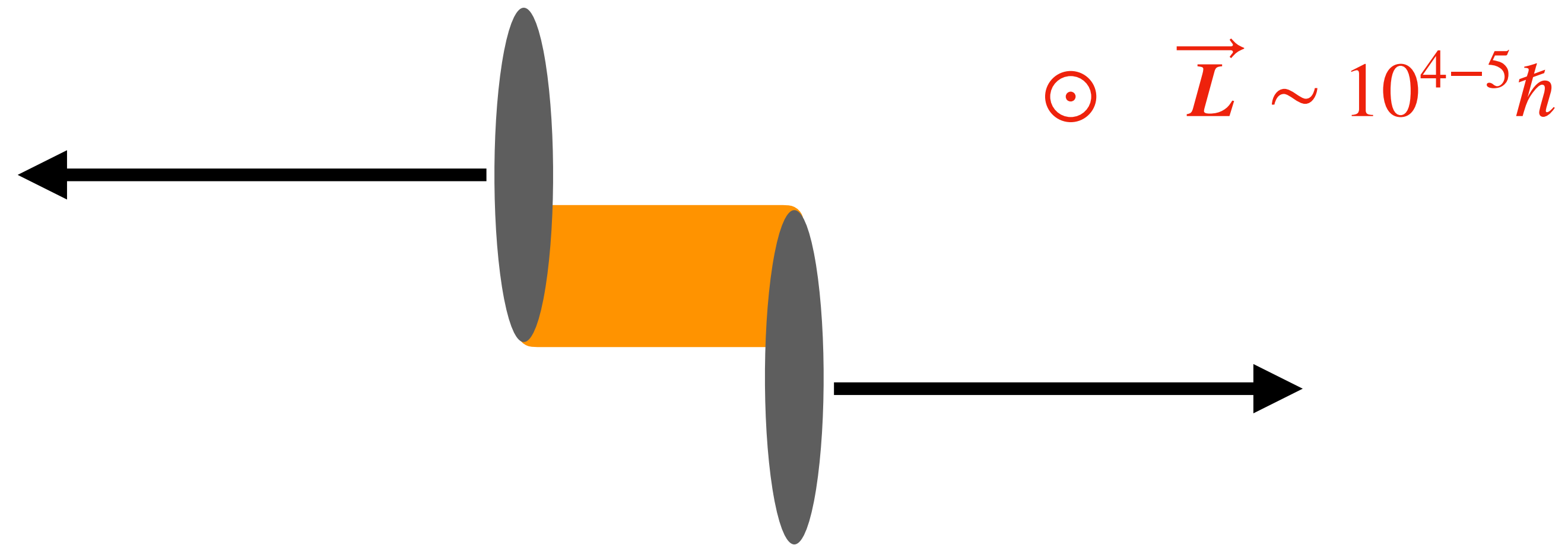


Spin and Chirality Transport Theories

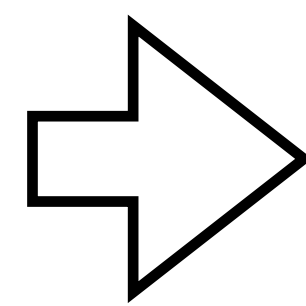
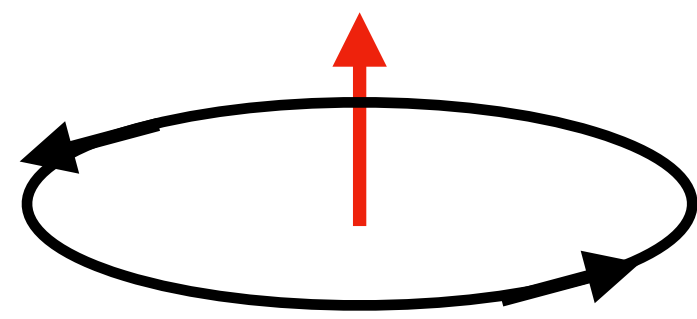
施舒哲(Shuzhe Shi)

清华大学(Tsinghua Univ.)

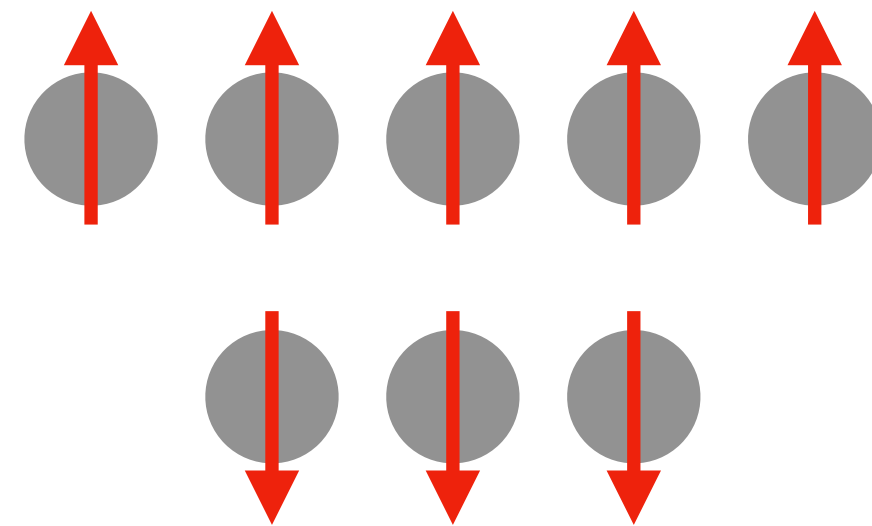


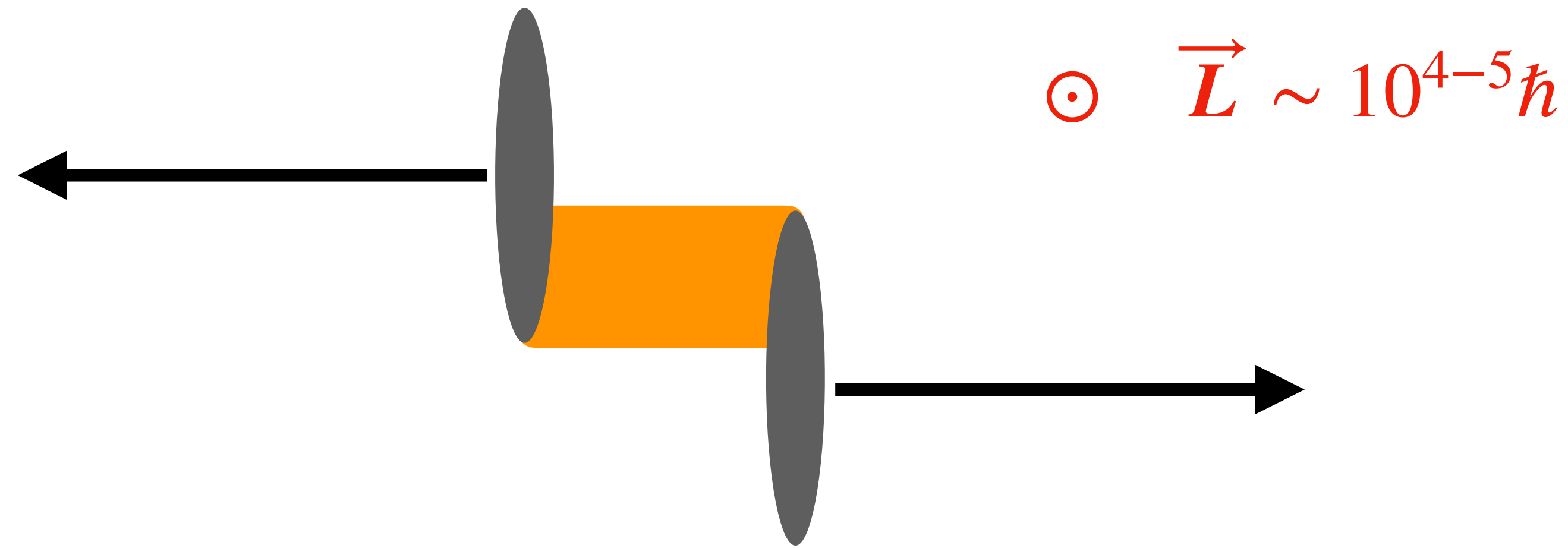
Angular Velocity: $\omega \sim 10^{21}$ Hz

rotation

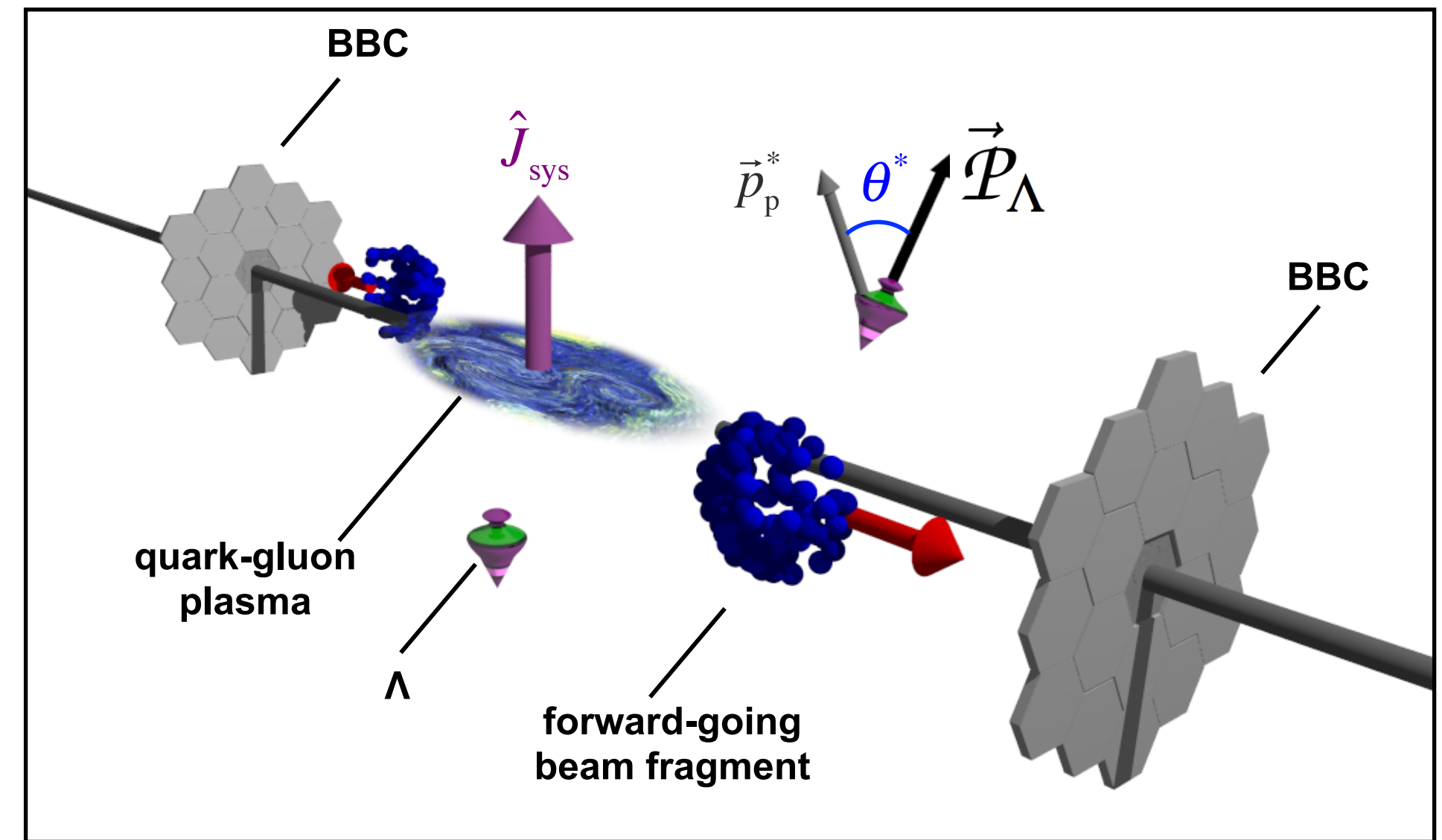


spin

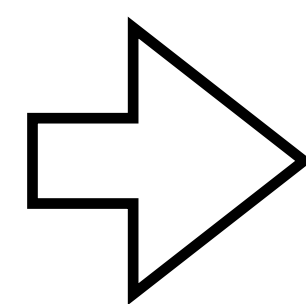
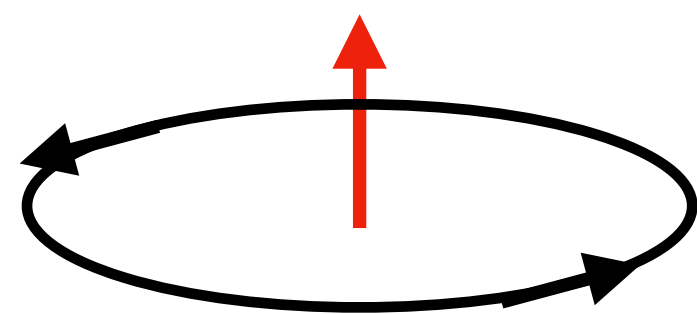




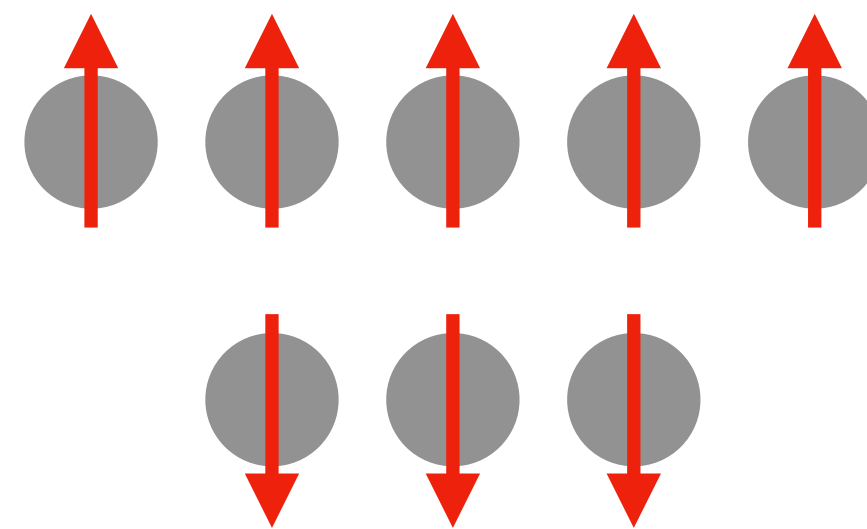
Angular Velocity: $\omega \sim 10^{21}$ Hz

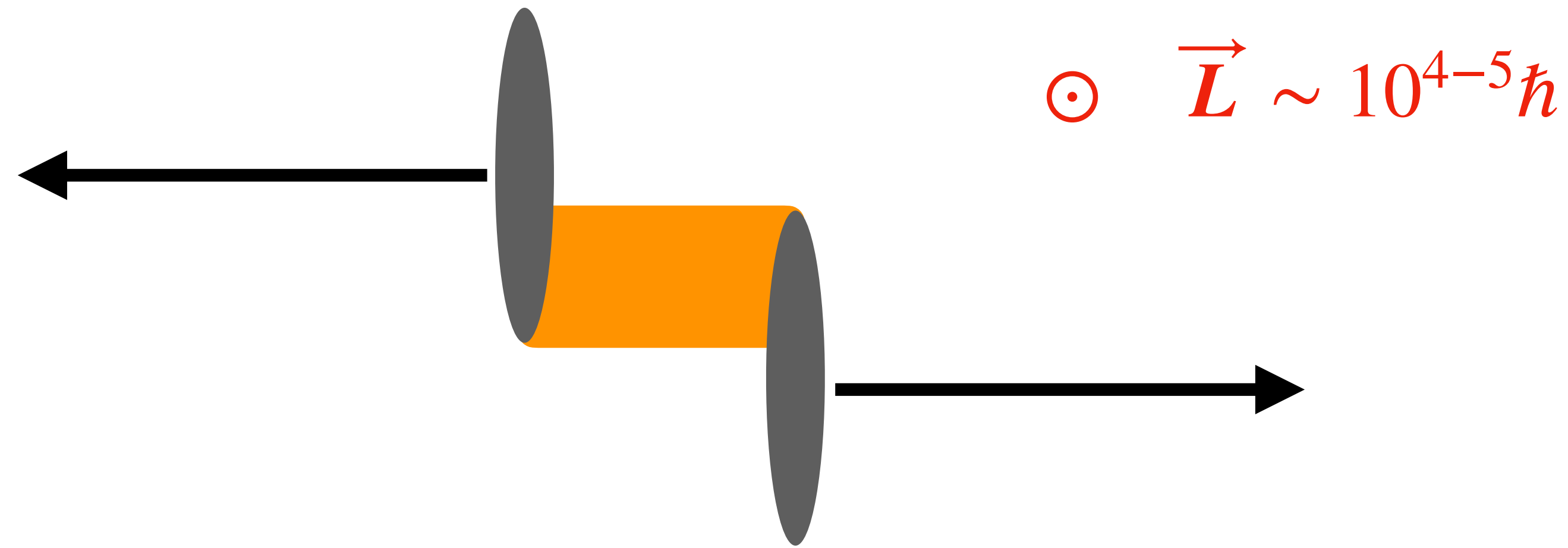


rotation

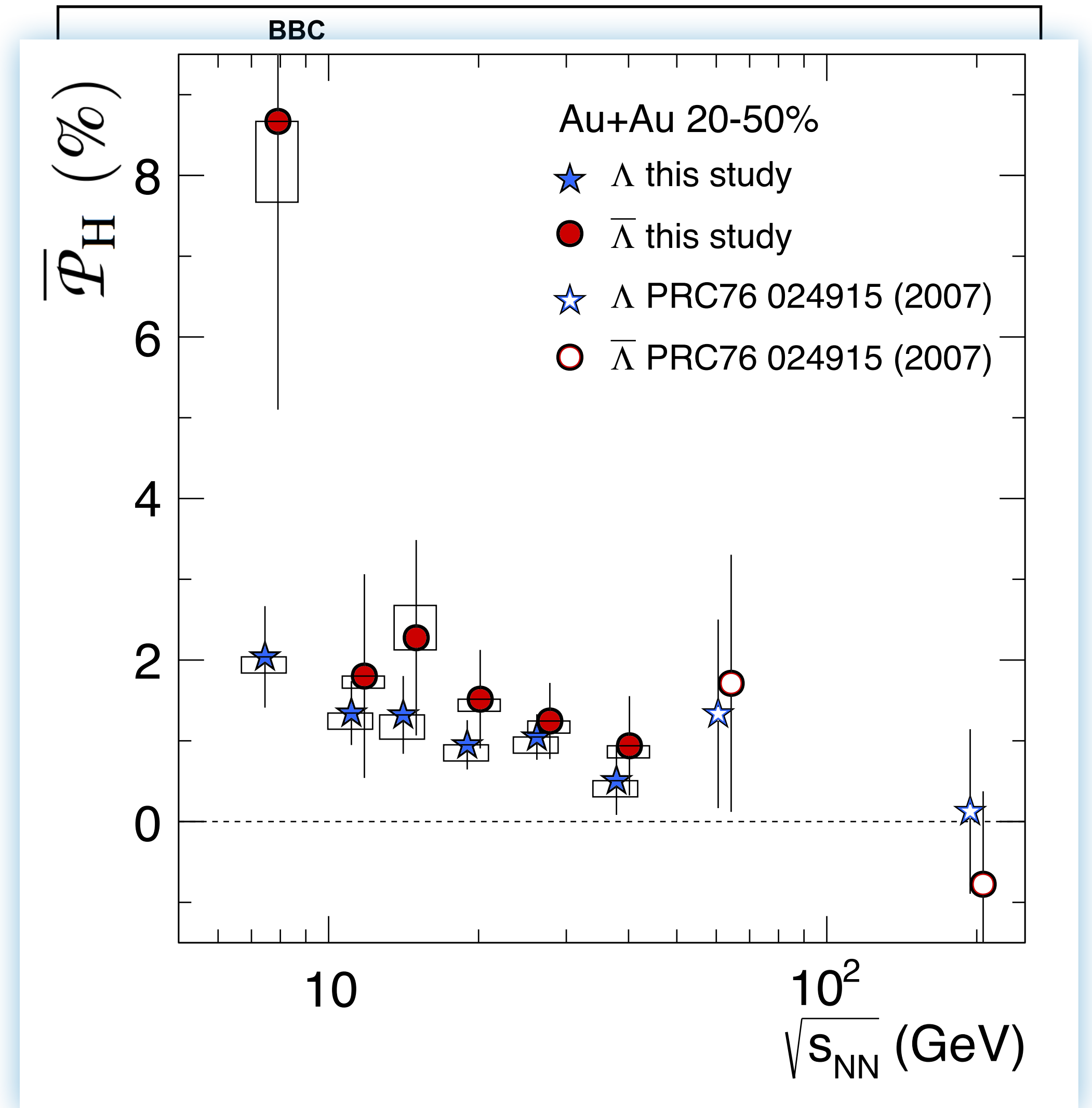
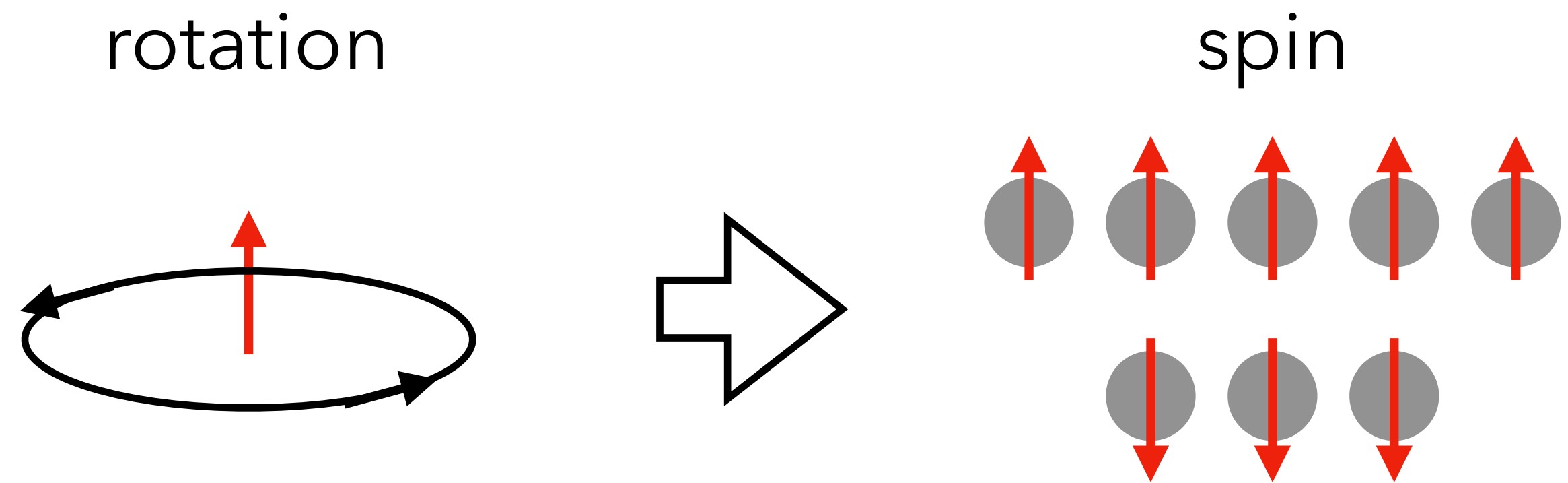


