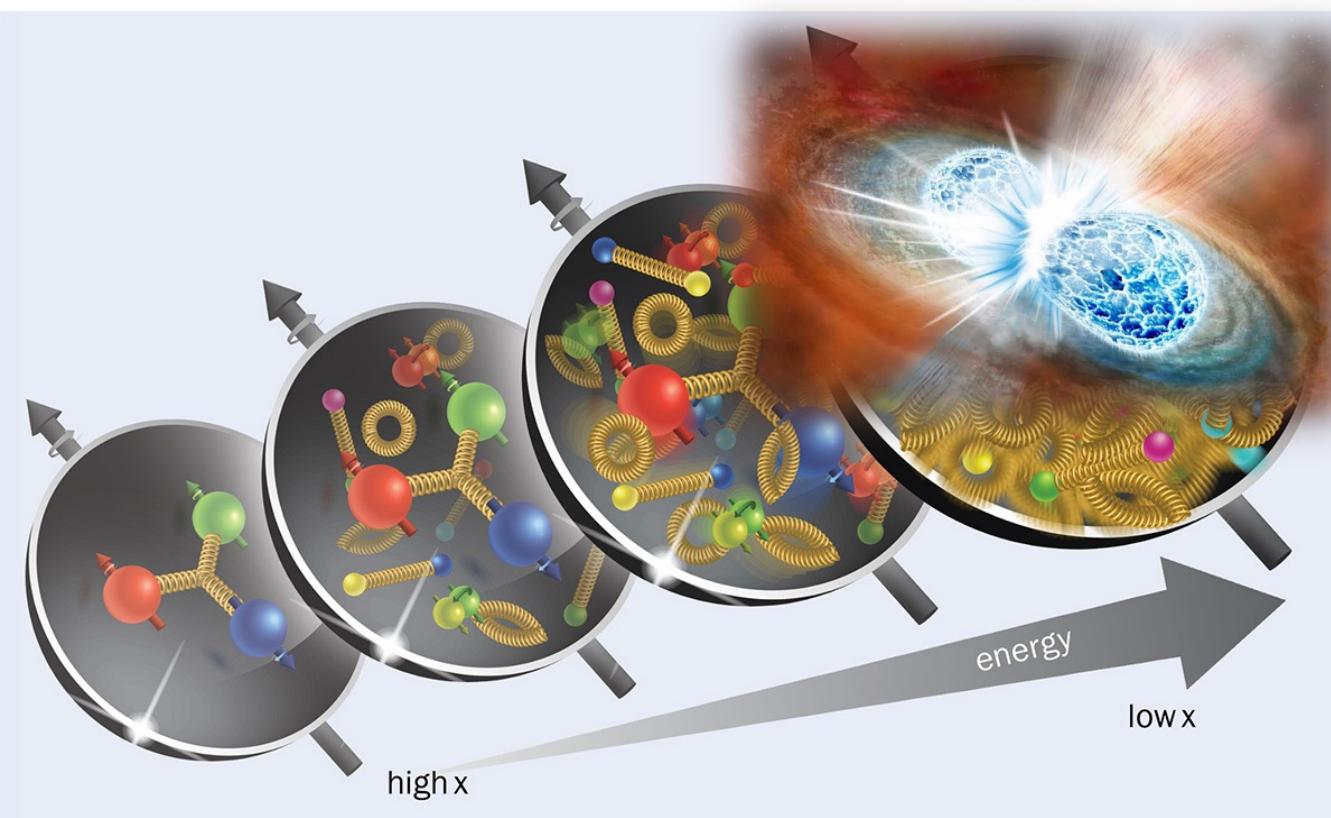


Production of baryon clusters of $B=\pm 1$ to ± 4 in relativistic heavy-ion collisions

Zhangbu Xu

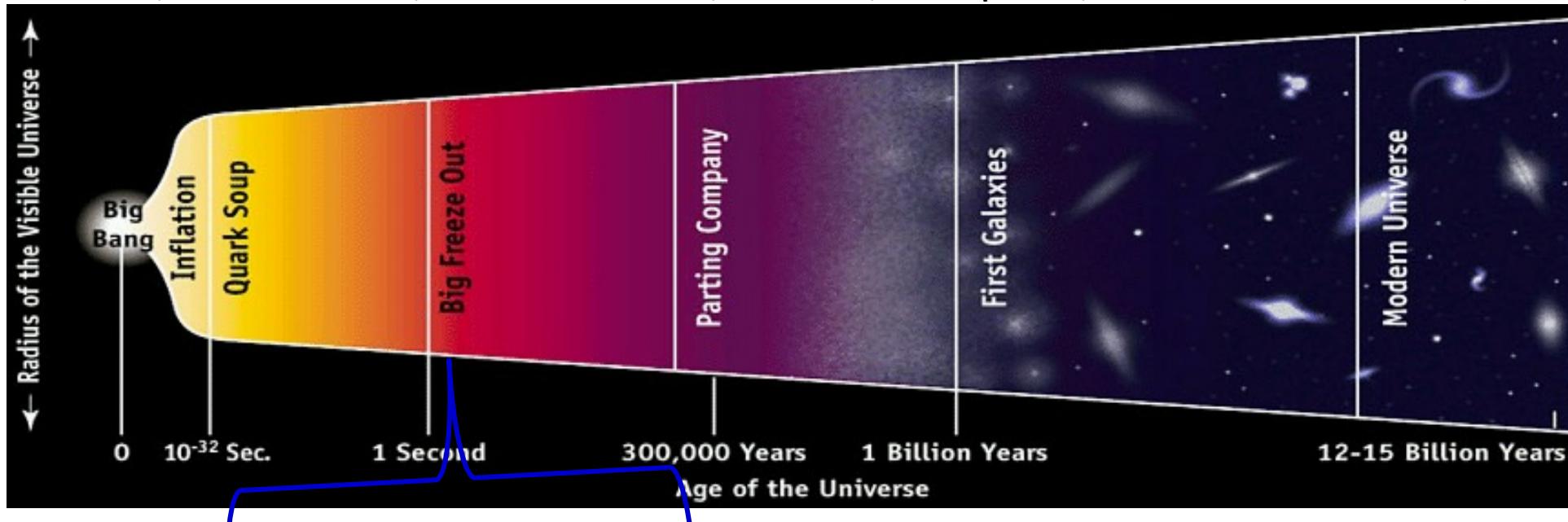
(Kent State University& BNL)

- $B=1$: simplest gluon topology in QCD
- $B=3,2$ nuclear yield ratio as probe of quantum wavefunction overlaps and density fluctuation and correlations
- $B=4,3$ Hypernuclear properties
- Discovery of $B=-4$ hypernucleus
- Future:
charmed hypernuclei

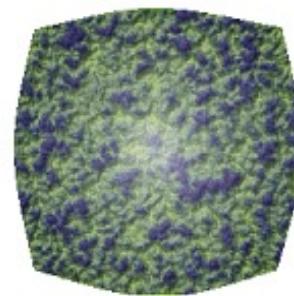
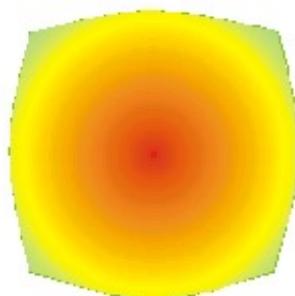


Little Big Bangs

BIG; All 4 forces at work; Gravitation dominates; QGP@ 10^{-6} s; Slow expansion; Antimatter-matter annihilate;



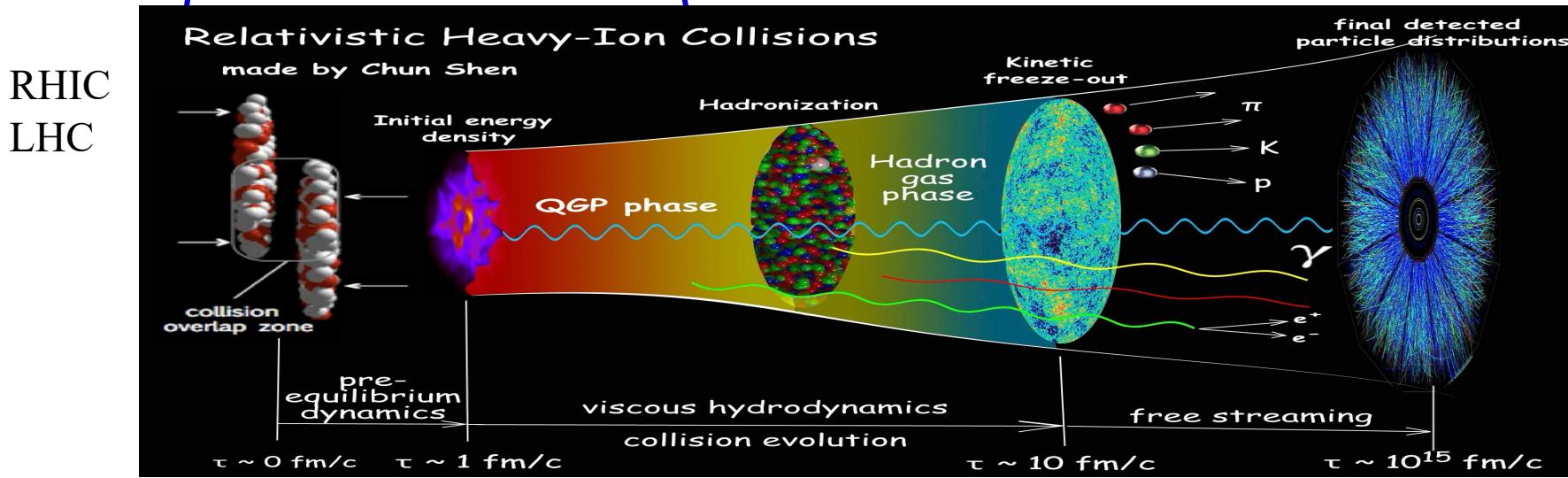
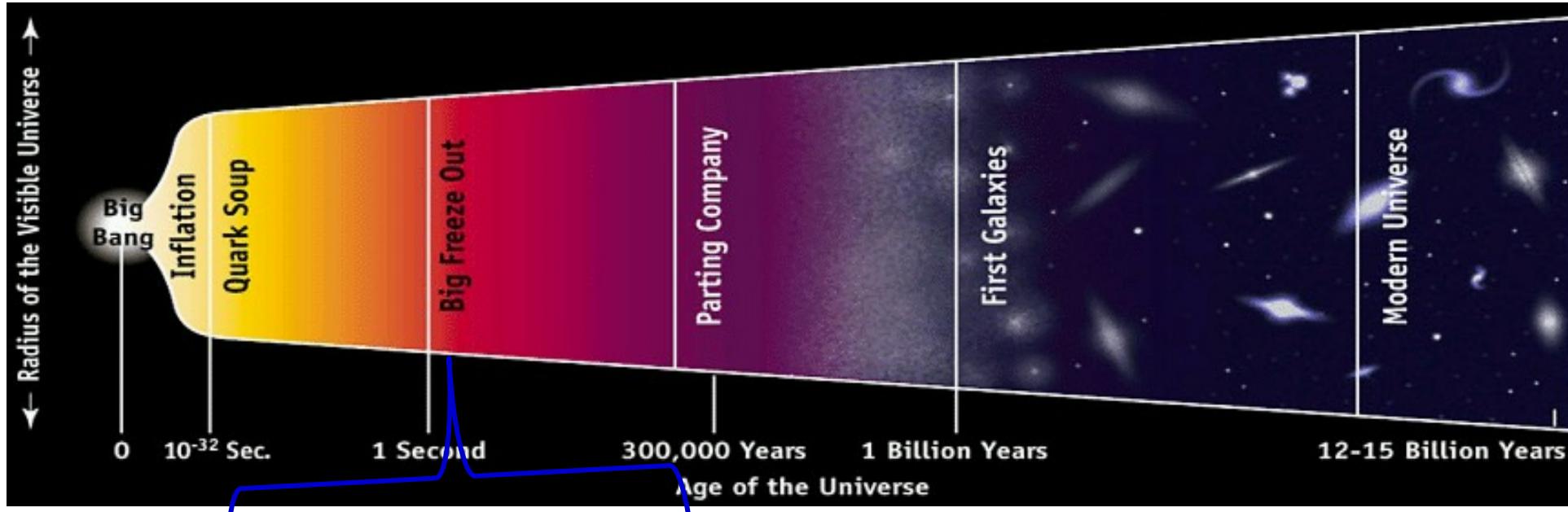
RHIC
LHC



TIME

Little Big Bangs

BIG; All 4 forces at work; Gravitation dominates; QGP@ 10^{-6} s; Slow expansion; Antimatter-matter annihilate;

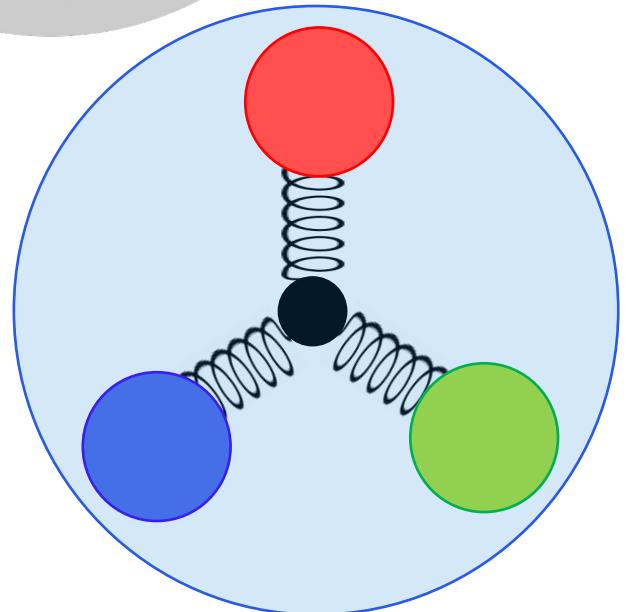
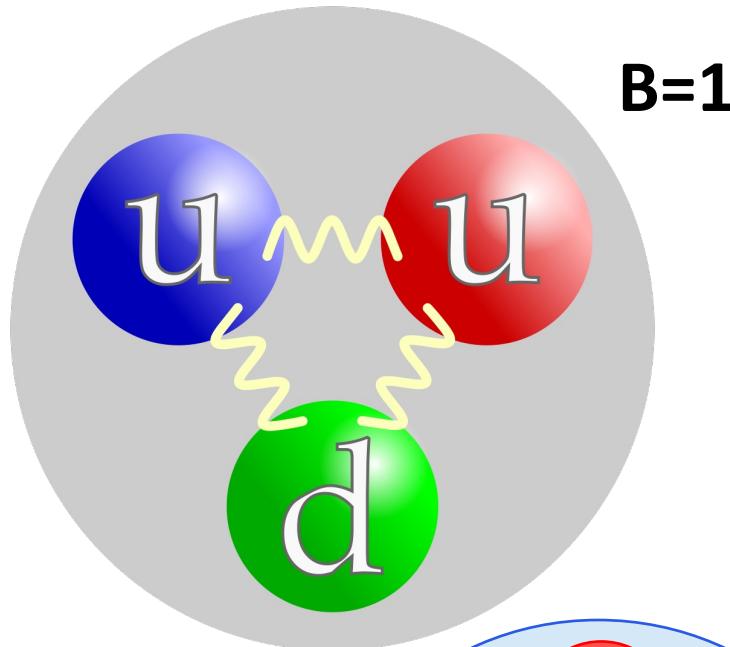


Little; Strong force at work; QGP@ 10^{-23} s; Fast expansion; Antimatter-matter decouple; repeat trillion times

TIME

Baryon Number (B) Carrier

- Textbook picture of a proton
 - Lightest baryon with strictly conserved baryon number
 - Each valence quark carries $1/3$ of baryon number
 - Proton lifetime $>10^{34}$ years
 - Quarks are connected by gluons
- Alternative picture of a proton
 - Proposed at the Dawn of QCD in 1970s
 - A Y-shaped gluon junction topology carries baryon number ($B=1$)
 - The topology number is the strictly conserved number
 - Quarks do not carry baryon number
 - Valence quarks are connected to the end of the junction always

