

Axion dark matter and the cosmic dipole problem

韩成成
中山大学

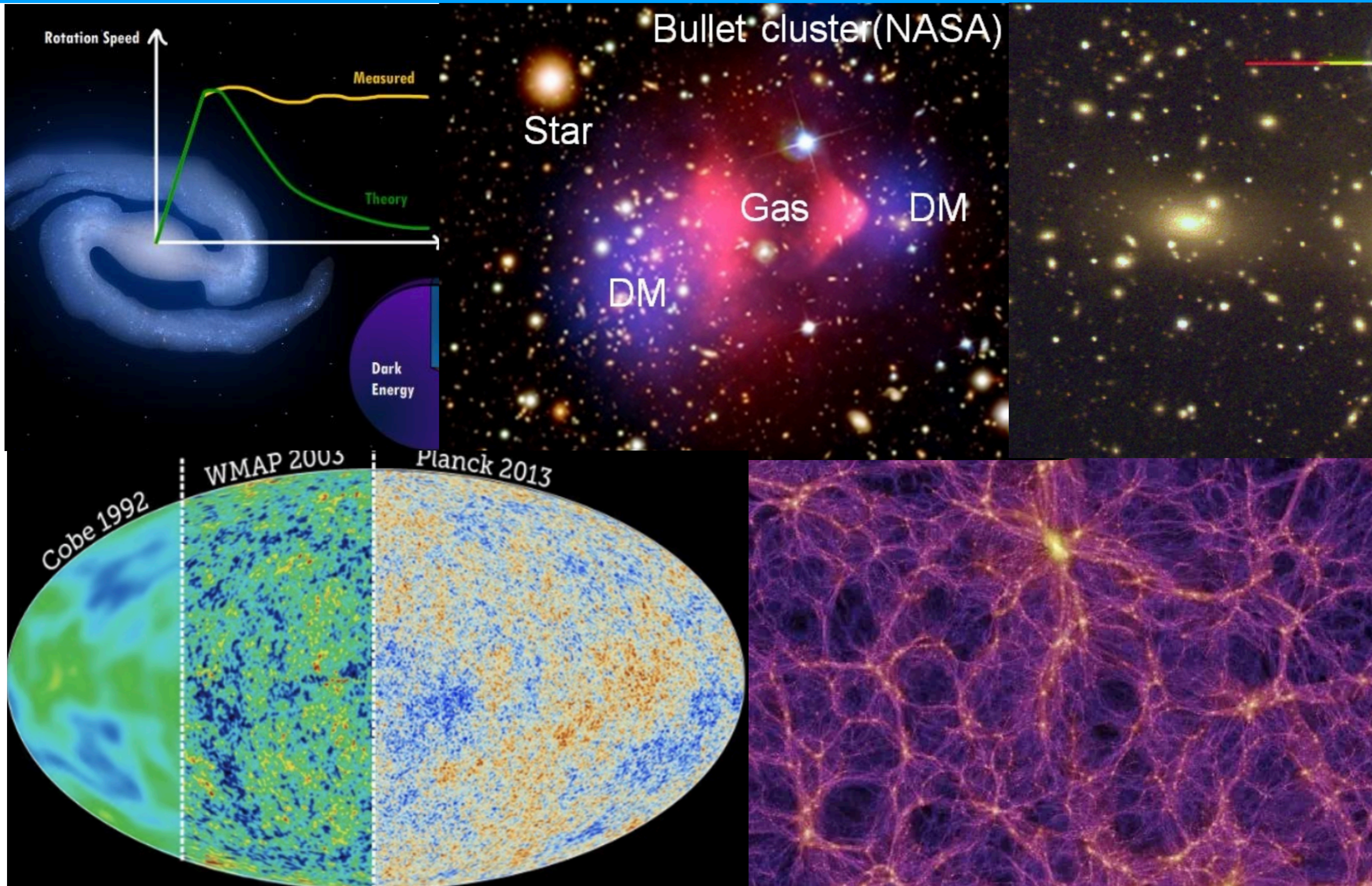
Based on arXiv: 2211.06912(PRD)

紫金山暗物质研讨会

南京师范大学

2023.12.29-2023.12.31

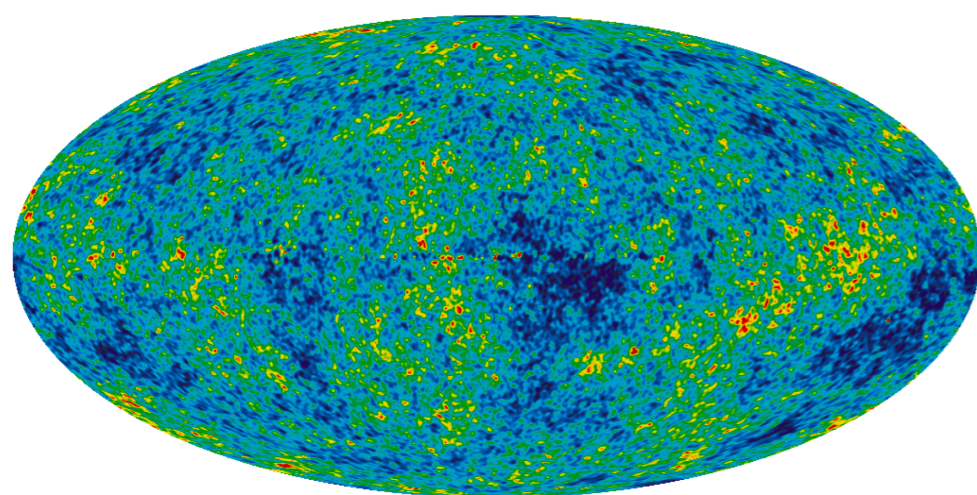
Evidence of dark matter at different length scales(kpc-Gpc)



Dark matter is important for structure formation

Quantum fluctuations from inflation

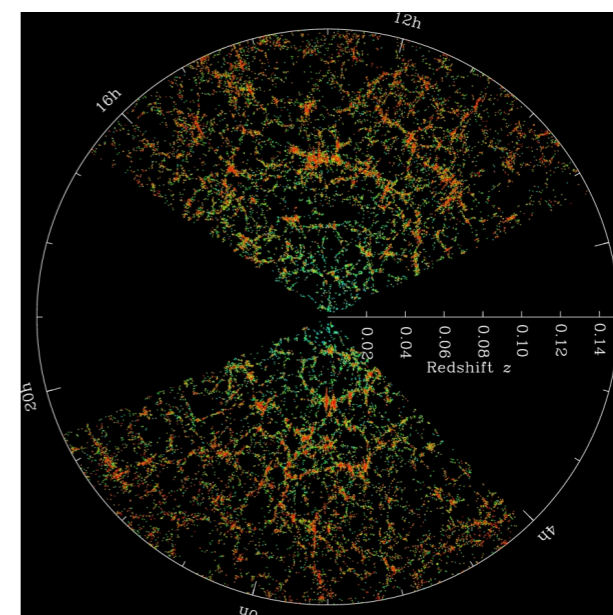
$$\delta\phi \sim H/2\pi \longrightarrow \frac{\delta\rho}{\rho} \sim 10^{-5}$$



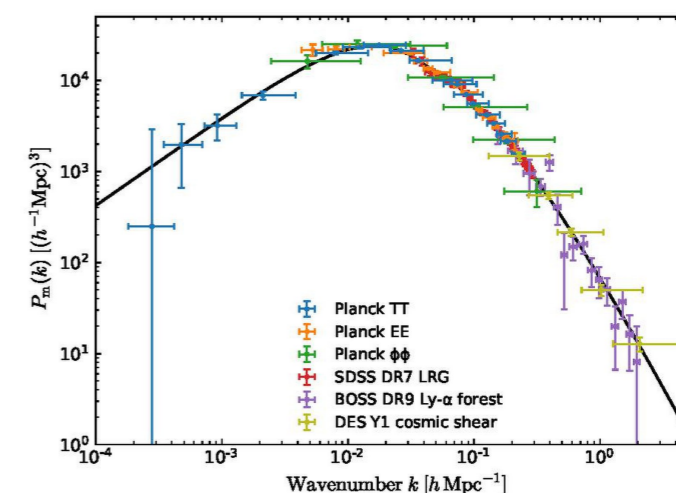
$$\frac{\delta T}{T} \sim 10^{-5}$$

dark matter

Void(空洞)
Filaments(丝)



通过对暗物质扰动的观测，我们能
获得暗物质性质的一些蛛丝马迹



Cosmological principles

Modern cosmology is based on the cosmological principle:

On a large enough scale, the Universe is homogeneous and isotropic

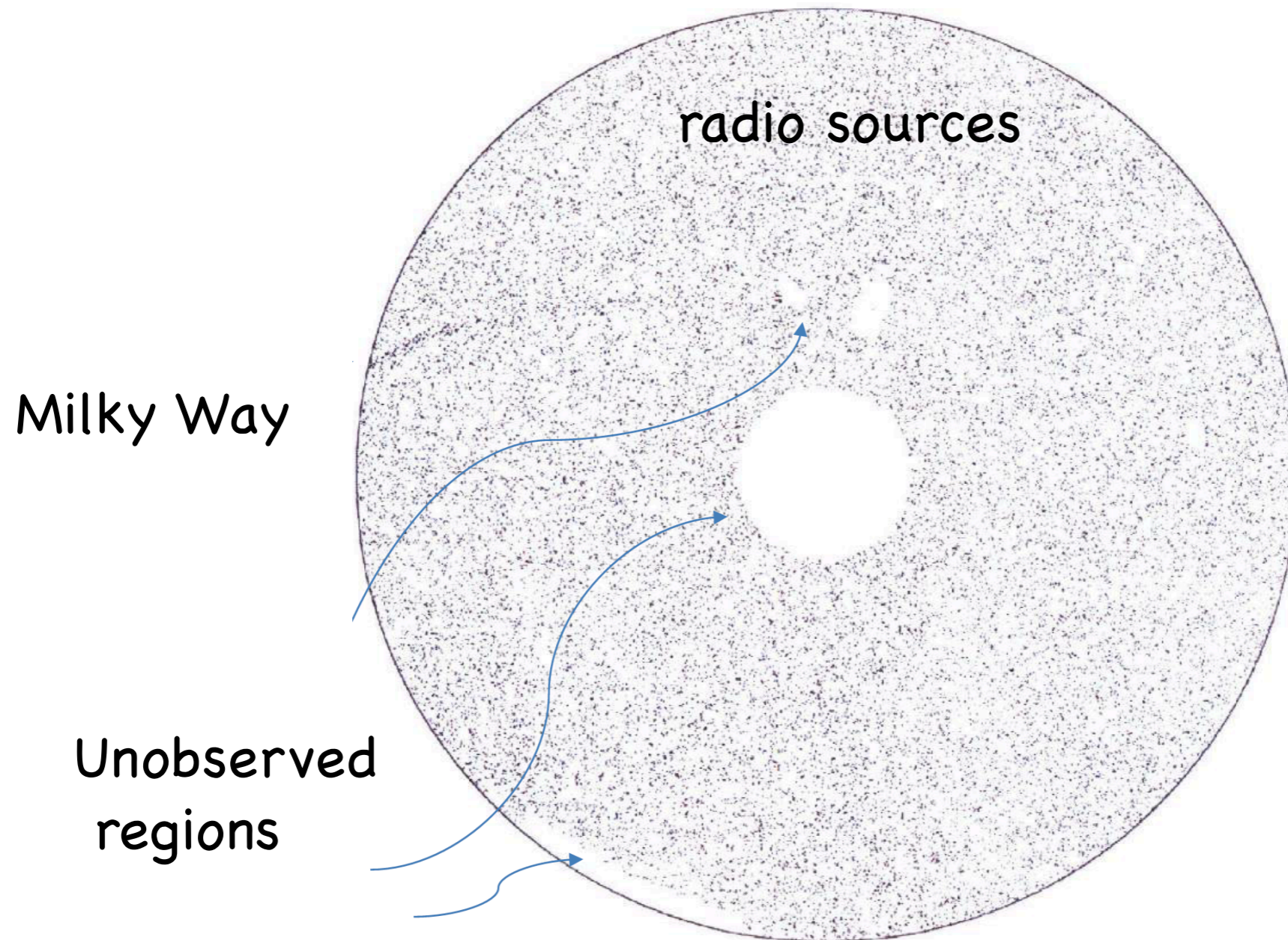


Friedmann–Robertson–Walker (FRW) metric

$$ds^2 = -dt^2 + a^2(t) \left(\frac{dr^2}{1 - kr^2} + r^2(d\theta^2 + \sin^2 \theta d\phi^2) \right)$$

