



Search for Lepton Number Violation Processes at BESIII

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Outline

• Introduction

- LNV Results at BESIII $\geq D^0 \rightarrow K^- \pi^- e^+ e^+, D^+ \rightarrow K_s^0 \pi^- e^+ e^+, D^+ \rightarrow K^- \pi^0 e^+ e^+ \text{(PRD99, 112002)}$ $\geq \Sigma^- \rightarrow p e^- e^-, \Sigma^- \rightarrow \Sigma^+ X \text{ (including } e^- e^-) \text{ (PRD103, 052011)}$ $\geq \phi \rightarrow \pi^+ \pi^+ e^- e^- \text{ via } J/\psi \rightarrow \phi \eta \text{ (arXiv: 2308.05490)}$
- Prospects
- Summary

Introduction

- Lepton number is conserved in the Standard Model.
- SM can't explain the finite mass of neutrinos.
- The nature of neutrinos (Dirac or Majorana), is still an open question.
 - If neutrinos are Majorana particles, they can be manifested through $\Delta L = 2$ processes.
 - $\Delta L = 2$ processes are widely viewed as the cleanest test of the Majorana neutrino.

Introduction

- The asymmetry of matter and antimatter in the universe \rightarrow possible BNV
 - Many theories believe that if there is BNV, there will also be LNV.
- LNV is also introduced in many New Physics models.
 - > 4th quark generation
 - ≻ SO(10) SUSY GUT
 - ≻Exotic Higgs

▶.....

 $D^{0(+)} \rightarrow hh'e^+e^+$

- $\Delta L = 2$ processes
- Decay diagrams proposed at BESIII [Chin.Phys. C39,013101(2015)]



 $D^{0(+)} \rightarrow hh'e^+e^+$

PRD99, 112002

- Data: 2.9 fb^{-1} @3.773 GeV
- Single tag analysis
- Two variables are used to identify the signal candidates:

$$M_{\rm BC} = \sqrt{E_{\rm beam}^2 - |\vec{p}_D|^2},$$
$$\Delta E = E_D - E_{\rm beam},$$

• Fit the ΔE distribution in signal MC samples.

Channel	$\Delta E \ ({ m MeV})$		
$D^0 \to K^- \pi^- e^+ e^+$	[-33.0, 19.7]		
$D^+ \to K^0_S \pi^- e^+ e^+$	[-30.6, 19.3]		
$D^+ \to K^- \pi^0 e^+ e^+$	[-54.8, 24.4]		

 $D^{0(+)} \rightarrow hh'e^+e^+$

PRD99, 112002



- Scanning likelihood value for different number of signal events.
- Likelihood curve convolved with a Gaussian function.

Channel	$\epsilon(\%)$	$N_{ m sig}^{ m UL}$	$\mathcal{B}_{ m sig}^{ m UL}(imes 10^{-6})$	
$D^0 \to K^- \pi^- e^+ e^+$	16.8	10.0	< 2.8	improved by 2 orders
$D^+ \to K^0_S \pi^- e^+ e^+$	11.5	4.4	< 3.3	
$D^+ \to K^- \pi^0 e^+ e^+$	10.6	14.8	< 8.5	first search

MC simulated shape convolved with a Gaussian function + ARGUS function

 $D^{0(+)} \rightarrow hh'e^+e^+$

PRD99, 112002

• Neutiono mass related UL for $D^0 \to K^- e^+ \nu_m(\pi^- e^+)$ and $D^+ \to K^0_s e^+ \nu_m(\pi^- e^+)$



$$\frac{\Gamma(m_{\nu_m}, V_{e\nu_m}(m_{\nu_m}))}{\Gamma(m_{\nu_m}, V'_{e\nu_m}(m_{\nu_m}))} = \frac{|V_{e\nu_m}(m_{\nu_m})|^4}{|V'_{e\nu_m}(m_{\nu_m})|^4}$$