spin

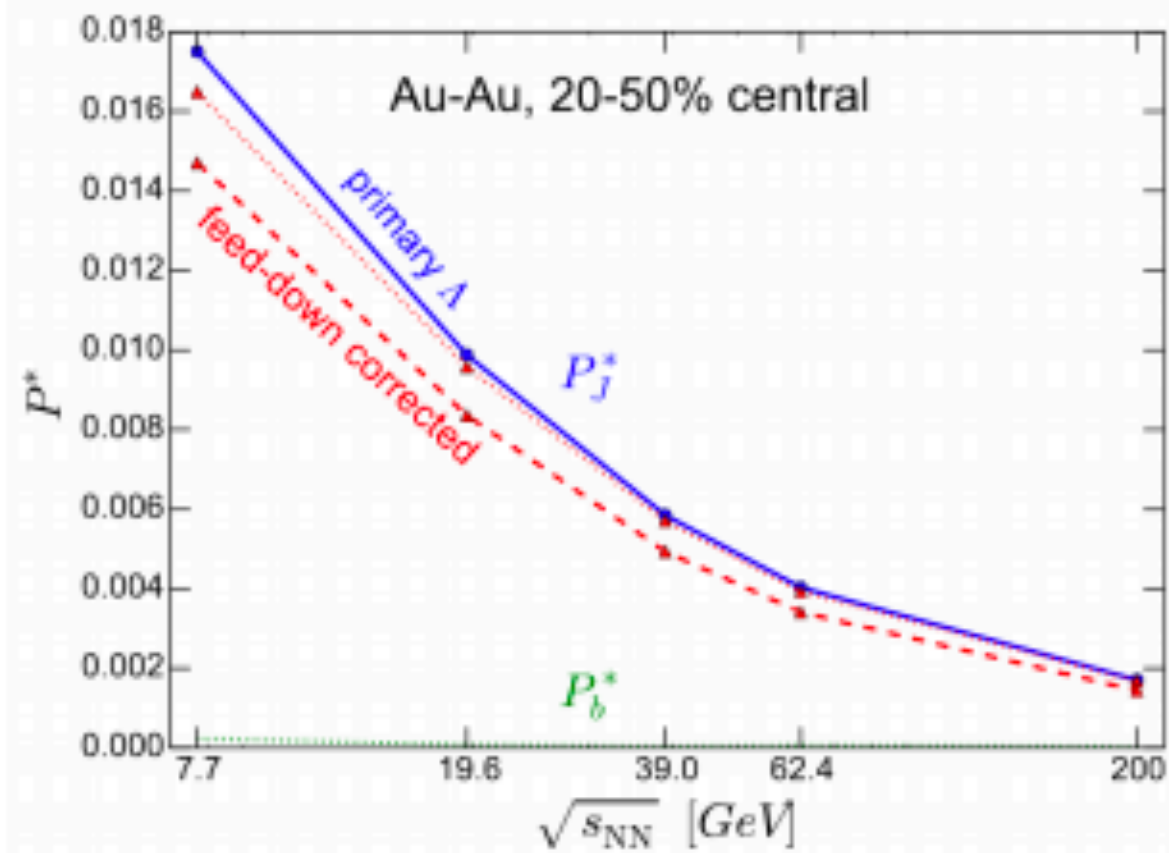




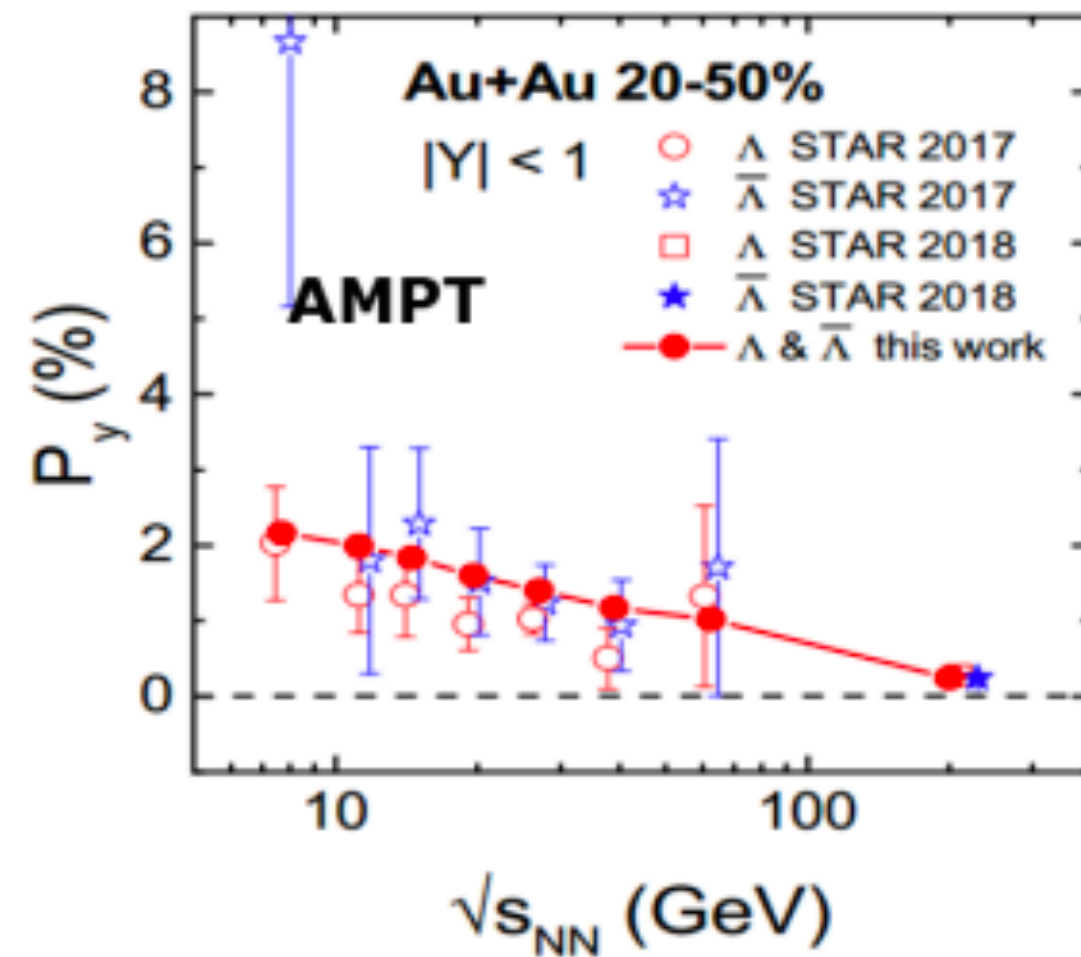
Angular Velocity: $\omega \sim 10^{21}$ Hz



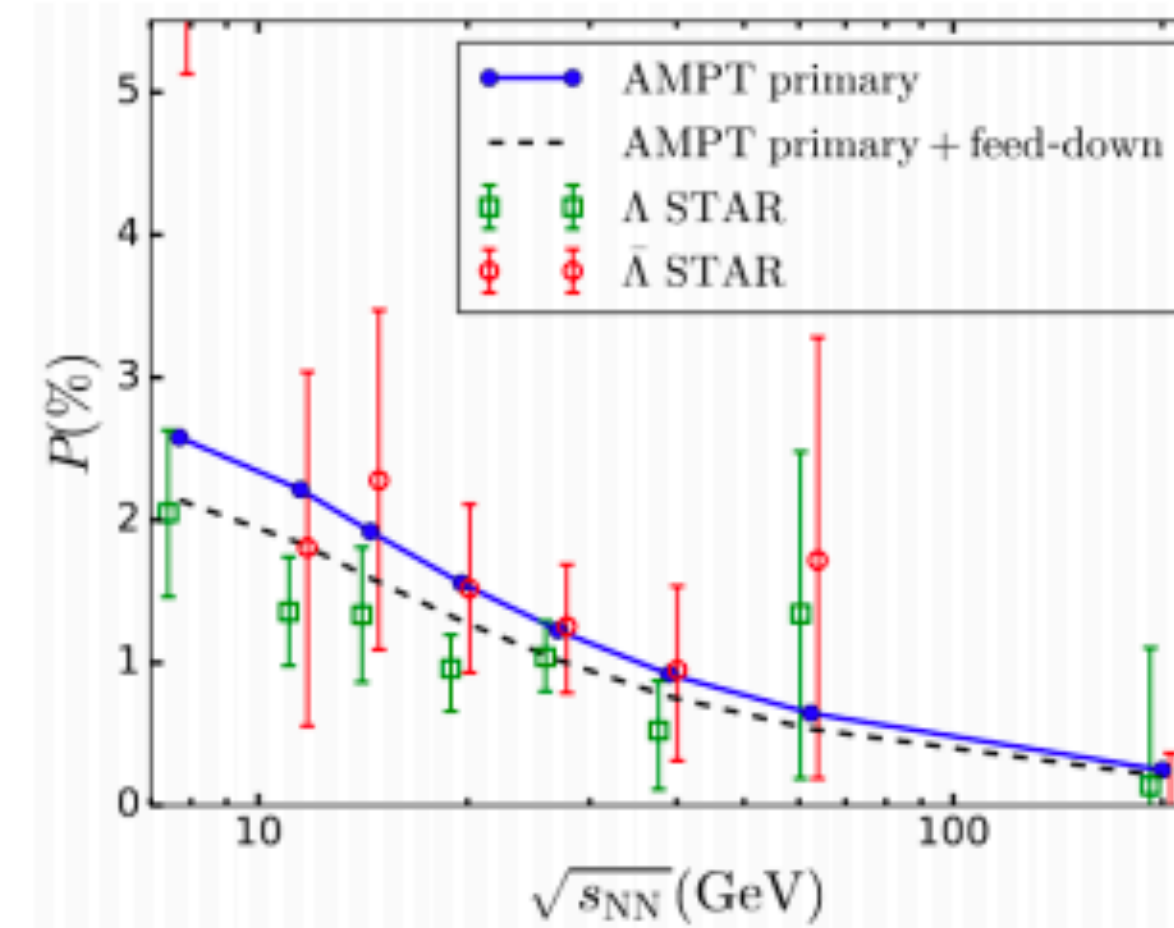
(Karpenko-Becattini EPJC2016)



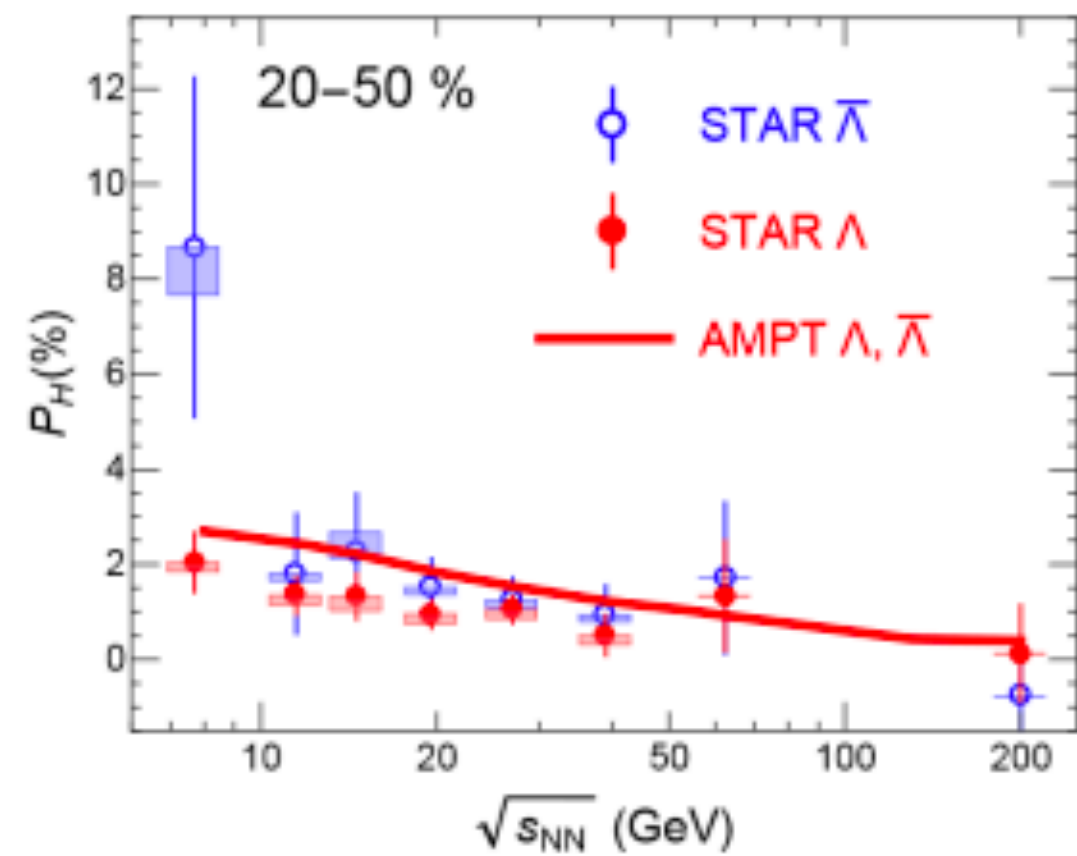
(Wei-Deng-XGH PRC2019)



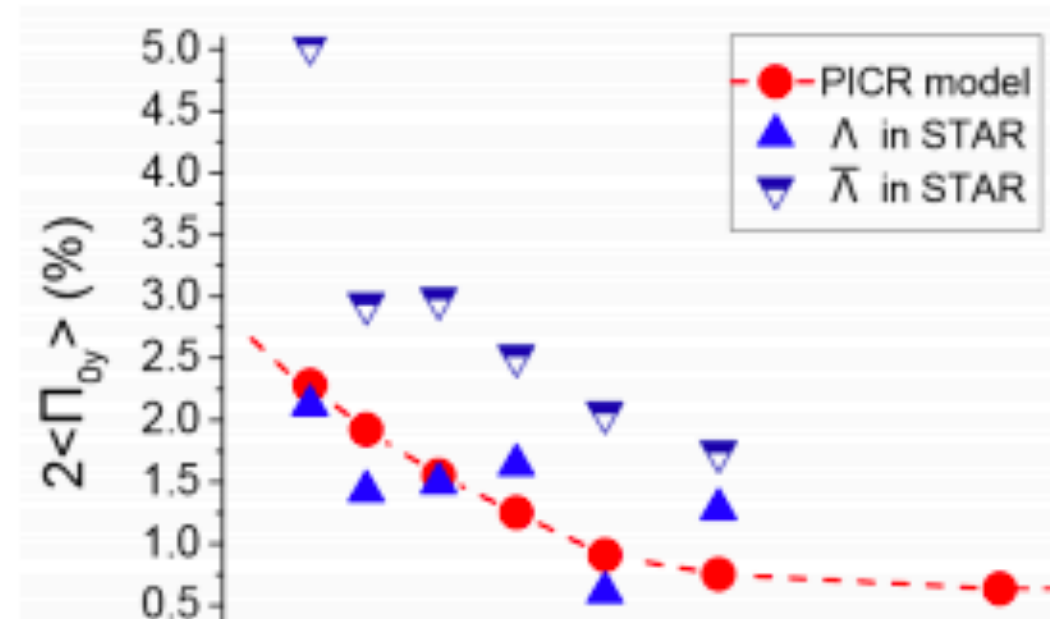
(Li-Pang-Wang-Xia PRC2017)



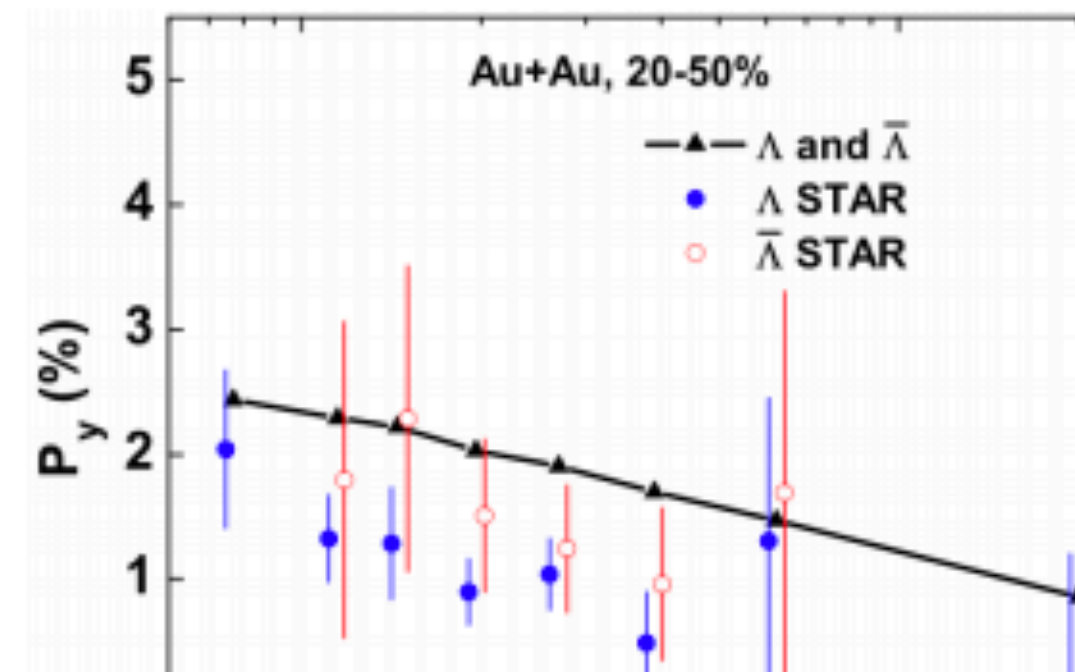
(Shi-Li-Liao PLB2018)



(Xie-Wang-Csernai PRC2017)



(Sun-Ko PRC2017)

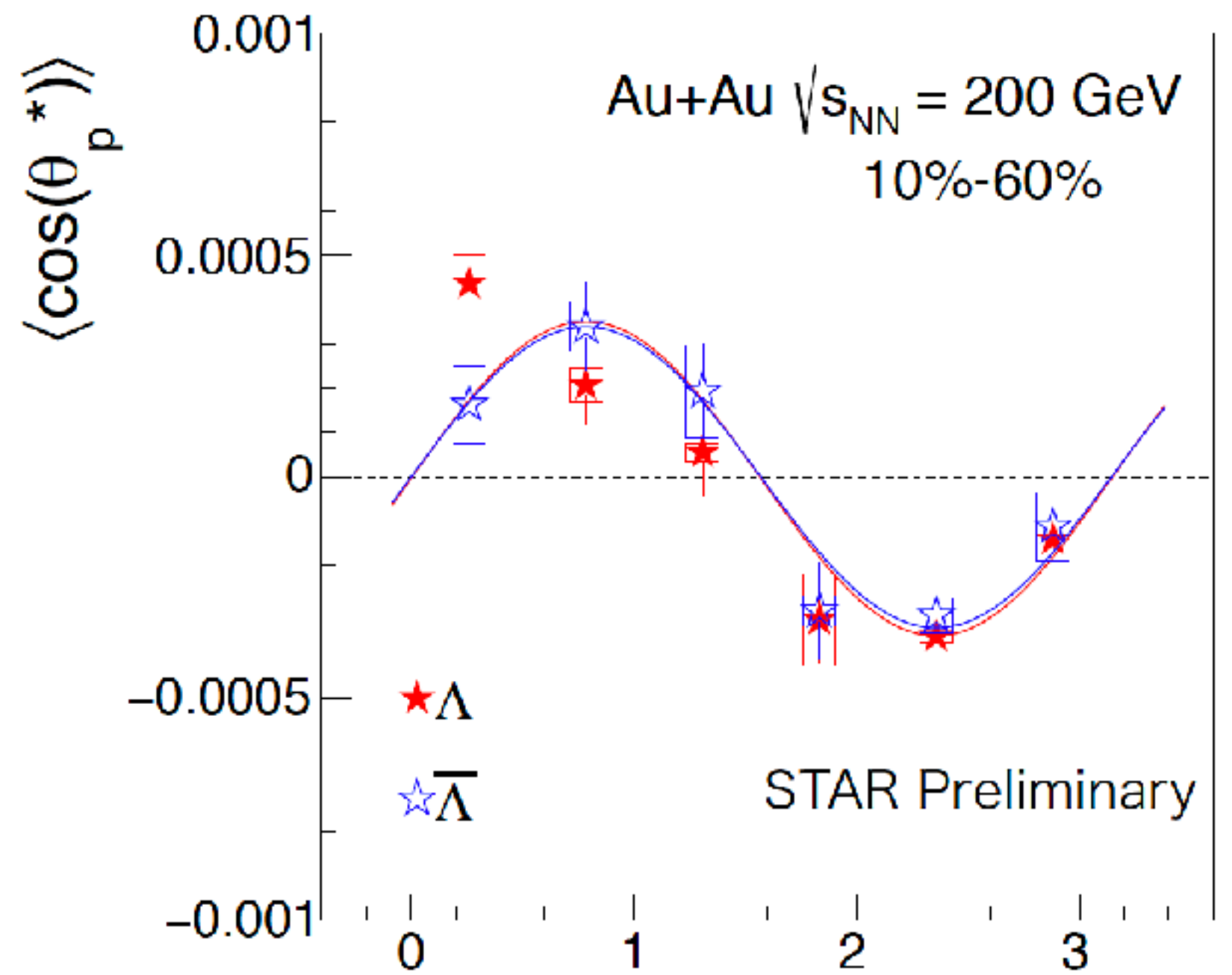


Slide from Xu Guang Huang's QM19 Plenary Talk

Assuming equilibrium of the spin degrees of freedom, global polarization rate can be well understood by theo. models.