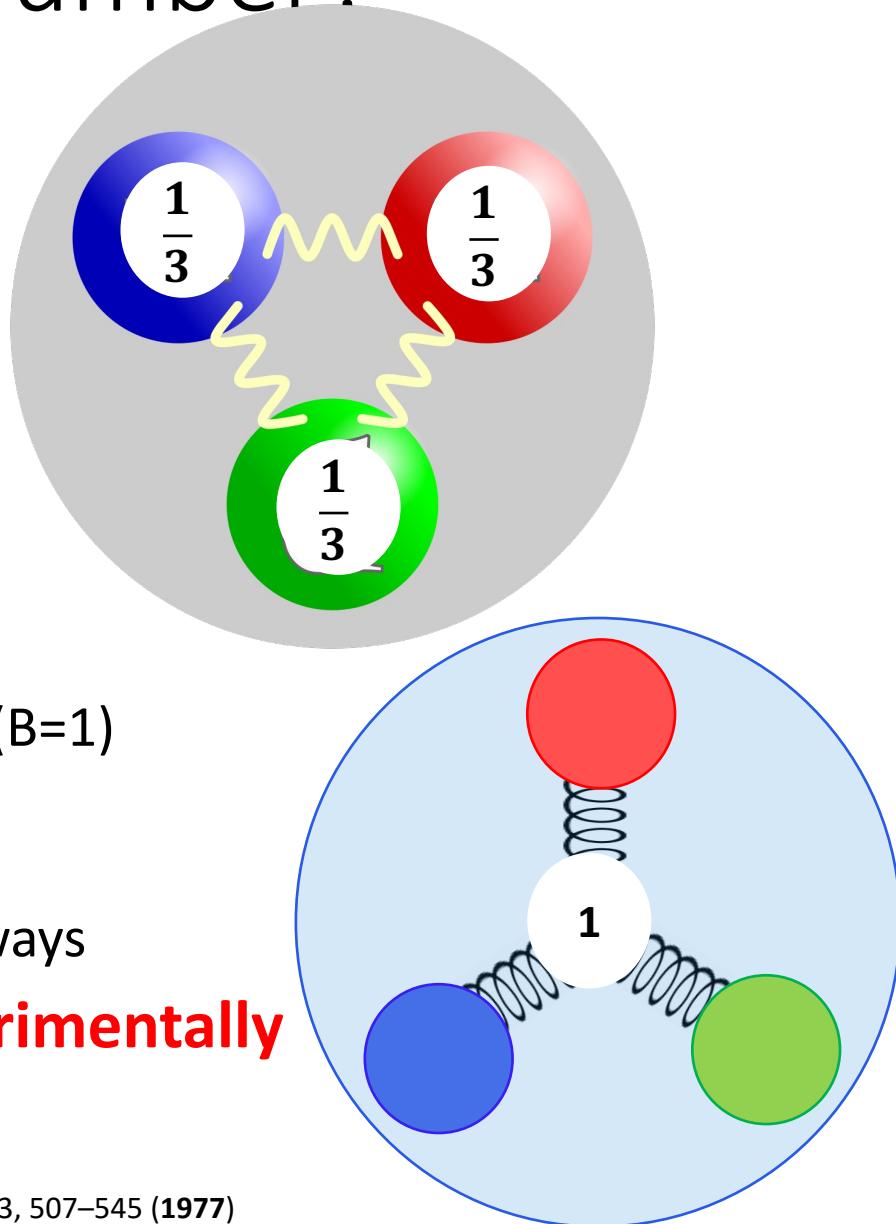


[1]: Artru, X.; String Model with Baryons: Topology, Classical Motion. Nucl. Phys. B 85, 442–460 (1975).

[2]: Rossi, G. C. & Veneziano, G. A; Possible Description of Baryon Dynamics in Dual and Gauge Theories. Nucl. Phys. B 123, 507–545 (1977)

Measurements of quark baryon number?

- Textbook picture of a proton
 - Lightest baryon with strictly conserved baryon number
 - Each valence quark carries $1/3$ of baryon number
 - Proton lifetime $>10^{34}$ years
 - Quarks are connected by gluons
- Alternative picture of a proton
 - Proposed at the Dawn of QCD in 1970s
 - A Y-shaped gluon junction topology carries baryon number ($B=1$)
 - The topology number is the strictly conserved number
 - Quarks do not carry baryon number
 - Valence quarks are connected to the end of the junction always
- **Neither of these postulations has been verified experimentally**



[1]: Artru, X.; String Model with Baryons: Topology, Classical Motion. Nucl. Phys. B 85, 442–460 (1975).

[2]: Rossi, G. C. & Veneziano, G. A; Possible Description of Baryon Dynamics in Dual and Gauge Theories. Nucl. Phys. B 123, 507–545 (1977)

Model implementations of baryons at RHIC

2003 RBRC Workshop on “Baryon Dynamics at RHIC”

FIRST WORKSHOP ON BARYON DYNAMICS
FROM RHIC TO EIC

- Many of the models used for heavy-ion collisions at RHIC (HIJING, AMPT, UrQMD) have implemented a nonperturbative baryon stopping mechanism

V. Topor Pop, *et al*, Phys. Rev. C **70**, 064906 (2004)

Zi-Wei Lin, *et al*, Phys. Rev. C **72**, 064901 (2005)

M. Bleicher, *et al*, J.Phys.G **25**, 1859-1896 (1999)

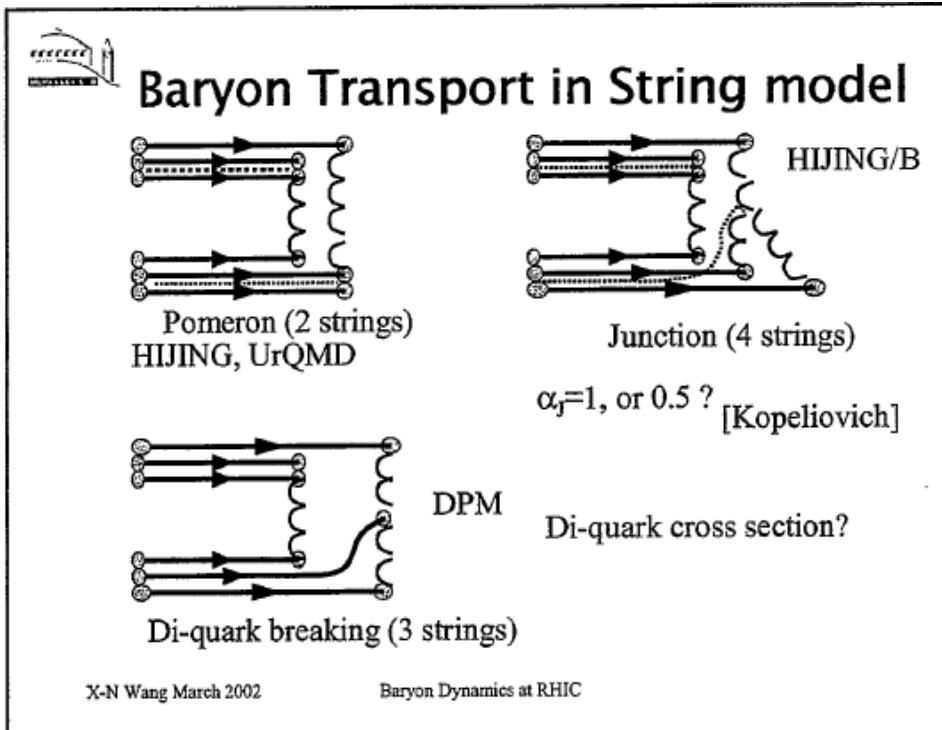
- Baryon Stopping
 - Theorized to be an effective mechanism of stopping baryons in $p\bar{p}$ and AA

D. Kharzeev, Physics Letters B **378**, 238-246 (1996)

- Specific rapidity dependence is predicted:

$$p = \sim e^{-\alpha_B y}$$

$$\alpha_B \sim= 0.5$$



There is only one way to construct a gauge-invariant state vector of a baryon from quarks and gluons

It is evident from the structure of (1) that the trace of baryon number should be associated not with the valence quarks, but with a non-perturbative configuration of gluon fields located at the point x - the “string junction” .

Dates: Jan 22 – 24, 2024
Location: Center for Frontiers in Nuclear Science (CFNS), Stony Brook University
Format: In-person & zoom
Participation: Invited Talks + Open Mic Discussion
Registration Deadline: Jan 15th, 2024
No registration fee - Limited student support available

Scientific Motivation:

This workshop aims to address fundamental questions such as what carries the baryon quantum number and how a baryon is stopped in high-energy collisions, which have profound implications for understanding the baryon structure. It also challenges our current knowledge of QCD and its non-perturbative aspects, such as baryon junctions and gluonic topology. The workshop will explore the origin and transport of baryons in high-energy collisions, from the AGS/SPS/RHIC/LHC to JLab F_n, HERA/EIC, and discuss the experimental and theoretical challenges and opportunities in this field.

Key Topics:

- Baryon junctions and gluonic topology
- Baryon and charge stopping in heavy-ion collisions
- Baryon transport in photon-induced processes
- Baryon-meson-transition in backward u-channel reaction
- Models of baryon dynamics and baryon-rich matter
- Novel experimental methods at EIC

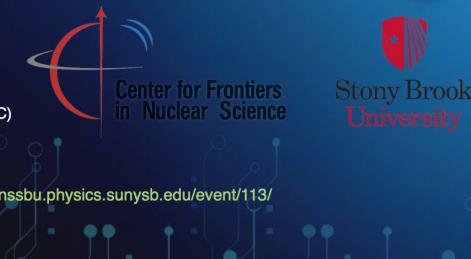
Keynote speaker: Gabriele Veneziano

Organizers:

- D. Kharzeev (SBU/BNL)
W. B. Li (SBU/CFNS)
N. Lewis (Rice)
J. Norohna Hostler (UIUC)
C. Shen (Wayne State/RBRC)
P. Tribedy (BNL)
Z. Xu (BNL)

Webpage: <https://indico.cfnsbu.physics.sunysb.edu/event/113/>

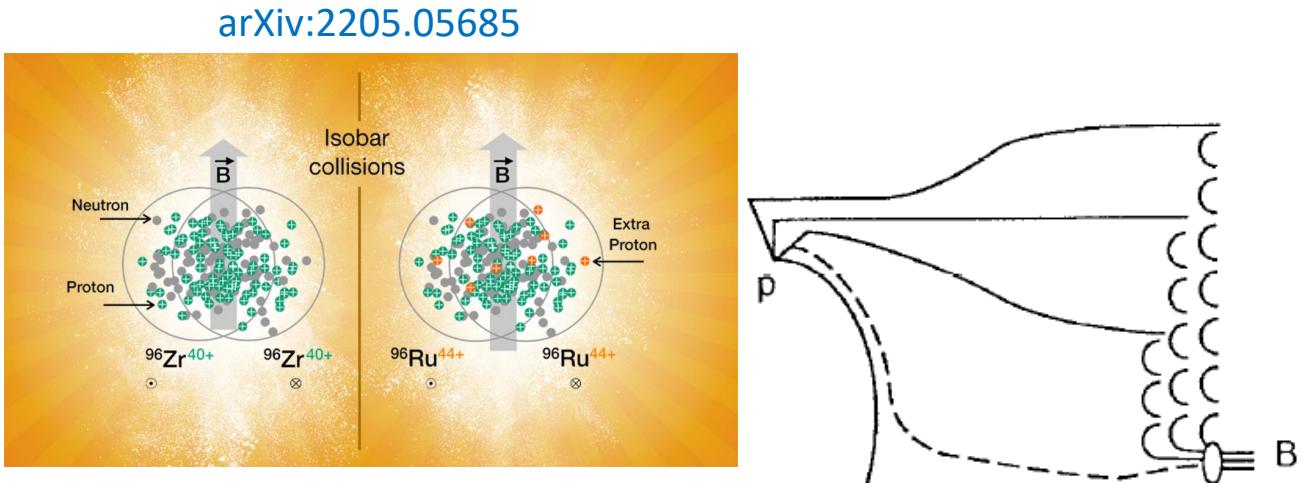
Contact: ptribedy@bnl.gov



Three approaches toward tracking the origin of the baryon number

1. STAR Method:

Charge (Q) stopping vs baryon (B) stopping:
if valence quarks carry Q and B,
 $Q=B$ at middle rapidity



2. Kharzeev-STAR Method:

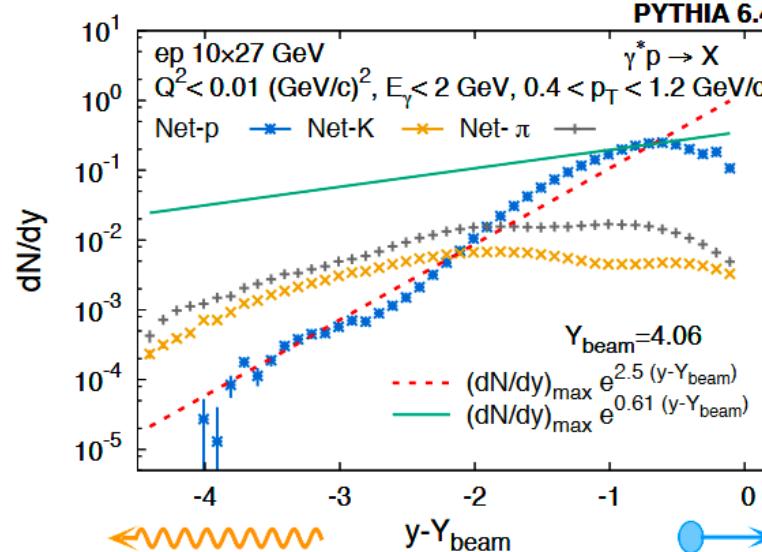
If gluon topology (J) carries B as one unit, it should show scaling according to Regge theory

