



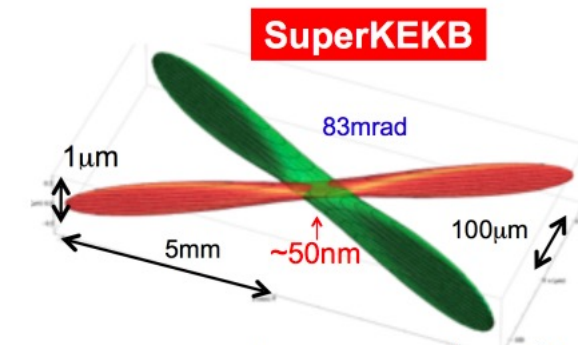
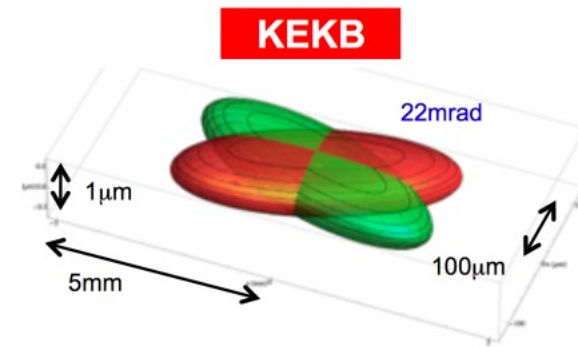
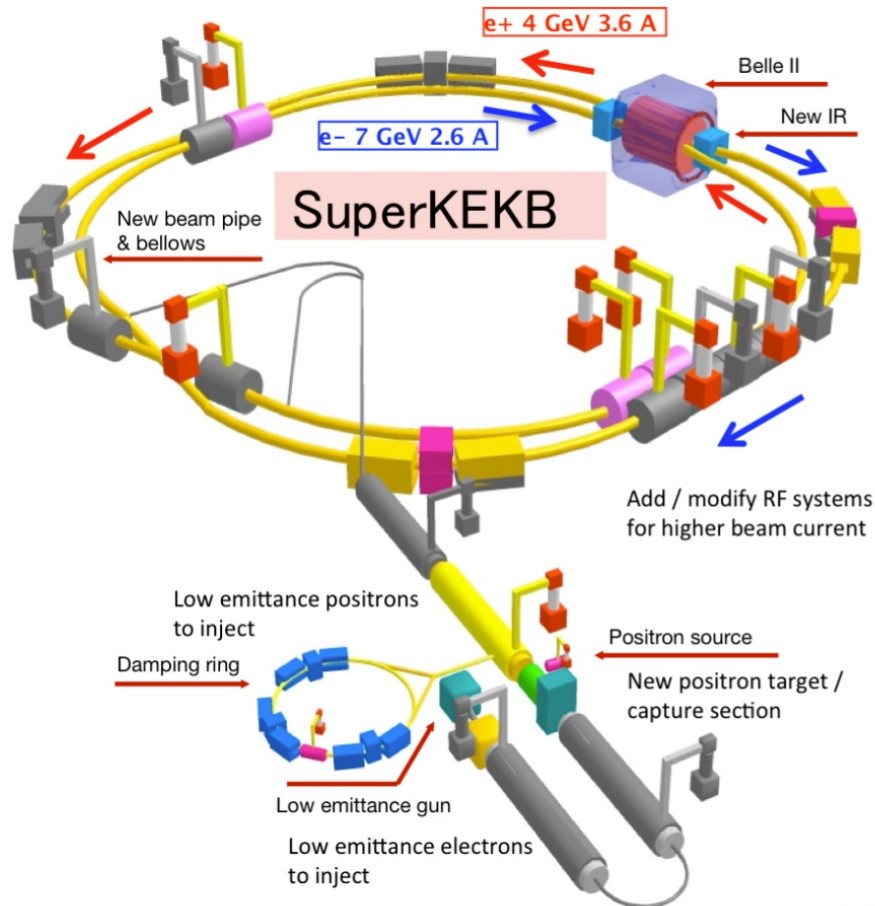
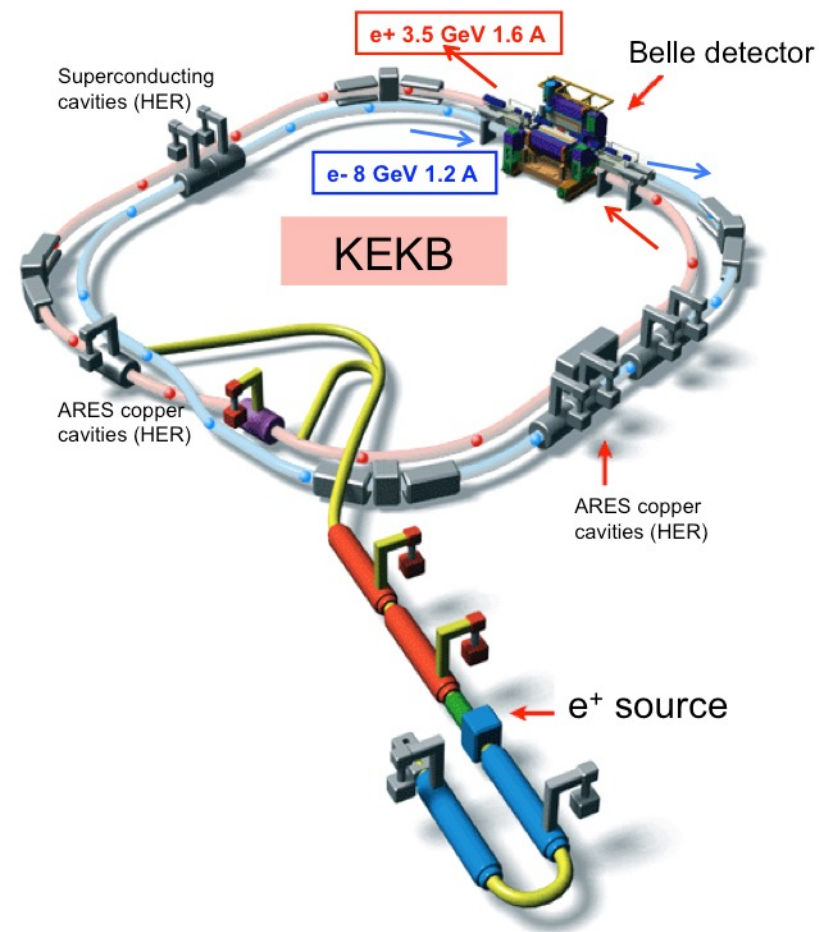
Quarkonium(-like) states at Belle and Belle II

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on behalf of the Belle and Belle II Collaboration

第八届XYZ粒子研讨会
2023年7月25 – 30 日, 长春

From KEKB to SuperKEKB



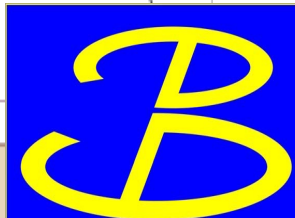
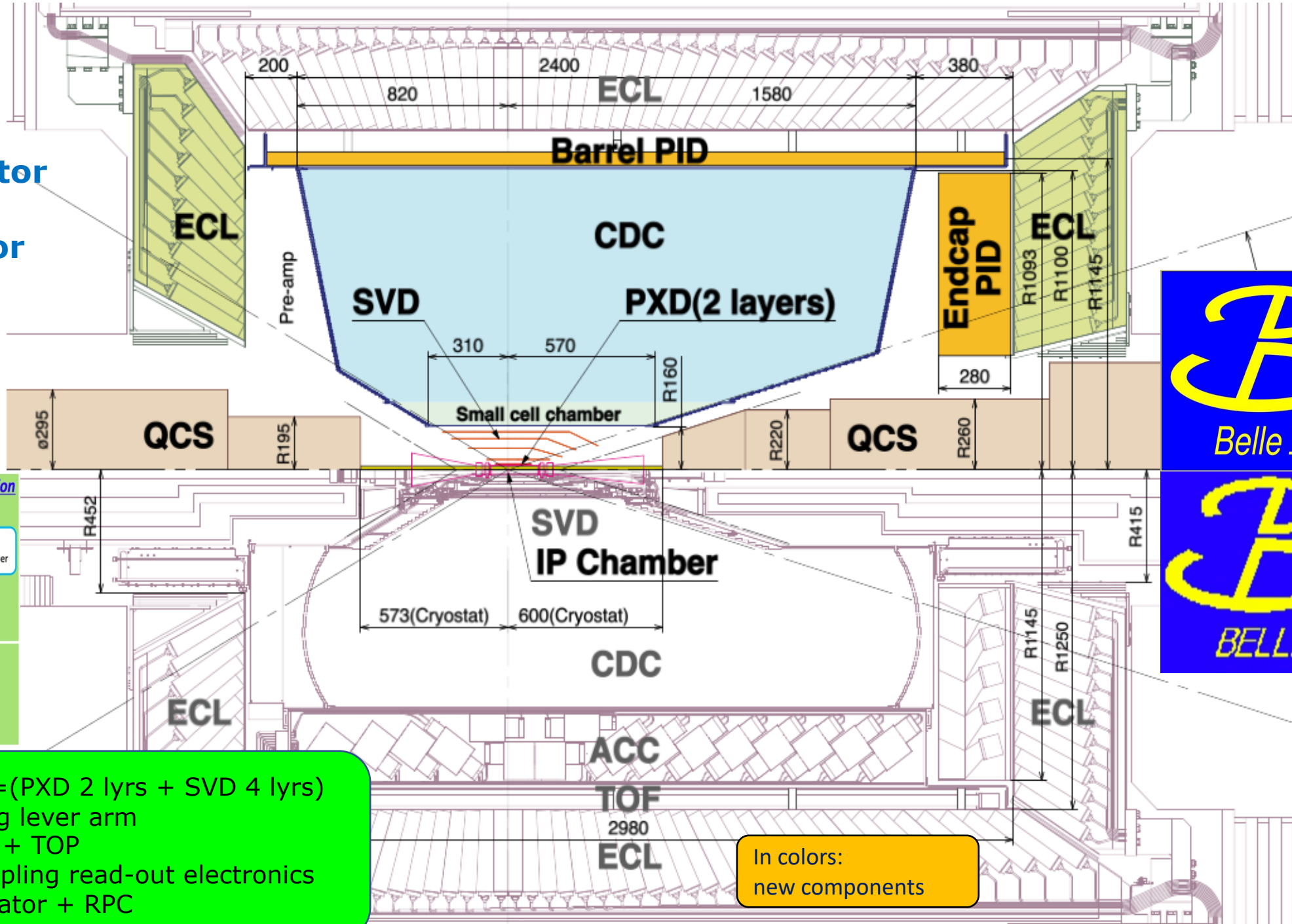
$$\mathcal{L} = \frac{\gamma_{\pm}}{2er_e} \left(1 + \frac{\sigma_y^*}{\sigma_x^*} \right) \frac{I_{\pm} \bar{\zeta}_{y\pm}}{\beta_{y\pm}^*} \left(\frac{R_L}{R_{\bar{\zeta}_{y\pm}}} \right)$$

- Double beam currents
- Squeeze beams @IP by $\sim 1/20$

$$\mathcal{L}_{II}^{\text{peak}} \approx 30 \times \mathcal{L}_I^{\text{peak}}$$

$$\int^{\text{goal}} \mathcal{L}_{II} dt = 50 \text{ ab}^{-1} \approx 50 \int \mathcal{L}_I dt$$

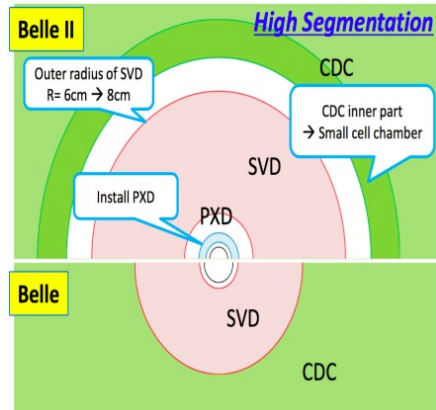
Belle II detector Vs. Belle detector



Belle II



BELLE

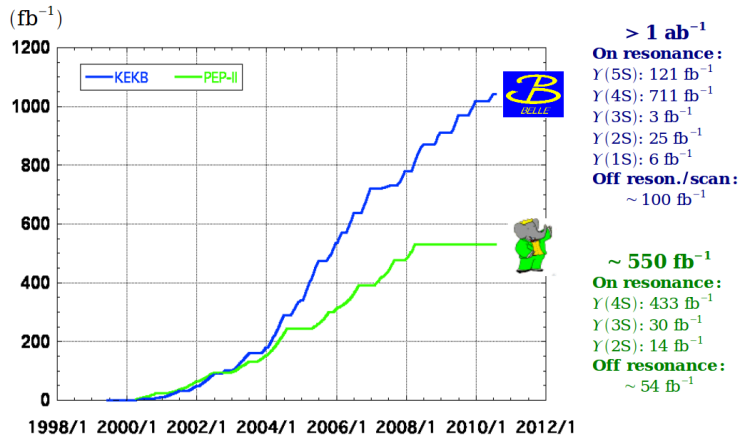


SVD: 4 lyrs → VXD=(PXD 2 lyrs + SVD 4 lyrs)
 CDC: small cell, long lever arm
 ACC+TOF → ARICH + TOP
 ECL: waveform sampling read-out electronics
 KLM: RPC → Scintillator + RPC

In colors:
 new components

Datasets and new quarkonia at Belle

Integrated luminosity of B factories

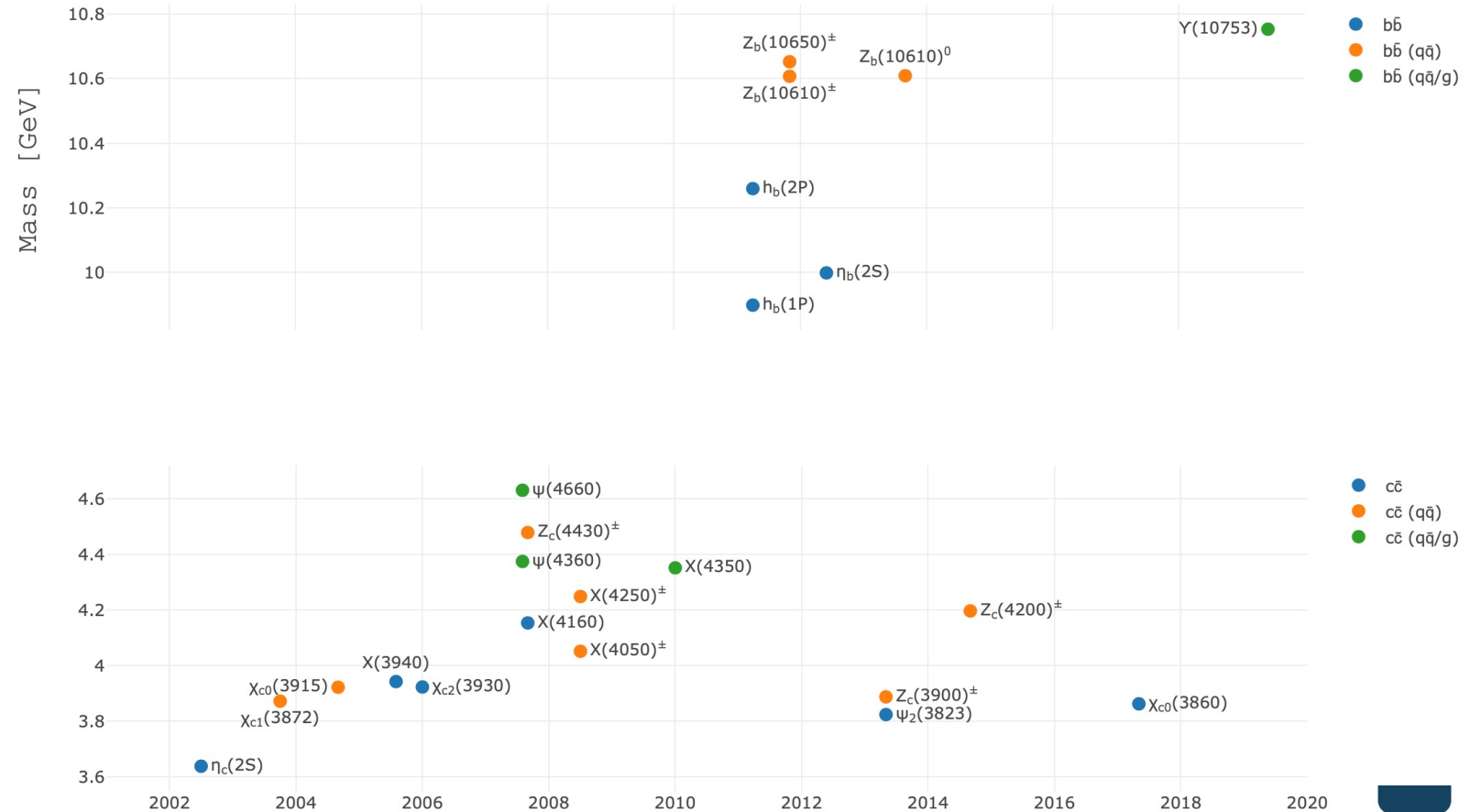


Data taking: 1999 – 2010

On/off/Scan $\Upsilon(nS)$ peaks

772M $B\bar{B}$ events @ $\Upsilon(4S)$

From QWG website:

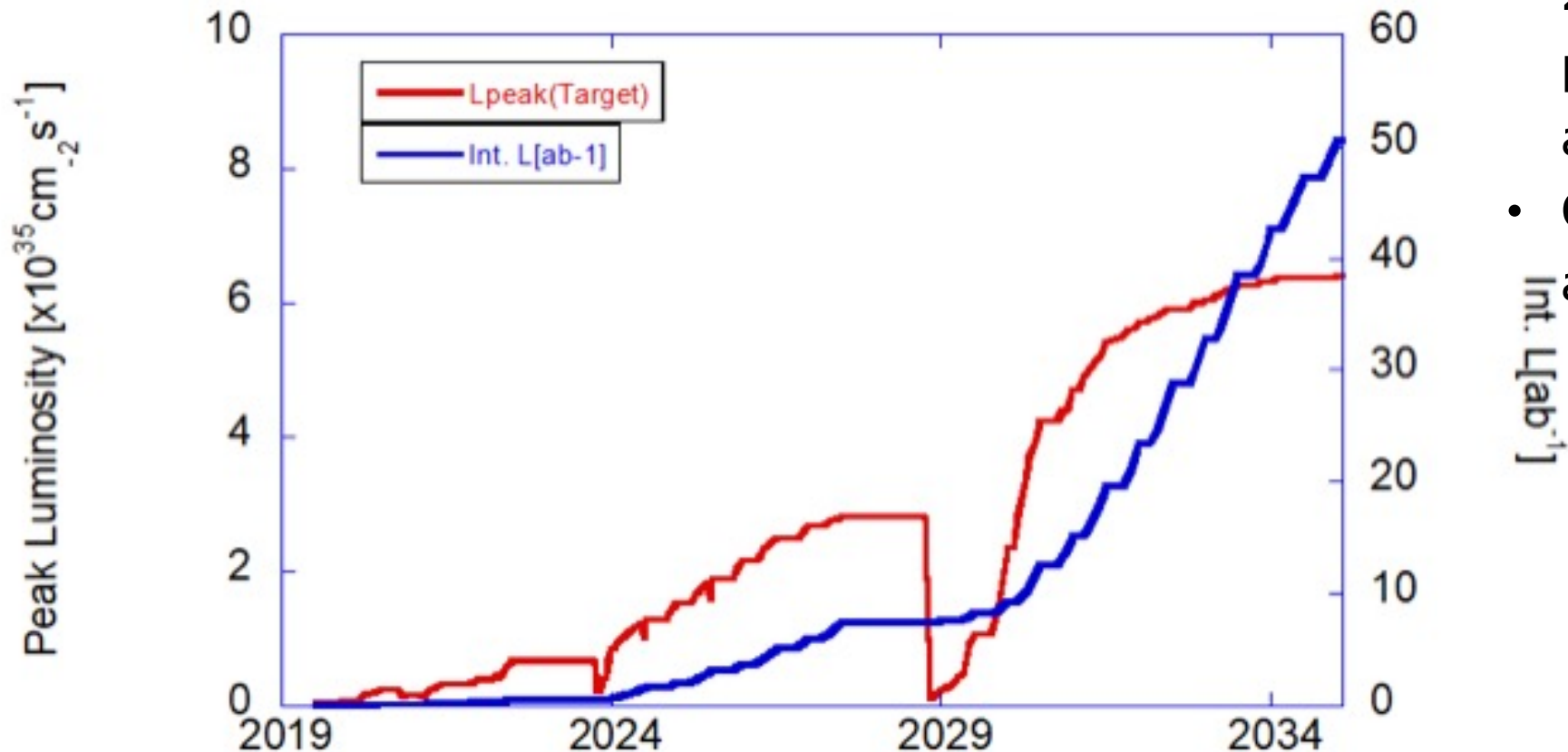


Datasets at Belle II

In June 2022



WORLD RECORD: $4.7 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$



- Collected $\sim 428 \text{ fb}^{-1}$ around $\Upsilon(4S)$ until now
- LS1 starts from summer 2022 to fully install the pixel detector and accelerator machine study
- Operation will be resumed around the end of 2023

