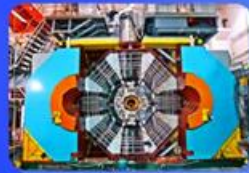


Prediction for Neutrino Interaction Cross Sections: **From Low to High Energies**

WWW.IHEP.CAS.CN



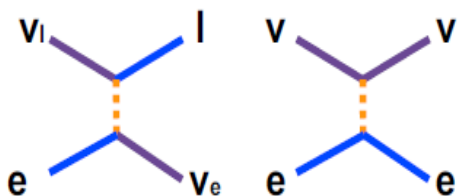
Yu-Feng Li (李玉峰)

Institute of High Energy Physics &

University of Chinese Academy of Sciences, Beijing

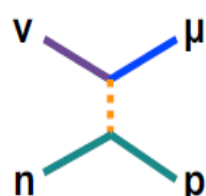
***The 23rd International Conference on Few-Body Problems in Physics,
Beijing, 27th Sep. 2024***

Neutrino interactions from low to high energies

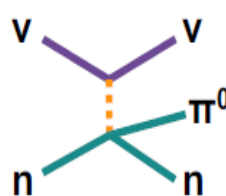


inverse muon (tau)
decay

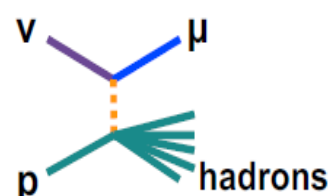
elastic electron
scattering



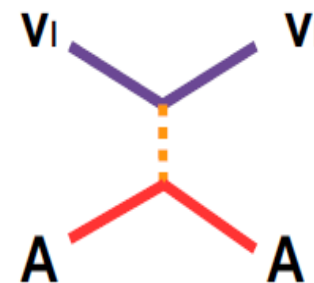
(quasi) - elastic
nucleon scattering



nuclear excitation and
resonant production



Deep inelastic scattering
and jet production



➤ Neutrino-lepton interactions:

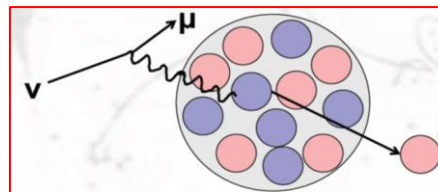
$$\bar{\nu}_\alpha + e^- \rightarrow \bar{\nu}_\alpha + e^- .$$

➤ Neutrino-nucleon interactions

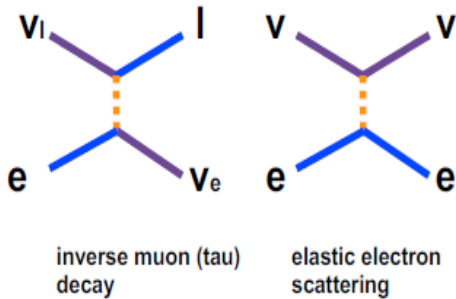
$$\nu_\ell + n \rightarrow p + \ell^-$$

$$\bar{\nu}_\ell + p \rightarrow n + \ell^+$$

➤ Neutrino-nucleus interactions



Neutrino-lepton interactions



➤ Neutrino-lepton interactions:

Pure leptonic process and easy to calculate (at tree level)

See radiative corrections in Bacall et al.,
PRD 51 (1995) 6146-6158

Accelerator $\nu\mu$: Observation of NC (Gargamelle, 1973)
Measurement of weak mixing angle (CHARM-II, 1994)

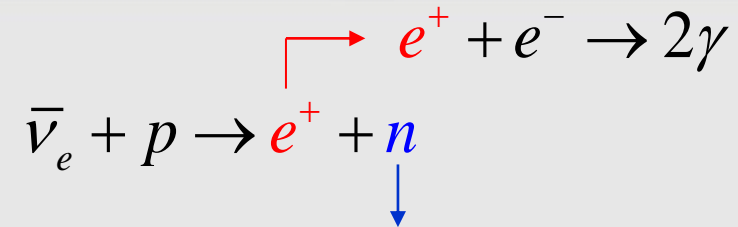
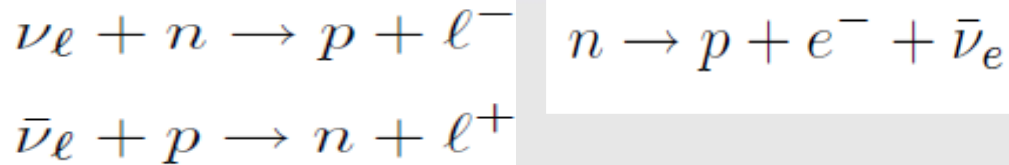
Solar neutrinos: Super-Kamiokande, Borexino, JUNO etc.
Dark Matter Direct Detection experiments

New physics: neutrino magnetic moment
GEMMA: $2.9 \times 10^{-11} \mu_B$ [Rev.Mod.Phys. 87 (2015) 531]

XENON-nT: $6.4 \times 10^{-12} \mu_B$ (2207.11330)

Neutrino-nucleon interactions

Prompt signal



Capture on H or Gd,
Delayed signal, 2.2, 8 MeV

- Famous inverse beta decay on free proton (in **Hydrogen rich** detectors)

- Hadron weak current: **induced currents**

$$\bar{u}_u(p_u) \gamma^\rho (1 - \gamma^5) u_d(p_d) \rightarrow \langle p(p_p) | h_W^\rho(0) | n(p_n) \rangle$$

- Isospin symmetry

- Correlated with free neutron decay

Dedicated calculations in :

Vogel & Beacom, 1999

Strumia & Vissani, 2003

Ricciardi, Vignaroli, Vissani, 2022

Radiative correction:

Kurylov, Ramsey-Musolf, Vogel, 2003

Uncertainty as small as **~0.2%**

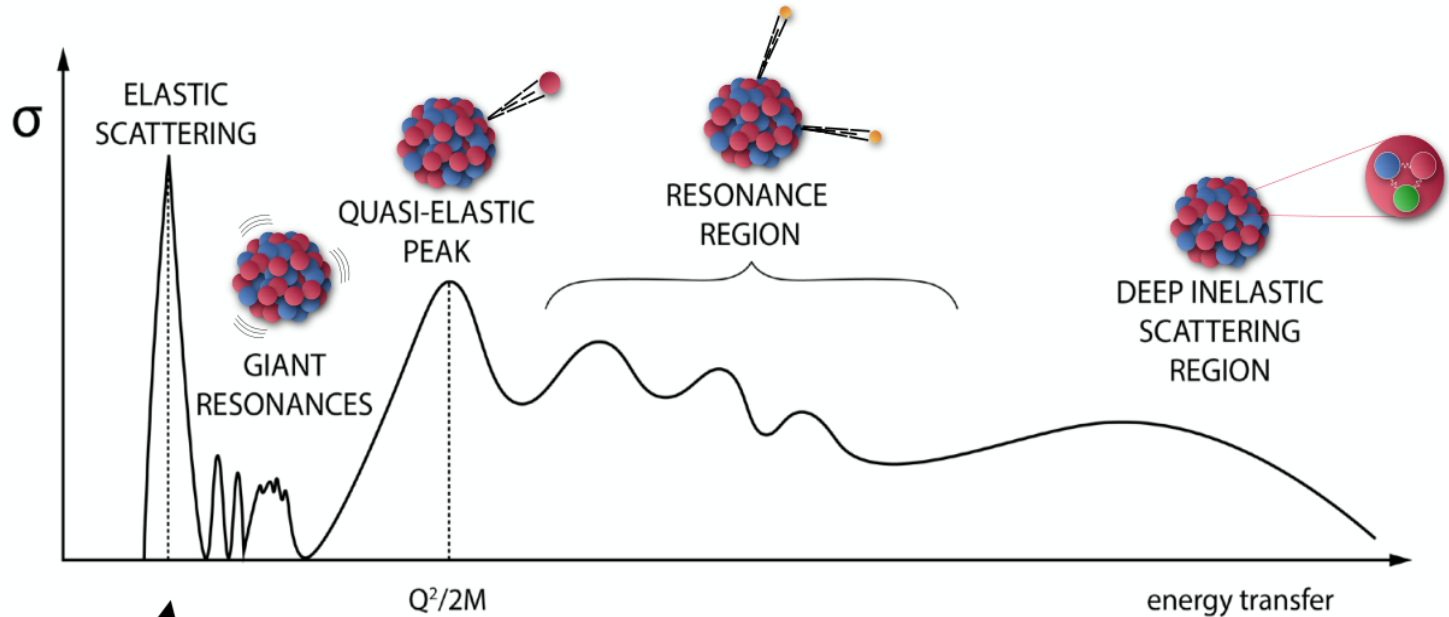
May affect the reactor
antineutrino anomaly

Giunti, YFL, Ternes, Xin PLB (2022)

$$\langle p(p_p) | v_W^\rho(0) | n(p_n) \rangle = \bar{u}_p(p_p) \left[\gamma^\rho F_1(Q^2) + \frac{i \sigma^{\rho\eta} q_\eta}{2m_N} F_2(Q^2) + \frac{q^\rho}{m_N} F_3(Q^2) \right] u_n(p_n)$$

$$\langle p(p_p) | a_W^\rho(0) | n(p_n) \rangle = \bar{u}_p(p_p) \left[\gamma^\rho \gamma^5 G_A(Q^2) + \frac{q^\rho}{m_N} \gamma^5 G_P(Q^2) + \frac{p_p^\rho + p_n^\rho}{m_N} \gamma^5 G_3(Q^2) \right] u_n(p_n)$$

Neutrino-nucleus interactions



Inelastic scattering

Elastic scattering: CEvNS
(coherent elastic neutrino-nucleus scattering)

e.g. Supernovae
neutrinos

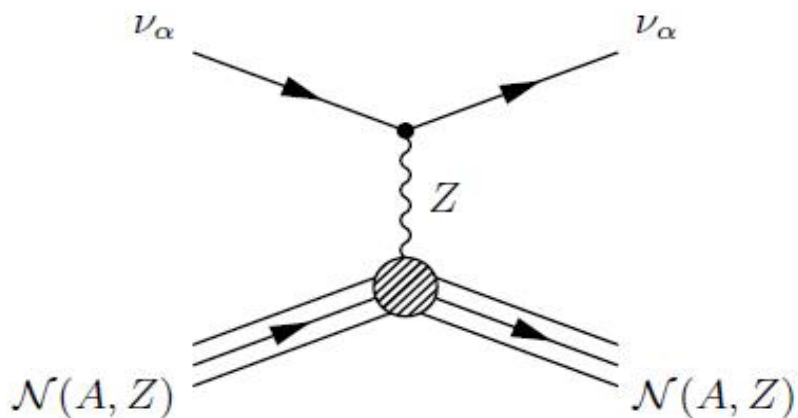
Long-baseline experiments
(DUNE, HyperK)

- Nuclear response at different energy transfer regions is crucial to make the predictions.

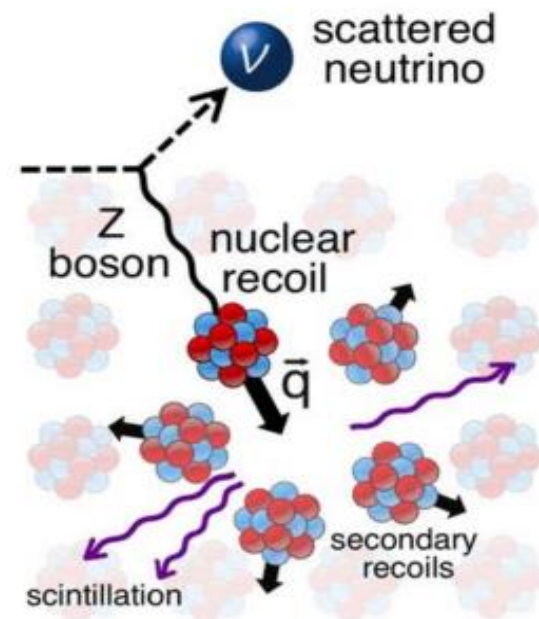
A: low energy NC: CEνNS

Coherent Elastic Neutrino-Nucleus Scattering

- ▶ $CE\nu NS$: pronounced “sevens”
- ▶ Weak Neutral-Current (NC) interaction:



- ▶ The nucleus $\mathcal{N}(A, Z)$ recoils as a whole!
- ▶ So what?



