

Probing the nuclear structure with relativistic heavy ion collisions

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Exploring nuclear physics across energy scales 2024

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04/21/2024

Landscape of nuclear physics

degrees of freedom

**quarks
& gluons**



quarks, gluons

Energy
(MeV)

940
neutron mass



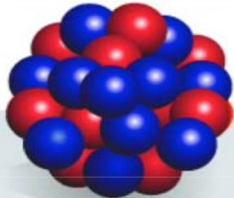
constituent quarks

hadrons



baryons, mesons

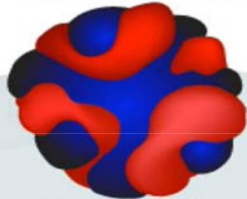
140
pion mass



protons, neutrons

8
proton separation
energy in lead

nuclei



nucleonic densities
and currents

1.32
vibrational
state in tin



collective coordinates

0.043
rotational
state in uranium

Landscape of nuclear physics

degrees of freedom

quarks
& gluons



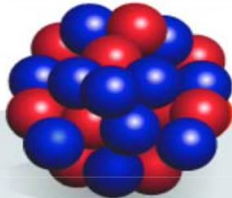
quarks, gluons



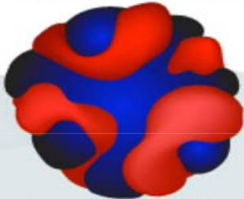
constituent quarks



baryons, mesons



protons, neutrons



nucleonic densities
and currents



collective coordinates

hadron
s

nuclei

Energy
(MeV)

940
neutron mass

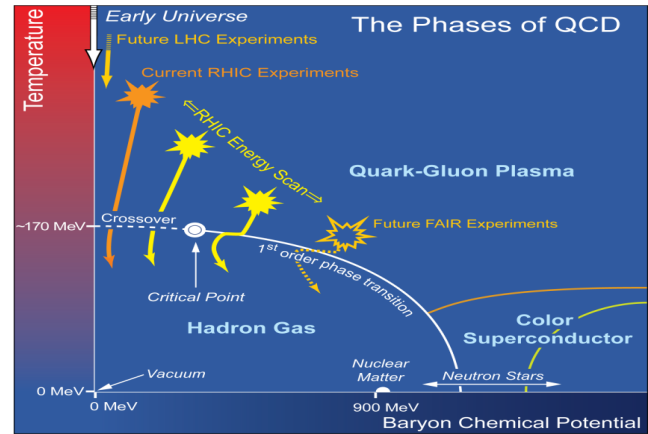
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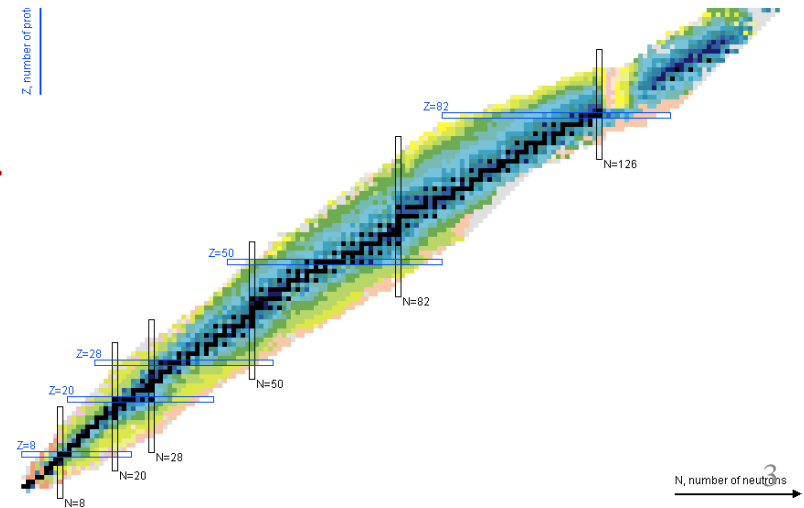
1.32
vibrational
state in tin

0.043
rotational
state in uranium

- intermediate and high
energy nuclear physics

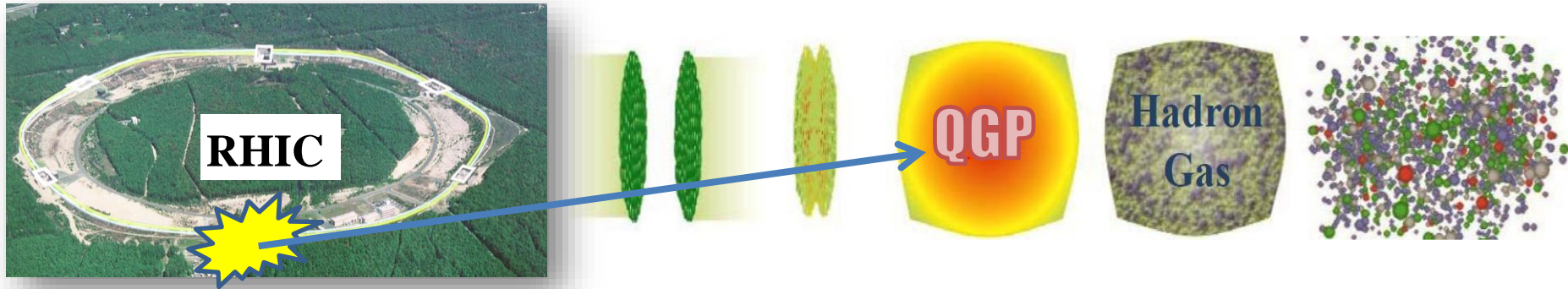


- nuclear structure



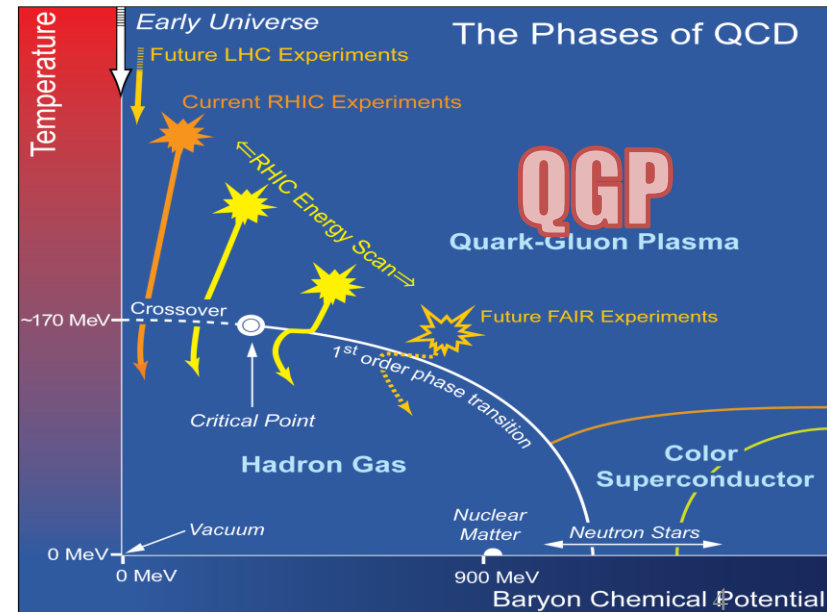
N, number of neutrons

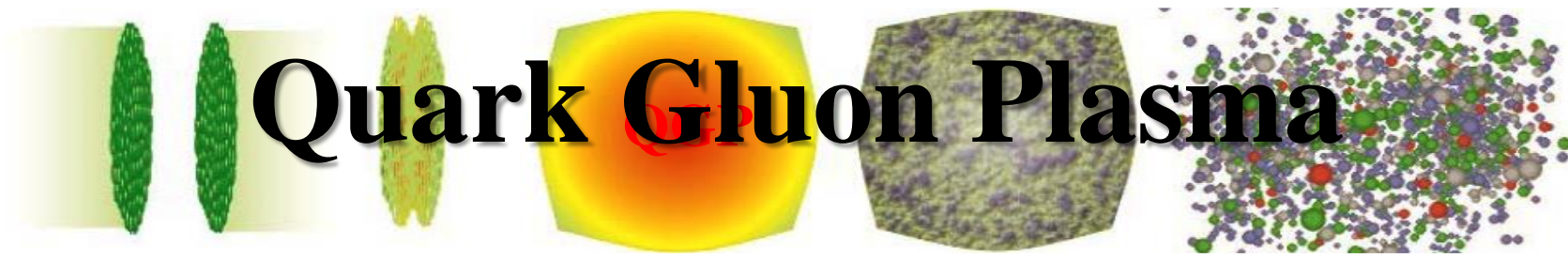
Relativistic Heavy Ion Collisions



Relativistic heavy ion collisions

- create and study QGP
- the QCD phase diagram
- the deconfinement & chiral phase transition
- the QCD vacuum



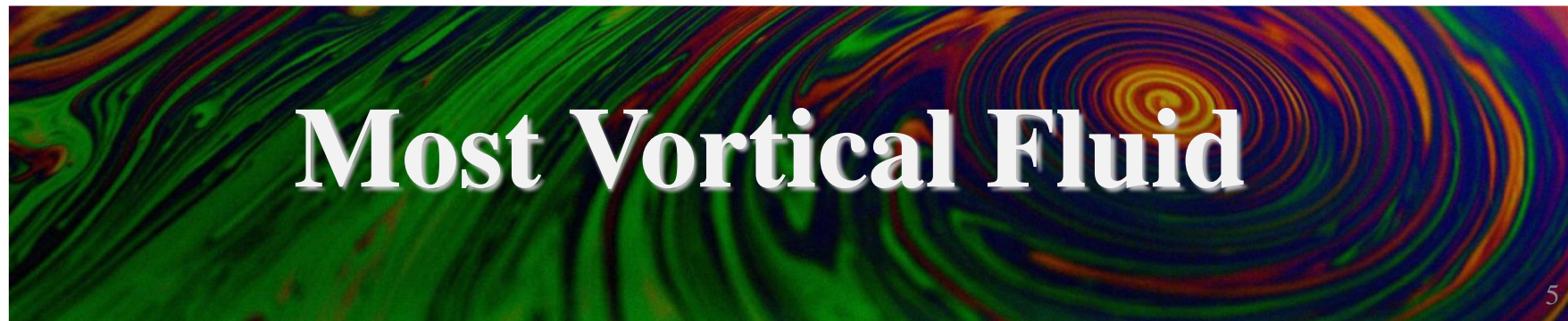


Quark Gluon Plasma

Hottest Matter on Earth

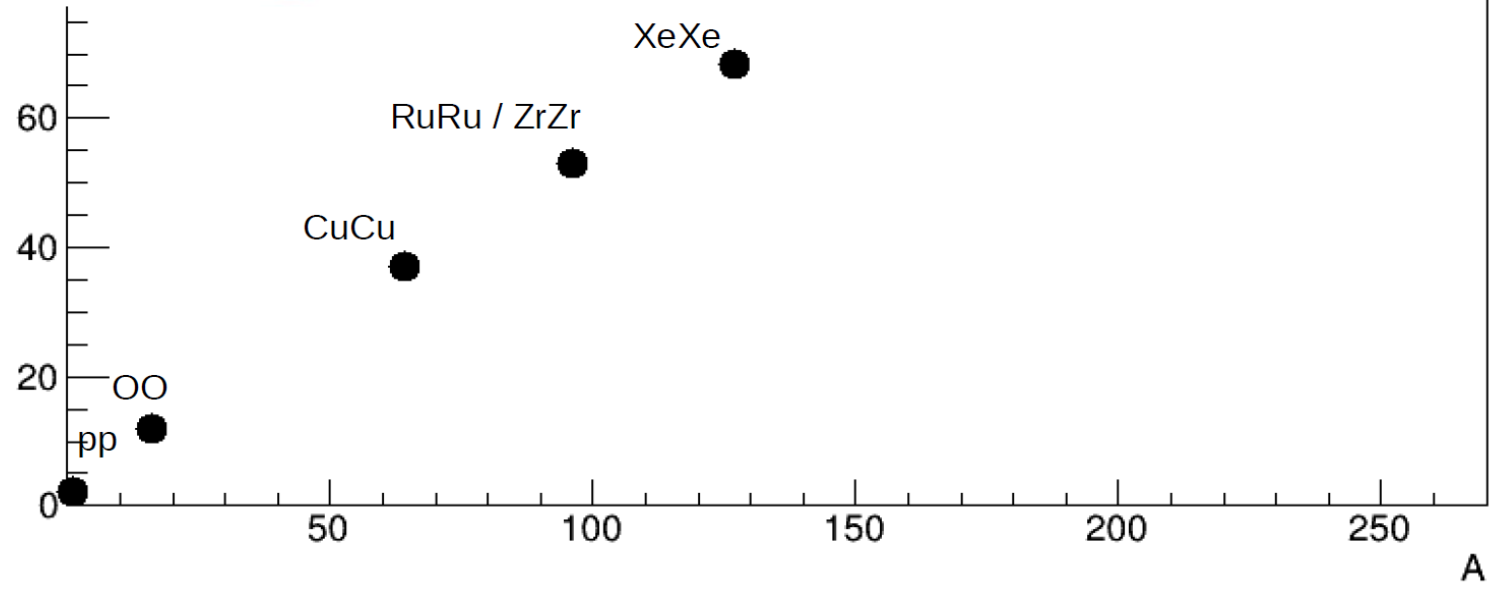
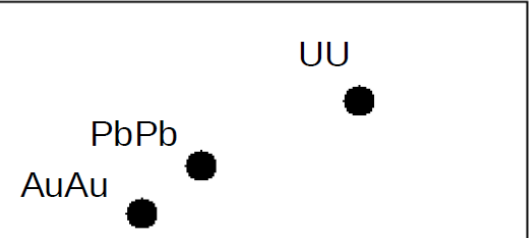
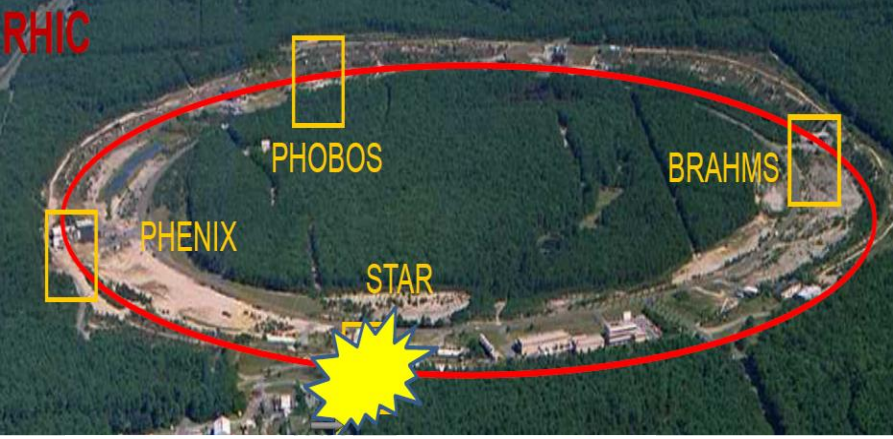


Most Perfect Liquid

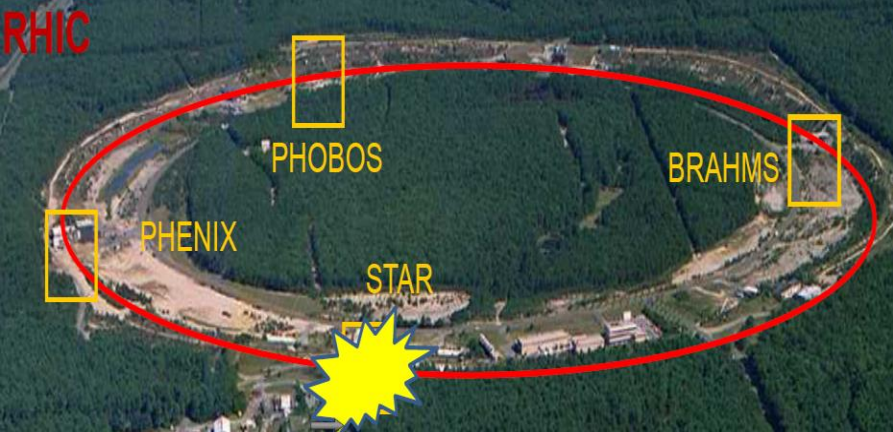


Most Vortical Fluid

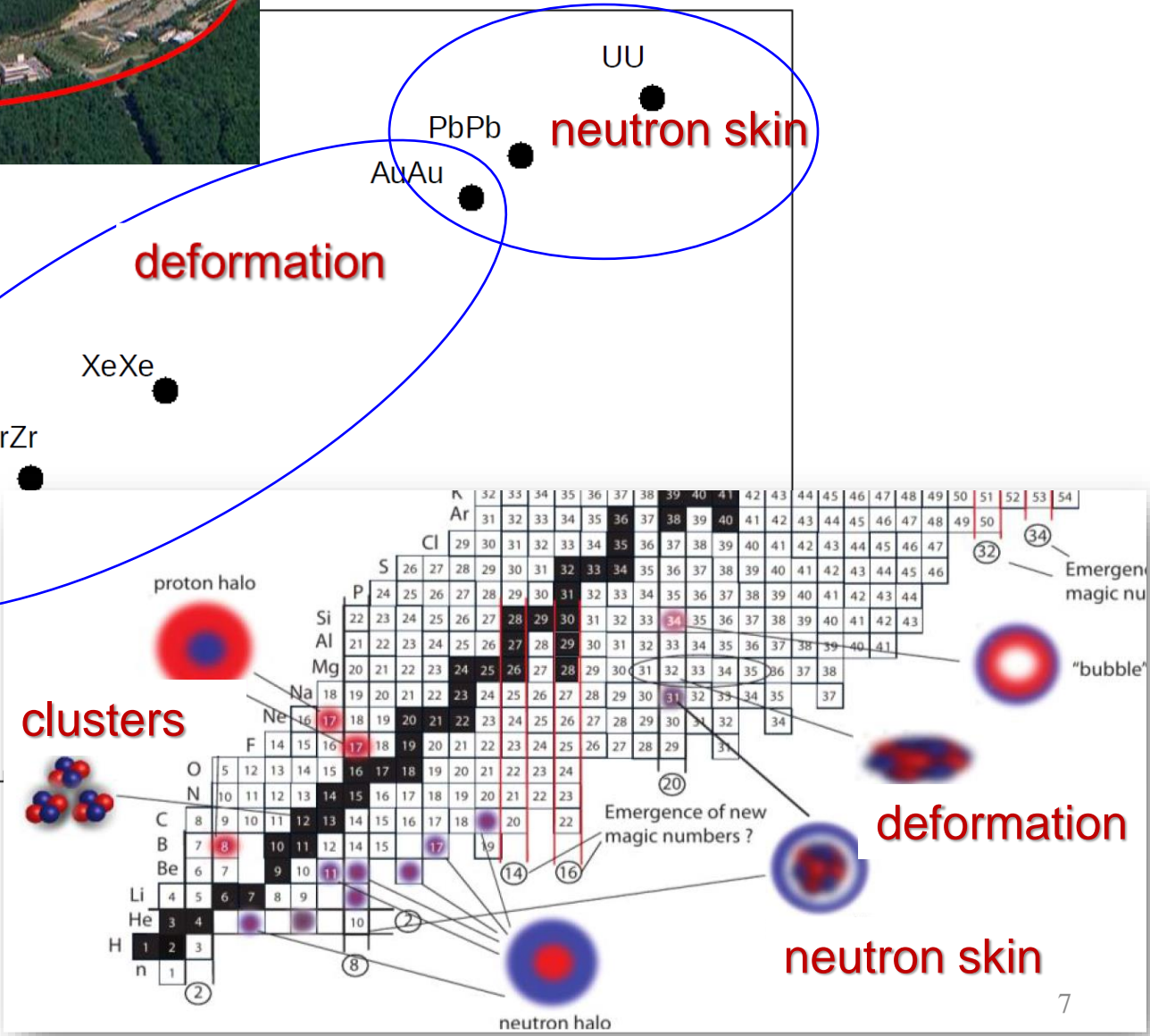
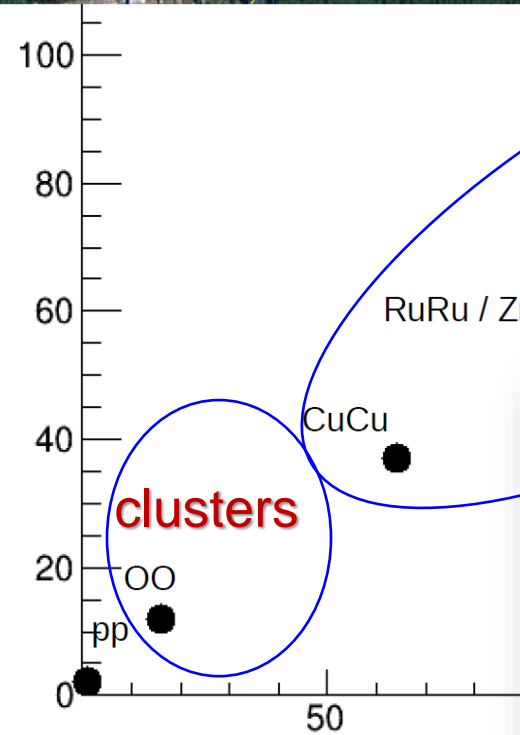
Rich A-A Collisions at RHIC & the LHC

