



河南省科学院
HENAN ACADEMY OF SCIENCES



復旦大學

Revisiting Vacuum Stability and CEPC Constraints on the Georgi-Machacek Model

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8/18/2023

Based on Xk Du, Zhuang Li, Fei Wang, Yingkai Zhang, 2204.05760;

Xk Du, Fei Wang, 2308.****

2023年8月18日 @ CEPC味物理-新物理和相关探测技术研讨会



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1、Introduction to the Georgi-Machacek Model

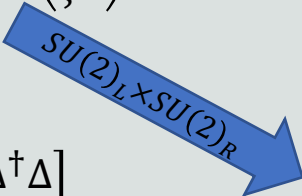
The Georgi-Machacek Model

$$\phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}, \quad \chi = \begin{pmatrix} \chi^{++} \\ \chi^+ \\ \chi^0 \end{pmatrix}, \quad \xi = \begin{pmatrix} \xi^+ \\ \xi^0 \\ \xi^- \end{pmatrix}$$

Fields	$U(1)_Y$	$SU(2)_L$	$SU(3)_C$
ϕ	1/2	2	1
χ	1	3	1
ξ	0	3	1

$$\mathcal{L}_{GM} = \mathcal{L}_{kin} + \mathcal{L}_Y + \mathcal{L}_\nu - V_H$$

$$\begin{aligned} V(\Phi, \Delta) = & \frac{1}{2} m_\Phi^2 \text{Tr}[\Phi^\dagger \Phi] + \frac{1}{2} m_\Delta^2 \text{Tr}[\Delta^\dagger \Delta] \\ & + \lambda_1 (\text{Tr}[\Phi^\dagger \Phi])^2 + \lambda_2 (\text{Tr}[\Delta^\dagger \Delta])^2 \\ & + \lambda_3 \text{Tr}[(\Delta^\dagger \Delta)^2] + \lambda_4 \text{Tr}[\Phi^\dagger \Phi] \text{Tr}[\Delta^\dagger \Delta] \\ & + \lambda_5 \text{Tr} \left[\Phi^\dagger \frac{\sigma^a}{2} \Phi \frac{\sigma^b}{2} \right] \text{Tr}[\Delta^\dagger T^a \Delta T^b] \\ & + \mu_1 \text{Tr} \left[\Phi^\dagger \frac{\sigma^a}{2} \Phi \frac{\sigma^b}{2} \right] (P^\dagger \Delta P)_{ab} \\ & + \mu_2 \text{Tr}[\Delta^\dagger T^a \Delta T^b] (P^\dagger \Delta P)_{ab} \end{aligned}$$



$$\Phi = \begin{pmatrix} \phi^{0*} & \phi^+ \\ \phi^- & \phi^0 \end{pmatrix}, \quad \Delta = \begin{pmatrix} \chi^{0*} & \xi^+ & \chi^{++} \\ \chi^- & \xi^0 & \chi^+ \\ \chi^{--} & \xi^- & \chi^0 \end{pmatrix}$$

$$\mathcal{L}_\nu \supset h_{ij} \overline{L}_L^{ic} i\tau_2 \chi L_L^j + h.c.$$

Nucl.Phys.B 262 (1985) 463-477, JHEP 01 (2013) 026, JHEP 01 (2016) 120

1、Introduction to the Georgi-Machacek Model

$$\boxed{SU(2)_L \times SU(2)_R} \xrightarrow[\langle \chi^0 \rangle = \langle \xi^0 \rangle = v_\Delta]{\langle \phi^0 \rangle = v_\phi / \sqrt{2}} \boxed{SU(2)_V}$$

$$v_{EW}^2 = \sum_i [4T_i(T_i + 1) - Y_i^2] |v_i|^2 c_i = v_\phi^2 + 4v_\chi^2 + 4v_\xi^2 = v_\phi^2 + 8v_\Delta^2 = \frac{1}{\sqrt{2}G_F} \approx (246\text{GeV})^2$$

$$c_{T,Y} = \begin{cases} 1, & (T, Y) \in \text{complex representation} \\ 1/2, & (T, Y = 0) \in \text{real representation} \end{cases}$$

$$\rho_{tree} \equiv \frac{m_W^2}{m_Z^2 \cos^2 \theta_W} = \sum_i [4T_i(T_i + 1) - Y_i^2] |v_i|^2 c_i / \sum_i 2Y_i |v_i|^2$$

$$\Delta\rho_{tree} \equiv \rho - 1 = \frac{v_\phi^2 + 4v_\chi^2 + 4v_\xi^2}{v_\phi^2 + 8v_\Delta^2} - 1 \approx \frac{4v_\chi^2 - 4v_\xi^2}{v_{EW}^2}$$

$$\tan \theta = 2\sqrt{2}v_\Delta / v_\phi$$



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2、Improved the Vacuum Stability Constraint on the GM Model

Vacuum stability

$$\begin{aligned}
 V(\Phi, \Delta) = & \frac{1}{2} m_{\Phi}^2 \text{Tr}[\Phi^\dagger \Phi] + \frac{1}{2} m_{\Delta}^2 \text{Tr}[\Delta^\dagger \Delta] + \lambda_1 (\text{Tr}[\Phi^\dagger \Phi])^2 + \lambda_2 (\text{Tr}[\Delta^\dagger \Delta])^2 \\
 & + \lambda_3 \text{Tr}[(\Delta^\dagger \Delta)^2] + \lambda_4 \text{Tr}[\Phi^\dagger \Phi] \text{Tr}[\Delta^\dagger \Delta] + \lambda_5 \text{Tr}\left[\Phi^\dagger \frac{\sigma^a}{2} \Phi \frac{\sigma^b}{2}\right] \text{Tr}[\Delta^\dagger T^a \Delta T^b] \\
 & + \mu_1 \text{Tr}\left[\Phi^\dagger \frac{\sigma^a}{2} \Phi \frac{\sigma^b}{2}\right] (P^\dagger \Delta P)_{ab} + \mu_2 \text{Tr}[\Delta^\dagger T^a \Delta T^b] (P^\dagger \Delta P)_{ab}
 \end{aligned}$$

Z_2 ?