Selected topics:

Quarkonium at Belle:

- $X(3872) \rightarrow D^0 \bar{D}^{*0}$ [PRD 107, 112011 (2023)]
- $e^+e^- \rightarrow \eta_c J/\psi$ [arXiv: 2305.17947]

Quarkonium at Belle II:

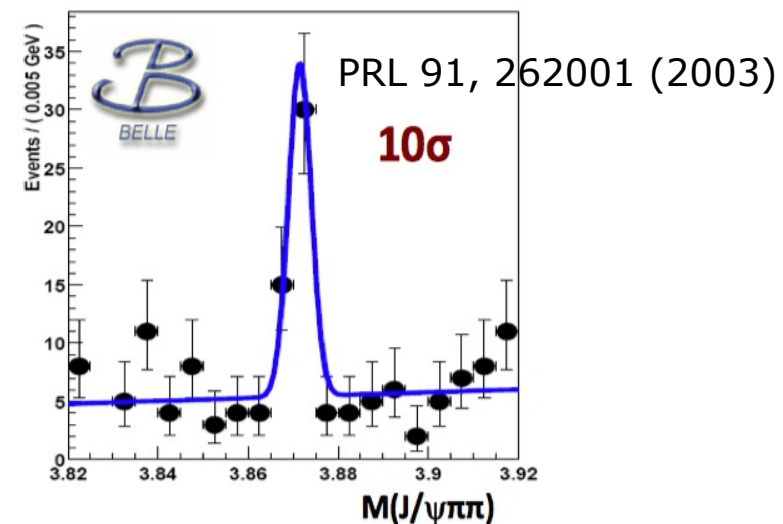
- $e^+e^- \rightarrow \omega \chi_{bJ}$ and $X_b \rightarrow \omega Y(1S)$ [PRL 130, 091902 (2023)]
- $e^+e^- \rightarrow B\bar{B}, B\bar{B}^*$ and $B^*\bar{B}^*$ [Preliminary]
- $e^+e^- \rightarrow \omega \eta_b(1S)$ and $e^+e^- \rightarrow \omega \chi_{b0}(1P)$ [Preliminary]

Quarkonium at Belle

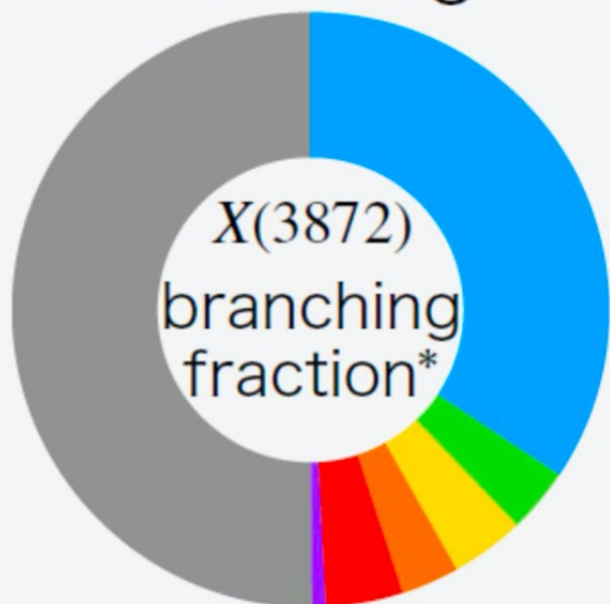
A new measurement of $X(3872) \rightarrow D^0 \bar{D}^{*0}$ at Belle

$X(3872)$ (aka $\chi_{c1}(3872)$) – very famous exotic states

- Narrow width
- Close to DD^* threshold
- No place in charmonium potential model
- $\pi\pi$ from ρ decays thus isospin-violating process



Branching fraction observed to date

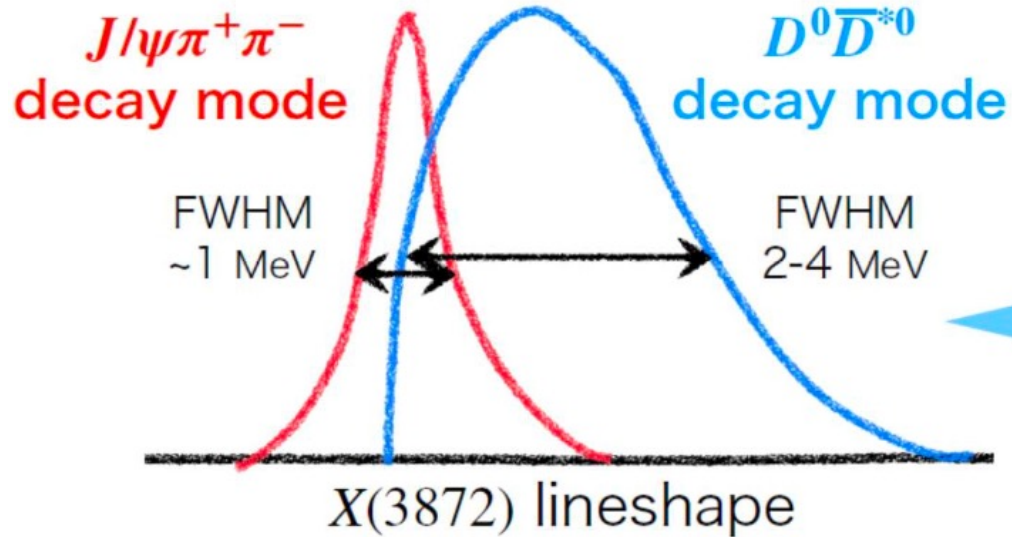


| | |
|---|-------------------|
| ● $D^0 \bar{D}^{*0}$ | $(37 \pm 9)\%$ |
| ● $J/\psi \rho (\rightarrow \pi^+ \pi^-)$ | $(3.8 \pm 1.2)\%$ |
| ● $J/\psi \omega$ | $(4.3 \pm 2.1)\%$ |
| ● $\chi_{c1} \pi^0$ | $(3.4 \pm 1.6)\%$ |
| ● $\psi(2S) \gamma$ | $(4.5 \pm 2.0)\%$ |
| ● $J/\psi \gamma$ | $(0.8 \pm 0.4)\%$ |
| ● Not seen | 46% |

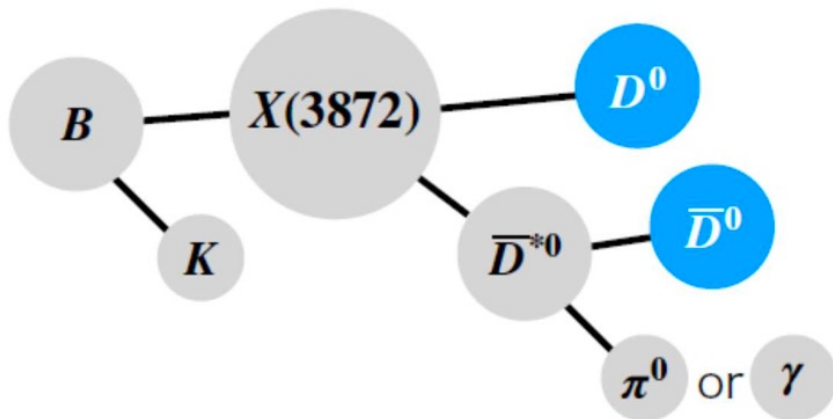
- The $X(3872) \rightarrow D^0 \bar{D}^{*0}$ decay has the largest BR.
- Determining $X(3872) \rightarrow D^0 \bar{D}^{*0}$ coupling strength is important to discuss the structure

Analysis strategy

$X(3872) \rightarrow \pi^+\pi^-J/\psi$ or $X(3872) \rightarrow D^0\bar{D}^{*0}$?



- **Wider lineshape**
because of phase space and threshold effect
- **Better mass resolution**
thanks to small Q-value (~ 100 keV, $\sim 1/20$ of $J/\psi\pi^+\pi^-$)
- **Belle experiment is suitable**
because $D^{*0} \rightarrow D^0\gamma, D^0\pi^0$ can be reconstructed.



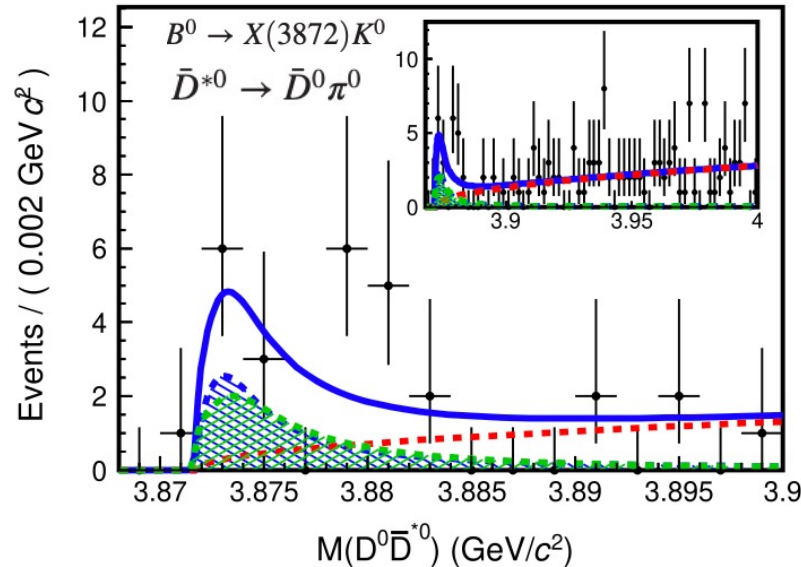
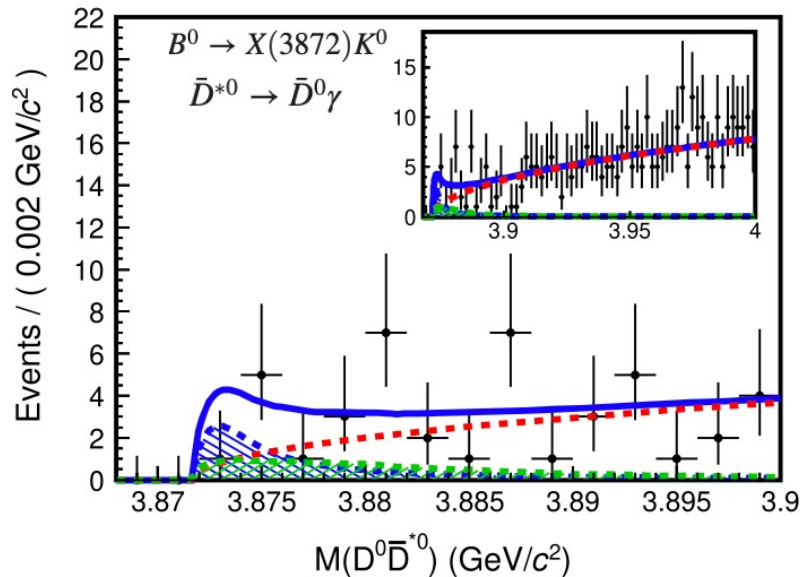
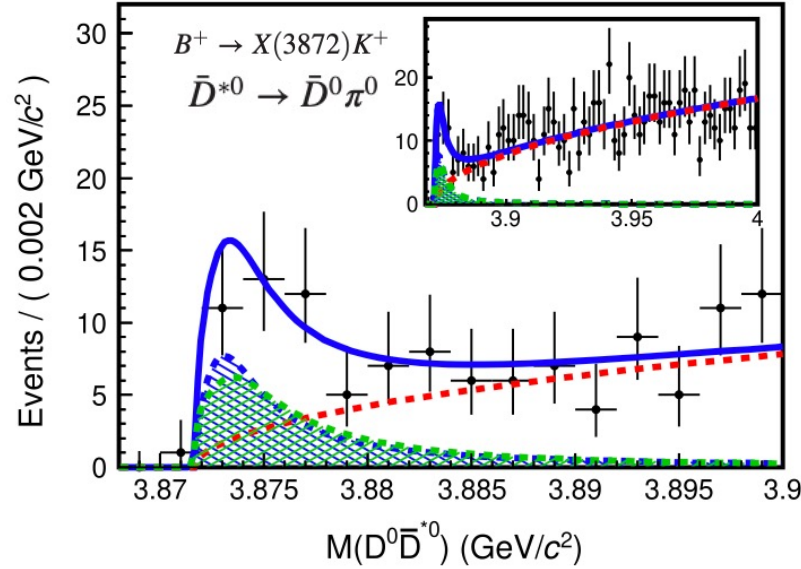
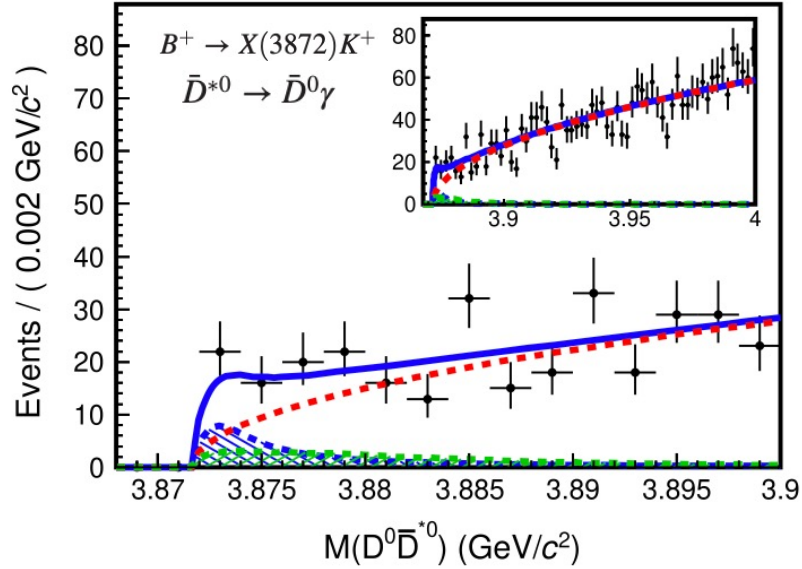
$$D^0 \rightarrow \begin{cases} K^-\pi^+ \\ K^-\pi^+\pi^0 \\ K^-\pi^+\pi^-\pi^+ \\ K_S^0\pi^+\pi^- \\ K^+K^- \\ K_S^0\pi^+\pi^-\pi^0 \end{cases}$$

Signal efficiency is improved by a factor of 1.9 compared to previous Belle measurement [PRD 81, 031103 (2010)].

Fits to $M(D^0\bar{D}^{*0})$ with a BW function

[PRD 107, 112011 (2023)]

▨ Signal
 ▨ Broken signal
 - - - Generic background
 — Total



Broken signal: at least one of the final states is wrongly assigned.

The yield ratio between signal and broken signal is fixed based on MC simulations.

Results:

- Significance: 7.5σ in total
- First observation from the B^0 decay (5.2σ)
- Branching ratios

$$\mathcal{B}(B^+ \rightarrow X(3872)K^+) \times \mathcal{B}(X(3872) \rightarrow D^0\bar{D}^{*0})$$

$$= (0.97_{-0.18}^{+0.21}(\text{stat}) \pm 0.10(\text{syst})) \times 10^{-4},$$

$$\mathcal{B}(B^0 \rightarrow X(3872)K^0) \times \mathcal{B}(X(3872) \rightarrow D^0\bar{D}^{*0})$$

$$= (1.30_{-0.31}^{+0.36}(\text{stat})_{-0.07}^{+0.12}(\text{syst})) \times 10^{-4}.$$

- Mass and width

$$m_{\text{BW}} = 3873.71_{-0.50}^{+0.56}(\text{stat}) \pm 0.13(\text{syst}) \text{ MeV}/c^2$$

$$\Gamma_{\text{BW}} = 5.2_{-1.5}^{+2.2}(\text{stat}) \pm 0.4(\text{syst}) \text{ MeV}$$

- All are consistent with Ref. [PRD 81, 031103 (2010)]

Flatté-like model

$$f(E) = \frac{gk_{D^0\bar{D}^{*0}}}{|E - E_f + \frac{i}{2}[\Gamma_0 + \Gamma_{J/\psi\rho}(E) + \Gamma_{J/\psi\omega}(E) + g(k_{D^0\bar{D}^{*0}} + k_{D^+D^{*-}})]|^2}$$

[PRD 76, 034007 (2007)]

Mass difference from $D^0\bar{D}^{*0}$ threshold

Partial widths for radiative, $J/\psi\rho$, and $J/\psi\omega$ decays

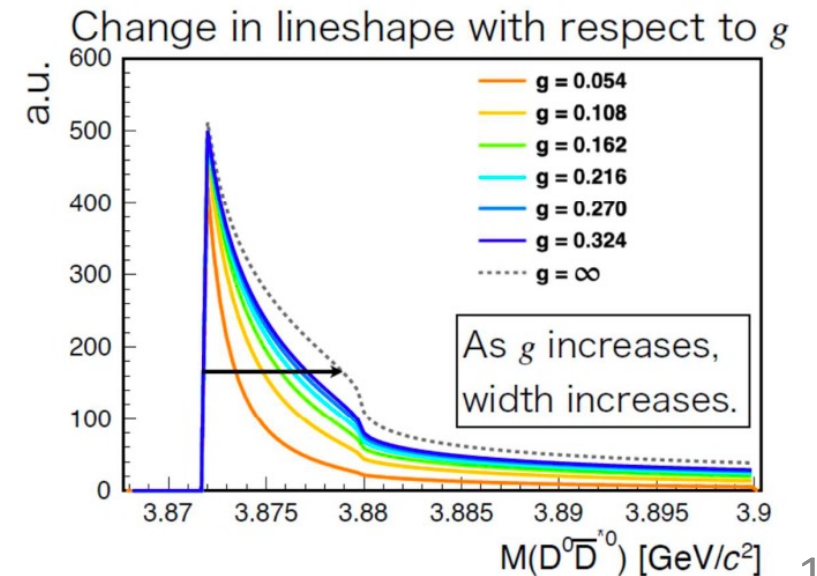
Coupling to $D\bar{D}^*$ channel
 ... g : Coupling constant to $D\bar{D}^*$ channel