$$p = \sim e^{-\alpha_B y}$$

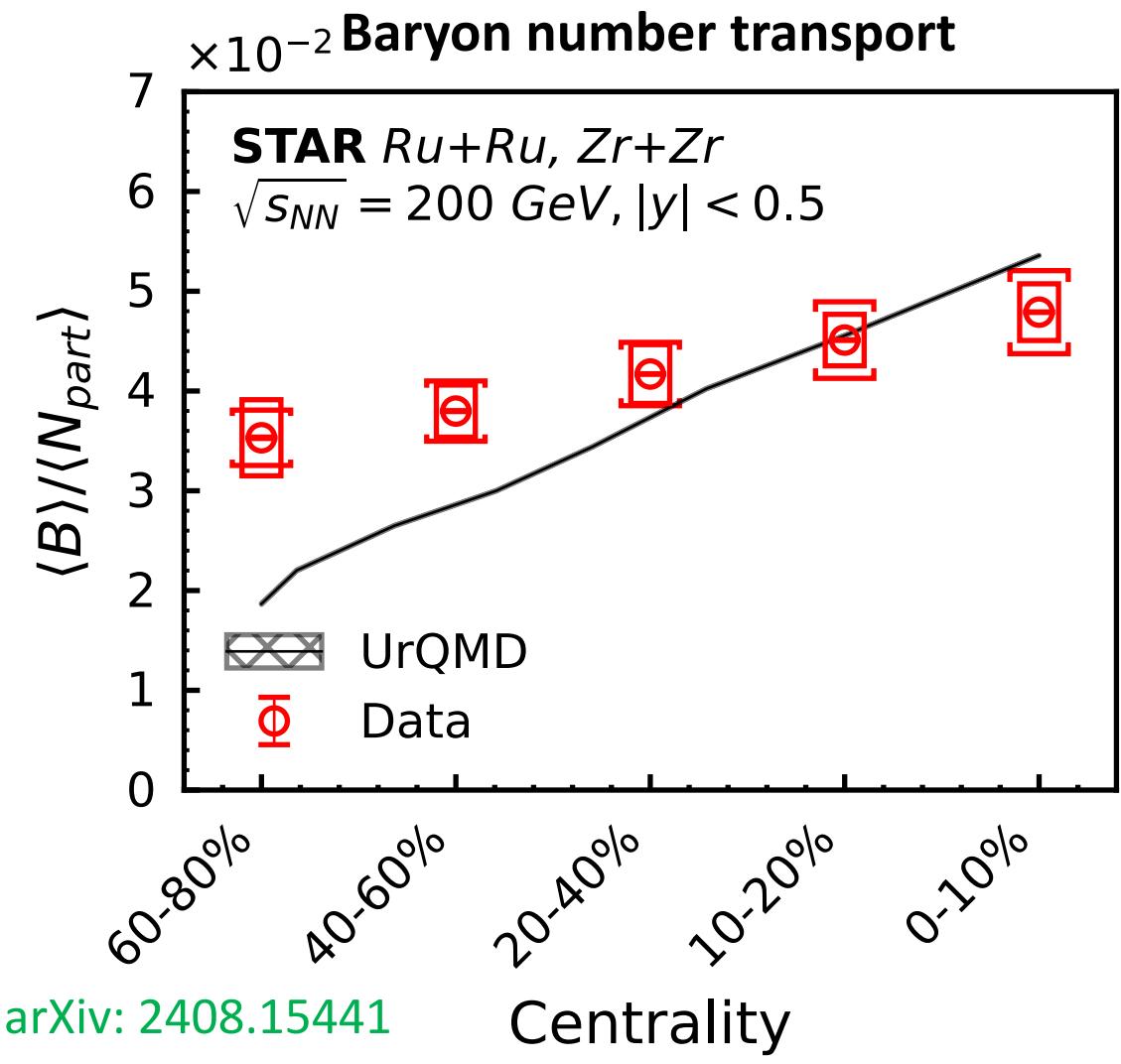
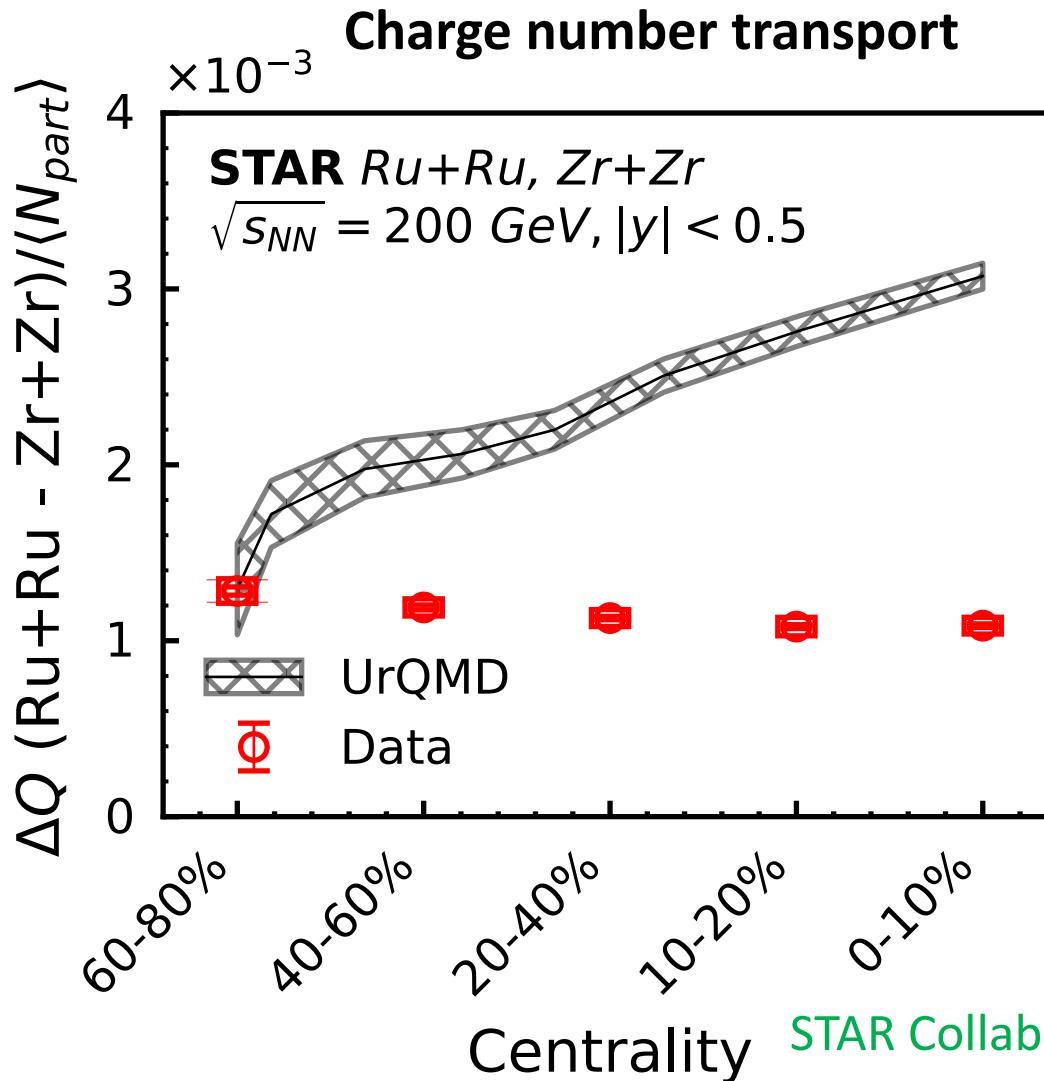
$$\alpha_B \approx 0.5$$

3. Artru Method:

In $\gamma+Au$ collision, rapidity asymmetry can reveal the origin



Separate charge and baryon transports



UrQMD matches data on charge stopping better in peripheral; better on baryon stopping in central
overpredicts charge stopping in central; underpredicts baryon stopping in peripheral

Three approaches toward tracking the origin of the baryon number

1. STAR Method:

Charge (Q) stopping vs baryon (B) stopping:
if valence quarks carry Q and B,
 $Q=B$ at middle rapidity

$$B/Q=2$$

2. Kharzeev-STAR Method:

If gluon topology (J) carries B as one unit,
it should show scaling according to
Regge theory

$$\alpha_B = 0.61$$

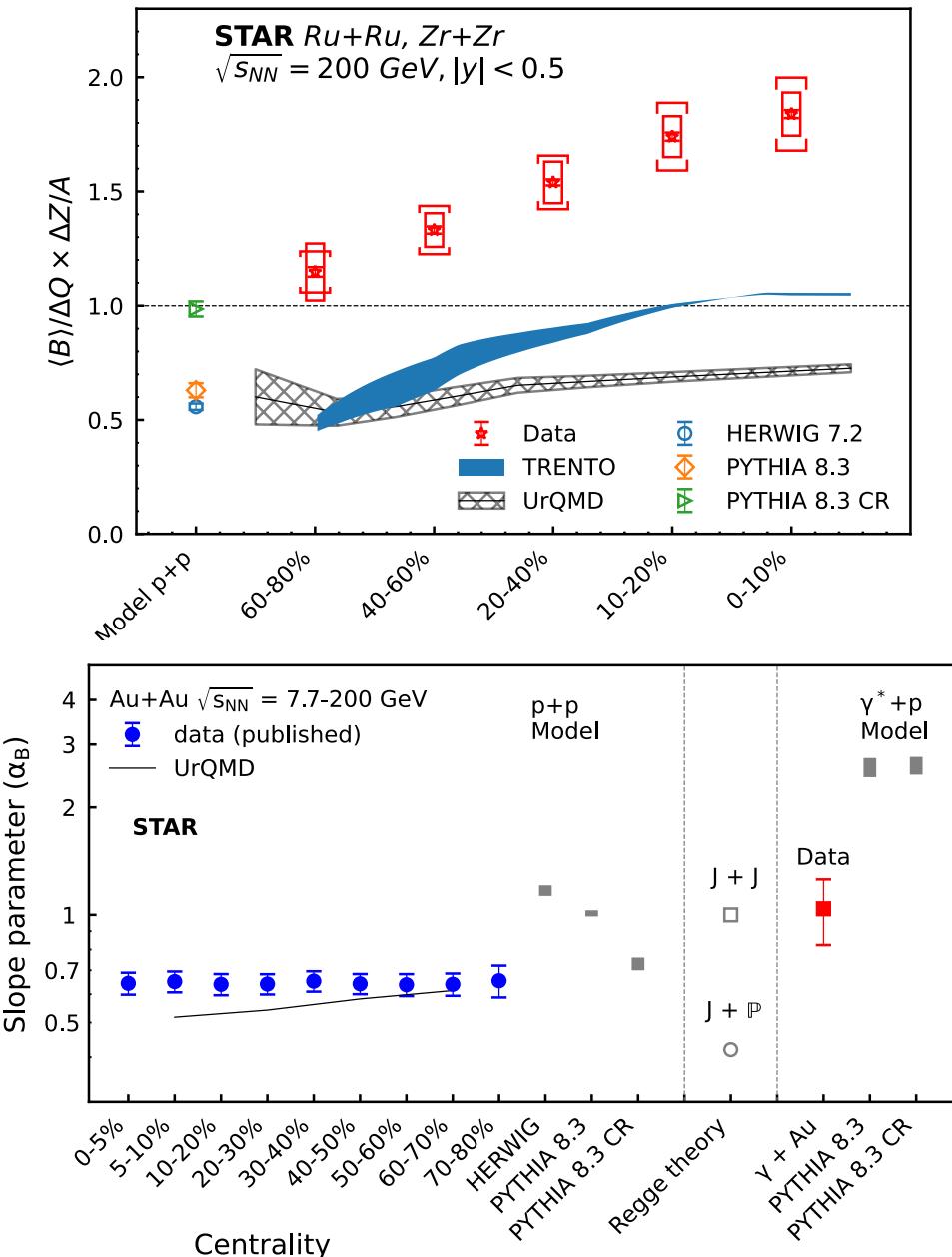
$$p = \sim e^{-\alpha_B y}$$

$$\alpha_B \approx 0.5$$

3. Artru Method:

In $\gamma+A$ collision, rapidity asymmetry can reveal the origin

$$\alpha_B(A+A)=0.61 < \alpha_B(\gamma+A)=1.1 < \alpha_B(\text{PYTHIA})$$



B=1,2,3 nuclear yield ratios

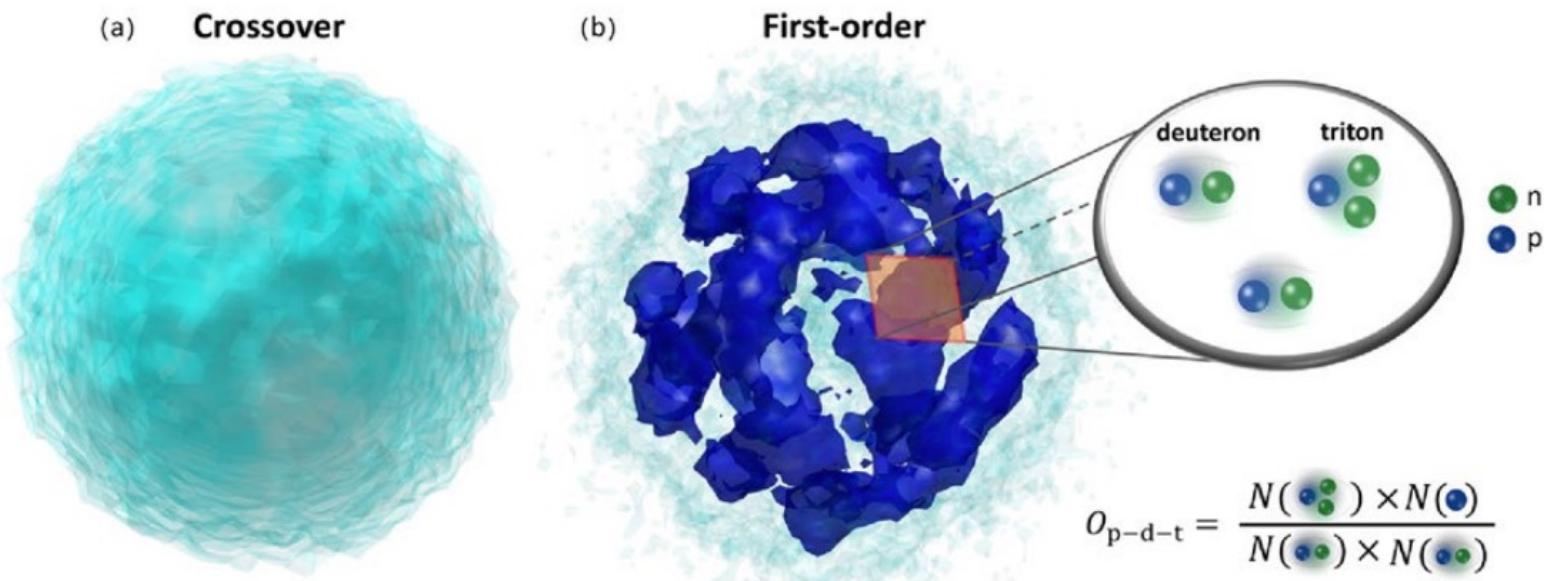


Fig. 1 (Color online) Density distribution of strongly interacting matter in a heavy ion collision after its expansion for the cases of crossover transition (panel **a**) and first-order chiral phase transition (panel **b**). Also shown for illustration of the latter case are deuterons and tritons produced from the density fluctuating hadronic matter and their yield ratio $\mathcal{O}_{p-d-t} = N_t N_p / N_d^2$, which depends on the magnitude of neutron density distribution as discussed in the text

- Light nuclei production as a probe of the QCD phase diagram
K.J. Sun, et al., PLB 781 (2018) 499

- Probing QCD critical fluctuations from light nuclei production in relativistic heavy-ion collisions

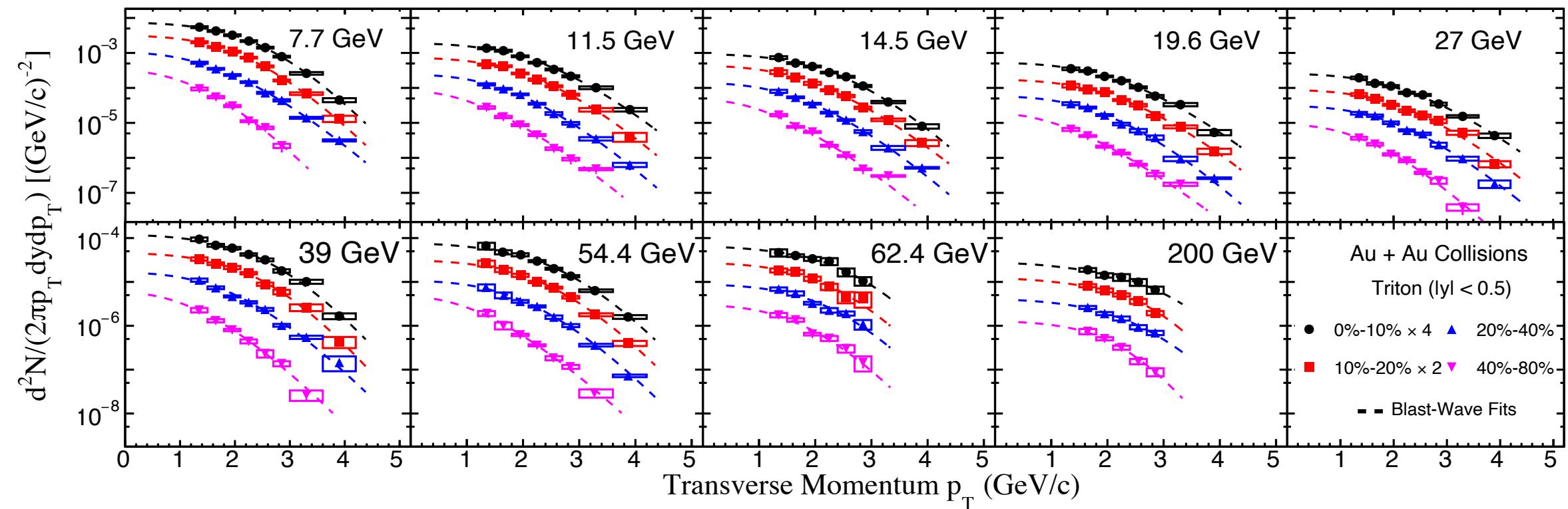
K.J. Sun, et al., PLB 774 (2017) 103

C.M. Ko, NST 34 (2023) 80

$$\mathcal{O}_{p-d-t} = \frac{N_{\text{triton}} N_{\text{deuteron}}}{N_{\text{triton}}^2} = g \frac{1 + (1 + 2\alpha) \Delta n}{(1 + \alpha \Delta n)^2},$$

