



Measurement of the electric dipole moment (EDM) of ^{171}Yb atoms in an optical dipole trap (ODT)

Tian Xia

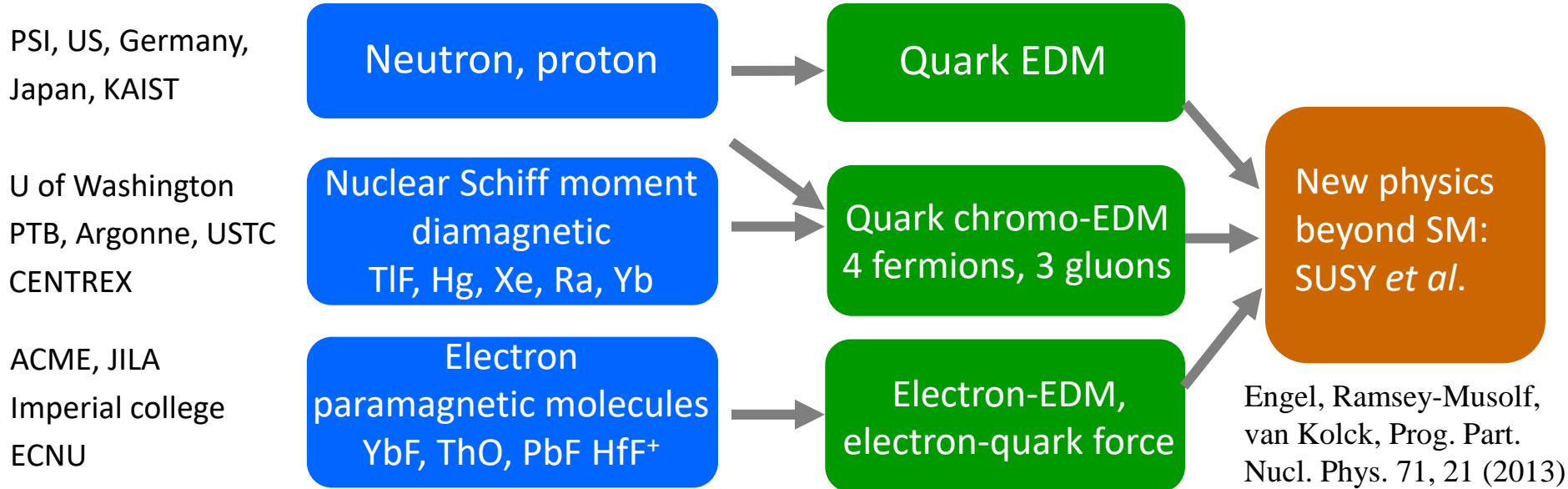
University of Science and Technology of China (USTC)

2023 CSNS NOPTREX workshop
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育天下英才
嚴濟慈
一九八八年五月



Searching for EDM in three different categories

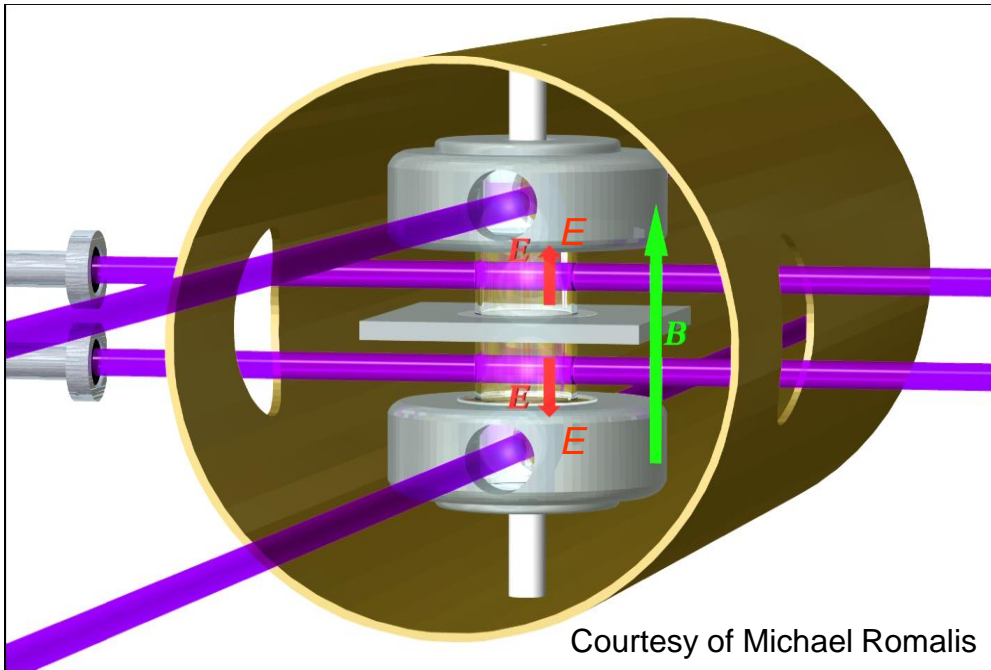


system	Upper limit (e-cm)	method	Value in Standard model (e-cm)
Electron	1×10^{-29}	Molecules – beam	10^{-38}
Neutron	2×10^{-26}	Neutrons – bottle	10^{-31}
¹⁹⁹ Hg	7×10^{-30}	Atoms – vapor cell	10^{-34}
¹⁷¹ Yb	This work	Atoms – trap	10^{-34}

The Seattle EDM Measurement

¹⁹⁹Hg

stable, high Z, groundstate 1S_0 , $I = \frac{1}{2}$, high vapor pressure



$$f_+ = \frac{2\mu B + 2dE}{h} \approx 15 \text{ Hz}$$

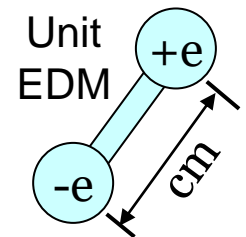
$$f_- = \frac{2\mu B - 2dE}{h} \approx 15 \text{ Hz}$$

$$|f_+ - f_-| < 25 \text{ pHz}$$

The best limit on atomic EDM

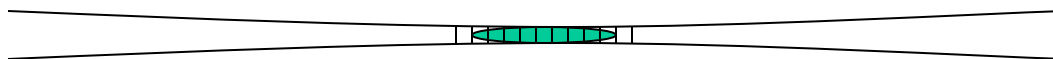
$$\text{EDM } (^{199}\text{Hg}) < 7 \times 10^{-30} \text{ e-cm}$$

Graner *et al.*, Phys Rev Lett (2016)



Measure EDM in an Optical Dipole Trap

M.V. Romalis and E.N. Fortson, Phys. Rev. A 59, 4547 (1999)