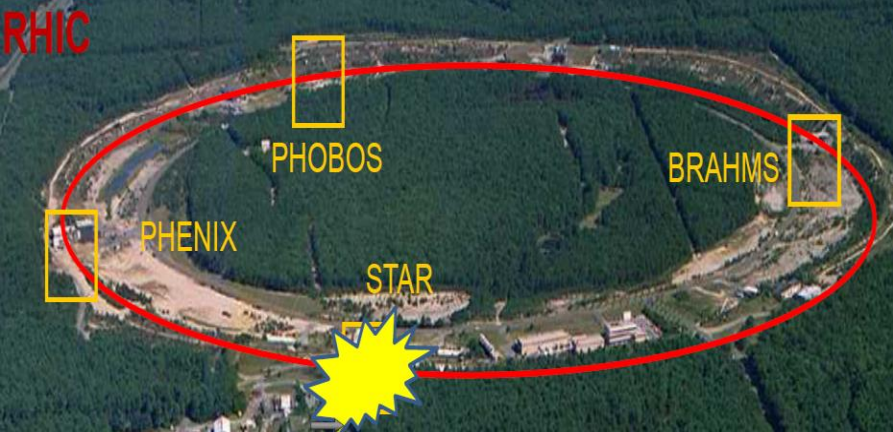


$^{197}\text{Au} + ^{197}\text{Au}$, $^{238}\text{U} + ^{238}\text{U}$, $^{208}\text{Pb} + ^{208}\text{Pb}$, $^{129}\text{Xe} + ^{129}\text{Xe}$, $^{96}\text{Zr} + ^{96}\text{Zr}$,
 $^{96}\text{Ru} + ^{96}\text{Ru}$, $^{64}\text{Cu} + ^{64}\text{Cu}$, $^{16}\text{O} + ^{16}\text{O}$, $p + ^{208}\text{Pb}$, $p + p$

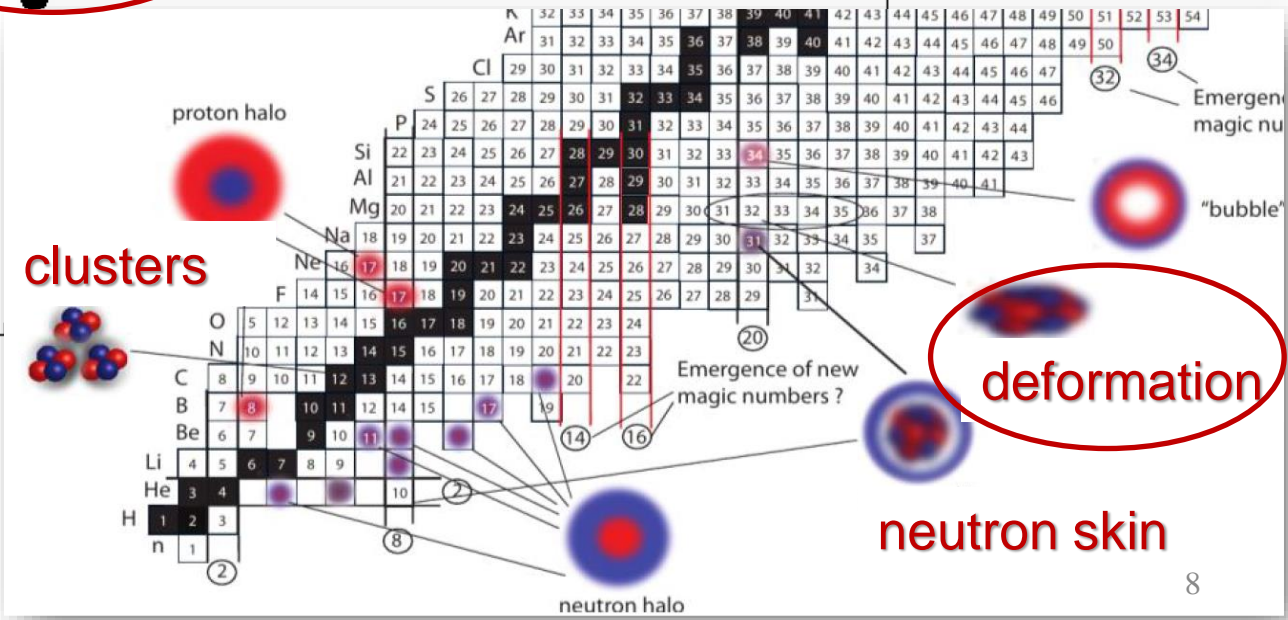
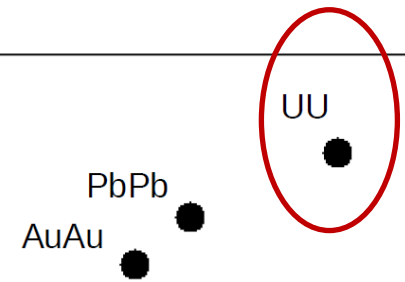
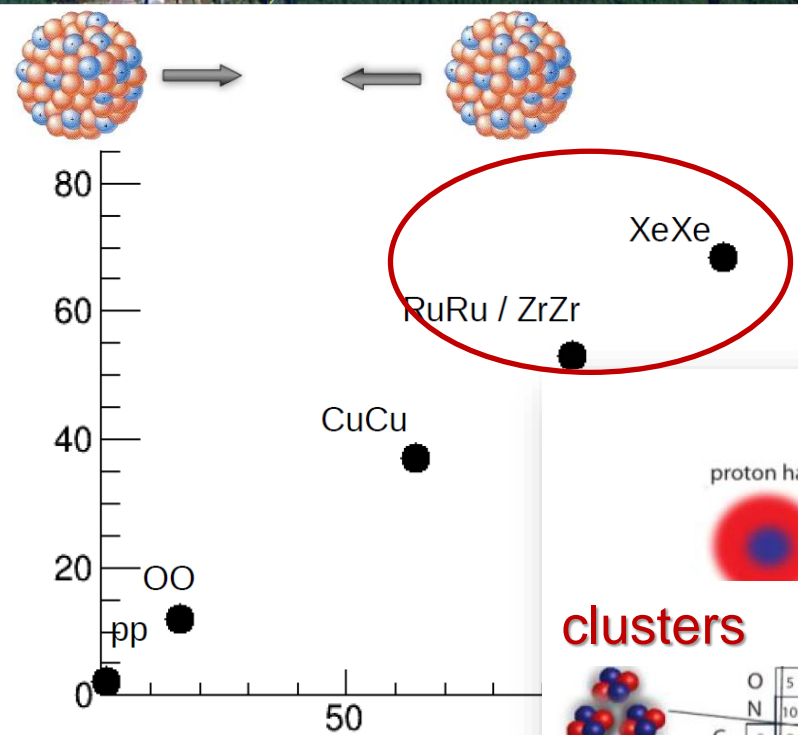


High energy nucleon nucleon collisions & nuclear structure

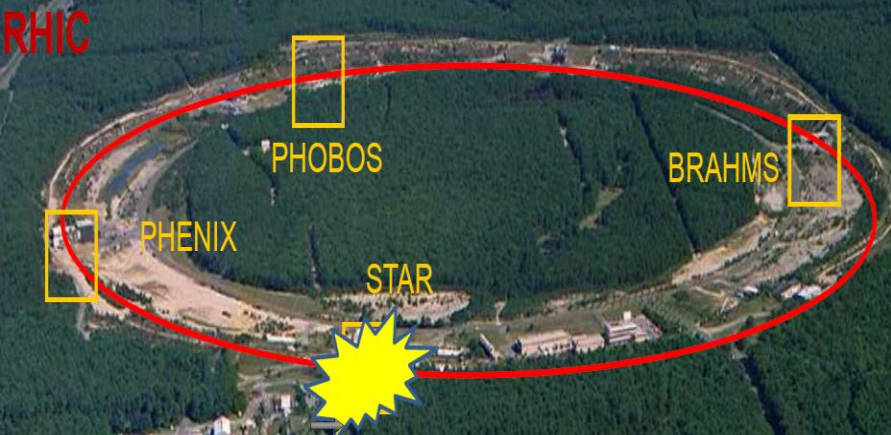




Probe the deformation of nuclei with relativistic heavy ion collisions

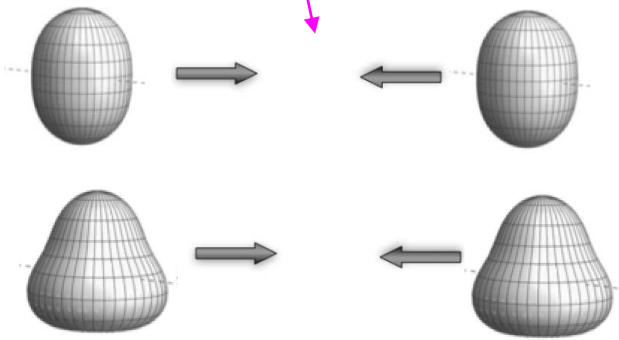


neutron skin

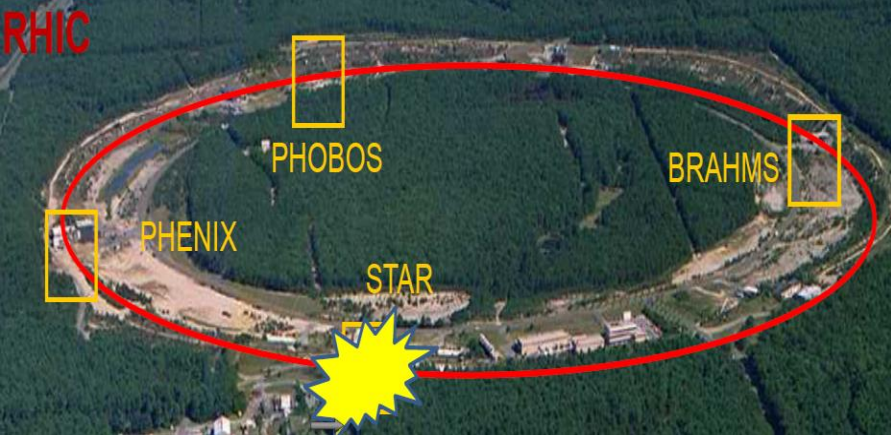


Probe the deformation of nuclei
with relativistic heavy ion collisions

- Relativistic heavy collisions
start from nuclei

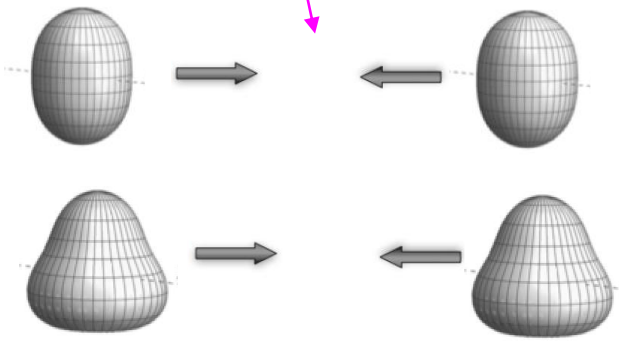


**initial conditions:
(with deformations)**

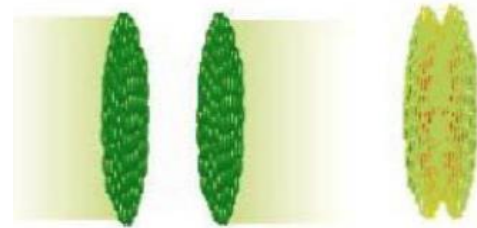


Probe the deformation of nuclei with relativistic heavy ion collisions

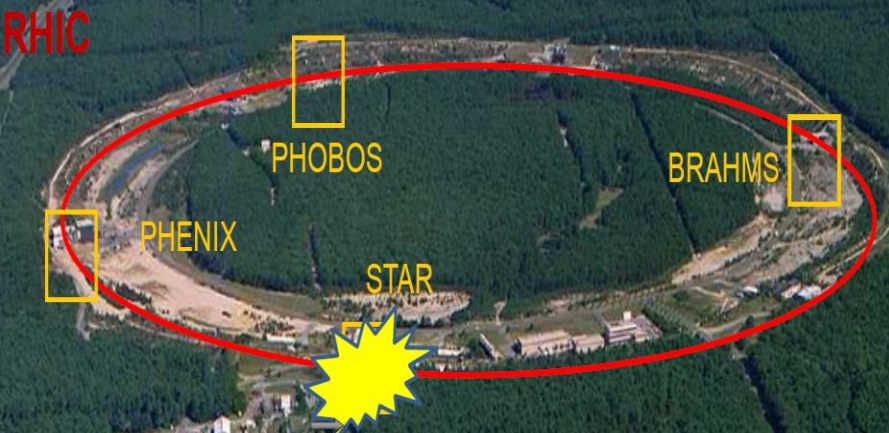
- Relativistic heavy collisions **start from nuclei**
- Collision time $< 10^{-24}$ s directly **probe the ground state of nuclei**



**initial conditions:
(with deformations)**

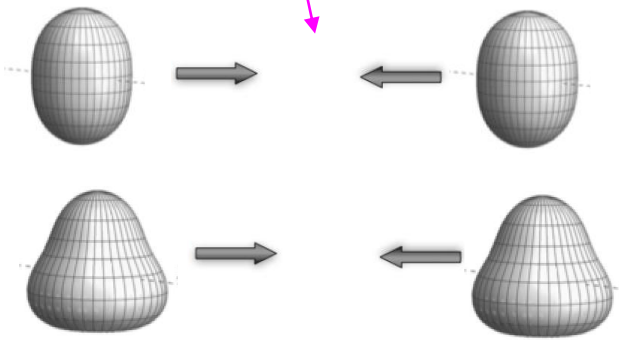


Collision time $< 10^{-24}$ s

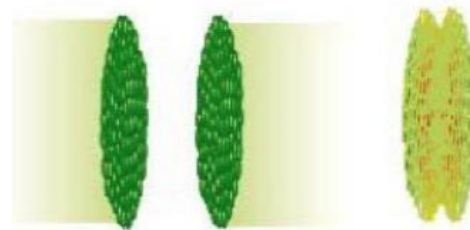


Probe the deformation of nuclei with relativistic heavy ion collisions

- Relativistic heavy collisions **start from nuclei**
- Collision time $< 10^{-24}$ s **directly probe the ground state of nuclei**

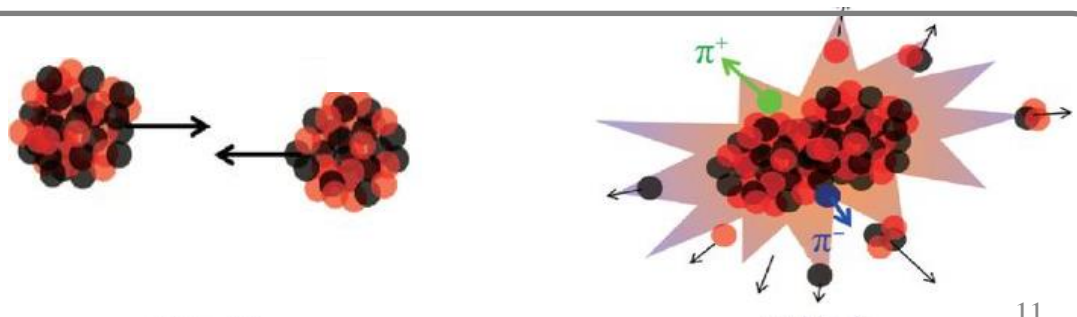


**initial conditions:
(with deformations)**



Collision time $< 10^{-24}$ s

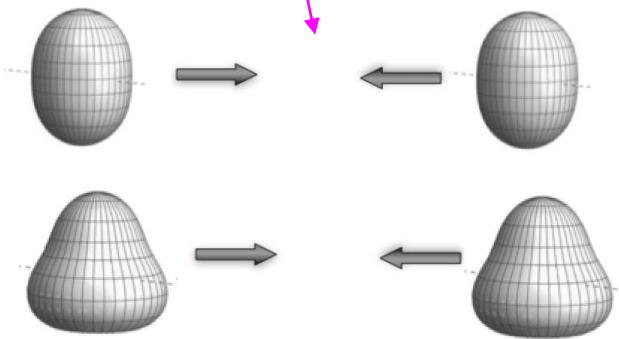
heavy ion collision at intermediate energies breaks up / excites nuclei during the collisions





Probe the deformation of nuclei with relativistic heavy ion collisions

- Relativistic heavy collisions **start from nuclei**
- Collision time $< 10^{-24}$ s directly **probe the ground state of nuclei**
- **Well calibrated calculations for QGP evolution**; to focus on the initial state



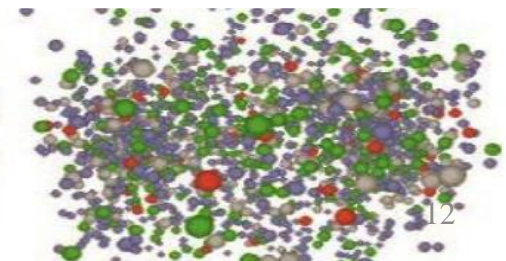
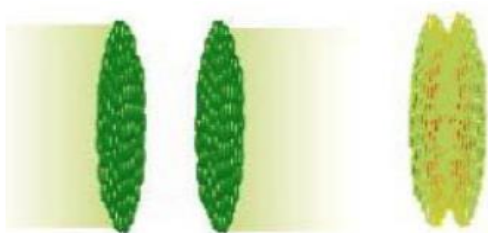
**initial conditions:
(with deformations)**

