



中國地質大學

# Investigating collective flow in high energy pp collisions with a multi-phase transport model

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Zhuhai Guangdong  
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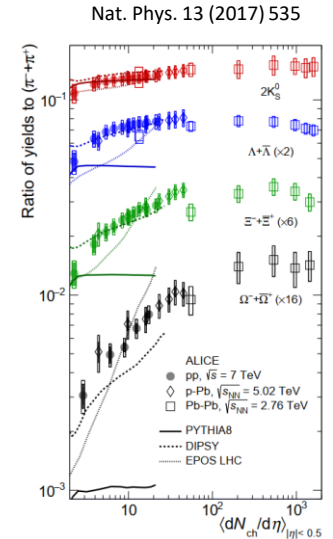
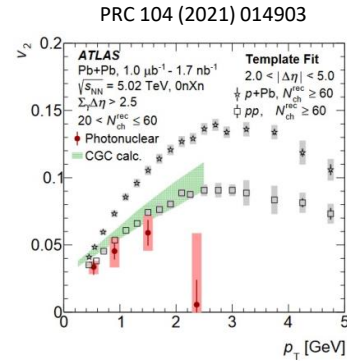
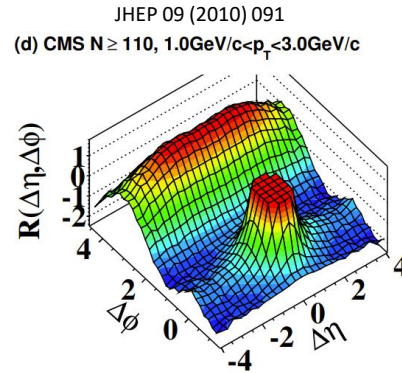
# Outline

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- Collectivity in small system
- AMPT with PYTHIA8 initial condition and sub-nucleon structures
- Near side long range correlations
- Collective flow in pp collisions

# Collectivity in small system

- Long range correlation in high multiplicity pp/pA: “ridge” structure
- Non-zero azimuthal collective flow even in UPC collisions
- Multiplicity dependent strangeness enhancement
- ...



- Hydrodynamic final state evolution

EPJC 76 (2016) 7, PLB 774 (2017) 351, PLB 780 (2018) 495, EPJC 80 (2020) 9, PRC 101 (2020) 2

- CGC Initial state correlations

PRL 108 (2012) 262001, PRD 87 (2013) 051502, PRL 117 (2016) 162301

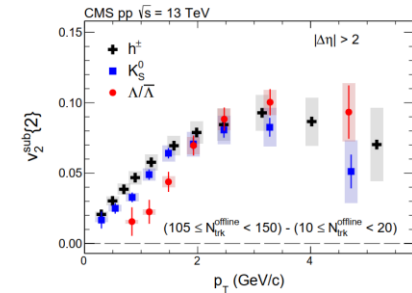
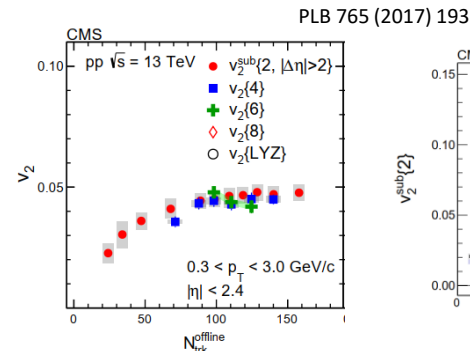
- Collective string effects

□ color reconnection: PRL 111 (2013) 042001, PRD 92 (2015) 094010

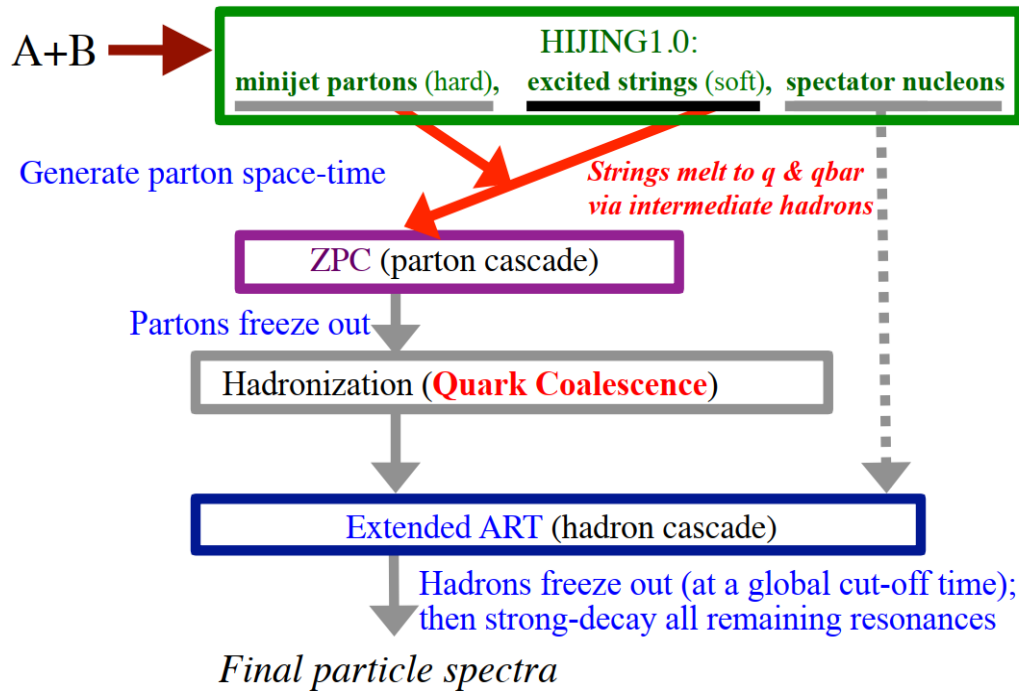
□ color rope: JHEP 03 (2015) 148, PLB 779 (2018) 58, JHEP 03 (2021) 270

