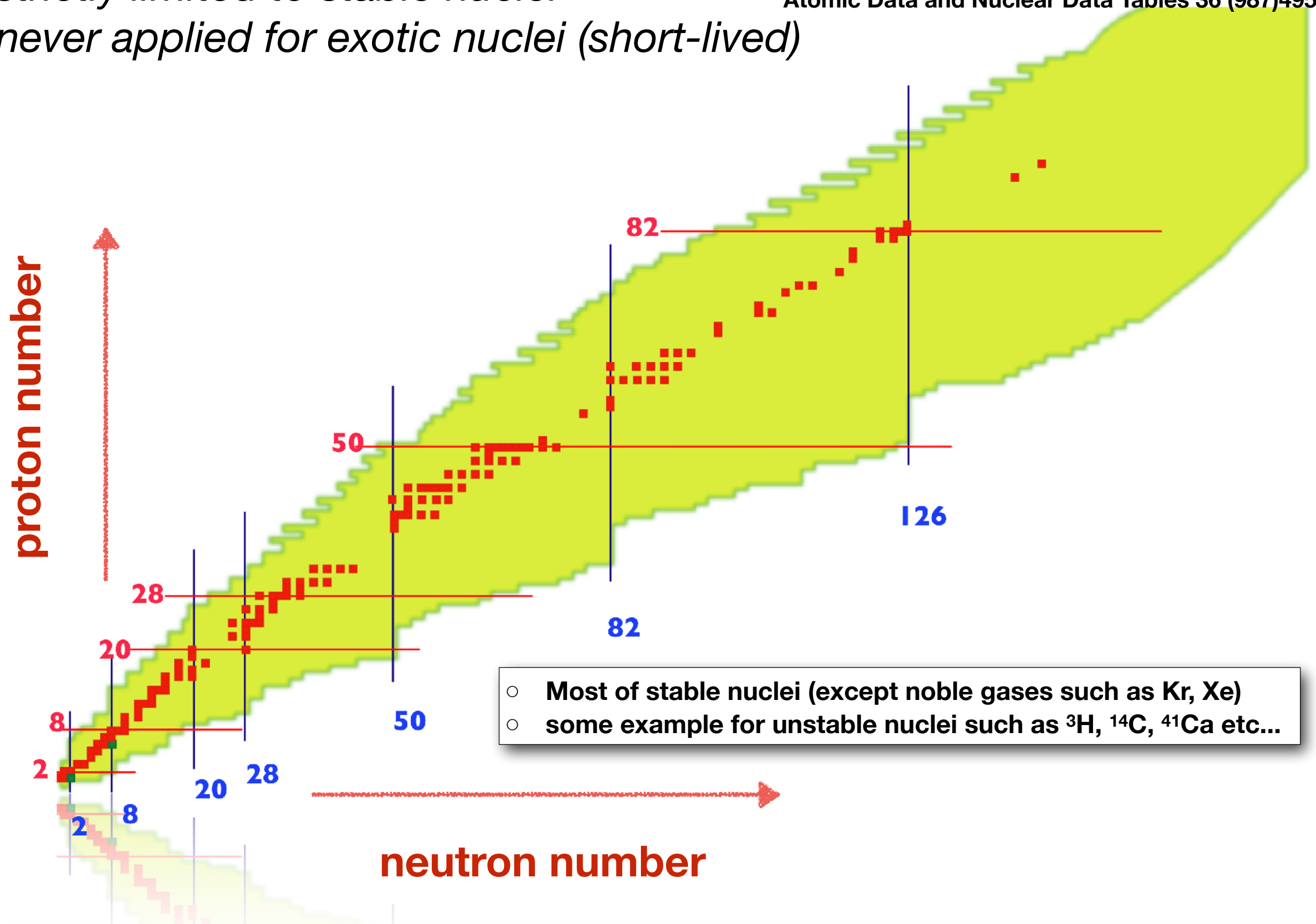


Elastic scattering

- 1) Charge density distributions
- 2) single-particle properties
- 3) deformed nuclei
- 4) valence neutron

- *strictly limited to stable nuclei*
- *never applied for exotic nuclei (short-lived)*

H.deVries, C. deJager and C. deVries
Atomic Data and Nuclear Data Tables 36 (1987)495



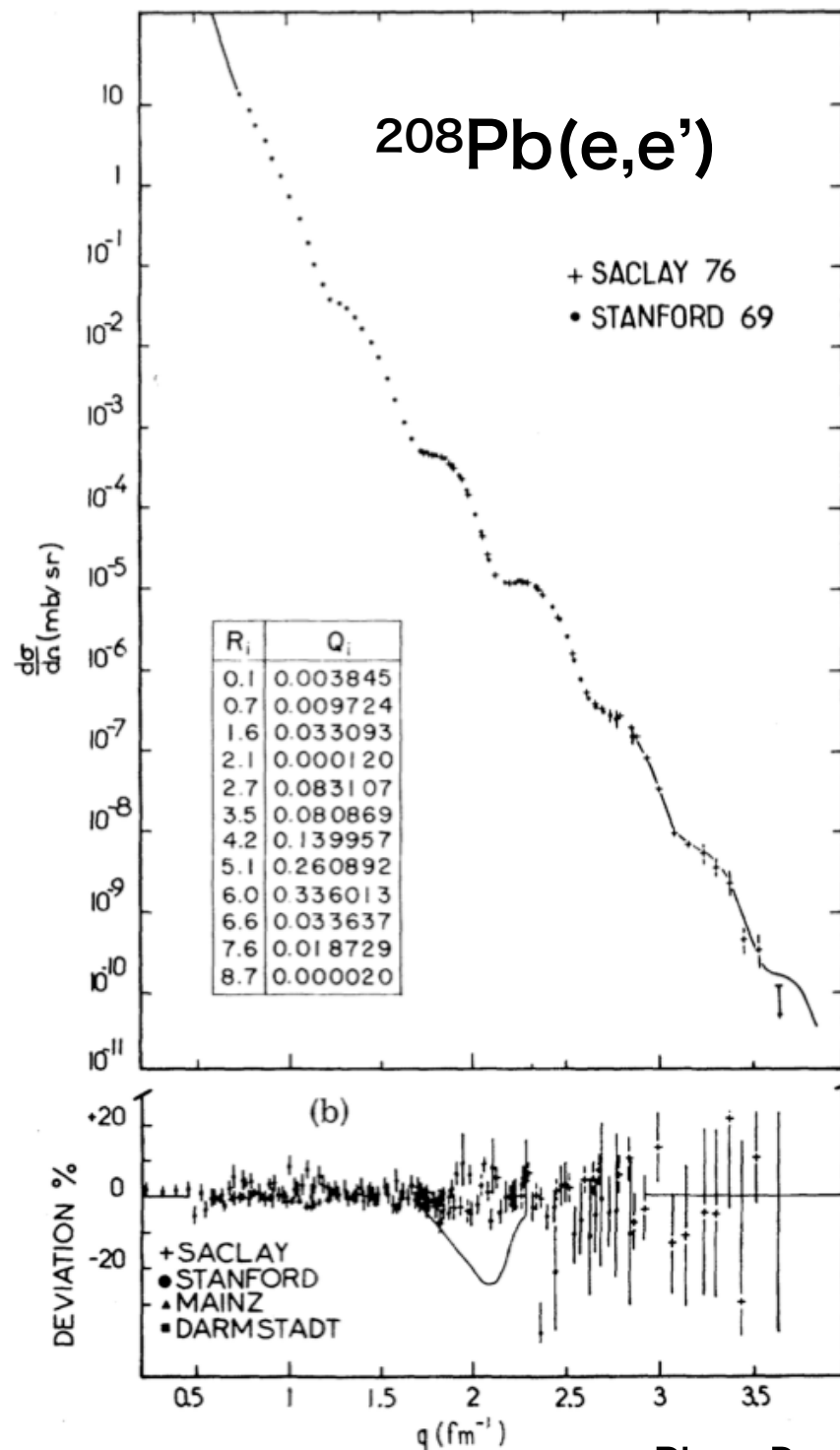
even-even nuclei

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma_{Mott}}{d\Omega} |F_c(q)|^2$$

$$F_c(q) = \int \rho(\vec{r}) e^{-i\vec{q}\cdot\vec{r}} d\vec{r}$$

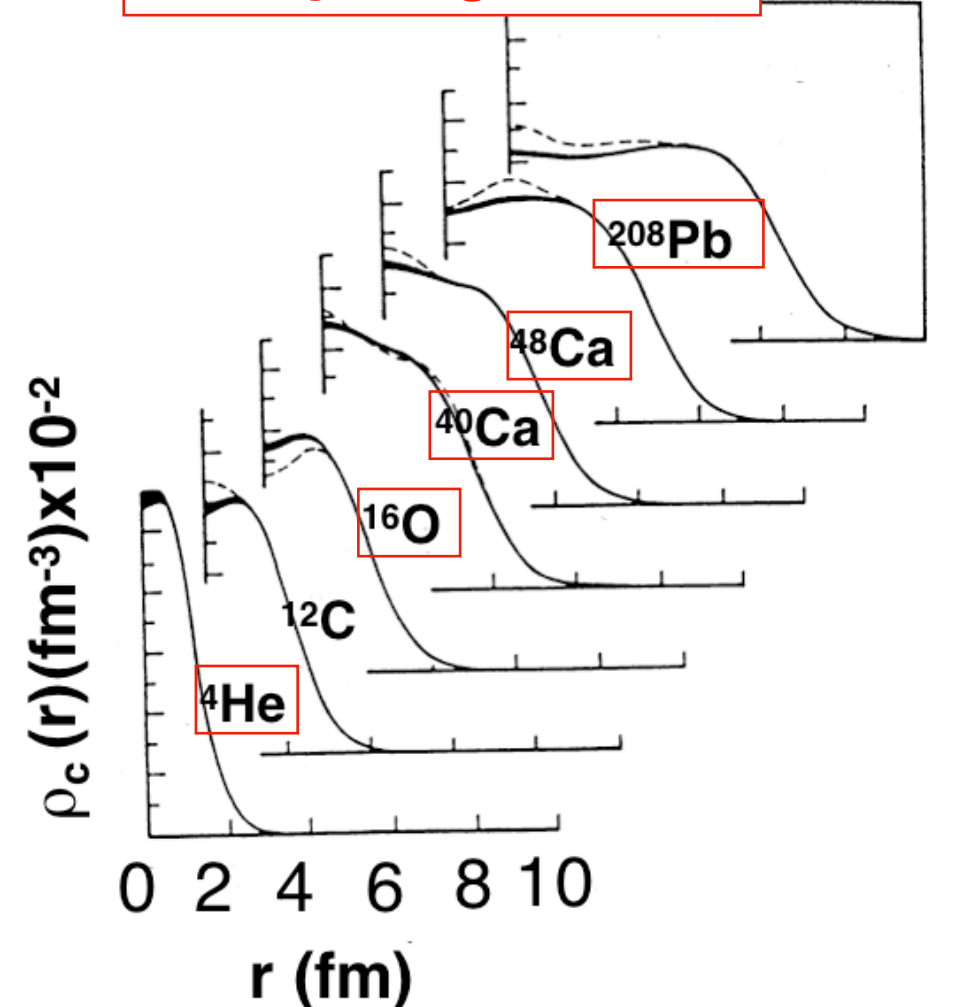
$$\Rightarrow \rho(r) = \int F_c(q) e^{i\vec{q}\cdot\vec{r}} d\vec{q}$$

$$\Rightarrow \rho(r) = \sum_{i=1}^Z |\phi_i(r)|^2$$

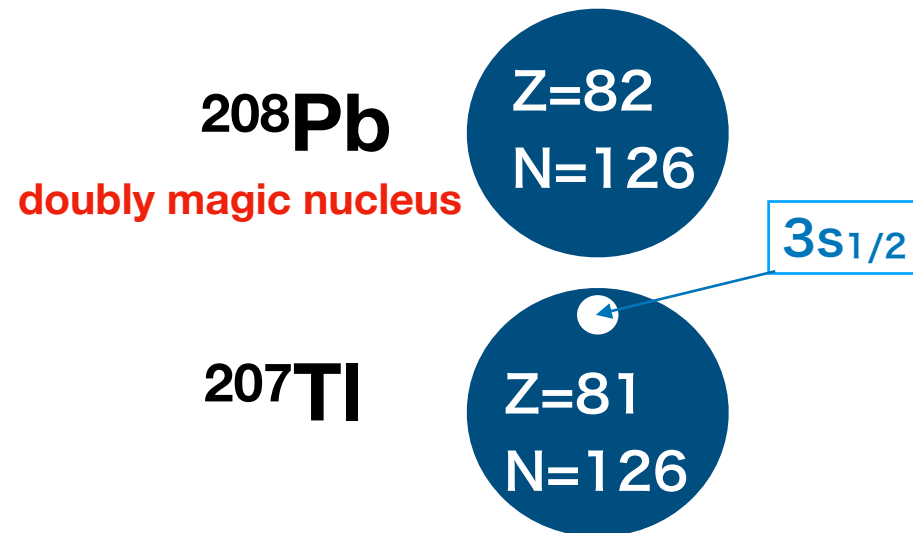


Phys. Rev. Lett. 38(1977)153.

doubly magic nuclei



let's take ^{208}Pb and ^{207}Tl as examples



$$\rho_{^{208}\text{Pb}}(r) = \sum_{i=1}^{82} |\psi_i(r)|^2 = \int F_{^{208}\text{Pb}}(q) e^{i\vec{q}\cdot\vec{r}} d\vec{q}$$

$$\rho_{^{207}\text{Tl}}(r) = \sum_{i=1}^{81} |\psi_i(r)|^2 = \int F_{^{207}\text{Tl}}(q) e^{i\vec{q}\cdot\vec{r}} d\vec{q}$$

under a naive independent particle picture

$$|\psi_{3s_{1/2}}(r)|^2 = \sum_{i=1}^{82} |\psi_i(r)|^2 - \sum_{i=1}^{81} |\psi_i(r)|^2$$

$$= \rho_{^{208}\text{Pb}}(r) - \rho_{^{207}\text{Tl}}(r)$$

elastic electron scattering
for ^{208}Pb and ^{207}Tl

$$= \int [F_{^{208}\text{Pb}}(q) - F_{^{207}\text{Tl}}(q)] e^{i\vec{q}\cdot\vec{r}} d\vec{q}$$

only electromagnetic interaction involved !!

