



東南大學·物理學院
SCHOOL OF PHYSICS SOUTHEAST UNIVERSITY

Strong decays of the $\Lambda_c(2910)$ and $\Lambda_c(2940)$ in D^*N molecular picture

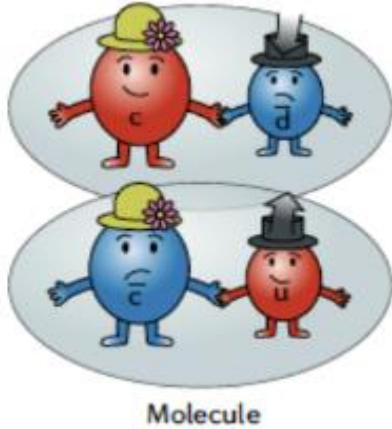
Reporter: Zi-Li Yue

2023.12.17

全国第十二届重味物理和CP破坏研讨会

Based on Zi-Li Yue, Dian-Yong Chen, ...to appear

- **Background**
- **Molecular Scenario**
- **The Strong Decay of the Λ_c**
- **Numerical Results and Discussions**
- **Summary**



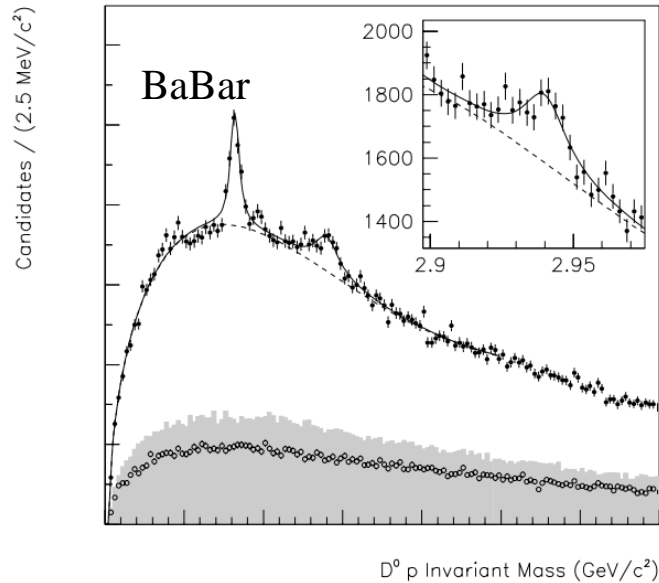
Conventional hadrons

Meson($q\bar{q}$), Baryon(qqq)

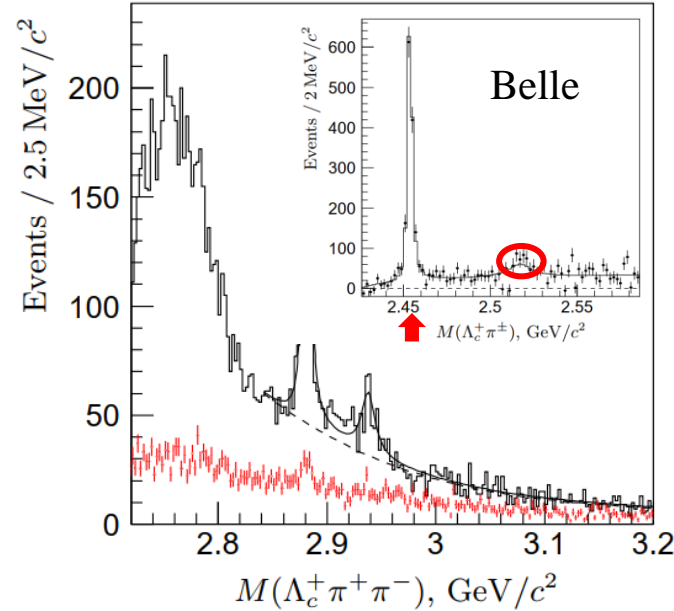
Exotic states

- **Molecular states:** loosely bound states composed of a pair of hadrons, bound by residual interactions of strong interactions. (Candidates: $X(3872)$, $D_{s1}(2460)$, $X(3915)$...)
- Tetraquarks: $qq\bar{q}\bar{q}$
- Pentaquarks: $qqqq\bar{q}$
- Hybrids: $q\bar{q}g$
- Glueballs: gg/ggg

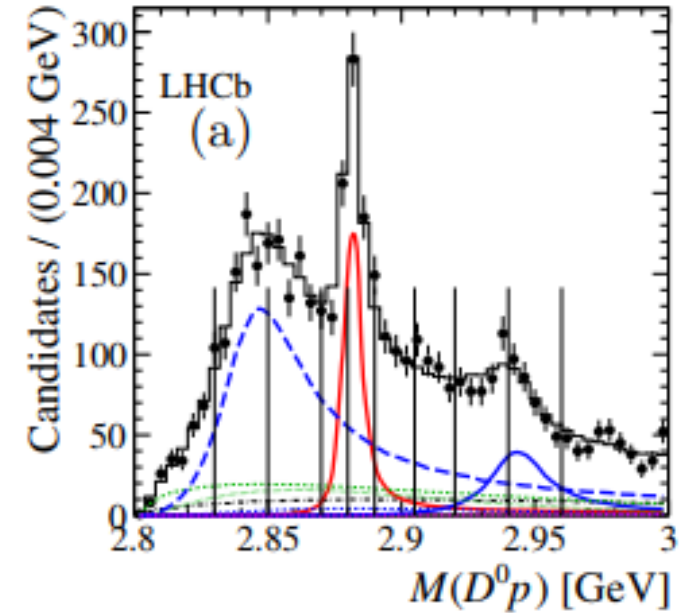
The Discovery of the $\Lambda_c(2940)^+$.



