

Observation of $\chi_{cJ} \rightarrow p\bar{p}\pi^+\pi^-\pi^0$, $p\bar{p}\omega$ and Evidence of $\eta_c(2S) \rightarrow p\bar{p}\pi^+\pi^-\pi^0$

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HFP Group Meeting

2023.06.20

outline

- Motivation
- Data sample
- Event selection
- Analysis
 - ▣ $\chi_{cJ}/\eta_c(2S) \rightarrow p\bar{p}\pi^+\pi^-\pi^0$
 - ▣ Investigate the intermediate process
- Summary & next to do

Motivation

- $R_1 = \frac{\Gamma_{h_c}^h}{\Gamma_{\eta_c(1S)}^h} = 0.010 \pm 0.001$ and 0.083 ± 0.018 based on pQCD and NRQCD calculation [1] .
- $R_2 = \frac{B(\eta_c(2S) \rightarrow h)}{B(\eta_c(1S) \rightarrow h)} = \frac{B(\eta_c(2S) \rightarrow \gamma\gamma)}{B(\eta_c(1S) \rightarrow \gamma\gamma)} \approx 0.52 \sim 1.56$ [2] .
- $h_c \rightarrow p\bar{p}\pi^+\pi^-\pi^0$ is observed with strong significance of 4.9σ based on $\sim 448\text{M}$ $\psi(3686)$ data sample at BESIII [3] .
- Search for $\eta_c(2S) \rightarrow p\bar{p}\pi^+\pi^-\pi^0$ is important to test the calculation above.
- Large blank room of χ_{cJ} decays need to be filled up [4] .
- $\sim 2.7\text{B}$ $\psi(3686)$ data sample at BESIII give us an excellent chance to investigate this topic via $\psi(3686) \rightarrow \gamma p\bar{p}\pi^+\pi^-\pi^0$ decay

[1] Phys. Rep.. 41, 1(1978); [2]Phys. Rev. Lett. 50, 569 (1983); [3] Journal of High Energy Physics. 05, 108 (2022)
[4] PDG2022, Prog. Theor. Exp. Phys. 2022, 083C01. (2022).

Data sample

Data set	Number of events	Boss version
09+12 $\psi(3686)$ data 2021 $\psi(3686)$ data	4.48×10^8 22.60×10^8	709
09+12 $\psi(3686)$ inclusive MC 2021 $\psi(3686)$ inclusive MC	5.06×10^8 $\sim 23 \times 10^8$	
PHSP MC	3 million (09+12+21)	

Event selection

Good charged track selection

- $|R_{xy}| < 1 \text{ cm}, |R_z| < 10 \text{ cm}$
- $|\cos\theta| < 0.93$
- $N = 4, N_m = N_p = 2, \Sigma Q = 0$

Particle identification

- $\text{Prob}(\pi) > \text{Prob}(p), \text{Prob}(\pi) > \text{Prob}(K), \text{Prob}(\pi) > 0.001$
- $\text{Prob}(p) > \text{Prob}(\pi), \text{Prob}(p) > \text{Prob}(K), \text{Prob}(p) > 0.001$
- $N_{\pi^+} = N_{\pi^-} = N_p = N_{\bar{p}} = 1$

Vertex Fit for $p\bar{p}\pi^+\pi^-$

Good photon selection

- $0 \leq TDC \leq 14$
- Barrel: $E > 0.025 \text{ GeV}, |\cos\theta| < 0.8$
- End cap: $E > 0.025 \text{ GeV}, 0.86 < |\cos\theta| < 0.92$
- $N_\gamma \geq 3$

1C for π^0 list

- $0.08 \text{ GeV}/c^2 < M(\pi^0) < 0.2 \text{ GeV}/c^2$
- $\chi^2_{1C}(\pi^0) < 200$
- $N_{\pi^0} \geq 1$

4C-kinematic fit

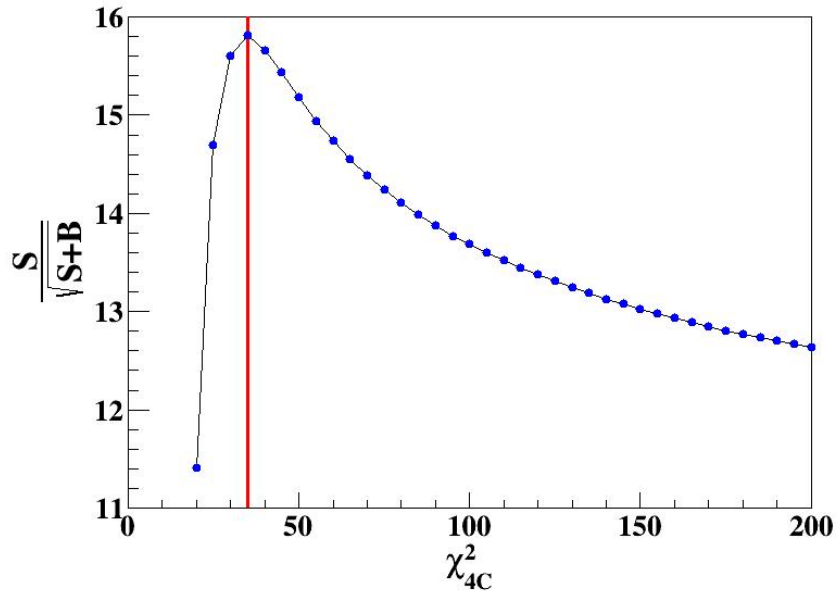
- choose the photons with least χ^2_{4C} .
- Additional 4C kinematic fit for $2\gamma p\bar{p}\pi^+\pi^-$, $3\gamma p\bar{p}\pi^+\pi^-$ and $4\gamma p\bar{p}\pi^+\pi^-$ final states.

3C-kinematic fit

- To separate the signal from background $\Psi(3686) \rightarrow p\bar{p}\pi^+\pi^-\gamma\gamma(\pi^0)$

Further Event selection

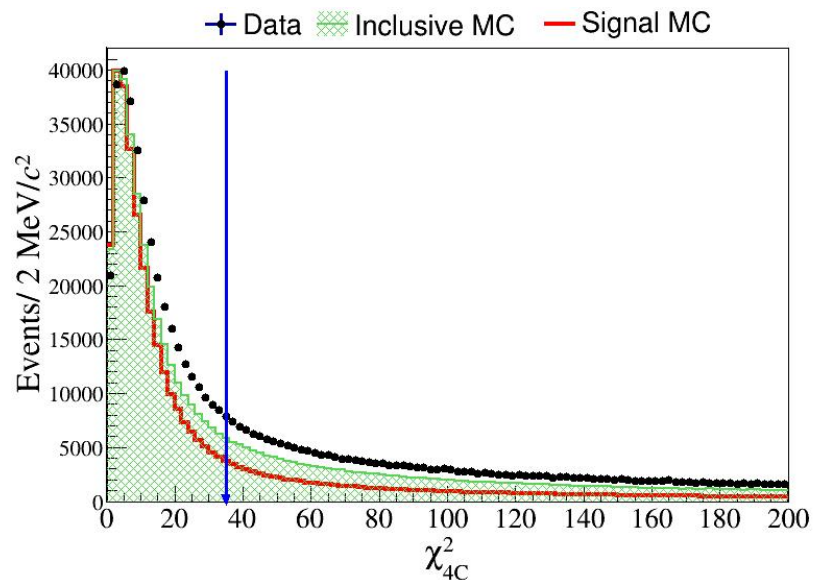
4C kinematic fit



➤ $\chi^2_{4C} < 35$

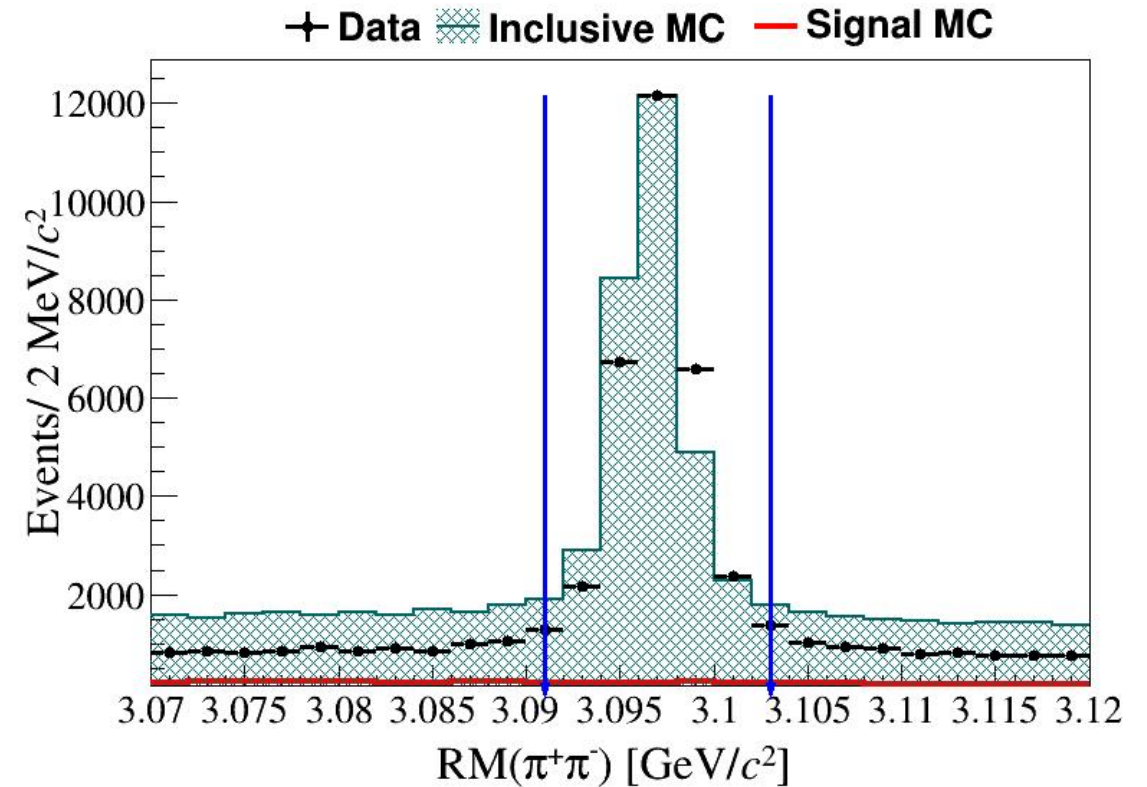
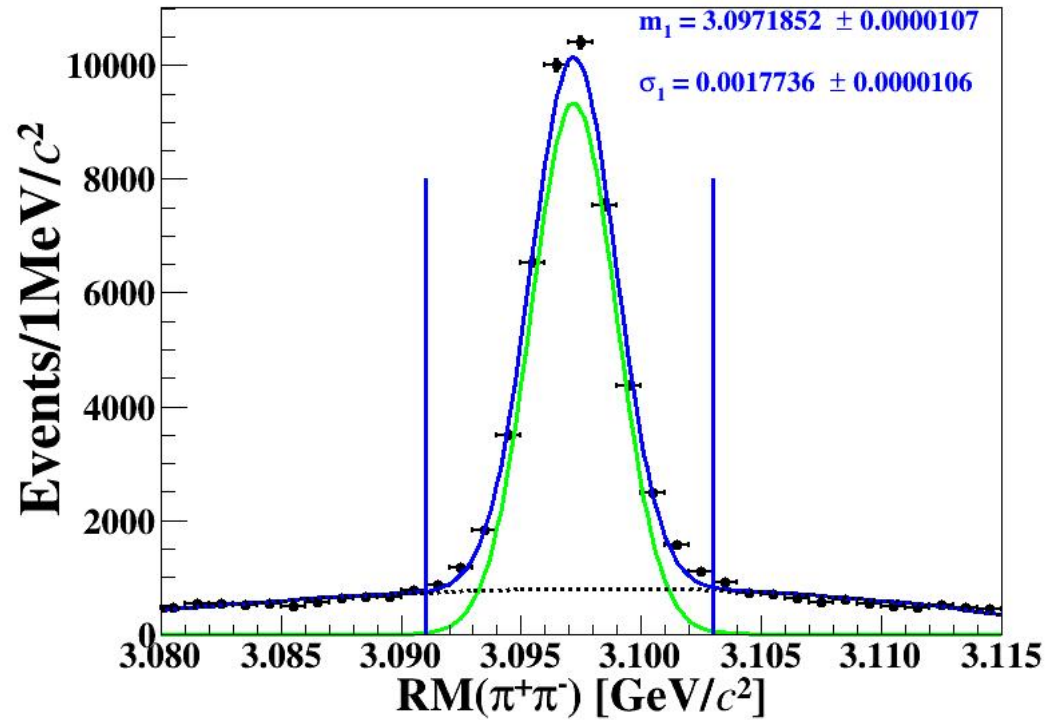
➤ Additional requirement:

