



北京航空航天大学
BEIHANG UNIVERSITY

宇宙学一级相变与物质、暗物质 和原初黑洞的起源

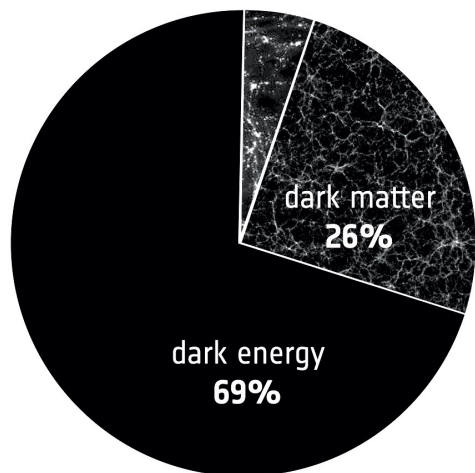
谢柯盼 (Ke-Pan Xie)

北京航空航天大学

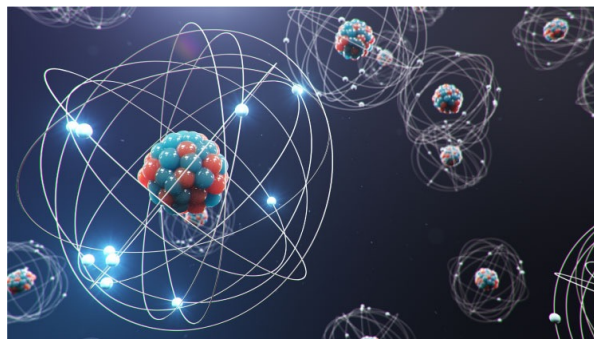
2023.4.26 @中科院高能所

我们生活的宇宙，充满了谜团

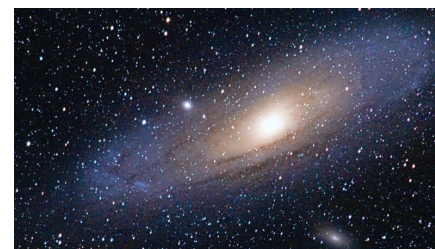
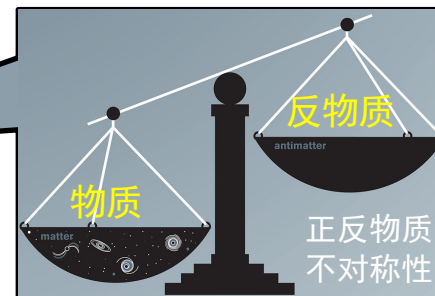
当今宇宙的能量分布



可见物质 5%



反物质失踪了!



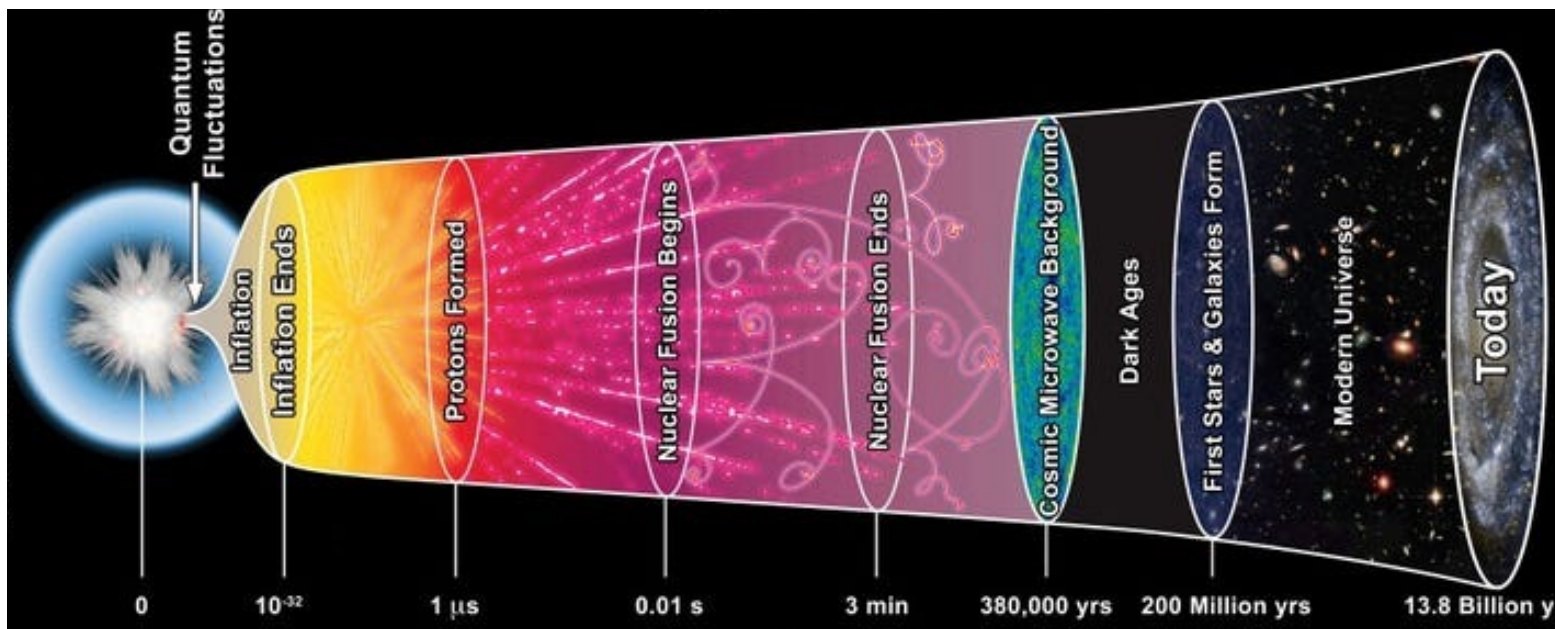
粒子物理标准模型——“已知的宇宙”

暗物质 26% [???

暗能量 69% [???

} 未知的宇宙.....

这些谜团存在已久

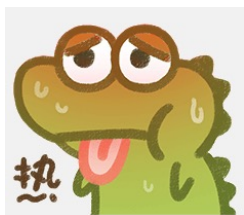


大爆炸后约0.1纳秒~1秒内，**正反物质不对称性**就已经产生，暗物质已经形成

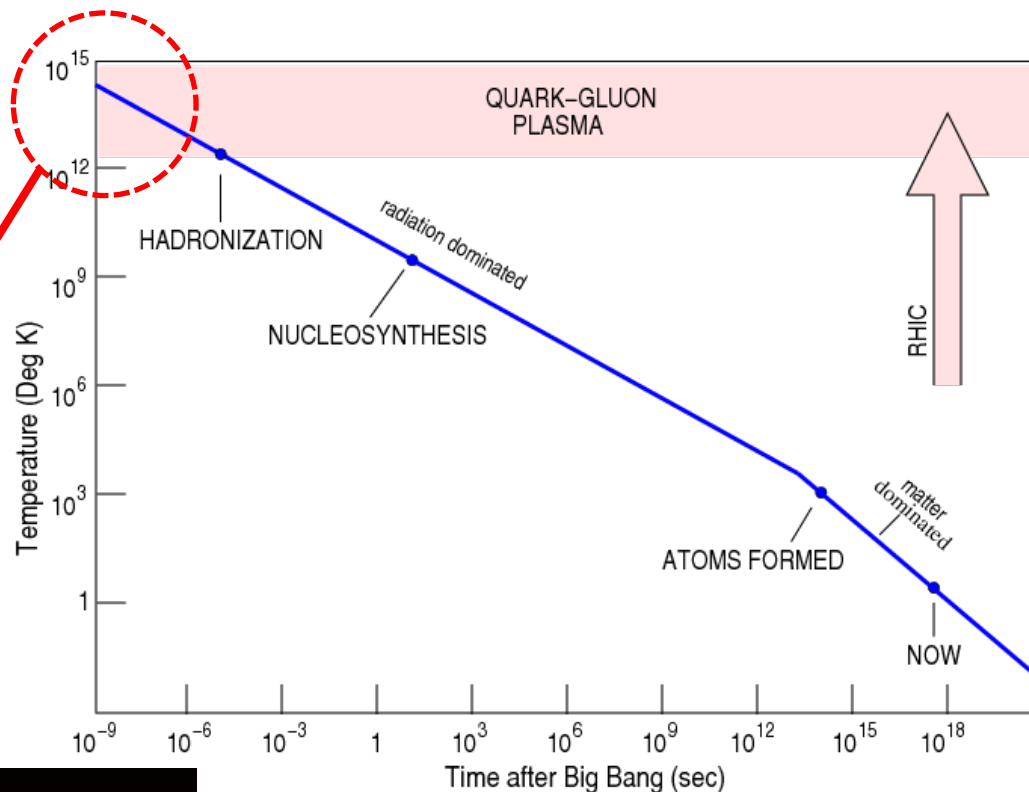
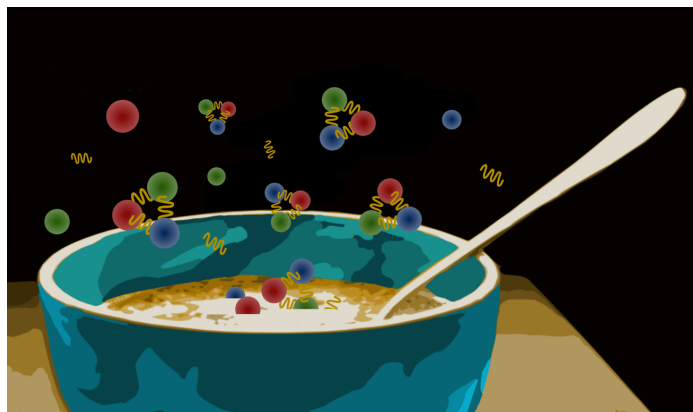
- 1920年代起，就陆续有暗物质存在的证据；至1970年代，暗物质成为被广泛承认的未解之谜，持续至今
- 正反物质不对称性于1960年代被注意到，一直没有妥善的解答

追溯早期宇宙的历史

早期宇宙的显著特点:



粒子宇宙学



基本粒子（电子、中微子、夸克……）以等离子体的形式存在

宇宙像一锅炽热的汤

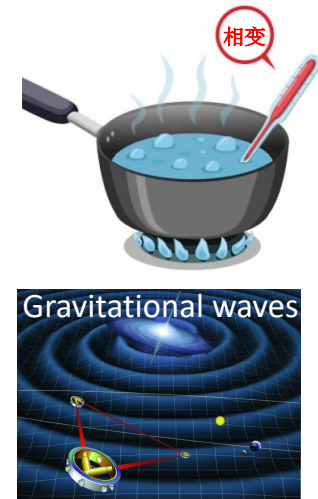
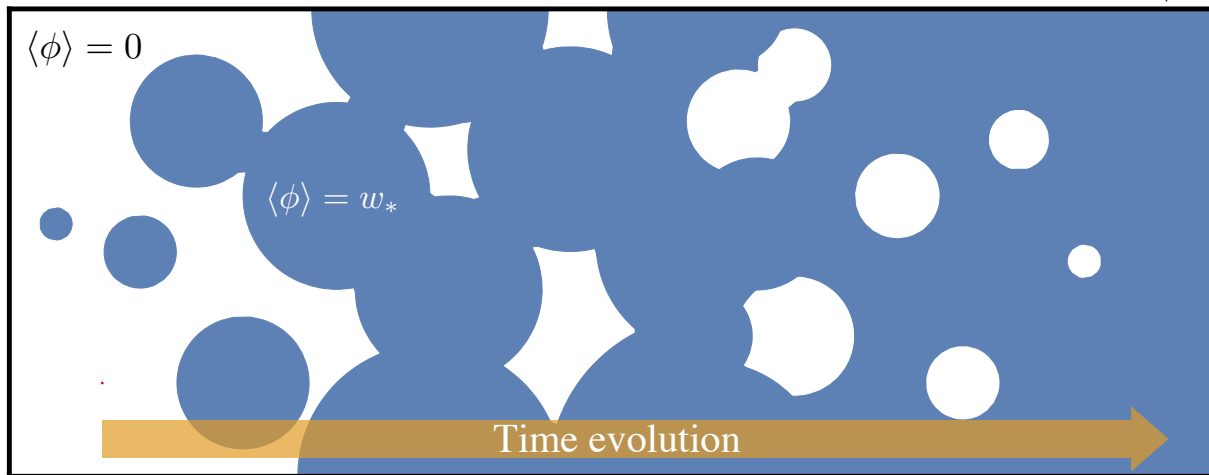
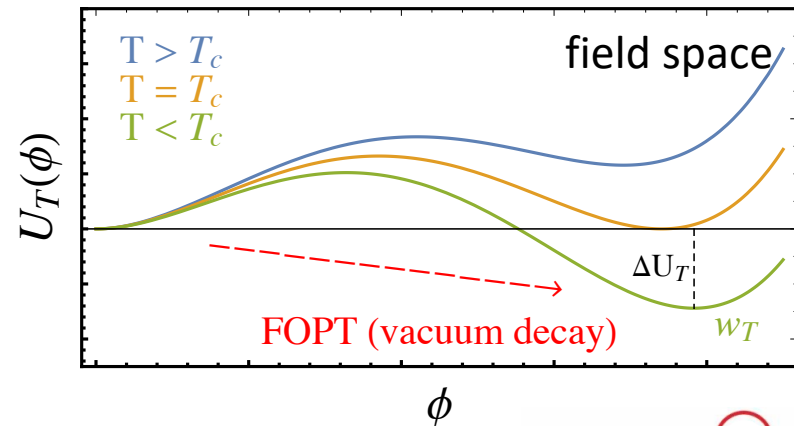
Solving the puzzles with a first-order phase transition

Vacuum decay

$$\mathcal{L} \supset \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - U(\phi)$$



Thermal corrections
 $U(\phi) \Rightarrow U_T(\phi, T)$



Boiling of the Universe, vacuum bubble nucleation and expansion

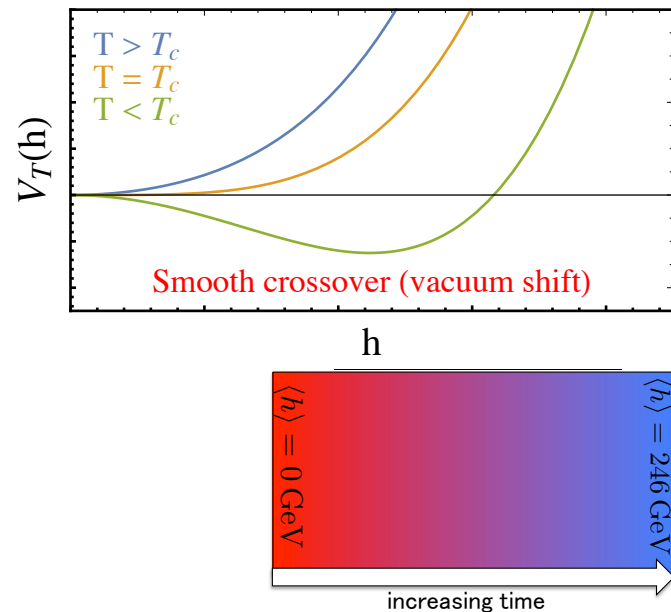
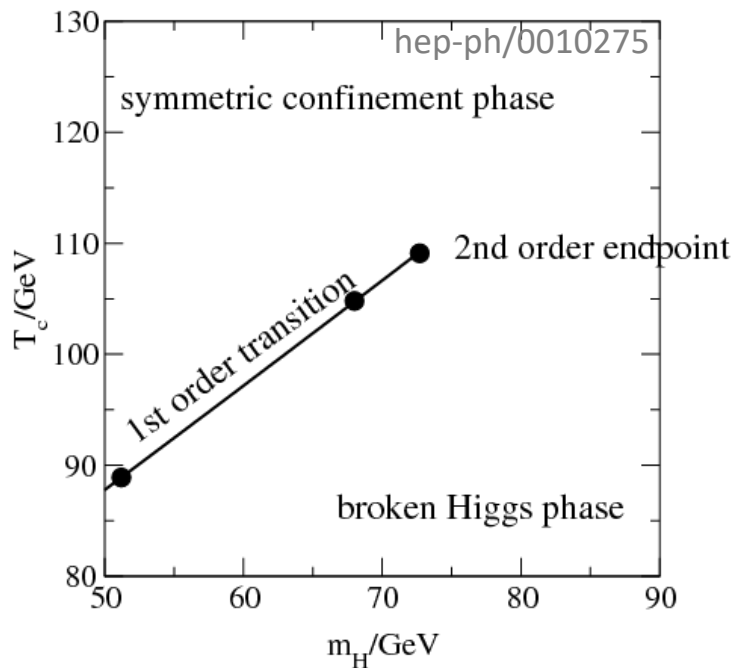
If $\phi = h$ (SM Higgs) – first-order electroweak phase transition