... k_a : Momentum for channel a

Fit does not converge w/o constraints due to poor statistics. Thus,

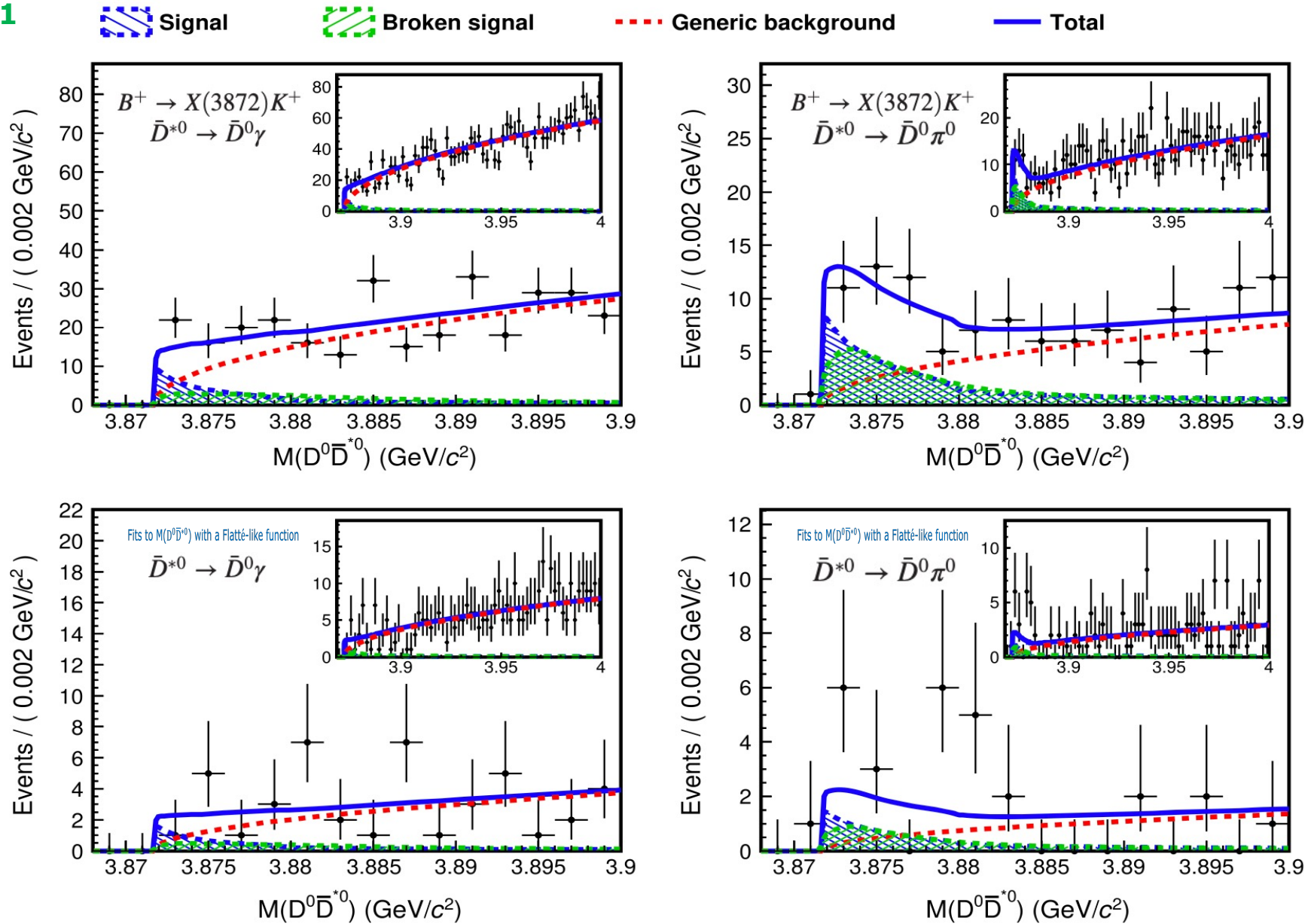
- $\mathcal{T}_{J/\psi\omega}$ is fixed by world-average BR
- E_f , $\mathcal{T}_{J/\psi\omega}$, and $\mathcal{T}_{J/\psi\rho}$ are fixed based on LHCb results [PRD 102, 092005 (2020)]

⇒ Only g is floated



Fits to $M(D^0\bar{D}^{*0})$ with a Flatté-like function

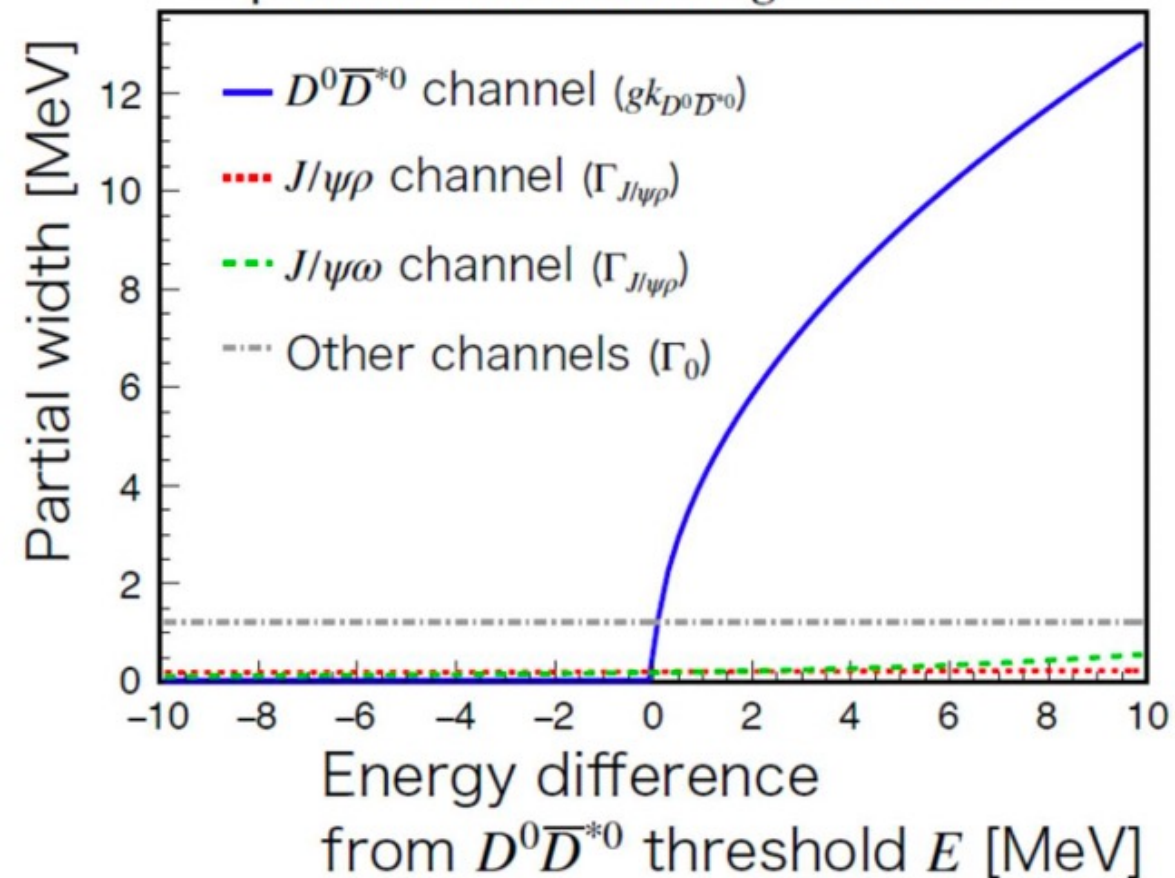
[PRD 107, 112011
(2023)]



Result and discussion

[PRD 107, 112011 (2023)]

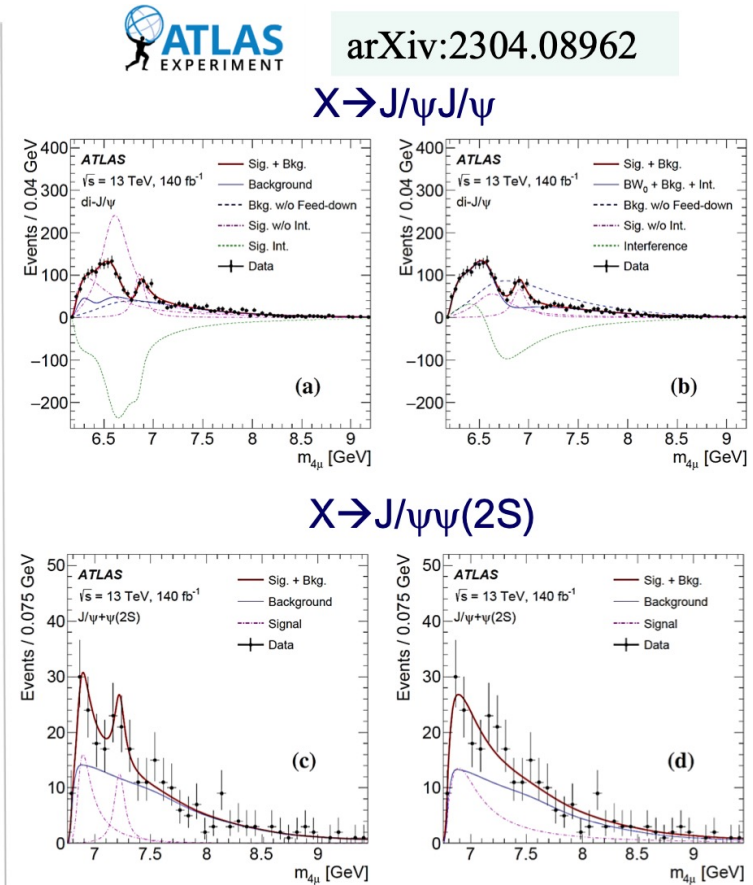
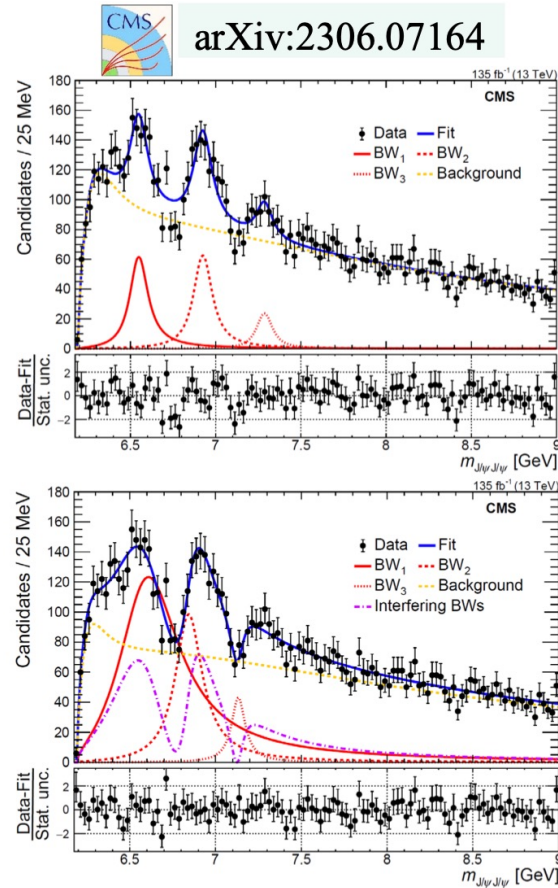
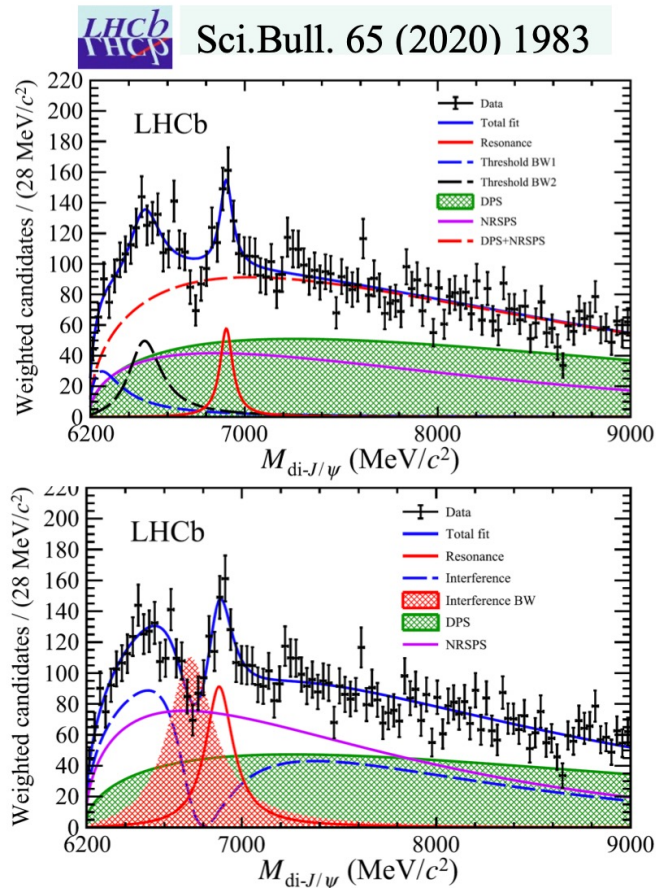
Energy-dependent
partial width for $g = 0.094$



- Fit result: $g = 0.29_{-0.15}^{+2.69}$ (stat. only)
⇒ Lower limit: $g > 0.094$ (90% C.L.)
including systematic uncertainty
- Partial width for $D^0\bar{D}^{*0}$ channel is rather large

Search for the double-charmonium state with $\eta_c J/\psi$ at Belle

- LHCb, CMS, and ATLAS observed new resonances in the $J/\psi J/\psi$ ($c\bar{c}c\bar{c}$) invariant mass distributions.

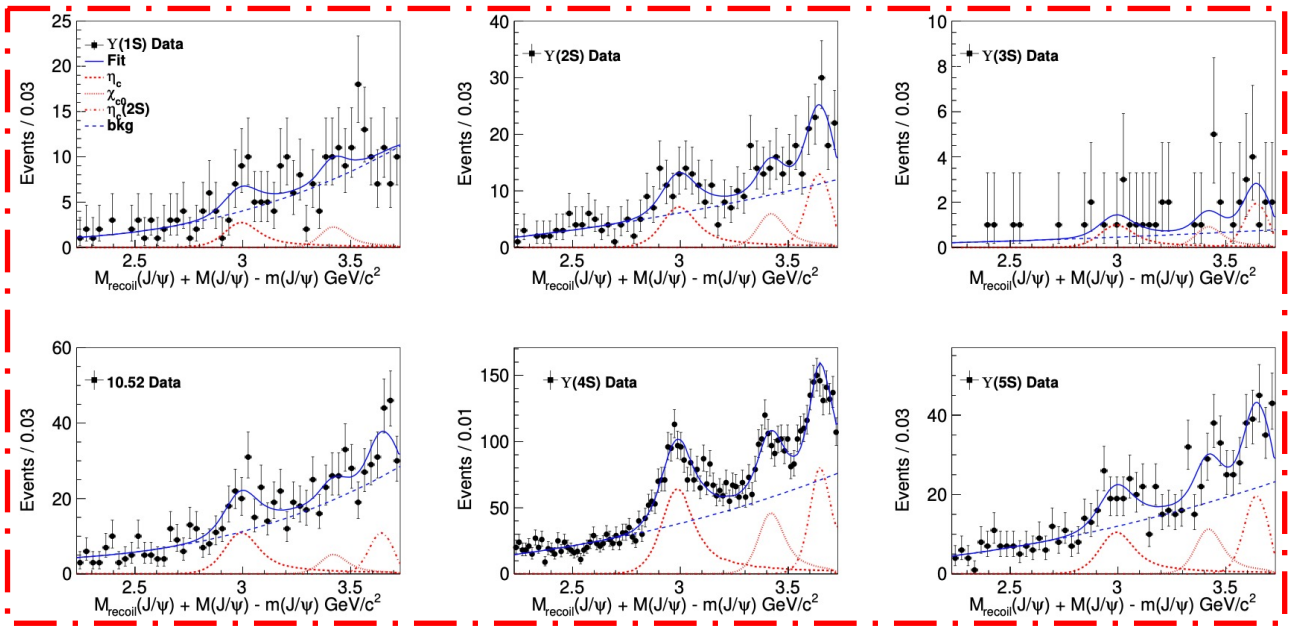
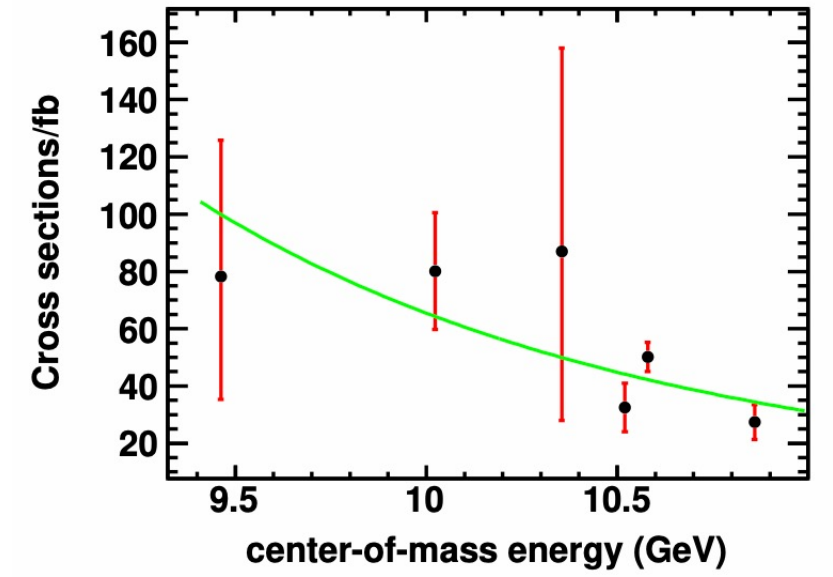
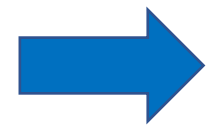
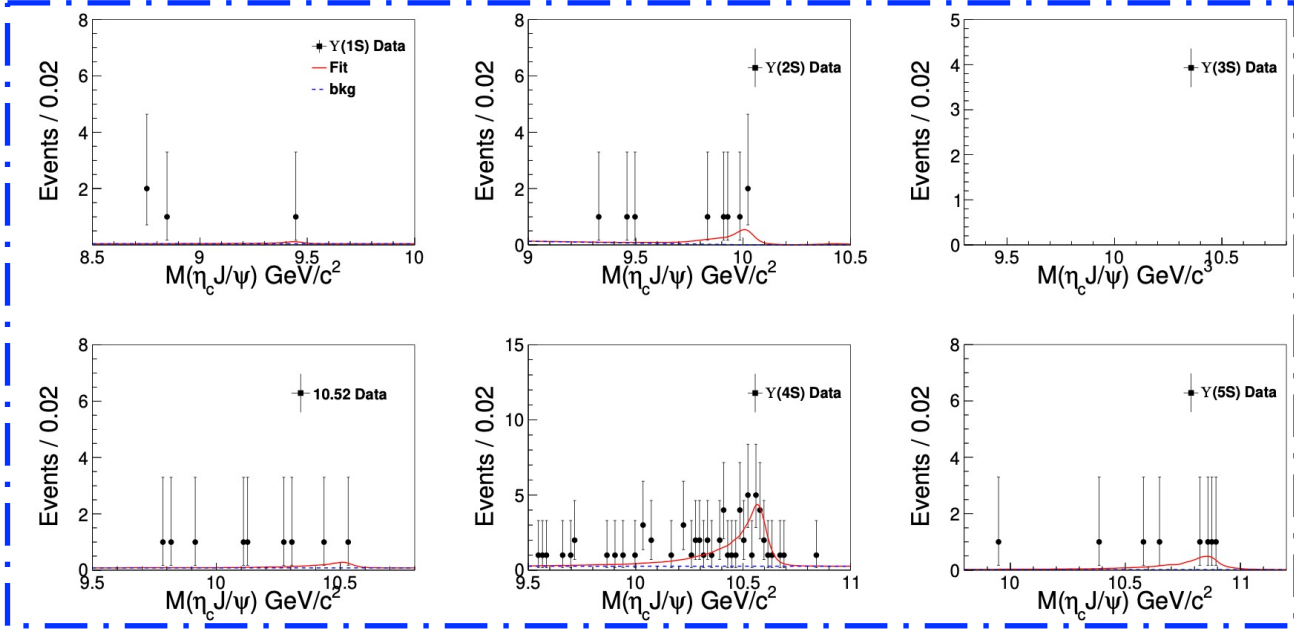


- The lowest mass combination of charmonia to which a vector $c\bar{c}c\bar{c}$ (Y_{cc}) could decay is $\eta_c J/\psi$, and this process may have a relative large branching fraction [Phys. Rev. D 73, 094510 (2006)].