The CEνNS kinematics

$$|\vec{q}| R \lesssim 1$$

- ▶ Heavy target nucleus $\mathcal{N}(A, Z)$:

$$A \sim 100 \quad M \sim 100 \text{ GeV}$$

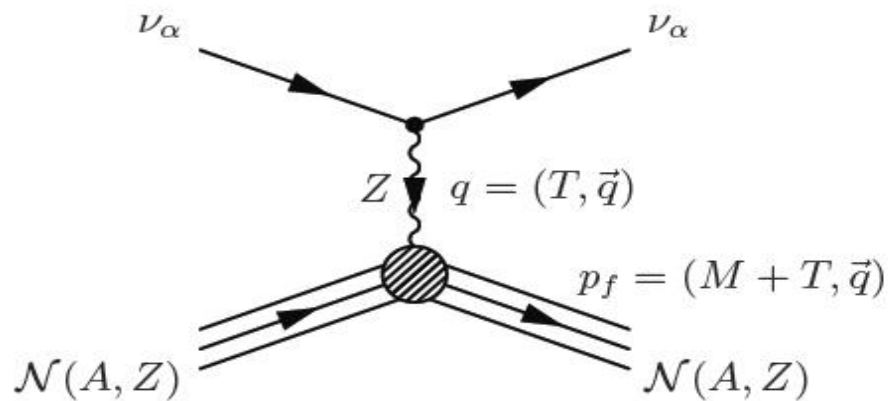
$$R \approx 1.2 A^{1/3} \text{ fm} \approx 5 \text{ fm}$$

- ▶ CEνNS for $|\vec{q}| \lesssim 40 \text{ MeV}$

- ▶ Non-Relativistic nuclear recoil:

$$|\vec{q}| \simeq \sqrt{2MT}$$

$$q^0 = T \leftarrow \text{Kinetic Energy}$$



- ▶ Observable nuclear recoil kinetic energy:

$$T \simeq \frac{|\vec{q}|^2}{2M} \lesssim 10 \text{ keV} \leftarrow \text{Very Small!}$$

The CEνNS Cross Section

Standard Model:
$$\frac{d\sigma_{\text{CE}\nu\text{NS}}}{dT}(E_\nu, T) = \frac{G_F^2 M}{4\pi} \left(1 - \frac{MT}{2E_\nu^2}\right) [Q_W(Q^2)]^2$$

- ▶ Weak charge of the nucleus \mathcal{N} :

$$|\vec{q}| = \sqrt{2MT}$$

$$Q_W(Q^2) = g_V^n N F_N(|\vec{q}|) + g_V^p Z F_Z(|\vec{q}|)$$

$$g_V^n = -\frac{1}{2} \quad g_V^p = \frac{1}{2} - 2 \sin^2 \vartheta_W(Q^2 \simeq 0) = 0.0227 \pm 0.0002$$

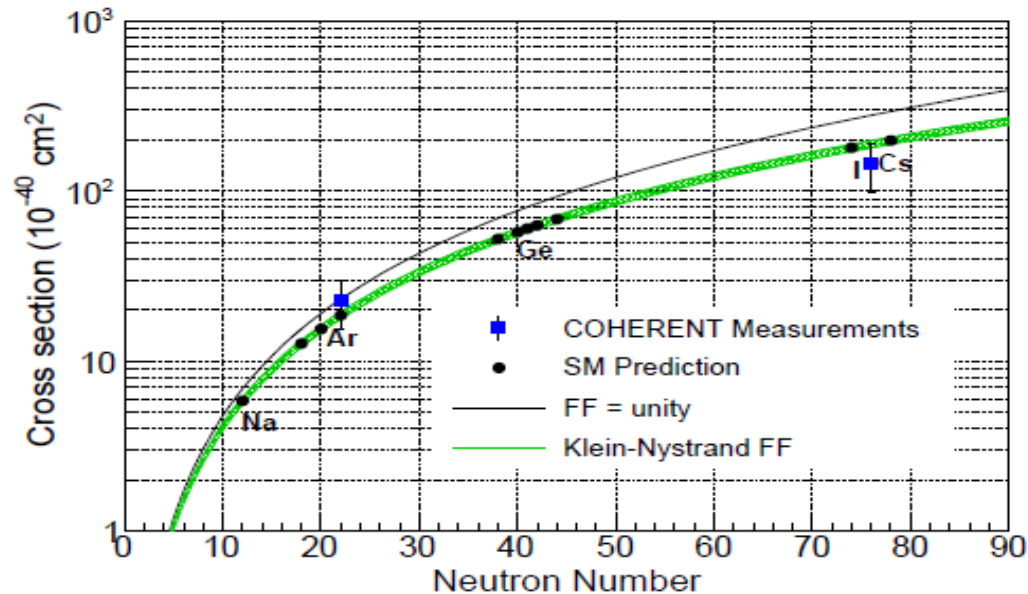
See the radiative correction in 2011.05960

The neutron contribution is dominant! $\implies \frac{d\sigma_{\text{CE}\nu\text{NS}}}{dT} \propto N^2$

[Freedman, PRD 9 (1974) 1389; Drukier, Stodolsky, PRD 30 (1984) 2295; Barranco, Miranda, Rashba, hep-ph/0508299]

- ▶ The coherent nuclear recoil gives a big cross section enhancement for heavy nuclei: $\sigma_{\text{NC}}^{\text{incoherent}} \propto N \implies \sigma_{\text{CE}\nu\text{NS}}/\sigma_{\text{NC}}^{\text{incoherent}} \propto N$
- ▶ The nuclear form factors $F_N(|\vec{q}|)$ and $F_Z(|\vec{q}|)$ describe the loss of coherence for $|\vec{q}|R \gtrsim 1$. [Patton et al, arXiv:1207.0693; Bednyakov, Naumov, arXiv:1806.08768; Papoulias et al, arXiv:1903.03722; Ciuffoli et al, arXiv:1801.02166; Canas et al, arXiv:1911.09831; Van Dessel et al, arXiv:2007.03658]

Neutron Form Factor



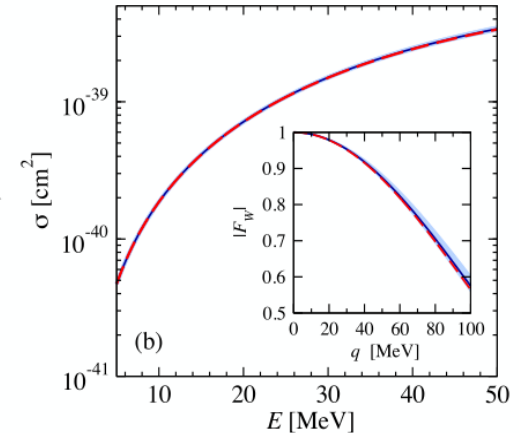
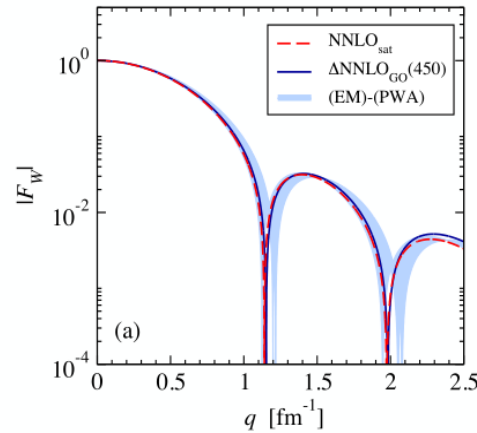
[COHERENT, arXiv:2003.10630]

- ▶ Partial coherency is described by the **nuclear neutron form factor** $F_N(|\vec{q}|)$
- ▶ Fourier transform of the **neutron distribution in the nucleus** $\rho_N(r)$:
$$F_N(|\vec{q}|) = \int e^{-i\vec{q}\cdot\vec{r}} \rho_N(r) d^3r$$
- ▶ Measurable parameter: the radius R_n of the nuclear neutron distribution

Neutron Form Factor: nuclear inputs

Weak form-factor of ^{40}Ar

- Ab initio method (coupled-cluster theory)
- various nuclear potentials

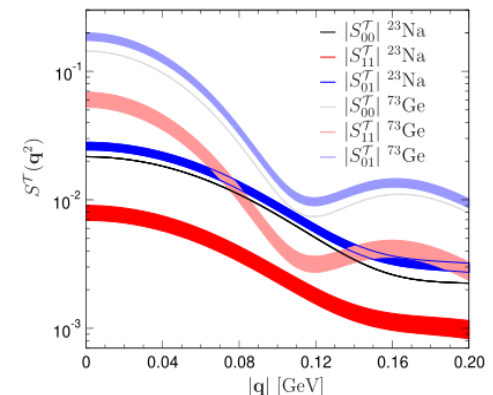
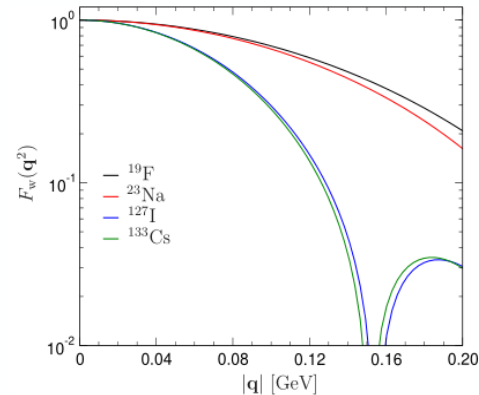


C. Payne et al.

Phys.Rev.C 100 (2019) 6, 061304

Shell model calculations

- Generalisation for beyond SM — new nuclear responses



M. Hoefliger, J. Menendez, A. Schwenk

Phys.Rev.D 102 (2020) 7, 074018

Neutron Form Factor: neutron radius

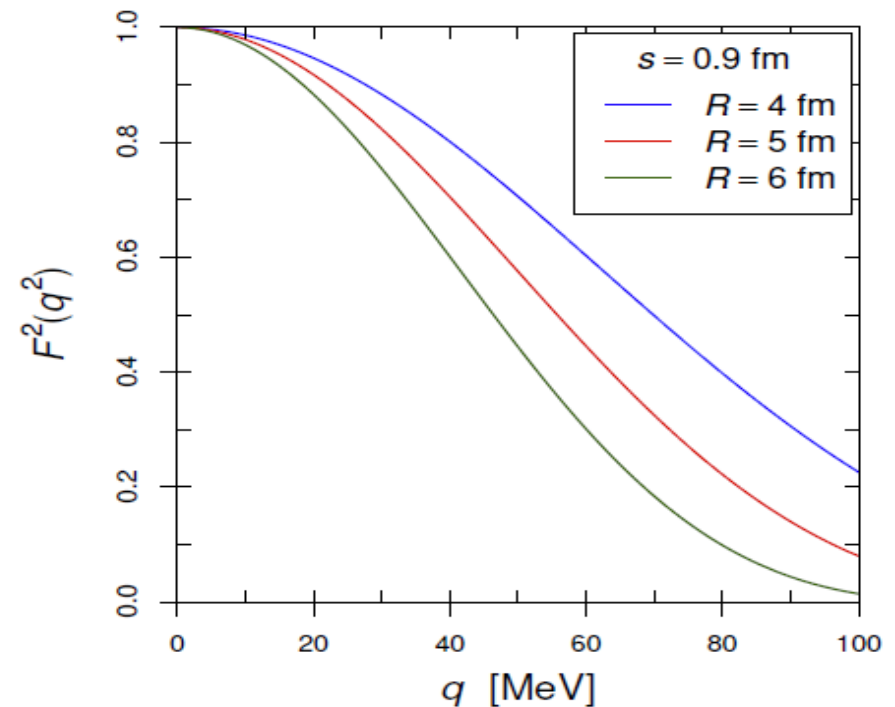
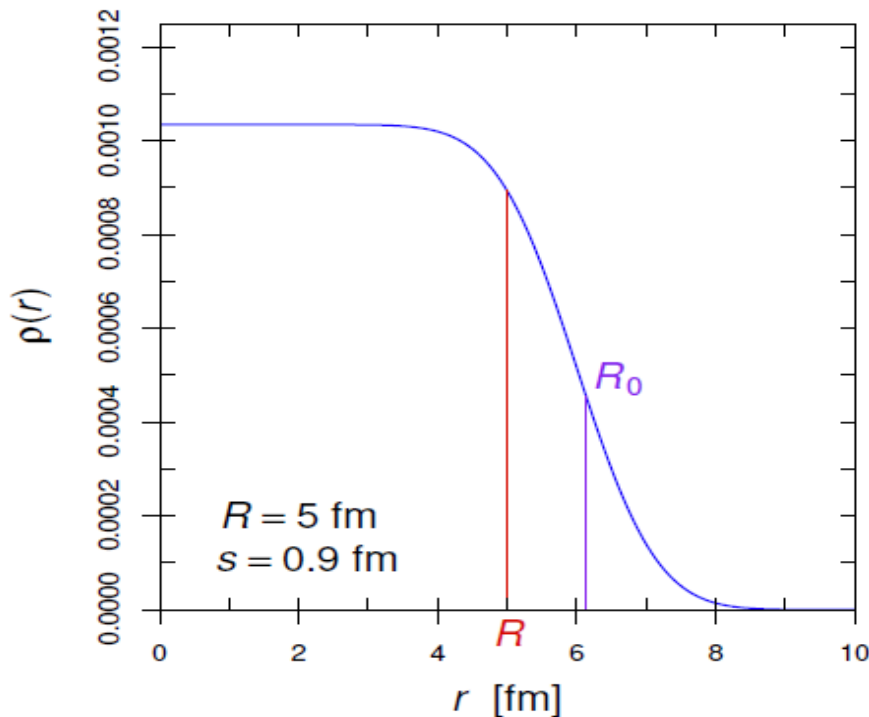
Helm form factor:
$$F_N^{\text{Helm}}(|\vec{q}|^2) = 3 \frac{j_1(|\vec{q}|R_0)}{|\vec{q}|R_0} e^{-|\vec{q}|^2 s^2/2}$$

Spherical Bessel function of order one: $j_1(x) = \sin(x)/x^2 - \cos(x)/x$

Obtained from the convolution of a sphere with constant density with radius R_0 and a gaussian density with standard deviation s