- Parton escape effects

PLB 753 (2016) 506

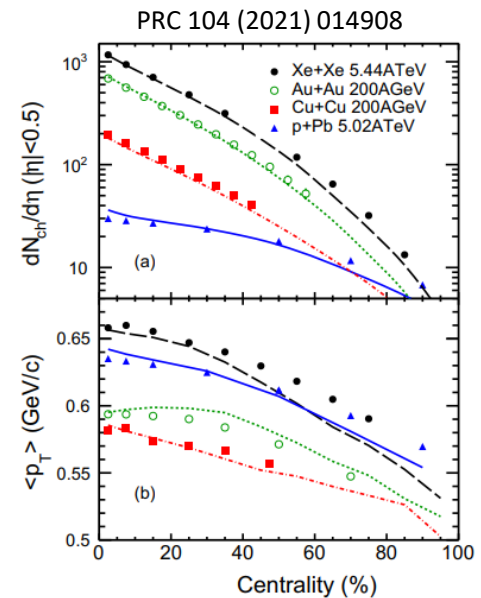


# A multi-phase transport model (AMPT)

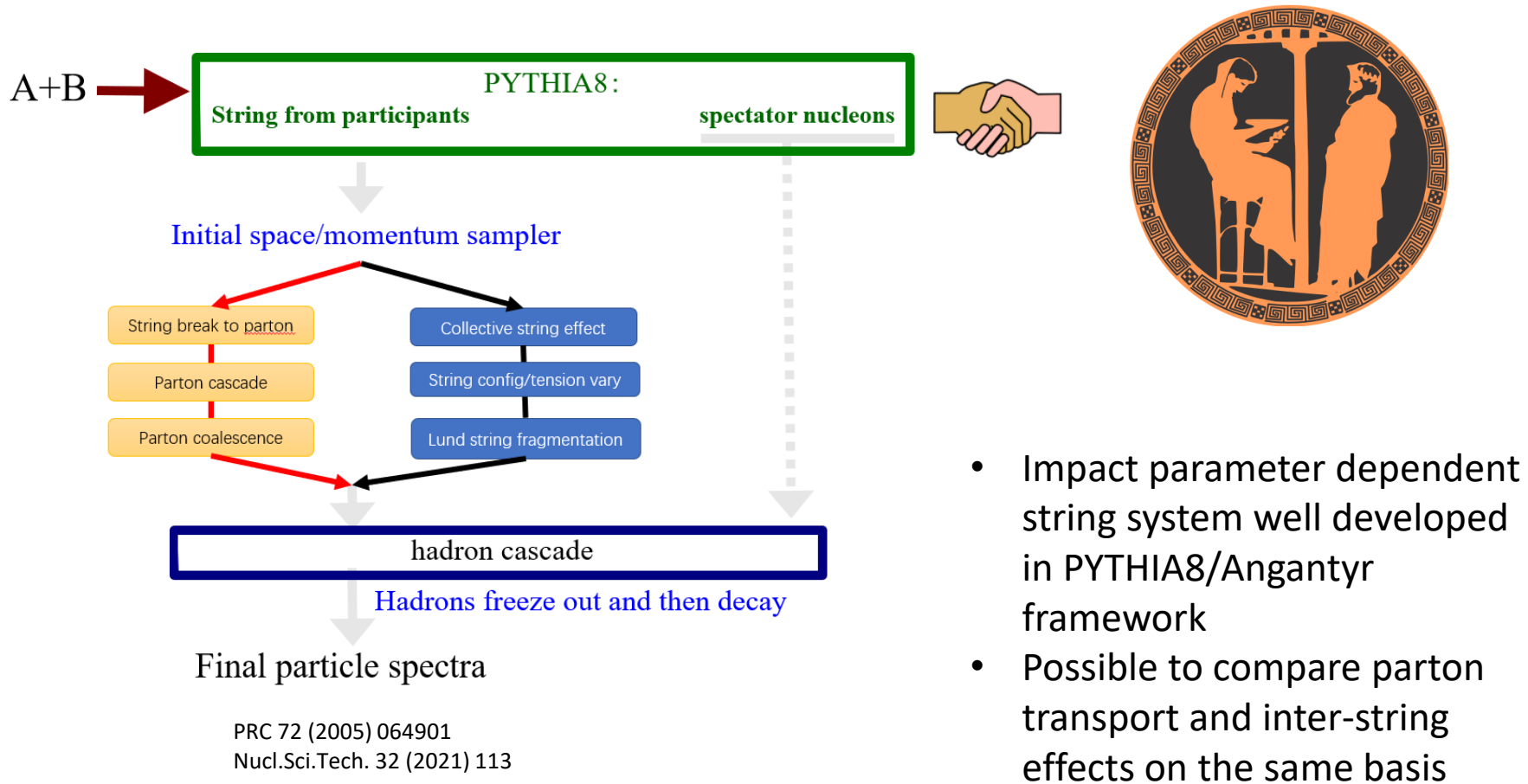


PRC 72 (2005) 064901  
 Nucl.Sci.Tech. 32 (2021) 113

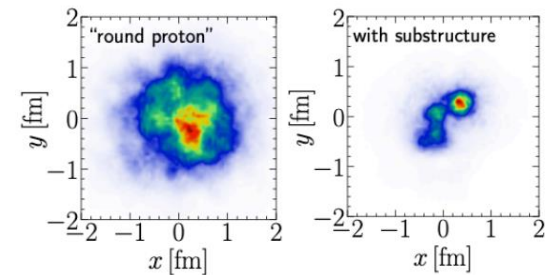
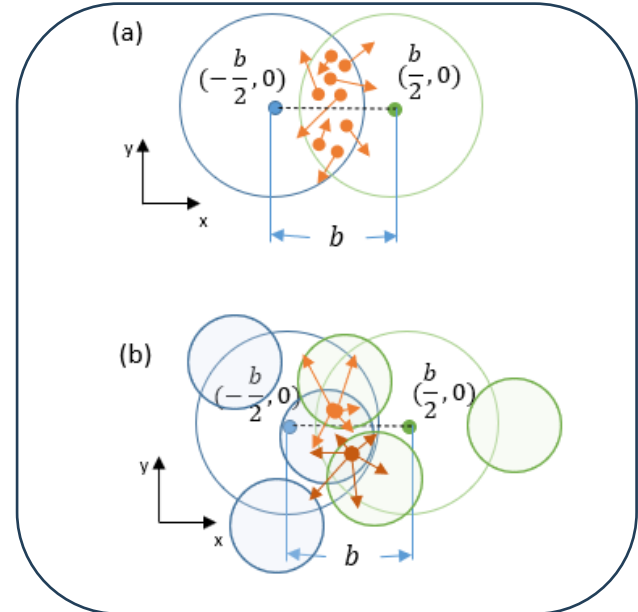
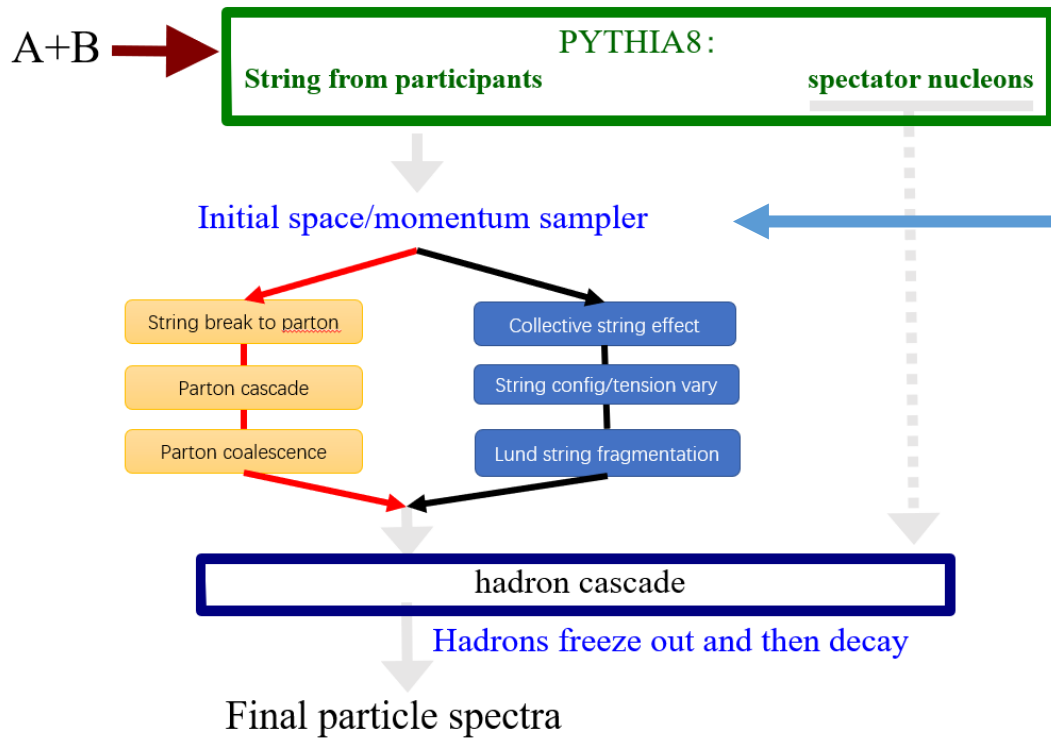
- AMPT has been extensively applied in collectivity study for large systems
- System size dependence has been improved recently



# AMPT x PYTHIA8



# Sub-nucleon structure

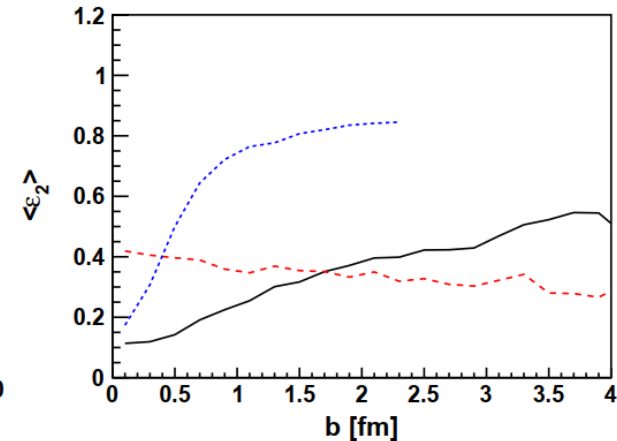
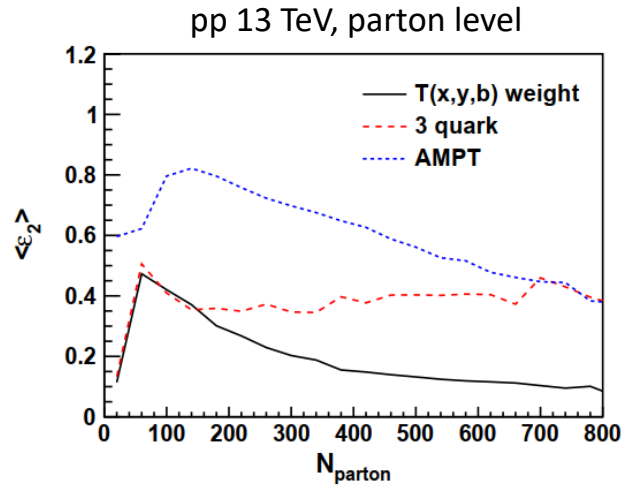
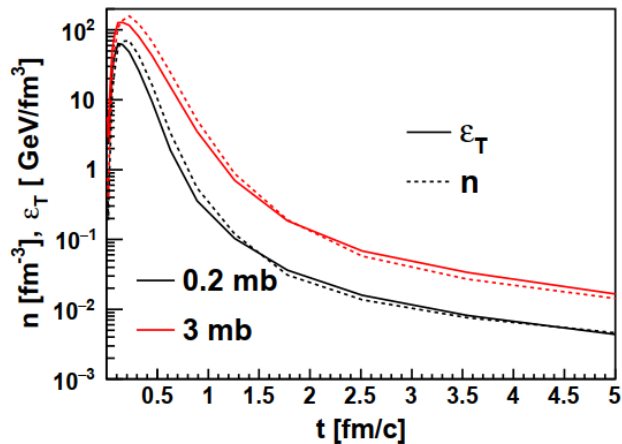


Rept.Prog.Phys. 84 (2021) 082301

PRC 72 (2005) 064901  
Nucl.Sci.Tech. 32 (2021) 113

# Space-time evolution

- Different features of spatial eccentricity distribution
- Rather flat eccentricity with sub-nucleon fluctuation (slightly increasing towards high multiplicity/central events)
- Energy density above  $\sim 1\text{GeV}/\text{fm}^3$ , within  $t < 1\text{ fm}/c$ .

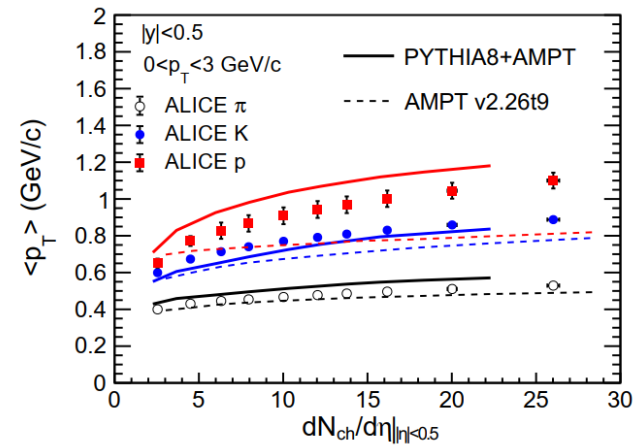
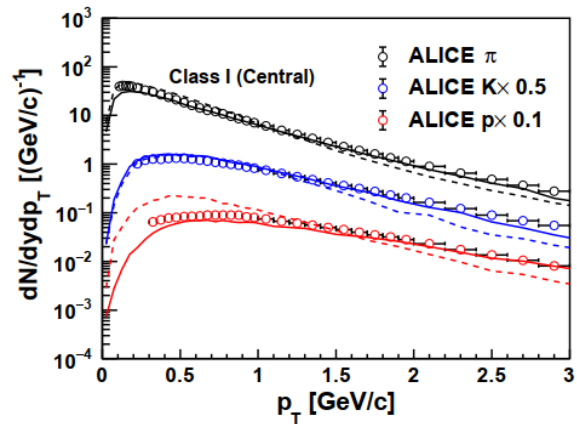
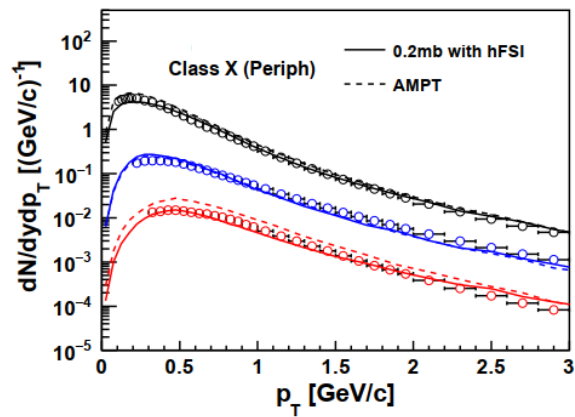