The energy dependency of the Born cross sections for $e^+e^- \rightarrow \eta_c J/\psi$

[arXiv: 2305.17947]

Exclusive process



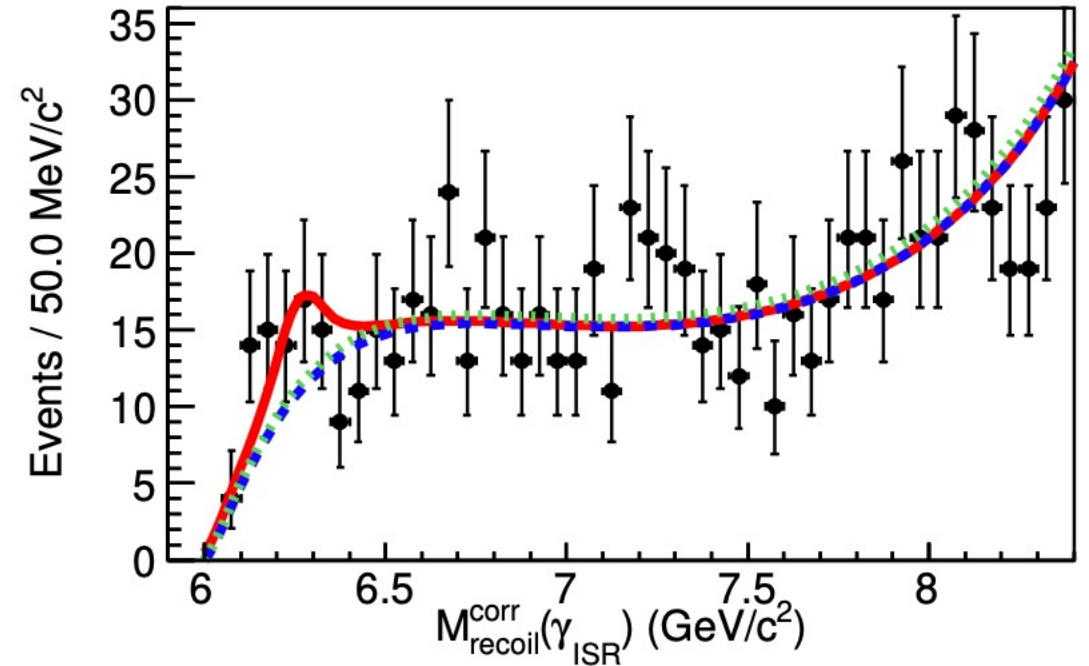
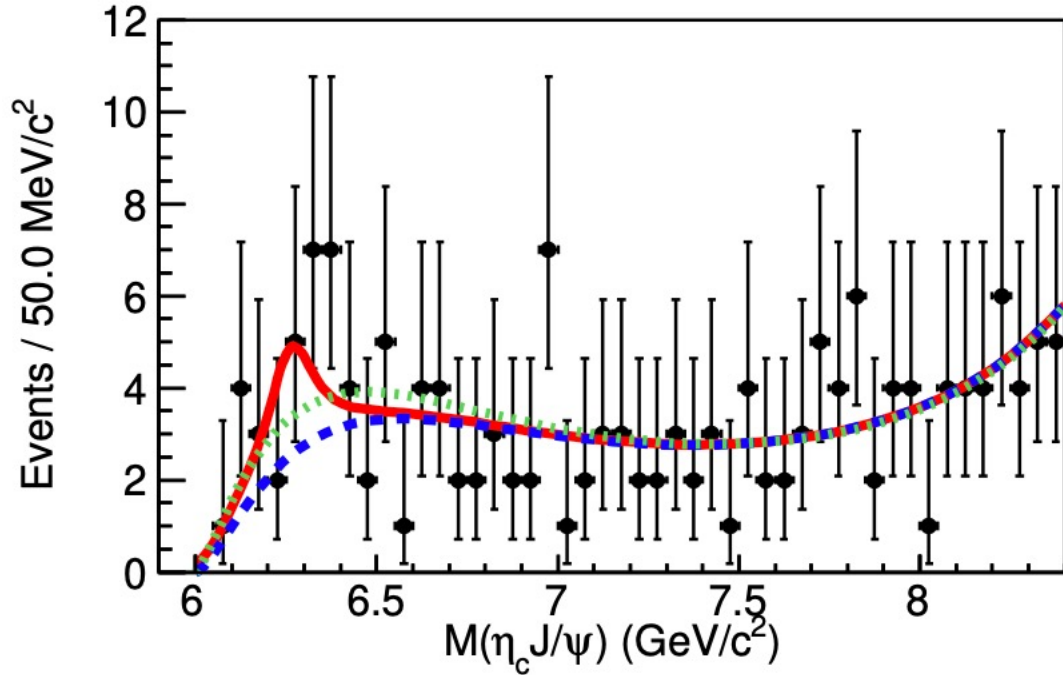
Inclusive process

$$\sigma = A \frac{\sqrt{2\mu\Delta M}}{\left(\frac{s}{s_0}\right)^n}$$

$$n = 4.5 \pm 1.3$$

$e^+e^- \rightarrow \eta_c J/\psi$ near threshold

[arXiv: 2305.17947]



- Red solid curve: best fits
- Blue dashed curve: the background component
- Green dotted curve: fits without the signal component

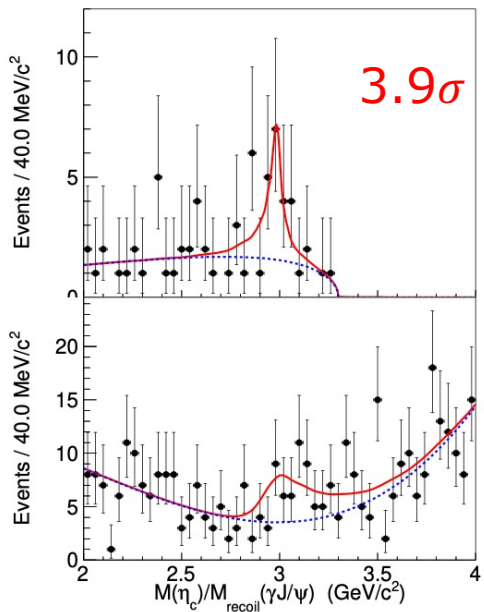
| Parameters | Exclusive | Inclusive |
|--------------|---------------------------------|-------------|
| Mass | $(6267 \pm 43) \text{ MeV}/c^2$ | |
| Width | $(121 \pm 72) \text{ MeV}$ | |
| Yield | 9 ± 4 | 23 ± 11 |
| Significance | 2.2σ | |

Measured cross sections of $e^+e^- \rightarrow \eta_c J/\psi$ near the threshold

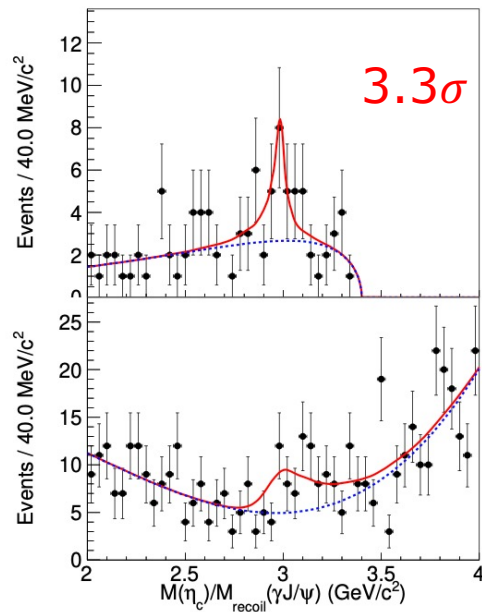
A simultaneous unbinned maximum likelihood fit for the η_c invariant mass and $\gamma_{\text{ISR}} J/\psi$ recoil mass:

[arXiv: 2305.17947]

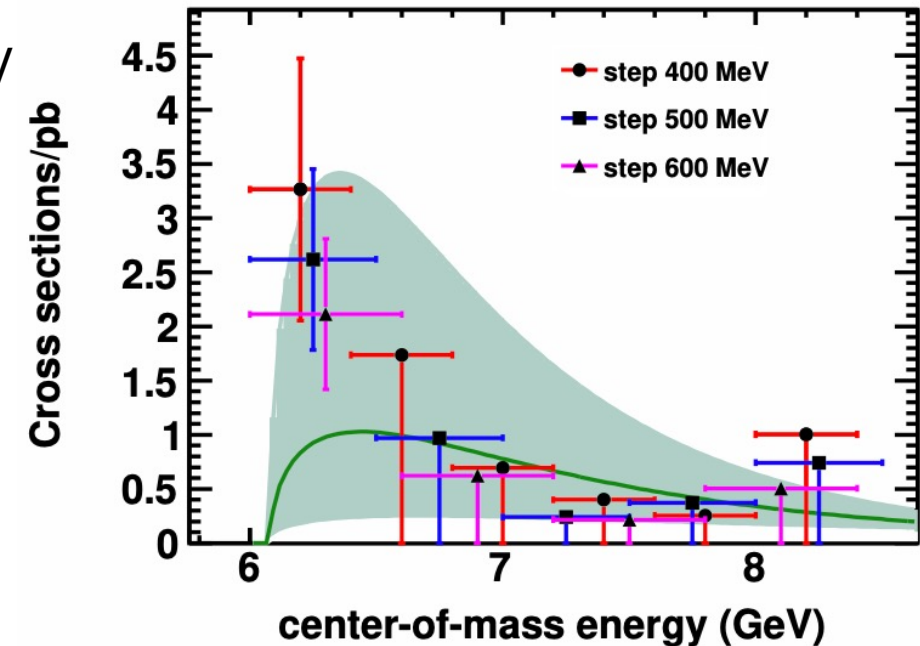
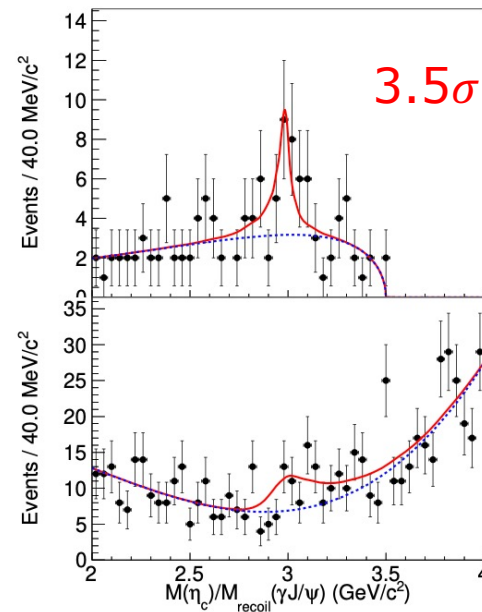
$\sqrt{s} = 6.0 - 6.4$ GeV



$\sqrt{s} = 6.0 - 6.5$ GeV



$\sqrt{s} = 6.0 - 6.6$ GeV

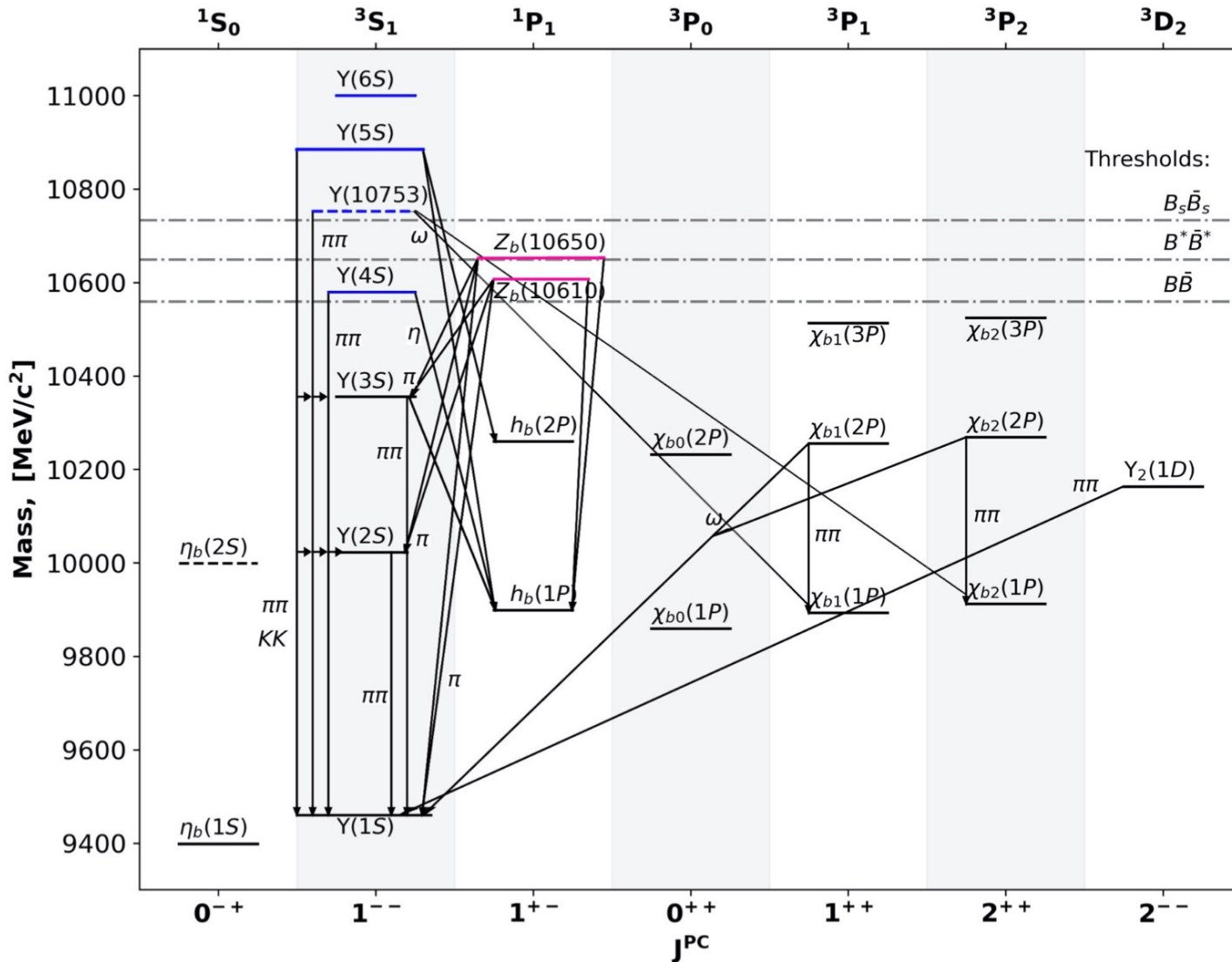


- Green solid curve: the fit
- Shadow area: $\pm 1\sigma$ region

The evidence for the enhancement of $e^+e^- \rightarrow \eta_c J/\psi$ near the threshold.

Quarkonium at Belle II

Bottomonium



- Below $B\bar{B}$ thresholds – bottomonia are well described by the potential models.

- Above $B\bar{B}$ thresholds – bottomonia express unexpected properties:

- Two charged Z_b^+ states are observed ($B^{(*)}\bar{B}^*$ molecular?)

- Hadronic transitions are strongly enhanced (OZI rule violation);

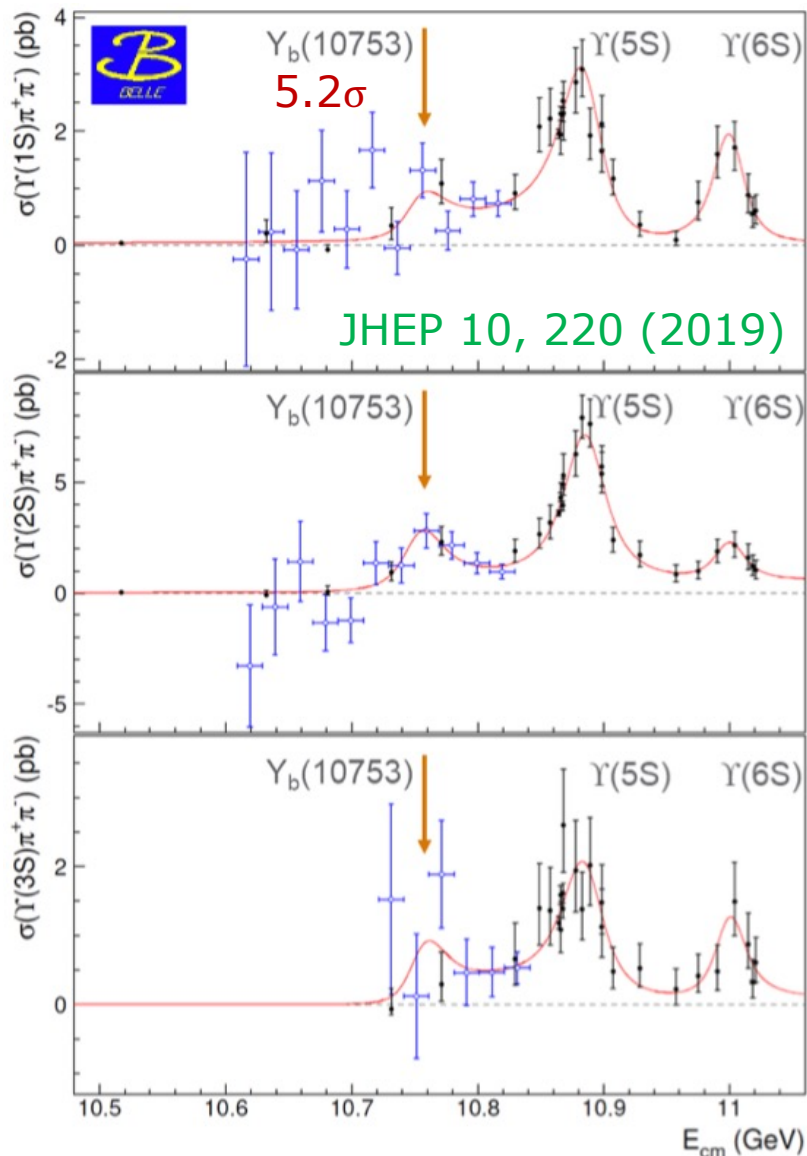
- η transitions are not suppressed compare to $\pi^+\pi^-$ transitions (heavy quark spin-symmetry violation);

Conventional bottomonium (pure $b\bar{b}$ states)

Bottomonium-like states (mix of $b\bar{b}$ and $B\bar{B}$)

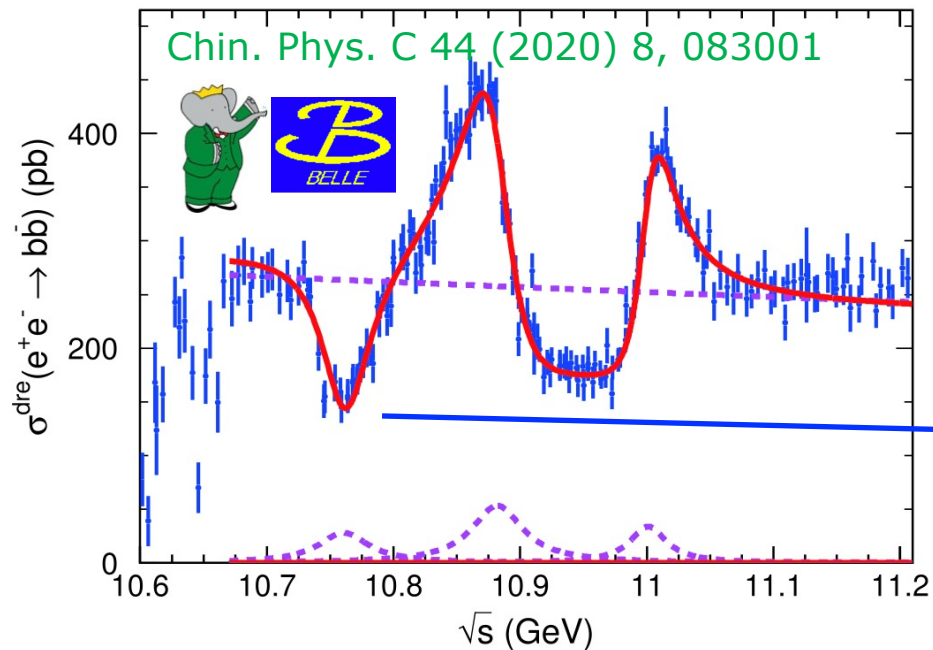
Exotic charged states (Z_b^+)

Discovery of $\Upsilon(10753)$



- Belle: several $\sim 1\text{fb}^{-1}$ scan points below $\Upsilon(5S)$
- New structure observed in $\pi^+\pi^-\Upsilon(nS)$ transitions

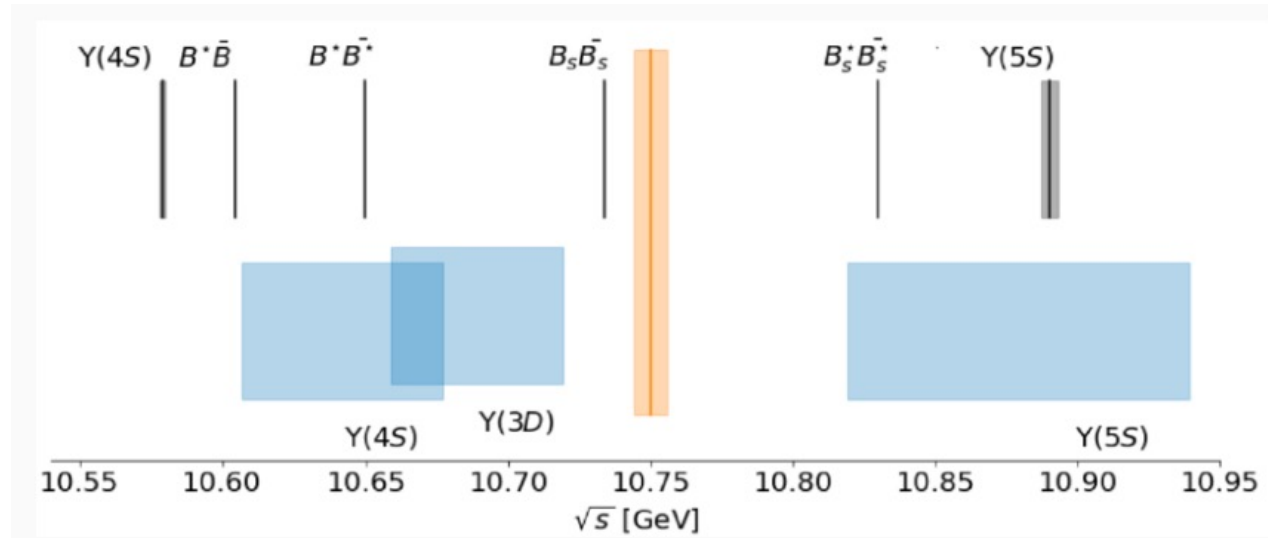
| | $\Upsilon(10860)$ | $\Upsilon(11020)$ | New structure |
|-------------------|---------------------------------------|--|---|
| M (MeV/ c^2) | $10885.3 \pm 1.5^{+2.2}_{-0.9}$ | $11000.0^{+4.0}_{-4.5} {}^{+1.0}_{-1.3}$ | $10752.7 \pm 5.9^{+0.7}_{-1.1}$ |
| Γ (MeV) | $36.6^{+4.5}_{-3.9} {}^{+0.5}_{-1.1}$ | $23.8^{+8.0}_{-6.8} {}^{+0.7}_{-1.8}$ | $35.5^{+17.6}_{-11.3} {}^{+3.9}_{-3.3}$ |



A dip at 10.75 GeV may correspond to $\Upsilon(10753)$.