- ✓ $\chi^2_{4C}(\mathbf{p\bar{p}} \gamma \gamma(\pi^0) \gamma) < \chi^2_{4C}(\mathbf{p\bar{p}} \gamma \gamma(\pi^0))$
- ✓ $\chi^2_{4C}(\mathbf{p\bar{p}} \gamma \gamma(\pi^0) \gamma) < \chi^2_{4C}(\mathbf{p\bar{p}} \gamma \gamma(\pi^0) 2\gamma)$
- ✓ $\chi^2_{4C}(\mathbf{3\gamma p\bar{p}}) < \chi^2_{4C}(\mathbf{2\gamma p\bar{p}})$
- ✓ $\chi^2_{4C}(\mathbf{3\gamma p\bar{p}}) < \chi^2_{4C}(\mathbf{4\gamma p\bar{p}})$



Background study

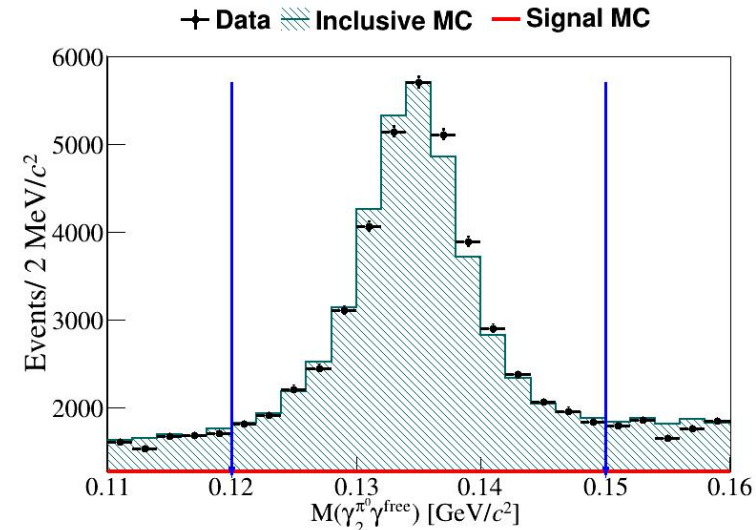
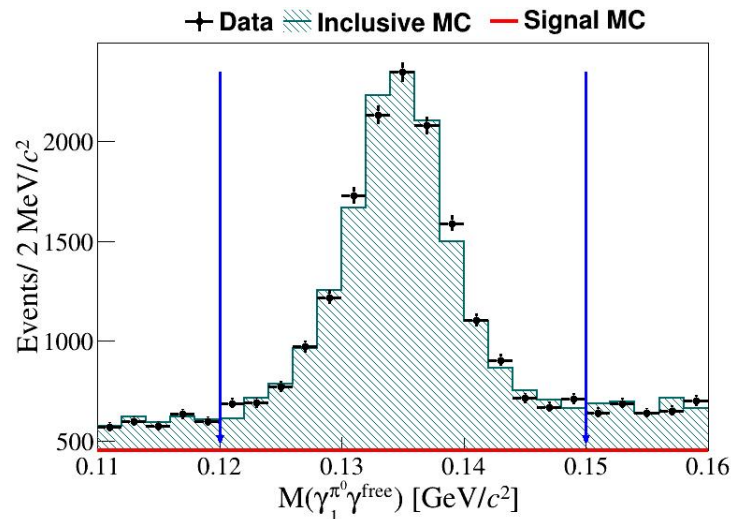
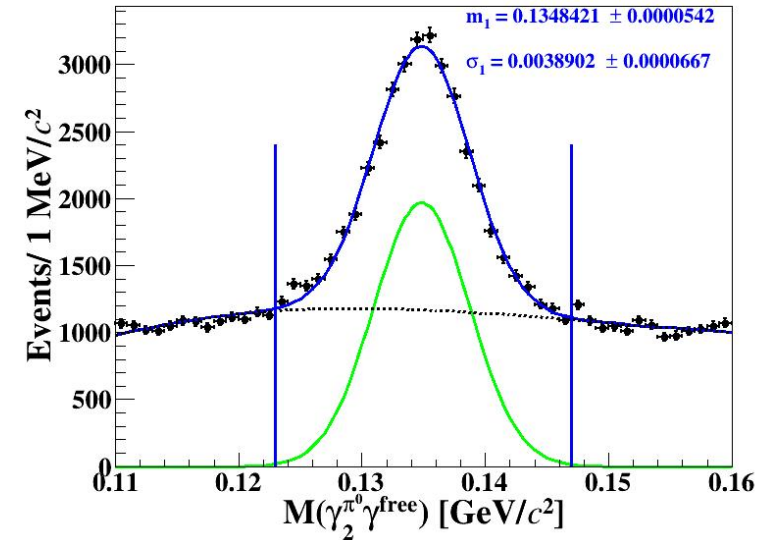
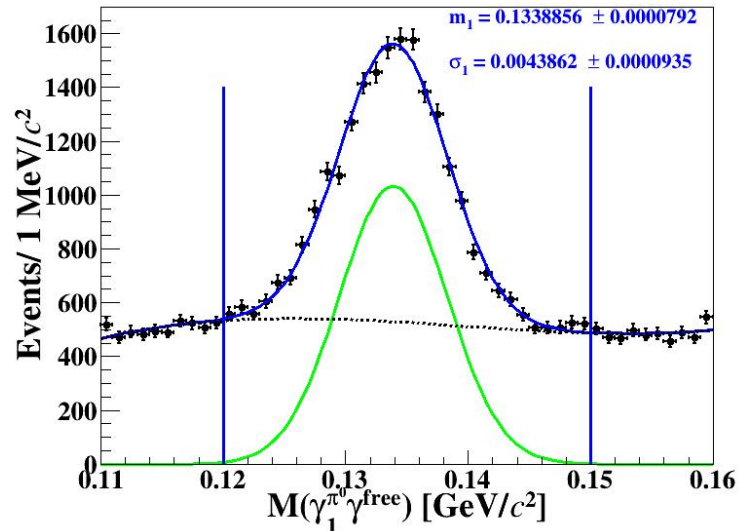
Veto on $\psi(3686) \rightarrow \pi^+ \pi^- J/\psi$



➤ $|\text{RM}(\pi^+\pi^-) - m(J/\psi)| > 6 \text{ MeV}/c^2$

Background study

Veto on $\pi^0 \rightarrow \gamma^{\pi^0} \gamma^{\text{free}}$ (Wrong Combination)

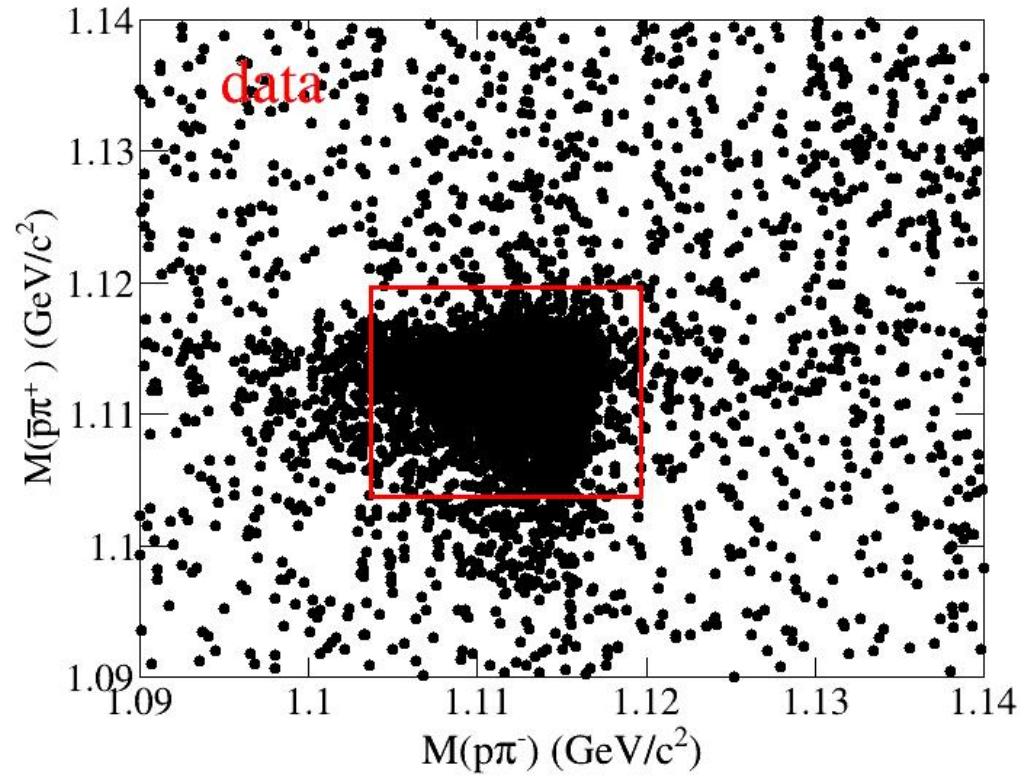


➤ $\left| M(\gamma_1^{\pi^0} \gamma^{\text{free}}) - m(\pi^0) \right| > 15 \text{ MeV}/c^2$

➤ $\left| M(\gamma_2^{\pi^0} \gamma^{\text{free}}) - m(\pi^0) \right| > 12 \text{ MeV}/c^2$

