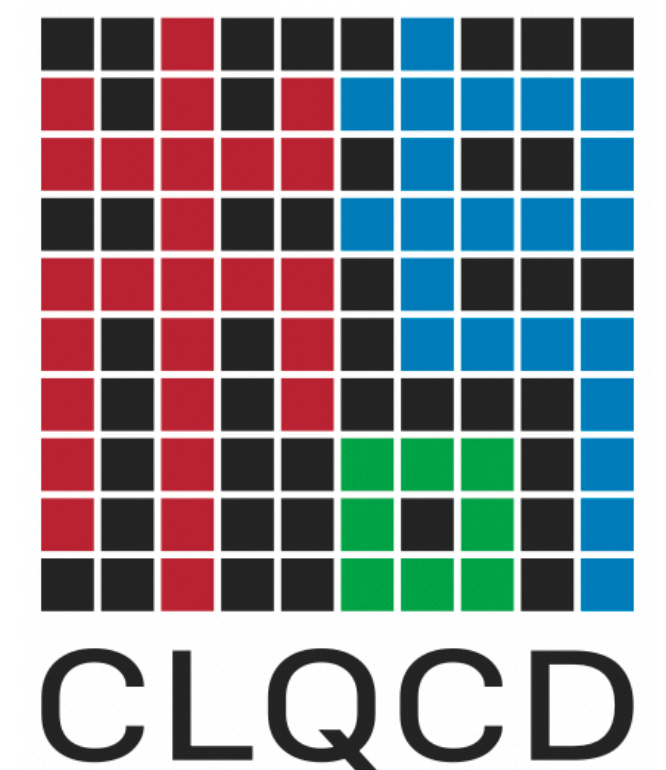


Lattice QCD study of the Exotic states



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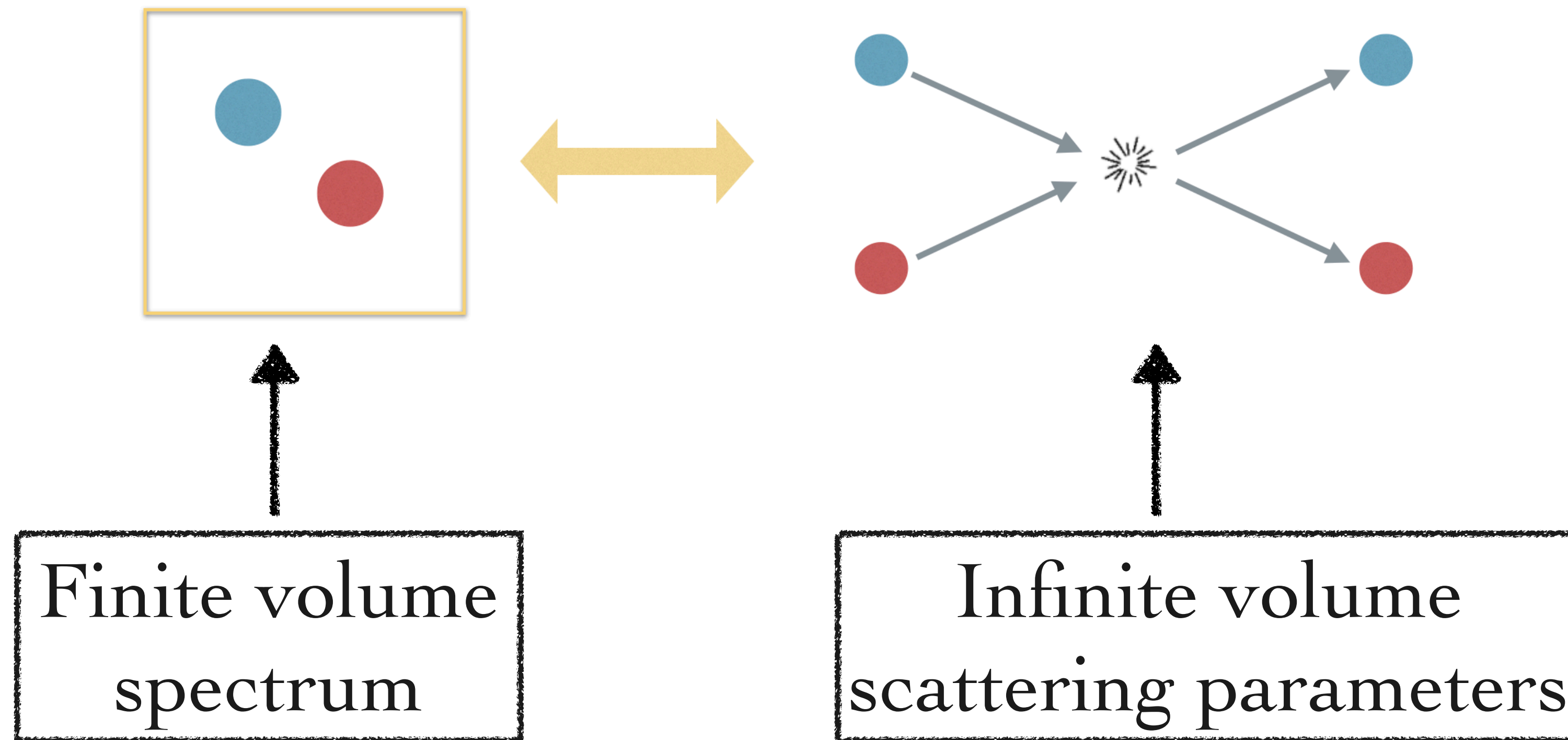


Introduction



- ◆ Since 2003, large amount of new hadronic states beyond the conventional quark model are discovered in experiments.
- ◆ Nearly all of the exotic candidates are close to the thresholds of two hadrons.
- ◆ Study of the hadron scattering is a main approach to probe the structure of the exotic hadrons.
 - Scattering on lattice.
 - Preliminary results on the hidden-charm pentaquarks.
 - Preliminary results on $D\pi$ scattering.

Lüscher's finite volume method: M. Lüscher, Nucl. Phys. B354, 531(1991)



- ◆ General Lüscher's formula for two-body scattering:

$$\det[\mathbf{1} + i\rho \cdot \mathbf{t} \cdot (1 + i\mathbf{M})] = 0$$

Diagonal matrix of
phase-space factors

$$\rho_{ij} = \delta_{ij} \frac{2k_i}{E_{cm}}$$

Infinite-volume
scattering matrix

Finite volume
information

$$M(E_{cm}, L)$$

- ◆ Resonances/bound states are formally defined as poles in scattering amplitudes.



Scattering on lattice



Finite volume spectrum:

- ◆ build large basis of operators $\{\mathcal{O}_1, \mathcal{O}_2, \dots\}$ with desired quantum numbers, construct the matrix of correlation function:

$$C_{ij} = \langle 0 | \mathcal{O}_i \mathcal{O}_j^\dagger | 0 \rangle = \sum_n Z_i^n Z_j^{n*} e^{-E_n t}$$

- ◆ Solve the generalized eigenvalue problem(GEVP): $C_{ij} v_j^n(t) = \lambda_n(t) C_{ij}^0 v_j^n(t)$
- ◆ Eigenvalues: $\lambda_n(t) \sim e^{-E_n t} (1 + e^{-\Delta E t})$
- ◆ Optimal linear combinations of the operators to overlap on the n'th state:

$$\Omega_n = \sum_i v_i^n \mathcal{O}_i$$

- ◆ Computational technique: distillation quark smearing.
 - Improve precision
 - Disconnected diagrams
 - Efficient for large numbers of ops

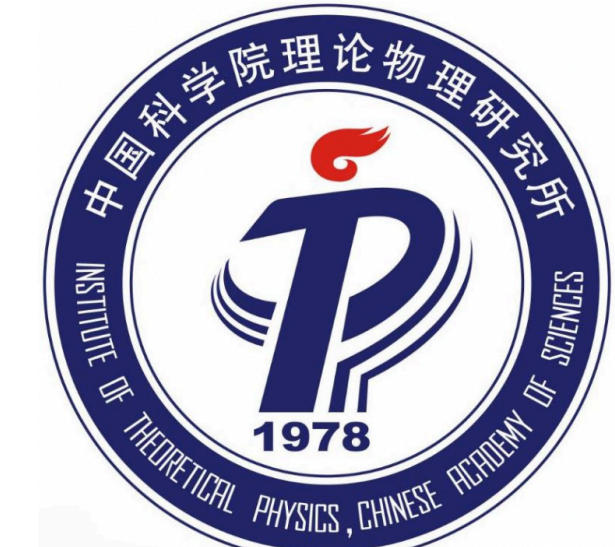


Lattice QCD configurations



- 2+1 flavor Wilson-clover configurations generated by CLQCD.

