## Pseudoscalar mesons & Emergent Mass

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**Pseudoscalar Mesons and Emergent Mass** 

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## **QCD: Emergent Phenomena**

- QCD is characterized by two emergent phenomena: confinement and dynamical generation of mass (DGM).
- Quarks and gluons not *isolated* in nature.
- → Formation of colorless bound states: "<u>Hadrons</u>"
- 1-fm scale size of hadrons?



 Emergence of hadron masses (EHM) from QCD dynamics





#### **QCD: Emergent Phenomena**

QCD is characterized by two emergent phenomena: confinement and dynamical generation of mass (DGM).





$$\mathcal{L}_{\text{QCD}} = \sum_{j=u,d,s,\dots} \bar{q}_j [\gamma_\mu D_\mu + m_j] q_j + \frac{1}{4} G^a_{\mu\nu} G^a_{\mu\nu},$$
$$D_\mu = \partial_\mu + ig \frac{1}{2} \lambda^a A^a_\mu,$$
$$G^a_{\mu\nu} = \partial_\mu A^a_\nu + \partial_\nu A^a_\mu - \underline{g} f^{abc} A^b_\mu A^c_\nu,$$

 Emergence of hadron masses (EHM) from QCD dynamics



## **QCD: Emergent Phenomena**

QCD is characterized by two emergent phenomena: confinement and dynamical generation of mass (DGM).

#### Can we trace them down to fundamental d.o.f?



 $\begin{aligned} \mathcal{L}_{\text{QCD}} &= \sum_{j=u,d,s,\dots} \bar{q}_j [\gamma_\mu D_\mu + m_j] q_j + \frac{1}{4} G^a_{\mu\nu} G^a_{\mu\nu}, \\ D_\mu &= \partial_\mu + ig \frac{1}{2} \lambda^a A^a_\mu, \\ G^a_{\mu\nu} &= \partial_\mu A^a_\nu + \partial_\nu A^a_\mu - \underline{g} f^{abc} A^b_\mu A^c_\nu, \end{aligned}$ 

 Emergence of hadron masses (EHM) from QCD dynamics



Gluon and quark running masses

#### Mass Budgets

 $M_{u/d} \approx 0.3 \,\mathrm{GeV}$ 

➤ What is the origin of EHM?

... its connection with *e.g.* **confinement** and **DCSB**?

- Most of the mass in the visible universe is contained within nucleons
  - Which remain pretty massive whether there is Higgs mechanism or not...





Proton and rho meson mass budgets are practically identical

#### Mass Budgets

 $m_s/m_u \sim 20$  $f_K/f_\pi \sim M_s/M_u \sim 1.2$ 

➤ What is the origin of EHM?

... its connection with e.g. **confinement** and **DCSB**?

- > What is the role of **NG-bosons**?
- Pion and Kaon would be massless in the absence of Higgs mass generation.
  - → And structurally alike.

$$m_{\pi} = 0.14 \,\mathrm{GeV} \neq M_u + M_d$$
  
 $m_K = 0.49 \,\mathrm{GeV} \neq M_u + M_s$ 



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#### Pion and Kaon

- Both quark-antiquark bound-states and NG bosons
  - Their mere existence is connected with mass generation in the SM



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The scrutiny of all pseudoscalars sheds light on the role of Mass Sources on the hadron structural properties.

Mass Budgets

 $m_s/m_u \sim 20$  $f_K/f_\pi \sim M_s/M_u \sim 1.2$ 

# Valence-quark distribution amplitudes (DAs) $f_M \phi_M^q(x) = \operatorname{tr} \int_{dk} \delta_n^x(k_M) \gamma_5 \gamma \cdot n \chi_M(k_-, P)$

Light-front momentum fraction

Written in terms of **BSWF** 

- 1-dimensional projection of the light-front wavefunction.
- Clear probe of EHM, related with hard exclusive processes, etc.

#### π-K DAs



#### 'Heavy' mesons DAs

' In systems with heavy quarks, the **DAs** become **narrow**.



Unlike the Kaon, heavy-light systems DAs are markedly skewed

+ The peaks located at:  $x_{\max}^{\pi, K, D, B} = 0.5, \, 0.4, \, 0.18, \, 0.1$ 

#### **Drawing boundaries**

 Systems with ss-bar components draw the line between strong and weak mass generation being dominant.



Х



M. Ding *et al.*, PRD 101 5, 054014 (2020) Z-F Cui *et al.*, EPJC 80 11, 1064 (2021)

# **Distribution functions (DFs)**





- Yields *e.g.* information on **momentum** distribution.
- Evolution disentangles valence, sea and gluon contributions.