Background study

Veto on $\Lambda(\bar{\Lambda})$

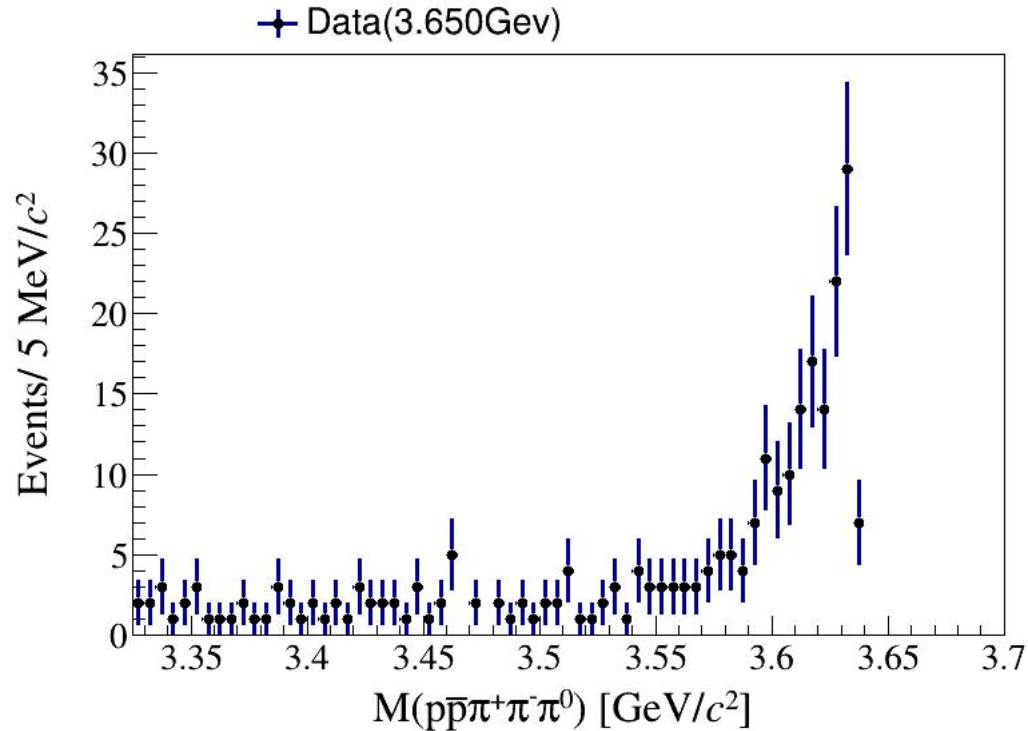


$M(p\pi^-)$ vs $M(\bar{p}\pi^+) \in (1.104, 1.120) \text{ GeV}/c^2$

Background study

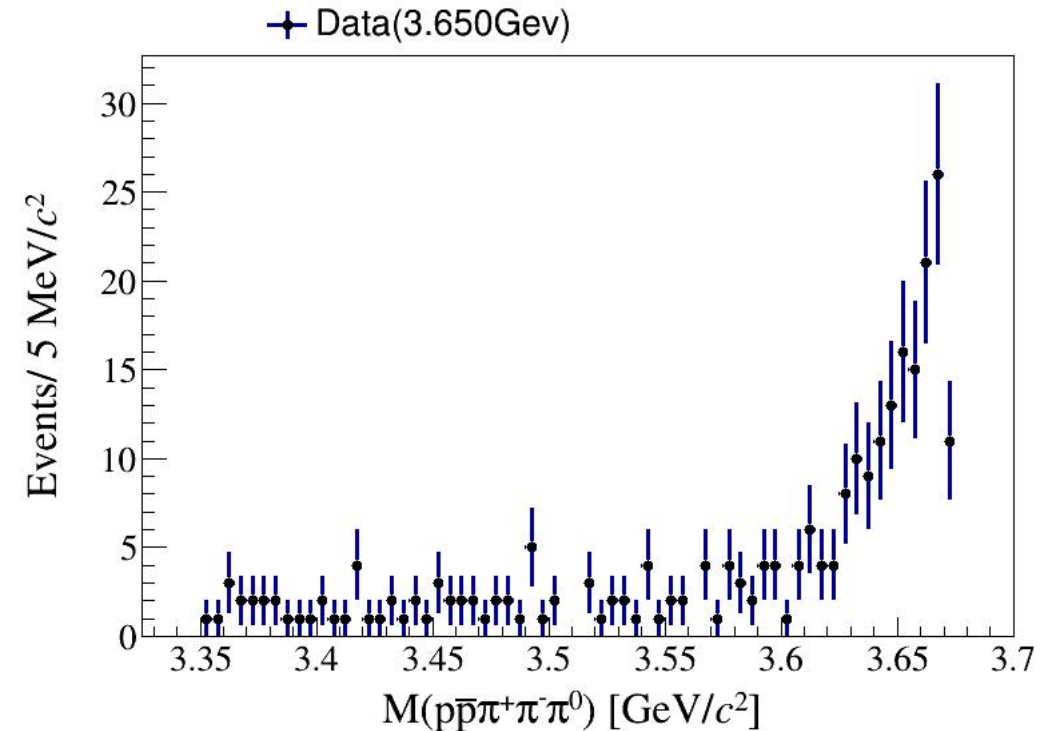
Background from continuum data

@3.650GeV(09+12+21): 434.95pb⁻¹ (online)



$$f_{\text{continuum}} = \frac{\mathcal{L}_{3.686}}{\mathcal{L}_{3.650}} \times \frac{\sigma_{3.686}}{\sigma_{3.650}} = \frac{3878.55/\text{pb}}{434.95/\text{pb}} \times \frac{3.650^2}{3.686^2}$$

$$N_{\text{continuum}} = N_{3.650} \times f_{\text{continuum}} \approx 2133.43$$



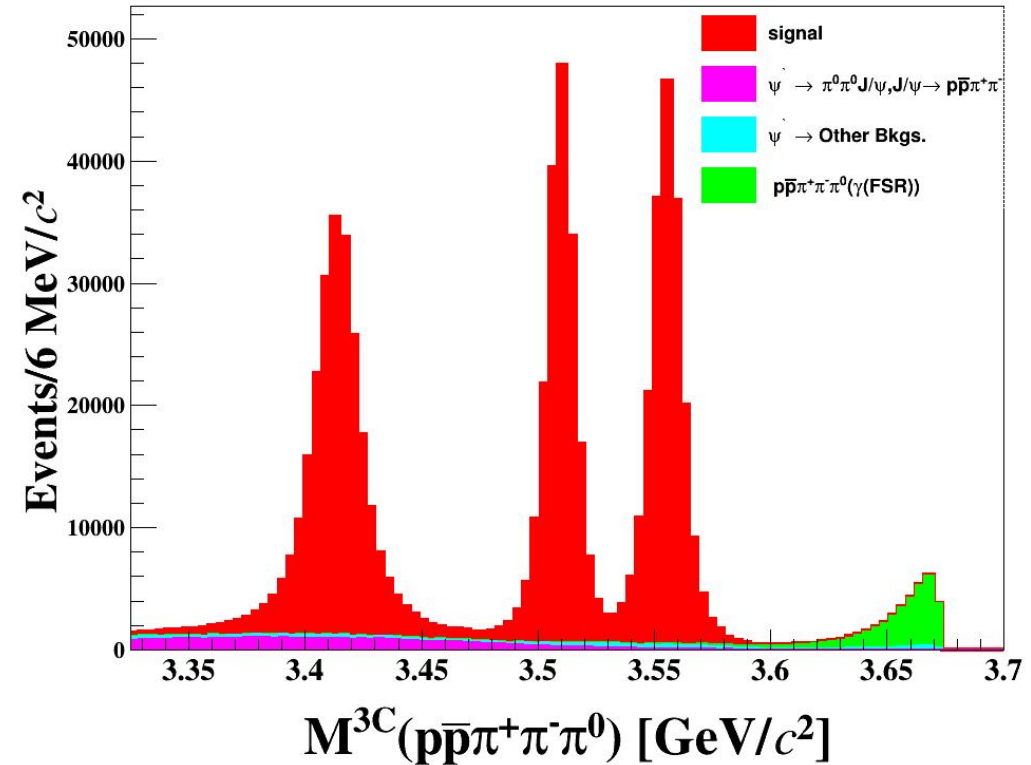
$$m_{\text{shift}} = a(m - m_0) + m_0$$

$$m_0: M_{p\bar{p}\pi^+\pi^-\pi^0}$$

$$a = \frac{(3.686 - m_0)}{(3.650 - m_0)}$$

Topology analysis

rowNo	decay tree	decay final state	iDcyTr	nEtr	nCEtr
1	$\psi' \rightarrow \chi_{c0}\gamma, \chi_{c0} \rightarrow \pi^0 \Delta^{++}, \Delta^{++} \rightarrow \pi^+ p, \Delta^{++} \rightarrow \pi^- \bar{p}$	$\pi^0 \pi^+ \pi^- p \bar{p}$	1	62625	62625
2	$\psi' \rightarrow \chi_{c1}\gamma, \chi_{c1} \rightarrow \pi^0 \Delta^{++}, \Delta^{++} \rightarrow \pi^+ p, \Delta^{++} \rightarrow \pi^- \bar{p}$	$\pi^0 \pi^+ \pi^- p \bar{p}$	9	43291	105816
3	$\psi' \rightarrow \chi_{c2}\gamma, \chi_{c2} \rightarrow \pi^0 \Delta^{++}, \Delta^{++} \rightarrow \pi^+ p, \Delta^{++} \rightarrow \pi^- \bar{p}$	$\pi^0 \pi^+ \pi^- p \bar{p}$	2	42935	148751
4	$\psi' \rightarrow \chi_{c0}\gamma, \chi_{c0} \rightarrow p \bar{p} h_1(1170), h_1(1170) \rightarrow \pi^+ \rho^-, \rho^- \rightarrow \pi^0 \pi^-$	$\pi^0 \pi^+ \pi^- p \bar{p}$	5	23554	172303
5	$\psi' \rightarrow \chi_{c0}\gamma, \chi_{c0} \rightarrow p \bar{p} h_1(1170), h_1(1170) \rightarrow \pi^- \rho^+, \rho^+ \rightarrow \pi^0 \pi^+$	$\pi^0 \pi^+ \pi^- p \bar{p}$	6	23446	195751
6	$\psi' \rightarrow \chi_{c0}\gamma, \chi_{c0} \rightarrow \pi^- \Delta^+, \Delta^+ \rightarrow \pi^0 p, \Delta^+ \rightarrow \pi^+ \bar{p}$	$\pi^0 \pi^+ \pi^- p \bar{p}$	4	23165	218916
7	$\psi' \rightarrow \chi_{c0}\gamma, \chi_{c0} \rightarrow \pi^- \Delta^0, \Delta^0 \rightarrow \pi^0 p, \Delta^0 \rightarrow \pi^- \bar{p}$	$\pi^0 \pi^+ \pi^- p \bar{p}$	42	22877	241793
8	$\psi' \rightarrow \chi_{c2}\gamma, \chi_{c2} \rightarrow p \bar{p} h_1(1170), h_1(1170) \rightarrow \pi^+ \rho^-, \rho^- \rightarrow \pi^0 \pi^-$	$\pi^0 \pi^+ \pi^- p \bar{p}$	19	22583	264376
9	$\psi' \rightarrow \chi_{c2}\gamma, \chi_{c2} \rightarrow p \bar{p} h_1(1170), h_1(1170) \rightarrow \pi^- \rho^+, \rho^+ \rightarrow \pi^0 \pi^+$	$\pi^0 \pi^+ \pi^- p \bar{p}$	15	22563	286939
10	$\psi' \rightarrow \chi_{c0}\gamma, \chi_{c0} \rightarrow p \bar{p} h_1(1170), h_1(1170) \rightarrow \pi^0 \rho^0, \rho^0 \rightarrow \pi^+ \pi^-$	$\pi^0 \pi^+ \pi^- p \bar{p}$	12	22117	309056
11	$\psi' \rightarrow \chi_{c1}\gamma, \chi_{c1} \rightarrow p \bar{p} h_1(1170), h_1(1170) \rightarrow \pi^- \rho^+, \rho^+ \rightarrow \pi^0 \pi^+$	$\pi^0 \pi^+ \pi^- p \bar{p}$	56	21095	330151
12	$\psi' \rightarrow \chi_{c2}\gamma, \chi_{c2} \rightarrow p \bar{p} h_1(1170), h_1(1170) \rightarrow \pi^0 \rho^0, \rho^0 \rightarrow \pi^+ \pi^-$	$\pi^0 \pi^+ \pi^- p \bar{p}$	43	20728	350879
13	$\psi' \rightarrow \chi_{c1}\gamma, \chi_{c1} \rightarrow p \bar{p} h_1(1170), h_1(1170) \rightarrow \pi^+ \rho^-, \rho^- \rightarrow \pi^0 \pi^-$	$\pi^0 \pi^+ \pi^- p \bar{p}$	7	20653	371532
14	$\psi' \rightarrow \pi^0 \pi^0 J/\psi, J/\psi \rightarrow \pi^+ \pi^- p \bar{p}$	$\pi^0 \pi^0 \pi^+ \pi^- p \bar{p}$	6	19603	391135
15	$\psi' \rightarrow \chi_{c1}\gamma, \chi_{c1} \rightarrow p \bar{p} h_1(1170), h_1(1170) \rightarrow \pi^0 \rho^0, \rho^0 \rightarrow \pi^+ \pi^-$	$\pi^0 \pi^+ \pi^- p \bar{p}$	11	19522	410657
16	$\psi' \rightarrow \chi_{c1}\gamma, \chi_{c1} \rightarrow \pi^- \Delta^+, \Delta^+ \rightarrow \pi^0 p, \Delta^+ \rightarrow \pi^+ \bar{p}$	$\pi^0 \pi^+ \pi^- p \bar{p}$	3	15248	425903
17	$\psi' \rightarrow \chi_{c1}\gamma, \chi_{c1} \rightarrow \pi^- \Delta^0, \Delta^0 \rightarrow \pi^0 p, \Delta^0 \rightarrow \pi^- \bar{p}$	$\pi^0 \pi^+ \pi^- p \bar{p}$	17	14970	440875
18	$\psi' \rightarrow \chi_{c2}\gamma, \chi_{c2} \rightarrow \pi^- \Delta^+, \Delta^+ \rightarrow \pi^0 p, \Delta^+ \rightarrow \pi^+ \bar{p}$	$\pi^0 \pi^+ \pi^- p \bar{p}$	20	14138	455013
19	$\psi' \rightarrow \chi_{c2}\gamma, \chi_{c2} \rightarrow \pi^- \Delta^0, \Delta^0 \rightarrow \pi^0 p, \Delta^0 \rightarrow \pi^- \bar{p}$	$\pi^0 \pi^+ \pi^- p \bar{p}$	39	14015	469028
20	$\psi' \rightarrow \chi_{c0}\gamma, \chi_{c0} \rightarrow \omega p \bar{p}, \omega \rightarrow \pi^0 \pi^+ \pi^-$	$\pi^0 \pi^+ \pi^- p \bar{p}$	25	11879	480907
21	$\psi' \rightarrow \pi^0 \pi^+ \pi^- p \bar{p}$	$\pi^0 \pi^+ \pi^- p \bar{p}$	26	11503	492410
22	$\psi' \rightarrow \chi_{c2}\gamma, \chi_{c2} \rightarrow \omega p \bar{p}, \omega \rightarrow \pi^0 \pi^+ \pi^-$	$\pi^0 \pi^+ \pi^- p \bar{p}$	21	11399	503809
23	$\psi' \rightarrow \pi^0 \Delta^{++} \Delta^{++}, \Delta^{++} \rightarrow \pi^+ p, \Delta^{++} \rightarrow \pi^- \bar{p}$	$\pi^0 \pi^+ \pi^- p \bar{p}$	30	9398	513207
24	$\psi' \rightarrow \pi^0 \pi^0 J/\psi, J/\psi \rightarrow \pi^+ p \Delta^{++}, \Delta^{++} \rightarrow \pi^- \bar{p}$	$\pi^0 \pi^0 \pi^+ \pi^- p \bar{p}$	42	8982	522189
25	$\psi' \rightarrow \chi_{c0}\gamma, \chi_{c0} \rightarrow \pi^- \Delta^+, \Delta^+ \rightarrow \pi^0 p, \Delta^+ \rightarrow \pi^+ \bar{p}$	$\pi^0 \pi^+ \pi^- p \bar{p}$	55	6384	528573
26	$\psi' \rightarrow \chi_{c0}\gamma, \chi_{c0} \rightarrow \pi^- \Delta^0, \Delta^0 \rightarrow \pi^0 p, \Delta^0 \rightarrow \pi^- \bar{p}$	$\pi^0 \pi^+ \pi^- p \bar{p}$	45	6333	534906
27	$\psi' \rightarrow \chi_{c0}\gamma, \chi_{c0} \rightarrow \rho^- \bar{p} \Delta^{++}, \rho^- \rightarrow \pi^0 \pi^-, \Delta^{++} \rightarrow \pi^+ p$	$\pi^0 \pi^+ \pi^- p \bar{p}$	18	5917	540823
28	$\psi' \rightarrow \chi_{c0}\gamma, \chi_{c0} \rightarrow \rho^+ p \Delta^{--}, \rho^+ \rightarrow \pi^0 \pi^+, \Delta^{--} \rightarrow \pi^- \bar{p}$	$\pi^0 \pi^+ \pi^- p \bar{p}$	44	5831	546653
29	$\psi' \rightarrow \chi_{c2}\gamma, \chi_{c2} \rightarrow \pi^0 \pi^- p \Delta^{++}, \Delta^{++} \rightarrow \pi^+ p$	$\pi^0 \pi^+ \pi^- p \bar{p}$	22	5510	552166
30	$\psi' \rightarrow \pi^0 \pi^0 J/\psi, J/\psi \rightarrow \pi^- \bar{p} \Delta^{++}, \Delta^{++} \rightarrow \pi^+ p$	$\pi^0 \pi^0 \pi^+ \pi^- p \bar{p}$	76	5360	557524
31	$\psi' \rightarrow \chi_{c2}\gamma, \chi_{c2} \rightarrow \pi^0 \pi^+ p \Delta^{--}, \Delta^{--} \rightarrow \pi^- \bar{p}$	$\pi^0 \pi^+ \pi^- p \bar{p}$	44	5184	562708
32	$\psi' \rightarrow \chi_{c1}\gamma, \chi_{c1} \rightarrow \omega p \bar{p}, \omega \rightarrow \pi^0 \pi^+ \pi^-$	$\pi^0 \pi^+ \pi^- p \bar{p}$	47	4878	567586
33	$\psi' \rightarrow \chi_{c0}\gamma, \chi_{c0} \rightarrow \pi^0 \pi^- p \Delta^{++}, \Delta^{++} \rightarrow \pi^+ p$	$\pi^0 \pi^+ \pi^- p \bar{p}$	61	4767	572353



