



北京大學  
PEKING UNIVERSITY

第十七届 TeV 物理工作组学术研讨会

# Long-lived Searches of Vector-like Lepton and Its Accompanying Scalar at Colliders

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Based on

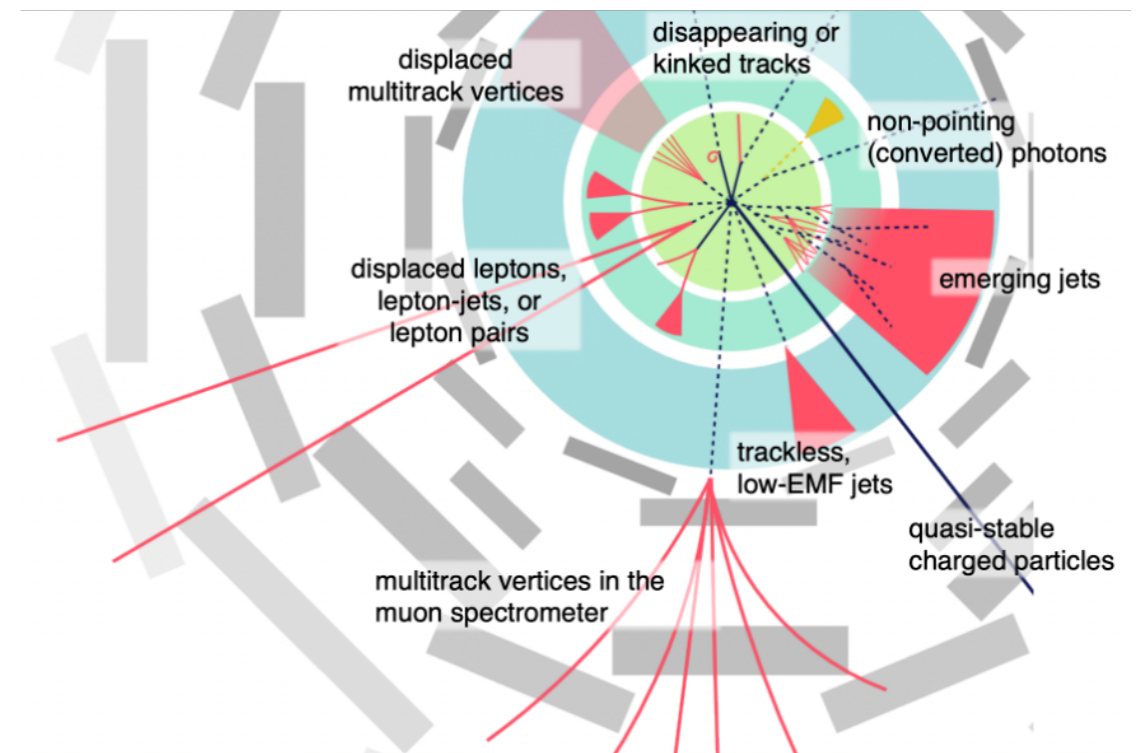
Qing-Hong Cao, Jinhui Guo, Jia Liu, Yan Luo, and Xiao-Ping Wang  
arXiv: 2311.12934

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# Introduction

- Vector-like fermion:  
non-chiral representation of SM  
gauge group
- More and more searches for  
vector-like lepton (VLL):  
but prompt decay 2202.08676. CMS Collaboration  
2303.05441. ATLAS Collaboration
- Small chiral mass mixing:  
natural long-lived particle
- Bypass the constraints from  
Lepton Flavor Violation



ATLAS Collaboration

We consider long-lived singlet VLL and its accompanying scalar.



# Models

- Vector-like Lepton Scenario (VLLS)

$$\mathcal{L}_{\text{eff}}^F \supset \bar{F}^0 i D_\mu \gamma^\mu F^0 + \bar{L}^0 i D_\mu \gamma^\mu L^0 + \bar{\ell}_R^0 i D_\mu \gamma^\mu \ell_R^0 - m_F^0 \bar{F}^0 F^0 - m_\ell^0 \bar{\ell}_R^0 \ell_R^0 - (\delta \bar{F}_L^0 \ell_R^0 + \text{h.c.})$$

	$F^\pm$	$\ell_R^0$
$SU(2)_L$	1	1
$U(1)_Y$	-2	-2

- diagonalizing the mass terms

$$\mathcal{L}_{\text{mass}}^F = (\bar{F}_L^0, \bar{\ell}_L^0) \begin{pmatrix} m_F^0 & \delta \\ 0 & m_\ell^0 \end{pmatrix} \begin{pmatrix} F_R^0 \\ \ell_R^0 \end{pmatrix} = (\bar{F}_L^0, \bar{\ell}_L^0) M \begin{pmatrix} F_R^0 \\ \ell_R^0 \end{pmatrix} = (\bar{F}_L, \bar{\ell}_L) U_L M U_R^\dagger \begin{pmatrix} F_R \\ \ell_R \end{pmatrix}$$

$$U_L M U_R^\dagger = \text{diag}(m_F, m_\ell) \quad \text{with} \quad U_L = \begin{pmatrix} \cos \theta_L & -\sin \theta_L \\ \sin \theta_L & \cos \theta_L \end{pmatrix} \quad U_R = \begin{pmatrix} \cos \theta_R & -\sin \theta_R \\ \sin \theta_R & \cos \theta_R \end{pmatrix}$$

mixing angle

$$\tan \theta_R \simeq -\frac{\delta}{m_F^0} \quad \tan \theta_L \simeq -\frac{m_\ell^0 \delta}{(m_F^0)^2} \simeq \frac{m_\ell^0}{m_F^0} \tan \theta_R$$

$$m_F \simeq m_F^0 + \frac{\delta^2}{2m_F^0} \simeq m_F^0 \quad m_\ell \simeq m_\ell^0 \left( 1 - \frac{1}{2} \left( \frac{\delta}{m_F^0} \right)^2 \right)$$