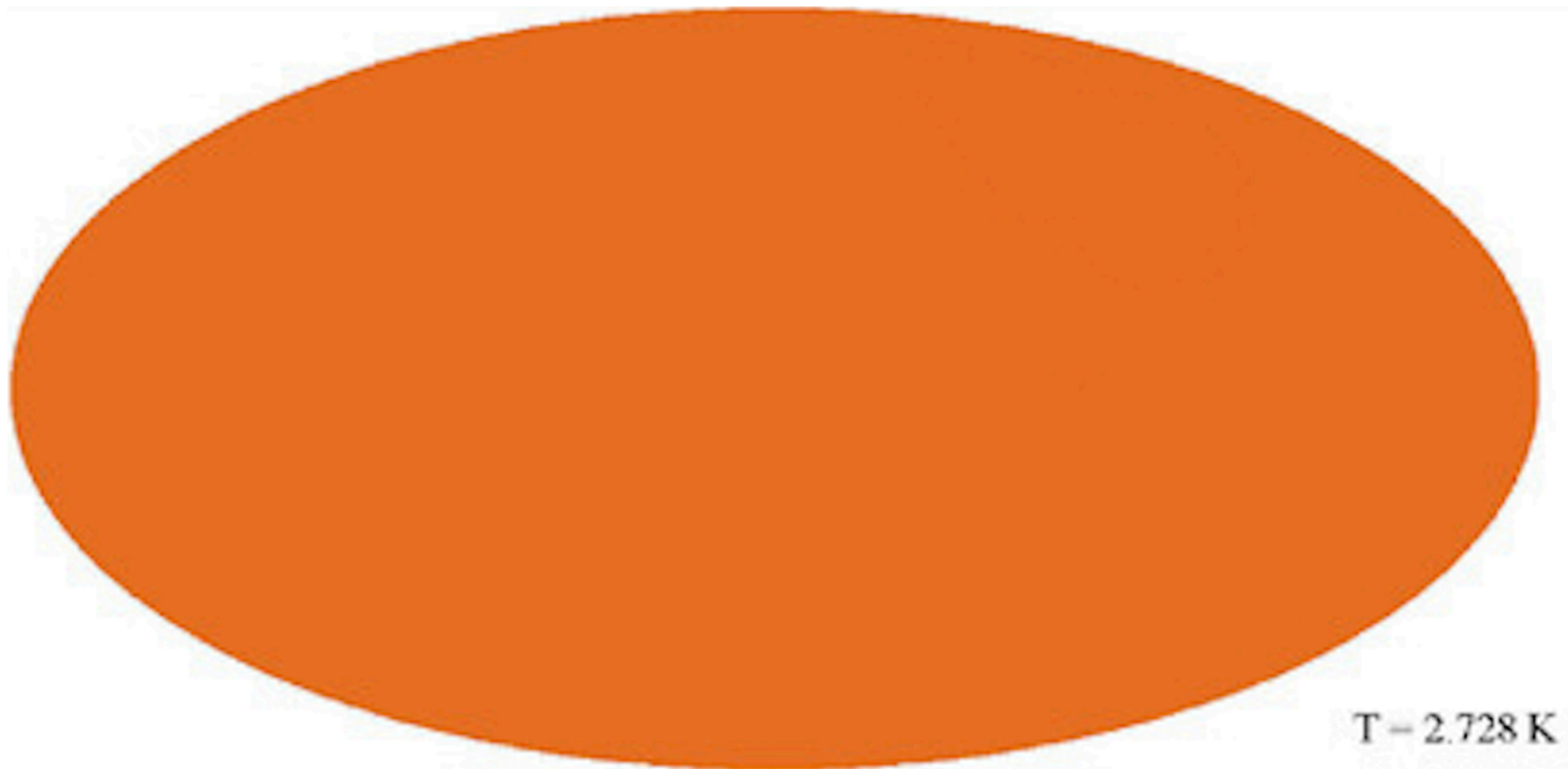
Cosmological principles

Observations support cosmological principle

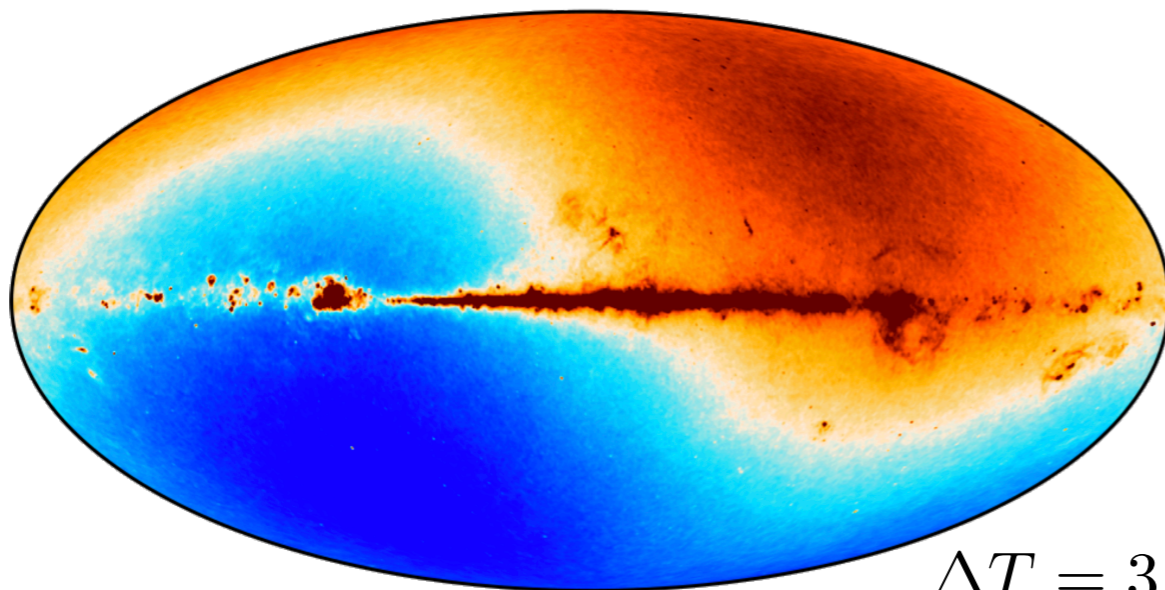


Peebles, Principles of Physical Cosmology, 1993

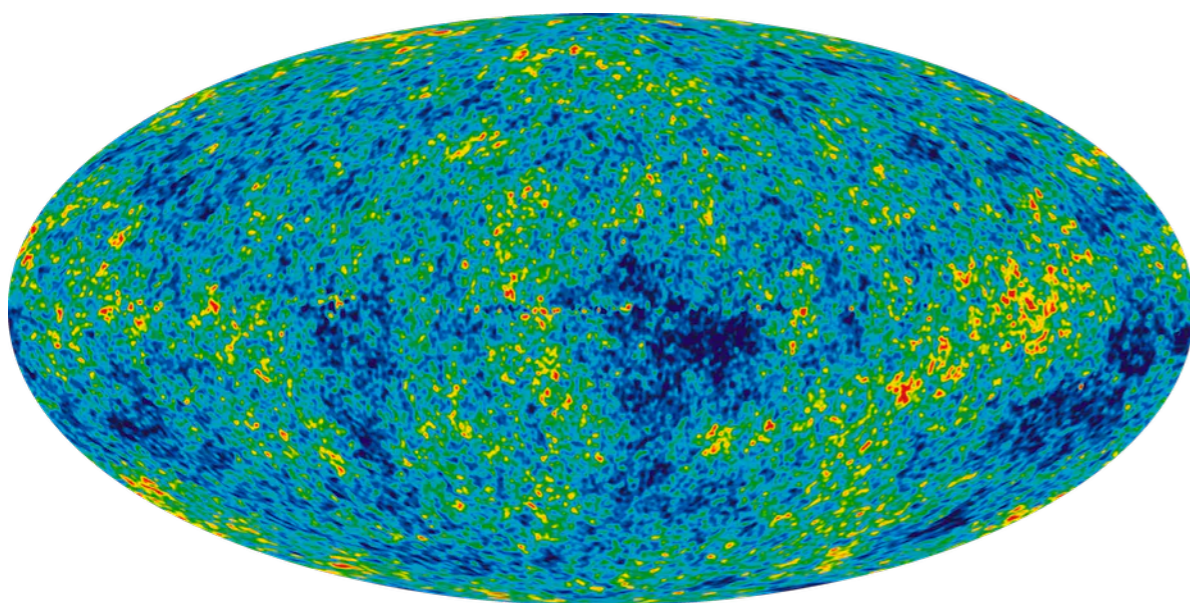
cosmic microwave background



Dipole from CMB



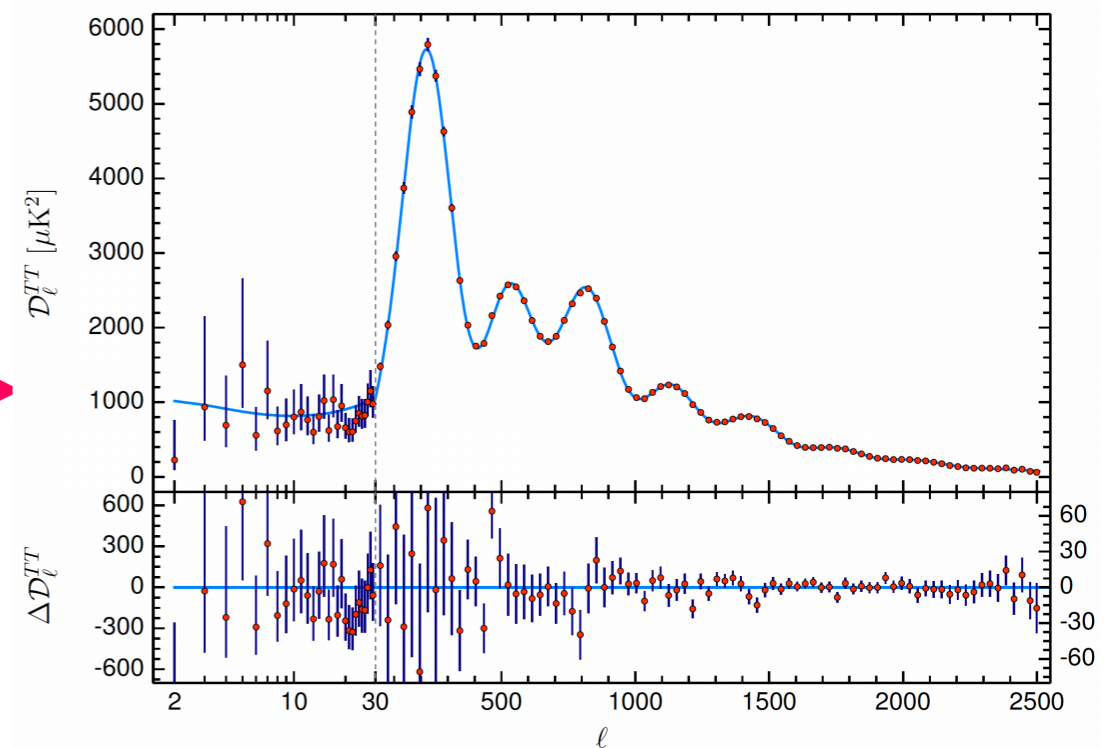
↓ Minus 370 km/s
plus inpainting



$$\frac{\Delta T}{T}(\hat{n}) = \sum_{l,m} a_{lm} Y_{lm}(\hat{n})$$



Relative velocity	Speed [km s ⁻¹]	l [deg]	b [deg]
Sun-CMB ^a	369.82 ± 0.11	264.021 ± 0.011	48.253 ± 0.005
Sun-LSR ^b	17.9 ± 2.0	48 ± 7	23 ± 4
LSR-GC ^c	239 ± 5	90	0
GC-CMB ^d	565 ± 5	265.76 ± 0.20	28.38 ± 0.28
Sun-LG ^e	299 ± 15	98.4 ± 3.6	-5.9 ± 3.0
LG-CMB ^d	620 ± 15	271.9 ± 2.0	29.6 ± 1.4



Dipole from CMB

Relative velocity	Speed [km s ⁻¹]	l [deg]	b [deg]
Sun-CMB ^a	369.82 ± 0.11	264.021 ± 0.011	48.253 ± 0.005



We are not the “rest” observer



Dipole in radio sources(distant galaxies)

Testing the cosmological principle

We should observe the dipole anisotropy of discrete objects (galaxies, quasars)

Ellis & Baldwin (1984): for sources in a flux-limited catalog

$$\frac{dN}{d\Omega}(S > S_*) \propto S_*^{-x}; \quad S \propto \nu^\alpha$$

Typical values $x = 0.7$ to 1.1 , $\alpha = -0.9$ to -0.7

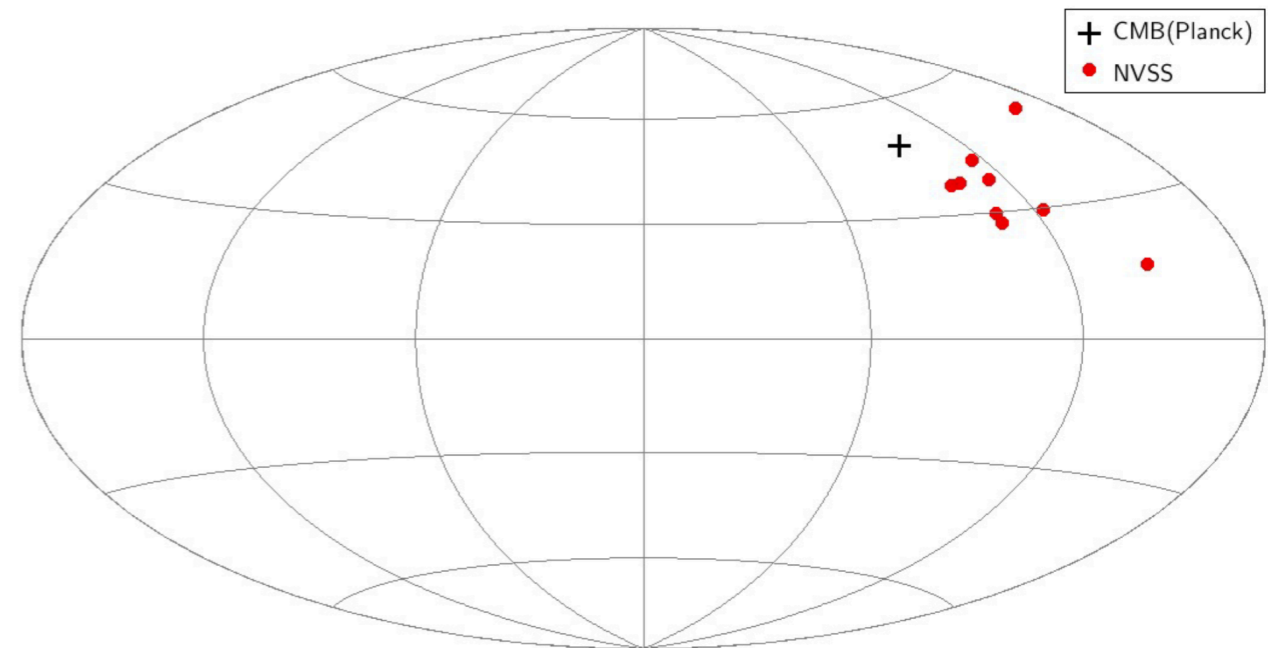
+ aberration & Doppler boosting

$$\left[\frac{dN}{d\Omega} \right]_{\text{obs}} = \left[\frac{dN}{d\Omega} \right]_{\text{com}} (1 + d_{\text{radio}} \cos \theta + \dots); \quad d_{\text{radio}} = [2 + x(1 - \alpha)] \frac{v}{c}$$

Testing the cosmological principle

NVSS - NRAO VLA Sky Survey Catalog

Source	d (10^{-2})	R.A. (deg)	decl. (deg)	Significance (σ)
Blake & Wall (2002)	0.8	148	+31	1.5
Singal (2011)	1.9	157	-12	3
Gibelyou & Huterer (2012)	2.7	214.5	+15.6	>2.3
Rubart & Schwarz (2013)	1.8	154	-2	3.5
Tiwari et al. (2015)	1.4	159	-14	2
Tiwari & Nusser (2016)	0.9	151	-6	2.1
Colin et al. (2017)	1.2	149.1	-15.7	3
Bengaly et al. (2018)	2.3	147.45	-17.54	2.9
Siewert et al. (2021)	1.8	140.02	-5.14	3.5
CMB expectation	0.46	167.942	-6.944	



Dipole \sim 2-3 times larger than expectation (0.0046)

Similar direction to the CMB dipole.