| Lattice spacing | Volume($L^3 \times T$) | M_π (MeV) | # of confs |
|-----------------|--------------------------|---------------|------------|
| ~0.108fm | $24^3 \times 72$ | 290 | 1000 |
| | $32^3 \times 64$ | 290 | 1000 |
| | $32^3 \times 64$ | 220 | 450 |
| | $48^3 \times 96$ | 220 | 200 |
| | $48^3 \times 96$ | 140 | 200 |
| ~0.080fm | $32^3 \times 96$ | 300 | 480 |
| | $48^3 \times 96$ | 300 | 200 |
| | $32^3 \times 64$ | 220 | 460 |
| | $48^3 \times 96$ | 220 | 200 |
| ~0.055fm | $48^3 \times 144$ | 300 | 200 |



arXiv: 2310.00814

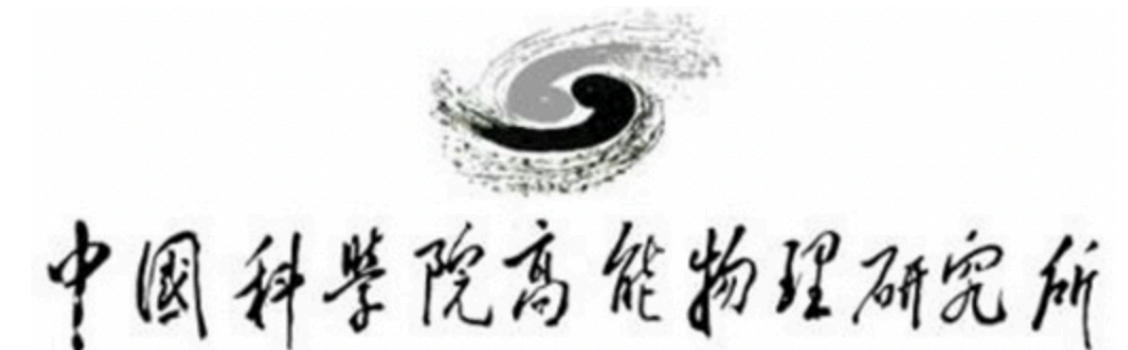
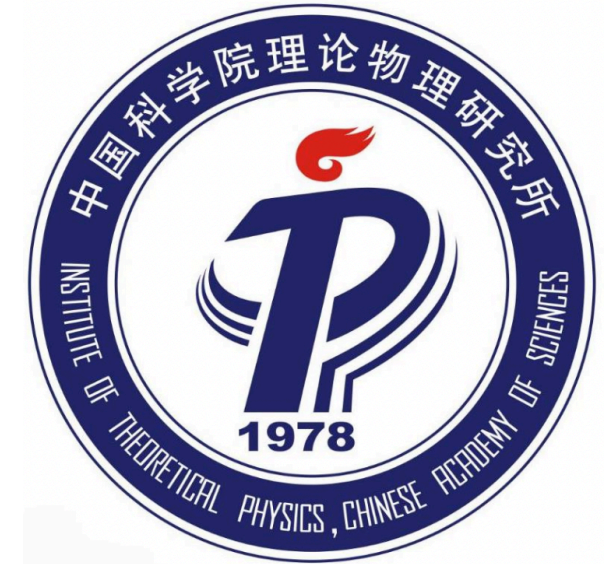


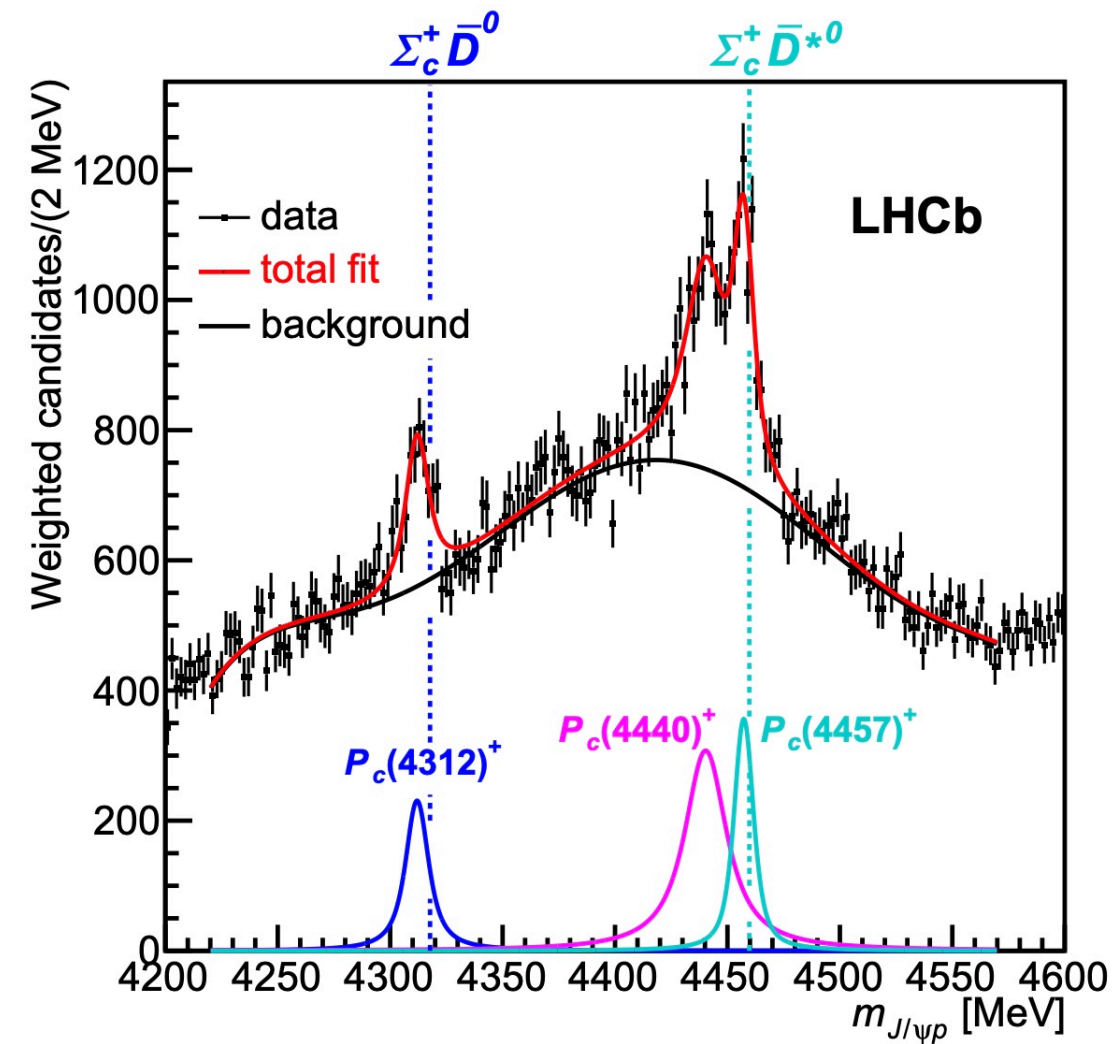
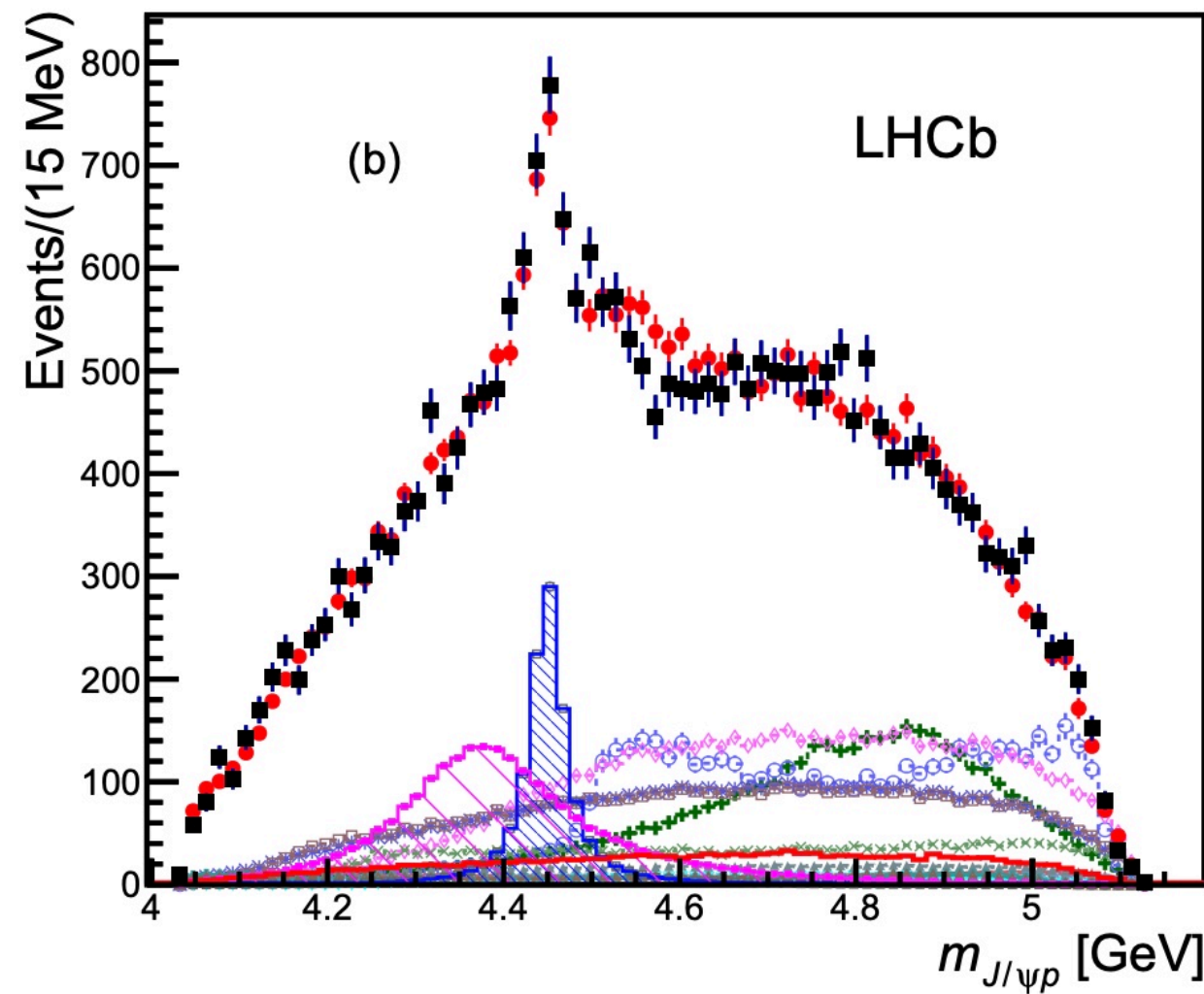
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$P_c(4380)$ \rightarrow $P_c(4312)$
 $P_c(4450)$ \rightarrow $P_c(4440)$
 $P_c(4450)$ \rightarrow $P_c(4457)$