$$\begin{aligned}
 r & \equiv \sqrt{\text{Tr}[\Phi^\dagger \Phi] + \text{Tr}[\Delta^\dagger \Delta]}, & r^2 \cos^2 \gamma & \equiv \text{Tr}[\Phi^\dagger \Phi], & r^2 \sin^2 \gamma & \equiv \text{Tr}[\Delta^\dagger \Delta], \\
 \zeta & \equiv \frac{\text{Tr}[(\Delta^\dagger \Delta)^2]}{(\text{Tr}[\Delta^\dagger \Delta])^2}, & \omega & \equiv \frac{\text{Tr}\left[\Phi^\dagger \frac{\sigma^a}{2} \Phi \frac{\sigma^b}{2}\right] \text{Tr}[\Delta^\dagger T^a \Delta T^b]}{\text{Tr}[\Phi^\dagger \Phi] \text{Tr}[\Delta^\dagger \Delta]}. \\
 r & \in [0, \infty), & \gamma & \in [0, \frac{\pi}{2}], & \zeta & \in [\frac{1}{3}, 1], & \omega & \in [-\frac{1}{4}, \frac{1}{2}].
 \end{aligned}$$

$$V(r, \tan \gamma, \zeta, \omega) = \frac{r^4}{(1 + \tan^2 \gamma)^2} [\lambda_1 + (\lambda_4 - \omega \lambda_5) \tan^2 \gamma + (\zeta \lambda_3 + \lambda_2) \tan^4 \gamma]$$

arXiv:1105.1925, 1404.2640, JHEP 01 (2016) 120

2、Improved the Vacuum Stability Constraint on the GM Model

Vacuum stability

$$\lambda_1 > 0,$$
$$\lambda_2 > \begin{cases} -\frac{1}{3}\lambda_3 & \text{for } \lambda_3 \geq 0, \\ -\lambda_3 & \text{for } \lambda_3 < 0, \end{cases}$$
$$\lambda_4 > \begin{cases} \frac{1}{2}\lambda_5 - 2\sqrt{\lambda_1\left(\frac{1}{3}\lambda_3 + \lambda_2\right)} & \text{for } \lambda_5 \geq 0 \text{ and } \lambda_3 \geq 0, \\ \omega_+(\zeta)\lambda_5 - 2\sqrt{\lambda_1(\lambda_3 + \lambda_2)} & \text{for } \lambda_5 \geq 0 \text{ and } \lambda_3 < 0 \\ \omega_-(\zeta)\lambda_5 - 2\sqrt{\lambda_1(\lambda_3 + \lambda_2)}. & \text{for } \lambda_5 < 0 \end{cases}$$

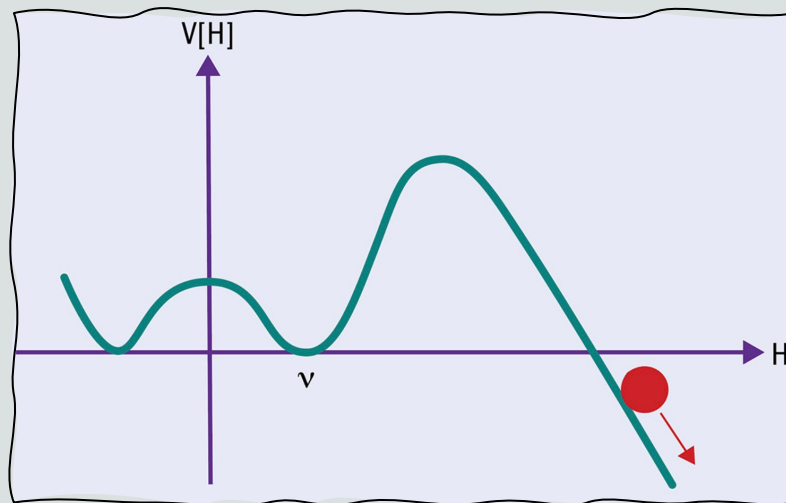
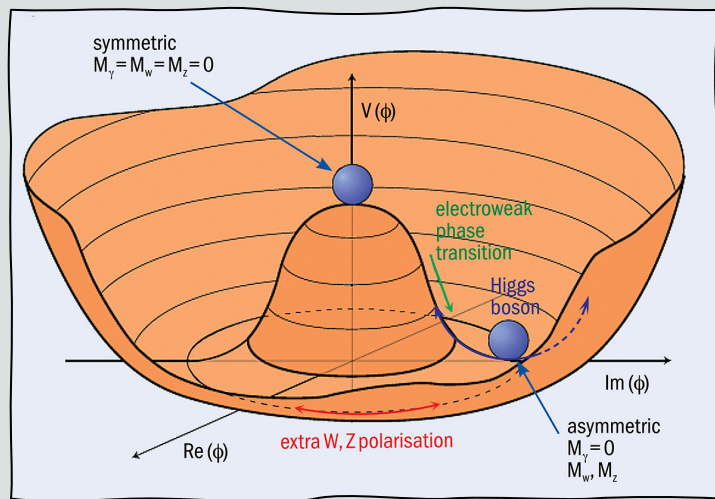
$$\lambda_1 > 0,$$
$$\lambda_2 + \lambda_3 > 0,$$
$$\lambda_2 + \frac{1}{2}\lambda_3 > 0,$$
$$-|\lambda_4| + 2\sqrt{\lambda_1(\lambda_2 + \lambda_3)} > 0,$$
$$\lambda_4 - \frac{1}{4}|\lambda_5| + \sqrt{2\lambda_1(2\lambda_2 + \lambda_3)} > 0.$$

Cheng-Wei Chiang etc, 1211.2658, 1511.00865

$$\omega_{\pm}(\zeta) = \frac{1}{6}(1 - B) \pm \frac{\sqrt{2}}{3} \left[(1 - B) \left(\frac{1}{2} + B \right) \right]^{1/2}, \quad B \equiv \sqrt{\frac{3}{2} \left(\zeta - \frac{1}{3} \right)} \in [0, 1], \quad \zeta \in \left[\frac{1}{3}, 1 \right].$$

Logan etc, arXiv:1404.2640, JHEP 01 (2016) 120

2、Improved the Vacuum Stability Constraint on the GM Model



Credit: J Ellis/M Neubauer

Large field approximation



High energy scale approximation



RGE

2、Improved the Vacuum Stability Constraint on the GM Model

$$\phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}, \quad \chi = \begin{pmatrix} \chi^+/\sqrt{2} & -\chi^{++} \\ \chi^0 & -\chi^+/\sqrt{2} \end{pmatrix}, \quad \xi = \begin{pmatrix} \xi^0/\sqrt{2} & -\xi^+ \\ -\xi^- & -\xi^0/\sqrt{2} \end{pmatrix}$$

$$V(\phi, \chi, \xi) = m_\phi^2 (\phi^\dagger \phi) + m_\chi^2 (\chi^\dagger \chi) + m_\xi^2 (\xi^\dagger \xi) + \lambda (\phi^\dagger \phi)^2 \\ + \rho_1 (\text{Tr}[\chi^\dagger \chi])^2 + \rho_2 \text{Tr}[\chi^\dagger \chi \chi^\dagger \chi] + \rho_3 \text{Tr}[\xi^4] \\ + \rho_4 \text{Tr}[\chi^\dagger \chi] \text{Tr}[\xi^2] + \rho_5 \text{Tr}[\chi^\dagger \xi] \text{Tr}[\xi \chi] + \sigma_1 \text{Tr}[\chi^\dagger \chi] \phi^\dagger \phi \\ + \sigma_2 \phi^\dagger \chi \chi^\dagger \phi + \sigma_3 \text{Tr}[\xi^2] \phi^\dagger \phi + \sigma_4 (\phi^\dagger \chi \zeta \phi^c + \text{h. c.})$$