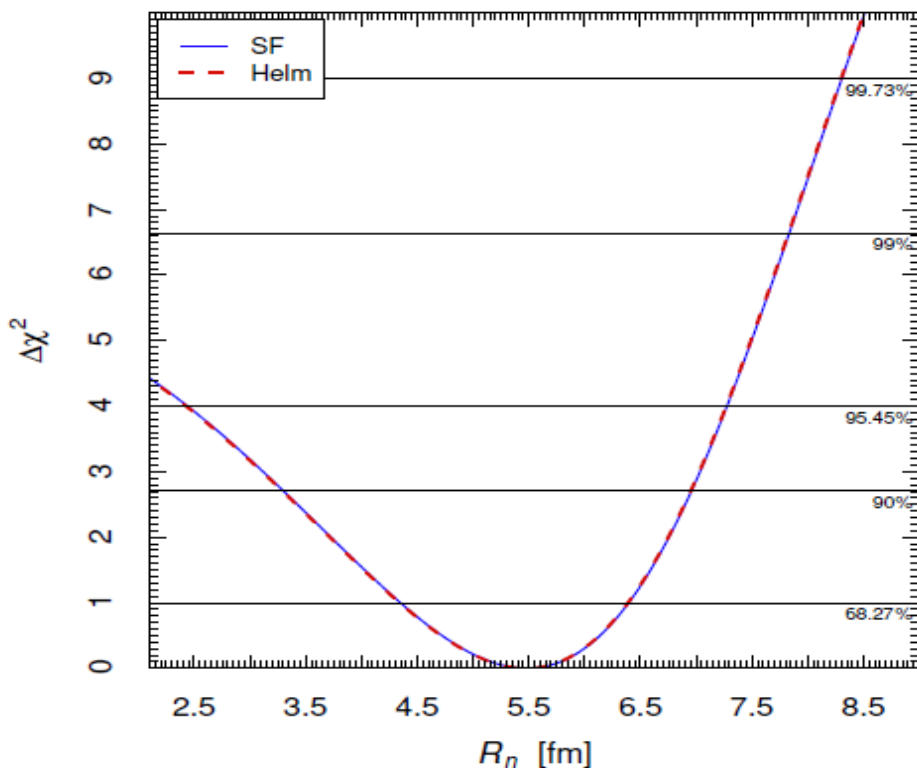
Rms radius: $R^2 = \langle r^2 \rangle = \frac{3}{5} R_0^2 + 3s^2$

Surface thickness: $s \simeq 0.9 \text{ fm}$



Neutron Distributions of Cs & I

- Fit of the 2017 COHERENT CsI data to get $R_n(^{133}\text{Cs}) \simeq R_n(^{127}\text{I})$:



First determination of R_n with neutrino-nucleus scattering:

$$R_n(\text{CsI}) = 5.5^{+0.9}_{-1.1} \text{ fm}$$

[Cadeddu, Giunti, Li, Zhang, arXiv:1710.02730]

- With new 2020 COHERENT CsI data:

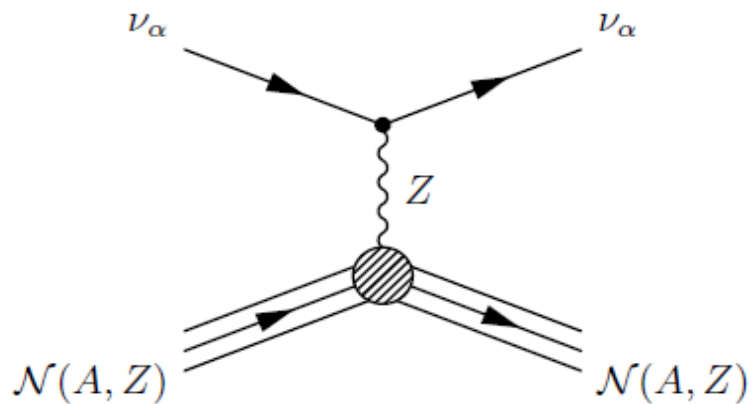
[Pershey @ Magnificent CE ν NS 2020]

$$R_n(\text{CsI}) = 5.55 \pm 0.44 \text{ fm}$$

[Cadeddu et al, arXiv:2102.06153]

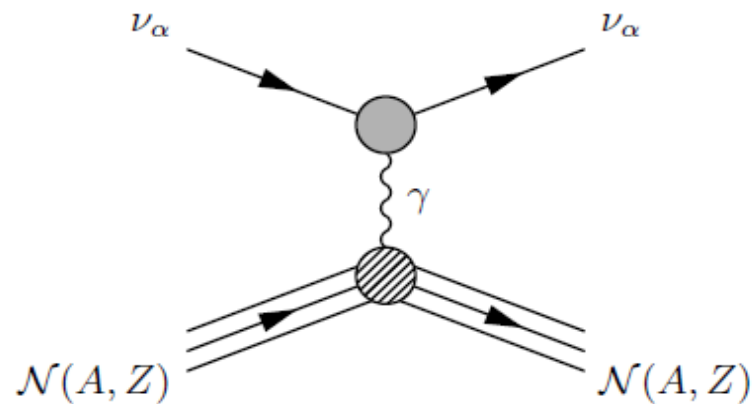
BSM Neutrino Interactions in CEvNS

Standard Model NC

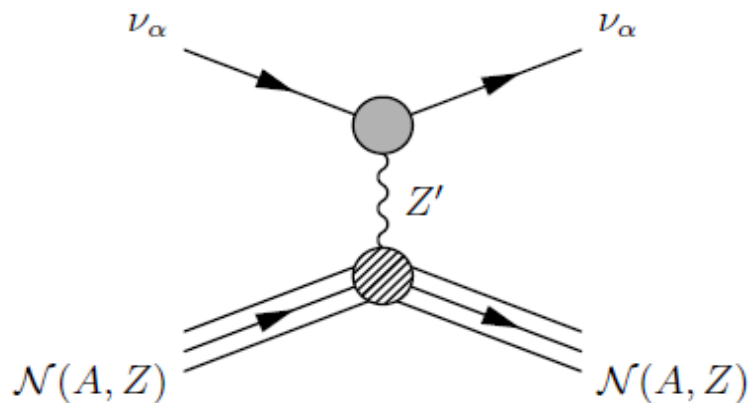


+

Electromagnetic Interactions

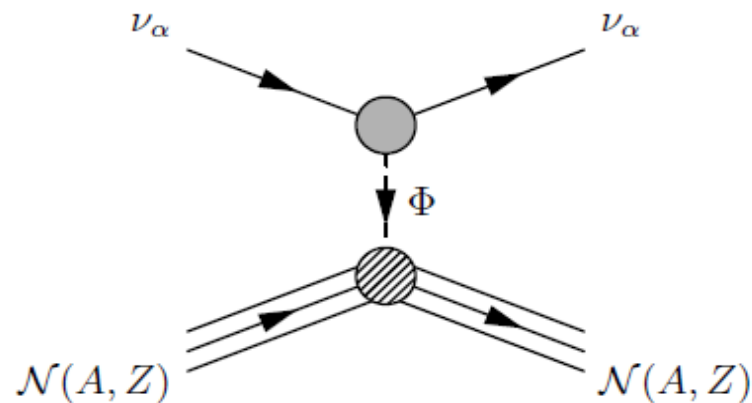


BSM Vector Mediator



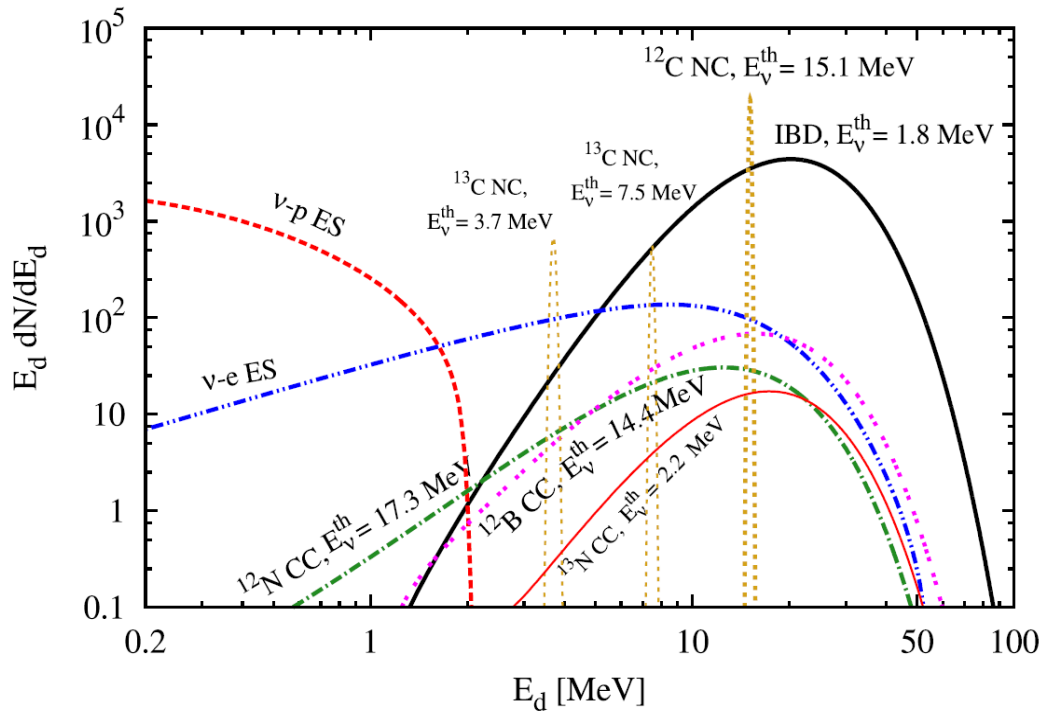
+

BSM Scalar Mediator



**B: low energy CC:
QE on the Target Nuclei**

(Quasi-)elastic ν -nucleus CC/NC interactions



Experiment	Nuclear Target	Reaction	σ_0 [10^{-46}cm^2]	ΔE_{nucl} [MeV] (no det. Thres.)
GALLEX/GNO SAGE	$^{71}\text{Ga}_{33}$	$\nu_e + ^{71}\text{Ga} \rightarrow e^- + ^{71}\text{Ge}$	$8.611 \pm 0.4\%$ (GT)	0.2327
HOMESTAKE	$^{37}\text{Cl}_{17}$	$\nu_e + ^{37}\text{Cl} \rightarrow e^- + ^{37}\text{Ar}$	1725 (F)	0.814
SNO	$^2\text{H}_1$	$\nu_e + ^2\text{H} \rightarrow e^- + p + p$	(GT)	1.442
DUNE, ICARUS, etc.	$^{40}\text{Ar}_{18}$	$\nu_e + ^{40}\text{Ar} \rightarrow e^- + ^{40}\text{K}^*$	148.58 (F) ... 44.367 (GT ₂) ... 41.567 (GT ₆) ...	1.505 +

From Kevin McFarland

➤ Important for solar & supernova neutrino detection

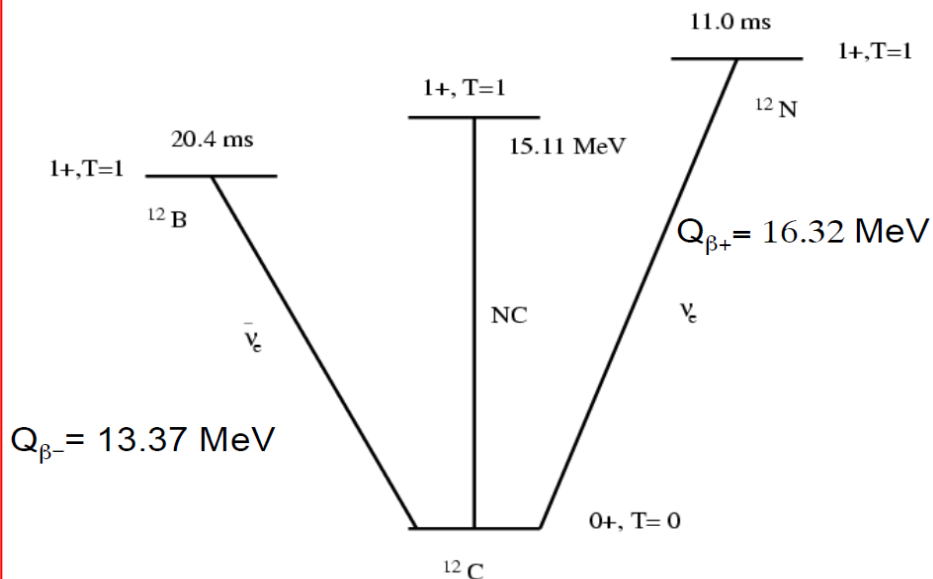
JUNO, Prog.Part.Nucl.Phys. 123 (2022) 103927

Channels	Threshold [MeV]	Signal	Event numbers	
			[200 kt×yrs]	after cuts
CC $\nu_e + ^{13}\text{C} \rightarrow e^- + ^{13}\text{N}(\frac{1}{2}^-; \text{gnd})$	2.2 MeV	$e^- + ^{13}\text{N}$ decay	3929	647
NC $\nu_x + ^{13}\text{C} \rightarrow \nu_x + ^{13}\text{C}(\frac{3}{2}^-; 3.685 \text{ MeV})$	3.685 MeV	γ	3032	738
ES $\nu_x + e \rightarrow \nu_x + e$	0	e^-	3.0×10^5	6.0×10^4

JUNO, Astrophys.J. 965 (2024) 2, 122

(Quasi-)elastic ν -nucleus CC/NC interactions

A=12 triad



Nuclear structure effects:

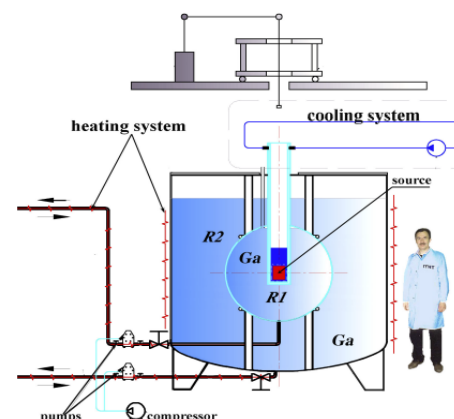
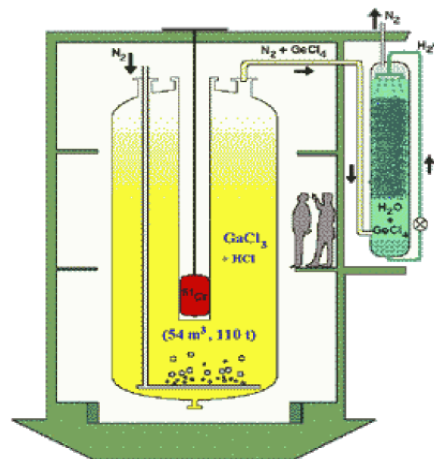
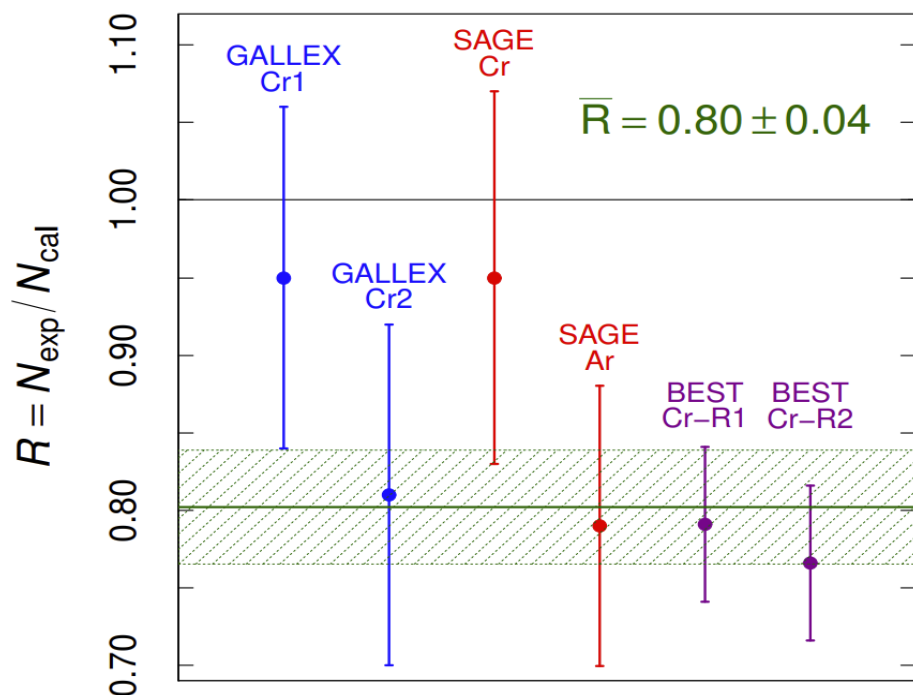
- beta (M1) decay calibration
- ν -energy \sim nuclear excitation energy: shell model
- giant resonance: CRPA or shell model
- $> 100 \text{ MeV}$, fermi Gas models or spectral function method
- DIS region: Parton

From Vogel, NPA 777 (2006) 340-355

	$^{12}\text{C}(\nu_e, e^-)^{12}\text{N}_{\text{gs}}$ decay at rest	$^{12}\text{C}(\nu_\mu, \mu^-)^{12}\text{N}_{\text{gs}}$ decay in flight	$^{12}\text{C}(\nu_e, e^-)^{12}\text{C}(15.11)$ decay at rest
Experiment [31]	$9.4 \pm 0.5 \pm 0.8$	–	$11 \pm 0.85 \pm 1.0$
Experiment [32]	$9.1 \pm 0.4 \pm 0.9$	$66 \pm 10 \pm 10$	–
Experiment [33]	$10.5 \pm 1.0 \pm 1.0$	–	–
Shell model [36]	9.1	63.5	9.8
CRPA [34,35]	8.9	63.0	10.5
EPT [37]	9.2	59	9.9

Gallium Anomaly

Gallium Radioactive Source Experiments: GALLEX, SAGE, BEST (2021)



$\approx 5\text{-}6\sigma$ deficit \implies Anomaly!

$\langle L \rangle_{\text{GALLEX}} \simeq 1.9 \text{ m}$ $\langle L \rangle_{\text{SAGE}} \simeq 0.6 \text{ m}$
 $\langle L \rangle_{\text{BEST}}^{\text{R1}} \simeq 0.7 \text{ m}$ $\langle L \rangle_{\text{BEST}}^{\text{R2}} \simeq 1.1 \text{ m}$
 $\Delta m_{\text{SBL}}^2 \gtrsim 1 \text{ eV}^2 \gg \Delta m_{\text{ATM}}^2$

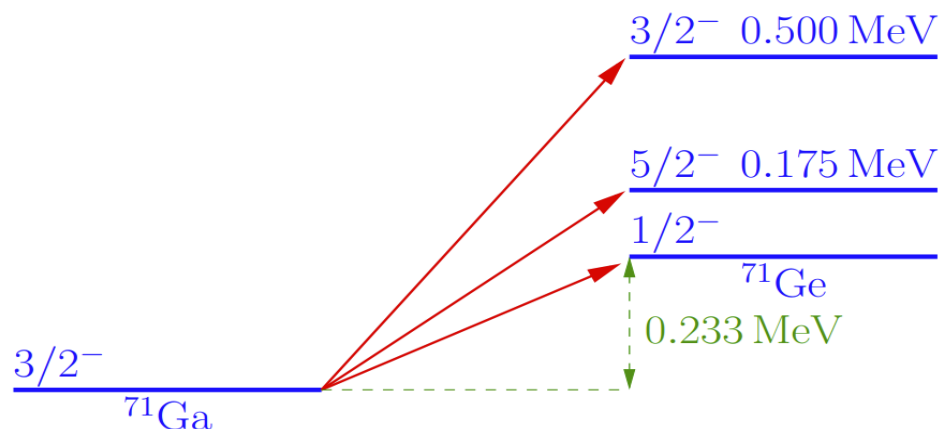
[SAGE, arXiv:nucl-ex/0512041, arXiv:0901.2200; Laveder et al, NPPS 168 (2007) 344, arXiv:hep-ph/0610352, arXiv:0711.4222, arXiv:1006.3244; Kostensalo et al, arXiv:1906.10980; BEST, arXiv:2109.11482, arXiv:2109.14654; Berryman et al, arXiv:2111.12530]

Cross section calculation

- ▶ A deficit could be due to an **overestimate** of

$$\sigma(\nu_e + {}^{71}\text{Ga} \rightarrow {}^{71}\text{Ge} + e^-)$$

- ▶ First calculation: Bahcall, PRC 56 (1997) 3391, hep-ph/9710491



- ▶ $\sigma_{\text{G.S.}}$ from $T_{1/2}({}^{71}\text{Ge}) = 11.43 \pm 0.03$ days [Hampel, Remsberg, PRC 31 (1985) 666]

$$\sigma_{\text{G.S.}}({}^{51}\text{Cr}) = (5.54 \pm 0.02) \times 10^{-45} \text{ cm}^2$$

- ▶ $\sigma({}^{51}\text{Cr}) = \sigma_{\text{G.S.}}({}^{51}\text{Cr}) \left(1 + 0.669 \frac{\text{BGT}_{175}}{\text{BGT}_{\text{G.S.}}} + 0.220 \frac{\text{BGT}_{500}}{\text{BGT}_{\text{G.S.}}} \right)$

- ▶ The contribution of **excited states** is only $\sim 5\%$!

[Bahcall, hep-ph/9710491]

Cross section calculation: excited states

$\nu_e + {}^{71}\text{Ga} \rightarrow {}^{71}\text{Ge} + e^-$ cross sections in units of 10^{-45} cm^2 :

		${}^{51}\text{Cr}$		${}^{37}\text{Ar}$		\bar{R}	GA
		σ_{tot}	δ_{exc}	σ_{tot}	δ_{exc}		
Ground State <small>[Phys.Atom.Nucl. 83 (2020) 1549]</small>	$T_{1/2}({}^{71}\text{Ge})$	5.539 ± 0.019	—	6.625 ± 0.023	—	0.844 ± 0.031	5.0σ
Bahcall <small>[hep-ph/9710491]</small>	${}^{71}\text{Ga}(p, n){}^{71}\text{Ge}$	5.81 ± 0.16	4.7%	7.00 ± 0.21	5.4%	0.802 ± 0.037	5.4σ
Kostensalo et al. <small>[arXiv:1906.10980]</small>	Shell Model	5.67 ± 0.06	2.3%	6.80 ± 0.08	2.6%	0.824 ± 0.031	5.6σ
Semenov <small>[Phys.Atom.Nucl. 83 (2020) 1549]</small>	${}^{71}\text{Ga}({}^3\text{He}, {}^3\text{H}){}^{71}\text{Ge}$	5.938 ± 0.116	6.7%	7.169 ± 0.147	7.6%	0.786 ± 0.033	6.6σ

Giunti, YFL, Ternes, Xin, arXiv: 2212.09722

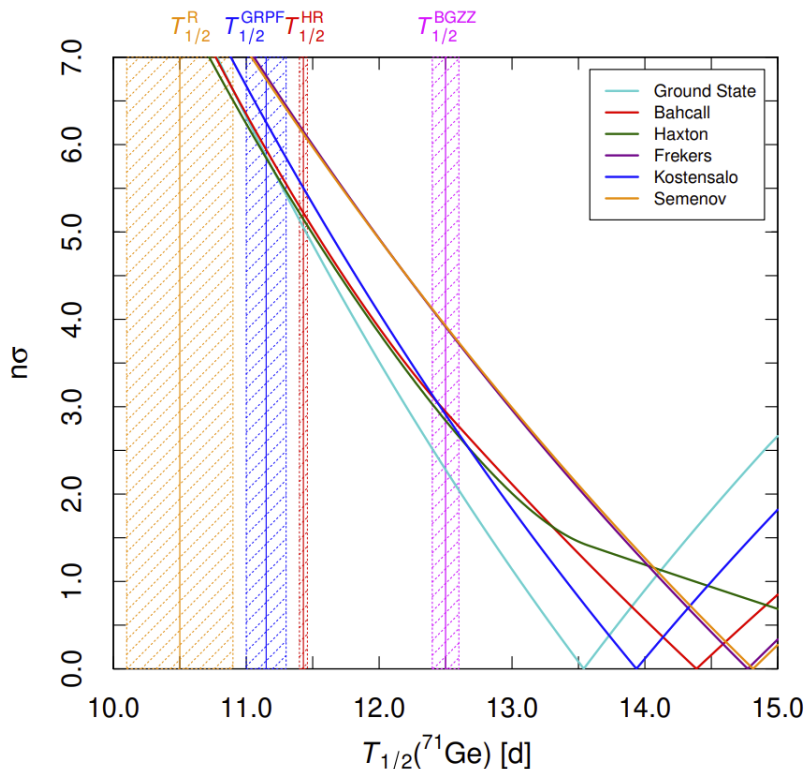
Cross section calculation: ground state

$$T_{1/2}^{\text{BGZZ}}(^{71}\text{Ge}) = 12.5 \pm 0.1 \text{ d} \quad (\text{Bisi, Germagnoli, Zappa, and Zimmer, 1955}) [39],$$

$$T_{1/2}^{\text{R}}(^{71}\text{Ge}) = 10.5 \pm 0.4 \text{ d} \quad (\text{Rudstam, 1956}) [40], \quad \text{Giunti, YFL, Ternes, Xin, arXiv: 2212.09722}$$

$$T_{1/2}^{\text{GRPF}}(^{71}\text{Ge}) = 11.15 \pm 0.15 \text{ d} \quad (\text{Genz, Renier, Pengra, and Fink, 1971}) [41],$$

$$T_{1/2}^{\text{HR}}(^{71}\text{Ge}) = 11.43 \pm 0.03 \text{ d} \quad (\text{Hempel and Remsberg, 1985}) [42].$$



➤ **An enlarged life-time will reduce or eliminate the anomaly.**

➤ **Triggered an active campaign of re-measurement !**

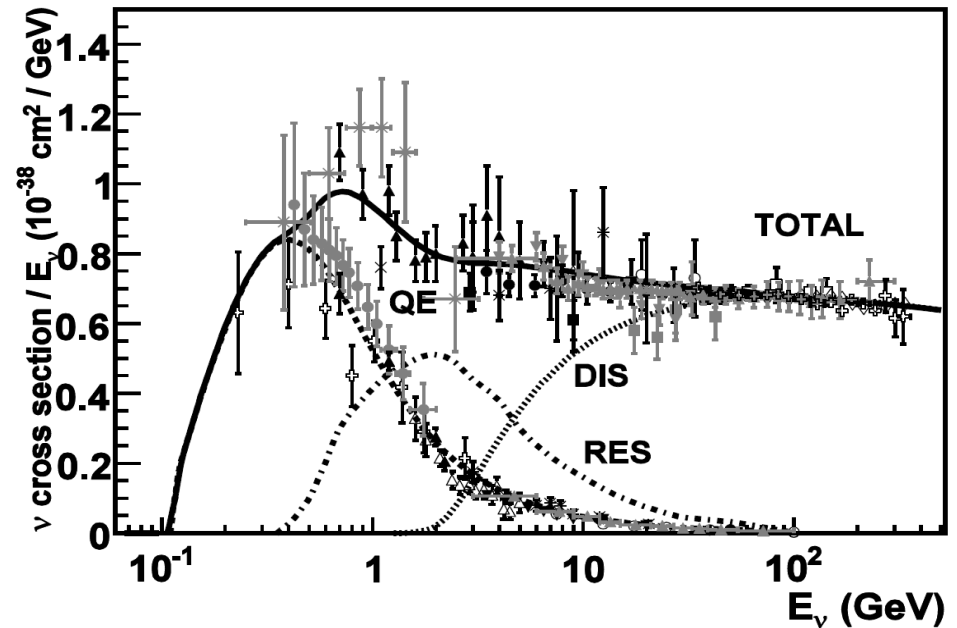
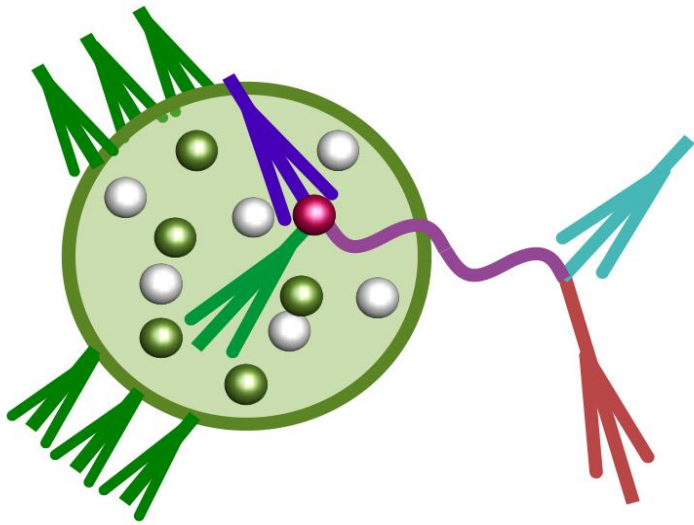
11.43 ± 0.03 days (2307.05353)

