

Muon Particle Physics Experiments

Yoshitaka Kuno
Osaka University, Japan

November 4th, 2023

MELODY2023

Chinese Spallation Neutron Source (CSNS),
Dongguan, China

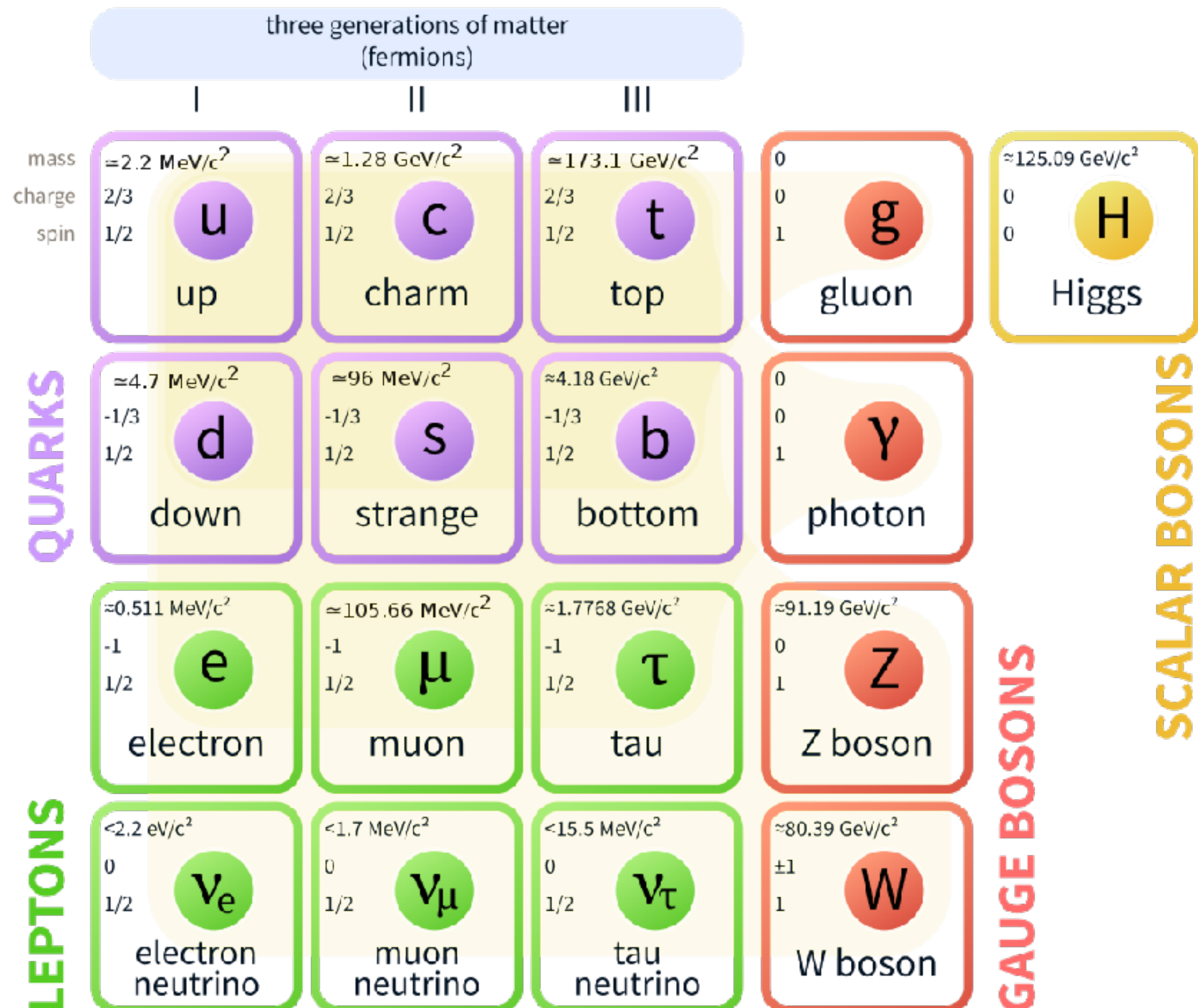


The Standard Model and Muons



The Standard Model and Muons

Standard Model of Elementary Particles

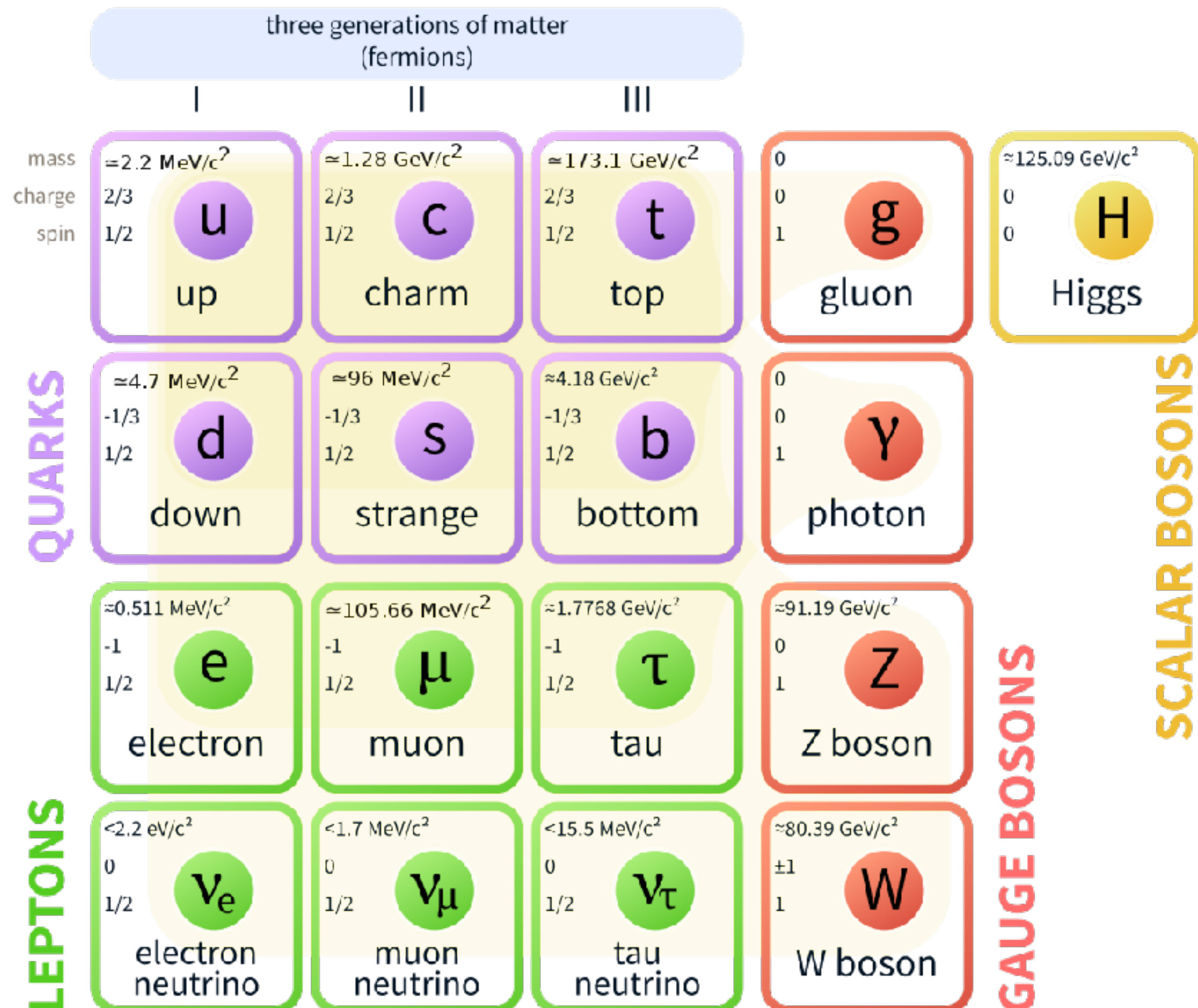


The Standard Model is considered to be incomplete.

example: origin of flavour
neutrino masses
dark matter,
baryogenesis,
dark energy,
strong CP and more.

The Standard Model and Muons

Standard Model of Elementary Particles



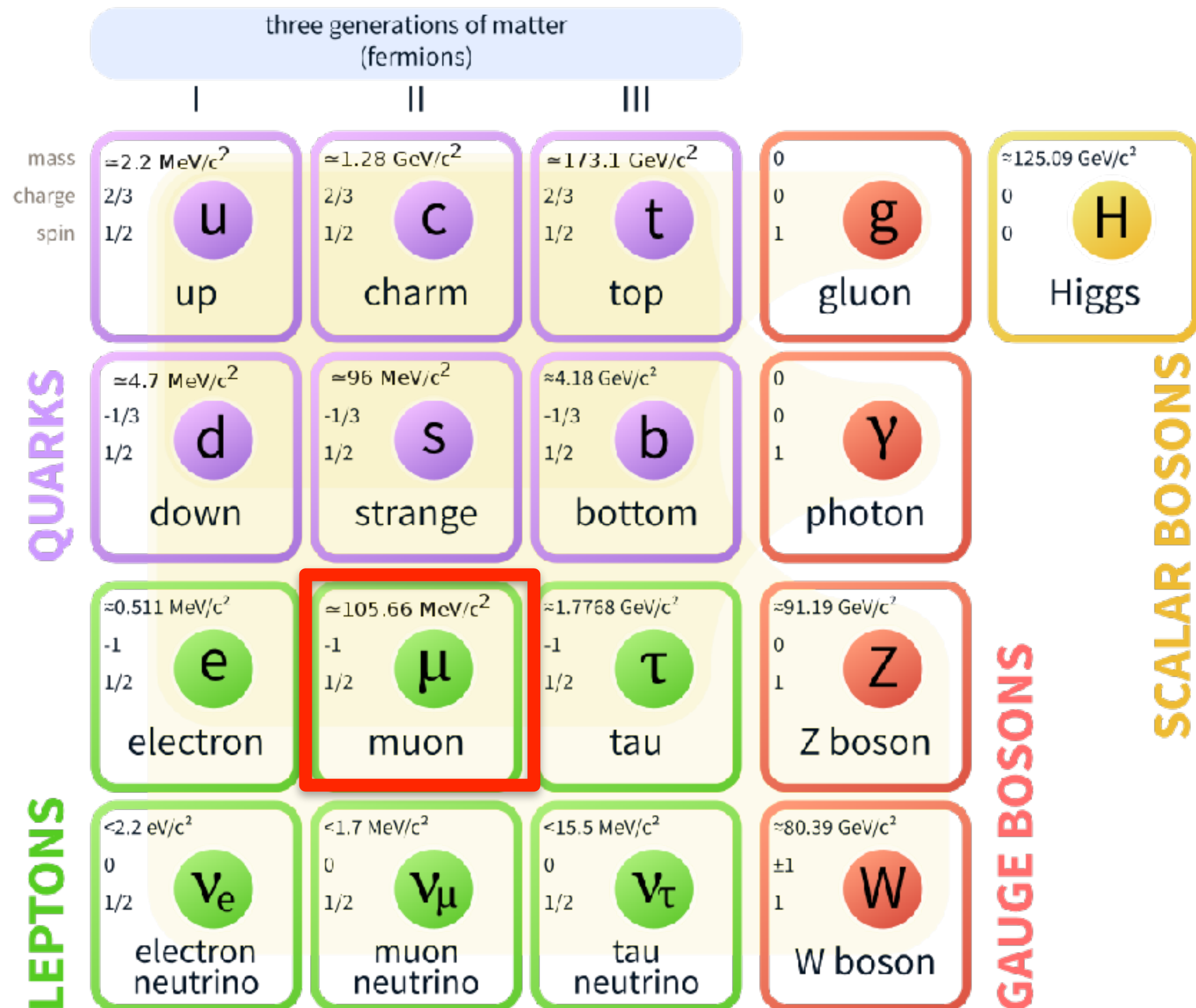
The Standard Model is considered to be incomplete.

example: origin of flavour
neutrino masses
dark matter,
baryogenesis,
dark energy,
strong CP and more.

New physics beyond the Standard Model (BSM) is needed.

The Standard Model and Muons

Standard Model of Elementary Particles



The Standard Model is considered to be incomplete.

example: origin of flavour
neutrino masses
dark matter,
baryogenesis,
dark energy,
strong CP and more.

New physics beyond the Standard Model (BSM) is needed.

→ Muons

Why Muons ?



Why Muons ?



Muons are many.

Why Muons ?



Muons are many.



Outline



- Why New physics needed in Particle Physics ?
- Muon's dipole moments
- Muon's charged lepton flavour violation (CLFV)
 - $\mu^+ \rightarrow e^+ \gamma, \mu^+ \rightarrow e^+ e^+ e^-, \mu^- \rightarrow e^-$ conversion
 - in the case of discoveries
- Muon's lepton number violation (LNV) and axion-like particles
 - $\mu^- \rightarrow e^+$ conversion, $\mu \rightarrow ea$
- Muonium to anti-muonium conversion
- Muon Collider
- Summary

Muon's Dipole Moments



Muon $g-2$ and Muon edm



Muon g-2 and Muon edm

Muon magnetic dipole moment (g-2)

- Charged, spin=1/2 particles have magnetic dipole moments.
- $g=2$ for a “pure” Dirac particle, and $g>2$ due to quantum radiative corrections.
- Deviation of the measurement from the Standard Model prediction indicates new physics beyond the SM.

$$\vec{\mu} = g \frac{q}{2m} \vec{S} \quad g = 2 \left(1 + \frac{\alpha}{2\pi} \dots \right)$$

Muon g-2 and Muon edm

Muon magnetic dipole moment (g-2)

- Charged, spin=1/2 particles have magnetic dipole moments.
- $g=2$ for a “pure” Dirac particle, and $g>2$ due to quantum radiative corrections.
- Deviation of the measurement from the Standard Model prediction indicates new physics beyond the SM.

$$\vec{\mu} = g \frac{q}{2m} \vec{S} \quad g = 2 \left(1 + \frac{\alpha}{2\pi} \dots \right)$$

See the talk by Liang Li
and Kim Siang Khaw

Muon g-2 and Muon edm

Muon magnetic dipole moment (g-2)

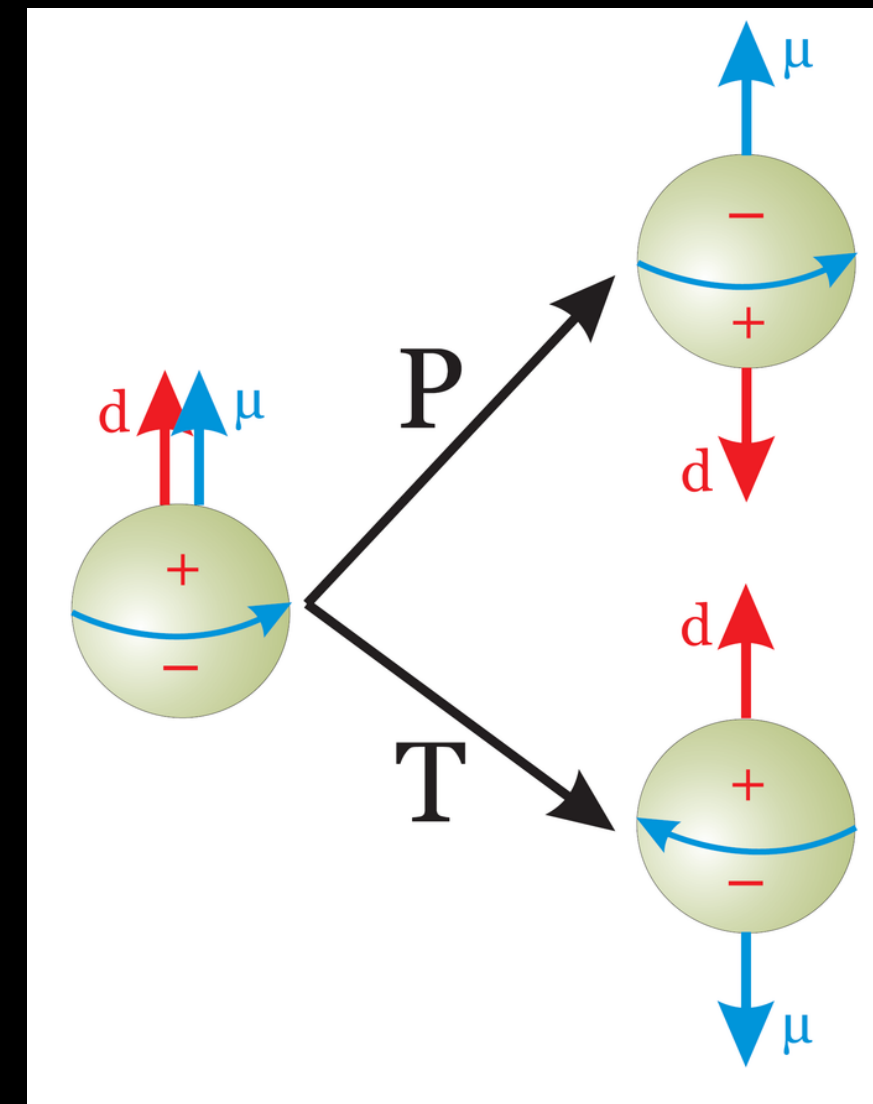
- Charged, spin=1/2 particles have magnetic dipole moments.
- $g=2$ for a “pure” Dirac particle, and $g>2$ due to quantum radiative corrections.
- Deviation of the measurement from the Standard Model prediction indicates new physics beyond the SM.

$$\vec{\mu} = g \frac{q}{2m} \vec{S} \quad g = 2 \left(1 + \frac{\alpha}{2\pi} \dots \right)$$

See the talk by Liang Li and Kim Siang Khaw

Muon electric dipole moment (edm)

- Permanent edm of an elementary particle violates T and P reversal invariance.
- The SM contribution is very small.
 $d_{\mu} \sim 10^{-(36-38)} e \cdot cm$
- Observation of edm indicates new physics beyond the SM.
- The current experimental limits $d_{\mu} < 1.5 \times 10^{-19} e \cdot cm$



Muon g-2 and Muon edm

Muon magnetic dipole moment (g-2)

- Charged, spin=1/2 particles have magnetic dipole moments.
- $g=2$ for a “pure” Dirac particle, and $g>2$ due to quantum radiative corrections.
- Deviation of the measurement from the Standard Model prediction indicates new physics beyond the SM.

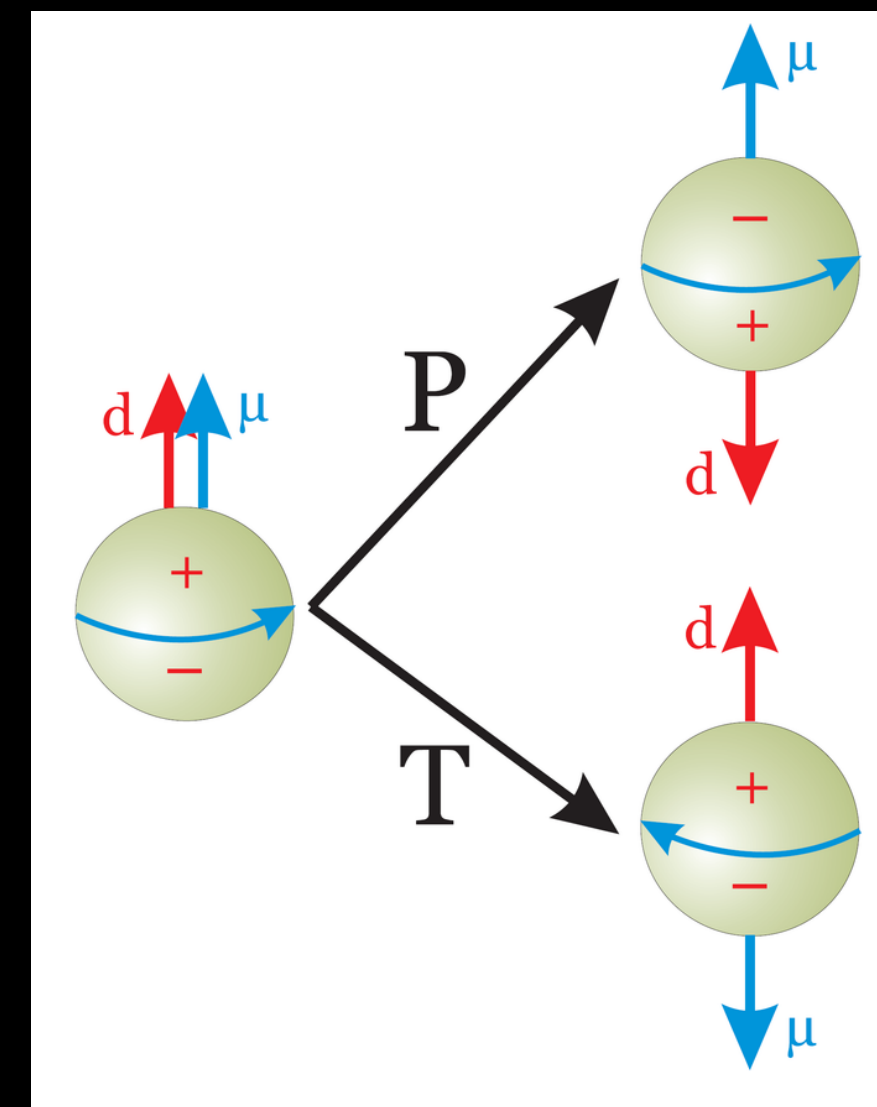
$$\vec{\mu} = g \frac{q}{2m} \vec{S} \quad g = 2 \left(1 + \frac{\alpha}{2\pi} \dots \right)$$

See the talk by Liang Li and Kim Siang Khaw

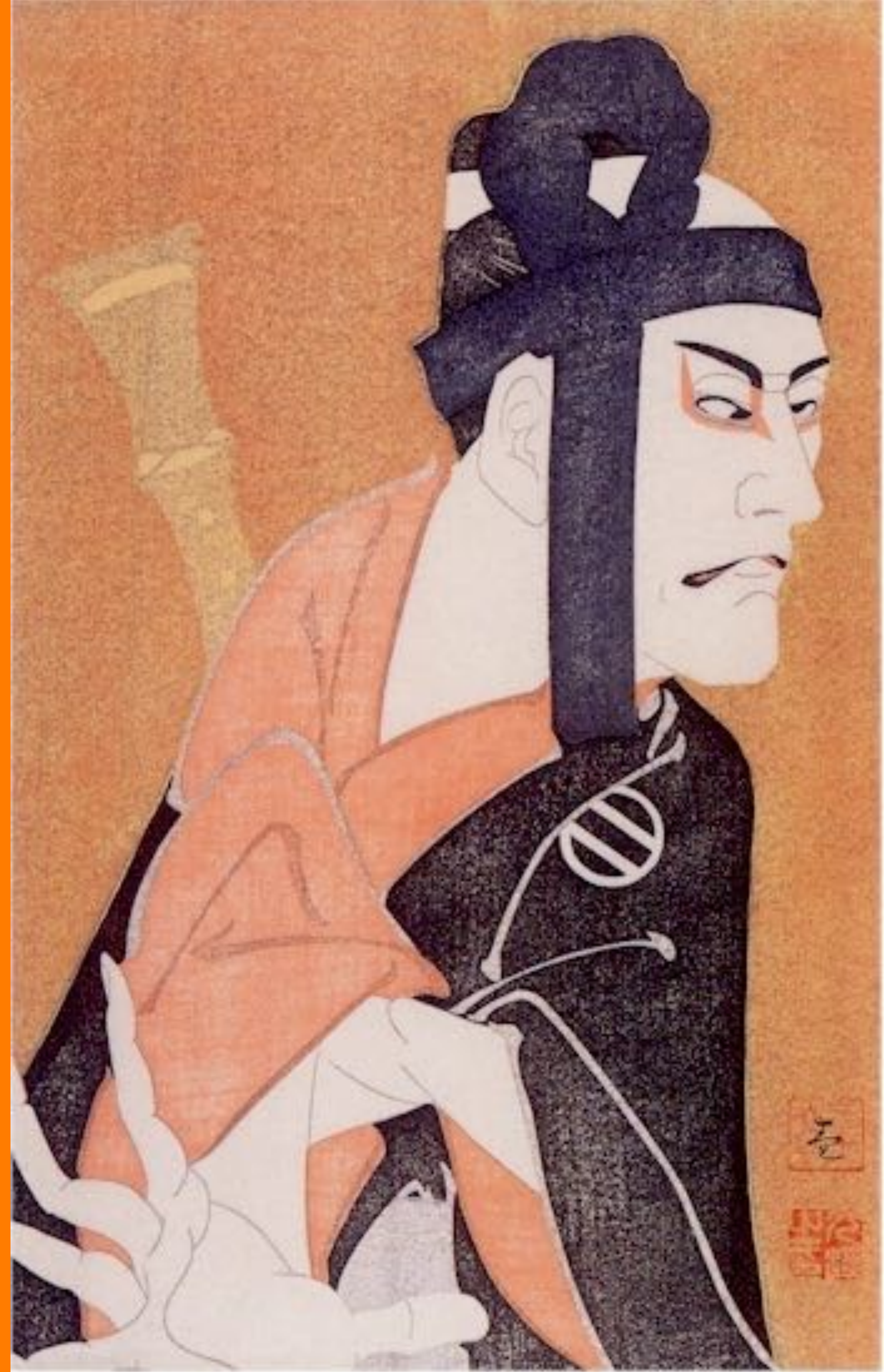
Muon electric dipole moment (edm)

- Permanent edm of an elementary particle violates T and P reversal invariance.
- The SM contribution is very small.
 $d_{\mu} \sim 10^{-(36-38)} e \cdot cm$
- Observation of edm indicates new physics beyond the SM.
- The current experimental limits $d_{\mu} < 1.5 \times 10^{-19} e \cdot cm$

See the talk by Kim Siang Khaw



Charged Lepton Flavor Violation with Muons



Lepton Flavours



Lepton Flavours



Lepton flavour quantum
numbers

Lepton Flavours

Lepton flavour quantum numbers

	electron flavour	muon flavour	tau flavour
e generation	1	0	0
μ generation	0	1	0
τ generation	0	0	1

Lepton Flavours

Lepton flavour quantum numbers

	electron flavour	muon flavour	tau flavour
e generation	1	0	0
μ generation	0	1	0
τ generation	0	0	1

In the (old) Standard Model, each of lepton flavour is additively conserved separately. Neutrino oscillation indicated violation of lepton flavour conservation.

How about charged lepton flavour violation (CLFV) like $\mu \rightarrow e\gamma$?

Lepton Flavours

Lepton flavour quantum numbers

	electron flavour	muon flavour	tau flavour
e generation	1	0	0
μ generation	0	1	0
τ generation	0	0	1

In the (old) Standard Model, each of lepton flavour is additively conserved separately. Neutrino oscillation indicated violation of lepton flavour conservation.

How about charged lepton flavour violation (CLFV) like $\mu \rightarrow e\gamma$?

CLFV for BSM search

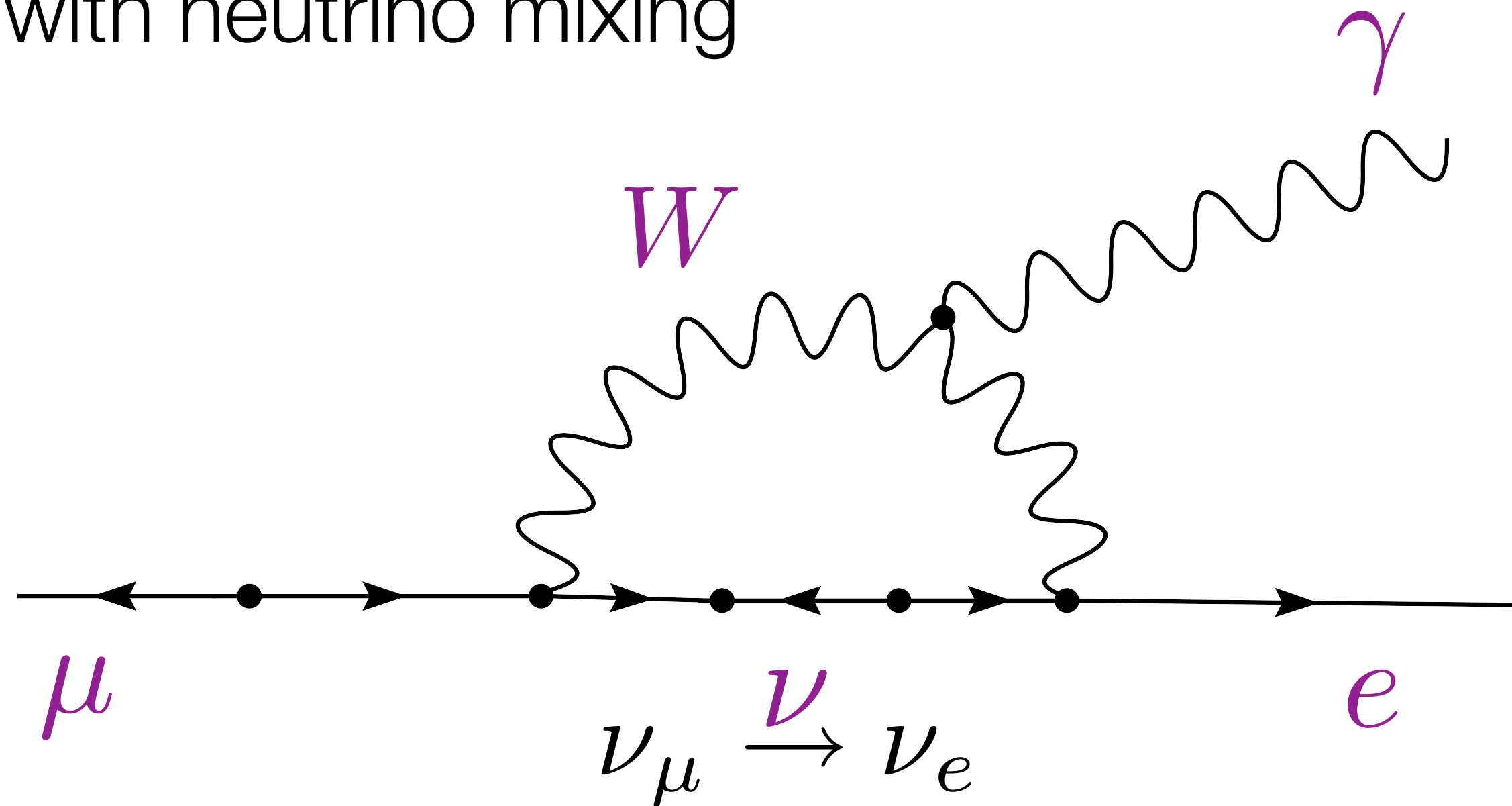
Two Things you must know on CLFV No.1



Two Things you must know on CLFV No.1

Standard Model Contribution to CLFV

with neutrino mixing

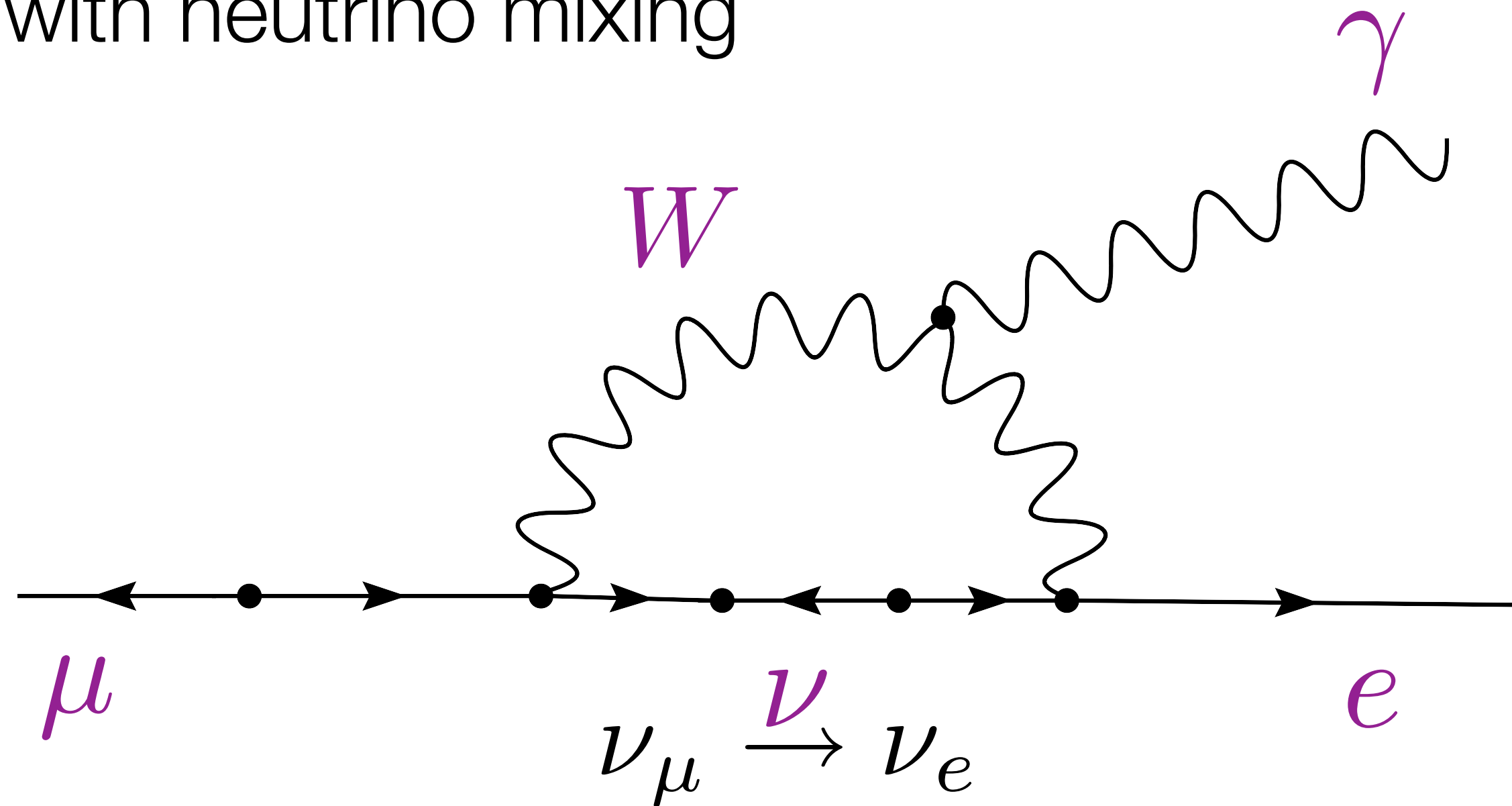


- S.T. Petcov, Sov.J. Nucl. Phys. 25 (1977) 340
- W.J. Marciano et al., Phys. Lett. B 67 (1977) 303
- B.W. Lee, et al., Phys. Rev. Lett. 38 (1977) 937
- B.W. Lee et al., Phys. Rev. D 16 (1977) 1444.

Two Things you must know on CLFV No.1

Standard Model Contribution to CLFV

with neutrino mixing

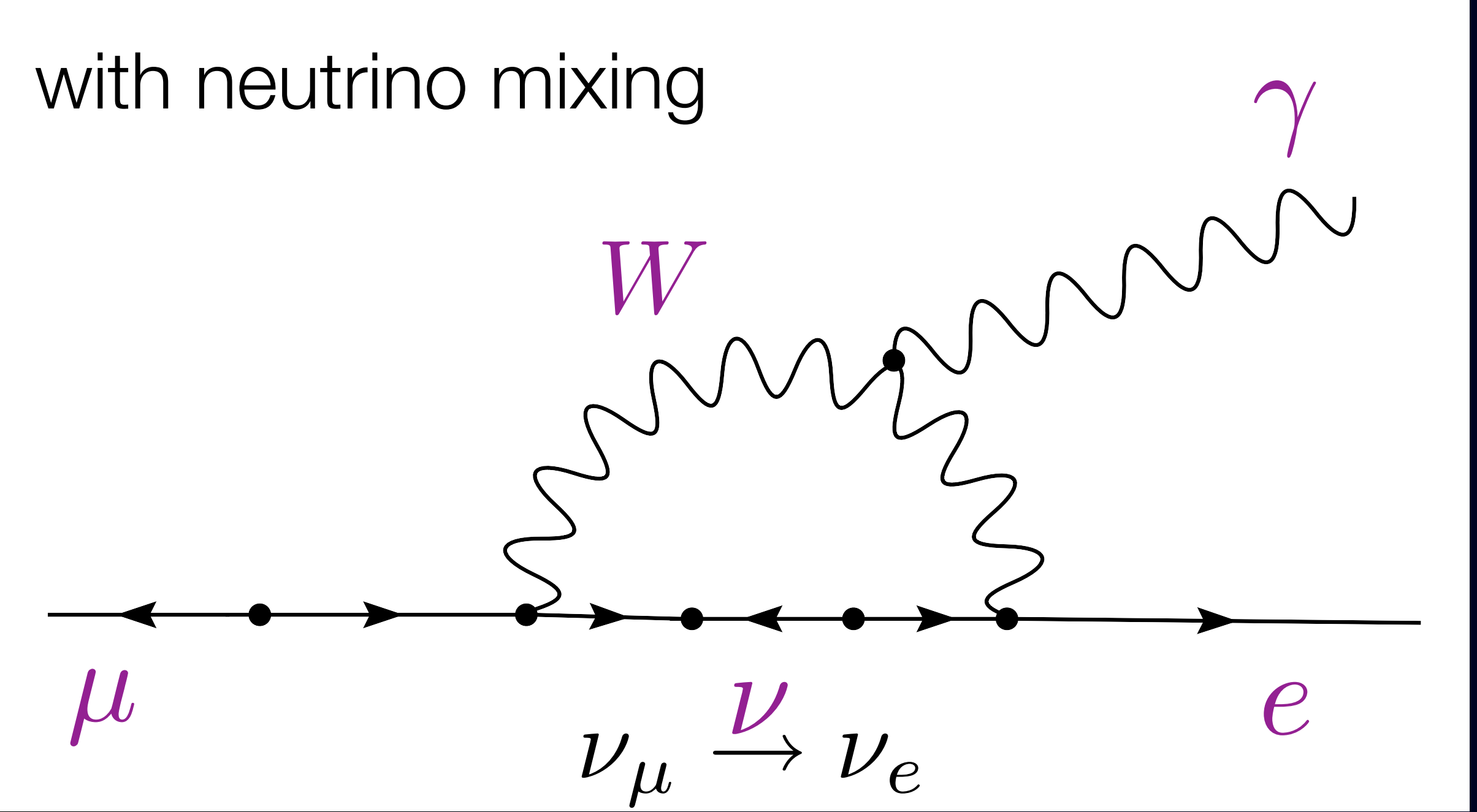


$$B(\mu \rightarrow e\gamma) = \frac{3\alpha}{32\pi} \left| \sum_l (V_{MNS})_{\mu l}^* (V_{MNS})_{el} \frac{m_{\nu_l}^2}{M_W^2} \right|^2$$

- S.T. Petcov, Sov.J. Nucl. Phys. 25 (1977) 340
- W.J. Marciano et al., Phys. Lett. B 67 (1977) 303
- B.W. Lee, et al., Phys. Rev. Lett. 38 (1977) 937
- B.W. Lee et al., Phys. Rev. D 16 (1977) 1444.

Two Things you must know on CLFV No.1

Standard Model Contribution to CLFV

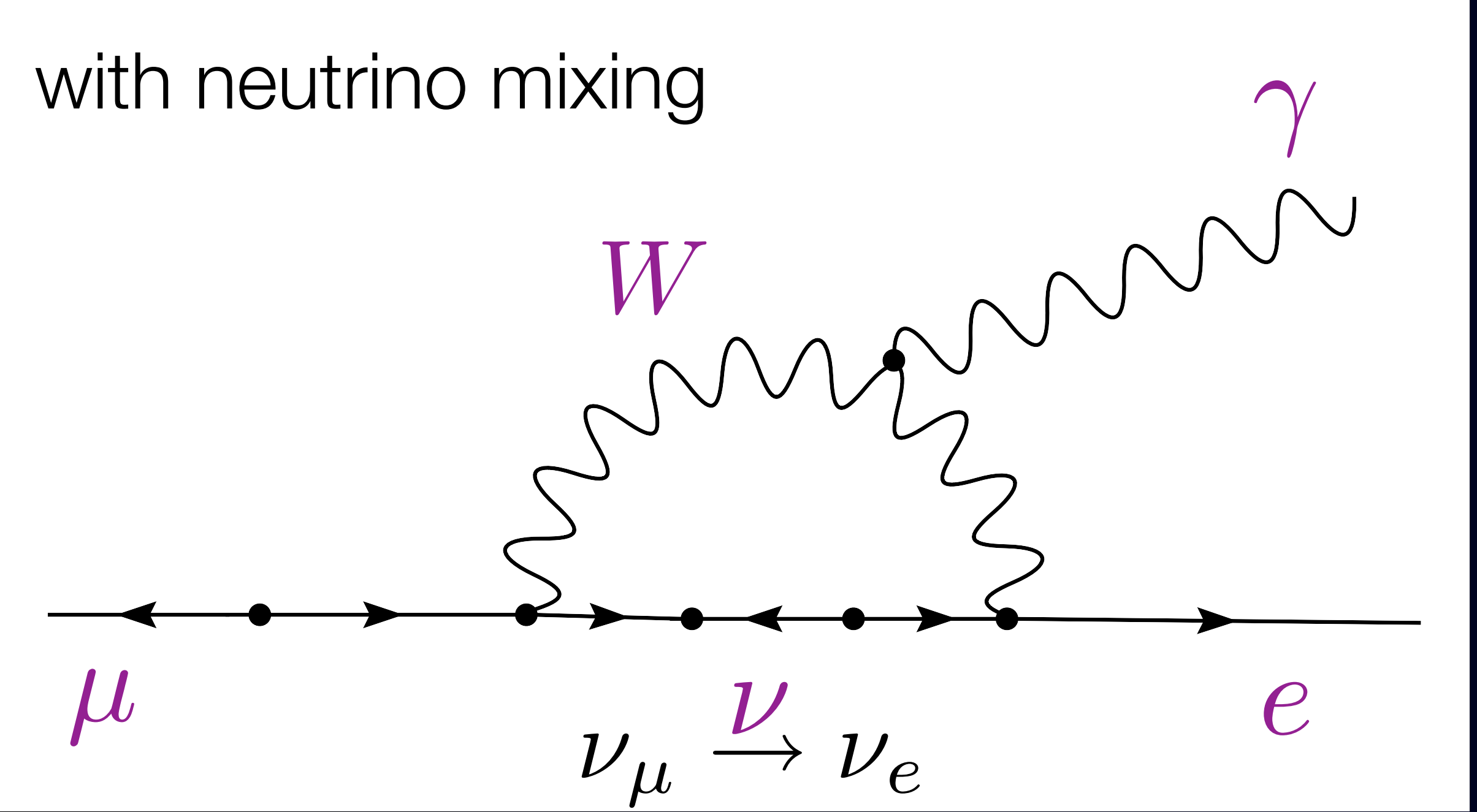


$$B(\mu \rightarrow e\gamma) = \frac{3\alpha}{32\pi} \left| \sum_l (V_{MNS})_{\mu l}^* (V_{MNS})_{el} \frac{m_{\nu_l}^2}{M_W^2} \right|^2$$

S.T. Petcov, Sov.J. Nucl. Phys. 25 (1977) 340
 W.J. Marciano et al., Phys. Lett. B 67 (1977) 303
 B.W. Lee, et al., Phys. Rev. Lett. 38 (1977) 937
 B.W. Lee et al., Phys. Rev. D 16 (1977) 1444.

Two Things you must know on CLFV No.1

Standard Model Contribution to CLFV



$$B(\mu \rightarrow e\gamma) = \frac{3\alpha}{32\pi} \left| \sum_l (V_{MNS})_{\mu l}^* (V_{MNS})_{el} \frac{m_{\nu_l}^2}{M_W^2} \right|^2$$

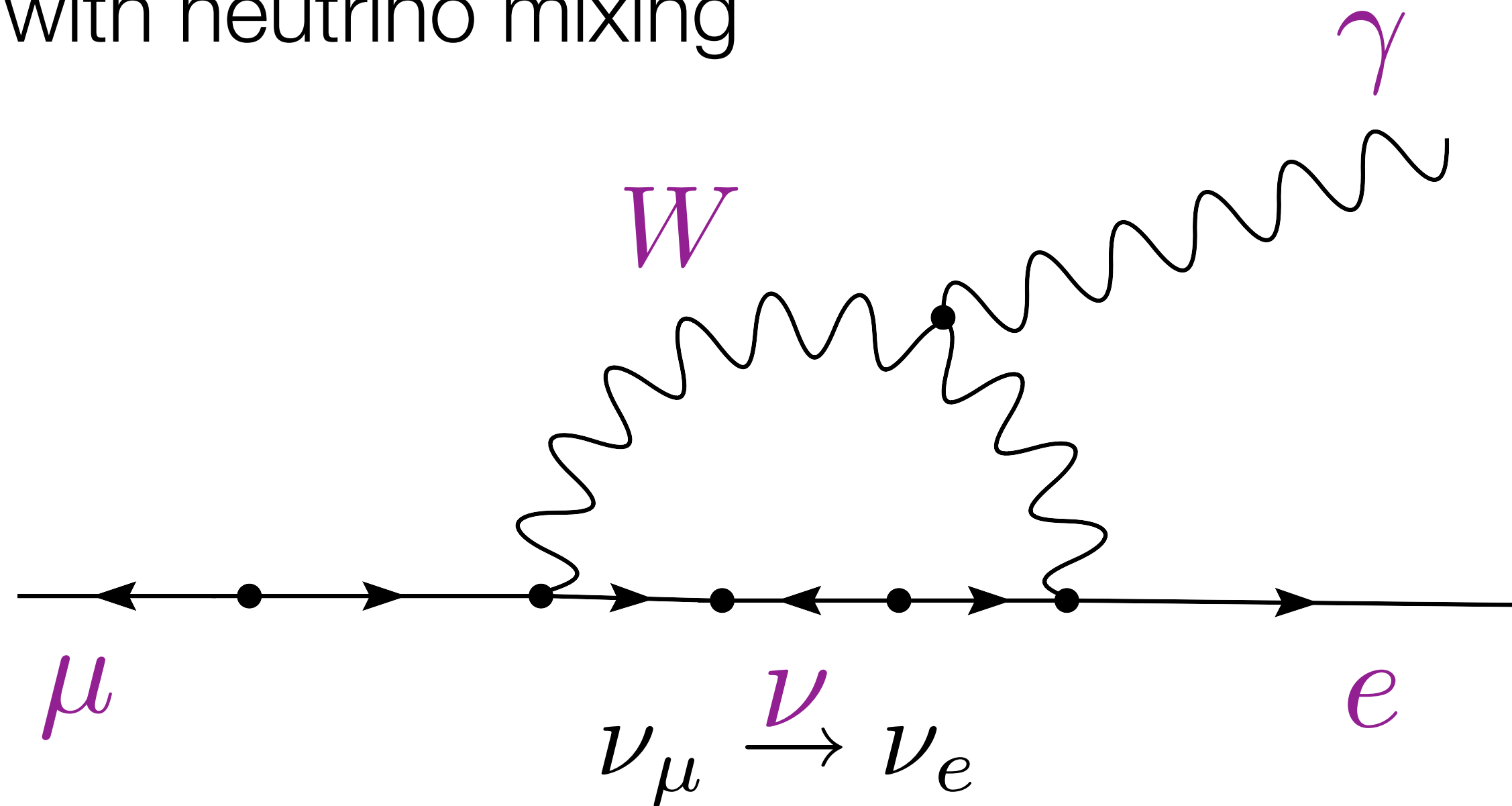
BR ~ O(10⁻⁵⁴)

S.T. Petcov, Sov.J. Nucl. Phys. 25 (1977) 340
 W.J. Marciano et al., Phys. Lett. B 67 (1977) 303
 B.W. Lee, et al., Phys. Rev. Lett. 38 (1977) 937
 B.W. Lee et al., Phys. Rev. D 16 (1977) 1444.

Two Things you must know on CLFV No.1

Standard Model Contribution to CLFV

with neutrino mixing



$$B(\mu \rightarrow e\gamma) = \frac{3\alpha}{32\pi} \left| \sum_l (V_{MNS})_{\mu l}^* (V_{MNS})_{el} \frac{m_{\nu_l}^2}{M_W^2} \right|^2$$

BR ~ O(10⁻⁵⁴)

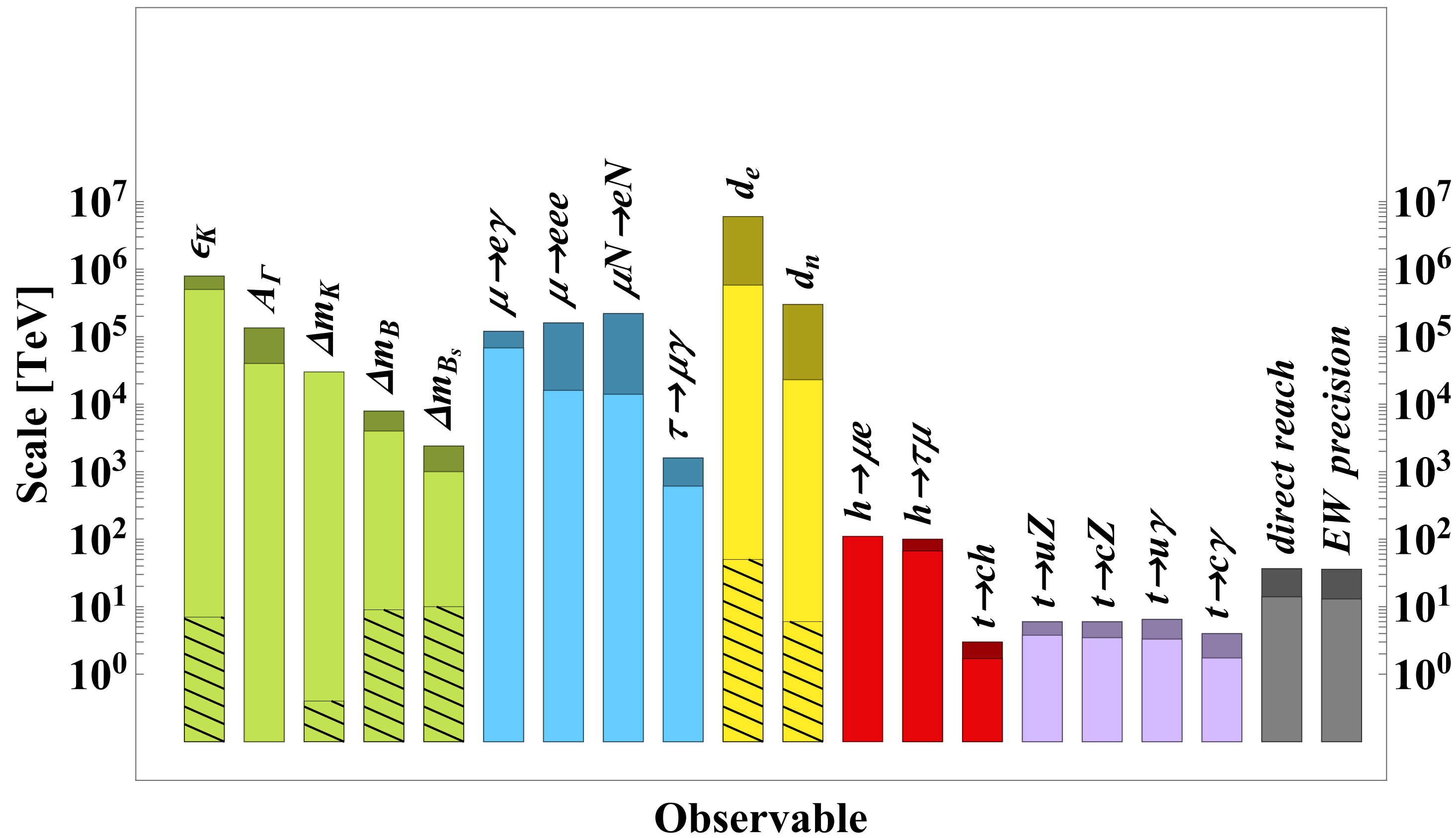
Searches for CLFV have clear signatures of new physics beyond the Standard Model **without the SM backgrounds** of massive neutrinos.

- S.T. Petcov, Sov.J. Nucl. Phys. 25 (1977) 340
- W.J. Marciano et al., Phys. Lett. B 67 (1977) 303
- B.W. Lee, et al., Phys. Rev. Lett. 38 (1977) 937
- B.W. Lee et al., Phys. Rev. D 16 (1977) 1444.

Two Things You must know on CLFV No.2



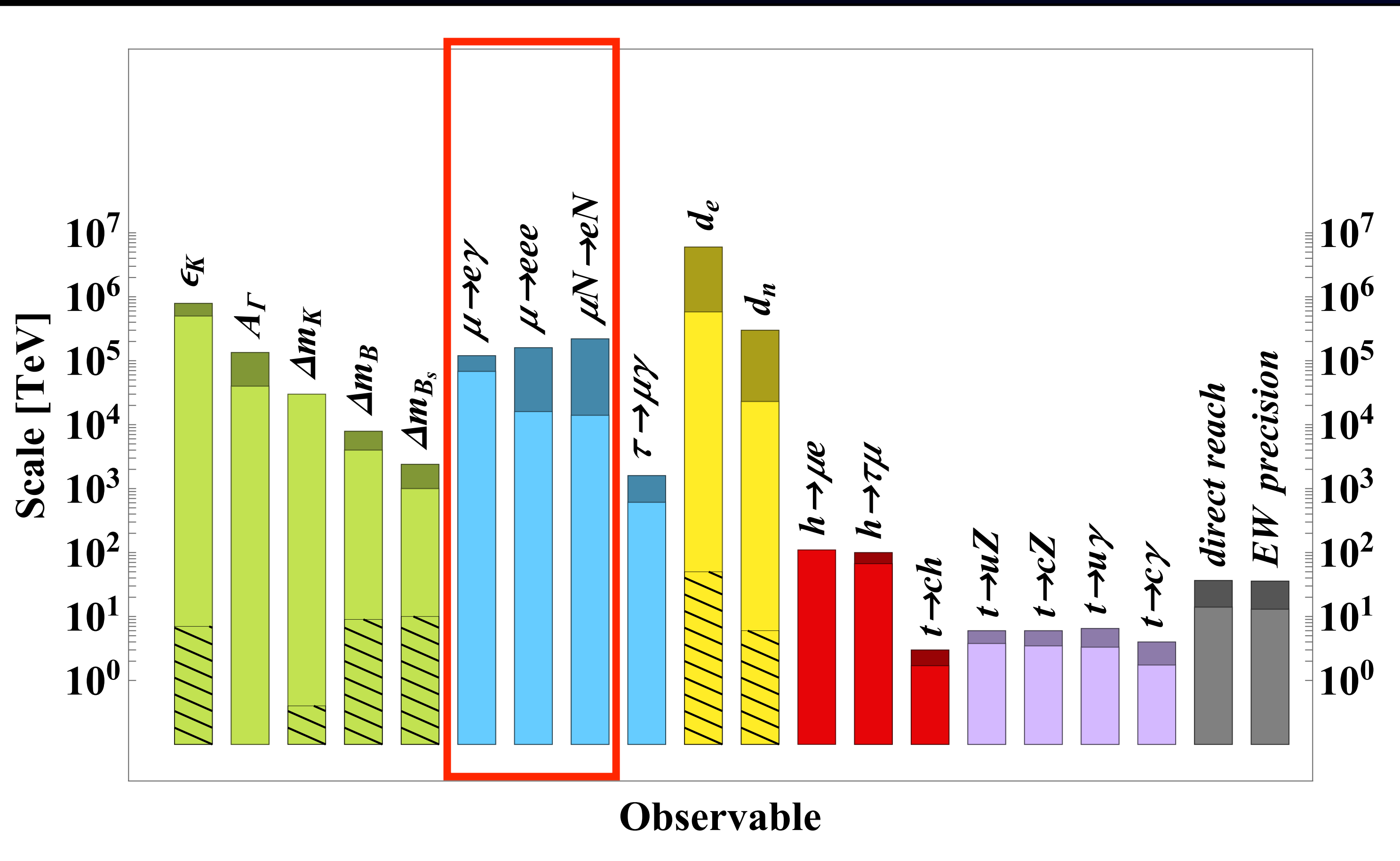
Two Things You must know on CLFV No.2



light colour: present

dark colour: future prospect

Two Things You must know on CLFV No.2

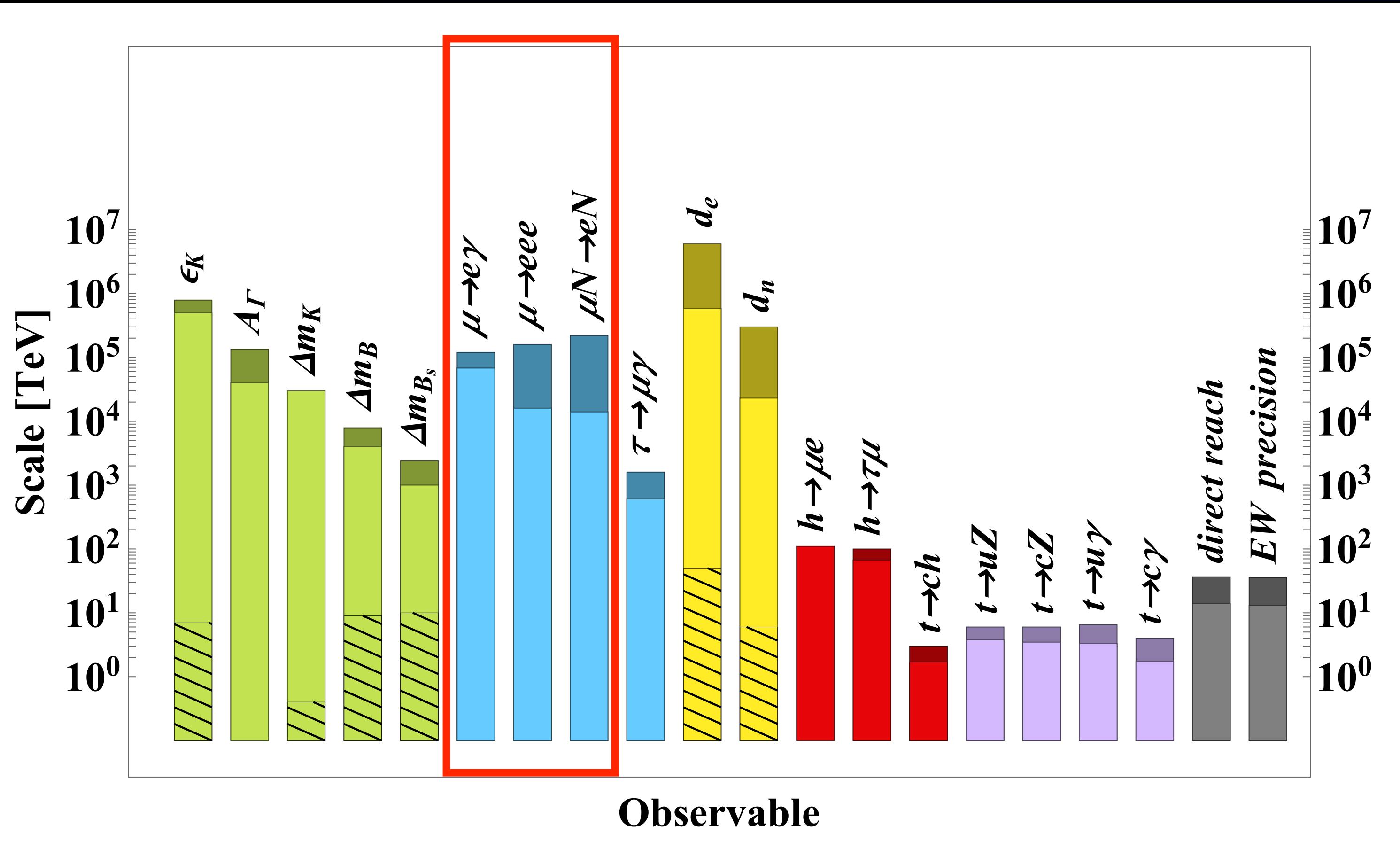


light colour: present
 dark colour: future prospect

Two Things You must know on CLFV No.2

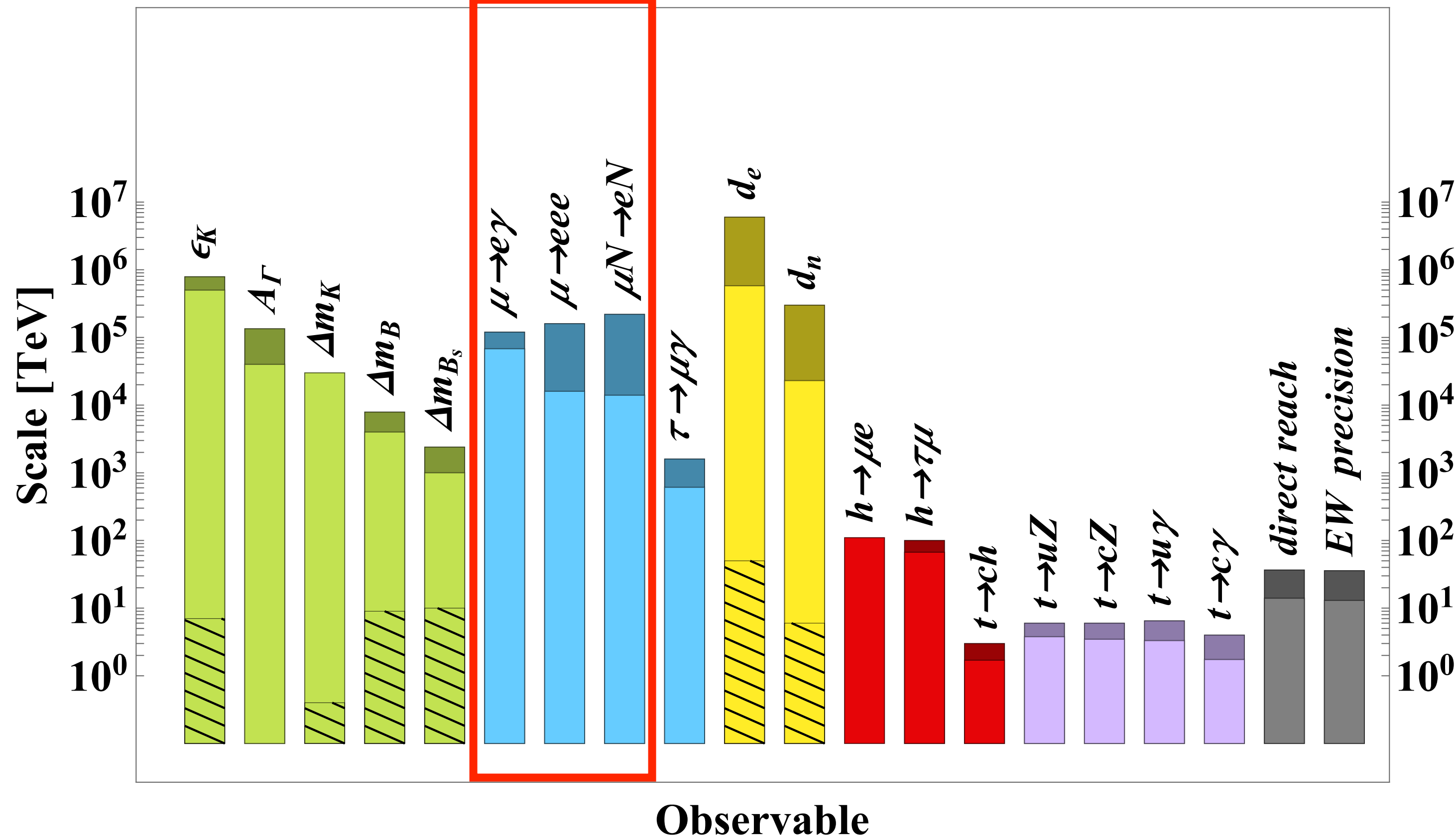
Present CLFV physics scales

Physics scale: $\Lambda = \mathcal{O}(10^3 - 10^4)$ TeV



light colour: present
dark colour: future prospect

Two Things You must know on CLFV No.2



Present CLFV physics scales

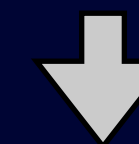
Physics scale: $\Lambda = \mathcal{O}(10^3 - 10^4)$ TeV

Prospect from present to future

x10 in physics scale

$$\text{SM forbidden rate} \propto \frac{c^2}{\Lambda^4}$$

in dimension 6 operators

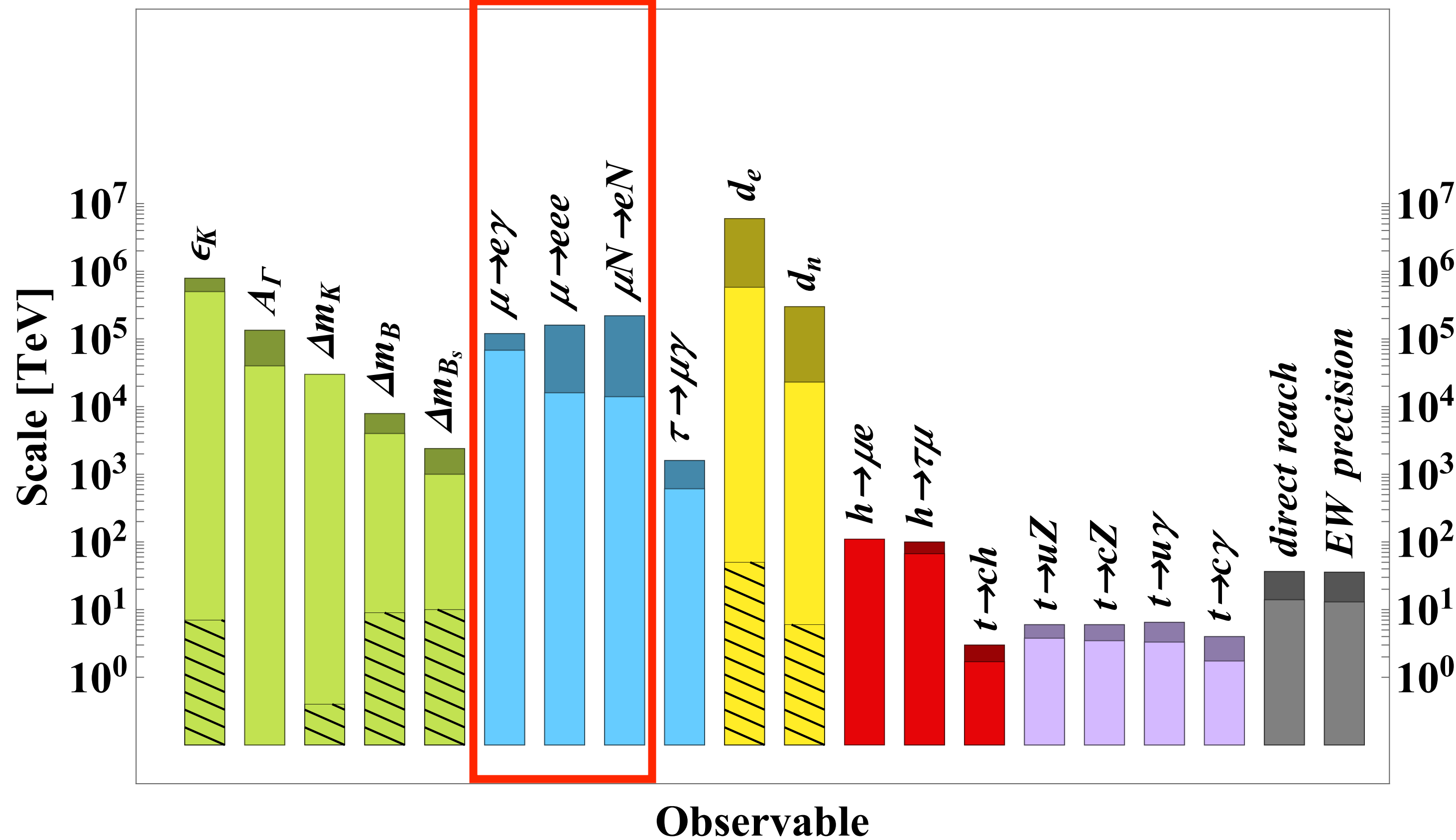


x10000 in experimental sensitivity

light colour: present

dark colour: future prospect

Two Things You must know on CLFV No.2



light colour: present
dark colour: future prospect

Present CLFV physics scales

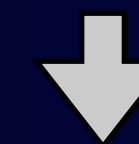
Physics scale: $\Lambda = \mathcal{O}(10^3 - 10^4)$ TeV

Prospect from present to future

x10 in physics scale

$$\text{SM forbidden rate} \propto \frac{c^2}{\Lambda^4}$$

in dimension 6 operators



x10000 in experimental sensitivity

CLFV probes very high energy scale of new physics.