²⁰⁸Pb - ²⁰⁷Tl

unstable
 $\tau \sim 5$ min.

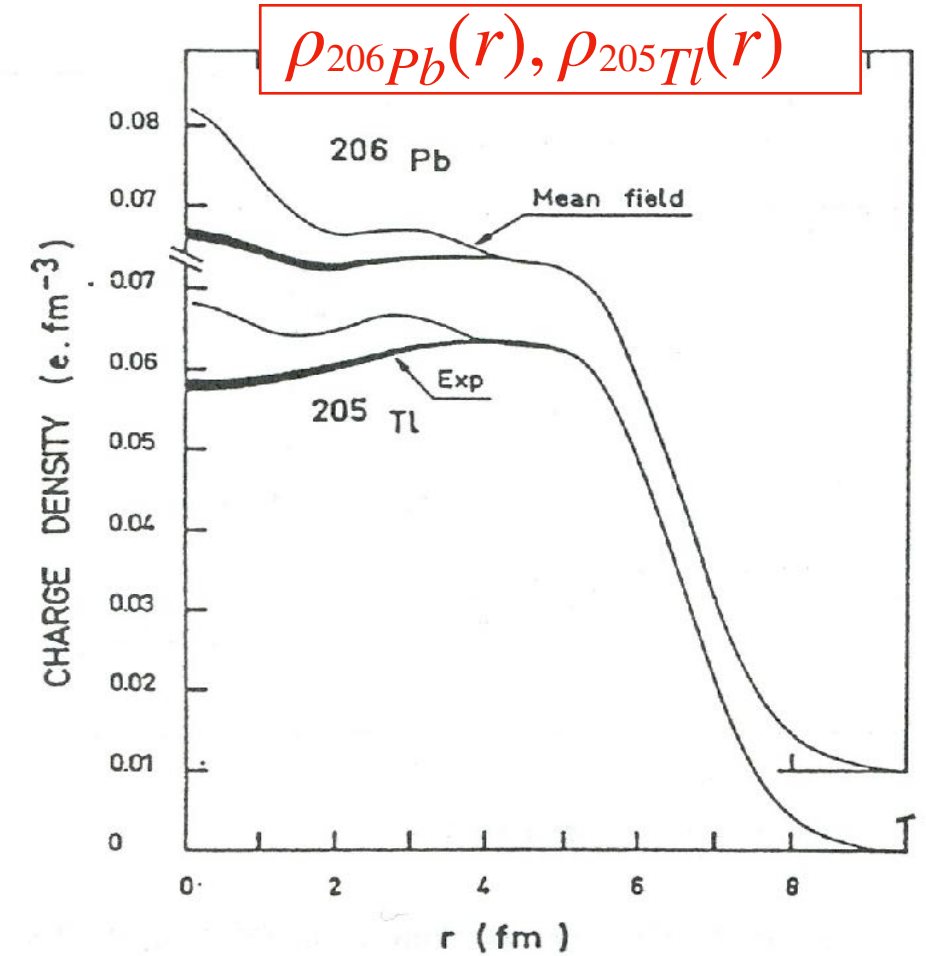
Po							
Bi	203	204	205	206	207	208	209
Pb	202	203	204	205	206	207	208
Tl	201	202	203	204	205	206	207
Hg	200	201	202	203	204	205	206
Au							

Z = 82

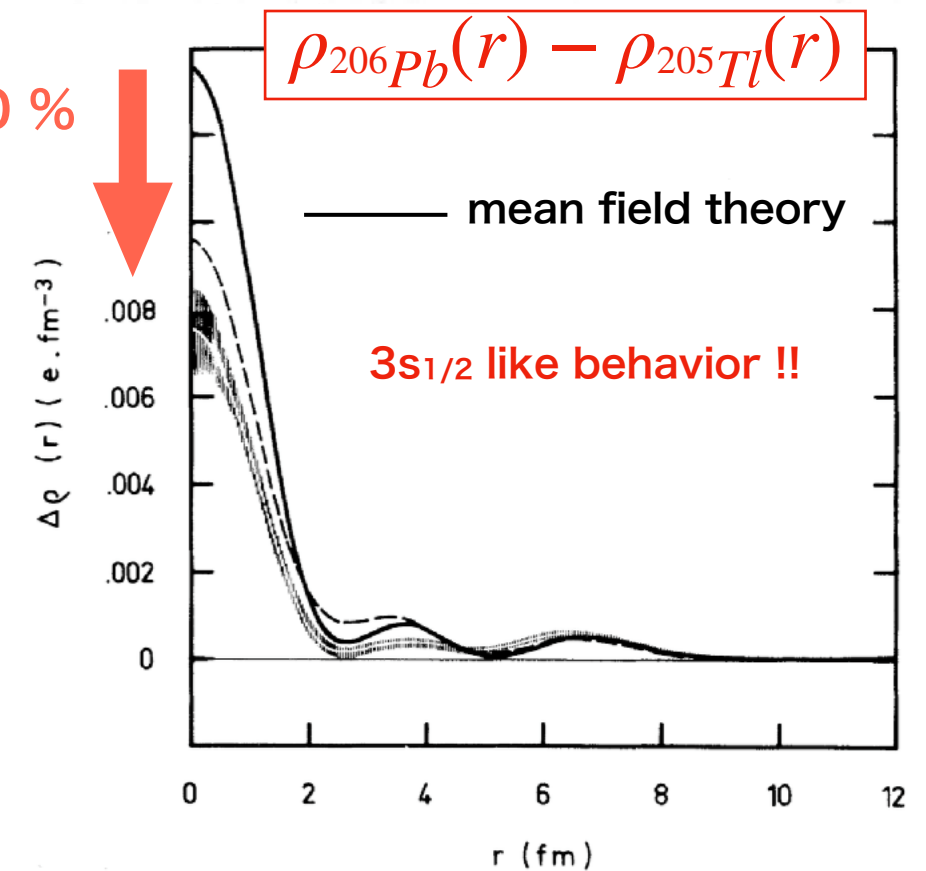
N = 126

instead

²⁰⁶Pb - ²⁰⁵Tl

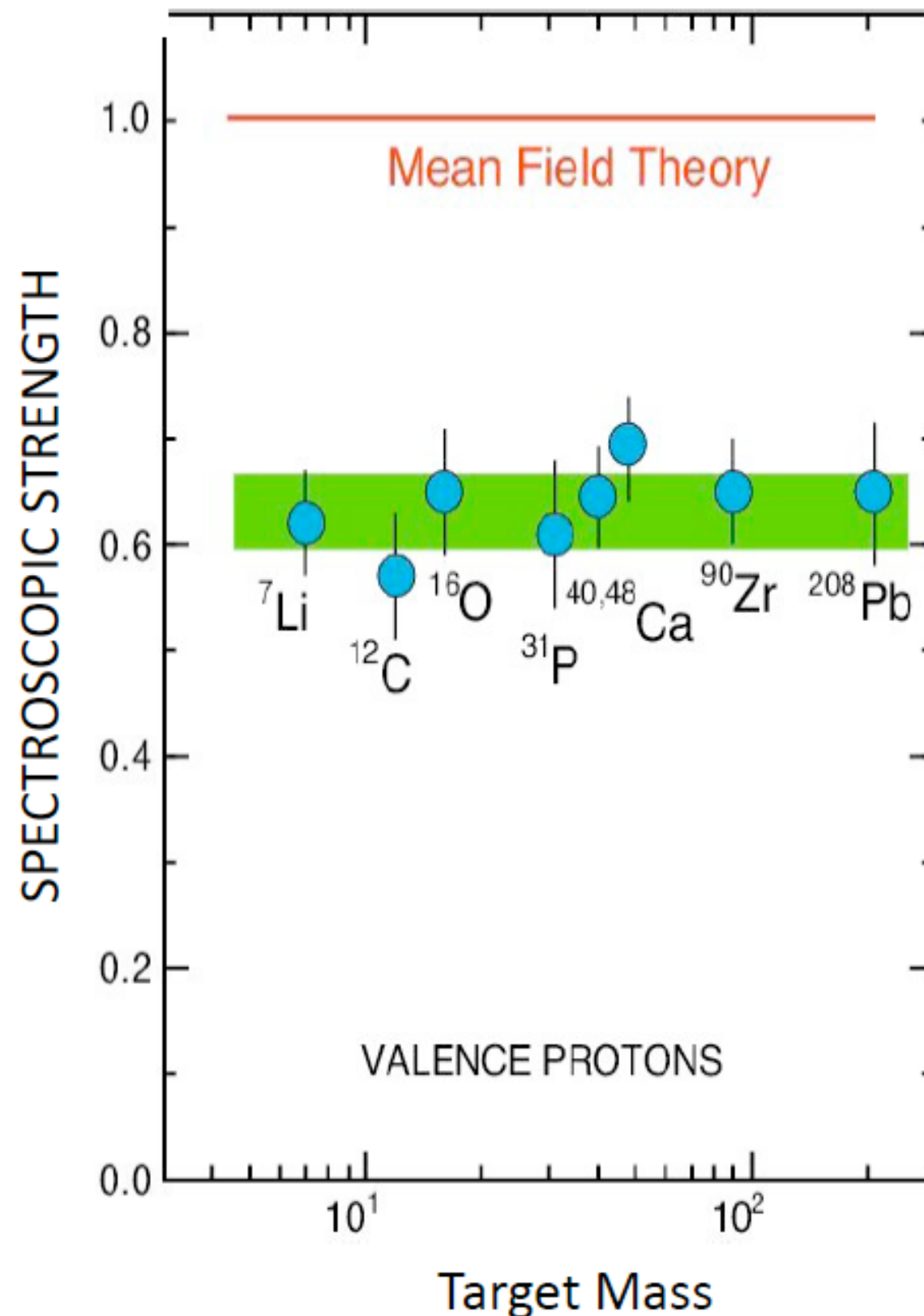


~60 %



spectroscopic factors deduced from quasi-elastic (e,e'p)

60-70 % of the occupation expected from the independent particle model



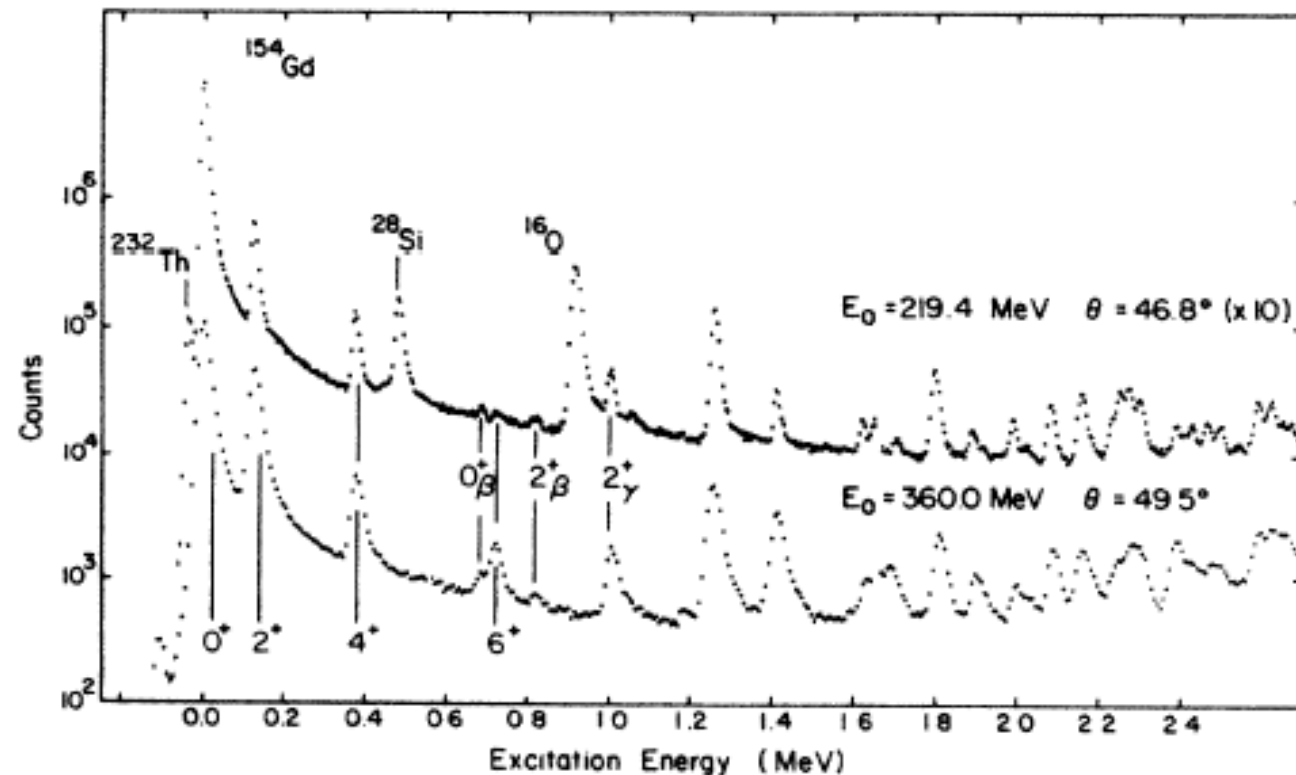
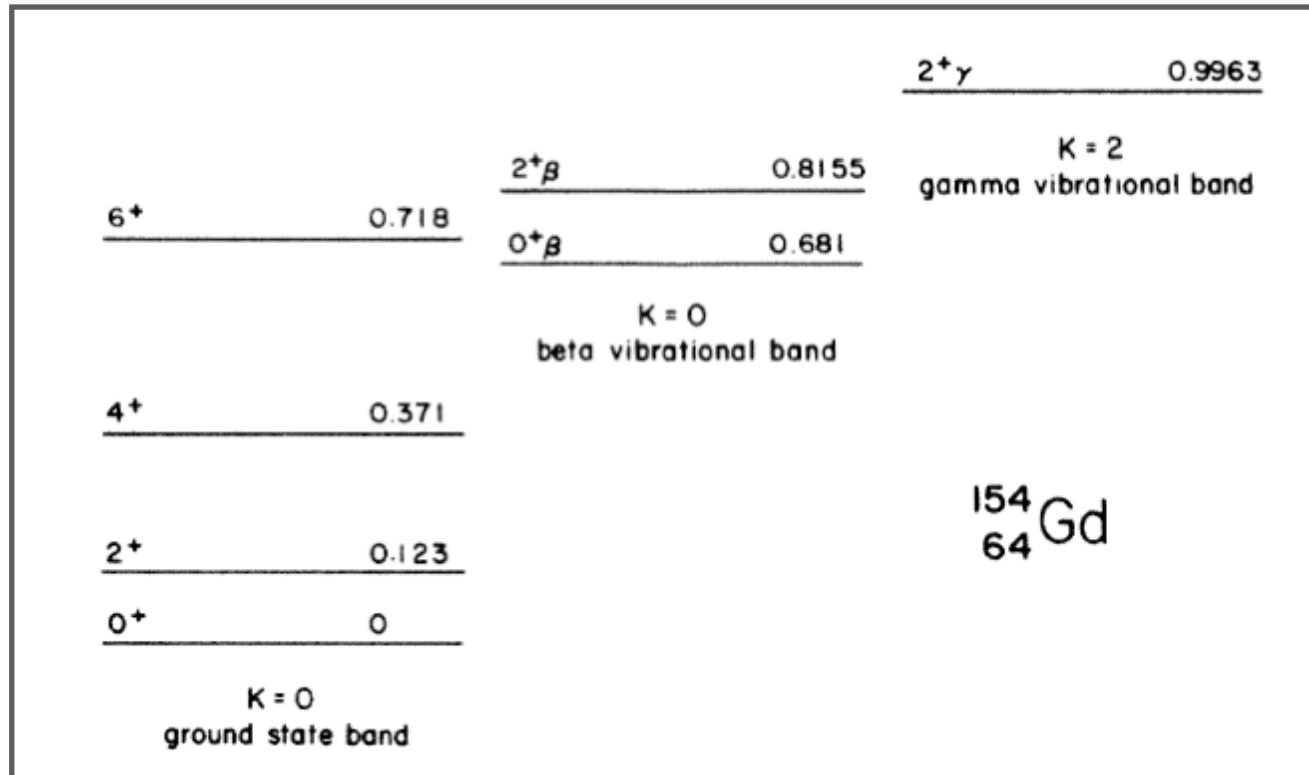


