



Interplay between Neutrino & Dark Matter

Shao-Feng Ge

葛韶锋

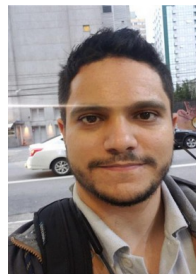
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紫金山暗物质研讨会
December 30, 2023

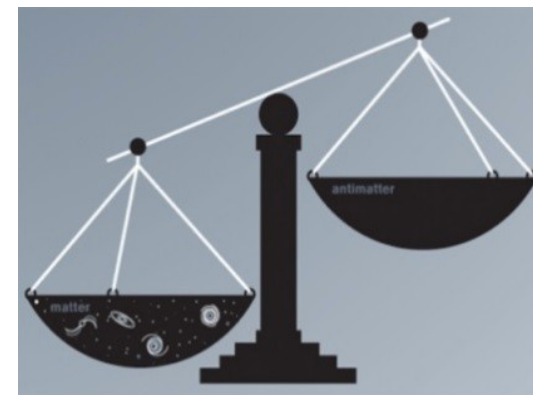
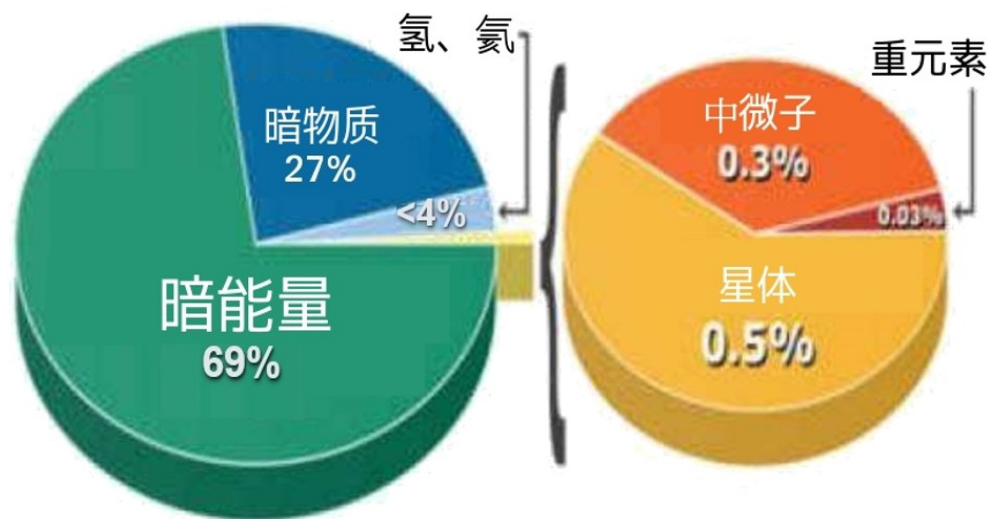
李政道研究所

Tsung-Dao Lee Institute

- 1) Overview
- 2) Sterile ν DM
- 3) Heavy Right-Handed Neutrino DM
- 4) Absorption DM @ Direct Detection Experiments
- 5) Environmental Mass with Dark NSI
- 6) ν Mass with Cosmic Gravitational Focusing
- 7) Summary

暗物质是什么？怎么产生的？

暗物质主导了物质世界的形成



可见物质的起源？
轻子生成机制

两者之间存在怎样的关联？

Georg G. Raffelt

Stars as Laboratories for Fundamental Physics

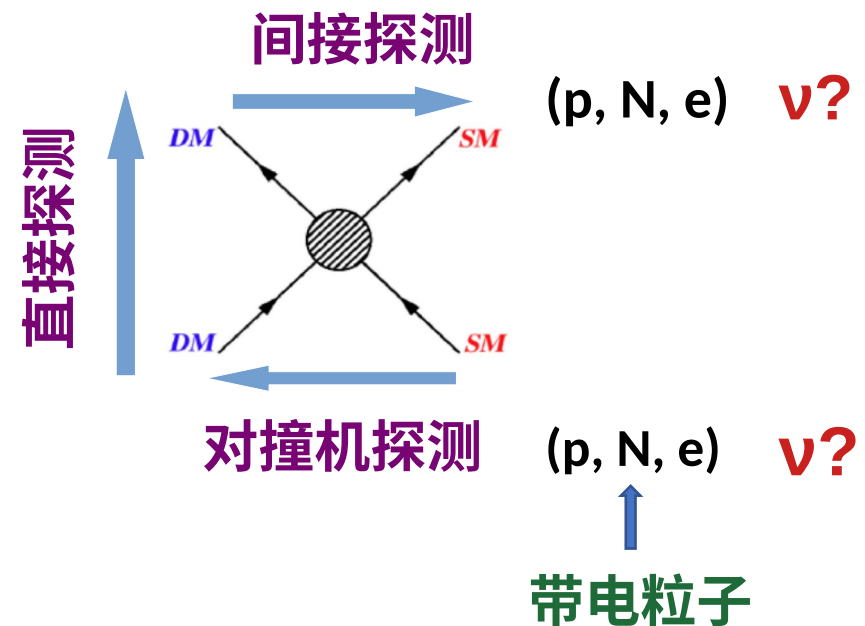
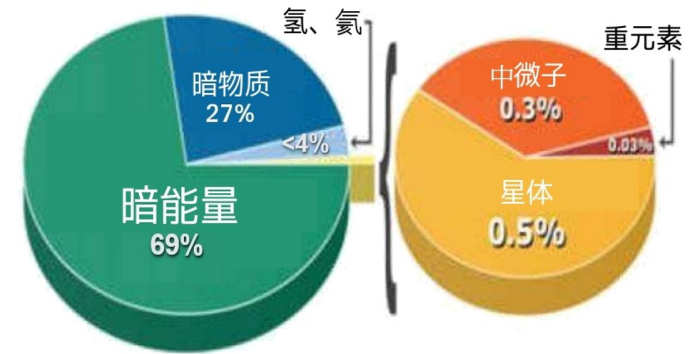
The Astrophysics of Neutrinos, Axions, and Other
Weakly Interacting Particles

In the standard model, neutrinos have been assigned the most minimal properties compatible with experimental data: zero mass, zero charge, zero dipole moments, zero decay rate, zero almost everything.

Almost the same for dark matter!

Neutrino vs Dark Matter

- 都对理解宇宙中的物质世界起关键作用
- 都是电中性物质
- 都和普通物质的相互作用很弱
- 都在宇宙早期就脱耦
- 背后是否有一个全新的分区?
- 两者之间是否有关联?
-



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Cosmological Lower Bound on Heavy-Neutrino Masses

Benjamin W. Lee^(a)

Fermi National Accelerator Laboratory,^(b) Batavia, Illinois 60510

and

Steven Weinberg^(c)

Stanford University, Physics Department, Stanford, California 94305

(Received 13 May 1977)

The present cosmic mass density of possible stable neutral heavy leptons is calculated in a standard cosmological model. In order for this density not to exceed the upper limit of 2×10^{-29} g/cm³, the lepton mass would have to be *greater* than a lower bound of the order of 2 GeV.

Lee & Weinberg, PRL77

Are Neutrinos Dark Matter?

Cosmological Upper Bound on Heavy-Neutrino Lifetimes

Duane A. Dicus^(a) and Edward W. Kolb^(a)

Center for Particle Theory, University of Texas at Austin, Austin, Texas 78712

and

Duane Dicus, Kolb & Teplitz, PRL77

Vigdor L. Teplitz^(b)

Department of Physics, Virginia Polytechnic Institute and State University, Blacksburg, Virginia 24061

(Received 31 May 1977)

An upper bound on the lifetime of a massive, neutral, weakly interacting lepton, ν_H , is derived from standard big-bang cosmology. Saturation of the bound and reasonable assumptions about the weak interaction of the ν_H then yield a prediction of approximately 10 MeV for its mass.