If $\phi =$ new physics field – general FOPT

3 main directions of research on a boiling Universe

1. The dynamics of FOPT;
2. FOPTs in new physics models;
3. New physics mechanisms based on FOPTs

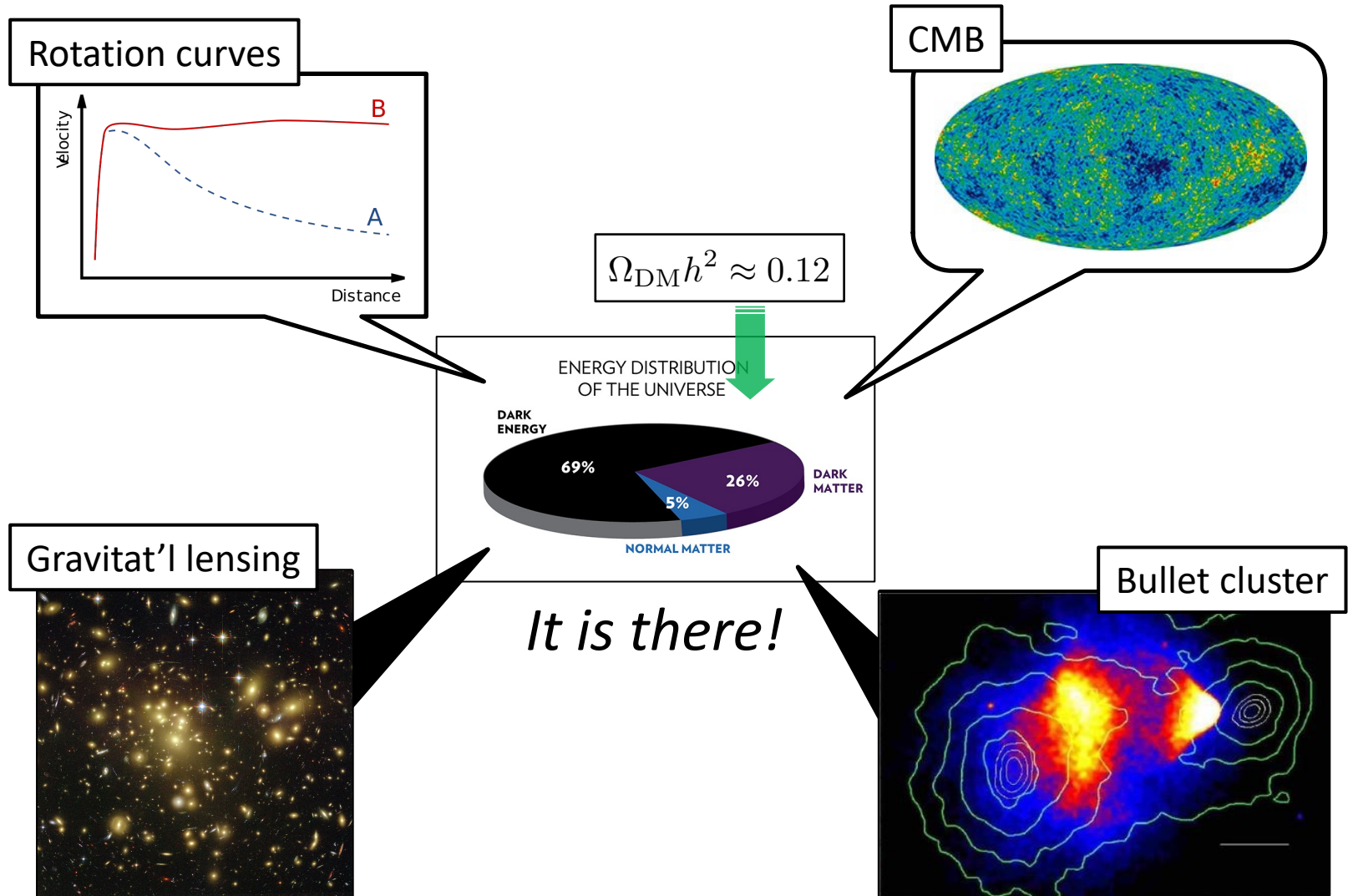
There is NO FOPT in the Standard Model



FOPT \Rightarrow new physics

First-order EW phase transition of h , or FOPT of a new field ϕ

The mystery of dark matter

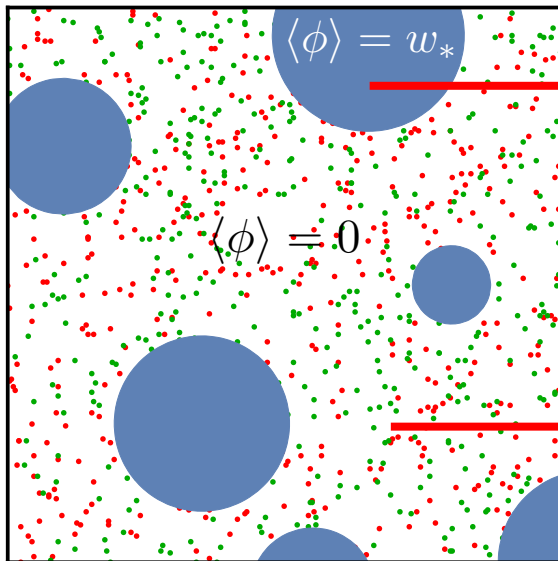


Dark matter & FOPTs

Relevant Lagrangian

$$\mathcal{L} \supset \underbrace{\frac{1}{2} \partial_\mu \phi \partial^\mu \phi - U(\phi)}_{\text{FOPT Lagrangian}} + \bar{\chi} (i\gamma^\mu \partial_\mu - m_0) \chi - \underbrace{y_\chi \phi \bar{\chi} \chi}_{\text{Portal coupling}}$$

Dark fermion



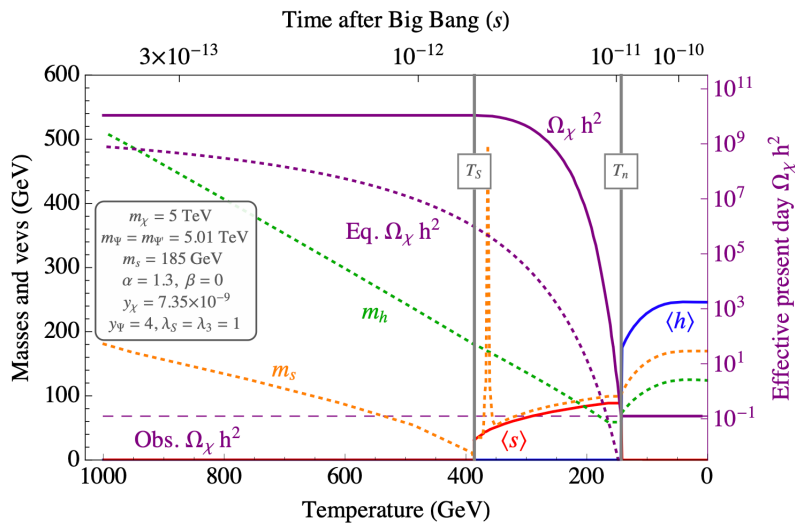
Inside the bubble:
 $m_\chi = m_0 + y_\chi w_*$

Outside the bubble:
 $m_\chi = m_0$

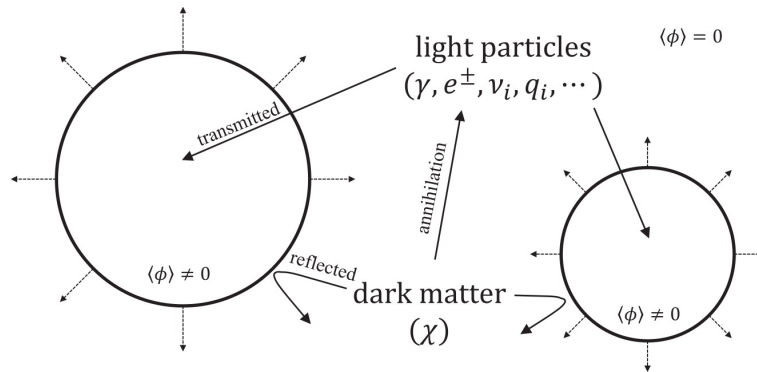
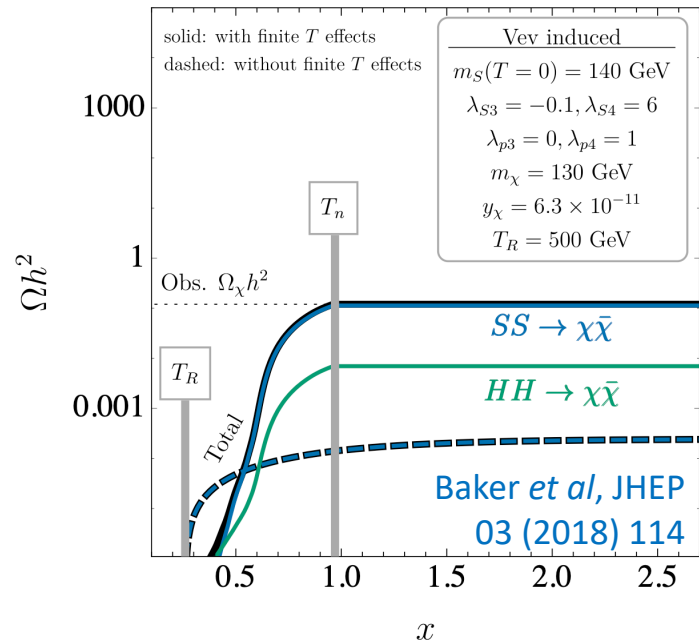


Mass discontinuity between two sides of the bubble wall

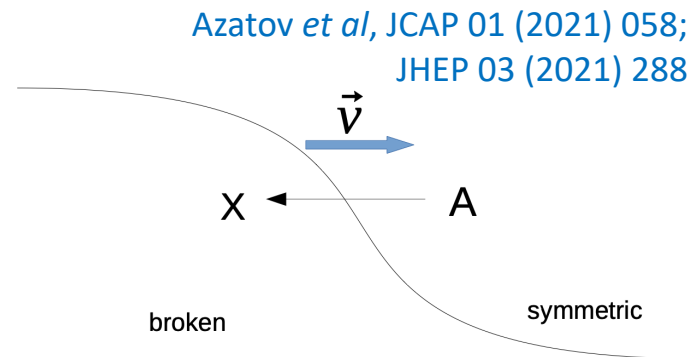
Dark matter mechanisms based on a FOPT



Baker et al, PRL 119 (2017) 6, 061801



Baker et al, PRL 125 (2020) 15, 151102
Chway et al, PRD 101 (2020) 9, 095019



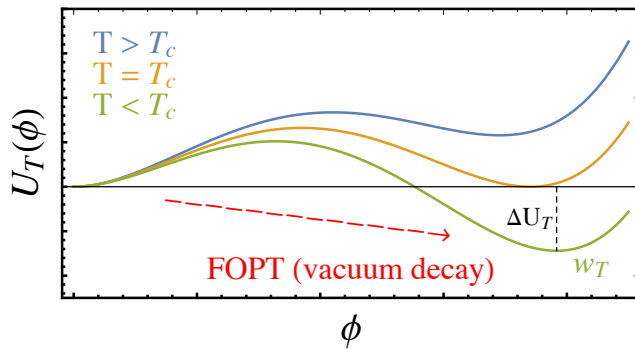
The fermion soliton dark matter from a FOPT

$$\mathcal{L} = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - U(\phi) + \bar{\chi} i \gamma^\mu \partial_\mu \chi - y_\chi \phi \bar{\chi} \chi$$

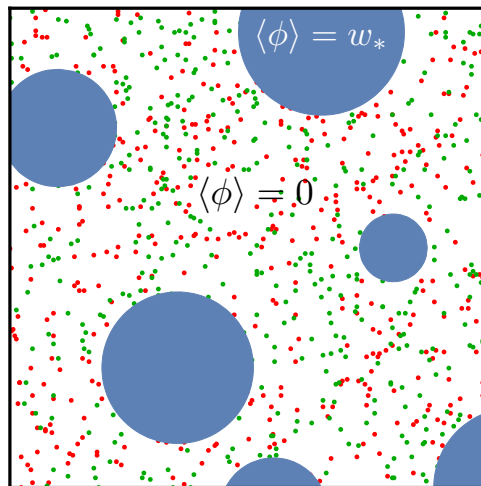
FOPT Scalar
Fermion

Potential for phase transition
Yukawa vertex

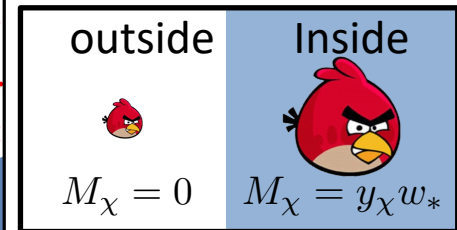
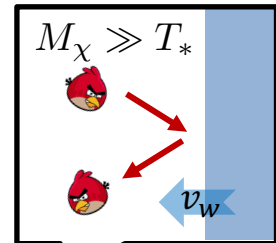
The interaction between bubble and fermion!



Bubble nucleation

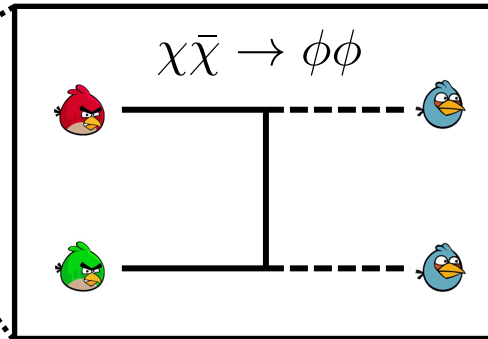
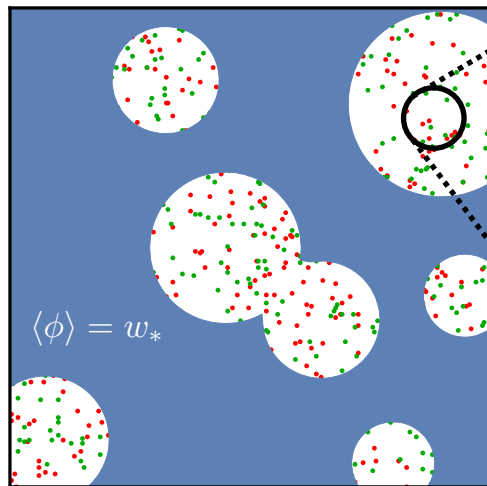


Reflected!