FIG. 5. Energy loss spectra of electrons scattered inelastically from ¹⁵⁴Gd at two incident energies. Spectrometer momentum resolution was 6×10^{-5} .

elastic + inelastic electron scattering

Transition form factors ($0^+ \rightarrow 2^+, 4^+, 6^+$)

→ their transition densities

$$\rho_\lambda(r) = \langle \lambda || \hat{\rho} || 0 \rangle = \frac{1}{\sqrt{2I+1}} \int \rho(r, \Omega) Y_{\lambda 0}(\Omega) d\Omega$$

$$\rho(r, \Omega) = \sum_{\lambda} \rho_\lambda(r) Y_{\lambda 0}(\Omega)$$

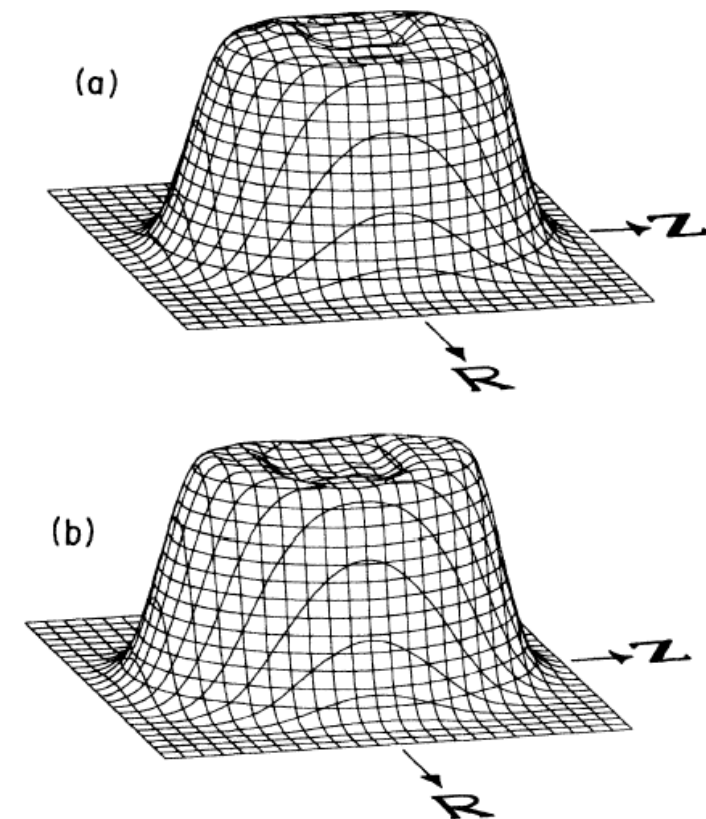


FIG. 3. Deformed charge density distributions in the body reference frame for (a) the calculated density, and (b) the density reconstructed from the experimental results. The symmetry axis and a perpendicular axis are denoted by Z and R, respectively.

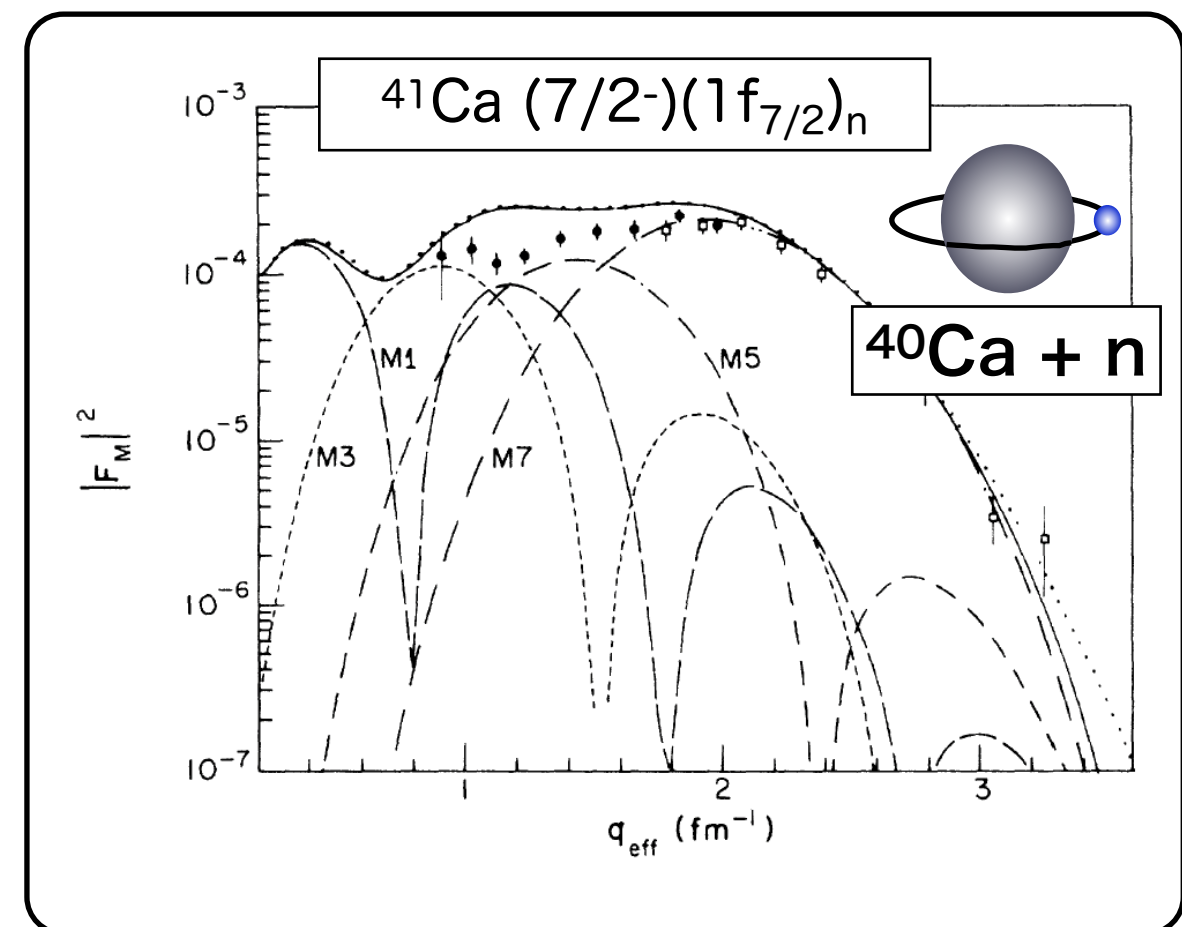
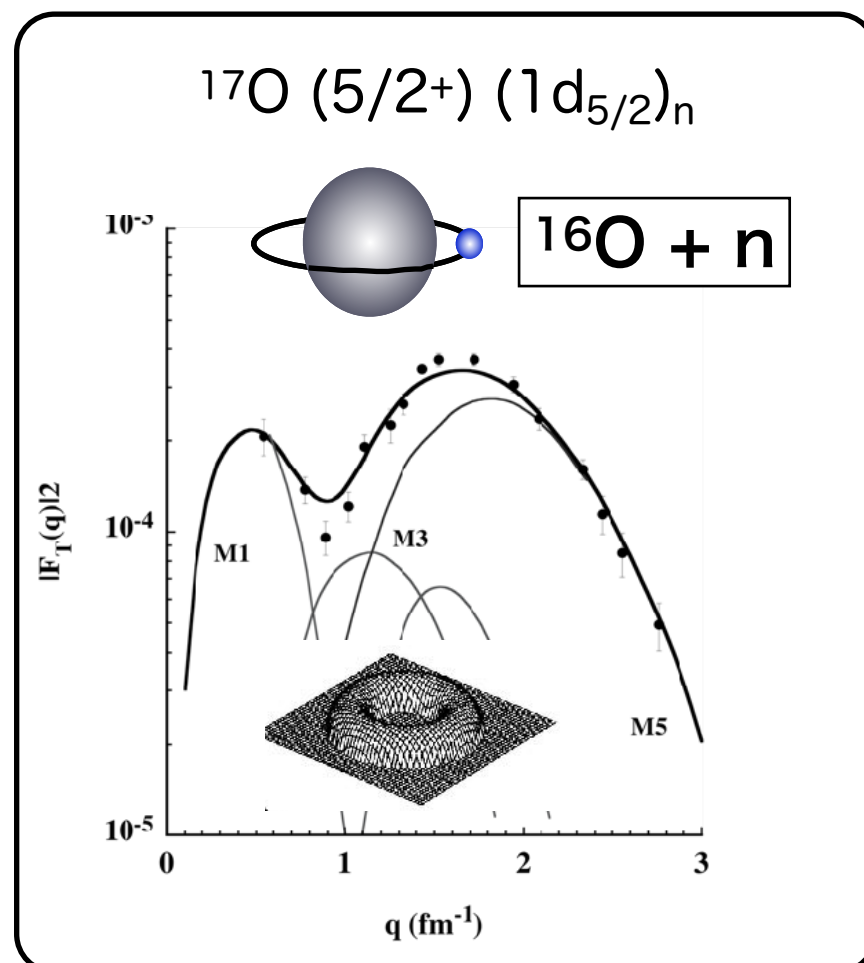
$$\frac{d\sigma_{Mott}}{d\Omega} = 0 \quad \leftarrow \quad \frac{d\sigma_{Mott}}{d\Omega} = \frac{Z^2\alpha^2}{4e^2} \frac{\cos^2 \frac{\theta}{2}}{\sin^4 \frac{\theta}{2}}$$

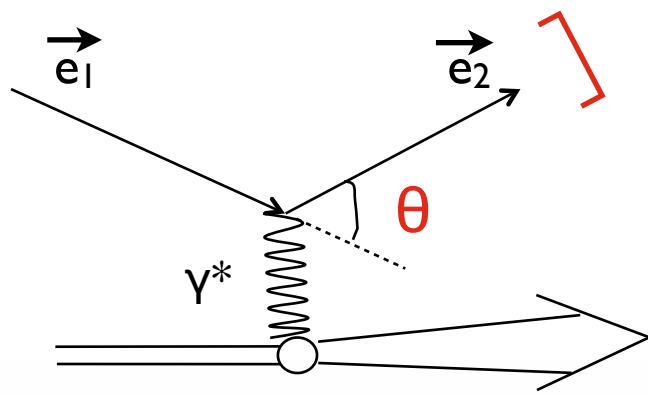
$$\frac{d\sigma}{d\Omega}(\theta = 180^\circ) \propto \frac{\cancel{\cos^2 \frac{\theta}{2}}}{\sin^4 \frac{\theta}{2}} \times \frac{\sin^2 \frac{\theta}{2}}{\cancel{\cos^2 \frac{\theta}{2}}} | \langle J_0 || T_J^M(q) || J_0 \rangle |^2$$