Well calibrated calculations

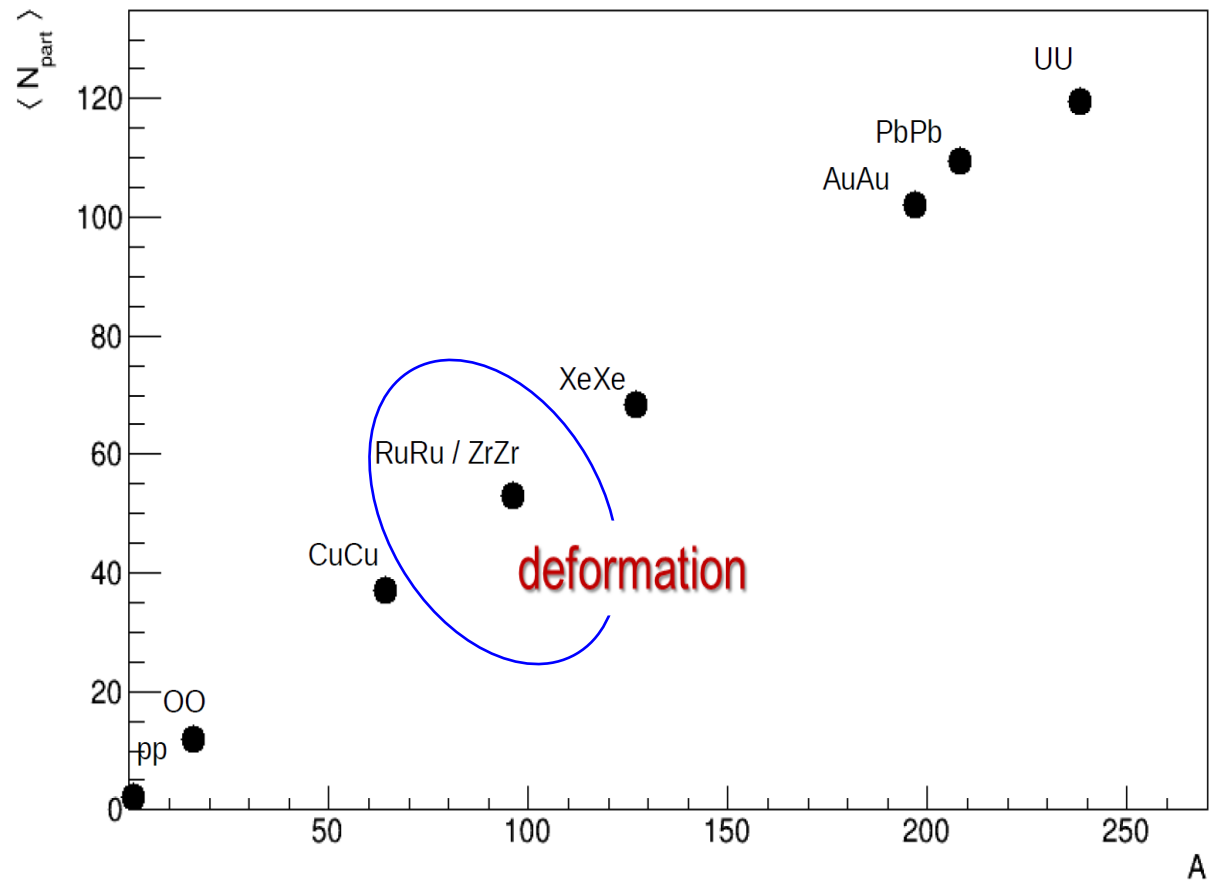
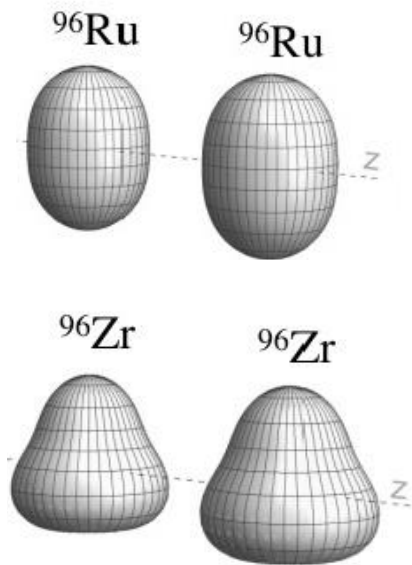
Initial conditions

viscous hydro

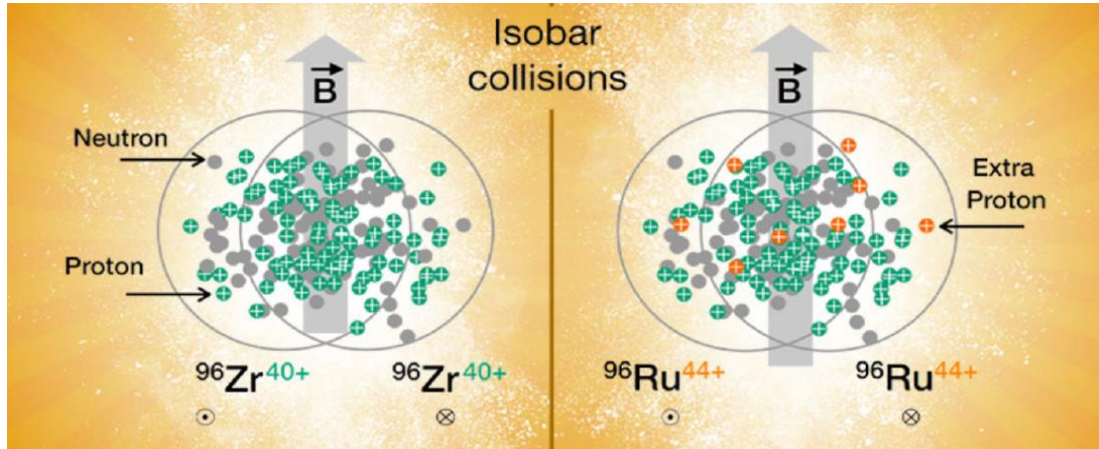
hadron cascade



Study the deformation of ^{96}Ru and ^{96}Zr at RHIC isobar run



$^{96}\text{Ru}+^{96}\text{Ru}$ and $^{96}\text{Zr}+^{96}\text{Zr}$ Collisions @ RHIC isobar run



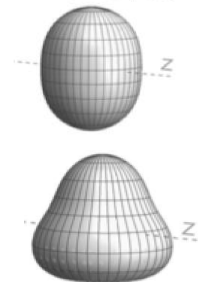
- To search the Chiral Magnetic Effect (CME)

- Obviously different early magnetic field for Ru+Ru and Zr+Zr collisions

Deformation of ^{96}Ru and ^{96}Zr

Conversion from $B(E_n)$ to β_n via: $\beta_2 = \frac{4\pi}{3ZR_2^2} \sqrt{\frac{B(E2)_{\uparrow}}{e^2}}$, $\beta_3 = \frac{4\pi}{3ZR_0^3} \sqrt{\frac{B(E3)_{\uparrow}}{e^2}}$

	β_2	$E_{2^+_1}$ (MeV)	β_3	$E_{3^-_1}$ (MeV)
^{96}Ru	0.154	0.83	-	3.08
^{96}Zr	0.062	1.75	0.202, 0.235, 0.27	1.90



ADNDT107,1 (2016) ADNDT80,35(2002)

Model calculation for Nuclear Deformation

General approach (DFT level)

Non-relativistic Schrodinger equation:
Skyrme and Gogny DFT

Relativistic Dirac equation:
covariant DFT (CDFT)

Range of interaction

Zero range - point coupling models
in CDFT (no mesons)
- Skyrme DFT

Finite range - meson exchange models
in CDFT
- Gogny DFT

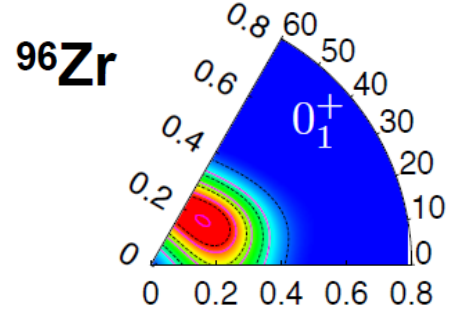
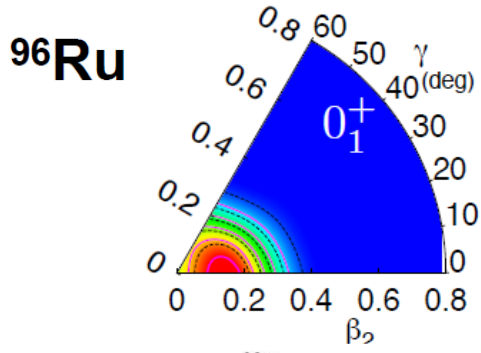
Effective density dependence

CDFT :

- explicit (DD-ME2, DD-PC1)
- non-linear (through the powers of mesons) (NL1, NL3*)

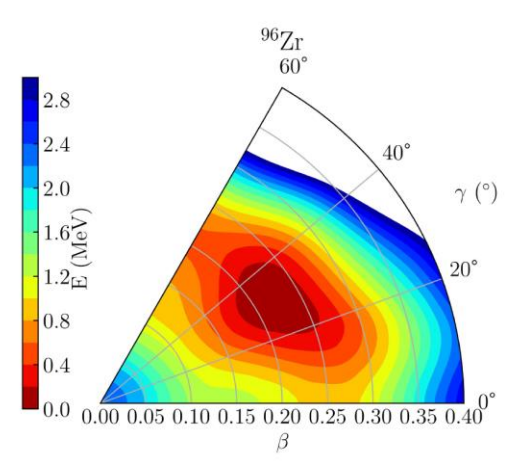
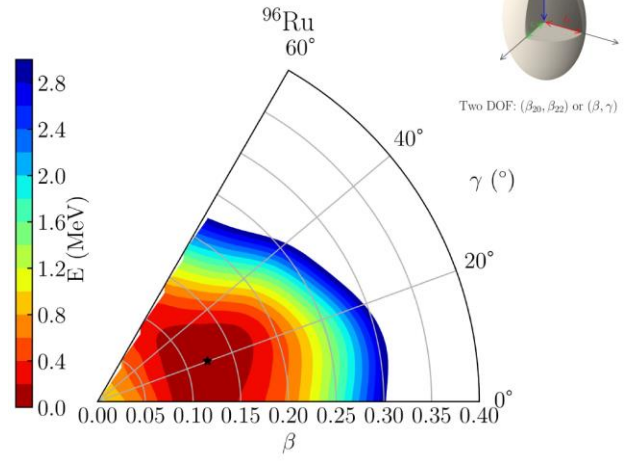
Skyrme and Gogny DFTs: different prescriptions for density dependence

Deformation of ^{96}Ru & ^{96}Zr — DFT calculations



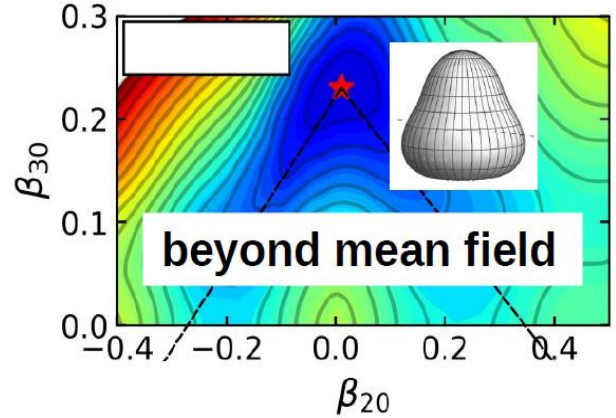
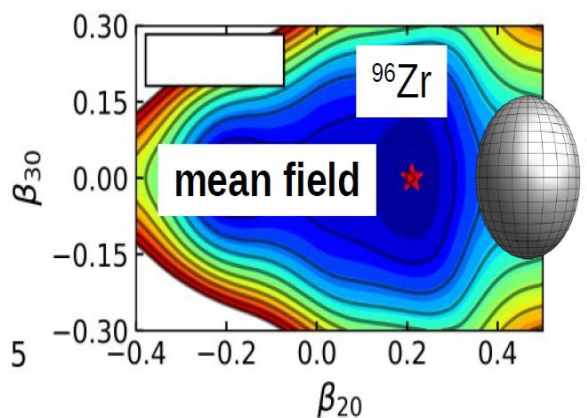
Gogny energy density functional (Tiaxial)

T R. Rodríguez EMMI RRTF 2022



Skyrme EDF (with rotational correction)

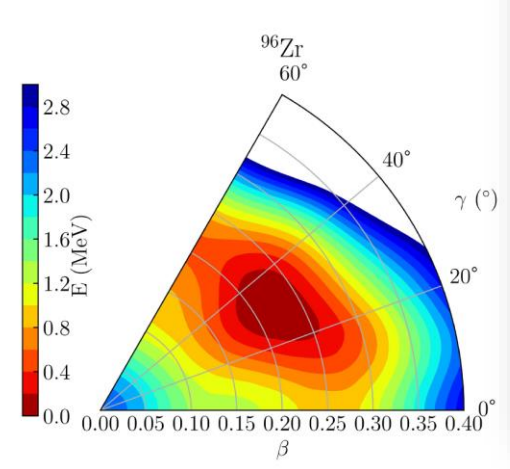
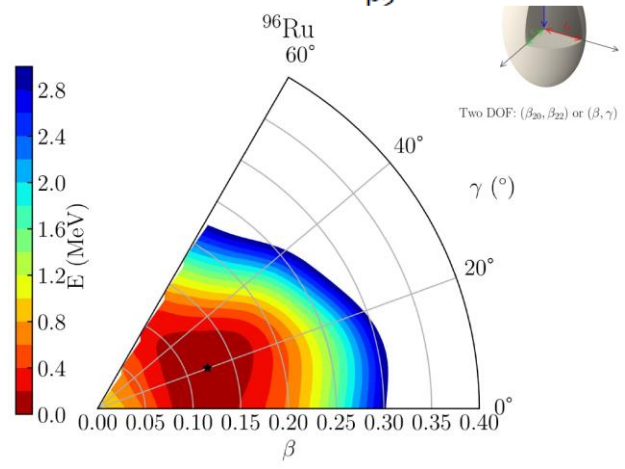
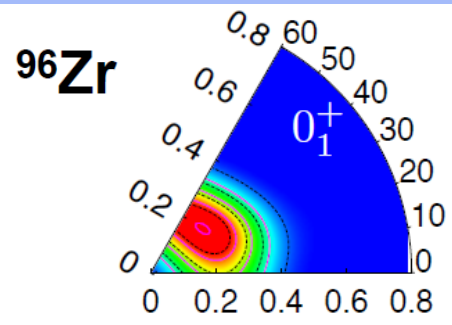
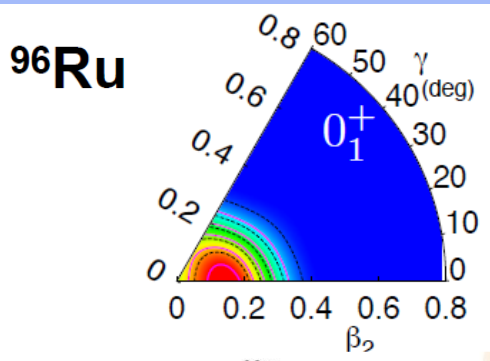
W Ryssens EMMI RRTF 2022



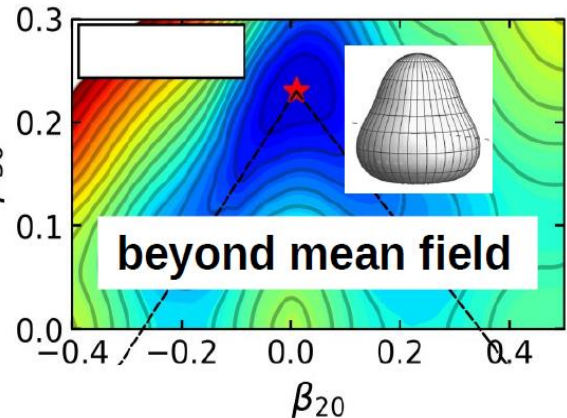
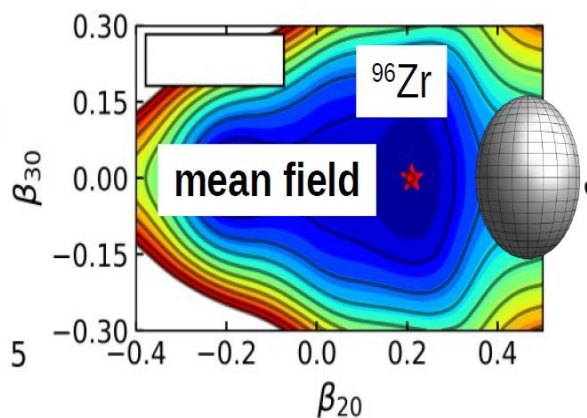
Beyond-mean-field correction is very important

Rong, Lu, arXiv:2201.02114

Deformation of ^{96}Ru & ^{96}Zr — personal comments

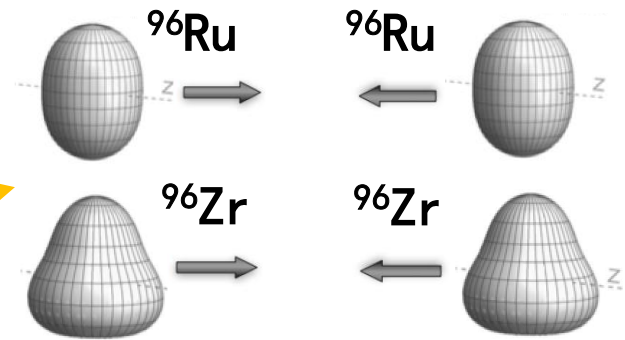
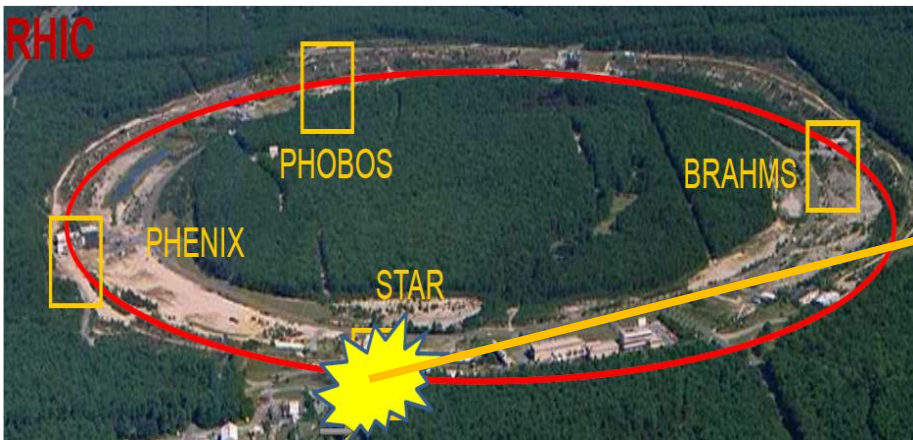


18^+	8205.7	15^-	6754.1
16^+	6441.6	13^-	5750.2
14^+	5680.7	11^-	4798.7
12^+	4418.3	9^-	3951.1
10^+	3817.2	7^-	3291.5
8^+	2950.4	5^-	2588.4
6^+	2149.7		
4^+	1518.1		
2^+	832.6		
0^+	0.0		



Nuclear structure physics obtain the deformation information from the spectrum with certain model calculations (not directly image the deformation in position space)

Probe the deformation (mass distributions) of ^{96}Ru & ^{96}Zr



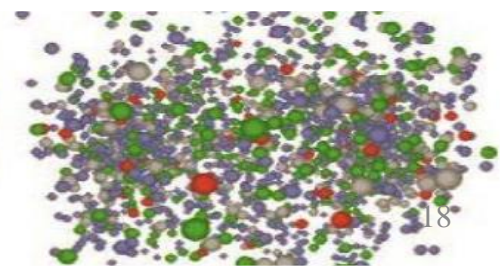
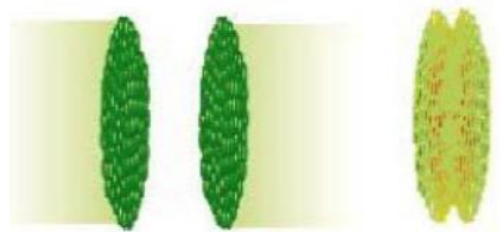
**initial conditions:
(deformation / mass distributions)**

Well calibrated calculations

Initial conditions

viscous hydro

hadron cascade



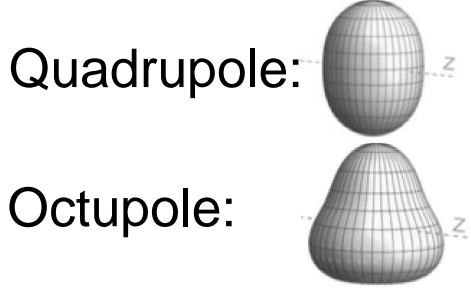
Hydrodynamic calculation with initially deformed nuclei

Initial conditions (TRENTO)

- Sample nucleon position in deformed nuclei with:

$$\rho(r, \theta, \phi) = \frac{\rho_0}{1 + e^{(r-R(\theta, \phi))/a_0}}$$

$$R(\theta, \phi) = R_0 \left(1 + \beta_2 [\cos \gamma Y_{2,0} + \sin \gamma Y_{2,2}] + \beta_3 \sum_{m=-3}^3 \alpha_{3,m} Y_{3,m} + \beta_4 \sum_{m=-4}^4 \alpha_{4,m} Y_{4,m} \right)$$

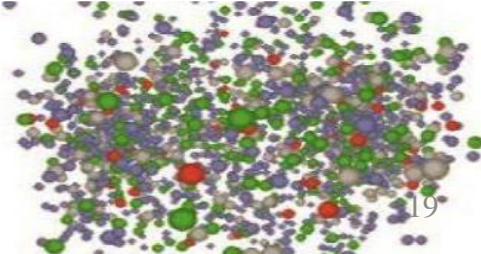
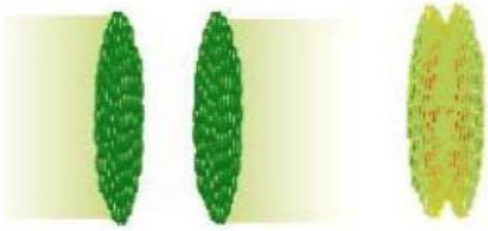


