# Muon collider signatures for a Z' with a maximal $\mu - \tau$ coupling in $U(1)_{L_{\mu}-L_{\tau}}$

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## Introduction



BNL g-2 FNAL g-2 + 4.2σ Standard Mode Experiment Average 18.0 18.5 21.0 17.5 19.0 19.5 20.0 20.5 21.5  $a_{\mu} \times 10^9 - 1165900$ 

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## Introduction

$$\Delta a_{\mu} = a_{\mu}^{exp} - a_{\mu}^{SM} = (251 \pm 59) \times 10^{-11} \qquad 4.2\sigma$$

Simplest  $U(1)_{L_{\mu}-L_{\tau}}$  model :

• First, second and third generations  $U(1)_{L_{\mu}-L_{\tau}}$  charge : 0,+1,-1

$$\begin{split} L_{\text{int}-Z'} &= -\tilde{g} \Big( \bar{\mu} \gamma^{\mu} \mu - \bar{\tau} \gamma^{\mu} \tau + \bar{\nu}_{\mu} \gamma^{\mu} L \nu_{\mu} - \bar{\nu}_{\tau} \gamma^{\mu} L \nu_{\tau} \Big) Z'_{\mu}, \\ \Delta a^{Z'}_{\mu} &= \frac{\tilde{g}^2}{8\pi^2} \frac{m^2_{\mu}}{m^2_{Z'}} \int_0^1 \frac{2x^2(1-x)dx}{1-x + (m^2_{\mu}/m^2_{Z'})x^2} \,. \end{split}$$

- In the limit  $m_{Z'} >> m_{\mu}, \Delta a_{\mu}^{Z'} = (\tilde{g}^2/12\pi^2)(m_{\mu}^2/m_{Z'}^2).$  $\tilde{g}^2/m_{Z'}^2 = (2.66 \pm 0.63) \times 10^{-5} \text{GeV}^{-2}$
- Will induce muon neutrino trident (MNT) process.

## **Introduction** :Neutrino trident process

$$v_{\mu}N \rightarrow v_{i}\mu\bar{\mu}N: | \text{arge } m_{Z'}: \text{ Contribution Proportional to } \tilde{g}^{2}/m_{Z'}^{2}$$

$$\frac{\sigma_{Z'}}{\sigma_{SM}}\Big|_{\text{trident}} = \frac{\left(1 + 4s_{W}^{2} + 8\tilde{g}^{2}m_{W}^{2}/g^{2}m_{Z'}^{2}\right)^{2} + 1}{1 + (1 + 4s_{W}^{2})^{2}}, \quad \text{Data:} \\1.58 \pm 0.57, \text{CHARM-II} \\0.82 \pm 0.28 \text{CFR} \\0.72^{\pm 1.73} \text{ NuTeV} \\ \bullet \quad \sigma_{Z'}/\sigma_{SM} = 5.86 \text{ for } \tilde{g}^{2}/m_{Z'}^{2} = (2.66 \pm 0.63) \times 10^{-5} \text{GeV}^{-2}$$

$$\frac{1}{\sigma_{\sigma}^{CCFR}} = 0.82 \pm 0.28 \frac{1}{\sigma_{\sigma}^{CCFR}} = 0.82 \pm 0.28 \frac{1}{\sigma_{Z'}^{CCFR}} = 0.82 \pm 0.28 \frac{1}{\sigma_{Z'$$

V

μ

\_\_\_\_ N

## **U(1) model for maximal coupling**

$$Z' \rightarrow -Z', H_1 \leftrightarrow H_1, H_2 \leftrightarrow H_3$$

$$L_{H} = -\tilde{g}(\bar{l}_{2}\gamma^{\mu}Ll_{2} - \bar{l}_{3}\gamma^{\mu}Ll_{3} + \bar{e}_{2}\gamma^{\mu}Re_{2} - \bar{e}_{3}\gamma^{\mu}Re_{3})Z'_{\mu}$$
  
$$-[Y_{11}^{l}\bar{l}_{1}Re_{1}H_{1} + Y_{22}^{l}(\bar{l}_{2}Re_{2} + \bar{l}_{3}Re_{3})H_{1} + Y_{23}^{l}(\bar{l}_{2}Re_{3}H_{2} + \bar{l}_{3}Re_{2}H_{3})] + H.C.$$