$\mu \rightarrow e$ CLFV Golden Processes and EFT



$\mu \rightarrow e$ CLFV Golden Processes and EFT

$$\mu^+ \rightarrow e^+ \gamma$$

$$\mu^+ \rightarrow e^+ e^+ e^-$$

$$\mu^- N \rightarrow e^- N$$

$\mu \rightarrow e$ CLFV Golden Processes and EFT

dipole interaction

contact interaction

Effective Field Theory (EFT)

$$\mu^+ \rightarrow e^+ \gamma$$

$$\mu^+ \rightarrow e^+ e^+ e^-$$

$$\mu^- N \rightarrow e^- N$$

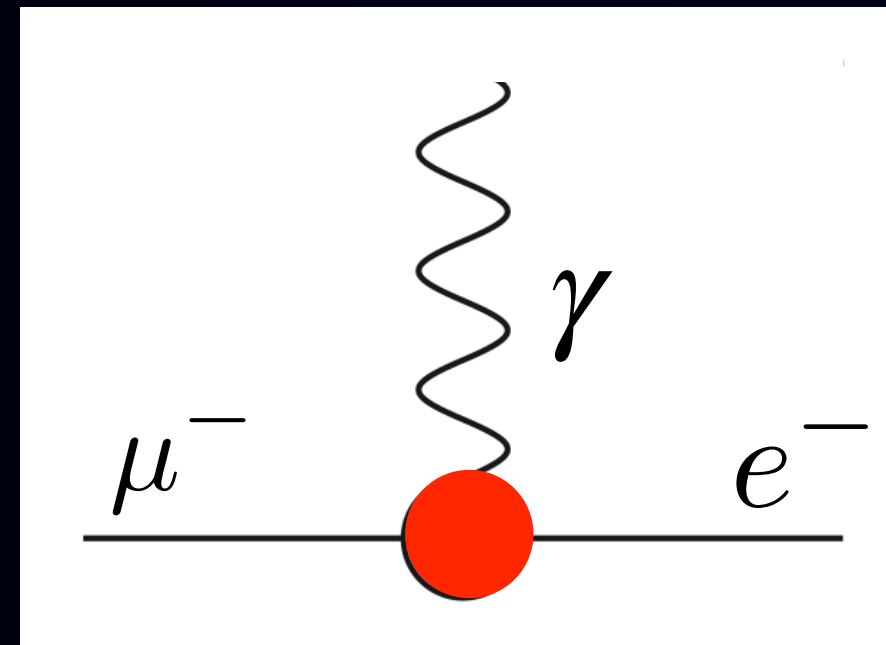
$\mu \rightarrow e$ CLFV Golden Processes and EFT

dipole interaction

contact interaction

Effective Field Theory (EFT)

$$\mu^+ \rightarrow e^+ \gamma$$



$$\mu^+ \rightarrow e^+ e^+ e^-$$

$$\mu^- N \rightarrow e^- N$$

- dipole operators (L/R)

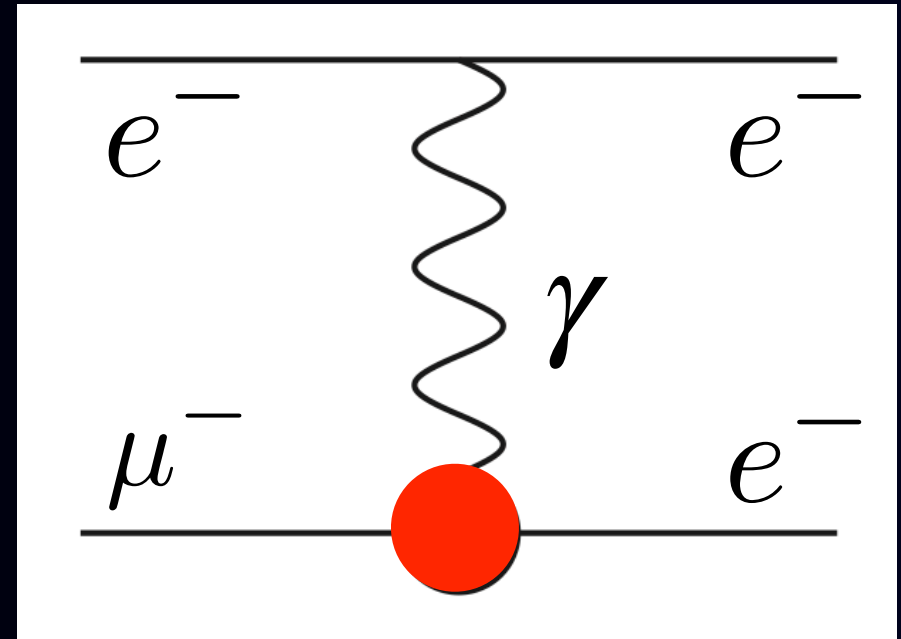
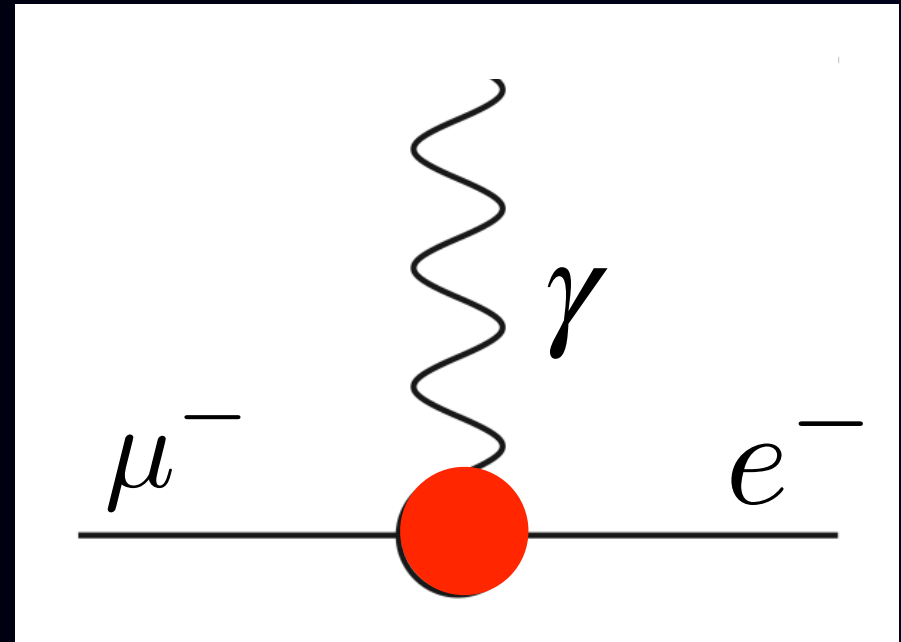
$\mu \rightarrow e$ CLFV Golden Processes and EFT

$$\mu^+ \rightarrow e^+ \gamma$$

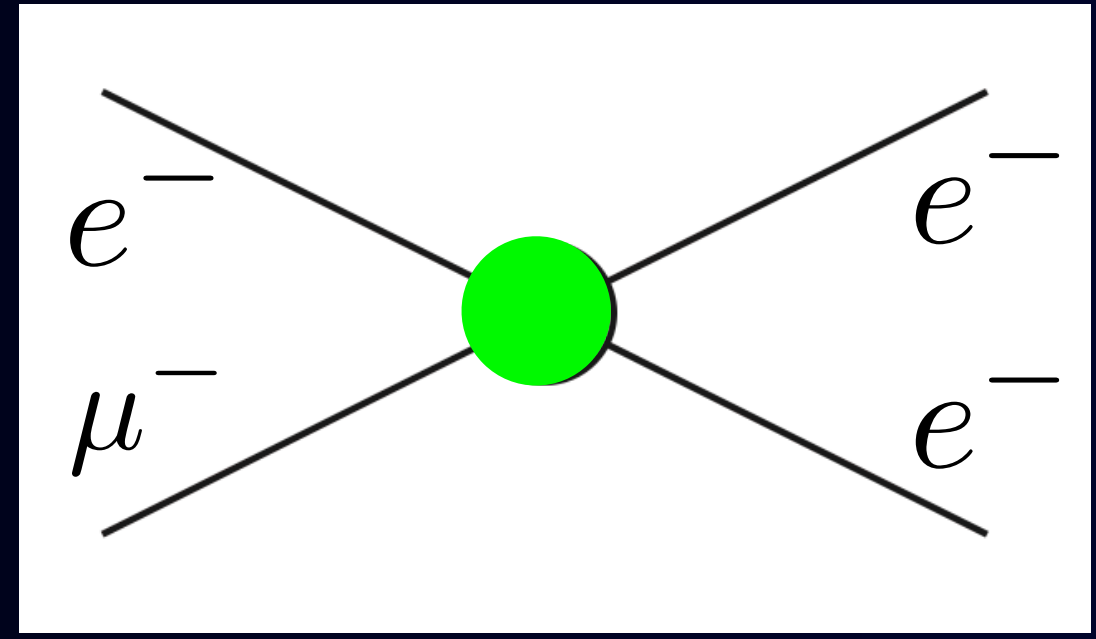
$$\mu^+ \rightarrow e^+ e^+ e^-$$

$$\mu^- N \rightarrow e^- N$$

dipole interaction



contact interaction



Effective Field Theory (EFT)

- dipole operators (L/R)
- dipole operators (L/R)
- 2 scalar contact and 4 vector contact operators

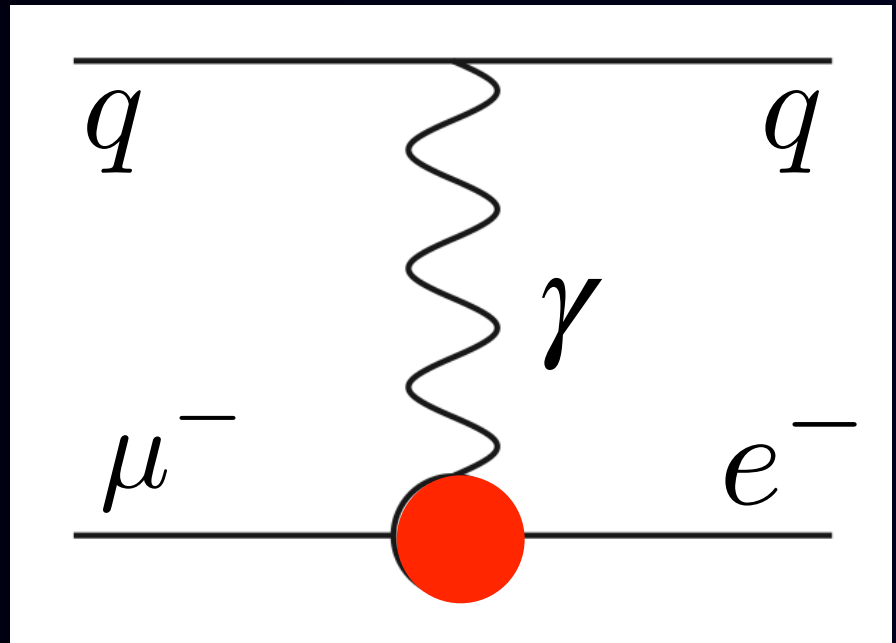
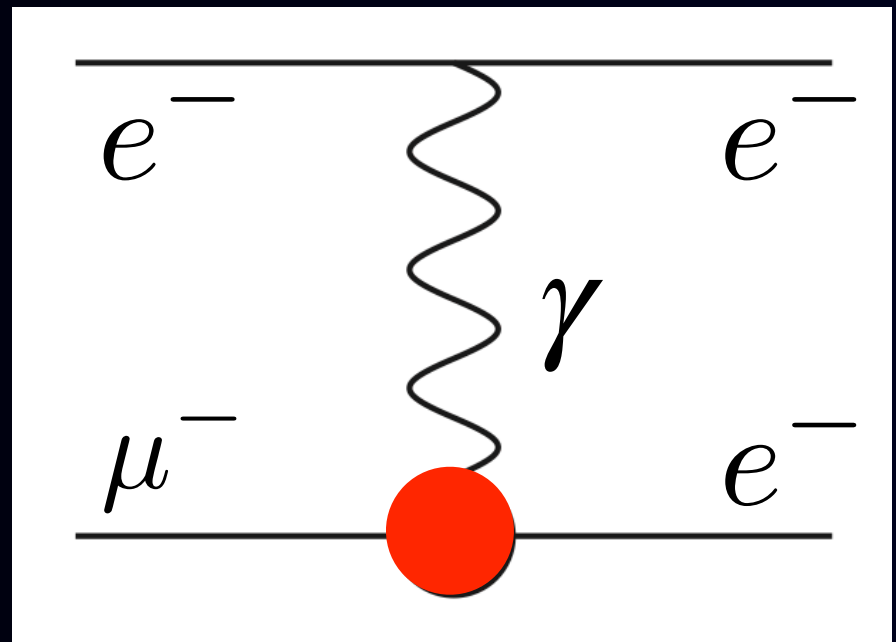
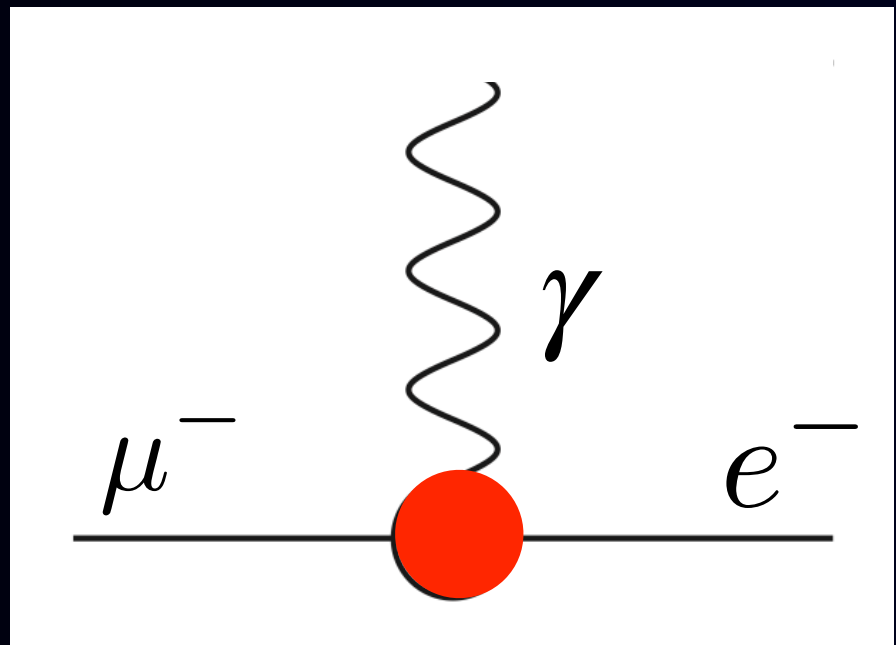
$\mu \rightarrow e$ CLFV Golden Processes and EFT

$$\mu^+ \rightarrow e^+ \gamma$$

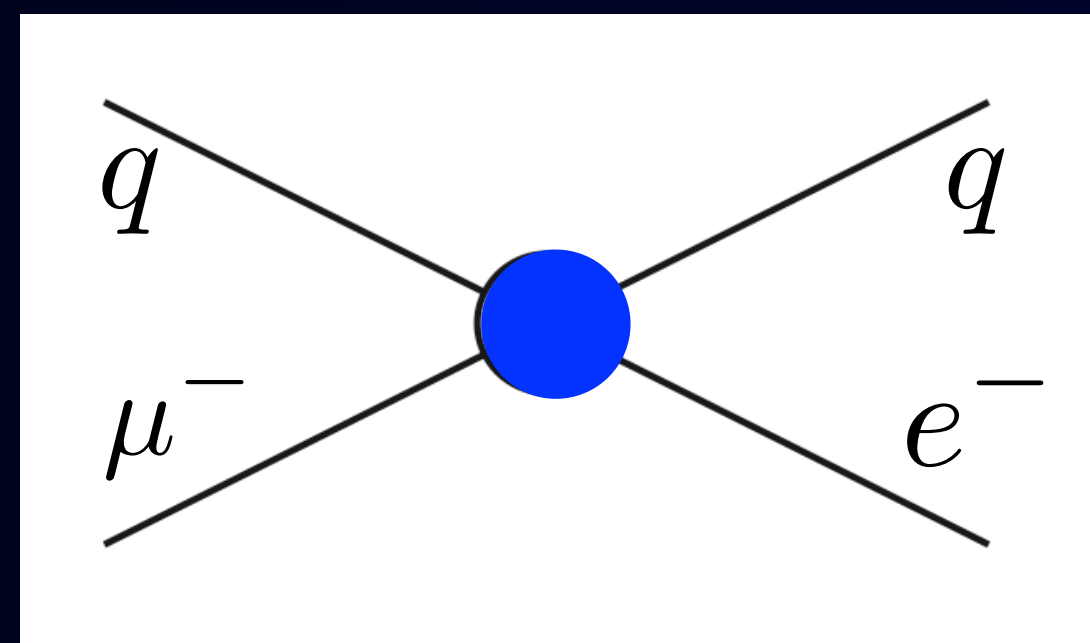
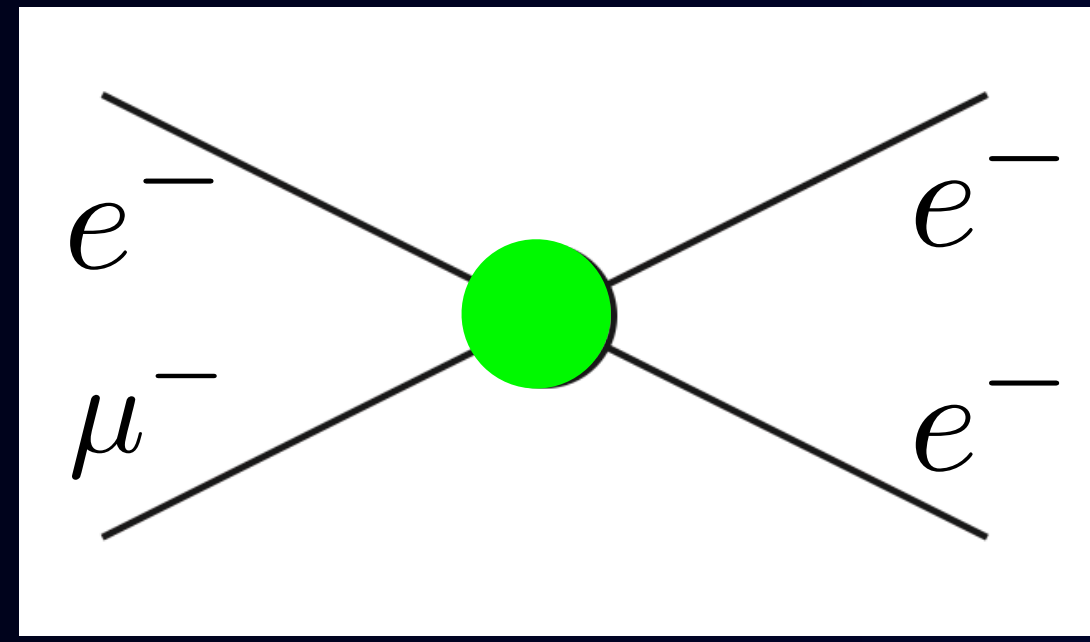
$$\mu^+ \rightarrow e^+ e^+ e^-$$

$$\mu^- N \rightarrow e^- N$$

dipole interaction



contact interaction



Effective Field Theory (EFT)

- dipole operators (L/R)
- dipole operators (L/R)
- 2 scalar contact and 4 vector contact operators
- dipole operators (L/R)
- many contact operators
- scalar, vector (spin independent)
- pseudoscalar, axial vector, tensor (spin dependent)



Present CLFV Upper Limits

from YK, Y. Okada,
RMP (2000)

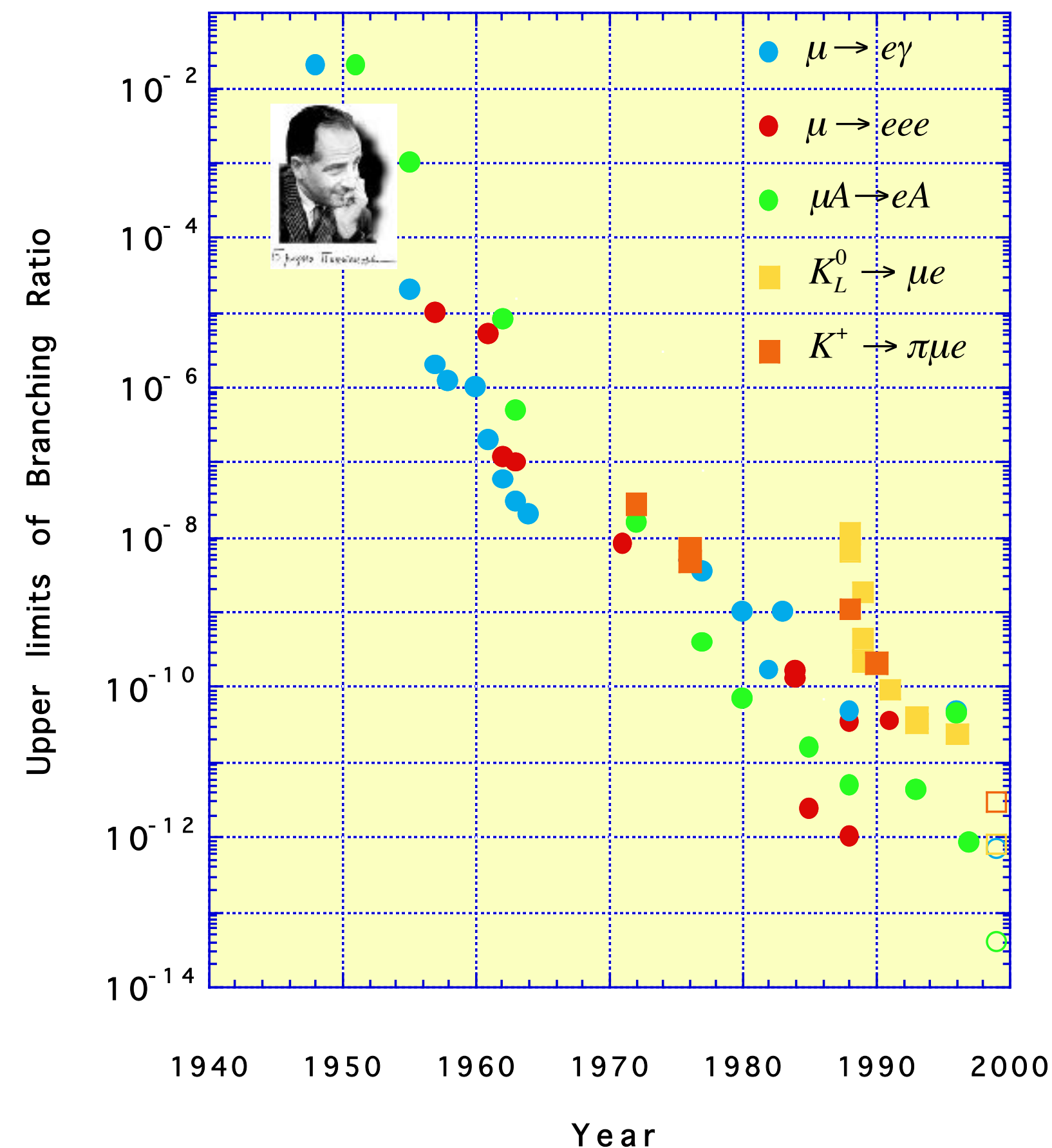
Present CLFV Upper Limits

process	present limit	future
$\tau \rightarrow \mu\eta$	$< 6.5 \times 10^{-8}$	$10^{-9} - 10^{-10}$
$\tau \rightarrow \mu\gamma$	$< 6.8 \times 10^{-8}$	
$\tau \rightarrow \mu\mu\mu$	$< 3.2 \times 10^{-8}$	
$\tau \rightarrow eee$	$< 3.6 \times 10^{-8}$	
$K_L \rightarrow e\mu$	$< 4.7 \times 10^{-12}$	
$K^+ \rightarrow \pi^+ e^- \mu^+$	$< 1.3 \times 10^{-11}$	
$B^0 \rightarrow e\mu$	$< 7.8 \times 10^{-8}$	
$B^+ \rightarrow K^+ e\mu$	$< 9.1 \times 10^{-8}$	
$Z^0 \rightarrow e\mu$	$< 7.5 \times 10^{-7}$	
$Z^0 \rightarrow e\tau$	$< 1.2 \times 10^{-5}$	
$Z^0 \rightarrow \mu\tau$	$< 9.8 \times 10^{-6}$	
$H^0 \rightarrow e\mu$	$< 3.5 \times 10^{-4}$	
$H^0 \rightarrow e\tau$	$< 3.7 \times 10^{-3}$	
$H^0 \rightarrow \mu\tau$	$< 2.5 \times 10^{-3}$	
$\mu^+ \rightarrow e^+ \gamma$	$< 4.2 \times 10^{-13}$	10^{-14} (MEG II)
$\mu^+ \rightarrow e^+ e^+ e^-$	$< 1.0 \times 10^{-12}$	10^{-16} (Mu3e)
$\mu^- \text{Au} \rightarrow e^- \text{Au}$	$< 7.0 \times 10^{-13}$	10^{-17} (COMET, Mu2e)
$\mu^- \text{Ti} \rightarrow e^+ \text{Ca}$	$< 3.6 \times 10^{-11}$	10^{-17} (COMET, Mu2e)
$\mu^+ e^- \rightarrow \mu^- e^+$	$< 8.3 \times 10^{-11}$	

from YK, Y. Okada,
RMP (2000)

Present CLFV Upper Limits

process	present limit	future
$\tau \rightarrow \mu\eta$	$< 6.5 \times 10^{-8}$	$10^{-9} - 10^{-10}$
$\tau \rightarrow \mu\gamma$	$< 6.8 \times 10^{-8}$	
$\tau \rightarrow \mu\mu\mu$	$< 3.2 \times 10^{-8}$	
$\tau \rightarrow eee$	$< 3.6 \times 10^{-8}$	
$K_L \rightarrow e\mu$	$< 4.7 \times 10^{-12}$	
$K^+ \rightarrow \pi^+ e^- \mu^+$	$< 1.3 \times 10^{-11}$	
$B^0 \rightarrow e\mu$	$< 7.8 \times 10^{-8}$	
$B^+ \rightarrow K^+ e\mu$	$< 9.1 \times 10^{-8}$	
$Z^0 \rightarrow e\mu$	$< 7.5 \times 10^{-7}$	
$Z^0 \rightarrow e\tau$	$< 1.2 \times 10^{-5}$	
$Z^0 \rightarrow \mu\tau$	$< 9.8 \times 10^{-6}$	
$H^0 \rightarrow e\mu$	$< 3.5 \times 10^{-4}$	
$H^0 \rightarrow e\tau$	$< 3.7 \times 10^{-3}$	
$H^0 \rightarrow \mu\tau$	$< 2.5 \times 10^{-3}$	
$\mu^+ \rightarrow e^+ \gamma$	$< 4.2 \times 10^{-13}$	10^{-14} (MEG II)
$\mu^+ \rightarrow e^+ e^+ e^-$	$< 1.0 \times 10^{-12}$	10^{-16} (Mu3e)
$\mu^- \text{Au} \rightarrow e^- \text{Au}$	$< 7.0 \times 10^{-13}$	10^{-17} (COMET, Mu2e)
$\mu^- \text{Ti} \rightarrow e^+ \text{Ca}$	$< 3.6 \times 10^{-11}$	10^{-17} (COMET, Mu2e)
$\mu^+ e^- \rightarrow \mu^- e^+$	$< 8.3 \times 10^{-11}$	

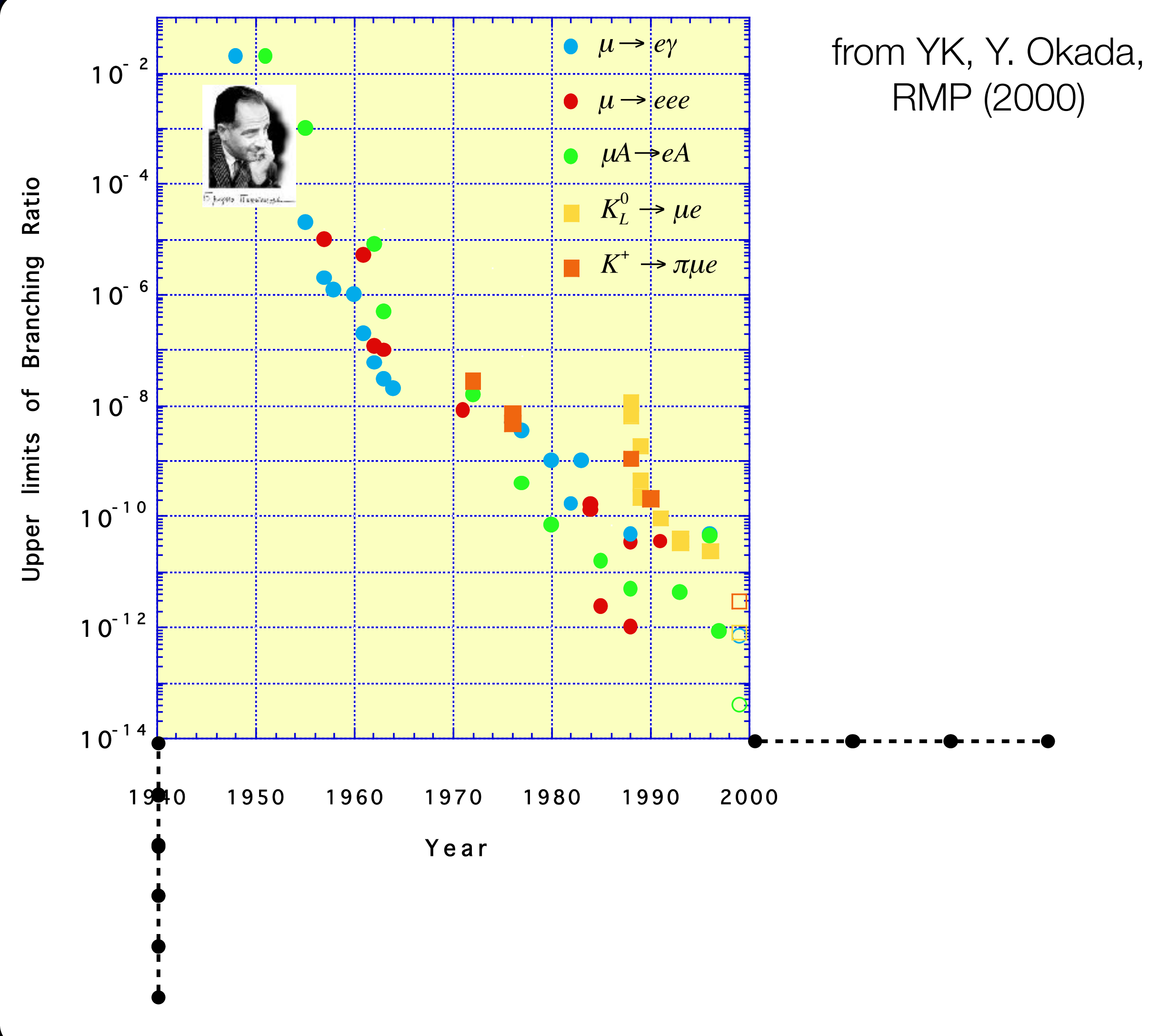


from YK, Y. Okada, RMP (2000)

Present CLFV Upper Limits



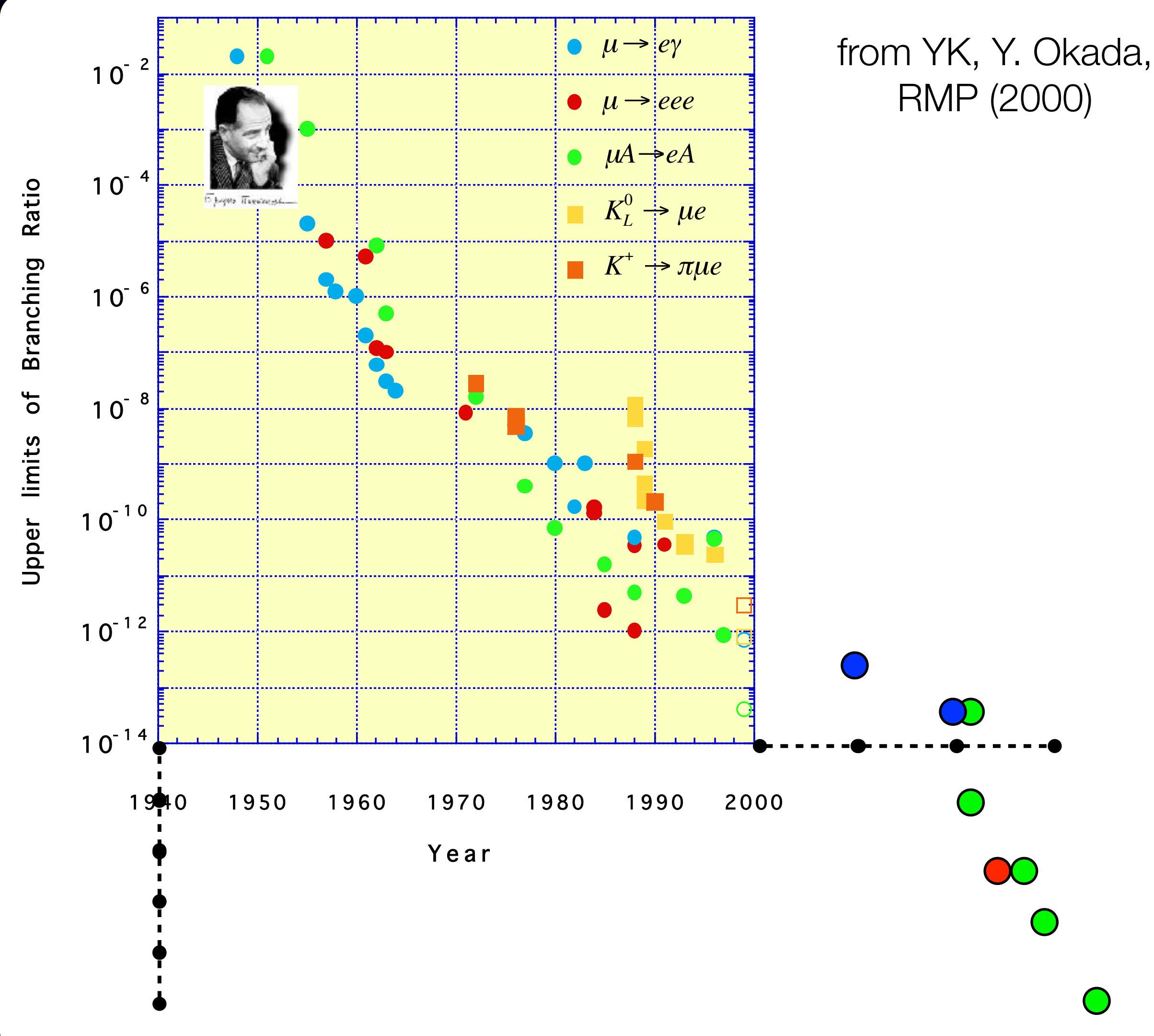
process	present limit	future
$\tau \rightarrow \mu\eta$	$< 6.5 \times 10^{-8}$	$10^{-9} - 10^{-10}$
$\tau \rightarrow \mu\gamma$	$< 6.8 \times 10^{-8}$	
$\tau \rightarrow \mu\mu\mu$	$< 3.2 \times 10^{-8}$	
$\tau \rightarrow eee$	$< 3.6 \times 10^{-8}$	
$K_L \rightarrow e\mu$	$< 4.7 \times 10^{-12}$	
$K^+ \rightarrow \pi^+ e^- \mu^+$	$< 1.3 \times 10^{-11}$	
$B^0 \rightarrow e\mu$	$< 7.8 \times 10^{-8}$	
$B^+ \rightarrow K^+ e\mu$	$< 9.1 \times 10^{-8}$	
$Z^0 \rightarrow e\mu$	$< 7.5 \times 10^{-7}$	
$Z^0 \rightarrow e\tau$	$< 1.2 \times 10^{-5}$	
$Z^0 \rightarrow \mu\tau$	$< 9.8 \times 10^{-6}$	
$H^0 \rightarrow e\mu$	$< 3.5 \times 10^{-4}$	
$H^0 \rightarrow e\tau$	$< 3.7 \times 10^{-3}$	
$H^0 \rightarrow \mu\tau$	$< 2.5 \times 10^{-3}$	
$\mu^+ \rightarrow e^+ \gamma$	$< 4.2 \times 10^{-13}$	10^{-14} (MEG II)
$\mu^+ \rightarrow e^+ e^+ e^-$	$< 1.0 \times 10^{-12}$	10^{-16} (Mu3e)
$\mu^- \text{Au} \rightarrow e^- \text{Au}$	$< 7.0 \times 10^{-13}$	10^{-17} (COMET, Mu2e)
$\mu^- \text{Ti} \rightarrow e^+ \text{Ca}$	$< 3.6 \times 10^{-11}$	10^{-17} (COMET, Mu2e)
$\mu^+ e^- \rightarrow \mu^- e^+$	$< 8.3 \times 10^{-11}$	

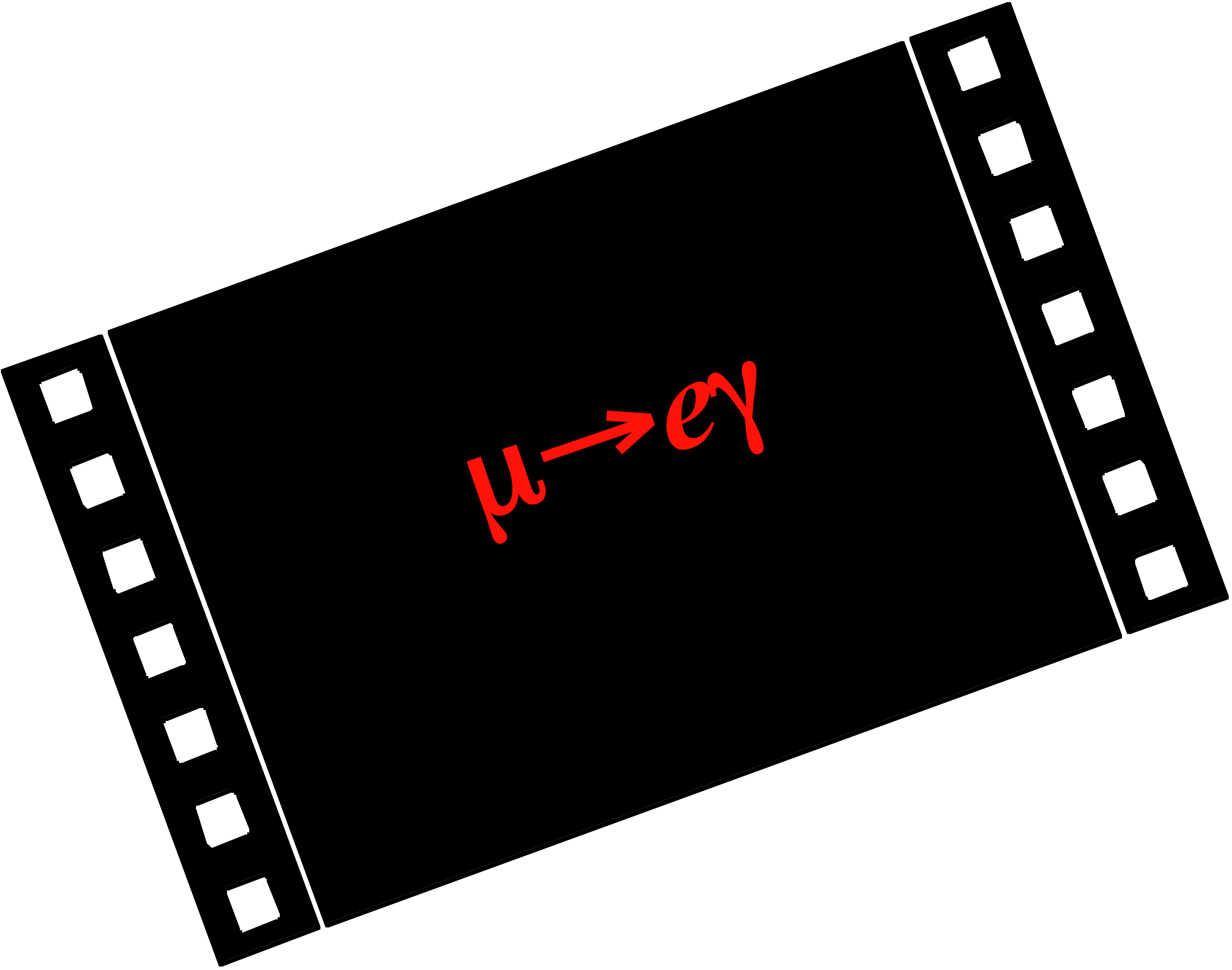


Present CLFV Upper Limits



process	present limit	future
$\tau \rightarrow \mu\eta$	$< 6.5 \times 10^{-8}$	$10^{-9} - 10^{-10}$
$\tau \rightarrow \mu\gamma$	$< 6.8 \times 10^{-8}$	
$\tau \rightarrow \mu\mu\mu$	$< 3.2 \times 10^{-8}$	
$\tau \rightarrow eee$	$< 3.6 \times 10^{-8}$	
$K_L \rightarrow e\mu$	$< 4.7 \times 10^{-12}$	
$K^+ \rightarrow \pi^+ e^- \mu^+$	$< 1.3 \times 10^{-11}$	
$B^0 \rightarrow e\mu$	$< 7.8 \times 10^{-8}$	
$B^+ \rightarrow K^+ e\mu$	$< 9.1 \times 10^{-8}$	
$Z^0 \rightarrow e\mu$	$< 7.5 \times 10^{-7}$	
$Z^0 \rightarrow e\tau$	$< 1.2 \times 10^{-5}$	
$Z^0 \rightarrow \mu\tau$	$< 9.8 \times 10^{-6}$	
$H^0 \rightarrow e\mu$	$< 3.5 \times 10^{-4}$	
$H^0 \rightarrow e\tau$	$< 3.7 \times 10^{-3}$	
$H^0 \rightarrow \mu\tau$	$< 2.5 \times 10^{-3}$	
$\mu^+ \rightarrow e^+ \gamma$	$< 4.2 \times 10^{-13}$	10^{-14} (MEG II)
$\mu^+ \rightarrow e^+ e^+ e^-$	$< 1.0 \times 10^{-12}$	10^{-16} (Mu3e)
$\mu^- \text{Au} \rightarrow e^- \text{Au}$	$< 7.0 \times 10^{-13}$	10^{-17} (COMET, Mu2e)
$\mu^- \text{Ti} \rightarrow e^+ \text{Ca}$	$< 3.6 \times 10^{-11}$	10^{-17} (COMET, Mu2e)
$\mu^+ e^- \rightarrow \mu^- e^+$	$< 8.3 \times 10^{-11}$	



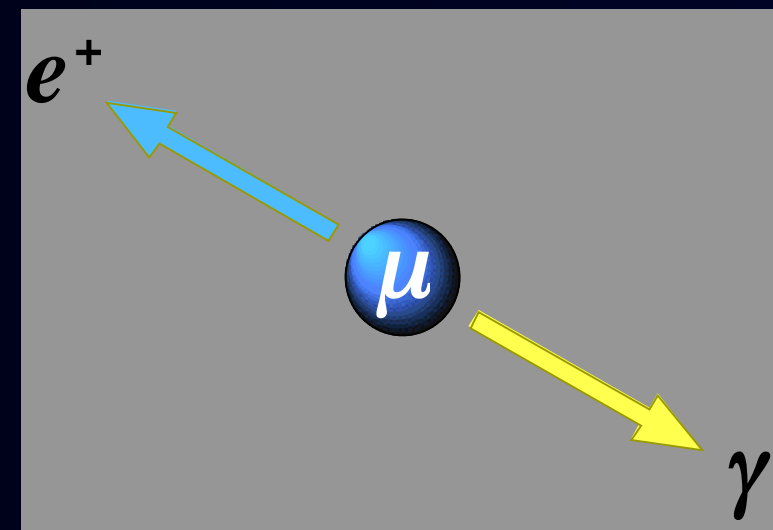


$\mu^+ \rightarrow e^+ \gamma$: MEG II



$\mu^+ \rightarrow e^+ \gamma$: MEG II

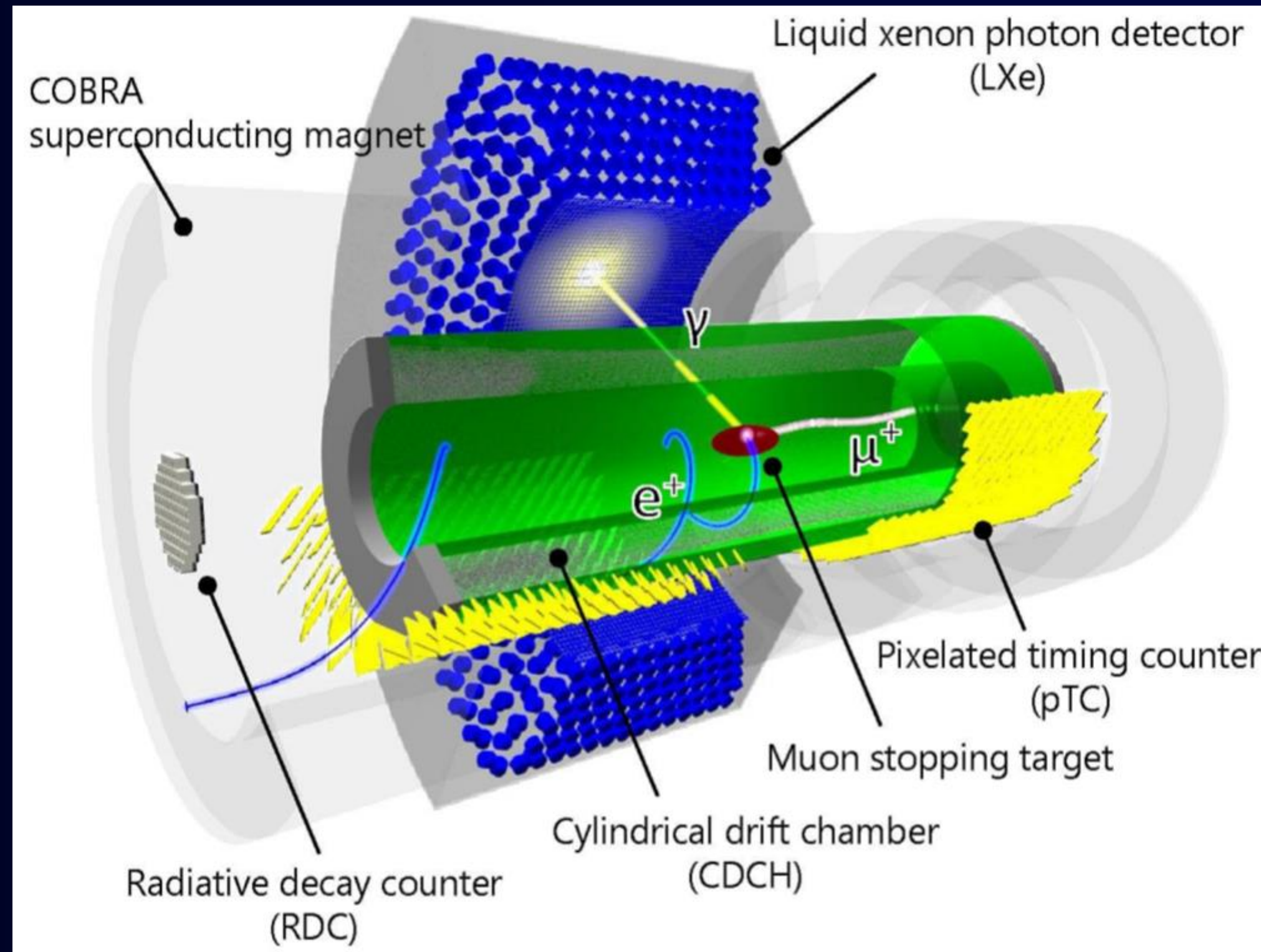
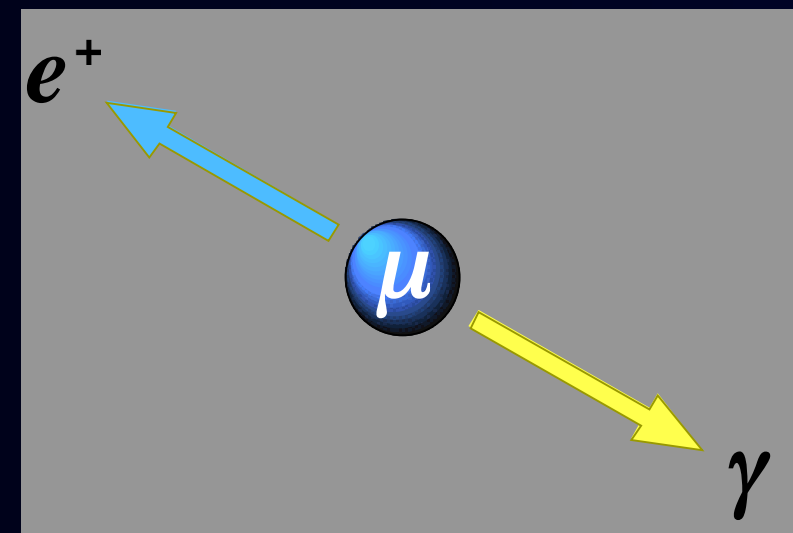
- Event Signature (μ^+ decay at rest)
 - mono energetic, $m_e = m_\gamma = m_\mu/2$ (=52.8 MeV)
 - angle $\theta_{e\gamma}=180$ degrees
 - time coincidence $\Delta t_{e\gamma}$
- Backgrounds
 - physics background, $\mu^+ \rightarrow e^+ \nu \bar{\nu} \gamma$
 - accidental background
 - e^+ in $\mu^+ \rightarrow e^+ \nu \bar{\nu}$ and γ in $\mu^+ \rightarrow e^+ \nu \bar{\nu} \gamma$
- MEG experiment at PSI
 - $B(\mu^+ \rightarrow e^+ \gamma) < 4.2 \times 10^{-13}$ (90 % C.L.)
- MEG II experiment at PSI
 - x10 times improvement



$\mu^+ \rightarrow e^+ \gamma$: MEG II

- **Event Signature (μ^+ decay at rest)**
 - mono energetic, $m_e = m_\gamma = m_\mu/2$ (=52.8 MeV)
 - angle $\theta_{e\gamma}=180$ degrees
 - time coincidence $\Delta t_{e\gamma}$
- **Backgrounds**
 - physics background, $\mu^+ \rightarrow e^+ \nu \bar{\nu} \gamma$
 - accidental background
 - e^+ in $\mu^+ \rightarrow e^+ \nu \bar{\nu}$ and γ in $\mu^+ \rightarrow e^+ \nu \bar{\nu} \gamma$
- **MEG experiment at PSI**
 - $B(\mu^+ \rightarrow e^+ \gamma) < 4.2 \times 10^{-13}$ (90% C.L.)
- **MEG II experiment at PSI**
 - x10 times improvement

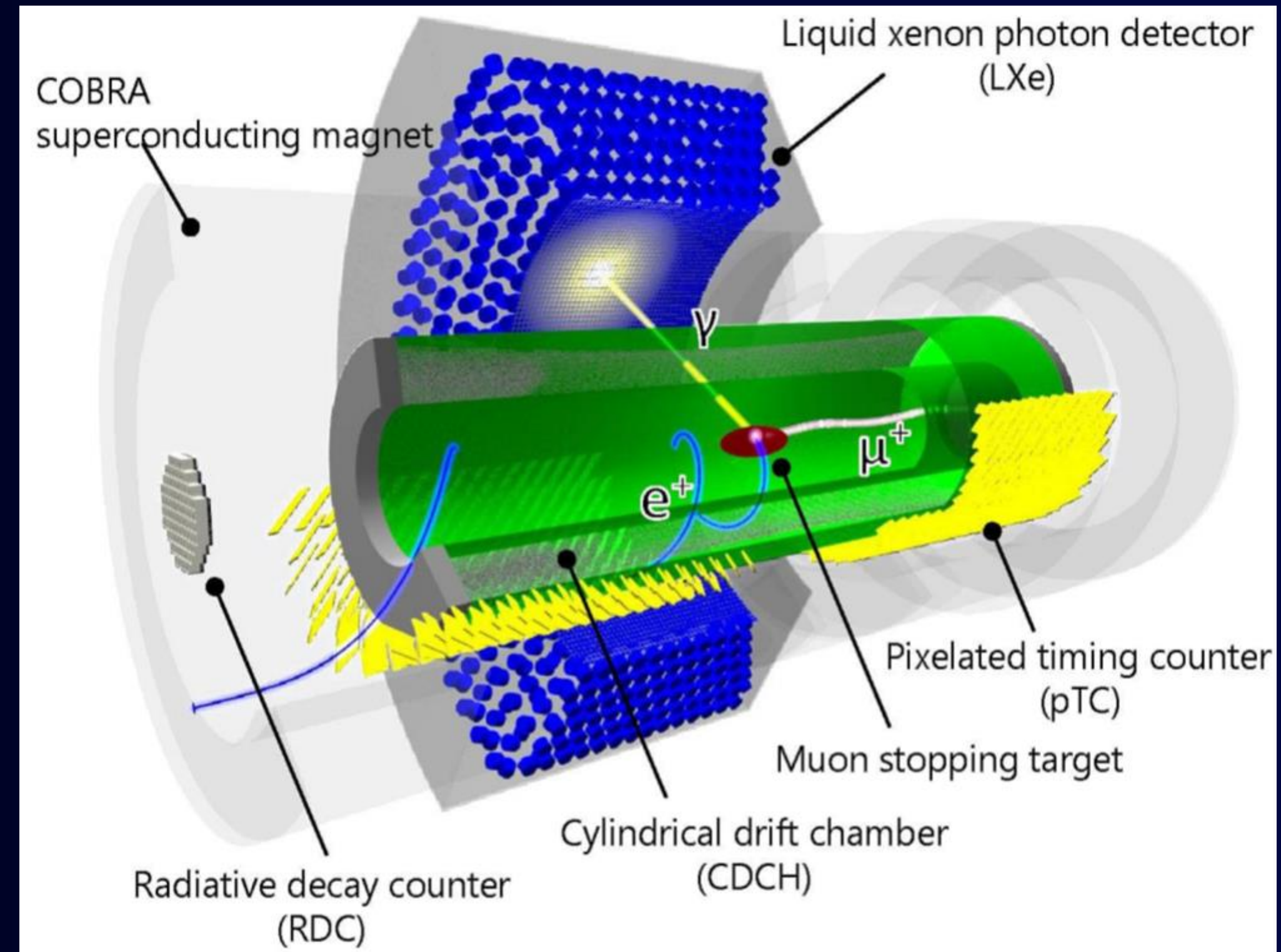
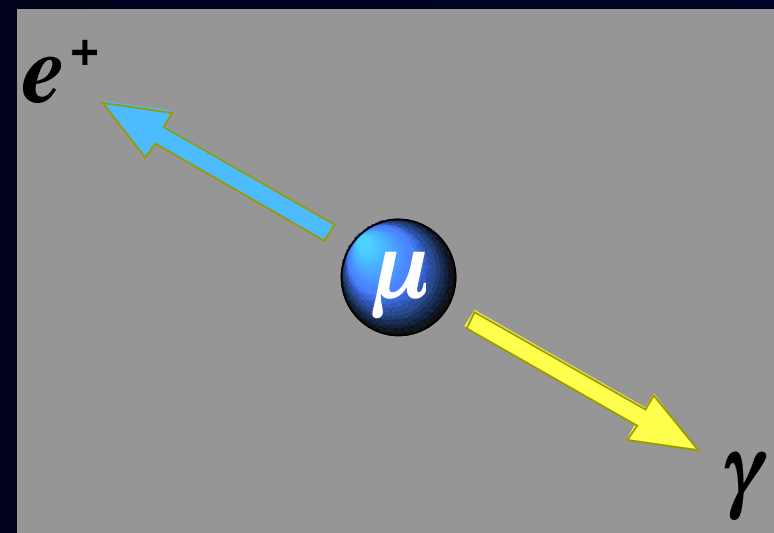
MEG-II



