CEPC味物理-新物理和相关探测技术研讨会 2023年8月16日-18日,复旦大学

Some theoretical aspects of hadronic τ decays



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Overview



Z exchange dominates at CEPC.

Number of taus produced at e⁺e⁻ colliders:

ALEPH: $\sim 3 \times 10^5$	BaBar / Belle: $\sim 1 \times 10^9$
Belle-II: $\sim 5 \times 10^{10}$	CEPC (Tera-Z factory): $\sim 3 \times 10^{10}$

Tau provides broad interests for particle physics:

- ✓ Precision tests for electroweak sector: V_{CKM} , lepton universality, g-2,
- \checkmark Stong interactions: α_s , hadron resonances, chiral symmetry,
- ✓ Possible discoveries for new physics: cLFV, CPV,

Sketch for hadronic tau decays (similar for leptonic decays by dropping QCD part)



Theoretical tools: SM EFT + Chiral EFT

- $$\begin{split} \mathbf{SM} \ \mathbf{EFT} & \rightarrow \mathbf{LEFT} \qquad \qquad [Cirigliano\ \mathbf{et}\ \mathbf{al},\ \mathbf{'10}] \dots \dots \\ \mathcal{L}_{\mathrm{eff}} &= -\frac{G_{\mu}V_{uD}}{\sqrt{2}} \bigg[\Big(1 + \epsilon_{L}^{D\ell} \Big) \bar{\ell} \gamma_{\mu} (1 \gamma_{5}) \nu_{\ell} \cdot \bar{u} \gamma^{\mu} (1 \gamma_{5}) D + \epsilon_{R}^{D\ell} \bar{\ell} \gamma_{\mu} (1 \gamma_{5}) \nu_{\ell} \cdot \bar{u} \gamma^{\mu} (1 + \gamma_{5}) D \\ & \quad + \bar{\ell} (1 \gamma_{5}) \nu_{\ell} \cdot \bar{u} \Big[\epsilon_{S}^{D\ell} \epsilon_{P}^{D\ell} \gamma_{5} \Big] D + \frac{1}{4} \hat{\epsilon}_{T}^{D\ell} \bar{\ell} \sigma_{\mu\nu} (1 \gamma_{5}) \nu_{\ell} \cdot \bar{u} \sigma^{\mu\nu} (1 \gamma_{5}) D \bigg] + \mathrm{h.c.}, \\ \mathbf{\epsilon_{X}} \ \mathbf{parameterize} \ \mathbf{various} \ \mathbf{new} \ \mathbf{physics} \ \mathbf{at} \ \mathbf{high} \ \mathbf{energy} \ \mathbf{scale} \end{split}$$
 - Chiral EFT $\mathcal{L}_{2} = \frac{F^{2}}{4} \langle u_{\mu}u^{\mu} + \chi_{+} \rangle$ $\begin{bmatrix} Gasser, Leutwyler, '83 '84 \end{bmatrix}$ $\mathcal{L}_{4}^{\chi PT} = L_{1} \langle u_{\mu}u^{\mu} \rangle^{2} + L_{2} \langle u_{\mu}u^{\nu} \rangle \langle u^{\mu}u_{\nu} \rangle + L_{3} \langle u_{\mu}u^{\mu}u_{\nu}u^{\nu} \rangle + L_{4} \langle u_{\mu}u^{\mu} \rangle \langle \chi_{+} \rangle$ $+L_{5} \langle u_{\mu}u^{\mu}\chi_{+} \rangle + L_{6} \langle \chi_{+} \rangle^{2} + L_{7} \langle \chi_{-} \rangle^{2} + \frac{L_{8}}{2} \langle \chi_{+}^{2} + \chi_{-}^{2} \rangle + \cdots$

Hadronic decays: a unique feature for tau lepton



Strong coupling of QCD: α_s





Hadronic tau decay: an invaluable source to test the QCD prediction of $\alpha_s(Q^2)$ below 2 GeV.

PDG

$$R_{\tau} = \frac{\Gamma(\tau \rightarrow \nu_{\tau} + \text{hadrons})}{\Gamma(\tau \rightarrow \nu_{\tau} e \bar{\nu}_{e})}$$

$$\Gamma_{\tau \rightarrow \nu_{\tau} + \text{had}} \sim \text{Im} \left\{ \begin{array}{c} & & \\ & &$$