Well calibrated calculations

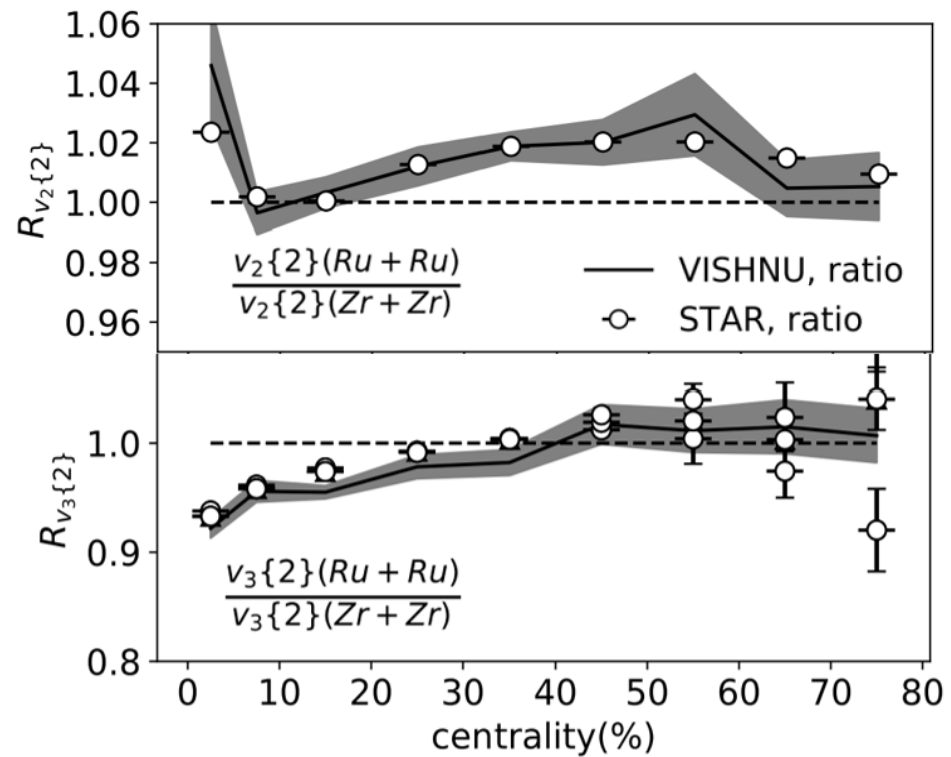
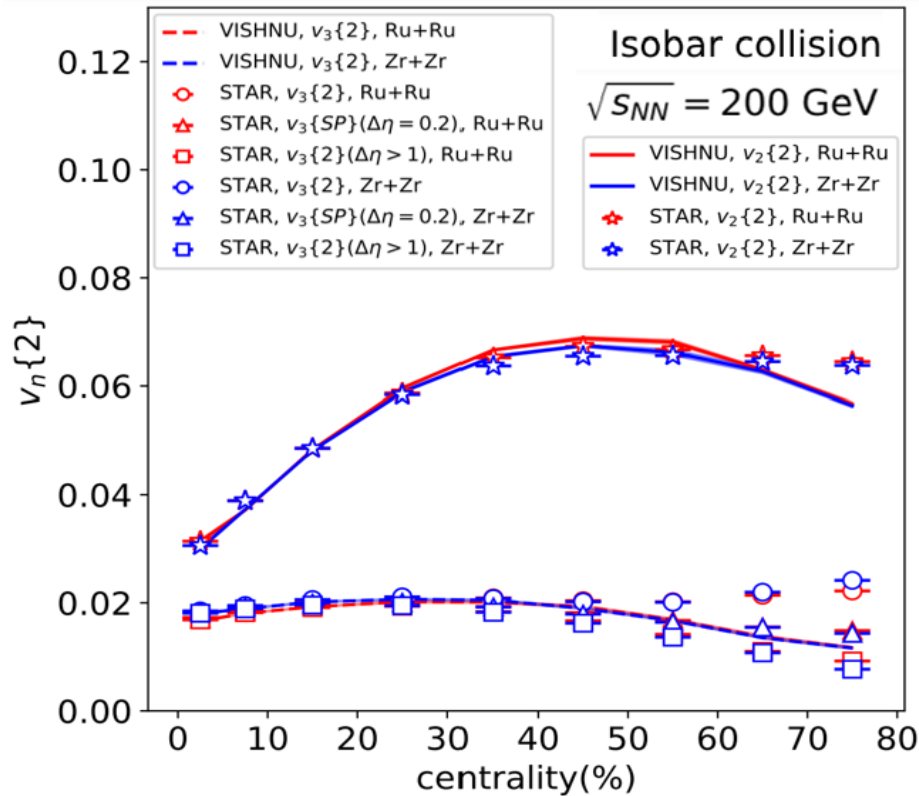
Initial conditions

viscous hydro

hadron cascade



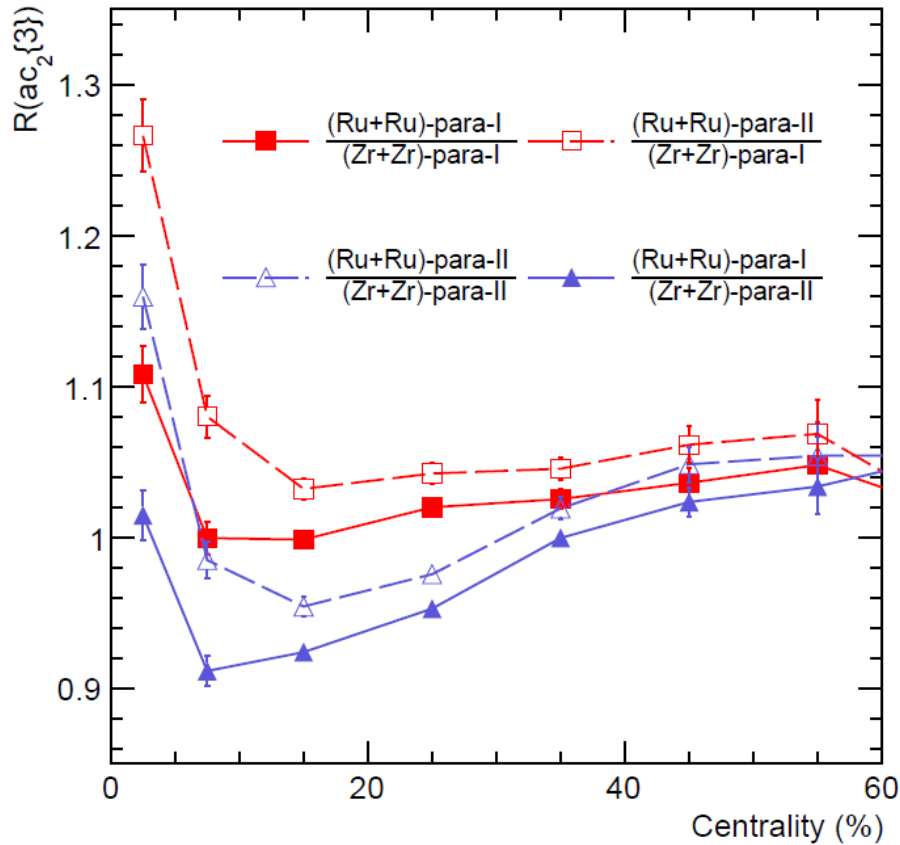
V_2 and V_3 for Ru+Ru and Zr+Zr collisions



- With fine tuning parameters, iEBE-VISHNU fits V_2 & V_3 for Ru+Ru collisions
- Using β_2 β_3 in table1, it “predicts” V_2 & V_3 for Zr+Zr collisions & the related ratio
- (the data are roughly described).

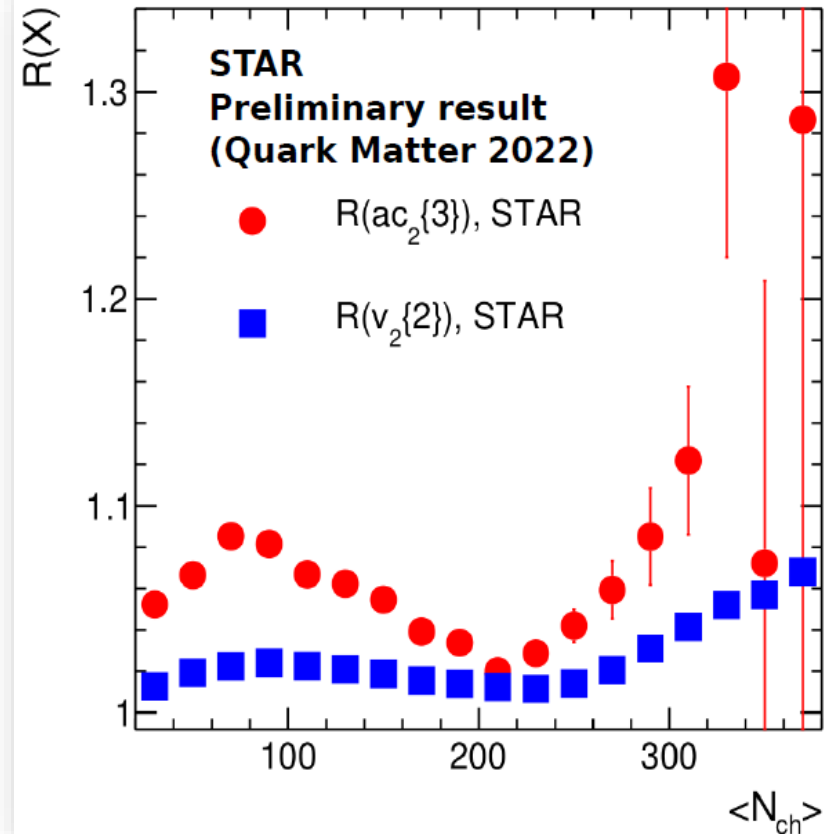
“standard”	Ru	Zr
a_0	0.46	0.52
β_2	0.162	0.060
β_3	0.00	0.200

$ac_2\{3\}$ for Ru+Ru and Zr+Zr collisions



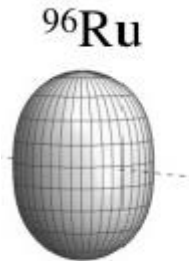
$ac_2\{3\}$ is sensitive to quadrupole and octupole deformations

$$ac_2\{3\} = \langle v_2^2 v_4 \cos 4(\Phi_2 - \Phi_4) \rangle,$$



Probe the deformation of ^{96}Ru and ^{96}Zr

-- a short summary

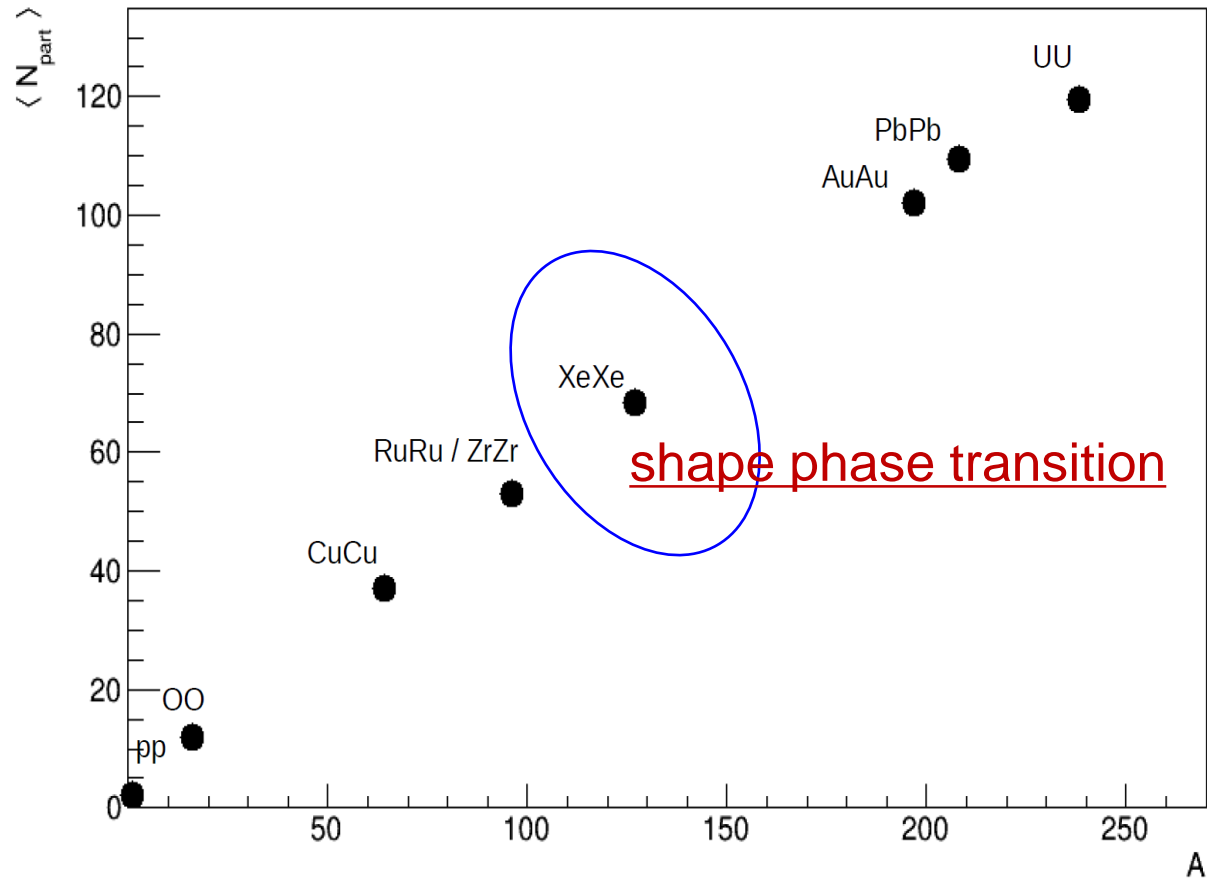


- ^{96}Ru and ^{96}Zr : two ideal nuclei for interdisciplinary research between relativistic heavy ion physics and nuclear structure

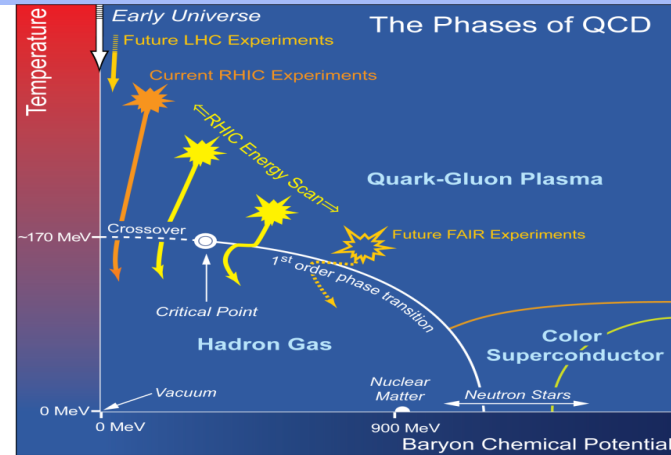
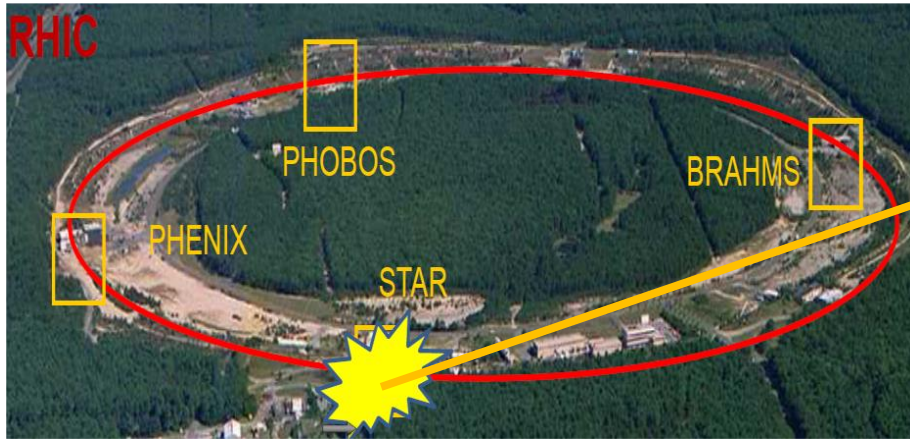
-isobar collisions provide rich and high statistical run data for various flow analysis, which could constrain the deformation of ^{96}Ru and ^{96}Zr from heavy ion physics side

-Need more efforts to study the deformation of ^{96}Ru & ^{96}Zr from both experimental and theoretical sides in nuclear structure

Probe the shape phase transition with Xe +Xe collisions



The Phase Transition



Relativistic heavy ion collisions

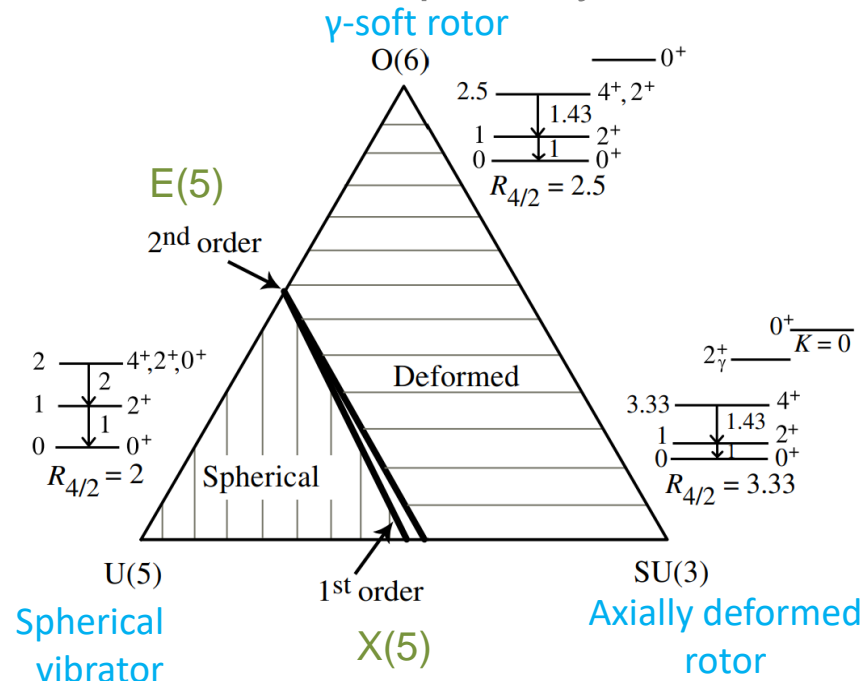
-mainly aim to explore QCD Phase Transition

-can also probe the shape phase transition of finite nuclei along certain isotope/ isotone chain

$^{129}\text{Xe} + ^{129}\text{Xe}$ collision

-explore the second-order shape phase transition occurring in the vicinity of $^{128-130}\text{Xe}$

The critical point symmetries



Shape phase transition for Xe isotopes

The shape phase transition:

- rapid structural change along certain isotope or isotone chains
- the dynamic interplay between the spherical-driving pairing interaction and the deformation-driving proton-neutron interaction

The shape phase transition for the Xe isotopes:

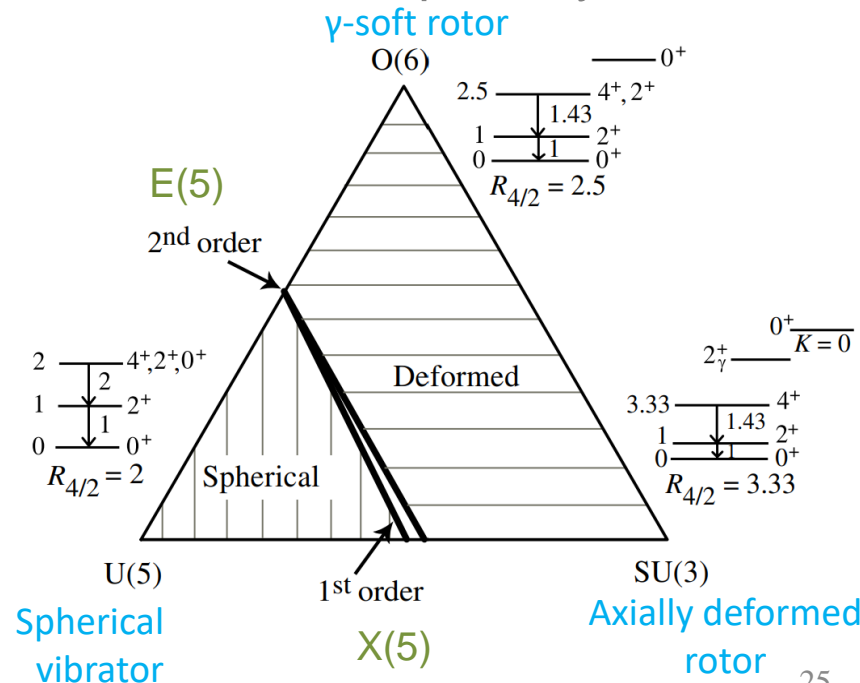
-Within the the framework of the interacting boson model (IBM), the Xe isotopes undergo a shape phase transition from a γ -soft rotor to a spherical vibrator

R. F. Casten, Nucl. Phys. A 439, 289 (1985). G. Puddu, O. Scholten, and T. Otsuka, Nucl. Phys. A 348, 109 (1980). R. F. Casten and P. Von Brentano, Phys. Lett. B 152, 22 (1985).