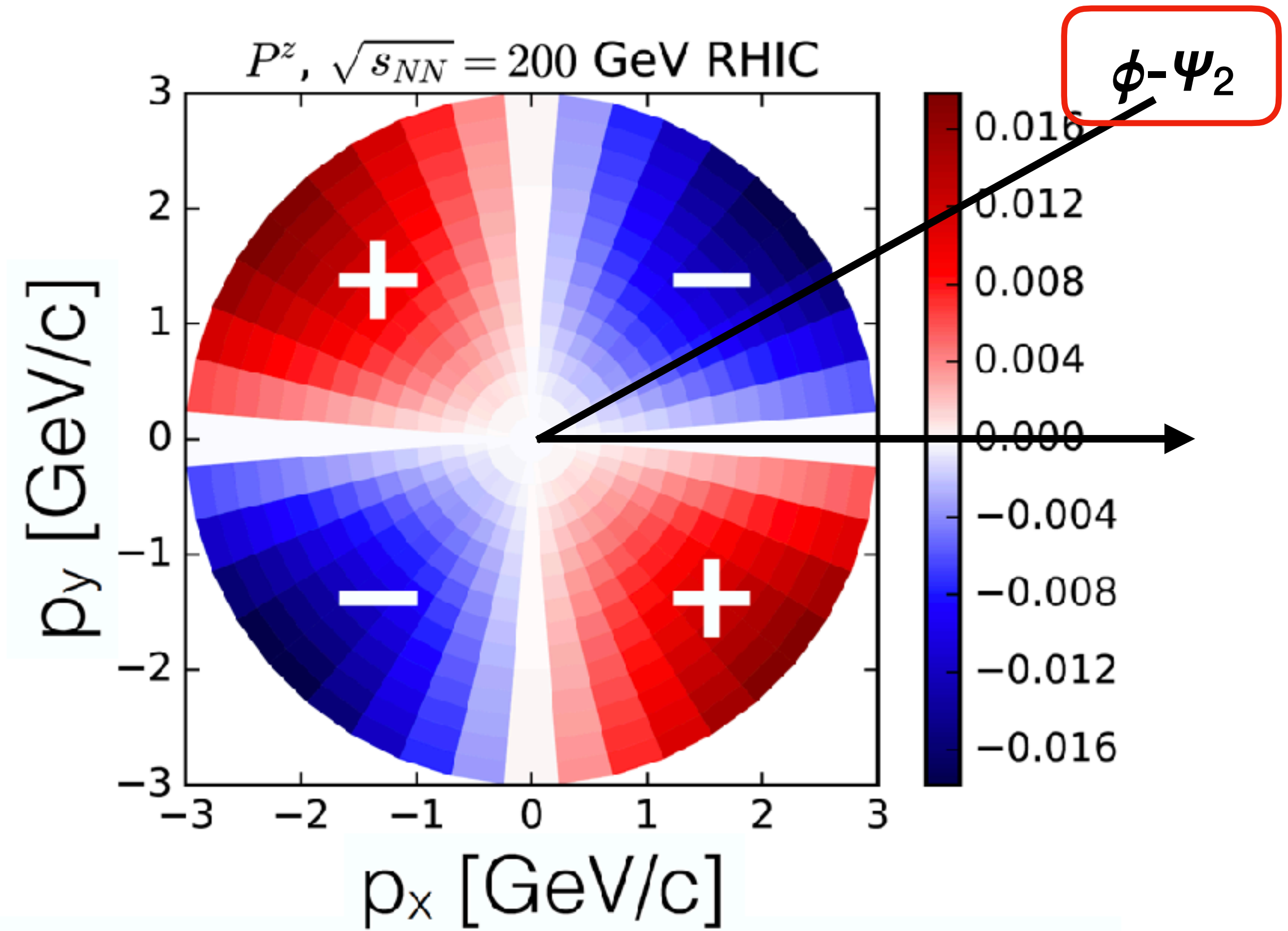
$$\varpi^{\mu\nu} \equiv \frac{1}{2} \left(\partial_\nu \frac{u_\mu}{T} - \partial_\mu \frac{u_\nu}{T} \right)$$

Local longitudinal polarization due to collective flow



$\phi - \Psi_2$ [rad]

STAR

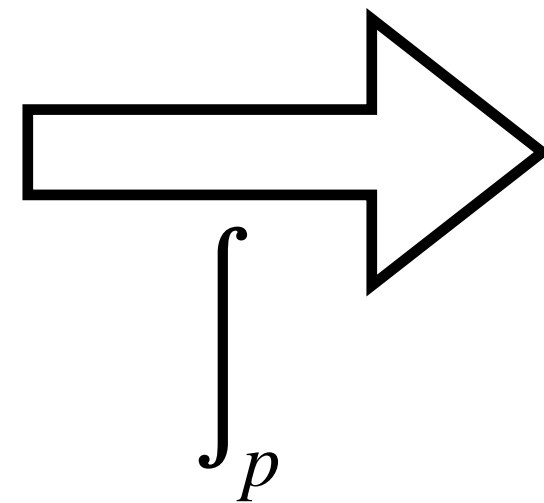


- (Hydro) F. Becattini & I. Karpenko, PRL 2018
 Similarly in other models

Microscopic:

distribution function

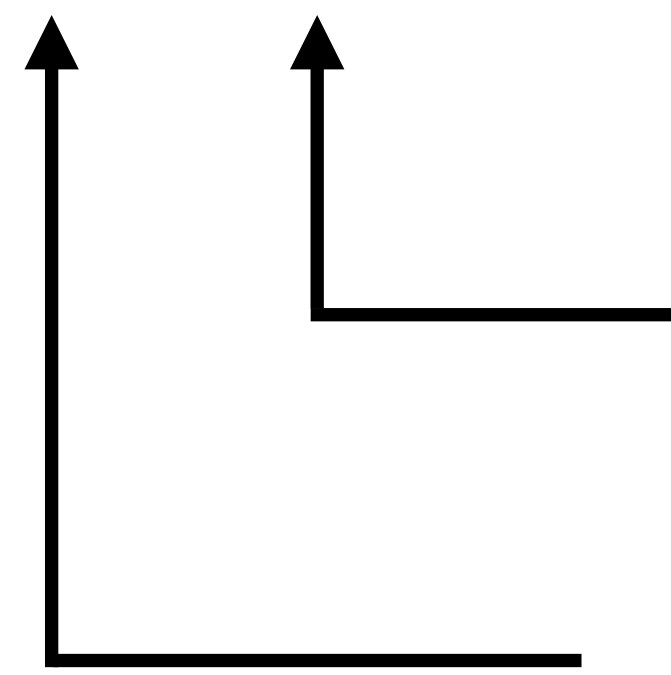
$$W_{ab}(x, p) \equiv \left\langle \int d^4y e^{i\frac{p \cdot y}{\hbar}} \bar{\psi}_b(x + \frac{y}{2}) \psi_a(x - \frac{y}{2}) \right\rangle,$$



Macroscopic:

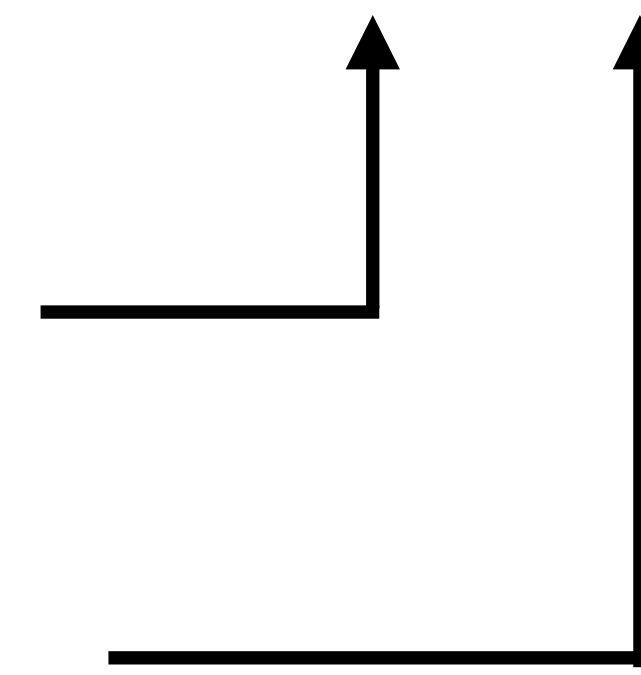
conserved (hydrodynamic) currents

$$J^\mu(x), \quad T^{\mu\nu}(x), \quad S^{\mu\nu\lambda}(x)$$



Vortical+Magnetic Effect

Non-Equilibrium Effect



Wigner Function

$$W_{ab}(x, p) \equiv \left\langle \int d^4y e^{i\frac{p \cdot y}{\hbar}} \bar{\psi}_b(x + \frac{y}{2}) \psi_a(x - \frac{y}{2}) \right\rangle,$$

a 4×4 matrix, decomposed in Clifford basis

$$W = \frac{1}{4} \left(\mathcal{F} + i\mathcal{P}\gamma^5 + \mathcal{V}_\mu \gamma^\mu + \mathcal{A}_\mu \gamma^5 \gamma^\mu + \frac{1}{2} \mathcal{L}_{\mu\nu} \sigma^{\mu\nu} \right),$$

Wigner Function

$$W_{ab}(x, p) \equiv \left\langle \int d^4y e^{i\frac{p \cdot y}{\hbar}} \bar{\psi}_b(x + \frac{y}{2}) \psi_a(x - \frac{y}{2}) \right\rangle,$$

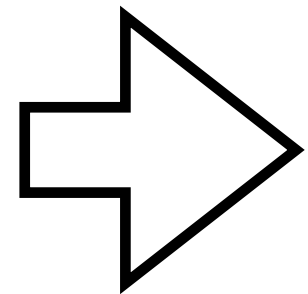
a 4x4 matrix, decomposed in Clifford basis

$$W = \frac{1}{4} \left(\mathcal{F} + i\mathcal{P}\gamma^5 + \mathcal{V}_\mu \gamma^\mu + \mathcal{A}_\mu \gamma^5 \gamma^\mu + \frac{1}{2} \mathcal{L}_{\mu\nu} \sigma^{\mu\nu} \right),$$

- Boltzmann equation
- Lorentz force
- chiral vortical effect + spin polarization
- chiral magnetic effect

Dirac Equation

$$\gamma^\mu (i\hbar \partial_\mu - QA_\mu) \psi(x) = 0$$



$$\mathcal{F}_\pm^\mu \equiv \frac{\mathcal{V}^\mu \pm \mathcal{A}^\mu}{2} = \left[p^\mu \delta(p^2) \pm \hbar \delta(p^2) \frac{\epsilon^{\mu\nu\lambda\rho} p_\nu n_\lambda}{2p \cdot n} \nabla_\rho \pm \hbar Q \delta'(p^2) \widetilde{F}^{\mu\nu} p_\nu \right] f_\pm,$$

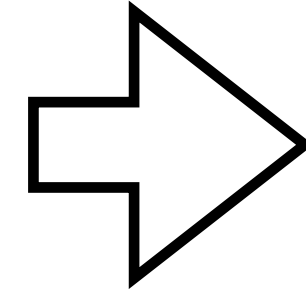
$$0 = \delta \left(p^2 \mp \hbar Q \frac{p \cdot B}{p \cdot n} \right) \left[p \cdot \partial - Q p^\mu F_{\mu\nu} \partial_p^\nu \pm \hbar \left(\partial_\mu \frac{\epsilon^{\mu\nu\lambda\rho} p_\nu n_\lambda}{2p \cdot n} \right) \nabla_\rho \right. \\ \left. \mp \hbar Q \frac{\epsilon^{\mu\nu\lambda\rho} E_\mu n_\nu p_\lambda}{2(p \cdot n)^2} \nabla_\rho \pm \hbar \frac{Q}{2p \cdot n} p_\lambda \left(\partial_\sigma \widetilde{F}^{\lambda\nu} \right) n_\nu \partial_p^\sigma \right] f_\pm.$$

Wigner Function

$$W_{ab}(x, p) \equiv \left\langle \int d^4y e^{i\frac{p \cdot y}{\hbar}} \bar{\psi}_b(x + \frac{y}{2}) \psi_a(x - \frac{y}{2}) \right\rangle,$$

a 4x4 matrix, decomposed in Clifford basis

$$W = \frac{1}{4} \left(\mathcal{F} + i\mathcal{P}\gamma^5 + \mathcal{V}_\mu \gamma^\mu + \mathcal{A}_\mu \gamma^5 \gamma^\mu + \frac{1}{2} \mathcal{L}_{\mu\nu} \sigma^{\mu\nu} \right),$$



Hydrodynamic Quantities:

$$J^\mu \equiv \langle \bar{\psi} \gamma^\mu \psi \rangle = \int \frac{d^4p}{(2\pi)^4} \mathcal{V}^\mu,$$

$$T^{\mu\nu} \equiv \langle \bar{\psi} (i\gamma^\mu \partial^\nu) \psi \rangle = \int \frac{d^4p}{(2\pi)^4} p^\mu \mathcal{V}^\nu,$$

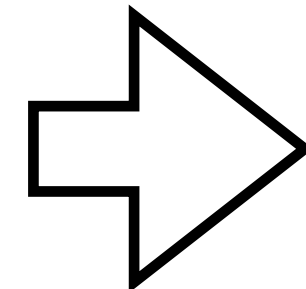
$$S^{\lambda\mu\nu} \equiv \frac{1}{4} \langle \bar{\psi} \{ \gamma^\lambda, \sigma^{\mu\nu} \} \psi \rangle = \frac{\epsilon^{\lambda\mu\nu\sigma}}{2} \int \frac{d^4p}{(2\pi)^4} \mathcal{A}_\sigma,$$

Wigner Function

$$W_{ab}(x, p) \equiv \left\langle \int d^4y e^{i\frac{p \cdot y}{\hbar}} \bar{\psi}_b(x + \frac{y}{2}) \psi_a(x - \frac{y}{2}) \right\rangle,$$

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$$W = \frac{1}{4} \left(\mathcal{F} + i\mathcal{P}\gamma^5 + \mathcal{V}_\mu \gamma^\mu + \mathcal{A}_\mu \gamma^5 \gamma^\mu + \frac{1}{2} \mathcal{L}_{\mu\nu} \sigma^{\mu\nu} \right),$$

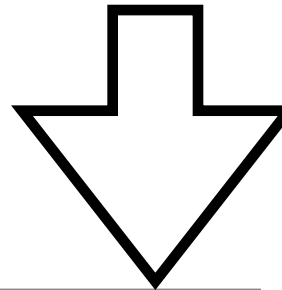


Hydrodynamic Quantities:

$$J^\mu \equiv \langle \bar{\psi} \gamma^\mu \psi \rangle = \int \frac{d^4p}{(2\pi)^4} \mathcal{V}^\mu,$$

$$T^{\mu\nu} \equiv \langle \bar{\psi} (i\gamma^\mu \partial^\nu) \psi \rangle = \int \frac{d^4p}{(2\pi)^4} p^\mu \mathcal{V}^\nu,$$

$$S^{\lambda\mu\nu} \equiv \frac{1}{4} \langle \bar{\psi} \{ \gamma^\lambda, \sigma^{\mu\nu} \} \psi \rangle = \frac{\epsilon^{\lambda\mu\nu\sigma}}{2} \int \frac{d^4p}{(2\pi)^4} \mathcal{A}_\sigma,$$



SHI, GALE, AND JEON

Substituting the distribution function in the definitions (13)–(15), we find the RH and LH particle currents and energy-momentum stress tensor,

$$T^{\mu\nu} = \epsilon u^\mu u^\nu - P \Delta^{\mu\nu} + \pi^{\mu\nu} + \frac{4\hbar}{5} \omega^\mu (v_+^\nu - v_-^\nu)$$

$$+ \frac{\hbar n_A}{4} (8\omega^\mu u^\nu + T \epsilon^{\mu\nu\sigma\lambda} \omega_{\sigma\lambda})$$

Chiral-viscous hydrodynamic equations

$$\begin{aligned} & \pm \frac{\hbar}{2} \epsilon^{\mu\rho\sigma\lambda} u_\rho \partial_\sigma \left(\frac{G_{4,1}^{(1),\pm}}{D_{3,1}^\pm} v_{\pm,\lambda} \right) \\ & \pm \frac{\hbar J_{2,2}^\pm}{4J_{4,2}^\pm} (\epsilon^{\mu\rho\sigma\lambda} u_\rho \sigma_\sigma^\xi \pi_{\lambda\xi} - \pi^{\mu\lambda} \omega_\lambda) \\ & \equiv n_\pm u^\mu + v_\pm^\mu + \hbar J_{\text{quantum},\pm}^\mu \end{aligned} \quad (67)$$

$$\begin{aligned} & + \frac{\hbar}{2} \epsilon^{\mu\rho\sigma\lambda} u_\rho u^\nu \partial_\sigma (v_\lambda^+ - v_\lambda^-) \\ & - \frac{\hbar}{10} \epsilon^{\mu\nu\rho\sigma} u_\rho (\partial_\sigma u^\lambda) (v_\lambda^+ - v_\lambda^-) \\ & + \frac{2\hbar}{5} \epsilon^{\mu\lambda\rho\sigma} u_\rho (\partial_\sigma u^\nu) (v_\lambda^+ - v_\lambda^-) \\ & \equiv \epsilon u^\mu u^\nu - \mathbf{SS}, \text{ Gale, and Jeon, Phys.Rev.C(2021)} \end{aligned}$$

Wigner Function

$$W_{ab}(x, p) \equiv \left\langle \int d^4y e^{i\frac{p \cdot y}{\hbar}} \bar{\psi}_b(x + \frac{y}{2}) \psi_a(x - \frac{y}{2}) \right\rangle,$$

Hydrodynamic Quantities:

$$J^\mu \equiv \langle \bar{\psi} \gamma^\mu \psi \rangle = \int \frac{d^4p}{(2\pi)^4} \mathcal{V}^\mu,$$

many other references:

SYSU, Fudan, USTC, Shandong Univ., Tsinghua, ...
Cracow, Frankfurt, Florence, ...

SHI, GALE, AND JEON

Substituting the distribution function in the definitions (13)–(15), we find the RH and LH particle currents and energy-momentum stress tensor,

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Chiral-viscous hydrodynamic equations

$$\pm \frac{\hbar}{2} \epsilon^{\mu\rho\sigma\lambda} u_\rho \partial_\sigma \left(\frac{G_{4,1}^{(1),\pm}}{D_{3,1}^\pm} v_{\pm,\lambda} \right)$$

$$\pm \frac{\hbar J_{2,2}^\pm}{4J_{4,2}^\pm} (\epsilon^{\mu\rho\sigma\lambda} u_\rho \sigma_\sigma^\xi \pi_{\lambda\xi} - \pi^{\mu\lambda} \omega_\lambda)$$

$$\equiv n_\pm u^\mu + v_\pm^\mu + \hbar J_{\text{quantum},\pm}^\mu$$

(67)

$$+ \left[\frac{\hbar}{2} \epsilon^{\mu\rho\sigma\lambda} u_\rho u^\nu \partial_\sigma (v_\lambda^+ - v_\lambda^-) - \frac{\hbar}{10} \epsilon^{\mu\nu\rho\sigma} u_\rho (\partial_\sigma u^\lambda) (v_\lambda^+ - v_\lambda^-) + \frac{2\hbar}{5} \epsilon^{\mu\lambda\rho\sigma} u_\rho (\partial_\sigma u^\nu) (v_\lambda^+ - v_\lambda^-) \right] \pi_{\lambda\xi}$$

$$+ \frac{\hbar}{2} \epsilon^{\mu\rho\sigma\lambda} u_\rho u^\nu \partial_\sigma (v_\lambda^+ - v_\lambda^-)$$

$$- \frac{\hbar}{10} \epsilon^{\mu\nu\rho\sigma} u_\rho (\partial_\sigma u^\lambda) (v_\lambda^+ - v_\lambda^-)$$

$$+ \frac{2\hbar}{5} \epsilon^{\mu\lambda\rho\sigma} u_\rho (\partial_\sigma u^\nu) (v_\lambda^+ - v_\lambda^-)$$

$$\equiv \varepsilon u^\mu u^\nu - \mathbf{SS}, \text{ Gale, and Jeon, Phys.Rev.C(2021)}$$

a different approach: conservation law + second law of thermodynamics

Fate of spin polarization in a relativistic fluid: An entropy-current analysis

Koichi Hattori, M.Hongo, Xu-Guang Huang, M.Matsuo, H.Taya, Phys.Lett.B 795 (2019) 100-106

Spin hydrodynamics and symmetric energy-momentum tensors – A current induced by the spin vorticity