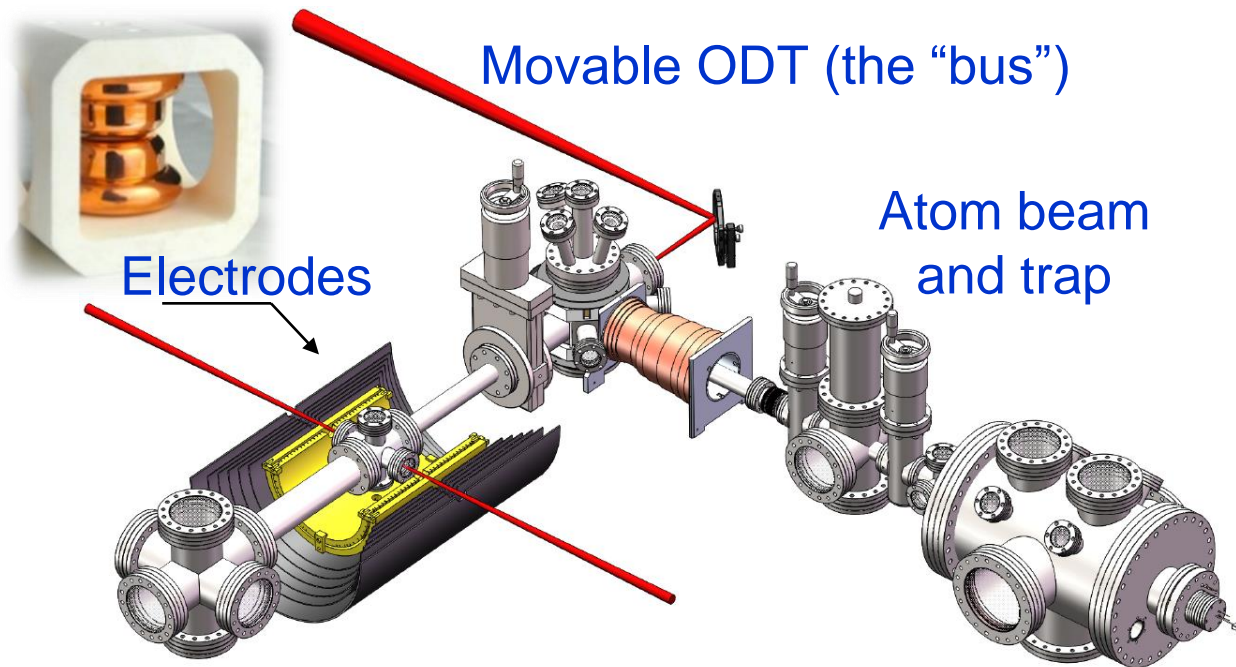
$$H = -\tilde{d}E = -\frac{1}{4}\alpha E_0^2$$

- Fiber laser: $\lambda = 1036 \text{ nm}$, Power = 10 Watts
- Focused to $50 \mu\text{m}$ \rightarrow trap depth $60 \mu\text{K}$

EDM in an optical dipole trap (ODT)

- $\mathbf{v} \times \mathbf{E}$, Berry's phase effects suppressed
- Cold scattering suppressed between cold Fermionic atoms
- Rayleigh scat. rate $\sim 10^{-1} \text{ s}^{-1}$; Raman scat. rate $\sim 10^{-12} \text{ s}^{-1}$
- Vector light shift $\sim \mu\text{Hz}$
- Parity mixing induced shift under control
- Conclusion: possible to reach 10^{-30} e cm for ^{199}Hg