 $\Sigma^- \rightarrow p e^- e^-, \Sigma^- \rightarrow \Sigma^+ X$

- Two down-type (d or s) quarks convert into two up-quarks
- Determined by local four-quark operators
- Underlying mechanism similar to that of $0\nu\beta\beta$ nuclear decay



• Based on MIT bag Model, predicted BF of $\Sigma^- \rightarrow pe^-e^-$ can reach 10⁻²³

$$\Sigma^- \to p e^- e^-, \Sigma^- \to \Sigma^+ X$$

PRD103, 052011

- Data: about 1 billion J/ψ events @3.097 GeV
- Double tag analysis $\mathcal{B}^{\text{sig}} = \frac{N^{\text{sig}} / \varepsilon^{\text{tag,sig}}}{N^{\text{tag}} / \varepsilon^{\text{tag}}}$ $J/\psi \to \overline{\Sigma}(1385)^{+}\Sigma^{-}$ tag side: $\overline{\Sigma}(1385)^{+} \to \overline{\Lambda}\pi^{+}$ Search for signal by looking at the recoiling mass of $\overline{\Sigma}(1385)^{+}$ Only reconstruct Σ^{+} for $\Sigma^{-} \to \Sigma^{+}X$

$$M_{\text{recoil}} = \sqrt{(E_{\mathrm{J}/\psi} - E_{\bar{\Lambda}} - E_{\pi^+})^2 - (\vec{\mathbf{p}}_{\mathrm{J}/\psi} - \vec{\mathbf{p}}_{\bar{\Lambda}} - \vec{\mathbf{p}}_{\pi^+})^2}$$

• Blind analysis



MC simulated shape convolved with a Gaussian function + 2nd order Chebyshev polynomial

$$\Sigma^- \rightarrow pe^-e^-, \Sigma^- \rightarrow \Sigma^+ X_{\text{PRD103, 052011}}$$

- Implemented by the package TROLKE
- UL @ 90% C.L.

$$\mathcal{B}(\Sigma^- \to p e^- e^-) < 6.7 \times 10^{-5},$$
$$\mathcal{B}(\Sigma^- \to \Sigma^+ X) < 1.2 \times 10^{-4}.$$

first search



 $\phi \rightarrow \pi^+ \pi^+ e^- e^- \operatorname{via} J/\psi \rightarrow \phi \eta$

• Extremely suppressed by several loops

arXiv: 2308.05490

• If any hint can be seen, definitely signal for new physics



$$\phi \rightarrow \pi^+ \pi^+ e^- e^- \operatorname{via} J/\psi \rightarrow \phi \eta$$

• Data: about 10 billion J/ψ events @3.097 GeV

arXiv: 2308.05490

• Relative measurement to $\phi \to K^+K^-$, lots of uncertainties cancelled

$$\mathcal{B}(\phi \to \pi \pi e e) = \mathcal{B}(\phi \to KK) \times \frac{\frac{N_{\pi \pi e e}^{\text{net}}}{\varepsilon_{\pi \pi e e}}}{\frac{N_{KK}^{\text{net}}}{\varepsilon_{KK}}} = \mathcal{B}(\phi \to KK) \times \frac{N_{\pi \pi e e}^{\text{net}} / \varepsilon_{\pi \pi e e}}{N_{KK}^{\text{net}} / \varepsilon_{KK}}$$



• Signal identified with two variables: $M_{\pi^+\pi^+e^-e^-} \& M_{\gamma\gamma}$

• Blind analysis

2023/10/15

 $\phi \rightarrow \pi^+ \pi^+ e^- e^- \operatorname{via} J/\psi \rightarrow \phi \eta$

arXiv: 2308.05490

$$N_{\rm obs}^{\rm MC} = N_{\phi \to K^+ K^-}^{\rm signal} - \frac{1}{2} \times N_{\phi \to K^+ K^-}^{\rm sideband}$$

$$\varepsilon_{KK} = \frac{N_{\text{obs}}^{\text{MC}}}{N_{\text{total}}^{\text{MC}}} = \frac{471297}{1000000} = 47.1\%$$

$$N_{K^+K^-}^{\text{net}} = N_{\text{signal}} - \frac{1}{2} \times N_{\text{sideband}} = 822665 \pm 1149$$