Only magnetic term survives at 180°

Core + valence neutron : ^{17}O , ^{41}Ca

electron scatters off the valence neutron via its anomalous magnetic moment





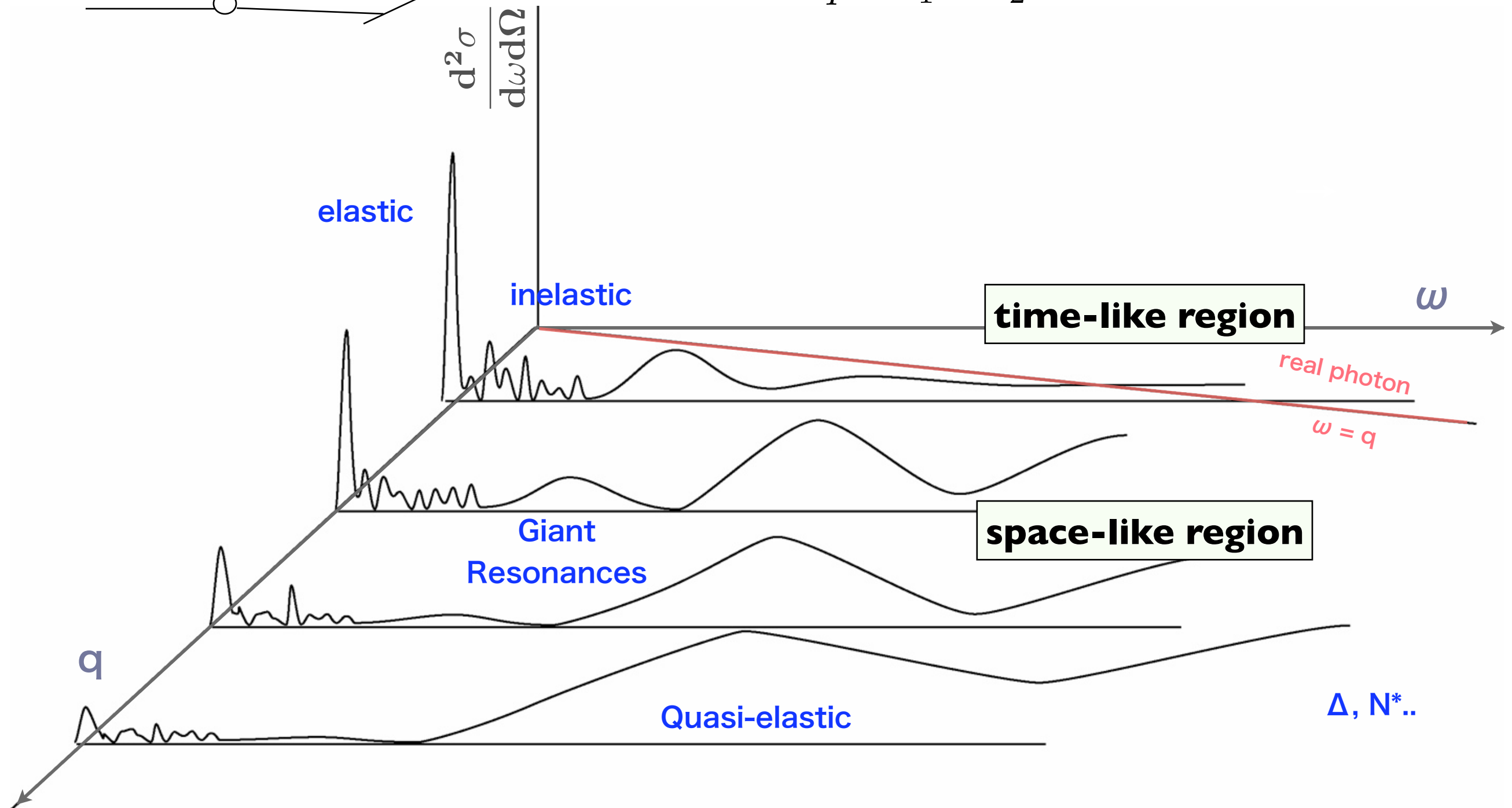
kinematical variables

$$\omega = e_1 - e_2$$

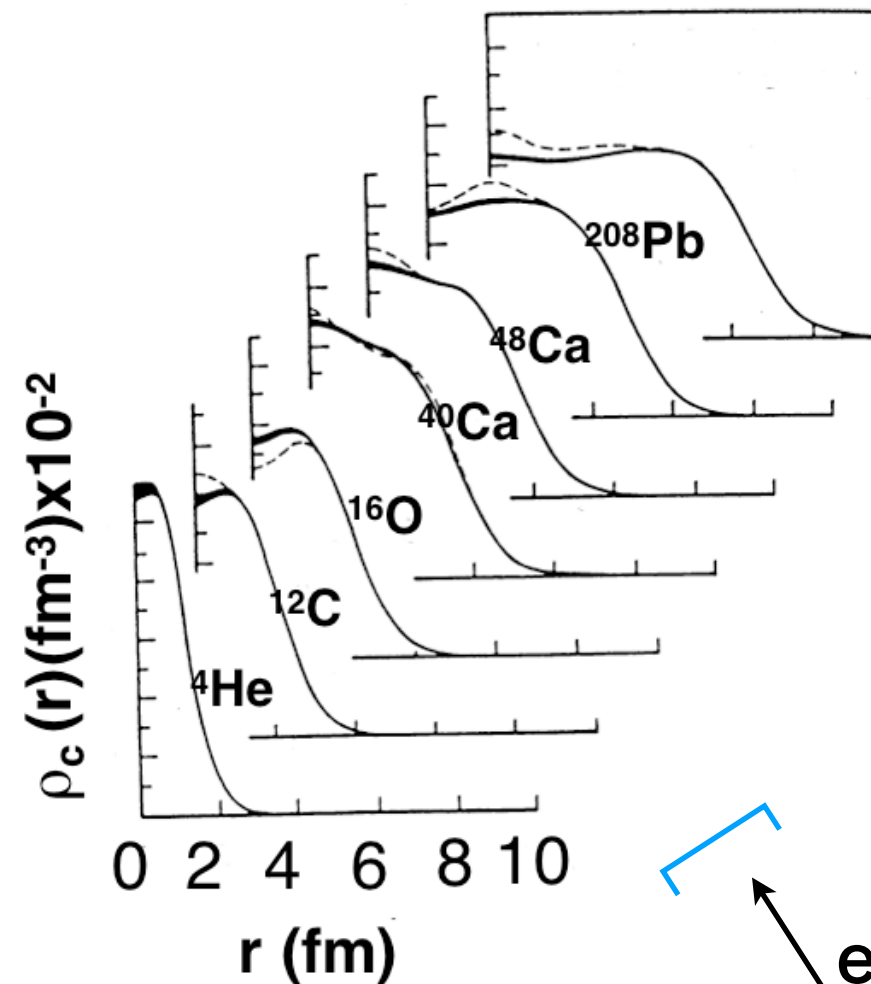
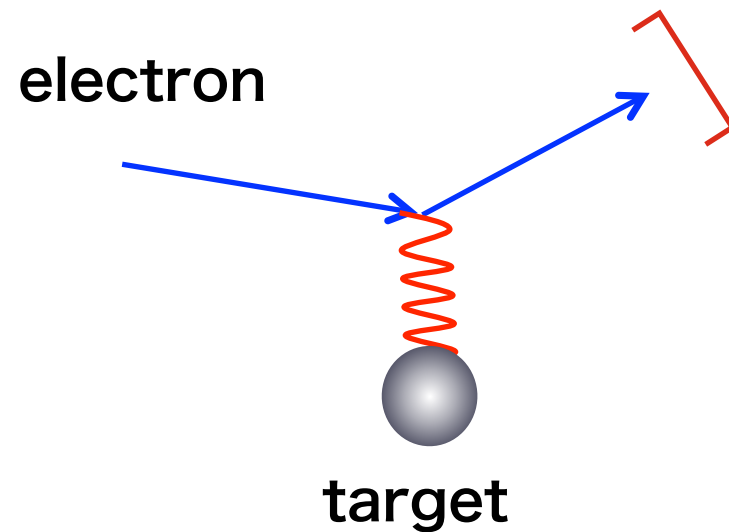
energy transfer

$$\vec{q} = \vec{e}_1 - \vec{e}_2$$

momentum transfer



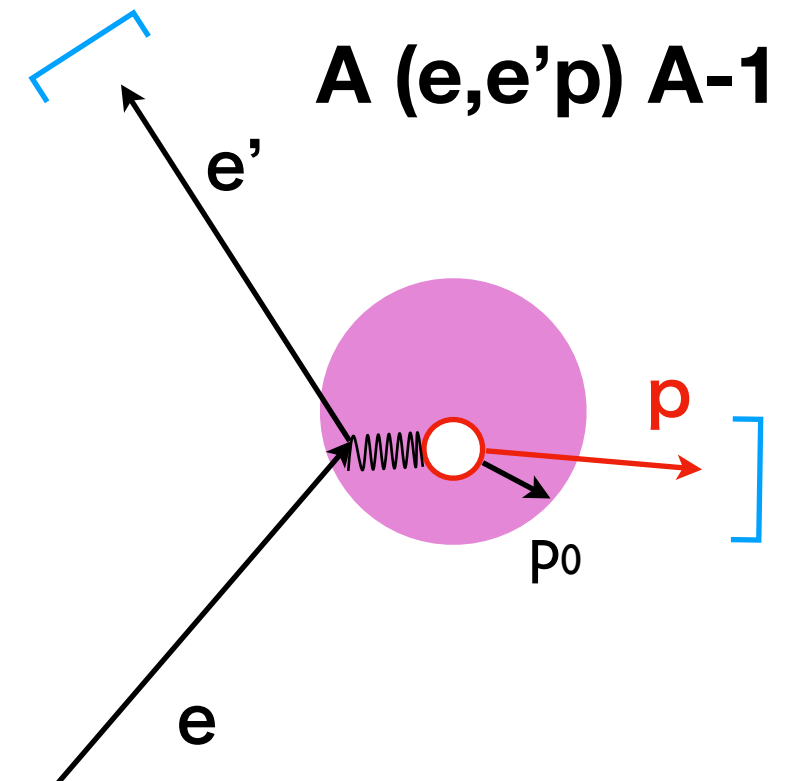
elastic scattering : $e+A$



quasi-elastic scattering

“elastic scattering” off proton in a nucleus

$e + \text{“p”} \rightarrow e + \text{“p”}$ (the residual is a spectator)



one studies protons inside a nucleus (momentum, energy ...)

under PWIA

$$\frac{d^6\sigma}{dk_e' d\Omega_{e'} dp_p d\Omega_p} = K \frac{d\sigma_{ep}}{d\Omega} S(E_m, \vec{p}_m)$$