$\mu^+ \rightarrow e^+ \gamma$: MEG II

- **Event Signature (μ^+ decay at rest)**
 - mono energetic, $m_e = m_\gamma = m_\mu/2$ (=52.8 MeV)
 - angle $\theta_{e\gamma}=180$ degrees
 - time coincidence $\Delta t_{e\gamma}$
- **Backgrounds**
 - physics background, $\mu^+ \rightarrow e^+ \nu \bar{\nu} \gamma$
 - accidental background
 - e^+ in $\mu^+ \rightarrow e^+ \nu \bar{\nu}$ and γ in $\mu^+ \rightarrow e^+ \nu \bar{\nu} \gamma$
- **MEG experiment at PSI**
 - $B(\mu^+ \rightarrow e^+ \gamma) < 4.2 \times 10^{-13}$ (90% C.L.)
- **MEG II experiment at PSI**
 - x10 times improvement

MEG-II



MEG II : $B(\mu^+ \rightarrow e^+ \gamma) < 6 \times 10^{-14}$

- x2 muon beam intensity
- x2 all detector resolution
- x2 efficiency

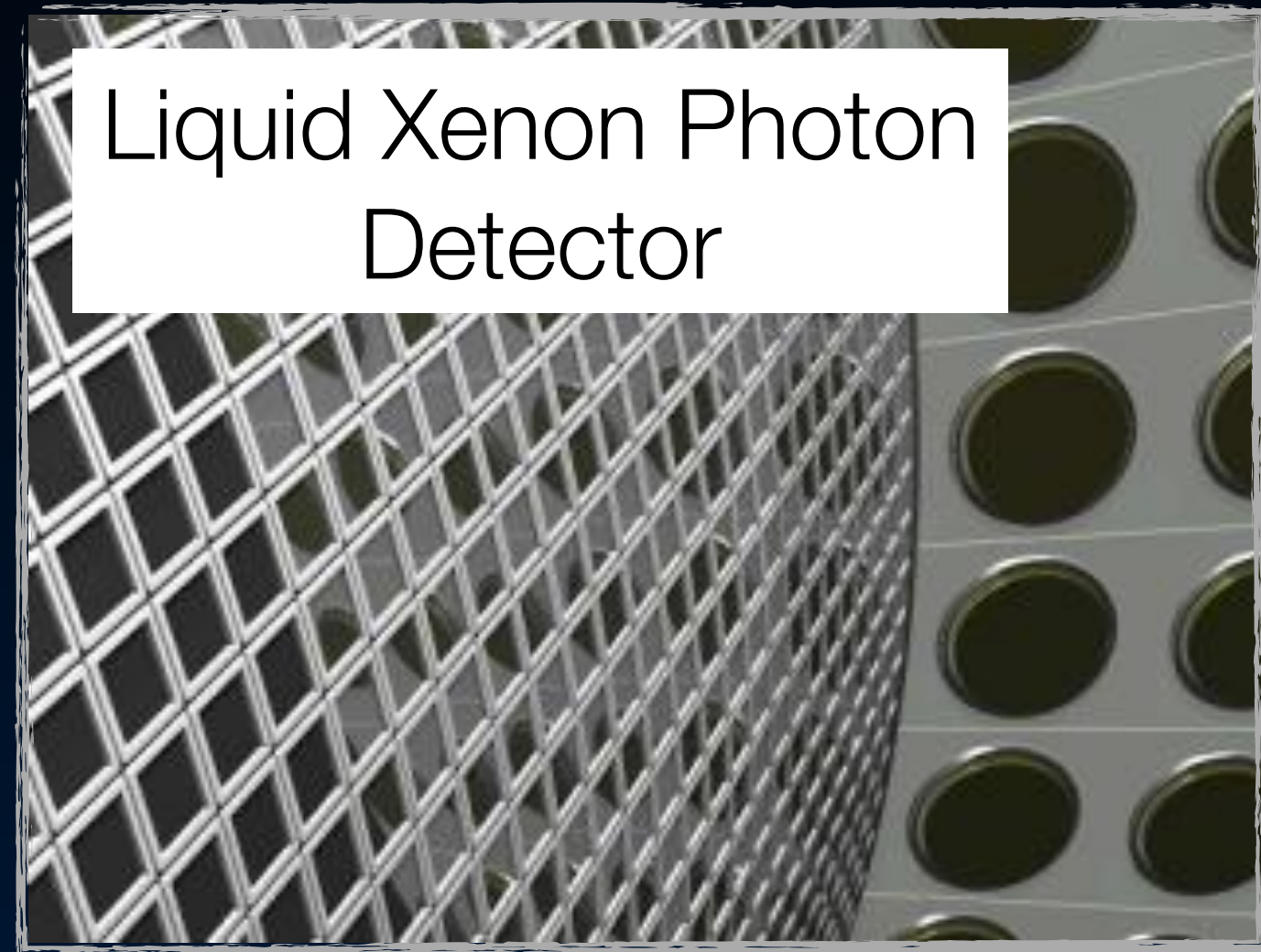


MEG II Detector Upgrade

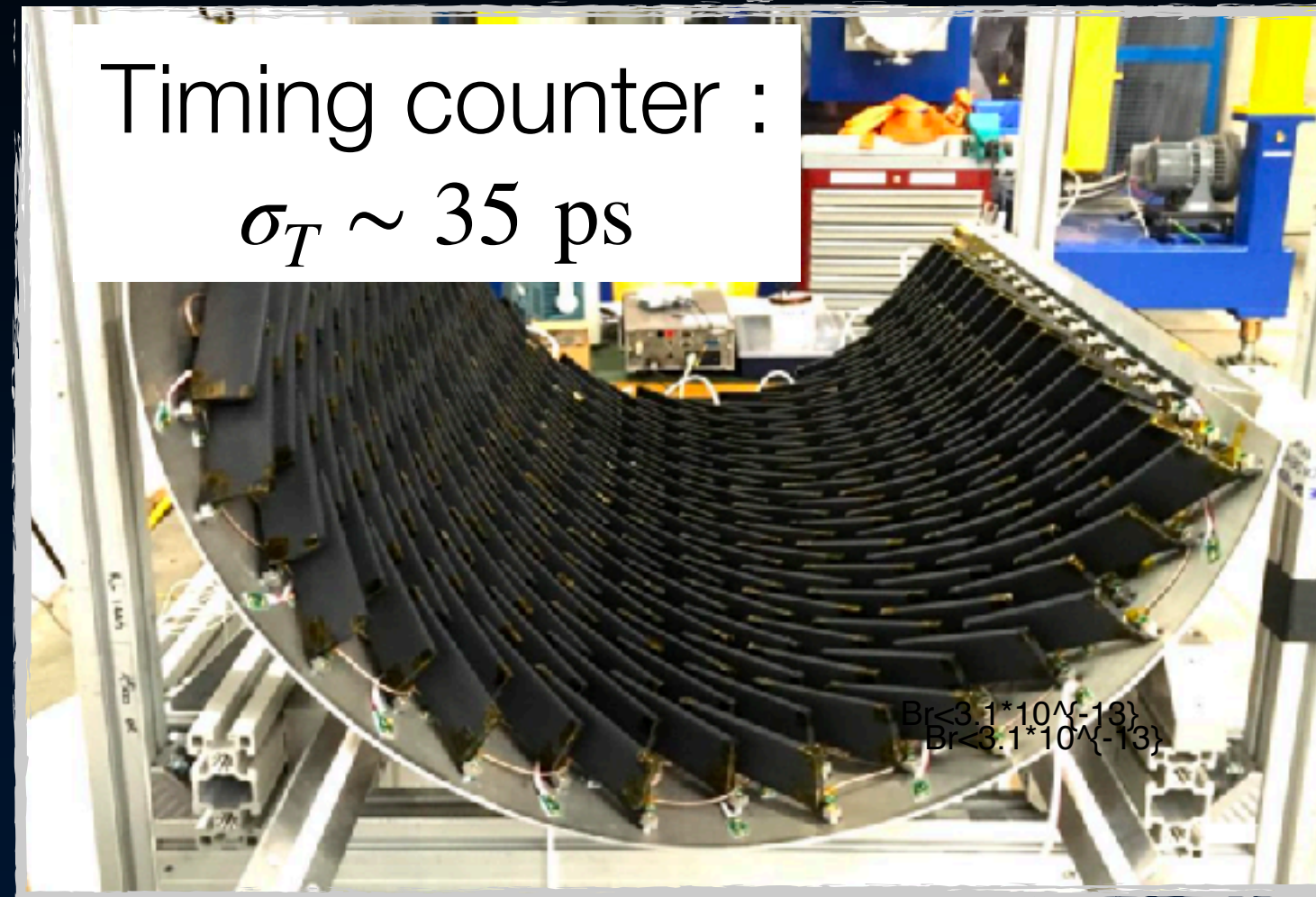
$Br < 3.1 \cdot 10^{-13}$
 $Br < 3.1 \cdot 10^{-13}$

MEG II Detector Upgrade

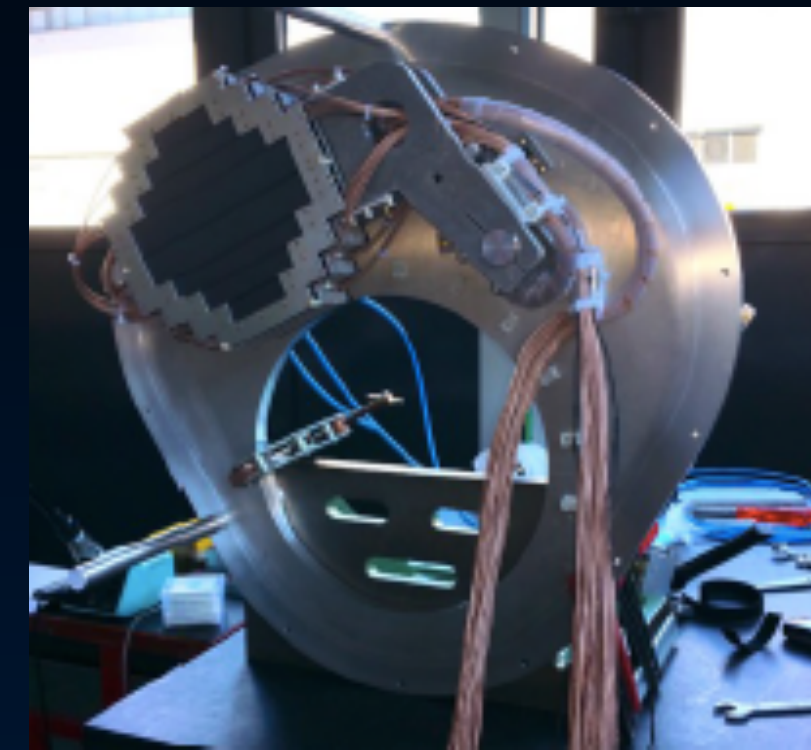
Liquid Xenon Photon Detector



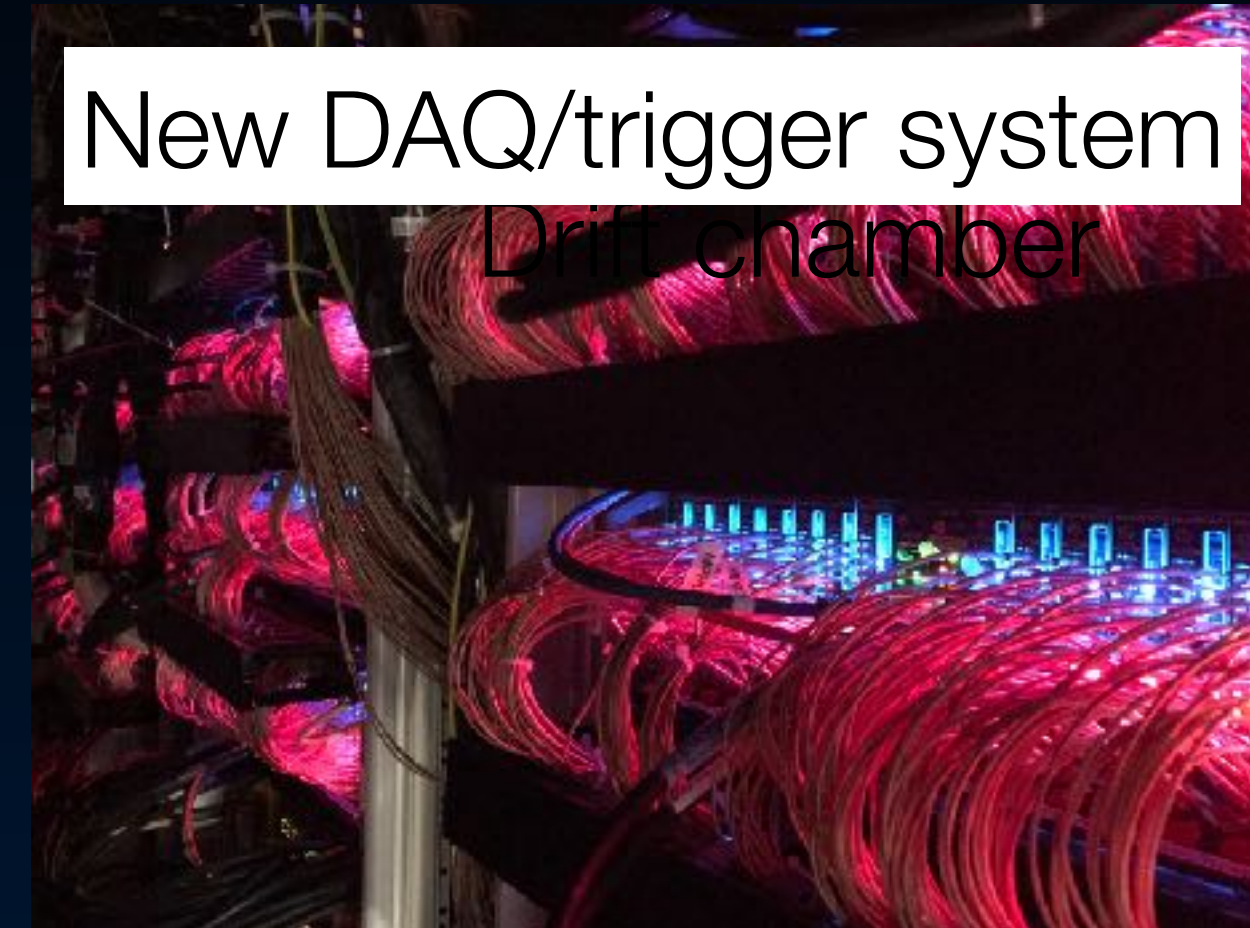
Timing counter :
 $\sigma_T \sim 35$ ps



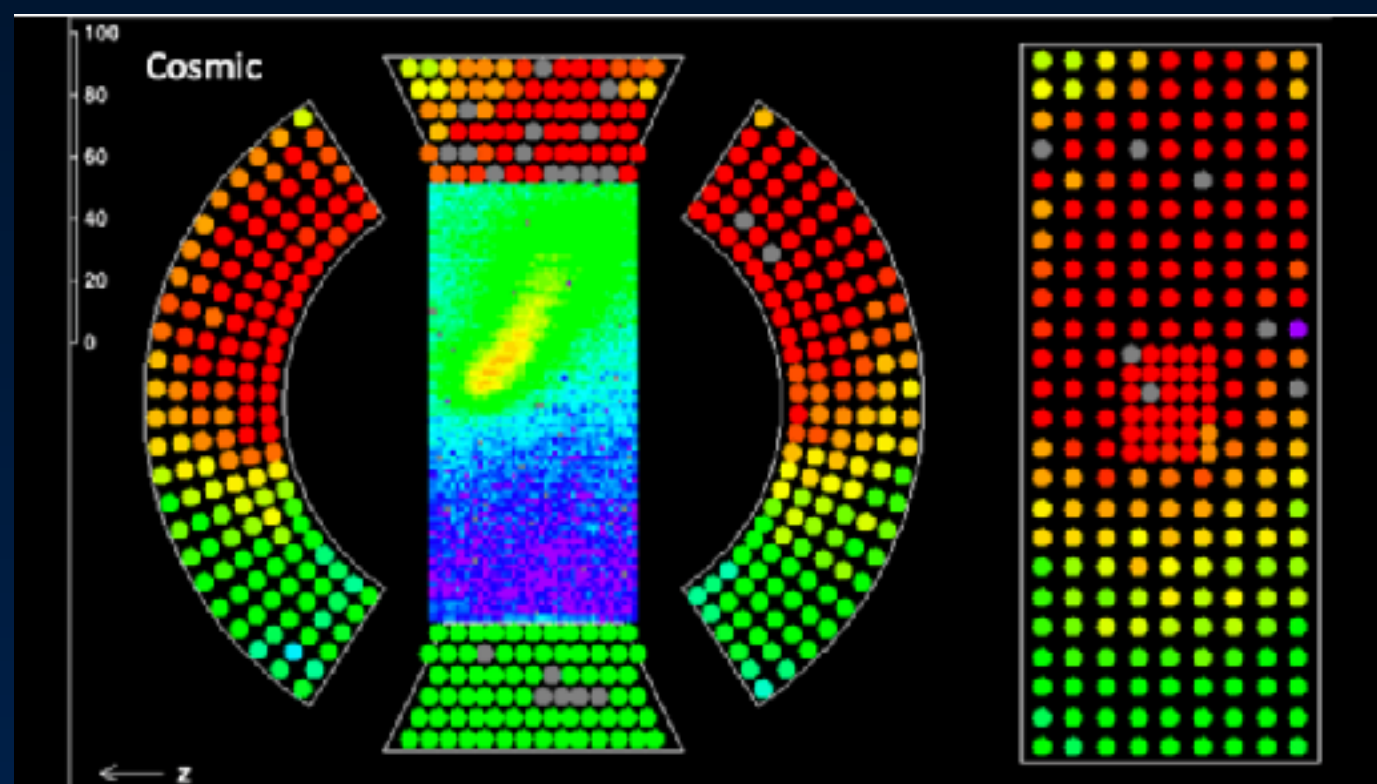
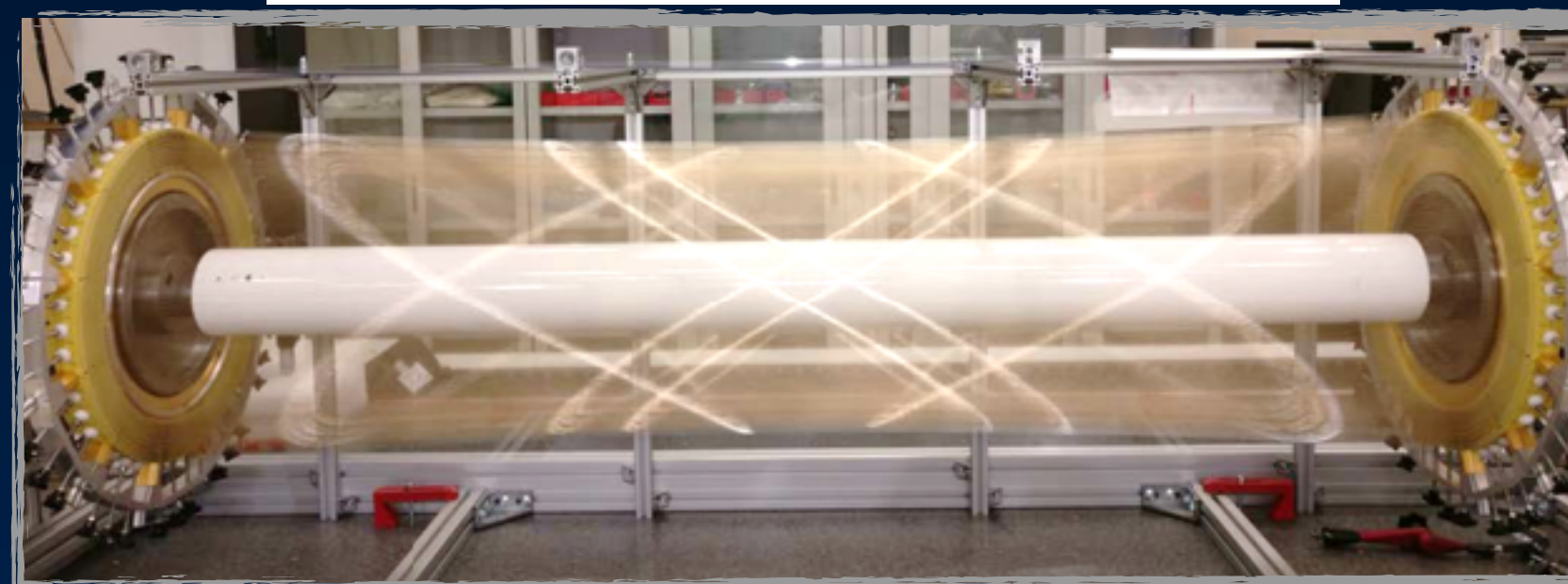
RDC detector



New DAQ/trigger system
Drift chamber

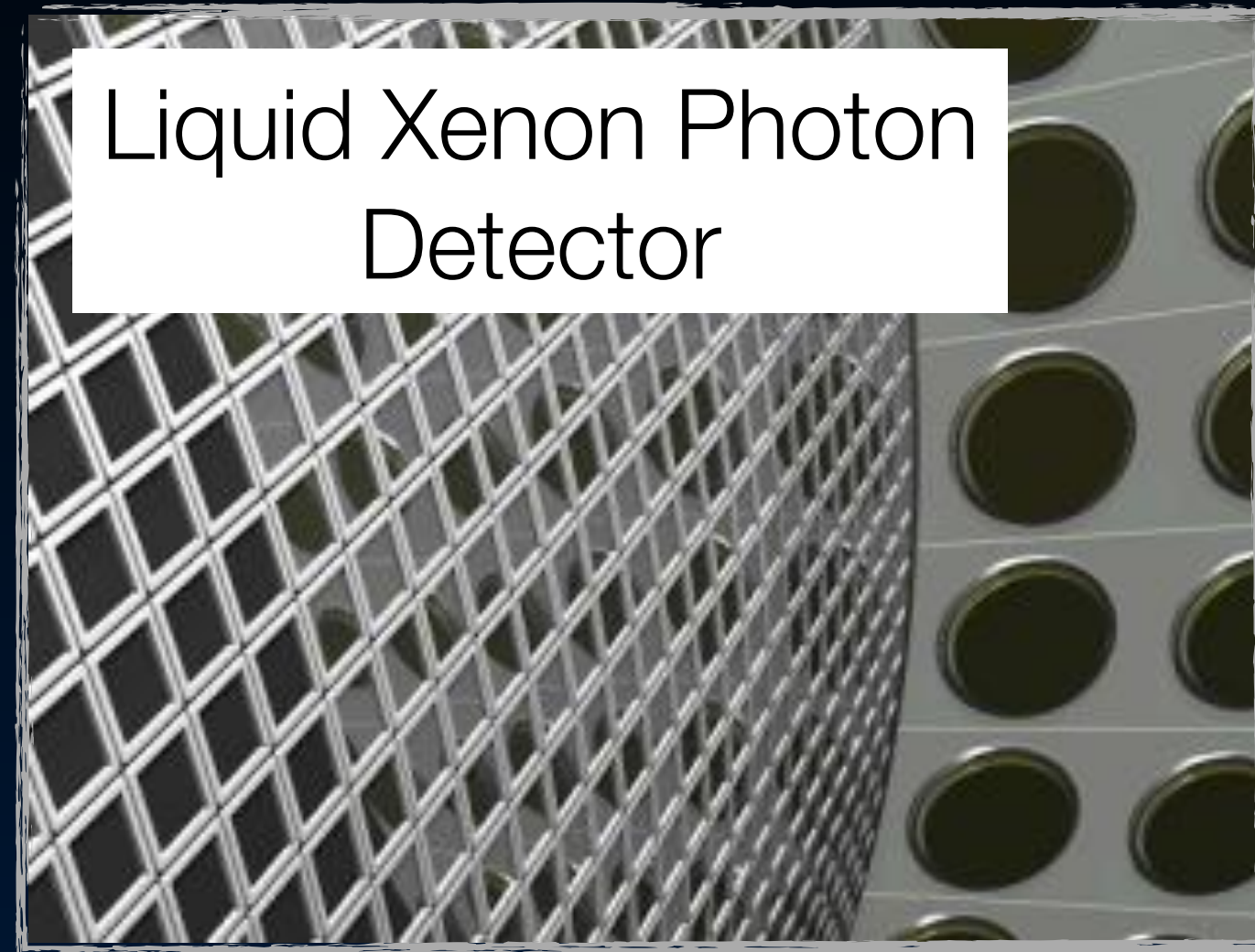


Cylindrical drift chamber

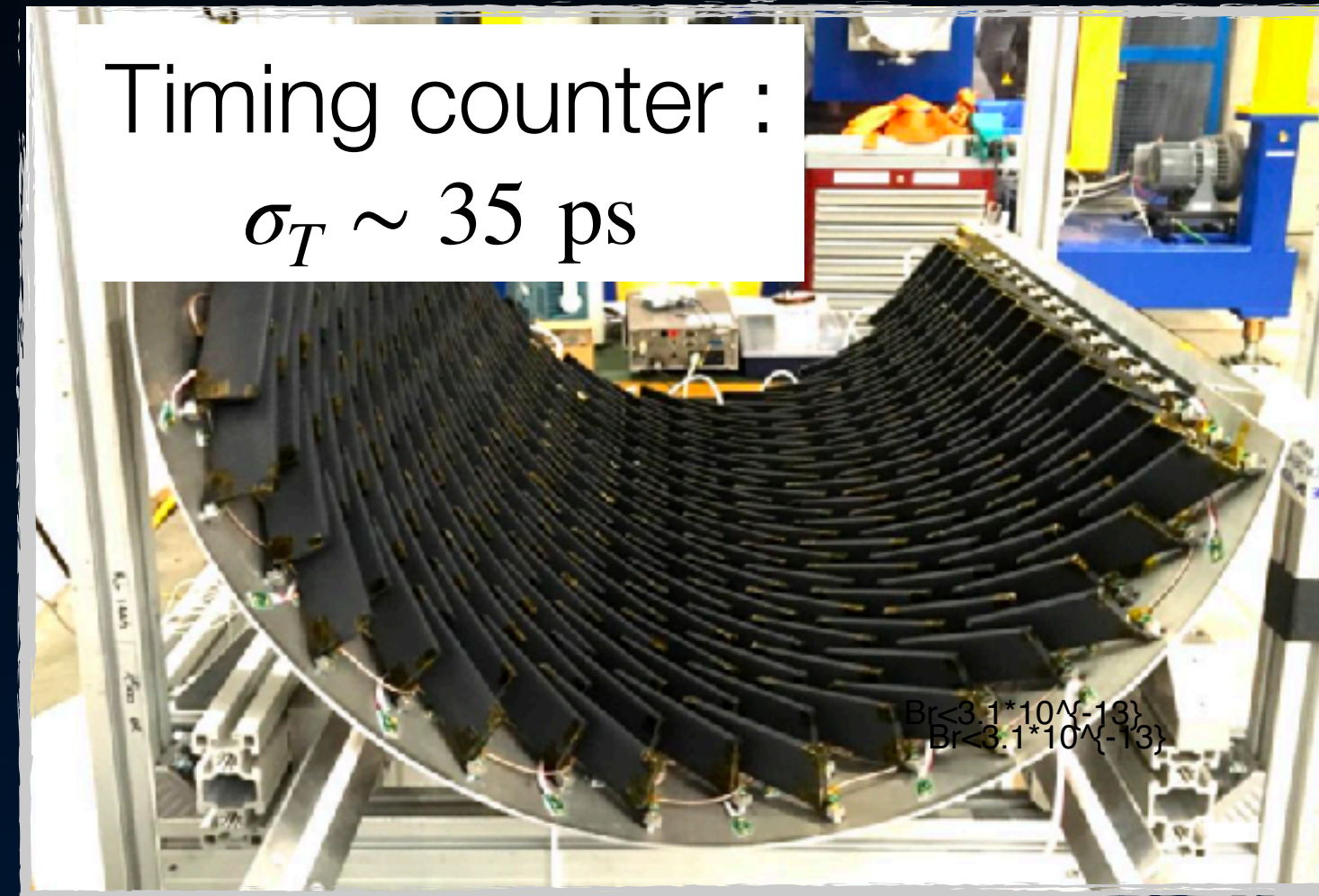


MEG II Detector Upgrade

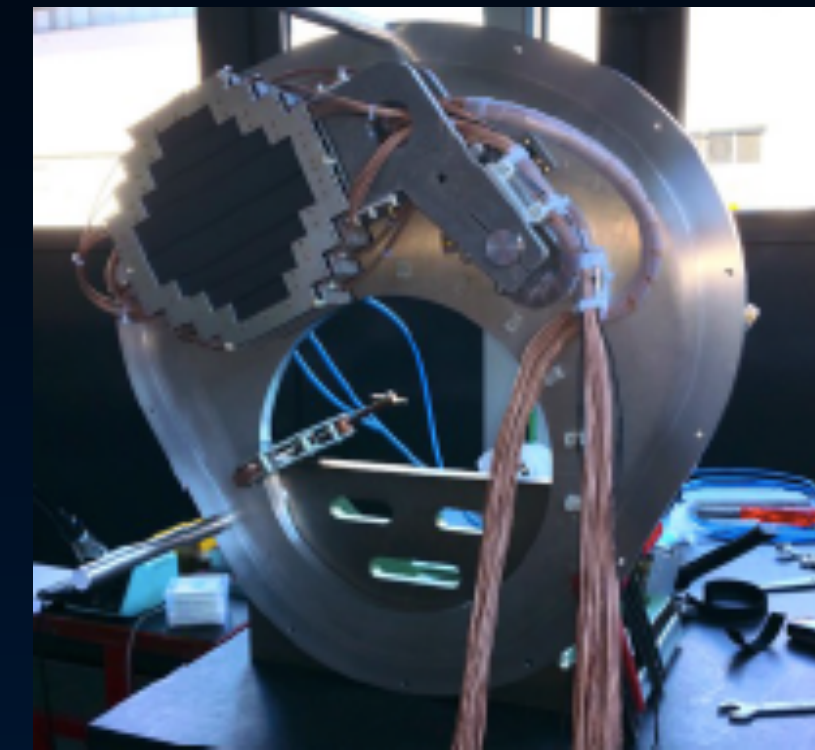
Liquid Xenon Photon Detector



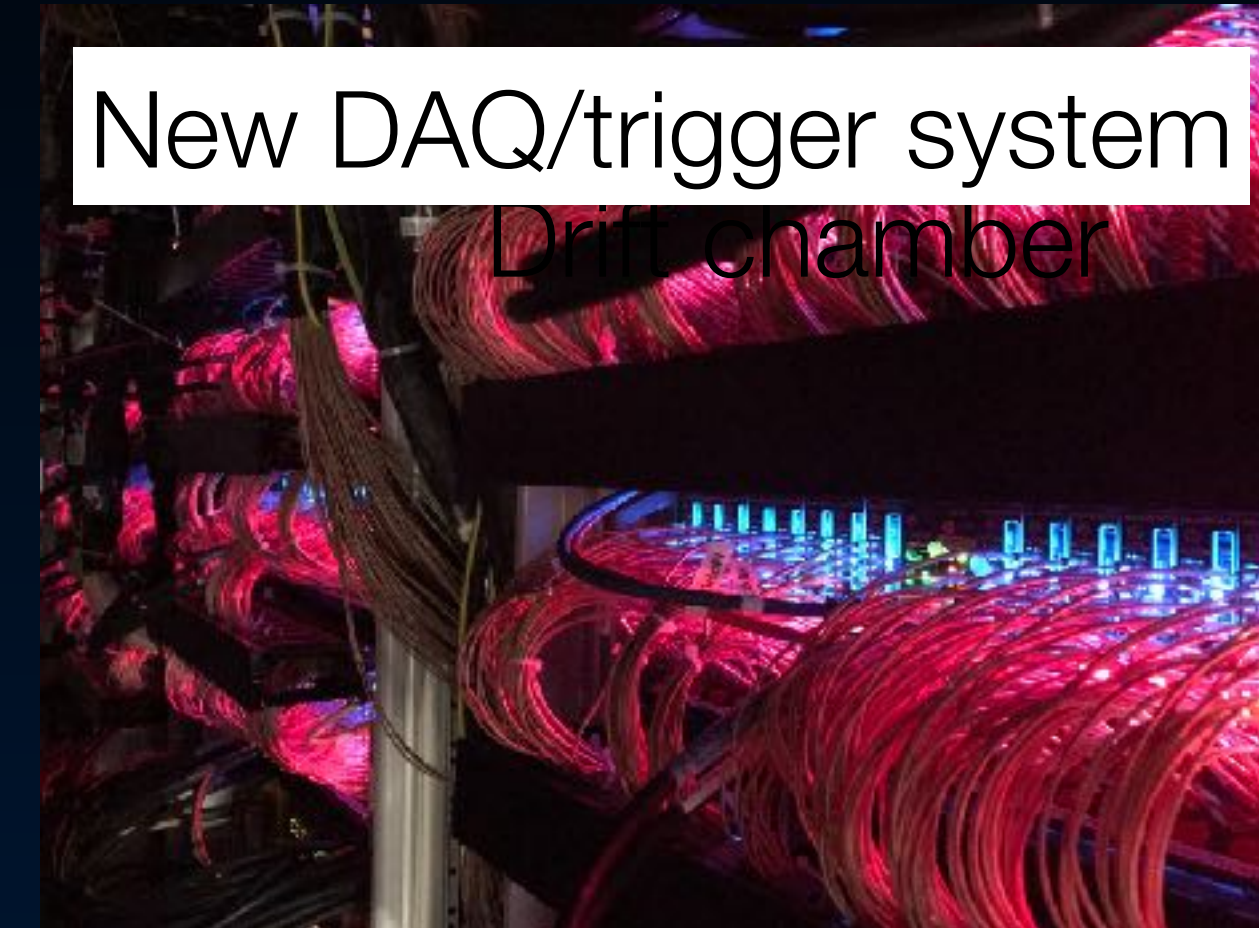
Timing counter :
 $\sigma_T \sim 35$ ps



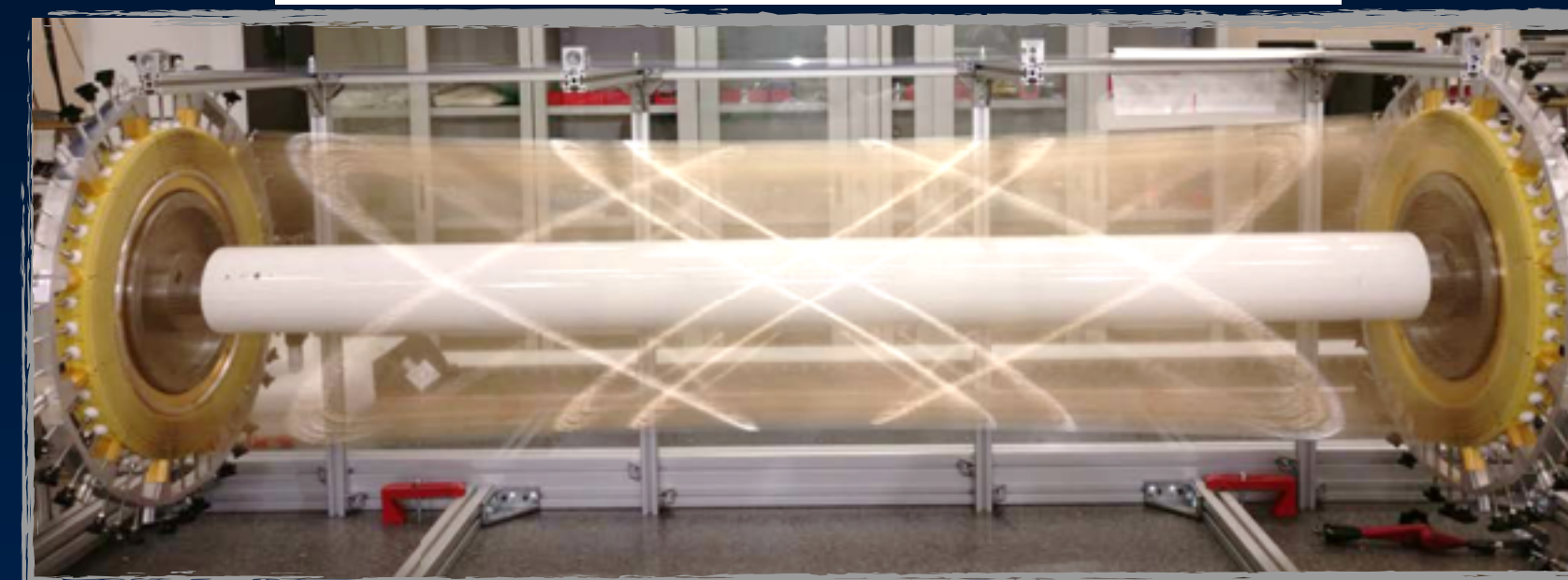
RDC detector



New DAQ/trigger system
Drift chamber

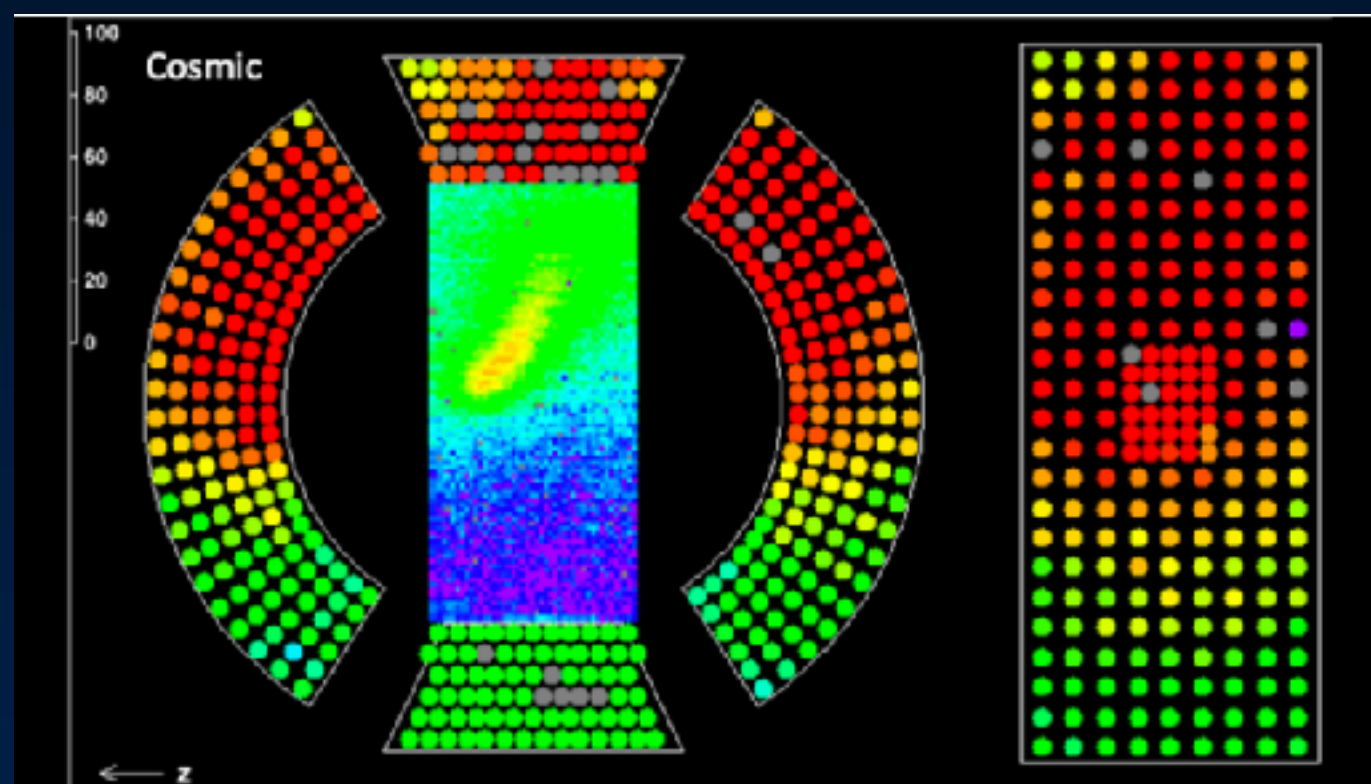


Cylindrical drift chamber



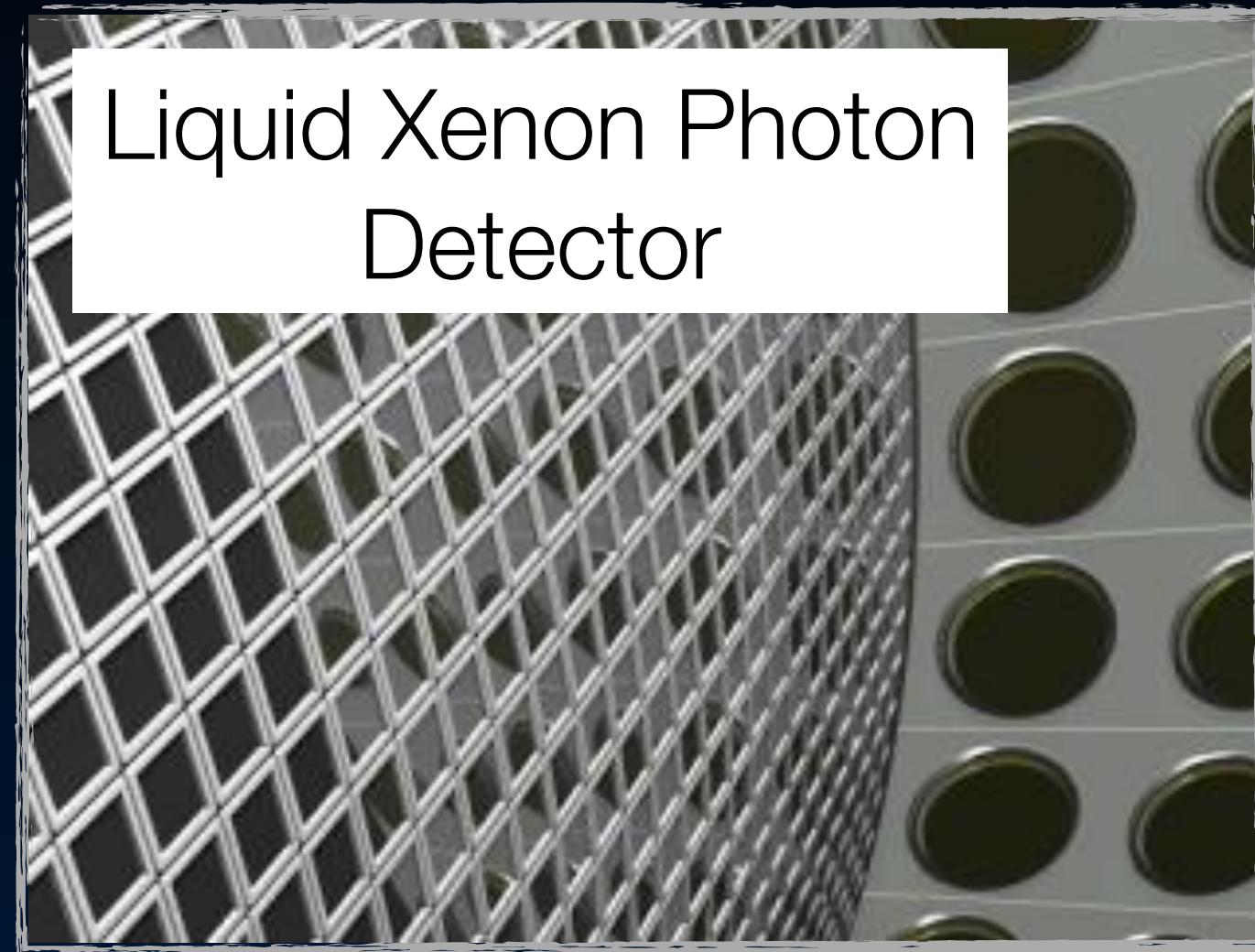
Timeline

- Physics runs in 2021, 2022, 2023 - 2025, (2026)
- PSI HIPA accelerator shutdown in 2027-2028

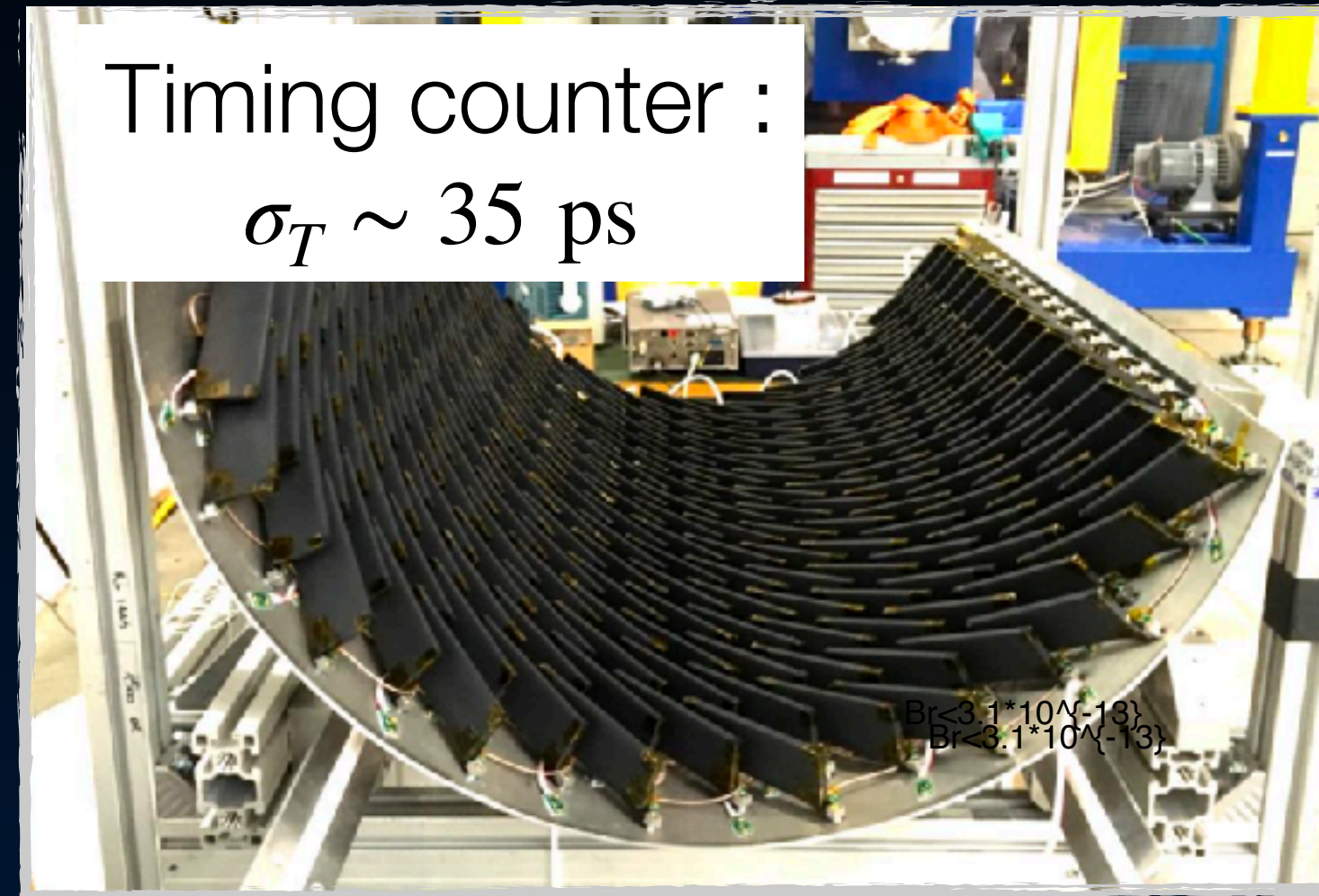


MEG II Detector Upgrade

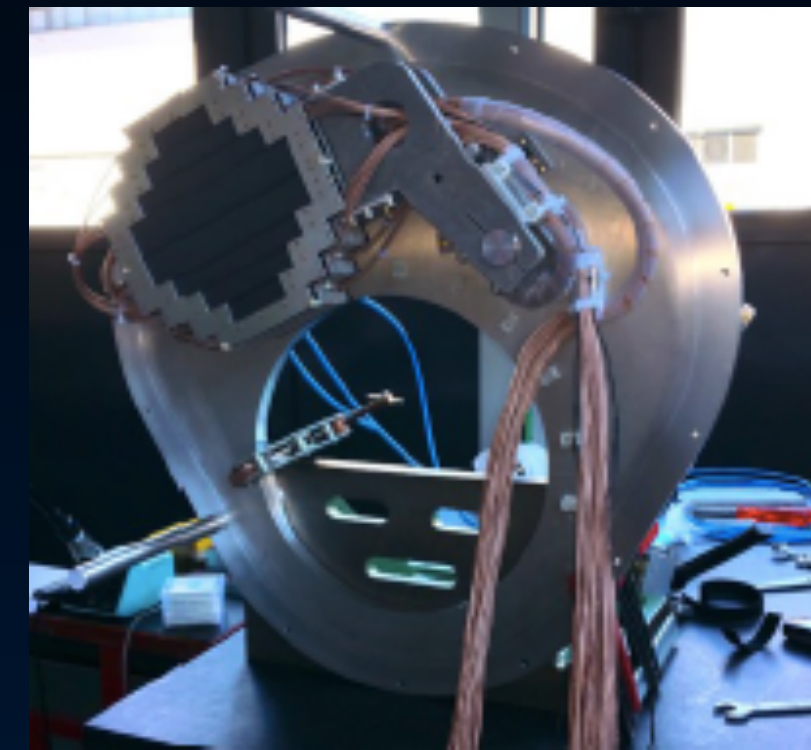
Liquid Xenon Photon Detector



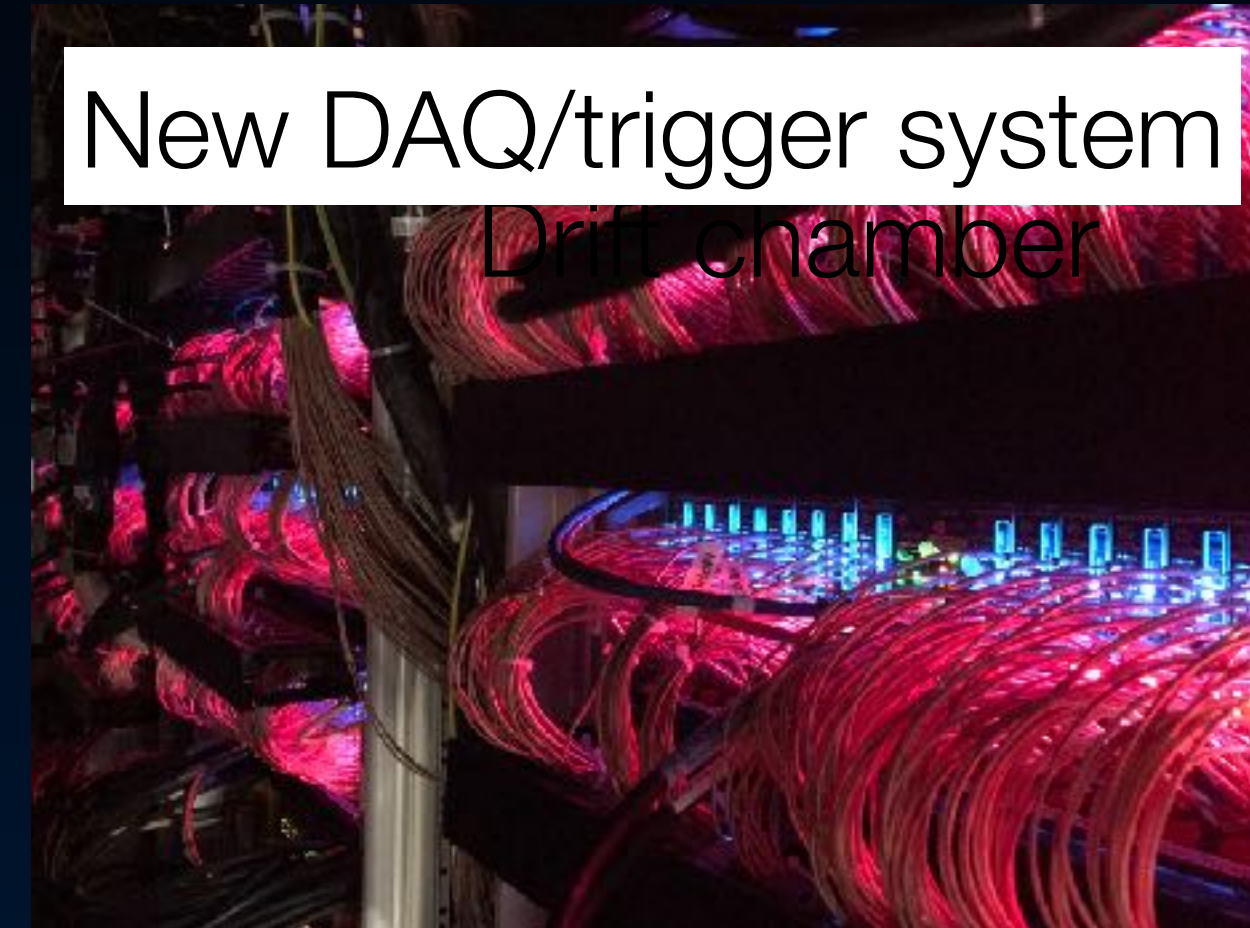
Timing counter :
 $\sigma_T \sim 35$ ps



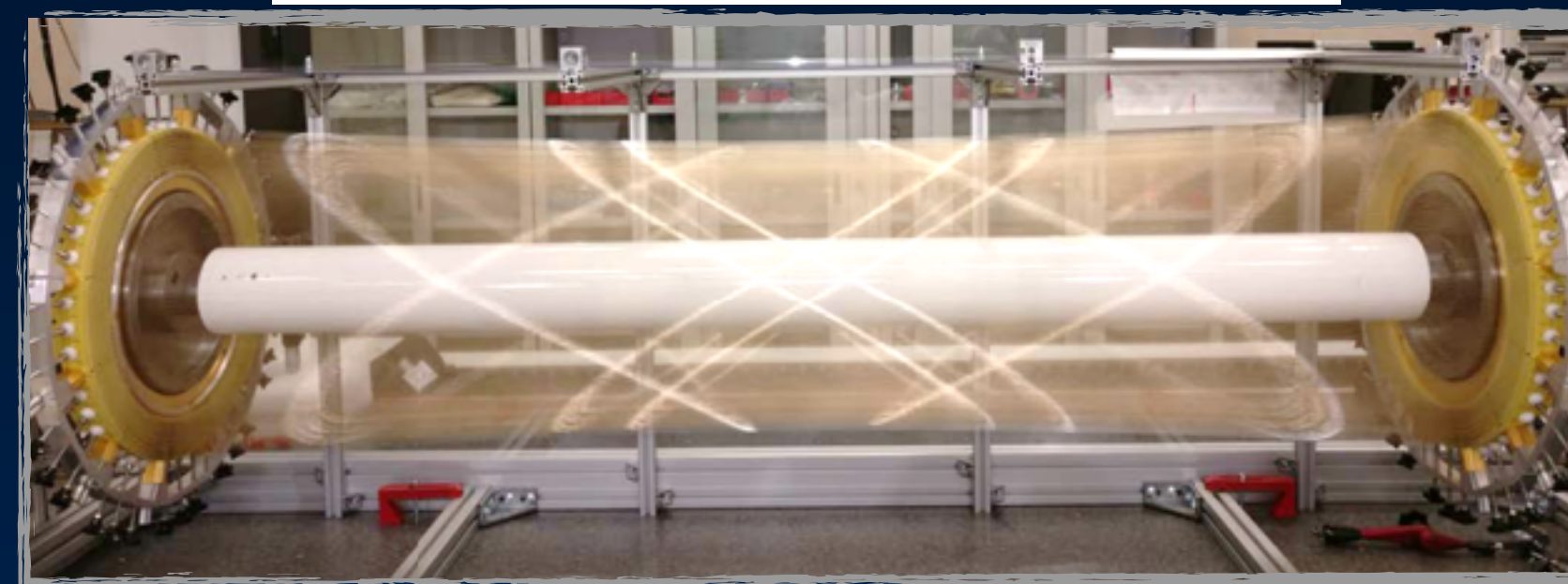
RDC detector



New DAQ/trigger system
Drift chamber

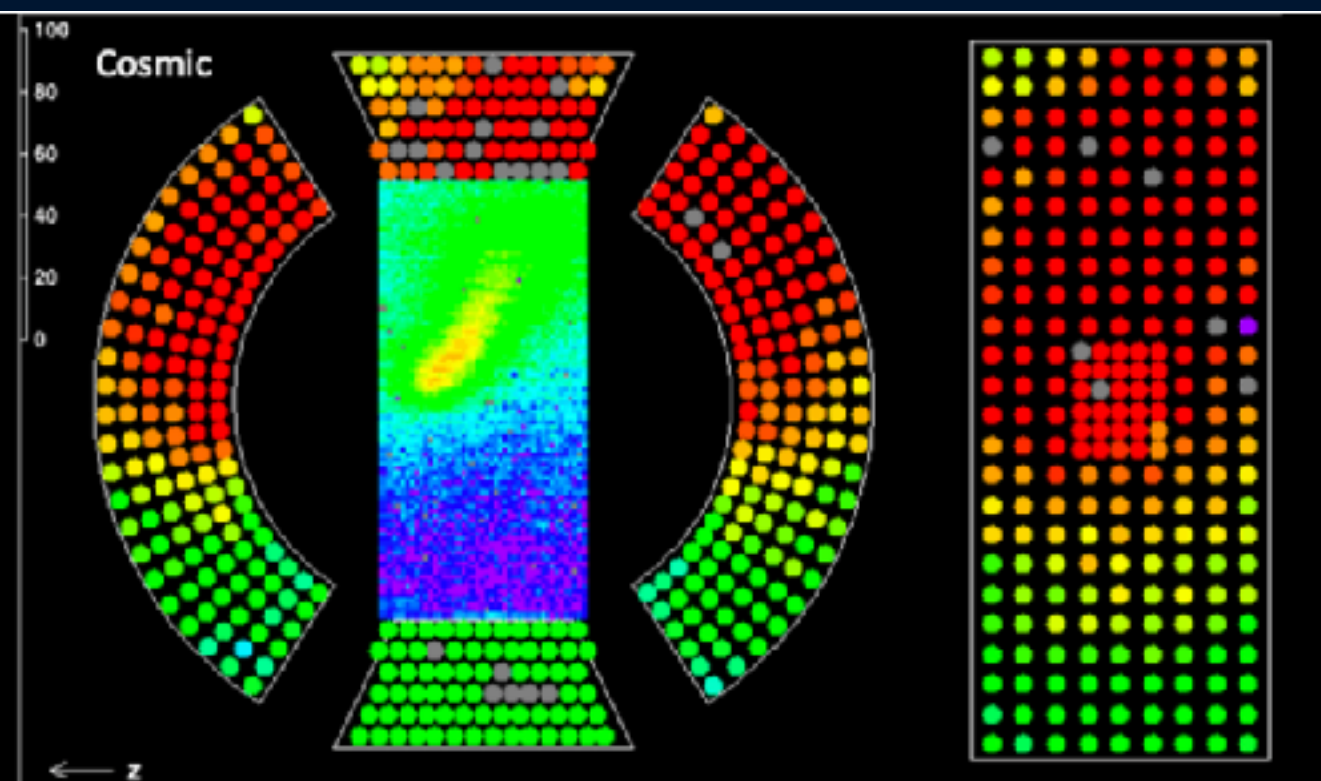


Cylindrical drift chamber



Timeline

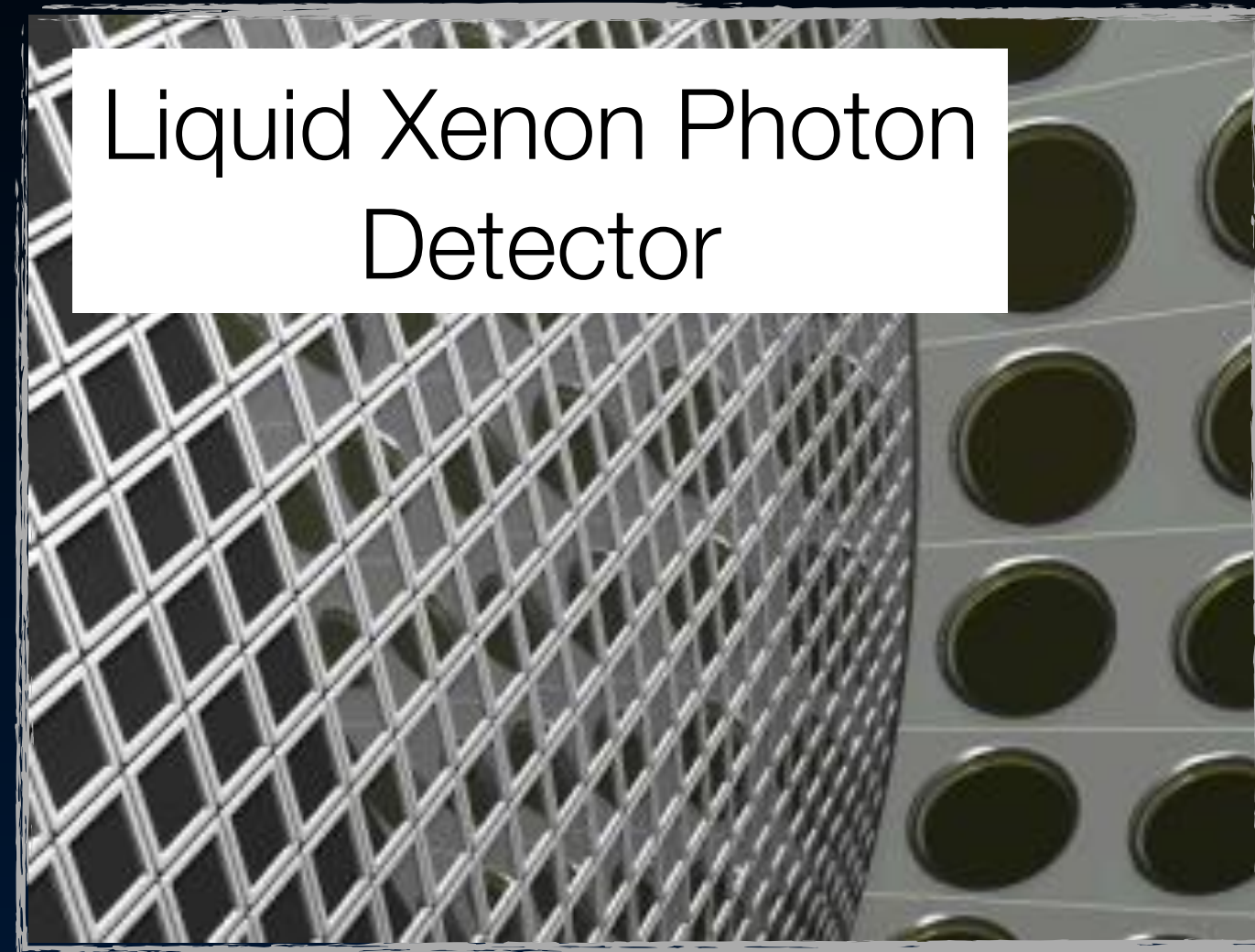
- Physics runs in 2021, 2022, 2023 - 2025, (2026)
- PSI HIPA accelerator shutdown in 2027-2028