# Models

- Vector-like Lepton Scenario (VLLS)

$$\mathcal{L}_{\text{eff}}^F \supset \bar{F}(i\partial_\mu - eA_\mu + e \tan \theta_W Z_\mu)\gamma^\mu F - m_F \bar{F}F - m_\ell \bar{\ell}\ell$$
$$+ \frac{1}{2} \frac{e}{\sin \theta_W \cos \theta_W} \theta_L Z_\mu (\bar{F}_L \gamma^\mu \ell_L + \text{h.c.}) - \frac{e}{\sqrt{2} \sin \theta_W} \theta_L (W_\mu^+ \bar{\nu}_L \gamma^\mu F_L + \text{h.c.})$$

Decay rate of  $F^\pm$ :

$$\Gamma(F^\pm \rightarrow \nu_\ell W^\pm) = \frac{\theta_L^2 g_W^2}{64\pi} \frac{(m_F^2 - m_W^2)^2 (m_F^2 + 2m_W^2)}{m_F^3 m_W^2}$$

$$\Gamma(F^\pm \rightarrow \ell^\pm Z) = \frac{\theta_L^2 g_Z^2}{64\pi} \frac{(m_F^2 - m_Z^2)^2 (m_F^2 + 2m_Z^2)}{m_F^3 m_Z^2}$$

small  $\theta_L \rightarrow$  long-lived  $F^\pm$



# Models

- Vector-like Lepton with Scalar (VLLWS)

$$\begin{aligned}\mathcal{L}_{\text{Int}}^\phi &\supset -y\phi\bar{F}_L^0\ell_R^0 + \text{h.c.} \\ &\simeq -y\phi(\bar{F}_L\ell_R + \bar{\ell}_R F_L + \theta_R\bar{F}F - \theta_L\bar{\ell}\ell)\end{aligned}$$

	$F^\pm$	$\phi$
$SU(2)_L$	1	1
$U(1)_Y$	-2	0

Decay rate of  $F^\pm$  and  $\phi$ :

$$\Gamma(F^\pm \rightarrow \phi\ell^\pm) = \frac{y^2(m_F^2 + m_\ell^2 - m_\phi^2)\sqrt{\lambda(m_F^2, m_\phi^2, m_\ell^2)}}{32\pi m_F^3}$$

$$\Gamma(\phi \rightarrow \ell^+\ell^-) = \frac{(y\theta_L)^2 m_\phi (1 - 4m_\ell^2/m_\phi^2)^{3/2}}{8\pi}$$

$$\tau(\phi) \simeq \left(\frac{3 \times 10^{-9}}{y\theta_L}\right)^2 \left(\frac{50 \text{ GeV}}{m_\phi}\right) \text{ ns}$$

$$\lambda(x, y, z) = x^2 + y^2 + z^2 - 2xy - 2yz - 2zx$$



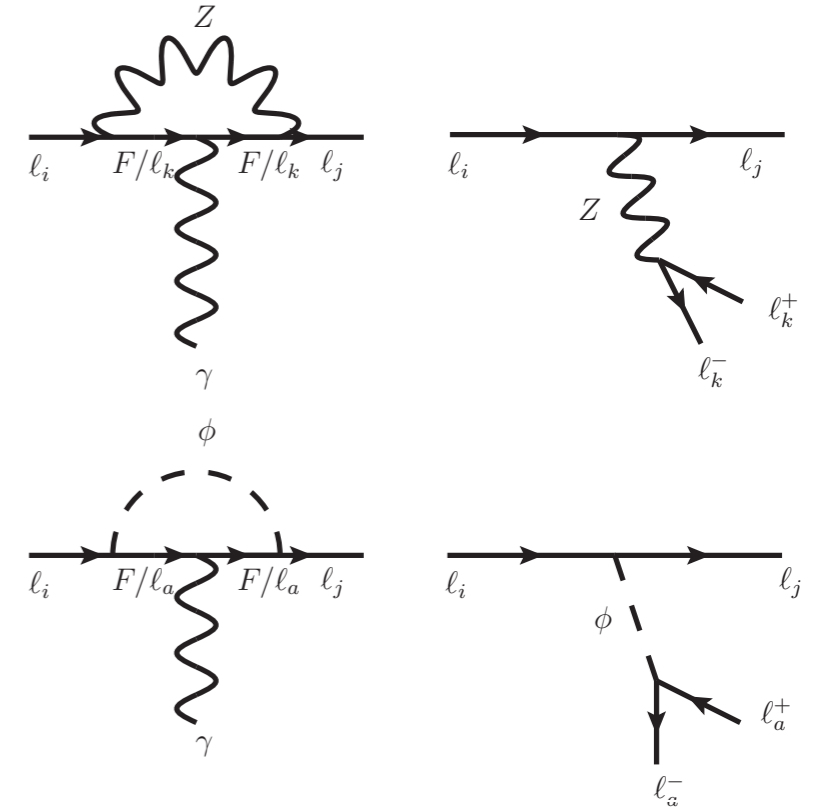
# Existing Constraints

- Constraints from Lepton Flavor Violations

Generally, the mass matrix is

$$x_i = U_{R,i2}^{\text{SM}}$$

$$M_{\text{mass}} = \begin{pmatrix} m_{ee} & m_{e\mu} & m_{e\tau} & 0 \\ m_{\mu e} & m_{\mu\mu} & m_{\mu\tau} & 0 \\ m_{\tau e} & m_{\tau\mu} & m_{\tau\tau} & 0 \\ 0 & \delta & 0 & m_F^0 \end{pmatrix} \longrightarrow M'_{\text{mass}} = \begin{pmatrix} m'_e & 0 & 0 & 0 \\ 0 & m'_\mu & 0 & 0 \\ 0 & 0 & m'_\tau & 0 \\ \delta \cdot x_1 & \delta \cdot x_2 & \delta \cdot x_3 & m_F^0 \end{pmatrix}$$