Example: trapping fraction = 98% for $M_\chi/T_* = 12$ and $v_w = 0.6$

What happens for the trapped fermions?



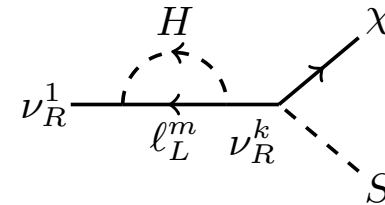
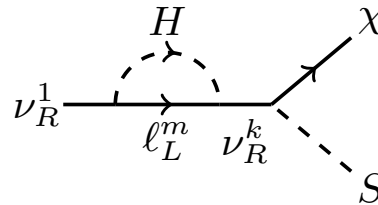
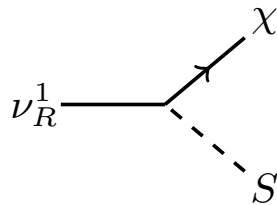
$$-y_\chi \phi \bar{\chi} \chi$$



Eventually to SM particles via portal coupling $\lambda_{H\phi} |H|^2 \phi^2$

To have a nontrivial result, there should be $N(\chi) \neq N(\bar{\chi})$

1. Thermal fluctuation; [Asadi et al, PRL 127 (2021) 21, 211101]
2. A baryogenesis-like asymmetry; [Shelton et al, PRD 82 (2010) 123512]

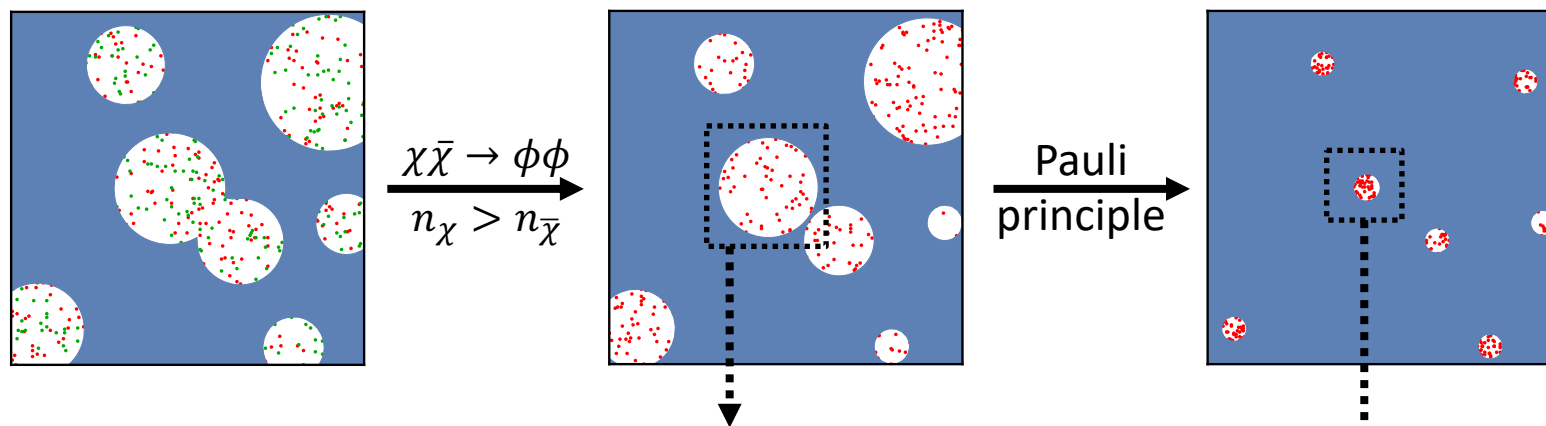


$$\Gamma(\nu_R^1 \rightarrow \chi S) > \Gamma(\nu_R^1 \rightarrow \bar{\chi} S)$$

$$\eta_\chi \equiv \frac{n_\chi - n_{\bar{\chi}}}{s}$$

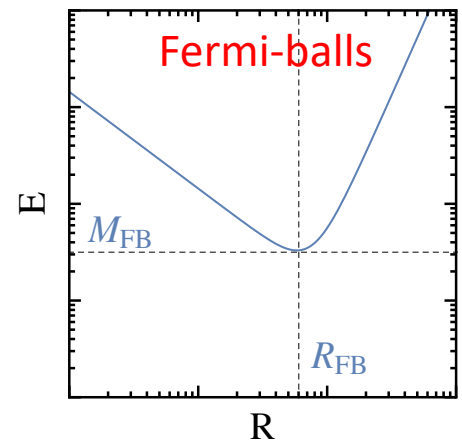
Similar to baryon asymmetry of the Universe

Formation of Fermi-ball solitons



Charge trapped: $Q_{\text{FB}} \approx \eta_{\chi} s_* V_*$ ← Remnant volume

Fermion asymmetry $\eta_{\chi} = \frac{n_{\chi} - n_{\bar{\chi}}}{s}$



$$E = \frac{3\pi}{4} \left(\frac{3}{2\pi}\right)^{2/3} \frac{Q_{\text{FB}}^{4/3}}{R} + \frac{4\pi}{3} R^3 U_0$$

Fermi-gas energy Volume energy

$$M_{\text{FB}} = Q_{\text{FB}} (12\pi^2 U_0)^{1/4};$$

$$R_{\text{FB}} = Q_{\text{FB}}^{1/3} \left[\frac{3}{16} \left(\frac{3}{2\pi}\right)^{2/3} \frac{1}{U_0} \right]^{1/4}$$

Hong, Jung and KPX, PRD 102 (2020) 7, 075028

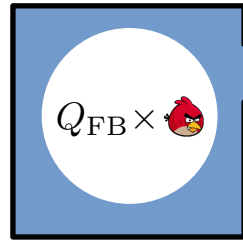
The Fermi-ball profile

$$M_{\text{FB}} = Q_{\text{FB}} (12\pi^2 U_0)^{1/4};$$

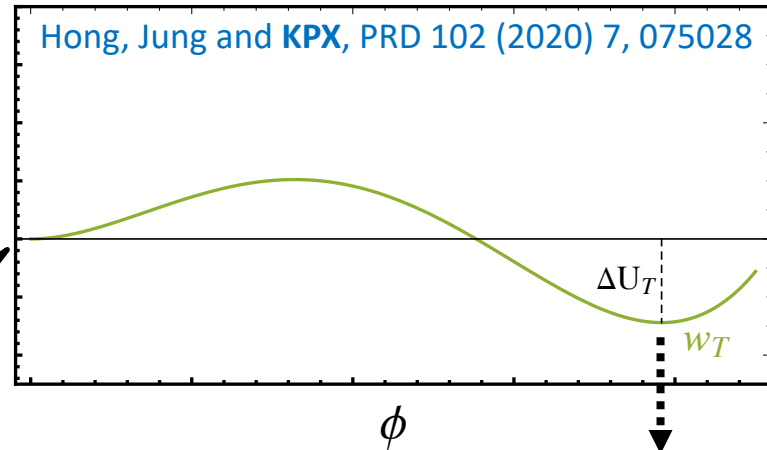
$$R_{\text{FB}} = Q_{\text{FB}}^{1/3} \left[\frac{3}{16} \left(\frac{3}{2\pi} \right)^{2/3} \frac{1}{U_0} \right]^{1/4}$$

Effective χ -mass
inside the ball:

$$M_{\text{eff}} = (12\pi^2 U_0)^{1/4}$$



$U_T(\phi)$



Mass at true vacuum $M_\chi = y_\chi w$

$$\text{Stability condition: } \frac{dM_{\text{FB}}}{dQ_{\text{FB}}} = M_{\text{eff}} < M_\chi$$



*Fission stability $d^2 M_{\text{FB}}/dQ_{\text{FB}}^2 < 0$ always satisfied with surface term $\propto Q_{\text{FB}}^{2/3}$

Estimates

$$Q_{\text{FB}} \approx 10^{42} \times v_w^3 \left(\frac{\eta_\chi}{10^{-3}} \right) \left(\frac{100}{g_*} \right)^{1/2} \left(\frac{100 \text{ GeV}}{T_*} \right)^3 \left(\frac{100}{\beta/H_*} \right)^3$$

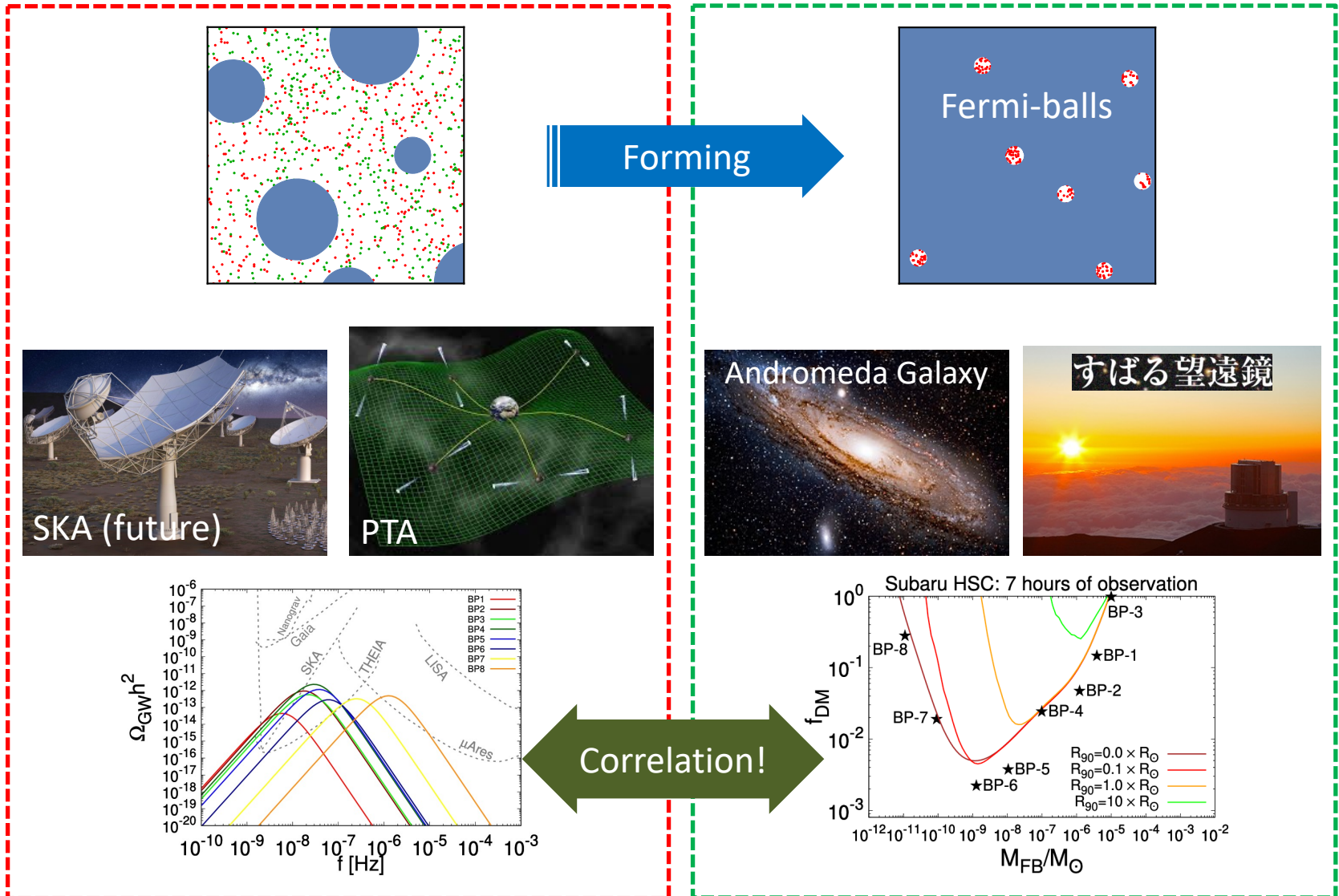
$$M_{\text{FB}} \approx 1.4 \times 10^{21} \text{ g} \times v_w^3 \left(\frac{\eta_\chi}{10^{-3}} \right) \left(\frac{100}{g_*} \right)^{1/4} \left(\frac{100 \text{ GeV}}{T_*} \right)^2 \left(\frac{100}{\beta/H_*} \right)^3 \alpha^{1/4}$$

$$R_{\text{FB}} \approx 4.8 \times 10^{-3} \text{ cm} \times v_w \left(\frac{\eta_\chi}{10^{-3}} \right)^{1/3} \left(\frac{100}{g_*} \right)^{5/12} \left(\frac{100 \text{ GeV}}{T_*} \right)^2 \left(\frac{100}{\beta/H_*} \right) \alpha^{-1/4}$$

Macroscopic Dark Matter!

An example of phenomenology

Marfatia *et al*, JHEP 11 (2021) 068



Big picture: fermion-type solitons

Proposed by T. D. Lee. [PRD.15.1694, PRD.16.1096]

Fermion-field nontopological solitons*

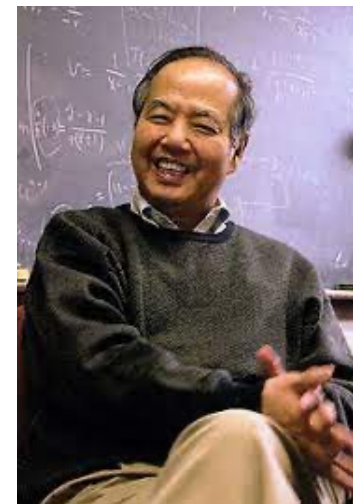
R. Friedberg

Barnard College and Columbia University, New York, New York 10027

T. D. Lee

Columbia University, New York, New York 10027

(Received 8 December 1976)