Theoretical interpretations

Godfrey and Moats, PRD 92, 054034 (2015)



- Mass does not match $Y(3D)$ theoretical predictions, and D-wave states are not seen in e^+e^- collisions.
- $Y(4S)$ - $Y(3D)$ mixing can be enhanced due to hadron loops.

□ Conventional bottomonium

Eur. Phys. J. C 80, 59 (2020)

Phys. Rev. D 101, 014020 (2020)

Phys. Rev. D 102, 014036 (2020)

Phys. Lett. B 803, 135340 (2020)

Phys. Rev. D 104, 034036 (2021)

Prog. Part. Nucl. Phys. 117, 103845 (2021)

Eur. Phys. J. Plus 137, 357 (2022)

Phys. Rev. D 105, 114041 (2022)

Phys. Rev. D 106, 094013 (2022)

Phys. Rev. D 105, 074007 (2022)

□ Hybrid

Phys. Rept. 873, 1 (2020)

Phys. Rev. D 104, 034019 (2021)

□ Tetraquark

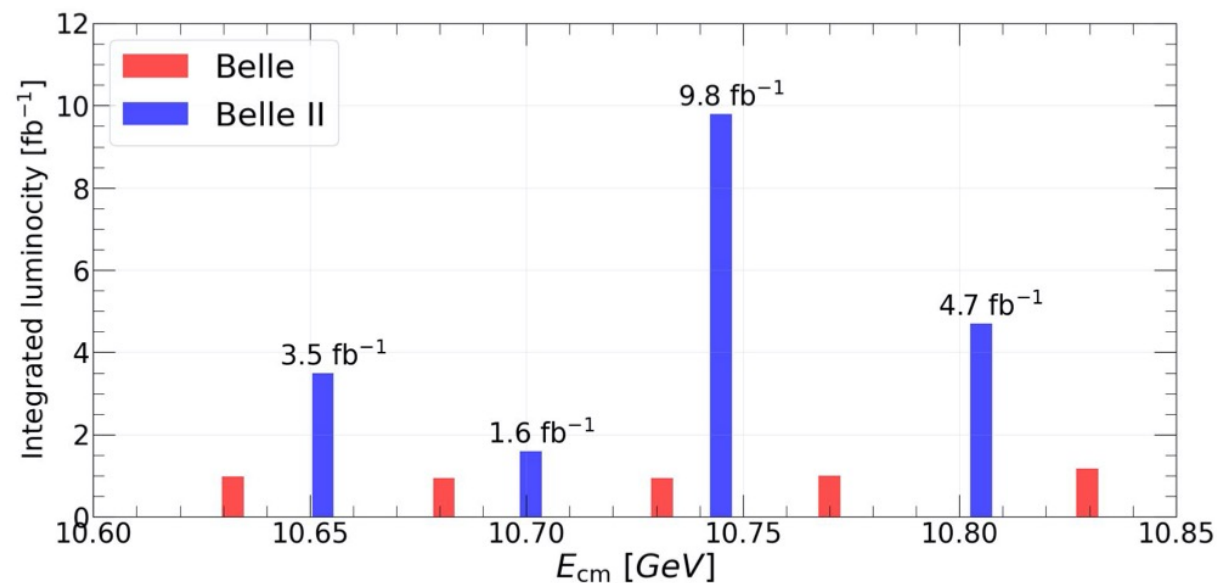
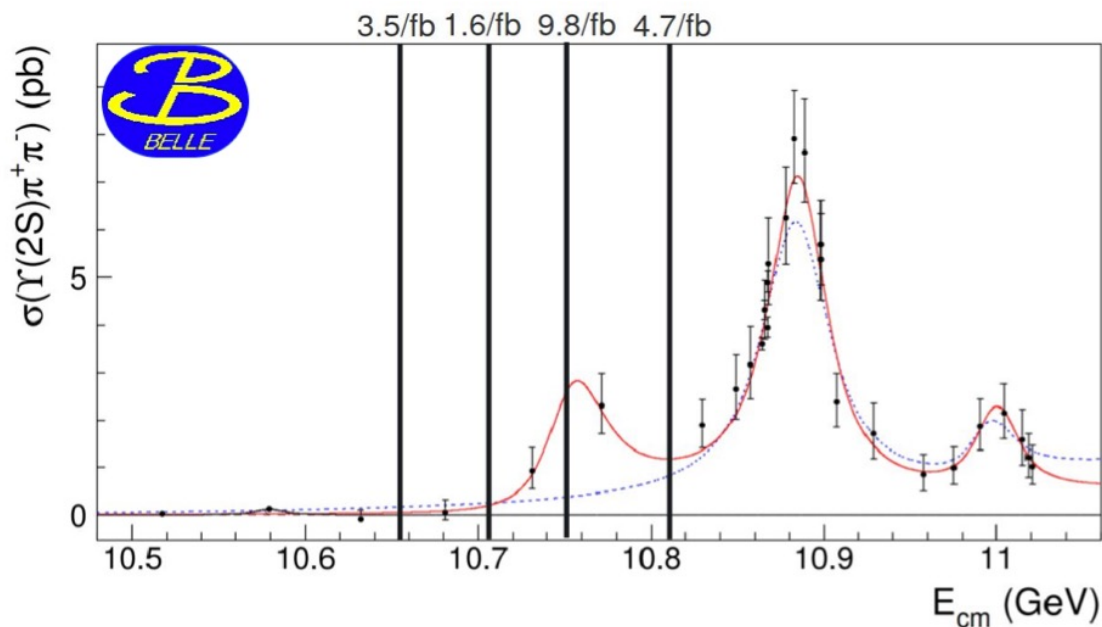
Chin. Phys. C 43, 123102 (2019)

Phys. Lett. B 802, 135217 (2020)

Phys. Rev. D 103, 074507 (2021)

Phys. Rev. D 107, 094515 (2023)

Unique scan data near $\sqrt{s} = 10.75$ GeV



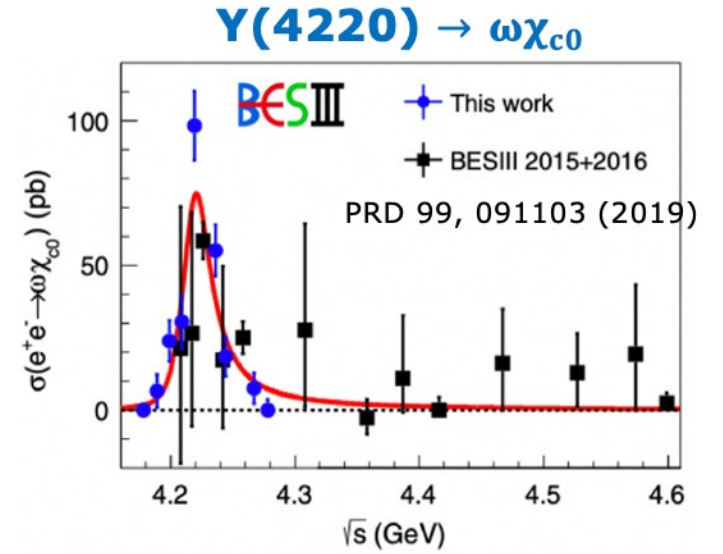
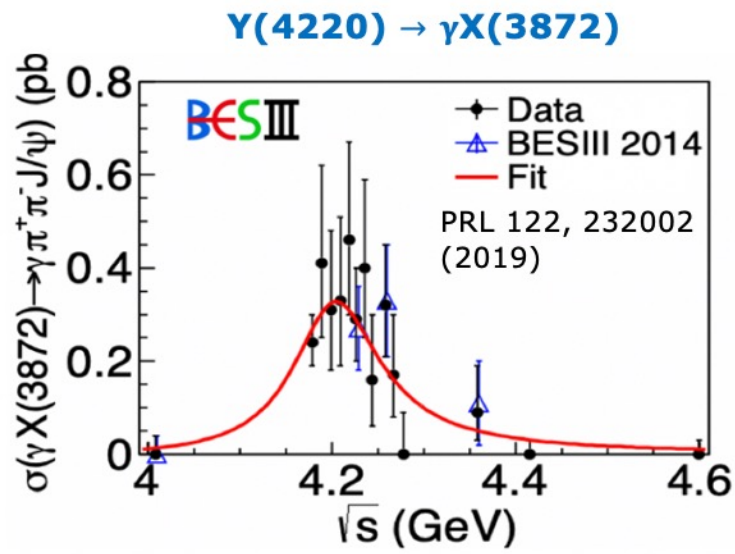
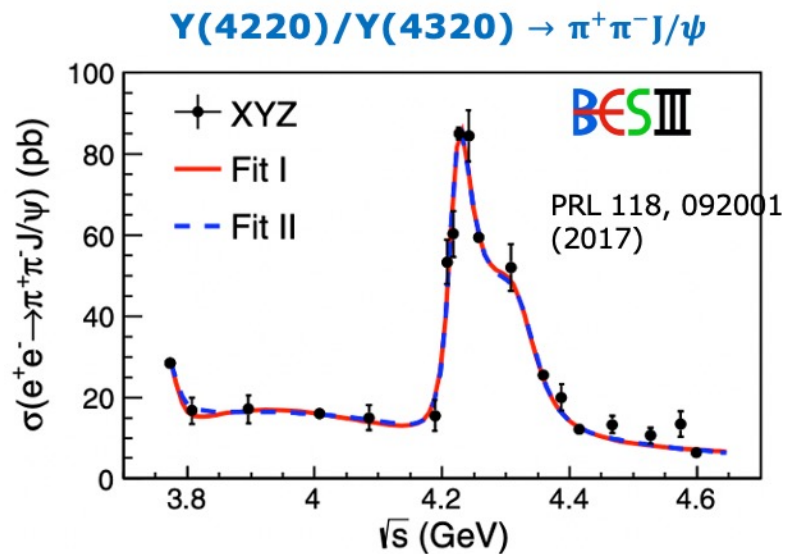
- In November 2021, Belle II collected 19 fb⁻¹ of unique data at energies above the $\Upsilon(4S)$: four energy scan points around 10.75 GeV.
- Belle II collected the data in the gaps between Belle energy scan points.
- Physics goal: understand the nature of the $\Upsilon(10753)$ energy region.

Motivation to search for $\Upsilon(10753) \rightarrow \omega\chi_{bJ}$

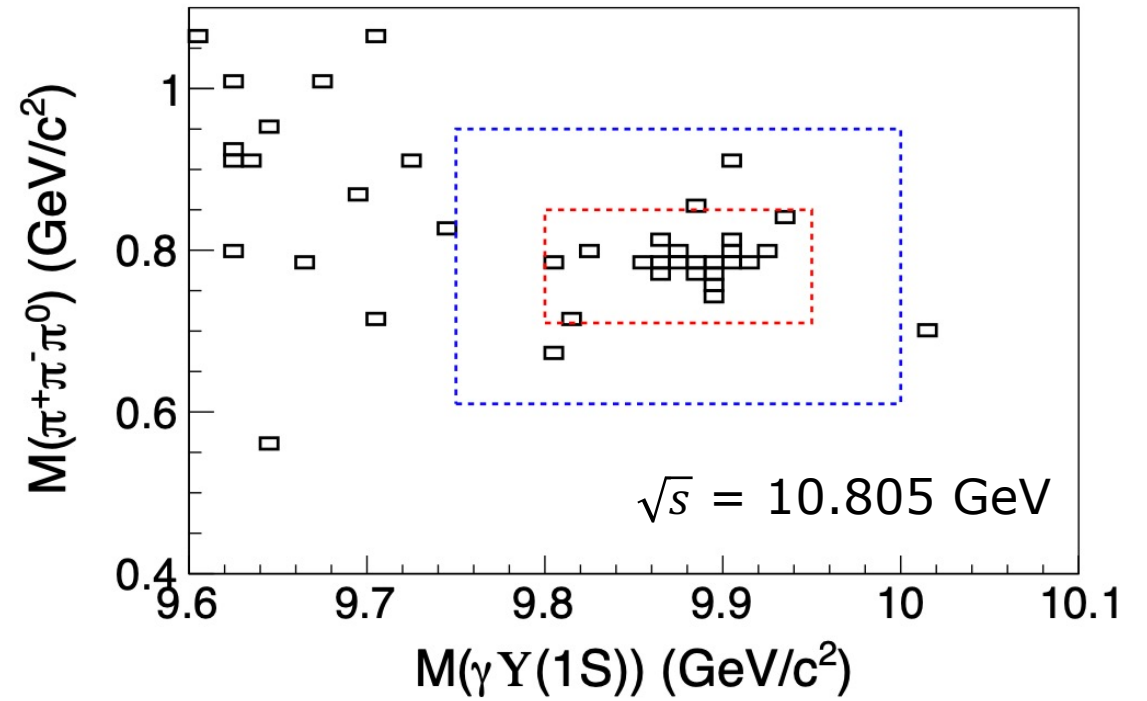
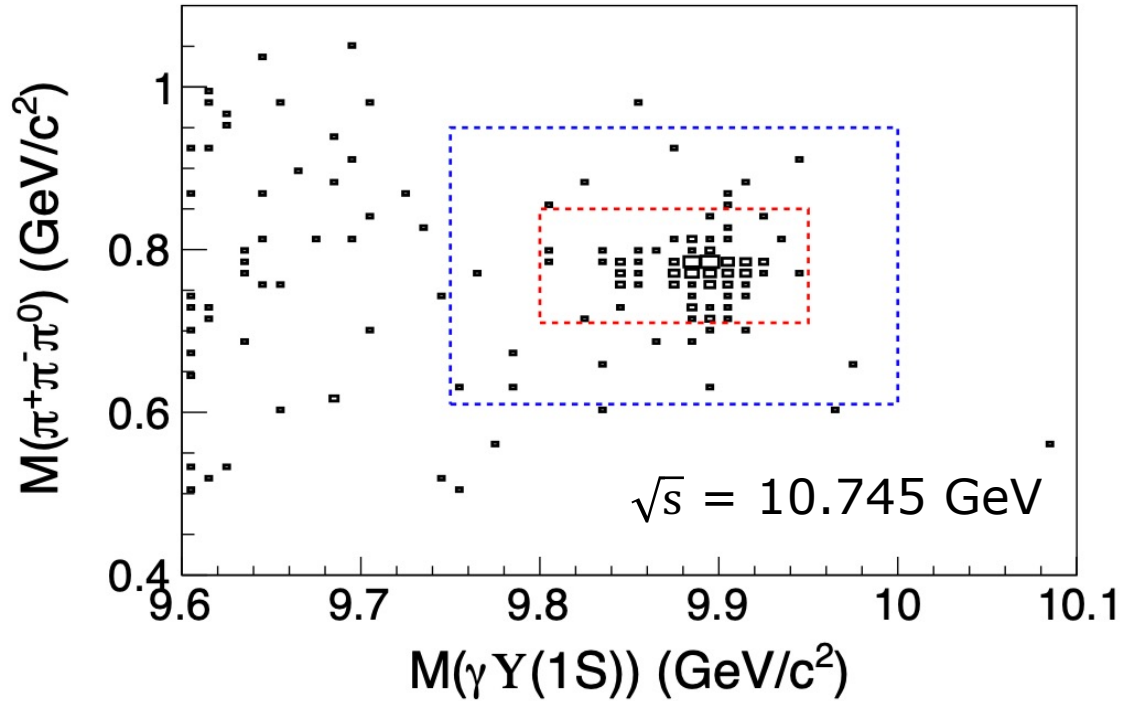
Theory: Branching fractions of 10^{-3} for $\Upsilon(10753) \rightarrow \omega\chi_{bJ}$ [PRD 104, 034036 (2021)] and $\Upsilon(10753) \rightarrow \pi^+\pi^-\Upsilon(nS)$ [PRD 105, 074007 (2022)] assuming $\Upsilon(4S) - \Upsilon(3D)$ mixing state for $\Upsilon(10753)$.

Charmonium sector:

- Two close peaks observed in the cross sections for $e^+e^- \rightarrow \pi^+\pi^-J/\psi$ by BESIII and $e^+e^- \rightarrow \pi^+\pi^-\Upsilon(nS)$ by Belle, respectively, may suggest similar nature.
- $\Upsilon(4220) \rightarrow \gamma X(3872)$ and $\omega\chi_{c0}$ observed by BESIII.
- So we expect the observations of $\Upsilon(10753) \rightarrow \gamma X_b$ and $\omega\chi_{bJ}$.



Mass distributions

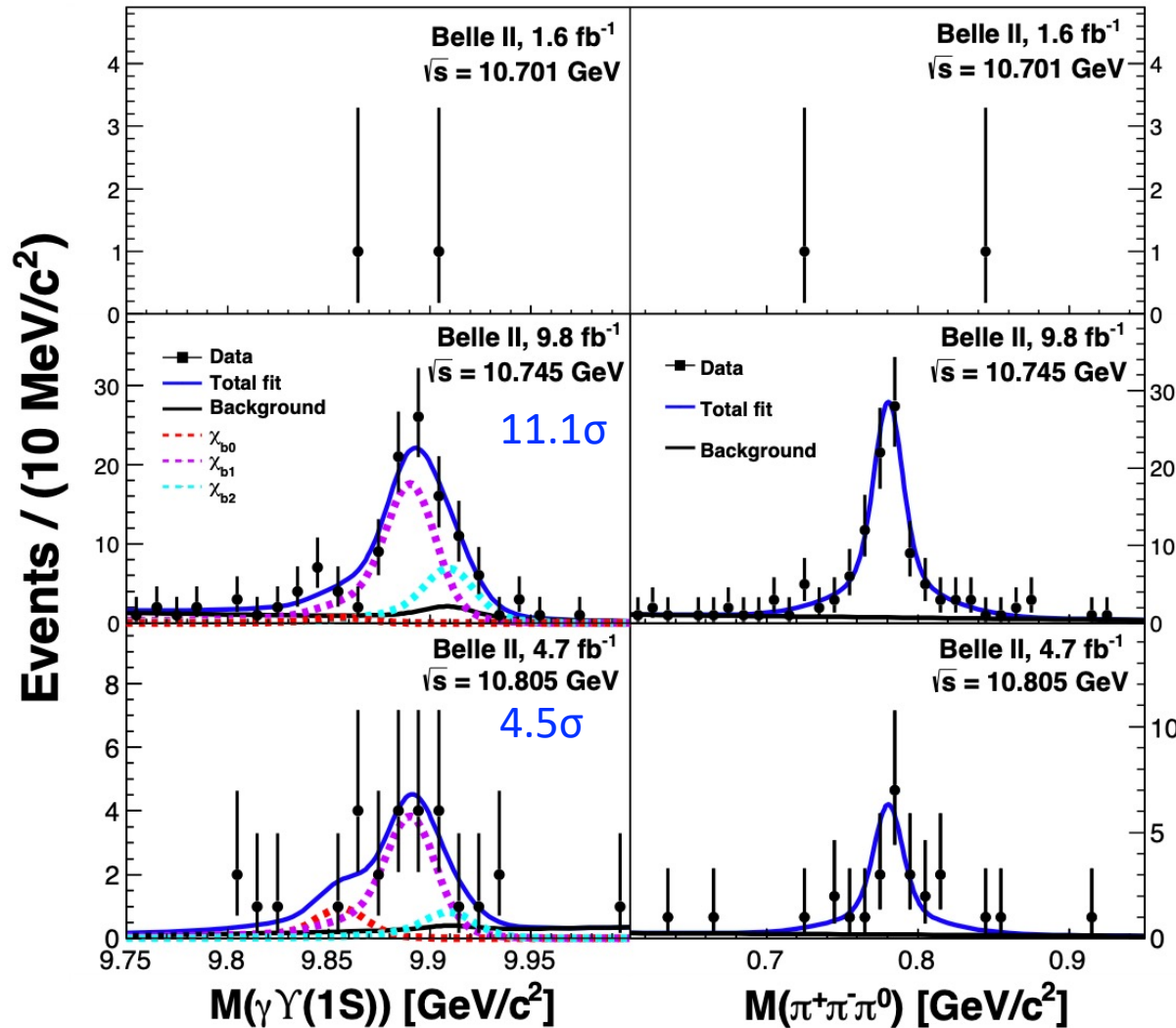


- Red boxes contains 95% of signals.
- Blue boxes show the fit ranges.