(a) $\phi \rightarrow K^{+}K^{-}$ Events / (1.1 MeV/c²) (×10³) 2 001 Signal region **2**F **(b)** $\phi \rightarrow K^{+}K^{-}$ Sideband region 0.5 1.05 1.1 $M_{K^+K^-}$ (GeV/ c^2)

MC simulated shape convolved with a double Gaussian function + inverted ARGUS function multiplied by a 4th order polynomial 2023/10/15

 $\phi \rightarrow \pi^+ \pi^+ e^- e^- \text{via } J/\psi \rightarrow \phi \eta$

arXiv: 2308.05490



• No signal was observed. UL @ 90% C.L. calculated with TROLKE

first search $\mathcal{B}(\phi \to \pi^+ \pi^+ e^- e^-) < 3.7 \times 10^{-6}$

Prospects

BESIII Collaboration, Future physics programme of BESIII, Chin. Phys. C 44 (2020), 040001

- BESIII has a potential to search for LNV transitions such as $D^+ \rightarrow l^+l^+X^-$ and $D_s^+ \rightarrow l^+l^+X^-$, due to the clean environment and low charge confusion rates.
- With the coming 20 fb⁻¹ of $\psi(3770)$ data, can perform more sensitive searches for LNV processes of D mesons.
 - could improve the best upper limit to 4.6×10^{-7} and 2.3×10^{-7} for $D^+ \rightarrow \pi^- e^+ e^+$ and $D^+ \rightarrow K^- e^+ e^+$.

Prospects

Averages of b-hadron, c-hadron, and τ -lepton properties as of 2021. (HFLAV Collaboration). PhysRevD.107.052008



- Many modes haven't been searched for.
- Four-body LNV decays of D_s; LNV processes of charmed baryons.

Prospects

- LNV processes of charmonium particles such as J/ψ
- of mesons such as ω , η/η' , K_S^0



Figure 1: One possible Feynman diagram of $J/\psi \to K^+K^+e^-e^- + c.c..$

BESIII Document 1091



Fig. 1: Possible Feynman diagrams of $\bar{K}^0 \rightarrow \pi^+ \pi^+ e^- e^- + c.c.$

BESIII BAM-00651

Summary

• In recent years, BESIII have searched for LNV processes.

Upper limits of many processes have been determined, providing stringent restrictions on the parameters of the future theoretical development.

• With the coming 20 fb⁻¹ data samples, more and better results for LNV process searching are expected.

Back Up

Chin.Phys. C39,013101

$$\mathcal{M} = 2G_{\rm F}^2 V_{\rm cs} V_{\rm ud} \langle K | \bar{s} \gamma^{\mu} (1 - \gamma_5) c | D \rangle \langle \pi | \bar{u} \gamma^{\nu} (1 - \gamma_5) d | 0 \rangle$$

$$\times \left[V_{\rm lN}^2 m_{\rm N} \bar{u}(p_2) \left(\frac{\gamma^{\mu} \gamma^{\nu}}{q_{\rm N}^2 - m_{\rm N}^2 + \mathrm{i} \Gamma_{\rm N} m_{\rm N}} + \frac{\gamma^{\nu} \gamma^{\mu}}{q_{\rm N}^{\prime 2} - m_{\rm N}^2 + \mathrm{i} \Gamma_{\rm N} m_{\rm N}} \right) P_{\rm R} v(p_3) \right], \qquad (4)$$

$$f_i = \frac{|\mathcal{M}_i|^2}{\sum_i |\mathcal{M}_i|^2} |\sum_i \mathcal{M}_i|^2,$$
$$\mathcal{M}|^2 = |\sum_i \mathcal{M}_i|^2 = \sum_i f_i.$$