R. Aaij et al. (LHCb), Phys. Rev. Lett. 115, 072001 (2015)

R. Aaij et al. (LHCb), Phys. Rev. Lett. 122, 222001 (2019)

Theory interpretations:

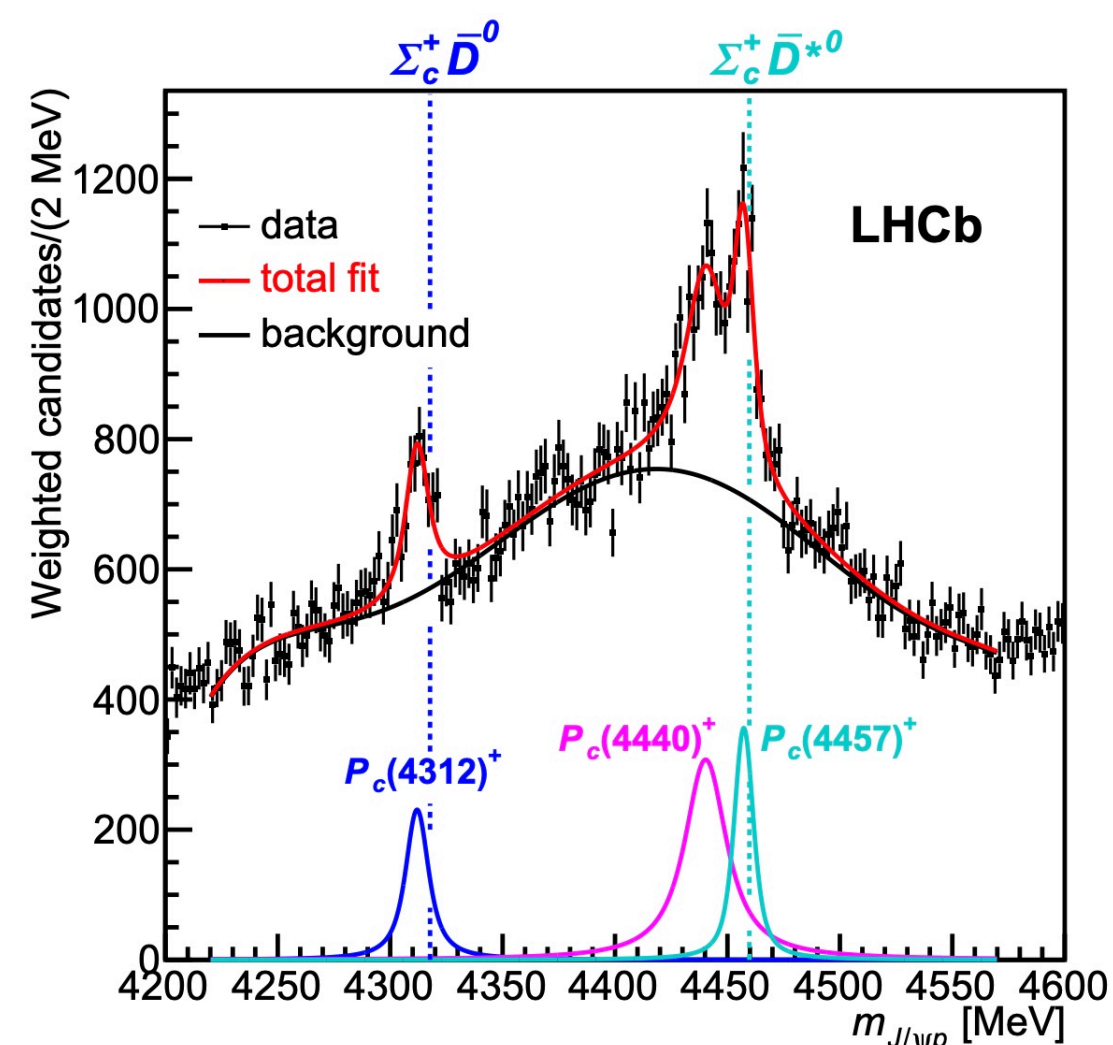
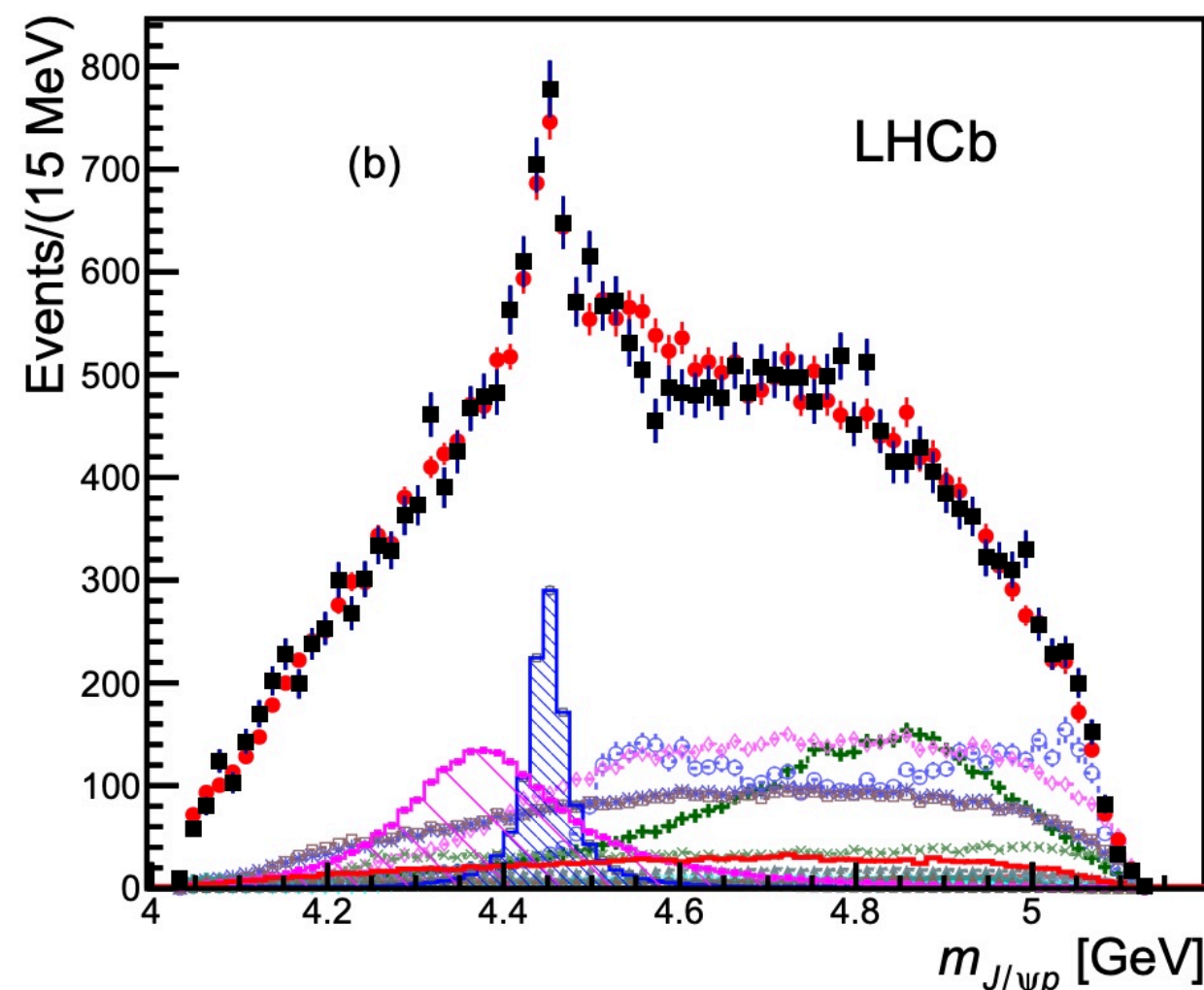
- Molecule bound states
- Compact pentaquark states
-

$\Sigma_c^{(*)} \bar{D}^{(*)}$ molecules:

$$\Sigma_c \bar{D}, J^P = \frac{1}{2}^-, P_c(4312)$$

$$\Sigma_c \bar{D}^*, J^P = \left(\frac{1}{2}^-, \frac{3}{2}^- \right), P_c(4440)/P_c(4457)$$

$$\Sigma_c^* \bar{D}, J^P = \frac{3}{2}^-, \quad \Sigma_c^* \bar{D}^*, J^P = \left(\frac{1}{2}^-, \frac{3}{2}^-, \frac{5}{2}^- \right),$$



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$\Sigma_c \bar{D}$ and $\Sigma_c \bar{D}^*$ scattering ($J^P = \frac{1}{2}^-$):

◆ Five operators:

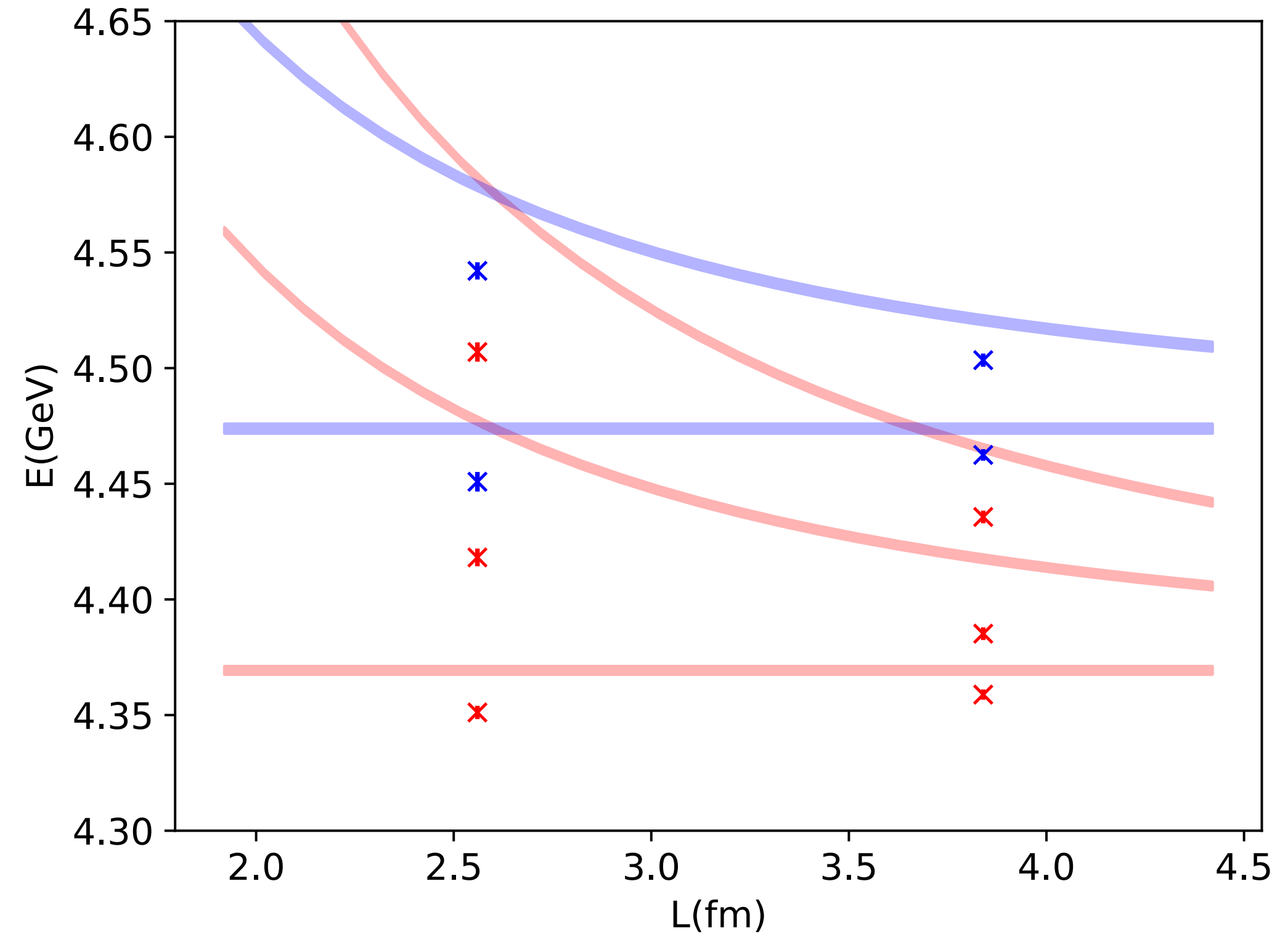
$$\mathcal{O}_1 = \Sigma_c(\mathbf{p}) \bar{D}(-\mathbf{p}) \quad (|\mathbf{p}| = 0)$$

$$\mathcal{O}_2 = \Sigma_c(\mathbf{p}) \bar{D}(-\mathbf{p}) \quad (|\mathbf{p}| = 1)$$

$$\mathcal{O}_3 = \Sigma_c(\mathbf{p}) \bar{D}(-\mathbf{p}) \quad (|\mathbf{p}| = \sqrt{2})$$

$$\mathcal{O}_4 = \Sigma_c(\mathbf{p}) \bar{D}^*(-\mathbf{p}) \quad (|\mathbf{p}| = 0)$$

$$\mathcal{O}_5 = \Sigma_c(\mathbf{p}) \bar{D}^*(-\mathbf{p}) \quad (|\mathbf{p}| = 1)$$



◆ The finite-volume energies lie below the free energies, indicating rather strong attractive interactions.

Scattering amplitude:

$$T \sim \frac{1}{p \cot \delta - ip}$$

Effective range expansion:

$$p \cot \delta(p) = \frac{1}{a_0} + \frac{1}{2} r_0 p^2 + \dots$$

Lüscher's formula:

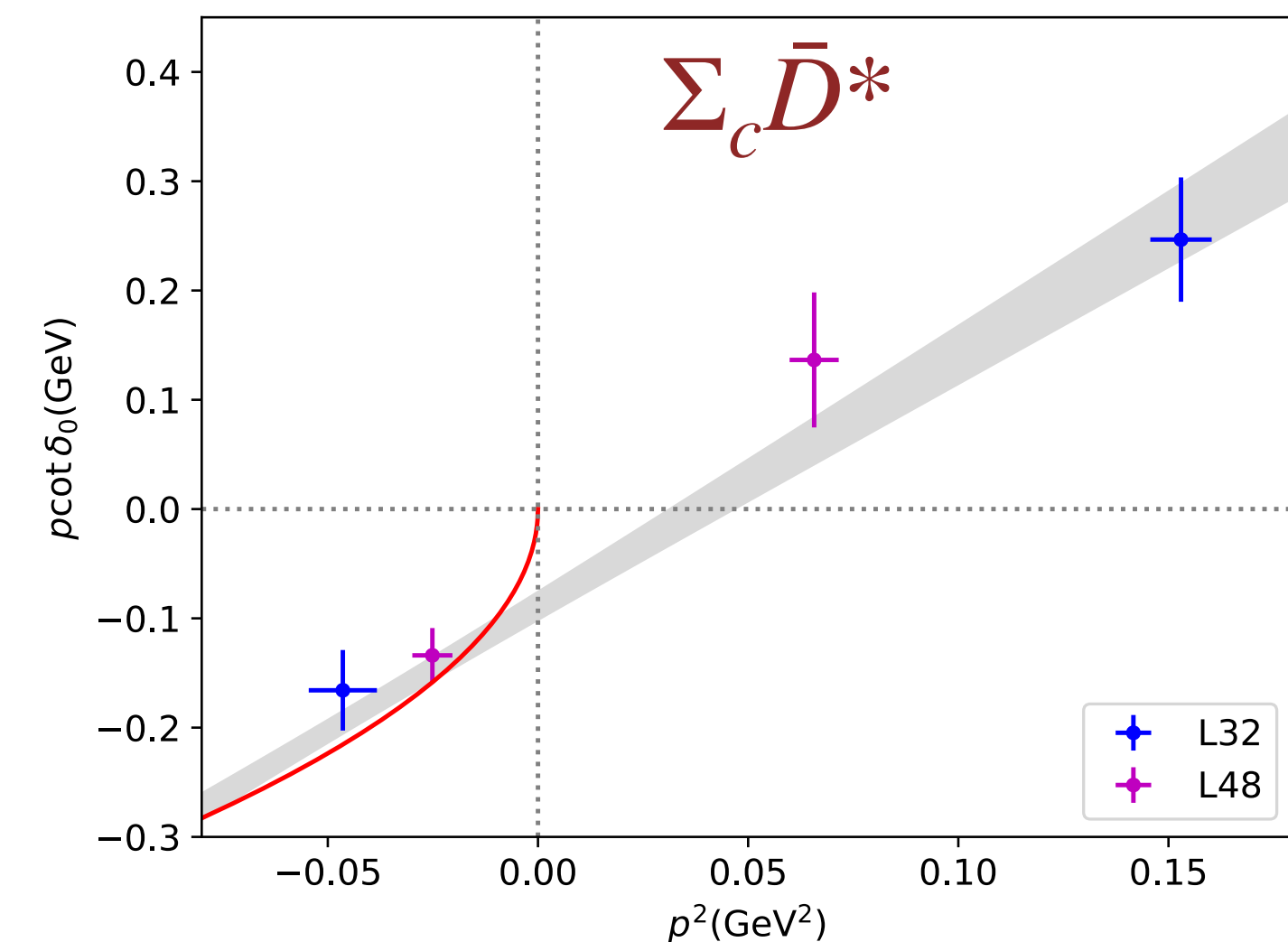
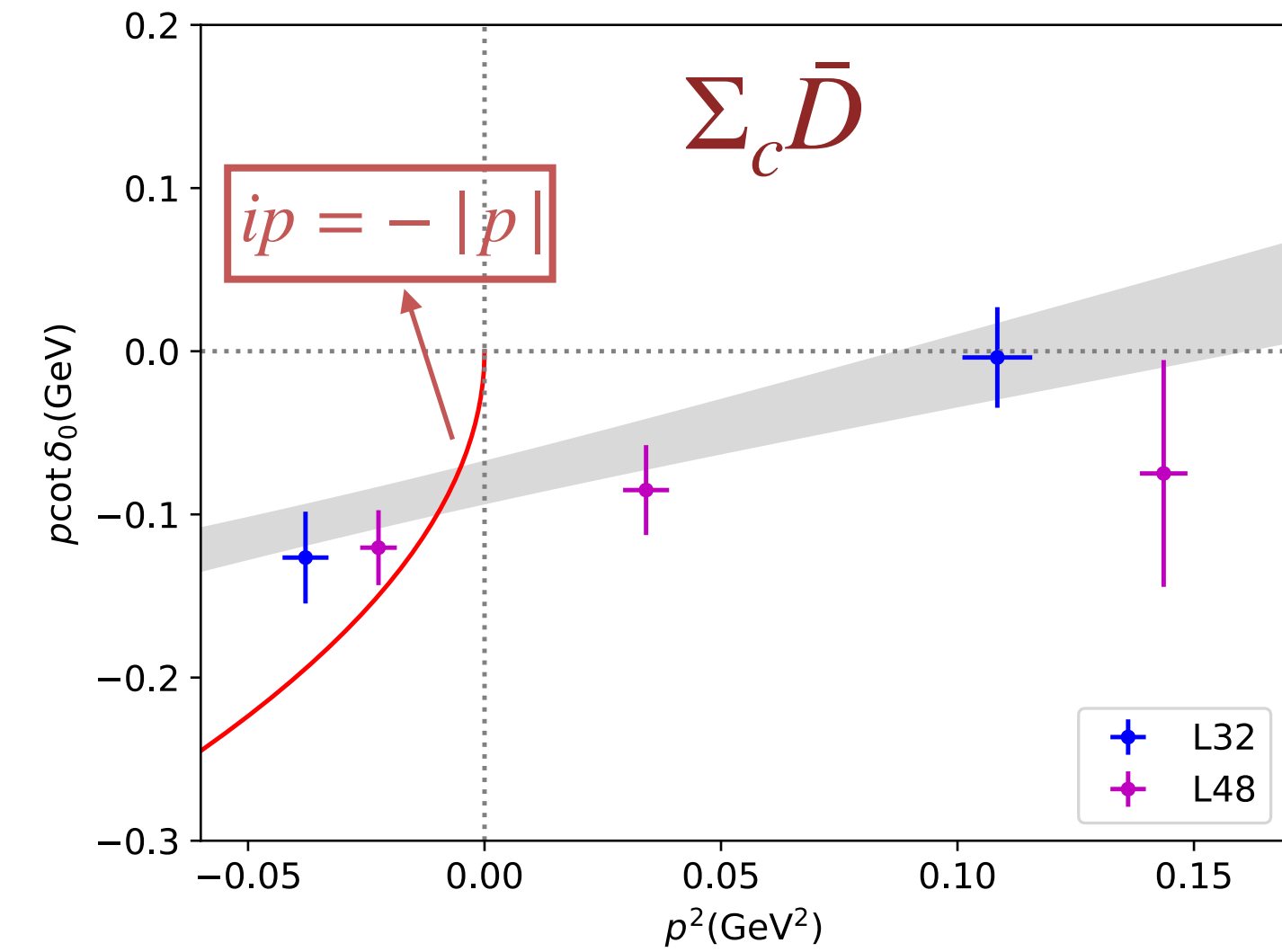
$$p \cot \delta(p) = \frac{2Z_{00}(1; (\frac{pL}{2\pi})^2)}{L\sqrt{\pi}}$$