Phys.Rev.Lett. 98 (2007) 012001



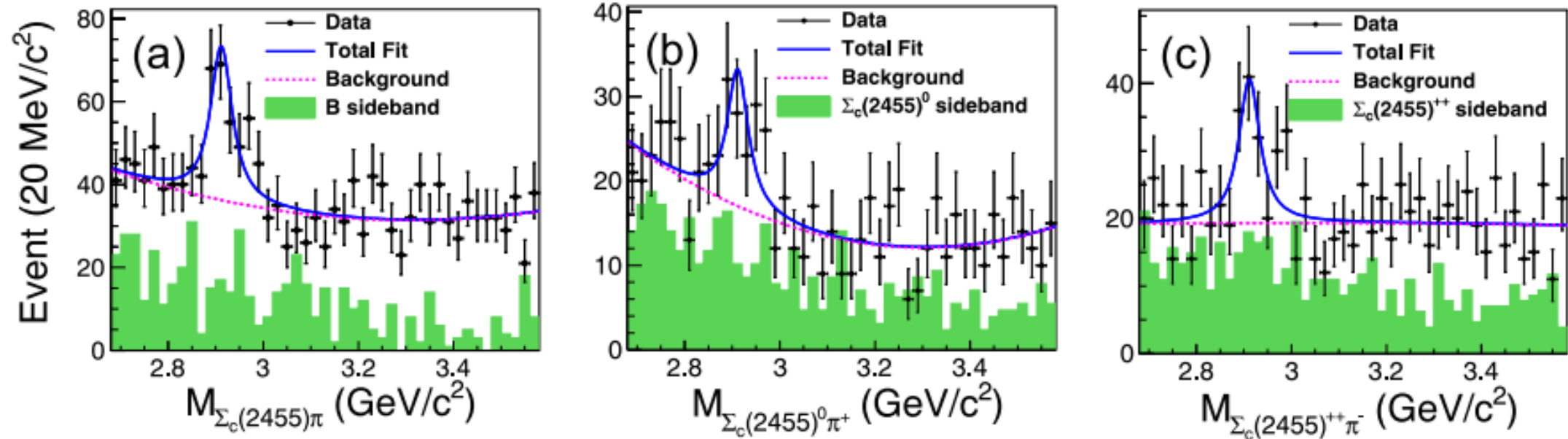
Phys.Rev.Lett. 98 (2007) 262001



JHEP 05 (2017) 030

Intrinsic parameters: $m_{\Lambda_c(2940)^+} = 2939.6_{-1.5}^{+1.3}$ MeV,
 $\Gamma_{\Lambda_c(2940)^+} = 20_{-5}^{+6}$ MeV.

$\Lambda_c(2910)^+$ in $\bar{B}^0 \rightarrow \Sigma_c(2455)^{0,++} \pi^\pm \bar{p}$ process.



Phys.Rev.Lett. 130 (2023) 3, 031901

Intrinsic parameters: $m_{\Lambda_c(2910)^+} = 2913.8 \pm 5.6 \pm 3.8$ MeV,
 $\Gamma_{\Lambda_c(2910)^+} = 51.8 \pm 20. \pm 18.8$ MeV.

Theoretical work on Λ_c states:

$m_{D^*} + m_N - m_{\Lambda_c(2910)} \sim 30 \text{ MeV}$; $m_{D^*} + m_N - m_{\Lambda_c(2940)} \sim 6 \text{ MeV}$;

D^*N molecular states:

- Eur.Phys.J.C 51 (2007) 883-889: D^*N molecular with $J^P = \frac{1}{2}^-$, decay,
- Phys.Rev.D 81 (2010) 014006,
- Phys.Rev.D 83 (2011) 094005: D^*N molecular with $J^P = \frac{1}{2}^+$, decay,
- Phys.Rev.D 107 (2023) 3, 034036: conventional baryons dressed with the D^*N channel,
- Eur.Phys.J.C 83 (2023) 6, 524: $\Lambda_c(2940)$ composed of D^*N with $J^P = \frac{3}{2}^-$...

Others explanation:

For $\Lambda_c(2910)$:

- Eur.Phys.J.C 82 (2022) 10, 920 $\Lambda_c(2P, 1/2^-)$ state,
- Phys.Rev.D 106 (2022) 7, 074020 Candidate of $\Lambda_c |J^P = 5/2^- \rangle_\rho$...

For $\Lambda_c(2940)$:

- Phys.Rev.D 75 (2007) 094017: D-wave charmed baryon with $J^P = \frac{1}{2}^+$ or $\frac{3}{2}^+$,
- Phys.Lett.B 659 (2008) 612-620: $\Lambda_c(2286)^+$ first orbital excitation...

Background



P_c states:

$P_c(4457)$: $\Sigma_c D^*$ molecular, $J^P = \frac{3}{2}^-$;

$P_c(4440)$: $\Sigma_c D^*$ molecular, $J^P = \frac{1}{2}^-$;

$P_c(4312)$: $\Sigma_c D$ molecular, $J^P = \frac{1}{2}^-$;

References:

Phys.Rev.D 100 (2019) 5, 056005,

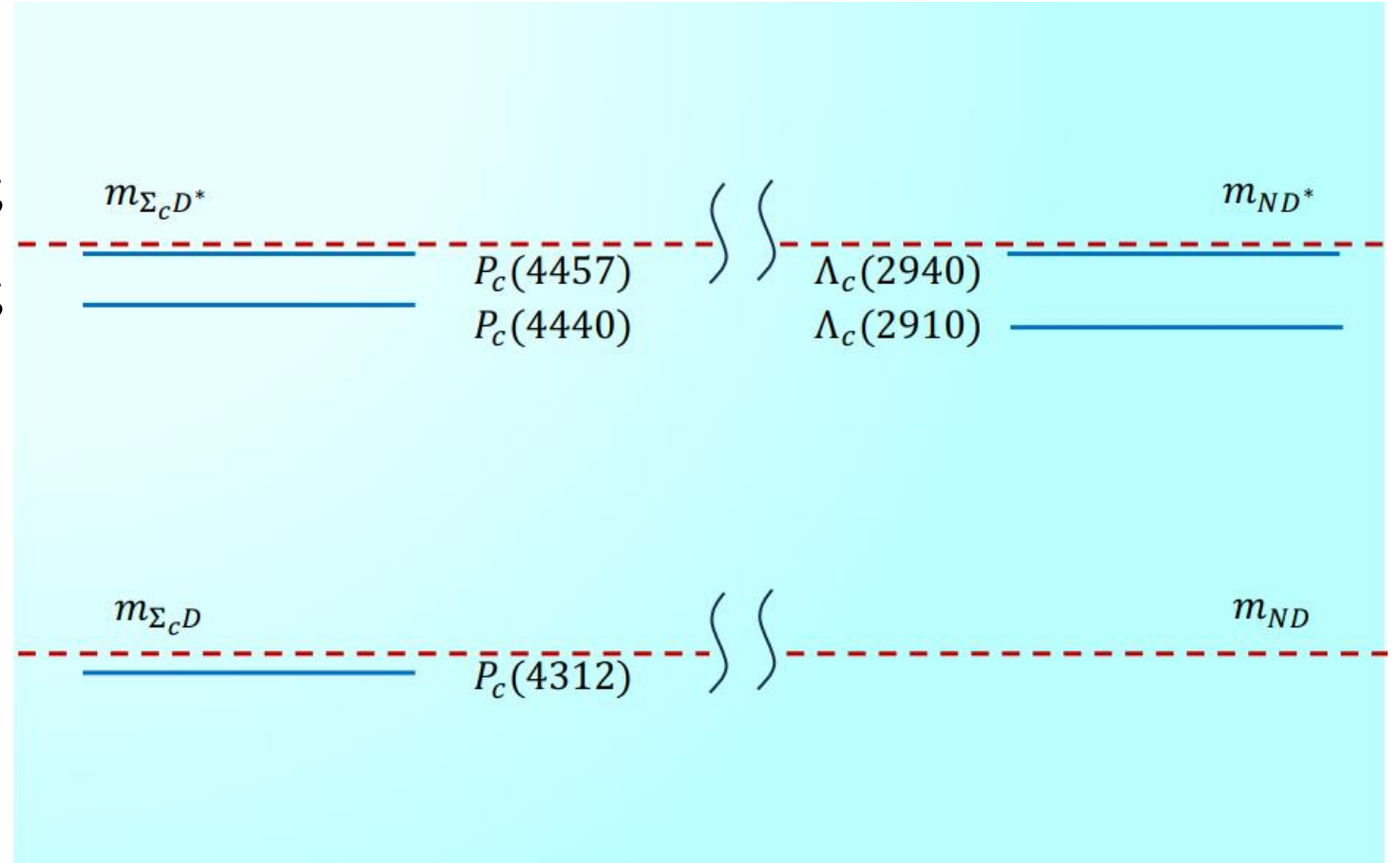
Phys.Rev.D 100 (2019) 1, 014022,

Phys.Rev.D 100 (2019) 11, 114002,

...

SetA: $\Lambda_c(2910) J^P = \frac{1}{2}^-$, $\Lambda_c(2940) J^P = \frac{3}{2}^-$.

SetB: $\Lambda_c(2940) J^P = \frac{1}{2}^-$, $\Lambda_c(2910) J^P = \frac{3}{2}^-$.



The mass spectrum of P_c states and Λ_c states.

● Effective Lagrangian:

$$\mathcal{L}_{\Lambda_c(1/2^-)}$$

$$= g_{\Lambda_c D^* N}^{1/2^-} \bar{\Lambda}_c(x) \gamma^\mu \gamma^5 \int dy \phi(y^2) [D_\mu^{*0}(x + \omega_{pD^*0} y) p(x - \omega_{D^*0} p y) + D_\mu^{*+}(x + \omega_{nD^*+} y) n(x - \omega_{D^*+} n y)]$$

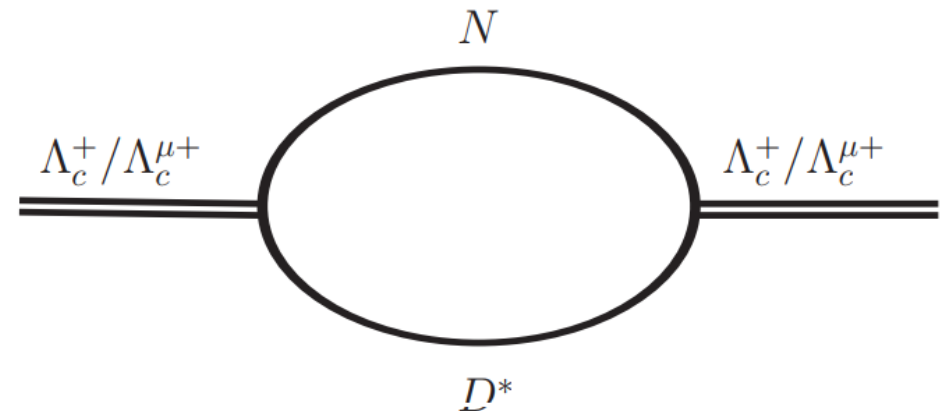
+ H. c.

$$\mathcal{L}_{\Lambda_c(3/2^-)}$$

$$= g_{\Lambda_c D^* N}^{3/2^-} \bar{\Lambda}_c^\mu(x) \int dy \phi(y^2) [D_\mu^{*0}(x + \omega_{pD^*0} y) p(x - \omega_{D^*0} p y) + D_\mu^{*+}(x + \omega_{nD^*+} y) n(x - \omega_{D^*+} n y)] + \text{H. c.}$$

Correlation function:

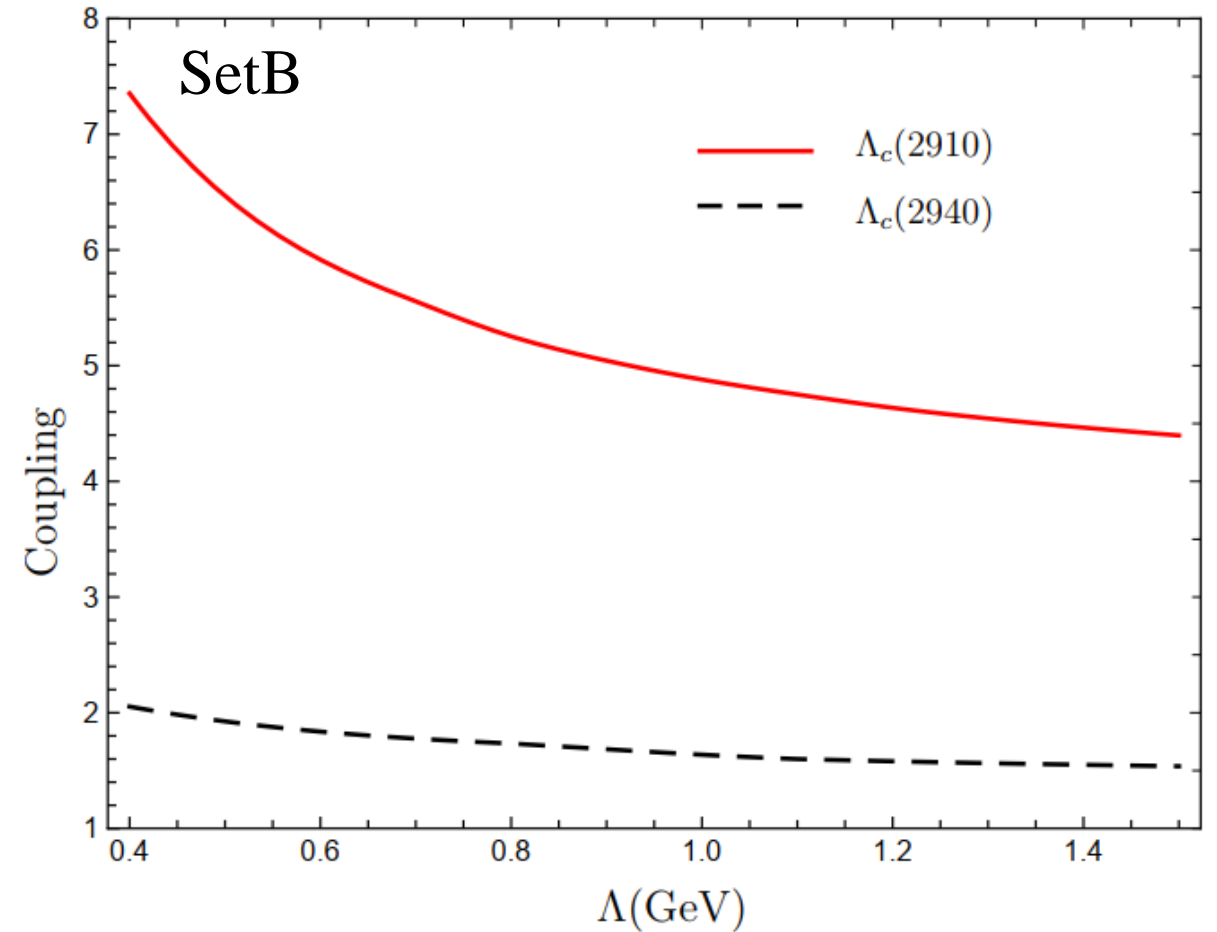
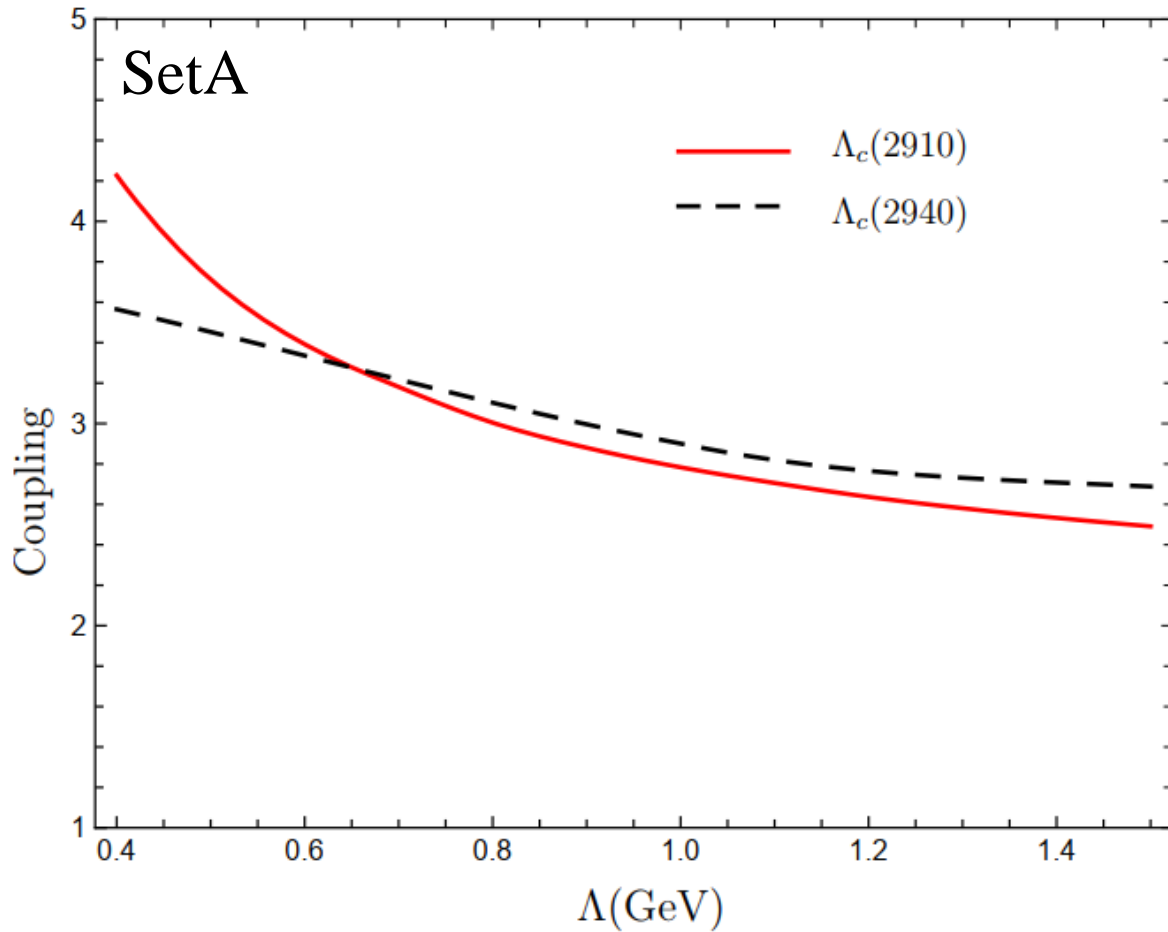
- ✓ Fourier transformation: $\phi(y^2) = \int \frac{d^4 p}{(2\pi)^4} e^{-ipy} \tilde{\phi}(-p^2)$,
- ✓ describe the interior structure of the molecular state,
- ✓ fall fast enough in the ultraviolet region,
- ✓ Gaussian form: $\tilde{\phi}(p_E^2) = \exp(-p_E^2/\Lambda^2)$.



Molecular scenario



- Compositeness condition: $Z = \langle Bare | Phys \rangle = 0$, $Z = 1 - \Pi' \left(m_{T_{C\bar{S}0}}^2 \right) = 0 \Rightarrow g_{\bar{\Lambda}_c}$.



