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Introduction: QGP physics with ALICE

-
- Quark-gluon plasma (QGP) is a deconfined state of
- hadronic matter which can form at high temperatures
- and/or net baryonic densities
- **ALICE at the LHC observes QGP produced in ultra**
	- relativistic collisions of heavy ions
- **Direct observation of QGP is impossible due to the short**
	- life of the deconfined phase
- **A** QGP is studied indirectly by means of a number of
	- The suppression of high $p_{\rm T}$ particles and jets
	- The enhancement of strange and multi–strange particles **F** Signatures of a collective motion of the medium

Introduction: quarkonium as a probe of QGP

Quarkonium production inside the QGP (min)

- Color screening and dissociations (**sequential melting**) vs **regeneration** mechanisms
- Parton energy loss (collisional vs radiative processes) IF

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Start of collision Development of QGP Hadronization RHIC LHC **POOL OD** OD

Quarkonium: bound state of a heavy quark and its corresponding antiquark Heavy quarks (charm and bottom): produced at the early stages of the heavy-ion collisions

Nature [448, 302-309 \(2007\)](http://volume%20448,%20pages302%E2%80%93309%20(2007))

Introduction: quarkonium in small collision systems

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Measurements in pp collisions:

Provide a baseline for the measurement in p-Pb and Pb-Pb collisions E Study of quarkonium production mechanisms: both perturbative (i.e. heavy-quark pair formation) and non-perturbative (quarkonium state formation) QCD processes involved Measurements of quarkonium production in pp allows to constrain QCD based models

Measurements in p-Pb collisions:

If Investigate cold nuclear matter (CNM) effects (shadowing, coherent parton energy loss,...) \leftrightarrow Help the interpretation of the measurements in Pb-Pb collisions

Measurements in high-multiplicity pp/p-Pb collisions: Study the role of *multiparton interactions* (MPIs) Observe possible *collective-like* effects

Small collision systems (pp and p-Pb collisions): no (or very tiny) QGP effect is expected

Introduction: quarkonium in small collision systems

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Measurements in high-multiplicity pp/p-Pb collisions: Study the role of *multiparton interactions* (MPIs) *E* Observe possible *collective-like* effects

Small collision systems (pp and p-Pb collisions): no (or very tiny) QGP effect is expected $(\binom{n}{\text{min}})$

A Large **I**on **C**ollider **E**xperiment

A Large **I**on **C**ollider **E**xperiment

- Event trigger
- Event characterization
- ↓ Background rejection

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- ITS layers)
- Vertex reconstruction
- Multiplicity estimation

A dedicated heavy-ion experiment at the LHC (min) **V. Silicon Pixel Detector** (SPD, the two innermost

VI. V0 Detectors $(-3.7 < \eta < -1.7$ and $2.8 < \eta < 5.1)$

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Multiplicity dependence of Υ production in pp collisions

Introduction (I)

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D-meson self-normalized yields present a **stronger than**

- **linear increase** as a function of self-normalized charged-
- particle multiplicity, both measured at **midrapidity**
- Models**,** including Percolation and EPOS with
- hydrodynamical expansion, describe qualitatively the data trend
- Hint of collective effects in high charged-particle multiplicity
- events?

Multiplicity dependence of D-meson production in pp collisions

What about the multiplicity dependence of quarkonium production in small collision system?

Υ production as a function of multiplicity

- Self-normalized $\Upsilon(1S)$ yield: *linear increase* with charged-particle multiplicity within uncertainties
- The linear trend results in a *flat trend* of the double ratio of the normalized $\Upsilon(1S)$ yield to the normalized multiplicity
- No firm conclusion can be drawn for the excited states $[\Upsilon(2S)$ and $\Upsilon(3S)]$ due to the limited data sample

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Y production as a function of multiplicity

Model comparison

- **We The linear increase trend is qualitatively** reproduced by Coherent Particle Production (CPP) within large uncertainties
- If $dN_{ch}/dy/(dN_{ch}/dy) < 4$: the linear increase behavior is qualitatively reproduced by PYTHIA 8.2 (with or without CR), and 3-pomeron CGC approach
- If $dN_{ch}/dy/(dN_{ch}/dy) > 4$: the theoretical computations diverge **12-3-pomeron CGC overestimates the observed** trend
	- \mathbb{P} PYTHIA 8.2 underestimates the data trend

Still lack of predictions for bottomonium studies!