^{171}Yb EDM Apparatus: Trapping + Science



$$\omega_+ = 2\mu B + 2dE$$

$$\omega_- = 2\mu B - 2dE$$

$$\delta d = \frac{\hbar}{2E\sqrt{\tau N \epsilon T}}$$

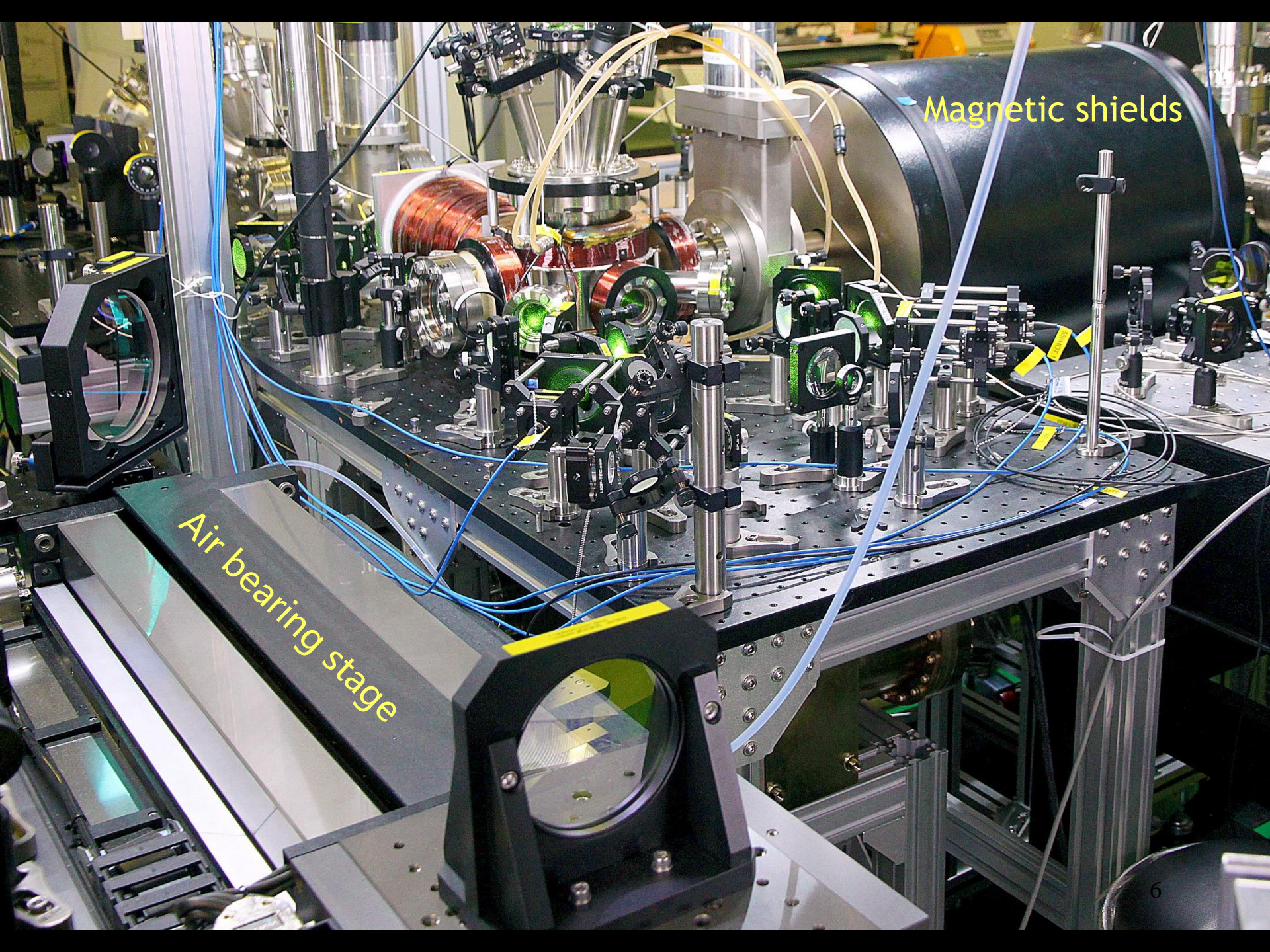
E : electric field

τ : precession time

N : atom number

ϵ : spin – state detection efficiency

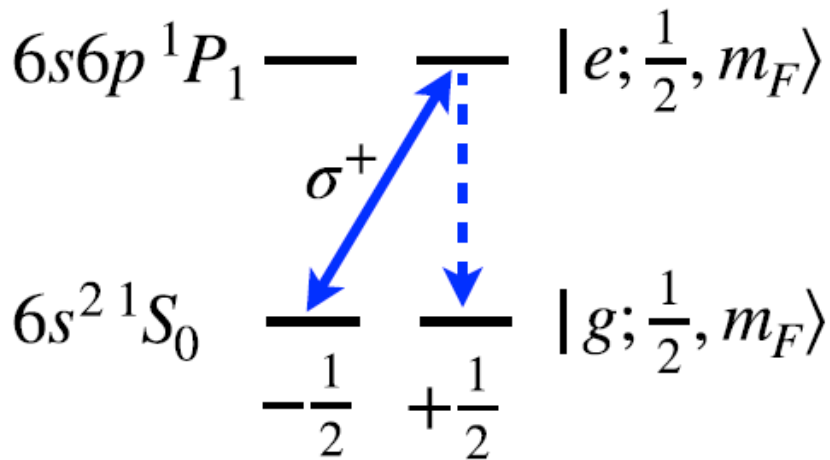
T : time of average



Magnetic shields

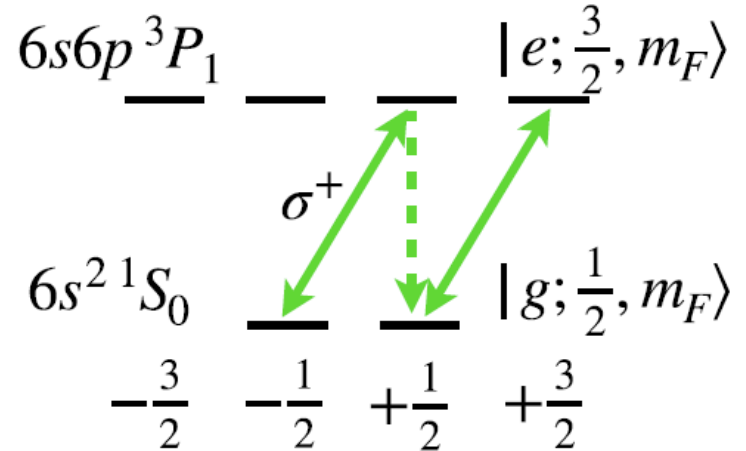
Air bearing stage

Spin-state detection – conventional method



Bright state:
~ 3 photons
before state
“demolished”

Dark state:
no photons,
state preserved

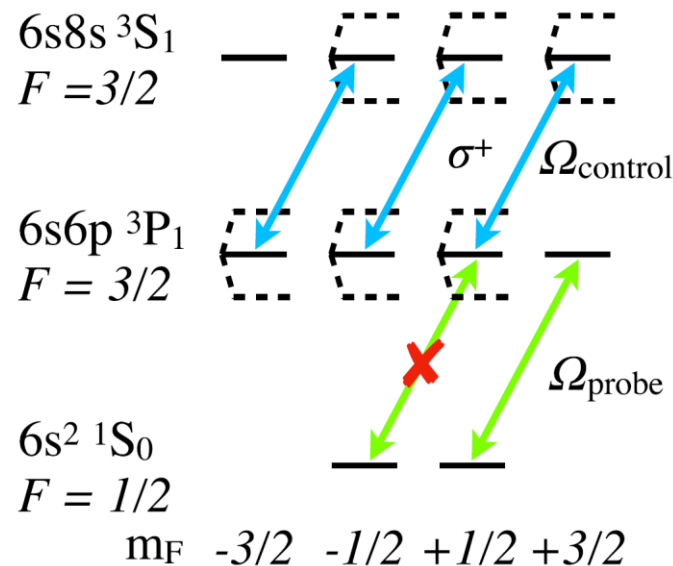
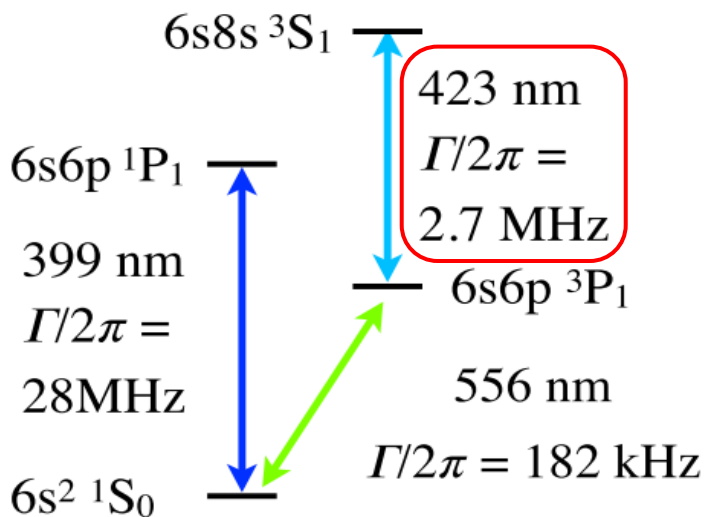


Bright state:
state
“demolished”

Brighter state:
cycling,
state preserved

Spin-state detection by a QND method

Quantum Non-Demolition



- $\Omega_{control} \sim 2\pi \times 40\text{ MHz}$
- $\Omega_{probe} \sim 2\pi \times 70\text{ kHz}$
- Spin flip suppressed by: $\Omega_c^2 / (\Gamma_e \Gamma_c) \sim 10^3$
- Need ODT to be at magic wavelength