signal

$\psi' \rightarrow \pi^0 \pi^0 J/\psi$

$\psi' \rightarrow p\bar{p}\pi^+\pi^-\pi^0(\gamma_{FSR})$

FSR correction

Background of $\psi(2S) \rightarrow p\bar{p}\pi^+\pi^-\pi^0(\gamma_{FSR})$:

FSR correction: we used the control sample $\psi(3686) \rightarrow \gamma\chi_{c0}, \chi_{c0} \rightarrow p\bar{p}\pi^+\pi^-(\gamma_{FSR})$ to estimate the difference of FSR factor between MC and data.

Good charged track selection

- $|R_{xy}| < 1 \text{ cm}, |R_z| < 10 \text{ cm}$
- $|\cos\theta| < 0.93$
- $N = 4, N_m = N_p = 2, \Sigma Q = 0$

Particle identification

- $\text{Prob}(\pi) > \text{Prob}(p), \text{Prob}(\pi) > \text{Prob}(K), \text{Prob}(\pi) > 0.001$
- $\text{Prob}(p) > \text{Prob}(\pi), \text{Prob}(p) > \text{Prob}(K), \text{Prob}(p) > 0.001$
- $N_{\pi^+} = N_{\pi^-} = N_p = N_{\bar{p}} = 1$

Vertex Fit for $p\bar{p}\pi^+\pi^-$

Good photon selection

- $0 \leq TDC \leq 14$
- Barrel: $E > 0.025 \text{ GeV}, |\cos\theta| < 0.8$
- End cap : $E > 0.025 \text{ GeV}, 0.86 < |\cos\theta| < 0.92$
- $N_\gamma \geq 2$

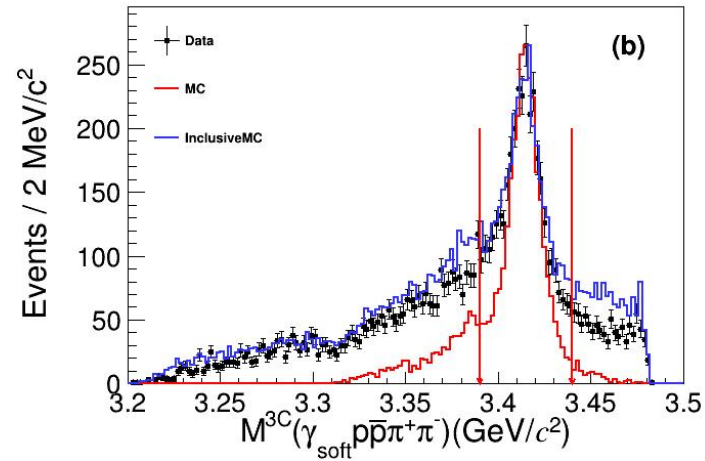
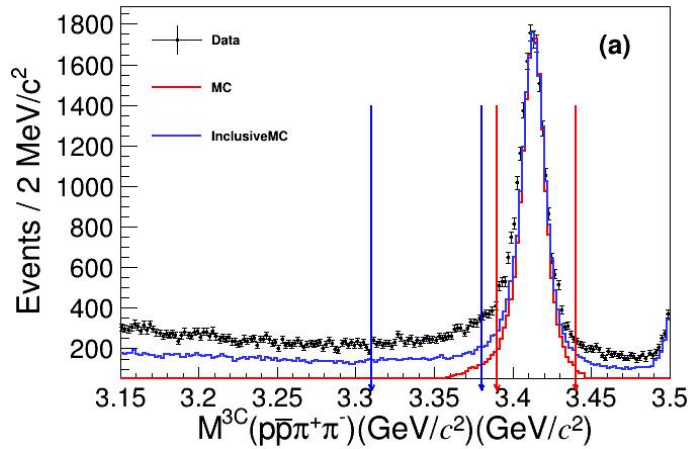
4C-kinematic fit

- choose the photons with least χ_{4C}^2 .

3C-kinematic fit

FSR correction

The FSR correction factor is defined as: $f = \frac{R_{FSR}^{data}}{R_{FSR}^{MC}}$ $R_{FSR} = \frac{N_{FSR}}{N_{nonFSR}}$

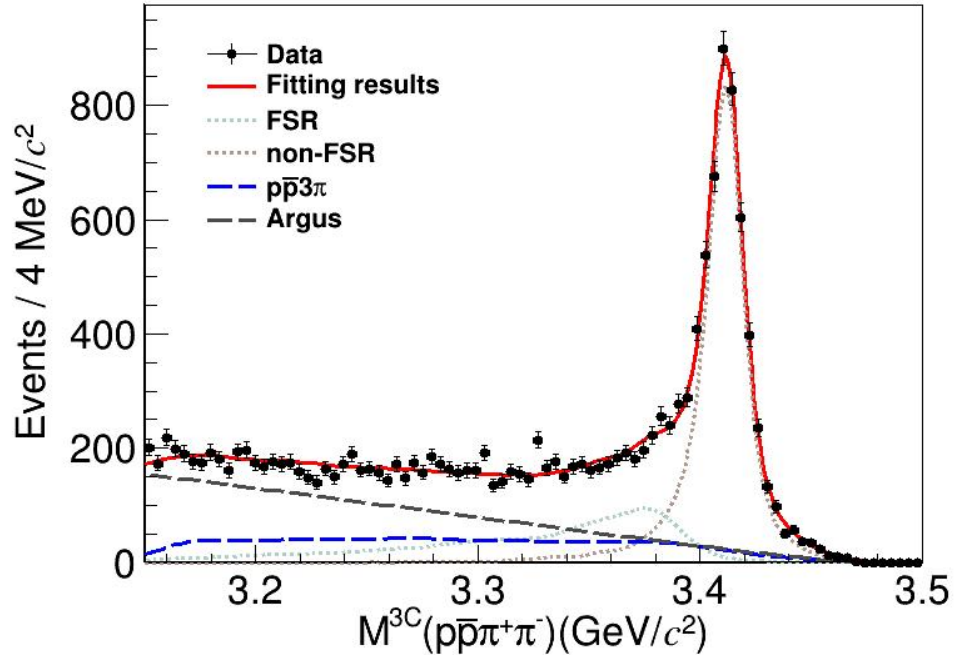


no FSR region : $M^{3C}(p\bar{p}\pi^+\pi^-) : 3.39\sim 3.44\text{GeV}/c^2$

FSR tail region : $M^{3C}(p\bar{p}\pi^+\pi^-) : 3.30\sim 3.38\text{GeV}/c^2$ && $M^{3C}((p\bar{p}\pi^+\pi^-)\gamma_{\text{soft}}) : 3.39\sim 3.44\text{GeV}/c^2$

FSR correction

We get the FSR ration of data sample by fitting.