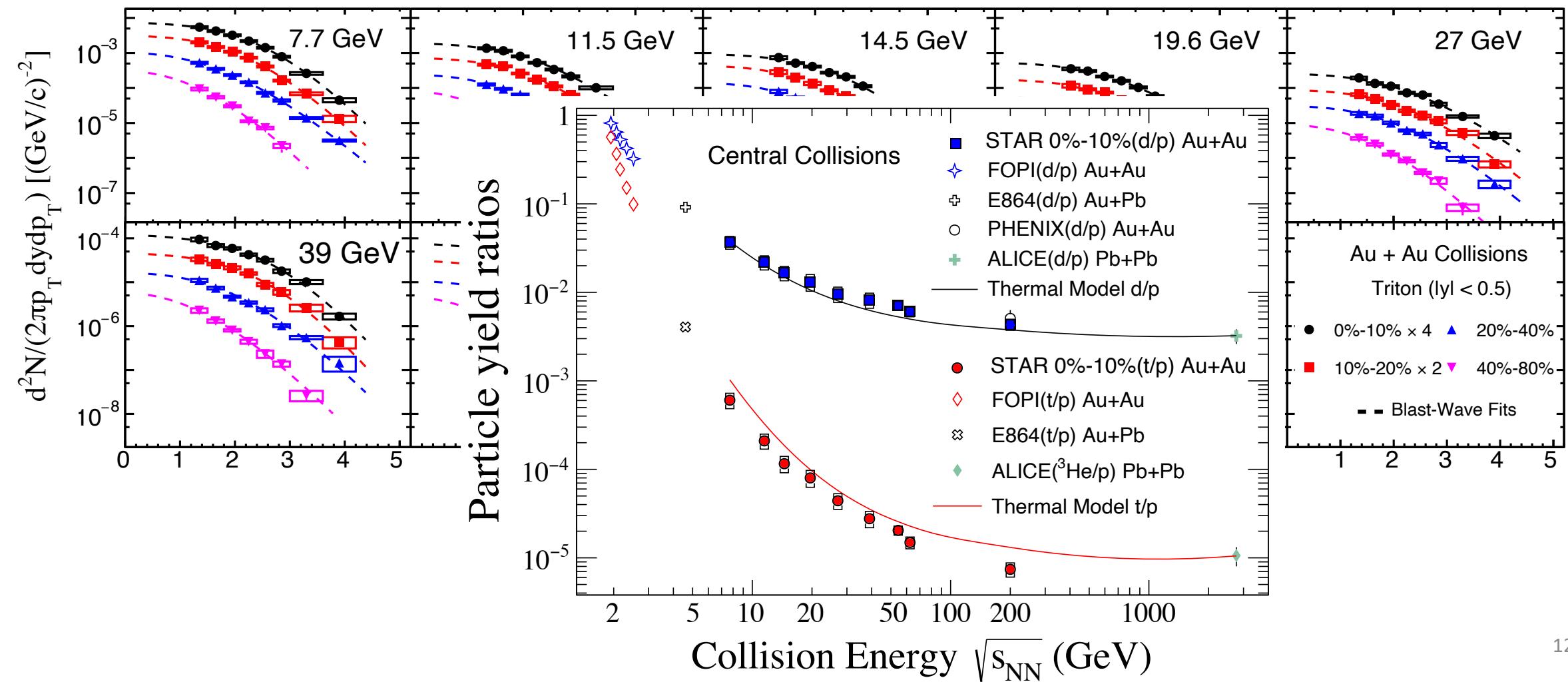
Spectra and two-particle ratios

STAR, Phys.Rev.Lett. 130 (2023) 202301



Spectra and two-particle ratios

STAR, Phys.Rev.Lett. 130 (2023) 202301



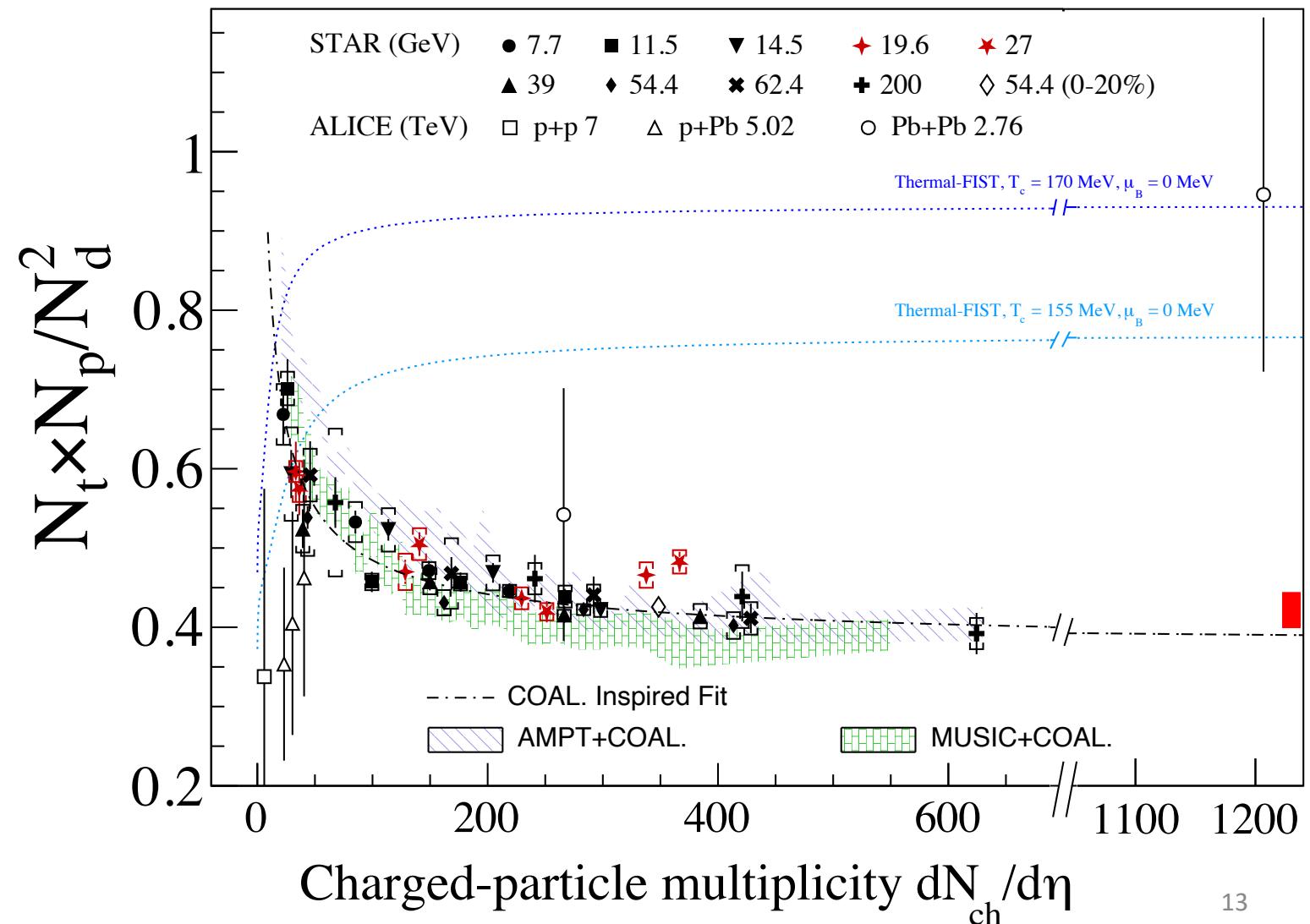
Quantum Wavefunction overlap efficiency

STAR, Phys.Rev.Lett. 130 (2023) 202301

RHIC+LHC data, C. Pinto, X.F. Luo, XZB,EMMI RRTF 04/24

$$\frac{N_t \times N_p}{N_d^2} = p_0 \times \left(\frac{R^2 + \frac{2}{3}r_d^2}{R^2 + \frac{1}{2}r_t^2} \right)^3$$

Coalescence wavefunction
overlap between nucleus
and nucleons



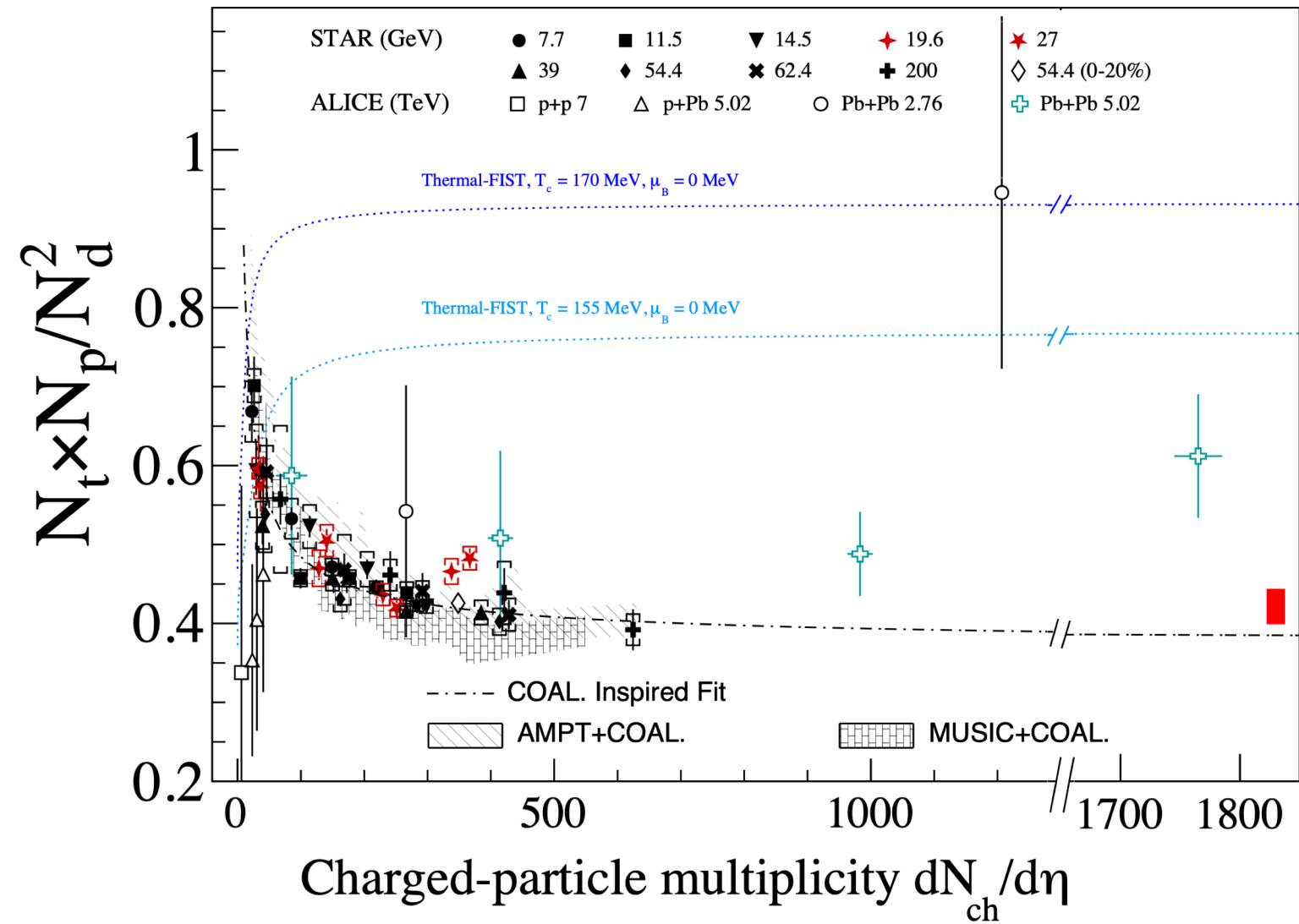
Quantum Wavefunction overlap efficiency

STAR, *Phys.Rev.Lett.* 130 (2023) 202301

RHIC+LHC data, C. Pinto, X.F. Luo, XZB,EMMI RRTF 04/24

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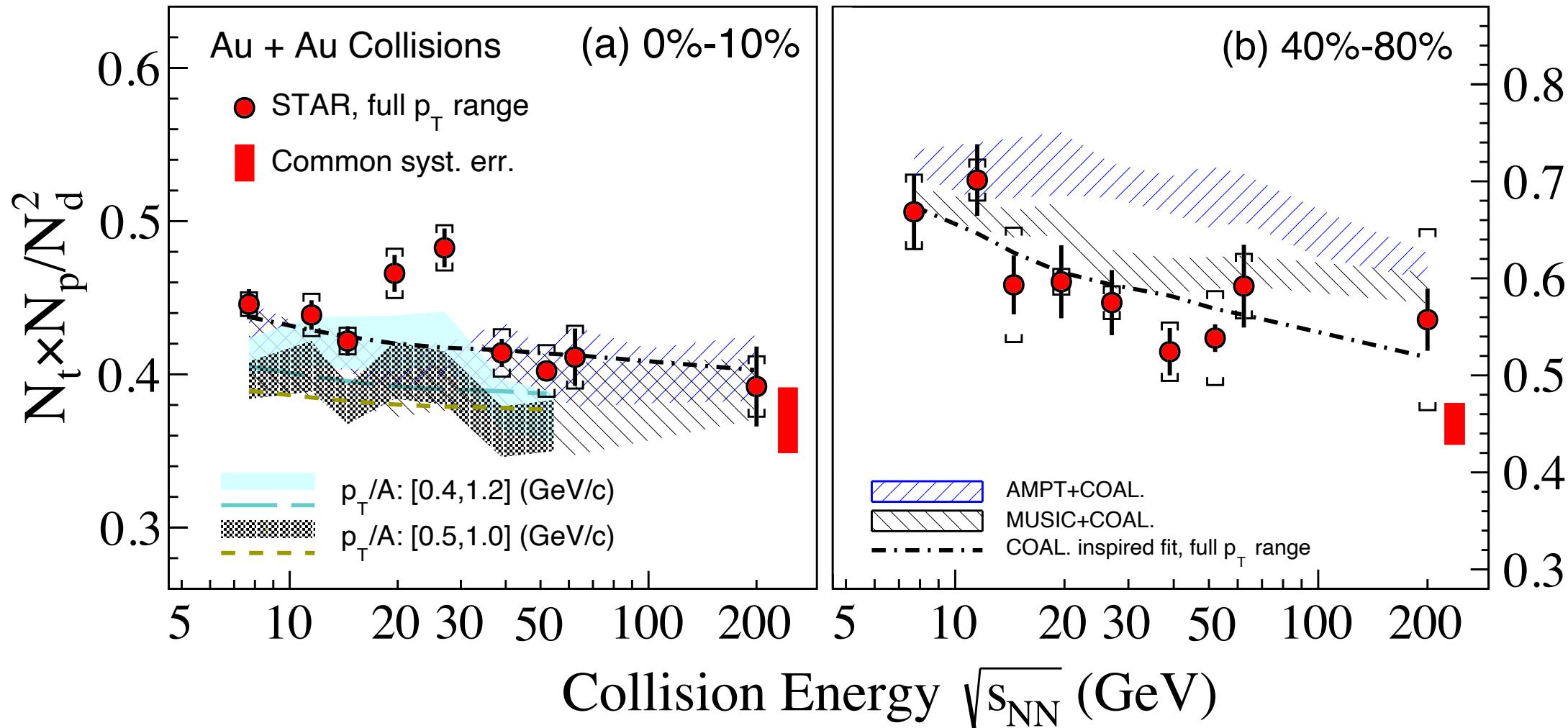
Coalescence wavefunction
overlap between nucleus
and nucleons