$$a_{s}(\mathbf{m}_{\tau}) \text{ from spectral functions} \qquad [\text{Bratten et al., '92]} \qquad [\text{Im}(s) \\ R_{\tau} = \frac{\Gamma(\tau \to \nu_{\tau} \text{ hadrons})}{\Gamma(\tau \to \nu_{\tau} e \nu_{e})} = \int_{0}^{m_{\tau}^{2}} \frac{ds}{m_{\tau}^{2}} \omega_{J}(s) \text{ Im}\Pi^{J}(s) \qquad [\mathbf{m}_{\tau}^{2}] \\ R_{\tau} = \oint_{|s|=m_{\tau}^{2}} \frac{ds}{m_{\tau}^{2}} \omega_{J}(s) \Pi^{J}(s) \qquad [\mathbf{m}_{\tau}^{2}] \\ \text{Along } |s|=m_{\tau}^{2}, \qquad \Pi^{J}(s) = \underbrace{\text{OPE}}_{D} \sum_{D} \frac{C_{D}^{J}(s, \alpha_{s}(\mu), \mu) \langle O_{D}(\mu) \rangle}{(-s)^{D/2}} \\ R_{\tau} = N_{C} S_{\text{EW}} \left(1 + \delta_{\text{P}} + \delta_{\text{NP}}\right) = R_{\tau,V} + R_{\tau,A} + R_{\tau,S} \\ S_{\text{EW}} = 1.0201 (3) \qquad ; \qquad \delta_{\text{NP}} = -0.0064 \pm 0.0013 \\ \text{Marciano-Sirlin, Braaten-Li, Erler} \qquad Fitted from data (Davier et al)} \\ \delta_{p} = a_{\tau} + 5.20 a_{\tau}^{2} + 26 a_{\tau}^{3} + 127 a_{\tau}^{4} + ... \approx 20\% \qquad ; \qquad a_{\tau} \equiv \alpha_{s}(m_{\tau})/\pi \\ \text{Baikov-Chetyrkin-Kühn} \\ Precision limited by FOPT vs CIPT discrepancy: \qquad Duality violation issue: \end{cases}$$

 $\alpha_s(m_\tau^2) = 0.319 \pm 0.014$ $\alpha_s(m_\tau^2) = 0.341 \pm 0.013$

Boito et al., Pich et al.,

Invariant-mass spectra for exclusive decays









s (GeV2)



[Davier et al., '08]

Exclusive case

$$\begin{split} \frac{\mathcal{B}(\tau^- \to K^- \nu_{\tau})}{\mathcal{B}(\tau^- \to \pi^- \nu_{\tau})} &= \frac{f_{K\pm}^2 |V_{us}|^2}{f_{\pi\pm}^2 |V_{ud}|^2} \frac{(m_{\tau}^2 - m_K^2)^2}{(m_{\tau}^2 - m_{\pi}^2)^2} (1 + \delta R_{\tau K/\tau \pi}) & \qquad \text{[Arroyo-Urena et al., '21 '22]} \\ \mathcal{B}(\tau^- \to K^- \nu_{\tau}) &= \frac{G_F^2}{16\pi\hbar} f_{K\pm}^2 |V_{us}|^2 \tau_{\tau} m_{\tau}^3 \left(1 - \frac{m_K^2}{m_{\tau}^2}\right)^2 S_{\rm EW}(1 + \delta R_{\tau K}) \end{split}$$

"Tension" of V_{us} from various determinations



Beyond SM tests in tau physics

> Lepton Universality (LU) test

$$\begin{array}{c} \begin{array}{c} & & \\$$

> Relevance to precise determination of a_{μ}

[Muon g-2, '23]

SM uncertainty dominated by



Dominated by $\pi\pi$ (> ~75%)









* Key problem in the matching: isospin breaking (IB) effects IB corrections to a_{μ} [Cirigliano et al., JHEP'02]





TABLE IV. Contributions to $\Delta a_{\mu}^{\text{HVP,LO}}$ in units of 10^{-11} using the dispersive representation of the form factor. From the two evaluations labeled $\mathcal{O}(p^4)$, the left (right) one corresponds to $F_V = \sqrt{2}F$ ($F_V = \sqrt{3}F$).

$[s_1, s_2]$	$\Delta a^{ m HVP,LO}_{\mu,{ m G}^{(0)}_{ m EM}}$	$\Delta a_{\mu,{ m SI}}^{ m HVP,LO}$	$\Delta a_{\mu,[\mathcal{O}(p^4)]}^{ ext{HVP,LO}}$	$\Delta a^{ ext{HVP,LO}}_{\mu,[\mathcal{O}(p^4)]}$	$\Delta a^{\mathrm{HVP,LO}}_{\mu,[SD]}$	$\Delta a^{\mathrm{HVP,LO}}_{\mu,[\mathcal{O}(p^6)]}$
$[4m_{\pi}^2, 1 \text{ GeV}^2]$	+17.8	-11.0	-11.3	-17.0	-32.4	-74.8 ± 44.0
$[4m_{\pi}^2, 2 \text{ GeV}^2]$	+18.3	-10.1	-10.3	-16.0	-31.9	-75.9 ± 45.5
$[4m_{\pi}^2, 3 \text{ GeV}^2]$	+18.4	-10.0	-10.2	-15.9	-31.9	-75.9 ± 45.6
$[4m_{\pi}^2, m_{\tau}^2]$	+18.4	-10.0	-10.2	-15.9	-31.9	-75.9 ± 45.6

