

The nuts and bolts of muon precession frequency measurements in the Fermilab Muon g-2 experiment

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Muon g-2 collaboration @ Liverpool

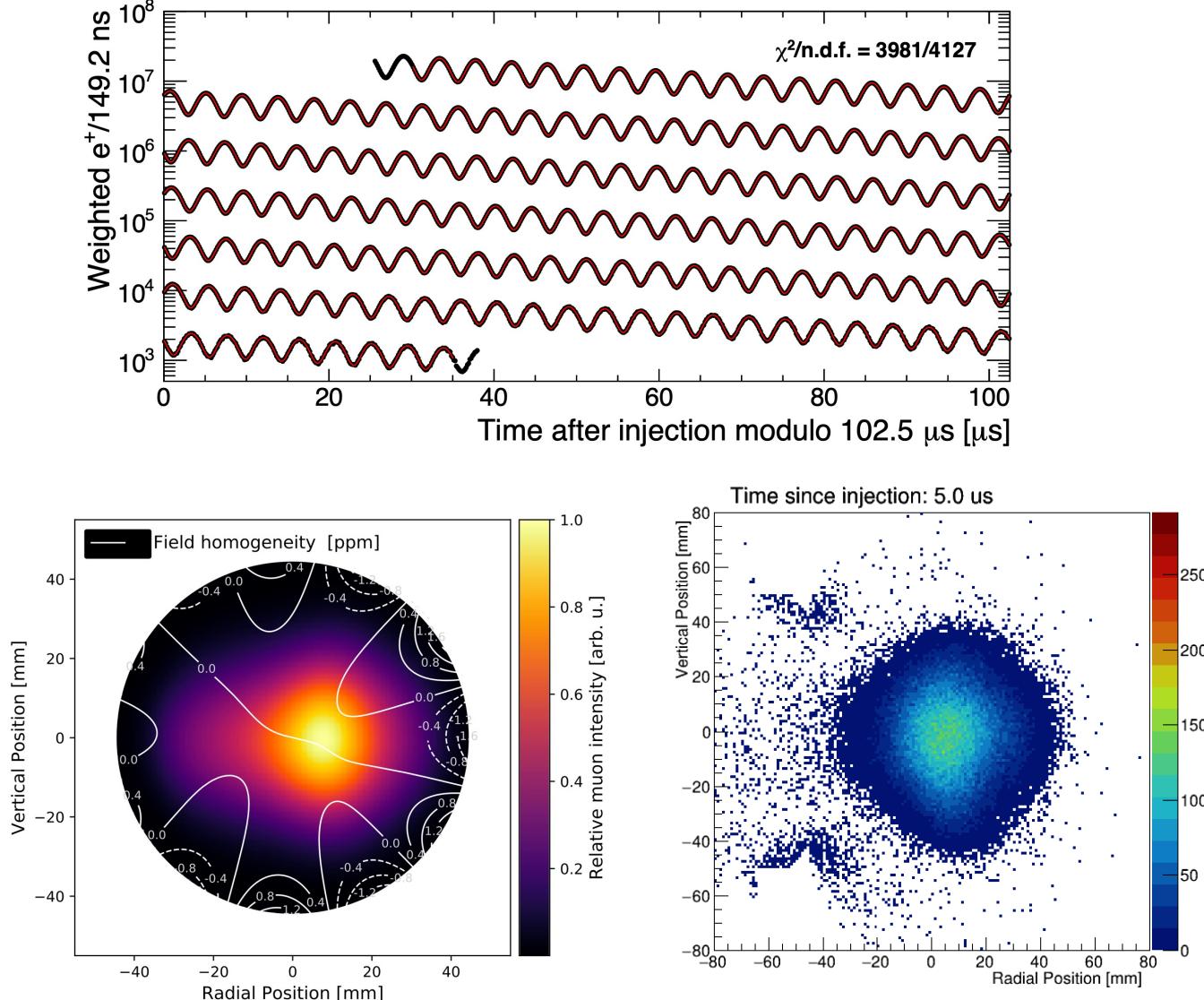


Muon g-2 Experiment in a Nutshell



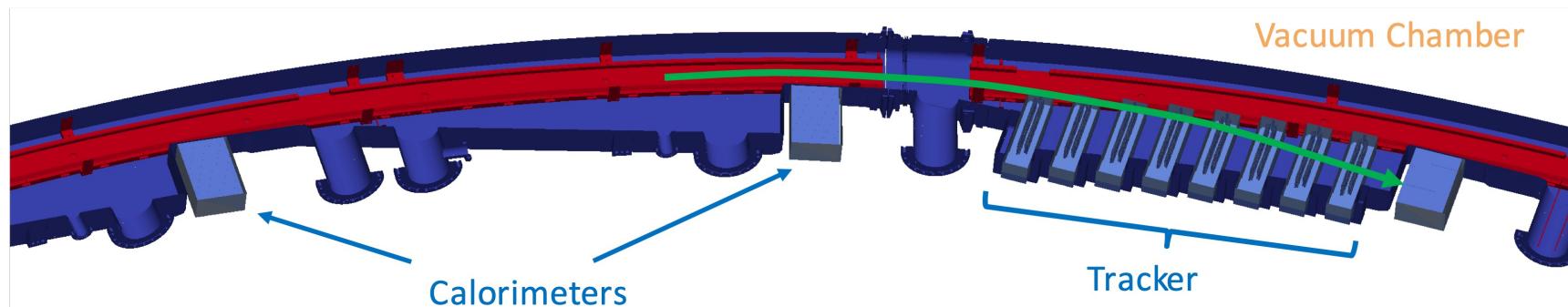
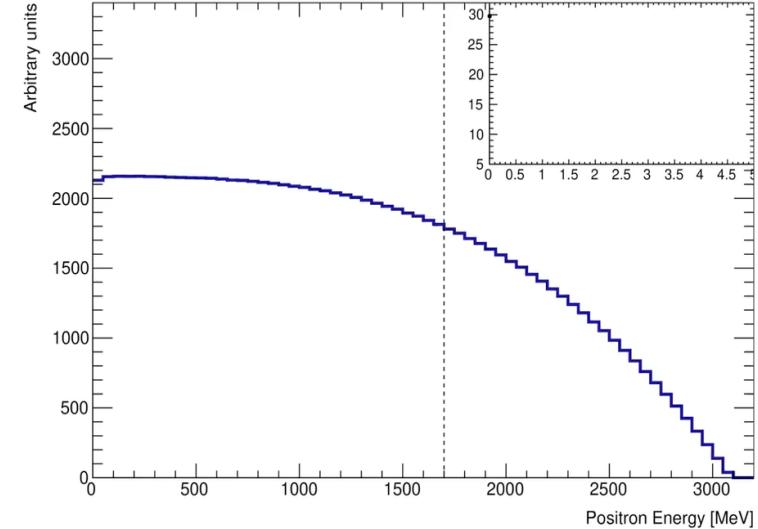
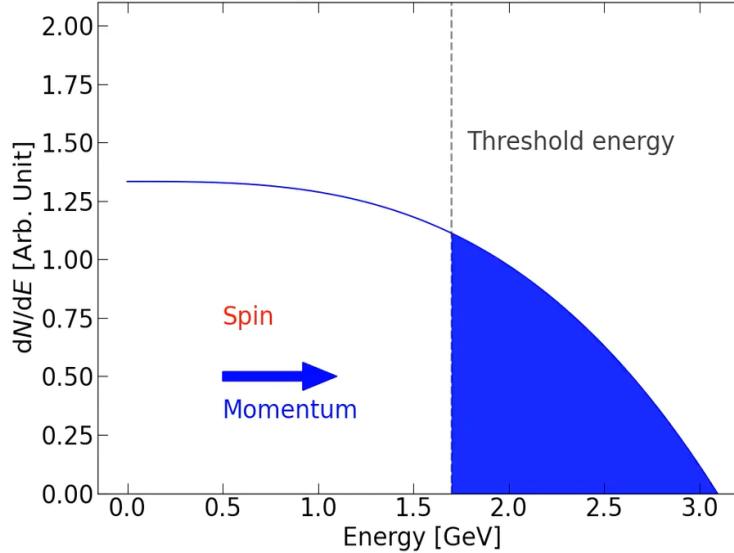
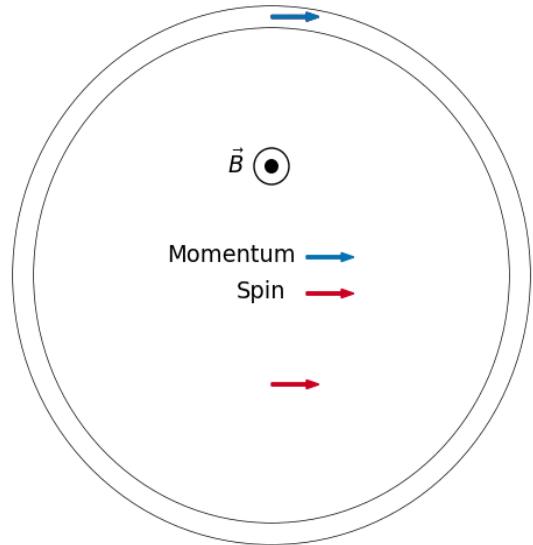
$$\frac{\omega_a}{\omega_p} = \frac{\omega_a^m}{\omega_p^m}$$