The strong decay of the Λ_c

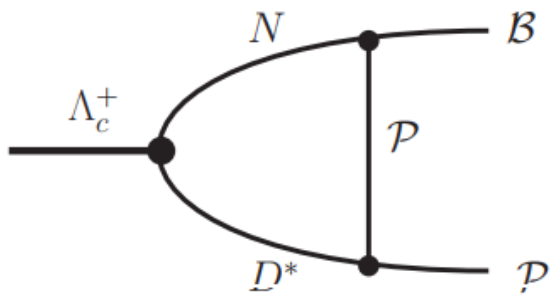


The possible decay mode:

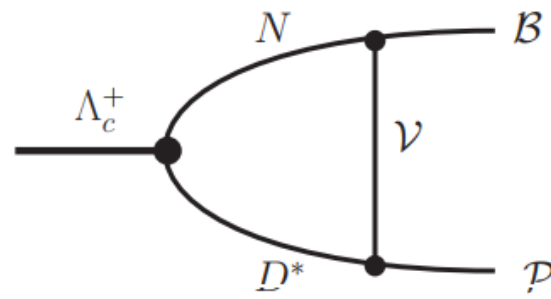
● $\Lambda_c \rightarrow (c\bar{u})(uud), (u\bar{u})(cud), (\bar{u}d)(cuu), (u\bar{d})(cdd) \dots$

	Final states	process	
$\frac{1^-}{2}$	ND	$1/2^- \rightarrow 0^- + 1/2^+$	$S - wave$
	$\Sigma_c \pi$	$1/2^- \rightarrow 0^- + 1/2^+$	$S - wave$
	$\Sigma_c^* \pi$	$1/2^- \rightarrow 0^- + 3/2^+$	$D - wave$
$\frac{3^-}{2}$	ND	$3/2^- \rightarrow 0^- + 1/2^+$	$D - wave$
	$\Sigma_c \pi$	$3/2^- \rightarrow 0^- + 1/2^+$	$D - wave$
	$\Sigma_c^* \pi$	$3/2^- \rightarrow 0^- + 3/2^+$	$S - wave$

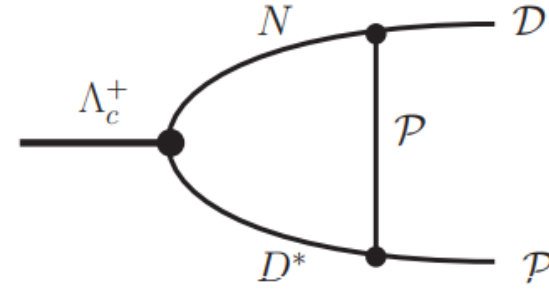
The strong decay of the Λ_c



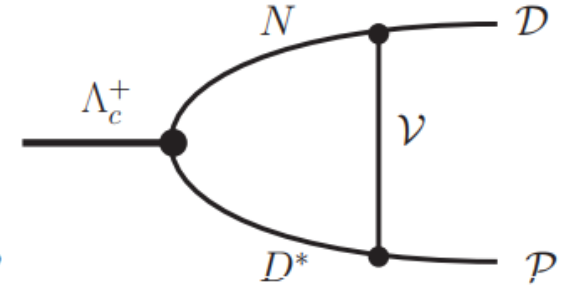
(a)



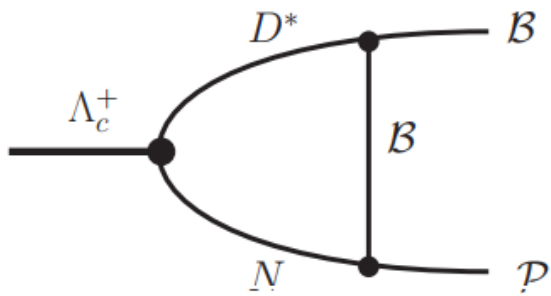
(b)



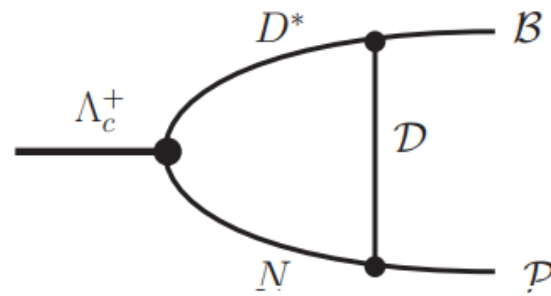
(e)



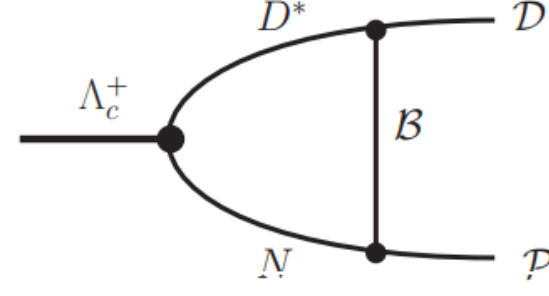
(f)



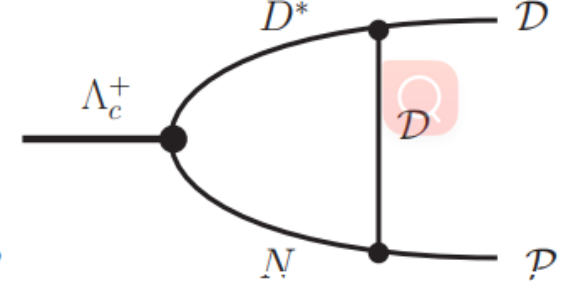
(c)



(d)



(g)



(h)

All the decay mode of the Λ_c at the hadron level.

The strong decay of the Λ_c



All the possible specific decay channels of the Λ_c

J^P	Final states	loops
$\frac{1}{2}^-$	ND	$[ND^*\pi], [ND^*\eta], [ND^*\rho], [ND^*\omega]$
	$\Sigma_c^{(*)+}\pi^0$	$[pD^{*0}\bar{D}^0], [pD^{*0}\bar{D}^{*0}], [nD^{*+}D^-], [nD^{*+}D^{*-}], [D^{*0}p\bar{\Delta}^+], [D^{*0}p\bar{p}], [D^{*+}n\bar{\Delta}^0], [D^{*+}n\bar{n}]$
	$\Sigma_c^{(*)++}\pi^-$	$[pD^{*0}D^-], [pD^{*0}D^{*-}], [D^{*+}n\bar{\Delta}^+], [D^{*+}n\bar{p}], [D^{*0}p\bar{\Delta}^{++}]$
	$\Sigma_c^{(*)0}\pi^+$	$[nD^{*+}\bar{D}^0], [nD^{*+}\bar{D}^{*0}], [D^{*0}p\bar{\Delta}^0], [D^{*0}p\bar{n}], [D^{*+}n\bar{\Delta}^-]$
$\frac{3}{2}^-$	ND	$[ND^*\pi], [ND^*\eta], [ND^*\rho], [ND^*\omega]$
	$\Sigma_c^{(*)+}\pi^0$	$[pD^{*0}\bar{D}^0], [pD^{*0}\bar{D}^{*0}], [nD^{*+}D^-], [nD^{*+}D^{*-}], [D^{*0}p\bar{\Delta}^+], [D^{*0}p\bar{p}], [D^{*+}n\bar{\Delta}^0], [D^{*+}n\bar{n}]$
	$\Sigma_c^{(*)++}\pi^-$	$[pD^{*0}D^-], [pD^{*0}D^{*-}], [D^{*+}n\bar{\Delta}^+], [D^{*+}n\bar{p}], [D^{*0}p\bar{\Delta}^{++}]$
	$\Sigma_c^{(*)0}\pi^+$	$[nD^{*+}\bar{D}^0], [nD^{*+}\bar{D}^{*0}], [D^{*0}p\bar{\Delta}^0], [D^{*0}p\bar{n}], [D^{*+}n\bar{\Delta}^-]$

