

# Quasicrystal in QCD

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# Low Energy Dense QCD

Spontaneous Chiral Symmetry Breaking

$$U(1)_V \times SU(N_f)_L \times SU(N_f)_R \rightarrow U(1)_V \times SU(N_f)_V$$

Nambu Goldstone (NG) boson

$$\Sigma = \exp(iT^a \pi^a(x) / f_\pi) \in SU(N_f)$$

Chiral Perturbation Theory (ChPT):

$$\mathcal{L}_{\text{chiral}} = \frac{f_\pi^2}{4} \text{Tr}(\partial^\mu \Sigma \partial_\mu \Sigma) - \frac{b}{2} \text{Tr}[M(\Sigma - 1) + \text{h.c.}]$$

$N_f = 2$  case with approximately  $m_u \approx m_d$

$$M = \begin{pmatrix} m_u & 0 \\ 0 & m_d \end{pmatrix} \quad m_u \approx m_d \equiv m$$

# Under Magnetic Field

Pions  $\xrightarrow{B}$  Chiral Soliton Lattice (CSL) /  $\pi^0$  domain wall (DW)

LLL:  $n = 0$ , Pion: Spin  $s = 0$

$$\varepsilon^2 = p_z^2 + \boxed{2|eB|(n + 1/2) + m^2 - 2seB}$$

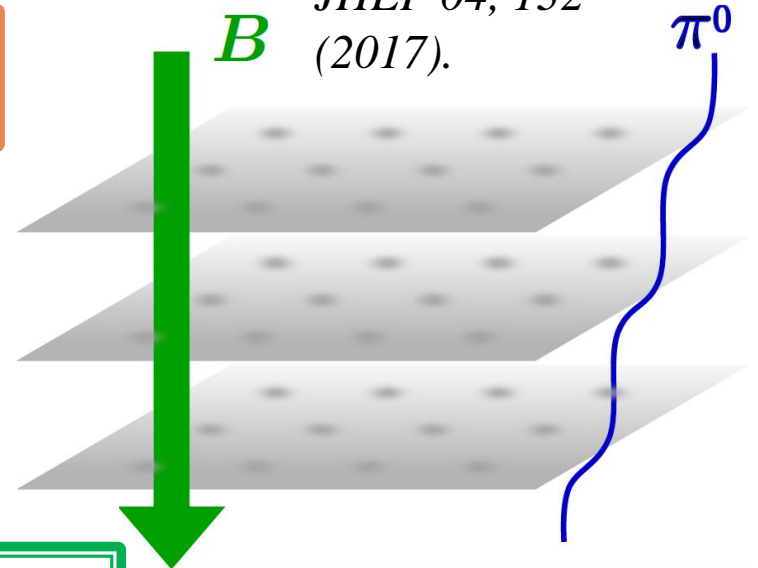
$\pi^\pm \xrightarrow{B\uparrow}$  “massive”,  $\Sigma \rightarrow \exp(i\tau^3 \pi^3)$

Transverse homogeneity  $\partial_{x,y} = 0$   
minimizes the Hamiltonian

$$\frac{f_\pi^2}{2} (\nabla \pi^0)^2 + m_\pi^2 f_\pi^2 (1 - \cos \pi^0) \boxed{-\frac{\mu \cdot B}{4\pi^2} \partial_z \pi^0}$$

Kinetic  $\partial_z$  + mass term  $\rightarrow$  SG soliton Winding  $\pi_1(U(1)) = \mathbb{Z}$

*Idea by T. Brauner  
and N. Yamamoto,  
JHEP 04, 132  
(2017).*



# WZW Term by Triangle Anomaly

Coupling of two U(1) via triangle anomaly to NG boson  $\phi_i$

$$\mathcal{L}_B = \frac{1}{8\pi^2} \epsilon^{\mu\nu\alpha\beta} \sum_i C_i \partial_\mu \phi_i A_\nu^B F_{\alpha\beta} \quad \text{D. T. Son, M. A. Stephanov, and A. R. Zhitnitsky, Phys. Rev. Lett. 86, 3955 (2001);}$$

Setup of chemical potential  $\mu$  and magnetic field  $\mathbf{B} = B\hat{z}$

$$U(1)_B : A_\nu^B = (\mu, \mathbf{0}). \quad U(1)_{\text{EM}} : A_\mu = (0, yB/2, -xB/2, 0)$$