Measured Values



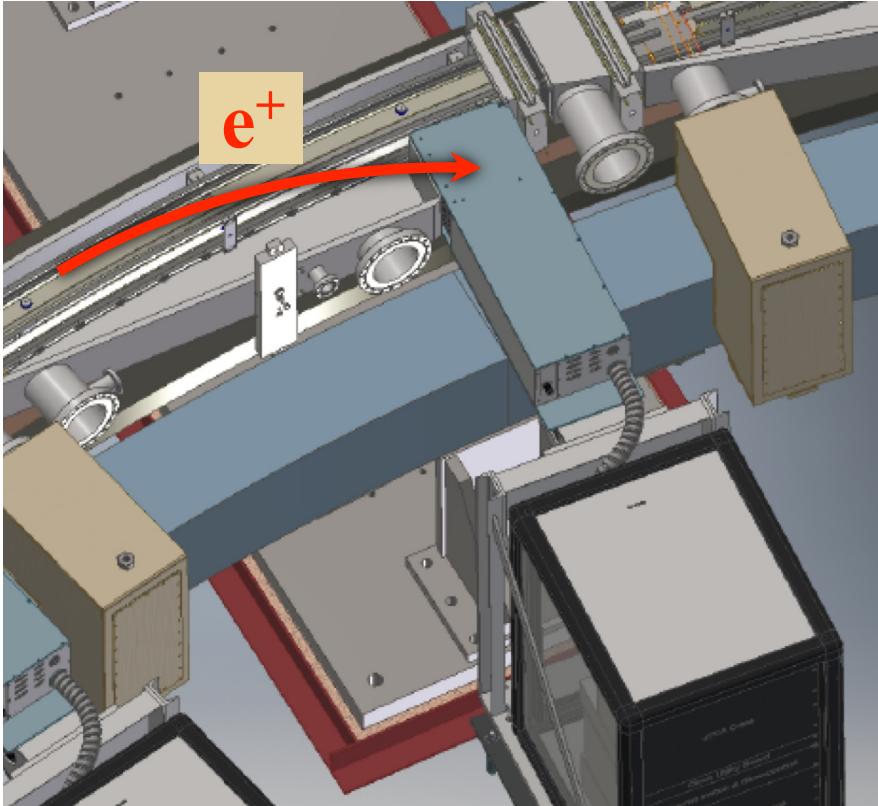
Muon precession frequency

$$N(t) = N_0 e^{-t/\tau} [1 + A_\mu \cos(\omega_a t + \phi)]$$



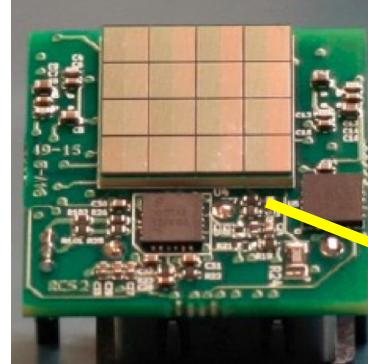
The SWISS KNIFE for g-2 experiment
3

Calo measure positron time and energy

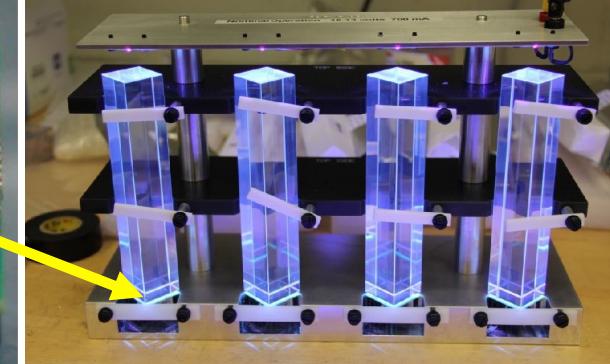


Decay positron curving in and
striking a calorimeter

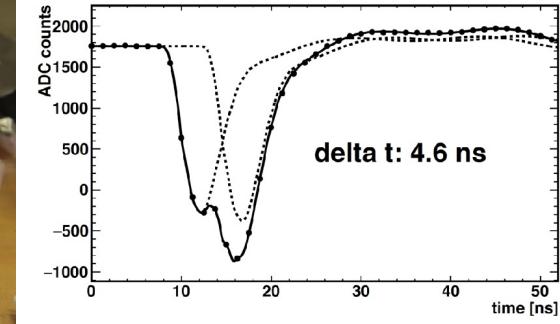
SiPM



PbF₂



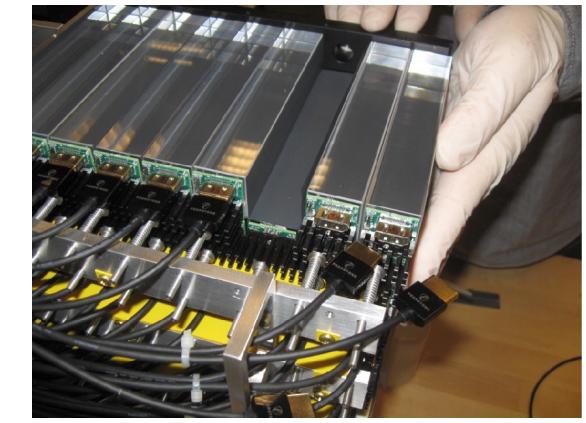
pileup separation



PMT-like signal, B-field operation, 100% separation > 2.5 ns

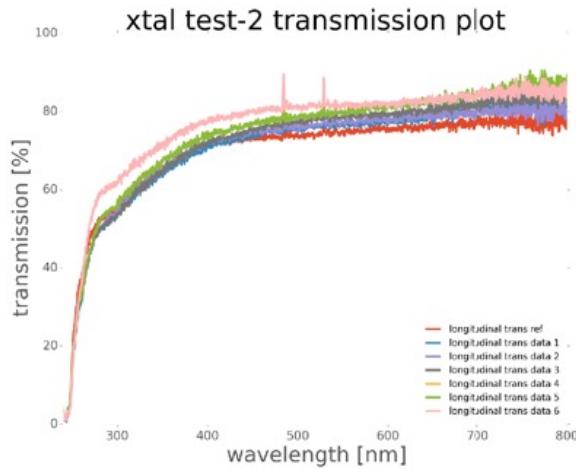


Opened up calorimeter

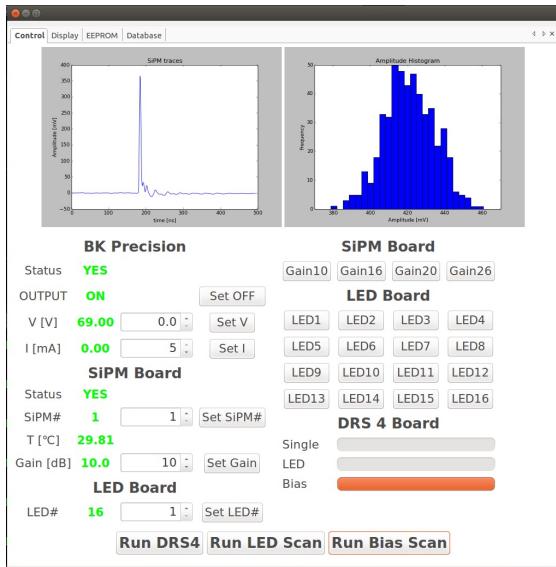
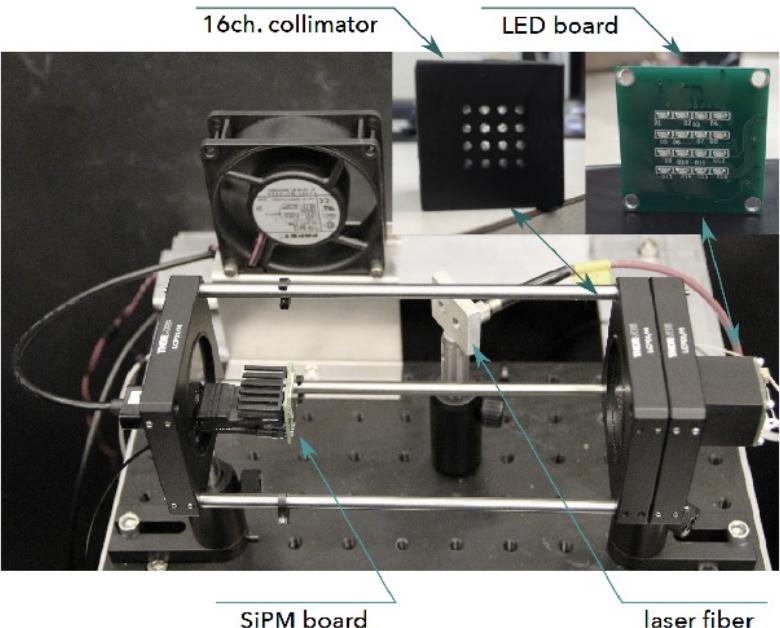


Stacking crystals

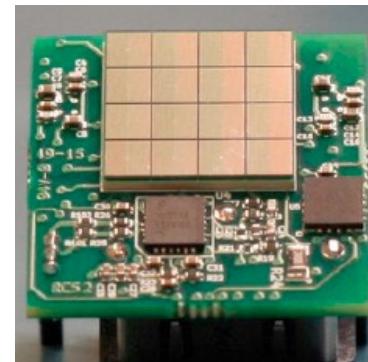
Calorimeter: Quality control systems



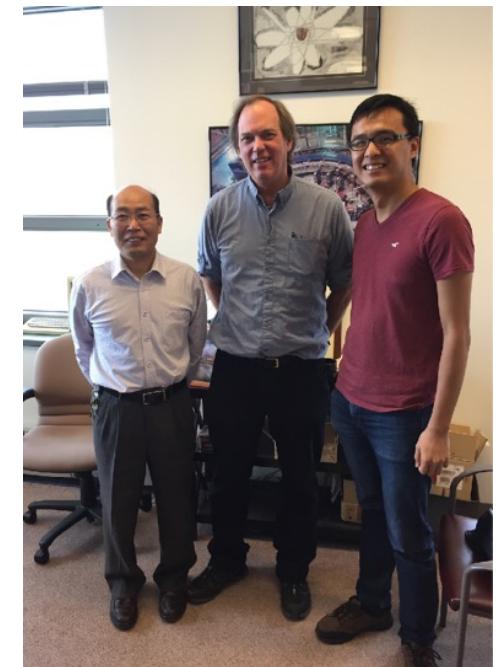
Inspection procedure and transmission measurement



Good transmission and in-range dimension needed! 1350 of them

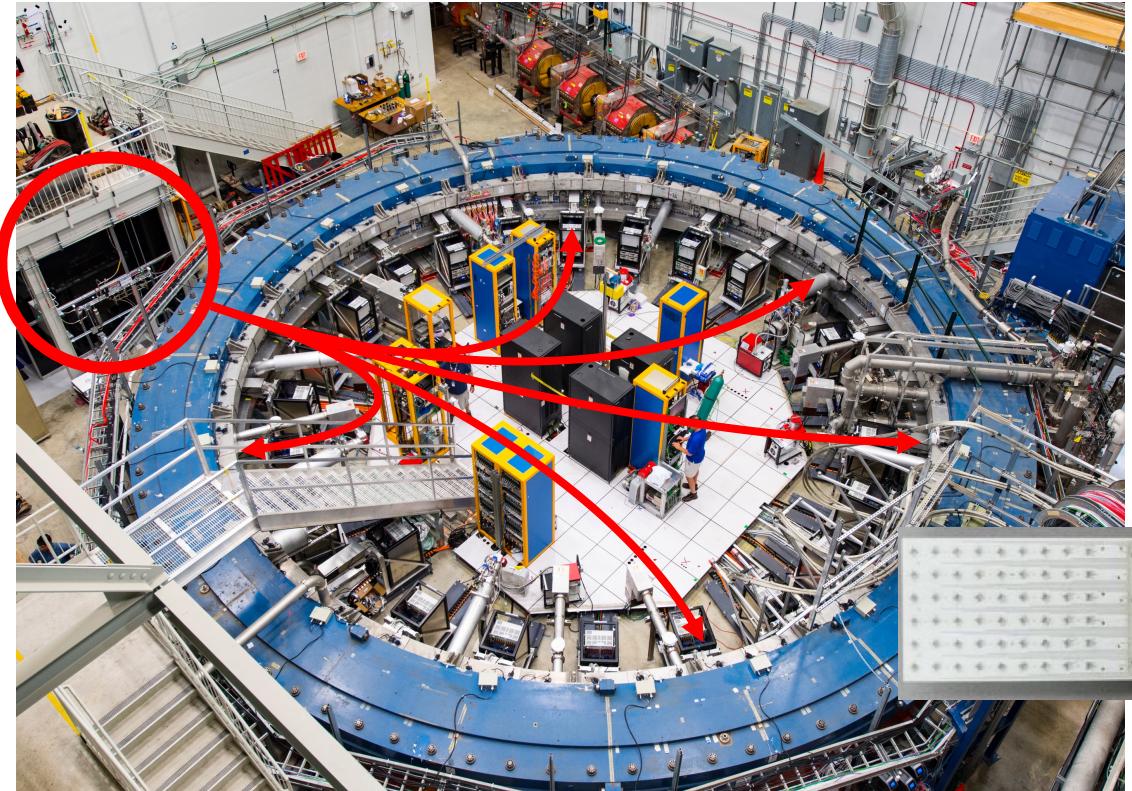


Guohao Ren
Head of Inorganic Scintillator
(SICCAS, 上海硅酸盐研究所)