- $R_{FSR}^{data} = \frac{N_{data}^{FSR}}{N_{data}^{nonFSR}} = 0.429 \pm 0.029$

- $R_{FSR}^{MC} = \frac{N_{MC}^{FSR}}{N_{MC}^{nonFSR}} = 0.253$

- $f = \frac{R_{FSR}^{data}}{R_{FSR}^{MC}} = 1.69 \pm 0.12$

We also listed the results for four similar studied channels. One finds that they are consistent with each other.

	Lina	Yaqian	GuangRui	Suxian
Decay channel	$\eta_c(2S) \rightarrow K^+ K^- \eta$	$\eta_c(2S) \rightarrow VV$	$\eta_c(2S) \rightarrow \pi^+ \pi^- \eta$	$\eta_c(2S) \rightarrow 3(\pi^+ \pi^-)$
f	1.61 ± 0.09	1.70 ± 0.10	1.62 ± 0.07	1.62 ± 0.13

Fit method for $M(p\bar{p}\pi^+\pi^-\pi^0)$

➤ $\eta_c(2S)$ Fit line shape : $(BW(m; M, \Gamma) \times E_\gamma^3 \times f_d(E_\gamma) \otimes DG(m_1, \sigma_1, m_2, \sigma_2, f) \times \epsilon(\mathbf{m})) \otimes G(\delta m_3, \delta \sigma_3)$

$\mathbf{m}$: $M_{p\bar{p}\pi^+\pi^-\pi^0}^{3C}$

$BW(m; M, \Gamma)$: Breit-Wigner function $\frac{1}{2\pi} \cdot \frac{\Gamma}{(m - m_0) + \Gamma^2/4}$

E_γ : transition photon energy in rest frame $\frac{M_{\psi(2S)}^2 - m^2}{2M_{\psi(2S)}}$

$\epsilon(\mathbf{m})$: argus function (fix)

$DG(m_1, \sigma_1, m_2, \sigma_2, f)$: Double Gaussian function, detector resolution (fixed)

$G(\delta m_3, \delta \sigma_3)$: Gaussian function, difference between data and MC

➤ χ_{cJ} line shape: MC shape $\otimes$ Gaussian function

➤ Background components : ① $\psi(2S) \rightarrow p\bar{p}\pi^+\pi^-\pi^0(\gamma_{FSR})$: Exclusive MC shape

② other background : Argus function

$\eta_c(2S)$ Efficiency curve: $\epsilon(m)$

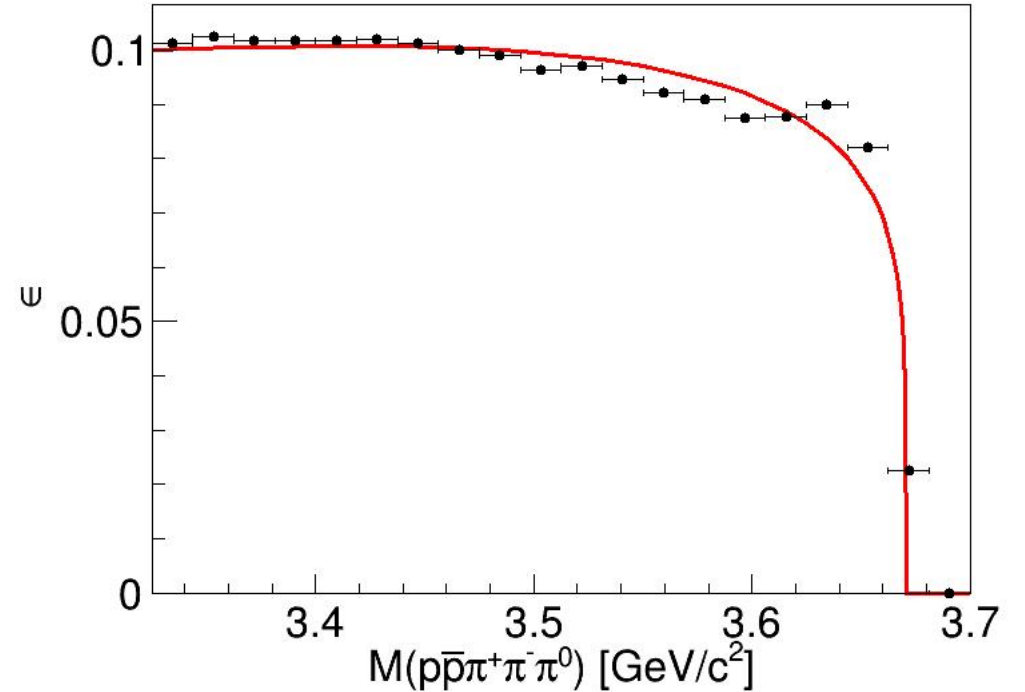
LSFLAT eta_c(2S)

ChangeMassMin eta_c(2S) 3.30

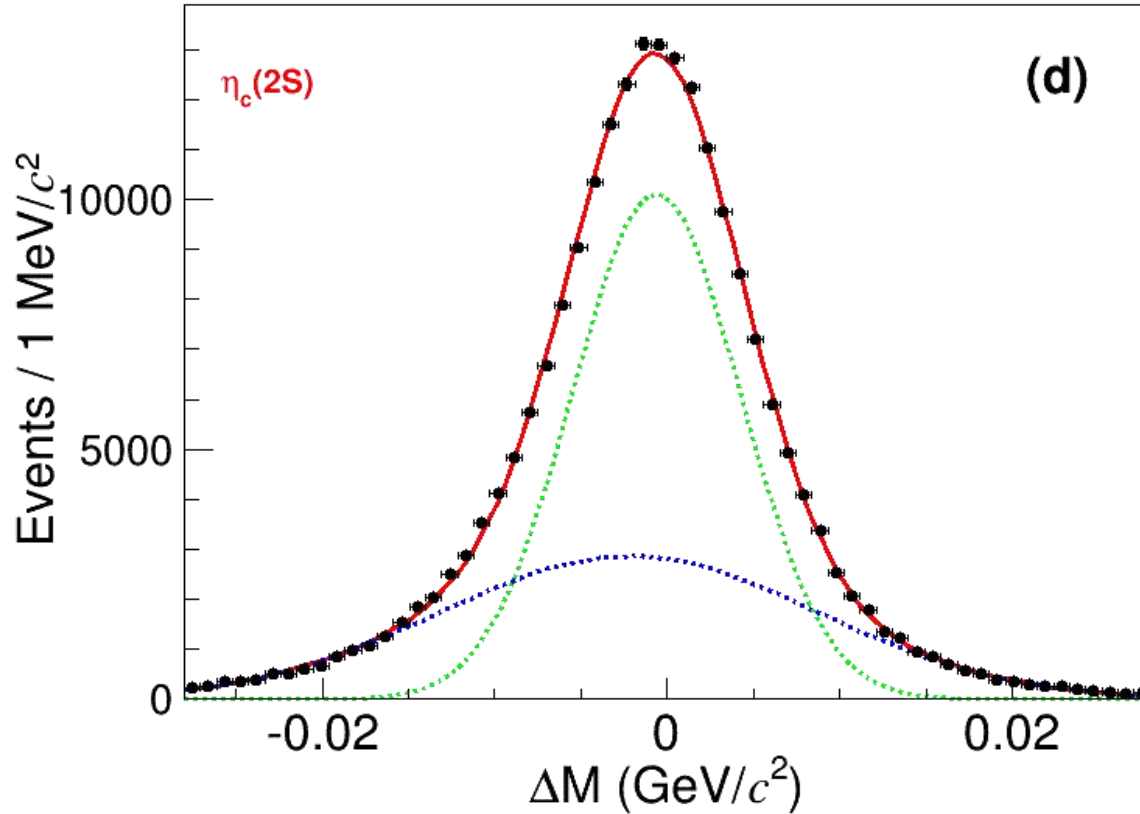
ChangeMassMax eta_c(2S) 3.75

argus function:

$$m(1 - (m/3.67)^2)^{0.169} \times e^{-0.6787(1-(m/3.67)^2)}$$



$\eta_c(2S)$ Mass resolution from signal MC

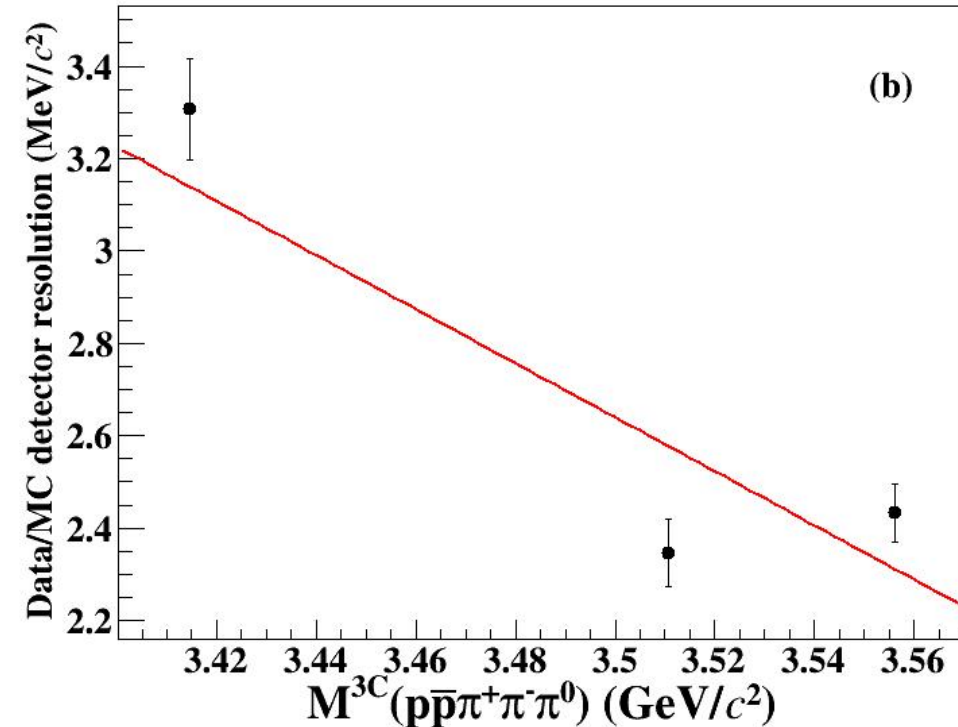
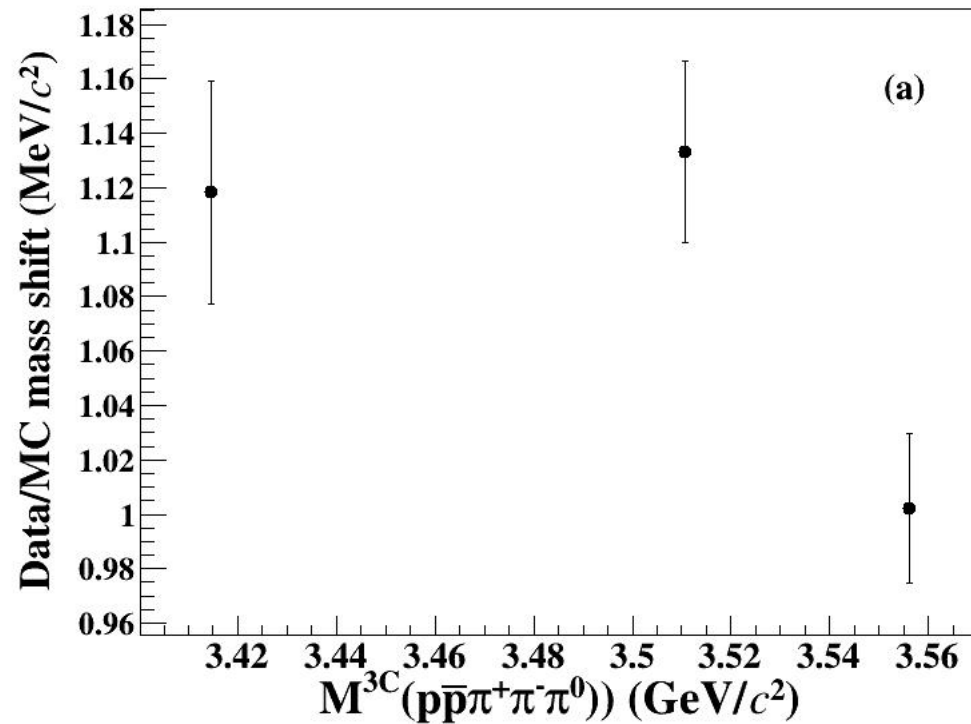


To obtain **Double Gauss**(m, σ) parameters, we fitted the mass difference ($\Delta M = M_{p\bar{p}\pi^+\pi^-\pi^0}^{3C} - M_{p\bar{p}\pi^+\pi^-\pi^0}^{truth}$) distribution from signal MC sample .