Possible sign of Density Fluctuation

4 σ effect, BES-II data x10 statistics

STAR, Phys.Rev.Lett. 130 (2023) 202301

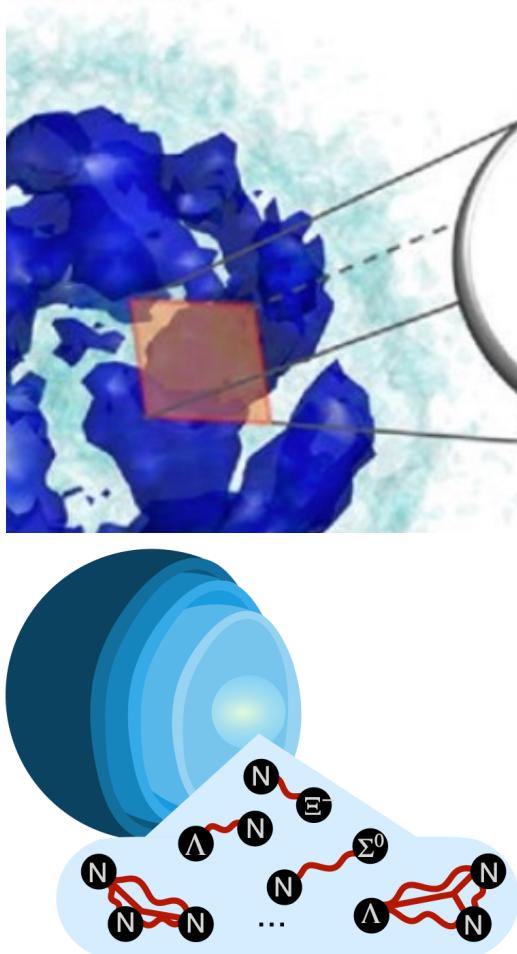


Few-body correlations at RHIC and LHC

Collision system size: 1-10fm

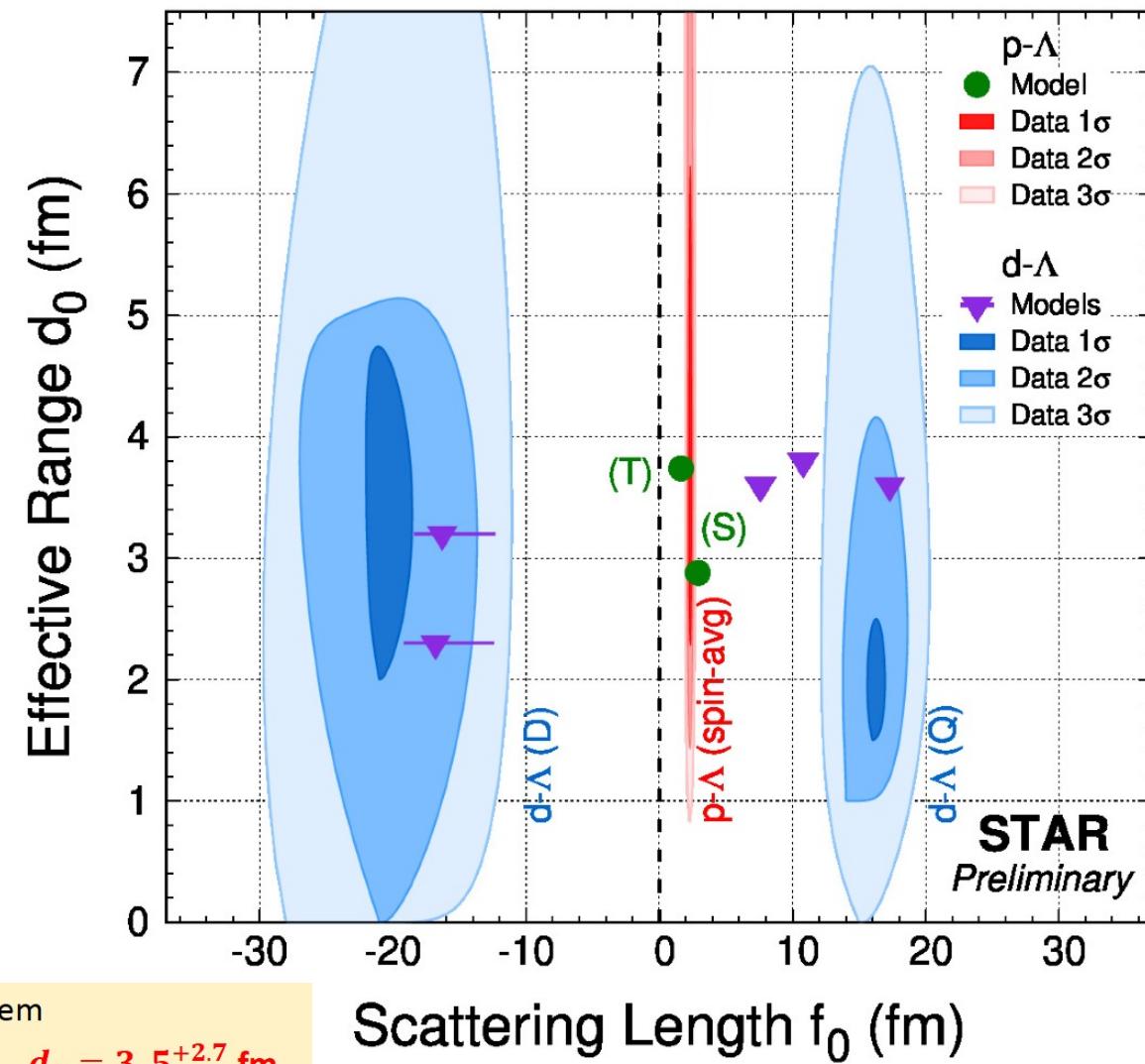
Scattering length: ~fm

d-p, n-p-p, p-p-p, p-p- Λ , d- Λ

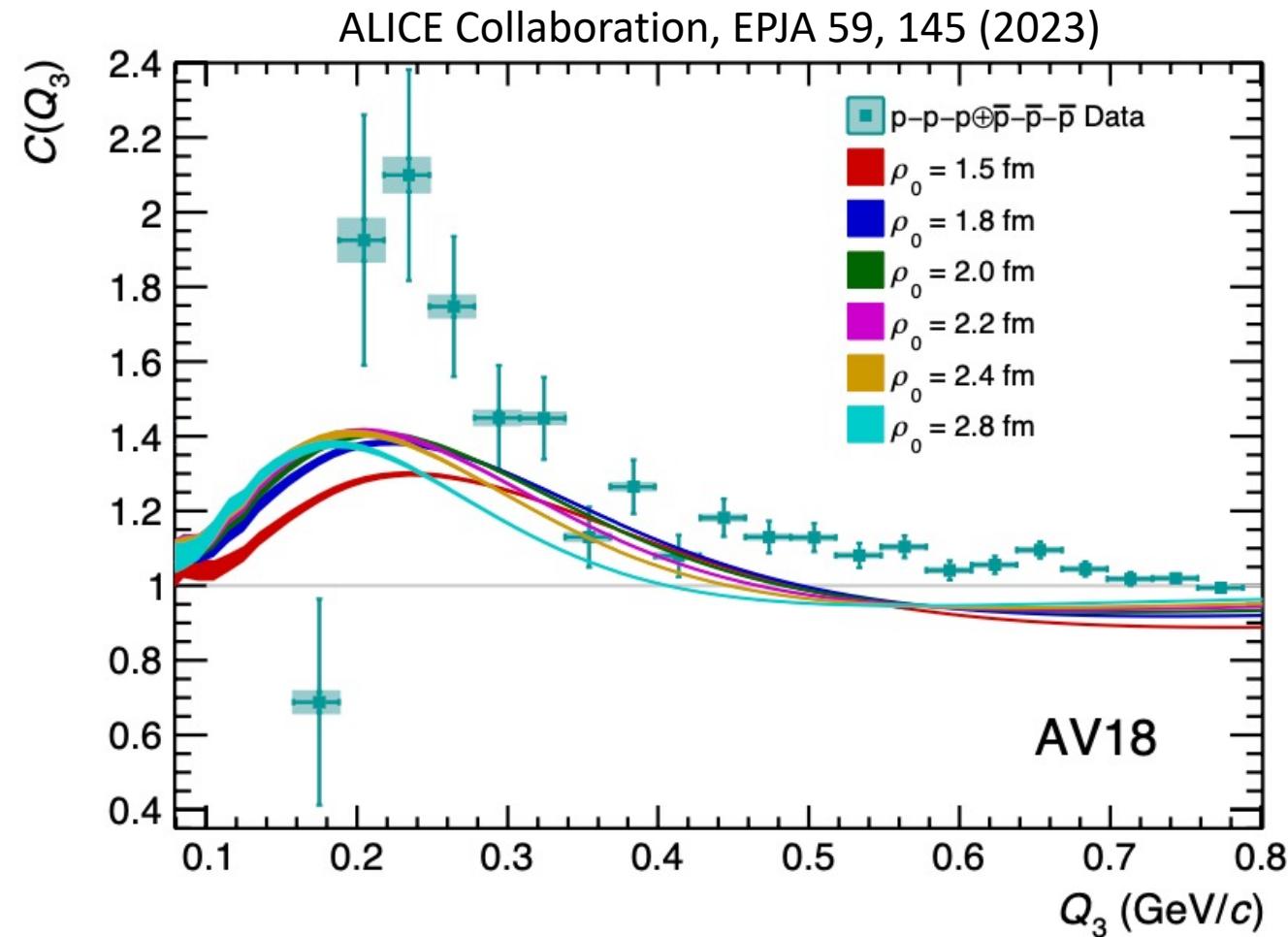
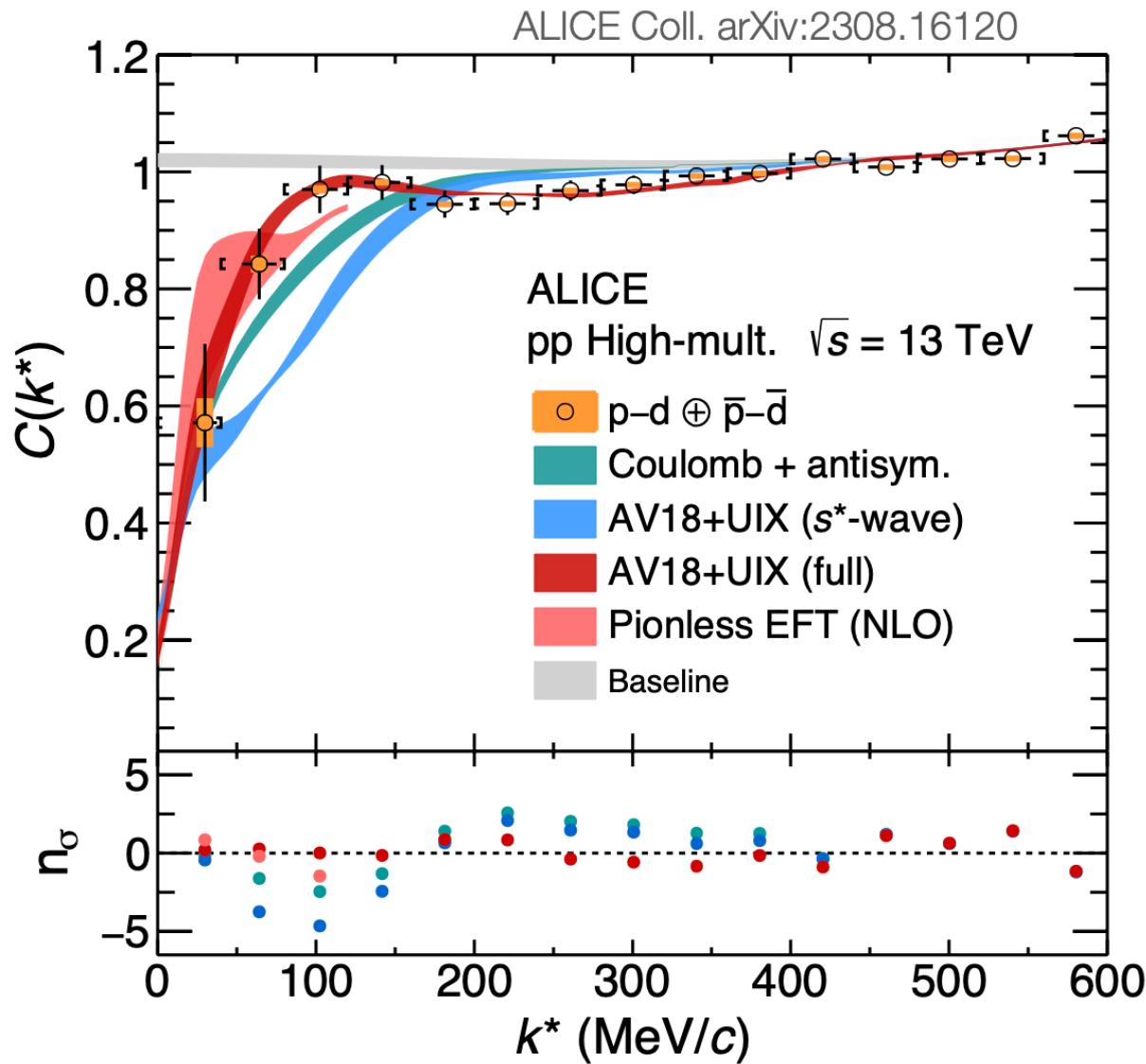


D. Lonardoni et al, PRL 114, 092301 (2015)
J. Schaffner-Bielich et al, NPA 835 (2010)

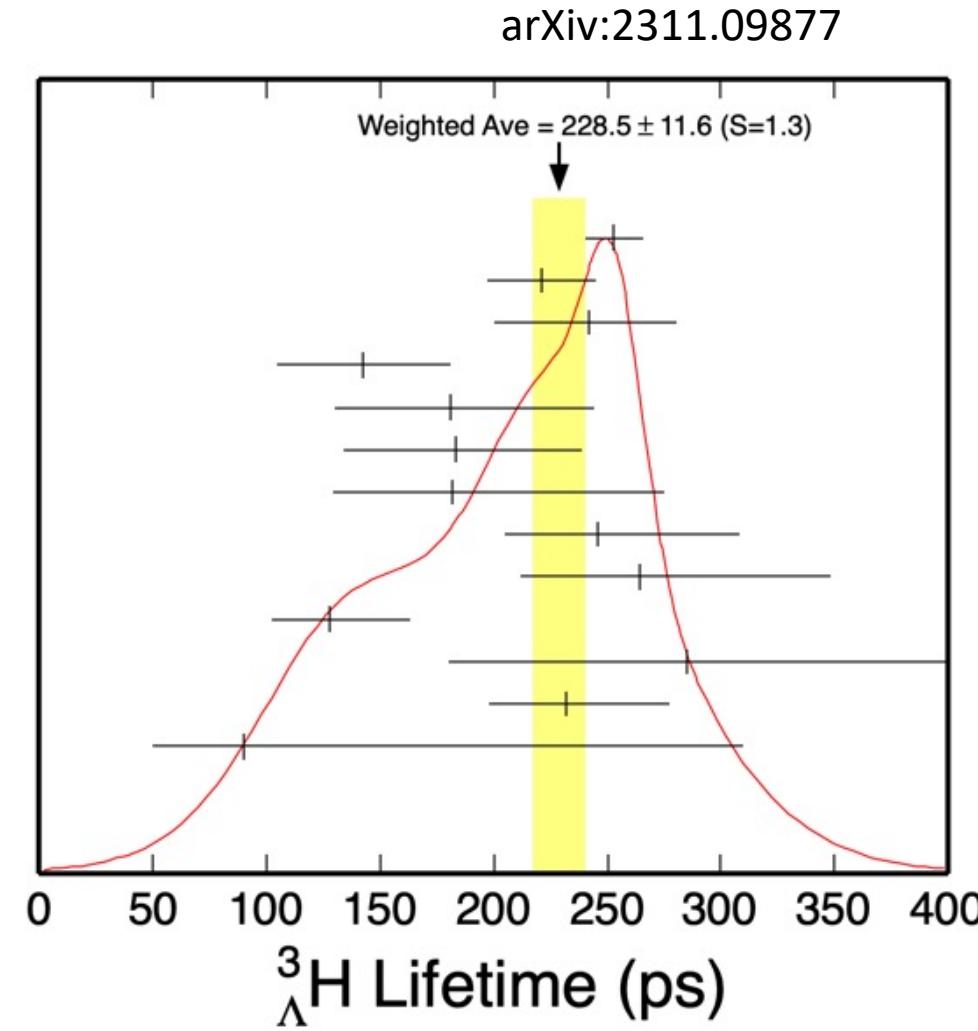
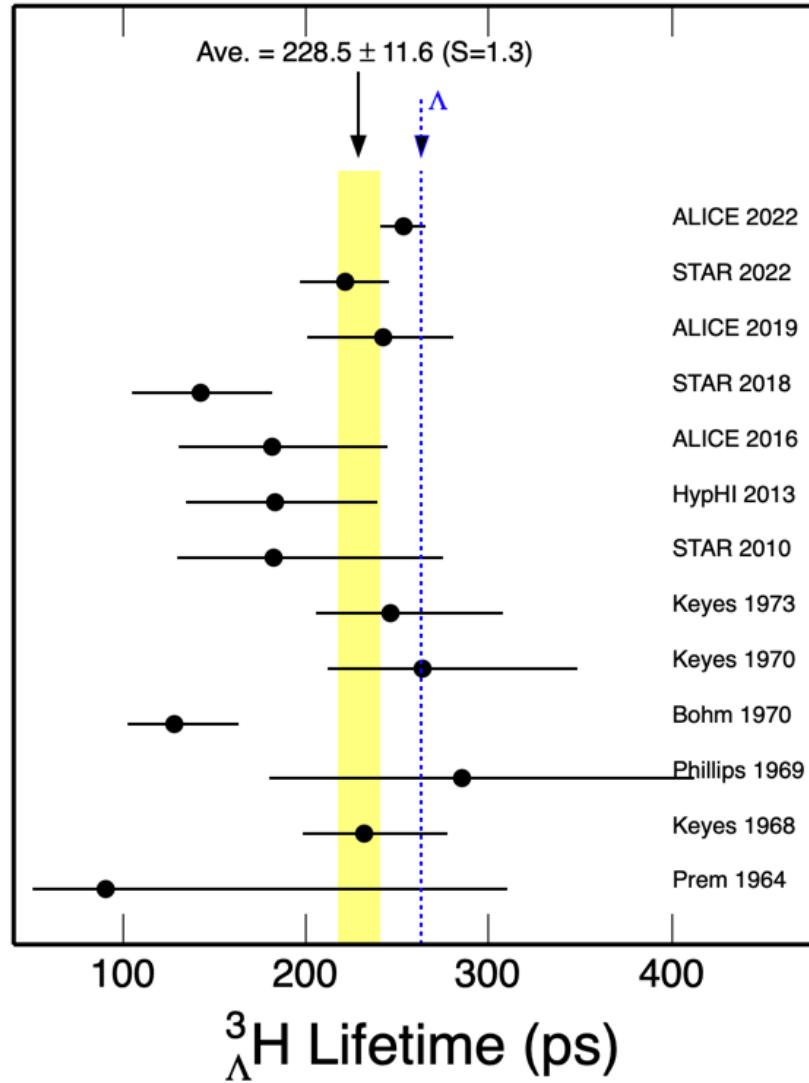
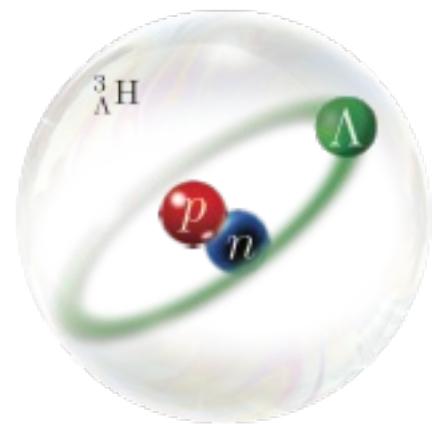
- ❖ Spin-avg for f_0 & d_0 p- Λ system
 $f_0 = 2.32^{+0.12}_{-0.11}$ fm $d_0 = 3.5^{+2.7}_{-1.3}$ fm
- ❖ Successfully separate two spin states in d- Λ
 $f_0(D) = -20^{+3}_{-3}$ fm $d_0(D) = 3^{+2}_{-1}$ fm
 $f_0(Q) = 16^{+2}_{-1}$ fm $d_0(Q) = 2^{+1}_{-1}$ fm