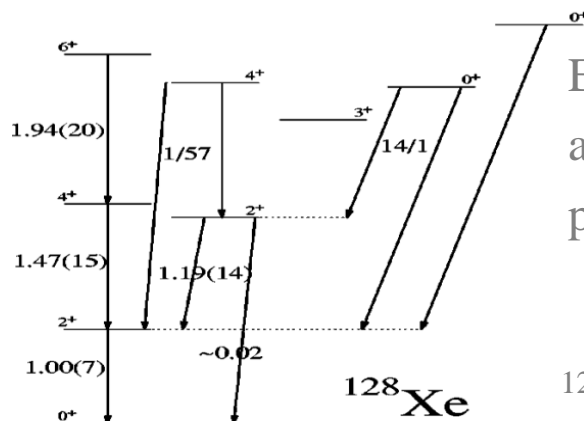
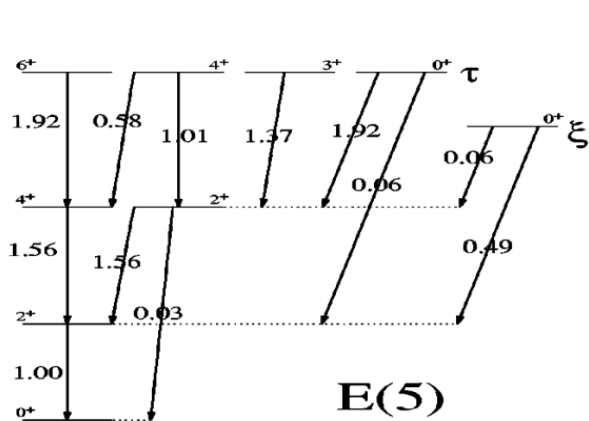
-the critical point is described by the $E(5)$ symmetry, associated with a 2nd order phase transition

F. Iachello, Phys. Rev. Lett. 87, 052502 (2001).
F. Iachello, Phys. Rev. Lett. 85, 3580 (2000).

The critical point symmetries



Exp evidence of E(5) symmetry for ^{128}Xe



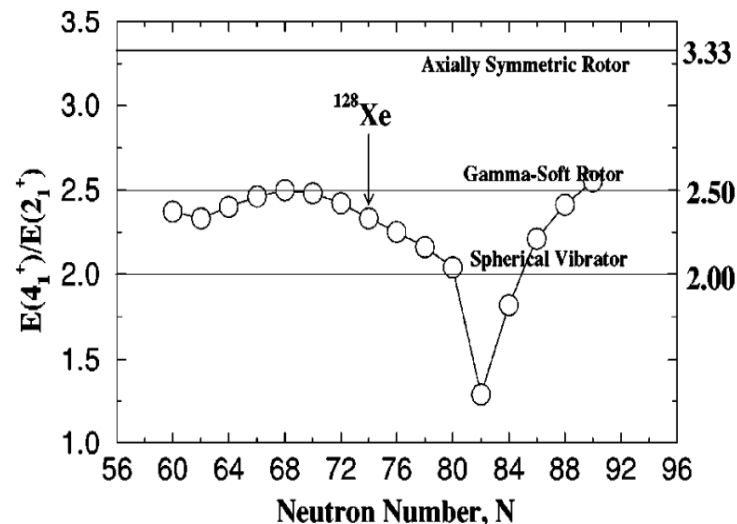
Energy spectroscopy: good agreement with E(5) prediction

^{128}Xe lies in between γ -soft rotor and spherical vibrator.

Nucleus	$E(4_1^+)/E(2_1^+)$	$E(0_2^+)/E(2_1^+)$	$E(0_3^+)/E(2_1^+)$
^{128}Xe	2.33	3.57	4.24
^{130}Xe	2.25	(3.35)	(3.76)
^{132}Xe	2.16		
^{134}Xe	2.04		

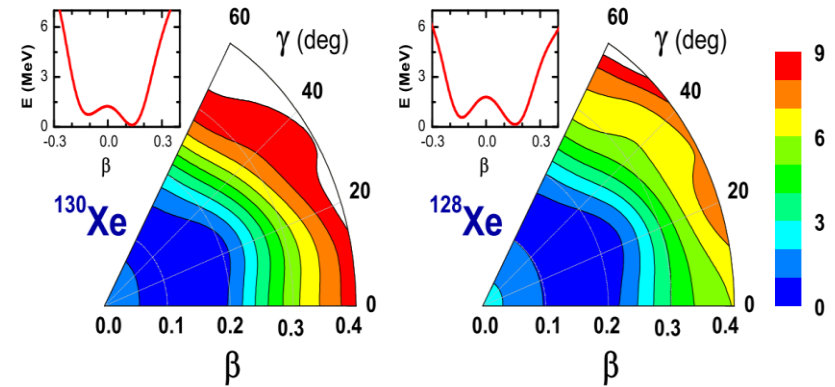
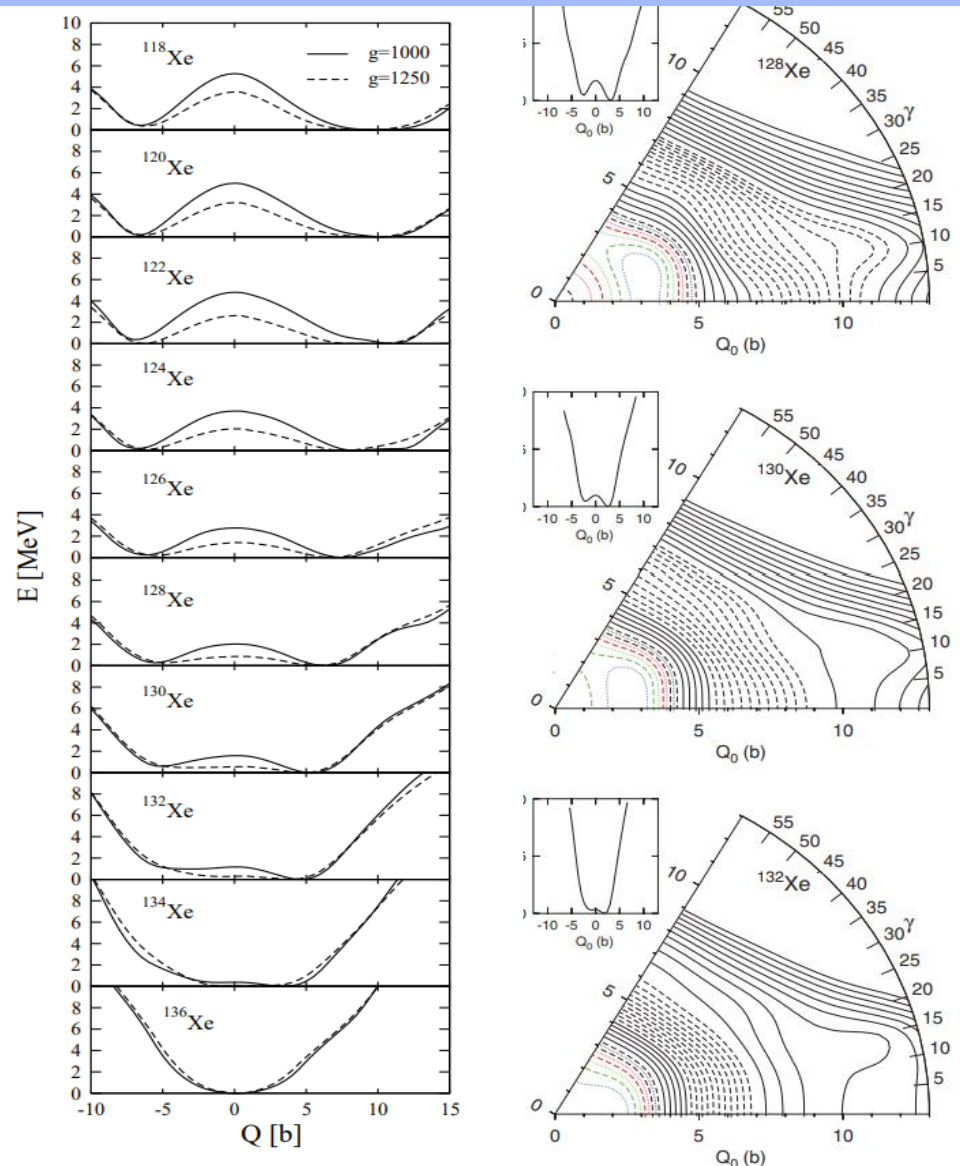
Evolution of $E(4_1^+)/E(2_1^+)$ ratio close to 2.2

Existence of two 0^+ states with $3 < E(0_n^+)/E(2_1^+) < 4$



the measured energy spectroscopy of ^{128}Xe agrees well with the E(5) predic. (the normalized transition strengths, the branching ratios, the energy ratios between different energy levels

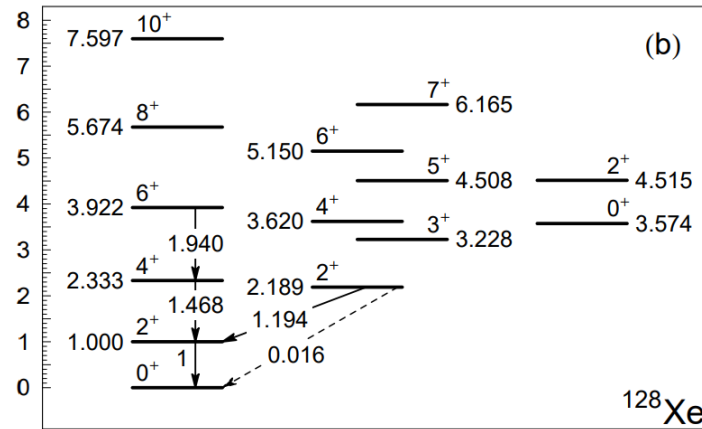
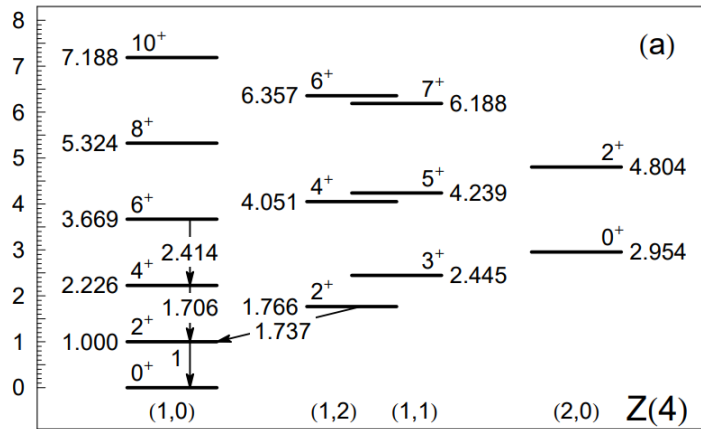
Theoretical predictions on E(5) symmetry near $^{128-130}\text{Xe}$



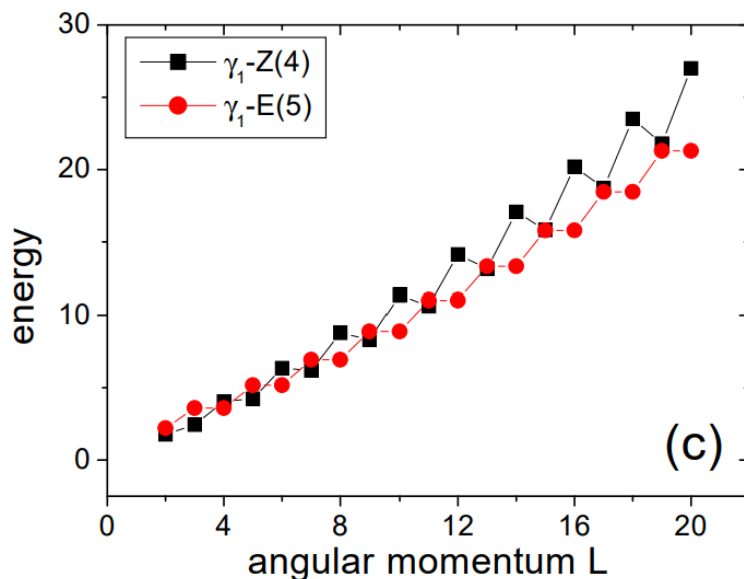
Z. P. Li, T. Niksic, D. Vretenar, and J. Meng (2010)

-Various theoretical calculations indicate a critical point of the second-order shape phase transition (E(5) symmetry) lies in the vicinity of $^{128-130}\text{Xe}$, associated with a γ -soft deformation

Z(4) or E(5) symmetry for $^{128-132}\text{Xe}$?



-Z(4) symmetry with a fixed γ at 30° can also describe the spectra and $B(E2)$ rates for $^{128,130,132}\text{Xe}$



The mean difference between E(5) and Z(4) is the pair order of energy levels in the γ band.

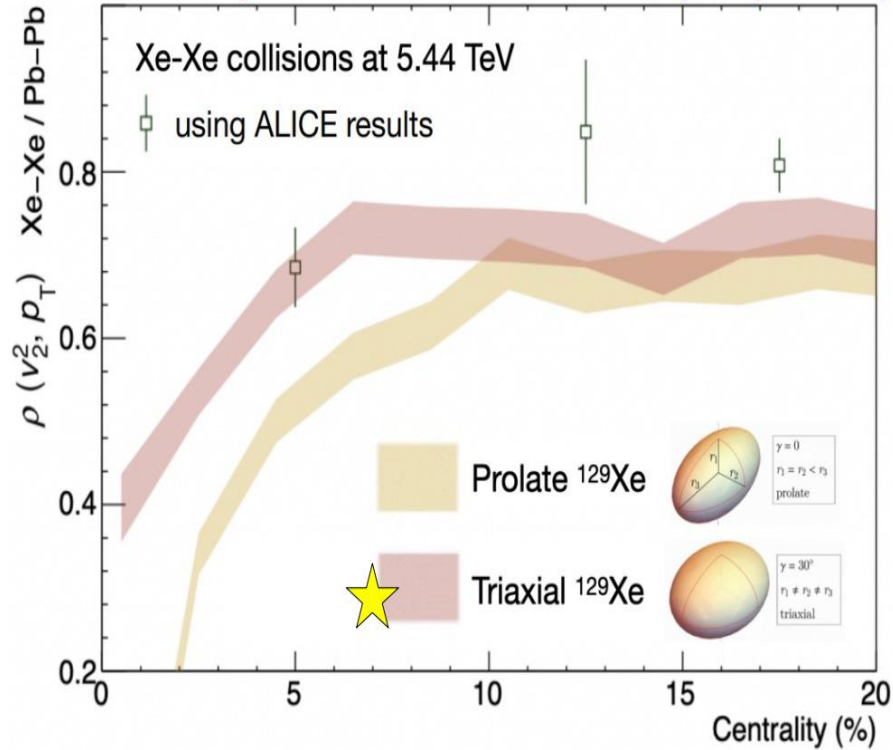
-It is hard to distinguish such E(5) and Z(4) symmetry in nuclear physics.

D. Bonatsos, D. Lenis, D. Petrellis, P. A. Terziev, and I. Yigitoglu, Phys. Lett. B 621, 102 (2005),

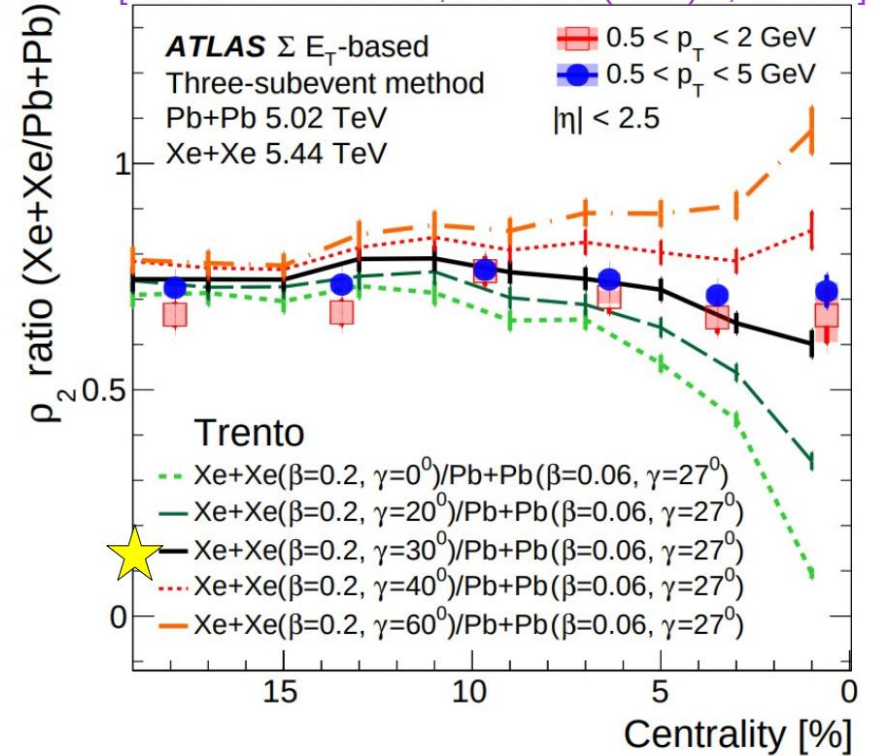
Probing triaxial deformation of ^{129}Xe with $^{129}\text{Xe}+^{129}\text{Xe}$ collisions

B. Bally, M. Bender, G. Giacalone, V. Somà, Phys. Rev. Lett. 128 (8) (2022) 082301

[ALICE Collaboration, PLB **834** (2022) 137393]



[ATLAS Collaboration, PRC **107** (2023) 5, 054910]



Distinguish triaxial and γ -soft deformation of ^{129}Xe with $^{129}\text{Xe}+^{129}\text{Xe}$ collisions

explore the possible 2nd order phase transition of Xe isotopes

Hydro initial condi. with triaxial or γ -soft deformation

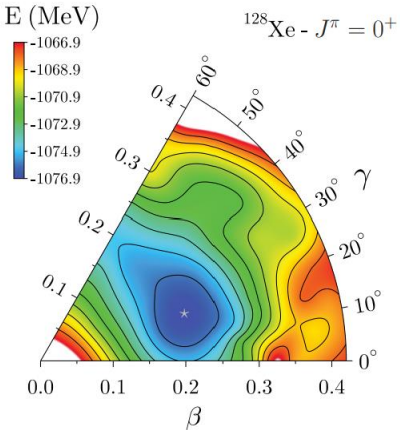
S. Zhao, H. Xu, Y. Zhou, Y. Liu, H. Song, arXiv: 2403.07441 [nucl-th]

Initial conditions (TRENTO)

- Sample nucleon position in deformed nuclei with:

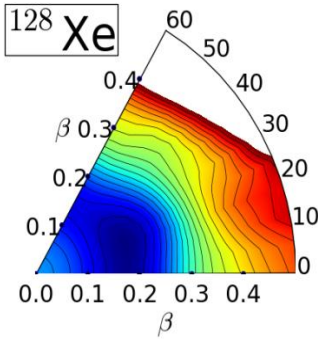
$$\rho(r, \theta, \phi) = \frac{\rho_0}{1 + e^{(r-R(\theta, \phi))/a_0}}$$

$$R(\theta, \phi) = R_0(1 + \beta_2[\cos \gamma Y_{2,0}(\theta, \phi) + \sin \gamma Y_{2,2}(\theta, \phi)]).$$



Rigid triaxial deformation ($\gamma=30^\circ$)

Bally et. al. Eur.Phys.J. A 58 (2022) 9, 187,



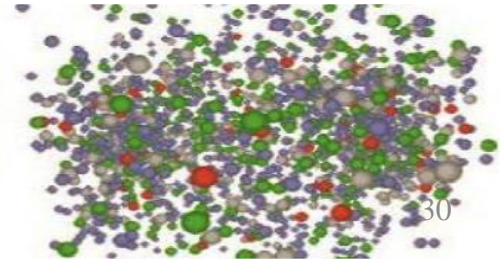
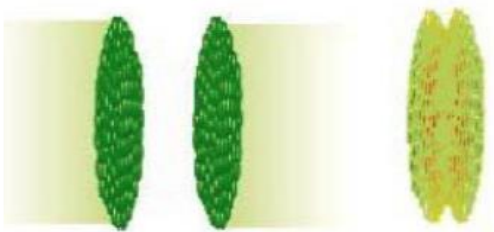
γ -soft (flat distribution in $0 \leq \gamma \leq 60^\circ$)