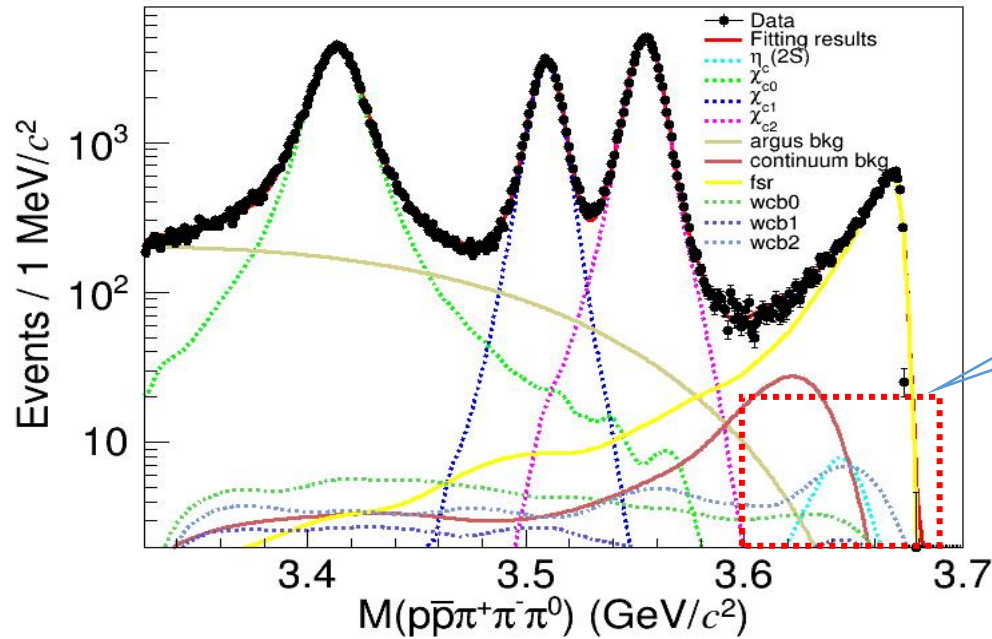
$\eta_c(2S)$ Detector resolution difference between data and MC

Gauss (m, σ)

	m	σ
χ_{c0}	$1.118 \pm 0.041 \text{ MeV}/c^2$	$3.307 \pm 0.109 \text{ MeV}$
χ_{c1}	$2.345 \pm 0.033 \text{ MeV}/c^2$	$2.347 \pm 0.072 \text{ MeV}$
χ_{c2}	$1.000 \pm 0.027 \text{ MeV}/c^2$	$2.433 \pm 0.064 \text{ MeV}$



Preliminary result

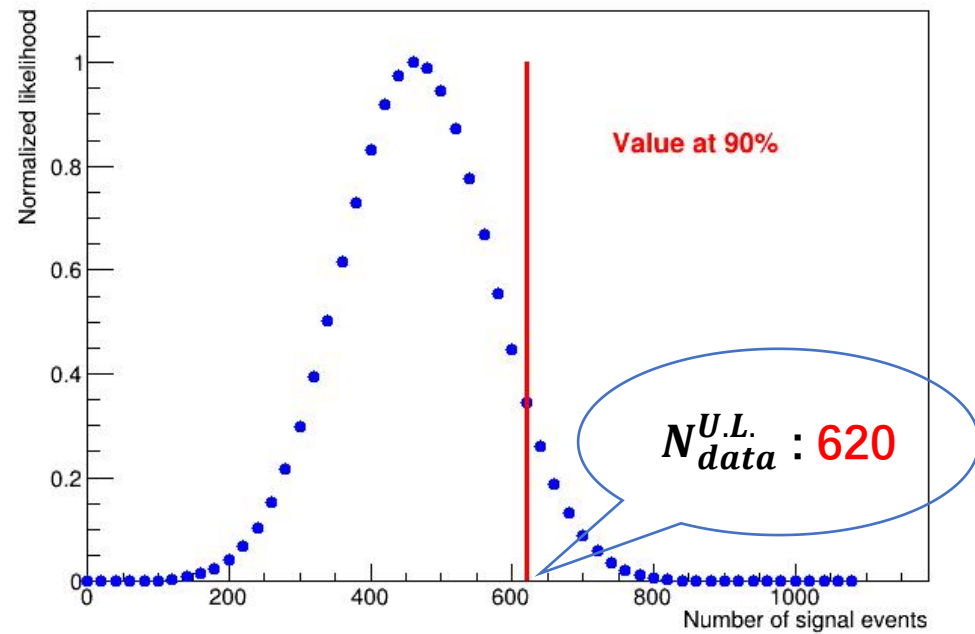


Statistical Significance: 3.16σ

$$B(\chi_{cJ}/\eta_c(2S) \rightarrow p\bar{p}\pi^+\pi^-\pi^0) = \frac{N_{\chi_{cJ}/\eta_c(2S)}^{obs}}{N_{\psi(2S)}^{tot} \cdot B(\psi(2S) \rightarrow \gamma\chi_{cJ}/\eta_c(2S)) \cdot B(\pi^0 \rightarrow \gamma\gamma) \cdot \epsilon}$$

channel	N^{obs}	$\epsilon(\%)$	Branching fraction
$\chi_{c0} \rightarrow p\bar{p}\pi^+\pi^-\pi^0$	$(1.180 \pm 0.004) \times 10^5$	9.23	$(4.880 \pm 0.017) \times 10^{-3}$
$\chi_{c1} \rightarrow p\bar{p}\pi^+\pi^-\pi^0$	$(5.346 \pm 0.026) \times 10^4$	9.70	$(2.112 \pm 0.011) \times 10^{-3}$
$\chi_{c2} \rightarrow p\bar{p}\pi^+\pi^-\pi^0$	$(8.447 \pm 0.032) \times 10^4$	8.79	$(3.772 \pm 0.014) \times 10^{-3}$
$\eta_c(2S) \rightarrow p\bar{p}\pi^+\pi^-\pi^0$	$(4.65 \pm 1.42) \times 10^2$	7.05	$(3.521 \pm 1.081) \times 10^{-3}$

Preliminary result



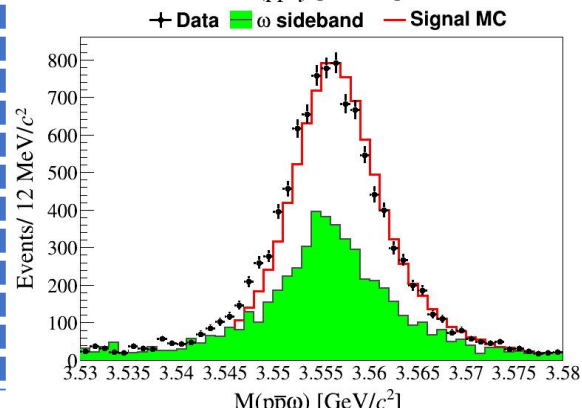
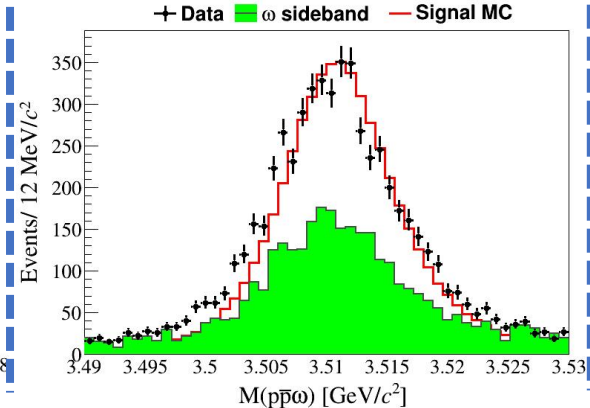
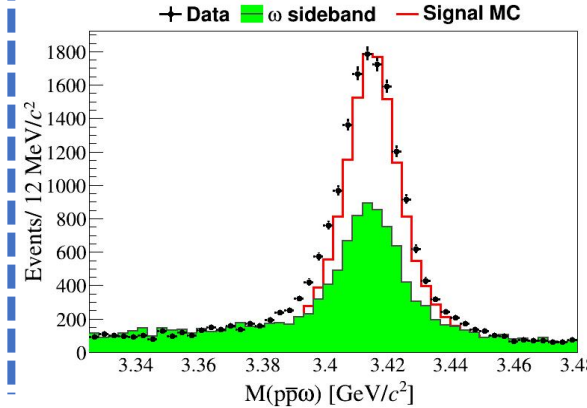
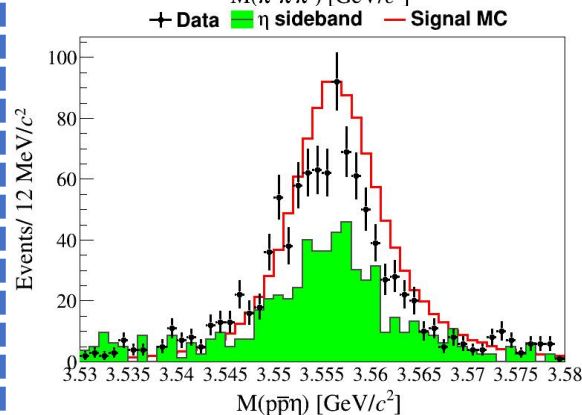
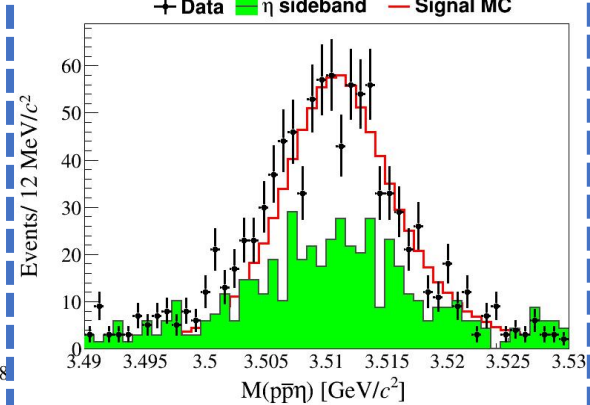
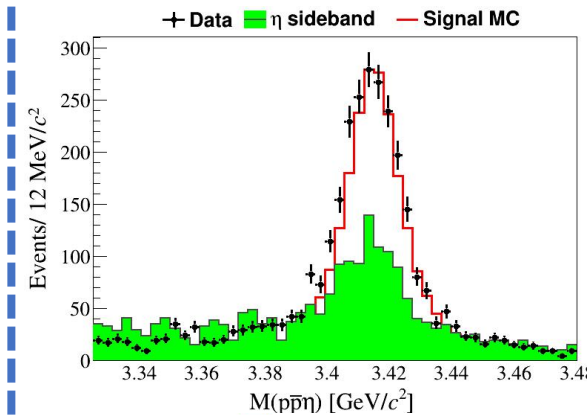
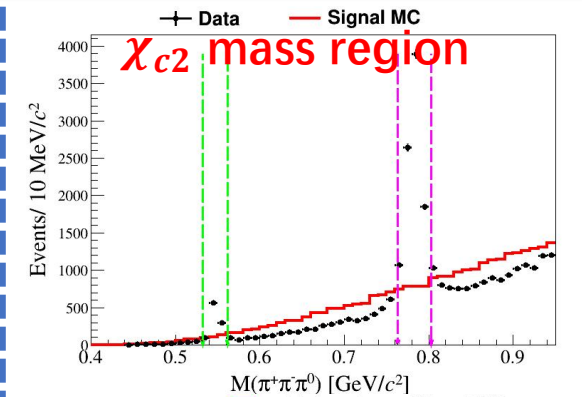
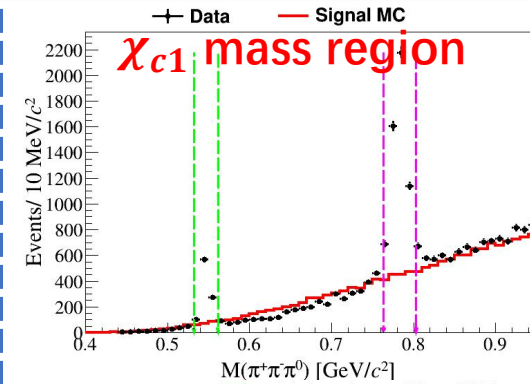
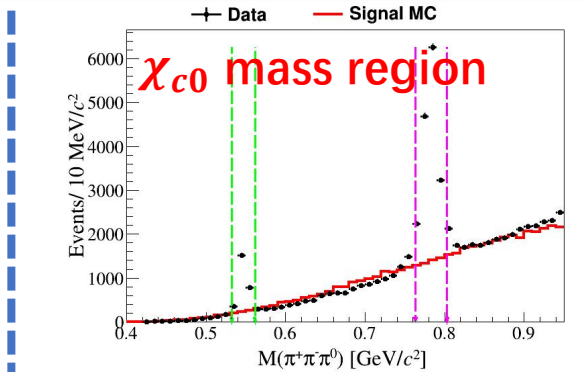
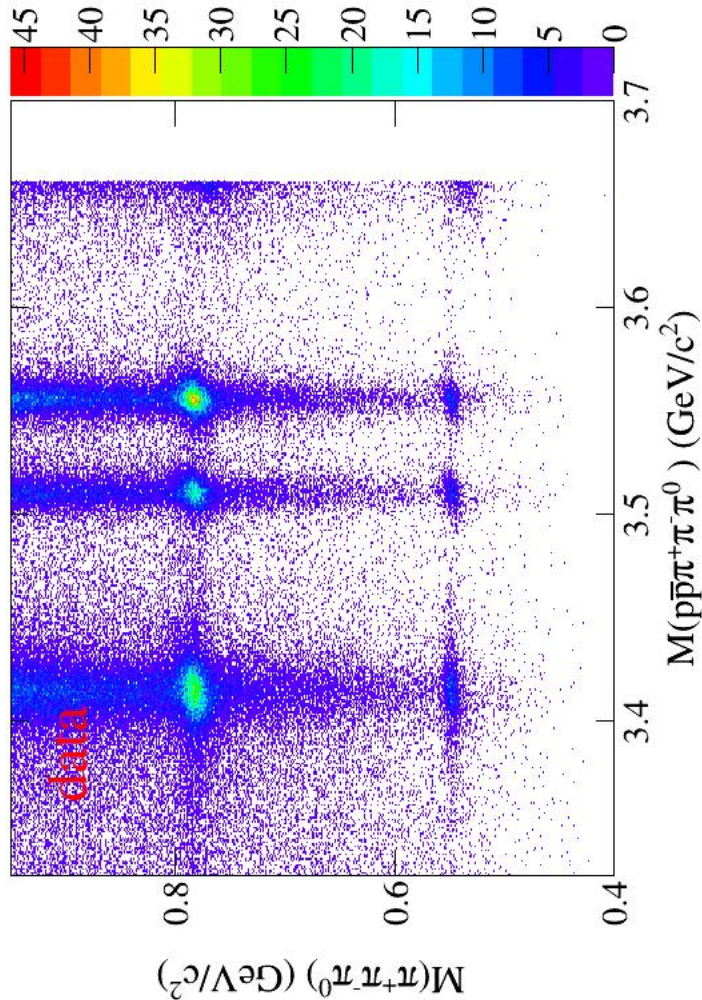
We calculate the upper limits of the branching fractions are :