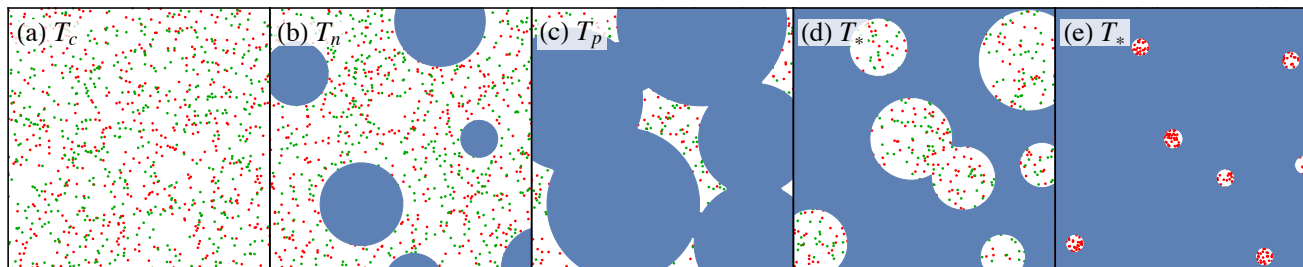


Fermion-soliton is possible with a scalar

Theorem 1. There exists a critical value N_S . For $N > N_S$, the lowest-energy state is a soliton, not the plane-wave solution. Furthermore, as $N \rightarrow \infty$,

$$E \leq \frac{4}{3}\pi\sqrt{2} N^{3/4} [U(-m/g)]^{1/4}. \quad (2.1)$$

Our work provides a dynamical mechanism to form such solitons

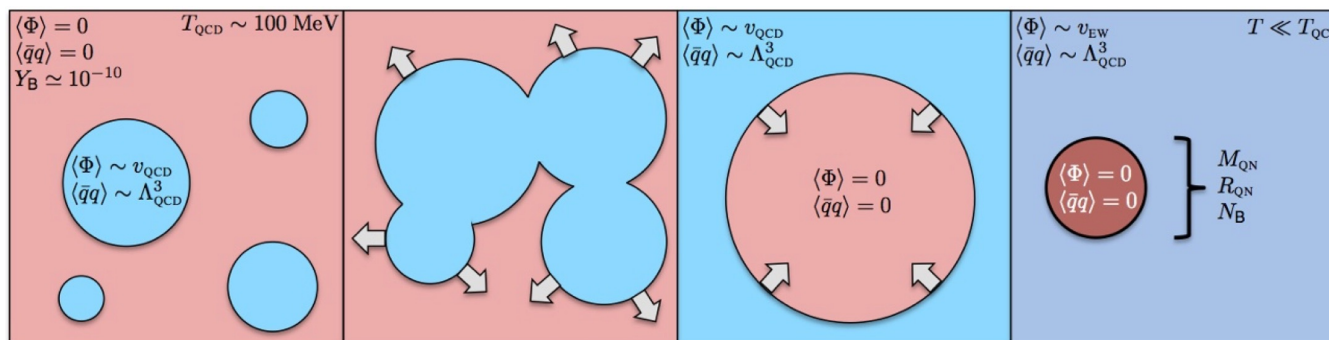


Hong, Jung and KPX, PRD 102 (2020) 7, 075028

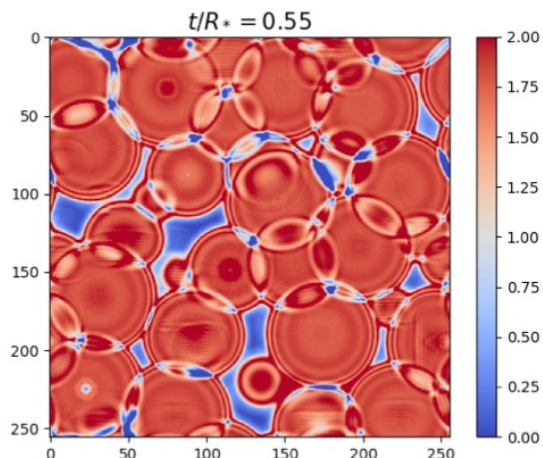
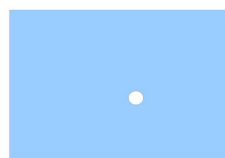
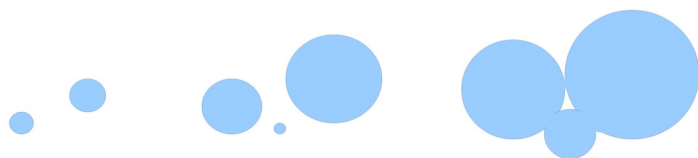
Big picture: solitons formation during a FOPT

Fermions trapped to form dark dwarfs [Gross *et al*, JHEP 09 (2021) 033]

Quark-nuggets in a QCD FOPT [Bai *et al*, JHEP 06 (2018) 072; PRD 99 (2019) 5, 055047]



Scalar Q-balls [Krylov *et al*, PRD 87 (2013) 8, 083528; Huang *et al*, PRD 96 (2017) 9, 095028]



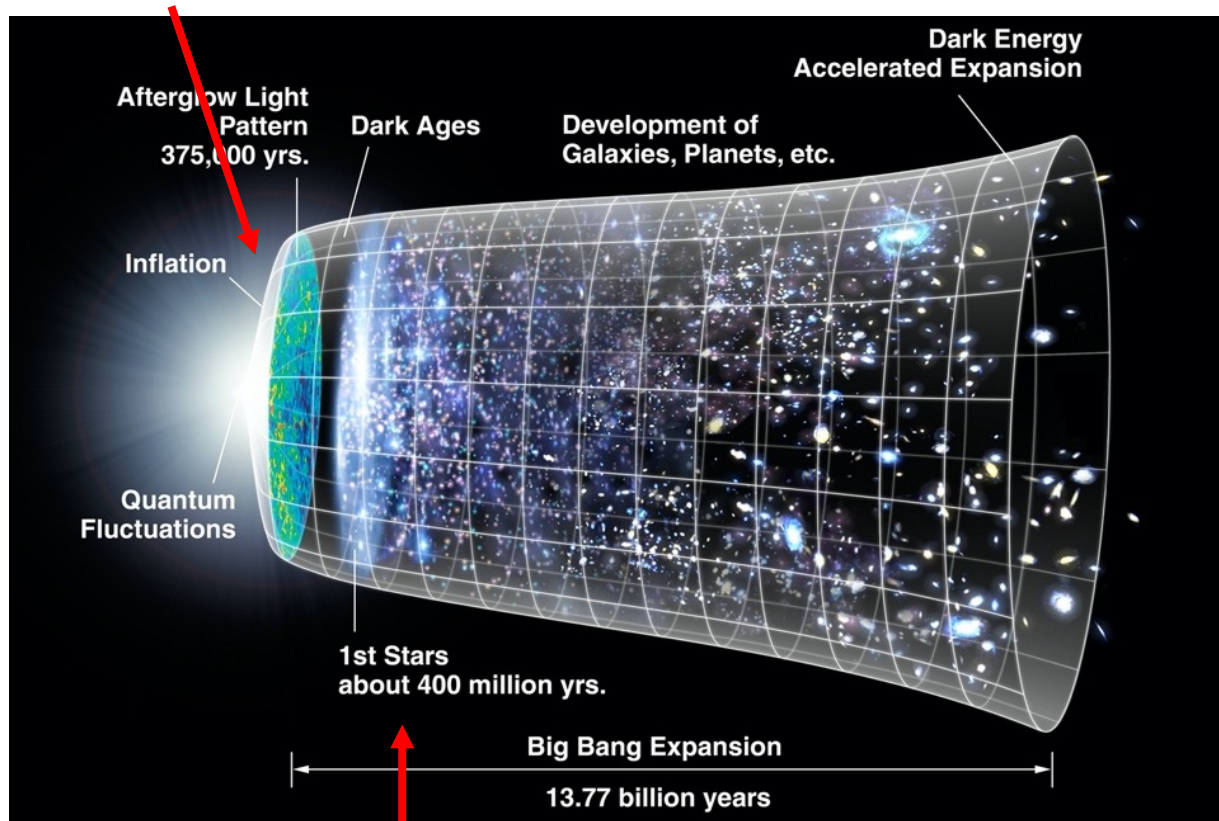
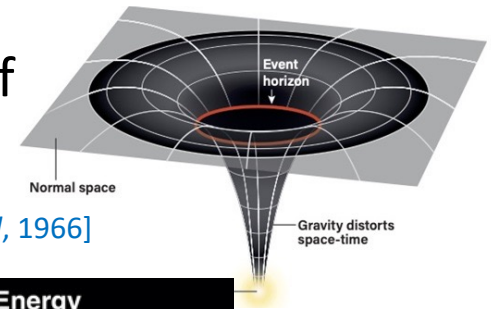
Topological solitons—

- Monopoles [Bian *et al*, PLB 839 (2023) 137822]
- Cosmic strings [Bian, Cai *et al*, 2204.04427]
- Domain walls [Jiang *et al*, 2208.07186]

From solitons to black holes

Almost all kinds of solitons have the possibility of **collapsing into black holes**

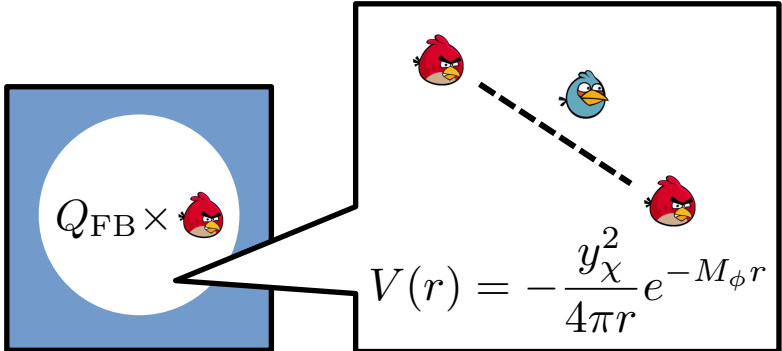
Primordial black holes (soon after Big Bang); [Zel'dovitch *et al*, 1966]



Astrophysical black holes (from stellar collapse)

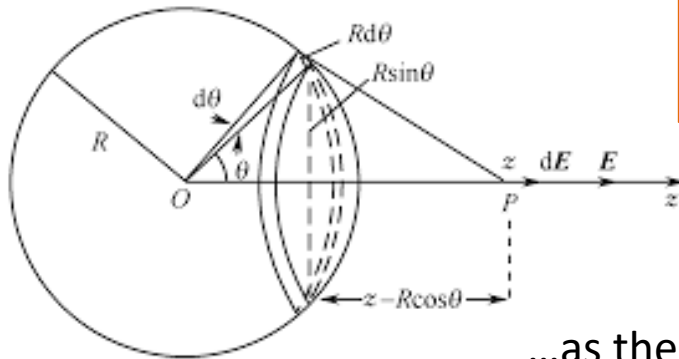
Attractive force inside a soliton

Yukawa force *inside* a Fermi-ball



The diagram shows a blue square containing a white circle representing a Fermi ball with charge $Q_{\text{FB}} \times$ and a red Angry Bird icon. A callout box shows a dashed line between two red Angry Birds and a blue Angry Bird, representing the Yukawa potential $V(r) = -\frac{y_\chi^2}{4\pi r} e^{-M_\phi r}$. To the right, text states "Originates from $\mathcal{L} \supset -y_\chi \phi \bar{\chi} \chi$ ".

A "textbook exercise" calculation...



$$E_{\text{Yuk}} \approx -\frac{15y_\chi^2}{40\pi} \frac{Q_{\text{FB}}^2}{R} \left(\frac{L_\phi}{R}\right)^2$$

$L_\phi = M_\phi^{-1}$: range of force

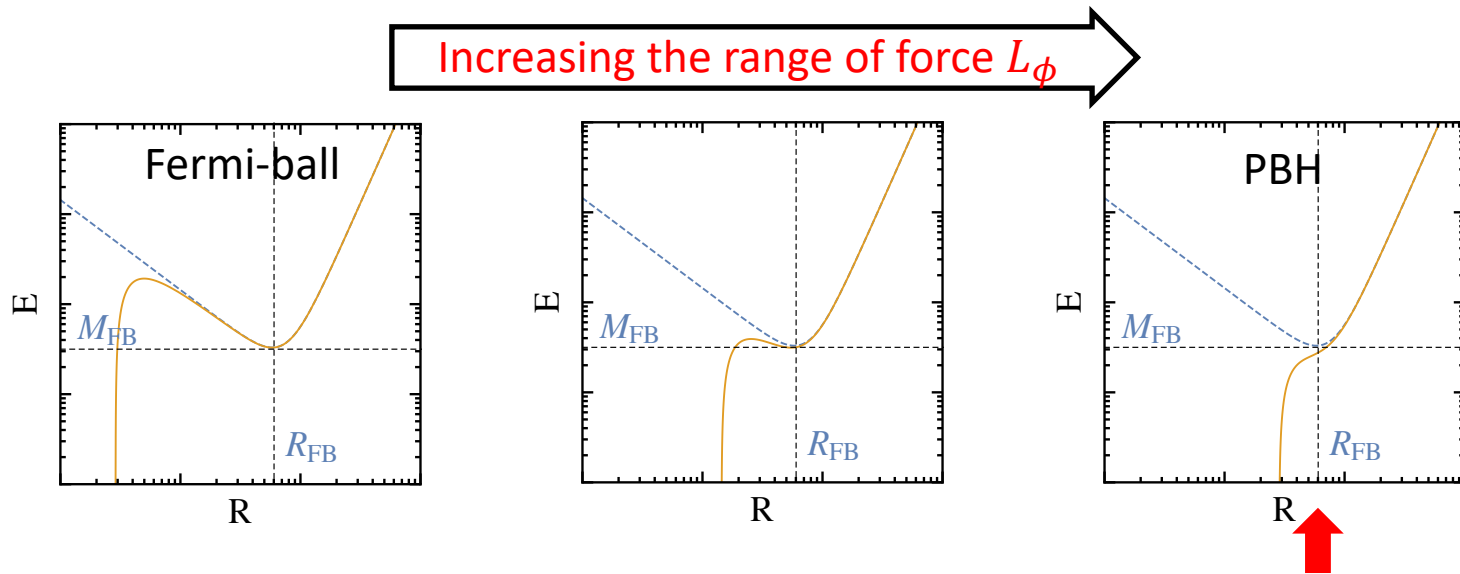
...as the leading-order approximation

The complete energy profile

The improved energy profile (when $L_\phi = M_\phi^{-1} \ll R$)

$$E \approx \frac{3\pi}{4} \left(\frac{3}{2\pi}\right)^{2/3} \frac{Q_{\text{FB}}^{4/3}}{R} + \frac{4\pi}{3} R^3 U_0 - \frac{15y_\chi^2}{40\pi} \frac{Q_{\text{FB}}^2}{R} \left(\frac{L_\phi}{R}\right)^2$$

What if the Yukawa energy dominates?