EPJC 81 (2021) 8, 755

# Compared to original AMPT

- Original AMPT fails to describe the multiplicity dependence
- Improvement may come from the MPI based PYTHIA8 initial condition

pp 13 TeV, hadron level



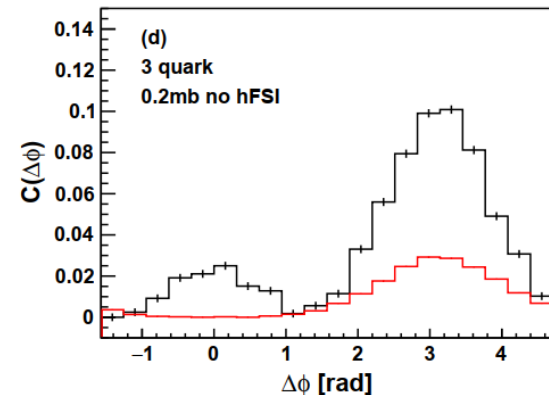
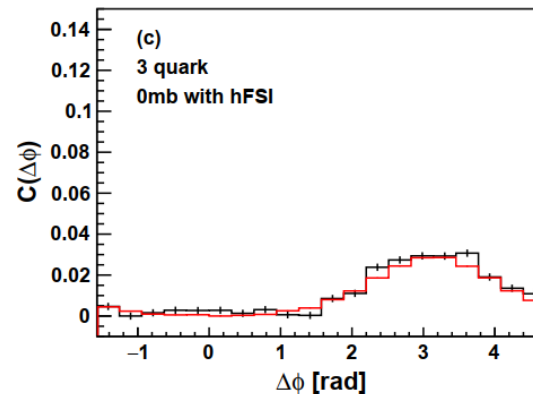
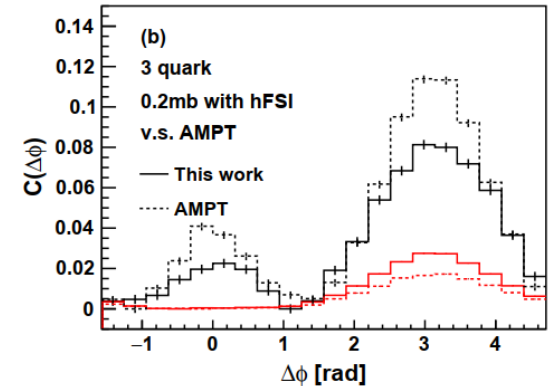
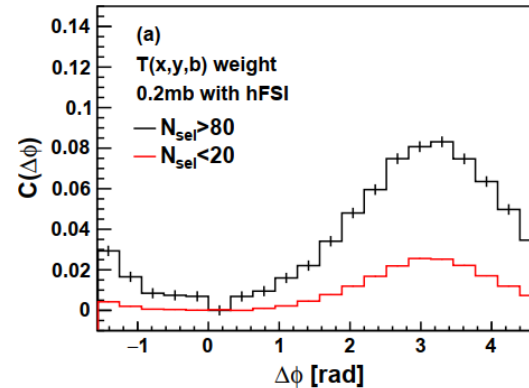
EPJC 81 (2021) 8, 755



# Two particle correlation

pp 13 TeV, charged hadron  
 $1 < p_T < 3$  GeV/c,  $|\Delta\eta| > 2$

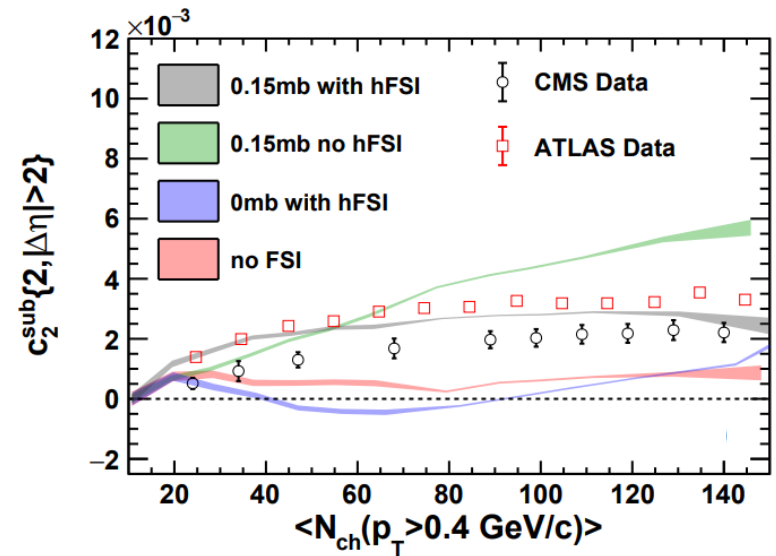
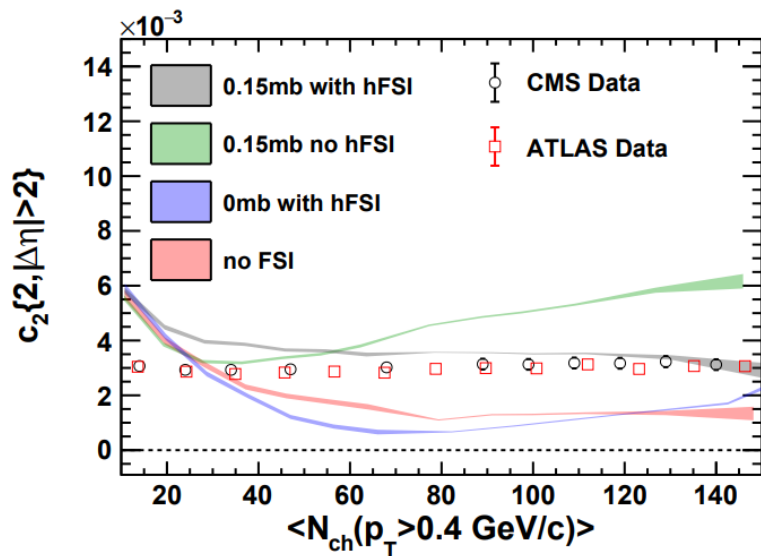
- Near side ridge appears with sub-nucleon fluctuation
- Significant ridge also shows in AMPT due to its large initial eccentricity
- Robust HM ridge developed in parton evolution stage



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# Two-particle cumulants

- Low multiplicity dominate by non-flow effects
- Final state parton rescattering leads to sizable  $c_2\{2\}$
- Non-flow effectively subtracted with low multiplicity events
- With all FSI, data behavior well captured



$$Q_n = \sum_{i=1}^M e^{in\phi_i} \quad c_n\{2, |\Delta\eta|\} = \langle\langle 2 \rangle\rangle_{\Delta\eta}$$

$$\langle 2 \rangle_{\Delta\eta} = \frac{Q_n^A \cdot Q_n^{B*}}{M_A \cdot M_B}$$

$$c_n^{sub}\{2, |\Delta\eta|\} = c_n\{2, |\Delta\eta|\} - k \cdot c_n^{low}\{2, |\Delta\eta|\}$$

$$k = \langle M \rangle^{low} / \langle M \rangle$$

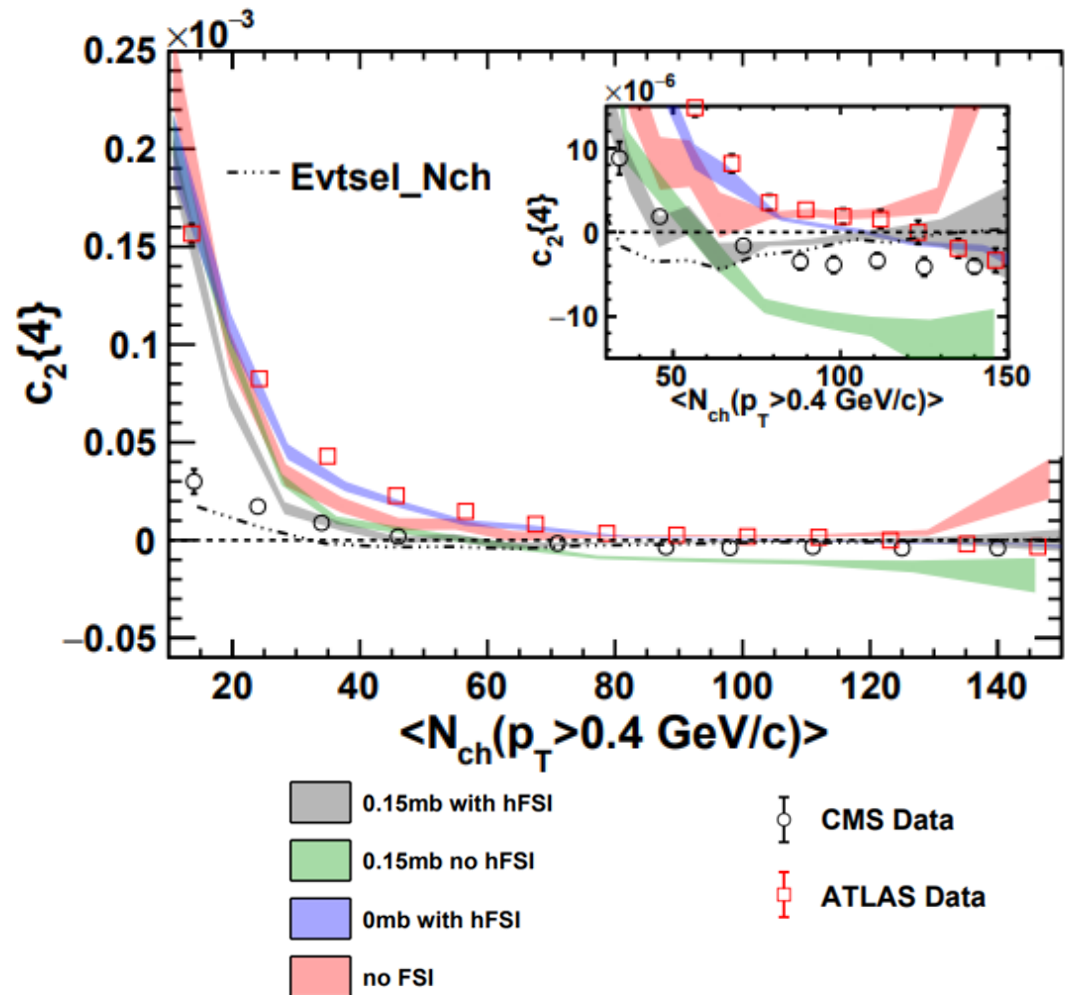
# Four-particle cumulants

- Parton rescattering leads to the most significant negative  $c_2\{4\}$
- Sign change dependent on multiplicity observed in the model calculations
- Data discrepancy due to different event averaging treatment well described in the model