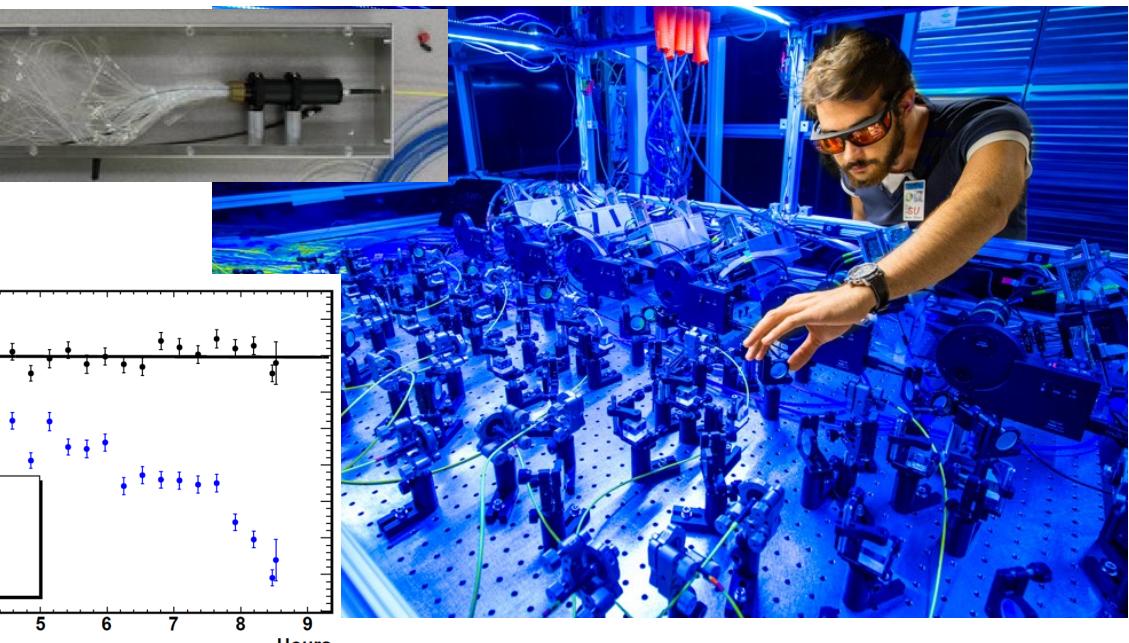


All 16 CHs, on-chip T-sensor, memory have to work! 1400 of them 5

Gain monitoring and calibration system



- Laser light @405 nm sent to all 1296 crystals with distribution panels
- Time synchronization at the ~50 ps
- Gain calibration of the SiPMs at the 10^{-4} level (< 20 ppb systematics in gain)



Real World Complications: Corrections



- We need to make corrections for seven small effects:

$$\frac{\omega_a}{\omega_p} = \frac{\omega_a^m}{\omega_p^m} \frac{1 + \overbrace{C_e + C_p}^{\text{E-field \& vertical motion: Spin precesses slower than in basic equation}} + \overbrace{C_{pa} + C_{dd} + C_{ml}}^{\text{Phase changes over each fill: Phase-Acceptance, Differential Decay, Muon Losses}}}{1 + \underbrace{B_k + B_q}_{\text{Transient Magnetic Fields: Quad Vibrations, Kicker Eddy Current,}}}$$

Measured Values

- ω_a corrections are small, but dominated Run-1 systematics: **500(93) ppb**
- For Run-2/3 they are **580(40) ppb**, dominated by C_e and C_p

Spin precession in B and E-field

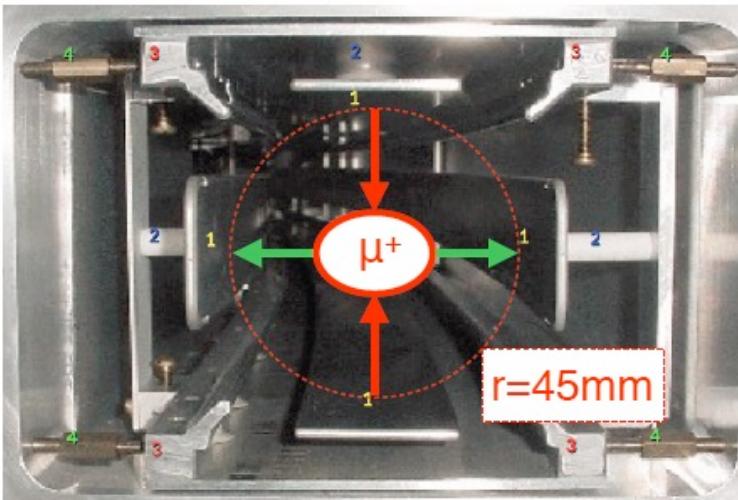
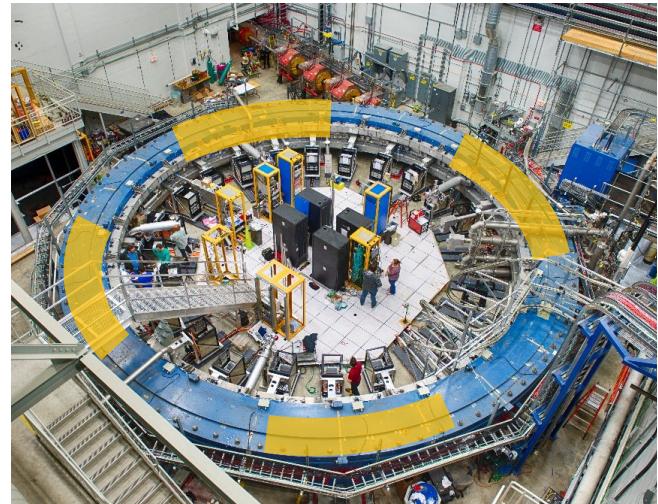
- Generic spin-motion is described by BMT equation:

$$\frac{d(\hat{\beta} \cdot \vec{S})}{dt} = -\frac{q}{m} \vec{S}_T \cdot \left[a_\mu \hat{\beta} \times \vec{B} + \beta \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{E}}{c} \right]$$

simplified equation

$$\omega_a = a_\mu \frac{eB}{m}$$

- Muons travel in **E-field** from focusing quadrupoles: experience a **motional magnetic field** in their rest frame



- The electrostatic quadrupoles (ESQ) are used to focus the beam vertically.
- Four sections cover 43% of the circumference.
- Each plate is charged to ± 18 kV.

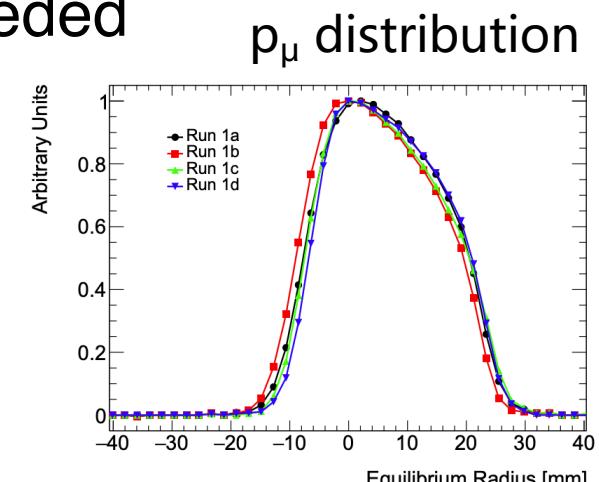
Finite width of the p-distribution

- Generic spin-motion is described by BMT equation:

$$\frac{d(\hat{\beta} \cdot \vec{S})}{dt} = -\frac{q}{m} \vec{S}_T \cdot \left[a_\mu \hat{\beta} \times \vec{B} + \beta \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{E}}{c} \right]$$

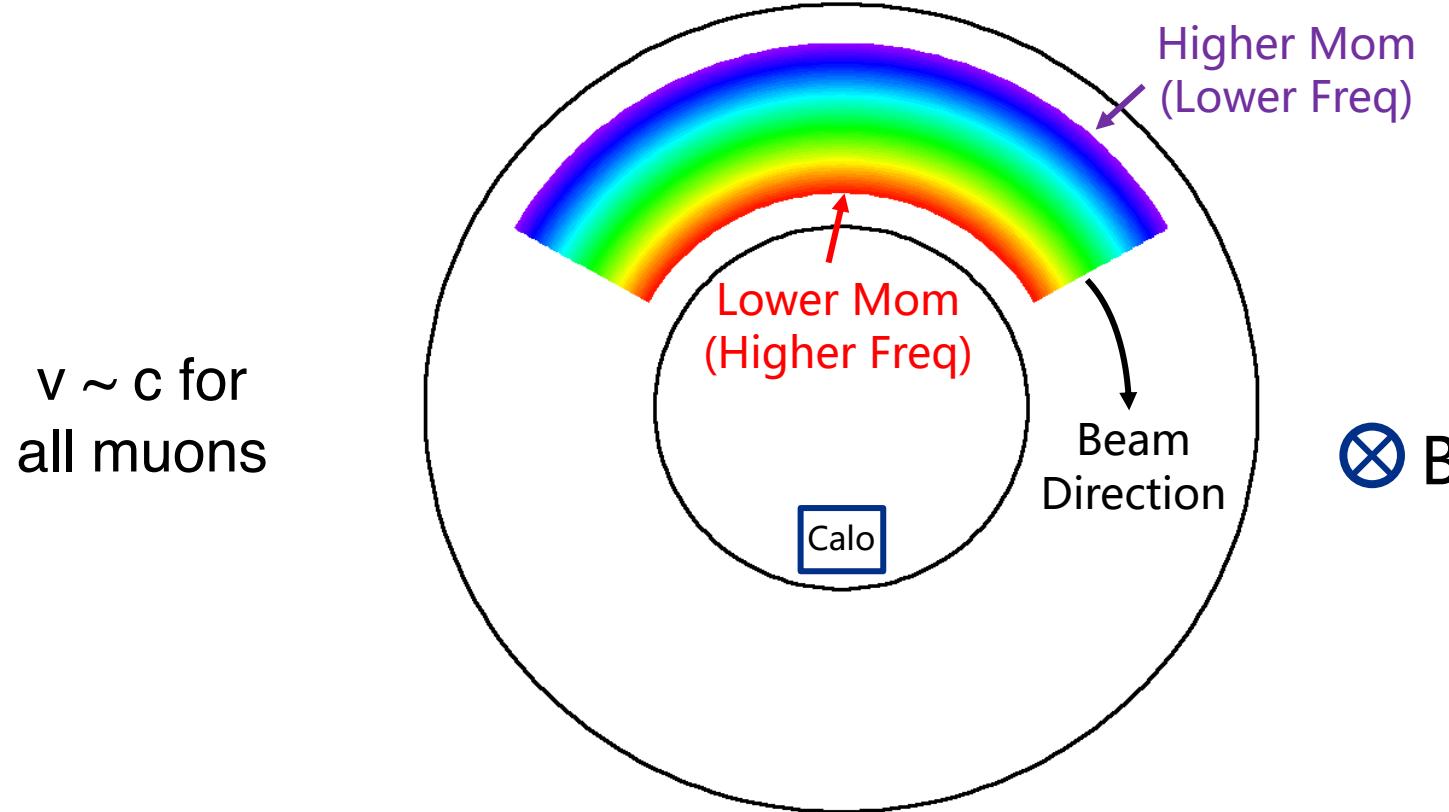
- Muons travel in **E-field** from focusing quadrupoles: experience a **motional magnetic field** in their rest frame
- Term vanishes at "magic" momentum ($p_\mu = 3.094 \text{ GeV}$)
- But not all muons are at p_{magic} , hence corrections are needed
- C_E comes from p_μ distribution measured using timing data from calorimeters

$$C_E = 451 \pm 32 \text{ ppb}$$



Fast rotation signal

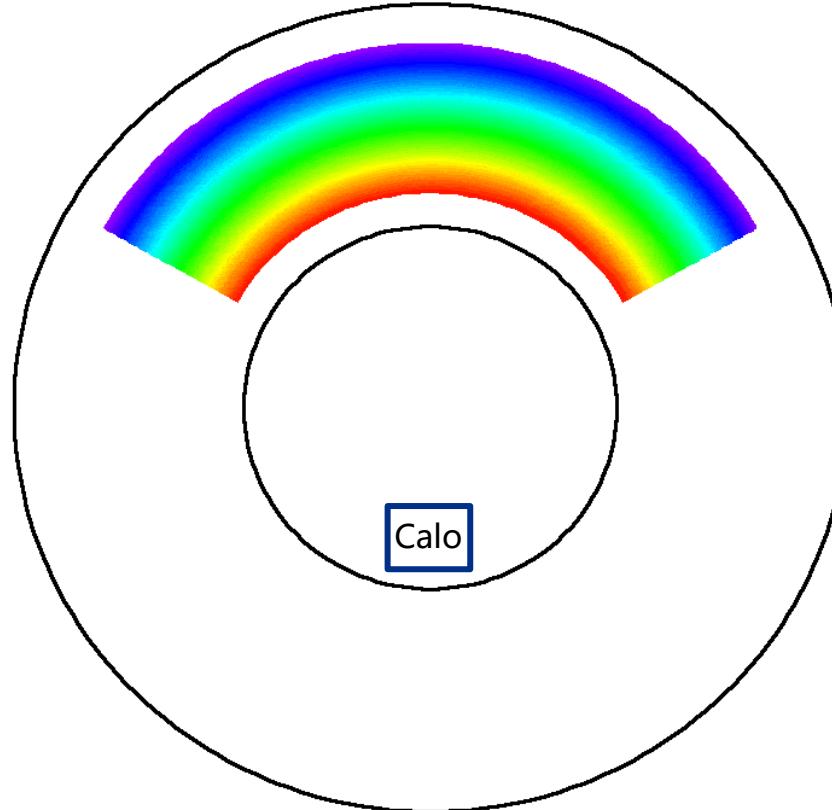
- Imagine injecting uniform momentum & time distributions:



- Higher momentum muons have further to travel, so have lower cyclotron frequencies

Fast rotation signal

- Over time, lower momentum muons will catch up with higher momentum muons:

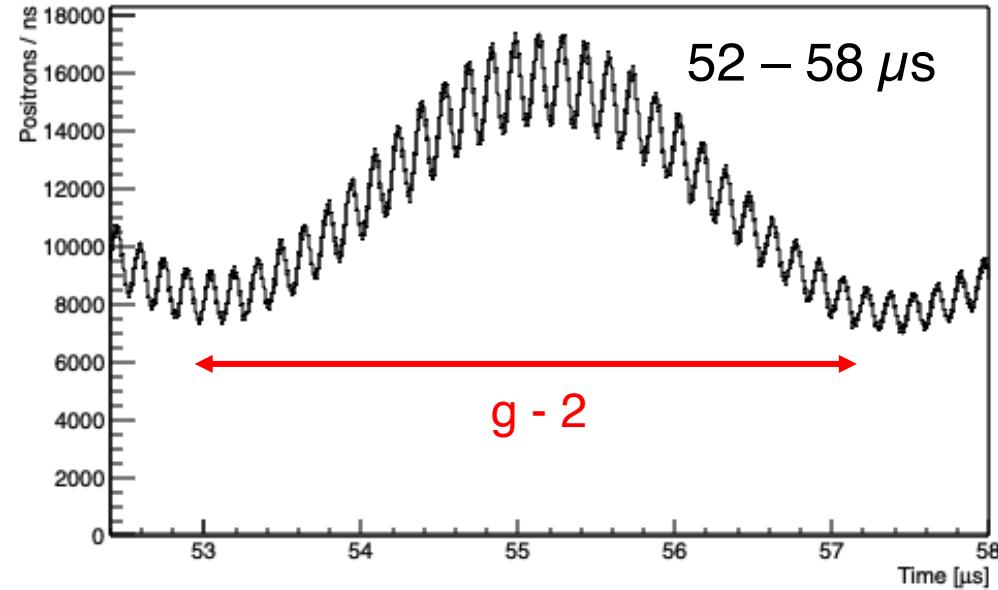
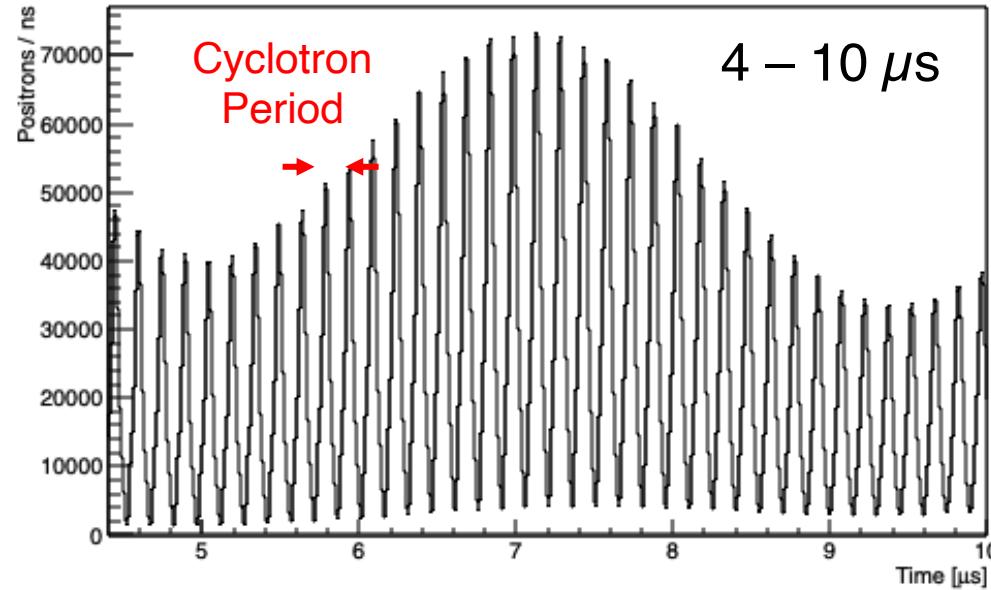


- The way that the gaps are filled in is related to the momentum distribution of the stored beam

Fast rotation signal in the calorimeter



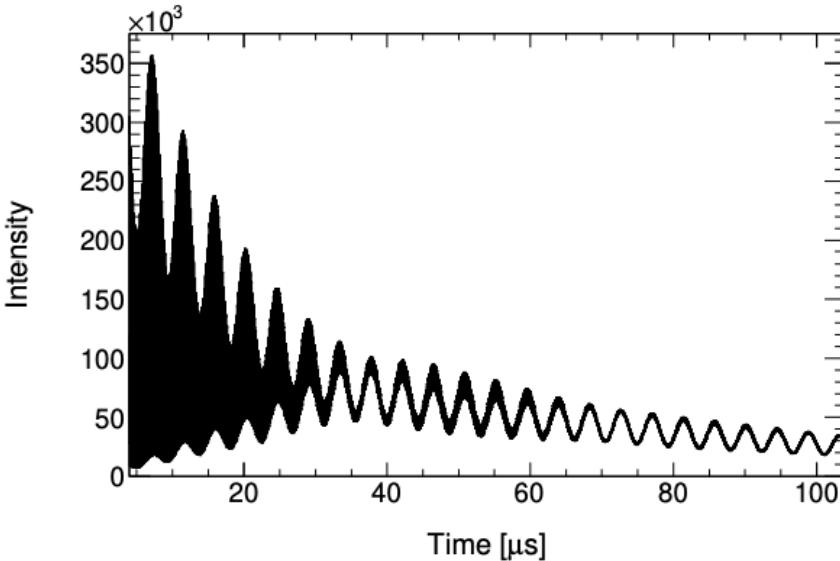
- Effect is a strong feature of the calorimeter timing data at early times:



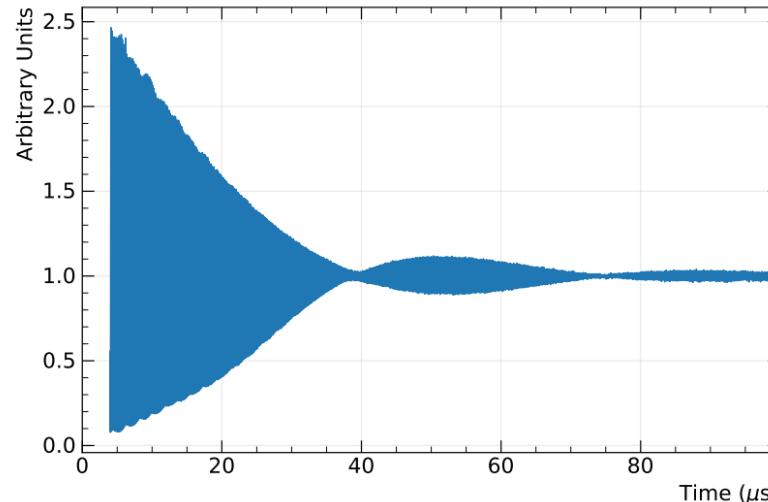
- Less pronounced when all calos are added together
- We utilized this feature to measure the muon momentum distribution or equilibrium position

E-field Correction: C_e

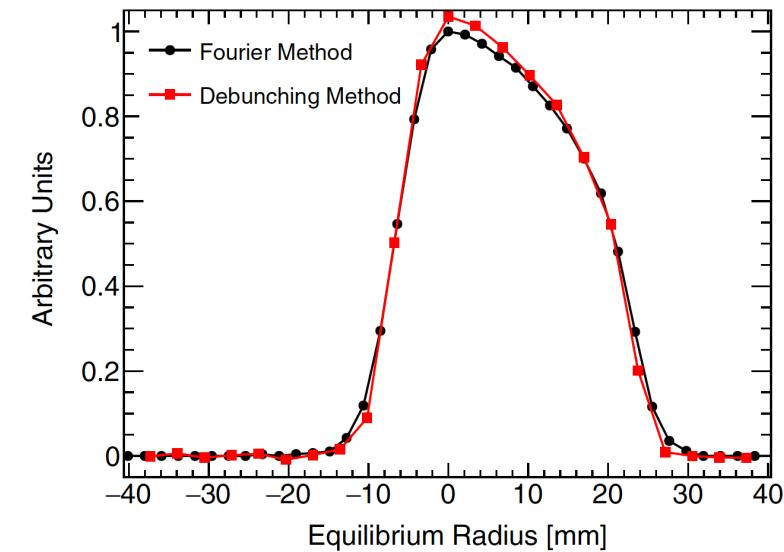
Finely-binned decay e^+ spectrum



Isolated FR signal



Reconstructed x_e distribution



Equilibrium position reconstruction approaches:

1. Fourier method (frequency domain)
2. Debunching method (time domain)

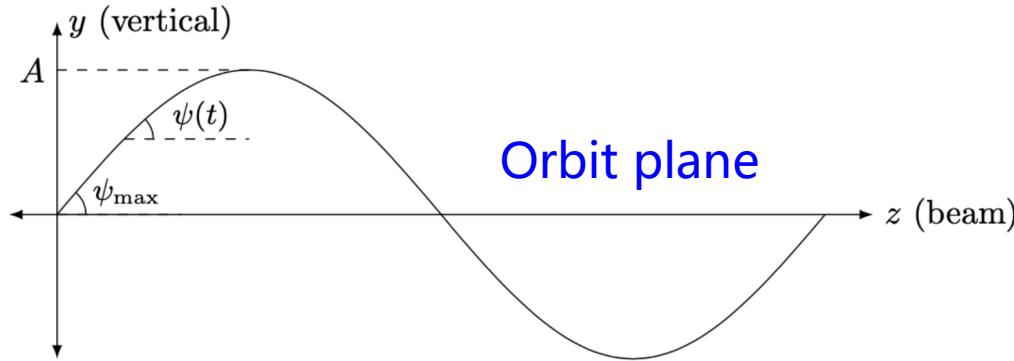
$$C_e \approx 2n(1-n)\beta_0^2 \frac{\langle x_e^2 \rangle}{R_0^2}$$

Beam vertical oscillation

- Generic spin-motion is described by BMT equation:

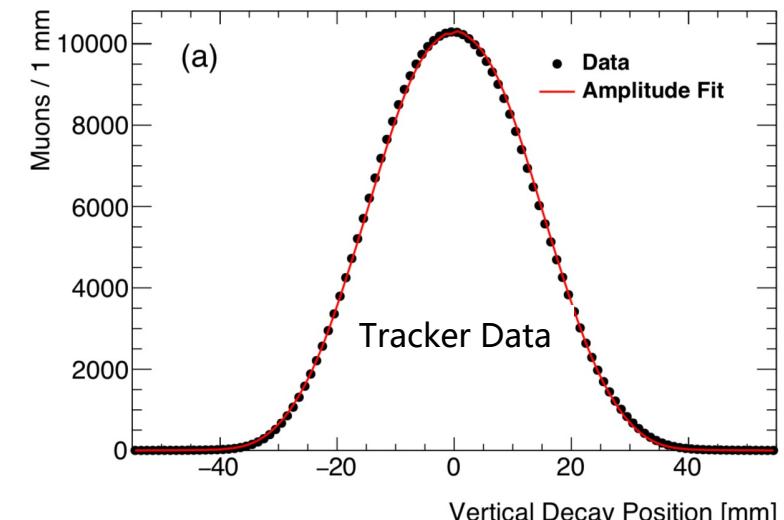
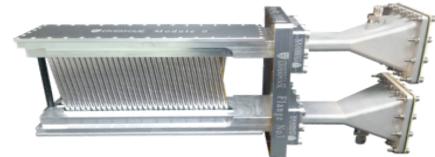
$$\frac{d(\hat{\beta} \cdot \vec{S})}{dt} = -\frac{q}{m} \vec{S}_T \cdot \left[a_\mu \hat{\beta} \times \vec{B} + \beta \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{E}}{c} \right]$$

- Muons oscillate vertically (**pitch**) so the $\hat{\beta} \times \vec{B}$ term is reduced



- C_P is extracted from vertical width measured by the trackers

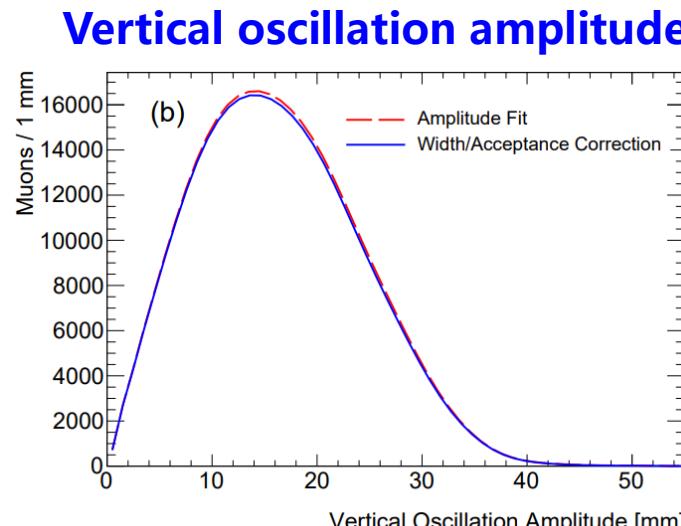
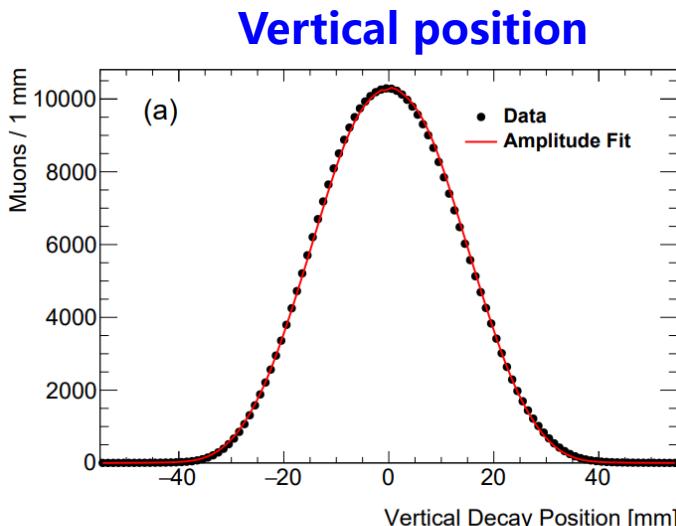
$$C_P = 170 \pm 10 \text{ ppb}$$



Pitch correction: C_p



- C_p is calculated from the amplitude distribution of the vertical motion
- The amplitude is reconstructed from the position distribution, including the acceptance correction.
- Dominant systematic error source: tracker alignment and reconstruction.
- Improvement after Run-1: Independent & different method analysis cross-check.



$$C_p = \frac{n}{2} \frac{\langle y^2 \rangle}{R_0^2} = \frac{n}{4} \frac{\langle A^2 \rangle}{R_0^2}$$

C_p	Correction [ppb]	Uncertainty [ppb]
Run-1	180	13
Run-2/3	170	10

Real World Complications: Corrections



- We need to make corrections for seven small effects:

$$\frac{\omega_a}{\omega_p} = \frac{\omega_a^m}{\omega_p^m} \frac{1 + \overbrace{C_e + C_p}^{\text{E-field \& vertical motion: Spin precesses slower than in basic equation}} + \overbrace{C_{pa} + C_{dd} + C_{ml}}^{\text{Phase changes over each fill: Phase-Acceptance, Differential Decay, Muon Losses}}}{1 + \overbrace{B_k + B_q}^{\text{Transient Magnetic Fields: Quad Vibrations, Kicker Eddy Current,}}}$$

Measured Values

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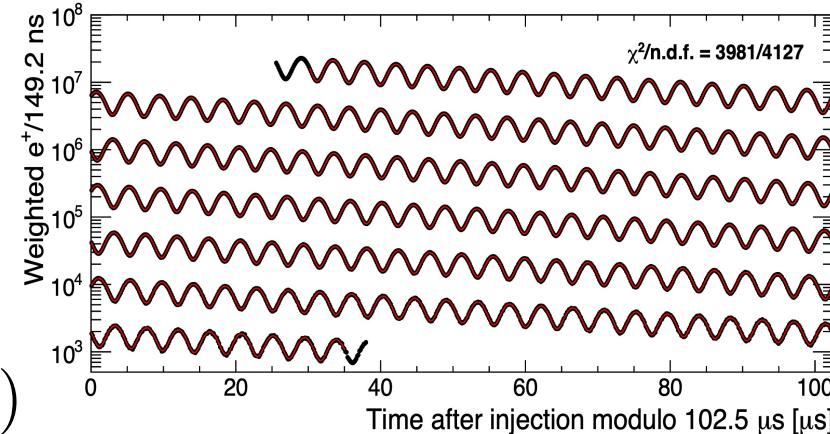
ω_a systematics due to phase shifts



- If there are effects that **change the g-2 phase** of the detected e^+ over time
- These make us mis-measure ω_a with no indications that we're getting it wrong

$$\begin{aligned} \cos(\omega_a t + \phi(t)) &= \cos(\omega_a t + \phi_0 + \phi' t + \dots) \\ &= \cos((\omega_a + \phi')t + \phi_0 + \dots) \end{aligned}$$

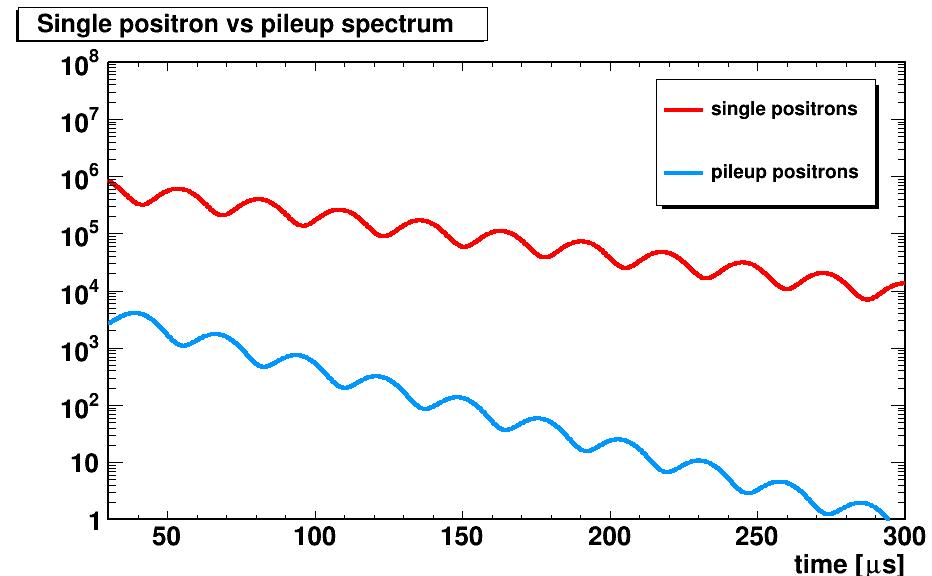
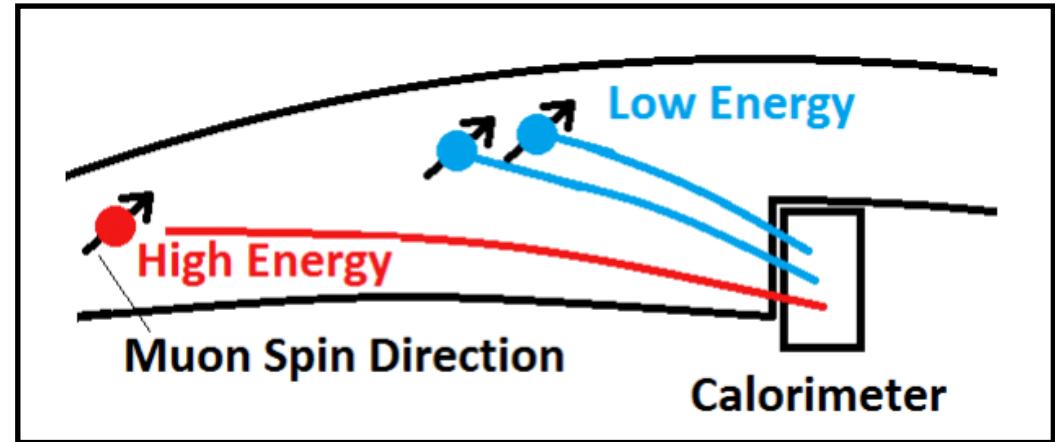
$$N(t) = N_0 e^{-t/\tau} \left[1 + A_\mu \cos(\omega_a t + \phi) \right]$$



- In general, anything that changes from **early-to-late** within each muon fill can be a cause of systematic error.
- Most phase shifts are eliminated by design or before fitting the data

ω_a systematics

- **Detector Effects**
 - Positron event pileups
 - Gain instability
- **Beam Dynamics**
 - Horizontal betatron motion
 - Vertical betatron motion
 - Beam de-bunching



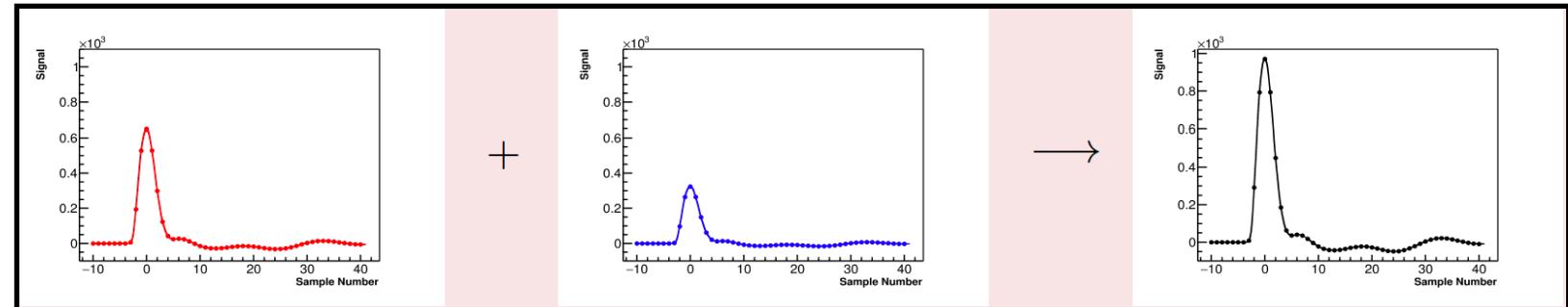
$$A_1 \cos(\omega t + \phi_1) + A_2 \cos(\omega t + \phi_2) = A_3 \cos(\omega t + \phi_3)$$