和 Lee-Weinberg 文章前后页

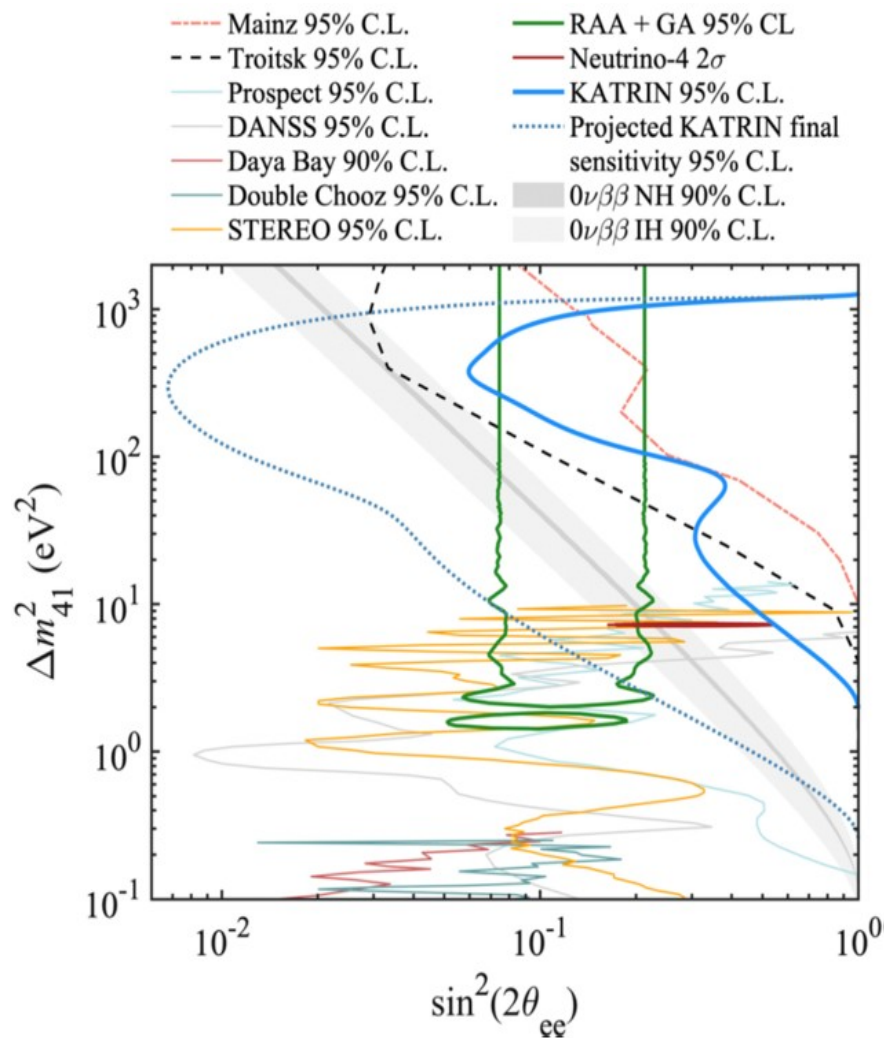


Duane Alfred Dicus

November 23, 1938– September 7, 2022

University of Texas

Sterile Neutrino

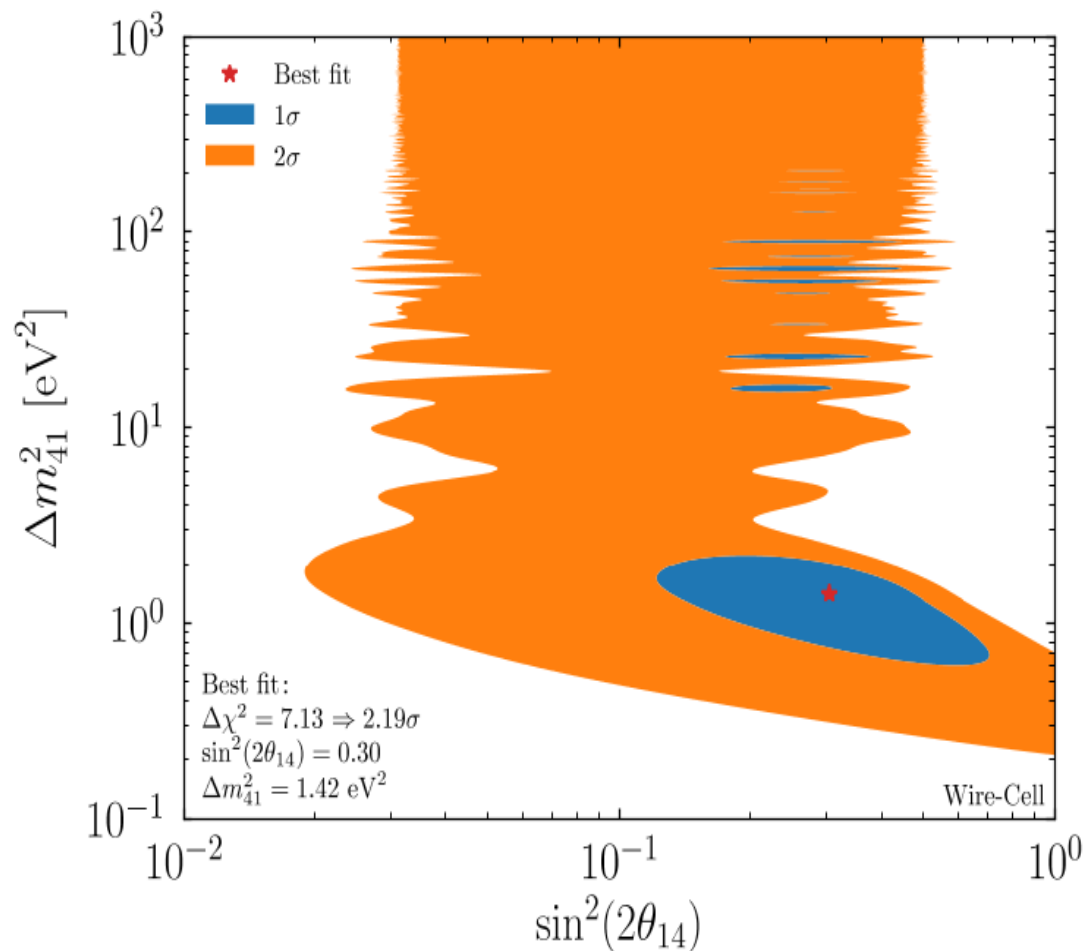


Schoppmann [2109.13541]

主要限制 eV sterile neutrino

Denton [2111.05793]

Latest result @ MicroBooNE



- If only neutrino oscillation => oscillation back & forth

$$\sin^2 2\theta \sin^2 \left(\delta m^2 \frac{L}{4E} \right) \quad \Rightarrow \quad \overline{P_{as}} = \frac{1}{2} \sin^2 2\theta_{as}$$

- Neutrino collision with nuclei = Quantum measurement

$$\left(\frac{\partial}{\partial t} - H E \frac{\partial}{\partial E} \right) f_s(E, t) = \left[\frac{1}{2} \sin^2(2\theta_M(E, t)) \Gamma(E, t) \right] f_a(E, t)$$

Hubble constant H denotes the influence of cosmological expansion.

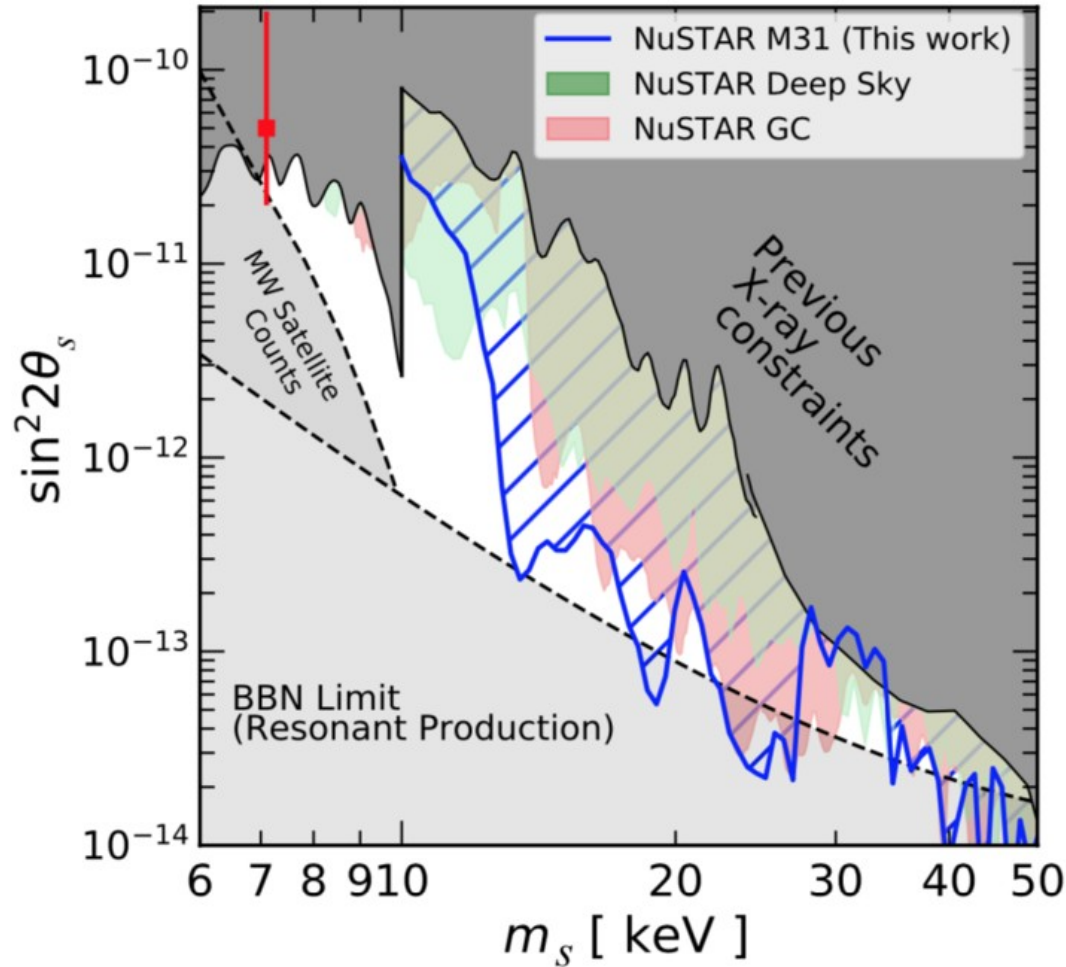
$f_s(E, t)$ and $f_a(E, t)$ are the time-dependent momentum distribution functions

$$f_a(E, t) = \frac{1}{e^{E/T} + 1}$$

$$\Gamma(E, t) \simeq \frac{7\pi}{24} G_F^2 T^4 E$$

不断转化为惰性中微子

active neutrino interaction rate

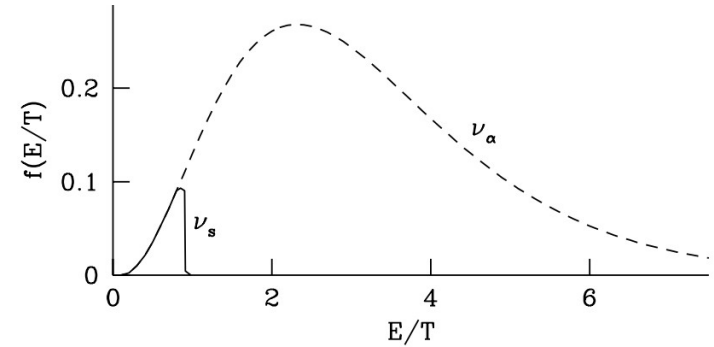


Ng, Roach, Perez, Beacom, Horiuchi, Krivonos, Wik [1901.01262]

- Resonant production

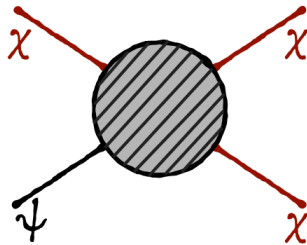
$$\sin^2 2\theta_{as}^m = \frac{\sin^2 2\theta_{as}}{\sin^2 2\theta_{as} + \left(\cos 2\theta_{as} \mp \frac{E_\nu V}{\Delta m_{as}^2} \right)^2}$$

Shi & Fuller, PRL99 [astro-ph/9810076]



- Pandemic Mechanism

Bringmann, Depta, Hufnagel, Ruderman, Schmidt-Hoberg [2103.16572]



$$n_\chi(t) \sim \frac{1}{a^3} e^{\langle \sigma v \rangle_{\text{tr}} n_\psi^{\text{eq}} t}$$

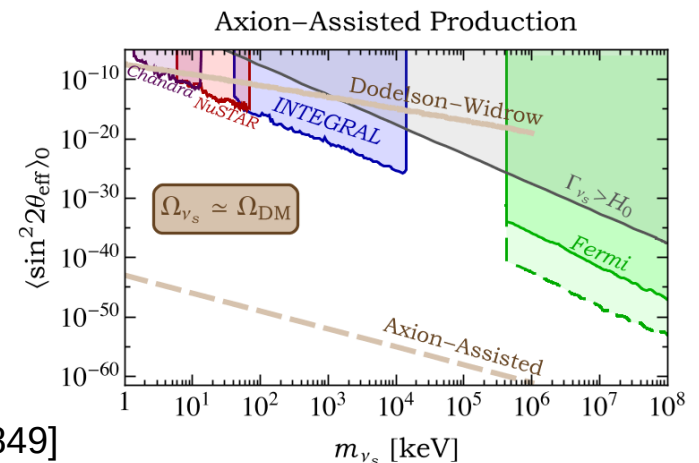
Bringmann, Depta, Hufnagel, Kersten, Ruderman, Schmidt-Hoberg [2206.10630]

- Axion Assisted Production

$$\sin^2 2\theta_m \simeq \frac{\Delta^2 \langle \sin^2 2\theta_{\text{eff}} \rangle}{\Delta^2 \langle \sin^2 2\theta_{\text{eff}} \rangle + (\Delta - V_T)^2}$$

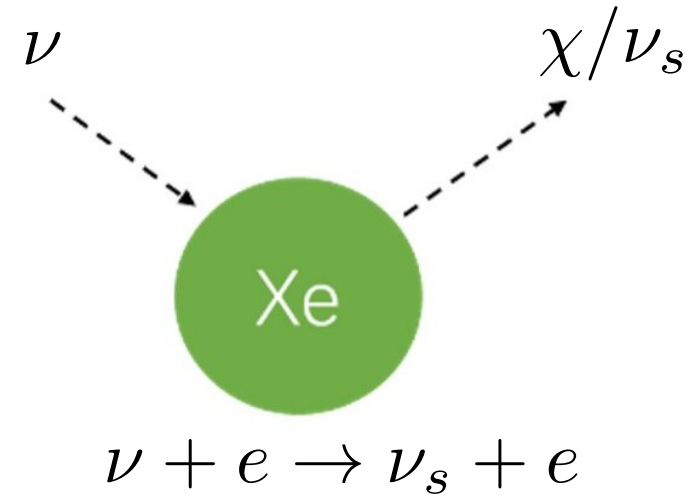
$$\Delta \equiv m_{\nu_s}^2 / 2E_\nu$$

Berlin & Hooper [1610.03849]



Recoil matrix element $|M|^2$ with a light mediator

$$(2m_e T_r + m_s^2) \frac{(y_S^\nu y_S^e)^2 (2m_e + T_r) + (y_P^\nu y_P^e)^2 T_r}{8\pi E_\nu^2 (2m_e T_r + m_\phi^2)^2}$$



- Active ν final state

$$m_s \rightarrow 0 \quad m_\phi \rightarrow 0$$

$$|\mathcal{M}|^2 \sim \begin{cases} \frac{1}{T_r^2} & \text{Scalar (S)} \\ \text{const} & \text{Pseudo-scalar (P)} \end{cases}$$

Pseudo-scalar mediator is also enhanced at low energy.