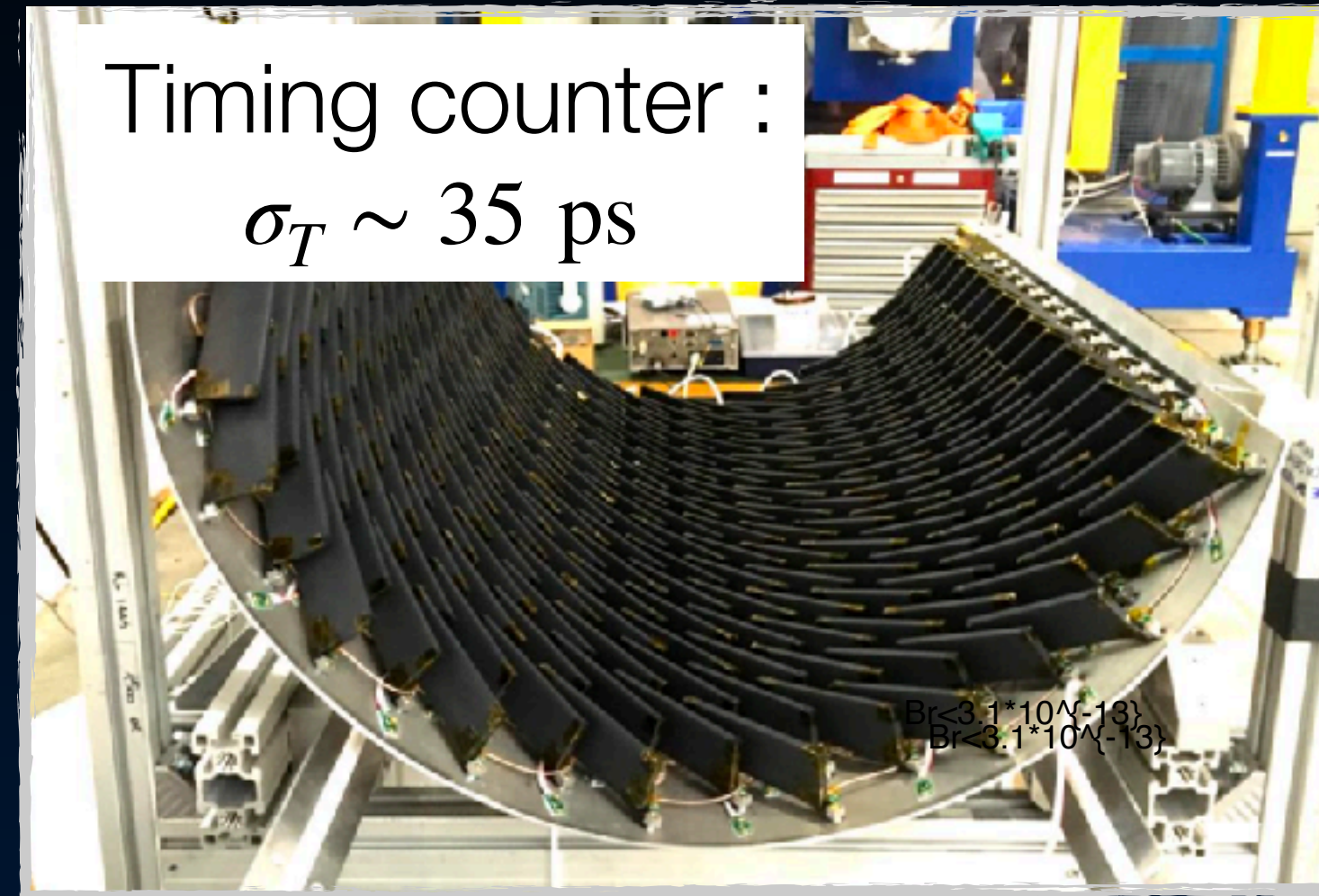
MEG II first result (10.2023): $BR < 7.5 \times 10^{-13}$

MEG II Detector Upgrade

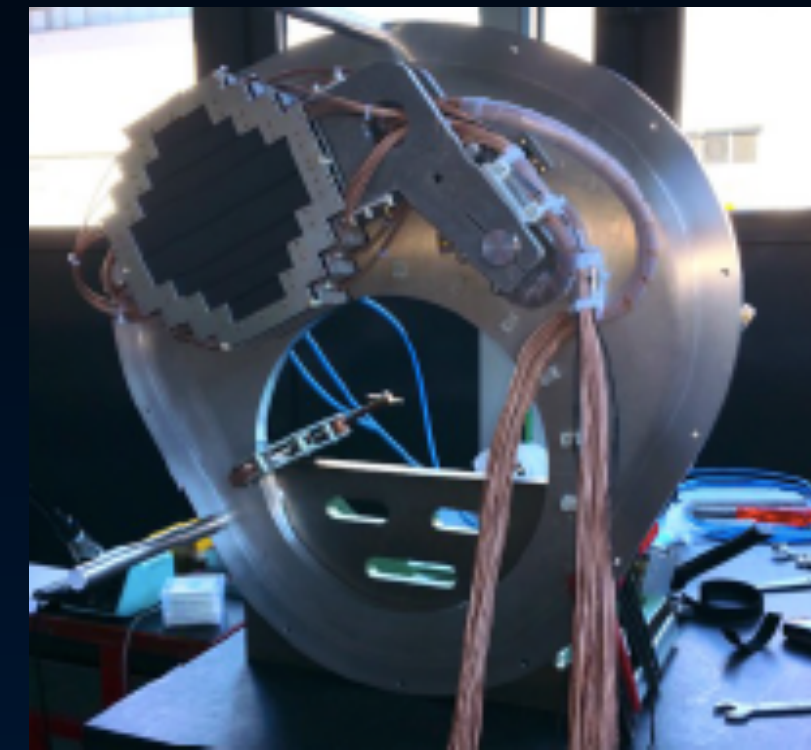
Liquid Xenon Photon Detector



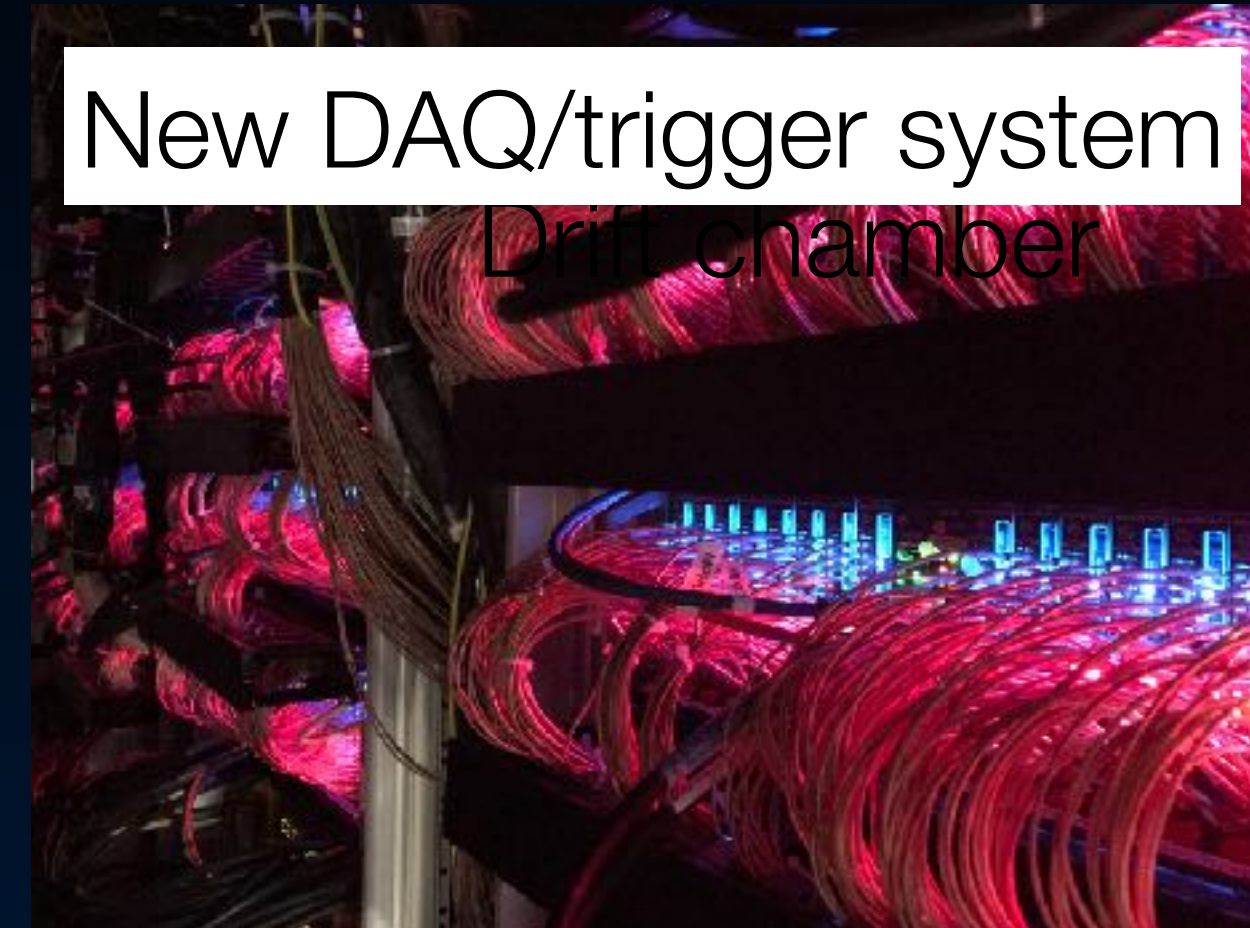
Timing counter :
 $\sigma_T \sim 35$ ps



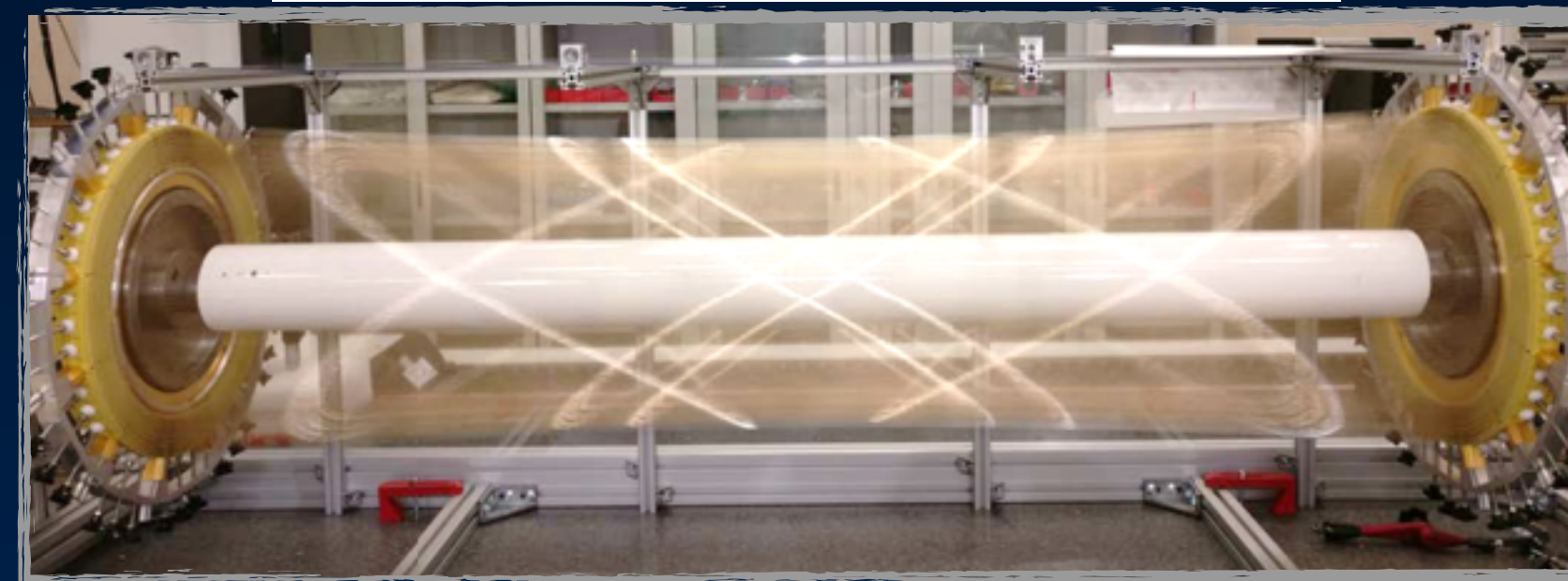
RDC detector



New DAQ/trigger system
Drift chamber

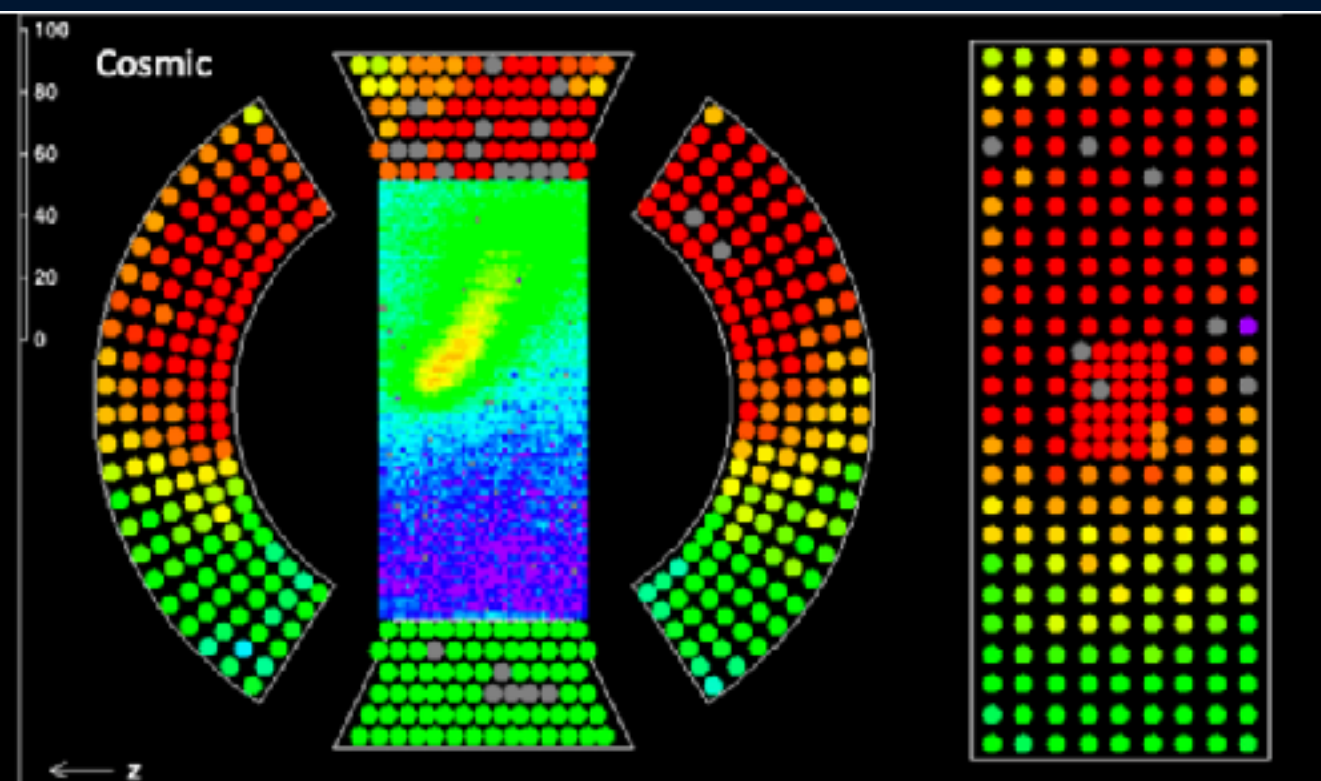


Cylindrical drift chamber



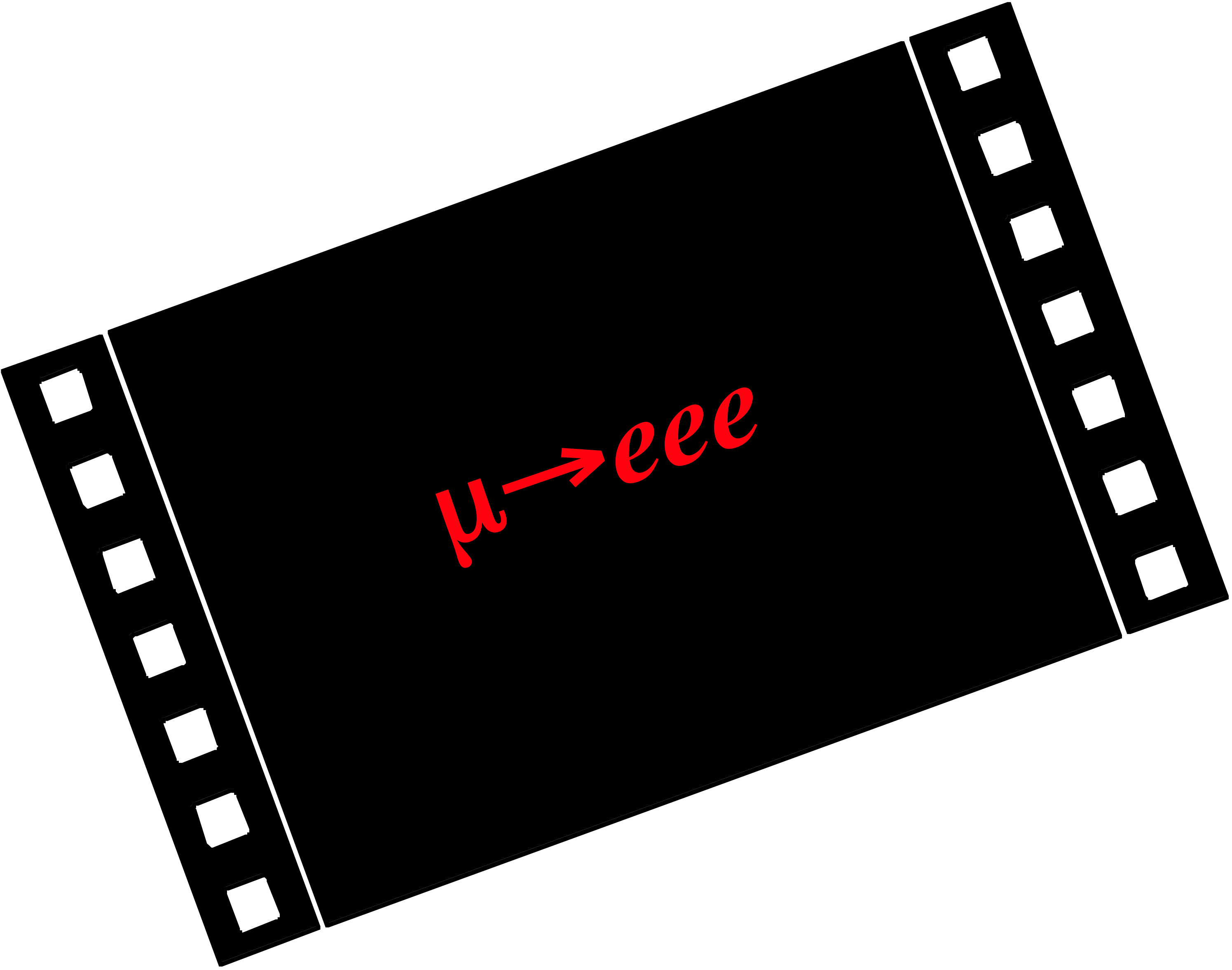
Timeline

- Physics runs in 2021, 2022, 2023 - 2025, (2026)
- PSI HIPA accelerator shutdown in 2027-2028



MEG II first result (10.2023): $BR < 7.5 \times 10^{-13}$

combined: $BR < 3.1 \times 10^{-13}$





$$\mu^+ \rightarrow e^+ e^+ e^- : \text{Mu3e}$$

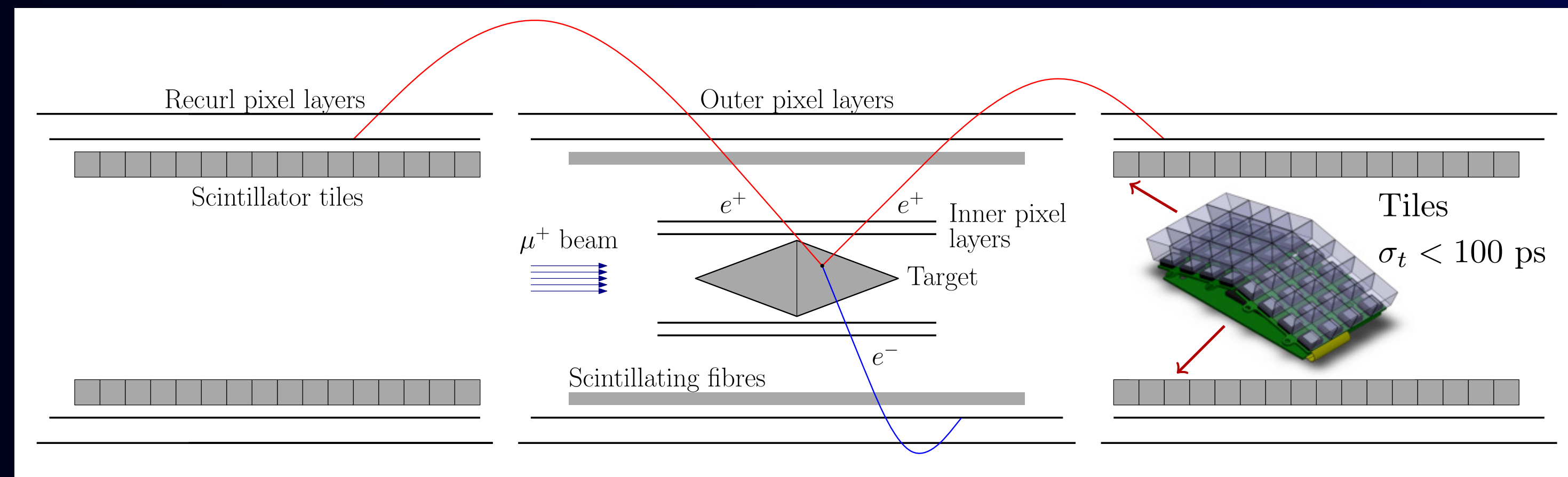


- Event Signature (μ^+ decay at rest)
 - $\Sigma E_e = m_\mu$ and $\Sigma P_e = 0$ (vector sum)
 - common vertex and time coincidence
- Backgrounds
 - Physics backgrounds, $\mu^+ \rightarrow e^+ \nu_e \bar{\nu}_\mu e^+ e^-$
 - Accidental backgrounds from Michel decays + Bhabha scattering
- Current limits (from SINDRUM at PSI)
 - $B(\mu \rightarrow eee) < 1.0 \times 10^{-12}$ (90% C.L.)
- Mu3e at PSI
 - Phase-I: $\mathcal{O}(10^{-15})$ with $10^8 \mu/s$ ($\pi E5$)
 - inner pixel detectors, scintillating fibers
 - Phase-II: $\mathcal{O}(10^{-16})$ with $10^9 \mu/s$ (HiMB)

$\mu^+ \rightarrow e^+ e^+ e^-$: Mu3e



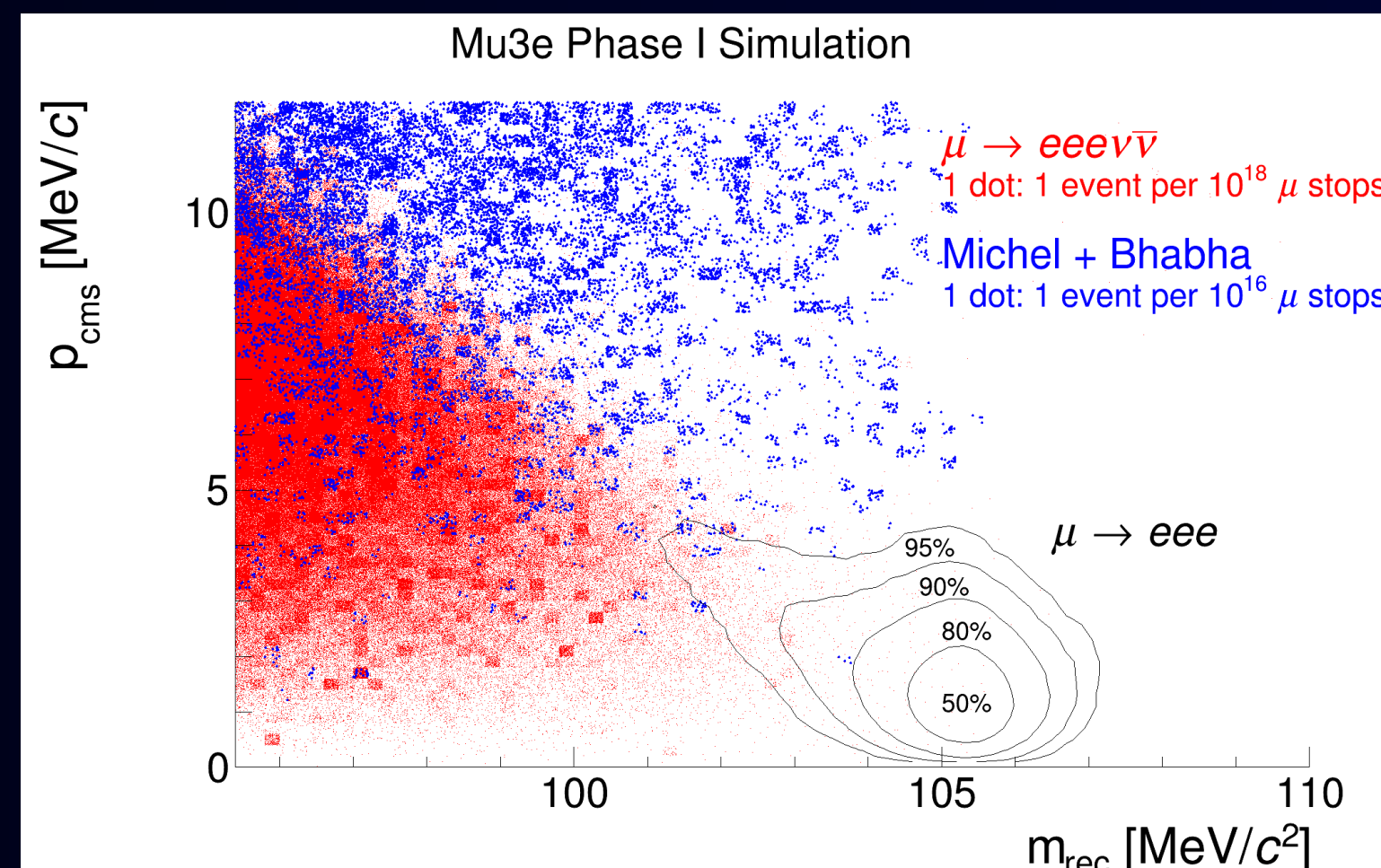
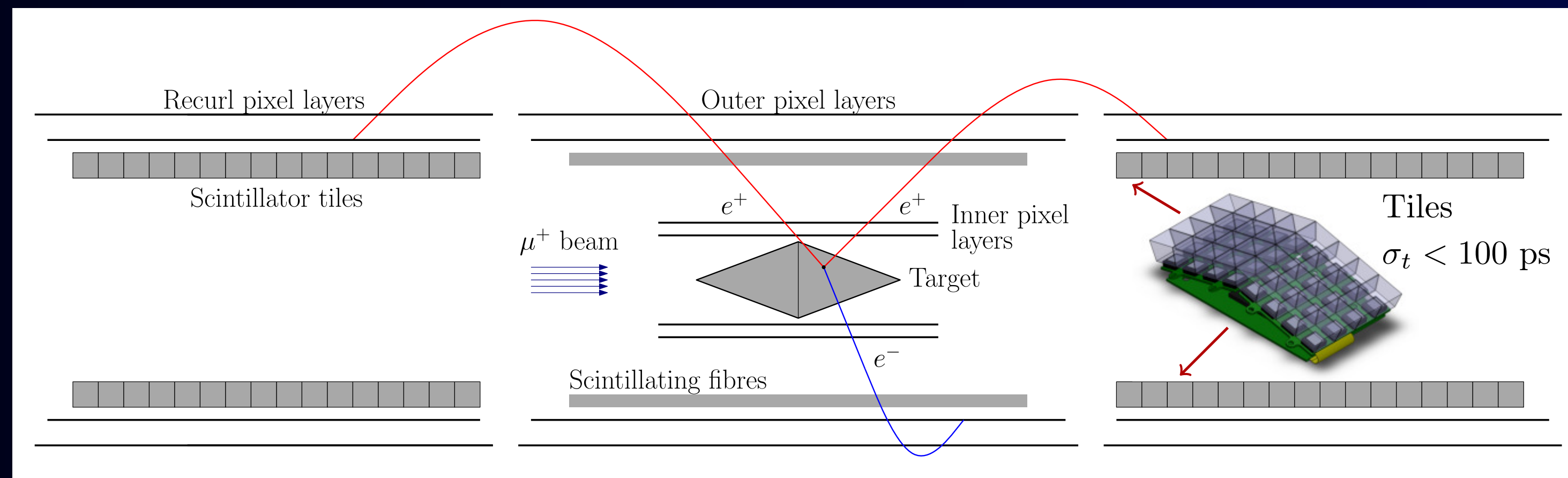
- **Event Signature (μ^+ decay at rest)**
 - $\Sigma E_e = m_\mu$ and $\Sigma P_e = 0$ (vector sum)
 - common vertex and time coincidence
- **Backgrounds**
 - Physics backgrounds, $\mu^+ \rightarrow e^+ \nu_e \bar{\nu}_\mu e^+ e^-$
 - Accidental backgrounds from Michel decays + Bhabha scattering
- **Current limits (from SINDRUM at PSI)**
 - $B(\mu \rightarrow eee) < 1.0 \times 10^{-12}$ (90% C.L.)
- **Mu3e at PSI**
 - Phase-I: $\mathcal{O}(10^{-15})$ with $10^8 \mu/s$ ($\pi E5$)
 - inner pixel detectors, scintillating fibers
 - Phase-II: $\mathcal{O}(10^{-16})$ with $10^9 \mu/s$ (HiMB)



$\mu^+ \rightarrow e^+ e^+ e^-$: Mu3e



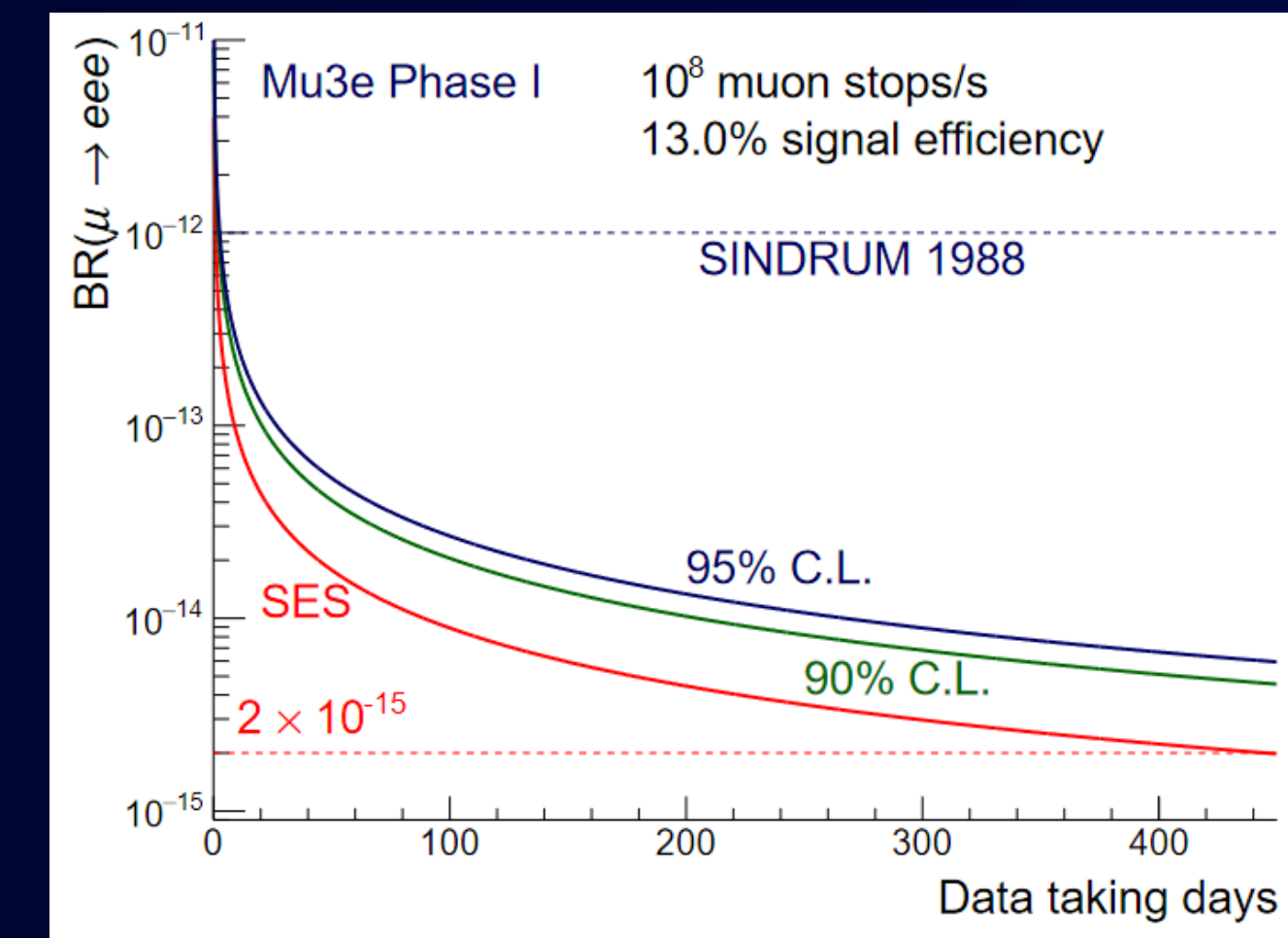
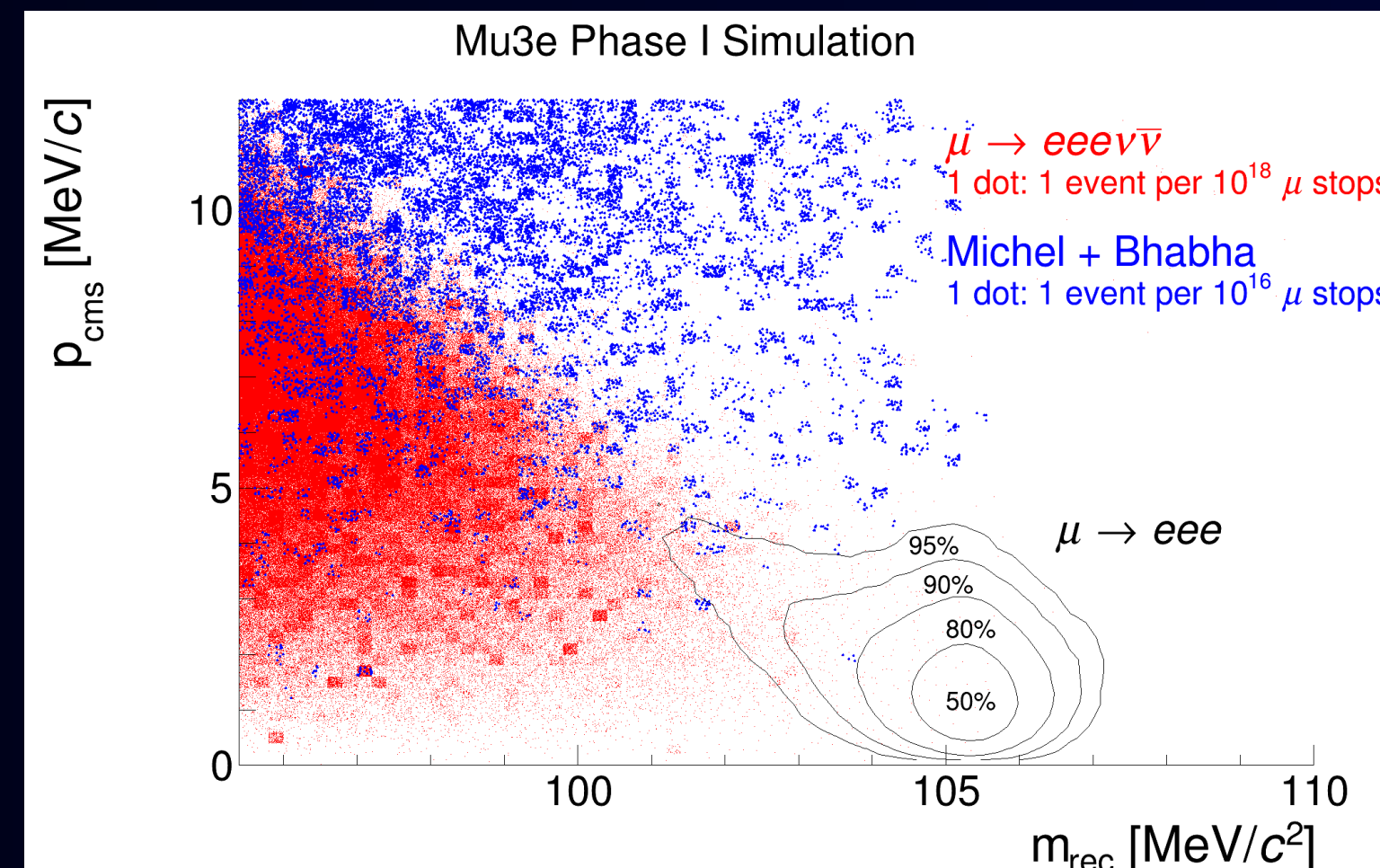
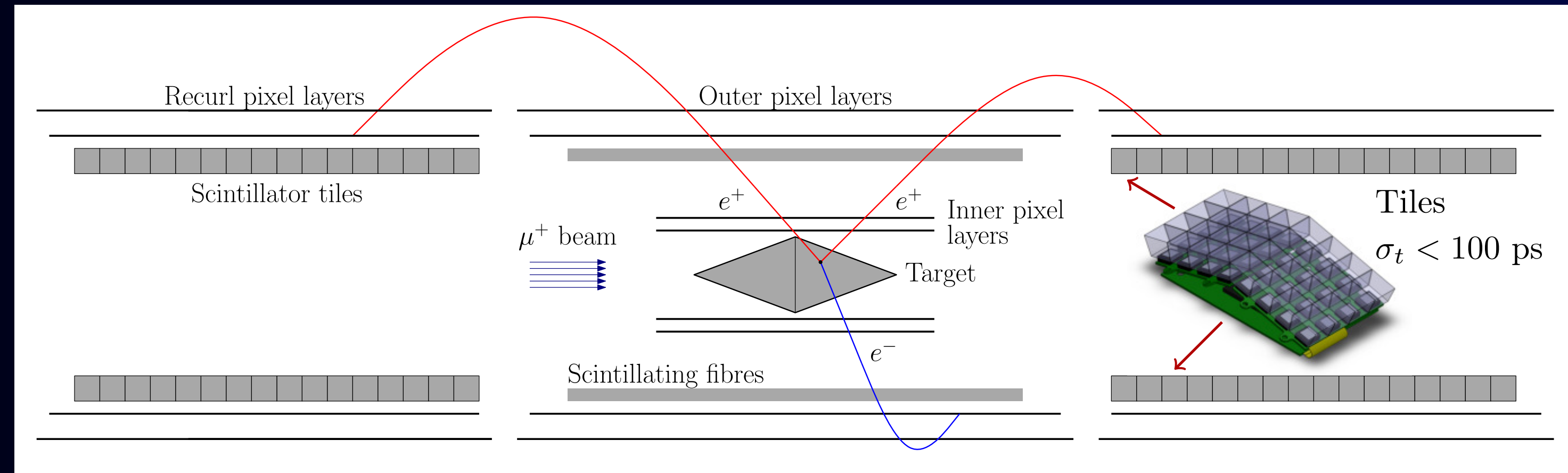
- **Event Signature (μ^+ decay at rest)**
 - $\Sigma E_e = m_\mu$ and $\Sigma P_e = 0$ (vector sum)
 - common vertex and time coincidence
- **Backgrounds**
 - Physics backgrounds, $\mu^+ \rightarrow e^+ \nu_e \bar{\nu}_\mu e^+ e^-$
 - Accidental backgrounds from Michel decays + Bhabha scattering
- **Current limits (from SINDRUM at PSI)**
 - $B(\mu \rightarrow eee) < 1.0 \times 10^{-12}$ (90% C.L.)
- **Mu3e at PSI**
 - Phase-I: $\mathcal{O}(10^{-15})$ with $10^8 \mu/s$ ($\pi E5$)
 - inner pixel detectors, scintillating fibers
 - Phase-II: $\mathcal{O}(10^{-16})$ with $10^9 \mu/s$ (HiMB)



$\mu^+ \rightarrow e^+ e^+ e^-$: Mu3e



- **Event Signature** (μ^+ decay at rest)
 - $\Sigma E_e = m_\mu$ and $\Sigma P_e = 0$ (vector sum)
 - common vertex and time coincidence
- **Backgrounds**
 - Physics backgrounds, $\mu^+ \rightarrow e^+ \nu_e \bar{\nu}_\mu e^+ e^-$
 - Accidental backgrounds from Michel decays + Bhabha scattering
- **Current limits (from SINDRUM at PSI)**
 - $B(\mu \rightarrow eee) < 1.0 \times 10^{-12}$ (90% C.L.)
- **Mu3e at PSI**
 - Phase-I: $\mathcal{O}(10^{-15})$ with $10^8 \mu/s$ ($\pi E5$)
 - inner pixel detectors, scintillating fibers
 - Phase-II: $\mathcal{O}(10^{-16})$ with $10^9 \mu/s$ (HiMB)



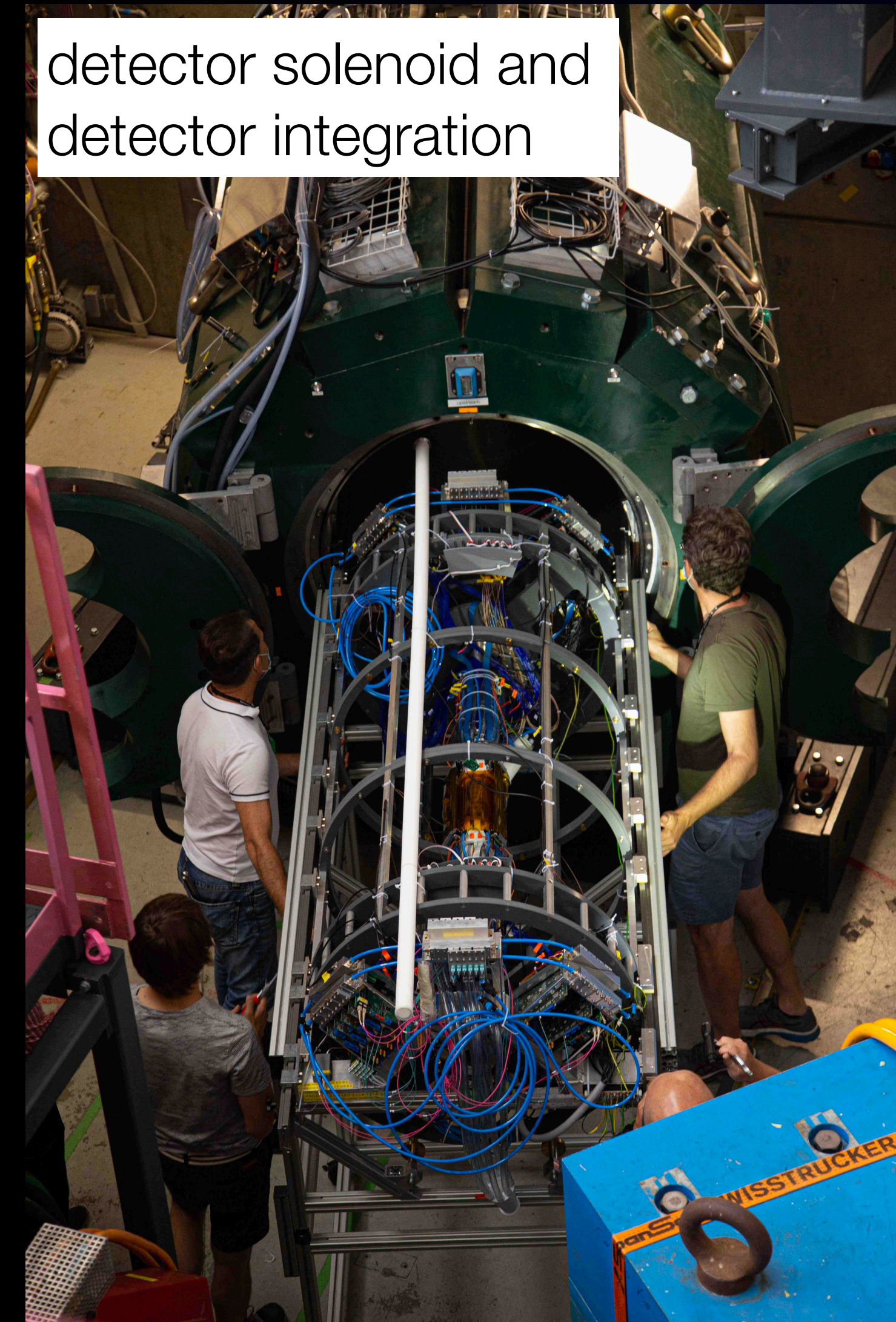
Mu3e Phase-I Preparation Status



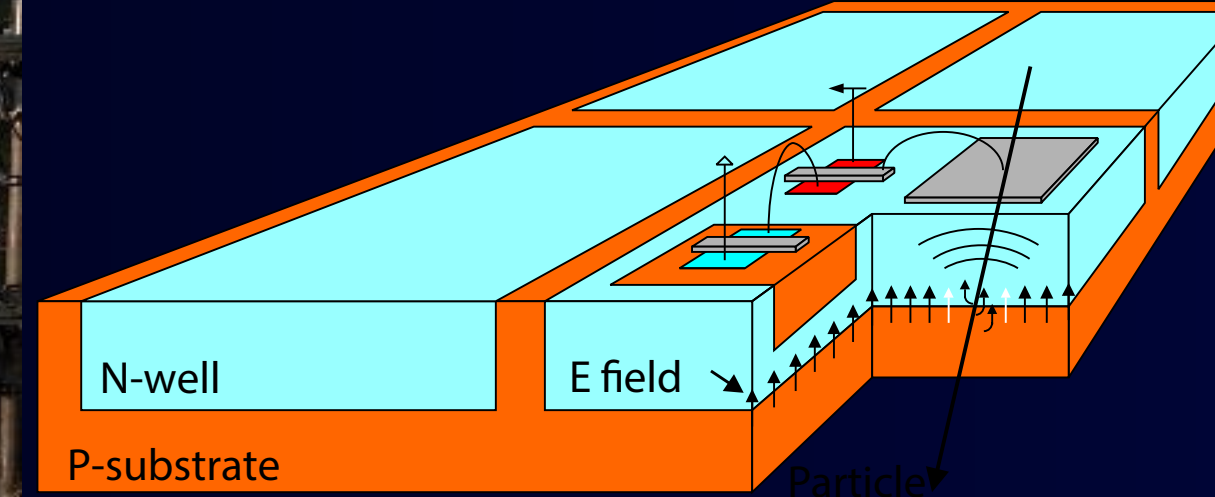
Mu3e Phase-I Preparation Status



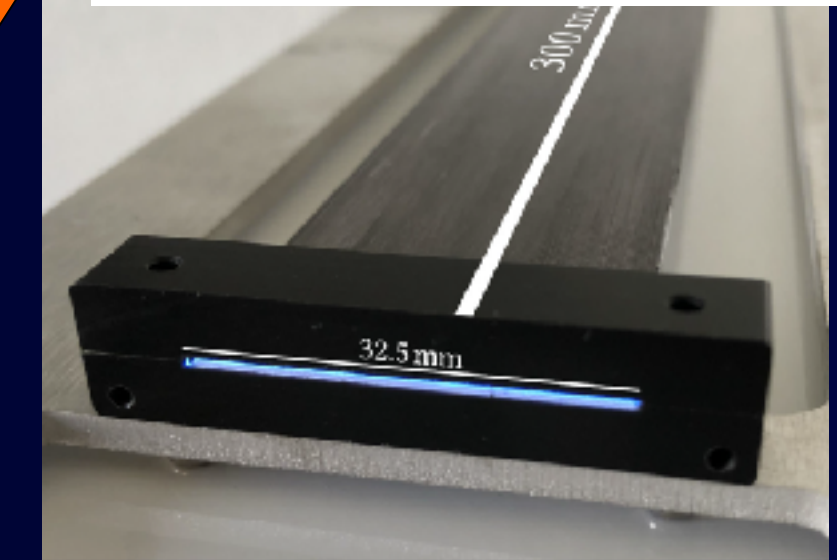
detector solenoid and detector integration



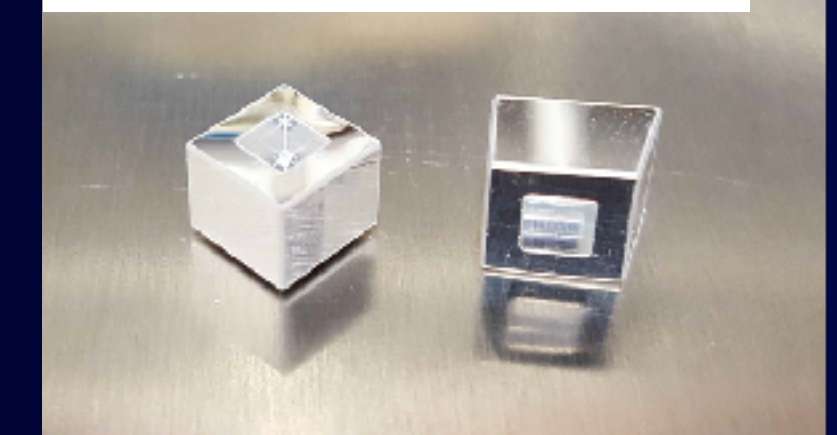
silicon pixel tracker



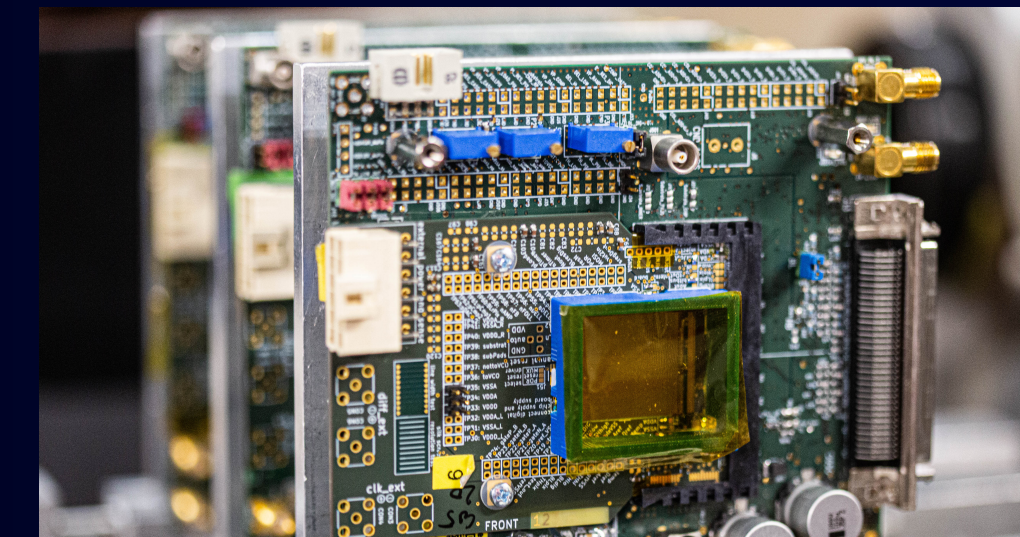
SciFi timing counter



scin. tile counter



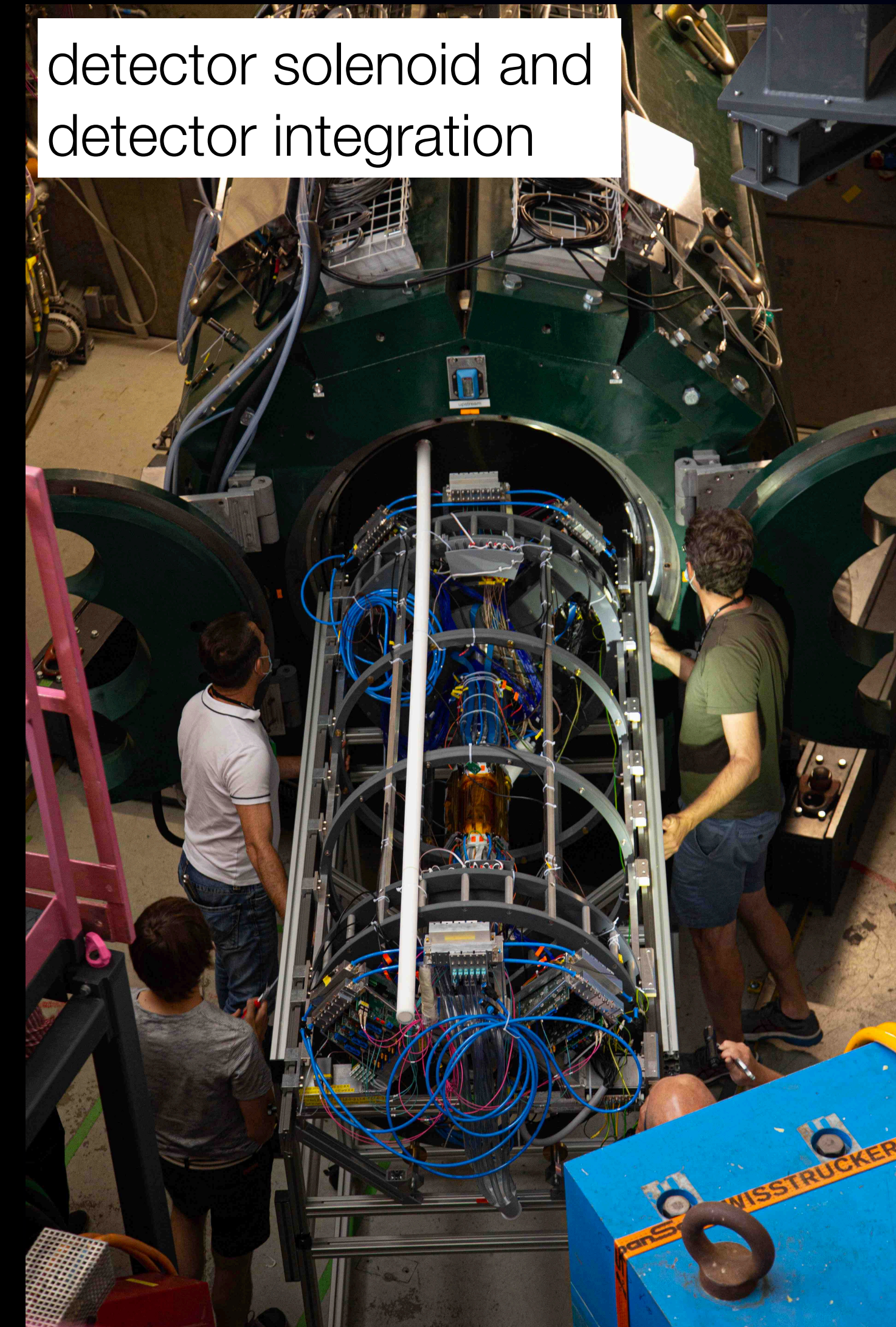
silicon pixel board (50 μm)



