

# AMPT simulation with clustering & nuclear structure effects

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(Fudan University)

[1] Xin-Li Zhao and GLM, Phys.Rev.C 106 (2022) , 034909, arXiv: 2203.15214

[2] Xin-Li Zhao, GLM, You Zhou, Zi-Wei Lin, Chao Zhang, arXiv:2404.09780

[3] Pei Li, Xin-Li Zhao, GLM, et al., in progress



# Outline

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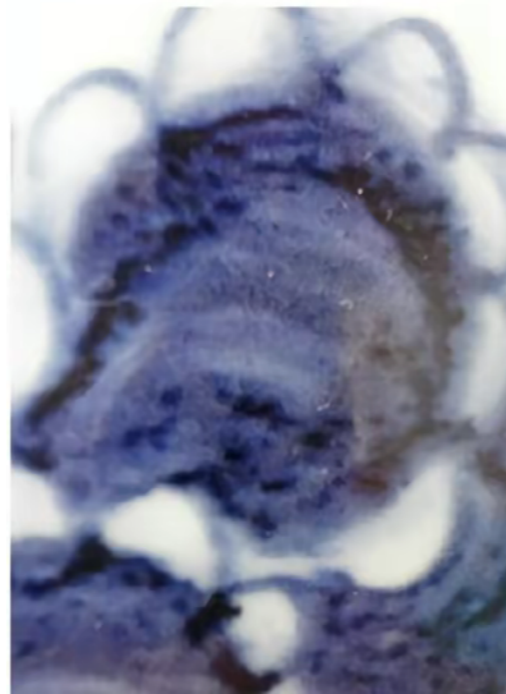
- Introduction
- AMPT simulation on isobar collisions at 200 GeV
- AMPT simulation on O+O collisions at 200 GeV
- Summary and outlook

# Nuclear structure vs porcelain design

梅瓶 (Meiping vase)



青花晕散 (Aqua Diffusion)



钧窑 (Jun kiln plate)

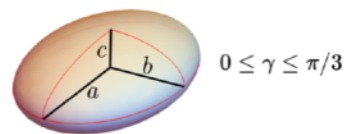


## deformation

$$\rho(r, \theta, \phi) = \frac{\rho_0}{1 + e^{(r-R(\theta, \phi))/a_0}}$$

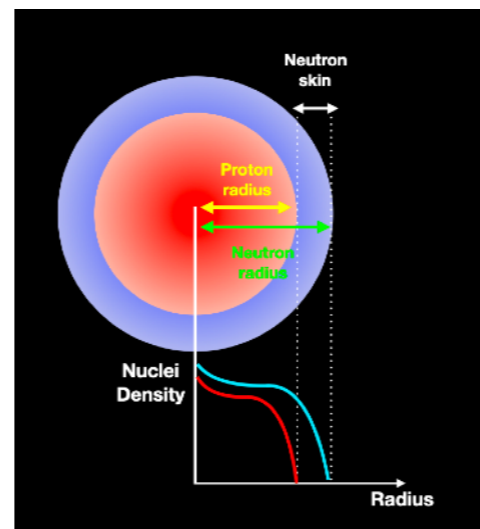
$$R(\theta, \phi) = R_0 \left( 1 + \beta_2 [\cos \gamma Y_{2,0} + \sin \gamma Y_{2,2}] + \beta_3 \sum_{m=-3}^3 \alpha_{3,m} Y_{3,m} + \beta_4 \sum_{m=-4}^4 \alpha_{4,m} Y_{4,m} \right)$$

- Quadrupole:  $1 + \beta_2 Y_{2,0}(\theta, \phi)$   
四极形变
- Octupole:  $1 + \beta_3 Y_{3,0}(\theta, \phi)$   
八极形变
- Hexadecapole:  $1 + \beta_4 Y_{4,0}(\theta, \phi)$   
16极形变

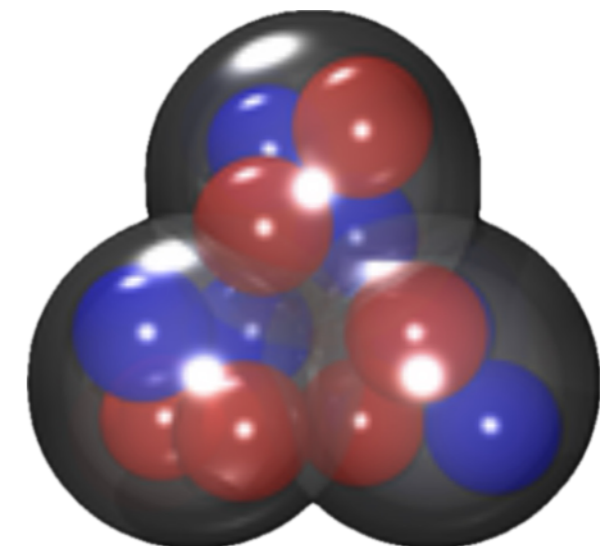


- Prolate:  $a=b < c \rightarrow \beta_2, \gamma=0$
- Oblate:  $a < b=c \rightarrow \beta_2, \gamma=\pi/3$
- Triaxial:  $a < b < c \rightarrow \beta_2, \gamma=\pi/6$

## neutron skin/halo



## $\alpha$ -cluster



# Study nuclear structure using heavy-ion experimental observables

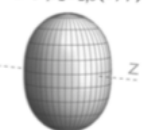
## deformation

$$\rho(r, \theta, \phi) = \frac{\rho_0}{1 + e^{(r-R(\theta, \phi))/a_0}}$$

$$R(\theta, \phi) = R_0 \left( 1 + \beta_2 [\cos \gamma Y_{2,0} + \sin \gamma Y_{2,2}] + \beta_3 \sum_{m=-3}^3 \alpha_{3,m} Y_{3,m} + \beta_4 \sum_{m=-4}^4 \alpha_{4,m} Y_{4,m} \right)$$

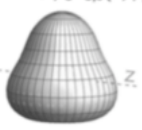
$$1 + \beta_2 Y_{2,0}(\theta, \phi)$$

Quadrupole:  
四极形变



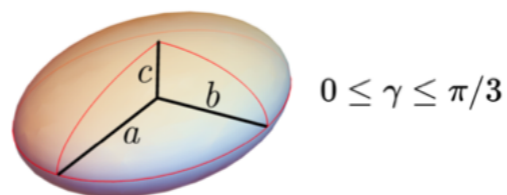
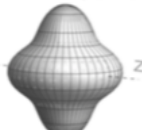
$$1 + \beta_3 Y_{3,0}(\theta, \phi)$$

Octupole:  
八极形变



$$1 + \beta_4 Y_{4,0}(\theta, \phi)$$

Hexadecapole:  
16极形变

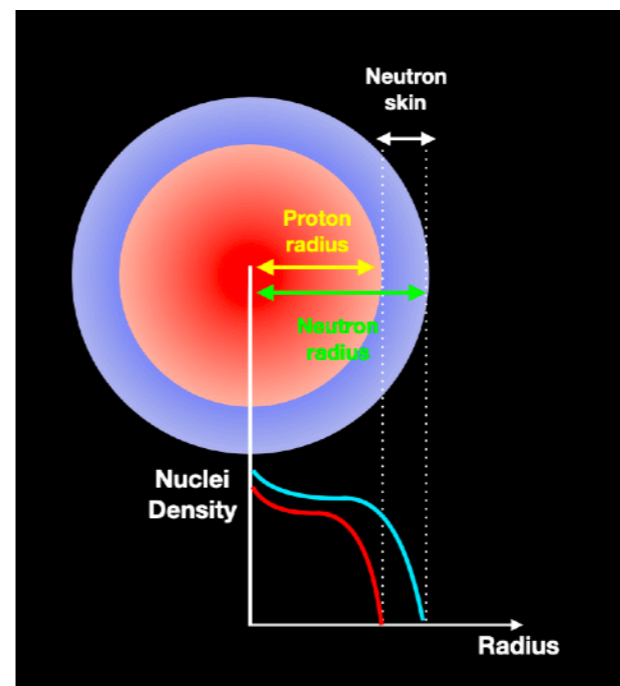


Prolate:  $a=b < c \rightarrow \beta_2, \gamma=0$

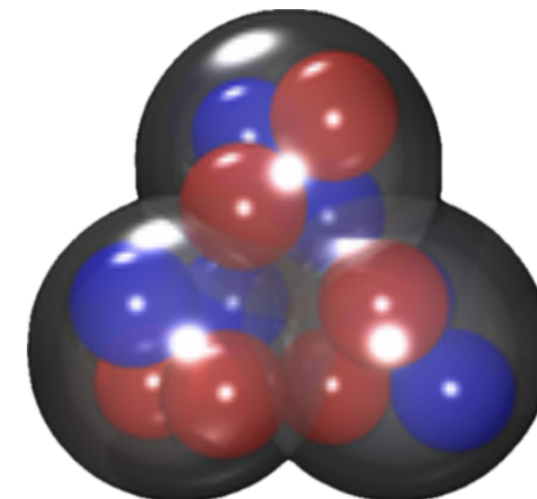
Oblate:  $a < b=c \rightarrow \beta_2, \gamma=\pi/3$

Triaxial:  $a < b < c \rightarrow \beta_2, \gamma=\pi/6$

## neutron skin

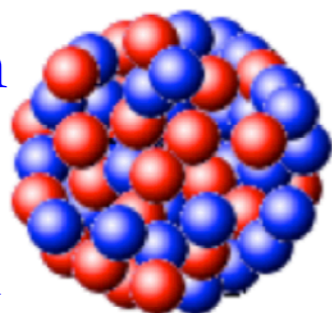


## $\alpha$ -cluster



## Nucleus

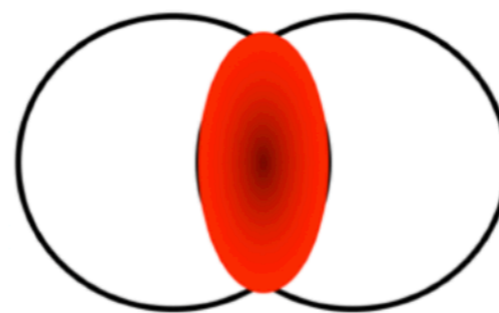
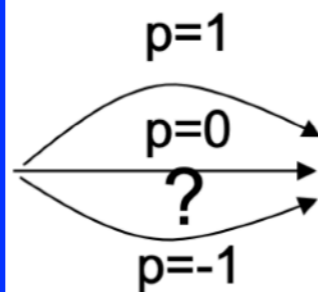
deformation  
 $\alpha$ -cluster  
neutron skin



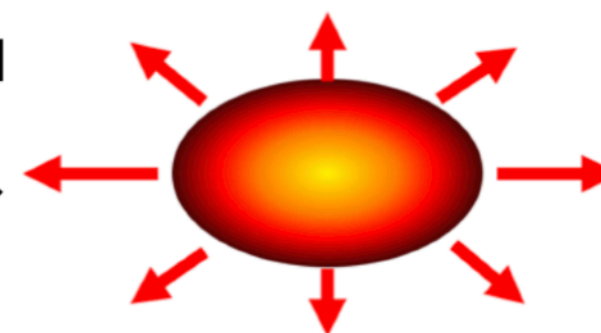
Initial condition

$$v_n = k_n \varepsilon_n$$

Final state



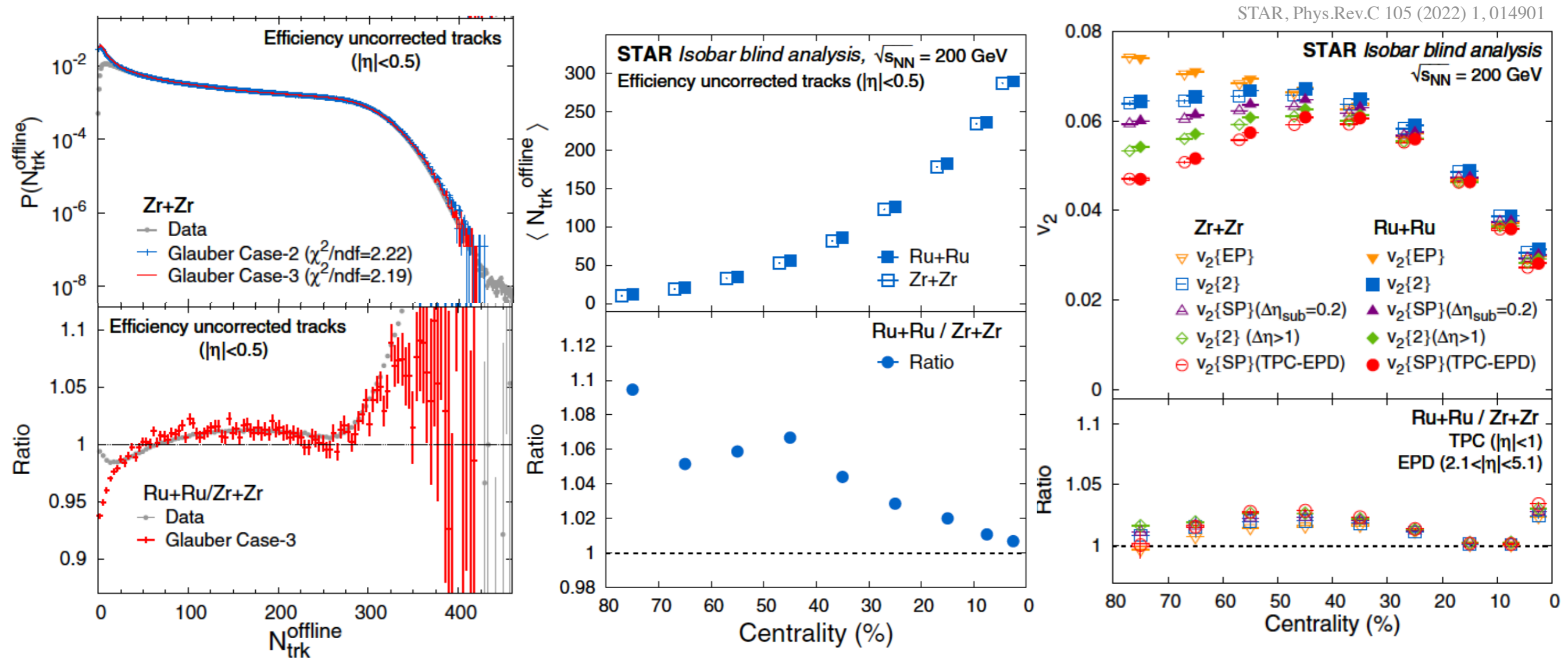
understood



Impact on initial geometry

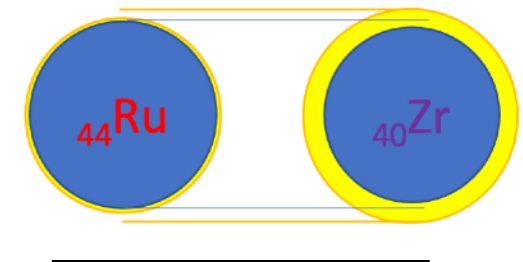
Reflect on collective flow

# Different CME backgrounds between isobar collisions



- Differences in multiplicity distribution,  $\langle N_{ch} \rangle$  and  $v_2$  between two isobar systems.

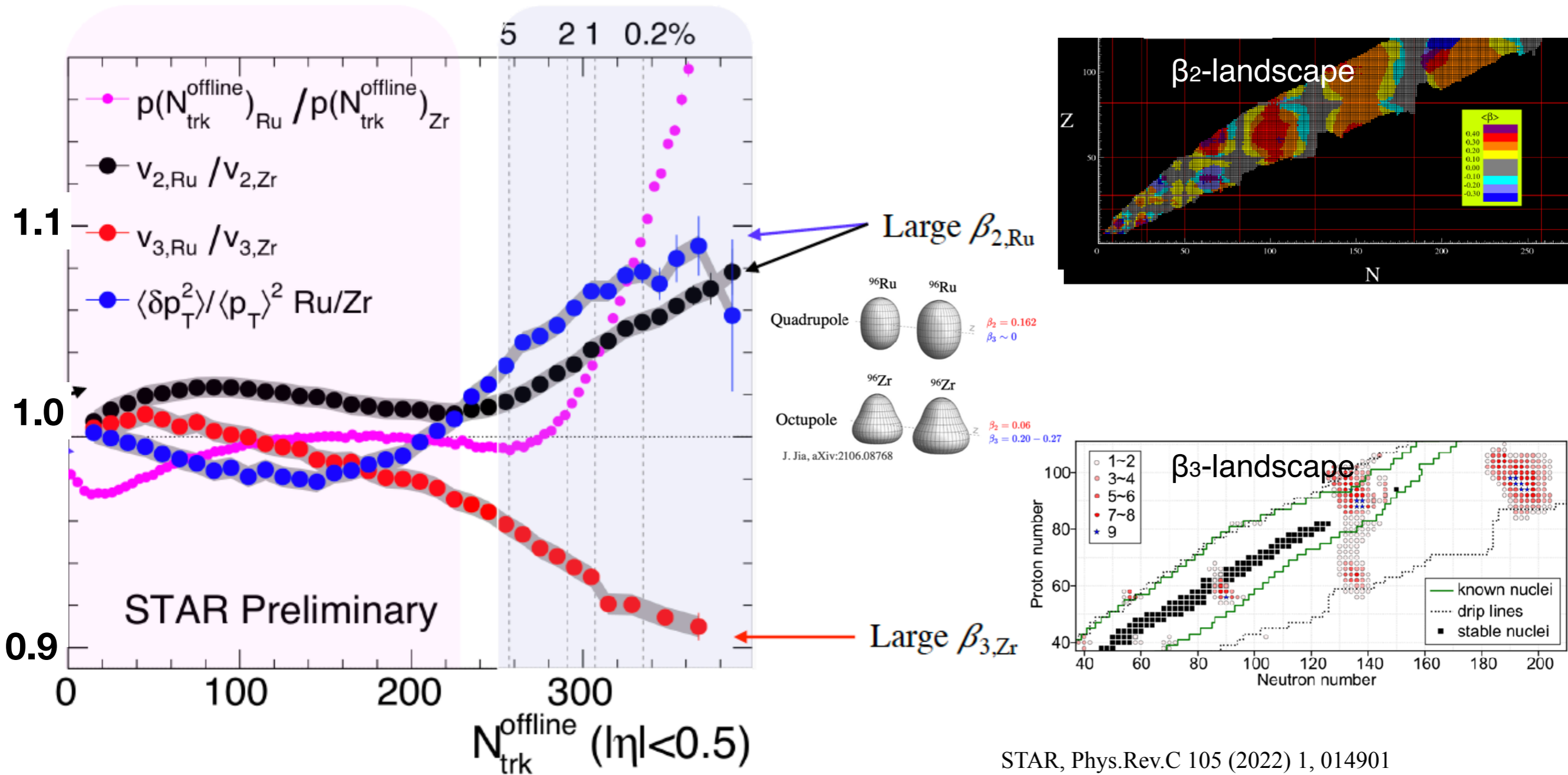
Nucleus	Case-1 [83]			Case-2 [83]			Case-3 [113]		
	$R$ (fm)	$a$ (fm)	$\beta_2$	$R$ (fm)	$a$ (fm)	$\beta_2$	$R$ (fm)	$a$ (fm)	$\beta_2$
$^{96}_{44}\text{Ru}$	5.085	0.46	0.158	5.085	0.46	0.053	5.067	0.500	0
$^{96}_{40}\text{Zr}$	5.02	0.46	0.08	5.02	0.46	0.217	4.965	0.556	0



- Related to nuclear deformation/structure.