11.468 ± 0.008 days (2401.15286)

➤ **The Gallium anomaly remains and deserves further theoretical and experimental investigations.**

**C: GeV range CC/NC:
accelerator and atmospheric neutrinos**

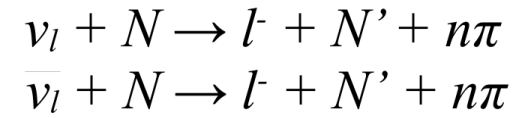
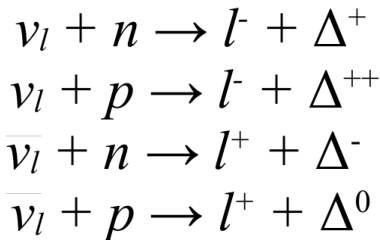
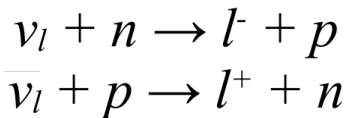
GeV neutrino interactions



quasielastic
scattering

resonance
production

deep-inelastic
scattering



**Fermi motion, binding
energy, M_A , 2p2h,**

Hadron production, FSI

Parton Model, FSI

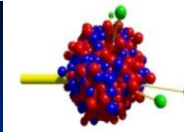
GeV neutrino interaction generators

Status overview

Marco Roda @ NUINT24



- Well established generator
 - Used by many experiments around the world
 - Main new addition is JUNO
 - Main generator for all the LAr experiments
- Two main efforts
 - Model development
 - Tuning
- Contacts, details and code are all available from our website: www.genie-mc.org/
- Latest release: version 3.04.02, released in April 2024
 - Previous release was 3.04.00, released in March 2023
 - <http://releases.genie-mc.org/>
- Recent publications
 - Neutrino-nucleon cross-section model tuning in GENIE v3 - *Phys.Rev.D 104 (2021) 7, 072009*
 - Hadronization model tuning in genie v3 - *Phys.Rev.D 105 (2022) 1, 012009*



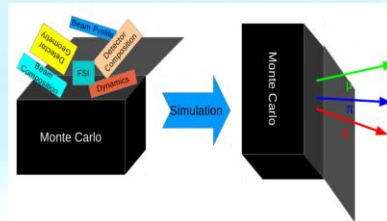
GiBUU

The Giessen Boltzmann-Uehling-Uhlenbeck Project
Ulrich Mosel and Kai Gallmeister @ NUINT24

- GiBUU is presently used to describe
 - Dilepton and pion production in heavy-ion collisions (HADES experiment at GSI)
 - Inelastic electron scattering at JLAB (and SLAC, MAMI)
 - Neutrino-nucleus reactions at Fermilab, T2K and FASER
- All with the same theory input and code!
- We provide the code for download from gibuu.hepforge.org,

NuWro - general information (1)

Jan T. Sobczyk @ NUINT24



- Monte Carlo generator of neutrino interactions
- Beginning ~ 2005 at the University of Wrocław
- Optimized for ~1 GeV
- Can handle all kind of targets, neutrino fluxes, equipped with detector interface
- Written in C++
- Output files in the ROOT format
- PYTHIA6 used for hadronization in DIS
- Open source code, repository: <https://github.com/NuWro/nuwro>



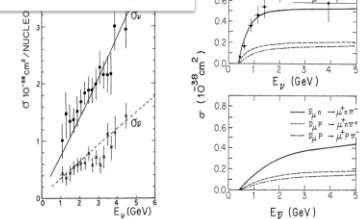
The NEUT neutrino interaction simulation program library

Yoshinari Hayato¹ and Luke Pickering² *The European Physical Journal Special Topics* volume 230, pages 4469–4481 (2021)

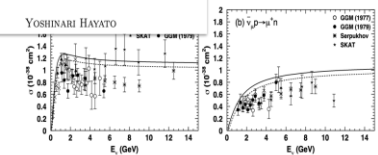
Patrick Stowell @ NUINT24

- MeV to TeV scale neutrino interaction generator originally created in the 70s to support neutrino backgrounds at Kamioka.
- Long history of development driven by evolving requirements of KamiokaNDE, Super-KamiokaNDE, and T2K.
- Currently the primary interaction generator for SK and T2K, used in all oscillation/cross-section analyses.
- See Laura, Stephen, Ulyesse, and Cesar's talks this NuINT!

Atmospheric Neutrino Background and Pion Nuclear Effect for KAMIOKA Nucleon Decay Experiment

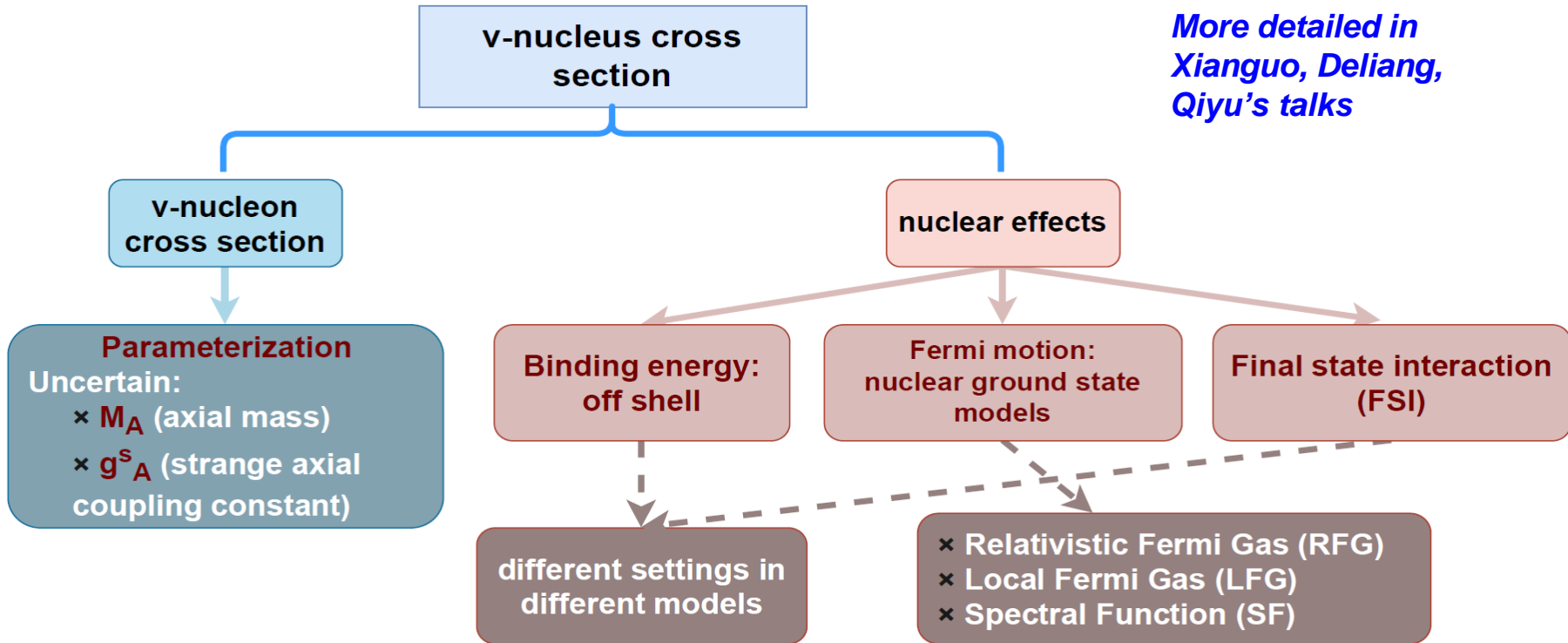


A NEUTRINO INTERACTION SIMULATION PROGRAM LIBRARY NEUT*



General components in generators

Brief summary of GeV neutrino interaction models



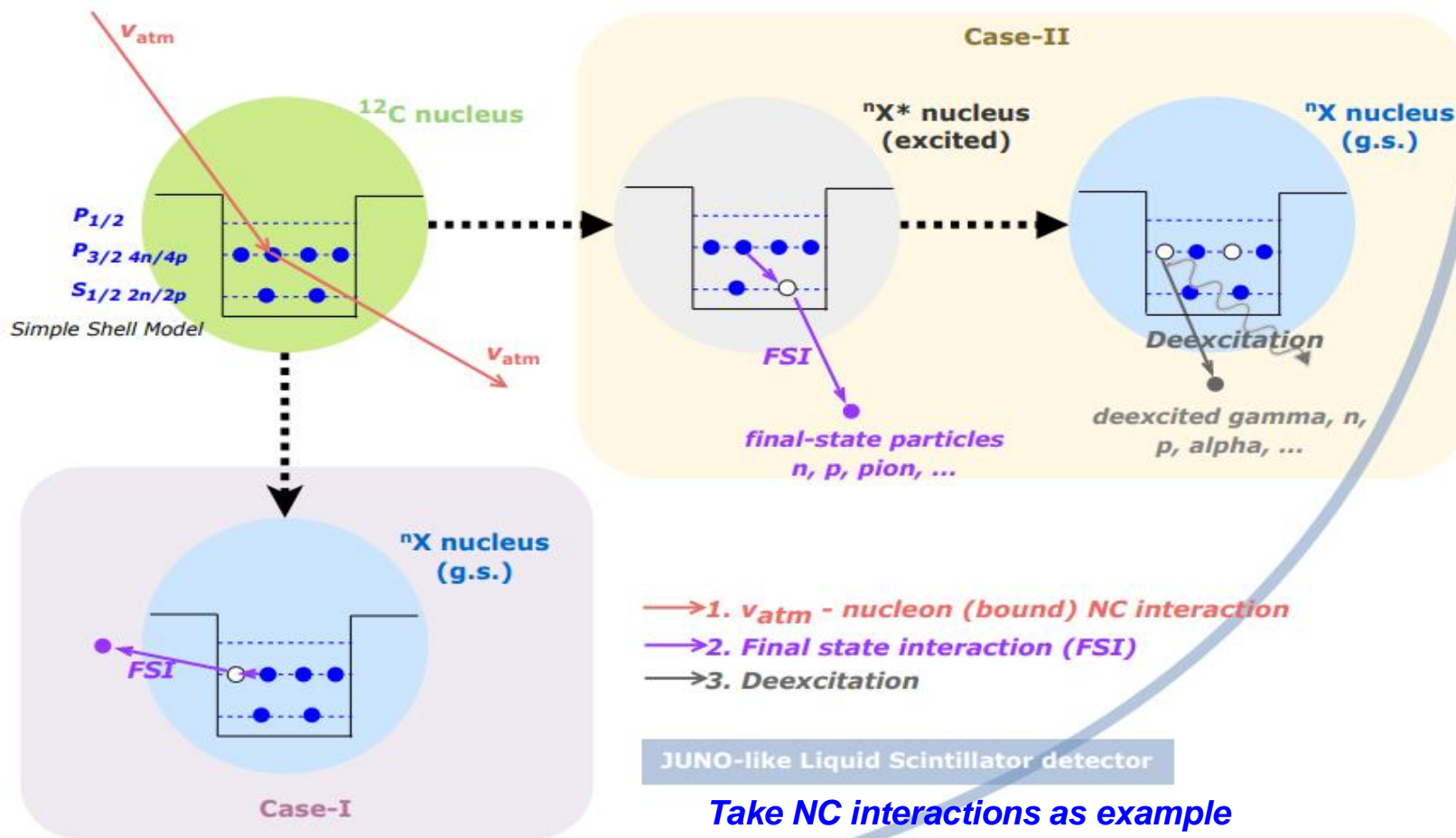
$$\frac{d\sigma_{\ell A}^{\text{IA}}}{d\omega d\Omega} = \sum_N \int d^3p dE P_{\text{hole}}^N(\mathbf{p}, E) \frac{M}{E_{\mathbf{p}}} \frac{d\sigma_{\ell N}^{\text{elem}}}{d\omega d\Omega} P_{\text{part}}^N(\mathbf{p}', \mathcal{T}')$$