$$\beta(g_3) = \frac{g_3^3}{16\pi^2} (-7), \quad \beta(g_2) = \frac{g_2^3}{16\pi^2} \left(-\frac{11}{6} \right), \quad \beta(g_1) = \frac{g_1^3}{16\pi^2} \frac{47}{6},$$

$$\beta(y_t) = \frac{1}{16\pi^2} \left[\frac{9}{2} y_t^3 + \frac{3}{2} y_b^3 - y_t \left(8g_3^2 + \frac{9}{4} g_2^2 + \frac{17}{12} g_1^2 \right) \right],$$

$$\beta(y_b) = \frac{1}{16\pi^2} \left[\frac{9}{2} y_b^3 + \frac{3}{2} y_t^3 - y_b \left(8g_3^2 + \frac{9}{4} g_2^2 + \frac{5}{12} g_1^2 \right) \right].$$

$$16\pi^2 \beta(\lambda) = \frac{3}{8} (3g_4^4 + 2g_2^2 g_1^2 + g_1^4) + 24\lambda^2 - 6(y_t^4 + y_b^4) + 3\sigma_1^2 + 3\sigma_1 \sigma_2 + \frac{5\sigma_2^2}{4} + 6\sigma_3^2 + 2\sigma_4^2 \\ - 3\lambda(g_1^2 + 3g_2^2 - 4y_t^2 - 4y_b^2), \quad (C4)$$

$$16\pi^2 \beta(\rho_1) = 15g_2^4 - 12g_1^2 g_2^2 + 6g_1^4 + 28\rho_1^2 + 24\rho_1 \rho_2 + 6\rho_2^2 + 6\rho_4^2 + 4\rho_4 \rho_5 + 3\rho_5^2 \\ + 2\sigma_1^2 + 2\sigma_1 \sigma_2 - 12\rho_1 (g_1^2 + 2g_2^2),$$

$$16\pi^2 \beta(\rho_2) = 24g_1^2 g_2^2 - 6g_2^4 + 24\rho_1 \rho_2 + 18\rho_2^2 - 2\rho_5^2 + \sigma_2^2 - 12\rho_2 (2g_2^2 + g_1^2), \quad (C6)$$

$$16\pi^2 \beta(\rho_3) = 2(3g_2^4 + 22\rho_3^2 + 3\rho_4^2 + 2\rho_4 \rho_5 + \rho_5^2 + 2\sigma_2^2 - 12g_2^2 \rho_3), \quad (C7)$$

$$16\pi^2 \beta(\rho_4) = 2 \left[3g_2^4 + \rho_4 (8\rho_1 + 6\rho_2 + 10\rho_3 + 4\rho_4) + 2\rho_5 (\rho_1 + \rho_2 + \rho_3) \right. \\ \left. + \rho_5^2 + 2\sigma_1 \sigma_3 + \sigma_2 \sigma_3 + \sigma_4^2 - 3\rho_4 (g_1^2 + 4g_2^2) \right], \quad (C8)$$

$$16\pi^2 \beta(\rho_5) = 2 \left[3g_2^4 + \rho_5 (2\rho_1 + 4\rho_3 + 8\rho_4 + 5\rho_5) - \sigma_4^2 - 3\rho_5 (4g_2^2 + g_1^2) \right], \quad (C9)$$

$$16\pi^2 \beta(\sigma_1) = 3g_1^4 - 6g_1^2 g_2^2 + 6g_2^4 + 2\sigma_1 (6\lambda + 8\rho_1 + 6\rho_2 + 2\sigma_1) + 2\sigma_2 (2\lambda + 3\rho_1 + \rho_2) \\ + (2(6\rho_4 \sigma_3 + 2\rho_5 \sigma_3 + \sigma_4^2) + \sigma_2^2 - \frac{3}{2} \sigma_1 (5g_1^2 + 11g_2^2 - 4y_t^2 - 4y_b^2)), \quad (C10)$$

$$16\pi^2 \beta(\sigma_2) = 12g_1^2 g_2^2 + 4\sigma_2 [\lambda + \rho_1 + 2(\rho_2 + \sigma_1) + \sigma_2] + 4\sigma_4^2 \\ - \frac{3}{2} \sigma_2 (5g_1^2 + 11g_2^2 - 4y_t^2 - 4y_b^2), \quad (C11)$$

$$16\pi^2 \beta(\sigma_3) = 3g_2^4 + 2\sigma_3 (6\lambda + 10\rho_3 + 4\sigma_3) + (3\rho_4 + \rho_5) (2\sigma_1 + \sigma_2) + 4\sigma_4^2 \\ - \frac{3}{2} \sigma_3 (g_1^2 + 11g_2^2 - 4y_t^2 - 4y_b^2), \quad (C12)$$

$$16\pi^2 \beta(\sigma_4) = \frac{\sigma_4}{2} [4(2\lambda + 2\rho_4 - \rho_5 + 2\sigma_1 + 2\sigma_2 + 4\sigma_3) - 3(3g_1^2 + 11g_2^2 - 4y_t^2 - 4y_b^2)], \quad (C13)$$

Simone Blasi et al, arXiv:1704.08512,

Logan et al, arXiv:1807.11511

$$\beta(X) \equiv \frac{d}{d \ln \mu} X$$

2、Improved the Vacuum Stability Constraint on the GM Model

$$\begin{aligned}
 V(\phi, \chi, \xi) = & m_\phi^2(\phi^\dagger\phi) + m_\chi^2(\chi^\dagger\chi) + m_\xi^2(\xi^\dagger\xi) + \lambda(\phi^\dagger\phi)^2 \\
 & + \rho_1(\text{Tr}[\chi^\dagger\chi])^2 + \rho_2\text{Tr}[\chi^\dagger\chi\chi^\dagger\chi] + \rho_3\text{Tr}[\xi^4] \\
 & + \rho_4\text{Tr}[\chi^\dagger\chi]\text{Tr}[\xi^2] + \rho_5\text{Tr}[\chi^\dagger\xi]\text{Tr}[\xi\chi] + \sigma_1\text{Tr}[\chi^\dagger\chi]\phi^\dagger\phi \\
 & + \sigma_2\phi^\dagger\chi\chi^\dagger\phi + \sigma_3\text{Tr}[\xi^2]\phi^\dagger\phi + \sigma_4(\phi^\dagger\chi\zeta\phi^c + \text{h.c.})
 \end{aligned}$$