Bound state pole:

$$p = i|p_B|$$

$$\begin{aligned} \Sigma_c \bar{D} : P_c(4312) ? \\ a_0 = -2.0(3)(5) \text{ fm} \\ E_B = 6(2)(2) \text{ MeV} \end{aligned}$$

$$\begin{aligned} \Sigma_c \bar{D}^* : P_c(4440)/P_c(4457) ? \\ a_0 = -2.3(5)(1) \text{ fm} \\ E_B = 7(3)(1) \text{ MeV} \end{aligned}$$





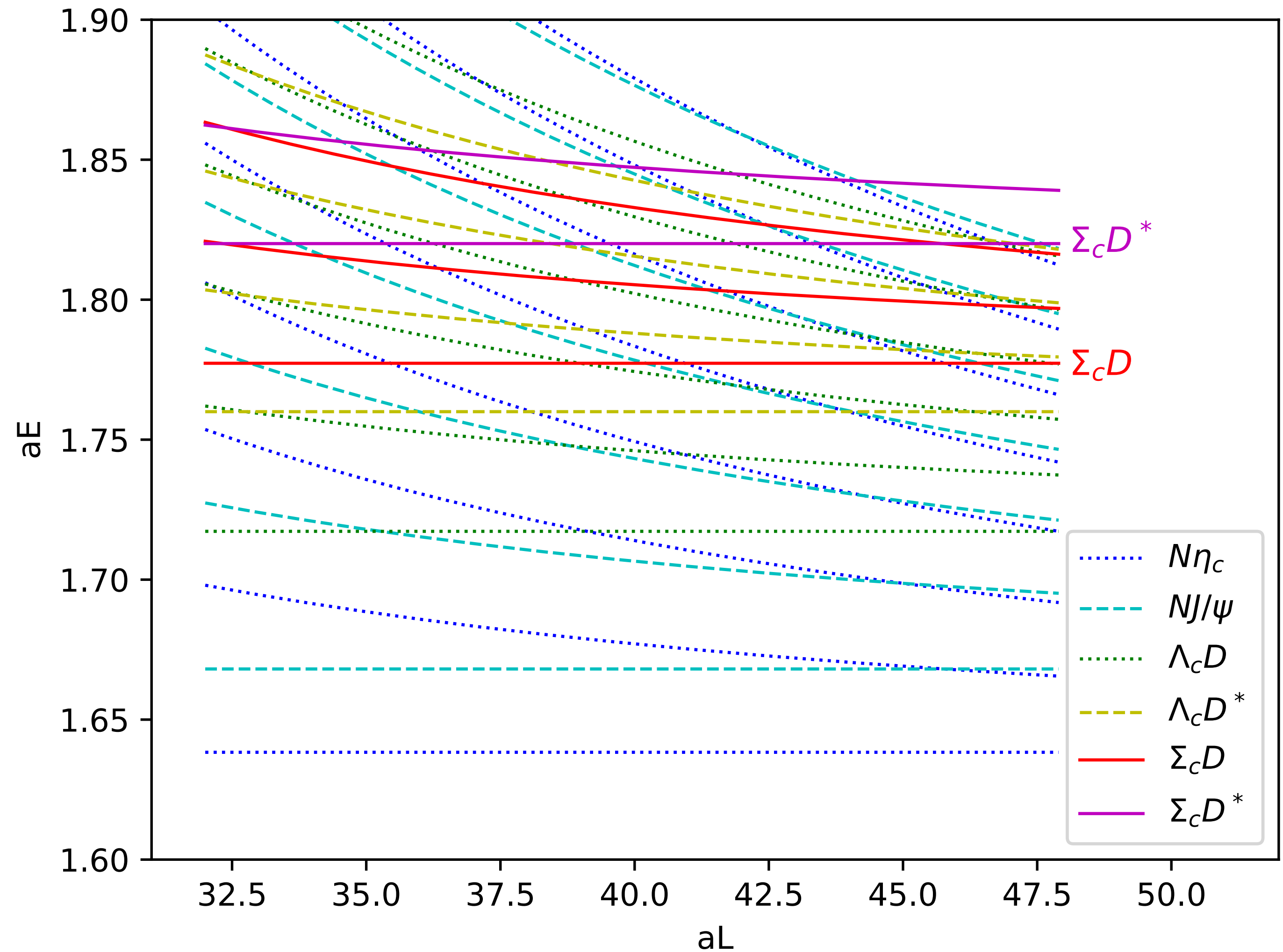
Coupled channels



Coupled channels: $\eta_c N, J/\psi N, \Lambda_c \bar{D}, \Lambda_c \bar{D}^*, \Sigma_c \bar{D}, \Sigma_c \bar{D}^*$

Non-interacting energy levels:

$$E_{free} = \sqrt{m_1^2 + p_1^2} + \sqrt{m_1^2 + p_2^2}$$





Coupled channels



Coupled channels: $\eta_c N, J/\psi N, \Lambda_c \bar{D}, \Lambda_c \bar{D}^*, \Sigma_c \bar{D}, \Sigma_c \bar{D}^*$

◆ 15 operators for the $L = 32$ ensemble:

$$\mathcal{O}_{1,2,3} = N(\mathbf{p})\eta_c(-\mathbf{p}) \quad (\mathbf{p}^2 = 0,1,2)$$

$$\mathcal{O}_{4,5} = N(\mathbf{p})J/\psi(-\mathbf{p}) \quad (\mathbf{p}^2 = 0,1)$$

$$\mathcal{O}_{6,7,8} = \Lambda_c(\mathbf{p})\bar{D}(-\mathbf{p}) \quad (\mathbf{p}^2 = 0,1,2)$$

$$\mathcal{O}_{9,10} = \Lambda_c(\mathbf{p})\bar{D}^*(-\mathbf{p}) \quad (\mathbf{p}^2 = 0,1)$$

$$\mathcal{O}_{11,12,13} = \Sigma_c(\mathbf{p})\bar{D}(-\mathbf{p}) \quad (\mathbf{p}^2 = 0,1,2)$$

$$\mathcal{O}_{14,15} = \Sigma_c(\mathbf{p})\bar{D}^*(-\mathbf{p}) \quad (\mathbf{p}^2 = 0,1)$$

◆ 23 operators for the $L = 48$ ensemble:

$$\mathcal{O}_{1,2,3,4,5} = N(\mathbf{p})\eta_c(-\mathbf{p}) \quad (\mathbf{p}^2 = 0,1,2,3,4)$$

$$\mathcal{O}_{7,8,9,10} = N(\mathbf{p})J/\psi(-\mathbf{p}) \quad (\mathbf{p}^2 = 0,1,2,3)$$

$$\mathcal{O}_{10,11,12,13,14} = \Lambda_c(\mathbf{p})\bar{D}(-\mathbf{p}) \quad (\mathbf{p}^2 = 0,1,2,3,4)$$

$$\mathcal{O}_{15,16,17,18} = \Lambda_c(\mathbf{p})\bar{D}^*(-\mathbf{p}) \quad (\mathbf{p}^2 = 0,1,2,3)$$