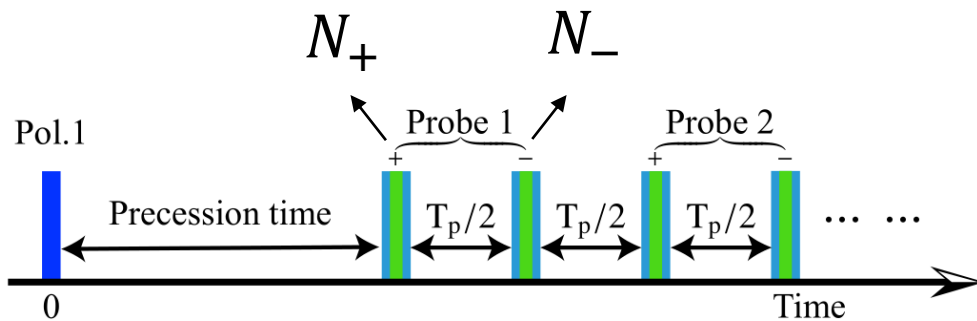
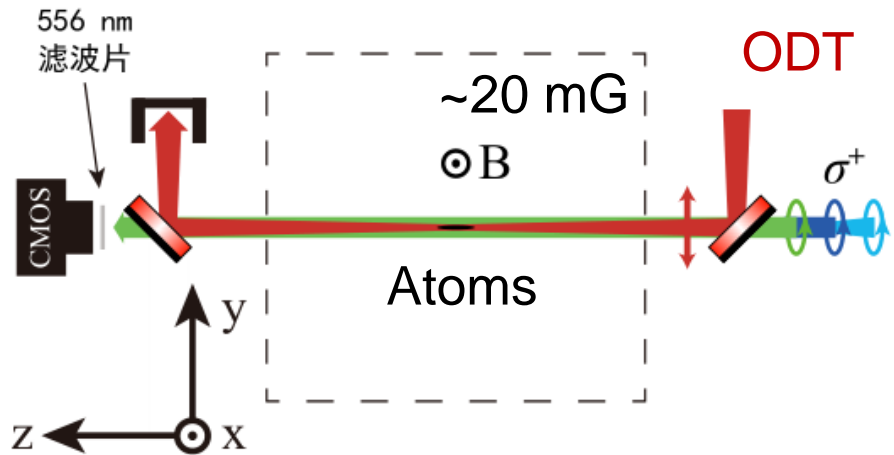
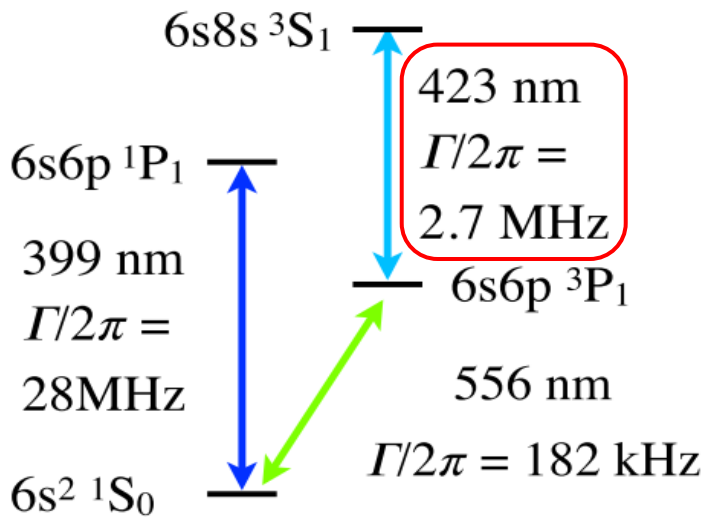
-- T. Zheng, M.S. Safronova *et al.*, PRA (2020)