Self-normalized yield ratio as a function of multiplicity

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Double ratios (self-normalized signal yield ratios) of $\Upsilon(2S)/\Upsilon(1S)$ and $\Upsilon(3S)/\Upsilon(1S)$. \mathbb{F} Compatible with unity within the large uncertainties **E.** PYTHIA 8.2, CPP and 3-pomeron CGC calculations show a almost flat trend **Comovers approach** predicts a **dissociation** of the excited states, leading to a **suppression** at high multiplicity, especially for the Y(3S)

Self-normalized yield ratio as a function of multiplicity

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The model computations are comparable with unity with uncertainties, indicating the **initial- and** $final{\text -}effects$ act on $\Upsilon(1S)$ and J/Ψ in a similar way

Double ratio of $\Upsilon(1S)/J/\Psi$:

 \mathbb{F} Compatible with unity within the current uncertainties

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Υ(1S) polarization in pp collisions

Introduction (II)

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Quarkonium polarization

Related to the particle spin-alignment with respect to a chosen direction Measured via anisotropies in the decay products angular distributions

Reference frames: (m)

$$
W(\cos\theta, \varphi) \propto \frac{1}{3 + \lambda_{\theta}} (1 + \widehat{\lambda_{\theta}} \cos^2 \theta + \widehat{\lambda_{\phi}} \sin^2 \theta \sin^2 \theta)
$$

- F Helicity (HE): the direction of quarkonium in the center-of-mass frame
- **E** Collins-Soper (CS) the bisector of the angle between the direction of one beam and the opposite of the other beam in the quarkonium rest frame

Introduction (II)

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Polarization in pp collisions: **constrain quarkonium production mechanisms**

- No sizeable polarization is observed for the J/ψ polarization measurement in pp collisions
- Theoretical calculations for J/ψ : **E** NLO NRQCD \rightarrow transverse (longitudinal) polarization in Helicity (Collins-Soper) frame **If NLO CSM** \rightarrow longitudinal (transverse) polarization in Helicity (Collins-Soper) frame

Υ(1S) polarization as a function of transverse momentum

- First ALICE $\Upsilon(1S)$ polarization measurement in pp collisions
	- λ_{θ} compatible with zero (maximum deviation of 1.5 *w.r.t* zero) in both **HE** and **CS** frames
	- λ_φ and $\lambda_{\theta\varphi}$ consistent with zero within uncertainties in both frames

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- First ALICE $\Upsilon(1S)$ polarization measurement in pp collisions
	- λ_θ , λ_φ and $\lambda_{\theta\varphi}$ consistent with zero within uncertainties in both **HE** and **CS** frames
	- Good agreement with LHCb pp data at $\sqrt{s} = 8$ TeV, in a similar rapidity range, within the large experimental uncertainties

LHCb [Collaboration,](https://link.springer.com/article/10.1007/JHEP12(2017)110) *JHEP* 12 (2017) 110

E Qualitatively described by **NLO NRQCD** calculations

[M. Butenschoen](https://journals.aps.org/prl/abstract/10.1103/PhysRevLett.108.172002) *et al.*, *Phys. Rev. Lett.* 108 (2012) 172002

Υ(1S) polarization as a function of transverse momentum

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Υ(1S) polarization as a function of transverse momentum

Improved Color Evaporation Model (ICEM): using the k_T factorization approach

- At low $p_{\rm T}$, the polarization is slightly transverse in the HX frame, while it is slightly longitudinal in the CS frame
- At high p_T , unpolarization is expected in both frames \to c<mark>onsistent with the LHCb data</mark>