Z. P. Li, et. al. Phys. Rev. C 81, 034316 (2010),

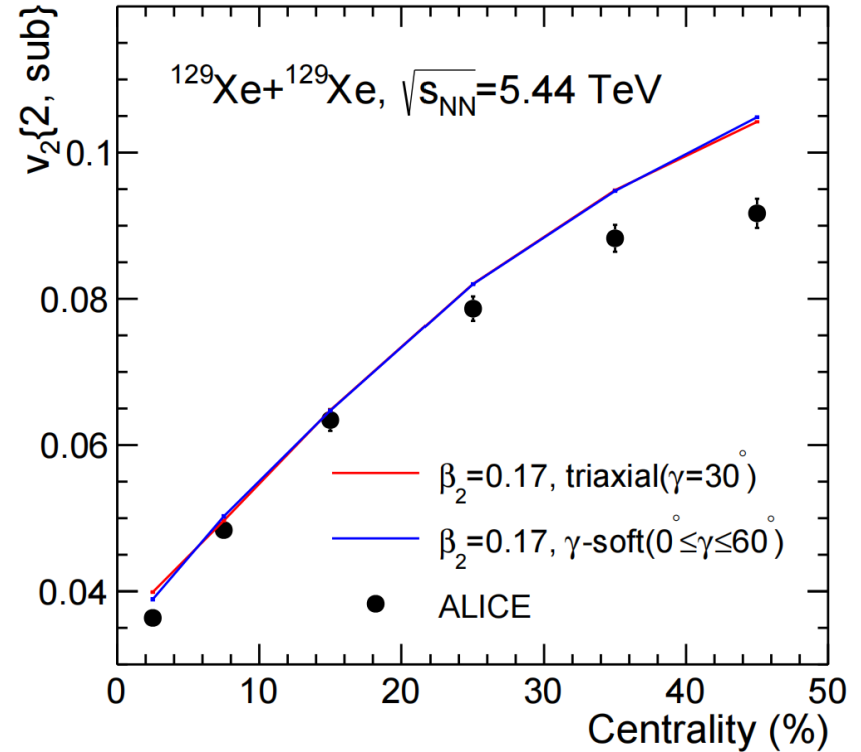
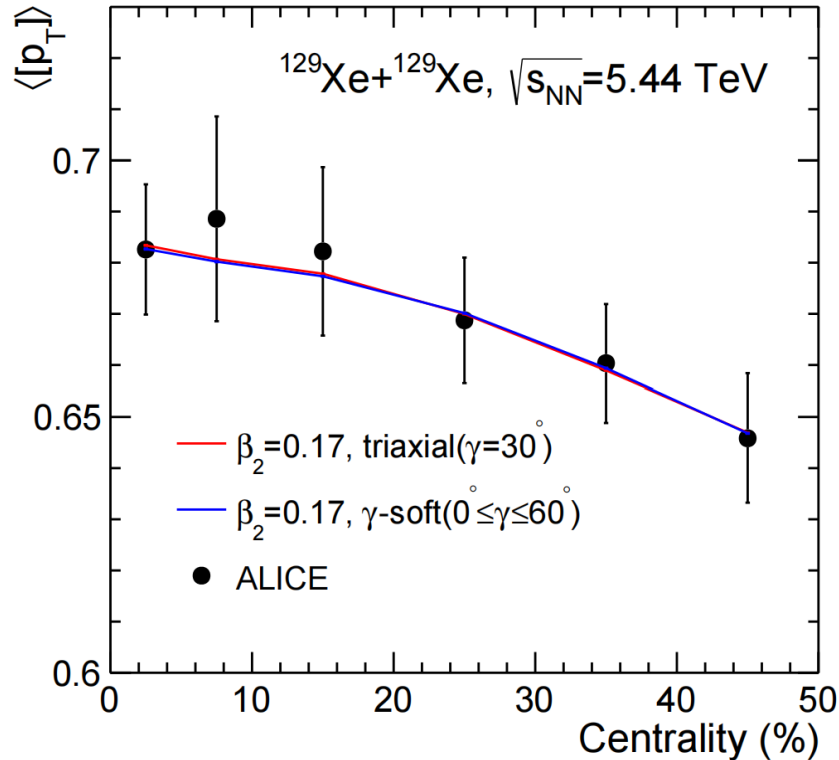
Initial conditions

viscous hydro

hadron cascade

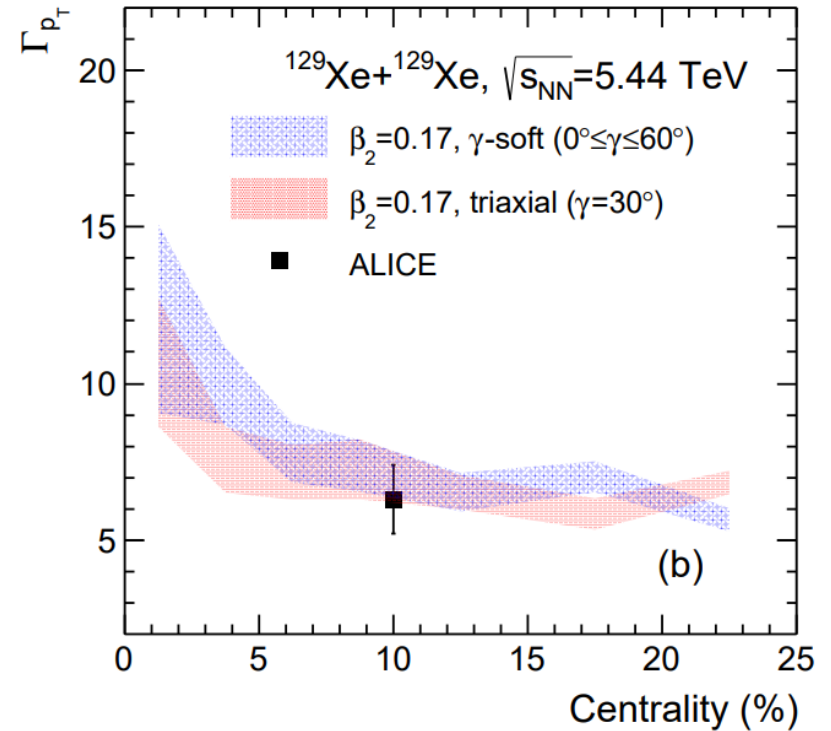
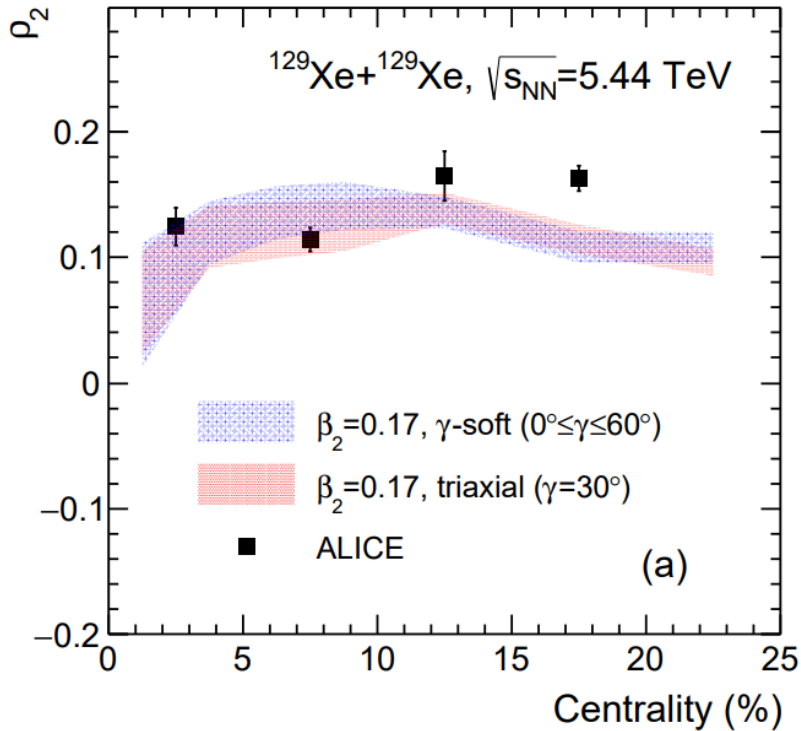


mean P_T & v_2



-With the parameters obtained from previous Bayesian analysis (Pb+Pb coll), our iEBE-VISHNU, with rigid triaxial or γ -soft deformation of ^{129}Xe , can describe most of the bulk observables in Xe+Xe collisions

3-particle correlation



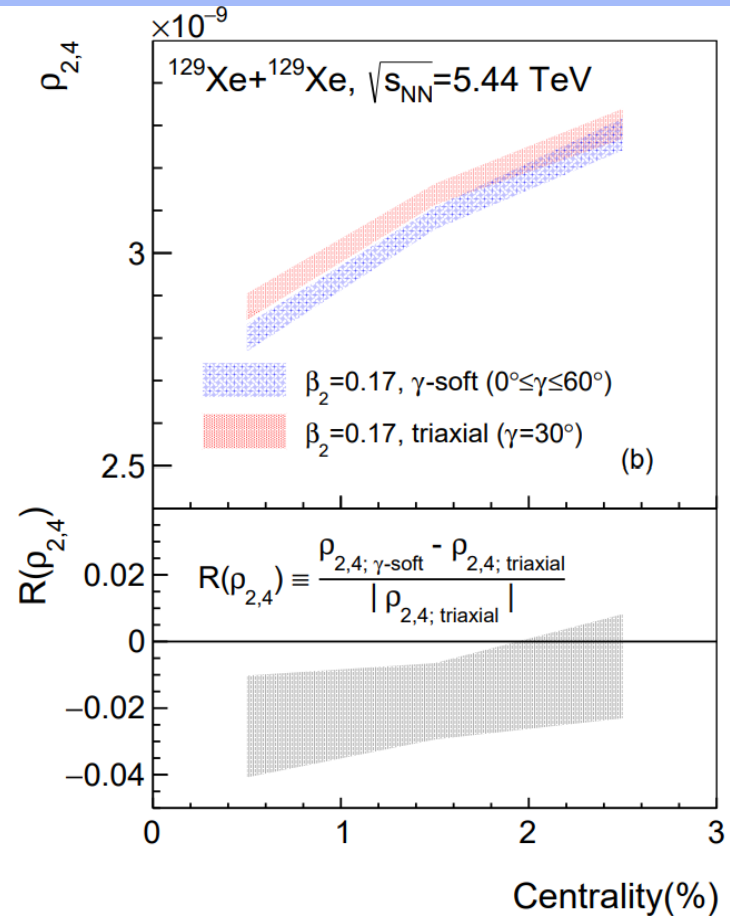
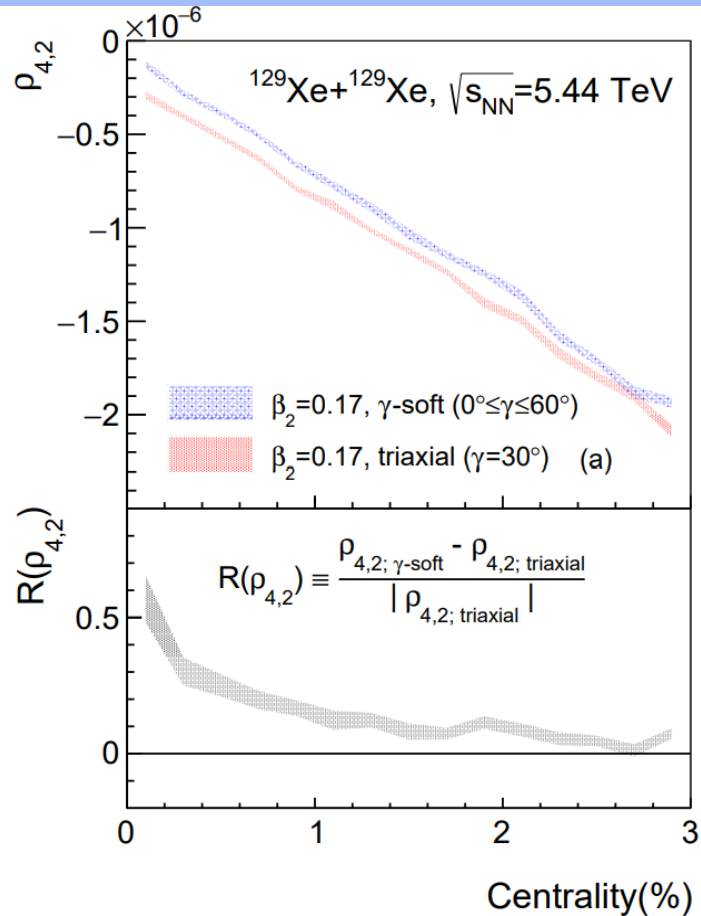
$$\rho_2 \equiv \frac{\text{cov}(v_2\{2\}^2, [p_T])}{\sqrt{\text{var}(v_2\{2\}^2)}\sqrt{\text{var}([p_T])}$$

$$\Gamma_{p_T} = \frac{\langle \delta p_{T,i} \delta p_{T,j} \delta p_{T,k} \rangle \langle [p_T] \rangle}{\langle \delta p_{T,i} \delta p_{T,j} \rangle^2},$$

-Our calculations with rigid triaxial or γ -soft deformation of ^{129}Xe can describe the measured ρ_2 and Γ_{p_T} equally well.

$\rho_2, \Gamma_{p_T} \propto \beta_2^3 \cos(3\gamma)$ insensitive to triaxial deformation $\gamma=30^\circ$ and γ -soft $0 \leq \gamma \leq 60^\circ$

6-particle correlations



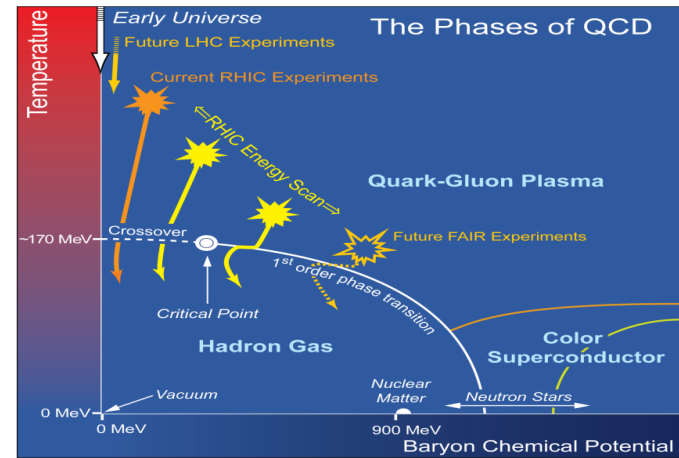
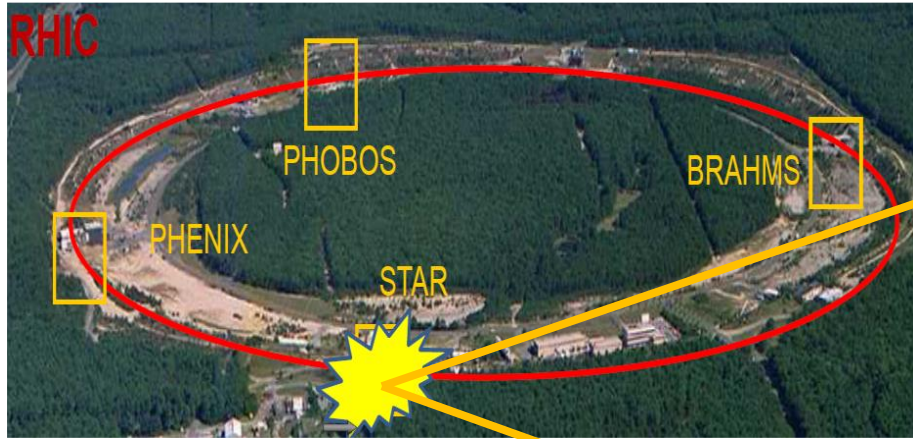
$$\rho_{4,2} \equiv \left(\frac{\langle \varepsilon_2^4 \delta d_\perp^2 \rangle}{\langle \varepsilon_2^4 \rangle \langle d_\perp \rangle^2} \right)_c$$

$$\rho_{2,4} \equiv \left(\frac{\langle \varepsilon_2^2 \delta d_\perp^4 \rangle}{\langle \varepsilon_2^2 \rangle \langle d_\perp \rangle^4} \right)_c \quad \langle \cos(6\gamma) \rangle$$

The γ -soft deformation of ^{129}Xe lead to a clear enhancement of 6-particle correlations $\rho_{4,2}$ in ultra-central Xe+Xe collisions

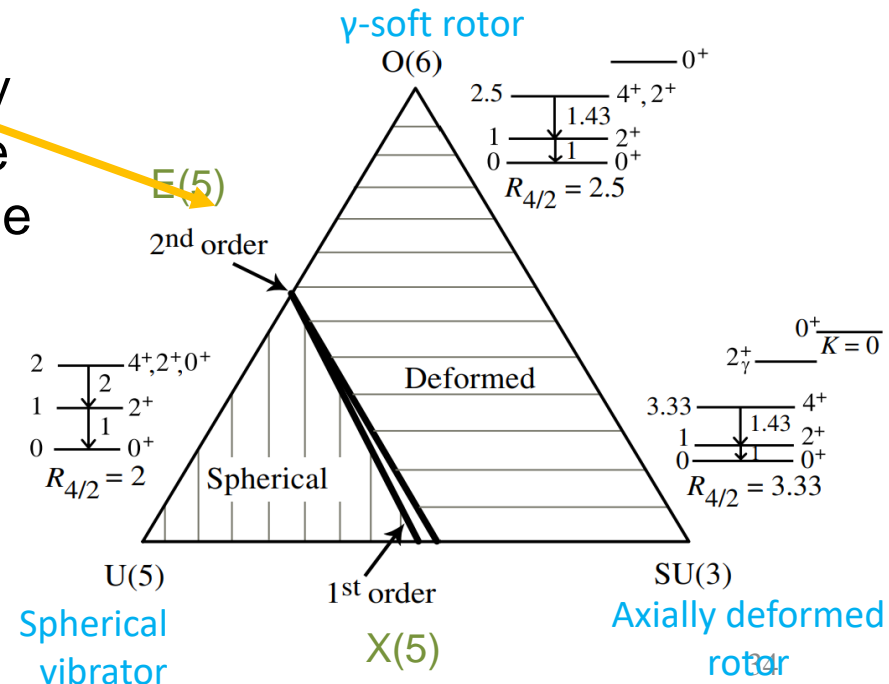
Exploring the Shape Phase Transition along Xe isotopes

-- a short summary

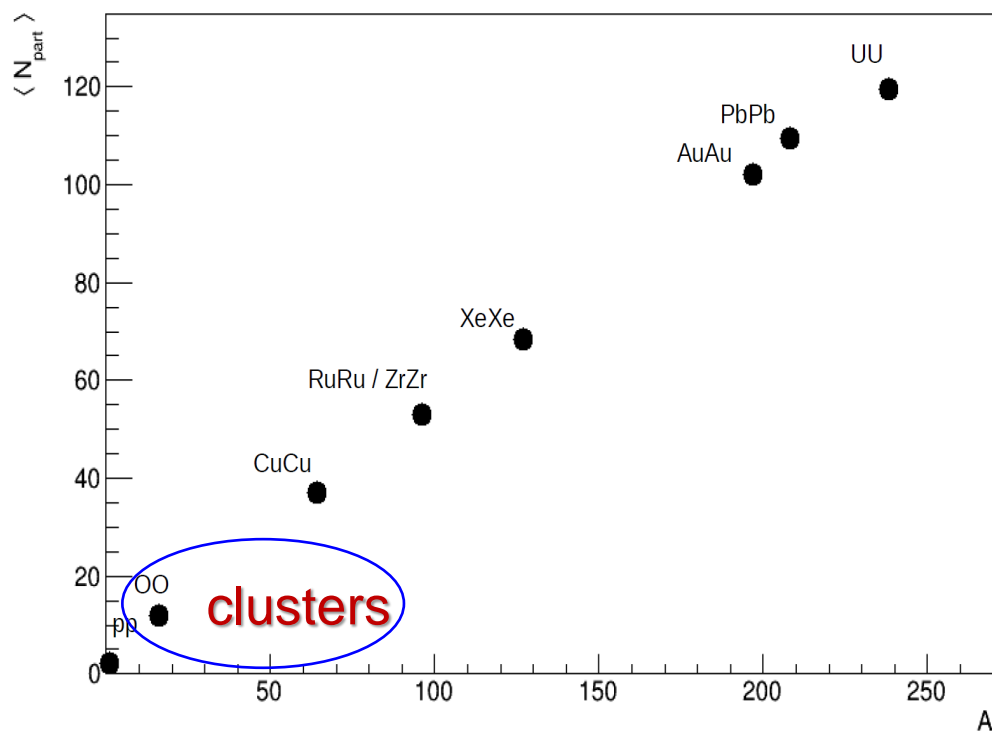


In nuclear structure physics, it is generally believed that Xe isotope chains experience a second-order shape phase transition in the vicinity of $^{128-130}\text{Xe}$.