$\phi_i$  be not only  $\pi^{\pm,0}$  but also  $\eta$  mesons (for larger density)

$$\phi_3 = \frac{\pi_3}{f_\pi} : \quad \mathcal{L}_{\text{WZW}} = \frac{\mu B}{4\pi^2} \partial_z \phi_3 \Rightarrow \text{CSL} \quad \text{D. T. Son and A. R. Zhitnitsky, Phys. Rev. D 70, 074018 (2004).}$$

$$\phi_0 \equiv \frac{\eta}{f_\eta} : \quad \mathcal{L}_\eta = \frac{\mu B}{12\pi^2} \partial_z \phi_0 \Rightarrow \eta\text{-CSL?}$$

Separately equivalent.  
But mix  $\eta$  and  $\pi^0$  ...  
new ground state?

# U(2) ChPT in Magnetic Field

$\pi^\pm$  decoupled,  $\eta$  and  $\pi^0$  remained:  $U = \exp \phi_0 \exp (i\tau_3 \phi_3)$   
Kinetic Anomaly

$$H = \frac{1}{2} \left[ \alpha (\partial_z \phi_3)^2 + (\partial_z \phi_0)^2 \right] - \frac{\gamma}{2\pi} \left( \partial_z \phi_3 + \frac{1}{3} \partial_z \phi_0 \right) + \sin \beta (1 - \cos 2\phi_0) + \cos \beta (1 - \cos \phi_0 \cos \phi_3) \text{ Mass}$$

Redefine parameters to have dimensionless quantities

$$\alpha \equiv \frac{f_\pi^2}{f_\eta^2}, \quad \beta \equiv \arctan \frac{a}{2mb}, \quad \gamma \equiv \frac{\mu B \left[ a^2 + (4mb)^2 \right]^{1/4}}{2\pi f_\eta}$$

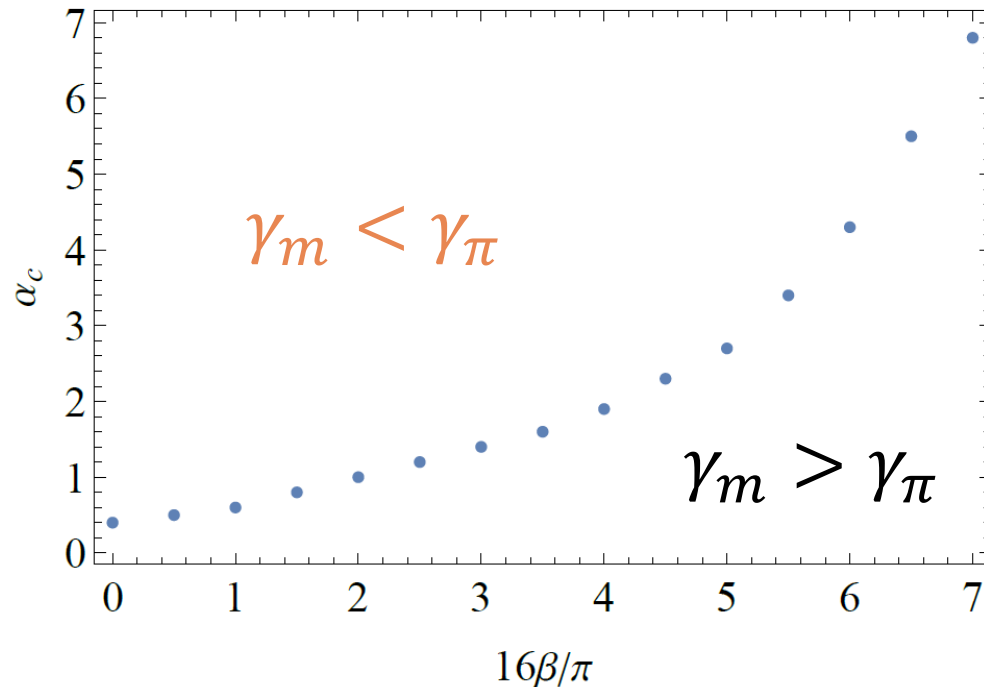
Boundary conditions to form Mixed Soliton Lattice

$$(\phi_3, \phi_0) \Big|_{z=0} = (0, 0), \quad (\phi_3, \phi_0) \Big|_{z=d} = (p\pi, q\pi), \quad \frac{p \pm q}{2} \in \mathbb{Z}$$

# Mixed Soliton Lattice

Irrelevant  $\gamma_\eta \gg \gamma_{\pi,m}$ . Duel is between  $\gamma_\pi$  and  $\gamma_m$ .