average over the initial nucleon state

nucleon cross section

final-state interactions

New Methodology: adding deexcitation



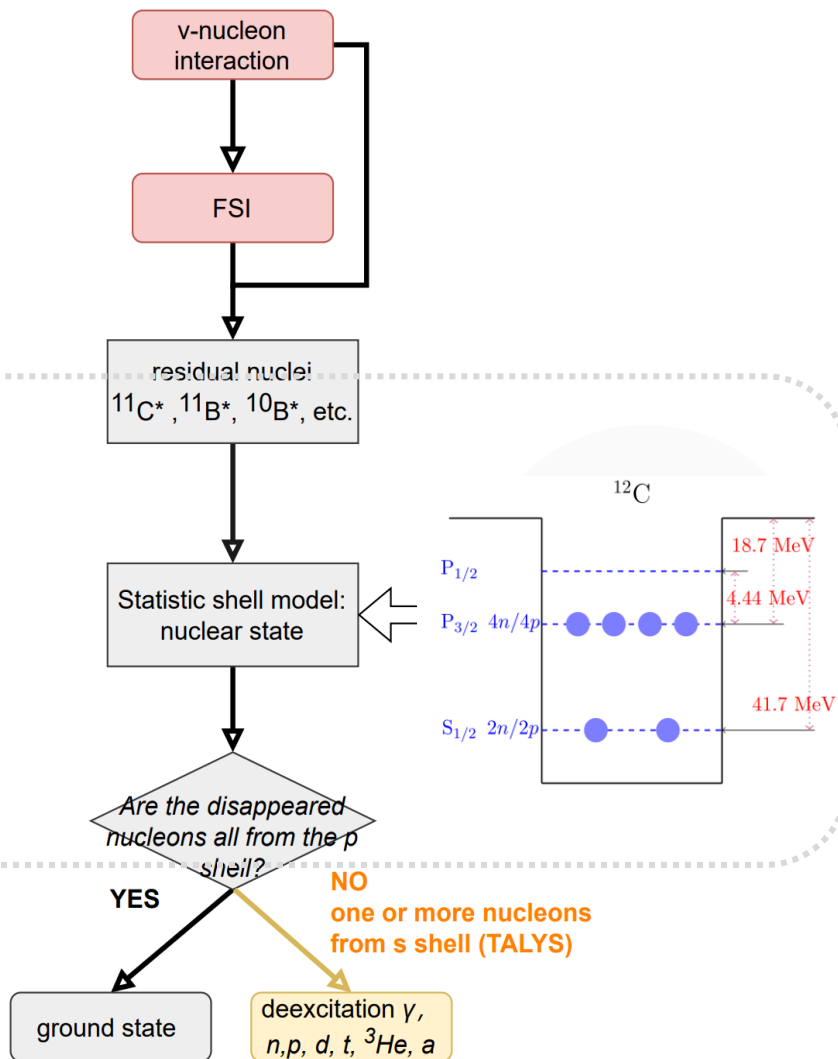
Take NC interactions as example

Cheng et al, *Phys. Rev. D* 103. 05001 (2021)

Cheng et al, arXiv 2404.07429

More detailed in Wanlei, Abe, Jie's talks

TALYS-based Deexcitation of Residual Nucleus



1. Simple shell model → Status of the residual nuclei

- All residual nuclei with $A > 5$ have been considered
- Taking $^{11}\text{C}^*$, $^{11}\text{B}^*$, $^{10}\text{C}^*$, $^{10}\text{Be}^*$ and $^{10}\text{B}^*$ for example

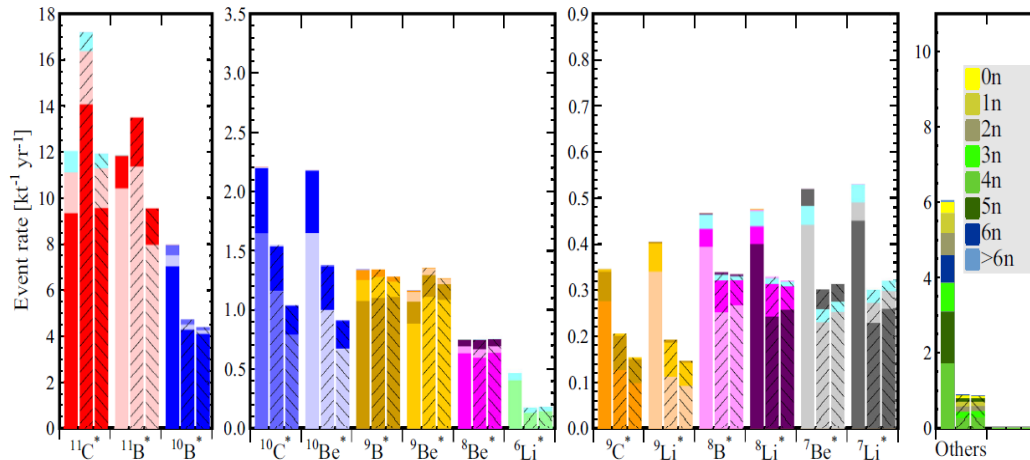
Daughter Nuclei	Shell Hole	Configuration Probability	Excitation Energy
$^{11}\text{C}^*$ or $^{11}\text{B}^*$	$s_{1/2}$	1/3	$E^* = 23$ MeV
	$p_{3/2}$	2/3	$E^* = 0$ MeV
$^{10}\text{C}^*$ or $^{10}\text{Be}^*$	$s_{1/2}$	1/15	$E^* = 46$ MeV
	$p_{3/2}$	6/15	$E^* = 0$ MeV
$^{10}\text{B}^*$	$s_{1/2}$ & $p_{3/2}$	8/15	$E^* = 23$ MeV
	$s_{1/2}$	1/9	$E^* = 46$ MeV
	$p_{3/2}$	4/9	$E^* = 0$ MeV
	$s_{1/2}$ & $p_{3/2}$	4/9	$E^* = 23$ MeV

Assume each neutron or proton has same possibility(1/6) to leave the shell.

More complicated shell information can be included.

Impact on exclusive cross sections

Before deexcitation



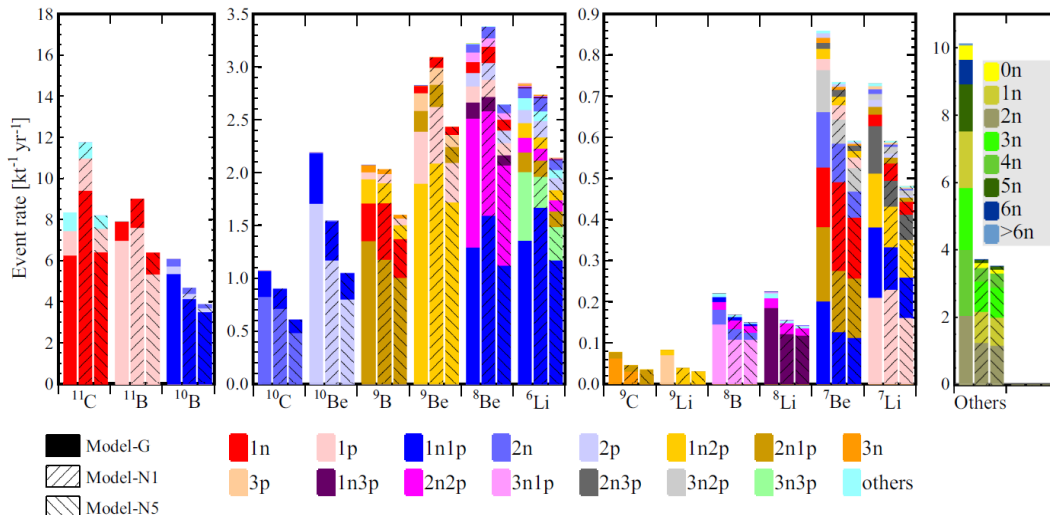
➤ ^{11}C , ^{11}B , ^{10}B reduced, and lighter nuclei increased; neutron multiplicity redistributed.

➤ Exclusive final-state information, such as **the neutron multiplicity, the charge pion multiplicity, the unstable nuclei**, is important for

(a) energy reconstruction

(b) tagging and reducing the backgrounds

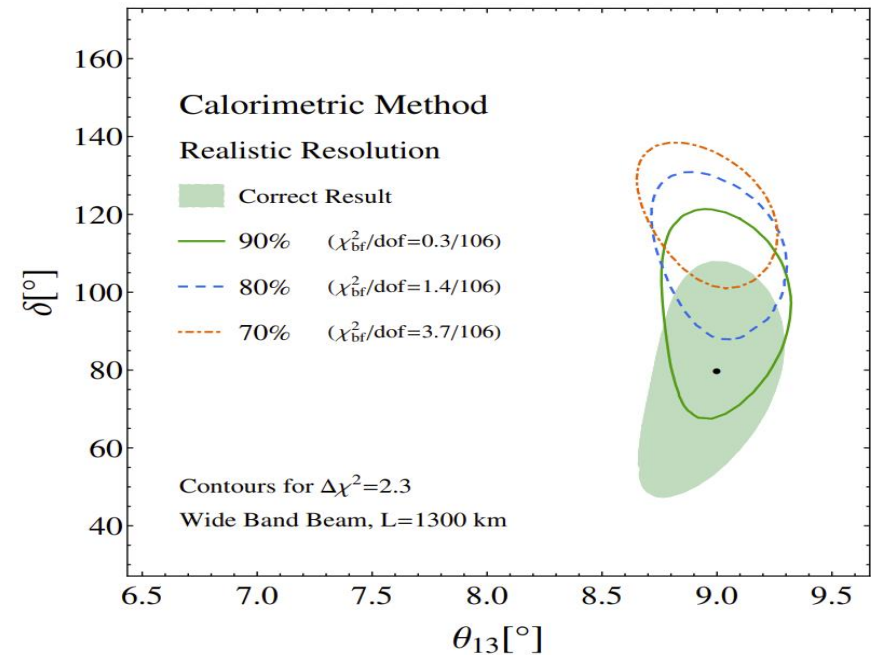
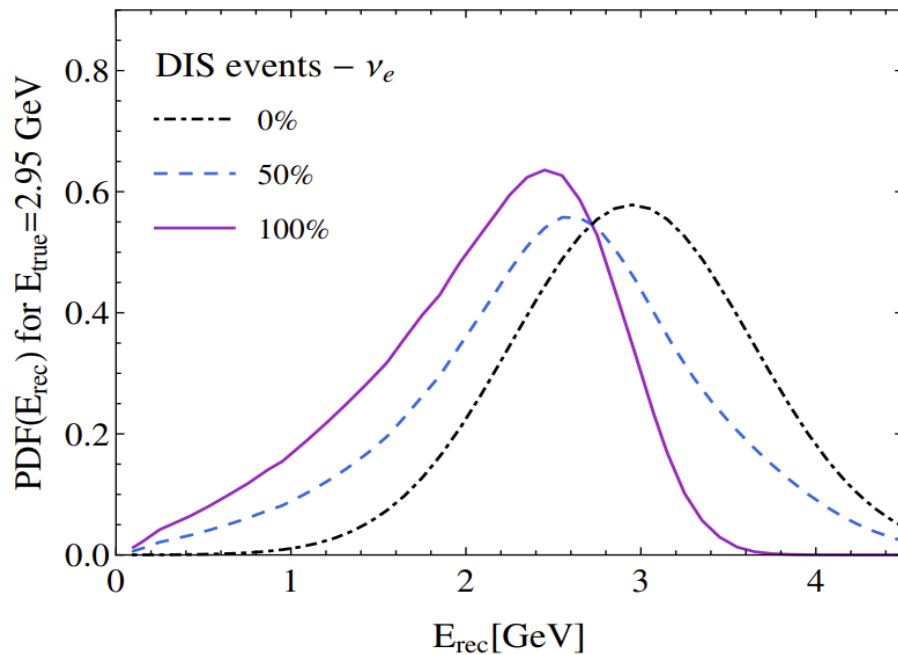
After deexcitation



Energy reconstruction in DUNE

$$E_\nu^{\text{cal}} = E_\ell + \sum_i T_i^N + \epsilon_n + \sum_j E_j,$$

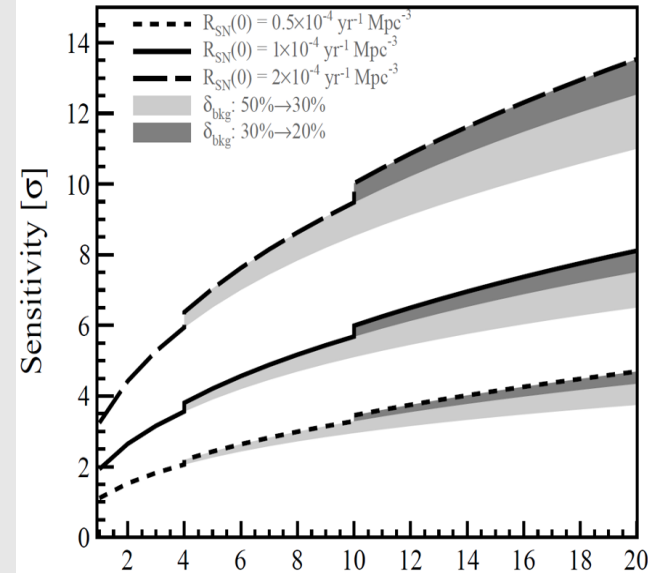
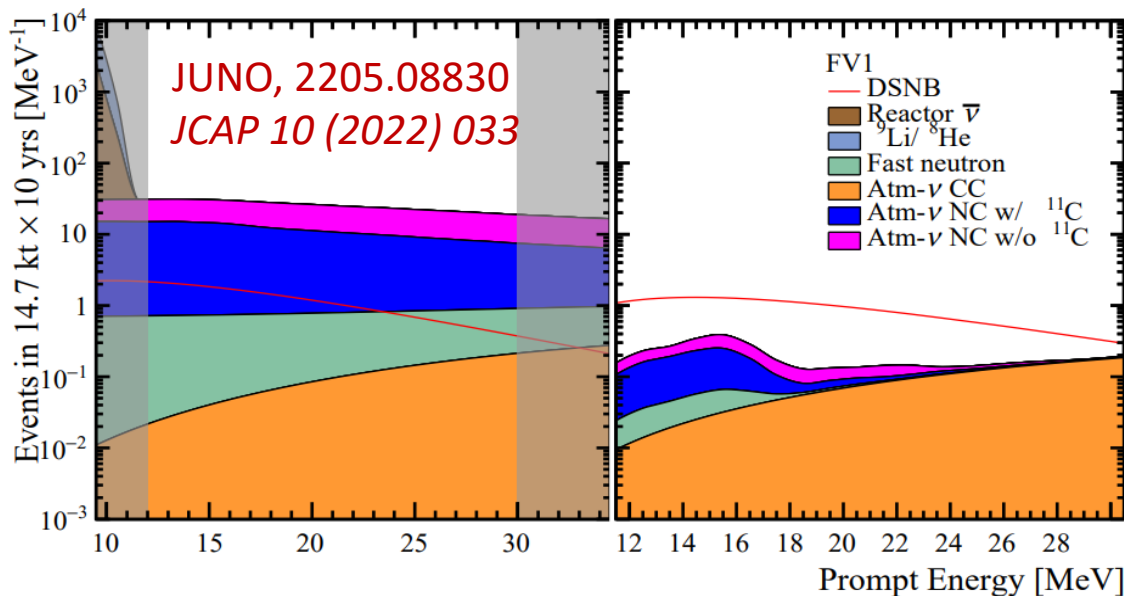
Ankowski, Coloma, Huber,
Mariani, Vagnoni, 1507.08561