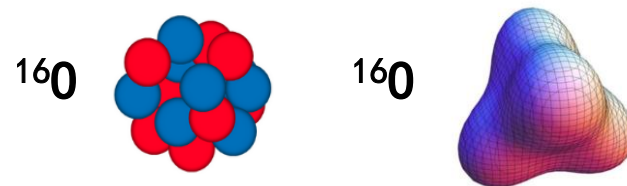
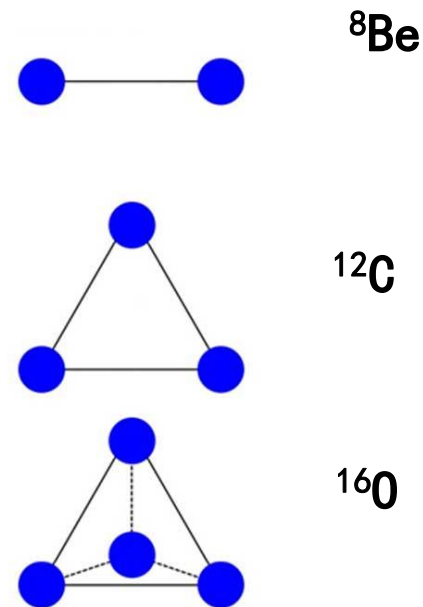
Relativistic heavy ion collisions (such as $^{129}\text{Xe} + ^{129}\text{Xe}$) can not only study the QCD phase transition, but also explore such shape phase transition of finite nuclei.



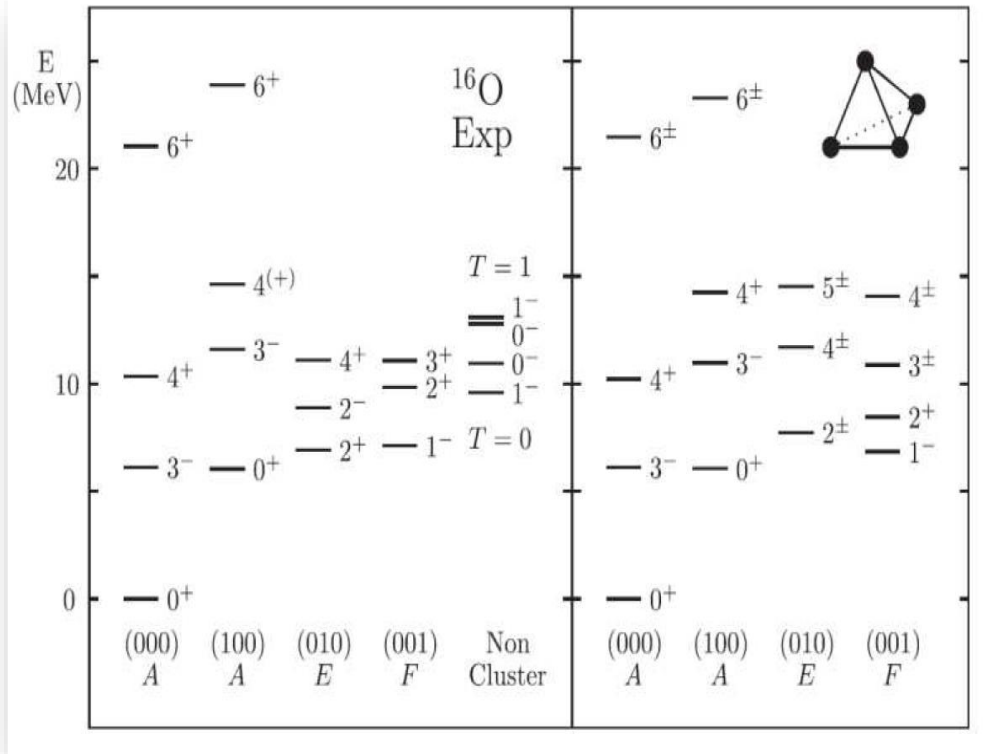
Probe the α -cluster of ^{16}O at RHIC and the LHC



$^{16}\text{O}+^{16}\text{O}$ collisions and $p+^{16}\text{O}$ collisions originally aim to study the possible formation of the QGP in small systems

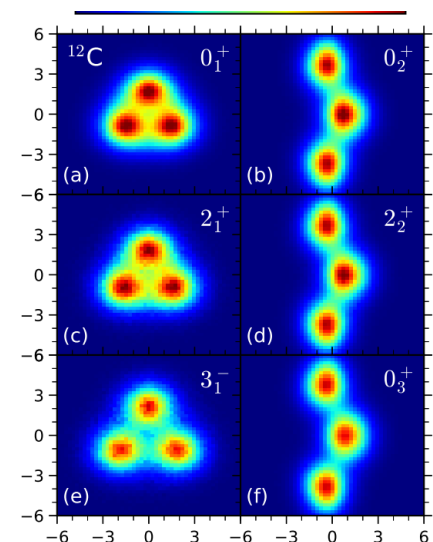
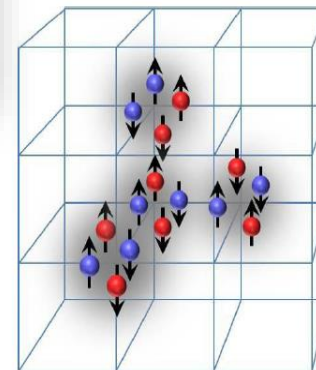


α -cluster of ^{16}O from nuclear structure



-ACM calculations show that the low-lying states of ^{16}O can be described as rotation-vibration of a 4α cluster with tetrahedral symmetry.

R.Bijker and F.Iachello, Phys. Rev. Lett. 112, no.15, 152501 (2014)



-ab initio lattice calculations demonstrate the nucleons are arranged in a tetrahedral alpha clusters in the ground state

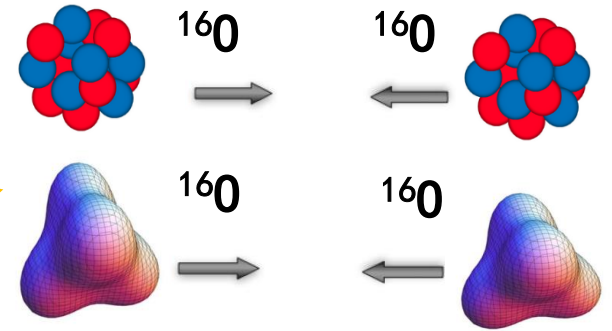
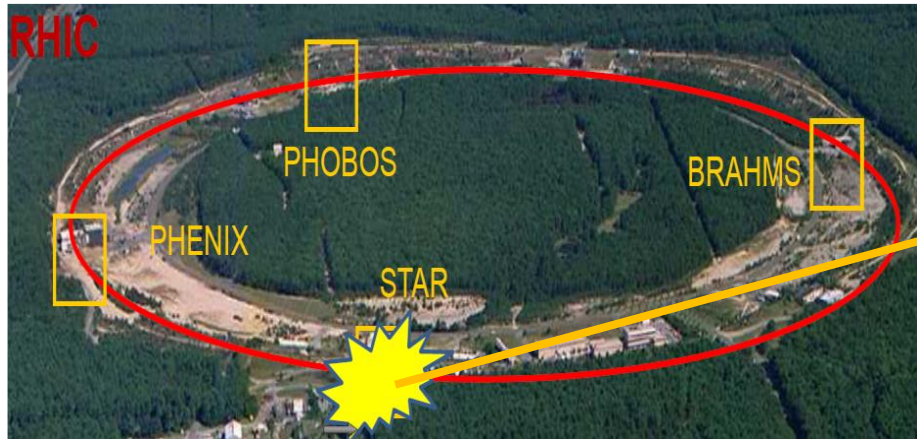
E. Epelbaum, et al Phys. Rev. Lett.112, no.10, 102501 (2014)

Recent NLEFT calculations for light nuclei: Intrinsic shape composed of alpha clusters

Shen, Elhatisari, Lahde, Lee, Lu, UGM, Nature Commun. **14** (2023) 2777

Dee Lee talk, today

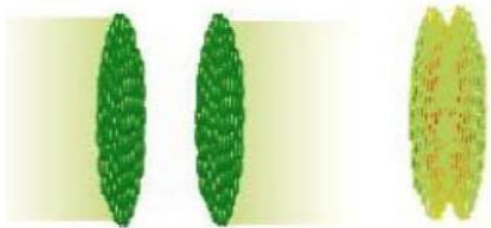
Relativistic heavy ion collision to probe the structure of ^{16}O



**initial conditions:
(with or without α -cluster)**

Well calibrated calculations

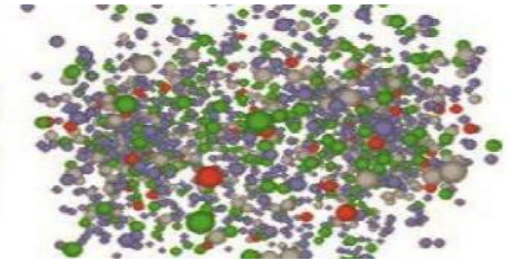
Initial conditions



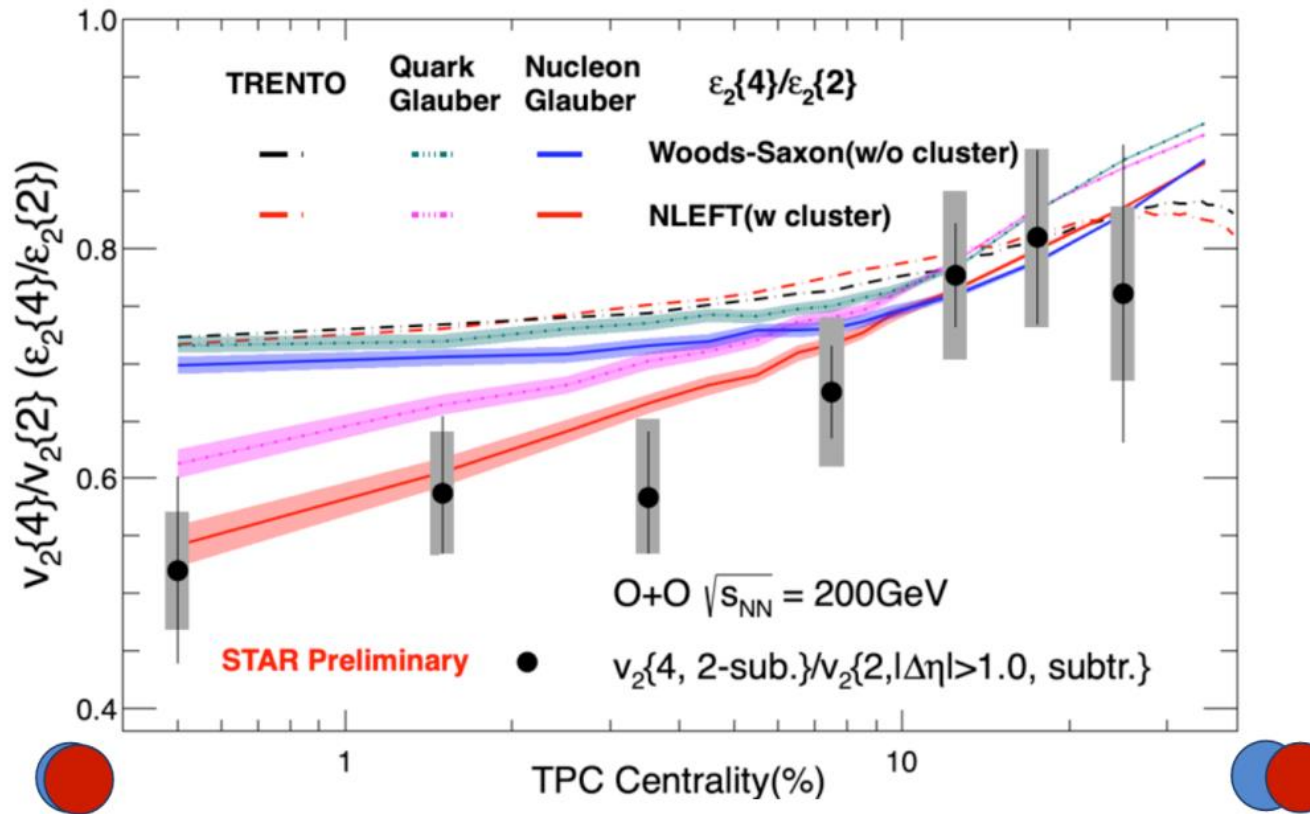
viscous hydro



hadron cascade



Measurement from $^{16}\text{O}+^{16}\text{O}$ collisions



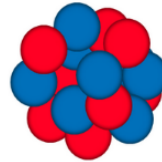
- $v_2\{4\}/v_2\{2\}$: enhanced fluctuations in ultra-central collisions
- heavy ion collision data hint alpha-clustering in ^{16}O

Hydrodynamic calculation w/wo clustering

Initial conditions (TRENTO)

-Woods-Saxon:

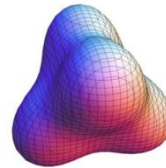
$$\rho(r, \theta, \phi) = \frac{\rho_0}{1 + e^{(r-R(\theta, \phi))/a_0}}$$



Spherical shape

-Alpha-Cluster:

$$f_i(\mathbf{r}) = A \exp \left[-\frac{3(\mathbf{r} - \mathbf{r}_i)^2}{2r_\alpha^2} \right]$$



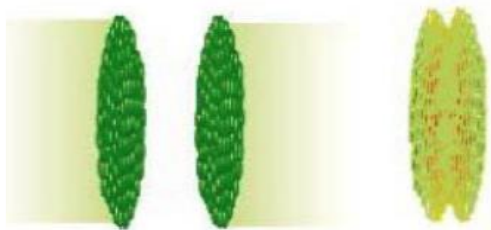
tetrahedral alpha clusters

Y. Wang, S. Zhao, B. Cao, H. Xu and H. Song
arXiv: 2401.15723 [nucl-th].

	distribution	l	r_α	r_α/l
I	Woods-Saxon			
II	α cluster	3.0	2.0	0.67
III	α cluster	3.6	1.6	0.44
IV	α cluster	4.0	1.2	0.30

Well calibrated calculations

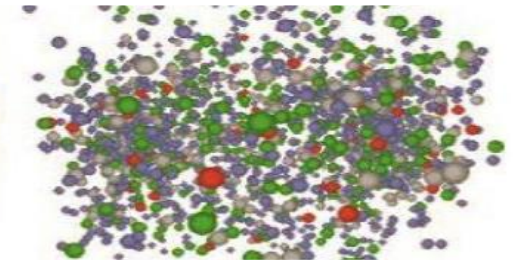
Initial conditions



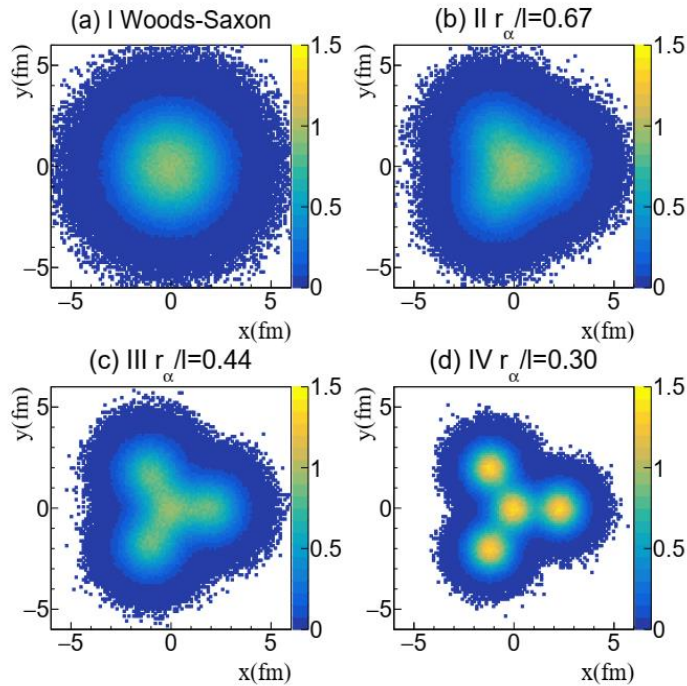
viscous hydro



hadron cascade

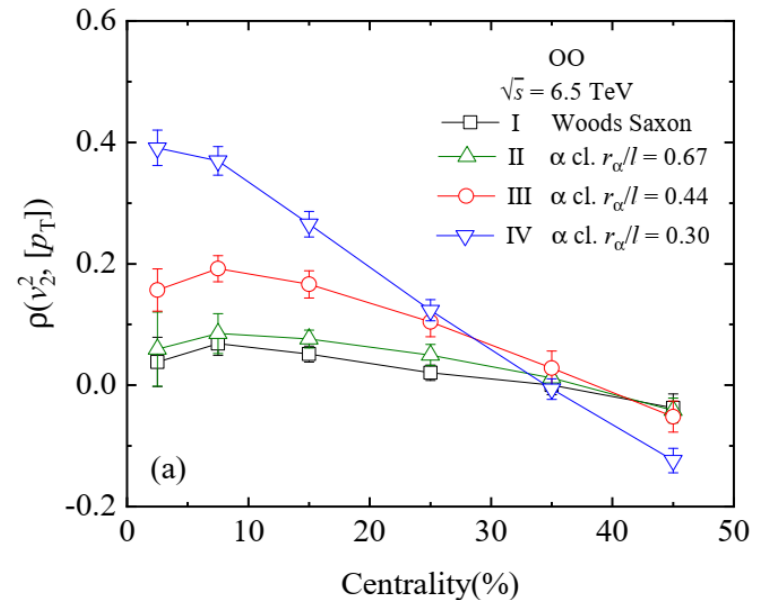
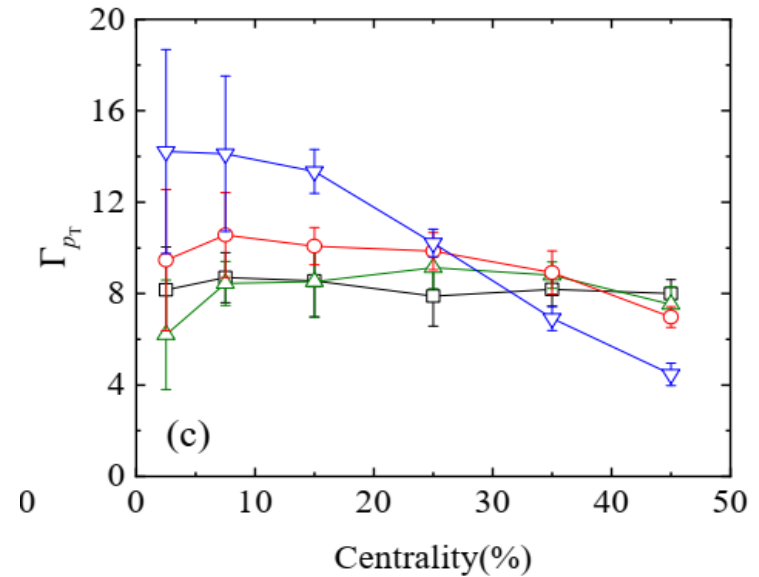


Sensitive observables for α -clustering

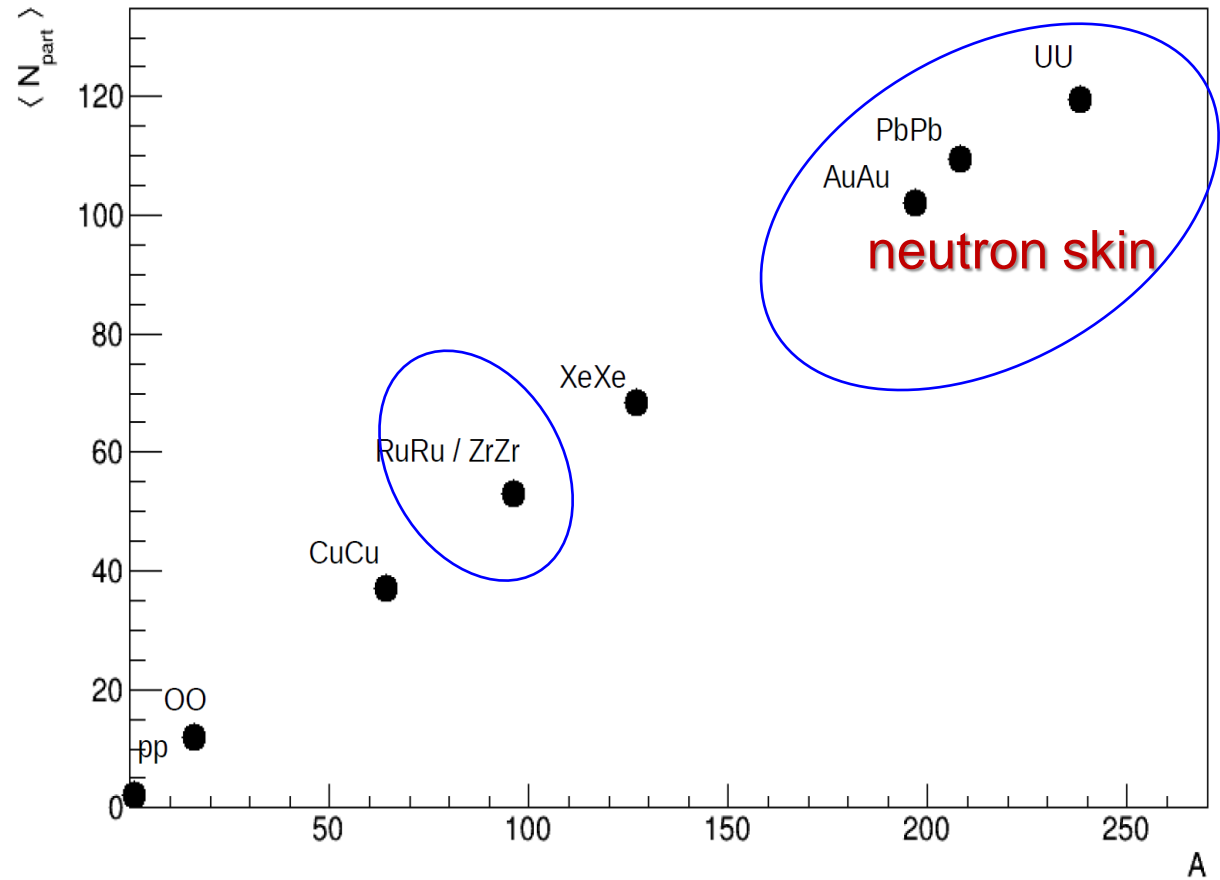
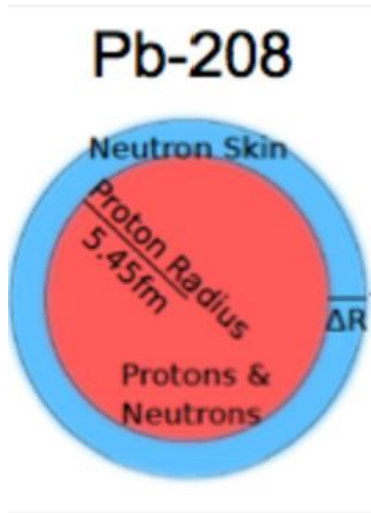


Several observables, such as the correlator Γ the $v_n - p_T$ correlations in $^{16}\text{O}+^{16}\text{O}$ collisions are sensitive to the compactness of the α cluster in the colliding nuclei, which can be used to constrain the detailed configurations of ^{16}O in the future.