$$\text{Br}(\mu \rightarrow eee) < 10^{-12}, \text{Br}(\tau \rightarrow \ell_i \ell_j \ell_j) \lesssim 2.7 \times 10^{-8}, \text{Br}(\ell_i \rightarrow \ell_j \gamma) \lesssim 4.2 \times 10^{-13}$$

can be easily satisfied in the parameter regions of our LLP study

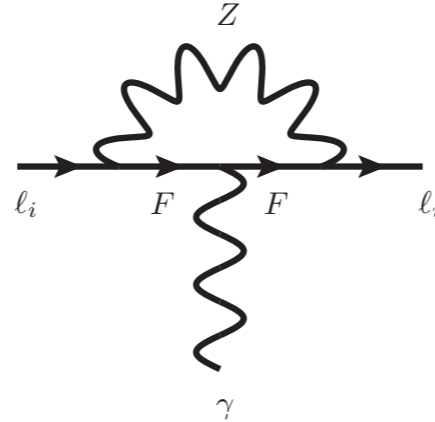
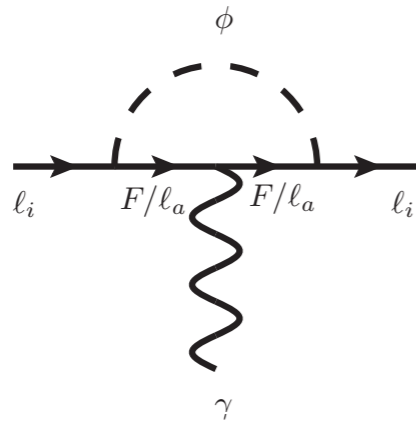
$$\left\{ y\theta_R < 10^{-2}, y\theta_L < 0.5 \times 10^{-3}, m_F > 200 \text{ GeV and } m_\phi > 20 \text{ GeV} \right\}$$

In the following phenomenology discussion, we assume  $\ell_0$  to be SM mass eigenstate.



# Existing Constraints

- Constraints from  $g-2$



$$(g - 2)_\mu^{\text{BSM}} < 2.49 \times 10^{-9}$$

$$(g - 2)_e^{\text{BSM}} < 9.8 \times 10^{-13}$$

0801.1134. D. Hanneke et al.

2308.06230. D. P. Aguillard et al.

the dominant contribution:

$$(g - 2)_{\ell_i}^{F,\phi} \simeq \begin{cases} 2.64 \times 10^{-10} \left(\frac{y}{1}\right)^2 \left(\frac{m_{\ell_i}}{0.1 \text{ GeV}}\right)^2 \left(\frac{200 \text{ GeV}}{m_F}\right)^2 (m_F > m_\phi \gg m_{\ell_i}), \\ 1.32 \times 10^{-10} \left(\frac{y}{1}\right)^2 \left(\frac{m_{\ell_i}}{0.1 \text{ GeV}}\right)^2 \left(\frac{200 \text{ GeV}}{m_F}\right)^2 (m_\phi = m_F \gg m_{\ell_i}), \end{cases}$$

we set  $y = 1$  in the following discussion

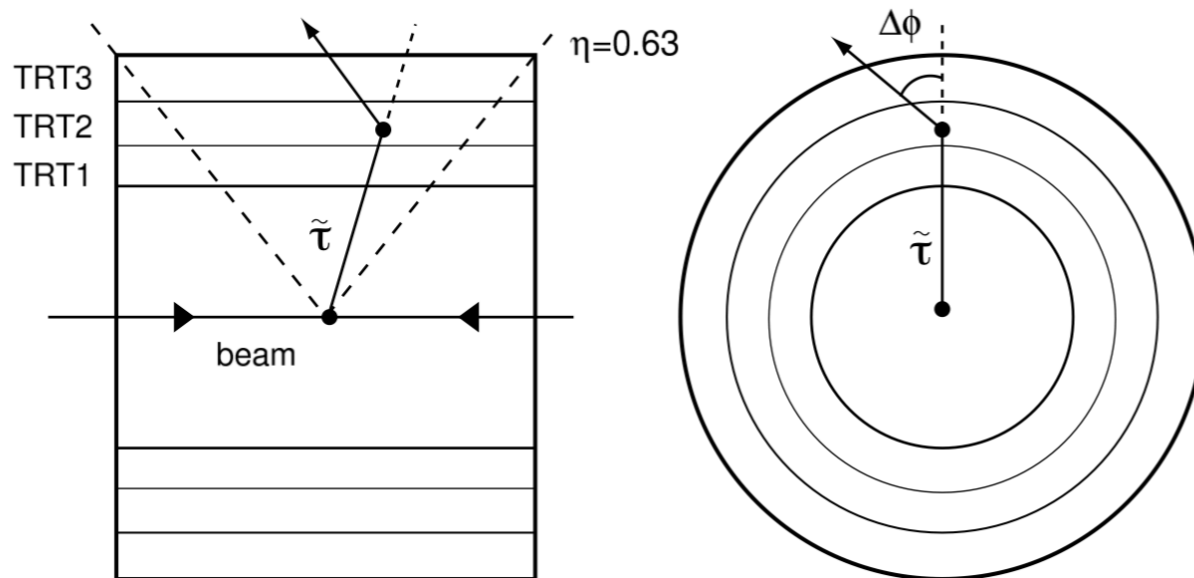
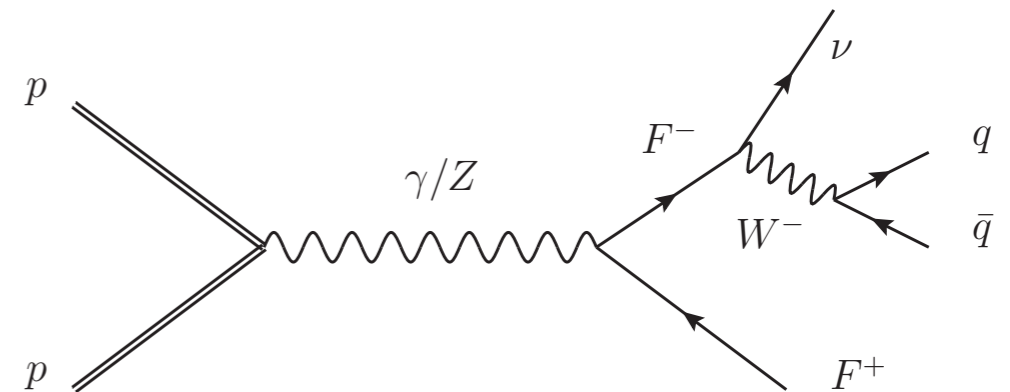
# LLP Signatures at Collider