Kenji Fukushima, Shi Pu, Phys.Lett.B 817 (2021) 136346

Cross effects in spin hydrodynamics: Entropy analysis and statistical operator

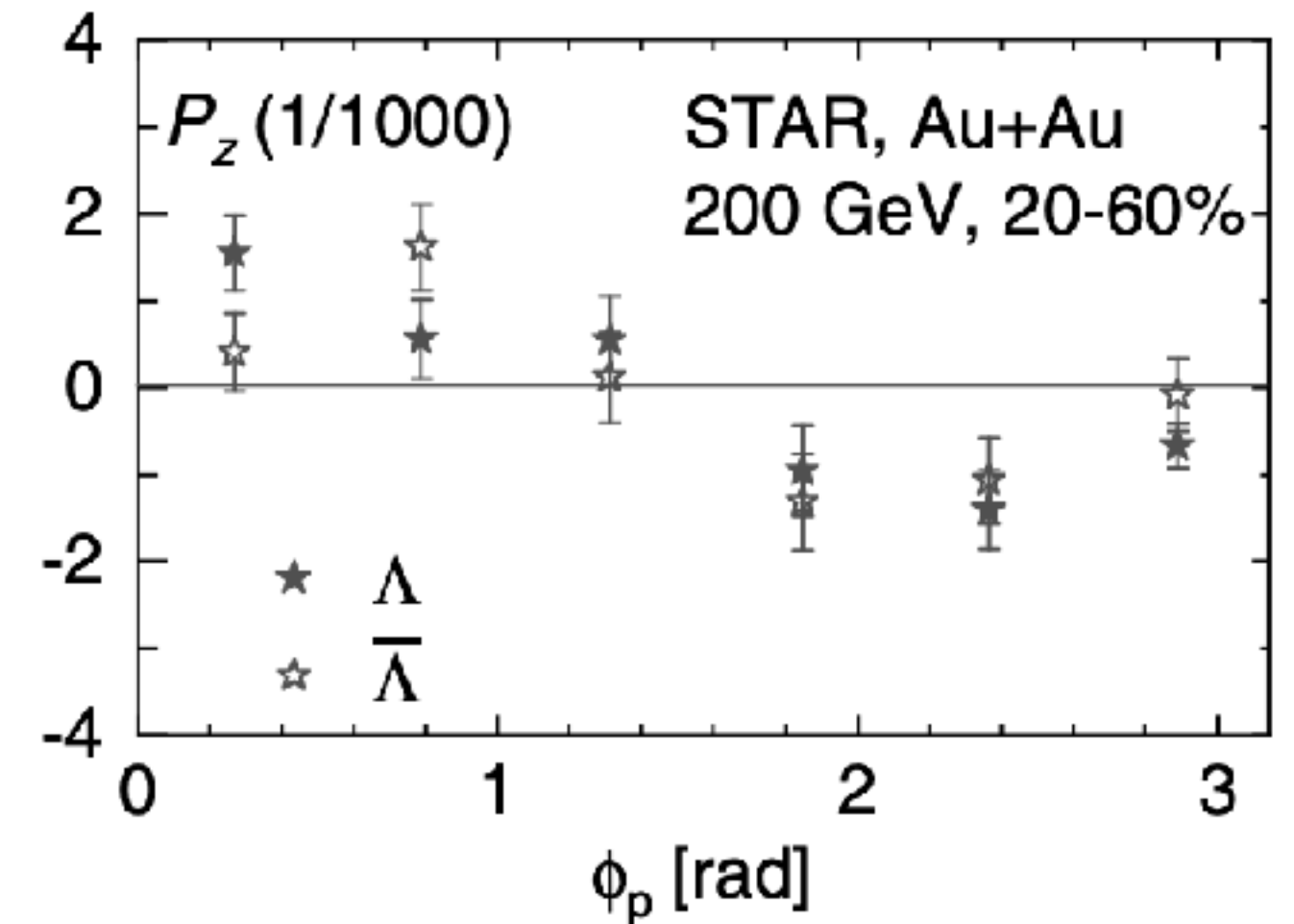
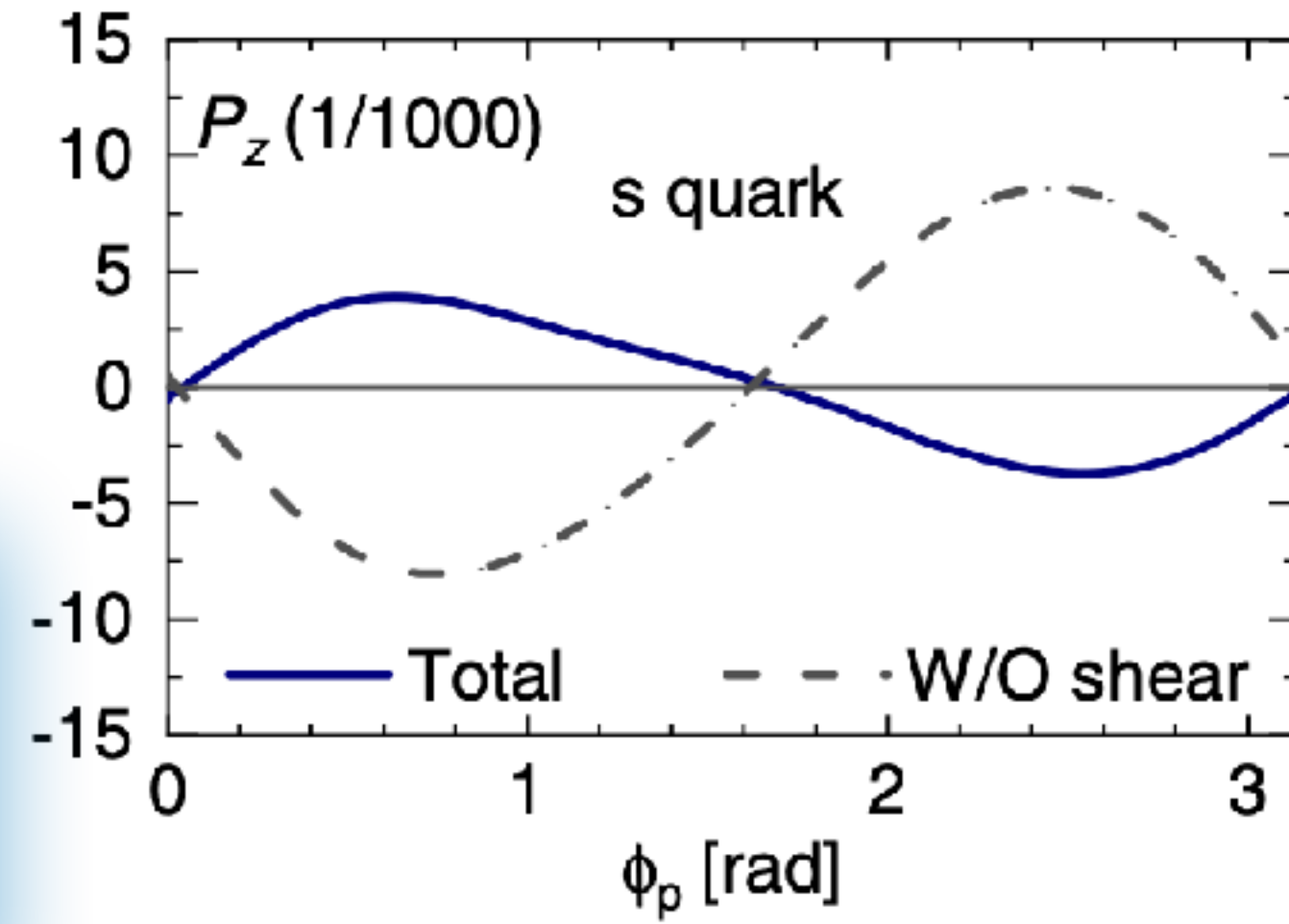
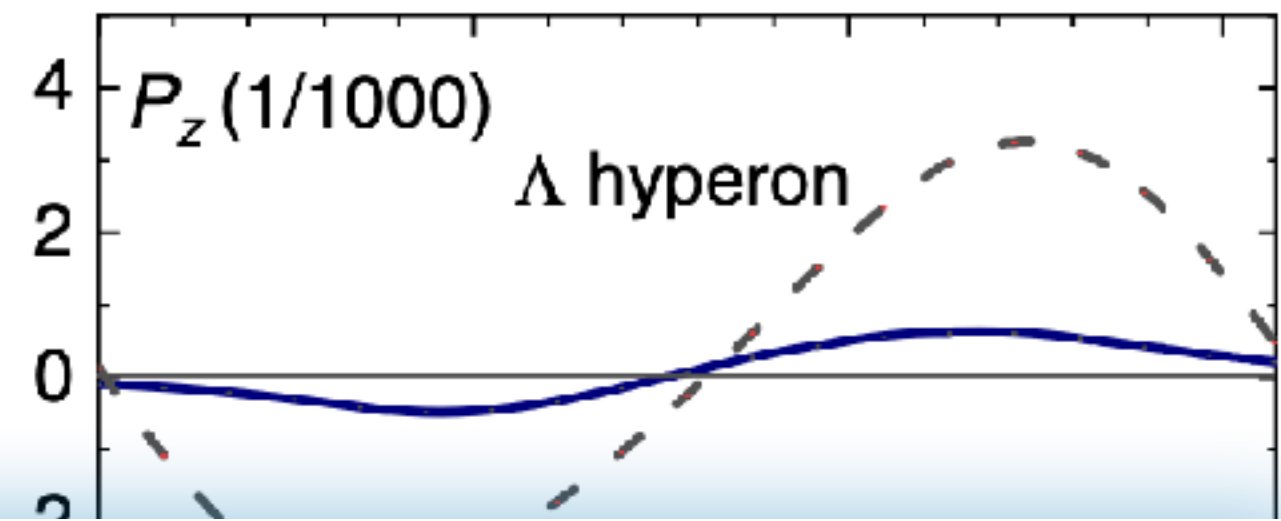
Jin Hu, PhysRevC.107.024915

12.17 afternoon, II, 国兴厅

17:30-17:50

Spin hydrodynamics, entropy principle and fluctuations

Lixin Yang



$$\mathcal{A}_{\text{SIP}}^\mu = -\beta n_0(1 - n_0) \frac{1}{\varepsilon_0} \epsilon^{\mu\nu\alpha\rho} u_\nu p_\rho p^\lambda \partial_{(\alpha}^\perp u_{\lambda)}$$

$$= -\beta n_0(1 - n_0) \frac{p_\perp^2}{\varepsilon_0} \epsilon^{\mu\nu\alpha\rho} u_\nu Q_\alpha^\lambda \sigma_{\rho\lambda}.$$

Spin polarization induced by the hydrodynamic gradients

Shuai Y.F. Liu, Yi Yin, JHEP 07 (2021) 188

Shear-Induced Spin Polarization in Heavy-Ion Collisions

Baochi Fu, Shuai Y.F. Liu, Long-Gang Pang, Huichao Song, Yi Yin, PhysRevLett.127.142301

Local Polarization and Isothermal Local Equilibrium in Relativistic Heavy Ion Collisions

F. Becattini, M. Buzzegoli, G. Inghirami, I. Karpenko, A. Palermo, PhysRevLett.127.272302

Spin-thermal shear coupling in a relativistic fluid

F. Becattini, M. Buzzegoli, A. Palermo, Phys.Lett.B 820 (2021) 136519

Experiment

12.17 afternoon, II, 国兴厅

15:20-15:40	Measurements of Global and Local Polarization of Hyperons in Heavy Ion Collisions from STAR	Xingrui Gou
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Theory

12.17 afternoon, II, 国兴厅

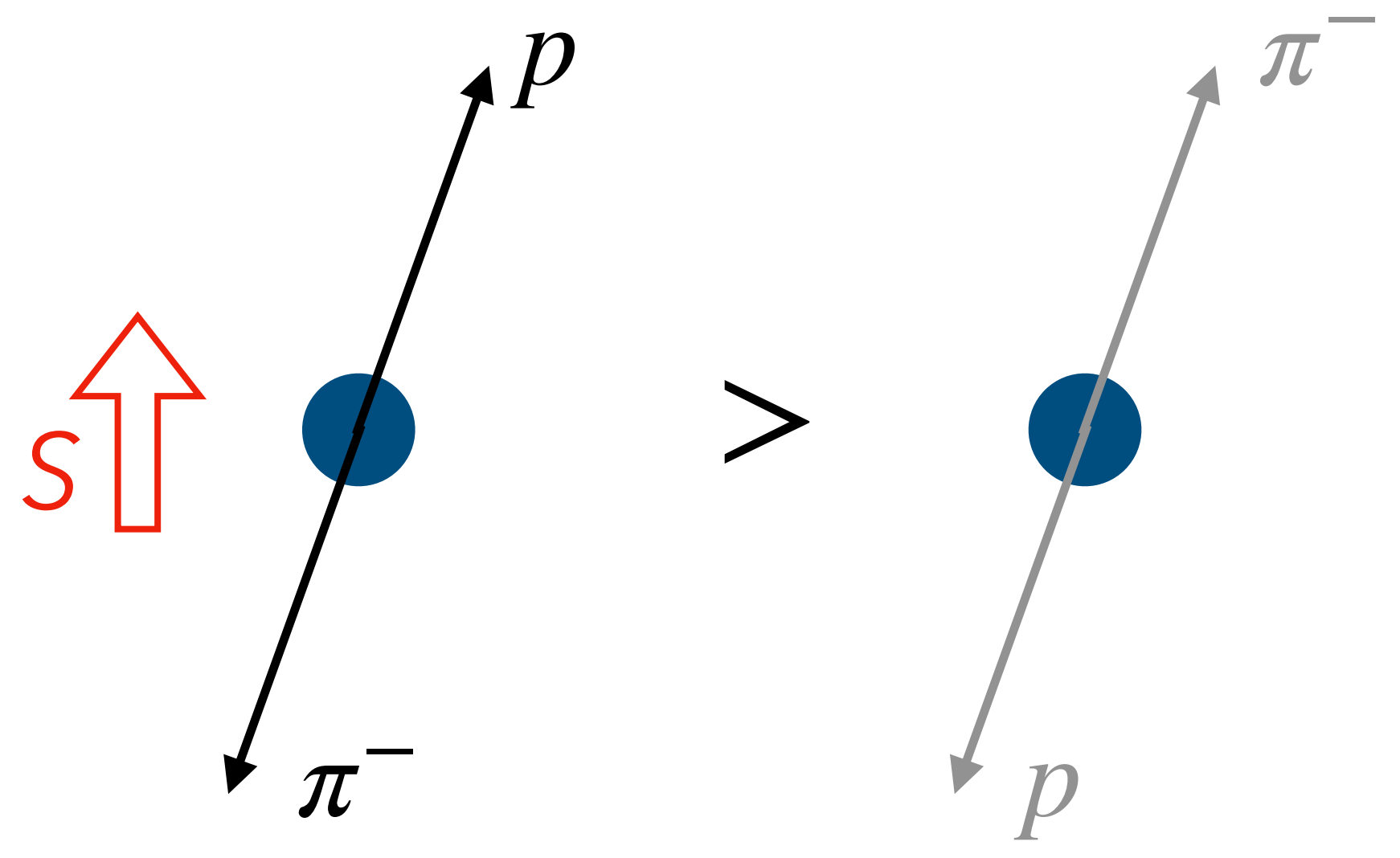
14:40-15:00	Spin Polarization, Anomalous Magnetic Moment and Transportation in the magnetized QCD background	Shengqin Feng
16:10-16:30	A topological realization of spin polarization through vortex formation in collisions of Bose-Einstein condensates	Jian Deng
16:30-16:50	Polarized hadron production in unpolarized high energy collisions	Shu-yi Wei

12.17 afternoon, III, 国旺厅

17:30-17:50	The relation of global polarization and directed flow in Au+Au collisions at RHIC	Ze-Fang Jiang
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weak decay (Λ)

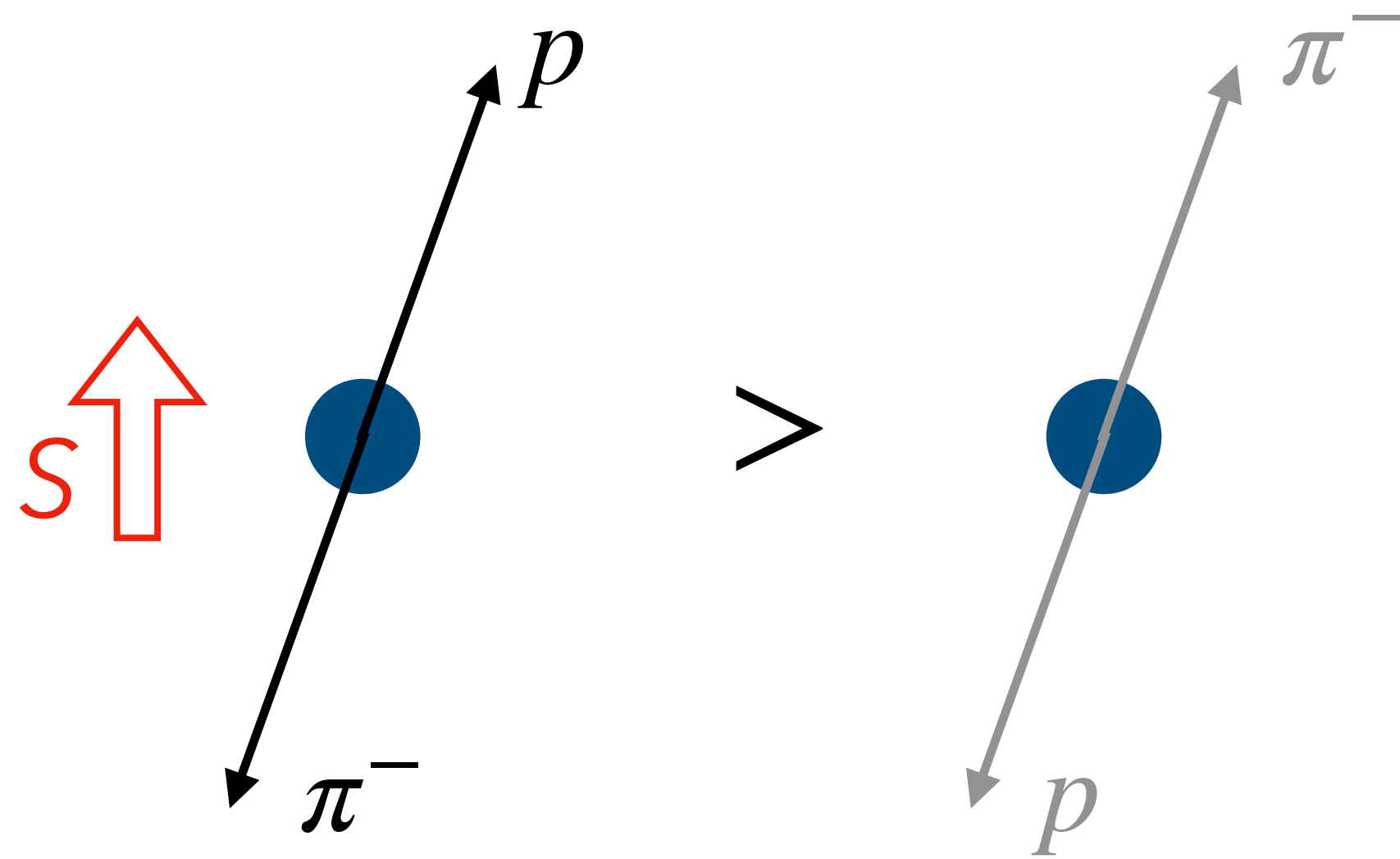
\not{p}



$$P \propto \frac{\omega}{T}$$

weak decay (Λ)

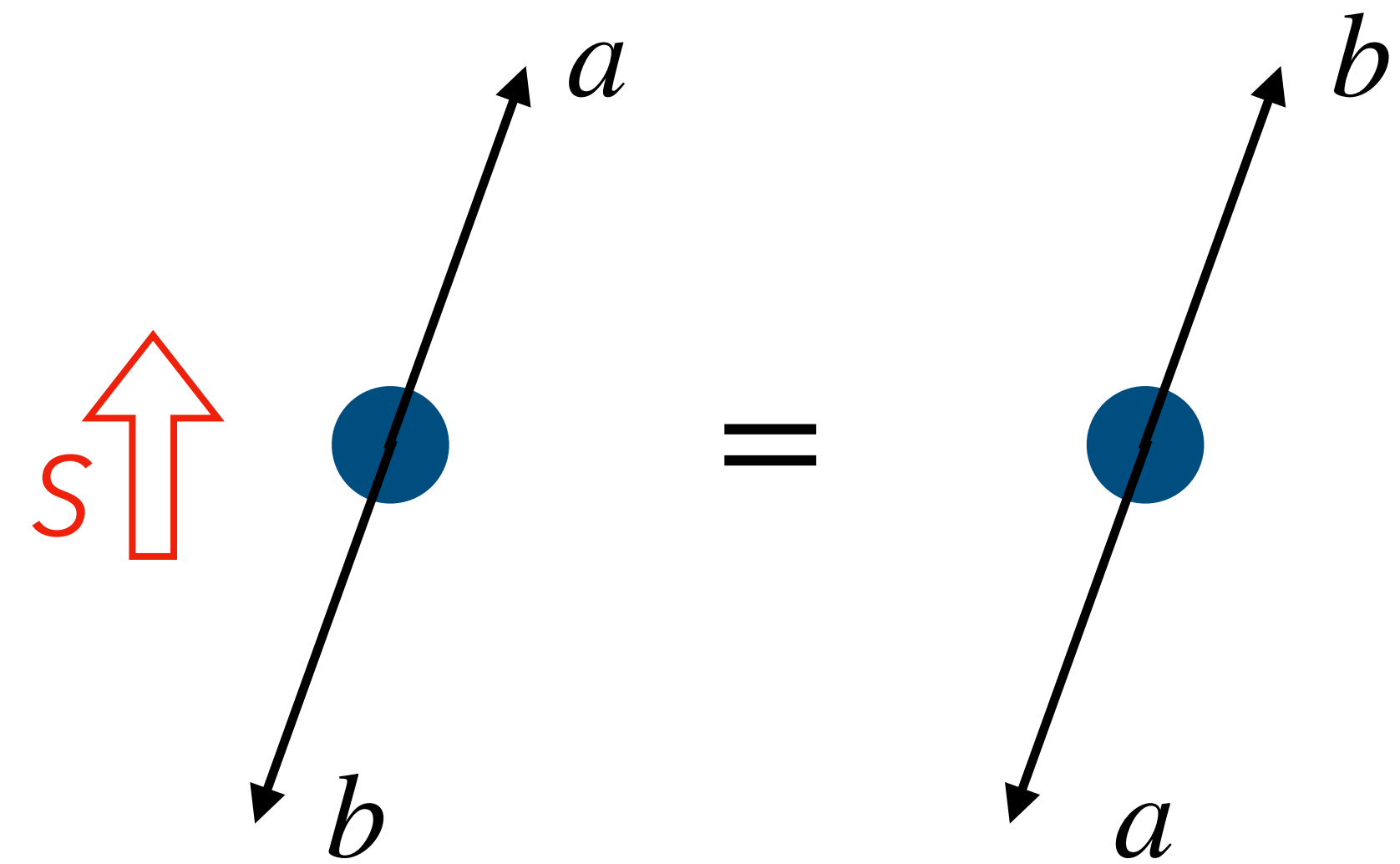
\not{P}



$$P \propto \frac{\omega}{T}$$

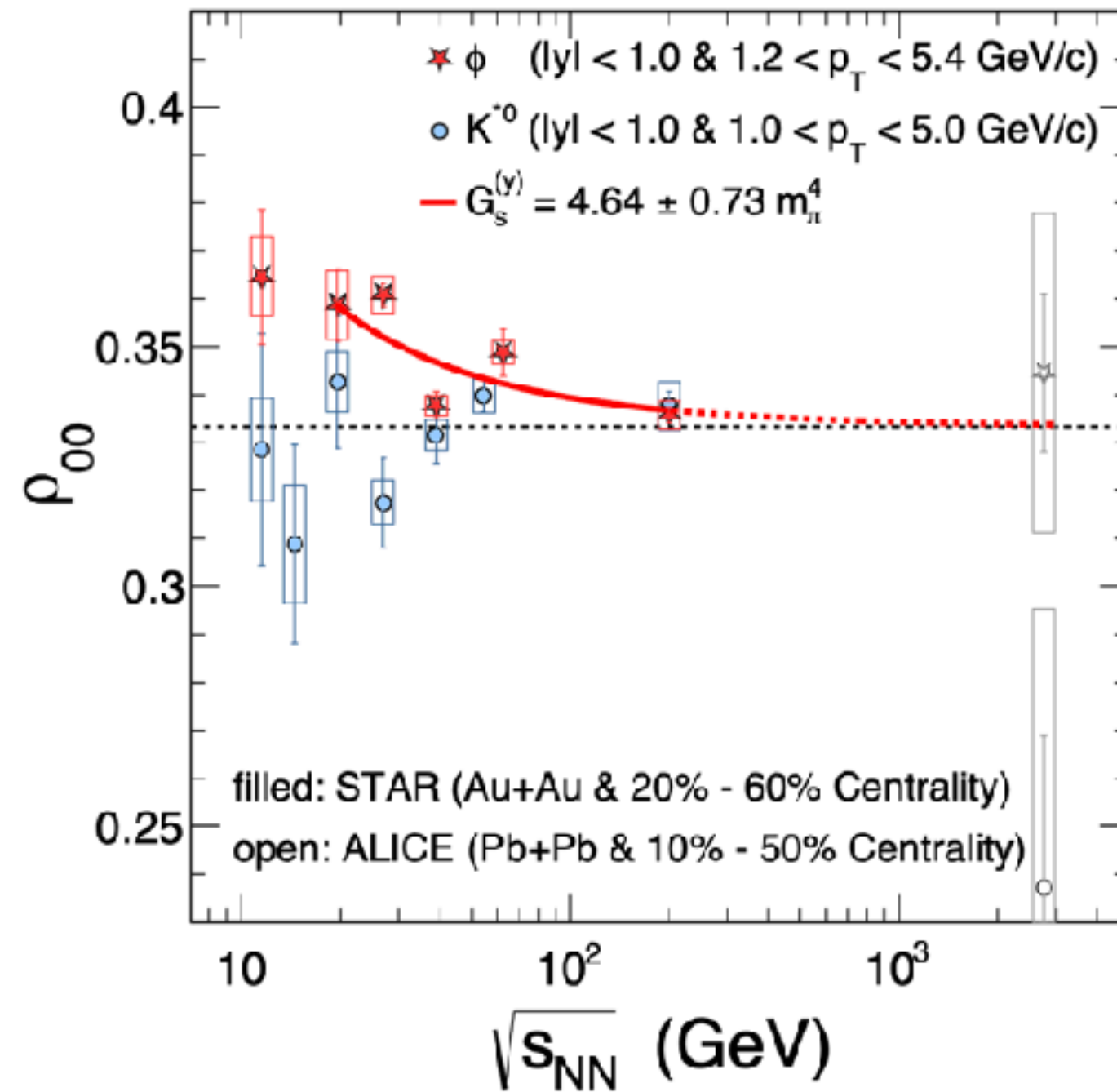
strong/EM decay ($\phi, \rho, K^*, J/\psi$)