Dark state:
no photons,
state preserved

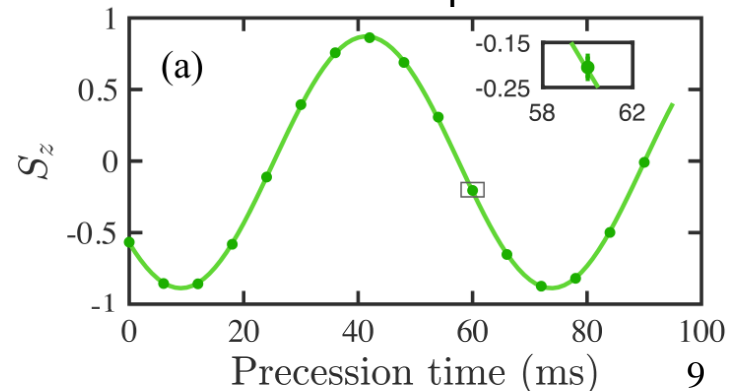
Bright state:
cycling,
state preserved

Spin-state detection by a QND method

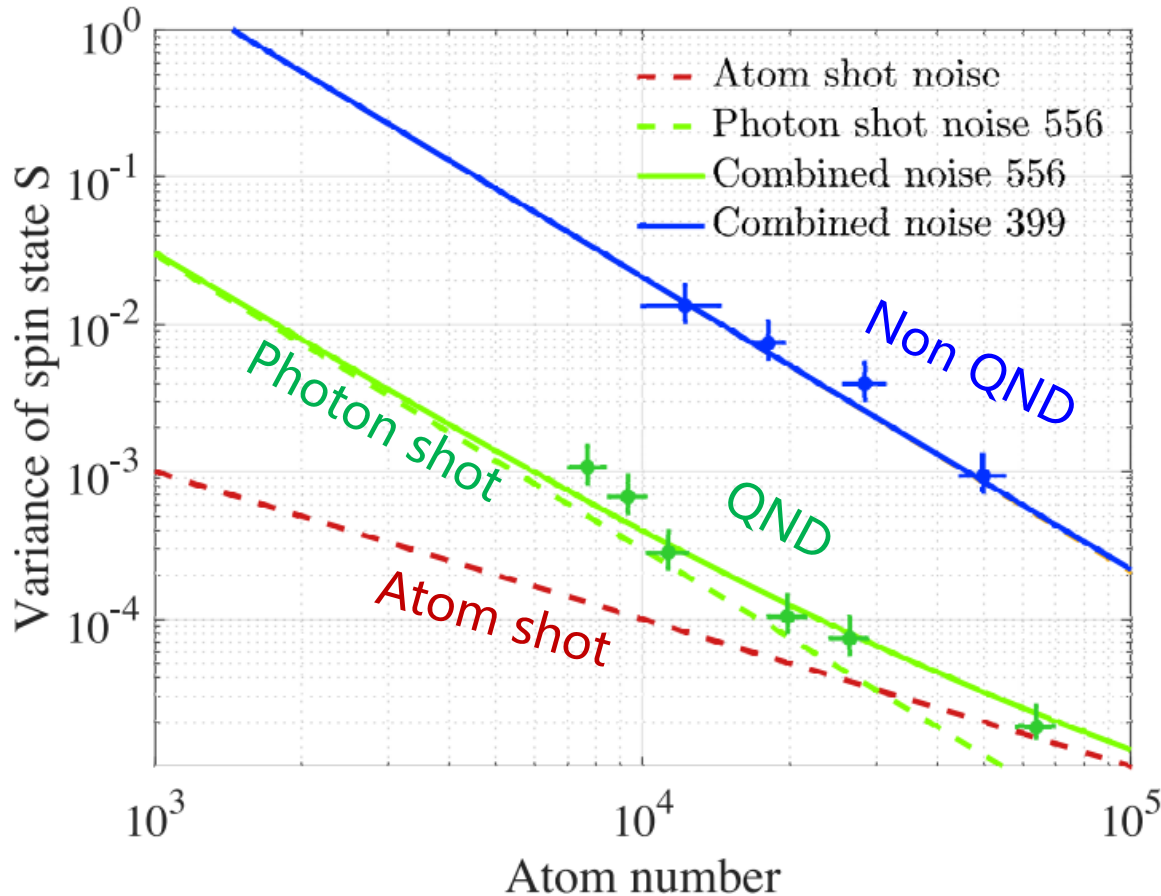
Quantum Non-Demolition



$$S_z = \frac{N_+ - N_-}{N_+ + N_-}$$



QND reduces measurement noise

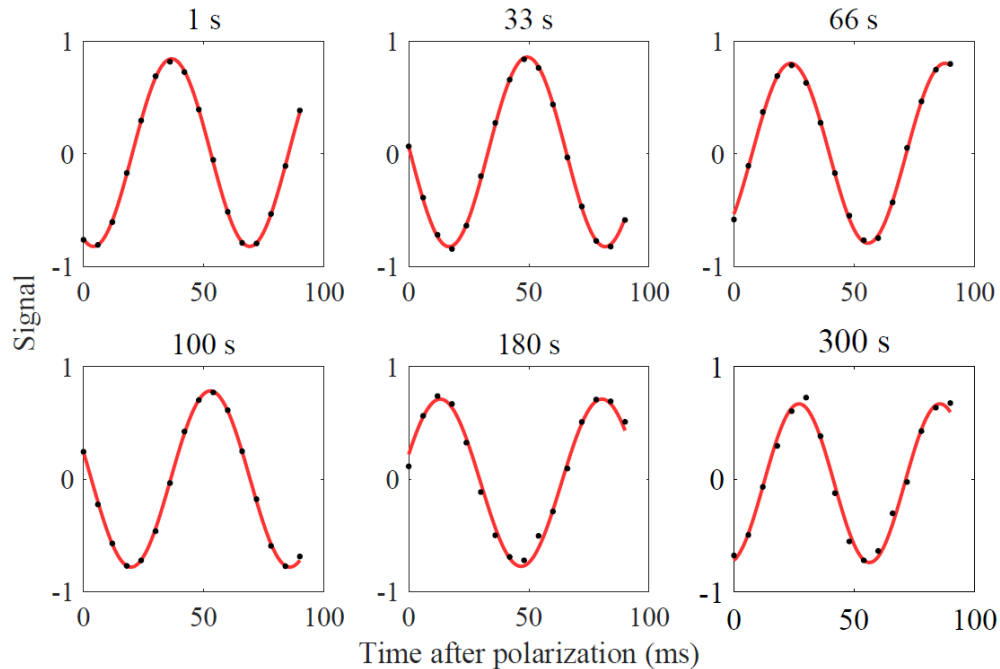


- Photons scattered non-QND: $\bar{n} \sim 2.5$
QND: $\bar{n} \sim 23$
- ~ 19 dB reduction of variance

QND measurement on the spin precession of laser-trapped ^{171}Yb atoms
Y.A. Yang *et al.*, Phys. Rev. Applied 19, 054015 (2023)

Long spin coherence time (T_2)

$$S_z = \frac{N_+ - N_-}{N_+ + N_-} = C \exp\left(-\frac{t}{T_2}\right) \cos(2\pi f t + \phi_0) + O$$



- Vacuum-limited trap lifetime: 75 s
- EDM measurement precession time: 96 s
- Observed precession to: 300 s
- T_2 : $(9 \pm 4) \times 10^3$ s (2.5 ± 1.1 hr)

Systematic errors

- effects correlated with E-field flipping

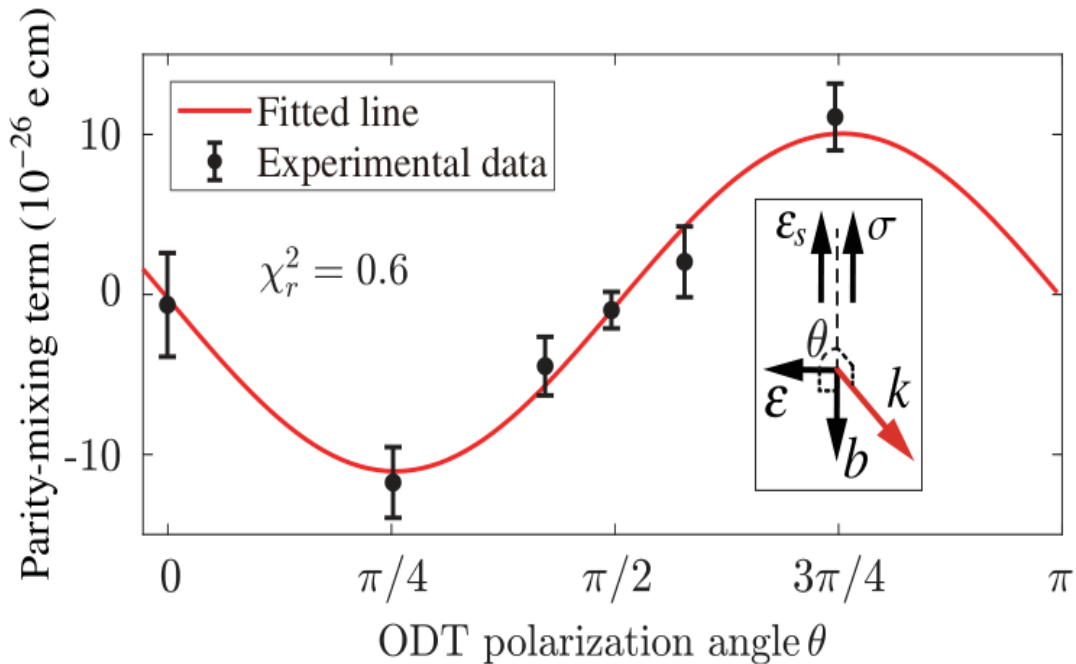
Source	Error
B-field correlations	1.00×10^{-27}
Parity mixing	0.59×10^{-27}
Leakage current	0.14×10^{-27}
ODT power effect	0.09×10^{-27}
E-squared effect	0.04×10^{-27}
Total	1.18×10^{-27}

Parity mixing

$$\Delta E = \frac{e^2 \mu_B E_0^2 E_S}{4 \hbar^2} \sum_{J', J'', m', m''} \left[\frac{\langle J', m' | \boldsymbol{\varepsilon} \cdot \mathbf{r} | J, m \rangle \langle J'', m'' | \mathbf{b}^* \cdot (\mathbf{L} + 2\mathbf{S}) | J', m' \rangle \langle J, m | \boldsymbol{\varepsilon}_S \cdot \mathbf{r} | J'', m'' \rangle}{(\omega_{J'} - \omega)(\omega_{J''} - \omega)} \right. \\ \left. + \frac{\langle J, m' | \mathbf{b} \cdot (\mathbf{L} + 2\mathbf{S}) | J, m \rangle \langle J', m' | \boldsymbol{\varepsilon}^* \cdot \mathbf{r} | J, m' \rangle \langle J, m | \boldsymbol{\varepsilon}_S \cdot \mathbf{r} | J', m' \rangle}{\omega_{J'}(\omega_{J'} - \omega)} + \text{perm.} + \text{c.r.} \right]$$

M.V. Romalis and E.N. Fortson, Phys. Rev. A 59, 4547 (1999)