- $\begin{array}{l} \mathcal{C}_{1} &= f_{\pi}^{2}|PN_{2}|^{2}G_{F}^{4}|V_{cs}|^{2}|V_{ud}|^{2} \Big\{ 8D_{0}^{2}(f_{0}-f_{1})^{2}m_{K}^{4}(-2m_{\pi}^{2}((p_{23}+p_{24}-p_{34})m_{1}^{2}+m_{1}^{4}-2p_{23}p_{34}) \\ &+ 4p_{34}^{2}(m_{1}^{2}+p_{23})+m_{\pi}^{4}p_{23})+m_{K}^{2}(16f_{1}(-D_{0}(f_{0}-f_{1})(-2m_{\pi}^{2}((p_{12}+p_{13}+p_{24}+p_{24}-p_{34})m_{1}^{2}) \\ &+ m_{1}^{4}-p_{14}p_{23}+p_{13}p_{24}-p_{12}p_{34}-2p_{23}p_{34})+4p_{34}((p_{14}+p_{34})m_{1}^{2}+p_{14}p_{23}-p_{13}p_{24}+p_{12}p_{34} \\ &+ p_{23}p_{34})+m_{\pi}^{4}p_{23})+2f_{1}m_{\pi}^{2}p_{23}-4f_{1}p_{24}p_{34})-16D_{0}^{2}(f_{0}-f_{1})^{2}m_{D}^{2}(-2m_{\pi}^{2}((p_{23}+p_{24}-p_{34})m_{1}^{2}) \\ &+ m_{1}^{4}-2p_{23}p_{34})+4p_{34}^{2}(m_{1}^{2}+p_{23})+m_{\pi}^{4}p_{23}) +8D_{0}^{2}(f_{0}-f_{1})^{2}m_{D}^{4}(-2m_{\pi}^{2}((p_{23}+p_{24}-p_{34})m_{1}^{2}) \\ &+ m_{1}^{4}-2p_{23}p_{34})+4p_{34}^{2}(m_{1}^{2}+p_{23})+m_{\pi}^{4}p_{23}) +16D_{0}(f_{0}-f_{1})f_{1}m_{D}^{2}(-2m_{\pi}^{2}((p_{12}+p_{13}+p_{13}+p_{23})) \\ &+ p_{24}-p_{34}m_{1}^{2}+m_{1}^{4}-p_{14}p_{23}+p_{13}p_{24}-p_{12}p_{34}-2p_{23}p_{34})+4p_{34}((p_{14}+p_{34})m_{1}^{2}+p_{14}p_{23} \\ &- p_{13}p_{24}+p_{12}p_{14}+p_{23}p_{34})+m_{\pi}^{4}p_{23}) +8f_{1}^{2}(-2m_{\pi}^{2}((2p_{12}+2p_{13}+p_{23}+p_{24}-p_{34})m_{1}^{2}+m_{1}^{4} \\ &+ 4p_{12}p_{13}-2p_{14}p_{23}+2p_{13}p_{24}-2p_{12}p_{34}-2p_{23}p_{34})+4p_{34}((2p_{14}+p_{34})m_{1}^{2}+4p_{12}p_{14}+2p_{14}p_{24} \\ &- 2p_{13}p_{24}+2p_{12}p_{34}+p_{23}p_{34})+m_{\pi}^{4}p_{23}) \Big\} \end{array}$
- $$\begin{split} \mathcal{C}_2 &= f_\pi^2 |PN_3|^2 G_F^4 |V_{cs}|^2 |V_{ud}|^2 \Big\{ 8 D_0^2 (f_0 f_1)^2 m_K^4 (-2m_\pi^2 ((p_{23} p_{24} + p_{34}) m_1^2 + m_1^4 2p_{23} p_{24}) \\ &+ 4p_{24}^2 (m_1^2 + p_{23}) + m_\pi^4 p_{23}) + m_K^2 (16f_1 (-D_0 (f_0 f_1) (-2m_\pi^2 ((p_{12} + p_{13} + p_{23} p_{24} + p_{34}) m_1^2 \\ &+ m_1^4 