Haojie Xu et al. PRL 121 (2018) 022301  
Hanlin Li et al. PRC 98 (2018) 054907

# Nuclear structure in relativistic heavy-ion collisions

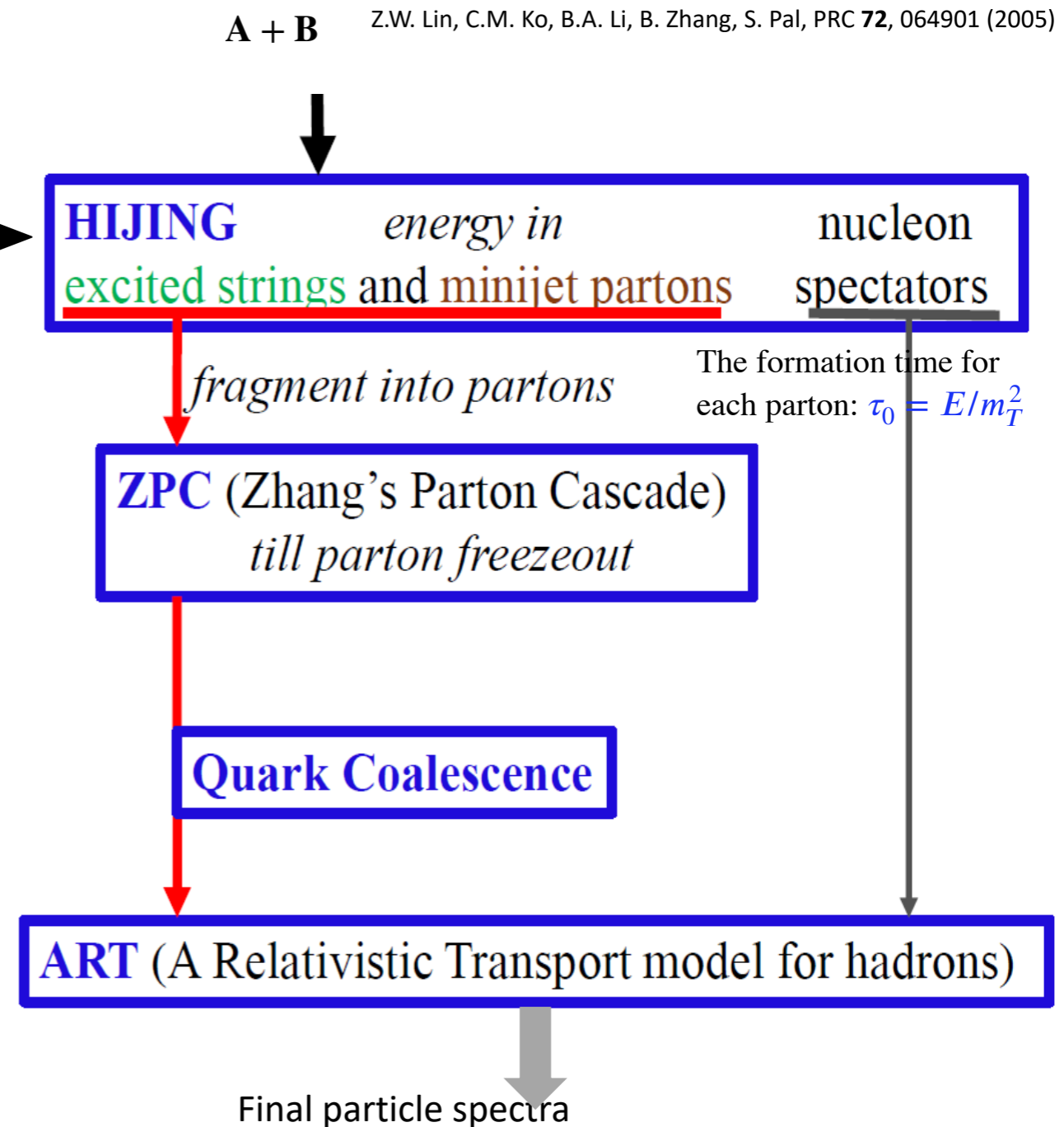
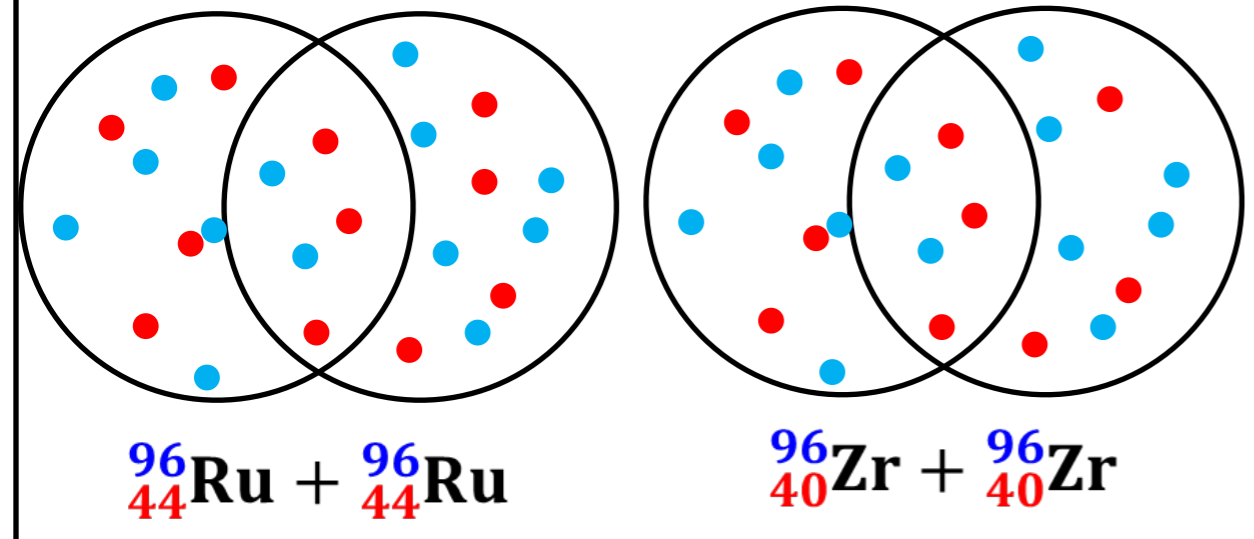


- The ratios of observables between two isobar systems are sensitive to nuclear deformation/structure

STAR, Phys.Rev.C 105 (2022) 1, 014901  
 H. Li et al., Phys.Rev.C 98 (2018) 5, 054907  
 H. J. Xu, Phys.Lett.B 819 (2021) 136453  
 J. Jia, Phys.Rev.C 105 (2022) 4, 044905  
 J. Jia, Phys.Rev.C 105, 014905 (2022)  
 C. Zhang and J. Jia, Phys.Rev.Lett. 128 (2022), 022301  
 G. Giacalone et al., Phys.Rev.Lett. 127 (2021), 242301  
 J. Jia and C. Zhang, Phys.Rev.C 107 (2023) 2, L021901

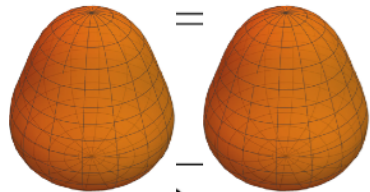
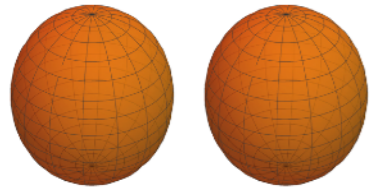
# Simulating isobar collisions using AMPT

## Initialization of isobar



# Initialization of isobar nuclei with geometry

## Woods-Saxon form of spatial distribution of nucleons:

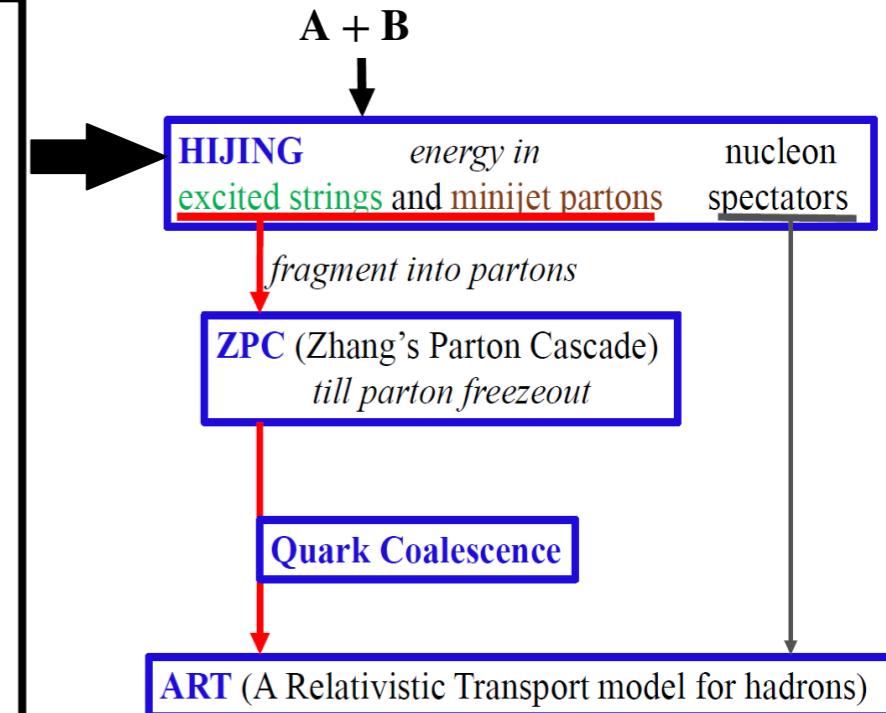


$$\rho(r, \theta, \phi) \propto \frac{1}{1 + e^{[r - R_0(1 + \beta_2 Y_2^0(\theta, \phi) + \beta_3 Y_3^0(\theta, \phi))]/a}}$$

	Case1	old Case2			Case1			Case2			Case3		
	$\beta_2$	$R_0$	a	$\beta_2$	$R_0$	a	$\beta_2$	$R_0$	a	$\beta_2$	$R_0$	a	$\beta_2$
${}^{96}_{44}\text{Ru}$	0.13	5.13	0.46	0.03	5.085	0.46	0.158	5.085	0.46	0.053	5.067	0.500	0
${}^{96}_{40}\text{Zr}$	0.06	5.06	0.46	0.18	5.02	0.46	0.080	5.02	0.46	0.217	4.965	0.556	0

	Case4				Case5				Case6				Case7			Case8		
	$R_0$	a	$\beta_2$	$\beta_3$	$R_0$	a	$\beta_2$	$\beta_3$	$R_0$	a	$\beta_2$	$\beta_3$	$R_0$	a	$\beta_2$	$R_0$	a	$\beta_2$
${}^{96}_{44}\text{Ru}$	5.09	0.46	0.162	0	5.09	0.46	0.162	0	5.09	0.52	0.154	0	5.065	0.485	0.16	5.085	0.523	0
${}^{96}_{40}\text{Zr}$	5.09	0.52	0.060	0.2	5.02	0.46	0.060	0.2	5.09	0.52	0.060	0.2	4.961	0.544	0.16	5.021	0.523	0

	Case9			Case10			Case11			skin-type			halo-type		
	$R_0$	a	$\beta_2$	$R_0$	a	$\beta_2$	$R_0$	a	$\beta_2$	$R_0$	a	$\beta_2$	$R_0$	a	$\beta_2$
${}^{96}_{44}\text{Ru}$ n	5.075	0.505	0	5.073	0.490	0.16	5.085	0.46	0.158	5.085	0.523	0	5.085	0.523	0
${}^{96}_{44}\text{Ru}$ p	5.060	0.493	0	5.053	0.480	0.16	5.085	0.46	0.158	5.085	0.523	0	5.085	0.523	0
${}^{96}_{40}\text{Zr}$ n	5.015	0.574	0	5.007	0.564	0.16	5.080	0.46	0	5.194	0.523	0	5.021	0.592	0
${}^{96}_{40}\text{Zr}$ p	4.915	0.521	0	4.912	0.508	0.16	5.080	0.34	0	5.021	0.523	0	5.021	0.523	0



Final particle spectra

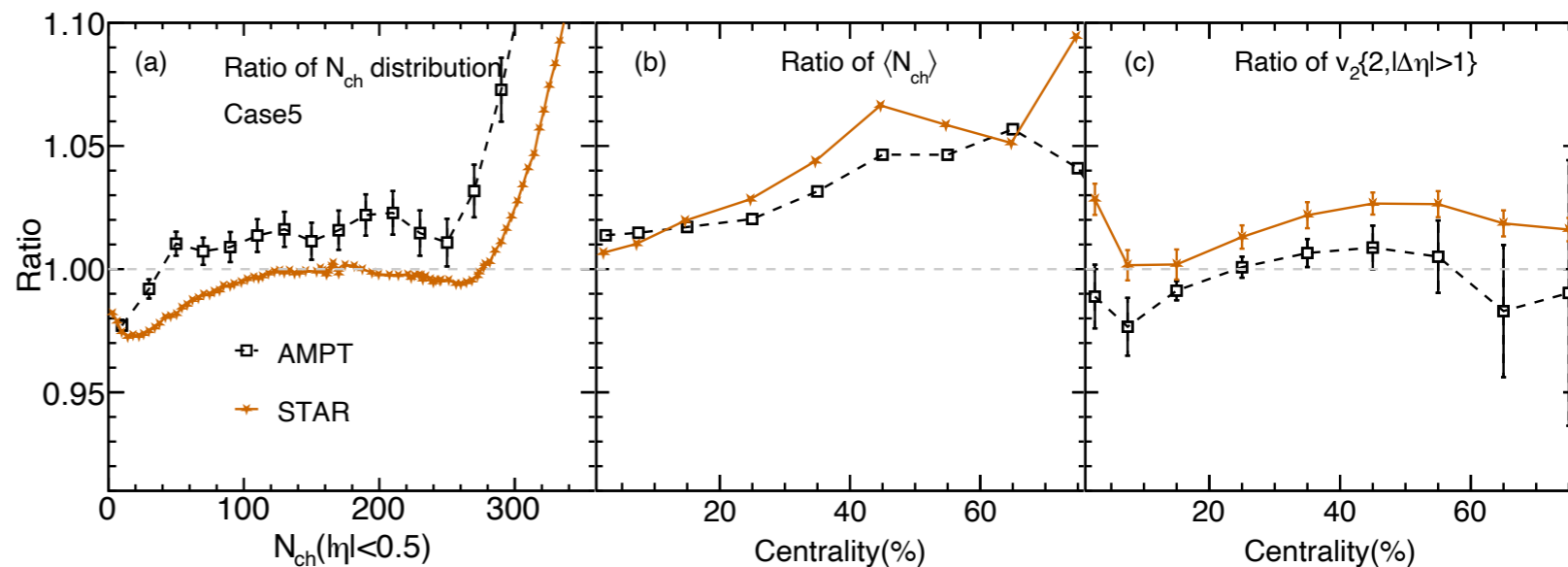
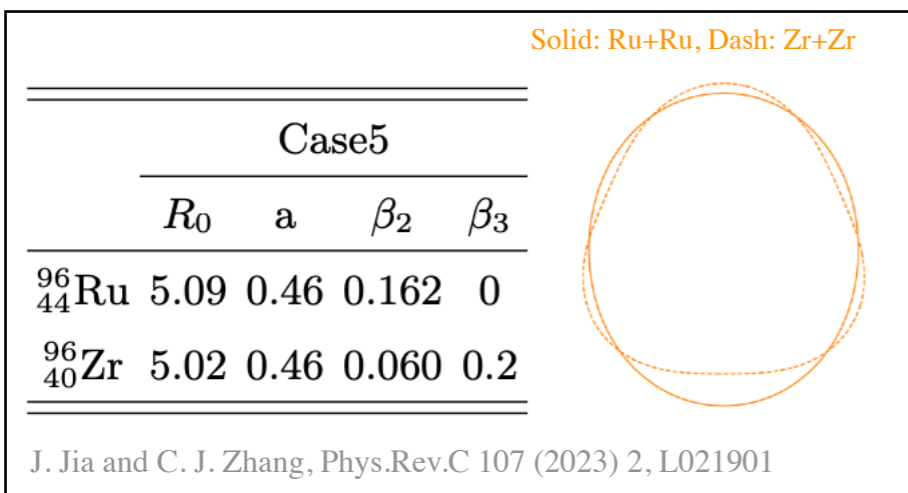
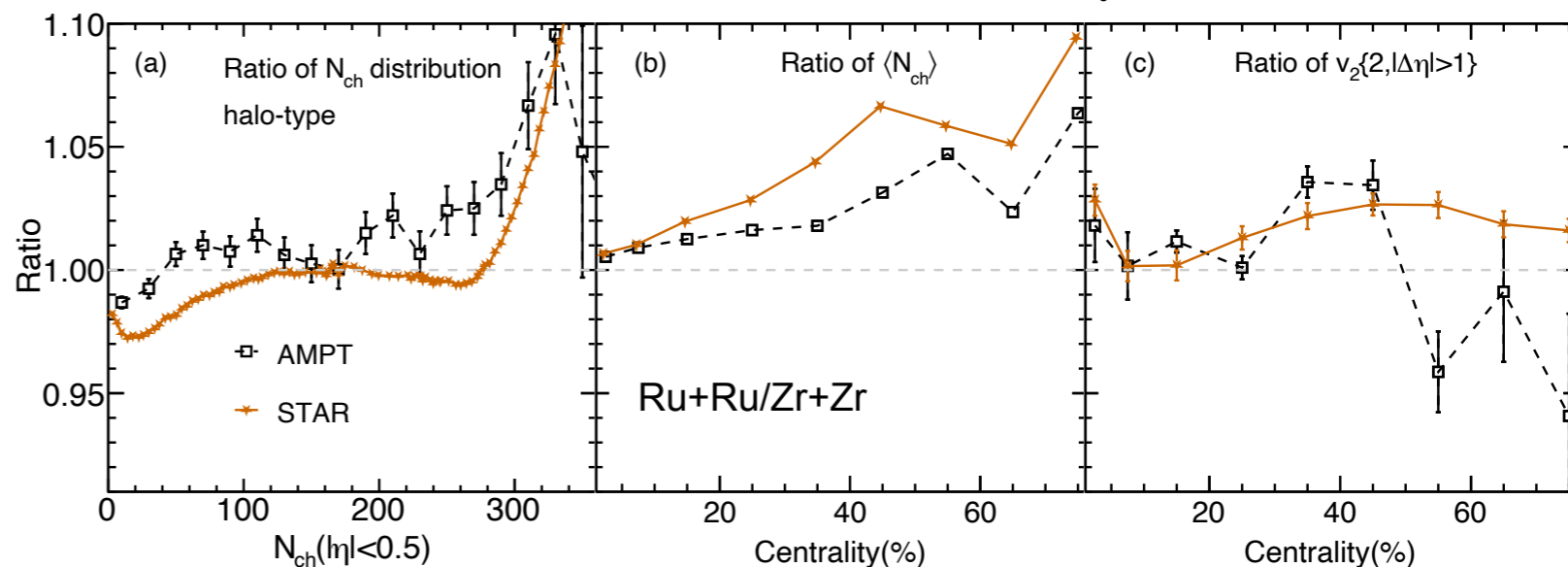
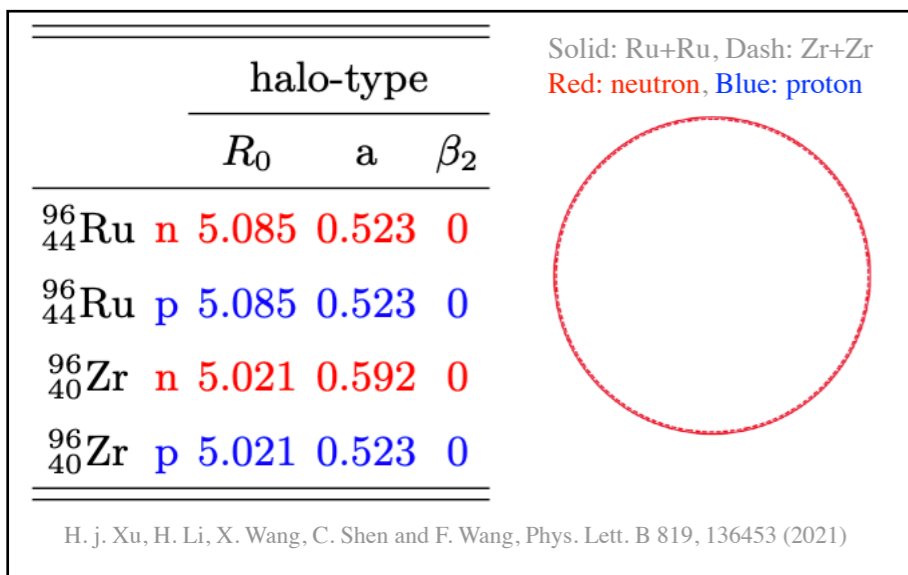
- Which is better or worse?
- Three ratios are our judging criteria:
  - 1) Mult. dist. ratio
  - 2)  $\langle N_{ch} \rangle$  ratio vs centrality
  - 3)  $v_2$  ratio vs centrality

\* 1 M events for each case are used to test if it can pass the criterion test.