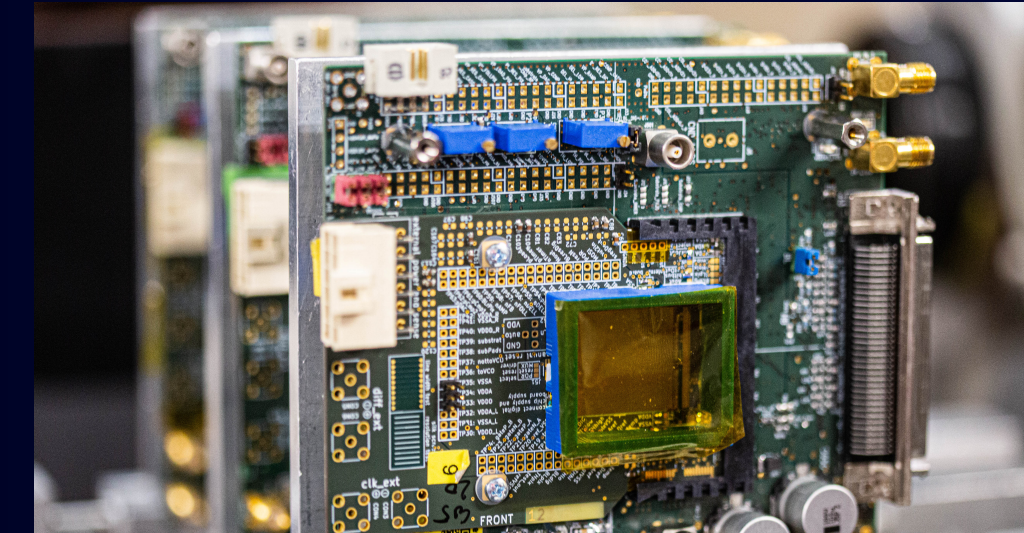
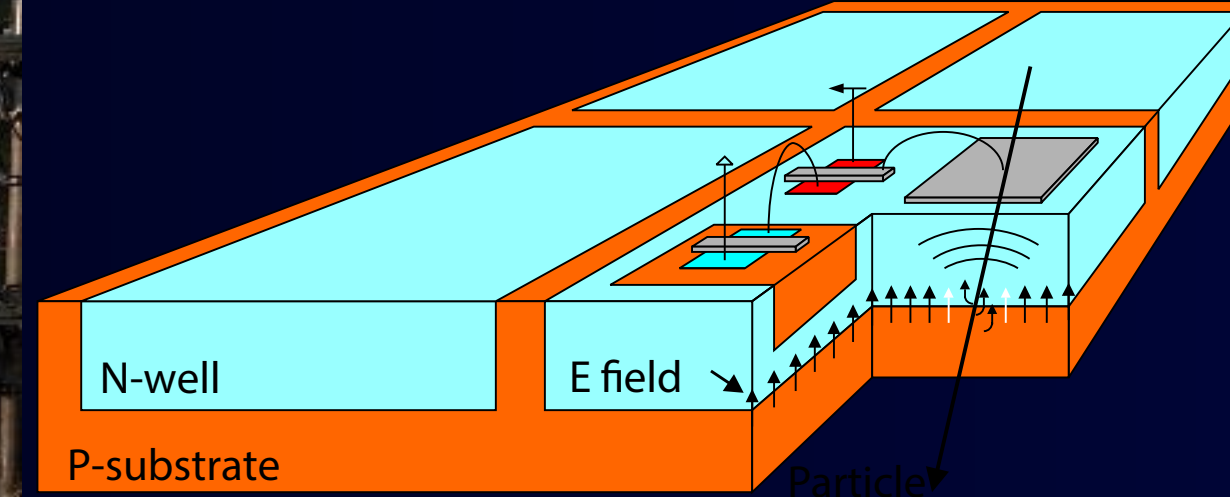
Mu3e Phase-I Preparation Status



detector solenoid and detector integration

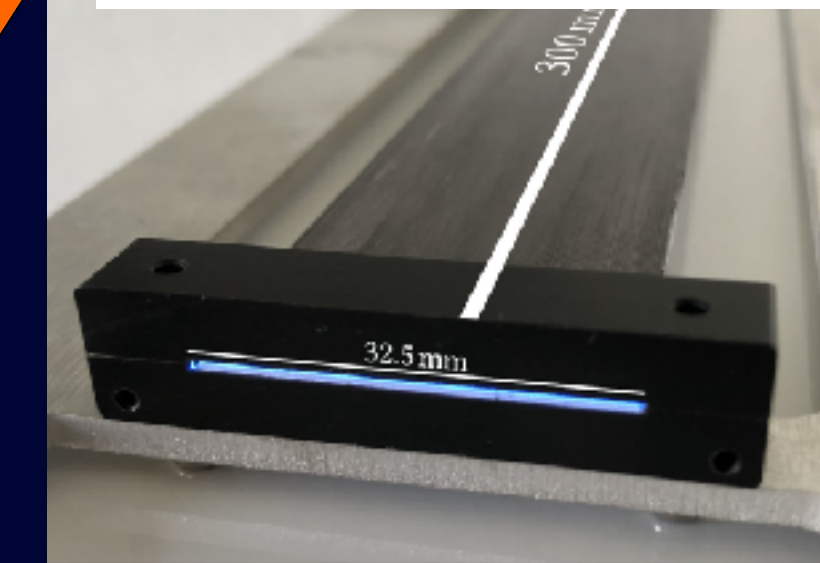


silicon pixel tracker

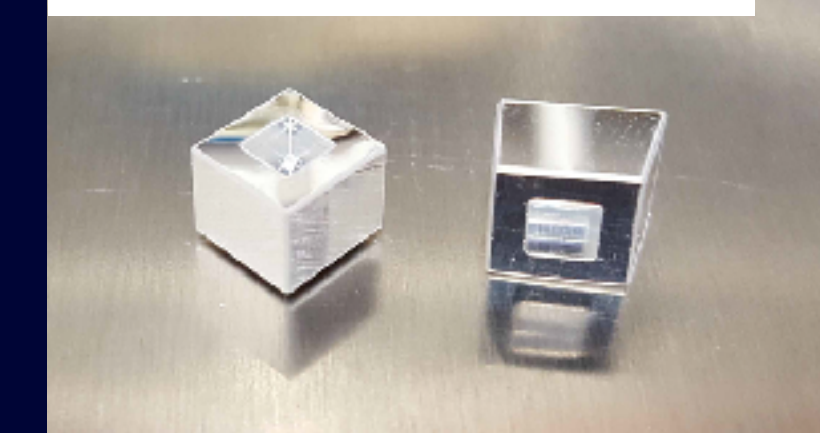


silicon pixel board (50 μm)

SciFi timing counter



scin. tile counter



- Timeline**
- First detector installation in 2023
 - Detector commissioning in 2024/2025
 - First data taking in 2025/2026

Mu3e Phase II



Mu3e Phase II

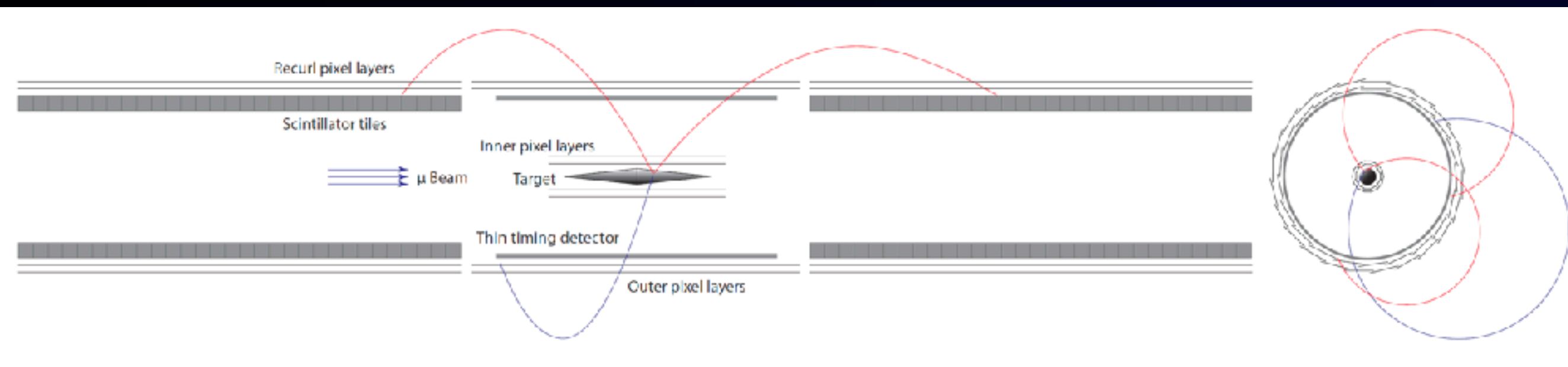


Mu3e Phase II

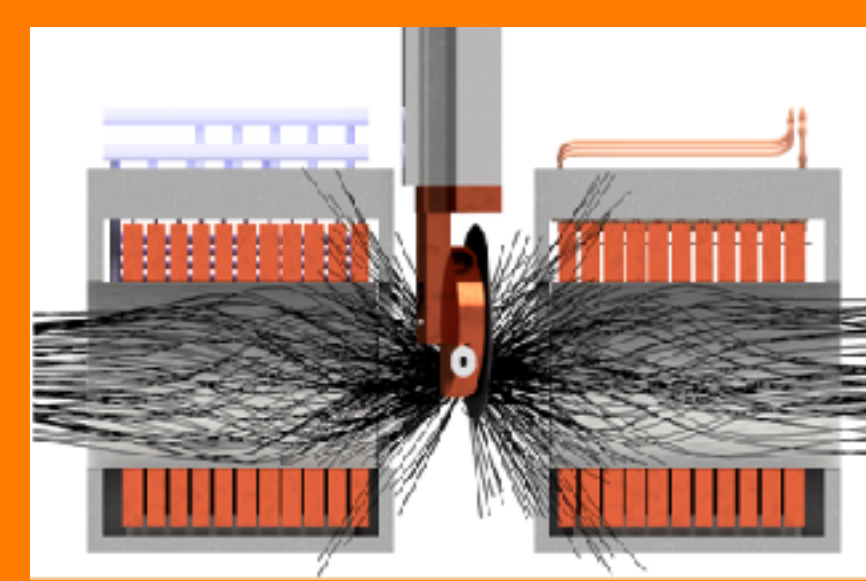
- Ultimate sensitivity goal of $BR < 1 \times 10^{-16}$
- muon intensity 2×10^9 /sec from **HiMB**
- Upgraded Mu3e detector
 - elongated pixel station for recurl tracks
 - muon target with smaller radius
 - thinner pixel detector
- scheduled after 2029

HiMB

- High Intensity Muon Beamline (HiMB) at PSI
- Surface muon (μ^+) beam, $\sim O(10^{10})$ /s
- New target and new capturing solenoids
- Installation in 2027-2028
- Planned to be operational in 2029

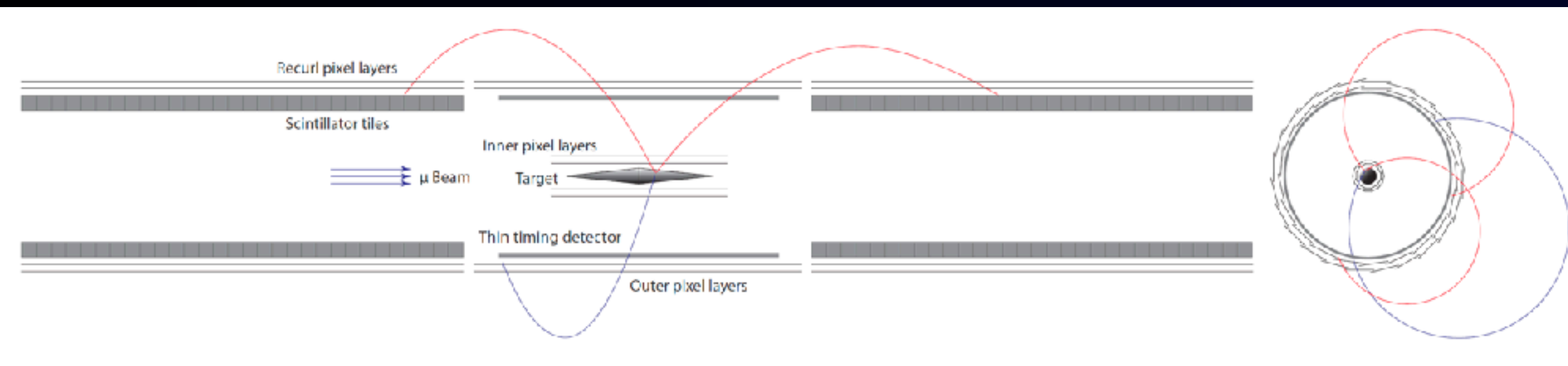


Mu3e Phase II



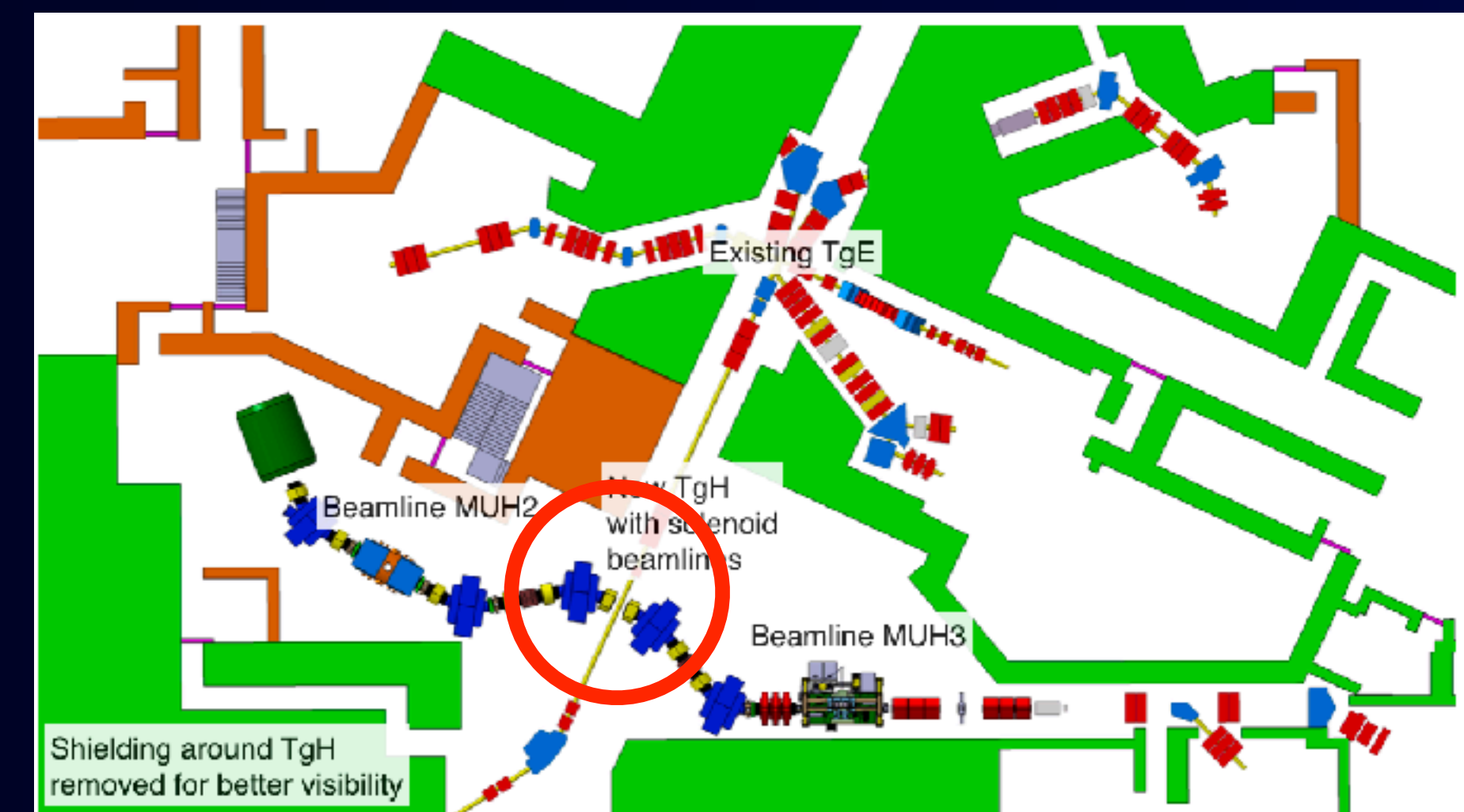
Mu3e Phase II


- Ultimate sensitivity goal of $BR < 1 \times 10^{-16}$
- muon intensity 2×10^9 /sec from **HiMB**
- Upgraded Mu3e detector
 - elongated pixel station for recurl tracks
 - muon target with smaller radius
 - thinner pixel detector
- scheduled after 2029



HiMB

- High Intensity Muon Beamline (HiMB) at PSI
- Surface muon (μ^+) beam, $\sim O(10^{10})$ /s
- New target and new capturing solenoids
- Installation in 2027-2028
- Planned to be operational in 2029





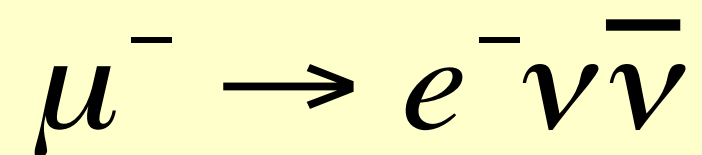
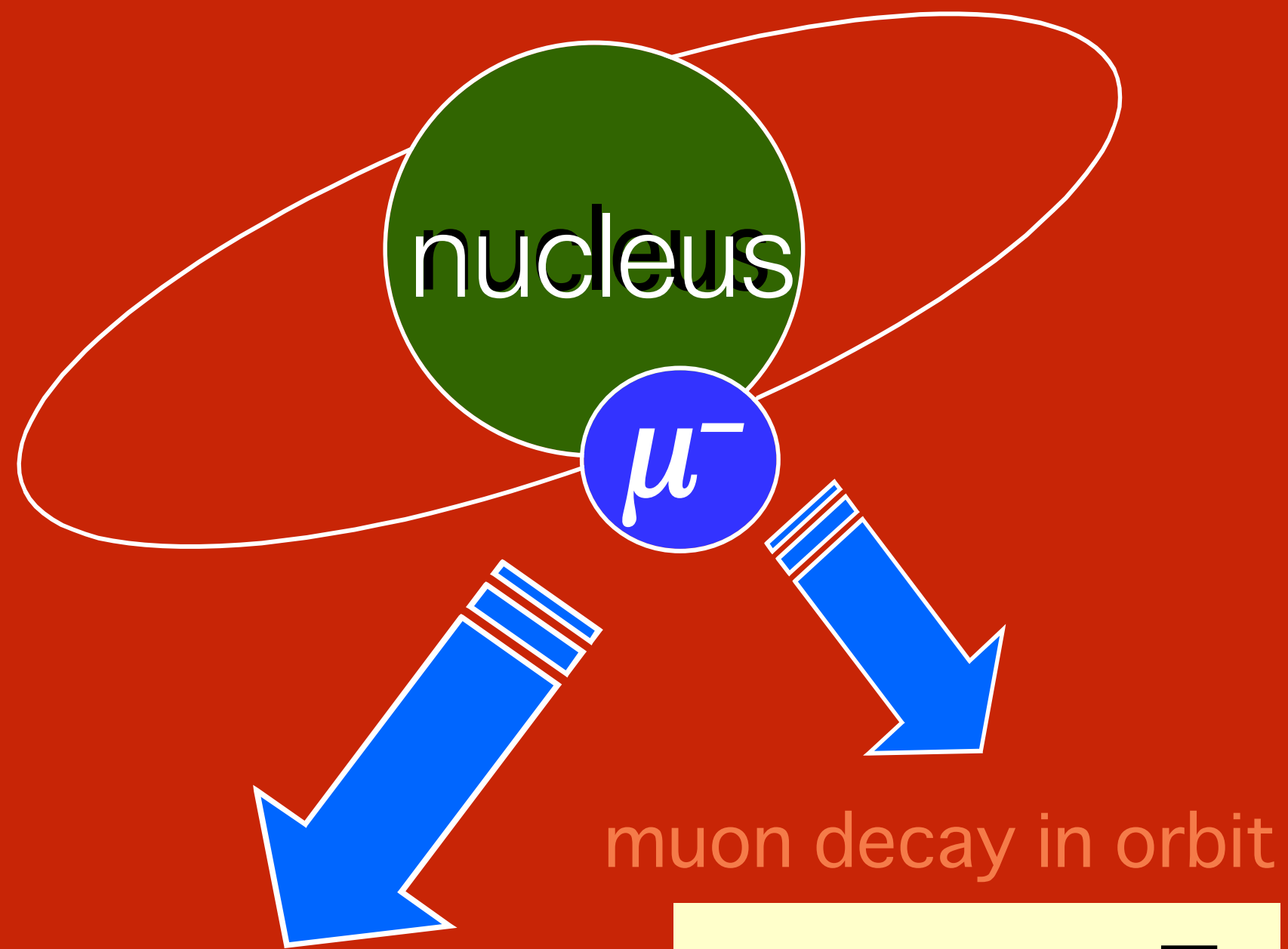
$\mu \rightarrow e$ conversion
in
a muonic atom

$\mu^- \rightarrow e^-$ Conversion

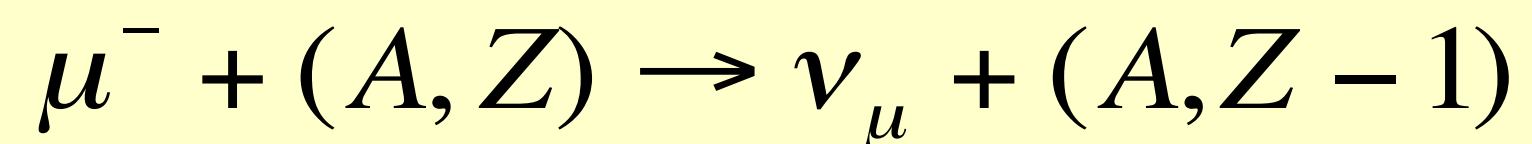


$\mu^- \rightarrow e^-$ Conversion

1s state in a muonic atom

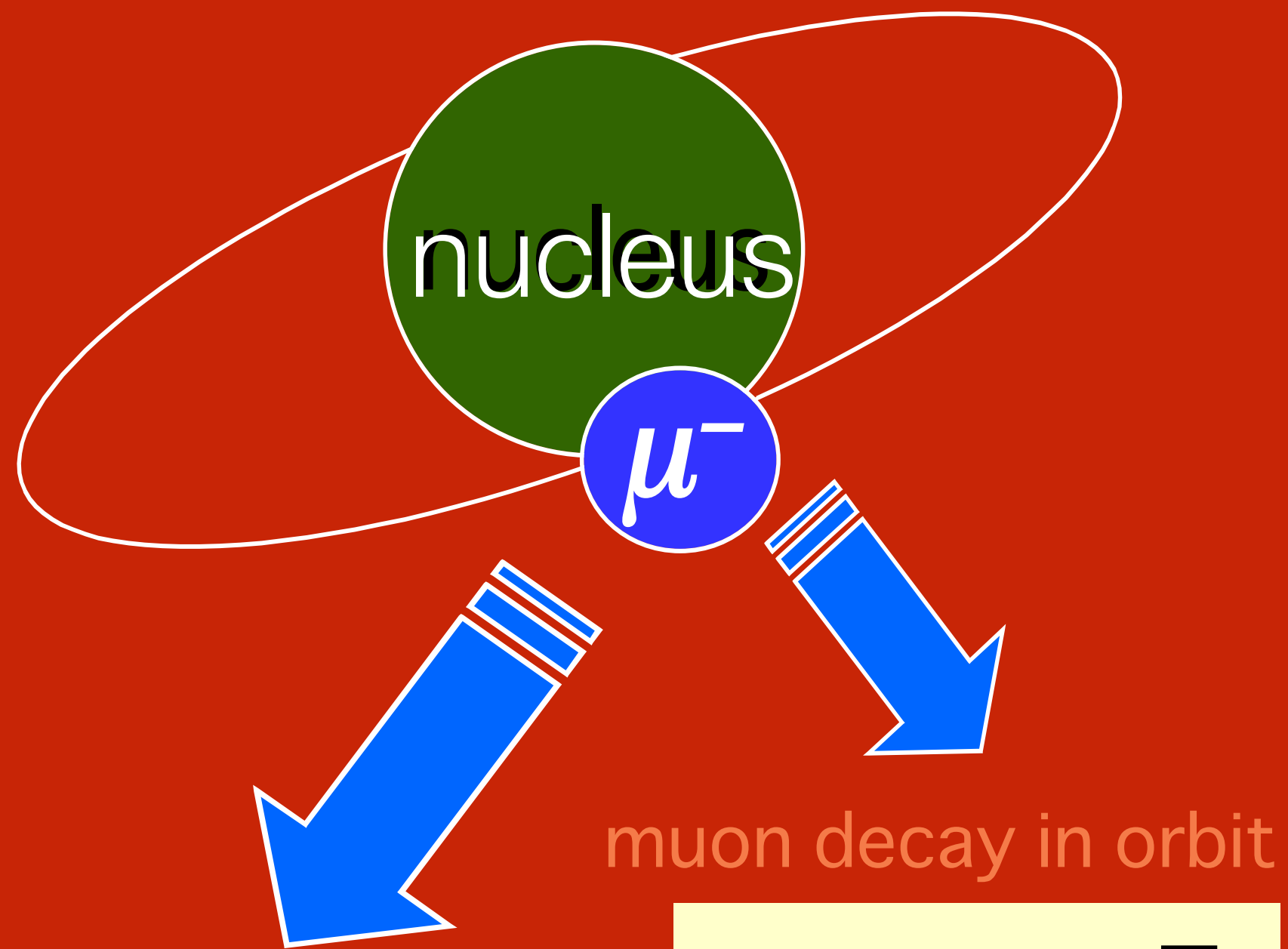


nuclear muon capture

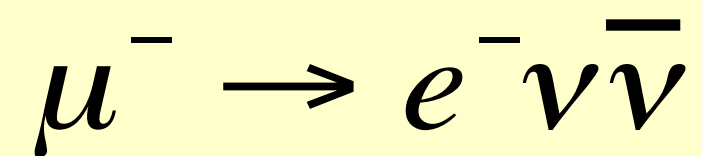


$\mu^- \rightarrow e^-$ Conversion

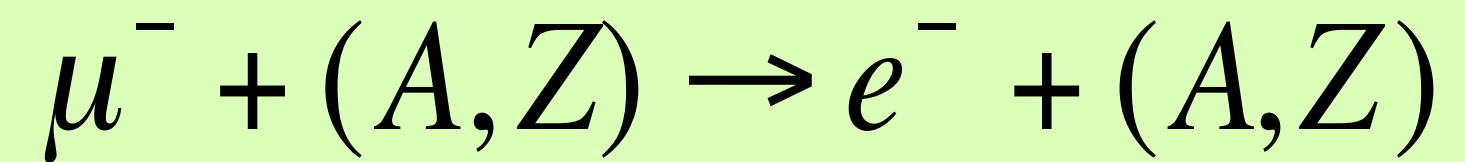
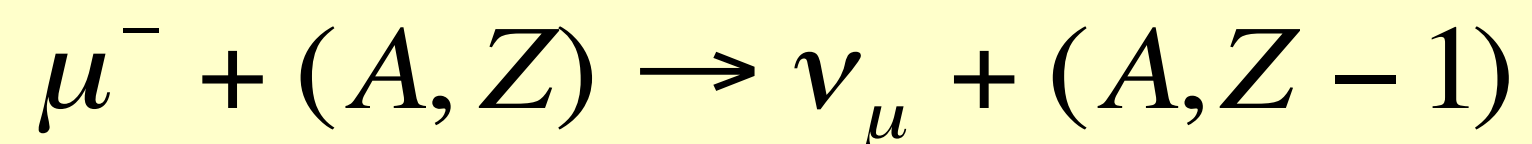
1s state in a muonic atom



muon decay in orbit

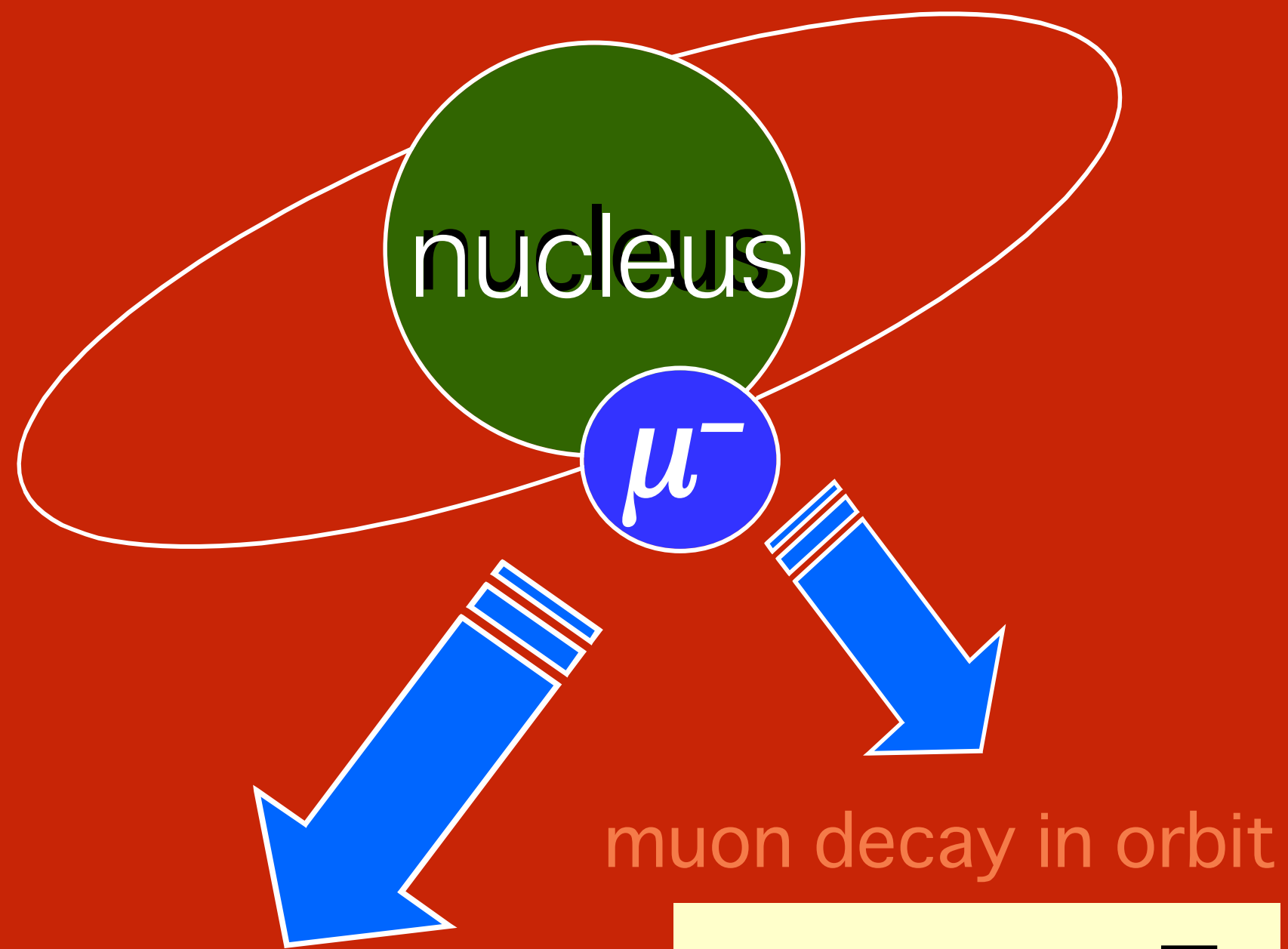


nuclear muon capture



$\mu^- \rightarrow e^-$ Conversion

1s state in a muonic atom



$$\mu^- \rightarrow e^- \nu \bar{\nu}$$

nuclear muon capture

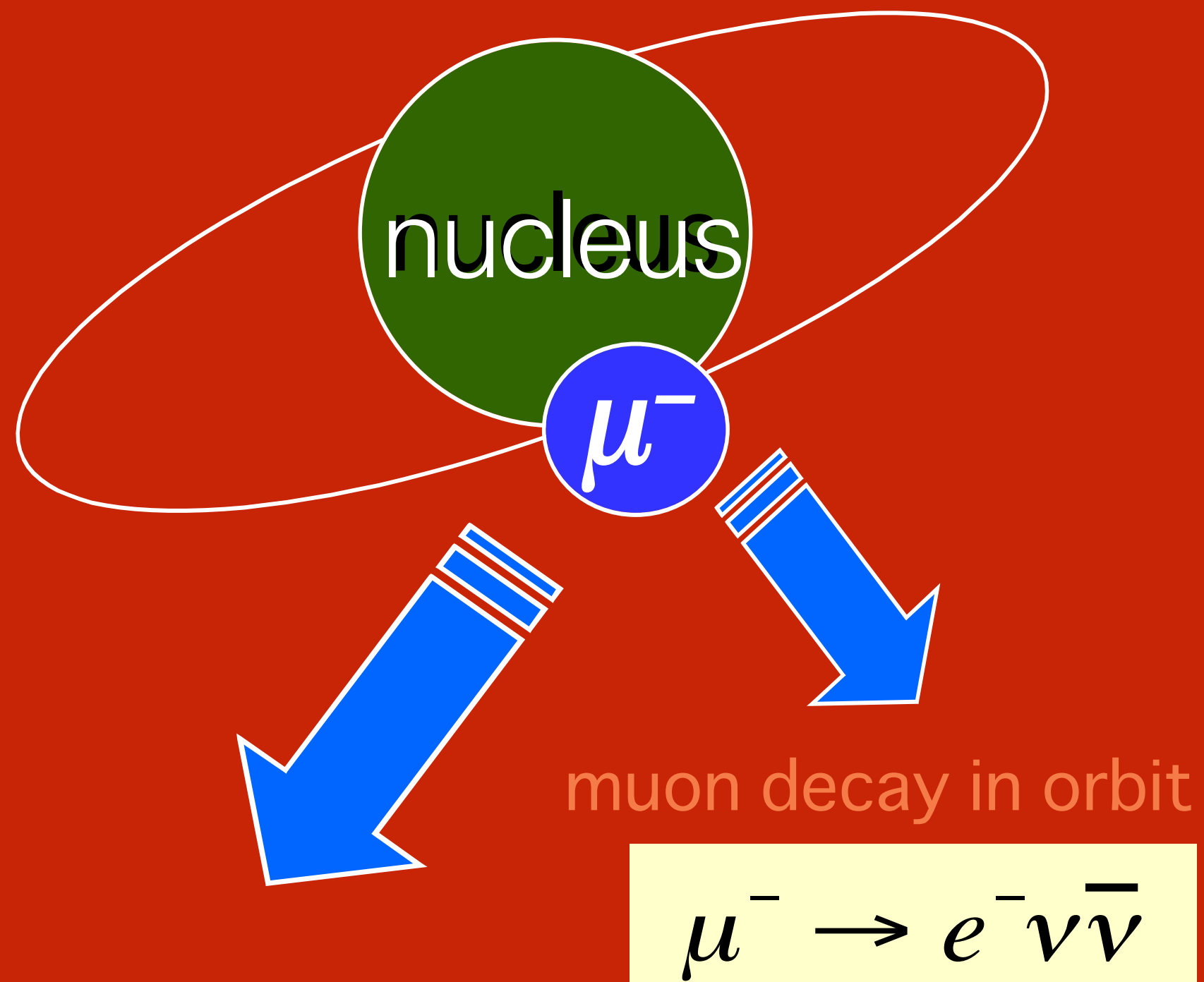
$$\mu^- + (A, Z) \rightarrow \nu_\mu + (A, Z - 1)$$

$$\mu^- + (A, Z) \rightarrow e^- + (A, Z)$$

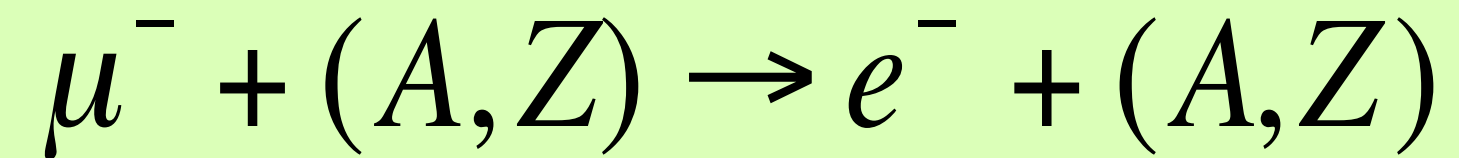
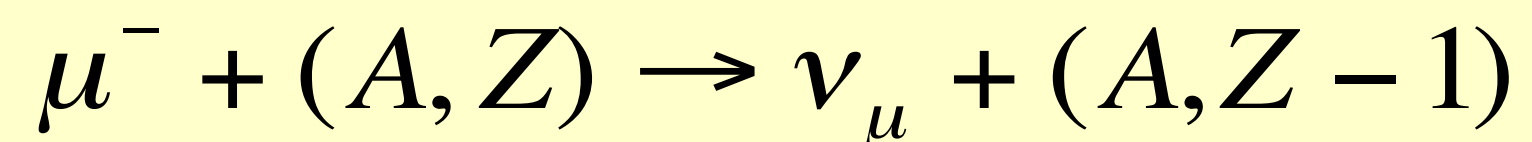
coherent process (for transition to ground state) $\propto Z^5$

$\mu^- \rightarrow e^-$ Conversion

1s state in a muonic atom



nuclear muon capture



coherent process (for transition to ground state) $\propto Z^5$

Event Signature :

a mono-energetic electron

(one particle measurement allows higher muon rates.)

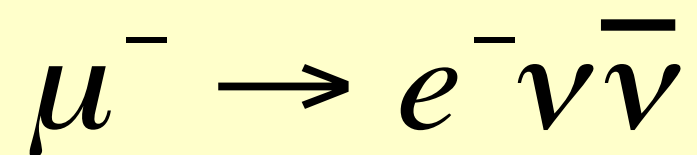
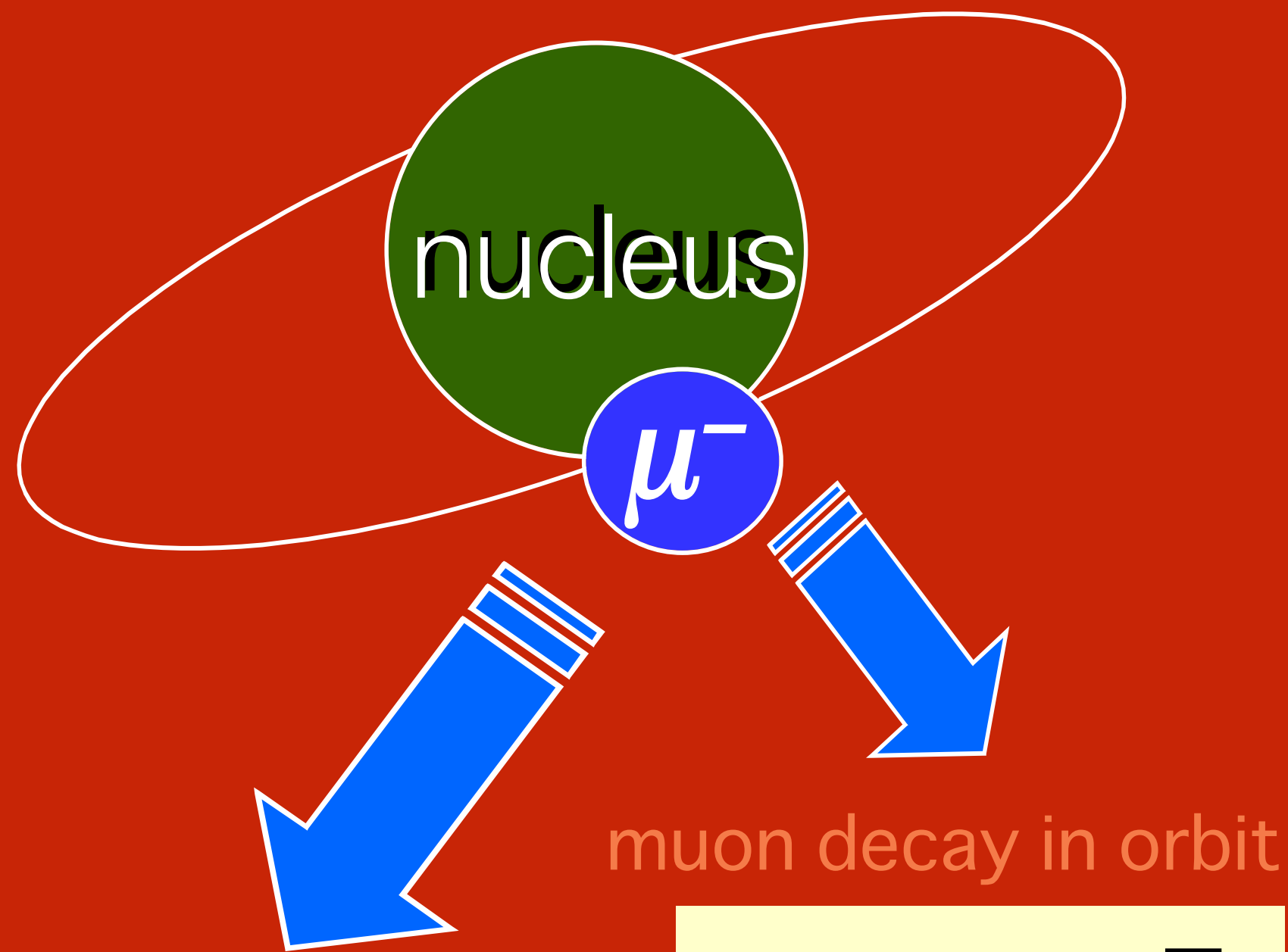
$$E_{\mu e} \approx m_\mu - E_{bound \mu} - E_{recoil} \approx 105 \text{ MeV}$$

Backgrounds:

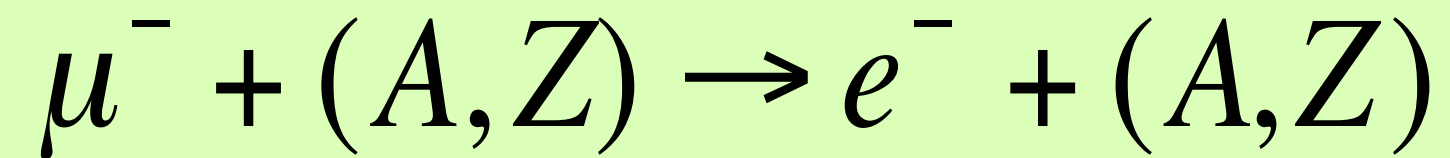
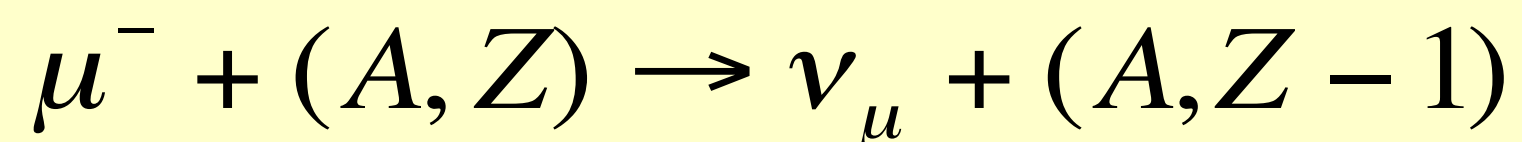
- (1) physics backgrounds
- (2) beam-related backgrounds
- (3) cosmic rays, false tracking

$\mu^- \rightarrow e^-$ Conversion

1s state in a muonic atom



nuclear muon capture



coherent process (for transition to ground state) $\propto Z^5$

Event Signature :

a mono-energetic electron

(one particle measurement allows higher muon rates.)

$$E_{\mu e} \approx m_\mu - E_{\text{bound } \mu} - E_{\text{recoil}} \approx 105 \text{ MeV}$$

Backgrounds:

- (1) physics backgrounds
- (2) beam-related backgrounds
- (3) cosmic rays, false tracking

Conversion rate:

$$\text{CR}(\mu^- \text{N} \rightarrow e^- \text{N}) \equiv \frac{\Gamma(\mu^- \text{N} \rightarrow e^- \text{N})}{\Gamma(\mu^- \text{N} \rightarrow \text{all})}$$

nucleus	Z	CR limit
sulfur	16	7×10^{-11}
titanium	22	4.3×10^{-12}
copper	39	1.6×10^{-8}
gold	79	7×10^{-13}
lead	82	4.6×10^{-11}

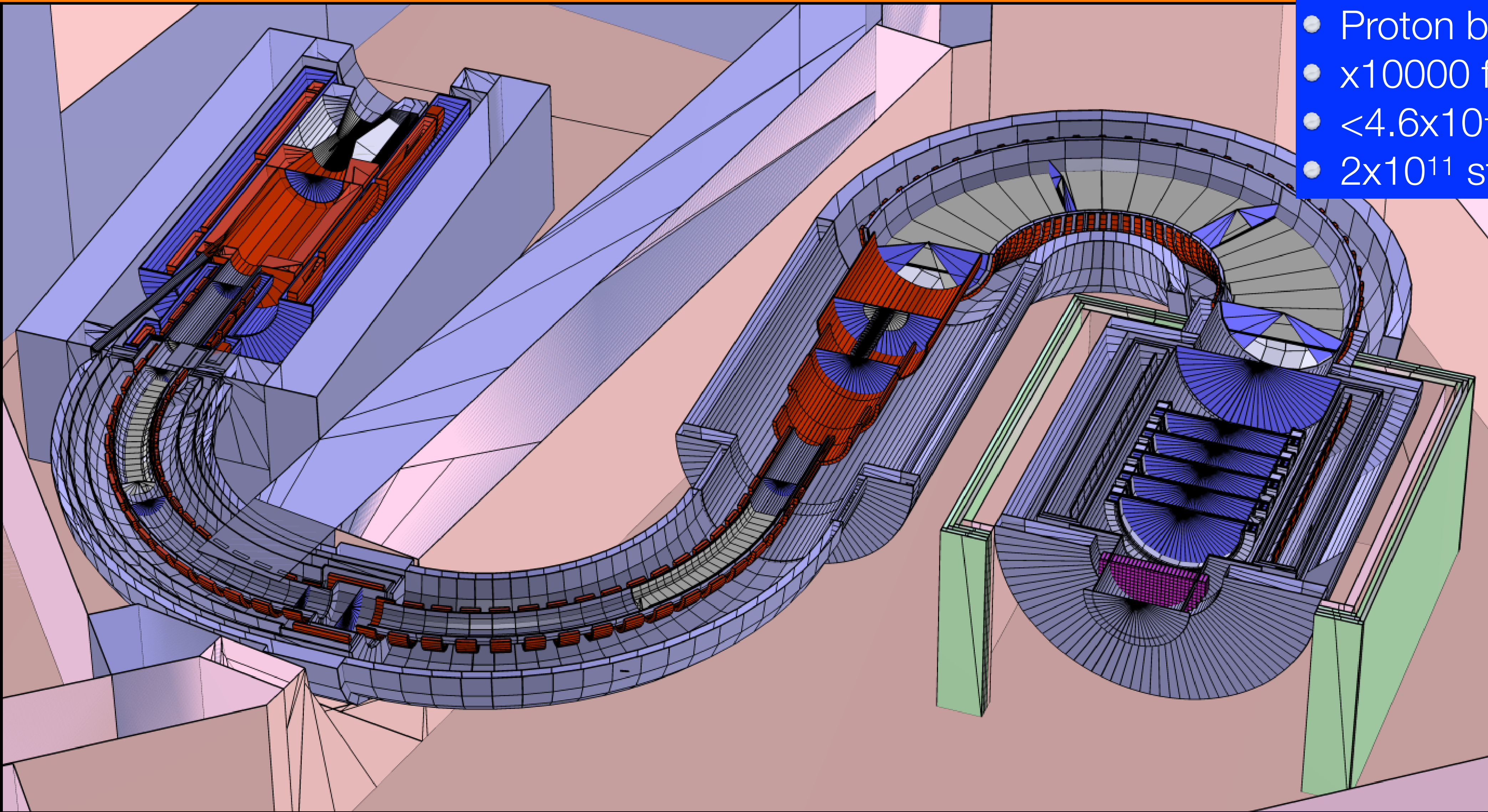
COMET at J-PARC



- COMET= COherent Muon to Electron Transition



COMET at J-PARC

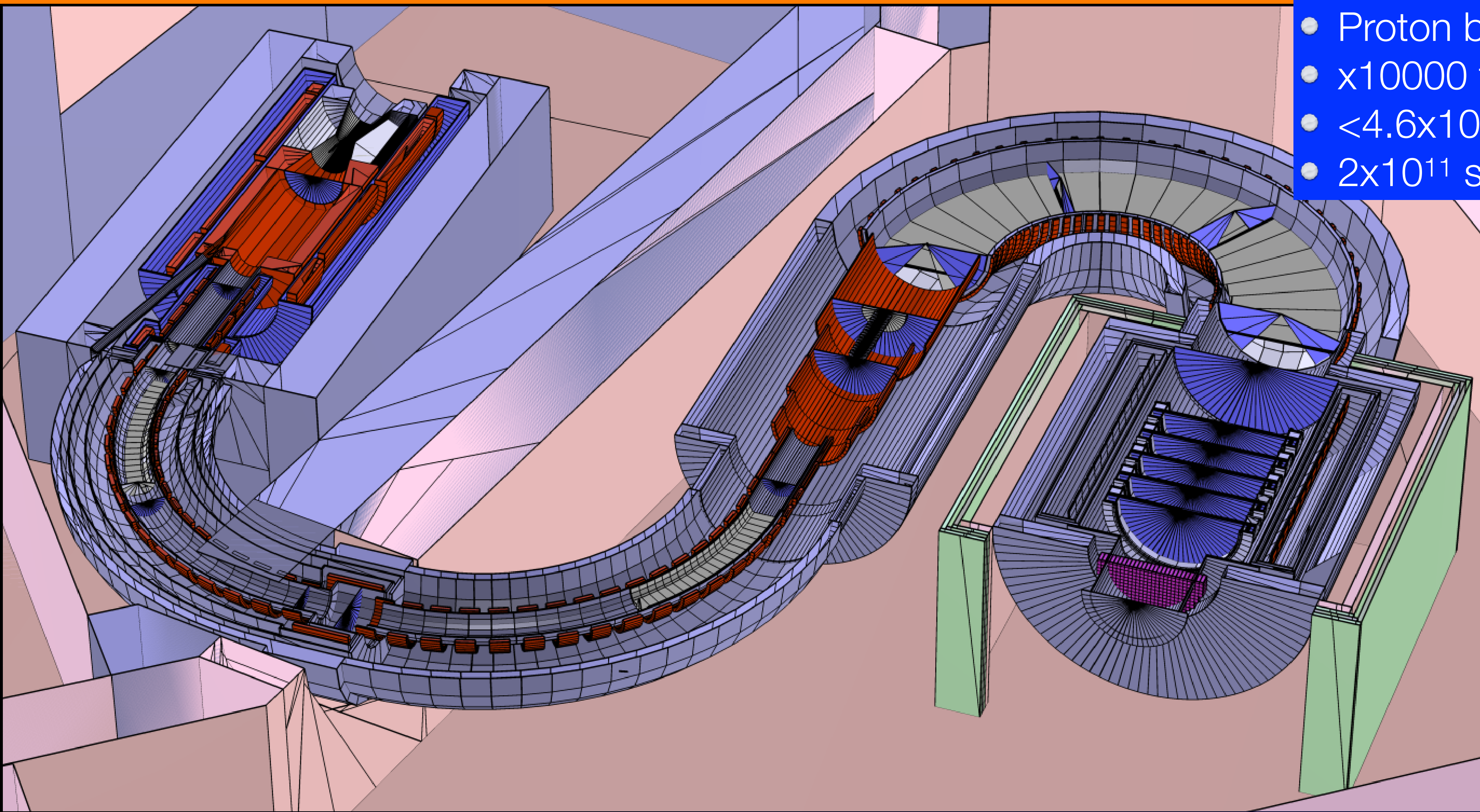


- Proton beam, 8 GeV, 56 kW
- $\times 10000$ from SINDRUM-II
- $< 4.6 \times 10^{-17}$ 90% C.L.
- 2×10^{11} stopped muons/s

- COMET= COherent Muon to Electron Transition



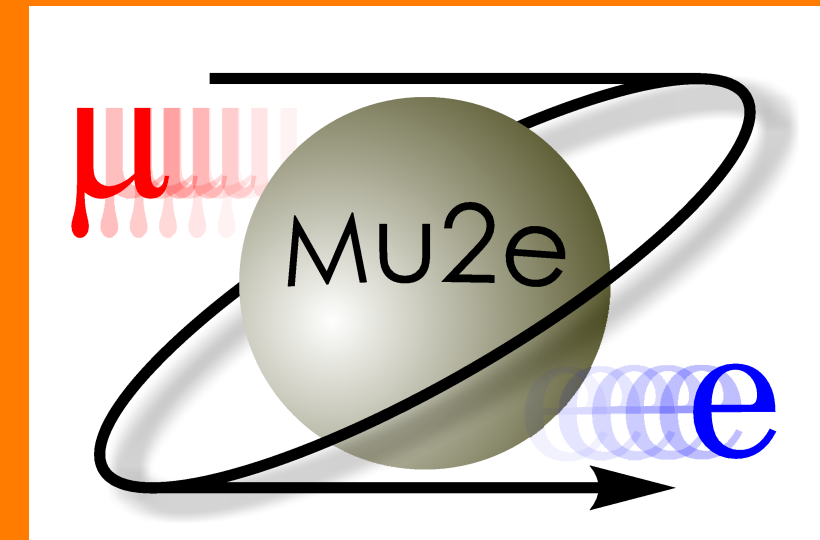
COMET at J-PARC



- Proton beam, 8 GeV, 56 kW
- $\times 10000$ from SINDRUM-II
- $< 4.6 \times 10^{-17}$ 90% C.L.
- 2×10^{11} stopped muons/s

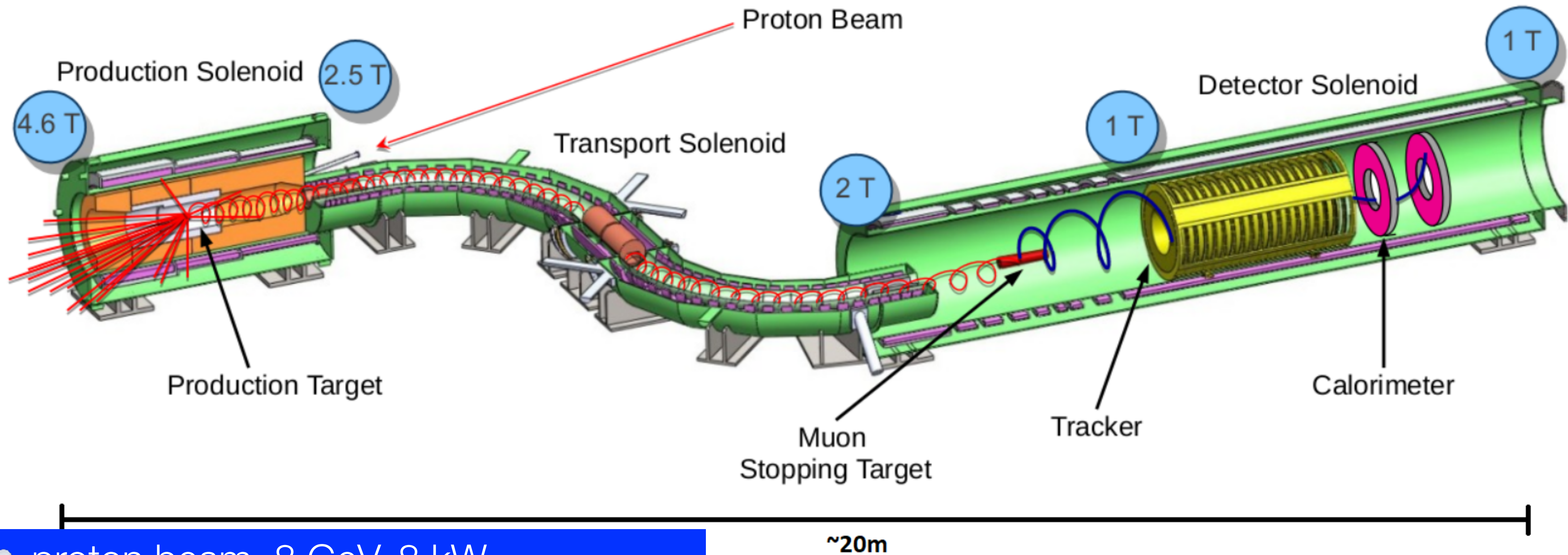
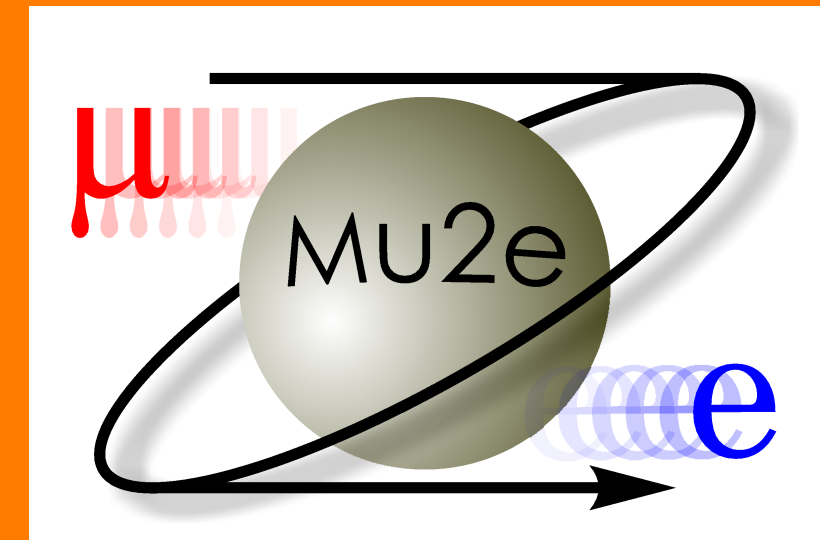
See the talk
by Chen Wu

Mu2e at Fermilab



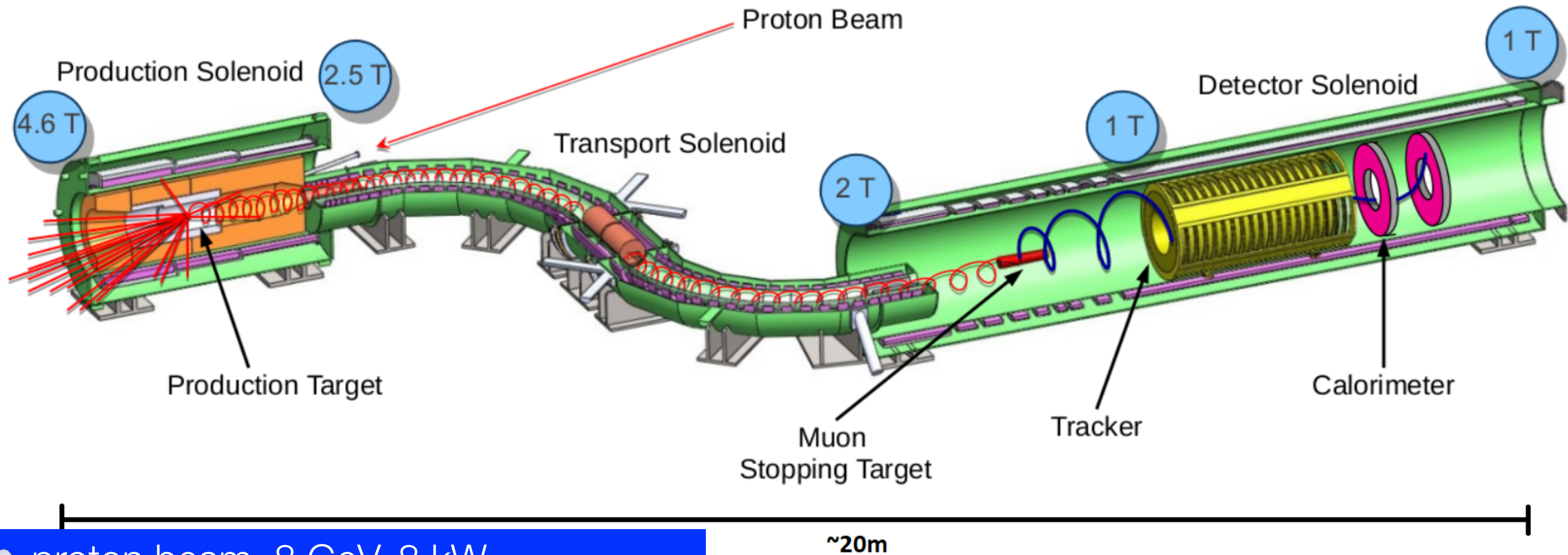
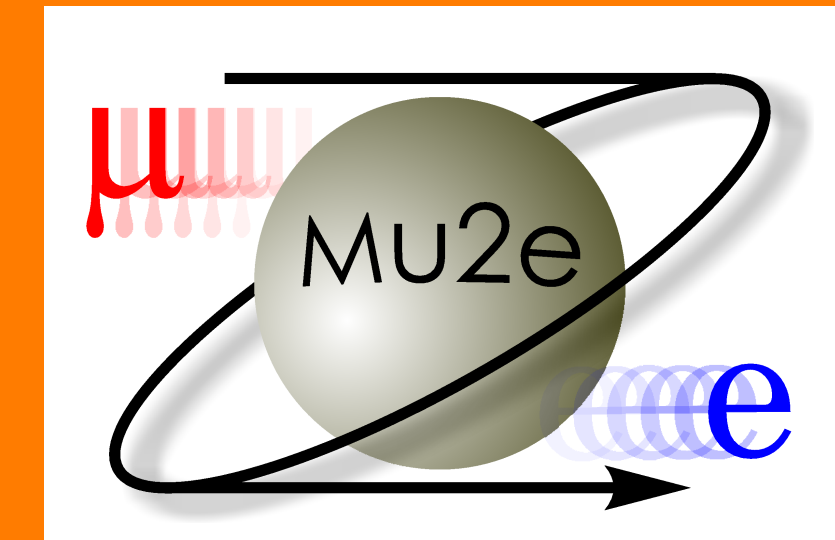
proton beam power = 8 kW

Mu2e at Fermilab



- proton beam, 8 GeV, 8 kW
- x10000 from SINDRUM-II
- $<8 \times 10^{-17}$ 90% C.L. or 5σ
discovery= 2×10^{-16}
- 0.11 ± 0.03 background events (Run 1)