$$B(\eta_c(2S) \rightarrow p\bar{p}\pi^+\pi^-\pi^0)$$

$$< \frac{N_{data}^{U.L.}}{N_{\psi(2S)}^{tot} \cdot B(\psi(2S) \rightarrow \gamma\eta_c(2S)) \cdot B(\pi^0 \rightarrow \gamma\gamma) \cdot \epsilon^{MC}}$$

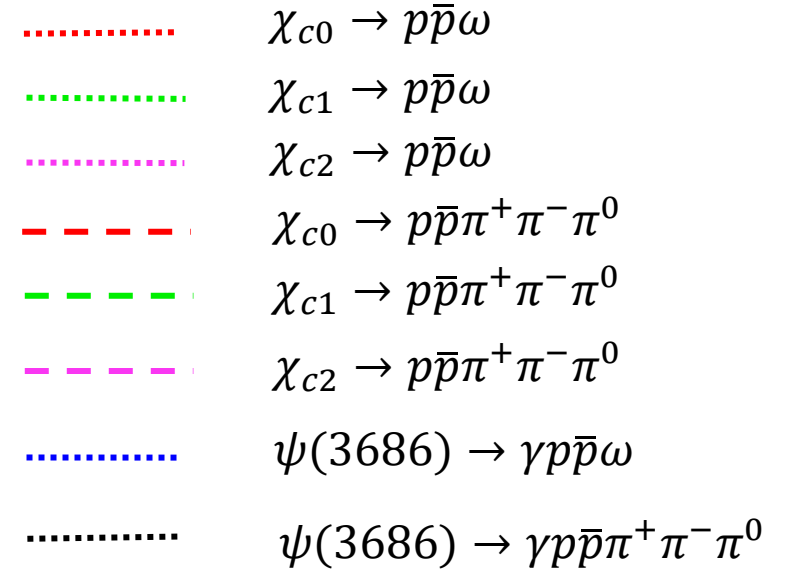
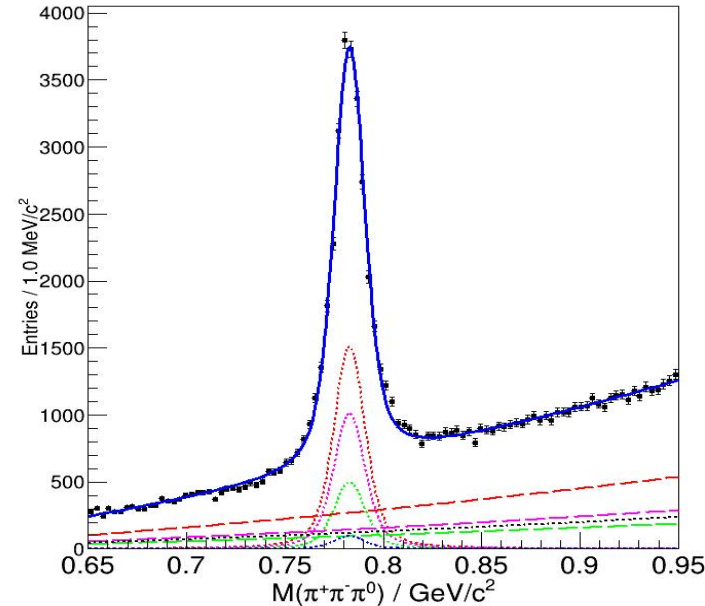
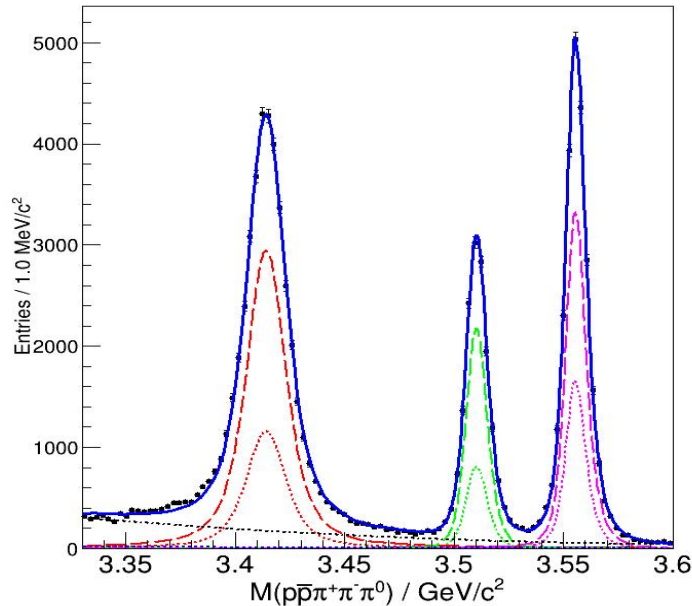
Channel	Branching fraction	PDG
$\eta_c(2S) \rightarrow p\bar{p}\pi^+\pi^-\pi^0$	$< 4.69 \times 10^{-3}$	-

Investigate the intermediate process



Pre. 2D fit to Investigate the intermediate process

Event type	$M(p\bar{p}\pi^+\pi^-\pi^0)$	$M(\pi^+\pi^-\pi^0)$
$\chi_{cJ} \rightarrow p\bar{p}\omega$	$\chi_{cJ}(\rightarrow p\bar{p}\omega)$ MC shape $\otimes G(\Delta m, \Delta\sigma)$	$BW \otimes G(m, \sigma)$
$\chi_{cJ} \rightarrow p\bar{p}\pi^+\pi^-\pi^0$	$\chi_{cJ}(\rightarrow p\bar{p}\pi^+\pi^-\pi^0)$ MC shape $\otimes G(\Delta m, \Delta\sigma)$	Polynomial function
$\psi(3686) \rightarrow \gamma p\bar{p}\omega$	Polynomial function	$BW \otimes G(m, \sigma)$
$\psi(3686) \rightarrow \gamma p\bar{p}\pi^+\pi^-\pi^0$	Polynomial function	Polynomial function



Preliminary result

$$B(\chi_{cJ} \rightarrow p\bar{p}\omega) = \frac{N_{\chi_{cJ}}^{obs}}{N_{\psi(2S)}^{tot} \cdot B(\psi(2S) \rightarrow \gamma\chi_{cJ}) \cdot B(\omega \rightarrow \pi^+\pi^-\pi^0) \cdot B(\pi^0 \rightarrow \gamma\gamma) \cdot \epsilon}$$

	N^{sig}	$\epsilon(\%)$	Br
$\chi_{c0} \rightarrow p\bar{p}\omega$	$(1.194 \pm 0.016) \times 10^4$	9.47	$(5.396 \pm 0.073) \times 10^{-4}$
$\chi_{c1} \rightarrow p\bar{p}\omega$	$(3.948 \pm 0.085) \times 10^3$	9.74	$(1.741 \pm 0.038) \times 10^{-4}$
$\chi_{c2} \rightarrow p\bar{p}\omega$	$(8.010 \pm 0.115) \times 10^3$	8.63	$(4.084 \pm 0.059) \times 10^{-4}$

Summary & next to do

- Summary:

Using about 2.7 billion $\psi(3686)$ events at BESIII,

- ✓ The evidence of $\eta_c(2S) \rightarrow p\bar{p}\pi^+\pi^-\pi^0$ is observed with 3.16σ for the first time, and its BR. is measured to be $(3.521 \pm 1.081) \times 10^{-3}$ and the upper limit of the BR. of $\eta_c(2S) \rightarrow p\bar{p}\pi^+\pi^-\pi^0 : < 4.69 \times 10^{-3}$.
- ✓ $\chi_{cJ} \rightarrow p\bar{p}\pi^+\pi^-\pi^0$ is observed for the first time, the BR. are determined to be $(4.880 \pm 0.017) \times 10^{-3}$, $(2.112 \pm 0.011) \times 10^{-3}$ and $(3.772 \pm 0.014) \times 10^{-3}$ for $J=0, 1$ and 2 , respectively.
- ✓ The intermediate process $\chi_{cJ} \rightarrow p\bar{p}\omega$ is also observed for the first time, their BR. are determined to be $(5.396 \pm 0.073) \times 10^{-4}$, $(1.741 \pm 0.038) \times 10^{-4}$, $(4.084 \pm 0.059) \times 10^{-4}$ for $J=0, 1$ and 2 , respectively.