ω_a systematics

- **Detector Effects**

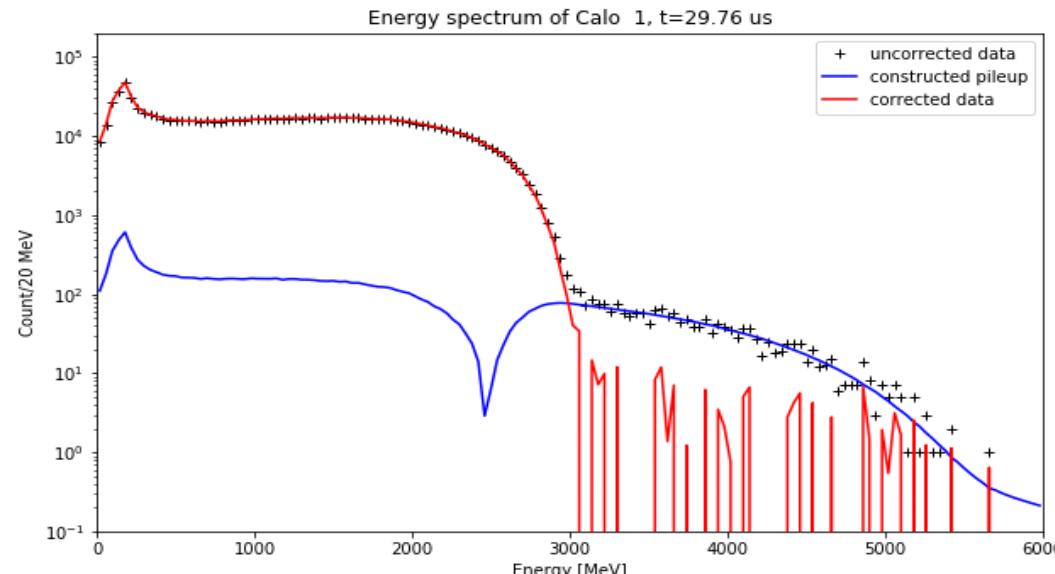
- Positron event pileups
- Gain instability

Artificially constructing pileup spectrum and removing it from the raw data



- **Beam Dynamics**

- Horizontal betatron motion
- Vertical betatron motion
- Beam de-bunching



ω_a systematics

- **Detector Effects**

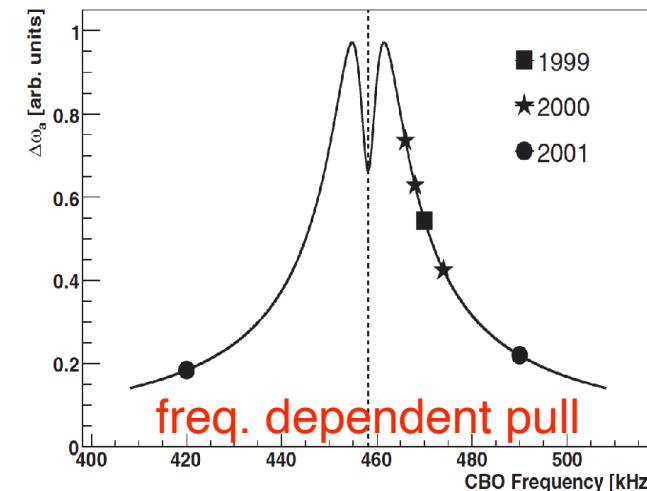
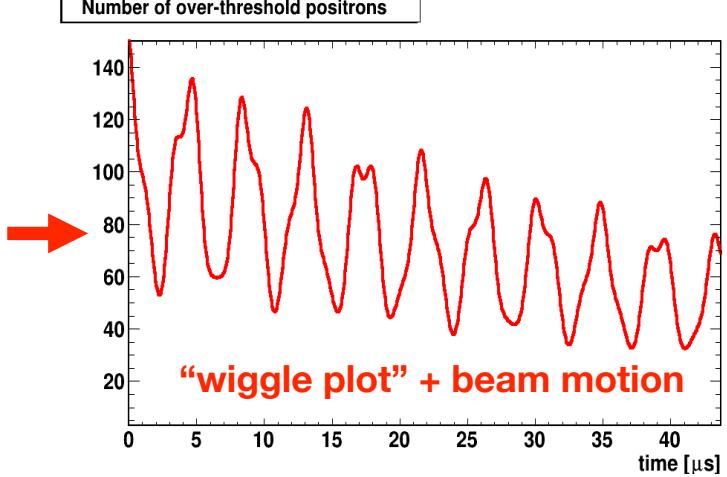
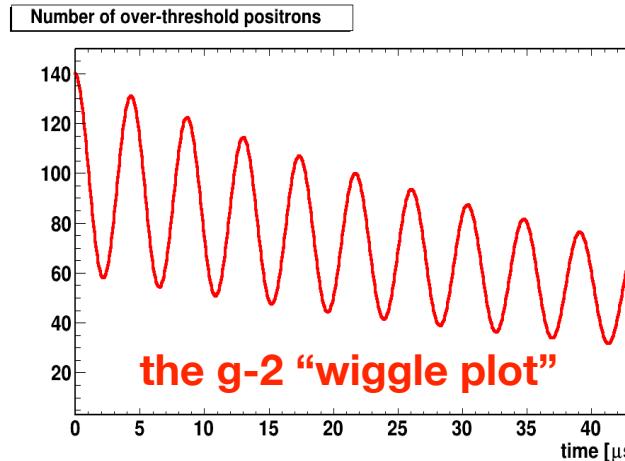
- Positron event pileups
- Gain instability

- **Beam Dynamics**

- Horizontal betatron motion
- Vertical betatron motion
- Beam de-bunching



Detector acceptance
effect



$$C(t) = 1 - e^{-t/\tau_{cbo}} A_1 \cos(\omega_{cbo} t + \phi_1)$$

Summary of ω_a systematics for Run 2/3

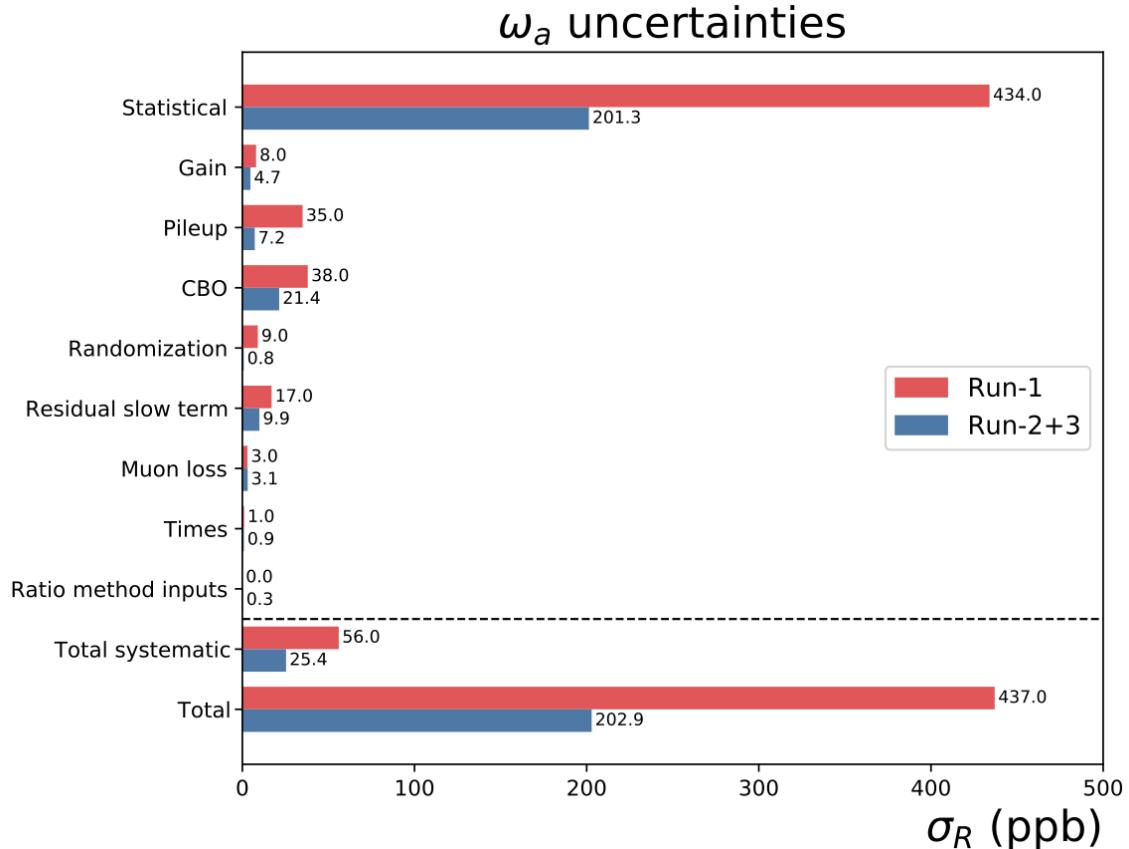
- With respect to Run-1:
 - Statistical uncertainty reduced by **2.2x**
 - Systematic uncertainty reduced by **2.2x**
- Statistical uncertainty still completely dominates the final Run-2/3 error

TDR (2015)		Run-2/3
Category	Syst goal [ppb]	Actual [ppb]
Gain	20	5
Pileup	40	7
CBO	<30	21
Lost muons	20	3
Other	-	10
Total	57	25

✓



KSK – Task force leader for reducing CBO systematics in Run 4/5/6



ω_a systematics due to phase shifts



- If there are effects that **change the g-2 phase** of the detected e^+ over time
- These make us mis-measure ω_a with no indications that we're getting it wrong

$$\begin{aligned} \cos(\omega_a t + \phi(t)) &= \cos(\omega_a t + \phi_0 + \phi' t + \dots) \\ &= \cos((\omega_a + \phi')t + \phi_0 + \dots) \end{aligned}$$