$$\langle 2 \rangle = \frac{|Q_n|^2 - M}{M(M-1)},$$

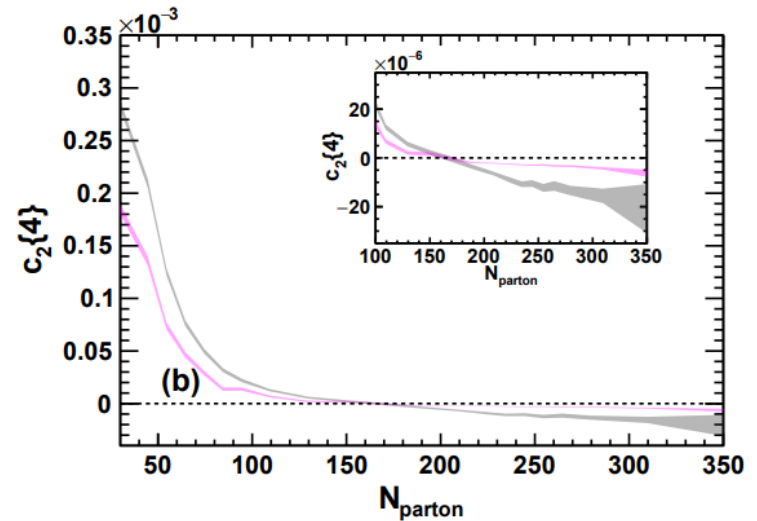
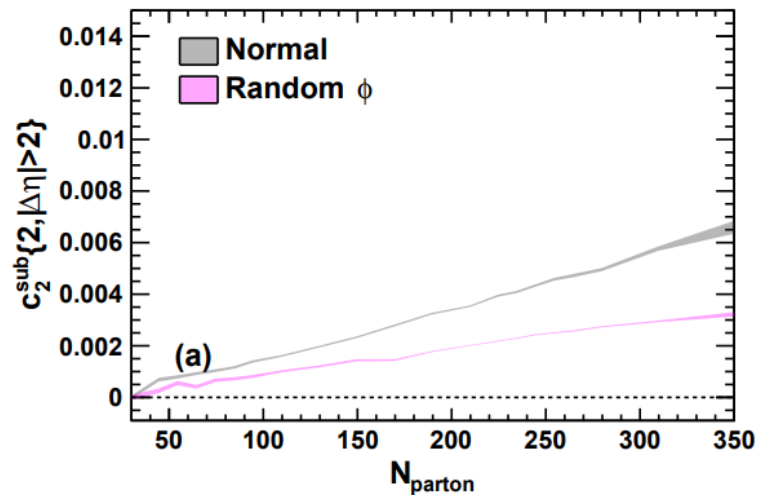
$$\langle 4 \rangle = \frac{|Q_n|^4 + |Q_{2n}|^2 - 2 \cdot \text{Re}[Q_{2n}Q_n^*Q_n^*]}{M(M-1)(M-2)(M-3)} - 2 \frac{2(M-2) \cdot |Q_n|^2 - M(M-3)}{M(M-1)(M-2)},$$

$$c_n\{4\} = \langle\langle 4 \rangle\rangle - 2 \cdot \langle\langle 2 \rangle\rangle^2,$$



# Parton escape effects

- Sizeable contribution from parton escape, decreases with event multiplicity
- Sign change of  $c_2\{4\}$  can be related to the parton escape mechanism



# $p_T$ differential elliptic flow for charge particle

- Sizable elliptic flow induced by parton level evolutions
- Hadronic interactions generally reduce the flow magnitude
- Negligible flow with only hadronic scattering or no FSI

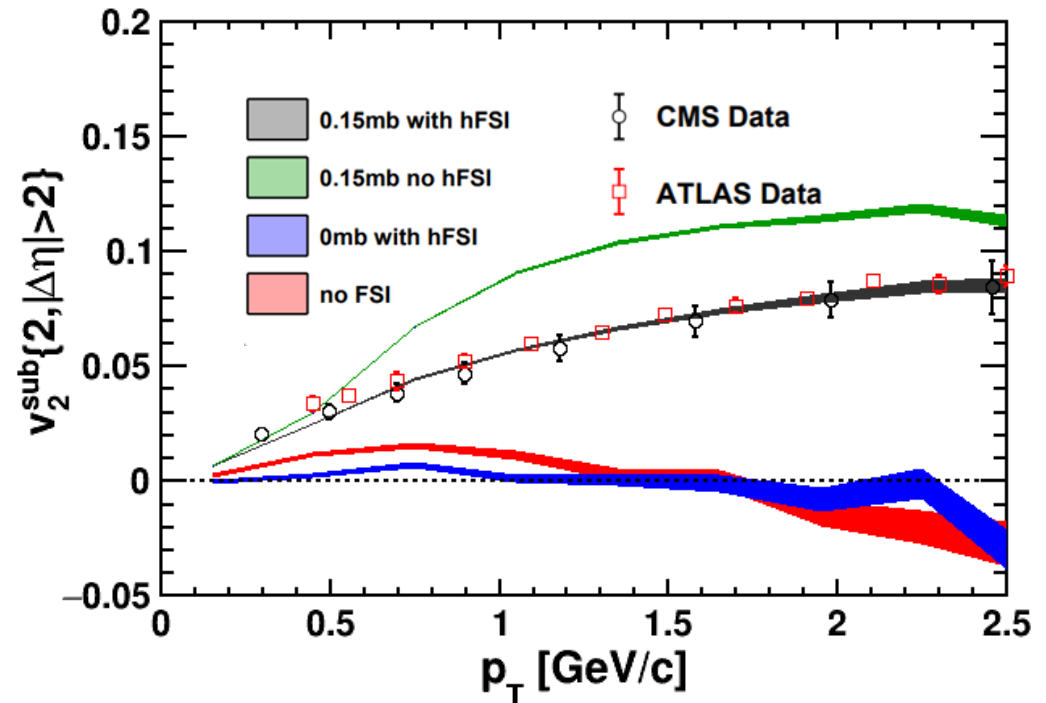
$$p_n = \sum_{i=1}^{m_p} e^{in\phi_i}$$

$$\langle 2' \rangle_{\Delta\eta} = \frac{p_{n,A} Q_{n,B}^*}{m_{p,A} M_B}$$

$$d_n\{2, |\Delta\eta|\} = \langle \langle 2' \rangle \rangle_{\Delta\eta}$$

$$v_n^{sub}\{2, |\Delta\eta|\} = \frac{d_n\{2, |\Delta\eta|\} - k \cdot d_n^{low}\{2, |\Delta\eta|\}}{\sqrt{c_n\{2, |\Delta\eta|\} - k \cdot c_n^{low}\{2, |\Delta\eta|\}}}$$

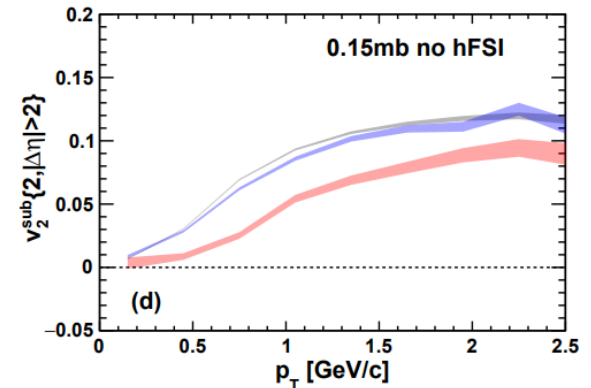
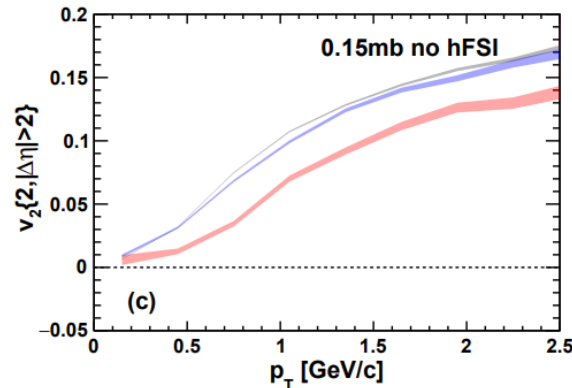
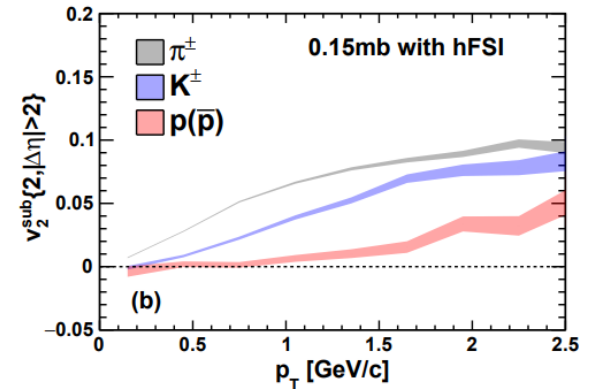
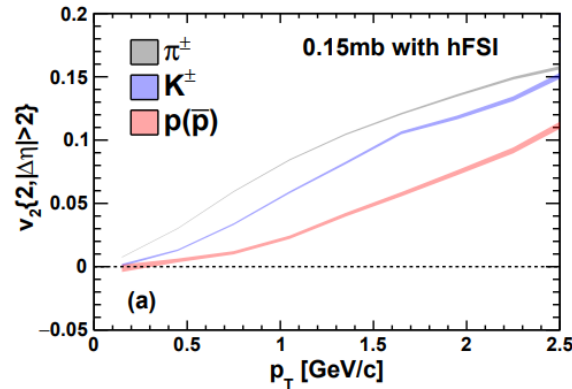
Elliptic flow in high multiplicity events ( $M > 100$ )



# Elliptic flow for identified particles

- Significant mass ordering at low  $p_T$  w/ or w/o low multiplicity subtraction
- Hadronic FSI coupled to parton evolutions further enlarge the mass ordering feature

Elliptic flow in high multiplicity events ( $M > 100$ )



# Summary

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- Collectivity in small systems has been stimulating for new theoretical and experimental developments in QGP studies.
- Further improvements have been implemented in the AMPT framework to reach better description for small system.
- Sub-nucleon fluctuations can be important to understand the collective effects developed in small system collisions.
- Collective flow in pp collisions may be related to the existence of final state parton evolutions.
- Strong mass ordering of identified hadron elliptic flow can be regarded as important signature to expose the parton collectivity in pp collisions.