Referenced value using the tau data to calculate a_{μ}

$$\Delta a_{\mu} = a_{\mu}^{\exp} - a_{\mu}^{\th} = (12.5 \pm 6.0) \times 10^{-10}$$

> CP violation in tau decays

$$A_{CP} = \frac{\Gamma(\tau^- \to \nu_\tau H) - \Gamma(\tau^+ \to \nu_\tau \bar{H})}{\Gamma(\tau^- \to \nu_\tau H) + \Gamma(\tau^+ \to \nu_\tau \bar{H})}$$

Intensive discussions on tau -> Ks pi nu

$$A_{\varrho} = \frac{\Gamma\left(\tau^{+} \to \pi^{+} K_{S}^{0} \overline{\nu}_{\tau}\right) - \Gamma\left(\tau^{-} \to \pi^{-} K_{S}^{0} \nu_{\tau}\right)}{\Gamma\left(\tau^{+} \to \pi^{+} K_{S}^{0} \overline{\nu}_{\tau}\right) + \Gamma\left(\tau^{-} \to \pi^{-} K_{S}^{0} \nu_{\tau}\right)}$$
$$\approx (0.36 \pm 0.01)\% \qquad (-0.36 \pm 0.23_{\text{stat}} \pm 0.11_{\text{syst}})\%$$

SM prediction

BaBar

[Bigi et al., PLB'05] [Grossman et al., JHEP'12] [Lees et al., PRD'12]
[Cirigliano et al., PRL'18] [Rendo et al., PRD'19] [Chen et al., PRD'19 JHEP'20]

Other types of CPV observables: T-odd triple-product asymmety

A typical T-odd kinematical variable:

$$\xi \equiv \varepsilon_{\mu\nu\rho\sigma} a^{\mu} b^{\nu} c^{\rho} d^{\sigma} \xrightarrow{\text{rest frame}}_{\text{of particle } a} \vec{b} \cdot (\vec{c} \times \vec{d}) m_a / s_a$$

a, b, c, d: either momentum or spin

T transformation $(t \to -t, \vec{p} \to -\vec{p}, \cdots): \bar{\xi} \to -\xi$

- When spin is involved, measurement of polarization is needed.
 [Nelson, et al., PRD'94] [Tsai, PRD'95] [Datta, PRD'07] ...
- ***** When focusing on the situation with four momenta, *i.e.*

$$\xi = \varepsilon_{\mu\nu\rho\sigma} p_1^{\mu} p_2^{\nu} p_3^{\rho} p_4^{\sigma} \xrightarrow{\text{rest frame}}_{\text{of particle 1}} \vec{p_2} \cdot (\vec{p_3} \times \vec{p_4}) m_1$$

In this case, there should be at least four particles in the final state!

> Pro: Strong phase is not necessary for a CPV phenomenon using TPA.

Con: TPA could also be caused by the final-state interactions!

Predicitons of the T-odd asymmetry distribution in $\tau \rightarrow \pi \pi \gamma v_{\tau}$



Variant predictions for the branching ratios (×10-4):

$E_{\gamma}^{\mathrm{cut}}$	\mathbf{SI}	CEN	Our-1A	Our-2A	Our-1B	Our-2B
$100 \mathrm{MeV}$	7.9	8.3	8.7/9.6/8.6/9.4	9.5/10/9.2/9.7	13/9.6/12/9.4	14/10/13/9.7
$300 \mathrm{MeV}$	1.5	1.8	2.4/3.0/2.3/2.8	2.9/3.3/2.6/3.0	5.6/3.0/5.2/2.8	6.3/3.3/5.5/3.0
$500 \mathrm{MeV}$	0.26	0.40	0.73/1.0/0.68/0.90	0.93/1.1/0.81/0.91	2.6/1.0/2.4/0.90	2.9/1.1/2.4/0.91

It has the good chance to be measured in STCF.