$SU(2)_V$
Limit

$$\begin{aligned}
 m_\phi^2 = 2m_\Phi^2, m_\chi^2 = 2m_\Delta^2, m_\xi^2 = m_\Delta^2, \mu_1 = -\frac{\bar{\mu}_1}{\sqrt{2}}, \mu_2 = -\frac{\bar{\mu}_1}{\sqrt{2}}, \\
 \mu_3 = 6\sqrt{2}\bar{\mu}_2, \lambda = 4\lambda_1, \rho_1 = 4\lambda_4 + 6\lambda_3, \rho_2 = -4\lambda_3, \rho_3 = 2(\lambda_4 + \lambda_3), \\
 \rho_4 = 4\lambda_4, \rho_5 = 4\lambda_3, \sigma_1 = 4\lambda_2 - \lambda_5, \sigma_2 = 2\lambda_5, \sigma_3 = 2\lambda_2, \sigma_4 = \sqrt{2}\lambda_5.
 \end{aligned}$$

$$\begin{aligned}
 m_\chi^2 = 2m_\xi^2, \quad \mu_1 = \sqrt{2}\mu_2, \\
 \rho_3 = \frac{1}{2}\rho_1 + \frac{1}{4}\rho_2, \quad \rho_4 = \rho_1 + \frac{3}{2}\rho_2, \quad \rho_5 = -2\rho_2 \\
 \sigma_3 = \frac{1}{2}\sigma_1 + \frac{1}{4}\sigma_2, \quad \sigma_4 = \frac{1}{\sqrt{2}}\sigma_2.
 \end{aligned}$$

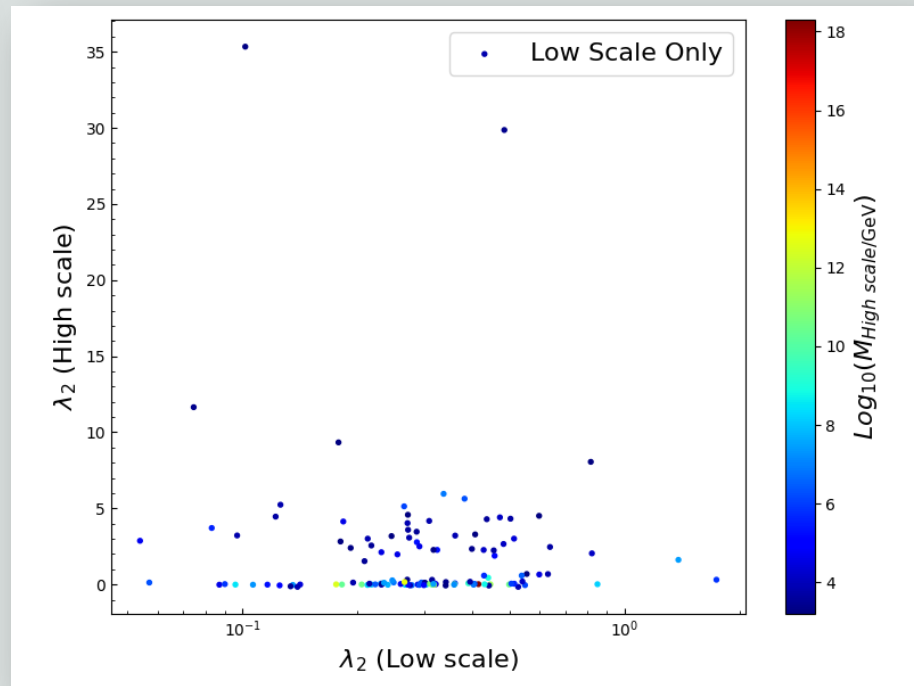
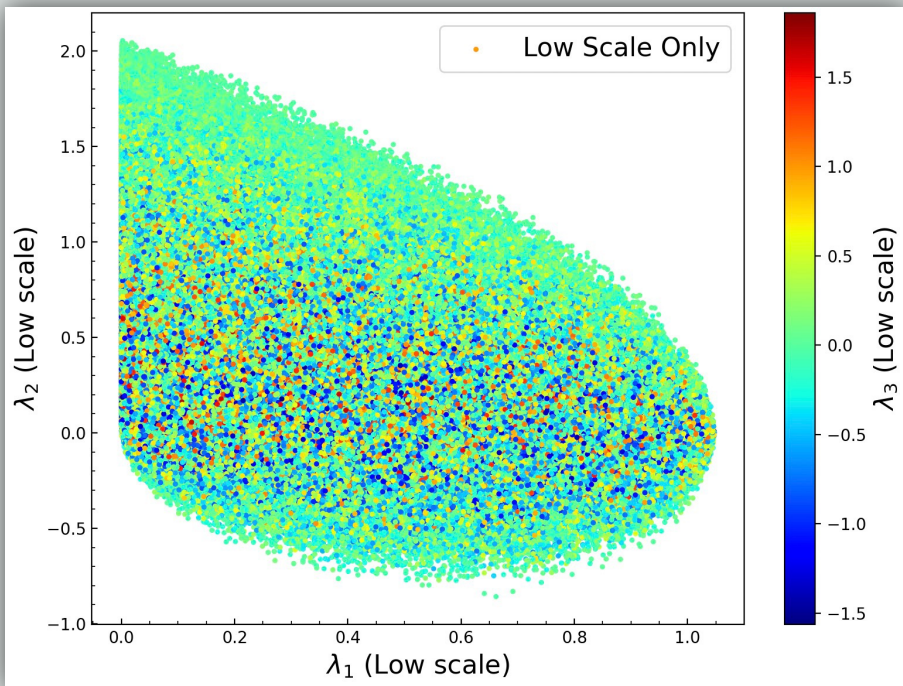
$$\begin{aligned}
 \delta_m^2 = m_\chi^2 - 2m_\xi^2, \quad \delta_\mu = \mu_1 - \sqrt{2}\mu_2, \\
 \delta_\rho^1 = \rho_3 - \frac{1}{2}\rho_1 - \frac{1}{4}\rho_2, \quad \dots \dots \\
 \delta_\sigma^1 = \sigma_3 - \frac{1}{2}\sigma_1 - \frac{1}{4}\sigma_2, \quad \dots \dots
 \end{aligned}$$