$$S(E_m, \vec{p}_m) = \rho(\vec{p}_m) n_\alpha \delta(E_m - E_{tr})$$

$$\rho(\vec{p}_m) n_\alpha = \int S(E_m, \vec{p}_m) dE_m$$

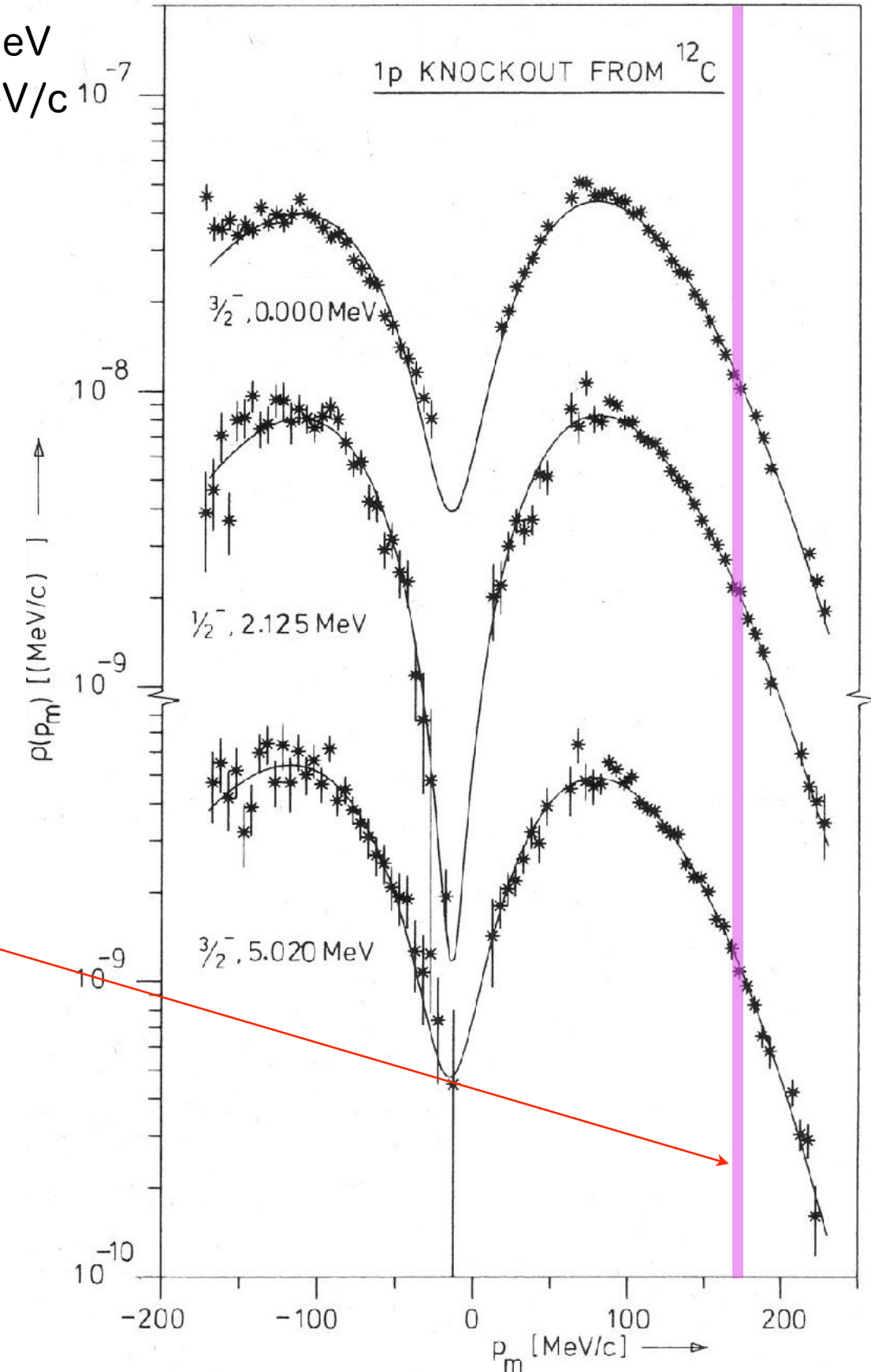
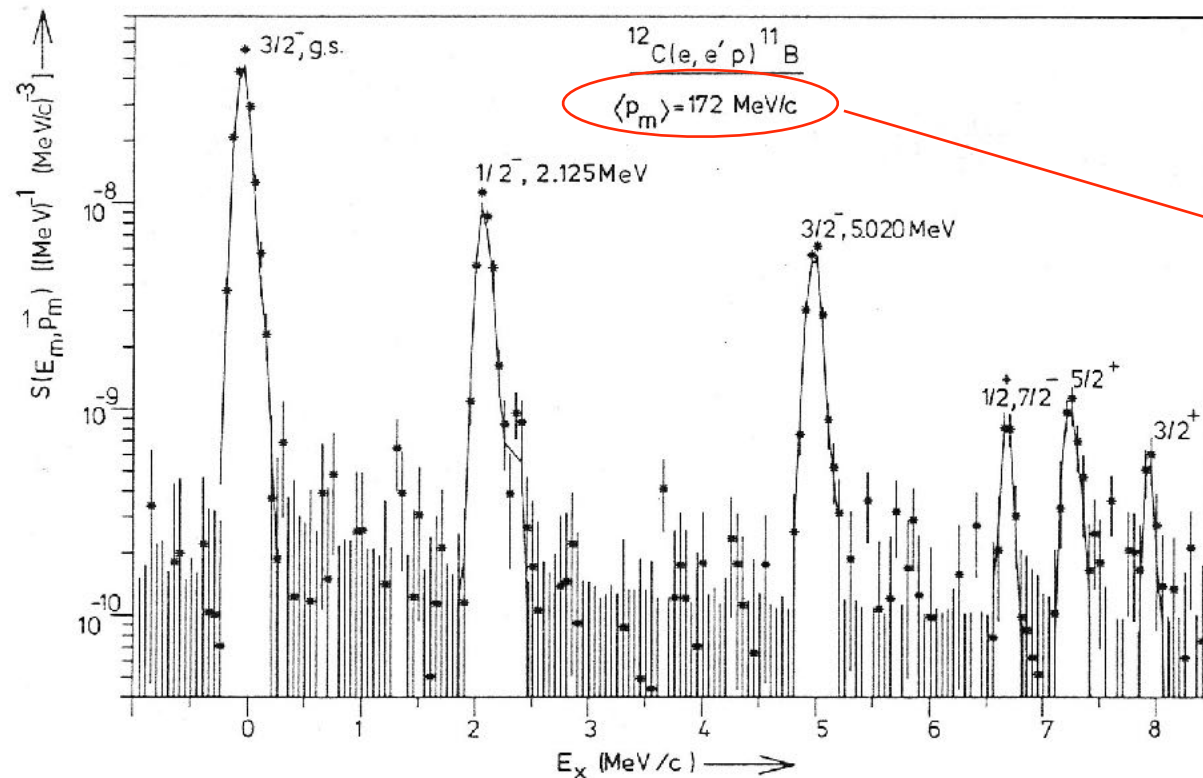
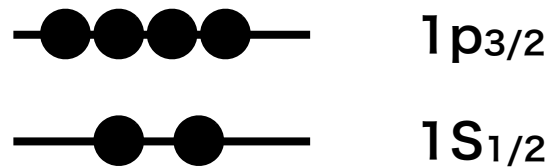
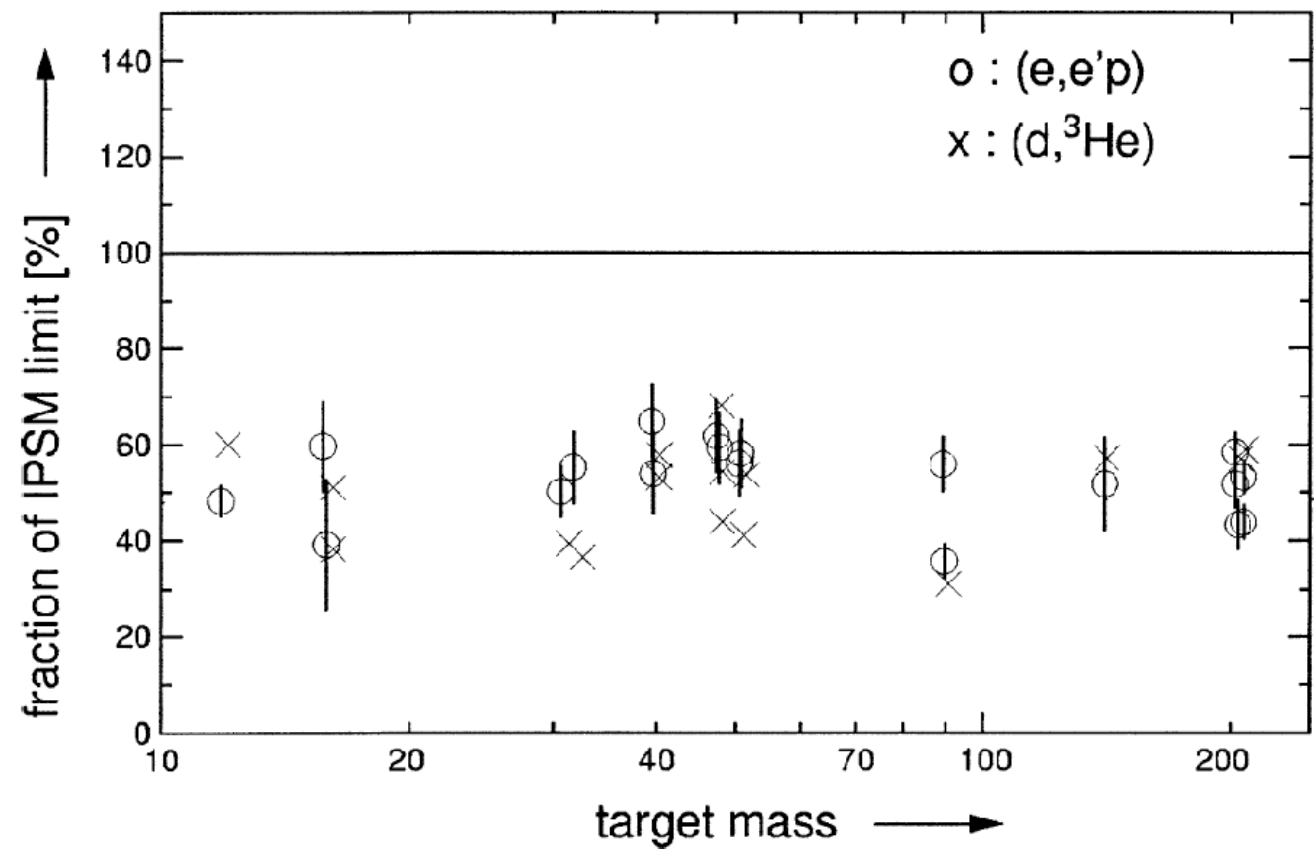
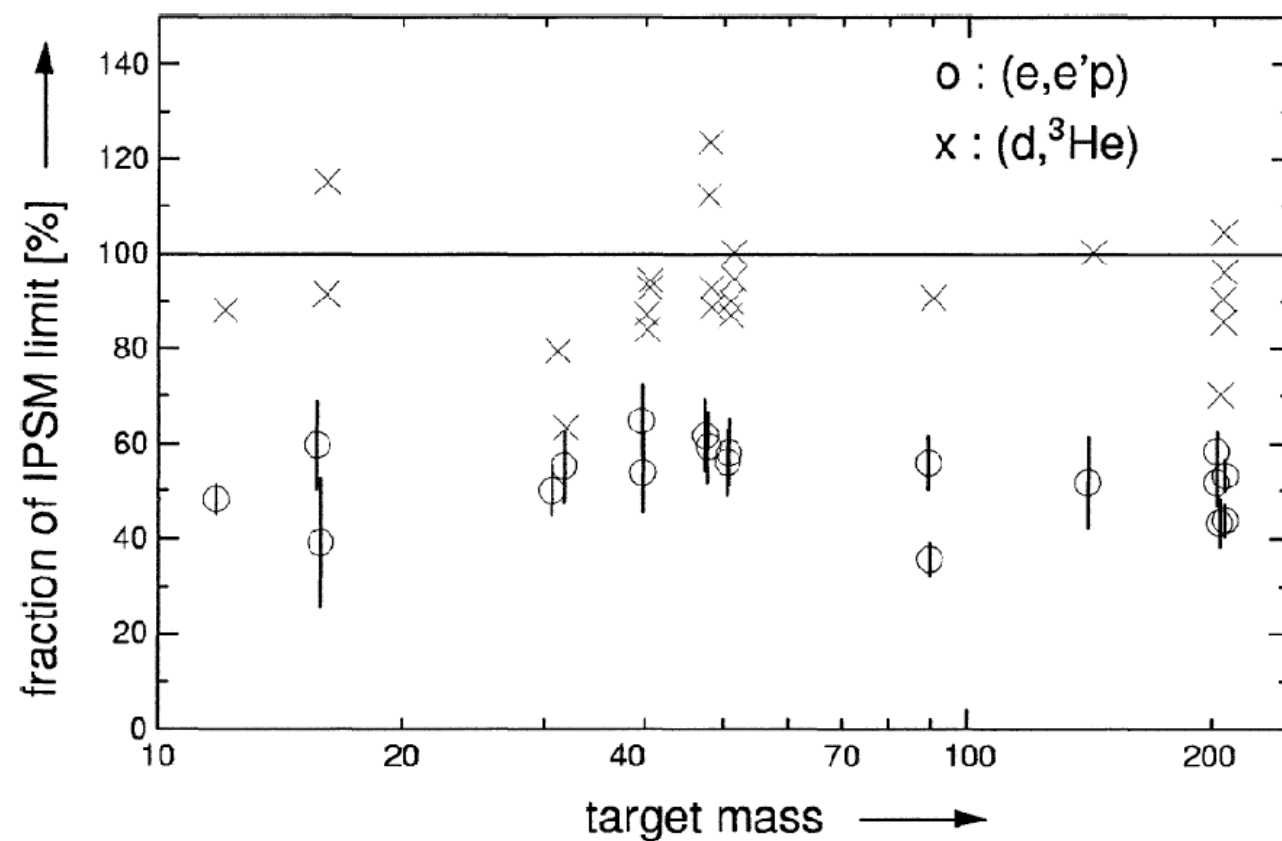


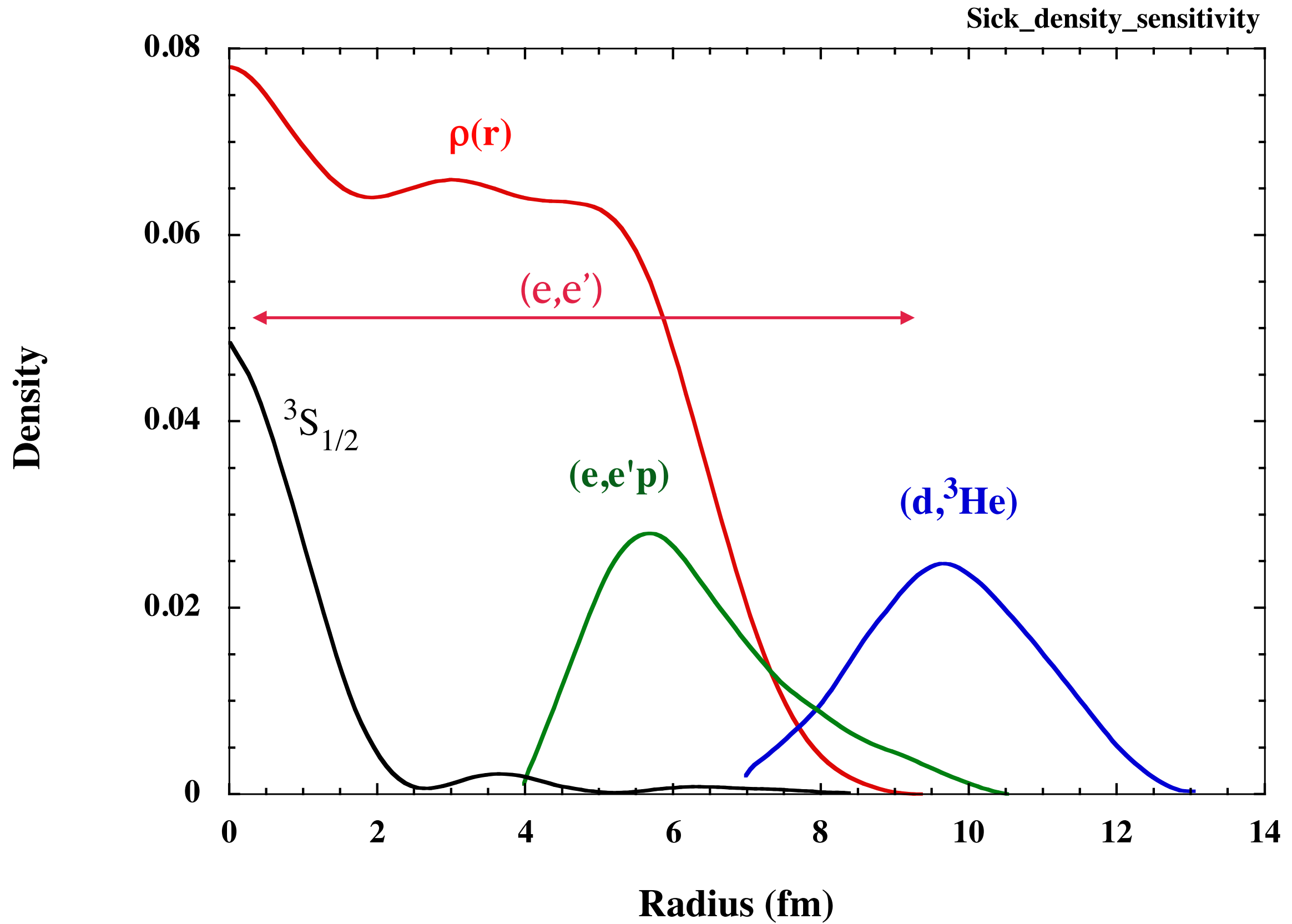
Fig. 1. Excitation-energy spectrum of the $^{12}\text{C}(e,e'p)^{11}\text{B}$ reaction at $E_0 = 284.5 \text{ MeV}$. The central value of the missing momentum is 172 MeV/c .