#### **π-K DFs: hadronic scale**

Fully-dressed valence quarks (quasiparticles)

- $\zeta_H$  : hadronic scale
- At this scale, **all properties** of the hadron are contained within their valence quarks.

 $(M_u = M_d)$ 

 Equally massive quarks symmetric distributions and equitable distribution of momentum fraction:

$$\langle x \rangle_{\pi}^{u} = 0.5, \ u_{\pi}(x; \zeta_{H}) = u_{\pi}(1-x; \zeta_{H})$$

The kaon distributions are only-shifted by a fewpercentage.

$$\langle x \rangle_K^u = 0.48, \ \langle x \rangle_K^{\bar{s}} = 0.52$$



Endpoint **behavior** is a reflection of the **underlying interaction** 

$$1/(k^2)^{\beta} \to (1-x)^{2\beta}$$

Farrar:1975yb

#### **π-K DFs: hadronic scale**



- > As the **DAs**, the  $\pi$ -K **DFs** are **dilated**.
- The bridge between DA and DF is the light-front wavefunction:

$$f_{\mathsf{P}}\varphi_{\mathsf{P}}^{u}(x,\zeta_{\mathcal{H}}) = \int \frac{dk_{\perp}^{2}}{16\pi^{3}}\psi_{\mathsf{P}}^{u}\left(x,k_{\perp}^{2};\zeta_{\mathcal{H}}\right)$$
$$u^{\mathsf{P}}(x;\zeta_{\mathcal{H}}) = \int \frac{d^{2}k_{\perp}}{16\pi^{3}} \left|\psi_{\mathsf{P}}^{u}\left(x,k_{\perp}^{2};\zeta_{\mathcal{H}}\right)\right|^{2}$$

 $\succ$  Such connection suggests:

$$\boldsymbol{u}^{P}(x;\zeta_{H}) \sim [\boldsymbol{\varphi}^{\boldsymbol{u}}_{P}(x;\zeta_{H})]^{2}$$

 Which is a fair approximation for integrated quantities of NG bosons



#### Pion DFs: Lattice & Experiment



At **5.2 GeV**, the experimental scale, our predictions matches that from Aicher *et al.* 



An agreement with novel **lattice** "Cross Section" results is also obtained.

At 2 GeV, the valence DF shows agreement with lattice moments:

$\zeta_2$	$\langle x \rangle_u^{\pi}$	$\langle x^2 \rangle_u^\pi$	$\langle x^3 \rangle_u^\pi$
Ref. [34]	0.24(2)	0.09(3)	0.053(15)
Ref. [35]	0.27(1)	0.13(1)	0.074(10)
Ref. [36]	0.21(1)	0.16(3)	
Herein	0.24(2)	0.098(10)	0.049(07)

#### □ The **Gluon DF** profiles matches **lattice** expectations:



#### Kaon DFs: Lattice & Experiment



## Pion vs Proton

- The (nearly) massless pion DFs differs vastly from the massive proton. For instance:
- Counting rules entail large-x behaviors (1-x)<sup>2</sup> and (1-x)<sup>3</sup> for the pion and proton, respectively.
- $\text{The momentum fractions at } \zeta_{\text{H}}: \quad (M_u = M_d) \\ \langle x \rangle_{u_p}^{\zeta_{\mathcal{H}}} = 0.687 , \ \langle x \rangle_{d_p}^{\zeta_{\mathcal{H}}} = 0.313 , \ \langle x \rangle_{u_{\pi}}^{\zeta_{\mathcal{H}}} = 0.5$

 $\Rightarrow u_V(x) \neq 2d_V(x)$  EHM induced diquark correlations inside the proton:

- No equitable distribution of momentum!
- Differences are preserved after evolution.



# **Electromagnetic** Elastic Form **Factors (EFFs)** $K_{\mu}F_M(Q^2) = N_c \operatorname{tr} \int_{\mathcal{A}^{h}} \chi_{\mu}(k+p_f,k+p_i) \Gamma_M(k_i;p_i) S(k) \gamma_M(k_f;-p_f)$ All can be written in terms of propagators and vertices

- Gives information on momentum/charge distribution.
- Pion EFF highly relevant for contemporary physics.

 $\Gamma_{\pi}(k_i;p_i)$ 

E

S(k)

E



#### **Elastic Form Factors**

- Clear probe of the hadron's structure.
  - Structure manifests
     in F(Q<sup>2</sup>) != constant
- Connected with the DA:



Factorization is a proof of the validity of QCD itself.



#### **Elastic Form Factors**

- Clear probe of the hadron's structure.
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 Factorization is a proof of the validity of QCD itself.