• The transformation between the lepton mass eigen- state and weak eigenstate basis

$$\begin{pmatrix} \mu \\ \tau \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} e_2 \\ e_3 \end{pmatrix} , \quad \begin{pmatrix} \nu_\mu \\ \nu_\tau \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} \nu_2 \\ \nu_3 \end{pmatrix} .$$

The conflict can also be avoided if Z' only has  $\bar{\mu}\gamma^{\mu}\tau Z'_{\mu}$  type of interaction.

$$L_{\text{int}-Z'} = -\tilde{g}(\bar{\mu}\gamma^{\mu}\tau + \bar{\tau}\gamma^{\mu}\mu + \bar{\nu}_{\mu}\gamma^{\mu}L\nu_{\tau} + \bar{\nu}_{\tau}\gamma^{\mu}L\nu_{\mu})Z'_{\mu}.$$

• Give addition contribution to  $\tau \rightarrow \mu \nu \bar{\nu}$  decay which is highly constrained the model.

$$L_{Z'+W} = -\frac{g^2}{2m_W^2} \bar{\nu}_\tau \gamma^\mu L \nu_\mu \bar{\mu} \gamma_\mu L \tau - \frac{\tilde{g}^2}{m_{Z'}^2} (\bar{\nu}_\tau \gamma^\mu L \nu_\mu + \bar{\nu}_\mu \gamma^\mu L \nu_\tau) \bar{\mu} \gamma_\mu \tau .$$
$$R_{\tau\mu} = \Gamma(\tau \to \mu \nu \bar{\nu}) / \Gamma_{SM}(\tau \to \mu \nu \bar{\nu}) = 1.0066 \pm 0.0041$$

## U(1) model for maximal coupling

• Introduce scalar triplet:  $\Delta_i$ : (1,3,1)

$$\Delta = \begin{pmatrix} \Delta^+ / \sqrt{2} & \Delta^{++} \\ \Delta^0 & -\Delta^+ / \sqrt{2} \end{pmatrix}$$

Yukawa terms in the mass-eigen basis  

$$L_{\Delta} = -(\bar{\nu}_{e}^{c}, \bar{\nu}_{\mu}^{c}, \bar{\nu}_{\tau}^{c})M(\Delta^{0})L\begin{pmatrix}\nu_{e}\\\nu_{\mu}\\\nu_{\tau}\end{pmatrix} + \sqrt{2}(\bar{\nu}_{e}^{c}, \bar{\nu}_{\mu}^{c}, \bar{\nu}_{\tau}^{c})M(\Delta^{+})L\begin{pmatrix}e\\\mu\\\tau\end{pmatrix} + (\bar{e}^{c}, \bar{\mu}^{c}, \bar{\tau}^{c})M(\Delta^{++})L\begin{pmatrix}e\\\mu\\\tau\end{pmatrix}$$

$$M(\Delta) = \begin{pmatrix}Y_{11}^{\nu}\Delta_{1} & 0 & 0\\ 0 & (Y_{22}^{\nu}(\Delta_{2} + \Delta_{3}) - 2Y_{23}^{\nu}\Delta_{1})/2 & Y_{22}^{\nu}(\Delta_{2} - \Delta_{3})/2\\ 0 & Y_{22}^{\nu}(\Delta_{2} - \Delta_{3})/2 & (Y_{22}^{\nu}(\Delta_{2} + \Delta_{3}) + 2Y_{23}^{\nu}\Delta_{1})/2 \end{pmatrix}$$

#### Simplification

1.  $Y_{11,23} < < Y_{22}, m_{\Delta 2} = m_{\Delta 3}$  2.  $m_{\Delta^{++}} = m_{\Delta^{+}} = m_{\Delta} > 420 \text{GeV}$ 3. Other heavier new degrees of freedom