T-even, P-even
yet EDM-like



- $\boldsymbol{\varepsilon}$: ac E-field of ODT
- \mathbf{b} : ac B-field of ODT
- $\boldsymbol{\varepsilon}_S$: Static E-field
- σ : Quantization axis
- Static B-field

measurement : $1.2(2) \times 10^{-25} \text{ e cm}$ ($\boldsymbol{\varepsilon} \cdot \boldsymbol{\varepsilon}_S = \boldsymbol{\varepsilon} \cdot \sigma = \sqrt{2}/2$)
 calculation : $2 \times 10^{-25} \text{ e cm}$ ($\boldsymbol{\varepsilon} \cdot \boldsymbol{\varepsilon}_S = \boldsymbol{\varepsilon} \cdot \sigma = \sqrt{2}/2$)

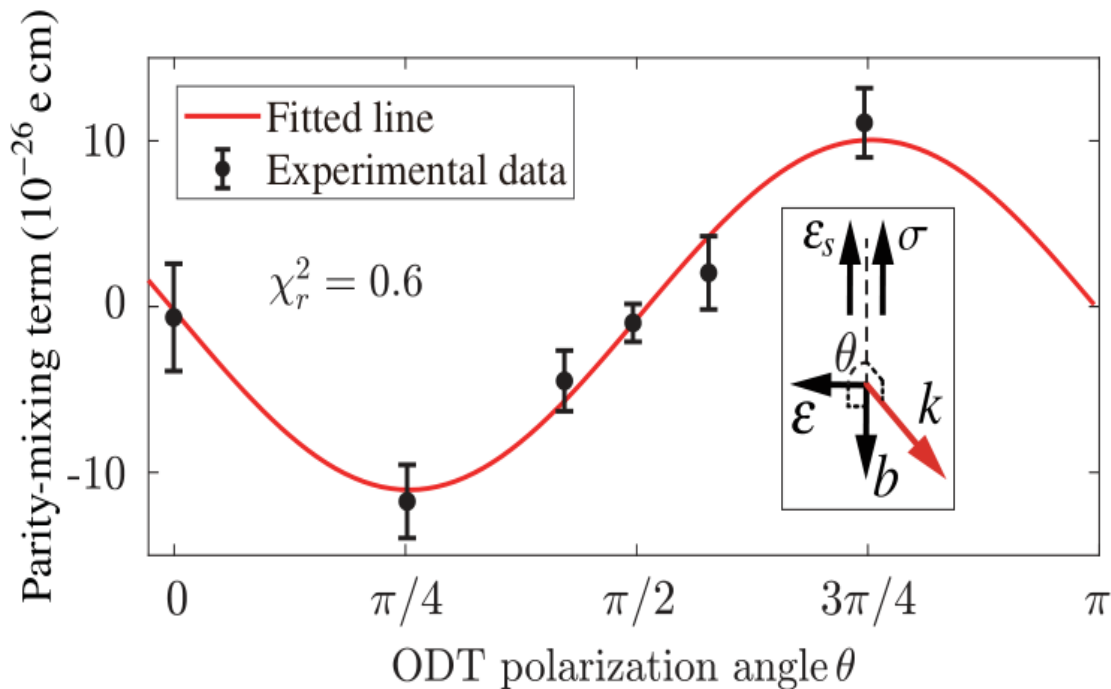
Parity mixing

$$\Delta \nu_{F=I \pm 1/2} = \mp [\nu_{MD}^1 (\hat{b} \cdot \hat{\sigma})(\hat{\epsilon} \cdot \hat{\epsilon}_s) + \nu_{MD}^2 (\hat{b} \cdot \hat{\epsilon}_s)(\hat{\epsilon} \cdot \hat{\sigma})] m$$

M.V. Romalis and E.N. Fortson, Phys. Rev. A 59, 4547 (1999)

$$\Delta \nu \approx -\frac{\nu_1 + \nu_2}{2} \sin(2\theta), \quad (0 \leq \theta < \pi)$$

T-even, P-even
yet EDM-like



ϵ : ac E-field of ODT

b : ac B-field of ODT

ϵ_s : Static E-field

σ : Quantization axis

Static B-field

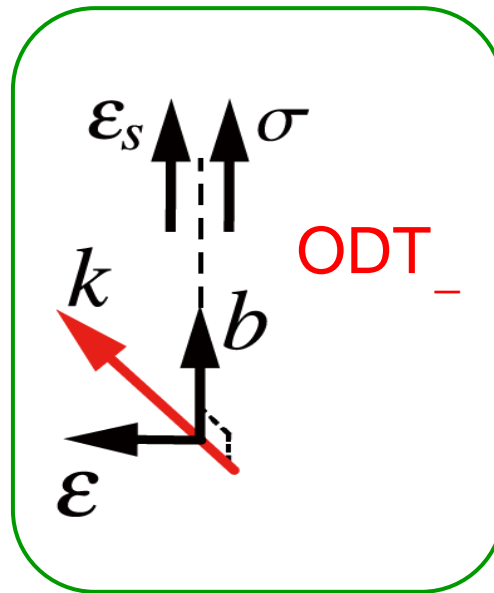
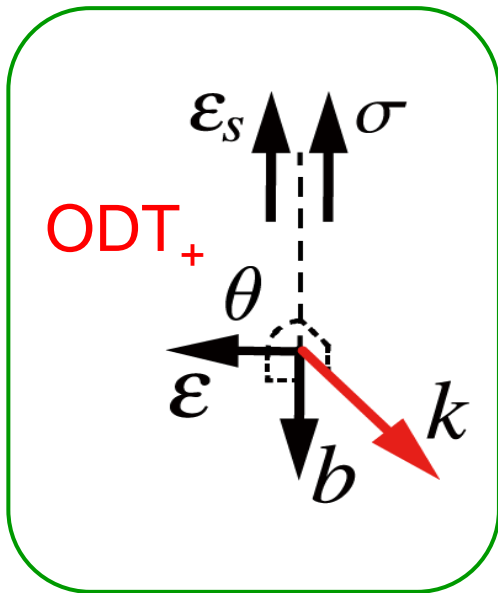
measurement : $1.2(2) \times 10^{-25}$ e cm ($\epsilon \cdot \epsilon_s = \epsilon \cdot \sigma = \sqrt{2}/2$)

calculation : 2×10^{-25} e cm ($\epsilon \cdot \epsilon_s = \epsilon \cdot \sigma = \sqrt{2}/2$)

Parity mixing

$$\Delta \nu_{F=I \pm 1/2} = \mp [\nu_{MD}^1 (\hat{\mathbf{b}} \cdot \hat{\boldsymbol{\sigma}}) (\hat{\boldsymbol{\epsilon}} \cdot \hat{\boldsymbol{\epsilon}}_s) + \nu_{MD}^2 (\hat{\mathbf{b}} \cdot \hat{\boldsymbol{\epsilon}}_s) (\hat{\boldsymbol{\epsilon}} \cdot \hat{\boldsymbol{\sigma}})] m$$

M.V. Romalis and E.N. Fortson, Phys. Rev. A 59, 4547 (1999)