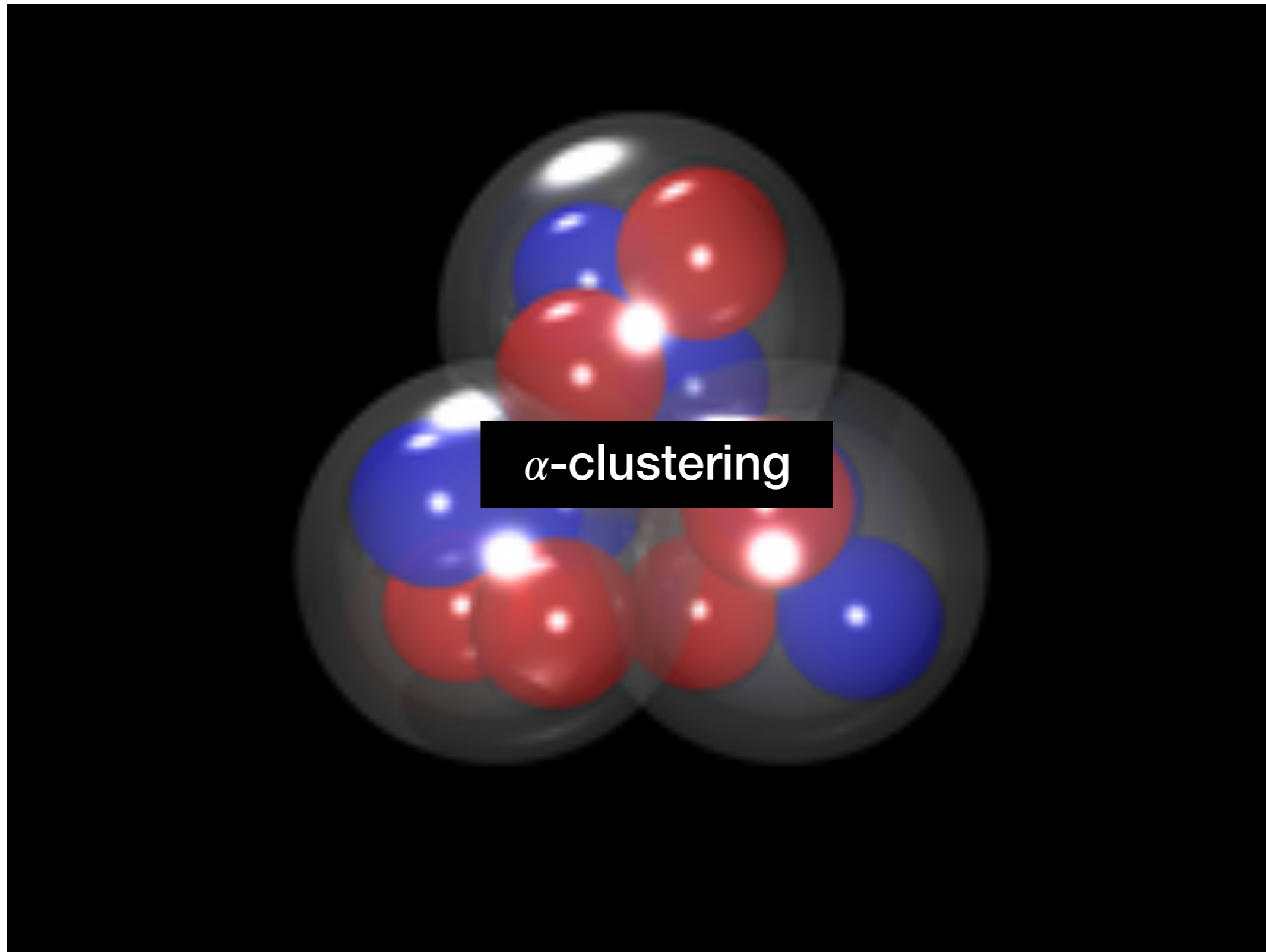
# AMPT performance: Halo-type vs deformation

Xin-Li Zhao and GLM, Phys. Rev. C 106, 034909 (2022)

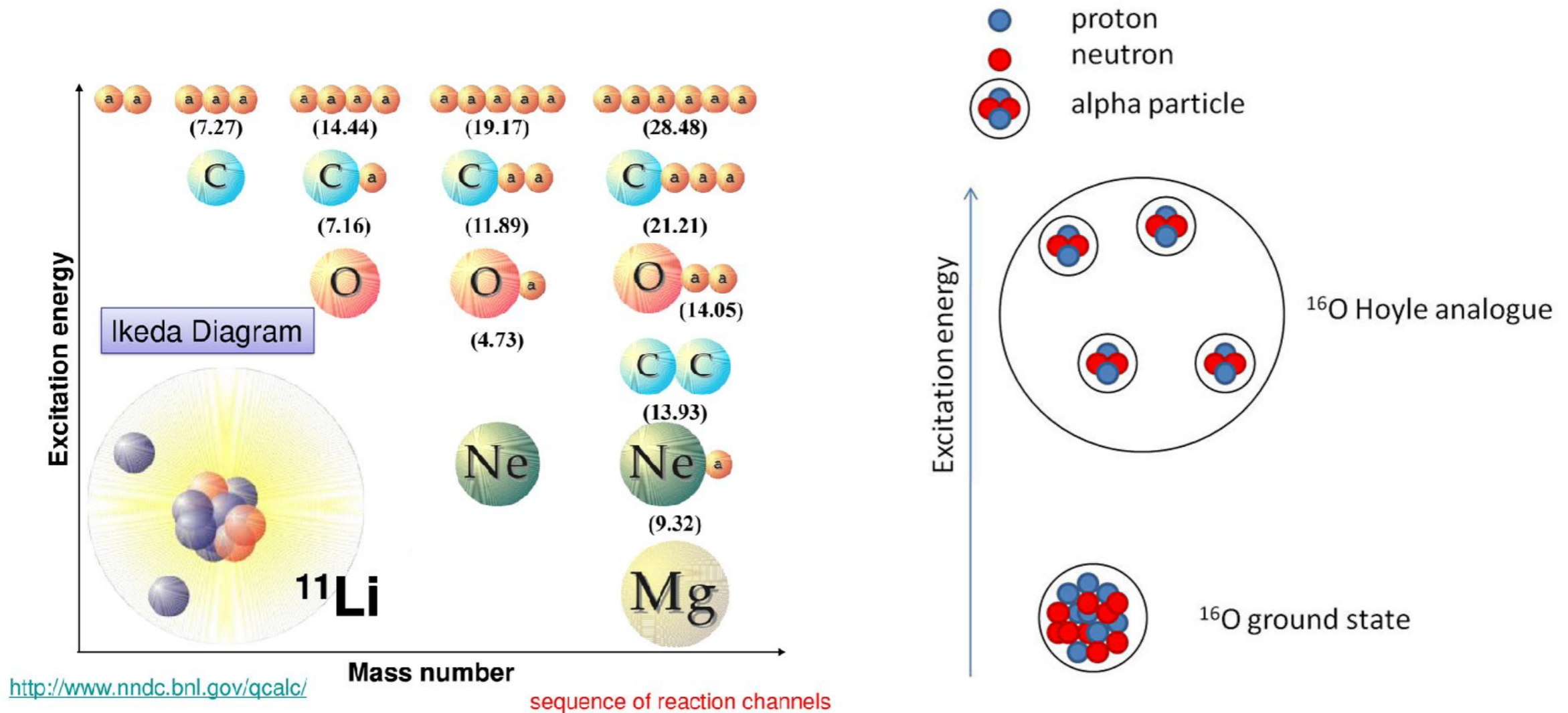


$\langle \chi^2 \rangle$	old Case 1	old Case 2	Case 1	Case 2	Case 3	Case 4
	0.204	0.682	0.255	0.400	0.097	0.053
$\langle \chi^2 \rangle$	Case 5	Case 6	Case 7	Case 8	Case 9	Case 10
	0.049	0.224	0.051	0.227	0.057	0.048
$\langle \chi^2 \rangle$	Case 11	skin-type	halo - type	Case 1 w/ $\beta_3$	Case 2 w/ $\beta_3$	Case 3 w/ $\beta_3$
	0.430	0.177	0.047	0.247	0.506	0.166

- Isobar nuclei are with halo-type neutron skin or deformed



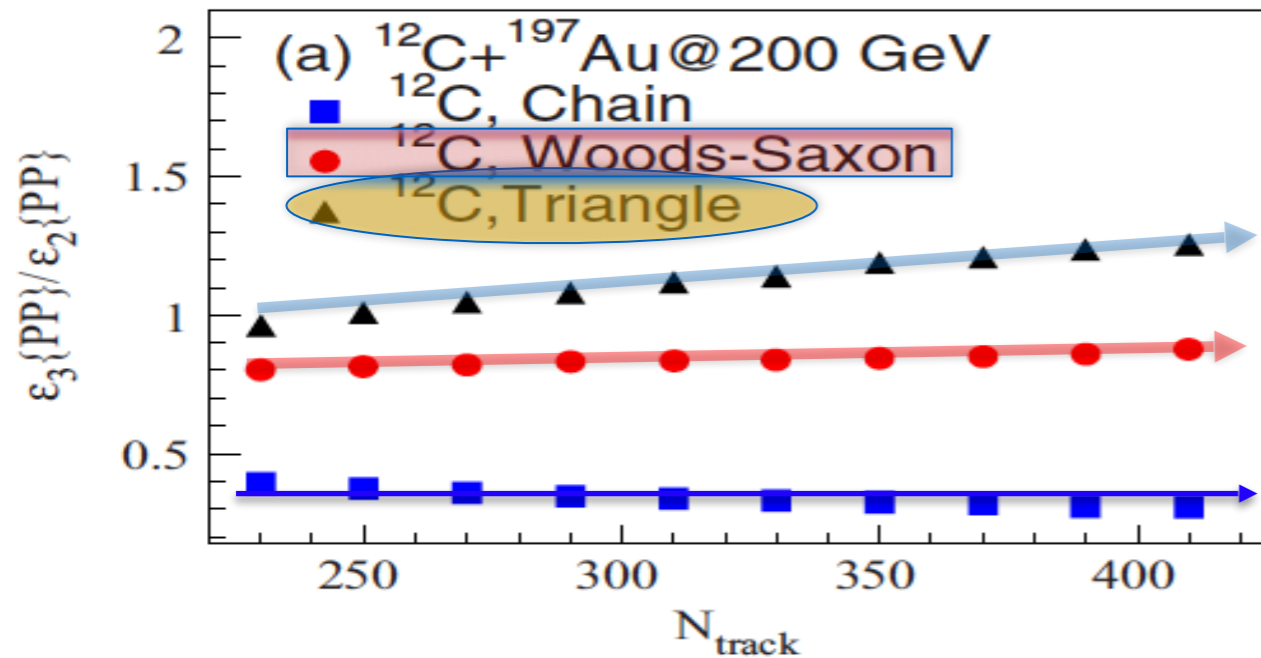
# Significance of $\alpha$ -clustering structure



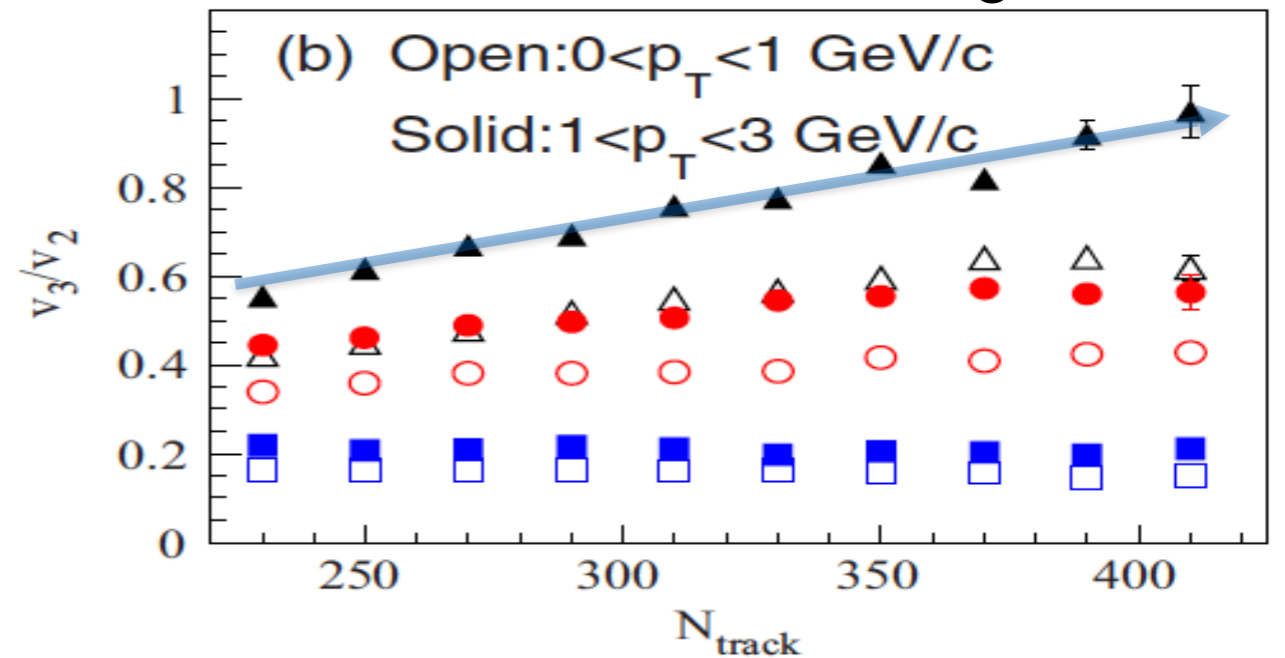
- $\alpha$ -clustering structure inside some nuclei
- The spatial distribution of  $\alpha$ -clustering nuclei is diffuse because condensed  $\alpha$ -particles are weakly bound
- Heavy-ion collisions are helpful to detect  $\alpha$ -clustering structures?

# $\alpha$ -cluster effect in C+Au collisions

Song Zhang, Yu-Gang Ma, et al., Phys. Rev. C 95 (2017) 064904.

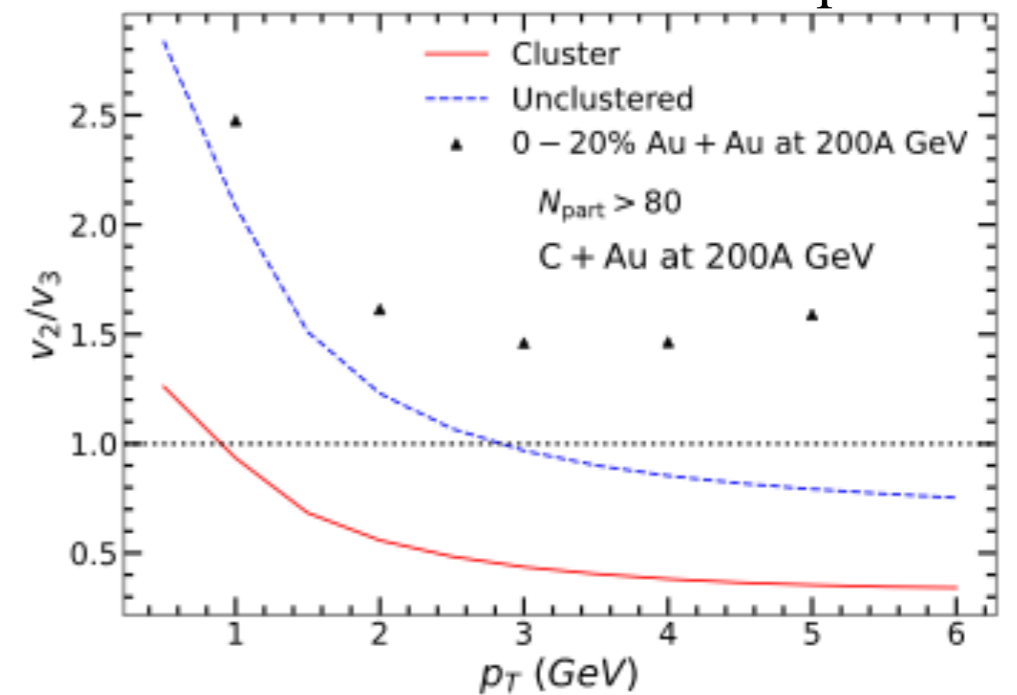
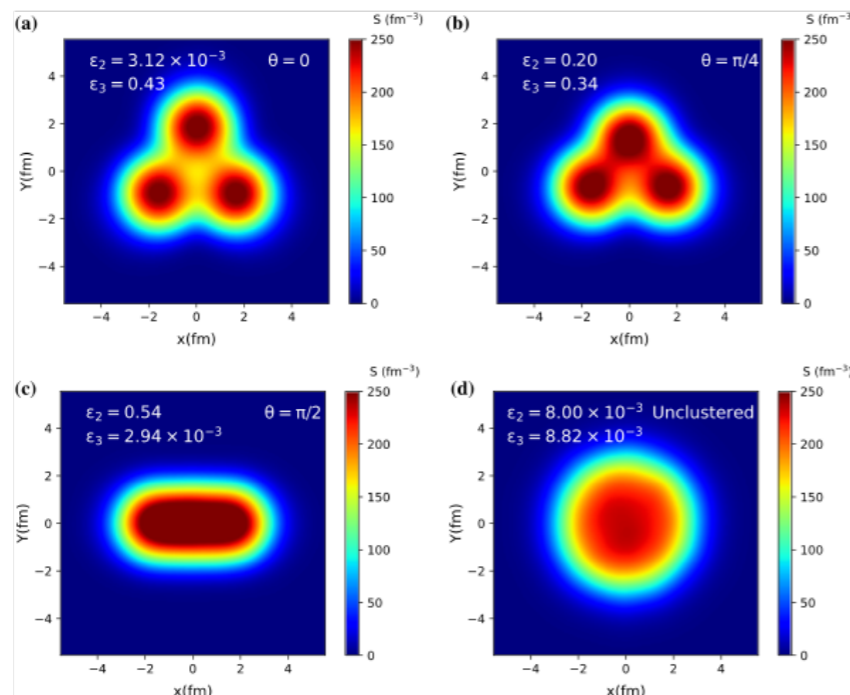


Charged hadron



P. Dasgupta, R. Chatterjee, GLM, et al., Phys.Rev.C 107 (2023), 044908

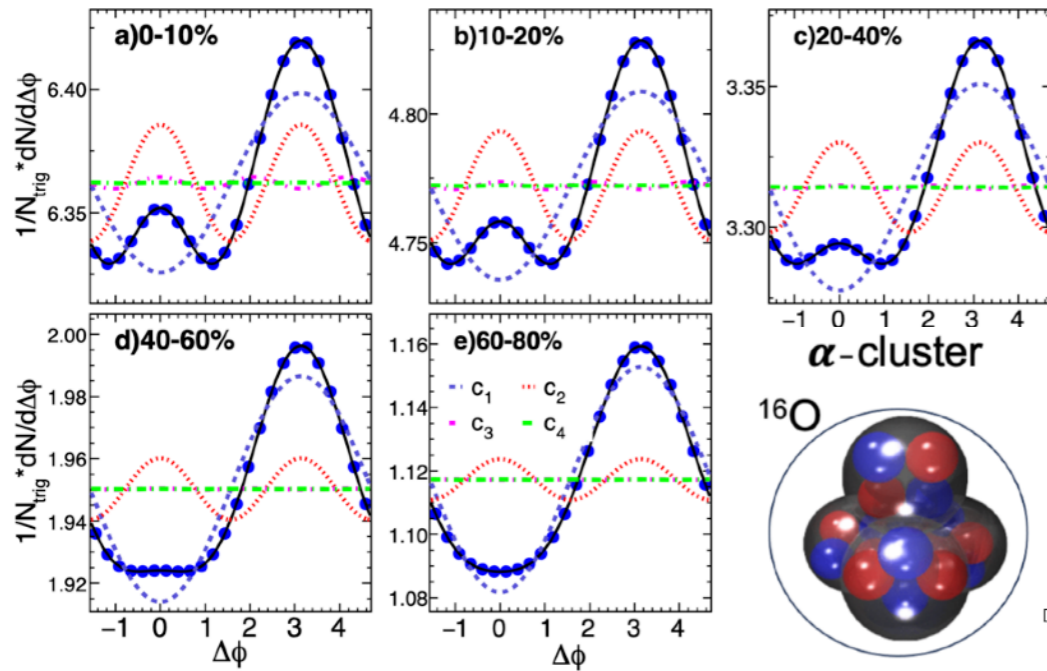
thermal photon



•  $\epsilon_3/\epsilon_2$  &  $v_3/v_2$  is sensitive to  $\alpha$ -cluster structure in C+Au collisions at 200GeV