- Long-lived VLL: Kink Signatures

Signal process:

$$pp \rightarrow F^- F^+, F^- \rightarrow W^- \nu_e, W^- \rightarrow q \bar{q}$$

$F$  mixed with  $\mu$



Kink Track in the TRT.

Inner Track
Pixel detector
Semi-conductor tracker (SCT)
Transition radiation tracker (TRT)
$563 \text{ mm} < R_1 < 694 \text{ mm}$
$697 \text{ mm} < R_2 < 860 \text{ mm}$
$863 \text{ mm} < R_3 < 1066 \text{ mm}$



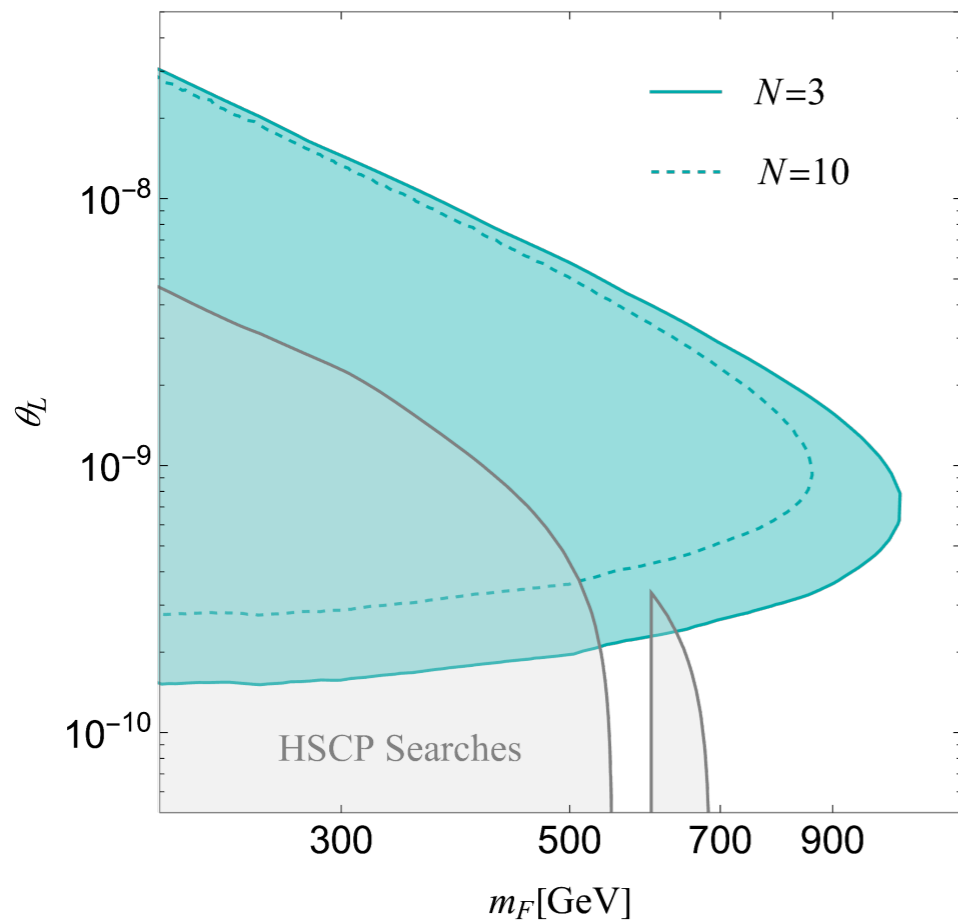


# LLP Signatures at Collider

- Long-lived VLL: Kink Signatures

free parameters

$$\{m_F, \theta_L\}$$



Selection criteria:

$$\text{KT} : p_T^F > 100 \text{ GeV}, |\eta_F| < 0.63, 0.1 < \Delta\phi < \pi/2,$$
$$563 \text{ mm} < r_F < 863 \text{ mm}, |z_F| < 712 \text{ mm},$$
$$p_T^q > 10 \text{ GeV}, |z_F/r_F| > |p_z^\mu/p_T^\mu|$$

$$\text{HL-LHC: } \sqrt{s} = 13 \text{ TeV}, \mathcal{L} = 3 \text{ ab}^{-1}$$

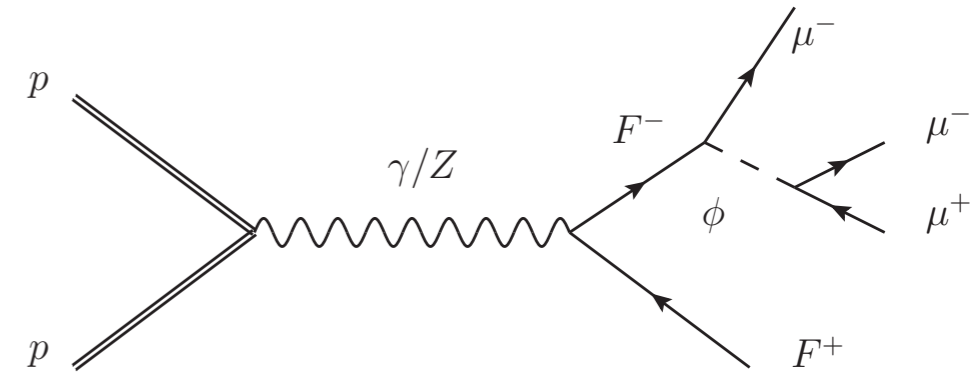
Constraints from the  
Heavy Stable Charged Particle searches.