Three-nucleon correlations at LHC



$|B|=3$ hypertriton lifetime



Potential discrepancy?

arXiv:2311.09877

Simultaneous fit to all heavy-ion data

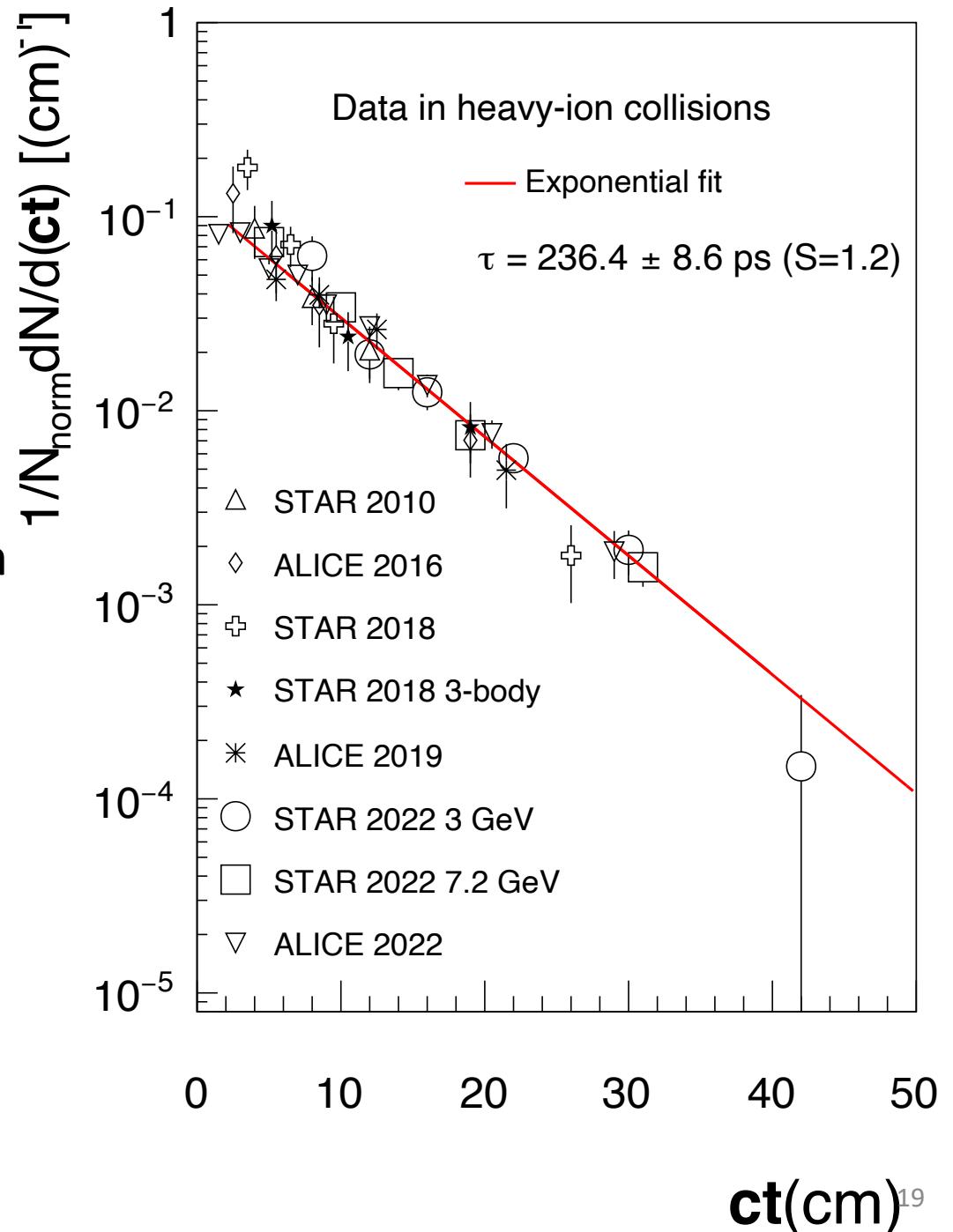
Scale yields to one common exponential function

Result consistent with other (average) methods

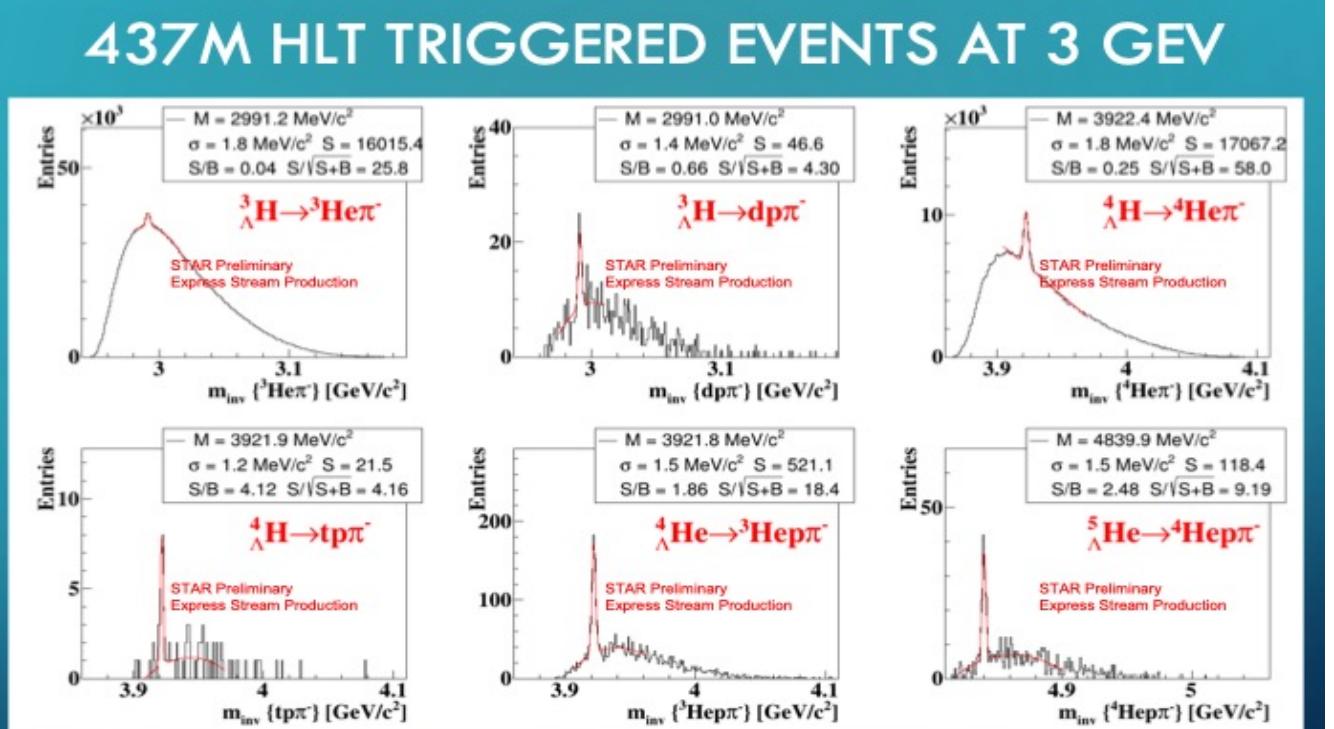
About 3σ smaller than Lambda lifetime

STAR 2018 first $c\tau$ point appears high

ALICE 2022 first $c\tau$ point appears low



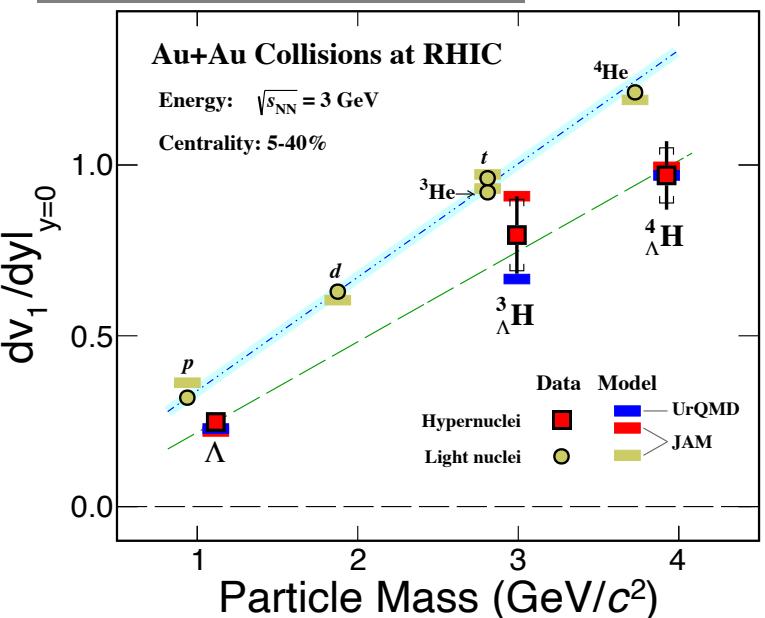
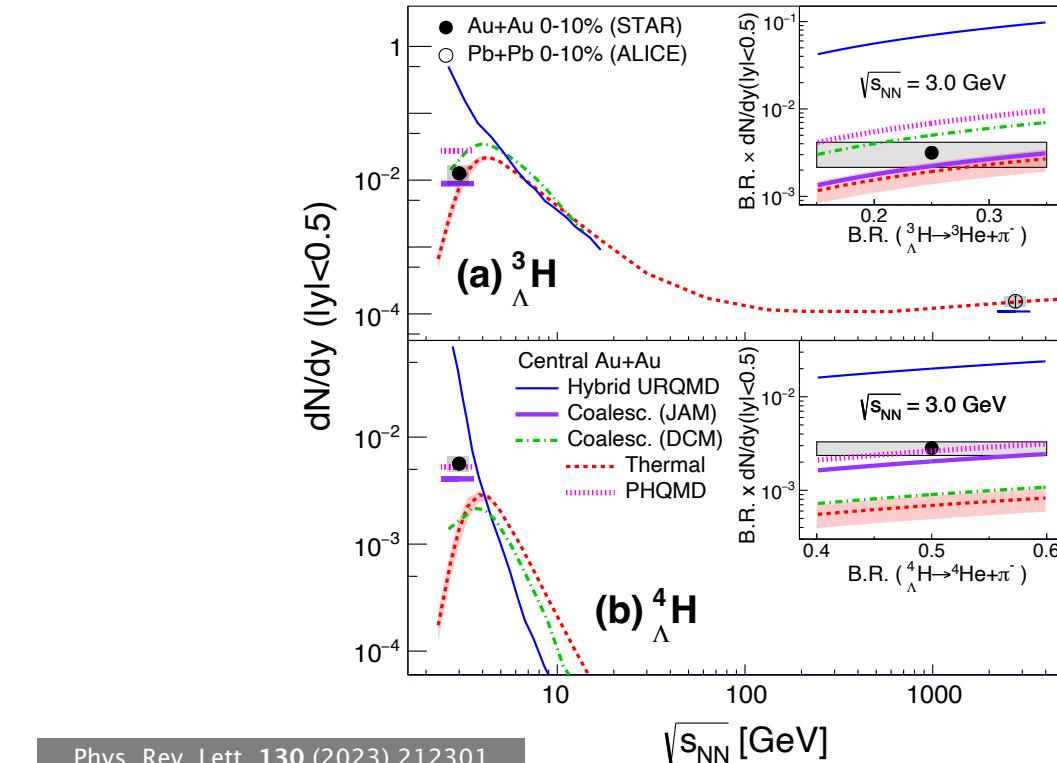
A zoo of hypernucleus measurements



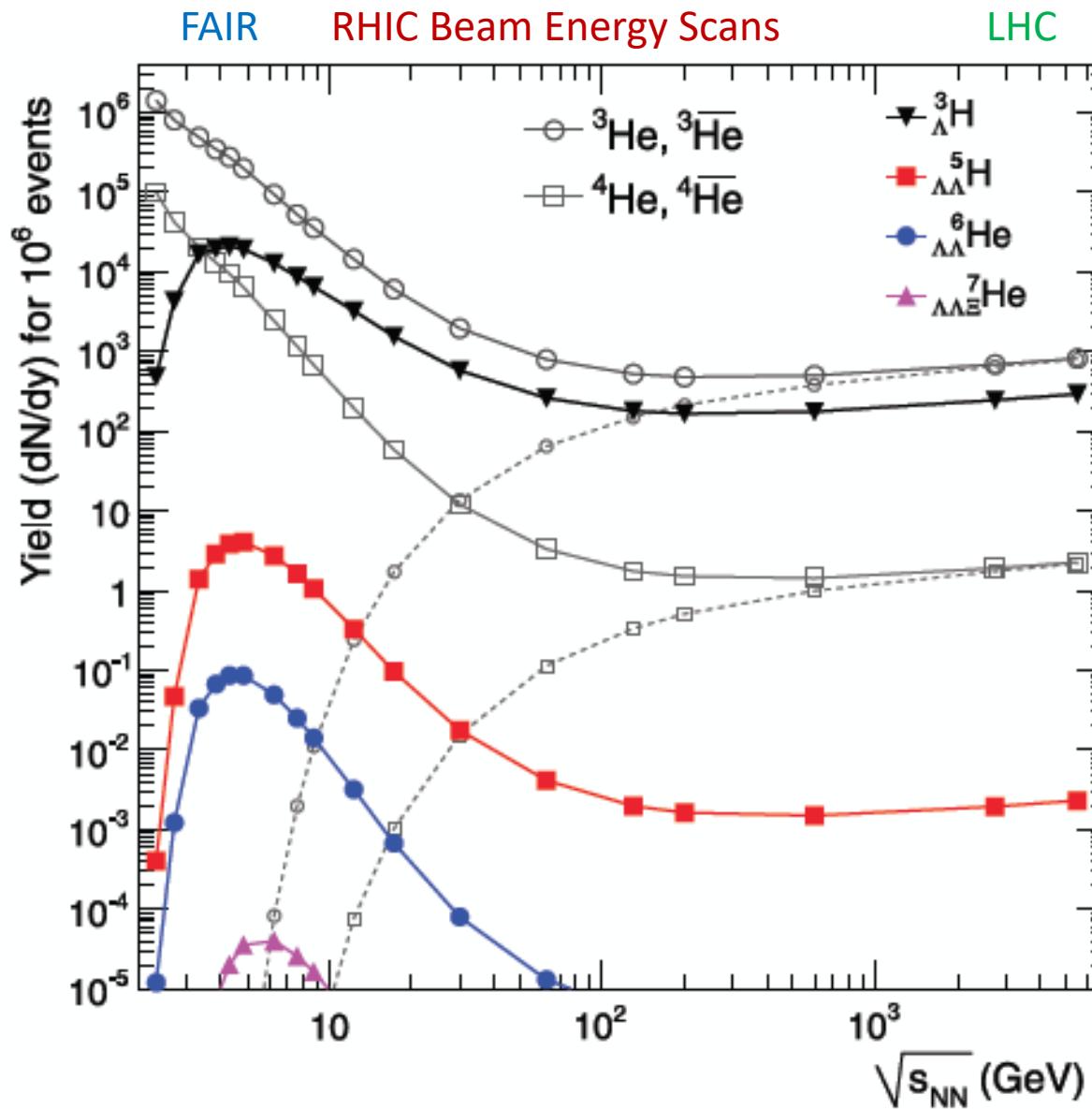
- With increased beam collision intensity in the Fixed Target mode HLT farm had not enough capacities to process all collected data online.