The strong decay of the Λ_c



- Effective Lagrangian:

$SU(4)$ symmetry:

$$\mathcal{L}_{BBP} = -\frac{g_{BBP}}{m_P} \bar{B} \gamma^\mu \gamma_5 \partial_\mu \mathcal{P} B$$

$$\mathcal{L}_{BBV} = -g_{BBV} \bar{B} \gamma_\mu \mathcal{V}^\mu B$$

$$\mathcal{L}_{BDP} = \frac{g_{BDP}}{m_P} (\bar{D}^\mu B - \bar{B} D^\mu) \partial_\mu \mathcal{P}$$

$$\mathcal{L}_{BDV} = -i \frac{g_{BDV}}{m_V} (\bar{D}^\mu \gamma^5 \gamma^\nu B - \bar{B} \gamma^5 \gamma^\nu D^\mu) (\partial_\mu \mathcal{V}_\nu - \partial_\nu \mathcal{V}_\mu)$$

$$\mathcal{L}_{BDV} = \bar{D}_\alpha \gamma^\nu \mathcal{V}_\nu D^\alpha$$

$$\mathcal{L}_{PPV} = -i g_{PPV} (\mathcal{P} \partial_\mu \mathcal{P} - \partial_\mu \mathcal{P} \mathcal{P}) \mathcal{V}^\mu$$

$$\mathcal{L}_{VVV} = g_{VVV} \epsilon_{\mu\nu\alpha\beta} \partial^\mu \mathcal{V}^\nu \partial^\alpha \mathcal{V}^\beta \mathcal{P}$$

The strong decay of the Λ_c



- Form factor:

$$\mathcal{F}(m_q, \Lambda_1) = \frac{\Lambda_1^4}{(m^2 - q^2)^2 + \Lambda_1^4},$$

- Loop Integral:

$$C_0 = \int \frac{d^4 q}{(2\pi)^4} \tilde{\Phi} \frac{1}{[p_1^2 - m_1^2][p_2^2 - m_2^2][q^2 - m_q^2]}$$

Schwinger parametrization:

$$\frac{1}{m^2 - k^2} = \int_0^\infty d\alpha e^{-\alpha(m^2 - k^2)}$$

- Decay width:

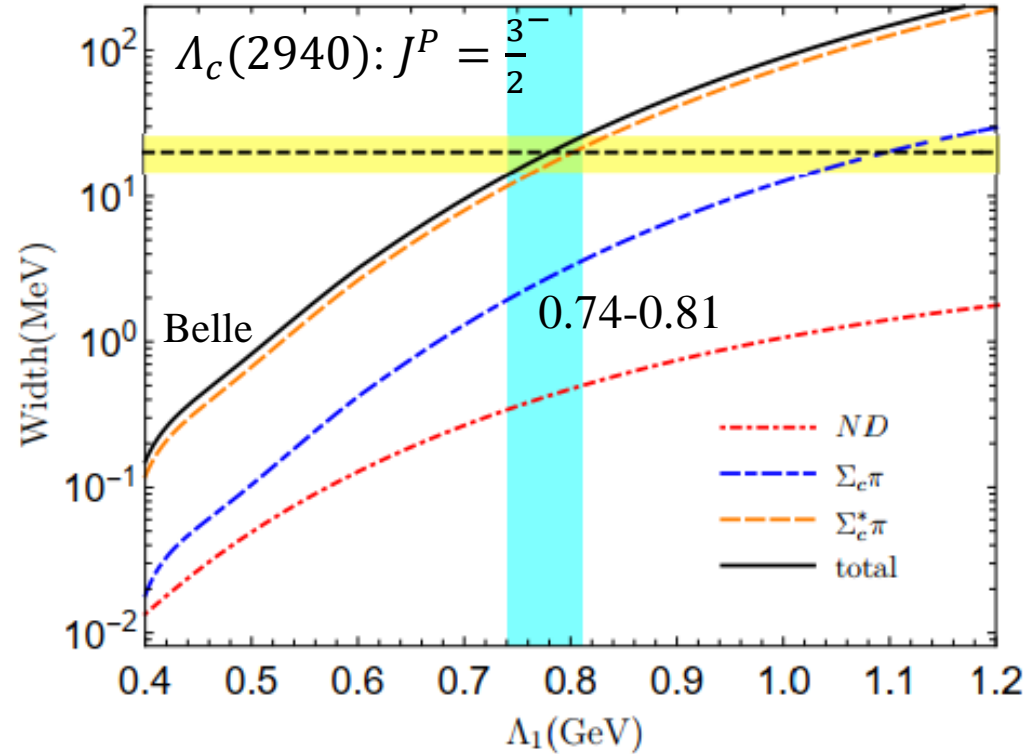
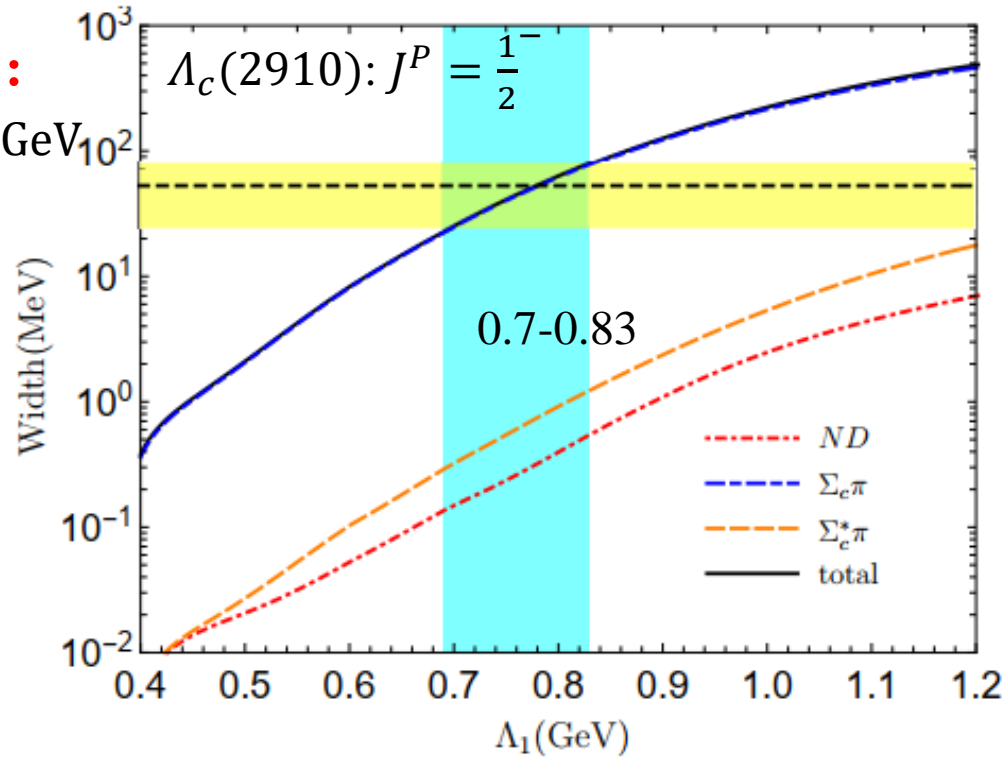
$$\Gamma_{\Lambda_c \rightarrow \dots} = \frac{1}{(2J + 1)8\pi m_{\Lambda_c}^2} |\overline{\mathcal{M}_{\Lambda_c \rightarrow \dots}}|^2$$