$(p,q)$	Soliton Type	Critical magnetic field
$(2,0)$	$\pi^0$	$\gamma_\pi$
$(0,2)$	$\eta$	$\gamma_\eta$
$(1,1)$	Mixed	$\gamma_m$



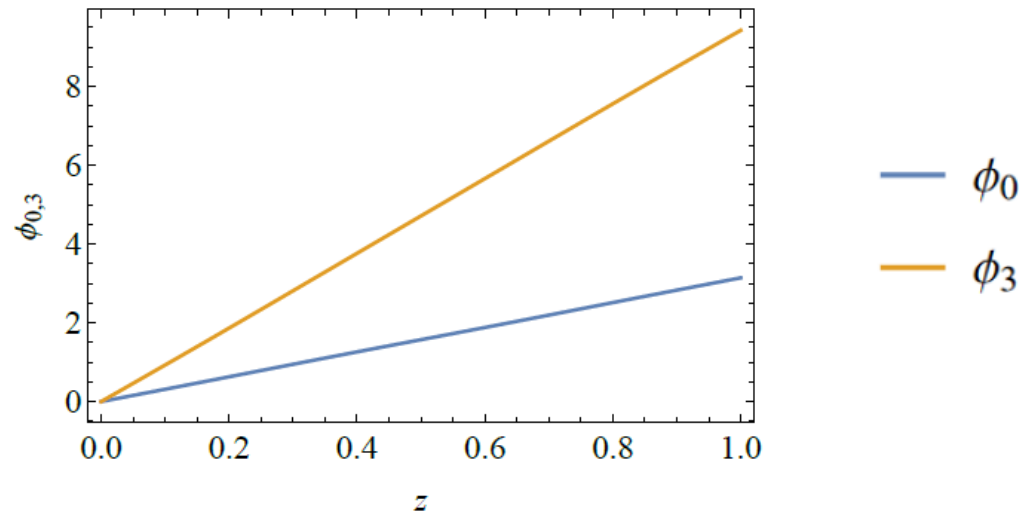
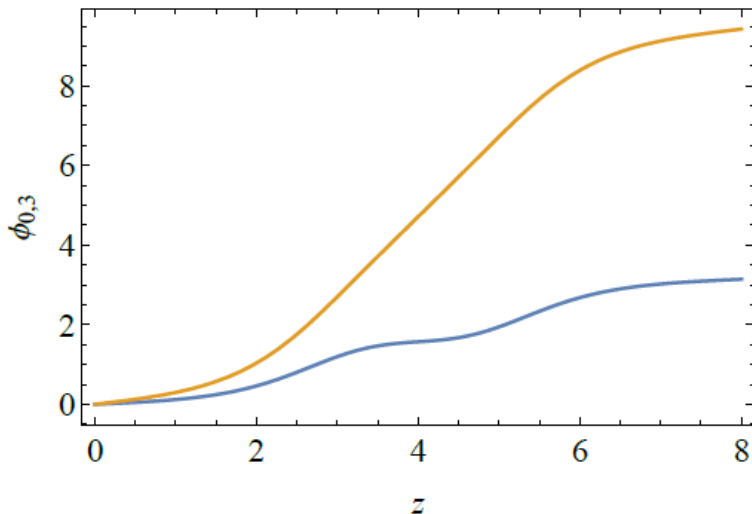
$\gamma_m$  can be lowest! for  $\alpha > \alpha_c(\beta)$

# Strong Magnetic Field Limit

In  $\gamma \rightarrow \infty$  limit, the ground state tends to take

$$r \equiv \frac{p}{q} = \frac{3}{\alpha}$$

Profiles:  $\gamma \uparrow$ ,  $d \downarrow$ ,  $\phi_{3(0)} \rightarrow p(q)\pi z/d$ ; (almost linear)



Indicate: derivative is large, mass term is subleading.