- Active-Sterile Conversion

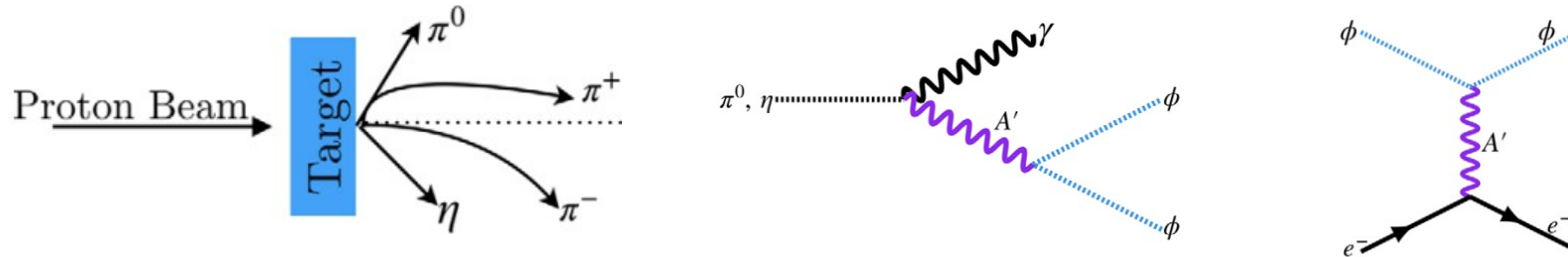
$$m_s \gtrsim \sqrt{2m_e T_r}$$

$$|\mathcal{M}|^2 \sim \frac{1}{T_r}$$

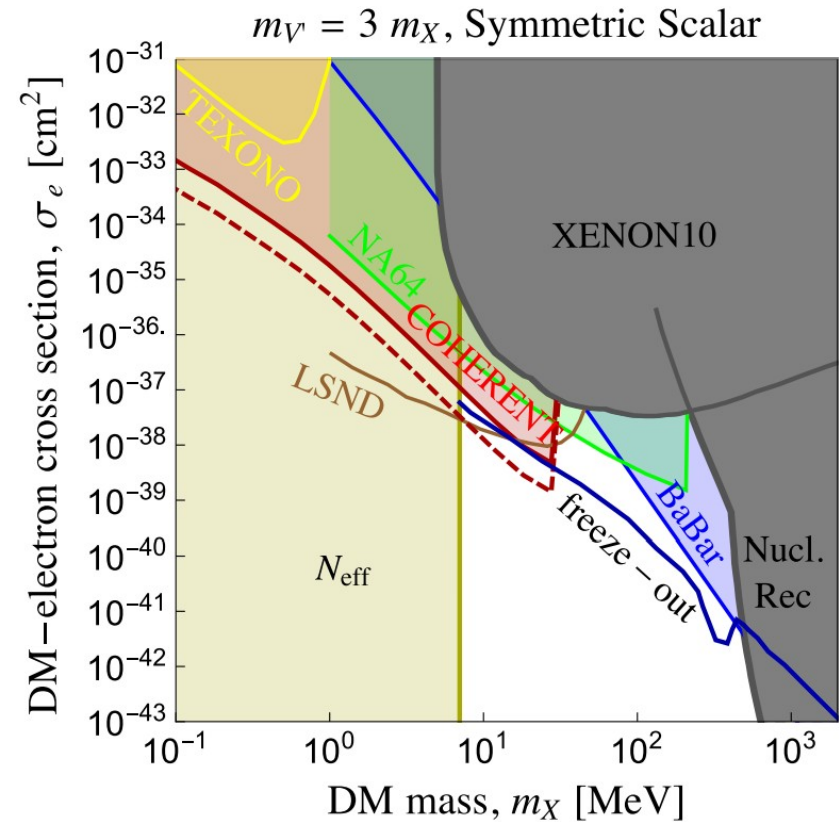
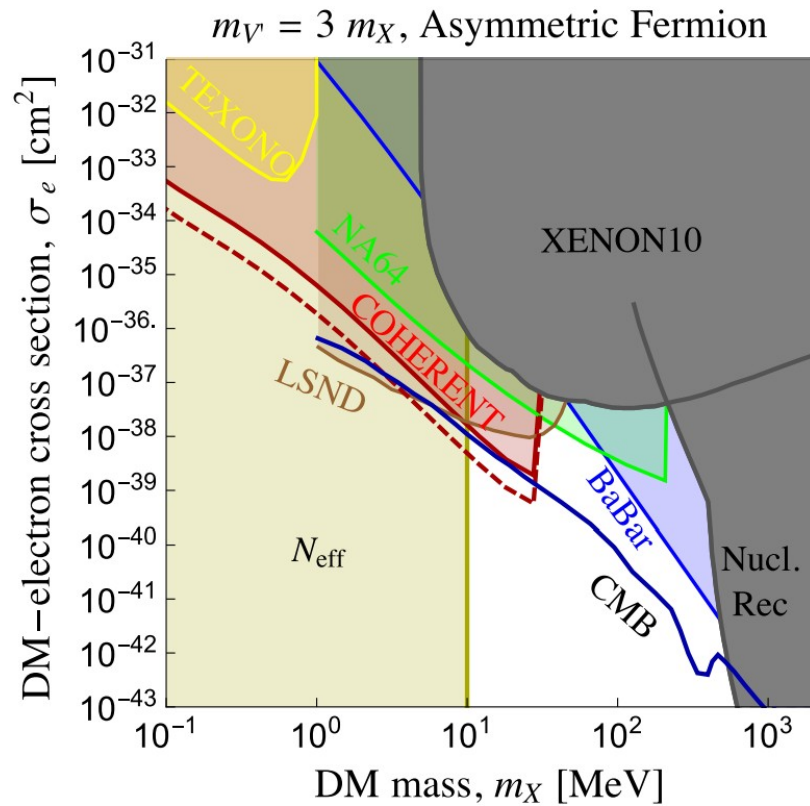
SFG, Pedro Pasquini, Jie Sheng [Phys.Lett.B 810 (2020) 135787]

SFG, Pedro Pasquini, Jie Sheng [JHEP 05 (2022) 088]

Testing light DM @ ν Exps



Romeri, Kelly & Machado, PRD 19 [arXiv:1903.10505]



SFG & Ian Shoemaker, JHEP 11 (2018) 066 [arXiv: 1710.10889]

Seesaw Mechanism (跷跷板机制)

- Heavy neutrinos (N)

$$\bar{\nu} M_D \mathcal{N} + h.c. + \bar{\mathcal{N}} M_N \mathcal{N} = \begin{pmatrix} \bar{\nu} & \bar{\mathcal{N}} \end{pmatrix} \begin{pmatrix} 0 & M_D \\ M_D^T & M_N \end{pmatrix} \begin{pmatrix} \nu \\ \mathcal{N} \end{pmatrix}$$

required by **Grand Unification Theory (大统一理论)**

- Seeaw Mechanism

The diagonalization of the full mass matrix

$$\begin{pmatrix} 0 & M_D \\ M_D^T & M_N \end{pmatrix} \Rightarrow M_\nu = -M_D \frac{1}{M_N} M_D^T$$

$$M_D \sim O(100) \text{ GeV}, \quad M_N \sim O(10^{15}) \text{ GeV} \quad \rightarrow \quad M_\nu \sim O(0.01) \text{ eV}$$

Light neutrino mass M_ν is suppressed by the heavy ones



柳田勉、何小刚：
我们来玩跷跷板吧

Seesaw with Right-Handed ν & DM

$$M_\nu = -M_D \frac{1}{M_N} M_D^T \quad \Rightarrow \quad \text{Det}(M_\nu) \propto \frac{1}{\text{Det}(M_N)}$$

(for NH)	-1σ	Best Value	$+1\sigma$
$\Delta m_s^2 \equiv \Delta m_{12}^2$ (10^{-5}eV^2)	7.37	7.56	7.75
$ \Delta m_a^2 \equiv \Delta m_{13}^2 $ (10^{-3}eV^2)	2.51	2.55	2.59

- The lightest neutrino mass could be zero: $m_{1,3} \rightarrow 0$
- Needs only 2 right-handed neutrinos N_R
- Effectively, the 3rd N_R could be infinitely heavy to decouple
- Could the 3rd N_R be DM?

P. Cox, C. Han, and T. T. Yanagida, "Right-handed Neutrino Dark Matter in a U(1) Extension of the Standard Model," [JCAP01\(2018\) 029](#), [[arXiv:1710.01585](#)]

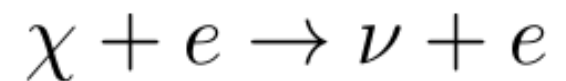
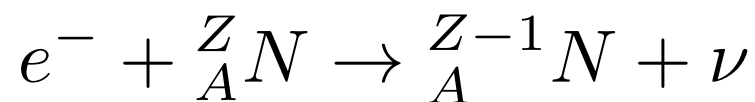
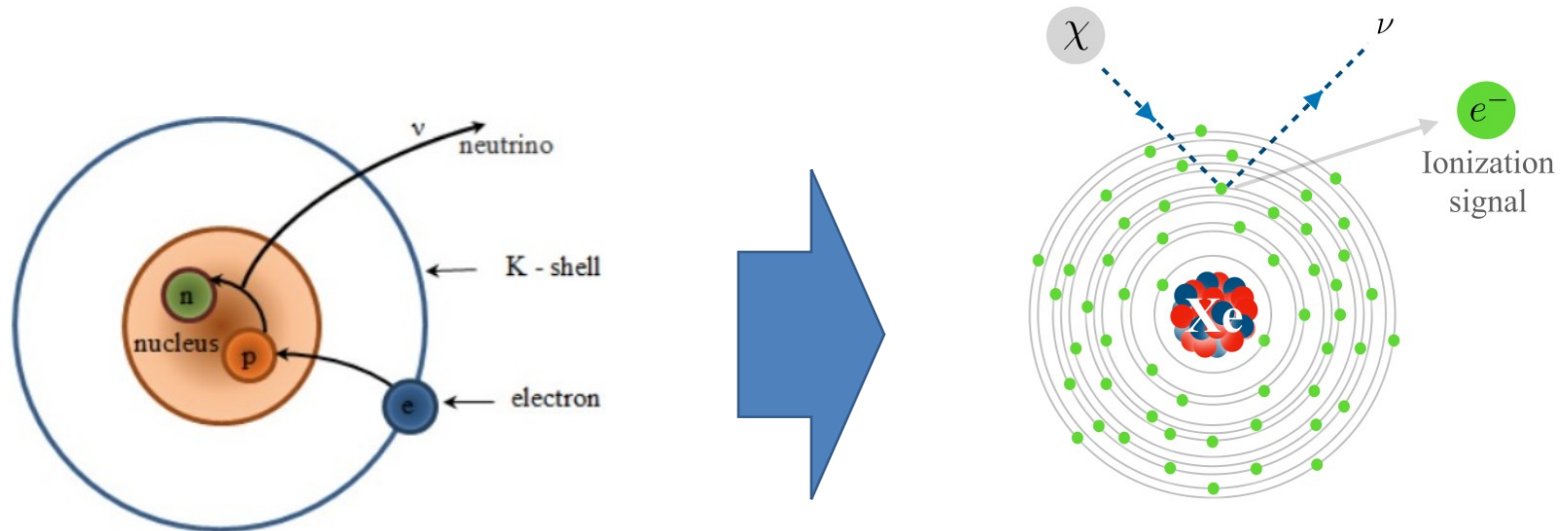
Cheng, Sheng & Yanagida [[arXiv:2312.15637](#)]