- Next to do:

- Perform the analysis of systematic uncertainty

Observation of $\eta_c(2S) \rightarrow 3(\pi^+\pi^-)$ and Measurements of $\chi_{cJ}(1P) \rightarrow 3(\pi^+\pi^-)$ in $\psi(3686)$ radiative transitions

Abstract

The hadronic decay $\eta_c(2S) \rightarrow 3(\pi^+\pi^-)$ is observed with a statistical significance of 9.3 standard deviations using $(448.1 \pm 2.9) \times 10^6$ $\psi(3686)$ events collected by the BESIII detector at the BEPCII collider. The measured mass and width of $\eta_c(2S)$ are $(3643.4 \pm 2.3 \text{ (stat)} \pm 4.4 \text{ (syst)}) \text{ MeV}/c^2$ and $(19.8 \pm 3.9 \text{ (stat)} \pm 3.1 \text{ (syst)}) \text{ MeV}$, respectively, which are consistent with the world average values within two standard deviations. The product branching fraction $\mathcal{B}[\psi(3686) \rightarrow \gamma\eta_c(2S)] \times \mathcal{B}[\eta_c(2S) \rightarrow 3(\pi^+\pi^-)]$ is measured to be $(9.2 \pm 1.0 \text{ (stat)} \pm 1.2 \text{ (syst)}) \times 10^{-6}$. Using $\mathcal{B}[\psi(3686) \rightarrow \gamma\eta_c(2S)] = (7.0^{+3.4}_{-2.5}) \times 10^{-4}$, we obtain $\mathcal{B}[\eta_c(2S) \rightarrow 3(\pi^+\pi^-)] = (1.31 \pm 0.15 \text{ (stat)} \pm 0.17 \text{ (syst)}^{+0.64}_{-0.47} \text{ (extr)}) \times 10^{-2}$, where the third uncertainty is from $\mathcal{B}[\psi(3686) \rightarrow \gamma\eta_c(2S)]$. We also measure the $\chi_{cJ} \rightarrow 3(\pi^+\pi^-)$ ($J = 0, 1, 2$) decays via $\psi' \rightarrow \gamma\chi_{cJ}$ transitions. The branching fractions are $\mathcal{B}[\chi_{c0} \rightarrow 3(\pi^+\pi^-)] = (2.080 \pm 0.006 \text{ (stat)} \pm 0.068 \text{ (syst)}) \times 10^{-2}$, $\mathcal{B}[\chi_{c1} \rightarrow 3(\pi^+\pi^-)] = (1.092 \pm 0.004 \text{ (stat)} \pm 0.035 \text{ (syst)}) \times 10^{-2}$, and $\mathcal{B}[\chi_{c2} \rightarrow 3(\pi^+\pi^-)] = (1.565 \pm 0.005 \text{ (stat)} \pm 0.048 \text{ (syst)}) \times 10^{-2}$.

Search for $\eta_c(2S) \rightarrow \omega\omega/\omega\phi$ in $\psi(3686) \rightarrow \gamma\eta_c(2S)$ decay and Measurements of $\chi_{cJ}(1P) \rightarrow \omega\omega/\omega\phi$ in $\psi(3686) \rightarrow \gamma\chi_{cJ}(1P)$ decay

Abstract

Using a sample of 2708.1×10^6 $\psi(3686)$ events collected with the BESIII detector at the BEPCII collider in 2009, 2012 and 2021, we search for $\eta_c(2S) \rightarrow \omega\omega$ and $\omega\phi$ via $\psi(3686) \rightarrow \gamma\eta_c(2S)$. The ω is reconstructed with $\pi^+\pi^-\pi^0$ and ϕ with K^+K^- final states. No obvious signal events are observed. The upper limits of the branching fractions at 90% confidence level are determined to be $Br(\psi(2S) \rightarrow \gamma\eta_c(2S), \eta_c(2S) \rightarrow \omega\omega) < 6.13 \times 10^{-7}$ and $Br(\psi(2S) \rightarrow \gamma\eta_c(2S), \eta_c(2S) \rightarrow \omega\phi) < 1.08 \times 10^{-7}$, respectively. We also update the branching fractions of $\chi_{cJ}(1P) \rightarrow \omega\omega$ decays and $\chi_{cJ}(1P) \rightarrow \omega\phi$ decays via $\psi(2S) \rightarrow \gamma\chi_{cJ}$ transition. The branching fractions are determined to be $Br(\chi_{c0} \rightarrow \omega\omega) = (9.42 \pm 0.09 \pm 0.38) \times 10^{-4}$, $Br(\chi_{c1} \rightarrow \omega\omega) = (5.93 \pm 0.07 \pm 0.27) \times 10^{-4}$, $Br(\chi_{c2} \rightarrow \omega\omega) = (8.02 \pm 0.08 \pm 0.35) \times 10^{-4}$, $Br(\chi_{c0} \rightarrow \omega\phi) = (1.19 \pm 0.03 \pm 0.06) \times 10^{-4}$, $Br(\chi_{c1} \rightarrow \omega\phi) = (2.04 \pm 0.14 \pm 0.15) \times 10^{-5}$, and $Br(\chi_{c2} \rightarrow \omega\phi) = (8.48 \pm 0.96 \pm 0.61) \times 10^{-6}$, where the first uncertainty is statistical and second is systematic.

Share

9.1 Lineshape of $\eta_c(2S)$

Before fitting to the mass spectrum, the lineshape of signal is defined firstly. For the $\eta_c(2S)$ signal, the lineshape is given below

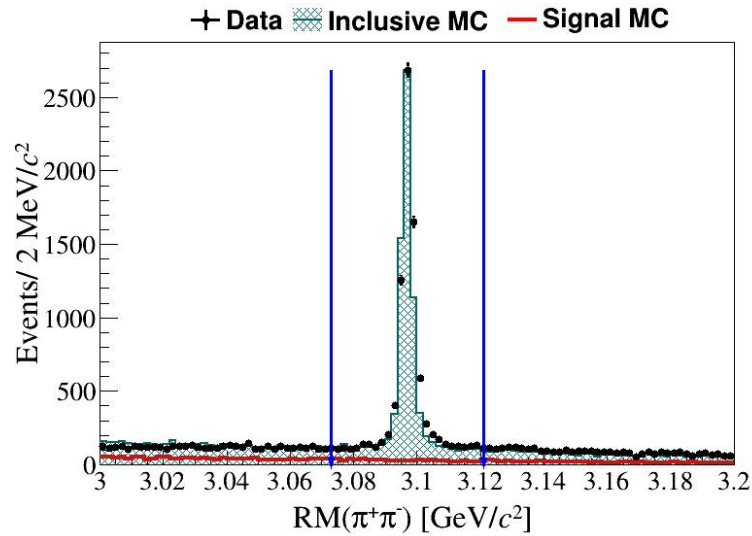
$$(BW(m; M, \Gamma) \times E_\gamma^3 \times f_d(E_\gamma) \otimes DG(m_1, \sigma_1, m_2, \sigma_2, f) \times \epsilon(m)) \otimes G(\delta m_3, \delta \sigma_3), \quad (13)$$

where m is the 3C kinematic fitted mass of $\omega\omega$ or $\omega\phi$, $BW(m; M, \Gamma)$ is the Breit-Wigner function for $\eta_c(2S)$ or χ_{cJ} , the form of the $BW(m)$ is $\frac{1}{2\pi} \cdot \frac{\Gamma}{(m-m_0)^2 + \Gamma^2/4}$ with M and Γ representing the mass and width of the resonance. $E_\gamma = \frac{M_{\psi(2S)}^2 - m^2}{2M_{\psi(2S)}}$ is the energy of the transition photon in the rest frame of $\psi(2S)$ ($M_{\psi(2S)}$ is the nominal mass of $\psi(2S)$). $f_d(E_\gamma)$ is the function to damp the diverging tail raised by E_γ^3 , the possible form of the damping function is used by KEDR for a similar process, which is $\frac{E_0^2}{E_\gamma E_0 + (E_\gamma - e_0)^2}$, where $E_0 = \frac{M_{\psi(2S)}^2 - M_{\eta_c(2S)}^2}{2M_{\psi(2S)}}$ is the peaking energy of the transition photon. $DG(m_1, \sigma_1, m_2, \sigma_2, f)$ is a double Gaussian function used to describe the mass resolution, the parameters are estimated by using signal MC samples. $\epsilon(m)$ is the MC efficiency as a function of m . $G(\delta m_3, \delta \sigma_3)$ is a single Gaussian function, which describes the difference of mass and resolution between data and MC sample.

Thanks

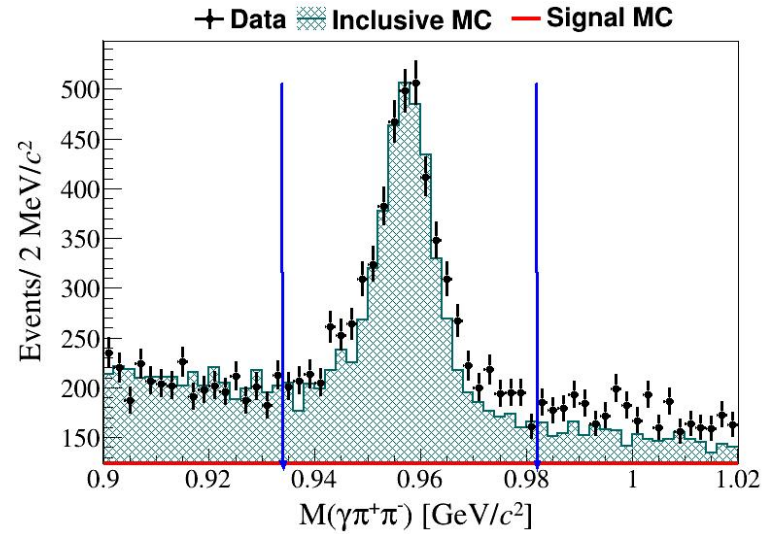
Back up

$\pi^+\pi^-J/\psi$ background



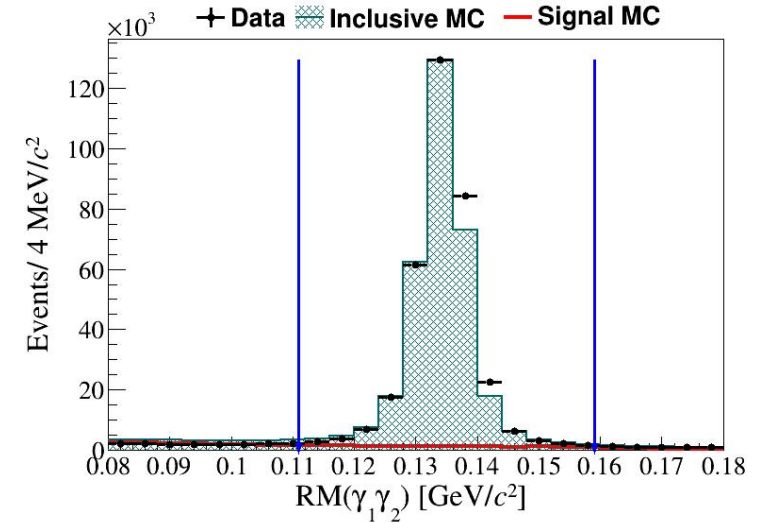
➤ $|\text{RM}(\pi^+\pi^-) - m(J/\psi)| > 24 \text{ MeV}/c^2$

η' background



➤ $|\text{M}(\gamma\pi^+\pi^-) - m(\eta')| > 24 \text{ MeV}/c^2$

π^0 background



➤ $|\text{M}(\gamma_1\gamma_2) - m(\pi^0)| > 24 \text{ MeV}/c^2$