# Semi-analytical Proof

$\gamma \rightarrow \infty$ , mass term becomes irrelevant:

$$H(\gamma \rightarrow \infty) \simeq \frac{1}{2} (\alpha \phi_3'^2 + \phi_0'^2) - \frac{\gamma}{2\pi} \left( \phi_3' + \frac{1}{3} \phi_0' \right) \equiv H_\infty.$$

Simple trick of total square finds us the minimum:

$$\begin{aligned} H_\infty &= \frac{1}{2} \left[ \left( \sqrt{\alpha} \phi_3' - \frac{\gamma}{2\sqrt{\alpha}\pi} \right)^2 + \left( \phi_0' - \frac{\gamma}{6\pi} \right)^2 \right] - \frac{\gamma^2}{8\pi^2} \left( \frac{1}{\alpha} + \frac{1}{9} \right) \\ &\geq -\frac{\gamma^2}{8\pi^2} \left( \frac{1}{\alpha} + \frac{1}{9} \right) \equiv E_{\min}. \end{aligned}$$

Lattice period  $d_L$  and ratio  $p/q$  that minimize the energy

$$d_L = \frac{2\pi^2}{\gamma} \cdot p\alpha = \frac{2\pi^2}{\gamma} \cdot 3q \quad \Rightarrow \quad \frac{p}{q} = \frac{3}{\alpha},$$



# Quasicrystal in QCD

Grey dots: mixed soliton lattice of different  $p, q$ .

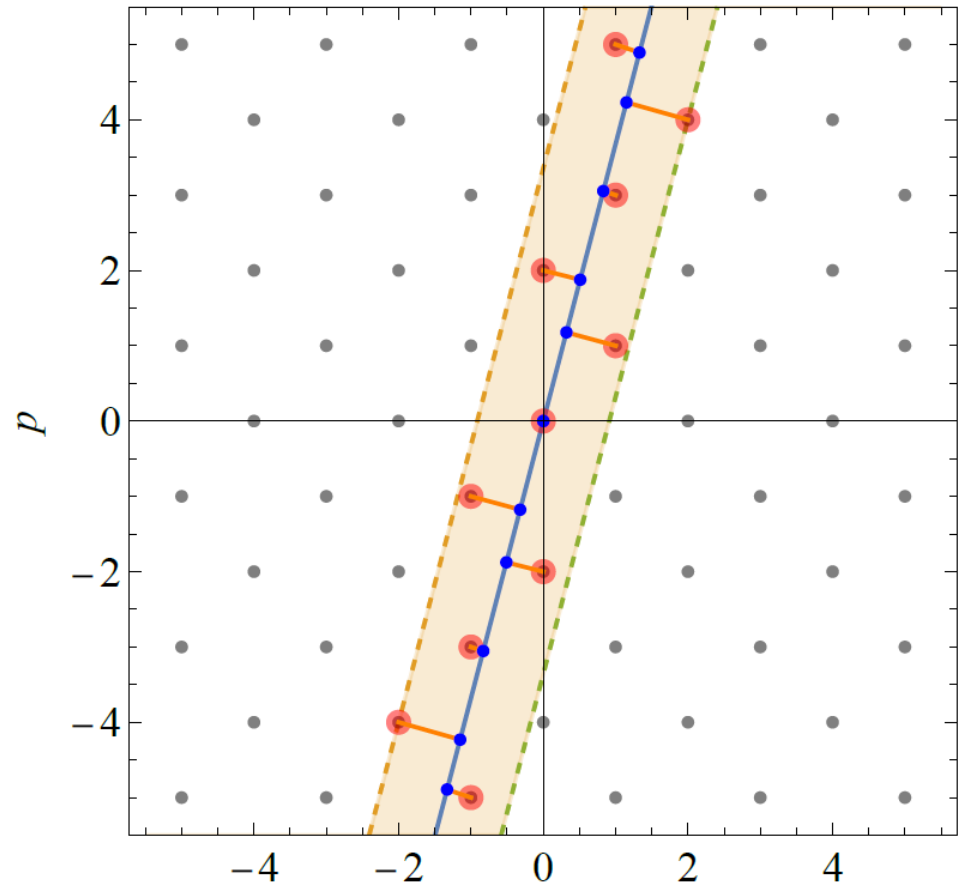
Blue solid line (example:  $p/q = 3f_\eta^2 / f_\pi^2$  (irrational))

Orange dots:  $(p, q)$  that adjacent to the blue line.