# LLP Signatures at Collider

- Long-lived Scalar Particle at HL-LHC

Signal process:  $pp \rightarrow F^-F^+, F^\pm \rightarrow \phi l^\pm, \phi \rightarrow l^+l^-$



Time delay:

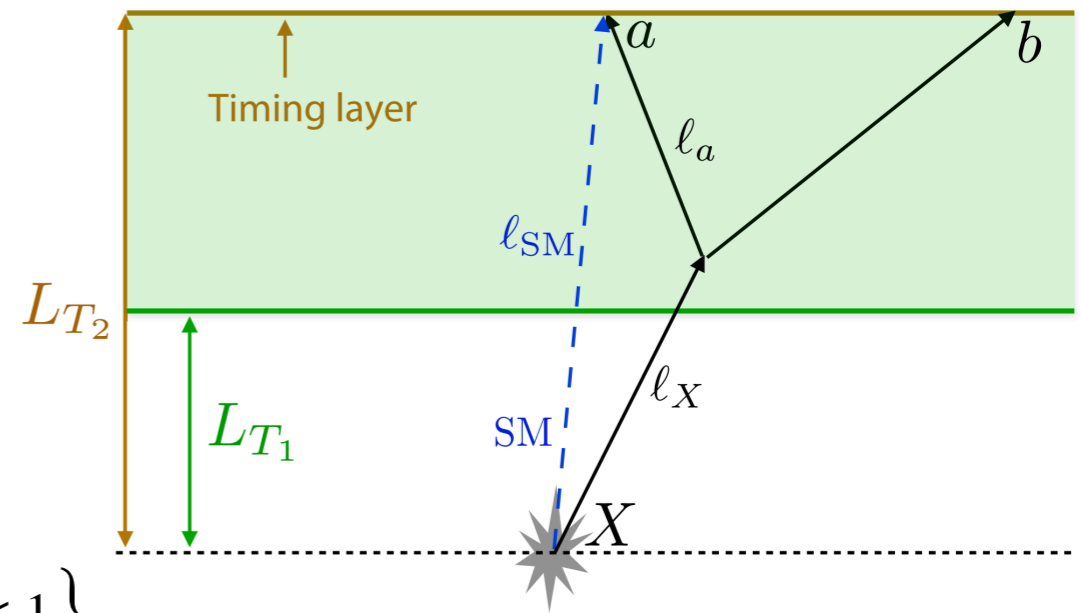
$$\Delta t_\ell = L_\phi/\beta_\phi + L_\ell/\beta_\ell - L_{SM}/\beta_{SM}$$

free parameters

$$\{m_F, m_\phi, y, \theta_L\}$$

Benchmark parameters setting:

- $\{m_F/m_\phi = 2, 5, 10, m_F > 200 \text{ GeV}, \theta_L \ll 1\}$
- $\{m_\phi \in [20 \text{ GeV}, m_F], m_F = 1 \text{ TeV}, 1.2 \text{ TeV}, \theta_L \ll 1\}$
- $\{m_\phi = 100 \text{ GeV}, 300 \text{ GeV}, 500 \text{ GeV}, m_F > 200 \text{ GeV}, \theta_L \ll 1\}$