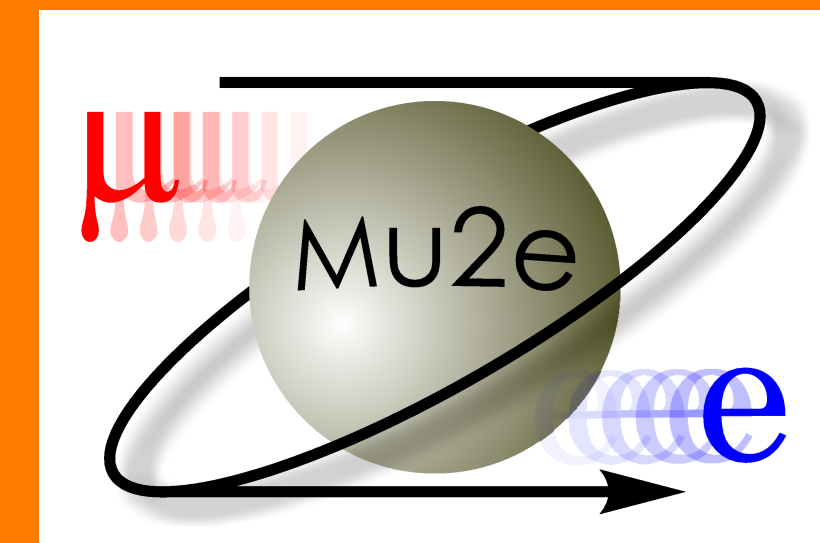
Mu2e at Fermilab



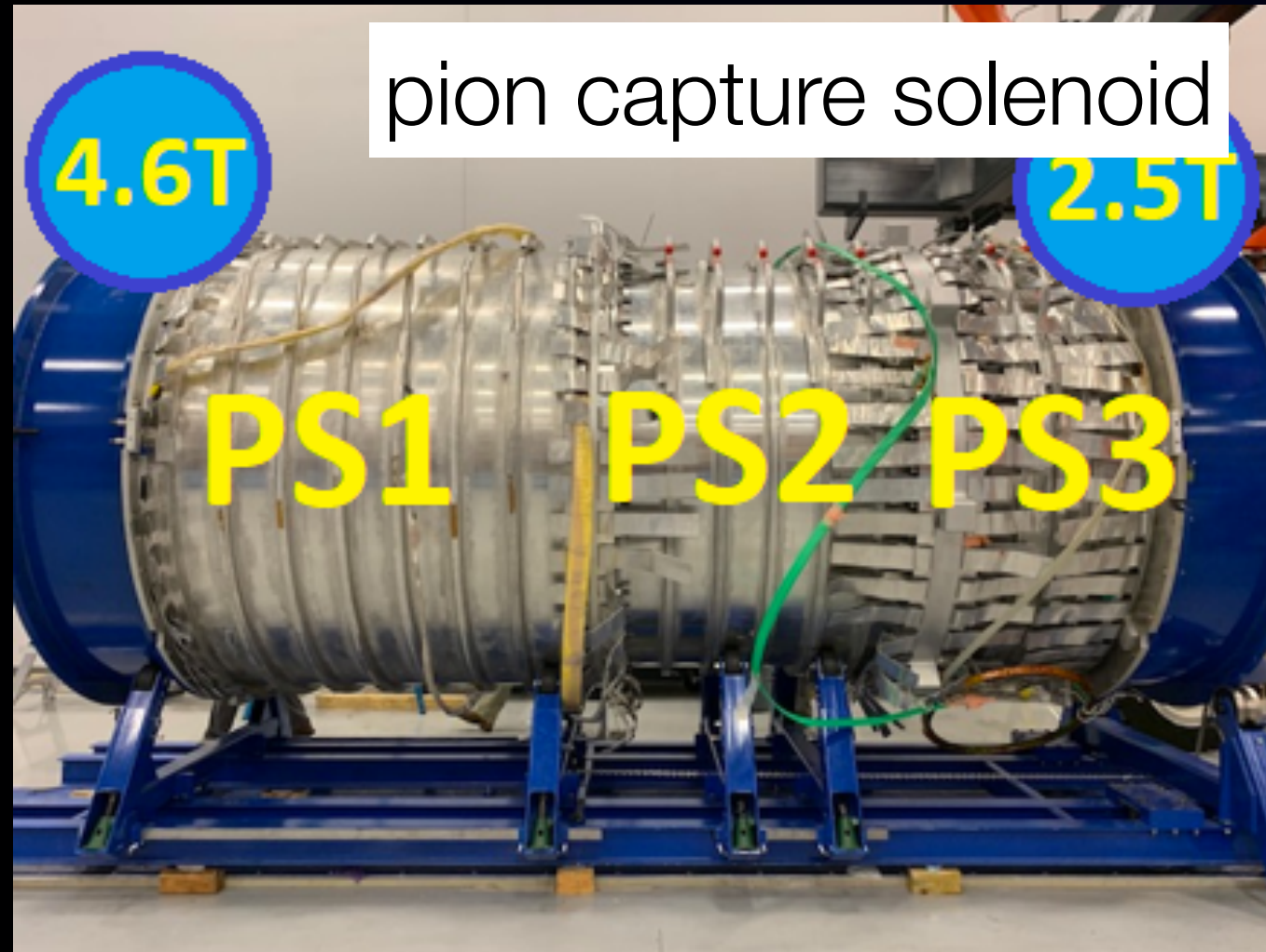
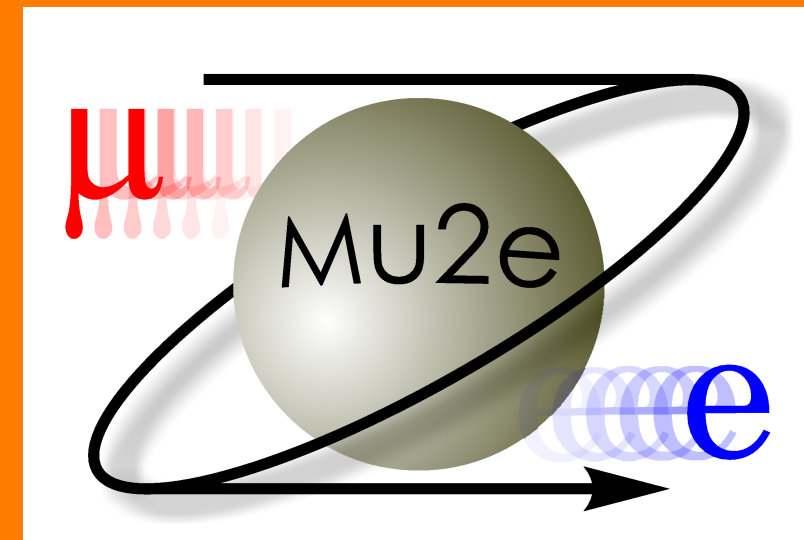
- proton beam, 8 GeV, 8 kW
- x10000 from SINDRUM-II
- $<8 \times 10^{-17}$ 90% C.L. or 5σ discovery= 2×10^{-16}
- 0.11 ± 0.03 background events (Run 1)

	muon transport	electron transport
Mu2e	s-shape curve	straight
COMET	c-shape curve	c-shape curve

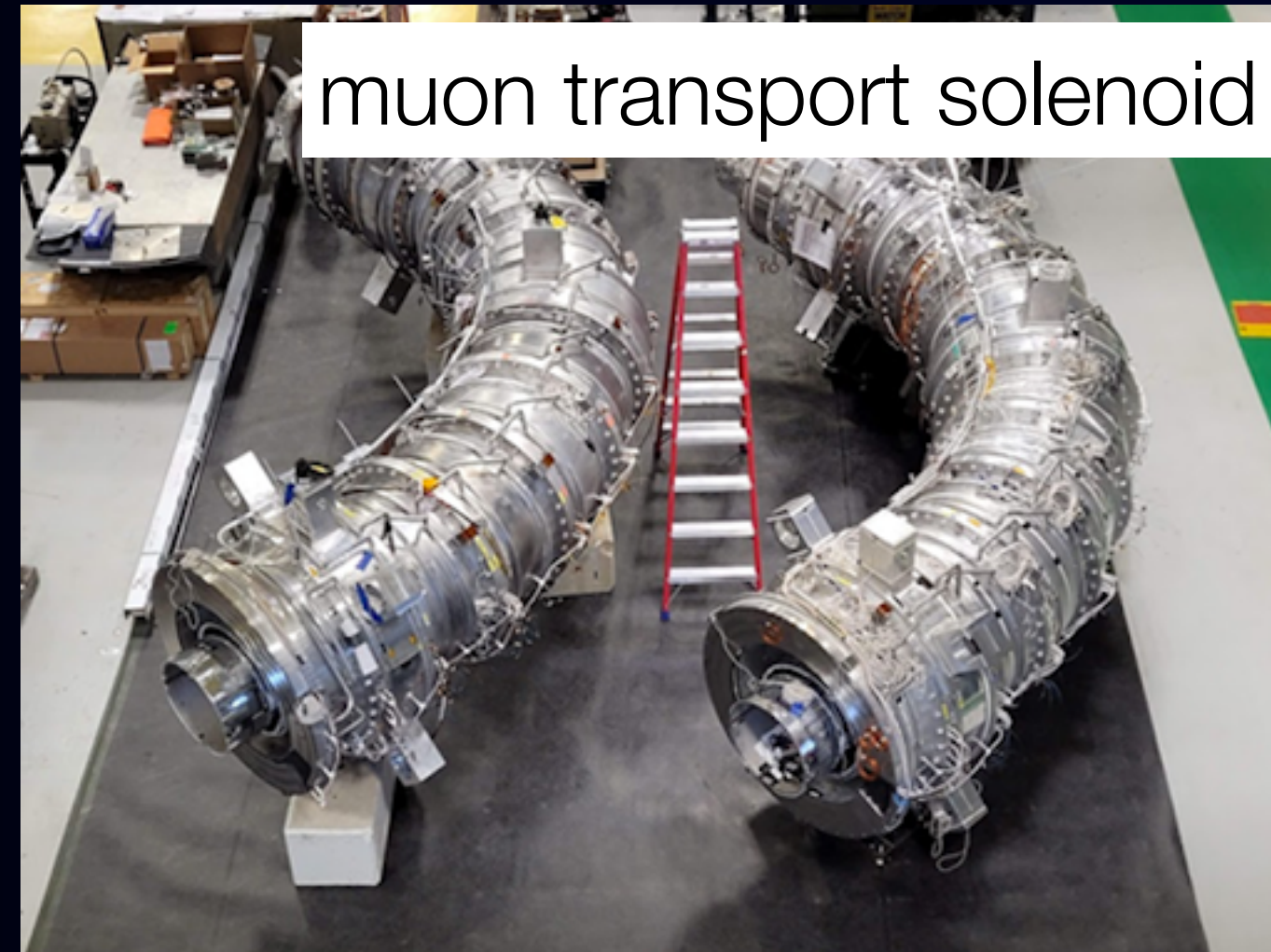
Mu2e at Fermilab



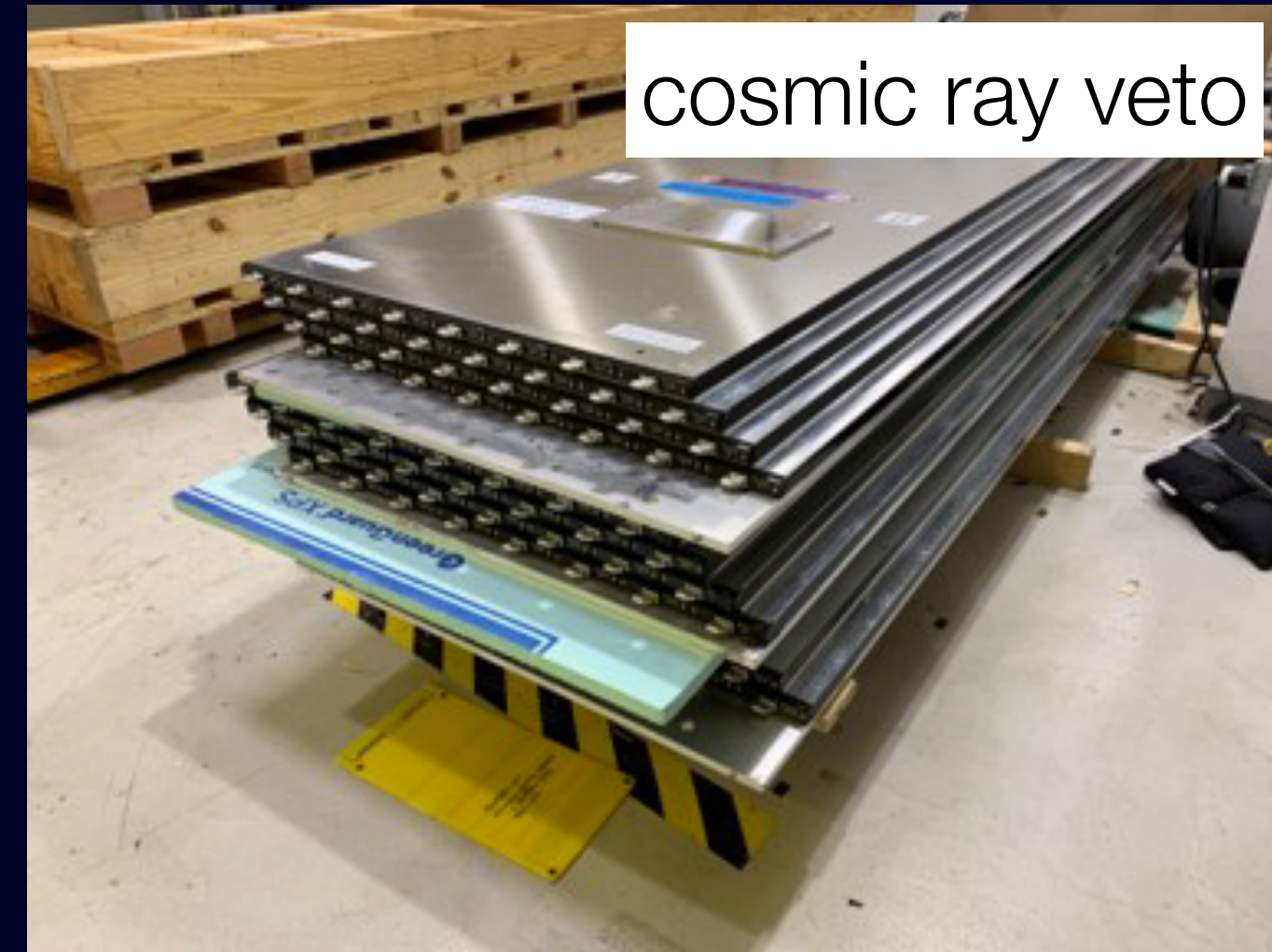
Mu2e at Fermilab



pion capture solenoid



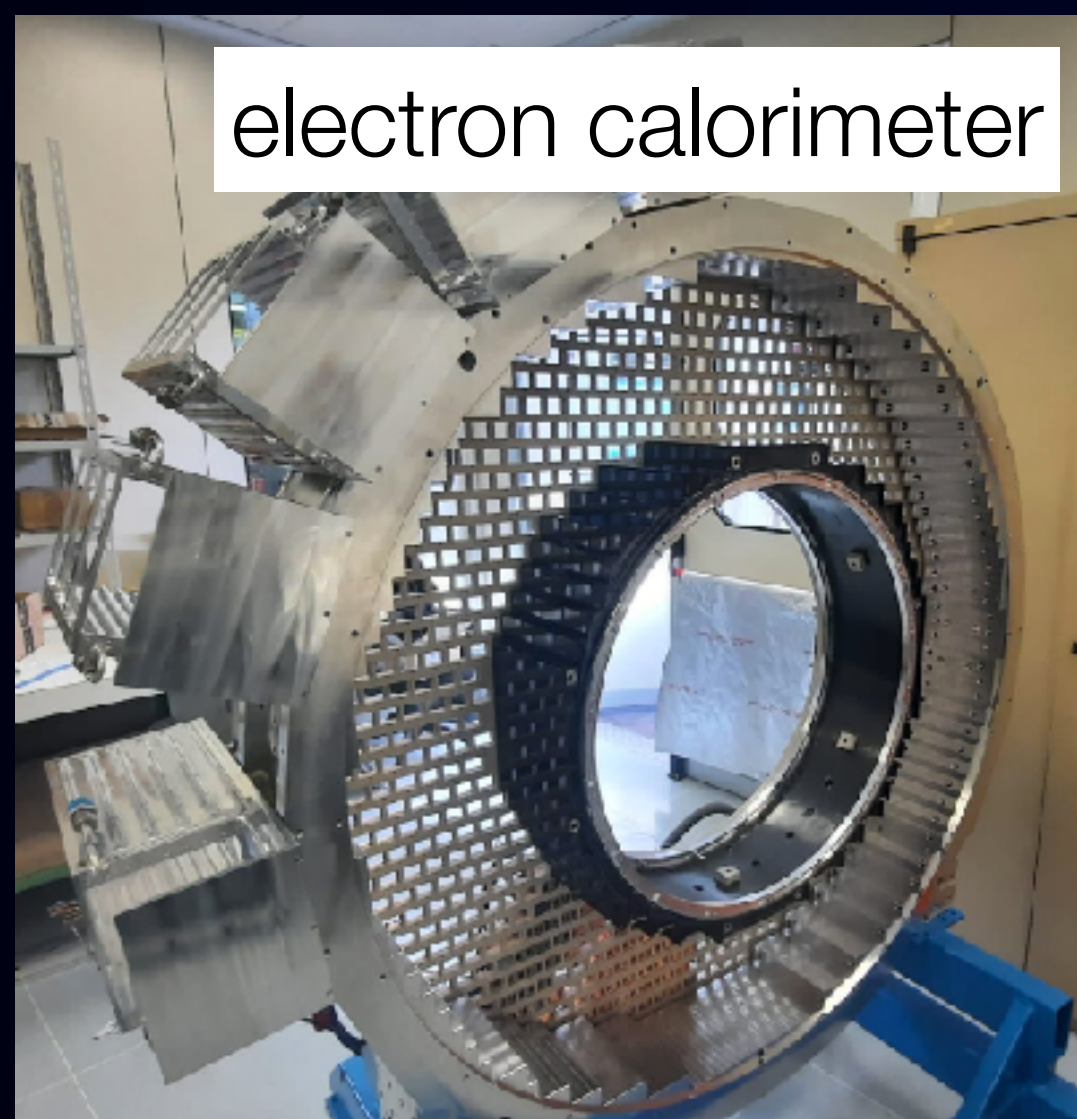
muon transport solenoid



cosmic ray veto

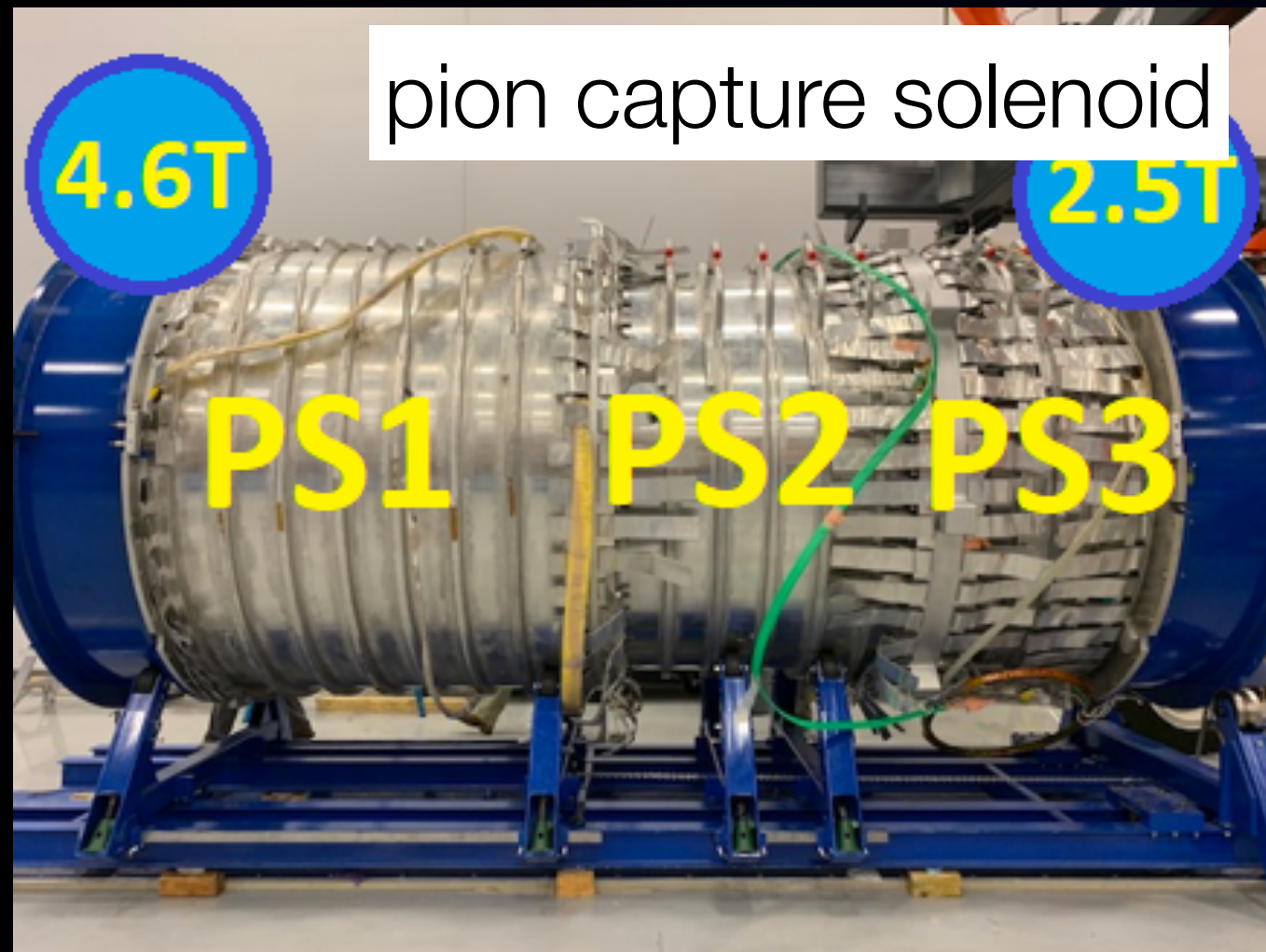
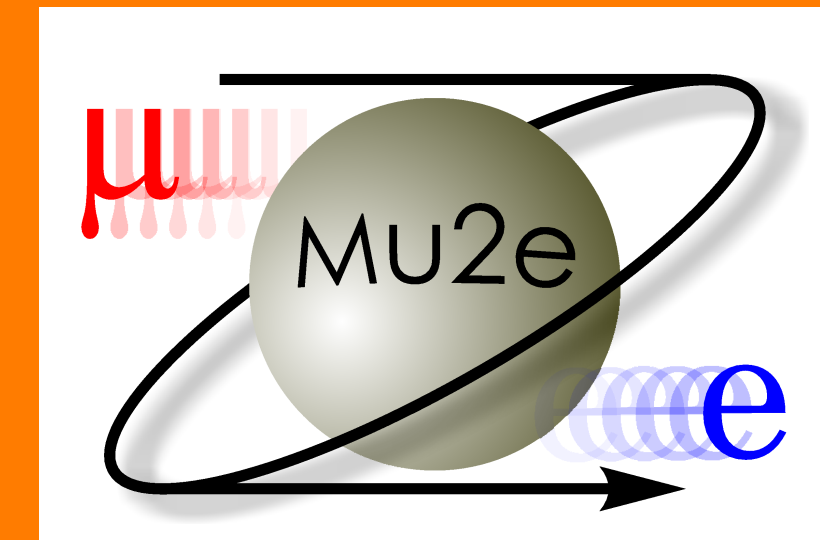


straw tracker



electron calorimeter

Mu2e at Fermilab

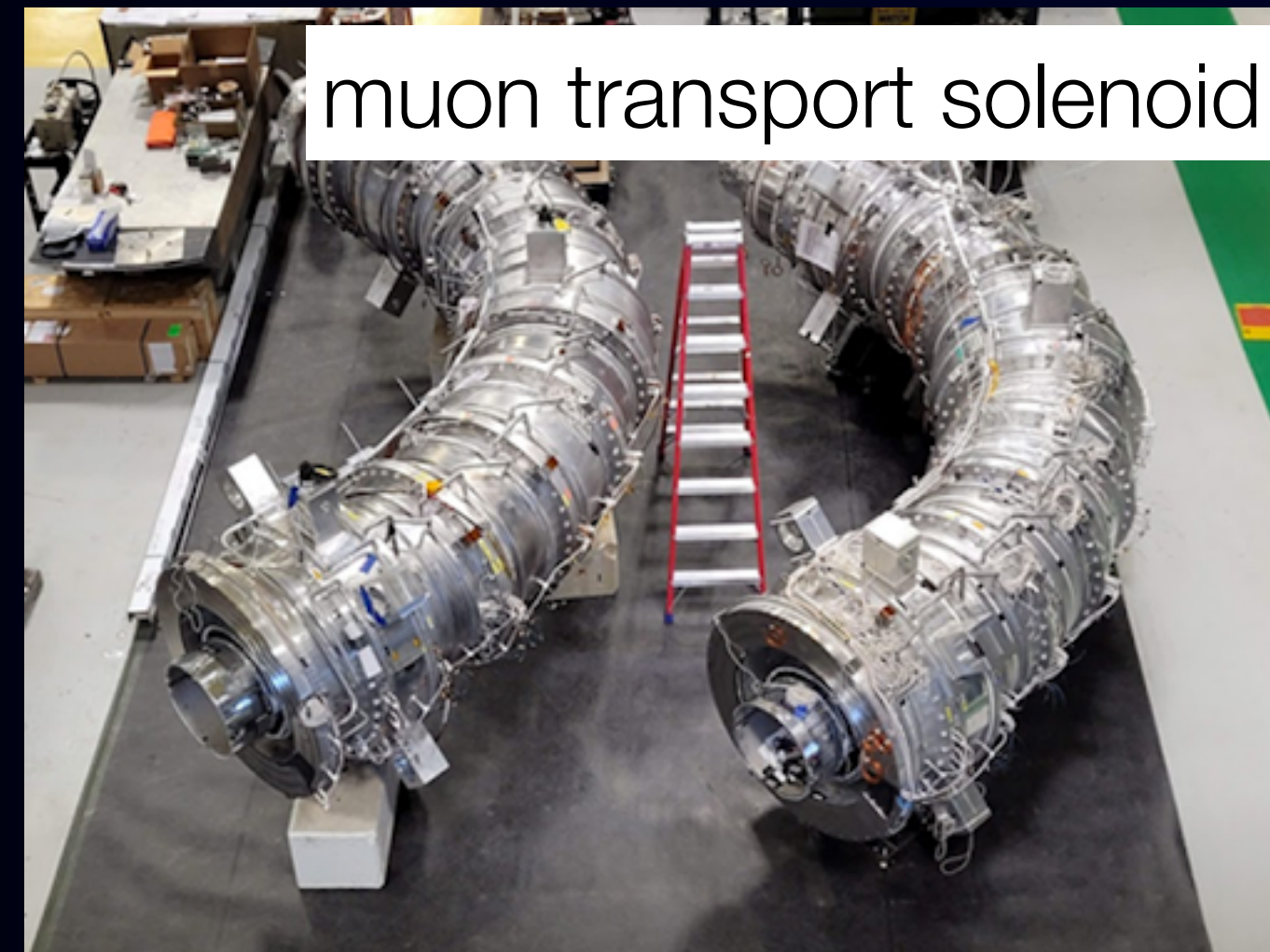


pion capture solenoid

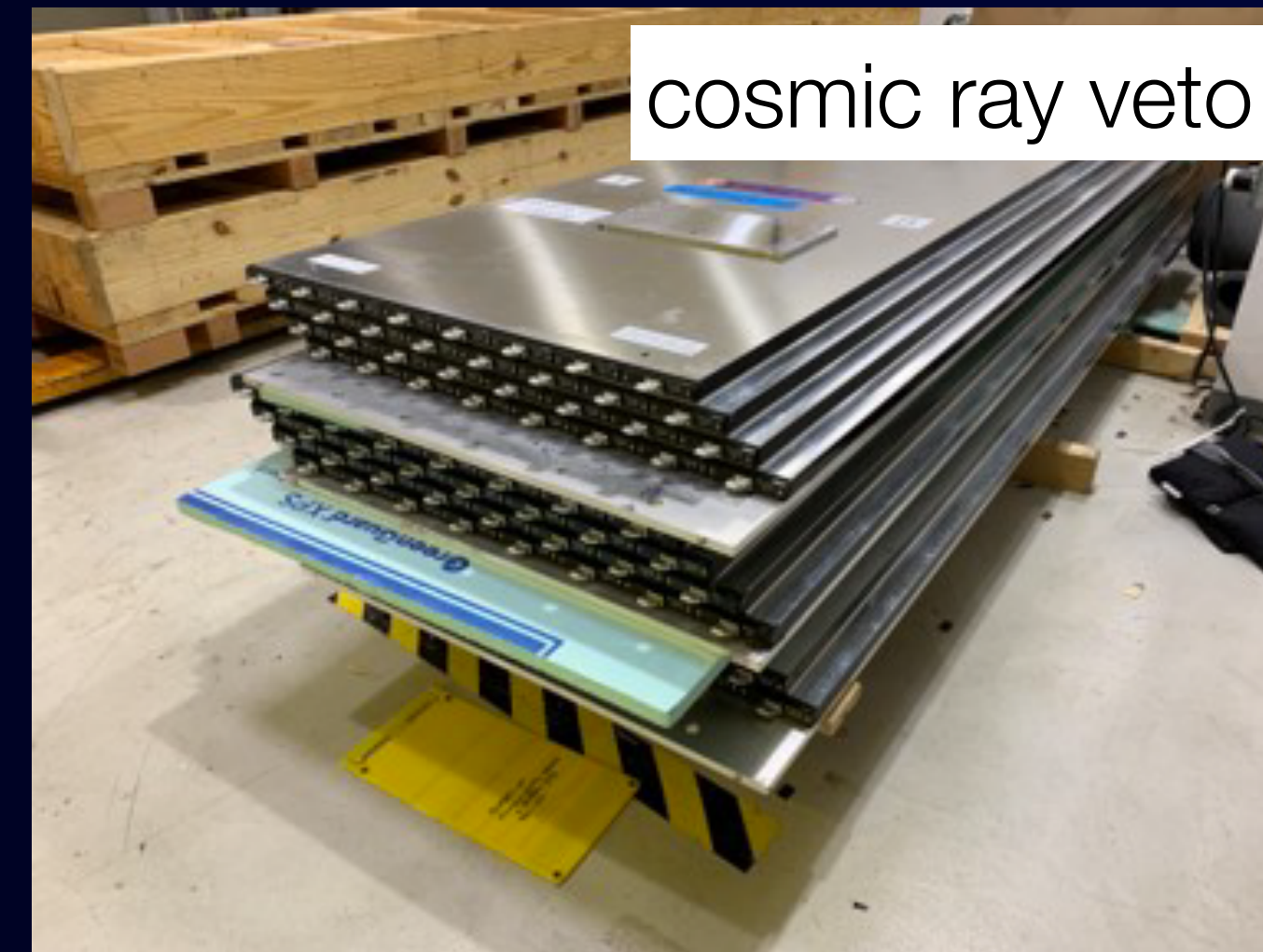
4.6T

2.5T

PS1 PS2 PS3



muon transport solenoid



cosmic ray veto



straw tracker



electron calorimeter

Timeline

- Detector commissioning from 2024
- Run 1 data taking in 2026 until LBNF/PIP-II shutdown (x1000 improvements)
- Run 2 data taking from 2029 after the shutdown

Mu2e-II and AMF at Fermilab

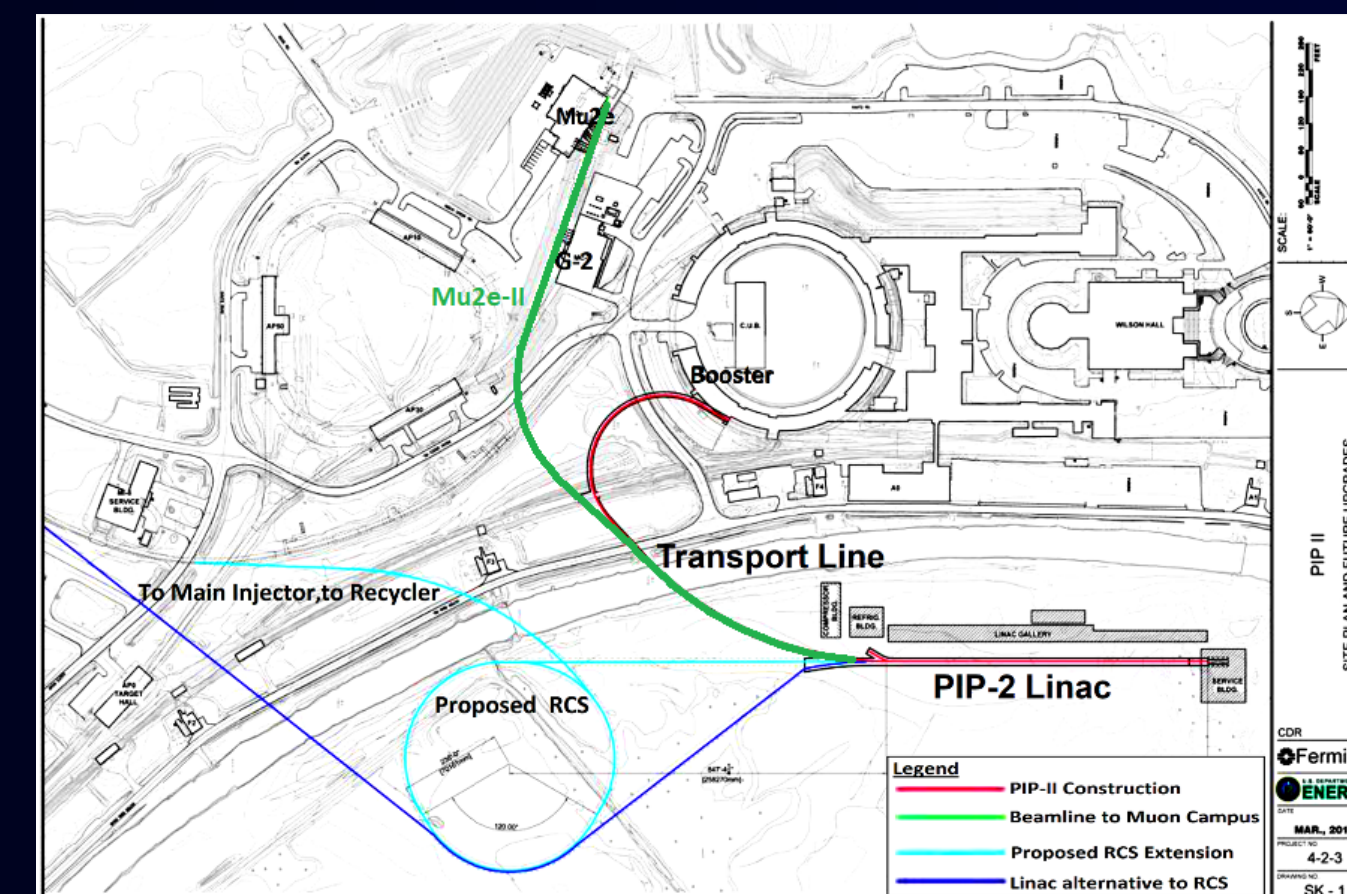


Mu2e-II and AMF at Fermilab



Mu2e-II

- 800 MeV, 100 kW from PIP-II Linac
- 5 years running (5.5×10^{19} stopped μ s)
- Goal: $< 6.4 \times 10^{-18}$ at 90% CL limit
 - x10 from Mu2e
 - 0.47 background events
- Detector refinements for Mu2e-II
 - proton target, tracker, calorimeter (BaF_2), cosmic ray veto
- Timeline ~ 2030s decade



PIP-II Project

Mu2e-II and AMF at Fermilab

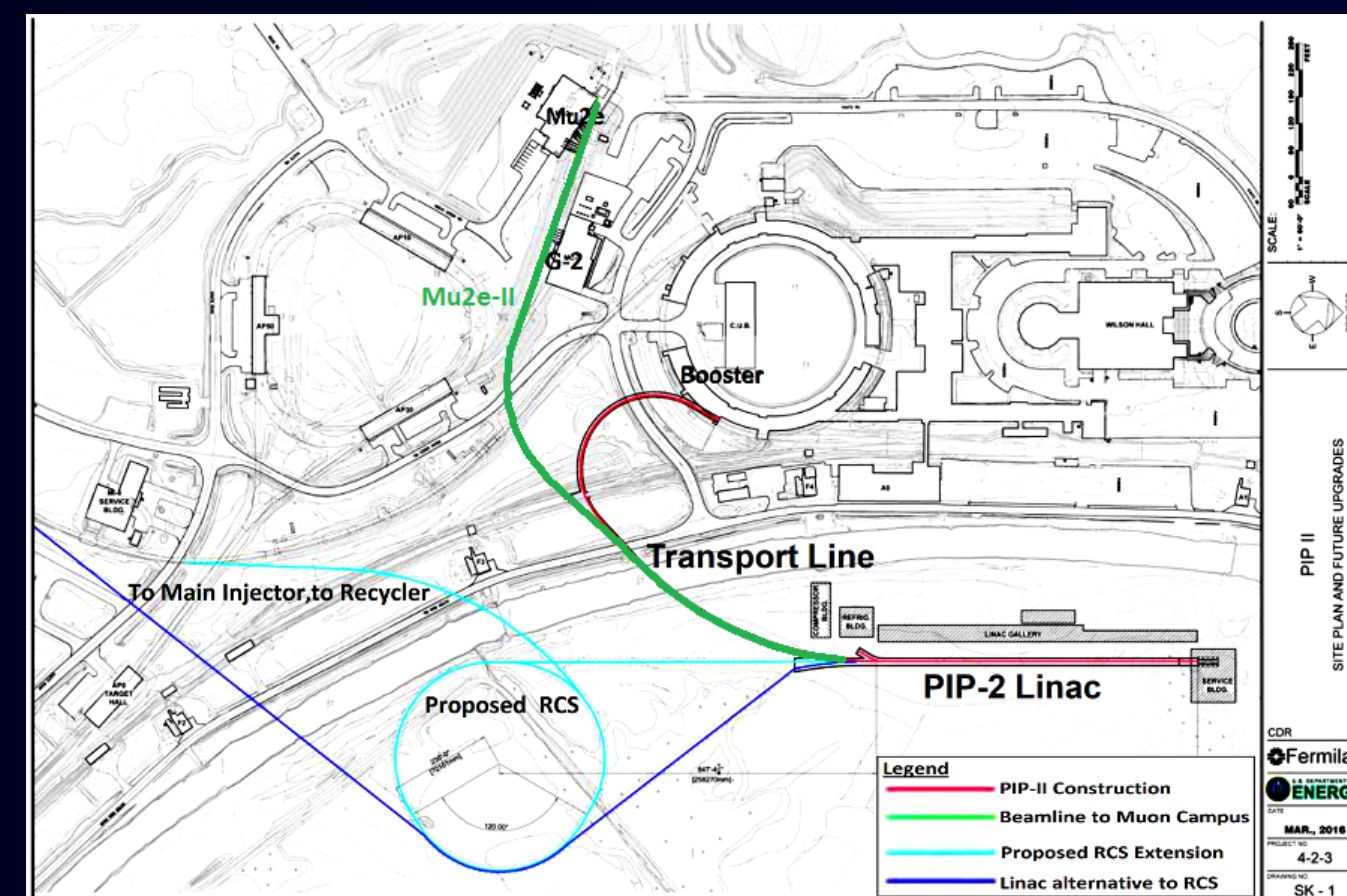


Mu2e-II

- 800 MeV, 100 kW from PIP-II Linac
- 5 years running (5.5×10^{19} stopped μ s)
- **Goal: $< 6.4 \times 10^{-18}$ at 90% CL limit**
 - x10 from Mu2e
 - 0.47 background events
- Detector refinements for Mu2e-II
 - proton target, tracker, calorimeter (BaF_2), cosmic ray veto
- Timeline ~ 2030s decade

AMF (Advanced Muon Facility)

- > 1 MW beam from PIP-II Linac
- aiming at searches for three muon CFLV processes
- **muon storage ring (FFA) for $\mu \rightarrow e$ conversion (Enigma@AMF)**
- ~2040

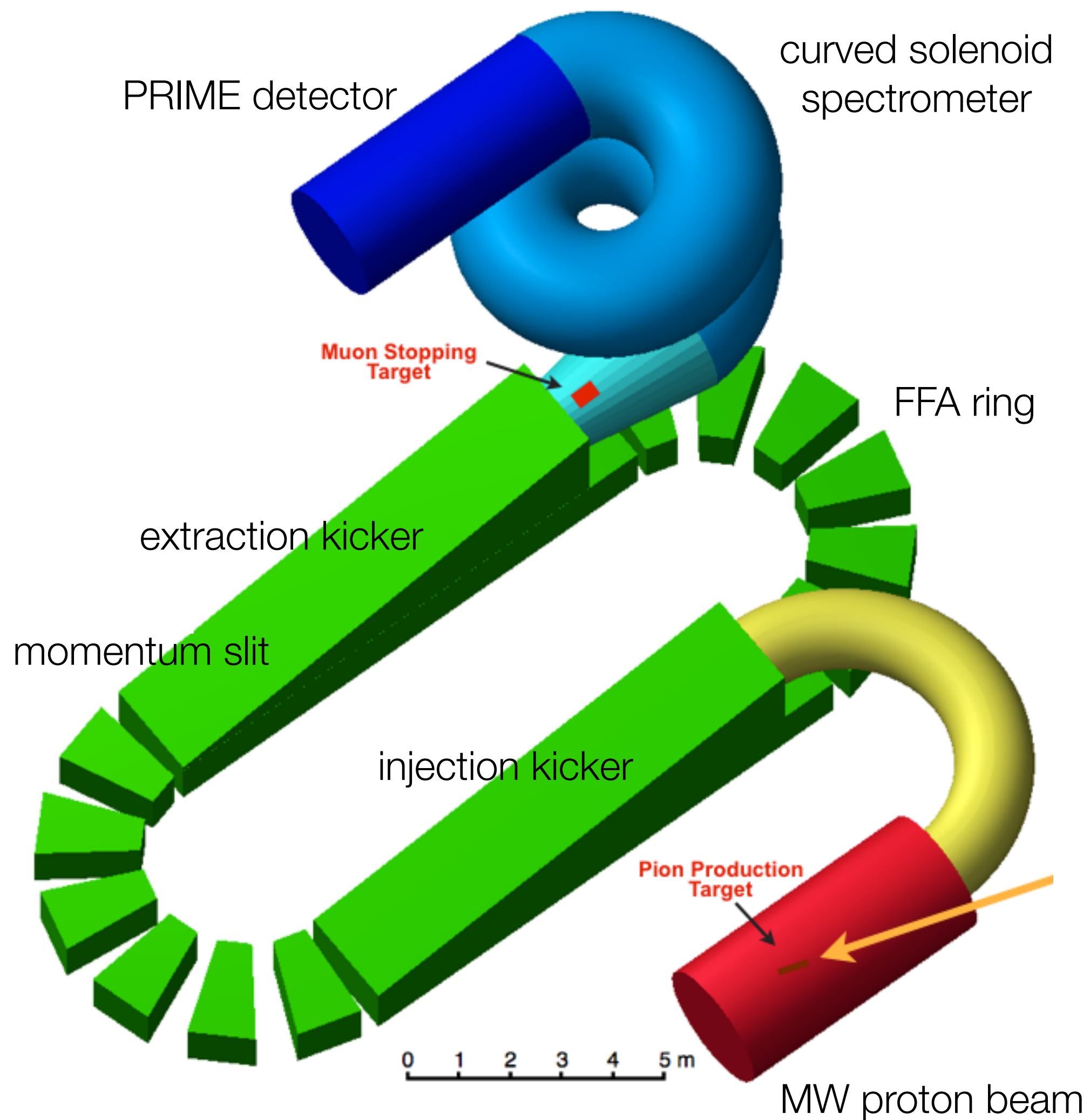


PIP-II Project

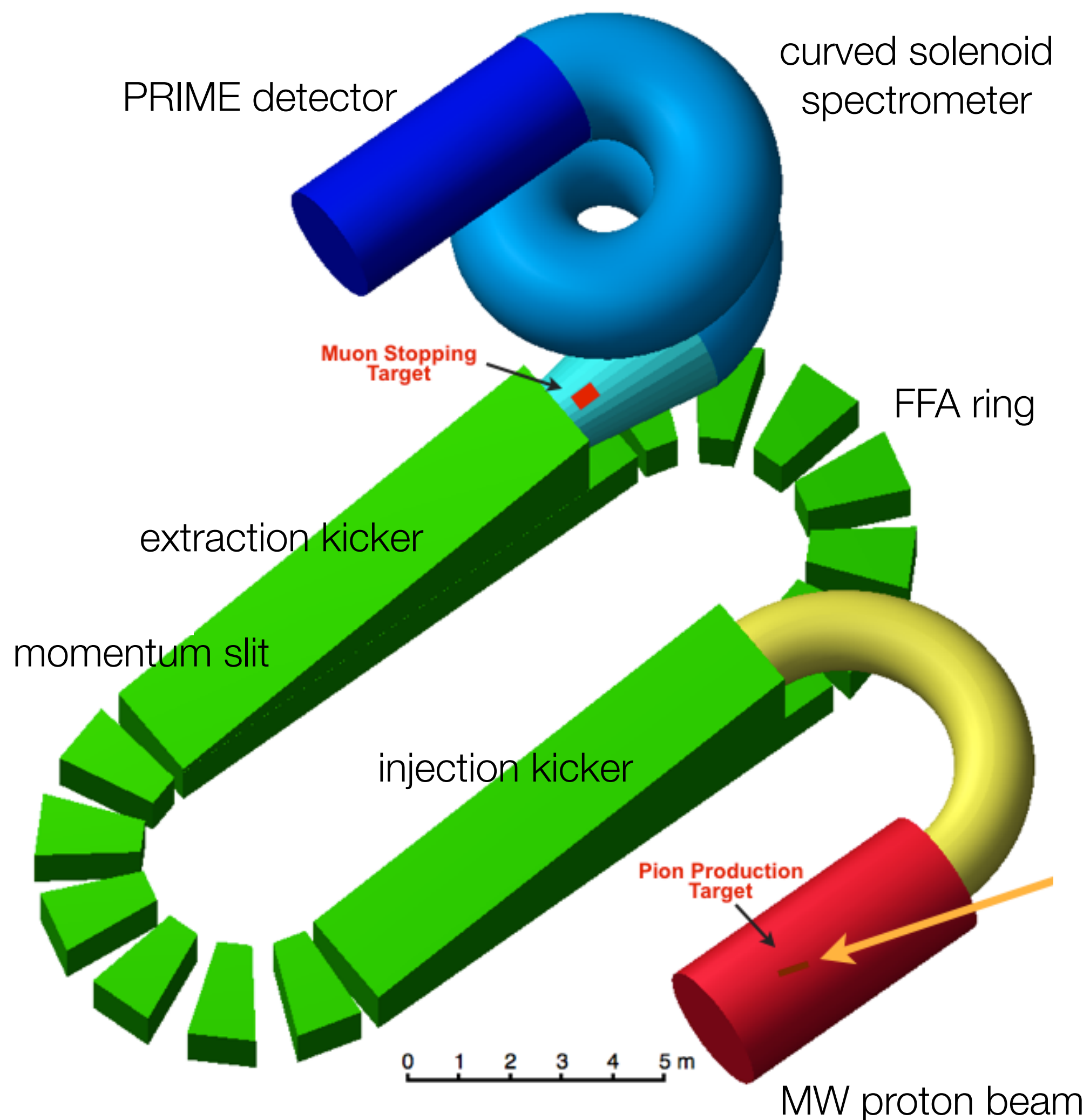
$\mu^- \rightarrow e^-$ Conversion at $<O(10^{-19})$
with Muon Storage Ring



$\mu^- \rightarrow e^-$ Conversion at $<O(10^{-19})$ with Muon Storage Ring

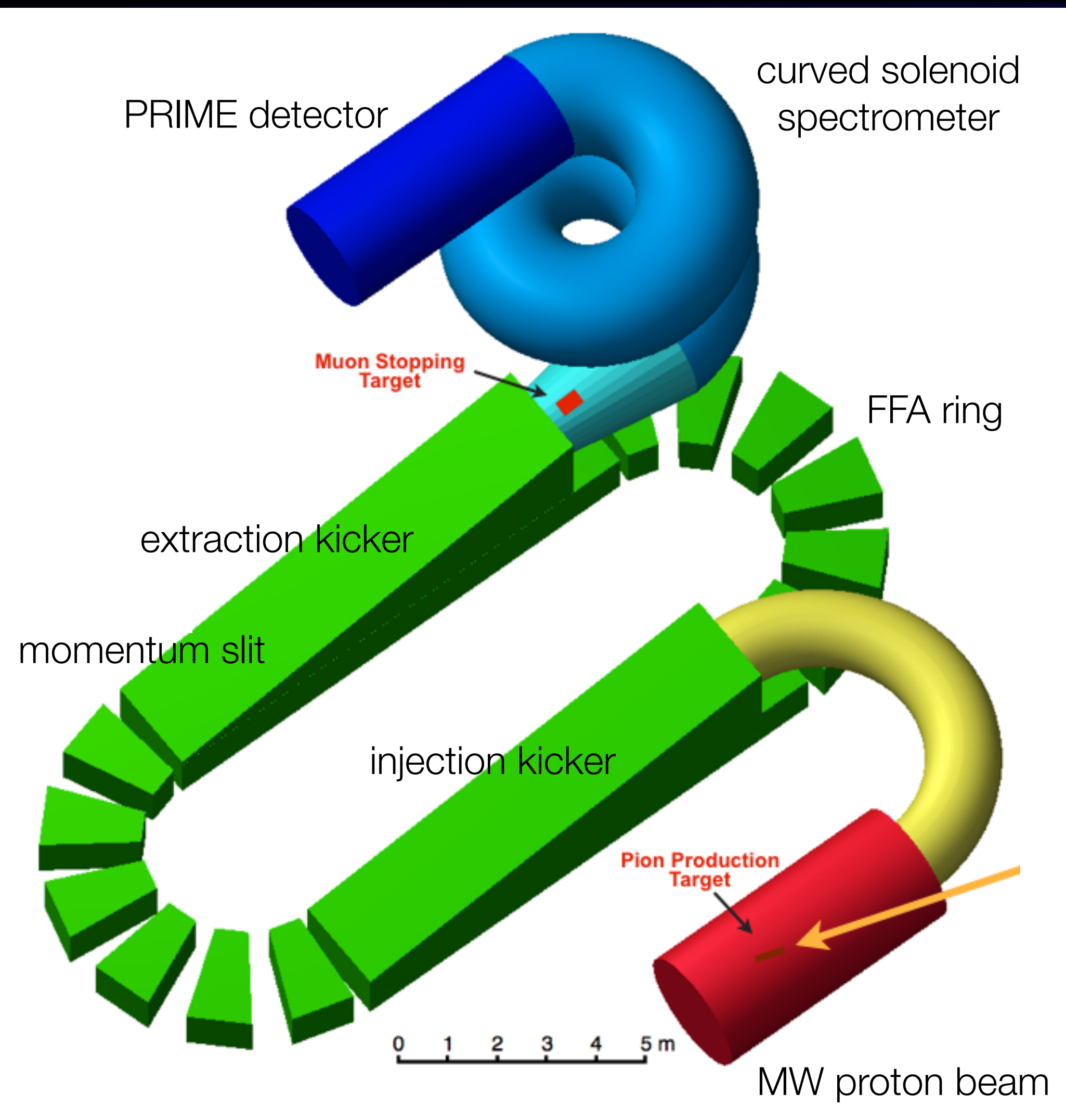


$\mu^- \rightarrow e^-$ Conversion at $<O(10^{-19})$ with Muon Storage Ring

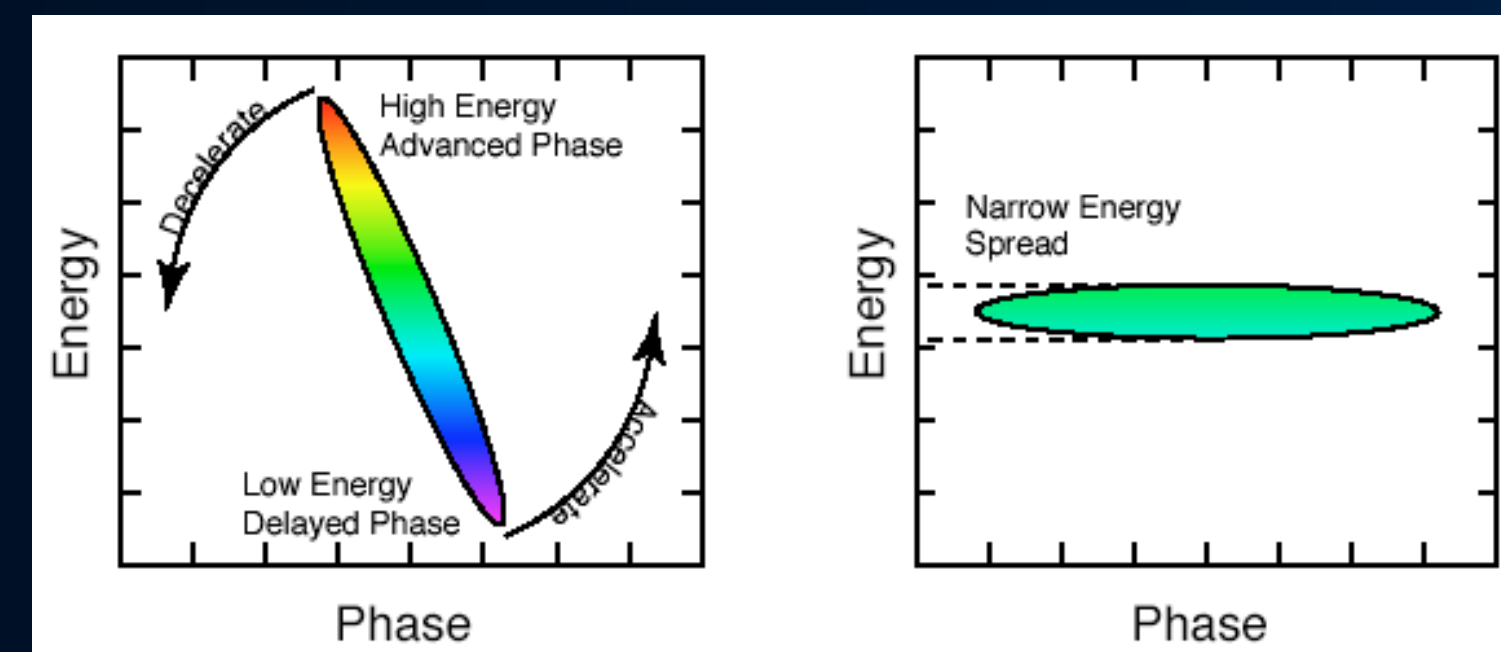


- Advantage of muon storage ring for $\mu \rightarrow e$ conversion
 - Removal of pions in a muon beam
 - long flight length in a ring to eliminate pions (pion reduction of 10^{-20} for 5 turns, 30 m circumference)
 - High-Z muon targets can be used
 - Phase rotation to narrow beam energy spread
 - accelerate slow muons and decelerate fast muons
 - thinner muon target for better signal acceptance

$\mu^- \rightarrow e^-$ Conversion at $<O(10^{-19})$ with Muon Storage Ring

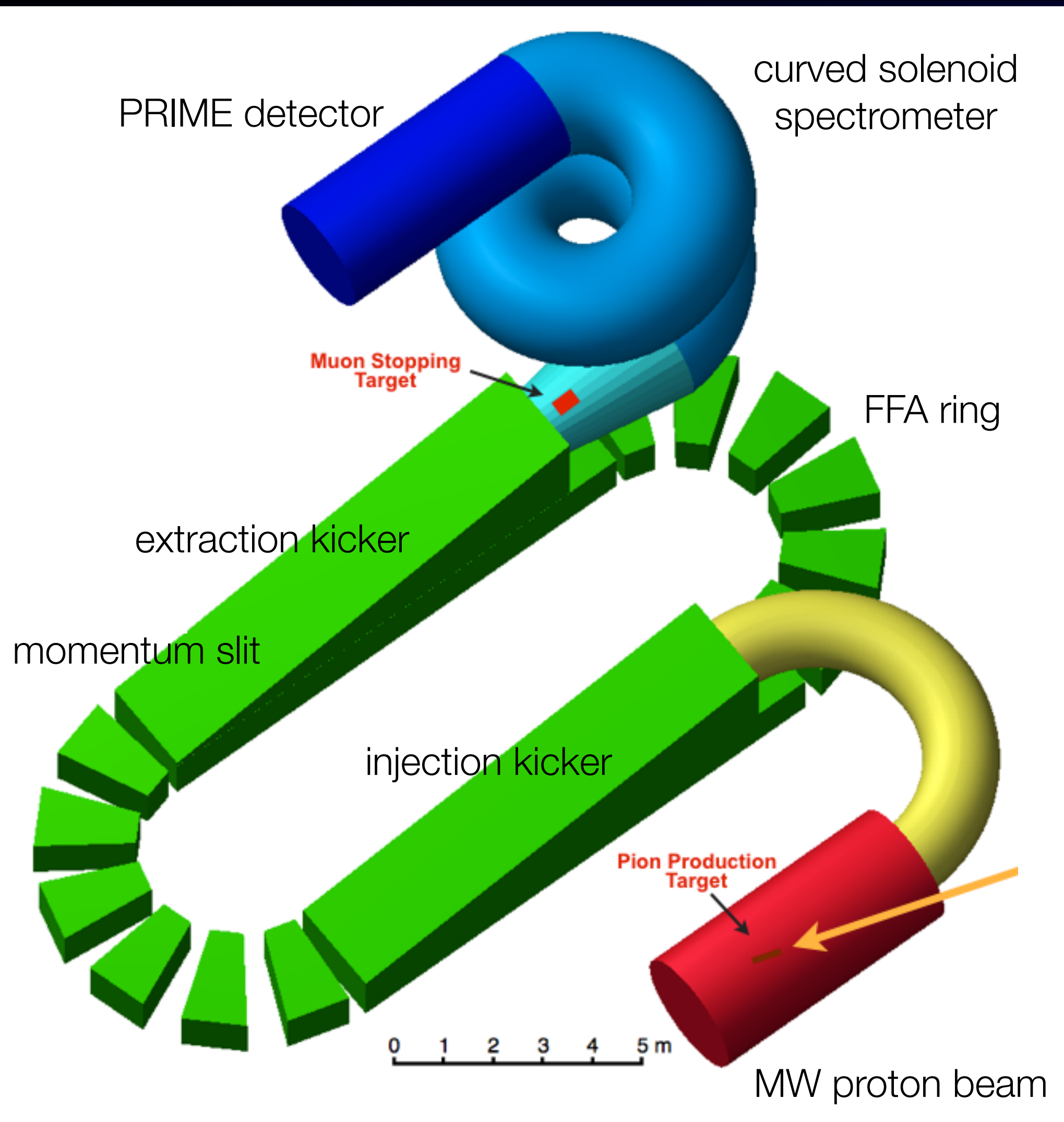


- Advantage of muon storage ring for $\mu \rightarrow e$ conversion
 - Removal of pions in a muon beam
 - long flight length in a ring to eliminate pions (pion reduction of 10^{-20} for 5 turns, 30 m circumference)
 - High-Z muon targets can be used
 - Phase rotation to narrow beam energy spread
 - accelerate slow muons and decelerate fast muons
 - thinner muon target for better signal acceptance

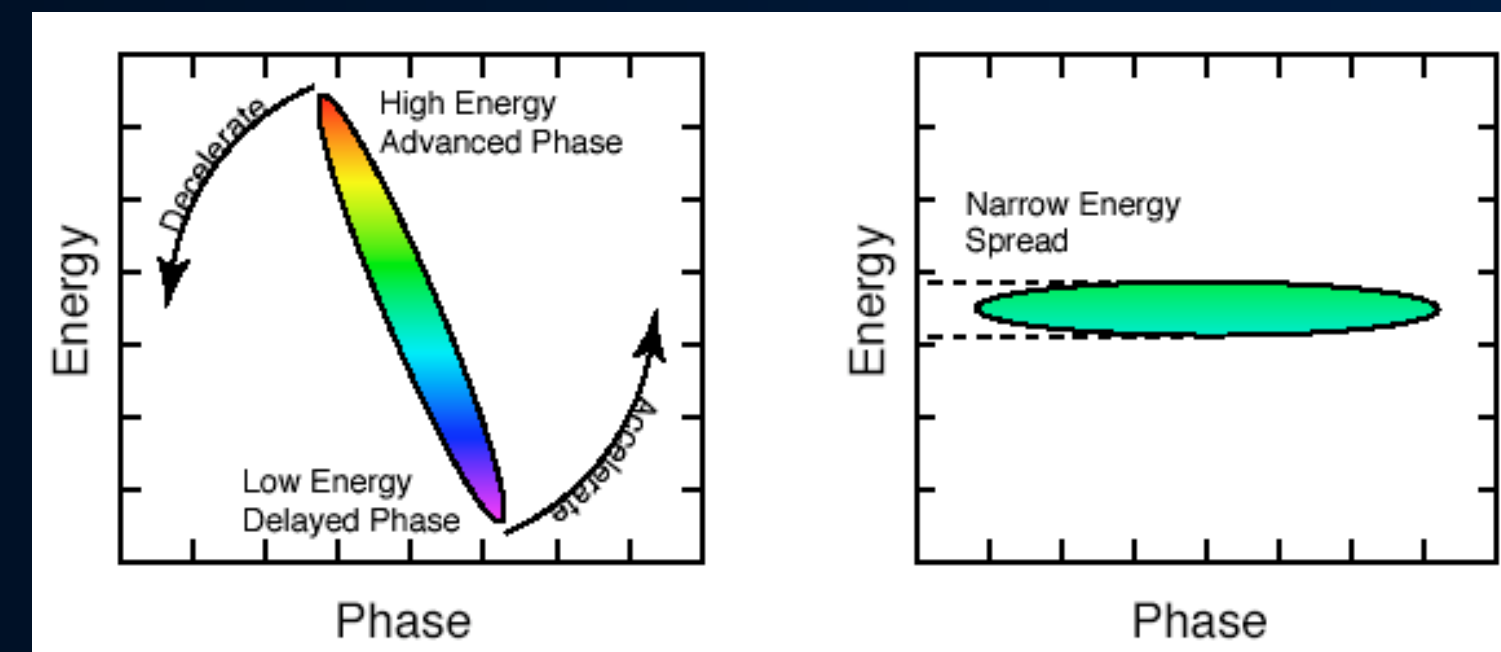


principle of phase rotation

$\mu^- \rightarrow e^-$ Conversion at $<O(10^{-19})$ with Muon Storage Ring



- Advantage of muon storage ring for $\mu \rightarrow e$ conversion
 - Removal of pions in a muon beam
 - long flight length in a ring to eliminate pions (pion reduction of 10^{-20} for 5 turns, 30 m circumference)
 - High-Z muon targets can be used
 - Phase rotation to narrow beam energy spread
 - accelerate slow muons and decelerate fast muons
 - thinner muon target for better signal acceptance



principle of phase rotation

- Future Prospects (2040~)
 - Enigma@AMF or PRISM@J-PARC

PRISM R&D at Osaka (2003-2007)