Numerical Results and discussions



Set A :

$\Lambda = 0.7 \text{ GeV}$



Conclusions:

- The estimations is consistent with the experiment measurement in parameter Λ_1 from 0.74 – 0.81;
- In this assignment, the dominant decay channel for $\Lambda_c(2910)$ is $\Sigma_c \pi$ and for $\Lambda_c(2940)$ is $\Sigma_c^* \pi$;
- Branching ratio: $Br(AB) = \Gamma_{AB} / \Gamma_{tot}$;

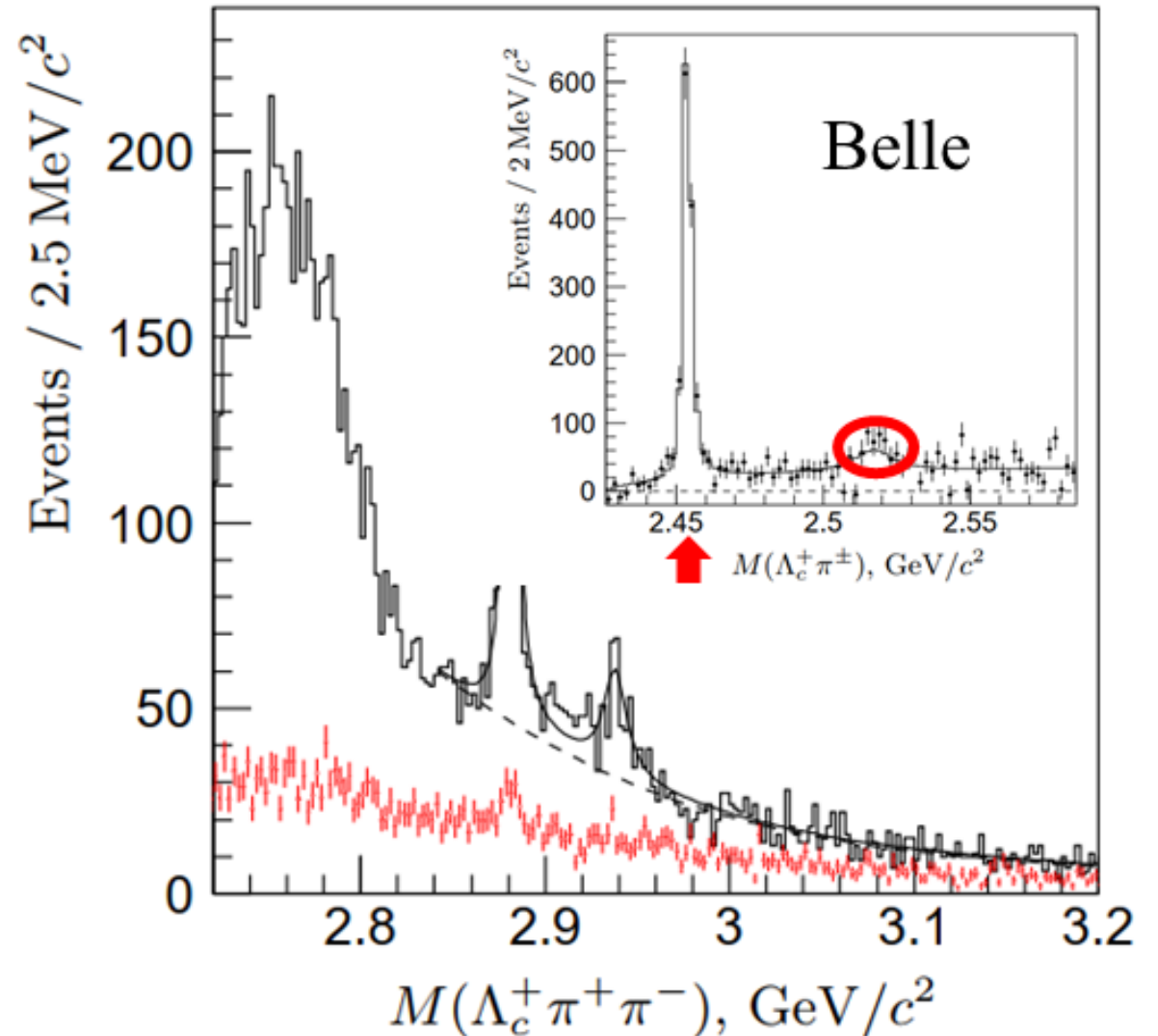
For $\Lambda_c(2910)$: $Br(ND) = 0.39 - 0.56\%$, $Br(\Sigma_c \pi) = 97.9 - 98.2\%$, $Br(\Sigma_c^* \pi) = 1.35 - 1.53\%$;

For $\Lambda_c(2940)$: $Br(ND) = 1.9 - 2.4\%$, $Br(\Sigma_c \pi) = 13.8 - 14\%$, $Br(\Sigma_c^* \pi) = 83.7 - 84\%$.

Numerical Results and discussions



- $\text{Br}(\Sigma_c^* \pi) = 83.7 - 84\%$
Large branching ratio for $\Lambda_c(2940)$!
- Experiment could look for $\Lambda_c(2940)$ in the $\Sigma_c^* \pi$ invariant mass spectrum.