Thank You!  
😊

Back up



# Elliptic flow in pp collisions

- Sub-nucleon structure provides more realistic multiplicity dependence of elliptic flow in pp collisions
- Negative 4-particle correlation appears with the transport model framework. What is the origin?
- Non-trivial final state cross section dependence revealed in  $c_2\{4\}$

pp 13 TeV, charged hadron,  $0.3 < p_T < 3$  GeV/c,  $|\Delta\eta| > 2$

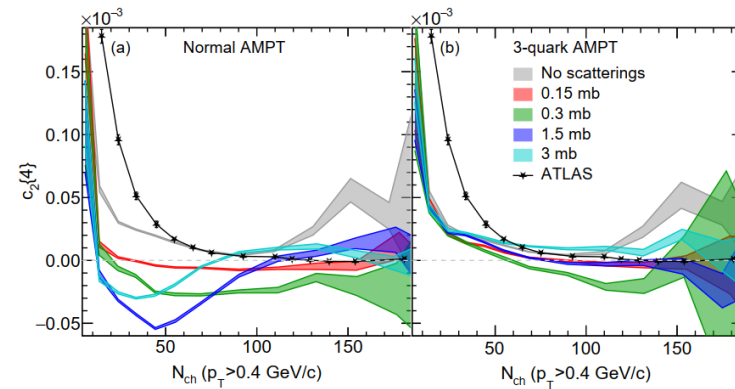
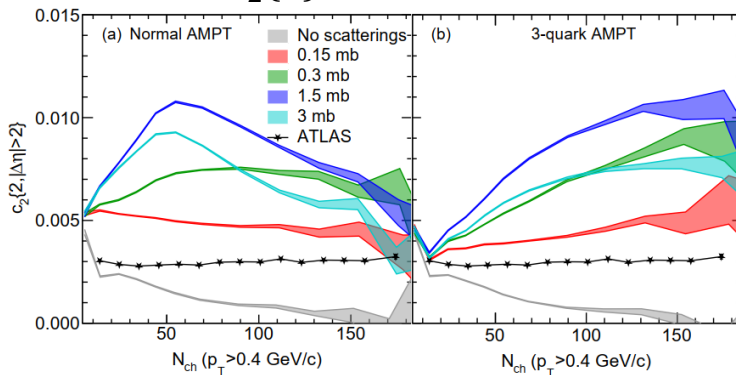
$$c_2\{4\} = -2\langle v_2^2 \rangle^2 + \langle v_2^4 \rangle$$

$$= -\langle v_2 \rangle^4 + 2\sigma_{v_2}^2 \langle v_2 \rangle^2 + \sigma_{v_2}^4$$

$$(v_2 = \langle v_2 \rangle + \sigma_{v_2})$$

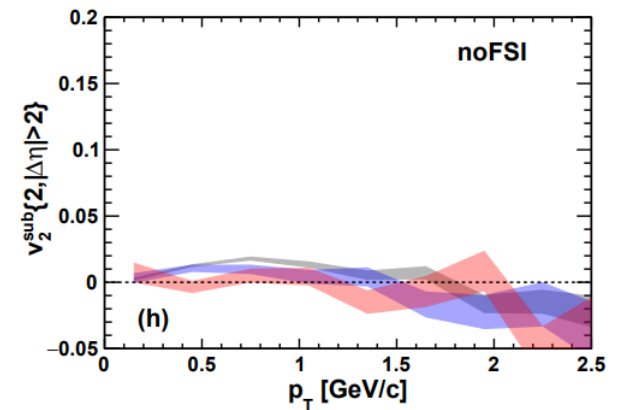
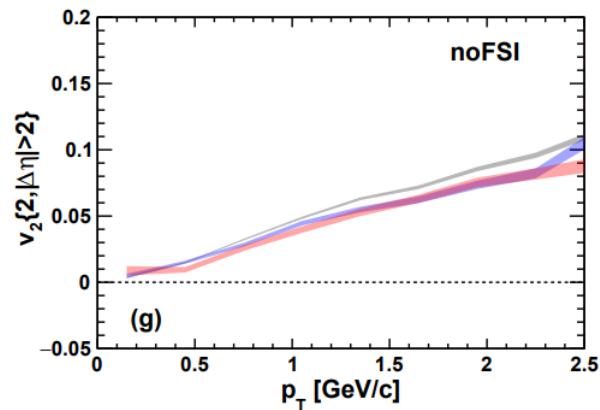
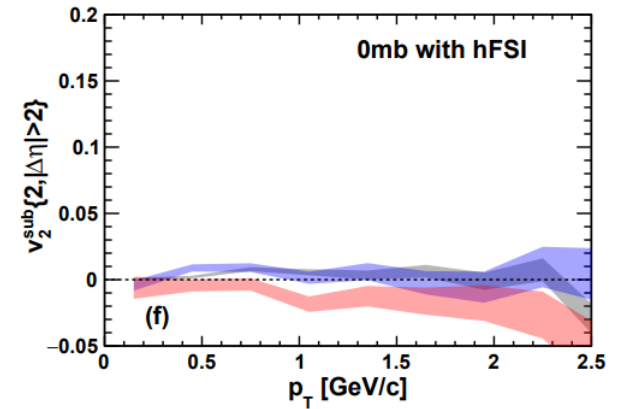
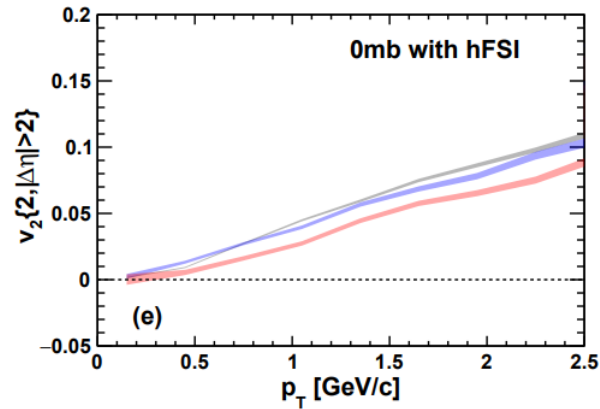
$$v_2\{4\} = (-c_2\{4\})^{\frac{1}{4}}$$

To get the real value of  $v_2\{4\}$ ,  $c_2\{4\}$  should be negative.



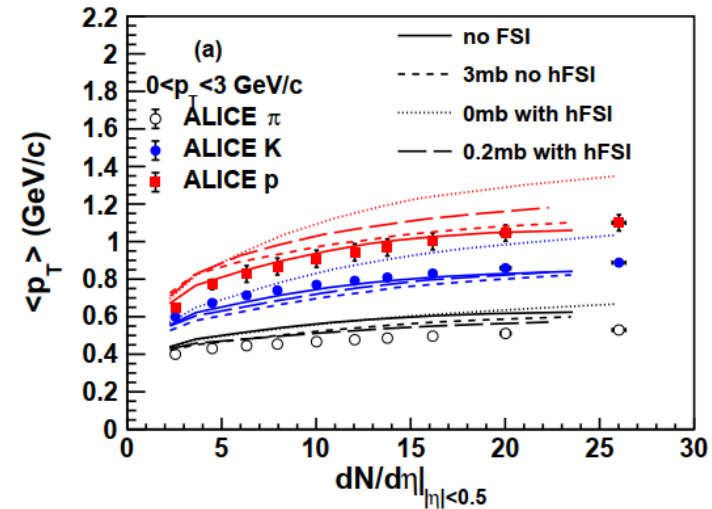
arXiv:2112.01232

# Elliptic flow for identified particles



# Final state interaction(FSI) in pp collisions

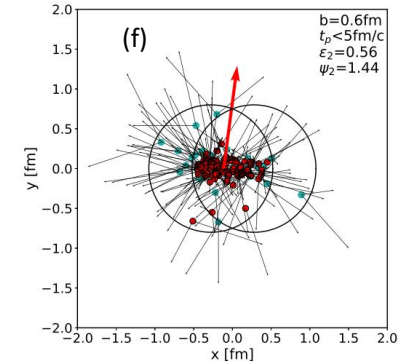
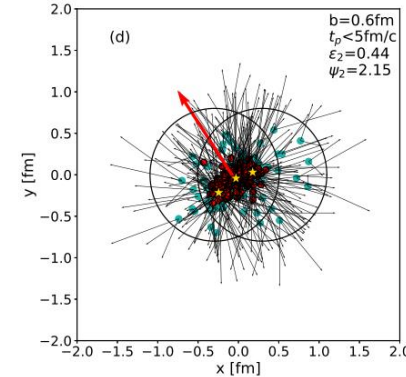
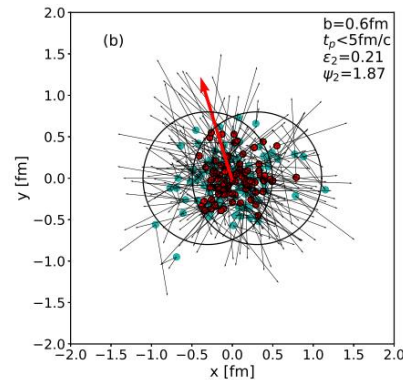
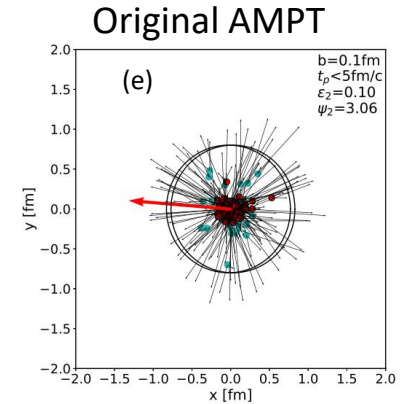
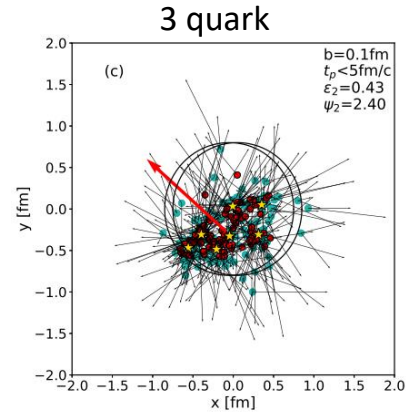
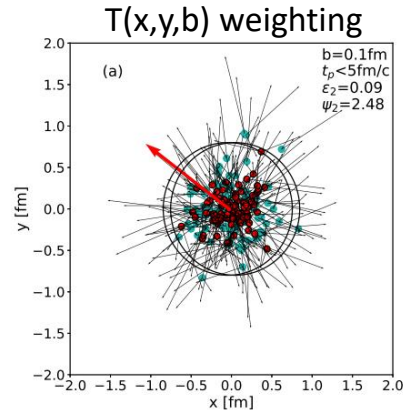
- Average transverse momentum increased due to hadron rescattering
- Proton receives stronger  $p_T$  boost compared to pi/K
- Parton rescattering slightly modifies the spectra