Beige band: possibly degenerate ground states, forming **Quasicrystal!**

\*valid range:

$$\mu B \ll \Lambda^3, \quad \Lambda \simeq 4\pi f_{\pi, \eta}; \quad \gamma \ll 5.6\pi^2$$



*Zebin Qiu, Muneto Nitta,  
JHEP 05, 170 (2023)*

# Rotational Counterpart

In parallel to WZW term, anomalous rotational ( $\mathbf{\Omega} = \Omega \hat{z}$ ) Lagrangian is the coupling  $j_{\text{CVE}}^5$  to axial rotation.

$$\mathcal{L}_\eta^\Omega = \frac{\mu^2 \Omega}{6\pi^2} \partial_z \phi_0, \quad \mathcal{L}_\pi^\Omega = \frac{\mu \mu_I \Omega}{2\pi^2} \partial_z \phi_3$$

*X.-G. Huang, K. Nishimura and N. Yamamoto, JHEP 02 (2018) 069*

We consider no isospin chemical potential  $\mu_I = 0$ . Recap:

	$\eta$	$\pi^0$
Rotation $\mathbf{\Omega}$	$(\partial_z \phi_0) \mu^2 \Omega / 6\pi^2$	0
Magnetic field $B$	$(\partial_z \phi_0) \mu B / 12\pi^2$	$(\partial_z \phi_3) \mu B / 4\pi^2$

In rotational case, the net winding number must be from  $\eta$  only. But  $\pi^0$  can have local “up and down”

# Non-Abelian (NA) CSL (Rotation)

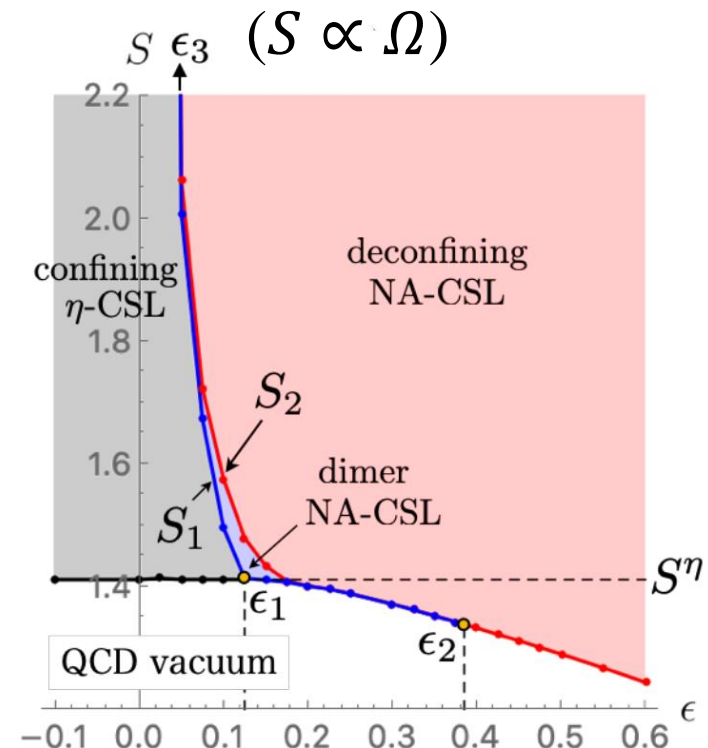
$\eta$ -CSL can split into “dimer”, i.e., “Up”  $\oplus$  “Down”

$$(\phi_3, \phi_0) \Big|_{z=0} = (0, 0);$$

$$(\phi_3, \phi_0) \Big|_{z=l} = \begin{cases} (0, 2\pi) & \eta\text{-CSL} \\ (\pi, \pi) & \text{”Up” branch} \\ (-\pi, \pi) & \text{”Down” branch} \end{cases}$$