Testing scaling violations...



# Two-photon Transition Form Factors (TFFs)

$$T_{\mu\nu}(k_1, k_2) = \epsilon_{\mu\nu\alpha\beta} k_{1\alpha} k_{2\beta} G_{M_5}(k_1^2, k_1 \cdot k_2, k_2^2) ,$$
  
$$T_{\mu\nu}(k_1, k_2) = \operatorname{tr} \int \frac{d^4 l}{(2\pi)^4} i \mathcal{Q}\chi_{\mu}(l, l_1) \Gamma_{M_5}(l_1, l_2) S(l_2) i \mathcal{Q}\Gamma_{\nu}(l_2, l)$$

All can be expressed in terms of **propagators** and **vertices** 

- Gives information on momentum/charge distribution.
- **Pion TFF** highly relevant for contemporary physics.

 $S(l_2)$ 

 $(\mathcal{Q}\Gamma_{
u}(l_2,l))$ 



- Clear probe of the hadron's structure.
  - Structure manifests
     in G(Q<sup>2</sup>) != constant
- Connected with the PDA:



(at sufficiently large Q<sup>2</sup>)

Factorization is a proof of the validity of QCD itself.



• The CSM prediction satisfies the Abelian anomaly,  $2f_{\pi}^{0}G_{\pi^{0}}^{0}(Q^{2}=0)=1$ 

... while faithfully recovering the **asymptotic limit**.

• A dilated+concave DA, at the hadronic scale, connects both pion EFF and TFF.



Precise agreement with <u>all</u> experimental data; except for <u>Babar</u> at large Q<sup>2</sup>.





- > **All** two-photon **TFFs** involving ground-state neutral pseudoscalars are within reach:
  - Invariably, agreement with the experimental data is found, with the exception of the large-Q<sup>2</sup> Babar data for the pion.



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- > **All** two-photon **TFFs** involving ground-state neutral pseudoscalars are within reach:
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- Clearly, the shape of M(k) echoes in TFFs and DAs.





- Gives information on mass/pressure distribution.
- A first step towards **nucleon GFFs**

#### **Gravitational form factors**

The expectation value of the energy-momentum tensor (EMT) in the pion defines the gravitational form factors:

$$\Lambda_{\mu\nu}(P,Q) = 2P_{\mu}P_{\nu}\theta_2(Q^2) + \frac{1}{2}\left(Q^2g_{\mu\nu} - Q_{\mu}Q_{\nu}\right)\theta_1(Q^2) + 2m_{\pi}^2g_{\mu\nu}\bar{c}(Q^2)$$

Where symmetry principles entail:



Momentum conservation



Soft-pion theorem



EMT conservation

Polyakov:2018zvc

(Other pseudoscalars are defined in analogy)

> The matrix element can be expressed in terms of **propagators** and **vertices**:

$$\Lambda_{\mu\nu}(P,Q) = N_c \int_{dk} \operatorname{Tr} \left[ \Gamma_{\pi} \left( k - \frac{Q}{4}, P - \frac{Q}{2} \right) S\left( k - \frac{P}{2} \right) \Gamma_{\pi} \left( k + \frac{Q}{4}, P + \frac{Q}{2} \right) + \text{beyond I.A.} \right]$$
$$S\left( k + \frac{P}{2} + \frac{Q}{2} \right) \Gamma_{\mu\nu} \left( k + \frac{P}{2}, Q \right) S\left( k + \frac{P}{2} - \frac{Q}{2} \right) \right]$$

Quark-tensor vertex (QTV): Interaction with a spin-2 probe

Partially constrained by its Ward-Green-Takahashi identity:

$$iQ_{\mu}\Gamma^{\mu\nu}(P,Q) = P_i^{\nu}S^{-1}(P_f) - P_f^{\nu}S^{-1}(P_i)$$

Those pieces escaping the WGTI, encode scalar and tensor meson poles.

## **π-K GFFs**

> It is found that:  $r_{\theta_1}^{\pi}$  (pressure) >  $r_F^{\pi}$  (charge) >  $r_{\theta_2}^{\pi}$  (mass)



