

Spinodal decomposition

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References

Supercritical Fluids
Fundamentals and Applications

Nuclear spinodal fragmentation

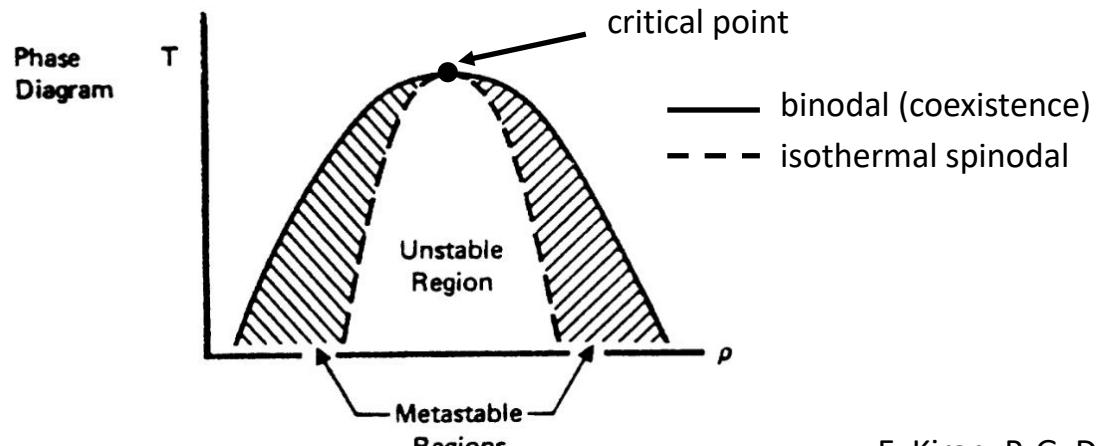
Philippe Chomaz^a, Maria Colonna^b, Jørgen Randrup^{c,*}

Spinodal decomposition during the hadronization stage at RHIC?

[hep-ph/0308271](#)