K-Shell Electron Capture for ν



王淦昌

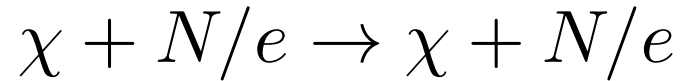
1942 – Kan Chang Wang proposed using K-shell electron capture for detecting neutrino



Kan Chang Wang, *A Suggestion on the Detection of the Neutrino*, Phys. Rev., 61, 97 (1942)

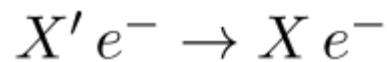
Kan Chang Wang, *Proposed Methods of Detecting the Neutrino*, Phys. Rev., 71, 645-646 (1947)

- Elastic Scattering



$$E_r \approx \frac{4m_\chi m_N}{(m_\chi + m_N)^2} T_\chi$$

- Exothermic

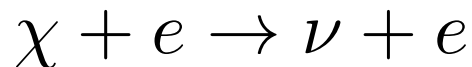


$$E_R \simeq \Delta m \left(1 - \frac{v_{\text{DM}}}{v_e} \cos \theta_e \right)$$

He, Wang & Zheng [JCAP21, 2007.04963]

Aboubrahim, Althueser, Klasen, Nath & Weinheimer [2207.08621]

- Fermionic absorption

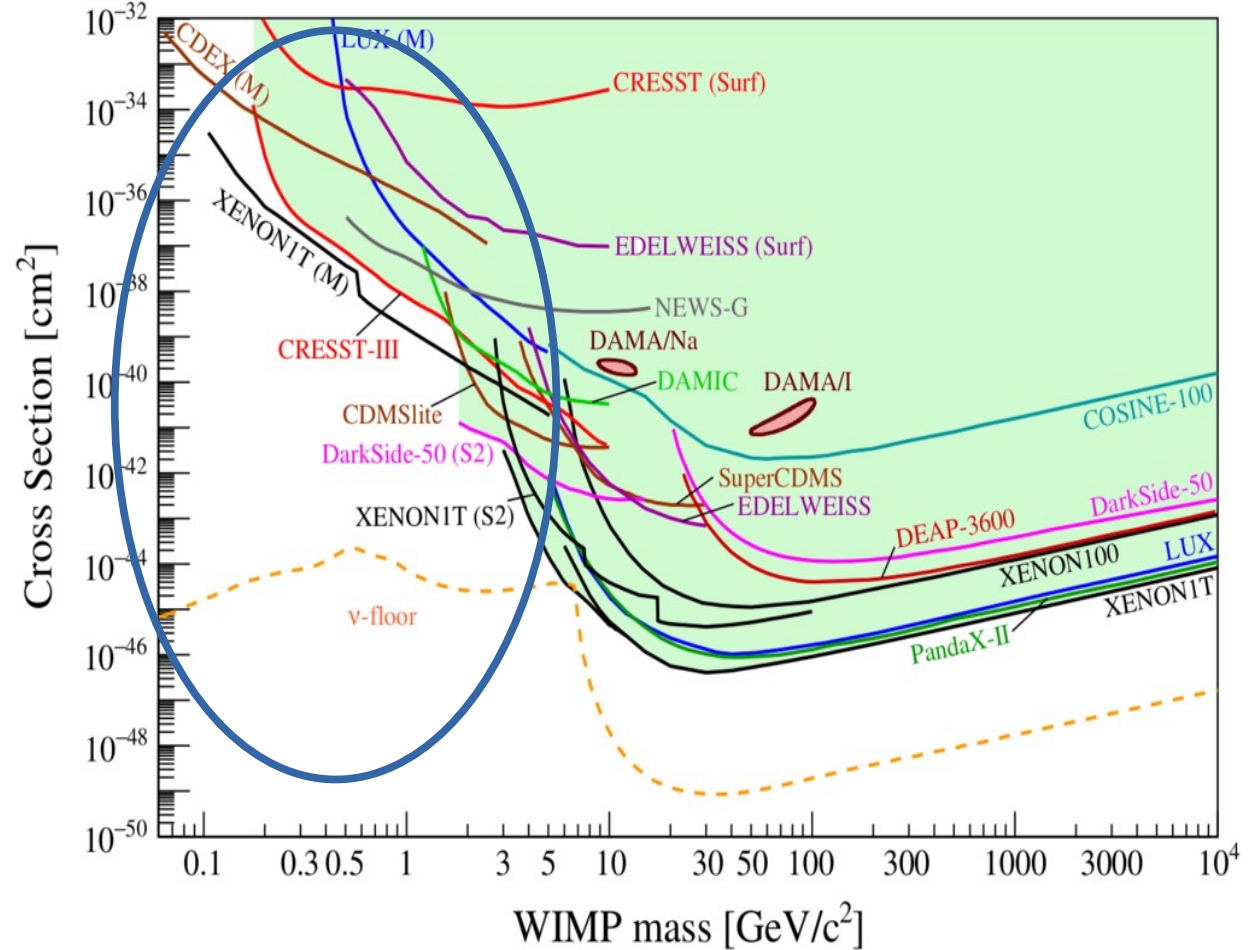
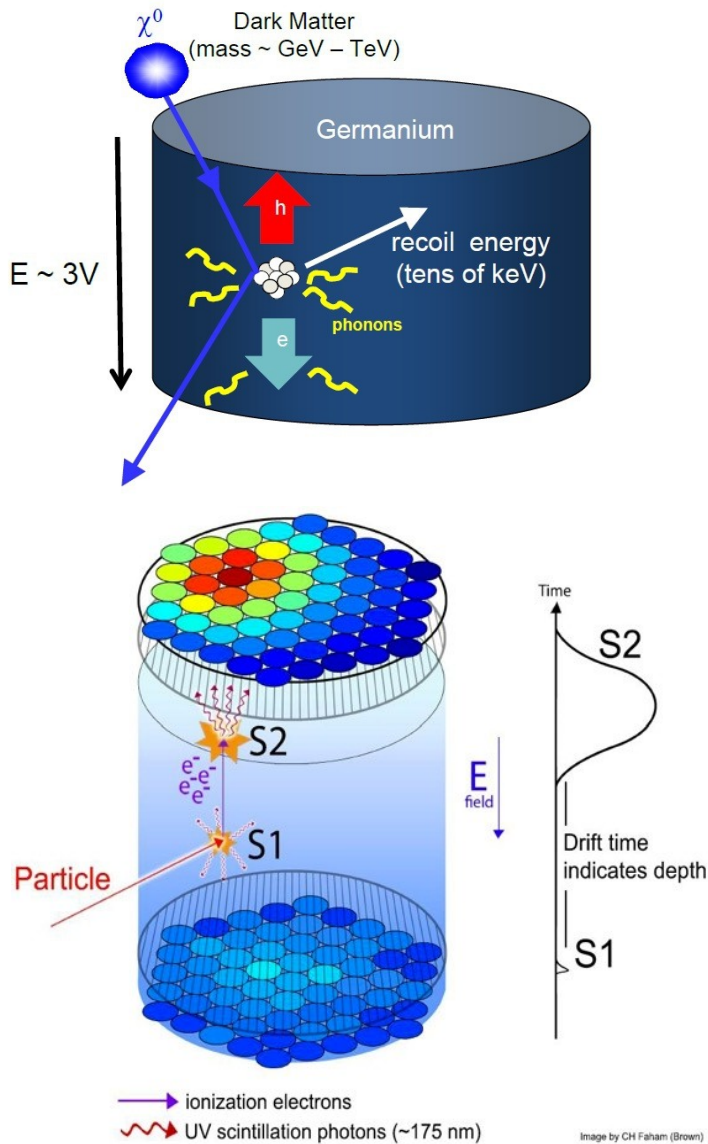


$$E_r \approx \frac{m_\chi^2}{2m_e} \quad m_\chi \sim \mathcal{O}(10) \text{ keV}$$

Dror, Elor, McGehee & Yu [2011.01940]

See also Dror, Elor & McGehee
[1905.12635, 1908.10861]

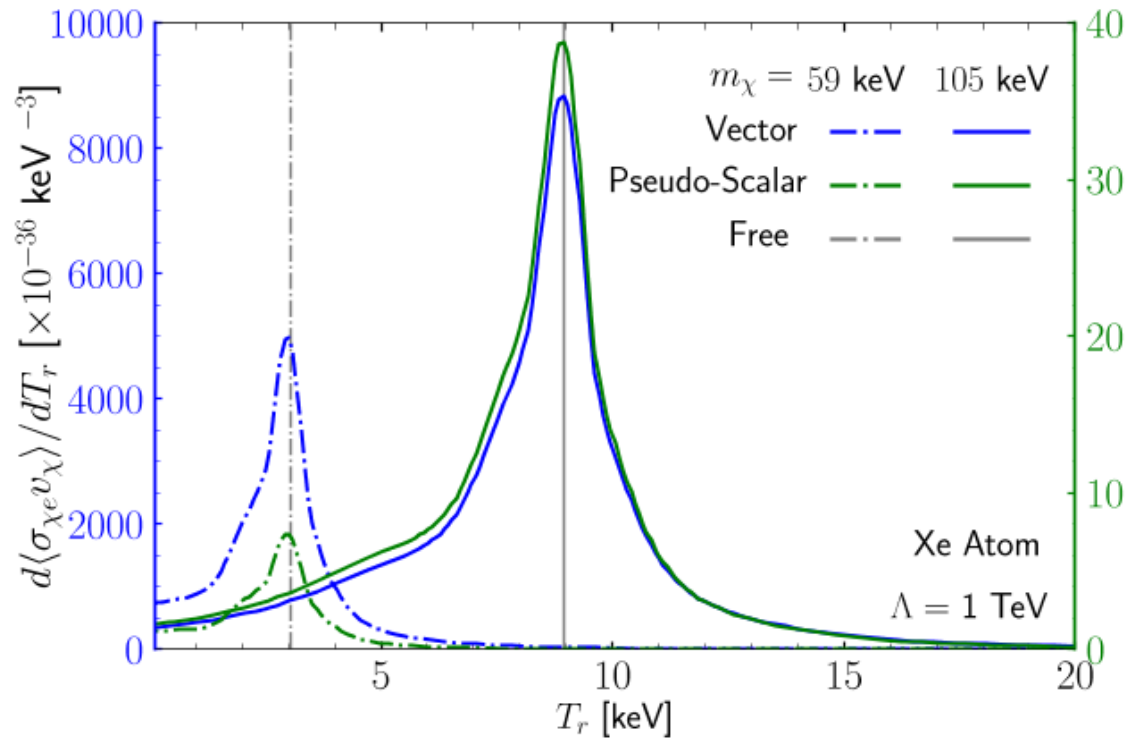
Direct Detection



APPEC Committee Report [2104.07634]