1805.05957. Jia Liu et al.

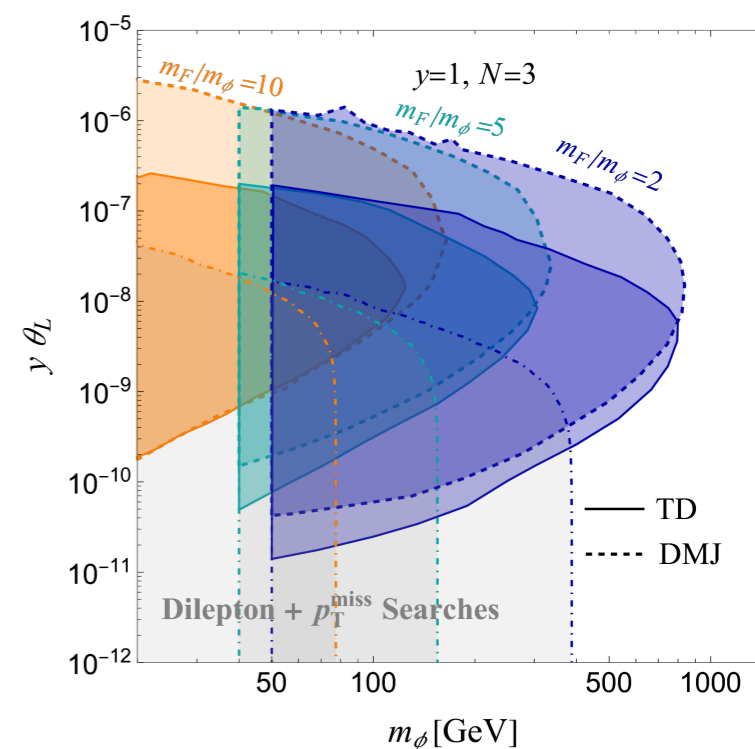
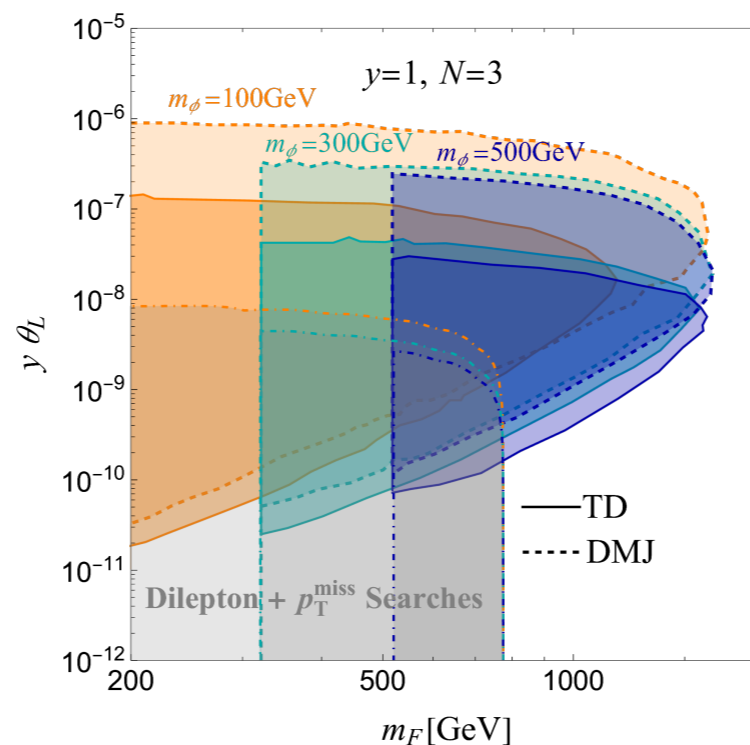
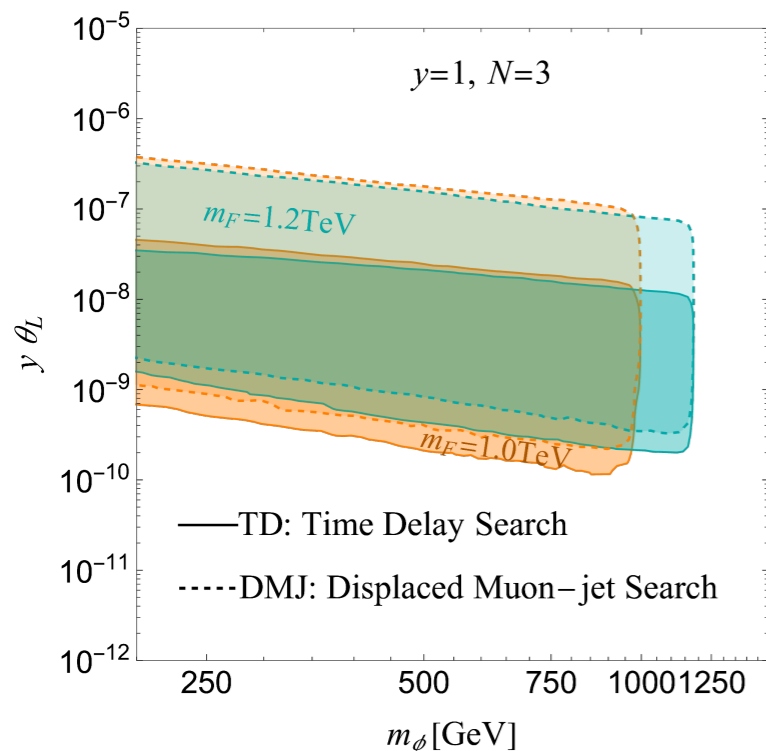
# LLP Signatures at Collider

- Long-lived Scalar Particle at HL-LHC

DMJ :  $p_T^{\mu,F} > 30 \text{ GeV}$ ,  $p_T^{\mu,\phi} > 5 \text{ GeV}$ ,  $r_\phi < 30 \text{ cm}$ ,  $d_0^{\mu,\phi} > 1 \text{ mm}$  (displaced muon-jets)

TD :  $p_T^{\mu,F} > 30 \text{ GeV}$ ,  $p_T^{\mu,\phi} > 5 \text{ GeV}$ ,  $|\eta_\mu| < 2.4$ ,  $\Delta t_{\mu,\phi} > 0.3 \text{ ns}$ ,  $5 \text{ cm} < r_\phi < 1.17 \text{ m}$ ,  $z_\phi < 3.04 \text{ m}$

CMS Minimum Ionizing Particle (MIP) timing detector



Constraints from dilepton plus missing transverse momentum searches.