# STAR results for O+O collisions

Shengli Huang [for STAR], arXiv:2312.12167

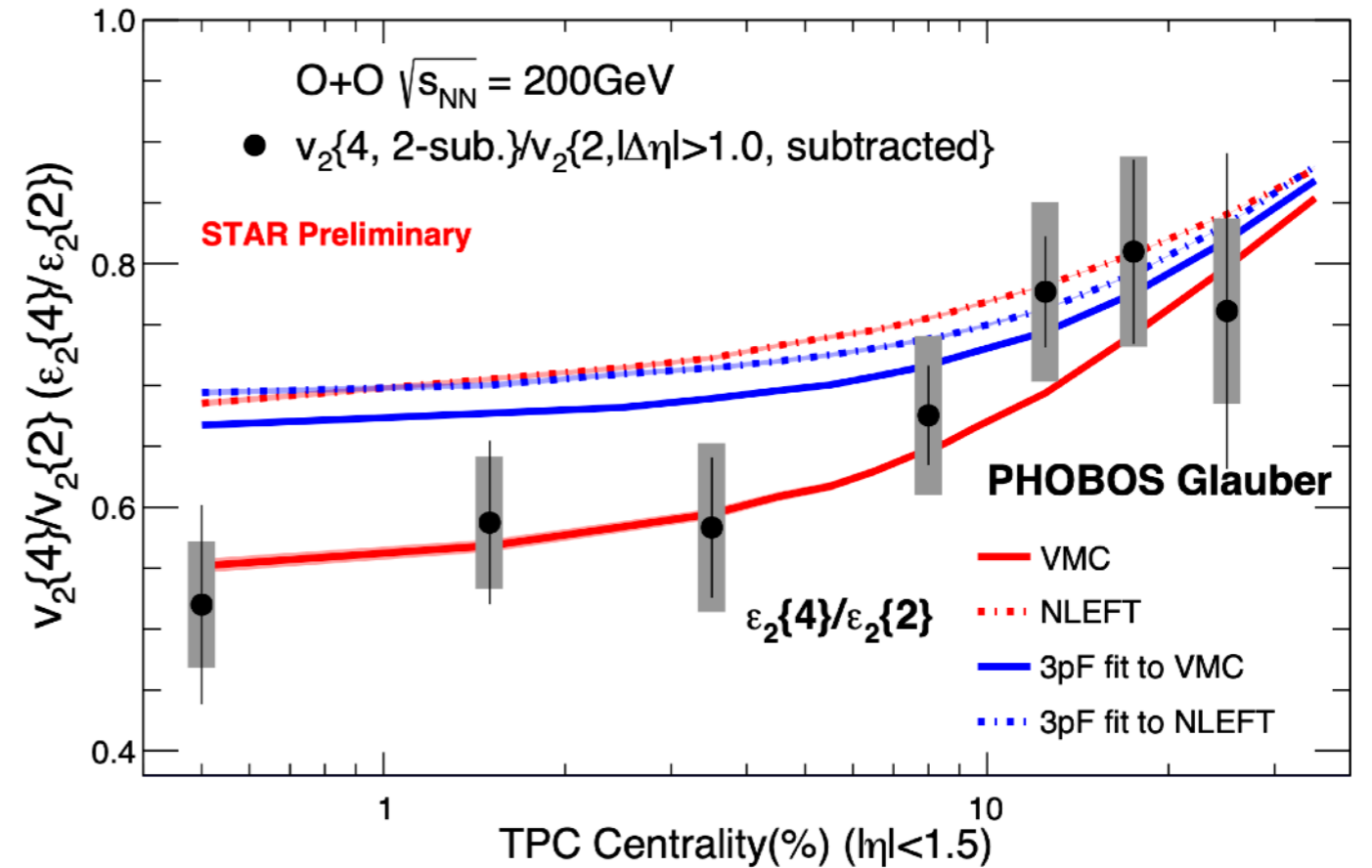
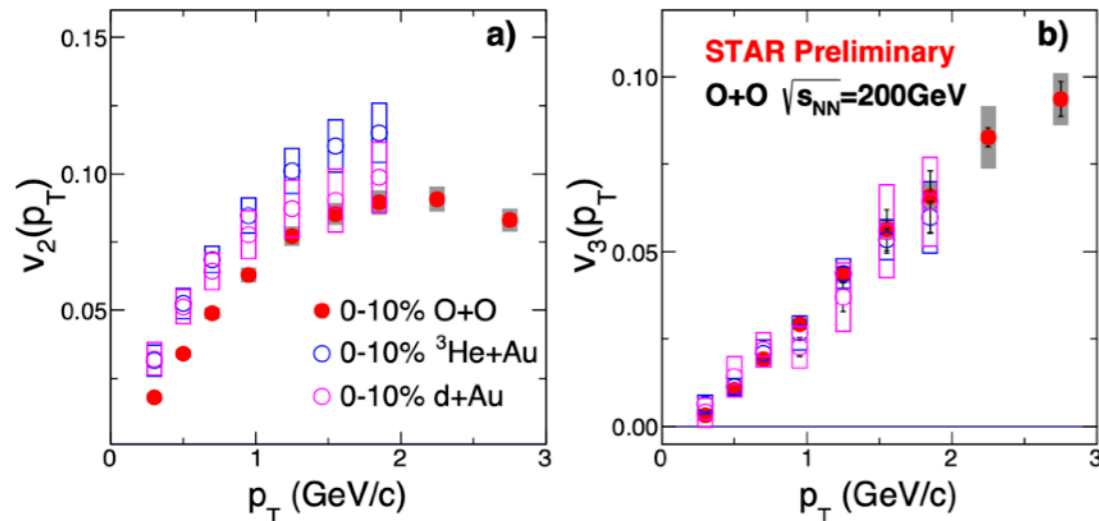


$$Y(\Delta\phi, p_T^{\text{trig}}) = c_0 \left( 1 + 2 \sum_{n=1}^{n=4} c_n \cos(n\Delta\phi) \right).$$

$$f = c_1^{\text{cent}} / c_1^{\text{peri}}$$

$$c_n^{\text{sub}} = c_n - c_n^{\text{non-flow}} = c_n^{\text{cent}} - c_n^{\text{peri}} \times f$$

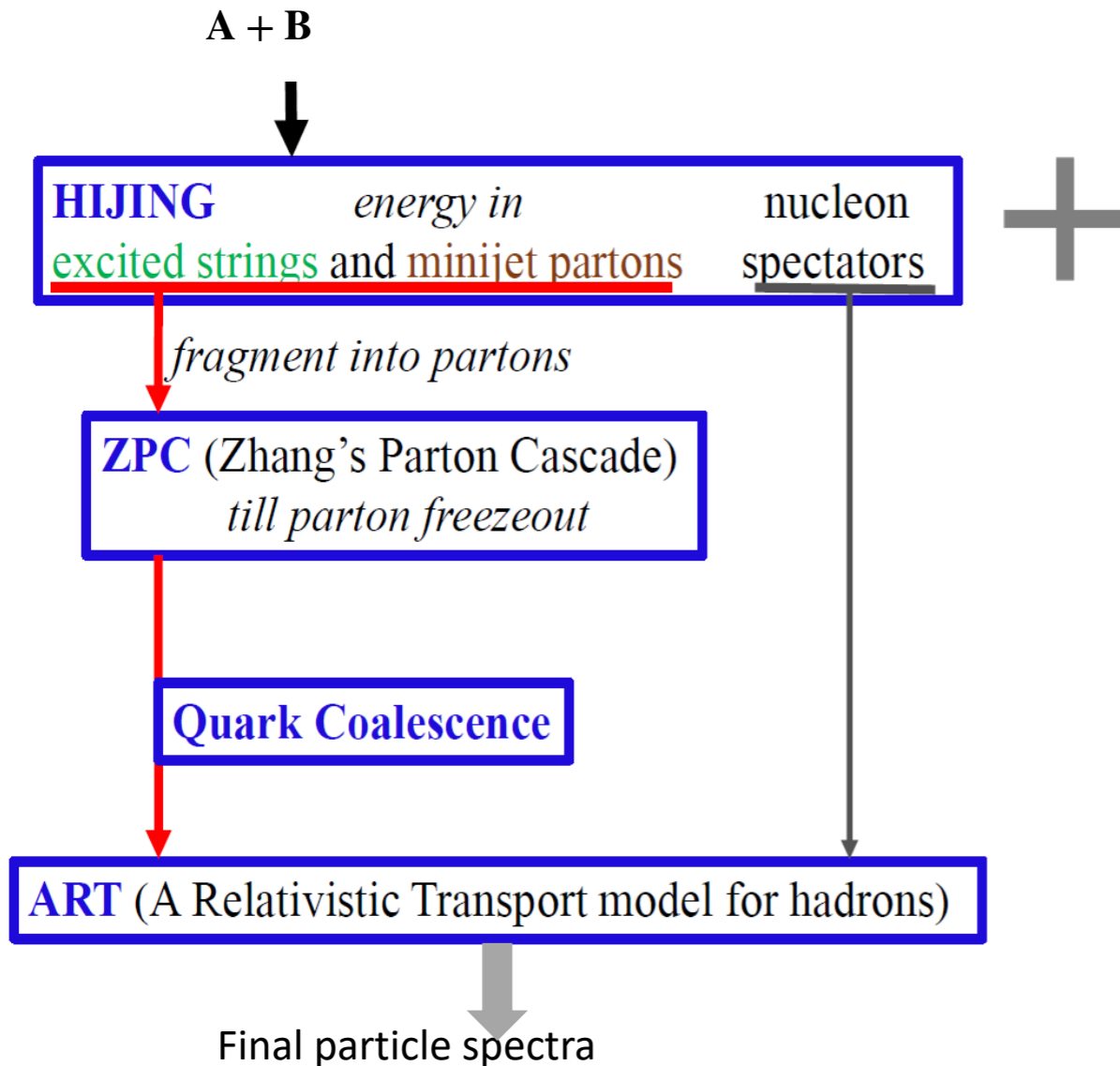
$$c_n = v_n^{\text{trig}} \times v_n^{\text{assoc}}$$



- $v_2\{4\}/v_2\{2\}$  serves as a powerful tool for studying nucleon-nucleon correlations in O+O collisions at 200 GeV.
- A hint on possible  $\alpha$ -cluster structure of oxygen nuclei?

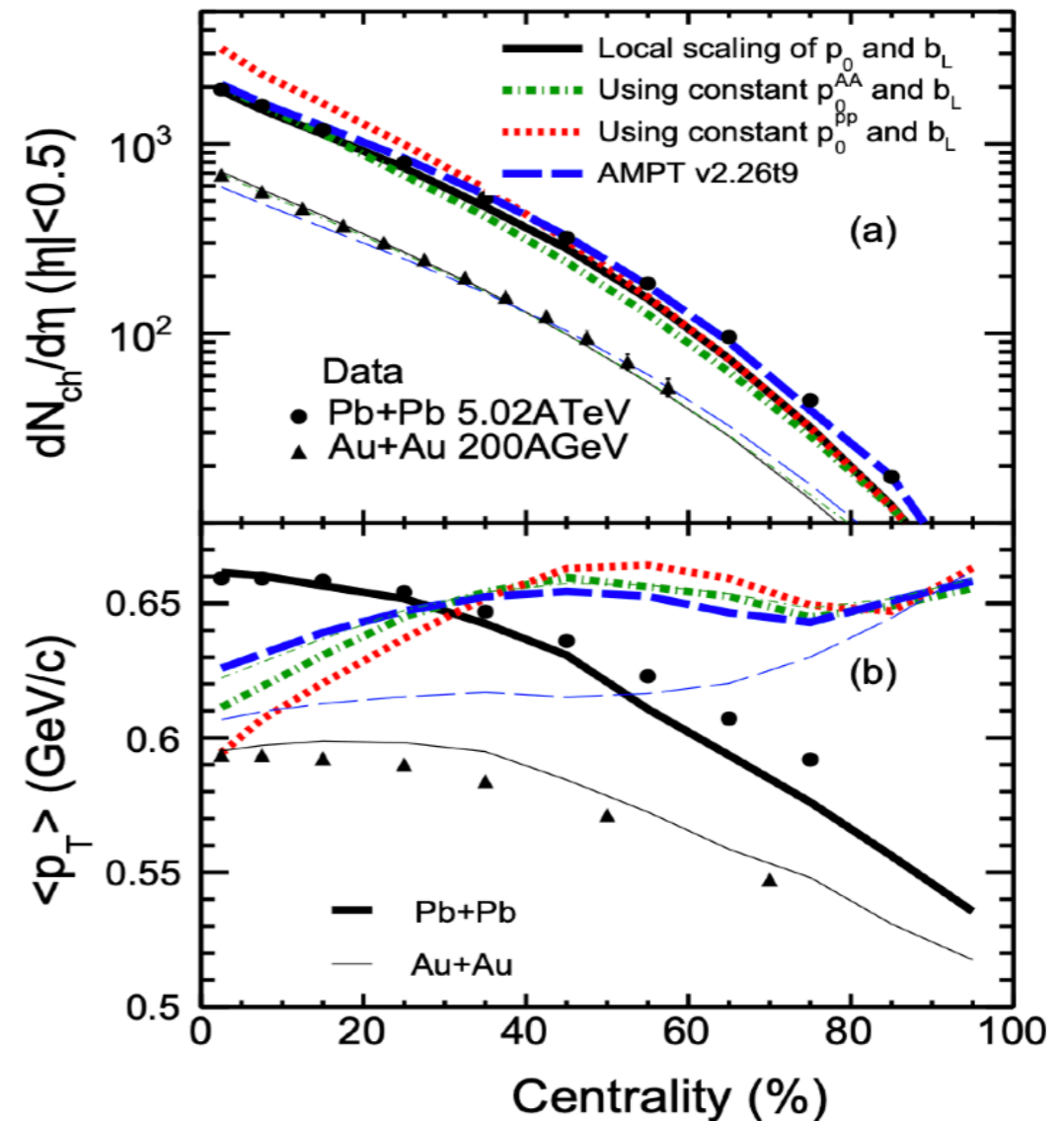
# Improved Version of String-Melting AMPT model

Z.W. Lin, C.M. Ko, B.A. Li, B. Zhang, S. Pal, PRC **72**, 064901 (2005)



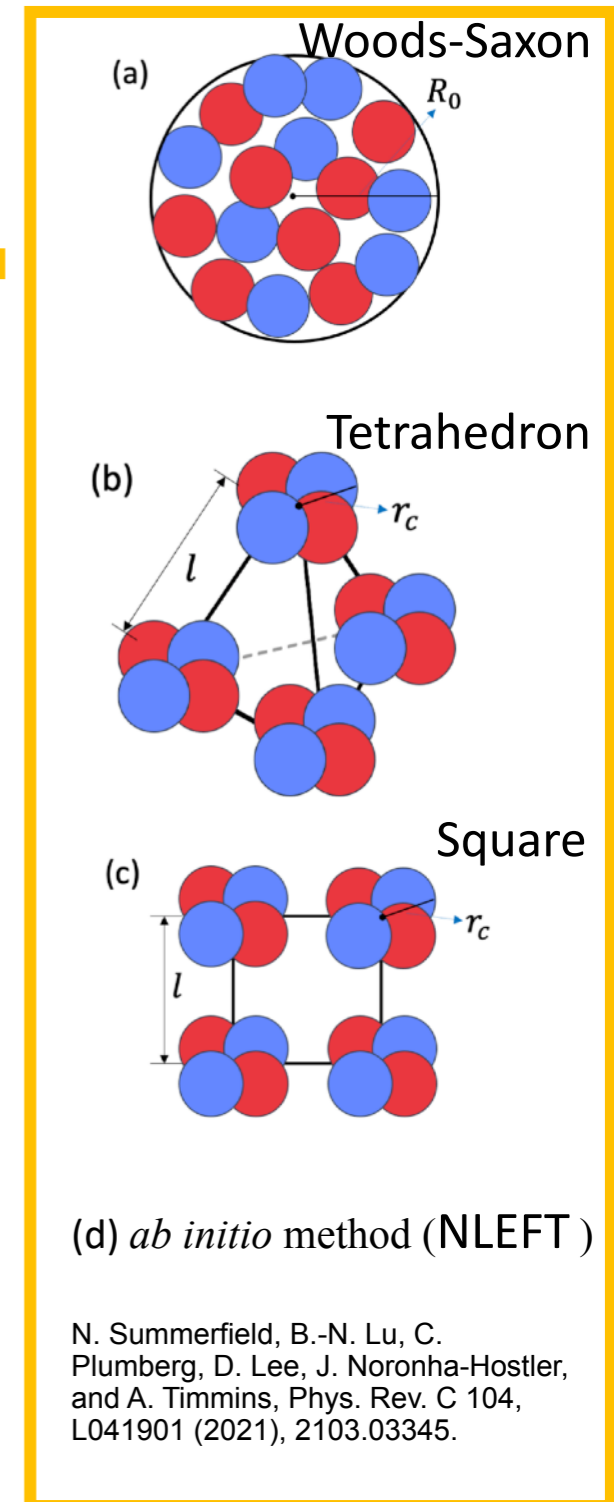
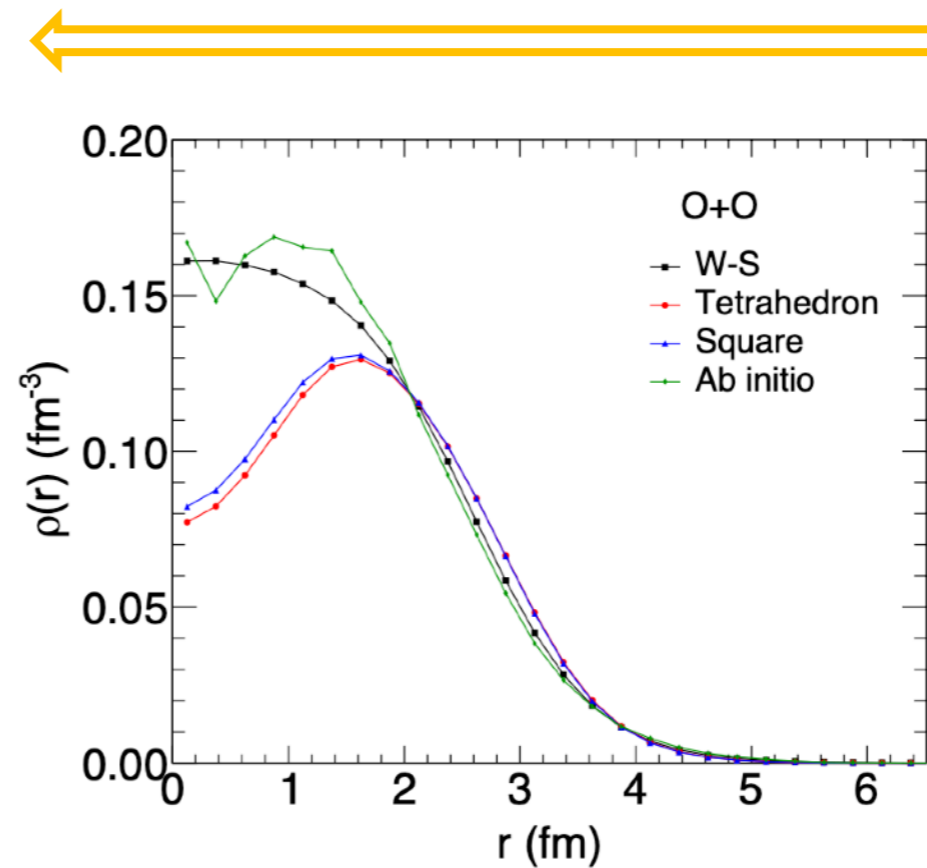
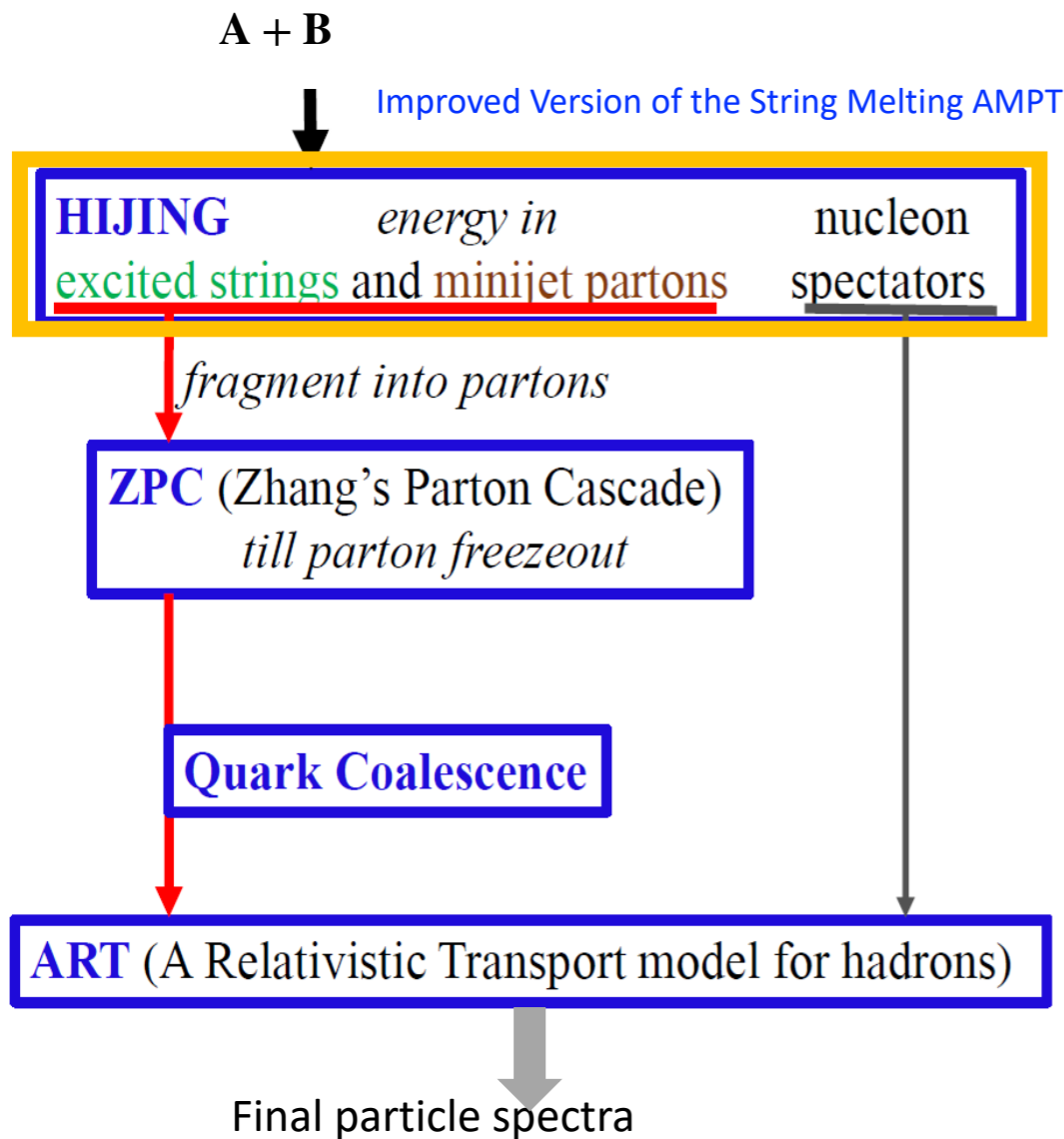
1. New quark coalescence model.
2. Improved heavy quark productions.
3. Modern set of parton distribution functions in proton and impact parameter-dependent nuclear shadowing.

➤ Improved AMPT model correctly describes the centrality dependences of  $\langle N_{ch} \rangle$  and  $\langle p_T \rangle$  rather well.