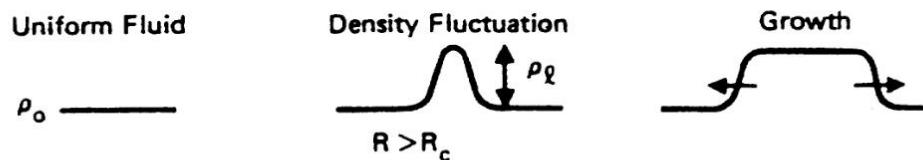
Jørgen Randrup

Two mechanisms of FOPT

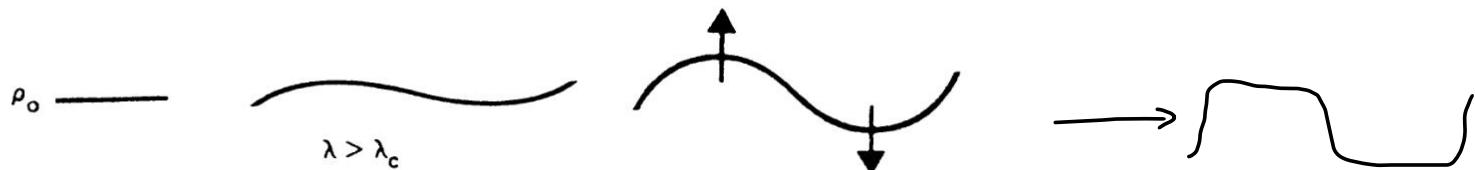


E. Kiran, P. G. Debenedetti, C. J. Peters
Supercritical Fluid

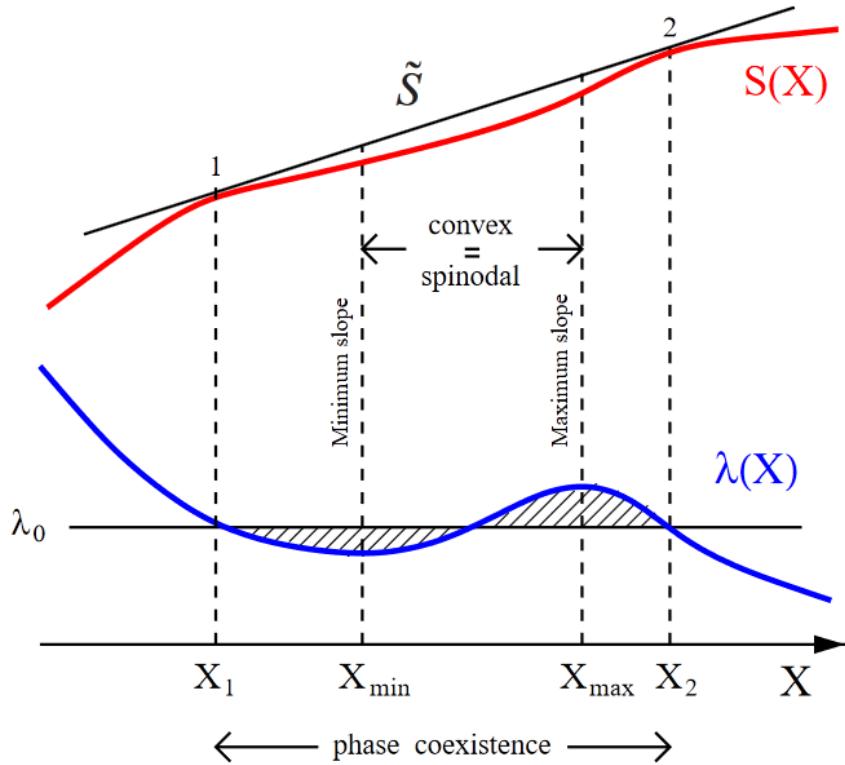
Metastable Region: Phase Transition by Nucleation



Unstable Region: Phase Transition by Spinodal Decomposition



Coexistence



$$S([X_1, X_2, \dots]) = \sum_i S_i(X_i),$$

$$0 \doteq \sum_i \delta S_i(X_i = \bar{X}_i) = \sum_{i\ell} \left(\frac{\partial S_i}{\partial X_i^\ell} \right)_{X_i = \bar{X}_i} \delta X_i^\ell = \sum_{i\ell} \lambda_i^\ell \delta X_i^\ell,$$

$$\lambda_i^\ell \equiv \partial S_i / \partial X_i^\ell \quad \delta X^\ell = \sum_i \delta X_i^\ell \doteq 0$$

e.g.

$$\lambda_i^E \equiv \frac{\partial S_i}{\partial E_i} = \frac{1}{T_i}, \quad \lambda_i^N \equiv \frac{\partial S_i}{\partial N_i} = -\frac{\mu_i}{T_i}, \quad \lambda_i^V \equiv \frac{\partial S_i}{\partial V_i} = \frac{P_i}{T_i},$$

$$\tilde{S} = N_1 \sigma(x_1) + N_2 \sigma(x_2) = \frac{N_1}{N} S(X_1) + \frac{N_2}{N} S(X_2) > S.$$

Maxwell construction

$$\int_{X_1}^{X_2} dX (\lambda(X) - \lambda_0) \doteq 0,$$

Spinodal instability

$$p_\kappa(m) = \kappa \frac{T^4}{2\pi^2} \sum_{n=1}^{\infty} \frac{\kappa^n}{n^4} \left(\frac{nm}{T}\right)^2 K_2\left(\frac{nm}{T}\right) \asymp \frac{\pi^2}{90} T^4, \quad (1)$$

- bag model

$$\begin{aligned} \varepsilon_\kappa(m) = \kappa \frac{T^4}{2\pi^2} \sum_{n=1}^{\infty} \frac{\kappa^n}{n^4} & \left[3 \left(\frac{nm}{T}\right)^2 K_2\left(\frac{nm}{T}\right) \right. \\ & \left. + \left(\frac{nm}{T}\right)^3 K_1\left(\frac{nm}{T}\right) \right] \asymp \frac{\pi^2}{30} T^4, \end{aligned} \quad (2)$$