Each branch is a “mixed soliton”  
 The interaction between branches  
 1. attractive:  $\eta$ -CSL; 2. null:  
 dimer; 3. repulsive: NA-CSL

$$\mu^2 \Omega_c \sim 1.3 \cdot N_c 2\pi^2 f_\eta^2 m_\eta$$



$$(\epsilon = 1 - f_\pi^2 / f_\eta^2 \sim 0.3)$$

*M. Eto, K. Nishimura and M. Nitta, JHEP 08 (2022) 305*

# Conclusion

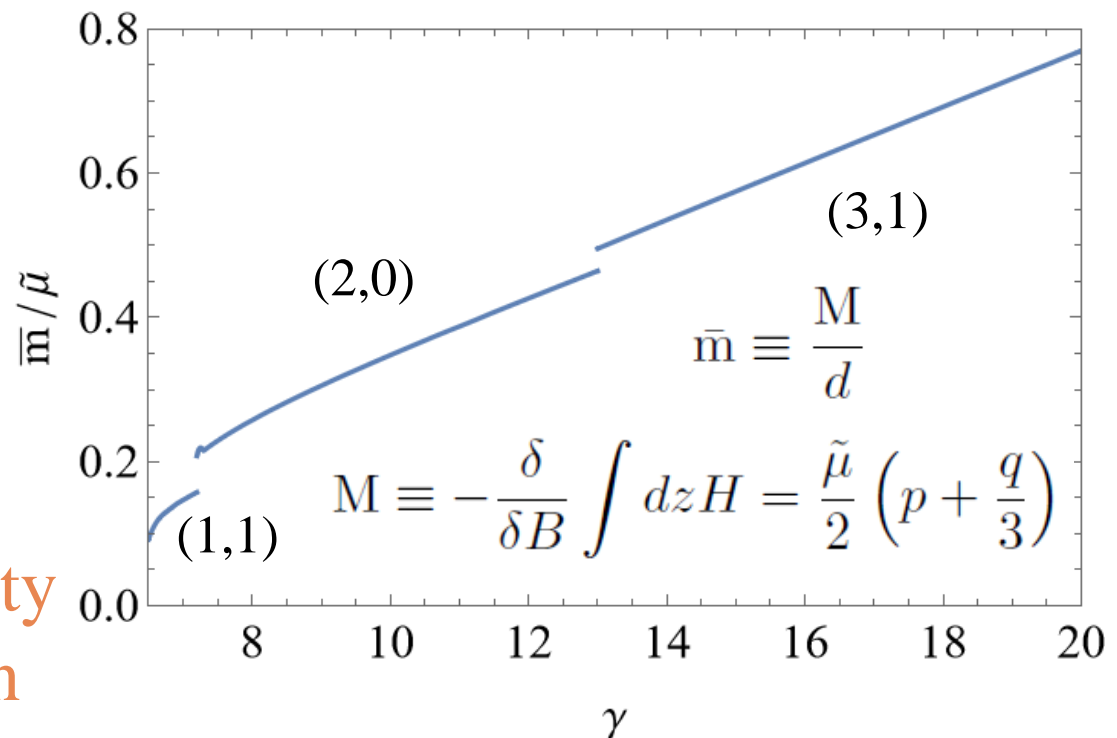
- Depending on decay constants and effective masses, the **mixed soliton lattice** of  $\eta$  and  $\pi^0$  can have lower energy / critical magnetic field than separate ones.
- In strong magnetic field / density limit, the ground state ratio  $p/q$  approaches  $3f_\eta^2 / f_\pi^2$  which is generally irrational. The result is a **mesonic quasicrystal**.
- In rotation case, the  $\eta$ -CSL can split into dimer and then **NA-CSL**, during increasing the angular velocity.

# Outlook: Intermediate $B$

Example: up to  $p+q=4$ , under  $\alpha = 0.7$  and  $\beta = \pi/16$ :  
 competitive configurations are (1,1), (2,0), and (3,1)  
 with critical  $\gamma_m = 6.5$ ,  $\gamma_\pi = 6.7$  and  $\gamma_{31} = 7.3$ .

**Ground  
State**

$$= \begin{cases} (1,1) & \gamma \in [6.4, 7.2) \\ (2,0) & \gamma \in [7.2, 13.0) \\ (3,1) & \gamma \in (13.0, \dots) \end{cases}$$



Magnetic moment density  
 quantifies the alternation  
 as a piece-wise function of  $\mu B$