## **π-K GFFs**



#### **π-K: Pressure profiles**

$$p_{K}^{u}(r) = \frac{1}{6\pi^{2}r} \int_{0}^{\infty} d\Delta \frac{\Delta}{2E(\Delta)} \sin(\Delta r) [\Delta^{2}\theta_{1}^{K_{u}}(\Delta^{2})],$$

$$s_{K}^{u}(r) = \frac{3}{8\pi^{2}} \int_{0}^{\infty} d\Delta \frac{\Delta^{2}}{2E(\Delta)} j_{2}(\Delta r) [\Delta^{2}\theta_{1}^{K_{u}}(\Delta^{2})],$$
"Pressure" Quark attraction/repulsion  
**CONFINEMENT**  
"Shear" Deformation QCD forces  

$$\int_{0.04}^{0.05} \int_{0.04}^{-\frac{-\overline{s}-in-K}{2E(\Delta)}} \frac{-\frac{-\overline{s}-in-K}{2E(\Delta)}}{r_{c}} = 0.39(1) \text{ fm}} r_{c}^{T} = 0.39(1) \text{ fm}} r_{c}^{T} = 0.26(1) \text{ fm}} \int_{0.05}^{0.04} \int_{0.05}^{0.05} \int_{0.00}^{0.05} \int_{0.00}^{0.$$

# Light-front wave functions (LFWF)

 $\psi_{\mathrm{M}}^{q}\left(x,k_{\perp}^{2}\right) = \mathrm{tr} \int_{dk_{\parallel}} \delta_{n}^{x}(k_{\mathrm{M}})\gamma_{5}\gamma \cdot n \chi_{\mathrm{M}}(k_{-},P)$ 



"One ring to rule them all"

Bethe-Salpeter wave function

- Intrinsic of the hadron's nature.
- Yields a variety of distributions.

#### The idea: Connect everything through the LFWF.



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#### **π-K: LFWFs & GPDs**



$$H_{\rm M}^q(x,\xi,t) = \int \frac{d^2 k_{\perp}}{16\pi^3} \psi_{\rm M}^{q*} \left(x^-, (\mathbf{k}_{\perp}^-)^2\right) \psi_{\rm M}^q \left(x^+, (\mathbf{k}_{\perp}^+)^2\right)$$





Likelihood of finding a valence-quark with momentum fraction x, at position b.

#### **Evolved IPS-GPD: Pion Case**

3.0

2.5 2.0 15 10

0.5

$$u^{\mathsf{P}}(x, b_{\perp}^{2}; \zeta_{\mathcal{H}}) = \int_{0}^{\infty} \frac{d\Delta}{2\pi} \Delta J_{0}(b_{\perp}\Delta) H^{u}_{\mathsf{P}}(x, 0, -\Delta^{2}; \zeta_{\mathcal{H}})$$



 Likelihood of finding a parton with LF momentum **x** at transverse position **b** 



#### **Evolved IPS-GPD: Kaon Case**



#### **Distributions: Charge & Mass**

$$\rho_{\rm P}(b) = \frac{1}{2\pi} \int_0^\infty d\Delta \,\Delta J_0(\Delta \, b) F_{\rm P}(\Delta^2)$$

 $F_{\rm P}^E(\Delta^2) \to \rho_{\rm P}^E(b)$ 

- Intuitively, we expect the meson to be localized at a finite space.
- **Charge** effect span over a **larger domain** than **mass** effects. ۶

More **massive** hadron More compressed



## **GPDs: Empirical determination**

#### **Pion GPD: Empirical determination**



#### **Pion GPD: Empirical determination**

- The connection of GPDs with <u>PDFs and EFFs</u> enable us to use existing data on those quantities to reconstruct the pion GPD.
- Using a chi^2-based probabilistic selection procedure, an ensemble of representations for the pion GPD is generated.
- The produced ensemble turns out to be in agreement with previous CSM predictions.



- Proving, once again, that  $\theta_2$  is harder than the EFE:
  - i.e. the **mass distribution** is **more compact** than the **charge** one.

The physical boundaries:

$$\frac{1}{\sqrt{2}} \le r_{\pi}^{\theta_2}/r_{\pi} \le 1$$

Xu, KR *et al.* Chin.Phys.Lett. 40 (2023) 4, 041201



## **Final Highlights**



# **Final Highlights**

- The emergent phenomena in QCD produces unique outcomes:
  - The degrees-of-freedom are not directly accessible, we get to observe hadrons (confinement).
  - Through their own mechanisms, **dynamical mass generation** is present in both **matter** and **gauge** sectors of QCD; the later yielding a running **coupling** that saturates at infrared momenta.
- > **Pseudoscalar** mesons are an ideal platform to inquire on these facets of **QCD**:
  - Their mere existence and properties are connected with the mass generation in the Standard Model and, potentially, confinement.
     A x = 0.0
  - Modern facilities are capable to address the properties of NG bosons and it's connection with QCD's emergent phenomena.
- Theory has evolved to the point where all sorts of parton distributions within <u>pseudoscalar</u> mesons are within reach.

→ Many of them connected via LFWF