$$\frac{d^6\sigma}{dk_{e'}d\Omega_{e'}dp_pd\Omega_p} = K \frac{d\sigma_{ep}}{d\Omega} S(E_m, \vec{p}_m)$$

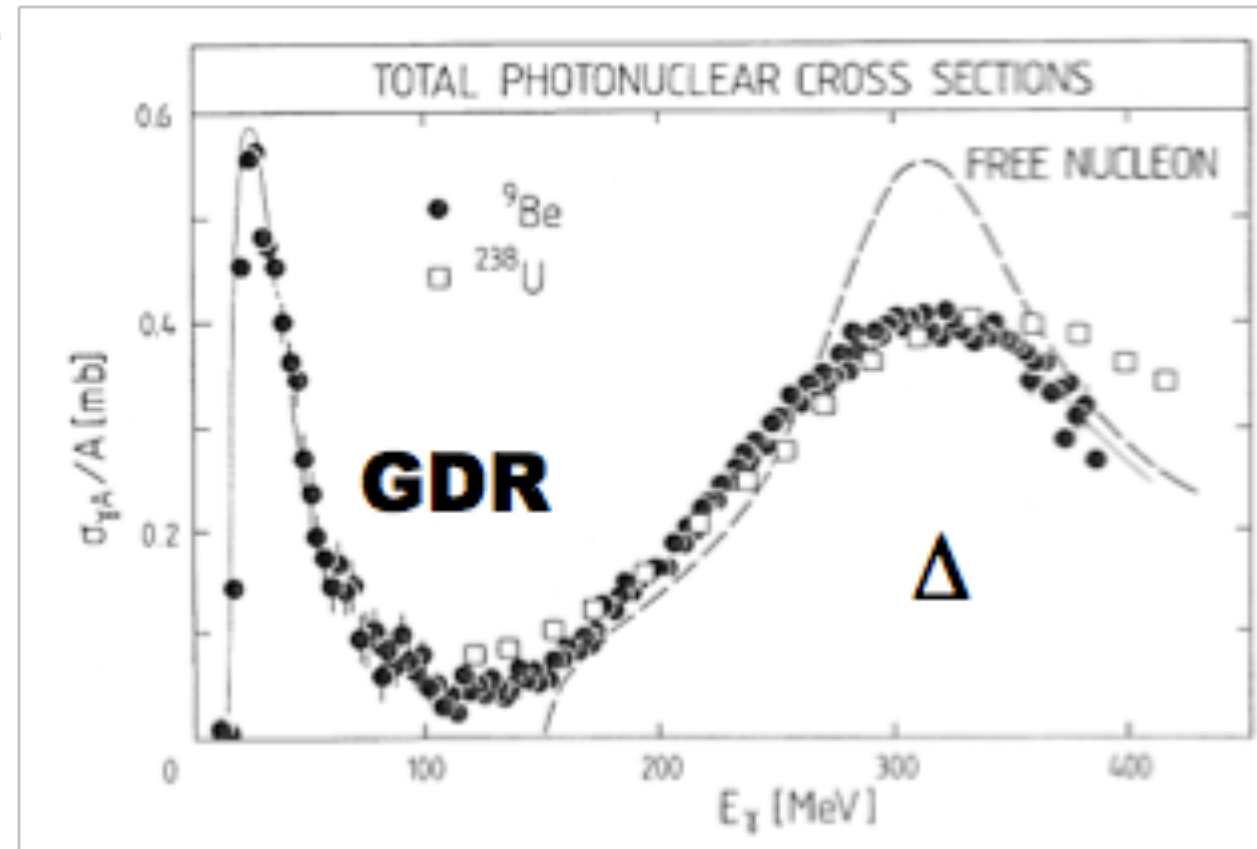
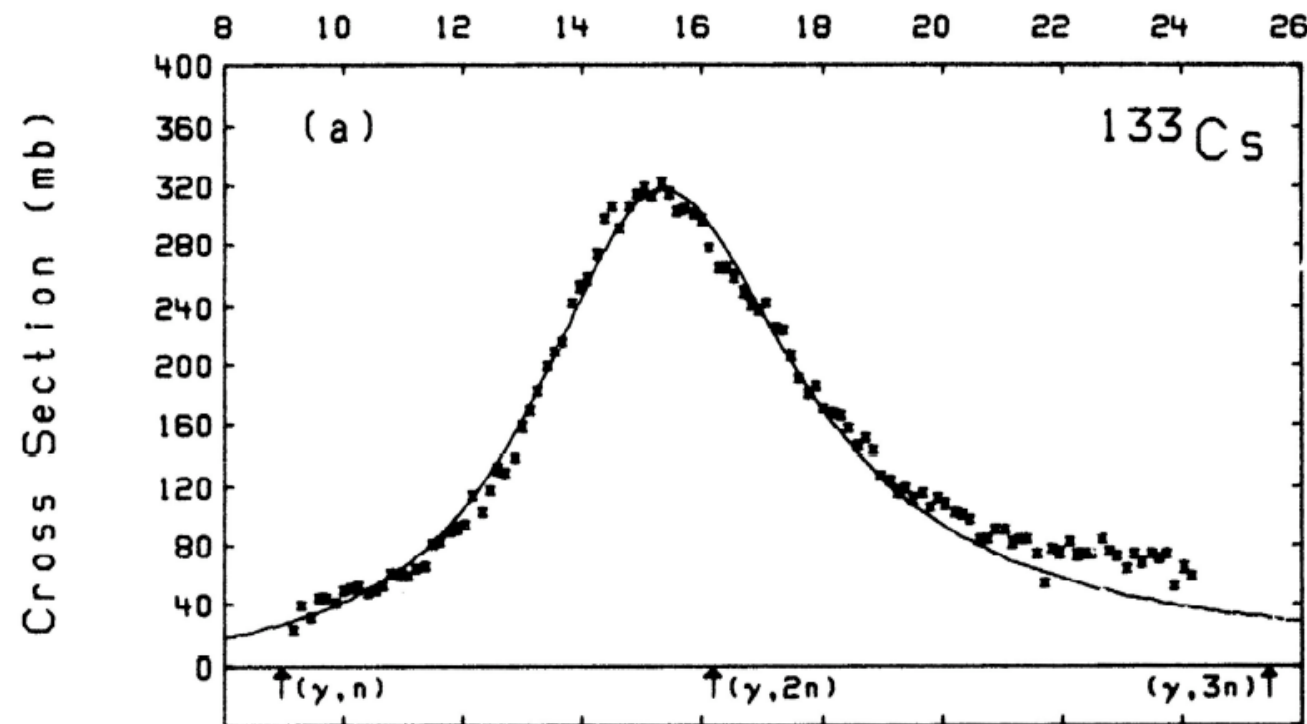
$$S(E_m, \vec{p}_m) = \rho(\vec{p}_m) n_\alpha \delta(E_m - E_{tr})$$

$$\rho(\vec{p}_m) n_\alpha = \int S(E_m, \vec{p}_m) dE_m$$





Photonuclear reaction



1) Response functions (operators : well-known)

2) Sum Rules

TRK sum rule

$$\int_0^\infty \sigma(E_\gamma) dE_\gamma = \frac{2\pi^2 e^2 \hbar}{M} \frac{NZ}{A} (1 + \kappa) = 60 \frac{NZ}{A} (1 + \kappa) \text{MeV} \cdot \text{mb}$$

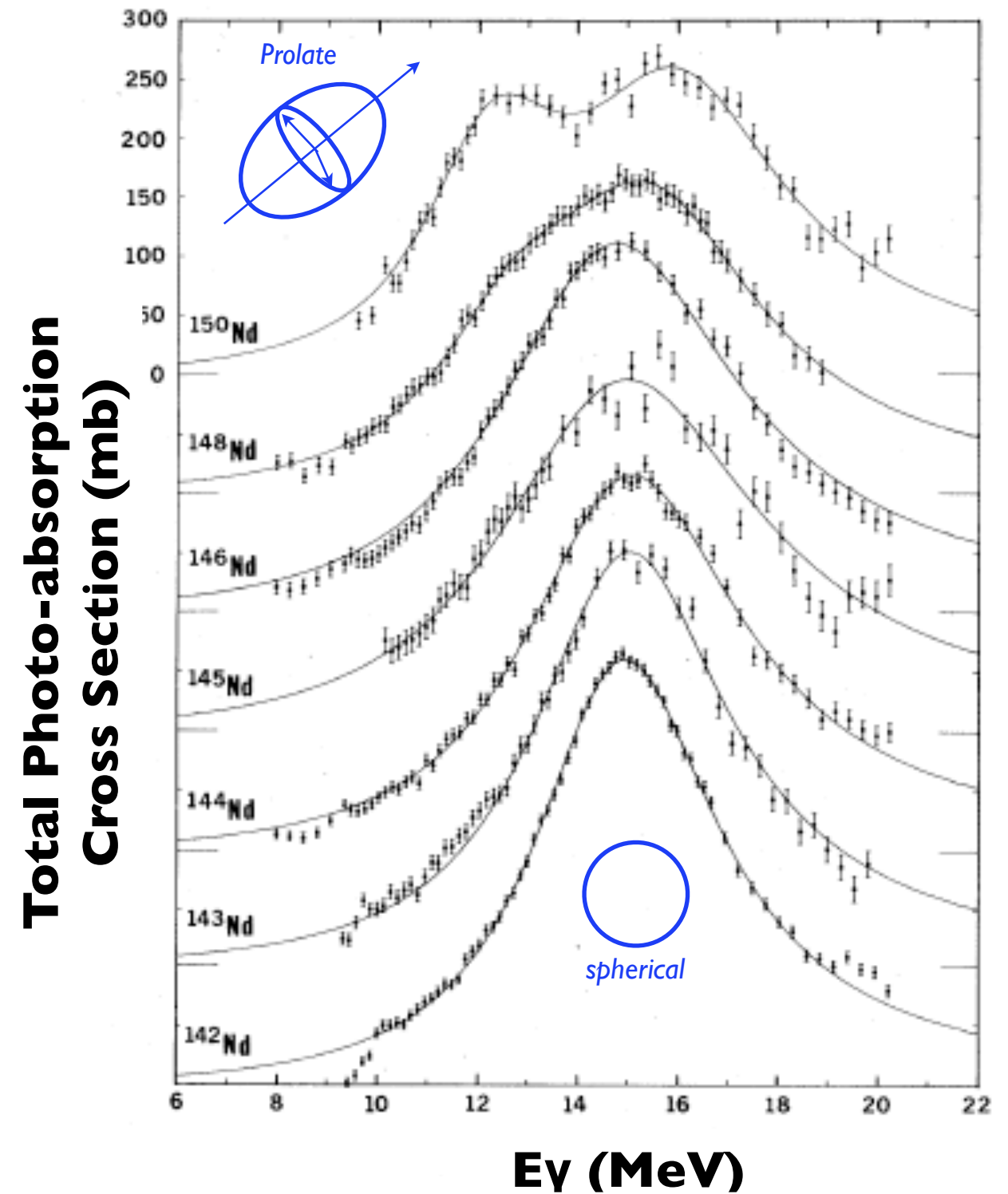
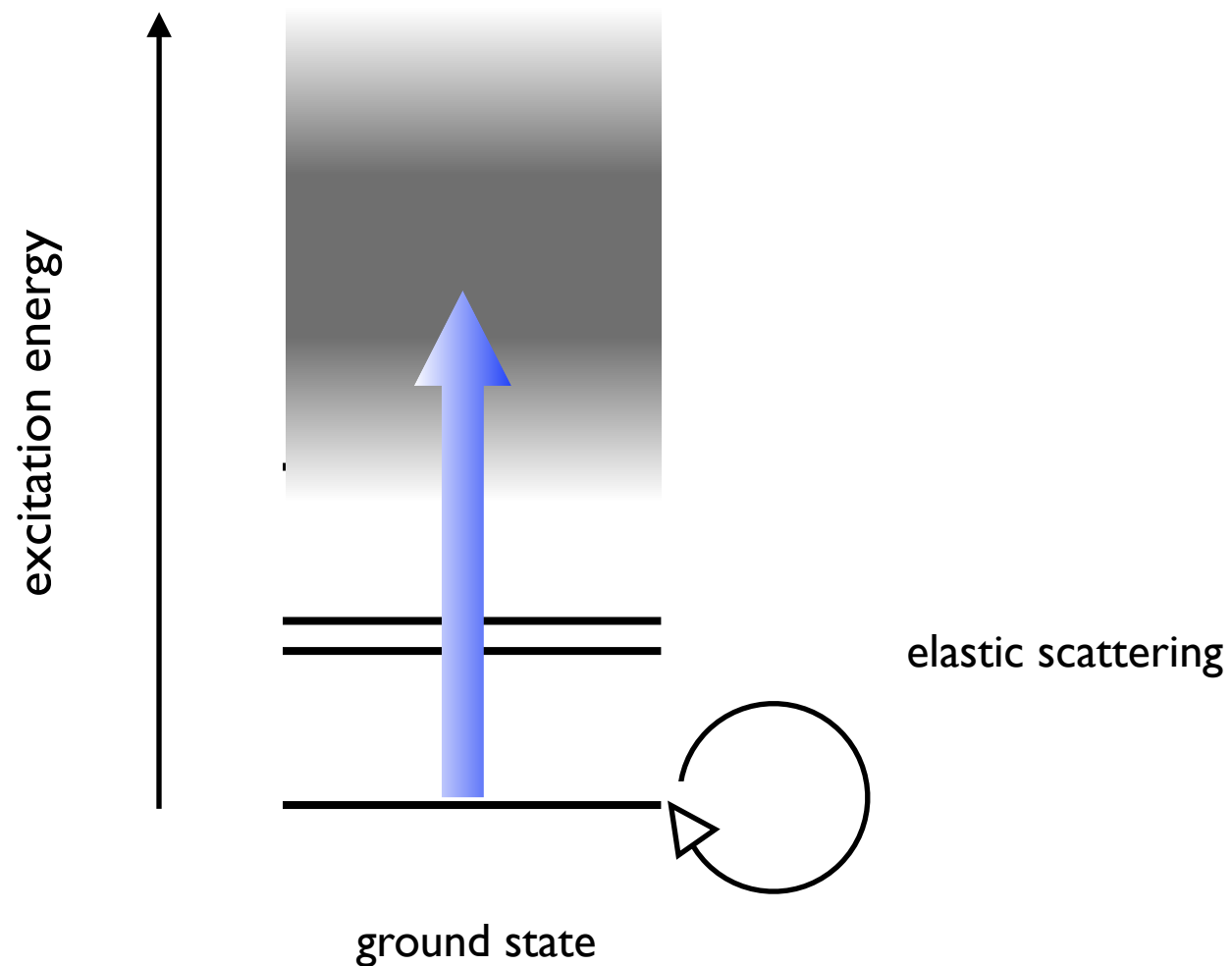
Bremmstrahlung sum rule