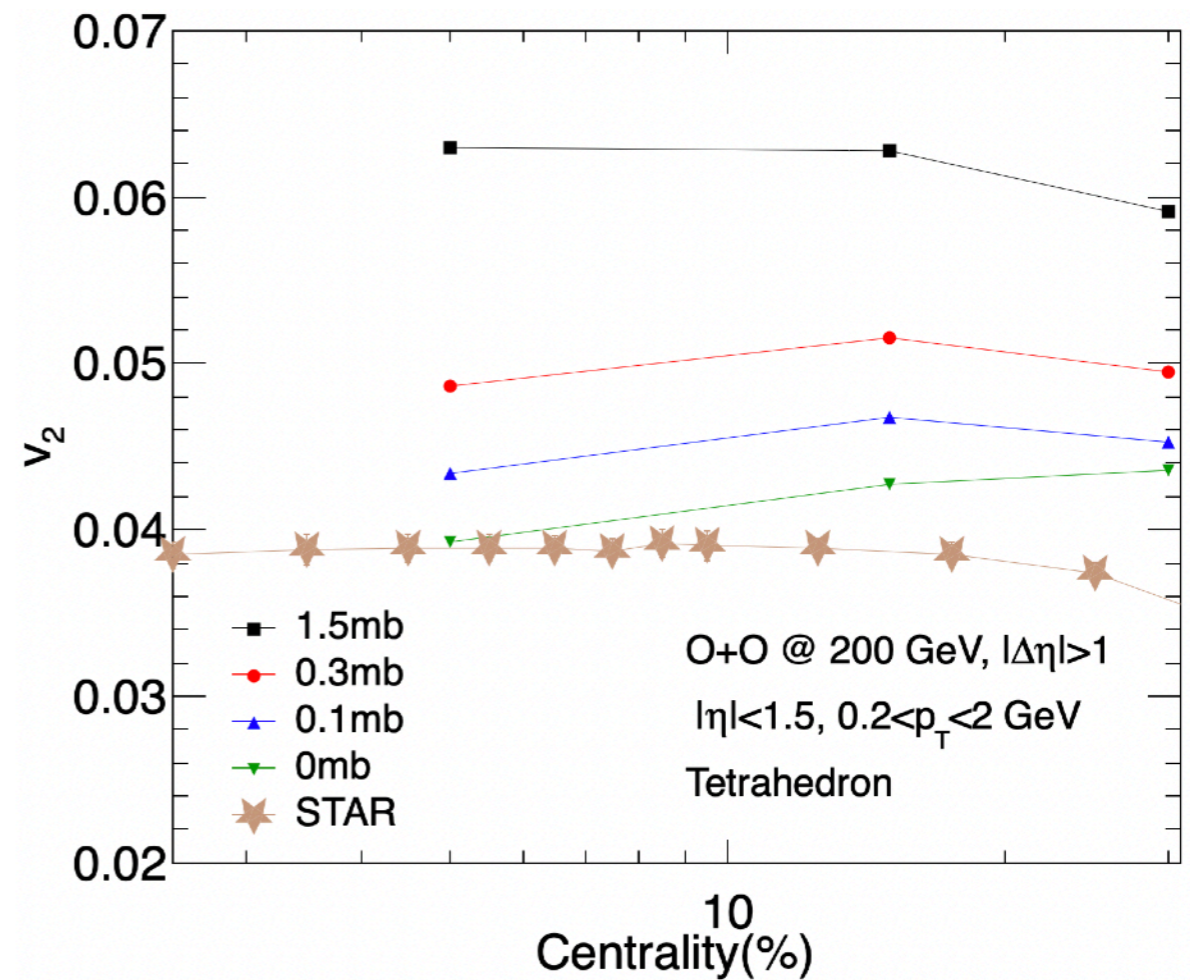
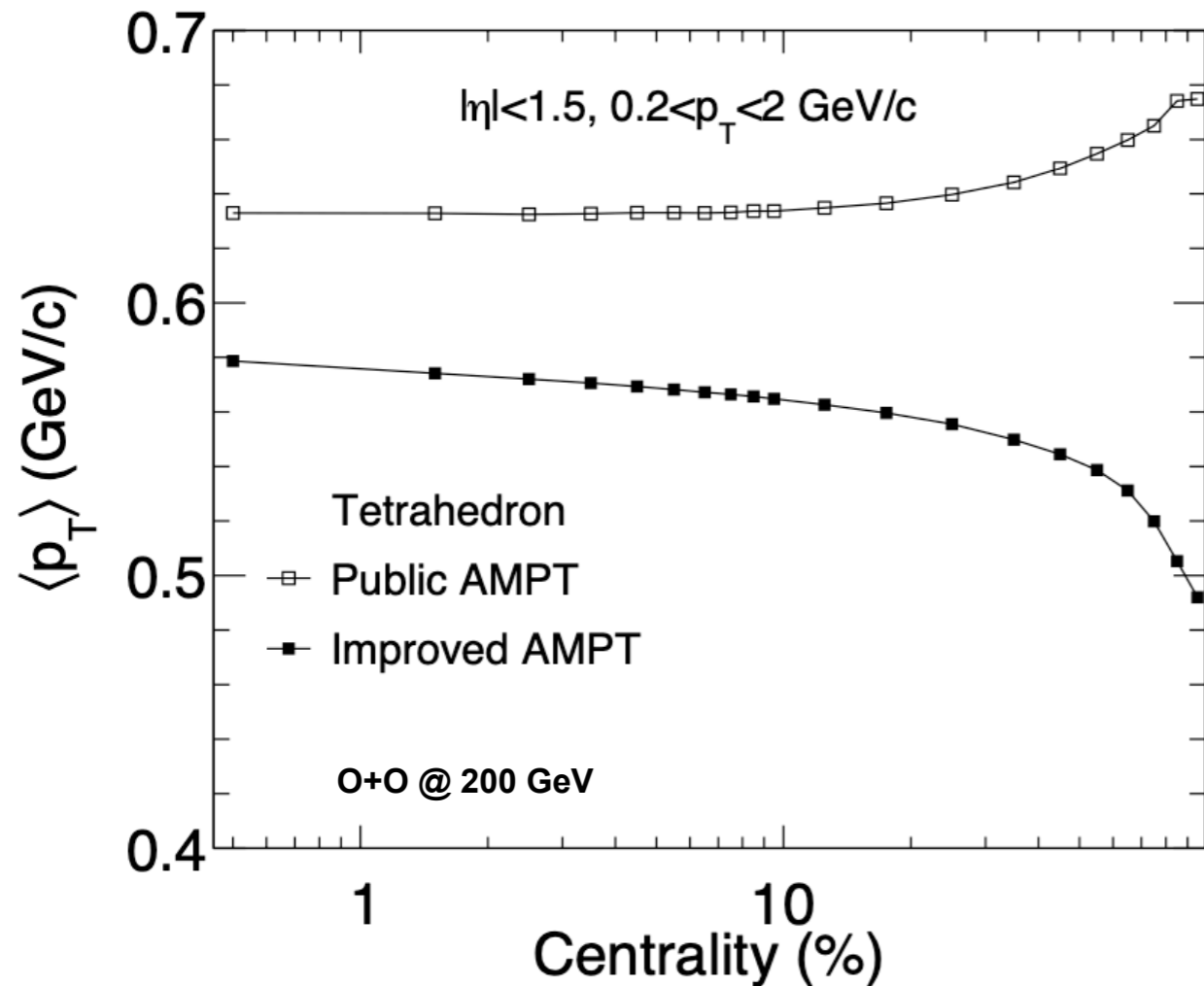
Chao Zhang, Liang Zheng, Shusu Shi, Zi-Wei Lin, PRC **104**, 014908 (2021)

# Introducing nuclear structures of $^{16}\text{O}$ to improved AMPT



Nuclear structure effect in O+O collisions  
at 200 GeV?

# Centrality dependence of $\langle p_T \rangle$ & $v_2$ in O+O collisions at 200 GeV



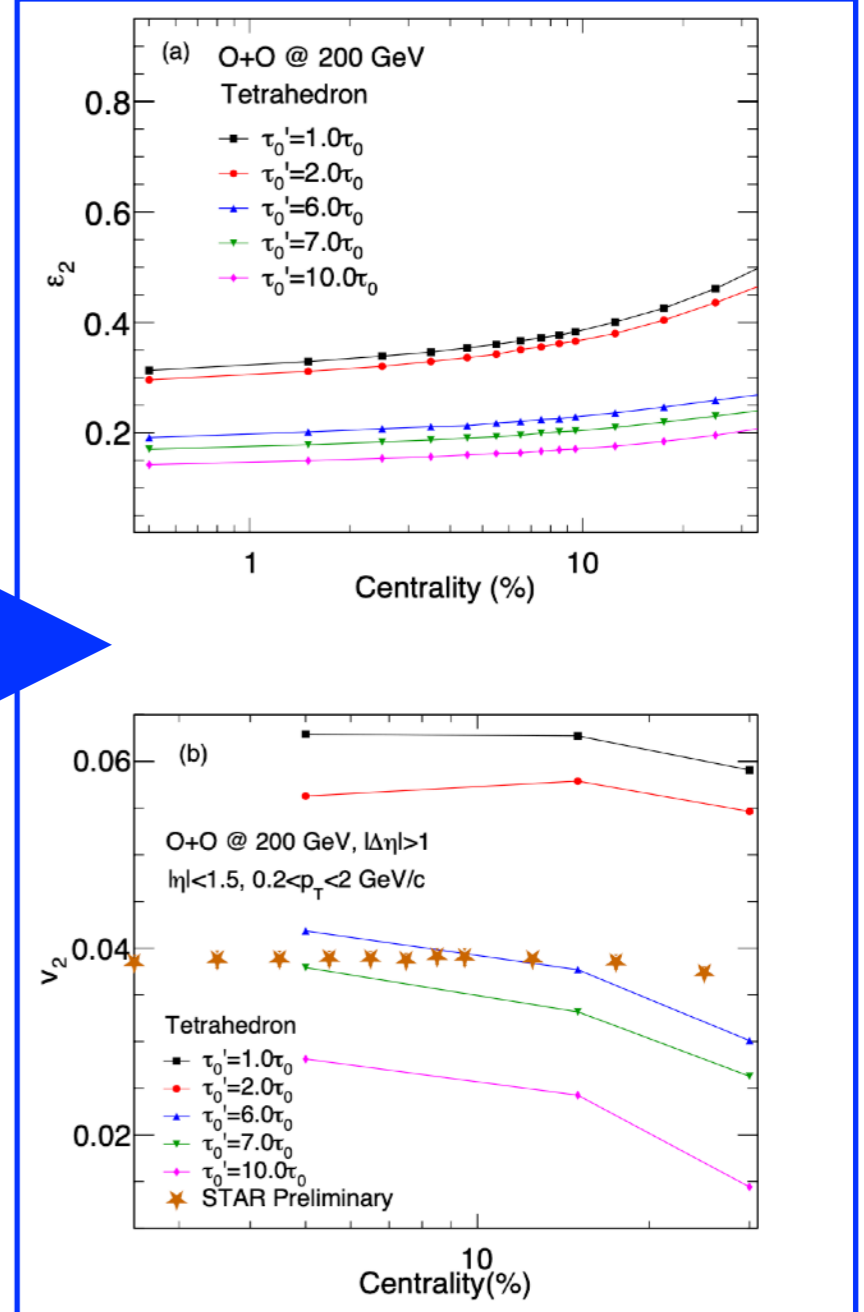
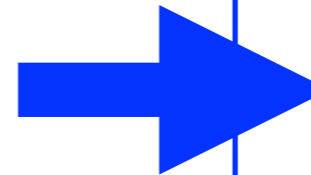
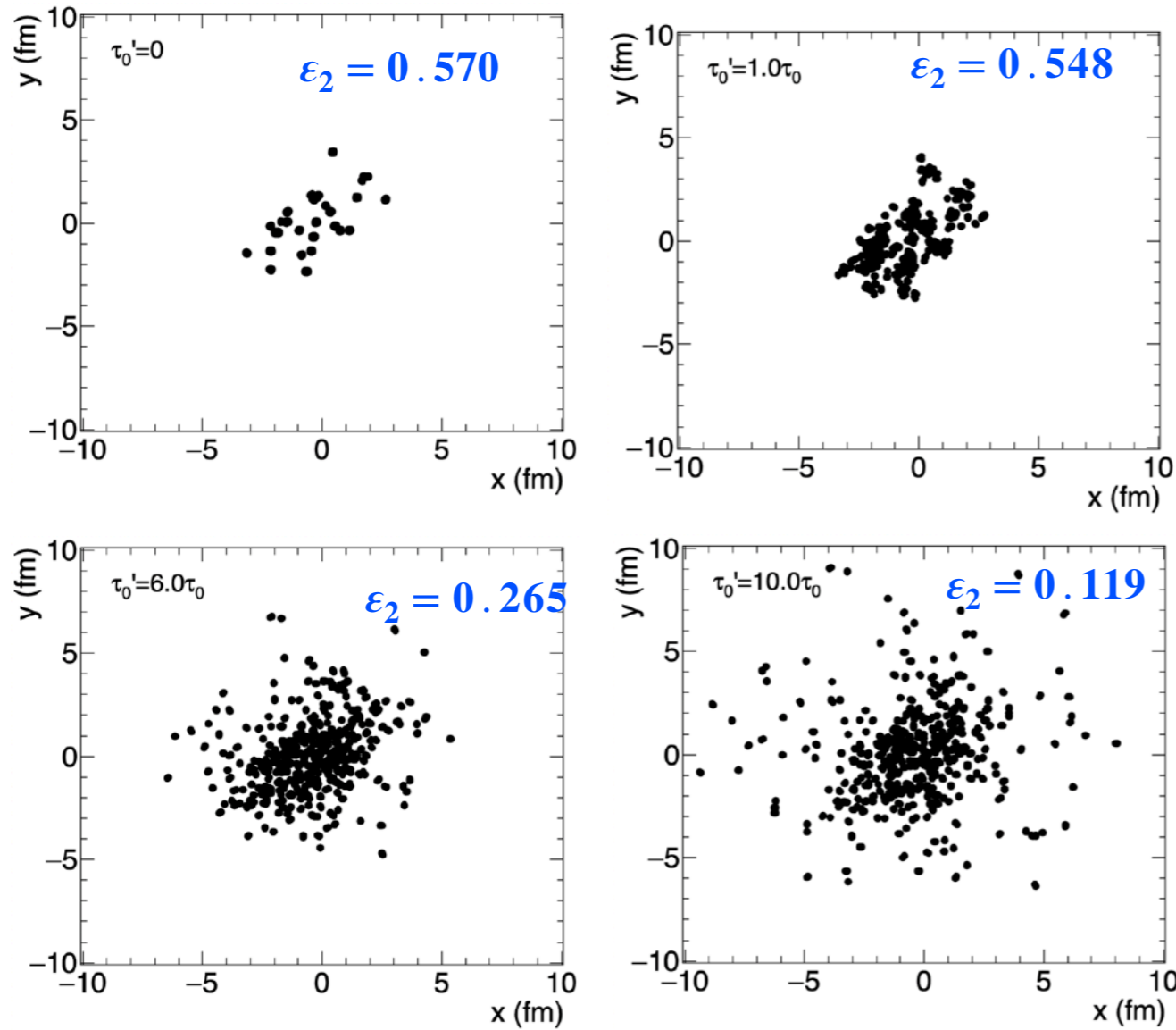
- The centrality dependence of  $\langle p_T \rangle$  is reasonable in improved AMPT version.
- Improved AMPT failed to reproduce data even when tuning off parton cascade
- Why the  $v_2$  problem? How to solve it?



# Parton formation time dependence of $\varepsilon_2$ & $v_2$

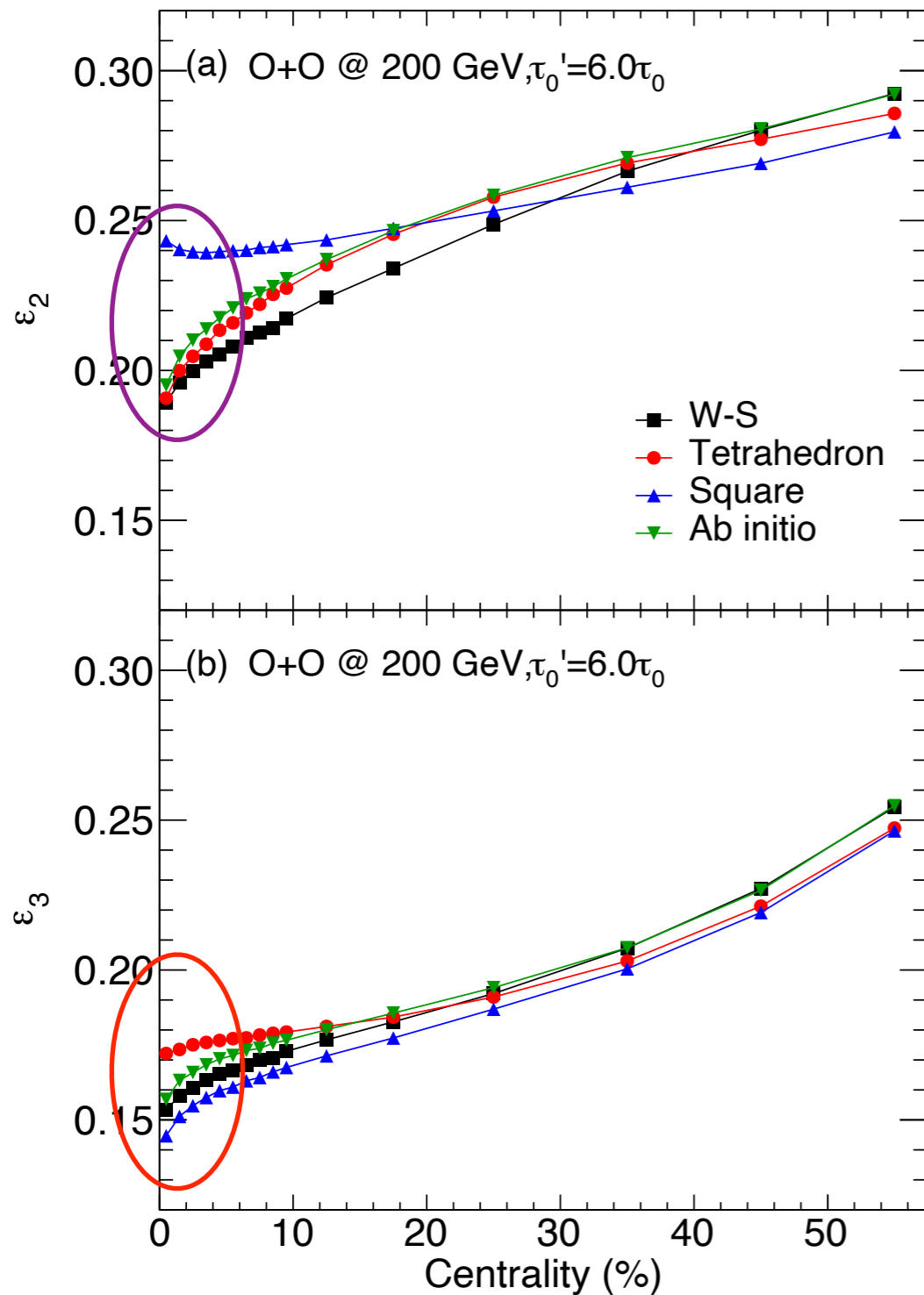
The formation time for each parton:

$$\tau'_0 = \text{const} \cdot E/m_T^2, \quad \tau_0 = E/m_T^2$$



- $\varepsilon_2$  &  $v_2$  decrease with increasing parton formation time,  $v_2$  at  $\tau'_0 = 6\tau_0$  is close to data. Why?
- Is spatial distribution diffuse and gaseous because condensed  $\alpha$ -particles are weakly bound?
- A similar concept/analog as the compactness from Huichao and Haojie et al.'s work (arXiv: 2401.15723).