# Space-time structure

pp 13 TeV,  $N_{\text{parton}} > 600$

- T(x, y, b) weighting  
Overlap geometry driven
- 3 quark  
Sub-nucleon fluctuation
- Original AMPT  
Parton sources aligned in a row

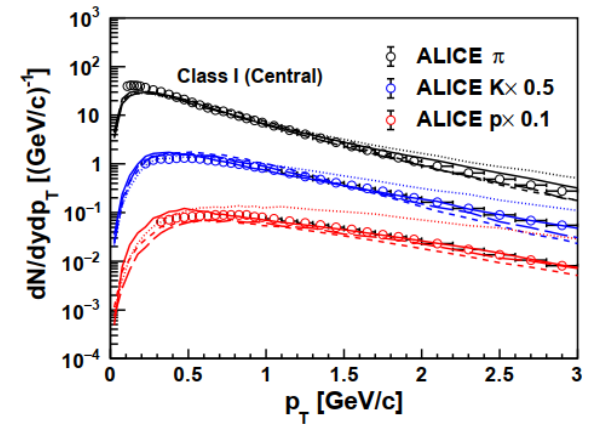
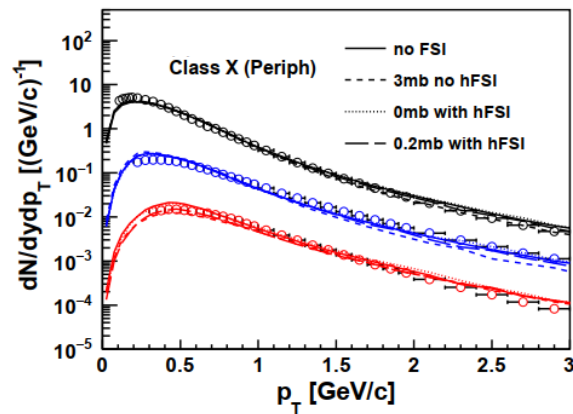
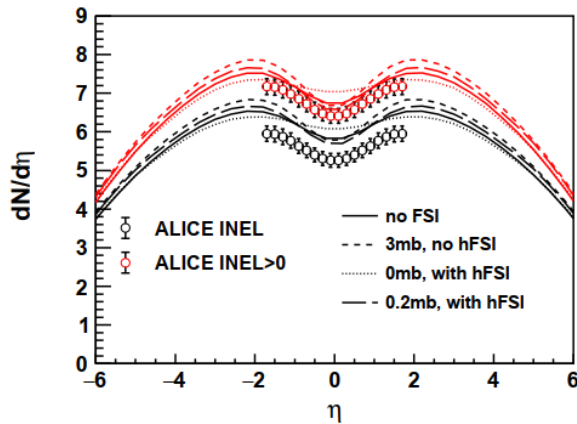


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# Final state interaction(FSI) in pp collisions

- Final state interaction effects weak in charge density
- Stronger FSI in high multiplicity events
- Radial flow like contribution from final state hadron interaction

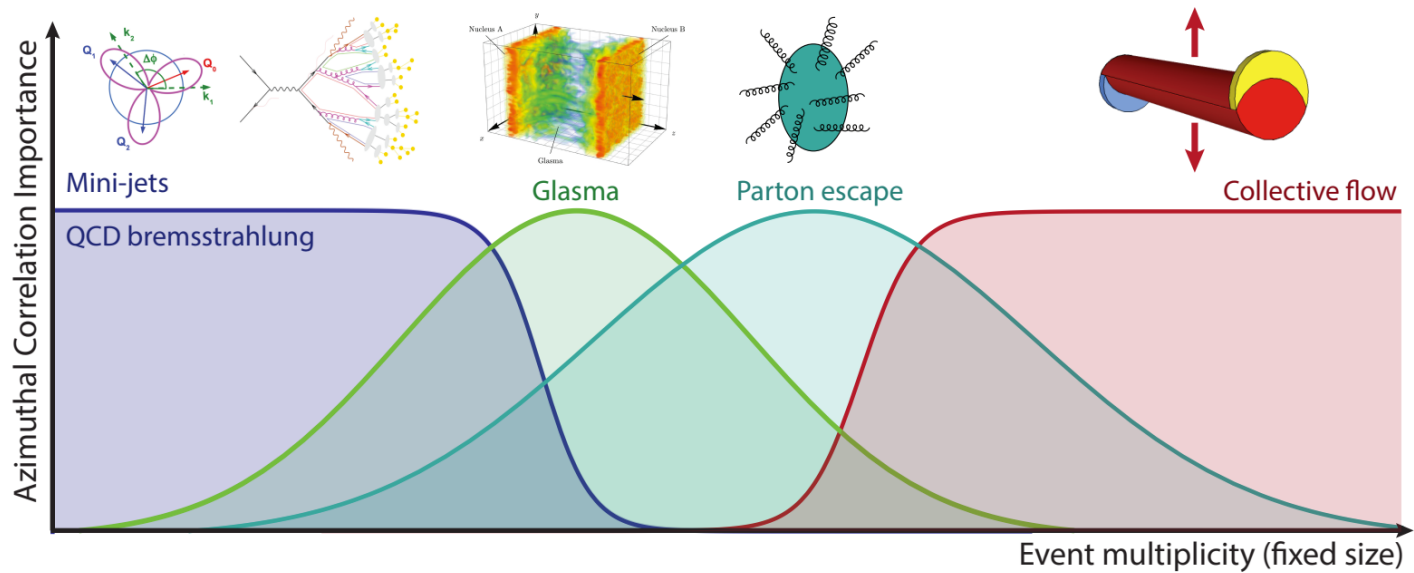
pp 13 TeV, hadron level



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# Collectivity in small system

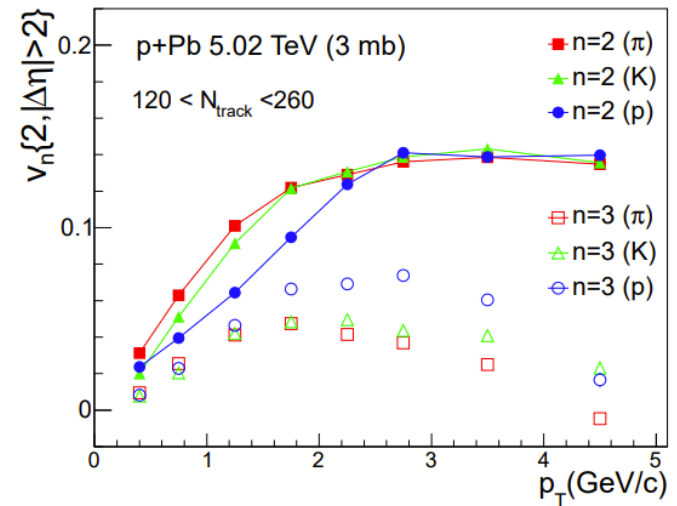
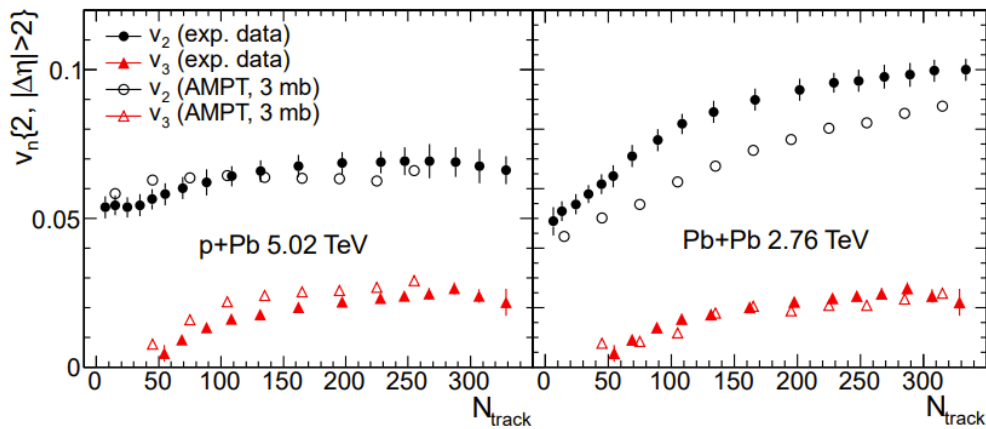
## Multiple sources of azimuthal anisotropy



NPA 982 (2019) 92  
arXiv:1807.07191

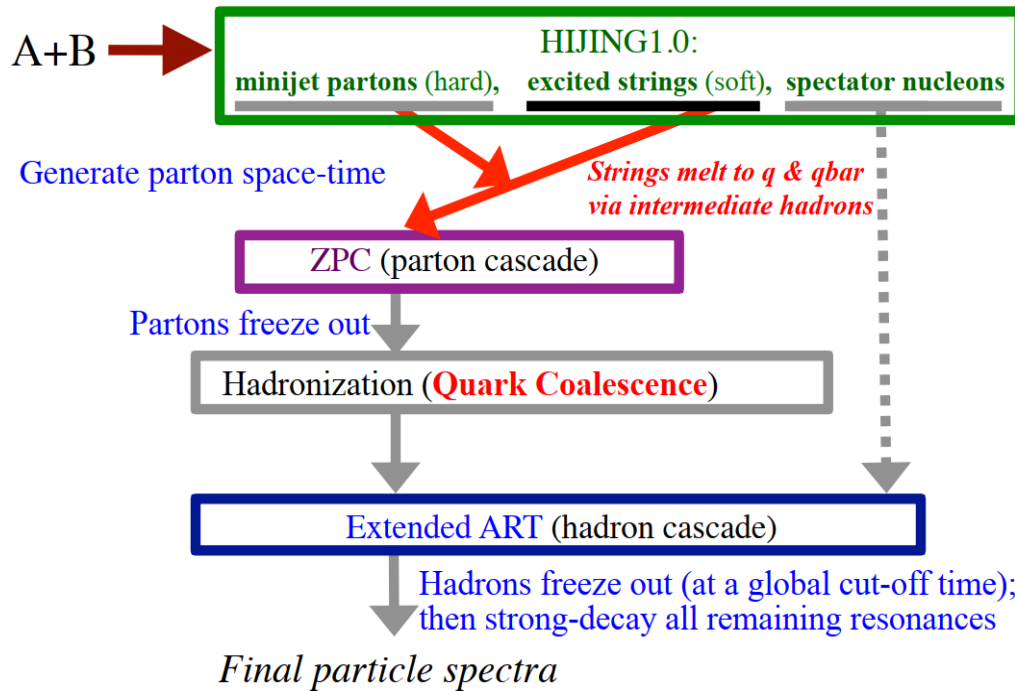
# Collectivity in small system with AMPT

pA collective flow well described by AMPT



PRL 113 (2014) 252301  
arXiv:1406.2804

# A multi-phase transport model (AMPT)



## Hadronization

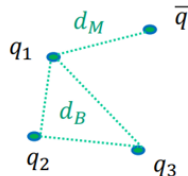
spatial coalescence

*A parton finds the closest partner (meson & baryon) in space to form a hadron (in rest frame):*

$d_M$ : closest anti-quark partner  
 $d_B$ : average distance among 3 quark with closest  $q_2$  &  $q_3$

If  $d_B < d_M * r_{BM}$ ,  $q_1$  will coalesce to a baryon; otherwise,  $q_1$  will coalesce to a meson.