Simone Blasi et al, arXiv:1704.08512

2、Improved the Vacuum Stability Constraint on the GM Model

Low: ✓

High: ✗

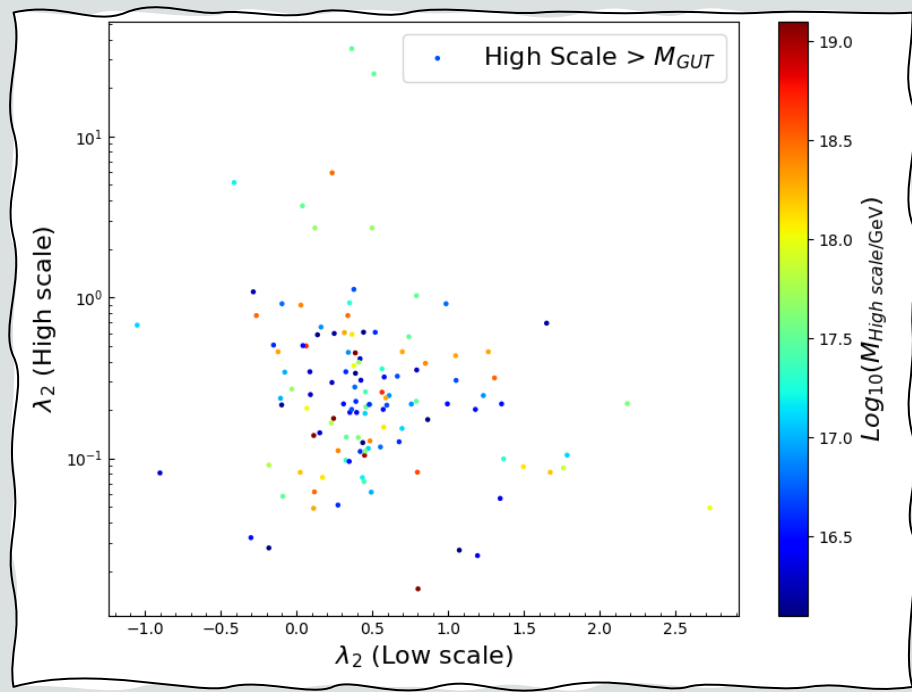
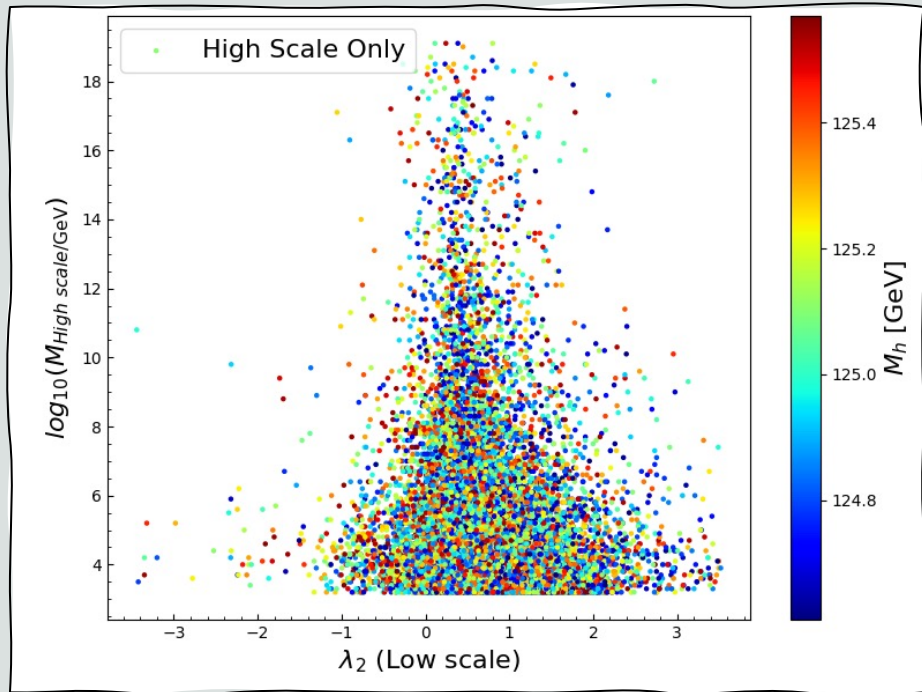


2、Improved the Vacuum Stability Constraint on the GM Model

High energy scale :



Low energy scale :

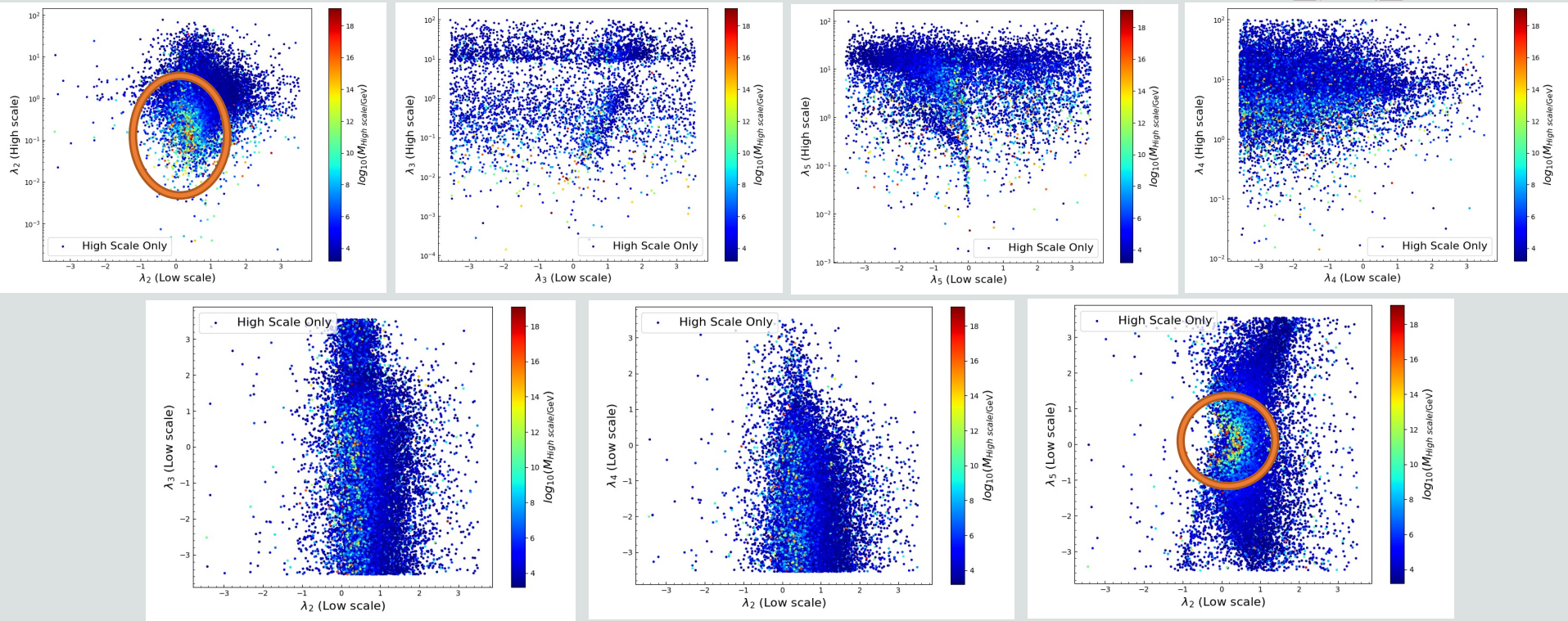


2、Improved the Vacuum Stability Constraint on the GM Model

High energy scale : ✓



Low energy scale : ✗





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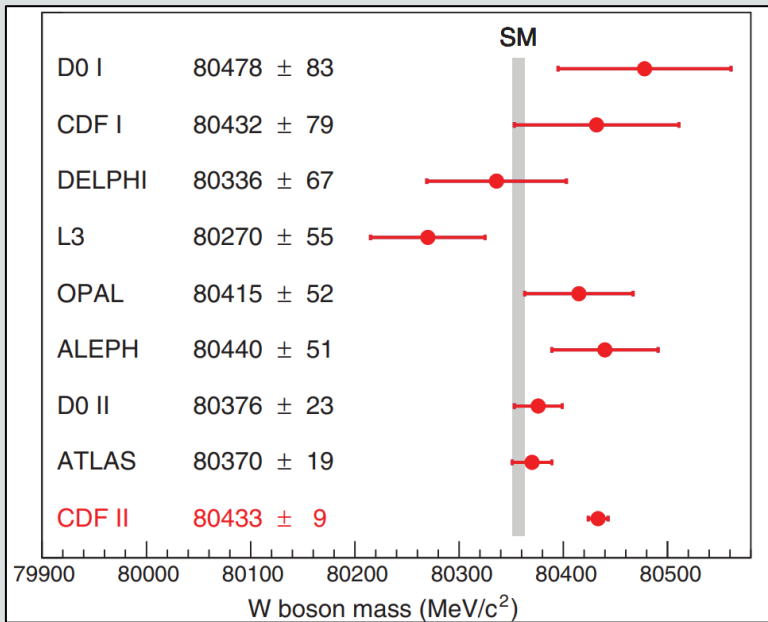
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3、W-Boson Mass and CEPC constraints on the the GM Model