# Nuclear structure effect on centrality dependence of $\epsilon_2$ and $\epsilon_3$



➤  $\epsilon_2$  and  $\epsilon_3$  are more sensitive to nuclear structure for more central collisions

➤ for central collisions:

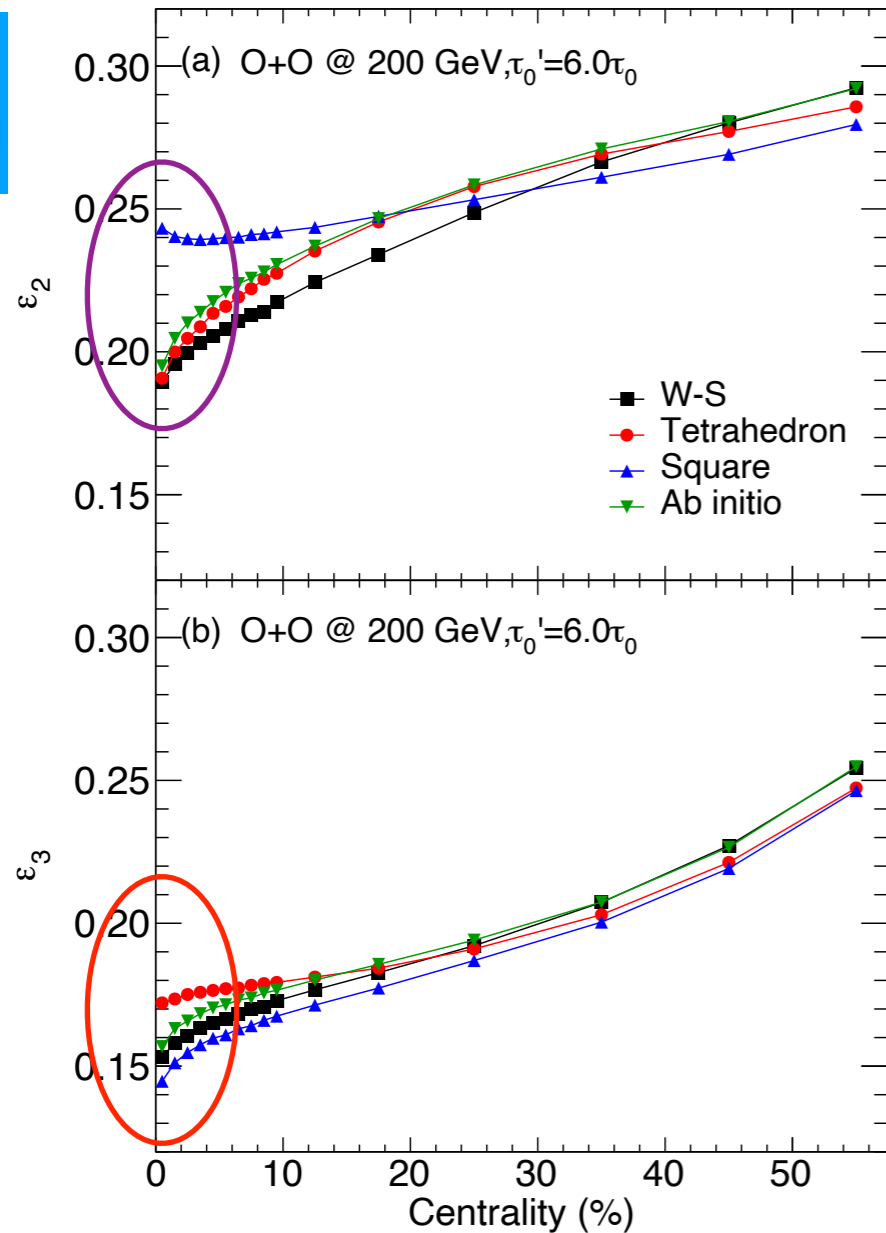
➤  $\epsilon_2$ : Square > Ab initio  $\approx$  Tetrahedron > W-S

➤  $\epsilon_3$ : Tetrahedron > Ab initio > W-S > Square

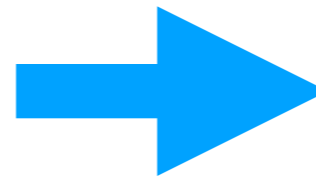
➔  $v_2$  and  $v_3$ ?

# Consistence between $v_n$ and $\epsilon_n$ in O+O collisions at 200 GeV

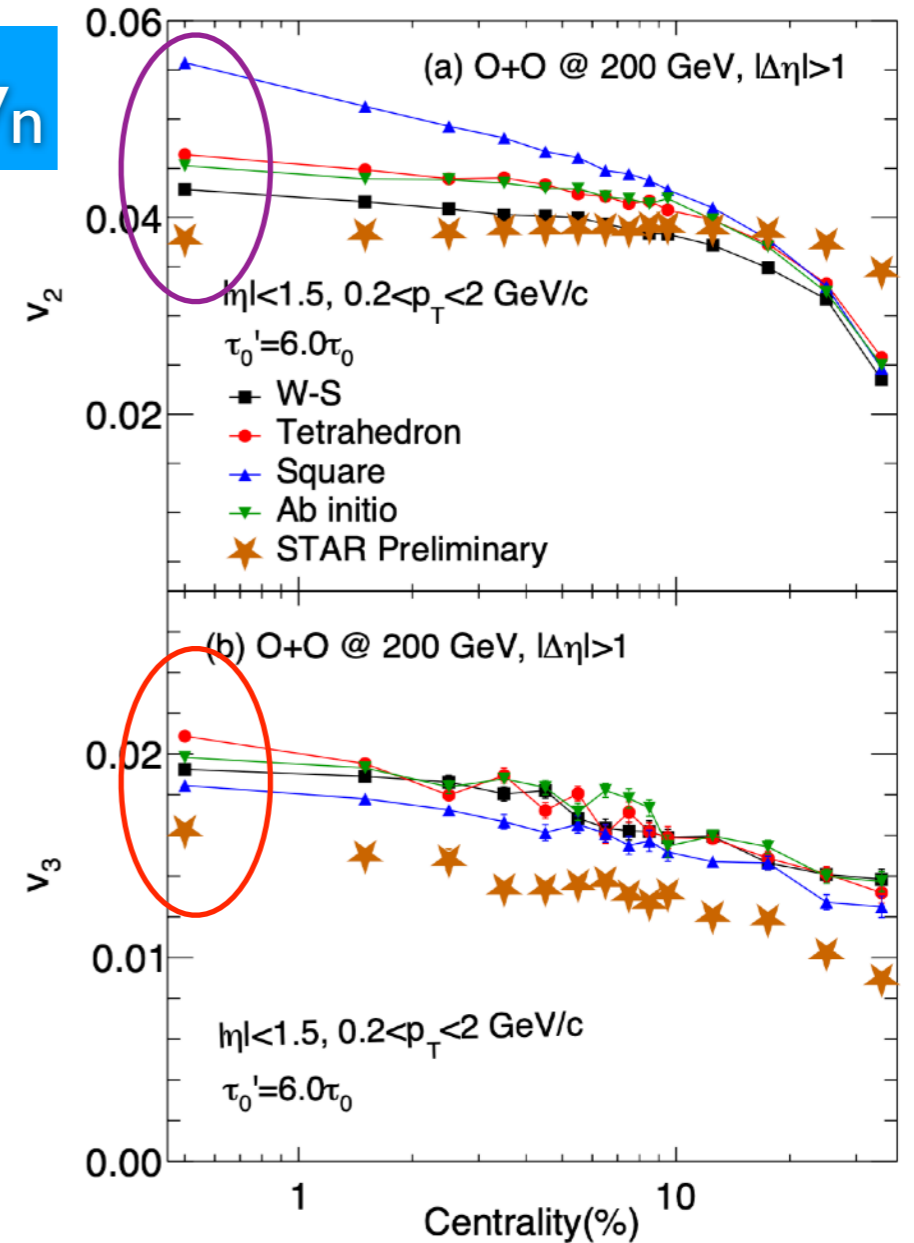
$\epsilon_n$



- $\epsilon_2$  and  $\epsilon_3$  are more sensitive to nuclear structure for more central collisions
- for central collisions:
  - $\epsilon_2$ : Square > Ab initio  $\approx$  Tetrahedron > W-S
  - $\epsilon_3$ : Tetrahedron > Ab initio > W-S > Square

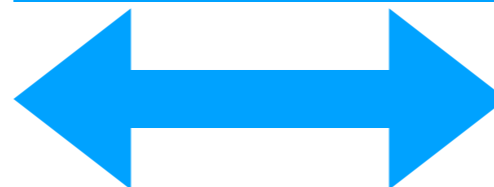


$v_n$

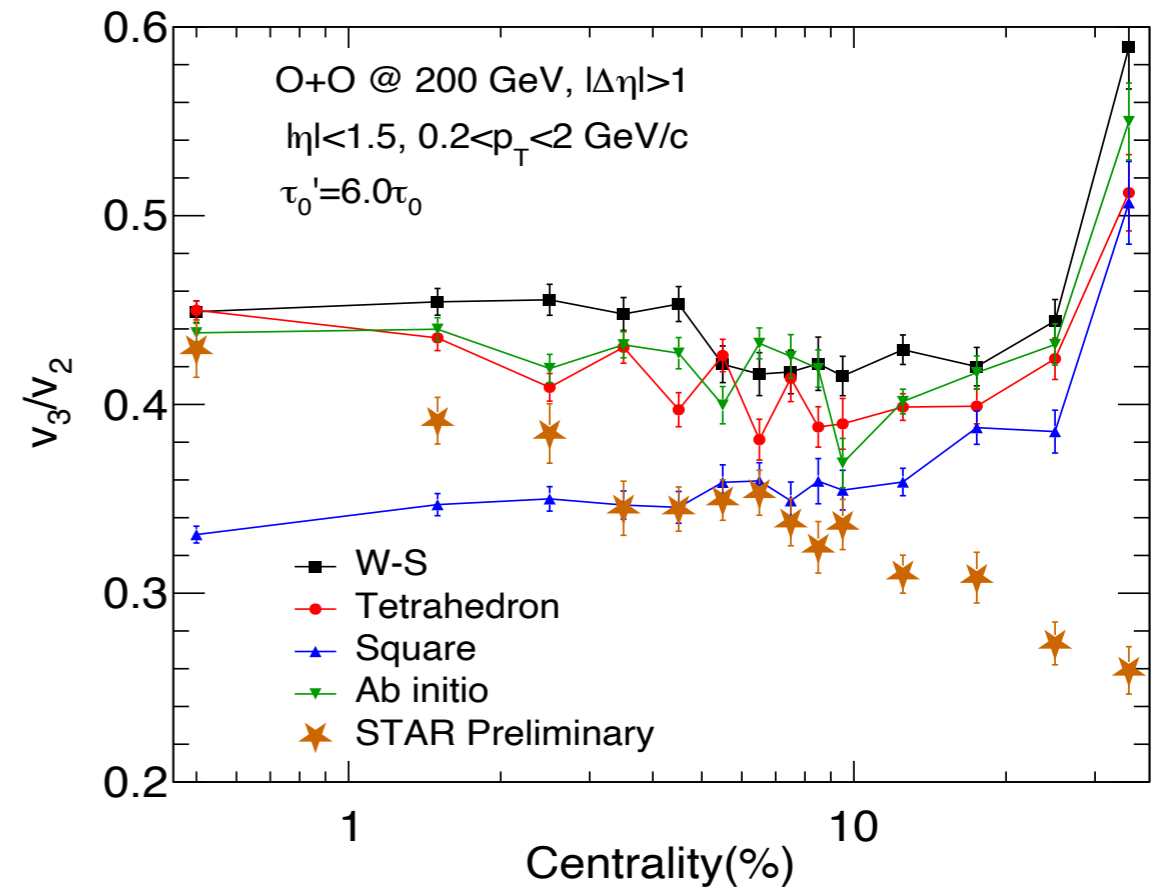
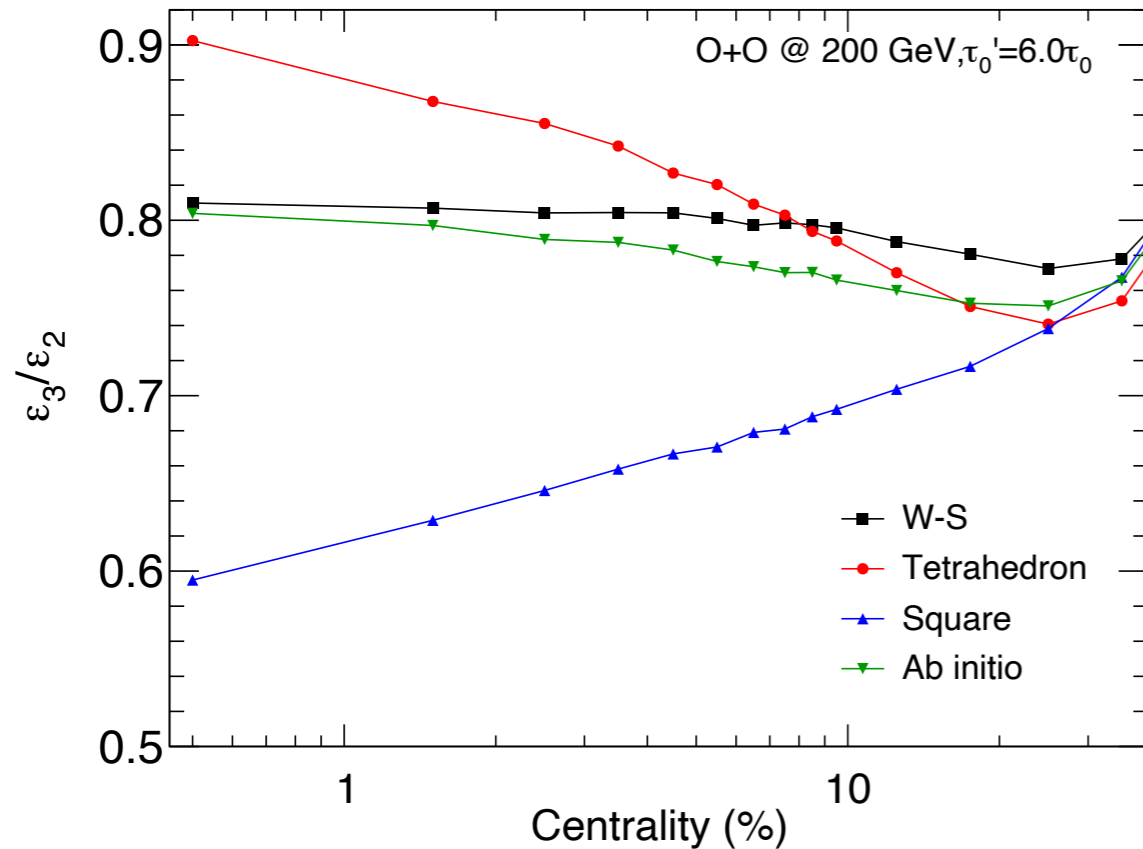


- $v_2$  and  $v_3$  are more sensitive to nuclear structure for more central collisions
- for central collisions:
  - $v_2$ : Square > Ab initio  $\approx$  Tetrahedron > W-S
  - $v_3$ : Tetrahedron > Ab initio > W-S > Square

Consistent

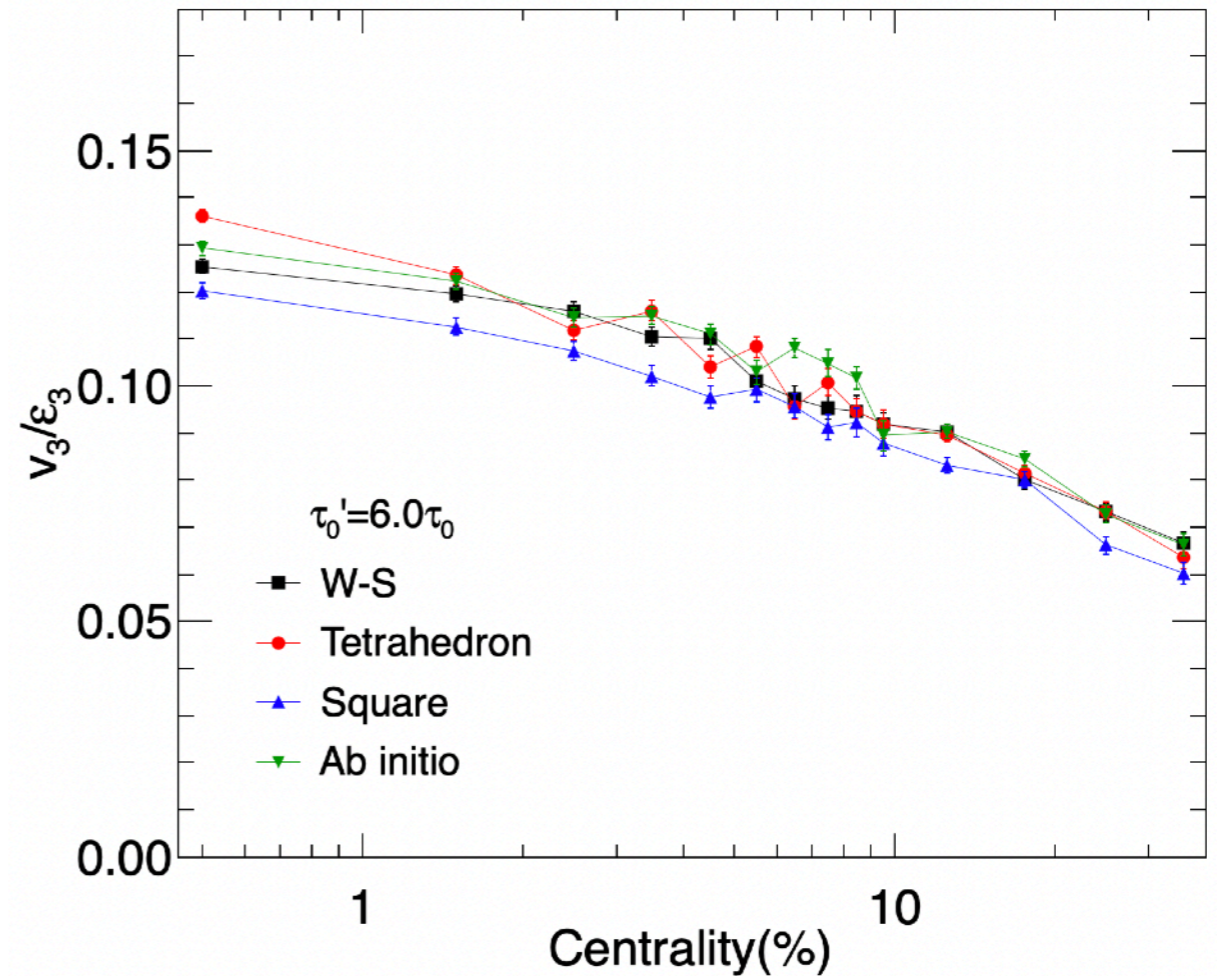
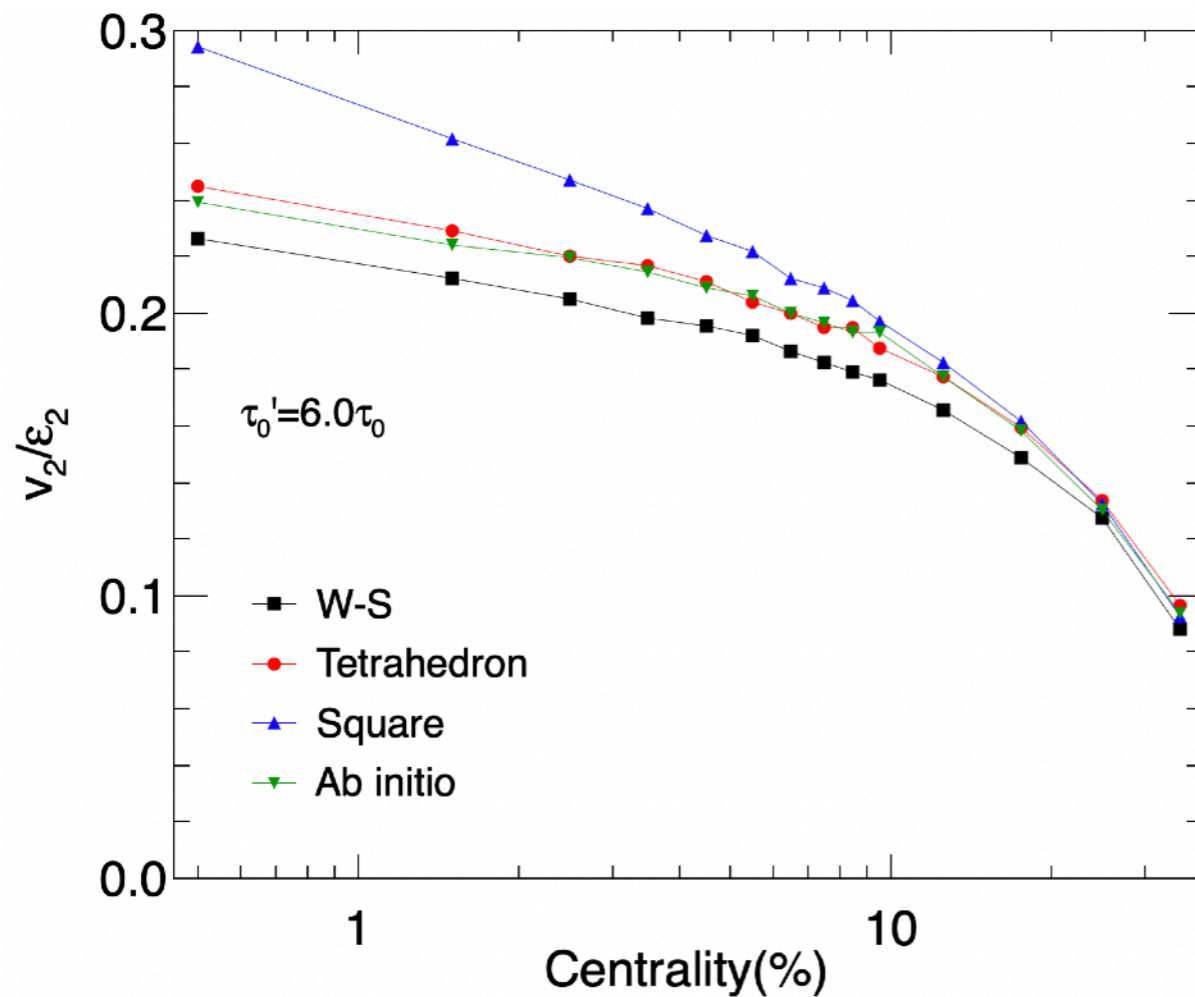


# $v_3/v_2$ ratio in O+O collisions at 200 GeV



- $\epsilon_3/\epsilon_2$  are more sensitive to nuclear structure for more central collisions
- $\epsilon_3/\epsilon_2$  for central collisions: Tetrahedron  $>$  W-S  $\approx$  Ab initio  $>$  Square
- $v_3/v_2$  for central collisions: Tetrahedron  $\approx$  W-S  $\approx$  Ab initio  $>$  Square.
- Why is  $v_3/v_2$  inconsistent with  $\epsilon_3/\epsilon_2$ ?