Kinematics of Absorption DM

$$\chi + e \rightarrow \nu + e \quad E_r = \frac{m_\chi^2}{2(m_e + m_\chi)} \approx \frac{m_\chi^2}{2m_e}$$

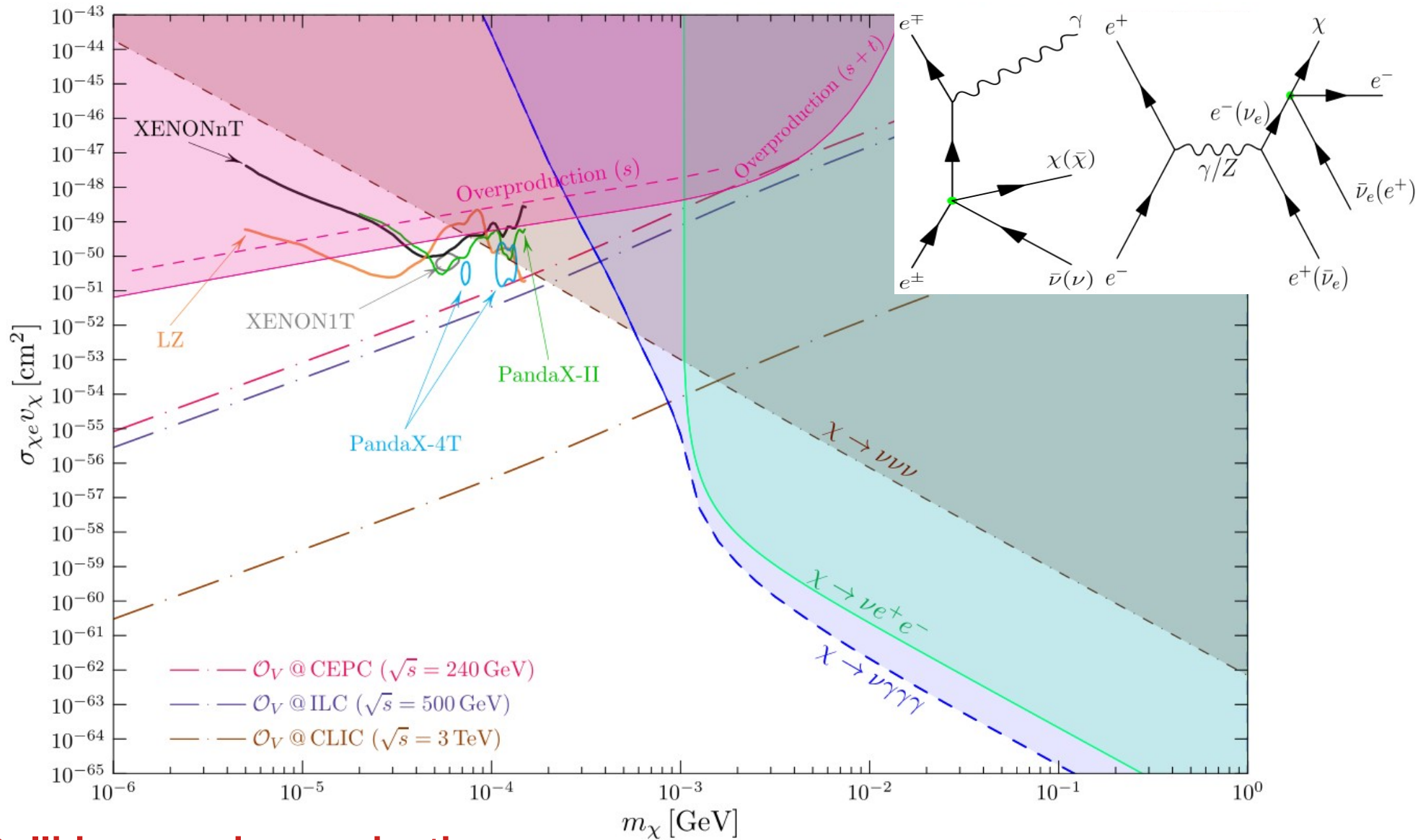


$$\frac{d\langle\sigma_{\chi e\nu\chi}\rangle}{dT_r} = \sum_{nl} (4l + 2) \frac{1}{T_r} \frac{m_\chi - \Delta E_{nl}}{16\pi m_e^2 m_\chi} |\mathcal{M}|^2(\mathbf{q}) K_{nl}(T_r, |\mathbf{q}|),$$

SFG, Xiao-Gang He, Xiao-Dong Ma, Jie Sheng [JHEP 05 (2022) 191]

SFG, Pedro Pasquini, Jie Sheng [JHEP 05 (2022) 088]

PandaX-4T & Lepton Colliders

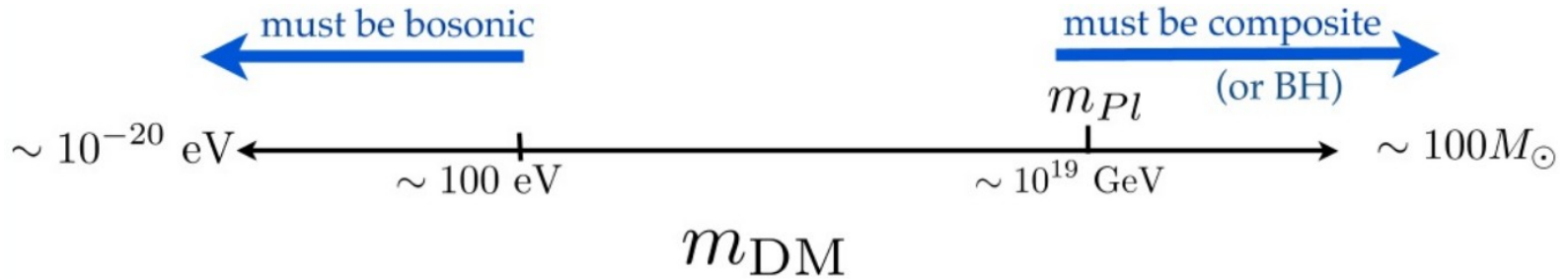


Collider searches probe the dark sector, not just DM

SFG, Kai Ma, Xiao-Dong Ma, Jie Sheng [arXiv:2306.00657]

Tiny-Varying Neutrino Mass

- Wave DM can be as light as 10^{-22} eV



$$-\mathcal{L} = \frac{1}{2} m_{\phi}^2 \phi^2 + \frac{1}{2} M_{\alpha\beta} \bar{\nu}_{\alpha} \nu_{\beta} + \underbrace{y_{\alpha\beta} \phi \bar{\nu}_{\alpha} \nu_{\beta}} + h.c.$$

- Bose-Einstein Condensation** $\langle \phi \rangle \neq 0$ **v Mass term**

$$\phi(x) \simeq \frac{\sqrt{2 \rho_{\text{DM}}(x)}}{m_{\phi}} \cos [m_{\phi} (t - \vec{v} \cdot \vec{x})]$$

Berlin [1608.01307]

Zhao [1701.02735]

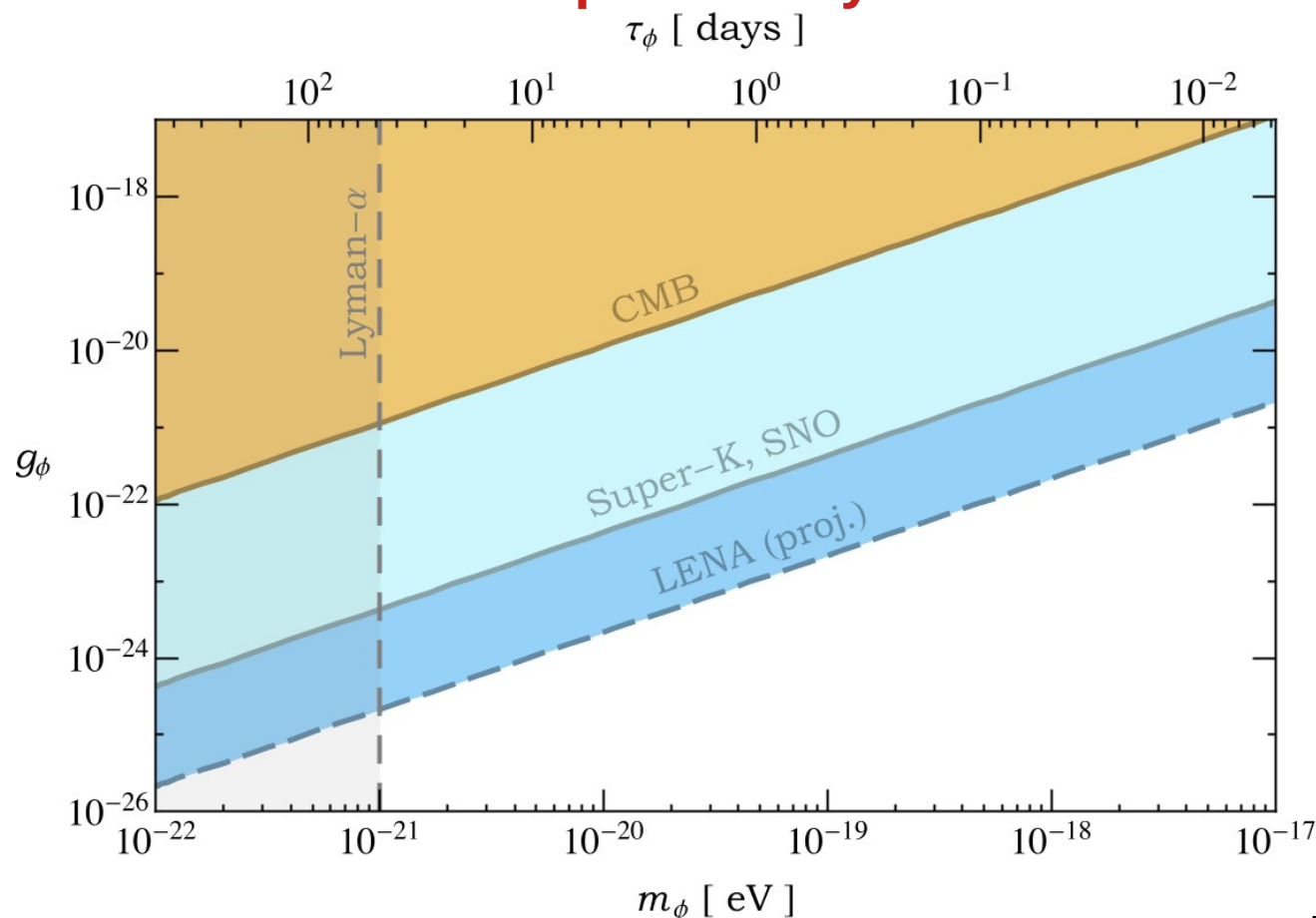
Brdar, Kopp, Liu, Prass & Wang [1705.09455]

Time-Varying oscillation probabilities!