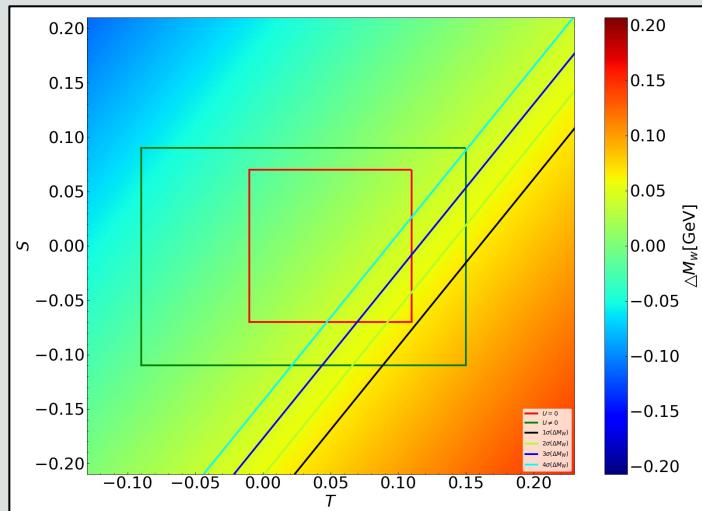
CDF-II Results on W Boson Mass



SM: $M_W = 80357 \pm 6$ MeV

CDF-II: $M_W = 80433.5 \pm 9.4$ MeV

Science 376, 170-176 (2022)



$S = 0.00 \pm 0.07,$
 $T = 0.05 \pm 0.06,$
 $U = 0$

$S = -0.01 \pm 0.10,$
 $T = 0.03 \pm 0.012,$
 $U = 0.02 \pm 0.11$

Phys. Rev. D 46 (1992) 381,
PTEP 2020 (2020) 8, 083C01

$$\Delta m_W = \frac{\alpha M_W}{2(c_W^2 - s_W^2)} \left(-\frac{1}{2} S + c_W^2 T + \frac{c_W^2 - s_W^2}{4s_W^2} U \right),$$

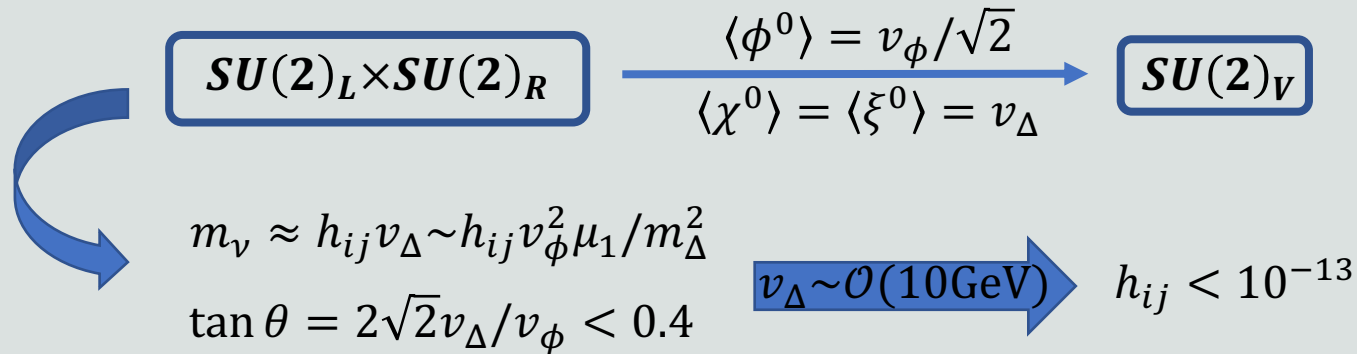
$$\alpha S = 4s_W^2 c_W^2 \left[\Pi'_{ZZ}(0) - \frac{c_W^2 - s_W^2}{s_W c_W} \Pi'_{Z\gamma}(0) - \Pi'_{\gamma\gamma}(0) \right],$$

$$\alpha T = \frac{\Pi_{WW}(0)}{m_W^2} - \frac{\Pi_{ZZ}(0)}{m_Z^2},$$

$$\alpha U = 4s_W^2 \left[\Pi'_{WW}(0) - c_W^2 \Pi'_{ZZ}(0) - 2s_W c_W \Pi'_{Z\gamma}(0) - s_W^2 \Pi'_{\gamma\gamma}(0) \right]$$

3、W-Boson Mass and CEPC constraints on the the GM Model

$$\mathcal{L}_{type-II} \supset h_{ij} \overline{L}_L^{ic} i\tau_2 \chi L_L^j + \mu_1 \text{tr} \left[\Phi^\dagger \frac{\sigma^a}{2} \Phi \frac{\sigma^b}{2} \right] (P^\dagger \Delta P)_{ab}$$



Vacuum stability

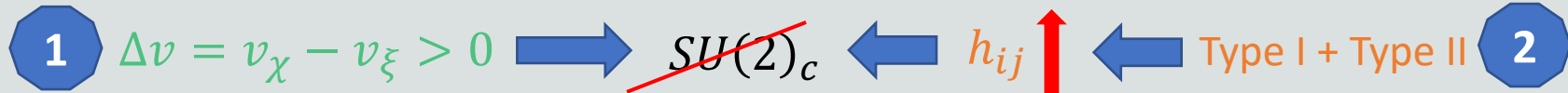
$$\begin{aligned} \lambda_1 > 0, \lambda_2 + \lambda_3 > 0, \lambda_2 + \frac{1}{2} \lambda_3 > 0, \\ -|\lambda_4| + 2\sqrt{\lambda_1(\lambda_2 + \lambda_3)} > 0, \\ \lambda_4 - \frac{1}{4} |\lambda_5| + \sqrt{2\lambda_1(2\lambda_2 + \lambda_3)} > 0. \end{aligned}$$

Perturbative unitarity

$$\begin{aligned} |\lambda_4 - \lambda_5| < 2\pi, \quad |2\lambda_3 + \lambda_2| < \pi, \\ |6\lambda_1 + 7\lambda_3 + 11\lambda_2| + \sqrt{(6\lambda_1 - 7\lambda_3 - 11\lambda_2)^2 + 36\lambda_4^2} < 4\pi, \\ |2\lambda_1 - \lambda_3 + 2\lambda_2| + \sqrt{(2\lambda_1 + \lambda_3 - 2\lambda_2)^2 + \lambda_5^2} < 4\pi. \end{aligned}$$

Nucl.Phys.B 262 (1985) 463-477, JHEP 01 (2013) 026, JHEP 01 (2016) 120

3、W-Boson Mass and CEPC constraints on the the GM Model



1. hypercharge gauge boson loops as a consequence of $SU(2)_L \times SU(2)_R$ breaking effects in the kinetic term.
2. Loop contributions contain ultra-violet (UV) divergences which cannot be cancelled by counterterms associated with the Vcst part alone.
3. ...