Testing the cosmological principle

Wide-field Infrared Survey Explorer (WISE) systematically independent quasar catalog

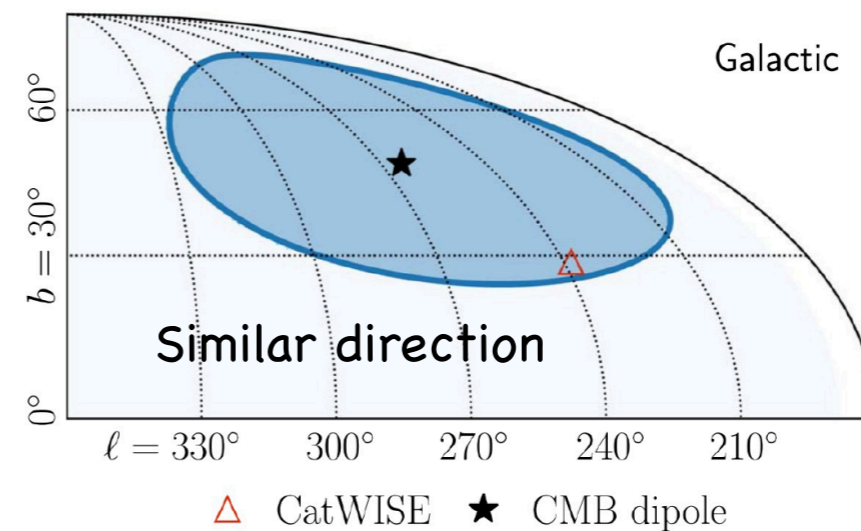
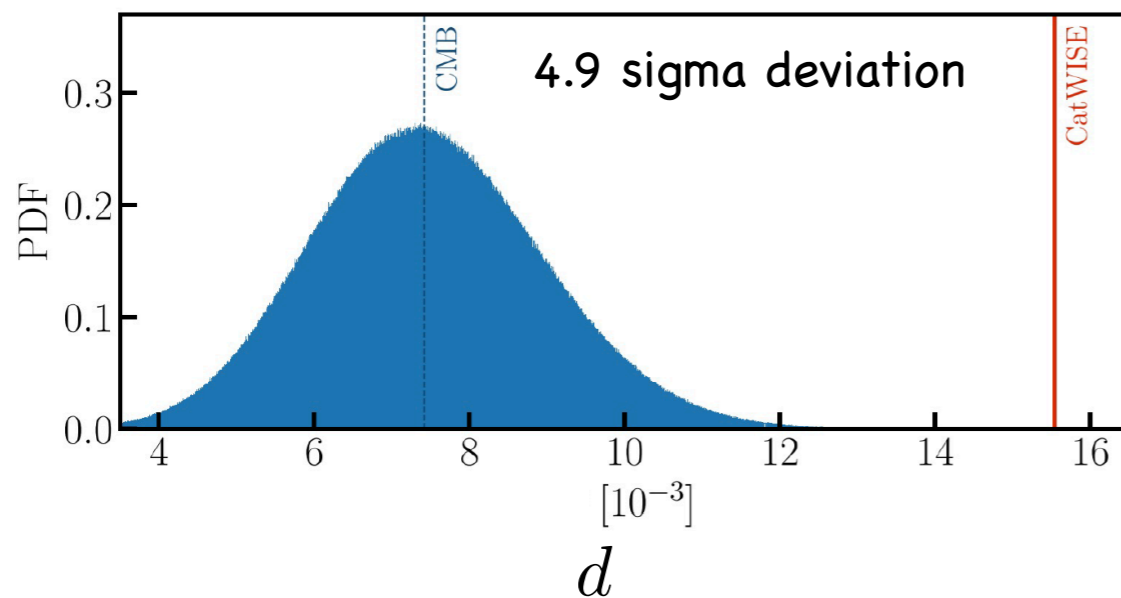
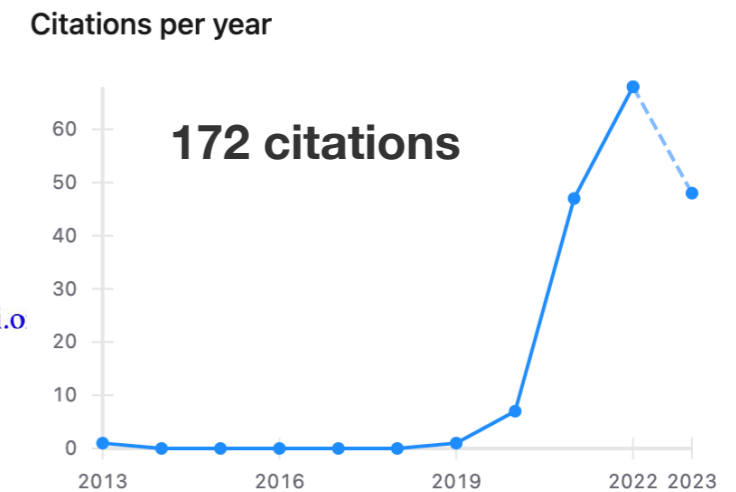
THE ASTROPHYSICAL JOURNAL LETTERS, 908:L51 (6pp), 2021 February 20

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A Test of the Cosmological Principle with Quasars

Nathan J. Secrest¹ , Sebastian von Hausegger^{2,3,4} , Mohamed Rameez⁵ , Roya Mohayaee³ , Subir Sarkar⁴ , and Jacques Colin³ 



$$n_i v_o^i = (2.66 \pm 0.29) \times 10^{-3} \Rightarrow 797 \pm 87 \text{ km/s}$$

更快的速度

arXiv > astro-ph > arXiv:2208.05018

Astrophysics > Cosmology and Nongalactic Astrophysics

[Submitted on 9 Aug 2022]

Anomalies in Physical Cosmology

[Phillip James E. Peebles](#)

I conclude that the present weight of the evidence from the other measures of the radio dipole and the WISE quasar dipole is that there is an anomalously large dipole common to distant radio galaxies and quasars.

How to explain the inconsistency?

我们生活在空洞附近?

arXiv > astro-ph > arXiv:2211.06857

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Astrophysics > Cosmology and Nongalactic Astrophysics

[Submitted on 13 Nov 2022]

Reconciling cosmic dipolar tensions with a gigaparsec void

Tingqi Cai, Qianhang Ding, Yi Wang

Recent observations indicate a 4.9σ tension between the CMB and quasar dipoles. This tension challenges the cosmological principle. We propose that if we live in a gigaparsec scale void, the CMB and quasar dipolar tension can be reconciled. This is because we are unlikely to live at the center of the void. And a 15% offset from the center will impact the quasars and CMB differently in their dipolar anisotropies. As we consider a large and thick void, our setup can also ease the Hubble tension.

宇宙轴心?

arXiv > astro-ph > arXiv:2209.14918

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Astrophysics > Cosmology and Nongalactic Astrophysics

[Submitted on 29 Sep 2022 (v1), last revised 8 Nov 2022 (this version, v2)]

Dipole Cosmology: The Copernican Paradigm Beyond FLRW

Chethan Krishnan, Ranjini Mondol, M. M. Sheikh-Jabbari

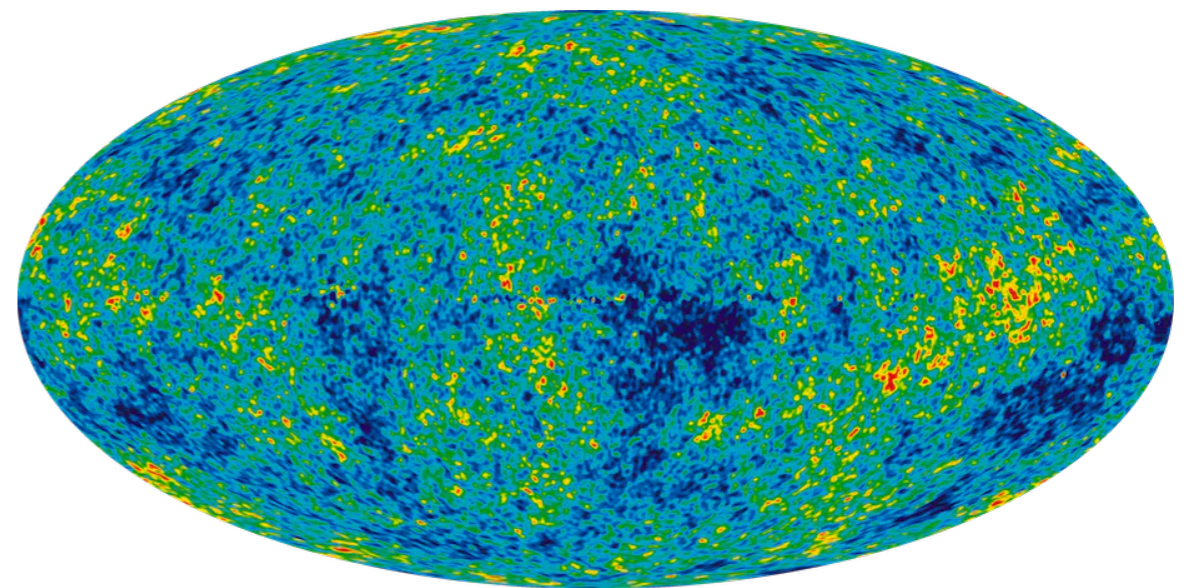
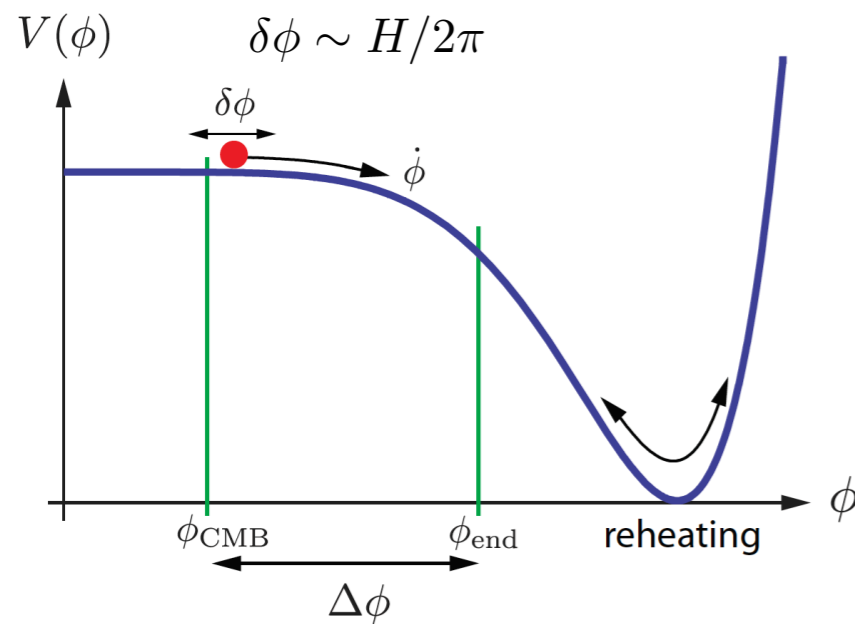
We introduce the *dipole cosmological principle*, the idea that the Universe is a maximally Copernican cosmology, compatible with a cosmic flow. It serves as the most symmetric paradigm that generalizes the FLRW ansatz, in light of the increasingly numerous (but still tentative) hints that have emerged in the last two decades for a non-kinematic component in the CMB dipole. Einstein equations in our "dipole cosmology" are still ordinary differential equations -- but instead of the two Friedmann equations, now we have four. The two new functions can be viewed as an anisotropic scale factor that breaks the isotropy group from $SO(3)$ to $U(1)$, and a "tilt" that captures the cosmic flow velocity. The result is an axially isotropic, tilted Bianchi V/VII_h cosmology. We assess the possibility of model building within the dipole cosmology paradigm, and discuss the dynamics of expansion rate, anisotropic shear and tilt, in various examples. A key observation is that the cosmic flow (tilt) can grow even while the anisotropy (shear) dies down. Remarkably, this can happen even in an era of late time acceleration.

How to explain the inconsistency?

Before giving up the cosmological principle, can we explain it from the **perturbed FRW**?



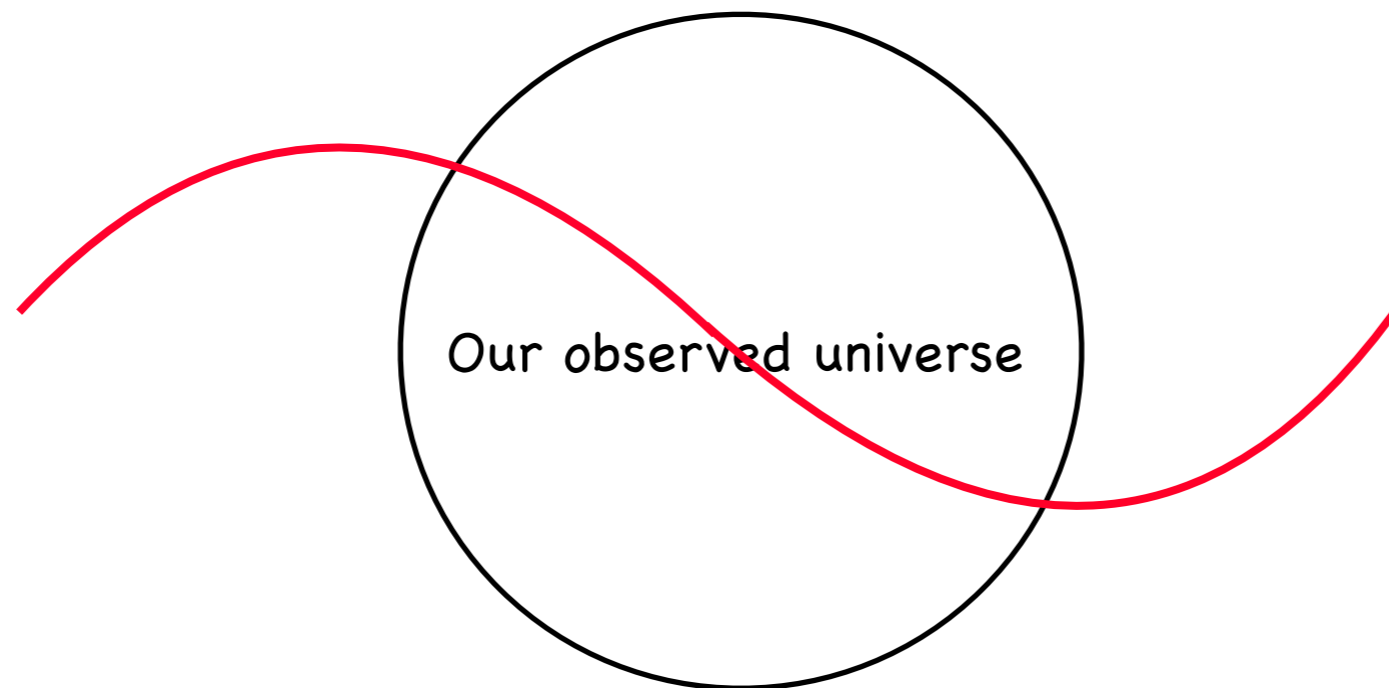
Anisotropies at CMB (commonly believed from the inflation)



Perturbations at super horizon scale

Inflation  Perturbations at super horizon scale

If we are living in a large super horizon mode, there may be a dipole



1785VA...22...125G

Long-wavelength perturbations of a Friedmann universe, and anisotropy of the microwave background radiation

L. P. Grishchuk and Ya. B. Zel'dovich

Shternberg Astronomical Institute, Moscow

(Submitted July 2, 1977)

Astron. Zh. **55**, 209–215 (March–April 1978)

PHYSICAL REVIEW D

VOLUME 44, NUMBER 12

15 DECEMBER 1991

倾斜的宇宙 **Tilted Universe** and other remnants of the preinflationary Universe

Michael S. Turner

*NASA/Fermilab Astrophysics Center, Fermi National Accelerator Laboratory, Batavia, Illinois 60510-0500
and Departments of Physics and Astronomy and Astrophysics, Enrico Fermi Institute, The University of Chicago,*