Therefore a trigger on He has been introduced to enhance hypernuclei.

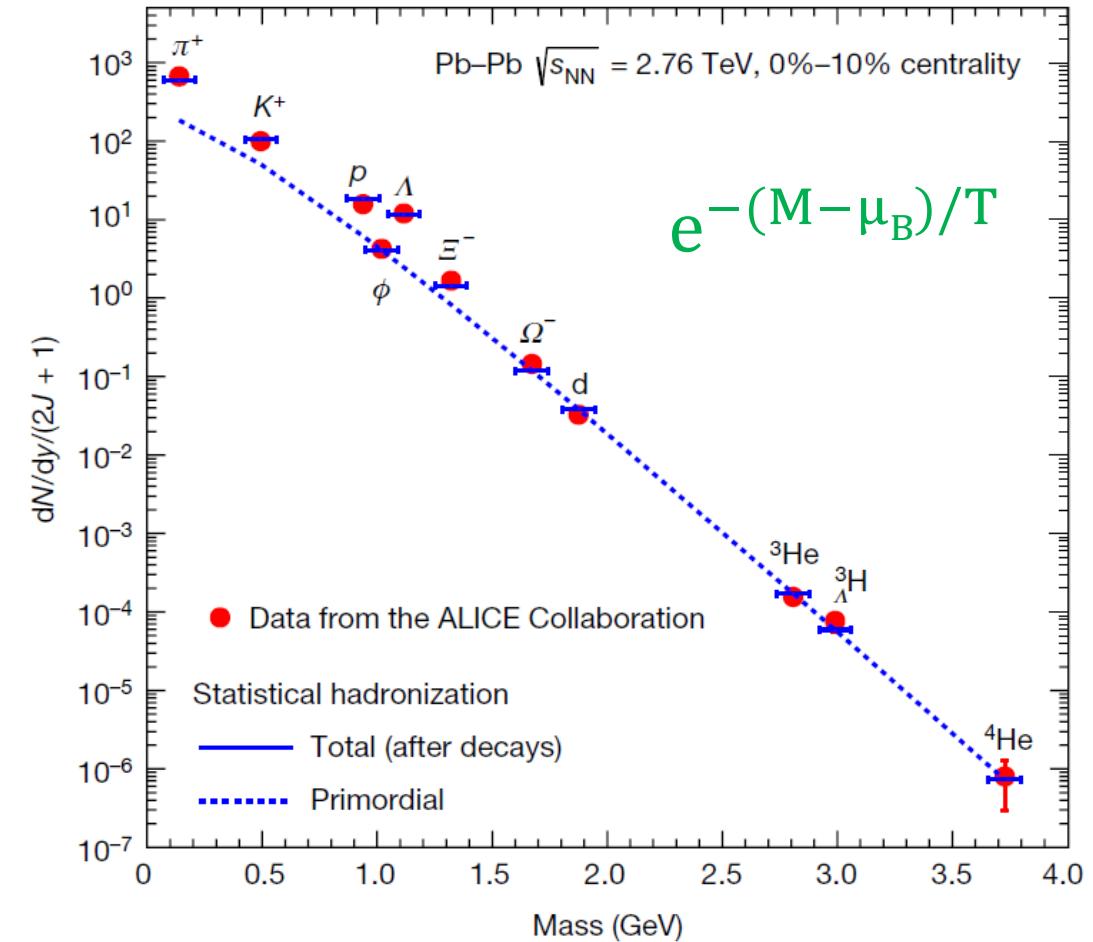
The collected statistics is enough to measure yields, lifetimes and spectra of these hypernuclei



Statistical Hadronization: powerful projections



Andronic et al., Nature 561, 321–330 (2018)



Search for heavy antimatter and baryon objects

Rapid Communication

Strangelet search

B. I. Abelev *et al.* (STAR Collaboration)
Phys. Rev. C **76**, 011901(R) – Published 24 April 2011

nature

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Published: 24 April 2011

Observation of the antimatter helium-4 nucleus

The STAR Collaboration

Nature **473**, 353–356 (2011) | Cite this article

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RESEARCH ARTICLE

Observation of an Antimatter Hypernucleus

THE STAR COLLABORATION, B. I. ABELEV, M. M. AGGARWAL, Z. AHAMMED, A. V. ALAKHVERDYANTS, I. ALEKSEEV, B. D. ANDERSON, D. ARK

Y. ZOULKARNEEVA +382 authors

Authors Info & Affiliations

SCIENCE • 4 Mar 2010 • Vol 328, Issue 5974 • pp. 58–62 • DOI: 10.1126/science.1183980

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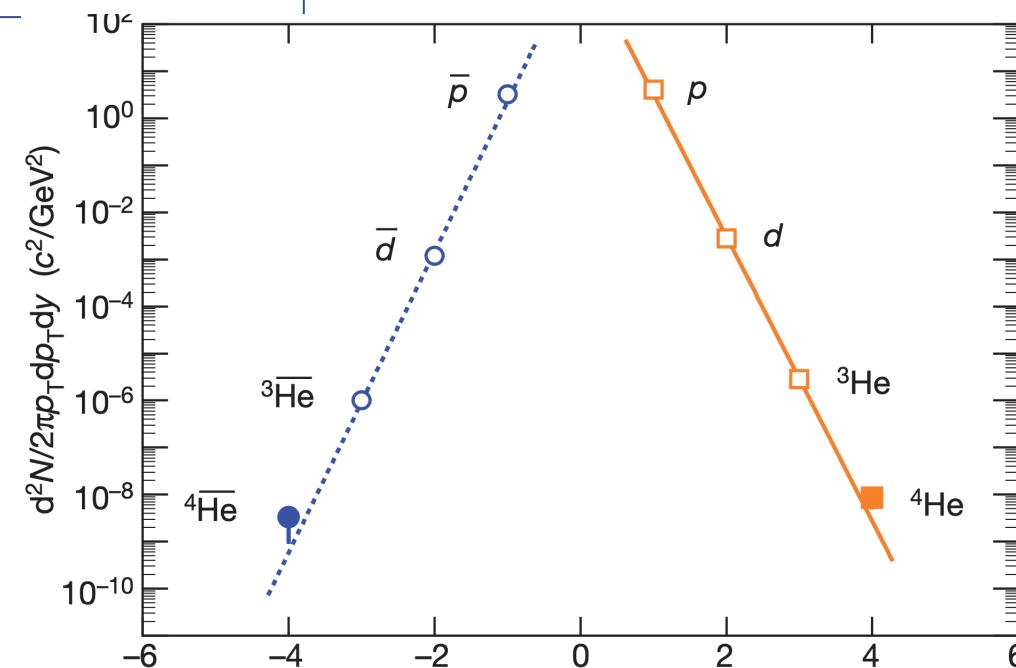
nature > articles > article

Article | Published: 21 August 2024

Observation of the antimatter hypernucleus $\bar{\Lambda}$ ${}^4\bar{H}$

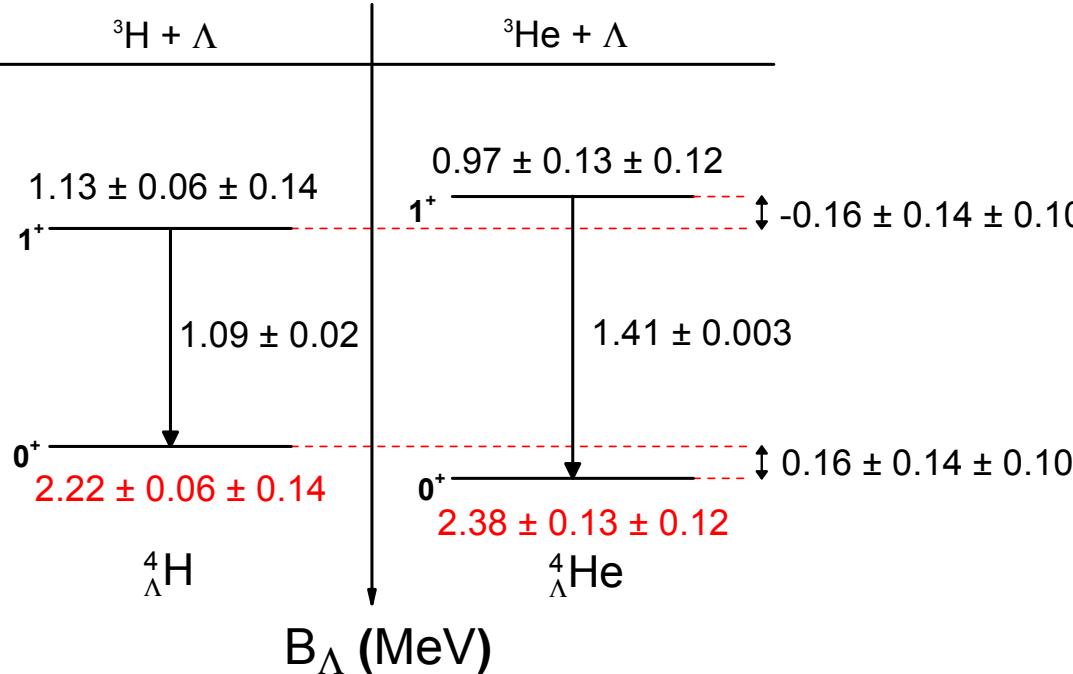
STAR Collaboration

Nature **632**, 1026–1031 (2024) | Cite this article



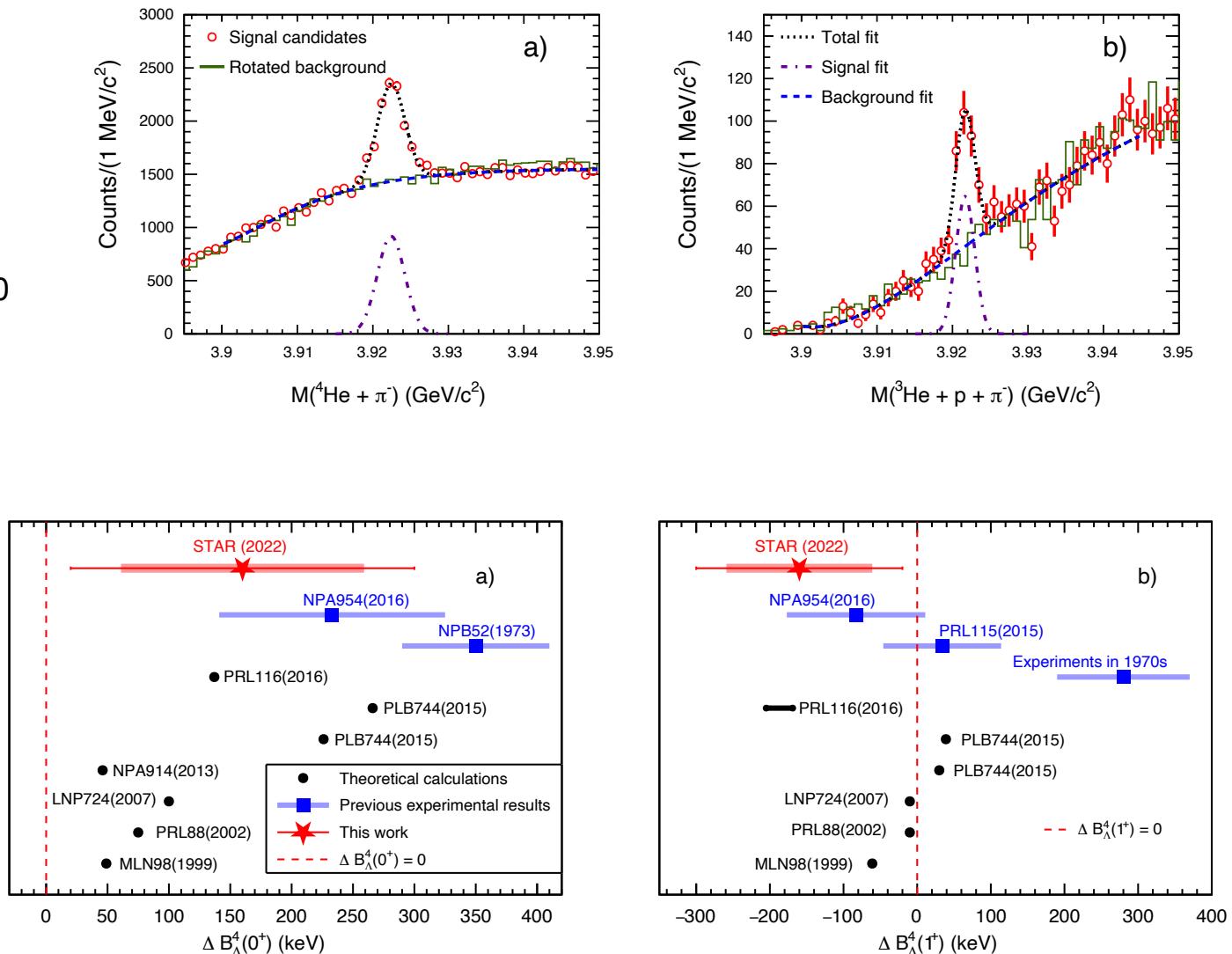
Charge Symmetry Breaking in B=4 hypernuclei

STAR, Phys. Lett. B 834 (2022) 137449



A puzzling CSB (70s):
both 0^+ and 1^+ large and positive ΔB

New measurements:
small and symmetric ΔB





Discovery potential at EIC

xzb

LBL Heavy-Ion Tea Seminar
05/2010

Heavy-flavor states

<http://belle.kek.jp/belle/talks/moriondQCD10/pakhlov.ppt>

Many (>10) states poorly consistent with quark model

State	M (MeV)	Γ (MeV)	J^{PC}	Decay Modes	Production Modes
$Y_s(2175)$	2175 ± 8	58 ± 26	1^{--}	$\phi f_0(980)$	e^+e^- (ISR) $J/\psi \rightarrow \eta Y_s(2175)$
$X(3872)$	3871.4 ± 0.6	< 2.3	1^{++}	$\pi^+\pi^- J/\psi, \gamma J/\psi, DD^*$	$B \rightarrow KX(3872), p\bar{p}$
$X(3915)$	3914 ± 4	23 ± 9	$0/2^{++}$	$\omega J/\psi$	$\gamma\gamma \rightarrow X(3915)$
$Z(3930)$	3929 ± 5	29 ± 10	2^{++}	$D\bar{D}$	$\gamma\gamma \rightarrow Z(3940)$
$X(3940)$	3942 ± 9	37 ± 17	$0^{?+}$	$D\bar{D}^*$ (not $D\bar{D}$ or $\omega J/\psi$)	$e^+e^- \rightarrow J/\psi X(3940)$
$Y(3940)$	3943 ± 17	87 ± 34	$?^{?+}$	$\omega J/\psi$ (not $D\bar{D}^*$)	$B \rightarrow KY(3940)$
$Y(4008)$	4008^{+82}_{-49}	226^{+97}_{-80}	1^{--}	$\pi^+\pi^- J/\psi$	e^+e^- (ISR)
$X(4160)$	4156 ± 29	139^{+113}_{-65}	$0^{?+}$	$D^*\bar{D}^*$ (not $D\bar{D}$)	$e^+e^- \rightarrow J/\psi X(4160)$
$Y(4260)$	4264 ± 12	83 ± 22	1^{--}	$\pi^+\pi^- J/\psi$	e^+e^- (ISR)
$Y(4350)$	4361 ± 13	74 ± 18	1^{--}	$\pi^+\pi^- \psi'$	e^+e^- (ISR)
$X(4630)$	4634^{+9}_{-11}	92^{+41}_{-32}	1^{--}	$\Lambda_c^+ \Lambda_c^-$	e^+e^- (ISR)
$Y(4660)$	4664 ± 12	48 ± 15	1^{--}	$\pi^+\pi^- \psi'$	e^+e^- (ISR)
$Z(4050)$	4051^{+24}_{-23}	82^{+51}_{-29}	$?$	$\pi^\pm \chi_{c1}$	$B \rightarrow KZ^\pm(4050)$
$Z(4250)$	4248^{+185}_{-45}	177^{+320}_{-72}	$?$	$\pi^\pm \chi_{c1}$	$B \rightarrow KZ^\pm(4250)$
$Z(4430)$	4433 ± 5	45^{+35}_{-18}	$?$	$\pi^\pm \psi'$	$B \rightarrow KZ^\pm(4430)$
$Y_b(10890)$	$10,890 \pm 3$	55 ± 9	1^{--}	$\pi^+\pi^- \Upsilon(1, 2, 3S)$	$e^+e^- \rightarrow Y_b$

observed last 6 years by B-factories

How about baryon states?