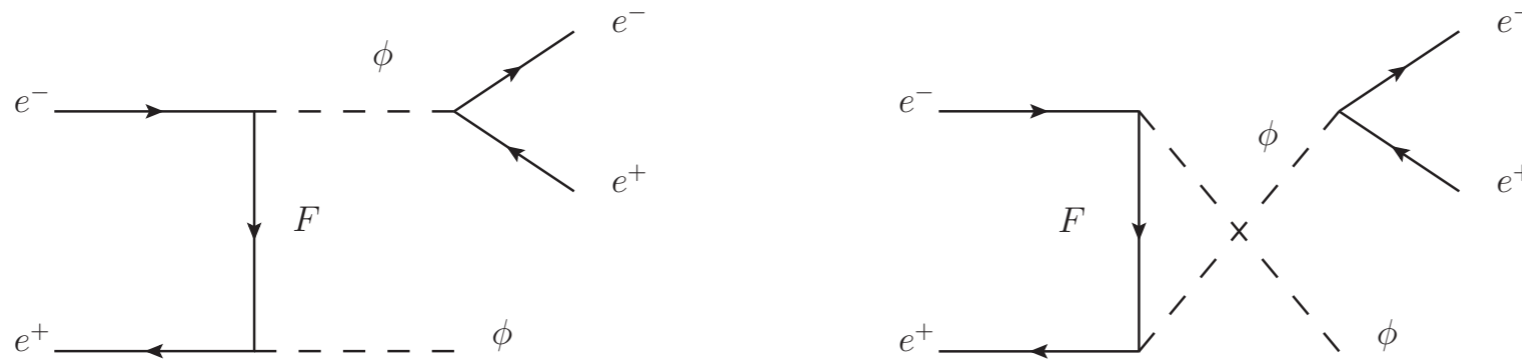
# LLP Signatures at Collider



- Long-lived Scalar Particle at CEPC

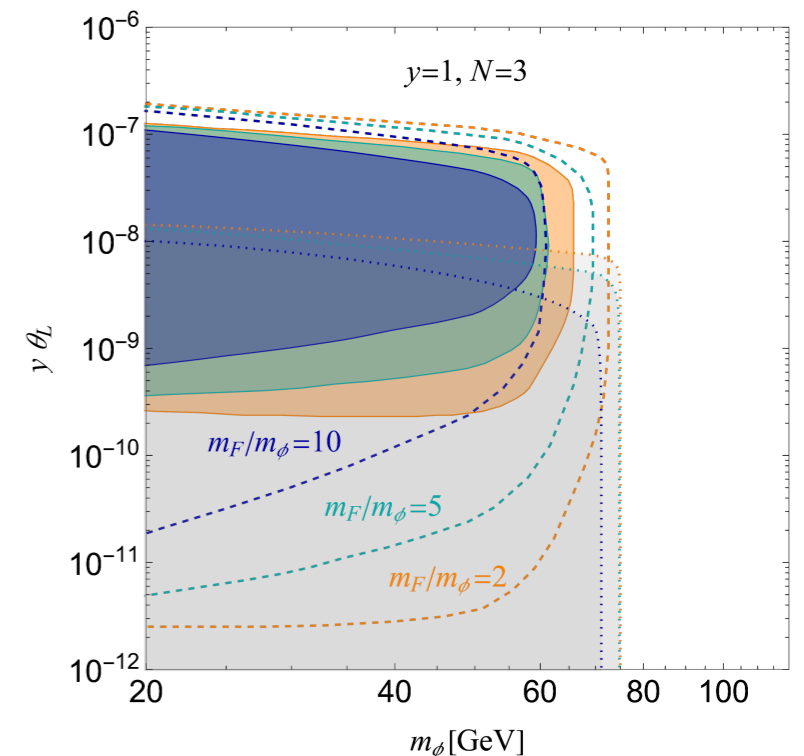
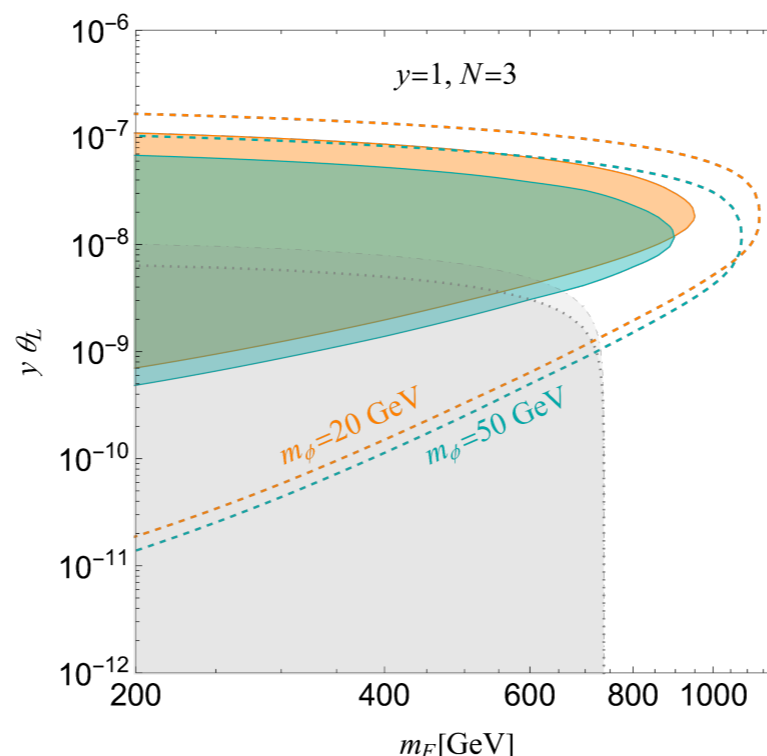
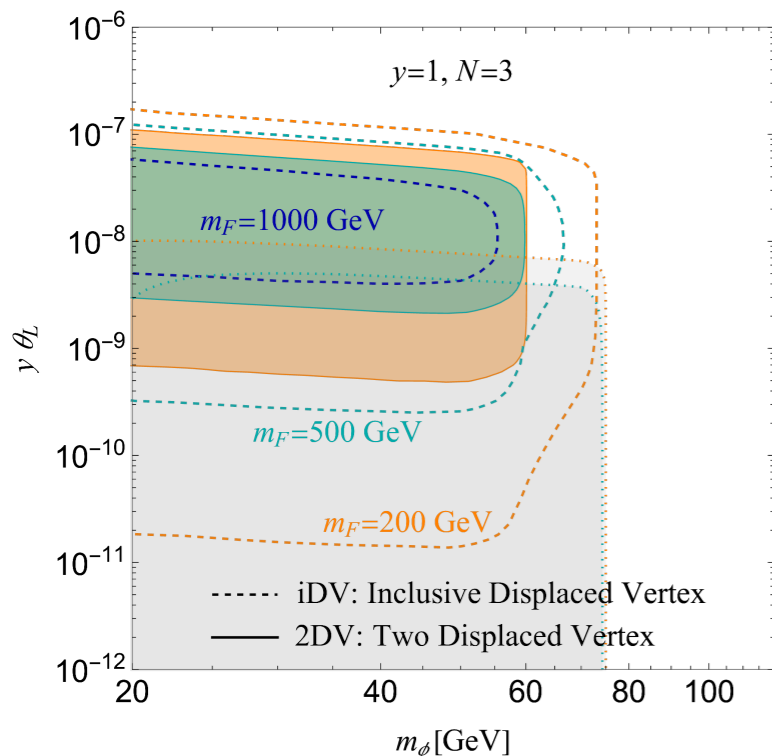
$F$  mixed with  $e$

(displaced vertex)



DV – CEPC :  $10 \text{ cm} < |d_\phi \cdot \sin \theta_i| < 1.8 \text{ m}$ ,  $|d_\phi \cdot \cos \theta_i| < 2.35 \text{ m}$ ,