Dipole Anisotropy from an Entropy Gradient 1996'

David Langlois^{1,2} and Tsvi Piran¹

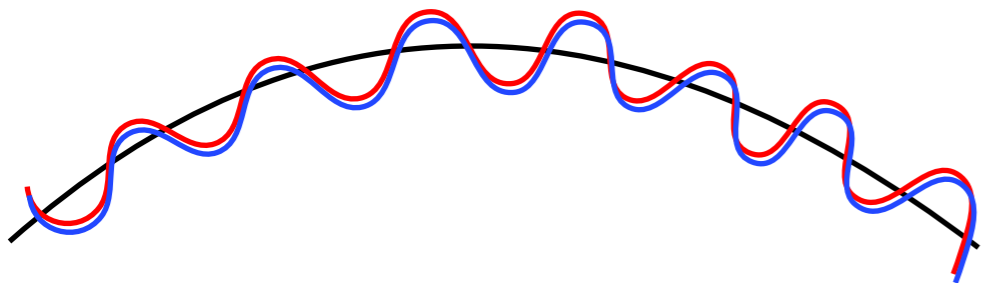
We can not observe this dipole from CMB if the perturbation is adiabatic

However, if there is entropy(isocurvature) mode at super horizon scale, an intrinsic dipole appears in CMB

Adiabatic/curvature vs entropy/isocurvature perturbation

dark matter and radiation share same fluctuation

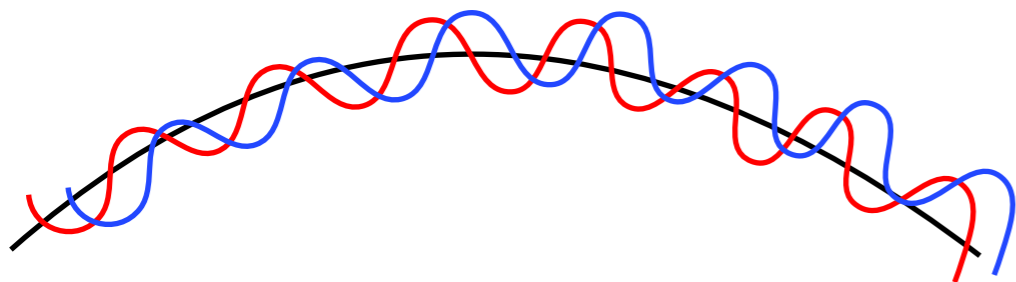
$$S = 0$$



$$S = \frac{3}{4} \frac{\delta\rho_r}{\rho_r} - \frac{\delta\rho_m}{\rho_m}$$

Single field inflation only generates adiabatic perturbation

$$S \neq 0$$



Thermal(for example, WIMP) dark matter only generates very tiny entropy perturbation

dark matter and radiation have different fluctuations

Comment on the isocurvature from thermal dark matter

Does freeze-in dark matter generate large isocurvature?

arXiv > astro-ph > arXiv:2210.15691v1

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Astrophysics > Cosmology and Nongalactic Astrophysics

[Submitted on 27 Oct 2022 (this version), latest version 3 Jul 2023 (v2)]

Dark matter freeze-in produces large post-inflationary isocurvature

Nicola Bellomo, Kim V. Berghaus, Kimberly K. Boddy

In this Letter, we show that the nonthermal nature of dark matter freeze-in production leads to large, totally correlated dark matter-photon isocurvature perturbations, which are imprinted in anisotropies of the cosmic microwave background (CMB). Isocurvature is typically expected from inflationary physics, but the isocurvature from freeze-in arises post inflation. We compute the freeze-in of millicharged dark matter, generated from electron-positron annihilations in the early Universe. We find that current CMB observations from `\textit{Planck}` exclude this scenario for dark matter masses between 1 MeV and 10 GeV at more than 2σ , whereas upcoming CMB experiments will have the sensitivity to reach at least the 4σ level. We anticipate any scenario in which dark matter is nonthermally produced to generically give rise to isocurvature. Our work opens a new avenue for exploring fundamental dark matter physics through its impact on cosmological observables.

被Riotto 和Strumia 写文反对

arXiv > hep-ph > arXiv:2211.08719

High Energy Physics – Phenomenology

[Submitted on 16 Nov 2022]

Freeze-in Dark Matter Perturbations are Adiabatic

Davide Racco, Antonio Riotto

arXiv > hep-ph > arXiv:2211.08359

High Energy Physics – Phenomenology

[Submitted on 15 Nov 2022]

Dark Matter from freeze-in and its inhomogeneities

Alessandro Strumia

We consider generic freeze-in processes for generation of Dark Matter, together with the consequent re-thermalization of the Standard Model fluid. We find that Dark Matter inherits the Standard Model adiabatic inhomogeneities on the cosmological scales probed by current observations, that were super-horizon during freeze-in. Thereby, freeze-in satisfies the bounds on iso-curvature perturbations.

Comment on the isocurvature from thermal dark matter

arXiv > astro-ph > arXiv:2311.17164

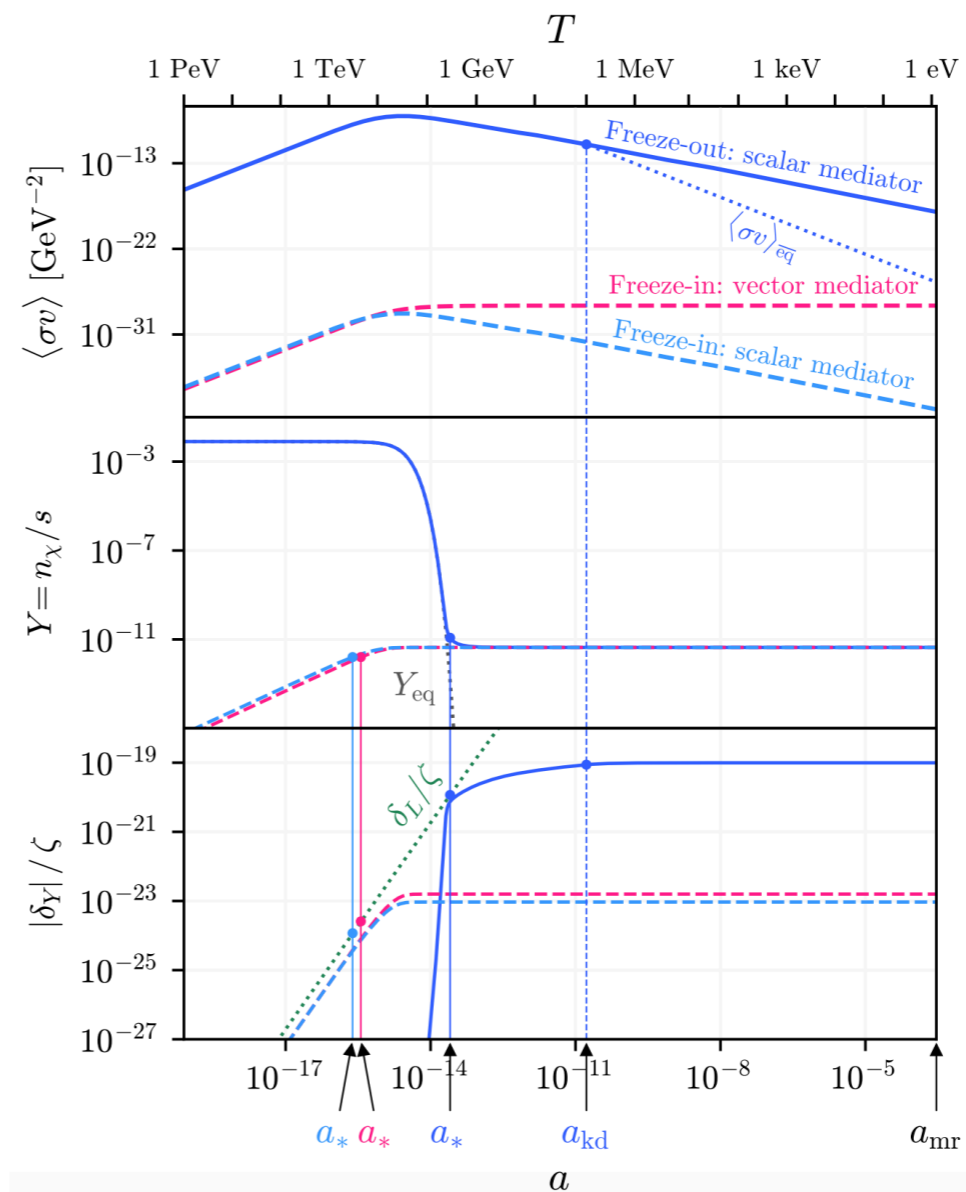
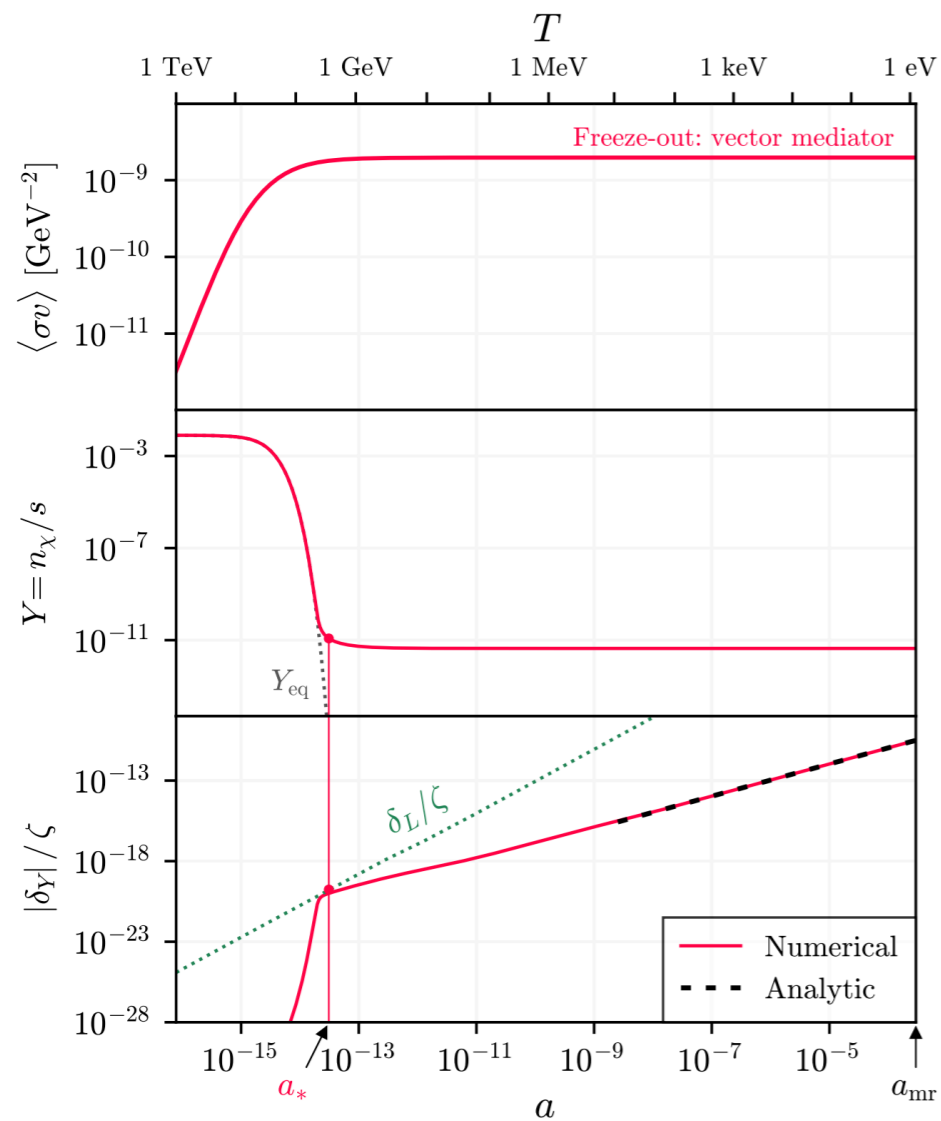
Astrophysics > Cosmology and Nongalactic Astrophysics

[Submitted on 28 Nov 2023]

Dark Matter Isocurvature from Curvature

Ian Holst, Wayne Hu, Leah Jenks

Thermal dark matter only generate very tiny isocurvature perturbation!



One solution to the cosmic dipole problem

CMB dipole

$$D_1^{\text{CMB}} = (1.23357 \pm 0.00036) \times 10^{-3}$$

$$n_i v_o^i = 369.82 \pm 0.11 \text{ km/s}$$

Galaxy number count dipole

$$d_{\mathcal{N}} = (15.54 \pm 1.7) \times 10^{-3}$$

$$n_i v_o^i = (2.66 \pm 0.29) \times 10^{-3} \Rightarrow 797 \pm 87 \text{ km/s}$$

If there is **intrinsic dipole** in CMB, it cancels part of kinematic dipole

$$d^{\text{CMB}} = d_{\text{kin}}^{\text{CMB}} + D_1^{\text{CMB}} = 1.23357 \times 10^{-3}$$

$D_1^{\text{CMB}} > 8 \times 10^{-4}$ to explain the cosmic dipole problem

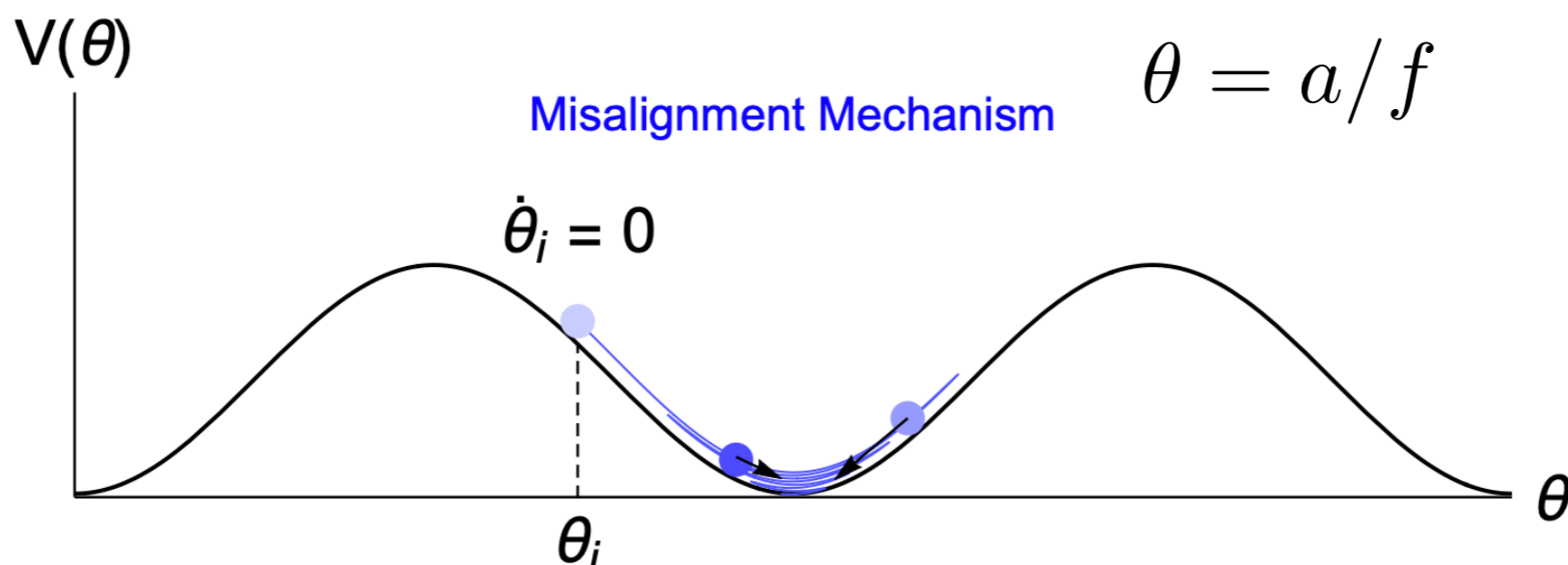
What is the origin of the isocurvature/entropy mode?