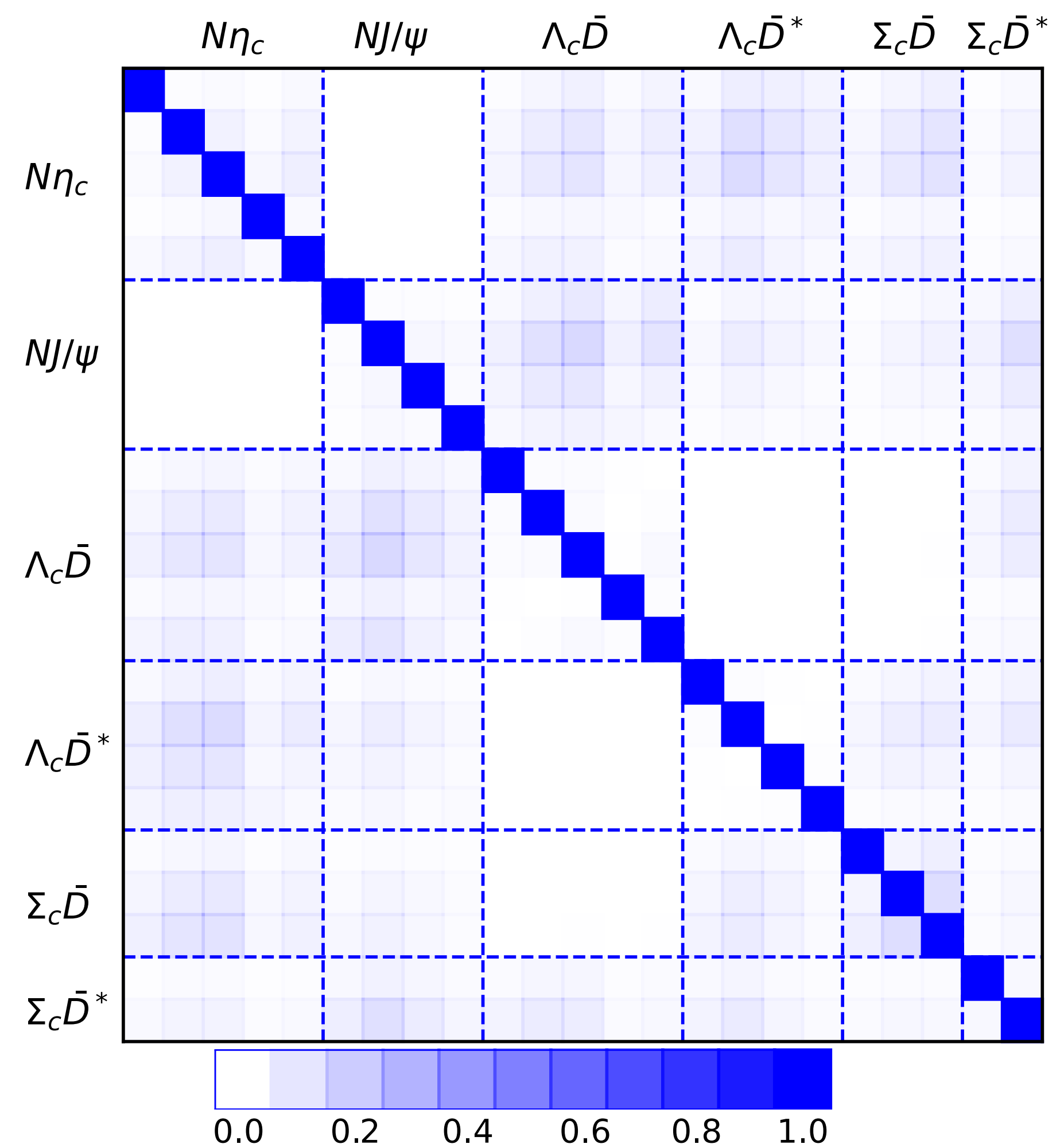
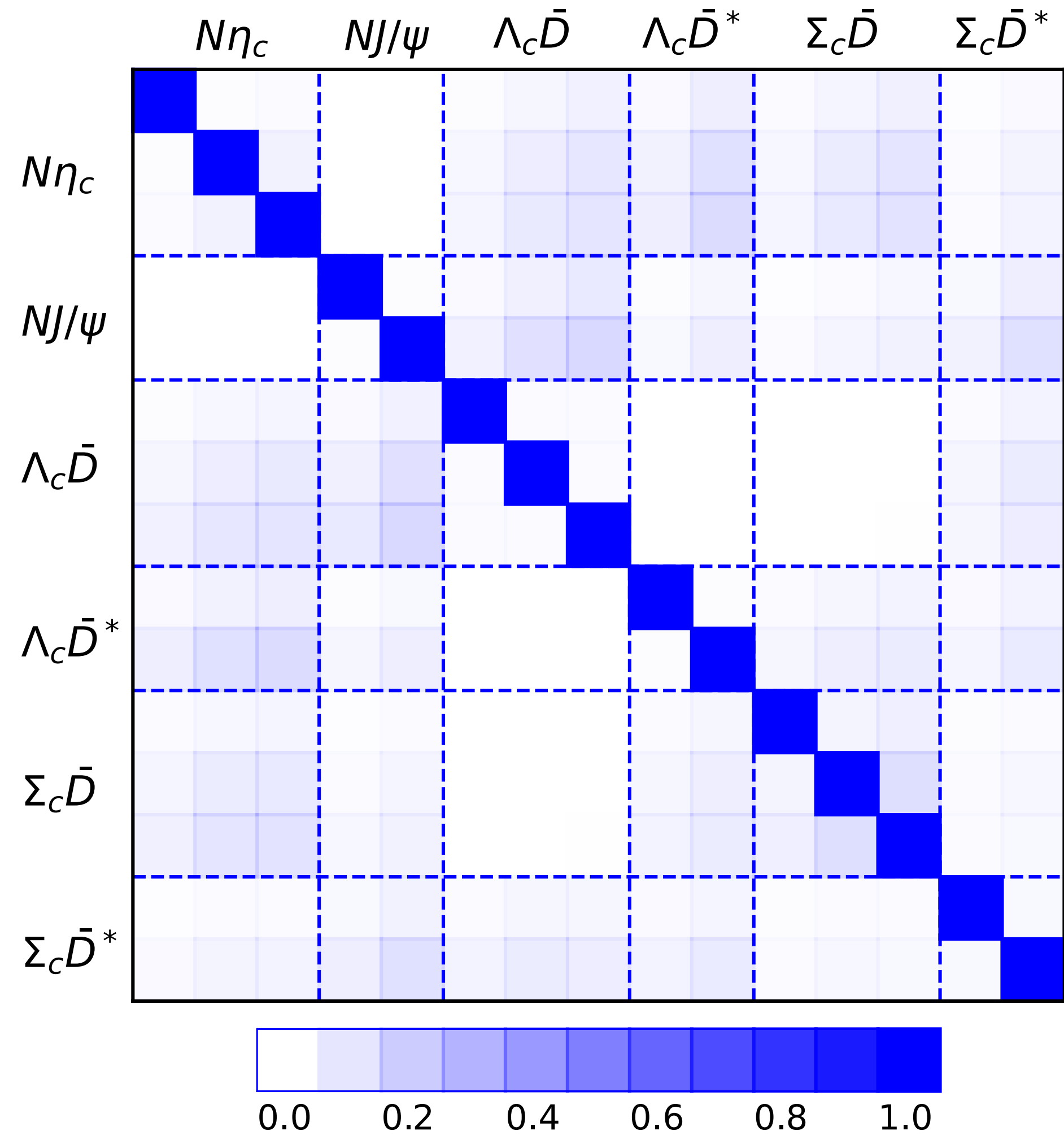
$$\mathcal{O}_{19,20,21} = \Sigma_c(\mathbf{p})\bar{D}(-\mathbf{p}) \quad (\mathbf{p}^2 = 0,1,2)$$