$p_T^{e_i} > 30 \text{ GeV}$ ,  $m_{e_1 e_2}(m_\phi) > 20 \text{ GeV}$ ,  $1 > \Delta R > 0.01$



CEPC:  $\sqrt{s} = 240 \text{ GeV}$ ,  $\mathcal{L} = 5.6 \text{ ab}^{-1}$

# Summary



- We explore the potential long-lived signatures of VLLs  $F^\pm$  or their subsequent decay products  $\phi$ .
- For the long-lived  $F^\pm$ , we use the kink track method at the ATLAS detector, which shows good sensitivities for  $m_F \in [200, 950]$  GeV, with coupling around  $10^{-10} < \theta_L < 10^{-7}$ .
- For the long-lived scalars  $\phi$ , we use the displaced muon-jet and time-delay methods, which show good sensitivities for  $m_F \in [200, 1200]$  GeV, with coupling around  $10^{-11} < y\theta_L < 10^{-6}$ . Moreover, CEPC has good performance for  $m_\phi < 70$  GeV and  $m_F < 1000$  GeV through the displaced vertex method.

**Thanks!**

# References



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# Backup

- More on LFV

The unitary matrix which diagonalizing the general mixed mass matrix

$$U'_L \simeq \begin{pmatrix} 1 - \frac{1}{2}r_{1a}^2\theta_L^2 & x_1x_2r_{12}\theta_R^2 & x_1x_3r_{13}\theta_R^2 & -x_1r_{1a}\theta_L \\ -x_1x_2r_{12}\theta_R^2 & 1 - \frac{1}{2}r_{2a}^2\theta_L^2 & x_2x_3r_{23}\theta_R^2 & -x_2r_{2a}\theta_L \\ -x_1x_3r_{13}\theta_R^2 & -x_2x_3r_{23}\theta_R^2 & 1 - \frac{1}{2}r_{3a}^2\theta_L^2 & -x_3r_{3a}\theta_L \\ x_1r_{1a}\theta_L & x_2r_{2a}\theta_L & x_3r_{3a}\theta_L & 1 - \frac{1}{2}\theta_L^2(r_{1a}^2x_1^2 + r_{2a}^2x_2^2 + r_{3a}^2x_3^2) \end{pmatrix}$$

$$U'_R \simeq \begin{pmatrix} 1 - \frac{1}{2}\theta_R^2 & x_1x_2r_{12}\theta_R^2 & x_1x_3r_{13}\theta_R^2 & -x_1\theta_R \\ -x_1x_2\theta_R^2 & 1 - \frac{1}{2}\theta_R^2 & x_2x_3r_{23}\theta_R^2 & -x_2\theta_R \\ -x_1x_3\theta_R^2 & -x_2x_3\theta_R^2 & 1 - \frac{1}{2}\theta_R^2 & -x_3\theta_R \\ x_1\theta_R & x_2\theta_R & x_3\theta_R & 1 - \frac{1}{2}\theta_R^2(x_1^2 + x_2^2 + x_3^2) \end{pmatrix}$$

$$r_{ij} \equiv m_{\ell_i}/m_{\ell_j}$$

$$r_{ia} \equiv m_{\ell_i}/m_{\ell_a}$$



# Backup

- More on LFV

- If  $\phi$  interaction comes from flavor eigenstates

$$\mathcal{L}_{\text{eff}}^{\phi} = y\phi(\bar{e}_L^0, \bar{\mu}_L^0, \bar{\tau}_L^0, \bar{F}_L^0) \begin{pmatrix} 0000 \\ 0000 \\ 0000 \\ 0100 \end{pmatrix} \begin{pmatrix} e_R^0 \\ \mu_R^0 \\ \tau_R^0 \\ F_R^0 \end{pmatrix} + \text{h.c.} = y\phi(\bar{e}'_L, \bar{\mu}'_L, \bar{\tau}'_L, \bar{F}'_L) \begin{pmatrix} 0000 \\ 0000 \\ 0000 \\ x_1x_2x_30 \end{pmatrix} \begin{pmatrix} e'_R \\ \mu'_R \\ \tau'_R \\ F'_R \end{pmatrix} + \text{h.c.}..$$

leads to

$$\text{Br}(\ell_i \rightarrow \ell_j \gamma) \simeq \frac{\pi^3 \alpha_{\text{EM}} (y^2 x_i x_j)^2}{3G_F^2} \frac{\left( m_F^6 - 6m_F^4 m_\phi^2 + 3m_F^2 m_\phi^4 + 2m_\phi^6 + 12m_F^2 m_\phi^4 \ln \left( \frac{m_F}{m_\phi} \right) \right)^2}{(m_F^2 - m_\phi^2)^8}$$

$$y < 6.8 \times 10^{-4} \left( \frac{1}{x_1} \right)^{1/2} \left( \frac{1}{x_2} \right)^{1/2} \left( \frac{m_F}{200 \text{ GeV}} \right)$$





# Backup

- More on LFV

- If  $\phi$  interaction comes from SM mass eigenstates

$$\text{Br}(\ell_a \rightarrow \ell_i \gamma) \simeq \frac{\pi^3 \alpha_{\text{EM}} (x_i x_a r_{ia} y^2 \theta_R^2)^2}{3G_F^2} \frac{\left( m_F^6 - 6m_F^4 m_\phi^2 + 3m_F^2 m_\phi^4 + 2m_\phi^6 + 12m_F^2 m_\phi^4 \ln \left( \frac{m_F}{m_\phi} \right) \right)^2}{\left( m_F^2 - m_\phi^2 \right)^8}$$
$$\simeq 3.5 \times 10^{-15} \left( \frac{x_i}{1} \right)^2 \left( \frac{x_a}{1} \right)^2 \left( \frac{r_{ia}}{10^{-3}} \right)^2 \left( \frac{y \theta_R}{10^{-2}} \right)^4 \left( \frac{200 \text{ GeV}}{m_F} \right)^4.$$