PRISM R&D at Osaka (2003-2007)



Fixed Field Alternating Gradient Synchrotron (FFA)

PRISM R&D at Osaka (2003-2007)



Fixed Field Alternating Gradient Synchrotron (FFA)

- Best accelerator ring for low-energy muons
 - Large beam acceptance
 - Fast beam acceleration
 - Synchrotron oscillation for phase rotation
- High field RF development
 - Finemet RF core, 150 kV/m

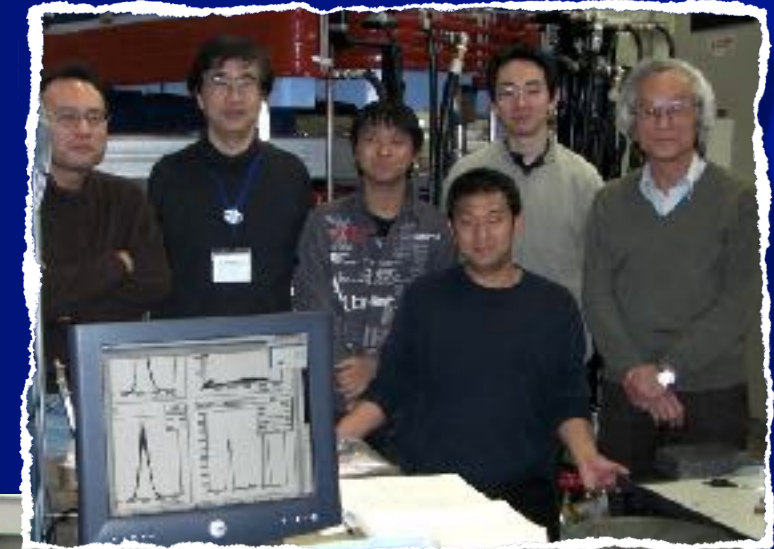
PRISM R&D at Osaka (2003-2007)



Fixed Field Alternating Gradient Synchrotron (FFA)

- Best accelerator ring for low-energy muons
 - Large beam acceptance
 - Fast beam acceleration
 - Synchrotron oscillation for phase rotation
- High field RF development
 - Finemet RF core, 150 kV/m

2003 - 2007 at Osaka University

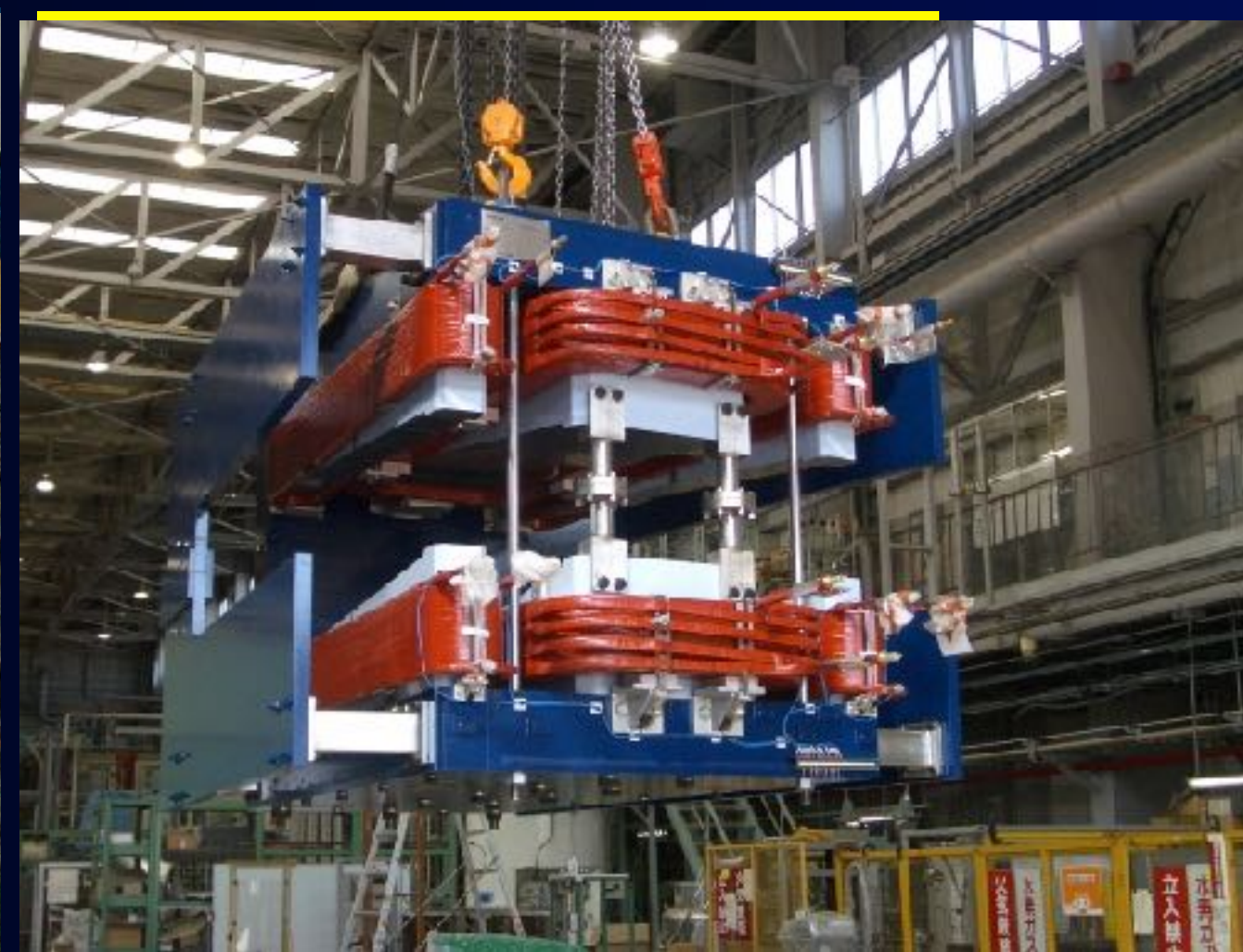
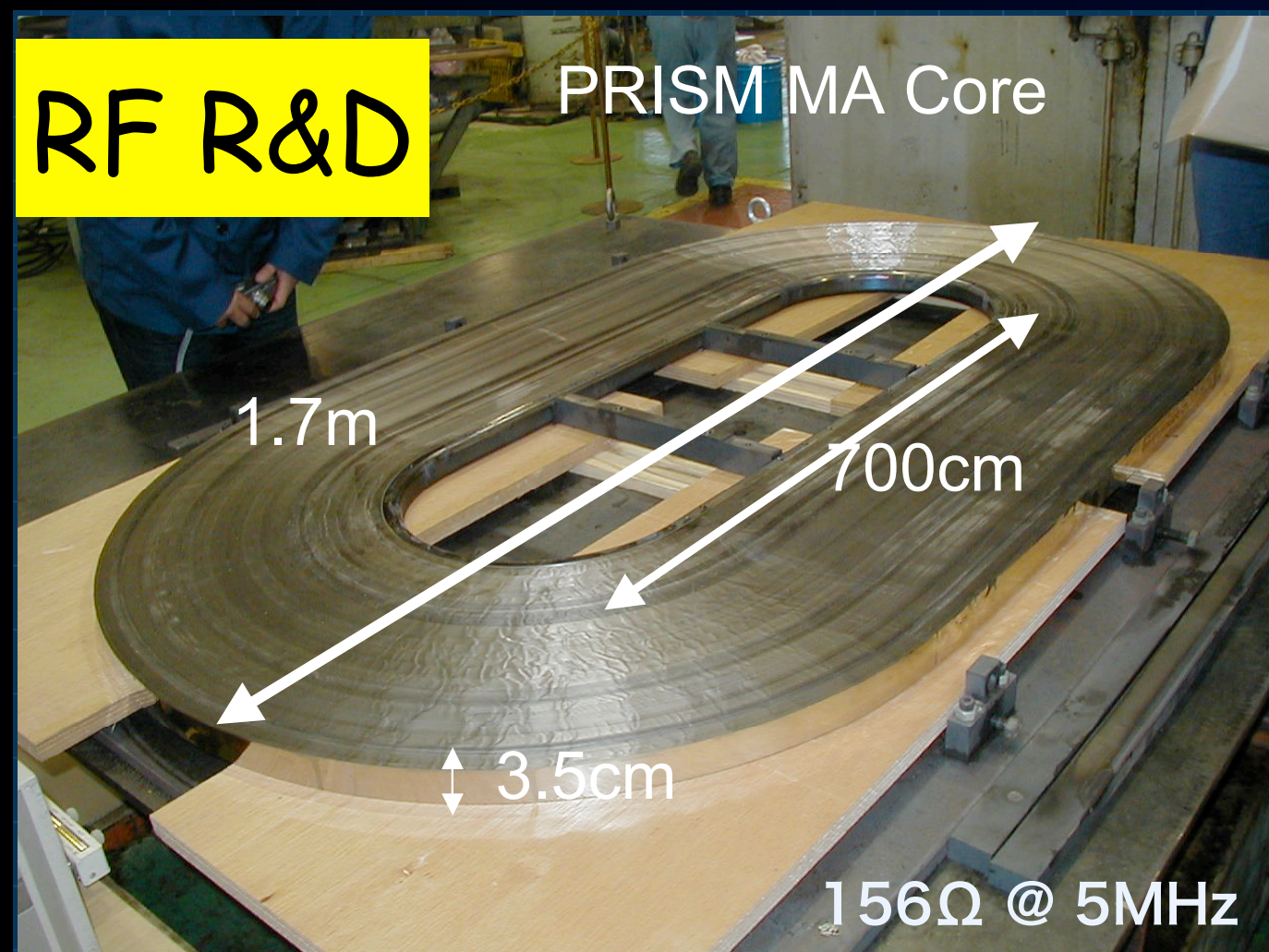


$$B(r) \propto \left(\frac{r}{r_0}\right)^5$$

PRISM-FFAG (6 sectors) in RCNP, Osaka



Demonstration of Phase Rotation

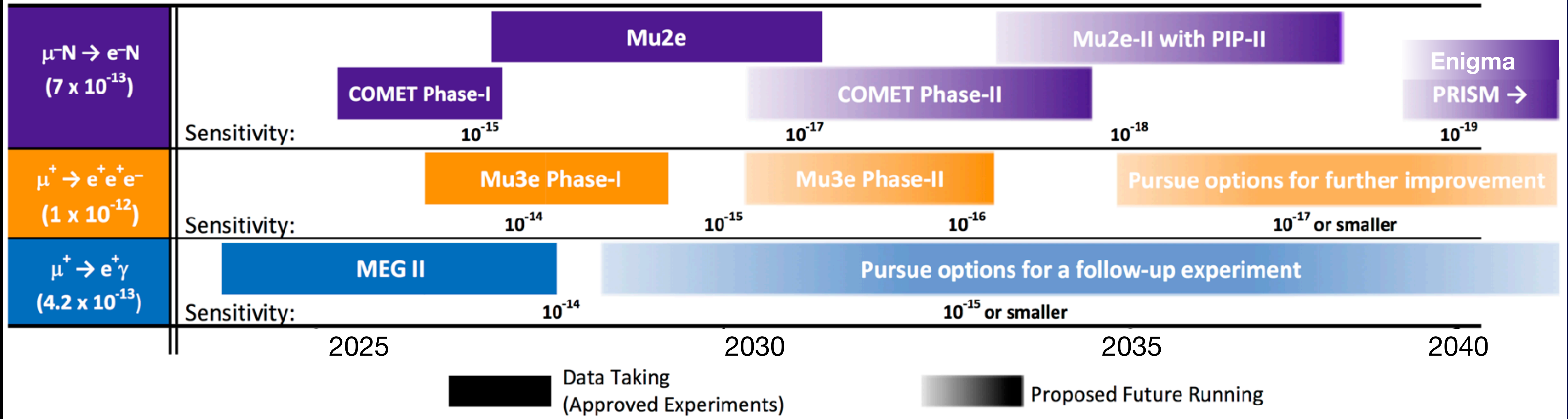


Global Timeline of Muon CLFV



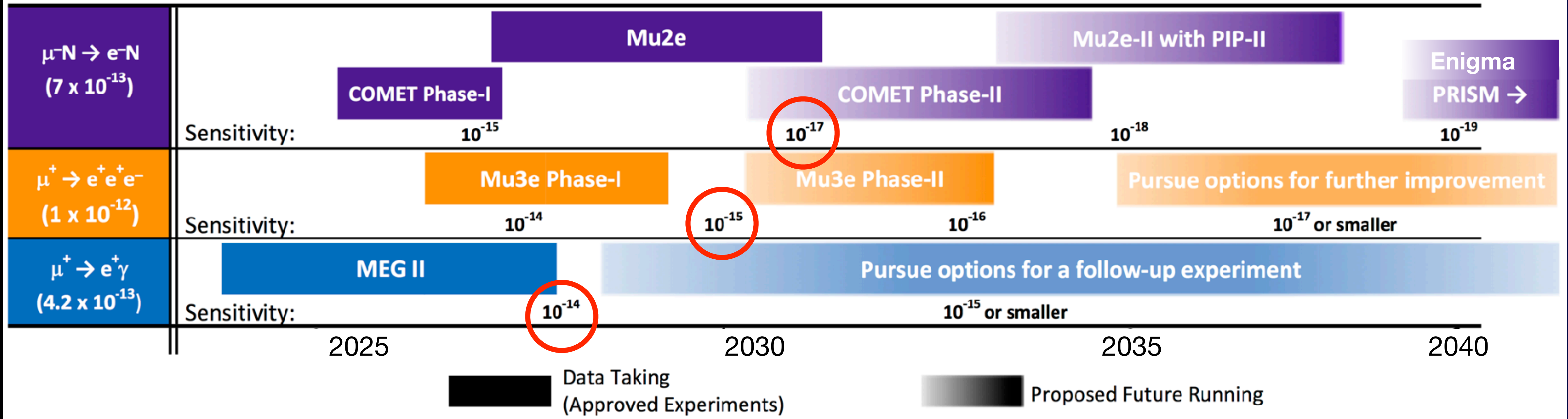
Global Timeline of Muon CLFV

Searches for Charged-Lepton Flavor Violation in Experiments using Intense Muon Beams

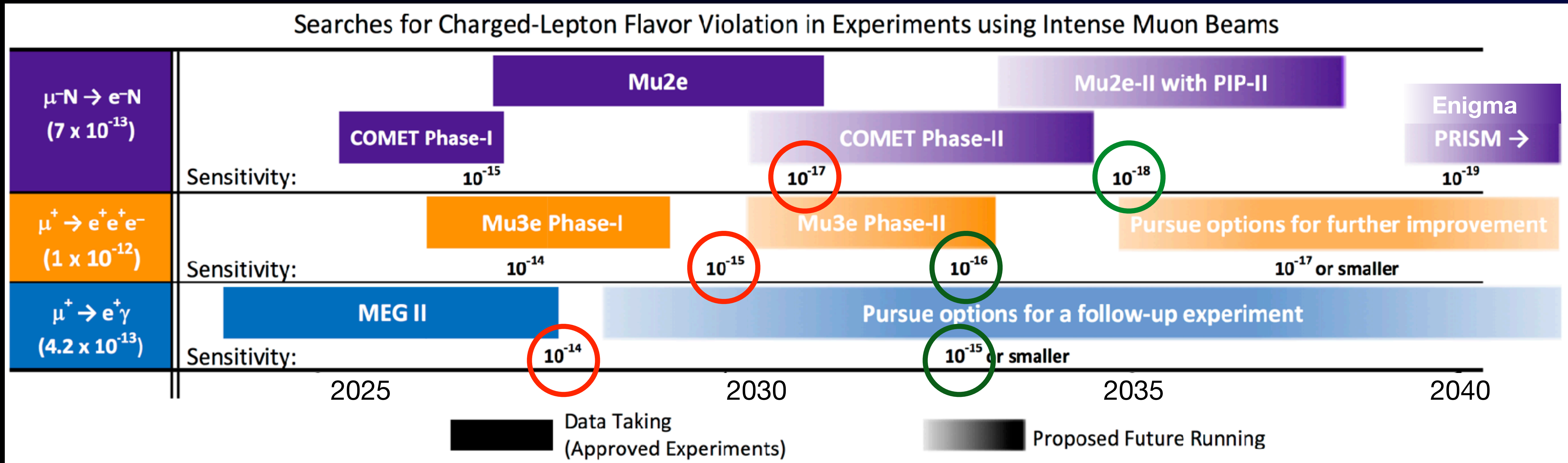


Global Timeline of Muon CLFV

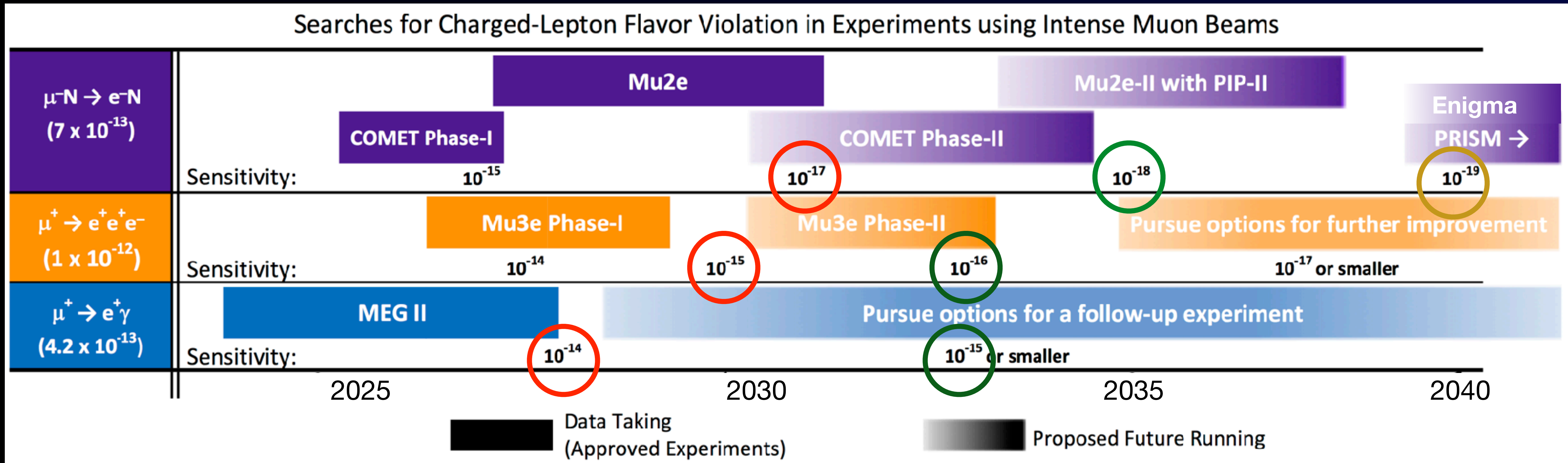
Searches for Charged-Lepton Flavor Violation in Experiments using Intense Muon Beams



Global Timeline of Muon CLFV



Global Timeline of Muon CLFV



In Case of Muon CLFV Discovery...



CLFV is an indirect search but,

In Case of Muon CLFV Discovery...

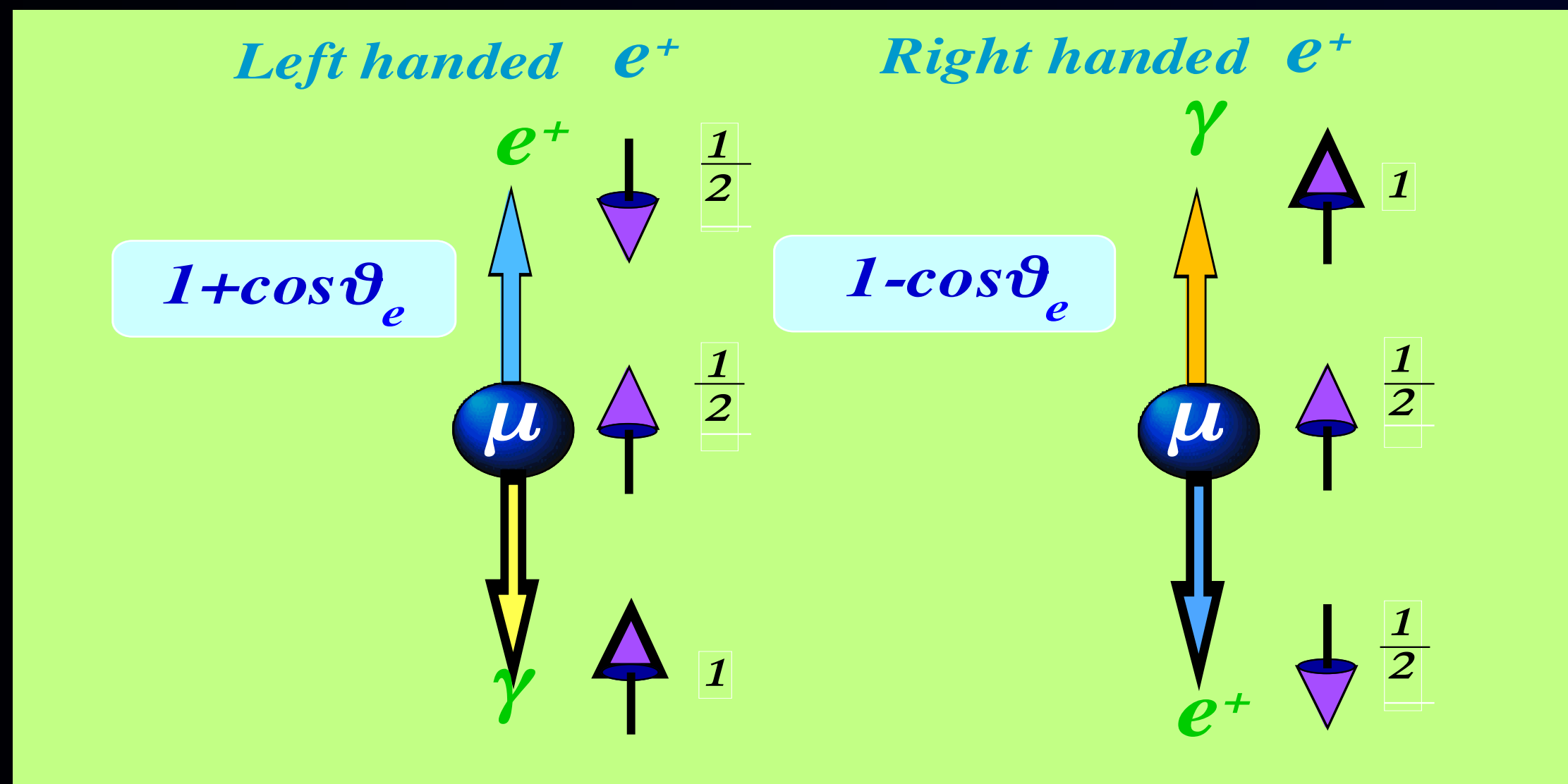


$$\mu^+ \rightarrow e^+ \gamma \text{ and } \mu^+ \rightarrow e^+ e^+ e^-$$



$$\mu^+ \rightarrow e^+ \gamma \text{ and } \mu^+ \rightarrow e^+ e^+ e^-$$

Angular distribution of polarized $\mu^+ \rightarrow e^+ \gamma$ Decay

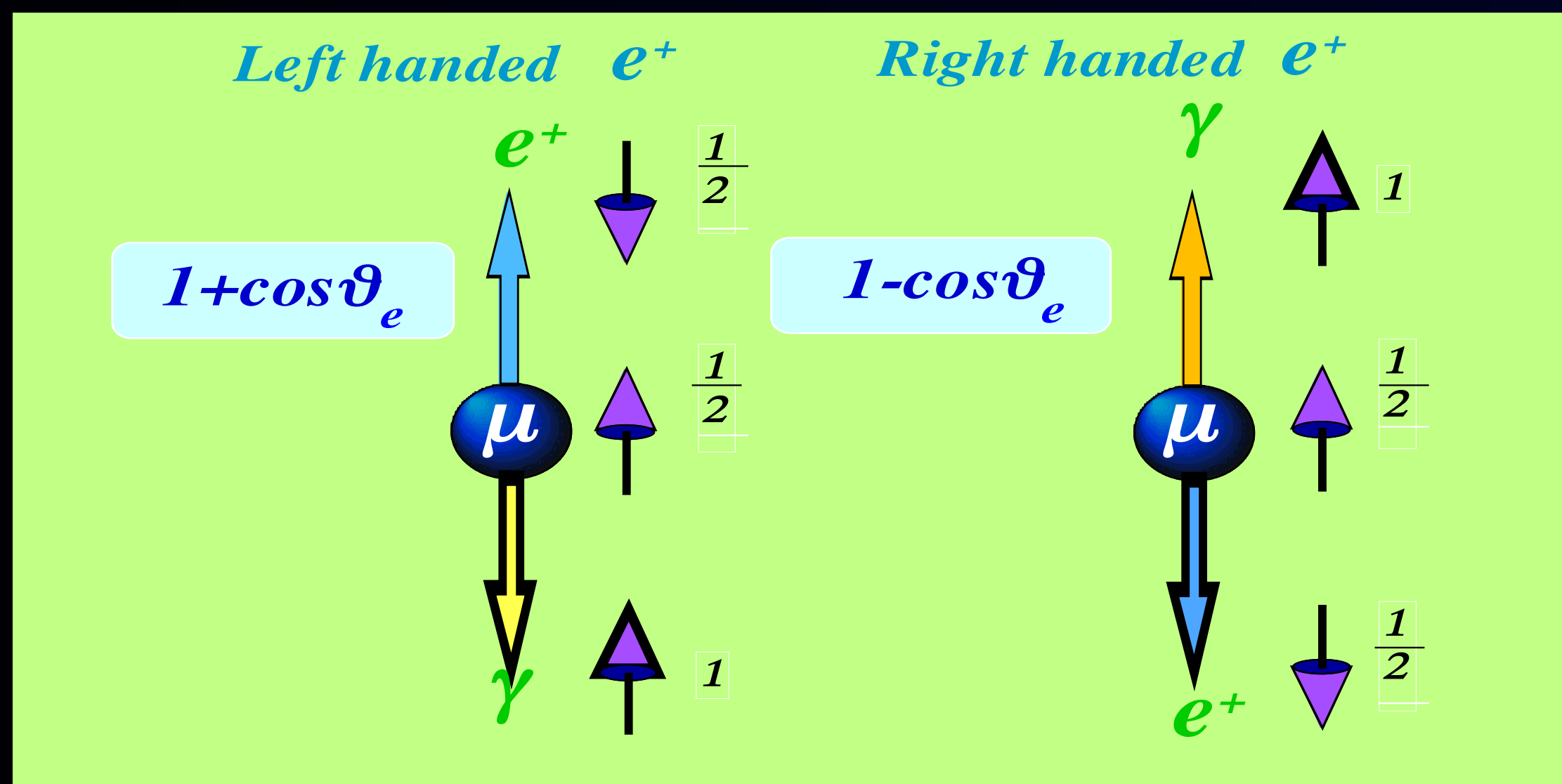


- discriminate models by chirality of positrons.
- surface muon beam is 100% polarized.
- polarized muon decays can be used to suppress backgrounds.

B.W. Lee, et al., Phys. Rev. Lett. 38 (1977) 937
 YK, Y. Okada, Phys. Rev. Lett. 77 (1996) 434
 YK. A. Maki and Y. Okada, Phys.Rev.D 55 (1997) 2517

$$\mu^+ \rightarrow e^+ \gamma \text{ and } \mu^+ \rightarrow e^+ e^+ e^-$$

Angular distribution of polarized $\mu^+ \rightarrow e^+ \gamma$ Decay



- discriminate models by chirality of positrons.
- surface muon beam is 100% polarized.
- polarized muon decays can be used to suppress backgrounds.

B.W. Lee, et al., Phys. Rev. Lett. 38 (1977) 937
 YK, Y. Okada, Phys. Rev. Lett. 77 (1996) 434
 YK. A. Maki and Y. Okada, Phys.Rev.D 55 (1997) 2517

$\mu^+ \rightarrow e^+ e^+ e^-$ Decay

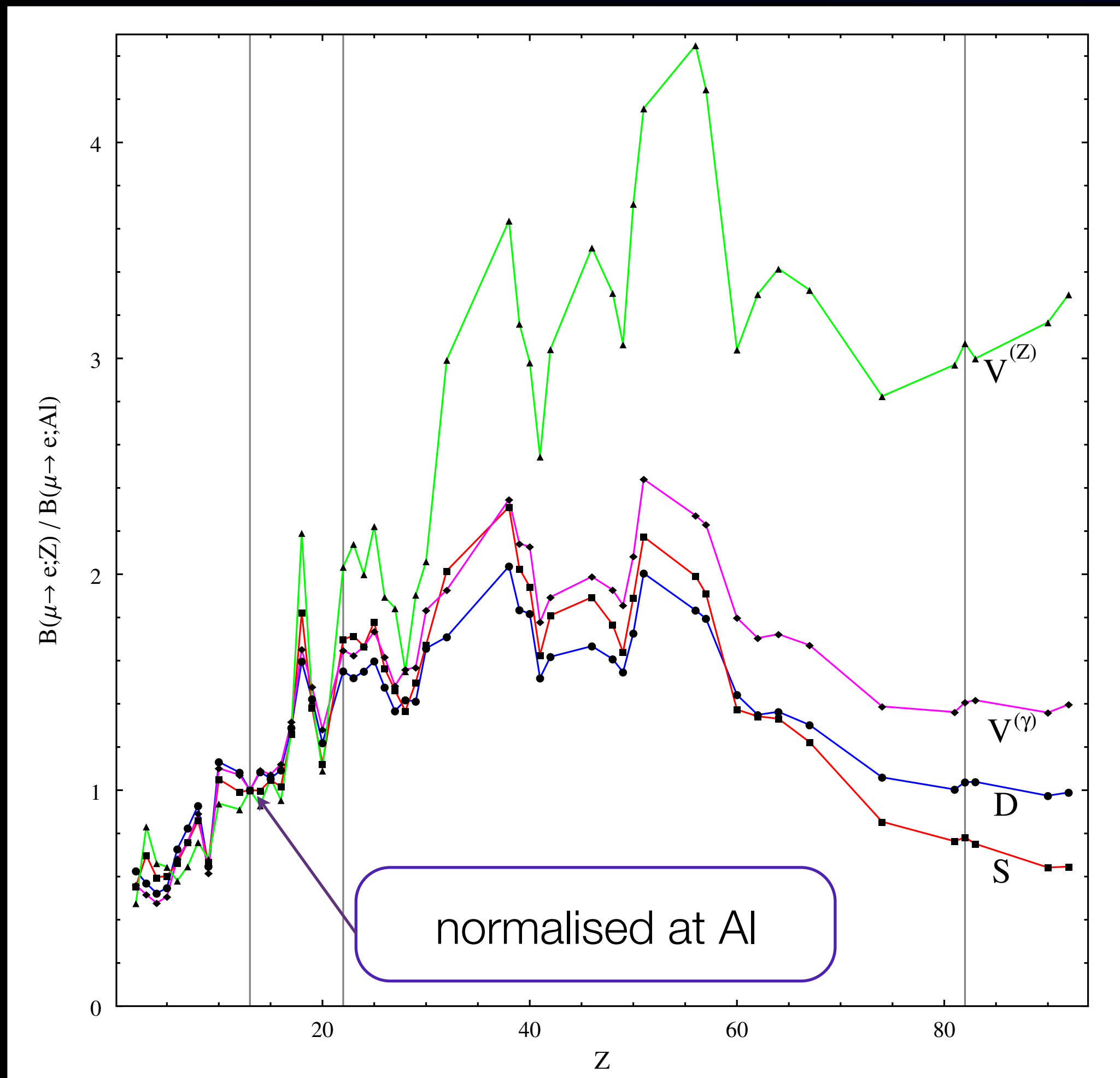
- Dalitz distribution of 3 body decay may be measured for model discrimination and correct rate estimation.

Model Discrimination in $\mu \rightarrow e$ conversion



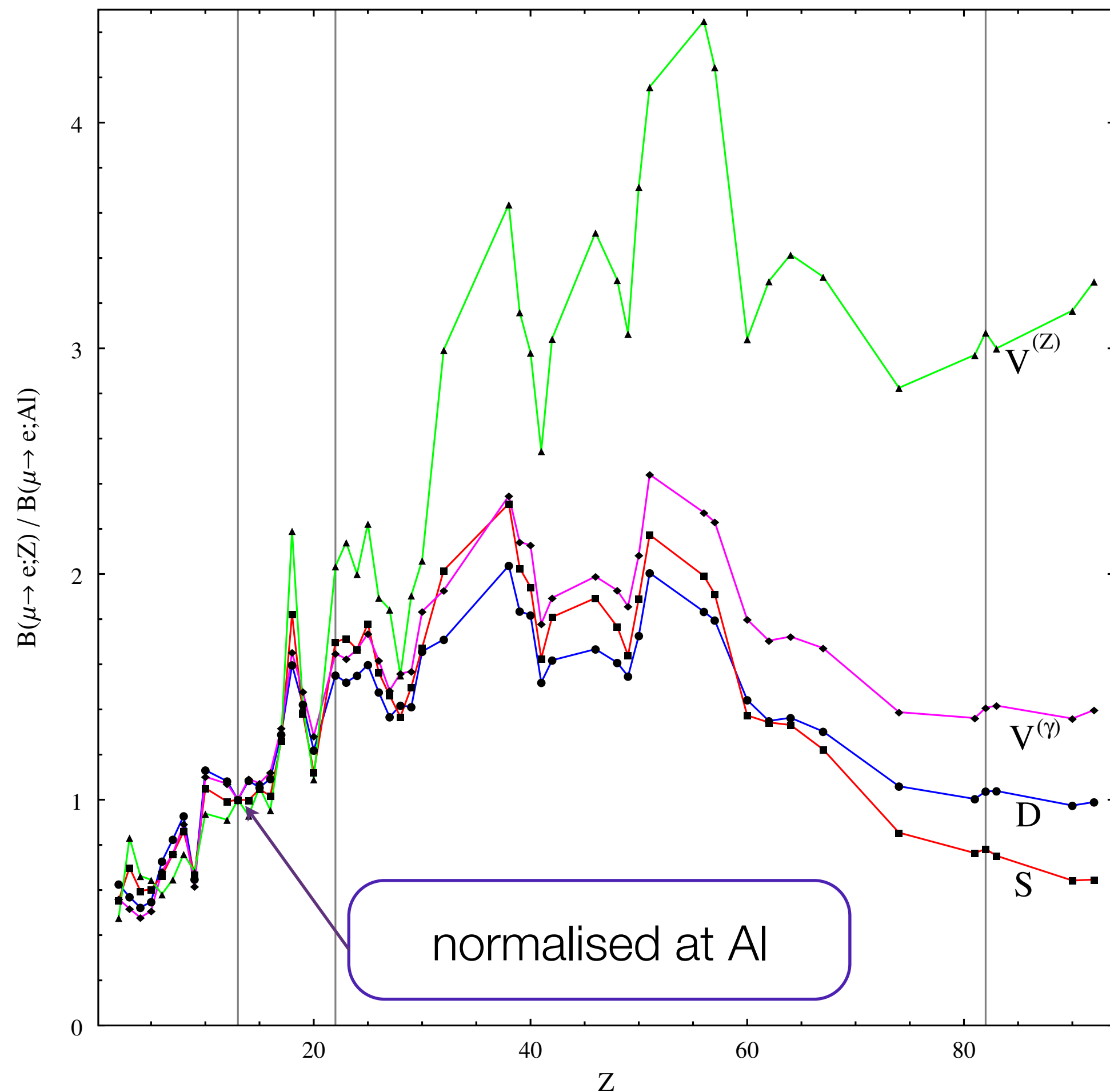
Model Discrimination in $\mu \rightarrow e$ conversion

Muon Target Dependence for $\mu \rightarrow e$ conversion



Model Discrimination in $\mu \rightarrow e$ conversion

Muon Target Dependence for $\mu \rightarrow e$ conversion



vector interaction
(with Z boson)

Left-Right symmetry, Type-III
seesaw, Leptoquark

vector interaction
(charge radius)

dipole interaction

SUSY GUT, SUSY seesaw

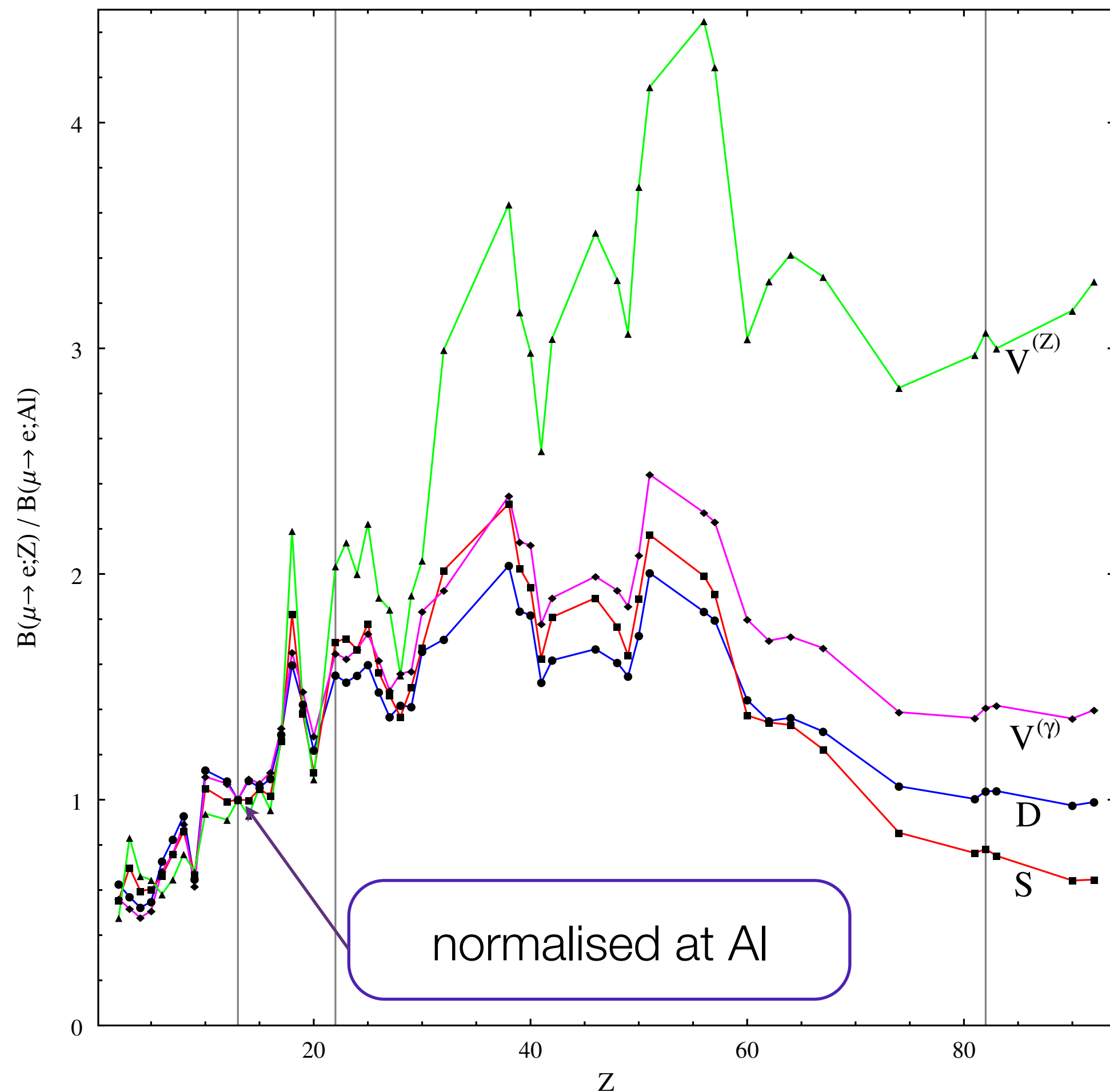
scalar interaction

RPV SUSY, Leptoquark

Model Discrimination in $\mu \rightarrow e$ conversion

Muon Target Dependence for $\mu \rightarrow e$ conversion

single operator analysis



vector interaction
(with Z boson)

Left-Right symmetry, Type-III
seesaw, Leptoquark

vector interaction
(charge radius)

dipole interaction

SUSY GUT, SUSY seesaw

scalar interaction

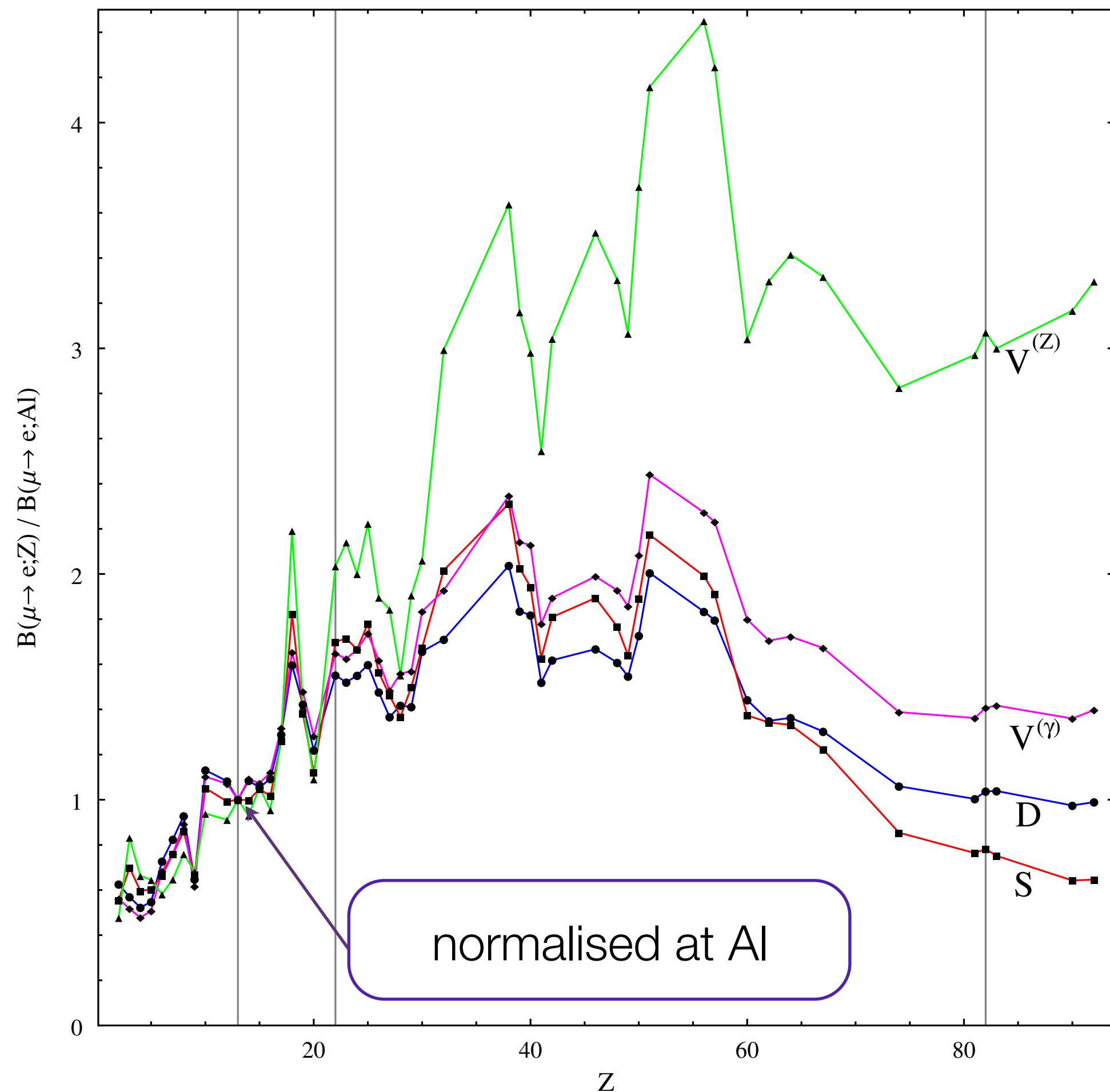
RPV SUSY, Leptoquark

Model Discrimination in $\mu \rightarrow e$ conversion



Muon Target Dependence for $\mu \rightarrow e$ conversion

single operator analysis



vector interaction
(with Z boson)

Left-Right symmetry, Type-III
seesaw, Leptoquark

vector interaction
(charge radius)

dipole interaction

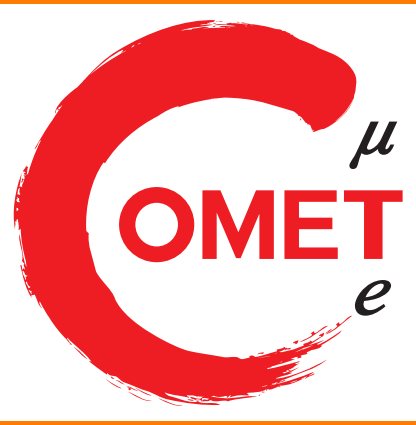
SUSY GUT, SUSY seesaw

scalar interaction

RPV SUSY, Leptoquark

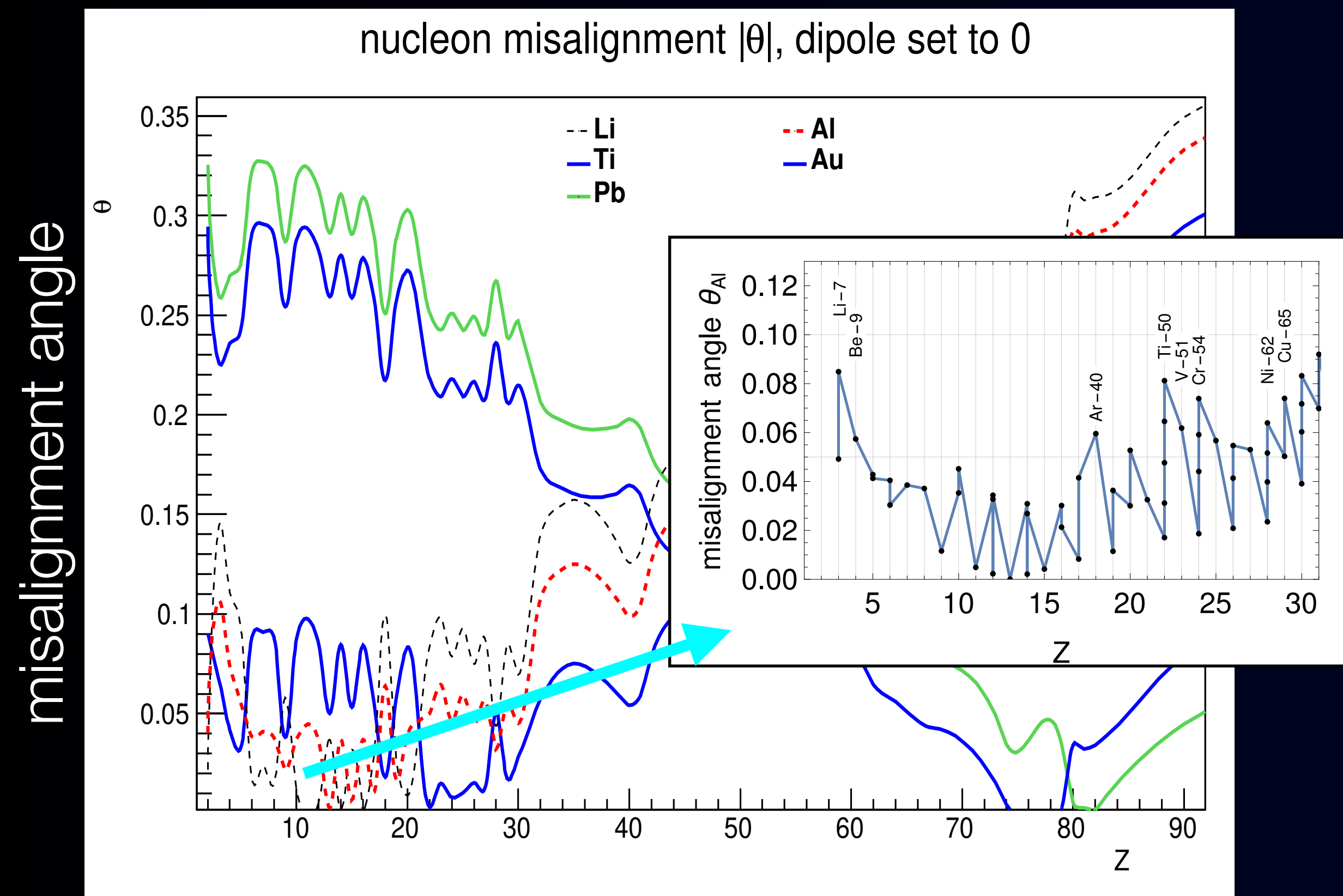
R. Kitano, M. Koike and Y. Okada, Phys. Rev. D 66, 096002 (2002)
 V. Cirigliano, R. Kitano et al., Phys. Rev. D 80, 013002 (2009)
 J. Heeck, R. Szafron, Y. Uesaka, Nucl. Phys. B 980 (2022) 115833
 W.C. Haxton, E. Rule, L. McElvain, et al., Phys. Rev. C 107, 3 (2023)
 E. Rule, W.C. Haxton, K. McElvain, Phys. Rev. Lett. 130, 131901 (2023)
 M. Hoferichter, J. Menendez and F. Noel, Phys. Rev. Lett. 130, 131902 (2023)

Disentanglement in $\mu \rightarrow e$ conversion



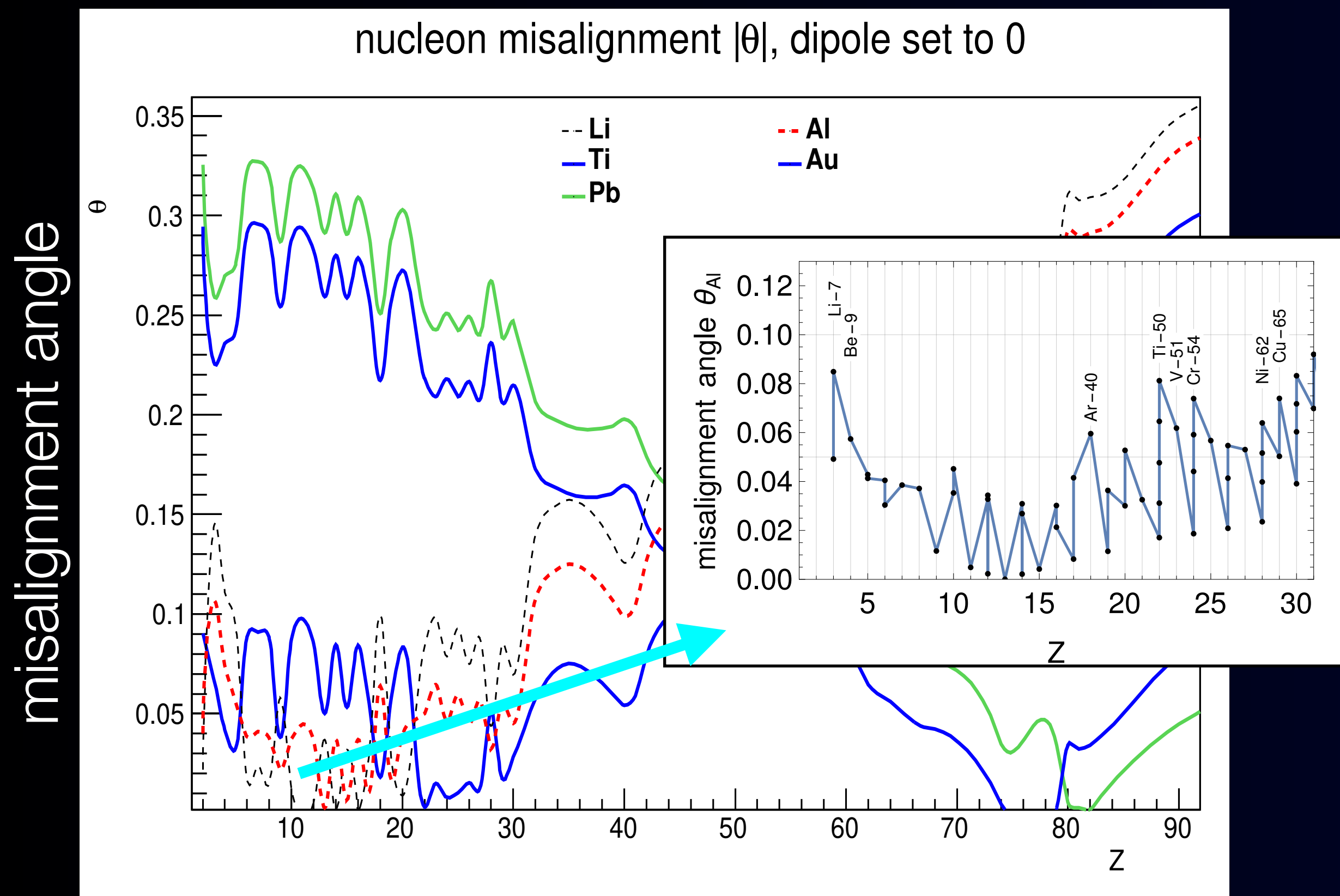
Disentanglement in $\mu \rightarrow e$ conversion

2nd Target selection with different information



Disentanglement in $\mu \rightarrow e$ conversion

2nd Target selection with different information



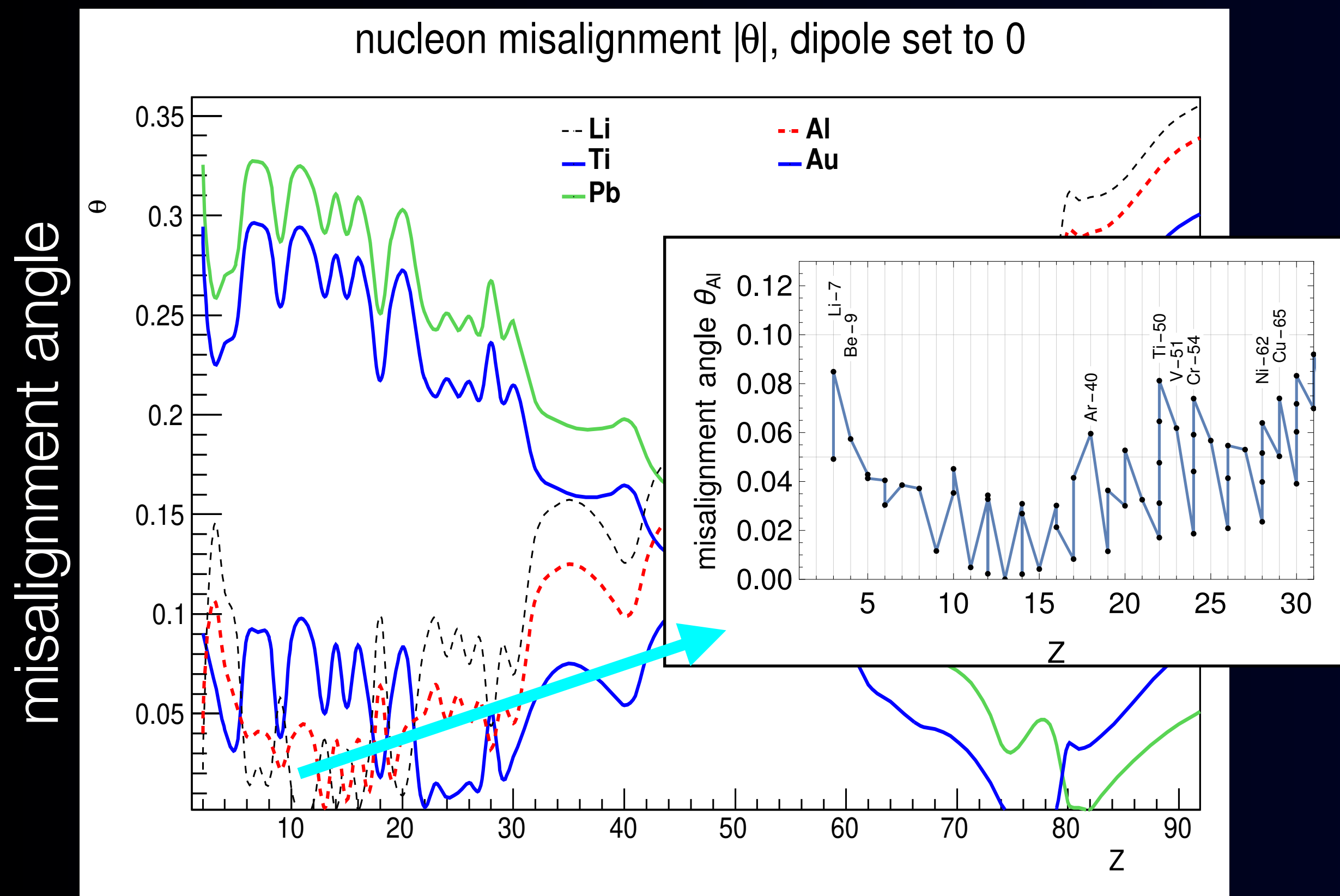
Spin Dependent $\mu \rightarrow e$ conversion

- Spin independent (SI) μ -e Conversion: coherent
- Spin dependent (SD) μ -e Conversion: incoherent
- Extract SD, by comparing zero-spin and non-zero spin targets
 - Effect may not be large by $1/N^2$

V. Cirigliano, S. Davidson, YK, Phys. Lett. B 771 (2017) 242
 S. Davidson, YK, A. Saporta, Eur. Phys. J. C78 (2018) 109

Disentanglement in $\mu \rightarrow e$ conversion

2nd Target selection with different information



Spin Dependent $\mu \rightarrow e$ conversion

- Spin independent (SI) μ -e Conversion: coherent
- Spin dependent (SD) μ -e Conversion: incoherent
- Extract SD, by comparing zero-spin and non-zero spin targets
 - Effect may not be large by $1/N^2$

V. Cirigliano, S. Davidson, YK, Phys. Lett. B 771 (2017) 242
 S. Davidson, YK, A. Saporta, Eur. Phys. J. C78 (2018) 109

Spin polarized $\mu \rightarrow e$ conversion

- Angular distribution with respect to the spin polarization determines the chirality of an electron.