p_{14} p_{23} p_{13} p_{24} 2p_{23} p_{24} + p_{12} p_{34}) + 4p_{24} ((p_{14} + p_{24}) m_1^2 + p_{14} p_{23} + p_{13} p_{24} + p_{23} p_{24} \\ &- p_{12} p_{34}) + m_\pi^4 p_{23}) + 2f_1 m_\pi^2 p_{23} 4f_1 p_{24} p_{34}) 16 D_0^2 (f_0 f_1)^2 m_D^4 (-2m_\pi^2 ((p_{23} p_{24} + p_{34}) m_1^2 \\ &+ m_1^4 2p_{23} p_{24}) + 4p_{24}^2 (m_1^2 + p_{23}) + m_\pi^4 p_{23})) + 8 D_0^2 (f_0 f_1)^2 m_D^4 (-2m_\pi^2 ((p_{22} p_{24} + p_{34}) m_1^2 \\ &+ m_1^4 2p_{23} p_{24}) + 4p_{24}^2 (m_1^2 + p_{23}) + m_\pi^4 p_{23}) + 16 D_0 (f_0 f_1) f_1 m_D^2 (-2m_\pi^2 ((p_{12} + p_{14} + p_{23} + p_{13} p_{24} + p_{23} p_{24} p_{12} p_{34}) + m_\pi^4 p_{23}) + 4p_{24} ((p_{14} + p_{24}) m_1^2 + p_{14} p_{23} \\ &+ p_{13} p_{24} + p_{23} p_{24} p_{12} p_{34}) + m_\pi^4 p_{23}) + 8f_1^2 (-2m_\pi^2 ((2p_{12} + 2p_{13} + p_{23} p_{24} + p_{34}) m_1^2 + m_1^4 \\ &- 2(p_{14} p_{23} + (p_{13} + p_{23}) p_{24}) + 2p_{12} (2p_{13} + p_{34})) + 4p_{24} ((2p_{14} + p_{24}) m_1^2 + 2p_{14} p_{23} + p_{23} p_{24} \\ &+ 2p_{13} (2p_{14} + p_{24}) 2p_{12} p_{24}) + m_\pi^4 p_{23}) \Big\} \end{split}$$
- $$\begin{split} \mathcal{C}_{3} &= \int_{\pi}^{2} PN_{2}^{*} PN_{3} G_{F}^{4} |V_{cs}|^{2} |V_{ud}|^{2} \Big\{ 8D_{0}^{2} (f_{0} f_{1})^{2} (p_{23} m_{\pi}^{4} + 2(m_{1}^{4} + p_{23} m_{1}^{2} + p_{23} (p_{24} + p_{34})) m_{\pi}^{2} \\ &\quad -2(m_{1}^{2} (p_{24}^{2} + p_{34}^{2}) 2p_{23} p_{24} p_{34})) m_{0}^{4} + 16D_{0} (f_{0} f_{1}) f_{1} (p_{23} m_{\pi}^{4} + 2(m_{1}^{4} + (p_{12} + p_{13} + p_{23}) m_{1}^{2} \\ &\quad + p_{23} (p_{14} + p_{24} + p_{34})) m_{\pi}^{2} 2((p_{24}^{2} + p_{34}^{2} + p_{14} (p_{24} + p_{34})) m_{1}^{2} + p_{13} p_{24}^{2} + p_{12} p_{34}^{2} p_{12} p_{24} p_{34} \\ &\quad p_{13} p_{24} p_{34} 2p_{23} p_{24} p_{34} p_{14} p_{23} (p_{24} + p_{34})) m_{\pi}^{2} + \zeta(-32 i (m_{\pi}^{2} + 2p_{14} + p_{24} + p_{34}) f_{1}^{2} \\ &\quad 32 i