$\langle H \rangle$

$M_\nu = -\langle H \rangle^2 Y_\nu M_R^{-1} Y_\nu^T$

+

$\langle H \rangle$

$M_\nu = \langle H \rangle^2 h_{ij} \mu / M_\Delta^2$

$$-\mathcal{L}_\nu \supset y_{ij}^N \bar{L}_{L,i} \phi N_{R,j} + \frac{1}{2} (M_R)_{ij} N_{R,i}^T C N_{R,j} + h_{ij} \bar{L}_L^{iC} i\tau_2 \chi L_L^j + h.c.$$

$$M_\nu = \begin{pmatrix} h_{ij} v_\Delta & (y_{ij}^N)^T v_\phi \\ y_{ij}^N v_\phi & (M_R)_{ij} \end{pmatrix} \quad M_R \gg v_\phi \gg v_\Delta$$

Phys.Rev.D 98 (2018) 1, 013008, Phys.Lett.B 774 (2017) 119-122

3、W-Boson Mass and CEPC constraints on the the GM Model



$$-\mathcal{L}_\nu \supset y_{ij}^N \bar{L}_{L,i} \phi N_{R,j} + \frac{1}{2} (M_R)_{ij} N_{R,i}^T C N_{R,j} + h_{ij} \bar{L}_L^{ic} i\tau_2 \chi L_L^j + h.c.$$

$$m_\nu \approx h_{ij} v_\Delta - v_\phi^2 (y_{ij}^N)^T M_{R,j}^{-1} (y_{ij}^N), \quad h_{ij} v_\Delta \approx (y^N v_\phi)^2 / M_{R,j}^{-1}$$

$$h_{ij} = 2\sqrt{2} (V_{PMNS}^T)^{-1} \left(\frac{v(1-s_H^2)}{s_H M_{R,i}} \right) \delta_{ij} (V_{PMNS})^{-1}, \quad y_{ij}^N = (V_{PMNS})^{-1}$$

$$\begin{aligned} BR(\mu \rightarrow e\gamma) &< 4.2 \times 10^{-13}, \\ BR(\mu \rightarrow 3e) &< 10^{-12} \end{aligned}$$

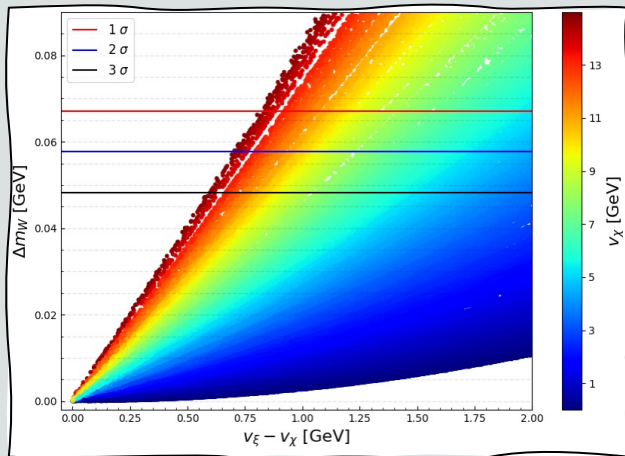
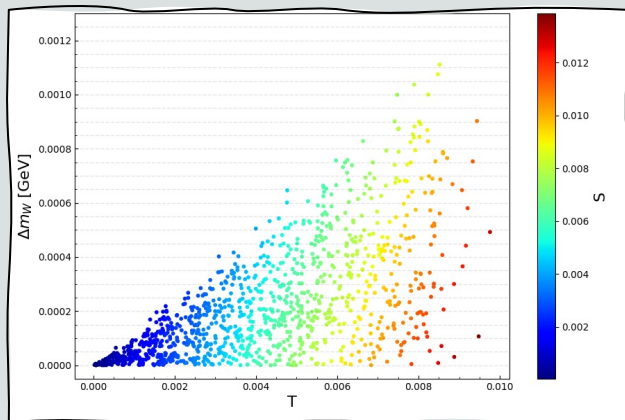
.....

$$BR(\mu \rightarrow e\gamma) \sim \frac{\alpha_{EM}}{192\pi} |h_{ij}|^4 \left(\frac{m_W}{M_{H^{++}}} \right)^4$$

$$h_{ij} < 10^{-2}, \quad M_R > 50\text{TeV}$$

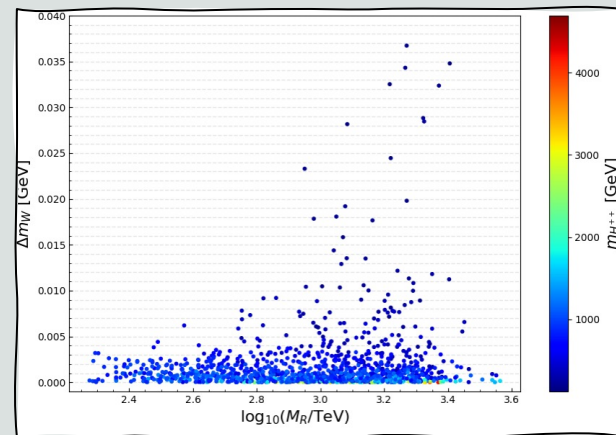
3、W-Boson Mass and CEPC constraints on the the GM Model

