# Long-lived dark photons at the LHC



## 第十七届TeV工作组学术研讨会

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- Enhanced Long-Lived Dark Photon Signals at the LHC Mingxuan Du, ZL, Van Que Tran JHEP 05 (2020) 055
   e-Print: 1912.00422 [hep-ph]
- Enhanced long-lived dark photon signals at lifetime frontier detectors Mingxuan Du, Rundong Fang, ZL, Van Que Tran *Phys.Rev.D* 105 (2022) 5, 055012
   e-Print: 2111.15503 [hep-ph]
- FACET: A new long-lived particle detector in the very forward region of the CMS experiment S. Cerci et al., JHEP 06 (2022) 110 • e-Print: 2201.00019 [hep-ex]





# Dark photon models & constraints Our new (long-lived) dark photon models Long-lived dark photon signals at the LHC



### Dark photon models & constraints



### Our new (long-lived) dark photon models

### Long-lived dark photon signals @ the LHC

# Dark photon models & constraints





# Hypercharge portal models $\Longrightarrow$ dark photon





### dark sector

### [Holdom 1986] [Foot & He 1991] [Kors & Nath 2004] [Feldman, ZL, Nath, <u>hep-ph/0702123</u>, 372 cites]





# Hypercharge portal models $\Longrightarrow$ dark photon



[Holdom 1986] [Foot & He 1991] [Kors & Nath 2004] [Feldman, ZL, Nath, <u>hep-ph/0702123</u>, 372 cites]





## $SU(3)_{c} \times SU(2)_{L} \times U(1)_{V} \times U(1)_{V}$

 $\mathscr{L} = -\frac{1}{4}B_{\mu\nu}B^{\mu\nu} - \frac{1}{4}X_{\mu\nu}X^{\mu\nu} + g_D X_\mu \bar{\chi}\gamma^\mu \chi - \frac{\delta}{2}B_{\mu\nu}X^{\mu\nu} - \frac{M_1^2}{2}(\partial_\mu \sigma + X_\mu + \tilde{\epsilon} B_\mu)^2$ 



6
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# $SU(3)_{c} \times SU(2)_{L} \times U(1)_{V} \times U(1)_{X}$

 $\mathscr{L} = -\frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4} X_{\mu\nu} X^{\mu\nu} + g_D X_\mu \bar{\chi} \gamma^\mu \chi - \frac{\delta}{2} B_{\mu\nu} X^{\mu\nu} - \frac{M_1^2}{2} (\partial_\mu \sigma + X_\mu + \tilde{\epsilon} B_\mu)^2$ 





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# $SU(3)_{c} \times SU(2)_{L} \times U(1)_{V} \times U(1)_{X}$

 $\mathscr{L} = -\frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4} X_{\mu\nu} X^{\mu\nu} + g_D X_\mu \bar{\chi} \gamma^\mu \chi - \frac{\tilde{\delta}}{2} B_{\mu\nu} X^{\mu\nu} - \frac{M_1^2}{2} (\partial_\mu \sigma + X_\mu + \tilde{\epsilon} B_\mu)^2$ mass mixing kinetic mixing







# $SU(3)_c \times SU(2)_L \times U(1)_Y \times U(1)_X$

[Feldman, ZL, Nath, <u>hep-ph/0702123</u>, **372** cites]

 $\mathscr{L} = -\frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4} X_{\mu\nu} X^{\mu\nu} + g_D X_\mu \bar{\chi} \gamma^\mu \chi - \frac{\tilde{\delta}}{2} B_{\mu\nu} X^{\mu\nu} - \frac{M_1^2}{2} (\partial_\mu \sigma + X_\mu + \tilde{\epsilon} B_\mu)^2$ mass mixing kinetic mixing

kinetic mixing  $\delta \delta$  mass mixing  $\tilde{\epsilon}$  are degenerate (w/o  $\chi$ ): only  $\epsilon \sim (\tilde{\epsilon} - \delta)$  is physical







[Feldman, ZL, Nath, <u>hep-ph/0702123</u>, **372** cites] [see also Fabbrichesi+, 2005.01515, Dark Photon Review]





$$\mathscr{L} = -\frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4} X_{\mu\nu} X^{\mu\nu} + g_D X_\mu \bar{\chi} \gamma^\mu \chi - \frac{\tilde{\delta}}{2} B_{\mu\nu} X^{\mu\nu} - \frac{M_1^2}{2} (\partial_\mu \sigma + X_\mu + \tilde{\epsilon} B_\mu)^2$$