P



vector mesons

$$\Delta\rho_{00} \propto \left(\frac{\omega}{T}\right)^2$$



STAR, Nature 614 (2023) 244

12.17 afternoon, II, 国兴厅

14:20-14:40	Vector meson polarization measurements in pp and Pb-Pb collisions with ALICE at the LHC	Xiaozhi Bai
15:00-15:20	Measurements of J/ψ polarization and global spin alignment in Ru+Ru and Zr+Zr collisions at $\sqrt{s_{NN}} = 200$ GeV from the STAR experiment	Qian Yang

meson field:

Spin Alignment of Vector Mesons in Heavy-Ion Collisions

Xin-Li Sheng, L. Oliva, Zuo-Tang Liang, Qun Wang, Xin-Nian Wang, PhysRevLett.131.042304

shear-induced modification of polarization tensor:

Tensor Polarization and Spectral Properties of Vector Meson in QCD Medium

Feng Li, Shuai Y.F. Liu, 2206.11890

12.17 afternoon, II, 国兴厅

14:00-14:20	Spin alignment for K mesons	Shi Pu
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12.18 morning, II, 国兴厅

9:00-9:20	Magneto-birefringence and axial Ward identity at finite temperature and density	Koichi Hattori
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12.17 afternoon, I, 国会厅

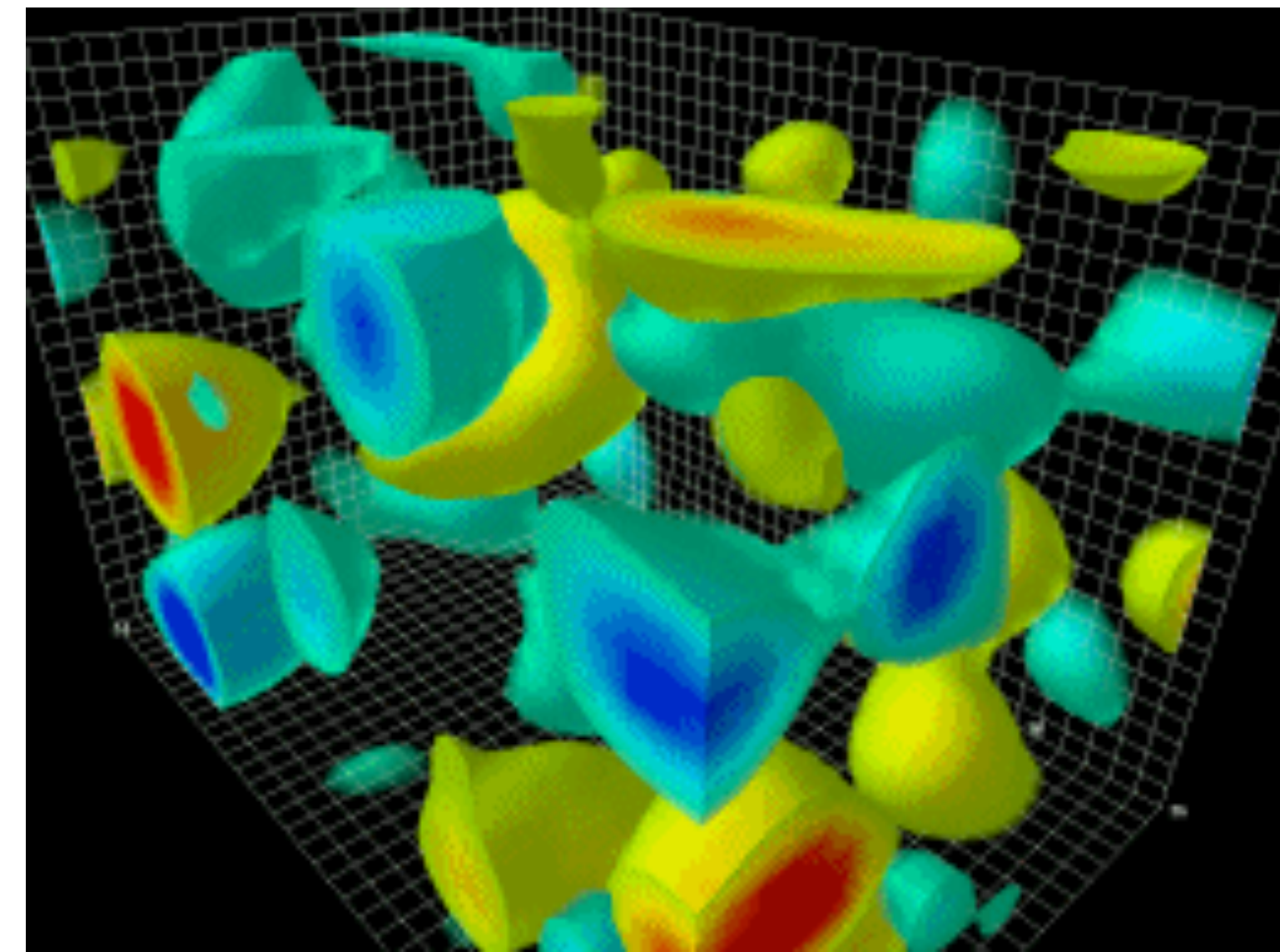
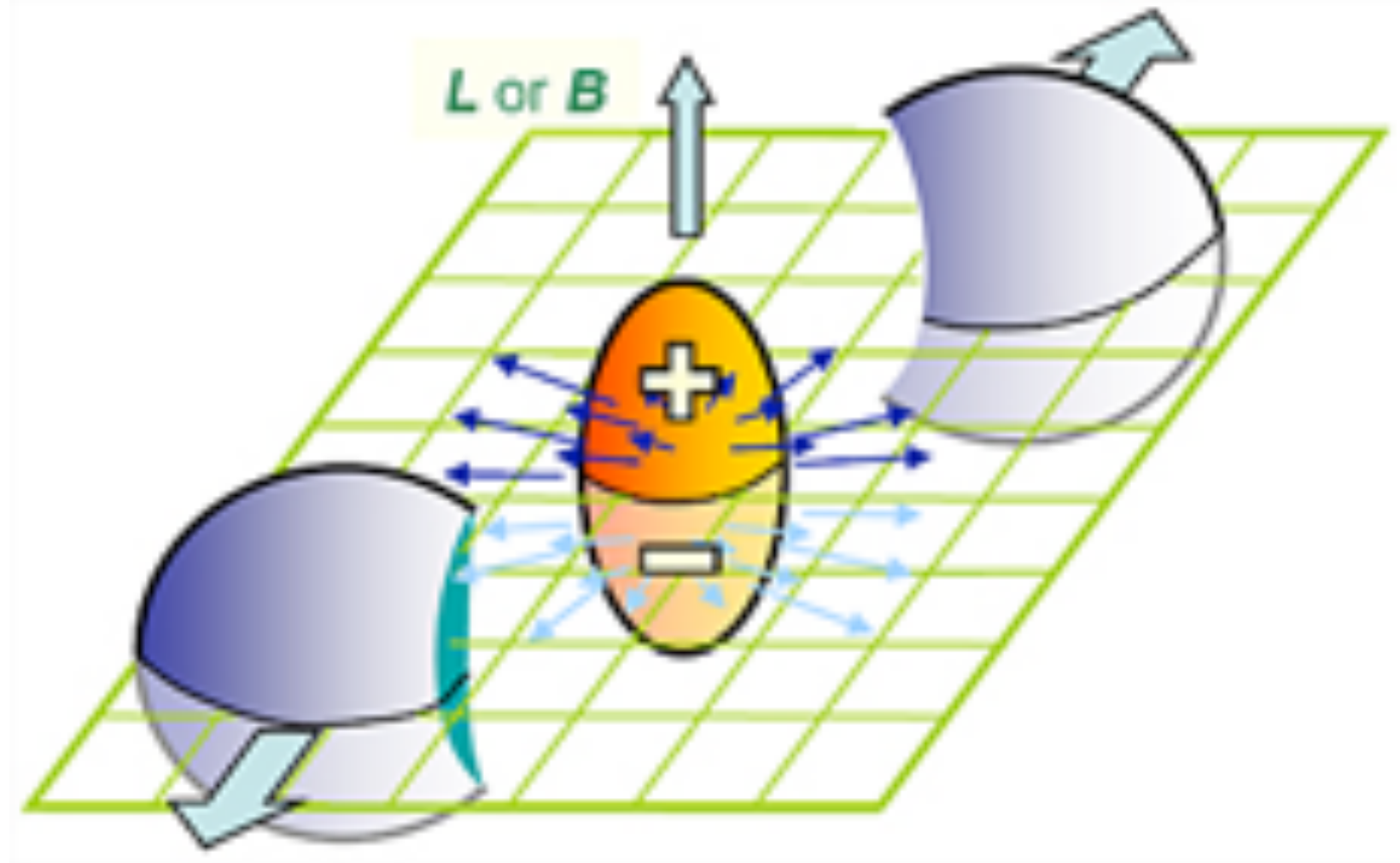
14:40-15:00	Fluctuations of conserved charges in strong magnetic fields in (2+1)- flavor QCD	Jinbiao Gu
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12.17 afternoon, II, 国兴厅

16:50-17:10	Thermodynamics for a Rotating Chiral Fermion System in the Uniform Magnetic Field	Ren-Hong Fang
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12.18 morning, II, 国兴厅

9:20-9:40	Chiral vortical catalysis	Yin Jiang
10:00-10:20	The role of neutral and charged pions in (inverse) magnetic catalysis and diamagnetism	Jie Mei



Magnetic Field

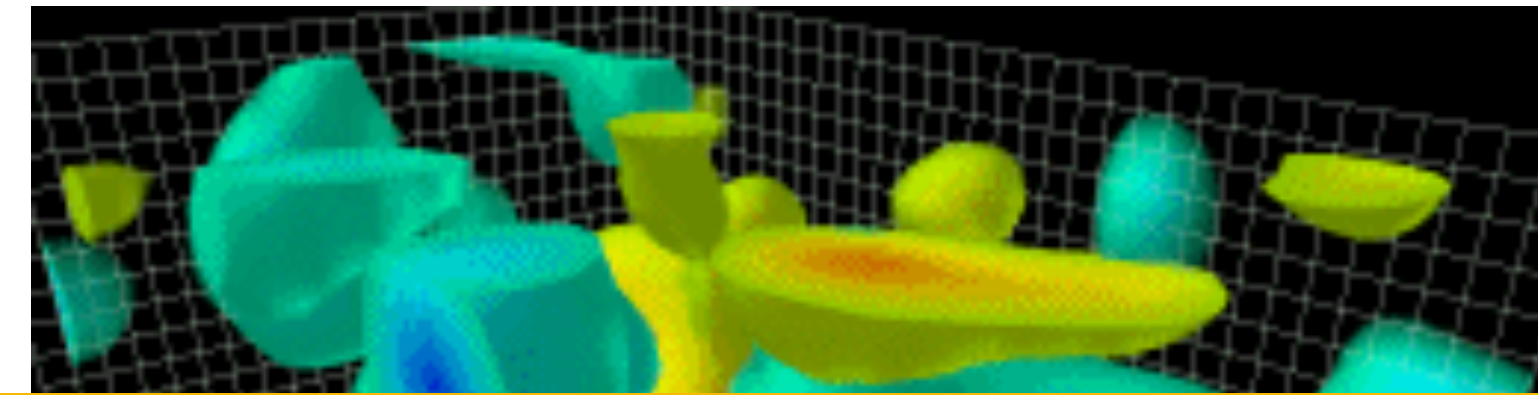
$$B \sim \gamma Z Q b / R^3$$

(10^{15} T @ 200 GeV Au-Au)



Topology
 ⊗ Chiral
 ⊗ Anomaly

$$Q_w = \frac{g^2}{32\pi^2} \int d^4x F_{\mu\nu}^a \tilde{F}_a^{\mu\nu}$$



The Chiral Magnetic Effect

Kenji Fukushima, Dmitri Kharzeev, H.Warringa, PhysRevD.78.074033

a recent review:

Chiral magnetic effect reveals the topology of gauge fields in heavy-ion collisions

Dmitri Kharzeev, Jinfeng Liao, Nature Review Physics 3 (2021) 1, 55-63

Magnetic Field

$$B \sim \gamma Z Q b / R^3$$

(10^{15} T @ 200 GeV Au-Au)



Topology

⊗ Chiral

⊗ Anomaly

$$Q_w = \frac{g^2}{32\pi^2} \int d^4x F_{\mu\nu}^a \tilde{F}_a^{\mu\nu}$$

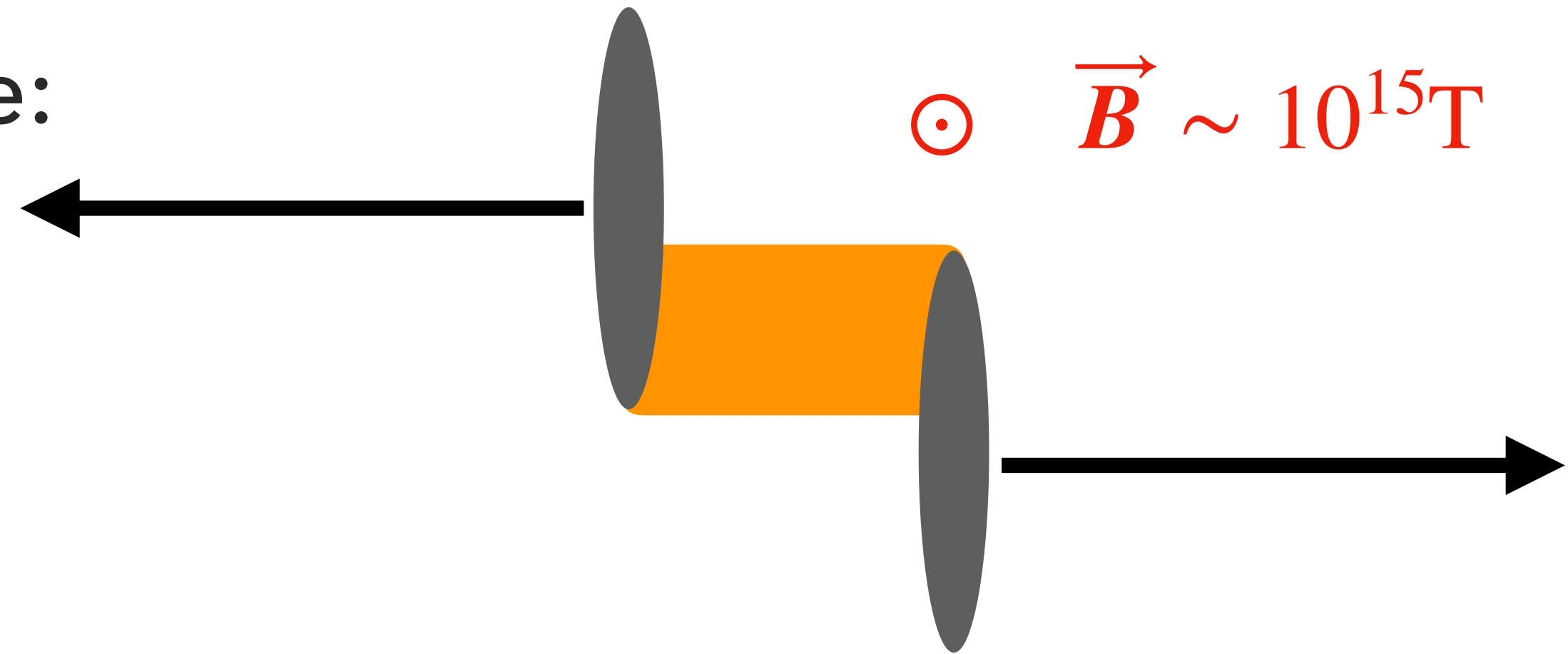
For right-handed particles w/ positive charge:

1. $\vec{p} \parallel \vec{S} \parallel \vec{\mu}$

2. Energy = $-\vec{\mu} \cdot \vec{B} \propto -\vec{p} \cdot \vec{B}$

⇒ lower energy if moving along B field direction

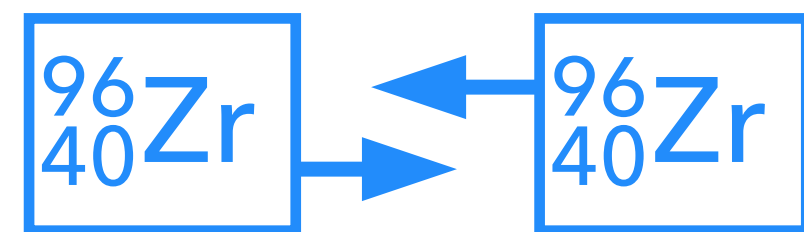
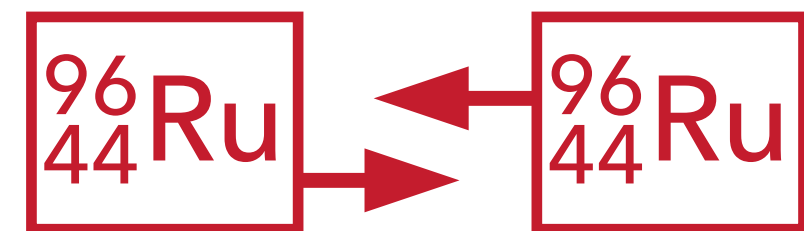
$$\vec{J} = \sigma_A \mu_A \vec{B}$$



expectation before the isobar collisions:

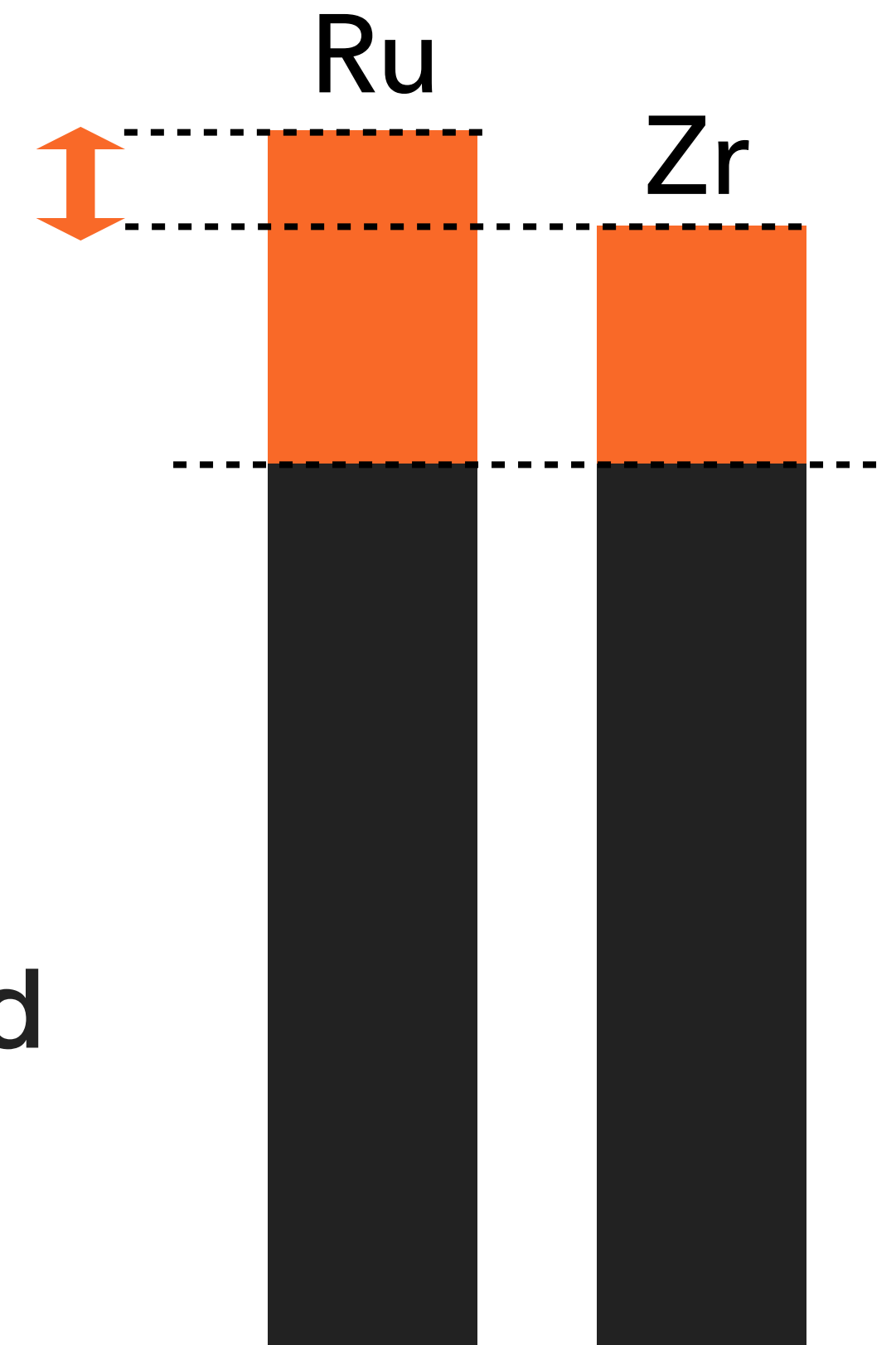
$Correlator[Ru] > Correlator[Zr]$ \longrightarrow **CME**