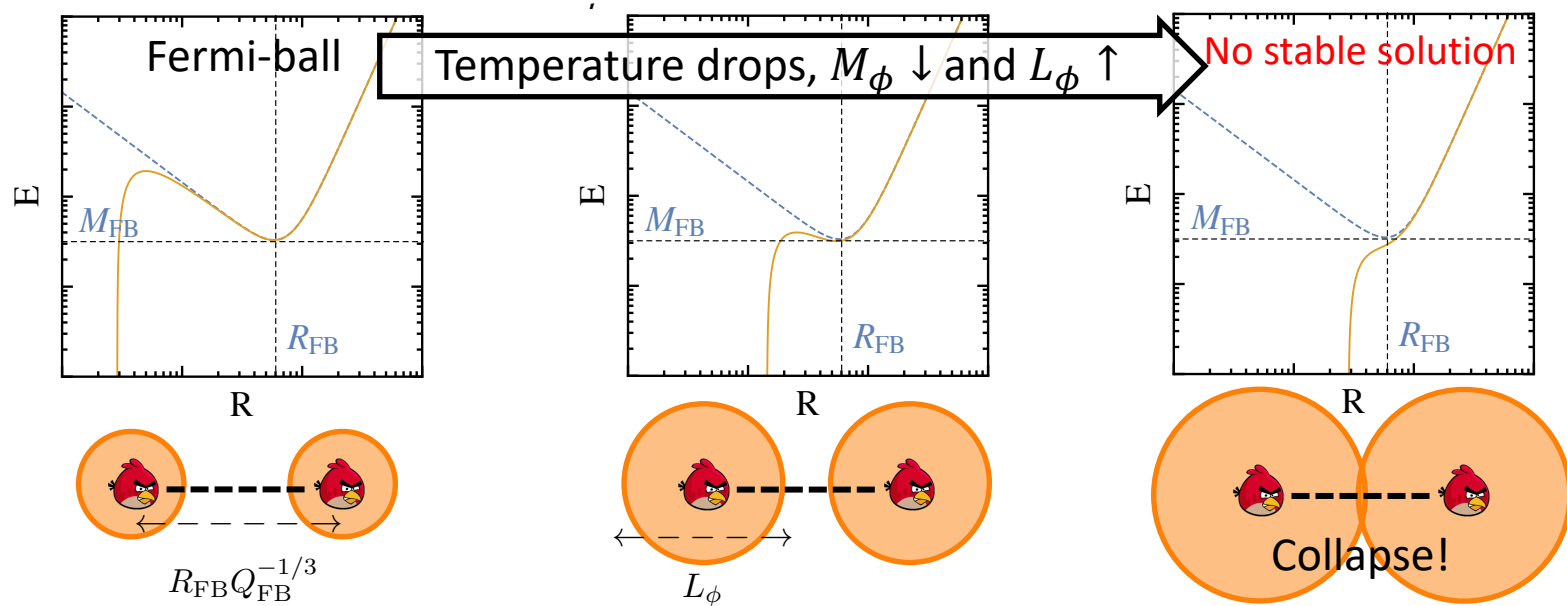
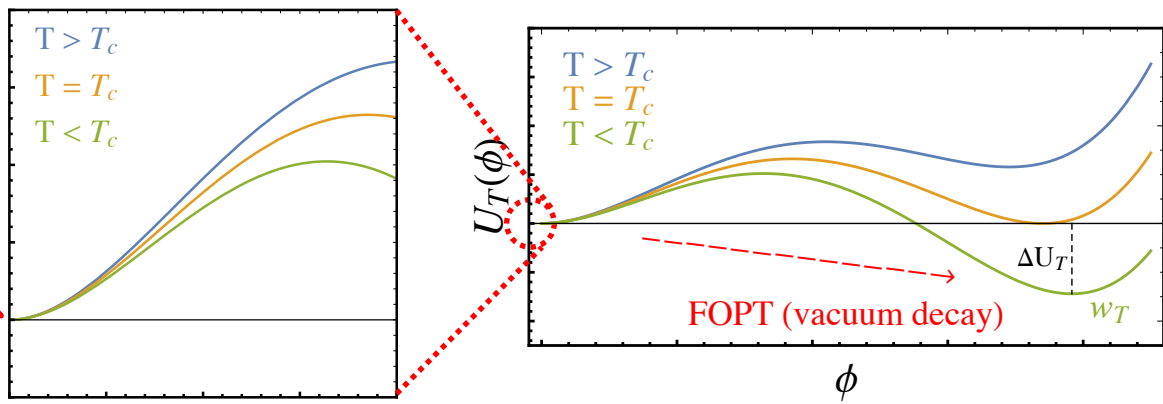
A Fermi-ball collapses in to a black hole!

$$L_\phi \sim R_{\text{FB}} Q_{\text{FB}}^{-1/3}$$

Fermi-balls may collapse when they cool down

The range of force **increases** as temperature drops!

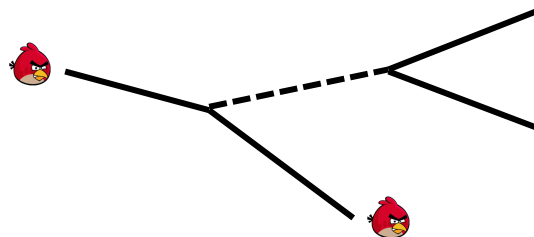
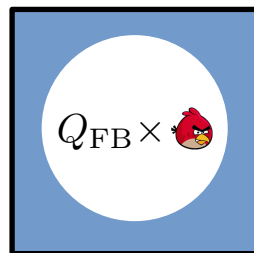
$$M_\phi = \sqrt{\mu^2 + cT^2}$$



Kawana and KPX, PLB 824 (2022) 136791

The cooling of Fermi-balls

Emitting SM light particles (black body radiation [\[Witten, PRD1984\]](#));



$e^+ e^-, \nu \bar{\nu}, \dots$

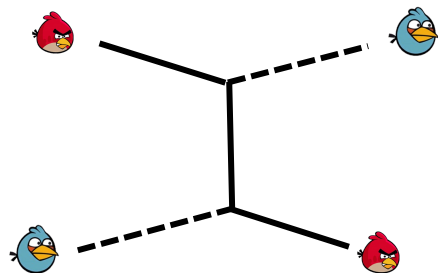
Energy decreasing rate

$$L(T) = \frac{7\pi^3 N_f}{240} R_{\text{FB}}^2 T^4$$

$$\tau_{\text{cool}} = \frac{240}{7\pi^2} \left(\frac{2\pi}{3} \right)^{1/3} \frac{Q_{\text{FB}}^{1/3} (12\pi^2 U_0)^{1/4}}{N_f T^2}.$$

Radiation cooling is very efficient: $\tau_{\text{cool}} \ll 1/H$

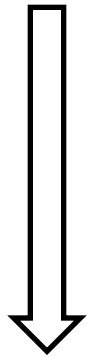
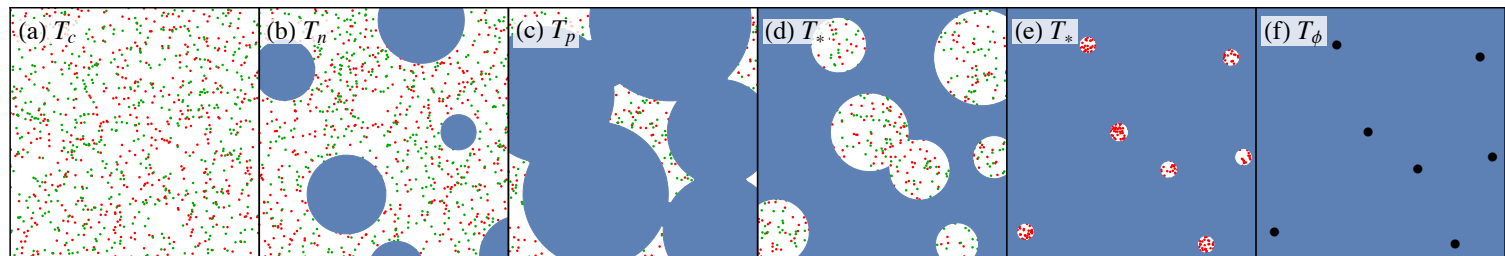
Scattering cooling: [\[Kawana, Lu and KPX, JCAP 10 \(2022\) 030\]](#)



In thermal bath via
 $\lambda_{H\phi} |H|^2 \phi^2$

In short: Fermi-balls can cool down!

PBHs from Fermi-ball collapse from a FOPT



Kawana and KPX, PLB 824 (2022) 136791

Advantages:

- Friendly to particle physicists;
- Can control the **formation time** of PBHs

Mass distribution: Lu, Kawana and KPX, PRD 105 (2022) 12, 123503

New soliton & PBH solutions: Kawana, Lu and KPX, JCAP 10 (2022) 030

EWPT & PBH DM: Huang and KPX, PRD 105 (2022) 11, 115033

PBH & gamma-rays: Marfatia et al, JHEP 08 (2022) 001

511 keV galactic line: Tseng et al, 2209.01552

PBHs after the CMB: Lu et al, 2210.16462

Boosted DM: Marfatia et al, JHEP 04 (2023) 006

Distinguishing different mechanisms: KPX, 2301.02352

GWs at PTA: Tseng et al, 2304.10084

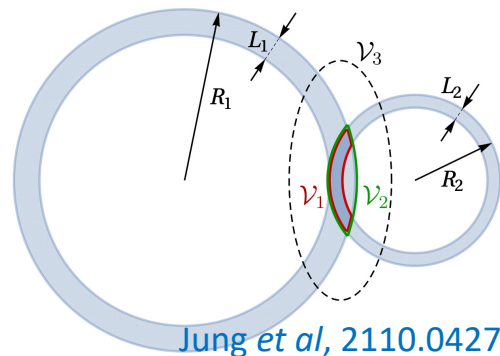
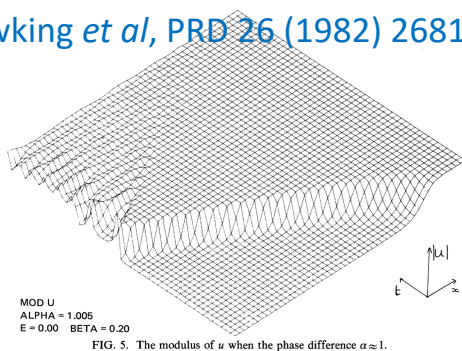
Mechanism

Application

Big picture: primordial black holes from FOPTs

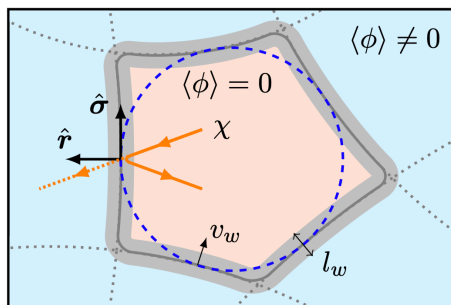
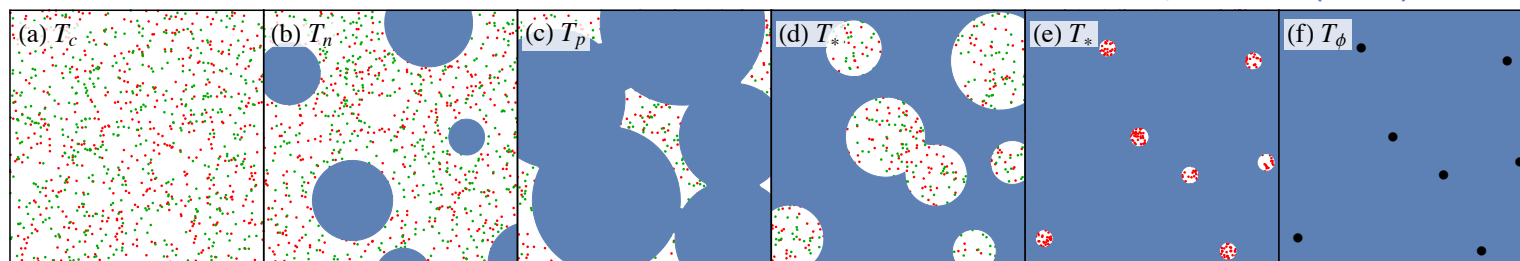
- Bubble collisions

Hawking *et al*, PRD 26 (1982) 2681



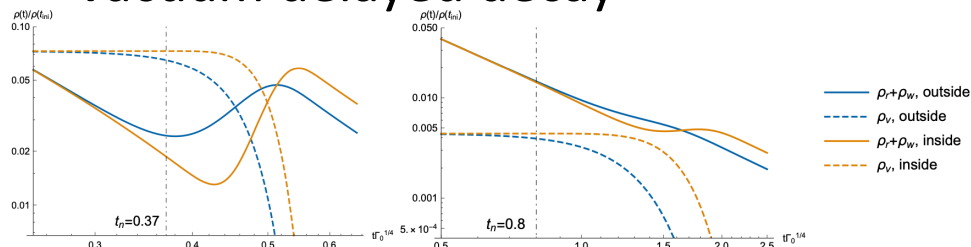
- Particle trapping

Kawana and KPX, PLB 824 (2022) 136791



Baker *et al*, 2105.07481

- Vacuum delayed decay



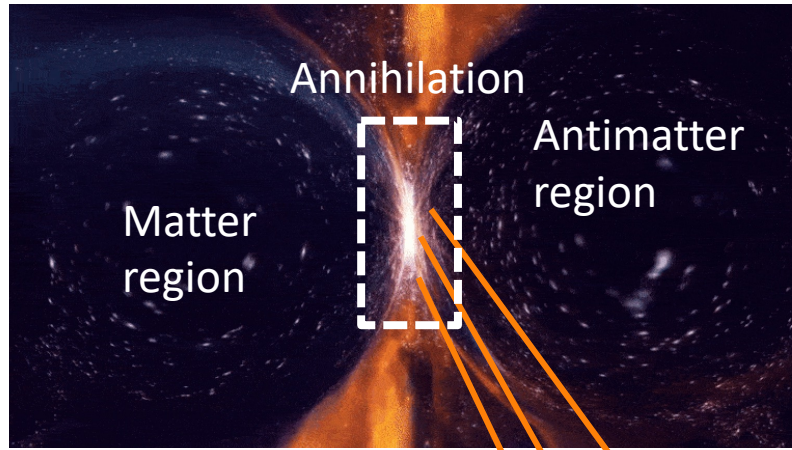
Liu, Bian, Cai, Guo and Wang PRD 105, L021303

The matter-antimatter asymmetry

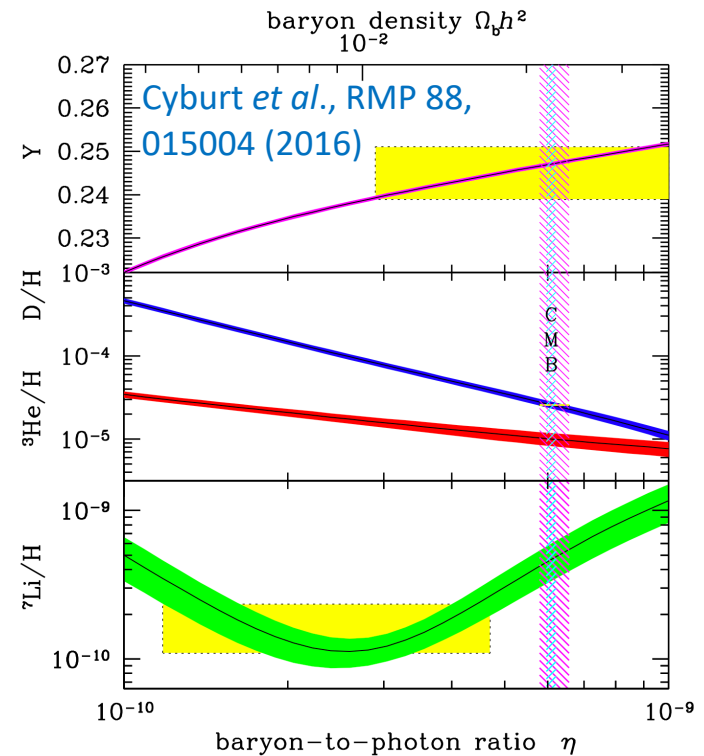
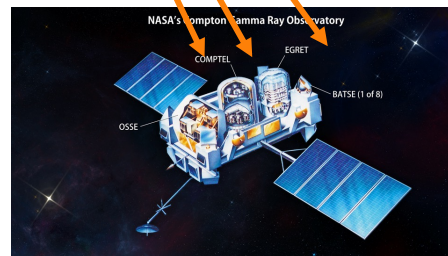
Also known as the *baryon asymmetry*

$$Y_B = \frac{n_B - n_{\bar{B}}}{s} \approx 10^{-10}$$

Two evidences:



Cohen *et al*,
Astrophys. J. 495, 539
(1998)



Generating the matter-antimatter asymmetry

Baryogenesis between inflationary reheating & BBN

Three conditions for baryogenesis: [Sakharov,1967]

1. Baryon number violation;
2. C/CP violation;
3. Departure from equilibrium.

SM cannot satisfy



A.Sakharov (1921-1989)

C&CP violation:
To avoid

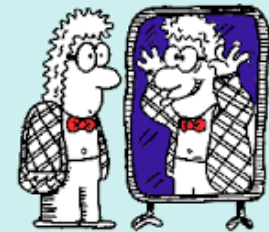
$$N(\text{matter}) = N(\text{antimatter})$$

$$\begin{aligned} N(\text{left-handed matter}) &= N(\text{right-handed antimatter}) \\ N(\text{right-handed matter}) &= N(\text{left-handed antimatter}) \end{aligned}$$