Observation of $\Upsilon(10753) \rightarrow \omega\chi_{bJ}$

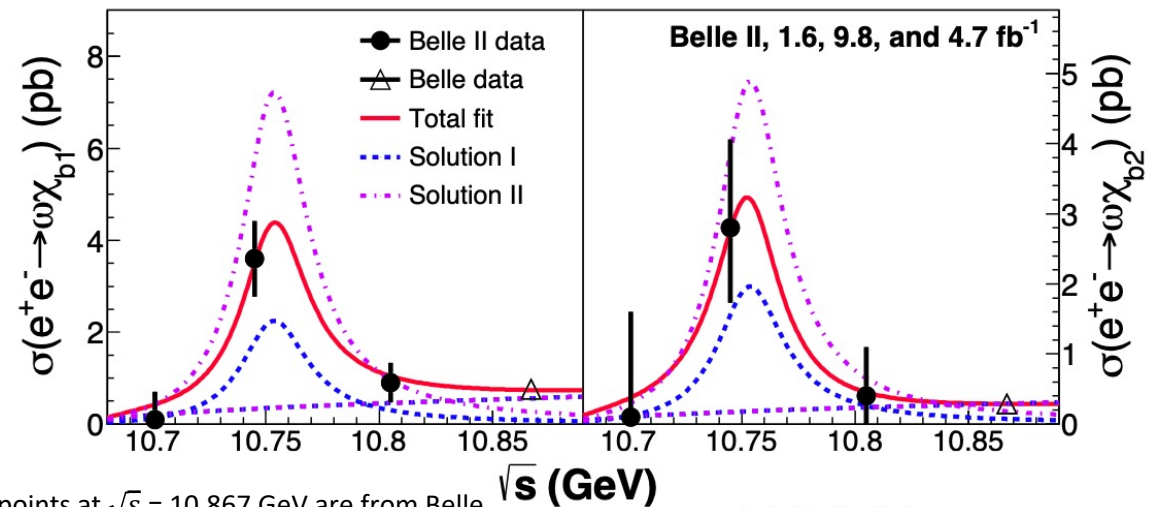
PRL 130, 091902 (2023)

Two dimensional unbinned maximum likelihood fits to the $M(\gamma\Upsilon(1S))$ and $M(\pi^+\pi^-\pi^0)$ distributions.



| Channel | \sqrt{s} (GeV) | N ^{sig} | $\sigma_{\text{Born}}^{(\text{UL})}$ (pb) |
|-------------------|------------------|------------------------|---|
| $\omega\chi_{b1}$ | 10.745 | $68.9_{-13.5}^{+13.7}$ | $3.6_{-0.7}^{+0.7} \pm 0.4$ |
| $\omega\chi_{b2}$ | | $27.6_{-10.0}^{+11.6}$ | $2.8_{-1.0}^{+1.2} \pm 0.5$ |
| $\omega\chi_{b1}$ | 10.805 | $15.0_{-6.2}^{+6.8}$ | 1.6 @90% C.L. |
| $\omega\chi_{b2}$ | | $3.3_{-3.8}^{+5.3}$ | 1.5 @90% C.L. |

The $e^+e^- \rightarrow \omega\chi_{bJ}$ ($J = 1, 2$) cross sections peak at $\Upsilon(10753)$.



The points at $\sqrt{s} = 10.867$ GeV are from Belle measurements [PRL 113, 142001 (2014)].

Discussion

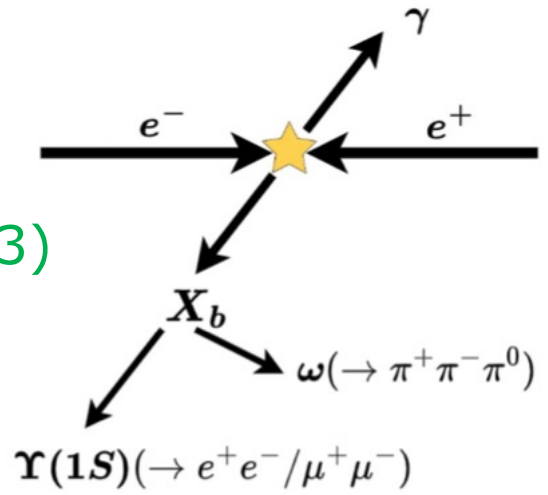
$$\frac{\sigma(e^+e^- \rightarrow \chi_{bJ}(1P)\omega)}{\sigma(e^+e^- \rightarrow Y(nS)\pi^+\pi^-)} \sim \begin{cases} \sim 1.5 \text{ at } \sqrt{s} = 10.745 \text{ GeV [PRL 130, 091902 (2023)]} \\ \sim 0.15 \text{ at } \sqrt{s} = 10.867 \text{ GeV [PRL 113, 142001 (2014)]} \end{cases}$$

- $Y(5S)$ and $Y(10753)$ have same quantum numbers and similar masses, but the difference on the above ratio is large. This may indicate **the difference in the internal structures of these two states**.

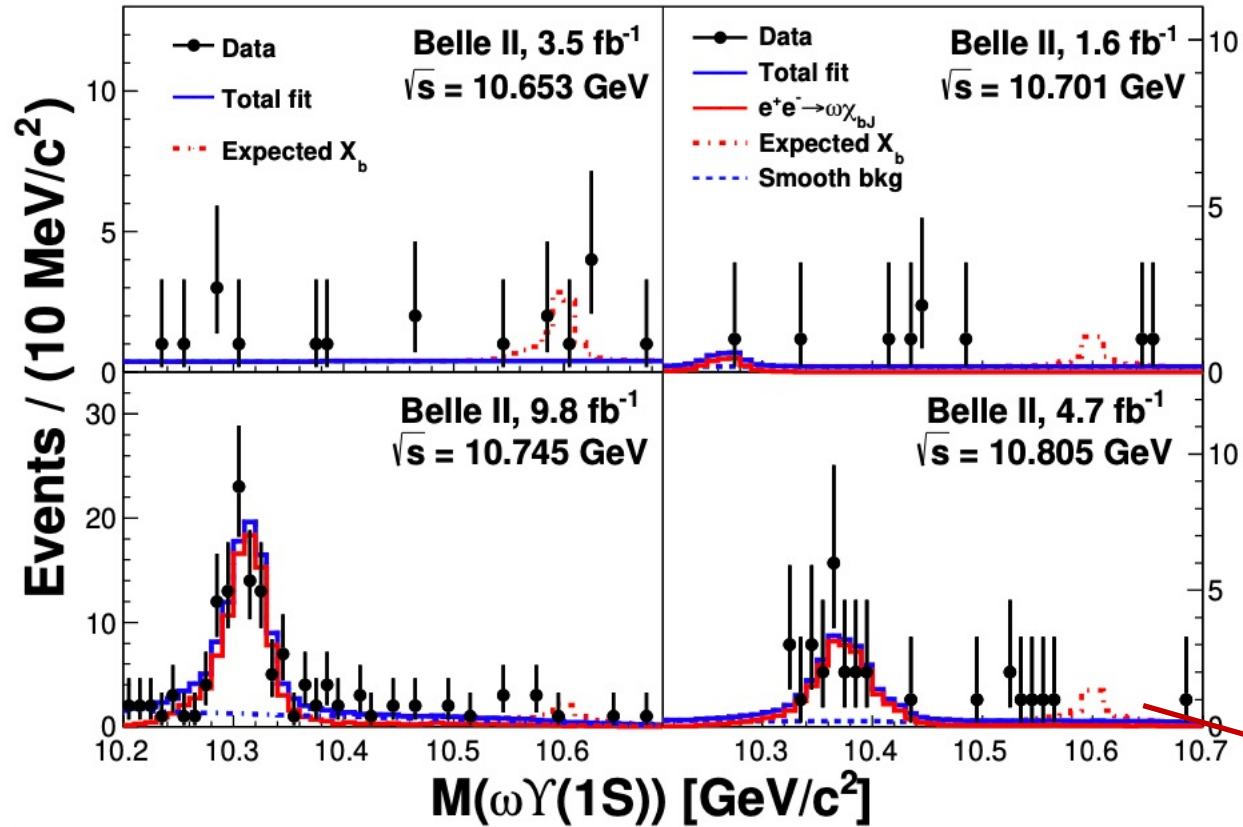
$$\frac{\sigma(e^+e^- \rightarrow \chi_{b1}(1P)\omega)}{\sigma(e^+e^- \rightarrow \chi_{b2}(1P)\omega)} = 1.3 \pm 0.6 \text{ at } \sqrt{s} = 10.745 \text{ GeV [PRL 130, 091902 (2023)]}$$

- **Contradicts the expectation for a pure D-wave bottomonium state of 15** [Phys. Lett. B 738, 172 (2014)]
- **An observation of 1.8σ difference with the prediction for a S-D-mixed state of 0.2** [Phys. Rev. D 104, 034036 (2021)]

Search for X_b



PRL 130, 091902 (2023)



- No significant X_b signal is observed.
- The peaks are the reflections of $e^+e^- \rightarrow \omega\chi_{bJ}$.

From simulated events with $m(X_b) = 10.6 \text{ GeV}/c^2$
The yield is fixed at the upper limit at 90% C.L.

| Upper limits at 90% C.L. on $\sigma_B(e^+e^- \rightarrow \gamma X_b) \cdot \mathcal{B}(X_b \rightarrow \omega\Upsilon(1S))$ (pb) | \sqrt{s} (GeV) | 10.653 | 10.701 | 10.745 | 10.805 |
|--|------------------|--------------|--------------|--------------|--------------|
| $m(X_b) = 10.6 \text{ GeV}/c^2$ | | 0.46 | 0.33 | 0.10 | 0.14 |
| $m(X_b) = (10.45, 10.65) \text{ GeV}/c^2$ | | (0.14, 0.55) | (0.25, 0.84) | (0.06, 0.14) | (0.08, 0.37) |

Measurement of the energy dependence of the $e^+e^- \rightarrow B\bar{B}$, $B\bar{B}^*$ and $B^*\bar{B}^*$ cross sections

$\sqrt{s} = 10.745 \text{ GeV}, 9.8 \text{ fb}^{-1}$

- The $B^{(*)}\bar{B}^{(*)}$ are expected to be dominant decay channels for excited bottomonium-like states. Their measurements are critical for understanding these states.

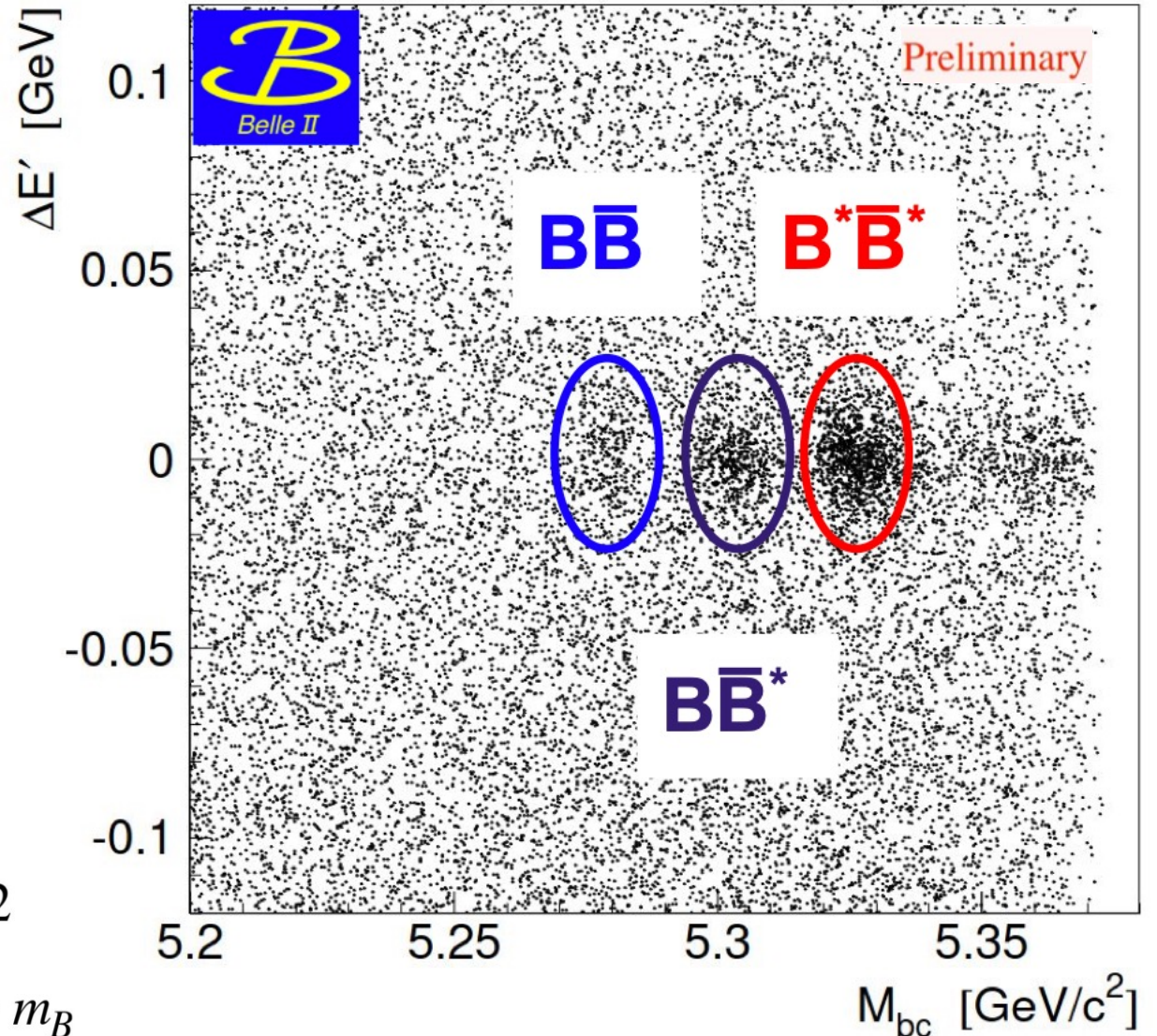
Method:

One B meson is reconstructed in hadronic channels, and signals are identified using

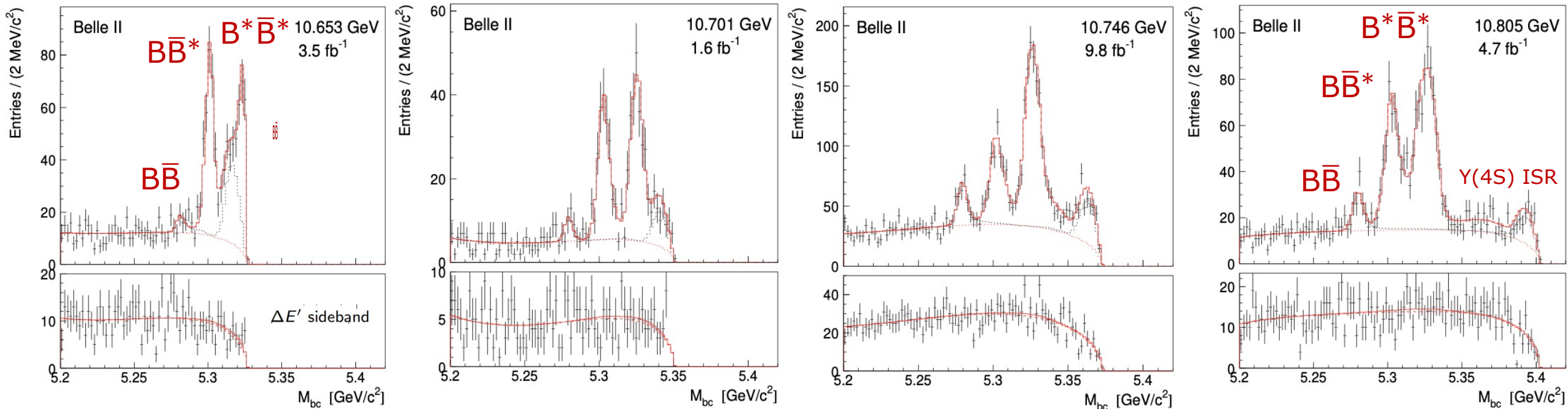
$$M_{bc} = \sqrt{(E_{cm}/2)^2 - p_B^2}$$

$$\Delta E = E_B - E_{cm}/2$$

$$\Delta E' = \Delta E + M_{bc} - m_B$$

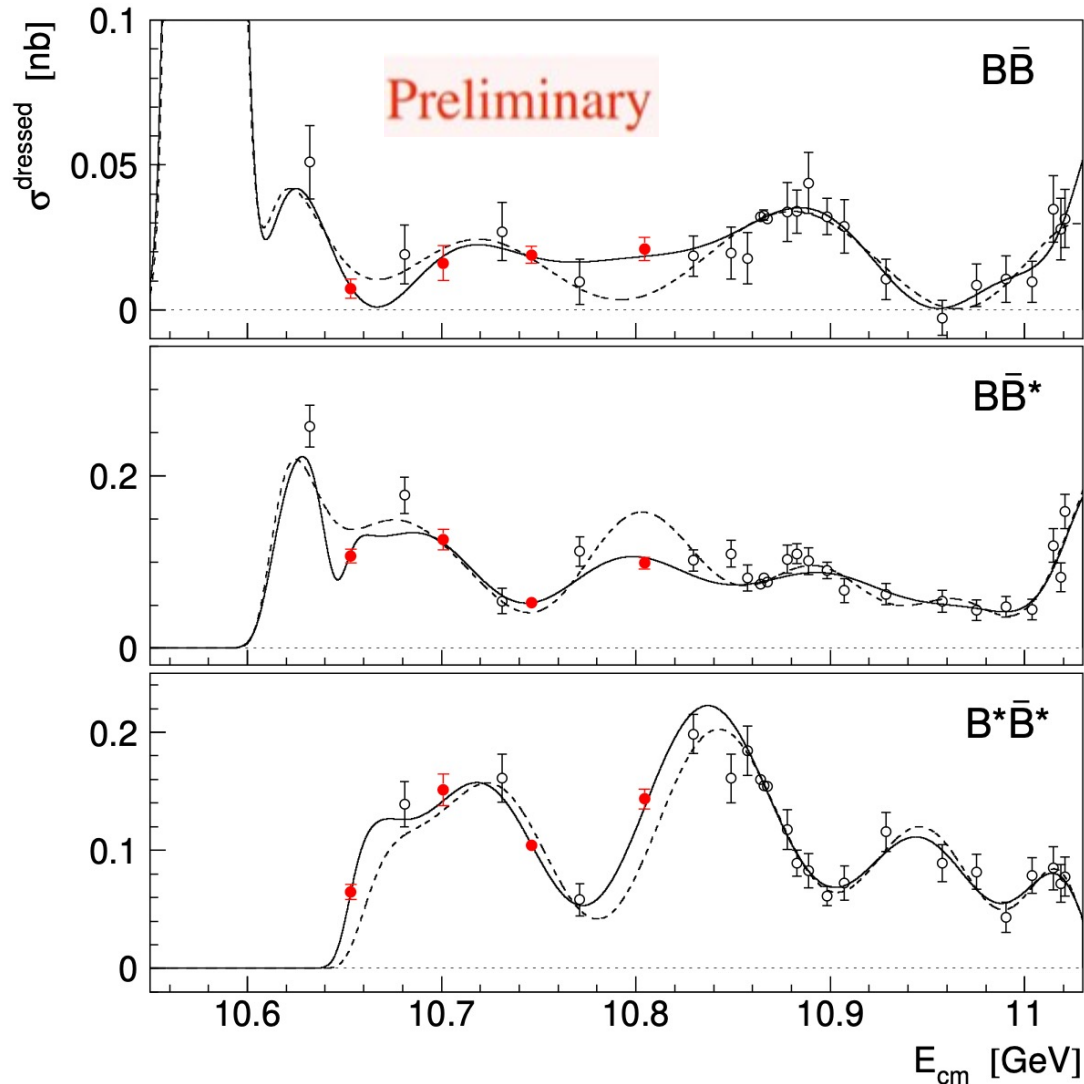


M_{bc} fit at scan energies



- $e^+e^- \rightarrow B\bar{B}, B\bar{B}^*$ and $B^*\bar{B}^*$ signals at $\sqrt{s} \sim 10.75$ GeV can be clearly observed
- Contribution of $Y(4S) \rightarrow B\bar{B}$ production via ISR is visible well (black dotted histograms)
- At $\sqrt{s} = 10.653$ GeV, the sharp cut of the data at right edge is due to threshold effect