Full theoretical description is still missing **Collaboration**, *JHEP* 12 (2017) 110

Multiplicity dependence of Υ production at forward rapidity has been measured with ALICE Self-normalized Υ yield shows a linear increase with increasing multiplicity **F** The results are compared to the model predictions (initial- and/or final-state effects) No significant polarization is observed for Υ (1S) in pp collisions

Midrapidity J/Ψ as a function of multiplicity

- multiplicity
	- *forward (V0)* rapidity
	- rapidity

J/*ψ* production in pp (ALICE)

J/*ψ* production in pp (ALICE)

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Percolation model: **color string interaction**s to describe p+p collisions

In a high-density environment, the coherence among the sources of the color strings leads to a reduction of their effective number. The total charged-particle multiplicity, which originates from soft sources, is more reduced than heavy-particle production for which the *sources have a smaller transverse size*

Multiplicity Dependent $\psi(2S)$ to J/ψ

- onset of QGP-like effects
- multiplicity dependence more or less consistent with unity

J/*ψ* production in pp (ALICE)

• Multiplicity-dependent studies in small systems provide a testing ground for examining the

• PHENIX ($\sqrt{s_{NN}}$ =200 GeV) and ALICE ($\sqrt{s_{NN}}$ =13 TeV) results consistent, with weak

• Note that ALICE results have charged particle multiplicity measured at mid-rapidity

J/*ψ* production in pp (STAR)

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J/*ψ* production in pp (PYTHIA)

Regions of azimuthal angle

Region for each track defined from angle φ between the track and J/ψ candidate

Regions constructed regarding a pair candidate J/ψ and not precisely a $J/\psi \rightarrow$ Would this influence the results?

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Isolate different effects in different regions:

- Toward: particles associated to J/ψ production
- Transverse: underlying event
- Away: recoil jet

J/*ψ* production in pp (PYTHIA)

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Pythia prediction (only one parsonic interaction activated) Eur. Phys. J. C 79, 36 (2019)

3. PHENIX Results

• PHENIX, STAR, ALICE (Measuring multiplicity at the same acceptance with J/ψ) → Similar multiplicity dependence despite different center-of-mass energy

J/*ψ* production in pp collisions

J/*ψ* production in pp (CMS)

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Cross section ratio

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J/*ψ* production in pp (CMS)

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Cross section ratio

J/*ψ* polarization in pp

J/*ψ* polarization in pp (STAR)

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 $\widehat{e_{i}}$ of

Upsilon polarization in pp (CMS)

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 $\widehat{e_{i0}}$

Figure 1.26: The polarization parameter $\lambda_{\vartheta}^{X_{c2}}$ values measured when the $\lambda_{\vartheta}^{X_{c1}}$ values are fixed to the unpolarized (left) or the NRQCD (right) scenarios as a function of p_T/M of the J/ ψ [140]. The purple band on the right is the NRQCD prediction for $\lambda_{\vartheta}^{\chi_{c2}}$ [95].

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 $\widehat{e_{i}}$ of β

$\chi_{\rm c2}$ production in pp (CMS)

J/*ψ* polarization in pp and Pb—Pb (ALICE)

ICEM with the collinear factorization approach for the direct ${\mathrm J}/\psi$ component, compared to the inclusive one

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No significant difference is observed between the Pb—Pb and pp collisions for the ICEM model