Time-Dependent Oscillation

$$\sin \theta_{12}(t) \simeq \sin \theta_{12} + \frac{\cos \theta_{12}}{\Delta m_{12}} \frac{g_{\phi} \sqrt{2} \rho_{\text{DM}}}{m_{\phi}} \cos m_{\phi} t$$

- **Neutrino oscillation probability varies with time!**

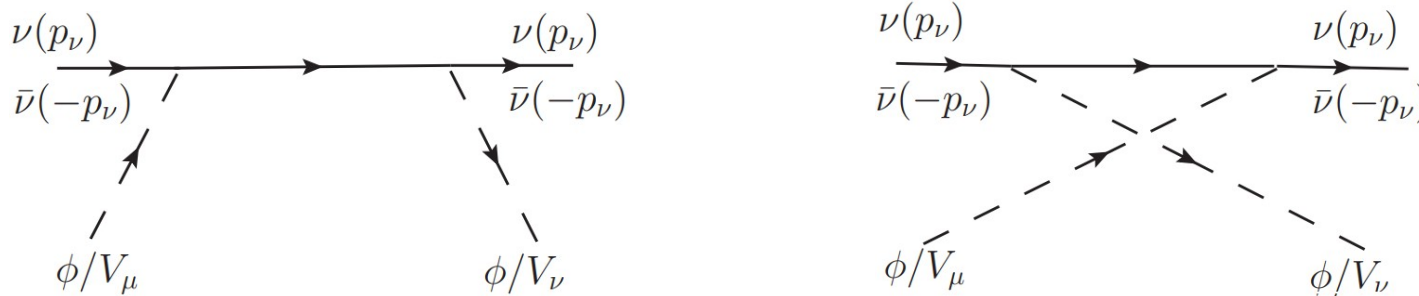


通过中微子振荡探测暗物质

Berlin [1608.01307]

$$-\mathcal{L} = \frac{1}{2}m_\phi^2\phi^2 + \frac{1}{2}M_{\alpha\beta}\bar{\nu}_\alpha\nu_\beta + y_{\alpha\beta}\phi\bar{\nu}_\alpha\nu_\beta + h.c.$$

- Forward scattering in DM medium



$$H = \frac{M^2}{2E_\nu} - \frac{1}{E_\nu} \sum_j y_{\alpha j} y_{j\beta}^* \frac{\rho_\chi}{m_\phi^2} \equiv \frac{M^2 + \delta M^2}{2E_\nu}$$

- With $1/E_\nu$ dependence, the dark potential is promoted to mass

Dark NSI can fake neutrino mass! Neutrino may not be genuine but environmental!

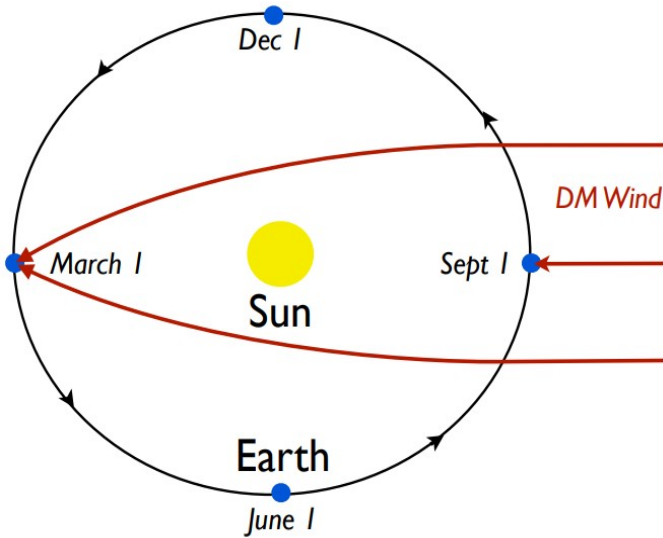
SFG, Murayama [1904.02518]

SFG, PoS NuFact2019 (2020) 108

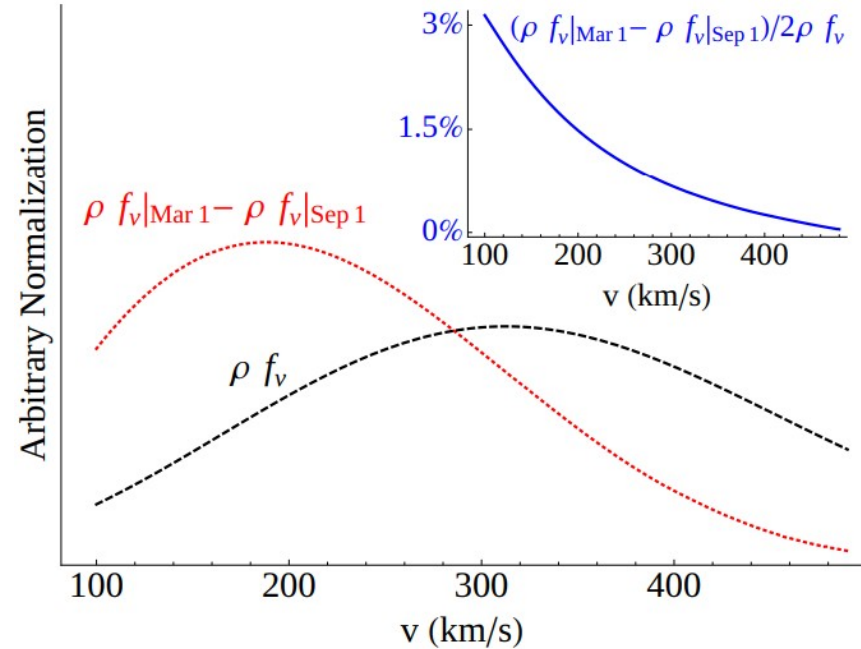
SFG, Chui-Fan Kong, Alexei Smirnov, to appear soon

SFG, J.Phys.Conf.Ser. 1468 (2020) 1, 012125

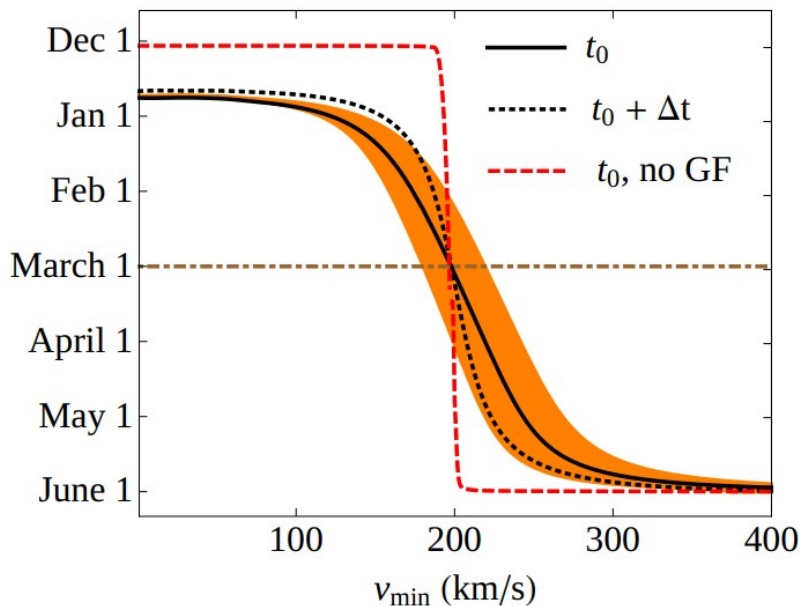
DM Gravitational Focusing



- Density enhancement



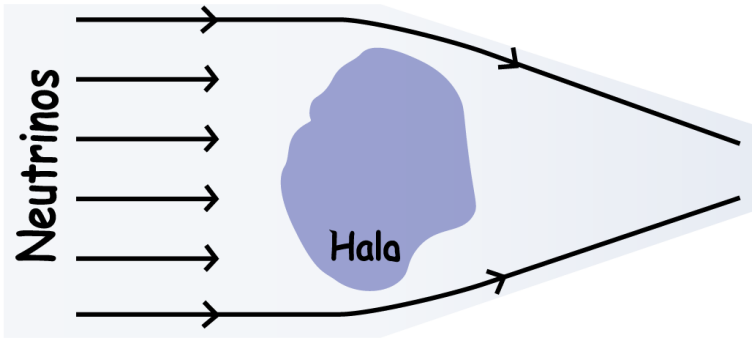
- Modulation phase shift



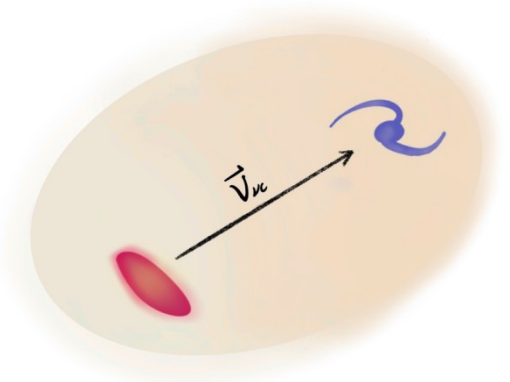
Lee, Lisanti, Peter, Safdi, Phys. Rev. Lett. 112, 011301 (2014) [arXiv:1308.1953]

Bozorgnia & Schwetz, JCAP 08 (2014) 013 [arXiv:1405.2340]

Cosmic Gravitational Focusing



$$\delta_m = \delta_{m0} + \sum_{i=1}^3 \frac{\delta\rho_{\nu i}}{\rho_m}$$



Density enhancement downwind!