$Correlator[Ru] = Correlator[Zr]$ \longrightarrow **no CME**



Different Proton # \longrightarrow Different CME Signal

Same Baryon # \longrightarrow Same Background

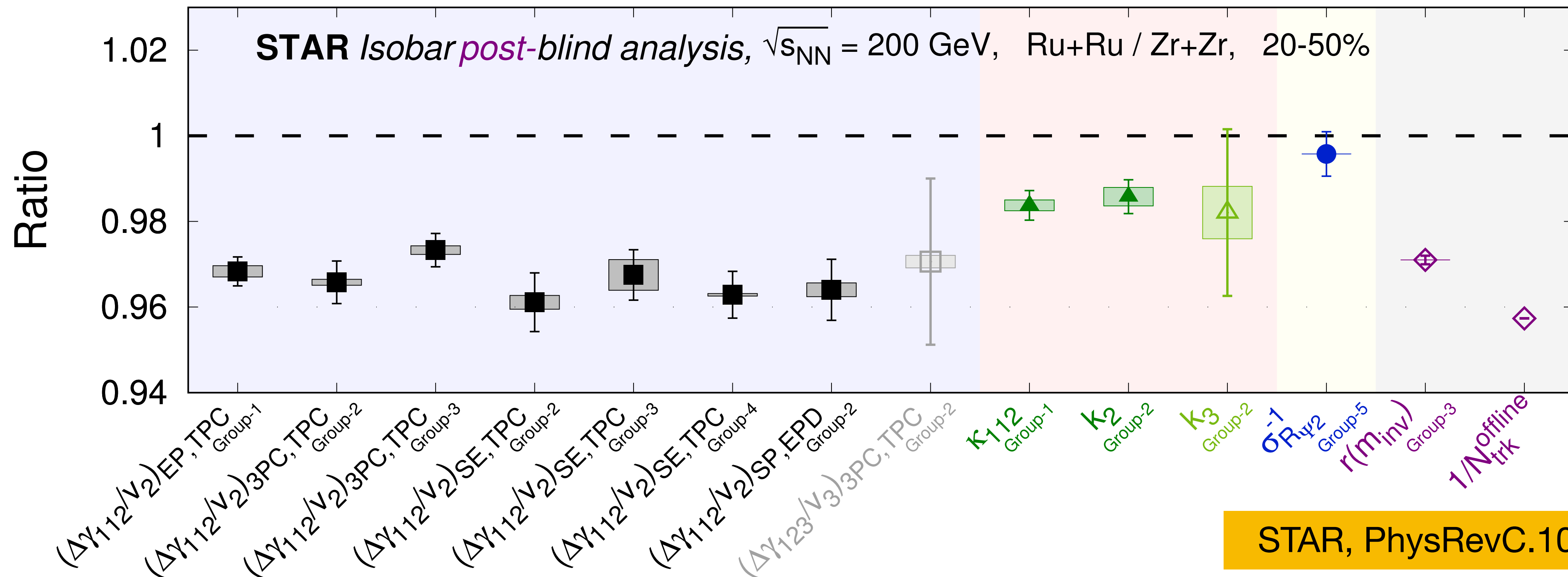


expectation before the isobar collisions:

$Correlator[Ru] > Correlator[Zr] \longrightarrow$ **CME**

$Correlator[Ru] = Correlator[Zr] \longrightarrow$ **no CME**

measurement: $Correlator[Ru] < Correlator[Zr]$

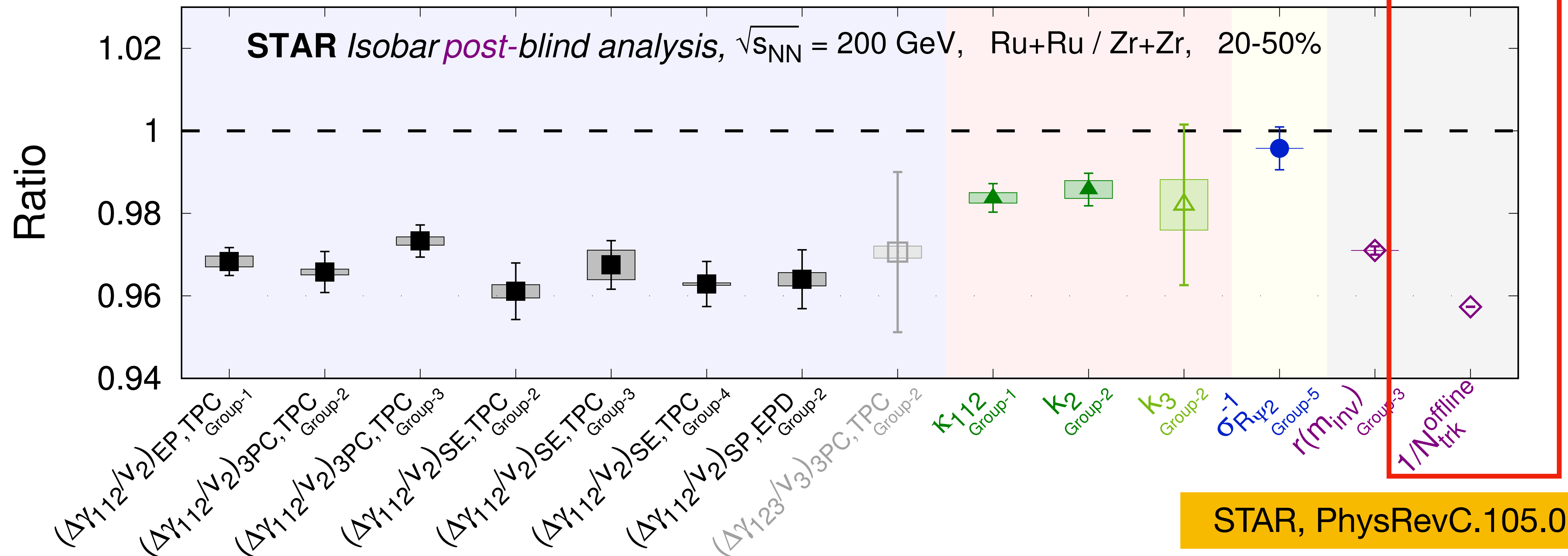


expectation before the isobar collisions:

$Correlator[Ru] > Correlator[Zr] \longrightarrow$ **CME**

$Correlator[Ru] = Correlator[Zr] \longrightarrow$ **no CME**

measurement: $Correlator[Ru] < Correlator[Zr]$

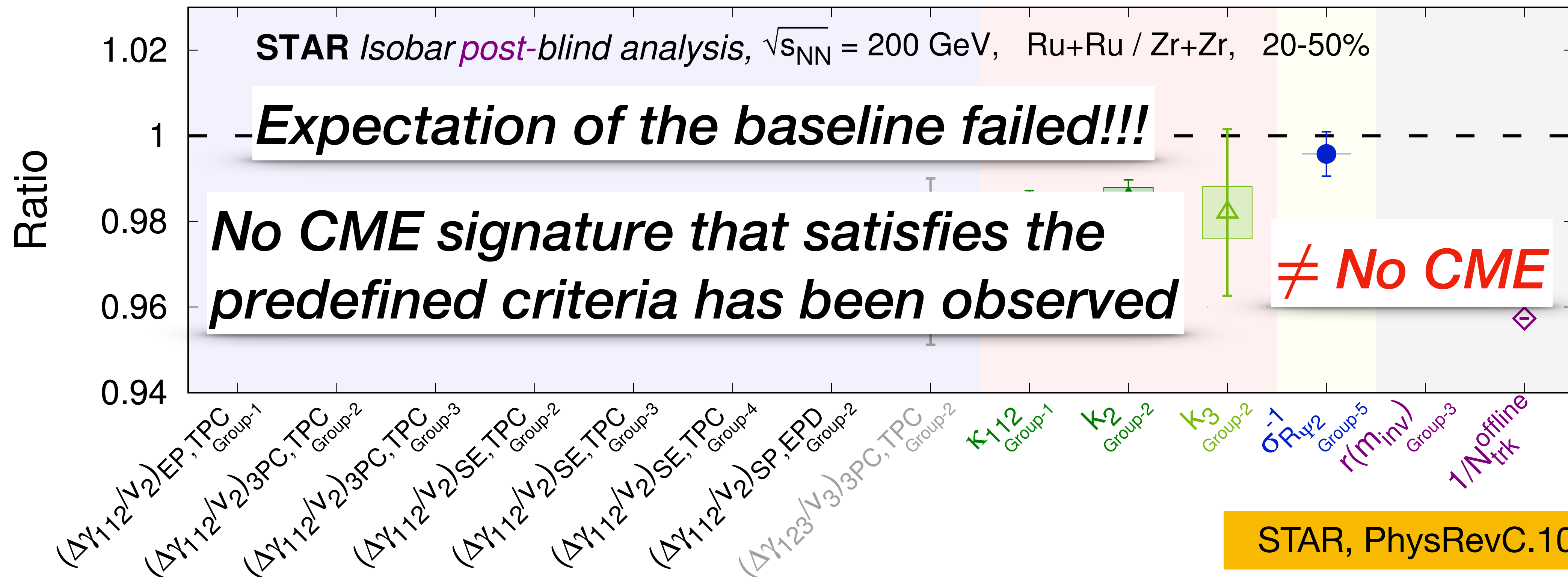


expectation before the isobar collisions:

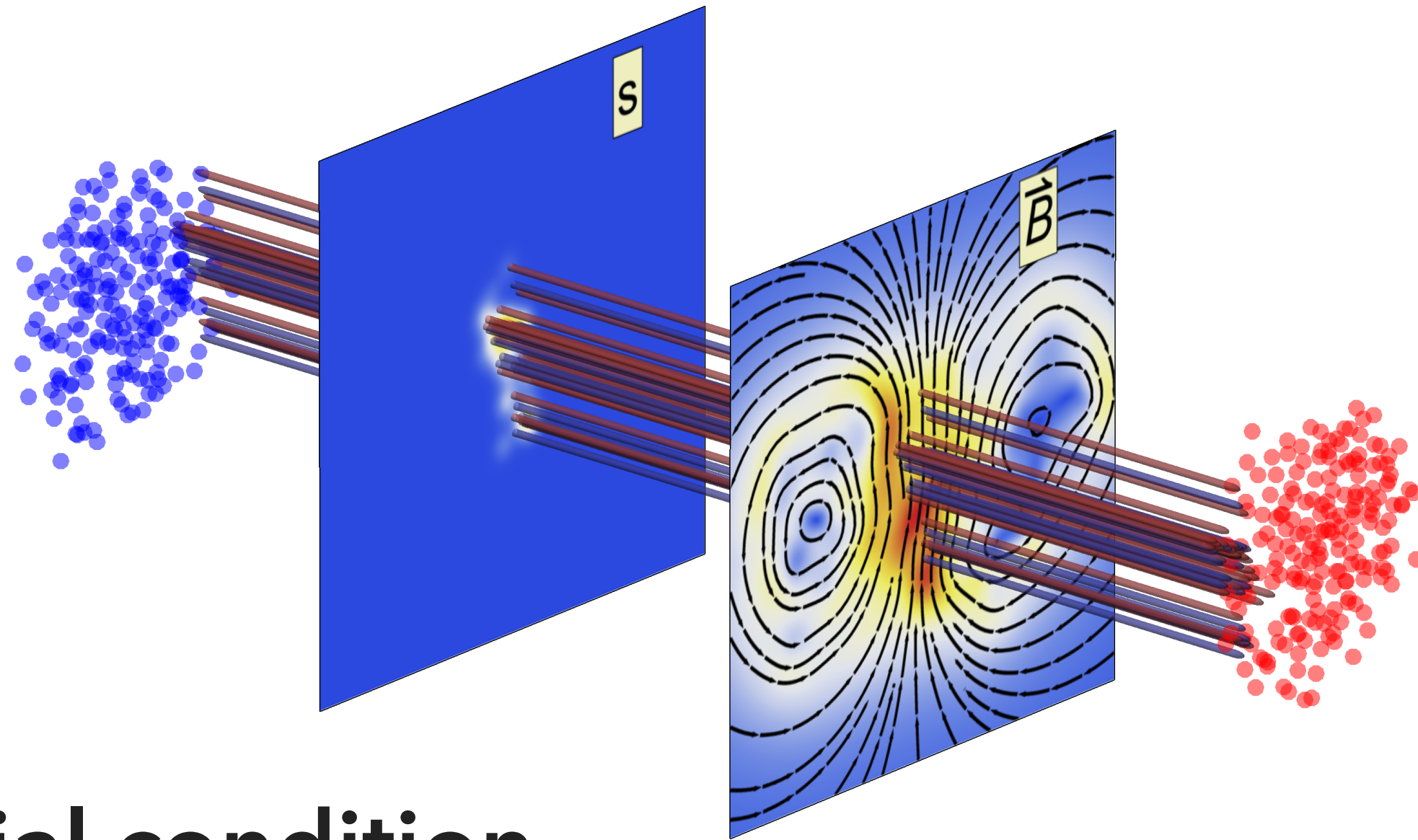
$Correlator[Ru] > Correlator[Zr] \longrightarrow$ **CME**

$Correlator[Ru] = Correlator[Zr] \longrightarrow$ **no CME**

measurement: $Correlator[Ru] < Correlator[Zr]$



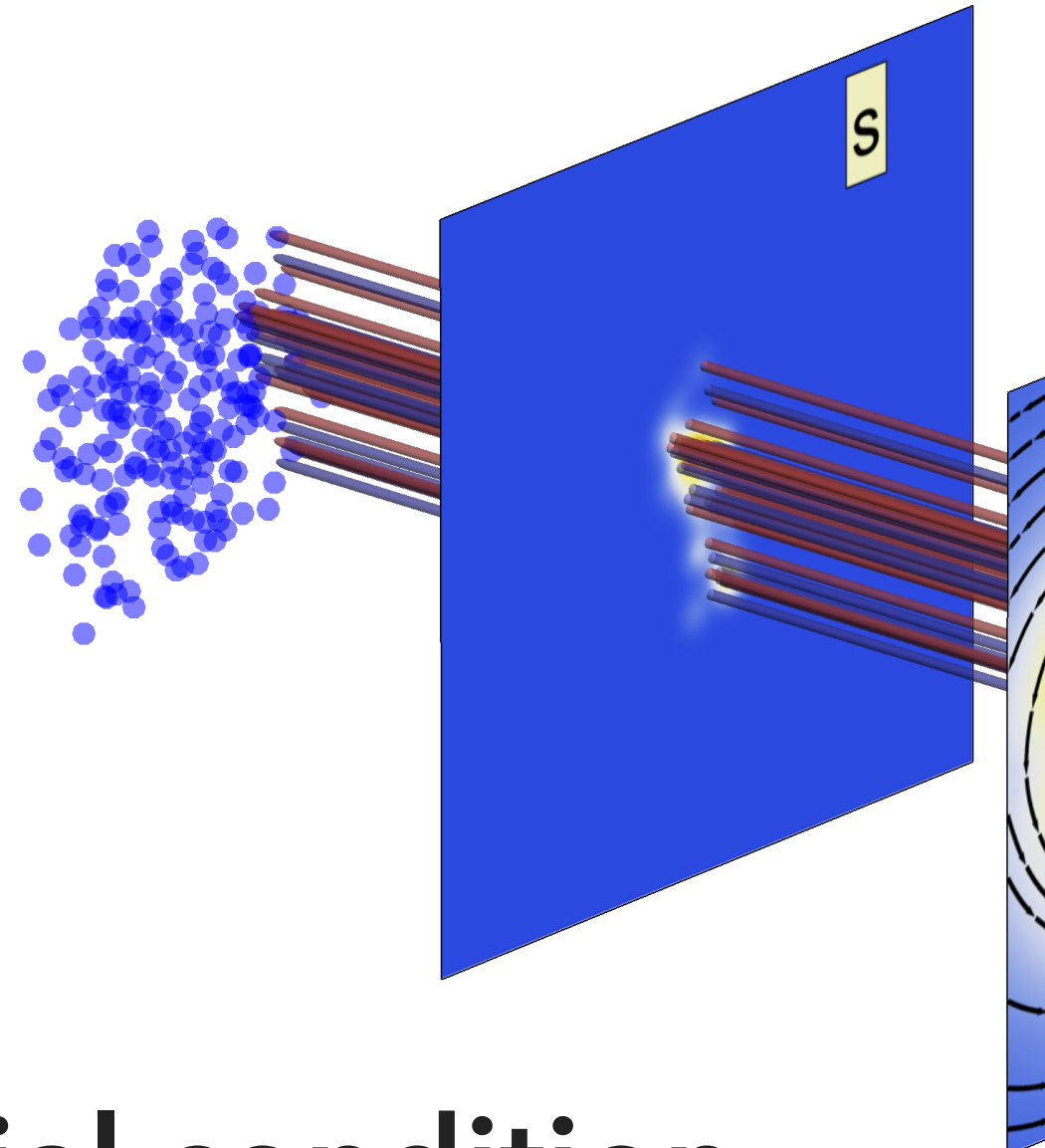
E-by-E IC of Bulk, n , $n5$, and B field



Initial condition
+ Driving force

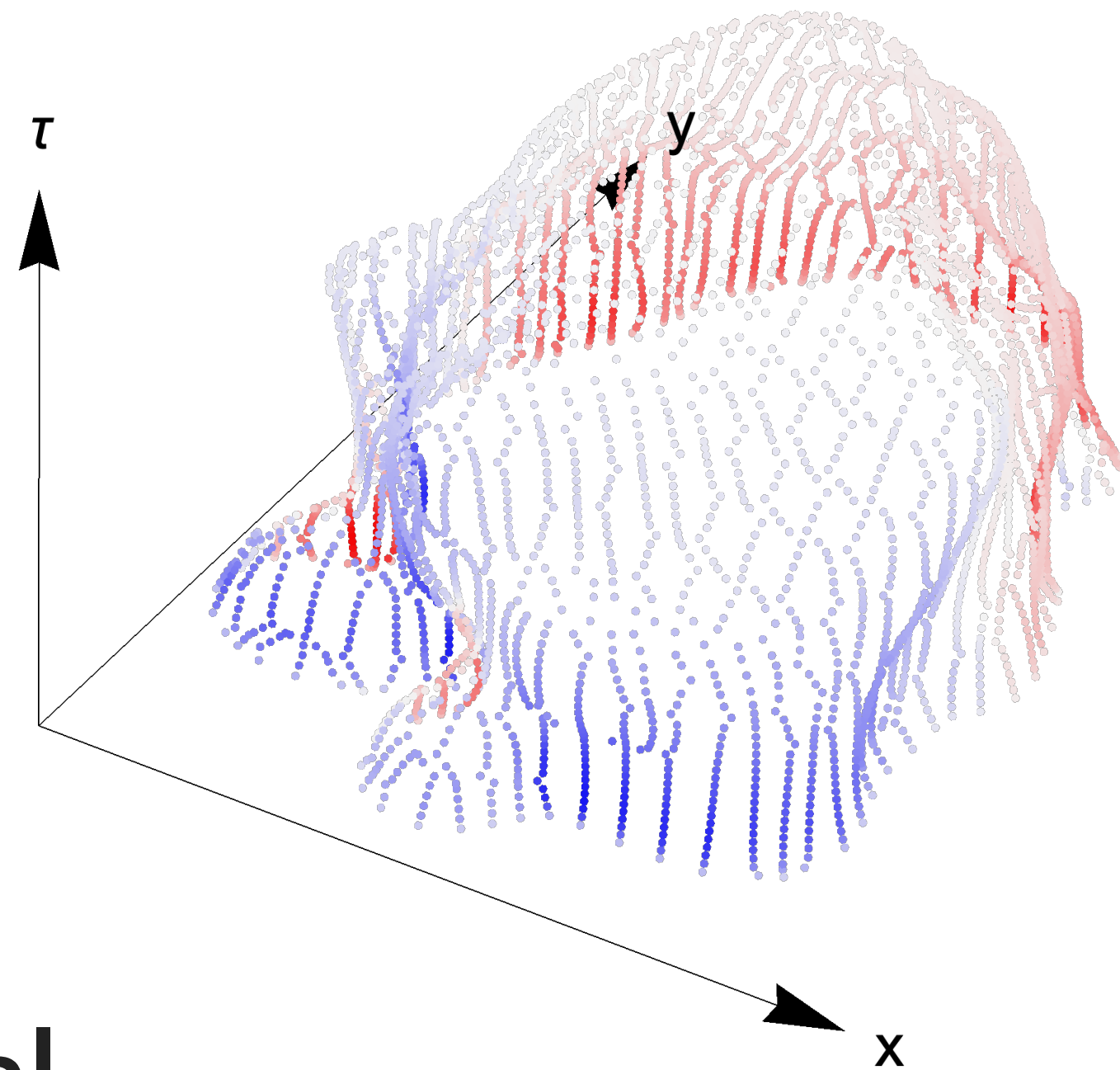
SS, Zhang, Hou, and Liao, Phys.Rev.Lett.(2020)
SS, Jiang, Lilleskov, and Liao, Annals Phys(2018)
Jiang, **SS**, Yin, and Liao, Chin.Phys.C(2018)