J/*ψ* polarization in pp and Pb—Pb collisions

ALICE measured J/ ψ polarization in Pb—Pb collisions

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- \mathbb{F} All polarization parameters are close to zero within uncertainties
	- λ_θ shows a maximum 2 σ deviation w.r.t zero in both HE and CS frames for $2 < p_T < 4$ GeV/ c
- **E.** Compatible with **ALICE** results in **pp** collisions within uncertainties *EPJC* [78 \(2018\) 562](https://link.springer.com/article/10.1140/epjc/s10052-018-6027-2)
- difference w.r.t **LHCb** in **pp** collisions in **HE** frame 3*σ* **LHCb** [Collaboration,](https://link.springer.com/article/10.1140/epjc/s10052-013-2631-3) *EPJC* 73 (2013) 11
- Difference due to **suppression**/**regeneration** effects in Pb—Pb w.r.t pp collisions?
- What is the role of the angular momentum (L) and the magnetic fields (B)? ⃗ ⃗ **ALI-PUB-490215**

Phys. Lett. B [815 \(2021\) 136146](https://www.sciencedirect.com/science/article/pii/S0370269321000861?via=ihub)

J/*ψ* polarization in theory

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 \mathscr{D} [Phys. Rev. Lett.](https://journals.aps.org/prl/abstract/10.1103/PhysRevLett.108.172002) 108 (2012) 172002

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J/*ψ* polarization in Pb—Pb

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EVENT-PLANE

Reference frame:

- \mathbb{F} Event-plane based frame (EP): axis orthogonal to the EP in the collision center-of-mass frame
- EP normal to *B* and *L*

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Light flavor hadrons (K^{*0}, Φ) polarization in Pb—Pb

- ρ_{00} measurement for light flavor hadrons in Pb—Pb collisions at
	- $\overline{s_\mathrm{NN}} = 2.76$ TeV and in pp collisions at \sqrt{s} = 13 TeV
	- $p_{\rm T}$ dependence
		- ρ_{00} < 1/3 for K^{*0} and Φ at low $p_{\rm T}$ (smaller central value for K^{*0}) in Pb—Pb collisions
		- ρ_{00} ~ 1/3 for:
			- $p_{\textrm{T}}^{K^{*0}}>2$ GeV/ c and $p_{\textrm{T}}^{\Phi}>0.8$ GeV/ c
		- A random event plane (**RP**)
			- K^{*0} and Φ in pp collisions

Zero spin hadron K_S^0 **: no spin alignment is observed** *S*

The spin-density matrix element ρ_{00} is determined from the distribution of the angle θ^* between the kaon decay daughter and the quantization axis in the decay rest frame [16], [B],

$$
\frac{dN}{d\cos\theta^*} \propto [1 - \rho_{00} + \cos^2\theta^* (3\rho_{00} - 1)]. \tag{1}
$$

 ρ_{00} measurement for light flavor hadrons in Pb–Pb collisions at $\sqrt{s_\text{NN}} = 2.76$ TeV and in pp collisions at \sqrt{s} = 13 TeV Gentrality dependence ρ_{00} deviates w.r.t 1/3 **at low** $p_{\rm T}$ in semi-central **collisions** No centrality dependence of ρ_{00} at high $p_{\rm T}$

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Light flavor hadrons (K^{*0}, Φ) polarization in Pb-Pb

Charm mesons polarization in pp

- Charmed vector meson (D^*+) polarization crucial to $(\binom{n}{\text{min}})$ complete the picture in HICs
- $\text{D}^{\ast+}$ polarization in pp collisions at $\sqrt{s} = 13$ TeV
	- ρ_{00} spin matrix element ($1/3$ means no polarization)
		- Prompt D^{*+} (c \rightarrow D^{*+}) unpolarized
		- Non-zero polarization for non-prompt D^{*+} (b $\rightarrow D^{*+}$)
		- Both well predicted by PYTHIA 8 + EVTGEN

Charm mesons polarization in pp

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PYTHIA 8: only considering the direct component

Introduction: quarkonium as a probe of QGP

Charmonia nuclear modification factor R_{AA}

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Transport model (TAMU) is in good agreement with results as a function of $p_{\rm T}$ and centrality

Statistical hadronization model (SHMc) tends to underestimate $ψ(2S)$ in central collisions