$$\int_0^\infty \frac{\sigma(E_\gamma)}{E_\gamma} dE_\gamma = \frac{4\pi^2 e^2}{3\hbar} \frac{NZ}{A-1} \langle r^2 \rangle$$

Migdal sum rule

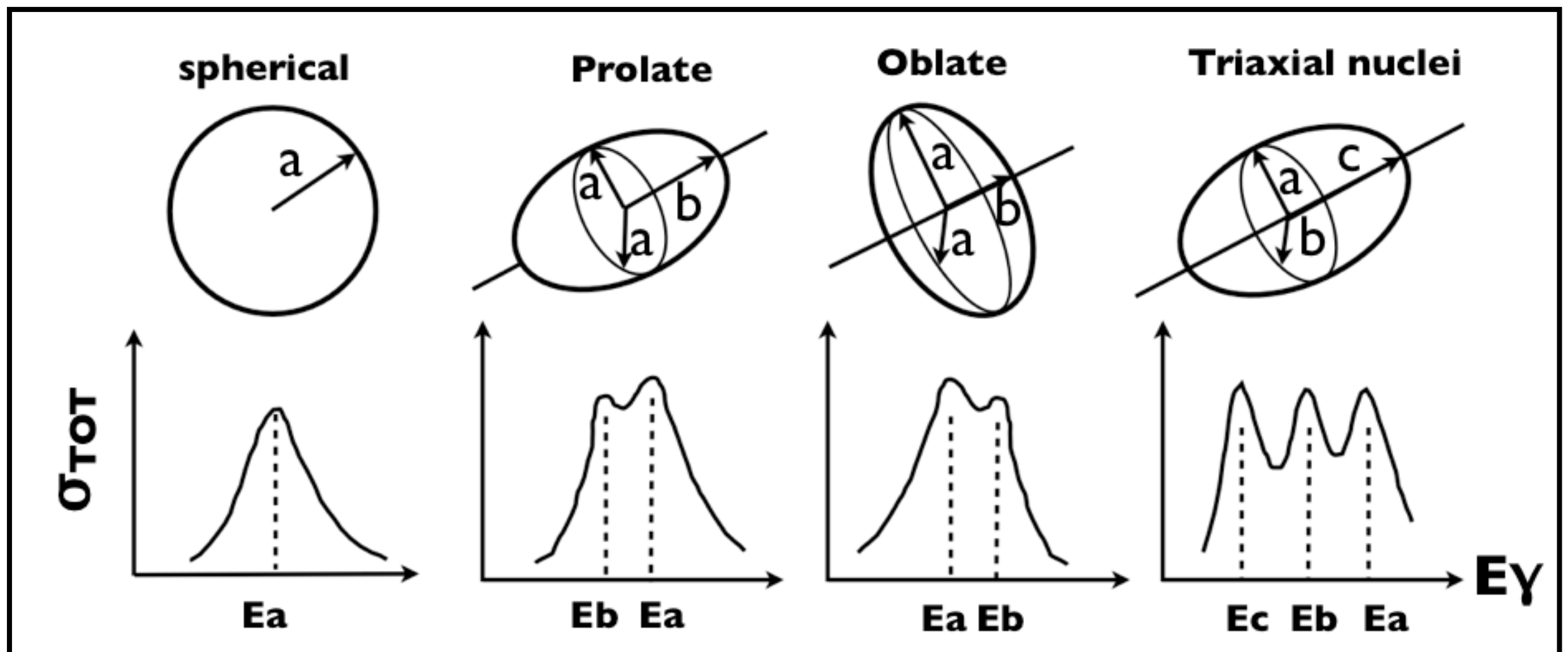
$$\int_0^\infty \frac{\sigma(E_\gamma)}{E_\gamma^2} dE_\gamma = \frac{2\pi^2}{\hbar} P$$

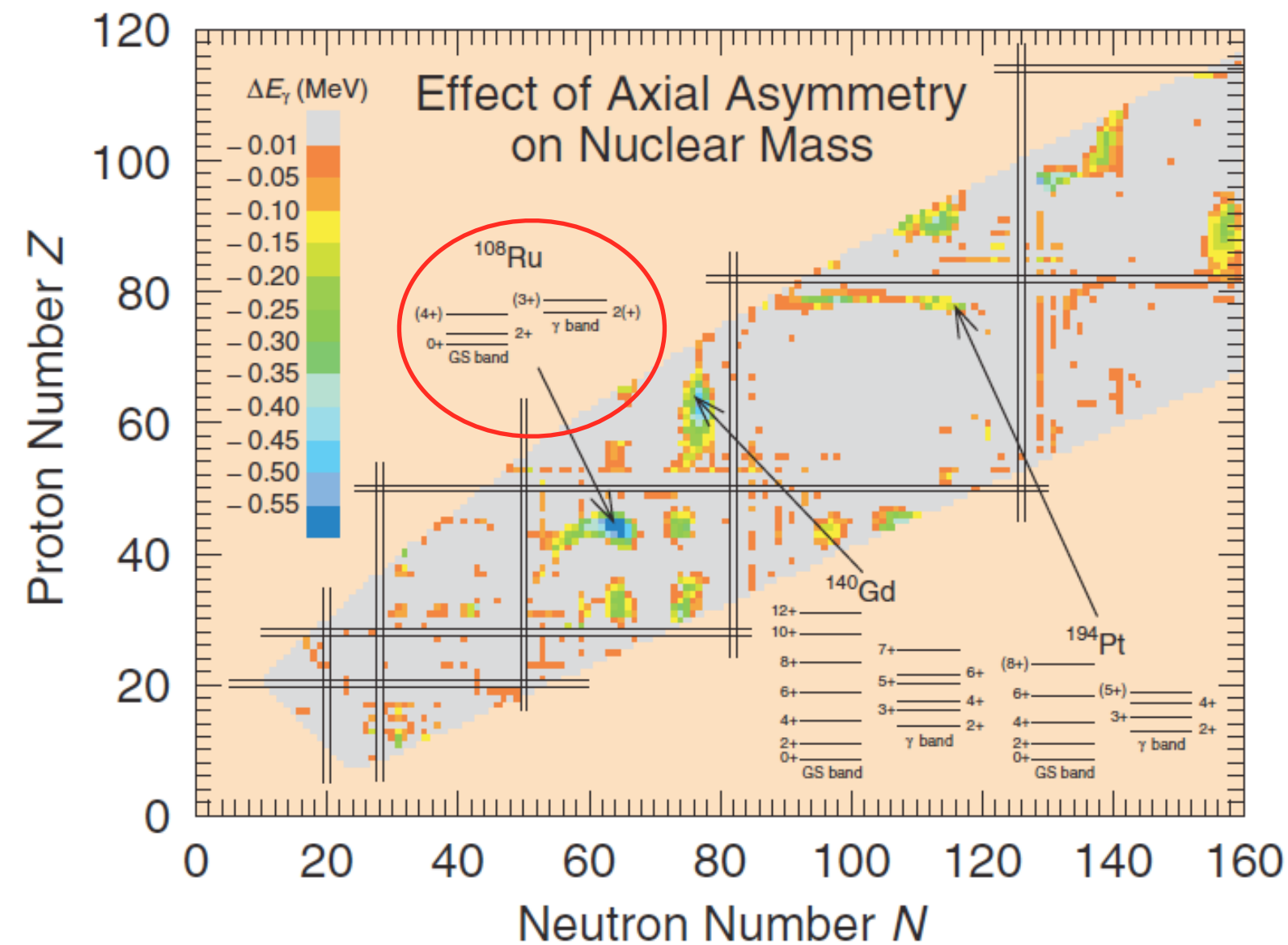
P : polarizability



Total photo-absorption cross section covering GDR

$$\gamma + A \rightarrow X$$

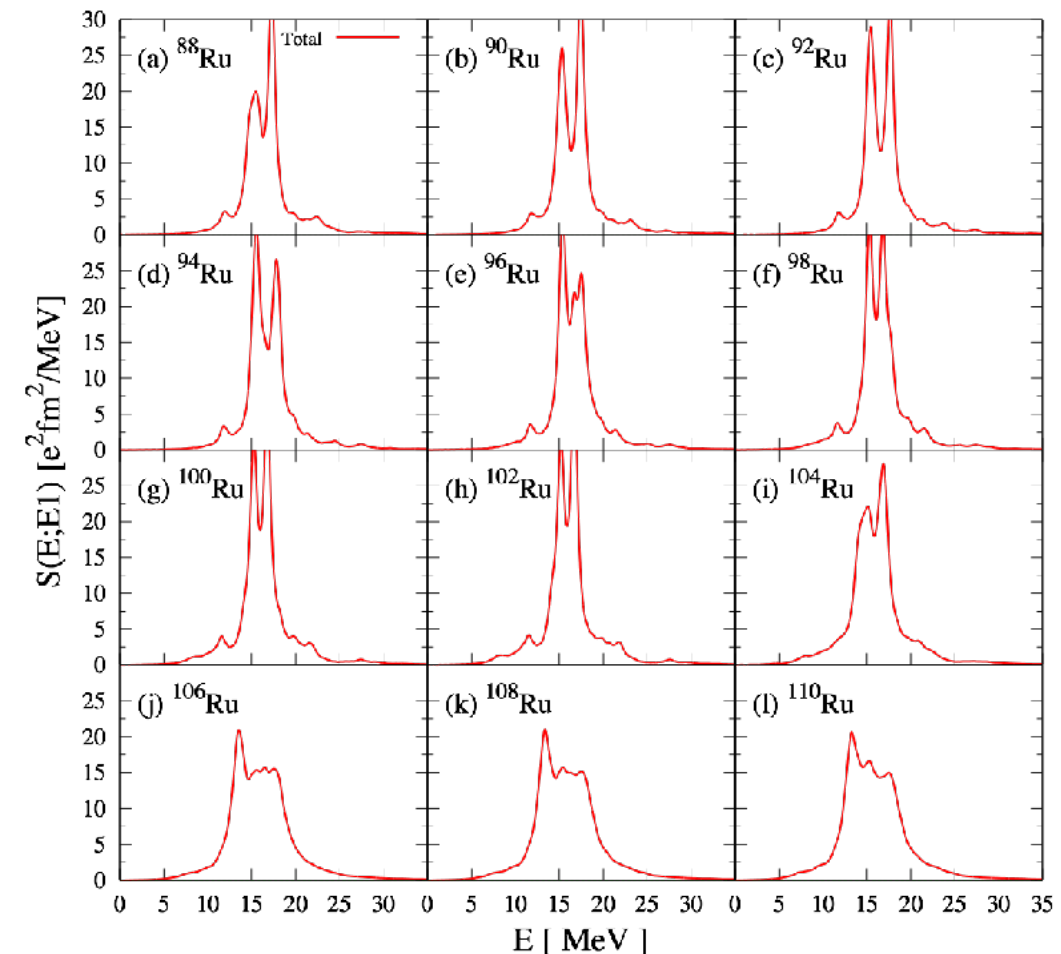
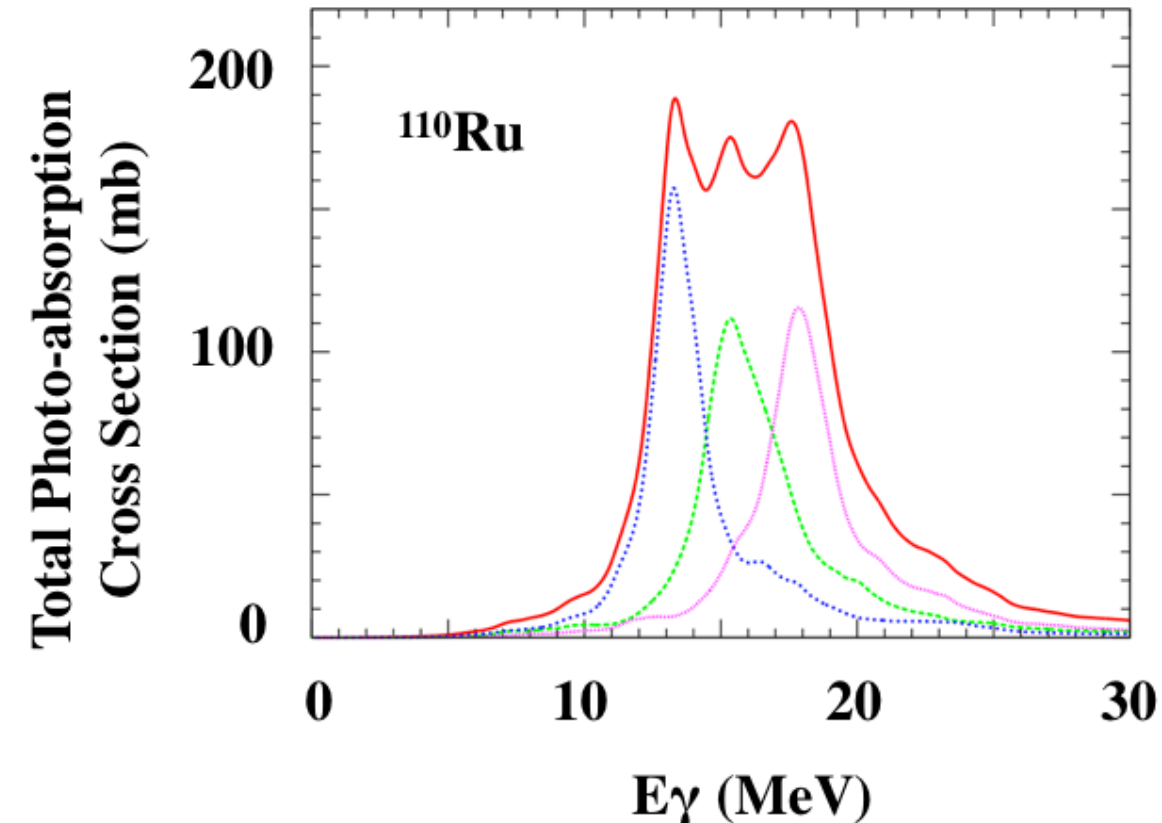