•  $X_{\mu} \Longrightarrow A'_{\mu}$  (dark photon), if  $M_1 \ll M_Z$  $\epsilon e Q_f A'_{\mu} \bar{f} \gamma^{\mu} f$  (SM sector) and  $g_D A'_{\mu} \bar{\chi} \gamma^{\mu} \chi$  (dark sector)



### [Feldman, ZL, Nath, <u>hep-ph/0702123</u>, 372 cites]

### [see also Fabbrichesi+, 2005.01515, Dark Photon Review]



$$\mathscr{L} = -\frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4} X_{\mu\nu} X^{\mu\nu} + g_D X_\mu \bar{\chi} \gamma^\mu \chi - \frac{\tilde{\delta}}{2} B_{\mu\nu} X^{\mu\nu} - \frac{M_1^2}{2} (\partial_\mu \sigma + X_\mu + \tilde{\epsilon} B_\mu)^2$$

- $X_{\mu} \Longrightarrow A'_{\mu}$  (dark photon), if  $M_1$  $\epsilon e Q_f A'_{\mu} \bar{f} \gamma^{\mu} f$  (SM sector) and  $g_D A'_{\mu} \bar{\chi} \gamma^{\mu} \chi$  (dark sector)
- $X_{\mu} \Longrightarrow Z'_{\mu}$  (hypercharge-like), if  $M_1 \gg M_Z$

$$\ll M_Z$$

### [Feldman, ZL, Nath, <u>hep-ph/0702123</u>, **372** cites]

### [see also Fabbrichesi+, 2005.01515, Dark Photon Review]



$$\mathscr{L} = -\frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4} X_{\mu\nu} X^{\mu\nu} + g_D X_\mu \bar{\chi} \gamma^\mu \chi - \frac{\tilde{\delta}}{2} B_{\mu\nu} X^{\mu\nu} - \frac{M_1^2}{2} (\partial_\mu \sigma + X_\mu + \tilde{\epsilon} B_\mu)^2$$

- $X_{\mu} \Longrightarrow A'_{\mu}$  (dark photon), if  $M_1$
- $X_{\mu} \Longrightarrow Z'_{\mu}$  (hypercharge-like), if  $M_1 \gg M_Z$

If  $A'_{\mu}$  or  $Z'_{\mu}$  is massive,  $\chi$  is millicharged ( $\epsilon e A_{\mu} \bar{\chi} \gamma^{\mu} \chi$ ) only when  $\tilde{\epsilon} \neq 0$ 

$$\ll M_Z$$

## $\epsilon e Q_f A'_{\mu} \bar{f} \gamma^{\mu} f$ (SM sector) and $g_D A'_{\mu} \bar{\chi} \gamma^{\mu} \chi$ (dark sector)

### [Feldman, ZL, Nath, <u>hep-ph/0702123</u>, **372** cites] [see also Fabbrichesi+, 2005.01515, Dark Photon Review]



# Limits on dark photon with mass below 1 MeV



 $\epsilon e Q_f A'_\mu \bar{f} \gamma^\mu f$ 

### astro/cosmo probes



# Limits on dark photon with mass above 1 MeV



### accelerator probes



# Limits on dark photon with mass above 1 MeV









### dark photon interaction with SM particles

 $\epsilon e Q_f A'_\mu \bar{f} \gamma^\mu f$ 









long decay length  $\implies$  small coupling



### Dark photon models & constraints

### Our new (long-lived) dark photon models 2

### Our new (long-lived) dark photon models

### Long-lived dark photon signals @ the LHC







# Long-lived dark photon signals

LLDP  $\implies$  small coupling  $\implies$  suppressed collider signals

[1503.06770]





# Long-lived dark photon signals

- LLDP  $\implies$  small coupling  $\implies$  suppressed collider signals
- Go beyond the simple one U(1) picture BSM theories can predict multiple U(1)'s, e.g. SO(32) string theory [1503.06770]
- $SU(3)_{c} \times SU(2)_{L} \times U(1)_{V} \times U(1)_{X} \times U(1)_{C}$ Extend SM with 2 U(1)'s





# Long-lived dark photon signals

- LLDP  $\implies$  small coupling  $\implies$  suppressed collider signals
- Go beyond the simple one U(1) picture BSM theories can predict multiple U(1)'s, e.g. SO(32) string theory [1503.06770]
- $SU(3)_{c} \times SU(2)_{I} \times U(1)_{V} \times U(1)_{X} \times U(1)_{C}$ Extend SM with 2 U(1)'s

Dirac fermion  $\psi$  charged under both U(1)'s  $(g_F X_{\mu} + g_W C_{\mu}) \bar{\psi} \gamma^{\mu} \psi$ 