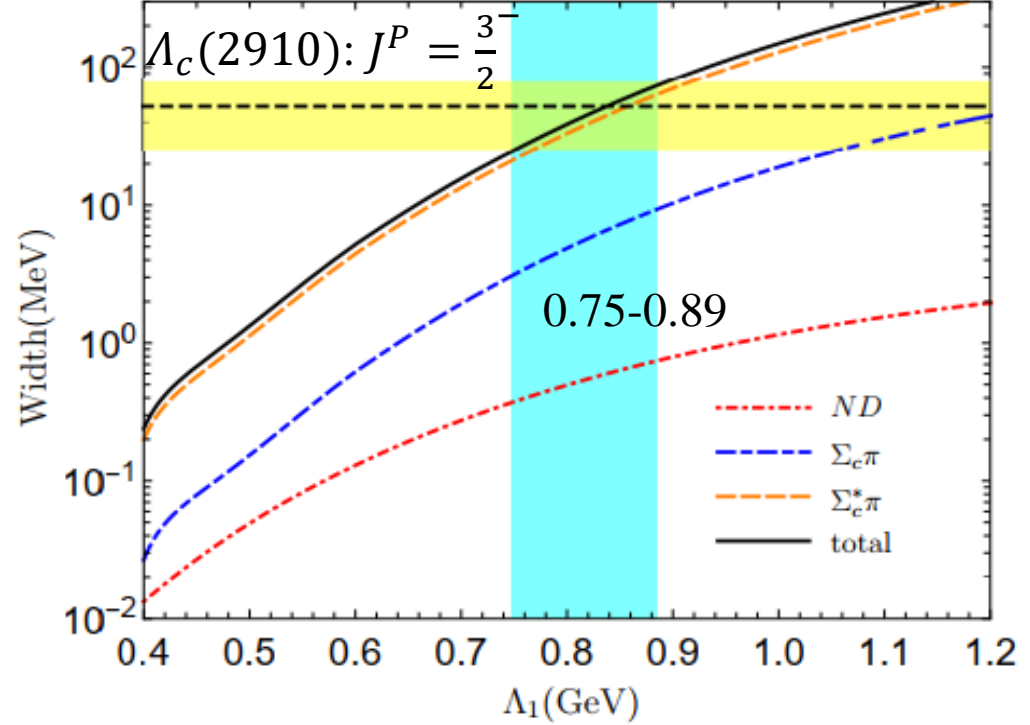
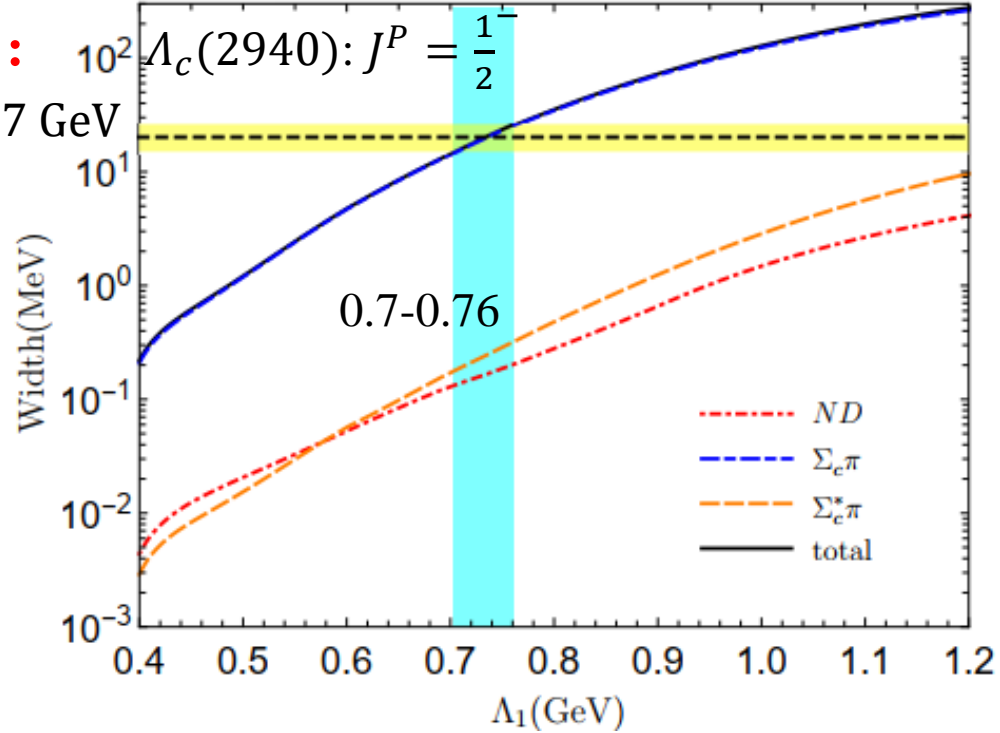


Numerical Results and discussions



Set B :

$\Lambda = 0.7 \text{ GeV}$



Conclusions:

- The estimations is consistent with the experiment measurement in parameter Λ_1 from 0.75 – 0.76;
- In this assignment, the dominant decay channel for $\Lambda_c(2910)$ is $\Sigma_c^*\pi$ and for $\Lambda_c(2940)$ is $\Sigma_c\pi$;
- Branching ratio: $Br(AB) = \Gamma_{AB} / \Gamma_{tot}$;

For $\Lambda_c(2910)$: $Br(ND) = 1.44 - 1.49\%$, $Br(\Sigma_c\pi) = 12.4 - 12.5\%$, $Br(\Sigma_c^*\pi) = 0.86\%$;

For $\Lambda_c(2940)$: $Br(ND) = 0.34 - 0.38\%$, $Br(\Sigma_c\pi) = 98.3\%$, $Br(\Sigma_c^*\pi) = 1.26 - 1.27\%$.

1. The present work investigate the strong decay behavior the $\Lambda_c(2940)$ and $\Lambda_c(2910)$ in D^*N molecular scenario by using the effective Lagrangian approach;
2. We estimate the partial decay width of $\Lambda_c \rightarrow ND$ and $\Lambda_c \rightarrow \Sigma_c^{(*)}\pi$;
3. By comparing our calculations with experiment measurements, the J^P for $\Lambda_c(2910)$ prefers $\frac{1}{2}^-$ and for $\Lambda_c(2940)$, the $J^P = \frac{3}{2}^-$ is weakly favored.

Thank you for your attention!