Why departure from equilibrium:

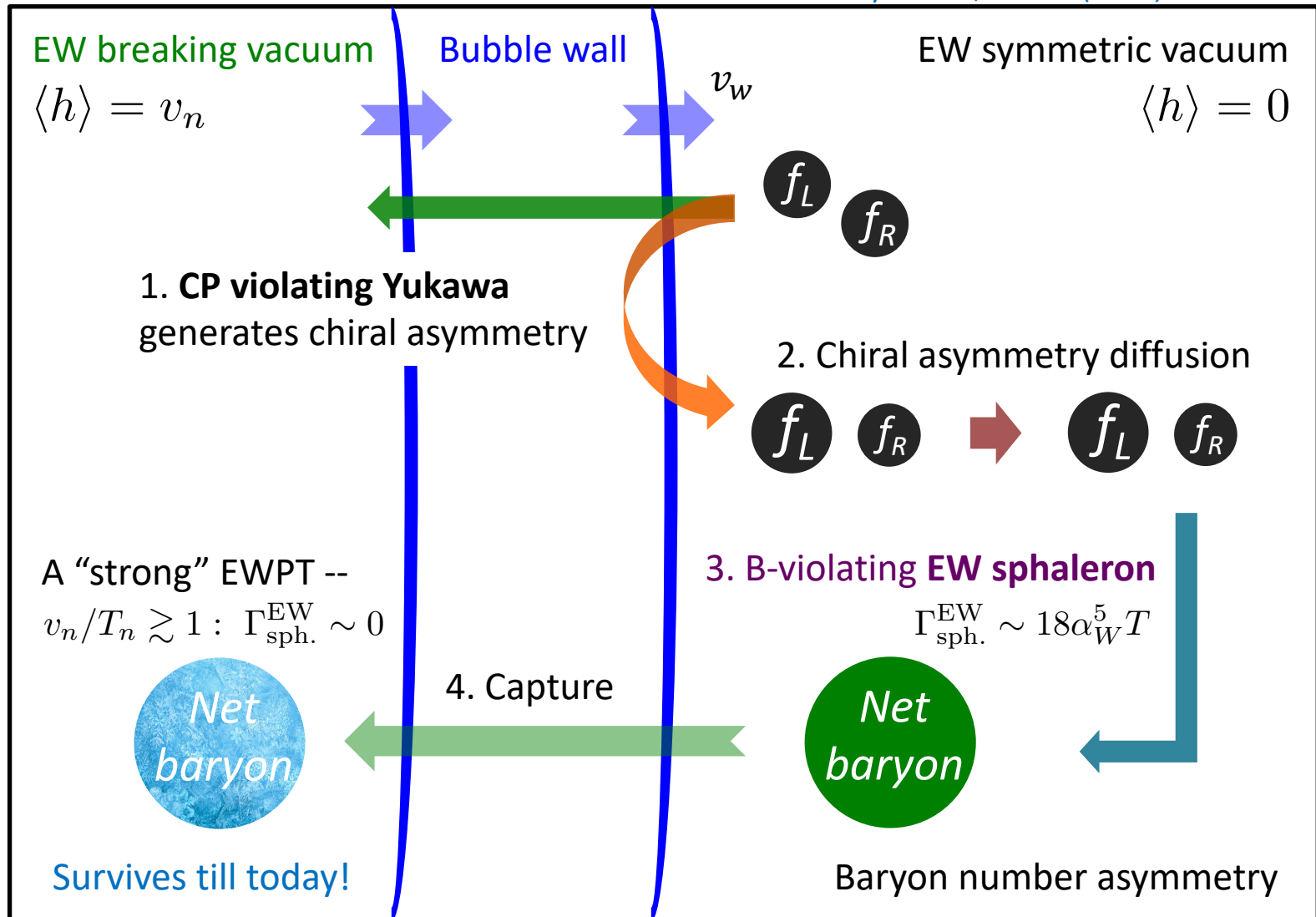
$$\begin{aligned} \langle B \rangle &= \text{tr}[B e^{-H/T}] \\ &= \text{tr}[(CPT)^{-1} B e^{-H/T} (CPT)] = -\text{tr}[B e^{-H/T}] \\ &= -\langle B \rangle = 0 \end{aligned}$$

CP violation



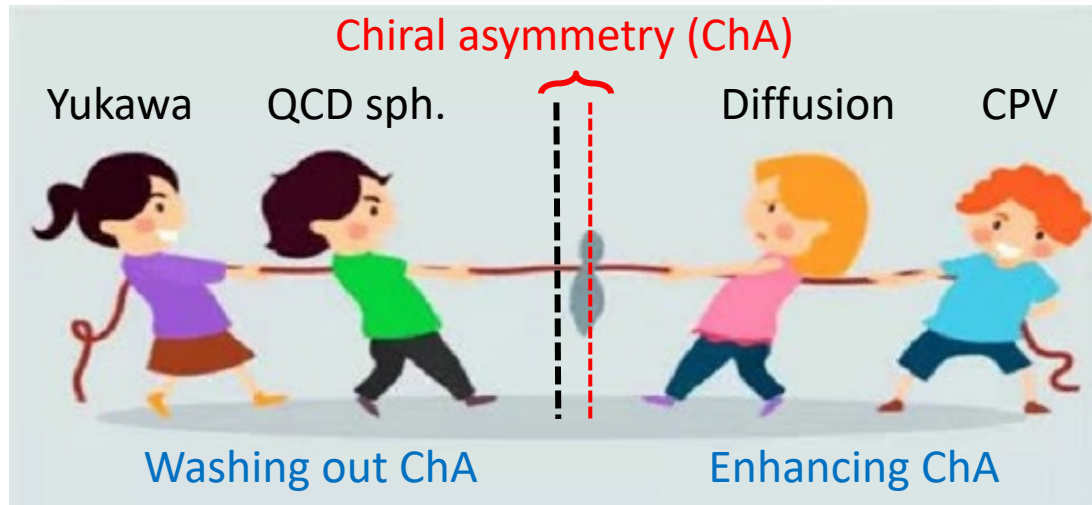
Electroweak baryogenesis

Joyce *et al*, PRL 75 (1995) 1695–1698



[Review] Morrissey *et al*, New J.Phys. 14 (2012) 125003

Which fermion?



Most popular choice

$f = t$ $y_t \approx 1$ Large $D_q \approx 6/T$ $y_t \approx 1$
 Strong washout v.s. Strong source

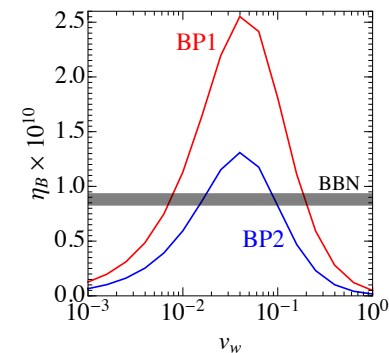
$f = \tau$ $y_\tau \approx 0.01$ Absent $D_\ell \approx 100/T$ $y_\tau \approx 0.01$
 Small washout v.s. Small source

τ -mediated [Chung et al, PRL 102 \(2009\) 061301](#); [Guo, Li, Liu, Ramsey-Musolf and Shu, PRD 96 \(2017\) 115034](#); [De Vries et al, JHEP 04 \(2019\) 024](#); [KPX, JHEP 02 \(2021\) 090](#)

b -mediated [\[Modak et al, PRD 99 \(2019\) 115022\]](#)

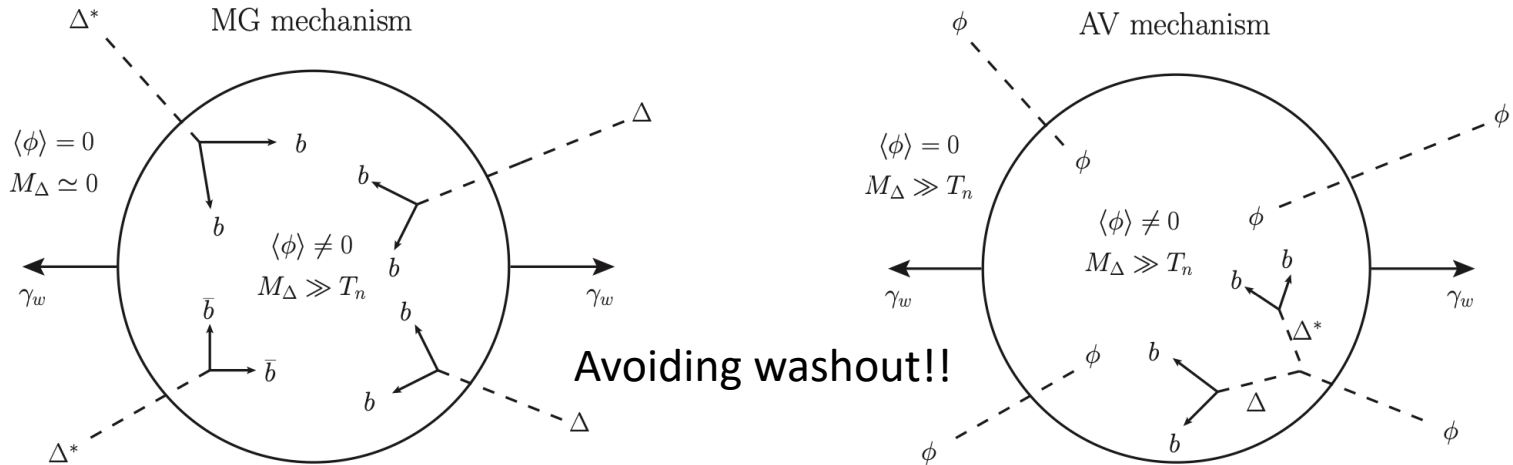
μ -mediated [\[Fuchs et al, PRL 124 \(2020\) 181801\]](#)

ν -mediated [\[Fernandez-Martinez et al, JHEP 10 \(2020\) 063\]](#)



Baryogenesis triggered by a FOPT

Relativistic bubbles [Baldes *et al*, PRD 104 (2021) 115029; Azatov *et al*, JHEP 10 (2021) 043]

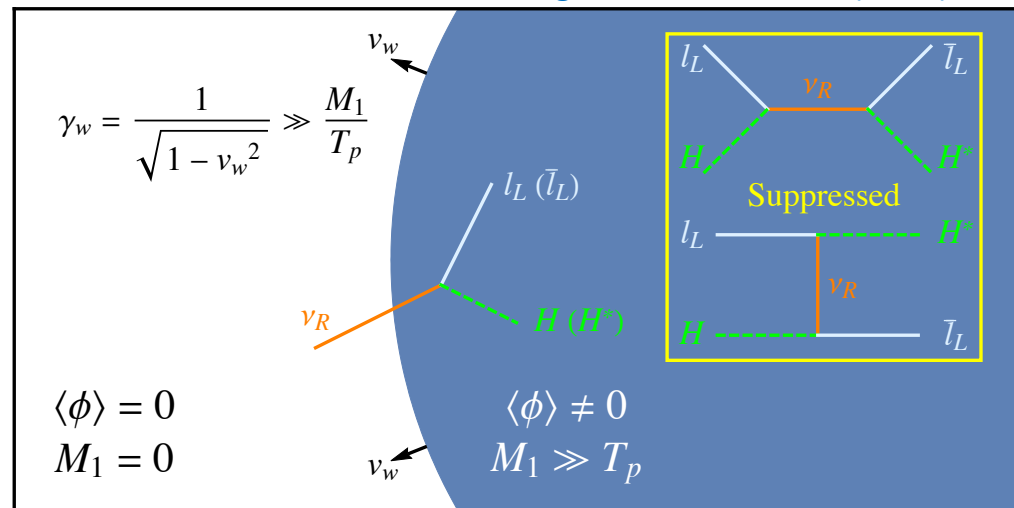


Huang and KPX, JHEP 09 (2022) 052

Apply to leptogenesis

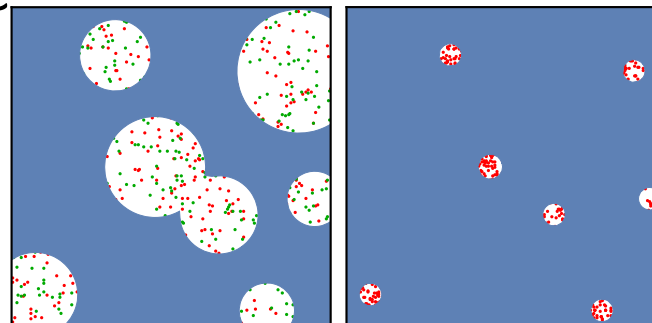
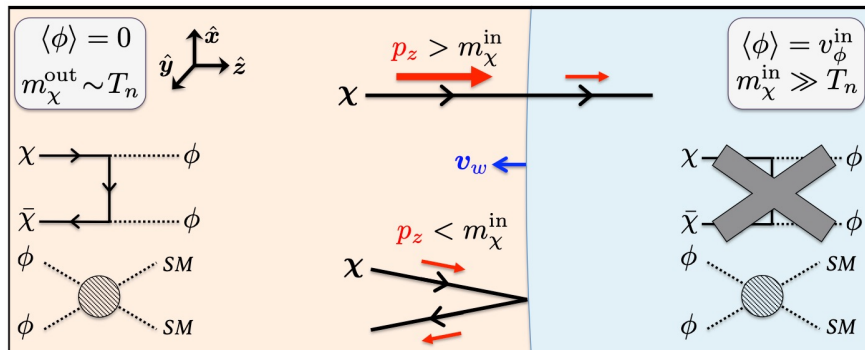
- $\nu_R \rightarrow \ell_L H$
- $\epsilon_1 = \frac{\Gamma_{\ell H} - \Gamma_{\bar{\ell} H^*}}{\Gamma_{\ell H} + \Gamma_{\bar{\ell} H^*}}$
- $Y_B \approx \epsilon_1 \frac{n_{\nu_R}}{s}$

No washout at all



A short summary on new physics mechanisms (1)

Slow-moving bubble walls ($v_w < 1$): particle filtering or trapping; diffusion; change of kinetic space, etc

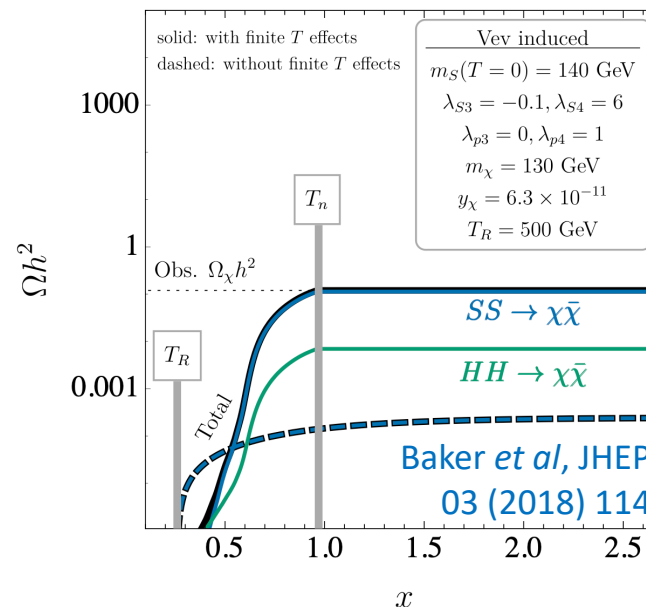
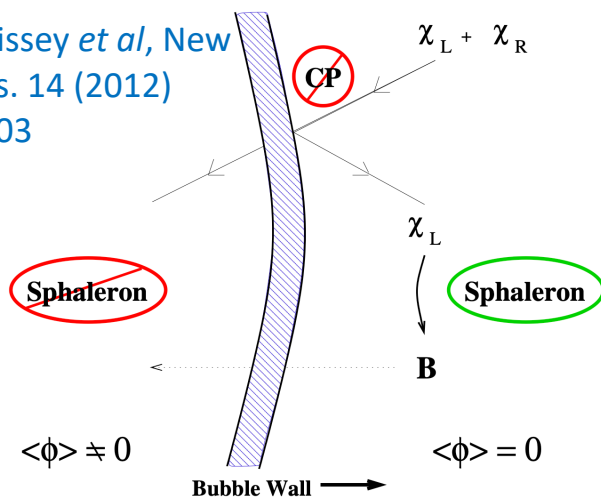