## $SU(3)_c \times SU(2)_L \times U(1)_Y \times U(1)_X \times U(1)_C$



 $\mathscr{L}_{F} = -\frac{1}{4} \frac{X_{\mu\nu}^{2}}{2} - \frac{1}{2} (\partial_{\mu}\sigma_{1} + m_{1}\epsilon_{1}B_{\mu} + m_{1}X_{\mu})^{2}$  $SU(3)_{c} \times SU(2)_{L} \times U(1)_{V} \times U(1)_{X} \times U(1)_{C}$ 



 $\mathscr{L}_{F} = -\frac{1}{A} \frac{X_{\mu\nu}^{2}}{2} - \frac{1}{2} (\partial_{\mu}\sigma_{1} + m_{1}\epsilon_{1}B_{\mu} + m_{1}X_{\mu})^{2}$  $SU(3)_{c} \times SU(2)_{L} \times U(1)_{V} \times U(1)_{X} \times U(1)_{C}$  $\mathscr{L}_{W} = -\frac{1}{4}C_{\mu\nu}^{2} - \frac{1}{2}(\partial_{\mu}\sigma_{2} + m_{2}\epsilon_{2}B_{\mu} + m_{2}C_{\mu})^{2}$ 



 $\mathscr{L}_F = -\frac{1}{\Delta} \frac{X_{\mu\nu}^2}{2} - \frac{1}{2} (\partial_\mu \sigma_1 + m_1 \epsilon_1 B_\mu + m_1 X_\mu)^2$  $SU(3)_{c} \times SU(2)_{L} \times U(1)_{V} \times U(1)_{X} \times U(1)_{C}$  $\mathscr{L}_W = -\frac{1}{4}C_{\mu\nu}^2 - \frac{1}{2}(\partial_\mu\sigma_2 + m_2\epsilon_2B_\mu)$ 



$$+m_2 C_{\mu})^2$$





 $\mathscr{L}_F = -\frac{1}{\Delta} \frac{X_{\mu\nu}^2}{2} - \frac{1}{2} (\partial_\mu \sigma_1 + m_1 \epsilon_1 B_\mu + m_1 X_\mu)^2$  $SU(3)_{c} \times SU(2)_{L} \times U(1)_{V} \times U(1)_{X} \times U(1)_{C}$  $\mathscr{L}_{W} = -\frac{1}{4}C_{\mu\nu}^{2} - \frac{1}{2}(\partial_{\mu}\sigma_{2} + m_{2}\epsilon_{2}B_{\mu} + m_{2}C_{\mu})^{2}$ 





$$m^{2} = \begin{pmatrix} m_{2}^{2} & 0 \\ 0 & m_{1}^{2} \\ m_{2}^{2}\epsilon_{2} & m_{1}^{2}\epsilon_{1} & m \\ 0 & 0 \end{pmatrix}$$

### [Du, ZL, Tran, 1912.00422]

4 by 4 mass square matrix in the basis of  $V = (C, X, B, A^3)$ 

$m_2^2\epsilon_2$	0
$m_1^2 \epsilon_1$	0
$m_1^2 \epsilon_1^2 + m_2^2 \epsilon_2^2 + \frac{{g'}^2 v^2}{4}$	$\frac{g'gv^2}{4}$
$\frac{g'gv^2}{4}$	$\frac{g^2v^2}{4}$

14

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$\frac{g'gv^2}{4}$	$\frac{g^2v^2}{4}$

mass eigenstates: E = (Z', A', Z, A) via  $V_i = O_{ij}E_j$ 

14

$$m^{2} = \begin{pmatrix} m_{2}^{2} & 0 \\ 0 & m_{1}^{2} \\ m_{2}^{2}\epsilon_{2} & m_{1}^{2}\epsilon_{1} & m \\ 0 & 0 \end{pmatrix}$$

determinant =  $0 \implies$  massless photon

[Du, ZL, Tran, 1912.00422]

4 by 4 mass square matrix in the basis of  $V = (C, X, B, A^3)$ 



mass eigenstates: E = (Z', A', Z, A) via  $V_i = O_{ij}E_j$ 

14

$$m^{2} = \begin{pmatrix} m_{2}^{2} & 0 \\ 0 & m_{1}^{2} \\ m_{2}^{2}\epsilon_{2} & m_{1}^{2}\epsilon_{1} & m \\ 0 & 0 \end{pmatrix}$$

- - $\epsilon_1 = 0 = \epsilon_2 \Longrightarrow \text{decouple}$

[Du, ZL, Tran, 1912.00422]