original
GM
Model



2

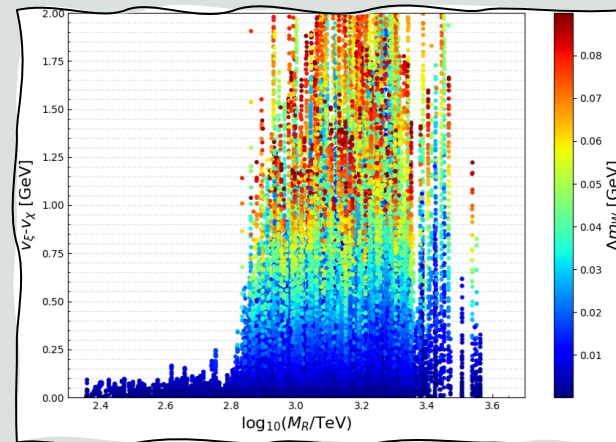
Type-I
+
Type-II



1

+

2

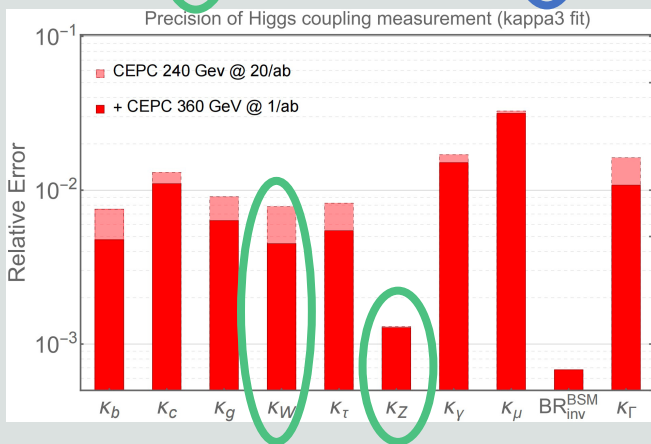
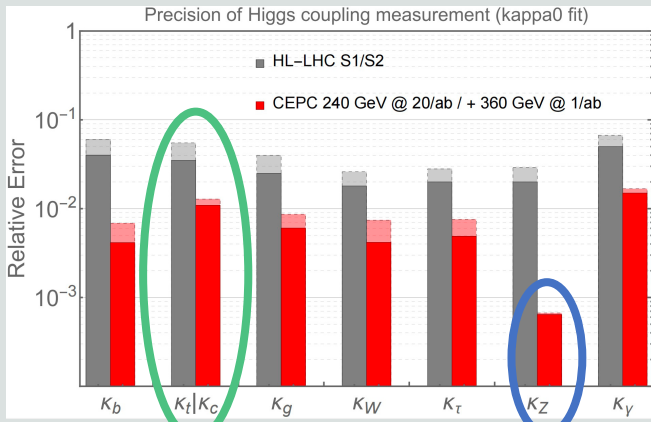


1

$$v_\chi \neq v_\xi$$

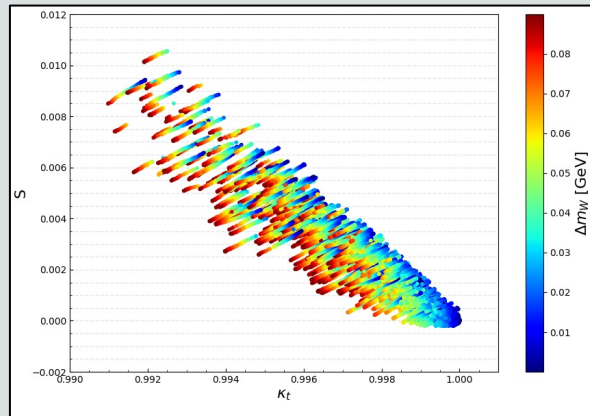
3、W-Boson Mass and CEPC constraints on the the GM Model

Precision of Higgs coupling measurement At CEPC



2

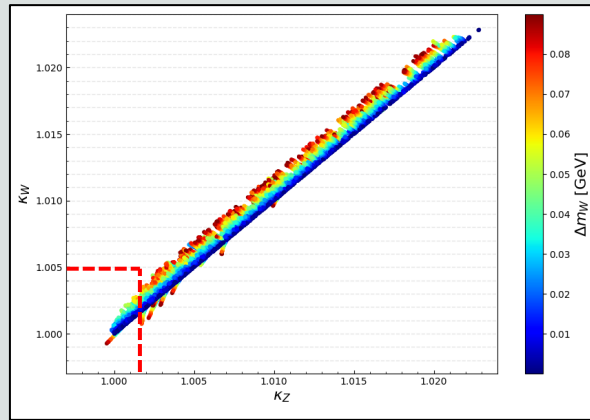
Type-I
+
Type-II



1

+

2



CEPC Snowmass report
arXiv:2205.08553

3、W-Boson Mass and CEPC constraints on the the GM Model

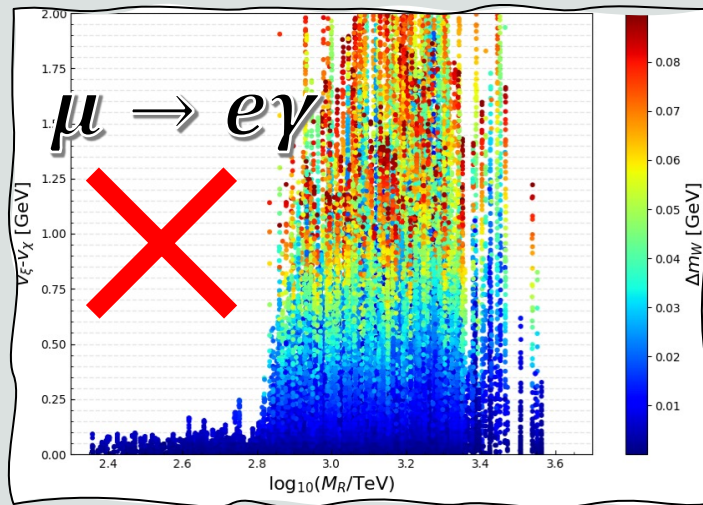
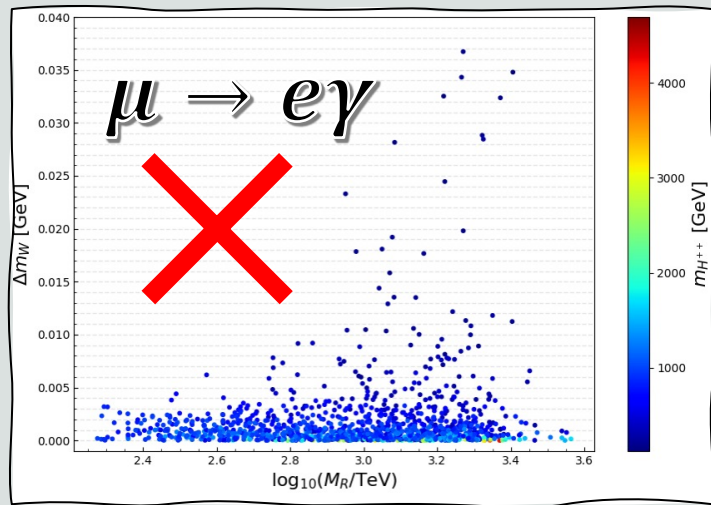
2

Type-I + Type-II

1

+

2



$$Br(\mu \rightarrow e\gamma) < 4.2 \times 10^{-13} \quad \text{MEG(2016)}$$

$$< 6 \times 10^{-14} \quad \text{MEG II}$$

$$Br(\tau \rightarrow e\mu\mu) < 2.7 \times 10^{-8} \quad \text{Bell(2010)}$$

$$< \sim 10^{-10} \quad \text{FCC/CEPC}$$



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**Summary
and Q&A**

4、 Summary and Q&A

1. Improved the Vacuum Stability Constraint on the GM model ;
2. Taken a general discussion about the contribution to W boson mass in the original GM Model; Explaining CDF-II results in the GM Extension Models;
3. Discussed the CEPC constraints on the GM model, especially the lepton flavor violation constraints.

THANKS

Q&A