TAMU: *Rapp et al, Nucl. Phys. A943 (2015) 147-158* **SHMc**: *Andronic et al, Phys. Lett. B 797 (2019) 134836*

In SHMc, charmonia are produced (regenerated) from *deconfined charm quarks* at the QCD phase boundary

whereas TAMU includes *regeneration* in the QGP phase

Introduction: quarkonium as a probe of QGP

Bottomonia nuclear modification factor R_{AA}

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- Stronger suppression for $\Upsilon(2S)$ w.r.t $\Upsilon(1S)$
- The various calculations (e.g. Transport model w/ or w/o $(\overbrace{(\overbrace{\cdots})})$ *regeneration*) reproduce the trend of the data within the corresponding uncertainties
- Weak regeneration effect expected in the bottom sector $(\binom{n}{k}$

ALICE Collaboration, *[Phys. Lett .B](https://arxiv.org/pdf/2011.05758.pdf) 822 (2021) 136579*

Dataset: collected in 2016, 2017 and 2018 in pp collisions at $\sqrt{s} = 13$ TeV

Data sample

- Event selection:
	- **E** Unlike-sign dimuon (CMUL7) trigger and minimum bias (MB) trigger \mathbb{F} Physics selection with default pileup rejection ${\sf SPD}$ vertex selection ($N_{\rm~contributors} > 0$, $\sigma_{z_{\rm vtx}} < 0.25$ cm, $|z_{\rm vtx}| < 10$ cm)
	-
	-
- In this analysis: $(\binom{n}{k}$
	- The signal Y is measured at forward rapidity (2.5<y<4.0), the charged-particle multiplicity is measured at central rapidity (|η|<1)
	- The Y yield ($\rm dN_Y/\rm dy$) and the pseudorapidity charged-particle multiplicity density ($\rm dN_{ch}/d\eta$) are both measured for $INEL > 0$ events

Analysis strategy

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Signal extraction

- IF Υ analysis is performed measuring the number of produced resonances
- \mathbb{R} After the standard muon and dimuon selections, invariant mass spectra muon pairs are built
- The invariant mass spectra are fitted with a LF combination of phenomenological functions for peaks and background
- The number of produced resonances is IF extracted through the integral of the peak

Data sample and analysis strategy

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Data sample: collected in 2016, 2017 and 2018 in pp collisions at \sqrt{s} = 13 TeV, dimuon trigger +

Signal extraction: the number of $\Upsilon(1S)$ candidates is obtained via a fit procedure on the dimuon invariant

Acceptance x efficiency correction: the raw number of $\Upsilon(1S)$ extracted from the fit procedure is corrected for a factor quantifying the geometrical acceptance and reconstruction efficiency effects,

Polarization parameters determination: the polarization parameters λ_θ , λ_φ and $\lambda_{\theta\varphi}$ are extracted by fitting the acceptance- and efficiency-corrected angular distributions of $\Upsilon(1S)$, in both the Helicity and

- quality criteria used to select the events
- Analysis procedure:
	- mass distribution in each angular interval considered in the analysis
		- estimated via a MC simulation
			-

Collins-Soper reference frames

Analysis strategy

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Polarization parameters determination

 λ_θ , λ_φ and $\lambda_{\theta\varphi}$ extracted by fitting the A \times ε -corrected $\Upsilon(1S)$ angular distributions in both frames simultaneously