S. Davidson, YK, M. Yamanaka, Phys. Lett. B790 (2019) 380-388
 J. Heeck, R. Szafron, and Y. Uesaka, Nucl.Phys.B 980 (2022) 115833

YK and Y. Okada, RMP 73 (2001) 151

Muon's Lepton Number Violation and Axion-like Particles



$\mu^- \rightarrow e^+$ conversion



$\mu^- \rightarrow e^+$ conversion



- Lepton number violation (LNV) and CLFV
- short distance TeV LNV Physics
- **Event Signature**
 - mono energetic positron (to the ground state)
 - $E_{e^+} = m_\mu - B_\mu - E_{rec} - M(A, Z - 2) + M(A, Z)$
- **Backgrounds**
 - Radiative muon capture (RMC) followed by photon conversion
- **Current limits (from SINDRUM at PSI)**

$\mu^- + \text{Ti} \rightarrow e^+ + \text{Ca}(\text{gs})$	1.7×10^{-12}
$\mu^- + \text{Ti} \rightarrow e^+ + \text{Ca}(\text{ex})$	3.6×10^{-11}

$\mu^- \rightarrow e^+$ conversion



- Lepton number violation (LNV) and CLFV
- short distance TeV LNV Physics

• Event Signature

- mono energetic positron (to the ground state)
- $E_{e^+} = m_\mu - B_\mu - E_{rec} - M(A, Z - 2) + M(A, Z)$

• Backgrounds

- Radiative muon capture (RMC) followed by photon conversion

• Current limits (from SINDRUM at PSI)

$\mu^- + \text{Ti} \rightarrow e^+ + \text{Ca}(\text{gs})$	1.7×10^{-12}
$\mu^- + \text{Ti} \rightarrow e^+ + \text{Ca}(\text{ex})$	3.6×10^{-11}

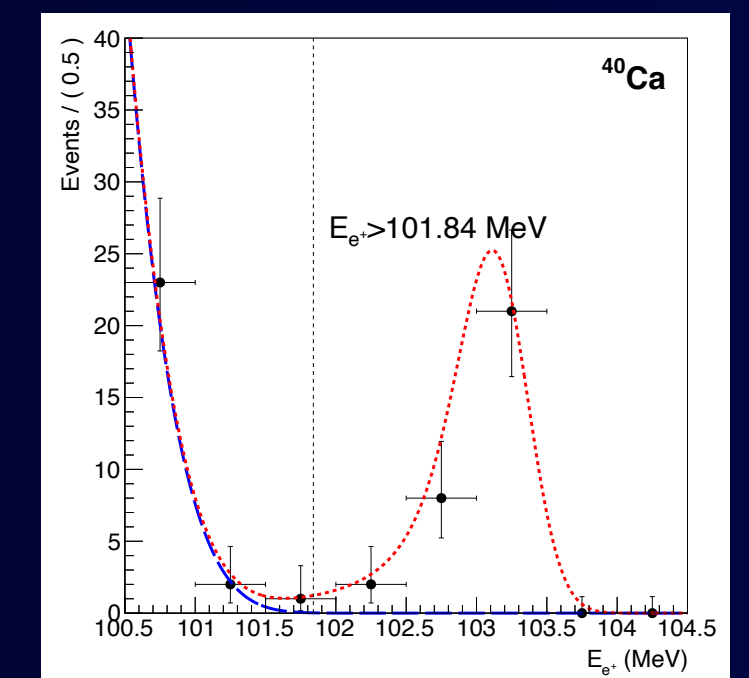
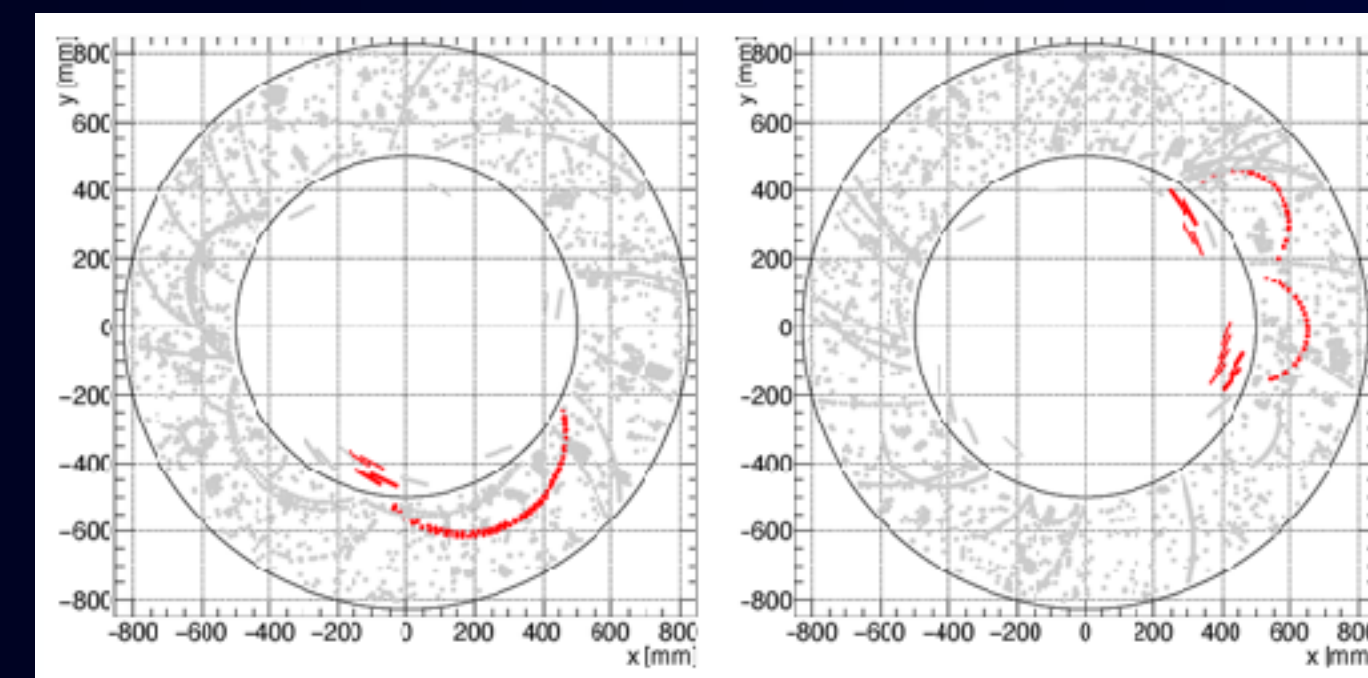
J. Kaulard et al. (SINDRUM-II), Phys. Lett. B422 (1998) 334.

- requirement of candidate target $N(A, Z)$

$$E_{\mu^-e^+} > E_{RMC}^{end} \longrightarrow M(A, Z - 1) > M(A, Z - 2)$$

B. Yeo, YK, M. Lee and K. Zuber, Phys. Rev. D96 (2017) 075027

- Mu2e and COMET are trying to measure at the same time of $\mu^- \rightarrow e^-$ conversion
 - x10000 improvement can be expected
 - RMC background spectrum, better to measured.



$\mu \rightarrow ea$ and $\mu \rightarrow eay$



$\mu \rightarrow ea$ and $\mu \rightarrow ea\gamma$

$$\mu^+ \rightarrow e^+ a$$

a is a light, invisible, neutral particle (ALP) with **LFV** coupling to leptons.

$\mu \rightarrow ea$ and $\mu \rightarrow ea\gamma$

$$\mu^+ \rightarrow e^+ a$$

a is a light, invisible, neutral particle (ALP) with **LFV** coupling to leptons.

$$\mathcal{L}_{all} = \frac{\partial^\mu a}{2f_a} (C_{ij}^V \bar{l}_i \gamma_\mu l_j + C_{ij}^A \bar{l}_i \gamma_\mu \gamma_5 l_j)$$

$$\Gamma(l_i \rightarrow l_j a) = \frac{1}{16\pi} \frac{m_{l_i}^3}{F_{ij}^2} \left(1 - \frac{m_a^2}{m_{l_i}^2}\right)^2 \quad F_{ij} \equiv \frac{2f_a}{\sqrt{|C_{ij}^V|^2 + |C_{ij}^A|^2}}$$

$\mu \rightarrow ea$ and $\mu \rightarrow ea\gamma$

$$\mu^+ \rightarrow e^+ a$$

a is a light, invisible, neutral particle (ALP) with **LFV** coupling to leptons.

$$\mathcal{L}_{all} = \frac{\partial^\mu a}{2f_a} (C_{ij}^V \bar{l}_i \gamma_\mu l_j + C_{ij}^A \bar{l}_i \gamma_\mu \gamma_5 l_j)$$

$$\Gamma(l_i \rightarrow l_j a) = \frac{1}{16\pi} \frac{m_{l_i}^3}{F_{ij}^2} \left(1 - \frac{m_a^2}{m_{l_i}^2}\right)^2 \quad F_{ij} \equiv \frac{2f_a}{\sqrt{|C_{ij}^V|^2 + |C_{ij}^A|^2}}$$

	upper limits	$F_{e\mu}$ (GeV)	
TRIUMF	$\text{BR}(\mu^+ \rightarrow e^+ a_{\text{RH}}) < 2.6 \times 10^{-6}$	$> 5.5 \times 10^9$	1988
online	$\text{BR}(\mu^+ \rightarrow e^+ a) < 5.8 \times 10^{-5}$	$> 1.2 \times 10^9$	2015
Los Alamos	$\text{BR}(\mu^+ \rightarrow e^+ a\gamma) < 1.1 \times 10^{-9}$	$> 9.8 \times 10^8$	1988
MEG-II forward	$\text{BR}(\mu^+ \rightarrow e^+ a_{\text{RH}}) < 10^{-7}$	$> 10^{10}$?
Mu3e-online	$\text{BR}(\mu^+ \rightarrow e^+ a) < 10^{-8}$	$> 10^{10}$?
MEG-II	$\text{BR}(\mu^+ \rightarrow e^+ a\gamma) < 10^{-10}$	$> 10^{10}$?

$\mu \rightarrow ea$ and $\mu \rightarrow ea\gamma$

$$\mu^+ \rightarrow e^+ a$$

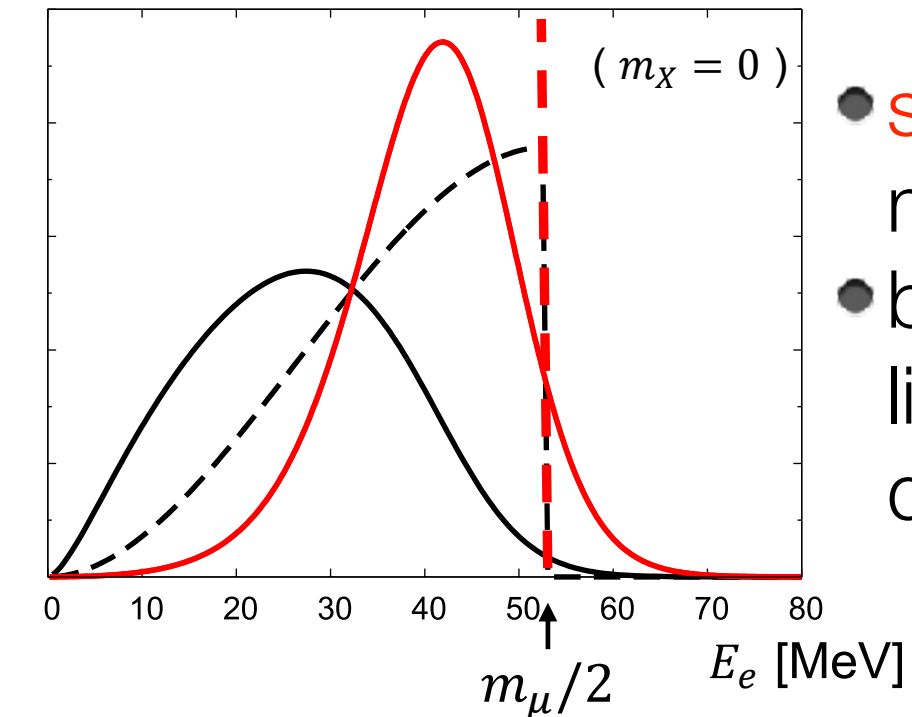
a is a light, invisible, neutral particle (ALP) with **LFV** coupling to leptons.

$$\mathcal{L}_{all} = \frac{\partial^\mu a}{2f_a} (C_{ij}^V \bar{l}_i \gamma_\mu l_j + C_{ij}^A \bar{l}_i \gamma_\mu \gamma_5 l_j)$$

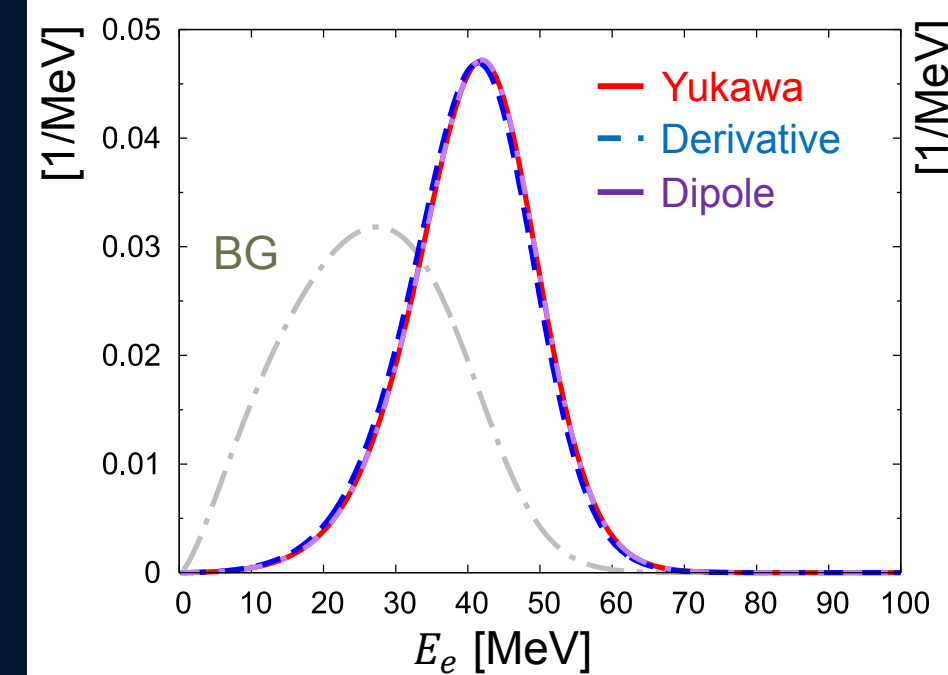
$$\Gamma(l_i \rightarrow l_j a) = \frac{1}{16\pi} \frac{m_{l_i}^3}{F_{ij}^2} \left(1 - \frac{m_a^2}{m_{l_i}^2}\right)^2 \quad F_{ij} \equiv \frac{2f_a}{\sqrt{|C_{ij}^V|^2 + |C_{ij}^A|^2}}$$

$$\text{bound } \mu^- \rightarrow e^- a$$

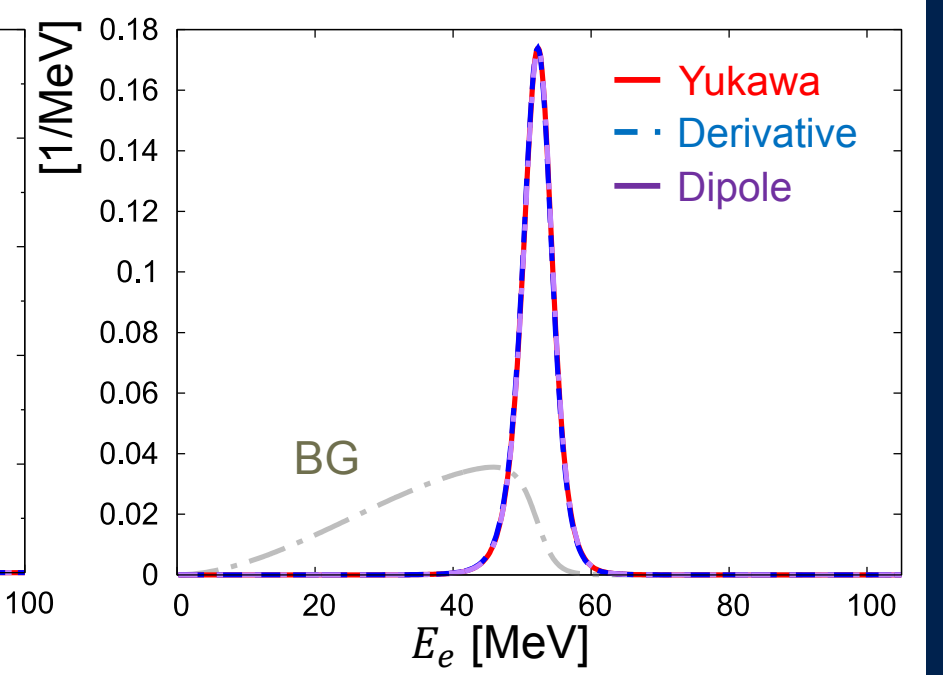
electron spectra (normalized by rate)



$\frac{1}{\Gamma} \frac{d\Gamma}{dE_e}$ ^{197}Au



^{27}Al



Y. Uesaka, Phys. Rev. D102, 095007 (2020)
Tianyu Xing et al., Chin. Phys. C47, 013108 (2023)

	upper limits	$F_{e\mu}$ (GeV)	
TRIUMF	$\text{BR}(\mu^+ \rightarrow e^+ a_{\text{RH}}) < 2.6 \times 10^{-6}$	$> 5.5 \times 10^9$	1988
online	$\text{BR}(\mu^+ \rightarrow e^+ a) < 5.8 \times 10^{-5}$	$> 1.2 \times 10^9$	2015
Los Alamos	$\text{BR}(\mu^+ \rightarrow e^+ a\gamma) < 1.1 \times 10^{-9}$	$> 9.8 \times 10^8$	1988
MEG-II forward	$\text{BR}(\mu^+ \rightarrow e^+ a_{\text{RH}}) < 10^{-7}$	$> 10^{10}$?
Mu3e-online	$\text{BR}(\mu^+ \rightarrow e^+ a) < 10^{-8}$	$> 10^{10}$?
MEG-II	$\text{BR}(\mu^+ \rightarrow e^+ a\gamma) < 10^{-10}$	$> 10^{10}$?

Muonium



Muonium to Antimuonium Conversion



Muonium to Antimuonium Conversion



$$\Delta L_\mu = 2, \Delta L_e = 2$$

Muonium to Antimuonium Conversion



$$\Delta L_\mu = 2, \Delta L_e = 2$$

- models: doubly-charged Higgs etc.

Muonium to Antimuonium Conversion

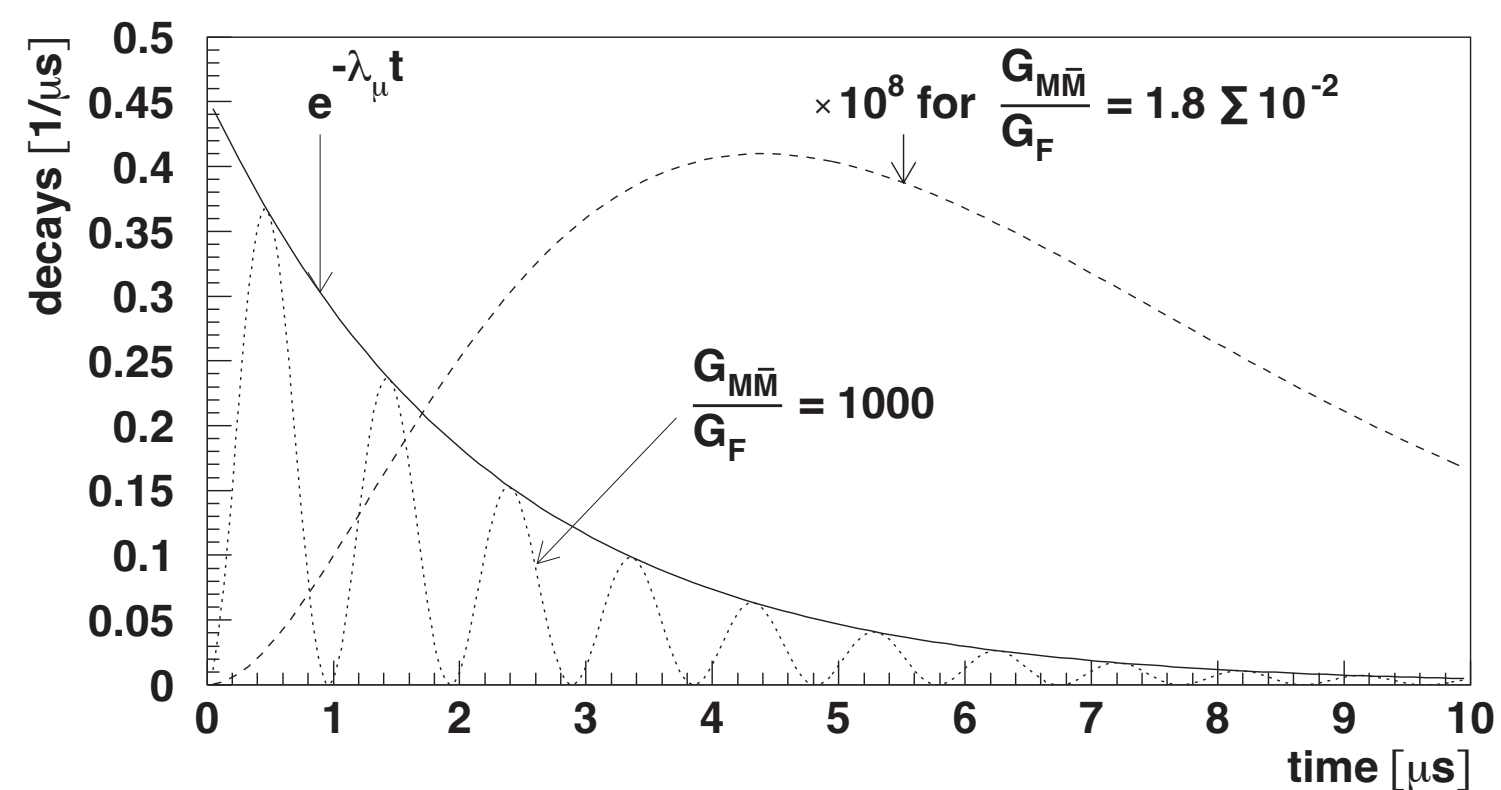


$$\Delta L_\mu = 2, \Delta L_e = 2$$

- models: doubly-charged Higgs etc.
- Oscillation probability

$$P_{\text{Mu}\overline{\text{Mu}}} = \sin^2\left(\frac{t \times \Delta}{2}\right) \exp\left[-\frac{t}{\tau_\mu}\right]$$

$$\Delta = 2 \langle \overline{\text{Mu}} | H_{\text{Mu}\overline{\text{Mu}}} | \text{Mu} \rangle$$



Oscillation
maximum
at $t = 2\tau_\mu$

Muonium to Antimuonium Conversion

$$\mu^+ e^- (\text{Mu}) \rightarrow \mu^- e^+ (\overline{\text{Mu}})$$

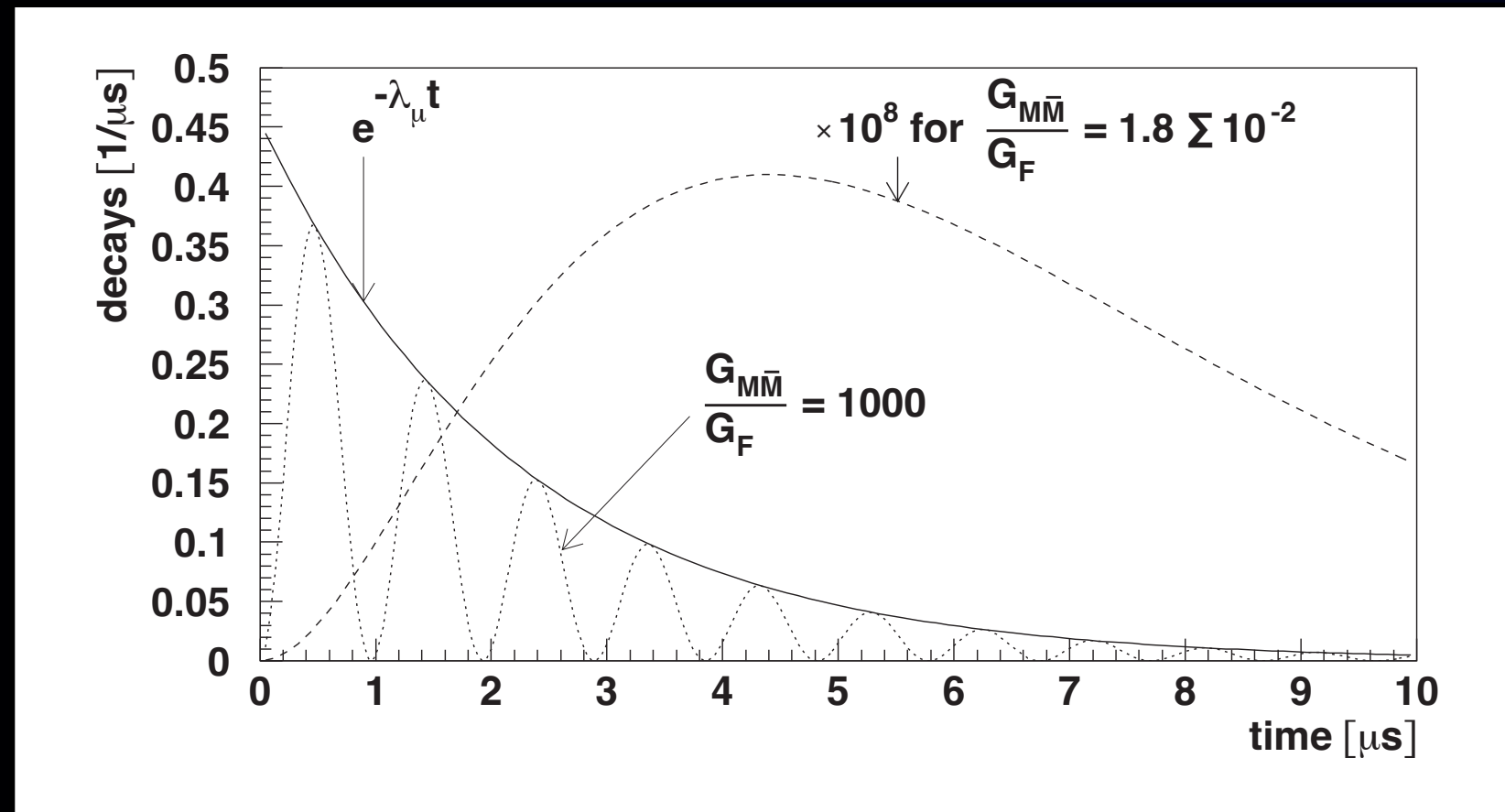
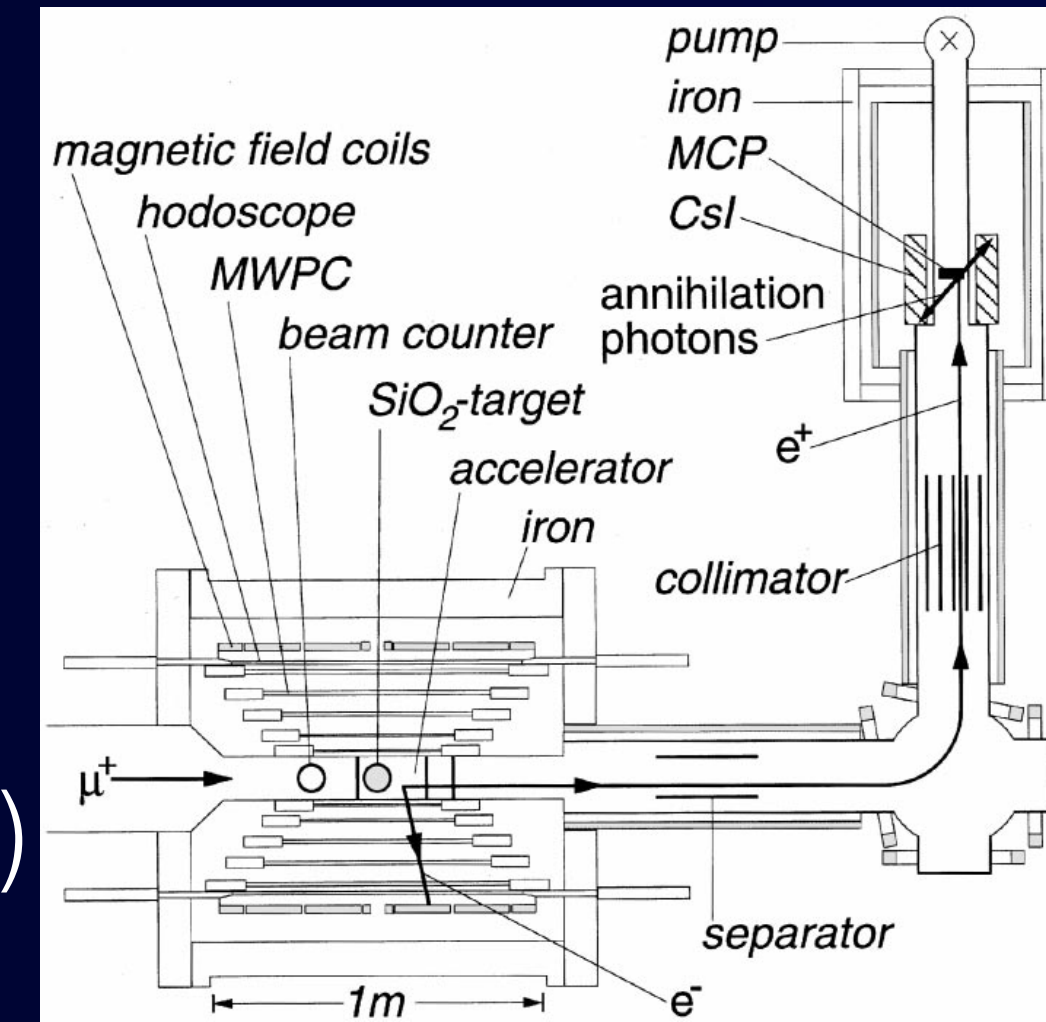
$$\Delta L_\mu = 2, \Delta L_e = 2$$

- models: doubly-charged Higgs etc.
- Oscillation probability

$$P_{\text{Mu}\overline{\text{Mu}}} = \sin^2\left(\frac{t \times \Delta}{2}\right) \exp\left[-\frac{t}{\tau_\mu}\right]$$

$$\Delta = 2 \langle \overline{\text{Mu}} | H_{\text{Mu}\overline{\text{Mu}}} | \text{Mu} \rangle$$

- Experimental methods
 - Production of Mu in vacuum
 - reduce residual EM fields
 - Detection of $\overline{\text{Mu}}$ decay
 - Current limit (from PSI, 1999)
 - $P_{\text{Mu}\overline{\text{Mu}}} \leq 8.3 \times 10^{-11}$ (90% C.L.)



Oscillation maximum at $t = 2\tau_\mu$

Muonium to Antimuonium Conversion

$$\mu^+ e^- (\text{Mu}) \rightarrow \mu^- e^+ (\overline{\text{Mu}})$$

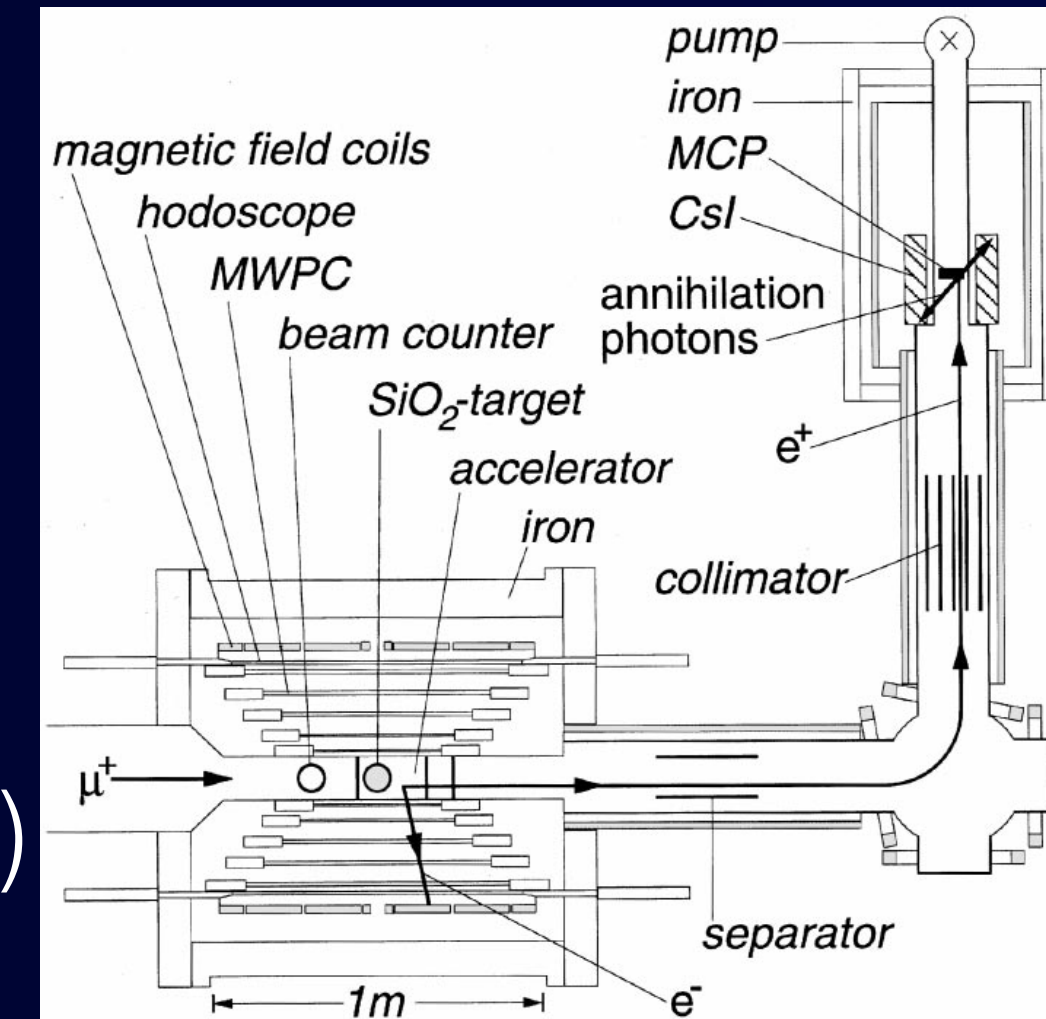
$$\Delta L_\mu = 2, \Delta L_e = 2$$

- models: doubly-charged Higgs etc.
- Oscillation probability

$$P_{\text{Mu}\overline{\text{Mu}}} = \sin^2\left(\frac{t \times \Delta}{2}\right) \exp\left[-\frac{t}{\tau_\mu}\right]$$

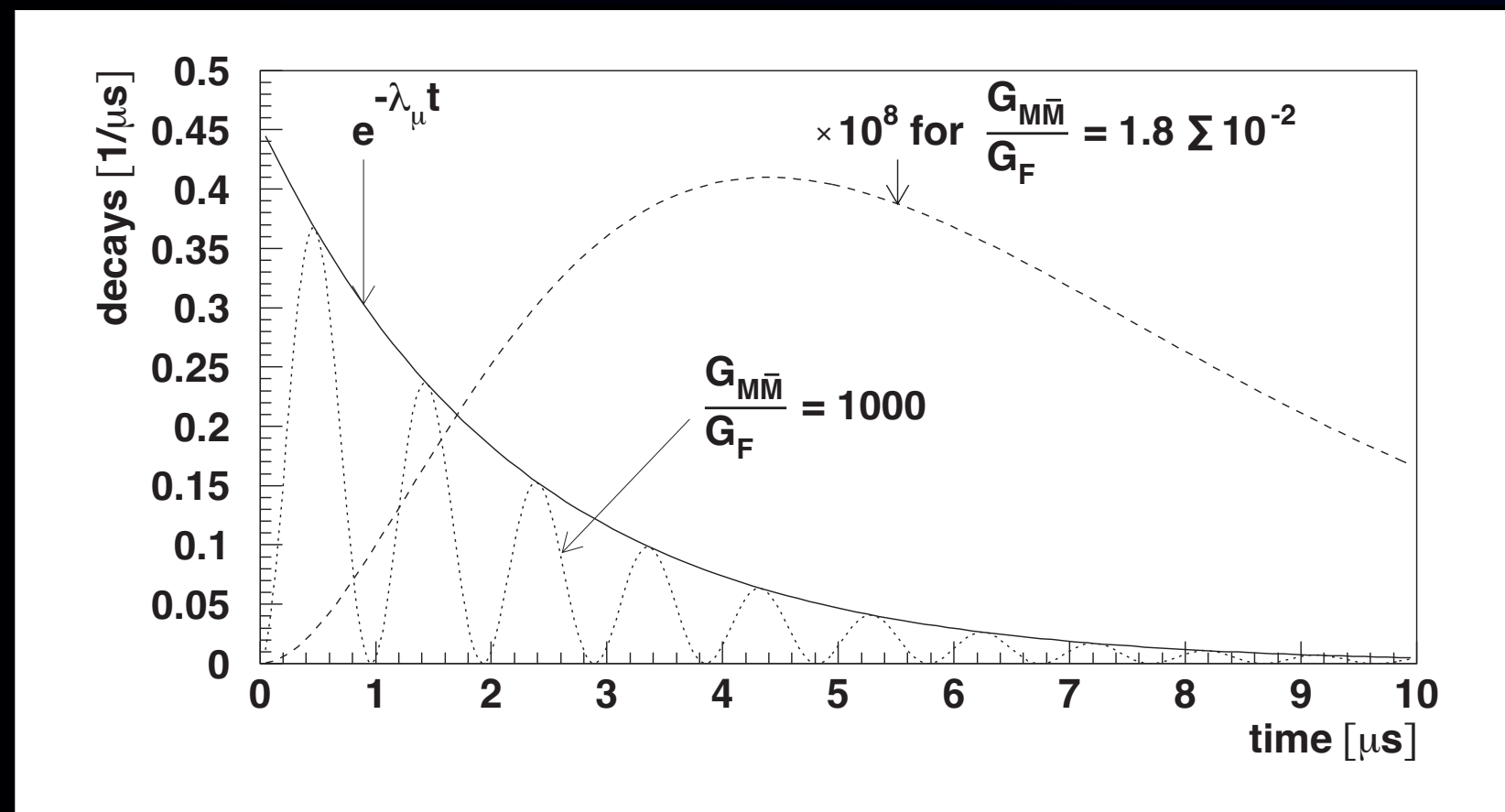
$$\Delta = 2 \langle \overline{\text{Mu}} | H_{\text{Mu}\overline{\text{Mu}}} | \text{Mu} \rangle$$

- Experimental methods
 - Production of Mu in vacuum
 - reduce residual EM fields
 - Detection of $\overline{\text{Mu}}$ decay
 - Current limit (from PSI, 1999)
 - $P_{\text{Mu}\overline{\text{Mu}}} \leq 8.3 \times 10^{-11}$ (90% C.L.)



- New proposals
 - **CSNS in China:** X(10-100) sensitivity

See the talk by Yuhang Guo



Oscillation maximum at $t = 2\tau_\mu$

Muon Colliders



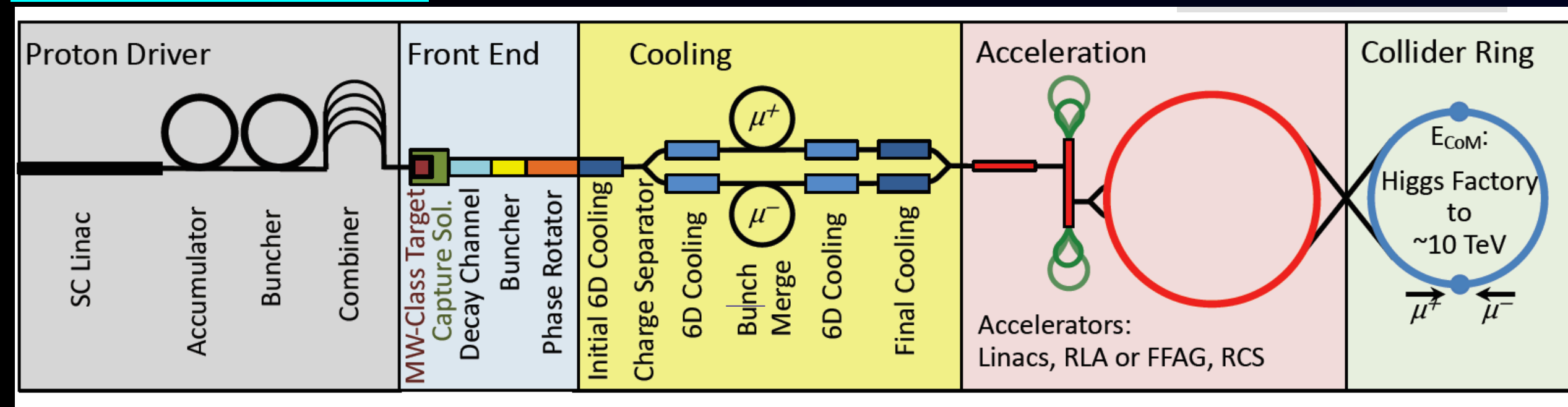
Muon Colliders



Muon Colliders



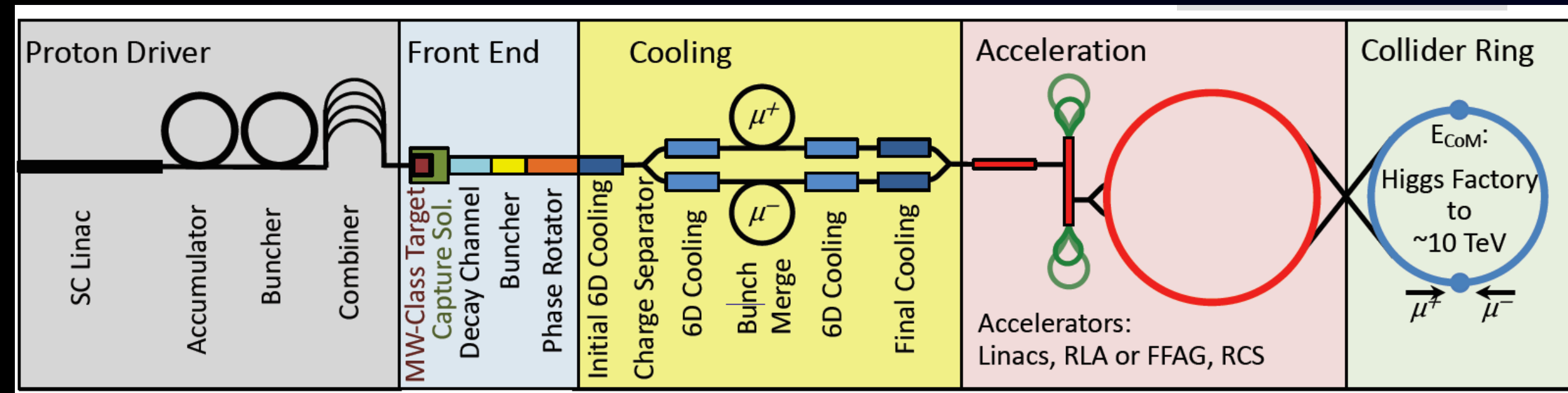
MAP (US)



Muon Colliders

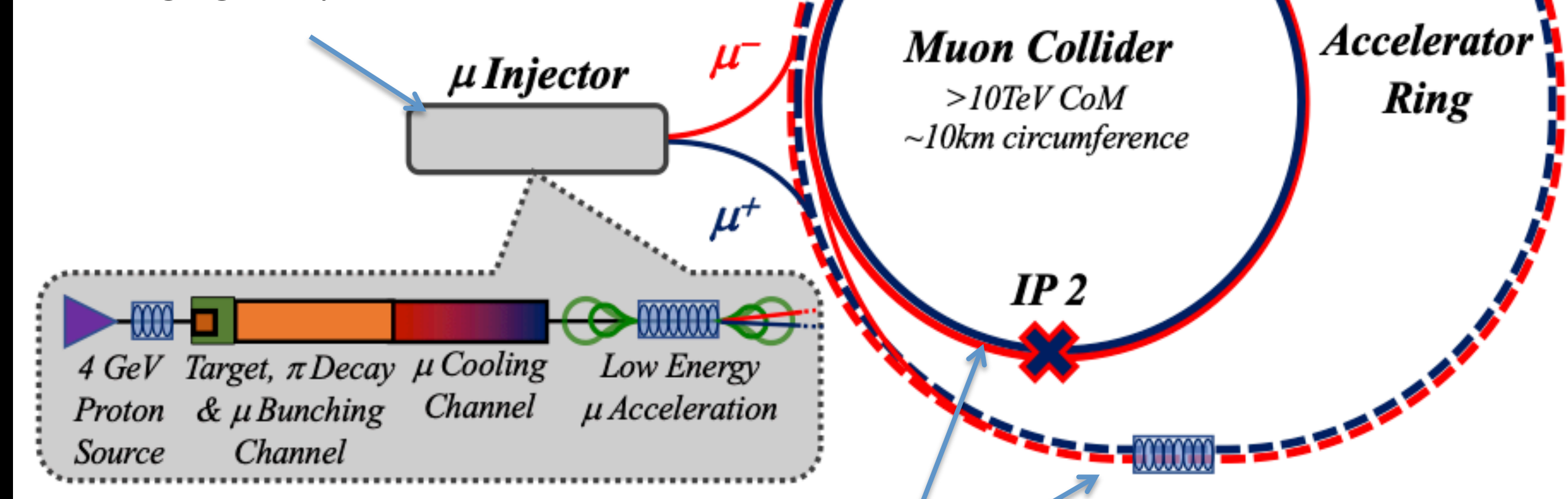


MAP (US)



CERN

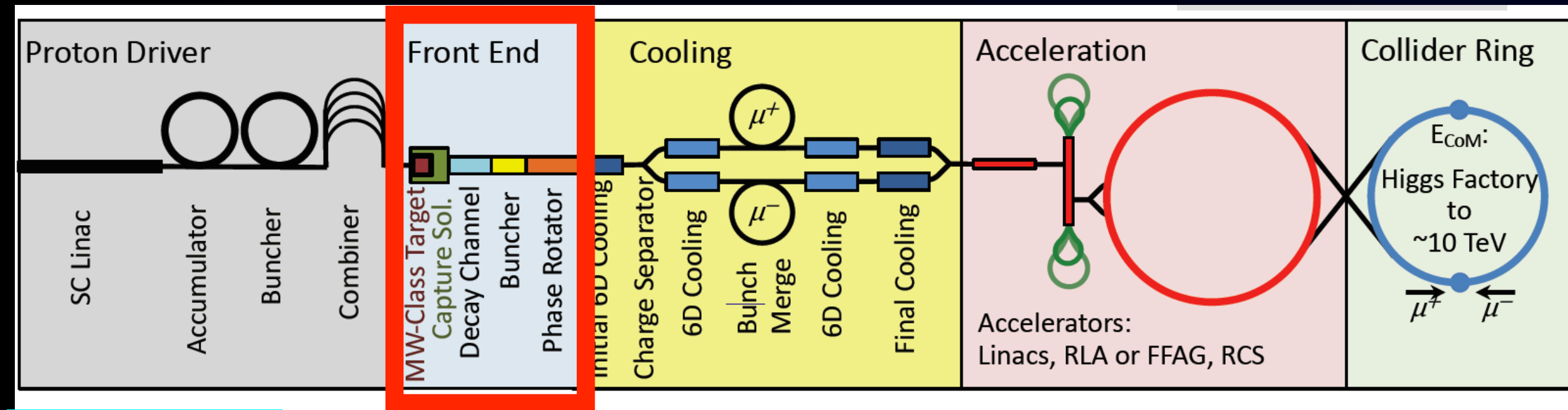
Drives the beam quality similar to MAP design
still challenging design with challenging components



Muon Colliders

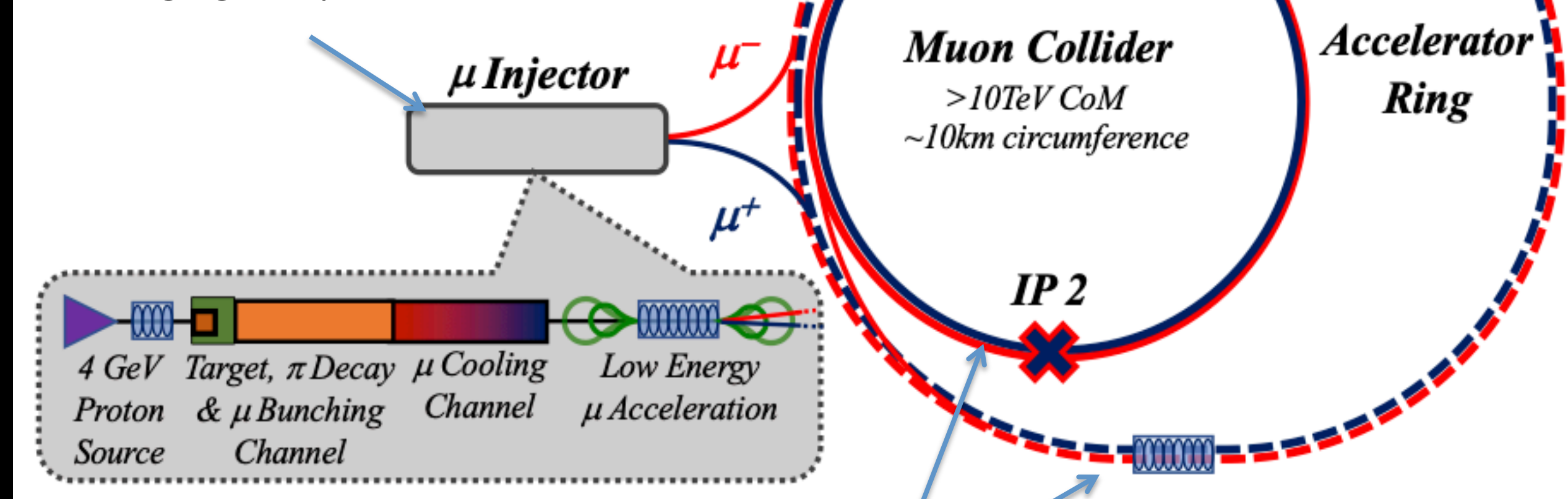


MAP (US)



CERN

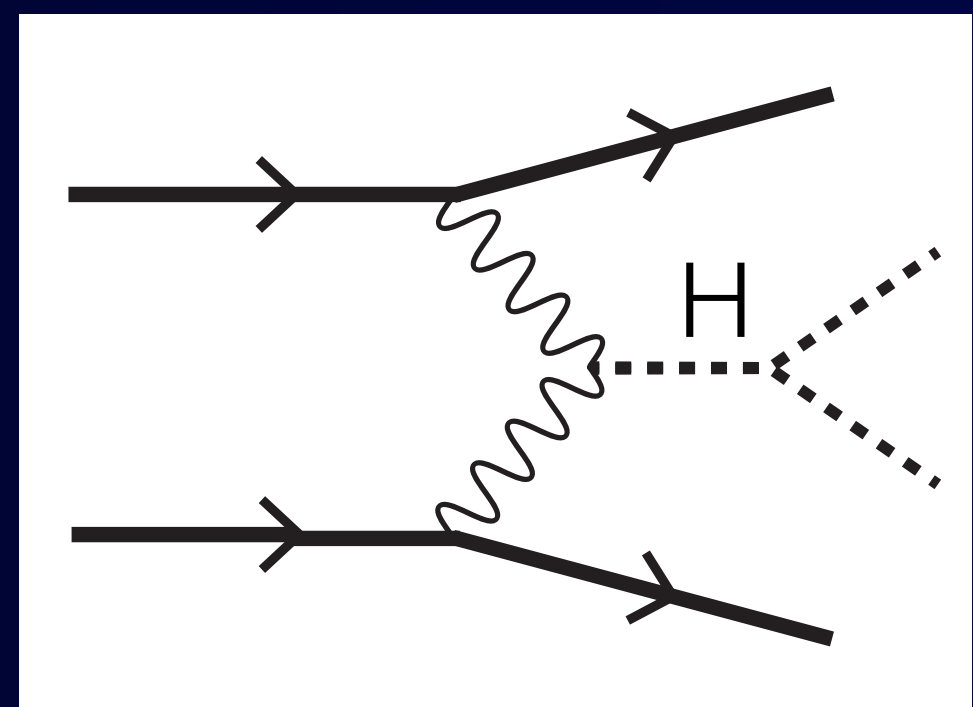
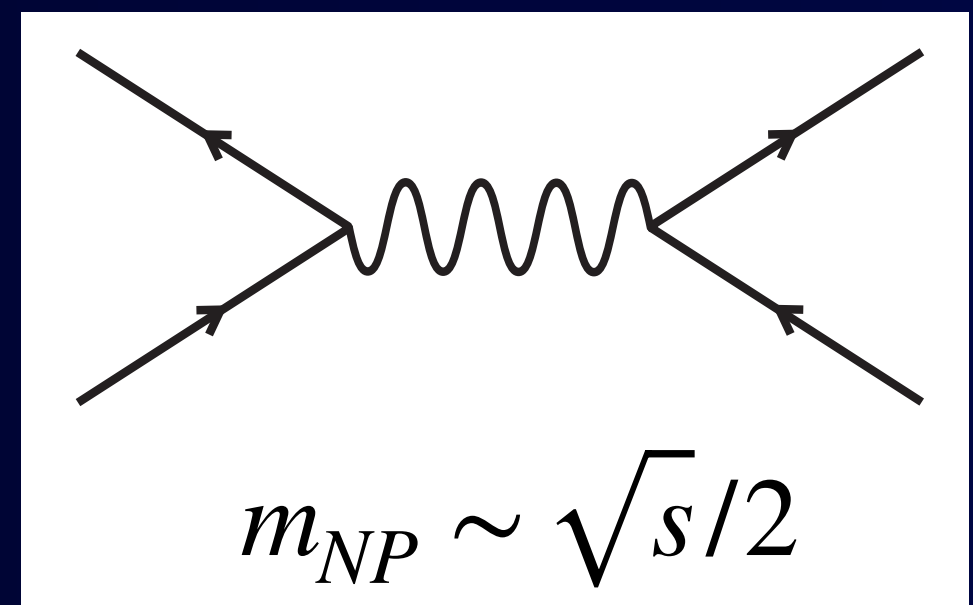
Drives the beam quality similar to MAP design still challenging design with challenging components



Muon collider physics

$$\sqrt{s} \sim 10+ \text{ TeV}$$

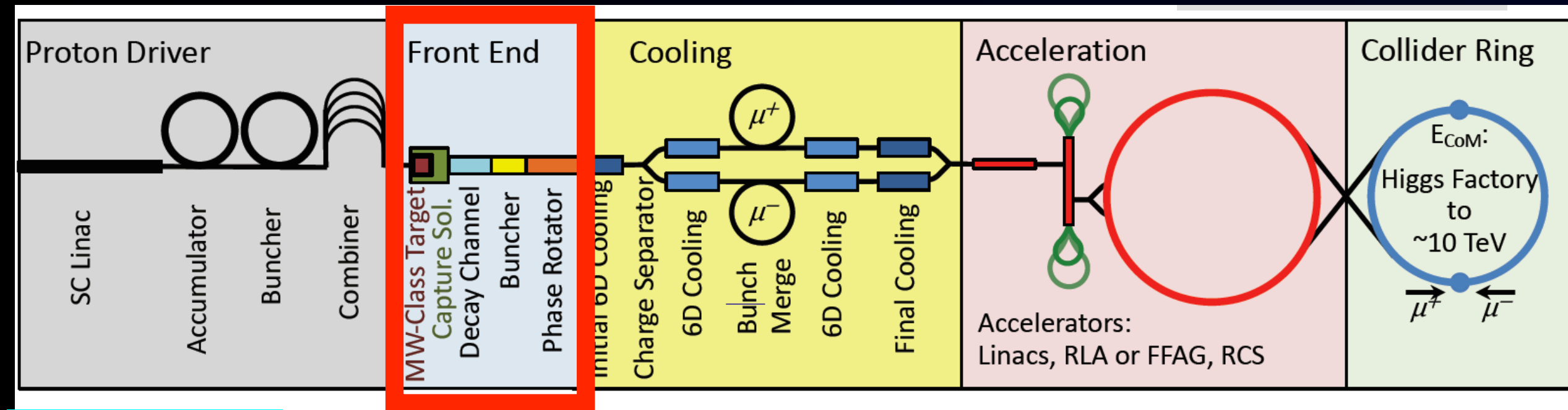
- Discovery
 - direct or indirect BSM production
- Precision
 - Higgs production with vector boson fusion
- μ Tristan at KEK
 - $\mu^+\mu^+$ or μ^+e^- @ 1 TeV



Muon Colliders

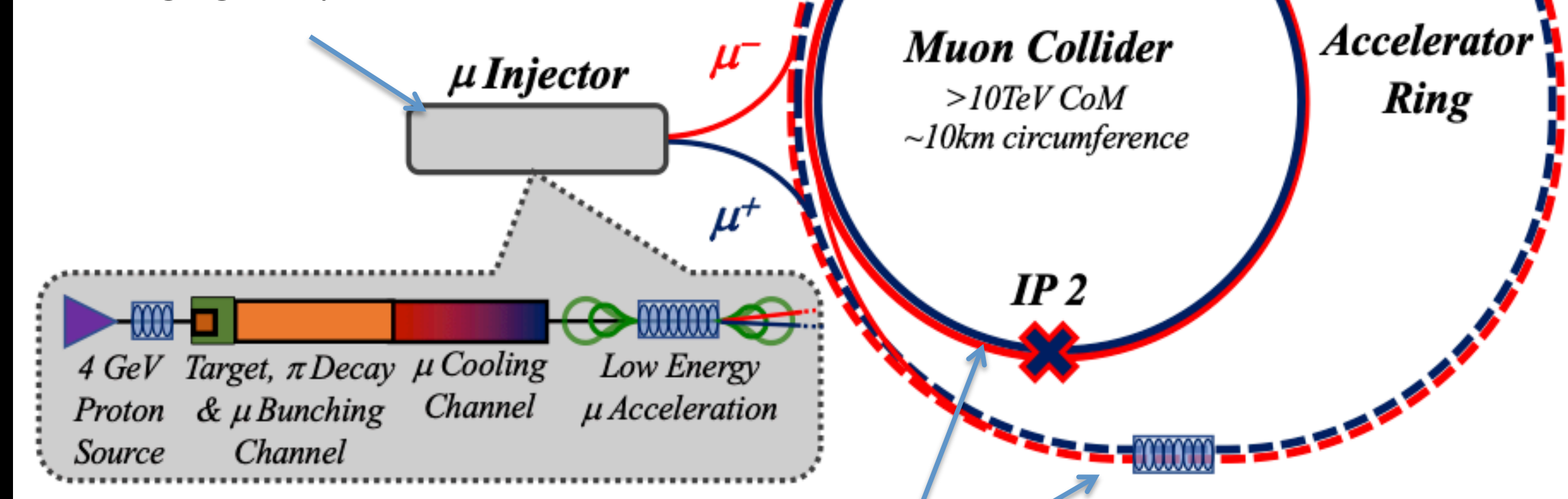


MAP (US)



CERN

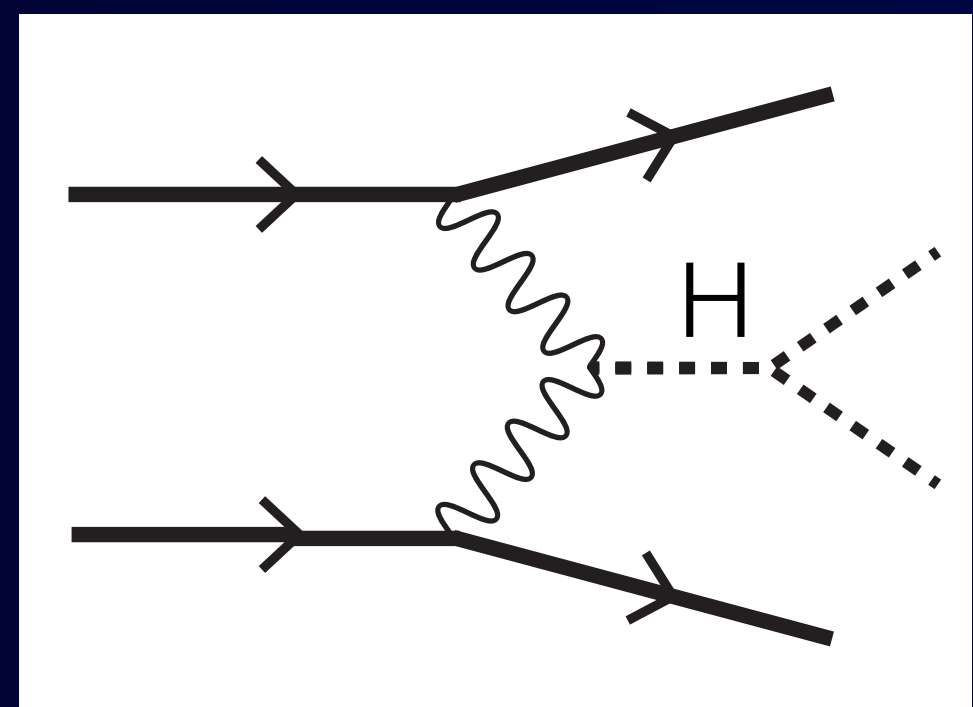
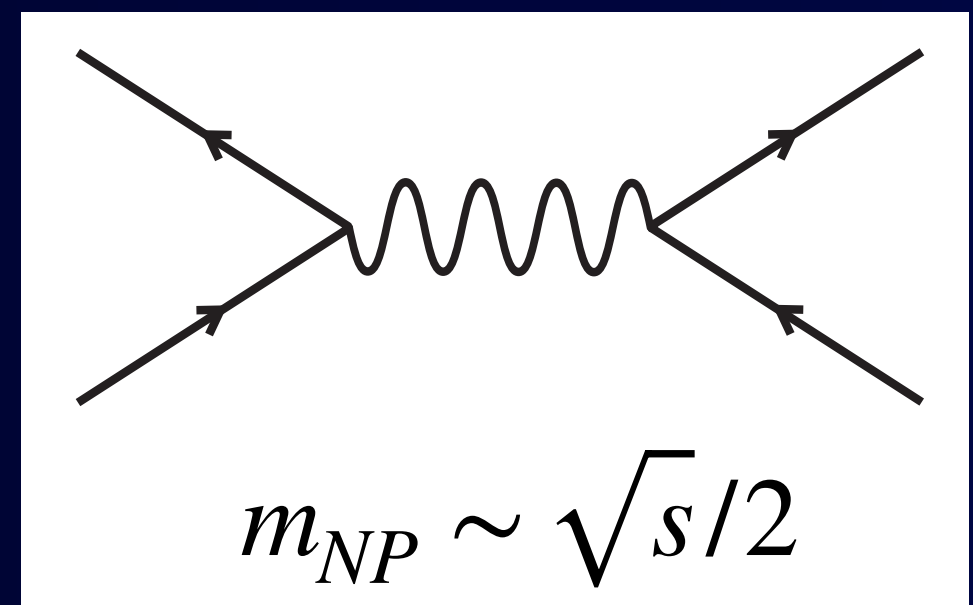
Drives the beam quality similar to MAP design still challenging design with challenging components



Muon collider physics

$$\sqrt{s} \sim 10+ \text{TeV}$$

- Discovery
 - direct or indirect BSM production
- Precision
 - Higgs production with vector boson fusion
- μ Tristan at KEK
 - $\mu^+\mu^+$ or μ^+e^- @ 1 TeV



Highly intense muon source

phase-rotation, muon cooling, etc.

Some More with Muons



Some More with Muons

Some of muon particle physics topics which were not covered are :

- Normal muon decay
 - precise measurements of Michel parameters
- CLFV
 - $\mu \rightarrow e\gamma\gamma$
 - $\mu^-e^- \rightarrow e^-e^-$ in a muonic atom
 - $\mu N \rightarrow \tau N$ DIS scattering
- Muonium
 - Muonium hyperfine splitting
 - Muonium 1s-2s spectroscopy

any missing !



Summary



Summary



Summary



- Muon particle physics could provide a unique discovery potential for physics beyond the Standard Model (BSM).
- Development of a highly intense muon source is an important key success factor for muon particle physics programs.

Summary

- Muon particle physics could provide a unique discovery potential for physics beyond the Standard Model (BSM).
- Development of a highly intense muon source is an important key success factor for muon particle physics programs.

Thank you for your attention!



Backup Slides