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Coupled channels

$L = 32$

$L = 48$





Coupled channels



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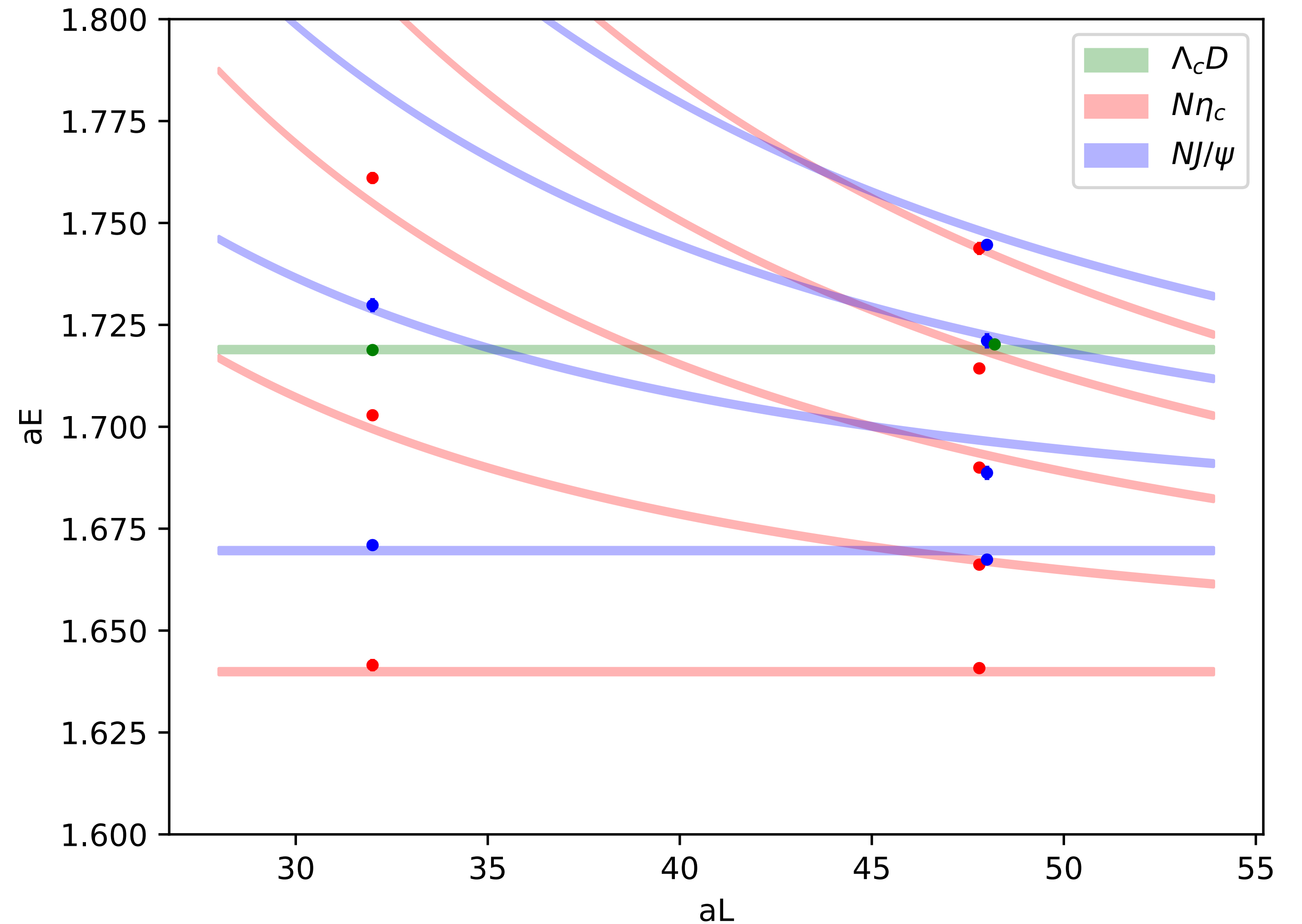
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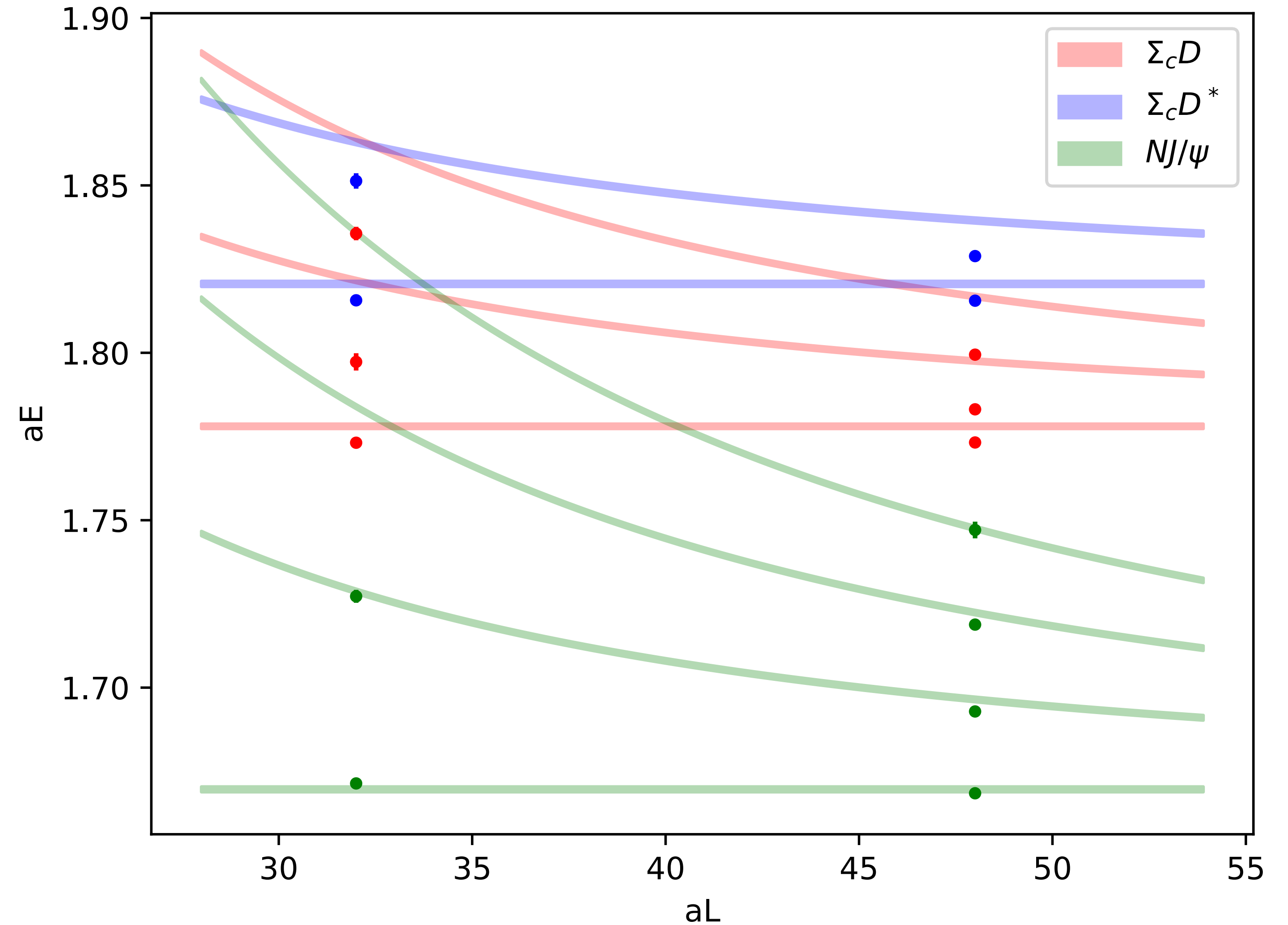
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$D\pi$ scattering



- A narrow exotic hadronic state with $J^{PC} = 0^{--}$ was predicted in
T. Ji, X.-K. Dong, F.-K. Guo and B.-S. Zou, PRL. 129, 102002 (2022)

$$\psi_0(4360) \longrightarrow D^* \bar{D}_1$$

- Scalar and axial vector charmed mesons:

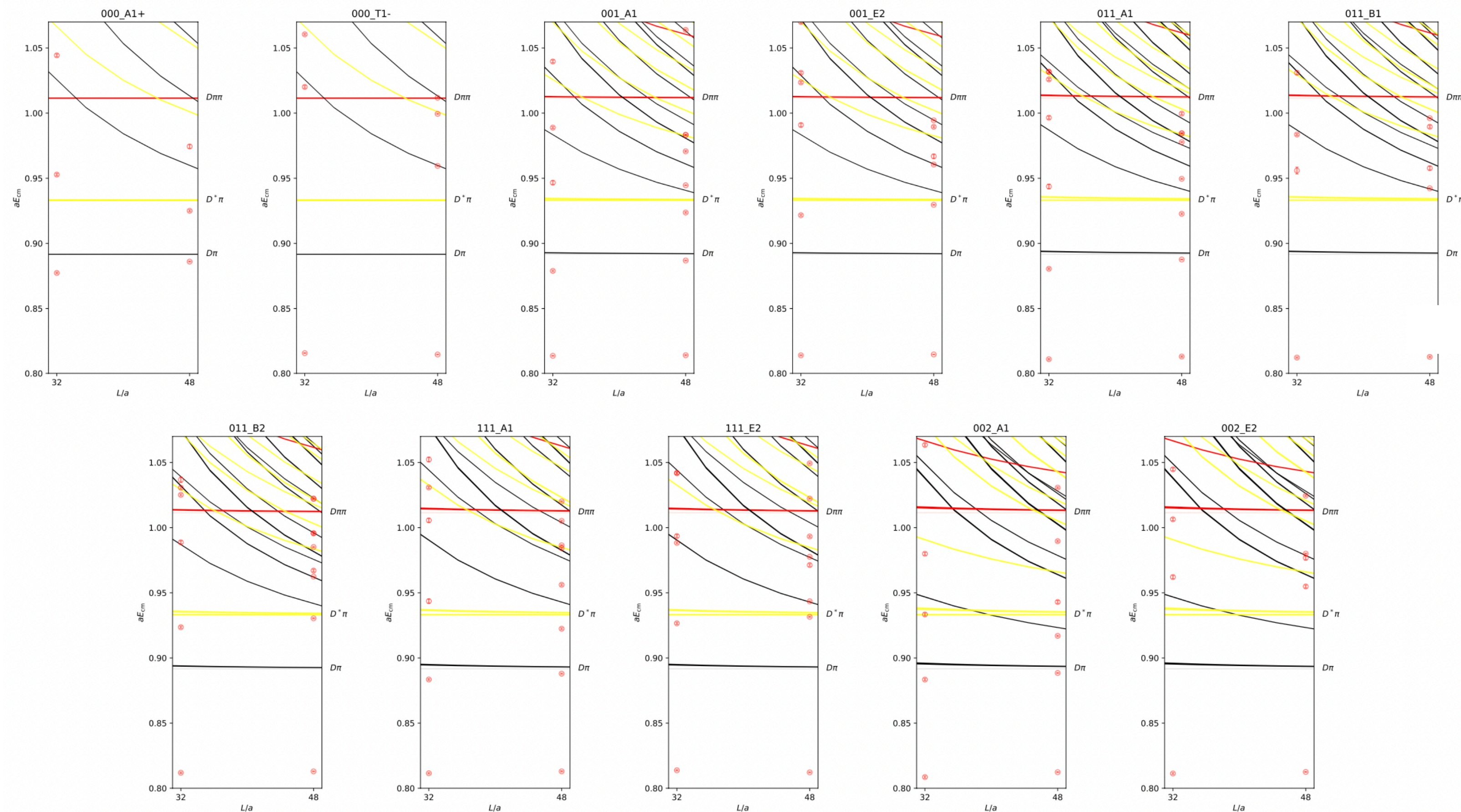
| 0^+ | 1^+ |
|--|------------------------------|
| $D_{s0}^* : 2317.8 \pm 0.5 \text{MeV}$ | $D_{s1}(2460), D_{s1}(2536)$ |
| $D_0^* : 2343 \pm 10 \text{MeV}$ | $D_1(2420), D_1(2430)$ |