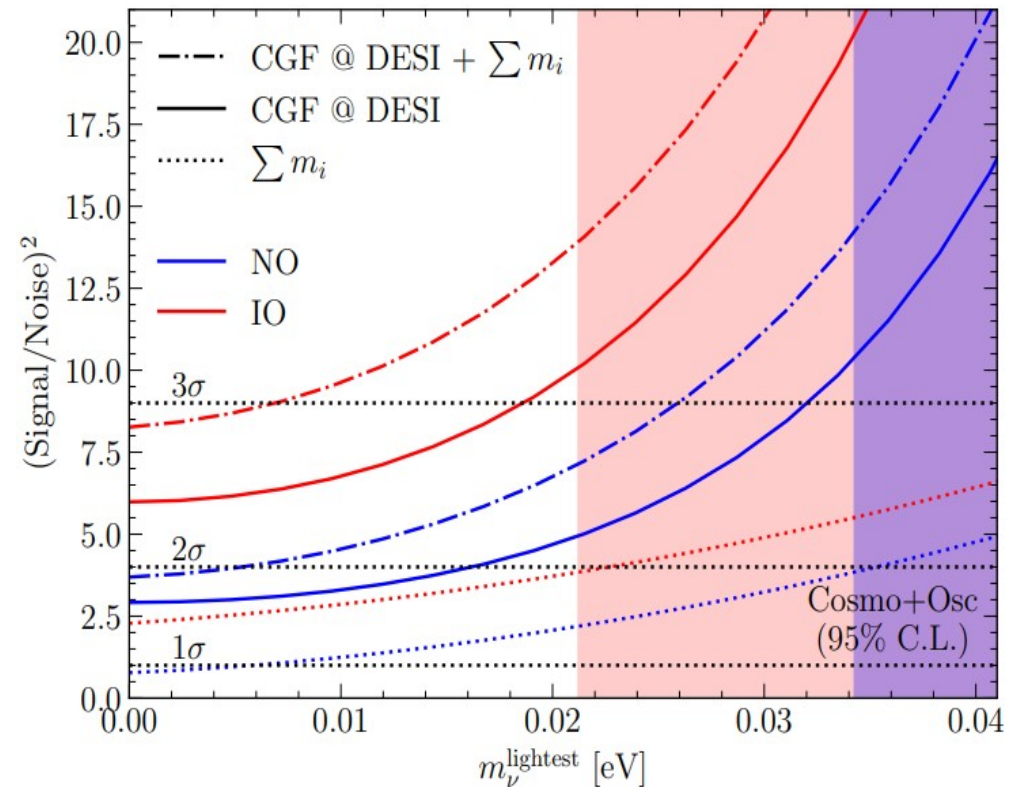
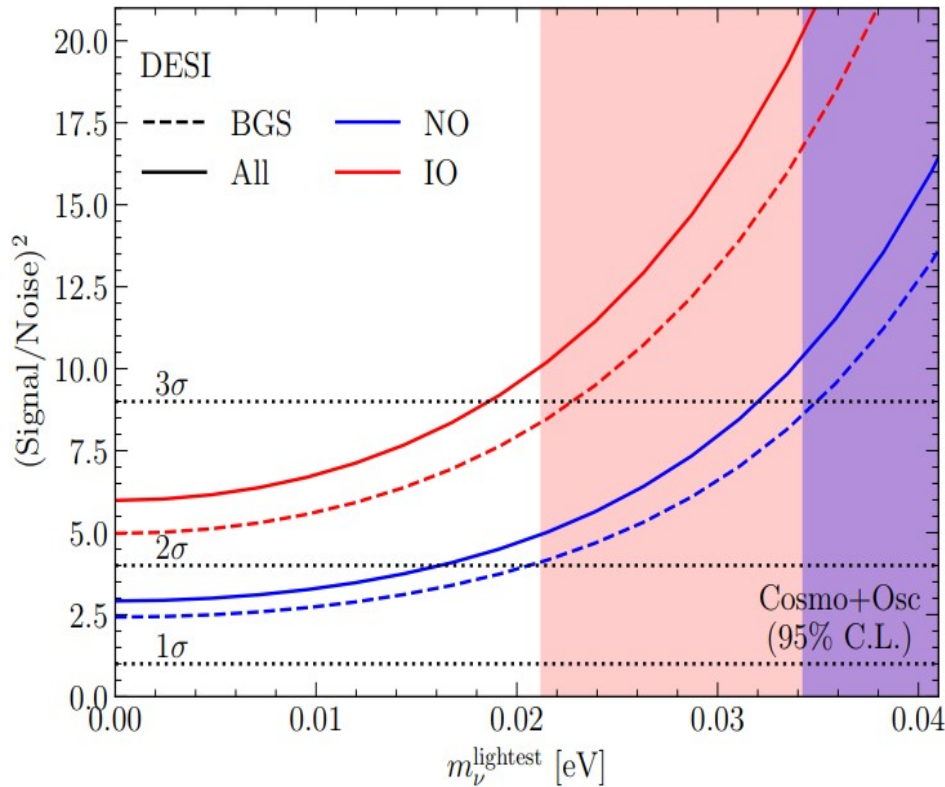
$$\delta\rho_\nu(-\mathbf{x}) = -\delta\rho_\nu(\mathbf{x})$$

$$\tilde{\delta}_m \equiv \tilde{\delta}_{m0}(1 + i\phi)$$

$$\begin{aligned} [\tilde{A}(\mathbf{k})]^* &= \int d\mathbf{x} e^{i\mathbf{k}\cdot\mathbf{x}} A(\mathbf{x}) = \int d\mathbf{x} e^{-i\mathbf{k}\cdot\mathbf{x}} A(-\mathbf{x}) \\ &= - \int d\mathbf{x} e^{-i\mathbf{k}\cdot\mathbf{x}} A(\mathbf{x}) = -\tilde{A}(\mathbf{k}) \end{aligned}$$

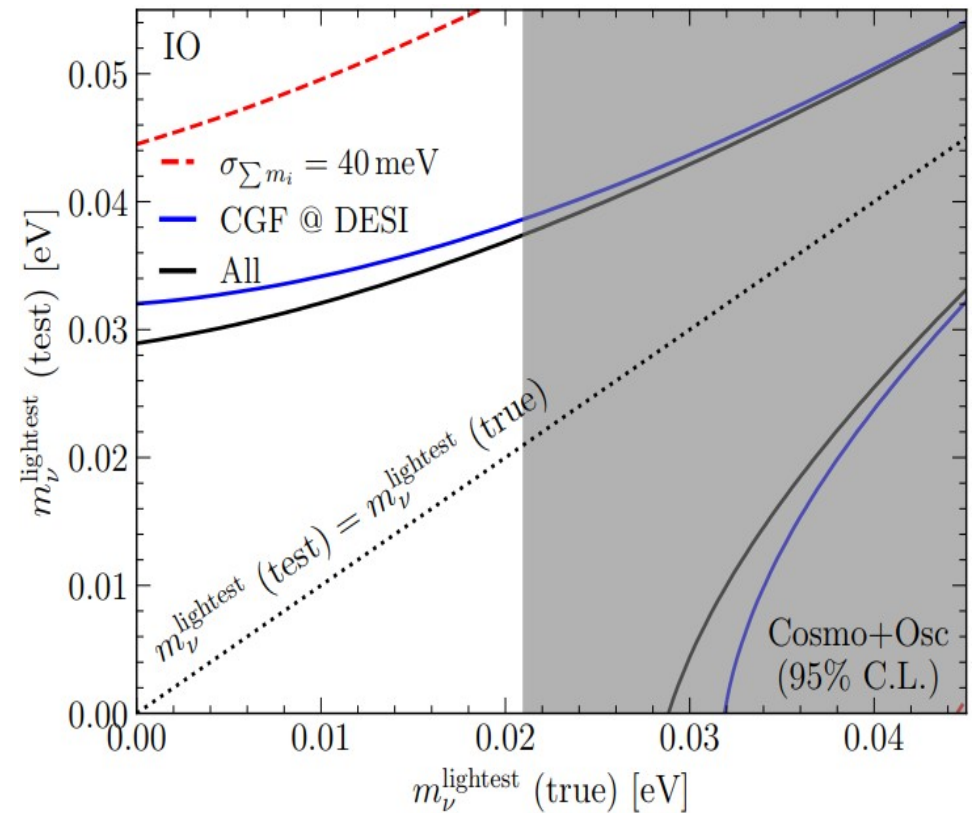
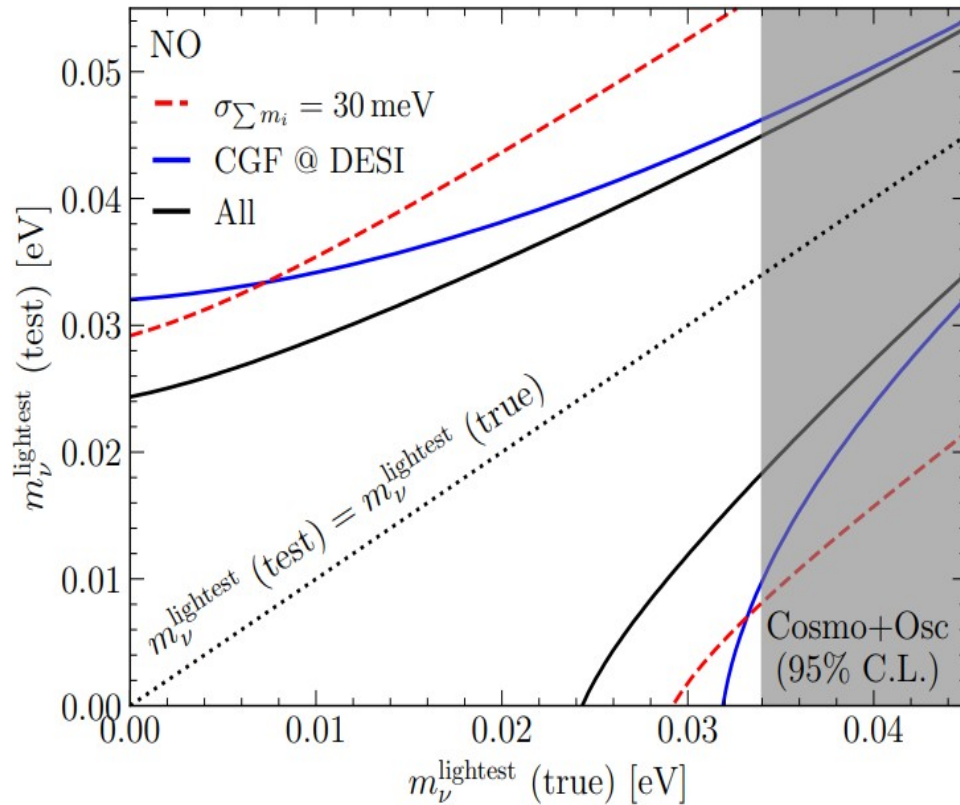
SFG, Pedro Pasquini, Liang Tan [arXiv:2312.16972]

$$\text{Im} \left[\tilde{\delta}_{g\alpha} \left(\tilde{\delta}_{g\beta} \right)^* \right] \propto \left(\mathbf{v}_{\nu c} \cdot \hat{\mathbf{k}} \right) \left(f_0 m_\nu^4 + f_1 m_\nu^2 T^2 + f_2 T^4 \right)$$



SFG, Pedro Pasquini, Liang Tan [arXiv:2312.16972]

$$\text{Im} \left[\tilde{\delta}_{g\alpha} \left(\tilde{\delta}_{g\beta} \right)^* \right] \propto \left(\mathbf{v}_{\nu c} \cdot \hat{\mathbf{k}} \right) \left(f_0 m_\nu^4 + f_1 m_\nu^2 T^2 + f_2 T^4 \right)$$



