Rapidity correlations and azimuthal correlations of strong color field

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References:

[1]/Yeyin Zhao, Mingmei Xu, Hengying Zhang, Yuanfang Wu, Nucl. Phys. A 955, 88 (2016).

[2] Hengying Zhang, Donghai Zhang, Yeyin Zhao, Mingmei Xu, Xue Pan and Yuanfang Wu, Phys. Rev. D, 034003 (2018).

[3] Donghai Zhang, Yeyin Zhao, Mingmei Xu, Xue Pan and Yuanfang Wu, Nucl. Phys. A 1011, 122201 (2021).

[4] Donghai Zhang, Yeyin Zhao, Mingmei Xu and Yuanfang Wu, Phys. Rev. D 107, 056017 (2023).

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OUTLINE

- Introduction
- Physical picture for long-range ridge correlations in CGC
- Rapidity correlations and dependences on p_t and $\sqrt{s_{NN}}$
- Fine structures of azimuthal correlations
- The features of two dimensional $\Delta y \Delta \varphi$ correlations

Ridge in small systems observed in experiments



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The physics underlying long range ridge correlations in small systems is inconclusive yet.

- / Final state effects, i.e. Hydrodynamics
- Initial correlations, CGC

CGC dynamics describe Initial state in high energy collisions



N. Armesto, L. McLerran, C. Pajares, NPA 781 (2007) 201.

Achievement of CGC on long range ridge correlations



, *Q*

p

e q

Glasma graph

- N_{ch} and p_t dependence of the long range azimuthal correlations in pPb collisions at 5.02 TeV at the LHC
- Describe CMS, ALICE and ATLASdata simultaneously with a common setof parameters

K. Dusling, R. Venugopalan, Phys. Rev. D 87 (2013) 051502(R). K. Dusling, R. Venugopalan, Phys. Rev. D 87 (2013) 054014. K. Dusling, R. Venugopalan, Phys. Rev. D 87 (2013) 094034.

(1) Only focus on azimuthal correlations at large |Δη|
(2) Lack a clear physical picture of correlation structure in η direction
We will focus on η direction in this study.





Observable:

$$C(\mathbf{p}_{\perp}, y_p, \mathbf{q}_{\perp}, y_q) = \frac{\frac{\mathrm{d}N_2}{\mathrm{d}^2 \mathbf{p}_{\perp} \mathrm{d}y_p \mathrm{d}^2 \mathbf{q}_{\perp} \mathrm{d}y_q}}{\frac{\mathrm{d}N}{\mathrm{d}^2 \mathbf{p}_{\perp} \mathrm{d}y_p} \frac{\mathrm{d}N}{\mathrm{d}^2 \mathbf{q}_{\perp} \mathrm{d}y_q}} - 1 = \frac{\frac{\mathrm{d}N_2^{\mathrm{corr.}}}{\mathrm{d}^2 \mathbf{p}_{\perp} \mathrm{d}y_p \mathrm{d}^2 \mathbf{q}_{\perp} \mathrm{d}y_q}}{\frac{\mathrm{d}N}{\mathrm{d}^2 \mathbf{p}_{\perp} \mathrm{d}y_p} \frac{\mathrm{d}N}{\mathrm{d}^2 \mathbf{q}_{\perp} \mathrm{d}y_q}}$$

where "p" and "q" are used to mark the two gluons.

(1) Physical Picture of Long Range Ridge Correlations





- Origin: strong correlations between source gluons and radiated gluons
- The same origin in pp and AA
- "W" shape

 $= 0 \begin{cases} \text{solid} & (r) \quad C_{rr}: \text{ platform, approx. boost invariance} \\ \text{dash} & (m) \quad C_{rm}: \text{ ridge-like corr. appears} \\ \text{dot-dash(s)} \quad C_{rs}: \text{ ridge-like corr. fully develop} \quad 10 \end{cases}$

(2) Rapidity correlations and dependences on p_t and $\sqrt{s_{NN}}$)



- Maximum corr. at $p_{\perp} = q_{\perp} \sim Q_{sA} + Q_{sB} = 2Q_{sp} \approx$ 1.5 GeV
- At lower energy, no small x in the central rapidity region, no ridge

(3) Fine structures of azimuthal correlations

To explore the contributions of different x degrees of freedom

$$C(\Delta \phi) = \int d\phi_p \int d\phi_q \, \delta(\phi_q - \phi_p - \Delta \phi) C(\mathbf{p}_{\perp}, y_p, \mathbf{q}_{\perp}, y_q)$$