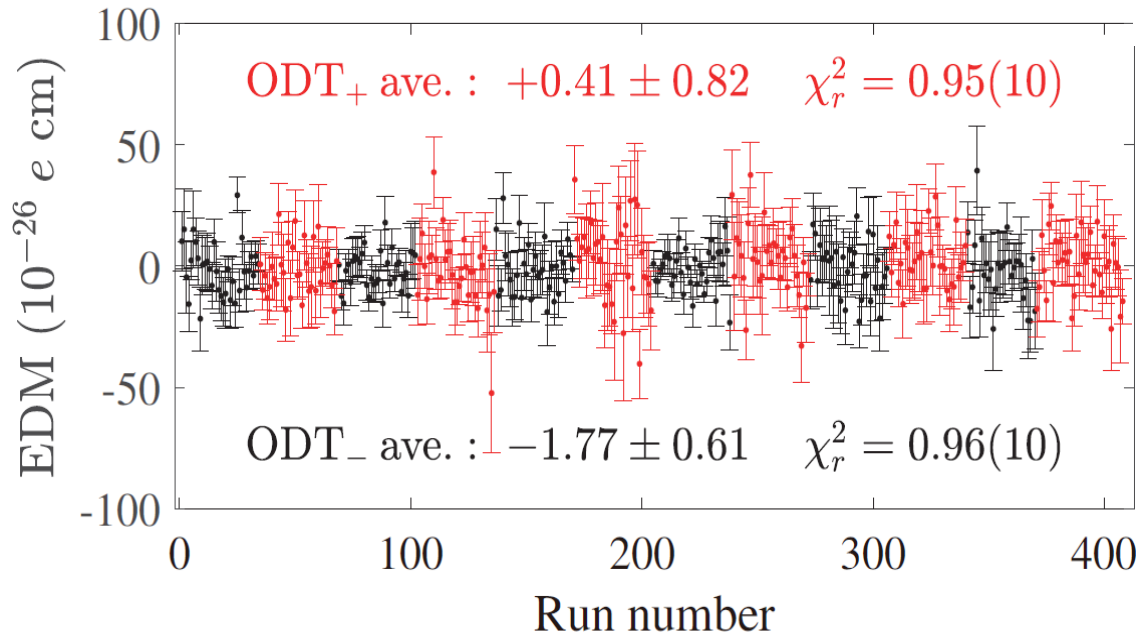
Control:

$$\theta = \pi/2 \pm 80 \text{ mrad}$$

$$(\theta_+ - \theta_-) = 0 \pm 7 \text{ mrad}$$

- Average between ODT+ and ODT-
- Residual parity mixing effect: 6×10^{-28} e cm
- Future: better k+ and k- balance in an optical cavity

^{171}Yb EDM measurement results



$$\delta d = \frac{\hbar}{2E\tau\sqrt{n}} \sqrt{\frac{1}{N_a\epsilon_d} + \sigma_{\phi(\delta B)}^2}$$

$$E: 73 \text{ kV/cm}$$

$$\tau: 96 \text{ s}$$

$$N: 6 \times 10^4$$

$$\epsilon_d: \sim 50\%$$

$$T: 510 \text{ h (22 d)}$$

$$\sigma_{\phi(\delta B)}^2 \approx 4 \times \frac{1}{N_a\epsilon_d}$$

$$d(^{171}\text{Yb}) = (-6.8 \pm 5.1_{\text{stat}} \pm 1.2_{\text{syst}}) \times 10^{-27} \text{ e cm}$$

$$\text{Upper limit (95\%)} : |d(^{171}\text{Yb})| < 1.5 \times 10^{-26} \text{ e cm}$$

Measurement of the EDM of ^{171}Yb atoms in an ODT
T.A. Zheng *et al.*, Phys. Rev. Lett. 129, 083001 (2022)

Upper limits on Schiff moments

Diamagnetic system	$^{205}\text{Tl}^{19}\text{F}$	^{199}Hg	^{129}Xe	^{225}Ra	^{171}Yb
Upper limit of EDM (E-26 e cm)	6500 [Yale-1991]	7.4×10^{-4} [U of Washington-2016]	0.14 [PTB-2019]	1400 [Argonne-2016]	1.5 [USTC-2022]
Upper limit of Schiff moment (e fm ³)* 10 ¹⁰	8.8	2.6×10^{-3}	5.2	1700	7.9
Calculation of Schiff moment (e fm ³)* 10 ⁸	$1.2 \eta_{pp} - 1.4 \eta_{pn}$ [Flambaum-86]	$-1.4 \eta_{np}$ [Flambaum-86]	$1.75 \eta_{np}$ [Flambaum-86]	$300 \eta_n,$ 1100η [Flambaum-03, Auerbach-96]	$\sim -1.4 \eta_{np}$ [Dzuba-07]

- The first ^{171}Yb EDM result contributes to constraining BSM physics: on the same order of as ^{129}Xe , ^{225}Ra , $^{205}\text{Tl}^{19}\text{F}$, although lagging ^{199}Hg .
- The global analysis of EDM results requires different systems with complementary sensitivities to BSM parameters, rather than ^{199}Hg alone.

Outlook #1

Upgrade: ^{171}Yb EDM precision improves into E-28 e-cm

- ❖ Larger E-field
- ❖ Less B-field noise
- ❖ Longer trap lifetime
- ❖ Optical cavity

E-field

- This work: 73 kV/cm, copper electrodes
- Test setup: 200 kV/cm, copper
- Other works: 500 kV/cm, niobium - Ready, et al. NIM A (2021)

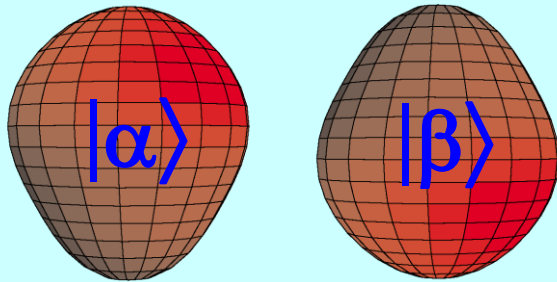
B-field

- This work: ~ 3 pT over 100 s, Seebeck effect due to ODT heating
- Environmental noise: ~ 1 pT, use magnetometers
- Johnson noise: ~ 0.6 pT, use co-magnetometer (^{173}Yb)

EDM of ^{225}Ra enhanced and more reliably calculated

- Closely spaced parity doublet – Haxton & Henley, PRL (1983)
- Large Schiff moment due to octupole deformation – Auerbach, Flambaum & Spevak, PRL (1996)
- Relativistic atomic structure ($^{225}\text{Ra} / ^{199}\text{Hg} \sim 3$) – Dzuba, Flambaum, Ginges, Kozlov, PRA (2002)

Parity doublet



$$\Psi^- = (|\alpha\rangle - |\beta\rangle)/\sqrt{2}$$

$$\Psi^+ = (|\alpha\rangle + |\beta\rangle)/\sqrt{2}$$

55 keV

$$\text{Schiff_moment} = \sum_{i \neq 0} \frac{\langle \psi_0 | \hat{S}_z | \psi_i \rangle \langle \psi_i | \hat{H}_{PT} | \psi_0 \rangle}{E_0 - E_i} + c.c.$$

Enhancement Factor: EDM (^{225}Ra) / EDM (^{199}Hg)

	Isoscalar	Isovector
Skyrme SIII	300	4000
Skyrme SkM*	300	2000
Skyrme SLy4	700	8000

Schiff moment of ^{225}Ra , Dobaczewski, Engel, PRL (2005)
Schiff moment of ^{199}Hg , Dobaczewski, Engel et al., PRC (2010)

“[Nuclear structure] calculations in Ra are almost certainly more reliable than those in Hg.”