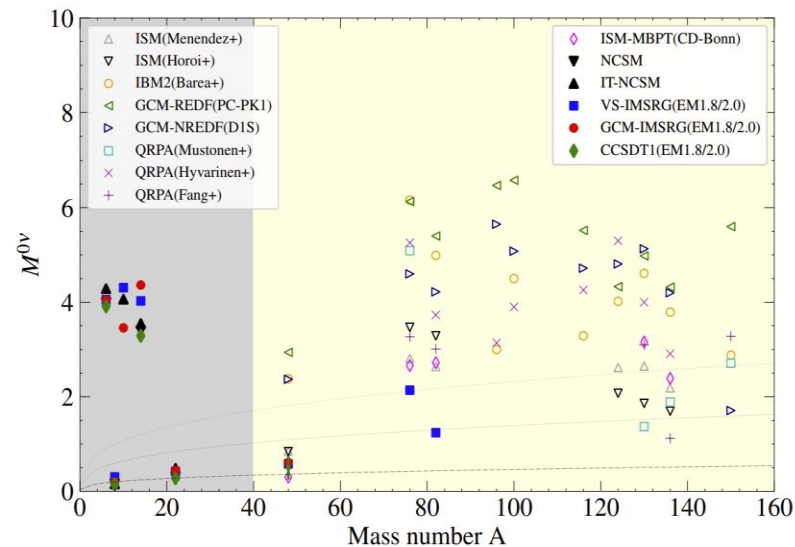
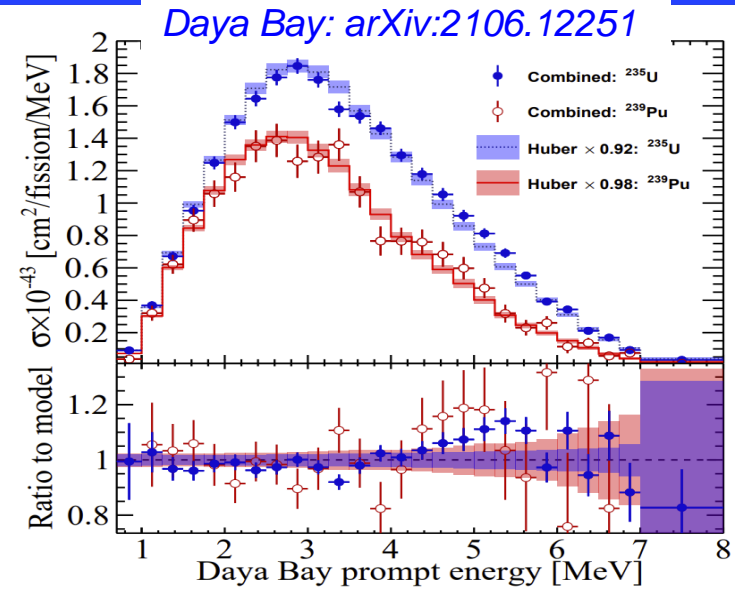
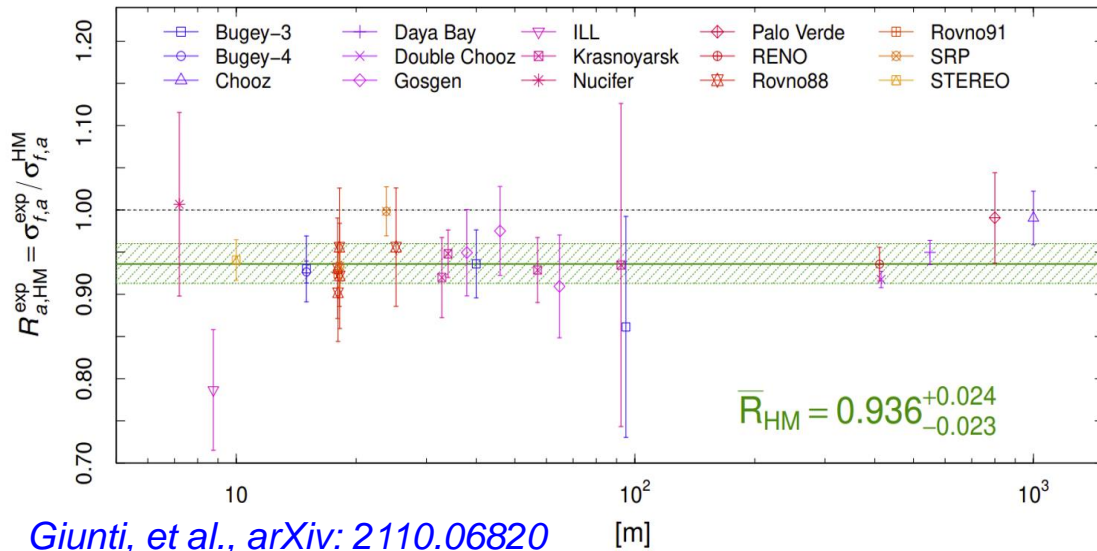
- The energy is reconstructed with calorimetric method.
- Missing neutrons (pions) may bias the energy and then result in wrong oscillation parameters.

Rare searches in JUNO

- **Diffuse Supernova Neutrino Background via IBD process: 2-4 events in JUNO per year**
- **Dominant backgrounds are from atmospheric neutrino NC interactions (20 times larger than the signal).**
- **A precise exclusive NC cross section is crucial (with neutron, ^{11}C) !**
- **Also pion and kaon production is important for proton decay search.**



Neutrino-nuclear connection **beyond cross sections**



➤ **Solution to Reactor antineutrino anomaly requires **precise energy spectrum of (forbidden) β decays!****

➤ **Search for neutrinoless double beta decay needs to **better control of the nuclear matrix element****

Yao et al., arXiv: 2111.15543

Conclusion

Neutrino interaction cross sections are **important prerequisite** to study **neutrino properties and new physics**.

Neutrino-lepton and neutrino (free-)nucleon interactions are relatively simple and widely used in low energy neutrino detection.

Electron scattering, & IBD of free protons

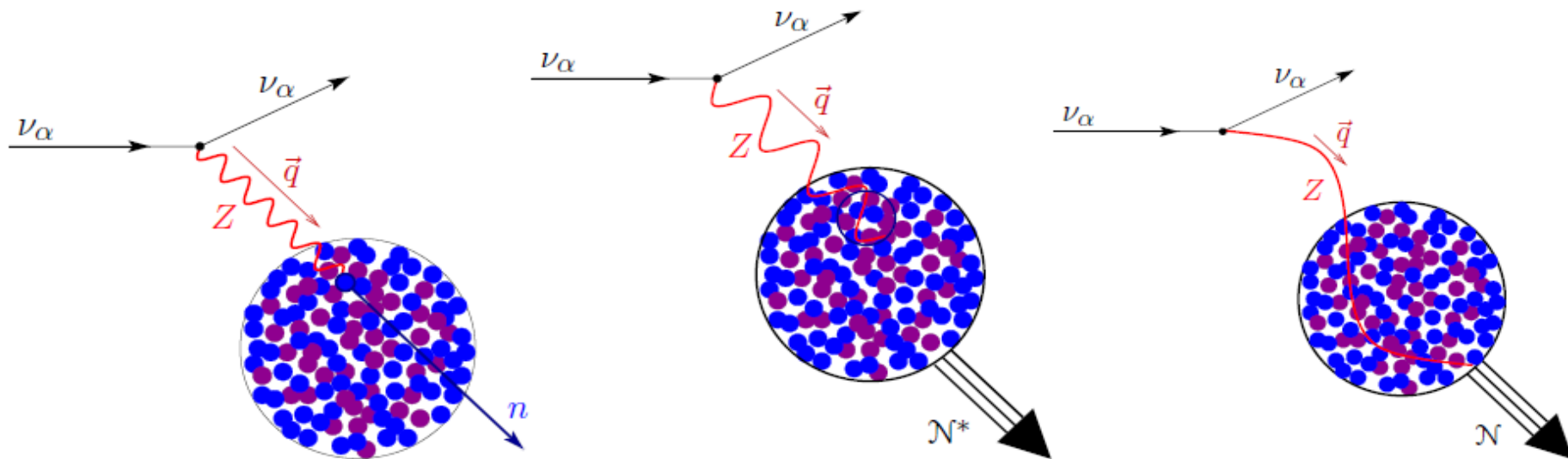
Neutrino-nucleus interactions:

From low to high energies, depend on different aspects of hadron and nuclear physics. **(Shell structure, Binding energy, Fermi Motion, Final state interactions, Deexcitation, Parton properties, etc.)**

Thank you !

Backup

Why Coherent ?



Inelastic Incoherent

$$\lambda_Z \ll R$$

Elastic Incoherent

$$\lambda_Z \lesssim R$$

Elastic Coherent

$$\lambda_Z \gtrsim 2R$$

$$\lambda_Z = 2\pi \frac{\hbar}{|\vec{q}|} \implies \text{CE}\nu\text{NS for } |\vec{q}| R \lesssim \hbar$$

$$|\vec{q}| R \lesssim 1$$

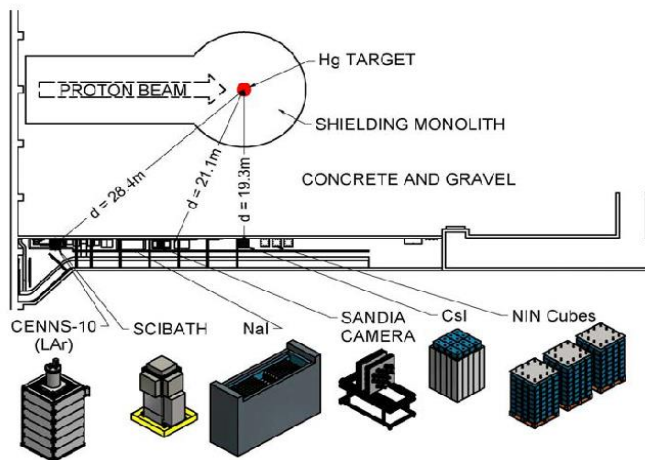
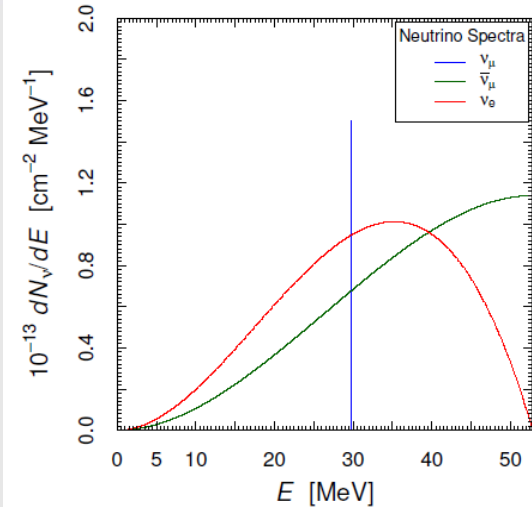
← Natural Units

The COHERENT experiment

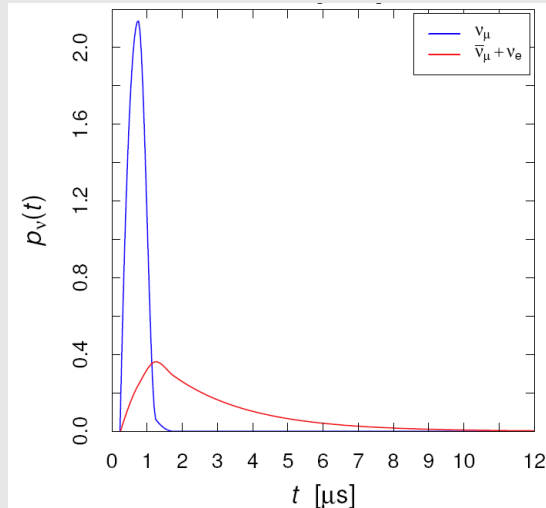
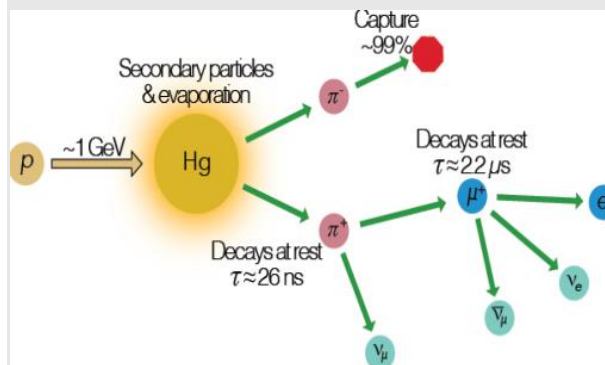
- 14.6 kg CsI scintillating crystal and 24 kg LAr detector.
- Prompt monochromatic ν_μ from stopped pion decays:

$$\pi^+ \rightarrow \mu^+ + \nu_\mu$$
- Delayed $\bar{\nu}_\mu$ and ν_e from the subsequent muon decays:

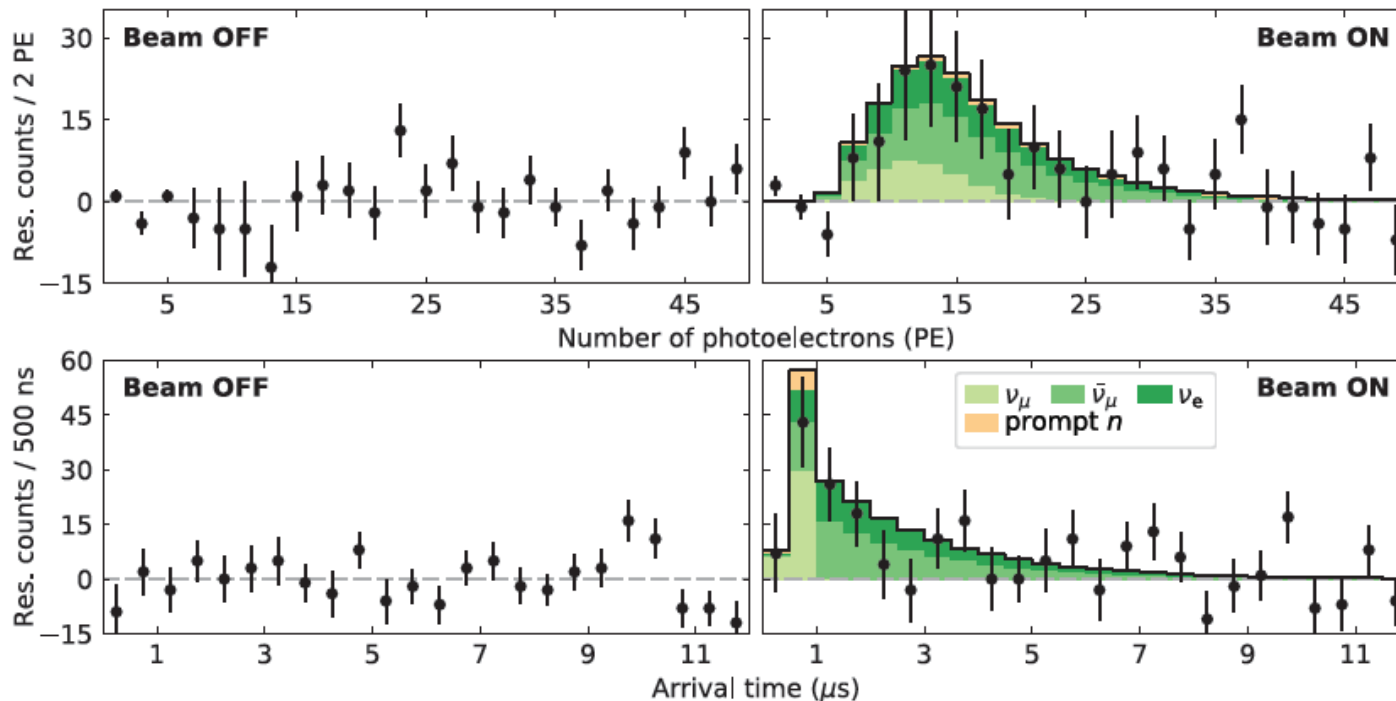
$$\mu^+ \rightarrow e^+ + \bar{\nu}_\mu + \nu_e$$
- The COHERENT energy and time information allow us to distinguish the interactions of ν_e , ν_μ and $\bar{\nu}_\mu$



[COHERENT, arXiv:1803.09183]



First observation of CEvNS at CsI (2017)

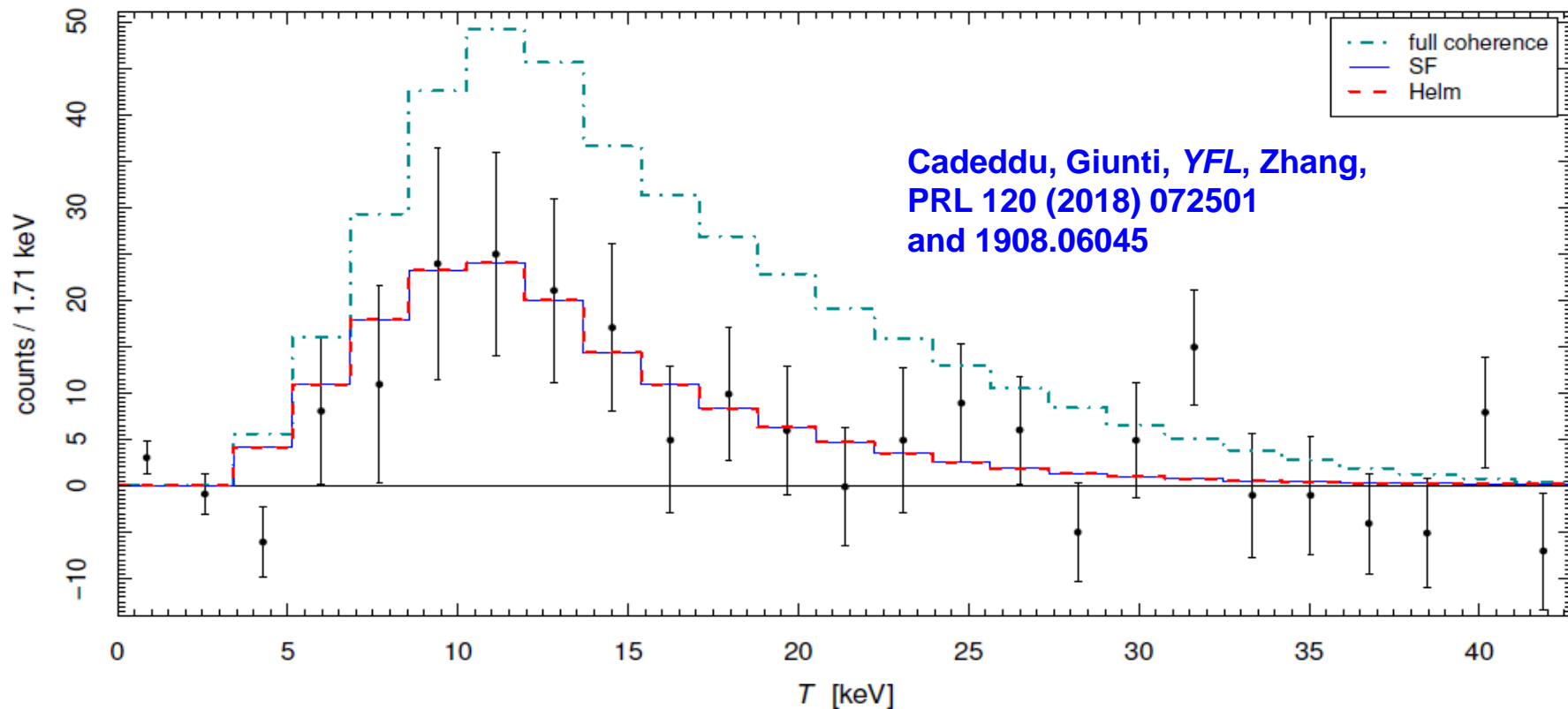


Akimov et al. *Science*
Vol 357, Issue 6356
15 September 2017

- Data are beam coincident and anti-coincident residuals during SNS operation, “On”, and during SNS shutdown periods, “Off”.
- Excess in light yield and timing distributions only for Beam on.

$^{133}_{55}\text{Cs}_{78}$ and $^{127}_{53}\text{I}_{74}$ ← Heavy nuclei well suited for CE ν NS

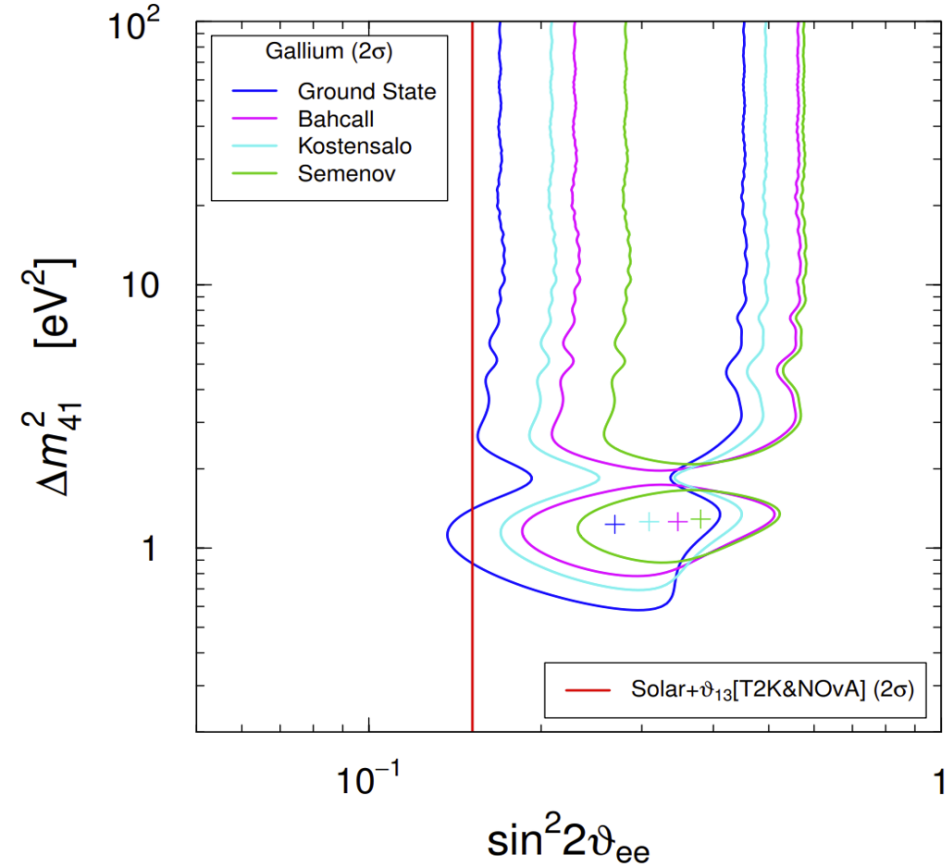
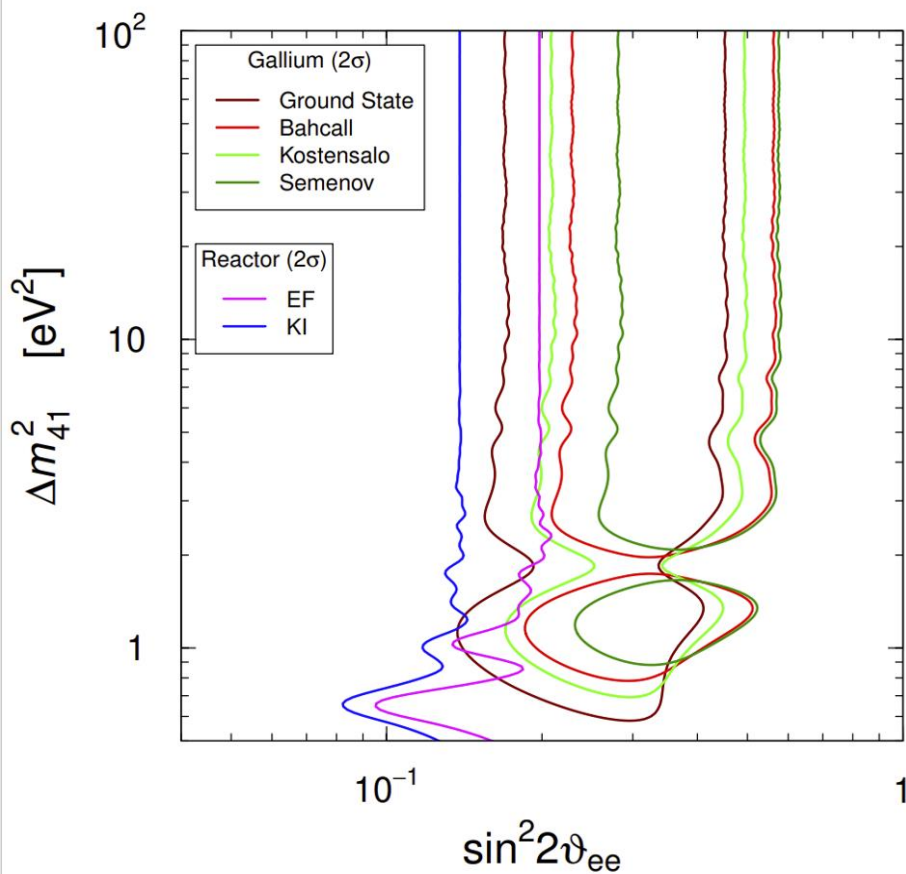
Test of Coherency



(1) Full coherence $\rightarrow F(\text{proton}) = F(\text{neutron}) = 1$.

(2) COHERENT data show **3.7** sigma evidence of the nuclear structure suppression of the full coherence

3+1 mixing ?



- **No 3+1 neutrino mixing and oscillation solution**
- **No CPT violation solution**

Giunti, YFL, Ternes, Xin, arXiv: 2209.00916

➔ **Source problem?**