Hong, Jung and KPX, PRD 102 (2020) 7, 075028

Baker *et al*, PRL 125 (2020) 15, 151102;

Chao *et al*, JCAP 06 (2021) 038

Morrissey *et al*, New J.Phys. 14 (2012) 125003



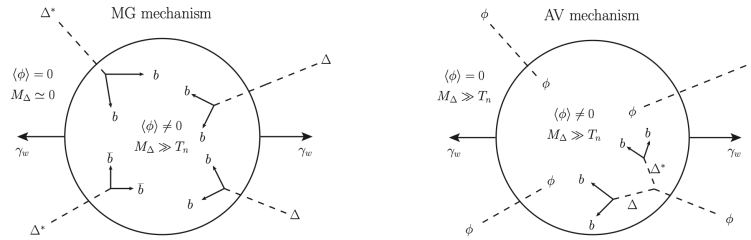
x

Baker *et al*, JHEP 03 (2018) 114

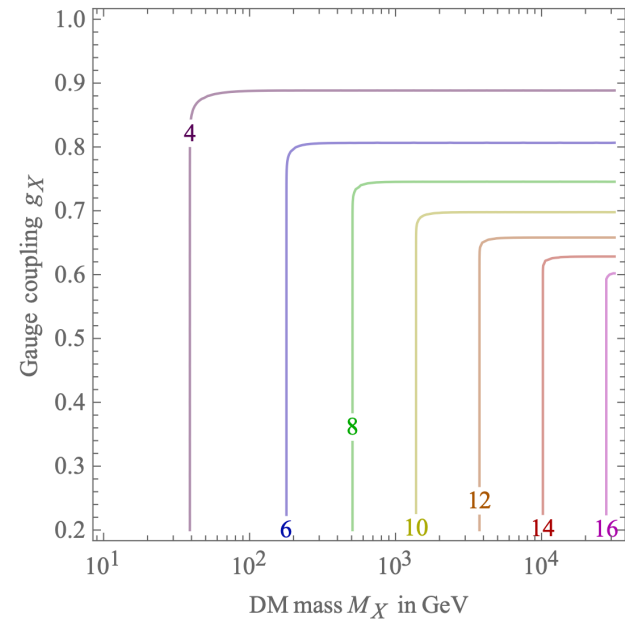
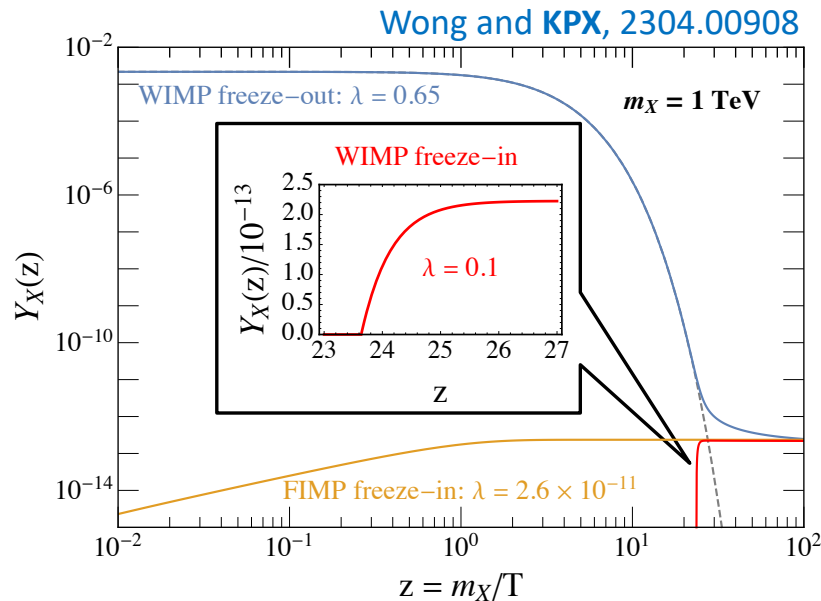
A short summary on new physics mechanisms (2)

Relativistic bubble walls ($v_w \approx 1$): heavy particle production; entropy injection, etc

Baldes *et al*, PRD 104 (2021) 115029
 Azatov *et al*, JHEP 10 (2021) 043
 Huang and **KPX**, JHEP 09 (2022) 052



Hambye *et al*, JHEP 08 (2018) 188;
 Baldes *et al*, JHEP 07, 084 (2022) Number of e^- -folds N



Models that can have FOPTs

SM + singlet scalar (xSM or cxSM)

Cline *et al*, JCAP 01 (2013) 012; Alanne *et al*, NPB 889 (2014) 692; Chiang *et al*, PLB 789 (2019) 154; Jiang, Bian, Huang and Shu, PRD 93 (2016) 6, 065032; Alves *et al*, JHEP 04 (2019) 052, JHEP 12 (2018) 070, JHEP 03 (2020) 053, PLB 818 (2021) 136377; Carena *et al*, JHEP 08 (2020) 107; Liu and **KPX**, JHEP 04 (2021) 015; Huang and **KPX**, PRD 105 (2022) 11, 115033, Liu *et al*, PRD 105 (2022) 11, 115040; etc

Two-Higgs-doublet model

Cline *et al*, JHEP 11 (2011) 089; Dorsch *et al*, JHEP 10 (2013) 029; Basler *et al*, JHEP 02 (2017) 121; Dorsch *et al*, JHEP 12 (2017) 086; Bian, Jiang *et al*, JHEP 05 (2018) 151; Wang, Yang, Zhang and Zhang, PLB 788 (2019) 519; Wang, Huang, Zhang, PRD 101 (2020) 015015; Su *et al*, JHEP 04 (2021) 219; etc

Left-right model

Brdar, Graf and Xu, JCAP 12 (2019) 027; Li, Yan, Zhang, Zhao, JHEP 03 (2021) 267; etc

Georgi-Machacek model

Zhou, Cheng, Deng, Bian and Wu, JHEP 01 (2019) 216; etc

Supersymmetric model

Lee *et al*, PRD 71 (2005) 075010; Balazs *et al*, PRD 71 (2005) 075002; Huang, Kang, Shu, Wu and Yang, PRD 91 (2015) 2, 025006; Bi, Bian, Huang, Shu and Yin, PRD 92 (2015) 023507; Bian, Guo, Shu, CPC 42 (2018) 9, 093106; Athron *et al*, JHEP 11 (2019) 151; Wang, **KPX**, Wu and Yang, EPJC 82 (2022) 12, 1120; etc

Composite Higgs model

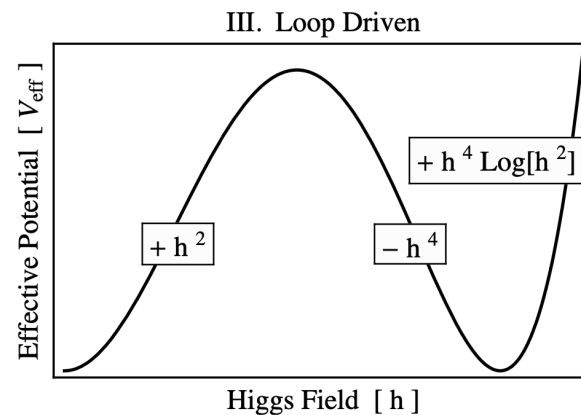
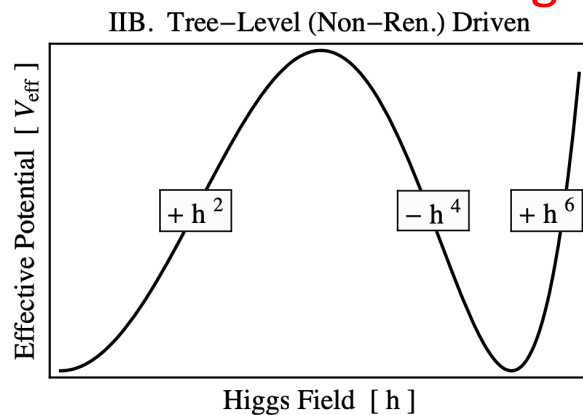
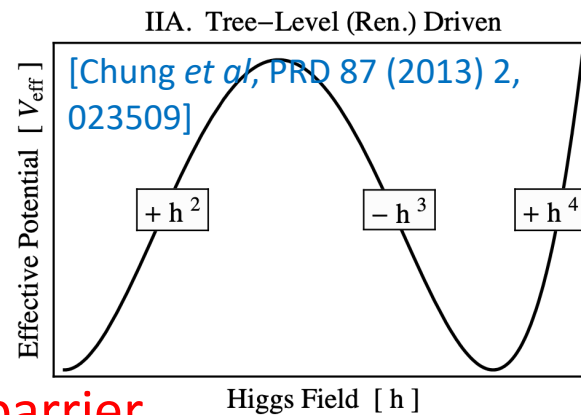
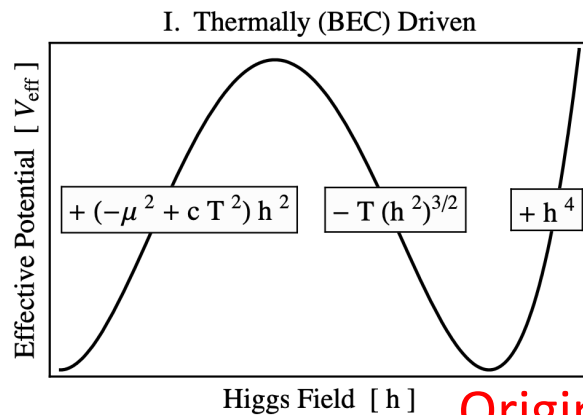
Espinosa *et al*, JCAP 01 (2012) 012; Bian, Wu and **KPX**, JHEP 12 (2019) 028, JHEP 12 (2020) 047; De Curtis *et al*, JHEP 12 (2019) 149; Angelescu *et al*, JHEP 10 (2022) 019; etc

Calculating the FOPTs (1): potential

The finite temperature potential [\[Quiros, hep-ph/9901312\]](#)

$$U_T(\phi, T) = U_0(\phi) + U_1(\phi) + U_{1,T}(\phi, T) + U_{\text{daisy}}(\phi, T)$$

↑ Tree level
 ↑ 1-loop CW
 ↑ 1-loop thermal
 ↑ Daisy resummation



Origin of barrier

IIA & III:
Hot now

Calculating the FOPTs (2): dynamics

Vacuum decay rate [\[Coleman et al, PRD 15 \(1977\) 2929-2936; PRD 16 \(1977\) 1762-1768\]](#)

At finite temperature [\[Linde, NPB 216 \(1983\) 421\]](#) $\Gamma(T) \sim \left(\frac{S_3}{2\pi T}\right) T^4 e^{-S_3/T}$

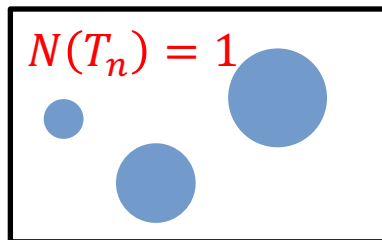
S_3 -- Action of $O(3)$ -symmetric bounce solution ↖ Bounce solution

$$S_3 = \int_0^\infty 4\pi r^2 dr \left[\frac{1}{2} \left(\frac{d\hat{\phi}}{dr} \right)^2 + U_T(\hat{\phi}, T) \right]$$

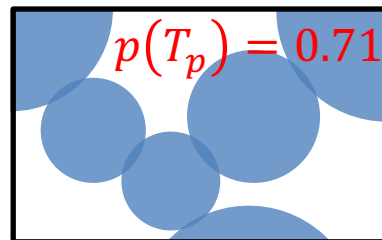
Nucleation rate $N(T) = \int_T^{T_c} \frac{dT'}{T'} \frac{\Gamma(T')}{H^4(T')}$

False vacuum fraction [\[Guth et al, PRD23 \(1981\) 876\]](#) $p(T) = e^{-I(T)}$;

$$I(T) = \frac{4\pi}{3} \int_T^{T_c} dT' \frac{\Gamma(T')}{T'^4 H(T')} \left(\int_T^{T'} \frac{v_w d\tilde{T}}{H(\tilde{T})} \right)^3$$



1. Nucleation T_n



2. Percolation T_p



3. Completed

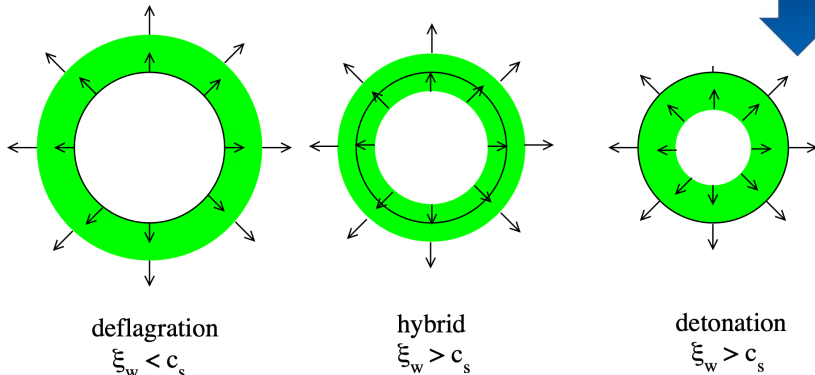
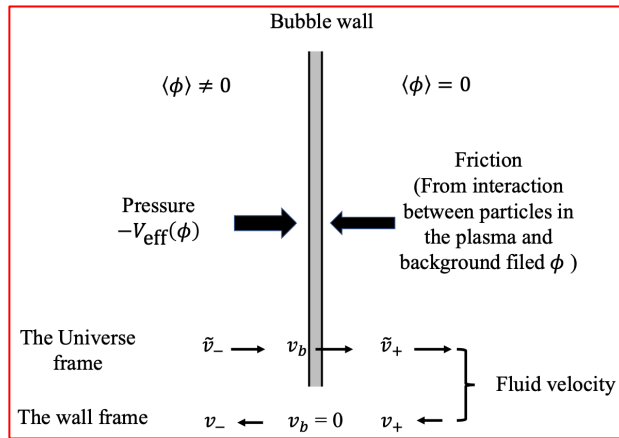
FOPT criterion

Very rough: $\frac{S_3}{T} \sim 4 \ln \frac{M_{Pl}}{T}$

Strict: $3 + T_p \left. \frac{dI}{dT} \right|_{T_p} < 0$

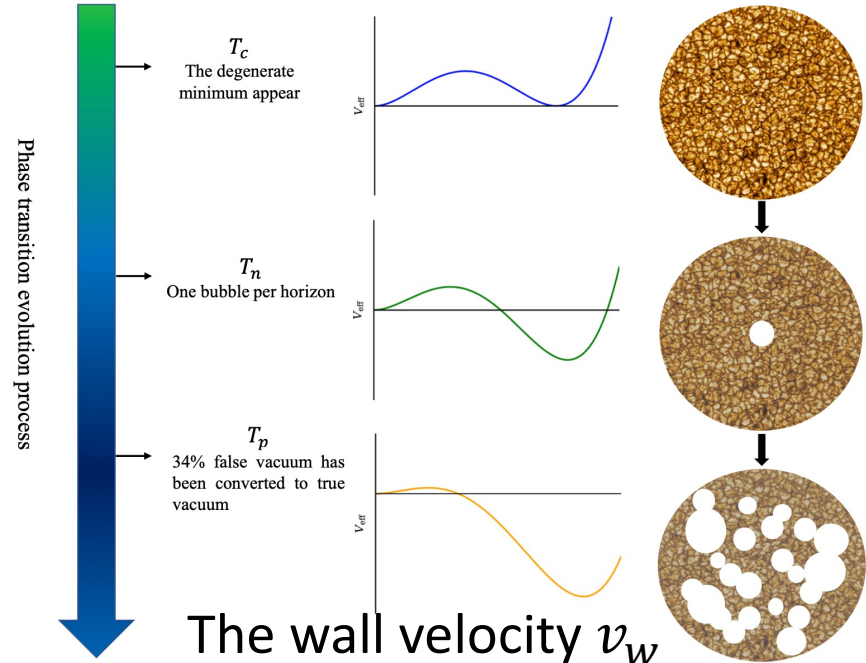
Dynamics of the FOPT

Energy budget: how much energy is devoted into wall or plasma motion?