## **U(1) model for maximal coupling**

### **Model parameters**

Four parameters  $U(1)_{L_u-L_\tau}$ :  $\tilde{g}, m_{Z'} + Y = 1$  triplet :  $Y_{22}, m_{\Delta}$ Revisiting type-II see-saw: present limits and Neutrino trident future prospects at LHC  $\begin{bmatrix} \tilde{g}^2 \\ m_{\pi'}^2 \end{bmatrix} \in (3,8) \times 10^{-7} \text{GeV}^{-2}$ Saiyad Ashanujjaman 🖂 & Kirtiman Ghosh Journal of High Energy Physics **2022**, Article number: 195 (2022) Cite this article  $|Y_{22}^{\nu}|^2/m_{\Delta}^2$  $\Delta m = 0$  with  $v_{\Lambda} \sim O(\text{GeV})$  $\tau \rightarrow \mu \bar{\nu} \nu$  $10^{-7}$ 

 $m_{\Lambda} > 420 \text{GeV}$ 



Highlight Z' effects and suppress the triplet contribution  $m_{\Lambda} = 450 \text{GeV}, |Y_{22}| = 0.117$ 







TABLE III: The cross section  $\tau^{\pm}\tau^{\mp}$  for  $U(1)_{L_{\mu}-L_{\tau}}$  model with Y = 1 triplet at  $\sqrt{s} = 3$  TeV for fixing  $m_{Z'} = 100$  GeV and  $\tilde{g} = 0.055$ .

$U(1)_{L_{\mu}-L_{\tau}}$ with triplet model	$m_{\Delta} = 800 \text{ GeV}$		$m_{\Delta} = 500 \text{ GeV}$		$m_{\Delta} = 450 \text{ GeV}$	
	$Y_{22} = 0.208$	$Y_{22} = 1.136$	$Y_{22} = 0.13$	$Y_{22} = 0.71$	$Y_{22} = 0.117$	$Y_{22} = 0.639$
cross section (pb)	0.0069	2.6156	0.0019	0.5180	0.0014	0.3569
luminosity (fb <sup>-1</sup> ) with $3\sigma$	4.00	0.0035	39.287	0.018	72.2335	0.026
Events ( $\mathcal{L} = 1ab^{-1}$ )	6900	2615600	1900	518000	1400	356900





t-channel pair production can easily be distinguished at more than 5\sigma level from the s-channel production in SM

 $\mu^+\mu^- \rightarrow \mu^+\mu^+\tau^-\tau^-$ 

#### This signal is very clean and effectively background-free.





Luminosity for  $\mu^+\mu^- \rightarrow \mu^+\mu^+\tau^-\tau^-$ 



(a). The required luminosity for  $3\sigma$  and  $5\sigma$  discovery with  $\tilde{g}/M_{Z'} = 0.55 \times 10^{-3} \text{GeV}^{-1}$ .





- ► We study in detail the maximal off-diagonal Z' interaction in  $U(1)_{L_{\mu}-L_{\tau}}$  at a muon collider
- > A Z' with off-diagonal mixing leads to very distinctive signatures, such as t-channel  $\mu^+\mu^- \rightarrow \tau^+\tau^-$  and  $\mu^+\mu^- \rightarrow \mu^+\mu^+\tau^-\tau^$ smoking gun
- > With a 3TeV muon collider with O(fb<sup>-1</sup>) luminosity, the above two processes can be distinguished at  $5\sigma$  level. Thanks !



## **Cross section of** $\mu^+\mu^- \rightarrow \tau^+\tau^-$



TABLE II: The cross section  $\tau^{\pm}\tau^{\mp}$  for  $U(1)_{L_{\mu}-L_{\tau}}$  model with Y = 1 triplet at  $\sqrt{s} = 3$  TeV for fixing  $m_{\Delta} = 450$  GeV and  $Y_{22} = 0.117$ .

$U(1)_{L_{\mu}-L_{\tau}}$ with triplet model	$m_{Z'} = 500 \text{ GeV}$		$m_{Z'} = 200 \text{ GeV}$		$m_{Z'} = 100 \text{ GeV}$	
	$\tilde{g} = 0.275$	$ \tilde{g} = 0.445$	$\tilde{g} = 0.11$	$\tilde{g} = 0.178$	$\tilde{g} = 0.055$	$  ilde{g} = 0.089 $
cross section (pb)	0.274	2.08	0.017	0.154	0.0014	0.01
luminosity (fb <sup>-1</sup> ) with $3\sigma$	0.034	0.004	0.97	0.063	72.2335	2.189
Events ( $\mathcal{L} = 1ab^{-1}$ )	274000	2080000	17000	154000	1400	10000