P. Moller et al., PRL97, 162502 (2006)

^{108}Ru 273 sec

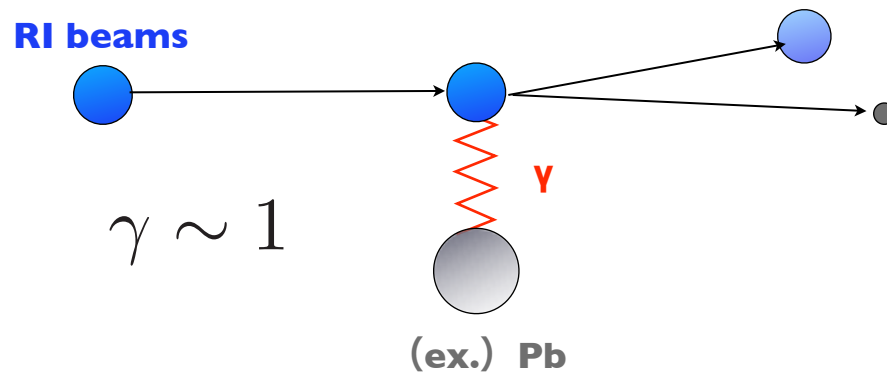
^{110}Ru 15 sec



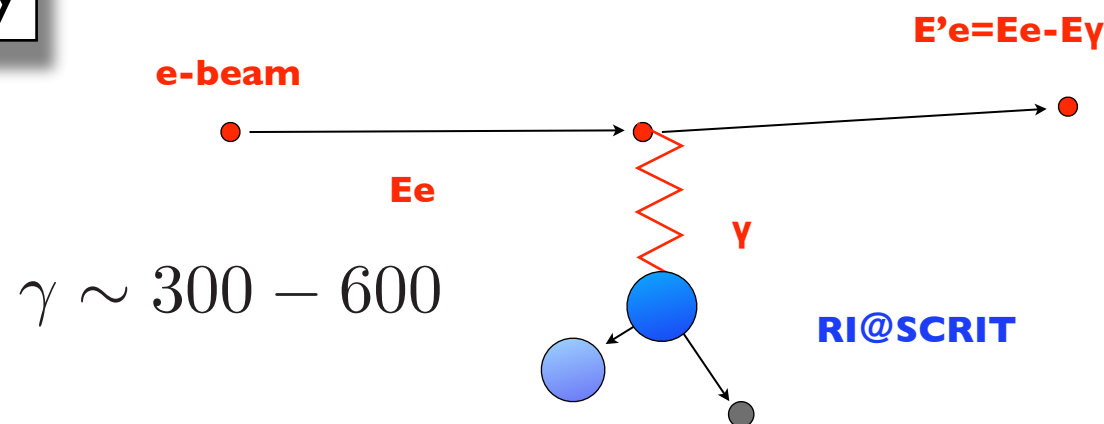
S. Ebata : private communication

so far

only way : Coulomb excitation in heavy ion reaction

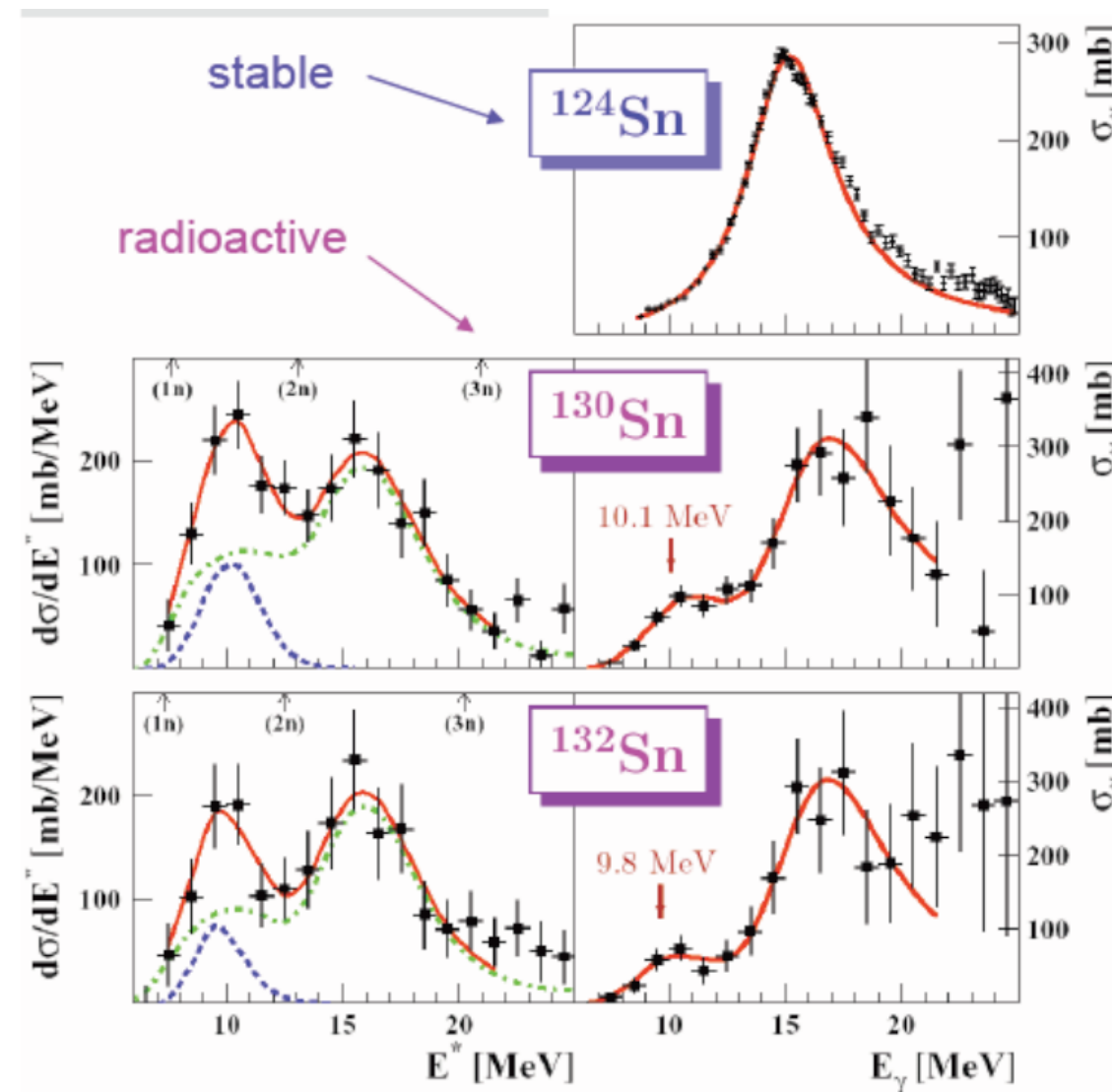


SCRIT facility



purely EM probe
well under control
negligible multi-stop
ultra-forward
electron scattering

$^{132}\text{Sn} + \text{Pb} \rightarrow ^{131}\text{Sn} + n + X$ @ GSI



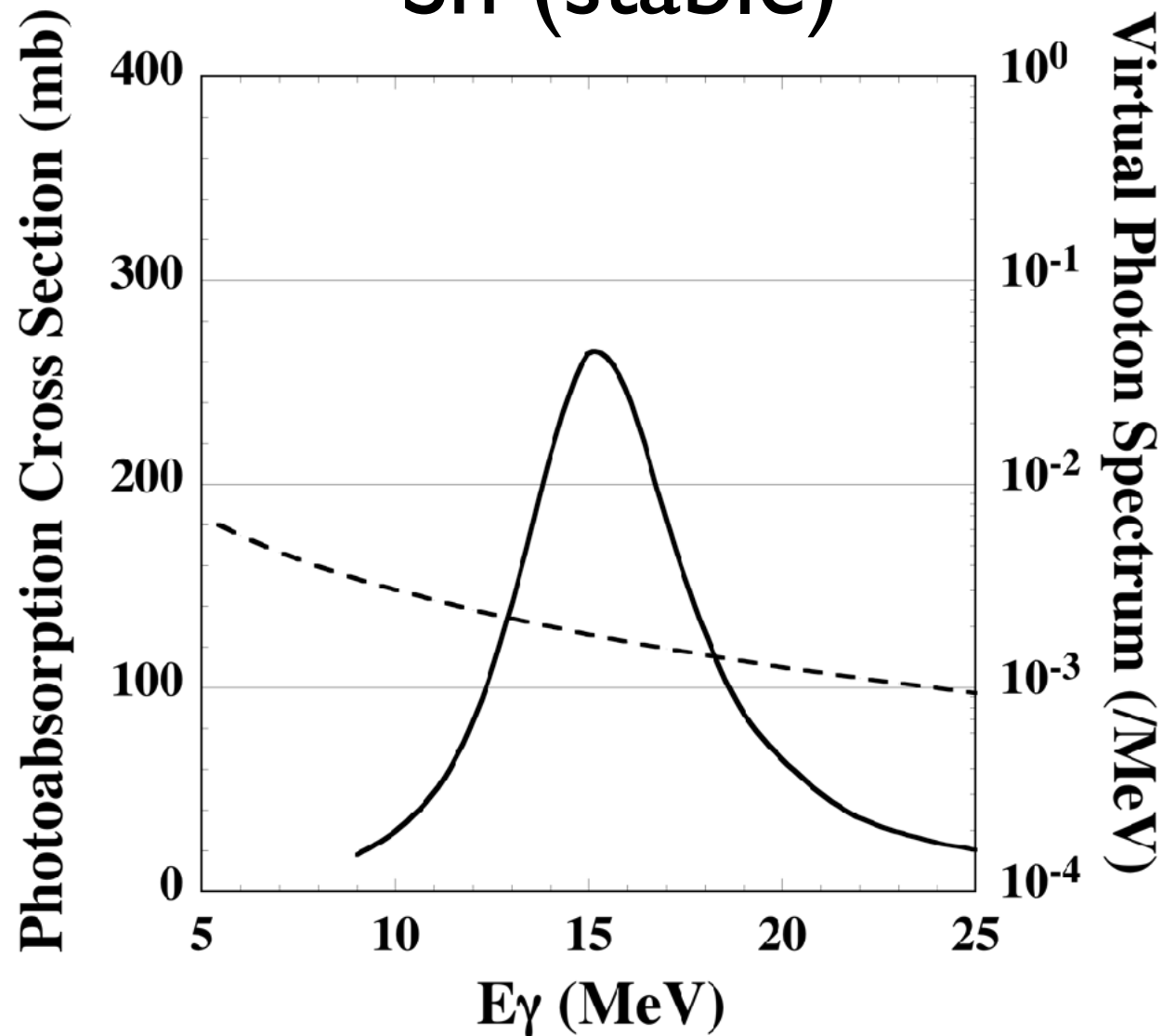
Virtual Photon flux

$$\frac{d^2\sigma}{dE_e d\Omega} = \sum \frac{d^2 N_e^{EL}(E, E_\gamma, \theta)}{dE_\gamma d\Omega} \cdot \sigma_\gamma^{EL}(E_\gamma)$$

virtual photon theory

$$\frac{dN}{dE_\gamma} = L \cdot \int d\Omega \frac{d^2 N_e^{E1}(E, E_\gamma, \theta)}{dE_\gamma d\Omega} \cdot \sigma_\gamma^{E1}(E_\gamma)$$

^{120}Sn (stable)



P. Durgapal and D.S. Onley
Comp. Phys. Comm. 32 (1984) 291