• Finer binning in
$$p_{\perp}$$

• /Dependences on specific y, not Δy

 $pp@\sqrt{s} = 7 \text{ TeV}$



- $p_{\perp}=1.0$ GeV: a valley
- $p_{\perp}=1.5$ GeV: a single bump
- $p_{\perp}=1.8$ GeV: a single bump
- $p_{\perp}=2.0$ GeV: flat
- $p_{\perp}=2.5$ GeV: a valley

- $p_{\perp}=1.5$ GeV: a single bump at $\Delta \phi = \pi/2$
- $p_{\perp}=1.8$ GeV: double bumps at $\Delta \phi \approx 1$, 2
- $p_{\perp}=2.0$ GeV: double shoulders at $\Delta \phi \approx 1$, 2

Fine structures is related to harmonic components of $\cos n\Delta\phi$. ¹³

Fine structures in azimuthal corr. show up when:

- Finer binning in p_{\perp} (integration over p_{\perp} will smear it)
- Near $p_{\perp} \sim 2Q_{sp} = 1.8$ GeV, associated with the saturation momentum of colliding particles
- At least one gluon located at small rapidity, specific to small-x region

(4) Two dimensional Δy - $\Delta \varphi$ correlations



$$\frac{1}{N_{\text{Trig}}} \frac{d^2 N^{\text{pair}}}{d\Delta y d\Delta \phi} \qquad S(\Delta y, \Delta \phi) = \frac{1}{N_{\text{Trig}}} \frac{d^2 N^{\text{same}}}{d\Delta y d\Delta \phi},$$
$$B(\Delta y, \Delta \phi) = \frac{1}{N_{\text{Trig}}} \frac{d^2 N^{\text{mixed}}}{d\Delta y d\Delta \phi},$$
$$Y(\Delta y, \Delta \phi) = B(0, 0) \frac{S(\Delta y, \Delta \phi)}{B(\Delta y, \Delta \phi)}$$
Approximate normalization factor

Approximate normalization factor assuming a boost invariant rapidity distribution *K. Dusling, R. Venugopalan, Phys. Rev. D* 87 (2013) 094034.

$$\frac{B(\Delta y, \Delta \phi)}{B(0,0)} = 1 - \frac{|\Delta y|}{y^{\max} - y^{\min}}$$

Correct normalization scheme use real single-gluon distribution





The rebound is more easily observed in the differential correlation function ---- a more obvious "W" shape

PHENIX and STAR collaborations report different values of v2 / v3. Their detectors cover forward and central rapidity regions.

Gluon dynamics is a function of rapidity. Due this, there may be no discrepancy between PHENIX and STAR measurements.

SUMMARY

- Physical Picture: Different rapidity regions has different degrees of freedom of gluon
- Calculate ridge corr. at large $|\Delta y|$
- Structures in Δy direction consistent with CMS measurements
- These features can be used to directly test the CGC dynamics

Thank you!

- the leading logarithmic accuracy in x
- the leading graphs of $\frac{p_{\perp}}{Q_s}$
- $p_{\perp}, q_{\perp} \gtrsim Q_s$
- one-loop approximation of BK eq.
- AAMQS initial condition
- $Q_{s0p}^2 = 0.168 \text{ GeV}^2$ pp collision at 7 TeV, $Q_{s0A}^2 = 0.504 \text{ GeV}^2$
- the leading order Balitsky-Kovchegov (BK) equation with a running coupling kernel, Balitsky's prescription

dipole forward scattering amplitude $\Phi_{A(B)}(x, \mathbf{k}_{\perp}) = \frac{N_c \mathbf{k}_{\perp}^2}{4\alpha_s} \int d^2 \mathbf{r}_{\perp} e^{i\mathbf{k}_{\perp} \cdot \mathbf{r}_{\perp}} [1 - \mathcal{N}_{ad.}(\mathbf{r}_{\perp}, Y)]$ unintegrated gluon dis. (uGD) • a phenomenological extrapolation of uGD at large x $\Phi(x, \mathbf{k}_{\perp}) = \left(\frac{1-x}{1-x_0}\right)^{\beta} \Phi(x_0, \mathbf{k}_{\perp}) \text{ for } x > x_0$