Obvious thermal dark matter can not work

Axion dark matter is one of the candidate

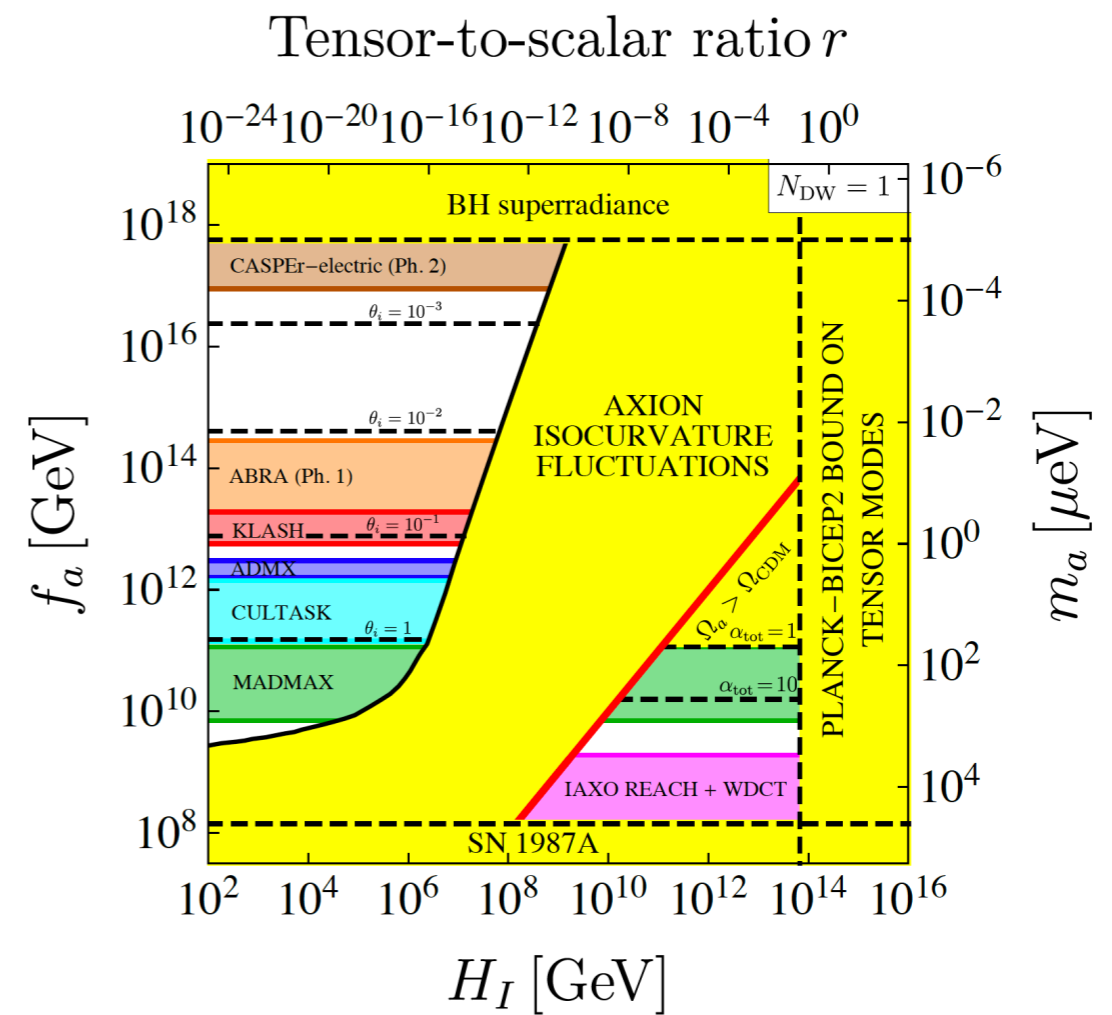
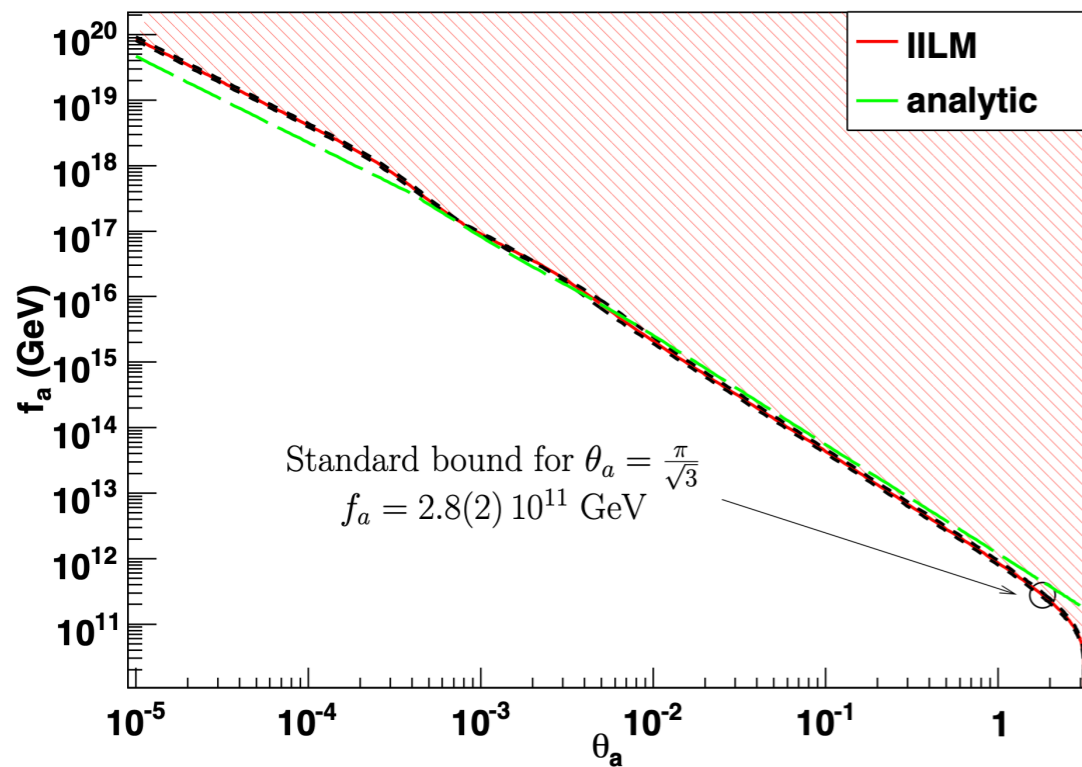
$$V(\Phi) = \lambda(\Phi\Phi^\dagger - f^2/2)^2 \quad \Phi = \frac{1}{\sqrt{2}}\varphi \exp(i\frac{a}{f})$$

$$\rho = \frac{1}{2}m_a^2 f^2 \theta_0^2 \quad V = \Lambda^4(1 - \cos \frac{a}{f})$$



What is the origin of the isocurvature mode?

For theta around $O(0.1-1)$ and axion be the dark matter $f_a \sim 10^{11-14}$ GeV



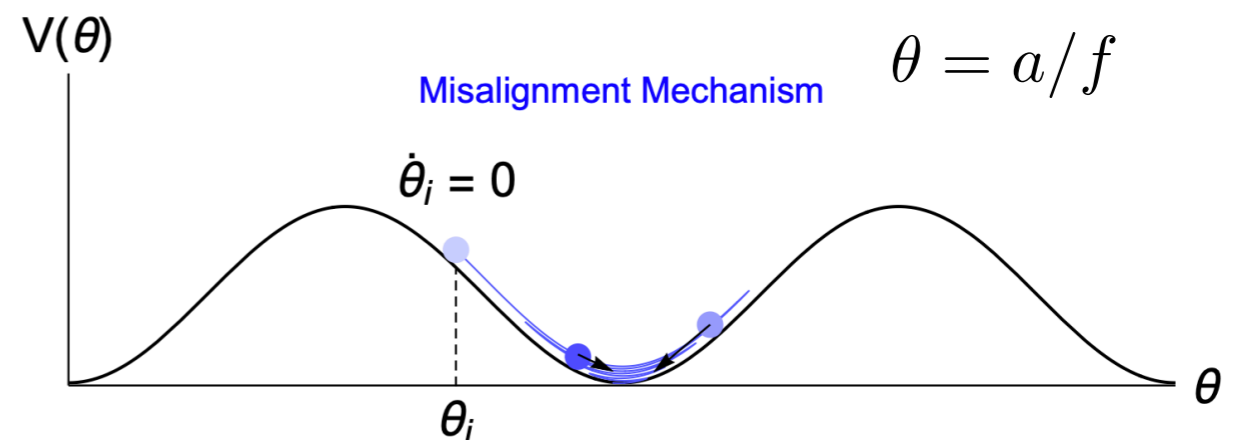
The landscape of QCD axion models, L. Luzio, M. Giannotti, E. Nardi, L. Visinelli

Axion dark matter

During inflation

$$\delta\theta = \frac{H}{2\pi\varphi} \quad \rho = \frac{1}{2}m_a^2 f^2 \theta_0^2$$

$$\delta\rho/\rho = \frac{H}{\pi\varphi\theta_0}$$



Limit on the large isocurvature from CMB for theta $O(1)$

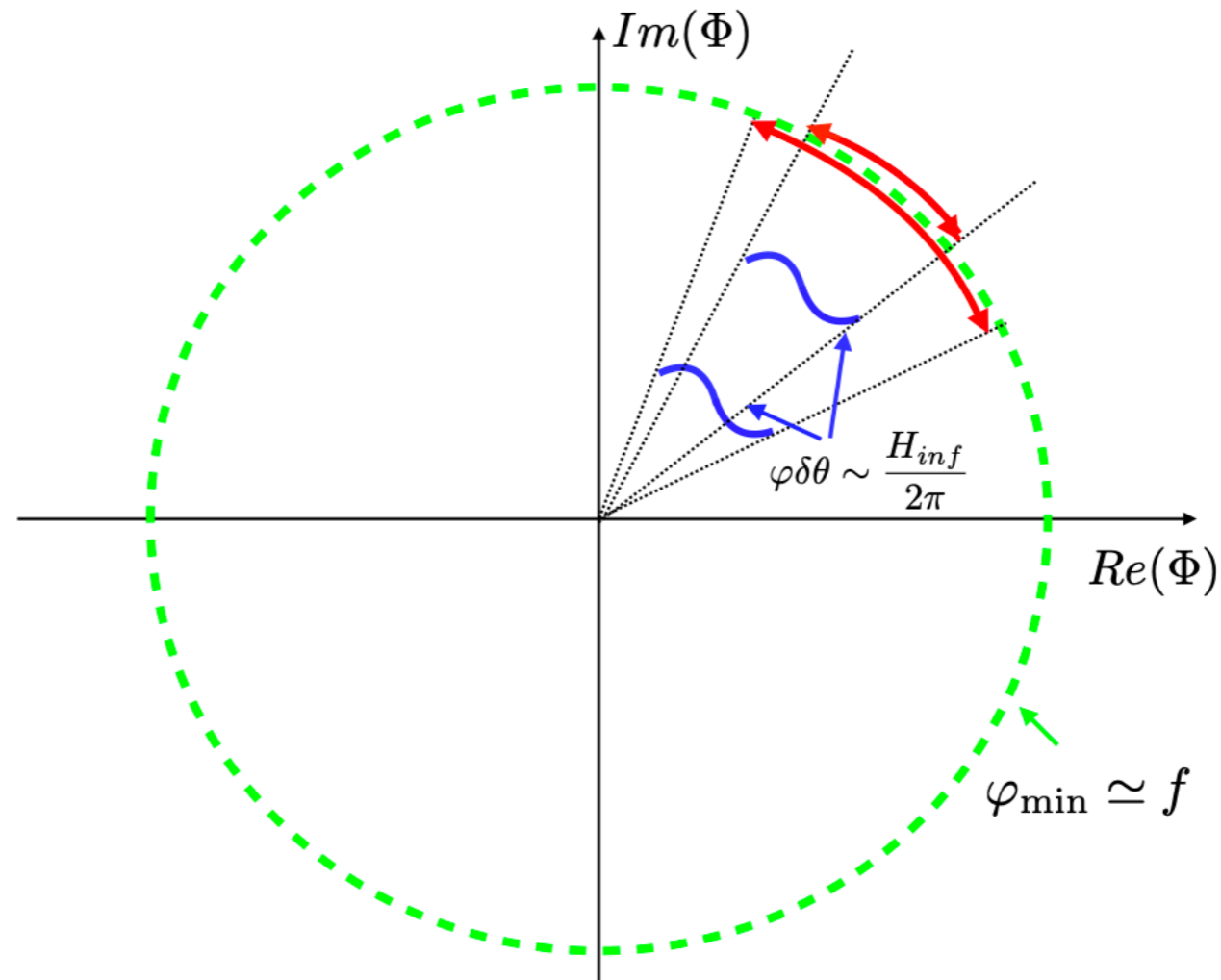
$$\frac{H^2}{\pi^2 f_a^2} < 10^{-10} \quad H/f_a < 10^{-5}$$