Y. Wang, S. Zhao, B. Cao, H. Xu and H. Song. arXiv: 2401.15723 [nucl-th].



Probe neutron skin at RHIC and the LHC



Neutron skin & neutron star

EOS of nuclear matter

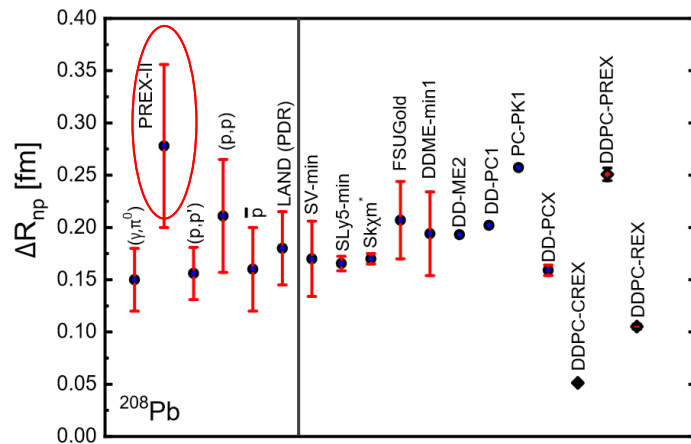
$$\epsilon(\rho, \alpha) = [\epsilon_{SNM}(\rho_0) + S(\rho_0)\alpha^2] + \alpha^2 L \frac{\rho - \rho_0}{3\rho_0} + \frac{1}{2}(K_0 + \alpha^2 K_{sym}) \left(\frac{\rho - \rho_0}{3\rho_0}\right)^2$$

L : the first order term in EOS; symmetry energy; Large L thick neutron skin

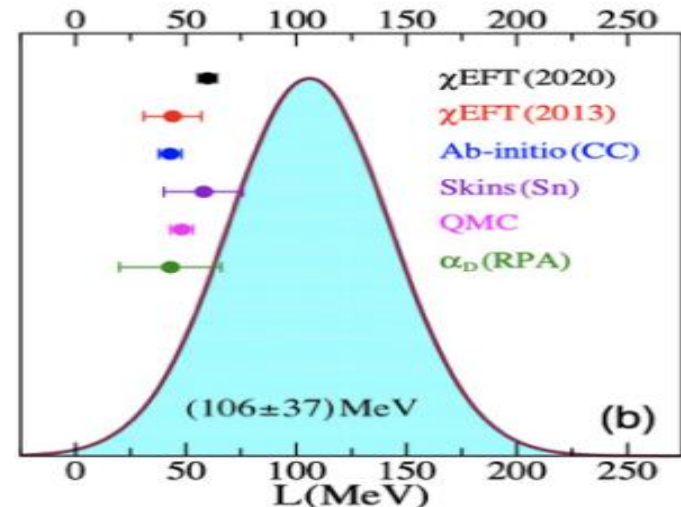
Probe the Neutron Skin at low energy nuclear physics

Parity-Violating Electron Scattering in Jefferson Lab

$$R_{skin}^{208} = 0.278_{-0.078}^{+0.078} fm \quad \text{Phys. Rev. Lett. 126, 172502, (2021)}$$

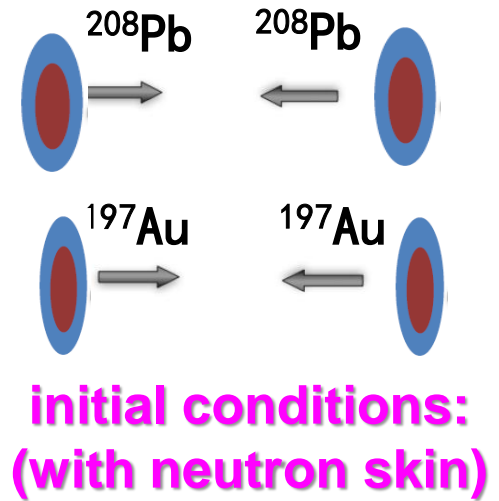
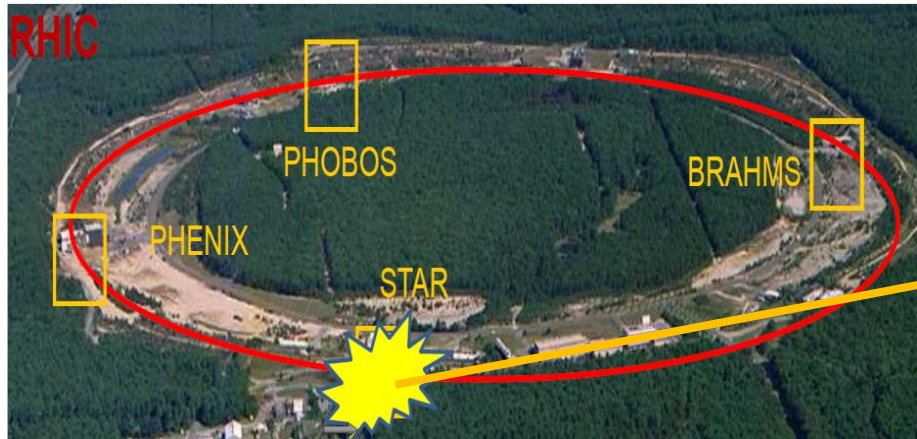


arXiv:2206.06527 Esra Yüksel and Nils Paar



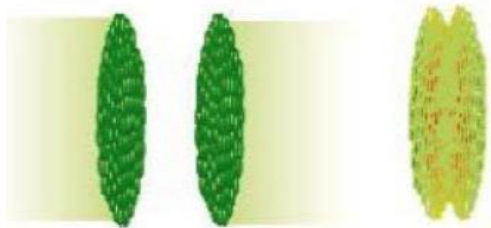
Phys. Rev. Lett. 126, 172502 D. Adhikari et al

Relativistic heavy ion collision to probe the neutron skin



Well calibrated calculations

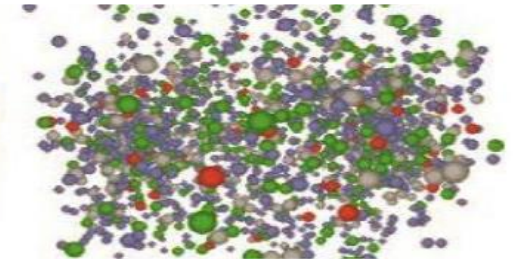
Initial conditions



viscous hydro

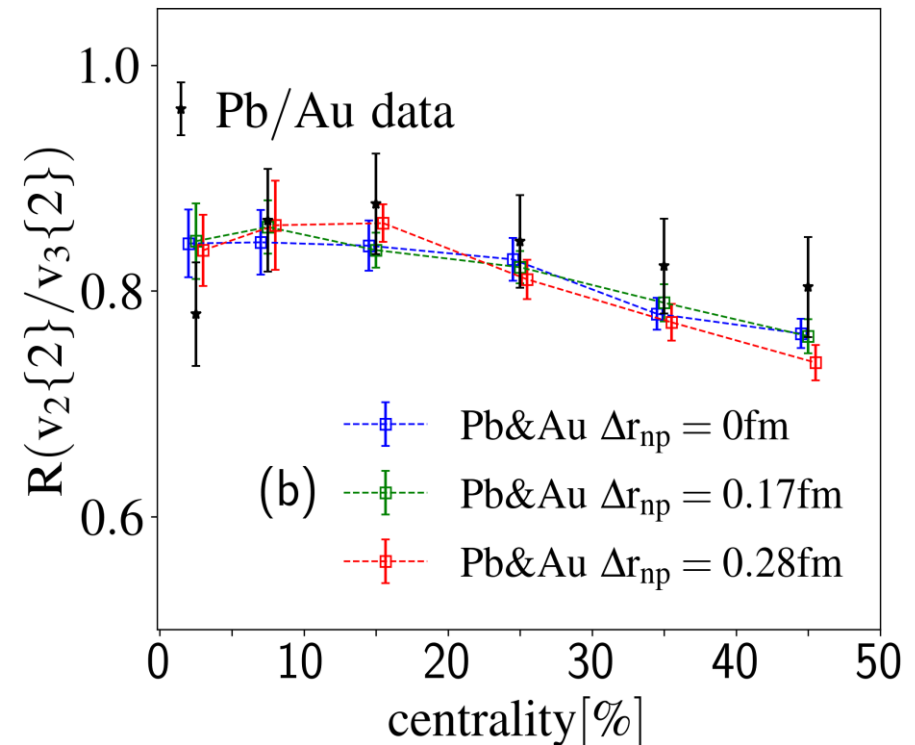
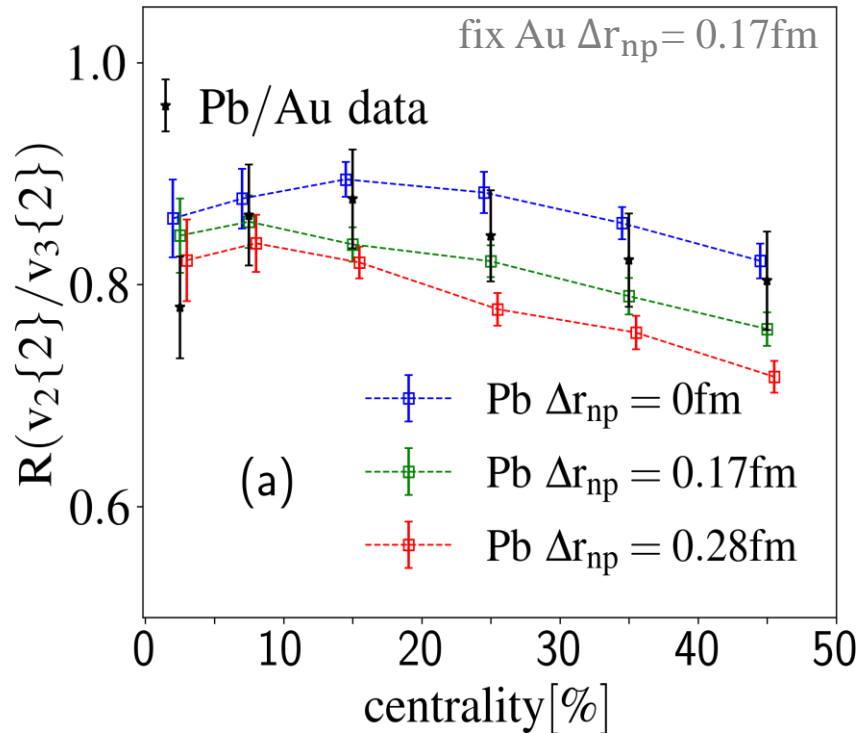


hadron cascade



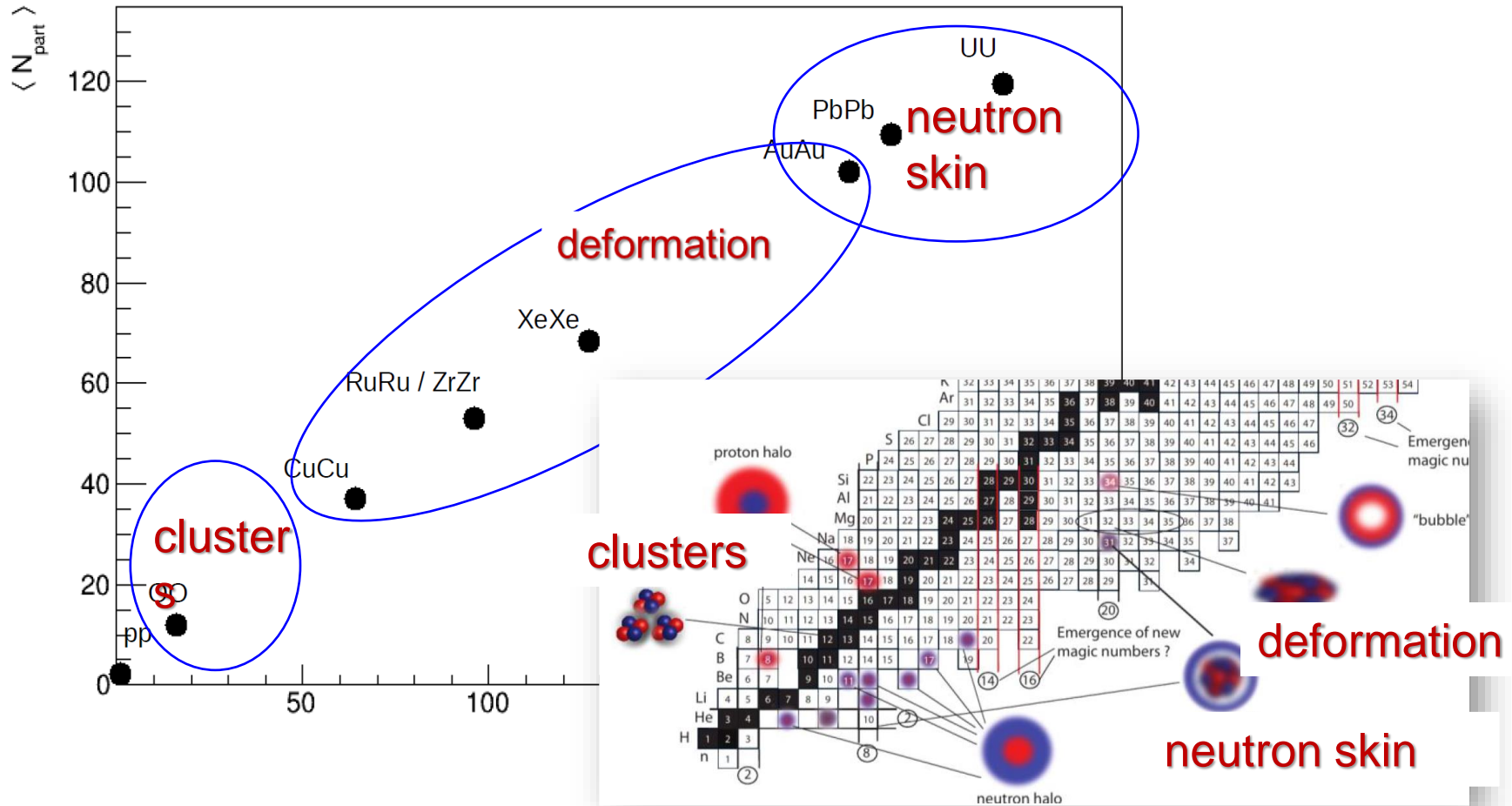
Probing the neutron skin of ^{197}Au and ^{208}Pb

semi-isobaric double ratio



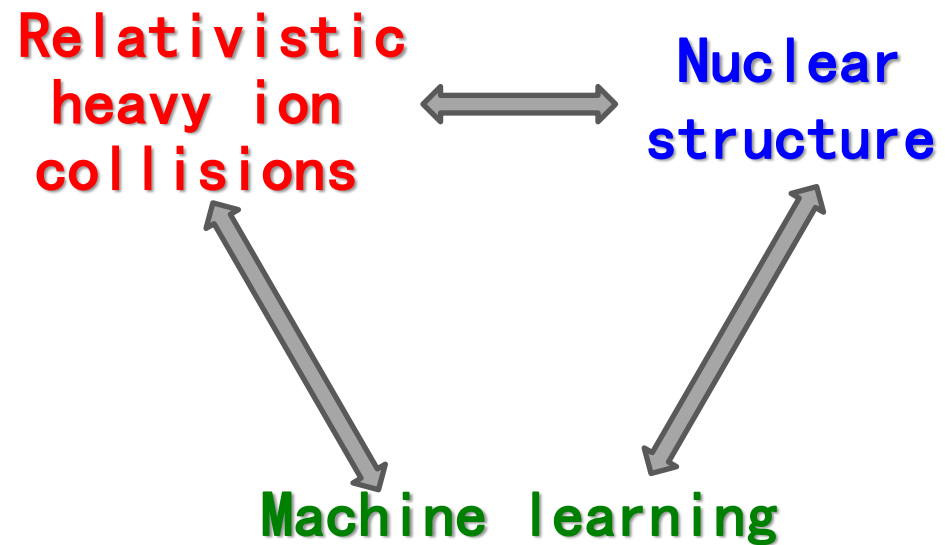
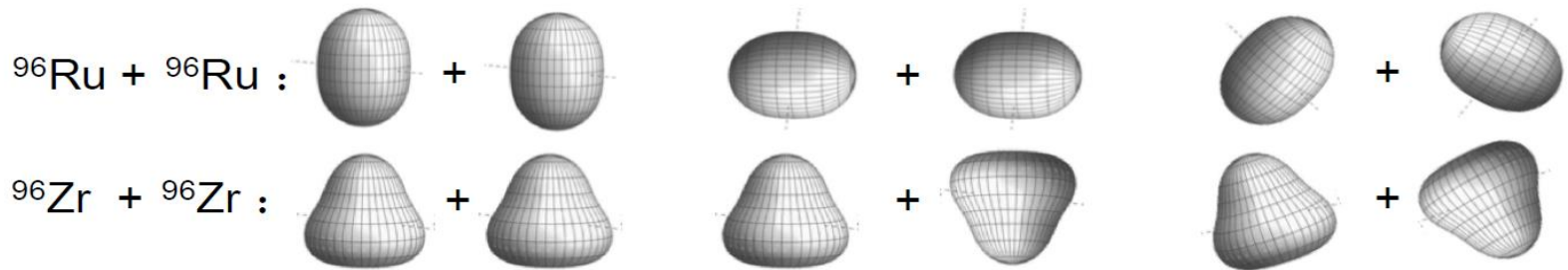
A scaling behavior was found in double ratio of $v_2\{2\}/v_3\{2\}$ when Au and Pb have the neutron skins of the same size,
The measured flow harmonics at various centrality suggest Au and Pb have similar neutron skin

Summary



Relativistic heavy ion collisions have already provide rich collision systems to study various aspects in nuclear structure, there are lots of things to explore !

- Nuclear structure: deformation, cluster, neutron skin; shape coexistence, γ -soft (shape phase transition)
- Rich configurations for QGP initial conditions



Weiyao Ke talk on April 23 Long-Gang Pang talk during the program

Backup