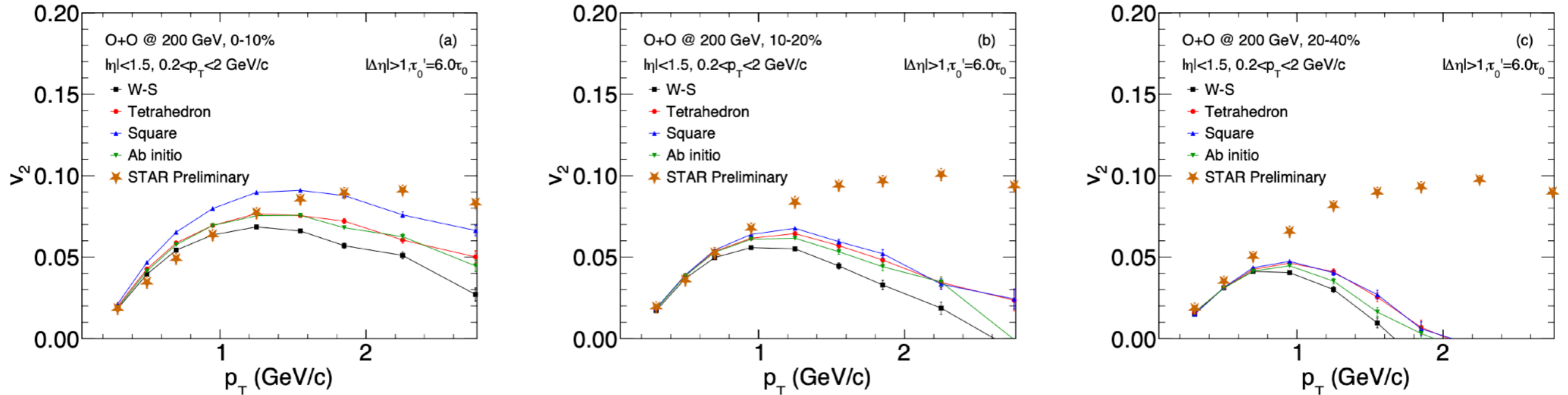
# Conversion efficiency ( $v_n/\epsilon_n$ ) from $\epsilon_n$ to $v_n$



➤ Because  $v_2/\epsilon_2 > v_3/\epsilon_3$

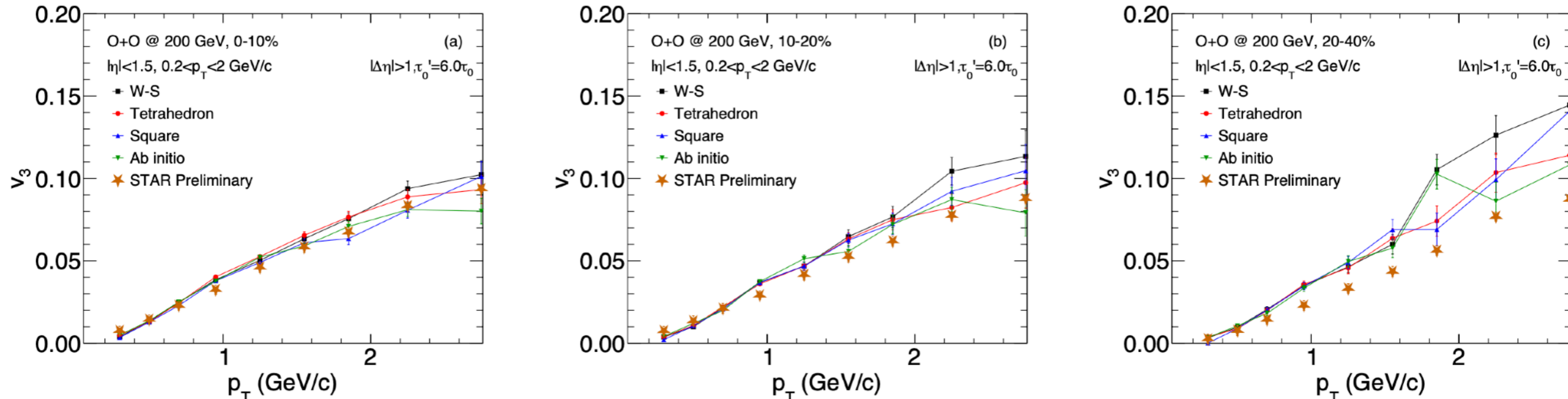
➤ The ordering of  $v_n/\epsilon_n$  in different configurations is consistent with the ordering of  $v_n$  or  $\epsilon_n$  in central collisions

# $p_T$ dependences of $v_2$ and $v_3$



➤  $v_2(p_T)$ : The AMPT results are consistent with the data for low  $p_T$ .

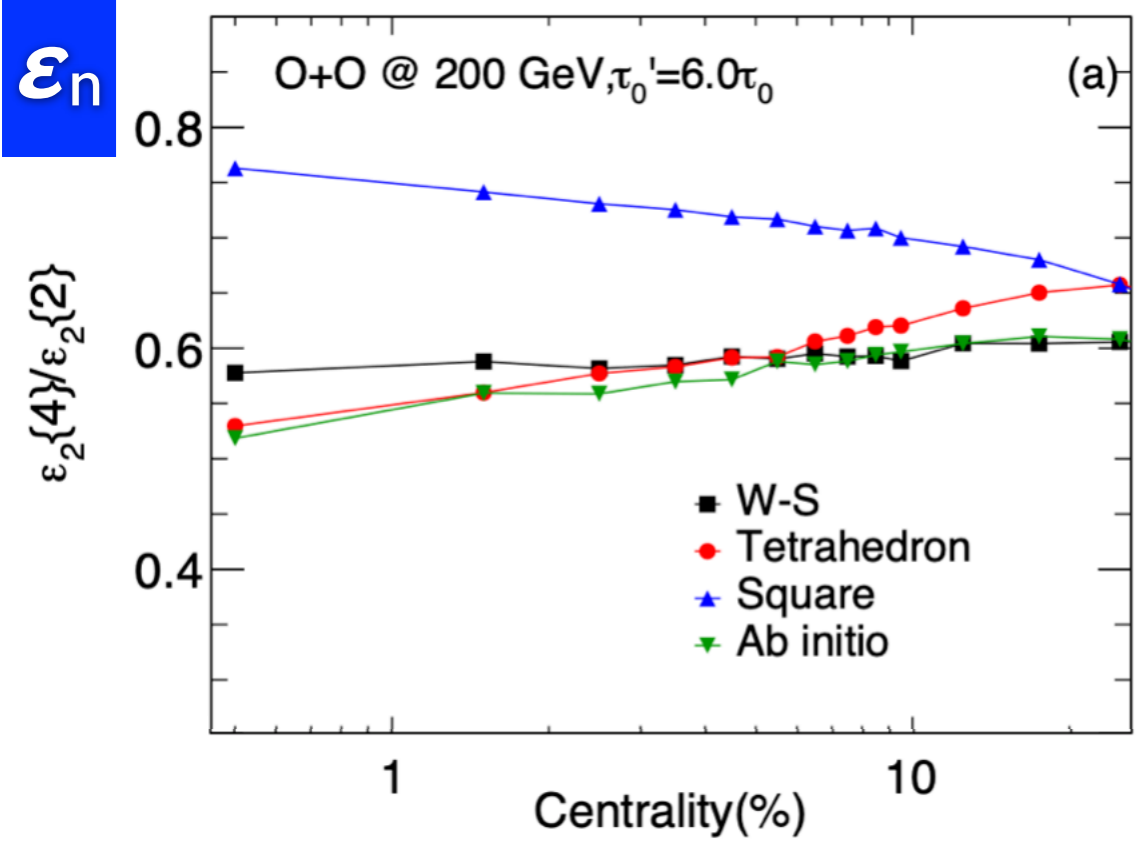
➤ Ours are lower than the data for high  $p_T$ . Is due to non-flow effect different from the data?



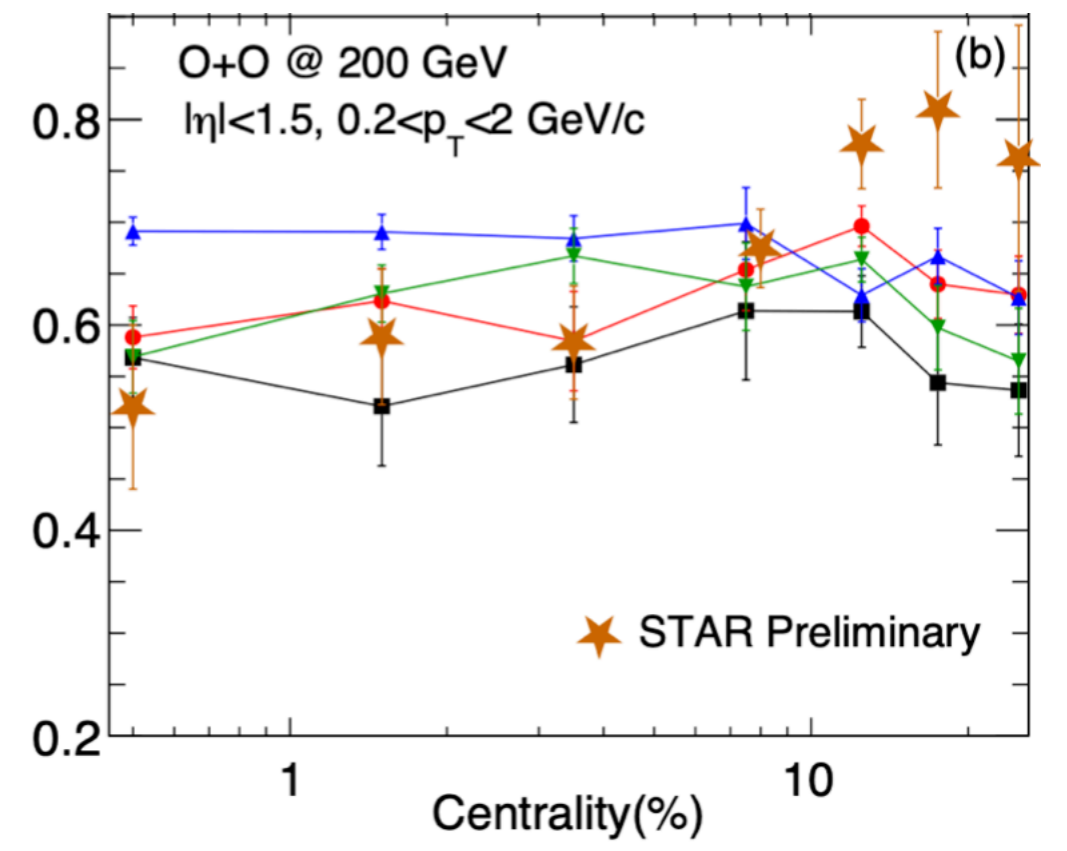
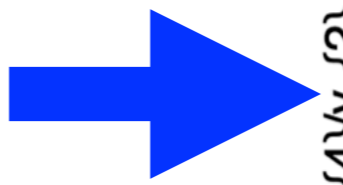
➤  $v_3(p_T)$ : The AMPT results are close to data and less sensitive to nuclear structure than  $v_2(p_T)$ .

# $\epsilon_2$ and $v_2$ fluctuations

$\epsilon_n$



$V_n$



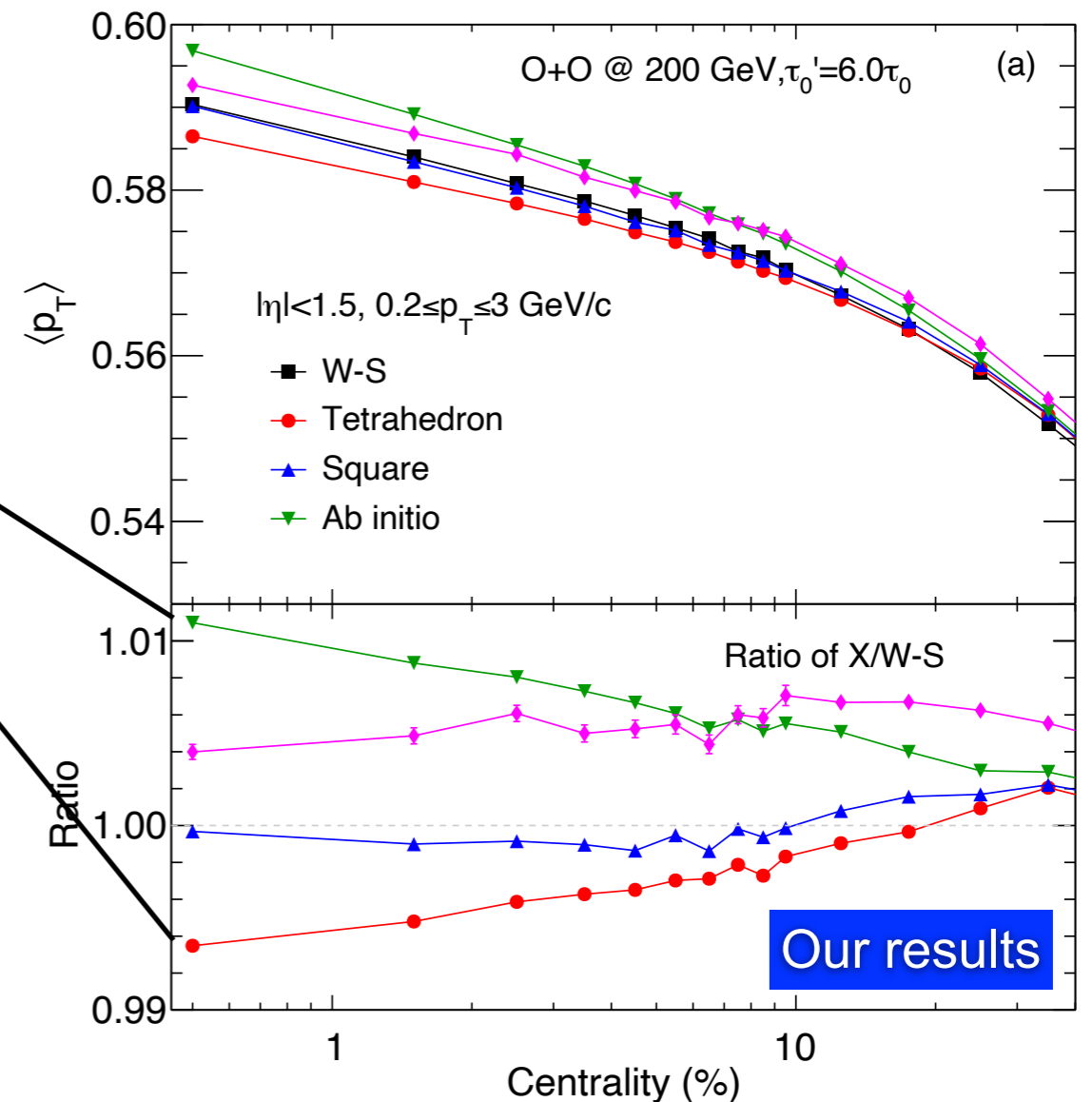
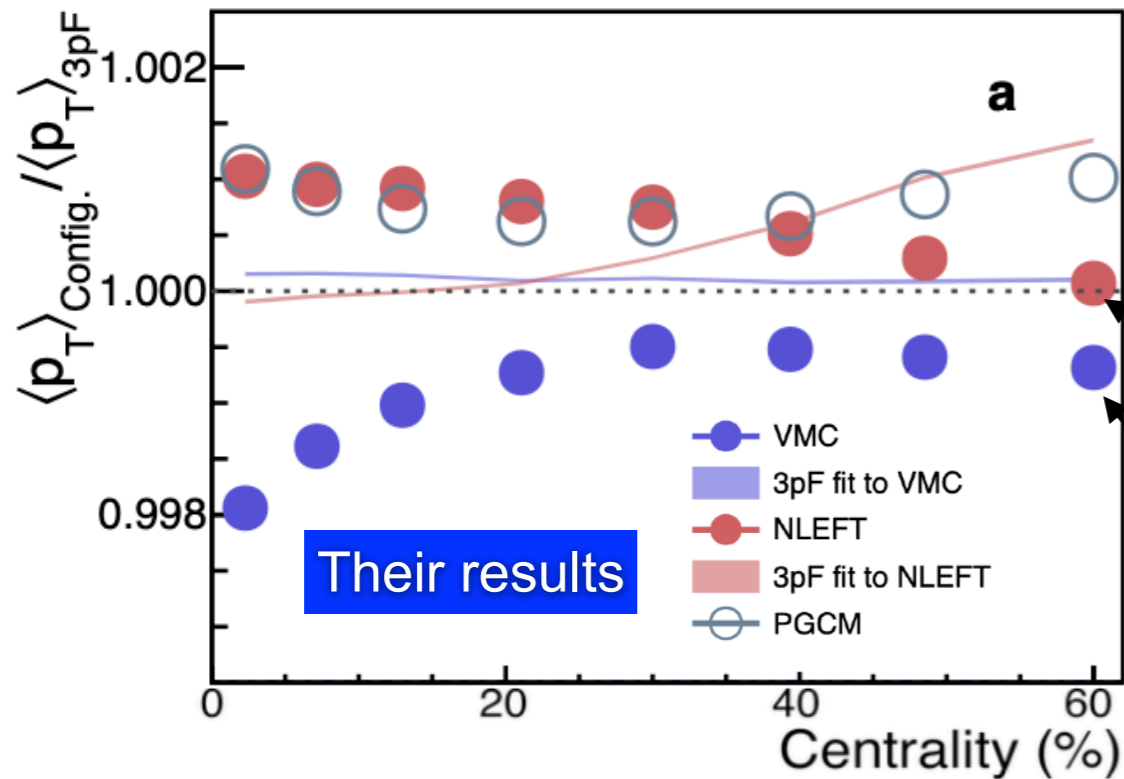
$$\epsilon_2\{2\}^2 = \langle \epsilon_2^2 \rangle = \langle \epsilon_2 \rangle^2 + \sigma_{\epsilon_2}^2,$$

$$\epsilon_2\{4\}^2 = (-\langle \epsilon_2^4 \rangle + 2\langle \epsilon_2^2 \rangle^2)^{1/2} \approx \langle \epsilon_2 \rangle^2 - \sigma_{\epsilon_2}^2$$

- $\epsilon_2$  and  $v_2$  fluctuations are consistent with each other.
- Larger  $v_2$  fluctuations are observed for Tetrahedron and Ab initio cases, which seems more consistent with the STAR data.