It is too small to explain the dipole problem by axion

Axion dark matter

If the radial mode vary in the early universe(during inflation)

$$\mathcal{P}_S^{1/2} = \frac{H_{inf}}{\pi\varphi(k)\theta}$$

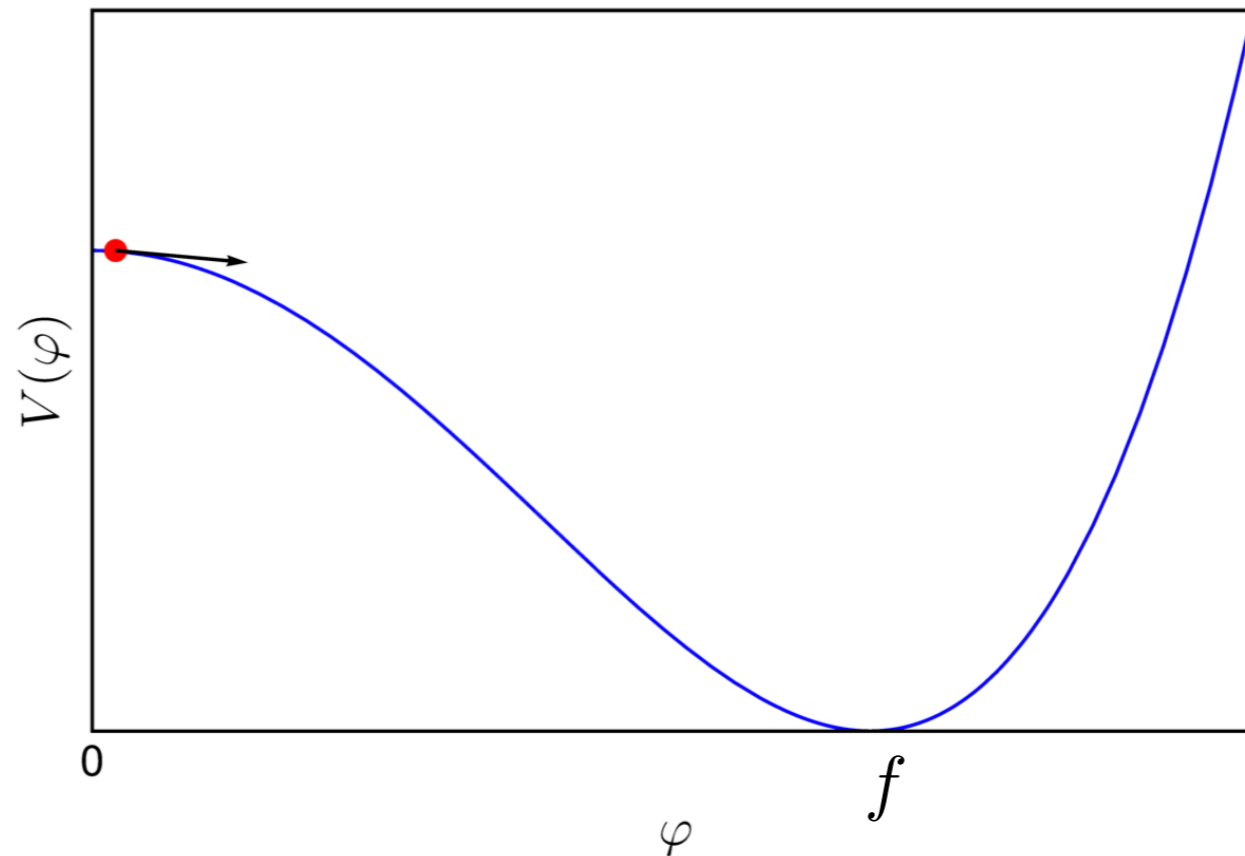


φ from a small value around H to a large value f

Potential

$$V(\Phi) = \lambda(\Phi\Phi^\dagger - f^2/2)^2$$

$$\Phi = \frac{1}{\sqrt{2}}\varphi \exp(i\frac{a}{f})$$



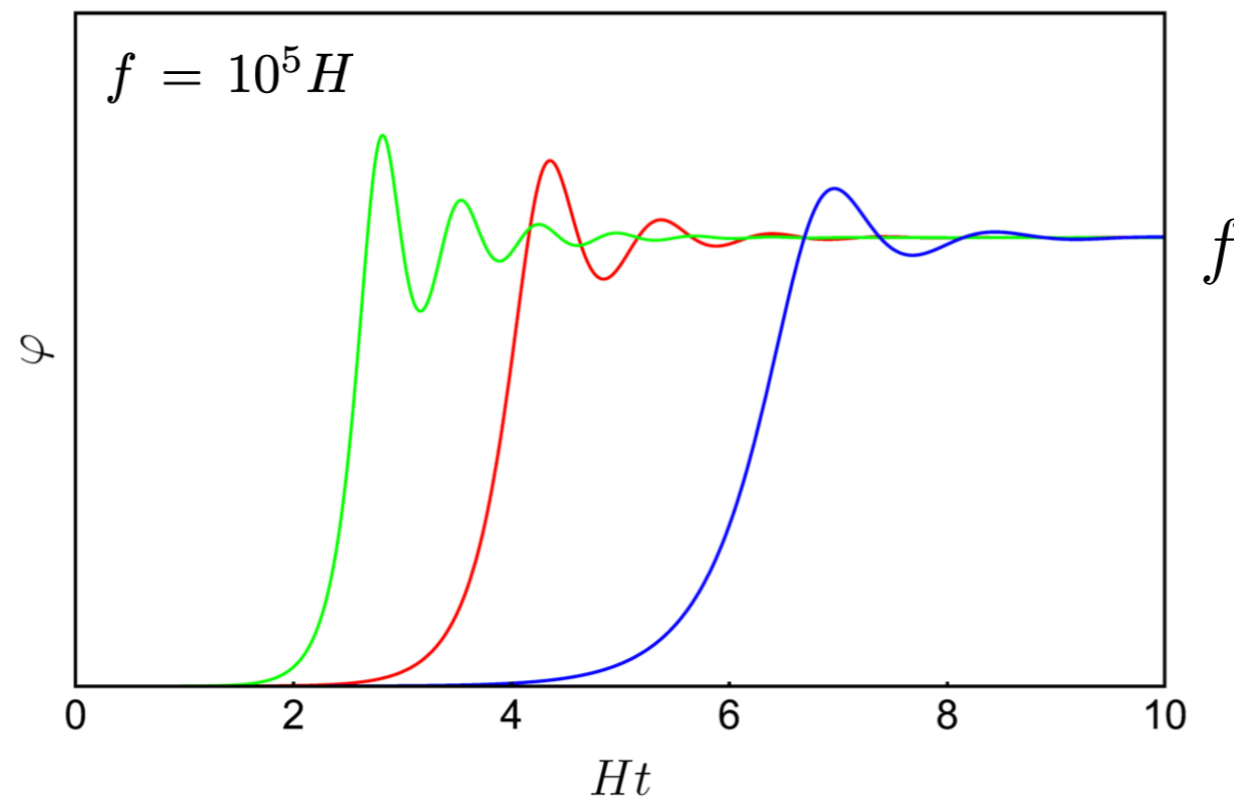
$$\delta\rho/\rho = \frac{H}{\pi\varphi}$$

initial phi should around H

A model of large isocurvature

$$\ddot{\varphi} + 3H\dot{\varphi} + V'(\varphi) = 0$$

$$\lambda = 10^{-9}, 2 \times 10^{-9}, 4 \times 10^{-9}$$

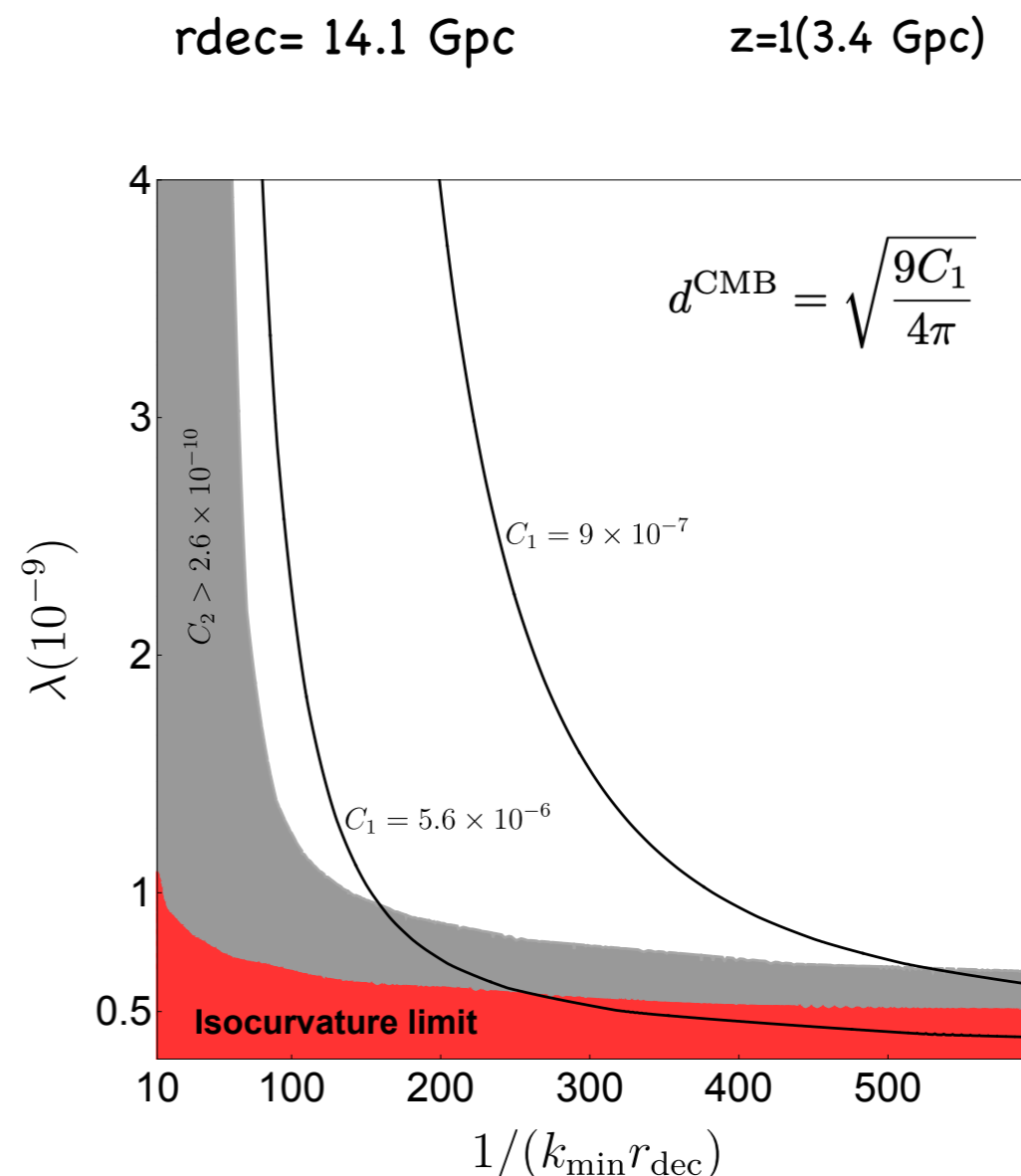


Small lambda, longer stay at small value

Axion model to explain dipole problem

Only two parameters

- (1) when to start?(super horizon size)
- (2) how long it takes?(lambda)



$$V(\Phi) = \lambda(\Phi\Phi^\dagger - f^2/2)^2 \quad 7 \times 10^{-10} < \lambda < 6.6 \times 10^{-5}$$

Axion model to explain dipole problem

THE ASTROPHYSICAL JOURNAL LETTERS, 908:L51 (6pp), 2021 February 20

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<https://doi.org/10.3847/2041-8213/abdd40>



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A Test of the Cosmological Principle with Quasars

Nathan J. Secrest¹ , Sebastian von Hausegger^{2,3,4} , Mohamed Rameez⁵ , Roya Mohayaee³ , Subir Sarkar⁴ , and Jacques Colin² 

Request to give a webinar over zoom on your recent work 2211.06912

Rameez 发送给 韩成成

Dear Prof Chengcheng Han

Your recent paper on "QCD axion dark matter and the cosmic dipole anomaly" is very interesting. I would be very grateful if you could spare some time to tell us about it over a zoom webinar.

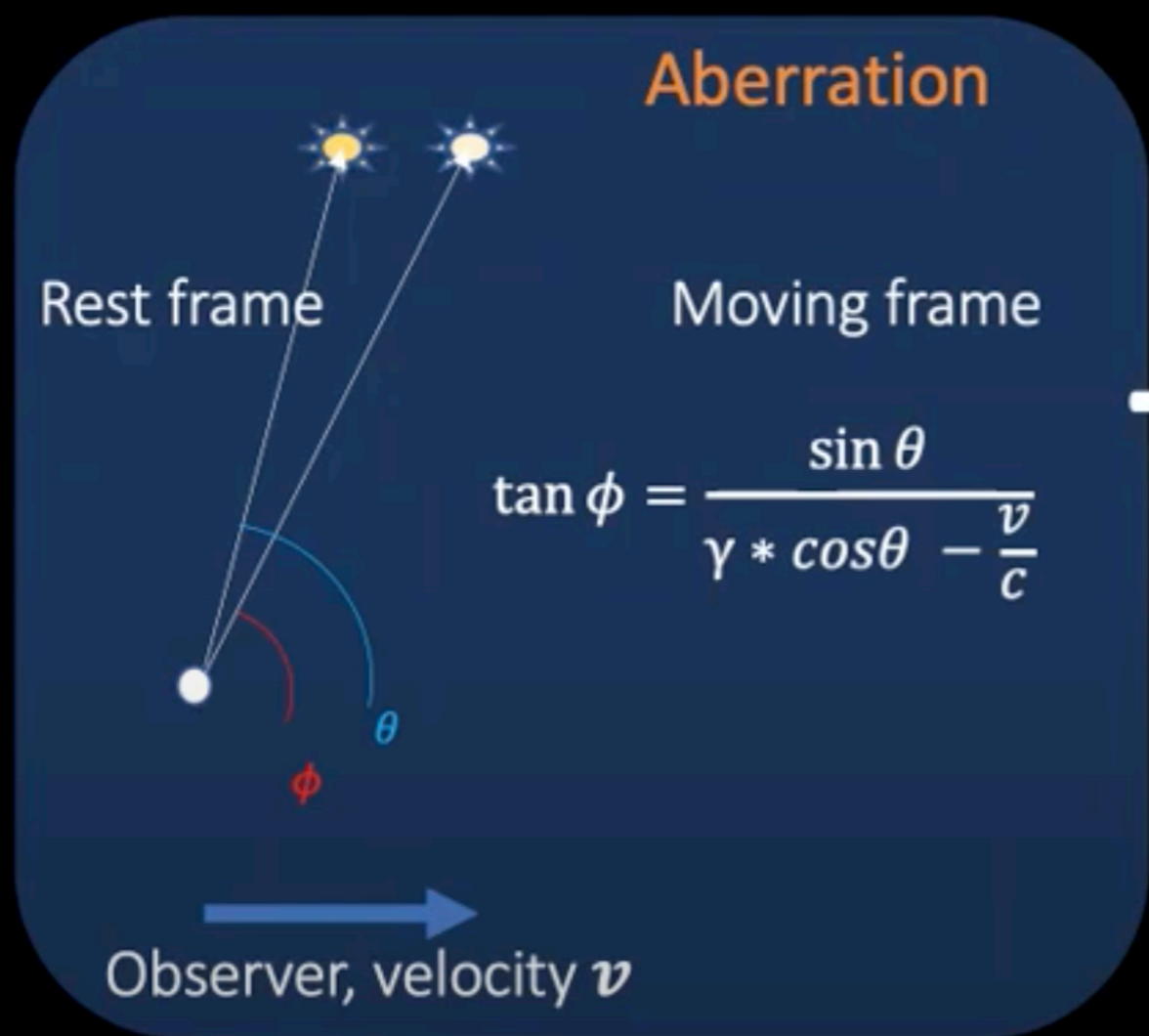
- Recently a cosmic dipole problem is reported
- If it originates from dark matter isocurvature, thermal dark matter is not favored
- The dipole may point to the first evidence of axion

The real reason, though, for our adherence here to the Cosmological Principle is not that it is surely correct, but rather, that it allows us to make use of the extremely limited data provided to cosmology by observational astronomy. If we make any weaker assumptions, as in the anisotropic or hierarchical models, then the metric would contain so many undetermined functions (whether or not we use the field equations) that the data would be hopelessly inadequate to determine the metric. On the other hand, by adopting the rather restrictive mathematical framework described in this chapter, we have a real chance of confronting theory with observation. If the data will not fit into this framework, we shall be able to conclude that either the Cosmological Principle or the Principle of Equivalence is wrong. Nothing could be more interesting.