– Engel, Ramsey-Musolf, van Kolck, Prog. Part. Nucl. Phys. (2013)

Constraining parameters in a global EDM analysis.

– Chupp, Ramsey-Musolf, arXiv1407.1064 (2014)

Outlook #2

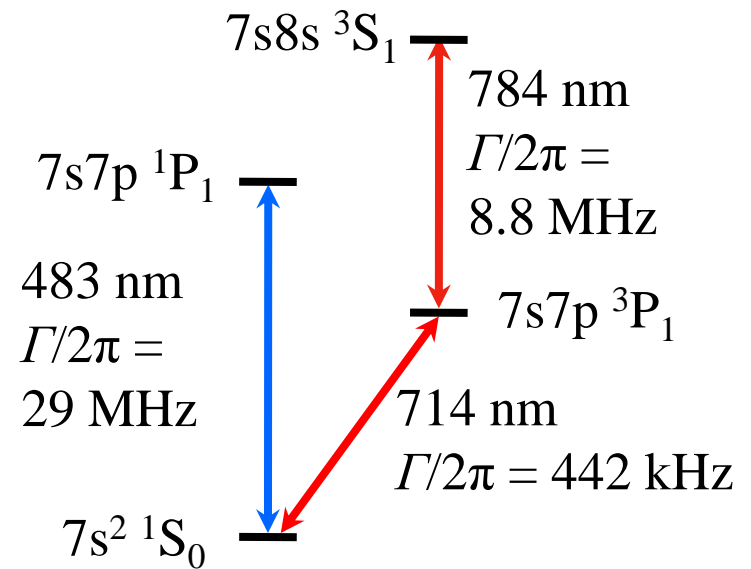
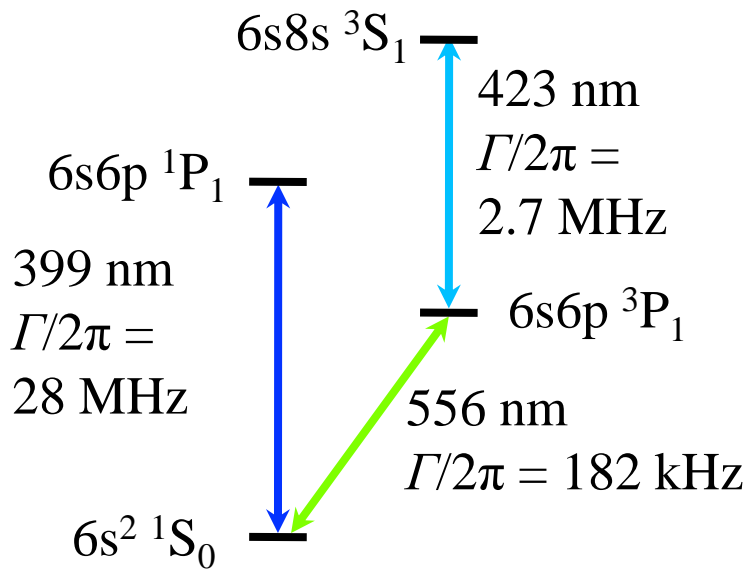
^{171}Yb

2-electron system
Nuclear spin $I = \frac{1}{2}$
Stable and abundant



^{225}Ra :

Group II
Nuclear spin $I = \frac{1}{2}$
Radioactive $t_{1/2} = 15 \text{ d}$

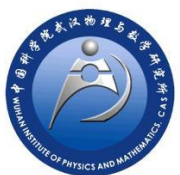
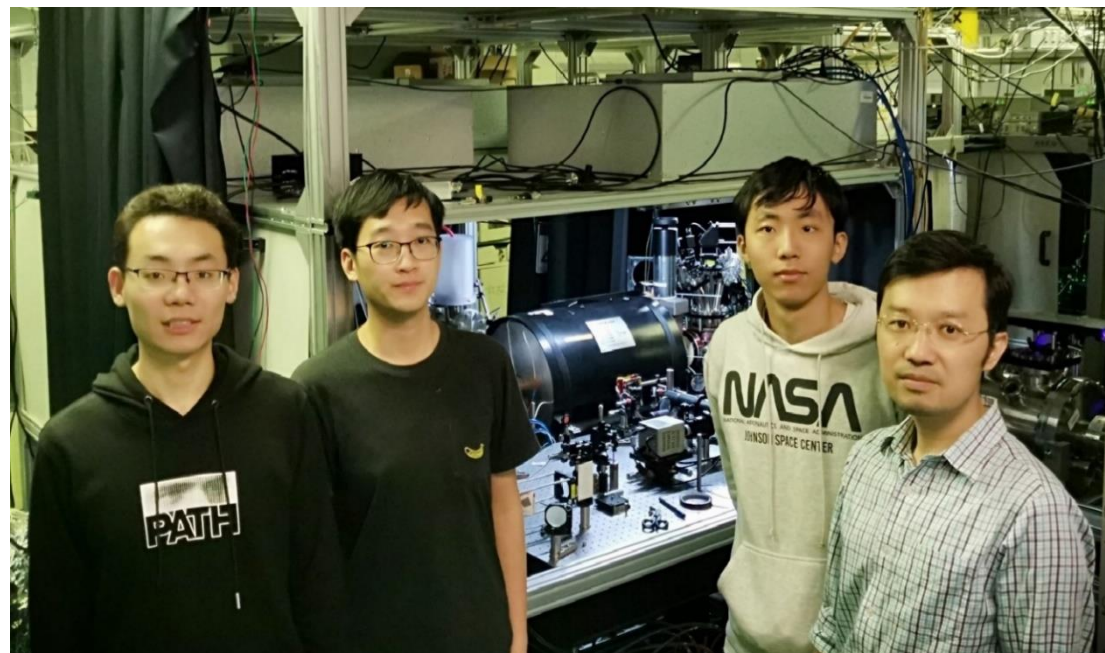


Collaboration and Support



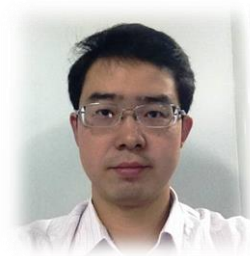
USTC

Z.-T. Lu, Tian Xia,
Dong Sheng, Tao Zheng,
Yang Yang, S.-Z. Wang



CAS Institute, Wuhan

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Jaideep Singh