# $\langle p_T \rangle$ Ratio between with and without nuclear structure

Chunjian Zhang, Jinhui Chen, Giuliano Giacalone, Shengli Huang, Jiangyong Jia and Yu-Gang Ma, arXiv:2404.08385

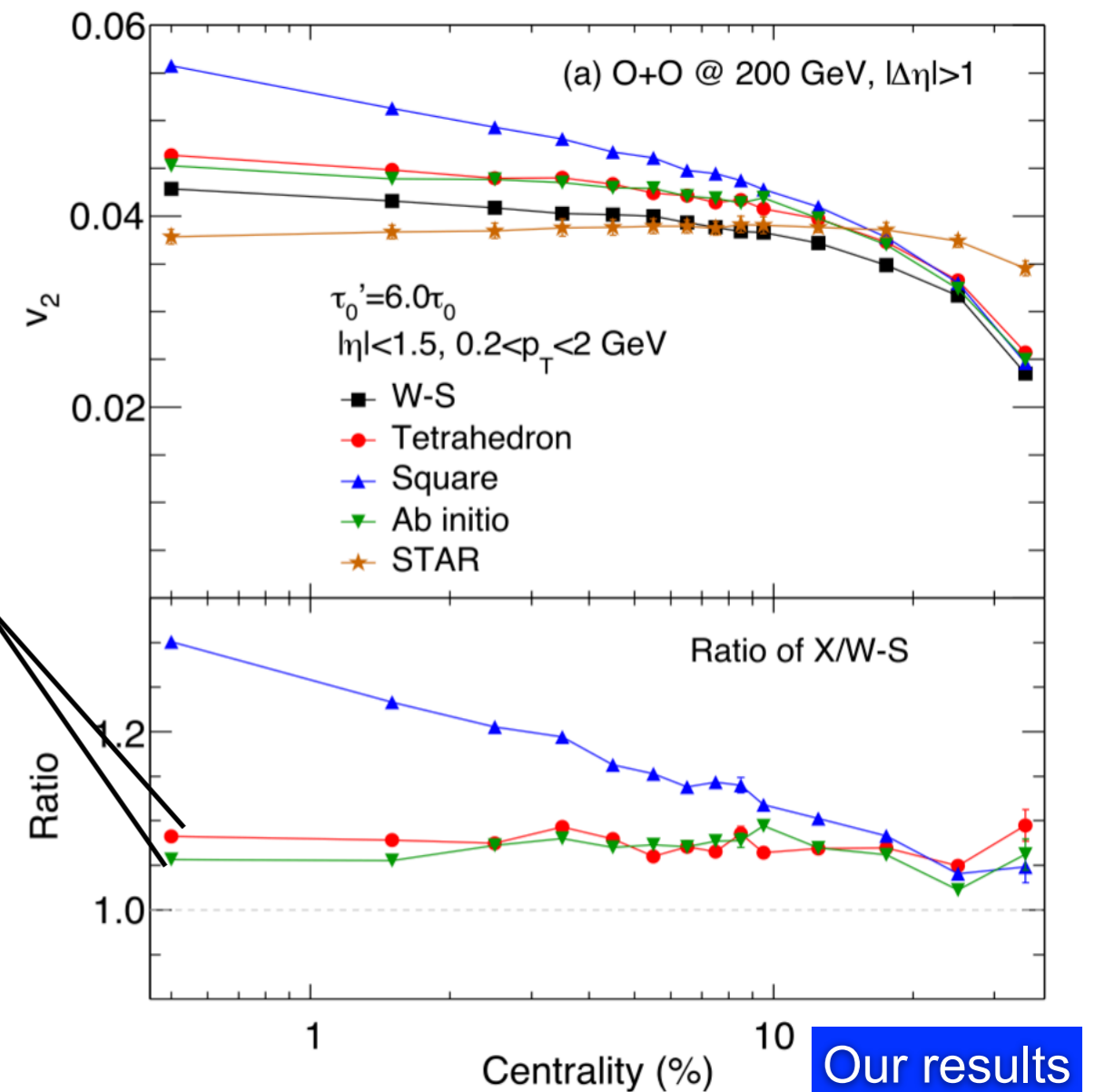
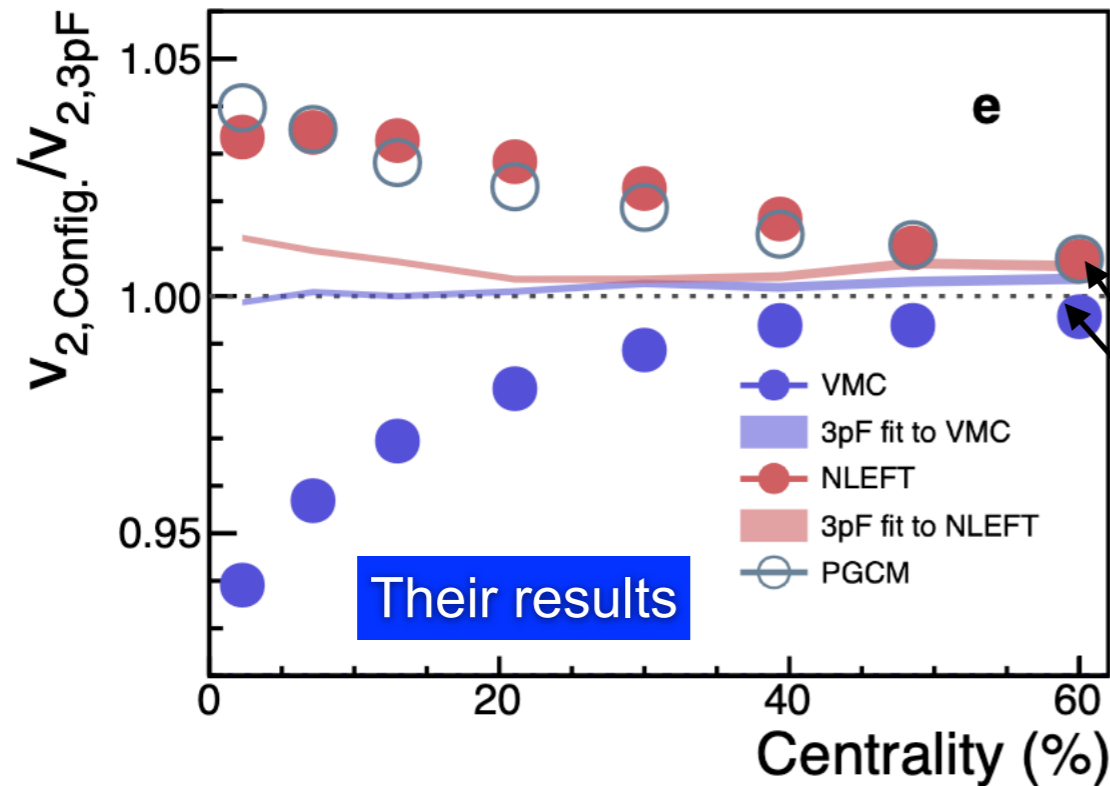


- Our  $\langle p_T \rangle$  ratio for Ab initio (NLEFT) is consistent with their NLEFT.
- Our  $\langle p_T \rangle$  ratio for Tetrahedron is consistent with their VMC.



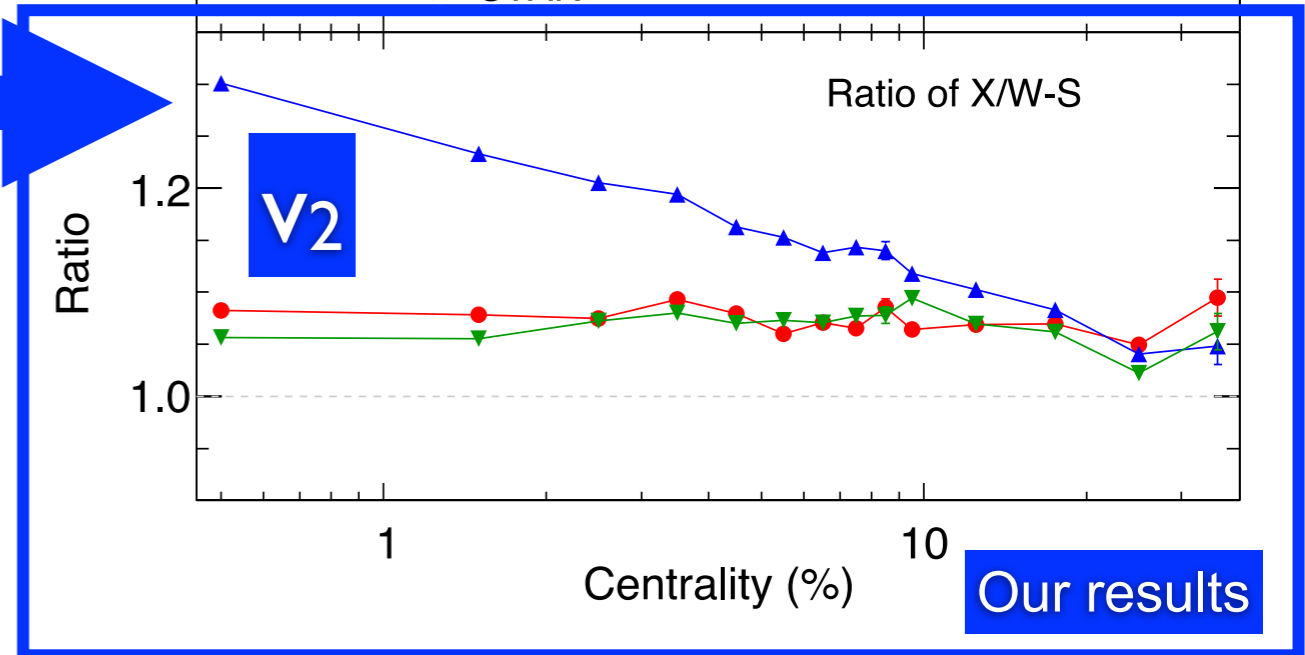
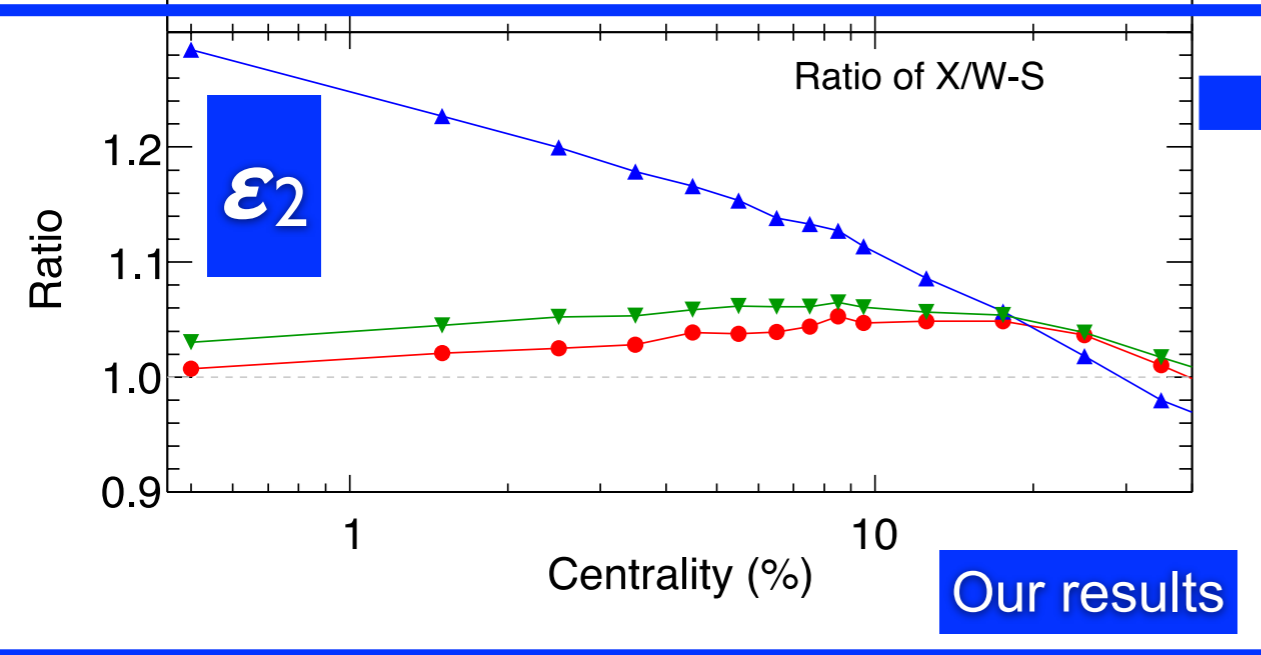
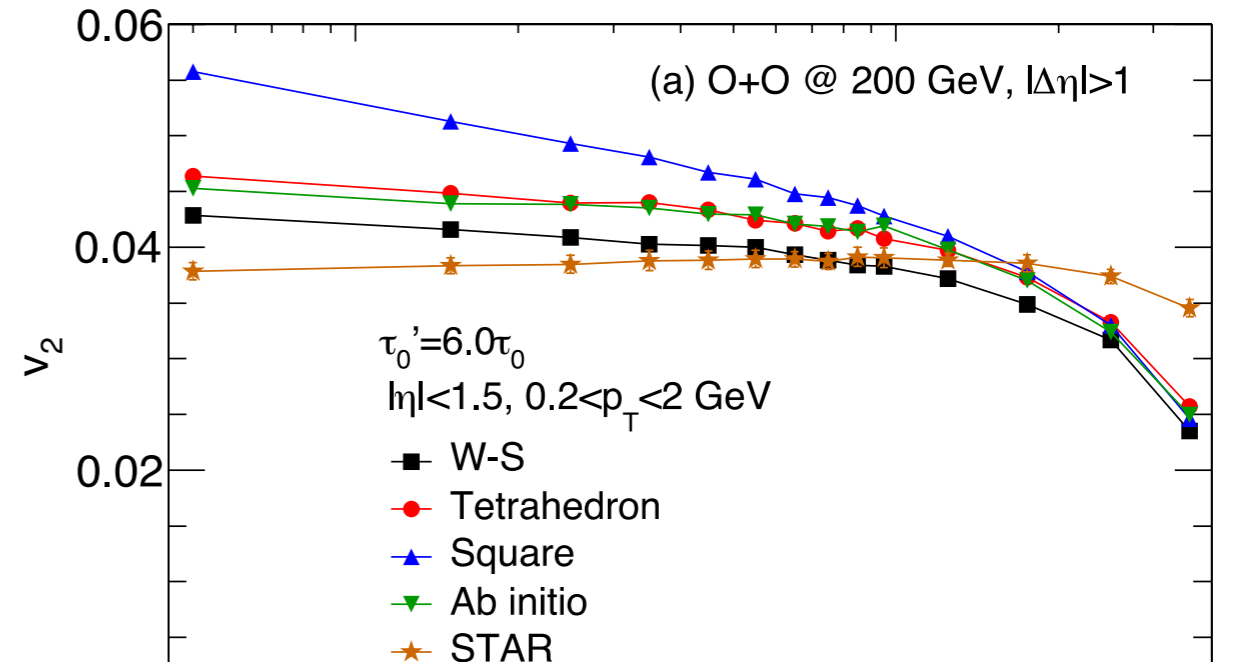
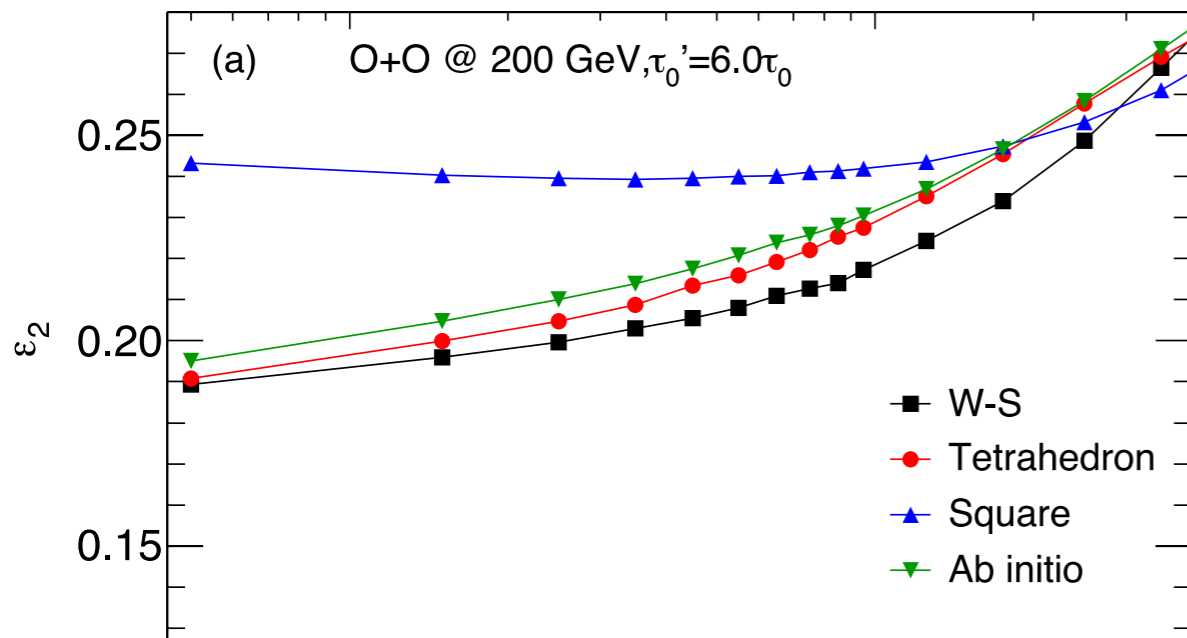
# v2 Ratio between with and without nuclear structure

Chunjian Zhang, Jinhui Chen, Giuliano Giacalone, Shengli Huang, Jiangyong Jia and Yu-Gang Ma, arXiv:2404.08385



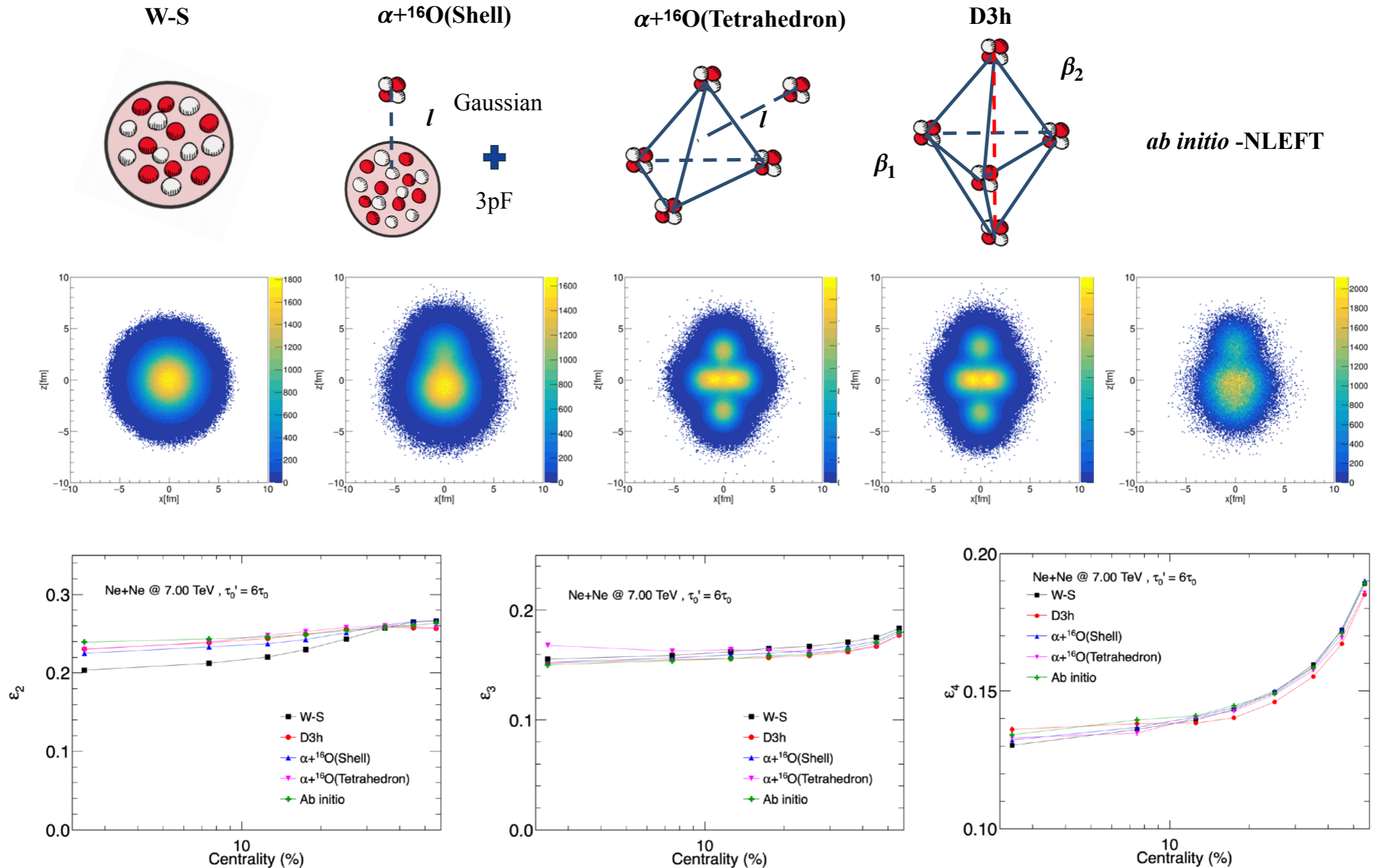
- Our  $v_2$  ratio for Ab initio (NLEFT) is consistent with their NLEFT.
- Our  $\langle p_T \rangle$  ratio for Tetrahedron is **inconsistent** with their VMC. **Why?**

# Ratio of $v_2$ and $\epsilon_2$ between with and without nuclear structure



➤ Our  $v_2$  ratio is consistent with our  $\epsilon_2$  ratio. Other reason?

# $^{20}\text{Ne} + ^{20}\text{Ne}$ at 7 TeV

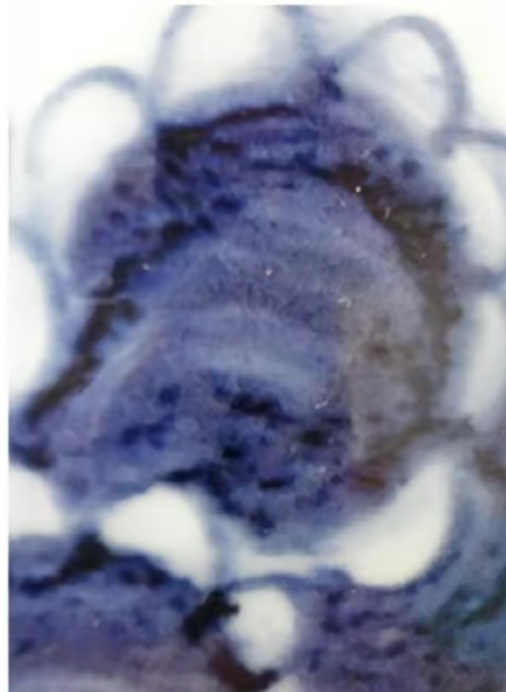


# Summary && Outlook

deformation



neutron skin/halo



$\alpha$ -cluster



- Different types of nuclear structures as the inputs of the AMPT model
- **Isobar collisions at RHIC**: Halo-type neutron skin or deformation?
- **O+O collisions at RHIC**: an extended effective parton formation time indicating a possible  $\alpha$ -clustering structure inside oxygen nuclei?
- More studies are in progress for Ne+Ne and O+O collisions at LHC
- Enjoy the beauty of nuclear porcelain!

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*Thank you!*