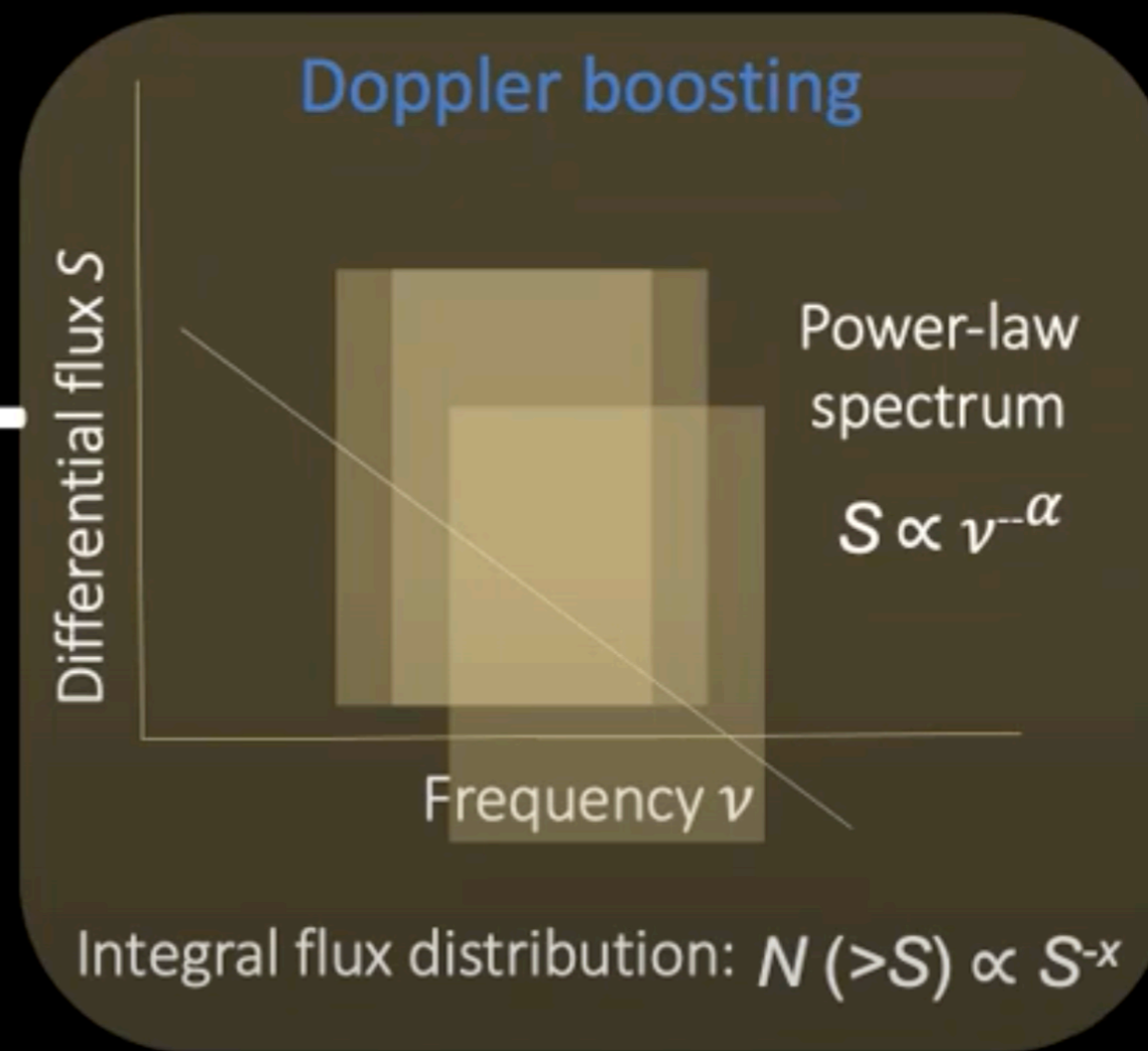
Steven Weinberg, Gravitation and Cosmology (1972)

IF THE DIPOLE IN THE CMB IS DUE TO OUR MOTION WRT THE 'CMB FRAME'
 THEN WE SHOULD SEE *SIMILAR* DIPOLE IN THE DISTRIBUTION OF DISTANT SOURCES

$$\sigma(\theta)_{obs} = \sigma_{rest} \left[1 + \left[2 + x(1 + \alpha) \right] \frac{v}{c} \cos(\theta) \right]$$



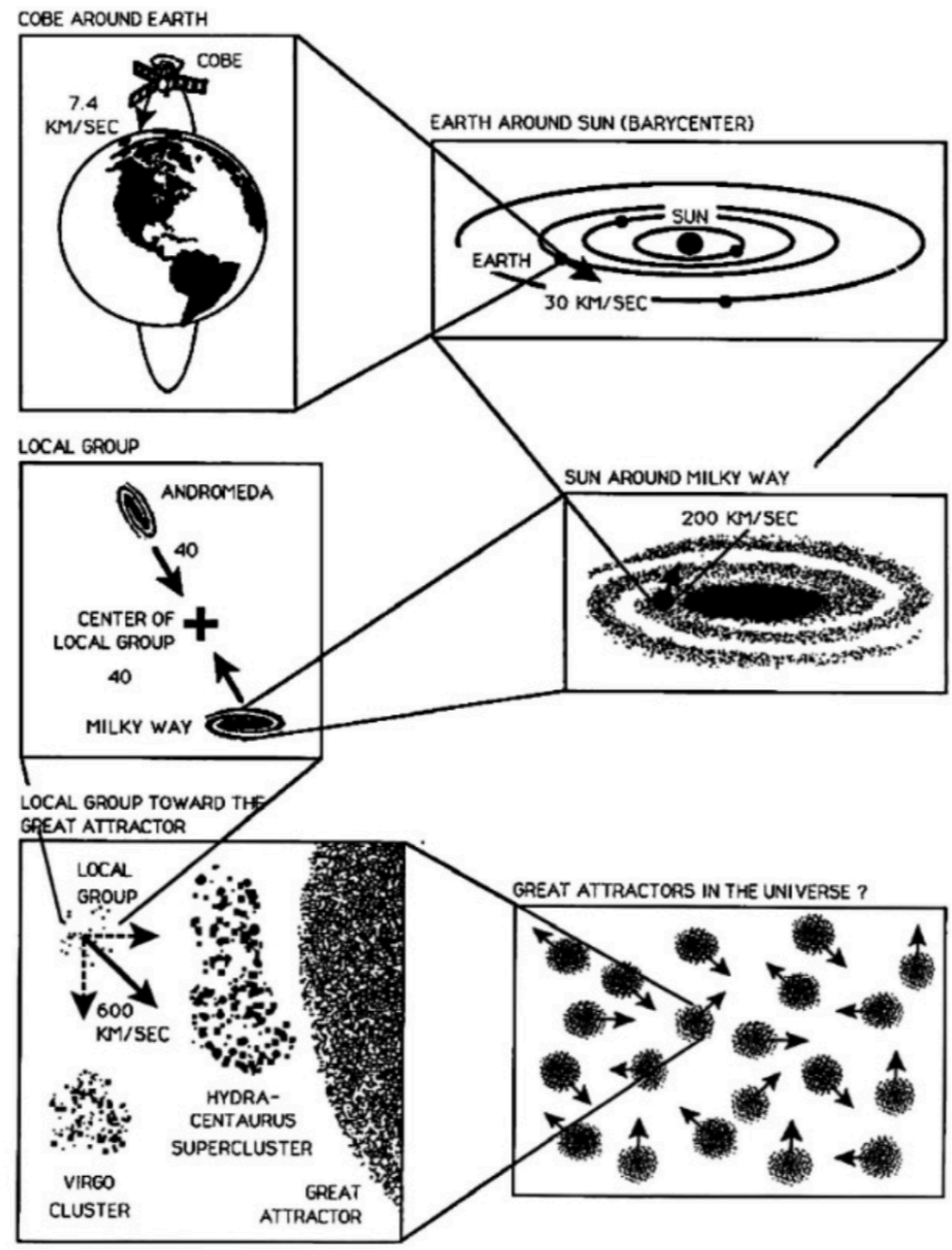
+

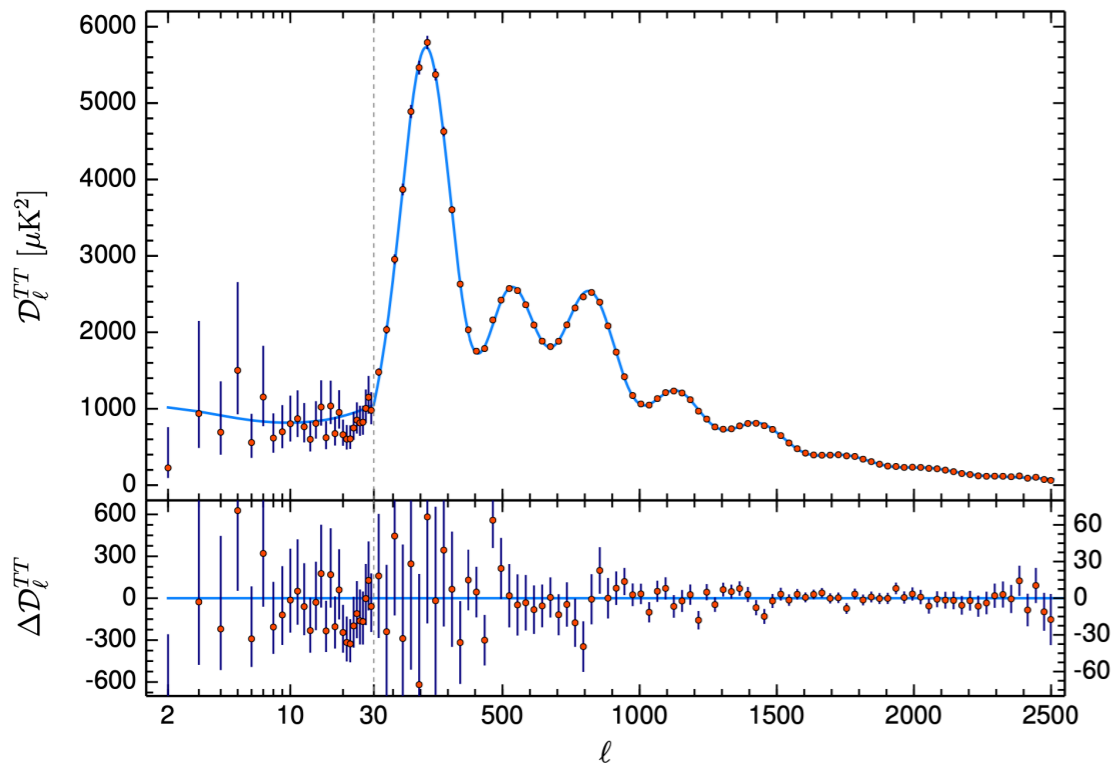
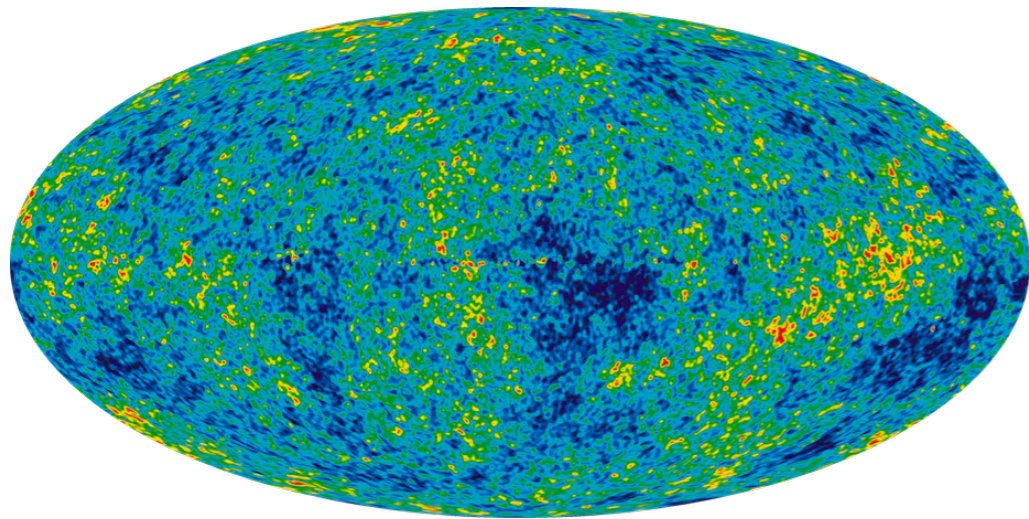


Flux-limited catalog \rightarrow *more* sources in direction of motion

Ellis & Baldwin,
 MNRAS 206:377,1984

VELOCITY COMPONENTS OF THE OBSERVED CMB DIPOLE





$$\frac{\Delta T}{T}(\hat{n}) = \sum_{l,m} a_{lm} Y_{lm}(\hat{n})$$

$$a_{lm} = \int d\Omega \frac{\Delta T}{T}(\hat{n}) Y_{lm}^*(\hat{n})$$

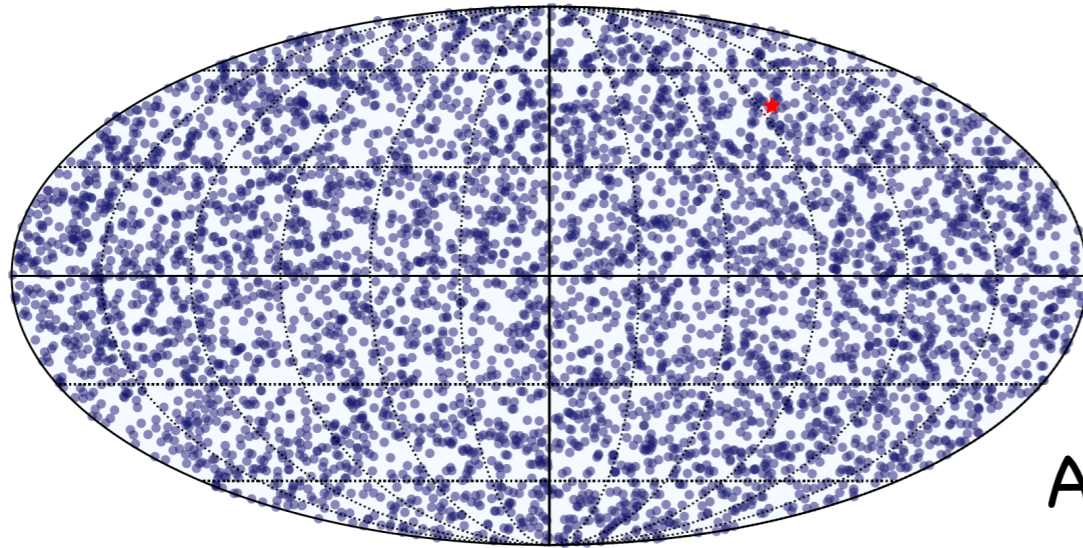
$$C_l = \frac{1}{2l+1} \sum_m \langle a_{lm}^* a_{lm} \rangle$$