4 by 4 mass square matrix in the basis of  $V = (C, X, B, A^3)$ 



mass eigenstates: E = (Z', A', Z, A) via  $V_i = O_{ij}E_j$ 

determinant =  $0 \implies$  massless photon

14

# vector and axial-vector couplings of neutral bosons

vector and axial-vector couplings between bosons & fermions

 $\bar{f}\gamma_{\mu}(v_{i}^{f}-\gamma_{5}a_{i}^{f})fE_{i}^{\mu}+v_{i}^{\psi}\bar{\psi}\gamma_{\mu}\psi E_{i}^{\mu}$ 





# vector and axial-vector couplings of neutral bosons

vector and axial-vector couplings between bosons & fermions

 $\bar{f}\gamma_{\mu}(v_{i}^{f}-\gamma_{5}a_{i}^{f})fE_{i}^{\mu}+v_{i}^{\psi}\bar{\psi}\gamma_{\mu}\psi E_{i}^{\mu}$  $v_i^f = (g\mathcal{O}_{4i} - g'\mathcal{O}_{3i})T_f^3/2 + g'\mathcal{O}_{3i}Q_f$  $a_i^f = (g\mathcal{O}_{4i} - g'\mathcal{O}_{3i})T_f^3/2$ 

SM fermion v

SM fermion a





# vector and axial-vector couplings of neutral bosons

vector and axial-vector couplings between bosons & fermions

 $\bar{f}\gamma_{\mu}(v_{i}^{f}-\gamma_{5}a_{i}^{f})fE_{i}^{\mu}+v_{i}^{\psi}\bar{\psi}\gamma_{\mu}\psi E_{i}^{\mu}$  $v_i^f = (g\mathcal{O}_{4i} - g'\mathcal{O}_{3i})T_f^3/2 + g'\mathcal{O}_{3i}Q_f$  $a_i^f = (g\mathcal{O}_{4i} - g'\mathcal{O}_{3i})T_f^3/2$  $v_i^{\psi} = g_W \mathcal{O}_{1i} + g_F \mathcal{O}_{2i}$ 

SM fermion v

SM fermion a

dark sector









# hidden radiation channel







# millicharged dark matter constraints



[see. e.g. Kovetz+ 1807.11482, Boddy+ 1808.00001, Puttter+ 1805.11616]

<sup>1</sup> millicharged DM abundance < 0.4%





# experimental constraints



More recent constraints on millicharged particles from SENSEI, BEBC, Super-K etc [2305.04964] [2011.08153]

[2211.11469]



### Dark photon models & constraints

### Long-lived dark photon signals @ the LHC 3

### Our new (long-lived) dark photon models

### Long-lived dark photon signals @ the LHC





# lifetime frontier detectors

### 3 types detectors: (1) far forward, (2) far transverse, (3) near timing

Detector	η	Distance from IP (m)	Decay volume $(m^3)$	LHC ru
FACET [3,4]	[6, 7.2]	100 (upstream)	12.3	Run 4 (2
FASER [5–9]	>9	480 (downstream)	0.047	Run 3 (2
FASER2 [9,10]	>6.87	480 (downstream)	15.7	HL-LH
AL3X [11]	[0.9, 3.7]	5.25 (upstream)	915.2	Run 5 (2
MoEDAL-MAPP [12]	~3.1	55 (upstream)	~150	Run 3 (2
CODEX-b [18,19]	[0.14, 0.55]	26 (transverse)	10 <sup>3</sup>	Run 4 (2
MATHUSLA [13–17]	[0.64, 1.43]	60 (transverse)	$2.5 \times 10^{5}$	HL-LH
ANUBIS [20]	[0.06, 0.21]	24 (transverse)	$\sim 1.3 \times 10^4$	HL-LH
CMS-MTD [21]	[-3,3]	1.17 (barrel), 3.04 (endcaps)	25.4	HL-LH
ATLAS-HGTD [22]	[2.4, 4]	3.5 (endcaps)	8.7	HL-LH
LHCb-TORCH [23,24]	[1.6, 4.9]	9.5 (beam direction)	• • •	HL-LH





# Timing detectors at ATLAS/CMS/LHCb

CMS-TMD
ATLAS-HGTD
HC-TORCH







[https://cds.cern.ch/record/2296612/files/LHCC-P-009.pdf]

## between tracker & calorimeter δt = 30 ps; pileup reduction & LLP





# Time delay for NR long-lived particles

D





$$\Delta t = \frac{\ell_X}{\beta_X} + \frac{\ell_a}{\beta_a} - \frac{\ell_{\rm SM}}{\beta_{\rm SM}}$$

[Liu, Liu, Wang, 1805.05957]



# Transverse detectors

CODEX-b
MATHUSLA
ANUBIS



# MATHUSLA



[MATHUSLA, 2009.01693]