DESI, Euclid, Subaru PFS, **CSST**

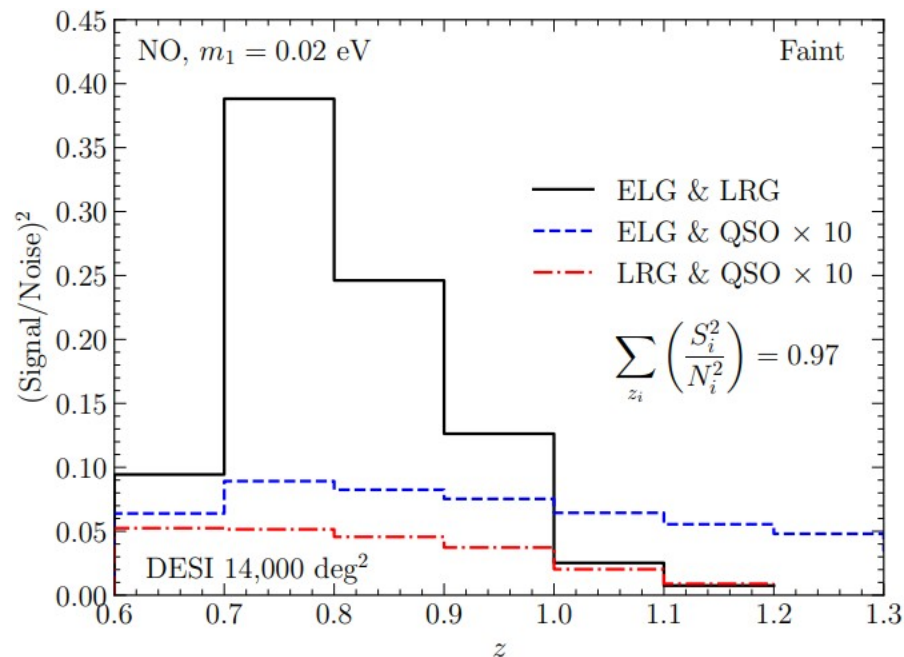
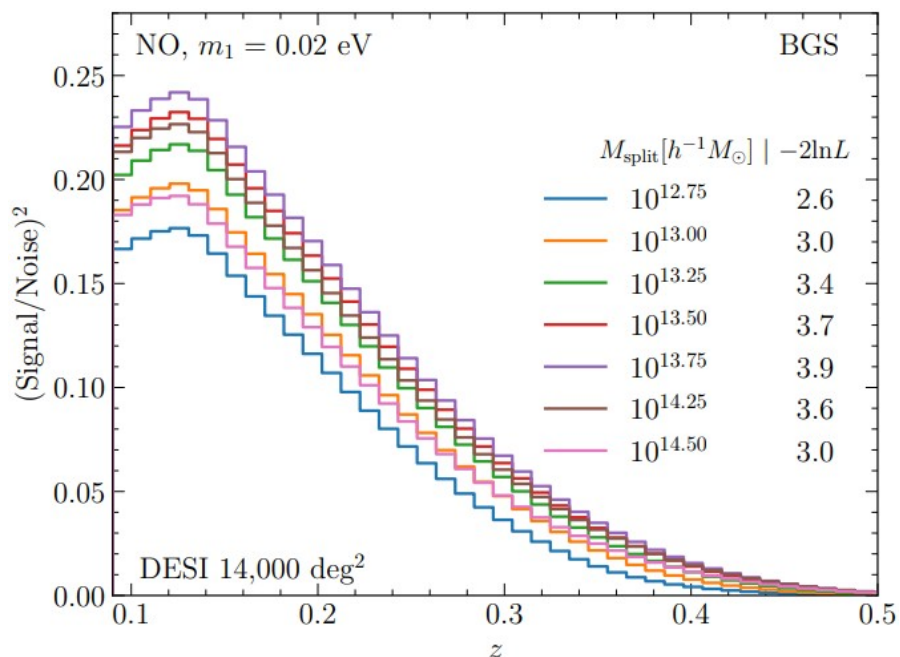
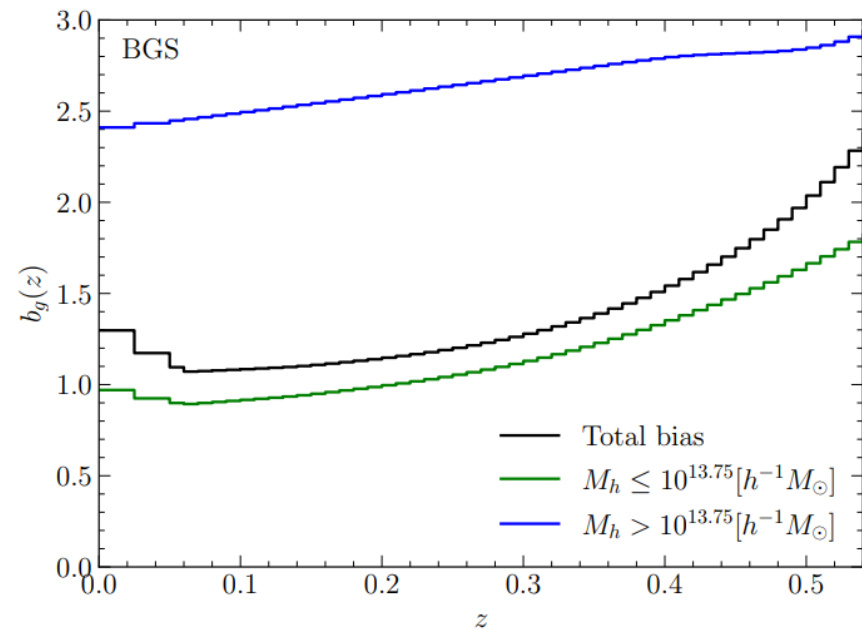
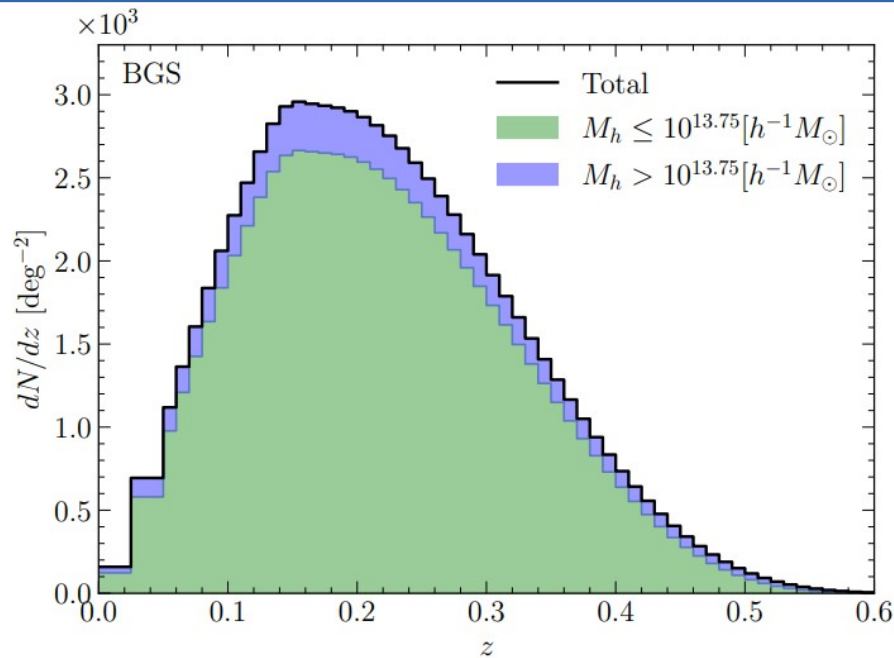
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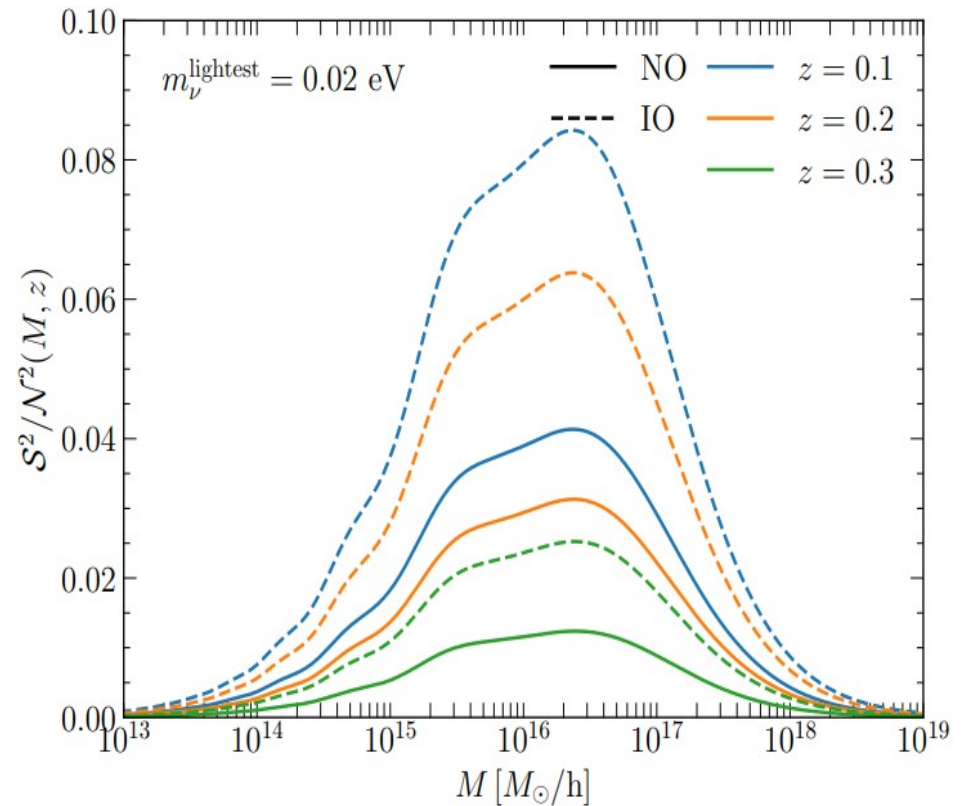
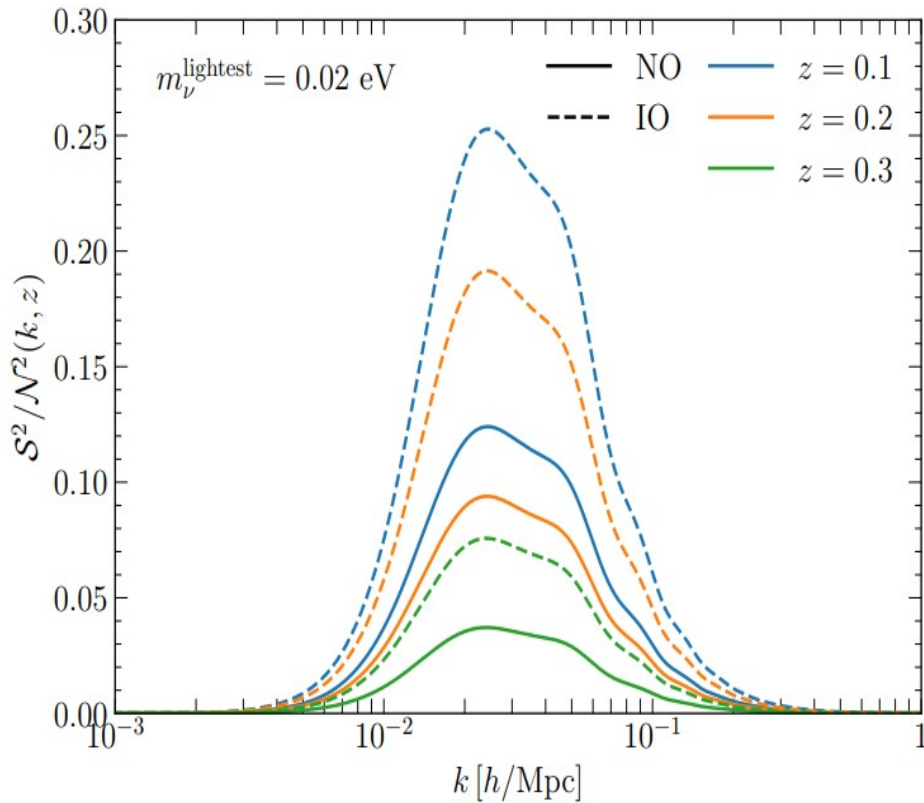
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Thank You

Backup Slides

Galaxy Categories





SFG, Pedro Pasquini, Liang Tan [arXiv:2312.16972]