$\kappa = +$ for bosons and $\kappa = -$ for fermions

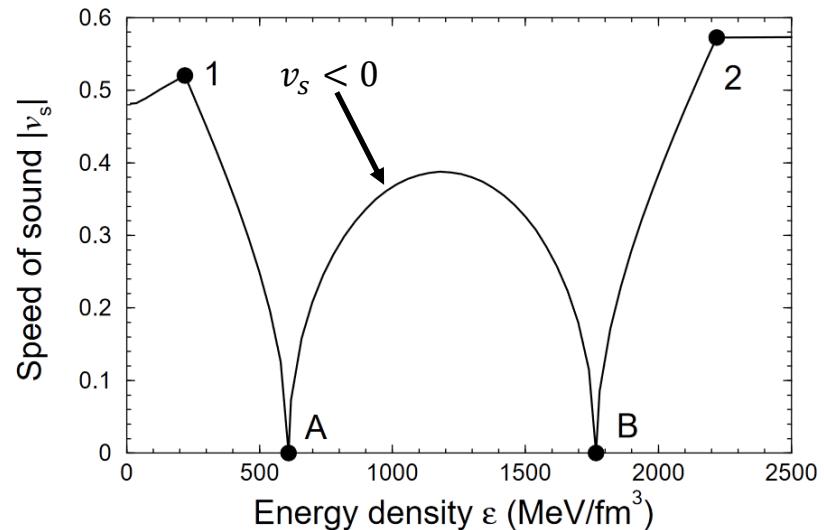
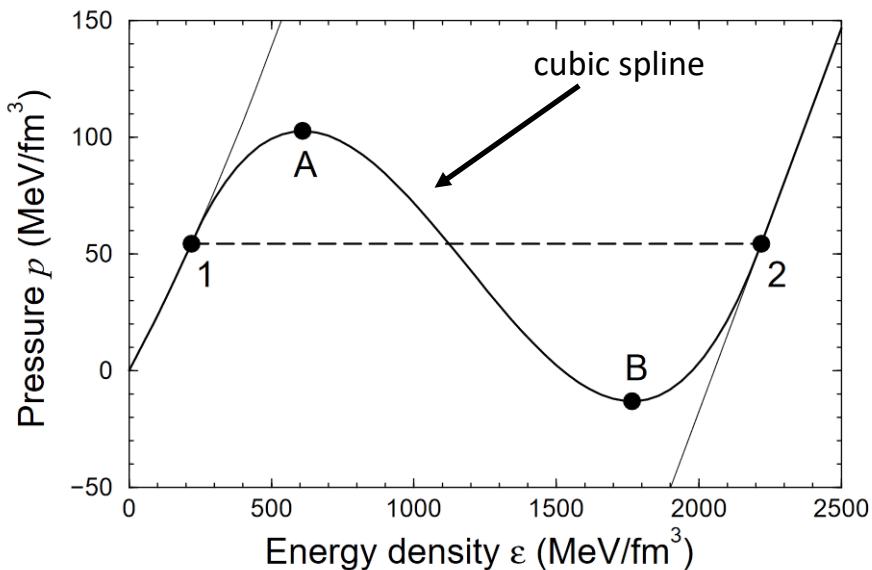
$$p_{\text{had}}(T) = \sum_{i=\pi}^{\Delta} g_i p_{\kappa_i}(m_i), \quad \varepsilon_{\text{had}}(T) = \sum_{i=\pi}^{\Delta} g_i \varepsilon_{\kappa_i}(m_i). \quad (3) \quad i = \pi, K, \eta, \rho, \dots, N, \Lambda, \Sigma, \Delta$$

$$p_{\text{qgp}}(T) = [g_g + \frac{7}{8}g_q] \frac{\pi^2}{90} T^4 + g_s p_-(m_s) - B, \quad (4) \quad g_g = 16, g_q = 24, \text{ and } g_s = 12.$$

$$\begin{aligned} \varepsilon_{\text{qgp}}(T) = [g_g + \frac{7}{8}g_q] \frac{\pi^2}{30} T^4 + g_s \varepsilon_-(m_s) + B. \quad (5) \quad m_g, m_q = 0, \quad m_s = 150 \text{ MeV} \\ T_c = 170 \text{ MeV} \quad B^{1/4} = 250 \text{ MeV} \end{aligned}$$