$$
W(\cos \theta) \propto \frac{1}{3 + \lambda_{\theta}} (1 + \lambda_{\theta} \cos^{2} \theta)
$$
\n
$$
W(\varphi) \propto 1 + \frac{2\lambda_{\varphi}}{3 + \lambda_{\theta}} \cos 2\varphi
$$
\n
$$
W(\tilde{\varphi}) \propto 1 + \frac{\sqrt{2}\lambda_{\theta\varphi}}{3 + \lambda_{\theta}} \cos 2\tilde{\varphi}
$$
\n
$$
\left[\begin{array}{cc}\n\tilde{\varphi} = \varphi - 3\pi/4, & \cos \theta < 0 \\
\varphi = \varphi - \pi/4, & \cos \theta > 0\n\end{array}\right]
$$
\n
$$
\left[\begin{array}{cc}\n\tilde{\varphi} = \varphi - \pi/4, & \cos \theta < 0 \\
\varphi = \varphi - \pi/4, & \cos \theta > 0\n\end{array}\right]
$$
\n
$$
\left[\begin{array}{cc}\n\tilde{\varphi} = \varphi - \pi/4, & \cos \theta > 0 \\
\frac{1}{\varphi} = \frac{1}{\sqrt{2}} \cos \theta - \frac{1}{\sqrt{2}} \cos \theta\n\end{array}\right]
$$
\n
$$
\left[\begin{array}{cc}\n\tilde{\varphi} = \varphi - \pi/4, & \cos \theta > 0 \\
\frac{1}{\varphi} = \frac{1}{\sqrt{2}} \cos \theta - \frac{1}{\sqrt{2}} \cos \theta\n\end{array}\right]
$$
\n
$$
\left[\begin{array}{cc}\n\tilde{\varphi} = \varphi - \pi/4, & \cos \theta > 0 \\
\frac{1}{\varphi} = \frac{1}{\sqrt{2}} \cos \theta - \frac{1}{\sqrt{2}} \cos \theta\n\end{array}\right]
$$

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Prospects for Run 3:

Significantly higher statistical data sample compared to Run 1 and 2 is expected ($L_{\rm int} = 10\,\,{\rm nb^{-1}}$ in Pb—Pb, $200\,{\rm ~nb^{-1}}$ in pp)

During the LS2 in 2019 - 2021, all detectors have been upgraded (min)

ALICE 2 - LS2 upgrade

Muon Forward Tracker (MFT, 2.5 < η < 3.6)

- \mathbb{F} Installed in the forward area near the interaction point
	- ↓ Matching muon track with MFT tracks
	- High-precision vertexing capabilities

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- **E** New silicon chips with MAPS technology (ALPIDE)
	- Pixel size = 27 x 27 μm^2
	- Position resolution ~ 5 *μ*m \blacktriangleright
- **Physics motivations**
	- ↓ Prompt and non-prompt charmonia disentangling
	- Precise measurement of low-mass dimuon

Both Central China Normal University and [Institut de Physique des 2 Infinis](https://mmm.cern.ch/owa/redir.aspx?C=ygHdjav-LCVmdLlPIsOrAeMq1kp8fVebxwN7giriYXSyW9DjuJ7YCA..&URL=https%3a%2f%2falice-glance.cern.ch%2falice%2fmembership%2finstitutes%2fdetails.php%3fid%3d1239%26view%3dabout) [de Lyon](https://mmm.cern.ch/owa/redir.aspx?C=ygHdjav-LCVmdLlPIsOrAeMq1kp8fVebxwN7giriYXSyW9DjuJ7YCA..&URL=https%3a%2f%2falice-glance.cern.ch%2falice%2fmembership%2finstitutes%2fdetails.php%3fid%3d1239%26view%3dabout) are involved in this MFT project

Ultra-relativistic heavy-ion collisions

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ALICE data taking history

Charm mesons polarization in pp

$$
:1^{--}]
$$

$$
\langle J, J_z \rangle = b_{+1} | 1, +1 \rangle + b_0 | 1, 0 \rangle + b_{-1} | 1, -1 \rangle
$$

Spin-alignment \Leftrightarrow decay products angular distribution
EPIC 69 (657-673), 2010, Faccoli et al.

Dilepton decay angular distribution

$$
\phi) \propto \frac{1}{3 + \lambda_{\theta}} \cdot (1 + \lambda_{\theta} \cos^2 \theta + \lambda_{\phi} \sin^2 \theta \cos 2\phi + \lambda_{\theta \phi} \sin 2\theta \cos \phi)
$$