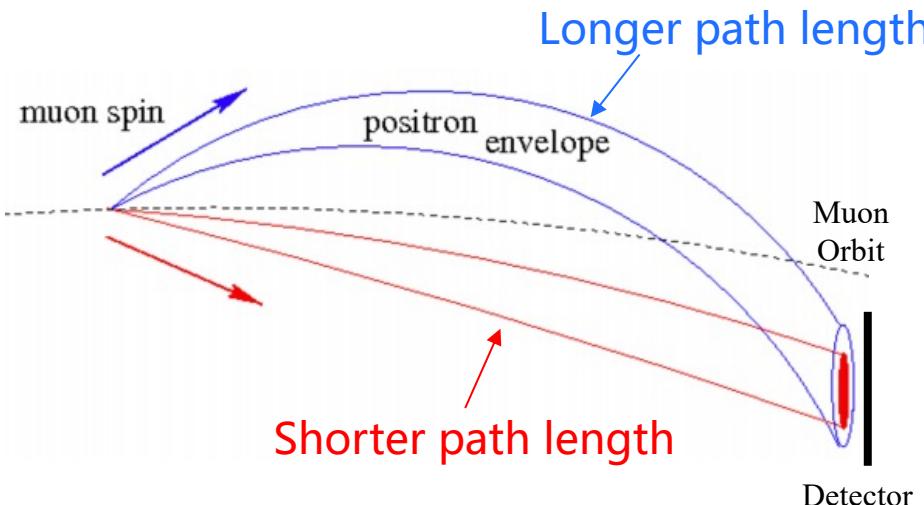
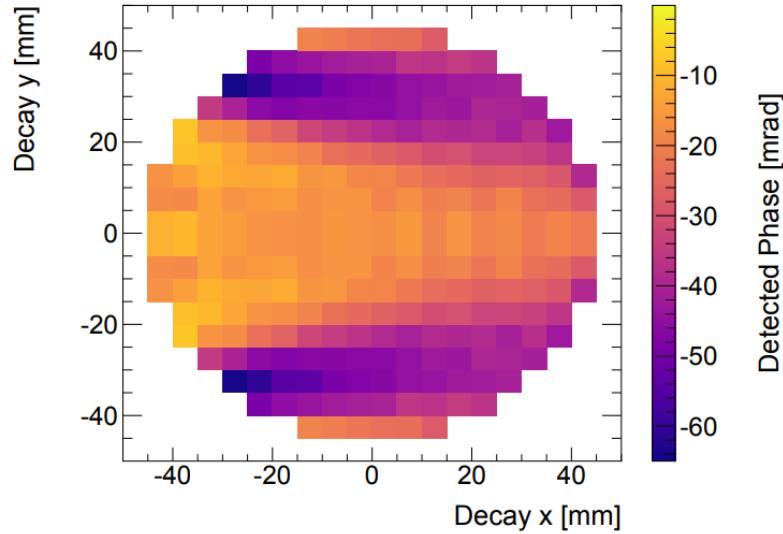
- In general, anything that changes from **early-to-late** within each muon fill can be a cause of systematic error.
- Most phase shifts are eliminated by design or before fitting the data, but we must correct for three effects (C_{pa} , C_{ml} , C_{dd})

$$\frac{\omega_a}{\omega_p} = \frac{\omega_p^m}{\omega_a^m} \frac{1 + \underbrace{C_e + C_p}_{\substack{\text{Measured Values} \\ \text{E-field \& vertical motion:} \\ \text{Spin precesses slower} \\ \text{than in basic equation}}}}{1 + \underbrace{B_k + B_q}_{\substack{\text{Transient Magnetic Fields:} \\ \text{Quad Vibrations,} \\ \text{Kicker Eddy Current,}}} + \underbrace{C_{pa} + C_{dd} + C_{ml}}_{\substack{\text{Phase changes over each fill:} \\ \text{Phase-Acceptance, Differential} \\ \text{Decay, Muon Losses}}}}$$

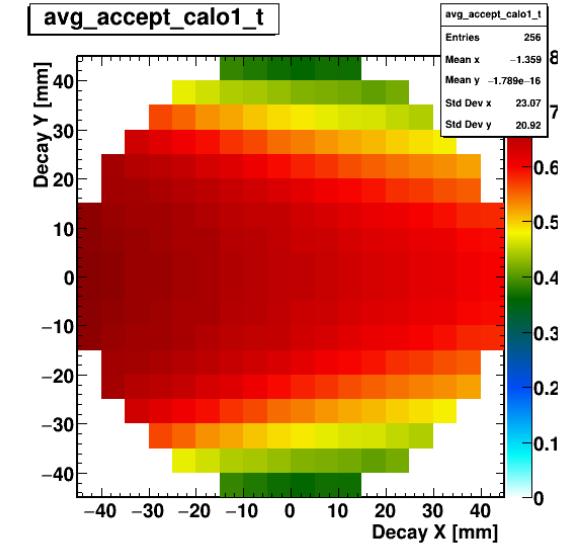
g-2 phase and detector acceptance

- Average detected g-2 phase changes with decay position:

Phase map (Geant4)



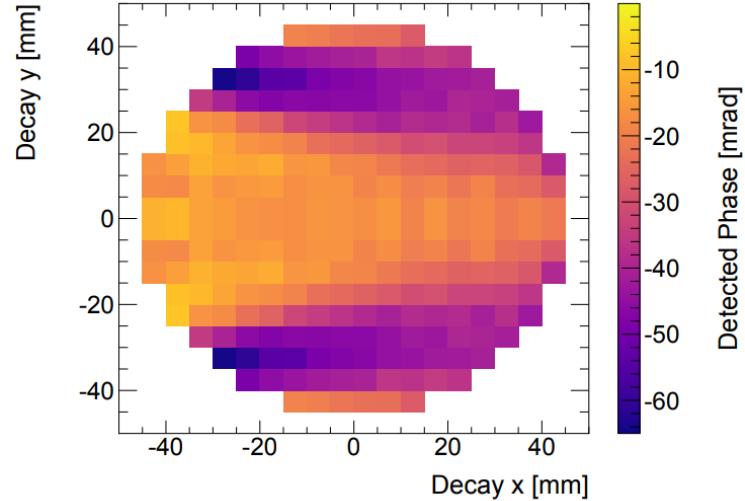
Acceptance map (Geant4)



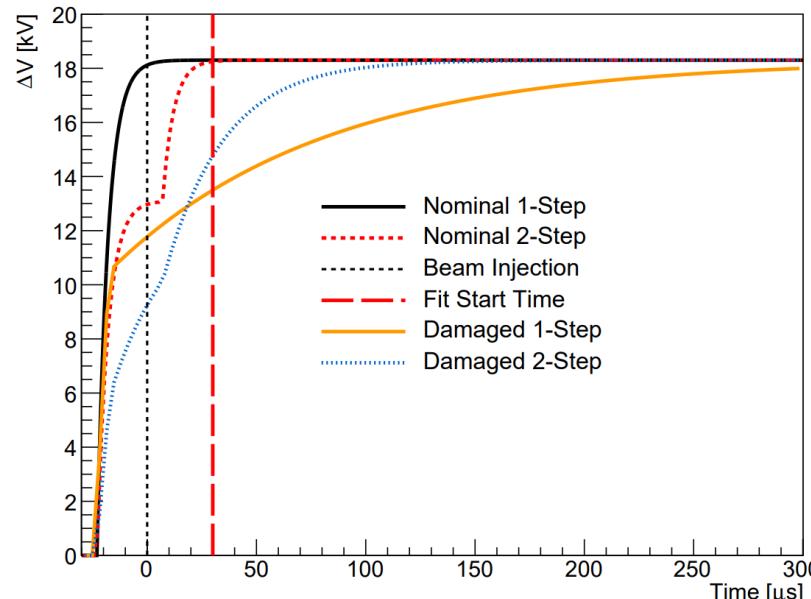
- Origin is acceptance: if e^+ decays outwards then it will have a longer path length to a detector
- We see fewer events from top/bottom of storage region as they miss the detectors vertically

Beam instability and phase-acceptance

Phase map (Geant4)



- Due to acceptance, ϕ depends on muon decay position (x, y)
- Not an issue if the muon distribution doesn't change shape over a fill
- But in Run 1, equipment failure led to beam instability
 - 2/32 quad HV resistors died
→ Focusing E-field changed
→ Beam width changed

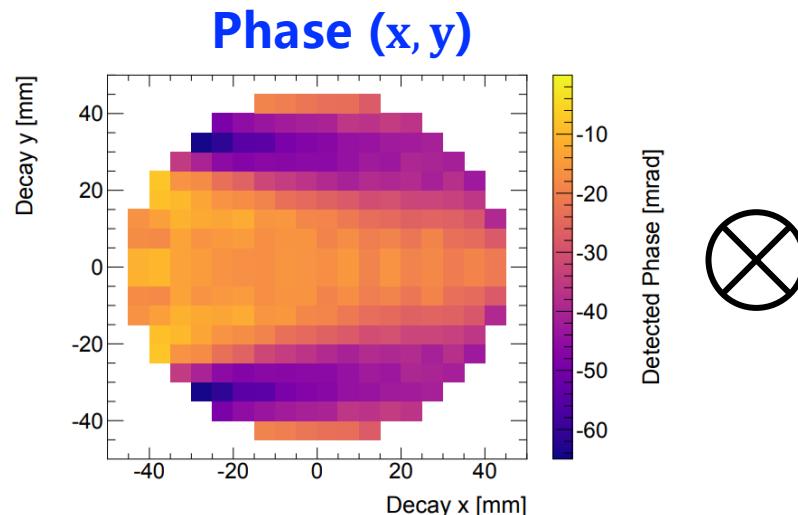


Phase versus time

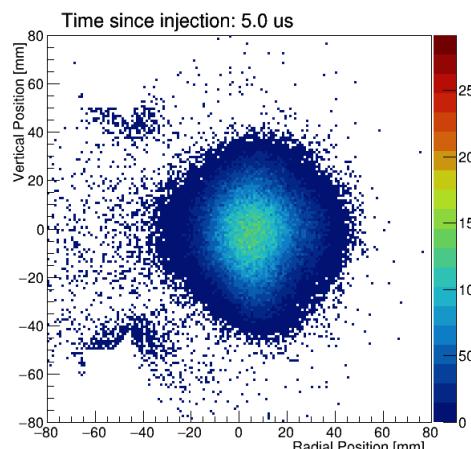
$$\varphi_{pa}^c(t) = \arctan \frac{\sum_{ij} M^c(x_i, y_j, t) \cdot \varepsilon^c(x_i, y_j) \cdot A^c(x_i, y_j) \cdot \sin(\varphi_{pa}^c(x_i, y_j))}{\sum_{ij} M^c(x_i, y_j, t) \cdot \varepsilon^c(x_i, y_j) \cdot A^c(x_i, y_j) \cdot \cos(\varphi_{pa}^c(x_i, y_j))}$$

Acceptance map Phase map

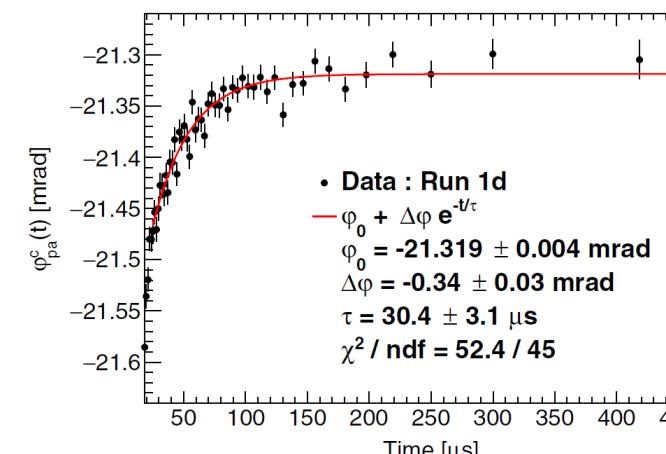
- The (azimuthally-averaged) muon position distribution does change over time.