Energy dependence of the cross sections

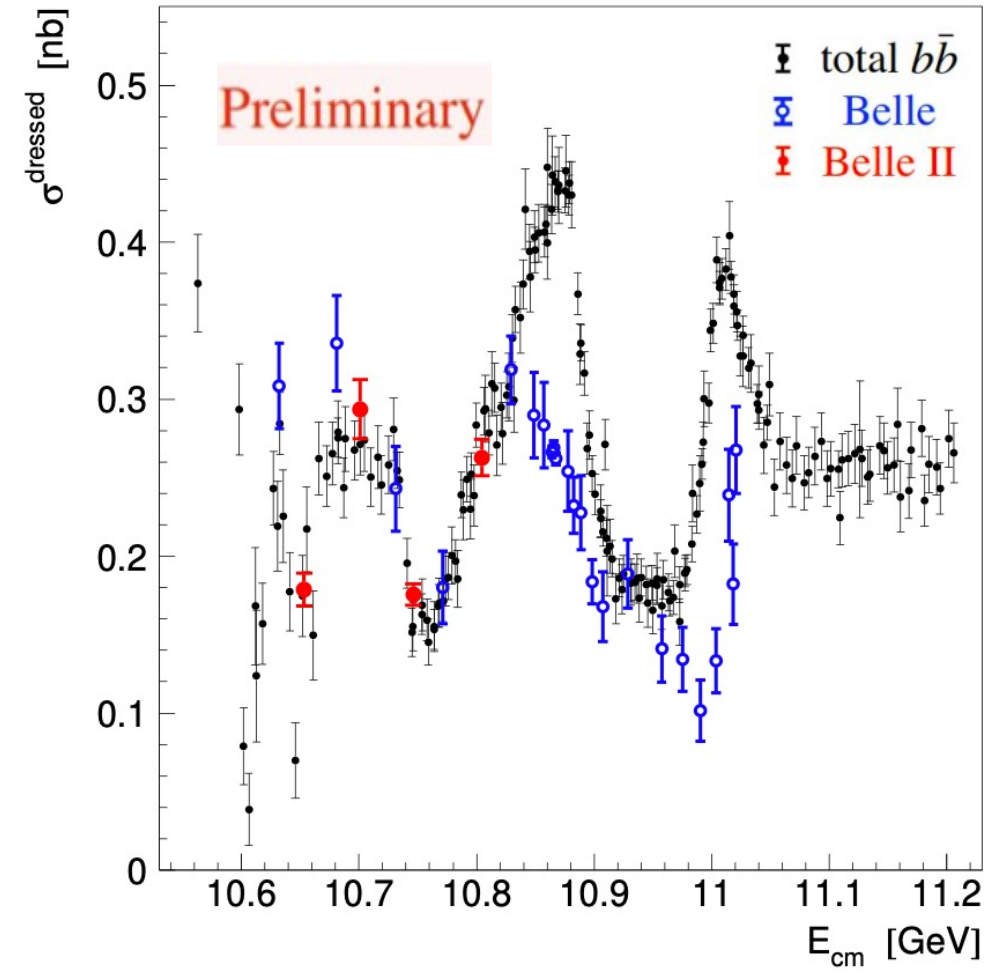
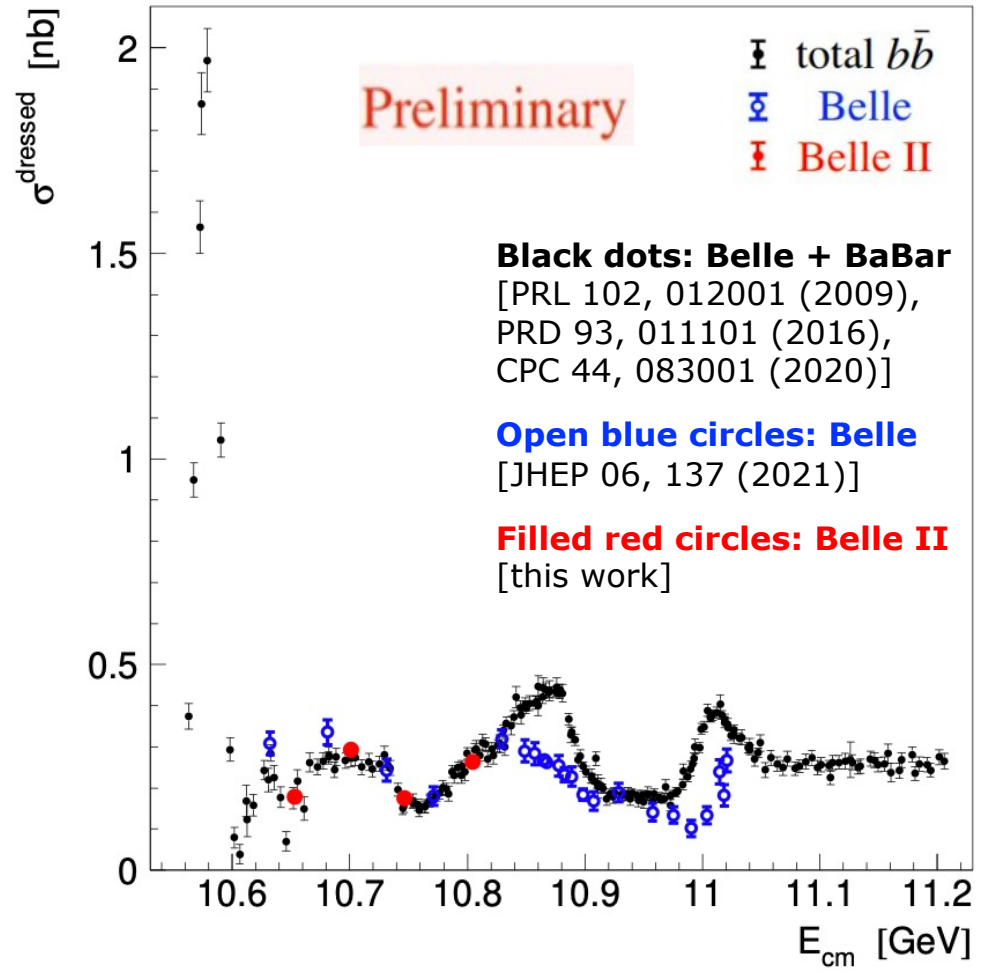


New: rapid increase of $\sigma_{B^*\bar{B}^*}$ above the threshold

- Similar behaviour was seen for $D^*\bar{D}^*$ cross section (PRD 97, 012002 (2018))
- Possible interpretation: **resonance or bound state** ($B^*\bar{B}^*$ or $b\bar{b}$) near threshold (MPL A 21, 2779 (2006))
- Also explains a narrow dip in $\sigma(e^+e^- \rightarrow B\bar{B}^*)$ near $B^*\bar{B}^*$ threshold by destructive interference between $e^+e^- \rightarrow B\bar{B}^*$ and $e^+e^- \rightarrow B^*\bar{B}^* \rightarrow B\bar{B}^*$
- Inelastic channels [$\pi^+\pi^-\Upsilon(nS)$ and $h_b(1P)\eta$] could also be enhanced (PRD 87, 094033 (2013))

Solid curve – combined Belle + Belle II data fit
Dashed curve – Belle data fit only

Comparison of $\sigma_{b\bar{b}}$ and $\sigma_{B\bar{B}} + \sigma_{B\bar{B}^*} + \sigma_{B^*\bar{B}^*}$



- Agreement at low energy
- Departure at high energy is due to $B_s^{(*)}\bar{B}_s^{(*)}$, multi-body $B^{(*)}\bar{B}^{(*)}\pi(\pi)$, and bottomonia

Search for $e^+e^- \rightarrow \omega\eta_b(1S)$ and $e^+e^- \rightarrow \omega\chi_{b0}(1P)$

- Tetraquark (diquark-antidiquark) interpretation of this state predicts **enhancement of $Y(10753) \rightarrow \omega\eta_b(1S)$ transition** [Chin. Phys. C 43, no.12, 123102 (2019)].

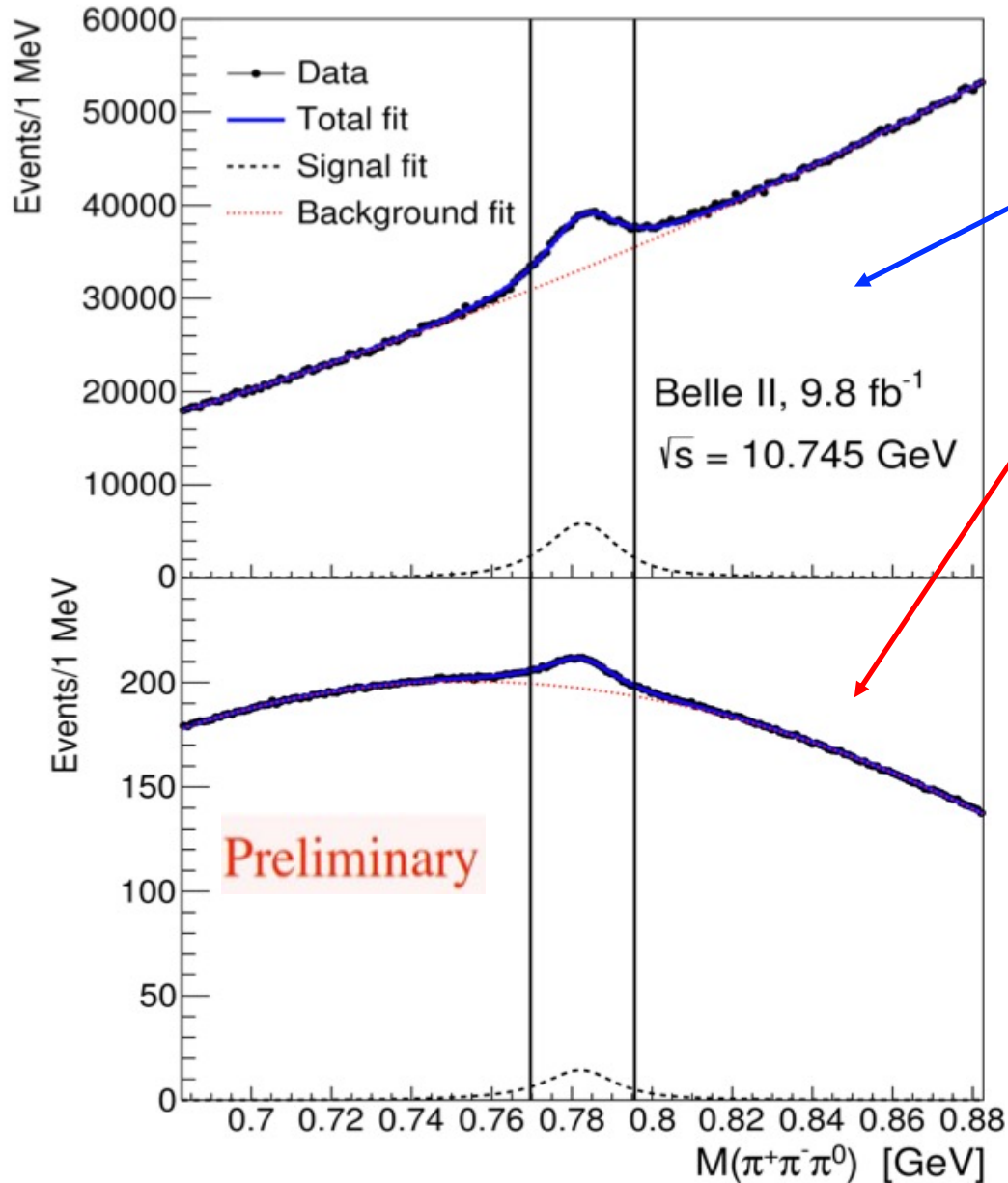
$$\frac{\Gamma(\eta_b \omega)}{\Gamma(\Upsilon \pi^+ \pi^-)} \sim 30$$

- The $e^+e^- \rightarrow \omega\chi_{bJ}(1P)$ ($J = 1, 2$) was found to be enhanced at $\sqrt{s} = 10.745$ GeV (PRL 130, 091902 (2023)). The $e^+e^- \rightarrow \omega\chi_{b0}(1P)$ transition was not observed due to low $\mathcal{B}[\chi_{b0}(1P) \rightarrow \gamma Y(1S)] = (1.94 \pm 0.27)\%$.

- We reconstruct only $\omega \rightarrow \pi^+ \pi^- \pi^0$ and use **its recoil mass to identify the signal**.

$$M_{\text{recoil}}(\pi^+ \pi^- \pi^0) = \sqrt{\left(\frac{E_{\text{c.m.}} - E^*}{c^2}\right)^2 - \left(\frac{p^*}{c}\right)^2}$$

Invariant mass distribution of $\pi^+\pi^-\pi^0$

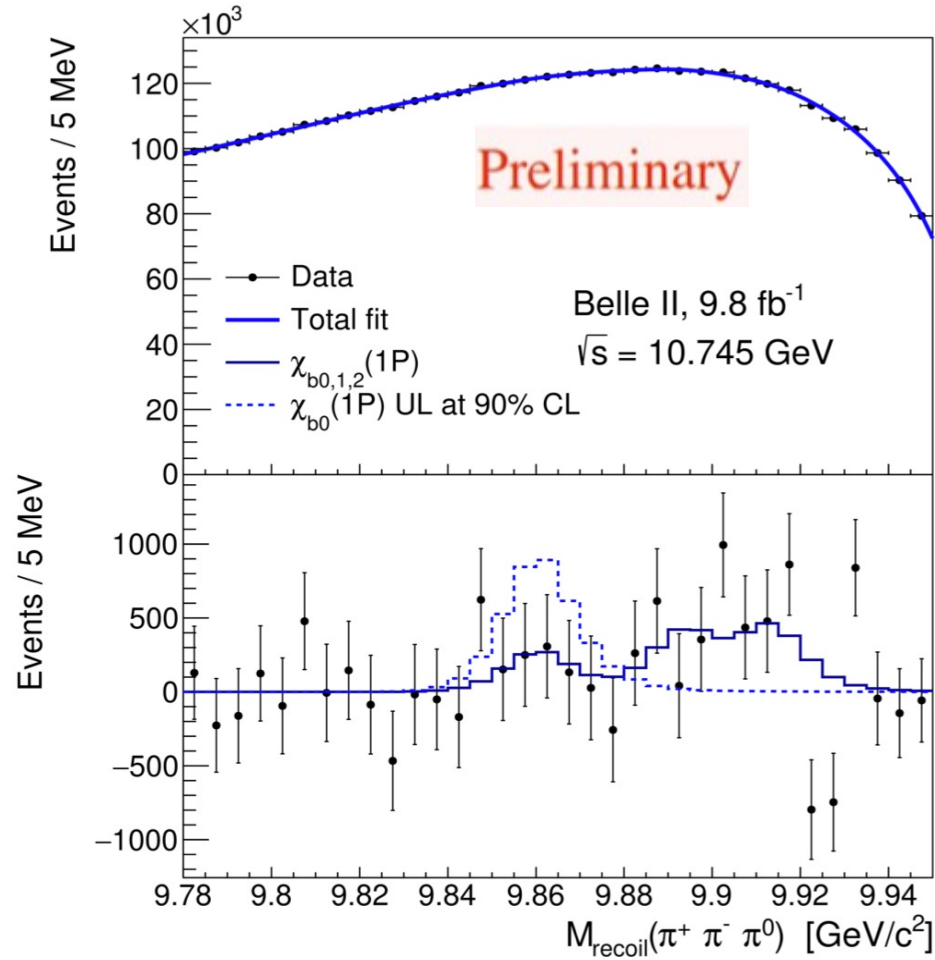
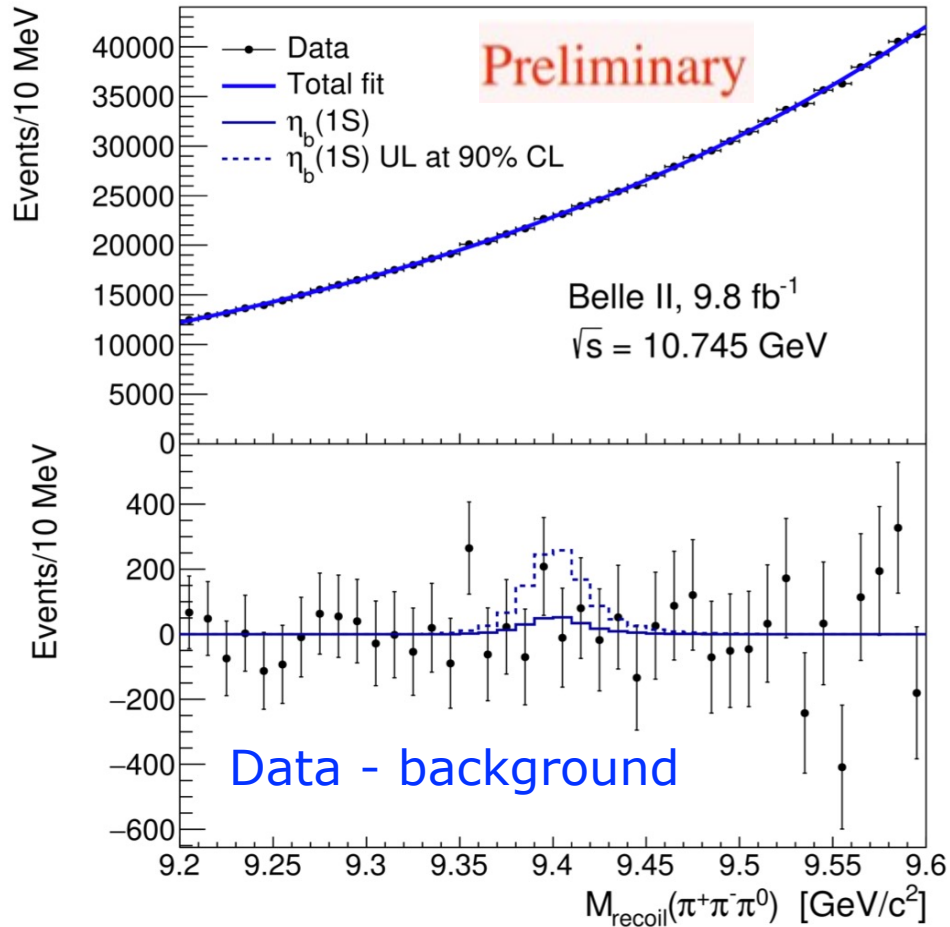


$9.2 < M_{\text{rec}}(\pi^+\pi^-\pi^0) < 9.6 \text{ GeV}/c^2$
($\eta_b(1S)$ included)

$9.78 < M_{\text{rec}}(\pi^+\pi^-\pi^0) < 9.95 \text{ GeV}/c^2$
($\chi_{bJ}(1P)$ included)

- A double-sided Crystal Ball + a Gaussian for ω signal
- 2nd or 3rd order Chebyshev polynomials for backgrounds
- The purities of ω -meson signals are 12.9% for $\eta_b(1S)$ and 5.3% for $\chi_{bJ}(1P)$

Recoil mass spectra of $\pi^+\pi^-\pi^0$



- A 3rd polynomial for $\eta_b(1S)$
- A product of a 4th polynomial and a square root function for $\chi_{b0}(1P)$
- Polynomial orders are chosen with maximum p-values
- The yields for $\chi_{b1}(1P)$ and $\chi_{b2}(1P)$ are fixed [PRL 130, 091902 (2023)].

| Channel | $e^+e^- \rightarrow \eta_b(1S)\omega$ | $e^+e^- \rightarrow \chi_{b0}(1P)\omega$ |
|---------|---------------------------------------|--|
|---------|---------------------------------------|--|

| | | |
|-------|---------------------------------------|------------------------------------|
| Yield | $(0.23 \pm 0.49 \pm 0.25) \cdot 10^3$ | $(1.2 \pm 1.4 \pm 0.9) \cdot 10^3$ |
|-------|---------------------------------------|------------------------------------|

No clear $\eta_b(1S)$ and $\chi_{b0}(1P)$ signals are observed.