Spinodal instability

- bag model



$$T^{\mu\nu} = (\varepsilon + p)u^\mu u^\nu - pg^{\mu\nu}$$

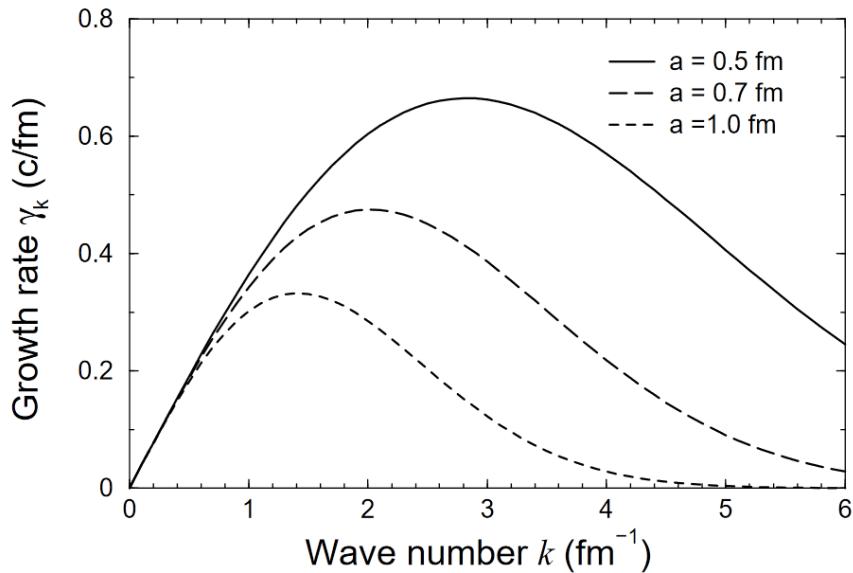
$$\varepsilon(t, \mathbf{r}) = \varepsilon_0 + \delta\varepsilon(t, \mathbf{r}), \quad \delta\varepsilon(t, \mathbf{r}) \ll \varepsilon_0.$$

$$0 \approx \partial_t \varepsilon - \nabla \cdot [(\varepsilon + p)\mathbf{v}] \approx \partial_t \delta\varepsilon - (\varepsilon_0 + p_0) \nabla \cdot \mathbf{v}, \quad (14)$$

$$\mathbf{0} \approx \partial_t [(\varepsilon + p)\mathbf{v}] - \nabla p \approx (\varepsilon_0 + p_0) \partial_t \mathbf{v} - v_s^2 \nabla \delta\varepsilon, \quad (15)$$

$$\partial_t^2 \delta\varepsilon(t, \mathbf{r}) = \frac{\partial p_0}{\partial \varepsilon_0} \nabla^2 \delta\varepsilon(t, \mathbf{r}). \quad (16)$$

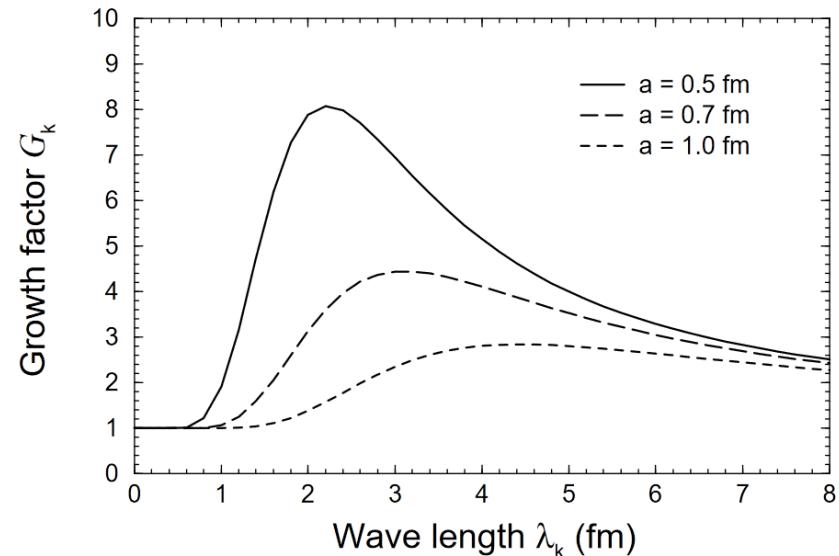
Spinodal instability



Consider finite range of interaction

$$p(\mathbf{r}_1) \approx \int d^3\mathbf{r}_2 g(|\mathbf{r}_1 - \mathbf{r}_2|) p(\varepsilon(\mathbf{r}_2)) . \quad (17)$$