→ Single coalescence parameter  $r_{BM}$



PRC 72 (2005) 064901  
 Nucl.Sci.Tech. 32 (2021) 113

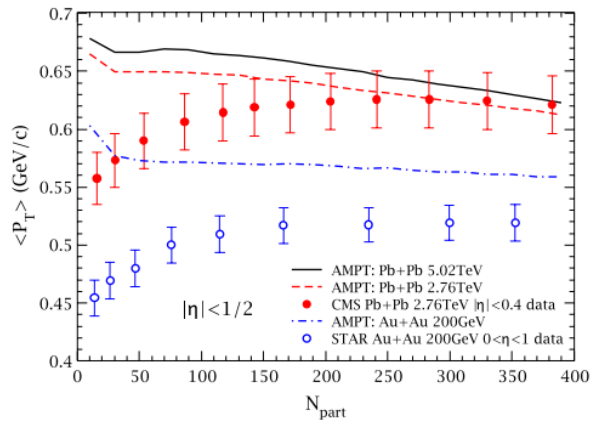


# System size dependence in AMPT

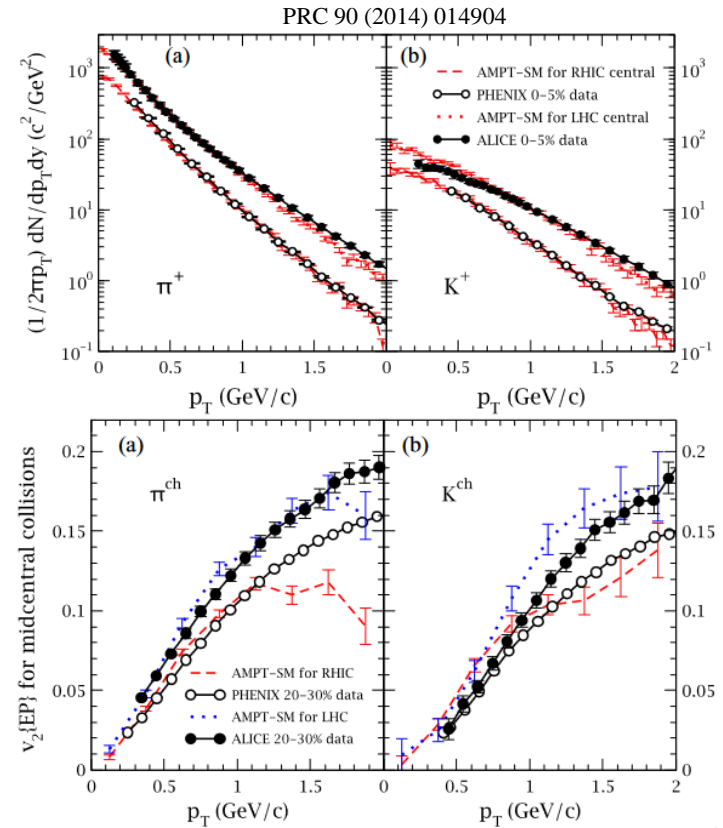
Key model parameters related to initial conditions depend on energy and system size

RHIC 200GeV  $a=0.55$ ,  $b=0.15$

LHC 2.76/5.02TeV  $a=0.3$ ,  $b=0.15$



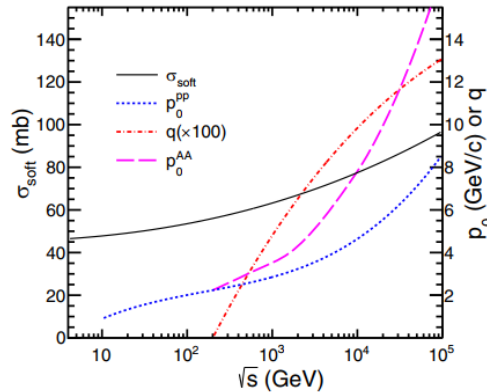
PRC 93 (2016) 054911



PRC 90 (2014) 014904

# System size dependence in AMPT

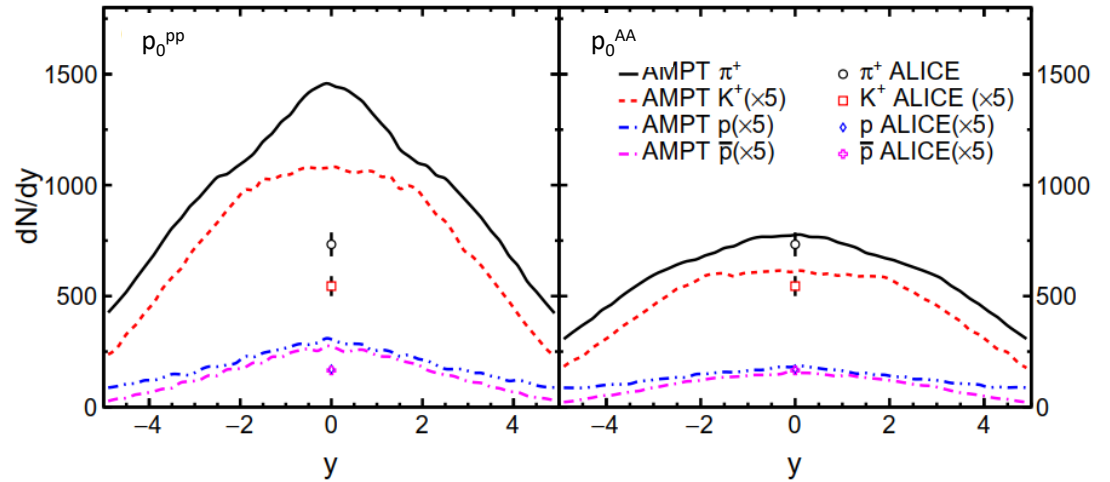
- Unified energy dependence parameterized for pp and central AA collisions



$$p_0^{AA}(s) = p_0^{pp}(s) A^{q(s)}$$

$$Q_s \propto A^{1/6}$$

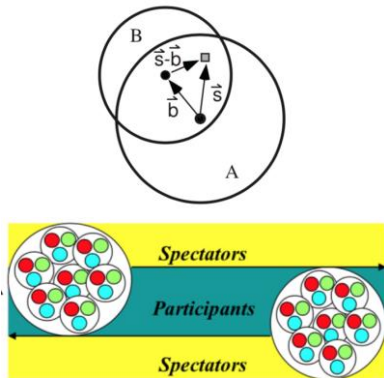
PRC 99 (2019) 6, 064906



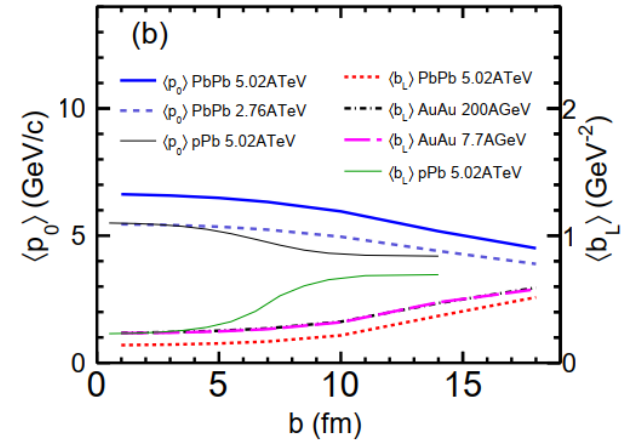
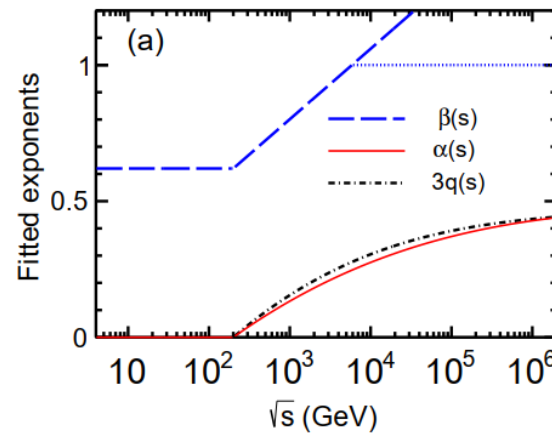
$$q(s) = 0.0334 \ln\left(\frac{\sqrt{s}}{E_0}\right) - 0.00232 \ln^2\left(\frac{\sqrt{s}}{E_0}\right) + 0.0000541 \ln^3\left(\frac{\sqrt{s}}{E_0}\right), \text{ for } \sqrt{s} \geq E_0 = 200 \text{ GeV}.$$

# System size dependence in AMPT

- Local thickness scaling



PRC 104 (2021) 1, 014908



$$p_0(s_A, s_B, s) = p_0^{pp}(s) \left[ \sqrt{T_A(s_A)T_B(s_B)}/T_p \right]^{\alpha(s)}$$

$$b_L(s_A, s_B, s) = \frac{b_L^{pp}}{\left[ \sqrt{T_A(s_A)T_B(s_B)}/T_p \right]^{\beta(s)}}$$