E-by-E IC of Bulk, n , $n5$, and B field



Initial condition
+ Driving force

Bulk => n , $n5$ evolution



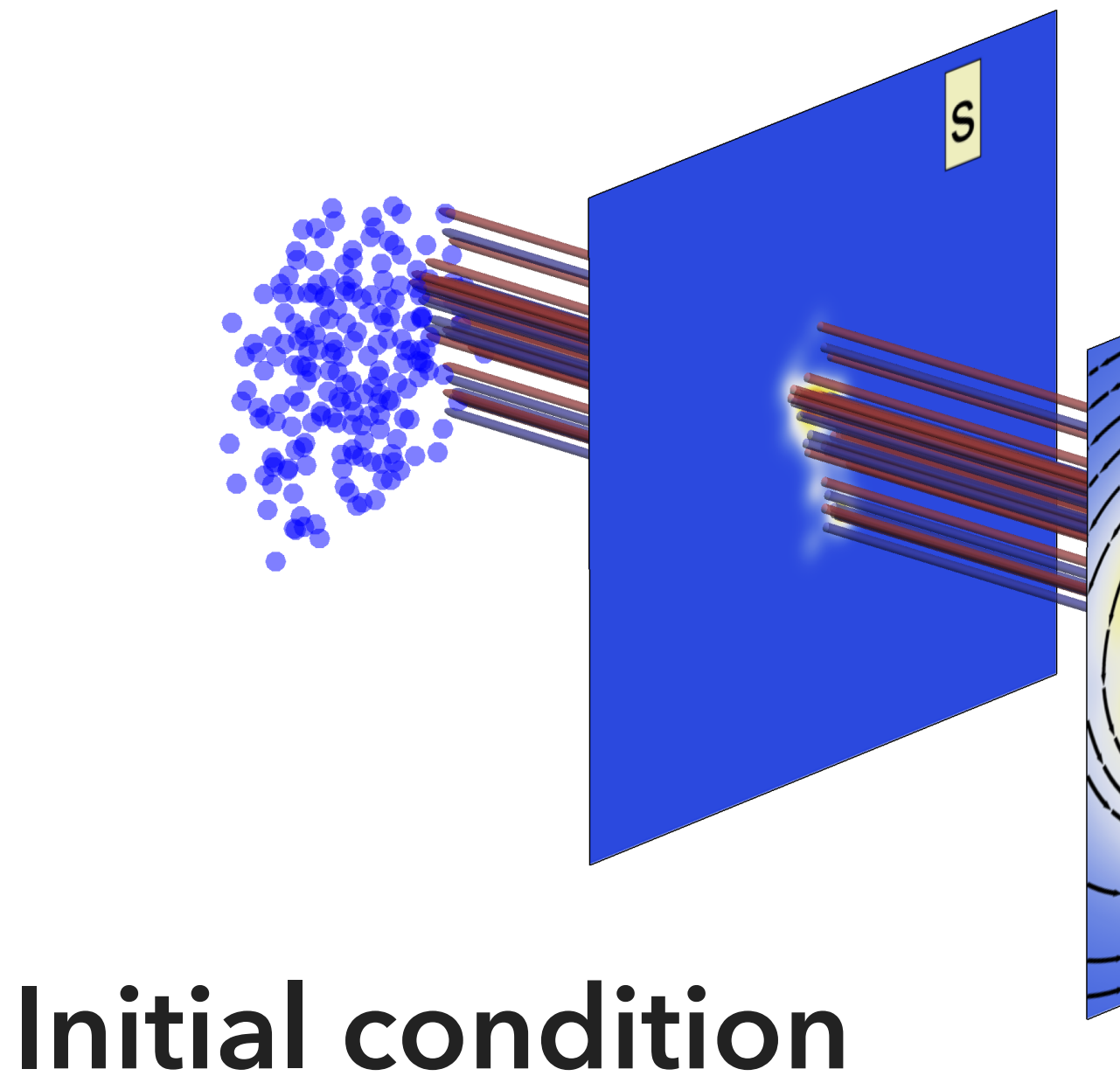
Dynamical
evolution

$\mu_Q @ FO$

Anomalous-**V**iscous
Fluid **D**ynamics

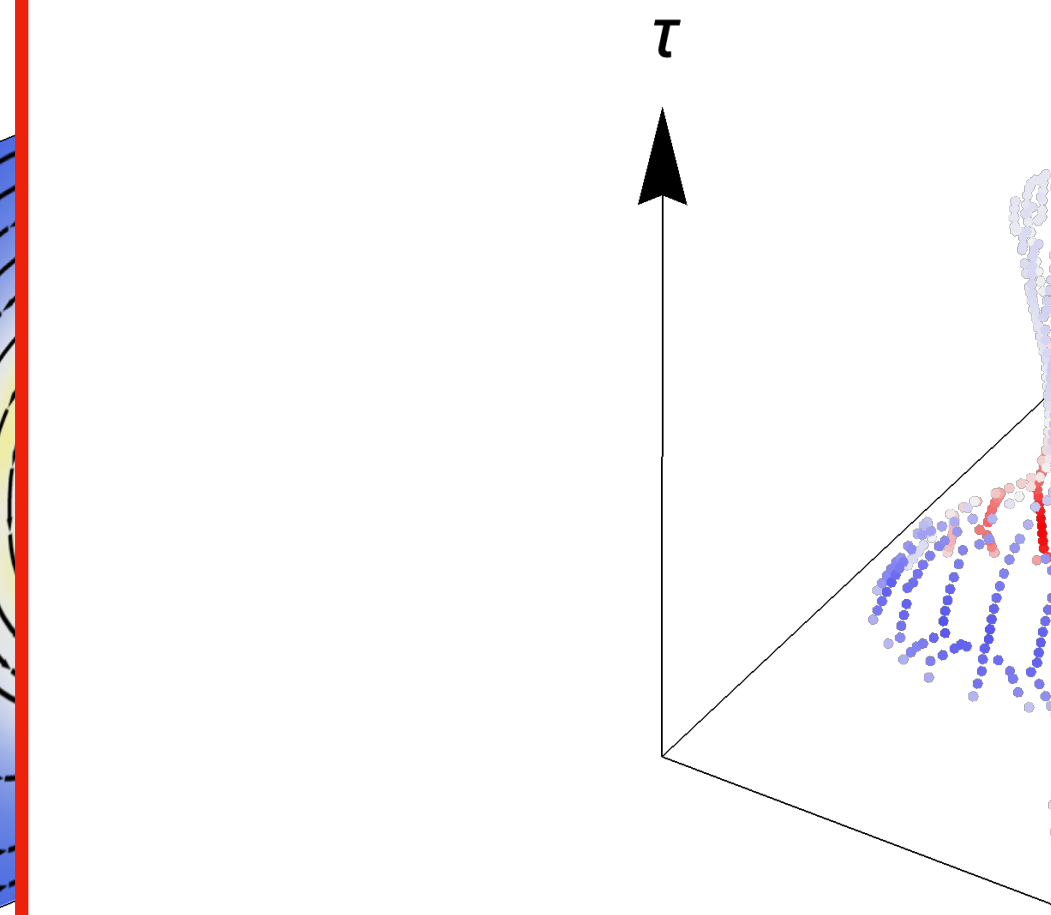
SS, Zhang, Hou, and Liao, Phys.Rev.Lett.(2020)
SS, Jiang, Lilleskov, and Liao, Annals Phys(2018)
Jiang, **SS**, Yin, and Liao, Chin.Phys.C(2018)

E-by-E IC of Bulk, n , $n5$, and B field



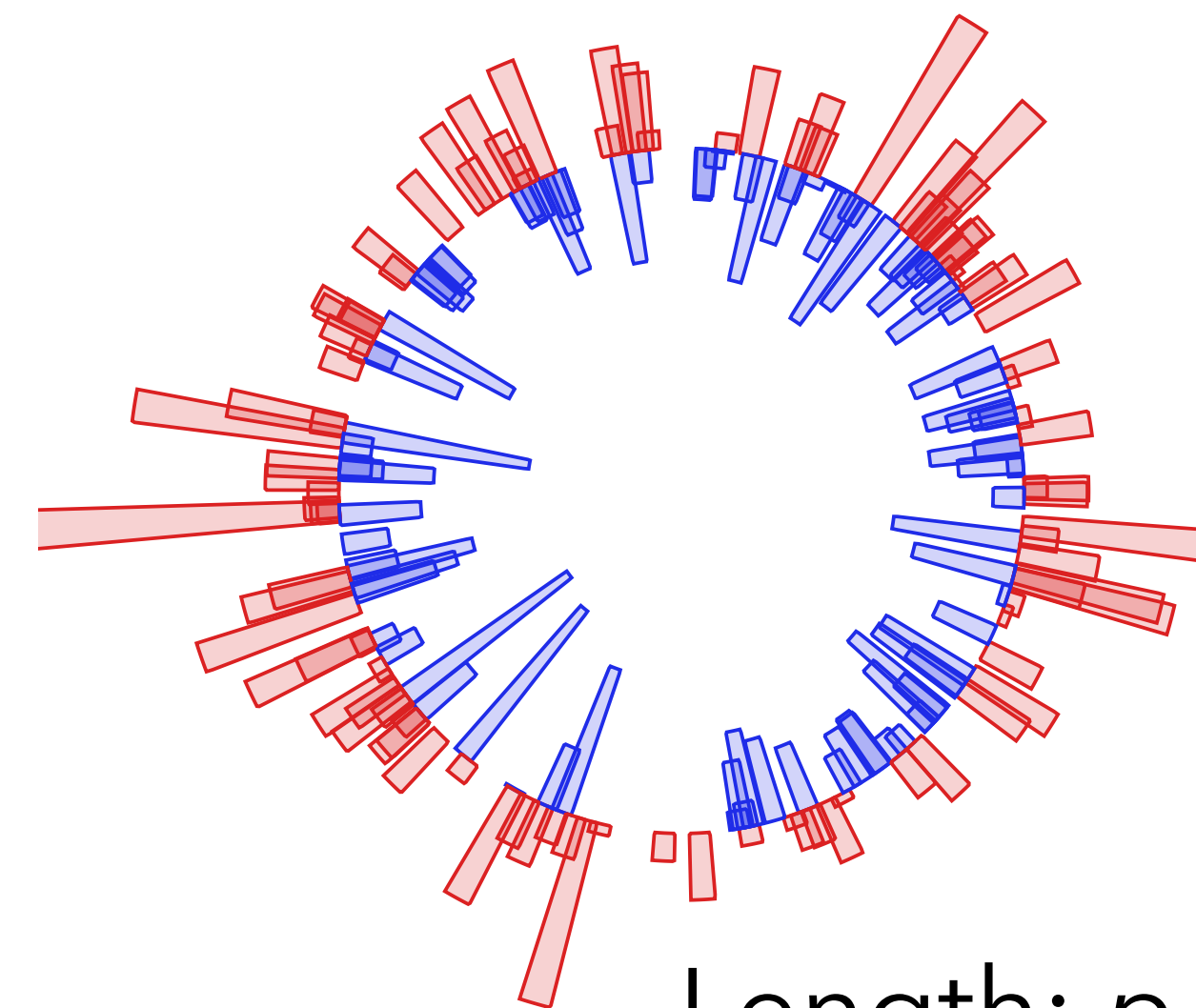
Initial condition
+ Driving force

Bulk => n , $n5$ evolution



Dynamical
evolution

Freeze-Out (LCC) + Hadron Cascade

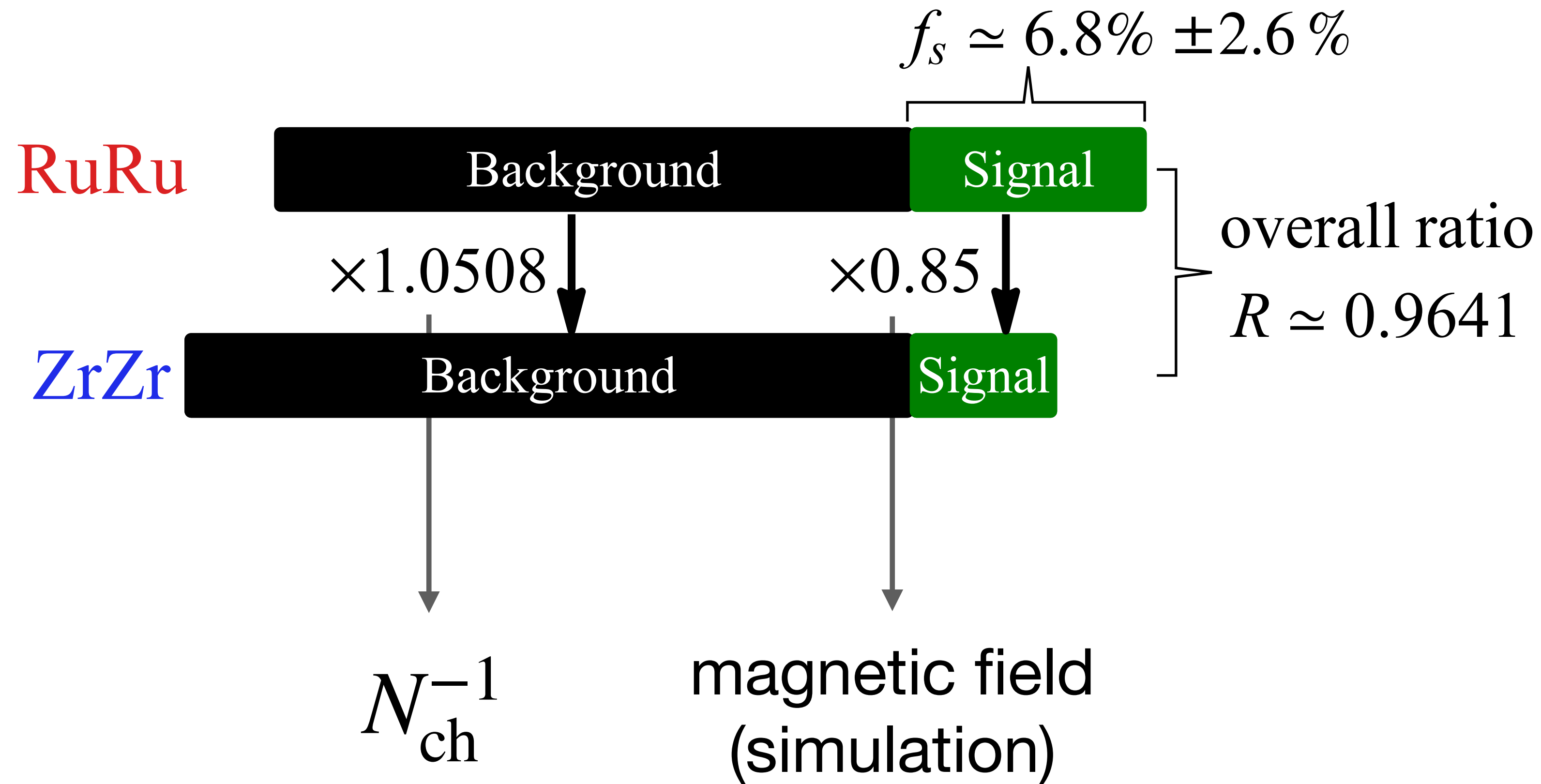


Hadron
distribution

Length: p_T , Angel: φ
Distribution of (+) / (-) ch.

Anomalous-Viscous
Fluid Dynamics

SS, Zhang, Hou, and Liao, Phys.Rev.Lett.(2020)
 SS, Jiang, Lilleskov, and Liao, Annals Phys(2018)
 Jiang, SS, Yin, and Liao, Chin.Phys.C(2018)



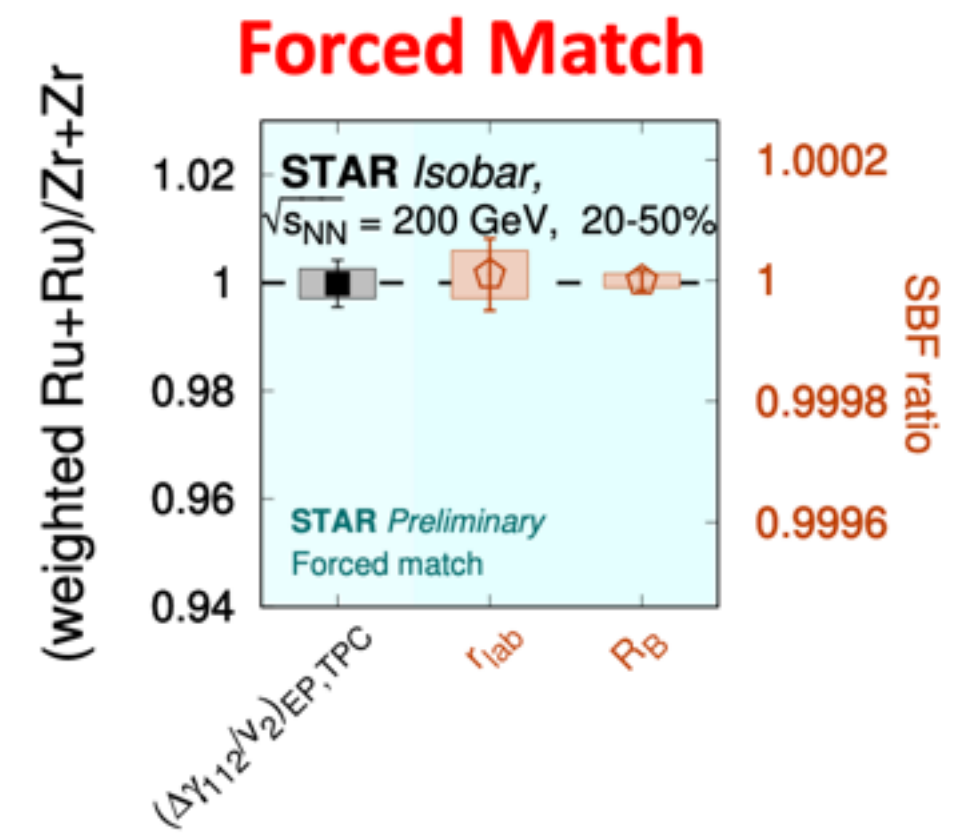
updates on experiment:

12.18 morning, II, 国兴厅

9:40-10:00	Search for the Chiral Magnetic Effect with Forced Match of Multiplicity and Elliptic Flow in Isobar Collisions at STAR	Yufu Lin
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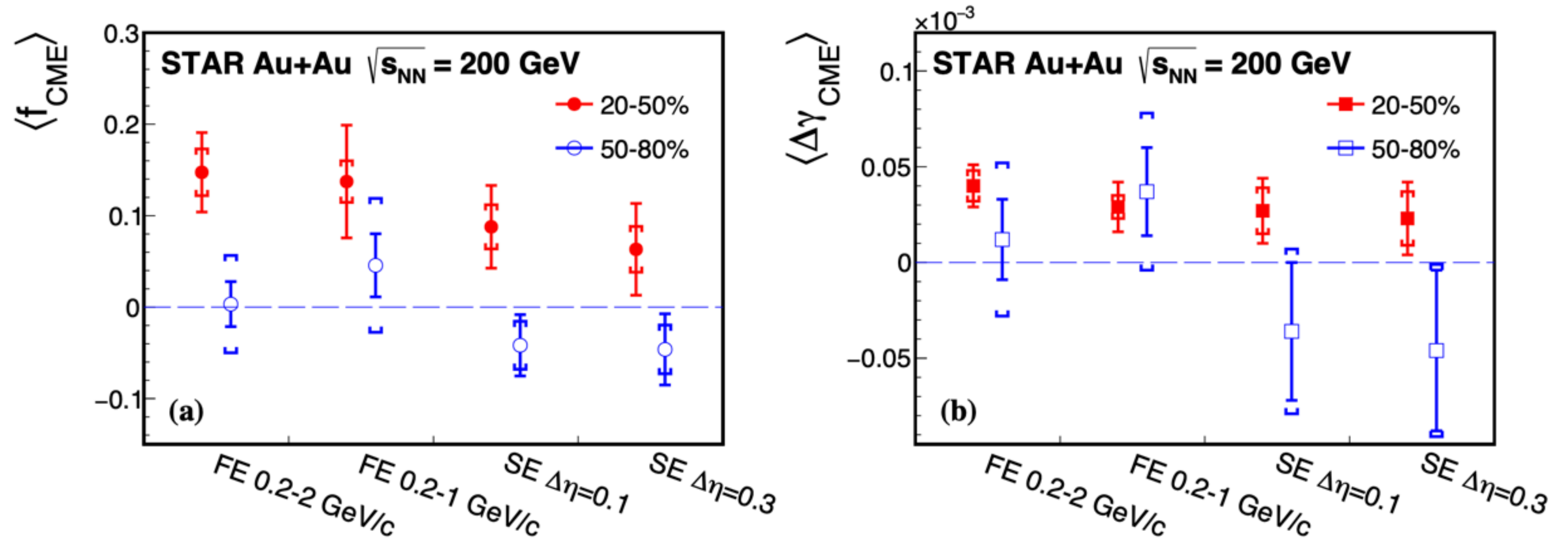
N_{ch}^{-1}

magnetic field
(simulation)

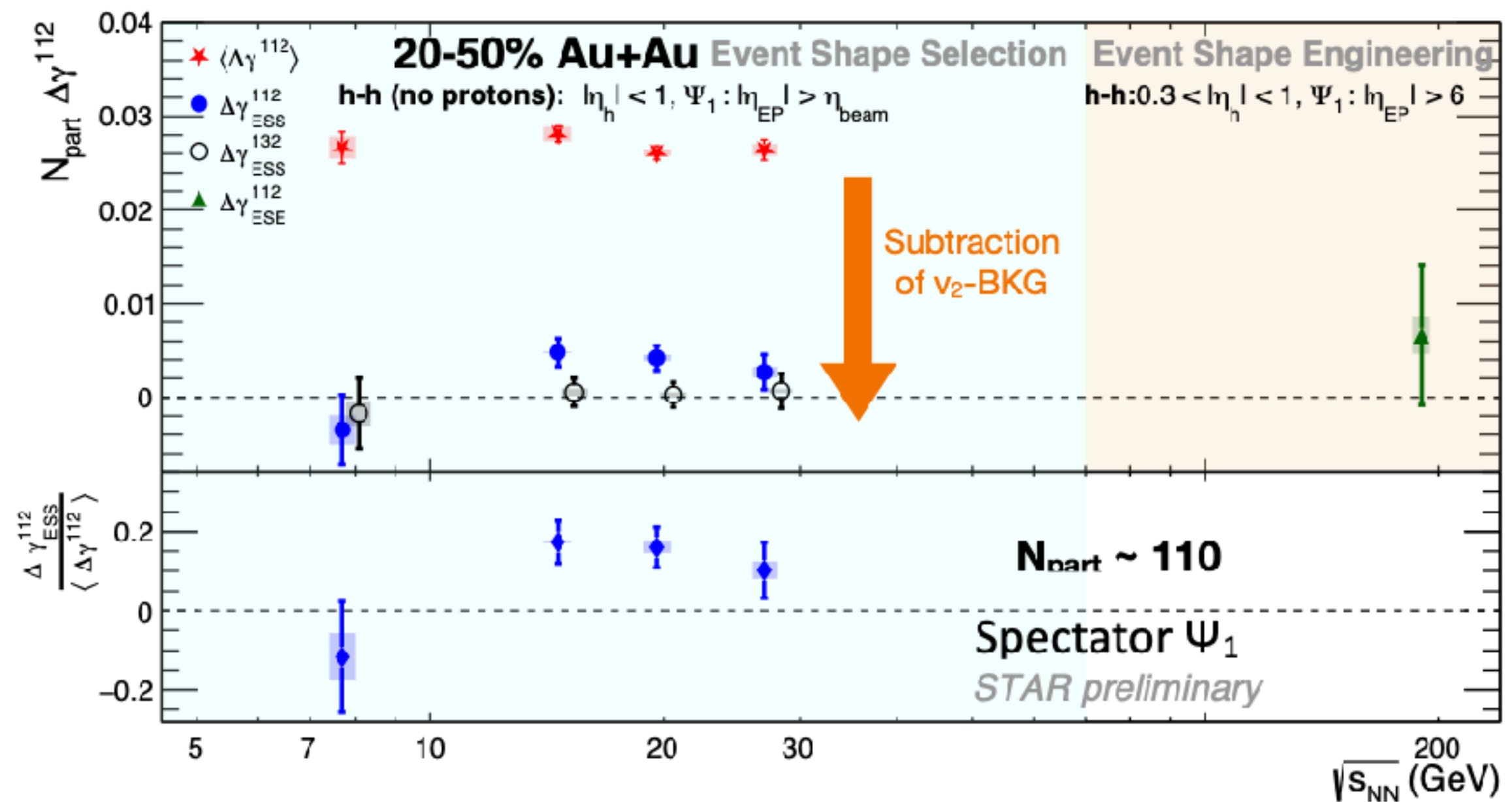


new method

STAR, PhysRevLett.128.092301



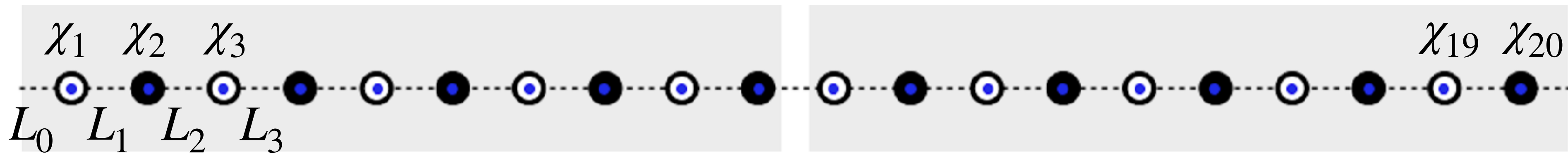
lower energy



1+1D massive Schwinger model

$$H = \int \left(\frac{E^2}{2} - \bar{\psi}(i\gamma^1 \partial_x - g\gamma^1 A - m)\psi \right) dx.$$

discretize and matrix(gate) representation:



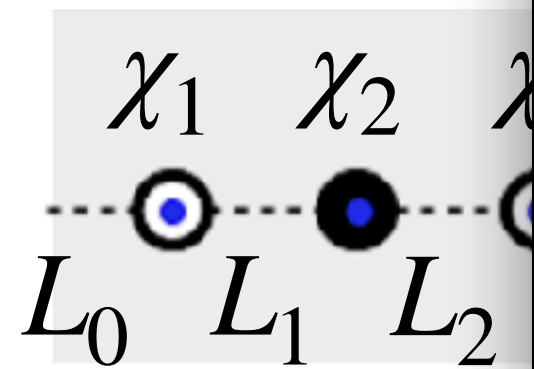
$$H = \frac{1}{4a} \sum_{n=1}^{N-1} (X_n X_{n+1} + Y_n Y_{n+1}) + \frac{m}{2} \sum_{n=1}^N (-1)^n Z_n + \frac{a g^2}{2} \sum_{n=1}^{N-1} L_n^2.$$

Pauli matrices: X, Y, Z

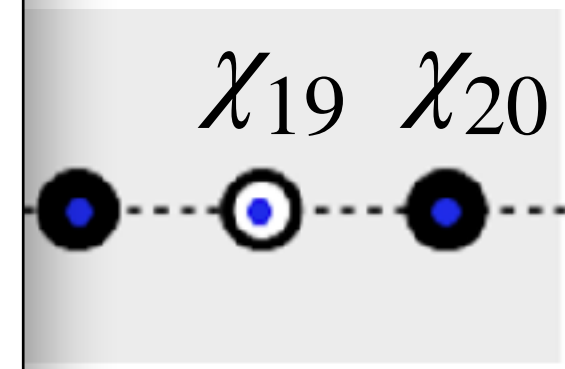
1+1D massive Schwinger model

$$H = \int \left(\frac{E^2}{2} - \bar{\psi}(i\gamma^1 \partial_x - g\gamma^1 A - m)\psi \right) dx.$$

discretize and map to time-dependent Schroedinger equation:

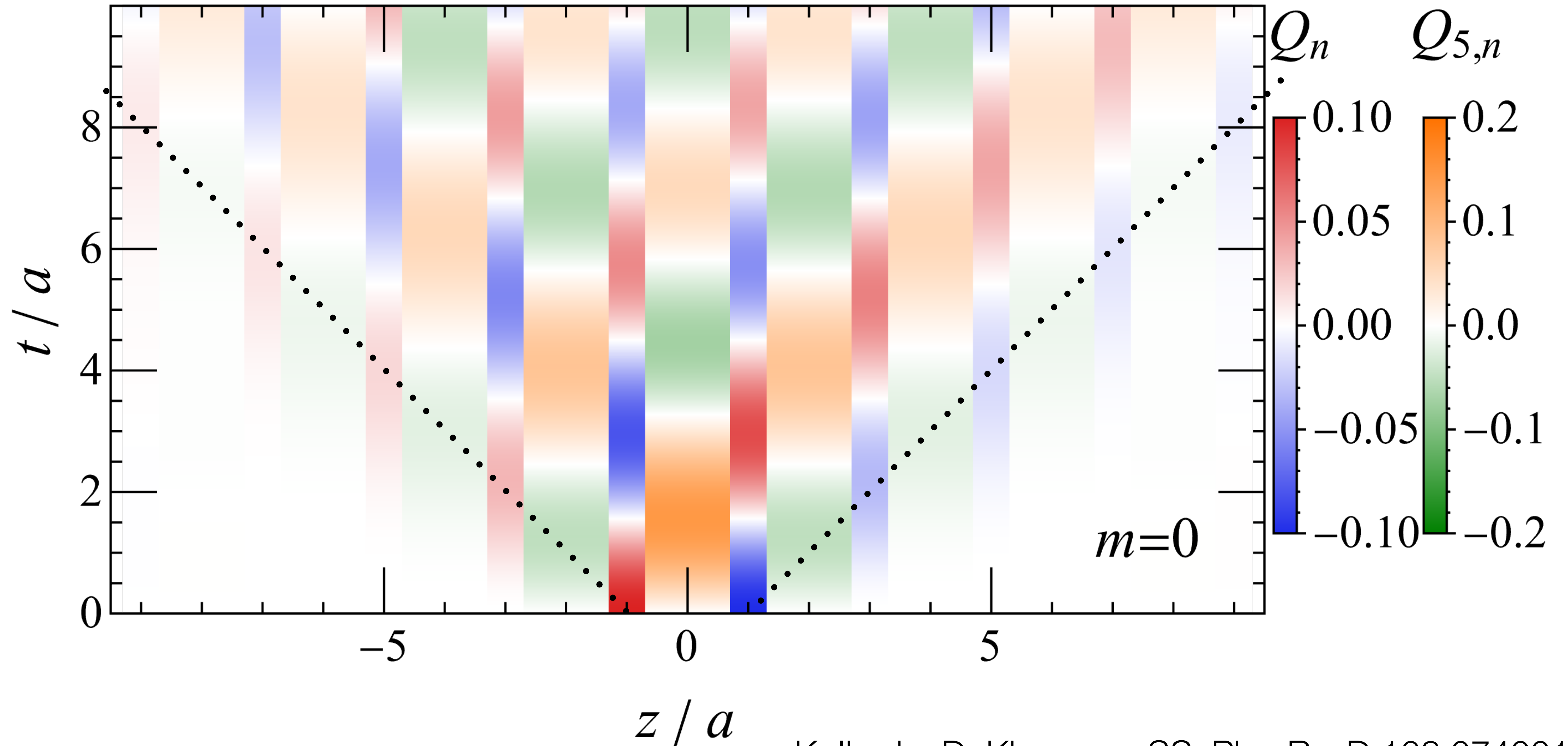


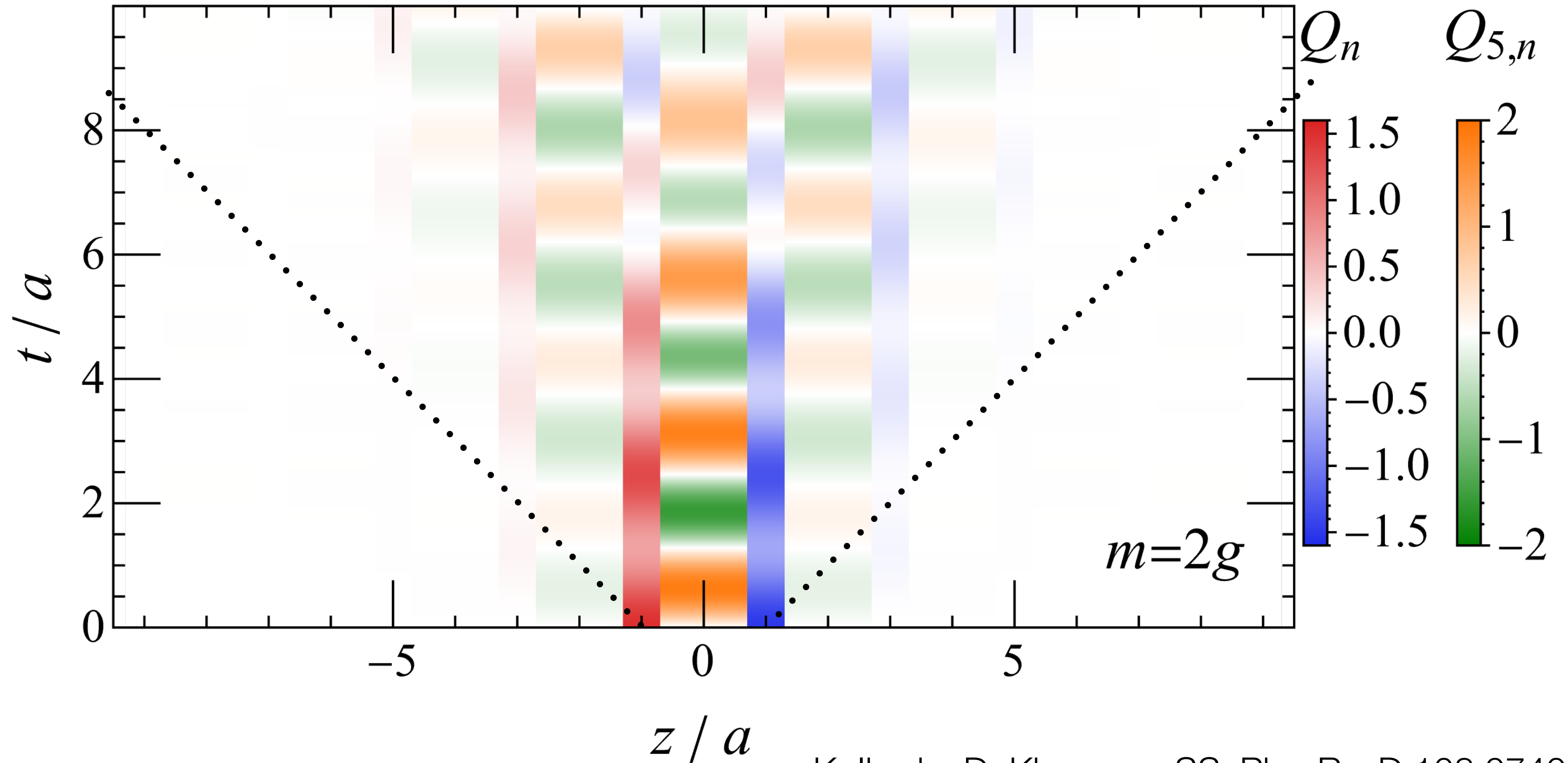
$$\frac{\partial}{\partial t} |\psi(t)\rangle = -i H |\psi(t)\rangle$$



$$H = \frac{1}{4a} \sum_{n=1}^{N-1} (X_n X_{n+1} + Y_n Y_{n+1}) + \frac{m}{2} \sum_{n=1}^N (-1)^n Z_n + \frac{a g^2}{2} \sum_{n=1}^{N-1} L_n^2.$$

Pauli matrices: X, Y, Z







More on quantum computation:

12.18 morning, II, 国兴厅

11:50-12:10

Quantum computing for parton fragmentation functions

Tianyin Li

12.18 afternoon

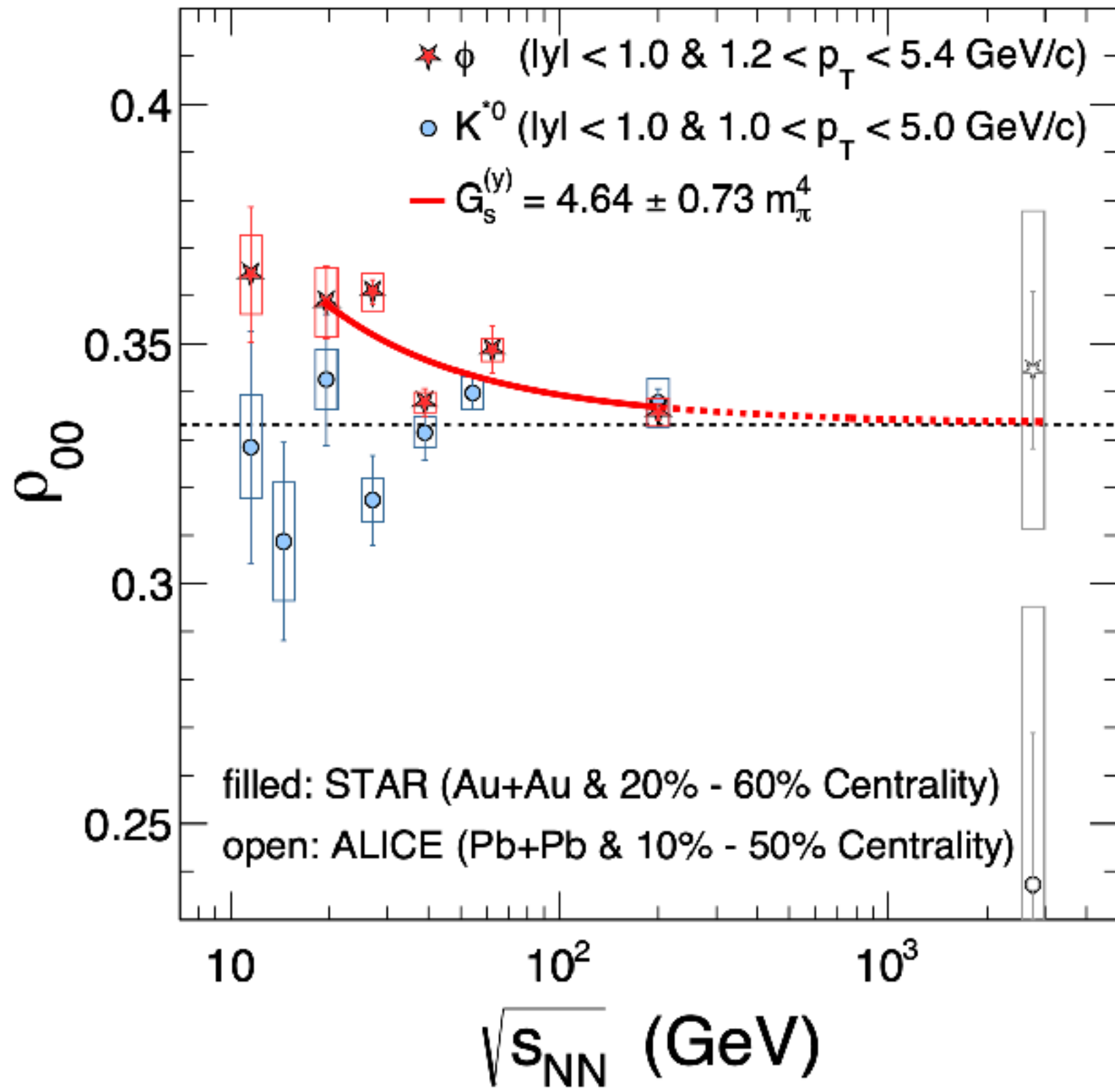
15:00-15:30

Quantum computing in QCD

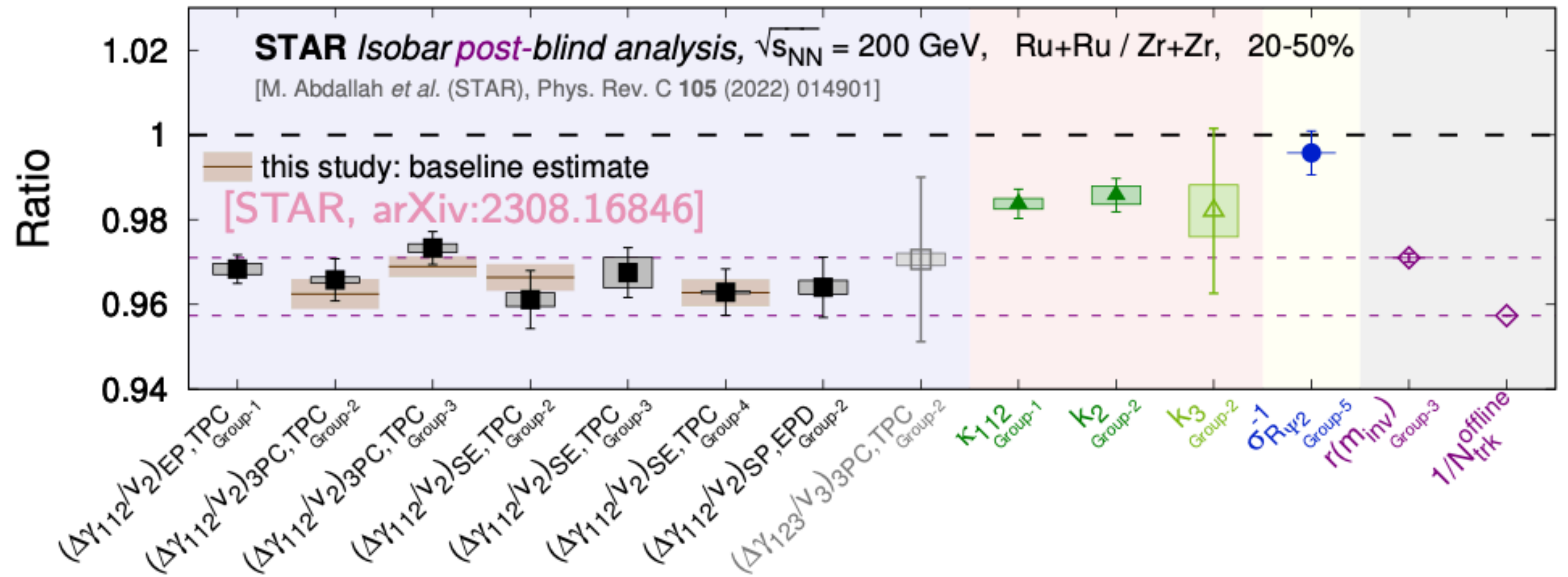
郭星雨

z / a

open questions

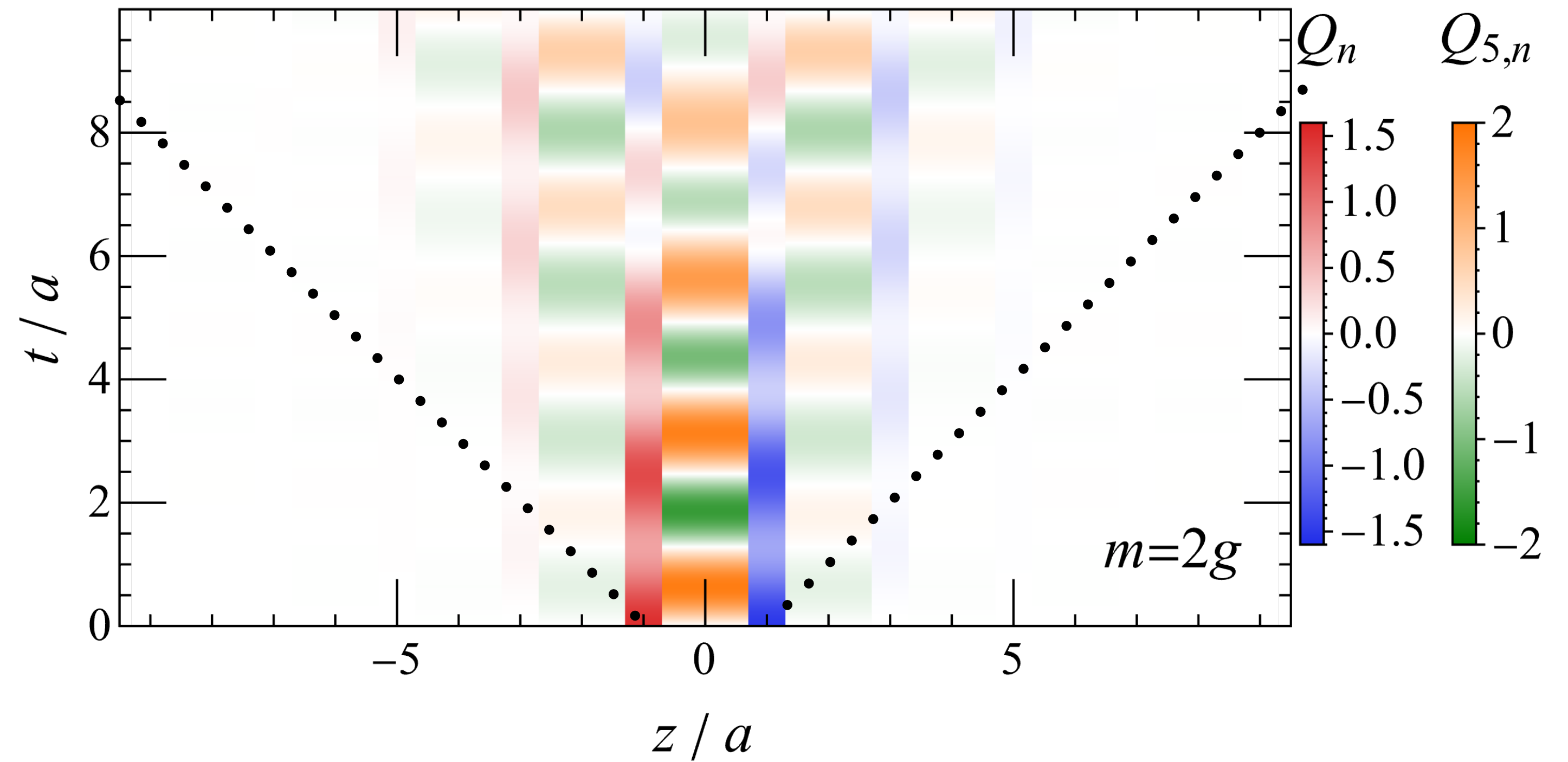


systematic understanding of
 vector meson spin alignment
 background of CME across beam energy



new opportunities using new computation techniques

quantum computation:



deep learning:

12.18 afternoon

14:30-15:00

Machine learning in QCD

王凌霄