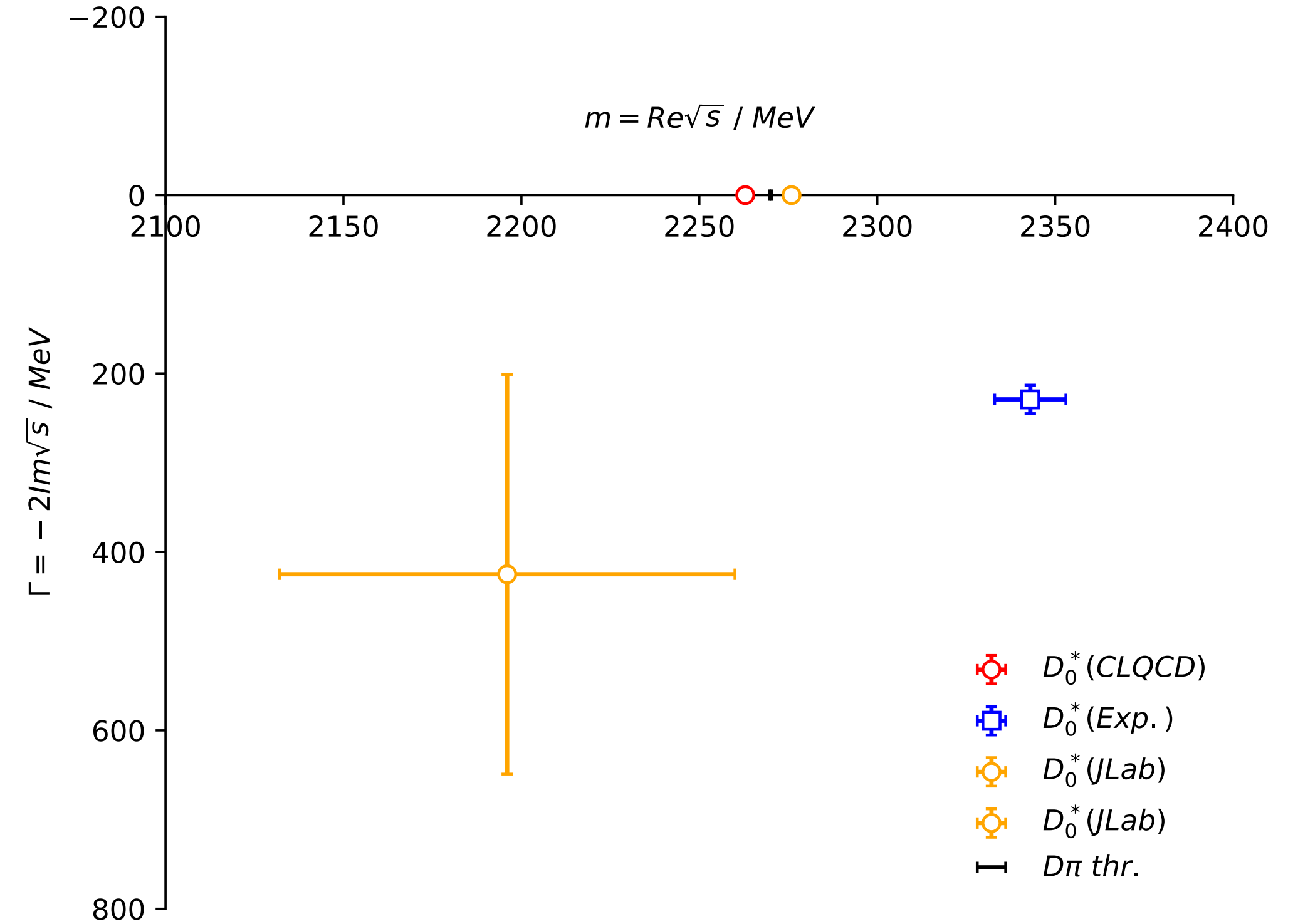
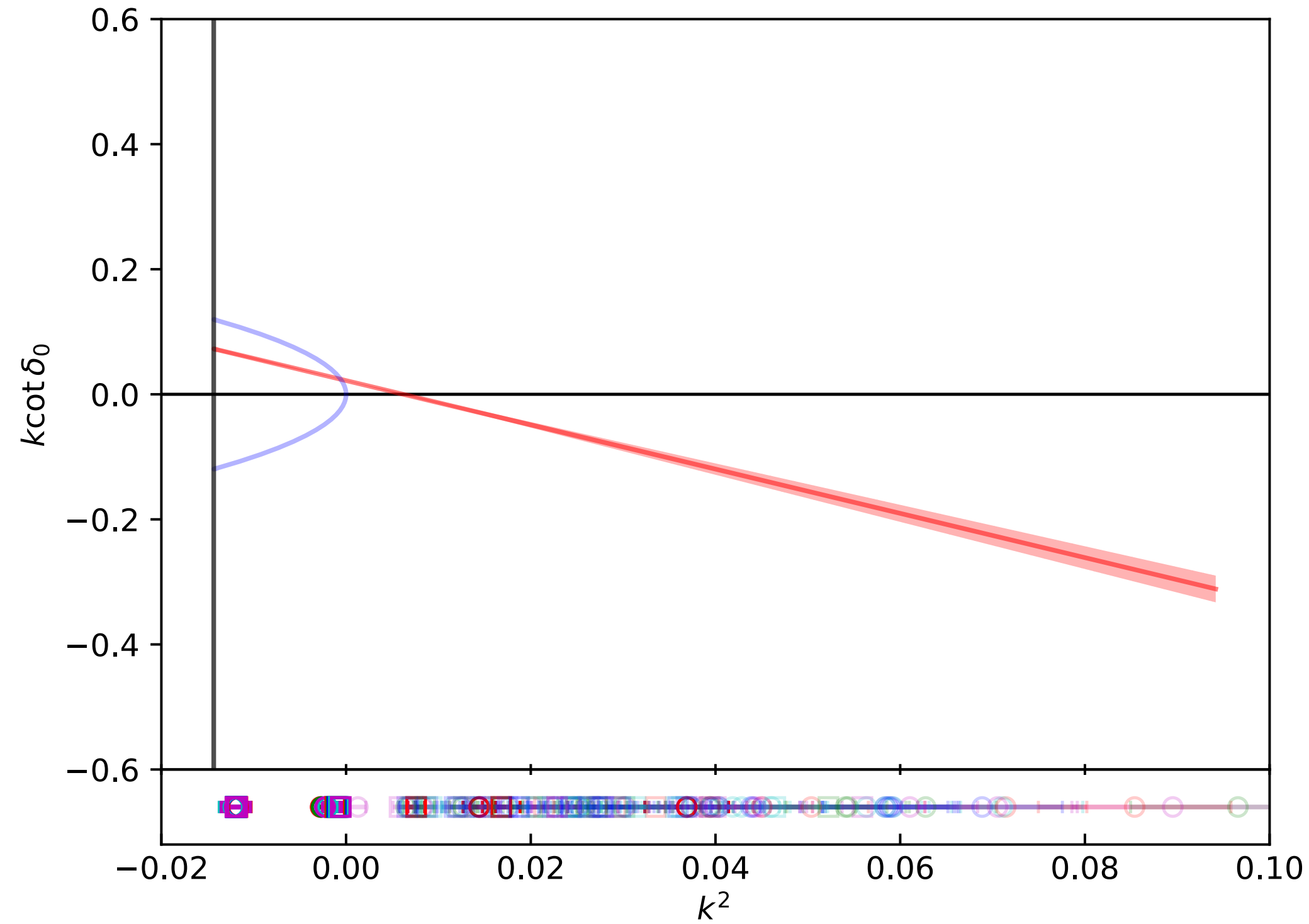
- $D^{(*)}\pi, D^{(*)}K$ scattering. Hao-bo Yan(PKU)

- Many interpolating operators are used, including both rest frame and moving frames

$$P = \frac{2\pi}{L} \{(0,0,0), (0,0,1), (0,1,1), (1,1,1), (0,0,2)\}.$$

- The finite volume spectra:







Summary



- ◆ Single channel analysis indicates bound states in $\Sigma_c \bar{D}$ and $\Sigma_c \bar{D}^*$ channel at $m_\pi \sim 300\text{MeV}$.
- ◆ Coupled channels:
 - Need robust determination of the spectrum with a complete set of interpolating operators.
 - Coupled channel scattering analysis.
- ◆ At $m_\pi \sim 300\text{MeV}$, a virtual bound state is found in $I = \frac{1}{2}$ S-wave $D\pi$ scattering.