Born cross sections

$$\sigma_B[e^+e^- \rightarrow X\omega] = \frac{N \cdot |1 - \Pi|^2}{\varepsilon \cdot \mathcal{L} \cdot (1 + \delta_{\text{ISR}}) \cdot \mathcal{B}_{\text{int}}}$$

Preliminary

| Channel | $e^+e^- \rightarrow \eta_b(1S)\omega$ | $e^+e^- \rightarrow \chi_{b0}(1P)\omega$ |
|------------------------------|---------------------------------------|--|
| Yield (10^3) | $0.23 \pm 0.49 \pm 0.25$ | $1.2 \pm 1.4 \pm 0.9$ |
| Born section section (pb) | $0.5 \pm 1.1 \pm 0.6$ | $2.6 \pm 3.1 \pm 2.1$ |
| Upper limit at 90% C.L. (pb) | <2.5 | <8.7 |

Upper limits at the 90% CL are set using the Feldman-Cousins method [Phys. Rev. D 57, 3873 (1998)]

Tetraquark model in Ref. [CPC 43, 123102 (2019)]:

$$\Gamma(\Upsilon(10753) \rightarrow \eta_b(1S)\omega) = 2.64_{-1.69}^{+4.70} \text{ MeV}$$

$$\Gamma(\Upsilon(10753) \rightarrow \Upsilon\pi^+\pi^-) = 0.08_{-0.06}^{+0.20} \text{ MeV}$$

This measurement and JHEP 10, 220 (2019):

$$\sigma^B(\Upsilon(10753) \rightarrow \eta_b(1S)\omega) < 2.5 \text{ pb}$$

$$\sigma^B(\Upsilon(10753) \rightarrow \Upsilon(2S)\pi^+\pi^-) \approx (3 \pm 1) \text{ pb}$$

Our results do not support the prediction within the tetraquark model that the $\Upsilon(10753) \rightarrow \omega\eta_b(1S)$ decay is enhanced.

Quarkonium prospects at Belle II

Production of Charmonium(-like) states at B-factory

◆ B decay ($B \rightarrow KX_{c\bar{c}}$)

- ✓ CKM favored process, large branching fractions $10^{-3} \sim 10^{-4}$
- ✓ $J^{PC} = 0^{-+}, 1^{--}, 1^{++}, \dots$

◆ Initial-state radiation (ISR)

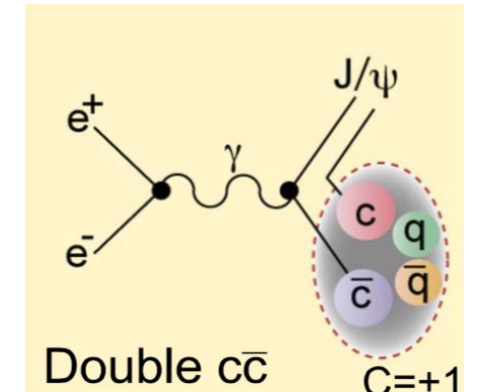
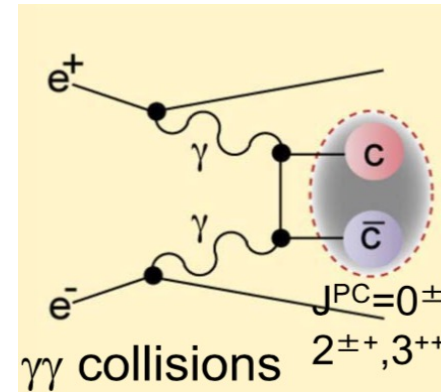
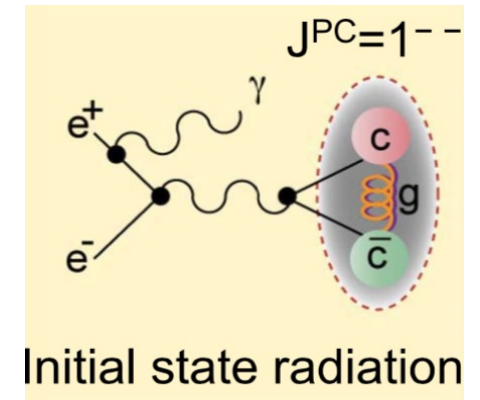
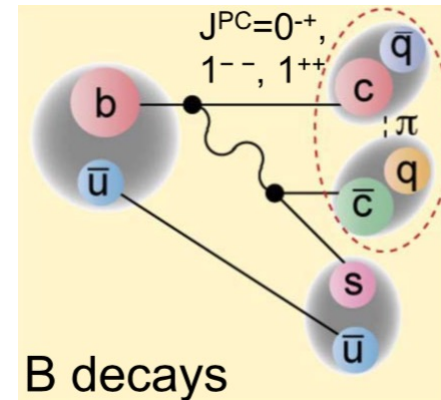
- ✓ $J^{PC} = 1^{--}$

◆ Two-photon process

- ✓ $J^{PC} = 0^{-+}, 0^{++}, 2^{++}, 2^{-+}, \dots$

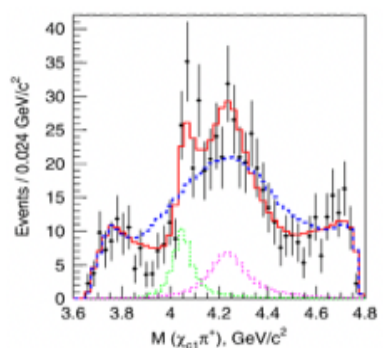
◆ Double charmonium

- ✓ e.g. $e^+e^- \rightarrow J/\psi X(3940)$ [PRL 98,082001(2007)]

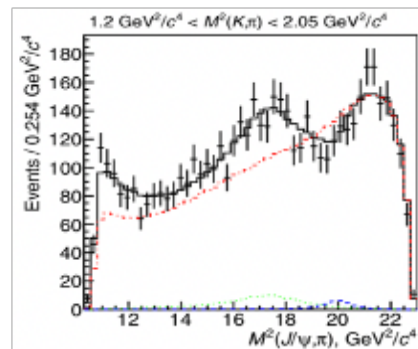


Charmonium(-like) states via B decays

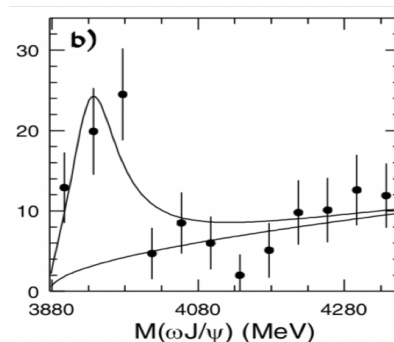
- Confirm Z_c states $Z(4050)^+$, $Z(4200)^+$, $Z(4250)^+$ and search for neutral partners
- Full amplitude analysis to $B \rightarrow K\omega J/\psi$ and $B \rightarrow K\omega\chi_{c1}$ to determine the spin-parities of $X(3915)$, $Z(4050)^+$ and $Z(4250)^+$.
- Confirmation of $X(3872)$ width measurement with $D^0\bar{D}^0\pi^0$ mode, search for more open-flavor decay modes, e.g., $B \rightarrow K(D\bar{D})$, $B \rightarrow K(D\bar{D}^*)$, $B \rightarrow K(D^*\bar{D}^*)$, $B \rightarrow K(D^*\bar{D}^{**})$... with more B mesons
- Absolute branching fractions are unique for Belle II
- Systematic investigations of charmonium plus light hadron final states: $B \rightarrow K(c\bar{c}+h)$



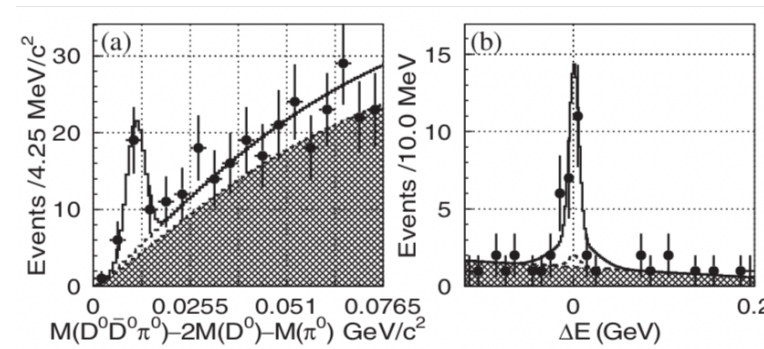
$Z(4050/4250) \rightarrow \pi^+\chi_{c1}$
PRD 78,072004



$Z(4200/4430) \rightarrow \pi^+J/\psi$
PRD 90,112009



$X(3915) \rightarrow \omega J/\psi$
PRL 94,182002



$B \rightarrow KX(3872) \rightarrow D^0\bar{D}^0\pi^0$
PRL 97,162002

Bottomonium(-like) prospects at Belle II

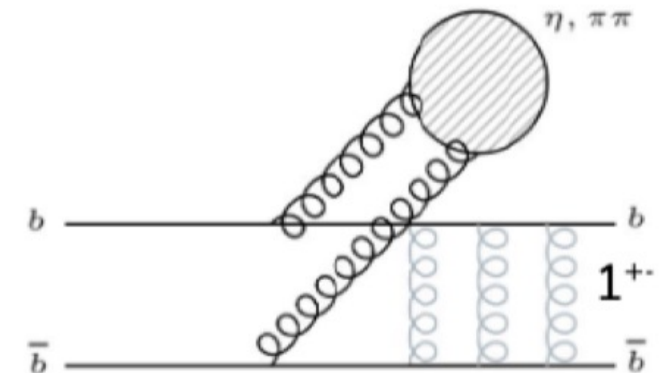
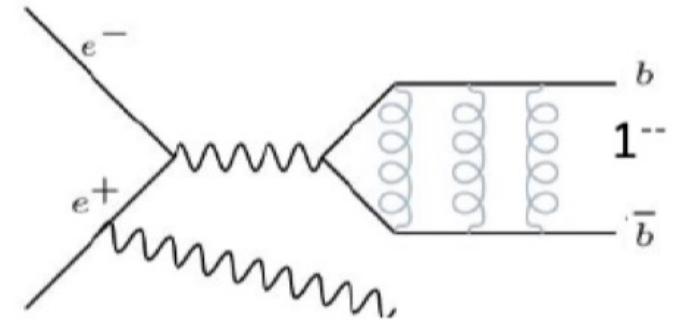
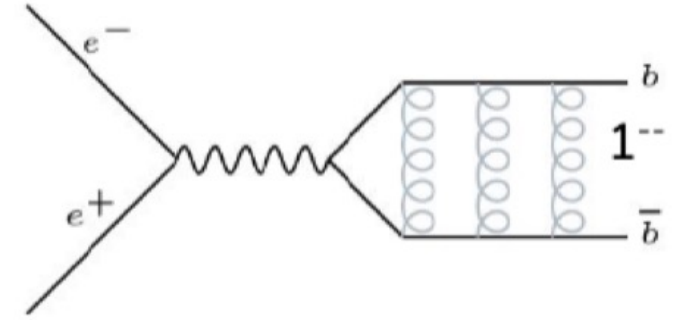
Four ways to access bottomonia:

- **Direct production from e^+e^-** : $J^{PC} = 1^{--}$: $\Upsilon(nS)$
- **ISR production**: $J^{PC} = 1^{--}$: $\Upsilon(nS)$
- **Hadronic transitions** from $\Upsilon(nS)$ through $\eta, \pi\pi, \dots$

$J^{PC} = 0^{-+}, 1^{--}, 1^{+-} \dots$: $\Upsilon(nS), \eta_b(nS), h_b(nS), \dots$

- **Radiative transitions** from $\Upsilon(nS)$

$J^{PC} = 0^{-+}, 0^{++}, 1^{++}, 2^{++}$: $\eta_b(nS), \chi_b(nP)$



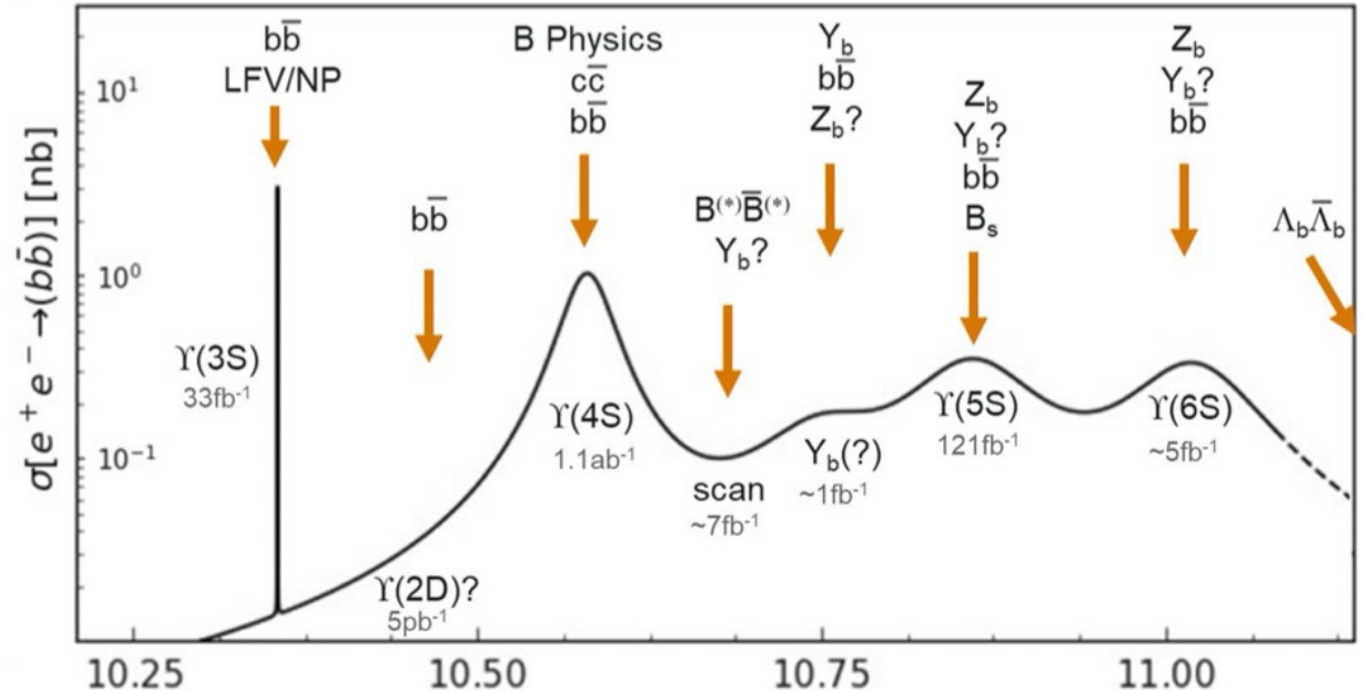
Bottomonium(-like) prospects at Belle II

Run at $Y(6S)$ and $Y(5S)$ and high energy scan:

- Search for new missing bottomonia $\eta_b(3S)$, $h_b(3P)$, $\Upsilon(D)$, exotic states Y_b , Z_b , etc
- Improve precision of already known processes and states, e.g., Z_b
- Measure the effect of the coupled channel contribution
- Study $B^{(*)}\bar{B}^{(**)}$ and $B_S^{(*)}B_S^{(**)}$ threshold regions (challenging for Super-KEKB)

Run at $Y(3S)$ and $Y(2S)$:

- Search for missing $\pi\pi/\eta$ transitions in inclusive decays to constrain further models
- Search for new physics: LFV, LFU, light Higgs, ...



Summary

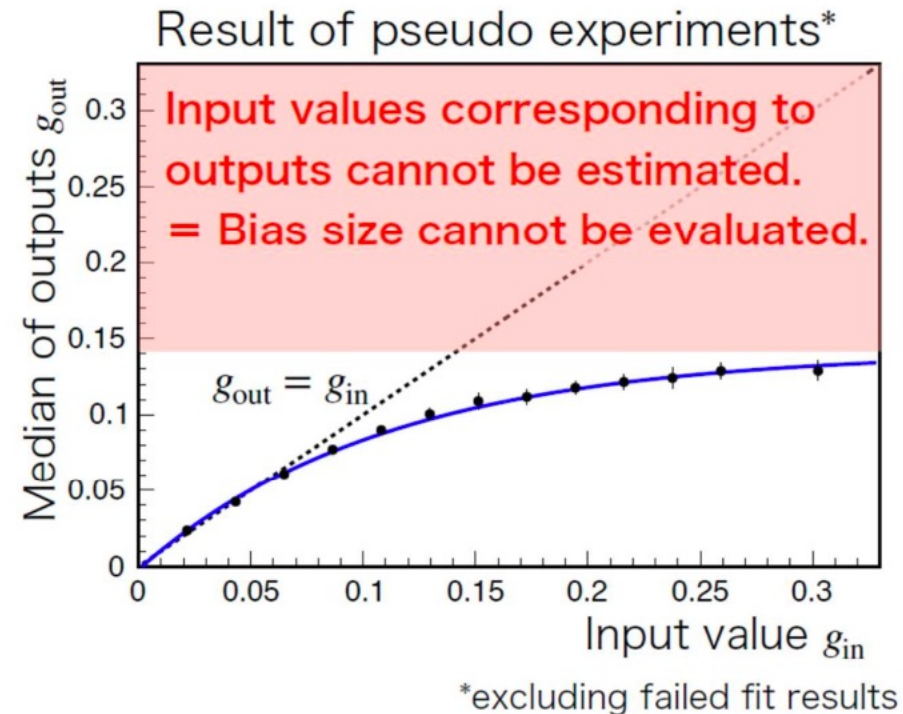
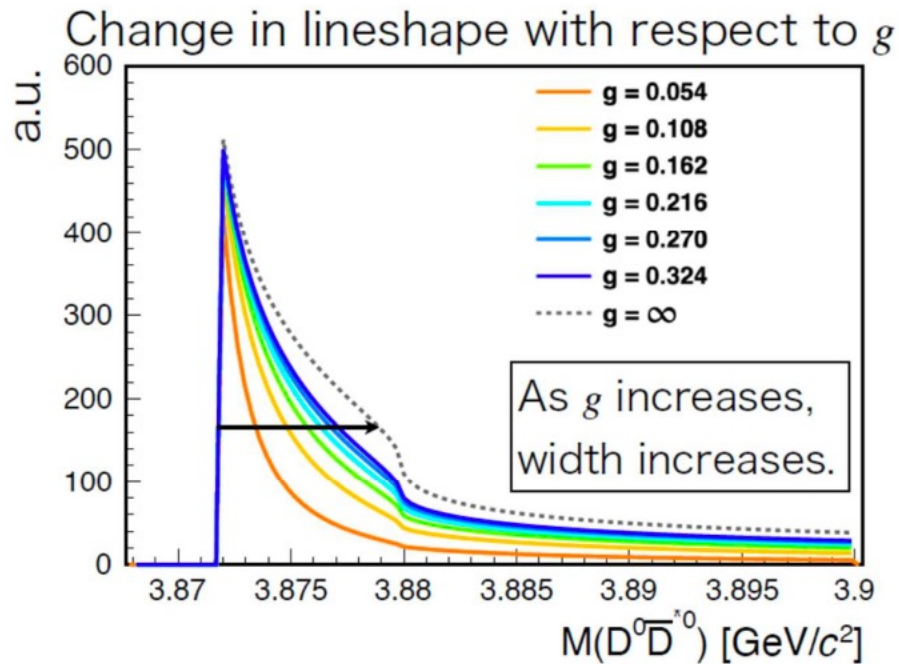
- The $g > 0.094$ (90% C.L.) is determined for $X(3872) \rightarrow D^0 \bar{D}^{*0}$
- Evidence for $e^+e^- \rightarrow \eta_c J/\psi$ near the threshold
- New decay modes of $\Upsilon(10753) \rightarrow \omega \chi_{bJ}$ are observed
- The rapid increase of $\sigma_{B^* \bar{B}^*}$ above the threshold may imply a resonance of $B^* \bar{B}^*$ or $b\bar{b}$
- The stringent upper limit is set for the $e^+e^- \rightarrow \omega \eta_b(1S)$ at $\sqrt{s} = 10.745$ GeV
- We expect more promising quarkonium results in the near future at Belle II with larger data samples

Thanks for your attention!

Backup slides

Fit bias

- Lineshape converges to a fixed form for large g



→ Only lower limit can be obtained for large g