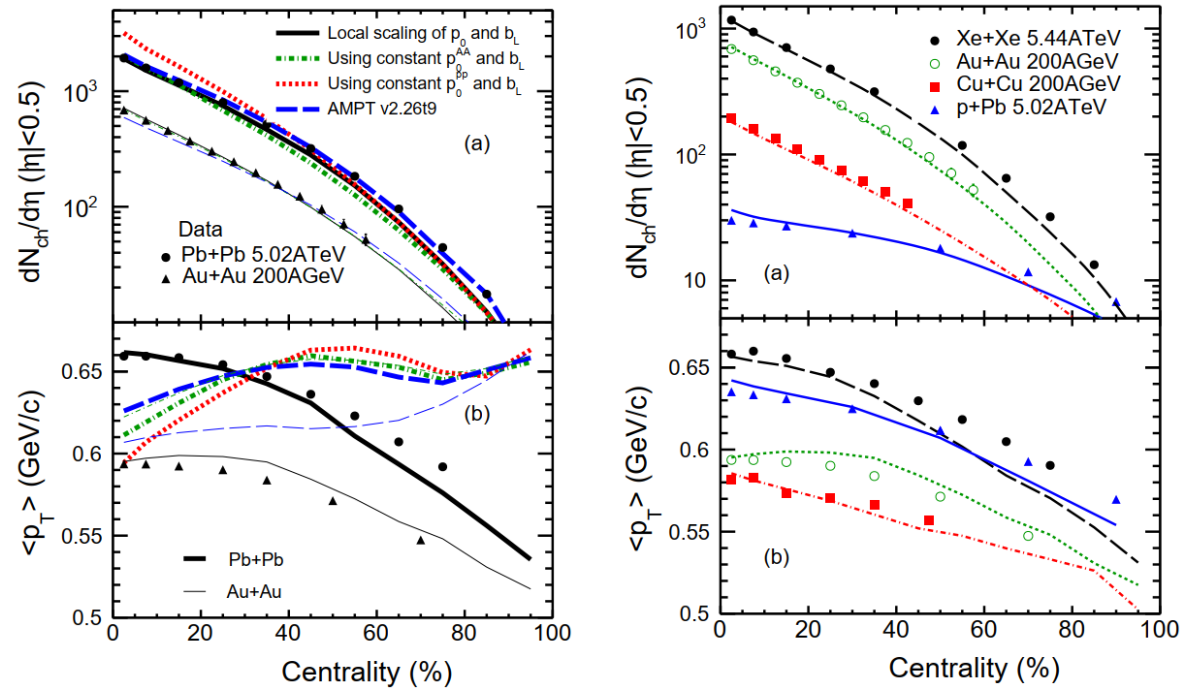
$$\alpha(s) = 0.0918 \ln \left( \frac{\sqrt{s}}{E_0} \right) - 0.00602 \ln^2 \left( \frac{\sqrt{s}}{E_0} \right) + 0.000134 \ln^3 \left( \frac{\sqrt{s}}{E_0} \right), \text{ for } \sqrt{s} \geq E_0,$$

$$\beta(s) = 0.620 + 0.112 \ln \left( \frac{\sqrt{s}}{E_0} \right) \Theta(\sqrt{s} - E_0)$$

# System size dependence in AMPT

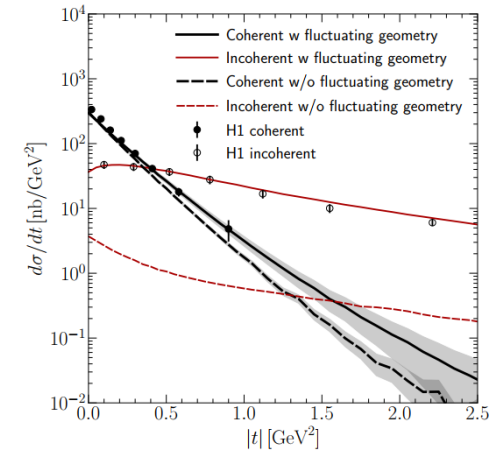
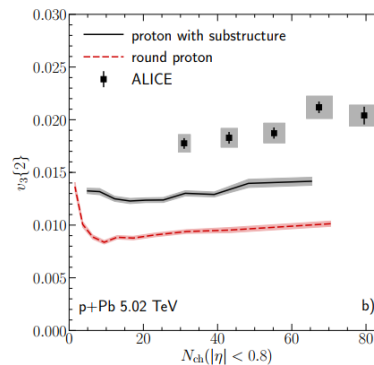
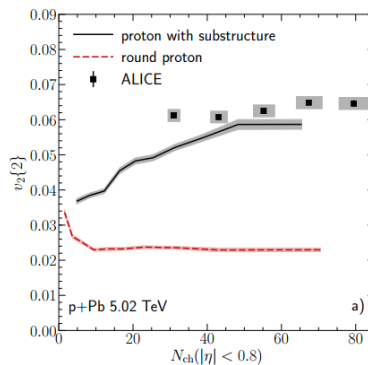
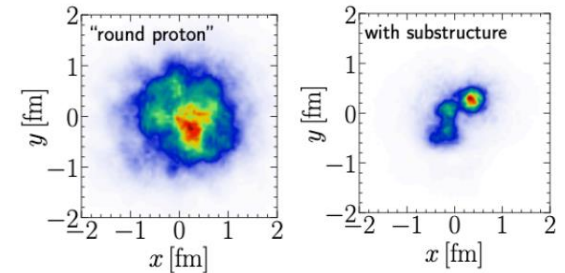
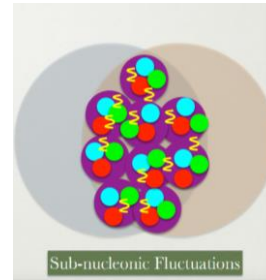
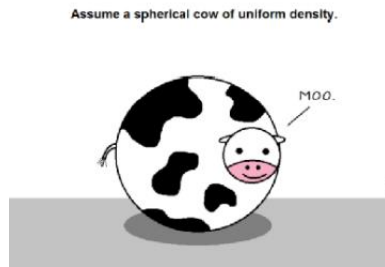
- System size dependence from pA to AA implemented via the local thickness scaling
- Charged hadron productions across different systems well described
- No more system dependent free model parameters

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# Sub-nucleon spatial fluctuations

- Modeling proton spatial structure with constituent quark picture
- Sub-nucleon fluctuation can be important for small system collectivity



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# Initial geometry configuration for pp

**Proton substructure**  $\rho(r) = \frac{1}{8\pi R^3} e^{-r/R}$

- Overlap function  $T(x,y,b)$  weighting method

$$T(x,y,b) = \int \rho_{p,1}(x - b/2, y, z) \rho_{p,2}(x + b/2, y, z) dz$$

- Constituent quark method (3 quark)

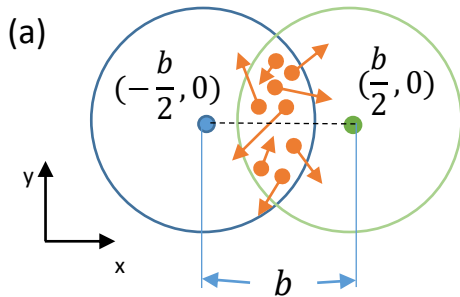
Glauber modeling with 3 quark participants

Collision criteria  $d < \sqrt{\sigma_{cc}/\pi}$

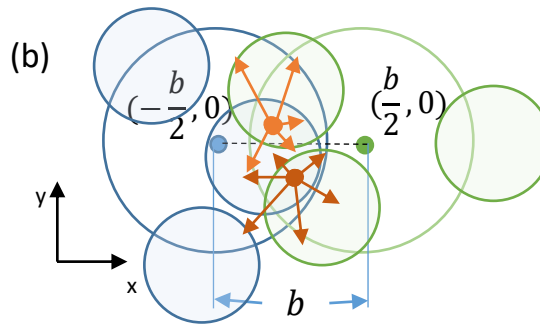
**Original AMPT spatial structures**

- Target/projectile string  
Target/projectile position
- Independent string  
Mid-point position

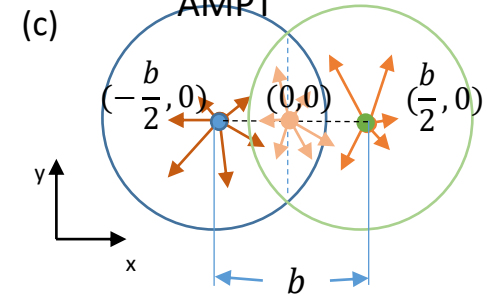
$T(x,y,b)$  weighting



3 quark constituent

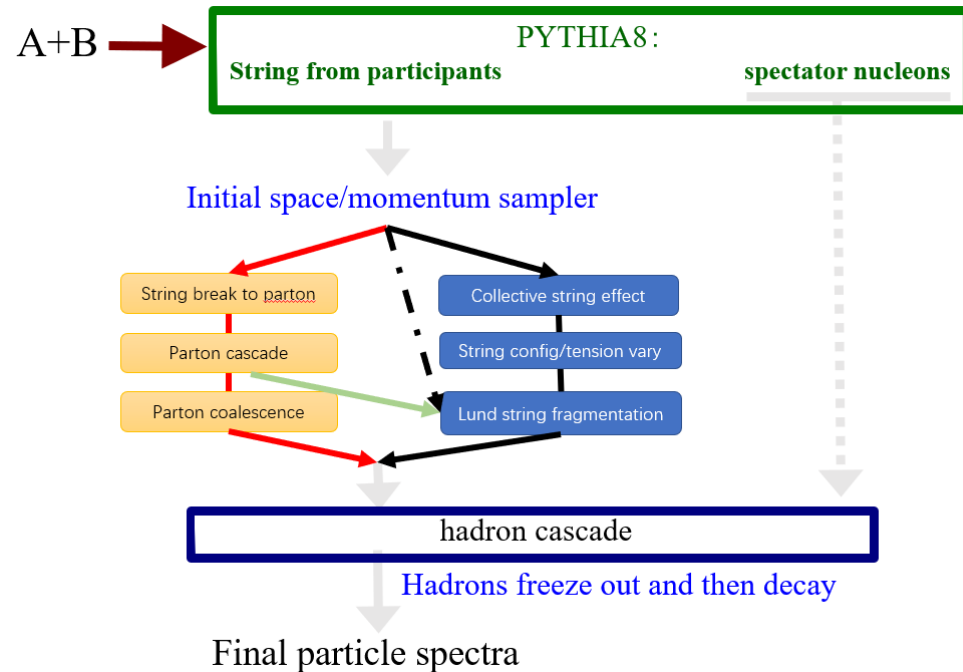


Original AMPT



# PYTHIA8 initial conditions

- PYTHIA8 initial conditions: modern parton distribution function, multiple parton interaction (MPI) based string system
- Possible to compare parton transport and inter-string effects on the same basis



# Parton escape effects