$l=0,1,2,3\dots$ monopole, dipole, quadrupole...

$$\mathcal{D}_l = \frac{l(l+1)}{2\pi} C_l$$

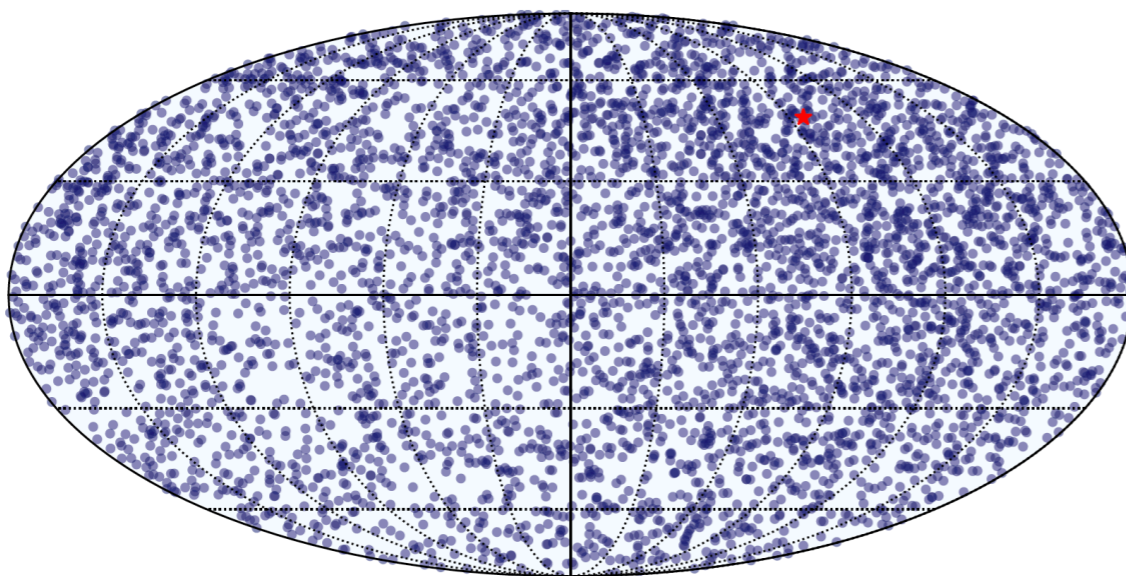
Aberration & Doppler boosting

Galaxies / quasars in CMB "rest frame"

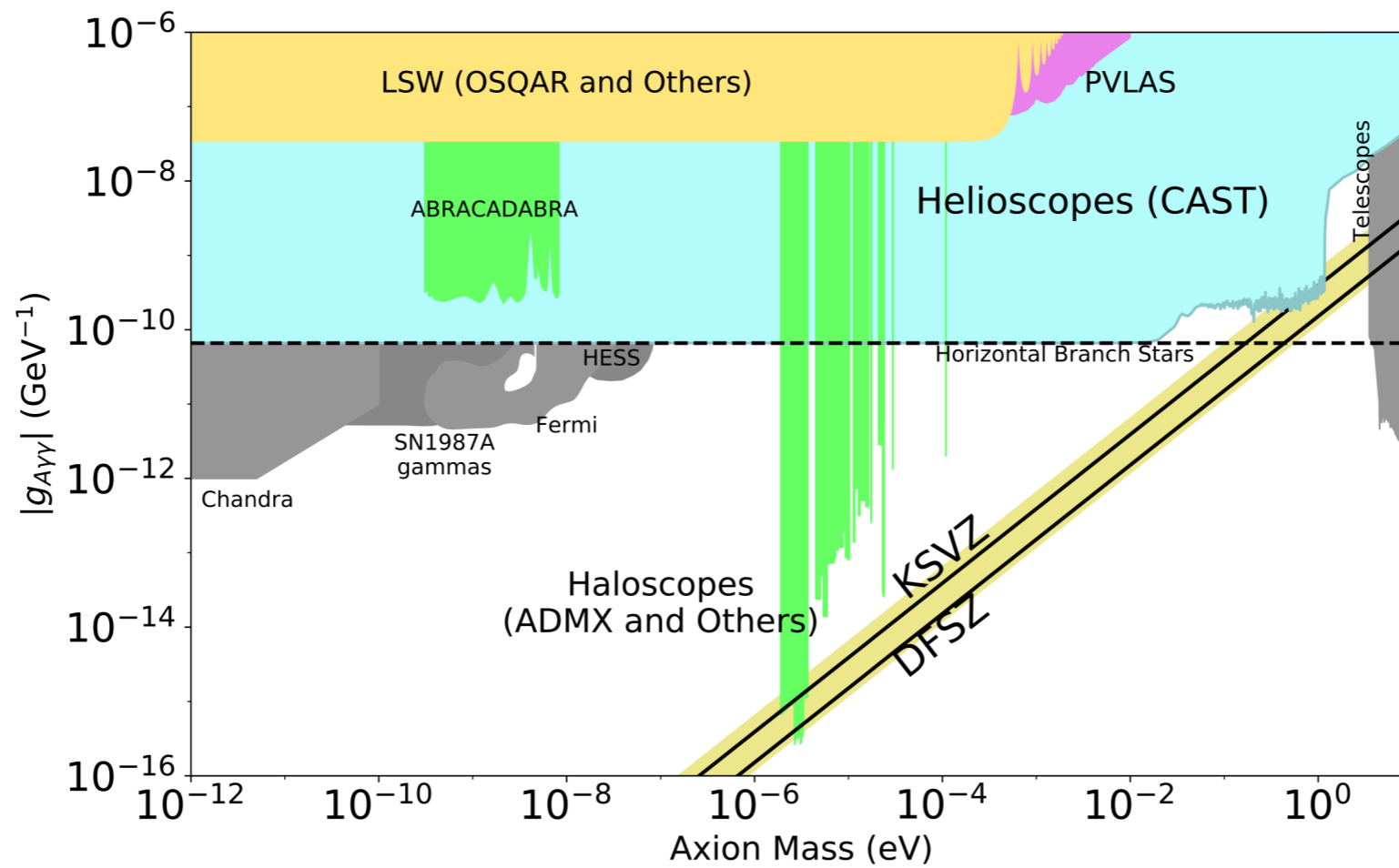


Aberration: object positions compressed in direction of motion

Doppler boosting: too-faint objects boosted into catalog flux limit



From Nathan Secrest



$$g_{A\gamma\gamma} = \frac{\alpha}{2\pi f_A} \left(\frac{E}{N} - 1.92(4) \right) \quad m_A = 5.691(51) \left(\frac{10^9 \text{ GeV}}{f_A} \right) \text{ meV}$$

One solution to the dipole problem

arXiv > astro-ph > arXiv:2207.01569

Astrophysics > Cosmology and Nongalactic Astrophysics

[Submitted on 4 Jul 2022]

Galaxy number-count dipole and superhorizon fluctuations JCAP 10 (2022) 019

Guillem Domènech, Roya Mohayaee, Subodh P. Patil, Subir Sarkar

Initial conditions Size of mode q	Adiabatic discrete mode	Isocurvature discrete mode
Superhorizon ($q < \mathcal{H}_0$)	No CMB dipole* [41] No NC dipole* Cannot solve dipole tension	Intrinsic CMB dipole [41] No NC dipole* Might resolve dipole tension**
Slightly subhorizon ($\mathcal{H}_0 \lesssim q \lesssim \mathcal{H}_{\text{dec}}$)	Amplitude $\lesssim 8 \times 10^{-5}$ (CMB [79]) $\mathcal{O}(10^{-3})$ maximum NC dipole Cannot solve dipole tension	Amplitude $\lesssim 10\%$ of adiabatic [79] $\mathcal{O}(10^{-4})$ maximum NC dipole Cannot solve dipole tension
Subhorizon ($q \gtrsim \mathcal{H}_{\text{dec}}$)	Amplitude $\sim 5 \times 10^{-5}$ [79] Cannot solve dipole tension [20]	Amplitude $\lesssim 10\%$ of adiabatic [79] Cannot solve dipole tension

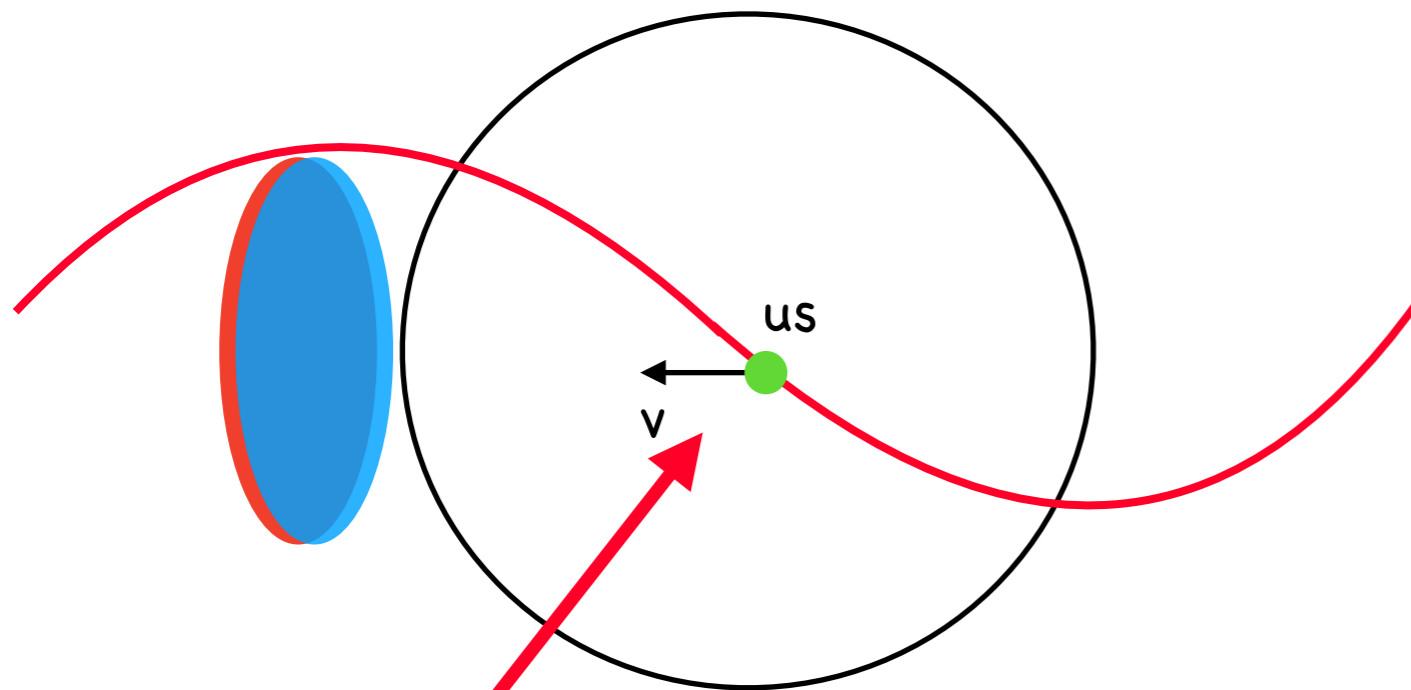
Considering a single mode of isocurvature to avoid multipole limit

The isocurvature mode should be large $\mathcal{O}(0.1-1)$

Perturbations at super horizon scale

Adiabatic perturbation

Our observed universe



CMB: $T < T_{\text{average}}$

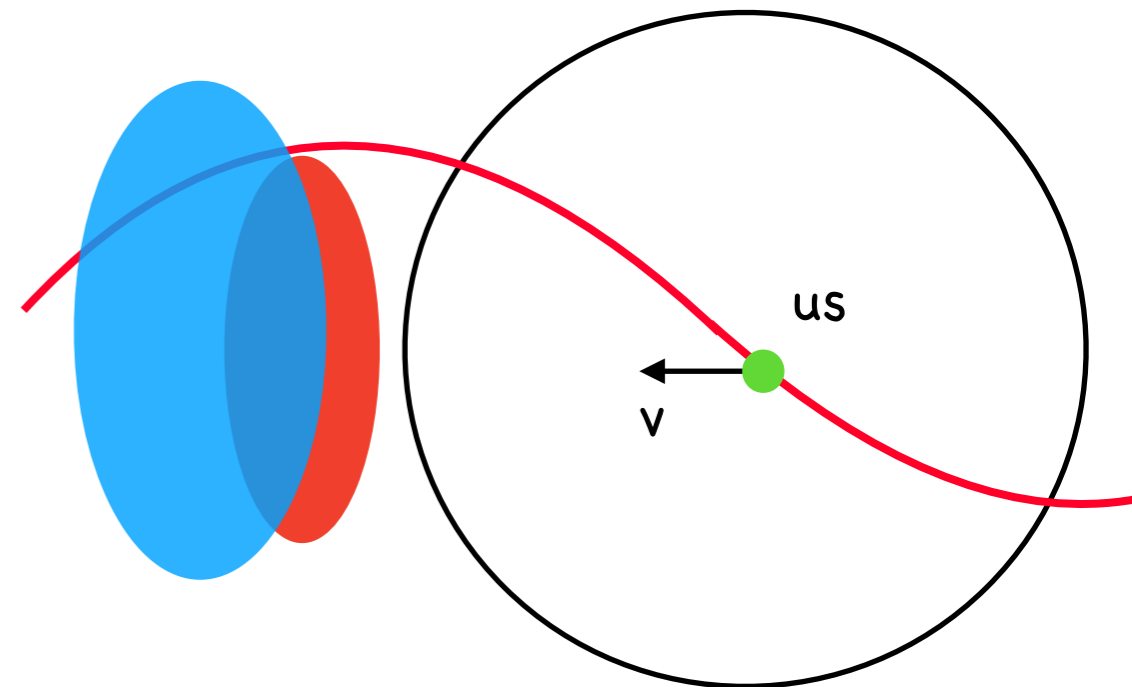
Pulling us, we see $T > T_{\text{average}}$

Two effect just cancel exactly!

Cancellation also happen for galaxy number count!

Entropy perturbation

Our observed universe



Can not cancel exactly in CMB

No effect on galaxy number count

Intrinsic dipole in CMB