Heavy-flavor hypernuclei

- Predicted to exist (70's)
- Cannot be produced in pp, ep collisions
- Cannot be detected in fixed target experiment
- Only solution: **EIC**
- EIC enough energy for charm and bottom hypernuclei
- Vertex detector at Fragmentation region
Displace vertex: 3cm

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Phys. Rev. Lett. 39, 1506–1509 (1977)

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Possibility of Charmed Hypernuclei

Abstract References Citing Articles (17) Page Images

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C. B. Dover and S. H. Kahana
Brookhaven National Laboratory, Upton, New York 11973



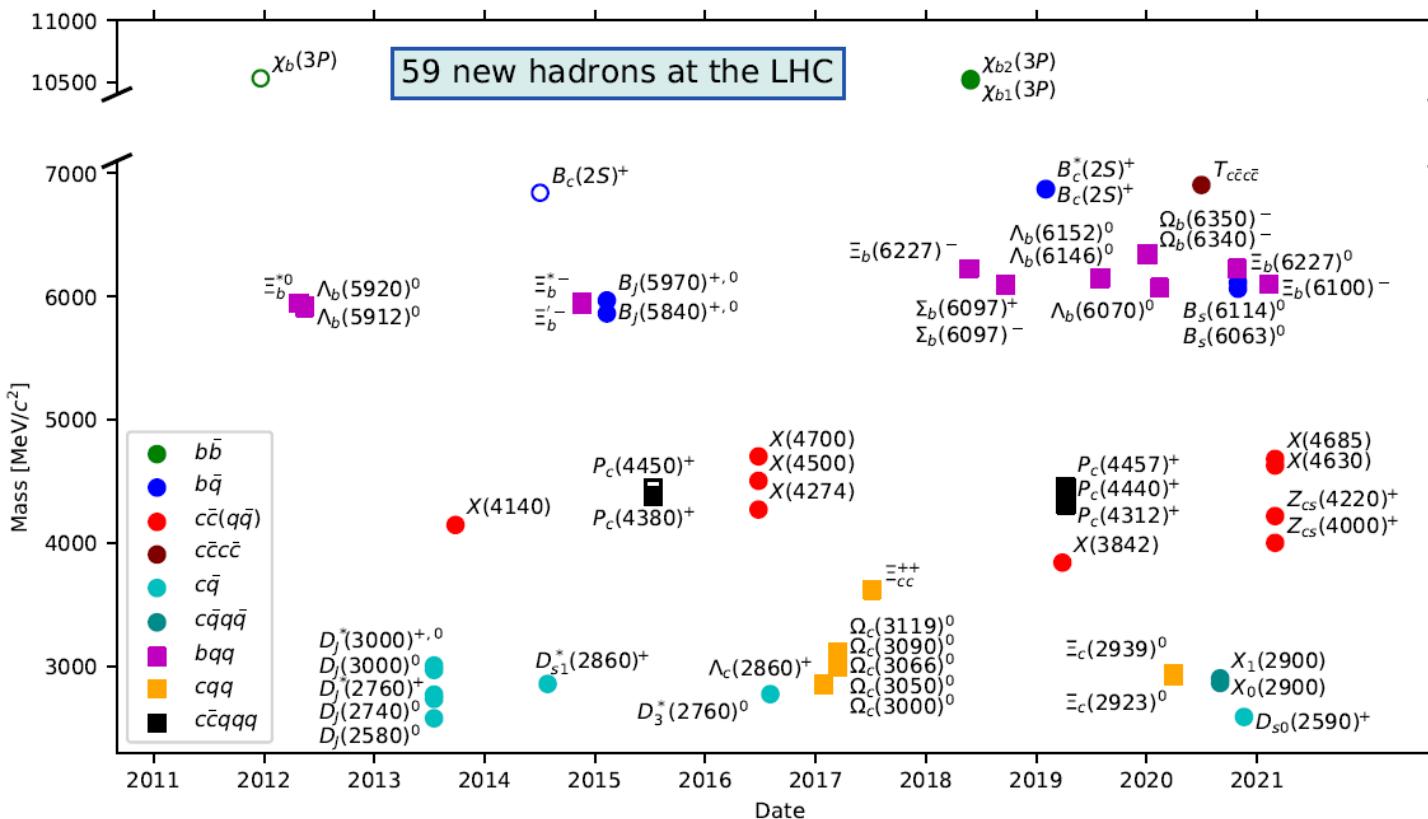
Discovery potential at EIC

xzb

LBL Heavy-Ion Tea Seminar
05/2010

Heavy-flavor states

<http://belle.kek.jp/belle/talks/moriondQCD10/pakhlov.ppt>



How about baryon states?

Heavy-flavor hypernuclei

- Predicted to exist (70's)
- Cannot be produced in pp, ep collisions
- Cannot be detected in fixed target experiment
- Only solution: **EIC**
- EIC enough energy for charm and bottom hypernuclei
- Vertex detector at Fragmentation region
Displace vertex: 3cm

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Phys. Rev. Lett. 39, 1506–1509 (1977)

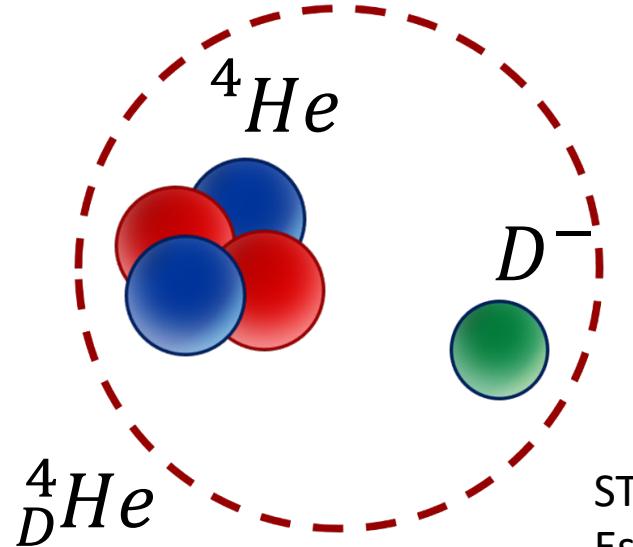
Possibility of Charmed Hypernuclei

Abstract References Citing Articles (17) Page Images

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Search for Stable Charmed Mesic Nucleus D^- ^4He in Heavy-Ion and EIC



Stable and existence due to Coulomb force

PYTHIA: $D^-/n \sim 5 \times 10^{-4}$ p+p collisions
at AGS and RHIC forward kinematics

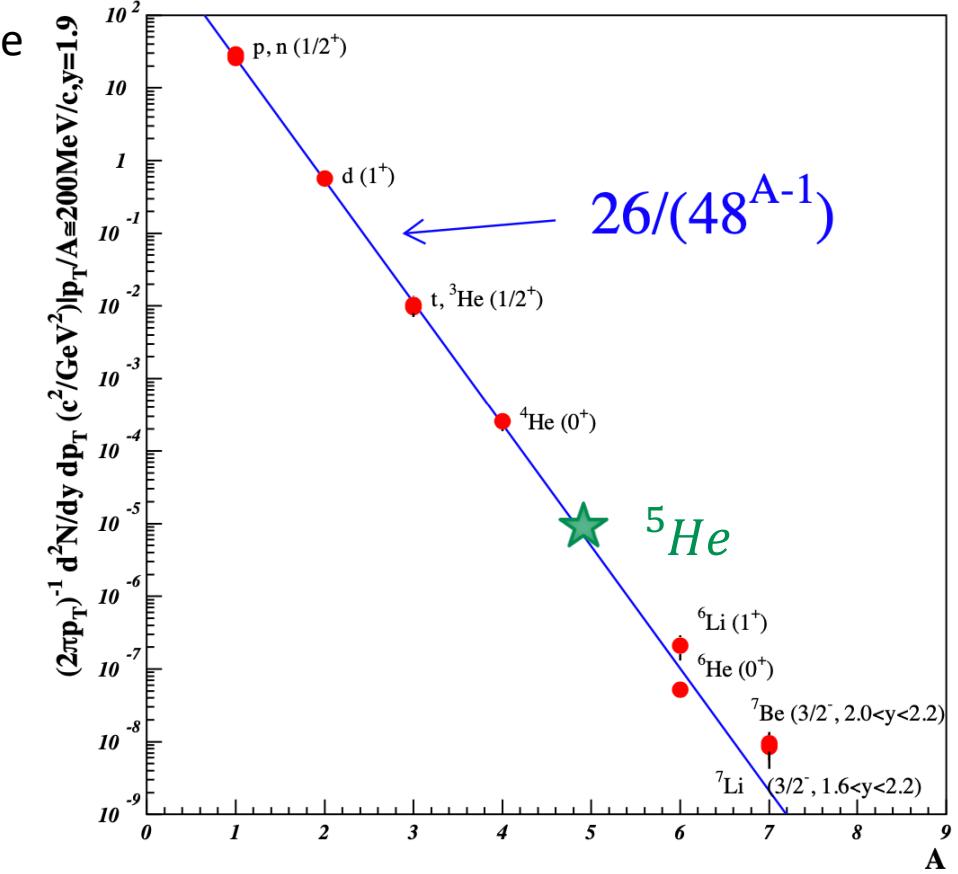
D^- ^4He yield 10^{-8} per collision

STAR@RHIC:
Estimate 1×10^5 /year in forward acceptance
But without vertex detector

CBM@FAIR high baryon, good vertex
LHCb@LHC forward with good vertex

EIC ion forward direction:
clean environment with good vertex
Nuclear cluster

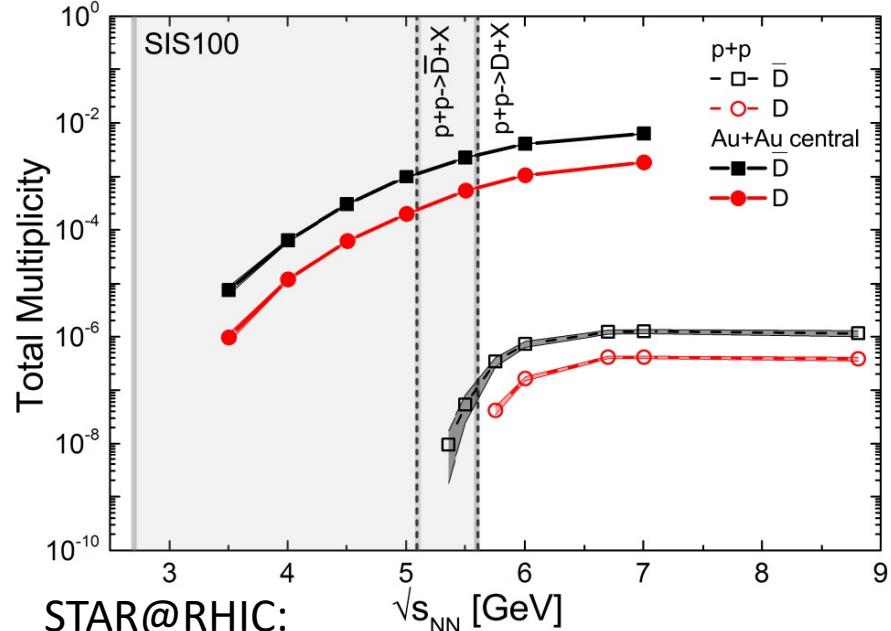
Zhangbu Xu (BNL)
Cheng-Wei Lin, Yi Yang (NCKU)
DNP (2022), EMMI (2023)



E864@AGS Phys. Rev. C **61**, 064908

Search for Stable Charmed Mesic Nucleus D^- 4He with CBM

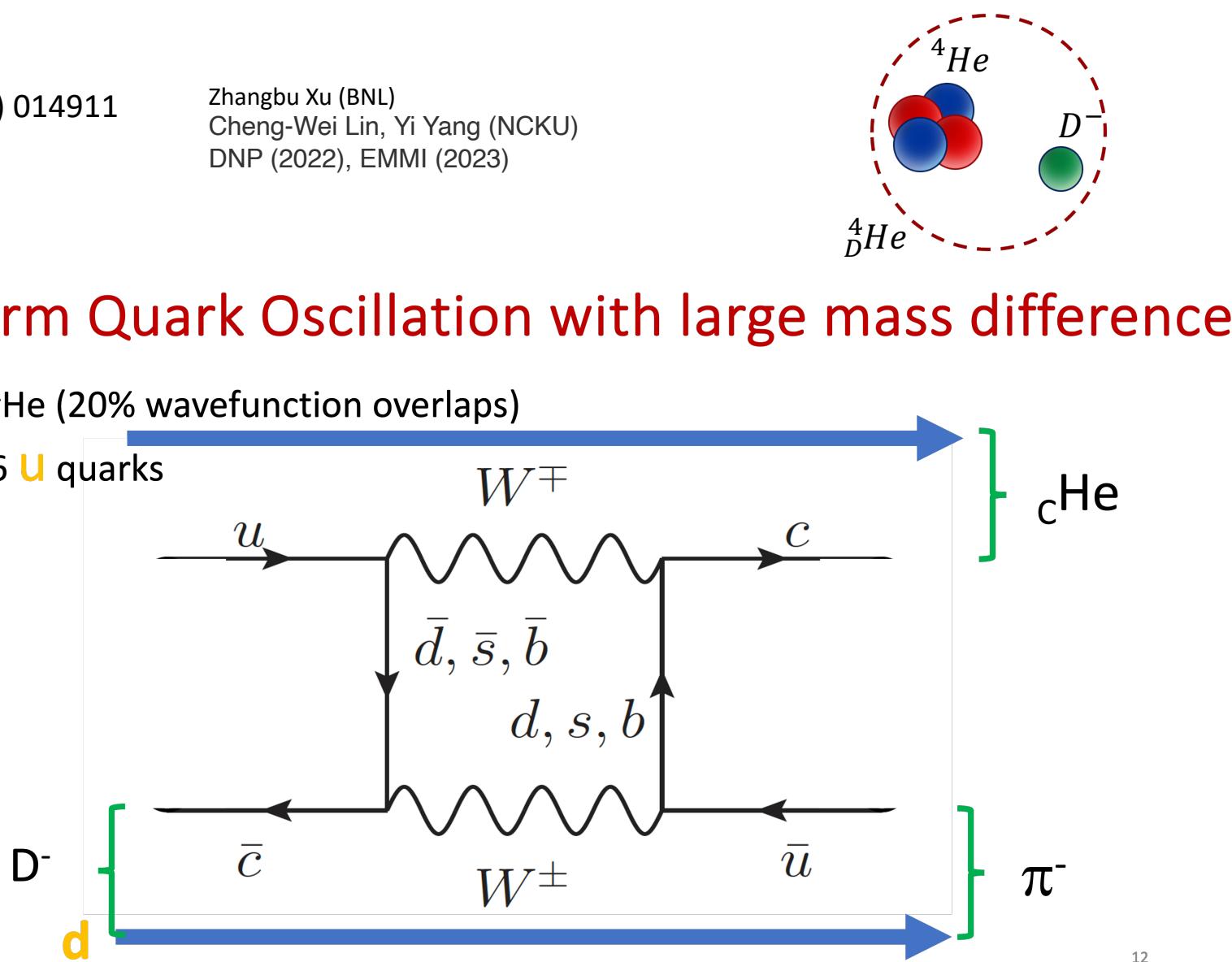
J. Steinheimer, A. Botvina, M. Bleicher, PRC 95 (2017) 014911



CBM@FAIR high baryon, good vertex
LHCb@LHC forward with good vertex

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Conclusions

- Discovery of the heaviest antimatter nuclear cluster (hyperhydrogen 4)
- Continue to improve our measurements on hypernuclear lifetime and binding energy (CSB)
- Use nuclear yields to study production mechanism, quantum wavefunction overlap: thermal vs coalescence model
- Use nuclear yield ratios as a sensitive probe of nucleon density fluctuation
- Baryon number is a strictly conserved quantum number, keeps the Universe as is; use baryon transport to study its tracer
- Explore other signatures
- Two and three-nucleon correlations sensitive to scattering length and bound states
- Charmed hypernuclei (EIC, LHC, FAIR)