D_{0} (f_{0} f_{1}) m_{\pi}^{2} (m_{\pi}^{2} + p_{24} + p_{34}) f_{1} + 32 i D_{0} (f_{0} f_{1}) m_{K}^{2} (m_{\pi}^{2} + p_{24} + p_{34}) f_{1} \\ &\quad + 8 D_{0}^{2} (f_{0} f_{1})^{2} m_{K}^{4} (p_{24} m_{1}^{4} + p_{23} m_{1}^{2} + p_{23} (p_{24} + p_{34})) m_{\pi}^{2} 2(m_{1}^{2} (p_{24}^{2} + p_{34}^{2}) 2p_{23} p_{24} + p_{34}) m_{\pi}^{2} 2(m_{1}^{2} (p_{24}^{2} + p_{34}^{2}) m_{\pi}^{2} 2(p_{24}^{2} + p_{34}^{2}) m_{1}^{2} 4p_{12}^{2} p_{13} + p_{23} (p_{24} + p_{34})) m_{\pi}^{2} \\ &\quad -2((p_{24}^{2} + p_{34}^{2}) m_{1}^{2} 4p_{14}^{2} p_{23} + 2p_{24} p_{24} p_{24} + p_{24} + p_{34})) m_{\pi}^{2} \\ &\quad -2(p_{13} (p_{4} + p_{24} p_{24}) p_{4} + p_{14} (p_{24} + p_{34}) p_{12} p_{24} p_{24} 2p_{23} p_{24} p_{34})) m_{\pi}^{2} \\ &\quad -4(m_{1}^{4} + p_{23} m_{1}^{2} + p_{23} (p_{24} + p_{34})) m_{\pi}^{2} 2(m_{1}^{2} (p_{24}^{2} + p_{34}^{2} p_{22} p_{24} p_{24}) m_{\pi}^{2} \\ &\quad -4(m_{1}^{4} + p_{23} m_{1}^{2} + p_{13} p_{24}^{2} + p_{12} p_{24}^{2} p_{22} p_{24} p_{24} p_{34}) m_{\pi}^{2} \\ &\quad -2((p_{24}^{2} + p_{34}^{2} + p_{14} (p_{24} + p_{34})) m_{\pi}^{2} 2(m_{1}^{2} (p_{24}^{2} + p_{34}^{2} p_{12} p_{24} p_{24} p_{13} p_{24} p_{24} 2p_{23} p_{24} p_{24} p_{34} \\ &\quad -4(p_{24}^{2} + p_{34} + p_{14} (p_{24} + p_{34})) m_{\pi}^{2} + p_{13} p_{24}^{2} + p_{23} (p_{24} + p_{2$$

 $|\mathcal{M}|^2 = C_1 + C_2 + C_3$

PRD87,036010

The most general form of the low-energy effective Lagrangian that is relevant for LNV semileptonic hyperon decays

$$-\mathcal{L}_{\beta\beta} = \frac{G_F^2}{\Lambda_{\beta\beta}} \{ c_1(\bar{u}\Gamma_i d)(\bar{u}\Gamma_j d) + c_2[(\bar{u}\Gamma_i d)(\bar{u}\Gamma_j s) + (\bar{u}\Gamma_i s)(\bar{u}\Gamma_j d)] + c_3(\bar{u}\Gamma_i s)(\bar{u}\Gamma_j s) \} \{ d_1(\bar{e}\Gamma_k e^c) + d_2(\bar{\mu}\Gamma_k \mu^c) + d_3(\bar{e}\Gamma_k \mu^c + \bar{\mu}\Gamma_k e^c) \}.$$
(1)