Prospects of revealing the genuine CPV signals

• CPV signals can be probed by taking the differences of A_{ξ} in $\tau \to \pi - \pi^0 \gamma v_{\tau}$ and $\tau^+ \to \pi^+ \pi^0 \gamma v_{\tau}$

$$A_{\xi} = \frac{\Gamma_{+} - \Gamma_{-}}{\Gamma_{+} + \Gamma_{-}} \qquad \qquad \overline{A}_{\overline{\xi}} = \frac{\overline{\Gamma}_{+} - \overline{\Gamma}_{-}}{\overline{\Gamma}_{+} + \overline{\Gamma}_{-}}$$

$$\overline{\Gamma}_{+} = \frac{(2\pi)^4}{2m_{\tau}} \int_{\overline{\xi}>0} \mathrm{d}\Phi \left(\overline{\hat{M}}_0 + \overline{\xi}\overline{\hat{M}}_1\right), \qquad \overline{\Gamma}_{-} = \frac{(2\pi)^4}{2m_{\tau}} \int_{\overline{\xi}<0} \mathrm{d}\Phi \left(\overline{\hat{M}}_0 + \overline{\xi}\overline{\hat{M}}_1\right)$$

$$\mathcal{M} = e \, G_F \, V_{ud}^* \epsilon^{*\mu}(k) \left\{ \left(1 + \mathbf{g}_{\mathbf{V}} \right) F_{\nu} \bar{u}(q) \gamma^{\nu} (1 - \gamma_5) (m_{\tau} + \not\!\!\!P - \not\!\!\!k) \gamma_{\mu} u(P) \right. \\ \left. + \left[\left(1 + \mathbf{g}_{\mathbf{V}} \right) V_{\mu\nu} - \left(1 - \mathbf{g}_{\mathbf{A}} \right) A_{\mu\nu} \right] \bar{u}(q) \gamma^{\nu} (1 - \gamma_5) u(P) \right\}$$

 $\mathcal{A}_{\xi} = A_{\xi} - \overline{A}_{\bar{\xi}} \supset \operatorname{Im}(\mathbf{g}_{\mathbf{V}}^* \mathbf{g}_{\mathbf{A}}) \operatorname{Re}[F_V(t/u)^* A_i], \ \operatorname{Im}(\mathbf{g}_{\mathbf{V}}^* \mathbf{g}_{\mathbf{A}}) \operatorname{Re}(V_j^* A_i)$

- Generally speaking, sizable hadronic contributions are also expected to enhance the CPV signals in $\tau \to \pi \pi \gamma v_{\tau}$.
- TPA in other types of τ decays could be also possible.

> Charged lepton flavor violation in tau decays

90% C.L. upper limits for LFV τ decays



- Not only statistic but also systematic uncertainties are important in $\tau \rightarrow l \gamma$
- Clean backgroud makes τ → l l'l" one of the best channels to search for LFV signals.
- $\tau \rightarrow l + hadrons$ provides a different laboratory to probe different LFV origins, comparing with the pure leptonic processes.

- > Proposals to search for second-class currents in tau decays
- First class of hadron currents : J^{PG}= 0⁺⁺, 0⁻⁻, 1⁺⁻, 1⁻⁺
- Second class of hadron currents : J^{PG}= 0⁺⁻, 0⁻⁺, 1⁺⁺, 1⁻⁻, which are usually suppressed at the level of 10⁻⁶ ~ 10⁻⁵ and are not observed yet in EXP.



BR: (0.3~2.0) ×10⁻⁵

BR: $10^{-9} \sim 10^{-6}$

[Escribano et al., PRD'16]

> Powerful tool to constrain new physics: combination of hadronic tau data + LHC data

$$\mathcal{L}_{\text{eff}} = -\frac{G_{\mu}V_{uD}}{\sqrt{2}} \left[\left(1 + \epsilon_L^{D\ell} \right) \bar{\ell} \gamma_{\mu} (1 - \gamma_5) \nu_{\ell} \cdot \bar{u} \gamma^{\mu} (1 - \gamma_5) D + \epsilon_R^{D\ell} \bar{\ell} \gamma_{\mu} (1 - \gamma_5) \nu_{\ell} \cdot \bar{u} \gamma^{\mu} (1 + \gamma_5) D \right. \\ \left. + \bar{\ell} (1 - \gamma_5) \nu_{\ell} \cdot \bar{u} \left[\epsilon_S^{D\ell} - \epsilon_P^{D\ell} \gamma_5 \right] D + \frac{1}{4} \hat{\epsilon}_T^{D\ell} \bar{\ell} \sigma_{\mu\nu} (1 - \gamma_5) \nu_{\ell} \cdot \bar{u} \sigma^{\mu\nu} (1 - \gamma_5) D \right] + \text{h.c.},$$
[Cirigliano et al., PRL'18]



Summary

Tau offers a laboratory for a broad range of interesting topics:

- > Precision tests of SM: CKM, α_s , m_{τ} , lepton universality,
- Hadron interactions: light-flavor resonances, chiral symmetry, form factors, second-class currents,
- **BSM tests:**

CPV (rate asym., triple-product asym.)

LFV (lepton/radiative decays, hadron decays)

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Thanks for your patience!