Muon distribution



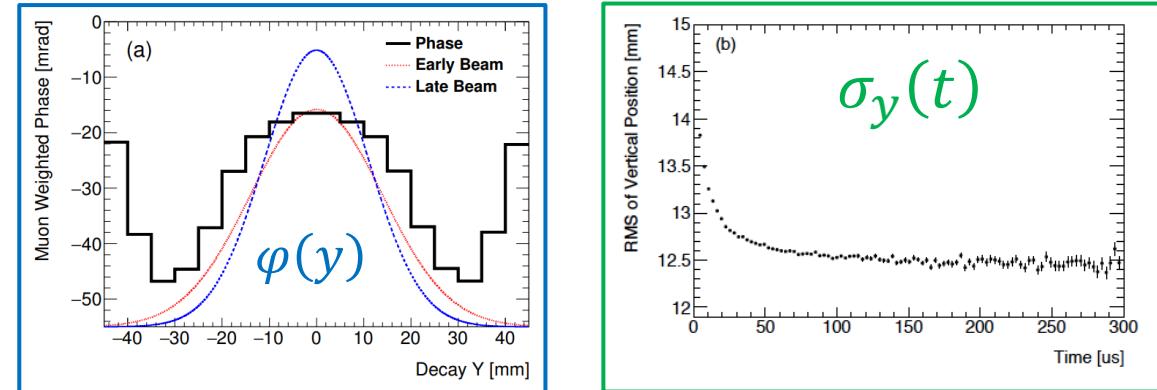
Phase (t)



Phase-acceptance correction

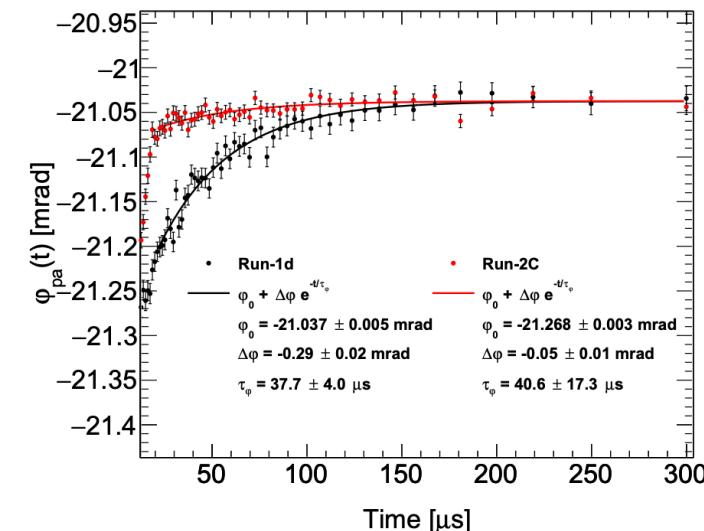
- The dominant effect from Run-1 came from the early-to-late vertical distribution change.

$$\frac{d\varphi}{dt} = \frac{d\varphi}{d\langle\sigma_y^2\rangle} \frac{d\langle\sigma_y^2\rangle}{dt}$$



- The damaged resistors were replaced before Run-2.
- It significantly improved the beam early-to-late stability, and hence reduced the correction and its uncertainty

C_{pa}	Correction [ppb]	Uncertainty [ppb]
Run-1	-158	75
Run-2/3	-27	13



Spin/Phase-Momentum Correlations

- Phase changes due to the coupled effects from ϕ - $\langle p \rangle$ correlation & $\langle p \rangle$ - t correlation.

$$\frac{d\phi}{dt} = \frac{d\phi}{d\langle p \rangle} \frac{d\langle p \rangle}{dt}$$

- Each correlation can be decomposed as follows:

$$\left(\frac{d\phi}{d\langle p \rangle} \right)_{bml}$$

Beamline

Upstream delivery ring dipole bending magnet

$$\left(\frac{d\phi}{d\langle p \rangle} \right)_{p-x}$$

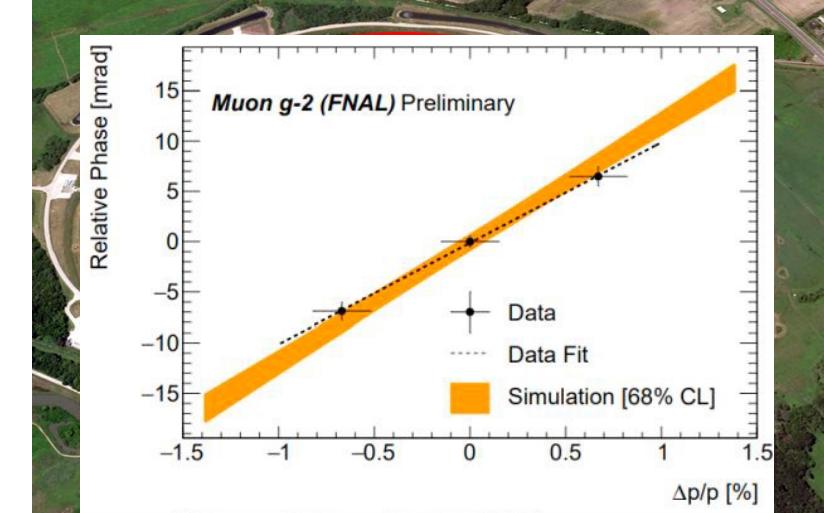
p - x/x'

Spin-Transverse coordinate correlation

$$\left(\frac{d\phi}{d\langle p \rangle} \right)_{p-t_0}$$

p - t₀

Head-to-tail phase difference &
Head-to-tail stored momentum distribution.



Momentum-Time Correlations

- Phase changes due to the coupled effects from ϕ - $\langle p \rangle$ correlation & $\langle p \rangle$ - t correlation.

$$\frac{d\phi}{dt} = \frac{d\phi}{d\langle p \rangle} \frac{d\langle p \rangle}{dt}$$

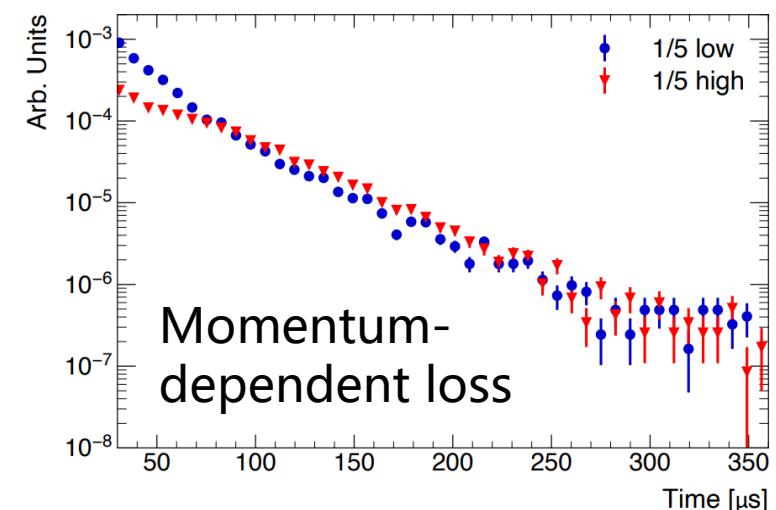
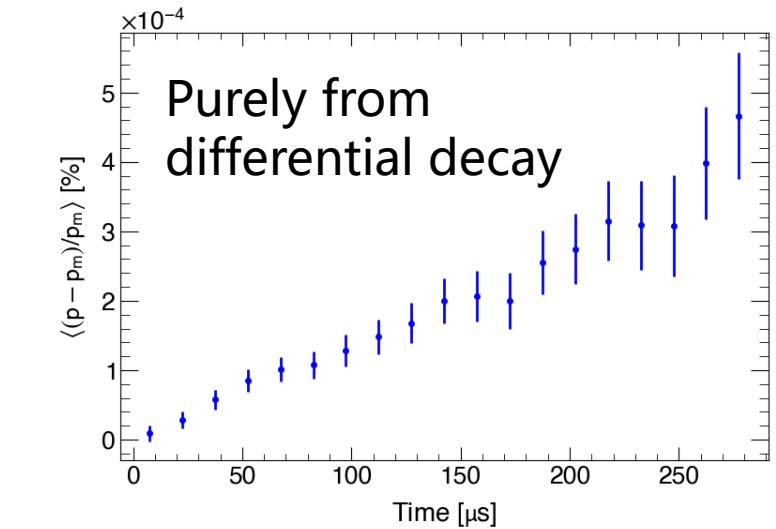
- Each correlation can be decomposed as follows:

$$\left(\frac{d\langle p \rangle}{dt} \right)_{dd}$$

C_{dd} : Differential decay
Muons have different lifetimes depending on their energies.

$$\left(\frac{d\langle p \rangle}{dt} \right)_{ml}$$

C_{ml} : Momentum-dependent loss
Muon loss spectrum depends on the momentum.



C_{ml} and C_{dd} – Muon Losses Correction



- In Run-1, we neglected C_{dd} .
- We were at the early stage of understanding the $p-x/x'$ and $p-t_0$ effects
- The beamline C_{dd} was negligible compared to C_{ml} which was enhanced due to the damaged resistors
- Dominant systematics comes from the bunch-by-bunch deviations in $\left(\frac{d\phi}{d\langle p \rangle}\right)_{p-t_0}$

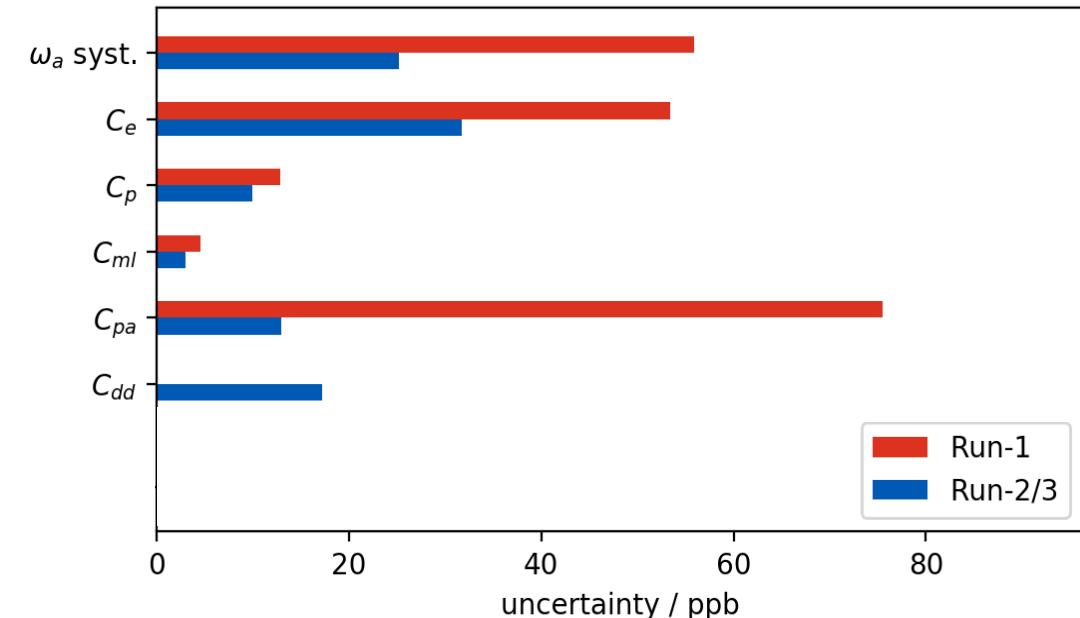
C_{dd}	Correction [ppb]	Uncertainty [ppb]
Run-1	0	0
Run-2/3	-15	17

C_{ml}	Correction [ppb]	Uncertainty [ppb]
Run-1	-11	5
Run-2/3	0	3

Summary of BD corrections

- Beam dynamics corrections to the anomalous spin precession frequency ω_a^m were reviewed
- Improvement to experimental conditions have significantly reduced the major systematics in these corrections
- The net uncertainty of the BD corrections was reduced by more than a factor of 2 in Run-2/3.

Corrections [ppb]	Run-1	Run-2/3
C_e	489 ± 53	451 ± 32
C_p	180 ± 13	170 ± 10
C_{pa}	-158 ± 75	-27 ± 13
C_{dd}	-	-15 ± 17
C_{ml}	-11 ± 5	0 ± 3
Sum	500 ± 93	580 ± 40



Summary and outlook



- Muon g-2 is now measured to 200 ppb using Run 1 to Run 3 of our collected data
 - Just published recently in PRL 131 (2023) 16, 161802
- Current systematics in precession frequency measurement have exceeded our expected goal of 70 ppb
 - 25 ppb vs 57 ppb in TDR
- There are still rooms for improvement for the beam dynamics corrections
 - 40 ppb vs 30 ppb in TDR
- Beam dynamics systematics are the limiting factor down the road
 - A very important aspect to be considered and studied in details for future g-2 experiments using magic momentum approach (HIAF?)

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