$$g(r) \sim \exp(-r^2/2a^2),$$



Consider finite traverse time in spinodal region

$$G_k = \exp(\Gamma_k)$$

$$\Gamma_k \equiv \int_{t_B}^{t_A} \gamma_k(\varepsilon(t)) dt \approx \gamma_k(\tilde{\varepsilon}) \Delta t_{\text{eff}} .$$

$$t_B - t_A \approx 4 \text{ fm}/c$$

following work

Phase transition dynamics for baryon-dense matter

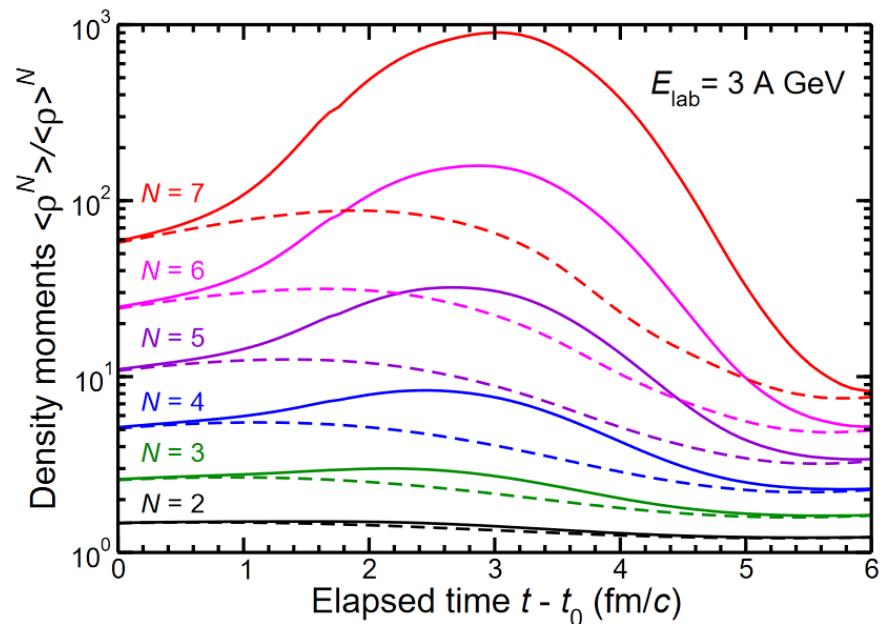
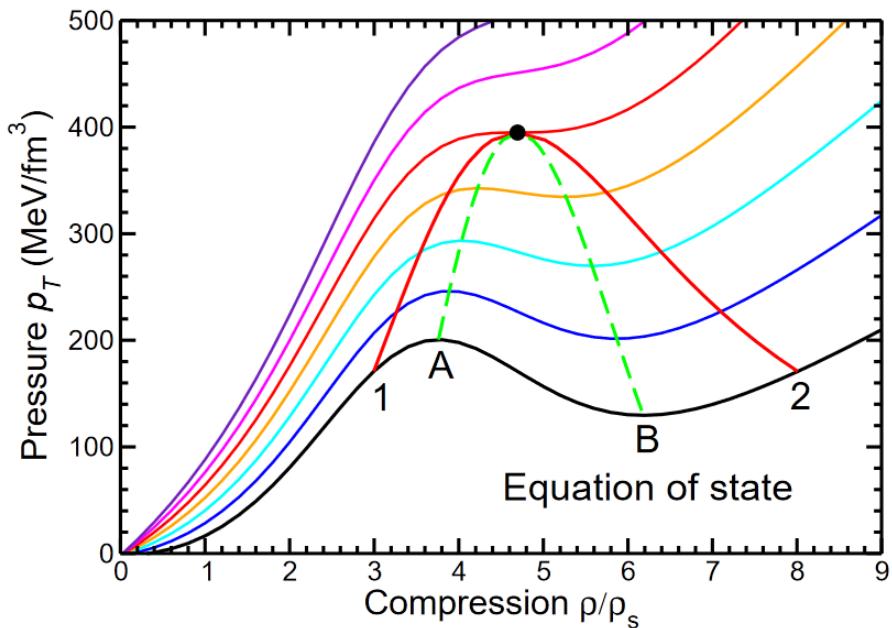
0903.4736

Jørgen Randrup

Spinodal amplification of density fluctuations
in fluid-dynamical simulations of relativistic nuclear collisions

1209.2462

Jan Steinheimer and Jørgen Randrup



Thanks for your listening!

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