Espinosa *et al*, JCAP 1006 (2010) 028;
Ellis *et al*, JCAP 06 (2019) 024, JCAP 11 (2020) 020

Wang, Huang and Zhang, JCAP 05 (2020) 045



Konstandin *et al*, JCAP 09 (2014) 028;
Bodeker *et al*, JCAP 0905 (2009) 009, JCAP 05 (2017) 025; Höche *et al*, JCAP 03 (2021) 009; Cai and Wang, JCAP 03 (2021) 096; Lewicki *et al*, JHEP 02 (2022) 017; Ai *et al*, JCAP 03 (2022) 03, 015; Gouttenoire (DESY and Tel Aviv U.), Ryusuke Jinno *et al*, JHEP 05 (2022) 004; Wang *et al*, PRD 107 (2023) 2, 023501

Gravitational waves & lattice simulation

Sources: [Caprini *et al*, JCAP 1604 (2016) 001]

1. Bubble collision;
2. Sound waves;
3. Turbulence

Normally, sound wave dominates

Espinosa *et al*, JCAP 1006 (2010) 028; Ellis *et al*, JCAP 04 (2019) 003; Wang, Huang and Zhang, JCAP 05 (2020) 045; Guo *et al*, JCAP 01 (2021) 001; etc

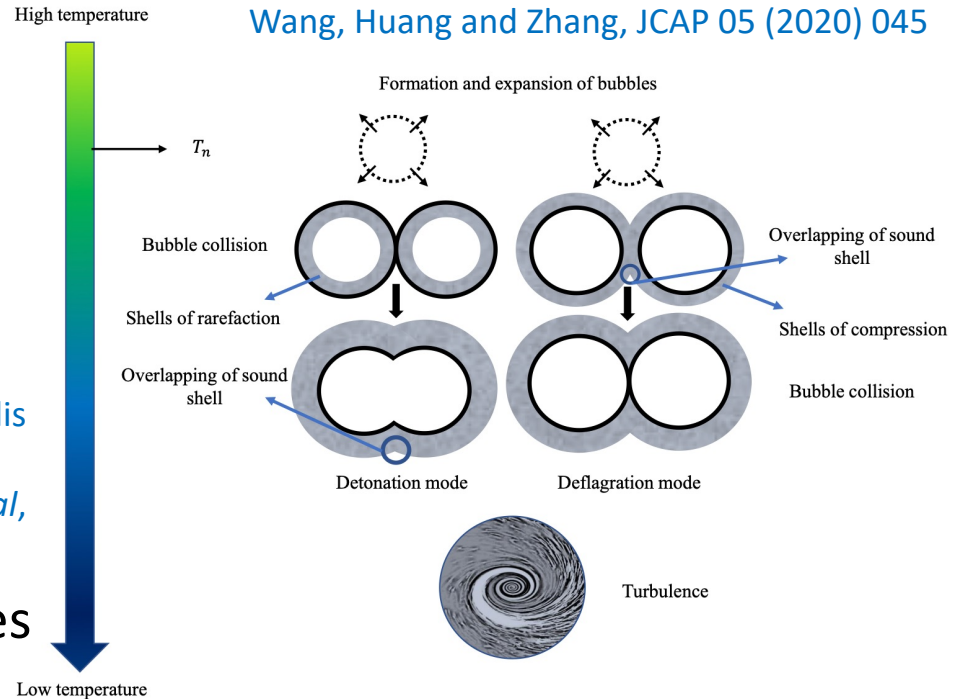
Bubble collision dominates for relativistic bubbles

Ellis *et al*, JCAP 06 (2019) 024, JCAP 11 (2020) 020; etc

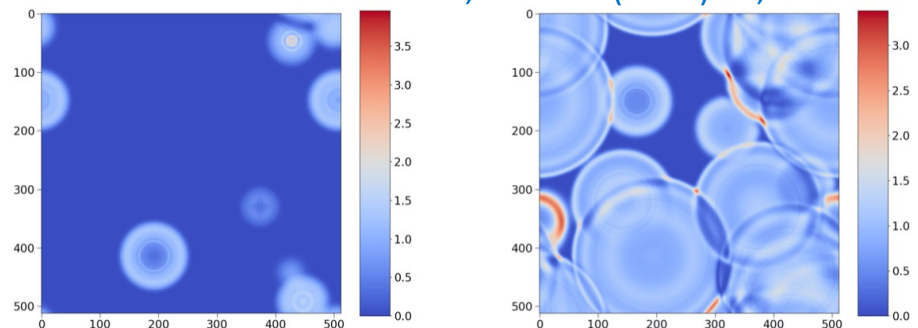
Lattice simulations

Di, Wang, Zhou, Bian, Cai and Liu, PRL 126 (2021) 25, 251102; Zhao, Di, Bian and Cai, 2204.04427; Li, Bian and Jia, 2304.05220; etc

Wang, Huang and Zhang, JCAP 05 (2020) 045

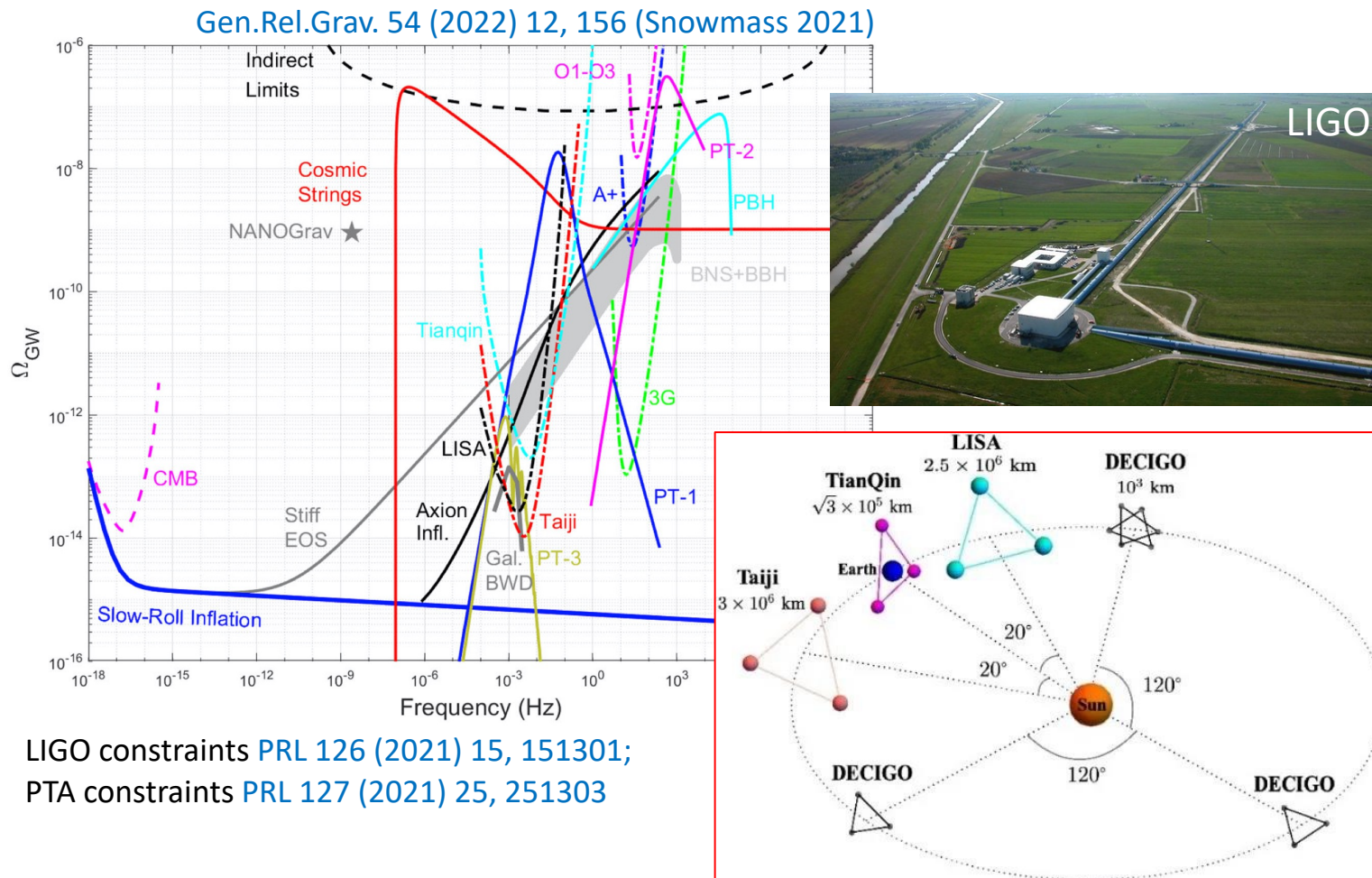


Di *et al*, PRL 126 (2021) 25, 251102



Detecting FOPTs (1)

Current and future GW detectors

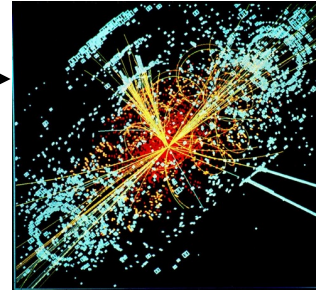
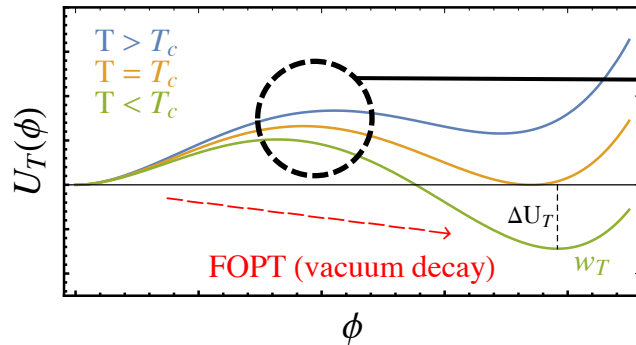


LIGO constraints PRL 126 (2021) 15, 151301;

PTA constraints PRL 127 (2021) 25, 251303

Detecting FOPTs (2)

Collider experiments complementarity



LHC → HL-LHC
HE-LHC? LHeC?
CEPC? ILC? FCC?
Muon collider?

Di-Higgs production

Huang, Gu, Yu and Zhang, PRD 93 (2016) 10, 103515; Alves, Guo, *et al*, JHEP 04 (2019) 052, JHEP 12 (2018) 070, JHEP 03 (2020) 053, PLB 818 (2021) 136377; Liu and **KPX**, JHEP 04 (2021) 015; etc

Exotic Higgs decay

Kozaczuk *et al*, PRD 101 (2020) no. 11, 115035; Carena, **KPX et al**, 2203.08206 (Snowmass 2021); Liu *et al*, PRD 105 (2022) 11, 115040; etc

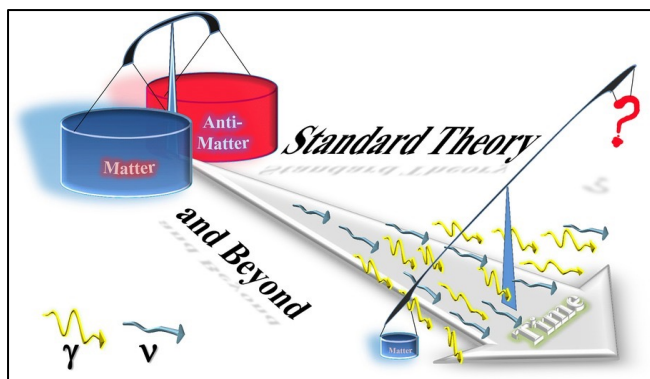
Higgs triple coupling, EW precision measurements

Huang, Gu, Yu and Zhang, PRD 93 (2016) 10, 103515; Cao, Huang, **KPX** and Zhang, CPC 42 (2018) 2, 023103; Su, Williams and Zhang, JHEP 04 (2021) 219; Song, Su and Zhang, JHEP 10 (2022) 048; etc

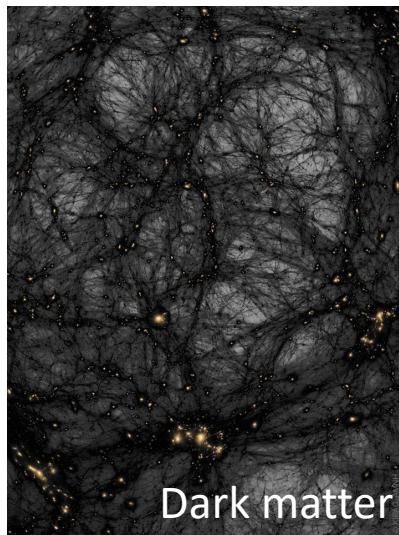
And other signals (very model-dependent)

Closing remarks

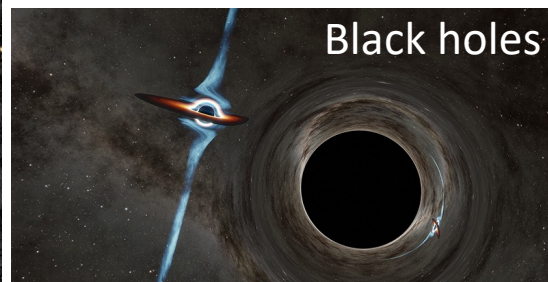
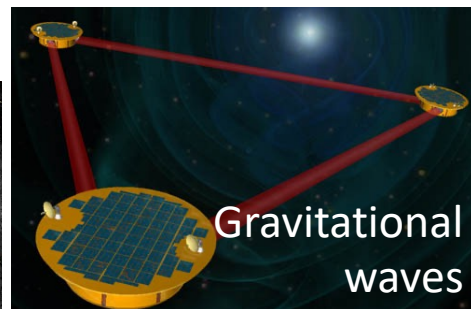
A lot of fun in a boiling Universe!



Matter-antimatter asymmetry



Dark matter



Thank you!