# Far detectors

# FASER and FASER2 FACET



# FASER





[see 2308.05587 for FASER's new results]

### [FASER 1901.00468]









# FACET: a new long-lived particle detector in the very forward region of the CMS experiment

- S. Cerci,<sup>*a*,1</sup> D. Sunar Cerci,<sup>*a*,1</sup> D. Lazic,<sup>*b*</sup> G. Landsberg,<sup>*c*,2</sup> F. Cerutti,<sup>*d*</sup>

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## 100 meter from CMS R = 50 cm pipe (vacuum)

### [FACET, JHEP, 2201.00019]

[see Albrow's talk for variants]



# Sensitivities from lifetime frontier detectors

### Minimal DP









distance d



### volume V





## probability of decaying at distance L: (suppressed at large L, but NOT at small L)

![](_page_56_Picture_4.jpeg)

![](_page_57_Figure_1.jpeg)

probability of decaying at distance L: (suppressed at large L, but NOT at small L)

ratio between 2 far detectors (assuming isotropic)

![](_page_57_Picture_4.jpeg)

![](_page_57_Picture_6.jpeg)

![](_page_58_Figure_1.jpeg)

probability of decaying at distance L: (suppressed at large L, but NOT at small L)

ratio between 2 far detectors (assuming isotropic)

larger & closer to IP is better (geometrical, w/o BG)

![](_page_58_Picture_5.jpeg)

- Minimal dark photon models have suppressed long-lived dark photon signals at the LHC
- We build a new model by extending the SM with two U(1) gauge fields in the dark sector, where the dark photon is produced via another gauge boson and is not suppressed
- We study a new dark photon production channel at the LHC: hidden radiation, in additional to the proton bremsstrahlung and meson decay channels
- Far detectors (FASER & FACET) and timing detectors (CMS, ATLAS, LHCb) are sensitive to long-lived dark photons with small and large masses, respectively

![](_page_59_Picture_5.jpeg)

![](_page_59_Picture_6.jpeg)

# additional slides

![](_page_60_Picture_1.jpeg)

![](_page_61_Figure_1.jpeg)

### probability of decaying at distance L:

probability of decaying in (d, d + h):

LLDP decays inside volume V

# LLDP decays inside a far detector

![](_page_61_Picture_6.jpeg)

![](_page_61_Picture_7.jpeg)

# **High-Granularity Timing Detector**

- ATLAS upgrade detector for the high luminosity LHC
- uses LGAD sensors to measure time with  $\sigma_t \sim 30-50$  ps per track until end of HL-LHC
- covers range  $2.4 \leq |\eta| \leq 4.0$
- two disks positioned at  $z = \pm 3.5m$  from the interaction point

![](_page_62_Picture_5.jpeg)

![](_page_62_Picture_6.jpeg)

# [taken from Leopold's talk] 2,2m 12,5 cm

![](_page_62_Picture_9.jpeg)

![](_page_62_Picture_10.jpeg)

![](_page_62_Picture_11.jpeg)

# LHCb-TORCH

![](_page_63_Figure_1.jpeg)

Figure 4.9: Schematic of TORCH detector for LHCb: (a) Front-on view of full detector; (b) View of single module showing focusing block and photodetector plane.

The precision of each track in the TORCH system is 15 ps

### [CERN-LHCC-2017-003]

![](_page_63_Figure_5.jpeg)

![](_page_63_Picture_6.jpeg)

### Mike Albrow 20231025

### **FACET FUTURE OPTIONS**

FACET – A (baseline) with *large* pipe as JHEP paper but with beam pipe optimized for background reduction:

![](_page_64_Figure_3.jpeg)

(Too late for Run4: Plan for Run 5 subject to feasibility from more studies) FACET – B (variant) with *standard* pipe, detectors as FACET-A (full azimuth): (Possible for Run4 – but would need modifications to LHC, e.g pipe supports) FACET – C (variant) *above standard* pipe : (smaller acceptance, no mods to LHC, easiest option. Area shown nominal example)

![](_page_64_Figure_5.jpeg)

![](_page_64_Figure_6.jpeg)

Contributions mainly: Mike Albrow Deniz Sunar Cerci Salim Cerci Suat Ozkorucuklu Aldo Penzo Burak Hacisahinoglu Orhan Aydilek

# FACET

### [Albrow's talk]

# TRACK CALO

![](_page_64_Picture_12.jpeg)

![](_page_64_Picture_13.jpeg)

![](_page_65_Figure_0.jpeg)

# FASER

### [2308.05587]

![](_page_65_Picture_3.jpeg)

![](_page_65_Picture_4.jpeg)